



HAL
open science

Decay and absorption for the Vlasov-Navier-Stokes system with gravity in a half-space

Lucas Ertzbischoff

► **To cite this version:**

Lucas Ertzbischoff. Decay and absorption for the Vlasov-Navier-Stokes system with gravity in a half-space. 2021. hal-03281075

HAL Id: hal-03281075

<https://hal.science/hal-03281075v1>

Preprint submitted on 7 Jul 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Decay and absorption for the Vlasov-Navier-Stokes system with gravity in a half-space

Lucas Ertzbischoff *

July 5, 2021

Abstract

This paper is devoted to the large time behavior of weak solutions to the three-dimensional Vlasov-Navier-Stokes system set on the half-space, with an external gravity force. This fluid-kinetic coupling arises in the modeling of sedimentation phenomena. We prove that the local density of the particles and the fluid velocity enjoy a convergence to 0 in large time and at a polynomial rate. In order to overcome the effect of the gravity, we rely on a fine analysis of the absorption phenomenon at the boundary. We obtain a family of decay estimates for the moments of the kinetic distribution, provided that the initial distribution function has a sufficient decay in the phase space.

Contents

1	Introduction	2
2	Main result	7
3	Conditional large time behavior of the fluid velocity	10
4	Preliminaries for the bootstrap procedure	11
4.1	Characteristic curves for the Vlasov equation: representation formula and change of variable . . .	12
4.2	Local in time estimates	14
4.3	Strong existence times and higher order energy estimates	20
5	Exit geometric condition and absorption	22
6	The bootstrap argument	26
6.1	Initialization of the bootstrap procedure	26
6.2	Absorption and decay in time of the moments	28
6.3	Weighted in time estimates	34
6.4	End of the proof of Theorem 2.1	37
A	Appendix	39
A.1	DiPerna-Lions theory in $\mathbb{R}_+^3 \times \mathbb{R}^3$	39
A.2	The Cauchy problem for the Vlasov-Navier-Stokes system	40
A.3	Gagliardo-Nirenberg-Sobolev inequality on \mathbb{R}_+^3	42
A.4	Maximal L^pL^q regularity for the Stokes system on \mathbb{R}_+^3	43
A.5	Conditional decay of the energy: proof of Theorem 3.1	44
A.6	Parabolic regularization for the Navier-Stokes system on \mathbb{R}_+^3	47
	References	50

*Centre de Mathématiques Laurent Schwartz (UMR 7640), Ecole polytechnique, Institut Polytechnique de Paris, 91128 Palaiseau Cedex, France (lucas.ertzbischoff@polytechnique.edu)

1 Introduction

Fluid-kinetic systems aim at modelling the collective motion of a dispersed phase of small particles immersed within a fluid. In such systems, also called spray models, the dispersed phase is described at a mesoscopic level by a distribution function solving a kinetic equation while the evolution of macroscopic quantities for the fluid is governed by fluid mechanics equations.

Among the wide family of fluid-kinetic systems (see the pioneering works [59, 54]), one can consider the so-called *thin spray* models where the particles volume fraction is small compared to that of the surrounding fluid. In this context, an interesting prototype is the incompressible Vlasov-Navier-Stokes system, coupling in a nonlinear way the fluid and kinetic equations through a drag term. The later one depends on the fluid unknowns and on the density function and allows for an exchange of momentum between the fluid and the particles. Beyond its mathematical interest, this system also appears in the study of the transport and deposition of a therapeutic aerosol in the airflows contained in the human upper airways (see [9]).

In this paper, we are interested in the following Vlasov-Navier-Stokes system set in $\mathbb{R}_+^3 \times \mathbb{R}^3$:

$$\partial_t f + v \cdot \nabla_x f + \operatorname{div}_v [(u - v)f + Gf] = 0, \quad (t, x, v) \in \mathbb{R}_+^* \times \mathbb{R}_+^3 \times \mathbb{R}^3, \quad (1.1)$$

$$\partial_t u + (u \cdot \nabla)u - \Delta u + \nabla p = \int_{\mathbb{R}^3} f(t, x, v)(v - u(t, x)) dv, \quad (t, x) \in \mathbb{R}_+^* \times \mathbb{R}_+^3, \quad (1.2)$$

$$\operatorname{div} u = 0, \quad (t, x) \in \mathbb{R}_+^* \times \mathbb{R}_+^3. \quad (1.3)$$

Here $\mathbb{R}_+^3 := \mathbb{R}^2 \times (0, +\infty)$ is the tridimensional half-space while $G := (0, 0, -g) \in \mathbb{R}^3$ is a given vector, with $g > 0$ the gravitational acceleration. In these equations, $u = u(t, x) \in \mathbb{R}^3$ and $p = p(t, x) \in \mathbb{R}$ stand for the velocity field and pressure of the fluid, while $f = f(t, x, v) \in \mathbb{R}^+$ is the distribution function of the particles in the phase space $\mathbb{R}_+^3 \times \mathbb{R}^3$. Here, the particles undergo the friction force produced by the surrounding fluid, as well as the effect of gravity. Thus, using the Stokes law, the resultant force exerted on the particles is the sum of the drag and weight/buoyancy, that is

$$u(t, x) - v + G,$$

and the Vlasov equation (1.1) is thus coupled to the Navier-Stokes equations (1.2)-(1.3). Note that we have implicitly considered spherical particles of same radius and that the mass density of the particles is greater than that of the fluid (because of the positive coefficient before the vector G - see e.g. [14]). Unlike the *thick spray* case (see e.g. [50, 22]), collisions between particles are neglected here, as well as the equation on their volume fraction.

A coupling term is also added in the Navier-Stokes equations (1.2)-(1.3), where a forcing term appears in the r.h.s and stems from the retroaction of the particles on the fluid. This source term is usually called the *Brinkman force* and can be rewritten as

$$\int_{\mathbb{R}^3} f(v - u) dv = j_f - \rho_f u, \quad (1.4)$$

where

$$\begin{aligned} \rho_f(t, x) &:= \int_{\mathbb{R}^3} f(t, x, v) dv, \\ j_f(t, x) &:= \int_{\mathbb{R}^3} v f(t, x, v) dv. \end{aligned}$$

Note that in the previous Navier-Stokes equations (1.2)-(1.3), the density and viscosity of the fluid are assumed to be constant and both chosen equal to 1, while the external gravity force $g = \nabla \Phi$, with $\Phi = gz$, has been absorbed in the pressure term. In short, the equations (1.1)-(1.2)-(1.3) account for the description of a cloud of fine particles sedimenting in an ambient incompressible viscous fluid.

The Vlasov-Navier-Stokes system has been studied in different directions over the past two decades.

The Cauchy problem. Concerning the existence theory of global weak solutions to the system (1.1)-(1.2)-(1.3) (without the gravity force), different settings and boundary conditions for the distribution function have

been addressed, depending on the spatial domain: a fixed bounded domain with specular reflexion in [1], the flat torus \mathbb{T}^3 in [8], time-dependent domain with absorption boundary condition in [10, 11], or a 2D rectangle with partly absorbing boundary condition in [28]. Local strong solutions can also be considered as in [18] (for inhomogeneous fluid equations), as well as blow-up in finite time of classical solutions in [17].

Note that the additional term involving the gravity has only been taken into account for the Vlasov-Stokes system on bounded domain with specular reflexion in [33], or for the Vlasov equation coupled to the stationary Stokes system with a regular and compactly supported initial distribution function on \mathbb{R}^3 in [45].

The 2D case. In two dimensions, more results are available for this fluid-kinetic system, essentially because of the study of Navier-Stokes system which is more favorable in this context: for instance, uniqueness of 2D-global weak solutions is proven for the whole space or the torus case in [36]. The controllability of the system is also explored in dimension 2 in [53].

Link to other models. In the spirit of Hilbert’s 6th problem of axiomatization of physics, fluid-kinetic models can be linked to other systems of ODEs and PDEs. Deriving rigorously the Vlasov-Navier-Stokes system from "first laws" appears as an important issue, but remains essentially an outstanding open problem for the whole system. Two main strategies have been proposed so far. The *mean-field* limit of a N-solid particle system coupled with a fluid equation has been considered in [23, 42, 43, 13] and allows one to recover the Brinkman force in a quasi-static framework thanks to homogenization techniques. Some partial results are also known for a dynamical but only macroscopic equation for the particles, in some dilute regime where they have no inertia (see [44, 51]). Another direction has been taken in [4, 5] where, in a formal way, the coupling between two mixtures provides the derivation of the Vlasov-Navier-Stokes system in the same fashion as the hydrodynamic limits of the Boltzmann equation [3, 2]. However, a full and rigorous justification of this program based on kinetic theory is still open.

Through *hydrodynamic limits* of the Vlasov-Navier-Stokes system, one can also seek to derive some systems involving only averaged quantities: more precisely, high friction regimes of the system have been shown to lead to Navier-Stokes type systems. These asymptotic regimes have been first considered in [30, 31] for the Vlasov-Fokker-Planck-Navier-Stokes equations, where the effect of Brownian motion is added in the equation of the distribution function.

Without diffusion in velocity in the Vlasov equation, one of these limits has been handled in [45] for the Vlasov-(steady)Stokes system with gravity in the whole space. Very recently, the question raised by these different regimes has been addressed in [35] for the full system (1.1)-(1.2)-(1.3) on the torus (without the gravity force).

Large time behavior. The *large time dynamics* of the Vlasov-Navier-Stokes system, which is the main issue of this article, has very recently received a particular attention. This natural question is studied for the first time by Jabin in [46] for a reduced kinetic model. In the absence of dissipative mechanism in the Vlasov equation (like a Fokker-Planck operator allowing to consider smooth equilibria, see [29]), the sole effect of the drag force in the system should lead to nontrivial equilibria which are singular. More precisely, one expects a monokinetic behavior for the distribution function of the particles (that is to say, a convergence towards a Dirac mass in velocity). A conditional result accounting for this phenomenon in the Vlasov-Navier-Stokes system has been provided by Choi and Kwon in [18]. In short, it requires a global bound in time which is not *a priori* satisfied by global weak solutions to the system. More recently, this extra assumption has been removed for small initial data: the first complete result stems from the article [37] of Han-Kwan, Moussa and Moyano where the authors work in a periodic setting and in a framework *à la* Fujita-Kato. In the same spirit, such a monokinetic behavior of weak solutions has been obtained for the whole space case by Han-Kwan in [34] and then extended to the case of bounded domains with absorption boundary conditions in [24]. In short, these results are all based on a remarkable energy-dissipation inequality satisfied by weak solutions to the system.

Unlike this series of works, the existence and stability of regular equilibria has been obtained by Glass, Han-Kwan and Moussa in [28] for a particular 2D bounded domain with partly absorbing and injection boundary conditions. Finally, let us emphasize the fact that the high friction limit tackled by Han-Kwan and Michel in [35] is closely linked to the monokinetic behavior we have mentioned just before, and in particular to the techniques used in [37, 34].

Main contribution of this paper. In the continuation of these previous works, the main goal of this article is the study of the large time dynamics of global weak solutions to the system (1.1)-(1.2)-(1.3). Its originality lies in dealing with a fluid-kinetic system on an unbounded domain with boundary, where one considers absorption boundary conditions for the distribution function, together with an additional gravity force term. Loosely speaking, the presence of a gravity force may ruin the decay of the energy of the system. At first sight, this prevents the use of exactly the same techniques as in [37, 34, 24]. However, it is actually possible to take advantage of the absorption at the boundary to analyse the large time behavior of global weak solutions starting close to equilibrium.

Before going further, we give several definitions and set notations about the system that we will consider in this article. In what follows, we will sometimes refer to (1.1)-(1.2)-(1.3) as the *VNS system*. Along this paper, we will make a constant use of the notation

$$\mathbb{R}_+^3 := \mathbb{R}^2 \times (0, +\infty).$$

First, the VNS system is supplemented with the following initial conditions for u and f :

$$u|_{t=0} = u_0 \text{ in } \mathbb{R}_+^3, \quad (1.5)$$

$$f|_{t=0} = f_0 \text{ in } \mathbb{R}_+^3 \times \mathbb{R}^3. \quad (1.6)$$

We prescribe the following Dirichlet boundary conditions for the fluid:

$$u(t, \cdot) = 0, \text{ on } \partial\mathbb{R}_+^3 = \mathbb{R}^2 \times \{0\}. \quad (1.7)$$

We also need to introduce the following outgoing/incoming phase-space boundaries:

$$\Sigma^\pm := \{(x, v) \in \partial\mathbb{R}_+^3 \times \mathbb{R}^3 \mid \pm v \cdot n(x) > 0\}, \quad (1.8)$$

$$\Sigma_0 := \{(x, v) \in \partial\mathbb{R}_+^3 \times \mathbb{R}^3 \mid v \cdot n(x) = 0\}, \quad (1.9)$$

$$\Sigma := \Sigma^+ \sqcup \Sigma^- \sqcup \Sigma_0 = \partial\mathbb{R}_+^3 \times \mathbb{R}^3, \quad (1.10)$$

where $n(x)$ stands for the normal vector to the boundary $\partial\mathbb{R}_+^3$ at point x . We observe that

$$\Sigma^\pm = \{(x, v) \in \mathbb{R}^3 \times \mathbb{R}^3 \mid x_3 = 0, \pm v_3 < 0\},$$

$$\Sigma_0 = \{(x, v) \in \mathbb{R}^3 \times \mathbb{R}^3 \mid x_3 = 0, v_3 = 0\}.$$

Then, we prescribe the following absorption boundary conditions for the distribution function:

$$f(t, \cdot, \cdot) = 0, \text{ on } \Sigma^-. \quad (1.11)$$

Several functionals play an important role in the study of the VNS system. We introduce the following ones.

Definition 1.1. 1. The **kinetic energy** of the Vlasov-Navier-Stokes system is defined for all $t \geq 0$ as:

$$\mathbf{E}(t) := \frac{1}{2} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v|^2 dx dv. \quad (1.12)$$

2. The **dissipation** of the Vlasov-Navier-Stokes system (without gravity) is defined for all $t \geq 0$ as:

$$\mathbf{D}(t) := \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |u(t, x) - v|^2 dx dv + \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2. \quad (1.13)$$

3. The **dissipation with gravity** of the Vlasov-Navier-Stokes system is defined for all $t \geq 0$ as:

$$\mathbf{D}_G(t) := \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |u(t, x) - v|^2 dx dv + \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 - \int_{\mathbb{R}_+^3} G \cdot j_f(t, x) dx. \quad (1.14)$$

At a formal level, the VNS system enjoys an energy-dissipation structure involving the previous functionals. More precisely, smooth solutions to the system satisfy the following *a priori* estimate

$$\frac{d}{dt}E(t) + D_G(t) \leq 0.$$

Thus, there is a variation of energy coming from the dissipation inside of the fluid and because of the friction between the particles and the fluid, but the gravity force induces an additional term leading to a potential non-decay of the kinetic energy E . Indeed, we expect the term $G \cdot j_f > 0$ to be positive at least after some time because the particles should ultimately fall in the same direction than the gravity.

We denote by $\mathcal{D}_{\text{div}}(\mathbb{R}_+^3)$ the set of smooth \mathbb{R}^3 valued divergence free vector-fields having compact support in \mathbb{R}_+^3 . The closures of $\mathcal{D}_{\text{div}}(\mathbb{R}_+^3)$ in $L^2(\mathbb{R}_+^3)$ and in $H^1(\mathbb{R}_+^3)$ are respectively denoted by $L_{\text{div}}^2(\mathbb{R}_+^3)$ and by $H_{0,\text{div}}^1(\mathbb{R}_+^3)$. We write $H_{\text{div}}^{-1}(\mathbb{R}_+^3)$ for the dual of the later.

We now define the class of admissible initial data for the VNS system.

Definition 1.2 (Initial condition). We shall say that a couple (u_0, f_0) is an admissible initial condition if:

$$u_0 \in L_{\text{div}}^2(\mathbb{R}_+^3), \quad (1.15)$$

$$f_0 \in L^1 \cap L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3), \quad (1.16)$$

$$f_0 \geq 0, \quad \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f_0 \, dx \, dv = 1, \quad (1.17)$$

$$(x, v) \mapsto f_0(x, v)|v|^2 \in L^1(\mathbb{R}_+^3 \times \mathbb{R}^3). \quad (1.18)$$

We then introduce some notations about the moments of any phase-space distribution function g .

Definition 1.3. For any $\alpha \geq 0$, and any measurable function $g : \mathbb{R}^+ \times \mathbb{R}_+^3 \times \mathbb{R}^3 \rightarrow \mathbb{R}^+$, we set

$$m_\alpha g(t, x) := \int_{\mathbb{R}^3} |v|^\alpha g(t, x, v) \, dv,$$

$$M_\alpha g(t, x) := \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} |v|^\alpha g(t, x, v) \, dv \, dx = \int_{\mathbb{R}_+^3} m_\alpha g(t, x) \, dx.$$

In our approach, we shall rely on some decay assumptions satisfied by the initial distribution function f_0 . We thus introduce the following quantities.

Definition 1.4. For any $q > 0$, $m > 0$ and $r \geq 1$, we set

$$N_q(f_0) := \|(1 + |v|^q)f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}, \quad (1.19)$$

$$K_{q,r}(f_0) := \left\| (1 + |v|^q) \|f_0(\cdot, v)\|_{L_x^r(\mathbb{R}_+^3)} \right\|_{L_v^\infty(\mathbb{R}^3)}, \quad (1.20)$$

$$H_{q,m}(f_0) := \left\| (1 + |v|^q) \|(1 + x_3^m)f_0(\cdot, v)\|_{L_x^\infty(\mathbb{R}_+^3)} \right\|_{L_v^1(\mathbb{R}^3)}, \quad (1.21)$$

$$F_{q,m,r}(f_0) := \left\| (1 + |v|^q) \|(1 + x_3^m)f_0(\cdot, v)\|_{L_x^r(\mathbb{R}_+^3)} \right\|_{L_v^1(\mathbb{R}^3)}. \quad (1.22)$$

We will consider weak solutions to the Vlasov equation with gravity force (1.1), with the boundary conditions (1.11) and the previous initial condition, which are defined as follows.

Definition 1.5 (Weak solutions to the Vlasov equation). Let $U \in L_{\text{loc}}^1(\mathbb{R}^+ \times \mathbb{R}_+^3)$. Consider also an initial distribution f_0 satisfying (1.16)-(1.18). We say that a nonnegative function $f \in L_{\text{loc}}^\infty(\mathbb{R}^+; L^1 \cap L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3))$ is a weak solution to the Vlasov equation (1.1) with force field U , with boundary condition (1.11) and with initial condition f_0 if, for all $\Psi \in \mathcal{D}([0, T] \times \mathbb{R}_+^3 \times \mathbb{R}^3)$ with $\Psi(T, \cdot) = 0$ and vanishing on $\mathbb{R}^+ \times (\Sigma^+ \cup \Sigma_0)$, one has

$$\int_0^T \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) [\partial_t \Psi + v \cdot \nabla_x \Psi + (U - v) \cdot \nabla_v \Psi](t, x, v) \, dx \, dv \, dt = - \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f_0(x, v) \Psi(0, x, v) \, dx \, dv.$$

Weak solutions to the Vlasov equations enter in the framework of the DiPerna-Lions theory for transport equations (in the phase space $\mathbb{R}_+^3 \times \mathbb{R}^3$). We refer to Section A.1 in the Appendix for more details, where we recall in particular the classical stability property of renormalized solutions which we will constantly use throughout this article.

We will also consider Leray solutions for the full Vlasov-Navier-Stokes with the boundary conditions (1.7)-(1.11) and the initial conditions described in (1.2), in the following sense.

Definition 1.6 (Leray solutions with strong energy inequality for the VNS system). Consider an admissible initial condition (u_0, f_0) in the sense of Definition 1.2. A global weak solution to the Vlasov-Navier-Stokes system with boundary condition (1.7)-(1.11) and with initial condition (u_0, f_0) is a pair (u, f) such that:

$$u \in L_{\text{loc}}^\infty(\mathbb{R}^+; L_{\text{div}}^2(\mathbb{R}_+^3)) \cap L_{\text{loc}}^2(\mathbb{R}^+; H_{0,\text{div}}^1(\mathbb{R}_+^3)), \quad (1.23)$$

$$\text{div}_x u = 0, \quad (1.24)$$

$$f \in L_{\text{loc}}^\infty(\mathbb{R}^+; L^1 \cap L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)), \quad (1.25)$$

$$j_f - \rho_f u \in L_{\text{loc}}^2(\mathbb{R}^+; H_{\text{div}}^{-1}(\mathbb{R}_+^3)), \quad (1.26)$$

and such that the following holds. The distribution function f is a weak solution to the Vlasov equation with force field $G + u$ with initial condition f_0 in the sense of Definition 1.5 and the velocity field u is a Leray solution to the Navier-Stokes equations with initial condition u_0 , that is for all $\Phi \in \mathcal{C}^1([0, T]; \mathcal{D}_{\text{div}}(\mathbb{R}_+^3))$ such that $\Phi(T, \cdot) = 0$, we have

$$\begin{aligned} & \int_0^T \int_{\mathbb{R}_+^3} [u \cdot \partial_t \Phi + (u \otimes u) : \nabla_x \Phi - \nabla_x u : \nabla_x \Phi](t, x) dx dt \\ &= - \int_{\mathbb{R}_+^3} u_0(x) \cdot \Phi(0, x) dx - \int_0^T \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) [v - u(t, x)] \cdot \Phi(t, x) dx dv dt, \end{aligned} \quad (1.27)$$

and the strong energy inequality holds for the Navier-Stokes equations: for any $t \geq 0$ and almost every $0 \leq s \leq t$ (including $s = 0$)

$$\|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + 2 \int_s^t \|\nabla u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 d\tau \leq \|u(s)\|_{L^2(\mathbb{R}_+^3)}^2 + 2 \int_s^t \int_{\mathbb{R}_+^3} (j_f(\tau, x) - \rho_f u(\tau, x)) \cdot u(\tau, x) dx d\tau. \quad (1.28)$$

Furthermore, the following energy estimate holds for the Vlasov-Navier-Stokes system: for any $t \geq 0$ and almost every $0 \leq s \leq t$ (including $s = 0$)

$$E(t) + \int_s^t D(\tau) d\tau \leq E(s) + \int_s^t \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} G \cdot v f(\tau, x, v) dx dv d\tau, \quad (1.29)$$

where the energy E and dissipation D have been defined in (1.14).

Note that the last integral in the r.h.s of the inequality (1.28) actually makes sense because of Sobolev embedding and the fact that $j_f - \rho_f u \in L_{\text{loc}}^2(\mathbb{R}^+; L^{6/5}(\mathbb{R}_+^3))$ (see Section A.2 in the Appendix).

Such global weak solutions can be obtained through an approximation procedure which seems to be classical by now (see e.g. [10, 28, 11]). Since the half-space/gravity framework has not been explicitly treated in the former literature, we provide some rather sketchy elements of proof about the Cauchy problem in Section A.2, with a particular insight on the obtention of the strong energy inequalities (1.28) and (1.29).

In Section 2, we present the main result obtained in this paper. As will be explained later on, the approach that we will use to prove this result shares some similar features with the ones introduced in [37, 34, 24, 28]. We will detail the strategy which has been set up in these works and in comparison, we will describe the method we need in our case.

2 Main result

The notations we use here are introduced in Definition 1.4 and in Section A.4. If $p, s > 0$ are given, we set

$$\begin{aligned}\mathcal{E}_{p,s} &:= W_0^{1,\frac{9}{7}}(\mathbb{R}_+^3) \cap D_3^{\frac{1}{2},2}(\mathbb{R}_+^3) \cap D_3^{1-\frac{1}{s},s}(\mathbb{R}_+^3) \cap D_p^{1-\frac{1}{p},p}(\mathbb{R}_+^3), \\ \mathcal{E}(0) &:= \|u_0\|_{H^1 \cap D_p^{1-\frac{1}{p},p}(\mathbb{R}_+^3)} + E(0) + \|u_0\|_{L^1(\mathbb{R}_+^3)} + M_6 f_0,\end{aligned}$$

and for $q, m > 0$

$$\mathcal{N}_{q,m}(f_0) := N_q(f_0) + H_{q,m}(f_0) + \max_{r \in \{1,3\}} \{K_{m,r}(f_0) + F_{q,m,r}(f_0)\}.$$

The main result of this paper reads as follows.

Theorem 2.1. *There exists $\varepsilon > 0$, $s \in (2, 3)$, $p \in (3, 3 + \varepsilon)$, $q > 3$ and $m > 0$ such that the following holds. Let (u_0, f_0) be an admissible initial condition in the sense of Definition 1.2 satisfying*

$$\begin{aligned}u_0 &\in H_{0,\text{div}}^1(\mathbb{R}_+^3) \cap L^1(\mathbb{R}_+^3) \cap \mathcal{E}_{p,s}, \\ M_6 f_0 &< \infty, \\ \mathcal{N}_{q,m}(f_0) &< \infty\end{aligned}\tag{2.1}$$

There exists a constant $C_0 > 0$ and an onto nondecreasing function $\Phi_{C_0} : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ such that the following holds. If the initial data are small enough in the sense that

$$\Phi_{C_0} \left(1 + \mathcal{E}(0) + \mathcal{N}_{q,m}(f_0) \right) \times \left(\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 + N_q(f_0) + H_{0,m}(f_0) \right) < C_0,\tag{2.2}$$

then there exists a continuous nonnegative function Ψ cancelling at 0 such that any global Leray solution (u, f) to the Vlasov-Navier-Stokes system with initial data (u_0, f_0) (in the sense of Definition 1.6) satisfies

$$\|u(t)\|_{L^2(\mathbb{R}_+^3)} \leq \frac{\Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right)}{(1+t)^{3/4}}.$$

Furthermore, if $k_1 \in \mathbb{R}^+$ with $q > k_1 + 3$ and $k_2 \in \mathbb{R}^+$, then for all $t \geq 0$

$$\begin{aligned}\rho_f(t, x) &\lesssim \frac{N_q(f_0)}{(1+t)^{k_1}} + \frac{H_{0,k_2}(f_0)}{(1+t)^{k_2}}, \\ \|\rho_f(t)\|_{L^r(\mathbb{R}_+^3)} &\lesssim \frac{K_{q,r}(f_0)}{(1+t)^{k_1}} + \frac{F_{0,k_2,r}(f_0)}{(1+t)^{k_2}}.\end{aligned}$$

Here, \lesssim only depends on k_1, k_2, q and g .

Remark 2.2. One can actually replace the quantity ρ_f in the end of the theorem by any moment $m_\ell f$, with $q > k + \ell + 3$ and an estimate involving $H_{\ell,k_2}(f_0)$ and $F_{\ell,k_2,r}(f_0)$.

Remark 2.3. As a byproduct of our analysis, we will prove along the way that if f_0 is compactly supported in velocity and in the third direction in space, that is

$$\text{supp } f_0 \subset (\mathbb{R}^2 \times (0, L)) \times B(0, R),$$

for some finite $L, R, > 0$, then there exists a finite time $T = T(L, R, g)$ such that $f(t) \equiv 0$ for every $t > T$.

Remark 2.4. In view of the results of [37, 34], one could hope for a monokinetic behavior of the distribution function in large time: in other words, $f(t)$ should concentrate in velocity to a Dirac mass supported at g . However, because of the absorption at the boundary, the result of Theorem 2.1 shows in particular that $\|f(t)\|_{L^1(\mathbb{R}_+^3 \times \mathbb{R}^3)} \xrightarrow[t \rightarrow +\infty]{} 0$ so that the previous singular behavior does not occur.

Let us explain the main strategy that has been already devised and used to study the large time behavior of the VNS system in [37, 34, 24] (which are gravity-less cases). Roughly speaking, under some smallness assumption on the initial data that we will detail below, [37, 34, 24] have proven that the fluid velocity u homogenizes to a constant when $t \rightarrow +\infty$, while the distribution function converges towards a Dirac mass in velocity. The later weak convergence is in particular measured thanks to the 1-Wasserstein distance on the phase space. All of these works rely heavily on the decay of a well-chosen energy functional which essentially controls the convergence of u and f . The choice of such a functional may depend on the domain. As explained before, two principal spatial frameworks have been explored.

- The case of bounded domains: as already mentioned in the introduction, Han-Kwan, Moussa and Moyano have tackled the large time behavior of small data global weak solutions to the system set on the torus \mathbb{T}^3 in [37], while the case of bounded domains with absorption boundary condition has been studied in [24]. The main strategy is the following. In the torus, a crucial quantity which plays a crucial role in the large time dynamics is the so-called *modulated energy* introduced by Choi and Kwon in [18]. It is defined as

$$\mathcal{E}(t) := \frac{1}{2} \int_{\mathbb{T}^3 \times \mathbb{R}^3} f(t, x, v) |v - \langle j_f(t) \rangle|^2 dx dv + \frac{1}{2} \int_{\mathbb{T}^3} |u(t, x) - \langle u(t) \rangle|^2 dx + \frac{1}{4} |\langle j_f(t) \rangle - \langle u(t) \rangle|^2,$$

where $\langle \cdot \rangle$ stands for the spatial average on \mathbb{T}^3 . On a smooth bounded domain $\Omega \subset \mathbb{R}^3$, the key functional is the kinetic energy E_Ω itself, defined as in (1.12) where \mathbb{R}_+^3 is replaced by Ω . In both cases, the decay of the energies \mathcal{E} and E_Ω is based on the following formal energy-dissipation identities

$$\frac{d}{dt} \mathcal{E}(t) + D_{\mathbb{T}^3}(t) = 0, \quad \frac{d}{dt} E_\Omega(t) + D_\Omega(t) = 0, \quad (2.3)$$

where $D_{\mathbb{T}^3}$ or D_Ω are defined as in (1.14) where \mathbb{R}_+^3 is replaced by \mathbb{T}^3 or Ω . Under the assumption that the global bound $\rho_f \in L^\infty(\mathbb{R}^+; L^\infty(Q))$ holds, where $Q = \mathbb{T}^3$ or $Q = \Omega$, one can show that an exponential decay of the energy is satisfied, namely

$$\forall t \geq 0, \quad \mathcal{E}(t) \lesssim e^{-\lambda t} \mathcal{E}(0), \quad E_\Omega(t) \lesssim e^{-\lambda t} E_\Omega(0), \quad (2.4)$$

for some $\lambda > 0$. This mainly comes from Poincaré(-Wirtinger) inequality. Then, a straightening change of variable in velocity shows that a sufficient condition for obtaining the previous bound on ρ_f is an estimate on the Lipschitz seminorm of u , that is

$$\int_0^\infty \|\nabla u(s)\|_{L^\infty(Q)} ds \ll 1.$$

A bootstrap procedure has to be performed in order to ensure such a global control: the main idea is to interpolate the pointwise conditional decay (2.4) with higher order parabolic regularity estimates for the fluid velocity. In short, the previous approach requires a smallness assumption on the initial data of the type

$$E_Q(0) + \|u_0\|_{H^1(Q)} \ll 1,$$

and one can even replace the previous norm for u_0 by $\|u_0\|_{H^{1/2}(\mathbb{T}^3)}$ in the torus case. Furthermore, one can describe the structure of the final spatial density, which depends on the whole evolution of the system. Note that in the case of a bounded domain with absorption boundary condition for the kinetic distribution, an initial support of the distribution function containing only high enough velocities leads to total absorption in finite time (namely, $f = 0$ after some finite time).

- The case of the whole space: in this case studied in [34], Han-Kwan has shown that the crucial functional to consider is the kinetic energy $E_{\mathbb{R}^3}$ itself, defined as in (1.12) where \mathbb{R}_+^3 is replaced by \mathbb{R}^3 . This energy satisfies the same formal energy identity (2.3) as above. A bound of the type $\rho_f \in L^\infty(\mathbb{R}^+; L^\infty(\mathbb{R}^3))$ now provides a decay of the form

$$\forall t \geq 0, \quad E_{\mathbb{R}^3}(t) \leq \frac{\varphi_\alpha(E_{\mathbb{R}^3}(0) + \|u_0\|_{L^1(\mathbb{R}^3)})}{(1+t)^\alpha}, \quad \text{for all } \alpha \in]0, 3/2[,$$

for some function φ_α . Here, the polynomial decay is the best one can hope for because of the absence of a Poincaré inequality (with respect to the Lebesgue measure) on domains unbounded in any direction. In short, it corresponds to the decay of the solutions to the heat equation on the whole space, with almost the same rate. As before, the second step of the analysis is a bootstrap analysis whose aim is to obtain the control on ρ_f , relying on the same sufficient control of $\|\nabla u\|_{L^1(\mathbb{R}^+; L^\infty(\mathbb{R}^3))}$. Since the energy decay is only polynomial and since the Brinkman force is not decaying, the use of dissipation functionals of higher order is needed¹. It essentially leads to the study of weighted in time estimates for the second order derivatives (in space) of the fluid velocity. Here, this procedure is applied for initial data close to equilibrium, in the sense that

$$E_{\mathbb{R}^3}(0) + \|u_0\|_{H^1(\mathbb{R}^3)} + \|f_0\|_{L^1_v(\mathbb{R}^3; L^\infty_x(\mathbb{R}^3))} \ll 1.$$

Let us emphasize the fact that the case of a bounded domain with specific boundary conditions has been considered by Glass, Han-Kwan and Moussa in [28]. More precisely, the VNS system is set on a bidimensional rectangle $\Omega := (-L, L) \times (-1, 1)$: the fluid velocity satisfies a Dirichlet boundary condition corresponding to a Poiseuille flow while the distribution function obeys to mixed absorption/injection boundary conditions. Compared to the previous state of the art, this particular framework leads to a somewhat other type of asymptotic behavior. Indeed, relying on a key geometric control condition (the so-called *exit geometric condition*), it has been shown that one can get the existence and asymptotic stability of non-trivial smooth equilibria for the system.

Main strategy. Let us explain the main approach used in this article. As mentioned before, the study of the large time dynamics of the system (1.1)-(1.2)-(1.3) is in the same spirit than that of the previous works [37, 34, 24]. However, a main obstacle comes from the presence of the additional gravity force term in the Vlasov equation (1.1), which creates an extra term in the r.h.s of the energy inequality (1.29). This breaks one of the main structural tools of the analysis of the gravity-less case because it rules out the decay of the total kinetic energy E . We thus need to base our study upon an additional mechanism which is at stake in the system. What comes into play here is the absorption of the particles at the boundary, coming from the boundary condition (1.11) for the distribution function f . To understand the crucial role of this phenomenon, we will use the Lagrangian structure of the Vlasov equation. We shall define the characteristic curves (X, V) for the Vlasov equation as the solutions of the following differential system:

$$\begin{cases} \dot{X}(s; t, x, v) = V(s; t, x, v), \\ \dot{V}(s; t, x, v) = u(s, X(s; t, x, v)) + G - V(s; t, x, v), \end{cases}$$

with $(X(t; t, x, v), V(t; t, x, v)) = (x, v)$ and where u has been extended by 0 outside the half-space. Introducing the following backward exit time

$$\tau^-(t, x, v) := \inf \{s \leq t \mid \forall \sigma \in [s, t], X(\sigma; t, x, v) \in \mathbb{R}_+^3\},$$

the method of characteristics shall provide the following representation formula:

$$f(t, x, v) = e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} f_0(X(0; t, x, v), V(0; t, x, v)). \quad (2.5)$$

In view of this expression, a certain decay in time of the moments of f should be satisfied along the evolution of the system, provided that f_0 enjoys some decay in the phase space.

In order to take advantage of the absorption, we shall rely on an *exit geometric condition*, reminiscent of the work of Glass, Han-Kwan and Moussa in [28]. In short, we ask that all the characteristic curves starting from a compact set leave the half-space before a fixed time. The main idea to propagate this condition will be to compare the coupled Vlasov equation to the Vlasov equation without fluid velocity and only governed by the gravity. The characteristic curves (X^g, V^g) for this simplified equation are defined as the solutions of

$$\begin{cases} \dot{X}^g(s; t, x, v) = V^g(s; t, x, v), \\ \dot{V}^g(s; t, x, v) = G - V^g(s; t, x, v), \end{cases}$$

¹This family of identities has also found a powerful application in the study of the hydrodynamic limits of the VNS system in the torus performed by Han-Kwan and Michel in [35].

with $(X(t; t, x, v), V(t; t, x, v)) = (x, v)$. We will then show that the exit geometric condition holds for all times for the VNS system, by using the particular geometry of the domain and the simple form of (X^g, V^g) . This will essentially require a control of the form

$$\int_0^\infty \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds \ll 1. \quad (2.6)$$

By combining the previous absorption phenomenon with a decay of the initial distribution function f_0 itself, we shall be able to obtain decay in time estimates of the moments of f . The argument will be based on the representation formula (2.5) and on a change of variable in velocity, namely $v \mapsto V(0; t, x, v)$ (which was already used in [37, 34, 24]). This procedure will be allowed if we can ensure a control of the Lipschitz seminorm of the fluid velocity, that is

$$\int_0^\infty \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds \ll 1. \quad (2.7)$$

Thus, it turns out that the presence of the gravity force is eventually favorable for our purpose and leads to the decay estimates of the moments stated in Theorem 2.1.

Thanks to the specific form of the Brinkman force $j_f - \rho_f u$ in the Navier-Stokes equations, we shall obtain pointwise in time estimates in various norms for this source term. Hence, as a solution of the Navier-Stokes equations with a sufficiently decaying forcing term, the fluid velocity u will enjoy a polynomial convergence towards 0. This is essentially the result bearing on u in Theorem 2.1.

Therefore, our main guiding line will be the obtention of decay in times estimates for the moments of f . We will base our proof on a bootstrap argument, mainly directly taken from [34].

Outline of the paper. According to the previous strategy, let us describe how this paper is organised.

In view of the arguments above, it makes sense to first consider the Navier-Stokes system having a source term (i.e. the Brinkman force (1.4)) with pointwise decay in time. It turns out that this assumption falls within the scope of the work of Wiegner in [58] for the Navier-Stokes equations with a decaying source term on \mathbb{R}^3 . *Modulo* an adaptation to the half-space case, this entails a polynomial decay of the L^2 norm of the fluid velocity. This conditional Theorem 3.1 is contained in Section 3.

Before going further, we shall need to state some preliminaries gathered in Section 4. They are necessary for a local in time analysis as well as for the subsequent bootstrap argument. We obtain rough bounds on the moments of f , ensuring short time controls. We also derive some H^1 estimates for the fluid velocity and define the notion of *strong existence time* for the Navier-Stokes system.

As explained before, the absorption of the particles at the boundary will be the key effect leading to global decay in time for the moments of f . In Section 5, we introduce the aforementioned *exit geometric condition* and analyse its effect on the system. In particular, this enables us to track which proportion of the support of the initial distribution has disappeared from the system at any given time.

Section 6 is dedicated to the bootstrap argument, which aim is to reach the global controls (2.6) and (2.7). Along the bootstrap, the previous absorption phenomenon is shown to lead to the desired decay estimates of the moments. We first show that the Brinkman force satisfies the desired pointwise decay. Polynomial weighted in time estimates for the fluid velocity are then obtained. This will allow to close the bootstrap argument thanks to an interpolation procedure.

3 Conditional large time behavior of the fluid velocity

The main goal of this short section is to show some conditional results about the polynomial decay of the L^2 norm of any Leray solution to the Navier-Stokes equation with a source term F . This decay holds if F satisfies a conditional pointwise decay in $L^2(\mathbb{R}_+^3)$ and is somewhat imposed by the decay of the Stokes semigroup on \mathbb{R}_+^3 . The main results reads as follows.

Theorem 3.1. *Let $u_0 \in L_{\text{div}}^2(\mathbb{R}_+^3) \cap L^1(\mathbb{R}_+^3)$ and $F \in L_{\text{loc}}^2(\mathbb{R}^+; L^2(\mathbb{R}_+^3))$. Let u be a global Leray solution to the Navier-Stokes system with strong energy inequality, associated to the initial data u_0 and the source term*

F. Let $T > 0$ and assume that

$$\forall t \in [0, T], \quad \|F(t)\|_{L^2(\mathbb{R}_+^3)} \leq \frac{C}{(1+t)^{7/4}}, \quad (3.1)$$

for some constant $C > 0$ independent of T . Then there exists a continuous nonnegative function Ψ_C cancelling at 0 and independent of T (and increasing in its parameter C), such that

$$\forall t \in [0, T], \quad \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim \frac{\Psi_C \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right)}{(1+t)^{3/2}}. \quad (3.2)$$

Here, we will follow Wiegner [58] and Borchers and Miyakawa [6], relying on the Fourier splitting method of Schonbeck [56]. More precisely, assuming a decay of the type (3.1) for the source term F means that the desired result for the large time behavior of the fluid velocity enters in the framework of [58]. Since we work in an unbounded domain with boundaries, we shall adapt this method written in the whole space case, thanks to a spectral decomposition of the Stokes operator and the use of its fractional powers, as in [6]. This will entail a polynomial decay of the fluid velocity similar to the one without a source term in the equations, and whose rate is roughly speaking the same as that of the unsteady Stokes equations.

Coming back to the Vlasov-Navier-Stokes system, we shall consider the Brinkman force $j_f - \rho_f u$ as a fixed source term in the Navier-Stokes equations. This means that we shall use Theorem 3.1 with $F = j_f - \rho_f u$. Note that compared to [34] where the conditional decay of the energy E was related to the whole VNS system, our result concerns the decay of u as a solution of the Navier-Stokes equations only and is independent of the coupling with the Vlasov equation: we only use the strong energy inequality (1.28) for the Navier-Stokes system with a given source term (even if it may depend on u and f). Therefore, the assumption (3.1) on the decay of this source term makes the situation simpler. We will be able to prove that this strong decay does occur thanks to the absorption phenomenon along a bootstrap procedure in Section 6. Here, we do not require a bound of the type $\rho_f \in L^\infty(0, T; L^\infty(\mathbb{R}_+^3))$ as in [34] and the rate of convergence is slightly better.

The combination of [58] and [6] for the proof of Theorem 3.1 may appear as a classical result for the Navier-Stokes system: for the reader's convenience, we only write the proof in Section A.5 of the Appendix.

In view of the conditional Theorem 3.1, obtaining decay in time for the moments of f will be the main goal of the rest of this paper.

4 Preliminaries for the bootstrap procedure

Thanks to a bootstrap argument, we shall prove that there exists $C > 0$ such that

$$\forall t \in \mathbb{R}^+, \quad \|j_f(t) - \rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} \leq \frac{C}{(1+t)^{7/4}}.$$

Then, by the conditional Theorem 3.1, this will imply the first statement of Theorem 2.1. Along the way, we shall rely on the following bounds and estimates for the local density ρ_f :

$$\begin{aligned} \|\rho_f\|_{L^\infty(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3))} &\lesssim 1, \\ \forall t \in \mathbb{R}^+, \forall r \geq 1, \quad \|\rho_f(t)\|_{L^r(\mathbb{R}_+^3)} &\lesssim \frac{1}{(1+t)^k}, \quad \text{for some } k > 0 \end{aligned}$$

which will essentially lead to the second part of Theorem 2.1.

In this section, we collect several useful information in order to be able to set up a bootstrap procedure in Section 6. We start by recalling some basic facts about the Lagrangian structure of the Vlasov equation (1.1). A careful analysis of the characteristic curves will indeed be required to deal with the absorption at the boundary. It also enables us to consider the aforementioned straightening change of variable in velocity. Then, we derive local in time estimates for the moments ρ_f and j_f , as well as for the fluid velocity in $L^1 L^\infty$. This will offer short time controls on these quantities. Thanks to a smoothing property of the Navier-Stokes system, we finally obtain $L^2 H^2 \cap L^\infty H^1$ estimates for the fluid velocity. This requires the introduction of the so-called *strong existence times* and eventually entails a local in time $L^1 W^{1, \infty}$ regularity for the fluid velocity.

4.1 Characteristic curves for the Vlasov equation: representation formula and change of variable

Given a time-dependent vector field u on $\mathbb{R}^+ \times \mathbb{R}_+^3$, a time $t \in \mathbb{R}^+$ and a point $(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3$, we define the characteristic curves $s \in \mathbb{R}^+ \mapsto (X(s; t, x, v), V(s; t, x, v)) \in \mathbb{R}^3 \times \mathbb{R}^3$ for the Vlasov equation (associated to u) as the solution of the following system of ordinary differential equations

$$\begin{cases} \dot{X}(s; t, x, v) = V(s; t, x, v), \\ \dot{V}(s; t, x, v) = (Pu)(s, X(s; t, x, v)) + G - V(s; t, x, v), \\ X(t; t, x, v) = x, \\ V(t; t, x, v) = v. \end{cases} \quad (4.1)$$

where the dot means derivative along the first variable. Here, P is the linear extension operator continuous from $L^\infty(\mathbb{R}_+^3)$ to $L^\infty(\mathbb{R}^3)$ and from $W_0^{1,\infty}(\mathbb{R}_+^3)$ to $W^{1,\infty}(\mathbb{R}^3)$ defined by

$$\forall x \in \mathbb{R}^d, \quad (Pw)(x) := \begin{cases} w(x) & \text{if } x \in \mathbb{R}_+^3, \\ 0 & \text{if } x \in \mathbb{R}^3 \setminus \mathbb{R}_+^3, \end{cases} \quad (4.2)$$

and which satisfies

$$\forall w \in L^\infty(\mathbb{R}_+^3), \quad \|Pw\|_{L^\infty(\mathbb{R}^3)} = \|w\|_{L^\infty(\mathbb{R}_+^3)}, \quad (4.3)$$

$$\forall w \in W_0^{1,\infty}(\mathbb{R}_+^3), \quad \|\nabla(Pw)\|_{L^\infty(\mathbb{R}^3)} \leq \|\nabla w\|_{L^\infty(\mathbb{R}_+^3)}. \quad (4.4)$$

We refer to [24, Appendix] for a simple proof. Also, we will use the convention

$$(Pu)(t, \cdot) = P(u(t, \cdot)),$$

as well as the notation

$$X_t^s(x, v) := X(s; t, x, v), \quad V_t^s(x, v) := V(s; t, x, v).$$

Let $T > 0$ be fixed and suppose

$$u \in L^2(0, T; H_0^1(\mathbb{R}_+^3)) \cap L^1(0, T; W^{1,\infty}(\mathbb{R}_+^3)).$$

We can apply the Cauchy-Lipschitz theorem to show the following proposition.

Proposition 4.1. *Given $(x, v) \in \mathbb{R}^3 \times \mathbb{R}^3$ and a time $t \in [0, T]$, the system (4.1) admits a unique solution $s \mapsto Z_t^s(x, v) \in \mathbb{R}^3 \times \mathbb{R}^3$ on $[0, T]$ and*

$$Z_t^s : \begin{cases} \mathbb{R}^3 \times \mathbb{R}^3 & \longrightarrow \mathbb{R}^3 \times \mathbb{R}^3 \\ (x, v) & \longmapsto Z_t^s(x, v) := (X(s; t, x, v), V(s; t, x, v)) \end{cases}$$

is a (bi-Lipschitz) diffeomorphism of $\mathbb{R}^3 \times \mathbb{R}^3$ whose inverse is given by $(Z_t^s)^{-1} = Z_s^t$.

In this context, the characteristic curves for the Vlasov equation are classically defined (at least) until time T and are given for all $s, t \in [0, T]$ by

$$\begin{cases} X(s; t, x, v) = x + (1 - e^{-s+t})v + (s - t + e^{-s+t} - 1)G + \int_t^s (1 - e^{\tau-s})(Pu)(\tau, X(\tau; t, x, v)) d\tau, \\ V(s; t, x, v) = e^{-s+t}v + (1 - e^{-s+t})G + \int_t^s e^{\tau-s}(Pu)(\tau, X(\tau; t, x, v)) d\tau. \end{cases} \quad (4.5)$$

Starting from a point $(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3$ at time t , the curve $X(s; t, x, v)$ remains during a certain interval of time in the half-space \mathbb{R}_+^3 . This naturally leads to the following definitions, already considered in [24].

Definition 4.2. For $(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3$ and for any $t \geq 0$, we set

$$\tau^+(t, x, v) := \sup \{s \geq t \mid \forall \sigma \in [t, s], X(\sigma; t, x, v) \in \mathbb{R}_+^3\}, \quad (4.6)$$

$$\tau^-(t, x, v) := \inf \{s \leq t \mid \forall \sigma \in [s, t], X(\sigma; t, x, v) \in \mathbb{R}_+^3\}, \quad (4.7)$$

with the convention $X(\sigma; t, x, v) = X(0; t, x, v)$ for all $\sigma < 0$ and $t \geq 0$.

We state two basic results whose proof can be found in [24, Appendix]. The second one is a representation formula for the weak solution to the Vlasov equation (1.1), where one has to take into account the absorption boundary condition (1.11) satisfied by the distribution function.

Lemma 4.3. For $z = (x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3$, if $\tau^+(0, z) < +\infty$ then we have

$$Z_0^{\tau^+(0, z)}(z) = \left(X_0^{\tau^+(0, z)}(z), V_0^{\tau^+(0, z)}(z) \right) \in \Sigma^+ \cup \Sigma^0, \quad (4.8)$$

where Σ^+ and Σ^0 are defined in (1.8).

Proposition 4.4. Let f be the weak solution to the initial boundary value problem for the Vlasov equation, associated to a velocity field $u \in L_{\text{loc}}^2(\mathbb{R}^+; H_0^1(\mathbb{R}_+^3)) \cap L_{\text{loc}}^1(\mathbb{R}^+; W^{1, \infty}(\mathbb{R}_+^3))$ with initial condition f_0 and with absorption boundary condition.

$$f(t, x, v) = e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} f_0(Z_t^0(x, v)) \quad a.e. \quad (4.9)$$

Recall that we aim at obtaining a sufficient decay in time of the Brinkman force $j_f - \rho_f u$ and more generally of the moments of f . The representation formula (4.9) will be our main starting point: elaborating on the same strategy than that of [37, 34, 24, 35], we shall rely on a straightening change of variable in velocity (and then in space) in this formula. Nevertheless, such a procedure requires a smallness assumption on the quantity $\|\nabla u\|_{L^1(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3))}$ and obtaining this control will be at the core of Section 6. We emphasize the fact that we also need the help of the absorption at the boundary in order to recover the desired decay in time of the moments.

In view of the following formulas

$$\begin{aligned} V(0; t, x, v) &= e^t v + (1 - e^t)G - \int_0^t e^\tau (Pu)(\tau, X(\tau; t, x, v)) d\tau, \\ D_v V(0; t, x, v) - e^t \text{Id} &= - \int_0^t e^\tau \nabla (Pu)(\tau, X(\tau; t, x, v)) D_v X(\tau; t, x, v) d\tau, \end{aligned}$$

and following closely the arguments of [37], we infer several statements which read as follows.

Lemma 4.5. Suppose $u \in L_{\text{loc}}^2(\mathbb{R}^+; H_0^1(\mathbb{R}_+^3)) \cap L_{\text{loc}}^1(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3))$. Fix $\delta > 0$ satisfying $\delta e^\delta < 1/9$. Then, for all times $t \in \mathbb{R}^+$ satisfying

$$\int_0^t \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \delta, \quad (4.10)$$

and for all $x \in \mathbb{R}_+^3$, the map

$$\Gamma_{t, x} : v \mapsto V(0; t, x, v),$$

is a global \mathcal{C}^1 -diffeomorphism from \mathbb{R}^3 to itself satisfying furthermore

$$\forall v \in \mathbb{R}^3, \quad |\det D_v \Gamma_{t, x}(v)| \geq \frac{e^{3t}}{2}. \quad (4.11)$$

Thanks to

$$|\Gamma_{t, x}^{-1}(w)| \leq e^{-t} \left[|w| + (e^t - 1)|G| + \int_0^t e^\tau \|Pu(\tau)\|_{L^\infty(\mathbb{R}^3)} d\tau \right], \quad (4.12)$$

$$|v| \leq |V(0; t, x, v)| + (1 - e^{-t})|G| + \int_0^t \|Pu(\tau)\|_{L^\infty(\mathbb{R}^3)} d\tau. \quad (4.13)$$

we also have the following result.

Lemma 4.6. *Let $t_0 > 0$. If $N_q(f_0) < \infty$ and if $u \in L^1_{\text{loc}}(\mathbb{R}_+; H^1_0 \cap L^\infty(\mathbb{R}^3_+))$ then $N_q(f(t_0)) < \infty$ with*

$$N_q(f(t_0)) \lesssim e^{3t_0} (1 + |G|^q + \|u\|_{L^1(0,t_0;L^\infty(\mathbb{R}^3_+))}^q) N_q(f_0).$$

As we shall see later, we will also perform a change of variable in space after the previous change of variable in velocity.

Lemma 4.7. *Consider the same assumptions as in Lemma 4.5. For any $t \in \mathbb{R}^+$ satisfying*

$$\int_0^t \|\nabla u(s)\|_{L^\infty(\mathbb{R}^3_+)} ds \leq \delta,$$

and for any $v \in \mathbb{R}^3$, the mapping

$$\Lambda_{t,v} : x \longrightarrow X_t^0(x, \Gamma_{t,x}^{-1}(v)) \tag{4.14}$$

is a global \mathcal{C}^1 -diffeomorphism from \mathbb{R}^3 to itself satisfying

$$\forall x \in \mathbb{R}^3, \quad |\det D_x \Lambda_{t,v}(x)| \geq \frac{1}{2}. \tag{4.15}$$

Proof. From

$$\Gamma_{t,x}^{-1}(w) = e^{-t}w + (1 - e^t)G + \int_0^t e^{\tau-t}(Pu)(\tau, X_t^\tau(x, \Gamma_{t,x}^{-1}(w))) d\tau,$$

we infer that

$$\Lambda_{t,w}(x) = x + (e^{-t} - 1)w + (1 - t + e^{-t})G + \int_0^t (e^{\tau-t} - 1)(Pu)(\tau, X_t^\tau(x, \Gamma_{t,x}^{-1}(w))) d\tau.$$

We then refer to [35, Lemma 3.26], the proof of which can be exactly adapted to the case of the characteristic curves with an additional gravity term. \square

4.2 Local in time estimates

In this subsection, we derive some local in time estimates for the moments and velocity field. Recalling the quantities of Definition 1.4, we introduce the following useful notations, which allows us to track down the dependency on the initial data in the later estimates. In view of the smallness condition (2.2), this will enable us to set up a bootstrap strategy in the proof of Theorem 2.1.

Notation 4.8. The notation $m \lesssim_0 M$ means that there exists some exponents $q > 3$ and $m > 0$ such that

$$m \leq \varphi \left(1 + \|u_0\|_{H^1 \cap D_p^{1-\frac{1}{p},p}(\mathbb{R}^3_+)} + \mathbf{E}(0) + \|u_0\|_{L^1(\mathbb{R}^3_+)} + N_q(f_0) + H_{0,q}(f_0) + \max_{r \in \{1,3\}} \{K_{m,r}(f_0) + F_{q,m,r}(f_0)\} + M_6 f_0 \right) M, \tag{4.16}$$

where $\varphi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is onto, continuous and nondecreasing. Here, $D_p^{1-\frac{1}{p},p}(\mathbb{R}^3_+)$ refers to the space defined in (A.7), with the exponent p given in the statement of Theorem 2.1. Note that in the end of the bootstrap argument, we shall be able to consider only the largest exponents q and m which are involved in the estimates.

Until the end of this Section 4, we consider a fixed global Leray solution (u, f) to the Vlasov-Navier-Stokes system, in the sense of Definition 1.6 and associated to an admissible initial data (u_0, f_0) .

We first state the following lemma, claiming that the total mass of the system is nonincreasing along its evolution, because of the absorption boundary condition (1.11). The proof is similar to that of [24, Lemma 2.1].

Lemma 4.9. For all $t \geq s \geq 0$, we have

$$\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) \, dx \, dv \leq \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(s, x, v) \, dx \, dv.$$

An application of Hölder inequality implies the following result.

Lemma 4.10. Let $p > 1$. For all $t > 0$, we have

$$\|(j_f - \rho_f u)(t)\|_{L^p(\mathbb{R}_+^3)} \leq \|\rho_f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v - u(t, x)|^p \, dx \, dv \right)^{1/p}. \quad (4.17)$$

Along the way, we will also need the simple following sublinear Grönwall lemma.

Lemma 4.11. Let $y \in \mathcal{C}([0, T]; \mathbb{R}^+)$ and $h \in \mathcal{C}([0, T]; \mathbb{R}^+)$ such that for all $t \in [0, T]$

$$y(t) \leq y_0 + \int_0^t h(s) y(s)^\beta \, ds,$$

where $\beta \in (0, 1)$ and $y_0 \in \mathbb{R}^+$. Then for all $t \in [0, T]$

$$y(t) \leq \left(y_0^{1-\beta} + (1-\beta) \int_0^t h(s) \, ds \right)^{\frac{1}{1-\beta}}.$$

We now state several rough bounds on the moments of the distribution function f .

Lemma 4.12. We have for all $t \geq 0$

$$\|j_f\|_{L^1(0, t; L^1(\mathbb{R}_+^3))} \lesssim E(0)^{\frac{1}{2}} t + t^2,$$

In particular, the map $t \mapsto \|j_f\|_{L^1(0, t; L^1(\mathbb{R}_+^3))}$ belongs to $L_{\text{loc}}^\infty(\mathbb{R}^+)$.

Proof. We use the energy inequality (1.29) and the Cauchy-Schwarz inequality to write that for all $t \geq 0$

$$E(t) \leq E(0) + g \int_0^t \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} |v| f(s, x, v) \, dx \, dv \, ds \leq E(0) + g \int_0^t E(s)^{\frac{1}{2}} \, ds,$$

thanks to Lemma 4.9. The sub-linear version of Grönwall Lemma 4.11 with $\beta = 1/2$ implies

$$E(t) \leq \left(E(0)^{\frac{1}{2}} + \frac{g}{2} t \right)^2.$$

The Cauchy-Schwarz inequality again reads

$$\int_0^t \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} |v| f(s, x, v) \, dx \, dv \, ds \leq \int_0^t E(s)^{\frac{1}{2}} \, ds \leq t E(0)^{\frac{1}{2}} + g \frac{t^2}{4},$$

and this concludes the proof. \square

Lemma 4.13. If $F := j_j - \rho_f u$ then for all $t \geq 0$, we have

$$\int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 \, ds \leq \sup_{s \in [0, t]} \|\rho_f(s)\|_{L^\infty(\mathbb{R}_+^3)} \left[E(0) + E(0)^{\frac{1}{2}} t + t^2 \right],$$

where \lesssim only depends on g .

Proof. Using Lemma 4.10 with $p = 2$, we obtain for all $s \in [0, t]$

$$\|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 \leq \|\rho_f(s)\|_{L^\infty(\mathbb{R}_+^3)} D(s),$$

where the dissipation D has been defined in (1.14). Then, we integrate the last inequality between 0 and t and use the energy inequality (1.29) to obtain

$$\int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \leq \sup_{s \in [0, t]} \|\rho_f(s)\|_{L^\infty(\mathbb{R}_+^3)} \left[E(0) + g \int_0^t \int_{\mathbb{R}_+^3} |j_f| d\tau dx \right],$$

which concludes the proof, thanks to Lemma 4.12. \square

As any weak (renormalised) solution to the Vlasov equation, the distribution function f also satisfies the following maximum principle.

Lemma 4.14. *We have for all $t \geq 0$*

$$\|f(t)\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)} \leq e^{3t} \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}. \quad (4.18)$$

We now recall standard interpolation estimates on the moments (see e.g. [33]), where we use the notations introduced in Definition 1.3.

Proposition 4.15. *Let $k > 0$ and let h be a nonnegative function in $L^\infty(\mathbb{R}^+ \times \mathbb{R}_+^3 \times \mathbb{R}^3)$. Then the following estimates hold for any $\ell \in [0, k]$ and a.e. (t, x)*

$$\|m_\ell h(t)\|_{L^{\frac{k+3}{\ell+3}}(\mathbb{R}_+^3)} \leq C_{k, \ell} \|h(t)\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{k-\ell}{k+3}} M_k h(t)^{\frac{\ell+3}{k+3}}.$$

We now provide a pointwise estimate for the moments of f .

Lemma 4.16. *Consider an integer $\alpha \geq 2$ such that $u \in L_{\text{loc}}^1(\mathbb{R}^+; L^{\alpha+3} \cap W^{1,1}(\mathbb{R}_+^3))$ and $M_\alpha f_0 < \infty$ for all $a \in [2, \alpha]$. Then for all $a \in [2, \alpha]$ and for all $t > 0$, $M_a f(t) < \infty$ and $M_a f \in L_{\text{loc}}^\infty(\mathbb{R}^+)$.*

Furthermore, if $T > 0$ then for all $t \in [0, T]$

$$M_\alpha f(t) \leq \left(\left[M_\alpha f_0 + \alpha g \int_0^T M_{\alpha-1} f(s) ds \right]^{\frac{1}{\alpha+3}} + e^{\frac{3t}{\alpha+3}} \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{\alpha+3}} \int_0^t \|u(s)\|_{L^{\alpha+3}(\mathbb{R}_+^3)} ds \right)^{\alpha+3}. \quad (4.19)$$

Proof. As in [24, Lemma 2.1], we rely on the strong stability results from DiPerna-Lions theory about transport equations on $\mathbb{R}_+^3 \times \mathbb{R}^3$ (see A.1 in the Appendix). In short, it allows to prove the desired estimates for a sequence of smooth distributions (f_n) associated to a sequence of regularized initial data $(f_n^0)_n$ and an approximating sequence of fluid velocities (u_n) : the strong stability property of renormalized solutions to the Vlasov equation is then used to recover the estimate for the original solution (u, f) . We do not detail the argument and we write the proof as if u and f were smooth.

We argue by induction on α . For $\alpha = 2$, we multiply the Vlasov equation by $|v|^2$ and we integrate on $\mathbb{R}_+^3 \times \mathbb{R}^3$ to get

$$\frac{d}{dt} M_2 f(t) + \alpha M_2 f(t) + \int_{\mathbb{R}^3} \int_{\partial \mathbb{R}_+^3} |v|^2 (\gamma f) n(x) \cdot v d\sigma(x) dv = 2 \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f v \cdot [u + G] dx dv.$$

As the boundary term of the l.h.s is non-negative (because of (1.11)), we integrate on the interval $(0, t)$ to obtain

$$M_2 f(t) \leq M_2 f_0 + 2g \int_0^t M_1 f(s) ds + 2 \int_0^t \|u(s)\|_{L^5(\mathbb{R}_+^3)} \|m_1 f(s)\|_{L^{\frac{5}{4}}(\mathbb{R}_+^3)} ds,$$

thanks to Hölder inequality. Furthermore, by Proposition 4.15 with $(\ell, k) = (1, 2)$ and the rough control provided by Lemma 4.14, we get

$$\|m_1 f(s)\|_{L^{\frac{5}{4}}(\mathbb{R}_+^3)} \lesssim \|f(s)\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{5}} M_2 f(s)^{\frac{4}{5}} \leq \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{5}} e^{\frac{3s}{5}} M_2 f(s)^{\frac{4}{5}},$$

where \lesssim is independent of s . We thus infer that for all $t \in [0, T]$

$$M_2 f(t) \leq M_2 f_0 + 2g \int_0^T M_1 f(s) ds + 2\|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{5}} \int_0^t e^{\frac{3s}{5}} \|u(s)\|_{L^5(\mathbb{R}_+^3)} M_2 f(s)^{\frac{4}{5}} ds.$$

Unsinf the Grönwall lemma stated in Proposition 4.11 with $\beta = 4/5$ entails

$$M_2 f(t) \lesssim \left(\left[M_2 f_0 + 2g \int_0^T M_1 f(s) ds \right]^{\frac{1}{5}} + \frac{2}{5} e^{\frac{3t}{5}} \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{5}} \int_0^t \|u(s)\|_{L^5(\mathbb{R}_+^3)} ds \right)^5.$$

This yields the result for $\alpha = 2$ (indeed, note that the previous r.h.s is finite because of Lemma 4.12). Now, if $M_{\alpha-1}$ satisfies the desired inequality, we can perform the same analysis as before. In short, multiplying the Vlasov equation by $|v|^\alpha$ and integrating on $\mathbb{R}_+^3 \times \mathbb{R}^3$, we obtain

$$M_\alpha f(t) \leq M_\alpha f_0 + \alpha g \int_0^t M_{\alpha-1} f(s) ds + \alpha \int_0^t \|u(s)\|_{L^{\alpha+3}(\mathbb{R}_+^3)} \|m_{\alpha-1} f(s)\|_{L^{\frac{\alpha+3}{\alpha+2}}(\mathbb{R}_+^3)} ds.$$

As before, we use Proposition 4.15 with $(\ell, k) = (\alpha - 1, \alpha)$ and Lemma 4.14 to get

$$\|m_{\alpha-1} f(s)\|_{L^{\frac{\alpha+3}{\alpha+2}}(\mathbb{R}_+^3)} \lesssim \|f_0\|_{L^\infty}^{\frac{1}{\alpha+3}} e^{\frac{3s}{\alpha+3}} M_\alpha f(s)^{\frac{\alpha+2}{\alpha+3}}.$$

This yields for all $t \in [0, T]$

$$M_\alpha f(t) \leq M_\alpha f_0 + \alpha g \int_0^t M_{\alpha-1} f(s) ds + \alpha \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{\alpha+3}} \int_0^t e^{\frac{3s}{\alpha+3}} \|u(s)\|_{L^{\alpha+3}(\mathbb{R}_+^3)} M_\alpha f(s)^{\frac{\alpha+2}{\alpha+3}} ds.$$

Thank to Proposition 4.11 with $\beta = \frac{\alpha+2}{\alpha+3}$, we obtain the conclusion. \square

Lemma 4.17. *We have $M_3 f \in L_{\text{loc}}^\infty(\mathbb{R}^+)$. Moreover, for all finite $T > 0$, there exists a continuous nonnegative and nondecreasing function $\varphi_{\mathbb{E}(0), M_3 f_0, T, \|f_0\|_{L_{x,v}^\infty}}$ (increasing in all its parameters) such that for all $t \in [0, T]$*

$$\|\rho_f(t)\|_{L^2(\mathbb{R}_+^3)} + \|j_f(t)\|_{L^{3/2}(\mathbb{R}_+^3)} \leq \varphi_{\mathbb{E}(0), M_3 f_0, T, \|f_0\|_{L_{x,v}^\infty}}(t).$$

Proof. Since $M_2 f_0 < \infty$ and $M_6 f_0 < \infty$ by (2.1), we have $M_3 f_0 \lesssim M_2 f_0 + M_6 f_0 < \infty$. Furthermore, u is a Leray solution so that by Sobolev embedding, we can apply Lemma 4.16 with $\alpha = 3$ and we deduce that $M_3 f \in L_{\text{loc}}^\infty(\mathbb{R}^+)$. Furthermore, the estimate (4.19) yields for all $t \in [0, T]$

$$M_3 f(t) \lesssim \left(\left[M_3 f_0 + 3g \int_0^T M_2 f(s) ds \right]^{\frac{1}{6}} + \frac{e^{\frac{t}{2}}}{2} \|f_0\|_{L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)}^{\frac{1}{6}} \int_0^t \|u(s)\|_{L^6(\mathbb{R}_+^3)} ds \right)^6.$$

Sobolev embedding on $H_0^1(\mathbb{R}_+^3)$ and the Cauchy-Schwarz inequality lead to

$$\int_0^t \|u(s)\|_{L^6(\mathbb{R}_+^3)} ds \lesssim_g t^{1/2} \left(\int_0^t \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right)^{1/2} \leq T^{1/2} \left[\mathbb{E}(0) + \mathbb{E}(0)^{\frac{1}{2}} T + T^2 \right]^{1/2},$$

where we have used the energy inequality (1.29) and Lemma 4.12. For the same reasons, we also have

$$\int_0^T M_2 f(s) ds \lesssim_g \mathbb{E}(0) + \mathbb{E}(0)^{\frac{1}{2}} T + T^2.$$

This implies that there exists a continuous nonnegative and nondecreasing function $\varphi = \varphi_{E(0), M_3 f_0, T, \|f_0\|_{L^\infty_v}}$ (increasing in all its parameters) such that for all $t \in [0, T]$, we have

$$M_3 f(t) \leq \varphi_{E(0), M_3 f_0, T, \|f_0\|_{L^\infty_v}}(t). \quad (4.20)$$

Using Proposition 4.15 on interpolation of moments of the distribution function f with $k = 3$ and $\ell \in \{0, 1\}$, together with Lemma 4.14, we get

$$\begin{aligned} \|\rho_f(t)\|_{L^2(\mathbb{R}_+^3)} &\lesssim \|f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{1/2} M_3 f(t)^{1/2} \leq \|f_0\|_{L^\infty(\mathbb{R}_+^3)}^{1/2} e^{3t/2} M_3 f(t)^{1/2}, \\ \|j_f(t)\|_{L^{3/2}(\mathbb{R}_+^3)} &\lesssim \|f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{1/3} M_3 f(t)^{2/3} \leq \|f_0\|_{L^\infty(\mathbb{R}_+^3)}^{1/3} e^t M_3 f(t)^{2/3}, \end{aligned}$$

which yields the conclusion, thanks to the bound (4.20). \square

We now prove that the source term in the Navier-Stokes equations, namely the Brinkman force $j_f - \rho_f u$, belong to $L^2 L^2$ (locally in time). The strategy of proof is very similar to that of [24, Lemma 4.7] (with a minor adaptation to the half-space case) and details are thus omitted. Let us only emphasize that the result mainly follows from the combination of Lemma 4.12, Lemma 4.17, Lemma 4.15 and the maximal regularity property for the Stokes system (see Section A.4 in the Appendix) which can be applied in that case because $u_0 \in W_0^{1,9/7}(\mathbb{R}_+^3)$ (see the assumption (2.1)).

Lemma 4.18. *We have*

$$j_f - \rho_f u \in L_{\text{loc}}^2(\mathbb{R}^+; L^2(\mathbb{R}_+^3)).$$

We are then in position to state the following local in times integrability results of the Leray solution (u, f) .

Proposition 4.19. *We have*

$$u \in L_{\text{loc}}^1(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3)), \quad (4.21)$$

and if $N_q(f_0) < \infty$ for some $q > 3$, then

$$\begin{aligned} \rho_f &\in L_{\text{loc}}^\infty(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3)), \\ j_f &\in L_{\text{loc}}^\infty(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3)). \end{aligned}$$

More precisely, there exists a continuous nondecreasing function $\eta : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ such that

$$\|u\|_{L^1(0,t; L^\infty(\mathbb{R}_+^3))} \lesssim \eta(t), \quad (4.22)$$

$$\|\rho_f\|_{L^\infty(0,t; L^\infty(\mathbb{R}_+^3))} + \|j_f\|_{L^\infty(0,t; L^\infty(\mathbb{R}_+^3))} \lesssim N_q(f_0) \eta(t). \quad (4.23)$$

Proof. Let $T > 0$. The proof of the fact that $u \in L^1(0, T; L^\infty(\mathbb{R}_+^3))$ is mostly directly taken from the arguments used in [24, Proposition 4.8] and mainly relies on the theory of epochs of regularity for the Leray solutions to the Navier-Stokes equations. Owing to [39, Theorem 8] and [40, Remark 4] (which are valid since the strong energy inequality (1.28) is satisfied by the weak solutions that we consider), we know there exists a subset $\sigma_T \subset [0, T]$ of full measure in $[0, T]$ with $\sigma_T = \bigsqcup_i]a_i, b_i[$ (the union being countable) and for which $u \in L_{\text{loc}}^\infty(a_i, b_i; H^1(\mathbb{R}_+^3)) \cap L_{\text{loc}}^2(a_i, b_i; H^2(\mathbb{R}_+^3))$ and $\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)} \xrightarrow[t \rightarrow b_i^-]{} +\infty$ for all i . Furthermore, the function $t \mapsto \|\nabla u(t)\|^2$ is absolutely continuous on each interval $]a_i, b_i[$ (see e.g. [55]).

Now, we can take the $L^2(\mathbb{R}_+^3)$ inner product of (1.2) with Au on each interval $]a_i, b_i[$, where A stands for the Stokes operator on $L_{\text{div}}^2(\mathbb{R}_+^3)$, to obtain

$$\frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + 2\|Au\|_{L^2(\mathbb{R}_+^3)}^2 + 2\langle \mathbb{P}(u \cdot \nabla)u, Au \rangle = 2\langle \mathbb{P}(j_f - \rho_f u), Au \rangle \quad \text{on }]a_i, b_i[,$$

where we have dropped the time variable. In order to estimate the term $|\langle \mathbb{P}(u \cdot \nabla)u, Au \rangle|$, we use the Gagliardo-Nirenberg-Sobolev inequality for the function ∇u with the exponents $(p, j, r, m, \alpha) = (3, 0, 2, 1, 1/2)$, which reads as

$$\|\nabla u\|_{L^3(\mathbb{R}_+^3)} \lesssim \|D^2 u\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^{1/2},$$

and we combine this inequality with the Hölder inequality to write

$$\begin{aligned} |\langle \mathbb{P}(u \cdot \nabla)u, Au \rangle| &\leq \|u\|_{L^6(\mathbb{R}_+^3)} \|\nabla u\|_{L^3(\mathbb{R}_+^3)} \|Au\|_{L^2(\mathbb{R}_+^3)} \\ &\lesssim \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^{3/2} \|Au\|_{L^2(\mathbb{R}_+^3)}^{3/2}. \end{aligned}$$

Thanks to the Young inequality, we infer that

$$\frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + \|Au\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim \|j_f - \rho_f u\|_{L^2(\mathbb{R}_+^3)}^2 + \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^6, \quad (4.24)$$

on each interval $]a_i, b_i[$, where \lesssim is independent of the time variable and independent of i . Dividing this inequality by $(1 + \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2)^2$, we get

$$\frac{1}{(1 + \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2)^2} \frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{(1 + \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2)^2} \|Au\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim \|j_f - \rho_f u\|_{L^2(\mathbb{R}_+^3)}^2 + \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2,$$

on each interval $]a_i, b_i[$. Integrating and summing over the previous epoch of regularities, and using in particular the fact that $\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)} \xrightarrow[t \rightarrow b_i^-]{} +\infty$ for all i , we can perform the same exact computations as in the proof of [24, Proposition 4.8] and end up with

$$\int_0^T \frac{\|Au(s)\|_{L^2(\mathbb{R}_+^3)}^2}{(1 + \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2)^2} ds \leq \int_0^T \left(\|j_f(s) - \rho_f u(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 \right) ds + 1,$$

from which we infer that

$$\int_0^T \|Au(s)\|_{L^2(\mathbb{R}_+^3)}^{2/3} ds \leq \left(\int_0^T \frac{\|Au(s)\|_{L^2(\mathbb{R}_+^3)}^2}{(1 + \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2)^2} ds \right)^{1/3} \left(\int_0^T (1 + \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2) ds \right)^{2/3} < \infty,$$

due to Proposition 4.18 and to the fact that u is a Leray solution to the Navier-Stokes equations. Using the Gagliardo-Nirenberg-Sobolev inequality with exponents $(p, j, r, m, \alpha) = (\infty, 0, 2, 2, 1/2)$ and Sobolev embedding, we deduce that

$$\int_0^T \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds \lesssim \int_0^T \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds + \int_0^T \|Au(s)\|_{L^2(\mathbb{R}_+^3)}^{2/3} ds < \infty,$$

from the same reasons as before, therefore this proves (4.21). The last estimate (4.23) is eventually obtained by observing that

$$\begin{aligned} \|\rho_f\|_{L^\infty(0,t;L^\infty(\mathbb{R}_+^3))} + \|j_f\|_{L^\infty(0,t;L^\infty(\mathbb{R}_+^3))} &\lesssim N_q(f(t)) \\ &\lesssim e^{3t} (1 + |G|^q + \|u\|_{L^1(0,t;L^\infty(\mathbb{R}_+^3))}^q) N_q(f_0) \\ &\lesssim N_q(f_0) \eta(t), \end{aligned}$$

thanks to Lemma 4.6 and (4.21). □

4.3 Strong existence times and higher order energy estimates

Along the bootstrap procedure, we shall need H^1 energy estimates for the fluid velocity u , which is *a priori* only a Leray solution to the Navier-Stokes equations. In order to consider higher regularity for this solution, we rely on a parabolic smoothing property of the (Vlasov-)Navier-Stokes system. We will be able to propagate this regularity if the contribution of the source term, that is the Brinkman force $F := j_f - \rho_f u$, and the initial data, are small enough.

Proposition 4.20. *There exists a universal constant C_\star such that the following holds. Assume that for some $T > 0$, one has*

$$\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds + \int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)} ds < C_\star. \quad (4.25)$$

Then one has

$$u \in L^\infty(0, T; H^1(\mathbb{R}_+^3)) \cap L^\infty(0, T; H^2(\mathbb{R}_+^3)),$$

and for all $t \in [0, T]$

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_0^t \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds, \quad (4.26)$$

where \lesssim only depends on C_\star .

Proof. The estimate is a direct consequence of the parabolic regularization for the Navier-Stokes system with source $F = j_f - \rho_f u$, that we state in Theorem A.8 in Section A.6 of the Appendix. \square

Remark 4.21. By Lemma 4.13 and the local estimates (4.23), the r.h.s of (4.26) is finite. In particular, if the condition (4.25) holds for some T , then for all $t \in [0, T]$

$$\|u\|_{L^\infty(0, t; L^6(\mathbb{R}_+^3))}^2 \lesssim \|\nabla u\|_{L^\infty(0, t; L^2(\mathbb{R}_+^3))}^2 \lesssim \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \sup_{s \in [0, t]} \|\rho_f(s)\|_{L^\infty(\mathbb{R}_+^3)} \left[E(0) + E(0)^{\frac{1}{2}} t + t^2 \right].$$

In order to ensure that the smallness condition (4.25) is satisfied for all times, we now introduce the following terminology, which has been already used in [37, 34] to take advantage of the parabolic regularization for the fluid.

Definition 4.22 (Strong existence time). A real number $T \geq 0$ is a *strong existence* time whenever (4.25) holds.

In the remaining part of this Section, we state a local in time $L^1 W^{1, \infty}$ regularity result for the fluid velocity. Note that for the moment, we are only interested in obtaining non-uniform in time estimates. Of course, quantitative and uniform in time estimates based on the polynomial decay of the kinetic energy will require an additional analysis.

Corollary 4.23. *For any finite strong times $t > 0$ and for any $p \in [1, 6]$, we have*

$$j_f - \rho_f u \in L^p(0, t; L^p(\mathbb{R}_+^3)).$$

Proof. Thanks to Lemma 4.17 and (4.23), we have

$$j_f \in L^\infty(0, t; L^{3/2}(\mathbb{R}_+^3)) \cap L^\infty(0, t; L^\infty(\mathbb{R}_+^3)) \hookrightarrow L^r(0, t; L^r(\mathbb{R}_+^3)),$$

for any $r \in [3/2, +\infty]$ by interpolation. For the same reasons, we also have

$$\rho_f \in L^\infty(0, t; L^2(\mathbb{R}_+^3)) \cap L^\infty(0, t; L^\infty(\mathbb{R}_+^3)) \hookrightarrow L^r(0, t; L^r(\mathbb{R}_+^3)),$$

for any $r \in [2, +\infty]$. On the other hand, by (6.16), we have $u \in L^\infty(0, t; L^q(\mathbb{R}_+^3))$ for any $q \in [2, 6]$, because t is a strong existence time. So, the Hölder inequality yields

$$\rho_f u \in L^p(0, t; L^p(\mathbb{R}_+^3)), \quad \frac{1}{p} = \frac{1}{r} + \frac{1}{q}, \quad r \in [2, +\infty], \quad q \in [2, 6].$$

This leads to the condition $p \in [1, 6]$ and then to the conclusion. \square

The next results is similar to that of [35, Lemma 3.28]. We detail the proof for the sake of completeness.

Proposition 4.24. *For any finite strong times $t \geq 0$ and for the exponent p given in the assumption (2.1), we have*

$$(u \cdot \nabla)u \in L^p(0, t; L^p(\mathbb{R}_+^3)). \quad (4.27)$$

Proof. For any $(a, b, r_1, r_2) \in (1, +\infty)^4$, we can use interpolation inequalities to write

$$\|(u \cdot \nabla)u\|_{L^a(0, t; L^b(\mathbb{R}_+^3))} \leq \|u\|_{L^\infty(0, t; L^6(\mathbb{R}_+^3))} \|\nabla u\|_{L^\infty(0, t; L^2(\mathbb{R}_+^3))}^{1-\frac{r_1}{a}} \|\nabla u\|_{L^{r_1}(0, t; L^{r_2}(\mathbb{R}_+^3))}^{\frac{r_1}{a}}, \quad (4.28)$$

provided that $r_1 \leq a$, $2 \leq b \leq r_2$ and

$$\frac{1}{b} = \frac{2}{3} + \frac{r_1}{a} \left(\frac{1}{r_2} - \frac{1}{2} \right). \quad (4.29)$$

Taking $(a, b, r_1, r_2) = (2, 3, 2, 6)$ in (4.28), Sobolev embedding and Proposition 4.20 imply

$$(u \cdot \nabla)u \in L^2(0, t; L^3(\mathbb{R}_+^3)).$$

Owing to the maximal regularity for the Stokes system (see Section A.4 in the Appendix) and to Corollary 4.23 which gives $j_f - \rho_f u \in L^2(0, t; L^3(\mathbb{R}_+^3))$, as well as on the assumption (2.1), we obtain

$$u \in L^2(0, t; W^{2,3}(\mathbb{R}_+^3)).$$

So, by Sobolev embedding and since $\nabla u \in L^2(0, t; L^2(\mathbb{R}_+^3))$, we infer that for all $r \in [2, +\infty)$

$$\nabla u \in L^2(0, t; L^r(\mathbb{R}_+^3)).$$

Coming back to the inequality (4.28) with $(b, r_1) = (3, 2)$, $r_2 \geq 2$ (which means $a = \tilde{a} = 3(r_2 - 2)r_2^{-1} \in [2, 3)$), we now get

$$(u \cdot \nabla)u \in L^{\tilde{a}}(0, t; L^3(\mathbb{R}_+^3)).$$

By Corollary 4.23, we also have $j_f - \rho_f u \in L^{\tilde{a}}(0, t; L^3(\mathbb{R}_+^3))$, therefore an other application of the maximal regularity for the Stokes system implies that for all $\tilde{a} \in [2, 3)$

$$u \in L^{\tilde{a}}(0, t; W^{2,3}(\mathbb{R}_+^3)),$$

if $u_0 \in D_3^{1-\frac{1}{\tilde{a}}, \tilde{a}}(\mathbb{R}_+^3)$. So, under this assumption, Sobolev embedding implies that for all $\tilde{a} \in [2, 3)$ and for all $r \in [2, +\infty]$, we have

$$\nabla u \in L^{\tilde{a}}(0, t; L^r(\mathbb{R}_+^3)).$$

This allows to apply the estimate (4.28) with $a = b = p > 3$ and $\tilde{a} = r_1 \in [2, 3)$ (and also $p \leq r_2$) so that

$$(u \cdot \nabla)u \in L^p(0, t; L^p(\mathbb{R}_+^3)).$$

The relation (4.29) reads as

$$r_2 = \frac{6r_1}{3r_1 + 6 - 4p}, \quad p \leq r_2,$$

therefore this turns into

$$p < \frac{3(r_1 + 2)}{4}.$$

Since $r_1 = 2$ leads to the limiting case $p < 3$, we see that taking $p > 3$ close enough to 3 and $\tilde{a} = r_1 \in (2, 3)$ ensures the previous condition. Thus, the assumption (2.1) is compatible with the previous analysis and this concludes the proof. \square

As in [37, 34, 24], we finally deduce the following non-uniform in time result.

Corollary 4.25. *For any finite strong existence time $t \geq 0$, one has*

$$\nabla u \in L^1(0, t; L^\infty(\mathbb{R}_+^3)).$$

Proof. Let $t > 0$ be a finite strong existence time. We consider the exponent $p > 3$ given in the assumption (2.1) and invoke the Gagliardo-Nirenberg-Sobolev inequality (see Theorem A.3 in the Appendix) which yields

$$\|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} \lesssim \|D^2 u(s)\|_{L^p(\mathbb{R}_+^3)}^{\beta_p} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^{1-\beta_p}, \quad s \in [0, t],$$

where $\beta_p := \frac{5p}{7p-6}$. Combining the energy inequality (1.29) and Lemma 4.12, we get

$$\begin{aligned} \int_0^t \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds &\lesssim \int_0^t \|D^2 u(s)\|_{L^p(\mathbb{R}_+^3)}^{\beta_p} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^{1-\beta_p} ds \\ &\leq \|u\|_{L^\infty(0, t; L^2(\mathbb{R}_+^3))}^{1-\beta_p} \int_0^t \|D^2 u(s)\|_{L^p(\mathbb{R}_+^3)}^{\beta_p} ds \\ &\leq \left(E(0) + E(0)^{\frac{1}{2}} t + t^2 \right)^{\frac{(1-\beta_p)}{2}} \int_0^t \|D^2 u(s)\|_{L^p(\mathbb{R}_+^3)}^{\beta_p} ds \\ &\leq \left(E(0) + E(0)^{\frac{1}{2}} t + t^2 \right)^{\frac{(1-\beta_p)}{2}} t^{1-\frac{p}{\beta_p}} \|D^2 u\|_{L^p(0, t; L^p(\mathbb{R}_+^3))}^{\beta_p}, \end{aligned}$$

where we have used the Hölder inequality in the last line. Furthermore, thanks to the maximal $L^p L^p$ regularity for the Stokes system, we have

$$\|D^2 u\|_{L^p(0, t; L^p(\mathbb{R}_+^3))} \lesssim \|u_0\|_{D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)} + \|j_f - \rho_f u\|_{L^p(0, t; L^p(\mathbb{R}_+^3))} + \|(u \cdot \nabla) u\|_{L^p(0, t; L^p(\mathbb{R}_+^3))} < \infty,$$

thanks to the assumption (2.1), Corollary 4.23 and Proposition 4.27. This allows to conclude the proof. \square

Remark 4.26. By Corollary 6.4 and Corollary 4.25, we get $u \in L^1(0, t; W_0^{1, \infty}(\mathbb{R}_+^3))$ when $t > 0$ is a strong existence time, so that the characteristic curves for the Vlasov equation are classically defined on $(0, t)$ and the representation formula (4.9) can be applied.

5 Exit geometric condition and absorption

The main goal of this section is to describe precisely the effect of the absorption boundary condition (1.11) satisfied by the distribution function f solution to the Vlasov equation (1.1). In short, we will study the time of absorption when one starts from a compact support for the initial distribution function f_0 . The simple geometry of the flat boundary will allow us to base our study upon the characteristics curves for the Vlasov equation.

As explained in the introduction, we rely on different ideas mainly taken from the work of Glass, Han-Kwan and Moussa in [28] (but which lead to different types of results). In some sense, the adaptation to the half-space case is less tedious because we only deal with a flat boundary. Here, the set $B_v(R)$ refers to $B(0, R) \subset \mathbb{R}^3$ with $R \in [0, +\infty]$.

We first introduce the so-called *exit geometric condition*.

Definition 5.1. Let $L, R > 0$. We say that a vector field $U : \mathbb{R}^+ \times \mathbb{R}_+^3 \rightarrow \mathbb{R}^3$ satisfies the exit geometric condition (EGC) in time T with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(R)$ if

$$\sup_{(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(R)} \tau_U^+(0, x, v) < T, \quad (5.1)$$

where τ_U^+ refers to Definition (4.2) for the characteristic curves (X_U, V_U) of the Vlasov equation associated to a velocity field U in (4.1).

A direct consequence of an EGC satisfied by a velocity field is the following.

Proposition 5.2. *Suppose that a velocity field $u \in L_{\text{loc}}^1(\mathbb{R}^+; W_0^{1,\infty}(\mathbb{R}_+^3))$ satisfies an EGC in time $T > 0$ with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(R)$ for some fixed $L, R > 0$. Then, if f is the solution to the Vlasov equation associated to u , we have for almost every $(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3$ and any $t > T$*

$$\begin{aligned} f(t, x, v) &= e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| > R} f_0(X_t^0(x, v), V_t^0(x, v)) \\ &\quad + e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| \leq R} \mathbf{1}_{X_t^0(x, v)_3 > L} f_0(X_t^0(x, v), V_t^0(x, v)). \end{aligned} \quad (5.2)$$

Proof. Let $t > T$. Observe that the regularity of u allows us to manipulate the characteristic curves for the Vlasov equation in a classical sense on $[0, t]$. From the representation formula (4.9), we have

$$f(t, x, v) = e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| \leq R} \mathbf{1}_{X_t^0(x, v)_3 \leq L} f_0(X_t^0(x, v), V_t^0(x, v)) + f^\sharp(t, x, v),$$

where $f^\sharp(t, x, v)$ denotes the expression in the r.h.s of (5.2). We thus have to prove that the first term of the previous equality vanishes. Using the change of variable $(x, v) \mapsto (X(0; t, x, v), V(0; t, x, v))$ together with the fact that

$$(X, V)(0; t, \mathcal{O}^t) = \{(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3 \mid \tau^+(0, x, v) > t\},$$

where

$$\mathcal{O}^t = \{(x, v) \in \mathbb{R}_+^3 \times \mathbb{R}^3 \mid \tau^-(t, x, v) < 0\},$$

we get

$$\begin{aligned} &\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| \leq R} \mathbf{1}_{X_t^0(x, v)_3 \leq L} f_0(X_t^0(x, v), V_t^0(x, v)) \, dx \, dv \\ &= \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} \mathbf{1}_{\tau^+(0, x, v) > t} \mathbf{1}_{|v| \leq R} \mathbf{1}_{x_3 \leq L} f_0(x, v) \, dx \, dv \\ &= \int_{x_3 \leq L, |v| \leq R} \mathbf{1}_{\tau^+(0, x, v) > t} f_0(x, v) \, dx \, dv. \end{aligned}$$

By definition of the EGC in time $T > 0$ with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(R)$, we have $\tau^+(0, x, v) < T$ for all $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(R)$ therefore the last integral is actually zero because $t > T$. Since it is true for all $t > T$, this concludes the proof. \square

The main idea that we follow now is to compare the Vlasov equation with velocity field u (solution to the Navier-Stokes equations) to the "free" Vlasov equation without coupling. We thus consider the following characteristic curves (X^g, V^g) for the Vlasov equation associated with the vector field $(x, v) \mapsto (v, G - v)$, namely

$$\begin{cases} \dot{X}^g(s; t, x, v) = V^g(s; t, x, v), \\ \dot{V}^g(s; t, x, v) = G - V^g(s; t, x, v), \\ X^g(t; t, x, v) = x, \\ V^g(t; t, x, v) = v. \end{cases} \quad (5.3)$$

These are the equations satisfied by the characteristic curves associated to a trivial velocity field (i.e $u = 0$) in the Vlasov equation and where the particles only undergo the effect of the gravity force G , without being coupled to a surrounding fluid. We have the formulas

$$\begin{cases} X^g(t; 0, x, v) = x + (1 - e^{-t})v + (t + e^{-t} - 1)G, \\ V^g(t; 0, x, v) = e^{-t}v + (1 - e^{-t})G, \end{cases} \quad (5.4)$$

and in particular, because $G = (0, 0, -g)$, we have

$$\begin{cases} X^g(t; 0, x, v)_3 = x_3 + (1 - e^{-t})v_3 - (t + e^{-t} - 1)g, \\ V^g(t; 0, x, v)_3 = -g + e^{-t}(v_3 + g). \end{cases}$$

Remark 5.3. In view of the property (4.8), we observe that for the characteristic curves $Z^g := (X^g, V^g)$, we have

$$\left\{ (x, v) \in (\mathbb{R}^2 \times (0, L)) \times \mathbb{R}^3 \mid \tau_g^+(0, x, v) = \infty \text{ or } Z^g(\tau^+(0, x, v); 0, x, v) \in \Sigma^0 \right\} = \emptyset,$$

where τ_g^+ refers to the forward exit time associated to the curves (X^g, V^g) . In addition, if $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times \mathbb{R}^3$ is fixed and if $V^g(t; 0, x, v)_3 < 0$ then $V^g(s; 0, x, v)_3 < 0$ for all $s \geq t \geq 0$.

As we shall see in the end of the current section, quantitative information about the EGC are easily available for the trivial velocity field. The following stability result, which is directly inspired by [28], will thus enable us to show that any velocity field solution to the Navier-Stokes equations satisfies the EGC in some finite time, provided that its L^1L^∞ norm is small enough.

Lemma 5.4. *Let $\alpha > 0$. There exists a constant $\kappa_\alpha > 0$ such that the following holds. Suppose that the trivial vector field (related to (X^g, V^g)) satisfies the EGC with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$ in time $T > 0$, where $L, R > 0$ are given. Then, any vector field $U \in L^1_{\text{loc}}(\mathbb{R}^+; W^{1,\infty}_+(\mathbb{R}^3_+))$ such that*

$$\int_0^{T+\alpha} \|U(s)\|_{L^\infty(\mathbb{R}^3_+)} ds \leq \kappa_\alpha, \quad (5.5)$$

satisfies the EGC in time $T + \alpha$ with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$.

Proof. For all $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$ and $(s, t) \in \mathbb{R}^+ \times \mathbb{R}^+$, we consider $(X_U(s, t, x, v), V_U(s, t, x, v))$ (resp. $(X^g(s, t, x, v), V^g(s, t, x, v))$) the characteristic curves associated to U (resp. to the trivial velocity field). We first set

$$(Y, W) := (X_U - X^g, V_U - V^g),$$

which satisfy the following equations

$$\begin{cases} \dot{Y}(s; 0, x, v) = W(s; 0, x, v), \\ \dot{W}(s; 0, x, v) = (PU)(s, X_U(s; 0, x, v)) - W(s; 0, x, v), \\ Y(0; 0, x, v) = 0, \\ W(0; 0, x, v) = 0. \end{cases}$$

We now fix $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$. We observe that we have for all $t \in \mathbb{R}^+$

$$Y(\tau; 0, x, v) = \int_0^\tau (1 - e^{\tau-t})(PU)(t, X_U(t, 0, x, v)) dt,$$

so that for all $t \in [0, T + \alpha]$

$$|Y(t; 0, x, v)| \leq \int_0^t (1 - e^{\tau-t}) \|(PU)(\tau)\|_{L^\infty(\mathbb{R}^3)} d\tau \leq \int_0^{T+\alpha} \|U(\tau)\|_{L^\infty(\mathbb{R}^3_+)} d\tau, \quad (5.6)$$

because of the property (4.3).

Furthermore, thanks to the EGC satisfied by the trivial velocity field in time T and Remark 5.3, we have $X^g(T; 0, x, v)_3 < 0$ and $V^g(T; 0, x, v)_3 < 0$. Hence,

$$\begin{aligned} X^g(T + \alpha; 0, x, v)_3 &= X^g(T + \alpha; T, X^g(T; 0, x, v)_3, V^g(T; 0, x, v)_3)_3 \\ &= X^g(T; 0, x, v)_3 + (1 - e^{-\alpha})(V^g(T; 0, x, v)_3 + g) - \alpha g \\ &< (1 - \alpha - e^{-\alpha})g < 0, \end{aligned} \quad (5.7)$$

because $\alpha > 0$. If we set $\eta_\alpha := (e^{-\alpha} + \alpha - 1)g > 0$ and $\kappa_\alpha := \eta_\alpha/2$, we see that if U satisfies the condition

$$\int_0^{T+\alpha} \|U(s)\|_{L^\infty(\mathbb{R}^3_+)} ds \leq \kappa_\alpha,$$

then the estimate (5.6) turns into

$$\sup_{t \in [0, T + \alpha]} |Y(t; 0, x, v)| \leq \frac{\eta_\alpha}{2}. \quad (5.8)$$

In view of (5.7) and Remark 5.3, we get the existence of $t_0 = t_0(x, v) \in (0, T + 1]$ such that $X^g(t_0; 0, x, v) \notin \{h \in \mathbb{R}^3 \mid h_3 > -\eta_\alpha\}$. From (5.8), we deduce that $X_U(t_0; 0, x, v) \notin \mathbb{R}_+^3$ and therefore $\tau_U^+(0, x, v) < T + \alpha$. As it is true for any $(\mathbb{R}^2 \times (0, L)) \times B_v(R)$, this means by definition that the EGC is satisfied for U in time $T + \alpha$ with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$. \square

Then, in view of the simple and explicit form of the characteristic curves (X^g, V^g) for the Vlasov equation associated to a trivial velocity field, we can easily obtain precise information on the EGC satisfied for this velocity field. Indeed, for all $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$ and for all $t \geq 0$, we have

$$\begin{aligned} X^g(t; 0, x, v)_3 &= x_3 + (1 - e^{-t})(v_3 + g) - tg \\ &< L + |v| + g - tg \\ &< L + R + g - tg. \end{aligned}$$

It naturally leads to the following definition and properties about the EGC for the trivial vector field and the characteristic curves (X^g, V^g) .

Definition 5.5. We set

$$t_0(L, R) := \frac{L + R + g}{g}, \quad (5.9)$$

$$t_0 := t_0(1, 1). \quad (5.10)$$

Lemma 5.6. *If L, R are given, the trivial vector field $U \equiv 0$ (associated to (X^g, V^g)) satisfies the EGC in time $t_0(R)$ with respect to $(\mathbb{R}^2 \times (0, L)) \times B_v(R)$.*

Proof. Assume that $X^g(t, 0, x, v)_3 > 0$ for all $t \in [0, t_0(L, R)]$ and $(x, v) \in (\mathbb{R}^2 \times (0, L)) \times B_v(0, R)$. Then

$$\begin{aligned} 0 &< x_3 + (1 - e^{-t})(v_3 + g) - tg \\ &< L + R + g - tg \\ &= g(t_0(L, R) - t), \end{aligned}$$

so that we get a contradiction by taking $t = t_0(L, R)$. \square

Remark 5.7. We note that $t_0 = t_0(1, 1)$ does not depend on the initial data. This time only has to be seen as a reference time after which we will use the absorption phenomenon. The subsequent analysis could have been performed by replacing $t_0(1, 1)$ by $t_0(\kappa, \kappa)$ for any $\kappa > 0$.

The following result is, in some sense, of reverse nature: given a time $t > t_0$, we describe which proportion of the initial velocities will lead to absorption before time t . More precisely, we state the following lemma.

Lemma 5.8. *There exist some continuous increasing functions $L_g : [t_0, +\infty) \rightarrow \mathbb{R}^+$ and $R_g : [t_0, +\infty) \rightarrow \mathbb{R}^+$ such that for all $t > t_0$ the trivial velocity vector field (associated to (X^g, V^g)) satisfies the EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1 + L_g(t))) \times B_v(0, 1 + R_g(t))$. Furthermore, there exists $C_g, \underline{C}_g, \overline{C}_g > 0$ such that for all $s > t_0$*

$$\begin{aligned} \frac{1}{1 + L_g(s)} &\leq \frac{C_g}{1 + s}, \\ \frac{\underline{C}_g}{1 + s} &\leq \frac{1}{1 + R_g(s)} \leq \frac{\overline{C}_g}{1 + s}. \end{aligned}$$

Proof. Since the characteristic curve X^g is a vertical line, a necessary and sufficient condition ensuring an EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1 + L_g(t))) \times B_v(0, 1 + R_g(t))$ (for some positive functions L_g and R_g to be determined) is that for all $(x, v) \in (\mathbb{R}^2 \times (0, 1 + L_g(t))) \times B_v(0, 1 + R_g(t))$, the following inequality holds:

$$X^g(t; 0, x, v)_3 = x_3 + (1 - e^{-t})(v_3 + g) - tg < 0. \quad (5.11)$$

Indeed, one can easily show that the function $s \mapsto X^g(s; 0, x, v)_3$ is decreasing on $[\bar{t}, t]$, for some time $\bar{t} \geq t_0$ which is independent of (x, v) . We then set

$$L_g(s) := \frac{1}{2}(sg - (1 - e^{-s})g) - 1, \quad R_g(s) := \frac{1}{2} \left(\frac{sg}{1 - e^{-s}} - g \right) - 1.$$

Now, if $0 < x_3 < 1 + L_g(t)$ and $|v| < 1 + R_g(t)$, we observe that

$$x_3 + (1 - e^{-t})v_3 \leq x_3 + (1 - e^{-t})|v| < tg - (1 - e^{-t})g,$$

therefore we have $X^g(t; 0, x, v)_3 < 0$. It remains to show that L_g and R_g are positive on $[t_0, +\infty[$. As they are nondecreasing functions, we only have to prove that $L_g(t_0) > 0$ and $R_g(t_0) > 0$: recalling the explicit Definition 5.5 of the time $t_0 = 1 + 2g^{-1}$, we have

$$\begin{aligned} L_g(t_0) &= \frac{1}{2} \left(\left(1 + \frac{2}{g}\right)g - (1 - e^{-(1+\frac{2}{g})})g \right) - 1 = \frac{1}{2}e^{-(1+\frac{2}{g})}g > 0, \\ R_g(t_0) &= \frac{1}{2} \left(\frac{(1 + \frac{2}{g})g}{1 - e^{-(1+\frac{2}{g})}} - g \right) - 1 = \frac{2}{1 - e^{-(1+\frac{2}{g})}} - 1 + g \left(\frac{1}{1 - e^{-(1+\frac{2}{g})}} - \frac{1}{2} \right) > 1 + \frac{g}{2} > 0, \end{aligned}$$

which is the desired claim. Concerning the last part, a direct computation shows that for all $s \geq t_0$, we have

$$\frac{1+s}{1+L_g(s)} \leq \frac{1+t_0}{1+L_g(t_0)}, \quad \frac{1+s}{1+R_g(s)} \leq \frac{1+t_0}{1+R_g(t_0)},$$

and that $s \mapsto \frac{1+R_g(s)}{1+s}$ is bounded from above on $[t_0, +\infty)$. The proof is then complete. \square

6 The bootstrap argument

In this section, we provide a proof of Theorem 2.1, relying on the absorption effect highlighted in Section 5. Our strategy is based on a bootstrap argument reminiscent of the ideas of [37, 34]. Roughly speaking, we will prove that one can propagate the controls (4.10) and (5.5) for the velocity field u .

6.1 Initialization of the bootstrap procedure

In order to set up a bootstrap procedure, we introduce the following quantities.

Definition 6.1. We set $T_0 := t_0 + 1$ where t_0 is given in Definition 5.5.

Definition 6.2. We consider $\delta_0 > 0$ which satisfies $\delta_0 e^{\delta_0} < 1/9$ and $\delta_0 < \kappa_{1/2}$, where $\kappa_{1/2}$ is given in Lemma 5.4.

Let (u, f) be a global Leray solution to the Vlasov-Navier-Stokes system with initial condition (u_0, f_0) . We start with the following lemma.

Lemma 6.3. *We have*

$$\int_0^{T_0} \|j_f(s) - \rho_f u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim N_q(f_0)\eta(T_0) [E(0) + E(0)T_0 + T_0^2], \quad (6.1)$$

for some nondecreasing positive continuous function η . Furthermore, under the smallness assumption (2.2), the time T_0 is a strong existence time, in the sense of Definition 4.22.

Proof. We first set $F := j_f - \rho_f u$. From Lemma 4.13 and (4.23) in Proposition 4.19, we infer that

$$\begin{aligned} \int_0^{T_0} \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds &\leq \sup_{s \in [0, T_0]} \|\rho_f(s)\|_{L^\infty(\mathbb{R}_+^3)} [\mathbf{E}(0) + \mathbf{E}(0)T_0 + T_0^2] \\ &\lesssim N_q(f_0)\eta(T_0) [\mathbf{E}(0) + \mathbf{E}(0)T_0 + T_0^2]. \end{aligned}$$

Thanks to the Cauchy-Schwarz inequality, we also have

$$\int_0^{T_0} \|F(s)\|_{L^2(\mathbb{R}_+^3)} ds \leq T_0^{1/2} \left(\int_0^{T_0} \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right)^{1/2},$$

therefore the inequality (6.1) now entails

$$\begin{aligned} \|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \int_0^{T_0} \left[\|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \|F(s)\|_{L^2(\mathbb{R}_+^3)} \right] ds &\lesssim \|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + N_q(f_0)\eta(T_0) [\mathbf{E}(0) + \mathbf{E}(0)T_0 + T_0^2] \\ &\quad + N_q(f_0)^{1/2} T_0^{1/2} \eta(T_0)^{1/2} [\mathbf{E}(0) + \mathbf{E}(0)T_0 + T_0^2]^{1/2}. \end{aligned}$$

If C_\star refers to the universal constant given in Proposition 4.20, we can thus use the smallness assumption (2.2) to ensure that

$$\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \int_0^{T_0} \left[\|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \|F(s)\|_{L^2(\mathbb{R}_+^3)} \right] ds < C_\star,$$

which means that T_0 is a strong existence time. \square

Corollary 6.4. *Under the smallness assumption (2.2), we have*

$$\int_0^{T_0} \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta_0}{2}. \quad (6.2)$$

Proof. By Lemma 6.3, we know that T_0 is a strong existence time therefore the parabolic regularization of the Navier-Stokes equations stated in Proposition 4.20 holds for u on $[0, T_0]$. Namely, we get

$$u \in L^\infty(0, T_0; H_{\text{div}}^1(\mathbb{R}_+^3)) \cap L^2(0, T_0; H^2(\mathbb{R}_+^3)),$$

and there exists $\tilde{C} > 0$ such that for all $t \in [0, T_0]$

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \leq \tilde{C} \left(\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \|j_f - \rho_f u\|_{L^2(0, T_0; L^2(\mathbb{R}_+^3))}^2 \right). \quad (6.3)$$

We then use the Gagliardo-Nirenberg-Sobolev inequality (see Theorem A.3 in the Appendix) with $(p, r, q, j, m) = (\infty, 2, 6, 0, 2)$ so that there exists a universal constant C_0 such that

$$\|u(s)\|_{L^\infty(\mathbb{R}_+^3)} \leq C_0 \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|u(s)\|_{L^6(\mathbb{R}_+^3)}^{1/2}.$$

By Sobolev embedding, we infer from (6.3) that

$$\begin{aligned} \int_0^{T_0} \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds &\leq C_0 \int_0^{T_0} \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^{1/2} ds \\ &\leq C_0 T_0^{3/4} \|\nabla u\|_{L^\infty(0, T_0; L^2(\mathbb{R}_+^3))}^{1/2} \left(\int_0^{T_0} \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right)^{1/4} \\ &\leq C_0 T_0^{3/4} \tilde{C}^{1/16} \left(\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \|j_f - \rho_f u\|_{L^2(0, T_0; L^2(\mathbb{R}_+^3))}^2 \right)^{1/16} \\ &\lesssim C_0 T_0^{3/4} \tilde{C}^{1/16} \left(\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + N_q(f_0)\eta(T_0) [\mathbf{E}(0) + \mathbf{E}(0)T_0 + T_0^2] \right)^{1/16}, \end{aligned}$$

thanks to (6.1). We can now proceed exactly in the same way as in the proof of Lemma 6.3: since T_0 is fixed, we can use the smallness condition (2.2) to ensure that

$$\int_0^{T_0} \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta_0}{2}.$$

This concludes the proof of the Corollary. \square

In order to set up the bootstrap argument, we introduce

$$t^* := \sup \left\{ \text{strong existence times } t \geq T_0 \text{ such that } \int_{T_0}^t \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \delta_0, \int_{T_0}^t \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta_0}{2} \right\}. \quad (6.4)$$

Our main goal is now to show that $t^* = +\infty$.

Lemma 6.5. *We have $t^* > T_0$.*

Proof. In view of Lemma 6.3, we see that using one more time the smallness assumption (2.2), the time $t = T_0 + \varepsilon_0$ is a strong existence time, if $\varepsilon_0 > 0$ is fixed. Now, we observe that a continuity argument, Proposition 4.25 and (4.21) ensure that

$$\int_{T_0}^{T_0+\varepsilon_0} \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \delta_0, \quad \int_{T_0}^{T_0+\varepsilon_0} \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta_0}{2},$$

if ε_0 is small enough. By definition of t^* , this concludes the proof. \square

We then argue by contradiction and shall assume from now on that $t^* < \infty$.

6.2 Absorption and decay in time of the moments

We eventually explain how one can take advantage of the absorption effect at the boundary. Recall first the concept of EGC introduced in Definition 5.1. In view of the bounds satisfied before t^* , the vector field u shall satisfy an EGC at time $t < t^*$ with respect to some compact depending on t . This will imply a (polynomial) decay in time of the moments of f and then of the Brinkman force $j_f - \rho_f u$.

Recalling the notations introduced in Definition 1.3 and Definition 1.4, we state the following fundamental lemma which highlights the precise link between absorption and decay.

Lemma 6.6. *Let $t \in (T_0, t^*)$. Let $k_1, k_2, \ell \in \mathbb{R}^+$ and $q > 3$ such that*

$$q > k_1 + \ell + 3,$$

Suppose that the velocity field u satisfies an EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1 + L(t))) \times B_v(0, 1 + R(t))$ for some nondecreasing functions $L : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ and $R : \mathbb{R}^+ \rightarrow \mathbb{R}^+$. Then we have for almost every $x \in \mathbb{R}_+^3$

$$m_\ell f(t, x) \lesssim \frac{N_q(f_0)}{(1 + R(t))^{k_1}} + \frac{H_{\ell, k_2}(f_0)}{(1 + L(t))^{k_2}}, \quad (6.5)$$

$$\|m_\ell f(t)\|_{L^r(\mathbb{R}_+^3)} \lesssim \frac{K_{q, r}(f_0)}{(1 + R(t))^{k_1}} + \frac{F_{\ell, k_2, r}(f_0)}{(1 + L(t))^{k_2}}, \quad (6.6)$$

where \lesssim only depends on $k_1, k_2, q, \ell, g, \delta_0$.

Proof. The whole proof is based on the fact that

$$f(t, x, v) = f^{\natural}(t, x, v) + f^{\flat}(t, x, v), \quad (6.7)$$

where

$$f^{\natural}(t, x, v) := e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| > 1 + R(t)} f_0(X_t^0(x, v), V_t^0(x, v)), \quad (6.8)$$

$$f^{\flat}(t, x, v) := e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| \leq 1 + R(t)} \mathbf{1}_{X_t^0(x, v)_3 > 1 + L(t)} f_0(X_t^0(x, v), V_t^0(x, v)). \quad (6.9)$$

Indeed, because of the EGC satisfied by u in time t with respect to $(\mathbb{R}^2 \times (0, 1 + L(t))) \times \mathbb{B}_v(0, 1 + R(t))$, we can apply Proposition 5.2 so that (6.7) holds. We then study separately the decay of the moments of f^{\natural} and f^{\flat} .

• For f^{\natural} defined in (6.8), we use the admissible change of variable $v \mapsto V_t^0(x, v) = \Gamma_{t,x}(v)$ given by Lemma 4.5, whose Jacobian inverse is bounded by $2e^{-3t}$. We get

$$\begin{aligned} m_{\ell} f^{\natural}(t, x) &= \int_{\mathbb{R}^3} |v|^{\ell} f^{\natural}(t, x, v) \, dv \\ &= \int_{\mathbb{R}^3} |v|^{\ell} e^{3t} \mathbf{1}_{\tau^-(t, x, v) < 0} \mathbf{1}_{|V_t^0(x, v)| > 1 + R(t)} f_0(X_t^0(x, v), V_t^0(x, v)) \, dv \\ &= \int_{\mathbb{R}^3} |\Gamma_{t,x}^{-1}(w)|^{\ell} e^{3t} \mathbf{1}_{\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0} \mathbf{1}_{|w| > 1 + R(t)} f_0(\Lambda_{t,w}(x), w) |\det D_w \Gamma_{t,x}^{-1}(w)| \, dw \\ &\lesssim \frac{1}{(1 + R(t))^{k_1}} \int_{\mathbb{R}^3} |w|^{k_1} |\Gamma_{t,x}^{-1}(w)|^{\ell} \mathbf{1}_{\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0} f_0(\Lambda_{t,w}(x), w) \, dw, \end{aligned}$$

where we have used the notations of Lemma 4.7, namely $\Lambda_{t,w}(x) = X_t^0(x, \Gamma_{t,x}^{-1}(w))$. Thanks to the inequality (4.12), we also have

$$\begin{aligned} |\Gamma_{t,x}^{-1}(w)| &\leq |w| + (1 - e^{-t})g + \int_0^t e^{\tau-t} \|Pu(\tau)\|_{L^{\infty}(\mathbb{R}^3)} \, d\tau, \\ &\leq |w| + g + \delta_0, \end{aligned}$$

and because of the Definition (4.7) of $\tau^-(t, x, \Gamma_{t,x}^{-1}(w))$, we see that if $\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0$ then $\Lambda_{t,w}(x) \in \mathbb{R}_+^3$. Therefore we obtain

$$\begin{aligned} m_{\ell} f^{\natural}(t, x) &\lesssim \frac{1}{(1 + R(t))^{k_1}} \int_{\mathbb{R}^3} [|w|^{k_1 + \ell} + |w|^{k_1}] \mathbf{1}_{\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0} f_0(\Lambda_{t,w}(x), w) \, dw \\ &= \frac{1}{(1 + R(t))^{k_1}} \int_{\mathbb{R}^3} \frac{|w|^{k_1 + \ell} + |w|^{k_1}}{1 + |w|^q} \mathbf{1}_{\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0} (1 + |w|^q) f_0(\Lambda_{t,w}(x), w) \, dw \\ &\leq \frac{N_q(f_0)}{(1 + R(t))^{k_1}} \int_{\mathbb{R}^3} \frac{|w|^{k_1 + \ell} + |w|^{k_1}}{1 + |w|^q} \, dw \\ &\lesssim \frac{N_q(f_0)}{(1 + R(t))^{k_1}}, \end{aligned}$$

for $q > k_1 + \ell + 3$. This entails the contribution from f^{\natural} in the inequality (6.5). Concerning its contribution to the inequality (6.6), we combine the previous change of variable in velocity with the Minkowski's integral

inequality (see [38, Theorem 202]) and the bound (4.12) to write with the notations of Lemma 4.7

$$\begin{aligned}
\|m_\ell f^\sharp(t)\|_{L^r(\mathbb{R}_+^3)} &= \left[\int_{\mathbb{R}_+^3} \left(\int_{\mathbb{R}^3} |v|^\ell f^\sharp(t, x, v) \, dv \right)^r dx \right]^{1/r} \\
&\lesssim \frac{1}{(1 + \mathbf{R}(t))^{k_1}} \left[\int_{\mathbb{R}_+^3} \left(\int_{\mathbb{R}^3} |w|^{k_1} |\Gamma_{t,x}^{-1}(w)|^\ell \mathbf{1}_{\tau^-(t,x,\Gamma_{t,x}^{-1}(w)) < 0} f_0(\Lambda_{t,w}(x), w) \, dw \right)^r dx \right]^{1/r} \\
&\leq \frac{1}{(1 + \mathbf{R}(t))^{k_1}} \int_{\mathbb{R}^3} |w|^{k_1} \left(\int_{\mathbb{R}_+^3} |\Gamma_{t,x}^{-1}(w)|^{r\ell} \mathbf{1}_{\tau^-(t,x,\Gamma_{t,x}^{-1}(w)) < 0} f_0(\Lambda_{t,w}(x), w)^r dx \right)^{1/r} dw \\
&\leq \frac{1}{(1 + \mathbf{R}(t))^{k_1}} \int_{\mathbb{R}^3} [|w|^{k_1+\ell} + |w|^{k_1}] \left(\int_{\mathbb{R}_+^3} \mathbf{1}_{x \in \mathbb{R}_+^3} \mathbf{1}_{\tau^-(t,x,\Gamma_{t,x}^{-1}(w)) < 0} f_0(\Lambda_{t,w}(x), w)^r dx \right)^{1/r} dw.
\end{aligned}$$

Performing the admissible change of variable $x \mapsto \Lambda_{t,w}(x)$ in the interior integral, with a Jacobian inverse bounded by 2 (see Lemma 4.7), we get

$$\begin{aligned}
\|m_\ell f^\sharp(t)\|_{L^r(\mathbb{R}_+^3)} &\lesssim \frac{1}{(1 + \mathbf{R}(t))^{k_1}} \int_{\mathbb{R}^3} [|w|^{k_1+\ell} + |w|^{k_1}] \left(\int_{\mathbb{R}^3} \mathbf{1}_{\Lambda_{t,w}^{-1}(x) \in \mathbb{R}_+^3} \mathbf{1}_{\tau^-\left(t, \Lambda_{t,w}^{-1}(x), \Gamma_{t, \Lambda_{t,w}^{-1}(x)}^{-1}(w)\right) < 0} f_0(x, w)^r dx \right)^{1/r} dw \\
&\lesssim \frac{1}{(1 + \mathbf{R}(t))^{k_1}} \int_{\mathbb{R}^3} \frac{|w|^{k_1+\ell} + |w|^{k_1}}{1 + |w|^q} (1 + |w|^q) \|f_0(\cdot, w)\|_{L^r(\mathbb{R}_+^3)} dw \\
&\leq \frac{K_{q,r}(f_0)}{(1 + \mathbf{R}(t))^{k_1}} \int_{\mathbb{R}^3} \frac{|w|^{k_1+\ell} + |w|^{k_1}}{1 + |w|^q} dw \\
&\lesssim \frac{K_{q,r}(f_0)}{(1 + \mathbf{R}(t))^{k_1}},
\end{aligned}$$

for $q > k_1 + \ell + 3$.

• For the term f^b defined in (6.9), we also perform the change of variable $v \mapsto V_t^0(x, v) = \Gamma_{t,x}(v)$ which entails

$$\begin{aligned}
m_\ell f^b(t, x) &= \int_{\mathbb{R}^3} |v|^\ell f^b(t, x, v) \, dv \\
&= \int_{\mathbb{R}^3} |v|^\ell e^{3t} \mathbf{1}_{\tau^-(t,x,v) < 0} \mathbf{1}_{|V_t^0(x,v)| \leq 1 + \mathbf{R}(t)} \mathbf{1}_{X_t^0(x,v)_3 > 1 + \mathbf{L}(t)} f_0(X_t^0(x, v), V_t^0(x, v)) \, dv \\
&= \int_{\mathbb{R}^3} |\Gamma_{t,x}^{-1}(w)|^\ell e^{3t} \mathbf{1}_{\tau^-(t,x,\Gamma_{t,x}^{-1}(w)) < 0} \mathbf{1}_{|w| \leq 1 + \mathbf{R}(t)} \mathbf{1}_{\Lambda_{t,w}(x)_3 > 1 + \mathbf{L}(t)} f_0(\Lambda_{t,w}(x), w) |\det D_w \Gamma_{t,x}^{-1}(w)| \, dw \\
&\lesssim \frac{1}{(1 + \mathbf{L}(t))^{k_2}} \int_{\mathbb{R}^3} (|w|^\ell + 1) \mathbf{1}_{|w| \leq 1 + \mathbf{R}(t)} \mathbf{1}_{\tau^-(t,x,\Gamma_{t,x}^{-1}(w)) < 0} |\Lambda_{t,w}(x)_3|^{k_2} f_0(\Lambda_{t,w}(x), w) \, dw,
\end{aligned}$$

so that

$$m_\ell f^b(t, x) \lesssim \frac{H_{\ell, k_2}(f_0)}{(1 + \mathbf{L}(t))^{k_2}}.$$

Here, we have again used the fact that if $\tau^-(t, x, \Gamma_{t,x}^{-1}(w)) < 0$ then $\Lambda_{t,w}(x) \in \mathbb{R}_+^3$. This provides the contribution of f^b to the estimate (6.5). To get its contribution to the inequality (6.6), we perform exactly the same

computations as for f^{\flat} to write

$$\begin{aligned}
\|m_{\ell}f^{\flat}(t)\|_{L^r(\mathbb{R}_+^3)} &= \left[\int_{\mathbb{R}_+^3} \left(\int_{\mathbb{R}^3} |v|^{\ell} f^{\flat}(t, x, v) \, dv \right)^r dx \right]^{1/r} \\
&\lesssim \frac{1}{(1+L(t))^{k_2}} \int_{|w| \leq 1+R(t)} [|w|^{\ell} + 1] \\
&\quad \left(\int_{\mathbb{R}^3} \mathbf{1}_{\Lambda_{t,w}^{-1}(x) \in \mathbb{R}_+^3} \mathbf{1}_{\tau^{-}\left(t, \Lambda_{t,w}^{-1}(x), \Gamma_{t, \Lambda_{t,w}^{-1}(x)}^{-1}(w)\right) < 0} |x_3|^{k_2 r} f_0(x, w)^r dx \right)^{1/r} \\
&\lesssim \frac{1}{(1+L(t))^{k_2}} \int_{|w| \leq 1+R(t)} [|w|^{\ell} + 1] \|x_3^{k_2} f_0(\cdot, w)\|_{L^r(\mathbb{R}_+^3)} dw,
\end{aligned}$$

so that

$$\|m_{\ell}f^{\flat}(t)\|_{L^r(\mathbb{R}_+^3)} \lesssim \frac{F_{\ell, k_2, r}(f_0)}{(1+L(t))^{k_2}}.$$

Gathering all the pieces together, we have proven inequalities (6.5) and (6.6). This achieves the proof. \square

Of course, the previous lemma is only available if we could ensure that a certain EGC is satisfied by the velocity field u at time t . Thanks to Lemma 5.8 stated in Section 5, we shall be able to obtain such a condition.

Proposition 6.7. *For any $t \in (T_0, t^*)$, the velocity field u satisfies an EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1+L(t)) \times B_v(0, 1+R(t)))$ for some nondecreasing continuous functions $L : [T_0, +\infty) \rightarrow \mathbb{R}^+$ and $R : [T_0, +\infty) \rightarrow \mathbb{R}^+$ satisfying for all $s > T_0$*

$$\frac{1}{1+L(s)} \lesssim \frac{1}{1+s}, \quad \frac{1}{1+R(s)} \lesssim \frac{1}{1+s},$$

where \lesssim only depends on T_0 and g .

In particular, if $k_1, k_2, \ell \in \mathbb{R}^+$ and $q > 3$ satisfy

$$q > k_1 + \ell + 3,$$

then

$$m_{\ell}f(t, x) \lesssim \frac{N_q(f_0)}{(1+t)^{k_1}} + \frac{H_{\ell, k_2}(f_0)}{(1+t)^{k_2}}, \quad (6.10)$$

$$\|m_{\ell}f(t)\|_{L^r(\mathbb{R}_+^3)} \lesssim \frac{K_{q,r}(f_0)}{(1+t)^{k_1}} + \frac{F_{\ell, k_2, r}(f_0)}{(1+t)^{k_2}}, \quad (6.11)$$

where \lesssim only depends on $k_1, k_2, q, \ell, g, \delta_0$ and T_0 .

Proof. Thanks to Lemma 5.8, we know that if $t - 1/2 > t_0$, the trivial velocity field satisfies an EGC in time $t - 1/2$ with respect to $(\mathbb{R}^2 \times (0, 1 + \tilde{L}(t - 1/2)) \times B_v(0, 1 + \tilde{R}(t - 1/2)))$ for some continuous nondecreasing functions $\tilde{L} : [t_0, +\infty) \rightarrow \mathbb{R}^+$ and $\tilde{R} : [t_0, +\infty) \rightarrow \mathbb{R}^+$ such that for any $s > t_0$

$$\frac{1}{1+\tilde{L}(s)} \lesssim \frac{1}{1+s}, \quad \frac{1}{1+\tilde{R}(s)} \lesssim \frac{1}{1+s}, \quad (6.12)$$

where \lesssim depends on t_0 and g . Furthermore, by definition of δ_0 , we have

$$\int_0^t \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \delta_0 < \kappa_{1/2},$$

so that, owing to Lemma 5.4, we get the fact that the velocity field u satisfies an EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1 + \tilde{L}(t - 1/2)) \times B_v(0, 1 + \tilde{R}(t - 1/2)))$. We have $t > T_0 = t_0 + 1 > t_0 + 1/2$ so that if we set $L(s) := \tilde{L}(s - 1/2)$ and $R(s) := \tilde{R}(s - 1/2)$, we observe that the velocity field u satisfies an EGC in time t with respect to $(\mathbb{R}^2 \times (0, 1 + L(t)) \times B_v(0, 1 + R(t)))$. Furthermore, by (6.12), for all $s > t_0 + 1/2$

$$\frac{1}{1 + L(s)} \lesssim \frac{1}{1 + s}, \quad \frac{1}{1 + R(s)} \lesssim \frac{1}{1 + s}.$$

The last part of the statement readily comes from the previous estimates and the inequalities (6.5)-(6.6) in Lemma 6.6. \square

We now fix $T \in (T_0, t^*)$. Since the decay provided by the absorption can only be considered after time T_0 , we will split the study of the decay estimates between $[0, T_0]$ and $[T_0, T]$. In particular, the local in time estimates of Section 4 will help us to treat the part on $[0, T_0]$.

In view of the assumptions (2.1) and (2.2) on the initial data, we will not be very precise about the exact range of exponents which are used when the quantities $N_q(f_0), H_{q,m}(f_0), K_{q,r}(f_0), F_{q,m,r}(f_0)$ of Definition 1.4 may appear. Note that they are nondecreasing when the parameters q and m increase. In the last part of the current section, we will consider the maximum of the finite family of indices we have used previously. Moreover, these indices may sometimes depend on the exponent γ that we use to quantify the decay in time but this one will be somehow fixed in the last step of the bootstrap. Thus, all the following results must be understood as if we take the involved exponents large enough.

Corollary 6.8. *There exists $\varepsilon > 0$ small enough such that the following holds. Let $p \in (3, 3 + \varepsilon)$. There exists nondecreasing positive functions φ_2 and φ_p vanishing at 0 such that for any $k > 0$, we have*

$$\forall t \in [0, T], \quad \|j_f(t) - \rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} \lesssim_0 \frac{\varphi_2(N_q(f_0) + H_{0,k}(f_0))}{(1 + t)^{\frac{k}{2}}}, \quad (6.13)$$

and if $\|u\|_{L^\infty(0, T; L^6(\mathbb{R}_+^3))} \lesssim_0 1$, we have

$$\forall t \in [0, T], \quad \|j_f(t) - \rho_f u(t)\|_{L^p(\mathbb{R}_+^3)} \lesssim_0 \frac{\varphi_p(N_q(f_0) + H_{0,k}(f_0))}{(1 + t)^{k \frac{p-1}{p}}}, \quad (6.14)$$

whenever $q > k + 3$.

Proof. We separate the study between $[0, T_0]$ and $[T_0, T]$.

- If $t \in [T_0, T]$, we use Lemma 4.10 and get

$$\begin{aligned} \| (j_f - \rho_f u)(t) \|_{L^p(\mathbb{R}_+^3)} &\leq \| \rho_f(t) \|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v - u(t, x)|^p dx dv \right)^{1/p} \\ &\leq \| \rho_f(t) \|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v|^p dx dv + \int_{\mathbb{R}_+^3} \rho_f(t, x) |u(t, x)|^p dx \right)^{1/p}. \end{aligned}$$

We first focus on the estimate (6.13). We have

$$\begin{aligned} \| (j_f - \rho_f u)(t) \|_{L^2(\mathbb{R}_+^3)} &\leq \| m_0 f(t) \|_{L^\infty(\mathbb{R}_+^3)}^{\frac{1}{2}} \left(\| m_2 f(t) \|_{L^1(\mathbb{R}_+^3)} + \| m_0 f(t) \|_{L^\infty(\mathbb{R}_+^3)} E(t) \right)^{1/2} \\ &\lesssim \| m_0 f(t) \|_{L^\infty(\mathbb{R}_+^3)}^{\frac{1}{2}} \left(\| m_2 f(t) \|_{L^1(\mathbb{R}_+^3)} \right. \\ &\quad \left. + \| m_0 f(t) \|_{L^\infty(\mathbb{R}_+^3)} \left[E(0) + E(0)^{1/2} T_0 + T_0^2 + \int_{T_0}^T \| j_f(\tau) \|_{L^1(\mathbb{R}_+^3)} d\tau \right] \right)^{1/2}, \end{aligned}$$

where we have used the energy inequality 1.29 and Lemma 4.12. The estimates (6.10) and (6.11) of Proposition 6.7 then imply

$$\begin{aligned} \|(j_f - \rho_f u)(t)\|_{L^2(\mathbb{R}_+^3)} &\leq \left(\frac{N_q(f_0)}{(1+t)^k} + \frac{H_{0,k}(f_0)}{(1+t)^k} \right)^{1/2} \\ &\quad \times \left(\frac{K_{q_1,1}(f_0)}{(1+t)^{k_1}} + \frac{F_{2,k_1,2}(f_0)}{(1+t)^{k_1}} + \left(\frac{N_{q_1}(f_0)}{(1+t)^{k_1}} + \frac{H_{0,k_1}(f_0)}{(1+t)^{k_1}} \right) \right. \\ &\quad \left. \times \left[E(0) + E(0)^{1/2}T_0 + T_0^2 + \int_{T_0}^T \left(\frac{K_{q_2,1}(f_0)}{(1+\tau)^{k_2}} + \frac{F_{1,k_2,1}(f_0)}{(1+\tau)^{k_2}} \right) d\tau \right] \right)^{1/2}, \end{aligned}$$

provided that $q > k + 3$, $q_1 > k_1 + 4$ and $q_2 > k_2 + 4$. Thus, we obtain for all $t \in [T_0, T]$

$$\|(j_f - \rho_f u)(t)\|_{L^2(\mathbb{R}_+^3)} \lesssim_0 \frac{N_q(f_0)^{1/2}}{(1+t)^{k/2}} + \frac{H_{0,k}(f_0)^{1/2}}{(1+t)^{k/2}}.$$

We now treat the estimate (6.14), assuming $\|u\|_{L^\infty(0,T;L^6(\mathbb{R}_+^3))} \lesssim_0 1$. By Hölder inequality, we have

$$\begin{aligned} \|\mathbb{P}(j_f - \rho_f u)(t)\|_{L^p(\mathbb{R}_+^3)} &\leq \|\rho_f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v - u(t, x)|^p dx dv \right)^{1/p} \\ &\leq \|\rho_f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) |v|^p dx dv + \int_{\mathbb{R}_+^3} \rho_f(t, x) |u(t, x)|^p dx \right)^{1/p} \\ &\leq \|m_0 f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{\frac{p-1}{p}} \left(\|m_p f(t)\|_{L^1(\mathbb{R}_+^3)} + \|u\|_{L^\infty(0,T;L^6(\mathbb{R}_+^3))}^p \|m_0 f(t)\|_{L^{r_\varepsilon}(\mathbb{R}_+^3)} \right)^{1/p}, \end{aligned}$$

where $2 < r_\varepsilon < 6/(3 - \varepsilon)$. Thanks to $\|u\|_{L^\infty(0,T;L^6(\mathbb{R}_+^3))} \lesssim_0 1$ and the estimates (6.10)-(6.11) of Proposition 6.7, we obtain in the same way as before that

$$\|(j_f - \rho_f u)(t)\|_{L^p(\mathbb{R}_+^3)} \lesssim_0 \frac{N_q(f_0)^{\frac{p-1}{p}}}{(1+t)^{k\frac{p-1}{p}}} + \frac{H_{0,k}(f_0)^{\frac{p-1}{p}}}{(1+t)^{k\frac{p-1}{p}}},$$

for $q > k + 3$ and $t \in [T_0, T]$.

• If $t \in [0, T_0]$, and for the proof of (6.13), we use the triangular inequality and the Hölder inequality to write

$$\begin{aligned} \|j_f(t) - \rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} &\leq \|j_f(t)\|_{L^2(\mathbb{R}_+^3)} + \|\rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} \\ &\leq \|j_f(t)\|_{L^2(\mathbb{R}_+^3)} + \|\rho_f(t)\|_{L^\infty(\mathbb{R}_+^3)} \|u(t)\|_{L^2(\mathbb{R}_+^3)}. \end{aligned}$$

Hence, by interpolation, Lemma 4.12 and Lemma 4.17, we get

$$\begin{aligned} \|j_f(t) - \rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} &\leq \|j_f(t)\|_{L^{3/2}(\mathbb{R}_+^3)}^{3/4} \|j_f(t)\|_{L^\infty(\mathbb{R}_+^3)}^{1/4} + \|\rho_f(t)\|_{L^\infty(\mathbb{R}_+^3)} [E(0) + E(0)^{1/2}T_0 + T_0^2] \\ &\leq \phi(T_0) \|j_f\|_{L^\infty(0,T_0;L^\infty(\mathbb{R}_+^3))}^{1/4} + \|\rho_f\|_{L^\infty(0,T_0;L^\infty(\mathbb{R}_+^3))} [E(0) + E(0)^{1/2}T_0 + T_0^2]. \end{aligned}$$

The local estimates (4.23) then provides

$$\begin{aligned} \|j_f(t) - \rho_f u(t)\|_{L^2(\mathbb{R}_+^3)} &\lesssim \phi(T_0) (N_q(f_0) \eta(T_0))^{1/4} + (N_q(f_0) \eta(T_0)) [E(0) + E(0)^{1/2}T_0 + T_0^2] \\ &\lesssim_0 N_q(f_0)^{1/4} + N_q(f_0) \\ &\lesssim_0 \frac{\varphi(N_q(f_0))}{(1+T_0)^{k/2}}, \end{aligned}$$

for some nonincreasing function φ vanishing at 0, provided that $N_q(f_0)$ is taken small enough by the smallness assumption (2.2). Concerning the estimate (6.14) under the assumption $\|u\|_{L^\infty(0,T;L^6(\mathbb{R}_+^3))} \lesssim_0 1$, we proceed as before by writing

$$\begin{aligned} \|j_f(t) - \rho_f u(t)\|_{L^p(\mathbb{R}_+^3)} &\leq \|j_f(t)\|_{L^p(\mathbb{R}_+^3)} + \|\rho_f u(t)\|_{L^p(\mathbb{R}_+^3)} \\ &\leq \|j_f(t)\|_{L^p(\mathbb{R}_+^3)} + \|\rho_f(t)\|_{L^{\tilde{r}_\varepsilon}(\mathbb{R}_+^3)} \|u(t)\|_{L^6(\mathbb{R}_+^3)}, \end{aligned}$$

where $6 < \tilde{r}_\varepsilon < \frac{18+6\varepsilon}{3-\varepsilon}$. By interpolation, this turns into

$$\begin{aligned} \|(j_f - \rho_f u)\|_{L^p(0,T_0;L^p(\mathbb{R}_+^3))}^p &\leq \|j_f\|_{L^\infty(0,T_0;L^{3/2}(\mathbb{R}_+^3))}^{3/2} \|j_f\|_{L^\infty(0,T_0;L^\infty(\mathbb{R}_+^3))}^{p-3/2} \\ &\quad + \|\rho_f\|_{L^\infty(0,T_0;L^1(\mathbb{R}_+^3))}^{p/\tilde{r}_\varepsilon} \|\rho_f\|_{L^\infty(0,T_0;L^\infty(\mathbb{R}_+^3))}^{p-p/\tilde{r}_\varepsilon} \|u\|_{L^\infty(0,T_0;L^6(\mathbb{R}_+^3))}^p \\ &\lesssim_0 N_q(f_0)^{p-3/2} + N_q(f_0)^{p-p/\tilde{r}_\varepsilon}, \end{aligned}$$

thanks to the Lemma 4.17, Lemma 4.9, the bound (4.23) and $\|u\|_{L^\infty(0,T;L^6(\mathbb{R}_+^3))} \lesssim_0 1$. By taking $N_q(f_0)$ small enough, we get for all $t \in [0, T_0]$

$$\|j_f(t) - \rho_f u(t)\|_{L^p(\mathbb{R}_+^3)} \lesssim_0 \frac{\varphi(N_q(f_0))}{(1+T_0)^{k\frac{p-1}{p}}}.$$

Combining the estimates on $[0, T_0]$ and on $[T_0, T]$, we finally obtain the result. \square

Corollary 6.9. *For all $0 \leq t \leq T$ and for any $q > k + 3$ and $k > 1$, we have*

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_0^t \|\mathbb{D}^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim_0 \|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \varphi(N_q(f_0) + H_{0,k}(f_0)). \quad (6.15)$$

Furthermore, the following estimate holds for all $0 \leq t \leq T$

$$\|u\|_{L^\infty(0,t;L^6(\mathbb{R}_+^3))}^2 \lesssim \|\nabla u\|_{L^\infty(0,t;L^2(\mathbb{R}_+^3))}^2 \lesssim_0 \|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \varphi(N_q(f_0) + H_{0,k}(f_0)), \quad (6.16)$$

and in particular, the estimate (6.14) holds true.

Proof. As we have $T \in (T_0, t^*)$ and in view of (6.1), we only treat the case where $t \in [T_0, T]$. Since the time t is a strong existence time (see Definition 4.22), Proposition 4.20 entails the following estimate

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_0^t \|\mathbb{D}^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|j_f - \rho_f u\|_{L^2(0,t;L^2(\mathbb{R}_+^3))}^2.$$

Applying Corollary 6.8 with $k > 1$ leads to

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_0^t \|\mathbb{D}^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim_0 \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \varphi(N_q(f_0) + H_{0,k}(f_0)),$$

which is the first part of the result. The other statement is then a consequence of Sobolev embedding. \square

6.3 Weighted in time estimates

We recall that we have fixed $T \in (T_0, t^*)$. As in [34], the guiding line in order to obtain polynomial weighted estimates for the fluid velocity u is to derive a Stokes system satisfied by a weighted version of u , in such a way that the maximal regularity result of Section A.4 can be applied. Note that since we do not have enough information about the regularity of $u(T_0)$ with respect to the interpolation spaces defined in (A.7), we rely on this result on the whole interval $[0, T]$.

In what follows, the function Ψ is a generic continuous positive function cancelling at 0 which may also depend on quantities involving the initial data but increasing when these quantities increase: for simplicity, this dependency will not be explicitly written down.

We first state the following L^2L^2 regularity result which is reminiscent of [34].

Lemma 6.10. *For all $\gamma \in (0, 3/4)$, we have*

$$\|(1+t)^\gamma \mathbb{D}^2 u\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)) + \Psi\left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2\right) + \|u_0\|_{H^1(\mathbb{R}_+^3)}. \quad (6.17)$$

Proof. Let $\gamma \in (0, 3/4)$ and set $U := (1+t)^\gamma u$. We observe that U satisfies the following Stokes system on $[0, T]$

$$\begin{cases} \partial_t U + A_2 U = S(u, f), \\ U|_{t=0} = u_0, \end{cases} \quad (6.18)$$

where A_2 stands for the Stokes operator on $L_{\text{div}}^2(\mathbb{R}_+^3)$ and where

$$S(u, f) := (1+t)^\gamma \mathbb{P}(j_f - \rho_f u) - (1+t)^\gamma \mathbb{P}(u \cdot \nabla) u + \gamma(1+t)^{\gamma-1} u.$$

According to the maximal regularity result in $L^2 L^2$ for the Stokes system (see Section A.4 in the Appendix), we get

$$\|\partial_t U\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} + \|\mathbb{D}^2 U\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} \lesssim \|S(u, f)\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} + \|u_0\|_{H^1(\mathbb{R}_+^3)},$$

where \lesssim is independent of T , u and f .

We now estimate each term in the modified source term $S(u, f)$. Owing to Corollary 6.8 with $k = 7/2$ and Theorem 3.1, we have

$$(1+t)^{\gamma-1} \|u(t)\|_{L^2(\mathbb{R}_+^3)} \leq \frac{\Psi\left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2\right)}{(1+t)^{3/4-\gamma+1}}, \quad (6.19)$$

which is bounded independently of T in $L^2(0, T)$ because $\gamma \in (0, 3/4)$. For the term coming from the Brinkman force $j_f - \rho_f u$, we use the continuity of the Leray projection on $L^2(\mathbb{R}_+^3)$ and the estimate (6.13) of Corollary 6.8 which entail

$$\begin{aligned} \|(1+t)^\gamma \mathbb{P}(j_f - \rho_f u)\|_{L^2(0,T;L^2(\mathbb{R}_+^3))}^2 &\lesssim_0 \int_0^T (1+s)^{2\gamma} \|j_f(s) - \rho_f u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \\ &\lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)) \int_0^T \frac{1}{(1+s)^{k-2\gamma}} ds \\ &\lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)), \end{aligned} \quad (6.20)$$

for $q > k + 3 > 4 + 3/2$. For the convection term $(u \cdot \nabla)u$, we proceed as in the proof of Theorem A.8 in the Appendix (more precisely, see (A.22)-(A.23)) to obtain, with the Young inequality, that

$$\|(u \cdot \nabla) u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \leq C\varepsilon^{-4} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^4 \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^4 + \frac{3\varepsilon^{4/3}}{4} \|\mathbb{D}^2 u(t)\|_{L^2(\mathbb{R}_+^3)}^2,$$

where $\varepsilon > 0$ can be taken as small as we want. We then use the Gagliardo-Nirenberg-Sobolev inequality to write

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)} \lesssim \|u(t)\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|\mathbb{D}^2 u(t)\|_{L^2(\mathbb{R}_+^3)}^{1/2},$$

therefore this yields

$$\|(1+t)^\gamma (u \cdot \nabla) u\|_{L^2(0,T;L^2(\mathbb{R}_+^3))}^2 \leq \int_0^T (1+t)^{2\gamma} \left[C\varepsilon^{-4} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^4 \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^4 + \frac{3\varepsilon^{4/3}}{4} \right] \|\mathbb{D}^2 u(t)\|_{L^2(\mathbb{R}_+^3)}^2 dt.$$

Furthermore, owing to Theorem 3.1 and the estimate (6.16), we can write

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^4 \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^4 &\leq \Psi\left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2\right)^4 \|\nabla u\|_{L^\infty(0,t;L^2(\mathbb{R}_+^3))}^4 \\ &\lesssim_0 \Psi\left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2\right)^4. \end{aligned}$$

Using the smallness assumption (2.2), we can first choose ε and then $\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2$ small enough so that

$$\|(1+t)^\gamma \mathbb{P}(u \cdot \nabla) u\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} \leq \frac{1}{2} \|(1+t)^\gamma \mathbb{D}^2 u\|_{L^2(0,T;L^2(\mathbb{R}_+^3))}. \quad (6.21)$$

As $\partial_t U = (1+t)^\gamma \partial_t u + \gamma(1+t)^{\gamma-1} u$, we combine (6.19), (6.22) and use an absorption argument stemming from (6.21) so that at the end of the day, we get

$$\|(1+t)^\gamma \mathbb{D}^2 u\|_{L^2(0,T;L^2(\mathbb{R}_+^3))} \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)},$$

which concludes the proof. \square

By interpolation the previous weighted L^2 maximal parabolic estimates allow us to obtain improved L^p estimates of the weighted source term $S(u, f)$ for $p > 3$. We refer to [34] where the proof of the two following results can be found in the whole space case and apply *mutadis mutandis* to the half-space case thanks to Theorem 3.1, Corollary 6.8 and the estimate (6.17).

Corollary 6.11. *Let $p > 3$. For all $\gamma \in (0, \frac{17}{8} - \frac{7}{4p})$, we have*

$$\|(1+t)^{\gamma-1} u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)}.$$

Corollary 6.12. *There exists $\varepsilon > 0$ such that for all $p \in (3, 3 + \varepsilon)$, the following holds. For any $t \in (T_0, T)$, we have for all $\gamma \geq 0$*

$$\|(1+t)^\gamma (u \cdot \nabla) u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \lesssim_0 \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) \|(1+t)^\gamma \mathbb{D}^2 u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))}.$$

We are now in position to provide some $L^p L^p$ estimates for the weighted source term

$$S(u, f) = (1+t)^\gamma \mathbb{P}(j_f - \rho_f u) - (1+t)^\gamma \mathbb{P}(u \cdot \nabla) u + \gamma(1+t)^{\gamma-1} u,$$

with $p > 3$ and $\gamma > 0$ large enough, thanks to the next lemma.

Lemma 6.13. *There exists $\varepsilon > 0$ such that for all $p \in (3, 3 + \varepsilon)$, the following holds. For all $\gamma \in (0, \frac{17}{8} - \frac{7}{4p})$, we have*

$$\|S(u, f)\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \lesssim_0 \left[\varphi(N_q(f_0) + H_{0,k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)} \right] \times \left[1 + \|(1+t)^\gamma \mathbb{D}^2 u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \right].$$

Proof. We first use Corollary 6.11 and Corollary 6.12 to estimate the contribution of $(1+t)^\gamma \mathbb{P}(u \cdot \nabla) u - \gamma(1+t)^{\gamma-1} u$, that is

$$\begin{aligned} & \|(1+t)^\gamma \mathbb{P}(u \cdot \nabla) u + \gamma(1+t)^{\gamma-1} u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \\ & \lesssim_0 \left[\varphi(N_q(f_0) + H_{0,k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)} \right] \\ & \quad \times \left[1 + \|(1+t)^\gamma \mathbb{D}^2 u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))} \right], \end{aligned}$$

for $p \in (3, 3 + \varepsilon)$, where we have also used the continuity of the Leray projection of $L^p(\mathbb{R}_+^3)$.

In order to treat the term coming from the Brinkman force, we use the estimate (6.14) of Corollary 6.8 (which is valid in view of (6.16)). We get

$$\begin{aligned} \|(1+t)^\gamma \mathbb{P}(j_f - \rho_f u)\|_{L^p(0,T;L^p(\mathbb{R}_+^3))}^p & \lesssim_0 \int_0^T (1+s)^{p\gamma} \|j_f(s) - \rho_f u(s)\|_{L^p(\mathbb{R}_+^3)}^p ds \\ & \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)) \int_0^T \frac{1}{(1+s)^{k(p-1)-p\gamma}} ds \\ & \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)), \end{aligned} \quad (6.22)$$

for k , and thus q large enough with respect to $\gamma \in (0, \frac{17}{8} - \frac{7}{4p})$.

Gathering all the pieces together, we end up with

$$\|S(u, f)\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \lesssim \left[\varphi(N_q(f_0) + H_{0, k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)} \right] \\ \times \left[1 + \|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \right],$$

and this concludes the proof. \square

Corollary 6.14. *There exists $\varepsilon > 0$ such that for all $p \in (3, 3 + \varepsilon)$, the following holds. For all $\gamma \in (0, \frac{17}{8} - \frac{7}{4p})$, we have*

$$\|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \lesssim 0 \ 1.$$

Proof. We apply the maximal $L^p L^p$ regularity for the Stokes system (see Section A.4 in the Appendix) satisfied by $U := (1+t)^\gamma u$:

$$\begin{cases} \partial_t U + A_p U = S(u, f), \\ U|_{t=0} = u_0, \end{cases}$$

where A_p stands for the Stokes operator in $L_{\text{div}}^p(\mathbb{R}_+^3)$. This entails

$$\|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \lesssim \|u_0\|_{D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)} + \|S(u, f)\|_{L^p(0, T; L^p(\mathbb{R}_+^3))},$$

where \lesssim involves a universal constant and where $D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)$ has been defined in (A.7). Recall the meaning of the notation \lesssim_0 from Definition 4.8. Using Lemma 6.13, we then have

$$\|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \leq \|u_0\|_{D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)} \\ + \psi \left(1 + \|u_0\|_{H^1 \cap D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)} + E(0) + \|u_0\|_{L^1(\mathbb{R}_+^3)} + N_q(f_0) + H_{0, q}(f_0) \right. \\ \left. + \max_{r \in \{1, 3\}} \{K_{m, r}(f_0) + F_{q, m, r}(f_0)\} + M_6 f_0 \right) \\ \times \left[\varphi(N_q(f_0) + H_{0, k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)} \right] \\ \times \left[1 + \|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))} \right].$$

Using the smallness assumption (2.2), we can ensure that

$$\psi \left(1 + \|u_0\|_{H^1 \cap D_p^{1-\frac{1}{p}, p}(\mathbb{R}_+^3)} + E(0) + \|u_0\|_{L^1(\mathbb{R}_+^3)} + N_q(f_0) + H_{0, q}(f_0) + \max_{r \in \{1, 3\}} \{K_{m, r}(f_0) + F_{q, m, r}(f_0)\} + M_6 f_0 \right) \\ \times \left[\varphi(N_q(f_0) + H_{0, k}(f_0)) + \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right) + \|u_0\|_{H^1(\mathbb{R}_+^3)} \right] < \frac{1}{2},$$

and this allows to absorb the term $\|(1+t)^\gamma D^2 u\|_{L^p(0, T; L^p(\mathbb{R}_+^3))}$ in the l.h.s. This concludes the proof. \square

6.4 End of the proof of Theorem 2.1

We are now able to close the bootstrap argument. Recall that $T \in (T_0, t^*)$, where t^* is given in Definition 6.4.

Corollary 6.15. *The following inequalities hold*

$$\int_{T_0}^T \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} \, ds \lesssim_0 \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right), \\ \int_{T_0}^T \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} \, ds \lesssim_0 \Psi \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right).$$

Proof. We use the Gagliardo-Nirenberg-Sobolev inequality twice (see Section A.3 in the Appendix) to write that for $p > 3$ (to be determined later)

$$\|u\|_{L^\infty(\mathbb{R}_+^3)} \lesssim \|D^2 u\|_{L^p(\mathbb{R}_+^3)}^{\alpha_p} \|u\|_{L^2(\mathbb{R}_+^3)}^{1-\alpha_p}, \quad \alpha_p := \frac{3p}{7p-6},$$

$$\|\nabla u\|_{L^\infty(\mathbb{R}_+^3)} \lesssim \|D^2 u\|_{L^p(\mathbb{R}_+^3)}^{\beta_p} \|u\|_{L^2(\mathbb{R}_+^3)}^{1-\beta_p}, \quad \beta_p := \frac{5p}{7p-6},$$

where \lesssim hides a universal constant. Thanks to the Hölder inequality in time, we then get for all $\gamma > 0$

$$\begin{aligned} \int_{T_0}^T \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds &\lesssim \int_0^T \|D^2 u(s)\|_{L^p(\mathbb{R}_+^3)}^{\alpha_p} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^{1-\alpha_p} ds \\ &\leq \|(1+t)^\gamma D^2 u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))}^{\alpha_p} \left(\int_0^T (1+t)^{-\gamma \frac{p\alpha_p}{p-\alpha_p}} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^{\frac{(1-\alpha_p)p}{p-\alpha_p}} ds \right)^{\frac{p-\alpha_p}{p}}, \end{aligned} \quad (6.23)$$

as well as

$$\int_{T_0}^T \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds \lesssim \|(1+t)^\gamma D^2 u\|_{L^p(0,T;L^p(\mathbb{R}_+^3))}^{\beta_p} \left(\int_0^T (1+t)^{-\gamma \frac{p\beta_p}{p-\beta_p}} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^{\frac{(1-\beta_p)p}{p-\beta_p}} ds \right)^{\frac{p-\beta_p}{p}}. \quad (6.24)$$

We also note that we have

$$\frac{p-\alpha_p}{p\alpha_p} = \frac{7}{3} - \frac{3}{p}, \quad \frac{p-\beta_p}{p\beta_p} = \frac{7}{5} - \frac{11}{5p}.$$

In view of Corollary 6.14, we thus choose $p > 3$ close to 3 (remembering the assumption $u_0 \in D_p^{1-\frac{1}{p}}(\mathbb{R}_+^3)$ in (2.1)) and take γ close to $\frac{17}{8} - \frac{7}{4p}$ (which is strictly greater than $\frac{7}{3} - \frac{3}{p}$ and $\frac{7}{5} - \frac{11}{5p}$ for p close to 3) so that

$$\gamma \frac{p\alpha_p}{p-\alpha_p} > 1, \quad \gamma \frac{p\beta_p}{p-\beta_p} > 1.$$

Observe that we can first choose such p and γ and then perform the whole analysis of the two previous subsections (using a finite number of quantities defined in Definition 1.4, as well as a finite number of times the smallness assumption (2.2), whether the estimates depend on γ or not). We then take the largest exponent involved in these quantities. We can now come back to (6.23) and (6.24), and use the uniform (in time) inequality $\|u(s)\|_{L^2(\mathbb{R}_+^3)} \lesssim \Psi\left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2\right)$ on $[0, T]$ with the previous exponents (ensuring uniform in time bound), therefore we get the desired results thanks to Corollary 6.14. \square

Recall that we have assumed that $t^* < \infty$.

End of the proof of Theorem 2.1. As before, we use the notation $F := j_f - \rho_f u$. Applying Corollary 6.8 with $k > 1$, we get

$$\int_0^{t^*} \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)),$$

while combining the Cauchy-Schwarz inequality with Corollary 6.8 with $k > 3$ yields

$$\begin{aligned} \int_0^{t^*} \|F(s)\|_{L^2(\mathbb{R}_+^3)} ds &\leq \left(\int_0^{t^*} \frac{ds}{(1+s)^2} \right)^{1/2} \left(\int_0^{t^*} (1+s)^2 \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right)^{1/2} \\ &\lesssim_0 \varphi(N_q(f_0) + H_{0,k}(f_0)). \end{aligned}$$

By the smallness assumption (2.2), we can thus ensure that

$$\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \int_0^{t^*} \left[\|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \|F(s)\|_{L^2(\mathbb{R}_+^3)} \right] ds < \frac{C_\star}{2}, \quad (6.25)$$

where C_\star refers to the universal constant given in Proposition 4.20. Since $F \in L_{\text{loc}}^2(\mathbb{R}^+; L^2(\mathbb{R}_+^3)) \hookrightarrow L_{\text{loc}}^1(\mathbb{R}^+; L^2(\mathbb{R}_+^3))$ by Proposition 4.18, a continuity argument shows that there exists a strong existence time strictly larger than t^* . Furthermore, by an appropriate choice of Φ in (2.2), we can use Corollary 6.15 so that

$$\int_{T_0}^{t^*} \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta}{4}, \quad \int_{T_0}^{t^*} \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta}{2}.$$

According to Proposition 4.19 and Proposition 4.25, this means that there exists a strong existence time $t > t^*$ such that

$$\int_{T_0}^t \|u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \frac{\delta}{2}, \quad \int_{T_0}^t \|\nabla u(s)\|_{L^\infty(\mathbb{R}_+^3)} ds < \delta.$$

This is a contradiction with the very definition of t^* therefore we necessarily have $t^* = +\infty$. The proof of Theorem 2.1 is now complete. \square

A Appendix

A.1 DiPerna-Lions theory in $\mathbb{R}_+^3 \times \mathbb{R}^3$

Theorem A.1. *Let $\chi \in \mathcal{C}^\infty(\mathbb{R})$ such that $|\chi(z)| \leq |z|$ and $\chi' \in L^\infty(\mathbb{R})$. Take $f_0 \in L^1 \cap L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)$, $G \in \mathbb{R}^3$ and a vector field $u \in L_{\text{loc}}^1(\mathbb{R}^+; W^{1,1}(\mathbb{R}_+^3))$. Consider the following boundary value problem on $\mathbb{R}_+^3 \times \mathbb{R}^3$.*

$$\begin{aligned} \partial_t f + v \cdot \nabla_x f + \text{div}_v(\chi(u-v)f + Gf) &= 0, \\ f|_{t=0} &= f_0, \\ f &= 0, \text{ on } \Sigma^-. \end{aligned}$$

Then we have, for all fixed $T > 0$

- Well-posedness: There exists a unique $f \in L_{\text{loc}}^\infty(\mathbb{R}^+; L^1 \cap L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3))$ which is a weak solution of the previous Cauchy problem. Furthermore,

$$f \in \mathcal{C}(\mathbb{R}^+; L_{\text{loc}}^p(\overline{\mathbb{R}_+^3 \times \mathbb{R}^3})),$$

for all $p \in [1, \infty)$ and the function f has a trace on $\partial\mathbb{R}_+^3 \times \mathbb{R}^3$ defined in the following sense: there exists a unique element $\gamma f \in L^\infty([0, T] \times \partial\mathbb{R}_+^3 \times \mathbb{R}^3)$ such that for any test function $\psi \in \mathcal{C}^\infty([0, T] \times \overline{\mathbb{R}_+^3 \times \mathbb{R}^3})$ with compact support in space and velocity, and for all $0 \leq t_1 \leq t_2 \leq T$

$$\begin{aligned} & \int_{t_1}^{t_2} \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) [\partial_t \psi + v \cdot \nabla_x \psi + (\chi(u(t, x) - v) + G) \cdot \nabla_v \psi](t, x, v) dx dv dt \\ &= \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t_2, x, v) \psi(t_2, x, v) dv dx - \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t_1, x, v) \psi(t_1, x, v) dx dv \\ & \quad + \int_{t_1}^{t_2} \int_{\partial\mathbb{R}_+^3 \times \mathbb{R}^3} [(\gamma f) \psi(t, x, v)] v \cdot n(x) d\sigma(x) dv dt. \end{aligned}$$

- Stability: If

$$u_n \longrightarrow u \text{ in } L_{\text{loc}}^1(\mathbb{R}^+; L^1(\mathbb{R}_+^3)) \text{ and } f_{0,n} \longrightarrow f_0 \text{ in } L_{\text{loc}}^1(\mathbb{R}_+^3 \times \mathbb{R}^3),$$

the corresponding sequence of solutions (f_n) satisfies for all $p < \infty$,

$$f_n \longrightarrow f \text{ in } L_{\text{loc}}^\infty(\mathbb{R}^+; L^p(\mathbb{R}_+^3 \times \mathbb{R}^3)).$$

Such a result can be found in [10, Theorem 3.2 - Proposition 3.2] for the well-posedness and renormalization properties and in [12, Theorem VI.1.9] for the stability property.

A.2 The Cauchy problem for the Vlasov-Navier-Stokes system

Since the existence of global Leray solutions to the Vlasov-Navier-Stokes system on a half-space with a gravity force has not been explicitly written in the literature, we provide some rather short elements of proof. We especially focus on the difficulties which may appear when obtaining the strong energy inequalities satisfied by these solutions (namely, the inequalities (1.28) and (1.29)).

Theorem A.2 (Existence of Leray solutions). *Consider (u_0, f_0) a pair of admissible initial conditions in the sense of Definition 1.2. Then there exists a Leray solution (u, f) in the sense of Definition 1.6 to the system (1.1)-(1.3) with boundary conditions (1.7)-(1.11) and with initial data (u_0, f_0) .*

The proof of this result follows a now classical method for those non linear coupled problems: first, we introduce an approximated problem of the whole system. For each fixed $n \in \mathbb{N}^*$, an appropriate fixed point procedure gives the existence of (f_n, u_n) which are solutions to the following regularized system on each interval $[0, T]$

$$\partial_t f_n + v \cdot \nabla_x f_n + \operatorname{div}_v (\chi_n (u_n - v) f_n + G f_n) = 0, \quad (\text{A.1})$$

$$\partial_t u_n + (J_n u_n \cdot \nabla) u_n - \Delta u_n + \nabla p_n = \int_{\mathbb{R}^3} f_n \chi_n (v - u_n) dv, \quad (\text{A.2})$$

$$\operatorname{div} u_n = 0, \quad (\text{A.3})$$

$$(f_n, u_n)|_{t=0} = (f_0^n, J_n u_0), \quad (\text{A.4})$$

where $T > 0$ is arbitrary and where we consider the following operator, regularizations and data, and their respective convergences when n goes to infinity:

- $\chi_n \xrightarrow{n \rightarrow +\infty} \operatorname{Id}_{\mathbb{R}}$, where for all $n \in \mathbb{N} \setminus \{0\}$, $\chi_n \in \mathcal{D}(\mathbb{R})$ is an odd, increasing and bounded function satisfying $|\chi(z)| \leq |z|$, $\|\chi'_n\|_{L^\infty(\mathbb{R})} \leq 1$ and $z \chi_n(z) \geq 0$. Above, χ_n is applied componentwise.
- $f_0^n \xrightarrow{n \rightarrow +\infty} f_0$ strongly in $L^p(\mathbb{R}_+^3 \times \mathbb{R}^3)$ for all $1 \leq p < \infty$ and weakly-* in $L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3)$, and where $f_0^n = \eta^n f_0$ with $(\eta^n)_n$ a family of positive functions compactly supported in velocity such that $0 \leq \eta^n \leq 1$ and increasing to 1
- for all $n \in \mathbb{N} \setminus \{0\}$, the operator J_n refers to the Yosida approximation of the identity defined by

$$J_n := (\operatorname{I} + n^{-1} A_2)^{-1}.$$

It acts on $L_{\operatorname{div}}^2(\mathbb{R}_+^3)$ and satisfies: for all $v \in L_{\operatorname{div}}^2(\mathbb{R}_+^3)$, $J_n v \in D(A_2)$, $\|J_n v\|_{L^2(\mathbb{R}_+^3)} \leq \|v\|_{L^2(\mathbb{R}_+^3)}$ and $J_n v \xrightarrow{n \rightarrow +\infty} v$ in $L^2(\mathbb{R}_+^3)$.

In the previous construction, u_n is a strong solution to the Navier-Stokes equations while f_n is a weak solution to the Vlasov equation which is compactly supported in velocity. Note that the presence of the gravity force is harmless in this procedure.

Then, we pass to the limit in the previous approximated problem by compactness arguments, in order to recover solution to the whole system of equations. To do so, we shall obtain energy estimates which are independent of n : all in all, it can be proven that $(u_n)_n$ is weakly compact in $L_{\operatorname{loc}}^\infty(\mathbb{R}^+; L^2(\mathbb{R}_+^3)) \cap L_{\operatorname{loc}}^2(\mathbb{R}^+; H^1(\mathbb{R}_+^3))$ while (f_n) is weakly-* compact in $L_{\operatorname{loc}}^\infty(\mathbb{R}^+; L^\infty(\mathbb{R}_+^3 \times \mathbb{R}^3))$. Furthermore, a standard application of the Aubin-Lions lemma ensures that $(u_n)_n$ is strongly compact in $L_{\operatorname{loc}}^2(\mathbb{R}^+; L_{\operatorname{loc}}^2(\mathbb{R}_+^3))$ and this also implies that, up to an additional extraction, $(u_n)_n$ converges almost everywhere. Note that the treatment of the gravity term G requires an application of the Grönwall lemma for the quantity $M_2 f_n$, which turns out to be bounded itself in $L_{\operatorname{loc}}^\infty(\mathbb{R}^+)$. These convergences are enough to pass to the limit in the weak formulation of the Vlasov equation and of the Navier-Stokes equations because each term converges in the sense of distributions (see e.g. [11] for a proof, in a more involved context). An additional diagonal extraction along increasing intervals of the type $[0, T_j]$ with $T_j \rightarrow +\infty$ allows to obtain global solutions.

We then explain how one can recover the strong energy inequality (1.28) for the Navier-Stokes equations, as well as the strong energy inequality (1.29) for the Vlasov-Navier-Stokes system: since the case $s = 0$ is

classical (at least for the Navier-Stokes equations with a given source term, see e.g [12]), we restrict ourselves to the case $s > 0$ and we pay attention to the particular treatment of the coupling terms.

We first write the strong energy inequality (1.28) for the Navier-Stokes equations. It seems that the construction of solutions satisfying this strong energy inequality (and not only the same inequality with $s = 0$ and $t \geq 0$) has long remained an open problem for general unbounded domains, due to the interaction of the pressure with the distant boundaries [41] (see however a general modern treatment in [25]). Nevertheless, the strategy applied in [52] for exterior domains actually extends for the half-space case. We thus refer to [52] where the proof of (1.28) for the Navier-Stokes equations with a forcing term is written down explicitly: namely, it consists in obtaining a localised energy inequality between almost any times $s < t$, by multiplying the regularized version of the Navier-Stokes equations by $2u_n\varphi_k$ (where $\varphi_k := \varphi(\cdot/k)$ with $\varphi \in \mathcal{D}(\mathbb{R}^3)$ satisfying $0 \leq \varphi \leq 1$, $\varphi(x) = 1$ for $|x| \leq 1$ and $\varphi(x) = 0$ for $|x| \geq 2$), integrating on $(s, t) \times \mathbb{R}_+^3$, performing integration by parts, and then understanding the behavior of the remaining terms, especially for the pressure ones. All in all, this reads as

$$\int_{\mathbb{R}_+^3} |u_n(t)|^2 \phi_k \, dx + 2 \int_s^t \int_{\mathbb{R}_+^3} |\nabla u_n|^2 \phi_k \, dx \, d\tau \leq R_{s,t}^{k,n} + \int_{\mathbb{R}_+^3} |u_n(s)|^2 \phi_k \, dx + 2 \int_s^t \int_{\mathbb{R}_+^3} \varphi_k u_n \cdot S_{\chi_n}(u_n, f_n) \, dx \, d\tau, \quad (\text{A.5})$$

with

$$S_{\chi_n}(u_n, f_n) := \int_{\mathbb{R}^3} f_n \chi_n(v - u_n) \, dv,$$

and where $R_{s,t}^{k,n}$ is a remaining term such that $\lim_{k \rightarrow +\infty} \limsup_{n \rightarrow +\infty} R_{s,t}^{k,n} = 0$ (see [52] for a proof).

Thus, the main difficulty essentially lies in the treatment of the last term in the r.h.s of the inequality (A.5). First, since $(u_n)_n$ is bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+; L^2(\mathbb{R}_+^3)) \cap L_{\text{loc}}^2(\mathbb{R}^+; L^6(\mathbb{R}_+^3))$, an interpolation argument shows that $(u_n)_n$ is bounded in $L_{\text{loc}}^r(\mathbb{R}^+; L^s(\mathbb{R}_+^3))$ if

$$\frac{2}{r} + \frac{3}{s} = \frac{3}{2}, \quad 2 \leq s \leq 6.$$

For every n , we have in particular $u_n \in L_{\text{loc}}^{20/9}(\mathbb{R}^+; L^5(\mathbb{R}_+^3)) \hookrightarrow L_{\text{loc}}^{19/9}(\mathbb{R}^+; L^5(\mathbb{R}_+^3))$. Furthermore, there exists $r_1 > 19/9$ and $s_1 > 5$ such that $(u_n)_n$ is bounded in $L_{\text{loc}}^{r_1}(\mathbb{R}^+; L^{s_1}(\mathbb{R}_+^3))$. Since (u_n) converges almost everywhere towards u , a corollary of Vitali convergence theorem shows that $u_n \xrightarrow[n \rightarrow +\infty]{} u$ in $L_{\text{loc}}^{19/9}(\mathbb{R}^+; L_{\text{loc}}^5(\mathbb{R}_+^3))$.

We also know that $(M_2 f_n)_n$ is bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+)$ so that Lemma 4.15 and the maximum principle ensure that $(j_{f_n})_n$ is bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+; L^{5/4}(\mathbb{R}_+^3))$ while $(\rho_{f_n})_n$ is bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+; L^{5/3}(\mathbb{R}_+^3)) \cap L_{\text{loc}}^\infty(\mathbb{R}^+; L^1(\mathbb{R}_+^3))$, and thus bounded in $L_{\text{loc}}^\infty(\mathbb{R}^+; L^{p_\theta}(\mathbb{R}_+^3))$ where $p_\theta = \frac{5}{3+2\theta}$ with $\theta \in [0, 1]$. As there exist $s_2 < 6$ close enough to 6 and $\theta \in [0, 1]$ such that $p_\theta^{-1} + s_2^{-1} = 4/5$ and $(u_n)_n$ is bounded in $L_{\text{loc}}^{r_2}(\mathbb{R}^+; L^{s_2}(\mathbb{R}_+^3))$ for some $r_2 > 2 > 19/10 = (19/9)'$, Hölder inequality shows that $(\rho_{f_n} u_n)_n$ is bounded in $L_{\text{loc}}^{19/10}(\mathbb{R}^+; L^{5/4}(\mathbb{R}_+^3))$, as well as $(j_{f_n})_n$. By the properties of χ_n , we observe that for all n

$$|S_{\chi_n}(u_n, f_n)| \leq j_{f_n} + |u_n| \rho_{f_n},$$

therefore the previous bound entails that $(S_{\chi_n}(u_n, f_n))_n$ converges weakly in $L_{\text{loc}}^{19/10}(\mathbb{R}^+; L^{5/4}(\mathbb{R}_+^3))$ to some limit S , up to extraction. Since $(S_{\chi_n}(u_n, f_n))_n$ also converges to $j_f - \rho_f u$ in the sense of distributions, we can identify $S = j_f - \rho_f u$.

All in all, combining the strong convergence of $(u_n)_n$ and the weak convergence of $(S_{\chi_n}(u_n, f_n))_n$ we have just obtained, we get for all fixed k

$$\int_s^t \int_{\mathbb{R}_+^3} \varphi_k u_n \cdot S_{\chi_n}(u_n, f_n) \, dx \, d\tau \xrightarrow[n \rightarrow +\infty]{} \int_s^t \int_{\mathbb{R}_+^3} \varphi_k u \cdot (j_f - \rho_f u) \, dx \, d\tau,$$

and by the dominated convergence theorem, we can then pass to the limit when $k \rightarrow +\infty$ in the inequality (A.5).

In this procedure, one has to first pick one $s > 0$ such that $u_n(s) \xrightarrow{n \rightarrow +\infty} u(s)$ in $L^2(K)$ for any compact $K \subset \mathbb{R}^3$ and the inequality (1.28) eventually holds for almost any $s > 0$ (including $s = 0$) and almost any $t \geq s$. Standard lower-semicontinuity arguments allow one to extend this inequality to any $t \geq s$.

Finally, we come back to the proof of the energy inequality (1.29) for the Vlasov-Navier-Stokes system: in view of (1.28), we only need to prove that for all $0 \leq s \leq t$

$$M_2 f(t) + 2 \int_s^t M_2 f(\tau) d\tau \leq M_2 f(s) + 2 \int_s^t \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(\tau, x, v) v \cdot [u(t, x) + G] dx dv d\tau. \quad (\text{A.6})$$

We proceed as follows: mimicking the beginning of the proof of Lemma 4.16 for $\alpha = 2$ (which does not use the energy inequality (1.29) but only the regularity of u), we obtain the fact that $(t, x, v) \mapsto |v|^2 f \in L_{\text{loc}}^\infty(\mathbb{R}^+; L^1(\mathbb{R}_+^3 \times \mathbb{R}^3))$. We now consider $\psi_n(v) := \theta_n(|v|^2)$ where $\theta_n(z) := n\theta(z/n)$ is defined thanks to a positive function $\theta \in \mathcal{D}(\mathbb{R}^+)$ equal to the identity function on $[0, 1]$, less than this function on $(1, +\infty)$ and with a derivative uniformly bounded by 1. The function ψ_n is an admissible test function in the weak formulation of the Vlasov equation with a trace (see Section A.1) and this implies

$$\begin{aligned} & \int_s^t \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(\tau, x, v) 2\theta' \left(\frac{|v|^2}{n} \right) v \cdot [(u(t, x) - v + G)] dx dv d\tau \\ &= \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) \theta_n(|v|^2) dv dx - \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(s, x, v) \theta_n(|v|^2) dx dv \\ & \quad + \int_s^t \int_{\Sigma^+} [(\gamma f) \theta_n(|v|^2)] v \cdot n(x) d\sigma(x) dv d\tau \\ & \geq \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(t, x, v) \theta_n(|v|^2) dv dx - \int_{\mathbb{R}_+^3 \times \mathbb{R}^3} f(s, x, v) \theta_n(|v|^2) dx dv, \end{aligned}$$

because of the boundary condition (1.11). Since $\theta_n(|v|^2) \xrightarrow{n \rightarrow +\infty} |v|^2$ with $\theta_n(|v|^2) \leq |v|^2$, and $\theta'(|v|^2/n) \xrightarrow{n \rightarrow +\infty} 1$ with $\theta'(|v|^2/n) \leq 1$, we can use the dominated convergence theorem to pass to the limit when $n \rightarrow +\infty$ in the previous inequality: indeed, $M_2 f \in L_{\text{loc}}^1(\mathbb{R}^+)$ while $u \in L_{\text{loc}}^2(\mathbb{R}^+; L^5(\mathbb{R}_+^3))$ and $m_1 f \in L_{\text{loc}}^2(\mathbb{R}^+; L^{5/4}(\mathbb{R}_+^3))$ by interpolation. Adding (1.28) to (A.6), we eventually get the energy inequality (1.29).

A.3 Gagliardo-Nirenberg-Sobolev inequality on \mathbb{R}_+^3

Theorem A.3. *Let $1 \leq p, q, r \leq \infty$ and $m \in \mathbb{N}$. Suppose $j \in \mathbb{N}$ and $\alpha \in [0, 1]$ satisfy the relations*

$$\begin{aligned} \frac{1}{p} &= \frac{j}{3} + \left(\frac{1}{r} - \frac{m}{3} \right) \alpha + \frac{1-\alpha}{q}, \\ \frac{j}{m} &\leq \alpha \leq 1, \end{aligned}$$

with the exception $\alpha < 1$ if $m - j - d/r \in \mathbb{N}$.

Then for all $g \in L^q(\mathbb{R}_+^3)$, if $D^m g \in L^r(\mathbb{R}_+^3)$, we have $D^j g \in L^p(\mathbb{R}_+^3)$ with the estimate

$$\|D^j g\|_{L^p(\mathbb{R}_+^3)} \lesssim \|D^m g\|_{L^r(\mathbb{R}_+^3)}^\alpha \|g\|_{L^q(\mathbb{R}_+^3)}^{1-\alpha},$$

where \lesssim refers to a constant only depending on the dimension d .

This result can be found in [16, Thm 1.5.2], at least for the case of the whole space \mathbb{R}^3 . In the case of the half-space \mathbb{R}_+^3 , we can rely on the existence of an extension operator mapping functions defined on \mathbb{R}_+^3 to function defined on \mathbb{R}^3 and which is continuous for the topology of Sobolev spaces. More precisely, for all $m \in \mathbb{N} \setminus \{0\}$ and $\ell \in [1, \infty)$, there exists an operator $E : W^{m, \ell}(\mathbb{R}_+^3) \rightarrow W^{m, \ell}(\mathbb{R}^3)$ (which is independent of ℓ) such that for all $v \in W^{m, \ell}(\mathbb{R}_+^3)$, we have

- $(Ev)|_{\mathbb{R}_+^3} = v$,

- for all $i \in \{0, \dots, m\}$, $\|D^i(Ev)\|_{L^\ell(\mathbb{R}^3)} \lesssim \|D^i v\|_{L^\ell(\mathbb{R}_+^3)}$.

This result can be found in (the proofs of) [21, Section 2.3.3] and be combined with the Gagliardo-Nirenberg-Sobolev inequality on \mathbb{R}^3 to deduce Theorem A.3.

A.4 Maximal $L^p L^q$ regularity for the Stokes system on \mathbb{R}_+^3

Let $1 < q < \infty$ be fixed. Given a vector field in $u \in L^q(\mathbb{R}_+^3)$, this can be uniquely decomposed as

$$\begin{aligned} u &= \tilde{u} + \nabla p, \\ \tilde{u} &\in L_{\text{div}}^q(\mathbb{R}_+^3), \quad p \in L^q(\mathbb{R}_+^3), \quad \nabla p \in L^q(\mathbb{R}_+^3). \end{aligned}$$

where $L_{\text{div}}^q(\mathbb{R}_+^3)$ stands for the closure in $L^q(\mathbb{R}_+^3)$ of $\mathcal{D}_{\text{div}}(\mathbb{R}_+^3)$. We recall that the projection $\mathbb{P}_q : u \mapsto \tilde{u}$ is continuous from $L^q(\mathbb{R}_+^3)$ to $L_{\text{div}}^q(\mathbb{R}_+^3)$.

For $1 < q < \infty$, we consider the following Stokes operator

$$A_q := -\mathbb{P}_q \Delta, \quad D(A_q) := L_{\text{div}}^q(\mathbb{R}_+^3) \cap W_0^{1,q}(\mathbb{R}_+^3) \cap W^{2,q}(\mathbb{R}_+^3).$$

We also set

$$D_q^{1-\frac{1}{s},s}(\mathbb{R}_+^3) := (D(A_q), L_{\text{div}}^q(\mathbb{R}_+^3))_{1/s,s}, \quad (\text{A.7})$$

where $(\cdot, \cdot)_{1/s,s}$ refers to the real interpolation space of exponents $(1/s, s)$. In the case of the Stokes operator A_q , which generates an analytic semigroup e^{-tA_q} , the quantity

$$\|u\|_{L^q(\mathbb{R}_+^3)} + \left(\int_0^\infty \|A_q e^{-tA_q} u\|_{L^q(\mathbb{R}_+^3)}^s dt \right)^{1/s} \quad (\text{A.8})$$

defines an equivalent norm on $D_q^{1-\frac{1}{s},s}(\mathbb{R}_+^3)$ (see [48, Chapter 5]).

The main result is the following and can be found with further references in [27].

Theorem A.4. *Consider $0 < T \leq \infty$ and $1 < q, s < \infty$. Then, for every $u_0 \in D_q^{1-\frac{1}{s},s}(\mathbb{R}_+^3)$ which is divergence free and $f \in L^s(0, T; L_{\text{div}}^q(\mathbb{R}_+^3))$, there exists a unique solution u of the Stokes system*

$$\begin{aligned} \partial_t u + A_q u &= f, \\ u(0, x) &= u_0(x), \end{aligned}$$

satisfying

$$u \in L^s(0, T'; D(A_q)) \text{ for all finite } T' \leq T,$$

and

$$\|\partial_t u\|_{L^s(0, T; L^q(\mathbb{R}_+^3))} + \|D^2 u\|_{L^s(0, T; L^q(\mathbb{R}_+^3))} \leq C \left(\|u_0\|_{D_q^{1-\frac{1}{s},s}(\mathbb{R}_+^3)} + \|f\|_{L^s(0, T; L^q(\mathbb{R}_+^3))} \right),$$

where $C = C(q, s) > 0$.

Furthermore, if $u_0 \in W_0^{1,q}(\mathbb{R}_+^3) \cap L_{\text{div}}^q(\mathbb{R}_+^3)$ and if $s \in (1, 2]$, the statement holds and we can replace $\|u_0\|_{D_q^{1-\frac{1}{s},s}(\mathbb{R}_+^3)}$ by $\|u_0\|_{W_0^{1,q}(\mathbb{R}_+^3)}$ in the right hand side of the previous inequality.

For the sake of completeness, we bring some precisions concerning the last statement of the theorem (even if a related fact can be found in [27, Remark 2.5]). Suppose that $u \in D(A_q^{\frac{1}{2}}) \cap L_{\text{div}}^q(\mathbb{R}_+^3)$. If $s \in (1, 2)$, we write

$\frac{1}{2} = 1 - \frac{1}{s} + \varepsilon$ with $\varepsilon = \frac{2-s}{2s} > 0$ (so that $1 - s\varepsilon < 1$). Since $A_q = A$ generates an analytic semigroup, we have by standard functional calculus manipulations (see e.g. [32, Chapter 3])

$$\begin{aligned} \int_0^\infty \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^s dt &= \int_0^1 \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^s dt + \int_1^\infty \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^s dt \\ &= \int_0^1 \|A^{\frac{1}{s}-\varepsilon}e^{-tA}A^{\frac{1}{2}}u\|_{L^q(\mathbb{R}_+^3)}^s dt + \int_1^\infty \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^s dt \\ &\lesssim \|A^{\frac{1}{2}}u\|_{L^q(\mathbb{R}_+^3)}^s \int_0^1 \frac{1}{t^{1-s\varepsilon}} dt + \|u\|_{L^q(\mathbb{R}_+^3)}^s \int_1^\infty \frac{1}{t^s} dt < \infty, \end{aligned}$$

where we have used the fact that the function $x \mapsto x^{\frac{1}{s}-\varepsilon}e^{-tx}$ and $x \mapsto xe^{-tx}$ are respectively bounded on \mathbb{R}^+ by (a constant times) $t^{\varepsilon-\frac{1}{s}}$ and t^{-1} .

In the case where $s = 2$, we proceed in a similar way and write

$$\begin{aligned} \int_0^\infty \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^2 dt &\leq \int_0^1 \|A^{\frac{1}{2}}e^{-tA}A^{\frac{1}{2}}u\|_{L^q(\mathbb{R}_+^3)}^2 dt + \int_1^\infty \|Ae^{-tA}u\|_{L^q(\mathbb{R}_+^3)}^2 dt \\ &\lesssim \|A^{\frac{1}{2}}u\|_{L^q(\mathbb{R}_+^3)}^2 \int_0^1 \frac{e^{-\frac{1}{t}}}{t} dt + \|u\|_{L^q(\mathbb{R}_+^3)}^2 \int_1^\infty \frac{1}{t^2} dt < \infty, \end{aligned}$$

where we have used the fact that the function $x \mapsto \sqrt{x}e^{-tx}$ is bounded by $t^{-1/2}e^{-\frac{1}{2t}}$.

Thus, the conclusion follows because $D(A_q^{\frac{1}{2}}) = W_0^{1,q}(\mathbb{R}_+^3) \cap L_{\text{div}}^q(\mathbb{R}_+^3)$ (see [6]).

A.5 Conditional decay of the energy: proof of Theorem 3.1

Let u be a Leray solution (with strong energy inequality) to the Vlasov-Navier-Stokes system with source term F and initial data u_0 . We provide some elements of proof concerning the conditional decay of the L^2 norm of u on $[0, T]$, provided that the forcing term F satisfies

$$\forall t \in [0, T], \quad \|F(t)\|_{L^2(\mathbb{R}_+^3)} \leq \frac{C}{(1+t)^{7/4}}, \quad (\text{A.9})$$

for some constant C . As explained in Section 3, this enters within the scope of the situation treated in [58], with the adaptation to the half-space case coming from [6].

We first recall the following L^rL^q estimates for the semigroup generated by the Stokes operator A_q on $L_{\text{div}}^q(\mathbb{R}_+^3)$ (see Section A.4 in the Appendix). A proof can be found in [6].

Lemma A.5. *Let $a \in L_{\text{div}}^2(\mathbb{R}_+^3) \cap L^r(\mathbb{R}_+^3)$ for some $r \in [1, \infty]$. Then for all $t \geq 0$*

$$\|e^{-tA_q}a\|_{L^q(\mathbb{R}_+^3)} \leq \frac{C}{t^{\frac{3}{2}(\frac{1}{r}-\frac{1}{q})}} \|a\|_{L^r(\mathbb{R}_+^3)}, \quad (\text{A.10})$$

where C is independent of a and t , provided either $1 < r \leq q < \infty$, or $1 \leq r < q \leq \infty$.

Proof of Theorem 3.1. Consider a given time-dependent smooth positive cut-off function $g : \mathbb{R}^+ \rightarrow (0, +\infty)$. Since the Stokes operator $A = A_2$ on $L_{\text{div}}^2(\mathbb{R}_+^3)$ is a self-adjoint and positive operator (see Section A.4 in the Appendix), we can use a spectral decomposition of this operator thanks to a resolution of the identity $(E_\lambda)_{\lambda \geq 0}$ (see [57, Section 3.2]), namely

$$A = \int_0^{+\infty} \lambda dE_\lambda.$$

By standard functional calculus rules and the Plancherel-Parseval theorem, we can write for all $\tau \in [0, T]$

$$\begin{aligned} \int_{\mathbb{R}_+^3} |\nabla u(\tau)|^2 dx &= \|A^{1/2}u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 = \int_0^{+\infty} \lambda \|dE_\lambda u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \\ &\geq \int_{\sqrt{\lambda} > g(\tau)} \lambda \|dE_\lambda u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \\ &\geq g^2(\tau) \|u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 - g^2(\tau) \int_0^{g^2(\tau)} \|dE_\lambda u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2. \end{aligned}$$

Then, we can combine the strong energy inequality (1.28) for the Navier-Stokes equations with the Cauchy-Schwarz inequality to get the following key inequality: for all $t \in [0, T]$ and almost every $s \in [0, t]$ (including $s = 0$)

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \int_s^t g^2(\tau) \|u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 d\tau &\leq \|u(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \int_s^t g^{-2}(\tau) \|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 d\tau \\ &\quad + 2 \int_s^t g^2(\tau) \int_0^{g^2(\tau)} \|dE_\lambda u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 d\tau. \end{aligned} \quad (\text{A.11})$$

Before going further, we introduce the notation $E_{\leq g^2(\tau)} := \mathbf{1}_{[0, g^2(\tau)]}(A)$ (in the sense of the bounded functional calculus). In order to estimate the last term in the r.h.s of (A.11), we fix $\tau \in [s, t]$ and follow [7]: we choose $\psi \in \mathcal{D}_{\text{div}}(\mathbb{R}_+^3)$ and then take $\varphi(\sigma) := e^{-(\tau-\sigma)A} E_{\leq g^2(\tau)} \psi$ as a test function in the weak formulation (1.27) of the Navier-Stokes equations between time $\underline{\tau}$ and τ , for some $\underline{\tau} \leq \tau$ (see e.g. [49] for a discussion about this available procedure). This yields

$$\int_{\underline{\tau}}^\tau \langle (u \cdot \nabla)u(\sigma), \varphi(\sigma) \rangle d\sigma = \int_{\underline{\tau}}^\tau \langle F(\sigma), \varphi(\sigma) \rangle d\sigma - \langle u(\tau), E_{\leq g^2(\tau)} \psi \rangle + \langle u(\underline{\tau}), e^{-(\tau-\underline{\tau})A} E_{\leq g^2(\tau)} \psi \rangle. \quad (\text{A.12})$$

We now take $\underline{\tau} = 0$ and by observing that $(u \cdot \nabla)u = \text{div}(u \otimes u)$, we have

$$\int_0^\tau \langle (u \cdot \nabla)u(\sigma), e^{-(\tau-\sigma)A} E_{\leq g^2(\tau)} \psi \rangle d\sigma = - \int_0^\tau \langle u(\sigma), u(\sigma) \cdot \nabla E_{\leq g^2(\tau)} e^{-(\tau-\sigma)A} \psi \rangle d\sigma,$$

therefore (A.12) yields

$$\begin{aligned} \left| \langle u(\tau), E_{\leq g^2(\tau)} \psi \rangle \right| &\leq \left| \langle u(\underline{\tau}), e^{-\tau A} E_{\leq g^2(\tau)} \psi \rangle \right| + \int_0^\tau \left| \langle F(\sigma), E_{\leq g^2(\tau)} e^{-(\tau-\sigma)A} \psi \rangle \right| d\sigma \\ &\quad + \int_0^\tau \left| \langle u(\sigma), u(\sigma) \cdot \nabla E_{\leq g^2(\tau)} e^{-(\tau-\sigma)A} \psi \rangle \right| d\sigma \\ &:= (\text{I}) + (\text{II}) + (\text{III}). \end{aligned}$$

For (I) and (II), we use Lemma A.5 and get

$$\begin{aligned} (\text{I}) &\lesssim \| \mathbf{1}_{[0, g^2(\tau)]}(A) e^{-\tau A} u_0 \|_{L^2(\mathbb{R}_+^3)} \| \psi \|_{L^2(\mathbb{R}_+^3)} \leq \| e^{-\tau A} u_0 \|_{L^2(\mathbb{R}_+^3)} \| \psi \|_{L^2(\mathbb{R}_+^3)}, \\ (\text{II}) &\leq \int_0^\tau \| F(\sigma) \|_{L^2(\mathbb{R}_+^3)} \| E_{\leq g^2(\tau)} e^{-(\tau-\sigma)A} \psi \|_{L^2(\mathbb{R}_+^3)} d\sigma \lesssim \| \psi \|_{L^2(\mathbb{R}_+^3)} \int_0^\tau \| F(\sigma) \|_{L^2(\mathbb{R}_+^3)} d\sigma, \end{aligned}$$

where \lesssim refers to a universal constant. For (III), we rely on the following lemma, which can be found in the proof of [6, Lemma 4.3].

Lemma A.6. *For any $w \in H_{\text{div}}^1(\mathbb{R}_+^3)$, $z \in L_{\text{div}}^2(\mathbb{R}_+^3)$ and $\gamma : \mathbb{R}^+ \rightarrow \mathbb{R}^+$, we have for all $\tau \geq 0$*

$$\left\langle w, w \cdot \nabla E_{\leq \gamma(\tau)} z \right\rangle \leq C \gamma^{\frac{5}{4}}(\tau) \|w\|_{L^2(\mathbb{R}_+^3)}^2 \|z\|_{L^2(\mathbb{R}_+^3)},$$

where $\langle \cdot, \cdot \rangle$ stands for the inner product on $L^2(\mathbb{R}_+^3)$ and where C is a universal constant.

We apply this Lemma with $w = u(\sigma)$, $z = e^{-(\tau-\sigma)A}\psi$ and $\gamma = g^2$ to obtain

$$(III) \lesssim \|\psi\|_{L^2(\mathbb{R}_+^3)} g^{5/2}(\tau) \int_0^t \|u(\sigma)\|_{L^2(\mathbb{R}_+^3)}^2 d\sigma.$$

Since all the previous estimates are valid for any $\psi \in \mathcal{D}_{\text{div}}(\mathbb{R}_+^3)$, we infer that

$$\int_0^{g^2(\tau)} \|dE_\lambda u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \leq C \|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 + C \left(\int_0^\tau \|F(r)\|_{L^2(\mathbb{R}_+^3)} dr \right)^2 + C g(\tau)^5 \left(\int_0^\tau \|u(r)\|_{L^2(\mathbb{R}_+^3)}^2 dr \right)^2,$$

where C stands for a universal constant. In view of (A.11), we end up with the following differential inequality: for almost all $s \geq 0$ and for all $t \geq s$

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \int_s^t g^2(\tau) \|u(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 d\tau &\leq \|u(s)\|_{L^2(\mathbb{R}_+^3)}^2 + 2 \int_s^t C g(\tau)^7 \left(\int_0^\tau \|u(r)\|_{L^2(\mathbb{R}_+^3)}^2 dr \right)^2 d\tau \\ &\quad + 2 \int_s^t C g^2(\tau) \left[\|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 + g^{-4}(\tau) \|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \right. \\ &\quad \left. + \left(\int_0^\tau \|F(r)\|_{L^2(\mathbb{R}_+^3)} dr \right)^2 \right] d\tau. \end{aligned}$$

We then rely on the following Grönwall-like inequality (see [7]).

Lemma A.7. *Let $t > 0$. Let $y \in L^\infty(0, t)$ and $\beta \in L^1(0, t)$ be nonnegative functions. Suppose that the following differential inequality holds: for almost all $0 < s < t$*

$$y(t) + \int_s^t g(\tau) y(\tau) d\tau \leq y(s) + \int_s^t \beta(\tau) d\tau,$$

where $\tau \mapsto g(\tau)$ is a smooth positive function on $[0, t]$. Then we have for almost all $0 < s < t$

$$\exp\left(\int_0^t g(\tau) d\tau\right) y(t) \leq \exp\left(\int_0^s g(\tau) d\tau\right) y(s) + \int_s^t \exp\left(\int_0^\tau g(r) dr\right) \beta(\tau) d\tau.$$

Applying this lemma, we finally obtain

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \exp\left(\int_0^t g^2(\tau) d\tau\right) &\leq \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 \\ &\quad + C \int_0^t g^2(\tau) \left[\|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 + g^{-4}(\tau) \|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \right] \exp\left(\int_0^\tau g^2(r) dr\right) d\tau \\ &\quad + C \int_0^t g^2(\tau) \left(\int_0^\tau \|F(\sigma)\|_{L^2(\mathbb{R}_+^3)} d\sigma \right)^2 \exp\left(\int_0^\tau g^2(r) dr\right) d\tau \\ &\quad + C \int_0^t g^7(\tau) \left(\int_0^\tau \|u(\sigma)\|_{L^2(\mathbb{R}_+^3)}^2 d\sigma \right)^2 \exp\left(\int_0^\tau g^2(r) dr\right) d\tau. \end{aligned} \tag{A.13}$$

We now take

$$g^2(t) := \frac{\alpha}{1+t}, \quad \exp\left(\int_0^t g^2(\tau) d\tau\right) = (1+t)^\alpha,$$

where $\alpha > 0$ and this means that for all $t \in [0, T]$

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 (1+t)^\alpha &\lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t (1+\tau)^{\alpha-1} \left[\|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 + (1+\tau)^2 \|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 \right] d\tau \\ &\quad + \int_0^t (1+\tau)^{\alpha-1} \left(\int_0^\tau \|F(\sigma)\|_{L^2(\mathbb{R}_+^3)} d\sigma \right)^2 d\tau \\ &\quad + \int_0^t (1+\tau)^{\alpha-7/2} \left(\int_0^\tau \|u(\sigma)\|_{L^2(\mathbb{R}_+^3)}^2 d\sigma \right)^2 d\tau, \end{aligned} \tag{A.14}$$

where \lesssim is independent of t . Using Lemma A.5, we get for all $\tau > 0$

$$\|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 \leq \frac{C}{(1+\tau)^{3/2}} \left(\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2 \right),$$

where C is independent of τ and u . Observe now that, according to the assumption (A.9), the contribution of the source term $g^{-4}(\tau)\|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 + \left(\int_0^\tau \|F(r)\|_{L^2(\mathbb{R}_+^3)} dr\right)^2$ is as $(1+\tau)^{-3/2}$ so that for all $t > 0$

$$\begin{aligned} \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t (1+\tau)^{\alpha-1} \left[\|e^{-\tau A} u_0\|_{L^2(\mathbb{R}_+^3)}^2 + g^{-4}(\tau)\|F(\tau)\|_{L^2(\mathbb{R}_+^3)}^2 + \left(\int_0^\tau \|F(r)\|_{L^2(\mathbb{R}_+^3)} dr\right)^2 \right] d\tau \\ \lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t \frac{1}{(1+\tau)^{1-\alpha+3/2}} d\tau, \\ \lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + (1+t)^{\alpha-3/2}, \end{aligned} \quad (\text{A.15})$$

provided that $\alpha \neq 3/2$, and where \lesssim depends on $\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \|u_0\|_{L^1(\mathbb{R}_+^3)}^2$.

Let us assume that on $[0, T]$, we have

$$\|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim \frac{1}{(1+t)^\beta}, \quad (\text{A.16})$$

for some $\beta > 0$, different from 1. From this estimate, we deduce the following inequality

$$\int_0^t (1+\tau)^{\alpha-7/2} \left(\int_0^\tau \|u(r)\|_{L^2(\mathbb{R}_+^3)}^2 dr \right)^2 d\tau \lesssim \begin{cases} 1 & \text{if } \beta < 1, \quad \alpha - 2\beta - 1/2 < 0, \\ (1+t)^{\alpha-2\beta-1/2} & \text{if } \beta < 1, \quad \alpha - 2\beta - 1/2 \geq 0, \\ 1 & \text{if } \beta > 1. \end{cases} \quad (\text{A.17})$$

Now, we start with $\beta = 0$: the a priori estimate (A.16) indeed holds with this choice of exponent because the energy inequality (1.28) can be rewritten, together with (A.9), as

$$\begin{aligned} \|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + 2 \int_0^t \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^t (1+s)^{-7/4} ds \\ \lesssim \|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + 1. \end{aligned}$$

We then take $\alpha = 1$ and use (A.14), (A.15) and (A.17) to obtain

$$(1+t)\|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim 1 + (1+t)^{-1/2} + (1+t)^{1/2},$$

therefore

$$\|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim (1+t)^{-1/2}.$$

This means that the a priori estimate (A.16) now holds with $\beta = 1/2$. We then take $\alpha = 2$ and combining again (A.14), (A.15) and (A.17) yields

$$\|u(t)\|_{L^2(\mathbb{R}_+^3)}^2 \lesssim (1+t)^{-2} + (1+t)^{-3/2} \lesssim (1+t)^{-3/2}.$$

Note that in view of (A.15) and (A.17), we cannot improve the exponent in the previous estimate. The proof is therefore complete. \square

A.6 Parabolic regularization for the Navier-Stokes system on \mathbb{R}_+^3

Theorem A.8. *For all $T > 0$, there exists a universal constant $C_\star > 0$ such that the following holds. Consider $u_0 \in H_{\text{div}}^1(\mathbb{R}_+^3)$ and $F \in L_{\text{loc}}^2(\mathbb{R}^+; L^2(\mathbb{R}_+^3))$ and $T > 0$ such that*

$$\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \int_0^T \left[\|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 + \|F(s)\|_{L^2(\mathbb{R}_+^3)} \right] ds \leq C_\star. \quad (\text{A.18})$$

Then, there exists on $[0, T]$ a unique Leray solution to the Navier-Stokes equations with initial data u_0 and source F . This solution u belongs to $L^\infty(0, T; H^1_{\text{div}}(\mathbb{R}_+^3)) \cap L^2(0, T; H^2(\mathbb{R}_+^3))$ and there exists $\tilde{C} > 0$ such that for almost every $t \in [0, T]$

$$\|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \int_0^t \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \leq \tilde{C} \left(\|u_0\|_{H^1(\mathbb{R}_+^3)}^2 + \|F\|_{L^2(0, T; L^2(\mathbb{R}_+^3))}^2 \right). \quad (\text{A.19})$$

Proof. We argue in two different steps. First, if such a Leray solution to the Navier-Stokes equations exists, it satisfies in particular $u \in L^4(0, T; H^1_{\text{div}}(\mathbb{R}_+^3))$, which is a well-known case of weak-strong uniqueness for this system (see e.g. [15, Theorem 3.3], which holds for general domains). Thus, such a solution will be equal to any Leray solution of the Navier-Stokes system with source F and initial value u_0 .

In a second step, we have to prove that such a solution indeed exists on $[0, T]$: to do so, we rely on a standard approximation procedure by regularizing the data F and u_0 and by smoothing the convection term in the momentum equation (see e.g. [52] where the Yosida operator is used, as in Section A.2). We essentially obtain a sequence $(u_N)_N$ of solutions to regularized Navier-Stokes systems. This sequence will classically be bounded in $L^\infty L^2 \cap L^2 H^1$ and will satisfy additional parabolic estimates coming from (A.18). A compactness argument allows to get the existence of a Leray solution in $L^\infty(0, T; H^1_{\text{div}}(\mathbb{R}_+^3)) \cap L^2(0, T; H^2(\mathbb{R}_+^3))$ and which satisfies the estimate (A.19) on $[0, T]$.

Hence in the following, we deal with a smooth solution u with smooth data u_0 and we are looking for parabolic estimates for the velocity field u . The following strategy is based on an idea which is well-known in the context of the inhomogeneous incompressible Navier-Stokes equations (see [47, 19]): we first apply the Leray projection \mathbb{P} to the Navier-Stokes equations and then multiply by $\partial_t u$ to obtain, after integration by parts

$$\frac{1}{2} \frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + \|\partial_t u\|_{L^2(\mathbb{R}_+^3)}^2 = -\langle (u \cdot \nabla)u, \partial_t u \rangle_{L^2(\mathbb{R}_+^3)} + \langle F, \partial_t u \rangle_{L^2(\mathbb{R}_+^3)},$$

where we have dropped the time variable. We then use the Young inequality twice in order to absorb $\|\partial_t u\|_{L^2(\mathbb{R}_+^3)}^2$ in the l.h.s, that is on $[0, T]$

$$\frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + \|\partial_t u\|_{L^2(\mathbb{R}_+^3)}^2 \leq 2 \int_{\mathbb{R}_+^3} |(u \cdot \nabla)u|^2 dx + 2 \int_{\mathbb{R}_+^3} |F|^2 dx. \quad (\text{A.20})$$

In order to estimate $D^2 u$, we rewrite the Navier-Stokes system satisfied by (u, p) as the following stationary Stokes system in the half-space \mathbb{R}_+^3 with source term $F - (u \cdot \nabla)u - \partial_t u$

$$\begin{cases} -\Delta u + \nabla p = F - (u \cdot \nabla)u - \partial_t u, \\ \operatorname{div} u = 0, \\ u|_{\partial \mathbb{R}_+^3} = 0. \end{cases}$$

According to [26, Theorem IV.3.2], there exists a universal constant $C > 0$ such that on $[0, T]$

$$\|D^2 u\|_{L^2(\mathbb{R}_+^3)} \leq C \|F - (u \cdot \nabla)u - \partial_t u\|_{L^2(\mathbb{R}_+^3)},$$

and thus, denoting by $C > 0$ another universal constant, we have on $[0, T]$

$$\|D^2 u\|_{L^2(\mathbb{R}_+^3)}^2 \leq C \left(\|F\|_{L^2(\mathbb{R}_+^3)}^2 + \|(u \cdot \nabla)u\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \|\partial_t u\|_{L^2(\mathbb{R}_+^3)}^2 \right).$$

Coming back to (A.20), we get

$$\frac{d}{dt} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2} \|\partial_t u\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{C} \|D^2 u\|_{L^2(\mathbb{R}_+^3)}^2 \leq 3 \int_{\mathbb{R}_+^3} |(u \cdot \nabla)u|^2 dx + 3 \int_{\mathbb{R}_+^3} |F|^2 dx. \quad (\text{A.21})$$

We now provide a bound for the convective term on $[0, T]$ by using consecutively the Cauchy-Schwarz inequality, interpolation inequality for Lebesgue spaces and the Gagliardo-Nirenberg-Sobolev inequality (see Theorem A.3) to write

$$\begin{aligned} \int_{\mathbb{R}_+^3} |(u \cdot \nabla)u|^2 dx &\leq \|u\|_{L^4(\mathbb{R}_+^3)}^2 \|\nabla u\|_{L^4(\mathbb{R}_+^3)}^2 \\ &\leq \|u\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|u\|_{L^6(\mathbb{R}_+^3)}^{3/2} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|\nabla u\|_{L^6(\mathbb{R}_+^3)}^{3/2} \\ &\lesssim \|u\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^{3/2} \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^{1/2} \|D^2 u\|_{L^2(\mathbb{R}_+^3)}^{3/2}. \end{aligned} \quad (\text{A.22})$$

Thanks to Young inequality, we deduce the existence of a universal constant $\bar{C} > 0$ such that

$$3\|(u \cdot \nabla)u\|_{L^2(\mathbb{R}_+^3)}^2 \leq \bar{C}\|u\|_{L^2(\mathbb{R}_+^3)}^2 \|\nabla u\|_{L^2(\mathbb{R}_+^3)}^8 + \frac{1}{2\bar{C}}\|D^2 u\|_{L^2(\mathbb{R}_+^3)}^2. \quad (\text{A.23})$$

We integrate (A.21) between 0 and $t \in [0, T]$ and get

$$\begin{aligned} \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2\bar{C}} \int_0^t \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \\ \leq \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + 3 \int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds + \bar{C} \int_0^t \|u(s)\|_{L^2(\mathbb{R}_+^3)}^2 \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^8 ds. \end{aligned}$$

If we set

$$x(t) := \|\nabla u(t)\|_{L^2(\mathbb{R}_+^3)}^2 + \frac{1}{2\bar{C}} \int_0^t \|D^2 u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds, \quad (\text{A.24})$$

$$h(t) := \|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + 3 \int_0^t \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds, \quad (\text{A.25})$$

$$g(s) := \bar{C} \|u(s)\|_{L^2(\mathbb{R}_+^3)}^2 \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^8, \quad (\text{A.26})$$

then the previous inequality can be written as

$$x(t) \leq h(t) + \int_0^t g(s)x(s)^3 ds,$$

for all $t \in [0, T]$. We are now in position to apply Bihari's Lemma, that we recall now (see [20] for a proof)

Lemma A.9. *Let x, g and h three positive continuous functions on \mathbb{R}^+ , such that h is increasing. Let w a continuous submultiplicative and nondecreasing function on \mathbb{R}^+ such that $w(u) > 0$ for all $u > 0$. Suppose that for all $t \geq 0$*

$$x(t) \leq h(t) + \int_0^t g(s)w(x(s)) ds. \quad (\text{A.27})$$

If the function

$$W(u) := \int_1^u \frac{ds}{w(s)}, \quad u \in (0, a) \text{ for some } a > 0,$$

is a bijection on $(0, a)$, then the following inequality

$$x(t) \leq h(t)W^{-1} \left(\int_0^t g(s) \frac{w(h(s))}{h(s)} ds \right) \quad (\text{A.28})$$

holds for all $t \geq 0$ such that $\int_0^t g(s) \frac{w(h(s))}{h(s)} ds \in (0, a)$.

We use this lemma with the function $w(s) = s^3$ associated to

$$W(u) = \int_1^u \frac{ds}{s^3} = \frac{1}{2} \left(1 - \frac{1}{u^2} \right), \quad u > 0,$$

and whose inverse, defined on $(0, 1/2)$, is $W^{-1}(v) = (1 - 2v)^{-1/2}$. The inequality (A.28) reads on $[0, T]$ as

$$x(t) \leq h(t) \left(1 - 2 \int_0^t g(s) h^2(s) ds \right)^{-1/2}, \quad (\text{A.29})$$

provided that for all $t \in [0, T]$, we have $1 - 2 \int_0^t g(s) h^2(s) ds > 0$. Thus, this is enough to ensure for instance that

$$1 - 2 \int_0^T g(s) h^2(s) ds > 1/2, \quad (\text{A.30})$$

is satisfied in order to get $x(t) \leq h(t)$ for all $t \in [0, T]$. We observe that

$$\int_0^T g(s) h^2(s) ds \leq \left[\|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + 3 \int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right] \int_0^T g(s) ds,$$

and

$$\int_0^T g(s) ds \leq \bar{C} \|u\|_{L^\infty(0, T; L^2(\mathbb{R}_+^3))}^2 \int_0^T \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds.$$

Furthermore, the following energy inequality for the Navier-Stokes system with source F and initial data u_0 is satisfied by u

$$\|u\|_{L^\infty(0, T; L^2(\mathbb{R}_+^3))}^2 + 2 \int_0^T \|\nabla u(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \leq \tilde{C} \left[\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \left(\int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)} ds \right)^2 \right],$$

where $\tilde{C} > 0$ is a universal constant. The condition (A.30) can thus be satisfied provided that the quantity

$$\left[\|\nabla u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)}^2 ds \right] \left[\|u_0\|_{L^2(\mathbb{R}_+^3)}^2 + \left(\int_0^T \|F(s)\|_{L^2(\mathbb{R}_+^3)} ds \right)^2 \right]^2$$

is small enough. We observe that an assumption of the type of (A.18) can indeed provide such a smallness. Now, in view of the definitions (A.24)-(A.25), the inequality $x(t) \leq h(t)$ for all $t \in [0, T]$ brings to the conclusion (A.19) and this completes the proof. \square

Acknowledgements

I would like to acknowledge my PhD advisors Daniel Han-Kwan and Ayman Moussa for their confidence, support and guidance, as well as their constructive criticism, scientific suggestions and careful rereadings during the preparation of this work. I would also like to thank Jean-Yves Chemin and Raphaël Danchin for discussions about Theorem A.8.

References

- [1] O. Anoshchenko and A. Boutet de Monvel-Berthier. The existence of the global generalized solution of the system of equations describing suspension motion. *Mathematical methods in the applied sciences*, 20(6):495–519, 1997.

- [2] C. Bardos, F. Golse, and C. D. Levermore. Fluid dynamic limits of kinetic equations. II. Convergence proofs for the Boltzmann equation. *Communications on pure and applied mathematics*, 46(5):667–753, 1993.
- [3] C. Bardos, F. Golse, and D. Levermore. Fluid dynamic limits of kinetic equations. I. Formal derivations. *Journal of statistical physics*, 63(1):323–344, 1991.
- [4] E. Bernard, L. Desvillettes, F. Golse, and V. Ricci. A derivation of the Vlasov-Navier-Stokes model for aerosol flows from kinetic theory. *Commun. Math. Sci.*, 15(6):1703–1741, 2017.
- [5] E. Bernard, L. Desvillettes, F. Golse, and V. Ricci. A derivation of the Vlasov-Stokes system for aerosol flows from the kinetic theory of binary gas mixtures. *Kinet. Relat. Models*, 11(1):43–69, 2018.
- [6] W. Borchers and T. Miyakawa. L2 decay for the Navier-Stokes flow in halfspaces. *Mathematische Annalen*, 282(1):139–155, 1988.
- [7] W. Borchers and T. Miyakawa. Algebraic L2 decay for Navier-Stokes flows in exterior domains. *Acta Mathematica*, 165(1):189–227, 1990.
- [8] L. Boudin, L. Desvillettes, C. Grandmont, and A. Moussa. Global existence of solutions for the coupled Vlasov and Navier-Stokes equations. *Differential Integral Equations*, 22(11-12):1247–1271, 2009.
- [9] L. Boudin, C. Grandmont, A. Lorz, and A. Moussa. Modelling and numerics for respiratory aerosols. *Communications in Computational Physics*, 18(3):723–756, 2015.
- [10] L. Boudin, C. Grandmont, and A. Moussa. Global existence of solutions to the incompressible Navier-Stokes-Vlasov equations in a time-dependent domain. *J. Differ. Equations*, 262:1317–1340, 2017.
- [11] L. Boudin, D. Michel, and A. Moussa. Global existence of weak solutions to the incompressible Vlasov-Navier-Stokes system coupled to convection-diffusion equations. *Math. Models Methods Appl. Sci.*, 2020.
- [12] F. Boyer and P. Fabrie. *Mathematical Tools For The Study Of The Incompressible Navier-Stokes Equations And Related Models*, volume 183. Springer Science & Business Media, 2012.
- [13] K. Carrapatoso and M. Hillairet. On the derivation of a Stokes–Brinkman problem from Stokes equations around a random array of moving spheres. *Communications in Mathematical Physics*, 373(1):265–325, 2020.
- [14] J. A. Carrillo and T. Goudon. Stability and asymptotic analysis of a fluid-particle interaction model. *Communications in Partial Differential Equations*, 31(9):1349–1379, 2006.
- [15] J.-Y. Chemin, B. Desjardins, I. Gallagher, and E. Grenier. *Mathematical Geophysics: An Introduction To Rotating Fluids And The Navier-Stokes Equations*, volume 32. Oxford University Press on Demand, 2006.
- [16] P. Cherrier and A. Milani. *Linear And Quasi-Linear Evolution Equations In Hilbert Spaces*. American Mathematical Society Providence, 2012.
- [17] Y.-P. Choi. Finite-time blow-up phenomena of Vlasov/Navier–Stokes equations and related systems. *Journal de Mathématiques Pures et Appliquées*, 108(6):991–1021, 2017.
- [18] Y.-P. Choi and B. Kwon. Global well-posedness and large-time behavior for the inhomogeneous Vlasov–Navier–Stokes equations. *Nonlinearity*, 28(9):3309, 2015.
- [19] R. Danchin and P. B. Mucha. The incompressible Navier-Stokes equations in vacuum. *Communications on Pure and Applied Mathematics*, 72(7):1351–1385, 2019.
- [20] F. M. Dannan. Integral inequalities of Gronwall-Bellman-Bihari type and asymptotic behavior of certain second order nonlinear differential equations. *Journal of mathematical analysis and applications*, 108(1):151–164, 1985.

- [21] F. Demengel, G. Demengel, and R. Ern e. *Functional Spaces For The Theory Of Elliptic Partial Differential Equations*. Springer, 2012.
- [22] L. Desvillettes. Some aspects of the modeling at different scales of multiphase flows. *Computer methods in applied mechanics and engineering*, 199(21-22):1265–1267, 2010.
- [23] L. Desvillettes, F. Golse, and V. Ricci. The mean-field limit for solid particles in a Navier-Stokes flow. *Journal of Statistical Physics*, 131(5):941–967, 2008.
- [24] L. Ertzbischoff, D. Han-Kwan, and A. Moussa. Concentration versus absorption for the Vlasov-Navier-Stokes system on bounded domains. *arXiv preprint arXiv:2101.05157*, 2021.
- [25] R. Farwig, H. Kozono, and H. Sohr. An Lq-approach to Stokes and Navier-Stokes equations in general domains. *Acta Mathematica*, 195(1):21–53, 2005.
- [26] G. Galdi. *An Introduction To The Mathematical Theory Of The Navier-Stokes Equations: Steady-State Problems*. Springer Science & Business Media, 2011.
- [27] Y. Giga and H. Sohr. Abstract Lp estimates for the Cauchy problem with applications to the Navier-Stokes equations in exterior domains. *Journal of functional analysis*, 102(1):72–94, 1991.
- [28] O. Glass, D. Han-Kwan, and A. Moussa. The Vlasov–Navier–Stokes system in a 2d pipe: Existence and stability of regular equilibria. *Archive for Rational Mechanics and Analysis*, 230(2):593–639, 2018.
- [29] T. Goudon, L. He, A. Moussa, and P. Zhang. The Navier–Stokes–Vlasov–Fokker–Planck system near equilibrium. *SIAM journal on mathematical analysis*, 42(5):2177–2202, 2010.
- [30] T. Goudon, P.-E. Jabin, and A. Vasseur. Hydrodynamic limit for the Vlasov-Navier-Stokes equations. part i: Light particles regime. *Indiana University mathematics journal*, pages 1495–1515, 2004.
- [31] T. Goudon, P.-E. Jabin, and A. Vasseur. Hydrodynamic limit for the Vlasov-Navier-Stokes equations. part ii: Fine particles regime. *Indiana University mathematics journal*, pages 1517–1536, 2004.
- [32] M. Haase. The functional calculus for sectorial operators. In *The Functional Calculus for Sectorial Operators*, pages 19–60. Springer, 2006.
- [33] K. Hamdache. Global existence and large time behaviour of solutions for the Vlasov-Stokes equations. *Japan journal of industrial and applied mathematics*, 15(1):51, 1998.
- [34] D. Han-Kwan. Large time behavior of small data solutions to the Vlasov–Navier–Stokes system on the whole space. *arXiv preprint arXiv:2006.09848*, 2020.
- [35] D. Han-Kwan and D. Michel. On hydrodynamic limits of the Vlasov-Navier-Stokes system. *arXiv preprint arXiv:2103.06668*, 2021.
- [36] D. Han-Kwan,  . Miot, A. Moussa, and I. Moyano. Uniqueness of the solution to the 2D Vlasov-Navier-Stokes system. *to appear in Revista Matem tica Iberoamericana*, Oct. 2017.
- [37] D. Han-Kwan, A. Moussa, and I. Moyano. Large time behavior of the Vlasov-Navier-Stokes system on the torus. *Archive for Rational Mechanics and Analysis*, 236(3):1273–1323, 2020.
- [38] G. H. Hardy, J. E. Littlewood, and G. P lya. *Inequalities*. Cambridge Mathematical Library. Cambridge University Press, Cambridge, 1988. Reprint of the 1952 edition.
- [39] J. G. Heywood. The Navier-Stokes equations: on the existence, regularity and decay of solutions. *Indiana University Mathematics Journal*, 29(5):639–681, 1980.
- [40] J. G. Heywood. Epochs of regularity for weak solutions of the Navier-Stokes equations in unbounded domains. *Tohoku Mathematical Journal, Second Series*, 40(2):293–313, 1988.

- [41] J. G. Heywood. Open problems in the theory of the Navier-Stokes equations for viscous incompressible flow. In *The Navier-Stokes Equations Theory and Numerical Methods*, pages 1–22. Springer, 1990.
- [42] M. Hillairet. On the homogenization of the Stokes problem in a perforated domain. *Archive for Rational Mechanics and Analysis*, 230(3):1179–1228, 2018.
- [43] M. Hillairet, A. Moussa, and F. Sueur. On the effect of polydispersity and rotation on the Brinkman force induced by a cloud of particles on a viscous incompressible flow. *Kinet. Relat. Models*, 12, 2019.
- [44] R. M. Höfer. Sedimentation of inertialess particles in Stokes flows. *Communications in Mathematical Physics*, 360(1):55–101, 2018.
- [45] R. M. Höfer. The inertialess limit of particle sedimentation modeled by the Vlasov-Stokes equations. *SIAM Journal on Mathematical Analysis*, 50(5):5446–5476, 2018.
- [46] P. Jabin. Large time concentrations for solutions to kinetic equations with energy dissipation. *Communications in Partial Differential Equations*, 25(3-4):541–557, 2000.
- [47] O. Ladyzhenskaya and V. Solonnikov. Unique solvability of an initial-and boundary-value problem for viscous incompressible nonhomogeneous fluids. *Journal of Soviet Mathematics*, 9(5):697–749, 1978.
- [48] A. Lunardi. *Interpolation Theory*, volume 16. Springer, 2018.
- [49] K. Masuda. Weak solutions of Navier-Stokes equations. *Tohoku Mathematical Journal, Second Series*, 36(4):623–646, 1984.
- [50] J. Mathiaud. *Etude de systemes de type gaz-particules*. PhD thesis, Cachan, Ecole normale supérieure, 2006.
- [51] A. Mecherbet. Sedimentation of particles in Stokes flow. *Kinet. Relat. Models*, 12, 2019.
- [52] T. Miyakawa and H. Sohr. On energy inequality, smoothness and large time behavior in L2 for weak solutions of the Navier-Stokes equations in exterior domains. *Mathematische Zeitschrift*, 199(4):455–478, 1988.
- [53] I. Moyano. Local null-controllability of the 2-d Vlasov-Navier-Stokes system. *arXiv preprint arXiv:1607.05578*, 2016.
- [54] P. J. O’Rourke. Collective drop effects on vaporizing liquid sprays. Technical report, Los Alamos National Lab., NM (USA), 1981.
- [55] J. C. Robinson, J. L. Rodrigo, and W. Sadowski. *The Three-Dimensional Navier-Stokes Equations: Classical Theory*, volume 157. Cambridge university press, 2016.
- [56] M. E. Schonbek. Large time behaviour of solutions to the Navier-Stokes equations. *Communications in partial differential equations*, 11(7):733–763, 1986.
- [57] H. Sohr. *The Navier-Stokes Equations: An Elementary Functional Analytic Approach*. Springer Science & Business Media, 2012.
- [58] M. Wiegner. Decay results for weak solutions of the Navier-Stokes equations on \mathbb{R}^n . *Journal of the London Mathematical Society*, 2(2):303–313, 1987.
- [59] F. A. Williams. *Combustion Theory*. Benjamin Cummings, second edition, 1985.