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Nickel(II) salicylaldiminates: Re-visiting a classic☆

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ABSTRACT

A series of Ni(II) trinuclear complexes, $[Ni_3(sala-X)_6]$ (sala-X = o -[(p-X-phenylimino)methyl]phenol; $X = Br 1-Br$, Me 1-Me and OMe 1-OMe), have been prepared and fully characterised. X-ray crystallographic studies reveal that the complexes are composed of three face-sharing octahedral Ni(II) metal ions bridged by the phenoxide oxygens of the sala-X ligands. Magnetic studies indicate that the magnetic Ni(II) centres are ferromagnetically coupled with the substituent group having only a minor impact on the magnitude of coupling. AC susceptibility studies show no evidence of single-molecule magnet behaviour. DFT calculations also support ferromagnetic coupling with smaller Ni-O-Ni angles leading to slightly larger coupling constants in line with previous studies.

1. Introduction

Schiff bases are common ligands in coordination chemistry largely due to their ease of synthesis and ready modifcation [1]. One of the most studied of the Schiff bases are salicylaldimines, typically prepared by the reaction of salicylaldehyde and a primary amine. Aryl amines are particularly diverse and their Ni(II) complexes have received considerable attention in the area of olefn oligomerisation and polymerisation [2–9]. The complexes have the general formula [Ni(sala-X)(R)(L)] (R = alkyl ligand, $L = co$ -ligand e.g. MeCN, pyridine; Scheme 1) and the tunability of both the phenyl and salicylaldimine rings permits substantial modifcation of catalytic behaviour.

In contrast, the simple $[Ni(sala-X)_2]$ complexes were first reported in 1962 [10] and over the course of a decade their basic spectroscopic data were collected and reported $[11-13]$. The pattern of substitution on the phenyl ring was found to be important with ortho-substituted complexes being diamagnetic in solution while para- or meta-substituted systems were paramagnetic. At the time, this was interpreted as the ortho complexes being square planar, while the other complexes were thought to associate in solution. The nature of the association was unclear, and although [Ni(sala-Me)₂] was reported to display a trinuclear octahedral structure in the solid state, [14] the structure is not present in the CSD. In seeking to defnitively answer the question of the structure of these complexes, we have prepared [Ni₃(sala-X)₆] (sala-X = o -[(p-X-phenylimino)methyl]phenol; $X = Br 1-Br$, Me 1-Me and OMe 1-OMe) determined their structures and explored their magnetic behaviour, with the results supported by DFT calculations.

2. Experimental

2.1. General remarks

All compounds were prepared with reagent-grade solvents. All other chemicals were purchased from Sigma-Aldrich or TCI Chemicals and used as received. Hsala-X were prepared according to literature procedures [10] using diisopropylether instead of MeOH as the solvent.

2.2. Synthesis of $[Ni_3(sala-X)_6]$

All complexes were synthesized according to the following general procedure:

To a solution of Hsala-X (2 mmol) in tetrahydrofuran (10 mL) was

Scheme 1. Common Ni(II) sala-X complexes used in olefin oligomerisation.

added NaOMe (0.1082 g, 2 mmol) and methanol (1 mL) to increase solubility. NiCl₂⋅6H₂O (1 mmol) was added and the solution was stirred overnight. The solution was evaporated to ca. 2 mL and fltered through a Buchner funnel. The solid was washed with ethanol (10 mL) and then recrystallized from CH_2Cl_2/h exane giving the compounds as green crystalline solids.

[Ni₃(sala-Br)₆] **1-Br**

Yield 49%. IR (cm^{−1}): 3022(m), 2899(m), 1610(sh). UV–Vis (ε, M^{-1} ·cm⁻¹) 629 (123).

[Ni₃(sala-Me) $_6$] 1-Me

Yield 29%. IR (cm^{−1}): 3001(m), 2938(m), 1615(sh). UV–Vis (ε, M^{-1} ·cm⁻¹) 600 (212).

 $[Ni_3(sala-OMe)_6]$ 1-OMe

Yield 23%. IR (cm^{−1}): 3002(m), 2833(m), 1614(sh). UV–Vis (ε, M^{-1} ·cm⁻¹) 615 (142).

2.3. Spectroscopic studies

Infrared spectra were recorded on a Bruker Tensor 27 FT-IR spectrometer with OPUS data collection program in the range of 400–4000 ${\rm cm^{-1}}$ (as KBr discs). UV–Vis spectra were recorded in Acetone, DCM and DMF at room temperature on an Avantes Fiber Optics Spectrometer with Avalight-DHC and Avaspec ULS2048XL-USB2 in the range of 200–1000 nm.

2.4. X-ray crystallographic studies

Crystallographic data for the basic complexes were collected using a Rigaku SuperNova diffractometer with a HyPix 3000 detector using Cu_{α} radiation ($\lambda = 1.54184$ Å). The data were scaled, integrated and an absorption correction applied in CrysAlisPro. The structures were solved in SHELXT [15] using intrinsic phasing and refned by full matrix least squares minimisation on F^2 using SHELXL [16]. All non-hydrogen atoms were refned anisotropically. Hydrogen atom positions were calculated geometrically and refned using the riding model. All pictures were generated with OLEX². [17] Powder X-ray diffraction data were measured on the same diffractometer using Cu_{α} radiation ($\lambda = 1.54184$) Å). The samples were ground and then suspended in Fomblim Y oil and the data collected between $2\theta = 5-80^\circ$.

2.5. Magnetic studies

Data on 1-Br and 1-OMe were collected with a Quantum Design MPMS 5 SQUID magnetometer under an applied feld of 1 T or 0.1 T over the temperature range 2–300 K. The powdered or polycrystalline samples were placed in gel capsules and care was taken to allow long thermal equilibration times at each temperature. DC magnetic susceptibility measurements on 1-Me were carried out with a Quantum Design SQUID magnetometer MPMS-XL. This magnetometer works between 1.8 and 400 K for dc applied fields ranging from -7 to 7 T. Ac susceptibility measurements were performed with a Quantum Design PPMS-9 susceptometer working with an oscillating ac felds of 1 up to 10 Oe with a frequency between 10 and 10,000 Hz. Consistent dc susceptibility at 0.1 T and in-phase ac susceptibility in zero-dc feld have been obtained between 1.85 and 15 K without detection of an out-of-phase ac signal up to 10,000 Hz, up to 2.8 T and above 1.9 K. These magnetic measurements were performed on a microcrystalline sample (13.58 mg) sealed in a polyethylene bag ($3 \times 0.5 \times 0.02$ cm; typical mass of 12.37 mg) and restrained in mineral oil (6.78 mg). Magnetic data were corrected for the sample holder, mineral oil, and diamagnetic contributions.

3. Results and discussion

3.1. Synthesis, structure and spectroscopic characterizations

The compounds $[Ni_3(sala-X)_6]$ $(X = Br 1-Br, Me 1-Me$ and OMe 1-OMe) were synthesized by reacting NiCl₂⋅H₂O with two equivalents of Na(sala-X), made from Hsala-X and NaOMe, as the reported procedure [10] did not give the compounds cleanly in our hands. The IR spectroscopic data is consistent with previous reports with an imine stretch observed between 1610 and 1615 $\rm cm^{-1}$. UV–Vis spectroscopic data were recorded in CH_2Cl_2 and reveal a band at 629, 600 and 615 nm for 1-Br, 1-Me and 1-OMe, respectively (Fig. S1). Interestingly, the weaker band at 1000 nm that is reported in CHCl3, and characteristic of octahedral geometry, is not observed but this may be due to the lower solubility of the compounds in $CH₂Cl₂$.

Recrystallisation of 1-Br, 1-Me and 1-OMe from $CH_2Cl_2/hexane$ gives green crystals. Single crystal X-ray crystallography reveals that the compounds belong in the Monoclinic $C2/c$ space group and have a trinuclear structure as shown in Fig. 1 (see Table 2 for full crystallographic details). Powder X-ray diffraction studies reveal that the crystal structures are reflective of the bulk materials (see Fig. S3). The compounds are isostructural and centrosymmetric with the Ni2 centre sitting on an inversion centre. The asymmetric nature of the sala-X ligand results in one of the Ni1 centres having N_3O_3 coordination sphere, while the central Ni2 has an O_6 donor set. The Ni(II) centres are linked through bridging phenoxides creating face-sharing octahedra. This arrangement is supported by weak π - π interactions, (centroid⋅⋅⋅centroid = 3.71 Å 1-Br; 3.71 Å 1-Me; 3.70 Å 1-OMe, Fig. $S4$ in the ESI) between the central phenyl rings and may be responsible for the trinuclear structure. The Ni-N/O bond lengths are typical of octahedral Ni(II), and while the bond lengths around Ni2 are similar to those observed at Ni1, the former is much more distorted as evidenced by the signifcantly higher *Σ* values $(\Sigma = \sum_{i=1}^{12} |90 - \alpha_i|)$, where α_i are the twelve cis N/O-Ni-N/O angles). The Ni1-Ni2 distance is virtually identical across the series at ca. 2.82 Å (Table 1) suggesting that magnetic communication is likely to be present and efficient across the face-shared bridge.

Fig. 1. Structure of $[Ni_3(sala-Br)_6]$ 1-Br with thermal ellipsoids drawn at 50%. For clarity, only the donor atoms in the asymmetric unit are labelled and the hydrogen atoms have been omitted (C = grey, N = blue, O = red, Br = redbrown, $Ni = green$).

Table 1 Selected bond lengths and angles for 1-Br, 1-Me and 1-OMe (Å,◦).

	1-Br	1-Me	1-OMe
$Ni1-O1$	2.037(3)	2.0560(10)	2.0475(16)
$Ni1-O2$	2.046(3)	2.0449(10)	2.0454(16)
$Ni1-O3$	2.050(3)	2.0487(10)	2.0546(17)
Ni1-N1	2.099(3)	2.0789(12)	2.088(2)
$Ni1-N2$	2.081(3)	2.0790(12)	2.080(2)
$Ni1-N3$	2.070(3)	2.0988(12)	2.082(2)
Ni1-Ni2	2.8116(6)	2.8248(2)	2.8180(4)
$Ni2-O1$	2.071(3)	2.0632(10)	2.0679(16)
$Ni2-O2$	2.045(3)	2.0554(10)	2.0631(16)
$Ni2-O3$	2.066(3)	2.0680(9)	2.0583(16)
Ni1-01-Ni2	86.13(10)	86.59(4)	86.43(6)
Ni1-02-Ni2	87.04(10)	87.09(4)	86.61(6)
Ni1-03-Ni2	86.16(10)	86.66(4)	86.49(6)
Σ Ni1 ^a	71.7	72.0	71.2
Σ Ni2 ^a	147	158	146

^a Calculated using OctaDist. [18].

3.2. Magnetic studies

The temperature dependence of the magnetic susceptibility of the three complexes were studied from 1.8 to 300 K (Fig. 2). In all cases, ac susceptibility studies in zero or applied feld show no evidence of single molecule magnet behaviour. The compounds exhibit almost identical behaviour and thus we discuss here only 1-Me in detail. At room temperature, the χT product is 4.3 cm³ K mol $^{-1}$, consistent with the expected value for the presence of three isolated $Ni(II)$ ions $(S = 1)$ with a Landé g-factor slight above 2. When the temperature is lowered, the χ T product increases continuously to reach a maximum of 6.7 $\text{cm}^3 \text{ K} \, \text{mol}^{-1}$ around 11 K, indicating dominant ferromagnetic interactions within the complex. Below 4 K, the χ T product drops down to 5.3 cm³ K mol⁻¹ at 1.85 K suggesting the presence of weak antiferromagnetic interactions, Zeeman depopulation effects and/or a significant magnetic anisotropy. The experimental $(\chi T)_{\text{max}}$ at 11 K agrees with a $S_T = 3$ total spin ground state of this trinuclear complex expected for three ferromagnetically arranged $S = 1$ Ni(II) spins. The $S_T = 3$ ground state is also supported by the field dependence of the magnetization at low temperatures, which reaches 6.7 μ B at 7 T and 1.8 K. The magnetization follows relatively well a simple $S = 3$ Brillouin function with $g = 2.20(2)$. This indicates that the intrinsic magnetic anisotropy of the Ni(II) metal ion is very small as also

Fig. 2. Temperature dependence of the χT product (where χ is the molar magnetic susceptibility that equals M/H per complex) collected in an applied dc magnetic feld of 0.1 T for 1-Br (blue dots), 1-Me (red dots) and 1-OMe (black dots). Inset: View of the low temperature region of the main figure below 50 K. The solid lines are the best fits to the model described in the text.

shown by the almost superposition of the M vs H/T data (see Figs. S5–S7 in the ESI). Therefore, antiferromagnetic intermolecular interactions are relevant and dominant at low temperature in this system to explain the decrease of the *χ*T product below 11 K.

Considering the structure, the magnetic susceptibility can be modelled using a linear Heisenberg tri-spin $S = 1$ model. In keeping with this structural motif, the Heisenberg spin Hamiltonian can be written as follows:

$$
\mathscr{H} = -2J_{12}\overrightarrow{S}_{\text{Ni2}}\hat{A}\cdot\left(\overrightarrow{S}_{\text{Ni1}a} + \overrightarrow{S}_{\text{Ni1}b}\right) \tag{1}
$$

where J is the exchange interactions within the trinuclear complex between adjacent Ni(II) through tris-phenoxido bridge; S_i is the spin operators for each center with $S_i = 1$. The application of the van Vleck equation [19] to the Kambe's vector coupling scheme [20], allows us to determine an analytical expression of the magnetic susceptibility in the weak field approximation:

$$
\chi_{Ni_3}T=\frac{2Ng^2\mu_B^2}{k_BT}\frac{14+5e^{\frac{-2I_{12}}{k_BT}}+e^{\frac{-4I_{12}}{k_BT}}+6e^{\frac{-6I_{12}}{k_BT}}+6e^{\frac{-6I_{12}}{k_BT}}}{7+5e^{\frac{-4I_{12}}{k_BT}}+3e^{\frac{-6I_{12}}{k_BT}}+8e^{\frac{-6I_{12}}{k_BT}}+e^{\frac{-8I_{12}}{k_BT}}+3e^{\frac{-10I_{12}}{k_BT}}}
$$

The above model reproduces the experimental data very well down to 20 K but below 20 K the introduction of intercomplex magnetic interactions in the frame of the mean-feld approximation is required:

$$
\chi = \frac{\chi_{Ni_3}}{1 - \frac{2z\Gamma}{Ng^2\mu_B^2}\chi_{Ni_3}}
$$

where χ_{Ni3} is the susceptibility of the non-interacting complexes, z the number of nearest neighbors and J*'* is the magnetic interactions between $Ni₃ complexes [21,22]$. A good fit of the data for **1-Me** has been achieved using $J_{12}/k_B = +11.4(2)$ K, $zJ'/k_B = -0.097(2)$ K and $g_{Ni} = 2.23(2)$, while the parameters obtained in the cases of 1-Br and 1-OMe are given in Table 3. These values are broadly comparable with a range of Ni(II) trinuclear complexes [23–25]. As mentioned before, the sign of the magnetic interactions implies that this complex possesses an $S_T = 3$ spin ground state. This result is consistent with the feld dependence shown in supporting information, that is very close to a $S = 3$ Brillouin function. Caution is needed regarding zJ*'* as it could phenomenologically contain the contribution from the Ni(II) magnetic anisotropy. While a numerical model considering only the local magnetic anisotropy is unable to reproduce the experimental data (using PHI) [26], a model considering the local magnetic anisotropy in addition to intermolecular interactions does not significantly improve the theory/experiment agreement.

3.3. Theoretical studies

Calculations based on Density Functional Theory provide an excellent estimation of the exchange coupling constants in polynuclear transition metal complexes despite the very small energy differences involved [27,28]. Detailed description of the computational procedure employed to calculate the exchange coupling constants in polynuclear complexes has been previously provided by some of us [29–30]. The calculations were performed using Gaussian16 B.01 version [31] with the B3LYP functional [32] and an all electron basis set (Def2TZV) [33]. For the studied symmetric trinuclear Ni(II) complexes, there are two intramolecular exchange interactions (see Table 4), and we employed three spin configurations to estimate these two J values: the high spin solution $(S = 3)$ and two $S = 1$ wavefunctions obtained with the spin inversion of the terminal or central nickel center. The magnetostructural correlation with the calculated J_{12} values shows the expected trend with larger ferromagnetic coupling when the Ni-O-Ni angle becomes smaller [34]. This fact is due to the well-known dependence, complexes with smaller bridging Ni-O-Ni angles have a weaker antiferromagnetic term. At the error bar of the experimental and theoretical data, the calculated values are in good agreement with the experimentally fitted parameters

Table 3

Magnetic ftting parameters for 1-Br, 1-Me and 1-OMe.

Table 4

DFT calculated intramolecular exchange coupling constants (J_{12} next- and J_{13} next-nearest neighbour interaction) for three trinuclear Ni(II) complexes using the experimental X-ray structures.

(Tables 2 and 3). The next-nearest neighbour interactions are weak and antiferromagnetic, as usual, but they do not follow a clear tendency with the Ni-O-Ni angle.

As an example, the spin density map of the $S = 3$ ground state of 1-Br system is represented in Fig. 3. As the coupling is ferromagnetic the main contributions in the three nickel atoms are positive due to the unpaired electrons in the eg orbitals. Consequently, the positive lobes are oriented towards the coordinated ligand atom. A small negative population in the metal atoms (see yellow lobes in Fig. 3) are due to the spin polarization mechanism $[35]$ of the non-bonding t_{2g} orbitals. In the atoms of the ligands corresponding the frst coordination sphere, the spin delocalization mechanism is predominant [36]. This is due to the antibonding nature of the e_g orbitals bearing the unpaired electrons in the metal ion with a large mixing with the ligand orbitals. Thus, positive densities are also found on the atoms of the frst coordination sphere.

4. Conclusions

To conclude, we have defnitively confrmed that the Ni(II) salicylaldiminate complexes with para-substituents exist as face-sharing trinuclear Ni(II) complexes in the solid state. In these complexes, the magnetic susceptibility measurements demonstrate that the Ni(II) spins are ferromagnetically coupled, consistent with the Ni-O-Ni angle of 86–87◦. DFT results confrm the ferromagnetic nature of the exchange interactions. Moreover, the calculated J_{12} values indicate the expected magnetostructural trend with the smaller bridging Ni-O-Ni angle showing stronger ferromagnetic coupling. Preliminary studies suggest that these trinuclear Ni(II) complexes can be used in the construction of 1D coordination polymers and these will be reported in due course.

CRediT authorship contribution statement

Supaporn Saechio: Methodology, Formal analysis, Investigation. Rodolphe Clérac: Methodology, Formal analysis, Investigation, Resources, Data curation, Writing - review & editing, Funding acquisition. Keith S. Murray: Data curation, Writing - review & editing, Funding acquisition. Wasinee Phonsri: Formal analysis, Investigation, Data curation. Eliseo Ruiz: Methodology, Formal analysis, Investigation, Resources, Data curation, Writing - review & editing, Funding acquisition. Phimphaka Harding: Conceptualization, Formal analysis, Writing review $\&$ editing, Visualization, Supervision, Funding acquisition. David J. Harding: Conceptualization, Formal analysis, Resources, Visualization, Writing - original draft, Writing - review & editing, Supervision, Funding acquisition.

Declaration of Competing Interest

The authors declare that they have no known competing fnancial

Fig. 3. Spin population map corresponding the $S = 3$ ground state of 1-Br calculated with the B3LYP method. The purple and yellow colors indicate positive and negative values, respectively. The isosurface corresponds to 0.14 eÅ $^{-3}$.

interests or personal relationships that could have appeared to infuence the work reported in this paper.

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Appendix A. Supplementary data

CCDC 2079808-2079810 contains the supplementary crystallographic data for 1-Br to 1-OMe, respectively. These data can be obtained free of charge via http://www.ccdc.cam.ac.uk/conts/retrieving.html, or from the Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: (+44) 1223-336-033; or e-mail: depos $it@cedc.cam.ac.uk.$

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Nickel(II) Salicylaldiminates: Re-visiting a classic

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Figure S1: IR spectra of **1-Br**, **1-Me** and **1-OMe** as KBr discs.

Figure S2: UV-Vis spectra of 1-Br, 1-Me and 1-OMe in CH₂Cl₂.

Crystal structures

Table S1 Crystal data and structure refinement for **1-Br**, **1-Me** and **1-OMe**.

	$1 - Br$	1-Me	1-OMe
Empirical formula	$C_{78}H_{54}Br_6N_6Ni_3O_6$	$C_{84}H_{72}N_6Ni_3O_6$	$C_{28}H_{24}N_{2}NiO_{4}$
Formula weight	1826.86	1437.60	511.20
Temperature/K	149.99(10)	200.01(10)	149.7(6)
Crystal system	monoclinic	monoclinic	monoclinic
Space group	C2/c	C2/c	C2/c
a/Å	23.5184(2)	22.88299(13)	23.96018(19)
b/Å	14.74717(14)	15.13748(11)	14.80924(11)
c/Å	20.75867(17)	20.57870(12)	20.44679(15)
α /°	90	90	90
$\beta/^\circ$	103.0924(8)	102.1199(6)	99.2758(7)
V°	90	90	90
Volume/Å ³	7012.58(11)	6969.39(8)	7160.31(9)
z	4	4	12
$p_{calc}g/cm^3$	1.730	1.370	1.423
μ /mm ⁻¹	5.418	1.420	1.484
F(000)	3624.0	3000.0	3192.0
Crystal size/mm ³	$0.459 \times 0.201 \times 0.1320.245 \times 0.115 \times 0.0590.301 \times 0.125 \times 0.091$		
Radiation		CuKa (λ = 1.54184) CuKa (λ = 1.54184) Cu Ka (λ = 1.54184)	
20 range for data collection/°	7.128 to 138.676	7.05 to 138.212	7.042 to 138.55
Index ranges	$17, -25 ≤ I ≤ 17$	17, $-24 \leq$ \leq 24	-28 ≤ h ≤ 28, -17 ≤ k ≤ -27 ≤ h ≤ 26, -18 ≤ k ≤ -26 ≤ h ≤ 29, -17 ≤ k ≤ $17, -24 \leq l \leq 22$
Reflections collected	35988	28751	37127
Independent reflections	6513 [R_{int} = 0.0383, $R_{\text{sigma}} = 0.0180$	6464 $[R_{int} = 0.0206]$ $R_{\text{sigma}} = 0.0138$	6638 [R_{int} = 0.0476, $R_{\text{sigma}} = 0.0231$
Data/restraints/parameters	6513/0/447	6464/0/450	6638/36/477
Goodness-of-fit on F ²	1.092	1.038	1.044
Final R indexes $[1 \ge 2\sigma(1)]$	$R_1 = 0.0517$ $wR_2 = 0.1402$	$R_1 = 0.0291$ $wR_2 = 0.0797$	$R_1 = 0.0518$ $wR_2 = 0.1443$
Final R indexes [all data]	$R_1 = 0.0524,$ $wR_2 = 0.1407$	$R_1 = 0.0306$ $wR_2 = 0.0807$	$R_1 = 0.0542$ $wR_2 = 0.1465$
Largest diff. peak/hole / e \AA ⁻³	$1.85/-1.43$	$0.21/-0.36$	$1.41/-0.68$
CCDC no.	2079808	2079809	2079810

Figure S3: View of the three intramolecular π - π interactions in 1-Br.

Magnetic studies

Figure S4: Field dependence of the magnetization, *M*, for **1-Me** below 8 K plotted as (left) *M* vs *H* and (right) *M* vs *H*/*T* plots. The solid blue line is the best fit of the *M* vs *H*/*T* plots to an *S* = 3 Brillouin function (with *g* = 2.20(2)) as discussed in the main text.

Figure S5: Field dependence of the magnetization, *M*, for **1-OMe** below 20 K plotted as (left) *M* vs *H* and (right) *M* vs *H*/*T* plots. The solid blue line is the best fit of the *M* vs *H*/*T* plots to an *S* = 3 Brillouin function (with *g* = 2.17(2)) as discussed in the main text.

Figure S6: Field dependence of the magnetization, *M*, for **1-Br** below 20 K plotted as (left) *M* vs *H* and (right) *M* vs *H*/*T* plots. The solid blue line is the best fit of the *M* vs *H*/*T* plots to an *S* $=$ 3 Brillouin function (with $g = 2.17(4)$) as discussed in the main text.