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# DYNAMICAL ZETA FUNCTIONS IN THE NONORIENTABLE CASE 

YONAH BORNS-WEIL AND SHU SHEN


#### Abstract

We use a simple argument to extend the microlocal proofs of meromorphicity of dynamical zeta functions to the nonorientable case. In the special case of geodesic flow on a connected non-orientable negatively curved closed surface, we compute the order of vanishing of the zeta function at the zero point to be the first Betti number of the surface.


## 1. Background

In this note we use a simple geometric argument to extend the results of Dyatlov and Zworski $[5,6]$ and of Dyatlov and Guillarmou $[3,4]$ to Axiom A flows with nonorientable stable and unstable bundles. It is classically known that on a closed manifold there are countably many closed orbits of such flows, and therefore one can define the Ruelle zeta function

$$
\zeta_{R}(\lambda)=\prod_{\gamma^{\sharp}}\left(1-e^{i \lambda T_{\gamma}^{\sharp}}\right)
$$

where the product is taken over all primitive closed geodesics $\gamma^{\sharp}$ with corresponding periods $T_{\gamma}^{\sharp}$. Note that by [3, Lemma 1.17] and [4, Section 3], this product converges for $\operatorname{Im}(\lambda) \gg 1$ large enough. The meromorphic continuation of $\zeta_{R}$ to all of $\mathbb{C}$ was conjectured by Smale [13], and proved by Fried [8] under analyticity assumptions. The case of smooth Anosov flows was first answered by Giulietti, Liverani and Policott [9] and then with microlocal methods by Dyatlov and Zworski [5] for manifolds with orientable stable and unstable bundles, and was extended to Axiom A flows by Dyatlov and Guillarmou [3, 4] under the same orientability assumptions. In [9, Appendix B], the authors also outlined ideas for removing the orientability assumptions.

We remove the orientability assumption and give a full proof for Axiom A flows. Specifically, we shall show

Theorem 1. If $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is an Axiom $A$ flow on a closed manifold, the Ruelle zeta function $\zeta_{R}$ extends to a meromorphic function on $\mathbb{C}$.

The definition of an Axiom A flow is given as Definition 1.3.
We then restrict to the case of contact Anosov flow on a 3-manifold, and study the order of vanishing of $\zeta_{R}$ at $\lambda=0$. An important example is when $M=S^{*} \Sigma$,
the cosphere bundle of a connected negatively curved closed surface $\Sigma$, and $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is geodesic flow [1]. This problem was treated in [6] in the case where the stable bundle is orientable, and it was shown that the order of vanishing is $b_{1}(M)-2$, where $b_{1}(M)$ is the first Betti number of $M$.

We shall show that for nonorientable stable bundle, the analogous result is the following:

Theorem 2. Let $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ be the Reeb flow on a connected contact closed 3-manifold. If $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is Anosov with nonorientable stable bundle $E_{s}$, the Ruelle zeta function has vanishing order at $\lambda=0$ equal to $b_{1}\left(\mathscr{O}\left(E_{s}\right)\right)$, the dimension of the first de Rham cohomology with coefficients in the orientation line bundle of $E_{s}$.

The orientation line bundle is reviewed in Definition 1.5.
In the special case of the geodesic flow on $M=S^{*} \Sigma$ with $\Sigma$ nonorientable, the vanishing order at $\lambda=0$ is given by $b_{1}(\Sigma)$, as is shown in Proposition 3.10. This is in contrast to the orientable case, in which it is $b_{1}(\Sigma)-2$.

More precisely, let $\chi^{\prime}(\Sigma)$ be the derived Euler characteristic of $\Sigma$, i.e.,

$$
\chi^{\prime}(\Sigma)=\sum_{i=0}^{2}(-1)^{i} i b_{i}(\Sigma)= \begin{cases}-b_{1}(\Sigma)+2, & \text { if } \Sigma \text { is orientable } \\ -b_{1}(\Sigma), & \text { otherwise }\end{cases}
$$

Corollary 3. If $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is the geodesic flow on the cosphere bundle of a connected negatively curved closed surface (orientable or not), the Ruelle zeta function has vanishing order at $\lambda=0$ equal to $-\chi^{\prime}(\Sigma)$.
1.1. Axiom A Flows. Let $M$ be a compact manifold without boundary of dimension $n$, and let $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ be a flow on $M$ generated by the vector field $V \in C^{\infty}(M ; T M)$.
Definition 1.1. A $\phi_{t}$-invariant set $K \subseteq M$ is called hyperbolic for the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ if $V$ does not vanish on $K$ and for each $x \in K$ the tangent space $T_{x} M$ can be written as the direct sum

$$
T_{x} M=E_{0}(x) \oplus E_{s}(x) \oplus E_{u}(x)
$$

where $E_{0}(x)=\operatorname{span}(V(x)), E_{s}, E_{u}$ are continuous $\phi_{t}$-invariant vector bundles on $K$, and for some Riemannian metric $|\cdot|$, there are $C, \theta>0$ such that for all $t>0$,

$$
\begin{align*}
\left|d \phi_{t}(x) v\right|_{\phi_{t}(x)} & \leq C e^{-\theta t}|v|_{x} & & v \in E_{s}(x) \\
\left|d \phi_{-t}(x) w\right|_{\phi_{-t}(x)} & \leq C e^{-\theta t}|w|_{x} & & w \in E_{u}(x) \tag{1}
\end{align*}
$$

In the important case where all of $M$ is hyperbolic, we call $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ an Anosov flow.
There is an analogous notion of hyperbolicity at fixed points.
Definition 1.2. A fixed point $x \in M$, i.e., $V(x)=0$, is called hyperbolic if the differential $D V(x)$ has no eigenvalues with vanishing real part.

A generalization of Anosov flows is the following, given first by Smale [13, II. 5 Definition 5.1]:

Definition 1.3. The flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is called Axiom $A$ if
(1) all fixed points of $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ are hyperbolic,
(2) the closure $\mathcal{K}$ of the union of all closed orbits of $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is hyperbolic,
(3) the nonwandering set ([4, Definition 2.2]) of $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is the disjoint union of the set of fixed points and $\mathcal{K}$.

We now recall the definition of a locally maximal set, given in [4, Definition 2.4].
Definition 1.4. A compact $\phi_{t}$-invariant set $K \subseteq M$ is called locally maximal for $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ if there is a neighborhood $V$ of $K$ such that

$$
K=\bigcap_{t \in \mathbb{R}} \phi_{t}(V)
$$

We may then state the key proposition, which generalises [4, Proposition 3.1] to the case where $E_{s}$ or $E_{u}$ is not necessarily orientable on $\mathcal{K}$.

Proposition 1.4. Let $K \subseteq M$ be a locally maximal hyperbolic set for $\left(\phi_{t}\right)_{t \in \mathbb{R}}$, and let $\zeta_{K}$ be defined as the Ruelle zeta function where we only take the product over trajectories in $K$. Then $\zeta_{K}$ has a continuation to a meromorphic function on all of $\mathbb{C}$.

Theorem 1 follows from Proposition 1.4, as we may remark that by [13, II. 5 Theorem 5.2] we can write $\mathcal{K}=K_{1} \sqcup \cdots \sqcup K_{N}$ with $K_{j}$ basic hyperbolic ${ }^{1}$. Then the product

$$
\zeta_{R}(\lambda)=\prod_{j=1}^{N} \zeta_{K_{j}}(\lambda)
$$

which a priori holds for $\operatorname{Im}(\lambda) \gg 1$, gives that $\zeta_{R}$ also has a meromorphic continuation to all of $\mathbb{C}$.

The goal of Section 2 is to prove Proposition 1.4.
1.2. The Orientation Bundle. To fix notation we recall the definition of transition functions of a vector bundle. Given a continuous real vector bundle $E$ of rank $k$ over a manifold $M$ with projection map $\pi$, let $U_{\alpha}, U_{\beta} \subseteq M$ be two small open sets with nonempty intersection, and let $\psi_{\alpha}: \pi^{-1} U_{\alpha} \rightarrow U_{\alpha} \times \mathbb{R}^{n}, \psi_{\beta}: \pi^{-1} U_{\beta} \rightarrow U_{\beta} \times \mathbb{R}^{n}$ be local trivializations. Then the map $\psi_{\alpha} \circ \psi_{\beta}^{-1}:\left(U_{\alpha} \cap U_{\beta}\right) \times \mathbb{R}^{n} \rightarrow\left(U_{\alpha} \cap U_{\beta}\right) \times \mathbb{R}^{n}$ is of the form

$$
\psi_{\alpha} \circ \psi_{\beta}^{-1}(p, v)=\left(p, \tau_{\alpha \beta}(p) v\right)
$$

[^0]where $\tau_{\alpha \beta} \in C^{0}\left(U_{\alpha} \cap U_{\beta}, \mathrm{GL}_{k}(\mathbb{R})\right)$ is called a transition function. If the local trivializations can be chosen such that $\tau_{\alpha \beta}$ are smooth, then $E$ is a smooth vector bundle. Similarly, if $\tau_{\alpha \beta}$ can be chosen to be locally constant functions, then $E$ is a flat vector bundle.

Furthermore, suppose we are given an open cover $\left(U_{\alpha}\right)_{\alpha \in A}$ of $M$ together with a set of continuous (resp. smooth, resp. locally constant) $\mathrm{GL}_{k}(\mathbb{R})$-valued functions $\left(\tau_{\alpha \beta}\right) \underset{\substack{\alpha, \beta \in A \\ U_{\alpha} \cap U_{\beta} \neq \emptyset}}{ }$ with $\tau_{\alpha \alpha}=I$ on $U_{\alpha}$. Then there exists a continuous (resp. smooth, resp. flat) vector bundle $E$ with transition functions $\tau_{\alpha \beta}$, provided the following triple product property holds:

$$
\tau_{\alpha \beta}(p) \tau_{\beta \gamma}(p) \tau_{\gamma \alpha}(p)=I
$$

for any $p \in U_{\alpha} \cap U_{\beta} \cap U_{\gamma}$.
Definition 1.5. If $E$ is a continuous (but not necessarily smooth) real vector bundle over $M$ with transition functions $\tau_{\alpha \beta}$, the orientation bundle of $E$ is a smooth flat line bundle $\mathscr{O}(E)$ with transition functions

$$
\sigma_{\alpha \beta}(p)=\operatorname{sgn} \operatorname{det}\left(\tau_{\alpha \beta}(p)\right)= \begin{cases}1 & \operatorname{det}\left(\tau_{\alpha \beta}(p)\right)>0 \\ -1 & \operatorname{det}\left(\tau_{\alpha \beta}(p)\right)<0\end{cases}
$$

Recall that if $f: M \rightarrow M$ is a map, we say $f$ lifts to a bundle map $F: E \rightarrow E$ if $\pi \circ F=f \circ \pi$.

Since $\mathscr{O}(E)$ is a flat vector bundle, using the associated flat connection, we can lift the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ to a flow $\left(\widetilde{\Phi}_{t}\right)_{t \in \mathbb{R}}$ on $\mathscr{O}(E)$. If the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ on $M$ lifts to a flow $\left(\Phi_{t}\right)_{t \in \mathbb{R}}$ on $E$, if $\psi, \eta$ are distinct trivializations of $E$ near $p, \phi_{t}(p)$ respectively, and $\widetilde{\psi}, \widetilde{\eta}$ are trivializations of $\mathscr{O}(E)$ near $p, \phi_{t}(p)$ respectively, we have for $p \in M$ and $l \in \mathscr{O}(E)_{p}:$

$$
\begin{equation*}
\widetilde{\Phi}_{t}(l)=\widetilde{\eta}^{-1}\left(\phi_{t}(p), \operatorname{sgn}\left(\left.\operatorname{det}\left(\eta \Phi_{t} \psi^{-1}\right)\right|_{p}\right) \operatorname{proj}_{2} \widetilde{\psi}(l)\right) \tag{2}
\end{equation*}
$$

where $\operatorname{proj}_{2}$ is the obvious projection to the second component.
1.3. Geodesic flows. Let $Z$ be a negatively curved closed Riemannian manifold. Let $M=S^{*} Z$ be the cosphere bundle on $Z$. It is classical that the geodesic flow on $M$ is Anosov [1].

Let $\pi: M \rightarrow Z$ be the canonical projection. For $x \in M$, we have a morphism of linear spaces

$$
\begin{equation*}
\pi_{*}: T_{x} M \rightarrow T_{\pi(x)} Z \tag{3}
\end{equation*}
$$

The following proposition is classical [1, Section 22] and [12, Proposition 6]. We include a proof for the sake of completeness.

Proposition 1.5. The morphism $\pi_{*}$ induces an isomorphism of continuous vector bundles on $M$,

$$
\begin{equation*}
E_{s} \oplus E_{0} \simeq \pi^{*}(T Z) \tag{4}
\end{equation*}
$$

Proof. Since both sides of (4) have the same dimension, it is enough to show that $\left.\pi_{*}\right|_{E_{s} \oplus E_{0}}$ is injective. We will show this using Jacobi fields. It is convenient to work on the sphere bundle $M^{\prime}=S Z$. We identify $M^{\prime}$ with $M$ via the Riemannian metric on $Z$.

We follow [7, Section II.H]. Let $\mathcal{M}$ be the total space of $T Z$. Denote still by $\pi$ : $\mathcal{M} \rightarrow Z$ the obvious projection. Let $T^{V} \mathcal{M} \subset T \mathcal{M}$ be the vertical subbundle of $T \mathcal{M}$. The Levi-Civita connection on $T Z$ induces a horizontal subbundle $T^{H} \mathcal{M} \subset T \mathcal{M}$ of $T \mathcal{M}$, so that

$$
\begin{equation*}
T \mathcal{M}=T^{V} \mathcal{M} \oplus T^{H} \mathcal{M} \tag{5}
\end{equation*}
$$

Since $T^{V} \mathcal{M} \simeq \pi^{*}(T Z)$ and $T^{H} \mathcal{M} \simeq \pi^{*}(T Z)$, by (5), we can identify the smooth vector bundles,

$$
T \mathcal{M}=\pi^{*}(T Z \oplus T Z)
$$

For $x=(z, v) \in \mathcal{M}$, let $\gamma_{x}$ be the unique geodesic on $Z$ such that $\left(\gamma_{x}(0), \dot{\gamma}_{x}(0)\right)=$ $(z, v)$. For $w \in T_{x} \mathcal{M}$, let $J_{x, w} \in C^{\infty}\left(\gamma_{x},\left.T Z\right|_{\gamma_{x}}\right)$ be the unique Jacobi field along $\gamma_{x}$ such that $\left(J_{x, w}(0), \dot{J}_{x, w}(0)\right)=w$, where $\dot{J}_{x, w}$ is the covariant derivation of $J_{x, w}$ in the direction $\dot{\gamma}_{x}$. Recall that a Jacobi field $J$ is called stable, if there is $C>0$ such that for all $t \geq 0$,

$$
|J(t)| \leq C .
$$

By [7, Proposition VI.A], given $x \in \mathcal{M}$, for any $Y_{1} \in T_{z} Z$, there exists one and only one stable Jacobi field $J$ along $\gamma_{x}$ such that $J(0)=Y_{1}$.

For $x=(z, v) \in M^{\prime}$, we have

$$
T_{x} M^{\prime}=\left\{\left(Y_{1}, Y_{2}\right) \in T_{z} Z \oplus T_{z} Z:\left\langle Y_{2}, v\right\rangle=0\right\}
$$

The morphism $\pi_{*}$ in (3) is just

$$
w \in T_{x} M^{\prime} \rightarrow J_{x, w}(0) \in T_{z} Z .
$$

By [7, Proposition VI.B], $w \in E_{s}(x) \oplus E_{0}(x)$ if and only if the Jacobi fields $J_{x, w}$ is stable. By the uniqueness of stable Jacobi fields, we see that $\left.\pi_{*}\right|_{E_{s} \oplus E_{0}}$ is injective.

Since $E_{0}$ is a trivial line bundle, our proposition implies immediately:
Corollary 1.6. We have the isomorphism of smooth flat line bundles

$$
\mathscr{O}\left(E_{s}\right) \simeq \pi^{*}(\mathscr{O}(T Z))
$$

## 2. Proof of Proposition 1.4

We use the notation of [5]. If $0 \leq k \leq n-1$, let $\mathcal{E}_{0}^{k} \subset \Lambda^{k}\left(T^{*} M\right)$ denote the subbundle of $k$-forms $\omega$ such that $\iota_{V} \omega=0$, where $\iota$ denotes interior multiplication.

Let $\widetilde{\mathcal{E}_{0}^{k}}=\mathcal{E}_{0}^{k} \otimes \mathscr{O}\left(E_{s}\right)$. We consider the pullback $\phi_{-t}^{*}$ on sections of $\widetilde{\mathcal{E}_{0}^{k}}$. Note that the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ lifts to a flow $\left(\Phi_{t}\right)_{t \in \mathbb{R}}$ on $\mathcal{E}_{0}^{k}$. Indeed, for $p \in M, \omega \in \mathcal{E}_{0, p}^{k}, \Phi_{t} \omega \in \mathcal{E}_{0, \phi_{t}(p)}^{k}$ is defined for $v_{1}, \cdots, v_{k} \in T_{\phi_{t}(p)} M$ by

$$
\begin{equation*}
\Phi_{t} \omega\left(v_{1}, \ldots, v_{k}\right)=\omega\left(\left(\left.d \phi_{t}\right|_{p}\right)^{-1} v_{1}, \cdots,\left(\left.d \phi_{t}\right|_{p}\right)^{-1} v_{k}\right) \tag{6}
\end{equation*}
$$

Note that from the above formula, it is easy to check that ${ }_{{ }_{V}} \Phi_{t} \omega=0$. Recall also that the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ lifts to a flow $\widetilde{\Phi}_{t}$ on $\mathscr{O}\left(E_{s}\right)$ (see (2)). For a section $s$ in $\widetilde{\mathcal{E}_{0}^{k}}$, we have

$$
\begin{equation*}
\phi_{-t}^{*} s(p)=\left(\Phi_{t} \otimes \widetilde{\Phi}_{t}\right)\left(s\left(\phi_{-t}(p)\right)\right) \tag{7}
\end{equation*}
$$

Let $\chi \in C^{\infty}(M)$ be a smooth function whose support is contained in a small neighborhood of $K$ such that $\chi(x)=1$ for all $x \in K$. We now invoke the Guillemin trace formula (see [11, pp. 501-502], [5, Appendix B], [3, (4.6)]) which says that the flat trace $\left.\operatorname{tr}^{b} \chi \phi_{-t}^{*} \chi\right|_{C^{\infty}\left(M ; \widetilde{\mathcal{F}_{0}^{k}}\right)}$ is a distribution on $(0, \infty)$ given by

$$
\begin{equation*}
\left.\operatorname{tr}^{b} \chi \phi_{-t}^{*} \chi\right|_{C^{\infty}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right)}=\sum_{\gamma \subset K} \frac{T_{\gamma}^{\sharp} \operatorname{tr}^{\widetilde{\mathcal{E}}_{0, y}^{k}}\left(\Phi_{T_{\gamma}} \otimes \widetilde{\Phi}_{T_{\gamma}}\right)}{\left|\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right)\right|} \delta_{t-T_{\gamma}}, \tag{8}
\end{equation*}
$$

where the sum is taken over all the periodic trajectories $\gamma$ in $K$ with period $T_{\gamma}$ and primitive period $T_{\gamma}^{\sharp}, y$ is any point on $\gamma$, and $\mathcal{P}_{\gamma}=\left.d \phi_{-T_{\gamma}}\right|_{\left(E_{s} \oplus E_{u}\right)_{y}}$ is the linearized Poincaré map at $y$. Note that as trace and determinant are invariant under conjugation, the right hand side does not depend on $y$.

By (6), the trace of $\Phi_{T_{\gamma}}$ on $\mathcal{E}_{0, y}^{k}$ is just $\operatorname{tr}\left(\bigwedge^{k} \mathcal{P}_{\gamma}\right)$. By (2), we may take trivializations $\psi, \widetilde{\psi}$ of $E_{s}, \mathscr{O}\left(E_{s}\right)$ in a neighborhood of $y$ and have the induced lifting on $\mathscr{O}\left(E_{s}\right)$ to be $\operatorname{sgn}\left(\operatorname{det}\left(\left.\psi d \phi_{T_{\gamma}}\right|_{E_{s, y}} \psi^{-1}\right) \mid\right)$. By definition we get this to be equal to

$$
\operatorname{sgn}\left(\left.\operatorname{det} d \phi_{T_{\gamma}}\right|_{E_{s, y}}\right)=\operatorname{sgn}\left(\left.\operatorname{det} d \phi_{-T_{\gamma}}\right|_{E_{s, y}}\right)=\operatorname{sgn} \operatorname{det}\left(\left.\mathcal{P}_{\gamma}\right|_{E_{s}}\right),
$$

and as it is a map between one dimensional spaces, the trace is given by that expression as well. By the above consideration, we can rewrite (8) as

$$
\begin{equation*}
\left.\operatorname{tr}^{b} \chi \phi_{-t}^{*} \chi\right|_{C \infty\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right)}=\sum_{\gamma \subset K} \frac{T_{\gamma}^{\sharp} \operatorname{tr}\left(\bigwedge^{k} \mathcal{P}_{\gamma}\right) \operatorname{sgn}\left(\operatorname{det} \mathcal{P}_{\gamma} \mid E_{s}\right)}{\left|\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right)\right|} \delta_{t-T_{\gamma}} . \tag{9}
\end{equation*}
$$

Let us follow [4, Section 3]. By [4, Lemma 3.2], we may and we will assume that near $K,\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is an open hyperbolic system in the sense of [3, Assumptions (A1)-(A4)].

By [3, Lemma 1.17], there is $C>0$ such that for all $t \geq 0$,

$$
\begin{equation*}
\mid\left\{\gamma \text { closed trajectory in } K: T_{\gamma} \leq t\right\} \mid \leq C e^{C t} . \tag{10}
\end{equation*}
$$

For $\operatorname{Im}(\lambda) \gg 1$ big enough, set

$$
\begin{equation*}
\zeta_{K, k}(\lambda)=\exp \left(-\sum_{\gamma \subset K} \frac{T_{\gamma}^{\sharp}}{T_{\gamma}} \frac{\operatorname{tr}\left(\bigwedge^{k} \mathcal{P}_{\gamma}\right) \operatorname{sgn}\left(\left.\operatorname{det} \mathcal{P}_{\gamma}\right|_{E_{s}}\right)}{\left|\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right)\right|} e^{i \lambda T_{\gamma}}\right) . \tag{11}
\end{equation*}
$$

Lemma 2.1. For $\operatorname{Im}(\lambda) \gg 1$ big enough, we have

$$
\begin{equation*}
\partial_{\lambda} \log \zeta_{K, k}(\lambda)=-\left.i \int_{0}^{\infty} e^{i \lambda t} \operatorname{tr}^{\mathrm{b}} \chi \phi_{-t}^{*} \chi\right|_{C^{\infty}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right)} d t \tag{12}
\end{equation*}
$$

The function $\zeta_{K, k}(\lambda)$ has a holomorphic extension to $\mathbb{C}$.
Proof. Let us first remark that by (9) and (10), the right hand side of (12) is well defined. Taking a logarithm and differentiating (11) and using Guillemin trace formula (9), we get (12). The last part of the lemma follows from the arguments of [3, Section 4].

Recall that for $\operatorname{Im}(\lambda) \gg 1$ big enough, we have

$$
\begin{equation*}
\zeta_{K}(\lambda)=\prod_{\gamma^{\sharp} \subset K}\left(1-e^{i \lambda T_{\lambda}^{\sharp}}\right)=\exp \left(-\sum_{\gamma \subset K} \frac{T_{\gamma}^{\sharp}}{T_{\gamma}} e^{i \lambda T_{\gamma}}\right) . \tag{13}
\end{equation*}
$$

Proposition 1.4 is a consequence of the following lemma. This lemma was stated in [2], but we restate and prove it for convenience.

Lemma 2.2. The following identity of meromorphic functions on $\mathbb{C}$ holds,

$$
\begin{equation*}
\zeta_{K}(\lambda)=\prod_{k=0}^{n-1}\left(\zeta_{K, k}(\lambda)\right)^{(-1)^{k+\operatorname{dim} E_{s}}} \tag{14}
\end{equation*}
$$

Proof. Following [5, (2.4)-(2.5)], since $\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right)=\sum_{k=0}^{n-1}(-1)^{k} \operatorname{tr}\left(\bigwedge^{k} \mathcal{P}_{\gamma}\right)$, by (11) and (13), it is enough to show

$$
\begin{equation*}
\left|\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right)\right|=(-1)^{\operatorname{dim} E_{s}} \operatorname{sgn}\left(\left.\operatorname{det} \mathcal{P}_{\gamma}\right|_{E_{s}}\right) \operatorname{det}\left(I-\mathcal{P}_{\gamma}\right) . \tag{15}
\end{equation*}
$$

Remark that

$$
\begin{align*}
\operatorname{det}\left(I-\mathcal{P}_{\gamma}\right) & =\operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}\right|_{E_{u}}\right) \operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}\right|_{E_{s}}\right) \\
& =(-1)^{\operatorname{dim} E_{s}} \operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}\right|_{E_{u}}\right) \operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}^{-1}\right|_{E_{s}}\right) \operatorname{det}\left(\left.\mathcal{P}_{\gamma}\right|_{E_{s}}\right) . \tag{16}
\end{align*}
$$

As time is running in the negative direction, we have by (1) that the eigenvalues $\lambda$ of $\left.\mathcal{P}_{\gamma}\right|_{E_{u}}$ have $|\lambda|<1$, and the eigenvalues $\mu$ of $\left.\mathcal{P}_{\gamma}^{-1}\right|_{E_{s}}$ have $|\mu|<1$. This gives any
eigenvalues of $I-\left.\mathcal{P}_{\gamma}\right|_{E_{u}}$ to be either $1-\lambda$ for $\lambda \in(-1,1)$ or conjugate pairs $1-\lambda, 1-\bar{\lambda}$ when $\lambda$ is not real. In any case, we get by multiplying that

$$
\operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}\right|_{E_{u}}\right)>0
$$

and similarly

$$
\operatorname{det}\left(I-\left.\mathcal{P}_{\gamma}^{-1}\right|_{E_{s}}\right)>0 .
$$

Then taking signs in (16), we get (15).
Remark. The key point of our argument is based on the smoothness of $\mathscr{O}\left(E_{s}\right)$. Thanks to this property, most of the analytic arguments in the proof of Proposition 1.4 are reduced to [3]. In [2, Section 2], Baladi and Tsujii used the orientation bundle in a different way for the flow with discrete time.

## 3. Vanishing Order at Zero on a Contact 3-Manifold

In this section, we assume that $M$ is a connected closed 3-manifold with a contact form $\alpha$, and that $V$ is the associated Reeb vector field. We suppose also that the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ of $V$ is Anosov. One such example would be when $M=S^{*} \Sigma$, the cosphere bundle of a connected closed surface $\Sigma$ with negative (variable) curvature, and $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is geodesic flow.

The following result was proven in [6]:
Theorem 3.1. If $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is a contact Anosov flow on a connected closed 3-manifold with orientable $E_{u}$ and $E_{s}$, the Ruelle zeta function has vanishing order at $\lambda=0$ equal to $b_{1}(M)-2$, where $b_{1}(M)$ denotes the first Betti number of $M$.

The goal of this section is to determine the order of vanishing of $\zeta_{R}$ at 0 in the case that $E_{s}, E_{u}$ are not orientable, and hence give a proof of Theorem 2. We remark that since a contact manifold is orientable, orientability of $E_{s}$ is equivalent to orientability of $E_{u}$.
3.1. The twisted cohomology. Let us recall some background and general facts on the twisted cohomology of a flat vector bundle. Let $X$ be a closed manifold. Let $F$ be a flat vector bundle on $X$ with flat connection $\nabla$. It induces a sheaf $\mathcal{F}$ on $X$ defined by locally constant sections, i.e., if $U \subset X$ is an open set, then

$$
\mathcal{F}(U)=\left\{s \in C^{\infty}\left(U ;\left.F\right|_{U}\right): \nabla s=0\right\}
$$

The twisted cohomology $H^{\bullet}(X ; F)$ is defined by the cohomology of the sheaf $\mathcal{F}[10$, Section II.4.4]. They are the algebraic invariants which describe the rigidity properties of the global flat sections of $F$. Let $b_{k}(F)$ be the twisted Betti number

$$
b_{k}(F)=\operatorname{dim} H^{k}(X ; F)
$$

If $F$ is the trivial line bundle, we get the classical de Rham cohomology with real coefficients.

To evaluate $H^{\bullet}(X ; F)$, one can use the twisted de Rham complex. Indeed, if we denote $F^{k}=\Lambda^{k}\left(T^{*} X\right) \otimes F$, the flat connection $\nabla$ extends to an operator $d_{k}$ : $C^{\infty}\left(X ; F^{k}\right) \rightarrow C^{\infty}\left(X ; F^{k+1}\right)$ by Leibniz rule: if $\alpha \in C^{\infty}\left(X ; \Lambda^{k}\left(T^{*} X\right)\right)$ and $s \in$ $C^{\infty}(X ; F)$, we have

$$
d_{k}(\alpha \cdot s)=d \alpha \cdot s+(-1)^{k} \alpha \wedge \nabla s
$$

By the flatness of $\nabla$, we have $d_{k+1} d_{k}=0$, so that $\left(C^{\infty}\left(X ; F^{\bullet}\right), d_{\bullet}\right)$ is a complex. By the de Rham isomorphism [10, Théorème II.4.7.1], we have

$$
\begin{equation*}
H^{k}(X ; F)=\operatorname{ker} d_{k} / \operatorname{Im} d_{k-1} . \tag{17}
\end{equation*}
$$

As an analogue of [6, Lemma 2.1], using the theory of elliptic operators, we can evaluate $H^{\bullet}(X ; F)$ using the complex of twisted currents, or more generally twisted currents with wavefront conditions.

More precisely, let $\Gamma \subset T^{*} X$ be a closed cone. We denote by $\mathcal{D}_{\Gamma}^{\prime}\left(X ; F^{k}\right)$ the space of $F^{k}$-valued distributions whose wavefront set is contained in $\Gamma$ (see [6, Section 2.1]). By microlocality, we have

$$
d_{k}: \mathcal{D}_{\Gamma}^{\prime}\left(X ; F^{k}\right) \rightarrow \mathcal{D}_{\Gamma}^{\prime}\left(X ; F^{k+1}\right)
$$

For simplicity, we will write $d$ sometimes.
Lemma 3.2. If $u \in \mathcal{D}_{\Gamma}^{\prime}\left(X ; F^{k}\right)$ and $d u \in C^{\infty}\left(X ; F^{k+1}\right)$, then there exist $v \in$ $C^{\infty}\left(X ; F^{k}\right)$ and $w \in \mathcal{D}_{\Gamma}^{\prime}\left(X ; F^{k-1}\right)$ such that

$$
u=v+d w
$$

In particular, if $u \in \mathcal{D}_{\Gamma}^{\prime}(X ; F)$ and $d u \in C^{\infty}\left(X ; F^{1}\right)$, then $u \in C^{\infty}(X ; F)$.
Proof. Take a Riemannian metric on $X$ and a Hermitian metric on $F$. Remark that these two metrics induce a fibrewise scalar product $\langle\cdot, \cdot\rangle$ on $F^{k}$. For $u, v \in C^{\infty}\left(X ; F^{k}\right)$, we can define the $L^{2}$-product by

$$
\begin{equation*}
\langle u, v\rangle_{L^{2}\left(X ; F^{k}\right)}=\int_{X}\langle u, v\rangle d \mathrm{vol}, \tag{18}
\end{equation*}
$$

where $d \mathrm{vol}$ is a volume form. Let $\delta_{k+1}: C^{\infty}\left(X ; F^{k+1}\right) \rightarrow C^{\infty}\left(X ; F^{k}\right)$ be the formal adjoint of $d$ with respect to the $L^{2}$-product (18). Define the twisted Hodge Laplacian by

$$
\Delta_{k}=d_{k-1} \delta_{k}+\delta_{k+1} d_{k}: C^{\infty}\left(X ; F^{k}\right) \rightarrow C^{\infty}\left(X ; F^{k}\right)
$$

Then $\Delta_{k}$ is an essentially self-adjoint second order elliptic differential operator. The remainder of the proof carries over identically from that of [6, Lemma 2.1].

Remark that if $F$ is the orientation bundle of certain vector bundle and $u \in$ $C^{\infty}(X ; F)$, then for $x \in X,|u(x)|^{2}$ is independent of the choice of trivializations. It defines a Hermitian metric on $F$.
3.2. Resonant State Spaces. Let $M$ be a connected 3-dimensional closed manifold with a contact form $\alpha \in C^{\infty}\left(M, T^{*} M\right)$. Let $V$ be the associated Reeb vector field. Then,

$$
\begin{equation*}
i_{V} \alpha=1, \quad \iota_{V} d \alpha=0 . \tag{19}
\end{equation*}
$$

We assume that the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ associated to $V$ is Anosov. Let $E_{u}^{*} \subset T^{*} M$ be the dual of $E_{s}$. We will apply the results of Section 3.1 to the case where $(X, F, \Gamma)=$ $\left(M, \mathcal{O}\left(E_{s}\right), E_{u}^{*}\right)$.

Since the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ is Anosov, we have $K=M$. For $0 \leq k \leq 2$, we write $\zeta_{k}=\zeta_{K, k}$. By (14), we have

$$
\begin{equation*}
\zeta_{R}(\lambda)=\frac{\zeta_{1}(\lambda)}{\zeta_{0}(\lambda) \zeta_{2}(\lambda)} \tag{20}
\end{equation*}
$$

We consider the operator $P_{k}=-i \mathcal{L}_{V}$, where $\mathcal{L}_{V}$ (in a slight abuse of notation) denotes the natural action on sections of $\widetilde{\mathcal{E}_{0}^{k}}$, given by the Lie derivative on sections of $\mathcal{E}_{0}^{k}$ tensored with the flat connection on $\mathscr{O}\left(E_{s}\right)$. For $\operatorname{Im} \lambda \gg 1$ large enough, the integral $R_{k}(\lambda)=i \int_{0}^{\infty} e^{i \lambda t} \phi_{-t}^{*} d t$ converges and defines a bounded operator on the $L^{2}$-space; this is nothing more than the resolvent operator of $P_{k}$. Then by [ $6, \operatorname{Section} 2.3$ ] we have that $R_{k}$ extends meromorphically to the entire complex plane,

$$
R_{k}(\lambda): C^{\infty}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right) \rightarrow \mathcal{D}^{\prime}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right)
$$

More precisely, near $\lambda_{0} \in \mathbb{C}$, we have

$$
R_{k}(\lambda)=R_{k, H}(\lambda)-\sum_{j=1}^{J\left(\lambda_{0}\right)} \frac{\left(P_{k}-\lambda\right)^{j-1} \Pi_{k}}{\left(\lambda-\lambda_{0}\right)^{j}}
$$

where $R_{k, H}$ is a holomorphic family defined near $\lambda_{0}, J\left(\lambda_{0}\right) \in \mathbb{N}$, and $\Pi_{k}$ has rank $m_{k}\left(\lambda_{0}\right)<\infty$. By the arguments at the end of [5], we have that at $\lambda_{0}$, the function $\zeta_{k}$ has a zero of order $m_{k}\left(\lambda_{0}\right)$.

We define the space of resonant states at $\lambda_{0}$ to be

$$
\operatorname{Res}_{k}\left(\lambda_{0}\right)=\left\{u \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right):\left(P_{k}-\lambda_{0}\right) u=0\right\} .
$$

Then a special case of [6, Lemma 2.2] gives the following:
Lemma 3.3. Suppose $P_{k}$ satisfies the semisimplicity condition:

$$
u \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \widetilde{\mathcal{E}_{0}^{k}}\right), \quad\left(P_{k}-\lambda_{0}\right)^{2} u=0 \quad \Longrightarrow \quad\left(P_{k}-\lambda_{0}\right) u=0
$$

Then $m_{k}\left(\lambda_{0}\right)=\operatorname{dim} \operatorname{Res}_{k}\left(\lambda_{0}\right)$.

Recall that we are trying to find the order at $\lambda=0$ of $\zeta_{R}$, which by (20) is simply

$$
\begin{equation*}
m_{R}(0)=m_{1}(0)-m_{0}(0)-m_{2}(0) \tag{21}
\end{equation*}
$$

We will compute each of these individually, by computing $\operatorname{dim}_{\operatorname{Res}}^{k}$ ( 0 ) and checking that the semisimplicity condition in Lemma 3.3 holds.

We begin with twisted "0-forms', which are just sections of the orientation bundle $\mathscr{O}\left(E_{s}\right)$.

Proposition 3.4. If $E_{s}$ is nonorientable, the space $\operatorname{Res}_{0}(0)$ is $\{0\}$.
Proof. Suppose $u \in \operatorname{Res}_{0}(0)$, i.e.,

$$
\begin{equation*}
P_{0} u=0 \tag{22}
\end{equation*}
$$

Since the flow $\left(\phi_{t}\right)_{t \in \mathbb{R}}$ preserves the contact volume form $\alpha \wedge d \alpha, P_{0}: C^{\infty}\left(M ; \mathscr{O}\left(E_{s}\right)\right) \rightarrow$ $C^{\infty}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ is a symmetric operator with respect to the $L^{2}$-product (18). By [6, Lemma 2.3], $u \in C^{\infty}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$. Using $\partial_{t}\left(\phi_{-t}^{*} u\right)=-\phi_{-t}^{*} \nabla_{V} u$ (where $\nabla$ is the flat connection), we see that $u$ is constant on the flow line: for all $t \in \mathbb{R}$,

$$
\begin{equation*}
u=\phi_{-t}^{*} u \tag{23}
\end{equation*}
$$

Let $(x, v) \in T M$. The pairing $\langle d u(x), v\rangle$ is an element of $\mathscr{O}\left(E_{s}\right)_{x}$. By (7) and (23), we have

$$
\langle d u(x), v\rangle=\left\langle\phi_{-t}^{*}(d u)(x), v\right\rangle=\widetilde{\Phi}_{t}\left\langle d u\left(\phi_{-t}(x)\right), d \phi_{-t}(x) v\right\rangle .
$$

If $v \in E_{u}(x)$, then sending $t \rightarrow \infty$ gives $\langle d u(x), v\rangle=0$ by (1). Similarly, if $v \in E_{s}(x)$, then sending $t \rightarrow-\infty$ gives $\langle d u(x), v\rangle=0$. This shows that

$$
\begin{equation*}
\left.d u\right|_{E_{s} \oplus E_{u}}=0 . \tag{24}
\end{equation*}
$$

By Cartan's formula and by (22), we have $\iota_{V} d u=0$, i.e.,

$$
\begin{equation*}
\left.d u\right|_{E_{0}}=0 . \tag{25}
\end{equation*}
$$

By (24) and (25), we have $d u=0$. So $u \in H^{0}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$. Since $E_{s}$ is nonorientable, we have $H^{0}\left(M ; \mathscr{O}\left(E_{s}\right)\right)=0$, so $u=0$ and $\operatorname{Res}_{0}(0)$ is trivial.

Corollary 3.5. If $E_{s}$ is nonorientable, the multiplicity for 0 -forms is $m_{0}(0)=0$.
Proof. If $P_{0}^{2}(u)=0$, then $P_{0} u \in \operatorname{Res}_{0}(0)$. By Proposition 3.4, $P_{0} u=0$, so $u \in \operatorname{Res}_{0}(0)$. This shows semisimplicity, so by Lemma 3.3 we see that $m_{0}(0)=\operatorname{dim} \operatorname{Res}_{0}(0)=0$.

Proposition 3.6. If $E_{s}$ is nonorientable, the space $\operatorname{Res}_{2}(0)$ is $\{0\}$.
Proof. We claim that

$$
\begin{equation*}
\alpha \wedge: \mathcal{E}_{0}^{2} \rightarrow \mathcal{E}^{3} \tag{26}
\end{equation*}
$$

is a bundle isomorphism. Indeed, using (19), it is easy to see that the inverse of (26) is given by $\iota_{V}$. Tensoring with $\mathscr{O}\left(E_{s}\right)$, we get a bundle isomorphism

$$
\begin{equation*}
\alpha \wedge: \widetilde{\mathcal{E}_{0}^{2}} \rightarrow \widetilde{\mathcal{E}^{3}} \tag{27}
\end{equation*}
$$

Let $u \in \operatorname{Res}_{2}(0)$. Since $\mathcal{E}^{3}$ is generated by $\alpha \wedge d \alpha$, by (27), there is $v \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ such that $\alpha \wedge u=v \alpha \wedge d \alpha$. Applying $\iota_{V}$ and using $\iota_{V} u=0$, we have $u=v d \alpha$. Then

$$
0=P_{2}(u)=\left(P_{0} v\right) d \alpha
$$

But this gives $P_{0} v=0$, so by Proposition 3.4 we have $v=0$. Therefore, $u=0$.
The following is then clear for the same reason as Corollary 3.5.
Corollary 3.7. If $E_{s}$ is nonorientable, the multiplicity for 2 -forms is $m_{2}(0)=0$.
We now turn to the case of $P_{1}$ acting on the space of twisted 1-form-valued distributions $\mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \widetilde{\mathcal{E}_{0}^{1}}\right)$. We can now state the analogous proposition for 1 -forms:

Proposition 3.8. If $E_{s}$ is nonorientable, the space $\operatorname{Res}_{1}(0)$ has dimension $b_{1}\left(\mathscr{O}\left(E_{s}\right)\right)$.
Proof. The proof is analogous to that of [6, Lemma 3.4], but slightly easier due to the holomorphy of the resolvent $R_{0}$ near 0 . Let $u \in \operatorname{Res}_{1}(0)$. Then $d u \in \operatorname{Res}_{2}(0)$ by Proposition 3.6, so $d u=0$. By Lemma 3.2 there is a $\phi \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ such that

$$
u-d \phi \in C^{\infty}\left(M ; \widetilde{\mathcal{E}^{1}}\right), \quad d(u-d \phi)=0
$$

We shall show that the map:

$$
\Theta: u \mapsto[u-d \phi] \in H^{1}\left(M ; \mathscr{O}\left(E_{s}\right)\right)
$$

is well-defined, linear and bijective, which is enough to prove the lemma.

## Well-Definedness and linearity:

Suppose there is another section $\psi \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ with $u-d \psi \in C^{\infty}\left(M ; \widetilde{\mathcal{E}^{1}}\right)$. Then subtracting gives $d(\phi-\psi) \in C^{\infty}\left(M ; \widetilde{\mathcal{E}^{1}}\right)$, so $\phi-\psi \in C^{\infty}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ by Lemma 3.2. This shows that the map $\Theta$ is well-defined. It is also easy to see that $\Theta$ is linear.

## Injectivity:

If $\Theta(u)=0$, then $u-d \phi$ is exact, so without loss of generality we can assume that $u=d \phi$. Combining with $\iota_{V} u=0$, we get $\phi \in \operatorname{Res}_{0}(0)$, so $\phi=0$ by Proposition 3.4. Therefore $u=0$, and this shows $\Theta$ to be injective.

## Surjectivity:

Let $v \in C^{\infty}\left(M ; \widetilde{\mathcal{E}^{1}}\right)$ with $d v=0$. Then as $m_{0}(0)=0$, the resolvent $R_{0}$ is holomorphic near 0 . Take $\phi=i R_{0}(0) \iota_{V} v \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$. Then $P_{0} \phi=i_{V} v$. This
rearranges to $\iota_{V}(v+d \phi)=0$, so $v+d \phi \in \operatorname{Res}_{1}(0)$. This gives that $\Theta$ is surjective, and completes the proof of our proposition.

Proposition 3.9. If $E_{s}$ is nonorientable, the multiplicity for 1 -forms is $m_{1}(0)=$ $b_{1}\left(\mathscr{O}\left(E_{s}\right)\right)$.

Proof. By Lemma 3.3, we must only check that the semisimplicity condition is satisfied. Take $u \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \widetilde{\mathcal{E}_{0}^{1}}\right)$ such that $\left(P_{1}\right)^{2} u=0$. Then $v=\iota_{V} d u \in \operatorname{Res}_{1}(0)$. It is enough to show that $v=0$.

Recall that in the proof of Proposition 3.8, we have seen that elements in $\operatorname{Res}_{1}(0)$ are closed. In particular,

$$
\begin{equation*}
d v=0 \tag{28}
\end{equation*}
$$

Note that $\alpha \wedge d u \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \widetilde{\mathcal{E}^{3}}\right)$. We claim that

$$
\begin{equation*}
\alpha \wedge d u=0 \tag{29}
\end{equation*}
$$

Indeed, there is some $a \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ such that

$$
\alpha \wedge d u=a \alpha \wedge d \alpha
$$

Since $\mathcal{L}_{V}(\alpha)=0$, by (28), we have

$$
\left(\mathcal{L}_{V} a\right) \alpha \wedge d \alpha=\alpha \wedge \mathcal{L}_{V}(d u)=\alpha \wedge d \iota_{V} d u=\alpha \wedge d v=0
$$

Then $\mathcal{L}_{V} a=0$, so $a=0$ by Proposition 3.4. This gives (29).
Since $\alpha(V)=1$, we have $(\alpha \wedge) \circ \iota_{V}+\iota_{V} \circ(\alpha \wedge)=$ id. By (29), we have

$$
\begin{equation*}
d u=\left((\alpha \wedge) \circ \iota_{V}+\iota_{V} \circ(\alpha \wedge)\right) d u=\alpha \wedge v . \tag{30}
\end{equation*}
$$

By Lemma 3.2 and by $(28)$, there are $w \in C^{\infty}\left(M ; \widetilde{\mathcal{E}^{1}}\right), \phi \in \mathcal{D}_{E_{u}^{*}}^{\prime}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$ such that

$$
\begin{equation*}
v=w+d \phi, \quad d w=0 \tag{31}
\end{equation*}
$$

Then

$$
\begin{equation*}
\iota_{V} w=\iota_{V}(v-d \phi)=-\mathcal{L}_{V} \phi . \tag{32}
\end{equation*}
$$

In particular, $\mathcal{L}_{V} \phi$ is smooth. We compute by Stokes' Theorem and by (30)-(32),

$$
\begin{aligned}
0 & =\operatorname{Re} \int_{M} d u \wedge \bar{w}=\operatorname{Re} \int_{M} \alpha \wedge d \phi \wedge \bar{w}=\operatorname{Re} \int_{M} \phi \bar{w} \wedge d \alpha \\
& =\operatorname{Re} \int_{M} \iota_{V}(\phi \bar{w}) \alpha \wedge d \alpha=-\operatorname{Re} \int_{M} \phi\left(\overline{\mathcal{L}_{V} \phi}\right) \alpha \wedge d \alpha=-\operatorname{Re}\left\langle\mathcal{L}_{V} \phi, \phi\right\rangle_{L^{2}\left(M ; \mathscr{O}\left(E_{s}\right)\right)},
\end{aligned}
$$

where the fourth equality comes from the fact the

$$
(\alpha \wedge) \circ \iota_{V}(\phi \bar{w} \wedge d \alpha)=\left((\alpha \wedge) \circ \iota_{V}+\iota_{V} \circ(\alpha \wedge)\right)(\phi \bar{w} \wedge d \alpha)=\phi \bar{w} \wedge d \alpha
$$

In the above formula, we use the fact that a product of two twisted forms is untwisted. By [6, Lemma 2.3] we have $\phi \in C^{\infty}\left(M ; \mathscr{O}\left(E_{s}\right)\right)$, so $v \in C^{\infty}\left(M ; \widetilde{\mathcal{E}_{0}^{1}}\right)$. Then by the same argument as in Proposition 3.4 (see [6, Lemma 3.5]) we have $v=0$.

Now Theorem 2 is a consequence of (21), Corollaries 3.5, 3.7, and Proposition 3.9.
Let $\Sigma$ be a connected negatively curved closed surface. Take $M=S^{*} \Sigma$. By Corollary 1.6, we have

$$
H^{1}\left(M ; \mathscr{O}\left(E_{s}\right)\right)=H^{1}\left(M ; \pi^{*} \mathscr{O}(T \Sigma)\right)
$$

Proposition 3.10. If $\Sigma$ is a connected negatively curved closed surface (oriented or not), we have

$$
\begin{equation*}
\operatorname{dim} H^{1}\left(M ; \pi^{*} \mathscr{O}(T \Sigma)\right)=\operatorname{dim} H^{1}(\Sigma) \tag{33}
\end{equation*}
$$

Proof. By the Gysin long exact sequence, we have the exact sequence
$0 \longrightarrow H^{1}(\Sigma ; \mathscr{O}(T \Sigma)) \xrightarrow{\pi^{*}} H^{1}\left(M ; \pi^{*} \mathscr{O}(T \Sigma)\right) \xrightarrow{\pi_{*}} H^{0}(\Sigma) \xrightarrow{e \wedge} H^{2}(\Sigma ; \mathscr{O}(T \Sigma)) \longrightarrow$,
where $\pi^{*}$ is the pullback, $\pi_{*}$ is the integration along the fibre of $M \rightarrow \Sigma$, and $e \in$ $H^{2}(\Sigma ; \mathscr{O}(T \Sigma))$ is the Euler class of $T \Sigma$.

We claim that the last map

$$
e \wedge: H^{0}(\Sigma) \rightarrow H^{2}(\Sigma ; \mathscr{O}(T \Sigma))
$$

in the Gysin exact sequence is an isomorphism. Indeed, since $\Sigma$ is connected, we have $\operatorname{dim} H^{0}(\Sigma)=1$, and by Poincaré duality, $\operatorname{dim} H^{2}(\Sigma ; \mathscr{O}(T \Sigma))=1$. It is enough to show that $e \in H^{2}(\Sigma ; \mathscr{O}(T \Sigma))$ is non zero, or equivalently $\int_{\Sigma} e \neq 0$. This is a consequence of the fact that $\Sigma$ has negative curvature, as $e=K \mu$ where $\mu$ is the Riemannian density and $K<0$ is the Gauss curvature.

Therefore, we get an isomorphism

$$
\begin{equation*}
\pi^{*}: H^{1}(\Sigma ; \mathscr{O}(T \Sigma)) \simeq H^{1}\left(M ; \pi^{*} \mathscr{O}(T \Sigma)\right) \tag{34}
\end{equation*}
$$

By Poincaré duality, we have

$$
\begin{equation*}
H^{1}(\Sigma ; \mathscr{O}(T \Sigma)) \simeq\left(H^{1}(\Sigma)\right)^{*} \tag{35}
\end{equation*}
$$

By (34) and (35), we get (33).

Now Corollary 3 is a consequence of Theorem 2 and Proposition 3.10.

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[^0]:    ${ }^{1}$ These are locally maximal hyperbolic by definition (see [4, Definition 2.5]).

