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# EXAMPLES OF FINITE TIME BLOW UP IN MASS DISSIPATIVE REACTION-DIFFUSION SYSTEMS WITH SUPERQUADRATIC GROWTH 

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#### Abstract

We provide explicit examples of finite time $L^{\infty}$-blow up for the solutions of $2 \times 2$ reaction-diffusion systems for which three main properties hold: positivity is preserved for all time, the total mass is uniformly controlled and the growth of the nonlinear reaction terms is superquadratic. They are obtained by choosing the space dimension large enough. This is to be compared with recent global existence results of uniformly bounded solutions for the same kind of systems with quadratic or even slightly superquadratic growth depending on the dimension. Such blow up may occur even with homogeneous Neumann boundary conditions. All these $L^{\infty}$-blowing up solutions may be extended as weak global solutions. Blow up examples are also provided in space dimensions one, two and three with various growths.


1. Introduction. The main goal of this paper is to provide examples of blow up in finite time for $2 \times 2$ reaction-diffusion systems of the form

$$
\left\{\begin{array}{l}
\partial_{t} u_{1}-d_{1} \Delta u_{1}=f_{1}\left(u_{1}, u_{2}\right) \text { on }(0, T) \times B_{N}  \tag{1}\\
\partial_{t} u_{2}-d_{2} \Delta u_{2}=f_{2}\left(u_{1}, u_{2}\right) \text { on }(0, T) \times B_{N} \\
\text { "good " boundary conditions on }(0, T) \times \partial B_{N} \\
u_{1}(0, \cdot)=u_{1}^{0} \geq 0, u_{2}(0, \cdot)=u_{2}^{0} \geq 0
\end{array}\right.
$$

where $B_{N}$ is the open unit ball in $\mathbb{R}^{N}, d_{1}, d_{2} \in(0,+\infty), f_{1}, f_{2}:[0,+\infty)^{2} \rightarrow \mathbb{R}$ are regular nonlinearities such that the nonnegativity of the solutions to (1) is preserved for all time which, as it is well-known, is equivalent to

$$
\begin{equation*}
\text { (P) quasipositivity: } f_{1}\left(0, s_{2}\right) \geq 0, f_{2}\left(s_{1}, 0\right) \geq 0, \forall s_{1}, s_{2} \in[0,+\infty) \tag{2}
\end{equation*}
$$

and such that the following 'mass dissipativity' condition holds

$$
\begin{equation*}
\text { (M) dissipativity: } f_{1}\left(s_{1}, s_{2}\right)+f_{2}\left(s_{1}, s_{2}\right) \leq 0, \forall s_{1}, s_{2} \in[0,+\infty) \tag{3}
\end{equation*}
$$

This last property implies that the $L^{1}\left(B_{N}\right)$-norm of the solutions $u_{1}(t), u_{2}(t)$ of (1) does not blow up in finite time ( $=$ control of the total mass for all time).

[^0]It is well-known that, for $u_{1}^{0}, u_{2}^{0} \in L^{\infty}\left(B_{N}\right)^{+}$, there exist $T \in(0,+\infty]$ and a classical solution to (1) on $[0, T)$. If all the $d_{i}$ are equal, then this solution can be extended to a global classical solution on $[0,+\infty)$. However, if the $d_{i}$ are different from each other, it is known that the two conditions $(\mathbf{P}),(\mathbf{M})$ are in general not sufficient to provide global classical solutions on [0, + ) as proved in [11], [12] where examples of solutions blowing up in $L^{\infty}\left(B_{N}\right)$ in finite time are described for this class of systems.

A main new result proved here (see Theorem 2.1) is that actually, for all $d=$ $2+\eta, \eta>0$, there exists a choice of $d_{1}, d_{2} \in(0,+\infty)$, of the space dimension $N$ and of the nonlinearities $f_{1}, f_{2}$ satisfying $(\mathbf{P}),(\mathbf{M})$ and the growth condition $(\mathbf{G})$ below, for which the solution to (1) blows up in $L^{\infty}\left(B_{N}\right)$ in finite time $T<+\infty$ :

$$
\begin{equation*}
\text { (G) d-growth: }\left|f_{i}\left(s_{1}, s_{2}\right)\right| \leq C_{0}+C_{1}\left(s_{1}+s_{2}\right)^{d}, \forall s_{1}, s_{2} \in[0,+\infty), i=1,2 \tag{4}
\end{equation*}
$$

where $C_{0}, C_{1} \in[0,+\infty)$. The point is that finite time blow up may occur for any superquadratic growth.

This result has to be analyzed in parallel with the recent results in [4], [5] saying that, for any dimension $N$, there exists $\eta>0$, depending on the dimension $N$, such that, if the growth of $f_{1}, f_{2}$ at infinity is at most $d=2+\eta$, then all solutions of (1) are global in time: in particular, the $L^{\infty}\left(B_{N}\right)$-norm of $u_{1}(t), u_{2}(t)$ is bounded on all intervals $[0, T]$. Note that these global existence results are proved in [4], [5] even for all $m \times m$ systems, $m \geq 2$, with the properties $(\mathbf{P}),(\mathbf{M}),(\mathbf{G})$. Theorem 2.1 shows that these results are in some sense optimal. Besides [4], [5], several papers had provided global existence results of classical solutions for close to quadratic systems like [7], [6], [2], [14]. More results may also be found in [3], [15], [8], [1], [13]. We refer to the introduction in [4] where a nice state of art on these questions is described.

Obviously, given $\eta>0$ and $f_{1}, f_{2}$ with a $2+\eta$ growth, to obtain blow up examples, it is necessary to work in space dimensions $N$ which are higher and higher when $\eta$ is smaller and smaller.

Four more facts are also interesting about these blow up examples.

1) First it turns out that, not only the $L^{\infty}$-norm blows up at time $t=T$, but also the $L^{m}$-norm for all $m \geq N(1+\eta) / 2, N$ large.
2) Next these $T$-blowing up solutions can nevertheless be extended as global weak solutions on $[0,+\infty)$ of the corresponding system (1) (see Section 7).
3) We also prove (see Theorem 2.2) ) that the same kind of solutions provide blow up examples for the following kind of systems:

$$
\left\{\begin{array}{l}
\partial_{t} u_{1}-d_{1} \Delta u_{1}=c_{1}(t, x) u_{1}^{\alpha} u_{2}^{\beta} \text { on }(0, T) \times B_{N}  \tag{5}\\
\partial_{t} u_{2}-d_{2} \Delta u_{2}=c_{2}(t, x) u_{1}^{\alpha} u_{2}^{\beta} \text { on }(0, T) \times B_{N} \\
" \text { good " boundary conditions on }(0, T) \times \partial B_{N}, \\
u_{1}(0, \cdot)=u_{1}^{0} \geq 0, u_{2}(0, \cdot)=u_{2}^{0} \geq 0
\end{array}\right.
$$

where $c_{1}+c_{2} \leq 0, c_{1}, c_{2} \in L^{\infty}\left(Q_{T}\right)$ and $\alpha+\beta=2+\eta$ with $\eta>0$ as small as we want.
4) The boundary conditions of the examples provided in Theorem 2.1 are nonhomogeneous boundary conditions. We also describe blowing up examples for systems with homogeneous Neumann boundary conditions, up to replacing $f_{i}, i=1,2$ by more general (still regular) nonlinearities $g_{i}=g_{i}\left(t, x, s_{1}, s_{2}\right)$. This is indicated in Theorem 2.3.

An optimal dependence of $N$ in terms of $\eta$ or $d$ is not explicitly known in the main result of Theorem 2.1. What comes out in the proof of this theorem [see (28)], is that one should choose $N$ large enough so that

$$
\frac{2 N-\theta}{N-\theta} \leq d=2+\eta \quad \Leftrightarrow \quad N \geq \theta\left(1+\eta^{-1}\right)
$$

where $\theta \in(4 / 3,+\infty)$. This means that we should have $N>4\left(1+\eta^{-1}\right) / 3$.
Another natural question is, given the dimension $N$, to find the optimal growth which provides global existence (or blow up in finite time) for nonlinearities satisfying the main properties (P), (M). As proved in [11], [12], examples of blow up may occur even in dimension $N=1$ by choosing a large enough growth for the nonlinearity. It is known that in dimension $N=1$, global existence is proved for at most cubic growth [see [6], [15]], but the optimal growth is not known.

To progress in this understanding, we give in Section 8 explicit blow up examples in the following cases :

$$
N=1, d=6 ; \quad N=2, d=7 / 2 ; \quad N=3, d=3 .
$$

2. Statement of the main results. Some notation. Here $B_{N}$ denotes the open euclidian unit ball in $\mathbb{R}^{N}, Q_{T}=(0, T) \times B_{N}, \Sigma_{T}=(0, T) \times \partial B_{N}$. We also denote by $C^{\infty}\left(A_{1}, A_{2}\right)$ the family of $C^{\infty}$-mappings from $A_{1}$ to $A_{2}$ where $A_{i} \subset \mathbb{R}^{N_{i}}, i=1,2$, and $N_{1}, N_{2}$ are positive integers. If $A_{2}=\mathbb{R}$, we simply denote $C^{\infty}\left(A_{1}\right)$.

Theorem 2.1. Let $d:=2+\eta, \eta \in(0,+\infty)$ and $T \in(0,+\infty)$. Then there exist $f_{1}, f_{2} \in$ $C^{\infty}\left([0,+\infty)^{2}\right)$ satisfying $(\mathbf{P})+(\mathbf{M})+(\mathbf{G}), d_{1}, d_{2} \in(0,+\infty)$, a dimension $N$ large enough, $u_{1}^{0}, u_{2}^{0} \in C^{\infty}\left(\overline{B_{N}}\right)^{+}, \alpha_{1}, \alpha_{2} \in C^{\infty}([0, T])^{+}$and $u_{1}, u_{2}$ nonnegative $C^{\infty}$-solutions of (1) with

$$
u_{1}(t, x)=\alpha_{1}(t), \quad u_{2}(t, x)=\alpha_{2}(t) \text { on } \Sigma_{T},
$$

and such that

$$
\lim _{t \rightarrow T^{-}}\left\|u_{1}(t)\right\|_{L^{m}\left(B_{N}\right)}=\lim _{t \rightarrow T^{-}}\left\|u_{2}(t)\right\|_{L^{m}\left(B_{N}\right)}=+\infty
$$

for $m \geq N(d-1) / 2$. Moreover, for $\lambda$ close to 1

$$
\begin{equation*}
\left(\mathbf{M}_{\lambda}\right) f_{1}\left(s_{1}, s_{2}\right)+\lambda f_{2}\left(s_{1}, s_{2}\right) \leq 0, \quad \forall s_{1}, s_{2} \in[0,+\infty) \tag{6}
\end{equation*}
$$

Remark 2.1. Despite the blow up of $u_{1}(t), u_{2}(t)$ at time $t=T$, it turns out that the solutions $u_{i}$ provided by Theorem 2.1 can be extended to the whole interval $[0,+\infty)$ as weak global solutions of the system (1) on the whole interval $[0,+\infty)$. By weak solution, we mean that the nonlinear terms $f_{i}\left(u_{1}, u_{2}\right), i=1,2$ are in $L^{1}\left((0, \tau) \times B_{N}\right)$ for all $\tau \in(0,+\infty)$ and $u_{1}, u_{2}$ are solutions in the sense of distributions or in the sense of the 'variation of constants' formula. This is a consequence of the fact that, in the blow up examples, not only $f_{1}+f_{2} \leq 0$ holds, but also $f_{1}+\lambda f_{2} \leq 0$ for some $\lambda \neq 1$. This provides the $L^{1}$-bound on the nonlinear terms and this is sufficient for the existence of weak global solutions. All this is made more precise in Section 7.
Theorem 2.2. Let $d:=2+\eta, \eta \in(0,+\infty), \alpha, \beta \in(1,+\infty)$ with $\alpha+\beta=d$ and $T \in$ $(0,+\infty)$. Then there exist a dimension $N$ large enough, $d_{1}, d_{2} \in(0,+\infty)$ and

$$
\left\{\begin{array}{l}
c_{1}, c_{2} \in C^{\infty}\left(\left[[0, T) \times \overline{B_{N}}\right) \cap L^{\infty}\left(Q_{T}\right),\right. \\
u_{1}^{0}, u_{2}^{0} \in C^{\infty}\left(\overline{B_{N}}\right)^{+}, \alpha_{1}, \alpha_{2} \in C^{\infty}([0, T])^{+},
\end{array}\right.
$$

and $u_{1}, u_{2}$ nonnegative $C^{\infty}$-solutions of the system (5) with $u_{i}=\alpha_{i}, i=1,2$ on $\Sigma_{T}$ and such that

$$
\begin{gathered}
c_{1}(t, x)+c_{2}(t, x) \leq 0, \quad \forall(t, x) \in \overline{Q_{T}} \\
\lim _{t \rightarrow T^{-}}\left\|u_{1}(t)\right\|_{L^{m}\left(B_{N}\right)}=\lim _{t \rightarrow T^{-}}\left\|u_{2}(t)\right\|_{L^{m}\left(B_{N}\right)}=+\infty
\end{gathered}
$$

for $m \geq N(d-1) / 2$.
Remark 2.2. The solutions $u_{1}, u_{2}$ appearing in this result are the same as those in Theorem 2.1. Only the interpretation of the nonlinear part changes. In particular, they blow up in $L^{m}\left(B_{N}\right)$ as $t \rightarrow T^{-}$for the same values of $m$.

It turns out that we can also use the same kind of solutions to construct new blowing up solutions for similar systems, but with homogeneous Neumann boundary conditions. Here the nonlinearities depend also on $(t, x)$. This is the purpose of the next theorem.

Theorem 2.3. Let $d:=2+\eta, \eta \in(0,+\infty)$ and $T \in(0,+\infty)$. Then there exist $\tau_{0} \in$ $[0, T), g_{1}, g_{2} \in C^{\infty}\left(\left[\tau_{0}, T\right) \times \overline{B_{N}} \times[0,+\infty)^{2}\right)$ satisfying $\left(\mathbf{P}^{\prime}\right)+\left(\mathbf{M}_{\lambda}^{\prime}\right)+\left(\mathbf{G}^{\prime}\right)$ below, $d_{1}, d_{2} \in$ $(0,+\infty)$, a dimension $N$ large, $v_{1}^{0}, v_{2}^{0} \in C^{\infty}\left(\overline{B_{N}}\right)^{+}$and $v_{1}, v_{2}$ nonnegative $C^{\infty}$-solutions of

$$
\left\{\begin{array}{l}
\partial_{t} v_{1}-d_{1} \Delta v_{1}=g_{1}\left(t, x, v_{1}, v_{2}\right) \text { in }\left(\tau_{0}, T\right) \times B_{N},  \tag{7}\\
\partial_{t} v_{2}-d_{2} \Delta v_{2}=g_{2}\left(t, x, v_{1}, v_{2}\right) \text { in }\left(\tau_{0}, T\right) \times B_{N}, \\
\partial_{\nu} v_{1}=0=\partial_{\nu} v_{2} \text { on }\left(\tau_{0}, T\right) \times B_{N} \\
v_{i}\left(\tau_{0}, \cdot\right)=v_{i}^{0} \geq 0
\end{array}\right.
$$

and such that

$$
\lim _{t \rightarrow T^{-}}\left\|v_{1}(t)\right\|_{L^{m}\left(B_{N}\right)}=\lim _{t \rightarrow T^{-}}\left\|v_{2}(t)\right\|_{L^{m}\left(B_{N}\right)}=+\infty
$$

4 for $m \geq N(d-1) / 2$.
Here the conditions $\left(\mathbf{P}^{\prime}\right)+\left(\mathbf{M}_{\lambda}^{\prime}\right)+\left(\mathbf{G}^{\prime}\right)$ are the same as $(\mathbf{P})+\left(\mathbf{M}_{\lambda}\right)+(\mathbf{G})$, but with a $(t, x)$-dependence, namely,

$$
\begin{align*}
& \forall(t, x) \in\left[\tau_{0}, T\right] \times \overline{B_{N}}, \forall\left(s_{1}, s_{2}\right) \in[0,+\infty)^{2}, \\
& \left(\mathbf{P}^{\prime}\right) g_{1}\left(t, x, 0, s_{2}\right) \geq 0, g_{2}\left(t, x, s_{1}, 0\right) \geq 0, \\
& \left(\mathbf{M}_{\lambda}^{\prime}\right) g_{1}\left(t, x, s_{1}, s_{2}\right)+\lambda g_{2}\left(t, x, s_{1}, s_{2}\right) \leq 0, \forall \lambda \text { close to } 1,  \tag{8}\\
& \left(\mathbf{G}^{\prime}\right)\left|g_{1}\left(t, x, s_{1}, s_{2}\right)\right|+\left|g_{2}\left(t, x, s_{1}, s_{2}\right)\right| \leq C_{0}+C_{1}\left(s_{1}+s_{2}\right)^{d} .
\end{align*}
$$

3. Steps of the proof of Theorem 2.1. The main idea in the proof of Theorem 2.1 is similar to the approach in [11], [12]. It consists in working with functions $u_{1}, u_{2}$ of the form

$$
\begin{equation*}
u_{i}(t, x)=\frac{a_{i}(T-t)+b_{i} r^{2}}{\left[(T-t)+r^{2}\right]^{\gamma}}, r=|x|, i=1,2, \tag{9}
\end{equation*}
$$

where $a_{1}, a_{2}, b_{1}, b_{2}$ will be well-chosen in $(0,+\infty)$, as well as $\gamma \in(1,2)$ and the dimension $N$ (large enough), in such a way that there exist $d_{1}, d_{2} \in(0,+\infty)$ and two functions $f_{1}, f_{2}$ as described in Theorem 2.1 for which

$$
\begin{equation*}
\partial_{t} u_{i}-d_{i} \Delta u_{i}=f_{i}\left(u_{1}, u_{2}\right), \quad \text { in }(0, T) \times B_{N}, \quad i=1,2 . \tag{10}
\end{equation*}
$$

The values of $u_{1}, u_{2}$ on $\partial B_{N}$ are obviously $C^{\infty}$-functions. Moreover the $L^{\infty}\left(B_{N}\right)$-norm of $u_{1}(t), u_{2}(t)$ blows up as $t \rightarrow T^{-}$since $\gamma>1$. More precisely, we directly check (see Lemma 3.1) that

$$
\begin{equation*}
\lim _{t \rightarrow T^{-}}\left\|u_{i}(t)\right\|_{L^{m}\left(B_{N}\right)}=+\infty, \quad \forall m \geq \frac{N}{2(\gamma-1)}, \quad i=1,2 \tag{11}
\end{equation*}
$$

Thus Theorem 2.1 will be proved if we can achieve (10) with regular functions $f_{1}, f_{2}$ satisfying $(\mathbf{P}),(\mathbf{M}),(\mathbf{G})$ with $d=2+\eta \geq \gamma /(\gamma-1)$. It will be a consequence of the three following lemmas.

First note that $u_{i}, i=1,2$, is of the following form:

$$
\begin{equation*}
u_{i}(t, x)=(T-t)^{1-\gamma} \rho_{i}\left(\frac{r^{2}}{T-t}\right), \rho_{i}(\sigma)=\frac{a_{i}+b_{i} \sigma}{(1+\sigma)^{\gamma}}, i=1,2 . \tag{12}
\end{equation*}
$$

We have the following technical lemma.
Lemma 3.1. Let $u_{1}$, $u_{2}$ be given by (9). For $i=1$, 2 , we have in $(0, T) \times B_{N}$

$$
\begin{equation*}
\partial_{t} u_{i}-d_{i} \Delta u_{i}=(T-t)^{-\gamma}(1+\sigma)^{-(\gamma+2)}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right], \quad \sigma=r^{2} /(T-t) \tag{13}
\end{equation*}
$$

where $A_{i}, B_{i}, C_{i}$ are real numbers given by

$$
\left\{\begin{array}{l}
A_{i}=(\gamma-1) a_{i}-2 d_{i} N\left(b_{i}-\gamma a_{i}\right)  \tag{14}\\
B_{i}=(\gamma-1)\left[2 a_{i}+b_{i}\left(1+2 d_{i} N+4 d_{i}\right)\right]+\left(b_{i}-\gamma a_{i}\right)\left[4 d_{i}(\gamma+1)-2 d_{i} N+1\right] \\
C_{i}=(\gamma-1)\left[a_{i}+b_{i}\left(1+2 d_{i} N-4 \gamma d_{i}\right)\right]+b_{i}-\gamma a_{i}
\end{array}\right.
$$

As a consequence

$$
\begin{equation*}
\partial_{t} u_{1}-d_{1} \Delta u_{1}+\partial_{t} u_{2}-d_{2} \Delta u_{2}=(T-t)^{-\gamma}(1+\sigma)^{-(\gamma+2)}\left[A+B \sigma+C \sigma^{2}\right] \tag{15}
\end{equation*}
$$

Notation. Writing a function $D: \mathbb{R} \mapsto \mathbb{R}$ as $D(N)=O\left(N^{-\alpha}\right), \alpha \in \mathbb{R}$, means that

$$
\limsup _{N \rightarrow+\infty}|D(N)| N^{\alpha}<+\infty
$$

Lemma 3.3. In the definition (9) of $u_{1}, u_{2}$, let us choose

$$
\left\{\begin{array}{l}
a_{1}=1 / N^{2}, a_{2}=2 N-N^{-2}, b_{1}=b_{2}=1 / \sqrt{N}  \tag{25}\\
d_{1}=\sqrt{N}, d_{2}=1 / N^{3}, \gamma=2-\theta / N, \theta \in(4 / 3,+\infty) \cap \mathbb{Q}
\end{array}\right.
$$

24

$$
\left\{\begin{array}{l}
A=-2 \theta+O\left(N^{-\frac{1}{2}}\right)<0, \quad C=-(8+2 \theta)+O\left(N^{-\frac{1}{2}}\right)<0,  \tag{26}\\
B^{2}-4 A C=64(4-3 \theta)+O\left(N^{-\frac{1}{2}}\right)<0 \\
A+B \sigma+C \sigma^{2} \leq A-B^{2} / 4 C=\frac{8(4-3 \theta)}{4+\theta}+O\left(N^{-\frac{1}{2}}\right)<0, \forall \sigma \in[0,+\infty)
\end{array}\right.
$$

Moreover (18) is satisfied and the functions $P_{1}, P_{2}$ can be chosen in Lemma 3.2 so that they satisfy (23) and even

$$
\begin{equation*}
P_{1}\left(s_{1}, s_{2}\right)+\lambda P_{2}\left(s_{1}, s_{2}\right) \leq 0, \quad \forall\left(s_{1}, s_{2}\right) \in(0,+\infty)^{2} \tag{27}
\end{equation*}
$$

for $\lambda$ close to 1 .
4. Proof of Theorem 2.1. Theorem 2.1 is a consequence of Lemmas 3.1, 3.2, 3.3. We will prove these lemmas below. But let us first show how these lemmas imply Theorem 2.1.

We consider the functions $u_{1}, u_{2}$ as defined in (9) with the choice of parameters $a_{1}, a_{2}, b_{1}, b_{2}, d_{1}, d_{2}, \gamma$ given in Lemma 3.3 with $\theta>4 / 3$ and $N$ large enough so that

$$
\begin{equation*}
\gamma^{\prime}=\frac{2 N-\theta}{N-\theta} \leq d=2+\eta \tag{28}
\end{equation*}
$$

We also impose $\theta \in \mathbb{Q}$. Consequently $\gamma \in \mathbb{Q}$ and also $\gamma \in(1,2)$ for $N>\theta$. For this choice of the parameters, $a_{2} b_{1}-a_{1} b_{2}=2 N^{1 / 2}+O\left(N^{-5 / 2}\right) \neq 0$ for $N$ large.

Thus by Lemma 3.2, there exist two $C^{\infty}$-functions $P_{1}, P_{2}:(0,+\infty)^{2} \rightarrow \mathbb{R}$ such that

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=P_{i}\left(u_{1}, u_{2}\right), \quad \text { in }(0, T) \times B_{N}
$$

$\left|P_{i}\left(s_{1}, s_{2}\right)\right| \leq K\left(s_{1}+s_{2}\right)^{\gamma^{\prime}}, \forall\left(s_{1}, s_{2}\right) \in(0,+\infty)^{2}, i=1,2$, for some $K \in(0+\infty)$.
By Lemma 3.3, we can even choose $P_{1}, P_{2}$ so that (27) holds for $\lambda$ close to 1 .
Note that there exists $m_{N} \in(0,+\infty)$ such that

$$
u_{i}(t, x) \geq m_{N}, \quad \forall(t, x) \in[0, T] \times \bar{B}_{N}, \quad i=1,2
$$

Indeed, we may choose $m_{N}=(T+1)^{1-\gamma} \min \left\{a_{1}, a_{2}, b_{1}, b_{2}\right\}$ as it can be seen by writing

$$
u_{i}(t, x)=\left(T-t+r^{2}\right)^{1-\gamma} \frac{a_{i}+b_{i} \sigma}{1+\sigma}
$$

Let us introduce $\varphi \in C^{\infty}([0,+\infty),[0,+\infty))$ such that

$$
0 \leq \varphi \leq 1, \varphi \equiv 0 \text { on }\left[0, m_{N} / 2\right], \varphi \equiv 1 \text { on }\left[m_{N},+\infty\right)
$$

Let us denote $f_{i}\left(s_{1}, s_{2}\right):=\varphi\left(s_{1}\right) \varphi\left(s_{2}\right) P_{i}\left(s_{1}, s_{2}\right), i=1,2$ on $[0,+\infty)^{2}$. Then

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=f_{i}\left(u_{1}, u_{2}\right), \text { in }(0, T) \times B_{N}
$$

and $f_{1}, f_{2}$ are exactly as described in Theorem 2.1. More precisely, and for future reference, we even have

$$
\begin{equation*}
f_{i}\left(s_{1}, s_{2}\right)=0, \forall\left(s_{1}, s_{2}\right) \in[0,+\infty)^{2} \backslash\left[m_{N} / 2,+\infty\right)^{2} \tag{29}
\end{equation*}
$$

Next we have $u_{i}^{0} \in C^{\infty}\left(\bar{B}_{N}\right)$ and $\alpha_{i} \in C^{\infty}([0, T])$ for $i=1,2$ since

$$
\left\{\begin{array}{l}
u_{i}^{0}(x)=\left(T+r^{2}\right)^{-\gamma}\left(a_{i} T+b_{i} r^{2}\right) \\
\alpha_{i}(t)=(T-t+1)^{-\gamma}\left(a_{i}(T-t)+b_{i}\right), i=1,2
\end{array}\right.
$$

Now, we know by Lemma 3.1, that the $L^{m}\left(B_{N}\right)$-norm of $u_{i}(t), i=1,2$ blows up if $m \geq N[2(\gamma-1)]^{-1}$. Here $\gamma$ is chosen $[$ see $(28)]$ so that $\gamma^{\prime} \leq d$, that is $(\gamma-1)^{-1} \leq d-1$. If $m \geq N(d-1) / 2$, then $m \geq N[2(\gamma-1)]^{-1}$ and the $L^{m}\left(B_{N}\right)$-norm of $u_{i}(t)$ blows up. This ends the proof of the main Theorem 2.1, assuming Lemmas 3.1, 3.2, 3.3.

Let us now prove these three lemmas.
Proof of Lemma 3.1. This result is derived by technical but easy computations. We may use $(\operatorname{see}(12)) u_{i}=(T-t)^{1-\gamma} \rho_{i}(\sigma), \sigma=r^{2} /(T-t)$ so that

$$
\begin{gathered}
\partial_{t} u_{i}=(\gamma-1)(T-t)^{-\gamma} \rho_{i}(\sigma)+\frac{r^{2}}{(T-t)^{2}}(T-t)^{1-\gamma} \rho_{i}^{\prime}(\sigma) \\
\partial_{r} u_{i}=(T-t)^{1-\gamma}\left(\frac{2 r}{T-t}\right) \rho_{i}^{\prime}\left(\frac{r^{2}}{T-t}\right)
\end{gathered}
$$

$$
\begin{gathered}
\partial_{r r} u_{i}=(T-t)^{1-\gamma}\left[\frac{4 r^{2}}{(T-t)^{2}} \rho_{i}^{\prime \prime}\left(\frac{r^{2}}{T-t}\right)+\frac{2}{T-t} \rho_{i}^{\prime}\left(\frac{r^{2}}{T-t}\right)\right] \\
\frac{N-1}{r} \partial_{r} u_{i}+\partial_{r r} u_{i}=(T-t)^{-\gamma}\left[2 N \rho_{i}^{\prime}(\sigma)+4 \sigma \rho_{i}^{\prime \prime}(\sigma)\right] . \\
\partial_{t} u_{i}-d_{i} \Delta u_{i}=(T-t)^{-\gamma}\left[(\gamma-1) \rho_{i}(\sigma)+\left(\sigma-2 d_{i} N\right) \rho_{i}^{\prime}(\sigma)-4 d_{i} \sigma \rho_{i}^{\prime \prime}(\sigma)\right] .
\end{gathered}
$$

With the choice of $\rho_{i}(\sigma)=\left(a_{i}+b_{i} \sigma\right) /(1+\sigma)^{\gamma}$, we have

$$
\begin{gathered}
\rho_{i}^{\prime}(\sigma)=\frac{b_{i}-\gamma a_{i}+b_{i}(1-\gamma) \sigma}{(1+\sigma)^{\gamma+1}}, \\
\rho_{i}^{\prime \prime}(\sigma)=\frac{b_{i}(1-\gamma)}{(1+\sigma)^{\gamma+1}}-\frac{(\gamma+1)\left[b_{i}-\gamma a_{i}+b_{i}(1-\gamma) \sigma\right]}{(1+\sigma)^{\gamma+2}} .
\end{gathered}
$$

where $A_{i}, B_{i}, C_{i}$ are given as stated in Lemma 3.1.
By summing for $i=1,2$ these two expressions, we obtain the claims (15)-(16)-(17) of Lemma 3.1. It is important to notice that this sum can be written only in terms of the four parameters $a, b, E, G$.

Finally we compute the $L^{m}\left(B_{N}\right)$-norm of $u_{i}(t)$. For some $C_{N} \in(0,+\infty)$

$$
\int_{B_{N}} u_{i}^{m}(t, x) d x=C_{N}(T-t)^{m(1-\gamma)} \int_{0}^{1} r^{N-1} \rho_{i}^{m}\left(r^{2} /(T-t)\right) d r
$$

By setting $\sigma=r^{2} /(T-t) \Leftrightarrow r=\sqrt{T-t} \sqrt{\sigma}$, we have

$$
\int_{B_{N}} u_{i}^{m}(t, x) d x=\frac{C_{N}}{2}(T-t)^{\frac{N}{2}+m(1-\gamma)} \int_{0}^{(T-t)^{-1}} \sigma^{\frac{N}{2}-1} \rho_{i}^{m}(\sigma) d \sigma
$$

As $\sigma^{\frac{N}{2}-1} \rho_{i}^{m}(\sigma)$ is equivalent to $\sigma^{\frac{N}{2}-1+m(1-\gamma)}$, up to a postive constant, as $\sigma \rightarrow+\infty$, we have

$$
\frac{N}{2}<m(\gamma-1) \Rightarrow \int_{0}^{+\infty} \sigma^{\frac{N}{2}-1} \rho_{i}^{m}(\sigma) d \sigma<+\infty
$$

As a consequence, in this case $\int_{B_{N}} u_{i}^{m}(t, x) d x$ behaves like $(T-t)^{\frac{N}{2}+(1-\gamma) m}$, up to a constant, as $t \rightarrow T^{-}$and therefore tends to $+\infty$. The conclusion is the same in the equality case $N=2 m(\gamma-1)$ since then $\int_{B_{N}} u^{m}(t, x) d x$ behaves like $|\log (T-t)|$ as $t \rightarrow T^{-}$.

Proof of Lemma 3.2. We are looking for two functions $P_{i}(\cdot, \cdot), i=1,2$ such that

$$
\begin{equation*}
\partial_{t} u_{i}-d_{i} \Delta u_{i}=P_{i}\left(u_{1}, u_{2}\right)=P_{i}\left((T-t)^{1-\gamma} \rho_{1}(\sigma),(T-t)^{1-\gamma} \rho_{2}(\sigma)\right) \tag{31}
\end{equation*}
$$

According to (13) in Lemma 3.1, and by the expected $\gamma^{\prime}$-homogeneity of $P_{i}$, this is equivalent to

$$
\begin{align*}
& (1+\sigma)^{-(\gamma+2)}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]=(1+\sigma)^{\gamma^{2} /(1-\gamma)} P_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right) \\
\Leftrightarrow & (1+\sigma)^{\frac{2-\gamma}{\gamma-1}}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]=P_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right), \forall \sigma \in[0,+\infty) \tag{32}
\end{align*}
$$

Let us write $\gamma^{\prime}=p / q \in(2,+\infty) \cap \mathbb{Q}$ with $p, q$ two coprime positive integers. Then

$$
\gamma=p /(p-q), \quad(2-\gamma) /(\gamma-1)=(p-2 q) / q>0
$$

2 Thus (32) may be rewritten

$$
\begin{equation*}
(1+\sigma)^{\frac{p-2 q}{q}}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]=P_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right) \tag{33}
\end{equation*}
$$

13
which also implies

$$
\begin{equation*}
(1+\sigma)^{p-2 q}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]^{q}=\left\{P_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right)\right\}^{q} \tag{34}
\end{equation*}
$$

Here $p>2 q$ and the left-hand side is a polynomial of degree at most $p$ in $\sigma$. It is easy to see that, since $a_{1} b_{2}-a_{2} b_{1} \neq 0$, the family of $p+1$ polynomials

$$
\left\{\left(a_{1}+b_{1} \sigma\right)^{k}\left(a_{2}+b_{2} \sigma\right)^{p-k}, k=0, \ldots, p\right\}
$$

is a linear basis of the linear space of polynomials of degree at most p. Indeed, let us analyze the relation

$$
\sum_{k=0}^{p} \mu_{k}\left(a_{1}+b_{1} \sigma\right)^{k}\left(a_{2}+b_{2} \sigma\right)^{p-k} \equiv 0, \mu_{k} \in \mathbb{R}, \forall k=0, \ldots, p
$$

Assume by contradiction that there exists $r \in\{0, \ldots, p\}$ such that

$$
\mu_{k}=0, \quad \forall k=0, \ldots, r-1, \mu_{r} \neq 0
$$

Then after dividing by $\left(a_{1}+b_{1} \sigma\right)^{r}$, the above relation may be rewritten

$$
\sum_{k=r}^{p} \mu_{k}\left(a_{1}+b_{1} \sigma\right)^{k-r}\left(a_{2}+b_{2} \sigma\right)^{p-k}=0
$$

If $r=p$, this obviously implies $\mu_{r}=0$ whence a contradiction. If $r<p$, choosing $\sigma=$ $-a_{1} / b_{1}$ leads to $\mu_{r}\left(a_{2}-b_{2} a_{1} / b_{1}\right)^{p-r}=0$, whence again $\mu_{r}=0$ thanks to the assumption (18) on $a_{i}, b_{i}, i=1,2$, and this is a contradiction.

We are now going to define the convenient function $P_{i}$. Going back to the polynomial of degree at most $p$ involved in (34), we can claim that there exist $\lambda_{k}^{i} \in \mathbb{R}, k=0, \ldots, p$ such that for all $\sigma \in \mathbb{R}$

$$
\begin{equation*}
(1+\sigma)^{p-2 q}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]^{q}=\sum_{k=0}^{p} \lambda_{k}^{i}\left(a_{1}+b_{1} \sigma\right)^{k}\left(a_{2}+b_{2} \sigma\right)^{p-k} \tag{35}
\end{equation*}
$$

7 Assume first that $q$ is odd. Let us define

$$
\begin{equation*}
\widetilde{P}_{i}\left(s_{1}, s_{2}\right):=\left\{\sum_{k=0}^{p} \lambda_{k} s_{1}^{k} s_{2}^{p-k}\right\}^{1 / q}, \quad\left(s_{1}, s_{2}\right) \in[0,+\infty)^{2} \tag{36}
\end{equation*}
$$

Note that $\widetilde{P}_{i}$ is homogeneous of degree $p / q$ and

$$
\begin{equation*}
(1+\sigma)^{p-2 q}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]=\widetilde{P}_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right) \tag{37}
\end{equation*}
$$

9
According to (30), this implies [see also (31), (32)]

$$
\begin{equation*}
\partial_{t} u_{i}-d_{i} \Delta u_{i}=\tilde{P}_{i}\left(u_{1}, u_{2}\right) \text { in } Q_{T} \tag{38}
\end{equation*}
$$

Moreover, we directly see on the definition (36) of $\widetilde{P}_{i}$ that for $\left(s_{1}, s_{2}\right) \in[0,+\infty)$

$$
\left|\widetilde{P}_{i}\left(s_{1}, s_{2}\right)\right| \leq\left\{\sum_{k=0}^{p}\left|\lambda_{k}\right|\right\}^{1 / q}\left(s_{1}+s_{2}\right)^{p / q}=C\left(s_{1}+s_{2}\right)^{\gamma^{\prime}} .
$$

Now, for the regularity of $\widetilde{P}_{i}$, it is obvious from its definition (36) that $\widetilde{P}_{i}$ is $C^{\infty}$ around each point $\left(s_{1}, s_{2}\right)$ for which $\widetilde{P}_{i}\left(s_{1}, s_{2}\right) \neq 0$. We are going to modify it into a new function $P_{i}$ which will be in $C^{\infty}\left((0,+\infty)^{2}\right)$ with the same main properties. For this, let us assume without loss of generality in (18) that $a_{1} b_{2}-a_{2} b_{1}>0$ or $b_{2} / b_{1}>a_{2} / a_{1}$ and let us consider the function

$$
s \in(0,+\infty) \rightarrow X:=\frac{a_{2}+b_{2} s}{a_{1}+b_{1} s} \in I:=\left(\frac{a_{2}}{a_{1}}, \frac{b_{2}}{b_{1}}\right)
$$

It is $C^{\infty}$ with derivative $\left(a_{1} b_{2}-a_{2} b_{1}\right) /\left(a_{1}+b_{1} s\right)^{2}>0$ on $(0,+\infty)$. Thus its inverse $\psi: I \mapsto(0,+\infty)$ is also $C^{\infty}$.

From (36), we may write

$$
\widetilde{P}_{i}\left(s_{1}, s_{2}\right)=s_{1}^{p / q} \widetilde{P}_{i}\left(1, s_{2} / s_{1}\right),
$$

12 and thanks to the relation (37) and the $p / q$ homogeneity of $\widetilde{P}_{i}$, we have

$$
\begin{equation*}
\widetilde{P}_{i}(1, X)=\left[a_{1}+b_{1} \psi(X)\right]^{-p / q}[1+\psi(X)]^{(p-2 q) / q}\left[A_{i}+B_{i} \psi(X)+C_{i} \psi(X)^{2}\right] \tag{39}
\end{equation*}
$$

Whence the $C^{\infty}$-property of $\left[X \in I \mapsto \widetilde{P}_{i}(1, X)\right]$ and therefore of $\widetilde{P}_{i}\left(s_{1}, s_{2}\right)$ on $\mathcal{I}$ where we set

$$
\mathcal{I}:=\left\{\left(s_{1}, s_{2}\right) \in(0,+\infty)^{2} ; s_{2} / s_{1} \in I\right\}
$$

Actually, $X \in \mathbb{R} \mapsto \widetilde{P}_{i}(1, X)$ is continuous and does not vanish in a neighborhood of the extremities of $I$, namely

$$
\begin{gathered}
X=a_{2} / a_{1} \Leftrightarrow \psi(X)=0, \widetilde{P}_{i}(1, X)=A_{i} a_{1}^{-p / q} \\
X=b_{2} / b_{1} \Leftrightarrow \psi(X)=+\infty, \widetilde{P}_{i}(1, X)=C_{i} b_{1}^{-p / q}
\end{gathered}
$$

1 Thus, as noticed before thanks to the definition (36), $X \mapsto \widetilde{P}_{i}(1, X)$ is $C^{\infty}$ even on some open interval $J$ containing $\bar{I}$. We now introduce a function $\chi \in C^{\infty}([0,+\infty))$ such that:

$$
\begin{equation*}
0 \leq \chi \leq 1, \chi \equiv 1 \text { on } I, \chi \equiv 0 \text { on }[0,+\infty) \backslash J \tag{40}
\end{equation*}
$$

And we set

$$
P_{i}\left(s_{1}, s_{2}\right):=s_{1}^{p / q} \chi\left(s_{2} / s_{1}\right) \widetilde{P}_{i}\left(1, s_{2} / s_{1}\right)=\chi\left(s_{2} / s_{1}\right) \widetilde{P}_{i}\left(s_{1}, s_{2}\right) .
$$

Thus $P_{i}$ is $C^{\infty}$ on all of $(0,+\infty)^{2}$ and coincides with $\widetilde{P}_{i}$ on $\mathcal{I}$. It follows that we may replace $\widetilde{P}_{i}$ by $P_{i}$ in the relations (37) and (38). This ends the proof of Lemma 3.2 when $q$ is odd.

Assume now that $q$ is even. We start again from the relation (35) from which we deduce

$$
\begin{equation*}
(1+\sigma)^{(p-2 q) / q}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right]=\widetilde{P}_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right) \operatorname{sign}\left\{A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right\} \tag{41}
\end{equation*}
$$

8 where we again set

$$
\begin{equation*}
\widetilde{P}_{i}\left(s_{1}, s_{2}\right):=\left\{\sum_{k=0}^{p} \lambda_{k}^{i} s_{1}^{k} s_{2}^{p-k}\right\}^{1 / q}=s_{1}^{p / q} \widetilde{P}_{i}\left(1, s_{2} / s_{1}\right), \forall\left(s_{1}, s_{2}\right) \in(0,+\infty)^{2}, \tag{42}
\end{equation*}
$$

and where

$$
\operatorname{sign}(r)=1, \forall r \in(0,+\infty), \operatorname{sign}(0)=0, \operatorname{sign}(r)=-1, \forall r \in(-\infty, 0)
$$

9 We deduce from (41) and the $p / q$ homogeneity of $\widetilde{P}_{i}$ that, for all $X \in I$,

$$
\left\{\begin{array}{l}
\widetilde{P}_{i}(1, X) \operatorname{sign}\left(A_{i}+B_{i} \psi(X)+C_{i} \psi(X)^{2}\right)  \tag{43}\\
=\left(a_{1}+b_{1} \psi(X)\right)^{-p / q}(1+\psi(X))^{(p-2 q) / q}\left[A_{i}+B_{i} \psi(X)+C_{i} \psi(X)^{2}\right]
\end{array}\right.
$$

and this function of $X$ is $C^{\infty}$ on $I$. Since $\widetilde{P}_{i}(1, X)$ does not vanish near the extremities of the interval $I$, it is also locally $C^{\infty}$ there. Moreover $\operatorname{sign}\left(A_{i}+B_{i} \psi(X)+C_{i} \psi(X)^{2}\right)$ is constant as $X$ tends to these extremities from inside $I$. It follows that the function of $X$ appearing in (43) may be extended in a $C^{\infty}$ - way to an open interval $J$ containing $\bar{I}$. Using the same function $\chi$ as above (see (40)), by setting

$$
P_{i}\left(s_{1}, s_{2}\right):=s_{1}^{p / q} \chi\left(s_{2} / s_{1}\right) \widetilde{P}_{i}\left(1, s_{2} / s_{1}\right) \operatorname{sign}\left(A_{i}+B_{i} \psi\left(s_{2} / s_{1}\right)+C_{i} \psi^{2}\left(s_{2} / s_{1}\right)\right)
$$

we define a $C^{\infty}$-function on $(0,+\infty)^{2}$ which is $p / q$-homogeneous and satisfies

$$
P_{i}\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right)=(1+\sigma)^{(p-2 q) / q}\left[A_{i}+B_{i} \sigma+C_{i} \sigma^{2}\right], \forall \sigma \in[0,+\infty)
$$

and therefore (see (31), (32), (33) )

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=P_{i}\left(u_{1}, u_{2}\right) \text { on } Q_{T} .
$$

This ends the proof of Lemma 3.2.
Proof of Lemma 3.3. As a consequence of the choice of the parameters in Lemma 3.3, we obtain

$$
\begin{gathered}
a=a_{1}+a_{2}=2 N, b=b_{1}+b_{2}=2 / \sqrt{N}, \\
E=d_{1} b_{1}+d_{2} b_{2}=\sqrt{N}(1 / \sqrt{N})+N^{-3}(1 / \sqrt{N})=1+O\left(N^{-7 / 2}\right), \\
G=d_{1} a_{1}+d_{2} a_{2}=\sqrt{N} N^{-2}+N^{-3}\left[2 N-\left(1 / N^{2}\right)\right]=O\left(N^{-3 / 2}\right) .
\end{gathered}
$$

It follows from the formulas (16) in Lemma 3.1 that for $N$ large

$$
\begin{gathered}
A=[1-\theta / N] 2 N-2 N\left[1+O\left(N^{-7 / 2}\right)\right]+2 N[2-\theta / N] O\left(N^{-3 / 2}\right)=-2 \theta+O\left(N^{-1 / 2}\right), \\
C=-2 N+[2-\theta / N] 2 / \sqrt{N}+2[1-\theta / N][N-4+2 \theta / N]\left[1+O\left(N^{-7 / 2}\right)\right]=-(8+2 \theta)+O\left(N^{-1 / 2}\right),
\end{gathered}
$$

$$
\begin{aligned}
& B=\left\{\begin{array}{l}
{[-\theta / N] 2 N+[2-\theta / N][2 / \sqrt{N}]+[2 N(-\theta / N)+8(2-\theta / N)]\left[1+O\left(N^{-7 / 2}\right)\right]} \\
\\
-2[2-\theta / N][2(3-\theta / N)-N]\left[O\left(N^{-3 / 2}\right)\right]=16-4 \theta+O\left(N^{-1 / 2}\right) .
\end{array}\right. \\
& B^{2}-4 A C=16[4-\theta]^{2}-8 \theta(8+2 \theta)+O\left(N^{-1 / 2}\right)=64[4-3 \theta]+O\left(N^{-1 / 2}\right) .
\end{aligned}
$$

With the choice of $\theta>4 / 3$, we obtain $A<0, C<0, B^{2}-4 A C<0$ for $N$ large as claimed in Lemma 3.3.

Since the maximum of $\sigma \in \mathbb{R} \mapsto A+B \sigma+C \sigma^{2}$ is reached for $\sigma=-B / 2 C$, we deduce

$$
\begin{equation*}
A+B \sigma+C \sigma^{2} \leq A-B^{2} / 4 C=-K_{\theta}+O\left(N^{-\frac{1}{2}}\right), \forall \sigma \in[0,+\infty) \tag{44}
\end{equation*}
$$

4 with $K_{\theta}:=8(3 \theta-4) /(4+\theta)>0$.
On the other hand, $a_{2} b_{1}-a_{1} b_{2}=2 N^{1 / 2}+O\left(N^{-5 / 2}\right) \neq 0$ and $\gamma \in(1,2) \cap \mathbb{Q}$ so that (18) is satisfied and we may apply Lemma 3.2. Using now the relation (21) in this Lemma 3.2 together with (44), we obtain that

$$
\left(P_{1}+P_{2}\right)\left(a_{1}+b_{1} \sigma, a_{2}+b_{2} \sigma\right) \leq-K_{\theta}+O\left(N^{-\frac{1}{2}}\right), \forall \sigma \in[0,+\infty)
$$

By the $p / q$-homogeneity of $P_{1}+P_{2}$, we deduce that

$$
\left(P_{1}+P_{2}\right)(1, X) \leq-K_{\theta} a_{1}^{-p / q}+O\left(N^{-\frac{1}{2}}\right), \forall X \in I
$$

where $I$ is defined as in the previous proof of Lemma 3.2. Since $I$ is bounded, so is $P_{2}(1, X)$ for $X \in I$. It follows that for $\lambda$ close to 1 and $N$ large enough, we may claim that

$$
\left(P_{1}+\lambda P_{2}\right)(1, X) \leq-\frac{K_{\theta} a_{1}^{-p / q}}{2}, \forall X \in I
$$

Up to reducing $J$ and the support of the function $\chi$ in (40) of the proof of Lemma 3.2, we may claim that

$$
\left(P_{1}+\lambda P_{2}\right)(1, X) \leq-\frac{K_{\theta} a_{1}^{-p / q}}{4}, \forall X \in J
$$

By homogeneity, this implies that

$$
\left(P_{1}+\lambda P_{2}\right)\left(s_{1}, s_{2}\right) \leq 0, \forall\left(s_{1}, s_{2}\right) \in(0,+\infty)^{2}
$$

5 This ends the proof of Lemma 3.3.
5. Proof of Theorem 2.2. We use the same functions $u_{1}, u_{2}$ as those introduced in the proof of Theorem 2.1 and defined in (9), and we choose $\gamma$ as in (28), namely such that

$$
\begin{equation*}
\gamma^{\prime} \leq d \Leftrightarrow \gamma \geq d /(d-1) \tag{45}
\end{equation*}
$$

We deduce from (30) that

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=\frac{A_{i}(T-t)^{2}+B_{i}(T-t) r^{2}+C_{i} r^{4}}{\left(T-t+r^{2}\right)^{\gamma+2}}, i=1,2 .
$$

Given $\alpha, \beta$ as indicated in the statement of Theorem 2.2, we may understand this expression as being of the form

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=c_{i}(t, x) u_{1}^{\alpha} u_{2}^{\beta},
$$

where we define

$$
\begin{gathered}
c_{i}(t, x):=\frac{A_{i}(T-t)^{2}+B_{i}(T-t) r^{2}+C_{i} r^{4}}{\left(T-t+r^{2}\right)^{\gamma+2}} u_{1}^{-\alpha} u_{2}^{-\beta}, \\
\Rightarrow \quad c_{i}(t, x)=\frac{\left(T-t+r^{2}\right)^{\gamma(d-1)-2}\left[A_{i}(T-t)^{2}+B_{i}(T-t) r^{2}+C_{i} r^{4}\right]}{\left[a_{1}(T-t)+b_{1} r^{2}\right]^{\alpha}\left[a_{2}(T-t)+b_{2} r^{2}\right]^{\beta}} .
\end{gathered}
$$

Obviously $c_{i}$ is $C^{\infty}$ on $[0, T) \times \overline{B_{N}}$. Let us check that it is in $L^{\infty}\left(Q_{T}\right)$ as follows. Let $H:=T-t+r^{2}$. We use

$$
T-t \leq H, r^{2} \leq H, a_{i}(T-t)+b_{i} r^{2} \geq \min \left\{a_{i}, b_{i}\right\} H
$$

Thus for some $K_{1}, K_{2} \in(0,+\infty)$, and using that $\gamma(d-1)-2 \geq \gamma(d-1)-d \geq 0$, we have

$$
\left|c_{i}(t, x)\right| \leq \frac{H^{\gamma(d-1)-2}\left[\left[A_{i}\left|+\left|B_{i}\right|+\left|C_{i}\right|\right] H^{2}\right.\right.}{K_{1} H^{\alpha+\beta}} \leq K_{2} H^{\gamma(d-1)-d} \leq K_{2}(T+1)^{\gamma(d-1)-d} .
$$

Finally, we see that the sign of $c_{1}+c_{2}$ is the same as the sign of $A(T-t)^{2}+B(T-t) r^{2}+C r^{4}$, that is $c_{1}+c_{2} \leq 0$ since $A<0, C<0, B^{2}-4 A C<0$ by Lemma 3.3. This ends the proof of Theorem 2.2.
6. Proof of Theorem 2.3. Let $u_{1}, u_{2}$ be the solutions obtained in Theorem 2.1. Their normal derivatives on $\partial B_{N}$ are given as follows

$$
\begin{align*}
& \partial_{r} u_{1}(t, 1)=\frac{2\left[\left(b_{1}-\gamma a_{1}\right)(T-t)-(\gamma-1) b_{1}\right]}{(T-t+1)^{\gamma+1}}, \partial_{r} u_{2}(t, 1)=\frac{2\left[\left(b_{2}-\gamma a_{2}\right)(T-t)-(\gamma-1) b_{2}\right]}{(T-t+1)^{\gamma+1}}, \\
& b_{1}=N^{-1 / 2}=b_{2}, b_{1}-\gamma a_{1}=N^{-1 / 2}-2 N^{-2}+O\left(N^{-3}\right),  \tag{46}\\
& b_{2}-\gamma a_{2}=-4 N+O(1)
\end{align*}
$$

6 Obviously for $N$ large

$$
\left\{\begin{array}{l}
\partial_{r} u_{1}(t, 1) \leq 0, \forall t \in\left[\tau_{0}, T\right], \tau_{0}:=(T-\gamma+1)^{+}(<T),  \tag{47}\\
\partial_{r} u_{2}(t, 1) \leq 0, \forall t \in[0, T] .
\end{array}\right.
$$

7 We now introduce the solutions $\beta_{i}, i=1,2$ of

$$
\left\{\begin{array}{l}
\partial_{t} \beta_{i}-d_{i} \Delta \beta_{i}=0 \text { in }\left(\tau_{0}, T\right) \times B_{N}  \tag{48}\\
\partial_{\nu} \beta_{i}=-\partial_{r} u_{i}(t, 1) \text { on }\left(\tau_{0}, T\right) \times \partial B_{N} \\
\beta_{i}\left(\tau_{0}, \cdot\right) \equiv 0
\end{array}\right.
$$

By the $C^{\infty}$-property of $\partial_{r} u_{i}(t, 1)$, these solutions are also in $C^{\infty}\left(\overline{Q_{T}}\right)$. Moreover, thanks to (47), and by maximum principle, we have $\beta_{i}(t, x) \geq 0$ on $\left[\tau_{0}, T\right) \times B_{N}$. We now introduce for $i=1,2$

$$
v_{i}(t, x):=u_{i}(t, x)+\beta_{i}(t, x) \geq 0 \text { so that } \partial_{\nu} v_{i}=\partial_{r} v_{i}(t, 1)=0 \text { on } \Sigma_{T} .
$$

And according to Theorem 2.1 and the definition of $\beta_{i}$, we have

$$
\partial_{t} v_{i}-d_{i} \Delta v_{i}=\partial_{t} u_{i}-d_{i} \Delta u_{i}=f_{i}\left(u_{1}, u_{2}\right)=g_{i}\left(t, x, v_{1}, v_{2}\right) \text { in } Q_{T}, i=1,2,
$$

where we set

$$
g_{i}\left(t, x, s_{1}, s_{2}\right):=f_{i}\left(s_{1}-\beta_{1}(t, x), s_{2}-\beta_{2}(t, x)\right) .
$$

Here the functions $f_{i}$ are extended by 0 outside $[0,+\infty)^{2}$ and with the help of (29), we see that they are $C^{\infty}$ on $\mathbb{R}^{2}$. The functions $g_{i}$ defined in this way satisfy $\left(\mathbf{P}^{\prime}\right),\left(\mathbf{M}_{\lambda}^{\prime}\right),\left(\mathbf{G}^{\prime}\right)$ since, by nonnegativity of $\beta_{i}$ :

$$
\begin{aligned}
& g_{1}\left(t, x, 0, s_{2}\right)=f_{1}\left(-\beta_{1}(t, x), s_{2}-\beta_{2}(t, x)\right)=0, \\
& g_{2}\left(t, x, s_{1}, 0\right)=f_{2}\left(s_{1}-\beta_{1}(t, x),-\beta_{2}(t, x)\right)=0, \\
& \left(g_{1}+\lambda g_{2}\right)\left(t, x, s_{1}, s_{2}\right)=\left(f_{1}+\lambda f_{2}\right)\left(s_{1}-\beta_{1}(t, x), s_{2}-\beta_{2}(t, x)\right) \leq 0[\text { using }(6)], \\
& \left|g_{i}\left(t, x, s_{1}, s_{2}\right)\right|=\left|f_{i}\left(s_{1}-\beta_{1}(t, x), s_{2}-\beta_{2}(t, x)\right)\right| \\
& \leq C_{0}+C_{1}\left[\left(s_{1}-\beta_{1}(t, x)\right)^{+}+\left(s_{2}-\beta_{2}(t, x)^{+}\right]^{d}\right. \\
& \Rightarrow\left|g_{i}\left(t, x, s_{1}, s_{2}\right)\right| \leq C_{0}^{\prime}+C_{1}^{\prime}\left|s_{1}+s_{2}\right|^{d},
\end{aligned}
$$

where $C_{0}^{\prime}, C_{1}^{\prime}$ depend on $C_{0}, C_{1}, d,\left\|\beta_{i}\right\|_{L^{\infty}\left(Q_{T}\right)}$. Finally, the $L^{m}\left(B_{N}\right)$-norm of $v_{i}(t)$ blows up as $t \rightarrow T^{-}$together with the $L^{m}\left(B_{N}\right)$-norm of $u_{i}(t)$. This ends the proof of Theorem 2.3.
7. About global weak solutions. Let us prove here that the solutions $u_{1}, u_{2}$ obtained in Theorem 2.1 may be extended to global weak solutions of the same system. For this we will use the approach of the previous section and, as a first step, we extend the functions $v_{1}, v_{2}$ to weak solutions on $[0,+\infty)$. This will provide extensions of $u_{1}, u_{2}$ as well.

Let us start by extending the functions $\beta_{i}$ as the solutions on $[0,+\infty)$ of

$$
\left\{\begin{array}{l}
\partial_{t} \beta_{i}-d_{i} \Delta \beta_{i}=0 \text { in }\left(\tau_{0},+\infty\right) \times B_{N}  \tag{49}\\
\partial_{\nu} \beta_{i}=\left\{\begin{array}{l}
-\partial_{r} u_{i}(t, 1) \text { on }\left(\tau_{0}, T\right) \times \partial B_{N} \\
-\partial_{r} u_{i}(T, 1) \text { on }(T,+\infty) \times \partial B_{N} \\
\beta_{i}\left(\tau_{0}, \cdot\right) \equiv 0
\end{array}\right.
\end{array}\right.
$$

These extensions are $C^{\infty}$ on $\left[\tau_{0},+\infty\right) \times \overline{B_{N}}$. We also extend the nonlinear function $g_{i}$ as follows for $(t, x) \in\left(\tau_{0},+\infty\right) \times \overline{B_{N}}$

$$
g_{i}\left(t, x, s_{1}, s_{2}\right):=f_{i}\left(s_{1}-\beta_{1}(t, x), s_{2}-\beta_{2}(t, x)\right), \forall\left(s_{1}, s_{2}\right) \in[0,+\infty)^{2} .
$$

1 They satisfy the same properties $\left(\mathbf{P}^{\prime}\right),\left(\mathbf{M}_{\lambda}^{\prime}\right),\left(\mathbf{G}^{\prime}\right)$ on $\left[\tau_{0},+\infty\right)$. We now apply Theorem 5.9 in [10] to the system

$$
\left\{\begin{array}{l}
\partial_{t} v_{1}-d_{1} \Delta v_{1}=g_{1}\left(t, x, v_{1}, v_{2}\right) \text { in }\left(\tau_{0},+\infty\right) \times B_{N},  \tag{50}\\
\partial_{t} v_{2}-d_{2} \Delta v_{2}=g_{2}\left(t, x, v_{1}, v_{2}\right) \text { in }\left(\tau_{0},+\infty\right) \times B_{N}, \\
\partial_{\nu} v_{1}=0=\partial_{\nu} v_{2} \text { on }\left(\tau_{0},+\infty\right) \times B_{N}, \\
v_{i}\left(\tau_{0}, \cdot\right)=v_{i}^{0}=u_{i}\left(\tau_{0}, \cdot\right)+\beta\left(\tau_{0}, \cdot\right) .
\end{array}\right.
$$

Indeed, thanks to the property $\left(\mathbf{M}_{\lambda}^{\prime}\right)$, any good approximation of this system with bounded nonlinearities $g_{i}^{n}$ in place of $g_{i}$ and the same regular data at $t=\tau_{0}$, will provide approximate regular solutions $v_{i}^{n}$ such that, for all $S \in\left[\tau_{0},+\infty\right)$

$$
\sup _{n \geq 1} \int_{\left(\tau_{0}, S\right) \times B_{N}}\left|g_{i}^{n}\left(t, x, v_{1}^{n}(t, x), v_{2}^{n}(t, x)\right)\right|<+\infty, i=1,2,
$$

We may for instance choose

$$
g_{i}^{n}\left(t, x, s_{1}, s_{2}\right):=\frac{g_{i}\left(t, x, s_{1}, s_{2}\right)}{1+n^{-1}\left[\left(\left|g_{1}\right|+\left|g_{2}\right|\right)\left(t, x, s_{1}, s_{2}\right)\right]},
$$

so that $\left|g_{i}^{n}\right| \leq n$. These approximate functions $g_{i}^{n}$ satisfy also $\left(\mathbf{P}^{\prime}\right),\left(\mathbf{M}_{\lambda}^{\prime}\right)$ and it is easy to prove that these properties imply the above $L^{1}\left(\left(\tau_{0}, S\right) \times B_{N}\right)$-bounds (see e.g. Proposition 5.1 in [10]). As a consequence, and as proved in [10], these approximate solutions $v_{i}^{n}$ converge to a weak solution of (50) which means that

$$
v_{i}(t)=S_{i}\left(t-\tau_{0}\right) v_{i}^{0}+\int_{\tau_{0}}^{t} S_{i}(t-s) g_{i}\left(s, \cdot, v_{1}(s, \cdot), v_{2}(s, \cdot)\right) d s, i=1,2
$$

where $S_{i}(t)$ denotes the semigroup generated by the operator $-d_{i} \Delta$ on $B_{N}$ with homogeneous Neumann boundary conditions.

These weak solution coincides on the interval $\left[\tau_{0}, T\right)$ with the classical solution of (7) as found in Theorem 2.3. Indeed, since $v_{i}^{0} \in L^{\infty}\left(B_{N}\right)$, it is classical to prove that the approximate solution $\left(v_{1}^{n}, v_{2}^{n}\right)$ stays uniformly bounded, independently of $n$, on some interval $\left[\tau_{0}, \tau_{1}\right] \subset\left[\tau_{0}, T\right]$. To see it, we might for instance use the inequalities

$$
\partial_{t} v_{i}^{n}-d_{i} \Delta v_{i}^{n} \leq\left|g_{i}\left(t, x, v_{1}^{n}, v_{2}^{n}\right)\right| \leq C_{0}^{\prime}+C_{1}^{\prime}\left|v_{1}^{n}+v_{2}^{n}\right|^{d} .
$$

5 Thus the limit $\left(v_{1}, v_{2}\right)$ is therefore a classical solution on $\left[\tau_{0}, \tau_{1}\right]$. By uniqueness of classical solutions for the reaction-diffusion system (7), both solutions coincide at least on $\left[\tau_{0}, \tau_{1}\right]$, and subsequently on the whole interval $\left[\tau_{0}, T\right]$ by a classical continuity argument.

Finally, if we now set $u_{i}(t, x):=v_{i}(t, x)-\beta_{i}(t, x), i=1,2$ on the whole domain $\left[\tau_{0},+\infty\right) \times \overline{B_{N}}$, we find a global weak solution of the initial system (1) on $[0,+\infty)$ which extends the (classical) solution obtained on $[0, T)$ in Theorem 2.1. It remains to get convinced that $u_{i} \geq 0$ for all $t$. This can be seen by noticing that $u_{i}^{n}:=v_{i}^{n}-\beta_{i}$ is a classical solution on $\left[\tau_{0},+\infty\right)$ of

$$
\partial_{t} u_{i}^{n}-d_{i} \Delta u_{i}^{n}=g_{i}^{n}\left(t, x, v_{1}^{n}, v_{2}^{n}\right)=f_{i}^{n}\left(u_{1}^{n}, u_{2}^{n}\right),
$$

8 with boundary conditions $\partial_{\nu} u_{i}^{n}=-\partial_{r} \beta_{i} \geq 0$ with a quasi-positive nonlinearity $\left(f_{1}^{n}, f_{2}^{n}\right)$. 9 Thus $u_{i}^{n} \geq 0$ and this is preserved at the limit.
8. Some more explicit examples in small dimensions. We use here Lemma 3.1 and Lemma 3.2. All blow up examples we are going to describe here are defined as in (9) so that they blow up as $t \rightarrow T^{-}$. Moreover by Lemma 3.2, we know that we can then find $P_{1}, P_{2}$ so that (19) holds, namely

$$
\partial_{t} u_{i}-d_{i} \Delta u_{i}=P_{i}\left(u_{1}, u_{2}\right), \text { in } Q_{T}, \quad i=1,2 .
$$

1 The goal is then to choose all parameters involved in the definition (9) well enough so that (24) holds, namely

$$
\begin{equation*}
A<0, C<0, B^{2}-4 A C<0 \tag{51}
\end{equation*}
$$

where $A, B, C$ are defined in Lemma 3.1. Then, exactly as in the proof of Theorem 2.1, we prove that the previous inequalities implies that $P_{1}+P_{2} \leq 0$ and that $P_{1}, P_{2}$ can then slightly be modified into functions $f_{1}, f_{2}$ satisfying $(\mathbf{P}),(\mathbf{M})$ and (1). Moreover the growth at infinity of $f_{1}, f_{2}$ is at most $\gamma^{\prime}$.

The following technical proposition is very useful to help finding examples of blow up. Notations are as in Lemma 3.1.

Proposition 8.1. Let $a, b, E, G \in(0,+\infty)$. Then we can find $a_{i}, b_{i}, d_{i}, i=1,2 \in(0,+\infty)$ such that

$$
\begin{equation*}
a_{1}+a_{2}=a, b_{1}+b_{2}=b, b_{1} d_{1}+b_{2} d_{2}=E, a_{1} d_{1}+a_{2} d_{2}=G \tag{52}
\end{equation*}
$$

$$
\begin{equation*}
a_{1} b_{2}-a_{2} b_{1}=\frac{a E-b G}{d_{2}-d_{1}} \tag{53}
\end{equation*}
$$

Proof of Proposition 8.1. Let us choose $d_{1}, d_{2}$ such that

$$
0<d_{1}<\min \left\{E b^{-1}, G a^{-1}\right\} \leq \max \left\{E b^{-1}, G a^{-1}\right\}<d_{2}
$$

Then, given the four parameters $a, b, E, G$, a solution to (52) is given by

$$
a_{1}=\frac{a d_{2}-G}{d_{2}-d_{1}}, a_{2}=\frac{G-a d_{1}}{d_{2}-d_{1}}, b_{1}=\frac{b d_{2}-E}{d_{2}-d_{1}}, b_{2}=\frac{E-b d_{1}}{d_{2}-d_{1}}
$$

We easily check that, with this choice, the relations (52), (53) are satisfied and all parameters are positive. The relation (53) will be used to check (18).

Remark 8.1. In what follows, we give several explicit choices of parameters which lead to blow up examples according to the analysis of the previous sections. Thanks to the Proposition 8.1, it is sufficient to provide the values of $a, b, E, G$ such that (51) holds for the corresponding values of $A, B, C$.
8.1. Blow up with a cubic growth in dimension $N=3$. Here we choose with $\epsilon \in(0,1)$ small enough,

$$
\gamma=3 / 2, a=10, b=1, E=1, G=\epsilon
$$

Then, $a E-b G=10-\epsilon \neq 0$ and

$$
\begin{gathered}
A=(\gamma-1) a-6 E+6 \gamma G=5-6+9 \epsilon=-1+9 \epsilon<0 \\
C=-a+\gamma b+2(\gamma-1)(3-2 \gamma) E=-10+3 / 2+0=-17 / 2<0 \\
B=(\gamma-2) a+\gamma b+[6(\gamma-2)+8 \gamma] E-2 \gamma[2(\gamma+1)-3] G=11 / 2-6 \epsilon \\
B^{2}-4 A C=\frac{11^{2}}{4}-4 \frac{17}{2}+0(\epsilon)=-15 / 4+O(\epsilon)<0
\end{gathered}
$$

Thus (51) holds so that this provides a blowing up example with growth $\gamma^{\prime}=3$. By Lemma 3.1, for $m \geq 3$, the $L^{m}\left(B_{N}\right)$-norm of the solution blows up as $t \rightarrow T^{-}$. Note that since all inequalities are strict in this example, we could slightly modify it to obtain a blowing up example with growth $3-\sigma$ for some small $\sigma>0$.
8.2. Blow up with a $7 / 2$-growth in dimension $N=2$. An example with such a growth was already mentioned in [12]. We give here another one, easier to compute. Let us choose the following parameters with $\epsilon>0$ small enough.

$$
\gamma=7 / 5, a=4, b=\epsilon, E=1, G=\epsilon .
$$

then $a E-b G=4-\epsilon^{2} \neq 0$ and

$$
\begin{gathered}
A=\frac{2}{5} a-4 E+\frac{28}{5} G=\frac{8}{5}-4+O(\epsilon)=-\frac{12}{5}+O(\epsilon)<0, \\
C=-a+\frac{7}{5} b-\frac{16}{25} E=-4-\frac{16}{25}+O(\epsilon)=-\frac{116}{25}+O(\epsilon)<0, \\
B=-\frac{3}{5} a+\frac{7}{5} b+\frac{44}{5} E-\frac{14^{2}}{25} G=-\frac{12}{5}+\frac{44}{5}+O(\epsilon)=\frac{32}{5}+O(\epsilon) . \\
B^{2}-4 A C=\frac{32^{2}}{25}-4 \frac{12}{5} \frac{116}{25}+O(\epsilon)=\frac{1}{125}[5120-5568]+O(\epsilon)=-\frac{448}{125}+O(\epsilon)<0 .
\end{gathered}
$$

Thus this provides a blow up example with growth $\gamma^{\prime}=7 / 2$ and an $L^{m}\left(B_{N}\right)$-norm blowing up for $m \geq 5 / 2$. Again, since all inqualities are strict, we could find another example with growth $7 / 2-\sigma$ for some $\sigma>0$.

### 8.3. Blow up with a 6 -growth in dimension $N=1$. Here we choose

$$
\gamma=6 / 5, a=5, b=1 / 2, E=1, G=1 / 10 .
$$

then $a E-b G=5-1 / 20 \neq 0$ and

$$
\begin{gathered}
A=1-2+6 / 25=-19 / 25<0 \\
C=-5+3 / 5+2(1 / 5)(1-12 / 5)=-22 / 5-14 / 25=-124 / 25<0 \\
B=-4+3 / 5+8-12.17 / 250=473 / 125 \\
125^{2}\left[B^{2}-4 A C\right]=473^{2}-4.95 .620=223729-235600=-11871<0
\end{gathered}
$$

This provides a blow up example with growth $\gamma^{\prime}=6$ and with an $L^{m}\left(B_{N}\right)$-norm blowing up for $m \geq 5 / 2$. As above again, we could improve up it to a lower growth $6-\sigma$ for some $\sigma>0$.

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