



**HAL**  
open science

# Classical and microlocal analysis of the x-ray transform on Anosov manifolds

Sébastien Gouëzel, Thibault Lefeuvre

► **To cite this version:**

Sébastien Gouëzel, Thibault Lefeuvre. Classical and microlocal analysis of the x-ray transform on Anosov manifolds. *Analysis & PDE*, 2021, 14 (1), pp.301-322. 10.2140/apde.2021.14.301 . hal-03400842

**HAL Id: hal-03400842**

**<https://hal.science/hal-03400842>**

Submitted on 25 Oct 2021

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# CLASSICAL AND MICROLOCAL ANALYSIS OF THE X-RAY TRANSFORM ON ANOSOV MANIFOLDS

SÉBASTIEN GOUËZEL AND THIBAUT LEFEUVRE

ABSTRACT. We complete the microlocal study of the geodesic X-ray transform on Riemannian manifolds with Anosov geodesic flow initiated by Guillarmou in [Gui17] and pursued by Guillarmou and the second author in [GL18]. We prove new stability estimates and clarify some properties of the operator  $\Pi_m$  — the generalized X-ray transform. These estimates rely on a refined version of the Livsic theorem for Anosov flows, especially on a new quantitative finite time Livsic theorem.

## 1. INTRODUCTION

Let  $\mathcal{M}$  be a smooth closed  $(n + 1)$ -dimensional manifold endowed with a vector field  $X$  generating a complete flow  $(\varphi_t)_{t \in \mathbb{R}}$ . We assume that the flow  $(\varphi_t)_{t \in \mathbb{R}}$  is transitive and Anosov in the sense that there exists a continuous flow-invariant splitting

$$T_x(\mathcal{M}) = \mathbb{R}X(x) \oplus E_u(x) \oplus E_s(x), \quad (1.1)$$

where  $E_s(x)$  (resp.  $E_u(x)$ ) is the *stable* (resp. *unstable*) vector space at  $x \in \mathcal{M}$ , and a smooth Riemannian metric  $g$  such that

$$\begin{aligned} |d\varphi_t(x) \cdot v|_{\varphi_t(x)} &\leq C e^{-\lambda t} |v|_x, \quad \forall t > 0, v \in E_s(x) \\ |d\varphi_t(x) \cdot v|_{\varphi_t(x)} &\leq C e^{-\lambda|t|} |v|_x, \quad \forall t < 0, v \in E_u(x), \end{aligned} \quad (1.2)$$

for some uniform constants  $C, \lambda > 0$ . The norm, here, is  $|\cdot|_x := g_x(\cdot, \cdot)^{1/2}$ . The dimension of  $E_s$  (resp.  $E_u$ ) is denoted by  $n_s$  (resp.  $n_u$ ). As a consequence,  $n + 1 = 1 + n_s + n_u$  (where the 1 stands for the *neutral* direction, that is the direction of the flow). The case we will have in mind will be that of a geodesic flow on the unit tangent bundle of a smooth Riemannian manifold  $(M, g)$  with negative sectional curvature.

**1.1. X-ray transform on  $\mathcal{M}$ .** We denote by  $\mathcal{G}$  the set of closed orbits of the flow and for  $f \in C^0(\mathcal{M})$ , its X-ray transform  $If$  is defined by:

$$\mathcal{G} \ni \gamma \mapsto If(\gamma) := \langle \delta_\gamma, f \rangle = \frac{1}{\ell(\gamma)} \int_0^{\ell(\gamma)} f(\varphi_t x) dt,$$

where  $x \in \gamma$ ,  $\ell(\gamma)$  is the length of  $\gamma$ .

The Livsic theorem characterizes the kernel of the X-ray transform for a hyperbolic flow: the latter is reduced to the coboundaries, which are the functions of the form  $f = Xu$ , where  $u$  is a function defined on  $\mathcal{M}$  whose regularity is prescribed by that of  $f$ . This

result was initially proved by Livsic [Liv72] in Hölder regularity: if  $f \in C^\alpha(\mathcal{M})$  is such that  $If = 0$ , then there exists  $u \in C^\alpha(\mathcal{M})$ , differentiable in the flow direction, such that  $f = Xu$  and  $u$  is unique up to an additive constant. There is also a version of the Livsic theorem in smooth regularity due to De la Llave-Marco-Moriyon [dLLMM86]. Much more recently, Guillarmou [Gui17, Corollary 2.8] proved a version of the Livsic theorem in Sobolev regularity which implies the theorem of [dLLMM86].

It is also rather natural to expect other versions of the Livsic theorem to hold. For instance, if we modify the condition  $If = 0$  by  $If \geq 0$ , is it true that  $f \geq Xu$ , for some well-chosen function  $u$  (positive Livsic theorem)? And if  $\|If\|_{\ell^\infty} := \sup_{\gamma \in \mathcal{G}} |If(\gamma)| \leq \varepsilon$ , can one write  $f = Xu + h$ , where some norm of  $h$  is controlled by a power of  $\varepsilon$  (approximate Livsic theorem)? Eventually, what can be said if  $If(\gamma) = 0$  for all closed orbits  $\gamma$  of length  $\leq L$  (finite Livsic theorem)?

The positive Livsic theorem for Anosov flows was proved by Lopes-Thieullen [LT05] with an explicit control of a Hölder norm of the coboundary  $Xu$  in terms of a norm of  $f$ .

**Theorem 1.1** (Lopes-Thieullen). *Let  $0 < \alpha \leq 1$ . There exist  $0 < \beta \leq \alpha$ ,  $C > 0$  such that: for all functions  $f \in C^\alpha(\mathcal{M})$ , there exist  $u \in C^\beta(\mathcal{M})$ , differentiable in the flow-direction with  $Xu \in C^\beta(\mathcal{M})$  and  $h \in C^\beta(\mathcal{M})$ , such that  $f = Xu + h + m(f)$ , with  $h \geq 0$  and  $m(f) = \inf_{\gamma \in \mathcal{G}} If(\gamma)$ . Moreover,  $\|Xu\|_{C^\beta} \leq C\|f\|_{C^\alpha}$ .*

In this article, we prove a finite approximate version of the Livsic theorem, as follows.

**Theorem 1.2.** *Let  $0 < \alpha \leq 1$ . There exist  $0 < \beta \leq \alpha$  and  $\tau, C > 0$  with the following property. Let  $\varepsilon > 0$ . Consider a function  $f \in C^\alpha(\mathcal{M})$  with  $\|f\|_{C^\alpha(\mathcal{M})} \leq 1$  such that  $|If(\gamma)| \leq \varepsilon$  for all  $\gamma$  with  $\ell(\gamma) \leq \varepsilon^{-1/2}$ . Then there exist  $u \in C^\beta(\mathcal{M})$  differentiable in the flow-direction with  $Xu \in C^\beta(\mathcal{M})$  and  $h \in C^\beta(\mathcal{M})$ , such that  $f = Xu + h$ . Moreover,  $\|u\|_{C^\beta} \leq C$  and  $\|h\|_{C^\beta} \leq C\varepsilon^\tau$ .*

We note that a rather similar result had already been obtained by S. Katok [Kat90] in the particular case of a contact Anosov flow on a 3-manifold.

The assumptions of Theorem 1.2 hold in particular if  $\|If\|_{\ell^\infty} = \sup_{\gamma \in \mathcal{G}} |If(\gamma)| \leq \varepsilon$ . Under the assumptions of the theorem (only mentioning the closed orbits of length at most  $\varepsilon^{-1/2}$ ), the decomposition  $f = Xu + h$  also gives a global control on  $\|If\|_{\ell^\infty}$ , of the form

$$\|If\|_{\ell^\infty} \leq C\varepsilon^\tau. \quad (1.3)$$

Indeed, if one integrates  $f = Xu + h$  along a closed orbit of any length, the contribution of  $Xu$  vanishes and one is left with a bound  $\|h\|_{C^0} \leq C\varepsilon^\tau$ . The bound (1.3) holds in particular if  $If(\gamma) = 0$  for all  $\gamma$  with  $\ell(\gamma) \leq \varepsilon^{-1/2}$ . This statement illustrates quantitatively the fact that the quantities  $If(\gamma)$  for different  $\gamma$  are far from being independent.

*Remark 1.3.* In Theorem 1.2, the constants  $\beta, C, \tau$  depend on the Anosov flow under consideration, but in a locally uniform way: given an Anosov flow, one can find such parameters

that work for any flow in a neighborhood of the initial flow. The local uniformity can be checked either directly from the proof, or using a (Hölder-continuous) orbit-conjugacy between the initial flow and the perturbed one.

*Remark 1.4.* It could be interesting to extend the positive and the finite approximate Livsic theorems to other regularities like  $H^s$  spaces for  $s > \frac{n+1}{2}$  but we were unable to do so.

**1.2. X-ray transform for the geodesic flow.** If  $(M, g)$  is a smooth closed Riemannian manifold, we set  $\mathcal{M} := SM$ , the unit tangent bundle, and denote by  $X$  the geodesic vector field on  $SM$ . We will always assume that the geodesic flow is Anosov on  $SM$  and we say that  $(M, g)$  is an *Anosov Riemannian manifold*. It is a well-known fact that a negatively-curved manifold has Anosov geodesic flow. We will denote by  $\mathcal{C}$  the set of free homotopy classes on  $M$ : they are in one-to-one correspondence with the set of conjugacy classes of  $\pi_1(M, \bullet)$ . If  $(M, g)$  is Anosov, we know by [Kli74] that given a free homotopy class  $c \in \mathcal{C}$ , there exists a unique closed geodesic  $\gamma \in \mathcal{G}$  belonging to the free homotopy class  $c$ . In other words,  $\mathcal{G}$  and  $\mathcal{C}$  are in one-to-one correspondence. As a consequence, we will rather see the X-ray transform as a map  $I^g : C^0(SM) \rightarrow \ell^\infty(\mathcal{C})$  and we will drop the index  $g$  if the context is clear.

If  $f \in C^\infty(M, \otimes_S^m T^*M)$  is a symmetric tensor, then by §2, we can see  $f$  as a function  $\pi_m^* f \in C^\infty(SM)$ , where  $\pi_m^* f(x, v) := f_x(v, \dots, v)$ . The X-ray transform  $I_m$  of  $f$  is simply defined by  $I_m f := I \circ \pi_m^* f$ . In other words, it consists in integrating the tensor  $f$  along closed geodesics by plugging  $m$ -times the speed vector in  $f$ .

The natural operator of derivation of symmetric tensors is  $D := \sigma \circ \nabla$ , where  $\nabla$  is the Levi-Civita connection and  $\sigma$  is the operator of symmetrization of tensors (see §2). Any smooth tensor  $f \in C^\infty(M, \otimes_S^m T^*M)$  can be uniquely decomposed as  $f = Dp + h$ , where  $p \in C^\infty(M, \otimes_S^{m-1} T^*M)$  and  $h \in C^\infty(M, \otimes_S^m T^*M)$  is a *solenoidal tensor* i.e., a tensor such that  $D^*h = 0$ , where  $D^*$  is the formal adjoint of  $D$ . We say that  $Dp$  is the *potential part* of the tensor  $f$ . We will see that  $I_m(Dp) = 0$ . In other words, the potential tensors are always in the kernel of the X-ray transform. We will say that  $I_m$  is *solenoidal injective* or in short s-injective if injective when restricted to

$$C_{\text{sol}}^\infty(M, \otimes_S^m T^*M) := C^\infty(M, \otimes_S^m T^*M) \cap \ker(D^*)$$

Note that we will often add an index sol to a functional space on tensors to denote the fact that we are considering the intersection with  $\ker D^*$ .

It is conjectured that  $I_m$  is s-injective for all Anosov Riemannian manifolds, in any dimension and without any assumption on the curvature. Under the additional assumption that the sectional curvatures are non-positive, the *Pestov energy identity* allows to show injectivity (see [GK80] and [CS98] for the original proofs). Without any assumption on the curvature, this is still true for surfaces by [PSU14] and [Gui17]. In higher dimensions, it holds for  $m = 0, 1$  (see [DS03]) but remains an open question for higher

order tensors without any assumption on the curvature. However, it is already known that  $C_{\text{sol}}^\infty(M, \otimes_S^m T^*M) \cap \ker(I_m)$  is finite-dimensional.

We will prove the following stability estimate on  $I_m$ :

**Theorem 1.5.** *Assume  $I_m$  is  $s$ -injective. Then, for all  $\frac{n+1}{2} < s < r$ , there exists  $\theta := \theta(s, r) > 0$  and  $C := C(s, r) > 0$  such that: if  $f \in H_{\text{sol}}^r(M, \otimes_S^m T^*M)$  is a solenoidal symmetric  $m$ -tensor such that  $\|f\|_{H^r} \leq 1$ , then  $\|f\|_{H^s} \leq C\|I_m f\|_{\ell^\infty}^\theta$ .*

Actually, if  $I_m$  is not known to be injective, one still has the previous estimate by taking  $f$  solenoidal and orthogonal to the kernel of  $I_m$ . Combining this estimate with Theorem 1.2 (and more specifically (1.3)), we immediately obtain:

**Theorem 1.6.** *For all  $\frac{n+1}{2} < s < r$ , there exists  $C, \beta > 0$  such that for any  $L > 0$  large enough: if  $f \in H_{\text{sol}}^r(M, \otimes_S^m T^*M)$  is a solenoidal symmetric  $m$ -tensor such that  $\|f\|_{H^r} \leq 1$  and  $I_m f(\gamma) = 0$  for all closed geodesics  $\gamma \in \mathcal{C}$  such that  $\ell(\gamma) \leq L$ , then  $\|f\|_{H^s} \leq CL^{-\beta}$ .*

The X-ray transform  $I_m$  may appear in different contexts. In particular,  $I_2$  is well-known to be the differential of the marked length spectrum and it was studied in [GL18] to prove its rigidity, thus partially answering the conjecture of Burns-Katok [BK85].

**Acknowledgements:** We thank Colin Guillarmou for fruitful discussions. TL acknowledges partial support from ERC funding COG IPFLOW.

## 2. ON SYMMETRIC TENSORS

We describe elementary properties of symmetric tensors on Riemannian manifolds.

### 2.1. Definitions and first properties.

2.1.1. *Symmetric tensors in Euclidean space.* Let  $E$  be a Euclidean  $(n+1)$ -dimensional vector space endowed with a metric  $g$  and let  $(\mathbf{e}_1, \dots, \mathbf{e}_{n+1})$  be an orthonormal basis. We say that a tensor  $f \in \otimes^m E^*$  is symmetric if  $f(v_1, \dots, v_m) = f(v_{\tau(1)}, \dots, v_{\tau(m)})$ , for all  $v_1, \dots, v_m \in E$  and  $\tau \in \mathfrak{S}_m$ , the group of permutations of order  $m$ . We denote by  $\otimes_S^m T^*E$  the vector space of symmetric  $m$ -tensors on  $E$ . There is a natural projection  $\sigma : \otimes^m E^* \rightarrow \otimes_S^m E^*$  given by

$$\sigma(v_1^* \otimes \dots \otimes v_m^*) = \frac{1}{m!} \sum_{\tau \in \mathfrak{S}_m} v_{\tau(1)}^* \otimes \dots \otimes v_{\tau(m)}^*,$$

for all  $v_1^*, \dots, v_m^* \in E^*$ . The metric  $g$  induces a scalar product  $\langle \cdot, \cdot \rangle$  on  $\otimes^m E^*$  by declaring the basis  $(\mathbf{e}_{i_1}^* \otimes \dots \otimes \mathbf{e}_{i_m}^*)_{1 \leq i_1, \dots, i_m \leq n+1}$  to be orthonormal which yields

$$\langle u_1^* \otimes \dots \otimes u_m^*, v_1^* \otimes \dots \otimes v_m^* \rangle = \prod_{i=1}^m g^{-1}(u_i^*, v_i^*),$$

where  $g^{-1}$  is the dual metric, that is the natural metric on  $E^*$  which makes the musical isomorphism  $\sharp : E \rightarrow E^*$  an isometry. Since  $\sigma$  is self-adjoint with respect to this metric, it is an orthogonal projection. Let  $(g_{ij})_{1 \leq i, j \leq n+1}$  denote the metric  $g$  in the coordinates  $(x_1, \dots, x_{n+1})$ . Then the metric can be expressed as

$$\langle f, h \rangle = \sum_{i_1, \dots, i_m=1}^{n+1} f_{i_1 \dots i_m} h^{i_1 \dots i_m},$$

where  $h^{i_1 \dots i_m} = \sum_{j_1, \dots, j_m=1}^{n+1} g^{i_1 j_1} \dots g^{i_m j_m} h_{j_1 \dots j_m}$ . We define the trace  $\text{Tr}_g : \otimes_S^m E^* \rightarrow \otimes_S^{m-2} E^*$  of a symmetric tensor by

$$\text{Tr}_g(f) = \sum_{i=1}^{n+1} f(\mathbf{e}_i, \mathbf{e}_i, \cdot, \dots, \cdot).$$

In coordinates,  $\text{Tr}_g(f)(v_2, \dots, v_m) = \text{Tr}(g^{-1}f(\cdot, \cdot, v_2, \dots, v_m))$ . Its adjoint with respect to the scalar products is the map  $I : \otimes_S^{m-2} E^* \rightarrow \otimes_S^m E^*$  given by  $I(u) = \sigma(g \otimes u)$ .

Symmetric tensors can also be seen as homogeneous polynomials on the unit sphere of the Euclidean space. We denote by  $\mathbb{S}_E$  the  $n$ -dimensional unit sphere on  $(E, g)$  and by  $dS$  the Riemannian measure on the sphere induced by the metric  $g|_{\mathbb{S}_E}$ . We define  $\pi_m : (x, v) \mapsto (x, \otimes^m v)$  for  $v \in E$ ; it induces a canonical morphism  $\pi_m^* : \otimes_S^m E^* \rightarrow C^\infty(\mathbb{S}_E)$  given by  $\pi_m^* f(v) = f(v, \dots, v)$ . Its formal adjoint is  $\langle \pi_m^* f, h \rangle_{L^2(\mathbb{S}_E, dS)} = \langle f, \pi_{m*} h \rangle_{\otimes^m T^*M}$ , where  $f \in \otimes_S^m T^*M, h \in C^\infty(\mathbb{S}_E)$ . In coordinates,

$$(\pi_{m*} h)_{i_1 \dots i_m} := \pi_{m*} h(\partial_{i_1}, \dots, \partial_{i_m}) = \sum_{j_1, \dots, j_m=1}^{n+1} g_{i_1 j_1} \dots g_{i_m j_m} \int_{\mathbb{S}_E} h(v) v_{j_1} \dots v_{j_m} dS \quad (2.1)$$

Also remark that (2.1) can be rewritten intrinsically as

$$\forall u_1, \dots, u_m \in E, \quad \pi_{m*} h(u_1, \dots, u_m) = \int_{\mathbb{S}_E} h(v) g(v, u_1) \dots g(v, u_m) dv \quad (2.2)$$

The map  $\pi_{m*} \pi_m^*$  is an isomorphism we will study in the next paragraph. Also note that  $\pi_m^*(\sigma f) = \pi_m^* f$  (since all the antisymmetric parts of the tensor  $f$  vanish by plugging  $m$  times the same vector  $v$ ).

We denote by  $j_\xi$  the multiplication by  $\xi$ , that is  $j_\xi : f \mapsto \xi \otimes f$ , and by  $i_\xi$  the contraction, that is  $i_\xi : f \mapsto u(\xi^\sharp, \cdot, \dots, \cdot)$ . The adjoint of  $i_\xi$  on symmetric tensors with respect to the  $L^2$ -scalar product is  $\sigma j_\xi$ , that is

$$\forall f \in \otimes_S^{m-1} E^*, h \in \otimes_S^m E^*, \quad \langle \sigma j_\xi f, h \rangle = \langle f, i_\xi h \rangle$$

The space  $\otimes_S^m E^*$  can thus be decomposed as the direct sum

$$\otimes_S^m E^* = \text{ran} \left( \sigma j_\xi |_{\otimes_S^{m-1} E^*} \right) \oplus^\perp \ker \left( i_\xi |_{\otimes_S^m E^*} \right)$$

We denote by  $\pi_{\ker i_\xi}$  the projection onto the right space, parallel to the left space. We will need the following

**Lemma 2.1.** *For all  $f, h \in \otimes_S^m E^*$ ,*

$$C_{n,m} \int_{\langle \xi, v \rangle = 0} \pi_m^* f(v) \pi_m^* h(v) dS_\xi(v) = \langle \pi_{\ker i_\xi} \pi_{m*} \pi_m^* \pi_{\ker i_\xi} f, h \rangle,$$

where

$$C_{n,m} = \int_0^\pi \sin^{n-1+2m}(\varphi) d\varphi = \sqrt{\pi} \frac{\Gamma((n+2m)/2)}{\Gamma((n+1+2m)/2)},$$

and  $dS_\xi$  is the canonical measure induced on the  $n-1$  dimensional sphere  $\mathbb{S}_{E,\xi} := \mathbb{S}_E \cap \{\langle \xi, v \rangle = 0\}$ .

*Proof.* We can write  $h = \sigma j_\xi h_1 + h_2$  where  $h_1 \in \otimes_S^{m-1} E^*$ ,  $h_2 \in \ker(i_\xi|_{\otimes_S^m T_x^* M})$ . Note that  $\pi_m^*(\sigma j_\xi h_1)(v) = \pi_m^*(j_\xi h_1)(v) = \langle \xi, v \rangle \pi_{m-1}^* h_1(v)$  and this vanishes on  $\{\langle \xi, v \rangle = 0\}$  (and the same holds for  $f$ ). In other words,  $\pi_m^* h = \pi_m^* \pi_{\ker i_\xi}$  on  $\{\langle \xi, v \rangle = 0\}$ . We are thus left to check that for  $f, h \in \ker i_\xi$ ,

$$C_{n,m} \int_{\langle \xi, v \rangle = 0} \pi_m^* f(v) \pi_m^* h(v) dS_\xi(v) = \int_{\mathbb{S}_E} \pi_m^* f(v) \pi_m^* h(v) dS(v)$$

We will use the coordinates  $v' = (v, \varphi) \in \mathbb{S}_{E,\xi} \times [0, \pi]$  on  $\mathbb{S}_E$  which allow to decompose  $v' = \sin(\varphi)v + \cos(\varphi)\xi^\# / |\xi|$ . Then the measure on  $\mathbb{S}_E$  disintegrates as  $dS = \sin^{n-1}(\varphi) d\varphi dS_\xi(v)$ . Also remark that  $\pi_m^* f(v + \cos(\varphi)\xi^\# / |\xi|) = \pi_m^* f(v)$ . Then, if  $C_{n,m} := \int_0^\pi \sin^{n-1+2m}(\varphi) d\varphi$ , we obtain:

$$\begin{aligned} & \int_{\langle \xi, v \rangle = 0} \pi_m^* f(v) \pi_m^* h(v) dS_\xi(v) \\ &= C_{n,m}^{-1} \int_0^\pi \sin^{n-1+2m}(\varphi) d\varphi \int_{\langle \xi, v \rangle = 0} \pi_m^* f(v) \pi_m^* h(v) dS_\xi(v) \\ &= C_{n,m}^{-1} \int_0^\pi \int_{\langle \xi, v \rangle = 0} \pi_m^* f(\sin(\varphi)v + \cos(\varphi)\xi^\# / |\xi|) \\ & \quad \times \pi_m^* h(\sin(\varphi)v + \cos(\varphi)\xi^\# / |\xi|) \sin^{n-1}(\varphi) d\varphi dS_\xi(v) \\ &= C_{n,m}^{-1} \int_{\mathbb{S}_E} \pi_m^* f(v') \pi_m^* h(v') dS(v') \end{aligned}$$

□

**2.1.2. Spherical harmonics.** Let  $\Delta|_{\mathbb{S}_E} := \operatorname{div}_{\mathbb{S}_E} \nabla_{\mathbb{S}_E}$  be the Laplacian on the unit sphere  $\mathbb{S}_E$  induced by the metric  $g|_{\mathbb{S}_E}$  and  $\Delta$  be the usual Laplacian on  $E$  induced by  $g$ . Let

$$L^2(\mathbb{S}_E) = \bigoplus_{m=0}^{+\infty} \Omega_m$$

be the spectral break up in spherical harmonics, where  $\Omega_m := \ker(\Delta|_{\mathbb{S}_E} + m(m+n-1))$  are the eigenspaces of the Laplacian. We denote by  $E_m$  the vector space of trace-free symmetric  $m$ -tensors, where the trace is, as before, taken over the first two coordinates.



**Lemma 2.2.**  $\pi_m^* : E_m \rightarrow \Omega_m$  is an isomorphism and  $\pi_{m*}\pi_m^*|_{E_m} = \lambda_{m,n}1_{E_m}$ , for some constant  $\lambda_{m,n} \neq 0$ .

This also shows that, up to rescaling by the constant  $\lambda_{m,n}$ ,  $\pi_m^* : E_m \rightarrow \Omega_m$  is an isometry. One could be more accurate and actually show that the maps

$$\pi_m^* : \otimes_S^m E^* \rightarrow \bigoplus_{k=0}^{\lfloor m/2 \rfloor} \Omega_{m-2k}, \quad \pi_{m*} \bigoplus_{k=0}^{\lfloor m/2 \rfloor} \Omega_{m-2k} \rightarrow \otimes_S^m E^* \quad (2.3)$$

are isomorphisms, where  $\lfloor m/2 \rfloor$  stands for the integer part of  $m/2$ . This follows from the (unique) decomposition of a symmetric tensor into a trace-free part and a remainder (which lies in the image of the adjoint of  $\text{Tr}$ ). More precisely, by iterating this process, one can decompose  $u$  as  $u = \sum_{k=0}^{\lfloor m/2 \rfloor} I^k(u_k)$ , where  $I : \otimes_S^\bullet E^* \rightarrow \otimes_S^{\bullet+2} E^*$  is the adjoint of  $\text{Tr}$  with respect to the scalar products and  $u_k \in \otimes_S^{m-2k} E^*$ ,  $\text{Tr}(u_k) = 0$  and  $\pi_m^* I^k(u_k) \in \Omega_{m-2k}$ . Then (2.3) is an immediate consequence of the previous lemma. The map  $\pi_{m*}\pi_m^*$  acts by scalar multiplication on each component  $I^k(u_k)$  (but with a different constant though, so  $\pi_{m*}\pi_m^*$  is not a multiple of the identity). Since we will only need the fact that  $\pi_{m*}\pi_m^*$  is an isomorphism, we do not provide further details.

*2.1.3. Symmetric tensors on a Riemannian manifold. Decomposition in solenoidal and potential tensors.* We now consider the Riemannian manifold  $(M, g)$  and denote by  $d\mu$  the Liouville measure on the unit tangent bundle  $SM$ . All the previous definitions naturally extend to the vector bundle  $TM \rightarrow M$ . For  $f, h \in C^\infty(M, \otimes_S^m T^*M)$ , we define the  $L^2$ -scalar product

$$\langle f, h \rangle = \int_M \langle f_x, h_x \rangle_x d\text{vol}(x),$$

where  $\langle \cdot, \cdot \rangle_x$  is the scalar product on  $T_x M$  introduced in the previous paragraph and  $d\text{vol}$  is the Riemannian measure induced by  $g$ . The map  $\pi_m^* : C^\infty(M, \otimes_S^m T^*M) \rightarrow C^\infty(SM)$  is the canonical morphism given by  $\pi_m^* f(x, v) = f_x(v, \dots, v)$ , whose formal adjoint with respect to the two  $L^2$ -inner products (on  $L^2(SM, d\mu)$  and  $L^2(\otimes_S^m T^*M, d\text{vol})$ ) is  $\pi_{m*}$ , i.e.,  $\langle \pi_m^* f, h \rangle_{L^2(SM, d\mu)} = \langle f, \pi_{m*} h \rangle_{L^2(\otimes_S^m T^*M, d\text{vol})}$ .

If  $\nabla$  denotes the Levi-Civita connection, we set  $D := \sigma \circ \nabla : C^\infty(M, \otimes_S^m T^*M) \rightarrow C^\infty(M, \otimes_S^{m+1} T^*M)$ , the symmetrized covariant derivative. Its formal adjoint with respect to the  $L^2$ -scalar product is  $D^* = -\text{Tr}(\nabla)$  where the trace is taken with respect to the two first indices, like in 2.1.1. One has the following relation between the geodesic vector field  $X$  on  $SM$  and the operator  $D$ :

**Lemma 2.3.**  $X\pi_m^* = \pi_{m+1}^* D$

The operator  $D$  can be seen as a differential operator of order 1. Its principal symbol is given by  $\sigma(D)(x, \xi)f \mapsto \sigma(\xi \otimes f) = \sigma_j \xi_j f$  (see [Sha94, Theorem 3.3.2]).

**Lemma 2.4.**  $D$  is elliptic. It is injective on tensors of odd order, and its kernel is reduced to  $\mathbb{R}\sigma(g^{\otimes m/2})$  on even tensors.



When  $m$  is even, we will denote by  $K_m = c_m \sigma(g^{\otimes m/2})$ , with  $c_m > 0$ , a unitary vector in the kernel of  $D$ .

*Proof.* We fix  $(x, \xi) \in T^*M$ . For a tensor  $u \in \otimes_S^m T_x^*M$ , using the fact that the antisymmetric part of  $\xi \otimes u$  vanishes in the integral:

$$\langle \sigma(D)u, \sigma(D)u \rangle = \int_{\mathbb{S}_x^n} \langle \xi, v \rangle^2 \pi_m^* u^2(v) dS_x(v) = |\xi|^2 \int_{\mathbb{S}_x^n} \langle \xi/|\xi|, v \rangle^2 \pi_m^* u^2(v) dS_x(v) > 0,$$

unless  $u \equiv 0$ . Since  $\otimes_S^m T_x^*M$  is finite dimensional, the map

$$(u, \xi/|\xi|) \mapsto \langle \sigma(D)(x, \xi/|\xi|)u, \sigma(D)(x, \xi/|\xi|)u \rangle,$$

defined on the compact set  $\{u \in \otimes_S^m T_x^*M, |u|^2 = 1\} \times \mathbb{S}^n$  is bounded and attains its lower bound  $C^2 > 0$  (which is independent of  $x$ ). Thus  $\|\sigma(x, \xi)\| \geq C|\xi|$ , so the operator is uniformly elliptic and can be inverted (on the left) modulo a compact remainder: there exists pseudodifferential operators  $Q, R$  of respective order  $-1, -\infty$  such that  $QD = 1 + R$ .

As to the injectivity of  $D$ : if  $Df = 0$  for some tensor  $f \in \mathcal{D}'(M, \otimes_S^m T^*M)$ , then  $f$  is smooth and  $\pi_{m+1}^* Df = X \pi_m^* f = 0$ . By ergodicity of the geodesic flow,  $\pi_m^* f = c \in \Omega_0$  is constant. If  $m$  is odd, then  $\pi_m^* f(x, v) = -\pi_m^* f(x, -v)$  so  $f \equiv 0$ . If  $m$  is even, then, by §2.1.2,  $f = I^{m/2}(u_{m/2})$  where  $u_{m/2} \in \otimes_S^0 E^* \simeq \mathbb{R}$  so  $f = c' \sigma(g^{\otimes m/2})$ .  $\square$

By classical elliptic theory, the ellipticity of  $D$  implies that

$$H^s(M, \otimes_S^m T^*M) = D(H^{s+1}(M, \otimes_S^{m-1} T^*M)) \oplus \ker D^*|_{H^s(M, \otimes_S^m T^*M)}, \quad (2.4)$$

and the decomposition still holds in the smooth category and in the  $C^{k,\alpha}$ -topology for  $k \in \mathbb{N}, \alpha \in (0, 1)$ . This is the content of the following theorem:

**Theorem 2.5** (Tensor decomposition). *Let  $s \in \mathbb{R}$  and  $f \in H^s(M, \otimes_S^m T^*M)$ . Then, there exists a unique pair of symmetric tensors  $(p, h) \in H^{s+1}(M, \otimes_S^{m-1} T^*M) \times H^s(M, \otimes_S^m T^*M)$  such that  $f = Dp + h$  and  $D^*h = 0$ . Moreover, if  $m = 2l + 1$  is odd,  $\langle p, K_{2l} \rangle = 0$ .*

*Tensorial distributions.* The spaces  $H^s(M, \otimes_S^m T^*M)$  that have been mentioned so far are the  $L^2$ -based Sobolev spaces of order  $s \in \mathbb{R}$ . They can be defined in coordinates (each coordinate of the tensor has to be in  $H_{\text{loc}}^s(\mathbb{R})$ ) or more intrinsically by setting  $H^s(M, \otimes_S^m T^*M) := (1 + D^*D)^{-s/2} L^2(M, \otimes_S^m T^*M)$ . These two definitions are equivalent by [Shu01, Proposition 7.3], following the properties of the operator  $1 + D^*D$  (it is elliptic, invertible, positive). In the same fashion, the spaces  $L^p(M, \otimes_S^m T^*M)$ , for  $p \geq 1$  can be defined in coordinates. Note that the maps

$$\pi_m^* : H^s(M, \otimes_S^m T^*M) \rightarrow H^s(SM), \quad \pi_{m*} : H^s(SM) \rightarrow H^s(M, \otimes_S^m T^*M).$$

are bounded for all  $s \in \mathbb{R}$  (and they are bounded on  $L^p$ -spaces for  $p \geq 1$ ). The operator  $\pi_{m*}$  acts by duality on distributions, namely:

$$\pi_{m*} : C^{-\infty}(SM) \rightarrow C^{-\infty}(M, \otimes_S^m T^*M), \quad \langle \pi_{m*} f_1, f_2 \rangle := \langle f_1, \pi_m^* f_2 \rangle$$

where  $\langle \cdot, \cdot \rangle$  denotes the distributional pairing.

*The projection on solenoidal tensors.* When  $m$  is even, we denote by  $\Pi_{K_m} := \langle K_m, \cdot \rangle K_m$  the orthogonal projection on  $\mathbb{R}K_m$ . We define  $\Delta_m := D^*D + \varepsilon(m)\Pi_{K_m}$ , where  $\varepsilon(m) = 1$  for  $m$  even,  $\varepsilon(m) = 0$  for  $m$  odd. The operator  $\Delta_m$  is an elliptic differential operator of order 2 which is invertible: as a consequence, its inverse is also pseudodifferential of order  $-2$  (see [Shu01, Theorem 8.2]). We can thus define the operator

$$\pi_{\ker D^*} := 1 - D\Delta_m^{-1}D^*. \quad (2.5)$$

One can check that this is exactly the  $L^2$ -orthogonal projection on solenoidal tensors, it is a pseudodifferential operator of order 0 (as a composition of pseudodifferential operators).

Since  $\sigma(D)(x, \xi) = \sigma j_\xi$ , we know by §2.1.1 that given  $(x, \xi) \in T^*M$ , the space  $\otimes_S^m T_x^*M$  breaks up as the direct sum

$$\begin{aligned} \otimes_S^m T_x^*M &= \text{ran} \left( \sigma(D)(x, \xi)|_{\otimes_S^{m-1} T_x^*M} \right) \oplus \ker \left( \sigma(D^*)(x, \xi)|_{\otimes_S^m T_x^*M} \right) \\ &= \text{ran} \left( \sigma j_\xi|_{\otimes_S^{m-1} T_x^*M} \right) \oplus \ker \left( i_\xi|_{\otimes_S^m T_x^*M} \right) \end{aligned}$$

We recall that  $\pi_{\ker i_\xi}$  is the projection on  $\ker(i_\xi|_{\otimes_S^m T_x^*M})$  parallel to  $\text{ran}(\sigma j_\xi|_{\otimes_S^{m-1} T_x^*M})$ .

**Lemma 2.6.** *The principal symbol of  $\pi_{\ker D^*}$  is  $\sigma_{\pi_{\ker D^*}} = \pi_{\ker i_\xi}$ .*

*Proof.* First, observe that:

$$\begin{aligned} D\Delta_m^{-1}D^*D\Delta_m^{-1}D^* &= D\Delta_m^{-1}(\Delta_m - \varepsilon(m)\Pi_{K_m})\Delta_m^{-1}D^* \\ &= D\Delta_m^{-1}D^* - \varepsilon(m)D\Delta_m^{-1}\Pi_{K_m}\Delta_m^{-1}D^* \end{aligned}$$

The second operator is smoothing so at the principal symbol level

$$\sigma_{(D\Delta_m^{-1}D^*)^2} = \sigma_{D\Delta_m^{-1}D^*}^2 = \sigma_{D\Delta_m^{-1}D^*},$$

which implies that  $\sigma_{D\Delta_m^{-1}D^*}$  is a projection. Moreover,  $\sigma_{D\Delta_m^{-1}D^*} = \sigma_D\sigma_{\Delta_m^{-1}}\sigma_{D^*} = \sigma j_\xi\sigma_{\Delta_m^{-1}}i_\xi$ , so it is the projection onto  $\text{ran} \sigma j_\xi$  with kernel  $\ker i_\xi$ . Since  $\pi_{\ker D^*} = 1 - D\Delta_m^{-1}D^*$ , the result is immediate.  $\square$

### 3. ON LIVSIC-TYPE THEOREMS

We will denote by  $d : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$  the Riemannian distance on  $\mathcal{M}$  inherited from the Riemannian metric  $g$ . The  $\alpha$ -Hölder norm of  $f$  is defined by:

$$\|f\|_{C^\alpha} := \sup_{x \in \mathcal{M}} |f(x)| + \sup_{x, y \in \mathcal{M}, x \neq y} \frac{|f(x) - f(y)|}{d(x, y)^\alpha} = \|f\|_\infty + \|f\|_\alpha$$

In a series of inequalities, we will sometimes write  $A \lesssim B$  to denote the fact that there exists a universal constant  $C > 0$  such that  $A \leq C \cdot B$ . Note that a constant  $C > 0$  may still appear from time to time and, as usual, it may change from one line to another.

**3.1. Properties of Anosov flows.** We refer to the exhaustive [KH95] and the forthcoming book [HF] for an introduction to hyperbolic dynamics.

**3.1.1. Stable and unstable manifolds.** The *global stable* and *unstable manifolds*  $W^s(x), W^u(x)$  are defined by:

$$\begin{aligned} W^s(x) &= \{x' \in \mathcal{M}, d(\varphi_t(x), \varphi_t(x')) \rightarrow_{t \rightarrow +\infty} 0\} \\ W^u(x) &= \{x' \in \mathcal{M}, d(\varphi_t(x), \varphi_t(x')) \rightarrow_{t \rightarrow -\infty} 0\} \end{aligned}$$

For  $\varepsilon > 0$  small enough, we define the *local stable* and *unstable manifolds*  $W_\varepsilon^s(x) \subset W^s(x), W_\varepsilon^u(x) \subset W^u(x)$  by:

$$\begin{aligned} W_\varepsilon^s(x) &= \{x' \in W^s(x), \forall t \geq 0, d(\varphi_t(x), \varphi_t(x')) \leq \varepsilon\} \\ W_\varepsilon^u(x) &= \{x' \in W^u(x), \forall t \geq 0, d(\varphi_{-t}(x), \varphi_{-t}(x')) \leq \varepsilon\} \end{aligned}$$

For all  $\varepsilon > 0$  small enough, there exists  $t_0 > 0$  such that:

$$\forall x \in \mathcal{M}, \forall t \geq t_0, \varphi_t(W_\varepsilon^s(x)) \subset W_\varepsilon^s(\varphi_t(x)), \varphi_{-t}(W_\varepsilon^u(x)) \subset W_\varepsilon^u(\varphi_{-t}(x)) \quad (3.1)$$

And:

$$T_x W_\varepsilon^s(x) = E_s(x), \quad T_x W_\varepsilon^u(x) = E_u(x)$$

**3.1.2. Classical properties.** The main tool we will use to construct suitable periodic orbits is the following classical shadowing property of Anosov flows. Part of the proof can be found in [KH95, Corollary 18.1.8] and [HF, Theorem 5.3.2]. The last bound is a consequence of hyperbolicity and can be found in [HF, Proposition 6.2.4]. For the sake of simplicity, we will write  $\gamma = [xy]$  if  $\gamma$  is an orbit segment with endpoints  $x$  and  $y$ .

**Theorem 3.1.** *There exist  $\varepsilon_0 > 0$ ,  $\theta > 0$  and  $C > 0$  with the following property. Consider  $\varepsilon < \varepsilon_0$ , and a finite or infinite sequence of orbit segments  $\gamma_i = [x_i y_i]$  of length  $T_i$  greater than 1 such that for any  $n$ ,  $d(y_n, x_{n+1}) \leq \varepsilon$ . Then there exists a genuine orbit  $\gamma$  and times  $\tau_i$  such that  $\gamma$  restricted to  $[\tau_i, \tau_i + T_i]$  shadows  $\gamma_i$  up to  $C\varepsilon$ . More precisely, for all  $t \in [0, T_i]$ , one has*

$$d(\gamma(\tau_i + t), \gamma_i(t)) \leq C\varepsilon e^{-\theta \min(t, T_i - t)}. \quad (3.2)$$

*Moreover,  $|\tau_{i+1} - (\tau_i + T_i)| \leq C\varepsilon$ . Finally, if the sequence of orbit segments  $\gamma_i$  is periodic, then the orbit  $\gamma$  is periodic.*

**Remark 3.2.** In this theorem, we could also allow the first orbit segment  $\gamma_i$  to be infinite on the left, and the last orbit segment  $\gamma_j$  to be infinite on the right. In this case, (3.2) should be replaced by its obvious reformulation: assuming that  $\gamma_i$  is defined on  $(-\infty, 0]$  and  $\gamma_j$  on  $[0, +\infty)$ , we would get for some  $\tilde{\tau}_{i+1}$  within  $C\varepsilon$  of  $\tau_{i+1}$ , and all  $t \geq 0$

$$d(\gamma(\tilde{\tau}_{i+1} - t), \gamma_i(-t)) \leq C\varepsilon e^{-\theta t} \quad (3.3)$$

and

$$d(\gamma(\tau_j + t), \gamma_j(t)) \leq C\varepsilon e^{-\theta t}.$$

In particular, if  $\gamma_0$  is an orbit segment  $[xy]$  with  $d(y, x) \leq \varepsilon_0$ , then applying the above theorem to  $\gamma_i := \gamma_0$  for all  $i \in \mathbb{Z}$ , one gets a periodic orbit that shadows  $\gamma_0$ : this is the *Anosov closing lemma*. We will also use thoroughly the version with two orbit segments that are repeated to get a periodic orbit.

**3.1.3. Cover by parallelepipeds.** We will now fix  $\varepsilon_0$  small enough so that the previous propositions are guaranteed. For  $\varepsilon \leq \varepsilon_0$ , we define the set  $W_\varepsilon(x) := \bigcup_{y \in W_\varepsilon^u(x)} W_\varepsilon^s(x)$ . We can cover the manifold  $\mathcal{M}$  by a finite union of flowboxes  $\mathcal{U}_i := \bigcup_{t \in (-\delta, \delta)} \varphi_t(\Sigma_i)$ , where  $\Sigma_i := W_{\varepsilon_0}(x_i)$  and  $x_i \in \mathcal{M}$ .

We denote by  $\pi_i : \mathcal{U}_i \rightarrow \Sigma_i$  the projection by the flow on the transverse section and we define  $\mathfrak{t}_i : \mathcal{U}_i \rightarrow \mathbb{R}$  such that  $\pi_i(x) = \varphi_{\mathfrak{t}_i(x)}(x)$  for  $x \in \mathcal{U}_i$ . We will need the following lemma:

**Lemma 3.3.**  *$\pi_i, \mathfrak{t}_i$  are Hölder-continuous.*

*Proof.* This is actually a general fact related to the Hölder regularity of the foliation and the smoothness of the flow.

For the sake of simplicity, we drop the index  $i$  in this proof. Let us first prove the Hölder continuity for  $x$  close to  $\Sigma$  and  $x'$  close to  $x$ . We fix  $p \in \Sigma$  and choose smooth local coordinates  $\psi : B(p, \eta) \rightarrow \mathbb{R}^{n+1} = \mathbb{R} \times \mathbb{R}^{n_s} \times \mathbb{R}^{n_u}$  around  $p$  (and centered at 0) so that  $d\psi_p(X) = \partial_{x_0}$ . This choice guarantees that in a neighborhood of 0, the flow is transverse to the hyperplane  $\{0\} \times \mathbb{R}^{n_s+n_u}$ . We still denote by  $\Sigma_\eta$  its image  $\psi(\Sigma_\eta) \subset \mathbb{R}^{n+1}$ , which is a submanifold of Hölder regularity (the index  $\eta$  indicates that we consider the same objects intersected with the ball  $B(x, \eta)$ ). Moreover, there exists a Hölder-continuous homeomorphism  $\Phi : S \rightarrow \Sigma_\eta$ , where  $S \subset \{0\} \times \mathbb{R}^{n_s+n_u}$  (since  $\Sigma_\eta$  is a submanifold of  $M$  with Hölder regularity). We consider  $\hat{\varphi} : (-\delta, \delta) \times S \rightarrow \varphi_{(-\delta, \delta)}(S) =: V \supset \Sigma_\eta$  defined by  $\hat{\varphi}(t, z) = \varphi_t(0, z)$ , which is a smooth diffeomorphism. Remark that  $\mathfrak{t}$  satisfies for  $(0, z) \in S$ ,  $(\mathfrak{t}(z), z) = \hat{\varphi}^{-1}(\Phi(z))$ . So it is Hölder-continuous on  $S$ . Then  $z \mapsto \pi(0, z) = \varphi_{\mathfrak{t}(z)}(0, z)$  is Hölder-continuous on  $S$  too.

We denote by  $\pi_S : V \rightarrow S$  the projection and by  $\mathfrak{t}_S : V \rightarrow S$  the time such that  $\pi_S(x) = \varphi_{\mathfrak{t}_S(x)}(x)$ . These two maps are smooth by the implicit function theorem since the flow is transverse to  $S$ . Moreover, we have:  $\pi(x) = \pi|_S(\pi_S(x))$  so  $\pi$  is Hölder-continuous. And  $\mathfrak{t}(x) = \mathfrak{t}_S(x) + \mathfrak{t}|_S(\pi_S(x))$  so  $\mathfrak{t}$  is Hölder-continuous too. Note that by compactness of  $\Sigma$ , this procedure can be done with only a finite number of charts, thus ensuring the uniformity of the constants. Thus,  $\pi_i, \mathfrak{t}_i$  are Hölder-continuous in a neighborhood of  $\Sigma$ . Now, in order to obtain the continuity on the whole cube  $\mathcal{U}$ , one can repeat the same argument i.e., write the projection as the composition of a first projection on a smooth small section  $S$  defined in a neighborhood of  $\Sigma$  with the actual projection on  $\Sigma$ . This provides the sought result.

□

**3.2. Proof of the approximate Livsic theorem.** We now deal with the proof of Theorem 1.2.

3.2.1. *A key lemma.* The following lemma states that we can find a sufficiently dense and yet separated orbit in the manifold  $\mathcal{M}$ . The separation can only hold transversally to the flow direction, and is defined as follows. Recall that  $W_\varepsilon(x) := \bigcup_{y \in W_\varepsilon^u(x)} W_\varepsilon^s(x)$ . Then we say that a set  $S$  is  $\varepsilon$ -transversally separated if, for any  $x \in S$ , we have  $S \cap W_\varepsilon(x) = \{x\}$ .

**Lemma 3.4.** *Consider a transitive Anosov flow on a compact manifold. There exist  $\beta_s, \beta_d > 0$  such that the following holds. Let  $\varepsilon > 0$  be small enough. There exists a periodic orbit  $\mathcal{O}(x_0) := (\varphi_t x_0)_{0 \leq t \leq T}$  with  $T \leq \varepsilon^{-1/2}$  such that this orbit is  $\varepsilon^{\beta_s}$ -transversally separated and  $(\varphi_t x_0)_{0 \leq t \leq T-1}$  is  $\varepsilon^{\beta_d}$ -dense. If  $\kappa > 0$  is some fixed constant, then one can also require that there exists a piece of  $\mathcal{O}(x_0)$  of length  $\leq C(\kappa)$  which is  $\kappa$ -dense in the manifold.*

*Proof.* We could give a combinatorial construction in terms of Markov partitions and carefully chosen sequences of symbols in the symbolic dynamics representation of the flow. However, controlling rigorously the boundary effects on separation is delicate. Instead, we give a geometric construction solely based on the shadowing theorem. It is easy to obtain an  $\varepsilon^{\beta_d}$ -dense orbit by concatenating orbit segments thanks to the shadowing Theorem. However, separation is harder to enforce. In this proof, we introduce several constants, but none of them will depend on  $\varepsilon$ .

Let us fix two periodic points  $p_1$  and  $p_2$  with different orbits  $\mathcal{O}(p_1)$  and  $\mathcal{O}(p_2)$  of respective lengths  $\ell_1$  and  $\ell_2$ . By the shadowing theorem and transitivity, there exists an orbit  $\gamma_-$  which is asymptotic to  $\mathcal{O}(p_1)$  in negative time and to  $\mathcal{O}(p_2)$  in positive time. Also, there exists an orbit  $\gamma_+$  which is asymptotic to  $\mathcal{O}(p_2)$  in negative time and to  $\mathcal{O}(p_1)$  in positive time. On  $\gamma_-$ , fix a point  $z_0$ , and  $\rho_0 > 0$  small enough so that  $\gamma_- \cup \gamma_+$  meets  $W_{3\rho_0}(z_0)$  only at  $z_0$ , and  $\mathcal{O}(p_1)$  and  $\mathcal{O}(p_2)$  are at distance  $> 3\rho_0$  of  $z_0$ . Denote by  $C_0$  the constant  $C$  in the shadowing theorem 3.1. Reducing  $\rho_0$  if necessary, we can assume  $\rho_0 < \varepsilon_0$  where  $\varepsilon_0$  is given by Theorem 3.1. Let us also fix a large constant  $C_1$ , on which our construction will depend.

We truncate  $\gamma_-$  in positive time, stopping it at a fixed time where it is within distance  $\rho_0/(2C_0)$  of  $p_2$ , to get an orbit  $\gamma'_-$ . Let  $t_-$  be the largest time in  $(-\infty, -2C_1|\log \varepsilon|]$  where  $\gamma'_-(t)$  is within distance  $\varepsilon$  of  $p_1$ . As the orbit  $\gamma'_-$  converges exponentially quickly in negative time to  $\mathcal{O}(p_1)$  by hyperbolicity, one has  $d(\gamma'_-(t), \mathcal{O}(p_1)) \leq \varepsilon$  for  $t \leq -2C_1|\log \varepsilon|$ , if  $C_1$  is large enough. Hence, one needs to wait at most  $\ell_1$  before being  $\varepsilon$ -close to  $p_1$ . This shows that the time  $t_-$  satisfies  $t_- = -2C_1|\log \varepsilon| + O(1)$ .

In the same way, we truncate  $\gamma_+$  in negative time at a fixed time for which it is within distance  $\rho_0/(2C_0)$  of  $p_2$ , obtaining an orbit  $\gamma'_+$ . We denote by  $t_+$  the smallest time in  $[2C_1|\log \varepsilon|, +\infty)$  with  $d(\gamma'_+(t), p_2) \leq \varepsilon$ . It satisfies  $t_+ = 2C_1|\log \varepsilon| + O(1)$ .

As the flow is transitive, it has a dense orbit. Therefore, for any  $x, y$ , there exists an orbit  $\gamma_{x,y}$  starting from a point within distance  $\rho_0/(2C_0)$  of  $x$ , ending at a point within distance  $\rho_0/(2C_0)$  of  $y$ , and with length  $\in [1, T_0]$  where  $T_0$  is fixed and independent of  $x$  and  $y$ .

To any  $x$ , we associate an orbit as follows. Start with  $\gamma'_-$ , then follow  $\gamma_{p_2, \varphi_{-C_1|\log \varepsilon|x}}$ , then follow the orbit of  $x$  between times  $-C_1|\log \varepsilon|$  and  $C_1|\log \varepsilon|$ , then follow  $\gamma_{\varphi_{C_1|\log \varepsilon|x}, p_2}$ , then follow  $\gamma'_+$ . In this sequence, the distance between an endpoint of a piece and the starting point of the next one is always less than  $\rho_0/C_0$ . Hence, Theorem 3.1 applies and yields an infinite orbit  $\gamma'_x$ , that follows the above pieces of orbits up to  $\rho_0$ . If  $C_1$  is large enough, (3.2) implies that  $x$  is within distance at most  $\varepsilon$  of  $\gamma'_x$ . The inequality (3.3) shows that  $\gamma'_-(t_-)$  and the corresponding point  $x_-$  on  $\gamma'_x$  are within distance  $e^{-\theta t_-}$ . If  $C_1$  is large enough, this is bounded by  $\varepsilon$  since  $t_- = -2C_1|\log \varepsilon| + O(1)$ . Therefore,  $d(x_-, p_1) \leq 2\varepsilon$ . In the same way, the point  $x_+$  on  $\gamma'_x$  matching  $\gamma'_+(t_+)$  is within distance  $\varepsilon$  of  $\gamma'_+(t_+)$ , and therefore within distance  $2\varepsilon$  of  $p_1$ . Let us truncate  $\gamma'_x$  between  $x_-$  and  $x_+$ , to get an orbit segment  $\gamma_x$  of length  $6C_1|\log \varepsilon| + O(1)$ , starting and ending within  $2\varepsilon$  of  $p_1$ , and passing within  $\varepsilon$  of  $x$ .

Let  $\beta_d = 1/(3 \dim(\mathcal{M}))$ . We define a sequence of points of  $\mathcal{M}$  as follows. Let  $x_1$  be an arbitrary point for which the  $C(\kappa)$ -beginning of its orbit is  $\kappa/2$ -dense, to guarantee in the end that the last condition of the lemma is satisfied. If  $\gamma_{x_1}$  is not  $\varepsilon^{\beta_d}/2$ -dense, we choose another point  $x_2$  which is not in the  $\varepsilon^{\beta_d}/2$ -neighborhood of  $\gamma_{x_1}$ . Then  $\gamma_{x_1} \cup \gamma_{x_2}$  contain both  $x_1$  and  $x_2$  in their  $\varepsilon$ -neighborhood, and therefore in their  $\varepsilon^{\beta_d}/2$ -neighborhood. If  $\gamma_{x_1} \cup \gamma_{x_2}$  is still not  $\varepsilon^{\beta_d}/2$ -dense, then we add a third piece of orbit  $\gamma_{x_3}$ , and so on. By compactness, this process stops after finitely many steps, giving a finite sequence  $x_1, \dots, x_N$ .

As all  $\gamma_{x_i}$  start and end with  $p_1$  up to  $2\varepsilon$ , we can glue the sequence

$$\dots, \gamma_{x_N}, \gamma_{x_1}, \gamma_{x_2}, \dots, \gamma_{x_N}, \gamma_{x_1}, \dots$$

thanks to Theorem 3.1. We get a periodic orbit  $\gamma$  which shadows them up to  $2C_0\varepsilon$ . We claim this orbit satisfies the requirements of the lemma. We should check its length, its density, and its separation.

Let us start with the length. The points  $x_i$  are separated by at least  $\varepsilon^{\beta_d}/3$ . The balls of radius  $\varepsilon^{\beta_d}/6$  are disjoint, and each has a volume  $\geq c\varepsilon^{\beta_d \dim(\mathcal{M})} = c\varepsilon^{1/3}$ . We get that the number  $N$  of points  $x_i$  is bounded by  $C\varepsilon^{-1/3}$ . As each piece  $\gamma_{x_i}$  has length at most  $C|\log \varepsilon|$ , it follows that the total length of  $\gamma$  is bounded by  $C|\log \varepsilon|\varepsilon^{-1/3} \leq \varepsilon^{-1/2}$ .

Let us check the density. By construction, the union of the  $\gamma_{x_i}$  is  $\varepsilon^{\beta_d}/2$ -dense. As  $\gamma$  approximates each  $\gamma_{x_i}$  within  $2C_0\varepsilon$ , it follows that  $\gamma$  is  $2C_0\varepsilon + \varepsilon^{\beta_d}/2$  dense, and therefore  $\varepsilon^{\beta_d}$ -dense. In the statement of the lemma, we require the slightly stronger statement that if one removes a length 1 piece at the end of the orbit it remains  $\varepsilon^{\beta_d}$ -dense. Such a length 1 piece in  $\gamma_{x_N}$  consists of points that are within  $2\varepsilon$  of  $\mathcal{O}(p_1)$ . They are approximated within  $\varepsilon^{\beta_d}$  by the start and end of all the other  $\gamma_{x_i}$ .



Finally, let us check the more delicate separation, which has motivated the finer details of the construction as we will see now. Let  $\beta_s$  be suitably large. We want to show that any two points  $x, y$  of  $\gamma$  within distance  $\varepsilon^{\beta_s}$  are on the same local flow line. Since the expansion of the flow is at most exponential, for any  $t \leq 20C_1|\log \varepsilon|$ , we have  $d(\varphi_t x, \varphi_t y) \leq \varepsilon$  if  $\beta_s$  is large enough. In the piece of  $\gamma$  of length  $10C_1|\log \varepsilon|$  starting at  $x$ , there is an interval  $[t_1, t_2]$  of length  $4C_1|\log \varepsilon| + O(1)$  during which  $\varphi_t x$  is within distance at most  $\rho_0/2$  of  $\mathcal{O}(p_1)$ , corresponding to the junction between the orbits  $\gamma_{x_i}$  and  $\gamma_{x_{i+1}}$  where  $i$  is such that  $x$  belongs to the shadow of  $\gamma_{x_{i-1}}$ . For  $t \in [t_1, t_2]$ , one also has  $d(\varphi_t y, \mathcal{O}(p_1)) \leq \rho_0$  as the orbits follow each other up to  $\varepsilon$ . Note that in each  $\gamma_j$  the consecutive time spent close to  $\mathcal{O}(p_1)$  is bounded by  $2C_1|\log \varepsilon|$  as we have forced a passage close to  $p_2$  (and therefore far away from  $\mathcal{O}(p_1)$ ) after this time in the construction. It follows that also for  $y$  the time interval  $[t_1, t_2]$  has to correspond to a junction between two orbits  $\gamma_{x_j}$  and  $\gamma_{x_{j+1}}$ . Consider the smallest times  $t$  and  $t'$  after the junctions for which  $\varphi_t(x)$  and  $\varphi_{t'}(y)$  are  $2\rho_0$ -close to  $z_0$ . Since the orbit  $\gamma'_-$  meets  $W_{3\rho_0}(z_0)$  at the single point  $z_0$ , these times have to correspond to each other, i.e., the orbits are synchronized up to an error  $O(\varepsilon)$ . To conclude, it remains to show that  $i = j$ . Suppose by contradiction  $i < j$  for instance. The orbit of  $x$  follows  $\gamma_{x_i}$  up to  $2C_0\varepsilon$ , the orbit of  $y$  follows  $\gamma_{x_j}$  up to  $2C_0\varepsilon$ , and the orbits of  $x$  and  $y$  are within  $\varepsilon$  of each other. We deduce that  $\gamma_{x_i}$  and  $\gamma_{x_j}$  follow each other up to  $(4C_0 + 1)\varepsilon$ . Since  $x_j$  is within  $\varepsilon$  of  $\gamma_{x_j}$ , it follows that  $x_j$  is at within  $(4C_0 + 2)\varepsilon$  of  $\gamma_{x_i}$ . This is a contradiction with the construction, as we could have added the point  $x_j$  only if it was not in the  $\varepsilon^{\beta_s}$ -neighborhood of  $\gamma_{x_i}$ , and  $\varepsilon^{\beta_s} > (4C_0 + 2)\varepsilon$  if  $\varepsilon$  is small enough.  $\square$

**3.2.2. Construction of the approximate coboundary.** Let us now prove Theorem 1.2. The result is obvious if  $\varepsilon$  is bounded away from 0, by taking  $u = 0$  and  $h = f$ . Hence, we can assume that  $\varepsilon$  is small enough to apply Lemma 3.4, with  $\kappa = \varepsilon_0$ . On the orbit  $\mathcal{O}(x_0)$  given by this lemma, we define a function  $\tilde{u}$  by  $\tilde{u}(\varphi_t x_0) = \int_0^t f(\varphi_s x_0) ds$ . Note that it may not be continuous at  $x_0$ . As a consequence, we will rather denote by  $\mathcal{O}(x_0)$  the set  $(\varphi_t x_0)_{0 \leq t \leq T-1}$  (which satisfies the required properties of density and transversality) in order to avoid problems of discontinuity.

**Lemma 3.5.** *There exist  $\beta_1, C > 0$  independent of  $\varepsilon$  such that  $\|\tilde{u}\|_{C^{\beta_1}(\mathcal{O}(x_0))} \leq C$ .*

*Proof.* We first study the Hölder regularity of  $\tilde{u}$ , namely we want to control  $|\tilde{u}(x) - \tilde{u}(y)|$  by  $Cd(x, y)^{\beta_1}$  for some well-chosen exponent  $\beta_1$ , when  $d(x, y) \leq \varepsilon_0$  (where  $\varepsilon_0$  is the scale under which the shadowing theorem 3.1 holds). If  $x$  and  $y$  are on the same local flow line, then the result is obvious since  $f$  is bounded by 1, so we are left to prove that  $\tilde{u}$  is transversally  $C^{\beta_1}$ . Consider  $x = \varphi_{t_0} x_0 \in \mathcal{O}(x_0)$  and  $y = \varphi_{t_0+t} x_0 \in W_{\varepsilon_0}(x)$ . By transversal separation of  $\mathcal{O}(x_0)$ , these points satisfy  $d(x, y) \geq \varepsilon^{\beta_s}$ . We can close the segment  $[xy]$  i.e., we can find a periodic point  $p$  such that  $d(p, x) \leq Cd(x, y)$  with period  $t_p = t + \tau$ , where  $|\tau| \leq Cd(x, y)$



which shadows the segment. Then:

$$|\tilde{u}(y) - \tilde{u}(x)| \leq \underbrace{\left| \int_0^t f(\varphi_s x) ds - \int_0^{t_p} f(\varphi_s p) ds \right|}_{=(I)} + \underbrace{\left| \int_0^{t_p} f(\varphi_s p) ds \right|}_{=(II)}$$

The first term (I) is bounded by  $Cd(x, y)^{\beta'_1}$  for some  $\beta'_1 > 0$  depending on the dynamics, whereas the second term (II) is bounded — by assumption — by  $\varepsilon t_p$ . But  $\varepsilon t_p \lesssim \varepsilon t \lesssim \varepsilon T \lesssim \varepsilon^{1/2} \lesssim d(x, y)^{1/2\beta_s}$ . We thus obtain the sought result with  $\beta_1 := \min(\beta'_1, 1/2\beta_s)$ .

We now prove that  $\tilde{u}$  is bounded for the  $C^0$ -norm. We know that there exists a segment of the orbit  $\mathcal{O}(x_0)$  — call it  $S$  — of length  $\leq C$  which is  $\varepsilon_0$ -dense in  $\mathcal{M}$ . In particular, for any  $x \in \mathcal{O}(x_0)$ , there exists  $x_S \in S$  with  $d(x, x_S) \leq \varepsilon_0$ , and therefore  $|\tilde{u}(x) - \tilde{u}(x_S)| \leq Cd(x, x_S)^{\beta_1} \leq C\varepsilon_0^{\beta_1}$  thanks to the Hölder control of the previous paragraph. Using the same argument with  $x_0$ , we get as  $\tilde{u}(x_0) = 0$

$$|\tilde{u}(x)| = |\tilde{u}(x) - \tilde{u}(x_0)| \leq |\tilde{u}(x) - \tilde{u}(x_S)| + |\tilde{u}(x_S) - \tilde{u}((x_0)_S)| + |\tilde{u}(x_0) - \tilde{u}((x_0)_S)|.$$

The first and last term are bounded by  $C\varepsilon_0^{\beta_1}$ , and the middle one is bounded by  $C$  as  $S$  has a bounded length and  $\|f\|_{C^0} \leq 1$ .  $\square$

For each  $i$ , we extend the function  $\tilde{u}$  (defined on  $\mathcal{O}(x_0)$ ) to a Hölder function  $u_i$  on  $\Sigma_i$ , by the formula  $u_i(x) = \sup \tilde{u}(y) - \|\tilde{u}\|_{C^{\beta_1}(\mathcal{O}(x_0))} d(x, y)^{\beta_1}$ , where the supremum is taken over all  $y \in \mathcal{O}(x_0)$ . With this formula, it is classical that the extension is Hölder continuous, with  $\|u_i\|_{C^{\beta_1}(\Sigma_i)} \leq \|\tilde{u}\|_{C^{\beta_1}(\mathcal{O}(x_0))}$ . We then push the function  $u_i$  by the flow in order to define it on  $\mathcal{U}_i$  by setting for  $x \in \Sigma_i, \varphi_t x \in \mathcal{U}_i$ :  $u_i(\varphi_t x) = u_i(x) + \int_0^t f(\varphi_s x) ds$ . Note that by Lemma 3.3, the extension is still Hölder with the same regularity. We now set  $u := \sum_i u_i \theta_i$  and  $h := f - Xu = -\sum_i u_i X\theta_i$ . The functions  $X\theta_i$  are uniformly bounded in  $C^\infty$ , independently of  $\varepsilon$  so the functions  $u_i X\theta_i$  are in  $C^{\beta_1}$  with a Hölder norm independent of  $\varepsilon > 0$  and thus  $\|h\|_{C^{\beta_1}} \leq C$ .

**Lemma 3.6.**  $\|h\|_{C^{\beta_1/2}} \leq \varepsilon^{\beta_3/2}$

*Proof.* We claim that  $h$  vanishes on  $\mathcal{O}(x_0)$ : indeed, on  $\mathcal{U}_i \cap \mathcal{O}(x_0)$  one has  $u_i \equiv \tilde{u}$  and thus  $h = -\tilde{u} \sum_i X\theta_i = -\tilde{u} X \sum_i \theta_i = -\tilde{u} X \mathbf{1} = 0$ . Since  $\mathcal{O}(x_0)$  is  $\varepsilon^{\beta_d}$ -dense and  $\|h\|_{C^{\beta_1}} \leq C$ , we get that  $\|h\|_{C^0} \leq C\varepsilon^{\beta_1\beta_d} = C\varepsilon^{\beta_3}$ , where  $\beta_3 = \beta_1\beta_d$ . By interpolation, we eventually obtain that  $\|h\|_{C^{\beta_1/2}} \leq \varepsilon^{\beta_3/2}$ .  $\square$

*Proof of Theorem 1.2.* The previous lemma provides the sought estimate on the remainder  $h$ . This completes the proof of Theorem 1.2.  $\square$

#### 4. GENERALIZED GEODESIC X-RAY TRANSFORM

From now on, we will rather use the dual decomposition of the cotangent space  $T^*\mathcal{M} = E_0^* \oplus E_u^* \oplus E_s^*$ , where  $E_0^*(E_u \oplus E_s) = 0$ ,  $E_s^*(E_s \oplus \mathbb{R}X) = 0$ ,  $E_u^*(E_u \oplus \mathbb{R}X) = 0$ . If  $A^{-\top}$

denotes the inverse transpose of a linear operator  $A$ , then the dual estimates to (1.2) are:

$$\begin{aligned} |d\varphi_t^{-\top}(x) \cdot \xi|_{\varphi_t(x)} &\leq C e^{-\lambda t} |\xi|_x, \quad \forall t > 0, \xi \in E_s^*(x) \\ |d\varphi_t(x) \cdot \xi|_{\varphi_t(x)} &\leq C e^{-\lambda|t|} |\xi|_x, \quad \forall t < 0, \xi \in E_u^*(x), \end{aligned} \quad (4.1)$$

where  $|\cdot|_x$  is now  $g^{-1}$ , the dual metric to  $g$  (which makes the musical isomorphism  $\flat : T\mathcal{M} \rightarrow T^*\mathcal{M}$  an isometry). For the sake of simplicity, we now assume that  $X$  generates a contact Anosov flow; the results of this paragraph will be applied to the case of an Anosov geodesic flow. It would actually be sufficient to assume that the flow is Anosov, preserves a smooth measure and that it is mixing for this measure. Note that a contact Anosov flow is exponentially mixing by [Liv04]. We will denote by  $\mu$  the *normalized* volume form induced by the contact 1-form. In the case of a geodesic flow,  $\mu$  is nothing but the Liouville volume form. By  $L^2(\mathcal{M})$ , we will always refer to the space  $L^2(\mathcal{M}, d\mu)$ . The orthogonal projection on the constant function is denoted by  $\mathbf{1} \otimes \mathbf{1}$ .

**4.1. Resolvent of the flow at 0.** By [FS11], we know that the resolvents  $R_{\pm}(\lambda) := (X \pm \lambda)^{-1} : \mathcal{H}_{\pm}^s \rightarrow \mathcal{H}_{\pm}^s$  (initially defined for  $\Re(\lambda) > 0$ ) admit a meromorphic extension to the half-space  $\{\Re(\lambda) > -cs\}$  — where  $\mathcal{H}_{\pm}^s$  are *anisotropic Sobolev spaces* — and thus  $R_{\pm}(\lambda) : C^{\infty}(\mathcal{M}) \rightarrow \mathcal{D}'(\mathcal{M})$  admit a meromorphic extension to the whole complex plane. For  $\Re(\lambda) > 0$ ,  $R_{\pm}(\lambda) : L^2(\mathcal{M}) \rightarrow L^2(\mathcal{M})$  are bounded and the expression  $R_+(\lambda)$  is given by

$$R_+(\lambda) = (X + \lambda)^{-1} = \int_0^{+\infty} e^{-\lambda t} e^{-tX} dt, \quad (4.2)$$

where  $e^{-tX} f(x) = f(\varphi_{-t}(x))$  for  $f \in C^{\infty}(\mathcal{M}), x \in \mathcal{M}$ .

In a neighborhood of 0, we can thus write the Laurent expansions

$$R_+(\lambda) = R_0^+ + \frac{\mathbf{1} \otimes \mathbf{1}}{\lambda} + \mathcal{O}(\lambda), \quad R_-(\lambda) = R_0^- - \frac{\mathbf{1} \otimes \mathbf{1}}{\lambda} + \mathcal{O}(\lambda), \quad (4.3)$$

where  $R_0^+ : \mathcal{H}_+^s \rightarrow \mathcal{H}_+^s, R_0^- : \mathcal{H}_-^s \rightarrow \mathcal{H}_-^s$  are bounded. Since  $H^s \subset \mathcal{H}_{\pm}^s \subset H^{-s}$ , we obtain that  $R_0^{\pm} : H^s \rightarrow H^{-s}$  are bounded and thus  $(R_0^+)^* : H^s \rightarrow H^{-s}$  is bounded too. Moreover, it is easy to check that formally  $(R_0^+)^* = -R_0^-$  (i.e., the operators coincide on  $C^{\infty}(\mathcal{M})$ ), in the sense that for all  $f_1, f_2 \in C^{\infty}(\mathcal{M})$ ,  $\langle R_0^- f_1, f_2 \rangle_{L^2(\mathcal{M})} = \langle f_1, -R_0^+ f_2 \rangle_{L^2(\mathcal{M})}$ . Since  $C^{\infty}(\mathcal{M})$  is dense in  $H^s(\mathcal{M})$ , we obtain that  $(R_0^+)^* = -R_0^-$  on  $H^s(\mathcal{M})$ , in the sense that for all  $f_1, f_2 \in H^s(\mathcal{M})$ ,  $\langle R_0^- f_1, f_2 \rangle_{L^2(\mathcal{M})} = \langle f_1, -R_0^+ f_2 \rangle_{L^2(\mathcal{M})}$ .

Also remark that, as operators  $C^{\infty}(\mathcal{M}) \rightarrow \mathcal{D}'(\mathcal{M})$ , one has:

$$XR_0^+ = R_0^+X = 1 - \mathbf{1} \otimes \mathbf{1}, \quad XR_0^- = R_0^-X = 1 - \mathbf{1} \otimes \mathbf{1} \quad (4.4)$$

For the sake of simplicity, we will write  $R_0 := R_0^+$ . We introduce the operator

$$\Pi := R_0 + R_0^*, \quad (4.5)$$

the sum of the two holomorphic parts of the resolvent. An easy computation, using (4.3), proves that  $\Pi(\mathbf{1}) = 0$  and the image  $\Pi(C^\infty(\mathcal{M}))$  is orthogonal to the constants. We recall the

**Theorem 4.1.** [Gui17, Theorem 1.1] *For all  $s > 0$ , the operator  $\Pi : H^s(\mathcal{M}) \rightarrow H^{-s}(\mathcal{M})$  is bounded, selfadjoint and satisfies:*

- (1)  $\forall f \in H^s(\mathcal{M}), X\Pi f = 0,$
- (2)  $\forall f \in H^s(\mathcal{M})$  such that  $Xf \in H^s(\mathcal{M}), \Pi Xf = 0.$

*If  $f \in H^s(\mathcal{M})$  with  $\langle f, \mathbf{1} \rangle_{L^2} = 0$ , then  $f \in \ker \Pi$  if and only if there exists a solution  $u \in H^s(\mathcal{M})$  to the cohomological equation  $Xu = f$ , and  $u$  is unique modulo constants.*

There exists two other characterizations of the operator  $\Pi$  that are more tractable and which we detail in the next proposition. We set  $\Pi_\lambda := \mathbf{1}_{(-\infty, \lambda]}(-iX).$

**Proposition 4.2.** *For  $f_1, f_2 \in C^\infty(\mathcal{M})$  such that  $\langle f, \mathbf{1} \rangle_{L^2} = 0$ :*

- (1)  $\langle \Pi f_1, f_2 \rangle = 2\pi \partial_\lambda|_{\lambda=0} \langle \Pi_\lambda f_1, f_2 \rangle,$
- (2)  $\langle \Pi f_1, f_2 \rangle = \int_{-\infty}^{+\infty} \langle f_1 \circ \varphi_t, f_2 \rangle dt.$

*Proof.* (1) For  $f_1, f_2 \in C^\infty(\mathcal{M})$  such that  $\int_{\mathcal{M}} f_i d\mu = 0$ , we have using Stone's formula, for  $\delta > 0$ :

$$\begin{aligned} \langle \Pi_{\lambda+\delta} f_1, f_2 \rangle - \langle \Pi_{\lambda-\delta} f_1, f_2 \rangle &= \langle \mathbf{1}_{[\lambda-\delta, \lambda+\delta]} f_1, f_2 \rangle \\ &= \frac{1}{2\pi} \int_{\lambda-\delta}^{\lambda+\delta} \langle (R_+(-i\lambda) - R_-(i\lambda)) f_1, f_2 \rangle d\lambda \end{aligned}$$

Dividing by  $2\delta$  and passing to the limit  $\delta \rightarrow 0^+$ , we obtain  $\partial_\lambda|_{\lambda=0} \langle \Pi_\lambda f_1, f_2 \rangle = \frac{1}{2\pi} \langle (R_0^+ - R_0^-) f_1, f_2 \rangle = \frac{1}{2\pi} \langle \Pi f_1, f_2 \rangle.$

(2) Thanks to the exponential decay of correlations (see [Liv04]), one can apply Lebesgue's dominated convergence theorem in the limit  $\lambda \rightarrow 0^+$  in the following expression

$$\langle \Pi f_1, f_2 \rangle = \lim_{\lambda \rightarrow 0^+} \int_{-\infty}^{+\infty} e^{-\lambda|t|} \langle f_1 \circ \varphi_{-t}, f_2 \rangle dt,$$

and the result is then immediate. Note that a polynomial decay would have been sufficient.  $\square$

The quantity  $\langle \Pi f, f \rangle$  is sometimes referred to in the literature as the *variance* of the flow. In particular, it enjoys the following positivity property:

**Lemma 4.3.** *The operator  $\Pi : H^s(\mathcal{M}) \rightarrow H^{-s}(\mathcal{M})$  is positive in the sense of quadratic forms, namely  $\langle \Pi f, f \rangle_{L^2} \geq 0$  for all real-valued  $f \in H^s(\mathcal{M})$ .*

There are different ways of proving this lemma, related to the different characterizations of the operator  $\Pi$ . We only detail one of them which is in the dynamical spirit of this article. Another way could be to use the first item of Proposition 4.2 and the fact that the spectral measure  $\Pi_\lambda$  is non-decreasing.

*Proof.* By density, it is sufficient to prove the lemma for a real-valued  $f \in C^\infty(\mathcal{M})$ . We will actually show that for  $\lambda > 0$ :

$$\left\langle \left( R_+(\lambda) - \frac{\mathbf{1} \otimes \mathbf{1}}{\lambda} \right) f, f \right\rangle = \langle R_+(\lambda) f, f \rangle - \frac{1}{\lambda} \left( \int_{\mathcal{M}} f d\mu \right)^2 \geq 0$$

The same arguments being valid for  $R_-(\lambda)$ , we will deduce the result by taking the limit  $\lambda \rightarrow 0^+$ . By Parry' formula [Par88, Paragraph 3], we know that:

$$\langle R_+(\lambda) f, f \rangle = \lim_{T \rightarrow \infty} \frac{1}{N(T)} \sum_{\ell(\gamma) \leq T} e^{\int_\gamma J^u} \frac{1}{\ell(\gamma)} \int_0^{\ell(\gamma)} R_+(\lambda) f(\varphi_t z) f(\varphi_t z) dt, \quad (4.6)$$

where  $\gamma$  is a periodic orbit,  $z \in \gamma$ ,  $\ell(\gamma)$  is the length of  $\gamma$  and  $N(T) = \sum_{\ell(\gamma) \leq T} e^{\int_\gamma J^u}$  is a normalizing coefficient,

$$J^u : x \mapsto \partial_t \det d\varphi_t(x)|_{E_u(x)}|_{t=0}$$

is the unstable Jacobian (or the *geometric potential*). Let us fix a closed orbit  $\gamma$  and a base point  $z \in \gamma$ . We set  $\tilde{f}(t) := f(\varphi_t z)$  which we see as a smooth function,  $\ell$ -periodic on  $\mathbb{R}$  (with  $\ell := \ell(\gamma)$ ). Since  $R_+(\lambda)$  commutes with  $X$ ,  $R_+(\lambda)$  acts as a Fourier multiplier on functions defined on  $\gamma$ . As a consequence, if we decompose  $\tilde{f}(t) = \sum_{n \in \mathbb{Z}} c_n e^{2i\pi n t / \ell}$ , we have:

$$\begin{aligned} R_+(\lambda) \tilde{f}(t) &= \int_0^{+\infty} e^{-\lambda s} \tilde{f}(t+s) ds \\ &= \sum_{n \in \mathbb{Z}} c_n e^{2i\pi n t / \ell} \int_0^{+\infty} e^{-(\lambda - 2i\pi n / \ell) s} ds \\ &= \sum_{n \in \mathbb{Z}} \frac{c_n (\lambda + 2i\pi n / \ell)}{\lambda^2 + 4\pi^2 n^2 / \ell^2} e^{2i\pi n t / \ell} \end{aligned}$$

Then:

$$\langle R_+(\lambda) \tilde{f}, \tilde{f} \rangle_{L^2} = \frac{1}{\ell} \int_0^\ell R_+(\lambda) \tilde{f}(t) \tilde{f}(t) dt = \sum_{n \in \mathbb{Z}} \frac{|c_n|^2 (\lambda + 2i\pi n / \ell)}{\lambda^2 + 4\pi^2 n^2 / \ell^2} = \lambda \sum_{n \in \mathbb{Z}} \frac{|c_n|^2}{\lambda^2 + 4\pi^2 n^2 / \ell^2},$$

by oddness of the imaginary part of the sum. In particular:

$$\frac{1}{\ell} \int_0^\ell R_+(\lambda) \tilde{f}(t) \tilde{f}(t) dt \geq \frac{|c_0|^2}{\lambda} = \frac{1}{\lambda} \left( \frac{1}{\ell} \int_0^\ell \tilde{f}(t) dt \right)^2 \quad (4.7)$$

Inserting (4.7) into (4.6), then applying Jensen's convexity inequality:

$$\begin{aligned} \langle R_+(\lambda)f, f \rangle &\geq \lambda^{-1} \lim_{T \rightarrow \infty} \frac{1}{N(T)} \sum_{\ell(\gamma) \leq T} e^{\int_\gamma J^u} \left( \frac{1}{\ell(\gamma)} \int_0^{\ell(\gamma)} f(\varphi_t z) dt \right)^2 \\ &\geq \lambda^{-1} \lim_{T \rightarrow \infty} \left( \frac{1}{N(T)} \sum_{\ell(\gamma) \leq T} e^{\int_\gamma J^u} \frac{1}{\ell(\gamma)} \int_0^{\ell(\gamma)} f(\varphi_t z) dt \right)^2 = \frac{1}{\lambda} \left( \int_{SM} f d\mu \right)^2, \end{aligned}$$

where we used again Parry's formula in the last equality.  $\square$

**4.2. The normal operator.** We introduce

$$\Pi_m := \pi_{m*}(\Pi + \mathbf{1} \otimes \mathbf{1})\pi_m^* \quad (4.8)$$

Recall from §2.1.3 that given  $(x, \xi) \in T^*M$ , the space  $\otimes_S^m T_x^*M$  decomposes as the direct sum

$$\begin{aligned} \otimes_S^m T_x^*M &= \text{ran} \left( \sigma_D(x, \xi)|_{\otimes_S^{m-1} T_x^*M} \right) \oplus \ker \left( \sigma_{D^*}(x, \xi)|_{\otimes_S^m T_x^*M} \right) \\ &= \text{ran} \left( \sigma_{j\xi}|_{\otimes_S^{m-1} T_x^*M} \right) \oplus \ker \left( i_\xi|_{\otimes_S^m T_x^*M} \right) \end{aligned}$$

The projection on the right space parallel to the left space is denoted by  $\pi_{\ker i_\xi}$  and  $\text{Op}(\pi_{\ker i_\xi}) = \pi_{\ker D^*} + S$  by Lemma 2.6, where  $S \in \Psi^{-1}$  and  $\text{Op}$  is any quantization on  $M$  (see [Shu01, Section 6.4] for instance). Here,  $\Psi^m$  denotes the set of pseudodifferential operators of order  $m \in \mathbb{R}$ . Given  $P \in \Psi^m$ , we will denote by  $\sigma_m$  its principal symbol. The following structure theorem is crucial in the sequel. It can be seen as a more intrinsic version of [SSU05, Theorem 2.1].

**Theorem 4.4.**  $\Pi_m$  is a pseudodifferential operator of order  $-1$  with principal symbol

$$\sigma_m := \sigma_{\Pi_m} : (x, \xi) \mapsto \frac{2\pi}{C_{n,m}} |\xi|^{-1} \pi_{\ker i_\xi} \pi_{m*} \pi_m^* \pi_{\ker i_\xi},$$

with  $C_{n,m} = \int_0^\pi \sin^{n-1+2m}(\varphi) d\varphi$ .

*Proof.* The fact that  $\Pi_m$  is pseudodifferential was proved in [Gui17]. All is left to compute is the principal symbol of  $\Pi_m$ . According to the proof in [Gui17, Theorem 3.1], we can only consider the integral in time between  $(-\varepsilon, \varepsilon)$ . Namely, given  $\chi \in C_c^\infty(\mathbb{R})$  a smooth cutoff function around 0 whose support is contained in  $(-\varepsilon, \varepsilon)$ , one has:

$$\begin{aligned} \Pi_m &= \pi_{m*} \int_{-\varepsilon}^{\varepsilon} \chi(t) e^{-tX} dt \pi_m^* \\ &\quad - \pi_{m*} R_0^+ \int_0^{+\infty} \chi'(t) e^{-tX} dt \pi_m^* - \pi_{m*} R_0^- \int_{-\infty}^0 \chi'(t) e^{-tX} dt \pi_m^* \\ &\quad + \left( 1 - \int_{-\infty}^{+\infty} \chi(t) dt \right) \pi_{m*} \mathbf{1} \otimes \mathbf{1} \pi_m^* \end{aligned}$$

On the right-hand side, the last term is obviously smoothing. Following the same computations as in [Gui17, Theorem 3.1], one can prove that the second and the third terms are also smoothing (this stems from an argument on the wavefront set of the kernel of these operators, using the fact that there are no conjugate points in the manifold). Thus, the pseudodifferential behaviour of the operator  $\Pi_m$  is encapsulated by the first term whose kernel has a support living in a neighborhood of the diagonal in  $M \times M$ . In the following,  $\varepsilon > 0$  is chosen small enough (less than the injectivity radius at the point  $x$ ).

Let us consider a smooth section  $f_1 \in C^\infty(M, \otimes_S^m T^*M)$  defined in a neighborhood of  $x \in M$  and  $f_2 \in \otimes_S^m T_x^*M$ , then:

$$\begin{aligned} \langle \sigma_m(x_0, \xi) f_1, f_2 \rangle_{x_0} &= \lim_{h \rightarrow 0} h^{-1} e^{-iS(x_0)/h} \langle \Pi_m(e^{iS(x)/h} f_1), f_2 \rangle_{x_0} \\ &= \lim_{h \rightarrow 0} h^{-1} e^{-iS(x_0)/h} \langle \Pi \pi_m^*(e^{iS(x)/h} f_1), \pi_m^* f_2 \rangle_{L^2(S_{x_0}M)}, \end{aligned}$$

where  $\xi = dS(x) \neq 0$ . Here, it is assumed that  $\text{Hess}_x S$  is non-degenerate. We obtain:

$$\begin{aligned} &\langle \sigma_m(x, \xi) f_1, f_2 \rangle_{x_0} \\ &= \lim_{h \rightarrow 0} h^{-1} \int_{\mathbb{S}^n} \int_{-\varepsilon}^{+\varepsilon} e^{i/h(S(\gamma(t)) - S(x))} \pi_m^* f_1(\gamma(t), \dot{\gamma}(t)) \pi_m^* f_2(x_0, v) \chi(t) dt dv \\ &= \lim_{h \rightarrow 0} h^{-1} \int_{\mathbb{S}^{n-1}} \left( \int_0^\pi \int_{-\varepsilon}^{+\varepsilon} e^{i/h(S(\gamma(t)) - S(x))} \pi_m^* f_1(\gamma(t), \dot{\gamma}(t)) \pi_m^* f_2(x_0, v) \sin^{n-1}(\varphi) \chi(t) dt d\varphi \right) du \end{aligned}$$

where  $\chi$  is a cutoff function with support in  $(-\varepsilon, \varepsilon)$ ,  $\gamma$  is the geodesic such that  $\gamma(0) = x$ ,  $\dot{\gamma}(0) = v$  and we have decomposed  $v = \cos(\varphi)n + \sin(\varphi)u$  with  $n = \xi^\# / |\xi| = dS(x)^\# / |dS(x)|$ ,  $u \in \mathbb{S}^{n-1}$ . We apply the stationary phase lemma [Zwo12, Theorem 3.13] uniformly in the  $u \in \mathbb{S}^{n-1}$  variable. For fixed  $u$ , the phase is  $\Phi : (t, \varphi) \mapsto S(\gamma(t)) - S(x)$  so  $\partial_t \Phi(t, \varphi) = dS(\dot{\gamma}(t))$ . More generally if  $\tilde{\Phi} : (t, v) \mapsto S(\gamma(t)) - S(x)$  denotes the map defined for any  $v \in \mathbb{S}^n$ , then

$$\partial_v \tilde{\Phi}(t, v) \cdot V = d\pi(d\varphi_t(x, v) \cdot V), \quad \forall V \in \mathbb{V},$$

where  $\mathbb{V} = \ker d\pi_0$ , with  $\pi_0 : SM \rightarrow M$  the natural projection. Since  $(M, g)$  has no conjugate points,  $d\pi(d\varphi_t(x, v)) \cdot V \neq 0$  as long as  $t \neq 0$  and  $V \in \mathbb{V} \setminus \{0\}$ . And  $dS(\dot{\gamma}(0)) = dS(\cos(\varphi)n + \sin(\varphi)u) = \cos(\varphi)|dS(x)| = 0$  if and only if  $\varphi = \pi/2$ . So the only critical point of  $\Phi$  is  $(t = 0, \varphi = \pi/2)$ . Let us also remark that

$$\text{Hess}_{(0, \pi/2)} \Phi = \begin{pmatrix} \text{Hess}_x S(u, u) & -|dS(x)| \\ -|dS(x)| & 0 \end{pmatrix}$$

is non-degenerate with determinant  $-|\xi|^2$ , so the stationary phase lemma can be applied and we get:

$$\begin{aligned} & \int_0^\pi \int_{-\varepsilon}^{+\varepsilon} e^{i/h(S(\gamma(t))-S(x_0))} \pi_m^* f_1(\gamma(t), \dot{\gamma}(t)) \pi_m^* f_2(x_0, v) \sin^{n-1}(\varphi) dt d\varphi \\ & \sim_{h \rightarrow 0} 2\pi h |\xi|^{-1} \pi_m^* f_1(x_0, u) \pi_m^* f_2(x_0, u). \end{aligned}$$

Eventually, we obtain:

$$\langle \sigma_m(x, \xi) f_1, f_2 \rangle_{x_0} = \frac{2\pi}{|\xi|} \int_{\{(\xi, v)=0\}} \pi_m^* f_1(v) \pi_m^* f_2(v) dS_\xi(v),$$

where  $dS_\xi$  is the canonical measure induced on the  $n-1$ -dimensional sphere  $\mathbb{S}_x M \cap \{(\xi, v) = 0\}$ . The sought result then follows from Lemma 2.1.  $\square$

### 4.3. Ellipticity, injectivity on solenoidal tensors.

**Lemma 4.5.** *The operator  $\Pi_m$  is elliptic on solenoidal tensors, that is there exists pseudodifferential operators  $Q$  and  $R$  of respective order 1 and  $-\infty$  such that:*

$$Q\Pi_m = \pi_{\ker D^*} + R$$

*Proof.* We define

$$\tilde{q}(x, \xi) = \begin{cases} 0, & \text{on } \text{ran}(\sigma j_\xi) \\ |\xi|(\pi_{m^*} \pi_m^*)^{-1}, & \text{on } \ker(i_\xi) \end{cases}$$

and  $q(x, \xi) = (1 - \chi(x, \xi))\tilde{q}(x, \xi)$  for some cutoff function  $\chi \in C_c^\infty(T^*M)$  around the zero section. By construction,  $\text{Op}(q)\Pi_m = \pi_{\ker D^*} - R'$  with  $R' \in \Psi^{-1}$ . Let  $r' = \sigma_{R'}$  and define  $a \sim \sum_{k=0}^\infty r'^k$ . Then  $\text{Op}(a)$  is a microlocal inverse for  $1 - R'$  that is  $\text{Op}(a)(1 - R') \in \Psi^{-\infty}$ . Since  $R'D = 0$ , we obtain that  $R' = R'\pi_{\ker D^*}$  and thus

$$\underbrace{\text{Op}(a) \text{Op}(q)}_{=Q} \Pi_m = \text{Op}(a)(1 - R')\pi_{\ker D^*} = \pi_{\ker D^*} + R,$$

where  $R$  is a smoothing operator.  $\square$

From now on, we assume that the X-ray transform is injective on solenoidal tensors.

**Lemma 4.6.** *If  $I_m$  is solenoidal injective, then  $\Pi_m$  is injective on  $H_{\text{sol}}^s(M, \otimes_S^m T^*M)$ , for all  $s \in \mathbb{R}$ .*

*Proof.* We fix  $s \in \mathbb{R}$ . We assume that  $\Pi_m f = 0$  for some  $f \in H_{\text{sol}}^s(M, \otimes_S^m T^*M)$ . By ellipticity of the operator, we get that  $f \in C_{\text{sol}}^\infty(M, \otimes_S^m T^*M)$ . And:

$$\begin{aligned} \langle \Pi_m f, f \rangle_{L^2} &= \langle \Pi \pi_m^* f, \pi_m^* f \rangle_{L^2} + \left( \int_{SM} \pi_m^* f d\mu \right)^2 \\ &= \langle (1 + \Delta_m)^{-s} \Pi \pi_m^* f, \pi_m^* f \rangle_{H^s} + \left( \int_{SM} \pi_m^* f d\mu \right)^2 = 0. \end{aligned}$$



Here, the Laplacian  $\Delta_m$  is the one introduced in §2.1.3. The scalar product on  $H^s$  is  $\langle f, h \rangle_{H^s} := \langle (1 + \Delta_m)^{s/2} f, (1 + \Delta_m)^{s/2} h \rangle_{L^2}$ . By Lemma 4.3, since  $\langle \Pi \pi_m^* f, \pi_m^* f \rangle \geq 0$ , we obtain that  $\int_{SM} \pi_m^* f d\mu = 0$ . Moreover,  $(1 + \Delta_m)^{-s} \Pi$  is bounded and positive (hence selfadjoint) on  $H^s$  so there exists a square root  $R : H^s \rightarrow H^s$ , that is a bounded positive operator satisfying  $(1 + \Delta_m)^{-s} \Pi = R^* R$ , where  $R^*$  is the adjoint on  $H^s$ . Then:

$$\langle (1 + \Delta_m)^{-s} \Pi \pi_m^* f, \pi_m^* f \rangle_{H^s} = 0 = \|R \pi_m^* f\|_{H^s}^2$$

This yields  $(1 + \Delta_m)^{-s} \Pi \pi_m^* f = 0$  so  $\Pi \pi_m^* f = 0$ . By Theorem 4.1, there exists  $u \in C^\infty(SM)$  such that  $\pi_m^* f = Xu$  so  $f \in \ker I_m \cap \ker D^*$ . By  $s$ -injectivity of the X-ray transform, we get  $f \equiv 0$ .  $\square$

A direct consequence of Lemma 4.6 and Theorem 4.5 is the

**Theorem 4.7.** *If  $I_m$  is solenoidal injective, then there exists a pseudodifferential operator  $Q'$  of order 1 such that:  $Q' \Pi_m = \pi_{\ker D^*}$ .*

*Proof.* The operator  $\Pi_m$  is elliptic of order  $-1$  on  $\ker D^*$ , thus Fredholm as an operator  $H_{\text{sol}}^s(M, \otimes_S^m T^* M) \rightarrow H_{\text{sol}}^{s+1}(M, \otimes_S^m T^* M)$  for all  $s \in \mathbb{R}$ . It is selfadjoint on  $H_{\text{sol}}^{-1/2}(M, \otimes_S^m T^* M)$ , thus Fredholm of index 0 (the index being independent of the Sobolev space considered, see [Shu01, Theorem 8.1]), and injective, thus invertible on  $H_{\text{sol}}^s(M, \otimes_S^m T^* M)$ . We multiply the equality  $Q \Pi_m = \pi_{\ker D^*} + R$  on the right by  $Q' := \pi_{\ker D^*} \Pi_m^{-1} \pi_{\ker D^*}$ :

$$Q \Pi_m Q' = Q \underbrace{\Pi_m \pi_{\ker D^*}}_{=\Pi_m} \Pi_m^{-1} \pi_{\ker D^*} = Q \pi_{\ker D^*} = Q' + R Q'$$

As a consequence,  $Q' = Q \pi_{\ker D^*} + \text{smoothing}$  so it is a pseudodifferential operator of order 1. And  $Q' \Pi_m = \pi_{\ker D^*}$ .  $\square$

This yields the following stability estimate:

**Lemma 4.8.** *For all  $s \in \mathbb{R}$ , there exists a constant  $C := C(s) > 0$  such that:*

$$\forall f \in H_{\text{sol}}^s(M, \otimes_S^m T^* M), \quad \|f\|_{H^s} \leq C \|\Pi_m f\|_{H^{s+1}}$$

Eventually, we will need a last lemma, whose proof can be found in [Gui17].

**Lemma 4.9.**  *$\Pi \pi_m^* : H^{-s}(M, \otimes_S^m T^* M) \rightarrow H^{-s}(SM)$  is bounded, for any  $s > 0$ . By duality,  $\pi_{m*} \Pi : H^s(SM) \rightarrow H^s(M, \otimes_S^m T^* M)$  is bounded too, for any  $s > 0$ .*

**4.4. Stability estimates for the X-ray transform.** We can now deduce from the previous work a stability estimate for the geodesic X-ray transform.

**Theorem 4.10.** *For all  $\frac{n+1}{2} < s < r$ , there exists an exponent  $\theta := \theta(s, r) > 0$  and a constant  $C := C(s, r) > 0$  such that:*

$$\forall f \in H_{\text{sol}}^r(M, \otimes_S^m T^* M) \text{ with } \|f\|_{H^r} \leq 1, \quad \|f\|_{H^s} \leq C \|I_m f\|_{\ell^\infty}^\theta$$

*Proof.* We fix such constants  $\frac{n+1}{2} < s < r$  and consider  $f \in H_{\text{sol}}^r(M, \otimes_S^m T^*M)$  such that  $\|f\|_{H^r} \leq 1$ . There exists an exponent  $0 < \alpha < 1$  such that  $f \in C^\alpha(M, \otimes_S^m T^*M)$ . We write  $\pi_m^* f + \|I_m f\|_{\ell^\infty} = Xu + h$ , where the decomposition is given by Theorem 1.1, with  $h \geq 0$ ,  $\|h\|_{C^\beta} \lesssim \|f\|_{C^\alpha} + \|I_m f\|_{\ell^\infty} \lesssim 1$  since  $\|f\|_{C^\alpha} \lesssim 1$  ( $\beta$  is provided by the theorem applied with a  $C^\alpha$ -function). Also remark that:

$$\begin{aligned} \|h\|_{L^1} &= \int_{SM} h \, d\mu = \int_{SM} (\pi_m^* f + \|I_m f\|_{\ell^\infty} - Xu) \, d\mu \\ &= \int_{SM} (\pi_m^* f + \|I_m f\|_{\ell^\infty}) \, d\mu \lesssim \|I_m f\|_{\ell^\infty}, \end{aligned} \quad (4.9)$$

where the estimate  $|\int_{SM} \pi_m^* f \, d\mu| \lesssim \|I_m f\|_{\ell^\infty}$  follows from Parry's formula (4.6). We have:

$$\begin{aligned} \|f\|_{H^s} &\lesssim \|\Pi_m f\|_{H^{s+1}}, && \text{by Lemma 4.8} \\ &\lesssim \|\Pi_m f\|_{H^{\beta/3}}^{\theta'} \underbrace{\|\Pi_m f\|_{H^{r+1}}^{1-\theta'}}_{\lesssim \|f\|_{H^r} \lesssim 1}, && \text{by interpolation with } \theta' = \frac{r-s}{r+1-\beta/3} \\ &\lesssim \|\pi_{m*} \Pi(\pi_m^* f + \|I_m f\|_{\ell^\infty})\|_{H^{\beta/3}}^{\theta'} + \|I_m f\|_{\ell^\infty}^{\theta'} \\ &\lesssim \|\pi_{m*} \Pi h\|_{H^{\beta/3}}^{\theta'} + \|I_m f\|_{\ell^\infty}^{\theta'}, && \text{by Theorem 4.1} \\ &\lesssim \|h\|_{H^{\beta/3}}^{\theta'} + \|I_m f\|_{\ell^\infty}^{\theta'}, && \text{by Lemma 4.9} \\ &\lesssim \|h\|_{L^2}^{\theta'/2} \|h\|_{H^{2\beta/3}}^{\theta'/2} + \|I_m f\|_{\ell^\infty}^{\theta'}, && \text{by interpolation} \\ &\lesssim \|h\|_{L^1}^{\theta'/4} \underbrace{\|h\|_{L^\infty}^{\theta'/4} \|h\|_{H^{2\beta/3}}^{\theta'/2}}_{\lesssim \|h\|_{C^\beta}^{3\theta'/4} \lesssim 1} + \|I_m f\|_{\ell^\infty}^{\theta'}, && \text{by interpolation.} \end{aligned}$$

By using (4.9), we obtain  $\|f\|_{H^s} \lesssim \|I_m f\|_{\ell^\infty}^\theta$ , with  $\theta = \theta'/4$ .  $\square$

*Remark 4.11.* Actually, the proof also works by taking  $f \in C^r$  and  $\|f\|_{C^r} \leq 1$  (for some  $r \in (0, 1)$ ) but the new control one obtains is still in the form  $\|f\|_{H^s} \lesssim \|I_m f\|_{\ell^\infty}^\theta$ . It is very likely that one can obtain a control in  $C^s$  in the end by looking at the action of the pseudodifferential operator  $\Pi_m$  on Hölder-Zygmund spaces  $C^{k,\alpha}$  (see [Tay91]). One could also possibly better the exponent  $\theta$  (and make it arbitrarily close to 1) by working in Besov spaces but we did not investigate these details.

## REFERENCES

- [BK85] K. Burns and A. Katok. Manifolds with nonpositive curvature. *Ergodic Theory Dynam. Systems*, 5(2):307–317, 1985.
- [CS98] Christopher B. Croke and Vladimir A. Sharafutdinov. Spectral rigidity of a compact negatively curved manifold. *Topology*, 37(6):1265–1273, 1998.
- [dlLMM86] R. de la Llave, J. M. Marco, and R. Moriyón. Canonical perturbation theory of Anosov systems and regularity results for the Livšic cohomology equation. *Ann. of Math. (2)*, 123(3):537–611, 1986.
- [DS03] Nurlan S. Dairbekov and Vladimir A. Sharafutdinov. Some problems of integral geometry on Anosov manifolds. *Ergodic Theory Dynam. Systems*, 23(1):59–74, 2003.

- [FS11] Frédéric Faure and Johannes Sjöstrand. Upper bound on the density of Ruelle resonances for Anosov flows. *Comm. Math. Phys.*, 308(2):325–364, 2011.
- [GK80] V. Guillemin and D. Kazhdan. Some inverse spectral results for negatively curved 2-manifolds. *Topology*, 19(3):301–312, 1980.
- [GL18] C. Guillarmou and T. Lefeuvre. The marked length spectrum of Anosov manifolds. *ArXiv e-prints*, June 2018.
- [Gui17] Colin Guillarmou. Invariant distributions and X-ray transform for Anosov flows. *J. Differential Geom.*, 105(2):177–208, 2017.
- [HF] Boris Hasselblatt and Todd Fisher. *Hyperbolic flows*. In preparation, \*.
- [Kat90] Svetlana Katok. Approximate solutions of cohomological equations associated with some Anosov flows. *Ergodic Theory Dynam. Systems*, 10(2):367–379, 1990.
- [KH95] Anatole Katok and Boris Hasselblatt. *Introduction to the modern theory of dynamical systems*, volume 54 of *Encyclopedia of Mathematics and its Applications*. Cambridge University Press, Cambridge, 1995. With a supplementary chapter by Katok and Leonardo Mendoza.
- [Kli74] Wilhelm Klingenberg. Riemannian manifolds with geodesic flow of Anosov type. *Ann. of Math. (2)*, 99:1–13, 1974.
- [Liv72] A. N. Livšic. Cohomology of dynamical systems. *Izv. Akad. Nauk SSSR Ser. Mat.*, 36:1296–1320, 1972.
- [Liv04] Carlangelo Liverani. On contact Anosov flows. *Ann. of Math. (2)*, 159(3):1275–1312, 2004.
- [LT05] A. O. Lopes and Ph. Thieullen. Sub-actions for Anosov flows. *Ergodic Theory Dynam. Systems*, 25(2):605–628, 2005.
- [Par88] William Parry. Equilibrium states and weighted uniform distribution of closed orbits. In *Dynamical systems (College Park, MD, 1986–87)*, volume 1342 of *Lecture Notes in Math.*, pages 617–625. Springer, Berlin, 1988.
- [PSU14] Gabriel P. Paternain, Mikko Salo, and Gunther Uhlmann. Spectral rigidity and invariant distributions on Anosov surfaces. *J. Differential Geom.*, 98(1):147–181, 2014.
- [Sha94] V. A. Sharafutdinov. *Integral geometry of tensor fields*. Inverse and Ill-posed Problems Series. VSP, Utrecht, 1994.
- [Shu01] M. A. Shubin. *Pseudodifferential operators and spectral theory*. Springer-Verlag, Berlin, second edition, 2001. Translated from the 1978 Russian original by Stig I. Andersson.
- [SSU05] Vladimir Sharafutdinov, Michal Skokan, and Gunther Uhlmann. Regularity of ghosts in tensor tomography. *J. Geom. Anal.*, 15(3):499–542, 2005.
- [Tay91] Michael E. Taylor. *Pseudodifferential operators and nonlinear PDE*, volume 100 of *Progress in Mathematics*. Birkhäuser Boston, Inc., Boston, MA, 1991.
- [Zwo12] Maciej Zworski. *Semiclassical analysis*, volume 138 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 2012.

LABORATOIRE JEAN LERAY, CNRS UMR 6629, UNIVERSITÉ DE NANTES, 2 RUE DE LA HOUSSINIÈRE,  
44322 NANTES, FRANCE

*Email address:* `sebastien.gouezel@univ-nantes.fr`

LABORATOIRE DE MATHÉMATIQUES D’ORSAY, UNIV. PARIS-SUD, CNRS, UNIVERSITÉ PARIS-SACLAY,  
91405 ORSAY, FRANCE

*Email address:* `thibault.lefeuvre@u-psud.fr`