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## Selective Ion Acceleration by Intense Radiation Pressure

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We report on the selective acceleration of carbon ions during the interaction of ultrashort, circularly polarized and contrast-enhanced laser pulses, at a peak intensity of  $5.5 \times 10^{20}$  W/cm<sup>2</sup>, with ultrathin carbon foils. Under optimized conditions, energies per nucleon of the bulk carbon ions reached significantly higher values than the energies of contaminant protons (33 MeV/nucleon vs 18 MeV), unlike what is typically observed in laser-foil acceleration experiments. Experimental data, and supporting simulations, emphasize different dominant acceleration mechanisms for the two ion species and highlight an (intensity dependent) optimum thickness for radiation pressure acceleration; it is suggested that the preceding laser energy reaching the target before the main pulse arrives plays a key role in a preferential acceleration of the heavier ion species.

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Laser-driven ion beams have gained significant scientific interest owing to their unique characteristics, such as ultrashort duration, high flux, and ultralow emittance [1,2]; this makes them a promising candidate for use in a range of applications [3,4]. While sheath acceleration mechanisms such as target normal sheath acceleration (TNSA) offer these properties, they suffer from a relatively slow scaling of ion energies with laser intensity limiting the applicative potential of these beams [5]. Furthermore, the dominant acceleration of protons typically shields and prevents acceleration of heavier ions [6], such as carbon, which also have significant applicative interest. Recently, mechanisms acting during the irradiation of ultrathin foils, such as radiation pressure acceleration (RPA), or acceleration in relativistically induced transparency (RIT) regimes have attracted significant attention [7,8]. A number of numerical studies have predicted that RPA will intrinsically dominate at intensities greater than  $10^{22}$  W/cm<sup>2</sup> with potential to produce monoenergetic beams of bulk target ions [9–15]; in particular, the “light sail” (LS) mode of RPA offers quadratic scaling of ion energies with laser fluence or a reduced target areal density, highlighting it as a promising route towards the production of high energy, high quality ion beams [8,11]. Numerical and experimental work has shown that this regime can be accessed at currently

available intensities by using a circularly polarized (CP) pulse to reduce the electron temperature and allow RPA features to emerge from (or dominate over) sheath acceleration effects [13,16–22].

In this Letter, we report on the experimental verification of an (intensity dependent) optimum target thickness for bulk ion acceleration via LS-RPA, as well as on the selective acceleration, at the optimum thickness, of bulk (carbon) ions to significantly higher energies than contaminant protons, whose energy is instead seen to drop substantially at this optimum thickness. This is the opposite of what is normally observed in multispecies ion acceleration from thin foils, where protons typically gain energy more efficiently during the acceleration process, thanks to the favorable  $Z/A$  ratio [16,17,23,24]. Multidimensional particle in cell (PIC) simulations, carried out for the conditions of the experiment, indicate that by exploiting the pulse’s unavoidable preceding intensity (the coherent contrast, usually a result of imperfect pulse compression [25]), proton contaminants can be displaced from the mass-limited target allowing the remaining carbon dominated target bulk to be accelerated by LS-RPA. The understanding provided in this Letter highlights routes for controllable, selective acceleration of a single ion species in the LS-RPA regime by tailoring the coherent contrast or

by using a separate prepulse. This approach will be an important step towards the development of versatile, high energy ion sources, and will facilitate a range of applications requiring single-species ion beams accelerated from the target bulk (an example being high dose-rate carbon radiobiology [26]).

The experiment was carried out using the GEMINI laser facility at the Rutherford Appleton Laboratory U.K. A 40 fs FWHM laser pulse, with central wavelength  $\lambda = 800$  nm was used, and temporally cleaned by a double plasma mirror arrangement [27,28] producing a contrast on the order of  $10^{-6}$  at 1 ps and  $10^{-14}$  at 100 ps which is sufficient to prevent significant premature expansion of the targets before the plasma mirrors are activated [29]. The pulse was focused using an  $f/2$  parabola at normal incidence producing a  $3 \mu\text{m}$  FWHM spot with  $\sim 6$  J on target and a peak intensity of  $5.5 \times 10^{20} \text{ W/cm}^2$ .

The main diagnostics were three Thomson parabola spectrometers (TPSs) placed at  $-4^\circ$ ,  $0^\circ$ , and  $9^\circ$  with respect to the target normal, each coupled to a microchannel plate (MCP)-CCD system to instantaneously record the ion spectra. Amorphous carbon targets (density of  $\sim 2 \text{ g/cm}^3$ ,  $n_e = 350 n_c$  where  $n_e$  and  $n_c$  are the electron density and plasma critical density, respectively) with thickness in the range 2–100 nm were irradiated to match the optimal LS parameters; an optimal areal density ( $\sigma$ ) equivalent to a target thickness of  $\sim 10$  nm at this electron density is required for the dimensionless laser amplitude  $a_0 \approx 13$  [11]. The laser polarization could be controlled with a  $\lambda/4$  wave plate placed in the collimated beam between the plasma mirrors and the parabola.

The maximum energies for the two main species in the spectra ( $\text{C}^{6+}$  and  $\text{H}^+$ ) as recorded on the  $0^\circ$  TPS for CP shots are shown in Fig. 1. The data highlights the presence of an optimal thickness at 15 nm, where the  $\text{C}^{6+}$  energies reach up to 33 MeV/nucleon ( $\sim 400$  MeV). A notable feature of the data is the fact that, under the conditions where the maximum carbon energies are obtained (15 nm and CP), a local minimum (18 MeV) is obtained for the proton energies. The intensity dependence of the maximum energy for the two ion species at the optimum thickness of 15 nm is shown in Fig. 2(a). While  $\text{C}^{6+}$  scales favorably with intensity ( $E_{\text{C}^{6+}} \propto I^{1.2 \pm 0.2}$ ),  $\text{H}^+$  follows a much slower trend. This behavior is not observed for linear polarization where there is no clear peak in carbon energies and consistently higher proton energies are observed for the same target thicknesses.

The preferential acceleration of the heavier species over protons as observed here is atypical of the acceleration mechanisms known to act on multispecies targets. In acceleration schemes such as TNSA, electrons set up a sheath field which will preferentially accelerate protons over other ion species due to their higher  $Z/A$  ratio, always resulting in higher energies for protons; this has also been

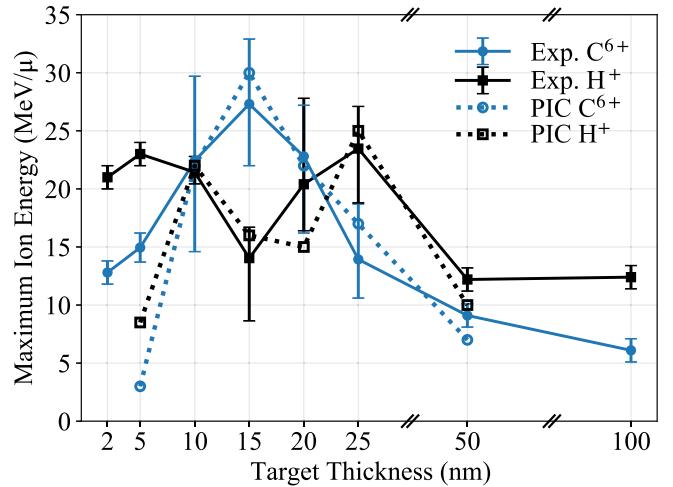


FIG. 1. Maximum ion energy from  $0^\circ$  TPS for  $\text{C}^{6+}$  (solid line and filled circles) and  $\text{H}^+$  (solid line and filled squares) for CP. 2D PIC simulations are dotted lines with the same markers but empty. Errors bars represent the fluctuations in the measured maximum energy due to shot to shot variation in laser energy for intensities greater than  $4 \times 10^{20} \text{ W/cm}^2$  throughout the experiment.

observed in the RIT regime [7,30,31] as well as in LS-RPA [13,16,17,21,23,32].

The observations at 15 nm therefore suggest that during the interaction for the optimized LS-RPA conditions, different acceleration processes act on the two species. In order to test this hypothesis, we have carried out supporting simulations with the EPOCH PIC code [33] for two cases: with and without the coherent contrast.

The first case considered a flat foil irradiated at normal incidence by an idealized 40 fs,  $3 \mu\text{m}$  pulse (both are Gaussian FWHM, with the simulation beginning at time  $t = -100$  fs so that no coherent contrast is considered). The target was composed of  $\text{C}^{6+}$  neutralized with electrons at a density of  $350 n_c$  with front and rear surface, 5 nm-thick  $10 n_c \text{ H}^+$  layers to simulate contaminant layers (where the  $\text{H}^+$  observed in experiments usually originate from). The simulations were performed on a grid with resolution  $\Delta x = 5 \text{ nm}$ ,  $\Delta y = 4\Delta x$ , and  $x$  range  $[-10 \mu\text{m}, 30 \mu\text{m}]$ . The target was initially located at  $x = 0$ , with  $y$  range  $[-5 \mu\text{m}, 5 \mu\text{m}]$  and was initialized with 200 particles per cell per species with a temperature of 10 keV. Collisions are not calculated in this case, as justified by their negligible influence at the temperature generated at this intensity. In these PIC simulations with an idealized pulse profile, the  $\text{H}^+$  energies scale similarly to carbon albeit they are  $\sim 30\%$  higher [triangles in Fig. 2(b)] and reproduce the  $I^2$  scaling associated with LS-RPA since an idealized target at this thickness (without considering any significant preexpansion) can enter the LS stage almost instantaneously and remain opaque for the majority of the pulse duration.

A better insight into the possible causes of the different behavior of protons and carbon ions (as observed in the

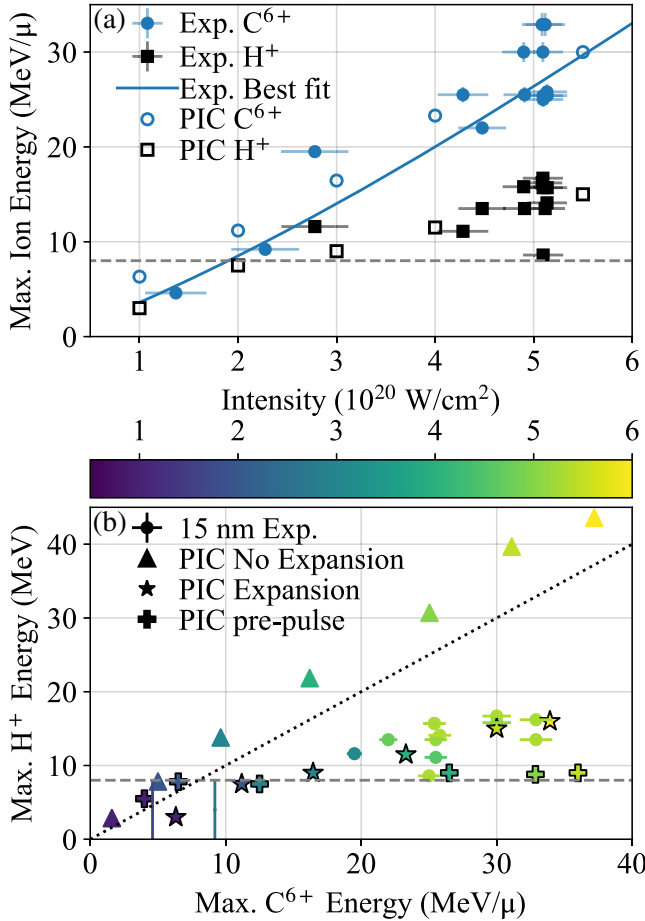


FIG. 2. (a) Maximum ion energy as a function of laser intensity for experimental measurements of C<sup>6+</sup> (filled blue circles) and H<sup>+</sup> (filled black squares). 2D PICs with expansion included are also shown with the same markers but empty. The solid line represents the best fit line to the C<sup>6+</sup> experimental data,  $\propto I^{1.2 \pm 0.2}$ . (b) Correlation between C<sup>6+</sup> and H<sup>+</sup> maximum energies for experimental and simulation data (marker colors indicate the laser intensity). Dotted line indicates where C<sup>6+</sup> and H<sup>+</sup> energies are equal. The dashed line at  $y = 8$  MeV in (a) and (b) represents the detection threshold for protons, and the maximum proton energy did not meet this threshold for the two lowest intensity shots. Crosses indicate simulations which included a prepulse with a peak intensity of  $10^{17}$  W/cm<sup>2</sup> arriving on target 2.5 ps before the peak of the main pulse.

experimental data of Figs. 1 and 2) can be obtained by introducing a more realistic temporal profile of the laser pulse into the simulations. The plasma mirrors, used to reduce target heating, are typically activated a few ps before the main pulse as the inherent contrast level decreases, meaning that the pulse is still preceded by a short pedestal and rising edge as shown in Fig. 3. This was considered in the second set of simulations. The intensity profile is based on a third order cross-correlator scan of the pulse before any contrast enhancement from the plasma mirrors (also shown in Fig. 3). Measurements on a similar system have

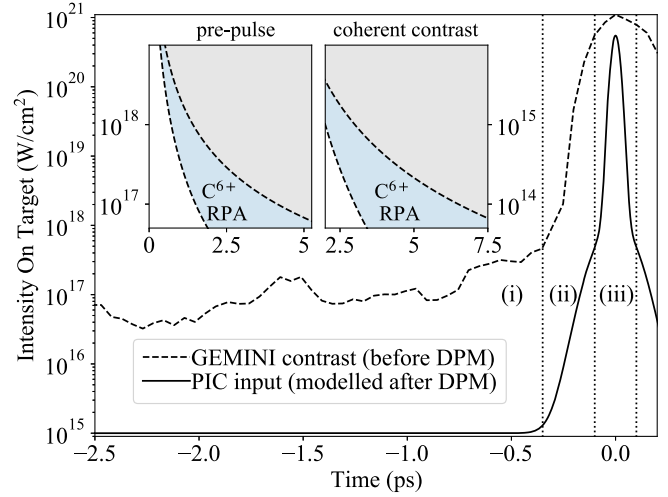


FIG. 3. Cross-correlator data of the GEMINI laser pulse profile without the double plasma mirror (DPM) arrangement in blue (dashed) alongside a modeled, contrast enhanced pulse profile used as a simulation input (solid black) and based on [34]. The pulse profiles have a peak intensity of  $5.5 \times 10^{20}$  W/cm<sup>2</sup> and  $1.1 \times 10^{21}$  W/cm<sup>2</sup> with and without plasma mirrors, respectively. Three annotated areas (divided by dotted black lines) mark the parts of the interaction highlighted in the text. The inset graphs show the parameter space for preferential species acceleration using a prepulse (left) and the coherent contrast (right) based on 2D PIC simulations. The intensities on the vertical axes indicate, respectively, the peak intensity of the prepulse (left frame) and the intensity of the flat pedestal (right frame), while the horizontal axes indicate the time separation between the prepulse and main pulse (left) and the start of the pedestal and the main pulse (right). The main pulse remains unchanged. The blue shaded areas indicate where preferential C<sup>6+</sup> acceleration by RPA is observed. The white area is where H<sup>+</sup> is favored by RPA and the grey area is where  $n_e < \gamma n_c$  and proton acceleration is favored in the RIT regime.

compared a similar cross-correlator scan with a much higher resolution technique both with and without a plasma mirror [34]. Based on this work, a truer representation of the experimental pulse profile is modeled in our simulations. The same grid is initialized in these simulations except now with ionization [35–37] and collisions [38] calculated and the target set as unionized and cold. Additional simulations with different contrast levels are shown in the Supplemental Material [39].

At the optimum thickness, the interaction can be thought of as consisting of three stages: (i) expansion, (ii) re-compression, and (iii) acceleration. The low intensity pedestal is responsible for the majority of the ionization and expansion. At this stage protons, with their lighter mass, will expand faster than the carbon ions and are mainly located in the front and rear underdense plasma. The target bulk has also expanded but it remains overdense. Furthermore, the expansion of the target inevitably means the mass of the target bulk partaking in the



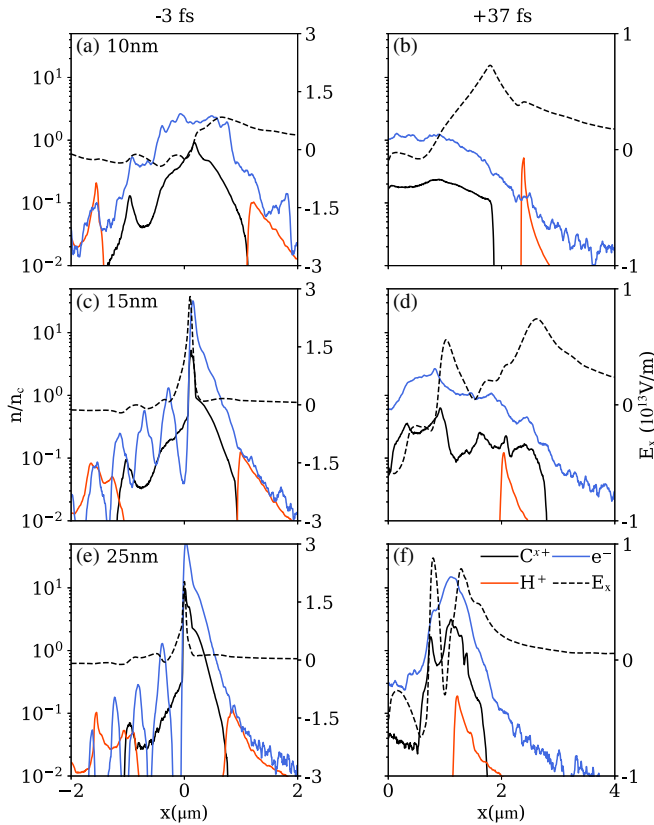


FIG. 4. Average density (left logarithmic axes) for  $\text{C}^{6+}$  (solid black), protons (red), and electrons (blue) for 10 nm [(a),(b)], 15 nm [(c),(d)], and 25 nm [(e),(f)] along the laser axis (averaged between  $y = \pm 100$  nm). The right axes display the cycle-averaged longitudinal electric field ( $E_x$ , dashed black line and linear axis). The left panels capture the interaction at close to the peak of the pulse ( $-3$  fs) while the panels on the right refer to a time towards the end of the pulse ( $+37$  fs).

subsequent LS-RPA stage is reduced explaining why the experimentally observed optimal thickness is higher than that predicted theoretically [11].

Less than 1 ps before the main pulse arrives, the intensity rises to the point where the thermal pressure in the target is overcome by the radiation pressure. The laser reflects off the overdense plasma, setting up a standing wave with the incoming radiation, which increases the local electric field, and pushes electrons to the points of lower laser intensity [corresponding to the density modulations in Figs. 4(c) and 4(e) for  $x < 0$ ]. The laser's rising edge begins to compress the front surface which further steepens the front density gradient, reversing some of the density decrease caused by the earlier expansion; this affects only the overdense carbon plasma, as the protons are not present in significant number in this central region. When the pulse peak arrives, it interacts with a recompressed (but reduced density) plasma which is close to the optimal areal density for the incident intensity (see Supplemental Material [39]). Here the laser accelerates the carbon bulk by LS-RPA whereas proton

energies will be determined by subsequent sheath acceleration, as shown in Figs. 4(c) and 4(d); as a result of these dynamics, the high energy carbon ions from the target bulk are accelerated much more efficiently than protons.

The LS-RPA phase is followed by an RIT phase which also has a role in accelerating ions. Close to the peak of the pulse, the plasma becomes relativistically transparent. Comparisons of simulations with and without this laser energy propagating through the target show that the RIT regime accounts for  $\sim 37\%$  of the final  $\text{C}^{6+}$  energy and  $\sim 63\%$  of the final proton energy.

The peak intensity was also varied in these simulations reproducing the experimentally observed species-dependent scaling of the cut-off energy with laser intensity (Fig. 2). The simulations show that a reduced intensity scaling for carbon compared to the ideal  $I^2$  dependence arises from the combination of two effects: (1) a delayed transition to LS since the initial hole boring phase must penetrate through the now expanded target; (2) an earlier transition to transparency for more intense pulses.

Achieving preferential carbon acceleration requires sufficient relative expansion of the proton species, such that protons are present only in the underdense preplasma, all while the target bulk remains relativistically overdense. As shown in the inset graphs of Fig. 3, these conditions can be accessed through the target heating induced by the coherent contrast reflected from the plasma mirrors, as in the experiment, but could also be reached by introducing a prepulse, of equal duration to the main pulse, a few ps before the pulse peak (representative of a pickoff from the main beam). The left inset panel in Fig. 3, summarizes results of 2D PIC simulations in which a 40 fs pulse, of varying peak intensity ( $y$  axis), irradiates the target at a time ( $x$  axis) before the main beam arrives, which allows the target to expand. The parameter region where preferential carbon acceleration was observed in simulations is shown in blue. Data points from a simulation employing a  $10^{17}$  W/cm<sup>2</sup> prepulse 2.5 ps before the main pulse are also shown in Fig. 2(b) as an example. Additional PIC simulations were performed to explore the effect of the intensity and duration of the coherent contrast on target (see right inset panel in Fig. 3). This was done by varying the amplitude of the flat pedestal [section (i) of the pulse profile in Fig. 3] and varying the arrival time of the peak ( $x$  axis of the right-hand inset in Fig. 3) showing that a window of 3–8 ps for  $10^{14}$  W/cm<sup>2</sup> as opposed to 2–2.5 ps for  $10^{15}$  W/cm<sup>2</sup> is required for preferential carbon acceleration. In general terms, if the laser contrast is better then one would need to activate the plasma mirrors earlier (by using a tighter focus on their surface) to achieve the same degree of target expansion. Both of these methods highlight routes to control the preferential acceleration of the carbon bulk.

The simulations also provide an explanation for the local minimum in the  $\text{H}^+$  energies and a comparison between

experiments and simulations is shown in Fig. 1. At the optimal thickness, the carbon bulk is accelerated much more efficiently by LS-RPA which, in turn, reduces the sheath potential at the target rear where the protons are located. The carbon ions will overtake the proton contaminants and are themselves further accelerated in the sheath [Fig. 4(d)].

For targets thicker than the optimum, the energy gain for the  $C^{6+}$  bulk is significantly reduced due to the increased areal density. While the transition to transparency is delayed, this does not overcome the reduced efficiency of the acceleration ( $E_{\max} \propto \sigma^{-2}$ ) [8]. Here, sheath acceleration dominates and favors protons since the carbon bulk remains behind the proton contaminants. For targets thinner than the LS optimum, by the time the main pulse arrives, the target has expanded so much that it is relativistically transparent (where  $n_e < \gamma n_c$ , where  $\gamma$  is the Lorentz factor, as seen in Fig. 4(a)). The  $C^{6+}$  ions can gain energy very quickly in an initial RPA phase but this is short lived as the target is too thin to sustain LS-RPA at this intensity and RIT sets in well before the peak of the pulse. This reduces the energy gain for carbons, but causes direct electron heating as the pulse propagates through the transparent plasma [Figs. 4(a) and 4(b)]. Since protons have expanded at the rear, they will experience the strongest field from this heating effect, screening the  $C^{6+}$  ions, and are subsequently accelerated to higher energies [31].

In conclusion, we have demonstrated, through experiments and simulations, that a moderate preexpansion of ultrathin targets can lead, via LS-RPA, to dominant acceleration of bulk carbon ions over contaminant protons. Furthermore, we present a more general case for controlling the expansion of a multispecies, ultrathin foil to selectively accelerate a single species. This can be achieved, in principle, by either modifying the activation time of the plasma mirror or by using a separate ps-timescale prepulse to controllably preheat the target as indicated by the simulations presented in this Letter. Additional opportunities for control may be achieved by modifying the pulse profile on a sub-ps timescale by tuning the third order dispersion [40] to further tailor the target density profiles prior to the arrival of the peak of the pulse. We believe that these considerations will be highly relevant and applicable to future investigations on multi-PW systems, where RPA acceleration will be a major topic of investigations. Our findings point to a promising route towards a versatile source of high energy ions with control over the preferential acceleration of a single species.

Data is available from [41].

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