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High temperature spin-driven multiferroicity in ludwigite chromocuprate Cu_2CrBO_5

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ABSTRACT

We report spin-driven multiferroicity above 100 K in the ludwigite Cu(II) oxyborate Cu_2CrBO_5 . Spontaneous polarization, which reaches $35 \mu\text{C}\cdot\text{m}^{-2}$ at 5 K, appears below 120 K, concomitantly with an incommensurate antiferromagnetic order and complex magnetodielectric effects. In magnetically induced ferroelectrics, multiferroicity usually appears at low temperature, because of the competing magnetic exchanges needed to stabilize a magnetic spiral : the remarkably high transition temperature observed in Cu_2CrBO_5 originates from the presence of strong Cu-O-Cu magnetic super-exchange interactions, which are not weakened by Cu/Cr cationic disorder. Our result provides an important contribution to the search for high temperature spin-driven multiferroics amongst low-dimensional cuprates.

Introduction

To understand and optimize the coupling between magnetic and dielectric properties in multiferroics is one of the enduring challenge of condensed matter physics [1], [2]. Amongst the most promising multiferroic compounds are spin-driven ferroelectrics, in which ferroelectricity emerges concomitantly with a non-collinear spiral magnetic order breaking inversion symmetry [3]. Different microscopic mechanisms are at play depending on whether this complex magnetic ordering is cycloidal [4] or helicoidal [5], but in both cases, a strong coupling between the ferroic parameters is expected, hence their high potential for applications [6]. The design of spiral magnetic order is not straightforward, however, as it often relies on the existence of competing (frustrated) magnetic interactions, which also means that such orders ordinarily occur at low temperature, typically below 40 K [7]. To overcome this issue, compounds with large magnetic super-exchange interactions have been studied, as, for instance, hexaferrites, in which magnetic super-exchange between spin blocks can be tuned to produce long-wavelength magnetic structures close to or above room temperature (RT), and large magnetoelectric effects [8], [9], [10]. Another example is cupric oxide CuO, with one of the largest Curie temperature (T_C) amongst spin-driven multiferroics so far, ~ 230 K [11], owing to the presence of a strong magnetic super-exchange $J \sim 700$ K. Yet, CuO is ferroelectric in a narrow temperature range of about 20 K only, which corresponds to the range of existence of its spiral magnetic order : this serious limitation calls for better candidates.

As rich as that of silicates or phosphates, borate chemistry has always provided interesting structural frameworks [12], with, sometimes, unique functionalities, especially in non-linear optics [13]. A notorious oxyborate crystal structure for transition metals (TM) is for instance the ludwigite, of formula $M_2M'BO_5$ (M and M' being divalent and trivalent TM, respectively). Its intricate framework of edge-sharing TM octahedra, prone to frustration, has been studied in the past for a wide variety of uncommon features, including mixed valence and dimerization [14], [15], [16], [17], magnetic sub-lattices decoupling [18], [19], ferroelectricity [20] and for

its potential use in lithium-ion battery [21]. Literature on heterometallic copper ludwigites $\text{Cu}_2\text{M}'\text{BO}_5$ ($\text{M}' = \text{Al}, \text{Fe}, \text{Ga}, \text{Mn}$) is in contrast somewhat scarcer. Indeed, Cu/M' ions non-random distribution over at least two out of the four crystallographically distinct TM sites of the ludwigite structure is always found to be an issue [22], [23], [24], [25], which, combined with the complex geometry of the magnetic exchange paths [26], is known to affect physical properties [27], [28] and hinder reproducibility [29].

To the best of our knowledge, ludwigite Cu_2CrBO_5 has only been investigated up to now for its redox activity [21]. In this article, we show that Cu_2CrBO_5 is a remarkable example amongst copper ludwigites of a perfect cationic order between $\text{Cu}(\text{II})$ and $\text{Cr}(\text{III})$ species, and the first instance in this family of a spin-driven multiferroic above 100 K with complex magnetodielectric properties. Perpendicular chains of strongly interacting magnetic spins, weakly coupled, are proposed to explain the complex non-centrosymmetric magnetic order at the origin of the spin driven multiferroic properties.

Experimental Methods

Cu_2CrBO_5 was prepared starting from a mixture of 2 $\text{CuO}/0.5 \text{Cr}_2\text{O}_3/0.5 \text{B}_2\text{O}_3$ (^{11}B enriched), heated in air at 900°C for 24 hrs. An impurity of CrBO_3 (~ 0.5 wt %) was identified.

Synchrotron X-ray powder diffraction (SXRPD) experiments were performed between 10 K and 260 K at BL04-MSPD at ALBA synchrotron (Barcelona, Spain) ($\lambda = 0.44296 \text{ \AA}$) [30].

Cu_2CrBO_5 powders were 20 μm sieved and filled in 0.5 mm \varnothing borosilicate capillaries.

Diffraction patterns were collected using the Position Sensitive Detector Mythen. Neutron powder diffraction (NPD) experiments were performed at room temperature on the D2B high-resolution diffractometer ($\lambda = 1.5946 \text{ \AA}$), at ILL (Grenoble, France), and between 1.5 K and 300 K on the G4.1 diffractometer ($\lambda = 2.426 \text{ \AA}$), at LLB-Orphée (Gif-sur-Yvette, France). Rietveld refinements were performed using the Fullprof Suite [31].

Magnetic susceptibility (χ) was derived from magnetization data recorded in a field of 1 T, on warming from 5 to 400 K after a zero-field cooling (MPMS, Quantum Design). Isothermal magnetic field (H) dependent magnetization (M) loops were recorded using the AC measurement system of a 9 T Physical Properties Measurement System (PPMS, Quantum Design). A homemade sample probe inserted in the PPMS was used to measure the dielectric permittivity (ϵ') and the electric polarization (P). For that purpose, square thin platelets (2 x 2 x 0.5 mm) were cut from the reacted bars. The electrical contacts were deposited on the opposite largest faces using silver paint. The electrode was soldered to the sample probe using Cu wires. ϵ' was measured as a function of T or H using a LCR meter (Agilent 4284A) whereas P was obtained by integrating the pyroelectric current collected via an electrometer (Keithley 6514A). For the P(T) curves, the data were collected upon warming at 4 K.min⁻¹.

Results

Cu₂CrBO₅ crystallizes at RT with the monoclinic space group $P2_1/c$ (#14, $a = 3.05487(2)$ Å, $b = 12.18070(8)$ Å, $c = 9.41432(7)$ Å and $\beta = 94.5602(4)^\circ$), like Cu-based ludwigites Cu₂FeBO₅ and Cu₂GaBO₅ [23]. Rietveld refinements are shown in Figure S1, and corresponding structural parameters are given in Tables S1, S2 and S3. Its framework can be described by TM oxygen polyhedra sharing edges to form zigzag walls, themselves connected to each other through apexes. These walls delimit triangular tunnels occupied by boron atoms (Figure 1a), with typical B-O bond lengths ~ 1.35 - $1.37(1)$ Å [23]. A few specifics of the Cu₂CrBO₅ structure, however, are noteworthy. A major structural difference lies in the shift of the Cu1 position to a $2c$ ($0 \frac{1}{2} 0$) site in Cu₂CrBO₅, which flattens the usually strongly corrugated Cu chains found in ludwigites. Moreover, no Cu/Cr mixing is observed on the TM sites, in contrast with other heterometallic copper ludwigites. In Cu₂CrBO₅, the oxygen coordination around all three Cu sites is (4 + 2), in agreement with a Jahn-Teller ion like Cu(II), with the axial Cu-O bonds being ~ 20 to ~ 50 % longer than the equatorial ones (Table S2). CrO₆ octahedra are very regular in

contrast, with the Cr-O bonds being equal within a few % (Table S2). Based on those considerations, the crystal structure of Cu_2CrBO_5 can be described as chains, running along b , of CuO_4 squares sharing edges or corners, and connected through their edges or corners to chains (parallel to a) of edge-sharing CrO_6 octahedra (Figure 1b). The shortest distance between TM corresponds to CuO_4 sharing edges ($\text{Cu1-Cu3} \sim 2.80 \text{ \AA}$, Table S3), forming trimer-like motifs (thick red segments on Figure 1a).

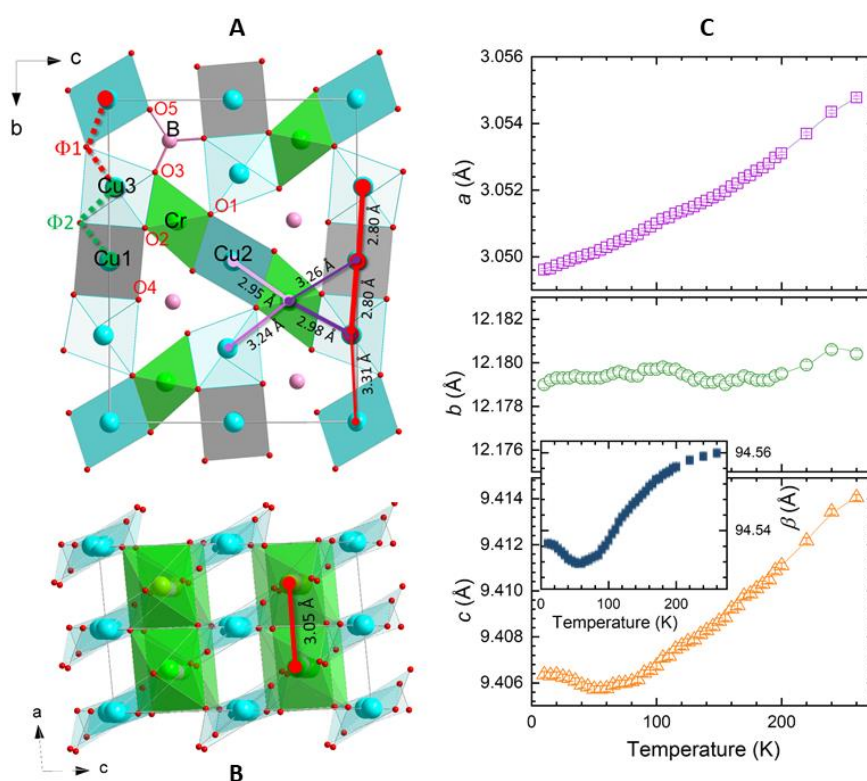


Figure 1. Projections of the monoclinic $P2_1/c$ crystal structure of Cu_2CrBO_5 along a (A) or along b (B). Cu atoms are in blue, Cr in green, B in pink and O in red. Distances between transition metals at RT are indicated, as well as the Cu-O-Cu angles Φ_1 and Φ_2 . C illustrates the temperature evolution of the cell parameters a , b , and c , and of the monoclinic angle β (inset) (from SXRPD Rietveld refinements).

The Cu_2CrBO_5 crystal structure is stable with decreasing temperature; it remains monoclinic $P2_1/c$ down to 10 K, the lattice parameters decreasing with a less than 0.1 % variation between

275 K and 10 K (Figure 1c). Slight increases of the c parameter and of the monoclinic distortion are observed, however, below 60 K (inset of Figure 1c).

Figure 2a shows the temperature evolution of the zero-field cooled (zfc) magnetic susceptibility of Cu_2CrBO_5 in 1 T. A susceptibility maximum is clearly seen at $T_N = 117$ K, suggesting a transition to an antiferromagnetic state. Above T_N , the susceptibility does not follow a Curie-Weiss law in the temperature range investigated ($T < 400$ K, inset of Figure 2a), which could be due to low-dimensional effects or short range magnetic ordering, and prevented us from extracting a meaningful Θ_{CW} . At 119 K, concomitantly with the magnetic transition, a λ -like peak is observed on the zero-field dielectric constant ϵ' (Figure 2b, left axis), and is accompanied by a step on the dielectric loss, whose low values indicate very small losses in the vicinity of the transition temperature (Figure 2b, right axis). This transition is frequency-invariant over the f -range studied ($f = 5$ to 100 kHz, not shown), thus excluding spurious experimental effects. Figure 2c shows the evolution with temperature of the electric polarization P obtained from measurements of the pyroelectric current. A finite P is measured in zero magnetic field below 122 K, and down to 5 K, where it reaches $\sim 35 \mu\text{C}\cdot\text{m}^{-2}$ for a poling field of $181.8 \text{ kV}\cdot\text{m}^{-1}$. P can be switched by inverting the sign of the poling electric field, unambiguously confirming the ferroelectric nature of Cu_2CrBO_5 in its magnetically ordered state. The $\epsilon'(T)_H$ curves (Figure 2d) reveal a clear magnetodielectric effect : as H increases, the dielectric peak shifts to lower temperature (-4 K between 0 and 9 T), with, up to 5 T, an increase of the peak amplitude. Above 5 T, the ϵ' peak starts to broaden markedly with H , its amplitude decreasing concomitantly. Further measurements versus magnetic field H were therefore performed to characterize this effect. Three different behaviors are identified, as illustrated by the magnetodielectric curves $\text{MDE} = ((\epsilon'_H - \epsilon'_0)/\epsilon'_0)(H)$ of Figure 2e ; just above T_N , there is no magnetodielectric effect ; just below T_N , a negative MDE is observed at 116 K, while a positive MDE is measured at 112 K. On the latter, a maximum of 0.6 % is reached at ± 5.9 T. This

characteristic magnetic field of 5.9 T at 112 K corresponds to a change of slope on the magnetization $M(H)$ curves, more clearly identified on the $dM/dH(H)$ curves (Figure S2), and is indicative of a metamagnetic transition in the vicinity of T_N , more generally suggesting for Cu_2CrBO_5 a rich multiferroic phase diagram to explore.

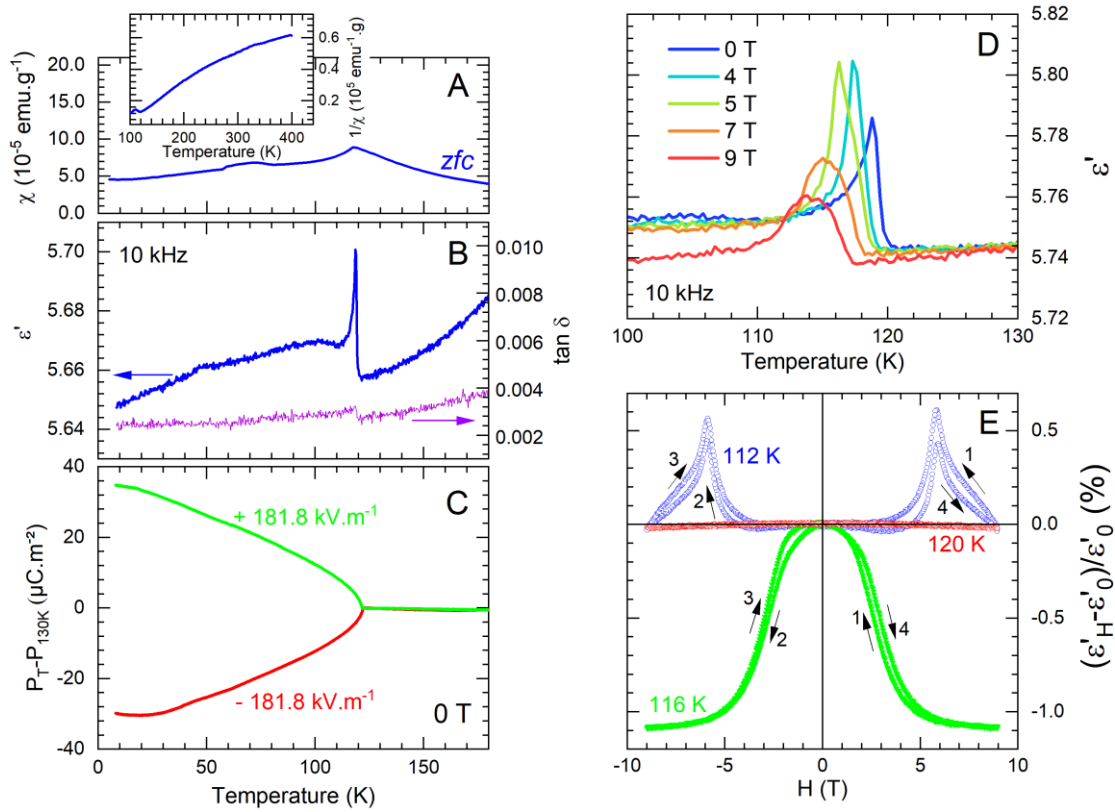


Figure 2. Temperature (T) dependence of the zfc magnetic susceptibility χ in 1 T (A). The small bump observed around 70 K is of unknown origin. The inset shows $1/\chi$ (T) in 0.01T, in the range 100 - 400 K. Temperature evolution of the dielectric constant ϵ' (B-right axis), of the corresponding loss ($\tan \delta$) (B-right axis) at 10 kHz in 0 T, and of the spontaneous polarization ΔP after +/- 181.8 kV.m^{-1} electric field poling (C). Magnetic field evolution of the dielectric constant $\epsilon'(T)$ measured at 10 kHz between 0 and 9 T (D). E shows the magnetodielectric coefficient MDE defined as $((\epsilon'_H - \epsilon'_0)/\epsilon'_0)$ versus H at 112 K, 116 K and 120 K.

A proper understanding of the microscopic mechanism behind the multiferroic properties of Cu_2CrBO_5 requires knowledge of its magnetic ordering. Cold neutron diffraction experiments were performed down to 1.5 K (Figure 3a). Below T_N , new Bragg peaks of magnetic origin

can be seen on the neutron diffractograms, whose incommensurate Q positions vary with temperature down to 30 K (inset of Figure 3a) and then remain invariant down to 1.5 K, but still incommensurate. Indexing of these peaks actually proved to be challenging, owing to the limited number of magnetic peaks. Several solutions with similar reliability factor could be obtained for the magnetic propagation vector \mathbf{k} . One of them, with only one incommensurate modulation along a , corresponds to $\mathbf{k} = (0.09\ 0\ 0.25)$ (illustrated on Figure S3), while all the other solutions involve at least two incommensurate modulations along b and c (or a and c). It is not possible to characterize further the magnetic ordering at that stage, because of the low symmetry of the crystal structure, the existence of four distinct magnetic sites, and the uncertainty on the number of modulations for \mathbf{k} . The fact that the magnetic state is ferroelectric nevertheless suggests a long wavelength helicoidal or cycloidal magnetic structure, rather than a spin amplitude modulation, arguably involving both Cr and Cu spins.

This incommensurate magnetic propagation vector confirms that the ludwigite framework of Cu_2CrBO_5 is host to frustration. Owing to its different crystal structure, the magnetic topology in Cu_2CrBO_5 is distinct from that of other ludwigites, and can be seen as two perpendicular sublattices made of spin chains. One chain is made of Cu spins and runs along b ; because it is made of an alternation of CuO_4 units sharing corners or edges, magnetic interactions in these chains are either of the super-exchange type (through the angle Φ_1 on Figure 1a) or direct-exchange like (through the angle Φ_2 on Figure 1a). This chain thus corresponds to an alternation of magnetic exchanges $-J_{\Phi_1}-J_{\Phi_1}-J_{\Phi_2}-J_{\Phi_2}-$ between Cu atoms (Figure 3b). The second magnetic sublattice is made of Cr chains. As CrO_6 octahedra only share edges along a , Cr-Cr interactions are through a unique direct-exchange J_a (Figure 3b, in which only nearest-neighbor (NN) interactions only are considered for simplicity). A minimum of six inter-sublattice interactions have to be taken into account in addition: they involve triangular configurations in the bc plane and along a (Figure 3b and Table S3), so that competing interactions are

conceivable in the three directions. This will be dwelt upon further in the discussion, after considerations on J_{Φ_1} and J_{Φ_2} .

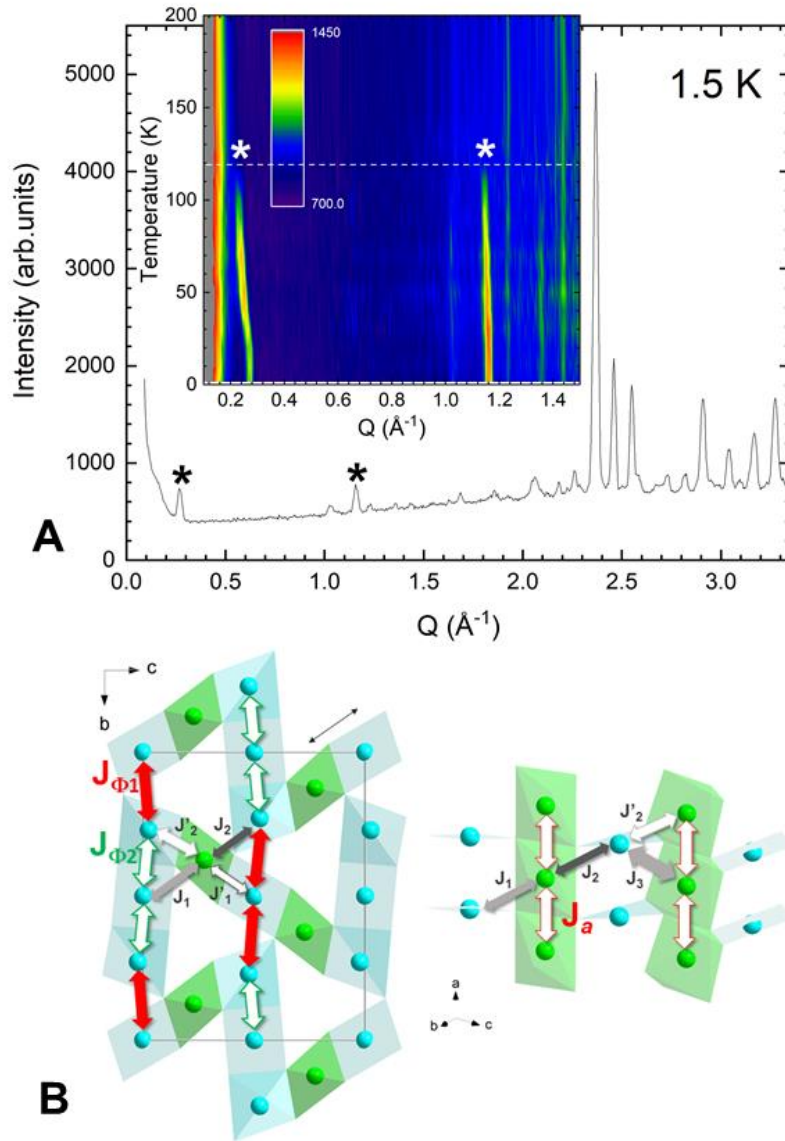


Figure 3. Cu_2CrBO_5 neutron diffraction data at 1.5 K (A), with the inset illustrating the corresponding temperature evolution. Asterisks indicate the incommensurate Bragg magnetic peaks appearing below $T_N = 120$ K. B schematizes the different magnetic exchange paths (as defined in Table S3). White (colored) arrows are for direct (super)-exchange configurations; intra-chains magnetic paths are in green or red, inter-sublattices exchange paths in dark/light grey. Cu/Cr atoms in blue/green for clarity. The thin arrow line on the left panel indicates the direction of the polyhedral block shown on the right.

The simultaneous occurrence of ferroelectricity and magnetic order designates copper ludwigite Cu_2CrBO_5 as a new member of the spin-driven multiferroic family. Compared with most known multiferroic compounds, Cu_2CrBO_5 exhibits in addition a high transition temperature of 120 K, amongst the highest known, but still behind tuned hexaferrites (295 K [32]), CuO (230 K [11]) or YBaCuFeO_5 (200 K [33], [34]). Along CuO , and non-oxide CuBr_2 [35] or Cu_2OCl_2 [36], Cu_2CrBO_5 confirms that low-dimensional divalent copper compounds have a potential to be high-temperature multiferroics. It is therefore of interest to rationalize our findings, using the extensive knowledge built on Cu(II) oxides of low dimensionality, since the discovery of high- T_c superconductivity in cuprates.

For more than 50 years, the Goodenough-Kanamori-Anderson (GKA) rules [37], [38] have been used to qualitatively relate magnetic exchange J with orbital overlap. In divalent copper oxide compounds, the NN Cu-Cu spin interaction J is thus expected to change from antiferromagnetic (AFM) to ferromagnetic (FM), when the angle Φ of the Cu-O-Cu bond goes from 180° to 90° . Such a correlation has indeed been shown in a series of low-dimensional Cu(II) compounds [39], [40], [41], including CuO [11], [42], and is recalled in Figure 4. In Cu_2CrBO_5 , Cu chains are made of edge-sharing trimer units (Φ_2 close to 90°), connected through corners by a Cu-O-Cu angle Φ_1 close to 120° (Figure 1a and Table S3). Using Figure 4, we can expect ferromagnetic Cu trimers (FM J_{Φ_2}), antiferromagnetically coupled through J_{Φ_1} (Figure 3b). J_{Φ_1} could be as high as 20 meV in this case, which is the right order of magnitude to account for $T_N = 120$ K (~ 10 meV).

For comparison with other copper ludwigites, two instances can be mentioned. In a very recent study on Cu_2MBO_5 ($M = \text{Ga}, \text{Al}$) [43], authors report for Cu_2GaBO_5 a complex non-collinear long-range magnetic order, with a propagation vector $\mathbf{k} = (0.45, 0, -0.7)$, below $T_N = 4.1$ K. Following the GKA rules for Cu-O-Cu bonds, they suggest that, since half of the TM sites (corresponding in this case to the Cu1 and Cr sites of Cu_2CrBO_5) are statistically occupied by

1/3 of the Cu^{2+} , this disordered sublattice of Cu orphan spins provides only a very weak FM coupling between corner-sharing AFM units. Cu_2MnBO_5 [24], in which a $\mathbf{k} = 0$ magnetic order below $T_N = 92$ K was found, is built on FM trimer units roughly coupled AFM. The overall scenario of a FM J_{Φ_2} and AFM J_{Φ_1} seems to hold in this case, but this resemblance has to be taken with caution, as Cu/Mn disorder, along with the Jahn-Teller nature of Mn^{3+} , will obviously alter magnetic exchanges.

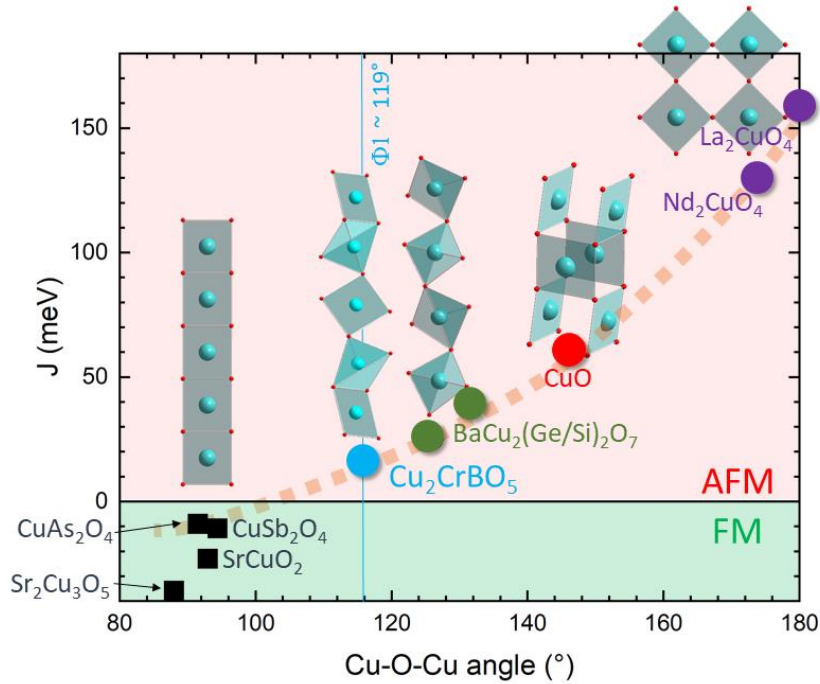


Figure 4. Evolution of the main magnetic exchange interaction J as a function of the Cu-O-Cu bond angle Φ in some cuprates, including Cu(II) oxyborate Cu_2CrBO_5 (drawn from data presented in [39], [41], [42], [44], obtained either from susceptibility measurements, neutron data or theoretical calculations). The dotted line is a guide to the eyes. The value of the predominant exchange J in Cu_2CrBO_5 can be estimated to be ~ 20 meV, from the Φ_1 value obtained from SXRPD refinements (see Figure 1a and Table S3).

The role of Cr chains becomes crucial because of the complex inter-sublattice pathways exemplified earlier. It is not possible to establish a hierarchy of exchange at this stage, but the modulated magnetic propagation vector tells us that one scenario is for J_2 and J'_1 to be both AF, or J'_2 and J_1 to be of opposite signs. Based on the fact that J_2 and J_1 , and J'_2 and J'_1 have very

similar geometries, it is most likely that frustration arises from the phase shift between Cu chains. It is a similar scenario, based on two magnetic sublattices with strong intra-sublattice interactions and weakly frustrated inter-sublattice interactions, which is invoked to explain spin frustration in CuO [45]. Arguably, magnetic anisotropy could also be an important ingredient in the stabilization of the spiral ground state. To compare with Cu₂MnBO₅, one could argue that in Cu₂CrBO₅ the Heisenberg character of the Cr spins favors incommensurate orderings, an effect that disappears when axial magnetic anisotropy directs the moment, such as in a Jahn-Teller ion like Mn³⁺.

Conclusion

Our results emphasize that the ludwigite network of copper oxyborates displays the right balance of low dimension, magnetic frustration, and strong magnetic exchange, for high temperature multiferroicity. Cu₂CrBO₅ is the first member of this family to be evidenced, but other heterometallic Cu(II) compounds remain to be investigated, especially if the disorder issue can be controlled. Density functional calculations should prove a really useful tool to further explore this system, in particular to better apprehend the importance of the inter-sublattice coupling, and understand how it can be used to further optimize its multiferroic properties.

AUTHOR CONTRIBUTIONS

Sample synthesis and structural characterizations by C.M. and J.S., SXRPD experiment by J.S. and F.F., physical properties measurements by C.M. and A.M., NPD experiments by E.S. and F.D. Manuscript writing by F.D. and J.S., with constant feedback from all co-authors.

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SUPPLEMENTARY MATERIAL

See supplementary material for Rietveld refinement of the Cu_2CrBO_5 SXRPD data at 260 K and NPD data at RT ; Magnetic field evolution of the magnetization and first field derivative of the magnetization of Cu_2CrBO_5 at different temperatures ; LeBail profile of the Cu_2CrBO_5 NPD data at 1.5 K, and results of the SXRPD Rietveld refinement of Cu_2CrBO_5 at 260 K (including M-O bond lengths and selected atomic distances and angles).

DATA AVAILABILITY STATEMENT

The data that supports the findings of this study are available within the article and its supplementary material.

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Supplementary Material

High temperature spin-driven multiferroicity in ludwigite chromocuprate Cu_2CrBO_5

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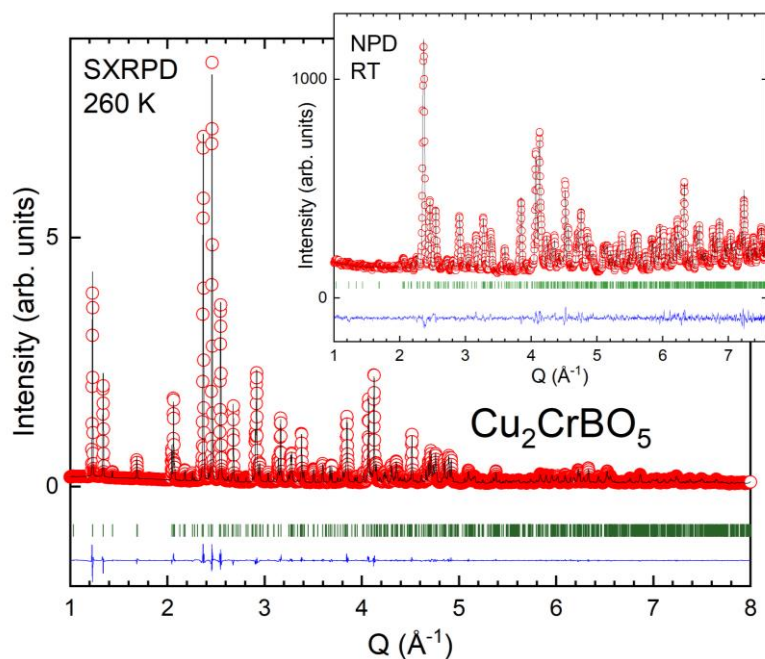


Figure S1. Rietveld refinement of the Cu_2CrBO_5 SXRPD data ($\lambda = 0.44296 \text{ \AA}$) at 260 K and NPD data ($\lambda = 1.5946 \text{ \AA}$) at RT (experimental data: open red circles, calculated profile: continuous line, allowed Bragg reflections: vertical green marks. The difference between the experimental and calculated profiles is displayed at the bottom of each graph as a blue line).

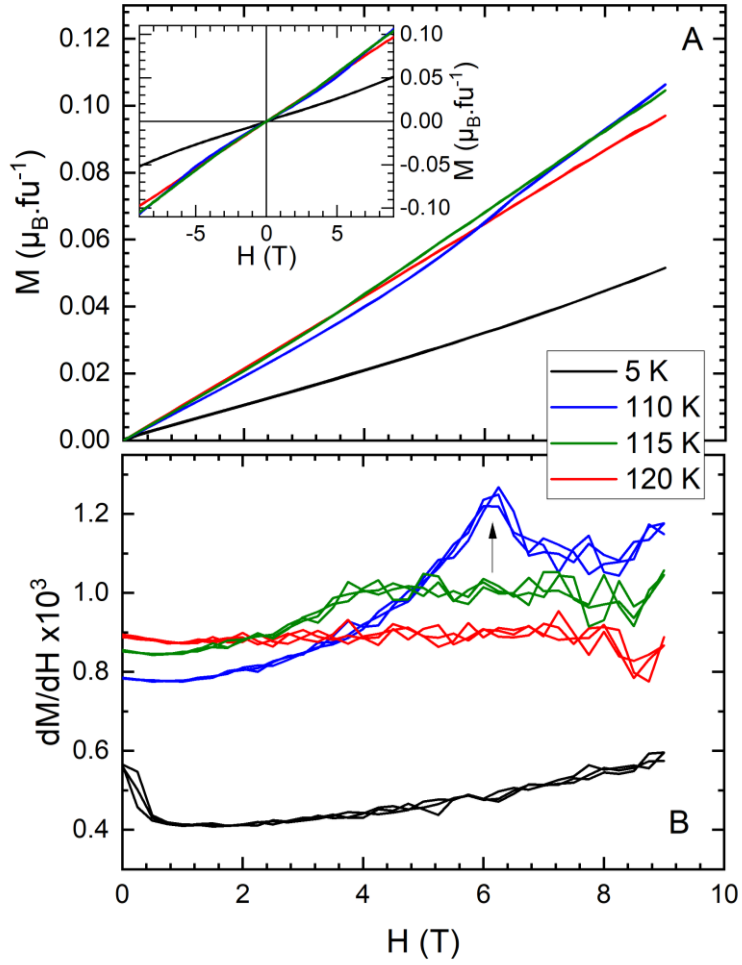


Figure S2. Magnetic field evolution (**A**) and corresponding first magnetic field derivative (**B**) of the magnetization M of Cu_2CrBO_5 at different temperatures (5 K, 110 K, 115 K and 120 K).

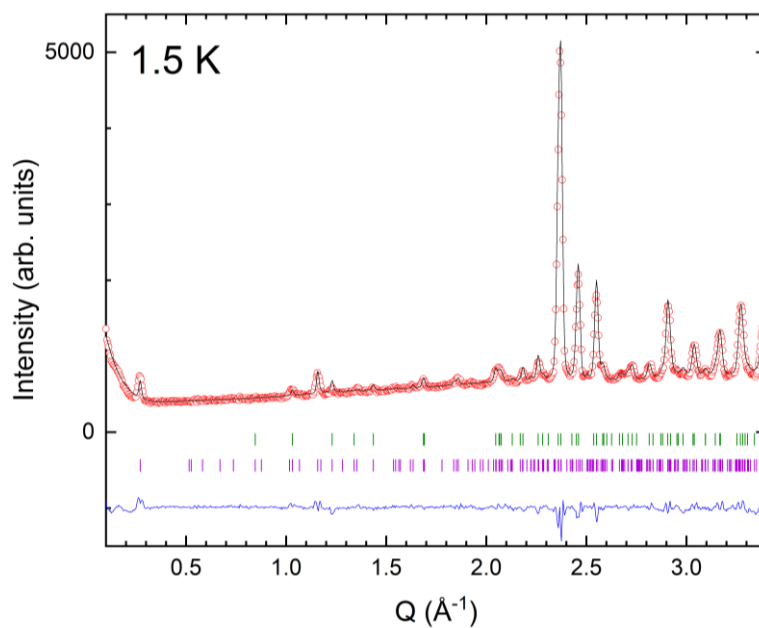


Figure S3. LeBail profile of the Cu_2CrBO_5 NPD data ($\lambda = 2.426 \text{ \AA}$) at 1.5 K. The chosen propagation vector in that case is $\mathbf{k} = (0.096(3), 0, 0.25)$ (experimental data: open red circles, calculated profile: continuous line, allowed Bragg (magnetic Bragg) reflections: vertical green (purple) ticks. The difference between the experimental and calculated profiles is displayed at the bottom of each graph as a blue line).

Table S1. Results of the SXRPD Rietveld refinement of Cu_2CrBO_5 at 260 K (SG: $P2_1/c$, $a = 3.05433(3)$ Å, $b = 12.1788(1)$ Å, $c = 9.4129(1)$ Å and $\beta = 94.5605(6)^\circ$. Cell volume $349.032(6)$ Å³. $R_{\text{Bragg}} = 1.8$ %; $\chi^2 = 17.6$). Atom labelling corresponds to that of Figure 1a. Thermal displacement parameters B_{iso} were constrained to be equal for all oxygen atoms. Anisotropic strain (strain_model = 2 [31]) was introduced in the refinement.

Atom	Site	x	y	z	B_{iso} (Å ²)
Cu1	$2c$	0	0.5	0	0.28(2)
Cu2	$2a$	0	0.5	0.5	0.20(2)
Cu3	$4e$	0.9979(3)	0.2709(5)	0.0244(7)	0.20(2)
Cr	$4e$	0.4959(4)	0.3750(1)	0.2738(1)	0.15(2)
B	$4e$	0.526(3)	0.3744(9)	0.7459(8)	0.4(1)
O1	$4e$	0.015(1)	0.1453(3)	0.9095(4)	0.03(3)
O2	$4e$	0.975(1)	0.3966(3)	0.1438(3)	0.03(3)
O3	$4e$	0.432(2)	0.2232(3)	0.1853(4)	0.03(3)
O4	$4e$	0.716(1)	0.3818(4)	0.8819(3)	0.03(3)
O5	$4e$	0.458(2)	0.0312(3)	0.1650(4)	0.03(3)

Table S2. M-O bond lengths in the TMO_6 polyhedra of Cu_2CrBO_5 at 260 K (from SXRPD refinement results). In italic are shown bond lengths > 2.35 Å (not taken into account in the average calculation).

	Cu1 ($2c$)	Cu2 ($2a$)	Cu3 ($4e$)	Cr ($4e$)
TM-O	O2:1.855(3) x2	O1:1.966(4) x2	O1:1.877(3)	O2:1.943(4)
distance (Å)	O4:1.977(4) x2	O5:2.042(4) x2	O2:1.904(3)	O1:1.972(4)
	<i>O4:2.911(4) x2</i>	<i>O5:2.388(5) x2</i>	O3:2.018(4)	O5:1.990(4)
			O4:2.045(4)	O2:1.998(4)
			<i>O3:2.455(5)</i>	O3:2.029(4)
			<i>O4:2.983(4)</i>	O1:2.036(4)
Average (Å)	1.916(4)	2.004(4)	1.961(4)	1.995(4)

Table S3. Selected atomic distances and angles in Cu_2CrBO_5 at 260 K (from SXRPD refinement results). Grey background indicates edge-sharing octahedra/direct-exchange paths.

Atom 1	Atom 2	Distance (Å)	Angle (°)	Angle (°)	J
Cu1	Cu3	2.797(1)	Cu1-O2-Cu3 : 96.3(2)	Cu1-O4-Cu3 : 88.2(2)	$J_{\Phi 2}$
Cu2	Cr	2.950(1)	Cu2-O1-Cr : 97.0(2)	Cu2-O5-Cr : 94.1(2)	J'_1
Cu3	Cr	2.977(1)	Cu3-O2-Cr : 101.3(2)	Cu3-O3-Cr : 94.7(1)	J'_2
Cr	Cr	3.054(2)	Cr-O1-Cr : 99.3(2)	Cr-O2-Cr : 101.5(2)	J_a
Cu1	Cr	3.261(1)	Cu1-O2-Cr : 118.2(2)		J_1
Cu2	Cu3	3.307(1)	Cu2-O1-Cu3 : 118.8(2)		$J_{\Phi 2}$
Cu2	Cr	3.108(1)	Cu2-O1-Cr : 101.9(2)		J_4
Cu3	Cr	3.169(1)	Cu3-O2-Cr : 108.6(2)		J_3
Cu3	Cr	3.240(1)	Cu3-O1-Cr : 114.7(2)		J_2