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Highly Coupled and Low Frequency Vibrational Energy Harvester Using Lithium Niobate on Silicon

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Abstract. LiNbO₃ has been little studied for the piezoelectric energy harvesting applications. Although, it is a cheap piezoelectric material without lead and toxic elements, and it beneficiates of technological maturity in single crystal fabrication for optical and acoustic applications. In this letter, we propose to investigate a (YXlt)/128°/90° LiNbO₃ cut as it offers a transverse piezoelectric coupling of $k_{23} = 0.49$ which is comparable to that of commonly used PZT ceramics. A flexible beam of 65 mm length and tip mass made of LiNbO₃ thick film bonded on silicon was studied under 0.1 g sinusoidal acceleration. The beam presented an open-circuit resonance frequency of 105.9 Hz and a displacement up to 1.5 mm. In the frame of single degree of freedom lumped model with rectifying bridge, the electromechanical coupling of the device (k_m^2) , the figure of merit $k_m^2 Q$ and the normalized average power density were compared to both Pb-based and Pb-free current devices. The generated power density by our device was 965 μ W/cm²/g², which is among the highest reported values compared to both Pb- and Pb-free vibrational harvesting devices.

Keywords: LiNbO₃, lead-free, piezoelectric material, vibrational energy harvesting

Energy harvesting and ultra-low power electronic are urgently needed for the implementation of selfpowered, battery- and maintenance-free sensors [1]. The integrated solutions with energy harvester are highly demanded in the market of the Internet of Things, health monitoring and more generally 4.0 industry [2]. Vibrational energy harvesting is frequently considered as an alternative solution to photovoltaics for this purpose. The specifications of vibrational sources available in vehicles or industrial environments are low acceleration levels (< 1 g = 9.81 m/s²) and frequencies (< 500 Hz). Pbcontaining piezoelectric materials such as $PbZr_{1-x}Ti_xO_3$ (PZT) ceramics are usually considered as the best performing electro-mechanical transducers. They are implemented in form bulk ceramic, fibers [3] or thin films. Recently, it has been shown that the thickness of the piezoelectric layer can be optimized [4], where a 57 µm PZT ceramic film was fabricated in unimorph configuration by means of phosphorbronze bonding and mechanical lapping, achieving 321 μ W at 100.8 Hz for high acceleration applications (up to 3 g). Kuo et al [5] have fabricated by aerosol deposition thick PZT films (15 μ m) in bimorph configuration on stainless steel, obtaining 423 μ W at 143.5 Hz for 1.5 g acceleration level. In other works, sputtered films were exploited [6] in form of bimorph harvesters on flexible Ni foils, obtaining higher power densities, 60 μ W at 70 Hz (0.5 g). Relaxor ferroelectrics in form of single crystals such as Pb[Mg_{1/3}Nb_{2/3}]O₃ - PbTiO₃ (PMN-PT) or Pb[Zn_{1/3}Nb_{2/3}]O₃ - PbTiO₃ (PZN-PT) are also considered for harvesting applications, although they are quite brittle and expensive. A comparison of PZT, PMN-PT and PZN-PT was presented [7], and the experimental results indicated that relaxor ferroelectrics can significantly outperform PZT ceramics energy harvesters, attaining 0.43 mW at 37.5 Hz at 0.3 g. Despite the outstanding results, these devices are still exploiting Pb-based materials that represent an environmental issue, and do not respect REACH and RoHS regulation in EU.

On the other hand, Pb-free materials have been also investigated for applications at low acceleration and low excitation frequency. One of the highest performances were obtained for AlN devices [8] capable of working at 58 Hz and generating 63 μ W at 0.7 g. Ferroelectric polymer PVDF has been also considered for low frequency applications (30.7 Hz) but has a limited power output (8 μ W at 0.5 g) [9]. The potassium and sodium niobate (KNN) family have shown promising results especially for MEMS applications. For instance, it was possible to attain a resonance frequency of 126 Hz and 7.7 μ W at 1.02 g by using KNN films on metal foils, deposited by hydrothermal method [10].

Among the Pb-free materials, very well-known lithium niobate (LiNbO₃) has drawn little attention as an electro-mechanical transducer in the literature [11], due to its low piezoelectric coefficients. LiNbO₃ is a ferroelectric material with extremely high Curie temperature (1100°C) that belongs to R3c space group at room temperature and that is commercially available at low price in form of high-quality poled single crystal wafers with diameter up to 6 inches (< 25 \$/wafer). LiNbO₃ elastic, piezoelectric and dielectric constants have been extensively studied [12] and its electro-mechanical coupling can be optimized as a function of crystal orientation [13]. Following IEEE standard on piezoelectricity [14], we can define the electro-mechanical coupling of LiNbO₃ (YXlt)/128°/90° crystal cut as $k_{23} = \sqrt{\frac{d_{23}^2}{s_{33}^E \epsilon_{22} T}} = 0.49$, taking the rotated piezoelectric coefficient $d_{23} = -27$ pC/N, dielectric constant $\epsilon_{22}T/\epsilon_0 = 50.5$, and elastic compliance $s_{33}^E = 6.91$ pm²/N. Therefore, the chosen orientation of LiNbO₃ has an expected stress-driven figure of merit (FoM⁻¹= $\frac{d_{23}^2}{\epsilon_{22}^T}$) of 0.61 MJ/cm³. This value is comparable with commercially available ceramics (PZT-5A, PZT-5H) and superior than other Pb-free materials (AlN, KNN, PVDF), where a detailed comparison of piezoelectric properties figure of merit is given in a recent work [15].

The implementation of $LiNbO_3$ in form of bulk crystals for purpose of energy harvesting, shows too low capacitance and results challenging for electronic interfacing and impedance matching [16]. Instead, the

thick films obtained by wafer on wafer technology [17], highly investigated in the field of acoustics and optics [18], allowed to address the capacitance issues and to show high vibrational energy harvesting efficiency at high acceleration (3 g) and frequencies in the kHz range [15]. In this paper, we have addressed the issue of the LiNbO₃-based vibrational harvester compatibility with vibrations available in the environment by lowering down the resonance frequency to ~ 105.9 Hz of LiNbO₃/Si beams, in order to attain high power densities under very low acceleration levels (0.1 g). The harvester was described by using the single degree of freedom lumped model and has been studied under resistive load near the resonance frequency.

The cantilevered beams were fabricated in the clean room by first bonding two substrates (Fig. 1a): a 500 µm thick (100) Si wafer and a 350 µm thick single crystal (YXlt)/128°/90° LiNbO3 wafer. For this purpose, a 30 nm Cr and 170 nm Au layers were sputtered on the polished surface of both wafers. The wafers were bonded by vacuum thermocompression at room temperature. The bonding interface quality was investigated by ultrasonic non-destructive testing in transmission [17]. The Fig. 1b presents in false colors the impedance matching in bonded wafers. The good quality of bonding corresponds to red color, while blue and green colors are related to bonding defects like air gap, roughness or dust particles at the interface. The thickness of LiNbO₃ was thinned down to approximately 30 μ m by mechanical lapping and micro-polishing. The typical Total Thickness Variation (TTV) achieved was 2 µm. The top electrodes (200 nm thick Cr/Au), deposited by electron-beam evaporation, were patterned by UV lithography and lift-off. The top electrode was covering 2/3 surface of the total final length. The devices were subsequently diced in rectangular beams of 65 mm length and 10 mm wide. The Si backside was protected with tape in order to leave 10 mm on each side. The unprotected beam was etched at the central part by deep reactive ion etching (DRIE) until thickness of 360 µm. The good bonding quality after dicing was confirmed by imaging the heterostructure cross-section by a scanning electron microscope (SEM-Apreo S, Fig. 1c). The final device was then clamped in between rigid aluminum bars. The top electrode was wire-bonded to a PCB while the bottom electrode was connected to the other branch of the circuit contacting the Au-Au layer using silver paste. The photo of the sample wire-bonded to the clamping system is given in the inset of Fig. 1d.

The setup for the vibrational measurements consisted of a shaker, providing the sinusoidal excitation for the vibrational tests, coupled with an accelerometer mounted at the base of the clamping system to monitor the magnitude of the input force (Fig. 1d). Meanwhile, the displacement of the tip was tracked by a laser interferometer with 50 μ m resolution, giving a relative error of 2.5% for the measurements at resonance. Typically, the acceleration magnitude range used was 0.1 ± 0.02 g. The voltage levels were measured both with an oscilloscope (LeCroy LT344) using a 10 M Ω attenuation probe, or with a Keithley 6517B electrometer after diode rectification, and connected to a National Instruments Labview interface. The voltage response at open-circuit frequency (105.9 Hz), measured experimentally, where the harvester showed a voltage peak of 15.1 V (RMS). This result is compared to the response simulated with Finite-Element modelling (using COMSOL® software), which shows the linear behavior of the harvester reaching a voltage peak of 16 V (RMS) (Fig. 2a). The modeled peak presents a shift towards higher frequencies (106.2 Hz) with respect to the experimental one. The experimental data is in good agreement with the simulations, even though the voltage levels seemed to be slightly overestimated (by 0.9 volt). This deviation was modeled in some recent works [19], and it was related to *softening* of the silicon substrate, which appeared whenever the restoring force decreases with the displacement. Nevertheless, the bending mode of the beam was attained in the expected frequency range and at low acceleration levels, which are compatible with environmental vibrations (< 1 g).

The electrical impedance of our device was investigated by using a network analyzer (Keysight E5061B) and by applying a Butterworth Van Dyke (BVD) equivalent model. The impedance was measured from 10 Hz to 100 kHz at room temperature. The samples showed very small dielectric losses ($\tan \delta = 0.24$ %) and a clamped capacitance of $C_0 = 5.8$ nF at 2 kHz measured away from resonance, confirming the good quality of the single crystal. The electrical resonance was observed at 104.53 Hz, and it is presented along with the BVD fitted parameters in Fig. 2b.

To investigate LiNbO₃ thick film on silicon as a competitive alternative for highly coupled energy harvesters, our results were analyzed through the widely used single degree of freedom model. The harvester operating close to its resonance frequency (ω_0) can be approximated as an effective spring (*K*), damper (*C*) and mass (*M*) system (Fig. 3a). The dynamic equations connecting the electrical and mechanical sub-systems are:

$$\begin{cases} M\ddot{y} = M\ddot{x} + C\dot{x} + Kx + \alpha V\\ I = \alpha \dot{x} - C_0 \dot{V} \end{cases}, \tag{1}$$

where the term $M\ddot{y}$ represents the external force acting on the harvester, while \dot{x} is the relative displacement of the mass. The piezoelectric element with clamped capacitance, C_0 , converts the kinetic energy given by the velocity, \dot{x} , in electrical current, *I*, by means of the force factor, $\alpha = \frac{C_0 V_{oc}}{x_M}$, which is defined by voltage (V_{oc}) and displacement magnitude (x_M) in open-circuit conditions.

In order to collect the output power provided by the harvester, a full-bridge rectifier with a smoothing capacitor, C_r , was used (Fig. 3b). The bridge was then connected to a resistive load, R_l , making a lumped circuit. For this system, the modified coupling coefficient is $k_m^2 = \frac{\alpha^2}{\kappa c_0}$, which is the ratio between electrostatic energy stored in the system and the elastic energy, and the quality factor, $Q = \frac{\sqrt{KM}}{c}$. When the system is driven by a constant-magnitude force, in the case of continuous excitation of the harvester

at resonance, the energy extraction circuit is attenuating the displacement of the beam. This effect is more evident for systems with high Q or k_m^2 . Thus, the normalized power can be expressed as [20]:

$$\bar{P} = 8 k_m^2 \frac{r\Omega^2}{\left(r\Omega + \frac{\pi}{2}\right)^2} |\bar{x}|^2, \qquad (2)$$

where the normalized frequency is $\Omega = \frac{\omega}{\omega_0}$, normalized load - $r = \omega_0 C_0 R_l$, and normalized displacement, \bar{x} . It is known [21] that the figure of merit, $k_m^2 Q$, can be used for the comparison of the harvester performance. If the harvester shows $k_m^2 Q > \pi$, the normalized power is fundamentally limited by the mechanical damping of the system, and therefore tends to reach the value of Q. In this scenario, we deal with a highly-coupled device showing two identical power peaks that correspond to the short-circuit (ω_0) and open-circuit ($\omega_0\sqrt{1+k_m^2}$) frequencies. Furthermore, it is possible to increase the bandwidth of the harvester by adjusting the resistive load to maximize the power output. The analytical model for standard rectification circuit was discussed in detail by Shu et al [22].

The transformer ratio, α , was measured in open-circuit condition. From the impedance spectrum investigations, we obtained the lumped-model parameters, summarized in Table 1. The measured shortand open-circuit frequencies were 104.5 Hz and 105.9 Hz, respectively. The experimentally identified quality factor was Q = 396 and the modified structural coupling $k_m^2 = 0.028$, leading to an estimated figure of merit of $k_m^2 Q = 11$, which indicated that our studied system was highly-coupled [22]. These values are typical for single crystal materials [23], but they represent a significant improvement compared to our previous results [15] ($k_m^2 Q = 0.3$). Indeed, in this study we had considered LiNbO₃ (YXlt)/128°/90°, which effectively increased the transverse electromechanical coupling [13] compared to (YXlt)/36°/90° orientation. Therefore, the choice of the crystal orientation is of essential importance in terms of performances.

To estimate the harvester output power, the DC voltage was measured after a full bridge rectification circuit. The circuit was essentially composed by four 1N4148 diodes and one smoothing capacitor, $C_r = 12 \ \mu\text{F}$, followed by a variable resistive load, R_l . The rectified voltage, V_{DC} , was used to estimate the power output, $P = \frac{V_{DC}^2}{R_l}$. In the case of highly coupled regime, we expect to have two power peaks for short- and open-circuit frequencies. A comparison between simulated (Eq. 2) and experimental power is given in Fig. 4a. The V_{DC} was measured as a function of the frequency for each resistive load, in order to investigate the behavior at both short- and open-circuit conditions. In terms of harvested power, a projection of the theoretical surface along with the data points is plotted in the plane to illustrate the dependency on the resistive load and frequency. Two power peaks were identified: one relative to the

short-circuit resonance (50 k Ω) and one at open-circuit (3 M Ω). Nevertheless, four fast-switching diodes were used in the rectification circuit, resulting in power losses for the short-circuit resistive load, where the voltage level was close to 1 V. Since diode losses were not considered in the model, the power at short-circuit resonance is lower compared to the simulation. Nevertheless, the optimal rectified power output for the short-circuit frequency was $P_{sc} = 29 \mu W$, while for open-circuit frequency was $P_{oc} = 41.5$ µW. Finally, considering the low acceleration level, we attain a considerably high normalized power density (NPD = 965 μ W/cm²/g²), which is among the highest values reported in literature for a Pb and Pb-free based energy harvester. Typically, other devices that are working in similar conditions can reach higher power levels but are implementing two piezoelectric layers in bimorph configuration. Even if LiNbO₃ bimorphs with inversion layers were fabricated by using bulk crystal wafers [16], their impedance matching issues with the electronic interfaces and limited flexibility were limiting their application. For this reason, we propose a thick film implementation, which results in a highly coupled device with rectified power response. Many authors tend to report the instantaneous power for vibrational harvesters which could be misleading, while in this work the rectified power was measured, confirming that it is possible to implement LiNbO3 thick films transducers as vibrational energy harvester with standard electronic interface.

The normalized power densities of low frequency energy harvesters, reported in literature, are compared to the one of our studied vibrational LiNbO₃ harvesters in Fig. 4b. The state of the art is a selection of works for both MEMS devices [24,25,26,27] exploiting materials that are deposited by means of microfabrication techniques, and meso-scale [28,29,30,31] devices exploiting commercially available ceramics or single crystals. Although they present different dimensions, their power is normalized by the active surface and input acceleration for a fair comparison. Finally, Pb-free MEMS devices show NPD up to 495 µW/cm²/g² for (MgZr)AlN/Si harvesters at high frequency (at 792 Hz) [32] and 71 µW/cm²/g² for LiNbO₃ (at 1140 Hz) [15]. Therefore, devices with large tip mass can attain higher power densities compared to simple beams as in our previous work, but they tend also be more fragile at high acceleration levels. In the case of low frequency devices PMN-PT highly coupled devices [31] can attain 9083 µW/cm²/g² at 28 Hz, while PZT [6] and LiNbO₃ harvesters can attain similar power densities in the order of $1039 \,\mu$ W/cm²/g² and 965 μ W/cm²/g² respectively. However, it is important to note that at low frequency, (YXlt)/128°/90° LiNbO₃/Si unimorph attains the highest performance among other leadfree materials. Even if relaxor ferroelectrics are showing higher normalized power density compared to our prototype, their expensive price and toxicity prevent them to be widely used in energy harvesting applications. Eventually, the performance of vibrational energy harvesters based on LiNbO₃ thick films can be further improved by implementing bimorphs on metal substrates [6].

In conclusion, a beam with tip mass, based on $LiNbO_3$ single-crystal film bonded on silicon, was investigated for the vibrational energy harvesting at low frequencies. The dynamic behavior was simulated by finite element method, whereas the single degree of freedom lumped-model was used to

simulate the power level of the harvester. Our device presented a short- and open-circuit resonance frequencies of respectively 104.5 and 105.9 Hz. The device electromechanical coupling was $k_m^2 = 2.8\%$, and the quality factor was Q = 396. The figure of merit ($k_m^2 Q = 11$) was indicating highly-coupled device behavior and the measured power density was 965 μ W/cm²/g², which is among the highest reported values compared to both Pb- and Pb-free vibrational harvesting devices. For future investigations, the fragility of the silicon-based beams can be overcome by the integration of LiNbO₃ thick films on more reliable or flexible substrates, such as metal foils.

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Table 1: Lumped model parameters, identified experimentally for $(YXlt)/128^{\circ}/90^{\circ}$ LiNbO₃/Si harvester.

M (g)	C (mN.s/m)	K (N/m)	α (mN/V)
0.416	0.691	179.54	0.17



Figure 1: a) Microfabrication flowchart for (YXlt)/128°/90° LiNbO₃ harvester: 1) Cr/Au sputtering on polished LiNbO₃ and Si wafers; 2) EVG bonding of Cr/Au interface layers; 3) lapping and polishing of LiNbO₃ wafer; 4) UV photolithography patterning and development; 5) deposition and lift-off of Cr/Au top electrodes; 6) mechanical dicing of cantilever structures with explicit rotation around X (in red) of the axes Y and Z; 7) deep reactive ion etching of Si substrate. Bonding interface quality between Si and (YXlt)/128°/90° LiNbO₃ wafers: b) ultrasonic imaging of bonded wafers surface, c) SEM image of cross section of (YXlt)/128°/90° LiNbO₃/Si heterostructure. d) Characterization setup for vibrational energy harvesters, (inset) (YXlt)/128°/90° LiNbO₃/Si energy harvester clamped and wire-bonded on PCB.



Figure 2: Frequency dependence of a) experimental and FEM voltage response (RMS) and b) impedance and phase response along with BVD simulation for (YXlt)/128°/90° LiNbO₃/Si harvester



Figure 3: Inertial energy harvester schematics: a) equivalent electro-mechanical lumped model for inertial energy harvester, b) full-bridge rectification circuit with resistive load.



Figure 4: a) Comparison of theoretical and experimental power output of $(YXIt)/128^{\circ}/90^{\circ}$ LiNbO₃/Si harvester in terms of frequency and resistive load. The transparent surface represents the simulated power using the proposed model, while the experimental maxima are depicted. b) Comparison of normalized power density for Pb-based (red) and Pb-free (green) MEMS and mesoscale devices, reported in literature and studied in this work.

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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