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Giuseppe Cavallaro, Giuseppe Lazzara, Frédéric Pignon, Leonardo Chiappisi, Erwan Paineau. Effect of Polymer Length on the Adsorption onto Aluminogermanate Imogolite Nanotubes. *Langmuir*, 2021, 37 (32), pp.9858-9864. 10.1021/acs.langmuir.1c01549 . hal-03365354

HAL Id: hal-03365354

<https://hal.science/hal-03365354v1>

Submitted on 5 Oct 2021

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1 Effect of Polymer Length on the Adsorption onto
2 Aluminogermanate Imogolite Nanotubes

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13

14 ABSTRACT.

15 This study evidences the adsorption of model non-ionic polymers onto aluminogermanate
16 imogolite nanotubes, an attractive porous nanofiller with potential molecular loading and
17 release applications. We resolve the underlying mechanisms between nanotubes and
18 polyethylene glycols with different molecular weight by means of nano-isothermal titration
19 calorimetry. The analysis of the results provides a direct thermodynamic characterization,
20 allowing to propose a detailed description of the energetics involved in the formation of
21 polymer/imogolite complexes. The affinity towards the nanotube surface is enthalpy-driven
22 and strongly depends on the polymer chain length, which significantly affects the polymer
23 configuration and the flow properties of the resulting complexes, probed by small-angle
24 neutron scattering and rheology, respectively. These findings open new avenues for the
25 rational design of these hybrid mixtures for advanced applications.

26

27 INTRODUCTION

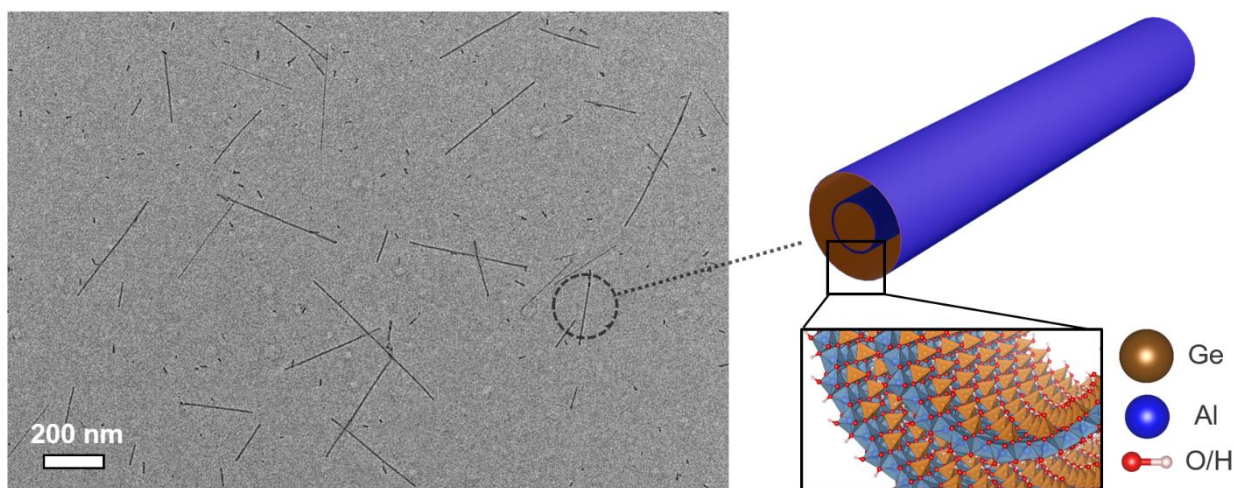
28 Imogolite nanotube (INT) is an aluminosilicate clay mineral, naturally occurring and
29 widespread in well-drained volcanic ash soils.¹ Unlike halloysite, another tubular clay
30 mineral,² INTs can be readily synthesized through hydrothermal methods.^{3,4} The structure of
31 these nanotubes is rather unique. The external walls consist of gibbsite-like sheet $\text{Al}(\text{OH})_3$
32 while the cavity interface is formed by isolated $(\text{SiO}_3)\text{OH}$ tetrahedron units connected upright
33 to the octahedral vacancies by covalent bonding.⁵ Silicon can be replaced by germanium
34 allowing to increase the diameter of the inner cavity from 1.5 to 3 nm for single-walled
35 INTs.^{6,7} Furthermore, modifications of the synthesis conditions offer a convenient way for
36 designing innovative INTs with well-defined morphologies (single or double-walled
37 structures),⁸ modular interfaces (e.g. hydrophilic or hydrophobic cavities)⁹⁻¹¹ and high
38 colloidal stability in aqueous media.¹²⁻¹⁴ All these properties have led in recent years to a
39 renewed interest in synthetic imogolite-type nanotubes as promising nano-reactors notably for
40 molecular confinement¹⁵⁻¹⁸ or photocatalytic applications.¹⁹⁻²³

41 INT offers similar characteristics to those of carbon nanotubes (CNTs) in terms of rigidity,
42 sizes and aspect ratios while retaining transparency and self-organization in highly dilute
43 conditions.^{24,25} Although they display larger band gaps than CNTs,²⁶ they are cheaper to
44 produce with monodisperse diameters without further sorting techniques. Consequently,
45 imogolite nanotubes represent an interesting alternative of the carbon nanotubes in numerous
46 applications including the fabrication of polymer/filler nanocomposites with multifunctional
47 characteristics. Within this, imogolite was successfully employed with various polymers as a
48 reinforced filler for functional nanocomposites with enhanced mechanical, optical, or
49 electrical properties, including self-healing capability.²⁷⁻³⁴ An interesting advantage of INTs
50 is the high density of hydroxyl groups on the outer surface ($\sim 18 \text{ OH}/\text{nm}^2$). It should promote
51 different possibilities of polymer adsorption,³⁵ or change in conformation as shown recently

52 by nuclear magnetic resonance.^{36,37} To date, however, the underlying mechanisms remain
53 largely unexplored.

54 One of the most used techniques to probe the polymer-nanoclay interaction mechanism is
55 isothermal titration calorimetry (ITC), which allows the experimentalist to determine the
56 thermodynamic parameters for the adsorption phenomenon through direct measurements. It
57 allowed clarifying the interactions in terms of hydrophobic/electrostatic forces and polymer
58 degree of freedom through the evaluation of the entropic contribution.^{38,39} Moreover, the
59 determination of the adsorption isotherm and the maximum amount of polymer adsorbed per
60 mass of nanoclay in comparison with the polymer gyration radius and geometric features of
61 the nanoparticle, provided an indirect representation on the polymer conformation at the
62 interface.⁴⁰ It should be noted that isothermal calorimetry is the only correct approach to
63 determine the enthalpy change of a reaction if the stoichiometry may change with temperature
64 as, in the latter case, the well-known van't Hoff approach cannot be used.⁴¹

65 Here, we investigated the adsorption of non-ionic polymers with variable molecular
66 weight on the outer surface of anisometric double-walled germanium-based imogolite (Ge-
67 DWINTs) with outer diameter of 4.5 nm and average length around 100 nm (Figure 1).¹⁴
68 Polyethylene glycols (PEGs) with variable molecular weight are selected as model
69 hydrosoluble non-ionic polymers that are frequently added with clay minerals.⁴²⁻⁴⁴ We
70 proposed the first detailed description on the energetics involved in the formation of
71 polymer/imogolite complexes by combining ITC measurements with small-angle neutron
72 scattering (SANS) and rheological experiments. Determination of thermodynamic parameters
73 allowed predicting the fraction of polymer adsorbed onto the nanotubes surface. The
74 adsorption mechanism is dependent of the PEG chain length, modifying the conformation and
75 the flow properties of the resulting mixtures.



76

77 **Figure 1.** Representative TEM image of aluminogermanate double-walled imogolite
 78 nanotubes (Ge-DWINTs). The sketch in the right illustrates the atomic structure of a Ge-
 79 DWINT with the curved $\text{Al}(\text{OH})_3$ sheet (in blue) and isolated $(\text{GeO}_3)\text{OH}$ units (in orange).
 80 Oxygen and hydrogen atoms are represented in red and white, respectively.

81 **EXPERIMENTAL SECTION**

82 **Materials.** Polyethylene glycols (PEGs) with various average molecular weight ($M_w = 200$,
 83 400 , $2k$, $3k$ and $20k \text{ g}\cdot\text{mol}^{-1}$), were purchased from Sigma Aldrich. Polymer solutions have
 84 been prepared by weight ($\pm 0.01 \text{ mg}$) and keeping under magnetic stirring overnight. Details
 85 on the synthesis and characterizations of Ge-DWINT clay nanotubes have been reported
 86 elsewhere.⁴⁵ After dialysis, as-obtained Ge-DWINTs form stable colloidal dispersions ($C =$
 87 $10 \text{ g}\cdot\text{L}^{-1}$), where the nanotubes are individually dispersed.²⁴ Mixtures of Ge-DWINTs with
 88 PEGs were prepared by adding the appropriate amount of polymer at different
 89 polymer/imogolite mass ratio ($R_{\text{P/INTs}}$) while keeping constant the imogolite concentration.

90 **Isothermal Titration Calorimetry (ITC).** ITC experiments were performed by using the
 91 ultrasensitive nano-ITC200 calorimeter (MicroCal). An amount of approximately $40 \mu\text{L}$ of
 92 the water/polymer mixture was injected into the thermally equilibrated ITC cell ($200 \mu\text{L}$)

93 containing the water/imogolite dispersion (concentration of 4.47 mg.L⁻¹). Each addition step
94 was 0.49 μL. Volumes have been calibrated by NaCl dilution experiment.⁴⁶ The concentration
95 of the polymer solutions was ca. 5 g.L⁻¹. The cell was thermally equilibrated at 25.000 ± 0.005
96 °C. The calorimeter sensitivity is at least 2 nanoW. The raw data were corrected for the
97 instrument time constant and an appropriate baseline was subtracted.

98 **Small-Angle Neutron Scattering (SANS).** SANS measurements were carried out at Institut
99 Laue-Langevin (ILL), Grenoble (France), on the instrument D11.⁴⁷ The experiments were
100 conducted at two different configurations with sample-to-detector (and collimation in
101 parenthesis) distances of 2 m (5 m), 8 m (8 m), using a wavelength λ of 6 Å and Full Width at
102 Half Maximum (FWHM) of 10%. The scattering vector q was calculated as $4\pi \sin(\theta/2)/\lambda$, θ
103 being the scattering angle. The two-dimensional patterns were corrected for the detector
104 efficiency using the scattering of a 1 mm H₂O sample and for the dark current signal. The
105 contribution from the H₂O/D₂O solvent was subtracted and finally, the patterns were radially
106 averaged, the scattering being isotropic. Data reduction was performed with BERSANS,⁴⁸
107 while SASfit 0.94.11 software was used for the analysis of SANS curves.⁴⁹ The experiments
108 were performed diluting the stack Ge-DWINTs dispersion with D₂O and the final solvent was
109 H₂O/D₂O with a mass ratio of 1/10. The density ($\pm 1 \times 10^{-6}$ g cm⁻³) of the Ge-DWINTs
110 dispersion was determined by using a DSA 5000 M (Anton Paar) at 25 °C (± 0.001 °C). The
111 specific volume of the INTs was calculated as:

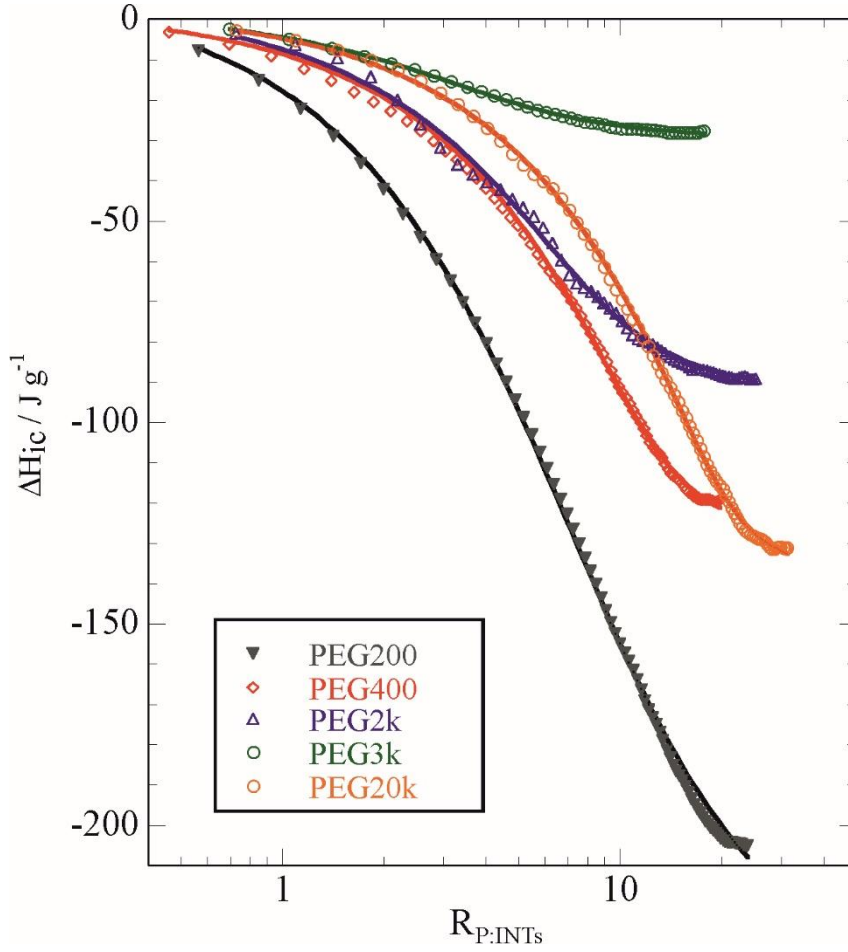
$$v_{sp} = (1/d) - 10^3 \frac{(d - d_0)}{w d d_0} \quad (1)$$

112 being d the experimental density for Ge-DWINTs dispersion at concentration w (expressed as
113 g of Ge-DWINTs/kg of solvent) and d_0 is the water density.

114 **Rheology.** Rheological measurements were performed on a controlled-stress rheometer
115 (Discovery HR-3, TA Instruments) using a cone-plate geometry (60 mm in diameter, 1° in
116 angle) and a solvent trap cover to prevent sample evaporation. Flow curves of PEGs/Ge-
117 DWINTs ($R_{P/INTs} = 10$) were measured under a controlled shear rate from 1 to 2200 s⁻¹.

118 **RESULTS AND DISCUSSION**

119 **Thermodynamic characterization of PEG/Ge-DWINTs dispersions.** The thermodynamics
120 of the polymer adsorption onto aqueous dispersions of Ge-DWINTs was investigated by ITC
121 through the stepwise injection method. The effects of dilution of the polymer and imogolite
122 were measured and subtracted from the heats of titration to obtain the thermal effects of the
123 polymer/imogolite interaction at each step. Based on the ITC data, we determined the
124 cumulative variation of enthalpy (ΔH_{ic}) at variable polymer/imogolite mass ratio ($R_{P/INTs}$) for
125 PEGs with different molecular weight (Figure 2). Details for the calculation of ΔH_{ic} are
126 presented in Supporting Information.



127

128 **Figure 2.** Isothermal titration data of Ge-DWINTs with PEGs of different molecular weight.

129 Lines correspond to best fits according to Langmuir adsorption model.

130 As for other polymer/clay systems,^{38,40} the ΔH_{ic} vs $R_{P:INTs}$ trends correspond to a single

131 adsorption process that were successfully described by the Langmuir adsorption model

132 (Figure 2), providing the enthalpy (ΔH_{ads}^0) and the equilibrium constant (K_{ads}) for the PEGs

133 adsorption onto imogolite surfaces. According to the K_{ads} and ΔH_{ads}^0 values, we estimated the

134 standard free energy (ΔG_{ads}^0) and the entropy (ΔS_{ads}^0) for the adsorption process as:

$$\Delta G_{ads}^0 = -RT \ln K_{ads} \quad (2)$$

$$T\Delta S_{ads}^0 = \Delta H_{ads}^0 - \Delta G_{ads}^0 \quad (3)$$

135 In addition, the fitting of ITC data (Figure 2) allowed us to determine the PEGs/INTs
136 stoichiometry (Z) at the saturation point. Namely, Z represents the maximum amount of PEGs
137 moles adsorbed per gram of imogolite nanotubes. The fitting analysis of the experimental
138 ΔH_{ic} vs $R_{P/INTs}$ functions are detailed in Supporting Information.

139 Table 1 collects the thermodynamic parameters and the stoichiometry values for the
140 adsorption of PEGs with variable molecular weights onto imogolite.

141 **Table 1.** Investigated polyethylene glycol (PEG) with various molecular weight and the corresponding adsorption parameters deduced from ITC
 142 measurements performed at 298 K.

Polymer	Molecular weight (g.mol ⁻¹)	K _{ads} (dm ³ mol ⁻¹)	Z ^a (mmol _P g _{INTs} ⁻¹)	ΔH ⁰ _{ads} (kJ mol ⁻¹)	ΔG ⁰ _{ads} (kJ mol ⁻¹)	ΔS ⁰ _{ads} (kJ mol ⁻¹ K ⁻¹)
PEG200	200	(9.5 ± 1.7) · 10 ⁴	37 ± 2	-2.15 ± 0.17	-28.8 ± 0.5	0.088 ± 0.007
PEG400	400	(1.6 ± 0.4) · 10 ⁶	30.6 ± 0.7	-2.18 ± 0.16	-35.4 ± 0.7	0.115 ± 0.004
PEG2k	2000	(3.3 ± 1.7) · 10 ⁶	2.5 ± 0.2	-9.65 ± 1.24	-37.2 ± 1.6	0.09 ± 0.01
PEG3k	3000	(8.0 ± 3.0) · 10 ⁶	4.8 ± 0.2	-9.3 ± 0.7	-39.4 ± 1.2	0.101 ± 0.010
PEG20k	20000	(1.4 ± 0.3) · 10 ⁸	1.00 ± 0.02	-55.5 ± 1.6	-46.5 ± 0.7	-0.030 ± 0.001

143 ^a Moles of polymer chain per gram of imogolite nanotubes

As evidenced from Table 1, the polymer molecular weight affects both the thermodynamic parameters and the maximum mass of PEG adsorbed onto Ge-DWINTs. The equilibrium constant displays a monotonic evolution according to the PEG chain length and it is related to different affinity. As shown for clay nanosheets, a limited affinity for the clay surface was observed for PEG with low molecular weight (1000-4000 g.mol⁻¹) whereas higher values resulted in an increase of the adsorbed mass of polymer in conjunction with a change of polymer conformation.^{50,51} The negative values of standard free energy indicate that PEG adsorption onto Ge-DWINTs surface is energetically favored. Both ΔG_{ads}° and ΔH_{ads}° decrease with PEG molecular weight, which indicates an enthalpy driven adsorption process for all the mixtures. Given that the enthalpy of adsorption per mole of polymer is always negative, PEG/imogolite interactions are the dominant forces in the adsorption process although the enthalpic contribution is much more relevant for polymers with a longer chain. As concerns ΔS_{ads}° , we determined positive values for all the PEGs used in this work (except for PEG20k). In general, the change in entropy is related to the following contrasting effects: (i) a loss of configuration freedom of the polymer in the adsorbed state, which induces a decrease of entropy ($\Delta S_{ads}^{\circ} < 0$) or (ii) a water release from the polymer ($\Delta S_{ads}^{\circ} > 0$). The positive entropy values highlight that de-hydration from the adsorbed PEG represents the dominant entropic factor during the adsorption. On the other hand, $\Delta S_{ads}^{\circ} < 0$ was observed for the highest molecular weight suggesting that the de-hydration phenomenon is less relevant for this mixture.

The maximum amount of PEG molecules adsorbed onto the imogolite surface decreases sharply with the polymer molecular weight (*Z* values in Table 1) and therefore the most hydrated structure is expected for PEG20k compared to the other polymers. Based on the thermodynamic parameters (Table 1), we can deduce the fraction of polymer adsorbed onto imogolite surfaces at

variable $R_{P/INTs}$. Based on the fitting parameters of ITC curves (see details in Supporting Information), we calculated the fraction of adsorbed polymer onto imogolite as a function of the $R_{P/INTs}$ for PEG2k/INTs and PEG20k/INTs complexes (Figure 3).

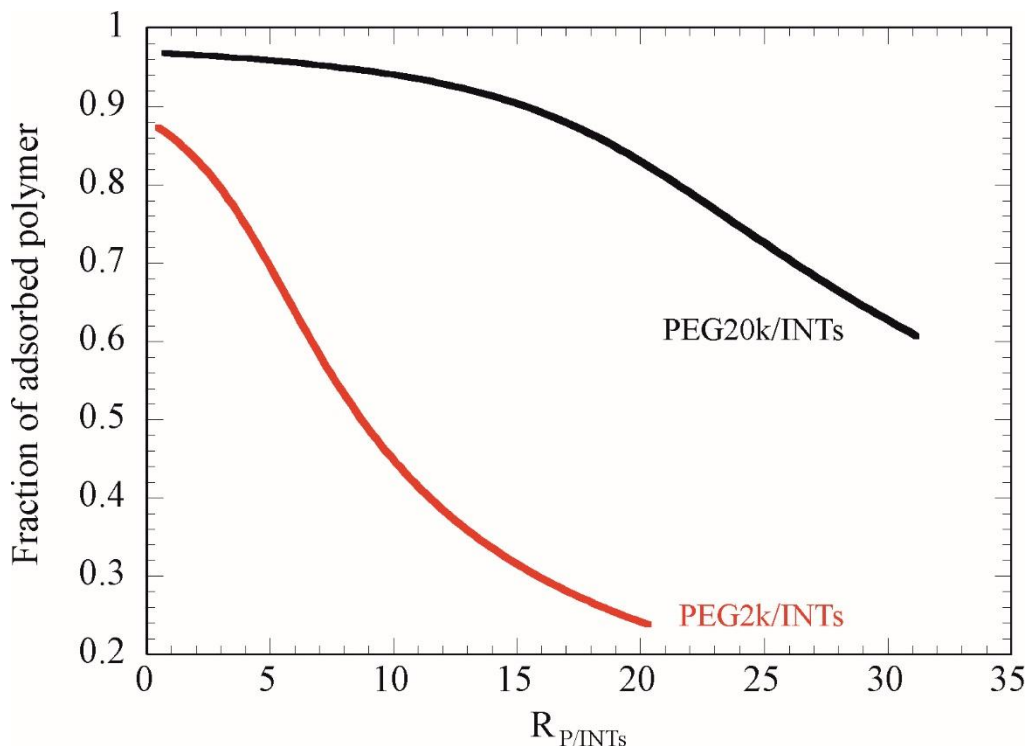


Figure 3. Calculated fraction of polymer adsorbed onto Ge-DWINTs at variable $R_{P/INTs}$ for Ge-DWINTs complexes with PEG2k (red) and PEG20k (black).

Within the investigated polymer/Ge-DWINTs range, we estimated a larger adsorbed fraction for PEG with the longest chain. As example, for $R_{P/INTs} = 10$, the fractions of polymer adsorbed onto Ge-DWINTs are ca. 0.4 and 0.95 for PEG2k and PEG20k, respectively. The different adsorption profiles are the result of the equilibrium constant, weighted by the maximum possible amount of polymer that can be allocated at the solid/liquid interface. In other words, there is more place for small polymers (expressed in total mass), but less tendency to fill that space as evidenced by the equilibrium constants.

The dependence of the affinity toward the Ge-DWINTs surface from the PEG molecular weight was further investigated by SANS for PEG2k and PEG20k. It should be noted that the scattering from dilute Ge-DWINTs dispersions (1 g.L⁻¹) in the same solvent has a negligible flat scattering intensity in the investigated scattering vector regime (Figure S1). As a general feature (Figure 4), the addition of Ge-DWINTs to PEG solutions does not influence the shape of the curves. All scattering curves were analysed with a model of a Gaussian coil⁵²:

$$I(q) = I_0 \frac{2[\exp(-q^2 R_g^2) + q^2 R_g^2 - 1]}{q^4 R_g^4} \quad (4)$$

with I_0 and R_g being the forward scattering intensity and the radius of gyration of the chain, respectively. The values are given in Table 2. The overlap concentration was calculated from the values of the radii of gyration, as:

$$c^* = \frac{Mw}{(4\pi/3)R_g^3 N_A} \quad (5)$$

Confirming that the experiment were performed in the dilute regime, the interactions between the polymer chains can be estimated between the ratio of the determined forward scattering and the predicted value,⁵³ calculated as:

$$I_0^{calc} = \frac{\phi_p Mw \Delta\rho^2}{d_p N_A} \quad (6)$$

In eq. 6, ϕ_p is the polymer concentration in volume fraction, d_p its density (a value of 1.199 g cm⁻³ was used)⁵⁴, and $\Delta\rho$ the scattering contrast. A ratio of 0.5 and 0.1 for PEG2k and PEG20k, respectively, between the experimental and the calculated forward scattering intensity indicates a

significant repulsive interaction between the polymer chains, in agreement with previous studies (Table 2).⁵³

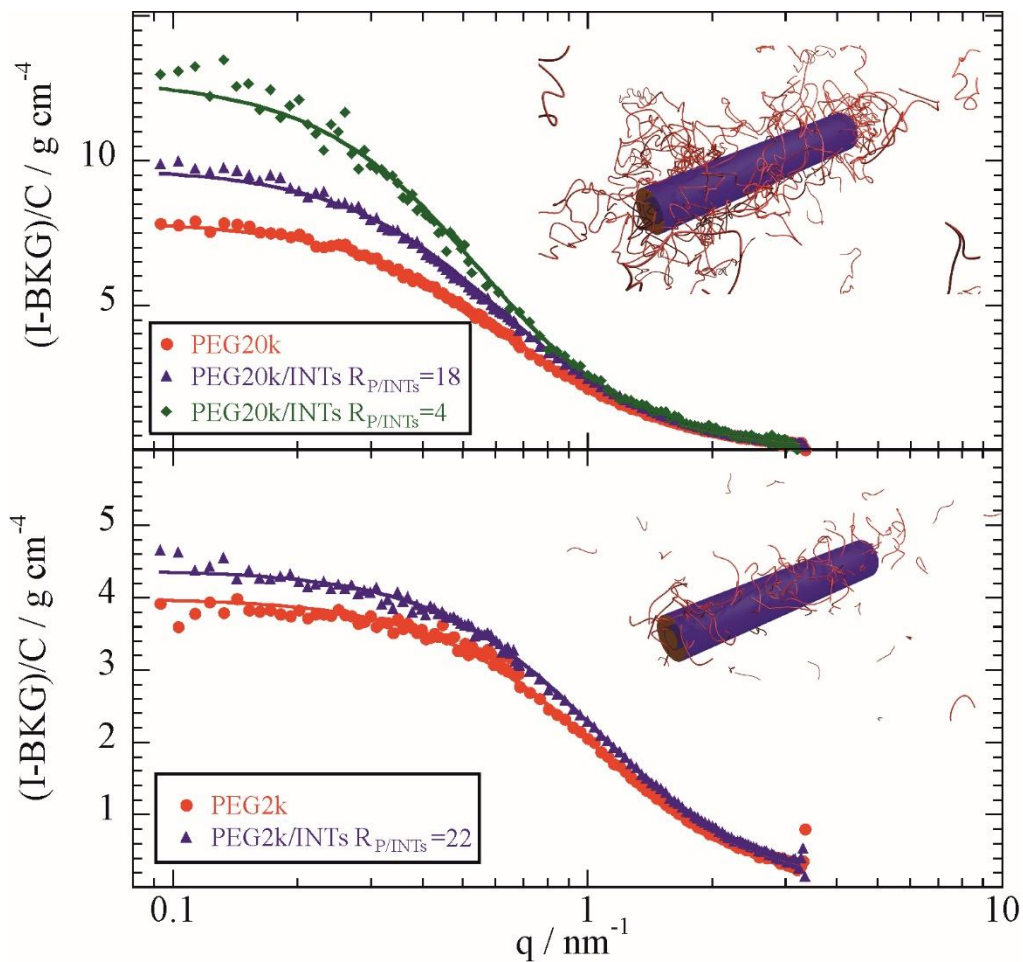


Figure 4. Neutron Scattering Intensity function after background (BKG) subtraction and polymer concentration normalization (C_p) of PEGs in the presence and absence of Ge-DWINTs (1 g.L^{-1}) at $25 \text{ }^\circ\text{C}$. Solvent was $\text{H}_2\text{O}/\text{D}_2\text{O}$ with a mass ratio of 1/10. The continuous lines correspond to the fit of the scattering curves (see text for details).

Table 2. Forward scattering intensity I_0 and radius of gyration R_g deduced from the modelling of SANS curves with a Gaussian coil model.

System	$R_{P/INTs}$	I_0^{exp} (cm^{-1})	R_g (nm)	I_0^{calc} (cm^{-1})	$\frac{I_0^{exp}}{I_0^{calc}}$
PEG2k	-	0.08 ± 0.01	1.56 ± 0.08	0.16	0.51
PEG2k/INTs	22	0.10 ± 0.02	1.54 ± 0.05	0.15	0.62
PEG20k	-	0.18 ± 0.02	2.49 ± 0.06	1.61	0.11
PEG20/INTs	18	0.17 ± 0.02	2.61 ± 0.06	1.22	0.14
	4	0.05 ± 0.01	3.02 ± 0.08	0.30	0.18

The analysis of the SANS data shows that the conformation of the polymer is not, or only minimally, affected by the presence of the clay nanotube, as deduced from constant value of radius of gyration obtained. In contrast, an increase in scattering intensity is observed upon addition of Ge-DWINTS to the dilute PEG solution. The origin of this increase has to be found in the attractive interaction between the nanoclay and the polymer. In fact, both the adsorption of PEG on the Ge-DWINTs and a reduced repulsion between the PEG chains in solution would result in an increase of the forward scattering intensity. The effect is more pronounced for the longer PEG chains for which the radius of gyration increases from 2.5 to 3 nm with decreasing $R_{P/INTs}$ (Table 2). This slight stretching of PEG20k chains can be ascribed to a crowding phenomenon due to the larger amount of polymer adsorbed onto Ge-DWINTs predicted by the ITC experiments, ca 0.96 compared to 0.87 for $R_{P/INTs} = 4$ and 18, respectively. The stretching of the polymer chain is an indication of a relatively compact organization at the solid/liquid interface. This is confirmed by the grafting density of ca. 1.6 chains nm^{-2} , estimated from Z values reported in Table 1 and by using the experimental specific volume of Ge-DWINTs ($0.3949 \text{ cm}^3 \text{ g}^{-1}$) and the characteristic sizes of the nanotubes.¹⁴ It should be noted that the fits of the scattering curves for PEG2k provided a radius of gyration of 1.5 nm with and without Ge-DWINTs in agreement with the thermodynamic prediction of negligible interactions in this case.

Therefore, we should expect different rheological and colloidal behaviors for these systems. In particular, depletion processes can play a role for PEG2k/Ge-DWINTs and steric colloidal stabilization for PEG20k/Ge-DWINTs (as a consequence of the large amount of polymer adsorbed onto Ge-DWINTs surfaces).

We investigated the flow properties of PEGs/Ge-DWINTs complexes (PEG2k and PEG20k) prepared at a fixed $R_{P/INTs} = 10$ while varying the concentration of imogolite nanotubes (C_{INTs}). In this case, the fraction of polymer adsorbed onto Ge-DWINTs is the same whatever the value of C_{INTs} but varies by a factor of 2 between the two series (Figure 3). This should allow us to probe how the difference in amount of PEG adsorbed onto the nanotube surface impact the viscosity of the suspensions. The flow curves (viscosity η vs. shear stress σ) exhibit a shear thinning behavior whatever the molecular weight or the nanotube concentration (Figure 5a,b). With increasing shear stress, hydrodynamic interactions tend to overcome electrostatic interaction between PEGs/Ge-DWINTs complexes that end up aligning in the flow. To go further, we reproduced flow curves by adjusting a viscosity model based on the Quemada's approach:^{55,56}

$$\eta(\sigma) = \eta_{\infty} \left(\frac{1 + \sigma/\sigma_c}{\chi + \sigma/\sigma_c} \right)^2 \quad (7a)$$

$$-1 \leq \chi \leq 1 \quad (7b)$$

where η_{∞} represents the viscosity for infinite shear stress and σ_c is a critical shear stress. Its value is close to σ when hydrodynamics effects are comparable to Brownian and interaction energies. χ is positive when a viscosity plateau at 0-shear is measured while it is negative when a yield stress σ_y occurs in the dispersion such as: $\sigma_y = -\chi\sigma_c$. By adjusting eq. 7a using the least-

square method, fair agreement was obtained between experimental and calculated flow curves for both PEG2k and 20k/Ge-DWINTs complexes (Figure 5a,b). The evolution of yield stress σ_y with Ge-DWINTs concentration depending the PEG molecular weight is represented in Figure 5c. The yield stress remains the same for low Ge-DWINTs concentrations but starts to increase when C_{INTs} is above 3 g.L^{-1} .

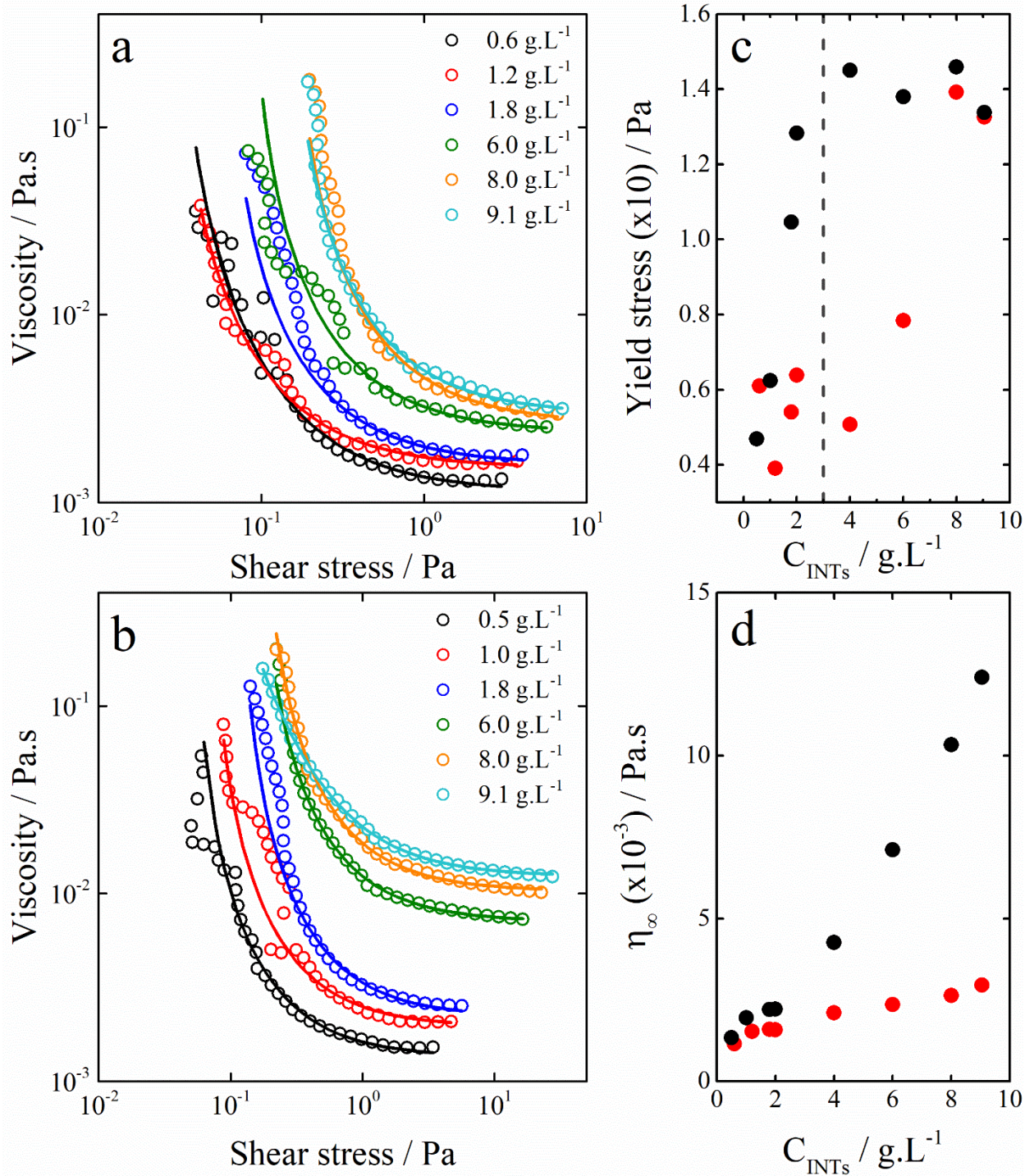


Figure 5. Flow curves for PEGs/imogolite mixtures prepared at different C_{INTs} and PEGs molecular weight: (a) $M_w = 2000$ g.mol⁻¹ and (b) $M_w = 20000$ g.mol⁻¹. Continuous lines

correspond to the best fit of eq. [7a]. Evolution of (c) yield stress and (d) infinite viscosity values with imogolite concentration for PEG2k (red) and PEG20k (black).

For the longest PEG chain, we observed a sharp increase of the yield stress right from the first concentrations investigated until reaching a constant value around 0.15 Pa. This difference is directly related to the larger amount of PEG adsorbed on the imogolite surface for PEG20k and the modification of the polymer conformation. This is particularly noticeable when looking at the evolution of infinite shear viscosity (Figure 5d). An almost linear increase is observed with varying the imogolite concentration but the values are always lower for PEG2k compared to PEG20k. Bridging interactions between the nanotubes seems unlikely since the average distance is at least higher than 60 nm for the most concentrated suspension ($\sim 9 \text{ g.L}^{-1}$) that is ~ 14 times the average diameter of the nanotubes.^{24,25} PEG/INTs mixtures rather look like a “decorated” nanotube with more or less amount of polymer as sketched in Figure 4. For PEG20k, steric interactions between the polymer and the nanotubes may form PEG/Ge-DWINTs complexes that are larger than the nanotubes themselves, which interfere more significantly thus raising the viscosity when shearing.

CONCLUSIONS

In summary, we investigated the adsorption mechanism of non-ionic polymers onto geo-inspired inorganic nanotubes in aqueous dispersions. Modeling nano-ITC measurements allowed us to quantify the standard variation of enthalpy, entropy and free energy during the adsorption of PEGs. Their affinity with the nanotube surface is dominated by favorable enthalpic contribution. We also evidenced that the adsorption mechanism depends on the length of the polymer molecular weight. Within the investigated polymer/Ge-DWINTs range, we estimated that more than 95% of polymer is adsorbed with the longest chain (PEG20k) against 40% for a 10-fold

lower polymer molecular weight. The role played by the molecular weight on the adsorption of PEG drastically impacts the high shear flow properties of the resulting complexes due to different interaction and conformation of the polymer coil at the interface with the nanotubes. These findings will open a fresh look on fundamental understanding of polymers with these modeled nanotubes but will also aid in an improved rational design of hybrid polymer/imogolite mixtures during the elaboration of nanocomposites.

ASSOCIATED CONTENT

Supporting Information. Details on the ITC data analysis and additional SANS data of Ge-DWINTs dispersion. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Funding Sources

Laboratoire Rhéologie et Procédés is part of LabEx Tec21 (Investissements d'Avenir - grant agreement no. ANR-11-LABX-0030), PolyNat Carnot Institut (Investissements d'Avenir - grant agreement #ANR-16-CARN-0025-01) and Glyco@Alps programme (Investissements d'Avenir - grant agreement #ANR-15-IDEX-02). G.L. and G.C. thank the University of Palermo for financial Support (FFR2020).

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENT

The present work has benefited from Imagerie-Gif core facility supported by l'Agence Nationale de la Recherche (ANR-11-EQPX-0029/Morphoscope; ANR-10-INBS-04/FranceBioImaging; ANR-11-IDEX-0003-02/Saclay Plant Sciences). SANS data (<https://doi.ill.fr/10.5291/ILL-DATA.9-12-557>) using D11 at Institut Laue Langevin took place with the approved proposal 9-12-557. The Partnership for Soft Condensed Matter (PSCM) is acknowledged for provided the laboratory infrastructure for sample preparation and pre-characterization.

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