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Pathways control in modification of solid surfaces induced by temporarily separated femtosecond laser pulses

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Reaction control by laser light is a unique method of the reaction dynamics mastery in the molecular chemistry. We provide evidence of phase control processing with femtosecond lasers in macroscopic solids. Rutile TiO₂ monocrystals with (001) and (100) surface orientations were irradiated with repetitive pulses of femtosecond KrF laser of variable fluences and a temporal delay between two superimposed linearly polarized beams. The appearance of three types of surface morphology was thoroughly analyzed: low-spatial frequency laser-induced periodic surface structures (LSFL), grooves and unusual featureless flat area (FFA). The interaction of light with the excited surface led to the onset or suppression of the subwavelength LSFL, depending on whether the temporal delay between laser beams is larger or smaller than the critical value of ~6 ps. By contrast, the suprawavelength grooves and FFA appeared at longer temporal delays. A strong decrease of the grooves onset energy was observed on the (001) oriented crystal after the delay of ~8 ps; the decrease of onset energy was also observed on the (100) oriented crystal, where FFA appeared instead of grooves. The critical delay is discussed in framework of a phenomenological model describing the energy evolution of excited transient states along the “reaction” coordinate.

Keywords: LIPSS; Electronic excitation; TiO₂; Femtosecond UV-laser irradiation.

1. Introduction

Femtochemistry is a rapidly developing topic of science since the 1990's, which became an effective method for realisation of the reaction dynamics control of single molecules in a gas phase, molecular beams and in more complex biological systems [1, 2]. In the same time, the control of modifications in solids faces severe restrictions because of the generally much shorter coherence time [3]. Recently, an ultra-high precision sub-femtosecond phase monitoring has been employed to validate the approach in condensed matter in several examples of delocalised vibration in solid para-hydrogen, collective phonon motion in bismuth single crystal [4] and electron transfer in CdSe quantum dots [5]. These systems are characterised by a relatively longer picoseconds vibrational dephasing time facilitating the analysis. The very exciting example of the metal-insulator structural phase transition in atomic indium wires on (111) silicon surface via specific phonon modes has been demonstrated only recently [6]. These studies provide some evidence that the interaction can be sensitive to a temporal delay at the femtosecond timescale between two laser beams superimposed on a sample, which affects greatly the quantum states of the system.

The control of the reaction pathways is connected to the transient state evolution along the reaction coordinate, which can lead to the specific bonds breaking, modifying reaction selectivity and final product. This kind of interaction, theoretically considered by Tannor et al. [7], can guide chemical transformations on the ground state potential energy surface of molecular species by intermediary excited states. In the same

time, there can be not pronounced similarity between the molecular structure and surface morphology of the crystalline solids [8]. The femtosecond excitation can involve relatively short-lived transient states spanning over the macroscopic volume and creating conditions for the phase control of the energy relaxation process directing the system to a particular surface morphology via the relatively slow melting and re-solidification stages taking place on a longer nanosecond time-scale. In such a case, the reaction pathways might be controlled by the laser fluence, instead of the photon energy assisting molecular transformations in a single molecule. Therefore, the laser fluence and temporal modulation of the absorbed energy are expected to control specific transient states and final reaction state via the density of electronic excitation of solids.

In most representative examples, a weak perturbation in solids was considered. By contrast, in conditions of the strong perturbation, the high density electronic excitation (HDEE) regime can alter the quantum state of a macroscopic crystal. In general, this requires the irradiation of solid surfaces with femtosecond pulses of duration shorter than the electron-phonon relaxation time, of photon energy above the band gap energy of the solid and of fluence above that necessary for producing Mott charge density, which marks the semiconductor-metal transition [9]. Such levels of excitation, when promotes ~10% electrons to the conduction band, can instantaneously trigger the displacement of cations from their equilibrium positions of the cooled lattice, stabilized by the non-perturbed electron clouds [10-13]. The characteristic time of this reorganisation can vary in the

range between 10^{-12} and 10^{-11} seconds. Acting on the electronic band structure during this time, may promote the overall macroscopic system to a new transition state with an unusual lattice dynamics, which may guide to a particular final morphology of the irradiated surface.

There is an abundance of experimental and theoretical studies about surface modifications of solids with the ultrashort-pulse laser irradiation. By varying the irradiation conditions of laser fluence (F), number of pulses (N), temporal delay between beams (Δt), etc. different laser-induced periodic surface structures (LIPSS) have been obtained and their formation mechanisms have been discussed [14-24]. It has been proposed that very short laser pulses can drive the irradiated matter into a transient “hyperthermal” state far from thermodynamic equilibrium, which is responsible for LIPSS formation [25]. The observation of coherence decay of this intermediate “soft” state with different electronic properties has been suggested for fs-laser ablation by Reif et al. [26] after first indications of an unusual “surface liquid” characterized by an increased absorption of laser photons [27]. An exciting interplay between suppression/amplification of the ablation effectiveness and surface structuring with the delay between laser pulses of the order of several picoseconds has been observed in silicon [28]. Recently, we have analysed the HDEE regime in rutile TiO_2 single crystal with two different (100) and (001) surface orientations and reported on their morphological surface maps (MSM) [29], which describe the domains of the appearance of nanoparticle-covered structures (NCS), low-spatial frequency LIPSS (LSFL), grooves and featureless flat area (FFA). The two types of structures, grooves and FFA, respectively appeared on the (001) and (100) oriented surfaces, have been explained by the crystalline anisotropy. The surface patterns in the experimental conditions of HDEE appeared to be connected to the density of electronic excitation (charge carriers density) [30].

In this communication, we analyzed the influence of the temporal delay between two intense femtosecond laser beams impinging the semiconductor surface on MSM and compare energy thresholds of the structures appearance. Rutile TiO_2 single crystals with (001) and (100) orientations of the crystalline planes with respect to the surface normal were irradiated. The laser photon energy (5.0 eV) was above the band gap energy of rutile TiO_2 ($E_g=3.0$ eV [31]) that enables one-photon transitions between the valence and conduction band and, consequently, permits obtaining the highest instantaneous density of the electron-hole plasma (EHP).

2. Experiment

The experiments were performed with rutile TiO_2 single crystals with (001) and (100) oriented polished surfaces, supplied by *Crystal-GmbH*, which were shaped in pallets of size $10 \times 10 \times 0.5$ mm³. The laser irradiation of crystals was carried out at normal incidence, ambient atmosphere and at room temperature at the Ultraviolet Laser Facility of IESL-FORTH (Heraklion, Greece) with a hybrid distributed feedback dye laser / KrF excimer system delivering linear polarized pulses of 10 mJ energy/pulse, $\tau=450$ fs (full width at half maximum) duration at $\lambda_L=248$ nm ($h\nu_L=5.0$ eV) and repetition rate 1-10 Hz, which assures effective dissipation of heat between laser pulses and prevents changes of sample temperature during repetitive illumination up to high integral doses, attained with a sequence of $10 \leq N \leq 10^3$ laser pulses. It is noted that the term ‘dose’ is related to the amount of

energy density deposited from the laser on the solid at each single or double pulse.

In pump-pump experiments, the laser beam was split in two beams, as depicted in Fig. 1a, and the temporal delay up to 420 ps (with the precision of 0.5 ps) was introduced with the Michelson interferometer. The “zero” delay correspondent to the temporal coincidence of the two beams was controlled by the appearance of interference fringes. The energy on the sample was varied by using two reflective attenuators and measured using an energy meter Moletron J04-09 with sensitivity 1.42 V/mJ. In two-beam irradiation experiment, the energies of the beams were E1 and E2 and the total energy sent to the sample (by one double pulse) was $E=E_1+E_2$. In one-beam irradiation experiment with the same energy E, the delay line (E1) was closed and the beam energy E2 was adjusted to that of the sum of the two beams $E=E_1+E_2$. The laser intensity distribution on sample, after the spatial filtering using a 200 mm focal distance lens, shown in Fig. 1b, was measured with a camera CM-140 GE-UV from JAI. This picture evidenced a quasi-perfect spatial superposition of the two split and delayed beams within the precision of one pixel (<5 μm). The beam-shape measurements, permanently conducted in all experimental series, evidenced a quasi-Gaussian spatial intensity profile on the sample, as shown in the x-axis and y-axis distributions in Fig. 1c. Accordingly, the irradiated spot had a characteristic diameter of ~ 30 μm (measured at $1/e$ of the maximum intensity), which corresponded to laser fluences on the samples varied between 50 and $2 \cdot 10^3$ mJ/cm².

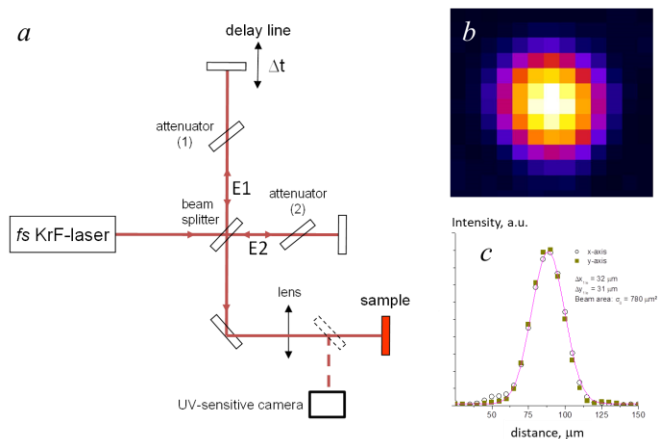


Fig. 1. Experimental schema (a), laser beams superposition (b) and spatial beam shape (c). The size of one pixel is 5×5 μm^2 (a).

The surface morphology after the irradiation was examined by a high-resolution Jeol Scanning Microscope JSM 7000F with a Schottky type field-emission gun (T-FE) operating in a vacuum of 10^{-8} Pa. The spatial resolution was 1.2 nm at 30 kV of the maximum accelerating voltage.

3. Results and discussion

SEM images of the (001) crystalline surface after one-beam irradiation with $N=30$, shown in Fig. 2, correspond to the characteristic MSM domains (F, N) described in Ref. [29]. At fluences above a particular value of the threshold $F^*=90$ mJ/cm², LSFL (ripples) were formed as shows image in Fig. 2b. With the

following increase of the fluence, 1D pattern transforms to 2D pattern (Fig. 2c). At very high fluences above 0.6 J/cm^2 ($N=100$), the suprawavelength structure that is termed grooves nucleates and becomes well-developed at $F \sim 1 \text{ J/cm}^2$. By contrast, at relatively low fluences below F^* , NCS (a) were observed (Fig. 2a). The onset of the ripples formation corresponds to a small material sublimation in the irradiated area and is in agreement with the theoretical model of Sipe et al. [32]. The (100) TiO_2 surface manifested a similar sequence of the SEM images, except for grooves, which were replaced by FFA.

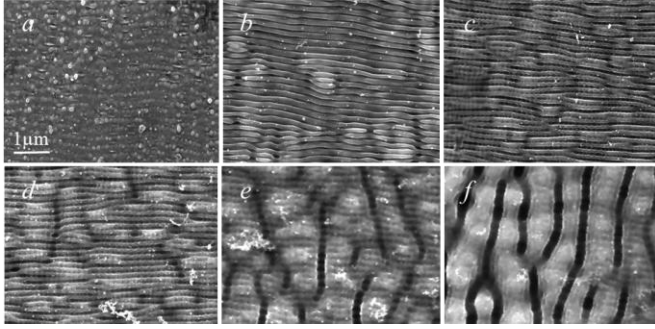


Fig. 2. Surface structuring of rutile TiO_2 (001) monocrystal after one-beam irradiation with KrF laser (248 nm, 450 fs, $N=100$): fluences are 75 (a), 150 (b), 230 (c), 440 (d), 640 (e) and 960 (f) mJ/cm^2 .

An exposition of the semiconductor surface with repetitive pairs of laser pulses with a temporal delay Δt between two beams in each pair can significantly modify MSM. A series of SEM images that correspond to the LSFL domain of MSM [29], shown in Fig. 3, were obtained with the temporal delay of $\Delta t=10 \text{ ps}$. Fig. 3a-b and 3d-e evidence the initiation of ripples at the periphery of the irradiation spot at (001) and (100) surfaces of rutile TiO_2 , which is similar to the one-beam experiment (Fig. 2). However, at a longer irradiation with $N=100$ the central part of the spot was covered with grooves on (001) surface, while FFA appeared on (100) surface, as respectively show images of Fig. 3c and 3f. The appearance of grooves and FFA in the low-energy domain of MSM ($F=F_1+F_2=300 \text{ mJ/cm}^2$) was totally unexpected, since these structures are commonly observed at much higher fluencies.

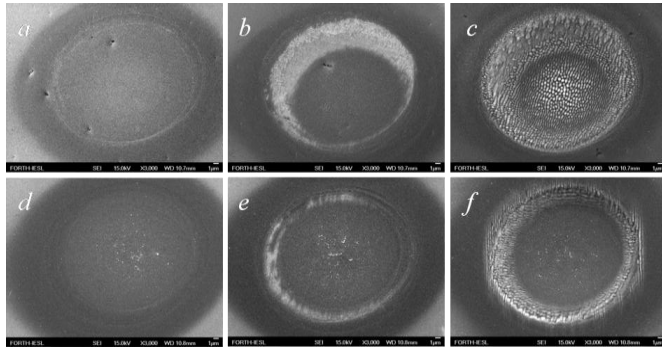


Fig. 3. Surface structuring of rutile TiO_2 (001) (a-c) and (100) (d-f) with repetitive pairs of pulses $N=5$ (a, d), 20 (b, e), 100 (c, f) of KrF laser (248 nm, 450 fs, two-beam with fluences $F_1=100 \text{ mJ/cm}^2$ and $F_2=200 \text{ mJ/cm}^2$ delayed by 10 ps).

The difference between the groove formation in one-beam and two-beam experiments is further explored with results shown in Fig. 4. In agreement with the commonly accepted mechanism, in one-beam experiments grooves were nucleated at the top of

ripples as the number of laser pulses N increased (Fig. 4a-c). In contrast, SEM images in Fig. 4d-e indicate nucleation of grooves in the central area of the spot in absence of any ripples. At large N , however, SEM pictures of crystals in both experiments looked quite similar (Fig. 4 c and f).

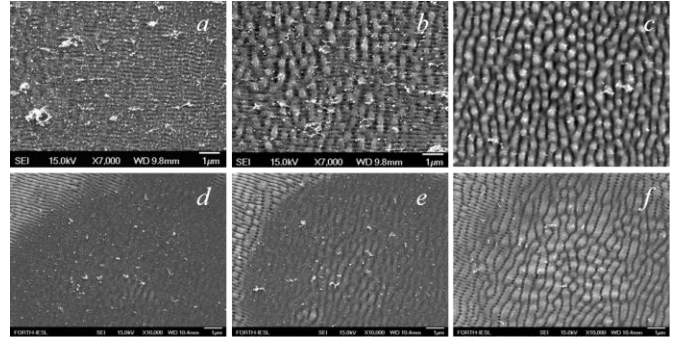


Fig. 4. Surface structuring of rutile TiO_2 (001) with KrF laser (248 nm, 450 fs) after one-beam irradiation with $F=1.1 \text{ J/cm}^2$ and $N=5$ (a), 10 (b) and 30 (c) and after two-beam irradiation with $F_1=100 \text{ mJ/cm}^2$ and $F_2=200 \text{ mJ/cm}^2$ delayed by 10 ps and number of pairs of pulses: $N=30$ (d), 50 (e), 100 (f).

At low fluences but above the threshold of LSFL formation ($F_{1,2} < F^*$ and $F_1+F_2 \geq F^*$) and sufficiently long irradiation ($N=100$), the surface area covered with ripples S_{ripples} progressively decreased with an increase of the delay time between two laser beams. The dependence of the normalized ratio $S_{\text{ripples}}/S_{\text{beam}}$ on the delay is presented in Fig. 5 for two differently oriented TiO_2 crystals. As these results show, the delay time of the ripples disappearance τ_r increased with an increase of the pulse energy, which is in agreement with the hypothesis about energy depletion of the common transient excited state. A critical energy delivered to this state (correspondent to the threshold laser fluence $F^*=90 \text{ mJ/cm}^2$) induced the surface morphology evolution towards ripples. The characteristic critical delay time of LSFL disappearance on (100) and (001) oriented surfaces is almost similar and fits the range of $\tau_r=6 \pm 1 \text{ ps}$. The overall picture apparently depends on the total energy density absorbed by the crystal to a common transient state, since the disproportioning of the two beams energies did not modify the critical delay time.

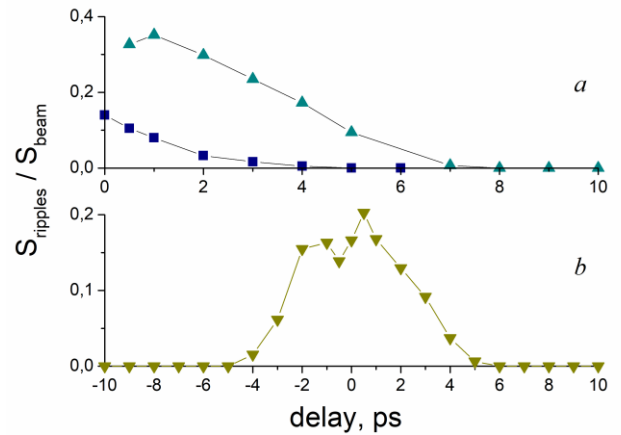


Fig. 5. Relative LSFL-covered area of rutile TiO_2 (001) (a) and (100) (b) with two-beam of equal fluences $E_1=E_2=50 \text{ mJ/cm}^2$ (\blacksquare), 75 mJ/cm^2 (\blacktriangle) and two beams of different fluences $E_1=50 \text{ mJ/cm}^2$ and $E_2=70 \text{ mJ/cm}^2$ (\blacktriangledown) (KrF laser: 248 nm, 450 fs, $N=100$). S_{beam} corresponds to the beam area with intensity above $1/e$ of the maximum.

In two-beam experiments with $F_{1,2} \geq F^*$, each laser pulse separately triggered the formation of LSF structures (ripples). Furthermore, two more types of surface modification could be observed in this case: grooves on (001) and FFA on (100) oriented rutile TiO_2 crystals, which dependence of $S_{\text{structure}}/S_{\text{beam}}$ on the temporal delay is shown in Fig. 6 for the two crystalline orientations. On the (001) oriented surface, ripples were replaced by grooves: they appeared in the central part of the irradiated spot after a critical delay time τ_g , which shortened from 7 ps to 3 ps with an increase of laser fluence and progressively spanned towards the spot periphery erasing ripples. This supports the previous conclusion of Fig. 4 about different mechanisms of their formation in one-beam and two-beam experiments. We also notice an asymmetry of the structuring area dependence on the delay time, in experiments with disproportioning energies of the two beams (Fig. 6). This point will be discussed below in section concerning possible mechanisms of the structuring.

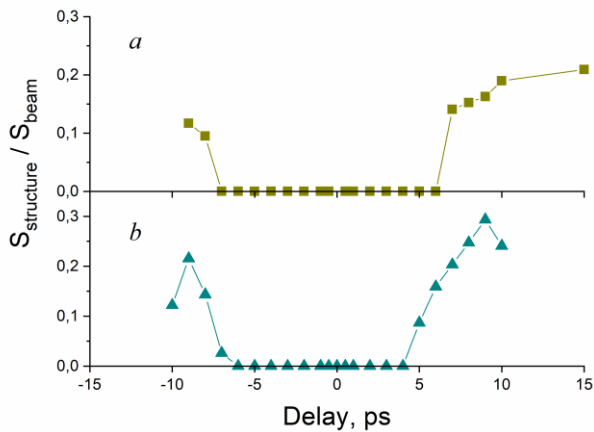


Fig. 6. Relative area of surface structuring of rutile TiO_2 (001) (a) and (100) (b) in pump-pump experiments with $F_1=120 \text{ mJ/cm}^2$ and $F_2=240 \text{ mJ/cm}^2$ versus delay between pulses (KrF laser: 248 nm, 450 fs, $N=100$): grooves (■) and FFA (▲) areas. S_{beam} corresponds to the beam area with intensity above $1/e$ of the maximum.

The precision of determination of the grooves area on the (100) surface was problematic because of a difficulty with its distinguishing from the area covered with short elongated precursor relief particles, which appeared at the grooves nucleation stage. In contrast, because the net border between structured and flat areas (see tilt SEM images in Fig. 7), (100) oriented crystal enabled precise measurements of the structuring area evolution with temporal delay. The data analysis showed that FFA appeared in the same range of laser fluencies and time delays, as grooves on the (001) surface. Furthermore, SEM images indicated that FFA replaced ripples in the central part of the irradiated spot and, similarly to grooves on the (001) oriented surface, was formed after the critical time delay of $\Delta t \approx \tau_g$ (Fig. 6) between two laser beams. Based on these observations, a general conclusion can be drawn about similar mechanisms responsible for the FFA and grooves formation. The critical delay time τ_g is apparently due to a transient state (different from that leading to the formation of ripples) excited by the first laser beam, which feeds with energy from the delayed beam.

Complementary to these measurements, Fig. 8 evidences that the critical delay required for the grooves / FFA formation shortens with an increase of the pulse energy. Indeed, fluence of

350 mJ/cm^2 set up the supra-wavelength structuring at $\Delta t_1=8 \text{ ps}$ that disappeared at $\Delta t_2=40 \text{ ps}$ by forming submicronic spike-like columns. When increasing fluence up to 600 mJ/cm^2 , the time-window of the grooves onset shortened and shifted to small delays of $3 \leq \Delta t \leq 10 \text{ ps}$. This indicates modifications of the transient energetic structure with time.

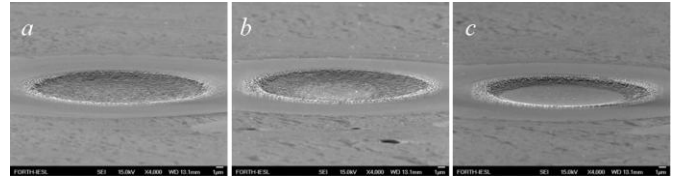


Fig. 7. Tilt SEM images of rutile TiO_2 (100) monocystal irradiated with two laser pulses (248 nm, 450 fs, $N=100$) of fluences $F_1=120 \text{ mJ/cm}^2$ and $F_2=240 \text{ mJ/cm}^2$ delayed by $\Delta t=6 \text{ ps}$ (a), 7 ps (b) and 15 ps (c).

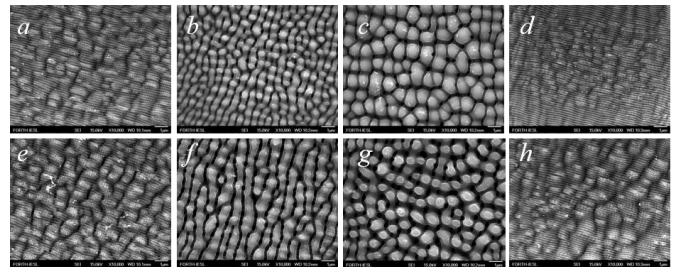


Fig. 8. Surface structuring of rutile TiO_2 (001) monocystal with two laser beams (248 nm, 450 fs, $N=100$) of total energy 350 mJ/cm^2 (a-d) and 600 mJ/cm^2 (e-h) delayed by $\Delta t=0$ (a, e), 8 ps (b), 40 ps (c), 3 ps (f), 10 ps (g) and 0.1 s (d, h).

The phenomenological schemas of the relevant structuring processes in the pump-pump experiments are depicted in Fig. 9. With respect to the processes which take place and contribute to the topography variation at increasing integral dose (i.e. increase of pulse number), it is noted that hydrodynamical factors play an important role along with electro-dynamical effects; more specifically, the latter determine to a great extent the orientation and periodicities of LIPSS while the former lead to the formation of a particular surface topography through transport of a fluid that has been produced following relaxation processes [22]. Fluid transport and resolidification of the material determine the surface morphological features that are produced within a few nanoseconds. The timescale of these processes has also been confirmed in previous reports [20, 33-37]. In contrast, the precursor state is activated during the ultrashort double-beam laser irradiation. This precursor state is characterized by the surface energy repartition and survives a time, during which it can be modified by an interaction with the second pulse. We could attribute an energetic structure to this state in terms of X_1 , X_2 etc. potentials. The mechanism of X_1 state excitation resulting in LSF consists in the absorption of 1st beam energy, which produces EHP continuum with the under-critical charges number density, which evolves along the reaction coordinate Q to result in the final NCS morphology. Furthermore, absorption of the 2nd beam energy after the temporal delay Δt , which defines a partial depletion $\Delta E_{\text{relaxation}}$ of the initially deposited energy density, can promote the system to X_1 state, which trigs LSF structuring. Considering fluences of both laser beams F_1 , F_2 to be smaller than energy density of the ripples onset $E_{g, \text{ripples}}$, the critical time delay $\Delta t = \tau_c$, after which ripples cannot be formed, can be defined by

$$F_1 - \Delta E_{\text{relaxation}}(\tau_r) + F_2 = E_{a,\text{ripples}} \quad (1)$$

The mechanism described by Eq. (1) is schematically depicted in Fig. 9a. In experiments with $F_{1,2} > F^*$, inhomogeneities of the surface energy distribution produced by the 1st laser beam can be maintained by the 2nd beam under a condition of the process in-phase contribution. The characteristic process deactivation time can be estimated as $\tau^* = \tau_r \approx 6$ ps at laser fluences close to the threshold F^* . Beyond this time, the additivity of two contributions is seemingly cancelled suppressing ripples formation, as evidences Fig. 4d. Even when the delayed beam by $\Delta t > \tau^*$ is capable producing ripples, its interaction with surface state prepared by the leading beam cancels the formation. Thus, two processes of the LSFL nucleation annihilate in these conditions.

The case of the grooves formation in the conditions of $F_{1,2} > F^*$ is more complex. On the one hand, grooves are generally considered to be formed on the top of ripples, which is confirmed by our measurements with one-beam irradiation (equivalent to zero delay $\Delta t = 0$). On the other hand, in two-beams irradiation experiments ($\Delta t \neq 0$), the grooves appeared on a flat surface, in absence of any sign of ripples, as evidence SEM images in Fig. 4 (e-f). Interestingly, grooves begin to appear at the time $\Delta t > \tau^*$, when LSFL vanishes. Since the formations of ripples and grooves at zero delay are connected (Fig. 4a-c), transient state X1 activated with F_1 on the X1 \leftarrow X0 transition participates in excitation of a higher lying transient state X2, which results in the grooves formation (Fig. 9b). We assume that the activation energy density $E_{a,\text{grooves}}$ of X2 decreased along the reaction coordinate Q of X1 towards the final surface morphology. If this decrease is stronger than energy losses of X1 state, X2 \leftarrow X0 transition becomes accessible with fluence F_2 at the critical time $\Delta t = \tau_g$ after the first beam absorption. Consequently, an increase of the laser energy would result in a shortening of the critical temporal delay between two laser beams permitting to reach X2 state, which is in agreement with Fig. 8 (b, f).

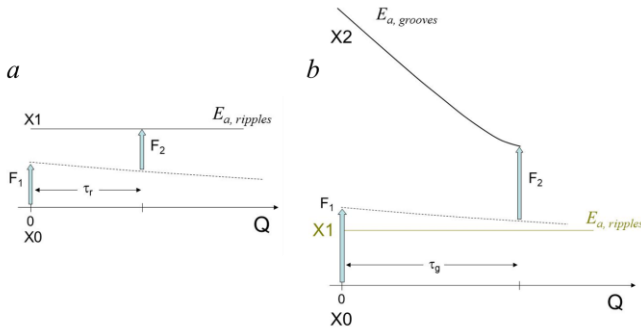


Fig. 9. Phenomenological schemas of structuring process in pump-pump experiments. F_1 and F_2 are laser fluences of two spatially superposed beams and temporarily delayed by τ , Q is reaction coordinate of the solid modification towards formation of NCS (a) and LSFL (b).

The asymmetry of FFA and grooves areas with respect to the temporal delay (positive $\Delta t > 0$ or negative $\Delta t < 0$) between two laser beams with different energies $E_1 \neq E_2$ supports the proposed schema in Fig. 9b. In our pump-pump experiments with disproportioning energies of two laser beams, the reference beam was that with a weaker fluence F_1 , while the stronger one with fluence $F_2 > F_1$ was sent before ($\Delta t < 0$) and after ($\Delta t > 0$) the reference beam. Because of the excited state energy depletion, the electronic energy density stored in the final state is $E(\tau) = E_i \cdot e^{-k|\Delta t|} + E_j$ ($i, j = 1$ or 2) and is stronger when $i=1$ and $j=2$, which

facilitates attaining the X2 state potential. This is in agreement with the experimental results reported in Figs 6a (grooves) and 6b (FFA), where critical time of the structures appearance is shorter for the positive delays $\Delta t > 0$, as compared with $\Delta t < 0$.

The proposed schema of the surface structures activation by lasers fits that proposed for controlling the reaction selectivity via the propagation time on the excited state potential energy surface, by applying delayed femtosecond pulses [7]. The case of grooves / FFA morphologies can be treated by considering the total energy evolution in the crystal. The interaction of light with an excited surface at intermediate reaction times can cancel LSFL (ripples), while maintain them at zero and infinite delays after the process beginning. Furthermore, an appropriately chosen delay time between femtosecond pulses permits to significantly decrease energy required for the surface modification: in case of FFA and grooves the energetic efficiency of the structuring process was increased by a factor of 3.

A theoretical study of physical processes triggered by HDEE regime in rutile TiO_2 crystals with (100) and (001) surface orientations has been recently reported [30]. This work has shown a strong enhancement of the absorption of the second laser pulse (delayed by Δt) in the lattice modified with the first pulse, creating necessary conditions for a significant weakening of onset energies for the grooves and FFA formation. In framework of the proposed schema, these energy states characterize the structure of superheated material and transient liquid phase, which remains largely unknown. These material states depend on the electronic excitation and lattice temperature, and the energy is a complex function of the charge carrier density N_e and interatomic distances r , which depend on the lattice temperature T_L of the solid. The relevant transient states are by consequence complex surfaces of energy density $E(r, N_e)$. A view of the electronic transitions in a frame $[E, r, N_e]$ leading to the formation of grooves is given in Fig. 10.

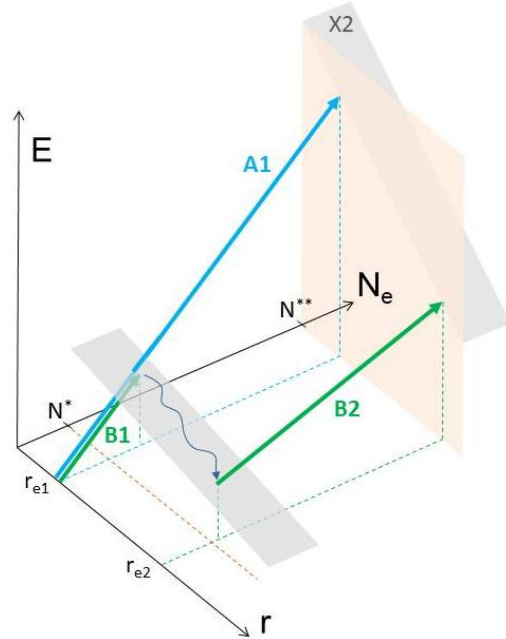


Fig. 10. View of electronic transitions in frame (E, r, N_e) leading to formation of grooves (see explanation in text). Laser fluence in one-beam is F_{A1} and laser fluences in two-beams are F_{B1} and F_{B2} . N_e is charge carrier density and r is characteristic interatomic distance.

The critical charge density N^{**} for the formation of grooves can be attained by applying one laser pulse A1 to TiO_2 ($T=300\text{K}$, $r=r_{e1}$): this corresponds to the vertical electronic transition in (E, r) plane and inclined transition in (E, N_e) plane. On the other hand, by applying first laser beam B1 ($T_1=300\text{K}$, $r=r_{e1}$) and second delayed B2 (after solid heating $T_1 \gg 300\text{K}$: $r_{e1} \rightarrow r_{e2}$) N^{**} can be attained at a smaller energy cost: $F_{B1}+F_{B2} < F_{A1}$. In this regard, a contribution of the electronically melt state to the critical delay τ_r may take place, as it has been previously reported in GaAs [11]. On longer times, the hypothetical state X2 can result in thermal gradients within the fluid melt characterized by large Prandtl numbers [37].

Despite of the tentative explanation has been given in Ref. [30], the qualitative agreement between theory and experiment using delayed laser beams is not yet attained. In particular, relatively narrow delay time-window for the initiation of supra-wavelength periodic surface structures and amplification/annihilation of LSFL with the temporal delay for $E_{1,2} > E^*$ remain unexplored. In these conditions, theoretical picture of the transient electronic states evolution on the picosecond timescale needs a better understanding.

4. Conclusions

In this study, we report on the application of delayed femtosecond laser beams to control the final surface morphology of solids in the regime of high density electronic excitation. Single rutile TiO_2 crystals with (001) and (100) surface orientations were irradiated with repetitive pulses of a femtosecond KrF laser (450 fs, 248 nm), which photon energy is higher than that of the interband transitions in solid: $h\nu > E_g = 3.0$ eV. This excitation created a huge number density of the electron-hole plasma that induced surface modifications resulting in nanoparticle-covered structures (NCS), low-spatial frequency LIPSS (LSFL), grooves and featureless flat area (FFA). The excitation beams quality (Gaussian), fluence and superposition on target were permanently monitored in these experiments.

The interaction of light with transient surface states can lead to the onset and suppression of LSFL, grooves and FFA structures, providing physical control of the surface modification. The process deactivation time was estimated to be $\tau^* = 6 \pm 1$ ps close to the threshold fluence. A strong decrease by a factor of 3 of the grooves onset energy was observed on the (001) oriented crystal after introduction of a temporal delay of ~ 8 ps between two laser beams. The last effect was confirmed on the (100) oriented crystal, however unusual FFA was observed instead of grooves, which is connected to the crystal anisotropy. The mechanism of the grooves and FFA formation in experiments with two delayed laser beams are different from that triggered by one laser pulse. We concluded that the both structures are promoted by the same sequence of transient energy states; this conclusion is supported by the analysis of variation of the critical temporal delay with laser fluence, required for the disappearance of LSFL and appearance of grooves and FFA. The experimental observations were rationalized in the framework of a phenomenological model describing evolution of transient states energy along the reaction coordinate.

As we stated, the critical delays between two pulses τ_r and τ_g separate two domains of ripples ($t \leq \tau_r$) and grooves/FFA ($t \geq \tau_g$) formation. We notice that one broad pulse of the shorter duration is expected to trig ripples. In contrast, by using one longer pulse the process coherence

will be disturbed, which would result in a poor-resolved picture of the surface structuring. Furthermore, recent theoretical modelling [30] predict in this case a loss of efficiency of the grooves/FFA formation.

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