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Three-body QED test and fifth-force constraint from vibrations and rotations of HD^+

M. Germann,^{1,*} S. Patra^{1,†} J.-Ph. Karr^{1,2,3} L. Hilico^{1,2,3} V. I. Korobov^{1,4} E. J. Salumbides,¹ K. S. E. Eikema¹,
W. Ubachs¹ and J. C. J. Koelemeij^{1,‡}

¹LaserLaB, Department of Physics and Astronomy, Vrije Universiteit Amsterdam, De Boelelaan 1081, 1081 HV Amsterdam, The Netherlands

²Laboratoire Kastler Brossel, Sorbonne Université, CNRS, ENS-PSL Research University, Collège de France,
4 place Jussieu, 75005 Paris, France

³Université d'Evry-Val d'Essonne, Université Paris-Saclay, Boulevard François Mitterrand, 91000 Evry, France

⁴Bogoliubov Laboratory of Theoretical Physics, Joint Institute for Nuclear Research, Dubna 141980, Russia



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We present a parts-per-million test of quantum electrodynamics (QED) in the HD^+ molecular hydrogen ion, improving on previous tests based on vibrational and rotational transitions by factors of 76 and 1.4, respectively. The test is performed following a unified statistical approach that also produces improved constraints on physics beyond the standard model. We furthermore show how individual constraints derived from the various degrees of freedom in HD^+ and antiprotonic helium can be combined to enhance the sensitivity, thus ruling out “fifth forces” on the Ångström scale that are 10^{11} times weaker than the electromagnetic interaction.

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The standard model (SM) of particle physics is challenged by open questions concerning matter-antimatter asymmetry and the nature of dark matter and dark energy, which motivate searches for new physics (NP) beyond the standard model. Besides experiments at ever increasing interaction energies, precision measurements on particles, atoms, and molecules can provide low-energy SM tests through comparison with comparably accurate predictions obtained by relativistic quantum mechanics and quantum electrodynamics (QED). Deviations from QED predictions may indicate hitherto unknown physics beyond the SM, while agreement with observations allows us to constrain putative new physics [1].

As QED has withstood stringent tests at ever higher precision for single particles (sub-ppb [2–4]) and atoms (0.2 ppm [5,6]), one naturally asks whether it accurately describes more complex forms of matter, such as molecules, as well. Molecular systems, with their additional degrees of freedom, also open up new windows on possible NP [1,7]. In this Letter, we address these concepts in the revised landscape formed by recent extremely precise measurements of rotational [8] and vibrational [9] transition frequencies of the molecular hydrogen ion HD^+ , and of the masses of its constituent particles [10–14]. These results allow us to improve theoretical predictions, and test QED in rotating and vibrating systems with

unprecedented sensitivity. We subsequently combine the new HD^+ results with data from another three-body system, $\bar{p}\text{He}$ [15,16], to establish more stringent bounds on the strength and range of a possible internuclear “fifth force” [7].

Most previous tests of QED in HD^+ were carried out using either rovibrational transitions [17–20] or rotational transitions [21], observed via spectroscopy at the ~ 1 -ppb level of trapped HD^+ ions cooled to 10 mK by laser-cooled beryllium ions. Such tests were based on a qualitative assessment of the level of agreement between the experimental frequencies and frequency values predicted using nonrelativistic QED (NRQED) theory [22,23]. In some of these and other studies involving spectroscopy of light molecular species, the absence of large frequency differences between theory and experiment was interpreted as a constraint on NP such as Ångström-range hadron-hadron interactions and compactified additional dimensions [7,8,19]. Fifth-force constraints exclude a region in the plane spanned by the range and coupling constant of the fifth force, and have so far been considered individually; i.e., no attempts have been made to combine data in order to increase the sensitivity and obtain an improved constraint.

A quantum leap in experimental accuracy was recently achieved in Doppler-free measurements in HD^+ of the $(v, L): (0, 0) \rightarrow (0, 1)$ fundamental rotational transition at 1.31 THz [8], the $(v, L): (0, 3) \rightarrow (9, 3)$ vibrational overtone transition at 415 THz [9], and the $(v, L): (0, 0) \rightarrow (1, 1)$ rotational-vibrational transition at 58.6 THz [24] (here, v and L denote vibrational and rotational quantum numbers, respectively). Of these, only the rotational transition was used to test QED and constrain NP [8], where it was furthermore noted that the sensitivity was limited by the uncertainty of the CODATA-18 value of the proton-electron mass ratio, m_p/m_e , used in the theoretical calculation [25]. This eventually led Alighanbari *et al.* [8] as well as Patra *et al.* [9] and Kortunov *et al.* [24] to assume QED theory to be correct, and exploit theory to

*Present address: Department of Physics, Umeå University, 901 87 Umeå, Sweden.

†Present address: Time and Frequency Division, National Institute of Standards and Technology, Boulder, Colorado 80305, USA.

‡j.c.j.koelemeij@vu.nl

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translate their precise measurement to an improved value of m_p/m_e .

Here, we present an improved test of molecular QED along with a more stringent fifth-force constraint that now also includes the $(v, L): (0, 3) \rightarrow (9, 3)$ transition which, at 3-ppt uncertainty, stands as the most accurate of all measurements performed on HD^+ so far (uncertainties of other measured spin-averaged frequencies are 13 ppt [8] and 15 ppt [24]). To this end we develop a unified statistical approach that ensures consistency between the QED test and the fifth-force constraint given prespecified confidence levels (CLs). Our test and constraint furthermore benefit from updated theoretical predictions reported in this work for both the rotational and vibrational transitions, taking advantage of strongly improved (and independent) values of m_p/m_e and the deuteron-proton mass ratio m_d/m_p [10–14] which reduce the uncertainty of the theoretical predictions by about 25%. The theoretical predictions also include previously unevaluated terms of order $R_\infty\alpha^5$ and $R_\infty\alpha^6$ (with R_∞ the Rydberg constant and α the fine-structure constant) [6,26–28], and a contribution from the deuteron polarizability [6,26]. These theory improvements are described in the Supplemental Material [29], where the resulting theoretical transition frequencies are listed together with the previous values. The new values have a relative uncertainty that approaches 1×10^{-11} , making them the most precise theoretical predictions for any molecular property published to date. We note that similar theory updates were reported very recently by Kortunov *et al.* for the $(v, L): (0, 0) \rightarrow (1, 1)$ vibrational transition [24]. However, in what follows we will focus primarily on the $(v, L): (0, 3) \rightarrow (9, 3)$ transition as a benchmark system for molecular vibrations, owing to its significantly higher experimental and theoretical accuracy (with the latter stemming from a lower sensitivity to the uncertainties of fundamental constants).

Figure 1 shows how the updated $(v, L): (0, 3) \rightarrow (9, 3)$ theoretical transition frequency, ν_{theo} , is broken down into a nonrelativistic (NR) contribution as well as a number of relativistic and QED corrections (including nuclear recoil). Purely relativistic corrections [30,31] appear at even powers in α ($R_\infty\alpha^2, R_\infty\alpha^4, R_\infty\alpha^6, \dots$, the latter contribution being negligibly small). QED corrections appear at order $R_\infty\alpha^3$ [32] and are considered up to order $R_\infty\alpha^6$ [22,30,33]. ν_{theo} also comprises a small correction to account for muonic and hadronic vacuum polarization. The NR term is calculated with negligible uncertainty by solving the three-body Schrödinger equation by a variational method [34,35]. Together with this part, we include a correction for the combined proton and deuteron finite nuclear size of $-830.6(5)$ kHz, computed using CODATA-18 values of the nuclear charge radii [25]. Besides the NR term, the leading-order relativistic ($R_\infty\alpha^2$) and QED ($R_\infty\alpha^3$) terms are also known with negligible uncertainty, and the theoretical uncertainty is caused almost entirely by the uncalculated $R_\infty\alpha^6$ -order QED corrections (see inset of Fig. 1). The largest contribution to the total theoretical uncertainty is from the uncertainty of physical constants (mostly m_p/m_e) [9].

Experimental values of the $(v, L): (0, 0) \rightarrow (0, 1)$ rotational and $(v, L): (0, 3) \rightarrow (9, 3)$ vibrational transition frequencies have previously been obtained by Doppler-free terahertz and laser spectroscopy, respectively (see Refs. [8,9]

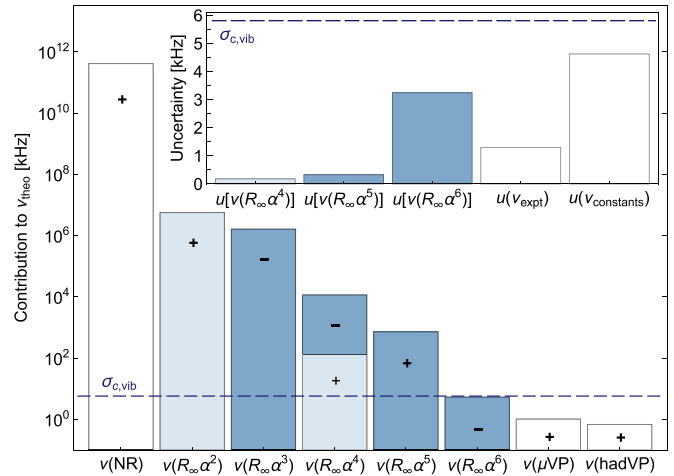


FIG. 1. Sizes and signs of contributions of order $R_\infty\alpha^n$ to the SM prediction of the $(v, L): (0, 3) \rightarrow (9, 3)$ transition frequency ν_{theo} . Light blue, pure relativistic terms; dark blue, QED corrections. Contributions from muonic and hadronic vacuum polarization are also shown. The combined uncertainty (dashed horizontal line) is comparable to the $R_\infty\alpha^6$ -order QED correction. Inset: Frequency uncertainties, u , of the contributions of orders $R_\infty\alpha^4$ through $R_\infty\alpha^6$ (the lower-order and NR terms are essentially exact). The experimental uncertainty, $u(\nu_{\text{expt}})$, is small enough to test state-of-the-art three-body QED theory at its very limit. The quadratic sum of all uncertainties shown here, including the contribution from the physical constants, $u(\nu_{\text{constants}})$, equals $\sigma_{c,vib}$.

for details). These values correspond to spin-averaged frequencies; i.e., they were obtained after correcting measured transition frequencies of various hyperfine components for hyperfine shifts, which were found theoretically with sufficient precision. We here apply updated theoretical hyperfine shifts [36,37] which shift the experimental $(v, L): (0, 3) \rightarrow (9, 3)$ transition frequency upward by 1.3 kHz, and increase the uncertainty slightly from 1.2 to 1.3 kHz [29]. For the $(v, L): (0, 0) \rightarrow (0, 1)$ transition frequency, the corresponding shift is negligibly small (less than 0.1σ), and we therefore retain the experimental frequency value reported by Alighanbari *et al.* [8]. These experimental values can be compared with the theoretical ones with a precision given by the combined uncertainty, σ_c , which comprises the uncertainties of the experiment, the physical constants used in the theory, and the uncertainties of the QED terms (Fig. 1). Given the relative uncertainties of the experimental rotational and vibrational overtone transition frequencies (13 and 3.1 ppt, respectively), molecular theory and experiment can now be compared at unprecedented relative combined uncertainties of 37 ppt for the $(v, L): (0, 0) \rightarrow (0, 1)$ and 14 ppt for the $(v, L): (0, 3) \rightarrow (9, 3)$ transition.

To test the HD^+ theory we go beyond the qualitative character of previous QED tests [8,17–21] and set up a binary hypothesis test characterized by a significance level, α_s , and statistical power, $1 - \beta$. α_s represents the probability of a type-I error (i.e., inferring that an effect beyond the theory exists while in reality it does not exist), and β the probability of a type-II error (i.e., judging the theory to be correct while a significant effect beyond the theory is present in reality).

As explained in the Supplemental Material [29], we perform the test for $\alpha_s = 0.0027$, thus checking for deviations between theory and experiment larger than $3\sigma_c$. The found differences (experiment minus theory) are $-2.1\sigma_{c,\text{rot}}$ and $-1.7\sigma_{c,\text{vib}}$ for the rotational and vibrational transitions, respectively [29], leading us to conclude that at 99.7% CL, the molecular theory is valid for rotations and vibrations at the level (3σ) of 110 and 42 ppt, respectively. For the theory of molecular vibrations, this result is 76 times as stringent as the previous best test [19].

The hypothesis test carried out above can be also be interpreted as a test of the QED part of the calculation, which we will do with special attention to the $(v, L): (0, 3) \rightarrow (9, 3)$ transition which so far has not been considered in this context. The QED terms contribute in their entirety -1.7 GHz to the vibrational transition frequency, and the uncertainty in the theory stems almost entirely from the QED terms (Fig. 1). The agreement between theory and experiment found above can therefore be seen as a test of high-order molecular QED, based on a vibrational transition, at the level (1σ) of $(5.8 \text{ kHz})/(1.7 \text{ GHz}) = 3.5$ ppm. These findings represent a remarkable 76-fold improvement over the previous most precise QED test for molecular vibrations by Biesheuvel *et al.* [19], and the sensitivity may be further reduced to ≈ 0.3 ppm if the precision of the theory and the constants are improved beyond the experimental uncertainty.

For the $(v, L): (0, 0) \rightarrow (0, 1)$ rotational transition in HD^+ [8], our improved theoretical prediction improves the stringency of the QED test from 7.0 to 5.1 ppm (at 1σ). It is worth noting that this test performed in a molecule exhibits qualitative differences with respect to those done, e.g., in hydrogenlike atoms, due to the existence of QED corrections specifically linked to the molecular nature of the system (see Supplemental Material [29] for details).

Our hypothesis test can also be used to derive bounds on possible frequency shifts caused by NP [29]. One advantage of this unified statistical approach is that the CL of such bounds is consistent with the CL specified for the QED test [29]. The location of the bound and the corresponding CL both depend on α_s and β , and for $\alpha_s = 0.0027$ and $\beta = 0.1$ we find a 90%-CL constraint at $4.3\sigma_c$ [29]. Our approach reveals that the CLs stated for some previous constraints were overestimated, with CLs being 50% rather than the claimed 90% [29]. In what follows we specifically consider hypothetical scalar hadron-hadron interactions at the $0.1\text{--}1 \text{ \AA}$ length scale, which is typical for three-body systems like HD^+ and pHe . This range lies in between the scales investigated in high-energy and nuclear physics experiments, and those studied with atoms in weakly bound dimers [38] and macroscopic Cavendish-type setups. Following Salumbides *et al.* [7], we model such a hypothetical ‘‘fifth force’’ phenomenologically through a Yukawa-type potential,

$$V_{\alpha_s, \lambda}(r) = \hbar c \alpha_5 A_1 A_2 \frac{e^{-r/\lambda}}{r}. \quad (1)$$

Here, r is the separation of the two nuclei containing the interacting hadrons, α_5 represents a dimensionless coupling constant, and \hbar and c are the reduced Planck constant and the speed of light, respectively. The range, λ , of the hypothetical interaction is related to the mass, m_5 , of the corresponding

hypothetical force carrier through $\lambda = \hbar/(m_5 c)$. We treat the nucleon numbers A_1 and A_2 as charges of the hypothetical interaction.

The interaction in Eq. (1) leads to a first-order frequency shift of a transition $(v, L) \rightarrow (v', L')$,

$$\begin{aligned} \Delta\nu_{\alpha_s, \lambda}^{vv'LL'} &= (\Delta E_{\alpha_s, \lambda}^{v'L'} - \Delta E_{\alpha_s, \lambda}^{vL}) / (2\pi \hbar) \\ &\equiv \alpha_5 A_1 A_2 \mathcal{Y}_{\lambda}^{vv'LL'}, \end{aligned} \quad (2)$$

where $\Delta E_{\alpha_s, \lambda}^{vL} = \langle vL | V_{\alpha_s, \lambda} | vL \rangle$, and $\mathcal{Y}_{\lambda}^{vv'LL'}$ is a normalized frequency shift that was computed on a dense grid of λ values for relevant transitions in HD^+ and pHe [7,29]. If theory and experiment are found to be in agreement, our hypothesis test constrains, for a given λ , the NP parameter $|\alpha_5|$ through the inequality

$$|\alpha_5| < \left| \frac{4.3\sigma_c}{A_1 A_2 \mathcal{Y}_{\lambda}^{vv'LL'}} \right|. \quad (3)$$

We apply Eq. (3) to the three recent measurements in HD^+ [8,9,24], as well as four transitions in p^3He and p^4He [15,16], which is justified as no deviations beyond 3σ are observed for these transitions. The resulting 90%-CL constraints are shown in Fig. 2.

The various transitions in Fig. 2 have diverse sensitivities with respect to λ . In the Supplemental Material [29] we show that in the (hypothetical) case of a detectable fifth-force frequency shift, the additional degrees of freedom offered by the various transitions and three-body systems studied here can be used to obtain separate estimates of the coupling constant as well as the range of a fifth force. It is also worth noting that the sensitivities $\mathcal{Y}_{\lambda}^{vv'LL'}$ of vibrational transitions change sign at shorter range (corresponding to the peaks in the constraints shown in Fig. 2), which can lead to an increased differential fifth-force sensitivity among different transitions, and thereby enhance the resolution at which the range of a hypothetical fifth force may be determined [29].

Further, the above data can be combined into a single, more severe constraint. To this end, we fit α_5 (for given λ) to multiple data simultaneously, through a least-squares optimization. At long range, the most constraining input data are the three recent HD^+ measurements. The corresponding constraints are partially correlated through the QED theory and the (correlated) physical constants used; all these correlations are taken into account [29]. To determine the combined constraint, we proceed in the same spirit as our treatment of single comparisons above. The least-squares fit produces an estimate $\hat{\alpha}_5$ with 1σ uncertainty $u(\hat{\alpha}_5)$. We again carry out a hypothesis test, this time based on the value of $\hat{\alpha}_5/u(\hat{\alpha}_5)$. We find no deviations at the 3σ level, and the corresponding 90%-CL ($4.3\sigma_c$) constraint is more stringent (by up to a factor of 3) than any of the individual constraints (Fig. 2). The impact of correlations is also visible in Fig. 2. The combined constraint clearly benefits from the multiple HD^+ and (uncorrelated) pHe data near $\lambda = 0.15 \text{ \AA}$, but the strong correlation between the HD^+ rotational and vibrational data prevents a significant improvement near $\lambda = 0.45 \text{ \AA}$. We expect our approach to become useful for future NP searches as it is generally applicable and may incorporate any number of constraints, obtained from any system.

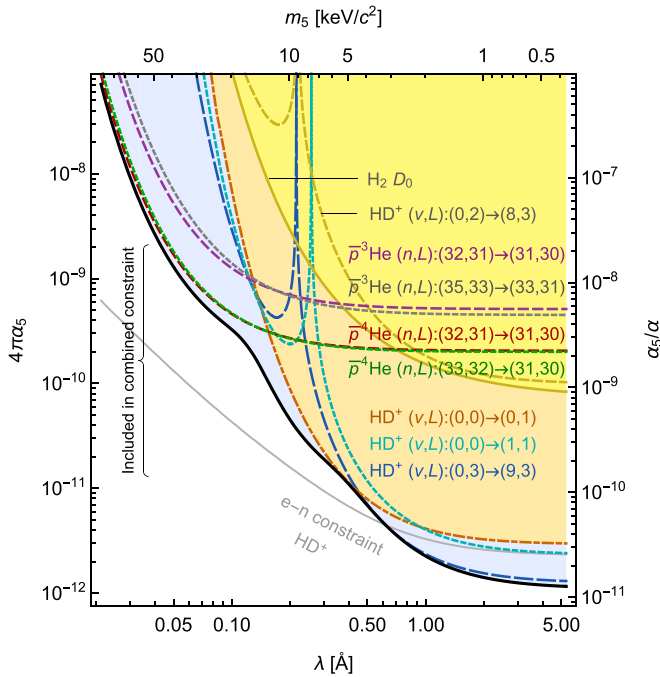


FIG. 2. 90%-CL constraints on scalar hadron-hadron interactions inferred from a hypothesis test with significance level 0.0027. Individual constraints are obtained from rotational and vibrational spectroscopy of HD^+ and from studies of $\bar{p}^3\text{He}$ and $\bar{p}^4\text{He}$ [8,9,15,16,24]. Black curve: Combined constraint obtained from a least-squares adjustment involving seven transitions. The upper yellow region indicates (λ, α_5) values excluded by the H_2 dissociation limit (solid yellow curve [39]) and the previous best constraint from vibrational spectroscopy of HD^+ (dashed yellow curve [19]). Below, the orange area represents the parameter space excluded by the work of Alighanbari *et al.* [8]. Our combined constraint expands the excluded area further into the light blue region. The solid gray curve illustrates the constraint on a hypothetical electron-neutron coupling based on the three recent HD^+ measurements [8,9,24].

At long range ($\lambda > 5 \text{ \AA}$), we find $4\pi\alpha_5 < 1 \times 10^{-12}$ at 90% CL. It is interesting to compare this with results from other NP searches on the atomic scale. One example is the search for Yukawa-type electron-neutron (e - n) couplings by analysis of data from isotope-shift spectroscopy and so-called King plots [40–43]. Such e - n constraints cannot be directly compared with our nucleon-nucleon constraint; however, we have also computed e - n Yukawa matrix elements similar to those in Eq. (2) by integrating $\hbar c \alpha_5 \exp(-r_{de}/\lambda)/r_{de}$, with r_{de} the deuteron-electron distance, over the nonrelativistic three-body wave functions. We use these matrix elements to derive

constraints from the three most recent HD^+ measurements, and inserting these constraints in a separate adjustment (assuming no electron-proton or neutron-proton coupling) leads to the e - n constraint shown in Fig. 2. This constraint is nearly two orders of magnitude more stringent than the recent King-plot results [42,43], and within one order of magnitude from the more stringent bound obtained from hydrogen-deuterium isotope shift spectroscopy [41]. However, it must be realized that our e - n constraint depends linearly on R_∞ , whose value is determined from spectroscopy of hydrogenlike atoms [25] which might be affected by the very e - n coupling we are trying to constrain. Note that this is not an issue for the Ångström-range hadron-hadron (or nucleon-nucleon) interactions considered in this work, since such interactions are absent in atomic systems. Future work in this direction should consider also the effect of NP on the physical constants used in the theory.

We may compare our nucleon-nucleon constraints with those inferred from neutron scattering [29,44,45], which are still two orders of magnitude more stringent than ours. However, a crucial difference is that our models for HD^+ and $\bar{p}\text{He}$, as search grounds for NP, are built up from first principles and do not rely on model assumptions except those of relativistic quantum mechanics and QED. A null result obtained with these systems can therefore constrain a wider class of possible NP phenomena. For example, King-plot searches for NP in atoms are necessarily restricted to electron-neutron interactions subject to certain model assumptions regarding atomic and nuclear structure [41–43], whereas our systems can be used to constrain electron-neutron as well as electron-proton and proton-neutron interactions in a model-independent way. We thus conclude that our work not only rules out scalar hadron-hadron interactions at the Ångström scale that are 1×10^{11} times weaker than the electromagnetic interaction (Fig. 2), but generally excludes NP that would modify the frequencies of molecular vibrations and rotations in HD^+ in the 11th decimal place, thereby providing evidence for the validity of the SM (and QED in particular) in molecules at the same level.

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