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Review of Geant4-DNA applications for micro and nanoscale simulations

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Abstract

Emerging radiotherapy treatments including targeted particle therapy, hadron therapy or radiosensitisation of cells by high-Z nanoparticles demand the theoretical determination of radiation track structure at the nanoscale. This is essential in order to evaluate radiation damage at the cellular and DNA level. Since 2007, Geant4 offers physics models to describe particle interactions in liquid water at the nanometre level through the Geant4-DNA Package. This package currently provides a complete set of models describing the event-by-event electromagnetic interactions of particles with liquid water, as well as developments for the modelling of water radiolysis.

Since its release, Geant4-DNA has been adopted as an investigational tool in kV and MV external beam radiotherapy, hadron therapies using protons and heavy ions, targeted therapies and radiobiology studies. It has been benchmarked with respect to other track structure Monte Carlo codes and, where available, against reference experimental measurements.

While Geant4-DNA physics models and radiolysis modelling functionalities have already been described in detail in the literature, this review paper summarises and discusses a selection of representative papers with the aim of providing an overview of a) geometrical descriptions of biological targets down to the DNA size, and b) the full spectrum of current micro- and nano-scale applications of Geant4-DNA.

1. Introduction

1 Mathematical and computational models describing the complex biophysical processes
2 associated with radiation induced cell death have been used since the early 1960s. In 1973
3 Chadwick [1] first presented a mathematical formula which accurately fitted experimental
4 data of cell survival as a function of absorbed dose. It was the first model that attempted to
5 consolidate theories of macroscopic dose deposition and micro/nanoscale damages caused
6 by ionising radiation. In macroscopic radiobiological models (such as Chadwick's), small
7 scale behaviour is consolidated into a set of analytical equations representing the large scale
8 behaviour of the system. While these models are fast in terms of computation time, they are
9 not robust enough to predict outcomes for a wide range of input parameters. As physical,
10 chemical and biological interactions of radiation within an organic medium are stochastic
11 processes, a stochastic type model is required for a more accurate description. As a result,
12 with improvements to the speed and general availability of computer hardware, a transition is
13 occurring from simple analytical models to more physically realistic stochastic (i.e. Monte
14 Carlo) models.
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19 In Monte Carlo codes based on a condensed history approach [2], such as EGS [3],
20 PENELOPE [4], GEANT4 [5-7] and MCNP [8], a large number of collision processes are
21 grouped together or "condensed" to a single "step". This approach made the Monte Carlo
22 simulation of charged particle transport computationally feasible, but it is intrinsically
23 inadequate to describe particle interactions at the nanometre scale. For more than twenty
24 years the ab-initio physical mechanism at the basis of radiation damage in biological
25 molecules has been investigated at the DNA and cellular level by means of other dedicated
26 Monte Carlo codes, referred to as track structure codes.
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31 A variety of track structure codes, such as PTra [9], PARTRAC [10], KURBUC [11], TRAX
32 [12], RITRACKS [13] have been developed to calculate the energy deposition at the
33 nanometre scale, modelling particle tracks interaction-by-interaction ("event-by-event"),
34 typically in gaseous medium or liquid water, to approximate biological systems (this will be
35 discussed further in section 2).
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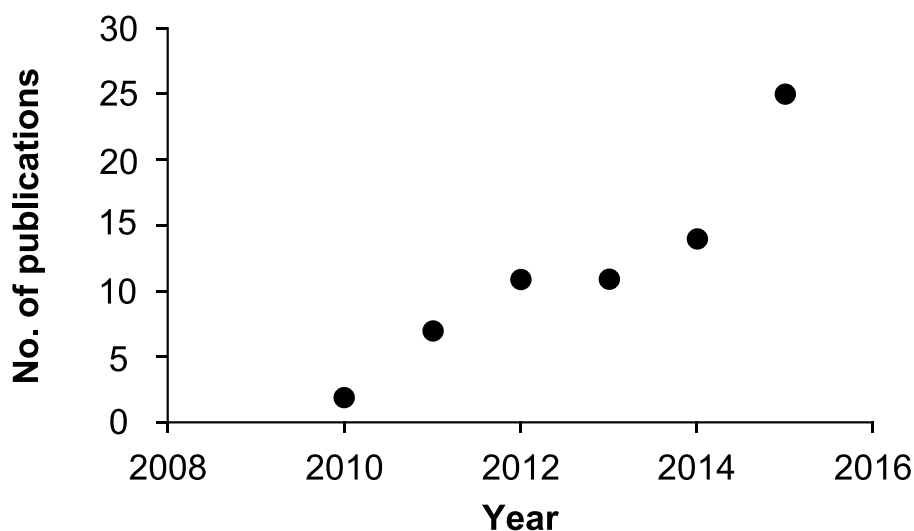
38 The success of the track structure codes resulted from insights from the observation that the
39 micro and nano scale pattern of radiation track and associated energy deposition has a crucial
40 role in the determination of the probability of formation of critical biological lesions [14].
41 The importance of modelling the electron interactions down to eV energies with nanoscale
42 resolution was shown by Nikjoo and Goodhead [15]. These authors estimated that
43 approximately 50% of all ionisations are produced by electrons with energy less than 1 keV
44 in the case of irradiation with a photon or a proton beam. These observations have a
45 significant impact in both radiation protection and radiotherapy studies.
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49 More recently, emerging radiotherapy treatments demand theoretical determination of
50 radiation track structure at the nanoscale, and the radiation damage at the cellular and DNA
51 level. For example, in Targeted Particle (alpha, beta, Auger) Therapy, the calculation of
52 energy deposition at cellular level is necessary because of the short path length of particles
53 emitted by the radionuclide and the spatial distribution of the radionuclide relative to the
54 small target volumes [16, 17]. The investigation of high-Z nanoparticles (NP) as
55 radiosensitisers has demonstrated the need to investigate the effect of radiation at nanoscale
56 [18]. The Local Effect Model (LEM) calculates the survival of cells exposed to a carbon ion
57 beam based on the assumption that the cell response to a given microscopic dose deposited
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1 by densely ionising radiation, such as carbon ions, is related to how the cell would respond to
2 a uniform dose of sparsely ionising radiation, such as X-rays, in a particular nanoscale
3 domain. The LEM is used clinically in the carbon ion beam Treatment Planning System at the
4 Heidelberg Heavy Ion Therapy centre [19]. Alternatively, at the Heavy Ion Medical
5 Accelerator (HIMAC) of NIRS in Chiba, Japan, the RBE₁₀ (RBE at 10% cell surviving
6 fraction in Human Salivary Gland tumour cells [20] was evaluated by means of the
7 Microdosimetric Kinetic Model (MKM), formulated by Hawkins [21, 22] which is based on
8 microdosimetric measurements.

9 Since its first release in 2007, the Geant4-DNA Package has been increasingly used for track
10 structure studies. The Geant4-DNA project, originally initiated by the European Space
11 Agency to study the effect of radiation in astronauts, was increasingly applied to other
12 application domains.
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15 Figure 1 shows the number of published journal articles using Geant4-DNA as an
16 investigation tool between 2010 and 2015. As far as the authors of this paper are aware, the
17 total number of journal articles published in this period is approximately 70. The
18 bibliographic research was performed by means of the Scopus (www.scopus.com) database.
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Figure 1. Number of journal articles published in the period between 2010-2015, using Geant4-DNA as an investigation tool according to Scopus database.

Given the large number of papers published, this review paper summarises and discusses a selection of representative papers with the aim of providing an overview of a) geometrical descriptions of biological targets down to the DNA size, and b) the full spectrum of current applications of Geant4-DNA.

2. Overview of physics and radiolysis modelling capabilities

2.1 Track structure codes

1 Track structure codes have been used for several decades in order to simulate mechanistically
2 the physical interactions of ionising particles with biological matter at small scale and low
3 energy. Some codes are limited to the simulation of physical interactions, while other include
4 the simulation of water radiolysis, realistic geometrical descriptions of biological targets as
5 well as biological repair processes. A detailed and exhaustive review as well as the inter
6 comparison of all past and existing track structure codes is out of the scope of the current
7 work. However, we provide a list of recent reviews on the subject, which the reader can
8 consult; these reviews illustrate the intense on-going development activities and underline the
9 need for such codes. The 2006 review by Nikjoo et al. [23] provides a description of suitable
10 codes for biophysical modelling at molecular level, including full references: these are
11 CPA100, DELTA, ETRACK, KURBUC, LEEPS, LEPHIST, LEAHIST, MC4, NOTRE
12 DAME, OREC, PARTRAC, PITS04, PITS99, SHERBROOKE, STBRGEN, TRION and
13 TRACEL. More recently El Naqa et al. published a topical review on Monte Carlo role in
14 radiobiological modelling of radiotherapy outcomes [24], selecting the representative cases of the
15 KURBUC, NOREC, PARTRAC and Geant4-DNA codes. Palmans et al. included in their recent work
16 [25] a comparison of different Monte Carlo codes for the calculation of microdosimetric spectra (not
17 limited to track structure codes), adding in particular a description of the PTRAN and TRAX codes.
18 Let us finally also add to this list of codes the LQD code [26, 27], NASIC [28], PTra [9] and the
19 RETRACKS/RITRACKS codes [29, 30].
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22 **2.2 Physics modelling**

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27 Since 2007, Geant4-DNA offers physics processes and models able to describe particle
28 interactions in liquid water at nanometre scale [31-33]. This Package currently provides a
29 complete set of models describing the event-by-event electromagnetic interactions of
30 particles (electrons, protons and neutral hydrogen atoms, alpha particles including their
31 charge states) with liquid water (as well as ionisation for a few ions – Li, Be, B, C, N, O, Si
32 and Fe).
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35 In brief, electron processes include ionisation, electronic excitation, elastic scattering,
36 vibrational excitation and molecular attachment. Inelastic interactions (ionisation, electronic
37 excitation) are derived from the model of the dielectric response function of liquid water
38 developed by Emfietzoglou, as recently recalled in [34]. In the most recent Geant4 release
39 (10.2), a new improved implementation of this model has been added in order to account for
40 binding-energy thresholds through a re-distribution of the oscillator strength and refinements
41 in the exchange and perturbation corrections to the Born approximation, as described in detail
42 in [35]. Three elastic scattering models are provided, either derived from a full partial wave
43 analysis for the liquid phase, or from the analytical Screened Rutherford theory with a low
44 energy screening parameter derived from data either in nitrogen gas or in gaseous water [31].
45 Finally, regarding sub-excitation electrons, vibrational excitation is based on the Michaud et
46 al. measurements scaled for the liquid phase, and molecular attachment based on
47 experiments by Melton [36].
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53 Interactions of protons, neutral hydrogen, neutral helium and charge states take into account
54 ionisation, electronic excitation, electron loss or capture and elastic scattering; the
55 corresponding models are described in detail in [32] and in all references cited therein. Low
56 energy ionisation for protons is described using the semi-empirical approach proposed by
57 Rudd et al. (<500 keV) and it is based on the dielectric formalism above this energy, as
58 proposed by Dingfelder et al. [37]. Low energy excitation for protons and neutral hydrogen
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uses a speed scaling from electrons proposed by Miller and Green (< 500 keV) and it is based on the Born and Bethe theory above this energy (for protons only). Electron loss and capture follow the semi-empirical approach also proposed in the work by Dingfelder et al. Finally, elastic scattering is modelled from a classical approach presented recently in [38]. These approaches are also used for the description of ionisation (semi-empirical model by Rudd et al.) and electronic excitation (speed scaling approach from protons by Miller and Green) by helium atoms and their charge states. Electron loss and capture follow the semi-empirical approach of Dingfelder et al. presented in [39]. Elastic scattering is modelled using the same method as for protons.

Finally, regarding ions heavier than helium, only ionisation is currently modeled, above about 1 MeV/amu, based on the Rudd approach and including relativistic extension and considering the effective charge of the ion [40].

The inclusion of new and/or alternative physics models in Geant4-DNA for liquid water or other target materials remains fully open, especially if developers of the above listed track structure codes are interested in providing their models in open access to the community.

2.3 Water radiolysis modelling

Since Geant4 release 10.1, Geant4-DNA includes models of free radical production, diffusion and chemical reactions, following the physical stage of interactions of particles with liquid water. Such functionalities allow the simulation of water radiolysis up to one microsecond after irradiation. These developments are described in detail in the recent papers by Karamitros et al. [41, 42], which include comparisons of Geant4-DNA simulated radiochemical yields to simulation and experimental literature data. They will be required for the simulation of indirect effects of ionising radiation to DNA. Such effects are known to be dominant especially in the low LET domain (see for example Hirayama et al. [43]).

Key parameters are based on PARTRAC as further explained in [41, 42]. In particular, the dissociation scheme is presented in Table 2 of ref. [31], starting from ionised or excited water molecules or dissociative attachment, simulated using the physics models previously presented. The list of molecular species currently considered in the simulation of water radiolysis is as follows: H_3O^+ , OH^- , $\bullet\text{OH}$, $\text{H}\bullet$, H_2 , e^-_{aq} , H_2O_2 . Their diffusion coefficients are listed in Table 1 and reaction and reactions rates are given in Table 2.

Table 1: list of available molecular species and corresponding diffusion coefficients.

Species	Diffusion coefficient ($10^{-9} \text{ m}^2 \text{ s}^{-1}$)
e^-_{aq}	4.9
$\bullet\text{OH}$	2.8
$\text{H}\bullet$	7.0
H_3O^+	9.0
H_2	5.0
OH^-	5.0
H_2O_2	1.4

Table 2: list of reactions and reaction rates.

Reaction	Reaction rate ($10^7 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$)
$\text{H}_3\text{O}^+ + \text{OH}^- \rightarrow 2 \text{H}_2\text{O}$	14.3
$\bullet\text{OH} + \text{e}^-_{\text{aq}} \rightarrow \text{OH}^-$	2.95
$\text{H}\bullet + \text{e}^-_{\text{aq}} + \text{H}_2\text{O} \rightarrow \text{OH}^- + \text{H}_2$	2.65
$\text{H}_3\text{O}^+ + \text{e}^-_{\text{aq}} \rightarrow \text{H}\bullet + \text{H}_2\text{O}$	2.11
$\text{H}\bullet + \bullet\text{OH} \rightarrow \text{H}_2\text{O}$	1.44
$\text{H}_2\text{O}_2 + \text{e}^-_{\text{aq}} \rightarrow \text{OH}^- + \bullet\text{OH}$	1.41
$\text{H}\bullet + \text{H}\bullet \rightarrow \text{H}_2$	1.20
$\text{e}^-_{\text{aq}} + \text{e}^-_{\text{aq}} + 2\text{H}_2\text{O} \rightarrow 2\text{OH}^- + \text{H}_2$	0.50
$\bullet\text{OH} + \bullet\text{OH} \rightarrow \text{H}_2\text{O}_2$	0.44

The following section will concentrate exclusively on Geant4 user applications adopting Geant4-DNA physics processes only, as research based on Geant4-DNA water radiolysis is not published yet. We briefly introduce in the last section of this paper two applications involving Geant4-DNA radiolysis simulation, one for validation purpose and the other one for nanoparticle sensitization study.

3. Geometrical models of DNA

One of the main goals of the Geant4-DNA Package is to be able to predict early damage to DNA [31, 44], taking into account both direct and indirect deleterious actions of ionising radiation on sensitive biological targets. In order to simulate damage to DNA, Geant4-DNA can be utilized following three approaches: (1) the estimation of damage using clustering algorithms, (2) the explicit geometrical modelling of the DNA double helix and associated biological structures of interest, and (3) a mixed approach combining the usage of clustering algorithms with geometrical modelling. These three approaches are described in the following paragraphs.

3.1 Clustering algorithms

The first approach consists of analysing the pattern of energy deposition in the irradiated medium using a clustering algorithm that is tuned in order to reproduce experimental data, e. g. on DNA strand break yields (single strand breaks (SSB), double strand breaks (DSB), complex strand breaks) and survival rates. Francis et al. [40, 45, 46] were the first to propose such an approach based on an adaption of the DBSCAN algorithm [47]. They could simulate SSB, DSB and DSB/SSB ratios as a function of incident proton energy close to other simulated and experimental data on cells and plasmids obtained from the literature. While their algorithm was not included in Geant4/Geant4-DNA, Perrot et al. followed a similar approach as Francis et al., adapting the algorithm from DBSCAN, and they included this clustering algorithm in the “clustering” Geant4-DNA “extended” example in December 2015 (Geant4 release 10.2), and is described in more detail in [31]. They could predict DSB/SSB ratios as a function of incident proton energy close to values predicted by the PARTRAC code [10] and reproduce the results by Francis et al.[46]. The “clustering” example can be conveniently controlled using Geant4 user interface and includes user commands for the

1 setting of clustering parameters, such as the minimum number interaction points (located in
2 space and associated to a local energy deposition) to form a cluster, the probability that a
3 point falls within a sensitive region of a target, the maximum distance separating two points
4 in a cluster and the probability to induce a strand break. Once clusters have been formed, the
5 code scans all clusters once and merge them if their barycenters are separated by less than the
6 maximum distance selected by the user (this is a simplification compared to the DBSCAN
7 algorithm and to the code developed by Francis et al.). The code can also be used with the
8 most recent Geant4-DNA physics models (see section 2), and it calculates single, complex
9 single and double strand breaks as well as cluster sizes and its performance is illustrated in
10 Bernal et al. [31].
11

12 Douglass et al. [48] also proposed independently their own clustering code, allowing in
13 particular the simulation of V79 cell survival after irradiation with MeV protons. The
14 algorithm is based on a hierarchical, geometric clustering process which attempts to group
15 closely spaced ionisation events into DSBs. This code is currently not available in Geant4-
16 DNA. While this approach is interesting for immediate applications (see e.g. [49]), it
17 currently does not take into account the mechanistic simulation of physico-chemical and
18 chemical interactions with sensitive biological targets, which may be required for detailed
19 simulations of damage induction.
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23 **3.2 Geometrical approach**

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25 Alternatively, another approach consists of developing geometrical models of biological
26 targets that can be combined with the simulation of physical, physico-chemical, chemical and
27 biological processes in order to mechanistically simulate damage, such as DNA strand breaks
28 and base oxidation. In the last four years, significant progress has been made in this direction
29 using Geant4-DNA. Two main categories of geometrical models of DNA have been
30 developed for the prediction of damage to DNA: (a) a cylindrical approach where sensitive
31 volumes are described using combinations of elementary cylinder volumes, and (b) a high-
32 resolution atomistic approach where component atoms are represented individually. These
33 geometric models are described below.
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38 *The cylinder volumes approach*

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40 Bernal et al. [50] proposed the first combination of Geant4-DNA physics modelling in liquid
41 water with a geometrical description of DNA based on previous simulations performed with
42 the PENELOPE Monte Carlo code [51, 52]. The DNA double helix is represented as a series
43 of slices in the B-DNA conformation, each of them including two phosphodiester groups
44 bound by a complementary base pair (bp). The base pairs are represented by cylindrical shells
45 and each slice is rotated from its neighbour in order to build a full double helix loop of 100
46 bp. Nucleosomes are modelled as cylindrical histones surrounded by two DNA loops and
47 they are assembled into 30-nm chromatin fibres, each containing a total of 3600 DNA bp.
48 Irradiating a total of 900 fragments of such chromatin fibres with incident protons and alpha
49 particles in the 0.5 – 10 MeV range, the authors could demonstrate that the total strand break
50 yield and the number of energy deposition events required to reach a certain absorbed dose
51 were found nearly independent of the type and energy of the incident ion. The SSB
52 generation process was found to be homogeneous, depending neither on the structure of the
53 DNA nor on the LET of the particles involved, while the DSB yield was found strongly
54 dependent on the LET of the incident radiation [50].
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1 Using the same geometrical description of DNA, Incerti et al. [53] proposed to construct
2 directly in a Geant4 application a simplified nucleus made of randomly oriented short
3 segments of chromatin fibres, for a total of 6×10^9 bp. This simplified nucleus was irradiated
4 with incident protons and the authors confirmed the weak dependence of the total strand
5 break yields with incident particle LET, as already observed in [50]. The DSB yield was
6 found to be strongly dependent on LET. These yields were also found to depend strongly on
7 the method used to account for a direct SSB when compared to other simulated or
8 experimental data.

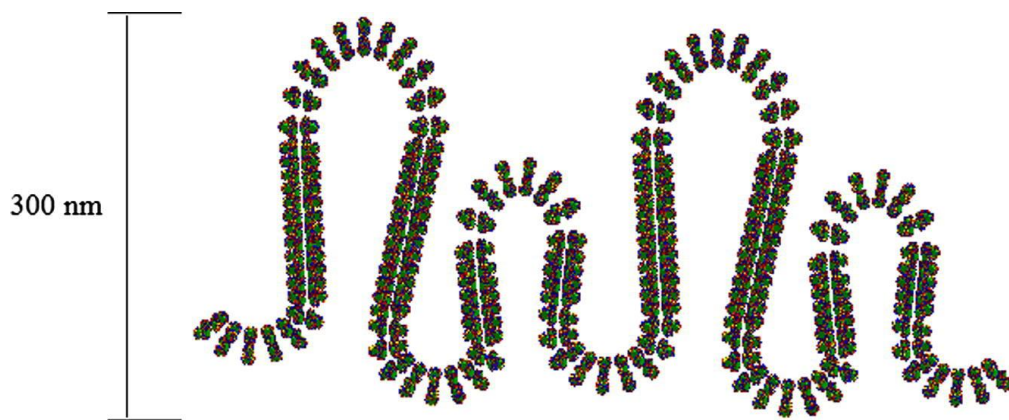
9
10 The B-DNA geometrical model was further completed with two other models, describing the
11 A- and Z- conformations of DNA [44]. It was found by the authors that the total direct strand
12 break yield for a given DNA form depends weakly on DNA conformation topology and this
13 yield is practically determined by the target volume of the DNA configuration. On the
14 contrary, the DSB yield was found to increase with the packing ratio of the DNA double
15 helix, and thus dependent on the DNA conformation.

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18 This cylindrical approach was also followed by Semsarha and colleagues for the prediction of
19 DNA damage in a variety of studies: from Auger emitter radionuclides (^{123}I and ^{125}I) [54],
20 from ^{60}Co gamma rays [55, 56] and from ultrasoft X-rays [57].

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23 Their study on iodine isotopes used two geometrical implementations of 41 base pairs in the
24 B-DNA conformation, the first being similar to the one proposed by Bernal et al. [50], the
25 second having a finer granularity. For the first time with Geant4-DNA, they also included a
26 method using a liquid water virtual cylinder around the DNA for the recording of indirect
27 effects, initially proposed by Pomplun [58] and using prototype radiolysis simulation
28 capabilities available in Geant4-DNA. The authors demonstrated the reasonable agreement
29 on the number of strand breaks per radionuclide decay between Geant4-DNA and theoretical
30 and experimental works from the literature, underlining the applicability of Geant4-DNA for
31 nanodosimetry. They also recommended the usage of accurate geometrical representations of
32 DNA for reliable predictions of indirect effects [54].

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37 The same groups performed two studies on DNA irradiation from ^{60}Co gamma rays. The first
38 study investigated the influence of DNA conformation of strand break yields [55]. The three
39 B-, A- and Z-DNA conformations were studied for 34 bp segments of DNA, including the B-
40 DNA description based on the earlier work by Charlton and Humm [59]. They also took into
41 account a liquid water hydration shell around the DNA for the recording of both direct and
42 indirect effects. From their simulations, the authors observed variations in strand break yields
43 compared to other experimental or simulation works for B-DNA (23% at maximum). They
44 concluded that the B-DNA conformation, the most common conformation of DNA in cells,
45 has the lowest sensitivity for both types of strand break damage. The A-DNA has the highest
46 SSB yields and the Z-DNA has the highest DSB yields. The second study on ^{60}Co gamma
47 rays irradiation focused exclusively on direct effects in a total of about 105 bp [56]. The
48 authors investigated direct strand breaking for several organization levels of DNA: double
49 helix, beads-on-a-string, solenoid (with and without histone proteins) and chromatin loop,
50 starting from the B-DNA geometrical model in [52]. An illustration of a 300 nm chromatin
51 fibre loop is shown in Figure 2, taken from [56]. They found reasonable agreement for DSB
52 yields compared to other theoretical and experimental works, but significant differences for
53 SSB yields of up to 68%. They suggested that the direct strand breaks yields depend mainly
54 on the primary double helix structure of the DNA and that the higher-order structures do not

1 have a noticeable effect on the direct DNA damage under ^{60}Co gamma rays irradiation. The
2 influence of histones on strand break induction was found to be negligible, while the
3 dependence with DNA structure volume was underlined.
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19 Figure 2. Example of 300 nm chromatin fibre loop implemented in Geant4 (Courtesy [56]).
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21 Finally, this group also investigated damage induction from ultrasoft X-rays. They used the
22 Charlton and Humm model [59] for the calculation of direct and indirect damage, for a total
23 of about 10^6 bp. They obtained predictions similar to theoretical and experimental works
24 demonstrating the efficiency of this simple approach. They observed a direct dependence of
25 DSB induction, relative biological effectiveness (RBE) for the induction of DSB and mean
26 lineal energy in strands, connecting nanodosimetry and microdosimetry observations. They
27 also underlined the influence of the SSB threshold induction on break yields, suggesting a
28 value of 10.79 eV for the first ionisation threshold in Geant4-DNA default models.
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32 A last study based on a similar geometrical approach has been proposed by Li et al. [60]: the
33 geometrical model by Charlton and Humm [59] has been used in combination with six
34 different sets of inelastic cross section data describing the transport of electrons and
35 calculated by the authors. These data were imported into their own Geant4-DNA application.
36 These cross sections data were based on Emfietzoglou's approach, as Geant4-DNA models,
37 but included two different optical datasets (by Heller or Hayashi) and three different
38 dispersion models (the extended-Drude model, the extended-oscillator-Drude model and
39 Ashley's delta-oscillator model). Further details on these cross section data are given in [60].
40 The geometrical model allowed the authors to demonstrate that the six inelastic cross sections
41 have a notable influence on the direct DNA strand break yields, underlining again the
42 necessity of accurate physics models for the simulation of electron physical interactions in
43 liquid water.
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49 *The high-resolution atomistic approach* 50

51 The works described in the previous paragraph (e.g. [54]) have underlined the influence of
52 DNA conformation on indirect damage yields induced by chemical species on DNA critical
53 sites. Bernal et al. recently proposed the first freely available stand-alone subroutine allowing
54 the construction of a fully atomistic geometrical description of B-DNA [61]. In this model,
55 the authors accounted for five-organization levels of the DNA, from the nucleotide pair up to
56 the 30 nm chromatin fibre. The provided subroutine is also capable of calculating the distance
57 from an arbitrary point in space to the closest atom, a requirement when one needs to
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1 calculate if a given energy deposition is located in a target volume of interest. This subroutine
2 is not available in Geant4-DNA but can be downloaded as described in [61]. This atomistic
3 model was then used in combination with Geant4-DNA physics modelling capabilities in
4 order to compute direct strand breaks [62]. The authors implemented a cell nucleus
5 containing about 6.5×10^9 base pairs; all atoms were explicitly represented as spheres with
6 corresponding Van der Waals radii. Damages were calculated for incident proton and alpha
7 particles, in the 0.5 – 10 MeV incident energy range. Total damage yields for protons were
8 found consistent with predictions by the PARTRAC software which also uses an atomistic
9 approach [63]. For alpha particles, SSB yields were close to experimental measurements,
10 while DSBs yields appeared larger than experimental values.
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12 In parallel, Xie et al. [64] compared estimations obtained using the previously described
13 cylinder volume approach (model from Charlton and Humm [59]) with an atomistic model
14 created from published DNA atomic coordinates; they also studied the effect of doubling Van
15 der Waals radii in the atomistic model in order to take into account DNA hydration shell. In
16 their results, the authors extracted simple and complex damage yields and demonstrated that
17 the yields obtained with the cylinder volume approach were the highest. They underlined the
18 necessity to take into account complex DSB damage. No comparisons to other calculations or
19 measurements were presented.
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23 A promising approach has been recently included in Geant4-DNA, allowing the
24 implementation of atomistic geometries of a large variety of macromolecules, directly from
25 the Protein Data Bank© (<http://www.rcsb.org>). This is described in the work of Delage et al.
26 [65] (<http://pdb4dna.in2p3.fr>) who developed a dedicated Geant4-DNA extended example,
27 the so-called “pdb4dna” example. This example also includes a cross-platform C++ library of
28 tools. With this code, users can easily implement macromolecules geometries available in the
29 Protein Data Bank© directly in their Geant4 application. In the case of DNA, the granularity
30 of the geometrical model can be selected easily at three levels: representation of individual
31 nucleotides, or representation of components of nucleotides (sugar, phosphate, base) or
32 representation of each individual atom for a full atomistic description (see Figure 3). An
33 algorithm capable of finding the closest atom to an energy deposition is also included, as well
34 as computation of DNA strand breaks.
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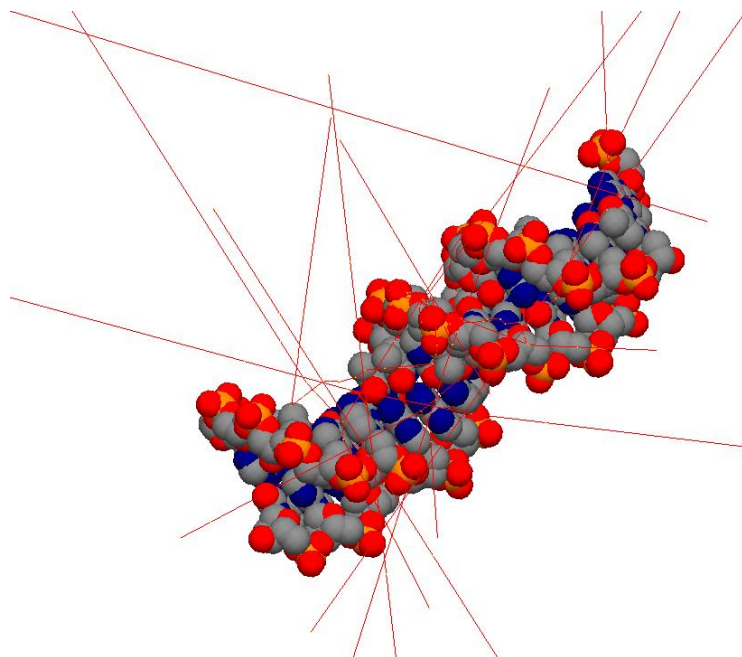


Figure 3. Example of irradiation of the 3BSE 16-base-pair B-DNA PDB© molecule by 10 incident 10 keV electrons, simulated with the Geant4-DNA "pdb4dna" example [65].

The combination of such an atomistic approach with the simulation of radiolysis will be required in order to simulate indirect damage accurately; such a feature is currently under development and is expected to be released in the near future in Geant4-DNA [31, 66].

The mixed approach

In parallel to the previously described applications, a mixed approach has been proposed by Dos Santos and colleagues from 2013 using Geant4-DNA [67-71].

They first proposed the geometrical implementation of two nuclei (including 6×10^9 bp) for a fibroblast and an endothelium cell in the G0/G1 phase. The fibroblast nucleus is included in Geant4-DNA as the so-called "wholeNuclearDNA" extended example [31]. They included five compaction levels: the DNA double helix, nucleosomes, chromatin fibres, chromatin fibre loops and chromosome territories. The reader is invited to look at illustrations available in the recent work of Bernal et al. [31]. In order to simulate direct damage from proton irradiation (0.5 – 50 MeV range), they applied the adapted DBSCAN [46, 71] clustering algorithm to energy deposition events simulated with Geant4-DNA physics processes and located in the sugar-phosphate group. The clustering algorithm is used for the estimation of potential SSBs and DSBs. With this study, the authors could underline the dependence of the quantity and complexity of clusters with increasing DNA density, as in endothelium cells. They extended this work to alpha particles (5 – 50 MeV range) [70] and confirmed that the quantity and the complexity of the potential direct damages in the endothelium cell nucleus are always higher than in the fibroblast cell nucleus. In addition, both studies showed that for a given LET, a proton irradiation induces more direct complex clustered damages than an alpha particle irradiation.

Using a similar approach, they also investigated the influence of chromatin condensation on direct DSBs induced by incident protons and alpha particles by comparing two types of chromosome territories corresponding to euchromatin (decondensed) or heterochromatin (condensed) regions [69, 70]. They reported that condensed chromatin could be the location of more severe radiation-induced lesions than decondensed chromatin.

The same group recently proposed to investigate the influence of the geometrical description of DNA on nanodosimetry parameters of track structures under proton an alpha particle irradiation (in the 1 to 162.5 keV/ μm LET range) [67, 68]. Geant4-DNA simulations were performed in a cylindrical target volume (equivalent to 18 kbp of DNA) containing either the previously described detailed implementation of the double helix [71] or randomly oriented 10-bp long individual homogeneous cylinders as in [72], an approach commonly used in experimental nanodosimetry. Counting ionisations in scoring volumes and using their adapted DBSCAN algorithm in the detailed implementation of the double helix, they could compare the topology of ionisation clusters in both geometries using several nanodosimetric quantities. From their simulations, the authors recommended the usage of the detailed DNA geometry and a scoring method without fixed boundaries in order to extract biologically relevant ionisation cluster size distributions [67, 68].

1 There is thus a large variety of works, which are proposing different geometrical approaches
2 for the modelling of DNA damage using Geant4-DNA. These studies suggest the need for
3 detailed geometrical models of DNA, which shall ultimately be used for the scoring of both
4 direct and indirect damage. In parallel, the recent development of experimental detection of
5 DNA repair proteins using fluorescence time lapse imaging [73, 74] will provide unique
6 opportunities to validate such simulations over different time scales. This will be needed
7 particularly for the inclusion of repair processes in Geant4-DNA, allowing damage estimation
8 well beyond the microsecond, the maximum time limit reachable by the simulation of water
9 radiolysis in Geant4-DNA.

10 11 12 13 **4. Geant4-DNA applications in radiation therapies**

14 Since 2010, Geant4-DNA has been adopted as investigation tool a) in external beam
15 radiotherapy using megavoltage photons produced in medical linear accelerators, b) in hadron
16 therapies using high-energy protons and heavy ions, c) for the study of radiosensitisation in
17 radiotherapy using NPs and d) in targeted therapies where tumour cells are directly targeted
18 by a vector with a radioactive substance attached, including targeted alpha or beta therapy
19 and Auger electron therapy. Some of these studies made use of the geometrical models
20 summarized in the previous section. We present below an overview of these applications in
21 radiation therapies.

22 23 24 25 26 27 28 **4.1 X-ray radiation therapy**

29 *Conventional radiation therapy*

30 For radiation therapy purposes, the interaction of radiation is often modelled in water
31 medium. However, there has been a trend to model semi-realistic cell media representing
32 tumours as well as healthy tissue. In the past, simplified spheroid models with different cell
33 and nucleus radii were used to model cells [75]. Recently, some more realistic cell and
34 tumour models were built [76-78]. Most of the works published to-date have modelled X-
35 rays, gamma-rays and low energy electrons and have investigated production of SSBs and
36 DSBs in order to evaluate the efficacy of various radiation therapy modalities.

37 For example, in the previously described study of Tajik et al. [57] the total yields of SSB and
38 DSB induced by monoenergetic electrons with energies of 0.28–4.55 keV, corresponding to
39 ultrasoft X-rays energies, have been modelled in the Charlton and Humm volume model
40 using the Geant4-DNA toolkit. They observed that in the low energy region, the yield of
41 SSBs remained fairly constant while the yield of DSBs increased with decreasing energy.
42 Moreover, a direct dependency between DSB induction, RBE value and the mean lineal
43 energy as a microdosimetry quantity was noticed.

44 In the other study of Tajik et al. [56], ^{60}Co irradiation was modelled using Geant4-DNA. The
45 yields of direct SSBs and direct DSBs yields induced by secondary electrons were obtained
46 for several DNA modelling structures. As the results obtained agreed well with the
47 experimental data, the authors concluded that Geant4-DNA is a useful tool for micro and
48 nano dosimetry calculations at cellular and sub-cellular levels. They also investigated the
49 importance of modelling the primary (i.e. photon) and secondary (i.e. electron) radiation
50 spectra. They concluded that the DSB yield depended primarily on the secondary electrons

and similar yields were obtained whether the primary photons were modelled or not. The SSB yields, however, decreased if only secondary electrons were considered.

Microbeam Radiation Therapy (MRT)

Geant4-DNA has also been adopted to characterise the secondary electron track structures for 30, 50 and 100 keV X-ray microbeams of width 20 μm , in a water phantom, in a first attempt to get an insight into the physical mechanism at the basis of MRT [79, 80]. The two studies demonstrate the importance of simulating track structures at nanoscale level to more realistically capture the full extent of radiation damage, while dose alone is inadequate for describing radiobiological effectiveness.

The study by McNamara et al. [81] compared the track structures of therapeutic MRT X-ray and proton beams in the attempt to find similarities between the two. If it can be shown that the ionisation cluster distributions are similar, the substitution or supplementation of photon microbeams for protons in targeted single cell irradiations could be feasible. The results of this study suggest that low energy X-rays could produce similar ionisation cluster distributions to MeV protons on the DNA scale at depths greater than $\sim 10 \mu\text{m}$ and at distances greater than $\sim 1 \mu\text{m}$ from the beam centre.

4.2 Hadron radiation therapy

Hadron therapy encompasses irradiations with neutrons, pions, protons and light ions such as carbon ions. Nowadays, proton therapy and carbon ion therapy are the most prevalent forms of hadron therapy and this section will focus on the Geant4-DNA applications to these particle beams.

While the sharpness of the Bragg peak is an important advantage of light ion beams for treatment conformality to the tumour, the rapidly changing stopping power at the end of the particle range means that the energy deposition characteristics of the beam also change with depth. LET increases with an associated increase in ionisation track structure density in the region of the Bragg peak. Quantifying the biological effect of this change in track structure has been the subject of a number of studies (see, for example, [82, 83]). In current clinical practice, the RBE of proton therapy is assumed to be a constant value of 1.1. This value is a result of experimental studies and clinical outcomes [84]. However, it is known that RBE is a complex quantity that depends on dose, LET, cell type, surrounding tissues, biological endpoint and other factors.

Radiobiological effects are handled somewhat differently in carbon ion therapy, where the varying RBE along the Bragg peak is accounted for with an LET-based radiobiological model. A non-uniform physical dose is delivered in an attempt to deliver a uniform RBE weighted dose [85]. Considering the importance of LET and ionisation track structure in hadron therapy treatment planning, a thorough understanding of the physical and radiobiological properties of irradiation by these beams is crucial.

Verification and validation of Geant4-DNA for hadron therapy applications

Several investigations have reported on the validity of the ionisation track structure calculated by Geant4-DNA for incident hadrons. Bug et al. [86] compared the Geant4-DNA angle-

1 dependent energy spectra of secondary electrons created by proton irradiation of amorphous
2 solid water with measurements performed by Toburen et al. [87]. The authors found that
3 simulated spectra were in good agreement with the measured spectra for electron energies
4 between 50 and 100 eV. Above 100 eV measured electron yields were underestimated by
5 Geant4-DNA by approximately 60%. The authors suggested this may have originated from
6 the abrupt limit of the low-energy excitation cross-section at this energy. Below 50 eV,
7 Geant4-DNA increasingly overestimated the yield of electrons, reaching a factor of 7 at 2 eV.
8 However, the authors demonstrated that inclusion of a potential step to model the surface
9 barrier potential of water improved the agreement of simulated and measured yield of
10 electrons below 50 eV, reducing the discrepancy to a factor of 3.
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12 Wang and Vassiliev [88] utilized Geant4-DNA to validate radial dose distributions along
13 single particle tracks for application in the Amorphous Track Model (ATM) commonly
14 applied in RBE models in hadron therapy. The ATM and its variations, including the Local
15 Effect Model, rely on a simplified model of a heavy charged particle track with a track
16 structure characterized by the radial dose distribution from the primary particle [89]. Wang
17 and Vassiliev compared Geant4-DNA calculated radial dose distributions with analytical
18 calculations for proton energies between 10 MeV and 100 MeV. The authors demonstrated
19 that their newly proposed formula could better match the Geant4-DNA radial dose
20 distribution than existing formulae.
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25 Francis et al. [90] compared stopping powers and ranges of electrons, protons and alpha
26 particles in liquid water with analytical models and ICRU data. They found that Geant4-DNA
27 resulted in lower proton stopping powers than ICRU 49 [91] below 10 keV and
28 correspondingly larger particle ranges below 100 keV, but were in good agreement above
29 these energies. This was further improved by including the elastic scattering process in
30 simulations [38]. Nanometric proton energy deposits calculated with Geant4-DNA with the
31 TRIOL code and experiments of Borak et al. [92] were also conducted by Francis et al. [36].
32 In addition, Francis et al. [40] also presented a Geant4-DNA study of energy depositions due
33 to protons, alpha particles and carbon ions of the same LET in liquid water. Energy
34 deposition spatial distributions were analysed using an adapted version of the DBSCAN
35 clustering algorithm [46]. The calculated yields of clustered and single damage and their
36 ratios were directly compared with experimental data on SSB and DSB obtained by the gel
37 electrophoresis technique on irradiated plasmid DNA and other DNA types. The simulation
38 showed good agreement with the experiments and also with the results obtained by the
39 PARTRAC calculations for protons.
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45 While the works mentioned above were all based on the irradiation of water, Champion et al.
46 extended the work of proton irradiations with Geant4-DNA in water to simulations in DNA
47 nucleobases [93, 94]. The authors concluded that material selection had a non-negligible
48 effect on energy deposit distributions in nanometre size targets of constant density. The
49 findings bring into question the validity of simulating biological systems as uniform water
50 environments.
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53 *Hadron therapy applications of Geant4-DNA*

54 Pham et al. [95] compared measured and simulated depth-dose distributions for therapeutic
55 proton beams when the Geant4-DNA Package was incorporated into the GATE [96] toolkit.
56 The authors also evaluated frequencies of energy depositions and DNA damage in the
57 simulated proton beams. DNA damage was evaluated using an algorithm to allocate energy
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1 depositions to atoms constituting DNA molecules represented by their Protein Data Bank
2 (PDB) description. The PDB format provides a standard representation for macromolecular
3 structure data obtained from X-ray diffraction and NMR studies. Proton and electron ranges
4 were calculated with Geant4-DNA implemented in GATE v7.0. No significant differences
5 were observed, validating the porting of Geant4-DNA physics models in GATE. The authors
6 also compared the calculated ranges with ICRU published data and found good agreement
7 with relative differences of less than 8.5% for proton energies between 100 keV and 100
8 MeV and less than 4.8% for electron energies between 10 keV and 100 keV. This work
9 provided an example of merging the popular macroscopic dosimetry toolkit GATE with sub-
10 cellular track structure modelling of Geant4-DNA.

11 Francis et al. [97] utilised Geant4-DNA to examine the track structure of therapeutic carbon
12 ion beams in the fragmentation region distal to the Bragg peak. The authors concluded that
13 the energy deposition clustering yields per event in the fragmentation region were
14 approximately three times less than those induced by the 400 MeV/amu carbon beam within a
15 few centimetres of the phantom surface. These calculations suggest an agreement with the
16 RBE ratio of 3 between the entrance beam and the Bragg peak tail as measured by Scholz et
17 al. [98] in Chinese hamster ovary cells.

18 Casiraghi and Schulte [99] utilized Geant4-DNA to calculate nanodosimetric quantities for
19 therapeutic proton and carbon ion beams in a simplified water environment. The
20 nanodosimetric quantities for each pencil beam were subsequently implemented in an inverse
21 optimization routine to generate beams of uniform nanodosimetric properties, rather than
22 physical dose properties.

23 **4.3 Nanoparticle radiosensitisation**

24 In the last ten years, the study of high atomic number (Z) nanoparticles (NPs) to enhance
25 radiotherapy treatment has been the subject of significant interest in the scientific literature,
26 as this technique shows very promising results in radiobiological experiments. Geant4 offers
27 the capability to select different physics models in different regions of the simulation set-up.
28 This allows the user to select condensed-history physics models, such as the “Low Energy”
29 or “Standard” electromagnetic categories, for modelling particle interactions in the NP, while
30 selecting the Geant4-DNA Package for the surrounding liquid water representing the
31 biological medium, as shown in Fig.4. A 40 MeV proton pencil beam is incident on a water
32 phantom with size equal to 200 μm , where the Geant4-DNA processes are active. A gold
33 nanoparticle is positioned in the centre of the water phantom. In this region the particle
34 interactions are described by means of the Geant4 condensed-random walk scheme physics
35 models. Fig.4 shows the energy deposition produced by the proton beam in a subregion of the
36 water phantom, containing the NP, with a lateral size of 0.2 μm . It can be observed that the
37 description of the proton track structure is much more refined in the liquid water medium
38 where the Geant4-DNA Package is used.

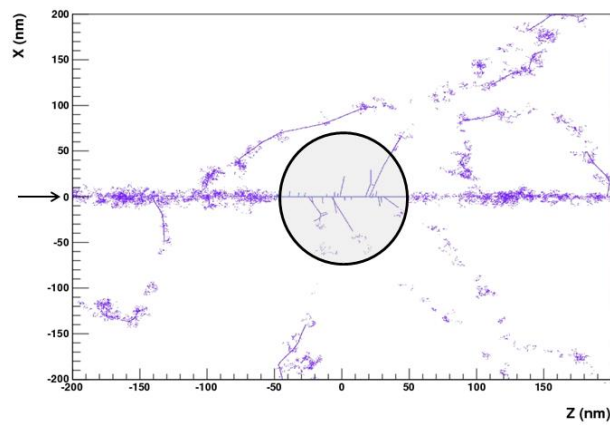


Fig. 4: Geant4-DNA simulation of a 40 MeV proton pencil beam incident on a liquid water phantom with size 200 μm . The direction of incidence of the beam is indicated by the black arrow on the left. A gold NP is placed in the water phantom (grey area enclosed by the black circle). Geant4-DNA processes are activated in the water phantom, while Geant4 condensed-history processes are used to simulate interactions in the NP. The plot shows the energy deposition topology (in arbitrary units) in a subregion of the water phantom containing the NP, with lateral size 0.2 μm . Image courtesy of S McKinnon and S Guatelli, Centre for Medical Radiation Physics, University of Wollongong [100].

McMahon et al. utilised Geant4-DNA models to simulate the dose distributions around GNPs (GNPs) on nanometre scales [18]. Geant4-DNA was selected to provide sufficient accuracy of the track structure in the water volume surrounding the GNPs. The biological effect of GNPs with diameters between 2 and 50 nm being irradiated by monoenergetic photons with energies between 20 and 150 keV was quantified using LEM. A single GNP of varying diameter (between 2 and 50 nm) was placed in the centre of a water cube of side length 200 μm and irradiated with monoenergetic photons between 20 and 150 keV. The rate at which ionisation events occurred in the gold, the spectrum of secondary electrons and the dose distribution in close proximity to the GNP surface were calculated. The study highlighted the role of Auger electrons in the enhancement of the dose close to the NP boundary. The higher energy photo or Compton electrons were found to travel long distances, depositing most of their energy away from the surface of the GNP, whereas Auger electrons deposited most of their energy in the vicinity of the GNP surface. It was also found that the dose deposited in the vicinity of the GNP following a single ionisation event in the gold was dependant on the size of the GNP. Smaller (2 nm) GNP showed the greatest dose deposition outside the NP due to its greater surface area to volume ratio. The Local Effect Model was applied to quantify the radiosensitisation effect of the GNPs. The model was applied to MDA-MB-231 breast cancer cells and it was found to agree with the cell survival curve obtained in radiobiological experiments. This study demonstrated that it is essential to take into account the dose inhomogeneities produced at nanoscale by NPs to be able to predict the experimental results obtained *in vitro* with more accuracy by means of the LEM.

Recently McMahon et al. [101] performed the first systematic study comparing the dose enhancement of alternative NP materials, spanning from silicon ($Z=14$) to mercury ($Z=80$), as possible NP contrast agents and radio sensitisers for imaging and radiotherapy applications. The Geant4-DNA Package was selected to model particle interactions in kV and MV photon beams. The purpose of the study was to investigate the underlying mechanism of

1 dose deposition due to ionising radiation interactions with NPs. A 20 nm diameter NP was
2 placed in a cube of water with side lengths of 10 μm . The “Livermore” set of low energy
3 electromagnetic physics models of Geant4 was used to simulate radiation interactions in the
4 NP and Geant4-DNA was utilised in the surrounding water medium. NPs were spheres
5 composed of pure elemental compositions ranging from $Z=14$ to $Z=80$ with densities and
6 isotope distribution based on NIST reference values. Only solid elements were considered in
7 this study (no gas or liquid phase elements). The NPs were irradiated with kV photon
8 energies and these were set depending on the material. The energy was set to be 20 keV
9 higher than the K-edge of the material being exposed. The energies ranged from 22 to 102
10 keV. The beam width was set to 20 nm so that the entire NP was exposed. NPs were also
11 exposed to a clinical 6 MV linac spectrum. For both spectra, all secondary particles emitted
12 from the NP were scored. The process leading to their emission was identified and the dose
13 deposited was scored in concentric 2 nm shells around the NP out to a range of 1 μm . The
14 LEM was then used to quantify the radiosensitisation effect. The results indicated that while
15 photo electrons contribute most strongly to the overall dose enhancement at kV energies, the
16 dose enhancement within 1 μm of the surface of high Z NPs is due primarily to Auger
17 electrons. Photoelectrons have a large range which results in much of the dose being
18 transported away from the NP. In the case of megavoltage photons, the total dose deposited is
19 dominated by secondary electrons from Compton scatter processes. The effects of Compton
20 electrons are mitigated due to their long range and a contribution from Auger electrons is still
21 seen at short range but this contribution is smaller compared with kV photons.
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27 It was concluded that there are significant variations in radiosensitisation effect for different
28 atomic number NPs. These differences are driven mainly by differences in Auger electron
29 spectra. At least for kV photons, material specific energy tuning is likely to be possible to
30 account for the characteristics of secondary electron energy deposition and NP design.
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33 Douglass et al. developed a randomised 3D tumour model in Geant4 [102, 103]. A cluster of
34 850 cells representing a small tumour was developed as a geometry for Geant4. The
35 translation and rotation of the membrane of each cell was parameterised using *Geant4*
36 *Parameterised Volumes*. Each membrane geometry contained the cytoplasm, reticulum,
37 nucleus and nucleolus. The cellular geometry was applied to investigate the radiosensitisation
38 effect of different sized GNPs from kV and MV photon radiation. Two gold geometries were
39 investigated; a 300 nm thick spherical shell of gold representing a cloud of NPs surrounding
40 the nucleus and a single 400 nm diameter gold sphere in the cytoplasm (Figure 5). These
41 geometries were selected to represent the extrema of possible internalised GNP geometries.
42 The simulations utilised both the Geant4 “Livermore” and Geant4-DNA models. A typical
43 clinical 80 kVp superficial and 6 MV linear accelerator energy spectrum was used to irradiate
44 the cells.
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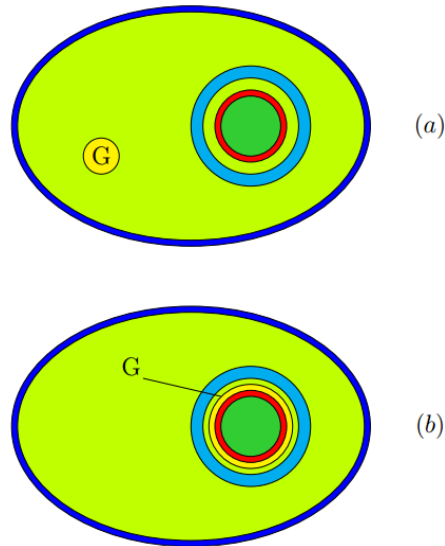


Figure 5. Illustration of cell cross-section with two types of simulation geometries. (a) 400 nm gold cluster (yellow) located randomly in the cytoplasm (lime) and (b) a thin (300 nm) layer of gold (yellow) is attached to the nucleus. Membrane: Dark Blue, Cytoplasm: lime, Reticulum: Cyan, Nucleus: Red, Nucleolus: Green, GNPs: Yellow (Courtesy [103]).

It was concluded that the largest contribution to the dose enhancement effect from GNPs was from increased photoelectron production due to a higher interaction cross section with gold at kV photon energies. It was also shown that Auger electrons may play an important role in the dose enhancement effect at small distances from the gold (up to 500 nm from the surface). Both findings were consistent with the results of McMahan et al. [18]. As a result, GNPs must be internalised close to, or inside, the nucleus in order to produce a significant dose enhancement effect.

Martínez-Rovira et al. [104] utilised GATE 7.0, to investigate the radioenhancement of gold and gadolinium NPs in proton therapy. The research aimed to investigate the effect of NP size, distance between source and NP on NP radiosensitisation. The energy deposited by a proton beam of 200 MeV incident on a NP was calculated in spherical shells around the NP. This work showed that, differently from the case of X-ray radiotherapy, the Auger electrons do not play a crucial role in the determination of the energy deposition enhancement and that the geometrical configuration of the simulation set-up plays a fundamental role when calculating the energy deposition. When adopting a more realistic set-up, the dose enhancement decreases significantly. The authors suggested that physical effects do not seem to be responsible for the NP induced proton radiosensitisation that has been shown in biological studies. This may indicate that it is a chemical or biological process that is responsible for the radiosensitisation effect. This study highlights the necessity to calculate the dose enhancement in more realistic scenarios, which are usually not adopted because of the extensive computing resources required.

In 2014 and 2015, Lin et al. [105, 106] published two articles on the subject of GNP radiosensitisation which encompassed kV and MV photons and therapeutic proton beams. In both papers, the dose deposited in water by secondary electrons through interactions between GNPs and ionising radiation (specifically photons and protons) was investigated.

1 The goal of the first paper [105] was to systematically compare the different mechanisms of
2 GNP radiosensitisation of kV, MV photons and protons using Geant4. The focus was on the
3 distribution of secondary electron dose distribution surrounding the NPs. The effect of energy
4 spectrum, depth of irradiation and radiation type on radiosensitisation was investigated. A
5 single GNP was irradiated with a particle source. The phase space file for the secondary
6 electrons produced by ionisation and excitation within the GNP was scored on the outer
7 surface of the GNP using Geant4 “Penelope” low energy electromagnetic physics models.
8 The track structure outside the NP was then simulated using Geant4-DNA. The most
9 interesting conclusion was that, for the same amount of energy absorbed in the GNP, the dose
10 deposited by secondary electrons within several nanometres of the gold surface differs by less
11 than 20% for kV, MV photons and protons. However, kV photons were shown to produce
12 secondary electrons with the longest range. The authors indicated that when GNPs are
13 internalised through endocytosis, most GNPs are accumulated in lysosomes and are not in
14 close proximity to the nucleus. In these cases, only kV photons (which have long range
15 secondary electrons) are capable of causing increased damage to the nucleus. If NPs could be
16 internalised close to the nucleus or internalised within the nucleus, protons may have a
17 similar radiosensitisation effect.
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21 In the second article by Lin et al. [106], the Monte Carlo models were extended to investigate
22 the biological effects of GNPs using the Local Effect Model. In this article, the same
23 radiation sources were used from the previous work but cell specific geometries were
24 designed. Breast cancer cells, prostate cancer cells and human brain tissue cells were
25 modelled to quantify the radiosensitisation effect. In addition, various GNP geometries were
26 used: GNPs randomly placed in the nucleus, GNPs randomly distributed in the whole cell,
27 GNPs randomly distributed in the cytoplasm, GNPs distributed in the extra cellular material
28 and GNPs distributed in both the cell and extra cellular material. The authors concluded that,
29 proton induced cellular damage is only enhanced by GNPs when they are internalised in the
30 nucleus. It was also discovered that for the same total mass of gold within the cell, 2 nm
31 diameter GNPs have the largest dose enhancement effect. This is due to the higher fraction
32 of low energy electrons escaping from the GNP and contributing to the local dose
33 enhancement. As a final note, it was concluded that if GNPs cannot be internalised within the
34 nucleus, no dose enhancement effect will be seen for proton treatments due to the short range
35 of secondary electrons. When gold is internalised within the cytoplasm, the dose
36 enhancement effect is of the same order for 6 MV X-rays and protons. kV photons showed
37 the highest dose enhancement effect and with the greatest range.
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47 **4.4 Targeted cellular therapy**

48 Geant4 can be adopted for simulation studies in targeted cellular therapy, thanks to its
49 capability to handle radioactive decay and a flexible modelling of radioactive sources both in
50 terms of geometry and generation of radiation field. The *General Particle Source* is
51 recognised to be a flexible, easy to use feature to model particle emission of radioactive
52 sources. Thanks to the atomic de-excitation package of Geant4, it is also possible to generate
53 and track fluorescence X-rays and Auger electrons.
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57 The Dose Point Kernels (DPK), which represent the radial energy deposition distribution
58 from point isotropic sources of electrons, are of great interest in targeted cellular therapy with
59 electron-emitting radiopharmaceuticals, because they are the basis to calculate the dose at
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cellular level, required to estimate the radiobiological effect of the treatment [107]. Another quantity of interest for targeted cellular therapy is the S-value, which represents the dose to the target per unit of cumulated activity in the source region [108].

DPKs and S-values have been calculated by means of Geant4-DNA and compared to other Monte Carlo codes. This approach is valid to verify the suitability of the physics models and is justified when experimental measurements to adopt as a reference are missing.

Dose Point Kernels

Simulations of DPKs with Geant4-DNA were first simulated in liquid water by [107] using mono-energetic point source incident electrons in the 10 keV – 100 keV range. Geant4-DNA was compared to other well-known Monte Carlo codes such as CPA100, EGSnrc, FLUKA, MCNPX and PENELOPE. The reader is invited to consult [107] for code version numbers, code-specific settings used in the simulations and full original references. Using a Kolmogorov-Smirnov statistical test, the authors showed that Geant4-DNA is in agreement with EGSnrc, PENELOPE, and FLUKA and thus provides reliable results in the whole range of incident energies (10 keV - 100 keV), and that significant differences are observed with CPA100 (for 30 keV and 50 keV incident energies) and with MCNPX (at all incident energies).

In order to further improve the accuracy of Geant4-DNA physics models for electrons in the very low energy domain, a new set of physics models has been recently added to Geant4-DNA in June 2015, based on the work by Kyriakou et al. and described in [34] and in [31]. DPKs were calculated with the new set of models at low incident electron energies (100 eV and 1 keV) considering either the Geant4-DNA default models or the newly proposed models. At 100 eV, it was observed that the default Geant4-DNA models lead to a long-tail DPK caused by the too low electronic excitation cross section available by default in Geant4-DNA, as underlined in [32, 48]. At 1 keV, a good agreement is observed between both sets of models.

S-values

First attempts to calculate S-values in liquid water with Geant4-DNA are described in [109]. In this work, the authors calculated S-values for incident mono-energetic electrons (100 eV to 20 keV) and for five iodine isotopes (^{131}I to ^{135}I – the spectra were obtained from previous MC simulations) in simple spherical geometrical targets: spheres of varying radius (10 nm to 1 micrometre) for mono-energetic electrons uniformly distributed in the spheres, or colloid and follicular cells of the thyroid for the five iodine isotopes (two concentric spheres of liquid water separated by 10 micrometres: the inner sphere contains the radionuclide uniformly distributed and its radius varies between 15 and 250 micrometres, representing colloids, whereas the surrounding 10-micrometre thick spherical region represents follicular cells). Results were compared to other Monte Carlo codes (CPA100, MC4V and PENELOPE for the mono-energetic case; CELLDOSE, EGSnrc, MCNP and PENELOPE for the iodine isotopes). The reader is invited to refer to [109] for full references on these codes. The Kolmogorov-Smirnov test was also used in this study to compare Geant4-DNA to the other Monte Carlo codes. For the mono-energetic case, global statistical agreement was found with the other codes, but differences were observed for the small spheres (10 nm radius). For the colloid and follicular cells, global statistical agreement was also observed. It should be however kept in mind that some large variations of relative differences between Geant4-DNA

and these codes can be observed in some extreme cases (e.g. up to 20% in the mono-energetic case and up to 75% in the thyroid case, see further details in [109]).

An attempt to simulate the irradiation of a simplified cell was made by Geng et al.[110]; they simulated S-values in a liquid water spherical cell (radius of 10 μm) composed of a nucleus (radius of 5 μm), a cytoplasm and a membrane. Incident particles were mono-energetic electrons (1 keV to 1 MeV) and 1 MeV alpha particles. S-values are calculated for different source-target combinations. Results obtained for the electron case (cell surface as source and whole cell as target, or cell surface as source and nucleus as target) are reported to be close to MIRD data. Unfortunately, alpha particle results were not compared to other data.

More recently, Fourie et al. [111] focused on ^{123}I in order to investigate the influence of the sub-cellular localization of the Auger emitter in spheres of liquid water. They used a more realistic geometrical model of a cell: a sphere for a whole cell (radius of 5 micrometer), a spherical nucleus (radius of 4 micrometre), a 40 nm thick nuclear membrane and a 7 nm thick cellular membrane. Liquid water was used everywhere except for the membranes where they used soft tissue material (not handled by Geant4-DNA but handled by Geant4). Decay of ^{123}I was simulated using the radioactive decay module of Geant4. S-values obtained with Geant4-DNA were compared with MIRD values and with other Monte Carlo codes (ETRAK and MC4 – refer to [111] for full references) for different source – target combinations: nucleus-nucleus, cytoplasm-nucleus and whole cell – whole cell. The authors found that the maximum dose to the nucleus is delivered when the radioactivity is distributed within the cell nucleus as opposed to the cytoplasm or the whole cell. They also observed that the Geant4-DNA S-values are generally lower than the values calculated by the other MC codes (up to about 24%). The authors explained that the observed differences are mainly due to the different particle emission spectra employed by the different codes, emphasizing the influence of the radionuclide spectra on dosimetry calculations.

Finally, Sefl et al. [112] performed a very detailed study of the influence of cellular geometry on S-values. They selected mono-energetic electrons (1 keV to 100 keV) and three Auger-emitter radionuclides ($^{99\text{m}}\text{Tc}$, ^{111}In and ^{125}I), which spectra were obtained from taken from the AAPM Report no. 2 including the correct spectrum for ^{125}I [113]. Three geometrical models were employed in order to simulate the cell geometry: one spherical model, two ellipsoid models and an irregular shape model. All geometries are fully described in [112]; they have the same cellular and nuclear volumes as the ellipsoid cell models proposed by MIRD and these volumes are equal to the volumes of a spherical cell with radius 5 micrometres and a spherical nucleus of radius 4 micrometres. Random and uniform sampling was presented in detail and was applied to all geometrical configurations, which is especially challenging in the case of non-spherical geometries. Regarding the spherical cell geometry and the mono-energetic electrons, the authors have shown that differences between Geant4-DNA and MIRD data are the largest when the electron range becomes comparable to the size of the sphere (e.g. up to 46% for the cell surface – nucleus source – target combination). In the case of radionuclides, differences with MIRD are in the 5 – 10% range, except for the (cell surface - nucleus) combination, up to 43% for ^{125}I . These differences are attributed to the neglect of delta rays and straggling in the MIRD approach. Regarding ellipsoid geometries and mono-energetic electrons, a very good agreement has been found between Geant4-DNA and MIRD values for S-value ratios of sphere-to-ellipsoid geometries with differences less than 3%. However, differences of about 100% are obtained for the cell surface – nucleus combination at 5 keV; the origin of this discrepancy remains unknown. Regarding radioemitters, it is

found that S-values are generally lower than in the spherical case (up to 32% for ^{111}In and for the cytoplasm – nucleus combination). Lastly, for the irregular cell shape, the largest difference compared to the spherical case has been obtained for the cell surface – nucleus combination while other combinations remain small (below 5%) or moderate (up to 30%).

These works illustrate the possibilities that Geant4-DNA offers to users for dosimetry in targeted cellular therapy. Thanks to the recent addition of alternative models describing more accurately the inelastic interactions of electrons in liquid water [34] and benefitting from advanced imaging techniques such as immunofluorescence staining of individual cells or groups of cells with sub-micrometre resolution [114-116], Geant4-DNA can provide users with a set of simulation features for internal dosimetry in very realistic geometrical models of cells for a variety of incident radiation qualities.

4.5 Geant4-DNA coupled with tracking in magnetic fields

With the advent of magnetic resonance imaging MRI-guided radiation therapy, it is becoming increasingly important to consider the potential influence of a magnetic field on ionising radiation at nanoscale as this can impact the radiobiological effectiveness of the therapeutic radiation beam. Bug et al, [117] and Lazarakis et al. [72] studied the effect of the magnetic field on the track structure of electrons, protons and alpha particles on the cluster size distribution in a DNA segment, modelled as a liquid water cylinder with nanometric sizes. The intensity of the magnetic field was varied from 0 to 14 T. The Geant4-DNA Package was adopted to model the physics interactions of particles in the set-up. The studies showed no significant impact of the magnetic field on the track structure, cluster size distribution and probability of producing a DSB. However, it is important to note that in Geant4 the presence of the magnetic field can impact the trajectory of charged particles only, and not the cross sections of the physical processes. It has been shown that molecules may align with magnetic field [118], while the Geant4-DNA cross sections assume that the hit direction is random. If the hit direction is not random then it could lead to significant differences in cross sections, particularly differential cross sections (kinetic energy of delta electrons etc.).

5. Future

5.1 Directions and Conclusions

One of the main achievements of Geant4-DNA is to extend the general purpose Geant4 Monte Carlo toolkit for the simulation of ionising radiation physical interactions with biological systems at the cellular and DNA level (i.e. micro and nano dosimetry).

As demonstrated in this review, the scope of potential applications includes X-ray radiotherapy, hadron therapy, NP radiosensitisation, targeted cellular therapy and other radiobiology applications. In addition, the extensive Geant4 functionality allowing the modelling of various geometries and the convenient user interface enabling the combination of physics models definitively encourages the use of Geant4-DNA in the medical physics community. Its free availability and open source access undoubtedly supported its widespread adoption.

There is still, however, scope and need for further development. For example, in order to use Geant4-DNA for early damage simulation in radiotherapy, it is pivotal to fully simulate radiation chemistry and the related production of indirect damage. As briefly explained, Geant4-DNA offers a full set of prototype features allowing the simulation of water

1 radiolysis from ionising radiation, up to one microsecond after irradiation. Pachnerova et al.
2 [119] performed the first validation of Geant4-DNA water radiolysis simulation capabilities
3 for incident 15-30 MeV protons. They found an acceptable agreement between simulated
4 radiochemical yields for hydroxyl radicals as a function of time versus experimental
5 measurements. Such validation studies are strongly needed to further evaluate Geant4-DNA
6 modelling accuracy for water radiolysis. Another recent application of Geant4-DNA,
7 including water radiolysis simulation, was recently presented by Tran et al. [120] for the
8 Geant4 simulation of molecular species production around a single 50 nm sphere of gold
9 irradiated by incident MeV protons. The authors estimated the Dose Enhancement Factor and
10 Radiolysis Enhancement Factor of the gold nanosphere compared to a liquid water
11 nanosphere, as a function of radial distance from the sphere; they showed preliminary results
12 demonstrating similar trends for the simulated incident proton energy range and underlined
13 current Geant4 limitations. More realistic mechanistic simulations will be required to fully
14 investigate the possible potential of NPs in radiotherapy, in particular, the development of
15 Geant4-DNA discrete models for the description of electrons interactions in gold NPs.
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19 These recent developments have clearly paved the way for a variety of new Geant4-DNA
20 applications at the nanometer scale requiring the simulation of physical, physicochemical and
21 chemical interactions.
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23 Geant4-DNA operates in a very low energy region that is governed by significant theoretical
24 and experimental complexity. Any theoretical calculations must take into account detailed
25 dielectric structure of the interacting material. Sometimes, for simplification purpose,
26 approximations, assumptions, semi-empirical models must be used. In parallel, experimental
27 measurements below 100 eV are needed to fully validate the models against experimental
28 data. However, such measurements are difficult and are subject to practical constraints. There
29 is a lack of experimental data in liquid water in this energy range and new measurements are
30 still needed. However, while all models have their own limitations, the Geant4-DNA models
31 provide reasonable agreement when compared to literature data and will continue to be
32 improved as new data (e.g. very low energy cross-sections) become available.
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37 In summary, Monte Carlo methods have been explored and used for years as a tool for
38 precise dosimetry as an alternative to analytical methods or for verification of experimental
39 measurements. Geant4-DNA is a very promising and exciting track structure simulation tool
40 for use in radiobiology and micro/nanodosimetry. It is also an excellent example of
41 multidisciplinary and international collaboration involving physicists, theoreticians,
42 radiobiologists, chemists and others from a large number of institutions worldwide. We can
43 be nothing but excited to see further development of this toolkit.
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