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# Riemann theta functions, Fredholm and wronskian representations of the solutions to the KdV equation

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#### Abstract

We degenerate the finite gap solutions of the KdV equation from the general formulation given in terms of abelian functions when the gaps tends to points, to get solutions to the KdV equation given in terms of Fredholm determinants and wronskians. For this we establish a link between Riemann theta functions, Fredholm determinants and wronskians. This gives the bridge between the algebro-geometric approach and the Darboux dressing method.

## 1 Introduction

We consider the KdV equation in the following normalization

$$u_t = 6uu_x - u_{xxx},\tag{1}$$

where the subscripts x and t as usual denote partial derivatives.

This equation (1) was introduced for the first time in 1895 by Korteweg and de Vries [56]. This equation appears in a wide range of physical problems and describes the propagation of waves with weak dispersion in various nonlinear media.

A method of resolution was given in 1967 by Gardner et al. [52]. It was proven that this equation is a complete integrable system by Zakharov and Faddeev in 1971 [76].

Solutions were constructed by Hirota in 1971 by using the bilinear method [53]. Its and Matveev present solutions in terms of Riemann theta functions [54] in 1975. Lax gives in the same year the expressions of periodic and almost periodic solutions [59]. A lot of works have been realized in the following years. We can mention for example Airault et al. in 1977 [3], Adler and Moser in 1978 [2], Ablowitz and Cornille in 1979 [1], Freeman and Nimmo in 1984 [5], Matveev in 1992 [64], Ma in 2004 [60], Kovalyov in 2005 [58] and more recently Ma in 2015 [61].

In the following, we are interested in the algebro-geometric approach given by Its and Matveev in 1975. We degenerate the solutions to the KdV equation given in terms of Riemann theta functions to get solutions in terms of Fredholm determinants. Then we give a representation in terms of wronskians. This gives the correspondence between the algebro-geometric approach and the Darboux dressing method.

# 2 The KdV equation and its solutions in terms of theta functions

We consider the Riemann surface  $\Gamma$  of the algebraic curve defined by

$$\omega^2 = \prod_{j=1}^{2g+1} (z - E_j)$$

with  $E_j \neq E_k$ ,  $j \neq k$ . Let *D* be some divisor  $D = \sum_{j=1}^{g} P_j$ ,  $P_j \in \Gamma$ . The so-called finite gap solution of the KdV equation

$$u_t = 6uu_x - u_{xxx} \tag{2}$$

can be expressed in the form [54]

$$u(x,t) = -2\partial_x^2 \left[\ln\theta(xg + tv + l)\right] + C.$$
(3)

We recall briefly, the notations. In (3),  $\theta$  is the Riemann function defined by

$$\theta(z) = \sum_{k \in \mathbf{Z}^g} \exp\{\pi i(Bk|k) + 2\pi i(k|z)\},\tag{4}$$

constructed from the matrix of the B-periods of the surface  $\Gamma$ , and the vectors g, v, l are defined by

$$g_j = 2ic_{j1},\tag{5}$$

$$v_j = 8i(\frac{c_{j1}}{2}\sum_{k=1}^{2g+1} E_k + c_{j2}), \tag{6}$$

$$l_j = -\sum_{k=1}^g \int_{\infty}^{P_k} dU_j + \frac{j}{2} - \frac{1}{2} \sum_{k=1}^g B_{kj},\tag{7}$$

$$C = \sum_{k=1}^{2g+1} E_k - 2\sum_{k=1}^g \int_{a_k} z dU_k,$$
(8)

the coefficients  $c_{jk}$  being relating with abelian differential  $dU_j$  by

$$dU_j = \frac{\sum_{k=1}^g c_{jk} z^{g-k}}{\sqrt{\prod_{k=1}^{2g+1} (z - E_k)}} dz,$$
(9)

and coefficients  $c_{jk}$  can be obtained by solving the system of linear equations

$$\int_{a_k} dU_j = \delta_{jk}, \quad 1 \le j \le g, \quad 1 \le k \le g.$$

# **3** Degeneracy of solutions

We suppose that  $E_j$  are real,  $E_m < E_j$  if m < j and try to evaluate the limits of all objects in formula (3) when  $E_{2m}$ ,  $E_{2m+1}$  tends to  $-\alpha_m$ ,  $-\alpha_m = -\kappa_m^2$ ,  $\kappa_m > 0$ , for  $1 \le m \le g$ , and  $E_1$  tends to 0 (these ideas were first presented by A. Its and V.B. Matveev, exposed for example in [4]).

#### 3.1 Degeneracy of the components of the solution

**3.1.1** Limit of  $P(z) = \prod_{j=1}^{2g+1} (z - E_j)$ 

The limit of  $P(z) = \prod_{j=1}^{2g+1} (z - E_j)$  is evidently equal to  $\tilde{P}(z) = z \prod_{j=1}^{g} (z + \alpha_j)^2$ 

**3.1.2** Limit of 
$$dU_m = \frac{\sum_{k=1}^g c_{mk} z^{g-k}}{\sqrt{\prod_{k=1}^{2g+1} (z - E_k)}} dz$$

The limit of  $dU_m$  is equal to  $d\tilde{U}_m = \frac{\varphi_m(z)}{\sqrt{z}\prod_{j=1}^g (z+\alpha_j)} dz$ , where  $\varphi_m(z) = \sum_{j=1}^g \tilde{z}_j = \frac{\varphi_m(z)}{z} dz$ 

 $\sum_{k=1}^{g} \tilde{c}_{mk} z^{g-k}$ . The normalization condition takes the form in the limit

$$\int_{a_k} dU_j \to \frac{2\pi i\varphi_j(-\alpha_k)}{-i\kappa_k \prod_{m \neq k} (-\alpha_k + \alpha_m)} = \delta_{kj},\tag{10}$$

which proves that the numbers  $-\alpha_m$ ,  $m \neq k$  are the zeros of the polynomials  $\varphi_k(z)$ , and so  $\varphi_k(z)$  can be written as  $\varphi_k(z) = \tilde{c}_{k1} \prod_{m \neq k} (z + \alpha_m)$ . By (10), we get in the limit

$$\tilde{c}_{k1} = -\frac{\kappa_k}{2\pi i}$$

 $\operatorname{So}$ 

$$d\tilde{U}_k = -\frac{\kappa_k}{2\pi i \sqrt{z}(z+\alpha_k)} dz$$

#### **3.1.3** Limit of $v_k$ and $g_k$

By identification of the powers of  $z^{g-2}$  in (11)

$$\tilde{\varphi_k}(z) = \tilde{c_{k1}} \prod_{l \neq k} (z + \alpha_l) = \sum_{j=1}^g \tilde{c_{kj}} z^{g-j}, \qquad (11)$$

we get in the limit

$$\tilde{c}_{k1} \sum_{l=1}^{g} \alpha_l = \tilde{c}_{k2}.$$

So we have the limit values of  $v_k$  and  $g_k$ :

$$\tilde{v}_k = \frac{4i\kappa_k^3}{\pi}$$

and

$$\tilde{g}_k = \frac{-i\kappa_k}{\pi}.$$

#### **3.1.4** Limit of $U_k(P)$ and $B_{mk}$

For  $\lambda_0 = -\alpha_m = -\kappa_m^2$ ,  $I = \int_{\lambda_0}^0 dU_k \to \frac{1}{2}\tilde{B}_{mk}$ . The integral I can be easily evaluate along the real axis on the upper sheet of surface  $\Gamma$  and we get

$$I \to \frac{i}{2\pi} \ln \left| \frac{\kappa_m + \kappa_k}{\kappa_m - \kappa_k} \right|$$

So we have the limit values of matrix  ${\cal B}$  :

$$\tilde{B}_{mk} = \frac{i}{\pi} \ln \left| \frac{\kappa_m + \kappa_k}{\kappa_m - \kappa_k} \right|.$$

Therefore  $iB_{kk}$  tends to  $-\infty$ . As previously, we have

$$\int_{\infty}^{P} dU_j \to -\frac{i}{2\pi} \ln \left| \frac{\kappa_j - \sqrt{z_P}}{\kappa_j + \sqrt{z_P}} \right|.$$
(12)

#### **3.1.5** Limit of argument of exponential in $\theta(p)$

Let us denote  $A_0$  the argument of exponential in  $\theta(p) = \sum_{k \in \mathbb{Z}^g} \exp\{\pi i(Bk|k) + 2\pi i(k|p)\}$ .

 $A_0$  can be rewritten in the form

$$A_0 = \pi i \sum_{j=1}^g B_{jj} k_j (k_j - 1) + 2\pi i \sum_{j>m} B_{mj} k_m k_j + \sum_{j=1}^g \pi i (2p_j + B_{jj}) k_j.$$
(13)

Using the inequality  $k_j(k_j - 1) \ge 0$  for all  $k \in \mathbb{Z}^g$  and the fact that  $iB_{kk}$  tends to  $-\infty$ , we can reduce the limit  $\tilde{\theta}$  of  $\theta(p)$  to a finite sum taken over vectors  $k \in \mathbb{Z}^g$  such that each  $k_j$  must be equal to 0 or 1.

So, if we denote A the argument of  $\theta(xg + tv + l)$ , it can be written in the form

$$A = \pi i \sum_{j=1}^{g} B_{jj} k_j (k_j - 1) + 2\pi i \sum_{j>m} B_{mj} k_m k_j + \sum_{j=1}^{g} k_j [2\pi i (g_j x + v_j t) -\pi i (-j + 2\sum_{k=1}^{g} \int_{\infty}^{P_k} dU_j + \sum_{m \neq j} B_{mj})].$$

In other words

$$A = \pi i \sum_{j=1}^{g} B_{jj} k_j (k_j - 1) + 2\pi i \sum_{j>m} B_{mj} k_m k_j + \sum_{j=1}^{g} k_j Q_j,$$

with

$$Q_j = 2\pi i (g_j x + v_j t) + \beta_j$$

and

$$\beta_j = -\pi i (-j + 2\sum_{k=1}^g \int_{\infty}^{P_k} dU_j + \sum_{m \neq j} B_{mj}).$$

The quantity  $\beta_j$  has a finite limit value  $\tilde{\beta}_j$  independent from x and t.

#### **3.1.6** Limit of $\theta(xg + vt + l)$

By means of the inequality  $k_j(k_j - 1) \ge 0$  for all  $k \in \mathbb{Z}^g$  and the previous relation  $iB_{kk}$  tends to  $-\infty$ , it turns out that the limit  $\tilde{\theta}$  of  $\theta(xg + tv + l)$  reduce to a finite sum taken over vectors  $k \in \mathbb{Z}^g$  with the property that each  $k_j$  must be equal to 0 or 1.

$$\tilde{\theta} = \sum_{k \in \mathbf{Z}^g, \, k_j = 0 \text{ or } 1} \exp\{\sum_{m > j} 2\ln\left|\frac{\kappa_m - \kappa_j}{\kappa_m + \kappa_j}\right| k_m k_j + (\sum_{j=1}^g 2\kappa_j x - 8\kappa_j^3 t + 2\kappa_j x_j + \pi j i + \sum_{m \neq j} \ln\left|\frac{\kappa_m + \kappa_j}{\kappa_m - \kappa_j}\right|) k_j\},$$

with

$$x_j = \frac{1}{2\kappa_j} \sum_{k=1}^g \ln \left| \frac{\sqrt{z_k} - i\kappa_j}{\sqrt{z_k} + i\kappa_j} \right|.$$

It can be rewritten as

$$\tilde{\theta} = \sum_{J \subset \{1,\dots,g\}} \prod_{j \in J} (-1)^j \prod_{j \in J \ k \notin J} \left| \frac{\kappa_j + \kappa_k}{\kappa_j - \kappa_k} \right| \exp \sum_{j \in J} 2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j).$$
(14)

#### **3.1.7** Limit of the coefficient C

The coefficient C is defined in (8) by

$$C = \sum_{k=1}^{2g+1} E_k - 2\sum_{k=1}^g \int_{a_k} z dU_k = C_1 + C_2.$$

can be evaluated as follows.

$$C_{2} = -2\sum_{k=1}^{g} \int_{a_{k}} z dU_{k} = -2\sum_{k=1}^{g} \int_{a_{k}} \frac{-\kappa_{k} z dz}{2\pi\sqrt{z}(z+\alpha_{k})} = \sum_{k=1}^{g} \frac{\kappa_{k}}{\pi} \int_{a_{k}} \frac{\sqrt{z} dz}{(z+\alpha_{k})}$$
$$= \sum_{k=1}^{g} \frac{\kappa_{k}}{\pi} 2i\pi(-i\kappa_{k}) = \sum_{k=1}^{g} 2\kappa_{k}^{2} = 2\sum_{k=1}^{g} \alpha_{k}.$$

Thus when the gaps tends to points,, the coefficient C tends to  $\tilde{C}$  equal to

$$\tilde{C} = 2\sum_{k=1}^{g} -\alpha_k + 2\sum_{k=1}^{g} \alpha_k = 0.$$

### 3.2 Degenerate solution to the KdV equation

Therefore, we have the following representation of the degenerate solution to the KdV equation

**Theorem 3.1** The fonction u defined by

$$u(x,t) = -2\partial_x^2 \ln\left(\sum_{J \subset \{1,\dots,g\}} \prod_{j \in J} (-1)^j \prod_{j \in J \ k \notin J} \left| \frac{\kappa_j + \kappa_k}{\kappa_j - \kappa_k} \right| \exp\left(\sum_{j \in J} 2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right)\right), (15)$$

with  $\kappa_i$ , and  $x_j$  arbitrary real parameters, is a solution to the KdV equation (1).

## 4 From theta function to Fredholm determinant

# 4.1 The link between the degenerate solution and the Fredholm determinant

In a recent paper, Kirillov and Van Diejen [72] have given formulas in terms of determinants for zonal spherical functions on hyperboloids. In particular, they compute  $\det(I + A)$ , where I is the unit matrix and  $A = (a_{jk})_{1 \le j,k \le m}$  the matrix defined as :

$$a_{jk} = \frac{2\epsilon_j \kappa_j}{\kappa_j + \kappa_k} \exp(-2\kappa_j x) \prod_{l \neq j} \left| \frac{\kappa_l + \kappa_j}{\kappa_l - \kappa_j} \right|,\tag{16}$$

where  $\epsilon_j \in \{-1; +1\}$  and  $\kappa_j > 0$  for  $1 \le j \le N$ . Then  $\det(I + A)$  has the following form

$$\det(I+A) = \sum_{J \subset \{1,\dots,N\}} \exp\left(-2x \sum_{j \in J} \kappa_j\right) \prod_{j \in J} \epsilon_j \prod_{j \in J \ k \notin J} \left|\frac{\kappa_j + \kappa_k}{\kappa_j - \kappa_k}\right|.$$
(17)

Using the same strategy, we can compute  $\det(I+A)$  where  $A=(a_{jk})_{1\leq j,k\leq m}$  is the matrix defined as :

$$a_{jk} = \frac{(-1)^j 2\kappa_j}{\kappa_j + \kappa_k} \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq j} \left|\frac{\kappa_l + \kappa_j}{\kappa_l - \kappa_j}\right|,$$

 $x_j$  being an arbitrary parameter. Then det(I + A) has the following form

$$\det(I+A) = \sum_{J \subset \{1,\dots,N\}} \exp\left(\sum_{j \in J} 2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right) \prod_{j \in J} (-1)^j \prod_{j \in J \ k \notin J} \left|\frac{\kappa_j + \kappa_k}{\kappa_j - \kappa_k}\right|, (18)$$

By the previous section,

$$\tilde{\theta} = \sum_{J \subset \{1,\dots,g\}} \prod_{j \in J} (-1)^j \prod_{j \in J \ k \notin J} \left| \frac{\kappa_j + \kappa_k}{\kappa_j - \kappa_k} \right| \exp\left(\sum_{j \in J} (2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j))\right).$$
(19)

If we compare the expression (18) to (19), we have clearly the equality with g = N

$$\tilde{\theta} = \det(I+A). \tag{20}$$

# 4.2 Solution to the KdV equation in terms of Fredholm determinant

So we have the following representation of the solutions to the KdV equation

**Theorem 4.1** The fonction u defined by

$$u(x,t) = -2\partial_x^2 \ln(\det(I+A)), \qquad (21)$$

with A the matrix defined by  $A = (a_{jk})_{1 \leq j,k \leq N}$ 

$$a_{jk} = \frac{(-1)^j 2\kappa_j}{\kappa_j + \kappa_k} \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq j} \left|\frac{\kappa_l + \kappa_j}{\kappa_l - \kappa_j}\right|,\tag{22}$$

and  $\kappa_j$ ,  $x_j$  arbitrary real parameters, is a solution to the KdV equation (1).

If we consider the matrix B defined by

$$b_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq k} \left|\frac{\kappa_l + \kappa_j}{\kappa_l - \kappa_k}\right|$$

it is easy to verify that  $\det(I + A) = \det(I + B)$ , and so we can give another representation of the solutions to the KdV equation. We get the following statement :

**Theorem 4.2** The fonction u defined by

$$u(x,t) = -2\partial_x^2 \ln(\det(I+B)), \tag{23}$$

with B the matrix defined by  $B = (b_{jk})_{1 \le j,k \le m}$ 

$$b_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq k} \left|\frac{\kappa_l + \kappa_j}{\kappa_l - \kappa_k}\right|,\tag{24}$$

and  $\kappa_j$ ,  $x_j$  arbitrary real parameters, is a solution to the KdV equation (1).

We can also consider the matrix C defined by

$$c_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \frac{\prod_{l \neq k} |\kappa_l + \kappa_j|}{\prod_{l \neq j} |\kappa_l - \kappa_j|}.$$

It is easy to check that det(I + A) = det(I + C), and so we can give a third representation of the solutions to the KdV equation :

**Theorem 4.3** The fonction u defined by

$$u(x,t) = -2\partial_x^2 \ln(\det(I+C)), \qquad (25)$$

with C the matrix defined by  $C = (c_{jk})_{1 \le j,k \le m}$ 

$$c_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \frac{\prod_{l \neq k} |\kappa_l + \kappa_j|}{\prod_{l \neq j} |\kappa_l - \kappa_j|},\tag{26}$$

and  $\kappa_j$ ,  $x_j$  arbitrary real parameters, is a solution to the KdV equation (1).

Another possibility is to choose the matrix D defined by

$$d_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq j} \left|\frac{\kappa_l + \kappa_k}{\kappa_l - \kappa_j}\right|.$$

It is also easy to check that  $\det(I + A) = \det(I + D)$ , and so we can give another representation of the solutions to the KdV equation :

**Theorem 4.4** The fonction u defined by

$$u(x,t) = -2\partial_x^2 \ln(\det(I+D)), \qquad (27)$$

with C the matrix defined by  $D = (d_{jk})_{1 \le j,k \le m}$ 

$$d_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq j} \left|\frac{\kappa_l + \kappa_k}{\kappa_l - \kappa_j}\right|,\tag{28}$$

and  $\kappa_j$ ,  $x_j$  arbitrary real parameters, is a solution to the KdV equation (1).

It remains to find the link between this Fredholm determinant and a certain wronskian.

## 5 From Fredholm determinants to wronskians

#### 5.1 Link between Fredholm determinants and wronskians

In this section, we consider the following functions

$$\phi_j(x) = \sinh(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j), \qquad (29)$$

where  $\kappa_j$  are real numbers such that  $\kappa_1 \leq \ldots \leq \kappa_N$ , and  $x_j$  an arbitrary constant independent of x.

We use the following notations :

 $\theta_j = (\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j).$ 

 $W = W(\phi_j, \ldots, \phi_N)$  is the classical Wronskian  $W = \det[(\partial_x^{j-1}\phi_i)_{i,j\in[1,\ldots,N]}]$ . We consider the matrix  $A = (a_{jk})_{j,k\in[1,\ldots,N]}$  defined by

$$a_{jk} = (-1)^j \exp\left[2(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)\right] \prod_{l \neq j} \left|\frac{\kappa_l + \kappa_k}{\kappa_l - \kappa_j}\right|.$$
 (30)

Then we have the following statement

#### Theorem 5.1

$$\det(I+A) = \frac{2^N (-1)^{\frac{N(N+1)}{2}} \exp(\sum_{j=1}^N \theta_j)}{\prod_{j=2}^N \prod_{i=1}^{j-1} (\kappa_j - \kappa_i)} W(\phi_1, \dots, \phi_N)$$
(31)

#### 5.2 Solutions to the KdV equation in terms of wronskians

From the previous subsection, we can give the following wronskian representation of the solutions to the KdV equation.

**Theorem 5.2** The function u defined by

$$u(x,t) = -2\partial_x^2 \left( \ln \left[ W(\phi_1, \dots, \phi_N) \right] \right), \tag{32}$$

where  $= W(\phi_1, \ldots, \phi_N) = \det[(\partial_x^{j-1}\phi_i)_{i,j\in[1,\ldots,N]}]$  is the wronskian of the functions  $\phi$  defined by  $\phi_j(x,t) = \sinh(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)$ ,  $\kappa_j$ ,  $x_j$  being real numbers, is a solution to the KdV equation (1).

It is relevant to note that we recover the result given by the Darboux dressing [64].

This realize the connection between the algebro-geometric approach and the Darboux dressing method.

**Remark 5.1** The choices of functions  $\phi_j$  are not unique. For example, we can choose :

 $\phi_j(x) = \cosh(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j), \text{ or } \phi_j(x) = \exp(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j), \text{ or } \phi_j(x) = \exp(-(\kappa_j x - 4\kappa_j^3 t + \kappa_j x_j)) \text{ or any combinations of these different last functions.}$ 

We can also choose the following functions :

 $\phi_j(x) = \sin(\kappa_j x + 4\kappa_j^3 t + \kappa_j x_j), \text{ or } \phi_j(x) = \cos(\kappa_j x + 4\kappa_j^3 t + \kappa_j x_j), \text{ or } \phi_j(x) = \exp(i(\kappa_j x + 4\kappa_j^3 t + \kappa_j x_j)), \text{ or } \phi_j(x) = \exp(-i(\kappa_j x + 4\kappa_j^3 t + \kappa_j x_j))) \text{ or any combinations of these different last functions.}$ 

## 6 Conclusion

In this paper, we succeed to construct different types of representations of the solutions to the KdV equation. First, it was essential to express the degenerate  $\theta$  function into an explicit Fredholm determinant. The second step was to get the transformation of the Fredholm determinant into a wronskian.

I have to mention a paper of Whitham [74] in connection with this work, and I would like to thank the referee about this information. The article [74] deals with equations as the KdV equation and the representation of solutions as sum of solitons, and also the relation of these solutions with Riemann theta functions in particular. It can be compared with the solutions given in the present work expressed in terms of Fredholm determinant and wronskians.

# References

- M.J. Ablowitz, H. Cornille, On solutions of the KdV equation, Phys. Lett., V. 72A, N. 4, 277-280, 1979
- [2] M. Adler, J. Moser, On a class of polynomials connected with the Kortewegde-Vries equation, Commun. Math. Phys., V. 61, 1-30, 1978
- [3] H. Airault, H.P. McKean, J. Moser, Rational and elliptic solutions of the KdV equation and a related many-body problem, Comm. Pure and Appl. Math., V. XXX, 2, 95-148, 1977
- [4] E.D. Belokolos, A.I. Bobenko, A.R. Its, V.Z. Enolskij and V.B. Matveev, Algebro-geometric approach to nonlinear integrable equations, Springer series in nonlinear dynamics, Springer Verlag, 1-360, 1994.
- [5] N. C Freeman, J.J.C. Nimmo Rational solutions of the KdV equation in wronskian form, Phys. Letters, V. 96 A, N. 9, 443-446, 1983
- [6] P. Gaillard, V.B. Matveev, Wronskian addition formula and its applications, Max-Planck-Institut für Mathematik, MPI 02-31, V. 161, 2002
- [7] P. Gaillard, A new family of deformations of Darboux-Pöschl-Teller potentials, Lett. Math. Phys., V. 68, 77-90, 2004
- [8] P. Gaillard, V.B. Matveev, New formulas for the eigenfunctions of the two-particle Calogero-Moser system, Lett. Math. Phys., V. 89, 1-12, 2009
- P. Gaillard, V.B. Matveev, Wronskian and Casorai determinant representations for Darboux-Pöschl-Teller potentials and their difference extensions, J. Phys A : Math. Theor., V. 42, 404409-1-16, 2009
- [10] P. Dubard, P. Gaillard, C. Klein, V. B. Matveev, On multi-rogue wave solutions of the NLS equation and positon solutions of the KdV equation, Eur. Phys. J. Spe. Top., V. 185, 247-258, 2010

- [11] P. Gaillard, Families of quasi-rational solutions of the NLS equation and multi-rogue waves, J. Phys. A : Meth. Theor., V. 44, 435204-1-15, 2011
- [12] P. Gaillard, Wronskian representation of solutions of the NLS equation and higher Peregrine breathers, J. Math. Sciences : Adv. Appl., V. 13, N. 2, 71-153, 2012
- [13] P. Gaillard, Degenerate determinant representation of solution of the NLS equation, higher Peregrine breathers and multi-rogue waves, J. Math. Phys., V. 54, 013504-1-32, 2013
- [14] P. Gaillard, Wronskian representation of solutions of NLS equation and seventh order rogue waves, J. Mod. Phys., V. 4, N. 4, 246-266, 2013
- [15] P. Gaillard, V.B. Matveev, Wronskian addition formula and Darboux-Pöschl-Teller potentials, J. Math., V. 2013, ID 645752, 1-10, 2013
- [16] P. Gaillard, Two-parameters determinant representation of seventh order rogue waves solutions of the NLS equation, J. Theor. Appl. Phys., V. 7, N. 45, 1-6, 2013
- [17] P. Gaillard, Deformations of third order Peregrine breather solutions of the NLS equation with four parameters, Phys. Rev. E, V. 88, 042903-1-9, 2013
- [18] P. Gaillard, Two parameters deformations of ninth Peregrine breather solution of the NLS equation and multi rogue waves, J. Math., V. 2013, 1-111, 2013
- [19] P. Gaillard, Six-parameters deformations of fourth order Peregrine breather solutions of the NLS equation, J. Math. Phys., V. 54, 073519-1-22, 2013
- [20] P. Gaillard, The fifth order Peregrine breather and its eight-parameters deformations solutions of the NLS equation, Commun. Theor. Phys., V. 61, 365-369, 2014
- [21] P. Gaillard, Ten parameters deformations of the sixth order Peregrine breather solutions of the NLS equation, Phys. Scripta, V. 89, 015004-1-7, 2014
- [22] P. Gaillard, Higher order Peregrine breathers, their deformations and multirogue waves, J. Of Phys. : Conf. Ser., V. 482, 012016-1-7, 2014
- [23] P. Gaillard, M. Gastineau, Eighteen parameter deformations of the Peregrine breather of order ten solutions of the NLS equation, Int. J. Mod. Phys. C, V. 26, N. 2, 1550016-1-14, 2014
- [24] P. Gaillard, Two parameters wronskian representation of solutions of nonlinear Schrödinger equation, eight Peregrine breather and multi-rogue waves, J. Math. Phys., V. 5, 093506-1-12, 2014

- [25] P. Gaillard, Hierarchy of solutions to the NLS equation and multi-rogue waves, J. Phys. : Conf. Ser., V. 574, 012031-1-5, 2015
- [26] P. Gaillard, Tenth Peregrine breather solution of the NLS, Ann. Phys., V. 355, 293-298, 2015
- [27] P. Gaillard, M. Gastineau, The Peregrine breather of order nine and its deformations with sixteen parameters solutions of the NLS equation Phys. Lett. A, V. 379, 1309-1313, 2015
- [28] P. Gaillard, Other 2N-2 parameters solutions to the NLS equation and 2N+1 highest amplitude of the modulus of the N-th order AP breather, J. Phys. A: Math. Theor., V. 48, 145203-1-23, 2015
- [29] P. Gaillard, Multi-parametric deformations of the Peregrine breather of order N solutions to the NLS equation and multi-rogue waves, Adv. Res., V. 4, N. 5, 346-364, 2015
- [30] P. Gaillard, Higher order Peregrine breathers solutions to the NLS equation, Jour. Phys. : Conf. Ser., V. 633, 012106-1-6, 2016
- [31] P. Gaillard, M. Gastineau Patterns of deformations of Peregrine breather of order 3 and 4, solutions to the NLS equation with multi-parameters, Journal of Theoretical and Applied Physics, V. 10,1-7, 2016
- [32] P. Gaillard, M. Gastineau Twenty parameters families of solutions to the NLS equation and the eleventh Peregrine breather, Commun. Theor. Phys, V. 65, N. 2, 136-144, 2016
- [33] P. Gaillard, Rational solutions to the KPI equation and multi rogue waves, Annals Of Physics, V. 367, 1-5, 2016
- [34] P. Gaillard, M. Gastineau Twenty two parameters deformations of the twelfth Peregrine breather solutions to the NLS equation, Adv. Res., V. 10, 83-89, 2016
- [35] P. Gaillard, Towards a classification of the quasi rational solutions to the NLS equation, Theor. And Math. Phys., V. 189, N. 1, 1440-1449, 2016
- [36] P. Gaillard, Fredholm and Wronskian representations of solutions to the KPI equation and multi-rogue waves, Jour. of Math. Phys., V. 57, 063505-1-13, doi: 10.1063/1.4953383, 2016
- [37] P. Gaillard, M. Gastineau Families of deformations of the thirteenth Peregrine breather solutions to the NLS equation depending on twenty four parameters, Jour. Of Bas. And Appl. Res. Int., V. 21, N. 3, 130-139, 2017
- [38] P. Gaillard, From Fredholm and Wronskian representations to rational solutions to the KPI equation depending on 2N2 parameters, Int. Jour. of Appl. Sci. And Math., V. 4, N. 3, 60-70, 2017

- [39] P. Gaillard, Families of rational solutions of order 5 to the KPI equation depending on 8 parameters, New Hor. in Math. Phys., V. 1, N. 1, 26-31, 2017
- [40] P. Gaillard, 6-th order rational solutions to the KPI Equation depending on 10 parameters, Jour. Of Bas. And Appl. Res. Int., V. 21, N. 2, 92-98, 2017
- [41] P. Gaillard, N-Order rational solutions to the Johnson equation depending on 2N - 2 parameters, Int. Jour. of Adv. Res. in Phys. Sci., V. 4, N. 9, 19-37, 2017
- [42] P. Gaillard, Families of rational solutions to the KPI equation of order 7 depending on 12 parameters, Int. Jour. of Adv. Res. in Phys. Sci., V. 4, N. 11, 24-30, 2017
- [43] P. Gaillard, Rational solutions to the Johnson equation and rogue waves, Int. Jour. of Inn. In Sci. and Math., V. 6, N. 1, 14-19, 2018
- [44] P. Gaillard, Multiparametric families of solutions of the KPI equation, the structure of their rational representations and multi-rogue waves, Theo. And Mat. Phys., V. 196, N. 2, 1174-1199, 2018
- [45] P. Gaillard, The Johnson equation, Fredholm and wronskian representations of solutions and the case of order three, Adv. In Math. Phys., V. 2018, 1-18, 2018
- [46] P. Gaillard, Families of Solutions of Order 5 to the Johnson Equation Depending on 8 Parameters, NHIMP, V. 2, N. 4, 53-61, 2018
- [47] P. Gaillard, Multiparametric families of solutions to the Johnson equation, J. Phys. : Conf. Series, V. 1141, 012102-1-10, 2018
- [48] P. Gaillard, Rational solutions to the Boussinesq equation, Fund. Jour. Of Math. And Appl., V., 109-112, 2019
- [49] P. Gaillard, Differential relations for the solutions to the NLS equation and their different representations, Comm. In Adv. Math. Sci., V. 2, N. 4, 1-4, 2019
- [50] P. Gaillard, Multi-parametric families of solutions of order N to the Boussinesq and KP equations and the degenerate rational case, UJMA, V. 3, N. 2, 44-52, 2020
- [51] P. Gaillard, The mKdV equation and multi-parameters rational solutions, Wave Motion, V. 100, 102667-1-9, 2021
- [52] C. S. Gardner, J. M. Green, M. D. Kruskall, R. M Miura, Method for solving the Korteweg-de Vries equation, Phys. Rev. Let., V. 19, N. 19, 1095-1097, 1967

- [53] R. Hirota, Exact solution of the Korteweg-de Vries equation for multiple collisions of solitons, Phys. Rev. Let., V. 27, N. 18, 1192-1194, 1971
- [54] A.R. Its, V.B. Matveev, Hill's operator with finitely many gaps, Funct. Anal. and Appl., V. 9, 69-70, 1975
- [55] M. Jaworski, J. Zagrodzinski, Positon and positon-like solution of the Korteweg-de Vries and Sine-Gordon equations, Chaos Solitons Fractals, V. 5, N. 12, 2229-2233, 1995.
- [56] D.J.Korteweg, G. de Vries, On the change of form of long wawes adwancing in a rectangular canal, and on a new type of long stationary waves, Phil. Mag., V. 39, 442-443, 1895
- [57] M. Kovalyov, M.H Ali Abadi, An explicit formula for a class of solutions of KdV equation, University of Alberta, Canada, 1997.
- [58] M. Kovalyov, On a class of solutions of KdV, Jour. Of Diff. Equ., V. 213, 1-80, 2005
- [59] P. D. Lax, Periodic solutions of the KdV equation, Comm. Pur. Applied Math., V. 28, 141-188, 1975
- [60] W.X. Ma, Y. You, Solving the KdV equation by its bilinear form wronskian solutions, Trans. Of The A. M. S., V. 357, N. 5, 1753-1778, 2004
- [61] W.X. Ma, Lump solutions to the KP equation, Physics Letters A, 1975-1978, 2015
- [62] V.B. Matveev, Darboux Transformation and explicit solutions of the Kadomtsev-Petviaschvily equation depending on functional parameters, Letters in Mathematical Physics, V. 3, 213-216, 1979
- [63] V.B. Matveev, Abelian functions and solitons, Preprint N. 313, Institute of Th. Physics, Univ. of Wroclow, 1976.
- [64] V.B. Matveev, Generalized Wronskian Formula for solutions of the KdV equation, Phys. Lett. A 166, 205-208, 1992
- [65] V.B. Matveev and M.A. Salle, Darboux transformations and solitons, Series in Nonlinear Dynamics, Springer-Verlag, Berlin, 1991
- [66] V.B. Matveev, Positon-positon and soliton-positon collisions : KdV case, Phys. Let. A, V. 166, 208-212, 1992.
- [67] V.B. Matveev, Asymptotics of mulpositon-solution  $\tau$  function of the Korteweg-de Vries equation and the supertransparency, J. Math. Phys., V. 35, N. 6, 1994.
- [68] V.B. Matveev, Positon : slowly decreasing analogues of solitons, Th. And Math. Phys. V. 131, N. 1, 2002.

- [69] C. Rasinariu, U. Sukhatme, A. Khare, Negaton and positon solutions of the KdV and the mKdV hierarchy, J. Phys. A, V. 29, 1803-1823, London, 1996.
- [70] J.F. Van Diejen, A.N. Kirillov, Formulas for q-Spherical functions using inverse scattering theory of reflectionless Jacobi operators, Commun. Math. Phys., V. 210, 335-369, 2000.
- [71] Van Diejen, J. F.: Integrability of difference Calogero-Moser systems, Jour. of Math. Physics, V. 35 N. 6, 2983-3004, 1999.
- [72] J.F. Van Diejen, A. N. Kirillov Determinantal formulas for zonal spherical functions on hyperboloids, Math. Ann., V. 319, 215-234, 2001.
- [73] H.D. Walquist, Bäcklund transformation of the potentials of the KdV type and the interaction of soliton with cnoïdal waves, Lecture Note, V. 515, 162-183, 1976
- [74] G.B. Whitham, Comments on Periodic Waves and Solitons, Journal of Applied Mathematics, V. 32, 353-366, 1984
- [75] A. Zabrodin, Finite-gap difference operators with elliptic coefficients and their spectral curves, in "Physical Combinatorics", M. Kashiwara and T. Miwa (eds), Birkhauser Ser. Progress in Mathematics, V. 191, ISBN 0-816-4175-0, Birkhauser-Boston, 301-317, 2000.
- [76] V.E. Zakharov, L.D. Faddeev, Korteweg-de Vries equation: A completely integrable Hamiltonian system, Func. Anal. and its Appl., V. 5, N. 4, 280-287, 1971