Tensor algebra and analysis

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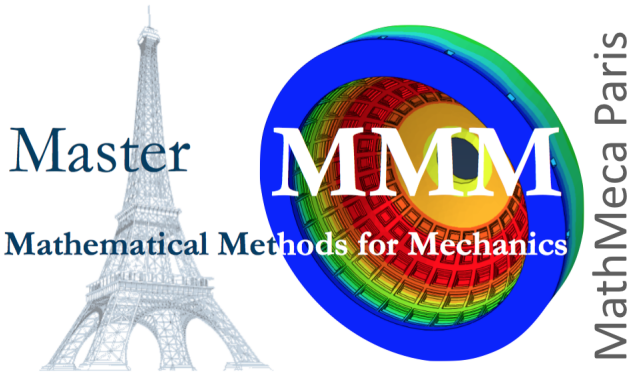
A short course on

Tensor algebra and analysis



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Preface

This text is a support for different courses of the master of Mechanics of the University Paris-Saclay.

The content of this text is an introduction, for graduate students, to modern tensor algebra and analysis, specially intended for applications in continuum mechanics.

Far from being exhaustive, the text focuses on some subjects, with the intention of providing the reader with the main algebraic tools necessary for a modern course in continuum mechanics.

Versailles, July 25, 2021

Contents

Chapter 1

Points and vectors

1.1 Points and vectors

We consider in the following a *point space* $\mathcal{E}$, whose elements are points p[1]. On $\mathcal{E}$ we admit the existence of an operation, the *difference* of any of its two elements:

$$q - p,\ p, q \in \mathcal{E}.$$

We associate to $\mathcal{E}$ a *vector space* $\mathcal{V}$ whose dimension is $\dim\mathcal{V} = 3$ and whose elements are vectors $\mathbf{v}$ representing *translations over* $\mathcal{E}$:

$$\forall p, q \in \mathcal{E},\ \exists!\ \mathbf{v} \in \mathcal{V} |\ q - p = \mathbf{v}.$$

Any element $\mathbf{v} \in \mathcal{V}$ is hence a transformation over $\mathcal{E}$ that can be written, using the previous definition, as :

$$\forall \mathbf{v} \in \mathcal{V},\ \mathbf{v} : \mathcal{E} \to \mathcal{E} |\ q = \mathbf{v}(p)\ \ \to\ \ q = p + \mathbf{v}.$$

To remark that the result of the application of the translation $\mathbf{v}$ depends upon the argument p:

$$q = p + \mathbf{v} \neq p' + \mathbf{v} = q',$$

whose geometric meaning is depicted in Fig. 1.1. Unlike difference, the sum of two points is not defined and is meaningless.

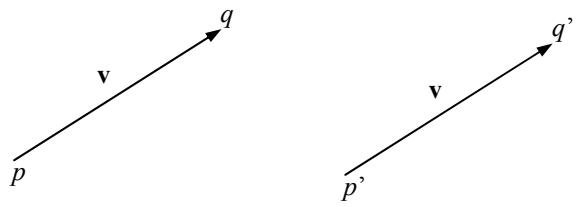

Figure 1.1: Same translation over two different points.

[1] $\mathcal{E}$ is to be identified with the *Euclidean three-dimensional space* in which the events of classical mechanics are intended to be set.

We define the *sum of two vectors* $\mathbf{u}$ and $\mathbf{v}$ as the vector $\mathbf{w}$ such that

$$(\mathbf{u}+\mathbf{v})(p) = \mathbf{u}(\mathbf{v}(p)) = \mathbf{w}(p)$$

This means that, if

$$q = \mathbf{v}(p) = p + \mathbf{v},$$

then

$$r = \mathbf{u}(q) = q + \mathbf{u} = \mathbf{w}(p),$$

see Fig. 1.2, which shows that the above definition actually coincides with the parallelogram rule and that

$$\mathbf{u} + \mathbf{v} = \mathbf{v} + \mathbf{u},$$

as obvious, for the sum over a vector space commutes. It is evident that the sum of more than two vectors can be defined iteratively, summing up a vector at time to the sum of the previous vectors.

The *null vector* $\mathbf{o}$ is defined as the difference of two coincident points:

$$\mathbf{o} = p - p \quad \forall p \in \mathcal{E};$$

$\mathbf{o}$ is unique and the only vector such that

$$\mathbf{v} + \mathbf{o} = \mathbf{v} \quad \forall \mathbf{v} \in \mathcal{V}.$$

In fact:

$$\forall p \in \mathcal{E}, \quad \mathbf{v} + \mathbf{o} = \mathbf{v} + p - p \;\rightarrow\; p + \mathbf{v} + \mathbf{o} = p + \mathbf{v} \;\iff\; \mathbf{v} + \mathbf{o} = \mathbf{v}.$$

A *linear combination* of n vectors $\mathbf{v}_i$ is defined as the vector[2]

$$\mathbf{w} = k_i \mathbf{v}_i, \quad k_i \in \mathbb{R}, \quad i = 1, ..., n.$$

The $n+1$ vectors $\mathbf{w}$, $\mathbf{v}_i$, $i = 1, ..., n$, are said to be *linear independent* if it does not exist a set of n scalars k_i such that the above equation is satisfied, *linear dependent* in the opposite case.

[2]We adopt here and in the following the *Einstein notation* for summations: all the times that an index is repeated in a monomial, then summation with respect to that index, called the *dummy index*, is understood. If a repeated index is underlined, then it is not a dummy index, i.e. there is no summation.

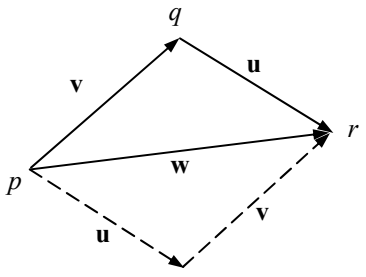

Figure 1.2: Sum of two vectors: the parallelogram rule.

1.2 Scalar product, distance, orthogonality

A *scalar product* on a vector space is a positive definite, symmetric, bilinear form. A *form* ω is a function

$$\omega : \mathcal{V} \times \mathcal{V} \to \mathbb{R},$$

i.e., ω operates on a couple of vectors to give a real number, a *scalar*. We will indicate the scalar product of two vectors $\mathbf{u}$ and $\mathbf{v}$ as

$$\omega(\mathbf{u}, \mathbf{v}) = \mathbf{u} \cdot \mathbf{v}.$$

The properties of *bilinearity* prescribe that, $\forall \mathbf{u}, \mathbf{v} \in \mathcal{V}$ and $\forall \alpha, \beta \in \mathbb{R}$,

$$\begin{aligned} \mathbf{u} \cdot (\alpha\mathbf{v} + \beta\mathbf{w}) &= \alpha\mathbf{u} \cdot \mathbf{v} + \beta\mathbf{u} \cdot \mathbf{w}, \\ (\alpha\mathbf{u} + \beta\mathbf{v}) \cdot \mathbf{w} &= \alpha\mathbf{u} \cdot \mathbf{w} + \beta\mathbf{v} \cdot \mathbf{w}, \end{aligned}$$

while *symmetry* implies that

$$\mathbf{u} \cdot \mathbf{v} = \mathbf{v} \cdot \mathbf{u} \quad \forall \mathbf{u}, \mathbf{v} \in \mathcal{V}.$$

Finally, the *positive definiteness* means that

$$\mathbf{v} \cdot \mathbf{v} > 0 \quad \forall \mathbf{v} \in \mathcal{V}, \quad \mathbf{v} \cdot \mathbf{v} = 0 \iff \mathbf{v} = \mathbf{o}.$$

Any two vectors are said to be *orthogonal* $\iff$

$$\mathbf{u} \cdot \mathbf{v} = 0.$$

Thanks to the properties of the scalar product, we can define the *Euclidean norm* of a vector $\mathbf{v}$ as the nonnegative scalar, denoted equivalently by v or $|\mathbf{v}|$,

$$v = |\mathbf{v}| = \sqrt{\mathbf{v} \cdot \mathbf{v}}$$

The norm of a vector has the following properties:

$$\begin{aligned} &|\mathbf{u} + \mathbf{v}| \leq u + v \quad \text{(Minkowski's triangular inequality)}; \\ &|\mathbf{u} \cdot \mathbf{v}| \leq u\, v \quad \text{(Schwarz's inequality)}; \\ &|k\mathbf{v}| = |k|v, \ k \in \mathbb{R}. \end{aligned}$$

We define *distance between two any points* p and $q \in \mathcal{E}$ the scalar

$$d(p, q) = |p - q| = |q - p|.$$

Similarly, the *distance between two any vectors* $\mathbf{u}$ and $\mathbf{v} \in \mathcal{V}$ is defined as

$$d(\mathbf{u}, \mathbf{v}) = |\mathbf{u} - \mathbf{v}| = |\mathbf{v} - \mathbf{u}|.$$

Two points or two vectors are *coincident* if and only if their distance is null.

The *unit sphere* $\mathcal{S}$ *of* $\mathcal{V}$ is defined as the set of all the vectors whose norm is one:

$$\mathcal{S} = \{\mathbf{v} \in \mathcal{V} |\ v = 1\}.$$

1.3 Basis of $\mathcal{V}$, expression of the scalar product

Generally speaking, a *basis* $\mathcal{B}$ of a vector space is any set of n linearly independent vectors, where n is equal to the dimension of the vector space. In the case of $\mathcal{V}$, $n = 3$, so that a basis $\mathcal{B}$ of $\mathcal{V}$ is any set

$$\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\},$$

of three linearly independent vectors $\mathbf{e}_i$. The concept of *basis of* $\mathcal{V}$ is useful for representing vectors: once a basis chosen, any vector $\mathbf{v} \in \mathcal{V}$ can be represented as a linear combination of the vectors of the basis, where the coefficients v_i of the linear combination are the *components* of $\mathbf{v}$:

$$\mathbf{v} = v_i\mathbf{e}_i = v_1\mathbf{e}_1 + v_2\mathbf{e}_2 + v_3\mathbf{e}_3.$$

Though the choice of the elements of a basis is completely arbitrary, the only condition being their linear independency, we will use in the following only *orthonormal bases*, that are bases composed by mutually orthogonal vectors of $\mathcal{S}$, i.e. satisfying

$$\mathbf{e}_i \cdot \mathbf{e}_j = \delta_{ij},$$

where the symbol δ_{ij} is the so-called *Kronecker's delta*:

$$\delta_{ij} = \begin{cases} 1 \text{ if } i = j, \\ 0 \text{ if } i \neq j. \end{cases}$$

The use of orthonormal bases has great advantages; for instance, it allows to give a very simple rule for the calculation of the scalar product:

$$\mathbf{u} \cdot \mathbf{v} = u_i\mathbf{e}_i \cdot v_j\mathbf{e}_j = u_iv_j\delta_{ij} = u_iv_i = u_1v_1 + u_2v_2 + u_3v_3.$$

In particular, it is

$$\mathbf{v} \cdot \mathbf{e}_i = v_k\mathbf{e}_k \cdot \mathbf{e}_i = v_k\delta_{ik} = v_i, \quad i = 1, 2, 3.$$

So, the *Cartesian components* of a vector are the projection of the vector on the three vectors of the basis $\mathcal{B}$; such quantities are the *director cosines* of $\mathbf{v}$ in the basis $\mathcal{B}$. In fact, if θ is the angle formed by two vectors $\mathbf{u}$ and $\mathbf{v}$, then

$$\mathbf{u} \cdot \mathbf{v} = u\ v\ \cos\theta.$$

This relation is used to define the angle between two vectors,

$$\theta = \arccos\frac{\mathbf{u} \cdot \mathbf{v}}{u\ v},$$

which can be proved easily: given two vectors $\mathbf{u}$ and $\mathbf{v}$, we look for $c \in \mathbb{R}$ such that the vector $\mathbf{u} - c\mathbf{v}$ be orthogonal to $\mathbf{v}$:

$$(\mathbf{u} - c\mathbf{v}) \cdot \mathbf{v} = 0 \iff c = \frac{\mathbf{u} \cdot \mathbf{v}}{\mathbf{v} \cdot \mathbf{v}} = \frac{\mathbf{u} \cdot \mathbf{v}}{v^2}.$$

Now, if $\mathbf{u}$ is inclined of θ on $\mathbf{v}$, its projection u_v on the direction of $\mathbf{v}$ is

$$u_v = u\ \cos\theta,$$

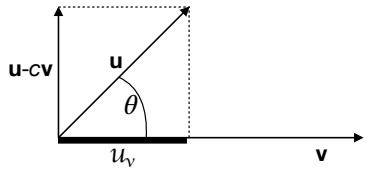

Figure 1.3: Angle between two vectors.

and, by construction (see Fig. 1.3), it is also

$$u_v = c\ v.$$

So

$$c = \frac{u}{v}\cos\theta \ \rightarrow\ \frac{u}{v}\cos\theta = \frac{\mathbf{u}\cdot\mathbf{v}}{v^2} \ \Rightarrow\ \cos\theta = \frac{\mathbf{u}\cdot\mathbf{v}}{u\ v}.$$

To remark that, while the scalar product, being an intrinsic operation, does not change for a change of basis, the components v_i of a vector are *not* intrinsic quantities, but they are basis-dependent: a change of the basis makes the components change. The way this change is done will be introduced in Sect. 2.11.

A *frame* $\mathcal{R}$ for $\mathcal{E}$ is composed by a point $o \in \mathcal{E}$, the *origin*, and a basis $\mathcal{B}$ of $\mathcal{V}$:

$$\mathcal{R} = \{o, \mathcal{B}\} = \{o; \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}.$$

The use of a frame for $\mathcal{E}$ is useful for determining the position of a point p, which can be done through its *Cartesian coordinates* x_i, defined as the components, in $\mathcal{B}$, of the vector $p - o$:

$$x_i = (p - o)\cdot\mathbf{e}_i, \quad i = 1, 2, 3.$$

Of course, the coordinates depend upon both the choices of o and of $\mathcal{B}$.

1.4 Exercices

1. Prove that the null vector is unique.
2. Prove that the norm of the null vector is zero.
3. Prove the inequality of Minkowski.
4. Prove the inequality of Schwarz.
5. Prove that
$$\mathbf{u}\cdot\mathbf{v} = 0 \iff |\mathbf{u} - \mathbf{v}| = |\mathbf{u} + \mathbf{v}| \ \ \forall \mathbf{u}, \mathbf{v} \in \mathcal{V}.$$
6. Prove the *linear forms representation theorem*: be $\psi : \mathcal{V} \rightarrow \mathbb{R}$ a linear function. Then, $\exists!\ \mathbf{u} \in \mathcal{V}$ such that
$$\psi(\mathbf{v}) = \mathbf{u}\cdot\mathbf{v} \ \ \forall \mathbf{v} \in \mathcal{V}.$$

Chapter 2

Second rank tensors

2.1 Second-rank tensors

A *second-rank tensor* $\mathbf{L}$ is any linear application from $\mathcal{V}$ to $\mathcal{V}$:

$$\mathbf{L} : \mathcal{V} \to \mathcal{V} \mid \mathbf{L}(\alpha_i \mathbf{u}_i) = \alpha_i \mathbf{L}\mathbf{u}_i \ \forall \alpha_i \in \mathbb{R},\ \mathbf{u}_i \in \mathcal{V},\ i = 1, ..., n.$$

Though here $\mathcal{V}$ indicates the vector space of translations over $\mathcal{E}$, the definition of tensor[1] is more general and in particular $\mathcal{V}$ can be any vector space.

Defining the *sum of two tensors* as

$$(\mathbf{L}_1 + \mathbf{L}_2)\mathbf{u} = \mathbf{L}_1\mathbf{u} + \mathbf{L}_2\mathbf{u} \quad \forall \mathbf{u} \in \mathcal{V}, \tag{2.1}$$

the *product of a scalar by a tensor* as

$$(\alpha \mathbf{L})\mathbf{u} = \alpha(\mathbf{L}\mathbf{u}) \quad \forall \alpha \in \mathbb{R}, \mathbf{u} \in \mathcal{V}$$

and the *null tensor* $\mathbf{O}$ as the unique tensor such that

$$\mathbf{O}\mathbf{u} = \mathbf{o}\ \forall \mathbf{u} \in \mathcal{V},$$

then the set of all the tensors $\mathbf{L}$ that operate on $\mathcal{V}$ forms a vector space, denoted by $Lin(\mathcal{V})$. We define the *identity tensor* $\mathbf{I}$ as the unique tensor such that

$$\mathbf{I}\mathbf{u} = \mathbf{u} \quad \forall \mathbf{u} \in \mathcal{V}.$$

Different possible operations can be defined for second-rank tensors. We consider all of them in the following Sections.

[1] We consider for the while only second-rank tensors, but we will see in the following how to introduce tensors of higher ranks.

2.2 Dyads, tensor components

For any couple of vectors $\mathbf{u}$ and $\mathbf{v}$, the *dyad*[2] $\mathbf{u} \otimes \mathbf{v}$ is the tensor defined by

$$(\mathbf{u} \otimes \mathbf{v})\mathbf{w} = \mathbf{v} \cdot \mathbf{w}\ \mathbf{u} \quad \forall \mathbf{w} \in \mathcal{V}.$$

The application defined above is actually a tensor because of the bi-linearity of the scalar product. The introduction of dyads allows for expressing any tensor as a linear combination of dyads. In fact, it can be proved that if $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ is a basis of $\mathcal{V}$, then the set of 9 dyads

$$\mathcal{B}^2 = \{\mathbf{e}_i \otimes \mathbf{e}_j,\ i, j = 1, 2, 3\},$$

is a basis of $Lin(\mathcal{V})$, i.e. $\dim(Lin(\mathcal{V})) = 9$. This implies that any tensor $\mathbf{L} \in Lin(\mathcal{V})$ can be expressed as

$$\mathbf{L} = L_{ij}\ \mathbf{e}_i \otimes \mathbf{e}_j, \quad i, j = 1, 2, 3,$$

where the L_{ij}s are the nine *Cartesian components* of $\mathbf{L}$ with respect to $\mathcal{B}^2$. The L_{ij}s can be calculated easily:

$$\mathbf{e}_i \cdot \mathbf{L}\mathbf{e}_j = \mathbf{e}_i \cdot L_{hk}\mathbf{e}_h \otimes \mathbf{e}_k\ \mathbf{e}_j = L_{hk}\mathbf{e}_i \cdot \mathbf{e}_h\ \mathbf{e}_k \cdot \mathbf{e}_j = L_{hk}\delta_{ih}\delta_{jk} = L_{ij}.$$

The above expression is sometimes called the *canonical decomposition of a tensor*. The components of a dyad can be computed easily:

$$(\mathbf{u} \otimes \mathbf{v})_{ij} = \mathbf{e}_i \cdot (\mathbf{u} \otimes \mathbf{v})\ \mathbf{e}_j = \mathbf{u} \cdot \mathbf{e}_i\ \mathbf{v} \cdot \mathbf{e}_j = u_i\ v_j. \tag{2.2}$$

The components of a vector $\mathbf{v}$ result of the application of a tensor $\mathbf{L}$ on a vector $\mathbf{u}$ can now be easily calculated:

$$\mathbf{v} = \mathbf{L}\mathbf{u} = L_{ij}(\mathbf{e}_i \otimes \mathbf{e}_j)(u_k\mathbf{e}_k) = L_{ij}u_k\delta_{jk}\mathbf{e}_i = L_{ij}u_j\mathbf{e}_i \ \rightarrow\ v_i = L_{ij}u_j. \tag{2.3}$$

Depending upon two indices, any second-rank tensor $\mathbf{L}$ can be represented by a matrix, whose entries are the Cartesian components of $\mathbf{L}$ in the basis $\mathcal{B}$:

$$\mathbf{L} = \begin{bmatrix} L_{11} & L_{12} & L_{13} \\ L_{21} & L_{22} & L_{23} \\ L_{31} & L_{32} & L_{33} \end{bmatrix};$$

because any $\mathbf{u} \in \mathcal{V}$, depending upon only one index, can be represented by a column vector, eq. (2.3) represents actually the classical operation of the multiplication of a 3×3 matrix by a 3×1 vector.

2.3 Tensor product

The *tensor product* of $\mathbf{L}_1$ and $\mathbf{L}_2 \in Lin(\mathcal{V})$ is defined as

$$(\mathbf{L}_1\mathbf{L}_2)\mathbf{v} = \mathbf{L}_1(\mathbf{L}_2\mathbf{v}) \quad \forall \mathbf{v} \in \mathcal{V}.$$

[2] In some texts, the dyad is also called *tensor product*; we prefer to use the term dyad because tensor product can be ambiguous, as used to denote the product of two tensors, see Sect. 2.3.

By linearity and eq. (2.1) we get

$$[\mathbf{L}(\mathbf{L}_1+\mathbf{L}_2)]\mathbf{v}=\mathbf{L}[(\mathbf{L}_1+\mathbf{L}_2)\mathbf{v}]=\mathbf{L}(\mathbf{L}_1\mathbf{v}+\mathbf{L}_2\mathbf{v})=$$
$$\mathbf{L}\mathbf{L}_1\mathbf{v}+\mathbf{L}\mathbf{L}_2\mathbf{v}=(\mathbf{L}\mathbf{L}_1+\mathbf{L}\mathbf{L}_2)\mathbf{v}\ \rightarrow\ \mathbf{L}(\mathbf{L}_1+\mathbf{L}_2)=\mathbf{L}\mathbf{L}_1+\mathbf{L}\mathbf{L}_2.$$

To remark that the tensor product is not symmetric:

$$\mathbf{L}_1\mathbf{L}_2\neq\mathbf{L}_2\mathbf{L}_1;$$

however, by the same definition of the identity tensor and of tensor product,

$$\mathbf{I}\mathbf{L}=\mathbf{L}\mathbf{I}=\mathbf{L}\ \forall\mathbf{L}\in Lin(\mathcal{V}).$$

The Cartesian components of a tensor $\mathbf{L}=\mathbf{A}\mathbf{B}$ can be easily calculated using eq. (2.3):

$$L_{ij}=\mathbf{e}_1\cdot(\mathbf{A}\mathbf{B})\mathbf{e}_j=\mathbf{e}_i\cdot\mathbf{A}(\mathbf{B}\mathbf{e}_j)=\mathbf{e}_i\cdot\mathbf{A}(B_{hk}(\mathbf{e}_j)_k\ \mathbf{e}_h)=B_{hk}\delta_{jk}\mathbf{e}_i\cdot\mathbf{A}\mathbf{e}_h$$
$$=B_{hk}\delta_{jk}\mathbf{e}_i\cdot(A_{pq}(\mathbf{e}_h)_q\ \mathbf{e}_p)=A_{pq}B_{hk}\delta_{jk}\delta_{qh}\delta_{ip}=A_{ih}B_{hj}.$$

The above result simply corresponds to the rule of the multiplication of rows by lines of two matrices. Using it, the following two identities can be readily shown:

$$\begin{aligned}&(\mathbf{a}\otimes\mathbf{b})(\mathbf{c}\otimes\mathbf{d})=\mathbf{b}\cdot\mathbf{c}(\mathbf{a}\otimes\mathbf{d})\ \ \forall\mathbf{a},\mathbf{b},\mathbf{c},\mathbf{d}\in\mathcal{V},\\&\mathbf{A}(\mathbf{a}\otimes\mathbf{b})=(\mathbf{A}\mathbf{a})\otimes\mathbf{b}\ \ \forall\mathbf{a},\mathbf{b}\in\mathcal{V},\ \mathbf{A}\in Lin(\mathcal{V}).\end{aligned}\tag{2.4}$$

Finally, the symbol $\mathbf{L}^2$ is normally used to denote in short the product $\mathbf{L}\mathbf{L}$, $\forall\mathbf{L}\in Lin(\mathcal{V})$.

2.4 Transpose, symmetric and skew tensors

For any tensor $\mathbf{L}\in Lin(\mathcal{V})$, it exists just one tensor $\mathbf{L}^\top$, called the *transpose* of $\mathbf{L}$, such that

$$\mathbf{u}\cdot\mathbf{L}\mathbf{v}=\mathbf{v}\cdot\mathbf{L}^\top\mathbf{u}\ \ \forall\mathbf{u},\mathbf{v}\in\mathcal{V}.$$

The transpose of the transpose of $\mathbf{L}$ is $\mathbf{L}$:

$$\mathbf{u}\cdot\mathbf{L}\mathbf{v}=\mathbf{v}\cdot\mathbf{L}^\top\mathbf{u}=\mathbf{u}\cdot(\mathbf{L}^\top)^\top\mathbf{v}\ \ \Rightarrow\ \ (\mathbf{L}^\top)^\top=\mathbf{L}.$$

The Cartesian components of $\mathbf{L}^\top$ are obtained swapping the indexes of the components of $\mathbf{L}$:

$$L^\top_{ij}=\mathbf{e}_i\cdot\mathbf{L}^\top\mathbf{e}_j=\mathbf{e}_j\cdot(\mathbf{L}^\top)^\top\mathbf{e}_i=\mathbf{e}_j\cdot\mathbf{L}\mathbf{e}_i=L_{ji}.$$

It is immediate to show that

$$(\mathbf{A}+\mathbf{B})^\top=\mathbf{A}^\top+\mathbf{B}^\top\ \ \forall\mathbf{A},\mathbf{B}\in Lin(\mathcal{V}),$$

while

$$\mathbf{u}\cdot(\mathbf{A}\mathbf{B})\mathbf{v}=\mathbf{B}\mathbf{v}\cdot\mathbf{A}^\top\mathbf{u}=\mathbf{v}\cdot\mathbf{B}^\top\mathbf{A}^\top\mathbf{u}\ \ \Rightarrow\ \ (\mathbf{A}\mathbf{B})^\top=\mathbf{B}^\top\mathbf{A}^\top.$$

Moreover,

$$\mathbf{u}\cdot(\mathbf{a}\otimes\mathbf{b})\mathbf{v}=\mathbf{a}\cdot\mathbf{u}\ \mathbf{b}\cdot\mathbf{v}=\mathbf{v}\cdot(\mathbf{b}\otimes\mathbf{a})\mathbf{u}\ \ \Rightarrow\ \ (\mathbf{a}\otimes\mathbf{b})^\top=\mathbf{b}\otimes\mathbf{a}.\tag{2.5}$$

A tensor $\mathbf{L}$ is *symmetric* $\Longleftrightarrow$

$$\mathbf{L} = \mathbf{L}^\top.$$

In such a case

$$L_{ij} = L_{ij}^\top = L_{ji} \iff L_{ij} = L_{ji}.$$

A symmetric tensor is hence represented, in a given basis, by a symmetric matrix and has just six independent Cartesian components. Applying eq. (2.4) to $\mathbf{I}$, it is immediately recognized that the identity tensor is symmetric: $\mathbf{I} = \mathbf{I}^\top$.

A tensor $\mathbf{L}$ is *antisymmetric* or *skew* $\Longleftrightarrow$

$$\mathbf{L} = -\mathbf{L}^\top.$$

In such a case (no summation on the index i, see footnote 2, Chap. 1)

$$L_{ij} = -L_{ij}^\top = -L_{ji} \iff L_{ij} = -L_{ji} \Rightarrow L_{\underline{ii}} = 0 \ \forall i = 1, 2, 3.$$

A skew tensor is hence represented, in a given basis, by an antisymmetric matrix whose components on the diagonal are identically null, in any basis; finally, a skew tensor just depends upon three independent Cartesian components.

If we denote by $Sym(\mathcal{V})$ the set of all the symmetric tensors and by $Skw(\mathcal{V})$ that of all the skew tensors, then it is evident that, $\forall \alpha, \beta, \lambda, \mu \in \mathbb{R}$,

$$\begin{gathered} Sym(\mathcal{V}) \cap Skw(\mathcal{V}) = \mathbf{O}, \\ \alpha\mathbf{A} + \beta\mathbf{B} \in Sym(\mathcal{V}) \ \forall \mathbf{A}, \mathbf{B} \in Sym(\mathcal{V}), \\ \lambda\mathbf{L} + \mu\mathbf{M} \in Skw(\mathcal{V}) \ \forall \mathbf{L}, \mathbf{M} \in Skw(\mathcal{V}), \end{gathered}$$

so $Sym(\mathcal{V})$ and $Skw(\mathcal{V})$ are vector subspaces of $Lin(\mathcal{V})$ with $\dim(Sym(\mathcal{V})) = 6$, while $\dim(Skw(\mathcal{V})) = 3$.

Any tensor $\mathbf{L}$ can be decomposed into the sum of a symmetric, $\mathbf{L}^s$, and an antisymmetric, $\mathbf{L}^a$, tensor:

$$\mathbf{L} = \mathbf{L}^s + \mathbf{L}^a,$$

with

$$\mathbf{L}^s = \frac{\mathbf{L} + \mathbf{L}^\top}{2} \quad \in Sym(\mathcal{V})$$

and

$$\mathbf{L}^a = \frac{\mathbf{L} - \mathbf{L}^\top}{2} \quad \in Skw(\mathcal{V}),$$

so that, finally,

$$Lin(\mathcal{V}) = Sym(\mathcal{V}) \oplus Skw(\mathcal{V}).$$

2.5 Trace, scalar product of tensors

It exists one and only one linear form

$$\mathrm{tr} : Lin(\mathcal{V}) \to \mathbb{R},$$

called the *trace*, such that

$$\mathrm{tr}(\mathbf{a}\otimes\mathbf{b}) = \mathbf{a}\cdot\mathbf{b} \quad \forall \mathbf{a},\mathbf{b}\in\mathcal{V}.$$

For its same definition, that has been given without making use of any basis of $\mathcal{V}$, the trace of a tensor is a *tensor invariant*, i.e. a quantity, extracted from a tensor, that does not depend upon the basis.

Linearity implies that

$$\mathrm{tr}(\alpha\mathbf{A}+\beta\mathbf{B}) = \alpha\mathrm{tr}\mathbf{A}+\beta\mathrm{tr}\mathbf{B} \quad \forall\alpha,\beta\in\mathbb{R},\ \mathbf{A},\mathbf{B}\in Lin(\mathcal{V}).$$

It is just the linearity to give the rule for calculating the trace of a tensor $\mathbf{L}$:

$$\mathrm{tr}\mathbf{L} = \mathrm{tr}(L_{ij}\mathbf{e}_i\otimes\mathbf{e}_j) = L_{ij}\mathrm{tr}(\mathbf{e}_i\otimes\mathbf{e}_j) = L_{ij}\ \mathbf{e}_i\cdot\mathbf{e}_j = L_{ij}\delta_{ij} = L_{ii}.$$

A tensor is hence an operator whose sum of the components on the diagonal,

$$\mathrm{tr}\mathbf{L} = L_{11}+L_{22}+L_{33},$$

is constant, regardless of the basis.

Following the same procedure above, it is readily seen that

$$\mathrm{tr}\mathbf{L}^\top = \mathrm{tr}\mathbf{L},$$

which implies, by linearity, that

$$\mathrm{tr}\mathbf{L} = 0 \quad \forall\mathbf{L}\in Skw(\mathcal{V}).$$

The *scalar product* of tensors $\mathbf{A}$ and $\mathbf{B}$ is the positive definite, symmetric bilinear form defined by

$$\mathbf{A}\cdot\mathbf{B} = \mathrm{tr}(\mathbf{A}^\top\mathbf{B}).$$

This definition implies that, $\forall\mathbf{L},\mathbf{M},\mathbf{N}\in Lin(\mathcal{V}),\ \alpha,\beta\in\mathbb{R}$,

$$\begin{aligned}
&\mathbf{L}\cdot(\alpha\mathbf{M}+\beta\mathbf{N}) = \alpha\mathbf{L}\cdot\mathbf{M}+\beta\mathbf{L}\cdot\mathbf{N},\\
&(\alpha\mathbf{L}+\beta\mathbf{M})\cdot\mathbf{N} = \alpha\mathbf{L}\cdot\mathbf{N}+\beta\mathbf{M}\cdot\mathbf{N},\\
&\mathbf{L}\cdot\mathbf{M} = \mathbf{M}\cdot\mathbf{L},\\
&\mathbf{L}\cdot\mathbf{L} > 0 \quad \forall\mathbf{L}\in Lin(\mathcal{V}),\quad \mathbf{L}\cdot\mathbf{L} = 0 \iff \mathbf{L} = \mathbf{O}.
\end{aligned}$$

Such properties give the rule for computing the scalar product of two tensors $\mathbf{A}$ and $\mathbf{B}$:

$$\begin{aligned}
\mathbf{A}\cdot\mathbf{B} &= A_{ij}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot B_{hk}(\mathbf{e}_h\otimes\mathbf{e}_k) = A_{ij}B_{hk}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot(\mathbf{e}_h\otimes\mathbf{e}_k)\\
&= A_{ij}B_{hk}\ \mathrm{tr}[(\mathbf{e}_i\otimes\mathbf{e}_j)^\top(\mathbf{e}_h\otimes\mathbf{e}_k)] = A_{ij}B_{hk}\ \mathrm{tr}[(\mathbf{e}_j\otimes\mathbf{e}_i)(\mathbf{e}_h\otimes\mathbf{e}_k)]\\
&= A_{ij}B_{hk}\ \mathrm{tr}[\mathbf{e}_i\cdot\mathbf{e}_h(\mathbf{e}_j\otimes\mathbf{e}_k)] = A_{ij}B_{hk}\ \mathbf{e}_i\cdot\mathbf{e}_h\ \mathbf{e}_j\cdot\mathbf{e}_k\\
&= A_{ij}B_{hk}\delta_{ih}\delta_{jk} = A_{ij}B_{ij}.
\end{aligned}$$

Like in the case of vectors, the scalar product of two tensors is equal to the sum of the products of the corresponding components. In the same manner, it is easily shown that, $\forall\mathbf{a},\mathbf{b},\mathbf{c},\mathbf{d}\in\mathcal{V}$,

$$(\mathbf{a}\otimes\mathbf{b})\cdot(\mathbf{c}\otimes\mathbf{d}) = \mathbf{a}\cdot\mathbf{c}\ \mathbf{b}\cdot\mathbf{d} = a_ib_jc_id_j,$$

while by the same definition of tensor scalar product,

$$\mathrm{tr}\mathbf{L} = \mathbf{I} \cdot \mathbf{L} \quad \forall \mathbf{L} \in Lin(\mathcal{V}).$$

Similarly to vectors, we define *Euclidean norm* of a tensor $\mathbf{L}$ the nonnegative scalar, denoted either by L or $|\mathbf{L}|$,

$$L = |\mathbf{L}| = \sqrt{\mathbf{L} \cdot \mathbf{L}} = \sqrt{\mathrm{tr}(\mathbf{L}^\top \mathbf{L})} = \sqrt{L_{ij} L_{ij}},$$

and the *distance* $d(\mathbf{L}, \mathbf{M})$ of two tensors $\mathbf{L}$ and $\mathbf{M}$ the norm of the tensor difference:

$$d(\mathbf{L}, \mathbf{M}) = |\mathbf{L} - \mathbf{M}| = |\mathbf{M} - \mathbf{L}|.$$

2.6 Spherical and deviatoric parts

Let $\mathbf{L} \in Sym(\mathcal{V})$; the *spherical part* of $\mathbf{L}$ is defined by

$$\mathbf{L}^{sph} = \frac{1}{3}\mathrm{tr}\mathbf{L}\mathbf{I},$$

and the *deviatoric part* by

$$\mathbf{L}^{dev} = \mathbf{L} - \mathbf{L}^{sph},$$

so that

$$\mathbf{L} = \mathbf{L}^{sph} + \mathbf{L}^{dev}.$$

To remark that

$$\mathrm{tr}\mathbf{L}^{sph} = \frac{1}{3}\mathrm{tr}\mathbf{L}\mathrm{tr}\mathbf{I} = \mathrm{tr}\mathbf{L} \;\Rightarrow\; \mathrm{tr}\mathbf{L}^{dev} = 0,$$

i.e. the deviatoric part is a traceless tensor. Let $\mathbf{A}, \mathbf{B} \in Lin(\mathcal{V})$; then

$$\mathbf{A}^{sph} \cdot \mathbf{B}^{dev} = \frac{1}{3}\mathrm{tr}\mathbf{A}\mathbf{I} \cdot \mathbf{B}^{dev} = \frac{1}{3}\mathrm{tr}\mathbf{A}\ \mathrm{tr}\mathbf{B}^{dev} = 0, \tag{2.6}$$

i.e. any spherical tensor is orthogonal to any deviatoric tensor.

The sets

$$Sph(\mathcal{V}) := \left\{ \mathbf{A}^{sph} \in Lin(\mathcal{V}) |\ \mathbf{A}^{sph} = \frac{1}{3}\mathrm{tr}\mathbf{A}\mathbf{I}\ \forall \mathbf{A} \in Lin(\mathcal{V}) \right\},$$
$$Dev(\mathcal{V}) := \left\{ \mathbf{A}^{dev} \in Lin(\mathcal{V}) |\ \mathbf{A}^{dev} = \mathbf{A} - \mathbf{A}^{sph}\ \forall \mathbf{A} \in Lin(\mathcal{V}) \right\}$$

form two subspaces of $Lin(\mathcal{V})$; the proof is left to the reader. For what proved above, $Sph(\mathcal{V})$ and $Dev(\mathcal{V})$ are two mutually orthogonal subspaces of $Lin(\mathcal{V})$.

2.7 Determinant, inverse of a tensor

The reader is probably familiar with the concept of determinant of a matrix. We show here that the determinant of a second rank tensor can be defined intrinsically and that it corresponds with the determinant of the matrix that represents it in any basis of $\mathcal{V}$. To this purpose, we need first to introduce a mapping:

$$\omega : \mathcal{V} \times \mathcal{V} \times \mathcal{V} \to \mathbb{R}$$

is a *skew trilinear form* if $\omega(\mathbf{u},\mathbf{v},\cdot),\omega(\mathbf{u},\cdot,\mathbf{v})$ and $\omega(\cdot,\mathbf{u},\mathbf{v})$ are linear forms on $\mathcal{V}$ and if

$$\omega(\mathbf{u},\mathbf{v},\mathbf{w}) = -\omega(\mathbf{v},\mathbf{u},\mathbf{w}) = -\omega(\mathbf{u},\mathbf{w},\mathbf{v}) = -\omega(\mathbf{w},\mathbf{v},\mathbf{u})\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V}. \tag{2.7}$$

After this definition, we can state the following

Theorem 1. *Three vectors are linearly independent if and only if every skew trilinear form on them is not null.*

Proof. In fact, be $\mathbf{u} = \alpha\mathbf{v} + \beta\mathbf{w}$; then, for any skew trilinear form ω,

$$\omega(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\alpha\mathbf{v}+\beta\mathbf{w},\mathbf{v},\mathbf{w}) = \alpha\omega(\mathbf{v},\mathbf{v},\mathbf{w}) + \beta\omega(\mathbf{w},\mathbf{v},\mathbf{w}) = 0$$

because of eq. (2.7), applied to the permutation of the positions of the two $\mathbf{u}$ and the two $\mathbf{w}$. □

It is evident that the set of all the skew trilinear forms is a vector space, that we denote by Ω, whose null element is the *null form* ω_0:

$$\omega_0(\mathbf{u},\mathbf{v},\mathbf{w}) = 0\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V}.$$

For a given $\omega(\mathbf{u},\mathbf{v},\mathbf{w}) \in \Omega$, any $\mathbf{L} \in Lin(\mathcal{V})$ induces another form $\omega_L(\mathbf{u},\mathbf{v},\mathbf{w}) \in \Omega$, defined as

$$\omega_L(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\mathbf{L}\mathbf{u},\mathbf{L}\mathbf{v},\mathbf{L}\mathbf{w})\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V}.$$

A key point[3] for the following developments is that $\dim\Omega = 1$.

This means that $\forall \omega_1,\omega_2 \neq \omega_0 \in \Omega, \exists \lambda \in \mathbb{R}$ such that

$$\omega_2(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda\omega_1(\mathbf{u},\mathbf{v},\mathbf{w})\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V}.$$

So, $\forall \mathbf{L} \in Lin(\mathcal{V})$, it must exist $\lambda_L \in \mathbb{R}$ such that

$$\omega(\mathbf{L}\mathbf{u},\mathbf{L}\mathbf{v},\mathbf{L}\mathbf{w}) = \omega_L(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda_L\ \omega(\mathbf{u},\mathbf{v},\mathbf{w})\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V}. \tag{2.8}$$

The scalar[4] λ_L is the *determinant of* $\mathbf{L}$ and in the following it will be denoted as $\det\mathbf{L}$. The determinant of a tensor $\mathbf{L}$ is an intrinsic quantity of $\mathbf{L}$, i.e. it does not depend upon the particular form ω, nor on the basis of $\mathcal{V}$. In fact, we have never introduced, so far, a basis for defining $\det\mathbf{L}$, hence it cannot depend upon the choice of a basis for $\mathcal{V}$, i.e. $\det\mathbf{L}$ is a *tensor invariant*.

Then, if ω^1 and $\omega^2 \in \Omega$, because $\dim\Omega = 1$, it exists $k \in \mathbb{R},\ k \neq 0$ such that

$$\begin{aligned}
&\omega^2(\mathbf{u},\mathbf{v},\mathbf{w}) = k\ \omega^1(\mathbf{u},\mathbf{v},\mathbf{w})\ \forall \mathbf{u},\mathbf{v},\mathbf{w} \in \mathcal{V} \Rightarrow \\
&\omega^2(\mathbf{L}\mathbf{u},\mathbf{L}\mathbf{v},\mathbf{L}\mathbf{w}) = k\ \omega^1(\mathbf{L}\mathbf{u},\mathbf{L}\mathbf{v},\mathbf{L}\mathbf{w}) \rightarrow \\
&\omega^2_L(\mathbf{u},\mathbf{v},\mathbf{w}) = k\ \omega^1_L(\mathbf{u},\mathbf{v},\mathbf{w}).
\end{aligned}$$

[3]The proof of this statement is rather articulated and out of our scope; the interested reader is addressed to the classical textbook of Halmos on linear algebra, §31 (see the bibiography). The theory of the determinants is developed in §53.

[4]More precisely, $\det\mathbf{L}$ is the function that associates a scalar with each tensor (Halmos, §53). We can however, for the sake of practice, identify $\det\mathbf{L}$ with the scalar associated to $\mathbf{L}$, without consequences for our purposes.

Moreover, by eq. (2.8) we get

$$\omega^1(\mathbf{Lu},\mathbf{Lv},\mathbf{Lw}) = \omega^1_L(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda^1_L\omega^1(\mathbf{u},\mathbf{v},\mathbf{w}),$$
$$\omega^2(\mathbf{Lu},\mathbf{Lv},\mathbf{Lw}) = \omega^2_L(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda^2_L\omega^2(\mathbf{u},\mathbf{v},\mathbf{w}),$$

so that

$$\lambda^2_L k\ \omega^1(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda^2_L\omega^2(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega^2_L(\mathbf{u},\mathbf{v},\mathbf{w}) =$$
$$k\ \omega^1_L(\mathbf{u},\mathbf{v},\mathbf{w}) = \lambda^1_L k\ \omega^1(\mathbf{u},\mathbf{v},\mathbf{w}) \iff \lambda^1_L = \lambda^2_L,$$

which proves that $\det\mathbf{L}$ does not depend upon the skew trilinear form, but only upon $\mathbf{L}$.

The definition given for $\det\mathbf{L}$ let us prove some important properties. First of all,

$$\det\mathbf{O} = 0;$$

in fact, $\forall\omega\in\Omega$,

$$\det\mathbf{O}\ \omega(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\mathbf{Ou},\mathbf{Ov},\mathbf{Ow}) = \omega(\mathbf{o},\mathbf{o},\mathbf{o}) = 0$$

because ω operates on three identical, i.e. linearly dependent, vectors. Then, if $\mathbf{L} = \mathbf{I}$,

$$\det\mathbf{I}\ \omega(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\mathbf{Iu},\mathbf{Iv},\mathbf{Iw}) = \omega(\mathbf{u},\mathbf{v},\mathbf{w})$$

if and only if

$$\det\mathbf{I} = 1. \tag{2.9}$$

A third property is that $\forall\mathbf{a},\mathbf{b}\in\mathcal{V}$,

$$\det(\mathbf{a}\otimes\mathbf{b}) = 0. \tag{2.10}$$

In fact, if $\mathbf{L} = \mathbf{a}\otimes\mathbf{b}$, then

$$\det\mathbf{L}\ \omega(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\mathbf{Lu},\mathbf{Lv},\mathbf{Lw}) = \omega((\mathbf{b}\cdot\mathbf{u})\mathbf{a},(\mathbf{b}\cdot\mathbf{v})\mathbf{a},(\mathbf{b}\cdot\mathbf{w})\mathbf{a}) = 0$$

because the three vectors on which $\omega\in\Omega$ operates are linearly dependent; being $\mathbf{u},\mathbf{v}$ and $\mathbf{w}$ arbitrary, this implies eq. (2.10).

An important result is the

Theorem 2. *(Theorem of Binet):* $\forall\mathbf{A},\mathbf{B}\in Lin(\mathcal{V})$

$$\det(\mathbf{AB}) = \det\mathbf{A}\det\mathbf{B}. \tag{2.11}$$

Proof. $\forall\omega\in\Omega$,

$$\lambda_{AB}\omega(\mathbf{u},\mathbf{v},\mathbf{w}) = \omega(\mathbf{ABu},\mathbf{ABv},\mathbf{ABw}) = \omega(\mathbf{A}(\mathbf{Bu}),\mathbf{A}(\mathbf{Bv}),\mathbf{A}(\mathbf{Bw})) =$$
$$\lambda_A\omega(\mathbf{Bu},\mathbf{Bv},\mathbf{Bw}) = \lambda_A\lambda_B\omega(\mathbf{u},\mathbf{v},\mathbf{w}) \iff \lambda_{AB} = \lambda_A\lambda_B,$$

which proves the theorem. □

A tensor $\mathbf{L}$ is called *singular* if $\det \mathbf{L} = 0$, otherwise it is *non-singular*.

Considering eq. (2.8), one can easily see that, if in a basis $\mathcal{B}$ of $\mathcal{V}$ it is $\mathbf{L} = L_{ij}\mathbf{e}_i \otimes \mathbf{e}_j$, then

$$\det \mathbf{L} = \sum_{\pi \in \mathcal{P}_3} \epsilon_{\pi(1),\pi(2),\pi(3)} L_{1,\pi(1)} L_{2,\pi(2)} L_{3,\pi(3)},$$

where $\mathcal{P}_3$ is the set of all the permutations π of $\{1,2,3\}$ and the $\epsilon_{i,j,k}$s are the components of the *Ricci alternator*:

$$\epsilon_{i,j,k} = \begin{cases} 1 & \text{if } (i,j,k) \text{ is an even permutation of } (1,2,3), \\ 0 & \text{if } (i,j,k) \text{ is not a permutation,} \\ -1 & \text{if } (i,j,k) \text{ is an odd permutation of } (1,2,3). \end{cases}$$

The above rule for $\det \mathbf{L}$ coincides with that for calculating the determinant of the matrix whose entries are the L_{ij}s. This shows that, once chosen a basis $\mathcal{B}$ for $\mathcal{V}$, $\det \mathbf{L}$ coincides with the determinant of the matrix representing it in $\mathcal{B}$, and finally that

$$\begin{aligned} \det \mathbf{L} = {} & L_{11}L_{22}L_{33} + L_{12}L_{23}L_{31} + L_{13}L_{32}L_{21} \\ & - L_{11}L_{23}L_{32} - L_{22}L_{13}L_{31} - L_{33}L_{12}L_{21}. \end{aligned} \tag{2.12}$$

This result shows immediately that $\forall \mathbf{L} \in Lin(\mathcal{V})$, and regardless of $\mathcal{B}$, it is

$$\det \mathbf{L}^\top = \det \mathbf{L}. \tag{2.13}$$

Using eq. (2.12), it is not difficult to show that, $\forall \alpha \in \mathbb{R}$,

$$\det(\mathbf{I} + \alpha \mathbf{L}) = 1 + \alpha I_1 + \alpha^2 I_2 + \alpha^3 I_3, \tag{2.14}$$

where I_1, I_2 and I_3 are the three *principal invariants* of $\mathbf{L}$:

$$I_1 = \mathrm{tr}\mathbf{L}, \quad I_2 = \frac{\mathrm{tr}^2\mathbf{L} - \mathrm{tr}\mathbf{L}^2}{2}, \quad I_3 = \det \mathbf{L}. \tag{2.15}$$

A tensor $\mathbf{L} \in Lin(\mathcal{V})$ is said to be *invertible* if there is a tensor $\mathbf{L}^{-1} \in Lin(\mathcal{V})$, called the *inverse* of $\mathbf{L}$, such that

$$\mathbf{L}\mathbf{L}^{-1} = \mathbf{L}^{-1}\mathbf{L} = \mathbf{I}. \tag{2.16}$$

If $\mathbf{L}$ is invertible, then $\mathbf{L}^{-1}$ is unique. By the above definition, if $\mathbf{L}$ is invertible, then

$$\mathbf{u}_1 = \mathbf{L}\mathbf{u} \Rightarrow \mathbf{u} = \mathbf{L}^{-1}\mathbf{u}_1.$$

Theorem 3. *Any invertible tensor maps triples of linearly independent vectors into triples of still linearly independent vectors.*

Proof. Be $\mathbf{L}$ an invertible tensor and $\mathbf{u}_1 = \mathbf{L}\mathbf{u}, \mathbf{v}_1 = \mathbf{L}\mathbf{v}, \mathbf{w}_1 = \mathbf{L}\mathbf{w}$, where $\mathbf{u}, \mathbf{v}, \mathbf{w}$ are three linearly independent vectors. Let us suppose that

$$\mathbf{u}_1 = h\mathbf{v}_1 + k\mathbf{w}_1, \ h, k \in \mathbb{R}.$$

Then, because $\mathbf{L}$ is invertible,

$$\mathbf{L}^{-1}\mathbf{u}_1 = \mathbf{L}^{-1}(h\mathbf{v}_1 + k\mathbf{w}_1) = h\mathbf{L}^{-1}\mathbf{v}_1 + k\mathbf{L}^{-1}\mathbf{w}_1 = h\mathbf{v} + k\mathbf{w},$$

which goes against the hypothesis. By consequence, $\mathbf{u}_1, \mathbf{v}_1$ and $\mathbf{w}_1$ are linearly independent. □

This result, along with the definition of determinant, eq. (2.8), and Theorem 1, proves the

Theorem 4. *(Invertibility theorem):* $\mathbf{L} \in Lin(\mathcal{V})$ *is invertible* $\iff \det \mathbf{L} \neq 0$.

Using the Theorem of Binet, 2, along with eqs. (2.9) and (2.16), we get

$$\det \mathbf{L}^{-1} = \frac{1}{\det \mathbf{L}}.$$

Equation (2.16) applied to $\mathbf{L}^{-1}$, along with the uniqueness of the inverse, gives immediately that

$$(\mathbf{L}^{-1})^{-1} = \mathbf{L},$$

while

$$\mathbf{B}^{-1}\mathbf{A}^{-1} = \mathbf{B}^{-1}\mathbf{A}^{-1}\mathbf{A}\mathbf{B}(\mathbf{A}\mathbf{B})^{-1} = (\mathbf{A}\mathbf{B})^{-1}.$$

The operations of transpose and inversion commute:

$$\mathbf{L}^\top(\mathbf{L}^\top)^{-1} = \mathbf{I} = \mathbf{L}^{-1}\mathbf{L} = \mathbf{I}^\top = (\mathbf{L}^{-1}\mathbf{L})^\top = \mathbf{L}^\top(\mathbf{L}^{-1})^\top \Rightarrow$$
$$(\mathbf{L}^{-1})^\top = (\mathbf{L}^\top)^{-1} := \mathbf{L}^{-\top}.$$

2.8 Eigenvalues and eigenvectors of a tensor

If it exists a $\lambda \in \mathbb{R}$ and a $\mathbf{v} \in \mathcal{V}$, except the null vector, such that

$$\mathbf{L}\mathbf{v} = \lambda\mathbf{v}, \tag{2.17}$$

then λ is an *eigenvalue* and $\mathbf{v}$ an *eigenvector*, relatif to λ, of $\mathbf{L}$. It is immediate to observe that, thanks to linearity, any eigenvector $\mathbf{v}$ of $\mathbf{L}$ is determined to within a multiplier, i.e., that $k\mathbf{v}$ is an eigenvector of $\mathbf{L}$ too, $\forall k \in \mathbb{R}$. Often, the multiplier k is fixed in such a way that $|\mathbf{v}| = 1$.

To determine the eigenvalues and eigenvectors of a tensor, we rewrite eq. (2.17) as

$$(\mathbf{L} - \lambda\mathbf{I})\mathbf{v} = \mathbf{o}. \tag{2.18}$$

The condition for this homogeneous system have a non null solution is

$$\det(\mathbf{L} - \lambda\mathbf{I}) = 0;$$

this is the so-called *characteristic* or *Laplace's equation*. In the case of a second rank tensor over $\mathcal{V}$, the Laplace's equation is an algebraic equation of degree 3 with real coefficients. The roots of the Laplace's equation are the eigenvalues of $\mathbf{L}$; because the components of $\mathbf{L}$, and hence the coefficients of the characteristic equation, are all real, then the eigenvalues of $\mathbf{L}$ are all real or one real and two complex conjugate.

For any eigenvalue $\lambda_i,\ i = 1, 2, 3$ of $\mathbf{L}$, the corresponding eigenvectors $\mathbf{v}_i$ can be found solving eq. (2.18), once put $\lambda = \lambda_i$.

The *proper space* of $\mathbf{L}$ relatif to λ is the subspace of $Lin(\mathcal{V})$ composed by all the vectors that satisfy eq. (2.18). The *multiplicity* of λ is the dimension of its proper space,

while the *spectrum* of $\mathbf{L}$ is the set composed by all of its eigenvalues, each one with its multiplicity.

$\mathbf{L}^\top$ has the same eigenvalues of $\mathbf{L}$, because the Laplace's equation is the same in both the cases:

$$\det(\mathbf{L}^\top - \lambda\mathbf{I}) = \det(\mathbf{L}^\top - \lambda\mathbf{I}^\top) = \det(\mathbf{L} - \lambda\mathbf{I})^\top = \det(\mathbf{L} - \lambda\mathbf{I}).$$

However, this is not the case for the eigenvectors, that, generally speaking, are different, as a simple example can show.

Developing the Laplace's equation, it is easy to show that it can be written as

$$\det(\mathbf{L} - \lambda\mathbf{I}) = -\lambda^3 + I_1\lambda^2 - I_2\lambda + I_3 = 0,$$

which is merely an application of eq. (2.14). If we denote $\mathbf{L}^3 = \mathbf{LLL}$, using eq. (2.15) one can prove the

Theorem 5. *(Cayley-Hamilton Theorem):* $\forall\mathbf{L} \in Lin(\mathcal{V})$,

$$\mathbf{L}^3 - I_1\mathbf{L}^2 + I_2\mathbf{L} - I_3\mathbf{I} = \mathbf{O}.$$

A *quadratic form* defined by $\mathbf{L}$ is any form $\omega : \mathcal{V} \times \mathcal{V} \to \mathbb{R}$ of the type

$$\omega = \mathbf{v} \cdot \mathbf{Lv};$$

if $\omega > 0\ \forall\mathbf{v} \in \mathcal{V},\ \omega = 0 \iff \mathbf{v} = \mathbf{o}$, then ω and $\mathbf{L}$ are said to be *positive definite*. The eigenvalues of a positive definite tensor are positive. In fact, if λ is an eigenvalue of $\mathbf{L}$, positive definite, and $\mathbf{v}$ its eigenvector, then

$$\mathbf{v} \cdot \mathbf{Lv} = \mathbf{v} \cdot \lambda\mathbf{v} = \lambda\mathbf{v}^2 > 0 \iff \lambda > 0.$$

Be $\mathbf{v}_1$ and $\mathbf{v}_2$ two eigenvectors of a symmetric tensor $\mathbf{L}$ relative to the eigenvalues λ_1 and λ_2, respectively, with $\lambda_1 \neq \lambda_2$; then

$$\lambda_1\mathbf{v}_1 \cdot \mathbf{v}_2 = \mathbf{Lv}_1 \cdot \mathbf{v}_2 = \mathbf{Lv}_2 \cdot \mathbf{v}_1 = \lambda_2\mathbf{v}_2 \cdot \mathbf{v}_1 \iff \mathbf{v}_1 \cdot \mathbf{v}_2 = 0.$$

Actually, symmetric tensors have a particular importance, specified by the

Theorem 6. *(Spectral Theorem): the eigenvectors of a symmetric tensor form a basis of* $\mathcal{V}$.

This theorem[5] is of the paramount importance in linear algebra: it proves that the eigenvalues of a symmetric tensor $\mathbf{L}$ are real valued and, remembering the definition of eigenvalues and eigenvectors, eq. (2.17), that it exists a basis $\mathcal{B}_N = \{\mathbf{u}_1, \mathbf{u}_2, \mathbf{u}_3\}$ of $\mathcal{V}$ composed by eigenvectors of $\mathbf{L}$, i.e. by vectors that are mutually orthogonal and that remain mutually orthogonal once transformed by $\mathbf{L}$. Such a basis is called the *normal basis*.

If $\lambda_i, i = 1, 2, 3$, are the eigenvalues of $\mathbf{L}$, then the components of $\mathbf{L}$ in $\mathcal{B}_N$ are

$$L_{ij} = \mathbf{u}_i \cdot \mathbf{Lu}_j = \mathbf{u}_i \cdot \lambda_j\mathbf{u}_j = \lambda_j\delta_{ij}$$

[5]The proof of the spectral theorem is omitted here; the interested reader can find a proof of it in the classical text of Halmos, page 155, see the bibliography.

so finally in $\mathcal{B}_N$ it is

$$\mathbf{L} = \lambda_i \mathbf{e}_i \otimes \mathbf{e}_i,$$

i.e. $\mathbf{L}$ is diagonal and is completely represented by its eigenvalues. In addition, it is easy to check that

$$I_1 = \lambda_1 + \lambda_2 + \lambda_3,\ I_2 = \lambda_1\lambda_2 + \lambda_2\lambda_3 + \lambda_3\lambda_1,\ I_3 = \lambda_1\lambda_2\lambda_3.$$

A tensor with a unique eigenvalue λ, of multiplicity 3, is said to be *spherical*; in such a case, any basis of $\mathcal{V}$ is $\mathcal{B}_N$ and

$$\mathbf{L} = \lambda \mathbf{I}.$$

Two tensors $\mathbf{A}$ and $\mathbf{B}$ are said to be *coaxial* if they have the same normal basis $\mathcal{B}_N$, i.e. if they share the same eigenvectors. Be $\mathbf{u}$ an eigenvector of $\mathbf{A}$, relative to the eigenvalue λ_A, and of $\mathbf{B}$, relatif to λ_B. Then,

$$\mathbf{ABu} = \mathbf{A}\lambda_B\mathbf{u} = \lambda_B\mathbf{Au} = \lambda_A\lambda_B\mathbf{u} = \lambda_A\mathbf{Bu} = \mathbf{B}\lambda_A\mathbf{u} = \mathbf{BAu},$$

which shows, on one hand, that also $\mathbf{Bu}$ is an eigenvector of $\mathbf{A}$, relative to the same eigenvalue λ_A; in the same way, of course, $\mathbf{Au}$ is an eigenvector of $\mathbf{B}$ relative to λ_B. In other words, this shows that $\mathbf{B}$ leaves unchanged any proper space of $\mathbf{A}$, and viceversa. On the other hand, we see that, at least for what concerns the eigenvectors, two tensors commute if and only if they are coaxial. Because any vector can be written as a linear combination of the vectors of $\mathcal{B}_N$, and for the linearity of tensors, we finally have proved the

Theorem 7. *(Commutation Theorem): two tensors commute if and only if they are coaxial.*

2.9 Skew tensors and cross product

Because $\dim(\mathcal{V}) = \dim(Skw(\mathcal{V})) = 3$, an isomorphism can be established between $\mathcal{V}$ and $Skw(\mathcal{V})$, i.e. between vectors and skew tensors. We establish hence a way to associate in a unique way a vector to any skew tensor and inversely. To this purpose, we first introduce the following

Theorem 8. *The spectrum of any tensor $\mathbf{W} \in Skw(\mathcal{V})$ is $\{0\}$ and the dimension of its proper space is 1.*

Proof. This theorem states that zero is the only real eigenvalue of any skew tensor and that its multiplicity is 1. In fact, be $\mathbf{w}$ an eigenvector of $\mathbf{W}$ relative to the eigenvector λ. Then

$$\begin{aligned}\lambda^2\mathbf{w}^2 &= \mathbf{Ww}\cdot\mathbf{Ww} = \mathbf{w}\cdot\mathbf{W}^\top\mathbf{Ww} = -\mathbf{w}\cdot\mathbf{WWw}\\ &= -\mathbf{w}\cdot\mathbf{W}(\lambda\mathbf{w}) = -\lambda\mathbf{w}\cdot\mathbf{Ww} = -\lambda^2\mathbf{w}^2 \iff \lambda = 0.\end{aligned}$$

Then, if $\mathbf{W} \neq \mathbf{O}$ its rank is necessarily 2, because $\det\mathbf{W} = 0\ \forall\mathbf{W} \in Skw(\mathcal{V})$; hence, the equation

$$\mathbf{Ww} = \mathbf{o} \tag{2.19}$$

has ∞^1 solutions, i.e. the multiplicity of λ is 1, which proves the theorem. □

The last equation gives also the way the isomorphism is constructed: in fact, using eq. (2.19) it is easy to check that if $\mathbf{w} = (a, b, c)$, then

$$\mathbf{w} = (a, b, c) \iff \mathbf{W} = \begin{bmatrix} 0 & -c & b \\ c & 0 & -a \\ -b & a & 0 \end{bmatrix}. \tag{2.20}$$

The proper space of $\mathbf{W}$ is called the *axis of* $\mathbf{W}$ and it is indicated by $\mathcal{A}(\mathbf{W})$:

$$\mathcal{A}(\mathbf{W}) = \{\mathbf{u} \in \mathcal{V} |\ \mathbf{W}\mathbf{u} = \mathbf{o}\}.$$

The consequence of what shown above is that $\dim \mathcal{A}(\mathbf{W}) = 1$. With regard to eq. (2.20), one can check easily that the equation

$$\mathbf{u} \cdot \mathbf{u} = \frac{1}{2}\mathbf{W} \cdot \mathbf{W} \tag{2.21}$$

is satisfied only by $\mathbf{w}$ and by its opposite $-\mathbf{w}$. Because both these vectors belong to $\mathcal{A}(\mathbf{W})$, choosing one of them corresponds to choose an orientation for $\mathcal{E}$, see below. We will always fix our choice according to eq. (2.20), which fixes once and for all the isomorphism between $\mathcal{V}$ and $Skw(\mathcal{V})$ that makes correspond any vector $\mathbf{w}$ with one and only one *axial tensor* $\mathbf{W}$ and vice-versa, any skew tensor $\mathbf{W}$ with a unique *axial vector* $\mathbf{w}$.

It is worth noting that the above isomorphism between the vector spaces $\mathcal{V}$ and $Skw(\mathcal{V})$ implies that to any linear combination of vectors $\mathbf{a}$ and $\mathbf{b}$ corresponds an equal linear combination of the corresponding axial tensors $\mathbf{W}_a$ and $\mathbf{W}_b$ and vice-versa, i.e. $\forall a, b \in \mathbb{R}$

$$\mathbf{w} = \alpha\mathbf{a} + \beta\mathbf{b} \iff \mathbf{W} = \alpha\mathbf{W}_a + \beta\mathbf{W}_b, \tag{2.22}$$

where $\mathbf{W}$ is the axial tensor of $\mathbf{w}$. Such a property is immediately checked using eq. (2.20).

It is useful, for further developments, to calculate the powers of $\mathbf{W}$:

$$\mathbf{W}^2 = \mathbf{W}\mathbf{W} = -\mathbf{W}^\top(-\mathbf{W}^\top) = (\mathbf{W}\mathbf{W})^\top = (\mathbf{W}^2)^\top \tag{2.23}$$

i.e. $\mathbf{W}^2$ is symmetric. Moreover:

$$\begin{aligned} \mathbf{W}^2\mathbf{u} = \mathbf{W}\mathbf{W}\mathbf{u} &= \mathbf{w} \times (\mathbf{w} \times \mathbf{u}) = \mathbf{w} \cdot \mathbf{u}\mathbf{w} - \mathbf{w} \cdot \mathbf{w}\mathbf{u} \\ &= -(\mathbf{I} - \mathbf{w} \otimes \mathbf{w})\mathbf{u} \ \Rightarrow\ \mathbf{W}^2 = -(\mathbf{I} - \mathbf{w} \otimes \mathbf{w}) \end{aligned} \tag{2.24}$$

So, applying recursively the previous results,

$$\begin{aligned} &\mathbf{W}^3 = \mathbf{W}\mathbf{W}^2 = -\mathbf{W}(\mathbf{I} - \mathbf{w} \otimes \mathbf{w}) = -\mathbf{W} + (\mathbf{W}\mathbf{w}) \otimes \mathbf{w} = -\mathbf{W} \\ &\mathbf{W}^4 = \mathbf{W}\mathbf{W}^3 = -\mathbf{W}^2 \\ &\mathbf{W}^5 = \mathbf{W}\mathbf{W}^4 = -\mathbf{W}^3 \\ &\textit{etc.} \end{aligned} \tag{2.25}$$

An important property of any couple axial tensor $\mathbf{W}$ - axial vector $\mathbf{w}$ is

$$\mathbf{W}\mathbf{W} = -\frac{1}{2}|\mathbf{W}|^2(\mathbf{I} - \mathbf{w} \otimes \mathbf{w}),$$

while eq. (2.21) can be generalized to any two axial couples $\mathbf{w}_1, \mathbf{W}_1$ and $\mathbf{w}_2, \mathbf{W}_2$:

$$\mathbf{w}_1 \cdot \mathbf{w}_2 = \frac{1}{2}\mathbf{W}_1 \cdot \mathbf{W}_2.$$

The proof of these two last properties is rather easy and left to the reader.

We define *cross product* of two vectors $\mathbf{a}$ and $\mathbf{b}$ the vector

$$\mathbf{a} \times \mathbf{b} = \mathbf{W}_a \mathbf{b},$$

where $\mathbf{W}_a$ is the axial tensor of $\mathbf{a}$. If $\mathbf{a} = (a_1, a_2, a_3)$ and $\mathbf{b} = (b_1, b_2, b_3)$, then by eq. (2.20) we get

$$\mathbf{a} \times \mathbf{b} = (a_2 b_3 - a_3 b_2, a_3 b_1 - a_1 b_3, a_1 b_2 - a_2 b_1).$$

The cross product is bilinear: $\forall \mathbf{a}, \mathbf{b}, \mathbf{u} \in \mathcal{V},\ \alpha, \beta \in \mathbb{R}$,

$$(\alpha\mathbf{a} + \beta\mathbf{b}) \times \mathbf{u} = \alpha\mathbf{a} \times \mathbf{u} + \beta\mathbf{b} \times \mathbf{u},$$
$$\mathbf{u} \times (\alpha\mathbf{a} + \beta\mathbf{b}) = \alpha\mathbf{u} \times \mathbf{a} + \beta\mathbf{u} \times \mathbf{b}.$$

In fact, the first equation above is a consequence of eq. (2.22), while the second one is a simple application to axial tensors of the same definition of tensor.

Three important results concerning the cross product are stated by the following theorems:

Theorem 9. *(Condition of parallelism): two vectors* $\mathbf{a}$ *and* $\mathbf{b}$ *are parallel, i.e.* $\mathbf{b} = k\mathbf{a},\ k \in \mathbb{R}, \iff$

$$\mathbf{a} \times \mathbf{b} = \mathbf{o}.$$

Proof. This property is actually a consequence of the fact that any eigenvalue of a tensor is determined to within a multiplier:

$$\mathbf{a} \times \mathbf{b} = \mathbf{W}_a \mathbf{b} = \mathbf{o} \iff \mathbf{b} = k\mathbf{a},\ k \in \mathbb{R},$$

for Theorem 8. □

Theorem 10. *(Orthogonality property):*

$$\mathbf{a} \times \mathbf{b} \cdot \mathbf{a} = \mathbf{a} \times \mathbf{b} \cdot \mathbf{b} = 0. \tag{2.26}$$

Proof.

$$\mathbf{a} \times \mathbf{b} \cdot \mathbf{a} = \mathbf{W}_a \mathbf{b} \cdot \mathbf{a} = \mathbf{b} \cdot \mathbf{W}_a^\top \mathbf{a} = -\mathbf{b} \cdot \mathbf{W}_a \mathbf{a} = -\mathbf{b} \cdot \mathbf{o} = 0,$$
$$\mathbf{a} \times \mathbf{b} \cdot \mathbf{b} = \mathbf{W}_a \mathbf{b} \cdot \mathbf{b} = \mathbf{b} \cdot \mathbf{W}_a^\top \mathbf{b} = -\mathbf{b} \cdot \mathbf{W}_a \mathbf{b} \iff \mathbf{a} \times \mathbf{b} \cdot \mathbf{b} = 0.$$

□

Theorem 11. $\mathbf{a} \times \mathbf{b}$ *is the axial vector of the tensor* $(\mathbf{b} \otimes \mathbf{a} - \mathbf{a} \otimes \mathbf{b})$.

Proof. First of all, by eq. (2.5) we see that

$$(\mathbf{b} \otimes \mathbf{a} - \mathbf{a} \otimes \mathbf{b}) \in Skew(\mathcal{V}).$$

Then,

$$(\mathbf{b} \otimes \mathbf{a} - \mathbf{a} \otimes \mathbf{b})(\mathbf{a} \times \mathbf{b}) = \mathbf{a} \cdot \mathbf{a} \times \mathbf{b}\ \mathbf{b} - \mathbf{b} \cdot \mathbf{a} \times \mathbf{b}\ \mathbf{a} = 0$$

for Theorem 10. □

Theorem 11 allows for showing that, unlike the scalar product, the cross product is anti-symmetric:

$$\mathbf{a} \times \mathbf{b} = -\mathbf{b} \times \mathbf{a}. \tag{2.27}$$

In fact, if $\mathbf{W}_1 = (\mathbf{b} \otimes \mathbf{a} - \mathbf{a} \otimes \mathbf{b})$ is the axial tensor of $\mathbf{a} \times \mathbf{b}$, $\mathbf{W}_2 = (-\mathbf{a} \otimes \mathbf{b} + \mathbf{b} \otimes \mathbf{a})$ is that of $-\mathbf{b} \times \mathbf{a}$. But, evidently, $\mathbf{W}_1 = \mathbf{W}_2$ which implies eq. (2.27) for the isomorphism between $\mathcal{V}$ and $Lin(\mathcal{V})$. This property and again Theorem 11 let us show the formula for the *double cross product*:

$$\mathbf{u} \times (\mathbf{v} \times \mathbf{w}) = -(\mathbf{v} \times \mathbf{w}) \times \mathbf{u} = -(\mathbf{w} \otimes \mathbf{v} - \mathbf{v} \otimes \mathbf{w})\mathbf{u} = \mathbf{u} \cdot \mathbf{w}\ \mathbf{v} - \mathbf{u} \cdot \mathbf{v}\ \mathbf{w}. \tag{2.28}$$

Another interesting result concerns the *mixed product*:

$$\mathbf{u} \times \mathbf{v} \cdot \mathbf{w} = \mathbf{W}_u \mathbf{v} \cdot \mathbf{w} = -\mathbf{v} \cdot \mathbf{W}_u \mathbf{w} = -\mathbf{v} \cdot \mathbf{u} \times \mathbf{w} = \mathbf{w} \times \mathbf{u} \cdot \mathbf{v}, \tag{2.29}$$

and similarly

$$\mathbf{u} \times \mathbf{v} \cdot \mathbf{w} = \mathbf{v} \times \mathbf{w} \cdot \mathbf{u}.$$

Using this last result, we can obtain a formula for the norm of a cross product; if $\mathbf{a} = a\ \mathbf{e}_a$ and $\mathbf{b} = b\ \mathbf{e}_b$, with $\mathbf{e}_a, \mathbf{e}_b \in \mathcal{S}$, are two vectors forming the angle θ, then

$$\begin{aligned}
&(\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{a} \times \mathbf{b}) = \mathbf{a} \times \mathbf{b} \cdot (\mathbf{a} \times \mathbf{b}) = (\mathbf{a} \times \mathbf{b}) \times \mathbf{a} \cdot \mathbf{b} = -\mathbf{a} \times (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{b} = \\
&(-\mathbf{a} \cdot \mathbf{b}\ \mathbf{a} + \mathbf{a}^2\ \mathbf{b}) \cdot \mathbf{b} = \mathbf{b} \cdot (\mathbf{a}^2 \mathbf{I} - \mathbf{a} \otimes \mathbf{a})\mathbf{b} = a^2\ \mathbf{b} \cdot (\mathbf{I} - \mathbf{e}_a \otimes \mathbf{e}_a)\mathbf{b} = \\
&a^2 b^2\ \mathbf{e}_b \cdot (\mathbf{I} - \mathbf{e}_a \otimes \mathbf{e}_a)\mathbf{e}_b = a^2 b^2 (1 - \cos^2\theta) = a^2 b^2 \sin^2\theta \rightarrow |\mathbf{a} \times \mathbf{b}| = ab \sin\theta.
\end{aligned}$$

So, the norm of a cross product can be interpreted, geometrically, as the area of the parallelogram spanned by the two vectors. As a consequence, the absolute value of the mixed product (2.29) measures the volume of the prism delimited by three non coplanar vectors.

Because the cross product is antisymmetric and the scalar one is symmetric, it is easy to check that the form

$$\beta(\mathbf{u}, \mathbf{v}, \mathbf{w}) = \mathbf{u} \times \mathbf{v} \cdot \mathbf{w}$$

is a skew trilinear form. Then, eq. (2.8), we get

$$\mathbf{L}\mathbf{u} \times \mathbf{L}\mathbf{v} \cdot \mathbf{L}\mathbf{w} = \det \mathbf{L}\ \mathbf{u} \times \mathbf{v} \cdot \mathbf{w}. \tag{2.30}$$

Following the interpretation given above for the absolute value of the mixed product, we can conclude that $|\det \mathbf{L}|$ can be interpreted as a coefficient of volume dilation. A geometrical interpretation can then be given to the case of a non invertible tensor, i.e. of $\det \mathbf{L} = 0$: it crushes a prism into a flat region (the three original vectors become coplanar, i.e. linearly dependent).

The *adjugate* of $\mathbf{L}$ is the tensor

$$\mathbf{L}^* := (\det \mathbf{L})\mathbf{L}^{-\top}.$$

From eq. (2.30) we get hence

$$\det \mathbf{L}\ \mathbf{u} \times \mathbf{v} \cdot \mathbf{w} = \mathbf{L}\mathbf{u} \times \mathbf{L}\mathbf{v} \cdot \mathbf{L}\mathbf{w} = \mathbf{L}^\top(\mathbf{L}\mathbf{u} \times \mathbf{L}\mathbf{v}) \cdot \mathbf{w} \ \ \forall \mathbf{w} \Rightarrow$$
$$\mathbf{L}\mathbf{u} \times \mathbf{L}\mathbf{v} = \mathbf{L}^*(\mathbf{u} \times \mathbf{v}).$$

2.10 Orientation of a basis

It is immediate to observe that a basis $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ can be oriented in two opposite ways[6]: e.g., once two unit mutually orthogonal vectors $\mathbf{e}_1$ and $\mathbf{e}_2$ chosen, there are two opposite unit vectors perpendicular to both $\mathbf{e}_1$ and $\mathbf{e}_2$ that can be chosen to form $\mathcal{B}$.

We say that $\mathcal{B}$ is *positively oriented* or *right-handed* if

$$\mathbf{e}_1 \times \mathbf{e}_2 \cdot \mathbf{e}_3 = 1,$$

while $\mathcal{B}$ is *negatively oriented* or *left-handed* if

$$\mathbf{e}_1 \times \mathbf{e}_2 \cdot \mathbf{e}_3 = -1.$$

Schematically, a right-handed basis is represented in Fig. 2.1, where a left-handed basis is represented too, with a dashed $\mathbf{e}_3$.

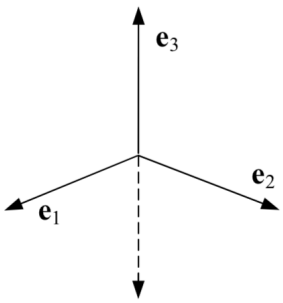

Figure 2.1: Right- and left-handed bases.

With a right-handed basis, by definition the axial tensors of the three vectors of the basis are

$$\mathbf{W}_{\mathbf{e}_1} = \mathbf{e}_3 \otimes \mathbf{e}_2 - \mathbf{e}_2 \otimes \mathbf{e}_3,$$
$$\mathbf{W}_{\mathbf{e}_2} = \mathbf{e}_1 \otimes \mathbf{e}_3 - \mathbf{e}_3 \otimes \mathbf{e}_1,$$
$$\mathbf{W}_{\mathbf{e}_3} = \mathbf{e}_2 \otimes \mathbf{e}_1 - \mathbf{e}_1 \otimes \mathbf{e}_2.$$

[6]It is evident that this is true also for one- and two-dimensional vector spaces.

2.11 Rotations

In the previous Chapter we have seen that the elements of $\mathcal{V}$ represent translations over $\mathcal{E}$. A *rotation*, i.e. a rigid rotation of the space, is an operation that transforms any two vectors $\mathbf{u}$ and $\mathbf{v}$ into two other vectors $\mathbf{u}'$ and $\mathbf{v}'$ in such a way that

$$u = u', \quad v = v', \quad \mathbf{u} \cdot \mathbf{v} = \mathbf{u}' \cdot \mathbf{v}', \tag{2.31}$$

i.e. it preserves norms and angles. Because a rotation is a transformation from $\mathcal{V}$ to $\mathcal{V}$, rotations are tensors, i.e. we can write

$$\mathbf{v}' = \mathbf{R}\mathbf{v},$$

with $\mathbf{R}$ the *rotation tensor* or simply the *rotation*.

Conditions (2.31) impose some restrictions on $\mathbf{R}$:

$$\mathbf{u}' \cdot \mathbf{v}' = \mathbf{R}\mathbf{u} \cdot \mathbf{R}\mathbf{v} = \mathbf{u} \cdot \mathbf{R}^\top \mathbf{R}\mathbf{v} = \mathbf{u} \cdot \mathbf{v} \iff \mathbf{R}^\top \mathbf{R} = \mathbf{I} = \mathbf{R}\mathbf{R}^\top.$$

A tensor that preserves angles belongs to $Orth(\mathcal{V})$, the subspace of *orthogonal tensors* (we leave to the reader the proof that actually $Orth(\mathcal{V})$ is a subspace of $Lin(\mathcal{V})$. Replacing in the above equation $\mathbf{v}$ with $\mathbf{u}$ shows immediately that an orthogonal tensor preserves also the norms. By the uniqueness of the inverse, we see that

$$\mathbf{R} \in Orth(\mathcal{V}) \iff \mathbf{R}^{-1} = \mathbf{R}^\top.$$

The above condition is not sufficient to characterize a rotation; in fact, a rotation must transform a right-handed basis into another right-handed basis, i.e. it must *preserve the orientation of the space*. This means that it must be

$$\mathbf{e}'_1 \times \mathbf{e}'_2 \cdot \mathbf{e}'_3 = \mathbf{R}\mathbf{e}_1 \times \mathbf{R}\mathbf{e}_2 \cdot \mathbf{R}\mathbf{e}_3 = \mathbf{e}_1 \times \mathbf{e}_2 \cdot \mathbf{e}_3.$$

By eq. (2.30) we get hence the condition[7]

$$\det \mathbf{R}(\mathbf{e}_1 \times \mathbf{e}_2 \cdot \mathbf{e}_3) = \mathbf{e}_1 \times \mathbf{e}_2 \cdot \mathbf{e}_3 \iff \det \mathbf{R} = 1.$$

The tensors of $Orth(\mathcal{V})$ that have a determinant equal to 1 form the subspace of *proper rotations* or simply *rotations*, indicated by $Orth(\mathcal{V})^+$ or also by $SO(3)$. Only tensors of $Orth(\mathcal{V})^+$ represent rigid rotations of $\mathcal{E}$[8].

Theorem 12. *: each tensor $\mathbf{R} \in Orth(\mathcal{V})$ has the eigenvalue ± 1, with $+1$ for rotations.*

Proof. Be $\mathbf{u}$ an eigenvector of $\mathbf{R} \in Orth(\mathcal{V})$ corresponding to the eigenvalue λ. Because $\mathbf{R}$ preserves the norm, it is

$$\mathbf{R}\mathbf{u} \cdot \mathbf{R}\mathbf{u} = \lambda^2 \mathbf{u}^2 = \mathbf{u}^2 \rightarrow \lambda^2 = 1.$$

[7] From the condition $\mathbf{R}^\top \mathbf{R} = \mathbf{I}$ and through eq.(2.13) and the Theorem of Binet, we recognize immediately that $\det \mathbf{R} = \pm 1 \quad \forall \mathbf{R} \in Orth(\mathcal{V})$.

[8] A tensor $\mathbf{S} \in Orth(\mathcal{V})$ such that $\det \mathbf{S} = -1$ represents a transformation that changes the orientation of the space, like mirror symmetries do, see Sect. 2.12.

We now must prove that it exists at least one real eigenvector λ. To this end, we consider the characteristic equation

$$f(\lambda) = \lambda^3 + k_1\lambda^2 + k_2\lambda + k_3 = 0,$$

whose coefficients k_i are real-valued, because $\mathbf{R}$ has real-valued components. It is immediate to recognize that

$$\lim_{\lambda\to\pm\infty} f(\lambda) = \pm\infty.$$

So, because $f(\lambda)$ is a real-valued continuous function, actually a polynomial of λ, it exists at least one $\lambda_1 \in \mathbb{R}$ such that

$$f(\lambda_1) = 0.$$

In addition, we already know that $\forall \mathbf{R} \in Orth(\mathcal{V}), \det \mathbf{R} = \pm 1$ and that, if $\lambda_i, i = 1, 2, 3$ are the eigenvalues of $\mathbf{R}$, then $\det \mathbf{R} = \lambda_1\lambda_2\lambda_3$. Hence, two are the possible cases:

i. $\lambda_1 \in \mathbb{R}$ and $\lambda_2, \lambda_3 \in \mathbb{C}$, with $\lambda_3 = \overline{\lambda}_2$, the complex conjugate of λ_2;

ii. $\lambda_i \in \mathbb{R}\ \forall i = 1, 2, 3$.

Let us consider the case of $\mathbf{R} \in Orth(\mathcal{V})^+$, i.e. a (proper) rotation $\rightarrow \det \mathbf{R} = 1$. Then, in the first case above,

$$\det \mathbf{R} = \lambda_1\lambda_2\overline{\lambda}_2 = \lambda_1[\Re^2(\lambda_2) + \Im^2(\lambda_2)].$$

But

$$\Re^2(\lambda_2) + \Im^2(\lambda_2) = 1$$

because it is the square of the modulus of the complex eigenvalue λ_2. So in this case

$$\det \mathbf{R} = 1 \iff \lambda_1 = 1.$$

In the second case, $\lambda_i \in \mathbb{R}\ \forall i = 1, 2, 3$, either $\lambda_1 > 0, \lambda_2, \lambda_3 < 0$, or all of them are positive. Because the modulus of each eigenvalue must be equal to 1, either $\lambda_1 = 1$ or $\lambda_i = 1\ \forall i = 1, 2, 3$ (in this case $\mathbf{R} = \mathbf{I}$).

Following the same steps, one arrives easily to show that $\forall \mathbf{S} \in Orth(\mathcal{V})$ with $\det \mathbf{S} = -1$, it exists at least one real eigenvalue $\lambda_1 = -1$. □

Generally speaking, a rotation tensor rotates the basis $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ into the basis $\mathcal{B}' = \{\mathbf{e}'_1, \mathbf{e}'_2, \mathbf{e}'_3\}$:

$$\mathbf{R}\mathbf{e}_i = \mathbf{e}'_i \ \ \forall i = 1, 2, 3 \ \Rightarrow\ R_{ij} = \mathbf{e}_i \cdot \mathbf{R}\mathbf{e}_j = \mathbf{e}_i \cdot \mathbf{e}'_j. \tag{2.32}$$

This result actually means that the j-th column of $\mathbf{R}$ is composed by the components in the base $\mathcal{B}$ of the vector $\mathbf{e}'_j$ of $\mathcal{B}'$. Because the two bases are orthonormal, such components are the director cosines of the axes of $\mathcal{B}'$ with respect to $\mathcal{B}$.

Geometrically speaking, any rotation is characterized by an *axis of rotation* $\mathbf{w}, |\mathbf{w}| = 1$ and by an *amplitude* φ, i.e. the angle through which the space is rotated about $\mathbf{w}$. By definition, $\mathbf{w}$ is the (only) vector that is left unchanged by $\mathbf{R}$, i.e.

$$\mathbf{R}\mathbf{w} = \mathbf{w},$$

or, in other words, it is the eigenvector corresponding to the eigenvalue $+1$.

The question is then: how a rotation tensor $\mathbf{R}$ can be expressed by means of its geometrical parameters, $\mathbf{w}$ and φ? To this end we have a fundamental theorem:

Theorem 13. *(Euler's rotation representation theorem):* $\forall \mathbf{R} \in Orth(\mathcal{V})^+$,

$$\mathbf{R} = \mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2 \tag{2.33}$$

with φ *the rotation's amplitude and* $\mathbf{W}$ *the axial tensor of the rotation axis* $\mathbf{w}$.

Proof. We observe preliminarily that

$$\mathbf{R}\mathbf{w} = \mathbf{I}\mathbf{w} + \sin\varphi\mathbf{W}\mathbf{w} + (1 - \cos\varphi)\mathbf{W}\mathbf{W}\mathbf{w} = \mathbf{I}\mathbf{w} = \mathbf{w} \tag{2.34}$$

i.e. that eq. (2.33) actually defines a transformation that leaves unchanged the axis $\mathbf{w}$, like a rotation about $\mathbf{w}$ must do, and that $+1$ is an eigenvalue of $\mathbf{R}$.

We need now to prove that eq. (2.33) actually represents a rotation tensor, i.e. we must prove that

$$\mathbf{R}\mathbf{R}^\top = \mathbf{I}, \quad \det\mathbf{R} = 1.$$

Through eq. (2.25) we get

$$\begin{aligned}
\mathbf{R}\mathbf{R}^\top &= (\mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2)(\mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2)^\top \\
&= (\mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2)(\mathbf{I} - \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2) \\
&= \mathbf{I} + 2(1 - \cos\varphi)\mathbf{W}^2 - \sin^2\varphi\mathbf{W}^2 + (1 - \cos\varphi)^2\mathbf{W}^4 \\
&= \mathbf{I} + 2(1 - \cos\varphi)\mathbf{W}^2 - \sin^2\varphi\mathbf{W}^2 - (1 - \cos\varphi)^2\mathbf{W}^2 = \mathbf{I}.
\end{aligned}$$

Then, through eq. (2.24) we obtain

$$\begin{aligned}
\mathbf{R} &= \mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2 \\
&= \mathbf{I} + \sin\varphi\mathbf{W} - (1 - \cos\varphi)(\mathbf{I} - \mathbf{w} \otimes \mathbf{w}) \\
&= \cos\varphi\mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{w} \otimes \mathbf{w}.
\end{aligned} \tag{2.35}$$

To go on, we need to express $\mathbf{W}$ and $\mathbf{w} \otimes \mathbf{w}$; if $\mathbf{w} = (w_1, w_2, w_3)$, then by eq. (2.20) we have

$$\mathbf{W} = \begin{bmatrix} 0 & -w_3 & w_2 \\ w_3 & 0 & -w_1 \\ -w_2 & w_1 & 0 \end{bmatrix}$$

and by eq. (2.2)

$$\mathbf{w} \otimes \mathbf{w} = \begin{bmatrix} w_1^2 & w_1w_2 & w_1w_3 \\ w_1w_2 & w_2^2 & w_2w_3 \\ w_1w_3 & w_2w_3 & w_3^2 \end{bmatrix},$$

that injected into eq. (2.35) gives

$$\mathbf{R} = \begin{bmatrix} \cos\varphi + (1 - \cos\varphi)w_1^2 & -w_3\sin\varphi + w_1w_2(1 - \cos\varphi) & w_2\sin\varphi + w_1w_3(1 - \cos\varphi) \\ w_3\sin\varphi + w_1w_2(1 - \cos\varphi) & \cos\varphi + (1 - \cos\varphi)w_2^2 & -w_1\sin\varphi + w_2w_3(1 - \cos\varphi) \\ -w_2\sin\varphi + w_1w_3(1 - \cos\varphi) & w_1\sin\varphi + w_2w_3(1 - \cos\varphi) & \cos\varphi + (1 - \cos\varphi)w_3^2 \end{bmatrix}. \tag{2.36}$$

This formula gives $\mathbf{R}$ as function exclusively of $\mathbf{w}$ and φ, the geometrical elements of the rotation. Then

$$\det \mathbf{R} = (w^2 + (1 - w^2)\cos\varphi)(\cos^2\varphi + w^2 \sin^2\varphi)$$

and because $w = 1, \det \mathbf{R} = 1$, which proves that eq. (2.33) actually represents a rotation.

We eventually need to prove that eq. (2.33) represents the rotation about $\mathbf{w}$ of amplitude φ. To this end, we choose an orthonormal basis $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ of $\mathcal{V}$ such that $\mathbf{w} = \mathbf{e}_3$, i.e. we analyze the particular case of a rotation of amplitude φ about $\mathbf{e}_3$. This is always possible, thanks to the arbitrariness of the basis of $\mathcal{V}$. In such a case, eq. (2.32) gives

$$\mathbf{R} = \begin{bmatrix} \cos\varphi & -\sin\varphi & 0 \\ \sin\varphi & \cos\varphi & 0 \\ 0 & 0 & 1 \end{bmatrix}. \tag{2.37}$$

Moreover,

$$\mathbf{W} = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \quad \mathbf{w} \otimes \mathbf{w} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix},$$
$$\mathbf{W}^2 = -(\mathbf{I} - \mathbf{w} \otimes \mathbf{w}) = \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix}.$$

Hence

$$\begin{aligned} \mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} + \sin\varphi \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + \\ + (1 - \cos\varphi) \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} \cos\varphi & -\sin\varphi & 0 \\ \sin\varphi & \cos\varphi & 0 \\ 0 & 0 & 1 \end{bmatrix} = \mathbf{R}. \end{aligned} \tag{2.38}$$

□

Equation (2.33) gives another result: to obtain the inverse of $\mathbf{R}$ it is sufficient to change the sign of φ. In fact, because $\mathbf{W} \in Skw(\mathcal{V})$ and through eq. (2.23)

$$\begin{aligned} \mathbf{R}^{-1} = \mathbf{R}^\top &= (\mathbf{I} + \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2)^\top = \mathbf{I} + \sin\varphi\mathbf{W}^\top + (1 - \cos\varphi)(\mathbf{W}^2)^\top \\ &= \mathbf{I} - \sin\varphi\mathbf{W} + (1 - \cos\varphi)\mathbf{W}^2 = \mathbf{I} + \sin(-\varphi)\mathbf{W} + (1 - \cos(-\varphi))\mathbf{W}^2. \end{aligned}$$

The knowledge of the inverse of a rotation allows also to perform the operation of *change of basis*, i.e. to determine the components of a vector or of a tensor in a basis $\mathcal{B}' = \{\mathbf{e}'_1, \mathbf{e}'_2, \mathbf{e}'_3\}$ rotated with respect to an original basis $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ by a rotation $\mathbf{R}$ (in the following equation, a prime indicates a quantity specified in the basis $\mathcal{B}'$). Considering that

$$\mathbf{e}_i = \mathbf{R}^{-1}\mathbf{e}'_i = \mathbf{R}^\top\mathbf{e}'_i = R^\top_{hk}(\mathbf{e}'_h \otimes \mathbf{e}'_k)\mathbf{e}'_i = R^\top_{hk}\delta_{ki}\mathbf{e}'_h$$

we get, for a vector $\mathbf{u}$,

$$\mathbf{u} = u_i\mathbf{e}_i = u_i R^\top_{ki}\mathbf{e}'_k$$

i.e.

$$u'_k = R^\top_{ki} u_i \;\rightarrow\; \mathbf{u}' = \mathbf{R}^\top \mathbf{u}.$$

We remark that, because $\mathbf{R}^\top = \mathbf{R}^{-1}$, the operation of change of basis is just the opposite one of the rotation of the space (and actually, we have seen that it is sufficient to take the opposite of φ in eq. (2.33) to get $\mathbf{R}^{-1}$).

For a second-rank tensor $\mathbf{L}$ we get

$$\mathbf{L} = L_{ij}\mathbf{e}_i \otimes \mathbf{e}_j = L_{ij} R^\top_{mi}\mathbf{e}'_m \otimes R^\top_{nj}\mathbf{e}'_n = R^\top_{mi} R^\top_{nj} L_{ij}\mathbf{e}'_m \otimes \mathbf{e}'_n,$$

i.e.

$$L'_{mn} = R^\top_{mi} R^\top_{nj} L_{ij} \;\rightarrow\; \mathbf{L}' = \mathbf{R}^\top \mathbf{L}\mathbf{R}.$$

We remark something that is typical of tensors: the components of a r-rank tensor in a rotated basis $\mathcal{B}'$ depend upon the r-th powers of the directors cosines of the axes of $\mathcal{B}'$, i.e. on the r-th powers of the components R_{ij} of $\mathbf{R}$.

If a rotation tensor is known through its Cartesian components in a given basis $\mathcal{B}$, it is easy to calculate its geometrical elements: the rotation axis $\mathbf{w}$ is the eigenvector of $\mathbf{R}$ corresponding to the eigenvalue 1, so it is found solving the equation

$$\mathbf{R}\mathbf{w} = \mathbf{w}$$

and then normalizing it, while the rotation amplitude φ can be found still using (2.33): because the trace of a tensor is an invariant, we get

$$\mathrm{tr}\mathbf{R} = 3 + (1-\cos\varphi)\mathrm{tr}(-\mathbf{I} + \mathbf{w}\cdot\mathbf{w}) = 1 + 2\cos\varphi \;\rightarrow\; \varphi = \arccos\frac{\mathrm{tr}\mathbf{R} - 1}{2}.$$

It is interesting to consider the geometrical meaning of eq. (2.33). To this purpose we apply eq. (2.33) to a vector $\mathbf{u}$, see Fig. 2.2

$$\begin{aligned}\mathbf{R}\mathbf{u} &= (\mathbf{I} + \sin\varphi\mathbf{W} + (1-\cos\varphi)\mathbf{W}^2)\mathbf{u} \\ &= \mathbf{u} + \sin\varphi\mathbf{w}\times\mathbf{u} + (1-\cos\varphi)\mathbf{w}\times(\mathbf{w}\times\mathbf{u})\end{aligned}$$

The rotated vector $\mathbf{R}\mathbf{u}$ is the sum of three vectors; in particular, $\sin\varphi\mathbf{W}\mathbf{u}$ is always

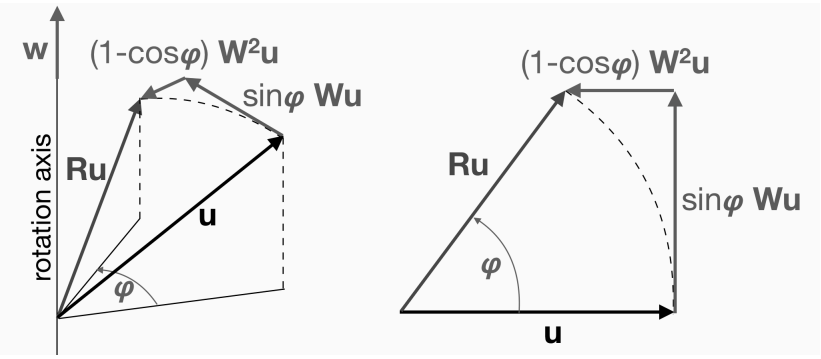

Figure 2.2: Rotation of a vector.

orthogonal to $\mathbf{u}$, $\mathbf{w}$ and $(1-\cos\varphi)\mathbf{W}^2\mathbf{u}$. If $\mathbf{u}\cdot\mathbf{w} = 0$, see the sketch on the right in Fig. 2.2, then $(1-\cos\varphi)\mathbf{W}^2\mathbf{u}$ is also parallel to $\mathbf{u}$.

Let us consider now a composition of rotations. In particular, let us imagine that a vector $\mathbf{u}$ is rotated first by $\mathbf{R}_1$, around $\mathbf{w}_1$ through φ_1, then by $\mathbf{R}_2$, around $\mathbf{w}_2$ through φ_2. So, first the vector $\mathbf{u}$ becomes the vector

$$\mathbf{u}_1 = \mathbf{R}_1\mathbf{u}.$$

Then, the vector $\mathbf{u}_1$ is rotated about $\mathbf{w}_2$ through φ_2 to become

$$\mathbf{u}_{12} = \mathbf{R}_2\mathbf{u}_1 = \mathbf{R}_2\mathbf{R}_1\mathbf{u}.$$

Let us now suppose to change the order of the rotations: $\mathbf{R}_2$ first and then $\mathbf{R}_1$. The final result will be the vector

$$\mathbf{u}_{21} = \mathbf{R}_1\mathbf{R}_2\mathbf{u}. \tag{2.39}$$

Because the tensor product is not symmetric (i.e., it has not the commutativity property), generally speaking[9]

$$\mathbf{u}_{12} \neq \mathbf{u}_{21}.$$

In other words, the order of the rotations matters: changing the order of the rotations leads to a different final result. An example is shown in Fig. 2.3.

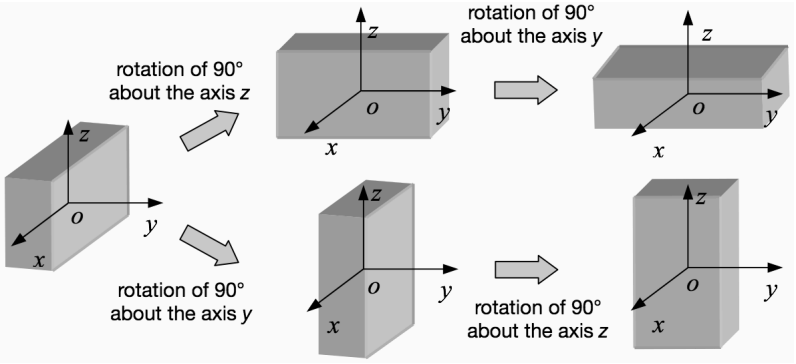

Figure 2.3: Non-commutativity of the rotations.

This is a fundamental difference between rotations and displacements, that commute, see Fig. 1.2, because the composition of displacements is ruled by the sum of vectors:

$$\mathbf{w} = \mathbf{u} + \mathbf{v} = \mathbf{v} + \mathbf{u} \tag{2.40}$$

This difference, which is a major point in physics, comes from the difference of the operators: vectors for the displacements, tensors for the rotations.

Any rotation can be specified by the knowledge of three parameters. This can be easily seen from eq. (2.33): the parameters are the three components of $\mathbf{w}$, that are not independent because

$$w = |\mathbf{w}| = \sqrt{w_1^2 + w_2^2 + w_3^2} = 1$$

and by the amplitude angle φ. The choice of the parameters by which to express a rotation is not unique. Besides the use of the Cartesian components of $\mathbf{w}$ and φ, other choices are possible, let us see three of them:

[9]We have seen, Theorem 7, that two tensors commute $\iff$ they are coaxial, i.e. if they have the same eigenvectors. Because the rotation axis is always a real eigenvector of a rotation tensor, if two tensors operate a rotation about different axes they are not coaxial. Hence, the rotation tensors about different axes never commute.

i. *physical angles*: the rotation axis $\mathbf{w}$ is given through its spherical coordinates ψ, the *longitude*, $0 \leq \psi < 2\pi$, and θ, the *colatitude*, $0 \leq \theta \leq \pi$, see Fig. 2.4, the third parameter being the *rotation amplitude* φ. Then

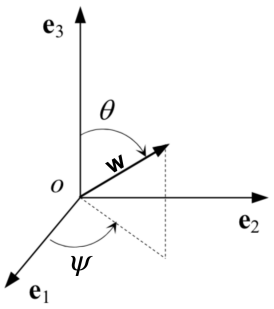

Figure 2.4: Physical angles.

$$\mathbf{w} = (\sin\theta\cos\psi, \sin\theta\sin\psi, \cos\theta) \;\rightarrow\; \theta = \arccos w_3,\; \psi = \arctan\frac{w_2}{w_1},$$

and, eq. (2.36),

$$\mathbf{R} = \begin{bmatrix} c\psi^2 s\theta^2 + c\varphi(c\theta^2 + s\psi^2 s\theta^2) & s\psi c\psi s\theta^2(1-c\varphi) - c\theta s\varphi & c\psi s\theta c\theta(1-c\varphi) + s\psi s\theta s\varphi \\ s\psi c\psi s\theta^2(1-c\varphi) + c\theta s\varphi & s\psi^2 s\theta^2 + c\varphi(c\theta^2 + c\psi^2 s\theta^2) & s\psi s\theta c\theta(1-c\varphi) - c\psi s\theta s\varphi \\ c\psi s\theta c\theta(1-c\varphi) - s\psi s\theta s\varphi & s\psi s\theta c\theta(1-c\varphi) + c\psi s\theta s\varphi & c\theta^2 + c\varphi(c\psi^2 s\theta^2 + s\psi^2 s\theta^2) \end{bmatrix},$$

where $c\psi = \cos\psi, s\psi = \sin\psi, c\theta = \cos\theta, s\theta = \sin\theta, c\varphi = \cos\varphi, s\varphi = \sin\varphi$. We remark that all the components of $\mathbf{R}$ so expressed depend upon the first powers of the circular functions of φ. Hence, for what said above, with this representation of the rotations, the components of a rotated r-rank tensor depend upon the r-th power of the circular functions of φ, i.e. of the physical rotation, but not on ψ nor on θ.

ii. *Euler's angles*: in this case the three parameters are the amplitude of three particular rotations into which the rotation is decomposed. Such parameters are the angles ψ, the *precession*, θ, the *nutation*, and φ, the *proper rotation*, see Fig. 2.5 These three

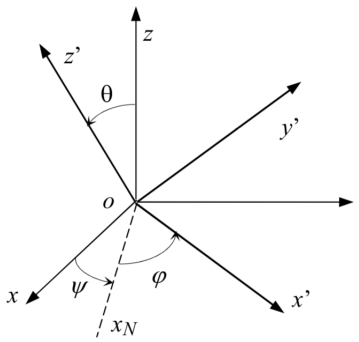

Figure 2.5: Euler's angles.

rotations are represented in Fig. 2.6. The first one, of amplitude ψ, is made about z to carry the axis x onto the *knots line* x_N, the line perpendicular to both the axes z

and z', and y onto $\overline{y}$; by eq. (2.32), in the frame $\{x, y, z\}$ it is

$$\mathbf{R}_\psi = \begin{bmatrix} \cos\psi & -\sin\psi & 0 \\ \sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix}.$$

The second one, of amplitude θ, is made about x_N to carry z onto z'; in the frame $\{x_N, \overline{y}, z\}$ it is

$$\mathbf{R}_\theta = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{bmatrix},$$

while in the frame $\{x, y, z\}$

$$\mathbf{R}_\theta^o = (\mathbf{R}_\psi^{-1})^\top \mathbf{R}_\theta \mathbf{R}_\psi^{-1} = \mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\psi^\top.$$

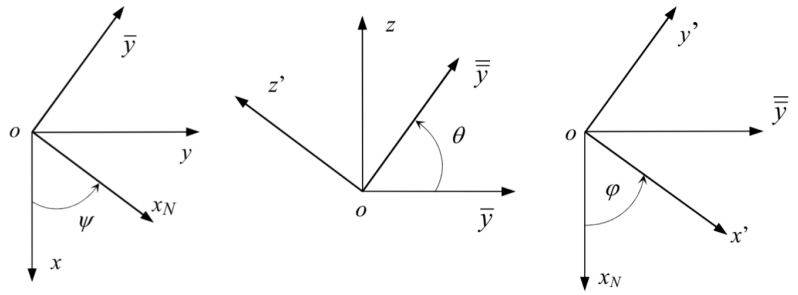

Figure 2.6: Euler's rotations, as seen from the respective axes of rotation.

The last rotation, of amplitude φ, is made about z' to carry x_N onto x' and $\overline{\overline{y}}$ onto y'; in the frame $\{x_N, \overline{\overline{y}}, z'\}$ it is

$$\mathbf{R}_\varphi = \begin{bmatrix} \cos\varphi & -\sin\varphi & 0 \\ \sin\varphi & \cos\varphi & 0 \\ 0 & 0 & 1 \end{bmatrix},$$

while in $\{x, y, z\}$

$$\mathbf{R}_\varphi^o = (\mathbf{R}_\psi^{-1})^\top (\mathbf{R}_\theta^{-1})^\top \mathbf{R}_\varphi \mathbf{R}_\theta^{-1} \mathbf{R}_\psi^{-1} = \mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\varphi \mathbf{R}_\theta^\top \mathbf{R}_\psi^\top.$$

Any vector $\mathbf{u}$ is transformed, by the global rotation, into the vector

$$\mathbf{u}' = \mathbf{R}\mathbf{u}.$$

But we can write also

$$\mathbf{u}' = \mathbf{R}_\varphi^o \overline{\overline{\mathbf{u}}},$$

where $\overline{\overline{\mathbf{u}}}$ is the vector transformed by rotation $\mathbf{R}_\theta^o$,

$$\overline{\overline{\mathbf{u}}} = \mathbf{R}_\theta^o \overline{\mathbf{u}},$$

and $\overline{\mathbf{u}}$ is the vector transformed by rotation $\mathbf{R}_\psi$,

$$\overline{\mathbf{u}} = \mathbf{R}_\psi \mathbf{u}.$$

Finally,

$$\mathbf{u}' = \mathbf{R}\mathbf{u} = \mathbf{R}^o_\varphi \mathbf{R}^o_\theta \mathbf{R}_\psi \mathbf{u} \;\rightarrow\; \mathbf{R} = \mathbf{R}^o_\varphi \mathbf{R}^o_\theta \mathbf{R}_\psi,$$

i.e. the global rotation tensor is obtained composing, in the opposite order of execution of the rotations, the three tensors all expressed in the original basis. However,

$$\mathbf{R} = \mathbf{R}^o_\varphi \mathbf{R}^o_\theta \mathbf{R}_\psi = \mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\varphi \mathbf{R}_\theta^\top \mathbf{R}_\psi^\top \mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\psi^\top \mathbf{R}_\psi = \mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\varphi,$$

i.e., the global rotation tensor is also equal to the composition of the three rotations, in the order of execution, if the three rotations are expressed in their own particular bases. This result is general, not bounded to the Euler's rotations nor to three rotations.

Performing the tensor multiplications we get

$$\mathbf{R} = \begin{bmatrix} \cos\psi\cos\varphi - \sin\psi\sin\varphi\cos\theta & -\cos\psi\sin\varphi - \sin\psi\cos\varphi\cos\theta & \sin\psi\sin\theta \\ \sin\psi\cos\varphi + \cos\psi\sin\varphi\cos\theta & -\sin\psi\sin\varphi + \cos\psi\cos\varphi\cos\theta & -\cos\psi\sin\theta \\ \sin\varphi\sin\theta & \cos\varphi\sin\theta & \cos\theta \end{bmatrix}.$$

The components of a vector $\mathbf{u}$ in the basis $\mathcal{B}'$ are then given by

$$\mathbf{u}' = \mathbf{R}^\top \mathbf{u} = \mathbf{R}^\top_\varphi \mathbf{R}^\top_\theta \mathbf{R}^\top_\psi \mathbf{u},$$

and those of a second-rank tensor

$$\mathbf{L}' = \mathbf{R}^\top \mathbf{L}\mathbf{R} = \mathbf{R}^\top_\varphi \mathbf{R}^\top_\theta \mathbf{R}^\top_\psi \mathbf{L}\mathbf{R}_\psi \mathbf{R}_\theta \mathbf{R}_\varphi.$$

iii. *coordinate angles*: in this case, the rotation $\mathbf{R}$ is decomposed into three successive rotations α, β, γ, respectively about the axes x, y and z of each rotation, i.e.

$$\mathbf{R} = \mathbf{R}_\alpha \mathbf{R}_\beta \mathbf{R}_\gamma$$

with

$$\mathbf{R}_\alpha = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\alpha & -\sin\alpha \\ 0 & \sin\alpha & \cos\alpha \end{bmatrix}, \mathbf{R}_\beta = \begin{bmatrix} \cos\beta & 0 & -\sin\beta \\ 0 & 1 & 0 \\ \sin\beta & 0 & \cos\beta \end{bmatrix}, \mathbf{R}_\gamma = \begin{bmatrix} \cos\gamma & -\sin\gamma & 0 \\ \sin\gamma & \cos\gamma & 0 \\ 0 & 0 & 1 \end{bmatrix},$$

so finally

$$\mathbf{R} = \begin{bmatrix} \cos\beta\cos\gamma & -\cos\beta\sin\gamma & -\sin\beta \\ \cos\alpha\sin\gamma - \sin\alpha\sin\beta\cos\gamma & \cos\alpha\cos\gamma + \sin\alpha\sin\beta\sin\gamma & -\sin\alpha\cos\beta \\ \sin\alpha\sin\gamma + \cos\alpha\sin\beta\cos\gamma & \sin\alpha\cos\gamma - \cos\alpha\sin\beta\sin\gamma & \cos\alpha\cos\beta \end{bmatrix}.$$

Let us now consider the case of *small rotations*, i.e. $|\varphi| \to 0$. In such a case,

$$\sin\varphi \simeq \varphi, \quad 1 - \cos\varphi \simeq 0$$

and

$$\mathbf{R} \simeq \mathbf{I} + \varphi\mathbf{W},$$

i.e. in the small rotations approximation, any vector $\mathbf{u}$ is transformed as

$$\mathbf{Ru} \simeq (\mathbf{I} + \varphi\mathbf{W})\mathbf{u} = \mathbf{u} + \varphi\mathbf{w} \times \mathbf{u}, \tag{2.41}$$

i.e. by a skew tensor, not by a rotation tensor. The term $(1 - \cos\varphi)\mathbf{W}^2\mathbf{u}$ has disappeared (i.e., it is a higher order infinitesimal quantity) and the term $\varphi\mathbf{w} \times \mathbf{u}$ is orthogonal to $\mathbf{u}$. Because $\varphi \to 0$, the arc is approximated by its tangent, the vector $\varphi\mathbf{w} \times \mathbf{u}$, see Fig. 2.7. Applying to eq. (2.41) the procedure already seen for the composition of finite amplitude

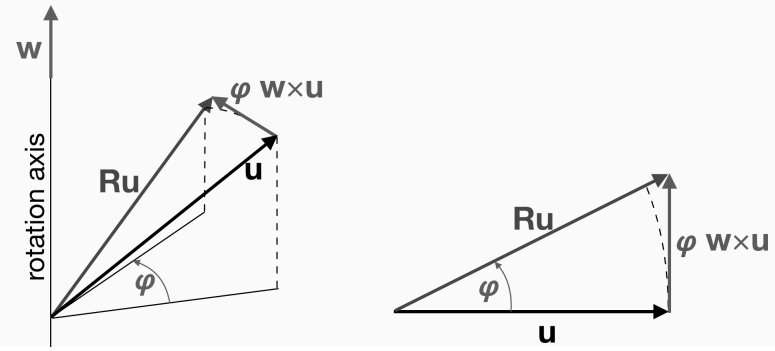

Figure 2.7: Small rotations.

rotations, we get

$$\begin{aligned}
\mathbf{u}_1 &= \mathbf{R}_1\mathbf{u} = (\mathbf{I} + \varphi_1\mathbf{W}_1)\mathbf{u} = \mathbf{u} + \varphi_1\mathbf{w}_1 \times \mathbf{u}, \\
\mathbf{u}_{21} &= \mathbf{R}_2\mathbf{u}_1 = (\mathbf{I} + \varphi_2\mathbf{W}_2)\mathbf{u}_1 = \mathbf{u}_1 + \varphi_2\mathbf{w}_2 \times \mathbf{u}_1 \\
&= \mathbf{u} + \varphi_1\mathbf{w}_1 \times \mathbf{u} + \varphi_2\mathbf{w}_2 \times \mathbf{u} \\
&\quad + \varphi_1\varphi_2\mathbf{w}_2 \times (\mathbf{w}_1 \times \mathbf{u}).
\end{aligned}$$

If the order of the rotations is changed, the last term becomes $\varphi_1\varphi_2\mathbf{w}_1 \times (\mathbf{w}_2 \times \mathbf{u})$, which is, in general, different from $\varphi_1\varphi_2\mathbf{w}_2 \times (\mathbf{w}_1 \times \mathbf{u})$: strictly speaking, also small rotations do not commute[10]. However, for small rotations, $\varphi_1\varphi_2$ is negligible with respect to φ_1 and φ_2: in this approximation, small rotations commute. To remark that the approximation (2.41) gives, for the displacements, a law which is quite similar to that of the velocities of the points of a rigid body:

$$\mathbf{v} = \mathbf{v}_0 + \boldsymbol{\omega} \times (p - o)$$

This is quite natural, because

$$\omega = \frac{d\varphi}{dt},$$

i.e. a small amplitude rotation can be seen as the rotation made with finite angular velocity $\boldsymbol{\omega}$ in a small time interval dt.

[10]This can happen for some vectors, all the times that $\mathbf{w}_1 \cdot \mathbf{u} = \mathbf{w}_2 \cdot \mathbf{u}$, like for the case of a vector $\mathbf{u}$ orthogonal to both $\mathbf{w}_1$ and $\mathbf{w}_2$; however, this is no more than a curiosity, it has no importance in practice.

2.12 Symmetries

Let us consider now tensors $\mathbf{S} \in Orth(\mathcal{V})$ that are not a rotation, i.e. such that $\det \mathbf{S} = -1$. Let us call $\mathbf{S}$ an *improper rotation*. A particular improper rotation, whose all eigenvalues are equal to -1, is the *inversion* or *reflexion tensor*

$$\mathbf{S}_I = -\mathbf{I}.$$

The effect of $\mathbf{S}_I$ is to transform any basis $\mathcal{B}$ into the basis $-\mathcal{B}$, i.e. with all the basis vectors changed of orientation (or, which is the same, to change the sign of all the components of a vector). In other words, $\mathbf{S}_I$ changes the orientation of the space. This is also the effect of any other improper rotation $\mathbf{S}$, that can be decomposed into a proper rotation $\mathbf{R}$ followed by the reflexion $\mathbf{S}_I$[11]:

$$\mathbf{S} = \mathbf{S}_I\mathbf{R}. \tag{2.42}$$

Be $\mathbf{n} \in \mathcal{S}$, then

$$\mathbf{S}_R = \mathbf{I} - 2\mathbf{n} \otimes \mathbf{n} \tag{2.43}$$

is the tensor that operates the transformation of symmetry with respect to a plane orthogonal to $\mathbf{n}$. In fact

$$\mathbf{S}_R\mathbf{n} = -\mathbf{n}, \quad \mathbf{S}_R\mathbf{m} = \mathbf{m}\ \forall \mathbf{m} \in \mathcal{V} : \ \mathbf{m} \cdot \mathbf{n} = 0.$$

$\mathbf{S}_R$ is an improper rotation; in fact, eqs. (2.4), (2.14) and exercice 11,

$$\begin{aligned}(\mathbf{I} - 2\mathbf{n} \otimes \mathbf{n})(\mathbf{I} - 2\mathbf{n} \otimes \mathbf{n})^\top &= (\mathbf{I} - 2\mathbf{n} \otimes \mathbf{n})(\mathbf{I} - 2\mathbf{n} \otimes \mathbf{n}) \\ &= \mathbf{I} - 2\mathbf{n} \otimes \mathbf{n} - 2\mathbf{n} \otimes \mathbf{n} + 4(\mathbf{n} \otimes \mathbf{n})(\mathbf{n} \otimes \mathbf{n}) = \mathbf{I},\end{aligned}$$

$$\det(\mathbf{I} - 2\mathbf{n} \otimes \mathbf{n}) = 1 - 2\mathrm{tr}(\mathbf{n} \otimes \mathbf{n}) + 4\frac{\mathrm{tr}^2(\mathbf{n} \otimes \mathbf{n}) - \mathrm{tr}(\mathbf{n} \otimes \mathbf{n})(\mathbf{n} \otimes \mathbf{n})}{2} - 8\det(\mathbf{n} \otimes \mathbf{n}) = -1.$$

Be $\mathbf{S} = \mathbf{S}_I\mathbf{R}$ an improper rotation; then

$$\begin{aligned}(\mathbf{S}\mathbf{u}) \times (\mathbf{S}\mathbf{v}) &= (\mathbf{S}_I\mathbf{R}\mathbf{u}) \times (\mathbf{S}_I\mathbf{R}\mathbf{v}) = \det(\mathbf{S}_I\mathbf{R}) \left[(\mathbf{S}_I\mathbf{R})^{-1}\right]^\top (\mathbf{u} \times \mathbf{v}) \\ &= \det \mathbf{S}_I \det \mathbf{R}(\mathbf{R}^{-1}\mathbf{S}_I^{-1})^\top(\mathbf{u} \times \mathbf{v}) = -(-\mathbf{R}^{-1}\mathbf{I})^\top(\mathbf{u} \times \mathbf{v}) = \mathbf{R}(\mathbf{u} \times \mathbf{v}).\end{aligned}$$

The transformation by $\mathbf{S}$ of any vector $\mathbf{u}$ gives

$$\mathbf{S}\mathbf{u} = \mathbf{S}_I\mathbf{R}\mathbf{u} = -\mathbf{R}\mathbf{u},$$

i.e. it changes the orientation of the rotated vector; this is not the case when the same improper rotations transforms the vectors of a cross product: the rotated vector result of the cross product does not change of orientation, i.e. the cross product is insensitive to a reflexion. That is why, strictly speaking, the result of a cross product is not a vector, but a *pseudo-vector*: it behaves like vectors apart for the reflexions. For the same reason a scalar result of a mixed product (scalar plus cross product of three vectors) is called a *pseudo-scalar*, because in this case the scalar result of the mixed product changes of sign under a reflexion, as it is easy to be seen.

[11] The application of the Binet's Theorem shows immediately that $\det \mathbf{S} = -1$, while $\mathbf{S}_I\mathbf{R}(\mathbf{S}_I\mathbf{R})^\top = \mathbf{S}_I\mathbf{R}\mathbf{R}^\top\mathbf{S}_I^\top = -\mathbf{I}(-\mathbf{I})^\top = \mathbf{I}$: the decomposition in eq. (2.42) actually gives an improper rotation.

2.13 Polar decomposition

Theorem 14. *(Square root theorem): be* $\mathbf{L} \in Sym(\mathcal{V})$ *and positive definite; then is exists a unique tensor* $\mathbf{U} \in Sym(\mathcal{V})$ *and positive definite such that*

$$\mathbf{L} = \mathbf{U}^2.$$

Proof. Existence: be $\mathbf{L}, \mathbf{U}, \mathbf{V} \in Sym(\mathcal{V})$ positive definite and

$$\mathbf{L} = \omega_i \mathbf{e}_i \otimes \mathbf{e}_i$$

a spectral decomposition of $\mathbf{L}$, $\omega_i > 0 \ \forall i$. Define $\mathbf{U}$ as

$$\mathbf{U} = \sqrt{\omega_i} \mathbf{e}_i \otimes \mathbf{e}_i;$$

then, by eq. $(2.4)_1$ we get

$$\mathbf{U}^2 = \mathbf{L}.$$

Uniqueness: suppose that also

$$\mathbf{V}^2 = \mathbf{L}$$

and be $\mathbf{e}$ an eigenvector of $\mathbf{L}$ corresponding to the (positive) eigenvalue ω. Then, if $\lambda = \sqrt{\omega}$,

$$\mathbf{O} = (\mathbf{U}^2 - \lambda \mathbf{I})\mathbf{e} = (\mathbf{U} - \lambda \mathbf{I})(\mathbf{U} - \lambda \mathbf{I})\mathbf{e},$$

and put

$$\mathbf{v} = (\mathbf{U} - \lambda \mathbf{I})\mathbf{e},$$

we get

$$\mathbf{U}\mathbf{v} = -\lambda \mathbf{v} \ \Rightarrow \ \mathbf{v} = \mathbf{o} \ \Rightarrow \ \mathbf{U}\mathbf{e} = \lambda \mathbf{e}$$

because $\mathbf{U}$ is positive definite and $-\lambda$ cannot be an eigenvalue of $\mathbf{U}$, because $\lambda > 0$. In the same way

$$\mathbf{V}\mathbf{e} = \lambda \mathbf{e} \ \Rightarrow \ \mathbf{U}\mathbf{e} = \mathbf{V}\mathbf{e}$$

for every eigenvector $\mathbf{e}$ of $\mathbf{L}$. Because, spectral theorem, it exists a basis of eigenvectors of $\mathbf{L}$, $\mathbf{U} = \mathbf{V}$. □

We symbolically write that

$$\mathbf{U} = \sqrt{\mathbf{L}}.$$

For any $\mathbf{F} \in Lin(\mathcal{V})$, both $\mathbf{F}\mathbf{F}^\top$ and $\mathbf{F}^\top\mathbf{F}$ clearly $\in Sym(\mathcal{V})$. If in addition $\det \mathbf{F} > 0$, then

$$\mathbf{u} \cdot \mathbf{F}^\top \mathbf{F}\mathbf{u} = (\mathbf{F}\mathbf{u}) \cdot (\mathbf{F}\mathbf{u}) \geq 0$$

and the zero value is obtained $\iff \mathbf{F}\mathbf{u} = \mathbf{o}$ and because $\mathbf{F}$ is invertible, $\iff \mathbf{u} = \mathbf{o}$. As a consequence, $\mathbf{F}^\top\mathbf{F}$ is positive definite. In the same way it can be proved that $\mathbf{F}\mathbf{F}^\top$ is also positive definite.

An important tensor decomposition is given by the

Theorem 15. *(Polar decomposition theorem):* $\forall \mathbf{F} \in Lin(\mathcal{V}) | \det \mathbf{F} > 0$ *exist and are uniquely determined two positive definite tensors* $\mathbf{U}, \mathbf{V} \in Sym(\mathcal{V})$ *and a rotation* $\mathbf{R}$ *such that*

$$\mathbf{F} = \mathbf{RU} = \mathbf{VR}.$$

Proof. Uniqueness: Be $\mathbf{F} = \mathbf{RU}$ a *right polar decomposition* of $\mathbf{F}$; because $\mathbf{R} \in Orth(\mathcal{V})^+$ and $\mathbf{U} \in Sym(\mathbf{V})$,

$$\mathbf{F}^\top \mathbf{F} = \mathbf{U}\mathbf{R}^\top \mathbf{R}\mathbf{U} = \mathbf{U}^2 \rightarrow \mathbf{U} = \sqrt{\mathbf{F}^\top \mathbf{F}}.$$

By the Square-root Theorem, tensor $\mathbf{U}$ is unique, and because

$$\mathbf{R} = \mathbf{F}\mathbf{U}^{-1},$$

$\mathbf{R}$ is unique too.

Be now $\mathbf{F} = \mathbf{VR}$ a *left polar decomposition* of $\mathbf{F}$; by the same procedure, we get

$$\mathbf{F}\mathbf{F}^\top = \mathbf{V}^2 \rightarrow \mathbf{V} = \sqrt{\mathbf{F}\mathbf{F}^\top},$$

so $\mathbf{V}$ is unique and also

$$\mathbf{R} = \mathbf{V}^{-1}\mathbf{F}.$$

Existence: be

$$\mathbf{U} = \sqrt{\mathbf{F}^\top \mathbf{F}}$$

so $\mathbf{U} \in Sym(\mathcal{V})$ and it is positive definite, and let

$$\mathbf{R} = \mathbf{F}\mathbf{U}^{-1}.$$

To prove that $\mathbf{F} = \mathbf{RU}$ is a right polar decomposition, we just have to show that $\mathbf{R} \in Orth(\mathcal{V})^+$. Since $\det \mathbf{F} > 0, \det \mathbf{U} > 0$ (the latter because all the eigenvalues of $\mathbf{U}$ are strictly positive), by the Theorem of Binet also $\det \mathbf{R} > 0$. Then

$$\mathbf{R}^\top \mathbf{R} = (\mathbf{F}\mathbf{U}^{-1})^\top (\mathbf{F}\mathbf{U}^{-1}) = \mathbf{U}^{-1}\mathbf{F}^\top \mathbf{F}\mathbf{U}^{-1} = \mathbf{U}^{-1}\mathbf{U}^2\mathbf{U}^{-1} = \mathbf{I} \Rightarrow \mathbf{R} \in Orth(\mathcal{V})^+.$$

Let now

$$\mathbf{V} = \mathbf{R}\mathbf{U}\mathbf{R}^\top;$$

then $\mathbf{V} \in Sym(\mathcal{V})$ and is positive definite, see exercice 16, and

$$\mathbf{VR} = \mathbf{R}\mathbf{U}\mathbf{R}^\top\mathbf{R} = \mathbf{RU} = \mathbf{F},$$

which completes the proof. □

2.14 Exercices

1. Prove that
$$\mathbf{Lo} = \mathbf{o} \quad \forall \mathbf{L} \in Lin(\mathcal{V}).$$

2. Prove that, if a straight line r has the direction of $\mathbf{u} \in \mathcal{S}$, then the tensor giving the projection of a vector $\mathbf{u} \in \mathcal{V}$ on r is $\mathbf{u} \otimes \mathbf{u}$, while the one giving the projection on a direction orthogonal to r is $\mathbf{I} - \mathbf{u} \otimes \mathbf{u}$.

3. For any $\alpha \in \mathbb{R}, \mathbf{a}, \mathbf{b} \in \mathcal{V}$ and $\mathbf{A}, \mathbf{B} \in Lin(\mathcal{V})$, prove that

$$(\alpha\mathbf{A})^\top = \alpha\mathbf{A}^\top, \quad (\mathbf{A}+\mathbf{B})^\top = \mathbf{A}^\top + \mathbf{B}^\top, \quad (\mathbf{a}\otimes\mathbf{b})\mathbf{A} = \mathbf{a}\otimes(\mathbf{A}^\top\mathbf{b}).$$

4. Prove that

$$\mathrm{tr}\mathbf{I} = 3, \quad \mathrm{tr}\mathbf{O} = 0, \quad \mathrm{tr}(\mathbf{AB}) = \mathrm{tr}(\mathbf{BA}) \quad \forall \mathbf{A}, \mathbf{B} \in Lin(\mathcal{V}).$$

5. Prove that, $\forall \mathbf{L}, \mathbf{M}, \mathbf{N} \in Lin(\mathcal{V})$,

$$\mathbf{L}^\top \cdot \mathbf{M}^\top = \mathbf{L}\cdot\mathbf{M}, \quad \mathbf{LM}\cdot\mathbf{N} = \mathbf{L}\cdot\mathbf{NM}^\top = \mathbf{M}\cdot\mathbf{L}^\top\mathbf{N}.$$

6. Prove that $Sym(\mathcal{V})$ and $Skw(\mathcal{V})$ are orthogonal, i.e. prove that

$$\mathbf{A}\cdot\mathbf{B} = 0 \quad \forall \mathbf{A} \in Sym(\mathcal{V}), \ \mathbf{B} \in Skw(\mathcal{V}).$$

7. For any $\mathbf{L} \in Lin(\mathcal{V})$, prove that, if $\mathbf{A} \in Sym(\mathcal{V})$, then

$$\mathbf{A}\cdot\mathbf{L} = \mathbf{A}\cdot\mathbf{L}^s,$$

   while if $\mathbf{B} \in Skw(\mathcal{V})$, then

$$\mathbf{B}\cdot\mathbf{L} = \mathbf{B}\cdot\mathbf{L}^a.$$

8. Express by components the second principal invariant I_2 of a tensor $\mathbf{L}$.
9. Prove that, if $\mathbf{a} = (a_1, a_2, a_3), \mathbf{b} = (b_1, b_2, b_3), \mathbf{c} = (c_1, c_2, c_3)$, then

$$\mathbf{a}\times\mathbf{b}\cdot\mathbf{c} = \det\begin{bmatrix} a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{bmatrix}.$$

10. Prove the uniqueness of the inverse tensor.
11. Prove that all the dyads are singular.
12. Prove that if $\mathbf{L}$ is invertible and $\alpha \in \mathbb{R} - \{0\}$ then

$$(\alpha\mathbf{L})^{-1} = \alpha^{-1}\mathbf{L}^{-1}.$$

13. Prove that any quadratic form defined by a tensor $\mathbf{L}$ can be written as a scalar product of tensors:

$$\mathbf{u}\cdot\mathbf{Lu} = \mathbf{L}\cdot\mathbf{u}\otimes\mathbf{u} \quad \forall \mathbf{u} \in \mathcal{V}, \mathbf{L} \in Lin(\mathcal{V}).$$

14. Prove that, if $\mathbf{W}$ is the axial tensor of $\mathbf{w}$, then

$$\mathbf{WW} = -\frac{1}{2}|\mathbf{W}|^2(\mathbf{I} - \mathbf{w}\otimes\mathbf{w}).$$

15. Prove that for any two axial couples $\mathbf{w}_1, \mathbf{W}_1$ and $\mathbf{w}_2, \mathbf{W}_2$, it is:

$$\mathbf{w}_1 \cdot \mathbf{w}_2 = \frac{1}{2}\mathbf{W}_1 \cdot \mathbf{W}_2.$$

16. Let $\mathbf{L} \in Sym(\mathcal{V})$ and positive definite and $\mathbf{R} \in Orth(\mathcal{V})^+$; then prove that $\mathbf{RLR}^\top \in Sym(\mathcal{V})$ and that it is positive definite.

17. Let $\mathbf{A}, \mathbf{B}, \mathbf{C}, \mathbf{D} \in Lin(\mathcal{V})$; prove that

$$\mathbf{A} \cdot \mathbf{BCD} = \mathbf{B}^\top \mathbf{A} \cdot \mathbf{CD} = \mathbf{AD}^\top \cdot \mathbf{BC}.$$

18. Prove that the spectrum of $\mathbf{L}^{sph}$ is composed by only

$$\lambda^{sph} = \frac{1}{3}\mathrm{tr}\mathbf{L},$$

and that any $\mathbf{u} \in \mathcal{S}$ is an eigenvector.

19. Prove that the eigenvalues λ^{dev} of $\mathbf{L}^{dev}$ are given by

$$\lambda^{dev} = \lambda - \lambda^{sph},$$

where λ is an eigenvalue of $\mathbf{L}$.

Chapter 3

Fourth rank tensors

3.1 Fourth-rank tensors

A *fourth-rank tensor* $\mathbb{L}$ is any linear application from $Lin(\mathcal{V})$ to $Lin(\mathcal{V})$:

$$\mathbb{L} : Lin(\mathcal{V}) \to Lin(\mathcal{V}) | \mathbb{L}(\alpha_i \mathbf{A}_i) = \alpha_i \mathbb{L}\mathbf{A}_i \ \forall \alpha_i \in \mathbb{R}, \ \mathbf{A}_i \in Lin(\mathcal{V}), \ i = 1, ..., n.$$

Defining the *sum of two fourth-rank tensors* as

$$(\mathbb{L}_1 + \mathbb{L}_2)\mathbf{A} = \mathbb{L}_1\mathbf{A} + \mathbb{L}_2\mathbf{A} \quad \forall \mathbf{A} \in Lin(\mathcal{V}),$$

the *product of a scalar by a fourth-rank tensor* as

$$(\alpha\mathbb{L})\mathbf{A} = \alpha(\mathbb{L}\mathbf{A}) \quad \forall \alpha \in \mathbb{R}, \mathbf{A} \in Lin(\mathcal{V})$$

and the *null fourth-rank tensor* $\mathbb{O}$ as the unique tensor such that

$$\mathbb{O}\mathbf{A} = \mathbf{O} \ \forall \mathbf{A} \in Lin(\mathcal{V}),$$

then the set of all the tensors $\mathbb{L}$ that operate on $Lin(\mathcal{V})$ forms a vector space, denoted by $\mathbb{L}\text{in}(\mathcal{V})$. We define the *fourth-rank identity tensor* $\mathbb{I}$ as the unique tensor such that

$$\mathbb{I}\mathbf{A} = \mathbf{A} \quad \forall \mathbf{A} \in Lin(\mathcal{V}).$$

It is apparent that the algebra of fourth-rank tensors is similar to that of second-rank tensors and in fact the operations with fourth-rank tensors can be introduced in almost the same way, in some sense shifting from $\mathcal{V}$ to $Lin(\mathcal{V})$ the operations. However, the algebra of fourth-rank tensors is richer than that of the second-rank ones and some care must be paid.

In the following sections, we consider some of the operations that can be done with fourth-rank tensors.

3.2 Dyads, tensor components

For any couple of tensors $\mathbf{A}$ and $\mathbf{B} \in Lin(\mathcal{V})$, the *(tensor) dyad* $\mathbf{A} \otimes \mathbf{B}$ is the fourth-rank tensor defined by

$$(\mathbf{A} \otimes \mathbf{B})\mathbf{L} := \mathbf{B} \cdot \mathbf{L} \ \mathbf{A} \quad \forall \mathbf{L} \in Lin(\mathcal{V}).$$

The application defined above is actually a fourth-rank tensor because of the bi-linearity of the scalar product of second-rank tensors. Applying this rule to the nine dyads of the basis $\mathcal{B}^2 = \{\mathbf{e}_i \otimes \mathbf{e}_j,\ i,j = 1,2,3\}$ of $Lin(\mathcal{V})$ let us introduce the 81 fourth-rank tensors

$$\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l := (\mathbf{e}_i \otimes \mathbf{e}_j) \otimes (\mathbf{e}_k \otimes \mathbf{e}_l)$$

that form a basis $\mathcal{B}^4 = \{\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l,\ i,j = 1,2,3\}$ for $\mathbb{L}\mathrm{in}(\mathcal{V})$. We remark hence that $\dim(\mathbb{L}\mathrm{in}(\mathcal{V})) = 81$. A useful result is that

$$(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l)(\mathbf{e}_p \otimes \mathbf{e}_q) = (\mathbf{e}_k \otimes \mathbf{e}_l) \cdot (\mathbf{e}_p \otimes \mathbf{e}_q)(\mathbf{e}_i \otimes \mathbf{e}_j) = \delta_{kp}\delta_{lq}(\mathbf{e}_i \otimes \mathbf{e}_j). \tag{3.1}$$

Any fourth-rank tensor can be expressed as the linear combination (the *canonical decomposition*)

$$\mathbb{L} = L_{ijkl}\ \mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l, \quad i,j = 1,2,3,$$

where the L_{ijkl}s are the 81 *Cartesian components* of $\mathbb{L}$ with respect to $\mathcal{B}^4$. The L_{ijkl}s are defined by the operation:

$$\begin{aligned}(\mathbf{e}_i \otimes \mathbf{e}_j) \cdot \mathbb{L}(\mathbf{e}_k \otimes \mathbf{e}_l) &= (\mathbf{e}_i \cdot \mathbf{e}_j) \cdot (L_{pqrs}\mathbf{e}_p \otimes \mathbf{e}_q \otimes \mathbf{e}_r \otimes \mathbf{e}_s)(\mathbf{e}_k \otimes \mathbf{e}_l) \\ &= (\mathbf{e}_i \otimes \mathbf{e}_j) \cdot (L_{pqrs}\delta_{rk}\delta_{sl}\mathbf{e}_p \otimes \mathbf{e}_q) \\ &= L_{pqrs}\delta_{rk}\delta_{sl}\delta_{ip}\delta_{jq} = L_{ijkl}.\end{aligned}$$

The components of a tensor dyad can be computed without any difficulty:

$$\begin{aligned}\mathbf{A} \otimes \mathbf{B} = (A_{ij}\mathbf{e}_i \otimes \mathbf{e}_j) \otimes (B_{kl}\mathbf{e}_k \otimes \mathbf{e}_l) = A_{ij}B_{kl}\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l \ \Rightarrow \\ (\mathbf{A} \otimes \mathbf{B})_{ijkl} = A_{ij}B_{kl},\end{aligned}$$

so that in particular

$$((\mathbf{a} \otimes \mathbf{b}) \otimes (\mathbf{c} \otimes \mathbf{d}))_{ijkl} = a_i b_j c_k d_l.$$

Concerning the identity of $\mathbb{L}\mathrm{in}(\mathcal{V})$,

$$\begin{aligned}I_{ijkl} := (\mathbf{e}_i \otimes \mathbf{e}_l) \cdot \mathbb{I}(\mathbf{e}_k \otimes \mathbf{e}_l) = (\mathbf{e}_i \otimes \mathbf{e}_j) \cdot (\mathbf{e}_k \otimes \mathbf{e}_l) = \mathbf{e}_i \cdot \mathbf{e}_k \mathbf{e}_j \cdot \mathbf{e}_l = \delta_{ik}\delta_{jl} \ \rightarrow \\ \mathbb{I} = \delta_{ik}\delta_{jl}(\mathbf{e}_i \otimes \mathbf{e}_l \otimes \mathbf{e}_k \otimes \mathbf{e}_l).\end{aligned}$$

The components of $\mathbf{A} \in Lin(\mathcal{V})$ result of the application of $\mathbb{L} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ on $\mathbf{B} \in Lin(\mathcal{V})$ can now be easily calculated:

$$\begin{aligned}\mathbf{A} = \mathbb{L}\mathbf{B} &= L_{ijkl}(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l)(B_{pq}\mathbf{e}_p \otimes \mathbf{e}_q) \\ &= L_{ijkl}B_{pq}\delta_{kp}\delta_{lq}(\mathbf{e}_i \otimes \mathbf{e}_j) \\ &= L_{ijkl}B_{kl}(\mathbf{e}_i \otimes \mathbf{e}_j) \ \rightarrow \ A_{ij} = L_{ijkl}B_{kl}.\end{aligned} \tag{3.2}$$

Moreover,

$$\begin{aligned}\mathbb{L}(\mathbf{A} \otimes \mathbf{B})\mathbf{C} = \mathbb{L}((\mathbf{A} \otimes \mathbf{B})\mathbf{C}) = \mathbb{L}(\mathbf{B} \cdot \mathbf{C}\mathbf{A}) = \mathbf{B} \cdot \mathbf{C}\ \mathbb{L}\mathbf{A} = ((\mathbb{L}\mathbf{A}) \otimes \mathbf{B})\mathbf{C} \ \Rightarrow \\ \mathbb{L}(\mathbf{A} \otimes \mathbf{B}) = (\mathbb{L}\mathbf{A}) \otimes \mathbf{B}.\end{aligned}$$

Using this result and eq. (3.1), we can determine the components of a product of fourth-rank tensors:

$$\begin{aligned}\mathbb{A}\mathbb{B} &= A_{ijkl}(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l)B_{pqrs}(\mathbf{e}_p \otimes \mathbf{e}_q \otimes \mathbf{e}_r \otimes \mathbf{e}_s) \\ &= A_{ijkl}B_{pqrs}(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l)(\mathbf{e}_p \otimes \mathbf{e}_q) \otimes (\mathbf{e}_r \otimes \mathbf{e}_s) \\ &= A_{ijkl}B_{pqrs}[(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l)(\mathbf{e}_p \otimes \mathbf{e}_q)] \otimes (\mathbf{e}_r \otimes \mathbf{e}_s) \\ &= A_{ijkl}B_{pqrs}[\delta_{kp}\delta_{lq}(\mathbf{e}_i \otimes \mathbf{e}_j)] \otimes (\mathbf{e}_r \otimes \mathbf{e}_s) \\ &= A_{ijkl}B_{klrs}(\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_r \otimes \mathbf{e}_s) \ \Rightarrow \ (\mathbb{A}\mathbb{B})_{ijrs} = A_{ijkl}B_{klrs}.\end{aligned} \tag{3.3}$$

Depending upon four indices, a fourth-rank tensor $\mathbb{L}$ cannot be represented by a matrix; however, we will see in Sect. 3.8 that a matrix representation of a fourth-rank tensor is still possible, and that it is currently used in some cases, e.g. in elasticity.

3.3 Conjugation product, transpose, symmetries

For any two tensors $\mathbf{A}, \mathbf{B} \in Lin(\mathcal{V})$ we call *conjugation product* the the tensor $\mathbf{A} \boxtimes \mathbf{B} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ defined by the operation

$$(\mathbf{A} \boxtimes \mathbf{B})\mathbf{L} := \mathbf{A}\mathbf{L}\mathbf{B}^\top \ \forall \mathbf{L} \in Lin(\mathcal{V}).$$

As a consequence, for the vectors of $\mathcal{B}$,

$$(\mathbf{e}_i \otimes \mathbf{e}_j) \boxtimes (\mathbf{e}_k \otimes \mathbf{e}_l) = \mathbf{e}_i \otimes \mathbf{e}_k \otimes \mathbf{e}_j \otimes \mathbf{e}_l, \tag{3.4}$$

so that

$$(\mathbf{A} \boxtimes \mathbf{B})_{ijkl} = A_{ik}B_{jl}.$$

Moreover, by the uniqueness of the identity $\mathbb{I}$, $\forall \mathbf{A} \in Lin(\mathcal{V})$,

$$(\mathbf{I} \boxtimes \mathbf{I})\mathbf{A} = \mathbf{I}\mathbf{A}\mathbf{I}^\top = \mathbf{A} \ \Rightarrow \ \mathbb{I} = \mathbf{I} \boxtimes \mathbf{I}.$$

The *transpose* of a fourth-rank tensor $\mathbb{L}$ is the unique tensor $\mathbb{L}^\top$ such that

$$\mathbf{A} \cdot (\mathbb{L}\mathbf{B}) = \mathbf{B} \cdot (\mathbb{L}^\top \mathbf{A}) \ \forall \mathbf{A}, \mathbf{B} \in Lin(\mathcal{V}).$$

By this definition, putting $\mathbf{A} = \mathbf{e}_i \otimes \mathbf{e}_j, \mathbf{B} = \mathbf{e}_k \otimes \mathbf{e}_l$ gives

$$(L^\top)_{ijkl} = L_{klij}.$$

A consequence is that

$$\mathbf{A} \cdot (\mathbb{L}\mathbf{B}) = \mathbf{B} \cdot (\mathbb{L}^\top \mathbf{A}) = \mathbf{A} \cdot (\mathbb{L}^\top)^\top \mathbf{B} \ \Rightarrow \ (\mathbb{L}^\top)^\top = \mathbb{L}.$$

Then, using

$$\begin{aligned}
\mathbf{M} \cdot (\mathbf{A} \otimes \mathbf{B})^\top \mathbf{L} &= \mathbf{L} \cdot (\mathbf{A} \otimes \mathbf{B})\mathbf{M} \\
&= \mathbf{L} \cdot \mathbf{A}\mathbf{M} \cdot \mathbf{B} = \mathbf{M} \cdot (\mathbf{B}\mathbf{A} \cdot \mathbf{L}) \\
&= \mathbf{M} \cdot (\mathbf{B} \otimes \mathbf{A})\mathbf{L}, \\
\mathbf{M} \cdot (\mathbf{A} \boxtimes \mathbf{B})^\top \mathbf{L} &= \mathbf{L} \cdot (\mathbf{A} \boxtimes \mathbf{B})\mathbf{M} \\
&= \mathbf{L} \cdot \mathbf{A}\mathbf{M}\mathbf{B}^\top = \mathbf{A}^\top \mathbf{L} \cdot \mathbf{M}\mathbf{B}^\top = \mathbf{M}^\top \mathbf{A}^\top \mathbf{L} \cdot \mathbf{B}^\top \\
&= (\mathbf{M}^\top \mathbf{A}^\top \mathbf{L})^\top \cdot (\mathbf{B}^\top)^\top = \mathbf{L}^\top \mathbf{A}\mathbf{M} \cdot \mathbf{B} = \mathbf{A}\mathbf{M} \cdot \mathbf{L}\mathbf{B} \\
&= \mathbf{M} \cdot \mathbf{A}^\top \mathbf{L}\mathbf{B} = \mathbf{M} \cdot (\mathbf{A}^\top \boxtimes \mathbf{B}^\top)\mathbf{L},
\end{aligned}$$

so that

$$\begin{aligned}
(\mathbf{A} \otimes \mathbf{B})^\top &= \mathbf{B} \otimes \mathbf{A}, \\
(\mathbf{A} \boxtimes \mathbf{B})^\top &= \mathbf{A}^\top \boxtimes \mathbf{B}^\top.
\end{aligned}$$

A tensor $\mathbb{L} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ is *symmetric* $\iff \mathbb{L} = \mathbb{L}^\top$. It is then evident that

$$\mathbb{L} = \mathbb{L}^\top \Rightarrow \; L_{ijkl} = L_{klij},$$

relations called *major symmetries*. These symmetries are 36 on the whole, so that a symmetric fourth-rank tensor has 45 independent components. Moreover,

$$\begin{aligned}
&\mathbf{A} \boxtimes \mathbf{B} = (\mathbf{A} \boxtimes \mathbf{B})^\top = \mathbf{A}^\top \boxtimes \mathbf{B}^\top \iff \mathbf{A} = \mathbf{A}^\top, \mathbf{B} = \mathbf{B}^\top, \\
&\mathbf{A} \otimes \mathbf{B} = (\mathbf{A} \otimes \mathbf{B})^\top = \mathbf{B} \otimes \mathbf{A} \iff \mathbf{B} = \lambda \mathbf{A}, \; \lambda \in \mathbb{R}.
\end{aligned}$$

Let us now consider the case of a $\mathbb{L} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ such that

$$\mathbb{L}\mathbf{A} = (\mathbb{L}\mathbf{A})^\top \; \forall \mathbf{A} \in Lin(\mathcal{V}).$$

Then, by eq. (3.2),

$$L_{ijkl} = L_{jikl},$$

relations called *left minor symmetries*: a tensor $\mathbb{L}$ having the left minor symmetries has values in $Sym(\mathcal{V})$. On the whole, the left minor symmetries are 27. Finally, consider the case of a $\mathbb{L} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ such that

$$\mathbb{L}\mathbf{A} = \mathbb{L}(\mathbf{A}^\top) \; \forall \mathbf{A} \in Lin(\mathcal{V});$$

then, again by eq. (3.2), we get

$$L_{ijlk} = L_{jilk},$$

relations called *minor right-symmetries*, whose total number is also 27. It is immediate to recognize that if $\mathbb{L}$ has the minor right-symmetries, then

$$\mathbb{L}\mathbf{W} = \mathbf{O} \; \forall \mathbf{W} \in Skw(\mathcal{V}).$$

We say that a tensor has the minor symmetries if it has both the right and left minor symmetries; the total number of minor symmetries is 45, because actually some of the left and right minor symmetries are the same, so a tensor with the minor symmetries has 36 independent components.

If $\mathbb{L} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ as the major and minor symmetries, then the independent symmetry relations are actually 60 (some minor and major symmetries coincide), so in such a case $\mathbb{L}$ depends upon 21 independent components only. This is the case of the elasticity tensor.

Finally, the 6 *Cauchy-Poisson symmetries* are those of the type

$$L_{ijkl} = L_{ikjl}.$$

A tensor having the major, minor and Cauchy-Poisson symmetries is *completely symmetric*, i.e. swapping any couple of indices gives an identical component. In that case, the number of independent components is of only 15.

3.4 Trace and scalar product of fourth-rank tensors

We can introduce the scalar product between fourth-rank tensors in the same way used for second-rank tensors. We first introduce the concept of *trace for fourth-rank tensors* once again using the dyad (here, the tensor dyad):

$$\mathrm{tr}_4\mathbf{A}\otimes\mathbf{B} := \mathbf{A}\cdot\mathbf{B}.$$

The easy proof that $\mathrm{tr}_4 : \mathbb{L}\mathrm{in}(\mathcal{V}) \to \mathbb{R}$ is a linear form is based upon the properties of scalar product of second-rank tensors and it is left to the reader. An immediate result is that

$$\mathrm{tr}_4\mathbf{A}\otimes\mathbf{B} = A_{ij}B_{ij},$$

Then, using the canonical decomposition, we have that

$$\mathrm{tr}_4\mathbb{L} = \mathrm{tr}_4(L_{ijkl}(\mathbf{e}_i\otimes\mathbf{e}_j)\otimes(\mathbf{e}_k\otimes\mathbf{e}_l)) = L_{ijkl}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot(\mathbf{e}_k\otimes\mathbf{e}_l) = L_{ijkl}\delta_{ik}\delta_{jl} = L_{ijij}$$

and that

$$\mathrm{tr}_4\mathbb{L}^\top = \mathrm{tr}_4(L_{klij}(\mathbf{e}_i\otimes\mathbf{e}_j)\otimes(\mathbf{e}_k\otimes\mathbf{e}_l)) = L_{klij}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot(\mathbf{e}_k\otimes\mathbf{e}_l) = L_{klij}\delta_{ik}\delta_{jl} = L_{ijij} = \mathrm{tr}_4\mathbb{L}.$$

Then, we define the *scalar product of fourth-rank tensors* as

$$\mathbb{A}\cdot\mathbb{B} := \mathrm{tr}_4(\mathbb{A}^\top\mathbb{B}).$$

By the properties of tr_4, the scalar product is a positive definite symmetric bilinear form:

$$\begin{aligned}
&\alpha\mathbb{A}\cdot\beta\mathbb{B} = \mathrm{tr}_4(\alpha\mathbb{A}^\top\beta\mathbb{B}) = \alpha\beta\mathrm{tr}_4(\mathbb{A}^\top\mathbb{B}) = \alpha\beta\mathbb{A}\cdot\mathbb{B},\\
&\mathbb{A}\cdot\mathbb{B} = \mathrm{tr}_4(\mathbb{A}^\top\mathbb{B}) = \mathrm{tr}_4(\mathbb{A}^\top\mathbb{B})^\top = \mathrm{tr}_4(\mathbb{B}^\top\mathbb{A}) = \mathbb{B}\cdot\mathbb{A},\\
&\mathbb{A}\cdot\mathbb{A} = \mathrm{tr}_4(\mathbb{A}^\top\mathbb{A}) = (\mathbb{A}^\top\mathbb{A})_{ijij} = A_{klij}A_{klij} > 0\ \forall\mathbb{A}\in\mathbb{L}in(\mathcal{V}), \mathbb{A}\cdot\mathbb{A} = 0 \iff \mathbb{A} = \mathbb{O}.
\end{aligned}$$

By components

$$\begin{aligned}
\mathbb{A}\cdot\mathbb{B} &= \mathrm{tr}_4((A_{klij}\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l)(B_{pqrs}\mathbf{e}_p\otimes\mathbf{e}_q\otimes\mathbf{e}_r\otimes\mathbf{e}_s))\\
&= \mathrm{tr}_4(A_{ijkl}B_{pqrs}\delta_{kp}\delta_{lq}(\mathbf{e}_i\otimes\mathbf{e}_j)\otimes(\mathbf{e}_r\otimes\mathbf{e}_s))\\
&= A_{ijkl}B_{pqrs}\delta_{kp}\delta_{lq}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot(\mathbf{e}_r\otimes\mathbf{e}_s) = A_{ijkl}B_{pqrs}\delta_{kp}\delta_{lq}\delta_{ir}\delta_{js} = A_{klij}B_{klij}.
\end{aligned}$$

The rule for computing the scalar product is hence always the same already seen for vectors and second-rank tensors: all the indexes are to be saturated.

In complete analogy with vectors and second-rank tensors, we say that $\mathbb{A}$ is *orthogonal to* $\mathbb{B}$ $\iff$

$$\mathbb{A}\cdot\mathbb{B} = 0$$

and we define the *norm of* $\mathbb{L}$ as

$$|\mathbb{L}| := \sqrt{\mathbb{L}\cdot\mathbb{L}} = \sqrt{\mathrm{tr}_4\mathbb{L}^\top\mathbb{L}} = \sqrt{L_{ijkl}L_{ijkl}}.$$

3.5 Projectors, identities

For the spherical part of any $\mathbf{A} \in Sym(\mathcal{V})$ we can write

$$\mathbf{A}^{sph} := \frac{1}{3}\mathrm{tr}\mathbf{A}\mathbf{I} = \frac{1}{3}\mathbf{I}\cdot\mathbf{A}\mathbf{I} = \frac{1}{3}(\mathbf{I}\otimes\mathbf{I})\mathbf{A} = \mathbb{S}^{sph}\mathbf{A},$$

where

$$\mathbb{S}^{sph} := \frac{1}{3}\mathbf{I}\otimes\mathbf{I}$$

is the *spherical projector*, i.e. the fourth-rank tensor that extracts from any $\mathbf{A} \in Lin(\mathcal{V})$ its spherical part. Moreover,

$$\mathbf{A}^{dev} := \mathbf{A} - \mathbf{A}^{sph} = \mathbb{I}\mathbf{A} - \mathbb{S}^{sph}\mathbf{A} = \mathbb{D}^{dev}\mathbf{A},$$

where

$$\mathbb{D}^{dev} := \mathbb{I} - \mathbb{S}^{sph}$$

is the *deviatoric projector*, i.e. the fourth-rank tensor that extracts from any $\mathbf{A} \in Lin(\mathcal{V})$ its deviatoric part. It is worth noting that

$$\mathbb{I} = \mathbb{S}^{sph} + \mathbb{D}^{dev}.$$

Moreover, about the components of $\mathbb{S}^{sph}$,

$$\begin{aligned} S^{sph}_{ijkl} &= (\mathbf{e}_i\otimes\mathbf{e}_j)\cdot\frac{1}{3}(\mathbf{I}\otimes\mathbf{I})(\mathbf{e}_k\otimes\mathbf{e}_l) = \frac{1}{3}(\mathbf{e}_i\otimes\mathbf{e}_j)\cdot\mathbf{I}(\mathbf{e}_k\otimes\mathbf{e}_l)\cdot\mathbf{I} \\ &= \frac{1}{3}\mathrm{tr}(\mathbf{e}_i\otimes\mathbf{e}_j)\mathrm{tr}(\mathbf{e}_k\otimes\mathbf{e}_l) = \frac{1}{3}\delta_{ij}\delta_{kl} \;\rightarrow\; \mathbb{S}^{sph} = \frac{1}{3}\delta_{ij}\delta_{kl}(\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l). \end{aligned}$$

To remark that

$$\mathbb{S}^{sph} = (\mathbb{S}^{sph})^{\top}.$$

We introduce now the tensor $\mathbb{I}^s$, *restriction of* $\mathbb{I}$ *to* $A \in Sym(\mathcal{V})$. It can be introduced as follows: $\forall A \in Sym(\mathcal{V})$

$$\mathbf{A} = \frac{1}{2}(\mathbf{A} + \mathbf{A}^{\top}),$$

and

$$\mathbf{A} = \mathbb{I}\mathbf{A} = \frac{1}{2}(\mathbb{I}\mathbf{A} + \mathbb{I}\mathbf{A}^{\top}) = \frac{1}{2}(I_{ijkl}A_{kl} + I_{ijkl}A_{lk})(\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l);$$

because $\mathbf{A} = \mathbf{A}^{\top}$ there is insensitivity to the swap of indexes k and l, so

$$\mathbf{A} = \frac{1}{2}(I_{ijkl}A_{kl} + I_{ijlk}A_{lk})(\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l) = \frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})A_{kl}(\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l).$$

Then, if we admit the interchangeability of indexes k and l, i.e. if we postulate the existence of the minor right-symmetries for $\mathbb{I}$, then $\mathbb{I} = \mathbb{I}^s$, with

$$\mathbb{I}^s = \frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})(\mathbf{e}_i\otimes\mathbf{e}_j\otimes\mathbf{e}_k\otimes\mathbf{e}_l).$$

It is apparent that

$$I^s_{ijkl} = I^s_{klij},$$

i.e. $\mathbb{I}^s = (\mathbb{I}^s)^\top$, but also that

$$I^s_{ijkl} = \frac{1}{2}(\delta_{il}\delta_{jk} + \delta_{ik}\delta_{jl}) = I^s_{jikl},$$

i.e., $\mathbb{I}^s$ has also the minor left-symmetries; in other words, $\mathbb{I}^s$ has the major and minor symmetries, like an elasticity tensor, while this is not the case for $\mathbb{I}$. In fact

$$I_{ijkl} = I_{jilk} = \delta_{ik}\delta_{jl} \neq \delta_{il}\delta_{jk} = I_{jikl} = I_{ijlk}.$$

Because $\mathbb{S}^{sph}$ and $\mathbb{D}^{dev}$ operate on $Sym(\mathcal{V})$, it is immediate to recognize that it is also

$$\mathbb{D}^{dev} = \mathbb{I}^s - \mathbb{S}^{sph} \;\Rightarrow\; \mathbb{I}^s = \mathbb{S}^{sph} + \mathbb{D}^{dev}.$$

It is worth noting that

$$(\mathbb{D}^{dev})^\top = (\mathbb{I}^s - \mathbb{S}^{spsh})^\top = (\mathbb{I}^s)^\top - (\mathbb{S}^{sph})^\top = \mathbb{I}^s - \mathbb{S}^{sph} = \mathbb{D}^{dev}.$$

We can now determine the components of $\mathbb{D}^{dev}$:

$$D^{dev}_{ijkl} = I^s_{ijkl} - S^{sph}_{ijkl} = \frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk}) - \frac{1}{3}\delta_{ij}\delta_{kl} \;\rightarrow$$
$$\mathbb{D}^{dev} = \left[\frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk}) - \frac{1}{3}\delta_{ij}\delta_{kl}\right](\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l).$$

To remark that the result (2.6) implies that $\mathbb{S}^{sph}$ and $\mathbb{D}^{dev}$ are *orthogonal projectors*, i.e. they project the same $\mathbf{A} \in Sym(\mathcal{V})$ into two orthogonal subspaces of $\mathcal{V}$, $Sph(\mathcal{V})$ and $Dev(\mathcal{V})$.

The tensor $\mathbb{T}^{trp} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ defined by the operation

$$\mathbb{T}^{trp}\mathbf{A} := \mathbf{A}^\top,$$

is the *transposition projector*, whose components are

$$T^{trp}_{ijkl} = (\mathbf{e}_i \otimes \mathbf{e}_j) \cdot \mathbb{T}^{trp}(\mathbf{e}_k \otimes \mathbf{e}_l) = (\mathbf{e}_i \otimes \mathbf{e}_j) \cdot (\mathbf{e}_l \otimes \mathbf{e}_k) = \delta_{il}\delta_{jk}.$$

The following operation defines the *symmetry projector* $\mathbb{S}^{sym} \in \mathbb{L}\mathrm{in}(\mathcal{V})$:

$$\mathbb{S}^{sym}\mathbf{A} = \frac{1}{2}(\mathbf{A} + \mathbf{A}^\top)\ \forall \mathbf{A} \in Lin(\mathcal{V}),$$

while the *antisymmetry projector* $\mathbb{W}^{skw} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ is defined by

$$\mathbb{W}^{skw}\mathbf{A} = \frac{1}{2}(\mathbf{A} - \mathbf{A}^\top)\ \forall \mathbf{A} \in Lin(\mathcal{V}).$$

Also $\mathbb{S}^{sym}$ and $\mathbb{W}^{skw}$ are orthogonal projectors, because they project the same $\mathbf{A} \in Lin(\mathcal{V})$ into two orthogonal subspaces of $Lin(\mathcal{V})$: $Sym(\mathcal{V})$ and $Skw(\mathcal{V})$, see exercice 6 of Chap. 2.

We prove now two properties of the projectors: $\forall \mathbf{A} \in Lin(\mathcal{V})$,

$$(\mathbb{S}^{sym} + \mathbb{W}^{dev})\mathbf{A} = \frac{1}{2}(\mathbf{A} + \mathbf{A}^\top) + \frac{1}{2}(\mathbf{A} - \mathbf{A}^\top) = \mathbf{A} = \mathbb{I}\mathbf{A} \;\Rightarrow\; \mathbb{S}^{sym} + \mathbb{W}^{skw} = \mathbb{I}. \quad (3.5)$$

Then,

$$(\mathbb{S}^{sym} - \mathbb{W}^{dev})\mathbf{A} = \frac{1}{2}(\mathbf{A} + \mathbf{A}^\top) - \frac{1}{2}(\mathbf{A} - \mathbf{A}^\top) = \mathbf{A}^\top = \mathbb{T}^{trp}\mathbf{A} \;\Rightarrow\; \mathbb{S}^{sym} - \mathbb{W}^{skw} = \mathbb{T}^{trp}. \quad (3.6)$$

3.6 Orthogonal conjugator

For any $\mathbf{U} \in Orth(\mathcal{V})$ we define its *orthogonal conjugator* $\mathbb{U} \in \mathbb{L}\mathrm{in}(\mathcal{V})$ as

$$\mathbb{U} := \mathbf{U} \boxtimes \mathbf{U}.$$

Theorem 16. *(orthogonality of $\mathbb{U}$): the orthogonal conjugator is an orthogonal tensor of $\mathbb{L}in(\mathcal{V})$, i.e. it preserves the scalar product between tensors:*

$$\mathbb{U}\mathbf{A} \cdot \mathbb{U}\mathbf{B} = \mathbf{A} \cdot \mathbf{B} \quad \forall \mathbf{A}, \mathbf{B} \in Lin(\mathcal{V}).$$

Proof. By the assertion in exercice 17 of Chap. 2, and because $\mathbf{U} \in Orth(\mathcal{V})$, we have

$$\begin{aligned}\mathbb{U}\mathbf{A} \cdot \mathbb{U}\mathbf{B} &= (\mathbf{U} \boxtimes \mathbf{U})\mathbf{A} \cdot (\mathbf{U} \boxtimes \mathbf{U})\mathbf{B} = \mathbf{U}\mathbf{A}\mathbf{U}^\top \cdot \mathbf{U}\mathbf{B}\mathbf{U}^\top \\ &= \mathbf{U}^\top\mathbf{U}\mathbf{A}\mathbf{U}^\top \cdot \mathbf{B}\mathbf{U}^\top = \mathbf{A}\mathbf{U}^\top \cdot \mathbf{B}\mathbf{U}^\top = \mathbf{A}\mathbf{U}^\top\mathbf{U} \cdot \mathbf{B} = \mathbf{A} \cdot \mathbf{B}.\end{aligned}$$

□

Just as for tensors of $Orth(\mathcal{V})$, it is also

$$\mathbb{U}\mathbb{U}^\top = \mathbb{U}^\top\mathbb{U} = \mathbb{I}.$$

In fact, see the assertion of exercice 3,

$$\mathbb{U}\mathbb{U}^\top = (\mathbf{U} \boxtimes \mathbf{U})(\mathbf{U}^\top \boxtimes \mathbf{U}^\top) = \mathbf{U}\mathbf{U}^\top \boxtimes \mathbf{U}\mathbf{U}^\top = \mathbf{I} \boxtimes \mathbf{I} = \mathbb{I}. \tag{3.7}$$

The orthogonal conjugators have also some properties in relation with projectors:

Theorem 17. *: $\mathbb{S}^{sph}$ is unaffected by any orthogonal conjugator, while $\mathbb{D}^{dev}$ commutes with any orthogonal conjugator.*

Proof. For any $\mathbf{L} \in Sym(\mathcal{V})$ and $\mathbf{U} \in Orth(\mathcal{V})$,

$$\begin{aligned}\mathbb{U}\mathbb{S}^{sph}\mathbf{L} &= (\mathbf{U} \boxtimes \mathbf{U})\left(\frac{1}{3}\mathbf{I} \otimes \mathbf{I}\right)\mathbf{L} = \frac{1}{3}(\mathrm{tr}\mathbf{L})(\mathbf{U} \boxtimes \mathbf{U})\mathbf{I} = \frac{1}{3}(\mathrm{tr}\mathbf{L})\mathbf{U}\mathbf{I}\mathbf{U}^\top \\ &= \frac{1}{3}(\mathrm{tr}\mathbf{L})\mathbf{I} = \frac{1}{3}\mathbf{I} \cdot \mathbf{L}\mathbf{I} = \frac{1}{3}(\mathbf{I} \otimes \mathbf{I})\mathbf{L} = \mathbb{S}^{sph}\mathbf{L}.\end{aligned}$$

Moreover,

$$\begin{aligned}\mathbb{S}^{sph}\mathbb{U}\mathbf{L} &= \left(\frac{1}{3}\mathbf{I} \otimes \mathbf{I}\right)(\mathbf{U} \boxtimes \mathbf{U})\mathbf{L} = \frac{1}{3}(\mathbf{I} \otimes \mathbf{I})(\mathbf{U}\mathbf{L}\mathbf{U}^\top) = \frac{1}{3}(\mathbf{I} \cdot \mathbf{U}\mathbf{L}\mathbf{U}^\top)\mathbf{I} \\ &= \frac{1}{3}\mathrm{tr}(\mathbf{U}\mathbf{L}\mathbf{U}^\top)\mathbf{I} = \frac{1}{3}\mathrm{tr}(\mathbf{U}^\top\mathbf{U}\mathbf{L})\mathbf{I} = \frac{1}{3}(\mathrm{tr}\mathbf{L})\mathbf{I} = \frac{1}{3}\mathbf{I} \cdot \mathbf{L}\mathbf{I} = \frac{1}{3}(\mathbf{I} \otimes \mathbf{I})\mathbf{L} = \mathbb{S}^{sph}\mathbf{L}.\end{aligned}$$

Thus, we have proved that

$$\mathbb{S}^{sph}\mathbb{U} = \mathbb{U}\mathbb{S}^{sph} = \mathbb{S}^{sph},$$

i.e. that the spherical projector $\mathbb{S}^{sph}$ is unaffected by any orthogonal conjugator. Further

$$\mathbb{D}^{dev}\mathbb{U}\mathbf{L} = (\mathbb{I}^s - \mathbb{S}^{sph})\mathbb{U}\mathbf{L} = \mathbb{I}^s\mathbb{U}\mathbf{L} - \mathbb{S}^{sph}\mathbb{U}\mathbf{L} = \mathbb{U}\mathbf{L} - \mathbb{S}^{sph}\mathbf{L} = (\mathbb{U} - \mathbb{S}^{sph})\mathbf{L}$$

and

$$\mathbb{U}\mathbb{D}^{dev}\mathbf{L} = \mathbb{U}(\mathbb{I}^s - \mathbb{S}^{sph})\mathbf{L} = \mathbb{U}\mathbb{I}^s\mathbf{L} - \mathbb{U}\mathbb{S}^{sph}\mathbf{L} = \mathbb{U}\mathbf{L} - \mathbb{S}^{sph}\mathbf{L} = (\mathbb{U} - \mathbb{S}^{sph})\mathbf{L},$$

so that

$$\mathbb{D}^{dev}\mathbb{U} = \mathbb{U}\mathbb{D}^{dev}.$$

□

3.7 Rotations and symmetries

We ponder now how to rotate a fourth-rank tensor, i.e., what are the components of

$$\mathbb{L} = L_{ijkl}\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l$$

in a basis $\mathcal{B}' = \{\mathbf{e}'_1, \mathbf{e}'_2, \mathbf{e}'_3\}$ obtained rotating the basis $\mathcal{B} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ by the rotation $\mathbb{R} = R_{ij}\mathbf{e}_i \otimes \mathbf{e}_j, \mathbb{R} \in Orth(\mathcal{V})^+$. The procedure is exactly the same already seen for vectors and second-rank tensors:

$$\begin{aligned}\mathbb{L} &= L_{ijkl}\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l = L_{ijkl}R^\top_{pi}\mathbf{e}'_p \otimes R^\top_{qj}\mathbf{e}'_q \otimes R^\top_{rk}\mathbf{e}'_r \otimes R^\top_{sl}\mathbf{e}'_s \\ &= R^\top_{pi}R^\top_{qj}R^\top_{rk}R^\top_{sl}L_{ijkl}\mathbf{e}'_p \otimes \mathbf{e}'_q \otimes \mathbf{e}'_r \otimes \mathbf{e}'_s,\end{aligned}$$

i.e.

$$L'_{pqrs} = R^\top_{pi}R^\top_{qj}R^\top_{rk}R^\top_{sl}L_{ijkl}.$$

We see clearly that the components of $\mathbb{L}$ in the basis $\mathcal{B}'$ are a linear combination of those in $\mathcal{B}$, the coefficients of the linear combination being fourth-powers of the director cosines, the R_{ij}s. The introduction of the orthogonal conjugator[1] of the rotation $\mathbf{R}$,

$$\mathbb{R} = \mathbf{R} \boxtimes \mathbf{R},$$

allows to give a compact expression for the rotation of second- and fourth-rank tensors (for completeness we recall also that of a vector $\mathbf{w}$);

$$\begin{aligned}\mathbf{w}' &= \mathbf{R}^\top\mathbf{w}, \\ \mathbf{L}' &= \mathbf{R}^\top\mathbf{L}\mathbf{R} = (\mathbf{R}^\top \boxtimes \mathbf{R}^\top)\mathbf{L} = \mathbb{R}^\top\mathbf{L}, \\ \mathbb{L}' &= (\mathbf{R}^\top \boxtimes \mathbf{R}^\top)\mathbb{L}(\mathbf{R} \boxtimes \mathbf{R}) = \mathbb{R}^\top\mathbb{L}\mathbb{R}.\end{aligned}$$

The check of the above relations with the orthogonal conjugator $\mathbb{R}$ is left to the reader. It is worth noting that actually these transformations are valid not only for $\mathbf{R} \in Orth(\mathcal{V})^+$, but more generally for any $\mathbf{U} \in Orth(\mathcal{V})$, i.e. also for symmetries.

If by $\mathbf{U}$ we denote the tensor of change of basis under any orthogonal transformation, i.e., if we put $\mathbf{U} = \mathbf{R}^\top$ for the rotations, then the above relations become

$$\begin{aligned}\mathbf{w}' &= \mathbf{U}\mathbf{w}, \\ \mathbf{L}' &= \mathbf{U}\mathbf{L}\mathbf{U}^\top = (\mathbf{U} \boxtimes \mathbf{U})\mathbf{L} = \mathbb{U}\mathbf{L}, \\ \mathbb{L}' &= (\mathbf{U} \boxtimes \mathbf{U})\mathbb{L}(\mathbf{U} \boxtimes \mathbf{U})^\top = \mathbb{U}\mathbb{L}\mathbb{U}^\top.\end{aligned} \tag{3.8}$$

[1]Here the symbol $\mathbb{R}$ indicates the orthogonal conjugator of $\mathbf{R}$, not the set of real numbers.

Finally, we say that $\mathbf{L} \in Lin(\mathcal{V})$ or $\mathbb{L} \in \mathbb{L}\text{in}(\mathcal{V})$ is *invariant under an orthogonal transformation* $\mathbf{U}$ if

$$\mathbf{U}\mathbf{L}\mathbf{U}^\top = \mathbf{L}, \quad \mathbb{U}\mathbb{L}\mathbb{U}^\top = \mathbb{L};$$

right multiplying both terms by $\mathbf{U}$ or by $\mathbb{U}$ and through eq. (3.7), we get that $\mathbf{L}$ or $\mathbb{L}$ are invariant under $\mathbf{U} \iff$

$$\mathbf{U}\mathbf{L} = \mathbf{L}\mathbf{U}, \quad \mathbb{U}\mathbb{L} = \mathbb{L}\mathbb{U},$$

i.e. $\iff$ $\mathbf{L}$ and $\mathbf{U}$, or $\mathbb{L}$ and $\mathbb{U}$ commute. This relation allows, e.g., the analysis of material symmetries in elasticity.

If a tensor is invariant under *any* orthogonal transformation, i.e. if the previous equations hold true $\forall \mathbf{U} \in Orth(\mathcal{V})$ then the tensor is said to be *isotropic*. A general result[2] is that a fourth-rank tensor $\mathbb{L}$ is isotropic $\iff$ exist two scalar functions λ, μ such that

$$\mathbb{L}\mathbf{A} = 2\mu\mathbf{A} + \lambda \text{tr}\mathbf{A}\ \mathbf{I} \quad \forall \mathbf{A} \in Sym(\mathcal{V}).$$

The reader is addressed to the book of Gurtin (see references) for the proof of this result and for a deeper insight in isotropic functions.

3.8 The Kelvin formalism

As already mentioned, though fourth-rank tensors cannot be organized in and represented by a matrix, a matrix formalism for these operators exists. Such formalism is due to Kelvin[3] and it is strictly related to the theory of elasticity, i.e. it concerns the Cauchy's stress tensor $\boldsymbol{\sigma}$, the strain tensor $\boldsymbol{\varepsilon}$ and the elasticity tensor $\mathbb{E}$. The relation between $\boldsymbol{\sigma}$ and $\boldsymbol{\varepsilon}$ is given by the celebrated (generalized) *Hooke's law*:

$$\boldsymbol{\sigma} = \mathbb{E}\boldsymbol{\varepsilon}.$$

Both $\boldsymbol{\sigma}, \boldsymbol{\varepsilon} \in Sym(\mathcal{V})$ while $\mathbb{E} = \mathbb{E}^\top$ and it has also the minor symmetries, so $\mathbb{E}$ has just 21 independent components[4]. In the Kelvin formalism, the six independent components of $\boldsymbol{\sigma}$ and $\boldsymbol{\varepsilon}$ are organized into column vectors and renumbered as follows

$$\{\sigma\} = \begin{Bmatrix} \sigma_1 = \sigma_{11} \\ \sigma_2 = \sigma_{22} \\ \sigma_3 = \sigma_{33} \\ \sigma_4 = \sqrt{2}\sigma_{23} \\ \sigma_5 = \sqrt{2}\sigma_{31} \\ \sigma_6 = \sqrt{2}\sigma_{12} \end{Bmatrix}, \quad \{\varepsilon\} = \begin{Bmatrix} \varepsilon_1 = \varepsilon_{11} \\ \varepsilon_2 = \varepsilon_{22} \\ \varepsilon_3 = \varepsilon_{33} \\ \varepsilon_4 = \sqrt{2}\varepsilon_{23} \\ \varepsilon_5 = \sqrt{2}\varepsilon_{31} \\ \varepsilon_6 = \sqrt{2}\varepsilon_{12} \end{Bmatrix}.$$

[2] Actually, this is a quite famous result in classical elasticity, the Lamé's equation, defining an isotropic elastic material.

[3] W. Thomson (Lord Kelvin): *Elements of a mathematical theory of elasticity.* Philos. Trans. R. Soc., 146, 481-498, 1856. Later, Voigt (W. Voigt: *Lehrbuch der Kristallphysik.* B. G. Taubner, Leipzig, 1910) gave another, similar matrix formalism for tensors, more widely known than the Kelvin one, but less effective.

[4] Actually, the Kelvin formalism can be extended without major difficulties also to tensors that do not possess all the symmetries.

The elasticity tensor $\mathbb{E}$ is reduced to a 6×6 matrix $[E]$, consequence of the minor symmetries induced by the symmetry of $\boldsymbol{\sigma}$ and $\boldsymbol{\varepsilon}$; this matrix is symmetric because $\mathbb{E} = \mathbb{E}^\top$:

$$[E] = \begin{bmatrix} E_{11} = E_{1111} & E_{12} = E_{1122} & E_{13} = E_{1133} & E_{14} = \sqrt{2}E_{1123} & E_{15} = \sqrt{2}E_{1131} & E_{16} = \sqrt{2}E_{1112} \\ E_{12} = E_{1122} & E_{22} = E_{2222} & E_{23} = E_{2233} & E_{24} = \sqrt{2}E_{2223} & E_{25} = \sqrt{2}E_{2231} & E_{26} = \sqrt{2}E_{2212} \\ E_{13} = E_{1133} & E_{23} = E_{2233} & E_{33} = E_{3333} & E_{34} = \sqrt{2}E_{3323} & E_{35} = \sqrt{2}E_{3331} & E_{36} = \sqrt{2}E_{3312} \\ E_{14} = \sqrt{2}E_{1123} & E_{24} = \sqrt{2}E_{2223} & E_{34} = \sqrt{2}E_{3323} & E_{44} = 2E_{2323} & E_{45} = 2E_{2331} & E_{46} = 2E_{2312} \\ E_{15} = \sqrt{2}E_{1131} & E_{25} = \sqrt{2}E_{2231} & E_{35} = \sqrt{2}E_{3331} & E_{45} = 2E_{2331} & E_{55} = 2E_{3131} & E_{56} = 2E_{3112} \\ E_{16} = \sqrt{2}E_{1112} & E_{26} = \sqrt{2}E_{2212} & E_{36} = \sqrt{2}E_{3312} & E_{46} = 2E_{2312} & E_{56} = 2E_{3112} & E_{66} = 2E_{1212} \end{bmatrix}.$$

In this way, the matrix product

$$\{\sigma\} = [E]\{\varepsilon\} \tag{3.9}$$

is equivalent to the tensor form of the Hooke's law and all the operations can be done by the aid of classical matrix algebra[5], e.g. the computation of the inverse of $\mathbb{E}$, the *compliance tensor*.

An important operation is the expression of tensor $\mathbb{U}$ in eq. (3.8) in the Kelvin formalism; some tedious but straightforward passages give the result:

$$[U] = \begin{bmatrix} U_{11}^2 & U_{12}^2 & U_{13}^2 & \sqrt{2}U_{12}U_{13} & \sqrt{2}U_{13}U_{11} & \sqrt{2}U_{11}U_{12} \\ U_{21}^2 & U_{22}^2 & U_{23}^2 & \sqrt{2}U_{22}U_{23} & \sqrt{2}U_{23}U_{21} & \sqrt{2}U_{21}U_{22} \\ U_{31}^2 & U_{32}^2 & U_{33}^2 & \sqrt{2}U_{32}U_{33} & \sqrt{2}U_{33}U_{31} & \sqrt{2}U_{31}U_{32} \\ \sqrt{2}U_{21}U_{31} & \sqrt{2}U_{22}U_{32} & \sqrt{2}U_{23}U_{33} & U_{23}U_{32} + U_{22}U_{33} & U_{33}U_{21} + U_{31}U_{23} & U_{31}U_{22} + U_{32}U_{21} \\ \sqrt{2}U_{31}U_{11} & \sqrt{2}U_{32}U_{12} & \sqrt{2}U_{33}U_{13} & U_{32}U_{13} + U_{33}U_{12} & U_{31}U_{13} + U_{33}U_{11} & U_{31}U_{12} + U_{32}U_{11} \\ \sqrt{2}U_{11}U_{21} & \sqrt{2}U_{12}U_{22} & \sqrt{2}U_{13}U_{23} & U_{12}U_{23} + U_{13}U_{22} & U_{11}U_{23} + U_{13}U_{21} & U_{11}U_{22} + U_{12}U_{21} \end{bmatrix}$$

With some work, it can be checked that

$$[U][U]^\top = [U]^\top[U] = [I],$$

i.e. that $[U]$ is an orthogonal matrix in $\mathbb{R}^6$. Of course,

$$[R] = [U]^\top$$

is the matrix that in the Kelvin formalism represents the tensor $\mathbf{R} = \mathbf{U}^\top$. The change of basis for $\boldsymbol{\sigma}$ and $\boldsymbol{\varepsilon}$ are hence done through the relations

$$\{\sigma'\} = [U]\{\sigma\}, \quad \{\varepsilon'\} = [U]\{\varepsilon\},$$

which applied to eq. (3.9) give

$$\{\sigma\} = [E]\{\varepsilon\} \;\rightarrow\; [U]^\top\{\sigma'\} = [E][U]^\top\{\varepsilon'\} \;\rightarrow\; \{\sigma'\} = [U][E][U]^\top\{\varepsilon'\}$$

i.e. in the basis $\mathcal{B}'$

$$\{\sigma'\} = [E']\{\varepsilon'\},$$

where

$$[E'] = [U][E][U]^\top = [R]^\top[E][R]$$

is the matrix representing $\mathbb{E}$ in $\mathcal{B}'$ in the Kelvin formalism. Though it is possible to give the expression of the components of $[E']$, they are so long that they are omitted here.

[5]Mehrabadi and Cowin have shown that the Kelvin formalism transforms second- and fourth-rank tensors on $\mathbb{R}^3$ into vectors and second-rank tensors on $\mathbb{R}^6$ (M. M. Mehrabadi, S. C. Cowin: *Eigentensors of linear anisotropic elastic materials.* Q. J. Mech. Appl. Math., 43, 15-41, 1990).

3.9 Exercices

1. Prove eq. (3.4).
2. Prove that
$$\mathbf{A} \otimes \mathbf{B}\mathbb{L} = \mathbf{A} \otimes \mathbb{L}^{\top}\mathbf{B}.$$
3. Prove that
$$(\mathbf{A} \boxtimes \mathbf{B})(\mathbf{C} \boxtimes \mathbf{D}) = \mathbf{AC} \boxtimes \mathbf{BD}.$$
4. Prove eq. (3.3) using the result of the previous exercice.
5. Prove that
$$(\mathbf{A} \otimes \mathbf{B})(\mathbf{C} \boxtimes \mathbf{D}) = \mathbf{A} \otimes ((\mathbf{C}^{\top} \boxtimes \mathbf{D}^{\top})\mathbf{B}).$$
6. Prove that
$$(\mathbf{A} \boxtimes \mathbf{B})(\mathbf{C} \otimes \mathbf{D}) = ((\mathbf{A} \boxtimes \mathbf{B})\mathbf{C}) \otimes \mathbf{D}.$$
7. Let $\mathbf{p} \in \mathcal{S}$ and $\mathbf{P} = \mathbf{p} \otimes \mathbf{p}$; then prove that
$$\mathbf{P} \boxtimes \mathbf{P} = \mathbf{P} \otimes \mathbf{P}.$$
8. Prove that, $\forall \mathbb{A} \in \mathbb{L}\mathrm{in}(\mathcal{V})$,
$$\mathbb{I}\mathbb{A} = \mathbb{A}\mathbb{I} = \mathbb{A}.$$
9. Show that
$$(\mathbf{A} \otimes \mathbf{B}) \cdot (\mathbf{C} \otimes \mathbf{D}) = \mathbf{A} \cdot \mathbf{C}\ \mathbf{B} \cdot \mathbf{D}.$$
10. Show that
$$\mathbb{S}^{sph} = \frac{\mathbf{I}}{|\mathbf{I}|} \otimes \frac{\mathbf{I}}{|\mathbf{I}|}.$$
11. Show that
$$dim(Sph(\mathcal{V})) = 1, \quad dim(Dev(\mathcal{V})) = 5.$$
12. Show the following properties of $\mathbb{S}^{sph}$ and $\mathbb{D}^{dev}$:
$$\mathbb{S}^{sph}\mathbb{S}^{sph} = \mathbb{S}^{sph},$$
$$\mathbb{D}^{dev}\mathbb{D}^{dev} = \mathbb{D}^{dev},$$
$$\mathbb{S}^{sph}\mathbb{D}^{dev} = \mathbb{D}^{dev}\mathbb{S}^{sph} = \mathbb{O}.$$
13. Prove the results in eqs. (3.5) and (3.6) using the components.
14. Show that
$$\mathbb{S}^{sph} \cdot \mathbb{S}^{sph} = 1,$$
$$\mathbb{D}^{dev} \cdot \mathbb{D}^{dev} = 5,$$
$$\mathbb{S}^{sph} \cdot \mathbb{D}^{dev} = 0.$$
15. Explicit the orthogonal conjugator $\mathbb{S}_R$ of the tensor $\mathbf{S}_R$ in eq. (2.43).

Chapter 4

Tensor analysis, curves, fields

4.1 Introduction

We generalize to the spaces $\mathcal{E}, \mathcal{V}, Lin(\mathcal{V})$ and $\mathbb{L}\mathrm{in}(\mathcal{V})$ some concepts already known for functions in $\mathbb{R}$, like continuity, differentiation and integration. Curves, fields, deformations are also introduced along with the differential operators used in continuum mechanics, also with their expression in cylindrical and spherical coordinates. Finally, we will introduce some concepts of differential geometry of curves with the aid of tensor algebra.

4.2 Curves of points, vectors and tensors

The scalar products in $\mathcal{V}, Lin(\mathcal{V})$ and $\mathbb{L}\mathrm{in}(\mathcal{V})$ allow us to define a *norm*, the *Euclidean norm*, so they automatically endow these spaces with a *metric*, i.e. we are able to measure and calculate a distance between two elements of such a space and in $\mathcal{E}$. This allows us to generalize the concepts of continuity and differentiability already known in $\mathbb{R}$, whose definition intrinsically makes use of a distance between real quantities.

Let $\pi_n = \{p_n \in \mathcal{E}, n \in \mathbb{N}\}$ a sequence of points in $\mathcal{E}$. We say that π_n *converges to* $p \in \mathcal{E}$ if

$$\lim_{n\to\infty} d(p_n - p) = 0.$$

A similar definition can be given for sequences of vectors or tensors of any rank. Through this definition of convergence we can now precise the concepts of continuity and of curve.

Let $[a, b]$ an interval of $\mathbb{R}$; the function

$$p = p(t) : [a, b] \to \mathcal{E}$$

is *continuous* at $t \in [a, b]$ if for each sequence $\{t_n \in [a, b], n \in \mathbb{N}\}$ that converges to t the sequence π_n defined by $p_n = p(t_n)\ \forall n \in \mathbb{N}$ converges to $p(t) \in \mathcal{E}$. The function $p = p(t)$ is a *curve in* $\mathcal{E}$ $\iff$ it is continuous $\forall t \in [a, b]$. In the same way we can define a curve

of vectors and of tensors:

$$\mathbf{v} = \mathbf{v}(t) : [a, b] \to \mathcal{V},$$
$$\mathbf{L} = \mathbf{L}(t) : [a, b] \to Lin(\mathcal{V}),$$
$$\mathbb{L} = \mathbb{L}(t) : [a, b] \to \mathbb{L}\mathrm{in}(\mathcal{V}).$$

Mathematically, a curve is a function that lets correspond to a real value t (the *parameter*) in an interval an element of a space, $\mathcal{E}, \mathcal{V}, Lin(\mathcal{V})$ or $\mathbb{L}(\mathcal{V})$.

4.3 Differention of curves

Let $\mathbf{v} = \mathbf{v}(t) : [a, b] \to \mathcal{V}$ a curve of vectors and $g = g(t) : [a, b] \to \mathbb{R}$ a scalar function. We say that $\mathbf{v}$ is *of the order o with respect to g in* $t_0 \iff$

$$\lim_{t \to t_0} \frac{|\mathbf{v}(t)|}{|g(t)|} = 0,$$

and we write

$$\mathbf{v}(t) = o(g(t)) \text{ for } t \to t_0.$$

A similar definition can be given for a curve of tensors of any rank. We then say that the curve $\mathbf{v}$ is *differentiable* in $t_0 \in]a, b[\iff \exists \mathbf{v}' \in \mathcal{V}$ such that

$$\mathbf{v}(t) - \mathbf{v}(t_0) = (t - t_0)\mathbf{v}' + o(t - t_0).$$

We call $\mathbf{v}'$ the *derivative of* $\mathbf{v}$[1], also indicated by $\dfrac{d\mathbf{v}}{dt}$. Applying the definition of derivative to $\mathbf{v}'$ we define the *second derivative* $\mathbf{v}''$ of $\mathbf{v}$ and recursively all the derivatives of higher orders. We say that $\mathbf{v}$ is of *class* C^n if it is continuous with its derivatives up to the order n; if $n \geq 1$, $\mathbf{v}$ is said to be *smooth*. Similar definitions can be given for curves in $\mathcal{E}, Lin(\mathcal{V})$ and $\mathbb{L}\mathrm{in}(\mathcal{V})$, so defining derivatives of points and tensors. We remark that the derivative of a curve in $\mathcal{E}$, defined as a difference of points, is a curve in $\mathcal{V}$ (we say, in short, that *the derivative of a point is a vector*). For what concerns tensors, the derivative of a tensor of rank r is a tensor of the same rank.

Let $\mathbf{u}, \mathbf{v}$ curves in $\mathcal{V}$, $\mathbf{L}, \mathbf{M}$ curves in $Lin(\mathcal{V})$, $\mathbb{L}, \mathbb{M}$ curves in $\mathbb{L}\mathrm{in}(\mathcal{V})$ and α a scalar function, all of them defined and at least of class C^1 on $[a, b]$. The same definition of derivative of a curve gives the following results, whose proof is let to the reader:

[1]The symbol $\dot{}$ is also used, but it is usually reserved, in physics, to the case where t is the time.

$$
\begin{aligned}
(\mathbf{u}+\mathbf{v})' &= \mathbf{u}'+\mathbf{v}',\\
(\alpha\mathbf{v})' &= \alpha'\mathbf{v}+\alpha\mathbf{v}',\\
(\mathbf{u}\cdot\mathbf{v})' &= \mathbf{u}'\cdot\mathbf{v}+\mathbf{u}\cdot\mathbf{v}',\\
(\mathbf{u}\times\mathbf{v})' &= \mathbf{u}'\times\mathbf{v}+\mathbf{u}\times\mathbf{v}',\\
(\mathbf{u}\otimes\mathbf{v})' &= \mathbf{u}'\otimes\mathbf{v}+\mathbf{u}\otimes\mathbf{v}',\\
(\mathbf{L}+\mathbf{M})' &= \mathbf{L}'+\mathbf{M}',\\
(\alpha\mathbf{L})' &= \alpha'\mathbf{L}+\alpha\mathbf{L}',\\
(\mathbf{L}\mathbf{v})' &= \mathbf{L}'\mathbf{v}+\mathbf{L}\mathbf{v}',\\
(\mathbf{L}\mathbf{M})' &= \mathbf{L}'\mathbf{M}+\mathbf{L}\mathbf{M}',\\
(\mathbf{L}\cdot\mathbf{M})' &= \mathbf{L}'\cdot\mathbf{M}+\mathbf{L}\cdot\mathbf{M}',\\
(\mathbf{L}\otimes\mathbf{M})' &= \mathbf{L}'\otimes\mathbf{M}+\mathbf{L}\otimes\mathbf{M}',\\
(\mathbf{L}\boxtimes\mathbf{M})' &= \mathbf{L}'\boxtimes\mathbf{M}+\mathbf{L}\boxtimes\mathbf{M}',\\
(\mathbb{L}+\mathbb{M})' &= \mathbb{L}'+\mathbb{M}',\\
(\alpha\mathbb{L})' &= \alpha'\mathbb{L}+\alpha\mathbb{L}',\\
(\mathbb{L}\mathbf{L})' &= \mathbb{L}'\mathbf{L}+\mathbb{L}\mathbf{L}',\\
(\mathbb{L}\mathbb{M})' &= \mathbb{L}'\mathbb{M}+\mathbb{L}\mathbb{M}',\\
(\mathbb{L}\cdot\mathbb{M})' &= \mathbb{L}'\cdot\mathbb{M}+\mathbb{L}\cdot\mathbb{M}'.
\end{aligned}
\tag{4.1}
$$

To remark that the derivative of any kind of product is made according to the usual rule of the derivative of a product of functions.

Be $\mathcal{R}=\{o;\mathcal{B}\}$ a reference frame of the euclidean space $\mathcal{E}$, composed by an *origin* o and a basis $\mathcal{B}=\{\mathbf{e}_1,\mathbf{e}_2,\mathbf{e}_3\}$ of $\mathcal{V}, \mathbf{e}_i\cdot\mathbf{e}_j=\delta_{ij}\forall i,j=1,2,3$ and let us consider a point $p(t)=(p_1(t),p_2(t),p_3(t))$. If the three *coordinates* $p_i(t)$ are three continuous functions over the interval $[t_1,t_2]\in\mathbb{R}$, then, by the definition given above, the mapping $p(t):[t_1,t_2]\to\mathcal{E}$ is a curve in $\mathcal{E}$ and the equation

$$
p(t)=(p_1(t),p_2(t),p_3(t)) \;\to\; \begin{cases} p_1=p_1(t)\\ p_2=p_2(t)\\ p_3=p_3(t)\end{cases}
$$

is the *parametric point equation of the curve*: to each value of $t\in[t_1,t_2]$ it corresponds a point of the curve in $\mathcal{E}$, see Fig. 4.1.

The vector function $\mathbf{r}(t)=p(t)-o$ is the *position vector* of point p in $\mathcal{R}$; the equation

$$
\mathbf{r}(t)=r_i(t)\mathbf{e}_i=r_1(t)\mathbf{e}_1+r_2(t)\mathbf{e}_2+r_3(t)\mathbf{e}_3 \;\to\; \begin{cases} r_1=r_1(t)\\ r_2=r_2(t)\\ r_3=r_3(t)\end{cases}
$$

is the *parametric vector equation* of the curve: to each value of $t\in[t_1,t_2]$ it corresponds a vector of $\mathcal{V}$ that determines a point of the curve in $\mathcal{E}$ through the operation $p(t)=o+\mathbf{r}(t)$.

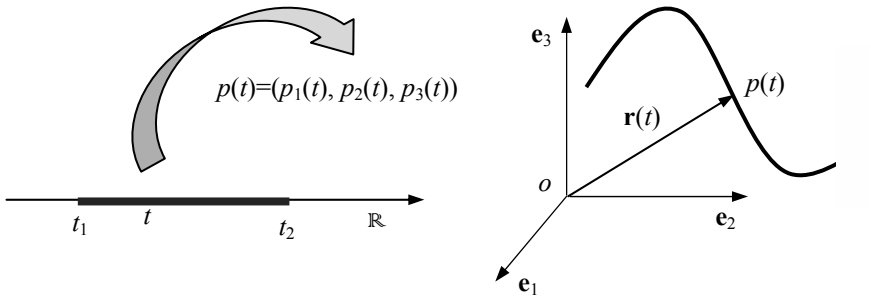

Figure 4.1: Mapping of a curve of points.

Similarly, if the components $L_{ij}(t)$ are continuous functions of a parameter t, the mapping $\mathbf{L}(t) : [t_1, t_2] \to Lin(\mathcal{V})$ defined by

$$\mathbf{L}(t) = L_{ij}(t)\mathbf{e}_i \otimes \mathbf{e}_j, \quad i, j = 1, 2, 3,$$

is a curve of tensors. In the same way we can give a curve of fourth-rank tensors $\mathbb{L}(t) : [t_1, t_2] \to \mathbb{L}\text{in}(\mathcal{V})$ by

$$\mathbb{L}(t) = L_{ijkl}(t)\mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l, \quad i, j, k, l = 1, 2, 3.$$

To be noticed that the choice of the parameter is *not unique*: the equation $p = p[\tau(t)]$ still represents the same curve $p = p(t)$, through the *change of parameter* $\tau = \tau(t)$.

The definition given above for the derivative of a curve of points $p = p(t)$ in $t = t_0$ is equivalent to the following one (probably more familiar to the reader)

$$\frac{dp(t)}{dt} = \lim_{\varepsilon \to 0} \frac{p(t_0 + \varepsilon) - p(t_0)}{\varepsilon};$$

represented in Fig. 4.2, it is apparent that $\mathbf{r}'(t) = \dfrac{dp(t)}{dt}$ is a vector.

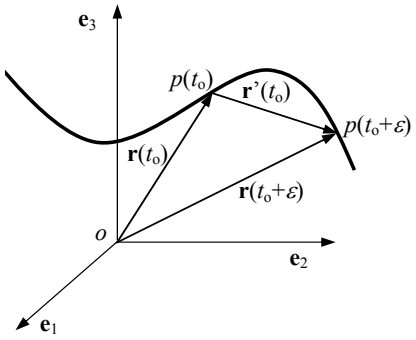

Figure 4.2: Derivative of a curve.

An important case is that of a vector $\mathbf{v}(t)$ whose norm $v(t)$ is constant $\forall t$:

$$(v^2)' = (\mathbf{v} \cdot \mathbf{v})' = \mathbf{v}' \cdot \mathbf{v} + \mathbf{v} \cdot \mathbf{v}' = 2\mathbf{v}' \cdot \mathbf{v} = 0 : \tag{4.2}$$

the derivative of such a vector is orthogonal to it $\forall t$. The contrary is also true, as immediately apparent.

Finally, using the above rules and assuming that the reference frame $\mathcal{R}$ is independent from t, we get easily that

$$\begin{aligned}
p'(t) &= p_i'(t)\ \mathbf{e}_i,\\
\mathbf{v}'(t) &= v_i'(t)\ \mathbf{e}_i,\\
\mathbf{L}'(t) &= L_{ij}'(t)\ \mathbf{e}_i \otimes \mathbf{e}_j,\\
\mathbb{L}'(t) &= L_{ijkl}'(t)\ \mathbf{e}_i \otimes \mathbf{e}_j \otimes \mathbf{e}_k \otimes \mathbf{e}_l,
\end{aligned}$$

i.e. that the derivative of a curve of points, vectors or tensors is simply calculated differentiating the coordinates or the components.

More involved is to prove that

$$\begin{aligned}
(\mathbf{L}^\top)' &= \mathbf{L}'^\top,\\
(\mathbb{L}^\top)' &= \mathbb{L}'^\top,\\
(\det \mathbf{L})' = \det \mathbf{L}\ \mathrm{tr}(\mathbf{L}'\mathbf{L}^{-1}) &= \det \mathbf{L}\ \mathbf{L}^{\top\prime} \cdot \mathbf{L}^{-1} = \det \mathbf{L}\ \mathbf{L}' \cdot \mathbf{L}^{-\top},
\end{aligned}$$

the reader is addressed to the book of Gurtin for the proof.

4.4 Integral of a curve of vectors

We define *integral of a curve of vectors* $\mathbf{r}(t)$ between a and $b \in [t_1, t_2]$ the curve that is obtained integrating each component of the curve:

$$\int_a^b \mathbf{r}(t)\ dt = \int_a^b r_i(t)\ dt\ \mathbf{e}_i.$$

If the curve is regular, we can generalize the second fundamental theorem of the integral calculus

$$\mathbf{r}(t) = \mathbf{r}(a) + \int_a^t \mathbf{r}'(t^*)\ dt^*.$$

Because

$$\mathbf{r}(t) = p(t) - o, \quad \mathbf{r}'(t) = (p(t) - o)' = p'(t),$$

we get also

$$p(t) = p(a) + \int_a^t p'(t^*)\ dt^*.$$

The integral of a vector function is the generalization of the vector sum, see Fig. 4.3.

A simple way to determine a point $p(t)$ on a curve is to fix a point p_0 on the curve and to measure the length $s(t)$ of the arc of curve between $p_0 = p(t = 0)$ and $p(t)$. This length $s(t)$ is called a *curvilinear abscissa* and it can be shown that

$$s(t) = \int_{t_0}^t |(p(t^*) - o)'| dt^* = \int_{t_0}^t |\mathbf{r}'(t^*)| dt^*, \tag{4.3}$$

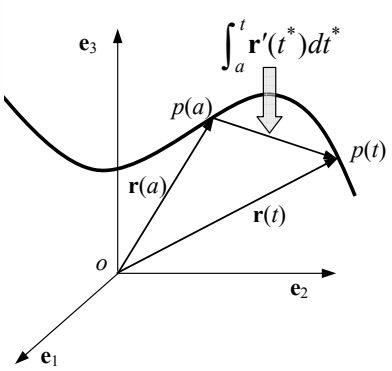

Figure 4.3: Integral of a vector curve.

so that the total length ℓ of a curve is

$$\ell = \int_a^b |\mathbf{r}'(t)|dt.$$

From eq. (4.3) we get

$$\frac{ds}{dt} = |\mathbf{r}'(t)| = \sqrt{\left(\frac{dr_1}{dt}\right)^2 + \left(\frac{dr_2}{dt}\right)^2 + \left(\frac{dr_3}{dt}\right)^2} > 0$$

so that $s(t)$ is an increasing function of t and the length of an infinitesimal arc is

$$ds = \sqrt{dr_1^2 + dr_2^2 + dr_3^2} = \sqrt{dx^2 + dy^2 + dz^2}.$$

For a plane curve $y = f(x)$, we can always put $t = x$, which gives the parametric equation

$$p(t) = (t, f(t)),$$

or in vector form

$$\mathbf{r}(t) = t\ \mathbf{e}_1 + f(t)\ \mathbf{e}_2,$$

from which we obtain

$$\frac{ds}{dt} = |\mathbf{r}'(t)| = |p'(t)| = \sqrt{1 + f'^2(t)}, \tag{4.4}$$

that gives the length of a plane curve between $t = x_0$ and $t = x$ as a function of the abscissa x:

$$s(x) = \int_{x_0}^{x} \sqrt{1 + f'^2(t)}dt.$$

4.5 The Frenet-Serret basis

We define the *tangent vector* $\boldsymbol{\tau}(t)$ to a regular curve $p = p(t)$ the vector

$$\boldsymbol{\tau}(t) = \frac{p'(t)}{|p'(t)|}.$$

By the definition of derivative, this unit vector is always oriented as the increasing values of t; the straight line tangent to the curve in $p_0 = p(t_0)$ has hence equation

$$q(\bar{t}) = p(t_0) + \bar{t}\ \boldsymbol{\tau}(t_0).$$

If the curvilinear abscissa s is chosen as parameter for the curve, through the change of parameter $s = s(t)$ we get

$$\boldsymbol{\tau}(t) = \frac{p'(t)}{|p'(t)|} = \frac{p'[s(t)]}{|p'[s(t)]|} = \frac{1}{s'(t)}\frac{dp(s)}{ds}\frac{ds(t)}{dt} = \frac{dp(s)}{ds} \ \rightarrow \boldsymbol{\tau}(s) = p'(s).$$

So, if the parameter of the curve is s, the derivative of the curve is $\boldsymbol{\tau}$, i.e. it is automatically a unit vector. The above equation, in addition, shows that the change of parameter does not change the direction of the tangent, because just a scalar, the derivative of the parameter's change, multiplies the vector. Nevertheless, generally speaking, a change of parameter can change the orientation of the curve.

Because the norm of $\boldsymbol{\tau}$ is constant, its derivative is a vector orthogonal to $\boldsymbol{\tau}$, see eq. (4.2). That is why we call *principal normal vector* to a curve the unit vector

$$\boldsymbol{\nu}(t) = \frac{\boldsymbol{\tau}'(t)}{|\boldsymbol{\tau}'(t)|}. \tag{4.5}$$

$\boldsymbol{\nu}$ is defined only on the points of the curve where $\boldsymbol{\tau}' \neq \mathbf{o}$ which implies that $\boldsymbol{\nu}$ is not defined on the points of a straight line. This simply means that there is not, among the infinite unit normal vectors to a straight line, a normal with special properties, a *principal* one, uniquely linked to $\boldsymbol{\tau}$.

Unlike $\boldsymbol{\tau}$, whose orientation changes with the choice of the parameter, $\boldsymbol{\nu}$ is an *intrinsic* local characteristic of the curve: *it is not affected by the choice of the parameter*. In fact, by its same definition, $\boldsymbol{\nu}$ does not depend upon the reference frame; then, because the direction of $\boldsymbol{\tau}$ is also independent upon the parameter's choice, the only factor that could affect $\boldsymbol{\nu}$ is the orientation of the curve, that depends upon the parameter. But a change of the orientation affects, in (4.5), both $\boldsymbol{\tau}$ and the sign of the increment dt, so that $\boldsymbol{\tau}'(t) = d\boldsymbol{\tau}/dt$ does not change, neither $\boldsymbol{\nu}$, which is hence an intrinsic property of the curve.

The vector

$$\boldsymbol{\beta}(t) = \boldsymbol{\tau}(t) \times \boldsymbol{\nu}(t)$$

is called the *binormal vector*; by construction, it is orthogonal to $\boldsymbol{\tau}$ and $\boldsymbol{\nu}$ and it is a unit vector. In addition, it is evident that

$$\boldsymbol{\tau} \times \boldsymbol{\nu} \cdot \boldsymbol{\beta} = 1,$$

so the set $\{\boldsymbol{\tau}, \boldsymbol{\nu}, \boldsymbol{\beta}\}$ forms a positively oriented othonormal basis that can be defined at any regular point of a curve with $\boldsymbol{\tau}' \neq \mathbf{o}$. Such a basis is called the *Frenet-Serret local basis*, local in the sense that it changes with the position along the curve. The plane $\boldsymbol{\tau} - \boldsymbol{\nu}$ is the *osculating plane*, the plane $\boldsymbol{\nu} - \boldsymbol{\beta}$ the *normal plane* and the plane $\boldsymbol{\beta} - \boldsymbol{\tau}$ the *rectifying plane*, see Fig. 4.4. The osculating plane is particularly important: if we consider a plane passing through three not aligned points of the curve, when these points

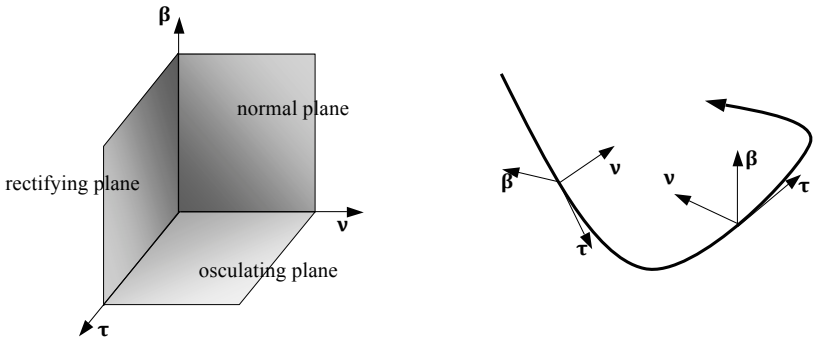

Figure 4.4: The Frenet-Serret basis.

become closer and closer, still remaining on the curve, the plane tends to the osculating plane: the osculating plane at a point of a curve is the plane that better approaches the curve near the point. A plane curve is entirely contained in the osculating plane, which is fixed.

The principal normal $\boldsymbol{\nu}$ is always oriented towards the part of the space, with respect to the rectifying plane, where the curve is; in particular, for a plane curve, $\boldsymbol{\nu}$ is always directed towards the concavity of the curve. To show it, it is sufficient to prove that the vector $p(t+\varepsilon) - p(t)$ forms with $\boldsymbol{\nu}$ an angle $\psi \leq \pi/2$, i.e. that $(p(t+\varepsilon) - p(t)) \cdot \boldsymbol{\nu} \geq 0$. In fact,

$$p(t+\varepsilon) - p(t) = \varepsilon\ p'(t) + \frac{1}{2}\varepsilon^2 p''(t) + o(\varepsilon^2)\ \rightarrow$$
$$(p(t+\varepsilon) - p(t)) \cdot \boldsymbol{\nu} = \frac{1}{2}\varepsilon^2 p''(t) \cdot \boldsymbol{\nu} + o(\varepsilon^2),$$

but

$$p''(t) \cdot \boldsymbol{\nu} = (\boldsymbol{\tau}'|p'| + \boldsymbol{\tau}|p'|') \cdot \boldsymbol{\nu} = (|\boldsymbol{\tau}'||p'|\boldsymbol{\nu} + \boldsymbol{\tau}|p'|') \cdot \boldsymbol{\nu} = |\boldsymbol{\tau}'||p'|,$$

so that, to within infinitesimal quantities of order $o(\varepsilon^2)$, we obtain

$$(p(t+\varepsilon) - p(t)) \cdot \boldsymbol{\nu} = \frac{1}{2}\varepsilon^2 |\boldsymbol{\tau}'||p'| \geq 0.$$

4.6 Curvature of a curve

It is important, in several situations, to evaluate how much a curve moves away from a straight line, in the neighborhood of a point. To do that, we calculate the angle formed by the tangents at two close points, determined by the curvilinear abscissae s and $s+\varepsilon$, and we measure the angle $\chi(s,\varepsilon)$ that they form, see Fig. 4.5.

We then define *curvature* of the curve in $p = p(s)$ the limit

$$c(s) = \lim_{\varepsilon \to 0} \left| \frac{\chi(s,\varepsilon)}{\varepsilon} \right|.$$

The curvature is hence a non-negative scalar that measures the rapidity of variation of the direction of the curve per unit length of the curve (that is why $c(s)$ is defined as a function of the curvilinear abscissa); by its same definition, the curvature is an *intrinsic*

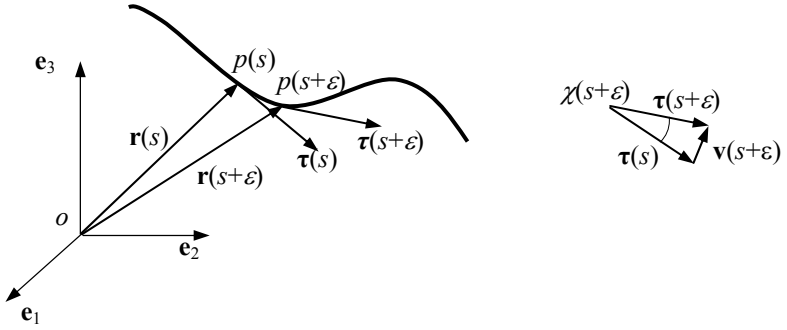

Figure 4.5: Curvature of a curve.

property of the curve, i.e. independent from the parameter's choice. For a straight line, the curvature is identically null everywhere.

The curvature is linked to the second derivative of the curve:

$$c(s) = \lim_{\varepsilon\to 0}\left|\frac{\chi(s,\varepsilon)}{\varepsilon}\right| = \lim_{\varepsilon\to 0}\left|\frac{\sin\chi(s,\varepsilon)}{\varepsilon}\right| = \lim_{\varepsilon\to 0}\left|\frac{2}{\varepsilon}\sin\frac{\chi(s,\varepsilon)}{2}\right| =$$
$$\lim_{\varepsilon\to 0}\left|\frac{\mathbf{v}(s,\varepsilon)}{\varepsilon}\right| = \lim_{\varepsilon\to 0}\left|\frac{\boldsymbol{\tau}(s+\varepsilon)-\boldsymbol{\tau}(s)}{\varepsilon}\right| = |\boldsymbol{\tau}'(s)| = |p''(s)|.$$

Another formula for the calculation of $c(s)$ can be obtained if we consider that

$$\frac{d\boldsymbol{\tau}[s(t)]}{dt} = \frac{d\boldsymbol{\tau}}{ds}\frac{ds}{dt} = \frac{d\boldsymbol{\tau}}{ds}|p'(t)| \;\rightarrow\; \frac{d\boldsymbol{\tau}}{ds} = \frac{1}{|p'(t)|}\frac{d\boldsymbol{\tau}}{dt},$$

so that

$$c(s) = |\boldsymbol{\tau}'(s)| = \frac{1}{|p'(t)|}\left|\frac{d\boldsymbol{\tau}}{dt}\right| = \frac{|\boldsymbol{\tau}'(t)|}{|p'(t)|}. \tag{4.6}$$

A better formula can be obtained as follows:

$$\frac{d\boldsymbol{\tau}}{ds} = \frac{1}{|p'(t)|}\frac{d\boldsymbol{\tau}}{dt} = \frac{1}{|p'(t)|}\frac{d}{dt}\frac{p'(t)}{|p'(t)|} = \frac{1}{|p'|}\frac{p''|p'| - p'\dfrac{p''\cdot p'}{|p'|}}{|p'|^2} =$$
$$\frac{p'' - \boldsymbol{\tau}\;p''\cdot\boldsymbol{\tau}}{|p'|^2} = (\mathbf{I}-\boldsymbol{\tau}\otimes\boldsymbol{\tau})\frac{p''}{|p'|^2}.$$

By consequence,

$$c(s) = \left|\frac{d\boldsymbol{\tau}(s)}{ds}\right| = \frac{1}{|p'|^2}|(\mathbf{I}-\boldsymbol{\tau}\otimes\boldsymbol{\tau})p''|.$$

Now, we use the following general formula expressing a skew tensor $\mathbf{W}$:

$$\mathbf{WW} = -\frac{1}{2}|\mathbf{W}|^2(\mathbf{I}-\mathbf{w}\otimes\mathbf{w});$$

if we use this formula for $\boldsymbol{\tau}$, so that $\mathbf{W}$ is the axial tensor of $\boldsymbol{\tau}$, we get

$$\mathbf{I}-\boldsymbol{\tau}\otimes\boldsymbol{\tau} = -2\frac{\mathbf{WW}}{|\mathbf{W}|^2} = -\mathbf{WW},$$

because if $\boldsymbol{\tau} = (\tau_1, \tau_2, \tau_3)$, then

$$|\mathbf{W}|^2 = \mathbf{W} \cdot \mathbf{W} = \begin{bmatrix} 0 & -\tau_3 & \tau_2 \\ \tau_3 & 0 & -\tau_1 \\ -\tau_2 & \tau_1 & 0 \end{bmatrix} \cdot \begin{bmatrix} 0 & -\tau_3 & \tau_2 \\ \tau_3 & 0 & -\tau_1 \\ -\tau_2 & \tau_1 & 0 \end{bmatrix} =$$
$$2(\tau_1^2 + \tau_2^2 + \tau_3^2) = 2.$$

So, recalling that for any skew tensor $\mathbf{W}$,

$$\mathbf{W}\ \mathbf{u} = \mathbf{w} \times \mathbf{u} \quad \forall \mathbf{u} \in \mathcal{V},$$

with $\mathbf{w}$ the axial vector of $\mathbf{W}$, we get

$$|(\mathbf{I} - \boldsymbol{\tau} \otimes \boldsymbol{\tau})p''| = | - \mathbf{W}\mathbf{W}p''| = | - \mathbf{W}(\boldsymbol{\tau} \times p'')| = | - \boldsymbol{\tau} \times (\boldsymbol{\tau} \times p'')| =$$
$$|\boldsymbol{\tau} \times (\boldsymbol{\tau} \times p'')| = |\boldsymbol{\tau} \times p''| = \frac{|p' \times p''|}{|p'|},$$

so that finally

$$c = \frac{|p' \times p''|}{|p'|^3}. \tag{4.7}$$

Applying this last formula to a plane curve $p(t) = (x(t), y(t))$, we get

$$c = \frac{|x'y'' - x''y'|}{(x'^2 + y'^2)^{\frac{3}{2}}}$$

and if the curve is given in the form $y = y(x)$, so that the parameter $t = x$, then we obtain

$$c = \frac{|y''|}{(1 + y'^2)^{\frac{3}{2}}}.$$

This last formula shows that if $|y'| \ll 1$, like in the infinitesimal theory of strain, then

$$c \simeq |y''|.$$

4.7 The Frenet-Serret formulae

From eqs. (4.5) for $t = s$ and (4.6) we get

$$\frac{d\boldsymbol{\tau}}{ds} = c\ \boldsymbol{\nu} \tag{4.8}$$

which is the *first Frenet-Serret Formula*, giving the variation of $\boldsymbol{\tau}$ per unit length of the curve. Such a variation is a vector whose norm is the curvature and that has as direction that of $\boldsymbol{\nu}$.

Let us now consider the variation of $\boldsymbol{\beta}$ per unit length of the curve; because $\boldsymbol{\beta}$ is a unit vector, we have

$$\frac{d\boldsymbol{\beta}}{ds} \cdot \boldsymbol{\beta} = 0,$$

and

$$\boldsymbol{\beta}\cdot\boldsymbol{\tau}=0\ \Rightarrow\ \frac{d(\boldsymbol{\beta}\cdot\boldsymbol{\tau})}{ds}=\frac{d\boldsymbol{\beta}}{ds}\cdot\boldsymbol{\tau}+\boldsymbol{\beta}\cdot\frac{d\boldsymbol{\tau}}{ds}=0.$$

Through eq. (4.8) and because $\boldsymbol{\beta}\cdot\boldsymbol{\nu}=0$ we get

$$\frac{d\boldsymbol{\beta}}{ds}\cdot\boldsymbol{\tau}=-c\,\boldsymbol{\beta}\cdot\boldsymbol{\nu}=0,$$

so that $\dfrac{d\boldsymbol{\beta}}{ds}$ is necessarily parallel to $\boldsymbol{\nu}$. We then put

$$\frac{d\boldsymbol{\beta}}{ds}=\vartheta\boldsymbol{\nu},$$

which is the *second Frenet-Serret formula*. The scalar $\vartheta(s)$ is called the *torsion of the curve* in $p=p(s)$. So, we see that the variation of $\boldsymbol{\beta}$ per unit length is a vector parallel to $\boldsymbol{\nu}$ and proportional to the torsion of the curve.

We can now find the variation of $\boldsymbol{\nu}$ per unit length of the curve:

$$\frac{d\boldsymbol{\nu}}{ds}=\frac{d(\boldsymbol{\beta}\times\boldsymbol{\tau})}{ds}=\frac{d\boldsymbol{\beta}}{ds}\times\boldsymbol{\tau}+\boldsymbol{\beta}\times\frac{d\boldsymbol{\tau}}{ds}=\vartheta\ \boldsymbol{\nu}\times\boldsymbol{\tau}+c\ \boldsymbol{\beta}\times\boldsymbol{\nu},$$

so finally

$$\frac{d\boldsymbol{\nu}}{ds}=-c\ \boldsymbol{\tau}-\vartheta\ \boldsymbol{\beta},$$

which is the *third Frenet-Serret formula*: the variation of $\boldsymbol{\nu}$ per unit length of the curve is a vector of the rectifying plane.

The three formulae of Frenet-Serret (discovered independently by J. F. Frenet in 1847 and by J. A. Serret in 1851) can be condensed in the symbolic matrix product

$$\begin{Bmatrix}\boldsymbol{\tau}'\\ \boldsymbol{\nu}'\\ \boldsymbol{\beta}'\end{Bmatrix}=\begin{bmatrix}0 & c & 0\\ -c & 0 & -\vartheta\\ 0 & \vartheta & 0\end{bmatrix}\begin{Bmatrix}\boldsymbol{\tau}\\ \boldsymbol{\nu}\\ \boldsymbol{\beta}\end{Bmatrix}.$$

4.8 The torsion of a curve

We have introduced the torsion of a curve in the previous section, with the second formula of Frenet-Serret. The torsion measures the deviation of a curve from flatness: if a curve is planar, it belongs to the osculating plane and $\boldsymbol{\beta}$, which is perpendicular to the osculating pane, is hence a constant vector. So, its derivative is null and by the Frenet-Serret second formula $\vartheta=0$.

Conversely, if $\vartheta=0$ everywhere, $\boldsymbol{\beta}$ is a constant vector and hence the osculating plane does not change and the curve is planar. So we have that *a curve is planar if and only if the torsion is null* $\forall p(s)$.

Using the Frenet-Serret formulae in the expression of $p'''(s)$ we get a formula for the torsion:

$$
\begin{aligned}
&p'(t) = |p'|\boldsymbol{\tau} = \frac{dp}{ds}\frac{ds}{dt} = s'\boldsymbol{\tau} \;\Rightarrow\; |p'| = s' \;\rightarrow \\
&p''(t) = s''\boldsymbol{\tau} + s'\boldsymbol{\tau}' = s''\boldsymbol{\tau} + s'^2\frac{d\boldsymbol{\tau}}{ds} = s''\boldsymbol{\tau} + c\; s'^2\boldsymbol{\nu} \;\rightarrow \\
&p'''(t) = s'''\boldsymbol{\tau} + s''\boldsymbol{\tau}' + (c\; s'^2)'\boldsymbol{\nu} + c\; s'^2\boldsymbol{\nu}' = \\
&\qquad s'''\boldsymbol{\tau} + s''s'\frac{d\boldsymbol{\tau}}{ds} + (c\; s'^2)'\boldsymbol{\nu} + c\; s'^3\frac{d\boldsymbol{\nu}}{ds} = \\
&\qquad s'''\boldsymbol{\tau} + s''s'c\boldsymbol{\nu} + (c\; s'^2)'\boldsymbol{\nu} - c\; s'^3(c\boldsymbol{\tau} + \vartheta\boldsymbol{\beta}) = \\
&\qquad (s''' - c^2s'^3)\boldsymbol{\tau} + (s''s'c + c's'^2 + 2c\; s's'')\boldsymbol{\nu} - c\; s'^3\vartheta\boldsymbol{\beta},
\end{aligned}
$$

so that, through eq. (4.7), we get

$$
\begin{aligned}
p' \times p'' \cdot p''' =& s'\boldsymbol{\tau} \times (s''\boldsymbol{\tau} + c\; s'^2\boldsymbol{\nu}) \cdot [(s''' - c^2s'^3)\boldsymbol{\tau} + \\
&(s''s'c + c's'^2 + 2c\; s's'')\boldsymbol{\nu} - c\; s'^3\vartheta\boldsymbol{\beta}] = \\
&- c^2s'^6\vartheta = -c^2|p'|^6\vartheta = -\frac{|p' \times p''|^2}{|p'|^6}|p'|^6\vartheta,
\end{aligned}
$$

so that, finally,

$$
\vartheta = -\frac{p' \times p'' \cdot p'''}{|p' \times p''|^2}.
$$

To remark that while the curvature is linked to the second derivative of the curve, the torsion is a function also of the third derivative.

Unlike curvature, which is intrinsically positive, the torsion can be negative. In fact, still using the Frenet-Serret formulae,

$$
\begin{aligned}
p(s+\varepsilon) - p(s) =& \varepsilon\; p' + \frac{1}{2}\varepsilon^2 p'' + \frac{1}{6}\varepsilon^3 p''' + o(\varepsilon^3) = \\
&\varepsilon\boldsymbol{\tau} + \frac{1}{2}\varepsilon^2 c\boldsymbol{\nu} + \frac{1}{6}\varepsilon^3(c\boldsymbol{\nu})' + o(\varepsilon^3) = \\
&\varepsilon\boldsymbol{\tau} + \frac{1}{2}\varepsilon^2 c\boldsymbol{\nu} + \frac{1}{6}\varepsilon^3(c'\boldsymbol{\nu} - c^2\boldsymbol{\tau} - c\;\vartheta\boldsymbol{\beta}) + o(\varepsilon^3) \;\rightarrow \\
&(p(s+\varepsilon) - p(s)) \cdot \boldsymbol{\beta} = -\frac{1}{6}\varepsilon^3 c\;\vartheta + o(\varepsilon^3).
\end{aligned}
$$

The above dot product determines if the point $p(s+\varepsilon)$ is located, with respect to the osculating plane, on the side of $\boldsymbol{\beta}$ or on the opposite one, see Fig. 4.6: if following the curve for increasing values of s, $\varepsilon > 0$, the point passes into the semi-space of $\boldsymbol{\beta}$ from the opposite one, because $1/6\; c\; \varepsilon^3 > 0$, it will be $\vartheta < 0$, while in the opposite case it will be $\vartheta > 0$.

This result is intrinsic, i.e. it does not depend upon the choice of the parameter, hence of the positive orientation of the curve; in fact, $\boldsymbol{\nu}$ is intrinsic, but changing the orientation of the curve, $\boldsymbol{\tau}$, and hence $\boldsymbol{\beta}$, change of orientation.

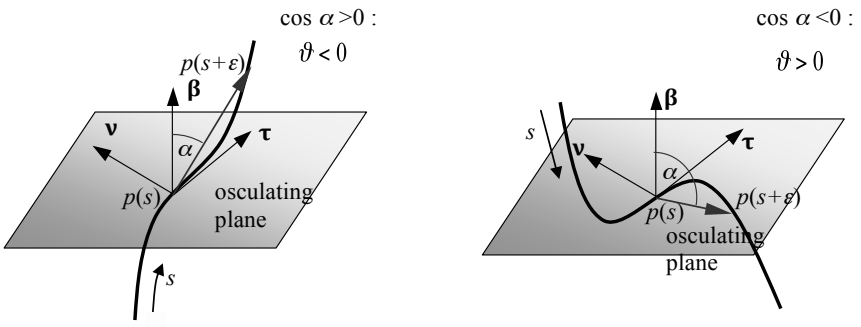

Figure 4.6: Torsion of a curve.

4.9 Osculating sphere and circle

The *osculating sphere*[2] to a curve at a point p is a sphere to which the curve tends to adhere in the neighborhood of p. Mathematically speaking, if q_s is the center of the sphere relative to point $p(s)$, then

$$|p(s+\varepsilon) - q_s|^2 = |p(s) - q_s|^2 + o(\varepsilon^3).$$

Using this definition, discarding the terms of order $o(\varepsilon^3)$ and using the Frenet-Serret formulae, we get:

$$\begin{aligned}
|p(s+\varepsilon) - q_s|^2 =& |p(s) - q_s + \varepsilon p' + \frac{1}{2}\varepsilon^2 p'' + \frac{1}{6}\varepsilon^3 p''' + o(\varepsilon^3)|^2 = \\
& |p(s) - q_s + \varepsilon \boldsymbol{\tau} + \frac{1}{2}\varepsilon^2 c\, \boldsymbol{\nu} + \frac{1}{6}\varepsilon^3 (c\boldsymbol{\nu})' + o(\varepsilon^3)|^2 = \\
& |p(s) - q_s|^2 + 2\varepsilon (p(s) - q_s)\cdot \boldsymbol{\tau} + \varepsilon^2 + \varepsilon^2 c (p(s) - q_s)\cdot \boldsymbol{\nu} + \\
& \frac{1}{3}\varepsilon^3 (p(s) - q_s)\cdot (c'\boldsymbol{\nu} - c^2\boldsymbol{\tau} - c\,\vartheta\boldsymbol{\beta}) + o(\varepsilon^3),
\end{aligned}$$

which gives

$$\begin{aligned}
(p(s) - q_s)\cdot \boldsymbol{\tau} &= 0, \\
(p(s) - q_s)\cdot \boldsymbol{\nu} &= -\frac{1}{c} = -\rho, \\
(p(s) - q_s)\cdot \boldsymbol{\beta} &= -\frac{c'}{c^2\vartheta} = \frac{\rho'}{\vartheta},
\end{aligned}$$

and finally

$$q_s = p + \rho\, \boldsymbol{\nu} - \frac{\rho'}{\vartheta}\boldsymbol{\beta}, \tag{4.9}$$

so the center of the sphere belongs to the normal plane; the sphere is not defined for a plane curve. ρ is the *radius of curvature* of the curve, defined as

$$\rho = \frac{1}{c}.$$

[2]The word osculating comes from the latin word *osculo* that means to kiss; an osculating sphere or circle or plane is a geometric object that is very close to the curve, as close as two lovers are in a kiss.

The radius of the osculating sphere is

$$\rho_s = |p - q_s| = \sqrt{\rho^2 + \left(\frac{\rho'}{\vartheta}\right)^2}.$$

The intersection between the osculating sphere and the osculating plane at a same point p is the *osculating circle*. This circle has the property to share the same tangent in p with the curve and its radius is the radius of curvature, ρ. From eq. (4.9) we get the position of the osculating circle center q:

$$q = p + \rho\, \boldsymbol{\nu}.$$

An example can be seen in Fig. 4.7, where the osculating plan, circle and sphere are shown for a point p of a conical helix.

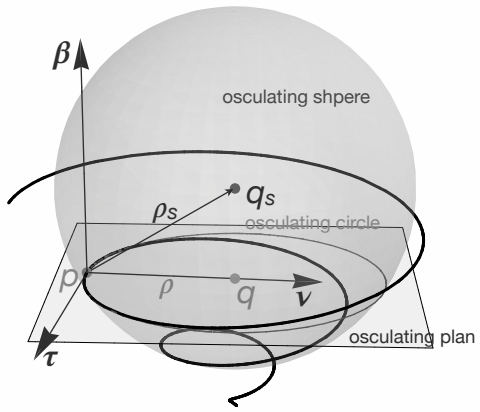

Figure 4.7: Osculating plan, circle and sphere for a point p of a conical helix.

The osculating circle is a diametral circle of the osculating sphere only when $q = q_s$, so if and only if

$$\frac{\rho'}{\vartheta} = -\frac{c'}{c^2\vartheta} = 0,$$

i.e. when the curvature is constant.

4.10 Scalar, vector and tensor fields

Let $\Omega \subset \mathcal{E}$ and $\mathbf{f} : \Omega \to \mathcal{V}$. We say that $\mathbf{f}$ is *continuous at* $p \in \Omega \iff \forall$ sequence $\pi_n = \{p_n \in \Omega, n \in \mathbb{N}\}$ that converges to $p \in \mathcal{E}$, the sequence $\{\mathbf{v}_n = \mathbf{f}(p_n), n \in \mathbb{N}\}$ converges to $\mathbf{f}(p)$ in $\mathcal{V}$. The function $\mathbf{f}(p) : \Omega \to \mathcal{V}$ is a *vector field on* Ω if it is continuous at each $p \in \Omega$. In the same way we can define a *scalar field* $\varphi(p) : \Omega \to \mathbb{R}$ and a *tensor field*, $\mathbf{L}(p) : \Omega \to Lin(\mathcal{V})$ or $\mathbb{L}(p) : \Omega \to \mathbb{L}\mathrm{in}(\mathcal{V})$.

A *deformation* is any continuous and bijective function $f(p) : \Omega \to \mathcal{E}$, i.e. any transformation of a region $\Omega \subset \mathcal{E}$ into another region of $\mathcal{E}$; bijectivity imposes that to any point $p \in \Omega$ corresponds one and only point in the transformed region, and vice-versa. This is a constraint imposed to a function from $\mathcal{E}$ to $\mathcal{E}$ ti represent a physical deformation of a body.

Finally, the basic difference between fields/deformations and curves, is that a field or a deformation is defined over a subset of $\mathcal{E}$, not of $\mathbb{R}$. In practice, this implies that the components of the field/deformation are functions of three variables, the coordinates x_i of a point $p \in \Omega$.

4.11 Differentiation of fields, differential operators

Let $\mathbf{f}(p) : \Omega \to \mathcal{V}$; we say that $\mathbf{f}$ is *differentiable in* $p_0 \in \Omega \iff \exists\ \text{grad}\mathbf{f} \in Lin(\mathcal{V})$ such that

$$\mathbf{f}(p_0 + \mathbf{u}) = \mathbf{f}(p_0) + \text{grad}\mathbf{f}(p_0)\ \mathbf{u} + o(\mathbf{u})$$

when $\mathbf{u} \to \mathbf{o}$. If $\mathbf{f}$ is differentiable $\forall p \in \Omega, \text{grad}\mathbf{f}$ defines a tensor field on Ω called the *gradient of* $\mathbf{f}$. It is also possible to define higher order differential operators, using higher order tensors, but this will not be done here. If $\mathbf{f}$ is continuous with grad$\mathbf{f}$ $\forall p \in \Omega$, then $\mathbf{f}$ is of class C^1 (*smooth*).

Let $\mathbf{f}$ a vector field of class C^1 on Ω. Then the *divergence* of $\mathbf{f}$ is the scalar field defined by

$$\text{div}\mathbf{f} := \text{tr}(\text{grad}\mathbf{f}),$$

while the *curl* of $\mathbf{f}$ is the unique vector field curl$\mathbf{f}$ that satisfies the relation

$$(\text{grad}\mathbf{f} - \text{grad}\mathbf{f}^\top)\mathbf{u} = (\text{curl}\mathbf{f}) \times \mathbf{u} \quad \forall \mathbf{u} \in \mathcal{V}.$$

The *divergence of a tensor field* $\mathbf{L}$ is the unique vector field div$\mathbf{L}$ that satisfies

$$(\text{div}\mathbf{L}) \cdot \mathbf{u} = \text{div}(\mathbf{L}^\top \mathbf{u}) \quad \forall \mathbf{u} \in \mathcal{V}.$$

Let $\varphi(p) : \Omega \to \mathbb{R}$ a scalar field over Ω. Similarly to what done for vector fields, we say that φ is *differentiable at* $p_0 \in \Omega \iff \exists\ \text{grad}\varphi \in \mathcal{V}$ such that

$$\varphi(p_0 + \mathbf{u}) = \varphi(p_0) + \text{grad}\varphi(p_0) \cdot \mathbf{u} + o(\mathbf{u})$$

when $\mathbf{u} \to \mathbf{o}$. If φ is differentiable $\forall p \in \Omega, \text{grad}\varphi$ defines a vector field on Ω called the *gradient of* φ. If gradφ is differentiable, its gradient is the tensor $\text{grad}^{II}\varphi$ called *second gradient* or *Hessian*. It is immediate to show that under continuity assumption,

$$\text{grad}^{II}\varphi = (\text{grad}^{II}\varphi)^\top.$$

A *level set* of a scalar field $\varphi(p)$ is the set $\mathcal{S}_L$ such that

$$\varphi(p) = const. \quad \forall p \in \mathcal{S}_L.$$

By the same definition of differentiability of $\varphi(p)$, we can prove that grad$\varphi(p)$ is a vector that is orthogonal to $\mathcal{S}_L$ at p. The curves of $\mathcal{E}$ that are tangent to grad$\varphi(p)$ $\forall p \in \Omega$ are the *streamlines of* φ; they have the property to be orthogonal to any $\mathcal{S}_L$ of φ $\forall p \in \Omega$.

gradφ allows to calculate the *directional derivative* of φ along any direction $\mathbf{n} \in \mathcal{S}$ as

$$\frac{d\varphi}{d\mathbf{n}} := \text{grad}\varphi \cdot \mathbf{n}.$$

The highest variation of φ is hence in the direction of $\mathrm{grad}\varphi$, and $|\mathrm{grad}\varphi|$ is the value of this variation; we remark also that $\mathrm{grad}\varphi$ is a vector directed as the increasing values of φ.

Let ψ a scalar of vector field of class C^2 at least. Then, the *laplacian* $\Delta\psi$ *of* ψ is defined by

$$\Delta\psi := \mathrm{div}(\mathrm{grad}\psi).$$

By the linearity of the trace, and hence of the divergence, we see easily that the laplacian of a vector field is the vector field whose components are the laplacian of each corresponding component of the field. A field is said to be *harmonic* on Ω if its laplacian is null $\forall p \in \Omega$.

The definitions given above for differentiable field, gradient and class C^1 can be repeated *verbatim* for a deformation $f(p) : \Omega \to \mathcal{E}$.

Let φ, ψ two scalar fields, $\mathbf{u}, \mathbf{v}, \mathbf{w}$ vector fields, $\mathbf{L}$ a tensor field and $\mathbf{W}$ the axial tensor of $\mathbf{w}$. Then, the following properties hold:

$$\mathrm{grad}(\varphi\psi) = \varphi\mathrm{grad}\psi + \psi\mathrm{grad}\varphi,$$

$$\mathrm{grad}(\varphi\mathbf{v}) = \varphi\mathrm{grad}\mathbf{v} + \mathbf{v} \otimes \mathrm{grad}\varphi,$$

$$(\mathrm{grad}\mathbf{v})\mathbf{v} = (\mathrm{curl}\mathbf{v}) \otimes \mathbf{v} + \frac{1}{2}\mathrm{grad}\mathbf{v}^2,$$

$$\mathrm{grad}(\mathbf{v} \cdot \mathbf{w}) = (\mathrm{grad}\mathbf{w})^\top\mathbf{v} + (\mathrm{grad}\mathbf{v})^\top\mathbf{w} = (\mathrm{grad}\mathbf{w})\mathbf{v} + (\mathrm{grad}\mathbf{v})\mathbf{w} + \mathbf{v} \times \mathrm{curl}\mathbf{w} + \mathbf{w} \times \mathrm{curl}\mathbf{v},$$

$$\mathrm{grad}(\mathbf{u} \cdot \mathbf{v}\ \mathbf{w}) = (\mathbf{u} \cdot \mathbf{v})\mathrm{grad}\mathbf{w} + (\mathbf{w} \otimes \mathbf{u})\mathrm{grad}\mathbf{v} + (\mathbf{w} \otimes \mathbf{v})\mathrm{grad}\mathbf{u},$$

$$\mathrm{grad}\mathbf{v} \cdot \mathrm{grad}\mathbf{v}^\top = \mathrm{div}((\mathrm{grad}\mathbf{v})\mathbf{v} - (\mathrm{div}\mathbf{v})\mathbf{v}) + (\mathrm{div}\mathbf{v})^2,$$

$$\mathrm{div}(\varphi\mathbf{v}) = \varphi\mathrm{div}\mathbf{v} + \mathbf{v} \cdot \mathrm{grad}\varphi,$$

$$\mathrm{div}(\mathbf{v} \otimes \mathbf{w}) = \mathbf{v}\mathrm{div}\mathbf{w} + (\mathrm{grad}\mathbf{v})\mathbf{w},$$

$$\mathrm{div}(\mathbf{L}^\top\mathbf{v}) = \mathbf{L} \cdot \mathrm{grad}\mathbf{v} + \mathbf{v} \cdot \mathrm{div}\mathbf{L},$$

$$\mathrm{div}(\varphi\mathbf{L}) = \varphi\mathrm{div}\mathbf{L} + \mathbf{L}\mathrm{grad}\varphi,$$

$$\mathrm{div}(\mathrm{grad}\mathbf{v}^\top) = \mathrm{grad}(\mathrm{div}\mathbf{v}),$$

$$\mathrm{div}((\mathrm{grad}\mathbf{v})\mathbf{v}) = \mathrm{grad}\mathbf{v} \cdot \mathrm{grad}\mathbf{v}^\top + \mathbf{v} \cdot \mathrm{grad}(\mathrm{div}\mathbf{v}),$$

$$\mathrm{div}(\mathbf{v} \times \mathbf{w}) = \mathbf{w} \cdot \mathrm{curl}\mathbf{v} - \mathbf{v} \cdot \mathrm{curl}\mathbf{w},$$

$$\mathrm{div}(\varphi\mathbf{L}\mathbf{v}) = \varphi\mathbf{L}^\top \cdot \mathrm{grad}\mathbf{v} + \varphi\mathbf{v} \cdot \mathrm{div}\mathbf{L}^\top + \mathbf{L}\mathbf{v} \cdot \mathrm{grad}\varphi,$$

$$\mathrm{div}(\mathrm{curl}\mathbf{v}) = 0,$$

$$\mathrm{curl}(\varphi\mathbf{v}) = \varphi\mathrm{curl}\mathbf{v} + \mathrm{grad}\varphi \times \mathbf{v},$$

$$\mathrm{curl}(\mathrm{curl}\mathbf{v}) = \mathrm{grad}(\mathrm{div}\mathbf{v}) - \mathrm{grad}\mathbf{v},$$

$$\mathrm{curl}(\mathrm{grad}\varphi) = \mathbf{o},$$

$$\mathrm{curl}(\mathbf{v} \otimes \mathbf{w}) = (\mathrm{grad}\mathbf{v})\mathbf{w} - (\mathrm{grad}\mathbf{w})\mathbf{v} + \mathbf{v}\mathrm{div}\mathbf{w} - \mathbf{w}\mathrm{div}\mathbf{v},$$

$$\mathrm{curl}\mathbf{w} = -\mathrm{div}\mathbf{W},$$

$$\Delta(\varphi\psi) = 2\mathrm{grad}\varphi \cdot \mathrm{grad}\psi + \varphi\Delta\psi + \psi\Delta\varphi,$$

$$\Delta(\mathbf{v} \cdot \mathbf{w}) = 2\mathrm{grad}\mathbf{v} \cdot \mathrm{grad}\mathbf{w} + \mathbf{v} \cdot \Delta\mathbf{w} + \mathbf{w} \cdot \Delta\mathbf{v}.$$

The proof of these properties is a good exercice for the reader (see also the book of Gurtin).

4.12 Theorems on fields

We recall here, without proof, some classical theorems on fields and operators.

Theorem 18. *(on harmonic fields): if* $\mathbf{v}(p)$ *is a vector field of class* $\geq \mathrm{C}^2$ *such that*

$$\mathrm{div}\mathbf{v} = 0, \quad \mathrm{curl}\mathbf{v} = \mathbf{o},$$

then $\mathbf{v}$ *is harmonic:* $\Delta\mathbf{v} = \mathbf{o}$.

Theorem 19. *(Potential theorem): let* $\mathbf{v}(p)$ *a vector field of class* $\geq \mathrm{C}^1$ *on a simply connected domain* $\Omega \subset \mathcal{E}$*; then*

$$\mathrm{curl}\mathbf{v} = \mathbf{o} \iff \mathbf{v} = \mathrm{grad}\varphi$$

with $\varphi(p)$ *a scalar field of class* $\geq \mathrm{C}^2$*, the potential.*

In what follows, Ω is a sufficiently regular region of $\mathcal{E}$, whose boundary is $\partial\Omega$ and the external normal $\mathbf{n} \in \mathcal{S}$.

Theorem 20. *(Divergence lemma): let* $\mathbf{v}(p)$ *a vector field of class* $\geq \mathrm{C}^1$ *on* Ω*; then*

$$\int_{\partial\Omega} \mathbf{v} \otimes \mathbf{n} \; ds = \int_{\Omega} \mathrm{grad}\mathbf{v} \; dv.$$

Theorem 21. *(Divergence or Gauss theorem): let* $\varphi, \mathbf{v}, \mathbf{L}$ *respectively a scalar, vector and tensor field on* Ω *of class* $\geq \mathrm{C}^1$*. Then*

$$\int_{\partial\Omega} \varphi\mathbf{n} \; ds = \int_{\Omega} \mathrm{grad}\varphi \; dv,$$
$$\int_{\partial\Omega} \mathbf{v} \cdot \mathbf{n} \; ds = \int_{\Omega} \mathrm{div}\mathbf{v} \; dv,$$
$$\int_{\partial\Omega} \mathbf{L}\mathbf{n} \; ds = \int_{\Omega} \mathrm{div}\mathbf{L} \; dv.$$

Theorem 22. *(Curl theorem): let* $\mathbf{v}(p)$ *a vector field of class* $\geq \mathrm{C}^1$ *on* Ω*; then*

$$\int_{\partial\Omega} \mathbf{n} \times \mathbf{v} \; ds = \int_{\Omega} \mathrm{curl}\mathbf{v} \; dv.$$

Theorem 23. *(Stokes theorem): let* $\mathbf{v}(p)$ *a vector field of class* $\geq \mathrm{C}^1$ *on* Ω *and be* Σ *an open surface whose support is the closed line* γ *and* $\mathbf{n} \in \mathcal{S}$ *the normal, see Fig. 4.8. Then*

$$\oint_{\gamma} \mathbf{v} \cdot d\ell = \int_{\Sigma} \mathrm{curl}\mathbf{v} \cdot \mathbf{n} \; ds.$$

The parametric equation of γ must be chosen in such a way that

$$p'(t_1) \times p'(t_2) \cdot \mathbf{n} > 0 \quad \forall t_2 > t_1.$$

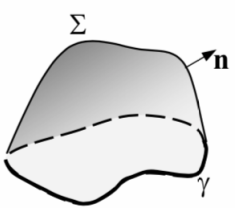

Figure 4.8: Scheme for the Stokes theorem.

Theorem 24. *(Green's formula): let $\varphi(p), \psi(p)$ two scalar fields on Ω of class $\geq \mathrm{C}^2$; then*

$$\int_{\partial\Omega} \left(\psi \frac{d\varphi}{d\mathbf{n}} - \varphi \frac{d\psi}{d\mathbf{n}} \right) ds = \int_{\Omega} (\psi \ \Delta\varphi - \varphi \ \Delta\psi) dv.$$

Theorem 25. *(Flux theorem): let $\mathbf{v}(p)$ a vector field of class $\geq \mathrm{C}^1$ on an open subset $\mathtt{R}$ of $\mathcal{E}$. Then*

$$\mathrm{div}\mathbf{v} = 0 \iff \int_{\partial\Omega} \mathbf{v} \cdot \mathbf{n} \ ds = 0 \quad \forall \Omega \subset \mathtt{R}.$$

4.13 Differential operators in Cartesian coordinates

In what follows, $f, \mathbf{v}, \mathbf{L}$ are respectively a scalar, vector and tensor field. The Cartesian components[3] of the differential operators are[4]

$$\begin{aligned}
&(\mathrm{grad}f)_i = f_i, \\
&(\mathrm{grad}\mathbf{v})_{ij} = v_{i,j}, \\
&div\mathbf{v} = v_{i,i}, \\
&(\mathrm{div}\mathbf{L})_i = L_{ij,j}, \\
&\Delta f = f_{,ii}, \\
&(\Delta\mathbf{v})_i = \Delta v_i = v_{i,jj}, \\
&\mathrm{curl}\mathbf{v} = (v_{3,2} - v_{2,3}, v_{1,3} - v_{3,1}, v_{2,1} - v_{1,2}).
\end{aligned}$$

The so-called *operator nabla* ∇:

$$\nabla := \frac{\partial \cdot}{\partial x_i}\mathbf{e}_i = \frac{\partial \cdot}{\partial x_1}\mathbf{e}_1 + \frac{\partial \cdot}{\partial x_2}\mathbf{e}_2 + \frac{\partial \cdot}{\partial x_3}\mathbf{e}_3$$

is often used to indicate the differential operators:

$$\begin{aligned}
&\mathrm{grad}f = \nabla f, \\
&\mathrm{div}\mathbf{v} = \nabla \cdot \mathbf{v}, \\
&\mathrm{curl}\mathbf{v} = \nabla \times \mathbf{v}, \\
&\Delta f = \nabla^2 f.
\end{aligned}$$

[3] In the following formulae, the Einstein summation rule holds.

[4] The comma indicates partial derivative, e.g. $f_{i,j} = \dfrac{\partial f_i}{\partial x_j}$.

4.14 Differential operators in cylindrical coordinates

The cylindrical coordinates ρ, θ, z of a point p, whose Cartesian coordinates in the (fixed) frame $\{o; \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ are $p = (x_1, x_2, x_3)$, are shown in Fig. 4.9. They are related together by

$$\begin{aligned} \rho &= \sqrt{x_1^2 + x_2^2}, \\ \theta &= \arctan \frac{x_2}{x_1}, \\ z &= x_3, \end{aligned}$$

or conversely

$$\begin{aligned} x_1 &= \rho \cos\theta, \\ x_2 &= \rho \sin\theta, \\ x_3 &= z. \end{aligned}$$

To notice that $\rho \geq 0$ and that the *anomaly* θ is bounded by $0 \leq \theta < 2\pi$.

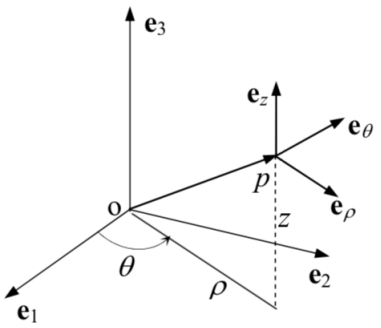

Figure 4.9: Cylindrical coordinates.

In the (local) frame $\{p; \mathbf{e}_\rho, \mathbf{e}_\theta, \mathbf{e}_z\}$, the differential operators are

$$\nabla f = \left(f_{,\rho}, \frac{1}{\rho} f_{,\theta}, f_{,z} \right),$$

$$\Delta f = \frac{1}{\rho}(\rho f_{,\rho})_{,\rho} + \frac{1}{\rho^2} f_{,\theta\theta} + f_{,zz},$$

$$\nabla \mathbf{v} = \begin{bmatrix} v_{\rho,\rho} & \frac{1}{\rho}(v_{\rho,\theta} - v_\theta) & v_{\rho,z} \\ v_{\theta,\rho} & \frac{1}{\rho}(v_{\theta,\theta} + v_\rho) & v_{\theta,z} \\ v_{z,\rho} & \frac{1}{\rho} v_{z,\theta} & v_{z,z} \end{bmatrix},$$

$$\mathrm{div}\mathbf{v} = v_{\rho,\rho} + \frac{1}{\rho}(v_{\theta,\theta} + v_\rho) + v_{z,z},$$

$$\mathrm{curl}\mathbf{v} = \left(\frac{1}{\rho} v_{z,\theta} - v_{\theta,z}, v_{\rho,z} - v_{z,\rho}, \frac{1}{\rho}((\rho v_\theta)_{,\rho} - v_{\rho,\theta}) \right),$$

$$\mathrm{div}\mathbf{L} = \begin{bmatrix} \frac{1}{\rho}((\rho L_{\rho\rho})_{,\rho} + L_{\rho\theta,\theta} - L_{\theta\theta}) + L_{\rho z,z} \\ L_{\theta\rho,\rho} + \frac{1}{\rho}(L_{\theta\theta,\theta} + L_{\rho\theta} + L_{\theta\rho}) + L_{\theta z,z} \\ \frac{1}{\rho}((\rho L_{z\rho})_{,\rho} + L_{z\theta,\theta}) + L_{zz,z} \end{bmatrix},$$

$$\Delta\mathbf{v} = \begin{pmatrix} \Delta v_\rho \\ \Delta v_\theta \\ \Delta v_z \end{pmatrix} = \begin{pmatrix} \frac{1}{\rho}(\rho v_{\rho,\rho})_{,\rho} + \frac{1}{\rho^2}v_{\rho,\theta\theta} + v_{\rho,zz} - \frac{1}{\rho^2}(v_\rho + 2v_{\theta,\theta}) \\ \frac{1}{\rho}(\rho v_{\theta,\rho})_{,\rho} + \frac{1}{\rho^2}v_{\theta,\theta\theta} + v_{\theta,zz} - \frac{1}{\rho^2}(v_\theta - 2v_{\rho,\theta}) \\ \frac{1}{\rho}(\rho v_{z,\rho})_{,\rho} + \frac{1}{\rho^2}v_{z,\theta\theta} + v_{z,zz} \end{pmatrix}.$$

4.15 Differential operators in spherical coordinates

The spherical coordinates r, φ, θ of a point p, whose Cartesian coordinates in the (fixed) frame $\{o; \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ are $p = (x_1, x_2, x_3)$, are shown in Fig. 4.10. They are related together by

$$\begin{aligned} r &= \sqrt{x_1^2 + x_2^2 + x_3^2}, \\ \varphi &= \arctan\frac{\sqrt{x_1^2 + x_2^2}}{x_3}, \\ \theta &= \arctan\frac{x_2}{x_1}, \end{aligned}$$

or conversely

$$\begin{aligned} x_1 &= r\cos\theta\sin\varphi, \\ x_2 &= r\sin\theta\sin\varphi, \\ x_3 &= r\cos\varphi. \end{aligned}$$

To notice that $r \geq 0$ and that the *anomaly* θ is bounded by $0 \leq \theta < 2\pi$ while the *colatitude* φ by $0 \leq \varphi \leq \pi$.

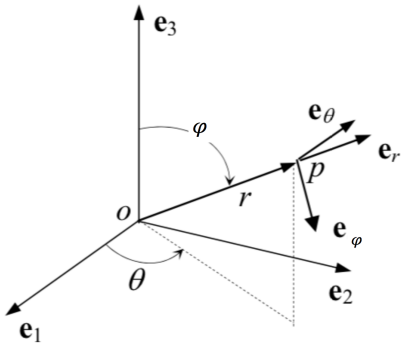

Figure 4.10: Spherical coordinates.

In the (local) frame $\{p; \mathbf{e}_r, \mathbf{e}_\varphi, \mathbf{e}_\theta\}$, the differential operators are

$$\nabla f = \left(f_{,r}, \frac{1}{r}f_{,\varphi}, \frac{1}{r\sin\varphi}f_{,\theta}\right),$$

$$\Delta f = \frac{1}{r^2}(r^2 f_{,r})_{,r} + \frac{1}{r^2 \sin\varphi}(f_{,\theta\theta} + (f_{,\varphi}\sin\varphi)_{,\varphi}),$$

$$\nabla \mathbf{v} = \begin{bmatrix} v_{r,r} & \frac{1}{r}(v_{r,\varphi} - v_\varphi) & \frac{1}{r}\left(\frac{1}{\sin\varphi} v_{r,\theta} - v_\theta\right) \\ v_{\varphi,r} & \frac{1}{r}(v_{\varphi,\varphi} + v_r) & \frac{1}{r}\left(\frac{1}{\sin\varphi} v_{\varphi,\theta} - v_\theta \cot\varphi\right) \\ v_{\theta,r} & \frac{1}{r} v_{\theta,\varphi} & \frac{1}{r}\left(\frac{1}{\sin\varphi} v_{\theta,\theta} + v_r + v_\varphi \cot\varphi\right) \end{bmatrix},$$

$$\mathrm{div}\mathbf{v} = \frac{1}{r^2}(r^2 v_r)_{,r} + \frac{1}{r\sin\varphi}(v_{\varphi,\varphi}\sin\theta + v_{\theta,\theta}),$$

$$\mathrm{curl}\mathbf{v} = \left(\frac{1}{r\sin\varphi}(v_{\theta,\varphi}\sin\theta - v_{\varphi,\theta}), \frac{1}{r\sin\varphi} v_{r,\theta} - \frac{1}{r}(r v_\theta)_{,r}, \frac{1}{r}((r v_\varphi)_{,r} - v_{r,\varphi})\right),$$

$$\mathrm{div}\mathbf{L} = \begin{pmatrix} \frac{1}{r^2}(r^2 L_{rr})_{,r} + \frac{1}{r} L_{r\varphi,\varphi} + \frac{1}{r\sin\varphi} L_{r\theta,\theta} - \frac{L_{\varphi\varphi} + L_{\theta\theta}}{r} + \frac{\cot\varphi}{r} L_{r\varphi} \\ \frac{1}{r^2}(r^2 L_{\varphi r})_{,r} + \frac{1}{r} L_{\varphi\varphi,\varphi} + \frac{1}{r\sin\varphi} L_{\varphi\theta,\theta} + \frac{1}{r} L_{r\varphi} + \frac{\cot\varphi}{r}(L_{\varphi\varphi} - L_{\theta\theta}) \\ \frac{1}{r^2}(r^2 L_{\theta r})_{,r} + \frac{1}{r} L_{\theta\varphi,\varphi} + \frac{1}{r\sin\varphi} L_{\theta\theta,\theta} + \frac{1}{r} L_{r\theta} + \frac{\cot\varphi}{r}(L_{\varphi\theta} + L_{\theta\varphi}) \end{pmatrix},$$

$$\Delta \mathbf{v} = \begin{pmatrix} v_{r,rr} + \frac{2v_{r,r}}{r} + \frac{v_{r,\varphi\varphi} - 2v_{\varphi,\varphi}}{r^2} + \frac{v_{r,\varphi} - 2v_\varphi}{r^2\tan\varphi} + \frac{1}{r^2\sin\varphi}\left(\frac{v_{r,\theta\theta}}{\sin\varphi} - 2v_{\theta,\theta}\right) - \frac{2v_r}{r^2} \\ v_{\varphi,rr} + \frac{2v_{\varphi,r}}{r} + \frac{v_{\varphi,\varphi\varphi} + 2v_{r,\varphi}}{r^2} + \frac{v_{\varphi,\varphi} - v_\varphi\cot\varphi}{r^2\tan\varphi} + \frac{1}{r^2\sin\varphi}\left(\frac{v_{\varphi,\theta\theta}}{\sin\varphi} - 2v_{\theta,\theta}\cot\varphi\right) - \frac{v_\varphi}{r^2} \\ v_{\theta,rr} + \frac{2v_{\theta,r}}{r} + \frac{v_{\theta,\varphi\varphi}}{r^2} + \left(v_{\theta,\varphi} + \frac{2v_{\varphi,\theta}}{\sin\varphi}\right)\frac{1}{r^2\tan\varphi} + \frac{1}{r^2\sin\varphi}\left(\frac{v_{\theta,\theta\theta}}{\sin\varphi} + 2v_{r,\theta}\right) - \frac{v_\theta}{r^2\sin^2\varphi} \end{pmatrix}.$$

4.16 Exercices

1. Using the same definition of derivative of a curve, prove the relations in eq. (4.1).
2. Prove the relations in eq. (4.3).
3. The curve whose polar equation is

$$r = a\ \theta, \quad a \in \mathbb{R},$$

   is an *Archimede's spiral*. Find its curvature, its length for $\theta \in [0, 2\pi)$ and prove that any straight line passing by the origin is divided by the spiral in segments of constant length $2\pi\ a$ (that is why it is used to record disks).

4. The curve whose polar equation is

$$r = a\ \mathrm{e}^{b\theta}, \quad a, b \in \mathbb{R},$$

   is the *logarithmic spiral*. Prove that the origin is an asymptotic point of the curve, find its curvature and the length of the segment in which a straight line by the origin is divided by two consecutive intersections with the spiral. Then prove that the curve is plane and its *equiangular property*: $(p(\theta) - o) \cdot \boldsymbol{\tau}(\theta) = const.$

5. The curve whose parametric equation is

$$p(\theta) = a(\cos\theta + \theta\sin\theta)\mathbf{e}_1 + a(\sin\theta - \theta\cos\theta)\mathbf{e}_2$$

with the parameter θ the angle formed by $p(\theta) - o$ with the x_1-axis is the *involute of the circle*. Find its curvature and length for $\theta \in [0, 2\pi)$ and prove that the geometrical set of the points $p(\theta) + \rho(\theta)\boldsymbol{\nu}(\theta)$ is exactly the circle of center o and radius a (that is why the involute of the circle is used to profile engrenages).

6. The curve whose parametric equation is

$$p(\theta) = a\cos\omega\theta\mathbf{e}_1 + a\sin\omega\theta\mathbf{e}_2 + b\omega\theta\mathbf{e}_3$$

is a *helix* that winds on a circular cylinder of radius a. Show that the angle formed by the helix and any generatrix of the cylinder is constant (a property that defines a helix in the general case). Then, find its length for $\theta \in [0, 2\pi)$, curvature, torsion and pitch (the distance, on a same generatrix, between two successive intersections with the helix). Prove then the *Bertrand's theorem*: a curve is a cylindrical helix if and only if the ratio $c/\vartheta = const$. Finally, prove that for the above circular helix there are two constants A and B such that

$$p' \times p'' = A\mathbf{u}(\theta) + B\mathbf{e}_3,$$

with

$$\mathbf{u} = \sin\omega\theta\mathbf{e}_1 - \cos\omega\theta\mathbf{e}_2;$$

find then A and B.

7. For the curve whose cylindrical equation is

$$\begin{cases} r = 1, \\ z = \sin\theta \end{cases}$$

find the highest curvature and determine whether or not it is planar.

8. Consider a rigid body $\mathtt{B}$, and a point $p_0 \in \mathtt{B}$. From the kinematics of rigid bodies, we now that the velocity of another point $p \in \mathtt{B}$ is given by

$$\mathbf{v}(p) = \mathbf{v}(p_0) + \boldsymbol{\omega} \times (p - p_0),$$

with $\boldsymbol{\omega}$ the angular velocity. Prove that

$$\boldsymbol{\omega} = \frac{1}{2}\mathrm{curl}\mathbf{v}.$$

9. Prove the relations at the end of Sect. 4.11.

10. Prove the three forms of the Gauss Theorem using the Divergence lemma.

11. Make use of the tensor form of the Gauss Theorem to prove the Curl Theorem.

12. Prove the following identities using the Gauss theorem:

$$\int_{\partial\Omega} \mathbf{v}\cdot\mathbf{L}\mathbf{n}\ ds = \int_{\Omega} (\mathbf{v}\cdot \mathrm{div}\mathbf{L} + \mathbf{L}\cdot\nabla\mathbf{v})dv,$$
$$\int_{\partial\Omega} (\mathbf{L}\mathbf{n})\otimes\mathbf{v}\ ds = \int_{\Omega} ((\mathrm{div}\mathbf{L})\otimes\mathbf{v} + \mathbf{L}(\nabla\mathbf{v}^{\top}))dv,$$
$$\int_{\partial\Omega} (\mathbf{w}\cdot\mathbf{n})\mathbf{v}\ ds = \int_{\Omega} (\mathbf{v}\mathrm{div}\mathbf{w} + (\nabla\mathbf{v})\mathbf{w})dv.$$

Suggested texts

1. P. R. Halmos: *Finite-dimensional vector spaces.* Van Nostrand Reynold, 1958.
2. M. E. Gurtin: *An introduction to continuum mechanics.* Academic Press, 1981.
3. P. Podio-Guidugli: *A primer in elasticity.* Journal of Elasticity, v. 58: 1-104, 2000.