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Determination of Electric Motor Losses and Critical Temperatures through an Inverse Approach

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ABSTRACT In this study, a practical numerical method is proposed to estimate losses of high specific power density electric motors, using few simulated temperature data. In such electric motors, these losses generate high heat fluxes inside the motor components that can be critically sensitive to temperature. Electromagnetic and mechanical friction phenomena are behind the occurring of these thermal dissipations. For both phenomena, losses could be difficult to compute with electrical or mechanical approaches. However, thermal management of electric motors requires a precise knowledge of those losses, in particular for high-performance motors such as those considered in future hybrid planes. To determine electric motor losses in a Permanent Magnet Synchronous Motor (PMSM) in real time, an inverse method using a Lumped Parameter Thermal Model (LPTM) is elaborated. In the first step, the dynamic profile of losses is determined through the inverse method, based on temperature data at easy-access points of the motor. In a second step, the identified losses are used to find temperatures at critical non-accessible hot spot points of the motor through forward LPTM. The method is applied for three useful cases, from the simplest case scenario, where only one type of losses has to be identified, to the most complicated case where all losses are simultaneously estimated. A global strategy for the choice of the number of future time steps used for regularization of the ill-posed problem is also proposed. Results show that this method enables adequate real-time supervision of the critical motor temperatures, mainly rotor and winding core.

Keywords: Electrical machines, Future time steps, Losses, Lumped Parameter Model, Thermal behavior, Regularization

I. INTRODUCTION

Nowadays, with the rising interest of electrification in the transportation sector, the study and conception of very high-specific power density electric motors is essential. Indeed, weight and volume are known to be classical constraints for electrified vehicles. Moreover, the recent use of electric motors for the propulsion of the new generation of ground and air vehicles requires the development of series of more efficient and higher performance electric motors [1]. This driving trend already exists in the automotive field for electric and hybrid-electric vehicles [2]. Recently, propulsion electrification is emerging into air vehicles design development. Numerous projects [3]-[6] are being conducted aiming to develop electrically propelled aircrafts and are targeting higher values of specific power density in electric motors. In general, these projects focus on hybrid aircrafts in which electric motors,

coupled to fans, would propel the plane, and the turbo-reactor will be mainly used as a turbine for electricity generation. The increase of specific power density of electric motors could be limited by electrical

Nomenclature							
Inomenciat	ure						
A(NN)	state matrix	Greek symbols					
$B_c(N)$	command vector relative to environment	Г	system boundary				
$B_P(N,n_n)$	command matrix relative to heat sources		time step, s				
C_n	specific heat, J.kg ⁻¹ .K ⁻¹	0	density, kg.m ⁻³				
$C(n_a,N)$	output matrix	σ_{U}	Mean quadratic error for <i>U</i> . W				
h	convective exchange coefficient, W.m ⁻² .K ⁻¹	σ_Y	Mean quadratic error for Y , °C				
Ι	identity matrix	χ	characteristic space function				
k	thermal conductivity, W.m ⁻¹ .K ⁻¹	Ω	system domain				
N	number of nodes						
nf	number of future times for specification						
	function						
n_p	number of heat sources	Subscripts					
n_q	number of outputs	ext	Exterior				
nt	number of time steps	k	time index				
Р	heat source vector, W						
U	vector of unknown heat sources						
V	vector of known inputs	Superscripts					
vol	volume, m ³	*	noisy temperature				
$Y(n_q)$	output vector	^	estimated value				
		Т	transposition sign				
		-1	inverse of a matrix				
Abbreviatio	ns						
FTS	Future Time Steps						
LPTM	Lumped Parameter Thermal Model						
PMSM	Permanent Magnet Synchronous Machine						
SM	Surface Mounted						

and thermal constraints. The multiple electrical limits and the approaches allowing an increase of the motors specific power density are not discussed here, since the related issues are not the main purpose of the paper and are associated to loadability concepts. The interested reader may refer to [7] for further comprehension of these issues.

Thermal limits in electric motors, majorly related to maximum temperature allowed in electric motors windings depending on the insulation thermal class, are at least of the same order of challenge - if not more - than the electrical limits [8-9]. Excessive temperature increase may irreversibly demagnetize magnets [10] or could lead to shortening motor lifetime due to overheated electric insulation material [11]. The temperature increase in electric machines results from electromagnetic and mechanical losses, generating heat in motor components. Generally, these losses are empirically estimated. Besides being hardly evaluated with analytical and some computational methods, it is believed that even with a motor prototype, the process of determination of the resulting losses and distinguishing their types or origins is quite complicated. Generally, such information can be barely obtained with long iterative methods.

Our study focuses on a Permanent Magnet Synchronous Machine (PMSM). In such machines, the temperatures are very sensitive to losses, specifically at hot spot zones. A possible solution to identify electric motor losses is to measure temperature inside the motor and then apply an inverse method. Such methods require a thermal model with quick solving process. Thus, a thermal nodal network could be a good choice as it is commonly used to carry out numerical simulations in the design process of high-performance electric machines.

Being mathematically ill-posed, the solution of the inverse problem may be unstable and not unique. Hence, several regularization techniques have been developed to ensure stable solutions [12]. Examples of inverse problems for electric machines can be found in the literature with different objectives: geometric, thermal or magnetic parameter estimation, loss mapping, source strength determination from temperature measurements. In [13], the authors propose an inverse problem in order to identify a set of 9 design parameters (geometry, magnetic parameters) for a permanent magnet brushless motor. The results show that the density of the magnetic flux in the air gap is increased, as well as the motor efficiency. Moreover, an inverse method for the identification of heat convection coefficients on coil end-windings from experimental data is presented by [14]. An interesting algorithm is developed and described for this nonlinear inverse heat conduction problem using a lumped parameter model. In [15], a low-order model is built in a first stage to estimate a set of thermal parameters. In a second stage, this model is used to identify thermal losses in the stator of an axial flux PM machine. In [16], the power loss density is estimated using a 2D finite-element thermal model associated to a conjugate gradient method from noisy temperature measurements. A prediction of the strength of the heat source field in an induction motor is also carried out in [17] using a steady state thermal model. Another example of reconstruction of thermal field sources can be found in [18]. The aim is to estimate a 3D source distribution using infrared data. The least square method is used in order to solve the inverse problem.

Concerning now the on-line identification of electromagnetic losses, a recursive parameter estimation method is implemented in [19]. Furthermore, a 3D inverse heat conduction problem is presented by [20] to estimate the time-dependent internal heat flux generated in rotor and stator for high speed electric motor.

This article deals with the temperature prediction in the electric machines at low accessibility locations through the identification of the heat sources and using an inverse technique. The originality of our numerical study lies in the fact that a low order lumped parameter thermal model that requires a very low computation time is used first to identify the unknown heat sources through an inverse procedure, and secondly to compute the temperature field in the electric machine at each time step in real time. The technique used to solve the mathematically ill-posed problem relies on Beck's contribution [21], [22], which takes into account the lagging and damping effects due to the diffusion process by using a sequential method associated to a function specification procedure. This technique acts as a regularization procedure, as it will be shown in this study.

After presenting the electric machine and associated heat transfer modeling, the inverse method applied to identify the heat sources dissipated inside the motor is detailed. Finally, results of parametric studies and corresponding analysis are presented.

II. THERMAL MODEL

A. DESCRIPTION OF THE ELECTRIC MOTOR

The electric motor illustrated in Fig. 1 is a Surface Mounted-Permanent Magnets Synchronous Motor (SM-PMSM) used for hybrid electric aircraft propulsion with specifications found in [23-24]. It is a high efficiency electric motor (more than 96%) and maximum losses resulting in thermal sources are around 16 kW. These losses are divided into winding Joule losses, Iron losses and mechanical losses.

Geometrically, the motor consists in two parts: the rotor (rotating in the center) and the stator around it.

As shown in Fig. 1, the rotor consists of a shaft (brown) supported by two bearings (grey), surrounded by the rotor lamination (yellow) and the permanent magnets (orange). In the studied motor, the considered heat sources are due to mechanical losses and electric losses. Mechanical losses are mainly due to friction inside the bearings ($P_{Bearings}$), and to aerodynamic heating resulting from very high velocity gradient in the air gap between the rotor and the stator (P_{Airgap}) and near the rotor end-caps (lateral surfaces of the rotor) termed P_{Rotor} .

The stator consists of the stator laminations (purple) and the winding (green) whose biggest part is embedded in the stator slots and the remaining part are the end-windings at the extremities of the stator, which are the complex parts in the cavities from geometrical and thermal viewpoints. Windings are heterogeneous elements made up of two types of materials: copper conductors and electric (and thermal) insulation material. The iron losses (P_{Iron}) and Joule losses (P_{Joule}) are thermal sources distributed in the stator laminations and winding respectively rising from the resistivity of copper in winding conductors. Iron losses are produced by magnetic effects (hysteresis and Foucault's current) and the Joule losses in windings are due to Joule effect.

It is assumed in this study that the thermo-physical properties are temperature independent.

The motor cooling system in this case is a combination of a stator water-jacket with water-glycol flowing around the stator, and a water-glycol axial flow at the center of the shaft for rotor cooling. Both types of liquid cooling are being used in electric motors for vehicular applications (examples can be found in [25]). After flowing through the motor, the water-glycol is cooled with a simple heat exchanger: a channel in contact with ambient air through a flat plate. Moreover, high conductivity resin potting of end-windings is used to increase the heat transfer in those critical regions (due to their high Joule losses). This technique is investigated in [26-27].



FIGURE 1. SM-PMSM schematic section with specific temperatures' locations.

These techniques have dynamic impact on motor temperatures and thus influence the extraction of the generated heat resulting from the losses. They must be taken into account in any proposed model of the motors.

B. HEAT TRANSFER EQUATIONS AND MODELING

Temperature distribution within the motor is determined from solving the heat equation associated with initial and boundary conditions. The equation governing heat diffusion and convection in the domain Ω is written as:

$$\rho(M)C_p(M)\left[\frac{\partial T(M,t)}{\partial t} + \vec{\nabla}(M).\vec{\nabla}T(M,t)\right] =$$

$$\vec{\nabla}.\left(k(M)\vec{\nabla}T(M,t)\right) + \sum_{j=1}^{n_p} \left[\frac{P_j(t)}{Vol_j}\chi_j(M)\right] \quad \forall M \in \Omega$$
(1)

Where T(M,t) is the temperature at a point M in the domain and $\vec{v}_{(M)}$ is the velocity field, ρ the density, C_p the specific heat and k the thermal conductivity. P_j is the time varying strength of the j^{th} volumetric heat source and Vol_j the corresponding volume. The characteristic function $\chi_i(M)$ is equal to 1 if M is located in source j, else $\chi_i(M) = 0$.

Heat exchanges between the motor and its environment are modeled through a global heat transfer coefficient *h*. Boundary conditions hence are written as follows:

$$k(M)\vec{\nabla}T(M,t).\vec{n} = h(T_{ext}(t) - T(M,t)) \quad \forall M \in \Gamma$$
⁽²⁾

Moreover, the initial condition is:

$$T(M,0) = T_{init} \quad \forall M \in \Omega \tag{3}$$

If the velocity field is not known in a particular region of domain Ω , the transport term in (1) may be modeled by a convective heat transfer as in (2). Then a convective coefficient *h* has to be evaluated from Nusselt correlations. Besides, the fluid temperature T_{ext} is required.

C. STATE SPACE REPRESENTATION

As analytical solutions for this 2D axisymmetric transient problem are not available, a Lumped Parameter Thermal Model (LPTM) is used to simulate motor thermal behavior according to (1)-(3). It consists in dividing the studied system into N discrete volumes, which are connected together by thermal conductances. For brevity reasons, details of the network are not given in this paper but can be found in [28]. It leads to a system of ordinary differential equations with a low number of parameters. In our study, N=58. This set of equations can be written as:

$$\dot{T}(t) = AT(t) + B_c T_{ext}(t) + B_P P(t)$$
(4)

Where T(t) (dimension N) is the vector of temperatures, function of time t, at the N discretization nodes, $\dot{T}(t)$ its derivative with respect to time t. Matrix A (dimension N,N) is the state matrix which connects temperatures at discretization nodes and contains diffusion and transport terms, as well as terms related to convective boundary conditions. The vector B_c is associated with these latter conditions and links corresponding nodes to environment nodes. Matrix B_p (dimension N,n_p) links discretization nodes to the internal heat sources gathered in vector function $P(t)=[P_{Joule} P_{Iron} P_{Rotor} P_{Airgap} P_{Bearings}]^T$ (dimension n_p).

In what follows, let us separate vector *P* into two sets: unknown heat sources included in *U* (dimension n_U) and known heat sources included in *K* (dimension (n_p-n_U)). The command matrix B_p is then also split into two matrices $B_p = [B_K B_U]$ Defining the vector $V = B_c T_{ext} + B_K K$, the state space representation takes the following form:

$$\begin{cases} \dot{T}(t) = AT(t) + V(t) + B_U U(t) \\ Y(t) = CT(t) \end{cases}$$
(5a)
(5b)

Where an observation matrix C (dimension n_q , N) allows to select n_q temperatures in the whole temperature field T(t) and to store them in vector function Y(t).

In this purely numerical study, a maximum of 9 temperature locations will be used as outputs. These temperatures are depicted in Fig. 1, and are located in the stator lamination (Ts1, Ts2 and Ts3), in the end-winding (Tw) at the surface of the end-winding (Tsw), on the surface of the stator in the airgap (Ta), in the cavity between the rotor end-cap and the frame (Tc), in the bearings (Tb), and in the magnets (Tm). These locations have been chosen given the heat sources but also considering plausible experimental measurement system positions. Temperature in the static part of the motor could be obtained using thermocouples (Ts1, Ts2, Ts3, Tsw, Ta, Tb and Tc). The rotor flange temperature could be obtained by infrared measurement. Tm and Tw are the critical temperatures difficult to obtain experimentally due to their locations, i.e. Tw inside the end-windings is confined and thought to be the hottest point and Tm requires a slip ring due to the rotating rotor.

III. INVERSE PROBLEM SOLUTION

The form of heat transfer model, developed in *subsection II.C*, is convenient to be used now in the inverse procedure. The inverse problem aims at finding an estimation \hat{u} of the input vector U from measured temperatures Y^* . Note that, in this study, real temperature measurements are not available. The temperatures included in vector Y^* derive from the direct problem with known heat sources. In order to simulate measurement errors, each temperature is altered with an additive Gaussian error, whose standard deviation is given by σ .

Subsequently, in order to have simplified and clear annotations, we will write U_k for $U(t_k)$, T_k for $T(t_k)$, Y_k for $Y(t_k)$. A sequential estimation method is adopted. Having the vector of noisy temperatures Y_{k+1}^* , and an estimated \hat{U}_k of U_k , an estimation of \hat{U}_{k+1} is then computed.

Using an Euler implicit scheme and Δt being the time step, (5a) becomes:

$$\dot{T} = \frac{T_{k+1} - T_k}{\Delta t} = AT_{k+1} + V_{k+1} + B_U U_{k+1}$$
(6)

Then:

$$T_{k+1} = (I - A\Delta t)^{-1} [T_k + \Delta t V_{k+1} + \Delta t B_U U_{k+1}]$$
(7)

Hence, according to (5b):

$$Y_{k+1} = C(I - A\Delta t)^{-1} [T_k + \Delta t V_{k+1} + \Delta t B_U U_{k+1}]$$
(8)

Due to the lagging and damping effects of heat diffusion and convection, a variation of heat source strength does not immediately affect the sensors. In order to take into account these effects, Future Time Steps (FTS) [29]-[32] are used. It means that output data at further times t_{k+2} , t_{k+3} , ... are used to correctly estimate \hat{U}_{k+1} .

If *nf* is the number of FTS, then for $1 \le f \le nf$, (8) is written for k+1+f instead of k+1. A temporary approximation of U_{k+1+f} is therefore needed to look for \widehat{U}_{k+1} . In this study, a constant value is chosen:

$$U_{k+1+f} = U_{k+1} \text{ for } 1 \le f \le nf$$
(9)

Hence, a global matrix formulation is obtained as in (10):

$$Y_{k+1}^* = SU_{k+1} + D_k$$
(10)

With:

$$\mathbf{Y}_{k+1}^{*} = \begin{bmatrix} Y_{k+1}^{*} \\ Y_{k+2}^{*} \\ \vdots \\ Y_{k+1+f}^{*} \\ \vdots \\ Y_{k+1+nf}^{*} \end{bmatrix}$$
(11)

$$S = \begin{bmatrix} C(I - A\Delta t)^{-1} B_U \Delta t \\ C[(I - A\Delta t)^{-1} + (I - A\Delta t)^{-2}] B_U \Delta t \\ \vdots \\ C\left[\sum_{j=1}^{f+1} (I - A\Delta t)^{-j}\right] B_U \Delta t \\ \vdots \\ C\left[\sum_{j=1}^{nf+1} (I - A\Delta t)^{-j}\right] B_U \Delta t \end{bmatrix}$$
(12)

And:

$$\mathsf{D}_{k} = \begin{bmatrix} C(I - A\Delta t)^{-1}[T_{k} + V_{k+1}\Delta t] \\ C(I - A\Delta t)^{-2}[T_{k} + V_{k+1}\Delta t] + C(I - A\Delta t)^{-1}V_{k+2}\Delta t \\ \vdots \\ C(I - A\Delta t)^{-(f+1)}[T_{k} + V_{k+1}\Delta t] + C\left[\sum_{j=1}^{f} (I - A\Delta t)^{-j}\right]V_{k+2+f-j}\Delta t \\ \vdots \\ C(I - A\Delta t)^{-(nf+1)}[T_{k} + V_{k+1}\Delta t] + C\left[\sum_{j=1}^{nf} (I - A\Delta t)^{-j}\right]V_{k+2+nf-j}\Delta t \end{bmatrix}$$
(13)

The size of vector Y^* is $(nf+1) \times n_q$ and matrix *S* is of size $((nf+1) \times n_q, n_U)$. The objective is to identify the pseudosolution \widehat{U}_{k+1} of the inverse problem, such that $Y^*-Y\approx 0$, where *Y* is the temperature vector computed by LPTM. As matrix *S* is not square due to the addition of future time steps, the least square method is used to solve (10) and leads to the sequential solution:

$$\widehat{U}_{k+1} = \left(\mathsf{S}^{T}\mathsf{S}\right)^{-1}\mathsf{S}^{T}\left(\mathsf{Y}_{k+1}^{*}\mathsf{-}\mathsf{D}_{k}\right)$$
(14)

In order to evaluate the accuracy of the estimations for the different studied cases, a mean quadratic discrepancy between exact heat sources (vector U^{exact}) and estimated ones (vector \hat{U}) is defined as follows:

$$\sigma_{U} = \left[\frac{1}{n_{U} \times (nt - nf)} \sum_{i=1}^{n_{U}} \sum_{k=1}^{nt - nf} \left(\widehat{U}_{i,k} - U_{i,k}^{exact}\right)^{2}\right]^{1/2}$$
(15)

Where *nt* is the number of time steps in the inverse problem. It is worth to remind that, in a practical application, this quantity is not reachable.

In a similar way, for each output Y_i , a mean quadratic error between measured temperatures (or artificially noisy temperatures) and those computed with the estimated set of heat sources strengths, is defined as follows:

$$\sigma_{Y_i} = \left[\frac{1}{(nt - nf)} \sum_{k=1}^{nt - nf} \left(\hat{Y}_{i,k} - Y_{i,k+1}^*\right)^2\right]^{1/2}$$
(16)

When the unknown heat sources \widehat{U}_{k+1} are identified, they will be used as inputs in the direct problem in order to predict temperature in inaccessible parts of the motor, for example critical temperatures *Tcrit* in windings or magnets. The complete procedure is illustrated in the block diagram of Fig. 2.



FIGURE 2. Block diagram: use of LPTM for both inverse and direct problems.

At each time-step, the model is solved in inverse problem to get motor losses. The inputs for inverse LPTM are the noisy temperature at accessible points (sensor temperature measurement on a real bench).

Then, forward LPTM is used to determine indicated critical temperatures.

IV. RESULTS AND ANALYSIS

A scenario of motor operation is chosen corresponding to a profile mission in an aircraft flight application. The corresponding profiles of heat sources generated in the motor are depicted in Fig. 3 during 3600 s. The different phases correspond to: taxi-out phase (0 to 450s), take off and climb (450 to 900s), cruise flight (900 to 2250s), descent and landing (2250s to 2700s) and taxi-in (2700s to 3600s).



FIGURE 3. PMSM heat sources evolution as function of time PMSM.

Rotor, Airgap, and Bearings losses remain relatively low and quite constant all along the mission whereas Joule and Iron losses change significantly. Those losses are particularly high during the take-off phase as they are linked to propulsion power. Airgap and Bearings losses depend principally on the rotor rotation speed, which varies slightly between take-off and cruise, and is very low during taxi phases.

Temperatures are then computed using LPTM and a time step equal to 1 s. Among these temperatures, a set of observables is chosen depending on the different case studies (see *subsection II.C*). A noise of standard deviation equal to $\sigma=0.1^{\circ}C$ is added to these observables so that they are used as temperature measurements in various inverse problems cases.

A. CASE 1: IDENTIFICATION OF IRON LOSSES

In this first case, among the five heat sources, four are supposed to be known ($P_{Joule} P_{Rotor} P_{Airgap} P_{Bearings}$) and only iron losses ($P_{Iron}(t)$) have to be identified. Fig. 4a shows the evolution of temperature Ts1 with respect to time. From this temperature, the inverse problem is solved in order to identify $P(t)=P_{Iron}(t)$ knowing $K(t)=[P_{Joule} P_{Rotor} P_{Airgap} P_{Bearings}]^T$. Hence, for this case, $n_q=n_U=1$.

The procedure is sequential, as mentioned before, expressed in (14), and depends on the number of future time steps. Results of identification are presented in Fig. 4b and Fig. 4c for respectively nf=2 and nf=10. It is remarkable that, without regularization (nf=0), the oscillations of $P_{Iron}(t)$ are huge. Hence, we have chosen to draw $P_{Iron}(t)$ for nf=2 in Fig. 4b.

It clearly appears that the number of future time steps filters the oscillations of P_{Iron} , what is confirmed by the quadratic criterion σ_U in Table 1. The main criterion, and the only one to be computed in a real case, is σ_Y . Its evolution with respect to *nf* is quite interesting. For *nf*=0, even if the problem is ill-posed, large oscillations in \hat{U} allow to follow exactly the evolution of the noisy temperature. Hence, as indicated in Table 1, $\sigma_Y \approx 0^\circ C$. When *nf* increases, regularization tends to damp the time evolution of \hat{U} , leading to a bias in the temperature computation. We choose then the number *nf* for which $\sigma_Y \geq \sigma$ to "get out" of the noise level. The identification of $P_{Iron}(t)$ for *nf*=10 is depicted in Fig. 4c and shows a very good agreement with the exact evolution of these losses drawn in Fig. 3.



FIGURE 4. Iron losses identification from node Ts1 with nf=2 and nf=10 and computation of critical temperatures Tw and Tm.

TABLE 1. Inversion results for iron losses identification from Ts1; bold lines correspond to results shown in Fig. 4.

nf	$\sigma_{U}(W)$	σ_{Ts1} (°C)
0	5303.1	3e-14
2	936.7	0.0871
4	451.3	0.0935
6	304.7	0.0964
8	249.4	0.0988
10	229.2	0.1009
12	224.0	0.1034

Once P_{Iron} is identified, the aim is to predict temperatures Tw and Tm in inaccessible parts of the electric motor. They have been computed for the different values of nf and are shown in Fig. 4d for nf=0. Note that, for nf=0, temperature prediction occurs without any delay whereas it is executed with a 10s delay with nf=10 since the time step Δt is equal to 1s. For any value of nf, the temperature evolution in critical zones is the same, due to the fact that oscillations in \hat{U} are damped by the thermal path and inertia of the system. Same as in (15), critical temperatures Tw and Tm are compared to the exact temperatures computed with the exact inputs. For nf=0, a quadratic error $\sigma_{Tcrit}=0.003^{\circ}C$ is obtained, which gives a very accurate prediction.

Now, a new test is carried out using a second measurement point Ts2 located in the stator. Graphical results are summarized in Fig. 5.

It appears obviously that the oscillations of $P_{Iron}(t)$ decreased compared to Fig. 4, which is confirmed by Table 2 where σ_U goes from 229 W to 200 W for nf=10. It is interesting to observe that in this case $(n_q=2, n_U=1)$ it is not possible to obtain $\sigma_T \approx 0^\circ C$ for each of the two outputs from only one heat source $P_{Iron}(t)$.



FIGURE 5. Iron losses identification from nodes Ts1 and Ts2 with nf=2 and nf=10 and computation of critical temperatures Tw and Tm.

TABLE 2.Inversion results for iron losses identification from Ts1 and Ts2; bold lines correspond to results shownin Fig. 5.

nf	$\sigma_{U}\left(W ight)$	σ_{Ts1} (°C)	σ_{Ts2} (°C)
0	3713.2	0.0690	0.0700
2	673.8	0.0921	0.0929
4	336.0	0.0957	0.0960
6	237.5	0.0976	0.0975
8	205.1	0.0992	0.0987
10	200.1	0.1011	0.1001
12	205.0	0.1034	0.1020

The temperature prediction in windings and magnets is also very accurate with an error $\sigma_{Tcrit}=0.002^{\circ}C$ for nf=0.

B. CASE 2: IDENTIFICATION OF JOULE AND IRON LOSSES

For the second case, Joule losses (P_{Joule}) are also unknown. The evolutions of $P_{Joule}(t)$ and $P_{Iron}(t)$ have to be identified. Firstly, outputs Ts1 and Tsw located in the stator and at the winding surface are used. It can be remarked in Fig. 6 that both heat sources are well identified when increasing the number of future time steps. In this case, nf=12 gives $\sigma_Y \ge \sigma$ for each output, as shown in Table 3.

The condition number of the matrix S^TS to invert is also indicated. It is noteworthy that this last one decreases when the number of future time steps increases, which confirms the regularization effect of the function specification, knowing that in case 1, S^TS is a scalar, hence its condition number is equal to 1.



FIGURE 6. Joule and iron losses identification and computation of critical temperatures from nodes Ts1 and Tsw with nf=2 and nf=12.

Compared to case 1, the winding temperature is quite sensitive to the oscillations in Joule losses identification, hence the error increases: $\sigma_{Tcrit}=1.161^{\circ}C$ for nf=0.

TABLE 3. Inversion results for Joule and iron losses identification from Ts1 and Tsw; bold lines correspond toresults shown in Fig. 6.

nf	$\sigma_{U}(W)$	σ_{Ts1} (°C)	σ_{Tsw} (°C)	$CN(S^{T}S)$
0	4125.2	3e-14	6e-15	3.74
2	741.2	0.0859	0.0869	3.64
4	359.2	0.0926	0.0931	3.54
6	246.4	0.0949	0.0971	3.45
8	209.3	0.0970	0.1007	3.36

10	198.9	0.0992	0.1051	3.28
12	200.4	0.1017	0.1104	3.20
14	207.9	0.1047	0.1168	3.12

In the other hand, by adding two other outputs Ts2 and Ts3, located in the stator, better inversion results are again obtained, as depicted in Fig. 7 and Table 4, since σ_U goes from 200 W to 180 W.



FIGURE 7. Joule and iron losses identification and computation of critical temperatures from nodes Ts1, Tsw, Ts2 and Ts3 with nf=2 and nf=10.

The addition of outputs acts as a filter in the inversion procedure and the condition number of matrix S^TS decreases also. The temperature prediction in windings and the magnets remains relatively accurate with an error $\sigma_{Tcrit}=1.158^{\circ}C$ for nf=0.

TABLE 4.Inversion results for Joule and iron losses identification from Ts1, Tsw, Ts2 and Ts3; bold linescorrespond to results shown in Fig. 7.

nf	$\sigma_{U}(W)$	σ_{Ts1} (°C)	σ _{Tsw} (°C)	σ _{Ts2} (°C)	σ _{Ts3} (°C)	$CN(S^{T}S)$
0	2849.0	0.0803	3e-5	0.0832	0.0816	1.26
2	518.4	0.0941	0.0867	0.0974	0.0952	1.23
4	262.6	0.0964	0.0937	0.0994	0.0974	1.21
6	194.5	0.0976	0.0975	0.1004	0.0987	1.18
8	179.2	0.0990	0.1012	0.1012	0.0999	1.16
10	180.3	0.1007	0.1055	0.1025	0.1016	1.13
12	188.3	0.1028	0.1109	0.1040	0.1038	1.11

C. CASE 3: IDENTIFICATION OF ALL LOSSES

In this third case, the five heat sources altogether are considered unknown and have to be identified simultaneously from five outputs Ts1, Tc, Tsw, Ta and Tb, whose temperature evolutions are given in Fig. 8a.

These nodes are located in the stator, in the cavity, on the surface of the windings, in the air-gap and on the bearings respectively.



FIGURE 8. Identification of all the heat sources and computation of critical temperatures from nodes Ts1, Tc, Tsw, Ta and Tb with *nf=4* and *nf=12*.

TABLE 5. Inversion results for all losses identification from Ts1, Tc, Tsw, Ta and Tb; bold lines correspond toresults shown in Fig. 8.

nf	$\sigma_{U}(W)$	σ_{Ts1} (°C)	σ_{Tc} (°C)	σ _{Tsw} (°C)	σ _{Ta} (°C)	σ_{Tb} (°C)	$CN(S^{T}S)$
0	11990.5	3e-14	1e-14	6e-15	2e-14	2e-14	2.46e6
2	1859.6	0.0881	0.0861	0.0842	0.1213	0.0876	4.48e5
4	859.3	0.0973	0.0930	0.0913	0.1639	0.0996	1.82e5
6	529.7	0.1038	0.0959	0.0952	0.1982	0.1129	9.90e4
8	382.1	0.1100	0.0975	0.0990	0.2282	0.1295	6.23e4
10	305.8	0.1169	0.0987	0.1037	0.2557	0.1490	4.85e4
12	263.9	0.1243	0.1000	0.1096	0.2816	0.1712	4.29e4
14	240.1	0.1326	0.1014	0.1167	0.3061	0.1950	3.82e4

It can be remarked that the temperature magnitude is very different according to the outputs locations.

As for previous cases, the quadratic errors and the condition number are given with respect to the number of future time steps. In this case, the condition number of the matrix S^TS is very high, hence, the noise in temperature measurements may be greatly amplified by the inverse problem. As in previous cases, nf=0 corresponds to a square case ($n_q=n_U=5$), thus, even if the problem is ill-posed, large oscillations in \hat{U} allow to follow exactly the noisy temperature evolutions. Hence, as indicated in Table 5, $\sigma_Y \approx 0^\circ C$. An increase of nf leads to $\sigma_Y \ge \sigma$ for all outputs for nf=12. Fig. 8c shows large oscillations in P_{Rotor} due to the fact that the outputs are located far from this heat source.

Concerning now the temperature prediction in windings and magnets (Fig. 8d), the error is completely acceptable with $\sigma_{Tcrit}=1.178^{\circ}C$ for nf=0.

V. CONCLUSION

The study presented in this paper deals with the determination of electric motors' losses or heat sources, and prediction of critical temperature at non-accessible parts. A low-order lumped parameter thermal model is used to solve both inverse and direct problems. The method is sequential and includes regularization with the function specification technique. Different configurations are tested according to the number of heat sources to be identified. For each case, the influence of the required number of outputs to be used to solve the inverse problem and the number of future time steps are analyzed. The technique gives accurate prediction of the temperature in unreachable critical parts of the motor and is computationally efficient: only 10⁻³ s at each time step is needed for identification of the heat sources and prediction of the motor temperatures.

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Declarations

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