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Qualitative analysis of optimisation problems with respect to non-constant Robin coefficients

Idriss Mazari* Yannick Privat^{† ‡}

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Abstract

Following recent interest in the qualitative analysis of some optimal control and shape optimisation problems, we provide in this article a detailed study of the optimisation of Robin boundary conditions in PDE constrained calculus of variations. Our main model consists of an elliptic PDE of the form $-\Delta u_\beta = f(x, u_\beta)$ endowed with the Robin boundary conditions $\partial_\nu u_\beta + \beta(x)u_\beta = 0$. The optimisation variable is the function β , which is assumed to take values between 0 and 1 and to have a fixed integral. Two types of criteria are under consideration; the first one is non-energetic criteria. In other words, we aim at optimising functionals of the form $\mathcal{J}(\beta) = \int_{\Omega \text{ or } \partial\Omega} j(u_\beta)$. We prove that, depending on the monotonicity of the function j , the optimisers may be of *bang-bang* type (in other words, the optimisers write 1_Γ for some measurable subset Γ of $\partial\Omega$) or, on the contrary, that they may only take values strictly between 0 and 1. This has consequence for a related shape optimisation problem, in which one tries to find where on the boundary Neumann ($\partial_\nu u = 0$) and constant Robin conditions ($\partial_\nu u + u = 0$) should be placed in order to optimise criteria. The proofs for this first case rely on new fine oscillatory techniques, used in combination with optimality conditions. We then investigate the case of compliance-type functionals. For such energetic functionals, we give an in-depth analysis and even some explicit characterisation of optimal β^* .

Keywords: Elliptic boundary values problems, Robin boundary conditions, Calculus of variations, Shape optimisation, Bilinear optimal control problems, Qualitative analysis of optimisation problems

AMS classification: 49J15, 49Q10.

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1 Introduction

1.1 Scope of the article, informal presentation of the problem

1.1.1 Informal presentation of the problem

The goal of this article is to provide a theoretical analysis of a class of PDE constrained optimisation problems which arise in many fields (for instance, in automation, in physics or in mathematical biology), and in which the aim is to minimise or maximise a certain criteria by acting on the coefficients of the Robin boundary conditions.

More specifically, we are working with *heterogeneous Robin boundary conditions*, in the following sense: the state equation of the phenomenon is supplemented, on the boundary of the domain, with a condition of the form

$$\frac{\partial u}{\partial \nu}(x) + \beta(x)u(x) = 0,$$

where β is a non-negative function on the boundary. Our goal is to optimise certain criteria with respect to β , under some natural constraints.

Context We provide bibliographical references in section 1.6 of this introduction, but let us point out that such problems have been the topic of a wide research activity. For instance, in [9, 10], several aspects of the optimisation of the natural energy of the underlying PDE or of some eigenvalues were tackled. Similarly, this type of question is very natural in the context of thermal insulation. In this case, a relevant query is to find the best way to coat a domain with an insulant in order to optimise certain criteria. This is the point of view chosen, for example, in [9, 10, 15]. Other authors have studied this problem in parabolic models, with applications to fluid dynamics [19], or in hyperbolic problems [20]. Let us finally mention that, in many of the aforementioned cases, the functionals to be optimised either derive from the natural energy of the PDE or are of "tracking-type" (*i.e.* the aim is to minimise the distance of the state to a reference state). However, many relevant optimisation problems do not fall in either category. This is for example the case in spatial ecology. One may consider, following [4, 25, 23, 26, 27, 28], the problem of maximising the population size in logistic models: how should one design the boundary of a domain in order to optimise the total population inside the domain? Although most of our analysis is, in the main proofs of this article, detailed in the case of linear models, we also provide in section 1.4 some extension to non-linear models.

We consider and analyse fairly general functionals, with a strong emphasis on the *qualitative properties* of optimisers. As we shall see, these properties are closely related to *existence results for shape optimisation problems*. To carry out our proofs, we introduce an oscillatory method which, although reminiscent of the one we introduced recently in [25], requires fine tuning to obtain our results.

One of our main contributions is the analysis of the influence of the type of functionals we wish to optimise (for instance, the *monotonicity of the functional* is crucial in the forthcoming analysis) on qualitative results. Let us mention that we consider two main types of functionals: energetic ones (in other words, functionals that are equal, up to a multiplicative constant, to the natural energy of the state equation), in which we may achieve an explicit characterisation of optimisers, and non-energetic ones. The latter case exhibits very different qualitative features; this is the main topic under consideration here.

Furthermore, let us underline that, from a mathematical perspective, our contributions can be read through the lens of *bilinear optimal control problems* set on the boundary. In this setting, we see β as the control. Likewise, such bilinear optimal control problems have been an active topic of research in the past years, and are not yet fully understood. We refer once again to section 1.6 of the introduction for further discussion of recent works in this field.

Paradigmatic formulation of our problem and motivation The most general version of our problem reads as follows: let Ω be a regular enough domain in \mathbb{R}^d and let, for any measurable subset $\Gamma \subset \partial\Omega$, u_α^Γ be the unique solution of

$$\begin{cases} -\nabla \cdot (A(x) \nabla u_\alpha^\Gamma(x)) = g_0(x, u_\alpha^\Gamma(x)) & x \in \Omega \\ \partial_\nu u_\alpha^\Gamma(x) + \alpha \mathbb{1}_\Gamma(x) u_\alpha^\Gamma(x) = 0 & x \in \partial\Omega. \end{cases} \quad (1.1)$$

In this formulation, $\alpha > 0$ is a fixed parameter, A is a matrix assumed to be uniformly elliptic and g_0 is a given non-linearity.

A possible interpretation of this equation is that it is an approximation of a mixed problem of the type

$$\begin{cases} -\nabla \cdot (A(x) \nabla v^\Gamma(x)) = g_0(x, v^\Gamma(x)) & x \in \Omega, \\ \partial_\nu v^\Gamma(x) = 0 & x \in \partial\Omega \setminus \Gamma, \\ v^\Gamma(x) = 0 & x \in \Gamma. \end{cases} \quad (1.2)$$

Indeed, under adequate assumptions, the solution to (1.1) converges in some sense to the one of (1.2) as $\alpha \rightarrow +\infty$. We refer to Appendix A for additional explanation in the case where $g_0 = g_0(x)$.

As a consequence, optimising criteria involving u_α^Γ with respect to Γ is closely related to optimising criteria involving v^Γ with respect to Γ .

Remark 1 (Comment on the methods used). *Let us stress the following fact: while it is plausible that the optimisers of the problems involving u_α^Γ and of the problems involving v^Γ have the same qualitative features, we believe that the tools necessary in order to analyse them are fundamentally different, as the proper convergence for the mixed Neumann-Dirichlet problem should be the Γ -convergence of sets while, for the problems involving u_α^Γ , the relevant topology is rather the weak $L^\infty - *$ one on the compactification of $\{\mathbb{1}_\Gamma, \mathcal{H}^{d-1}(\Gamma) = V_0\}$. We refer to section 1.5 for additional comments.*

The typical problems we consider in this paper are

$$\sup_{\substack{\Gamma \subset \partial\Omega \\ \mathcal{H}^{d-1}(\Gamma) = V_0}} / \inf \int_\Omega j(u_\alpha^\Gamma) \quad \text{or} \quad \sup_{\substack{\Gamma \subset \partial\Omega \\ \mathcal{H}^{d-1}(\Gamma) = V_0}} / \inf \int_{\partial\Omega} j(u_\alpha^\Gamma) \quad (1.3)$$

where $\mathcal{H}^{d-1}(\cdot)$ stands for the $(d-1)$ -Hausdorff dimensional measure of Γ and j is a smooth non-linearity. Stated as such, these problems are shape optimisation problems. We consider a relaxed version of this problem, where the term $\alpha \mathbb{1}_\Gamma$ in (1.1) is replaced with a function $\beta \in L^\infty(\partial\Omega)$ satisfying certain constraints; we explain later in this introduction how qualitative properties for the optimisation with respect to β are translated to (non-)existence results for the initial shape optimisation problem.

We also consider, for the sake of completeness, compliance-type problems: namely, the goal for this class of problems is to solve

$$\sup_{\substack{\Gamma \subset \partial\Omega \\ \mathcal{H}^{d-1}(\Gamma) = V_0}} / \inf \int_\Omega f u_\alpha^\Gamma \quad (1.4)$$

where f is the source term in the equation.

Since all the results we establish hereafter are valid whatever the value of $\alpha > 0$, we may without loss of generality fix $\alpha = 1$ in this formulation. We also underline that we first work with linear models and cover non-linear models in section 1.4.

1.1.2 Plan of the introduction

The introduction of this paper is long, and we thus take the liberty to give a plan to ease the reading. Subsection 1.1.3 contains the presentation of the state equation, as well as the definitions of the two types of functionals considered. In particular, it concludes with a presentation of the focal point of our analysis for non-energetic functionals, the bang-bang property. In subsection 1.1.4, we motivate the analysis of this property by linking it to existence properties for some shape optimisation problems. In sections 1.2-1.3, we give our main theorems in the linear case, first stating the ones dealing with non-energetic criteria, second presenting the ones relevant for energetic criteria. Section 1.4 contains the results for non-linear models. Section 1.5 is devoted to the technical context of our proofs. The introduction concludes with section 1.6, which contains a discussion of the relevant references.

1.1.3 Problem under consideration in this article

Relaxation of the problem and admissible class of coefficients β . We have mentioned we would consider a relaxed version of (1.3)-(1.4). In order to make the above statement about the relaxation of the problem precise, we define, for a fixed $V_0 \in (0; \mathcal{H}^{d-1}(\partial\Omega))$, the set $\mathcal{B}(\partial\Omega)$ as

$$\mathcal{B}(\partial\Omega) := \left\{ \beta \in L^\infty(\partial\Omega) : 0 \leq \beta \leq 1, \int_{\partial\Omega} \beta = V_0 \right\}. \quad (1.5)$$

This set corresponds to the closure of the set $\{\mathbb{1}_\Gamma, \mathcal{H}^{d-1}(\Gamma) = V_0\}$ for the weak-star topology of $L^\infty(\partial\Omega)$ [18, Proposition 7.2.17]. The set $\mathcal{B}(\partial\Omega)$ is the admissible class we consider throughout this paper. The link between existence properties for the shape optimisation problems of type (1.3) and the so-called *bang-bang property* for optimisation problems set in $\mathcal{B}(\partial\Omega)$ is investigated in section 1.1.4.

State equation For the sake of simplicity, we first focus in this paper on a simpler version of (1.2). This allows us to not dwell on existence and regularity issues, and we thus hope to provide a clear description of the type of arguments involved in the proofs of our results. We refer to section 1.4 for non-linear models.

Henceforth, Ω is a fixed open bounded connected subset of \mathbb{R}^d with a \mathcal{C}^2 boundary, and $f \in L^\infty(\Omega)$ is a fixed source term. We further assume that

$$f \text{ does not vanish identically and } f \geq 0 \text{ a.e. in } \Omega. \quad (\mathbf{H}_f)$$

For any $\beta \in \mathcal{B}(\partial\Omega)$, we denote by u_β the unique solution of the equation

$$\begin{cases} -\Delta u_\beta = f & \text{in } \Omega, \\ \frac{\partial u_\beta}{\partial \nu} + \beta u_\beta = 0 & \text{on } \partial\Omega. \end{cases} \quad (\mathbf{E}_\beta)$$

Alternatively, u_β is the unique minimiser in $W^{1,2}(\Omega)$ of the energy functional

$$\mathcal{E}_\beta : W^{1,2}(\Omega) \ni u \mapsto \frac{1}{2} \int_\Omega |\nabla u|^2 + \frac{1}{2} \int_{\partial\Omega} \beta u^2 - \int_\Omega f u. \quad (1.6)$$

As a consequence of the strong maximum principle, for any $\beta \in \mathcal{B}(\partial\Omega)$ and any $f \in L^\infty(\Omega)$ satisfying (\mathbf{H}_f) , we have

$$\inf_{\overline{\Omega}} u_\beta > 0.$$

First type of functional: energetic functionals Two natural optimisation problems that stem from (\mathbf{E}_β) are the problems of maximising and minimising the compliance. In other words, we shall tackle the problems

$$\max / \min_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\Omega} f u_{\beta} = \mathcal{F}(\beta) \text{ (energetic criteria).}$$

What is notable here is that the functional \mathcal{F} rewrites naturally using the energy defined in (1.6), whence the wording "energetic" to describe such criteria. Straightforward computations indeed lead to

$$\mathcal{F}(\beta) = -2\mathcal{E}_{\beta}(u_{\beta}) = -2 \min_{u \in W^{1,2}(\Omega)} \mathcal{E}_{\beta}(u).$$

This alternative formulation enables us to obtain a finer description of optimisers. Since this is not the central point of this paper, we state the relevant results last, in Theorem 8.

Second type of functional: non-energetic criteria We want to consider boundary and interior cost functionals. For both these criteria, we consider a fixed non-linearity $j : \mathbb{R} \rightarrow \mathbb{R}$. As we will see, *the monotonicity of j plays a crucial role in the qualitative analysis of optimisers*. Thus, we choose j to be monotone. Since we are dealing both with minimisation and maximisation problems, we may take j to be increasing. Overall, we assume that

$$j \in \mathcal{C}^2(\mathbb{R}_+) \text{ and } j' > 0 \text{ on } (0; +\infty). \quad (\mathbf{H}_j)$$

Using this non-linearity we define two functionals:

$$\begin{aligned} \mathcal{J}_{\partial\Omega} : \mathcal{B}(\partial\Omega) \ni \beta &\mapsto \int_{\partial\Omega} j(u_{\beta}) \quad (\text{boundary criterion}), \\ \mathcal{J}_{\Omega} : \mathcal{B}(\partial\Omega) \ni \beta &\mapsto \int_{\Omega} j(u_{\beta}) \quad (\text{distributed criterion}), \end{aligned} \quad (1.7)$$

Our focus is on the optimisation problems

$$\sup / \inf_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\partial\Omega} j(u_{\beta}), \quad (1.8)$$

where u_{β} solves (\mathbf{E}_{β}) . We refer again to section 1.6 of this introduction for a discussion of the history of these problems. Several features of the optimisers are relevant in such queries, among them the so-called bang-bang property: *do optimisers β^* write $\mathbb{1}_{\Gamma}$ for some subset Γ of $\partial\Omega$?* To justify why this is a relevant question, let us now discuss briefly the relationship between this bang-bang property and the existence of optimal shapes for the related shape optimisation problem.

1.1.4 Relationship between existence properties for the shape optimisation problem and the bang-bang property for the relaxed formulation.

Let us now explain a bit more in details the link between the initial shape optimisation problems

$$\sup / \inf_{\substack{\Gamma \subset \partial\Omega \\ \mathcal{H}^{d-1}(\Gamma) = V_0}} \int_{\partial\Omega} j(w^{\Gamma}), \quad (1.9)$$

where w^{Γ} solves

$$\begin{cases} -\Delta w^{\Gamma} = f & \text{in } \Omega, \\ \frac{\partial w^{\Gamma}}{\partial \nu} + \mathbb{1}_{\Gamma} w^{\Gamma} = 0 & \text{on } \partial\Omega, \end{cases} \quad (1.10)$$

and the relaxed problem (1.8).

The first thing that should be noted is that (1.8) has a solution β^* . We refer to Lemma 12 below and indicate that this follows from the direct method in the calculus of variations: the weak $L^\infty - *$ compactness of $\mathcal{B}(\partial\Omega)$ and the continuity for the $L^\infty - *$ topology of the functionals suffice to obtain this result. Obtaining such an existence property for (1.9) is much harder.

However, since $\mathcal{B}(\partial\Omega)$ is the compactification of the set $\{\mathbb{1}_\Gamma, \mathcal{H}^{d-1}(\Gamma) = V_0\}$ it follows that for every $\beta \in \mathcal{B}(\partial\Omega)$ there exists a sequence $\{\Gamma_k\}_{k \in \mathbb{N}}$ of measurable subsets of $\partial\Omega$ with Hausdorff measure V_0 such that

$$\mathbb{1}_{\Gamma_k} \xrightarrow[k \rightarrow \infty]{} \beta \text{ in the weak } L^\infty - * \text{ topology.}$$

Since $\mathcal{J}_{\partial\Omega}$ is continuous for this topology, we obtain

$$\mathcal{J}_{\partial\Omega}(\mathbb{1}_{\Gamma_k}) \xrightarrow[k \rightarrow +\infty]{} \mathcal{J}_{\partial\Omega}(\beta).$$

The set $\{\mathbb{1}_\Gamma, \mathcal{H}^{d-1}(\Gamma) = V_0\}$ corresponds exactly to the set of extreme points of the admissible set $\mathcal{B}(\partial\Omega)$. We call its elements *bang-bang functions*.

With these informations it is easy to obtain the following proposition/definition describing the relationships between the shape optimisation problem (1.9) and the bilinear optimal control problem (1.8).

Definition 2. 1. Problem (1.9) has a solution if, and only if there exists a bang-bang solution β^* of Problem (1.8). In this case, we say that Problem (1.8) satisfies the bang-bang property.

2. Alternatively, Problem (1.9) does not have a solution if, and only if any solution β^* of (1.8) satisfies $\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0$. In this case, we say that Problem (1.9) enjoys a relaxation property.

1.2 First case: boundary and distributed criteria

1.2.1 Existence results and bang-bang property for maximisation problems

Boundary criteria

We first tackle the maximisation problem

$$\max_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\partial\Omega} j(u_\beta), \quad (\mathbf{P}_{\max, \partial\Omega, \mathcal{B}})$$

where u_β denotes the unique solution to (\mathbf{E}_β) . We refer to Lemma 12 below for the existence of optimal profiles. We also state the related shape optimisation problem

$$\sup_{\substack{\Gamma \subset \partial\Omega \\ \mathcal{H}^{d-1}(\Gamma) = V_0}} \int_{\partial\Omega} j(w^\Gamma). \quad (\mathbf{P}_{\max, \partial\Omega, \Sigma})$$

Our first result states that maximisers of $(\mathbf{P}_{\max, \partial\Omega, \mathcal{B}})$ satisfy the bang-bang property. Following the discussion of Section 1.1.4, the shape optimisation problem $(\mathbf{P}_{\max, \partial\Omega, \Sigma})$ has a solution.

Theorem 3. Let Ω be a bounded open set of \mathbb{R}^d such that $\partial\Omega$ is \mathcal{C}^2 . Assume f satisfies (\mathbf{H}_f) and j satisfies (\mathbf{H}_j) . Any solution β^* of the optimisation problem $(\mathbf{P}_{\max, \partial\Omega, \mathcal{B}})$ is bang-bang: there exists $\Gamma^* \subset \partial\Omega$ such that $\beta^* = \mathbb{1}_{\Gamma^*}$. As a consequence, the shape optimisation problem $(\mathbf{P}_{\max, \partial\Omega, \Sigma})$ has a solution.

The proof of this Theorem is one of the central points of this paper. It is carried out using a high frequency analysis of the second order derivative of the functional. While this type of results is usually proved using convexity or concavity arguments, the fact that the problem is not energetic and that we are considering a bilinear control problems *a priori* prohibits obtaining a convexity property for the functional $\mathcal{J}_{\partial\Omega}$. We refer to section 1.5 for a discussion of the method.

Distributed criteria

Although we decided to start with boundary criteria as, to the best of our knowledge, they have received less attention in the literature, our methods naturally extend to the case of distributed criteria. In this case, the optimisation problem is

$$\max_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\Omega} j(u_{\beta}), \quad (\mathbf{P}_{\max, \Omega, \mathcal{B}})$$

where u_{β} denotes the unique solution to (\mathbf{E}_{β}) , and the related shape optimisation problem reads

$$\sup_{\Gamma \subset \partial\Omega, \mathcal{H}^{d-1}(\Gamma)=V_0} \int_{\Omega} j(w^{\Gamma}). \quad (\mathbf{P}_{\max, \Omega, \Sigma})$$

The main result is the following Theorem:

Theorem 4. *Let Ω be a bounded open set of \mathbb{R}^d such that $\partial\Omega$ is \mathcal{C}^2 . Assume f satisfies (\mathbf{H}_f) and j satisfies (\mathbf{H}_j) . Any solution β^* of the optimisation problem $(\mathbf{P}_{\max, \Omega, \mathcal{B}})$ is bang-bang: there exists $\Gamma^* \subset \partial\Omega$ such that $\beta^* = \mathbf{1}_{\Gamma^*}$. As a consequence, the shape optimisation problem $(\mathbf{P}_{\max, \Omega, \Sigma})$ has a solution.*

Since the proof is very similar to that of Theorem 3, we omit it in the main text of the article and only give it in Appendix D.

1.2.2 Non-existence and relaxation phenomenon for minimisation problems

Boundary criteria

Let us now consider the minimisation problem

$$\min_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\partial\Omega} j(u_{\beta}). \quad (\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$$

Once again, we refer to Lemma 12 for the existence of optimal profiles. As stated hereafter, we shall show that the related shape optimisation problem

$$\inf_{\Gamma \subset \partial\Omega, \mathcal{H}^{d-1}(\Gamma)=V_0} \int_{\partial\Omega} j(w^{\Gamma}) \quad (\mathbf{P}_{\min, \partial\Omega, \Sigma})$$

exhibits a relaxation phenomenon. It is interesting to notice that the main argument for showing the second part of the following result rests upon a low frequency analysis of the second order optimality conditions.

Theorem 5. *Let Ω be a bounded open set of \mathbb{R}^n such that $\partial\Omega$ is \mathcal{C}^2 . Assume f satisfies (\mathbf{H}_f) and j satisfies (\mathbf{H}_j) .*

(i) Any solution β^* of $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$ satisfies

$$\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0,$$

so that $(\mathbf{P}_{\min, \partial\Omega, \Sigma})$ does not have a solution and enjoys a relaxation phenomenon.

(ii) Furthermore, let us introduce

$$U_0(f) := \sup_{\beta \in \mathcal{B}(\partial\Omega)} \sup_{x \in \bar{\Omega}} u_\beta(x) \in (0, +\infty).$$

There exists $C > 0$ such that, if

$$j''(u) \leq -Cj'(0) \quad \text{on} \quad [0, U_0(f)] \quad \text{and} \quad j'(U_0(f)) > 0,$$

then for any solution β^* of $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$ we have

$$\mathcal{H}^{d-1}(\{\beta^* = 1\}) > 0.$$

Let us provide an example of function j satisfying the assumptions of (ii). Given $C > 0$ and $U_0(f) > 0$, the function j given by

$$j(u) = -\frac{1}{2}u^2 + \frac{1}{2}\left(U_0(f) + \frac{1}{C}\right)u$$

fulfills these conditions provided that $CU_0(f) < 1$. Since $U_0(f)$ does not depend on C , it suffices to chose C small enough.

Remark 6. One has $U_0(f) < +\infty$ since $f \in L^\infty(\Omega)$. More precisely, one has

$$\sup_{\beta \in \mathcal{B}(\partial\Omega)} \|u_\beta\|_{W^{1,p}(\Omega)} < +\infty$$

for any $p \in [1; +\infty)$, by standard elliptic regularity estimates which are detailed in Lemma 11. We conclude by the compactness of the embedding $W^{1,p}(\Omega) \hookrightarrow \mathcal{C}^0(\bar{\Omega})$ for p large enough.

We refer to section 1.5 for comments on the proof.

Distributed criteria

Here again, some of our methods naturally extend to the case of distributed criteria. The proof of the following result is very similar to that of Theorem 5, and we provide it in Appendix E. The minimisation problem under consideration is

$$\min_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\Omega} j(u_\beta). \tag{\mathbf{P}_{\min, \Omega, \mathcal{B}}}$$

Theorem 7. Let Ω be a bounded open set of \mathbb{R}^n such that $\partial\Omega$ is \mathcal{C}^2 . Assume f satisfies (\mathbf{H}_f) and j satisfies (\mathbf{H}_j) . Then, any solution β^* of $(\mathbf{P}_{\min, \Omega, \mathcal{B}})$ satisfies

$$\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0.$$

1.3 Second case: energetic criteria

Let us now tackle the two energetic optimisation problems

$$\min_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\Omega} f u_{\beta} \quad (\mathbf{Q}_{\min})$$

and

$$\max_{\beta \in \mathcal{B}(\partial\Omega)} \int_{\Omega} f u_{\beta}. \quad (\mathbf{Q}_{\max})$$

The existence of optimisers for (\mathbf{Q}_{\min}) – (\mathbf{Q}_{\max}) can be obtained by adapting the arguments of Lemma 12 below.

As we have noted earlier, a salient feature of these problems is that they can be rewritten in terms of the energy of the equation (\mathbf{E}_{β}) :

$$\int_{\Omega} f u_{\beta} = -2 \min_{u \in W^{1,2}(\Omega)} \frac{1}{2} \int_{\Omega} |\nabla u|^2 + \frac{1}{2} \int_{\partial\Omega} \beta u^2 - \int_{\Omega} f u. \quad (1.11)$$

Let us mention an important consequences of this fact: as an infimum of linear functionals is concave, the criterion is convex. This is why we can expect a more precise description of the optimisers of this problem.

Let us mention that two very related contributions to the study of this problem are [9, 10], in which several problems of minimising some energetic criteria are studied. The main difference with our case is that the authors of [9, 10] rather study the problem of optimising such criteria with respect to β for the boundary conditions $\beta \partial_{\nu} u_{\beta} + u_{\beta} = 0$, which significantly changes the behaviour of the functionals.

We sum up our results in the following Theorem:

Theorem 8. *Let Ω be a bounded open set of \mathbb{R}^n such that $\partial\Omega$ is \mathcal{C}^2 . Assume f satisfies (\mathbf{H}_f) and j satisfies (\mathbf{H}_j) .*

- (i) *Every solution β^* of the maximization problem (\mathbf{Q}_{\max}) is bang-bang: there exists $\Gamma^* \subset \partial\Omega$ such that $\beta^* = \mathbf{1}_{\Gamma^*}$.*
- (ii) *Let v_{Ω} denote the solution of the Dirichlet problem*

$$\begin{cases} -\Delta v_{\Omega} = f & \text{in } \Omega, \\ v_{\Omega} = 0 & \text{on } \partial\Omega, \end{cases} \quad (1.12)$$

and let V_0^{Ω} given by

$$V_0^{\Omega} = -\frac{1}{\|\partial_{\nu} v_{\Omega}\|_{L^{\infty}(\partial\Omega)}} \int_{\partial\Omega} \partial_{\nu} v_{\Omega} \in (0, \mathcal{H}^{d-1}(\partial\Omega)].$$

For every $V_0 \in (0, V_0^{\Omega})$, the minimization problem (\mathbf{Q}_{\min}) enjoys a relaxation property and has a unique solution β^ given by*

$$\beta^* = V_0 \frac{-\partial_{\nu} v_{\Omega}}{\int_{\partial\Omega} -\partial_{\nu} v_{\Omega}}. \quad (1.13)$$

- (iii) **Case where $f(\cdot) = 1$.** *Let us assume that $f \equiv 1$. The constant admissible profile $\beta^*_{V_0} \equiv \frac{V_0}{\mathcal{H}^{d-1}(\partial\Omega)}$ solves the minimization problem (\mathbf{Q}_{\min}) if, and only if Ω is a ball.*

1.4 Non-linear models: distributed criteria

One of our initial motivation is also to understand some optimal control problems that arise in some fields of applied mathematics, for instance in mathematical biology. For a detailed introduction to optimisation problems in mathematical biology we refer to [22, Introduction] and the references therein, or to [24]. In most applications of optimal control to mathematical biology, the goal is to optimise a criterion involving a resources term related to a population density modelled through a PDE. In a nutshell, optimizing such a criterion means in general looking for a resource term or a control that guarantees the “best” survival of the population in long time, or the largest population size. The underlying equation is typically non-linear and, as it models a population, the solutions of the equation should satisfy a non-negativity assumption. This motivates the upcoming analytic setting needed in order to state the relevant results. We give a concrete example after stating Theorem 9. It should be noted that this case is, at the notational level, heavier than the linear one, but that the methods are similar to that of the linear case. For this reason, we only give the proofs in Appendix F, and only state our theorem for distributed criteria.

Analytic set-up We fix a non-linearity $g = g(x, y)$, and we first assume

$$g \text{ is measurable in both variables and } \mathcal{C}^2 \text{ in the second variable.} \quad (\mathbf{H}_{NL})$$

Keeping in mind the aforementioned mathematical biology motivation, we are looking for non-negative solutions $y_\beta \in W^{1,2}(\Omega)$ of the equation

$$\begin{cases} -\Delta y_\beta = g(x, y_\beta), & \text{in } \Omega, \\ \frac{\partial y_\beta}{\partial \nu} + \beta y_\beta = 0 & \text{on } \partial\Omega, \\ y_\beta \geq 0, y_\beta \neq 0. \end{cases} \quad (1.14)$$

We of course assume:

$$\text{For any } \beta \in \mathcal{B}(\partial\Omega), \text{ there exists a unique solution } y_\beta \in W^{1,2}(\Omega) \text{ to (1.14).} \quad (\mathbf{H}_{WP})$$

Since we are working with optimality conditions, we need to be allowed to differentiate the map $\beta \mapsto y_\beta$. This is possible, granted the steady-states y_β are linearly stable: in other words, letting μ_β be the first eigenvalue of the linearised operator, *i.e.*

$$\mu_\beta := \inf_{\varphi \in W^{1,2}(\Omega), \int_\Omega \varphi^2 = 1} \left(\int_\Omega |\nabla \varphi|^2 - \int_\Omega \frac{\partial g}{\partial y}(\cdot, y_\beta) \varphi^2 + \int_{\partial\Omega} \beta \varphi^2 \right) \quad (1.15)$$

we must have

$$\forall \beta \in \mathcal{B}(\partial\Omega), \mu_\beta > 0. \quad (\mathbf{H}_{stab})$$

Finally, $W^{1,p}$ -estimates on y_β are crucial. We hence need to ensure that

$$\forall \beta \in \mathcal{B}(\partial\Omega), \quad \forall p \in [1; +\infty), \quad y_\beta \in W^{1,p}(\Omega). \quad (\mathbf{H}_{reg})$$

These are the only assumptions we need on g . After stating the theorem, we explain why a large class of monostable non-linearities satisfies these conditions.

Optimisation problem We assume g satisfies (\mathbf{H}_{NL}) -(\mathbf{H}_{WP})-(\mathbf{H}_{stab})-(\mathbf{H}_{reg}). We still work with a function j satisfying (\mathbf{H}_j) and define

$$\mathcal{R}_\Omega(\beta) := \int_\Omega j(y_\beta).$$

We consider the optimisation problems

$$\max_{\beta \in \mathcal{B}(\partial\Omega)} \mathcal{R}_\Omega(\beta) \quad (\mathbf{R}_{\max})$$

as well as

$$\min_{\beta \in \mathcal{B}(\partial\Omega)} \mathcal{R}_\Omega(\beta). \quad (\mathbf{R}_{\min})$$

Under the assumption that g satisfies (\mathbf{H}_{NL}) -(\mathbf{H}_{WP})-(\mathbf{H}_{stab})-(\mathbf{H}_{reg}), the existence of solutions to (\mathbf{R}_{\max}) -(\mathbf{R}_{\min}) is proved following the arguments of Lemma 12 below. Our main theorem is

Theorem 9. *Assume j satisfies (\mathbf{H}_j) and g satisfies (\mathbf{H}_{NL}) -(\mathbf{H}_{WP})-(\mathbf{H}_{stab})-(\mathbf{H}_{reg}). Then:*

1. *Any solution of (\mathbf{R}_{\max}) is of bang-bang type. In other words, for every maximiser β^* of (\mathbf{R}_{\max}) , there exists a measurable subset Γ^* of $\partial\Omega$ such that $\beta^* = \mathbf{1}_{\Gamma^*}$.*
2. *Any solution β^* of (\mathbf{R}_{\min}) satisfies*

$$\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0.$$

An application to the optimal design of barriers in logistic models Let us give an example of a class of non-linearities g satisfying (\mathbf{H}_{NL}) -(\mathbf{H}_{WP})-(\mathbf{H}_{stab})-(\mathbf{H}_{reg}). We consider any function $m \in L^\infty(\Omega)$ such that

$$\int_{\Omega} m > \int_{\partial\Omega} \beta = V_0. \quad (1.16)$$

Then the non-linearity

$$g_\beta : (x, y) \mapsto y(m(x) - y)$$

satisfies (\mathbf{H}_{NL}) -(\mathbf{H}_{WP})-(\mathbf{H}_{stab})-(\mathbf{H}_{reg}). To see why this is the case, we first observe that we are exactly in the context of monostable non-linearities, which, adapting the arguments of [6], yields the existence and uniqueness of a solution y_β of (1.14), that further satisfies $y_\beta \leq \|m\|_{L^\infty}$. This readily gives the regularity estimates of Assumption (\mathbf{H}_{reg}) . Finally, we observe that, as $y_\beta \neq 0$ is a non-negative eigenfunction, associated with eigenvalue 0, the monotonicity of the eigenvalue ensures that the first eigenvalue μ_β of $-\Delta - (m - 2y_\beta)$ is positive for any β .

This equation models a population density that can access resources, accounted for in this scenario by the function m . For more references on the modelling on such phenomena we refer to [22], as well as to section 1.6 below. If we take $j(x) = x$, the optimisation problem (\mathbf{R}_{\max}) (resp. (\mathbf{R}_{\min})) can be interpreted as: how should we design the features of the fence around the domain in order to maximise (resp. minimise) the population size? In this context, the bang-bang property has been deemed to be a relevant aspect of the qualitative analysis of optimisers [25, 28].

1.5 Comments on the proof of the bang-bang property and technical context

In this section, we comment upon our methods of proofs, in order to provide some context regarding the tools we introduce to analyse the bang-bang property for the bilinear optimal control problems under consideration.

Oscillatory techniques used here The bang-bang property is often proved *via* the following reasoning: one determines the so-called “switch function” of the control problem. This function, say ϕ , encodes the optimality conditions. If an optimiser is not bang-bang, this usually implies that ϕ has a level-set of positive measure. To conclude, unique continuation theorems are invoked, that prove that this can not be the case since ϕ often satisfies a certain well-behaved optimality system.

However, this strategy is mostly useful for energetic problems. For non-energetic problems, such as the ones considered here, the equation solved by the switching function is usually not tractable. We refer for instance to [23]. Instead, the method we introduce for Robin boundary conditions hinges on the one we introduced in [25]. In it, we show that, for certain distributed bilinear optimal control problems, the second order derivative of the functional to optimise writes as something very similar to a Rayleigh quotient of a certain operator \mathcal{L} , in terms of $\dot{u}_\beta[h]$. Here, $\dot{u}_\beta[h]$ denotes the derivative of $\beta \mapsto u_\beta$ in a direction h . Concretely, this means that, the functionals under consideration in Theorems 3, 4 and 9 being denoted generically \mathcal{J} and double dots standing for second-order derivatives in a direction h , we derive an estimate of the form

$$\ddot{\mathcal{J}} \gtrsim A \|\dot{u}_\beta[h]\|_{W^{1,2}(\Omega)}^2 - B \|\dot{u}_\beta[h]\|_{L^2(\Omega)}^2$$

for some constants A and B . The monotonicity of the functional enables to choose a positive A . We refer to Proposition 14 for a precise statement. Such an estimate allows to prove the bang-bang property: by assuming that a maximiser β^* is not bang-bang, we show that there exists a perturbation h that has a suitable support and such that $\dot{u}_\beta[h]$, has, in the spectral basis of \mathcal{L} , only high modes. Thus, the second order derivative of the functional can be made arbitrarily high, and we can reach a contradiction. This idea is one of the key points of the proofs of Theorems 3, 4 and 9. Of course, several points need to be refined in order to make this scheme suitable for the present context.

Second, and this is also a novelty of the article, we show that the same expression of the second order derivative in terms of Rayleigh quotients allows to prove relaxation phenomena for minimisation problem. Contrary to the bang-bang property, the method rests upon the use of low-eigenmodes, and it is at the center of Theorem 5 and 7.

Relationship with other existence theorems in shape optimisation A particularly crucial step in all our proofs is the monotonicity of the functionals we optimise. Concretely, consider the problem $(\mathbf{P}_{\max, \partial\Omega, \mathcal{B}})$. Then one of the central points of the proof of Theorem 3 is that the functional $\mathcal{J}_{\partial\Omega}(\beta) = \int_{\partial\Omega} j(u_\beta)$ is increasing. This implies that, from a shape optimization perspective, the functional $\Gamma \mapsto \int_{\partial\Omega} j(w^\Gamma)$ is increasing. In this context, it is tempting to use the seminal theorem of Buttazzo-DalMaso [11] to conclude that there exists a solution of $(\mathbf{P}_{\max, \partial\Omega, \Sigma})$. However, the topology on sets required by [11] is not suited to our type of problems: in [11], the convergence on sets is the Γ -convergence; unfortunately, it is not clear that this topology makes the functional we consider here continuous, which prohibits using this result. Furthermore, the strategies of proof are very different. Nevertheless, we wish to highlight the fact that, as in [11], the crux of the problem is the monotonicity of the functional.

1.6 Related works and bibliographical references

Optimisation of criteria involving the Robin coefficients Two lines of research coexist when it comes to optimisation in combination with Robin boundary conditions, one dedicated to the type of problems we consider here, where the domain Ω is fixed and β is the variable, and another one focused on optimising the domain Ω itself. Some contributions combine these two approaches.

Among the vast literature relevant in such queries, let us first single out [19] and [20]. Both papers deal with time-dependent equation, and focus on "tracking-type" functionals. In other words, in the framework of our paper, this would mean that the functional to optimise would involve a term of the form $\|u_\beta - u_{\text{ref}}\|_X^2$ for some norm X and some reference u_{ref} . Such criteria are known to be very relevant in many applications. Our paper, on the other hand, focuses on "free" functionals, and the methods used to analyse each problem are necessarily very different.

The aforementioned [9, 10] investigate the properties of the optimisers β that minimise the natural energy of the model or some related eigenvalues. Although Theorem 8 is closely related to these contributions, let us highlight the fact that they consider boundary conditions of the form $\beta \partial_\nu u + u = 0$, which changes the features of the problems. This contribution falls into the first category described above (energetic functionals), but it is noteworthy that the motivation in [9, 10] is an optimal insulation problem which was then deeply explored from the point of view of both optimal control and shape optimisation in the recent [14]. Notable in the latter is the fact that one of their main results, [14, Theorem 5.1] uses Talenti-like inequalities for Robin boundary conditions, which has been another very active line of research following [2]. Finally, let us mention [13], where a related spectral optimisation problem for quasilinear elliptic operators, which can be interpreted as a p -laplacian version of a Robin spectrum, is investigated. Of particular interest is that the authors of [13] study optimal designs under a volume constraint only (in our case, this would amount to only prescribing $\int_{\partial\Omega} \beta$) and obtain several results related to a relaxed formulation of the problem.

Bilinear control problems Let us first underline that the study of bilinear controllability of systems (*i.e.* trying to reach an exact state using a bilinear control) is a very active field. We refer, for instance, to [1, 5, 12].

On the other hand, bilinear optimal control problems, in which one rather aims at optimising a certain criterion, as is the case in the present paper, have received less attention. We have already mentioned several contributions related to the optimisation of Robin coefficients in order to minimise tracking-type functionals; in this broader context of bilinear optimisation, let us also refer to [16, 17], where bilinear optimal control problems for chemotaxis or chemorepulsion models are studied. The functional the authors wish to minimise is also of tracking type, but the control acts on the interior of the domain rather than on the boundary. Most of the emphasis is placed on deriving existence results and optimality conditions. Related to these contributions, we point, for another distributed bilinear optimal control problem, to [8], a contribution that focuses on a numerical multigrid analysis of the optimisation system.

However, the literature is scarcer when it comes to the qualitative analysis of bilinear optimal control problems when the functionals is not of tracking type. A paradigmatic example is that of the optimisation of the total population size in monostable models. In this problem, which originated in [21], the goal is to spread resources so as to maximise the integral of the solution of a reaction-diffusion equation. In the elliptic case, proving the bang-bang property for optimal resources distributions proves surprisingly difficult, and the problem exhibits a very intricate qualitative behaviour. We refer to [4, 7, 23, 26, 27, 28] and the references therein for an overview of this problem. It should be noted that the bang-bang property was only recently proved by the two authors in collaboration with a third one in [25], using a method that, as was explained, is linked to the one we develop here for Robin controls.

2 Preliminary material

We first give some existence results and regularity estimates for the solutions of (\mathbf{E}_β) . Since the proofs are standard, we only give them below in Appendix B. We begin with a uniform estimate:

Lemma 10. *Let Ω be a bounded domain in \mathbb{R}^d with a \mathcal{C}^2 boundary. There exists $C > 0$ such that*

$$\forall \beta \in \mathcal{B}(\partial\Omega), \quad \forall v \in W^{1,2}(\Omega), \quad C\|v\|_{W^{1,2}(\Omega)}^2 \leq \int_{\Omega} |\nabla v|^2 + \int_{\partial\Omega} \beta v^2.$$

Lemma 11. *Assume $f \in L^\infty(\Omega)$. Then the equation (\mathbf{E}_β) has a unique solution $u_\beta \in W^{1,2}(\Omega)$. Furthermore, for any $p \in [1; +\infty)$, $u_\beta \in W^{1,p}(\Omega)$,*

$$\sup_{\beta \in \mathcal{B}(\partial\Omega)} \|u_\beta\|_{W^{1,p}(\Omega)} < \infty$$

and

$$\forall \beta \in \mathcal{B}(\partial\Omega), \inf_{\Omega} u_\beta > 0.$$

In particular, as a consequence of Sobolev embeddings, we have the uniform estimates

$$0 < \inf_{\beta \in \mathcal{B}(\partial\Omega)} \inf_{\Omega} u_\beta \leq \sup_{\beta \in \mathcal{B}(\partial\Omega)} \sup_{\Omega} u_\beta < \infty.$$

This lemma is a standard consequence of Robin regularity results.

We finally state our existence result.

Lemma 12. *Let us assume that $\partial\Omega$ is \mathcal{C}^2 . Each of the problems $(\mathbf{P}_{\max, \partial\Omega, \mathcal{B}})$ - $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$ - $(\mathbf{P}_{\max, \Omega, \mathcal{B}})$ - $(\mathbf{P}_{\min, \Omega, \mathcal{B}})$ has a solution.*

3 Proofs of Theorems 3 and 5

Throughout this section, f and j are assumed to satisfy (\mathbf{H}_f) and (\mathbf{H}_j) respectively. We work only on boundary criteria; thus, to alleviate notation, we drop the subscript in $\mathcal{J}_{\partial\Omega}$ and simply write

$$\mathcal{J}(\beta) := \int_{\partial\Omega} j(u_\beta).$$

We begin by computing the derivatives of \mathcal{J} , as these derivatives are key in proving Theorems 3 and 5.

3.1 Computation of the derivatives of $\mathcal{J}_{\partial\Omega}$

The differentiability of the map $\mathcal{B}(\partial\Omega) \ni \beta \mapsto u_\beta \in W^{1,2}(\Omega)$ is standard. To specify the optimality conditions we use, let us introduce a bit of terminology.

Definition 13. *The cone of admissible perturbation at a given $\beta \in \mathcal{B}(\partial\Omega)$ is the set of functions $h \in L^\infty(\Omega)$ such that, for any sequence of positive real numbers ε_n decreasing to 0, there exists a sequence of functions $h_n \in L^\infty(\Omega)$ converging in $L^2(\Omega)$ to h as $n \rightarrow +\infty$, and $\beta + \varepsilon_n h_n \in \mathcal{B}(\partial\Omega)$ for every $n \in \mathbb{N}$. An element of the cone of admissible perturbations is called an admissible perturbation.*

Let us consider $\beta \in \mathcal{B}(\partial\Omega)$ and an admissible perturbation h at β . We denote with a single (resp. double) dot the first (resp. second) order Gateaux derivative of relevant quantities at β in the direction h . By differentiating (\mathbf{E}_β) , we see that \dot{u}_β solves

$$\begin{cases} -\Delta \dot{u}_\beta = 0 & \text{in } \Omega, \\ \frac{\partial \dot{u}_\beta}{\partial \nu} + \beta \dot{u}_\beta = -h u_\beta & \text{on } \partial\Omega, \end{cases} \quad (3.1)$$

while \ddot{u}_β satisfies

$$\begin{cases} -\Delta \ddot{u}_\beta = 0 & \text{in } \Omega, \\ \frac{\partial \ddot{u}_\beta}{\partial \nu} + \beta \ddot{u}_\beta = -2h\dot{u}_\beta & \text{on } \partial\Omega. \end{cases} \quad (3.2)$$

Existence and uniqueness of $W^{1,2}(\Omega)$ solutions to these equations are immediate. The derivatives of the criterion \mathcal{J} are, similarly, given by

$$\dot{\mathcal{J}}(\beta)[h] = \int_{\Omega} \dot{u}_\beta j'(u_\beta) \text{ and } \ddot{\mathcal{J}}(\beta)[h, h] = \int_{\partial\Omega} \ddot{u}_\beta j'(u_\beta) + \int_{\partial\Omega} (\dot{u}_\beta)^2 j''(u_\beta). \quad (3.3)$$

In order to make these derivatives more tractable, we introduce the adjoint state p_β as the unique solution in $W^{1,2}(\Omega)$ of

$$\begin{cases} -\Delta p_\beta = 0 & \text{in } \Omega, \\ \frac{\partial p_\beta}{\partial \nu} + \beta p_\beta = j'(u_\beta) & \text{on } \partial\Omega. \end{cases} \quad (3.4)$$

Since $\inf_{\Omega} u_\beta > 0$ and since $j' > 0$ on \mathbb{R}_+^* , it follows from the maximum principle that p_β is positive in $\overline{\Omega}$ and that we even have

$$\inf_{\overline{\Omega}} p_\beta > 0. \quad (3.5)$$

Multiplying (3.4) by \dot{u}_β and (3.1) by p_β , an integration by parts yields

$$\begin{aligned} 0 &= \int_{\Omega} \nabla \dot{u}_\beta \cdot \nabla p_\beta + \int_{\partial\Omega} \beta p_\beta \dot{u}_\beta - \int_{\partial\Omega} j'(u_\beta) \dot{u}_\beta \\ 0 &= \int_{\Omega} \nabla \dot{u}_\beta \cdot \nabla p_\beta + \int_{\partial\Omega} \beta p_\beta \dot{u}_\beta + \int_{\partial\Omega} h u_\beta p_\beta \end{aligned}$$

Hence, we obtain

$$\int_{\partial\Omega} j'(u_\beta) \dot{u}_\beta = - \int_{\partial\Omega} h u_\beta p_\beta,$$

by combining the two identities above. In particular, according to (3.3), we get

$$\dot{\mathcal{J}}(\beta)[h] = - \int_{\partial\Omega} h u_\beta p_\beta. \quad (3.6)$$

Multiplying (3.4) by \ddot{u}_β , and (3.2) by p_β , and integrating by parts also gives the following expression for the second-order Gateaux derivative of \mathcal{J} :

$$\ddot{\mathcal{J}}(\beta)[h, h] = -2 \int_{\partial\Omega} h \dot{u}_\beta p_\beta + \int_{\partial\Omega} j''(u_\beta) (\dot{u}_\beta)^2. \quad (3.7)$$

3.2 Proof of Theorem 3

We shall argue by contradiction: let us fix a maximiser $\beta^* \in \mathcal{B}(\partial\Omega)$ (which exists thanks to Lemma 12) such that the set

$$\omega^* := \{0 < \beta^* < 1\} \quad (3.8)$$

has positive measure:

$$\mathcal{H}^{d-1}(\omega^*) > 0.$$

It follows that for any admissible perturbation h at β^* supported in ω^* (in the sense of Definition 13), there holds

$$\dot{\mathcal{J}}(\beta^*)[h] = 0. \quad (3.9)$$

To reach a contradiction, it suffices to prove that there exists an admissible perturbation h supported in ω^* such that

$$\ddot{\mathcal{J}}(\beta^*)[h, h] > 0. \quad (3.10)$$

We start by recalling that, from (3.7), we have

$$\ddot{\mathcal{J}}(\beta^*)[h, h] = -2 \int_{\partial\Omega} h \dot{u}_{\beta^*} p_{\beta^*} + \int_{\partial\Omega} j''(u_{\beta^*}) (\dot{u}_{\beta^*})^2. \quad (3.11)$$

The first step of this proof is to obtain an expression of $\ddot{\mathcal{J}}$ that is reminiscent of a Rayleigh quotient. This is the purpose of the following proposition:

Proposition 14. *There exist three constants A, B, C with $A > 0$ such that*

$$\ddot{\mathcal{J}}(\beta^*)[h, h] \geq A \int_{\Omega} |\nabla \dot{u}_{\beta^*}|^2 - B \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)} - C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (3.12)$$

Proof of Proposition 14. We isolate the first part of (3.7), and define

$$\mathcal{W}(\beta^*)[h, h] := -2 \int_{\partial\Omega} h \dot{u}_{\beta^*} p_{\beta^*}. \quad (3.13)$$

Using the boundary condition of the equation (3.1) on \dot{u}_{β^*} this quantity rewrites

$$\mathcal{W}(\beta^*)[h, h] = 2 \int_{\partial\Omega} (\partial_{\nu} \dot{u}_{\beta^*} + \beta^* \dot{u}_{\beta^*}) \frac{p_{\beta^*}}{u_{\beta^*}} \dot{u}_{\beta^*}.$$

Let us introduce the function Ψ_{β^*} given by

$$\Psi_{\beta^*} := \frac{p_{\beta^*}}{u_{\beta^*}},$$

so that

$$\mathcal{W}(\beta^*)[h, h] = \int_{\partial\Omega} \Psi_{\beta^*} (\partial_{\nu} (\dot{u}_{\beta^*}^2) + 2\beta^* \dot{u}_{\beta^*}^2). \quad (3.14)$$

Since $\inf_{\bar{\Omega}} u_{\beta^*}, p_{\beta^*} > 0$ from (3.5) and lemma 11, and since $\sup_{\bar{\Omega}} u_{\beta^*}, p_{\beta^*} < \infty$, one has

$$\inf_{\bar{\Omega}} \Psi_{\beta^*} > 0.$$

Furthermore, for every function $z \in W^{2,2}(\Omega)$, one has

$$-\int_{\Omega} (\Delta \Psi_{\beta^*}) z + \int_{\Omega} (\Delta z) \Psi_{\beta^*} = \int_{\partial\Omega} (\partial_{\nu} z) \Psi_{\beta^*} - \int_{\partial\Omega} z (\partial_{\nu} \Psi_{\beta^*}). \quad (3.15)$$

We first compute

$$\frac{\partial \Psi_{\beta^*}}{\partial \nu} = \frac{\partial_{\nu} p_{\beta^*}}{u_{\beta^*}} - \Psi_{\beta^*} \frac{\partial_{\nu} u_{\beta^*}}{u_{\beta^*}} = -\beta^* \Psi_{\beta^*} + \beta^* \Psi_{\beta^*} + \frac{j'(u_{\beta^*})}{u_{\beta^*}} = \frac{j'(u_{\beta^*})}{u_{\beta^*}} > 0. \quad (3.16)$$

We want to take $z = \dot{u}_{\beta^*}^2$ in (3.15). We have

$$\Delta z = 2|\nabla \dot{u}_{\beta^*}|^2 + 2\dot{u}_{\beta^*} \Delta \dot{u}_{\beta^*} = 2|\nabla \dot{u}_{\beta^*}|^2 \text{ since } \Delta \dot{u}_{\beta^*} = 0.$$

Hence,

$$\int_{\partial\Omega} \Psi_{\beta^*} \partial_{\nu} (\dot{u}_{\beta^*}^2) = 2 \int_{\Omega} |\nabla \dot{u}_{\beta^*}|^2 \Psi_{\beta^*} - \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} + \int_{\partial\Omega} \frac{j'(u_{\beta^*})}{u_{\beta^*}} \dot{u}_{\beta^*}^2 \quad (3.17)$$

Thus, the second order derivative of \mathcal{J} rewrites

$$\ddot{\mathcal{J}}(\beta^*)[h, h] = 2 \int_{\Omega} \Psi_{\beta^*} |\nabla \dot{u}_{\beta^*}|^2 - \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} + \int_{\partial\Omega} \left(2\beta^* \Psi_{\beta^*} + \frac{j'(u_{\beta^*})}{u_{\beta^*}} + j''(u) \right) \dot{u}_{\beta^*}^2. \quad (3.18)$$

Since β^* and u_{β^*} belong to $L^\infty(\partial\Omega)$ and since j is \mathcal{C}^2 , there exists a constant C independent of h such that, for any admissible perturbation h ,

$$\int_{\partial\Omega} \left(2\beta^* \Psi_{\beta^*} + \frac{j'(u_{\beta^*})}{u_{\beta^*}} + j''(u) \right) (\dot{u}_{\beta^*})^2 \geq -C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (3.19)$$

Overall, we thus have the following estimate on $\ddot{\mathcal{J}}$:

$$\ddot{\mathcal{J}}(\beta^*)[h, h] \geq 2 \int_{\Omega} \Psi_{\beta^*} |\nabla \dot{u}_{\beta^*}|^2 - \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} - C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (3.20)$$

Since $\inf_{\Omega} \Psi_{\beta^*} > 0$, there exists $A > 0$ such that

$$\ddot{\mathcal{J}}(\beta^*)[h, h] \geq A \int_{\Omega} |\nabla \dot{u}_{\beta^*}|^2 - \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} - C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (3.21)$$

Direct computations yield

$$\begin{aligned} \Delta \Psi_{\beta^*} &= \frac{\Delta p_{\beta^*}}{u_{\beta^*}} - \Psi_{\beta^*} \frac{\Delta u_{\beta^*}}{u_{\beta^*}} - 2 \frac{\nabla u_{\beta^*} \cdot \nabla p_{\beta^*}}{u_{\beta^*}^2} + 2 \Psi_{\beta^*} \frac{|\nabla u_{\beta^*}|^2}{u_{\beta^*}^2} \\ &= \frac{f \Psi_{\beta^*}}{u_{\beta^*}} - 2 \frac{\nabla u_{\beta^*} \cdot \nabla p_{\beta^*}}{u_{\beta^*}^2} + 2 \Psi_{\beta^*} \frac{|\nabla u_{\beta^*}|^2}{u_{\beta^*}^2}. \end{aligned}$$

Besides, since u_{β^*} and p_{β^*} belong to $W^{1,p}(\Omega)$ for any $p \in [1; +\infty)$ and since $\inf_{\overline{\Omega}} u_{\beta^*} > 0$, it follows that, for any $p \geq 1$,

$$-2 \frac{\nabla u_{\beta^*} \cdot \nabla p_{\beta^*}}{u_{\beta^*}^2} + 2 \Psi_{\beta^*} \frac{|\nabla u_{\beta^*}|^2}{u_{\beta^*}^2} \in L^p(\Omega).$$

Therefore, $\Delta \Psi_{\beta^*}$ belongs to $L^p(\Omega)$ for any $p \geq 1$. Now, let us apply Hölder's inequality with three exponents: let $(q, r) \in (1; +\infty)^2$ be such that

$$\frac{1}{2} + \frac{1}{q} + \frac{1}{r} = 1 \text{ and } r \in (2; 2^*) \text{ with } 2^* := \frac{2d}{d-2}.$$

By the Hölder inequality, we obtain

$$\left| \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} \right| \leq \|\Delta \Psi_{\beta^*}\|_{L^q(\Omega)} \|\dot{u}_{\beta^*}\|_{L^r(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)} \quad (3.22)$$

and so, from the Sobolev embedding $W^{1,2}(\Omega) \hookrightarrow L^r(\Omega)$, there exists a constant $B > 0$ such that

$$\left| \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} \right| \leq B \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)}. \quad (3.23)$$

We thus infer that

$$\ddot{\mathcal{J}}(\beta^*)[h, h] \geq A \int_{\Omega} |\nabla \dot{u}_{\beta^*}|^2 - B \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)} - C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (3.24)$$

As a consequence, (3.12) is proved.

This concludes the proof of the Proposition. \square

We point that since the standard $W^{1,2}(\Omega)$ norm is equivalent to $W^{1,2}(\Omega) \ni u \mapsto (\int_{\Omega} |\nabla u|^2 + \int_{\partial\Omega} u^2)^{1/2}$ (we refer to Lemma 10 for instance), in what follows, we use

$$\|u\|_{W^{1,2}(\Omega)}^2 = \int_{\Omega} |\nabla u|^2 + \int_{\partial\Omega} u^2.$$

Up to multiplying B by a positive constant, (3.12) remains unchanged and we thus keep the notation B . We turn back to the proof of the Theorem.

The key point is to construct an admissible perturbation h , supported in ω^* (introduced in (3.8)) such that $\tilde{\mathcal{J}}(\beta^*)[h, h] > 0$. To this aim, we use (3.12) to say that it suffices to construct, for $\varepsilon > 0$ and $\delta > 0$ small enough, an admissible perturbation h supported in ω^* such that $\|h\|_{L^2(\partial\Omega)} = 1$, and that further satisfies

$$\|\dot{u}_{\beta^*}\|_{L^2(\Omega)} \leq \varepsilon \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \quad (3.25)$$

and

$$\|\dot{u}_{\beta^*}\|_{L^2(\partial\Omega)} \leq \sqrt{\delta} \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)} \quad (3.26)$$

Indeed, if h is non-zero and satisfies (3.25)-(3.26) then, according to (3.12), we obtain

$$\begin{aligned} \tilde{\mathcal{J}}(\beta^*)[h, h] &\geq A \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 - B \|\nabla \dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)} - C \|\dot{u}_{\beta^*}\|_{L^2(\partial\Omega)}^2 \\ &\geq (A - C\delta) \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 - B\varepsilon \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)}^2 \\ &\geq (A - C\delta) \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 - B\varepsilon \left(\|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 + \|\dot{u}_{\beta^*}\|_{L^2(\partial\Omega)}^2 \right) \\ &\geq (A - C\delta - B\varepsilon(1 + \delta)) \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2. \end{aligned}$$

In particular, to ensure that we have a positive right-hand side it suffices to pick ε and δ small enough.

It thus remains to prove that such a perturbation h exists. Let us highlight that we will obtain (3.25) and (3.26) by two different paths. We start with (3.26).

Regarding condition (3.26). Let us fix an arbitrary $\delta > 0$. To prove that we can choose an admissible perturbation h supported in ω^* such that $\|h\|_{L^2(\partial\Omega)} = 1$ and that satisfies (3.26), we rely on eigenvalues and eigenfunctions of a Robin-Steklov type operator. More precisely, we introduce the Hilbert basis $\{\phi_k\}_{k \in \mathbb{N}}$ of $L^2(\partial\Omega)$ given by

$$\begin{cases} -\Delta \phi_k = 0 & \text{in } \Omega, \\ \frac{\partial \phi_k}{\partial \nu} + \beta^* \phi_k = \sigma_k \phi_k & \text{on } \partial\Omega, \end{cases} \quad \text{where } 0 < \sigma_0 \leq \sigma_1 \leq \dots \leq \sigma_k \xrightarrow{k \rightarrow \infty} \infty \text{ and } \int_{\partial\Omega} \phi_k \phi_{k'} = \delta_{k,k'}. \quad (3.27)$$

We prove in Appendix C how such eigenpairs are defined; our arguments follow the classical [3, Section 11].

To see how this elements enable us to obtain (3.26), let us first observe that for any admissible h , the function $-hu_{\beta^*}$ belongs to $L^2(\partial\Omega)$, and thus expands as

$$-hu_{\beta^*} = \sum_{k=0}^{\infty} \alpha_k(h) \phi_k \quad (3.28)$$

where

$$\forall k \in \mathbb{N}, \alpha_k(h) = \int_{\partial\Omega} (-hu_{\beta^*}) \phi_k \quad (3.29)$$

Then, since \dot{u}_{β^*} solves

$$\begin{cases} -\Delta \dot{u}_{\beta^*} = 0 & \text{in } \Omega, \\ \frac{\partial \dot{u}_{\beta^*}}{\partial \nu} + \beta^* \dot{u}_{\beta^*} = \sum_{k=0}^{\infty} \alpha_k(h) \phi_k & \text{on } \partial\Omega, \end{cases}$$

we have

$$\dot{u}_{\beta^*} = \sum_{k=0}^{\infty} \frac{\alpha_k(h)}{\sigma_k} \phi_k, \quad (3.30)$$

which then allows us to compute

$$\|\dot{u}_{\beta^*}\|_{L^2(\partial\Omega)}^2 = \sum_{k=0}^{\infty} \frac{\alpha_k(h)^2}{\sigma_k^2} \quad \text{and} \quad \|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 + \int_{\partial\Omega} \beta^* \dot{u}_{\beta^*}^2 = \sum_{k=0}^{\infty} \frac{\alpha_k(h)^2}{\sigma_k}. \quad (3.31)$$

Therefore, since $0 \leq \beta^* \leq 1$, one has

$$\|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 \geq \sum_{k=0}^{\infty} \frac{\alpha_k(h)^2}{\sigma_k} - \int_{\partial\Omega} \dot{u}_{\beta^*}^2 \geq \sum_{k=0}^{\infty} \alpha_k(h)^2 \left(\frac{1}{\sigma_k} - \frac{1}{\sigma_k^2} \right).$$

Observe that if $K \in \mathbb{N}^*$ is chosen in such a way that

$$\forall k \in \llbracket 0, K-1 \rrbracket, \quad \alpha_k(h) = 0, \quad (3.32)$$

the previous estimates imply

$$\|\nabla \dot{u}_{\beta^*}\|_{L^2(\Omega)}^2 \geq \sum_{k=K}^{\infty} \alpha_k(h)^2 \left(\frac{1}{\sigma_k} - \frac{1}{\sigma_k^2} \right) = \sum_{k=K}^{\infty} \frac{\alpha_k(h)^2}{\sigma_k^2} (\sigma_k - 1) \geq (\sigma_K - 1) \|\dot{u}_{\beta^*}\|_{L^2(\partial\Omega)}^2. \quad (3.33)$$

Thus if we fix $K \in \mathbb{N}^*$ such that

$$\frac{1}{\sigma_K - 1} \leq \delta \quad (3.34)$$

and if we pick h such that (3.32) holds, we reach condition (3.26). We now prove that for any $K \in \mathbb{N}$, there exists an admissible perturbation $h \neq 0$ supported in ω^* such that (3.32) is satisfied. It will be convenient to observe that by the definition of ω^* and the fact that $\mathcal{H}^{d-1}(\omega^*) > 0$, there exists $\zeta > 0$ such that

$$\omega_{\zeta}^* := \{\zeta < \beta^* < 1 - \zeta\} \quad (3.35)$$

has positive measure. We will rather work on ω_{ζ}^* , for $\zeta > 0$ small enough, as any function $h \in L^\infty(\Omega)$ supported in ω_{ζ}^* that satisfies $\int_{\Omega} h = 0$ is then an admissible perturbation at β^* .

Let $K \in \mathbb{N}^*$ such that (3.34) is satisfied be fixed. According to the discussion above, we want to prove that there exists $h \in L^2(\partial\Omega)$ supported in ω_{ζ}^* such that

1. $\|h\|_{L^2(\partial\Omega)}^2 = 1$,
2. $\int_{\partial\Omega} h = 0$,
3. $\forall k \in \llbracket 0, K-1 \rrbracket, \alpha_k(h) = 0$.

Remember that $\mathcal{H}^{d-1}(\omega_{\zeta}^*) > 0$, so that $L^\infty(\omega_{\zeta}^*)$ is infinite dimensional. We introduce the following family of $(K+1)$ linear forms on $L^\infty(\omega_{\zeta}^*)$:

$$R : L^\infty(\omega_{\zeta}^*) \ni h \mapsto \int_{\omega_{\zeta}^*} h$$

$$T_k : L^\infty(\omega_\zeta^*) \ni h \mapsto \int_{\omega_\zeta^*} h u_{\beta^*} \phi_k$$

for all $k \in \llbracket 0, K-1 \rrbracket$. Since u_{β^*} belongs to $L^\infty(\partial\Omega)$, each T_k defines a continuous linear form and R is itself obviously continuous. As a consequence, the subspace

$$E_\delta := \ker(R) \cap \left(\bigcap_{k=0}^K \ker(T_k) \right) \quad (3.36)$$

is non-trivial: there exists, in particular, $h \in E_\delta \setminus \{0\}$. Up to normalising its L^2 -norm we can assume that $\|h\|_{L^2(\omega_\zeta^*)} = 1$. It suffices to extend h by 0 to $\partial\Omega$. As h is supported in ω_ζ^* and is bounded, it is an admissible perturbation at β^* that satisfies all the required conditions.

Satisfying both conditions (3.25) and (3.26). Thus, for a fixed $\delta > 0$ (not necessarily small), every $h \in E_\delta$ satisfies (3.26), where E_δ is defined in (3.36). We have also fixed $\varepsilon > 0$. Let us show that there exists $h \in E_\delta$ that satisfies (3.25), which suffices to conclude the proof. In other words, we will prove that

$$\forall C > 0, \quad \exists h \in E_\delta, \quad \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} > C \|\dot{u}_{\beta^*}\|_{L^2(\Omega)}. \quad (3.37)$$

To prove (3.37), let us argue by contradiction, assuming the existence of $C > 0$ such that

$$\forall h \in E_\delta, \quad \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \leq C \|\dot{u}_{\beta^*}\|_{L^2(\Omega)}. \quad (3.38)$$

In what follows, we will rather denote \dot{u}_{β^*} by $\dot{u}_{\beta^*}[h]$ to emphasize the dependency of this function in h . Let us introduce

$$X_\delta := \{\dot{u}_{\beta^*}[h], h \in E_\delta\}.$$

X_δ is a subspace of $W^{1,2}(\Omega)$ and since the map $E_\delta \ni h \mapsto \dot{u}_{\beta^*}[h]$ is an injection, X_δ is infinite dimensional. Consequently, there exists an $L^2(\Omega)$ -orthonormal family $\{v_k\}_{k \in \mathbb{N}}$ in X_δ . In particular, for any $k \in \mathbb{N}$, $\|v_k\|_{L^2(\Omega)} = 1$. Furthermore, according to the Parseval inequality, one has

$$v_k \xrightarrow[k \rightarrow \infty]{} 0 \text{ weakly in } L^2(\Omega).$$

However, should (3.38) hold, the family $\{v_k\}_{k \in \mathbb{N}}$ would be uniformly bounded in $W^{1,2}(\Omega)$ and thus, by the Rellich-Kondrachov theorem, converge strongly in L^2 to a closure point v_∞ (up to a subsequence). Since we already know it converges weakly to 0, one must have $v_\infty = 0$ on the one hand and $\|v_\infty\|_{L^2(\Omega)} = 1$ on the other hand, leading to a contradiction. The conclusion follows: there necessarily exists h such that (3.25)-(3.26) holds.

The proof of the Theorem is now complete.

3.3 Proof of Theorem 5

We prove each point of Theorem 5 separately. Once again, notational convenience leads us to dropping the $\partial\Omega$ subscript, and to just writing

$$\mathcal{J}(\beta) = \int_{\partial\Omega} j(u_\beta).$$

Proof of (i). Before we get to the core of the Theorem, let us point out a consequence of the expression of the first order derivative of the criterion given in (3.6). We set, for any $\beta \in \mathcal{B}(\partial\Omega)$,

$$\Phi_\beta := u_\beta p_\beta$$

where p_β is the solution of (3.4). Therefore, for any admissible perturbation h at a given β , there holds

$$\dot{\mathcal{J}}(\beta)[h] = - \int_{\partial\Omega} h \Phi_\beta. \quad (3.39)$$

Let β^* be a solution of $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$. Since we must have

$$-\dot{\mathcal{J}}(\beta^*)[h] = \int_{\partial\Omega} \Phi_{\beta^*} h \leq 0$$

for any admissible perturbation h at β^* , there exists a real number λ (necessarily positive as $\Phi_{\beta^*} > 0$) such that

1. $\{0 < \beta^* < 1\} \subset \{\Phi_{\beta^*} = \lambda\}$,
2. $\{\beta^* = 1\} \subset \{\Phi_{\beta^*} \geq \lambda\}$,
3. $\{\beta^* = 0\} \subset \{\Phi_{\beta^*} \leq \lambda\}$.

We shall now prove that for any solution β^* of $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$, such optimality conditions imply that

$$\mathcal{H}^{d-1}(\{\beta^* = 0\}) = 0, \quad (3.40)$$

which necessarily yields

$$\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0, \quad (3.41)$$

the required conclusion. Indeed, this follows from the volume constraint $\int_{\partial\Omega} \beta^* = V_0$.

To prove (3.40), we first consider the equation satisfied by the function $\Phi_{\beta^*} = u_{\beta^*} p_{\beta^*}$. By direct computation, we have

$$\nabla \Phi_{\beta^*} = u_{\beta^*} \nabla p_{\beta^*} + p_{\beta^*} \nabla u_{\beta^*}, -\Delta \Phi_{\beta^*} = -p_{\beta^*} \Delta u_{\beta^*} - u_{\beta^*} \Delta p_{\beta^*} - 2 \nabla u_{\beta^*} \cdot \nabla p_{\beta^*}. \quad (3.42)$$

Let us set

$$B := -2 \frac{\nabla p_{\beta^*}}{p_{\beta^*}}$$

and

$$V := \frac{f}{u_{\beta^*}} + 2 \frac{|\nabla p_{\beta^*}|^2}{p_{\beta^*}^2}.$$

We then note that

$$\nabla u_{\beta^*} = \frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{u_{\beta^*}}{p_{\beta^*}} \nabla p_{\beta^*} = \frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{\Phi_{\beta^*}}{p_{\beta^*}^2} \nabla p_{\beta^*}.$$

Plugging this expression into (3.42) yields

$$\begin{aligned} -\Delta \Phi_{\beta^*} &= f p_{\beta^*} - 2 \left(\frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{\Phi_{\beta^*}}{p_{\beta^*}^2} \nabla p_{\beta^*} \right) \cdot \nabla p_{\beta^*} \\ &= \Phi_{\beta^*} \frac{f}{u_{\beta^*}} + 2 \Phi_{\beta^*} \frac{|\nabla p_{\beta^*}|^2}{p_{\beta^*}^2} + \langle \nabla \Phi_{\beta^*}, B \rangle \end{aligned}$$

$$= V\Phi_{\beta^*} + \langle \nabla \Phi_{\beta^*}, B \rangle.$$

Since $\inf_{\Omega} \min\{u_{\beta^*}, p_{\beta^*}\} > 0$ it follows that

$$\inf_{\Omega} \Phi_{\beta^*} > 0. \quad (3.43)$$

As f satisfies (\mathbf{H}_f) , one has $V \geq 0$ according to the definition of V . Thus, Φ_{β^*} satisfies

$$-\Delta \Phi_{\beta^*} - \langle \nabla \Phi_{\beta^*}, B \rangle = V\Phi_{\beta^*} \geq 0. \quad (3.44)$$

In the last inequality we used the fact that $V \geq 0$ and (3.43). By the strong maximum principle we necessarily have either

$$\min_{\Omega} \Phi_{\beta^*} = \min_{\partial\Omega} \Phi_{\beta^*} \text{ and this minimum is never reached inside } \Omega \quad (3.45)$$

or

$$\Phi_{\beta^*} \text{ is constant in } \Omega. \quad (3.46)$$

We claim that (3.46) can not occur: if Φ_{β^*} is constant in Ω then (3.44) implies $V\Phi_{\beta^*} = 0$ in Ω whence, since $\Phi_{\beta^*} > 0$ in Ω , $V \equiv 0$ in Ω . In particular, $f \equiv 0$, a contradiction. Thus (3.45) holds. According to the maximum principle of Hopf, if we pick a minimum point $x^* \in \partial\Omega$ of Φ_{β^*} , there holds

$$\frac{\partial \Phi_{\beta^*}}{\partial \nu}(x^*) < 0. \quad (3.47)$$

We fix this point.

We now argue by contradiction and assume that $\mathcal{H}^{d-1}(\{\beta^* = 0\}) > 0$. Given that β^* satisfies first order optimality conditions, we can choose such a minimum point x^* that satisfies

$$x^* \in \{\beta^* = 0\}.$$

We finally compute the boundary conditions on Φ_{β^*} . Since

$$\frac{\partial \Phi_{\beta^*}}{\partial \nu} = \frac{\partial u_{\beta^*}}{\partial \nu} p_{\beta^*} + \frac{\partial p_{\beta^*}}{\partial \nu} u_{\beta^*} = -2\beta^* \Phi_{\beta^*} + j'(u_{\beta^*}) u_{\beta^*},$$

we get

$$\frac{\partial \Phi_{\beta^*}}{\partial \nu} + 2\beta^* \Phi_{\beta^*} = j'(u_{\beta^*}) u_{\beta^*}. \quad (3.48)$$

Going back to our minimum point $x^* \in \{\beta^* = 0\}$ and to (3.47) we should have $\frac{\partial \Phi_{\beta^*}}{\partial \nu}(x^*) \leq 0$ on the one hand, and $\frac{\partial \Phi_{\beta^*}}{\partial \nu} = -2\beta^* \Phi_{\beta^*} + j'(u_{\beta^*}) u_{\beta^*} = j'(u_{\beta^*}) u_{\beta^*} > 0$ on the other hand. The last condition comes from the fact that j satisfies (\mathbf{H}_j) . This is a contradiction, and the conclusion follows.

Proof of (ii). Let us use the same notation as in the proof of Theorem 3 in section 3.2. Let β^* denote a solution of Problem $(\mathbf{P}_{\min, \partial\Omega, \mathcal{B}})$, whose existence is guaranteed by Lemma 12. We use the Hilbert basis $\{\phi_k\}_{k \in \mathbb{N}}$ of $L^2(\partial\Omega)$ given by (3.27), associated to the sequence of eigenvalues $\{\sigma_k\}_{k \in \mathbb{N}}$. According to the first item of Theorem 5, one has $\mathcal{H}^{d-1}(\{0 < \beta^* < 1\}) > 0$.

Assume the existence of $C > 0$ such that

$$j''(u) \leq -C j'(0) = -C \sup_{u \in [0, U_0(f)]} j'(u).$$

Let us argue by contradiction, assuming that $\mathcal{H}^{d-1}(\{\beta^* = 1\}) = 0$. Let $h \in L^\infty(\partial\Omega)$ be an admissible perturbation of β^* in $\mathcal{B}(\partial\Omega)$. The function $-hu_{\beta^*} \in L^2(\partial\Omega)$ expands as

$$-hu_{\beta^*} = \sum_{k=0}^{\infty} \alpha_k(h) \phi_k \quad \text{with} \quad \alpha_k(h) = \int_{\partial\Omega} (-hu_{\beta^*}) \phi_k$$

for every $k \in \mathbb{N}$. Then, we also have

$$\dot{u}_{\beta^*} = \sum_{k=0}^{\infty} \frac{\alpha_k(h)}{\sigma_k} \phi_k.$$

Recall that, according to (3.7), we have

$$\begin{aligned} \ddot{\mathcal{J}}(\beta^*)[h, h] &= -2 \int_{\partial\Omega} h \dot{u}_{\beta^*} p_{\beta^*} + \int_{\partial\Omega} j''(u_{\beta^*}) (\dot{u}_{\beta^*})^2 \\ &= -2 \int_{\partial\Omega} \Psi_{\beta^*} h u_{\beta^*} \dot{u}_{\beta^*} + \int_{\partial\Omega} j''(u_{\beta^*}) (\dot{u}_{\beta^*})^2 \end{aligned}$$

where $\Psi_{\beta^*} := \frac{p_{\beta^*}}{u_{\beta^*}}$. Observe first that $p_{\beta^*} \leq z_{\beta^*} \sup_{u \in [0, U_0]} j'(u)$ where z_{β^*} is the unique solution to

$$\begin{cases} -\Delta z_{\beta^*} = 0 & \text{in } \Omega, \\ \frac{\partial z_{\beta^*}}{\partial \nu} + \beta^* z_{\beta^*} = 1 & \text{on } \partial\Omega. \end{cases}$$

Indeed, the function $P := p_{\beta^*} - z_{\beta^*} \sup_{u \in [0, U_0]} j'(u)$ is harmonic, and therefore reaches its maximal value on $\partial\Omega$. Furthermore, according to the Hopf maximum principle, $\partial_\nu P > 0$ at this point, which yields easily that $P(\cdot) < 0$ on $\partial\Omega$.

We isolate the following result, which follows from exactly the same arguments as in Lemma 11.

Lemma 15. *Let Ω be a bounded open set of \mathbb{R}^n such that $\partial\Omega$ is \mathcal{C}^2 . For $\beta \in \mathcal{B}(\partial\Omega)$, we define*

$$\mathcal{K}(\beta) = \frac{\max_{\overline{\Omega}} z_\beta}{\min_{\overline{\Omega}} u_\beta}.$$

One has

$$K := \sup_{\beta \in \mathcal{B}(\partial\Omega)} \mathcal{K}(\beta) < +\infty.$$

Using the notations of this Lemma, we thus have

$$|\Psi_{\beta^*}| \leq \frac{\max p_{\beta^*}}{\min u_{\beta^*}} \leq \sup_{[0; U_0(f)]} j'(u) K. \quad (3.49)$$

Let us introduce

$$h_{\alpha_1, \alpha_2} = \frac{\alpha_1 \phi_1 + \alpha_2 \phi_2}{u_{\beta^*}} \quad \text{with} \quad \alpha_1 = \int_{\partial\Omega} \frac{\phi_2}{u_{\beta^*}} \quad \text{and} \quad \alpha_2 = - \int_{\partial\Omega} \frac{\phi_1}{u_{\beta^*}},$$

so that $\int_{\Omega} h_{\alpha_1, \alpha_2} = 0$. Notice that h_{α_1, α_2} is admissible since $\mathcal{H}^{d-1}(\{\beta^* = 1\}) = \mathcal{H}^{d-1}(\{\beta^* = 0\}) = 0$. Indeed, consider a sequence $\{\varepsilon_n\}_{n \in \mathbb{N}}$ converging to 0 as $n \rightarrow \infty$. Defining, for any $n \in \mathbb{N}$,

$$h_n := \mathbb{1}_{\{2\|h\|_{L^\infty(\Omega)} \varepsilon_n < \beta^* < 1-2\|h\|_{L^\infty(\Omega)} \varepsilon_n\}} \left(h - \int_{\{2\|h\|_{L^\infty(\Omega)} \varepsilon_n < \beta^* < 1-2\|h\|_{L^\infty(\Omega)} \varepsilon_n\}} h \right)$$

it is clear that $h_n \xrightarrow{n \rightarrow \infty} h$ in $L^2(\Omega)$. Furthermore, for any $n \in \mathbb{N}$, $\int_{\Omega} h_n = 0$. Finally, for any $n \in \mathbb{N}$,

$$\beta^* + \varepsilon_n h_n \geq \mathbb{1}_{\{2\|h\|_{L^\infty(\Omega)}\varepsilon_n < \beta^* < 1-2\|h\|_{L^\infty(\Omega)}\varepsilon_n\}} (2\|h\|_{L^\infty(\Omega)}\varepsilon_n - 2\|h\|_{L^\infty(\Omega)}\varepsilon_n) = 0$$

and, similarly,

$$\beta^* + \varepsilon_n h_n \leq 1.$$

We deduce that h is admissible in the sense of Definition 13.

From the assumption on j and the Cauchy-Schwarz inequality, we have

$$\begin{aligned} \ddot{\mathcal{J}}(\beta^*)[h_{\alpha_1, \alpha_2}, h_{\alpha_1, \alpha_2}] &\leq \|\Psi_{\beta^*}\|_{\infty} \int_{\partial\Omega} |h_{\alpha_1, \alpha_2} u_{\beta^*}| |\dot{u}_{\beta^*}| - C \sup_{u \in [0, U_0]} j'(u) \int_{\partial\Omega} (\dot{u}_{\beta^*})^2 \\ &\leq \|\Psi_{\beta^*}\|_{\infty} \left(\int_{\partial\Omega} h_{\alpha_1, \alpha_2}^2 u_{\beta^*}^2 \right)^{1/2} \left(\int_{\partial\Omega} (\dot{u}_{\beta^*})^2 \right)^{1/2} - C \sup_{u \in [0, U_0(f)]} j'(u) \int_{\partial\Omega} (\dot{u}_{\beta^*})^2 \\ &\leq \sup_{u \in [0, U_0(f)]} j'(u) \left[K (\alpha_1^2 + \alpha_2^2)^{1/2} \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right)^{1/2} - C \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right) \right] \\ &\leq \sup_{u \in [0, U_0(f)]} j'(u) \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right)^{1/2} \left[K (\alpha_1^2 + \alpha_2^2)^{1/2} - C \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right)^{\frac{1}{2}} \right] \\ &\leq \sup_{u \in [0, U_0(f)]} j'(u) \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right)^{1/2} \left[K (\alpha_1^2 + \alpha_2^2)^{1/2} - C \left(\frac{\alpha_1^2 + \alpha_2^2}{\sigma_2^2} \right)^{\frac{1}{2}} \right] \\ &\leq \sup_{u \in [0, U_0(f)]} j'(u) (\alpha_1^2 + \alpha_2^2)^{1/2} \left(\frac{\alpha_1^2}{\sigma_1^2} + \frac{\alpha_2^2}{\sigma_2^2} \right)^{1/2} \left(K - \frac{C}{\sigma_2} \right). \end{aligned}$$

Furthermore, according to the Courant-Fisher principle, one has

$$\begin{aligned} \sigma_2 &= \min_{\substack{E_2 \subset W^{1,2}(\Omega) \\ \text{subspace of dim. 2}}} \max_{\substack{v \in E_2 \\ v \neq 0}} \frac{\int_{\Omega} |\nabla v|^2 + \int_{\partial\Omega} \beta^* v^2}{\int_{\partial\Omega} v^2} \\ &\leq \min_{\substack{E_2 \subset W^{1,2}(\Omega) \\ \text{subspace of dim. 2}}} \max_{\substack{v \in E_2 \\ v \neq 0}} \frac{\int_{\Omega} |\nabla v|^2 + \int_{\partial\Omega} v^2}{\int_{\partial\Omega} v^2} = \Lambda_2(\Omega) > 0, \end{aligned}$$

where $\Lambda_2(\Omega)$ denotes the second Robin-Steklov eigenvalue of the domain:

$$\begin{cases} -\Delta \varphi_k = 0 & \text{in } \Omega, \\ \frac{\partial \varphi_k}{\partial \nu} + \varphi_k = \Lambda_k(\Omega) \varphi_k & \text{on } \partial\Omega. \end{cases}$$

Therefore, if C is such that $\Lambda_2(\Omega)K < C$, one has $\ddot{\mathcal{J}}(\beta^*)[h_{\alpha_1, \alpha_2}, h_{\alpha_1, \alpha_2}] < 0$ and thus

$$\mathcal{J}(\beta^* + \varepsilon h_{\alpha_1, \alpha_2}) - \mathcal{J}(\beta^*) = \frac{\varepsilon^2}{2} \ddot{\mathcal{J}}(\beta^*)[h_{\alpha_1, \alpha_2}, h_{\alpha_1, \alpha_2}] + o(\varepsilon^2) < 0,$$

whenever $\varepsilon > 0$ is chosen small enough. This is in contradiction with the optimality of β^* , whence the result.

4 Proof of Theorem 8

We recall that we work with the energy functional

$$\mathcal{F}(\beta) := \int_{\Omega} f u_{\beta}.$$

- (i) **Proof of (i):** Let $\beta \in \mathcal{B}(\partial\Omega)$ and h denote an admissible perturbation at β (see Definition 13). By mimicking the computations of Section 3.1, one gets

$$\dot{\mathcal{J}}(\beta)[h] = - \int_{\partial\Omega} h u_{\beta}^2 \quad \text{and} \quad \ddot{\mathcal{J}}(\beta^*)[h, h] = -2 \int_{\partial\Omega} h u_{\beta^*} \dot{u}_{\beta^*},$$

where \dot{u}_{β} solves

$$\begin{cases} -\Delta \dot{u}_{\beta} = 0 & \text{in } \Omega, \\ \frac{\partial \dot{u}_{\beta}}{\partial \nu} + \beta \dot{u}_{\beta} = -h u_{\beta^*} & \text{on } \partial\Omega, \end{cases}$$

Let us use the Hilbert basis $\{\phi_k\}_{k \in \mathbb{N}}$ of $L^2(\partial\Omega)$ given by (3.27), associated with the sequence of eigenvalues $\{\sigma_k\}_{k \in \mathbb{N}}$. The function $-h u_{\beta} \in L^2(\partial\Omega)$ expands as

$$-h u_{\beta} = \sum_{k=0}^{\infty} \alpha_k(h) \phi_k \quad \text{with} \quad \alpha_k(h) = \int_{\partial\Omega} (-h u_{\beta}) \phi_k \quad \text{for every } k \in \mathbb{N}$$

and we also have

$$\dot{u}_{\beta} = \sum_{k=0}^{\infty} \frac{\alpha_k(h)}{\sigma_k} \phi_k.$$

As a consequence

$$\ddot{\mathcal{J}}(\beta^*)[h, h] = 2 \sum_{k=0}^{\infty} \frac{\alpha_k(h)^2}{\sigma_k}$$

and we easily infer that \mathcal{J} is strictly convex.

Let β^* be a solution of Problem (\mathbf{Q}_{\max}) . Let us assume by contradiction that the set $\mathcal{I} := \{0 < \beta^* < 1\}$ has positive measure. Let $\tilde{\beta}$ denote any element of $\mathcal{B}(\partial\Omega)$ equal to β on $\Omega \setminus \mathcal{I}$, such that $\tilde{\beta} \neq \beta$ a.e. on \mathcal{I} and $\int_{\mathcal{I}} \beta^* = \int_{\mathcal{I}} \tilde{\beta}$. Then, $\beta^* + \varepsilon h$ where $h = \tilde{\beta} - \beta^*$ is admissible and

$$\mathcal{F}(\beta^* + \varepsilon h) - \mathcal{F}(\beta^*) = \frac{\varepsilon^2}{2} \ddot{\mathcal{J}}(\beta^*)[h, h] + o(\varepsilon^2) > 0$$

whenever ε is small enough. We have thus reached a contradiction and it follows that $\mathcal{H}^{d-1}(\mathcal{I}) = 0$.

- (iii) **Proof of (ii)** According to the analysis above, the mapping $\mathcal{B}(\partial\Omega) \ni \beta \mapsto \mathcal{F}(\beta)$ is convex. Since we are dealing with a minimization problem, we get that β^* solves Problem (\mathbf{Q}_{\min}) if, and only if $\dot{\mathcal{J}}(\beta)[h] \geq 0$ for every admissible perturbation h . By using the expression of $\dot{\mathcal{J}}(\beta)[h]$ obtained previously, it is standard that the first order optimality conditions read as follows: there exists a positive real number λ such that

1. $\{0 < \beta^* < 1\} \subset \{u_{\beta^*}^2 = \lambda\},$
2. $\{\beta^* = 1\} \subset \{u_{\beta^*}^2 \geq \lambda\},$
3. $\{\beta^* = 0\} \subset \{u_{\beta^*}^2 \leq \lambda\}.$

Since these conditions are sufficient and necessary, it follows that β^* solves Problem (\mathbf{Q}_{\min}) if and only if u_{β^*} solves the overdetermined system given by (\mathbf{E}_{β}) complemented by the optimality conditions above. Therefore, to conclude, it is enough to check that the particular function β^* given by (1.13) satisfies the optimality system above.

Observe first that, according to the Hopf maximum principle, the function v_{Ω} reaches its minimal value a.e. on the boundary of Ω . Therefore, one has $\partial_{\nu} v_{\Omega} < 0$ on $\partial\Omega$, meaning that $\beta^* > 0$ on $\partial\Omega$. Moreover, because of the assumptions on V_0 , one has

$$\beta^* < V_0 \frac{\|\partial_{\nu} v_{\Omega}\|_{L^{\infty}(\partial\Omega)}}{\int_{\partial\Omega} \partial_{\nu} v_{\Omega}} = 1 \text{ a.e. in } \partial\Omega \quad \text{and} \quad \int_{\Omega} \beta^* = V_0,$$

so that $\beta^* \in \mathcal{B}(\partial\Omega)$ and $\{0 < \beta^* < 1\} = \partial\Omega$.

We set

$$\lambda = \frac{-\int_{\partial\Omega} \partial_{\nu} v_{\Omega}}{V_0}, \quad \text{and} \quad u_{\beta^*}^{\lambda} = \lambda + v_{\Omega}.$$

Straightforward computations show that $u_{\beta^*}^{\lambda}$ coincide with the solution u_{β^*} of (\mathbf{E}_{β}) and that

$$\{0 < \beta^* < 1\} = \{u_{\beta^*}^2 = \lambda\}.$$

The expected conclusion follows.

- (iii) **Proof of (iii):** Using the same arguments as in (ii), since $\{0 < \beta^*_{V_0} < 1\} = \partial\Omega$, one sees that $\beta^*_{V_0}$ solves Problem (\mathbf{Q}_{\min}) if, and only if there exists $\lambda \geq 0$ such that $u_{\beta^*_{V_0}}$ solves the overdetermined system

$$\begin{cases} -\Delta u_{\beta^*_{V_0}} = 1 & \text{in } \Omega, \\ u_{\beta^*_{V_0}} = \lambda & \text{on } \partial\Omega \\ \frac{\partial u_{\beta^*_{V_0}}}{\partial \nu} = -\frac{V_0}{|\partial\Omega|} \lambda & \text{on } \partial\Omega. \end{cases}$$

Setting $v_{\beta^*_{V_0}} = u_{\beta^*_{V_0}} - \lambda$, this overdetermined system is equivalent with

$$\begin{cases} -\Delta v_{\beta^*_{V_0}} = 1 & \text{in } \Omega, \\ v_{\beta^*_{V_0}} = 0 & \text{on } \partial\Omega \\ \frac{\partial v_{\beta^*_{V_0}}}{\partial \nu} = c_{V_0} & \text{on } \partial\Omega. \end{cases} \quad \text{with} \quad c_{V_0} = -\frac{V_0}{|\partial\Omega|} \lambda < 0.$$

We can now apply Serrin's theorem (see [30] for the original proof, [32] for a simpler proof that holds in the case $f \equiv 1$, and [29] for a survey of the proofs of this theorem), which fully characterises such overdetermined elliptic problems: this system has a solution if, and only if, Ω is a ball.

This concludes the proof of the Theorem.

Appendix

A Convergence of u_α^Γ towards v^Γ as $\alpha \rightarrow +\infty$

We investigate in this section the asymptotic behaviour of u_α^Γ as $\alpha \rightarrow +\infty$, in the simple case where $g_0 = g_0(x, u)$ does not depend on u . With a slight abuse of notation, we write $g_0 = g_0(x)$.

Proposition 16. *Let Ω be a connected bounded open set of class \mathcal{C}^1 . Let $\Gamma \subset \partial\Omega$, with $\mathcal{H}^{d-1}(\Gamma) > 0$, and let $g_0 \in L^2(\Omega)$. Assume there exists $\sigma_0 > 0$ such that $A \geq \sigma_0 \text{Id}$ a.e. in Ω in the sense of bilinear forms. The family $(u_\alpha^\Gamma)_{\alpha>0}$ converges to v^Γ , weakly in $W^{1,2}(\Omega)$ and strongly in $L^2(\Omega)$.*

Proof. It should be noted that the regularity assumptions on Ω is central as it guarantees the compactness of the trace operator $\text{Tr} : W^{1,2}(\Omega) \rightarrow L^2(\partial\Omega)$. We write Tr_Γ for the operator that maps u to $\mathbb{1}_\Gamma \text{Tr} u$. Multiplying the first equation of (1.1) by u_α^Γ and integrating by parts gives

$$\sigma_0 \int_\Omega |\nabla u_\alpha^\Gamma|^2 + \alpha \int_\Gamma (u_\alpha^\Gamma)^2 \leq \|g\|_{L^2(\Omega)} \|u_\alpha^\Gamma\|_{L^2(\Omega)}.$$

By continuity of the trace operator and by Lemma 10, there exists $C_0 = C_0(\sigma_0, \Omega) > 0$ such that

$$C_0 \|u_\alpha^\Gamma\|_{W^{1,2}(\Omega)}^2 \leq \sigma_0 \int_\Omega |\nabla u_\alpha^\Gamma|^2 + \alpha \int_\Gamma (u_\alpha^\Gamma)^2 \leq \|g\|_{L^2(\Omega)} \|u_\alpha^\Gamma\|_{W^{1,2}(\Omega)} \quad (\text{A.1})$$

Let $(\alpha_n)_{n \in \mathbb{N}}$ be an increasing sequence of positive number such that $\lim_{n \rightarrow +\infty} \alpha_n = +\infty$. As the family $\{u_{\alpha_n}^\Gamma\}_{n \in \mathbb{N}}$ is bounded in $W^{1,2}(\Omega)$ by (A.1), the Rellich-Kondrachov theorem, ensures that it converges, up to a subsequence, to a certain $\bar{u} \in W^{1,2}(\Omega)$ weakly in $W^{1,2}(\Omega)$ and strongly in $L^2(\Omega)$. With a slight abuse of notation, this subsequence is still written $\{u_{\alpha_n}^\Gamma\}_{n \in \mathbb{N}}$. Since the trace operator is compact, the sequence $\{\text{Tr}_\Gamma u_{\alpha_n}^\Gamma\}_{n \in \mathbb{N}}$ converges to $\text{Tr}_\Gamma \bar{u}$ in $L^2(\Gamma)$. As the sequence $\{\alpha_n \int_\Gamma (u_{\alpha_n}^\Gamma)^2\}_{n \in \mathbb{N}}$ is bounded as we must have $\text{Tr}_\Gamma \bar{u} = 0$ in $L^2(\Gamma)$.

Let us introduce the space $W_\Gamma^{1,2}(\Omega)$ as the subspace of functions φ in $W^{1,2}(\Omega)$ whose trace vanishes on Γ . Recall that u_α^Γ solves the minimization problem

$$\min_{u \in W^{1,2}(\Omega)} \mathcal{F}_\alpha(u) \quad \text{where} \quad \mathcal{F}_\alpha(u) = \frac{1}{2} \int_\Omega \langle A \nabla u, \nabla u \rangle + \frac{\alpha}{2} \int_\Gamma u^2 - \int_\Omega f u.$$

By minimality, one has, for any $n \in \mathbb{N}$,

$$\min_{u \in W^{1,2}(\Omega)} \mathcal{F}_{\alpha_n}(u) = \mathcal{F}_{\alpha_n}(u_{\alpha_n}^\Gamma) \leq \min_{u \in W_\Gamma^{1,2}(\Omega)} \mathcal{F}_{\alpha_n}(u) = \min_{u \in W_\Gamma^{1,2}(\Omega)} \mathcal{F}_0(u) = \mathcal{F}_0(v^\Gamma),$$

where v^Γ solves Problem (1.2) with mixed Dirichlet-Neumann boundary conditions. Furthermore, as A is uniformly positive in the sense of bilinear forms, the map $W^{1,2}(\Omega) \ni u \mapsto \int_\Omega \langle A \nabla u, \nabla u \rangle$ is convex, and, so, weakly lower semi-continuous. Hence, we have

$$\begin{aligned} \liminf_{n \rightarrow +\infty} \frac{1}{2} \int_\Omega \langle A \nabla u_{\alpha_n}^\Gamma, \nabla u_{\alpha_n}^\Gamma \rangle + \frac{\alpha_n}{2} \int_\Gamma (u_{\alpha_n}^\Gamma)^2 - \int_\Omega f u_{\alpha_n}^\Gamma &\geq \frac{1}{2} \int_\Omega \langle A \nabla \bar{u}, \nabla \bar{u} \rangle - \int_\Omega f \bar{u} \\ &\geq \min_{u \in W_\Gamma^{1,2}(\Omega)} \mathcal{F}_0(u). \end{aligned}$$

Combining both inequalities above, it follows that

$$\min_{u \in W_\Gamma^{1,2}(\Omega)} \frac{1}{2} \int_\Omega \sigma |\nabla u|^2 - \int_\Omega f u = \frac{1}{2} \int_\Omega \sigma |\nabla \bar{u}|^2 - \int_\Omega f \bar{u},$$

and by uniqueness of the minimiser of this last problem, we obtain that $\bar{u} = v^\Gamma$. Thus, the sequence $\{u_{\alpha_n}^\Gamma\}_{n \in \mathbb{N}}$ has a unique closure point v^Γ . It follows that the entire sequence $\{u_{\alpha_n}^\Gamma\}_{n \in \mathbb{N}}$ converges to v^Γ , weakly in $W^{1,2}(\Omega)$ and strongly in $L^2(\Omega)$. \square

B Proof of results stated in section 2

Proof of Lemma 10. Of course, it suffices to prove that there exists $C > 0$ such that

$$\forall \beta \in \mathcal{B}(\partial\Omega), \forall v \in W^{1,2}(\Omega), C\|v\|_{L^2(\Omega)}^2 \leq \int_{\Omega} |\nabla v|^2 + \int_{\partial\Omega} \beta v^2. \quad (\text{B.1})$$

To prove (B.1) we argue by contradiction: should no such constant C exist, there exists a sequence $\{v_n, \beta_n\}_{n \in \mathbb{N}} \in (W^{1,2}(\Omega) \times \mathcal{B}(\partial\Omega))^{\mathbb{N}}$ such that

$$\|v_n\|_{L^2(\Omega)}^2 = 1 \quad \text{and} \quad \int_{\Omega} |\nabla v_n|^2 + \int_{\partial\Omega} \beta_n v_n^2 < \frac{1}{n}$$

By the Rellich-Kondrachov theorem, there exists $\bar{v} \in W^{1,2}(\Omega)$ such that $\{v_n\}_{n \in \mathbb{N}}$ converges, up to a subsequence, to \bar{v} , weakly in $W^{1,2}(\Omega)$ and strongly in $L^2(\Omega)$. Denoting this subsequence by $\{v_n\}_{n \in \mathbb{N}}$ with a slight abuse of notation, it follows that

$$\|\bar{v}\|_{L^2(\Omega)}^2 = 1 \quad \text{and} \quad \int_{\Omega} |\nabla \bar{v}|^2 \leq \liminf_{n \rightarrow +\infty} \int_{\Omega} |\nabla v_n|^2 = 0.$$

Thus \bar{v} is a positive constant, say \bar{v}_0 . On the other hand, since $\mathcal{B}(\partial\Omega)$ is compact for the weak $L^\infty - *$ topology, there exists $\beta \in \mathcal{B}(\Omega)$ such that, still up to a subsequence,

$$\beta_n \xrightarrow[n \rightarrow \infty]{} \beta.$$

By compactness of the trace operator, $\{v_n\}_{n \in \mathbb{N}}$ converges strongly, in $L^2(\partial\Omega)$, to \bar{v} . Thus, passing to the limit in

$$\int_{\Omega} |\nabla v_n|^2 + \int_{\partial\Omega} \beta_n v_n^2 < \frac{1}{n}$$

yields

$$\int_{\Omega} |\nabla \bar{v}|^2 + \int_{\partial\Omega} \beta \bar{v}_n^2 = 0.$$

Since $\int_{\partial\Omega} \beta = V_0$, this is impossible as \bar{v} is a positive constant. □

Proof of Lemma 11. The proof of this result relies results for Neumann boundary conditions: consider the problem

$$\begin{cases} -\Delta u_{f,g}^N = f & \text{in } \Omega, \\ \frac{\partial u_{f,g}^N}{\partial \nu} = g & \text{on } \partial\Omega. \end{cases} \quad (\text{B.2})$$

The following regularity holds from, for example, [31, Theorem 4.4]: assume Ω has a \mathcal{C}^2 boundary. Let $f \in L^q(\Omega)$, $g \in W^{-\frac{1}{q},q}(\partial\Omega)$ satisfy the compatibility condition

$$\int_{\Omega} f = \int_{\partial\Omega} g.$$

Then, there exists a $W^{1,q}(\Omega)$ solution $u_{f,g}$ of (B.2). Furthermore, for any such solution,

$$\|\nabla u_{f,g}\|_{L^q(\Omega)} \leq C \left(\|f\|_{L^q(\Omega)} + \|g\|_{W^{-\frac{1}{q},q}(\partial\Omega)} \right). \quad (\text{B.3})$$

We turn back to the proof of Lemma 11: let us consider, for any $\beta \in \mathcal{B}(\partial\Omega)$, the energy functional

$$\mathcal{E}_{\beta,f} : W^{1,2}(\Omega) \ni u \mapsto \frac{1}{2} \int_{\Omega} |\nabla u|^2 + \frac{1}{2} \int_{\partial\Omega} \beta u^2 - \int_{\Omega} f u. \quad (\text{B.4})$$

By Lemma 10 this energy functional is coercive. As a consequence, it admits a minimiser. It is immediate to see that uniqueness holds for Equation (\mathbf{E}_{β}) . Thus, we have obtained a unique solution $u_{\beta} \in W^{1,2}(\Omega)$.

From the Sobolev embeddings $W^{1,2}(\Omega) \hookrightarrow W^{\frac{1}{2},2}(\partial\Omega) \hookrightarrow L^{2^*_{\partial\Omega}}(\partial\Omega)$ where $2^*_{\partial\Omega} = \frac{2(n-1)}{n-2}$ (if $n \geq 2$, the case $n = 1$ being trivial) and the fact that $\beta \in L^{\infty}(\partial\Omega)$ we obtain that u_{β} solves a Neumann problem with Neumann data $g := -\beta u_{\beta} \in L^{2^*_{\partial\Omega}}(\partial\Omega)$. Furthermore, for the constant C given by Lemma 10,

$$\|g\|_{L^{2^*_{\partial\Omega}}(\partial\Omega)} \leq \frac{1}{C} \|f\|_{L^2(\Omega)}.$$

Indeed, we have

$$C \|u_{\beta}\|_{W^{1,2}(\Omega)}^2 \leq \int_{\Omega} |\nabla u_{\beta}|^2 + \int_{\partial\Omega} \beta u_{\beta}^2 \leq \|f\|_{L^2(\Omega)} \|u_{\beta}\|_{W^{1,2}(\Omega)},$$

and it suffices to invoke the continuity of the trace application, and of the bound $|\beta| \leq 1$. Since Ω is \mathcal{C}^2 , it follows from the regularity for Neumann problems that

$$u_{\beta} \in W^{1,2^*_{\partial\Omega}}(\Omega).$$

We can then bootstrap this argument and obtain successively that

$$\forall k \in \mathbb{N}, u_{\beta} \in W^{1,q_k}(\Omega) \quad (\text{B.5})$$

where the sequence $\{q_k\}_{k \in \mathbb{N}}$ is defined, by recurrence, as

$$q_{k+1} := \frac{(n-1)q_k}{n-1-\frac{q_k}{2}} \text{ if } \frac{q_k}{2} < n-1, q_k+1 \text{ else.}$$

The conclusion follows: for any $\beta \in \mathcal{B}(\partial\Omega)$, for any $p \in [1; \infty)$,

$$u_{\beta} \in W^{1,p}(\Omega)$$

and furthermore

$$\sup_{\beta \in \mathcal{B}(\partial\Omega)} \|u_{\beta}\|_{W^{1,p}(\Omega)} < \infty.$$

To obtain the uniform estimate

$$\sup_{\beta \in \mathcal{B}(\partial\Omega)} \sup_{\Omega} u_{\beta} < \infty$$

(the symmetric estimate $\inf_{\beta \in \mathcal{B}(\partial\Omega)} \inf_{\Omega} u_{\beta^*} > 0$ is obtained in the same way) it suffices to take a maximising sequence $\{\beta_k\}_{k \in \mathbb{N}}$ for $\|u_{\beta}\|_{L^{\infty}(\Omega)}$. Up to a subsequence, $\{\beta_k\}_{k \in \mathbb{N}}$ weakly converges to $\beta \in \mathcal{B}(\partial\Omega)$. By Sobolev embeddings, $\{u_{\beta_k}\}_{k \in \mathbb{N}}$ is uniformly bounded in $\mathcal{C}^{0,\alpha}(\Omega)$ for some $\alpha > 0$. Hence, it converges, strongly in $\mathcal{C}^0(\Omega)$, to u_{β} , which concludes the proof. \square

Proof of Lemma 12. Let us first underline that the set $\mathcal{B}(\partial\Omega)$ defined in (1.5) endowed with the weak-star topology of $L^{\infty}(\Omega)$ is compact. To apply the direct method in the calculus of variations, it suffices to show that all the functionals that define problems $(\mathbf{P}_{\max,\partial\Omega,\mathcal{B}})$ - $(\mathbf{P}_{\min,\partial\Omega,\mathcal{B}})$ - $(\mathbf{P}_{\max,\Omega,\mathcal{B}})$ - $(\mathbf{P}_{\min,\Omega,\mathcal{B}})$ are continuous under this weak $L^{\infty} - *$ topology. All these maps write as $\int_{\Omega} j(u_{\beta})$ or

$\int_{\partial\Omega} j(u_\beta)$. Since the functions u_β are uniformly bounded from above by Lemma 11, the dominated convergence theorem implies that these functionals are continuous if the map

$$\mathcal{B}(\partial\Omega) \ni \beta \mapsto u_\beta \in W^{1,2}(\Omega)$$

is continuous for the weak $L^\infty - *$ topology on $\mathcal{B}(\partial\Omega)$ and the weak $W^{1,2}$ topology on $W^{1,2}(\Omega)$. Indeed, it then suffices to invoke the compactness of the embeddings $W^{1,2}(\Omega) \hookrightarrow L^2(\Omega)$ and $W^{1,2}(\Omega) \hookrightarrow L^2(\partial\Omega)$.

Let us then prove the continuity of $\beta \mapsto u_\beta$ for these weak topologies. Let $\{\beta_n\}_{n \in \mathbb{N}} \in \mathcal{B}(\partial\Omega)^\mathbb{N}$ be a weakly converging sequence in $\mathcal{B}(\partial\Omega)$. Let $\beta \in \mathcal{B}(\partial\Omega)$ be such that

$$\beta_n \xrightarrow[n \rightarrow \infty]{} \beta.$$

In order to alleviate notations, we define, for any $n \in \mathbb{N}$, u_n as the solution of (\mathbf{E}_β) associated with β_n . Our goal is to show

$$u_n \xrightarrow[n \rightarrow \infty]{} u_\beta \text{ in } W^{1,2}(\Omega). \quad (\text{B.6})$$

First of all, multiplying the main equation of (\mathbf{E}_β) by u_n and integrating by parts yields

$$\int_{\Omega} |\nabla u_n|^2 + \int_{\partial\Omega} \beta_n u_n^2 = \int_{\Omega} f u_n \leq \|f\|_{L^2(\partial\Omega)} \|u_n\|_{L^2(\Omega)}.$$

From Lemma 10, it follows that $\{u_n\}_{n \in \mathbb{N}}$ is uniformly bounded in $L^2(\Omega)$ and, in turn, in $W^{1,2}(\Omega)$. From the Rellich-Kondrachov theorem there exists $\bar{u} \in W^{1,2}(\Omega)$ such that $\{u_n\}_{n \in \mathbb{N}}$ converges, up to a subsequence, to \bar{u} weakly in $W^{1,2}(\Omega)$ and strongly in $L^2(\Omega)$. By the compactness of the trace operator, $\{u_n\}_{n \in \mathbb{N}}$ converges to \bar{u} strongly in $L^2(\partial\Omega)$. Passing to the limit in the weak formulation of (\mathbf{E}_β) , we obtain that \bar{u} is the solution of (\mathbf{E}_β) associated with β . This concludes the proof. \square

C Definition of Steklov eigenvalues and eigenfunctions

Namely, we consider the resolvent operator $T : L^2(\partial\Omega) \rightarrow L^2(\partial\Omega)$ defined for all $f \in L^2(\partial\Omega)$ by

$$T(f) = z_f|_{\partial\Omega} \text{ where } z_f \text{ is the unique solution of } \begin{cases} -\Delta z_f = 0 & \text{in } \Omega, \\ \frac{\partial z_f}{\partial \nu} + \beta^* z_f = f & \text{on } \partial\Omega. \end{cases}$$

By compactness of the trace operator and standard regularity estimates, T is a compact operator. It is furthermore self-adjoint since, for any $f, g \in L^2(\partial\Omega)$ there holds

$$\begin{aligned} \int_{\partial\Omega} T(f)g &= \int_{\partial\Omega} z_f \left(\frac{\partial z_g}{\partial \nu} + \beta^* z_g \right) = \int_{\partial\Omega} \beta^* z_f z_g + \int_{\Omega} z_f \Delta z_g - \int_{\Omega} z_g \Delta z_f + \int_{\partial\Omega} z_g \frac{\partial z_f}{\partial \nu} \\ &= \int_{\partial\Omega} z_g \left(\frac{\partial z_f}{\partial \nu} + \beta^* z_f \right) = \int_{\partial\Omega} T(g)f. \end{aligned}$$

Finally, T is a positive operator: for any $f \in L^2(\partial\Omega)$ we have

$$\int_{\partial\Omega} T(f)f = \int_{\partial\Omega} z_f \frac{\partial z_f}{\partial \nu} + \int_{\partial\Omega} \beta^* z_f^2 = \int_{\Omega} |\nabla z_f|^2 + \int_{\partial\Omega} \beta^* z_f^2.$$

According to the spectral decomposition Theorem, there exists a non-increasing sequence of positive eigenvalues $\{r_k\}_{k \in \mathbb{N}}$ converging to zero and an associated family $\{\phi_k\}_{k \in \mathbb{N}}$ of eigenfunctions

satisfying for all $k \in \mathbb{N}$, $T(\phi_k) = r_k \phi_k$. Furthermore, the family $\{\phi_k\}_{k \in \mathbb{N}}$ is a Hilbert basis of $L^2(\partial\Omega)$ and we have

$$\begin{cases} -\Delta \phi_k = 0 & \text{in } \Omega, \\ \frac{\partial \phi_k}{\partial \nu} + \beta^* \phi_k = \frac{1}{r_k} \phi_k & \text{on } \partial\Omega, \end{cases} \quad \text{and} \quad \int_{\partial\Omega} \phi_k \phi_{k'} = \delta_{k,k'}. \quad (\text{C.1})$$

for all $(k, k') \in \mathbb{N}^2$. Let us set $\sigma_k := \frac{1}{r_k}$, we have

$$\sigma_p \xrightarrow{p \rightarrow +\infty} +\infty \quad \text{and} \quad \begin{cases} -\Delta \phi_k = 0 & \text{in } \Omega, \\ \frac{\partial \phi_k}{\partial \nu} + \beta^* \phi_k = \sigma_k \phi_k & \text{on } \partial\Omega, \\ \int_{\partial\Omega} \phi_k^2 = 1, \end{cases} \quad (\text{C.2})$$

for all $k \in \mathbb{N}$. Moreover,

$$\forall k, k' \in \mathbb{N}, \quad \int_{\partial\Omega} \phi_k \phi_{k'} = \delta_{k,k'}. \quad (\text{C.3})$$

Alternatively, we can define, for any $k \in \mathbb{N}$, σ_k via the min-max formula

$$\sigma_k := \min_{\substack{S \text{ subspace of dim} \\ k+1 \text{ of } W^{1,2}(\Omega)}} \max_{v \in S \setminus \{0\}} \frac{\int_{\Omega} |\nabla v|^2 + \int_{\partial\Omega} \beta^* v^2}{\int_{\partial\Omega} v^2}. \quad (\text{C.4})$$

D Proof of Theorem 4

We omit the subscript in $\mathcal{J}_{\partial\Omega}$ and simply write

$$\mathcal{J}(\beta) := \int_{\Omega} j(u_{\beta}).$$

To prove Theorem 4, we simply need to obtain a lower estimate of the second order derivative of the type given in Proposition 14. Indeed, the rest of the proof can be adapted verbatim. We thus need the derivatives of the functional under consideration. The first and second order derivatives of $\beta \mapsto u_{\beta}$ are still denoted by \dot{u}_{β} and \ddot{u}_{β} . The equations on \dot{u}_{β} and \ddot{u}_{β} remain the same as in the proof of Theorem 3: \dot{u}_{β} solves (3.1), while \ddot{u}_{β} satisfies (3.2). We now compute the derivatives of \mathcal{J} ; they write:

$$\dot{\mathcal{J}}(\beta)[h] = \int_{\Omega} \dot{u}_{\beta} j'(u_{\beta}) \quad \text{and} \quad \ddot{\mathcal{J}}(\beta)[h, h] = \int_{\Omega} \ddot{u}_{\beta} j'(u_{\beta}) + \int_{\Omega} (\dot{u}_{\beta})^2 j''(u_{\beta}). \quad (\text{D.1})$$

We define the adjoint state p_{β} as the unique solution in $W^{1,2}(\Omega)$ of

$$\begin{cases} -\Delta p_{\beta} = j'(u_{\beta}) & \text{in } \Omega, \\ \frac{\partial p_{\beta}}{\partial \nu} + \beta p_{\beta} = 0 & \text{on } \partial\Omega. \end{cases} \quad (\text{D.2})$$

Since $\inf_{\Omega} u_{\beta} > 0$ and since $j' > 0$ on \mathbb{R}_+^* , the maximum principle entails

$$\inf_{\Omega} p_{\beta} > 0. \quad (\text{D.3})$$

If we multiply (D.2) by \dot{u}_{β} and (3.1) by p_{β} , integrating by parts leads to

$$0 = \int_{\Omega} \nabla \dot{u}_{\beta} \cdot \nabla p_{\beta} + \int_{\partial\Omega} \beta p_{\beta} \dot{u}_{\beta} - \int_{\Omega} j'(u_{\beta}) \dot{u}_{\beta}$$

$$0 = \int_{\Omega} \nabla \dot{u}_{\beta} \cdot \nabla p_{\beta} + \int_{\partial\Omega} \beta p_{\beta} \dot{u}_{\beta} + \int_{\partial\Omega} h u_{\beta} p_{\beta}$$

so that

$$\int_{\Omega} j'(u_{\beta}) \dot{u}_{\beta} = - \int_{\partial\Omega} h u_{\beta} p_{\beta}.$$

This leads to

$$\dot{\mathcal{J}}(\beta)[h] = - \int_{\partial\Omega} h u_{\beta} p_{\beta}. \quad (\text{D.4})$$

Similarly, we get

$$\ddot{\mathcal{J}}(\beta)[h, h] = -2 \int_{\partial\Omega} h \dot{u}_{\beta} p_{\beta} + \int_{\Omega} j''(u_{\beta}) (\dot{u}_{\beta})^2. \quad (\text{D.5})$$

We make a proof by contradiction: let $\beta^* \in \mathcal{B}(\partial\Omega)$ be a maximiser such that the set

$$\omega^* := \{0 < \beta^* < 1\} \quad (\text{D.6})$$

has positive measure:

$$\mathcal{H}^{d-1}(\omega^*) > 0.$$

Thus, for any admissible perturbation h at β^* supported in ω^* we must have

$$\dot{\mathcal{J}}(\beta^*)[h] = 0. \quad (\text{D.7})$$

We now prove that there exists an admissible perturbation h supported in ω^* such that

$$\ddot{\mathcal{J}}(\beta^*)[h, h] > 0. \quad (\text{D.8})$$

We aim at obtaining an expression of $\ddot{\mathcal{J}}$ that is similar to a Rayleigh quotient, since this is the main point of the proof. We can show the following adaptation of Proposition 14 there exist three constants A, B, C with $A > 0$ such that

$$\ddot{\mathcal{J}}(\beta^*)[h, h] \geq A \int_{\Omega} |\nabla \dot{u}_{\beta^*}|^2 - B \|\dot{u}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{u}_{\beta^*}\|_{L^2(\Omega)} - C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (\text{D.9})$$

To that effect, we set

$$\mathcal{W}(\beta^*)[h, h] := -2 \int_{\partial\Omega} h \dot{u}_{\beta^*} p_{\beta^*}. \quad (\text{D.10})$$

Using the same computations as in the proof of Proposition 14, if we set

$$\Psi_{\beta^*} := \frac{p_{\beta^*}}{u_{\beta^*}},$$

we get

$$\mathcal{W}(\beta^*)[h, h] = \int_{\partial\Omega} \Psi_{\beta^*} (\partial_{\nu} (\dot{u}_{\beta^*}^2) + 2\beta^* \dot{u}_{\beta^*}^2). \quad (\text{D.11})$$

We note that

$$\inf_{\Omega} \Psi_{\beta^*} > 0$$

and that

$$\frac{\partial \Psi_{\beta^*}}{\partial \nu} = \frac{\partial_{\nu} p_{\beta^*}}{u_{\beta^*}} - \Psi_{\beta^*} \frac{\partial_{\nu} u_{\beta^*}}{u_{\beta^*}} = -\beta^* \Psi_{\beta^*} + \beta^* \Psi_{\beta^*} = 0. \quad (\text{D.12})$$

We use identity (3.15) and get

$$\ddot{\mathcal{J}}(\beta^*)[h, h] = 2 \int_{\Omega} \Psi_{\beta^*} |\nabla \dot{u}_{\beta^*}|^2 - \int_{\Omega} \dot{u}_{\beta^*}^2 \Delta \Psi_{\beta^*} + \int_{\partial\Omega} (2\beta^* \Psi_{\beta^*} + j''(u)) \dot{u}_{\beta^*}^2. \quad (\text{D.13})$$

Since β^* and u_{β^*} belong to $L^\infty(\partial\Omega)$ and since $j \in \mathcal{C}^2$, there exists a constant C independent of h such that, for any admissible perturbation h

$$\int_{\partial\Omega} (2\beta^* \Psi_{\beta^*} + j''(u)) (\dot{u}_{\beta^*})^2 \geq -C \int_{\partial\Omega} \dot{u}_{\beta^*}^2. \quad (\text{D.14})$$

The rest of the proof follows exactly the same lines: indeed, the rest of the proof of Theorem 3 hinges upon the analysis of \dot{u}_{β^*} , not on the fact that the criterion to optimise is distributed. The equation on \dot{u}_{β^*} remains unchanged, and so does the rest of the analysis.

E Proof of Theorem 7

We define

$$\mathcal{J}(\beta) = \int_{\Omega} j(u_{\beta}).$$

We shall make use of the computations of Appendix D.

We recall that the first order derivative of the criterion is given in (D.1). We set, for any $\beta \in \mathcal{B}(\partial\Omega)$,

$$\Phi_{\beta} := u_{\beta} p_{\beta}$$

where p_{β} is the solution of (D.2). Therefore, for any admissible perturbation h at a given β , there holds

$$\dot{\mathcal{J}}(\beta)[h] = - \int_{\partial\Omega} h \Phi_{\beta}. \quad (\text{E.1})$$

Let β^* be a solution of $(\mathbf{P}_{\min, \Omega, \mathcal{B}})$. Since we must have

$$-\dot{\mathcal{J}}(\beta^*)[h] = \int_{\partial\Omega} \Phi_{\beta^*} h \leq 0$$

for any admissible perturbation h at β^* , there exists a real number λ (necessarily positive as $\Phi_{\beta} > 0$) such that

1. $\{0 < \beta^* < 1\} \subset \{\Phi_{\beta^*} = \lambda\}$,
2. $\{\beta^* = 1\} \subset \{\Phi_{\beta^*} \geq \lambda\}$,
3. $\{\beta^* = 0\} \subset \{\Phi_{\beta^*} \leq \lambda\}$.

As in the proof of Theorem 5, we show that these conditions imply

$$\mathcal{H}^{d-1}(\{\beta^* = 0\}) = 0. \quad (\text{E.2})$$

The required conclusion then follows.

To prove (E.2), let us first observe that

$$\nabla \Phi_{\beta^*} = u_{\beta^*} \nabla p_{\beta^*} + p_{\beta^*} \nabla u_{\beta^*}, -\Delta \Phi_{\beta^*} = -p_{\beta^*} \Delta u_{\beta^*} - u_{\beta^*} \Delta p_{\beta^*} - 2 \nabla u_{\beta^*} \cdot \nabla p_{\beta^*}. \quad (\text{E.3})$$

First, we set

$$B := -2 \frac{\nabla p_{\beta^*}}{p_{\beta^*}}$$

and

$$V := \frac{f}{u_{\beta^*}} + \frac{j'(u_{\beta^*})}{p_{\beta^*}} + 2 \frac{|\nabla p_{\beta^*}|^2}{p_{\beta^*}^2}.$$

Second, observe that

$$\nabla u_{\beta^*} = \frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{u_{\beta^*}}{p_{\beta^*}} \nabla p_{\beta^*} = \frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{\Phi_{\beta^*}}{p_{\beta^*}^2} \nabla p_{\beta^*}.$$

Plugging this expression into (E.3) yields

$$\begin{aligned} -\Delta \Phi_{\beta^*} &= f p_{\beta^*} + j'(u_{\beta^*}) u_{\beta^*} - 2 \left(\frac{\nabla \Phi_{\beta^*}}{p_{\beta^*}} - \frac{\Phi_{\beta^*}}{p_{\beta^*}^2} \nabla p_{\beta^*} \right) \cdot \nabla p_{\beta^*} \\ &= \Phi_{\beta^*} \left(\frac{f}{u_{\beta^*}} + \frac{j'(u_{\beta^*})}{p_{\beta^*}} \right) + 2 \Phi_{\beta^*} \frac{|\nabla p_{\beta^*}|^2}{p_{\beta^*}^2} + \langle \nabla \Phi_{\beta^*}, B \rangle \\ &= V \Phi_{\beta^*} + \langle \nabla \Phi_{\beta^*}, B \rangle. \end{aligned}$$

Since $\inf_{\Omega} \min\{u_{\beta^*}, p_{\beta^*}\} > 0$ it follows that $\inf_{\Omega} \Phi_{\beta^*} \geq 0$. We also have $V \geq 0$. We can then follow the proof of Theorem 5 verbatim.

F Proof of Theorem 9

We recall that

$$\mathcal{R}(\beta) := \int_{\Omega} j(y_{\beta}).$$

To prove Theorem 9, we need a lower estimate of the second order derivative of the type given in Proposition 14; in a second step, we will need a set of eigenfunctions different from the ones used in the proofs of Theorem 3. Although not straightforward, we show how this can be done.

We start by computing the derivatives of the criterion. The first and second order derivatives of $\beta \mapsto y_{\beta}$ are denoted by \dot{y}_{β} and \ddot{y}_{β} . By Assumptions (\mathbf{H}_{NL}) -(\mathbf{H}_{stab}), these derivatives exist. Furthermore, \dot{y}_{β} satisfies

$$\begin{cases} -\Delta \dot{y}_{\beta} = \frac{\partial g}{\partial u}(x, y_{\beta}) \dot{y}_{\beta} & \text{in } \Omega, \\ \frac{\partial \dot{y}_{\beta}}{\partial \nu} + \beta \dot{y}_{\beta} = -h y_{\beta} & \text{on } \partial\Omega, \end{cases} \quad (\text{F.1})$$

while \ddot{y}_{β} satisfies

$$\begin{cases} -\Delta \ddot{y}_{\beta} = \frac{\partial^2 g}{\partial u^2}(x, y_{\beta}) (\dot{y}_{\beta})^2 + \frac{\partial g}{\partial u}(x, y_{\beta}) \ddot{y}_{\beta} & \text{in } \Omega, \\ \frac{\partial \ddot{y}_{\beta}}{\partial \nu} + \beta \ddot{y}_{\beta} = -2h \dot{y}_{\beta} & \text{on } \partial\Omega. \end{cases} \quad (\text{F.2})$$

We now compute the derivatives of \mathcal{R} ; they write:

$$\dot{\mathcal{R}}(\beta)[h] = \int_{\Omega} \dot{y}_{\beta} j'(y_{\beta}) \text{ and } \ddot{\mathcal{R}}(\beta)[h, h] = \int_{\Omega} \ddot{y}_{\beta} j'(y_{\beta}) + \int_{\Omega} (\dot{y}_{\beta})^2 j''(y_{\beta}). \quad (\text{F.3})$$

We define the adjoint state p_{β} as the unique solution in $W^{1,2}(\Omega)$ of

$$\begin{cases} -\Delta p_{\beta} = \frac{\partial g}{\partial u}(x, y_{\beta}) p_{\beta} + j'(y_{\beta}) & \text{in } \Omega, \\ \frac{\partial p_{\beta}}{\partial \nu} + \beta p_{\beta} = 0 & \text{on } \partial\Omega. \end{cases} \quad (\text{F.4})$$

We need to check that a solution to this equation indeed exists.

Lemma 17. *There exists a unique solution $p_\beta \in W^{1,2}(\Omega)$ of (F.4). For any $p \in [1; +\infty)$, $p_\beta \in W^{1,p}(\Omega)$ and*

$$\inf_{\Omega} p_\beta > 0.$$

Proof of Lemma 17. By Assumption (**H_{stab}**), the energy functional

$$E : W^{1,2}(\Omega) \ni p \mapsto \frac{1}{2} \int_{\Omega} |\nabla p|^2 - \frac{1}{2} \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) p^2 - \int_{\Omega} j'(y_\beta) p + \int_{\partial\Omega} \beta p^2$$

can be bounded from below as

$$E(p) \geq \mu_\beta \int_{\Omega} p^2 - \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) p.$$

Hence, E is coercive, and so a solution p_{β^*} to (F.4) exists. To prove the uniqueness of this solution, we argue by contradiction: if two solutions p_β, q_β exist, then $z_\beta := p_\beta - q_\beta$ solves

$$\begin{cases} -\Delta z_\beta = \frac{\partial g}{\partial u}(x, y_\beta) z_\beta & \text{in } \Omega, \\ \frac{\partial z_\beta}{\partial \nu} + \beta z_\beta = 0 & \text{on } \partial\Omega. \end{cases} \quad (\text{F.5})$$

Hence, if $p_\beta \neq q_\beta$, $z_\beta \neq 0$ is an eigenfunction of the operator $-\Delta - \frac{\partial g}{\partial u}(x, y_\beta)$, associated with the eigenvalue 0. However, the lowest eigenvalue of this operator is $\mu_\beta > 0$, a contradiction. Uniqueness follows.

The $W^{1,p}$ -regularity of p_β is a consequence of the same arguments as in Lemma 11.

Finally, the positivity of p_β is a consequence of the following version of the maximum principle: as $\mu_\beta > 0$, should p_β not be positive, the negative part $p_\beta^- := -p_\beta \mathbb{1}_{p_\beta < 0}$ satisfies

$$\int_{\Omega} |\nabla p_\beta^-|^2 - \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) (p_\beta^-)^2 \leq - \int_{\Omega} j'(y_\beta) p_\beta^- \leq 0.$$

By the variational formulation (1.15) of $\mu_\beta > 0$, we necessarily have $p_\beta^- = 0$. So we first have $p_\beta^- \geq 0$. It then suffices to apply the classical maximum principle to conclude. \square

If we multiply (F.4) by \dot{y}_β and (F.1) by p_β , integrating by parts leads to

$$\begin{aligned} 0 &= \int_{\Omega} \nabla \dot{y}_\beta \cdot \nabla p_\beta - \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) p_\beta \dot{y}_\beta + \int_{\partial\Omega} \beta p_\beta \dot{y}_\beta - \int_{\Omega} j'(y_\beta) \dot{y}_\beta \\ 0 &= \int_{\Omega} \nabla \dot{y}_\beta \cdot \nabla p_\beta - \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) p_\beta \dot{y}_\beta + \int_{\partial\Omega} \beta p_\beta \dot{y}_\beta + \int_{\partial\Omega} h y_\beta p_\beta \end{aligned}$$

so that

$$\int_{\Omega} j'(y_\beta) \dot{y}_\beta = - \int_{\partial\Omega} h y_\beta p_\beta.$$

This leads to

$$\dot{\mathcal{R}}(\beta)[h] = - \int_{\partial\Omega} h y_\beta p_\beta. \quad (\text{F.6})$$

Similarly, we get

$$\ddot{\mathcal{R}}(\beta)[h, h] = -2 \int_{\partial\Omega} h \dot{y}_\beta p_\beta + \int_{\Omega} \frac{\partial^2 g}{\partial u^2}(x, y_\beta) p_\beta (\dot{y}_\beta)^2 + \int_{\Omega} j''(y_\beta) (\dot{y}_\beta)^2. \quad (\text{F.7})$$

We argue by contradiction: let $\beta^* \in \mathcal{B}(\partial\Omega)$ be a maximiser such that the set

$$\omega^* := \{0 < \beta^* < 1\} \quad (\text{F.8})$$

has positive measure:

$$\mathcal{H}^{d-1}(\omega^*) > 0.$$

Thus, for any admissible perturbation h at β^* supported in ω^* we must have

$$\dot{\mathcal{R}}(\beta^*)[h] = 0. \quad (\text{F.9})$$

We now prove that there exists an admissible perturbation h supported in ω^* such that

$$\ddot{\mathcal{R}}(\beta^*)[h, h] > 0. \quad (\text{F.10})$$

We aim at obtaining an expression of $\ddot{\mathcal{R}}$ that is similar to a Rayleigh quotient, since this is the main point of the proof. We show the following adaptation of Proposition 14: there exist $A > 0$ and two constants $B, C > 0$ such that

$$\ddot{\mathcal{R}}(\beta^*)[h, h] \geq A \int_{\Omega} |\nabla \dot{y}_{\beta^*}|^2 - B \|\dot{y}_{\beta^*}\|_{W^{1,2}(\Omega)} \|\dot{y}_{\beta^*}\|_{L^2(\Omega)} - C \int_{\partial\Omega} \dot{y}_{\beta^*}^2. \quad (\text{F.11})$$

To that effect, we set

$$\mathcal{W}(\beta^*)[h, h] := -2 \int_{\partial\Omega} h \dot{y}_{\beta^*} p_{\beta^*}. \quad (\text{F.12})$$

As before, if we set

$$\Psi_{\beta^*} := \frac{p_{\beta^*}}{u_{\beta^*}},$$

we obtain

$$\mathcal{W}(\beta^*)[h, h] = \int_{\partial\Omega} \Psi_{\beta^*} (\partial_{\nu} (\dot{y}_{\beta^*}^2) + 2\beta^* \dot{y}_{\beta^*}^2). \quad (\text{F.13})$$

We note that

$$\inf_{\Omega} \Psi_{\beta^*} > 0$$

and that

$$\frac{\partial \Psi_{\beta^*}}{\partial \nu} = \frac{\partial_{\nu} p_{\beta^*}}{u_{\beta^*}} - \Psi_{\beta^*} \frac{\partial_{\nu} u_{\beta^*}}{u_{\beta^*}} = -\beta^* \Psi_{\beta^*} + \beta^* \Psi_{\beta^*} = 0. \quad (\text{F.14})$$

From (3.15) we are led to

$$\ddot{\mathcal{R}}(\beta^*)[h, h] = 2 \int_{\Omega} \Psi_{\beta^*} |\nabla \dot{y}_{\beta^*}|^2 - \int_{\Omega} \dot{y}_{\beta^*}^2 \Delta \Psi_{\beta^*} + \int_{\partial\Omega} (2\beta^* \Psi_{\beta^*} + j''(u)) \dot{y}_{\beta^*}^2 + \int_{\Omega} \frac{\partial^2 g}{\partial u^2}(x, y_{\beta}) p_{\beta} (\dot{y}_{\beta})^2. \quad (\text{F.15})$$

Since β^* and u_{β^*} belong to $L^{\infty}(\partial\Omega)$, since $g \in \mathcal{C}^2$ in its second variable and since $j \in \mathcal{C}^2$, there exists a constant C independent of h such that, for any admissible perturbation h

$$\begin{aligned} \int_{\partial\Omega} (2\beta^* \Psi_{\beta^*} + j''(u)) (\dot{y}_{\beta^*})^2 + \int_{\Omega} \frac{\partial^2 g}{\partial u^2}(x, y_{\beta}) p_{\beta} (\dot{y}_{\beta})^2 &\geq -C \left(\int_{\partial\Omega} \dot{y}_{\beta^*}^2 + \int_{\Omega} \dot{y}_{\beta}^2 \right) \\ &\geq -C \int_{\partial\Omega} \dot{y}_{\beta}^2 - C \|\dot{y}_{\beta}\|_{L^2(\Omega)} \|\dot{y}_{\beta}\|_{W^{1,2}(\Omega)}. \end{aligned}$$

The rest of the proof follows exactly the same lines, except that, instead of considering the sequence of eigenpairs defined in (3.27), we rather need the following: we consider the eigenvalues

$$0 \leq \sigma_0 \leq \sigma_1 \leq \dots \leq \sigma_k \xrightarrow{k \rightarrow \infty} +\infty$$

where, for each k , the eigenvalue σ_k is associated with the eigenfunction ϕ_k solution of

$$\forall k \in \mathbb{N}, \begin{cases} -\Delta \phi_k = \frac{\partial g}{\partial u}(x, y_\beta) p_\beta & \text{in } \Omega, \\ \frac{\partial \phi_k}{\partial \nu} + \beta \phi_k = \sigma_k \phi_k & \text{on } \partial\Omega \end{cases} \quad \text{and, for any } k, k' \in \mathbb{N}, \int_{\partial\Omega} \phi_k \phi_{k'} = \delta_{k,k'}. \quad (\text{F.16})$$

The one thing that needs to be checked is that these eigenpairs are well defined. This is once again a consequence of the stability Assumption (**H_{stab}**): proceeding as in Appendix C, it suffices to show that the operator $T : L^2(\partial\Omega) \ni f \mapsto T(f)$ defined, for any $f \in L^2(\partial\Omega)$, as

$$z_f|_{(\partial\Omega)} \text{ where } z_f \text{ is the unique solution of } \begin{cases} -\Delta z_f - \frac{\partial g}{\partial u}(x, y_\beta) z_f = 0 & \text{in } \Omega, \\ \frac{\partial z_f}{\partial \nu} + \beta z_f = f & \text{on } \partial\Omega \end{cases}$$

is compact. First, we need to check that T is well-defined. However, this follows from the same arguments as in Lemma 17, considering this time the energy functional

$$E : W^{1,2}(\Omega) \ni z \mapsto \frac{1}{2} \int_{\Omega} |\nabla z|^2 - \frac{1}{2} \int_{\Omega} \frac{\partial g}{\partial u}(x, y_\beta) z^2 + \int_{\partial\Omega} \beta z^2 - \int_{\partial\Omega} z f.$$

Second, the compactness is a consequence of standard $W^{1,2}$ -estimates.

We can hence define the eigenpairs described in (F.16). The rest of the proof is adapted verbatim.

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References

- [1] F. Alabau-Boussouira, P. Cannarsa, and C. Urbani. Exact controllability to eigensolutions for evolution equations of parabolic type via bilinear control. *Nonlinear Differential Equations and Applications NoDEA*, 29(4):1–32, 2022.
- [2] A. Alvino, C. Nitsch, and C. Trombetti. A talenti comparison result for solutions to elliptic problems with robin boundary conditions. *arXiv: Analysis of PDEs*, 2019.
- [3] G. Auchmuty. Steklov eigenproblems and the representation of solutions of elliptic boundary value problems. *Numer. Funct. Anal. Optim.*, 25(3-4):321–348, 2004.
- [4] X. Bai, X. He, and F. Li. An optimization problem and its application in population dynamics. *Proc. Amer. Math. Soc.*, 144(5):2161–2170, 2016.
- [5] K. Beauchard and C. Laurent. Local controllability of 1D linear and nonlinear Schrödinger equations with bilinear control. *Journal de Mathématiques Pures et Appliquées*, 94(5):520–554, Nov. 2010.

- [6] H. Berestycki, F. Hamel, and L. Roques. Analysis of the periodically fragmented environment model : I – species persistence. *Journal of Mathematical Biology*, 51(1):75–113, 2005.
- [7] J. Bintz and S. Lenhart. Optimal resources allocation for a diffusive population model. *Journal of Biological Systems*, 28(04):945–976, Dec. 2020.
- [8] A. Borzi, E.-J. Park, and M. V. Lass. Multigrid optimization methods for the optimal control of convection–diffusion problems with bilinear control. *Journal of Optimization Theory and Applications*, 168(2):510–533, Aug. 2015.
- [9] D. Bucur, G. Buttazzo, and C. Nitsch. Symmetry breaking for a problem in optimal insulation. *Journal de Mathématiques Pures et Appliquées*, 107(4):451–463, 2017.
- [10] D. Bucur, G. Buttazzo, and C. Nitsch. Two optimization problems in thermal insulation. *Notices of the American Mathematical Society*, 64(08):830–835, Sept. 2017.
- [11] G. Buttazzo and G. Dal Maso. An existence result for a class of shape optimization problems. *Archive for Rational Mechanics and Analysis*, 122(2):183–195, Jun 1993.
- [12] P. Cannarsa, G. Floridia, and A. Y. Khapalov. Multiplicative controllability for semilinear reaction–diffusion equations with finitely many changes of sign. *Journal de Mathématiques Pures et Appliquées*, 108(4):425–458, Oct. 2017.
- [13] F. Della Pietra, N. Gavitone, and H. Kovařík. Optimizing the first eigenvalue of some quasi-linear operators with respect to boundary conditions. *ESAIM: Control, Optimisation and Calculus of Variations*, 23(4):1381–1395, sep 2017.
- [14] F. Della Pietra, C. Nitsch, R. Scala, and C. Trombetti. An optimization problem in thermal insulation with robin boundary conditions. *Communications in Partial Differential Equations*, pages 1–17, June 2021.
- [15] F. Della Pietra, C. Nitsch, and C. Trombetti. An optimal insulation problem. *Mathematische Annalen*, Aug. 2020.
- [16] K. Fister and C. McCarthy. Optimal control of a chemotaxis system. *Quarterly of Applied Mathematics*, 61(2):193–211, 2003.
- [17] F. Guillén-González, E. Mallea-Zepeda, and M. Á. Rodríguez-Bellido. Optimal bilinear control problem related to a chemo-repulsion system in 2d domains. *ESAIM: Control, Optimisation and Calculus of Variations*, 26:29, 2020.
- [18] A. Henrot and M. Pierre. *Shape variation and optimization. A geometrical analysis*, volume 28. EMS Tracts in Mathematics, European Mathematical Society (EMS), Zürich., 2018.
- [19] D. Hömberg, K. Krumbiegel, and J. Rehberg. Optimal control of a parabolic equation with dynamic boundary condition. *Appl. Math. Optim.*, 67(1):3–31, 2013.
- [20] S. Lenhart, V. Protopopescu, and J. Yong. Optimal control of a reflection boundary coefficient in an acoustic wave equation. *Appl. Anal.*, 68(1-2):179–194, 1998.
- [21] Y. Lou. On the effects of migration and spatial heterogeneity on single and multiple species. *Journal of Differential Equations*, 223(2):400–426, Apr. 2006.
- [22] I. Mazari. *Shape optimization and spatial heterogeneity in reaction-diffusion equations*. Theses, Sorbonne Université, July 2020.

- [23] I. Mazari, G. Nadin, and Y. Privat. Optimal location of resources maximizing the total population size in logistic models. *J. Math. Pures Appl. (9)*, 134:1–35, 2020.
- [24] I. Mazari, G. Nadin, and Y. Privat. Chapter 12 - some challenging optimization problems for logistic diffusive equations and their numerical modeling. In E. Trélat and E. Zuazua, editors, *Numerical Control: Part A*, volume 23 of *Handbook of Numerical Analysis*, pages 401–426. Elsevier, 2022.
- [25] I. Mazari, G. Nadin, and Y. Privat. Optimisation of the total population size for logistic diffusive equations: bang-bang property and fragmentation rate. *Communications in Partial Differential Equations*, 47(4):797–828, 2022.
- [26] I. Mazari and D. Ruiz-Balet. A fragmentation phenomenon for a nonenergetic optimal control problem: Optimization of the total population size in logistic diffusive models. *SIAM Journal on Applied Mathematics*, 81(1):153–172, Jan. 2021.
- [27] K. Nagahara, Y. Lou, and E. Yanagida. Maximizing the total population with logistic growth in a patchy environment. *Journal of Mathematical Biology*, 82(1-2), Jan. 2021.
- [28] K. Nagahara and E. Yanagida. Maximization of the total population in a reaction–diffusion model with logistic growth. *Calculus of Variations and Partial Differential Equations*, 57(3):80, Apr 2018.
- [29] C. Nitsch and C. Trombetti. The classical overdetermined serrin problem. *Complex Variables and Elliptic Equations*, 63(7-8):1107–1122, Dec. 2017.
- [30] J. Serrin. A symmetry problem in potential theory. *Arch. Rational Mech. Anal.*, 43:304–318, 1971.
- [31] C. G. Simader and H. Sohr. A new approach to the Helmholtz decomposition and the Neumann problem in l^q -spaces for bounded and exterior domains. In *Series on Advances in Mathematics for Applied Sciences*, pages 1–35. World Scientific, Aug. 1992.
- [32] H. F. Weinberger. Remark on the preceding paper of Serrin. *Archive for Rational Mechanics and Analysis*, 43(4):319–320, Jan. 1971.

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