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1 Standard units for ElectroChromic Shift (ECS) measurements in plant biology 2 Charlie Mathiot^{1, 2}, Jean Alric¹ 3 4 ¹ Aix Marseille Univ, CEA, CNRS, BIAM, UMR 7265, Photosynthesis and Environment, 5 6 Saint Paul-Lez-Durance, France F-13108 7 ² Groupe Biomasse 3G, CEA Tech, CEA Cadarache, 13108 Saint-Paul-Lez-Durance, France. 8 9 Highlight ElectroChromic Shift can be used to measure the light-response curves of electron transport 10 11 rate, the fraction of open photosystems and a transmission coefficient for the photosynthetic 12 chain. 13 14 **Abstract** 15 The absorbance shift of pigments is proportional to the membrane potential ($\Delta \psi$) in plants, 16 green algae, and many photosynthetic bacteria. It is currently denoted as ElectroChromic Shift 17 (ECS) at 515-520 nm for plant carotenoids. It is increasingly being used for phenotyping plants for traits related to photosynthesis or chloroplast metabolism because it is a non-18 19 invasive technique and also because more instruments are now commercially available from 20 various manufacturers. The ECS technique is currently used to monitor the post-illumination 21 decay of the proton-motive force (pmf), but it has a yet more general use for quantitative 22 studies on photosynthetic energy transduction. Here we briefly summarize the basic 23 knowledge on ECS, emphasize the full potential of this technique, and propose a quantitative 24 analysis of the photosynthetic performance with the definition of a transmission coefficient 25 for electrons along the photosynthetic chain. 26 27

Keywords

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29 Photosynthesis, chloroplast, thylakoid, cyclic electron flow, ATP-synthase

31 Abbreviations

- 32 ECS, ElectroChromic Shift
- 33 *pmf*, proton-motive force
- 34 Δw , membrane potential

- 35 *I*, initial electron flow rate
- 36 *J*, steady-state electron flow rate
- 37 PS_o, open photosystems

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Introduction

ECS signal was first reported as P515 (Duysens, 1954), an absorbance increase (ΔA) peaking around 515 nm. This signal was attributed to a change in absorption of carotenoids (Witt, 1979), proportional to the membrane potential ($\Delta \psi$). This signal was then used for the study of proton flow through ATPase (Junge, 1970) and the "slow" electrogenic phase of cyt $b_6 f$ complex (Joliot and Delosme, 1974). More recently, the ECS decay rate after a continuous illumination was related to photosynthetic electron flow, allowing this technique to measure a flux as for any other reaction intermediate in the photosynthetic chain (Sacksteder and Kramer, 2000). This concept deserves explanation: if the ECS signal is constant over time, as observed at steady-state under continuous illumination, then the reactions building the protonmotive force (light-induced electron and proton transport) are exactly compensated by those consuming the proton-motive force (by ATPase in the dark), so that the post-illumination decay rate of ECS (membrane conductivity g_H^+ , in s⁻¹) is exactly equal to the electron transfer rate through the photosystems. This contribution had a very broad impact on the photosynthesis and chloroplast research community, it allowed for semi-quantitative studies of energy-dependent non-photochemical chlorophyll fluorescence quenching (qE) in relation to electron and proton flow (Kanazawa and Kramer, 2002), of cyclic and linear electron flows by comparing WT and pgr5 plants (Avenson et al., 2005), of the parsing of the pmf between $\Delta \psi$ and ΔpH (Cruz et al., 2001) and the recent identification of the thylakoid proton antiporter KEA3 (Kunz et al., 2014).

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The rate of a reaction (in s⁻¹) defines enzyme kinetics. Yet, for the characterization of an integrated pathway like photosynthesis, the flux (in number of particles, per unit time per unit area) is a much more relevant parameter. Kramer and co-workers showed that the amplitude of the ECS decay is proportional to the proton motive force (*pmf*). Therefore, the efflux of protons through ATPase (v_H^+) can be related to its conductivity g_H^+ as follows: $v_H^+ \sim pmf \times g_H^+$ (Avenson et al., 2004). The units for this proton efflux are not always specified (often in arbitrary units) but they should be $\mu mol_H^+ m^{-2} s^{-1}$, like the electron transport rate (ETR) is expressed in $\mu mol e^- m^{-2} s^{-1}$. An uncertainty lies with these units as they depend on the amount of photons absorbed (and electrons or protons transferred) per leaf

area. Therefore we would like to follow up on an idea that was proposed previously (Joliot and Joliot, 2002): ECS can be calibrated on an internal standard to determine the number of electrons transferred per photosystem per unit time (e⁻ PS⁻¹ s⁻¹). Here the unit area is the functional unit of photosynthesis, the photosystems themselves.

Materials and Methods

Chlorella sorokiniana WT, and chloroplast ATPase-deficient mutant SL8 (Rappaport et al., 1999) were heterotrophically grown at 20°C, in continuous light (20 μmol_{photons} m⁻² s⁻¹), in Tris-acetate-phosphate medium. Prior to each experiment, cells were concentrated and resuspended in 10% Ficoll to a concentration of ~ 15 μg_{Chl} mL⁻¹, then adapted for >15 min in Erlenmeyer flasks stirred at 150 rpm prior to measurements. 3-(3,4-dichlorophenyl)-1,1-dimethyl-urea (DCMU) was used at a final concentration of 10 μM, hydroxylamine at 10

mM.

Absorption spectroscopy - The ElectroChromic Shift (ECS) of carotenoids was measured as an absorbance changes (ΔA) with a Joliot-type spectrophotometer using a Xenon flashlamp fitted to a Jobin-Yvon HL300 f/2 monochromator (Joliot et al., 1980) or to a white LED fitted to a 520 nm band pass dielectric filter (JTS-10, Biologic). Similar measurements can be done with Walz Dual-Pam 100 or MultispeQ. The detecting photodiodes were protected by a combination of BG39 (Schott) and CVI low-pass dielectric filters (600 nm). Action spectra were obtained with continuous light provided with a SPEX illuminator as excitation light (a Xenon arc lamp fitted to a Jobin Yvon 1681 Minimate f/4 monochromator) and computer-controlled shutter (Uniblitz). The JTS-10 is supplied with a detection filter centered at 546 nm that can be used for deconvolution of cytochrome f (554 nm) or b (563 nm) signals. We find that, while it is necessary to correct for cytochrome kinetics, it was not for ECS in the present study where

illuminations were rather short (< 10 s). For longer illuminations, inducing de-epoxidation of violaxanthin to zeaxanthin (Ilioaia et al., 2011), thylakoid membrane swelling and scattering changes or chloroplast migration in plant leaves, it would become necessary to extract pure ECS traces by subtracting 546 nm from 520 nm absorbance changes. For such applications, the technique of "continuous ECS-indicated recording" *via* the dual-wavelength (550–520 nm) accessory of the Dual-PAM-100 may be more appropriate (Klughammer et al., 2013).

Saturating flash - If a "single turnover saturating flash" is applied to the photosynthetic material, the ECS signal ($\Delta A_{520~nm}$) is proportional to the amount of active (or open) photosystems, PS_o. The technical requirements are the following: the flash must be short (typically $\leq 1~\mu s$, otherwise PSI can undergo a second turnover during the flash), and intense (typically $\geq 1~m J$, saturation is reached when twice as many photons have no significant effect (< 10%) to increase the signal). Pulsed lasers (flash-lamp pumped and frequency-doubled Nd:YAG lasers with dyes, $\leq 10~n s$ and $\geq 30~m J$) typically meet these criteria. Here flashes at 700 nm were delivered by an OPO pumped at 532 nm by a frequency-doubled Nd:YAG laser (Surelite II, Continuum). A saturating single-turnover flash gives the amount of "open" or "potentially active" photosystems: *i.e.* centers in the PA state if we call P the reduced form of the Primary donor and A the oxidized form of the Acceptor. The flash will close the centers, forming the P+A- state. In the dark, all centers are in the PA state, and during a continuous illumination some centers can be under the P+A or PA- states, which are other forms of closed centers. PS_o being the relative amplitude of the signal created by the flash is therefore the fraction of PA.

Data treatment - ECS kinetics were fitted by R-scripts. ECS rise and decay were respectively
fitted to a linear model (initial flow rate *I*) and first order exponential decay (pseudo steady state flow rate *J*), using the *nls* (Non-Linear Squares) function for the latter.

Results

Figure 1A shows the spectrum of light-induced absorbance changes (ΔA) 20 μs after a saturating flash inducing an increase in trans-thylakoid $\Delta \psi$. Similarly as in (Witt, 1979), the spectrum shows a peak at 518 nm corresponding to carotenoid absorption shifting to longer wavelengths. After the flash, ΔA decays to zero in about 100 ms, see Figure 1B (a). In the absence of ATPase, the ECS decay is much slower and the "slow electrogenic phase" of electron transfer in the low-potential chain (*b*-type hemes) of cytochrome $b_6 f$ complex is observed, see Figure 2A. Subsequent flashes increase $\Delta A_{520 \text{ nm}}$ by the same amount, showing ECS is proportional to $\Delta \psi$. The amplitude of the flash-induced signal $\Delta A_{20 \text{ }\mu s}$ is used as a simple calibration of 1 charge separated per photosystem (PS) and used for normalization of ΔA . Hereafter electron transfer rates (in s⁻¹) are converted into electron flow rates (in e⁻ PS⁻¹ s⁻¹).

Initial phase of ECS - A continuous illumination induces multiphasic changes in $\Delta A_{520 \text{ nm}}$ (Figure 1B). At the onset of illumination, the photosynthetic membrane is charged as a capacitor, the conversion of photons into electrons transferred across the membrane builds-up $\Delta \psi$. Figure 1B (left) shows that ECS increases rather linearly in the first tens of milliseconds. In the absence of ATPase (Figure 2B), $\Delta A_{520 \text{ nm}}$ follows a straight slope over a much larger time range. Therefore, in order to estimate in the WT the initial photochemical flow rate I we need to fit the slope of the ECS rise on the first three datapoints ($t \le 10$ ms). I is proportional to the light intensity (Figure 2C). Inhibition of PSII has contrasting effects on the initial phase: when PSII is blocked only on the acceptor side, the initial flow rate is identical to the non-treated conditions (DCMU blocks the QA to QB electron transfer, but trans-membrane charge separation $P_{680}^{+}Q_{A}^{-}$ still occurs); while when PSII is blocked at both donor and acceptor side (DCMU and hydroxylamine), the ECS rise is about twice slower (only PSI contributes to $\Delta \psi$). This method can be used to assess the relative antenna size and the relative amounts of photosystems I and II. It also provides a way to calibrate the actinic effect of different light sources (of different wavelengths or bandwidths), converting incident photons µmol_{photons} m⁻² s⁻¹ into effective photons e⁻ PS⁻¹ s⁻¹.

Pseudo steady-state and decay phase - After all photosystems were hit by photons (amplitude $\Delta A \sim 1$, at $t \sim 20$ ms in Figure 1B, see inset on the left), the increase of ECS slows down, and a typical shoulder-peak-and-slow-descent is observed. The shape and amplitude of this curve is dependent upon the light intensity, the redox state (aerobic/anaerobic) or the nature of the sample (plant leaf or algal suspension), it is mostly influenced by the activation status of the chloroplast ATPase, the inter-conversion between $\Delta \psi$ and ΔpH , and possibly the activation of cyclic electron flow. At $t \geq 3$ s after the onset of illumination, ECS is in its "slow descent" phase, *i.e.* it does not vary much over time and it can be defined as a pseudo steady-state. When illumination stops, a sharp decay of ECS is observed (see Figure 1B right). It reflects the discharge of the membrane (efflux) that corresponded to the charge flow rate in the light, we denote J this flow at pseudo steady-state.

A transmission coefficient for electrons - In one quick measurement, ECS can therefore monitor the initial (I) and pseudo steady-state (J) rate of photosynthesis. To our knowledge, the only other way to get to these parameters is using fast-response bare O_2 electrode with

modulated illumination (Haxo and Blinks, 1950; Joliot, 1965; Myers and French, 1960). Figure 3A shows that I increases linearly with irradiance. This is because when the photosystems are open, they are not limiting for electron transfer. In contrast J, here fitted to a Michaelis-Menten equation (solid line), reaches a plateau at $V_{\text{max}} = 88 \text{ e}^{-} \text{ s}^{-1} \text{ PS}^{-1}$. It shows that photosynthesis saturates at high light intensities and suggests that photosystems are getting inactive during continuous illumination. The closure of photosystems can be assessed directly by adding a single turnover saturating flash in the slow descent at t = 3 s (see Figure 1B). The flash gives a step increase with an amplitude equal to the fraction of open centers $0 \le PS_0 \le 1$. J is always smaller than I, and the ratio $0 \le J/I \le 1$ is the yield of electron transfer. In other words J/I is the probability for an electron, initially transferred in a photosystem, to be steadily transported through the chain. This dimensionless parameter is a transmission coefficient.

The slippery transmission of photosynthesis - A photosynthetic apparatus with perfect transmission would give J/I equal to PS_o. In other words, for a first order rate reaction, the steady-state flow would be equal to the maximal rate (all centers open) in proportion to the fraction of centers remaining open at steady-state. In Figure 3B where J/I is plotted against PS_o, data points do not align diagonally: the transmission coefficient is always smaller than the relative amount of open photosystems. It suggests that photosynthetic electron transfer is affected by the diffusion of the electron carriers. It is possible that a reduced PSII acceptor does not find its way towards an oxidized PSI donor before charge recombination occurs with a rate k_{rec} .

This problem of diffusion in the photosynthetic chain can be treated with a very economical model with only one parameter $k_{\rm rec}$, a rate of charge recombination (Lavergne, 2009). We found that $k_{\rm rec}$ values comprised between 200 and 700 s⁻¹ accounted for the gap between the measured transmission coefficient J/I and the fraction of centers remaining open in the light PS_o. The solid and dotted lines correspond to an intermediary value $k_{\rm rec} = 350 \, {\rm s}^{-1}$. A rate of 350 s⁻¹ ($t_{1/2} = 2 \, {\rm ms}$) is in the range of charge recombination in PSII or PSI. At very low light intensities, when PS_o approaches 1 in Figure 3B, the vertical error bars become rather large, this is because of the division by I whose values tend to zero. Within experimental accuracy, the solid lines fit experimental data more satisfactorily than the dotted lines in Figure 3B, and also in Figure 3A where J rises more slowly than I (see Figure 3A inset). It shows that the maximum transmission coefficient, $V_{\rm max}/K_{\rm m}$ is reached when irradiance tends to zero, here 88/106 = 0.83. In other words, by analogy with mechanical engineering, the electron transfer

machinery has >17 % energy losses, even at idle speed. If one considers Photosystem I and Photosystem II as transmission shafts, the coupling in the photosynthetic chain compares to a friction belt, rubbing on the pulleys and dissipating mechanical energy as heat. This energy dissipation is accounted for in the k_{rec} parameter: frictions may encompass unpaired PSI-PSII antenna sizes (state transitions), non-photochemical quenching, non-Q_B-reducing centers, as well as stochastic motion of electron carriers.

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Increasing the frictions – In the experiments shown in Figure 4, the wavelength of the actinic light was varied from 650 to 710 nm, in order to change the excitation balance between PSI and PSII. Yet, stronger light intensities were used at less absorbed wavelengths so that the actinic effect of the light was the same (same F_i , see Figure 4B). Figure 4A shows the ECS signal is much different under 680 nm excitation than at 710 nm. When both photosystems are excited, $\Delta \psi$ builds up to a much higher level. This is not because of a slower decay of $\Delta \psi$ because in this instance $\Delta \psi$ decays much faster, 30 e⁻ s⁻¹ PS⁻¹, than when only PSI is excited, $8 e^- s^{-1} PS^{-1}$ (Figure 4C). A "red drop" of the transmission coefficient J/I corresponds to the drop of the maximum quantum yield for O2 evolution (Emerson and Lewis, 1943). It is expected since these two parameters are homogenous: they both refer to a yield. Nevertheless, O₂ evolution drops more steeply than the transmission coefficient. It shows that far-red photons, absorbed by PSI and not by PSII, do not evolve O2 but sustain cyclic electron flow around PSI, although at a much slower rate than linear electron flow, as previously reported for DCMU treated samples (Alric, 2014, 2010; Maxwell and Biggins, 1976). Under such conditions P₇₀₀⁺ accumulates and PSI is closed while PSII centers stay potentially active and contribute to PS₀. When PSI and PSII excitation is unbalanced, the transmission coefficient in the chain decreases. In the long wave-range of photosynthesis, spinning Photosystem I transmission shaft while Photosystem II keeps still increases the frictions in the chain.

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Conclusion

Half-saturation of photosynthesis was observed for a time delay of 0.02 s between short saturating flashes (Emerson and Arnold, 1932). In modern terms, it corresponds to an electron flow of $1/0.02 = 50 \text{ e}^-\text{ s}^{-1}\text{ PS}^{-1}$. In the present work, ECS measurements give a very similar value for half-saturation of photosynthesis: $V_{\text{max}}/2 = 44 \text{ e}^-\text{ s}^{-1}\text{ PS}^{-1}$ (see Figure 3A). Calibration of photosynthesis measurements on single turnover saturating flashes is therefore a way to compare results from lab to lab and from instrument to instrument. It also offers to precisely quantify the yield of photosynthesis. These "standard units" are independent from

- the chlorophyll content (sample concentration, or thickness of the leaf) and are therefore also
- 237 "portable" between photosynthetic organisms.

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- 243 Author contribution
- 244 C.M. performed experiments and analyzed data, J.A. designed research, analyzed data and
- wrote the manuscript.

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- 247 Data Availability Statement
- All data supporting the findings of this study are available within the paper.

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331 Figure legends

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- Figure 1. ECS spectrum and kinetics in *Chlorella sorokiniana* intact cells, here measured with a Joliot-type spectrophotometer.
- 335 **(A)** Spectrum of absorbance changes (ΔA), peaking at 518 nm, detected 20 μ s after a single turnover saturating flash.
- 337 **(B)** Typical kinetics of ECS, induced by a ~ 5 s continuous red illumination 338 (140 μ mol_{photons} m⁻² s⁻¹, 638nm, white bar). The signal was normalized on the amplitude of the
- single turnover flash (upward arrow) in the dark (a). Insets show the fast phases of rise and
- decay in $\Delta \psi$ at the onset (b) and offset (d) of light. I is the initial electron flow rate; J is the
- 341 steady-state flow rate. The relative amount of photosystems remaining open during the light,
- 342 PS_o, is probed with a saturating flash (c).

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- Figure 2. Inactivating ATPase extends the range where $\Delta A_{520 \text{ nm}}$ increases with absorbed photons (SL8 mutant line).
- 346 **(A)** ECS following a train of 4 flashes (upward arrows). The ECS signal induced by the first flashes is of equal amplitude (see vertical bar: $\Delta \psi$ amplitude generated by the transfer of one
- electron per photosystem). The decay of the signal is very slow (low membrane conductivity).
- 349 (B) ECS increases linearly with light duration. The linear range is larger in the ATPase-less
- strain (solid lines) than in WT (dotted lines). The initial rates, calculated between 0 and 10 ms are identical.
- 352 (C) The initial rate of ECS increases linearly with light intensity. 10 µM DCMU does not
- block charge separation in PSII, but DCMU and hydroxylamine (10 mM) do, and decrease the
- 354 initial rate.

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Figure 3. ECS light-response curves

- 357 **(A)** Initial (*I*) and steady-state (*J*) electron fluxes plotted against light intensity. *I* is fitted to a straight-line, *J* to a Michaelis-Menten model, with $V_{max} = 88 \text{ e}^-\text{PS}^{-1} \text{ s}^{-1}$ and $K_M = 106 \text{ e}^-\text{PS}^{-1}$
- 359 ¹ s⁻¹ (solid line). The inset shows an expanded view at low-light conditions. The dashed line is
- 360 when forcing $K_M = V_{max}$ in the fit.
- 361 **(B)** Plot of the transmission coefficient J/I against the fraction of open photosystems PS_o.
- 362 The solid line corresponds to a diffusive model with intrinsic losses, considering charge
- recombination at a rate $k_{\text{rec}} = 350 \text{ s}^{-1}$. The dashed line simulates the same diffusive model with
- maximum efficiency at low light ($K_M = V_{max}$, *i.e.* same initial slope for I and J, see (D) inset).
- The dotted diagonal corresponds to an "ideal" photosynthetic apparatus with perfect transmission (for $k_{rec} = 0$).
- In (A), (B) and Figure 1B, open symbols represent the same dataset (obtained at 140 μ mol_{photons} m⁻² s⁻¹).

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- 370 Figure 4. ECS action spectra
- 371 **(A)** ECS kinetics induced by 680 nm and 710 nm continuous light excitation.
- 372 **(B)** Light intensity (stronger at 710 nm than 680 nm) was adjusted to give the same initial
- 373 rise.

- 374 **(C)** Steady-state electron flow is much slower under 710 nm light excitation specific to PSI.
- 375 **(D)** Photosynthesis action spectra: red drop in O₂ quantum yield (reproduced from (Emerson
- et al., 1957)) and transmission coefficient calculated from ECS (J/I). Under 710 nm
- 377 illumination P_{700}^+ is accumulated and PSI centers are closed.