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# CRYSTALLINITY OF THE HOMOGENIZED ENERGY DENSITY OF PERIODIC LATTICE SYSTEMS 

ANTONIN CHAMBOLLE AND LEONARD KREUTZ


#### Abstract

We study the homogenized energy densities of periodic ferromagnetic Ising systems. We prove that, for finite range interactions, the homogenized energy density, identifying the effective limit, is crystalline, i.e. its Wulff crystal is a polytope, for which we can (exponentially) bound the number of vertices. This is achieved by deriving a dual representation of the energy density through a finite cell formula. This formula also allows easy numerical computations: we show a few experiments where we compute periodic patterns which minimize the anisotropy of the surface tension.


## 1. Introduction

The study of discrete interfacial energies has attracted widespread attention in the mathematical community over last decades, with applications in various contexts such as computer vision [7], crystallization problems [8, fracture mechanics [6, 18, 33, or statistical physics [38, 39]. To give examples, in computer vision the understanding of these energies allows to investigate functional correctness of segmentation algorithms [22]. Whereas for crystallization problems it gives fluctuation estimates on the macroscopic shape of the crystal cluster of ground state configurations [27, 31, 32, 37].

In this work, we consider energies defined on discrete periodic sets $\mathcal{L} \subset \mathbb{R}^{d}$ and corresponding Ising systems. We refer to [1, 12, 20, 21, 29, 34, 35, 36] for an abundant literature on the derivation of continuum limits of such systems and their effective behavior. More precisely, we consider $\mathcal{L}$ satisfying the following two conditions (see Figure 1 )
(i) (Discreteness) There exists $c>0$ such that $\operatorname{dist}(x, \mathcal{L} \backslash\{x\}) \geq c$ for all $x \in \mathcal{L}$;
(ii) (Periodicity) There exists $T \in \mathbb{N}$ such that for all $z \in \mathbb{Z}^{d}$, it holds that $\mathcal{L}+T z=\mathcal{L}$;

To each function $u: \mathcal{L} \rightarrow\{0,1\}$ and each $A \subset \mathbb{R}^{d}$ we associate an energy

$$
\begin{equation*}
E(u, A)=\sum_{i \in \mathcal{L} \cap A} \sum_{j \in \mathcal{L}} c_{i, j}(u(i)-u(j))^{+} \tag{1}
\end{equation*}
$$

where $(z)^{+}$denotes the positive part of $z \in \mathbb{R}, c_{i, j}: \mathcal{L} \times \mathcal{L} \rightarrow[0,+\infty)$ are $T$-periodic, that is $c_{i+T z, j+T z}=c_{i, j}$ for all $i, j \in \mathcal{L}$ and $z \in \mathbb{Z}^{d}$ and satisfy the following decay assumption
(iii) (Decay of interactions) For all $i \in \mathcal{L}$ there holds

$$
\sum_{j \in \mathcal{L}} c_{i, j}|i-j|<+\infty
$$

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Figure 1. An example of the set $\mathcal{L}$
Assuming conditions (i)-(iii) (and some additional coercivity assumption) ensures that the asymptotic behavior of (1) is well described (in a variational sense) by a continuum perimeter energy. More precisely, let us introduce a scaling parameter $\varepsilon>0$. We consider the scaled energies

$$
E_{\varepsilon}(u)=\sum_{i, j \in \varepsilon \mathcal{L}} \varepsilon^{d-1} c_{i, j}^{\varepsilon}(u(i)-u(j))^{+},
$$

where $c_{i, j}^{\varepsilon}=c_{i / \varepsilon, j / \varepsilon}$ and $u: \varepsilon \mathcal{L} \rightarrow\{0,1\}$. By identifying $u$ with its piecewise constant interpolation taking the value $u(i)$ on the Voronoi cell centered at $i \in \varepsilon \mathcal{L}$ we may regard the energies as defined on $L_{\text {loc }}^{1}\left(\mathbb{R}^{d},\{0,1\}\right)$. Integral representation results [2, 3, 17] then guarantee that the energies $E_{\varepsilon} \Gamma$-converge (see [11, 28] for an introduction to that subject) with respect to the $L_{\text {loc }}^{1}\left(\mathbb{R}^{d}\right)$-topology to a continuum energy of the form

$$
E_{0}(u)=\int_{\partial^{*}\{u=1\}} \varphi\left(\nu_{u}(x)\right) \mathrm{d} \mathcal{H}^{d-1} \quad u \in B V_{\text {loc }}\left(\mathbb{R}^{d} ;\{0,1\}\right) .
$$

Here, $B V_{\text {loc }}\left(\mathbb{R}^{d} ;\{0,1\}\right)$ denotes the space of functions with (locally) bounded variation and values in $\{0,1\}, \partial^{*}\{u=1\}$ denotes the reduced boundary of the level set $\{u=1\}, \nu_{u}(x)$ its measure theoretic normal at the point $x \in \partial^{*}\{u=1\}$, and $\mathcal{H}^{d-1}$ denotes the ( $d-1$ )-dimensional Hausdorff measure, see [4 for the precise definitions of these notions. The energy density $\varphi: \mathbb{R}^{d} \rightarrow[0,+\infty)$ can be recovered via the asymptotic cell formula

$$
\begin{equation*}
\varphi(\nu):=\lim _{\delta \rightarrow 0} \lim _{S \rightarrow+\infty} \frac{1}{S^{d-1}} \inf \left\{E\left(u, Q_{S}^{\nu}\right): u: \mathcal{L} \rightarrow\{0,1\}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}\right\} \tag{2}
\end{equation*}
$$

where

$$
u_{\nu}(x)= \begin{cases}1 & \langle x, \nu\rangle \geq 0 \\ 0 & \langle x, \nu\rangle<0\end{cases}
$$

Here, $Q_{S}^{\nu}$ is a cube with side-length $S$ orientation in direction $\nu \in \mathbb{S}^{d-1}$. In the case $\mathcal{L}=\mathbb{Z}^{2}$, $c_{i, j}=1$ if $|i-j|=1$ and $c_{i, j}=0$ otherwise, we have that $\varphi(\nu)=2\|\nu\|_{1}$, see Figure 2 .

The goal of this article is to investigate the energy density $\varphi$. In particular we show, that for finite interaction range $c_{i, j}$, that is there exists $R>0$ such that $c_{i, j}=0$ if $|i-j|>R$, then $\varphi$ is crystalline. This means that the solution to

$$
\min \left\{\int_{\partial^{*} A} \varphi\left(\nu_{A}(x)\right) \mathrm{d} \mathcal{H}^{d-1}:|A|=1\right\}
$$



Figure 2. The energy density in the case of nearest neighbor interactions on $\mathbb{Z}^{2}$
is a convex polyhedron 30 . The finite range of interaction is crucial. Indeed, example 2.8 shows that for infinite range interactions this is in general not true. In [15, 16] it is shown that, as the periodicity $T$ of the interactions tends to $+\infty$, it is possible to approximate any norm as surface energy density satisfying suitable growth conditions. We refer to [2] for a random setting where it is shown that an isotropic energy density (and thus non-crystalline) can be obtained in the limit.

The proof of the crystallinity in the case of finite range interactions relies on the following alternative representation result of the density, proven in Proposition 2.6. Namely, we prove that

$$
\begin{equation*}
\varphi(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} \tag{3}
\end{equation*}
$$

This representation formula is reminiscent of the representation formula of the energy density of integral functionals obtained via homogenization of $T$-periodic integral functionals in $W^{1, p}$ [14. To motivate this, consider the positively 1-homogeneous extensions of $E_{\varepsilon}$ defined by

$$
F_{\varepsilon}(u)=\sum_{i, j \in \varepsilon \mathcal{L}} \varepsilon^{d-1} c_{i, j}^{\varepsilon}(u(i)-u(j))^{+},
$$

for $u: \varepsilon \mathcal{L} \rightarrow \mathbb{R}$. The $\Gamma$-limit $F_{0}$ of the above sequence is clearly positively 1-homogeneous and convex as the sequence of functionals satisfies these properties. Thus, $F_{0}$ admits an integral representation of the form

$$
F_{0}(u)=\int f_{0}(\nabla u) \mathrm{d} x+\int f_{0}\left(\frac{d D_{s} u}{d\left|D_{s} u\right|}\right) \mathrm{d}\left|D_{s} u\right|
$$

where $f_{0}: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ is convex and positively 1-homogeneous, see 13 . (We like to stress however, that this integral representation for the spin energies considered above is not proven in the literature) Here, the important point is that the density of the singular part and the density of absolutely continuous parts agree. In the continuous setting, in [23, 26] it has been shown that for continuous and convex densities, that satisfy a coarea formula, the $\Gamma$-convergence of sets of finite perimeter or in the space of $B V$-functions is equivalent. Thus also in their setting, the densities agree. The density of the absolutely continuous part can be calculated via (3). This property eventually allows us to express $\varphi$ via (3) since the density of the absolutely continuous part can be calculated via (3) and the density of the singular part agrees with the energy density in (2), see Proposition 2.6. Using convex duality (see 40) and using (3) we show in Theorem 2.7 that $\varphi$ is crystalline, and estimate an upper bound on the number of extreme points of the corresponding Wulff shape. We would like to stress that (3) is not only a useful tool in our proof but it can be used also for computational purposes as it is a finite and not an asymptotic cell formula.

The paper is organized as follows. In Section 2 we describe the mathematical setting and state the main theorems of our paper. In Section 3 we prove Proposition 2.6, the alternate representation formula for $\varphi$. In Section 4 we show that, in the case of finite range interactions, the density $\varphi$ is always crystalline. We present some numerical simulations of our findings in the last chapter.

## 2. Setting of the problem and statement of the main result

2.1. Notation. We denote by $\mathcal{B}\left(\mathbb{R}^{d}\right)$ the collection of all Borel-Sets in $\mathbb{R}^{d}$. For every $A \subset \mathbb{R}^{d}$ we denote by $|A|$ its $d$-dimensional Lebesgue measure. Given $r>0$, we denote by $(A)_{r}:=\{x \in$ $\left.\mathbb{R}^{d}: \operatorname{dist}(x, A)<r\right\}$ the $r$-neighbourhood of $A$. Given $\tau \in \mathbb{R}^{d}$, we set $A+\tau:=\{x+\tau: x \in A\}$. The set $\mathbb{S}^{d-1}:=\left\{\nu \in \mathbb{R}^{d}:|\nu|=1\right\}$ is the set of all $d$-dimensional unit vectors. For $v, w \in \mathbb{R}^{d}$ we denote by $\langle v, w\rangle$ their scalar product. We denote by $\left\{e_{1}, \ldots, e_{d}\right\} \subset \mathbb{R}^{d}$ the standard orthonormal basis of $\mathbb{R}^{d}$. Given $C \subset \mathbb{R}^{d}$ convex, we denote by extreme $(C)$ its extreme points. Given $\rho>0$, we denote by $Q_{\rho}:=[-\rho / 2, \rho / 2)^{d}$ the half open cube centred in 0 with side-length $\rho$. For $\nu \in \mathbb{S}^{d-1}$, we set $Q_{\rho}^{\nu}:=R^{\nu} Q_{\rho}$, where $R^{\nu}$ is a rotation such that $R^{\nu} e_{d}=\nu$. Furthermore, given $x \in \mathbb{R}^{d}$ we set $Q_{\rho}^{\nu}(x):=x+Q_{\rho}^{\nu}\left(\right.$ resp. $\left.Q_{\rho}(x)=x+Q_{\rho}\right)$. Given $x \in \mathbb{R}^{d}$ and $r>0$ we denote by $B_{r}(x)$ the open ball with radius $r>0$ and center $x$. We denote by $\omega_{d}$ the volume of the unit ball in $\mathbb{R}^{d}$. Given $\nu \in \mathbb{S}^{d-1}$ we define

$$
u_{\nu}(x):= \begin{cases}1 & \text { if }\langle\nu, x\rangle \geq 0  \tag{4}\\ 0 & \text { otherwise }\end{cases}
$$

For $z \in \mathbb{R}$ we denote by $(z)^{+}:=\max \{z, 0\}$ the positive part of $z$.
2.2. Discrete energies and homogenized surface energy density. In this paragraph we define the discrete energies we want to consider and the main object of homogenized surface energy density.

Let $\mathcal{L} \subset \mathbb{R}^{d}$ satisfy the following two conditions:
(S1) (Discreteness) There exists $c>0$ such that for all $x \in \mathcal{L}$ there holds

$$
\operatorname{dist}(x, \mathcal{L} \backslash\{x\}) \geq c
$$

(S2) (Periodicity) There exists $T \in \mathbb{N}$ such that for all $z \in \mathbb{Z}^{d}$ there holds

$$
\mathcal{L}+T z=\mathcal{L} .
$$

Note that the two assumptions (S1) and (S2) include multi-lattices, such as the hexagonal closed packing in three dimensions, and bravais lattices, such as $\mathbb{Z}^{d}$, or the face centred cubic lattice in three dimensions.

We consider interaction coefficients $c_{i, j}: \mathcal{L} \times \mathcal{L} \rightarrow[0,+\infty)$ and the corresponding (localized) ferromagnetic spin energies of the form

$$
\begin{equation*}
E(u, A):=\sum_{i \in \mathcal{L} \cap A} \sum_{j \in \mathcal{L}} c_{i, j}(u(i)-u(j))^{+} \tag{5}
\end{equation*}
$$

where $u: \mathcal{L} \rightarrow \mathbb{R}$ and $A \in \mathcal{B}\left(\mathbb{R}^{d}\right)$. If $A=\mathbb{R}^{d}$ we omit the dependence on the set and write $E(u):=E\left(u, \mathbb{R}^{d}\right)$. We want to remark that we are considering interactions on the directed graph instead of the undirected graph.

We introduce the following three hypothesis on the interaction coefficients $c_{i, j}: \mathcal{L} \times \mathcal{L} \rightarrow$ $[0,+\infty)$ hold true:
(H1) (Periodicity) There holds

$$
c_{i+T z, j+T z}=c_{i, j}
$$

for all $i, j \in \mathcal{L}, z \in \mathbb{Z}^{d}$.
(H2) (Decay of Interactions) For all $i \in \mathcal{L}$ there holds

$$
\sum_{j \in \mathcal{L}} c_{i, j}|i-j|<+\infty
$$

(H3) (Finite Range Interactions) There exists $R>1$ such that

$$
c_{i, j}=0
$$

for all $i, j \in \mathcal{L}$ such that $|i-j| \geq R$.
It is obvious, that hypothesis (H3) implies hypothesis (H2). Note that, if (H1) and (H2) are satisfied then

$$
\max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j|=\max _{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j|<+\infty
$$

and for all $R>0$, there exists $C_{R}>0$ such that $C_{R} \rightarrow 0$ as $R \rightarrow+\infty$ and

$$
\max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\|i-j| \geq R}} c_{i, j}|i-j| \leq C_{R}
$$

Definition 2.1. Let $c_{i, j}$ satisfy (H1) and (H2). We then define the homogenized surface energy density $\varphi: \mathbb{R}^{d} \rightarrow[0,+\infty)$ as the convex positively homogeneous function of degree one such that for all $\nu \in \mathbb{S}^{d-1}$ we have

$$
\begin{equation*}
\varphi(\nu):=\lim _{\delta \rightarrow 0} \lim _{S \rightarrow+\infty} \frac{1}{S^{d-1}} \inf \left\{E\left(u, Q_{S}^{\nu}\right): u: \mathcal{L} \rightarrow\{0,1\}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}\right\} \tag{6}
\end{equation*}
$$

with $u_{\nu}$ defined in (4).
Remark 2.2. The definition above can be interpreted as a passage from discrete to continuum description as follows. Given $\varepsilon>0$, we consider the scaled energies

$$
E_{\varepsilon}(u):=\sum_{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} \varepsilon^{d-1} c_{i, j}(u(\varepsilon i)-u(\varepsilon j))^{+}
$$

where $u: \varepsilon \mathcal{L} \rightarrow\{0,1\}$. Upon identifying $u$ with its piecewise-constant interpolation, we can regard these energies to be defined on $L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{d}\right)$. We know that their $\Gamma$-limit is infinite outside the space $B V_{\text {loc }}\left(\mathbb{R}^{d}\right)$, where it has the form

$$
E_{0}(u):=\int_{\partial^{*}\{u=1\}} \varphi(\nu) \mathrm{d} \mathcal{H}^{d-1}
$$

with $\varphi$ given by (6) 1 Here, $\partial^{*}\{u=1\}$ denotes the reduced boundary of the set $\{u=1\}$ and $\mathcal{H}^{d-1}$ denotes the $(d-1)$-dimensional Hausdorff measure in $\mathbb{R}^{d}$ (cf. [4], Chapters 2.8 and 3.5 ).
Remark 2.3. Testing with $u_{\nu}$ in (6), using (S1) and (H2), it is easy to see that $\varphi(\nu) \leq C$ for all $\nu \in \mathbb{S}^{d-1}$. Therefore, due to the convexity and the fact that it is a positively one homogeneous function of degree one, $\varphi$ is continuous.

[^0]2.3. Statement of the main result. In this section we state the main result.

Definition 2.4. Given $\varphi: \mathbb{R}^{d} \rightarrow[0,+\infty)$ convex, positively homogeneous of degree one, we define the Wulff set of $\varphi$ by

$$
W_{\varphi}:=\left\{\zeta \in \mathbb{R}^{d}:\langle\zeta, \nu\rangle \leq \varphi(\nu) \text { for all } \nu \in \mathbb{S}^{d-1}\right\} .
$$

We say that $\varphi$ is crystalline, if $W_{\varphi}$ is a polytope.
Remark 2.5. From the definition of the Wulff set, it is clear that

$$
\varphi(\nu)=\sup _{\zeta \in W_{\varphi}}\langle\nu, \zeta\rangle
$$

Furthermore, one can check, that if $\varphi$ is crystalline, then the set $\{\varphi \leq 1\}$ is a polytope.
The next proposition shows that with our proof, we obtain a finite cell formula in order to calculate $\varphi$ instead of the asymptotic one, given in (6). We think that this result in itself is interesting, since it allows for calculations on finite size systems in order to compute $\varphi$ for general Ising-systems. This result is in spirit very close to [9, 13], where convex and positively 1-homogenous continuum energies are considered. In this case, the surface energy density and the energy with respect to the absolutely continuous part coincide.
Proposition 2.6. Let $c_{i, j}: \mathcal{L} \times \mathcal{L} \rightarrow[0,+\infty)$ be interaction coefficients such that $(\mathrm{H} 1)$ and (H2) hold true. Then

$$
\begin{equation*}
\varphi(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is T-periodic }\right\} \tag{7}
\end{equation*}
$$

Theorem 2.7. Let $c_{i, j}: \mathcal{L} \times \mathcal{L} \rightarrow[0,+\infty)$ be interaction coefficients such that (H1) and (H3) hold true. Then, the homogenized surface energy density $\varphi$ is crystalline. Denote by

$$
N:=\#\left\{(i, j) \in \mathcal{L} \cap Q_{T} \times \mathcal{L}: c_{i, j} \neq 0\right\}
$$

Then,

$$
\# \operatorname{extreme}\left(W_{\varphi}\right) \leq 2^{N}
$$

The next example shows that without assumption (H2) Theorem 2.7 fails to hold true.
Example 2.8. To construct the example we first observe that if $f: \mathbb{R}^{d} \rightarrow[0,+\infty)$ is crystalline, then $D^{2} f$ is a Radon-measure with support contained in finitely many hyper-planes. To see this, note that if $f: \mathbb{R}^{d} \rightarrow[0,+\infty)$ is crystalline, then there exist $\left\{\xi_{k}\right\}_{k=1}^{N} \subset \mathbb{R}^{d}$ such that

$$
f(\nu)=\max _{1 \leq k \leq N}\left\langle\xi_{k}, \nu\right\rangle .
$$

Here, we assume that $\left\{\xi_{k}\right\}_{k=1}^{N}$ is chosen minimal, i.e. $\xi_{k} \neq \lambda \xi_{j}$ for some $\lambda>0$ and for some $j \neq k$. This assumption ensures that all the vectors $\xi_{k}$ play an active role in the definition of $f$. Now, $D f \in B V_{\text {loc }}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right)$ is given by

$$
D f(\nu)=\sum_{k=1}^{N} \chi_{V_{k}}(\nu) \xi_{k}, \text { where } V_{k}=\left\{\nu \in \mathbb{R}^{d}: f(\nu)=\left\langle\xi_{k}, \nu\right\rangle\right\}
$$

Then

$$
D^{2} f(\nu)=\sum_{1 \leq k<j \leq N}\left(\xi_{k}-\xi_{j}\right) \otimes \nu_{k j} \mathcal{H}^{d-1}\left\lfloor_{\partial V_{k} \cap \partial V_{j}}\right.
$$

where $\partial V_{k} \cap \partial V_{j}=\left\{\nu \in \mathbb{R}^{d}: f(\nu)=\left\langle\xi_{k}, \nu\right\rangle=\left\langle\xi_{j}, \nu\right\rangle\right\}$. and $\nu_{k j} \in \mathbb{S}^{d-1}$ denotes the normal pointing towards the set $V_{k}$.

Let now $\mathcal{L}=\mathbb{Z}^{d}$ and $c_{i, j}=c_{j-i}=c_{i-j}$ (in the following denoted by $\left\{c_{\xi}\right\}_{\xi \in \mathbb{Z}^{d}}$ ) be such that $c_{\xi}>0$ for all $\xi \in \mathbb{Z}^{d}$ and

$$
\sum_{\xi \in \mathbb{Z}^{d}} c_{\xi}|\xi|<+\infty
$$

It is then obvious that $c_{i, j}$ is 1-periodic, (H1) and (H2) hold true, but (H3) is violated. Therefore, due to Proposition 2.6. we have

$$
\varphi(\nu)=\sum_{\xi \in \mathbb{Z}^{d}} c_{\xi}|\langle\xi, \nu\rangle|
$$

This is true, since the only admissible functions in the minimum problem given by Proposition 2.6 are $u_{\nu}(i)=\langle\nu, i\rangle+c$ for some $c \in \mathbb{R}$. We claim that

$$
D \varphi(\nu)=\sum_{\xi \in \mathbb{Z}^{d}} \operatorname{sign}(\langle\xi, \nu\rangle) c_{\xi} \xi
$$

where sign: $\mathbb{R} \rightarrow \mathbb{R}$ is defined by

$$
\operatorname{sign}(t)= \begin{cases}1 & t \geq 0 \\ -1 & t<0\end{cases}
$$

Therefore

$$
\begin{equation*}
D^{2} \varphi=2 \sum_{\xi \in \mathbb{Z}^{d}} c_{\xi} \xi \otimes \frac{\xi}{|\xi|} \mathcal{H}^{d-1}\left\lfloor_{\{\nu:\langle\xi, \nu\rangle=0\}}\right. \tag{8}
\end{equation*}
$$

This can be seen by approximation. Consider $\varphi_{R}: \mathbb{R}^{d} \rightarrow \mathbb{R}$ defined by

$$
\varphi_{R}(\nu)=\sum_{\substack{\xi \in \mathbb{Z}^{d} \\|\xi| \leq R}} c_{\xi}|\langle\xi, \nu\rangle|, \quad D \varphi_{R}(\nu)=\sum_{\substack{\xi \in \mathbb{Z}^{d} \\|\xi| \leq R}} \operatorname{sign}(\langle\xi, \nu\rangle) c_{\xi} \xi
$$

Then

$$
D^{2} \varphi_{R}=2 \sum_{\substack{\xi \in \mathbb{Z}^{d} \\|\xi| \leq R}} c_{\xi} \xi \otimes \frac{\xi}{|\xi|} \mathcal{H}^{d-1} L_{\{\nu:\langle\xi, \nu\rangle=0\}}
$$

Now

$$
\left|D^{2} \varphi_{R}\right|\left(B_{r}\right) \leq C r^{d-1} \sum_{\xi \in \mathbb{Z}^{d}} c_{\xi}|\xi|
$$

so the total variation of $D^{2} \varphi_{R}$ is (locally) uniformly bounded with limiting measure $D^{2} \varphi$ and $D \varphi_{R} \rightarrow D \varphi$ in $L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right)$, actually weakly in $B V$. Hence, (8) is shown. Now, since $c_{\xi}>0$ for all $\xi \in \mathbb{Z}^{d}$ it is obvious that $D^{2} \varphi$ is not supported on finitely many hyper-planes. Thus $\varphi$ cannot be crystalline.

Note that $\varphi$ is differentiable in totally irrational directions $\Delta^{2}$ A similar property is known to hold, in the continuous setting [5, [24, for homgenized surface tensions. We can state a result showing that this is still the case in the discrete setting, under assumptions (H1) and (H2).

Proposition 2.9. Under the assumptions of Proposition 2.6, $\varphi$ is differentiable in any totally irrational direction.

[^1]It is expected that it should be, "in general", not differentiable in the other directions, at least whenever the minimizers $u$ in (7) are constant on an infinite set, however the proofs in [5, 24] rely on ellipticity properties of the problem and are less easy to transfer to the discrete case. The proof of Proposition 2.9, which mimicks the proof in [24], is postponed to Section 5, and relies on the dual representation 100 introduced later on.

## 3. Proof of Proposition 2.6

This section is devoted to the proof of Proposition 2.6. We assume throughout this section that assumptions (S1), (S2) and (H1), (H2) are satisfied. The proof consists in showing that $\varphi$ can be characterized by several (equivalent) cell-formulas and therefore passing from (6) to (7).

First, we will state and prove some elementary properties of $E$ that will be used throughout this section.

Lemma 3.1. Let $A \in \mathcal{B}\left(\mathbb{R}^{d}\right)$.
(i) Let $|A|>1$ and $\nu \in \mathbb{R}^{d}$. Then

$$
E(\langle\nu, \cdot\rangle, A) \leq C\left|\nu \|(A)_{c}\right|
$$

(ii) Let $u: \mathcal{L} \rightarrow \mathbb{R}$. For all $t \in \mathbb{R}, \lambda>0$ there holds

$$
E(\lambda u+t, A)=\lambda E(u, A)
$$

and $u \mapsto E(u, A)$ is convex. In particular,

$$
E(u+v, A) \leq E(u, A)+E(v, A)
$$

for all $u, v: \mathcal{L} \rightarrow \mathbb{R}$.
(iii) Let $u: \mathcal{L} \rightarrow \mathbb{R}$ and $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$ be such that $A \subset B$. Then

$$
E(u, A) \leq E(u, B)
$$

(iv) Let $u: \mathcal{L} \rightarrow \mathbb{R}$ and $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$ be such that $A \cap B=\emptyset$. Then

$$
E(u, A \cup B)=E(u, A)+E(u, B) .
$$

(v) There holds

$$
\#\{i \in \mathcal{L} \cap A\} \leq \frac{1}{c^{d} \omega_{d}}\left|(A)_{c}\right|
$$

where $c$ is given by (S1).
(vi) Let $u: \mathcal{L} \rightarrow \mathbb{R}$. Then, the function $\tau \mapsto E(u(\cdot-\tau), A+\tau)$ is $T$-periodic.

Proof. We start by proving (ii)-(iv) in Step 1, then (v) and (vi) in Step 2 and Step 3 respectively, and finally (i) in Step 4.

Step 1.(Proof of (ii) - (iv)) All the three statements are a direct consequence of (5) and the fact that $c_{i, j} \geq 0$.

Step 2. (Proof of (v)) Note that

$$
\bigcup_{i \in \mathcal{L} \cap A} B_{c}(i) \subset(A)_{c}
$$

and therefore, due to (S1),

$$
c^{d} \omega_{d} \#\{i \in \mathcal{L} \cap A\}=\left|\bigcup_{i \in \mathcal{L} \cap A} B_{c}(i)\right| \leq\left|(A)_{c}\right|
$$

This is the claim.
Step 3.(Proof of (vi)) Let $u: \mathcal{L} \rightarrow \mathbb{R}$ and $z \in \mathbb{Z}^{d}$. Then, using (H1) and (S1),

$$
\begin{aligned}
E(u(\cdot-T z), A+T z) & =\sum_{i \in \mathcal{L} \cap(A+T z)} \sum_{j \in \mathcal{L}} c_{i, j}(u(i-T z)-u(j-T z))^{+} \\
& =\sum_{i \in \mathcal{L} \cap A} \sum_{j \in(\mathcal{L}+T z)} c_{i+T z, j+T z}(u(i)-u(j))^{+} \\
& =\sum_{i \in \mathcal{L} \cap A} \sum_{j \in \mathcal{L}} c_{i, j}(u(i)-u(j))^{+}=E(u, A) .
\end{aligned}
$$

Step 4.(Proof of (i)) Let $S>1$ and $\nu \in \mathbb{R}^{d}$, then, due to (v), the fact that $S>1$, (S1), (S2),(H1), and (H2), we have

$$
E(\langle\nu, \cdot\rangle, A)=\sum_{i \in \mathcal{L} \cap A} \sum_{j \in \mathcal{L}} c_{i, j}|\langle\nu, i-j\rangle| \leq|\nu| \#\{i \in \mathcal{L} \cap A\} \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \leq C|\nu|\left|(A)_{c}\right|
$$

Lemma 3.2. Let $S>0$ and $\nu \in \mathbb{S}^{n-1}$. Then

$$
\begin{aligned}
& \inf \left\{E\left(u, Q_{S}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}\right\} \\
= & \inf \left\{E\left(u, Q_{S}^{\nu}\right): u: \mathcal{L} \rightarrow\{0,1\}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}\right\} .
\end{aligned}
$$

Proof. We divide the proof into two steps. We first show that the left hand side is smaller or equal than the right hand side and then we show the right hand side is smaller or equal than the left hand side.

Step 1. (Proof of ' $\leq$ ') This inequality is clear, since the infimum on the left hand side is taken over a larger class of functions.

Step 2. (Proof of ' $\geq^{\prime}$ ) Let us take $u: \mathcal{L} \rightarrow \mathbb{R}$ such that $u=u_{\nu}$ on $\mathbb{R}^{d} \backslash Q_{(1-\delta) S}^{\nu}$. It suffices to construct $u_{0}: \mathcal{L} \rightarrow\{0,1\}$ such that $u_{0}=u_{\nu}$ on $\mathbb{R}^{d} \backslash Q_{(1-\delta) S}^{\nu}$ and

$$
\begin{equation*}
E\left(u_{0}, Q_{S}^{\nu}\right) \leq E\left(u, Q_{S}^{\nu}\right) \tag{9}
\end{equation*}
$$

We prove this inequality by induction on $N=\# \operatorname{codomain}(u)$. If $N=2$, then we must have $\operatorname{codomain}(u)=\{0,1\}$ and there is nothing to prove. Suppose that $\# \operatorname{codomain}(u)=N+1$, with $N \geq 2$. Since

$$
E\left((u \wedge 1) \vee 0, Q_{S}^{\nu}\right) \leq E\left(u, Q_{S}^{\nu}\right)
$$

we can suppose that $\operatorname{codomain}(u) \subset[0,1]$ We can write $\operatorname{codomain}(u)=\left\{a_{k}: k=0, \ldots, N\right\}$ with $0=a_{0}<a_{1}<\ldots, a_{N}=1$. Let

$$
u=\sum_{k=1}^{N} a_{k} \chi_{E_{k}}, \text { with } \bigcup_{k=0}^{N} E_{k}=\mathbb{R}^{d}
$$

$\{x:\langle\nu, x\rangle \geq 0\} \backslash Q_{(1-\delta) S}^{\nu} \subset E_{N}$, and $\{x:\langle\nu, x\rangle<0\} \backslash Q_{(1-\delta) S}^{\nu} \subset E_{0}$ so that $u(i)=u_{\nu}(i)$ on $\mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}$. We define

$$
\begin{equation*}
u^{\prime}=\sum_{k=2}^{N} a_{k} \chi_{E_{k}}, \text { and } u^{\prime \prime}=\sum_{k=2}^{N} a_{k} \chi_{E_{k}}+a_{2} \chi_{E_{1}} \tag{10}
\end{equation*}
$$

Note that both $u^{\prime}(i)=u^{\prime \prime}(i)=u^{\nu}(i)$ on $\mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}$. Furthermore, we have that

$$
E\left(u^{\prime}, Q_{T}^{\nu}\right)=E\left(u, Q_{T}^{\nu}\right)+a_{1}\left(\sum_{k=2}^{N} \sum_{i \in \mathcal{L} \cap E_{k}} \sum_{j \in \mathcal{L} \cap E_{1}} c_{i, j}-\sum_{i \in \mathcal{L} \cap E_{0}} \sum_{j \in \mathcal{L} \cap E_{1}} c_{i, j}\right)
$$

and

$$
E\left(u^{\prime \prime}, Q_{T}^{\nu}\right)=E\left(u, Q_{T}^{\nu}\right)+\left(a_{1}-a_{2}\right)\left(\sum_{k=2}^{N} \sum_{i \in \mathcal{L} \cap E_{k}} \sum_{j \in \mathcal{L} \cap E_{1}} c_{i, j}-\sum_{i \in \mathcal{L} \cap E_{0}} \sum_{j \in \mathcal{L} \cap E_{1}} c_{i, j}\right)
$$

Hence, either $E\left(u^{\prime}, Q_{S}^{\nu}\right) \leq E\left(u, Q_{S}^{\nu}\right)$ or $E\left(u^{\prime \prime}, Q_{S}^{\nu}\right) \leq E\left(u, Q_{S}^{\nu}\right)$. Since, by 10, $u^{\prime}, u^{\prime \prime}: \mathcal{L} \rightarrow$ \#codomain $(u) \backslash\left\{a_{1}\right\}$ inequality (9) is true by induction. This concludes the proof.

Let $\phi: \mathbb{R}^{d} \rightarrow[0,+\infty]$ be defined by

$$
\begin{equation*}
\phi(\nu)=\lim _{\delta \rightarrow 0} \lim _{S \rightarrow+\infty} \frac{1}{S^{d}} \inf \left\{E\left(u, Q_{S}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}\right\} \tag{11}
\end{equation*}
$$

$\phi_{\mathrm{per}}: \mathbb{R}^{d} \rightarrow[0,+\infty]$ is defined by

$$
\begin{equation*}
\phi_{\mathrm{per}}(\nu)=\lim _{k \rightarrow+\infty} \frac{1}{(k T)^{d}} \inf \left\{E\left(u, Q_{k T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is }(k T) \text {-periodic }\right\} \tag{12}
\end{equation*}
$$

The following lemma uses a standard cutoff-argument. However, due to the infinite range of interactions, the finite scale arguments need to be adapted.

Lemma 3.3. Let $\nu \in \mathbb{R}^{d}$. Then

$$
\phi_{\mathrm{per}}(\nu)=\phi(\nu) .
$$

Proof. In order to prove the Lemma, we show first $\phi_{\text {per }}(\nu) \leq \phi(\nu)$ and then the reverse inequality. In order to do so, we modify competitors of the respective cell formulas in order to obtain a competitor for the other formula.

Due to the one homogeneity of both functions, we may assume that $\nu \in \mathbb{S}^{d-1}$.
Step 1. (Proof of $\left.{ }^{\prime} \leq '\right)$ Since, in the definition of $\phi_{\text {per }}(\nu)$, (resp. $\left.\phi(\nu)\right)$, the limit exists $\square^{3}$ we can assume without loss of generality that $S=k T$ for some $k \in \mathbb{N}$ with $k$ large. Let $\delta>0$ and let $u_{k}^{\delta}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{k}^{\delta}(i)=\langle\nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) k T}$ and

$$
\begin{equation*}
E\left(u_{k}^{\delta}, Q_{k T}\right)=\inf \left\{E\left(u, Q_{k T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) k T}\right\} \tag{13}
\end{equation*}
$$

We assume that

$$
\begin{equation*}
\left\|u_{k}^{\delta}\right\|_{L^{\infty}\left(Q_{(1+\delta) k T}\right)} \leq 2 k T . \tag{14}
\end{equation*}
$$

If that were not true we perform the following construction with $\tilde{u}_{k}^{\delta}(i)=\left(u_{k}^{\delta}(i) \vee(-2 k T)\right) \wedge(2 k T)$. Note that still $\tilde{u}_{k}^{\delta}(i)=\langle\nu, i\rangle$ on $Q_{(1+\delta) k T}$ for $\delta$ small enough. We define $v_{k}^{\delta}: \mathcal{L} \rightarrow \mathbb{R}$ by setting

$$
\begin{equation*}
v_{k}^{\delta}\left(i_{0}+k T z\right)=u_{k}^{\delta}\left(i_{0}\right)+\langle\nu, k T z\rangle \text { if } i_{0} \in Q_{k T}, z \in \mathbb{Z}^{d} \tag{15}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
\left|v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right| \leq C k T+|i-j| \tag{16}
\end{equation*}
$$

[^2]To see this note for $i=i_{0}+T k z, j=j_{0}+T k z^{\prime}, i_{0}, j_{0} \in Q_{k T}, z, z^{\prime} \in \mathbb{Z}^{d}$, due to 15) and 14, we have

$$
\begin{aligned}
\left|v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right| & =\left|u_{k}^{\delta}\left(i_{0}\right)-u_{k}^{\delta}\left(j_{0}\right)+\left\langle\nu, k T\left(z-z^{\prime}\right)\right\rangle\right| \leq\left|u_{k}^{\delta}\left(i_{0}\right)\right|+\left|u_{k}^{\delta}\left(j_{0}\right)\right|+\left|k T\left(z-z^{\prime}\right)\right| \\
& \leq C k T+|i-j|+\left|i_{0}\right|+\left|j_{0}\right| \leq C k T+|i-j|
\end{aligned}
$$

Clearly, $v_{k}^{\delta}(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic. Let us check that

$$
\begin{equation*}
v_{k}^{\delta}(i)=v_{k}^{\delta}(i) \text { for } i \in Q_{(1+\delta) k T} \tag{17}
\end{equation*}
$$

This holds true for $i \in Q_{k T}$. It remains to be checked for $i \in Q_{(1+\delta) k T} \backslash Q_{k T}$. Let $i \in$ $Q_{(1+\delta) k T} \backslash Q_{k T}$, i.e. $i=k T z+i_{0}$, where $i_{0} \in Q_{k T}$ and $\|z\|_{\infty}=1$. Let $n \in\{1, \ldots, d\}$ be such that $\left|i_{n}\right|=\|i\|_{\infty}$. Furthermore, let us assume for now that $i_{n} \geq 0$ and therefore, $z_{n}=1$. Since $i \in Q_{(1+\delta) T k}$ we have $i_{n}<(1+\delta) T k / 2$. Hence

$$
\left(i_{0}\right)_{n}=i_{n}-k T z_{n}<(1+\delta) T k / 2-k T=(-1+\delta) T k / 2
$$

Hence $i_{0} \in \mathcal{L} \backslash Q_{(1-\delta) k T}$ (The case that $i_{n}<0$ is done analogously - note that $Q_{k T}$ is defined as the half-open cube centred in 0 . Hence, we need to make this distinction.). Therefore, by 15 and the definition of $u_{k}^{\delta}$ in $\mathcal{L} \backslash Q_{(1-\delta) k T}$,

$$
v_{k}^{\delta}(i)=v_{k}^{\delta}\left(i_{0}+k T z\right)=v_{k}^{\delta}\left(i_{0}\right)+\langle\nu, k T z\rangle=u_{k}^{\delta}\left(i_{0}\right)+\langle\nu, k T z\rangle=\left\langle\nu, i_{0}\right\rangle+\langle\nu, k T z\rangle=u_{k}^{\delta}(i)
$$

Hence, 17) holds true. Additionally,

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{k T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is }(k T) \text {-periodic }\right\} \leq E\left(v_{k}^{\delta}, Q_{k T}\right) \tag{18}
\end{equation*}
$$

We are finished with Step 1 if we prove

$$
\begin{equation*}
E\left(v_{k}^{\delta}, Q_{k T}\right) \leq E\left(u_{k}^{\delta}, Q_{k T}\right)+\frac{C_{k}}{\delta}(k T)^{d} \tag{19}
\end{equation*}
$$

where $C_{k} \rightarrow 0$ as $k \rightarrow+\infty$. In fact, using (13), 18), 19), dividing by $(k T)^{d}$, and letting $k \rightarrow+\infty$, we obtain the claim. Let us prove (19). We have, using (17),

$$
\begin{aligned}
E\left(v_{k}^{\delta}, Q_{k T}\right) & =\sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{j \in \mathcal{L}} c_{i, j}\left(v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right)^{+} \\
& =\sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{j \in \mathcal{L} \cap Q_{(1+\delta) k T}} c_{i, j}\left(v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right)^{+}+\sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) k T}} c_{i, j}\left(v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right)^{+} \\
& \leq E\left(u_{k}^{\delta}, Q_{k T}\right)+\sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta k T / 2}} c_{i, j}\left|v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right| .
\end{aligned}
$$

Hence, in order to show (19), it remains to prove

$$
\begin{equation*}
\sum_{i \in Q_{k T}} \sum_{\substack{j \in \mathcal{L} \\|i-j| \geq \delta k T / 2}} c_{i, j}\left|v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right| \leq \frac{C_{k}}{\delta}(k T)^{d} \tag{20}
\end{equation*}
$$

Using 16, (H2), and Lemma 3.1(v), we have

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta k T / 2}} c_{i, j}\left|v_{k}^{\delta}(i)-v_{k}^{\delta}(j)\right| & \leq \sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta k T / 2}} c_{i, j}(C k T+|i-j|) \\
& \leq\left(\frac{C}{\delta}+1\right) \sum_{i \in \mathcal{L} \cap Q_{k T}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta k T / 2}} c_{i, j}|i-j| \\
& \leq \frac{C}{\delta} \#\left(\mathcal{L} \cap Q_{k T}\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta k T / 2}} c_{i, j}|i-j| \leq \frac{C_{k}}{\delta}(k T)^{d}
\end{aligned}
$$

where $C_{k} \rightarrow 0$ as $k \rightarrow+\infty$. This yields 20) and therefore the claim of Step 1.
Step 2. (Proof of ' $\geq^{\prime}$ ) Let $u_{k}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{k}(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic and

$$
E\left(u_{k}, Q_{k T}\right)=\inf \left\{E\left(u, Q_{k T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is }(k T) \text {-periodic }\right\}
$$

Fix $S \in \mathbb{N}$ such that $S=m k T \gg k T$ for some $m \in \mathbb{N}, m \gg 1$. Since $u_{k}(\cdot)-\langle\nu, \cdot\rangle$ is ( $k T$ )-periodic, we have

$$
E\left(u_{k}, Q_{k T}\left(x_{0}\right)\right)=E\left(u_{k}, Q_{k T}\right) \text { for all } x_{0} \in k T \mathbb{Z}^{d}
$$

and therefore

$$
\begin{equation*}
E\left(u_{k}, Q_{S}\right)=\frac{S^{d}}{(k T)^{d}} E\left(u_{k}, Q_{k T}\right) \tag{21}
\end{equation*}
$$

There exists a constant $C_{k}>0$ (we omit the dependence on $T$ ) such that, due to the fact that $u_{k}(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic, there holds

$$
\begin{equation*}
\max _{i \in \mathcal{L}}\left|u_{k}(i)-\langle\nu, i\rangle\right|=\max _{i \in \mathcal{L} \cap Q_{k T}}\left|u_{k}(i)-\langle\nu, i\rangle\right| \leq C_{k} \tag{22}
\end{equation*}
$$

Let $\varphi_{S} \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ be a cut-off function such that

$$
\varphi_{S}(x)=1 \text { for } x \in Q_{(1-3 \delta) S}, \varphi_{S}(x)=0 \text { for } x \in \mathbb{R}^{d} \backslash Q_{(1-2 \delta) S}, \text { and }\left\|\nabla \varphi_{S}\right\|_{\infty} \leq C_{k}^{-1}
$$

Define $u_{S}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
u_{S}(i)=\varphi_{S}(i) u_{k}(i)+\left(1-\varphi_{S}(i)\right)\langle\nu, i\rangle .
$$

Then, $u_{S}(i)=\langle\nu, i\rangle$ for $i \in \mathcal{L} \backslash Q_{(1-\delta) S}$ and therefore

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}\right\} \leq E\left(u_{S}, Q_{S}\right) \tag{23}
\end{equation*}
$$

For all $i, j \in \mathcal{L}$ there holds

$$
u_{S}(i)-u_{S}(j)=\varphi_{S}(i)\left(u_{k}(i)-u_{k}(j)\right)+\left(1-\varphi_{S}(i)\right)\langle\nu, i-j\rangle+\left(\varphi_{S}(i)-\varphi_{S}(j)\right)\left(u_{k}(i)-\langle\nu, i\rangle\right)
$$

which, together with 22 , implies for all $i, j \in \mathcal{L}$

$$
\begin{align*}
\left(u_{S}(i)-u_{S}(j)\right)^{+} & \leq\left(u_{k}(i)-u_{k}(j)\right)^{+}+|i-j|+C_{k}^{-1}\left|u_{k}(i)-\langle\nu, i\rangle \| i-j\right| \\
& \leq\left(u_{k}(i)-u_{k}(j)\right)^{+}+C|i-j| \tag{24}
\end{align*}
$$

and

$$
\begin{equation*}
\left(u_{S}(i)-u_{S}(j)\right)^{+}=\left(u_{k}(i)-u_{k}(j)\right)^{+} \text {for all } i, j \in Q_{(1-3 \delta) S} \tag{25}
\end{equation*}
$$

Using (21), (24), and 25), we obtain

$$
\begin{align*}
E\left(u_{S}, Q_{S}\right) \leq & \sum_{i \in \mathcal{L} \cap Q_{S}} \sum_{j \in \mathcal{L}} c_{i, j}\left(u_{k}(i)-u_{k}(j)\right)^{+}+C \sum_{i \in \mathcal{L} \cap Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L} \backslash Q_{(1-3 \delta) S}} c_{i, j}|i-j| \\
& +C \sum_{i \in \mathcal{L} \cap Q_{S} \backslash Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
= & \frac{S^{d}}{(k T)^{d}} E\left(u_{k}, Q_{k T}\right)+C \sum_{i \in \mathcal{L} \cap Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L} \backslash Q_{(1-3 \delta) S}} c_{i, j}|i-j|  \tag{26}\\
& +C \sum_{i \in \mathcal{L} \cap Q_{S} \backslash Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| .
\end{align*}
$$

We show that

$$
\begin{equation*}
\sum_{i \in \mathcal{L} \cap Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L} \backslash Q_{(1-3 \delta) S}} c_{i, j}|i-j| \leq C_{S} S^{d} \tag{27}
\end{equation*}
$$

where $C_{S} \rightarrow 0$ as $S \rightarrow+\infty$. Furthermore, we show that

$$
\begin{equation*}
\sum_{i \in \mathcal{L} \cap Q_{S} \backslash Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \leq C \delta S^{d} \tag{28}
\end{equation*}
$$

Note that from (27) and (28) we obtain the claim of Step 2 by using (23), 26), dividing by $S^{d}$, letting first $S \rightarrow+\infty$, then $k \rightarrow+\infty$ and lastly $\delta \rightarrow 0$.

We first prove (27). Note that, for $S$ big enough, due to (H2) and Lemma 3.1.v), we have

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L} \backslash Q_{(1-3 \delta) S}} c_{i, j}|i-j| & \leq \sum_{i \in \mathcal{L} \cap Q_{(1-6 \delta) S}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S}} c_{i, j}|i-j| \\
& \leq \#\left(\mathcal{L} \cap Q_{S}\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S}} c_{i, j}|i-j| \\
& \leq C_{S} S^{d}
\end{aligned}
$$

where $C_{S} \rightarrow 0$ as $S \rightarrow+\infty$. Next, we show 28). Using (H2), and Lemma 3.1(v), we obtain

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q S \backslash Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| & \leq \sum_{i \in \mathcal{L} \cap Q_{S} \backslash Q_{(1-6 \delta) S}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
& \leq \#\left(\mathcal{L} \cap Q_{S} \backslash Q_{(1-6 \delta) S}\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
& \leq C \delta S^{d} .
\end{aligned}
$$

This is 28 and hence the claim of Step 2.
The next Lemma shows that, using periodic boundary conditions, one can reduce from an asymptotic cell formula to a finite cell formula.

Lemma 3.4. Let $\nu \in \mathbb{S}^{1}$. For all $k \in \mathbb{N}$ there holds

$$
\begin{align*}
& \frac{1}{(k T)^{d}} \inf \left\{E\left(u, Q_{k T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is }(k T) \text {-periodic }\right\}  \tag{29}\\
& =\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is T-periodic }\right\}
\end{align*}
$$

In particular

$$
\begin{equation*}
\phi_{\text {per }}(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is T-periodic }\right\} \tag{30}
\end{equation*}
$$

Proof. We split the proof into two steps by first observing the (obvious) inequality that the right hand side in 29 is less than or equal to the left hand side. Then, we prove the converse inequality by using a superposition argument.

Step 1. (Proof of ' $\leq$ ') Let $u: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u(\cdot)-\langle\nu, \cdot\rangle$ is $T$-periodic and

$$
\begin{equation*}
E\left(u, Q_{T}\right)=\inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} \tag{31}
\end{equation*}
$$

Then, since $u(\cdot)-\langle\nu, \cdot\rangle$ is $T$-periodic, we have $E\left(u, Q_{T}(z)\right)=E\left(u, Q_{T}\right)$ for all $z \in \mathbb{Z}^{d}$. For $m \in\{0, \ldots, k-1\}^{d}$ set

$$
z_{m}^{k}=T\left(m-\frac{k-1}{2} \overrightarrow{1}\right)
$$

where $\overrightarrow{1}_{n}=1$ for all $n=1, \ldots, d$. It is easy to check that $Q_{T}\left(z_{m}^{k}\right) \subset Q_{k T}$ for all $m \in$ $\{0, \ldots, k-1\}^{d}$. Hence, using 31, we obtain

$$
\begin{aligned}
E\left(u, Q_{k T}\right) & \leq \sum_{m \in\{0, \ldots, k-1\}^{d}} E\left(u, Q_{T}\left(z_{m}^{k}\right)\right)=k^{d} E\left(u, Q_{T}\right) \\
& =\inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\}
\end{aligned}
$$

Noting that $u(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic we obtain the desired inequality.
Step 2. (Proof of ' $\geq$ ') Let $u_{k}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic. Define $u_{T}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
u_{T}(i)=\frac{1}{k^{d}} \sum_{z \in\{0, \ldots, k-1\}^{d}} u_{k}(i+T z) .
$$

Let us first check that $u_{T}(\cdot)-\langle\nu, \cdot\rangle$ is $T$-periodic. By definition, for $n \in\{1, \ldots, d\}$ there holds

$$
u_{T}\left(i+T e_{n}\right)-\nu\left(i+T e_{n}\right)=\frac{1}{k^{d}} \sum_{z \in\{0, \ldots, k-1\}^{d}} u_{k}\left(i+T z+T e_{n}\right)-\left\langle\nu, i+T e_{n}\right\rangle
$$

We now split the sum to obtain

$$
\begin{align*}
& u_{T}\left(i+T e_{n}\right)-\nu\left(i+T e_{n}\right) \\
= & \frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n} \in\{0, \ldots, k-2\}}} u_{k}\left(i+T z+T e_{n}\right)+\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=k-1}}\left(u_{k}\left(i+T z+T e_{n}\right)-k\left\langle\nu, T e_{n}\right\rangle\right)-\langle\nu, i\rangle . \tag{32}
\end{align*}
$$

For the first term in the sum, shifting the indices, we obtain

$$
\begin{equation*}
\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\ z_{n} \in\{0, \ldots, k-2\}}} u_{k}\left(i+T z+T e_{n}\right)=\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\ z_{n} \in\{1, \ldots, k-1\}}} u_{k}(i+T z) . \tag{33}
\end{equation*}
$$

For the second term in the sum, we obtain, due to the fact that $u(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic,

$$
\begin{aligned}
& \frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=k-1}}\left(u_{k}\left(i+T z+T e_{n}\right)-k\left\langle\nu, T e_{n}\right\rangle\right) \\
&= \frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=k-1}}\left(u_{k}\left(i+T\left(z^{\prime}, 0\right)+k T e_{n}\right)-\left\langle\nu, i+T\left(z^{\prime}, 0\right)+k T e_{n}\right\rangle+\left\langle\nu, i+T\left(z^{\prime}, 0\right)\right\rangle\right) \\
&= \frac{1}{k^{d}} \sum_{\substack{ \\
z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=k-1}} u_{k}\left(i+T\left(z^{\prime}, 0\right)\right)=\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=0}} u_{k}(i+T z) . \\
&
\end{aligned}
$$

Using this together with (32) and (33), we get

$$
\begin{aligned}
u_{T}\left(i+T e_{n}\right)-\nu\left(i+T e_{n}\right) & =\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n} \in\{1, \ldots, k-1\}}} u_{k}(i+T z)+\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d} \\
z_{n}=0}} u_{k}(i+T z)-\langle\nu, i\rangle \\
& =\frac{1}{k^{d}} \sum_{\substack{z^{\prime} \in\{0, \ldots, k-1\}^{d}}} u_{k}(i+T z)-\langle\nu, i\rangle=u_{T}(i)-\langle\nu, i\rangle .
\end{aligned}
$$

This shows that in fact $u_{T}(\cdot)-\langle\nu, \cdot\rangle$ is $T$-periodic. Hence

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} \leq E\left(u_{T}, Q_{T}\right) \tag{34}
\end{equation*}
$$

Now, using the convexity of $E$ (cf. Lemma 3.1(ii)), we obtain

$$
E\left(u_{T}, Q_{T}\right) \leq \frac{1}{k^{d}} \sum_{z \in\{0, \ldots, k-1\}^{d}} E\left(u_{k}(\cdot+T z), Q_{T}\right)=\frac{1}{k^{d}} E\left(u_{k}, Q_{k T}\right)
$$

Note that $u_{k}: \mathcal{L} \rightarrow \mathbb{R}$ is such that $u_{k}(\cdot)-\langle\nu, \cdot\rangle$ is $(k T)$-periodic and arbitrary. This together with (34) yields the claim. Equation (30) follows from Step 1, Step 2, and the definition of $\phi_{\text {per }}$, see (12).

Let $\psi: \mathbb{R}^{d} \rightarrow[0,+\infty]$ be defined as the positively homogeneous function of degree one that for $\nu \in \mathbb{S}^{d-1}$ is defined by

$$
\begin{equation*}
\psi(\nu)=\lim _{\delta \rightarrow 0} \lim _{S \rightarrow+\infty} \frac{1}{S^{d}} \inf \left\{E\left(u, Q_{S}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}\right\} \tag{35}
\end{equation*}
$$

Lemma 3.5. $\psi: \mathbb{R}^{d} \rightarrow[0,+\infty]$ satisfies the following properties:
(i) There exists $C>0$ such that $\psi(\nu) \leq C|\nu|$ for all $\nu \in \mathbb{R}^{d}$,
(ii) $\psi$ is a continuous function.

Proof. We divide the proof into two steps. We first prove (i) and then (ii). Throughout the proofs let $1 \ll S$.

Step 1.(Proof of (i)) Let $\nu \in \mathbb{S}^{d-1}$ it suffices to prove

$$
\psi(\nu) \leq C
$$

The general case then follows by one-homogeneity. In order to prove (i) we insert $u(i)=\langle\nu, i\rangle$ for all $i \in \mathcal{L}$ as a competitor in the cell formula. Using Lemma 3.1(i), we then have

$$
E\left(u, Q_{S}^{\nu}\right)=E\left(\langle\nu, \cdot\rangle, Q_{S}^{\nu}\right) \leq C\left|\nu \|\left(Q_{S}\right)_{c}\right| \leq C S^{d}
$$

Dividing by $S^{d}$ and letting $S \rightarrow+\infty$ yields the claim.

Step 2.(Proof of (ii)) Due to the one-homogeneity, it suffices to consider the case where $\nu_{1}, \nu_{2} \in \mathbb{S}^{d-1}$. Let $\eta>0$ and $\nu_{1}, \nu_{2} \in \mathbb{S}^{d-1}$ be such that $\left|\nu_{1}-\nu_{2}\right| \leq \eta$. Our goal is to prove that there exists $C>0$ independent of $\nu_{1}$ and $\nu_{2}$ such that

$$
\begin{equation*}
\left|\psi\left(\nu_{1}\right)-\psi\left(\nu_{2}\right)\right| \leq C \eta \tag{36}
\end{equation*}
$$

We only prove

$$
\begin{equation*}
\psi\left(\nu_{1}\right)-\psi\left(\nu_{2}\right) \leq C \eta \tag{37}
\end{equation*}
$$

since then (36) follows by exchanging $\nu_{1}$ and $\nu_{2}$ in 37. To this end let $\delta>0$ small enough, $S>0$ big enough, $u_{1}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{1}(i)=\left\langle\nu_{1}, i\right\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S}^{\nu}$ and

$$
\begin{equation*}
\frac{1}{S^{d}} E\left(u_{1}, Q_{S}^{\nu_{1}}\right) \leq \psi\left(\nu_{1}\right)+\eta \tag{38}
\end{equation*}
$$

We assume that

$$
\begin{equation*}
\left\|u_{1}\right\|_{L^{\infty}\left(Q_{S}^{\nu_{1}}\right)} \leq S \tag{39}
\end{equation*}
$$

If this were not the case, we consider

$$
\tilde{u}_{1}(i)= \begin{cases}\left(u_{1}(i) \wedge S\right) \vee(-S) & i \in Q_{2 S}^{\nu_{1}} \\ u_{1}(i) & \text { otherwise }\end{cases}
$$

Note that for $i, j \in Q_{2 S}^{\nu_{1}}$, due to truncation, $\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+} \leq\left(u_{1}(i)-u_{1}(j)\right)^{+}$, whereas in general there holds $\left|\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right| \leq C S+|i-j|$. From this, using Lemma 3.1(v) and (H2), we deduce

$$
\begin{aligned}
E\left(\tilde{u}_{1}, Q_{S}^{\nu_{1}}\right) & =\sum_{i \in \mathcal{L} \cap Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \cap Q_{2 S}^{\nu_{1}}} c_{i, j}\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+}+\sum_{i \in \mathcal{L} \cap Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \backslash Q_{2 S}^{\nu_{1}}} c_{i, j}\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+} \\
& \leq \sum_{i \in \mathcal{L} \cap Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \cap Q_{2 S}^{\nu_{1}}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+}+C \sum_{i \in \mathcal{L} \cap Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
& \leq E\left(u_{1}, Q_{S}^{\nu_{1}}\right)+C \#\left(\mathcal{L} \cap Q_{S}^{\nu_{1}}\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq S / 2}} c_{i, j}|i-j| \leq E\left(u_{1}, Q_{S}^{\nu_{1}}\right)+C_{S} S^{d}
\end{aligned}
$$

where $C_{S} \rightarrow 0$ as $S \rightarrow \infty$. In particular $C_{S} \leq \eta$ for $S$ big enough. Hence, we can assume (39). There exists $C>0$ such that for $\tilde{S}=(1+C \eta) S$ there holds $Q_{(1-\delta) \tilde{S}}^{\nu_{2}} \supset Q_{(1+\delta) S}^{\nu_{1}}$. We now define $u_{2}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
\begin{equation*}
u_{2}(i)=\left\langle\nu_{2}-\nu_{1}, i\right\rangle+u_{1}(i) \tag{40}
\end{equation*}
$$

First, note that $u_{2}(i)=\left\langle\nu_{2}, i\right\rangle$ for all $i \in \mathcal{L} \backslash Q_{(1-\delta) \tilde{S}}^{\nu_{2}}$ and therefore

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{\tilde{S}}^{\nu_{2}}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\left\langle\nu_{2}, i\right\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) \tilde{S}}^{\nu_{2}}\right\} \leq E\left(u_{2}, Q_{\tilde{S}}^{\nu_{2}}\right) \tag{41}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
E\left(u_{2}, Q_{\tilde{S}}^{\nu_{2}}\right) \leq E\left(u_{1} Q_{S}^{\nu_{1}}\right)+\frac{C_{S}}{\delta} S^{d}+C \eta S^{d}+C \delta S^{d} \tag{42}
\end{equation*}
$$

where $C_{S} \rightarrow 0$ as $S \rightarrow+\infty$. We postpone the proof of (42) and show first how it implies (37). Dividing (42) by $\tilde{S}^{d}$, letting $\tilde{S}$ (therefore also $S$ ) tend to $+\infty, \delta \rightarrow 0$, and using (41) as well as (38), we get

$$
\psi\left(\nu_{2}\right) \leq \psi\left(\nu_{1}\right)+C \eta \leq \psi\left(\nu_{1}\right)+C \eta
$$

This is (37). We now prove (42). Due to Lemma 3.1(ii), there holds

$$
\begin{equation*}
E\left(u_{2}, Q_{\tilde{S}}^{\nu_{2}}\right) \leq E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}}\right)+E\left(\left\langle\nu_{2}-\nu_{1}, \cdot\right\rangle, Q_{\tilde{S}}^{\nu_{1}}\right) \tag{43}
\end{equation*}
$$

Now, due to Lemma 3.1 (i) and the fact that $\tilde{S} \leq 2 S$, there holds

$$
\begin{equation*}
E\left(\left\langle\nu_{2}-\nu_{1}, \cdot\right\rangle, Q_{\tilde{S}}^{\nu_{1}}\right) \leq C\left|\nu_{2}-\nu_{1}\right| S^{d} \leq C \eta S^{d} \tag{44}
\end{equation*}
$$

Next, we prove

$$
\begin{equation*}
E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}}\right) \leq E\left(u_{1}, Q_{S}^{\nu_{1}}\right)+C \delta S^{d}+\frac{C_{S}}{\delta} S^{d} \tag{45}
\end{equation*}
$$

where $C_{S} \rightarrow 0$ as $S \rightarrow+\infty$. We use Lemma 3.1(iv), to obtain

$$
E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}}\right)=E\left(u_{1}, Q_{S}^{\nu_{1}}\right)+E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}\right) .
$$

In order to prove 45 it suffices to prove

$$
\begin{equation*}
E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}\right) \leq C \eta S^{d}+\frac{C_{S}}{\delta} S^{d} \tag{46}
\end{equation*}
$$

To see this we write

$$
\begin{align*}
E\left(u_{1}, Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}\right)= & \sum_{i \in \mathcal{L} \cap Q_{\tilde{S}_{2}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \cap Q_{(1-\delta) S}^{\nu_{1}}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+} \\
& +\sum_{i \in \mathcal{L} \cap Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu_{1}}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+} . \tag{47}
\end{align*}
$$

To estimate the first term, note that due to (39), we have $\left|u_{1}(i)-u_{1}(j)\right| \leq C S+|i-j|$, and therefore, up to changing $C$, using (H2), and Lemma 3.1(iv), we get

$$
\begin{align*}
\sum_{i \in \mathcal{L} \cap Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \cap Q_{(1-\delta) S}^{\nu_{1}}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+} & \leq \frac{C}{\delta} \sum_{i \in \mathcal{L} \cap Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S / 2}} c_{i, j}|i-j|  \tag{48}\\
& \leq \frac{C}{\delta} \#\left(\mathcal{L} \cap Q_{\tilde{S}_{2}}^{\nu_{2}}\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S / 2}} c_{i, j}|i-j| \leq \frac{C_{S}}{\delta} S^{d}
\end{align*}
$$

To estimate the first term, we use the fact that $u_{1}(i)=\left\langle\nu_{1}, i\right\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S}^{\nu_{1}}$, and Lemma 3.1 (i), to obtain
$\sum_{i \in \mathcal{L} \cap Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}} \sum_{j \in \mathcal{L} \backslash Q_{(1-\delta) S}^{\nu_{1}}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+} \leq E\left(\left\langle\nu_{1}, \cdot\right\rangle, Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}\right) \leq C\left|\nu_{1} \|\left(Q_{\tilde{S}}^{\nu_{2}} \backslash Q_{S}^{\nu_{1}}\right)_{c}\right| \leq C \eta S^{d}$.
This together with (47) and (48) implies (46) which in turn, together with (43) and (44) implies 42 and therefore the conclusion of Step 2.
Lemma 3.6. $\phi: \mathbb{R}^{d} \rightarrow[0,+\infty]$ satisfies the following properties:
(i) There exists $C>0$ such that $\phi(\nu) \leq C|\nu|$ for all $\nu \in \mathbb{R}^{d}$,
(ii) $\phi$ is a positively homogeneous function of degree one,
(iii) $\phi$ is a continuous function.

Proof. We divide the proof into two steps. We first prove (i) and then (ii). Throughout the proofs let $1 \ll S$.

Step 1. (Proof of (i) and (ii)) In order to prove (i) we insert $u(i)=\langle\nu, i\rangle$ for all $i \in \mathcal{L}$ as a competitor in the cell formula. Using Lemma 3.1(i), we then have

$$
E\left(u, Q_{S}^{\nu}\right)=E\left(\langle\nu, \cdot\rangle, Q_{S}^{\nu}\right) \leq C\left|\nu \|\left(Q_{S}\right)_{c}\right| \leq C S^{d}
$$

Dividing by $S^{d}$ and letting $S \rightarrow+\infty$ yields the claim. (ii) follows by using Lemma (3.1)(ii) to obtain $E\left(\lambda u, Q_{S}\right)=\lambda E\left(u, Q_{S}\right)$ for all $\lambda>0$ and by noting that, given $\nu \in \mathbb{R}^{d}$, if $u: \mathcal{L} \rightarrow \mathbb{R}$ satisfies $u(i)=\langle\nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S}$, then $\lambda u(i)=\langle\lambda \nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S}$. Employing this in 11) it is easy to see that $\phi$ is pos. one homogeneous function.

Step 2.(Proof of (iii)) In order to prove (ii), let $\nu_{1}, \nu_{2} \in \mathbb{R}^{d}$. We prove that

$$
\begin{equation*}
\left|\phi\left(\nu_{1}\right)-\phi\left(\nu_{2}\right)\right| \leq C\left|\nu_{1}-\nu_{2}\right| . \tag{49}
\end{equation*}
$$

Here, we only prove

$$
\begin{equation*}
\phi\left(\nu_{2}\right)-\phi\left(\nu_{1}\right) \leq C\left|\nu_{1}-\nu_{2}\right| \tag{50}
\end{equation*}
$$

since then follows by exchanging $\nu_{1}$ and $\nu_{2}$. To this end let $u: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{1}(i)=\left\langle\nu_{1}, i\right\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S}$. We then define $u_{2}: \mathcal{L} \rightarrow \mathbb{R}$ by $u_{2}(i)=u_{1}(i)+\left\langle\nu_{2}-\nu_{1}, i\right\rangle$. Clearly,

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S}\right\} \leq E\left(u_{2}, Q_{S}\right) \tag{51}
\end{equation*}
$$

Now, due to Lemma 3.1(i) and (ii),

$$
\begin{aligned}
E\left(u_{2}, Q_{S}\right) & =E\left(u_{1}+\left\langle\nu_{2}-\nu_{1}, \cdot\right\rangle, Q_{S}\right) \leq E\left(u_{1}, Q_{S}\right)+E\left(\left\langle\nu_{2}-\nu_{1}, \cdot\right\rangle, Q_{S}\right) \\
& \leq E\left(u_{1}, Q_{S}\right)+C\left|\nu_{1}-\nu_{2}\right|\left|\left(Q_{S}\right)_{c}\right| \leq E\left(u_{1}, Q_{S}\right)+C\left|\nu_{1}-\nu_{2}\right| S^{d}
\end{aligned}
$$

Using (51), noting that $u_{1}$ is arbitrary, dividing by $S^{d}$, letting first $S$ tend to $+\infty$, and then $\delta \rightarrow 0$, we obtain (50).

The next Lemma shows that the asymptotic cell-formula describing the surface energy density is equal to the asymptotic cell-formula with affine boundary conditions.
Lemma 3.7. Let $\nu \in \mathbb{R}^{d}$. Then

$$
\psi(\nu)=\varphi(\nu)
$$

Proof. Due the fact that both $\psi$ and $\phi$ are positively homogeneous functions of degree one, it suffices to consider the case where $\nu \in \mathbb{S}^{d-1}$. Furthermore, since both functions are continuous, see Lemma 3.5 (ii) and Remark 2.3, it suffices to prove the claim for $\nu \in \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$. For each such vector we can find $\nu_{1}, \ldots, \nu_{d-1} \in \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$ such that the set $\left\{\nu_{1}, \ldots, \nu_{d-1}, \nu\right\}$ forms an orthonormal basis of $\mathbb{R}^{d}$. Then, there exists $\lambda \in \mathbb{N}$ such that

$$
\begin{equation*}
\lambda \nu_{n}=z_{n} \text { for some } z_{n} \in \mathbb{Z}^{d} \text { for all } n \in\{1, \ldots, d\} \tag{52}
\end{equation*}
$$

Step 1.(Proof of ' $\leq$ ') Let $\left\{\nu_{1}, \ldots, \nu_{d-1}, \nu_{d}=\nu\right\} \subset \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$ be an orthonormal basis as previously described and let $1 \ll S_{1} \ll S_{2}$. We assume that $S_{1}=\lambda T$, where $\lambda$ satisfies (52) and $T$ is given by (H1). Note that if $\lambda$ satisfies (52), also $k \lambda$ satisfies (52) and therefore we can find a sequence $S_{k}=k \lambda T$ such that $S_{k} \rightarrow+\infty$ of the desired form. The existence of the limit in definition (6) of $\varphi$ permits us to assume that $S$ is of the specific form. Let $\delta>0$ and $u_{1}: \mathcal{L} \rightarrow\{0,1\}$ be such that $u_{1}(i)=u_{\nu}(i)$ on $\mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}$ and

$$
\begin{equation*}
E\left(u_{1}, Q_{S_{1}}^{\nu}\right)=\inf \left\{E\left(u, Q_{S_{1}}^{\nu}\right): u: \mathcal{L} \rightarrow\{0,1\}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}\right\} \tag{53}
\end{equation*}
$$

Due to the assumption on $S_{1}$ and Lemma 3.1(vi), we have

$$
\begin{equation*}
E\left(u_{1}(\cdot-z), Q_{S_{1}}^{\nu}(z)\right)=E\left(u_{1}, Q_{S_{1}}^{\nu}\right) \text { for all } z=\lambda T \sum_{n=1}^{d} k_{n} \nu_{n}, k \in \mathcal{L} \tag{54}
\end{equation*}
$$

Set (omitting the dependence on $S_{1}$ and $S_{2}$ )

$$
\mathcal{Z}=\left\{z=S_{1} \sum_{n=1}^{d} k_{n} \nu_{n}: k \in \mathcal{L}, Q_{S_{1}}^{\nu}(z) \subset Q_{S_{2}-2 R}^{\nu}\right\}
$$

We define $u_{2}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
u_{2}(i)= \begin{cases}S_{1}\left(u_{1}(i-z)-\frac{1}{2}\right)+\langle\nu, z\rangle & \text { if } z \in \mathcal{Z}, i \in Q_{S_{1}}^{\nu}(z) \\ \langle\nu, i\rangle & \text { otherwise }\end{cases}
$$

By the definition of $u_{2}$, it is clear that

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S_{2}}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{2}}^{\nu}\right\} \leq E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \tag{55}
\end{equation*}
$$

It remains to show that

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \leq \frac{S_{2}^{d}}{S_{1}^{d-1}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+C S_{2}^{d-1} S_{1}^{2}+\frac{S_{2}^{d}}{\delta} C_{S_{1}} \tag{56}
\end{equation*}
$$

In fact, once we have shown (56), Step 1 follows from 5 and (53) by dividing with $S_{2}^{d}$ and letting first tend $S_{2}$ to $+\infty$, then letting $S_{1}$ tend to $+\infty$, and lastly $\delta \rightarrow 0$. We are left to prove (56). In order to prove it we use Lemma 3.1(iv) to obtain

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}^{\nu}\right)=E\left(u_{2}, Q_{S_{2}-2 S_{1}}^{\nu}\right)+E\left(u_{2}, Q_{S_{2}}^{\nu} \backslash Q_{S_{2}-2 S_{1}}^{\nu}\right) \tag{57}
\end{equation*}
$$

and we estimate the two terms on the right hand side separately. We claim that

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}-2 S_{1}}^{\nu}\right) \leq \frac{S_{2}^{d}}{S_{1}^{d-1}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{S_{2}^{d}}{\delta} C_{S_{1}} \tag{58}
\end{equation*}
$$

Indeed, if $i \in Q_{S_{2}-2 S_{1}}^{\nu}$, then there exists $z=S_{1} \sum_{n=1}^{d} k_{n} \nu_{n} \in \mathcal{Z}$ such that

$$
\begin{equation*}
u_{2}(i)=S_{1}\left(u_{1}(i-z)-\frac{1}{2}\right)+\langle\nu, z\rangle \text { for all } i \in Q_{(1+\delta) S_{1}}^{\nu}(z) \tag{59}
\end{equation*}
$$

Due to 40, this is clearly true for $i \in Q_{S_{1}}^{\nu}(z)$, while for $i \in Q_{(1+\delta) S_{1}}^{\nu}(z) \backslash Q_{S_{1}}^{\nu}(z)$ we have $i \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right) \backslash Q_{(1-\delta) S_{1}}^{\nu}\left(z^{\prime}\right)$, for some $z^{\prime}=S_{1} \sum_{n=1}^{d} k_{n}^{\prime} \nu_{n}$ with $\left\|k-k^{\prime}\right\|_{\infty}=1$. Then, due to the boundary conditions of $u_{1}$, we have

$$
\begin{aligned}
u_{2}(i) & =S_{1}\left(u_{1}\left(i-z^{\prime}\right)-\frac{1}{2}\right)+\left\langle\nu, z^{\prime}\right\rangle=S_{1}\left(u_{\nu}\left(i-z^{\prime}\right)-\frac{1}{2}\right)+\langle\nu, z\rangle+\left\langle\nu, z^{\prime}-z\right\rangle \\
& =S_{1}\left(u_{\nu}(i-z)-\frac{1}{2}\right)+\langle\nu, z\rangle=S_{1}\left(u_{1}(i-z)-\frac{1}{2}\right)+\langle\nu, z\rangle
\end{aligned}
$$

Here, the third equality follows, from the fact that $\left\|k-k^{\prime}\right\|_{\infty}=1$ and therefore $\left\langle\nu, z^{\prime}-z\right\rangle \in$ $\left\{-S_{1}, 0, S_{1}\right\}$. To obtain the previous equality, we distinguish the following two cases:
$\left\langle\nu, z^{\prime}-z\right\rangle= \pm S_{1} \Longrightarrow u_{\nu}\left(i-z^{\prime}\right)-u_{\nu}(i-z)=\mp 1$ and $\left\langle\nu, z^{\prime}-z\right\rangle=0 \Longrightarrow u_{\nu}\left(i-z^{\prime}\right)=u_{\nu}(i-z)$.

Now (59) together with (54) implies for $z \in \mathcal{Z}$

$$
\begin{align*}
E\left(u_{2}, Q_{S_{1}}^{\nu}(z)\right)= & \sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \cap Q_{(1+\delta) S_{1}}^{\nu}(z)} c_{i, j}\left(u_{2}(i)-u_{2}(j)\right)^{+} \\
& +\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}(z)} c_{i, j}\left(u_{2}(i)-u_{2}(j)\right)^{+}  \tag{60}\\
\leq & E\left(u_{1}, Q_{S_{1}}^{\nu}(z)\right)+\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}(z)} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| \\
= & S_{1} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}(z)} u_{2}(i)-u_{2}(j) \mid .
\end{align*}
$$

We estimate the second term on the right hand side of 60 to obtain (58). In fact, here we claim that

$$
\begin{equation*}
\left|u_{2}(i)-u_{2}(j)\right| \leq \frac{C}{\delta}|i-j| \text { for all } i \in Q_{S_{1}}^{\nu}(z), j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}(z) \tag{61}
\end{equation*}
$$

for some $C>0$ independent of $S_{1}, S_{2}$ and $\delta$. If this is true, then we get

$$
\begin{aligned}
\sum_{i \in Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}(z)} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| & \leq \frac{C}{\delta} \sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S_{1} / 2}} c_{i, j}|i-j| \\
& \leq \#\left(\mathcal{L} \cap Q_{S_{1}}^{\nu}(z)\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S_{1} / 2}} c_{i, j}|i-j| \leq \frac{C_{S_{1}}}{\delta} S_{1}^{d}
\end{aligned}
$$

Hence, noting that for $z, z^{\prime} \in \mathcal{Z}$ such that $z \neq z^{\prime}$, we have $Q_{S_{1}}^{\nu}(z) \cap Q_{S_{1}}^{\nu}\left(z^{\prime}\right)=\emptyset$ and therefore $\# \mathcal{Z} \leq S_{2}^{d} / S_{1}^{d}$, we get

$$
\begin{aligned}
E\left(u_{2}, Q_{S_{2}-2 S_{1}}^{\nu}\right) \leq \sum_{\substack{z \in \mathcal{Z} \\
Q_{S_{1}}^{\nu}(z) \subset Q_{S_{2}-2 S_{1}}^{\nu}}} E\left(u_{2}, Q_{S_{1}}^{\nu}(S)\right) & \leq \# \mathcal{Z}\left(S_{1} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{1}^{d}\right) \\
& \leq \frac{S_{2}^{d}}{S_{1}^{d-1}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d}
\end{aligned}
$$

This is 58. It remains to prove 61. Note for $i \in Q_{S_{1}}^{\nu}(z), j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}$ we have $|i-j| \geq \delta S_{1} / 2$ as already used above. We claim that

$$
\begin{equation*}
\left|u_{2}(i)-u_{2}(j)\right| \leq C\left(S_{1}+|i-j|\right) \tag{62}
\end{equation*}
$$

and therefore (61) holds true. Let us prove 62). There are three cases to consider:
(a) $i=i_{0}+z, j=j_{0}+z^{\prime} i_{0} \in Q_{S_{1}}^{\nu}(z), j_{0} \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right), z, z^{\prime} \in \mathcal{Z}$,
(b) $i=i_{0}+z, i_{0} \in Q_{S_{1}}^{\nu}(z), z \in \mathcal{Z} j_{0} \notin Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for any $z^{\prime} \in \mathcal{Z}$,
(c) $i \notin Q_{S_{1}}^{\nu}(z)$ for any $z \in \mathcal{Z}$ and $j \notin Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for any $z^{\prime} \in \mathcal{Z}$.

Case (a): Note that in the case where $i=i_{0}+z, j=j_{0}+z^{\prime}$ for some $i_{0}, j_{0} \in Q_{S_{1}}^{\nu}$ and for some $z, z^{\prime} \in \mathcal{Z}$, we have

$$
\left|u_{2}(i)-u_{2}(j)\right| \leq\left|\left\langle\nu, z-z^{\prime}\right\rangle\right|+C S_{1} \leq\left|\left\langle\nu, z+i_{0}-z^{\prime}-j_{0}\right\rangle\right|+\left|i_{0}-j_{0}\right|+C S_{1} \leq|i-j|+C S_{1}
$$

and therefore 62 holds true.

Case (b): Note that in the case where $i=i_{0}+z, i_{0} \in Q^{\nu}, z \in \mathcal{Z}$ and $j \notin Q_{S_{1}}^{\nu}(z)$ for any $z \in \mathcal{Z}$, we have

$$
\left|u_{2}(i)-u_{2}(j)\right| \leq C S_{1}+|\langle\nu, z-j\rangle| \leq C S_{1}+|\langle\nu, i-j\rangle|+\left|i_{0}\right| \leq C S_{1}+|i-j| .
$$

Also here 62 holds true.
Case (c): In this case $u_{2}(i)=\langle\nu, i\rangle$ and $u_{2}(j)=\langle\nu, j\rangle$ and therefore 62 holds true.
Next, we prove

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}^{\nu} \backslash Q_{S_{2}-2 S_{1}}^{\nu}\right) \leq C S_{2}^{d-1} S_{1}^{2} \tag{63}
\end{equation*}
$$

Now, using 62, (H2), and Lemma 3.1.v), we get

$$
\begin{aligned}
E\left(u_{2}, Q_{S_{2}}^{\nu} \backslash Q_{S_{2}-2 S_{1}}^{\nu}\right) & \leq \sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu} \backslash Q_{S_{2}-2 S_{1}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| \\
& \leq C S_{1} \#\left(\mathcal{L} \cap Q_{S_{2}}^{\nu} \backslash Q_{S_{2}-2 S_{1}}^{\nu}\right)\left(1+\max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j|\right) \leq C S_{2}^{d-1} S_{1}^{2}
\end{aligned}
$$

This yields (63). Now (57), (58), and (63) give (56) and therefore the conclusion of Step 1.
Step 2.(Proof of ' $\geq$ ') Here, we proceed in two sub-steps. First we extend a competitor for $\psi\left(\right.$ for fixed $\left.S_{1}\right)$ periodically and then we perform a cut-off construction. Let $\left\{\nu_{1}, \ldots, \nu_{d-1}, \nu\right\}$ be an orthonormal basis as described at the beginning of the proof and let $1 \ll S_{1} \ll S_{2}$. As in Step 1, we assume that $S_{1}=\lambda T$, where $\lambda$ satisfies (52) and $T$ is given by (H1). Furthermore, we assume that $S_{2}=k S_{1}$ for some $k \in \mathbb{N}$. Let $\delta>0$ and $u_{1}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{1}(i)=\langle\nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S_{1}}$ and

$$
\begin{equation*}
E\left(u_{1}, Q_{S_{1}}^{\nu}\right)=\inf \left\{E\left(u, Q_{S_{1}}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}\right\} \tag{64}
\end{equation*}
$$

Due to the assumptions on $S_{1}$ and Lemma 3.1(vi), we have

$$
\begin{equation*}
E\left(u_{1}(\cdot-z), Q_{S_{1}}^{\nu}(z)\right)=E\left(u_{1}, Q_{S_{1}}^{\nu}\right) \text { for all } z=\lambda T \sum_{n=1}^{d} k_{n} \nu_{n}, k \in \mathcal{L} \tag{65}
\end{equation*}
$$

Set as in Step 1 (omitting the dependence on $S_{1}$ )

$$
\mathcal{Z}^{\prime}=\left\{z=S_{1} \sum_{n=1}^{d-1} k_{n} \nu_{n}: k \in \mathbb{Z}^{d-1}\right\}
$$

Step 2.1.(Periodic extension) We assume that

$$
\begin{equation*}
\left\|u_{1}\right\|_{L^{\infty}\left(Q_{S_{1}}^{\nu}\right)} \leq S_{1} / 2 \tag{66}
\end{equation*}
$$

since otherwise we consider $\hat{u}_{1}(i)=\left(u_{1}(i) \wedge\left(S_{1} / 2\right)\right) \vee\left(-S_{1} / 2\right)$. Let $\tilde{u}_{1}: \mathcal{L} \rightarrow \mathbb{R}$ be defined by

$$
\tilde{u}_{1}(i)= \begin{cases}S_{1}^{-1} u_{1}(i-z)+\frac{1}{2} & i \in Q_{S_{1}}^{\nu}(z), z \in \mathcal{Z}^{\prime} \\ u_{\nu}(i) & \text { otherwise }\end{cases}
$$

Note that, due to 66), we have

$$
\begin{equation*}
\left\|\tilde{u}_{1}\right\|_{L^{\infty}\left(Q_{S_{2}}^{\nu}\right)} \leq 1 \tag{67}
\end{equation*}
$$

We prove that

$$
\begin{equation*}
E\left(\tilde{u}_{1}, Q_{S_{2}}^{\nu}\right) \leq \frac{S_{2}^{d-1}}{S_{1}^{d}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1}+C \delta S_{2}^{d-1} \tag{68}
\end{equation*}
$$

for some $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. In order to see this, first observe that for all $z \in \mathcal{Z}^{\prime}$ there holds

$$
\begin{equation*}
\left|\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right| \leq S_{1}^{-1}\left|u_{1}(i-z)-u_{1}(j-z)\right| \text { for all } i, j \in Q_{(1+\delta) S_{1}}^{\nu}(z) \tag{69}
\end{equation*}
$$

Clearly this holds true if $i, j \in Q_{S_{1}}^{\nu}(z)$. On the other hand, if $i \in Q_{(1+\delta) S_{1}}^{\nu}(z) \cap Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$, then $i-z^{\prime} \in \mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}$. We therefore have $u_{1}(i-z)=\langle\nu, i-z\rangle=\langle\nu, i\rangle$ and with that

$$
\tilde{u}_{1}(i)=S_{1}^{-1} u_{1}\left(i-z^{\prime}\right)+\frac{1}{2}=S_{1}^{-1}\langle\nu, i\rangle+\frac{1}{2}=S_{1}^{-1} u_{1}(i-z)+\frac{1}{2} .
$$

Hence, it also holds true for $i, j \in\left(Q_{(1+\delta) S_{1}}^{\nu}(z) \cap Q_{S_{1}}^{\nu}\left(z^{\prime}\right)\right) \cup Q_{S_{1}}^{\nu}(z)$. If $i \in Q_{(1+\delta) S}^{\nu}(z) \backslash Q_{S}^{\nu}\left(z^{\prime}\right)$ for all $z^{\prime} \in \mathcal{Z}^{\prime}$. On the other hand if $i \notin Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for all $z^{\prime} \in \mathcal{Z}^{\prime}$ it suffices to note that $u_{\nu}(i)=S_{1}^{-1}\left(u_{1}(i) \wedge\left(S_{1} / 2\right)\right) \vee\left(-S_{1} / 2\right)$. Hence, also in this case 69 holds true. We first show

$$
E\left(\tilde{u}_{1}, Q_{S_{1}}^{\nu}(z)\right) \leq S_{1}^{-1} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1}
$$

for some $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Due to (65), (69), and Lemma 3.1(v) we have

$$
\begin{aligned}
& E\left(\tilde{u}_{1}, Q_{S_{1}}^{\nu}(z)\right)= \sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L}} c_{i, j}\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+} \leq S_{1}^{-1} \sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \cap Q_{(1+\delta) S_{1}}^{\nu}} c_{i, j}\left(u_{1}(i)-u_{1}(j)\right)^{+} \\
&+\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{\substack{ \\
|i-j| \geq \delta S_{1} / 2}} c_{i, j}\left|u_{1}(i)-u_{1}(j)\right| \\
& \leq S_{1}^{-1} E\left(u_{1}(\cdot-z), Q_{S_{1}}^{\nu}(z)\right)+\frac{C}{\delta S_{1}} \#\left(\mathcal{L} \cap Q_{S_{1}}^{\nu}(z)\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
&|i-j| \geq \delta S_{1} / 2 \\
& \leq S_{1}^{-1} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{1}^{d-1}
\end{aligned}
$$

for some $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Now note that $\# \mathcal{Z}^{\prime} \leq S_{2}^{d-1} / S_{1}^{d-1}$ and therefore, due to Lemma 3.1(iv) there holds

$$
\begin{equation*}
E\left(\tilde{u}_{1}, \bigcup_{z \in \mathcal{Z}^{\prime}} Q_{S_{1}}^{\nu}(z)\right) \leq \frac{S_{2}^{d-1}}{S_{1}^{d}} E\left(\tilde{u}_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1} \tag{70}
\end{equation*}
$$

for some $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Next, we show

$$
\begin{equation*}
E\left(\tilde{u}_{1}, Q_{S_{2}}^{\nu} \backslash \bigcup_{z \in \mathcal{Z}^{\prime}} Q_{S_{1}}^{\nu}(z)\right) \leq C \delta S_{2}^{d-1}+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1} \tag{71}
\end{equation*}
$$

To this end, we introduce

$$
H_{a, b}^{\nu}=\left\{x \in \mathbb{R}^{d}: b \leq\langle x, \nu\rangle<a\right\} .
$$

Note that, by Lemma 3.1(iv), there holds

$$
\begin{align*}
E\left(\tilde{u}_{1}, Q_{S_{2}}^{\nu} \backslash \bigcup_{z \in \mathcal{Z}^{\prime}} Q_{S_{1}}^{\nu}(z)\right)= & E\left(\tilde{u}_{1}, H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu}\right)+E\left(\tilde{u}_{1}, H_{S_{2} / 2,(1+\delta) S_{1} / 2}^{\nu}\right)  \tag{72}\\
& +E\left(\tilde{u}_{1}, H_{-S_{1} / 2,(1+\delta) S_{1} / 2}^{\nu}\right)+E\left(\tilde{u}_{1}, H_{-(1+\delta) S_{1} / 2,-S_{2} / 2}^{\nu}\right)
\end{align*}
$$

We only estimate the first to terms of 72 , the other two being analogous. We first show

$$
\begin{equation*}
E\left(\tilde{u}_{1}, H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu}\right) \leq C \delta S_{2}^{d-1}+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1} \tag{73}
\end{equation*}
$$

with $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Now, due to (H2), 67), 69), and Lemma 3.1(v), there holds

$$
\begin{aligned}
E\left(\tilde{u}_{1}, H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu}\right) & \sum_{i \in \mathcal{L} \cap H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L} \cap H_{(1+\delta) S_{1} / 2,(1-\delta) S_{1} / 2}^{\nu} \cap Q_{(1+\delta) S_{2}}^{\nu}} c_{i, j}\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+} \\
& +\sum_{i \in \mathcal{L} \cap H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}} c_{i, j}\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)_{H_{(1+\delta) S_{1 / 2,(1-\delta) S_{1} / 2}^{\nu} \cap Q_{(1+\delta) S_{2}}^{\nu}}} \leq S_{1}^{-1} \#\left(\mathcal{L} \cap H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu}\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
& +\frac{1}{\delta S_{1}} \#\left(\mathcal{L} \cap H_{(1+\delta) S_{1} / 2, S_{1} / 2}^{\nu}\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
\leq & C \delta S_{2}^{d-1}+\frac{C_{S_{1}}}{\delta} S_{2}^{d-1} .
\end{aligned}
$$

This is $(73)$. We now show that the second term in 72 is estimated by

$$
\begin{equation*}
E\left(\tilde{u}_{1}, H_{S_{2} / 2,(1+\delta) S_{1} / 2}^{\nu}\right) \leq C_{S_{1}} S_{2}^{d-1} \tag{74}
\end{equation*}
$$

where $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Note that $\tilde{u}_{1}(i)=u_{\nu}(i)=1$ for all $i \in \mathcal{L}$ such that $\langle\nu, i\rangle \geq S_{1} / 2$. Therefore, due to (H1), Lemma 3.1(v), 67), for $S_{1}$ big enough, there holds

$$
\begin{aligned}
& E\left(\tilde{u}_{1}, H_{S_{2} / 2,(1+\delta) S_{1} / 2}^{\nu}\right) \leq 2 \sum_{i \in H_{S_{2} / 2,(1+\delta) S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L} \cap\left\{\langle\nu, i\rangle \leq S_{1} / 2\right\}} c_{i, j} \\
& \leq C \sum_{i_{0}, j_{0} \in \mathcal{L} \cap Q_{T}} \sum_{z \in T \mathbb{Z}^{d} \cap H_{S_{2},(1+2 / 3 \delta) S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}} \sum_{\substack{z^{\prime} \in T \mathbb{Z}^{d} \\
\left\langle\nu, z^{\prime}\right\rangle \leq(1+1 / 3 \delta) S_{1} / 2}} c_{i_{0}+z, j_{0}+z^{\prime}} \sum_{i_{0}, j_{0}+\zeta} \\
& \leq C \max _{i_{0}, j_{0} \in \mathcal{L} \cap Q_{T}} \sum_{\substack{\zeta \in T \mathbb{Z}^{d} \\
|\zeta| \geq 1 / 6 \delta S_{1}}} \sum_{z \in T \mathbb{Z}^{d} \cap H_{S_{2},(1+2 / 3 \delta) S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}}^{\langle\zeta+z, \nu\rangle \leq(1+1 / 3 \delta) S_{1} / 2} \\
& \leq C \max _{i_{0}, j_{0} \in \mathcal{L} \cap Q_{T}} \sum_{\substack{\zeta \in T \mathbb{Z}^{d}}} c_{i_{0}, j_{0}+\zeta|\zeta| S_{2}^{d-1}} \sum_{|\zeta| \geq 1 / 6 \delta S_{1}} \\
& \leq C \max _{i \in \mathcal{L}} \sum_{\substack{ } \mathcal{L}} c_{i, j}|j-i| S_{2}^{d-1} \leq C_{S_{1}} S_{2}^{d-1} \\
&|i-j| \geq 1 / 12 \delta S_{1}
\end{aligned}
$$

This is (74). Now (73), (74) imply (71). (71) together with (70) implies (68) and with that the claim.

Step 2.2. (Cut-off construction) Let $\varphi \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$

$$
\varphi(x)=1 \text { on } Q_{(1-2 \delta) S_{2}}^{\nu}, \varphi(x)=0 \text { on } \mathbb{R}^{d} \backslash Q_{(1-\delta) S_{2}}^{\nu}, \text { and }\|\nabla \varphi\|_{\infty} \leq C \delta^{-1} S_{2}^{-1}
$$

Let

$$
u_{2}(i)=\varphi(i) \tilde{u}_{1}(i)+(1-\varphi(i)) u_{\nu}(i) .
$$

Clearly $u_{2}(i)=u_{\nu}(i)$ for $i \in \mathcal{L} \backslash Q_{(1-\delta) S_{2}}^{\nu}$ and therefore, due to Lemma 3.2, there holds

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S_{2}}^{\nu}\right): u: \mathcal{L} \rightarrow\{0,1\}, u(i)=u_{\nu}(i) \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{2}}^{\nu}\right\} \leq E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \tag{75}
\end{equation*}
$$

We prove that

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \leq E\left(\tilde{u}_{1}, Q_{S_{2}}^{\nu}\right)+C \delta S_{2}^{d-1}+\frac{C_{S_{2}}}{\delta} S_{2}^{d-1} \tag{76}
\end{equation*}
$$

where $C_{S_{2}} \rightarrow 0$ as $S_{2} \rightarrow+\infty$. Note that, using $\sqrt{64}$, (68), and 67 , this concludes Step 2 by dividing with $S_{2}^{d-1}$ letting first $S_{2} \rightarrow+\infty, S_{1} \rightarrow+\infty$, and then $\delta \rightarrow 0$. It remains to prove (76). Clearly $u_{2}(i)=u_{\nu}(i)$ on $\mathcal{L} \backslash Q_{(1-\delta) S_{2}}^{\nu}$, and using $\sqrt[67]{ }$, it is easy to see that

$$
\left(u_{2}(i)-u_{2}(j)\right)^{+} \leq\left(\tilde{u}_{1}(i)-\tilde{u}_{1}(j)\right)^{+}+\frac{C}{\delta S_{2}}\left|i-j\left\|\tilde{u}_{1}(j)-u_{\nu}(j)\left|+\left|1-\varphi(i) \| u_{\nu}(i)-u_{\nu}(j)\right|\right.\right.\right.
$$

Therefore,

$$
\begin{align*}
E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \leq & E\left(\tilde{u}_{1}, Q_{S_{2}}^{\nu}\right)+\frac{C}{\delta S_{2}} \sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|i-j \| \tilde{u}_{1}(j)-u_{\nu}(j)\right|  \tag{77}\\
& +\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu} \backslash Q_{(1-2 \delta) S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|u_{\nu}(i)-u_{\nu}(j)\right|
\end{align*}
$$

We show that

$$
\begin{equation*}
\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|i-j \| \tilde{u}_{1}(j)-u_{\nu}(j)\right| \leq \frac{C_{S_{1}}}{\delta} S_{2}^{d}+\frac{C}{\delta} S_{1} S_{2}^{d-1} \tag{78}
\end{equation*}
$$

To see this, we split the sum in two terms by writing

$$
\begin{aligned}
& \sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|i-j \| \tilde{u}_{1}(j)-u_{\nu}(j)\right| \leq \sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|i-j \| \tilde{u}_{1}(j)-u_{\nu}(j)\right| \\
&+\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}}^{|i-j| \leq S_{1}} \substack{\begin{subarray}{c}{j \in \mathcal{L} \\
|i-j| \geq S_{1}} }} \\
& c_{i, j}\left|i-j \| \tilde{u}_{1}(j)-u_{\nu}(j)\right|
\end{aligned}
$$

Now note that $\tilde{u}_{1}(j) \neq u_{\nu}(j)$ only if $j \in H_{S_{1} / 2,-S_{1} / 2}^{\nu}$ therefore, using Lemma 3.1(v), (H2), and 67), we can estimate

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \leq S_{1}}} c_{i, j}|i-j|\left|\tilde{u}_{1}(j)-u_{\nu}(j)\right| & \leq 2 \sum_{i \in \mathcal{L} \cap H_{3 S_{1} / 2,-3 S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \\
& \leq 2 \#\left(\mathcal{L} \cap H_{3 S_{1} / 2,-3 S_{1} / 2}^{\nu} \cap Q_{S_{2}}^{\nu}\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}|i-j| \leq C S_{1} S_{2}^{d-1}
\end{aligned}
$$

On the other hand, using again Lemma 3.1(v), (H2), and (67), we obtain

$$
\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu}} \sum_{\substack{j \in \mathcal{L} \\|i-j| \geq S_{1}}} c_{i, j}|i-j|\left|\tilde{u}_{1}(j)-u_{\nu}(j)\right| \leq 2 \#\left(\mathcal{L} \cap Q_{S_{2}}^{\nu}\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\|i-j| \geq S_{1}}} c_{i, j}|i-j| \leq C S_{2}^{d} C_{S_{1}}
$$

The previous two inequalities yield $\sqrt[78)]{ }$. Lastly we show

$$
\begin{equation*}
\sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu} \backslash Q_{(1-2 \delta) S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|u_{\nu}(i)-u_{\nu}(j)\right| \leq C \delta S_{2}^{d-1}, \tag{79}
\end{equation*}
$$

Using (S1) and (H2), we have

$$
\begin{aligned}
& \sum_{i \in \mathcal{L} \cap Q_{S_{2}}^{\nu} \backslash Q_{(1-2 \delta) S_{2}}^{\nu}} \sum_{j \in \mathcal{L}} c_{i, j}\left|u_{\nu}(i)-u_{\nu}(j)\right| \\
\leq & \sum_{i_{0}, j_{0} \in \mathcal{L} \cap Q_{T}} \sum_{z \in T \mathbb{Z}^{d} \cap Q_{(1+\delta) S_{2}}^{\nu} \backslash Q_{(1-3 \delta)}^{\nu}} \sum_{z^{\prime} \in T \mathbb{Z}^{d}} c_{i_{0}+z, j_{0}+z^{\prime} \mid} u_{\nu}\left(i_{0}+z\right)-u_{\nu}\left(j_{0}+z^{\prime}\right) \mid \\
\leq & \sum_{i_{0}, j_{0} \in \mathcal{L} \cap Q_{T}} \sum_{\zeta \in T \mathbb{Z}^{d}} c_{i_{0}, j_{0}+\zeta} \\
\leq & C \delta S_{2}^{d-1} \max _{i_{0}, j_{0} \in T \mathbb{Z}^{d} \cap Q_{(1+\delta) S_{2}}^{\nu} \backslash Q_{(1-3 \delta)}^{\nu}} \sum_{-|\zeta|-\sqrt{d} T \leq\langle z, \nu\rangle \leq|\zeta|+\sqrt{d} T} \sum_{\zeta \in T \mathbb{Z}^{d}} c_{i_{0}, j_{0}+\zeta}(|\zeta|+\sqrt{d} T) \\
\leq & C \delta S_{2}^{d-1} \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}(|i-j|+1) \leq C \delta S_{2}^{d-1} .
\end{aligned}
$$

This implies (79). Now (77)-79 imply (76) and with that the conclusion of Step 2.2.
In the next Lemma we show that, assuming affine boundary conditions, the calculation of the asymptotic cell formula with respect to the coordinate cube and the calculation of the asymptotic cell formula with respect to the rotated cube are equivalent.

Lemma 3.8. Let $\nu \in \mathbb{R}^{d}$. Then

$$
\psi(\nu)=\phi(\nu) .
$$

Proof. Before we start the proof, we would like to point out that the various steps of the proof are very similar to the steps of proof of Lemma 3.7. However, we decided to include them here for completeness.

First, note that fact that both $\psi$ and $\phi$ are positively homogeneous functions of degree one (cf. (35) and Lemma 3.6 (ii)) it suffices to consider the case where $\nu \in \mathbb{S}^{d-1}$. Thanks to Lemma 3.5 (ii) and Lemma 3.6 (iii) both functions are continuous. Thus it suffices to prove the claim for $\nu \in \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$. For each such vector we can find $\left\{\nu_{1}, \ldots, \nu_{d-1}\right\} \in \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$ such that the set $\left\{\nu_{1}, \ldots, \nu_{d-1}, \nu_{d}=\nu\right\}$ forms an orthonormal basis of $\mathbb{R}^{d}$. For such an orthonormal basis, it is clear that there exists $\lambda \in \mathbb{N}$ such that

$$
\begin{equation*}
\lambda \nu_{n}=z_{n} \text { for some } z_{n} \in \mathbb{Z}^{d} \text { for all } n \in\{1, \ldots, d\} \tag{80}
\end{equation*}
$$

Step 1. (Proof of ' $\geq$ ') Let $\left\{\nu_{1}, \ldots, \nu_{d-1}, \nu_{d}=\nu\right\} \subset \mathbb{S}^{d-1} \cap \mathbb{Q}^{d}$ be the orthonormal basis described previously and let $1 \ll S_{1} \ll S_{2}$. We assume that $S_{1}=\lambda T$, where $\lambda$ satisfies (80) and $T$ is given by (H1). Let $\delta>0$ and $u_{1}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{1}(i)=\langle\nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}$ and

$$
\begin{equation*}
E\left(u_{1}, Q_{S_{1}}^{\nu}\right)=\inf \left\{E\left(u, Q_{S_{1}}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{1}}^{\nu}\right\} \tag{81}
\end{equation*}
$$

Due to the assumption on $S_{1}$ and Lemma 3.1(vi), we have that

$$
\begin{equation*}
E\left(u_{1}(\cdot-z), Q_{S_{1}}^{\nu}(z)\right)=E\left(u_{1}, Q_{S_{1}}^{\nu}\right) \text { for all } z=S_{1} \sum_{n=1}^{d} k_{n} \nu_{n}, k \in \mathcal{L} . \tag{82}
\end{equation*}
$$

We can assume that

$$
\begin{equation*}
\left\|u_{1}\right\|_{L^{\infty}\left(Q_{S_{1}}^{\nu}\right)} \leq S_{1} \tag{83}
\end{equation*}
$$

since otherwise we consider $\hat{u}_{1}(i)=\left(u_{1}(i) \wedge S_{1}\right) \vee\left(-S_{1}\right)$.

We set

$$
\mathcal{Z}=\left\{z=S_{1} \sum_{n=1}^{d} k_{n} \nu_{n}, k \in \mathcal{L}, Q_{S_{1}}^{\nu}(z) \subset Q_{(1-\delta) S_{2}}\right\}
$$

and we define $u_{2}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
u_{2}(i)= \begin{cases}u_{1}(i-z)+\langle\nu, z\rangle & \text { if } z \in \mathcal{Z}, i \in Q_{S_{1}}^{\nu}(z)  \tag{84}\\ \langle\nu, i\rangle & \text { otherwise } .\end{cases}
$$

By the definition of $u_{2}$ it is clear that

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S_{2}}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{2}}\right\} \leq E\left(u_{2}, Q_{S_{2}}\right) \tag{85}
\end{equation*}
$$

We conclude the proof of Step 1 by showing that

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}\right) \leq \frac{S_{2}^{d}}{S_{1}^{d}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d}+C S_{2}^{d-1} S_{1}^{2} \tag{86}
\end{equation*}
$$

where $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow \infty$. Once this is shown, Step 1 is proven. In fact, using it together with (81) and (85) and dividing by $S_{2}^{d}$ and letting first $S_{2}$ tend to $+\infty$, then letting $S_{1}$ tend to $+\infty$, and finally letting $\delta \rightarrow 0$, we obtain the claim. We are left to prove 86 . In order to do so, we employ Lemma 3.1(iv) to obtain

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}}\right)=E\left(u_{2}, \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right)+E\left(u_{2}, Q_{S_{2}} \backslash \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right) \tag{87}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
E\left(u_{2}, \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right) \leq \frac{S_{2}^{d}}{S_{1}^{d}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d} \tag{88}
\end{equation*}
$$

In order to see this, note that for all $z \in \mathcal{Z}$, we have

$$
\begin{equation*}
u_{2}(i)=u_{1}(i-z)+\langle\nu, z\rangle \text { for all } i \in Q_{(1+\delta) S_{1}}^{\nu}(z) \tag{89}
\end{equation*}
$$

Consulting 84, this is clearly true for $i \in Q_{S_{1}}^{\nu}(z)$, while for $i \in Q_{(1+\delta) S_{1}}^{\nu}(z) \backslash Q_{S_{1}}^{\nu}(z)$ there are two cases to check:
(a) $i \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for some $z^{\prime} \in Z$,
(b) otherwise.
(a): If $i \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$, then $i \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right) \backslash Q_{(1-\delta) S_{1}}^{\nu}\left(z^{\prime}\right)$ for some $z^{\prime} \in Z$. Then, due to the boundary conditions of $u_{1}$ we have

$$
u_{2}(i)=u_{1}\left(i-z^{\prime}\right)+\langle\nu, z\rangle=\left\langle\nu, i-z^{\prime}\right\rangle+\left\langle\nu, z^{\prime}\right\rangle=\langle\nu, i\rangle=\langle\nu, i-z\rangle+\langle\nu, z\rangle=u_{1}(i-z)+\langle\nu, z\rangle
$$

and 89) follows.
(b): If $i \notin Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for any $z^{\prime} \in Z$ then

$$
u_{2}(i)=\langle\nu, i\rangle=\langle\nu, i-z\rangle+\langle\nu, z\rangle=u_{1}(i-z)+\langle\nu, z\rangle
$$

and also here 89 follows. In order to proceed, note that

$$
\begin{equation*}
\left|u_{2}(i)-u_{2}(j)\right| \leq C\left(S_{1}+|i-j|\right) . \tag{90}
\end{equation*}
$$

To see this, we need to distinguish between three cases:
(a) $i \in Q_{S_{1}}^{\nu}(z), j \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for $z, z^{\prime} \in \mathcal{Z}$,
(b) $i \in Q_{S_{1}}^{\nu}(z)$ for $z \in \mathcal{Z}, j \notin Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$ for any $z^{\prime} \in \mathcal{Z}$,
(c) otherwise.

Case (a): In order to see this, note that, we have $\left|z-z^{\prime}\right|=\left|z-i-z^{\prime}+j+i-j\right| \leq|i-j|-C S_{1}$. Therefore, using 83), we have for $i \in Q_{S_{1}}^{\nu}(z), j \in Q_{S_{1}}^{\nu}\left(z^{\prime}\right)$

$$
\left|u_{2}(i)-u_{2}(j)\right| \leq 2 S_{1}+\left|z-z^{\prime}\right| \leq C S_{1}+|i-j|
$$

Case (b): Again, due to 83), we get

$$
\left|u_{2}(i)-u_{2}(j)\right|=\left|u_{1}(i-z)+\langle\nu, z-j\rangle\right| \leq S_{1}+|i-z|+|j-i| \leq C S_{1}+|i-j|
$$

Case (c): This case is trivially true, due to the definition of $u_{2}$.
We are now in the position to prove 88 ). To this end, we prove

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{1}}^{\nu}(z)\right) \leq E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{1}^{d} \tag{91}
\end{equation*}
$$

where $C_{S_{1}} \rightarrow 0$ as $S_{1} \rightarrow+\infty$. Using this, Lemma 3.1(iv), and noting that for $z, z^{\prime} \in \mathcal{Z}$ such that $z \neq z^{\prime}$ we have $Q_{S_{1}}^{\nu}(z) \cap Q_{S_{1}}^{\nu}\left(z^{\prime}\right)=\emptyset$ and therefore $\# \mathcal{Z} \leq S_{2}^{d} / S_{1}^{d}$, we get

$$
E\left(u_{2}, \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right)=\sum_{z \in \mathcal{Z}} E\left(u_{2}, Q_{S_{1}}^{\nu}(z)\right) \leq \# \mathcal{Z}\left(E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{1}^{d}\right) \leq \frac{S_{2}^{d}}{S_{1}^{d}} E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\frac{C_{S_{1}}}{\delta} S_{2}^{d}
$$

This is 88. Now, let us prove (91). Using (89) and (82), we have

$$
\begin{align*}
E\left(u_{2}, Q_{S_{1}}^{\nu}(z)\right) & =\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \cap Q_{(1+\delta) S_{1}}^{\nu}} c_{i, j}\left(u_{2}(i)-u_{2}(j)\right)^{+}+\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}} c_{i, j}\left(u_{2}(i)-u_{2}(j)\right)^{+} \\
& \leq E\left(u_{1}, Q_{S_{1}}^{\nu}\right)+\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| \tag{92}
\end{align*}
$$

Now, we are in a position to estimate the second term on the right hand side of (92). Due to (90) and (H2), we have

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L} \backslash Q_{(1+\delta) S_{1}}^{\nu}} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| & \leq C \sum_{i \in \mathcal{L} \cap Q_{S_{1}}^{\nu}(z)} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S_{1} / 2}} c_{i, j}\left(S_{1}+|i-j|\right) \\
& \leq C\left(1+\frac{1}{\delta}\right) \#\left(\mathcal{L} \cap Q_{S_{1}}^{\nu}(z)\right) \max _{i \in \mathcal{L}} \sum_{\substack{j \in \mathcal{L} \\
|i-j| \geq \delta S_{1} / 2}} c_{i, j}|i-j| \\
& \leq \frac{C_{S_{1}}}{\delta} S_{1}^{d}
\end{aligned}
$$

This together with 92 implies 91 .
Next, we prove

$$
\begin{equation*}
E\left(u_{2}, Q_{S_{2}} \backslash \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right) \leq C S_{2}^{d-1} S_{1}^{2} \tag{93}
\end{equation*}
$$

To see this, we use 90 , to obtain

$$
\begin{aligned}
E\left(u_{2}, Q_{S_{2}} \backslash \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right) & \leq \sum_{i \in \mathcal{L} \cap Q_{S_{2}} \backslash \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)} \sum_{j \in \mathcal{L}} c_{i, j}\left|u_{2}(i)-u_{2}(j)\right| \\
& \leq C S_{1} \#\left(\mathcal{L} \cap Q_{S_{2}} \backslash \bigcup_{z \in \mathcal{Z}} Q_{S_{1}}^{\nu}(z)\right) \max _{i \in \mathcal{L}} \sum_{j \in \mathcal{L}} c_{i, j}(|i-j|+1) \leq C S_{2}^{d-1} S_{1}^{2}
\end{aligned}
$$

This is 93 . Now 87), 88), and 93 imply 86 .
Step 2. (Proof of ' $\leq$ ') Let $1 \ll S_{1} \ll S_{2}$. We assume that $S_{1}=\lambda T$ for some $\lambda \in \mathbb{N}$ and $T$ given by (H1). Let $\delta>0$ and $u_{1}: \mathcal{L} \rightarrow \mathbb{R}$ be such that $u_{1}(i)=\langle\nu, i\rangle$ on $\mathcal{L} \backslash Q_{(1-\delta) S_{1}}$ and

$$
\begin{equation*}
E\left(u_{1}, Q_{S_{1}}\right)=\inf \left\{E\left(u, Q_{S_{1}}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{1}}\right\} \tag{94}
\end{equation*}
$$

We set

$$
\mathcal{Z}^{\prime}=\left\{z=S_{1} k: k \in \mathbb{Z}^{d}, Q_{S_{1}}(z) \subset Q_{(1-\delta) S_{2}}^{\nu}\right\}
$$

and we define $u_{2}: \mathcal{L} \rightarrow \mathbb{R}$ by

$$
u_{2}(i)= \begin{cases}u_{1}(i-z)+\langle\nu, z\rangle & \text { if } z \in \mathcal{Z}^{\prime}, i \in Q_{S_{1}}(z) \\ \langle\nu, i\rangle & \text { otherwise }\end{cases}
$$

It is clear that

$$
\begin{equation*}
\inf \left\{E\left(u, Q_{S_{2}}^{\nu}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(i)=\langle\nu, i\rangle \text { on } \mathcal{L} \backslash Q_{(1-\delta) S_{2}}^{\nu}\right\} \leq E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \tag{95}
\end{equation*}
$$

Similar to proving (86) in Step 1, one can show that

$$
E\left(u_{2}, Q_{S_{2}}^{\nu}\right) \leq \frac{S_{2}^{d}}{S_{1}^{d}} E\left(u_{1}, Q_{S_{1}}\right)+C S_{2}^{d-1} S_{1}^{2}+\frac{C_{S_{1}}}{\delta} S_{2}^{d}
$$

Using (94) and (95), this implies the conclusion of Step 2. This can be seen by dividing by $S_{2}^{d}$ and letting first $S_{2}$ tend to $+\infty$, then $S_{1}$ to $+\infty$, and lastly $\delta \rightarrow 0$.

Proof of Proposition 2.6. Our goal is to prove

$$
\begin{equation*}
\varphi(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} \tag{96}
\end{equation*}
$$

for all $\nu \in \mathbb{R}^{d}$. Due to Lemma 3.8, Lemma 3.7, and Lemma 3.3, we have

$$
\begin{equation*}
\varphi(\nu)=\psi(\nu)=\phi(\nu)=\phi_{\text {per }}(\nu) . \tag{97}
\end{equation*}
$$

Additionally, Lemma 3.4 ensures that

$$
\phi_{\text {per }}(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} .
$$

This shows (96) and concludes the proof.

## 4. CRystallinity of the homogenized surface energy density

This section is devoted to the proof of Theorem 2.7. We assume throughout this section that assumptions (S1), (S2) and (H1), (H3) are satisfied.

We define the set of edges $\mathcal{E}$ by

$$
\begin{equation*}
\mathcal{E}=\left\{(i, j) \in\left(\mathcal{L} \cap Q_{T}\right) \times \mathcal{L}: c_{i, j} \neq 0\right\} \text { and } N=\# \mathcal{E} \tag{98}
\end{equation*}
$$

Proof of Theorem 2.7. We divide the proof into three steps. First, we derive a dual representation of $\varphi$. Then, using this representation, we show that $\varphi$ is crystalline.

Step 1. (Dual representation) We define

$$
\begin{align*}
& \mathcal{C}=\left\{\left(\alpha_{i, j}\right)_{i, j}: 0 \leq \alpha_{i, j} \leq c_{i, j}, \alpha_{i+T z, j+T z}=\alpha_{i, j} \text { for all } z \in \mathcal{L}\right. \\
& \left.\qquad \sum_{j \in \mathcal{L}}\left(\alpha_{j, i}-\alpha_{i, j}\right)=0 \text { for all } i \in Q_{T} \cap \mathcal{L}\right\} \tag{99}
\end{align*}
$$

Our goal is to prove

$$
\begin{equation*}
\varphi(\nu)=\frac{1}{T^{d}} \sup _{\left(\alpha_{i, j}\right)_{i, j} \in \mathcal{C}}\left\langle\nu, \sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j}(i-j)\right\rangle . \tag{100}
\end{equation*}
$$

Let $\nu \in \mathbb{R}^{d}$. Due to Proposition, 2.6 there holds

$$
\begin{aligned}
\varphi(\nu)=\phi_{\mathrm{per}}(\nu) & =\frac{1}{T^{d}} \inf \left\{E\left(u, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot)-\langle\nu, \cdot\rangle \text { is } T \text {-periodic }\right\} \\
& =\frac{1}{T^{d}} \inf \left\{E\left(u+\langle\nu, \cdot\rangle, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot) \text { is } T \text {-periodic }\right\}
\end{aligned}
$$

It therefore suffices to prove that the function

$$
\phi_{\mathrm{per}}(\nu)=\frac{1}{T^{d}} \inf \left\{E\left(u+\langle\nu, \cdot\rangle, Q_{T}\right): u: \mathcal{L} \rightarrow \mathbb{R}, u(\cdot) \text { is } T \text {-periodic }\right\}
$$

is crystalline. Note that, we can write

$$
\begin{equation*}
\phi_{\text {per }}(\nu)=\frac{1}{T^{d}} \inf _{\substack{u: \mathcal{L} \rightarrow \mathbb{R} \\ u(\cdot) T-\operatorname{per}}} \sup _{\substack{0 \leq \alpha_{i, j} \leq c_{i, j} \\ \alpha_{i+T z, j+T z}=\alpha_{i, j}}} \sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j}(u(i)-u(j)+\langle\nu, i-j\rangle) \tag{101}
\end{equation*}
$$

Given $0 \leq \alpha_{i, j} \leq c_{i, j}$ such that $\alpha_{i+T z, j+T z}=\alpha_{i, j}$ for all $z \in \mathbb{Z}^{d}$, and $u: \mathcal{L} \rightarrow \mathbb{R} T$-periodic, we have

$$
\begin{aligned}
\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j}(u(i)-u(j)) & =\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j} u(i)-\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j} u(j) \\
& =\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j} u(i)-\sum_{j \in \mathcal{L} \cap Q_{T}} \sum_{z \in \mathbb{Z}^{d}} \sum_{i \in \mathcal{L} \cap Q_{T}} \alpha_{i, j+T z} u(j+T z) \\
& =\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j} u(i)-\sum_{j \in \mathcal{\mathcal { L } \cap Q _ { T }}} \sum_{z \in \mathbb{Z}^{d}} \sum_{i \in \mathcal{L} \cap Q_{T}} \alpha_{i-T z, j} u(j) \\
& =\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j} u(i)-\sum_{j \in \mathcal{L} \cap Q_{T}} \sum_{i \in \mathcal{L}} \alpha_{i, j} u(j) \\
& =\sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}}\left(\alpha_{i, j}-\alpha_{j, i}\right) u(i) .
\end{aligned}
$$

Note that, since in all steps the sum over $i$ and, due to to (H3), the sum over $j$ runs over a finite index set, the order of summation can be changed without changing the value of the various sums. This implies that, given $0 \leq \alpha_{i, j} \leq c_{i, j}$ such that $\alpha_{i+T z, j+T z}=\alpha_{i, j}$ for all $z \in \mathcal{L}$, we have

$$
\inf _{\substack{u: \mathcal{L} \rightarrow \mathbb{R}  \tag{102}\\ u(\cdot) T-\operatorname{per}}} \sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i j}(u(i)-u(j))= \begin{cases}0 & \text { if } \sum_{j \in \mathcal{L}}\left(\alpha_{i, j}-\alpha_{j, i}\right)=0 \text { for all } i \in Q_{T} \cap \mathcal{L} \\ -\infty & \text { otherwise }\end{cases}
$$

Hence, using (97), 99, 101), and 102, , we obtain 100.
Step 2.(Crystallinity) By Remark 2.5, we have

$$
\varphi(\nu)=\frac{1}{T^{d}} \sup _{\zeta \in W_{\varphi}}\langle\nu, \zeta\rangle
$$

So that, by 100

$$
W_{\varphi}=\left\{\sum_{i \in \mathcal{\mathcal { L } \cap Q _ { T }}} \sum_{j \in \mathcal{L}} \alpha_{i, j}(i-j):\left(\alpha_{i, j}\right)_{i, j} \in \mathcal{C}\right\},
$$

with $\mathcal{C}$ given in 99 . Recall $N$ and $\mathcal{E}$ defined in (98). Define $L: \mathbb{R}^{N} \rightarrow \mathbb{R}^{d}$ by

$$
\begin{equation*}
L\left(\alpha_{i, j}\right)_{(i, j) \in \mathcal{E}}=\sum_{(i, j) \in \mathcal{E}} \alpha_{i, j}(i-j) \tag{103}
\end{equation*}
$$

We then find

$$
\begin{equation*}
W_{\varphi}=L\left(\left[0, c_{i, j}\right]^{N} \cap V\right) \tag{104}
\end{equation*}
$$

where $V \subset \mathbb{R}^{N}$ is a linear subspace of co-dimension $T^{d}-1$ given by

$$
\begin{equation*}
V=\left\{\alpha_{i, j} \in \mathbb{R}^{N}: \sum_{j \in \mathcal{L}}\left(\alpha_{i, j}-\alpha_{j, i}\right)=0 \text { for all } i \in Q_{T} \cap \mathcal{L}\right\} \tag{105}
\end{equation*}
$$

Hence, due to $\sqrt[104]{ }, W_{\varphi}$ is the image of the linear map $L$, given in 103), of a $N$-dimensional cube $\left[0, c_{i, \xi}\right]^{N}$ intersected with the linear subspace $V$, given in 105 . The intersection of a cube with a linear subspace is a polytope, and thus also its image through a linear map. This proves that $\varphi$ is crystalline.

Step 3. (Estimate on the number of vertices) Our goal is to prove that

$$
\begin{equation*}
\# \operatorname{extreme}\left(W_{\varphi}\right) \leq 2^{N} \tag{106}
\end{equation*}
$$

where we recall $N$ defined in (98). Let us note that, due to the Krein-Milman Theorem (cf. [19], Theorem 1.13) and (104), it is easy to see that there holds

$$
\# \operatorname{extreme}\left(W_{\varphi}\right)=\# \operatorname{extreme}\left(L\left(\left[0, c_{i, j}\right]^{N} \cap V\right)\right) \leq \# \operatorname{extreme}\left(\left[0, c_{i, j}\right]^{N} \cap V\right)
$$

In order to show 106), it remains to show

$$
\begin{equation*}
\text { \#extreme }\left(\left[0, c_{i, j}\right]^{N} \cap V\right) \leq 2^{N} \tag{107}
\end{equation*}
$$

In order to obtain this estimate we construct a (non necessarily orthogonal) projection $P: \mathbb{R}^{N} \rightarrow$ $\mathbb{R}^{N}$ such that

$$
\begin{equation*}
P\left(\left[0, c_{i, j}\right]^{N}\right)=\left[0, c_{i, j}\right]^{N} \cap V \tag{108}
\end{equation*}
$$

By the Krein-Milman Theorem, it then follows that

$$
\# \operatorname{extreme}\left(\left[0, c_{i, j}\right]^{N} \cap V\right)=\# \operatorname{extreme}\left(P\left(\left[0, c_{i, j}\right]^{N}\right)\right) \leq \# \operatorname{extreme}\left(\left[0, c_{i, j}\right]^{N}\right)=2^{N}
$$

In order to construct $P$ denote by $k=\operatorname{dim}(V)$ and let $\left\{v_{1}, \ldots, v_{k}\right\}$ be a basis of $V$. Add vectors $\left\{e_{i_{1}}, \ldots, e_{i_{N-k}}\right\}$ from the standard orthonormal basis of $\mathbb{R}^{N}$ in order to form a basis of $\mathbb{R}^{N}$. For every $x \in \mathbb{R}^{n}$ we can write

$$
x=x_{v}+x_{c} \text { with } x_{v}=\sum_{j=1}^{k} \lambda_{j} v_{j} \text { and } x_{c}=\sum_{j=1}^{N-k} \mu_{j} e_{i_{j}}
$$

where $\lambda_{k}, \mu_{k} \in \mathbb{R}^{N}$ for all $j \in\{1, \ldots, k\}$ and $j \in\{1, \ldots, N-k\}$ respectively. We define $P: \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ by

$$
P x=x_{v} .
$$

It is easy to see that 108 holds true. In fact on the one hand, we have $P\left(\left[0, c_{i, j}\right]^{N}\right) \supseteq$ $P\left(\left[0, c_{i, j}\right]^{N} \cap V\right)=\left[0, c_{i, j}\right]^{N} \cap V$. On the other hand given $x \in\left[0, c_{i, j}\right]^{N}$, we have

$$
x=\sum_{j=1}^{N} \sigma_{j} e_{j}=\sum_{j=1}^{k} \lambda v_{j}+\sum_{j=1}^{N-k} \mu_{j} e_{i_{j}}=x_{v}+x_{c} .
$$

Now, it is easy to see that $\mu_{j}=\sigma_{i_{j}}$ and therefore

$$
x_{v}=\sum_{j=1}^{k} \lambda v_{j}=\sum_{\substack{j=1 \\ j \notin\left\{i_{j}: j \in\{1, \ldots, N-k\}\right.}}^{N} \sigma_{j} e_{j} .
$$

This implies that $x_{v} \in\left[0, c_{i, j}\right]^{N}$ and clearly $x_{v} \in V$. This shows that $P\left(\left[0, c_{i, j}\right]^{N}\right) \subseteq\left[0, c_{i, j}\right]^{N} \cap V$ and therefore 107 ). This concludes Step 3.

## 5. Differentiability of the effective surface tension

In this Section, we prove Proposition 2.9 which states that $\varphi$ is differentiable in totally irrational directions. It is a corollary of the two lemmas which we state and prove below.

Lemma 5.1. Let $\nu \in \mathbb{S}^{1}$, let $u$ be a minimizer in (7) and assume that for any $s \in \mathbb{R}$, the set $\{u=s\}$ is finite. Then $\varphi$ is differentiable in $\nu$.

Proof. The expression 100 shows that $\varphi$ is a convex, one-homogeneous function with subgradient at $\nu$ given by

$$
\partial \varphi(\nu)=\left\{\frac{1}{T^{d}} \sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j \in \mathcal{L}} \alpha_{i, j}(i-j): \alpha=\left(\alpha_{i, j}\right)_{i, j} \in \mathcal{C} \text { maximizer in } 100\right\}
$$

It is differentiable at $\nu$ if and only if the above set has exactly one element.
Let $\alpha, \alpha^{\prime} \in \mathcal{C}$ be two maximizers in 100 . Classical optimality conditions guarantee that for any $i, j$, if $u(i) \neq u(j)$, then:

$$
\alpha_{i, j}=\alpha_{i, j}^{\prime}= \begin{cases}c_{i, j} & \text { if } u(i)-u(j)>0  \tag{109}\\ 0 & \text { if } u(i)-u(j)<0\end{cases}
$$

Let us denote by $p, p^{\prime} \in \partial \varphi$ the subgradients given by the dual variables, respectively, $\alpha$ and $\alpha^{\prime}$, we claim that $p=p^{\prime}$. One has:

$$
\begin{equation*}
p-p^{\prime}=\frac{1}{T^{d}} \sum_{i \in \mathcal{L} \cap Q_{T}} \sum_{j: u(j)=u(i)}\left(\alpha_{i, j}-\alpha_{i, j}^{\prime}\right)(i-j) \tag{110}
\end{equation*}
$$

Let $s \in \mathbb{R}, i_{0} \in \mathcal{L} \cap Q_{T}$ with $u\left(i_{0}\right)=s$ and such that the finite set $J_{s}:=\{j: u(j)=s\}$ has more than one element. For any $i, j$, let $\beta_{i, j}:=\alpha_{i, j}-\alpha_{i, j}^{\prime}$. Then

$$
\begin{aligned}
\sum_{i \in J_{s}} \sum_{j \in J_{s}} \beta_{i, j}(i-j) & =\sum_{z \in \mathbb{Z}^{d}} \sum_{i \in J_{s} \cap\left(T z+Q_{T}\right)} \sum_{j \in J_{s}} \beta_{i, j}(i-j) \\
& =\sum_{z \in \mathbb{Z}^{d}} \sum_{i \in\left(J_{s}-T z\right) \cap Q_{T}} \sum_{j \in J_{s}-T z} \beta_{i, j}(i-j)
\end{aligned}
$$

where for the last line we have substituted $(i, j)$ with $(i-T z, j-T z)$ and used that $\beta$ is $Q_{T^{-}}$ periodic. In addition, we have that $u(i)=u(j)$ if and only if $u(i-T z)=u(j-T z)$ so that this can be rewritten:

$$
\sum_{i \in J_{s}} \sum_{j \in J_{s}} \beta_{i, j}(i-j)=\sum_{z \in \mathbb{Z}^{d}} \sum_{i \in\left(J_{s}-T z\right) \cap Q_{T}} \sum_{j: u(j)=u(i)} \beta_{i, j}(i-j)
$$

By assumption, the sets $\left(J_{s}-T z\right) \cap Q_{T}, z \in \mathbb{Z}^{d}$ are all disjoint. Otherwise, there would be $i, z$ with $s=u(i-T z)=u(i)+T\langle\nu, z\rangle=s$, yielding in particular that $\langle\nu, z\rangle=0$, and one would deduce that $i-k T z \in J_{s}$ for all $k \in \mathbb{Z}$, a contradiction since we assumed $J_{s}$ was finite. As a consequence, showing that (110) vanishes is equivalent to showing that

$$
\begin{equation*}
\sum_{i \in J_{s}} \sum_{j \in J_{s}} \beta_{i, j}(i-j)=0 \tag{111}
\end{equation*}
$$

for any $s \in \mathbb{R}$ (such that $J_{s}$ is not empty and contains more than one point). Obviously, the expression in 111) is also

$$
\sum_{i \in J_{s}} \sum_{j \in J_{s}}\left(\beta_{i, j}-\beta_{j, i}\right) i
$$

Thanks to the definition (99) of $\mathcal{C}$, one has for any $i$ that $\sum_{j} \beta_{i, j}-\beta_{j, i}=0$, so that:

$$
\sum_{i \in J_{s}} \sum_{j \in J_{s}}\left(\beta_{i, j}-\beta_{j, i}\right) i=\sum_{i \in J_{s}} \sum_{j \notin J_{s}}\left(\beta_{j, i}-\beta_{i, j}\right) i=0
$$

thanks to 109 . Hence, 111 holds and we deduce $p=p^{\prime}$, which shows the lemma.
Lemma 5.2. Let $\nu \in \mathbb{S}^{1}$ be totally irrational and let $u$ be a minimizer in 77. Then for any $s \in \mathbb{R}$, the set $\{u=s\}$ is finite.

Proof. Recalling the notation in the previous proof, let $s \in \mathbb{R}$ and consider the set $J_{s}:=\{u=s\}$. For $z \in \mathbb{Z}$, let $J_{s}^{z}=J_{s} \cap\left(Q_{T}+T z\right)-T z \subset Q_{T}$. For $i \in J_{s}^{z}, u(i)=s+T\langle z, \nu\rangle$. Since $\nu$ is totally irrational, we deduce that $J_{s}^{z} \cap J_{s}^{z^{\prime}}=\emptyset$ for any $z \neq z^{\prime}$, showing that all sets $J_{s}^{z}$ but a finite number must be empty. Hence $J_{s}$ is finite.

## 6. Numerical illustration

6.1. A simplified framework. In this section, we address, as an illustrative experiment, the following issue. We consider a basic $2 D$ cartesian graph $\{(i, j): 0 \leq i \leq M-1,0 \leq j \leq N-1\}$, representing for instance the pixels of an image, and we want to approximate on this discrete grid the two-dimensional total variation $\int_{\Omega}|D u|, u \in B V(\Omega)$. Here it is assumed that $\Omega \subset \mathbb{R}^{2}$ is a rectangle and that $\{0, \ldots, M-1\} \times\{0, \ldots, N-1\}$ is a discretization of $\Omega$ at a length scale $\sim 1 / N \sim 1 / M$.

There are of course many ways to do this, but we propose here to consider a family of discrete "graph" total variations, defined for a $\left(u_{i, j}\right)_{i, j} \in \mathbb{R}^{M \times N}$ by:

$$
\begin{align*}
J(u)=\sum_{i, j} c_{i+\frac{1}{2}, j}^{+} & \left(u_{i+1, j}-u_{i, j}\right)^{+}+c_{i+\frac{1}{2}, j}^{-}\left(u_{i, j}-u_{i+1, j}\right)^{+}  \tag{112}\\
& +c_{i, j+\frac{1}{2}}^{+}\left(u_{i, j+1}-u_{i, j}\right)^{+}+c_{i, j+\frac{1}{2}}^{-}\left(u_{i, j}-u_{i, j+1}\right)^{+}
\end{align*}
$$

and which involves only nearest-neighbour interactions in horizontal and vertical directions.

We assume in addition that the weight $c^{ \pm}$are $T$-periodic for some $T \in \mathbb{N}, T>0$, that is, $C_{a+k T, b+l T}^{ \pm}=c_{a, b}^{ \pm}$for any $(k, l) \in \mathbb{Z}^{2},(a, b)=\left(i+\frac{1}{2}, j\right)$ or $\left(i, j+\frac{1}{2}\right)$, as long as the points fall inside the grid.

For $T=1, c_{a, b}^{ \pm} \equiv 1$, it is standard that 112 approximates, in the continuum limit, the anisotropic total variation $\int_{\Omega}\left|\partial_{1} u\right|+\left|\partial_{2} u\right|$, which, if used for instance as a regularizer for image denoising or reconstruction, may produce undesired artefacts (although hardly visible on standard applications, see Figure 8).

A standard way to mitigate this issue (besides, of course, resorting to numerical analysis based on finite differences or elements in order to define more refined discretizations), is to add to 112 diagonal interactions, with appropriate weights, in order to improve the isotropy of the limit (see for instance [10), with the drawback of complexifying the graph and the optimization. We show here that a similar effect can be attained by homogenization. To illustrate this, let us first consider the simplest situation, for $T=2$.


Figure 3. The alternating 2-periodic coefficients yielding the smallest anisotropy
In that case, one can explicitly build coefficients $c_{a, b}^{ \pm}$, taking two values $\alpha, \beta$ (see Figure 3), which will yield the homogenized surface tension

$$
\begin{equation*}
\varphi(\nu)=(\sqrt{2}-1)\left(\left|\nu_{1}\right|+\left|\nu_{2}\right|+\frac{\left|\nu_{1}+\nu_{2}\right|}{\sqrt{2}}+\frac{\left|\nu_{1}-\nu_{2}\right|}{\sqrt{2}}\right) \tag{113}
\end{equation*}
$$

whose 1-level set (or Frank diagram) is shown in Figure 4 Observe that this is the same


Figure 4. The Frank diagram $\{\nu: \varphi(\nu) \leq 1\}$ given by 113
anisotropy which would be obtained by using constant coefficients and adding interactions along the edges $((i, j),(i+1, j+1))$ and $((i, j),(i+1, j-1))$.

In order to obtain 113 , one needs to tune $\alpha, \beta$ so that a vertical edge and a diagonal edge, in the most favorable position, have the same length (with a $\sqrt{2}$ factor for the diagonal, whose intersection with the periodicity cell is of course longer). This is ensured if $\alpha+\beta=4 \alpha / \sqrt{2}$, that is, $\beta=(2 \sqrt{2}-1) \alpha$. We find that choosing

$$
\left\{\begin{array}{l}
\alpha=\frac{1}{4 \sqrt{2}} \approx 0.1768  \tag{114}\\
\beta=(2 \sqrt{2}-1) \alpha \approx 0.3232
\end{array}\right.
$$

yields 113 , as an effective homogenized anisotropy.
For larger periodicity cells, it seems difficult to do a similar analysis, first of all, because one should not expect the optimal minimizers, in most directions (if not all), to be given by straight lines, but rather by periodic perturbations of straigth lines. We propose an optimization process in order to compute the optimal weights $c_{a, b}^{ \pm}$.
6.2. The optimization method. The effective surface tension is obtained by solving the cell problem:

$$
\begin{gather*}
\phi(\nu)=\min _{u}\left\{\sum_{(i, j) \in Y} c_{i+\frac{1}{2}, j}^{+}\left(u_{i+1, j}-u_{i, j}\right)^{+}+c_{i+\frac{1}{2}, j}^{-}\left(u_{i, j}-u_{i+1, j}\right)^{+}\right. \\
+c_{i, j+\frac{1}{2}}^{+}\left(u_{i, j+1}-u_{i, j}\right)^{+}+c_{i, j+\frac{1}{2}}^{-}\left(u_{i, j}-u_{i, j+1}\right)^{+}:  \tag{115}\\
\left.u_{i, j}-\nu \cdot\binom{i}{j} Y \text {-periodic }\right\}
\end{gather*}
$$

where $Y=\mathbb{Z}^{2} \cap([0, T) \times[0, T))$ is the periodicity cell. This is easily solved, for instance by a saddle-point algorithm [25] which aims at finding a solution to:

$$
\begin{aligned}
\phi(\nu)= & \min _{v Y \text {-periodic }} \max _{0 \leq w \pm} \sum_{(i, j) \in Y}\left(w_{i+\frac{1}{2}, j}^{+} c_{i+\frac{1}{2}, j}^{+}-w_{i+\frac{1}{2}, j}^{-} c_{i+\frac{1}{2}, j}^{-}\right)\left(v_{i+1, j}-v_{i, j}+\nu_{1}\right) \\
& +\left(w_{i, j+\frac{1}{2}}^{-} c_{i, j+\frac{1}{2}}^{-}-w_{i, j+\frac{1}{2}}^{+} c_{i, j+\frac{1}{2}}^{+}\right)\left(v_{i, j+1}-v_{i, j}+\nu_{2}\right)
\end{aligned}
$$

where we have replaced the variable $u$ with the periodic vector $v_{i, j}=u_{i, j}-\nu \cdot(i, j)^{T}$. For technical reasons, we need to "regularize" slightly this problem in order to make it differentiable with respect to the coefficients $\mathbf{c}=\left(c_{\bullet}^{ \pm}\right)$. This is done by introducing $\varepsilon>0$ a (very) small parameter and adding to the previous objective the penalization

$$
-\frac{\varepsilon}{2} \sum_{(i, j) \in Y}\left(w_{i+\frac{1}{2}, j}^{+}\right)^{2}+\left(w_{i+\frac{1}{2}, j}^{-}\right)^{2}+\left(w_{i, j+\frac{1}{2}}^{-}\right)^{2}+\left(w_{i, j+\frac{1}{2}}^{+}\right)^{2}+\frac{\varepsilon}{2} \sum_{(i, j) \in Y} v_{i, j}^{2}
$$

which makes the problem strongly convex/concave and the solutions $w, v$ unique. We call $\phi_{\varepsilon}(\nu)[\mathbf{c}]$ the corresponding value. The advantage of this regularization is that one can easily show that $\mathbf{c} \mapsto \phi_{\varepsilon}(\nu)[\mathbf{c}]$ is locally $C^{1,1}$, with a gradient given by:

$$
\begin{aligned}
\lim _{t \rightarrow 0} \frac{\phi_{\varepsilon}(\nu)[\mathbf{c}+t \mathbf{d}]-\phi_{\varepsilon}(\nu)[\mathbf{c}]}{t}=\sum_{(i, j) \in Y}( & \left.w_{i+\frac{1}{2}, j}^{+} d_{i+\frac{1}{2}, j}^{+}-w_{i+\frac{1}{2}, j}^{-} d_{i+\frac{1}{2}, j}^{-}\right)\left(v_{i+1, j}-v_{i, j}+\nu_{1}\right) \\
& +\left(w_{i, j+\frac{1}{2}}^{-} d_{i, j+\frac{1}{2}}^{-}-w_{i, j+\frac{1}{2}}^{+} d_{i, j+\frac{1}{2}}^{+}\right)\left(v_{i, j+1}-v_{i, j}+\nu_{2}\right)
\end{aligned}
$$

where $(w, v)$ solves the saddle-point problem which defines $\phi_{\varepsilon}(\nu)[\mathbf{c}]$.

Then, to find coefficients which ensure that $\phi$ is as "isotropic" as possible, one fixes a finite set of directions $\left(\nu_{1}, \ldots, \nu_{k}\right)$ (typically, $(\cos (2 \ell \pi / k), \sin (2 \ell \pi / k))$ for $\left.\ell=1, \ldots, k\right)$, and use a first order gradient descent algorithm to optimize:

$$
\mathcal{L}(\mathbf{c})=\sum_{\ell=1}^{k}\left(\phi_{\varepsilon}\left(\nu_{\ell}\right)[\mathbf{c}]-1\right)^{2}
$$

The problem is easily solved for $Y=\{0,1\} \times\{0,1\}, k=8$ and $\left(\nu_{\ell}\right)_{\ell=1}^{8}$ given as above. For larger periodicity cells and more directions, it easily gets trapped in local minima and we use a random initialization in order to be able to find satisfactory solutions. We then test the result by computing the un-regularized surface tension $\phi$ with the resulting coefficients $\mathbf{c}$. We show some results in the next section. Of course, taking a large value of $\varepsilon$ will make the problem easier to solve, but the learned coefficients will not allow to reconstruct a satisfactory surface tension: we need to choose $\varepsilon$ small, an order of magnitude below the error which we expect on the anisotropy of $\phi$.
6.3. Numerical results. We show the outcome of the optimization, in the periodicity cell $Y=\{0, \ldots T-1\} \times\{0, \ldots, T-1\}$ for $T=2,4,6,8$. We plot first the set $\{\varphi \leq 1\}$ or Frank diagram for the effective surface tensions. Figure 5 shows the diagram obtained, for $T=2,4,8$.


Figure 5. Frank diagrams of the effective anisotropies for $T=2,4,8$.
For $T=2$, the optimization yields the same anisotropy as our theoretical proposition (compare with Fig. 4). However, except when initialized with the values in $\sqrt{114}$ ), the algorithm usually outputs different values with the same effective anisotropy, see Fig. 6] (the values in (114) are in some sense better, as for instance a vertical edge will always have the same effective energy with these values, while with the comptuted values displayed in Fig. 6, it will need to pass through the edges in the second column of the cell in order to get the minimal energy).

For $T=4$, one sees that the behaviour is almost isotropic, while for $T=8$, the relative error with the perfect unit disk is about $1 \%$. We illustrate this on an "inpainting" example, which consists in finding the minimal line in a given direction. We consider as an example the direction $(\cos 3 \pi / 8, \sin 3 \pi / 8)$, which is irrational, so that there cannot be a fully periodic solution. The figure 7 displays several minimal half-planes in this orientation. Observe that for this orientation, the results for $T=4$ or 6 look nicer than the result obtained for $T=8$.

We also show a denoising example based on the "ROF" method (which consists simply in minimizing the total variation (defined by the surface tension $\varphi$ ) of an image with a quadratic penalization of the distance to a noisy data, in order to produce a denoised version, see [41]) with the anisotropic tension $\varphi(\nu)=\left|\nu_{1}\right|+\left|\nu_{2}\right|$ ("T = 1") and the optimized homogenized surface tension for $T=4$. The original image is degraded with a Gaussian noise with $10 \%$ standard deviation (with respect to the range of the values). Here, the difference between the


Figure 6. An example of optimized 2-periodic coefficients yielding the same anisotropy as the choice 114


Figure 7. A minimal half-plane in the orientation $(\cos 3 \pi / 8, \sin 3 \pi / 8)$. Top left, boundary datum, the region where the perimeter is minimized is in gray. Top, middle: $\varphi(\nu)=\left|\nu_{1}\right|+\left|\nu_{2}\right|$. Top, right: optimal effective $\varphi$ for $T=2$. Bottom: for $T=4,6,8$.
two regularizers is hardly perceptible (since the data term strongly influences the position of the discontinuities), yet a close-up (bottom row) allows to see a slight difference, for instance on the cheek where the $T=1$ anisotropy produces block structures.

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Figure 8. "ROF" denoising example. Left: noisty image. Middle, denoised with $\varphi(\nu)=\left|\nu_{1}\right|+\left|\nu_{2}\right|$. Right: with the effective tension computed for $T=4$.

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[^0]:    ${ }^{1}$ Actually, the integral representation for the $\Gamma$-limit has only been shown for undirected graphs. However, a slight modification of the proof shows that it is still true for directed graphs.

[^1]:    ${ }^{2} p$ is totally irrational if there is no $q \in \mathbb{Z}^{d} \backslash\{0\}$ such that $\langle q, p\rangle=0$.

[^2]:    ${ }^{3}$ The existence of the limit of $\phi_{\text {per }}$ is a direct consequence of Lemma 3.4 whereas for $\phi$ it follows by classical subadditivity arguments.

