# Supplementary Information 

## For

# Large Perpendicular Magnetic Anisotropy in $\mathbf{T a} / \mathbf{C o F e B} / \mathbf{M g O}$ on Full-Coverage Monolayer MoS $\mathbf{I}_{2}$ and First Principles Study of Its Electronic Structure 

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Note1: Procedure to extract effective anisotropy energy from in-plane and out-of-plane $M-H$ curves.
The effective anisotropy constant ( $K_{\text {eff }}$ ) is an important parameter for characterizing the anisotropy of the sample, which can be calculated by $K_{\text {eff }}=M_{\text {s-eff }} \times \Delta / \mu_{0}$, where $M_{\text {s-eff }}$ is the effective saturation magnetization per volume, $\Delta$ is the difference in the area of the in-plane (IP) and out-ofplane (OOP) $M$ - $H$ curves in the first quadrant, as shown in Figure $\mathbf{S 1}$ by the shaded area (blue lines).

From the fitting curve of thickness dependent magnetization (Figure 3 b ), $M_{\mathrm{s} \text {-eff }}$ has been determined to be $1216 \mathrm{emu} / \mathrm{cm}^{3} . \Delta$ is extracted to be 0.09716 T in Figure S1. Furthermore, $\mu_{0}=4 \pi \times 10^{-}$ ${ }^{7} \mathrm{H} / \mathrm{m}$ and $1 \mathrm{emu} / \mathrm{cm}^{3}=4 \pi \times 10^{-4} \mathrm{~T}$. By injecting all parameters into the formula, the value of $K_{\text {eff }}$ can be calculated to be: $K_{\text {eff }}=M_{\text {s-eff }} \times \Delta / \mu_{0}=\left(1216 \times 4 \pi \times 10^{-4}\right) \times(0.09716) \div\left(4 \pi \times 10^{-7}\right)=118136.844 \mathrm{~J} / \mathrm{m}^{3}$.


Figure S1: Out-of-plane and In-plane $M-H$ curves for 1.2 nm CoFeB sample measured at 10 K .

Note2: Details for extracting the effective anisotropy energy for the samples with different thickness of CoFeB .


Figure S2: (a-g) Out-of-plane and in-plane $M-H$ curves for samples with different CoFeB thicknesses with $T_{\mathrm{an}}=250^{\circ} \mathrm{C}$ measured at 10 K . (a) 1 nm , (b) 1.1 nm , (c) 1.2 nm , (d) 1.3 nm , (e) 1.4 nm , (f) 1.5 nm , (g) 1.6 nm . (h) $t_{\text {eff }}$ dependence of the product of $K_{\text {eff }}$ and $t_{\text {eff }}$, where the intercept to the vertical axis of the linear extrapolation corresponds to the interface anisotropy $K_{\mathrm{i}}$.

Note3: Details for extracting the effective anisotropy energy for the samples with different annealing temperature.


Figure S3: (a-e) Out-of-plane and in-plane $M-H$ curves for the sample with 1.2 nm CoFeB before and after different temperature annealing measured at 10 K . (a) no annealing, (b) $200^{\circ} \mathrm{C}$ annealing, (c) $250^{\circ} \mathrm{C}$ annealing, (d) $300^{\circ} \mathrm{C}$ annealing and (e) $350^{\circ} \mathrm{C}$ annealing. (f) $K_{\text {eff }}$ as a function of annealing temperature.

Note4: Element distribution profile and B diffusion in MgO for the sample capped with Pt.


Figure S4: (a) The maximum-normalized element concentration profiles from EELS drawn perpendicularly to the $\mathrm{MoS}_{2}$ layer shows that B diffuses completely inside the MgO barrier. (b) EELS spectrum of $\mathrm{B}_{\mathrm{K}}$ edge extracted from spectrum image in the MgO layer. The energy loss near edge structure shows the fingerprint of B III, indicating the oxidation of boron simultaneously to its diffusion.

Note5: Band structure of the primitive (hexagonal) and conventional (rectangular) unit cells of a $\mathrm{MoS}_{2} \mathrm{ML}$.

Figs S5a and S5b show the band structure of an isolated $\mathrm{MoS}_{2}$ monolayer, when the crystal structure is respectively described with a hexagonal primitive cell and with a rectangular conventional cell similar to that used to study the $\mathrm{Fe} / \mathrm{MgO} / \mathrm{MoS}_{2}$ multilayer. The same $\Gamma-\mathrm{M}-\mathrm{K}-\Gamma-\mathrm{K}$ '- $\mathrm{M}-\Gamma$ path has been used in the two figures, except that this path connects high symmetry directions shown with red letters in Figure 6a for the former case, and with black letters for the latter case. Comparison of the two figures shows how the band structure of $\mathrm{MoS}_{2}$ is folded when the rectangular conventional cell is used. In particular, the electron states corresponding to some of the M points are folded at $\Gamma$ in this representation.


Figure S5: (a) Band structure of a $\mathrm{MoS}_{2}$ monolayer calculated with a primitive hexagonal unit cell. (b) Band structure of a $\mathrm{MoS}_{2}$ monolayer calculated with the conventional rectangular unit cell. The color bar shows the value of the spin component $s_{\mathrm{z}}$ for each of the Bloch states forming these bands.

Note6: First principle calculations for the $\mathrm{MoS}_{2} / 3 \mathrm{MLs} \mathrm{MgO} / 7 \mathrm{MLs}$ Fe structure


Figure S6: First principle calculations for the $\mathrm{MoS}_{2} / 3 \mathrm{MLs} \mathrm{MgO} / 7$ MLs Fe structure. (a) Majority and minority spin DOS integrated over atomic spheres belonging to the successive atomic layers of the supercell. (b) $z$-variations of the $x y$-averaged electrostatic potential. (c) The red curve shows the variations of the electron charge density when the supercell is built from the $\mathrm{MoS}_{2} \mathrm{ML}$, the 3 ML thick MgO layer and the 7 ML-thick Fe layer. Each blue dot corresponds to the $z$-integration of the red curve between the two mid-perpendicular planes separating an atomic layer and its neighbors on the right and left sides. (d) $z$-dependence of the DOS at the Fermi level, calculated for majority (black line) and minority spin electrons (red line). (e) Profile of spin polarization at the Fermi level, calculated from the spin-resolved DOS.

