



HAL
open science

Search for dark matter annihilation signals from unidentified Fermi-LAT objects with H.E.S.S

H. Abdalla, F. Aharonian, F. Ait Benkhali, E.O. Angüner, C. Arcaro, C. Armand, T. Armstrong, H. Ashkar, M. Backes, V. Baghmany, et al.

► **To cite this version:**

H. Abdalla, F. Aharonian, F. Ait Benkhali, E.O. Angüner, C. Arcaro, et al.. Search for dark matter annihilation signals from unidentified Fermi-LAT objects with H.E.S.S. *Astrophys.J.*, 2021, 918, pp.17. <10.3847/1538-4357/abff59>. <hal-03262638>

HAL Id: hal-03262638

<https://hal.science/hal-03262638v1>

Submitted on 6 Aug 2025

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons CC BY 4.0 - Attribution - International License



Search for Dark Matter Annihilation Signals from Unidentified Fermi-LAT Objects with H.E.S.S.

H. Abdalla¹, F. Aharonian^{2,3,4}, F. Ait Benkhali³, E. O. Angüner⁵, C. Arcaro⁶, C. Armand⁷, T. Armstrong⁸, H. Ashkar⁹, M. Backes^{1,6}, V. Baghmanyant¹⁰, V. Barbosa-Martins¹¹, A. Barnacka¹², M. Barnard⁶, Y. Becherini¹³, D. Berge¹¹, K. Bernlöhr³, B. Bi¹⁴, M. Böttcher⁶, C. Boisson¹⁵, J. Bolmont¹⁶, M. de Bony de Lavergne⁷, M. Breuhaus³, R. Brose², F. Brun⁹, T. Bulik¹⁷, T. Bylund¹³, F. Cangemi¹⁶, S. Caroff¹⁶, S. Casanova^{3,10}, P. Chambery¹², J. Catalano¹⁸, T. Chand⁶, A. Chen¹⁹, G. Cotter⁸, M. Curyło¹⁷, H. Dalgleish^{1,6}, J. Damascene Mbarubucye¹¹, I. D. Davids¹, J. Davies⁸, J. Devin²⁰, A. Djannati-Atai²⁰, A. Dmytriiev¹⁵, A. Donath³, V. Doroshenko¹⁴, L. Dreyer⁶, L. du Plessis⁶, C. Duffy²¹, K. Egberts²², S. Einecke²³, G. Emery¹⁶, J.-P. Ernenwein⁵, K. Feijen²³, S. Fegan²⁴, A. Fiasson⁷, G. Fichet de Clairfontaine¹⁵, G. Fontaine²⁴, S. Funk¹⁸, M. Füssling¹¹, S. Gabici²⁰, Y. A. Gallant²⁵, S. Ghafourizade²⁶, G. Giavitto¹¹, L. Giunti^{9,20}, D. Glawion^{18,39}, J. F. Glicenstein⁹, M.-H. Grondin²⁷, S. Hattingh⁶, M. Haupt¹¹, G. Hermann³, J. A. Hinton³, W. Hofmann³, C. Hoischen²², T. L. Holch¹¹, M. Holler²⁸, M. Hörbe⁸, D. Horns²⁹, Z. Huang³, D. Huber²⁸, M. Jamroz¹², D. Jankowsky¹⁸, F. Jankowsky²⁶, V. Joshi¹⁸, I. Jung-Richardt¹⁸, E. Kasai¹, K. Katarzyński³⁰, U. Katz¹⁸, D. Khangulyan³¹, B. Khélifi²⁰, S. Klepser¹¹, W. Kluźniak³², Nu. Komin¹¹, R. Konno⁹, K. Kosack⁹, D. Kostunin¹¹, M. Kreter⁶, G. Kukec Mezek¹³, A. Kundu⁶, G. Lamanna⁷, S. Le Stum⁵, A. Lemièrè²⁰, M. Lemoine-Goumard²⁷, J.-P. Lenain¹⁶, F. Leuschner¹⁴, C. Levy¹⁶, A. Luashvili¹⁵, T. Lohse³³, I. Lypova¹¹, J. Mackey², J. Majumdar¹¹, D. Malyshev^{14,39}, D. Malyshev¹⁸, V. Marandon³, P. Marchegiani¹⁹, A. Marcowith²⁵, A. Mares²⁷, G. Martí-Devesa²⁸, R. Marx^{3,26}, G. Maurin⁷, P. J. Meintjes³⁴, M. Meyer¹⁸, A. Mitchell³, R. Moderski³², L. Mohrmann¹⁸, A. Montanari^{9,39}, C. Moore²¹, P. Morris⁸, E. Moulin^{9,39}, J. Müller²⁴, T. Murach¹¹, K. Nakashima¹⁸, A. Nayerhoda¹⁰, M. de Naurois²⁴, H. Ndiayvala⁶, J. Niemiec¹⁰, A. Noel¹², L. Oberholzer⁶, P. O'Brien²¹, S. Ohm¹¹, L. Olivera-Nieto³, E. de Ona Wilhelmi¹¹, M. Ostrowski¹², M. Panter³, S. Panny²⁸, R. D. Parsons³³, G. Peron³, S. Pita²⁰, V. Poireau⁷, D. A. Prokhorov³⁵, H. Prokoph¹¹, G. Pühlhofer¹⁴, M. Punch^{13,20}, A. Quirrenbach²⁶, P. Reichherzer⁹, A. Reimer²⁸, O. Reimer²⁸, Q. Remy³, M. Renaud²⁵, F. Rieger³, C. Romoli³, G. Rowell²³, B. Rudak³², H. Rueda Ricarte⁹, E. Ruiz-Velasco³, V. Sahakian³⁶, S. Sailer³, H. Salzmann¹⁴, D. A. Sanchez⁷, A. Santangelo¹⁴, M. Sasaki¹⁸, J. Schäfer¹⁸, F. Schüssler⁹, H. M. Schutte⁶, U. Schwanke³³, M. Senniappan¹³, A. S. Seyffert⁶, J. N. S. Shapopi¹, K. Shiningayamwe¹, R. Simoni³⁵, A. Sinha²⁰, H. Spackman⁸, H. Sol¹⁵, A. Specovius¹⁸, S. Spencer⁸, M. Spir-Jacob²⁰, Ł. Stawarz¹², L. Sun³⁵, R. Steenkamp¹, C. Stegmann^{11,22}, S. Steinmassl³, C. Steppa²², T. Takahashi³⁷, T. Tanaka³⁸, T. Tavernier⁹, A. M. Taylor¹¹, R. Terrier²⁰, C. Thorpe-Morgan¹⁴, J. H. E. Thiersen⁶, M. Tluczykont²⁹, L. Tomankova¹⁸, M. Tsirou²⁵, M. Tsujii³¹, R. Tuffs³, Y. Uchiyama³¹, D. J. van der Walt⁶, C. van Eldik¹⁸, C. van Rensburg¹, B. van Soelen³⁴, G. Vasileiadiis²⁵, J. Veh¹⁸, C. Venter⁶, P. Vincent¹⁶, A. Viana³⁸, J. Vink³⁵, H. J. Völk³, S. J. Wagner²⁶, F. Werner³, R. White³, A. Wierzcholska^{10,26}, Yu Wun Wong¹⁸, H. Yassin⁶, A. Yusafzai¹⁸, M. Zacharias^{6,15}, R. Zanin³, D. Zargaryan^{2,4}, A. A. Zdziarski³², A. Zech¹⁵, S. J. Zhu¹¹, A. Zmija¹⁸, J. Zorn³, S. Zouari²¹, and N. Żywucka⁶

(H.E.S.S. Collaboration)

¹ University of Namibia, Department of Physics, Private Bag 13301, Windhoek 10005, Namibia² Dublin Institute for Advanced Studies, 31 Fitzwilliam Place, Dublin 2, Ireland³ Max-Planck-Institut für Kernphysik, P.O. Box 103980, D-69029 Heidelberg, Germany⁴ High Energy Astrophysics Laboratory, RAU, 123 Hovsep Emin St Yerevan 0051, Armenia⁵ Aix Marseille Université, CNRS/IN2P3, CPPM, Marseille, France⁶ Centre for Space Research, North-West University, Potchefstroom 2520, South Africa⁷ Laboratoire d'Annecy de Physique des Particules, Univ. Grenoble Alpes, Univ. Savoie Mont Blanc, CNRS, LAPP, F-74000 Annecy, France⁸ University of Oxford, Department of Physics, Denys Wilkinson Building, Keble Road, Oxford OX1 3RH, UK⁹ IRFU, CEA, Université Paris-Saclay, F-91191 Gif-sur-Yvette, France¹⁰ Instytut Fizyki Jądrowej PAN, ul. Radzikowskiego 152, 31-342 Kraków, Poland¹¹ DESY, D-15738 Zeuthen, Germany¹² Obserwatorium Astronomiczne, Uniwersytet Jagielloński, ul. Orla 171, 30-244 Kraków, Poland¹³ Department of Physics and Electrical Engineering, Linnaeus University, 351 95 Växjö, Sweden¹⁴ Institut für Astronomie und Astrophysik, Universität Tübingen, Sand 1, D-72076 Tübingen, Germany¹⁵ Laboratoire Univers et Théories, Observatoire de Paris, Université PSL, CNRS, Université de Paris, F-92190 Meudon, France¹⁶ Sorbonne Université, Université Paris Diderot, Sorbonne Paris Cité, CNRS/IN2P3, Laboratoire de Physique Nucléaire et de Hautes Energies, LPNHE, 4 Place Jussieu, F-75252 Paris, France¹⁷ Astronomical Observatory, The University of Warsaw, Al. Ujazdowskie 4, 00-478 Warsaw, Poland¹⁸ Friedrich-Alexander-Universität Erlangen-Nürnberg, Erlangen Centre for Astroparticle Physics, Erwin-Rommel-Str. 1, D-91058 Erlangen, Germany¹⁹ School of Physics, University of the Witwatersrand, 1 Jan Smuts Avenue, Braamfontein, Johannesburg, 2050 South Africa²⁰ Université de Paris, CNRS, Astroparticule et Cosmologie, F-75013 Paris, France²¹ Department of Physics and Astronomy, The University of Leicester, University Road, Leicester, LE1 7RH, UK²² Institut für Physik und Astronomie, Universität Potsdam, Karl-Liebknecht-Strasse 24/25, D-14476 Potsdam, Germany²³ School of Physical Sciences, University of Adelaide, Adelaide 5005, Australia²⁴ Laboratoire Leprince-Ringuet, École Polytechnique, CNRS, Institut Polytechnique de Paris, F-91128 Palaiseau, France²⁵ Laboratoire Univers et Particules de Montpellier, Université Montpellier, CNRS/IN2P3, CC 72, Place Eugène Bataillon, F-34095 Montpellier Cedex 5, France²⁶ Landessternwarte, Universität Heidelberg, Königstuhl, D-69117 Heidelberg, Germany

²⁷ Université Bordeaux, CNRS/IN2P3, Centre d'Études Nucléaires de Bordeaux Gradignan, F-33175 Gradignan, France

²⁸ Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität Innsbruck, A-6020 Innsbruck, Austria

²⁹ Universität Hamburg, Institut für Experimentalphysik, Luruper Chaussee 149, D-22761 Hamburg, Germany

³⁰ Institute of Astronomy, Faculty of Physics, Astronomy and Informatics, Nicolaus Copernicus University, Grudziadzka 5, 87-100 Torun, Poland

³¹ Department of Physics, Rikkyo University, 3-34-1 Nishi-Ikebukuro, Toshima-ku, Tokyo 113-0033, 171-8501, Japan

³² Nicolaus Copernicus Astronomical Center, Polish Academy of Sciences, ul. Bartycka 18, 00-716 Warsaw, Poland

³³ Institut für Physik, Humboldt-Universität zu Berlin, Newtonstr. 15, D-12489 Berlin, Germany

³⁴ Department of Physics, University of the Free State, PO Box 339, Bloemfontein 9300, South Africa

³⁵ GRAPPA, Anton Pannekoek Institute for Astronomy, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands

³⁶ Yerevan Physics Institute, 2 Alikhanian Brothers St., 375036 Yerevan, Armenia

³⁷ Kavli Institute for the Physics and Mathematics of the Universe (WPI), The University of Tokyo Institutes for Advanced Study (UTIAS), The University of Tokyo, 5-1-5 Kashiwa-no-Ha, Kashiwa, Chiba, 277-8583, Japan

³⁸ Now at Instituto de Física de São Carlos, Universidade de São Paulo, Av. Trabalhador São-carlense, 400—CEP 13566-590, São Carlos, SP, Brasil

Received 2021 February 22; revised 2021 April 23; accepted 2021 May 8; published 2021 August 30

Abstract

Cosmological N -body simulations show that Milky Way–sized galaxies harbor a population of unmerged dark matter (DM) subhalos. These subhalos could shine in gamma-rays and eventually be detected in gamma-ray surveys as unidentified sources. We performed a thorough selection among unidentified Fermi-Large Area Telescope Objects (UFOs) to identify them as possible tera-electron-volt-scale DM subhalo candidates. We search for very-high-energy ($E \gtrsim 100$ GeV) gamma-ray emissions using H.E.S.S. observations toward four selected UFOs. Since no significant very-high-energy gamma-ray emission is detected in any data set of the four observed UFOs or in the combined UFO data set, strong constraints are derived on the product of the velocity-weighted annihilation cross section $\langle\sigma v\rangle$ by the J factor for the DM models. The 95% confidence level observed upper limits derived from combined H.E.S.S. observations reach $\langle\sigma v\rangle J$ values of 3.7×10^{-5} and 8.1×10^{-6} GeV² cm⁻² s⁻¹ in the W^+W^- and $\tau^+\tau^-$ channels, respectively, for a 1 TeV DM mass. Focusing on thermal weakly interacting massive particles, the H.E.S.S. constraints restrict the J factors to lie in the range 6.1×10^{19} – 2.0×10^{21} GeV² cm⁻⁵ and the masses to lie between 0.2 and 6 TeV in the W^+W^- channel. For the $\tau^+\tau^-$ channel, the J factors lie in the range 7.0×10^{19} – 7.1×10^{20} GeV² cm⁻⁵ and the masses lie between 0.2 and 0.5 TeV. Assuming model-dependent predictions from cosmological N -body simulations on the J -factor distribution for Milky Way–sized galaxies, the DM models with masses >0.3 TeV for the UFO emissions can be ruled out at high confidence level.

Unified Astronomy Thesaurus concepts: [Dark matter \(353\)](#); [High energy astrophysics \(739\)](#); [Gamma-ray sources \(633\)](#); [Gamma-ray telescopes \(634\)](#)

1. Introduction

The presence of dark matter (DM) is suggested by a wealth of astrophysical and cosmological measurements; however, its underlying nature is yet unknown. Among the most promising candidates are weakly interacting massive particles (WIMPs): particles thermally produced in the early universe with mass and coupling strength at the electroweak scale predict a present relic density (Steigman et al. 2012) consistent with that observed today (Adam et al. 2016). WIMP self-annihilation would produce Standard Model particles including gamma-rays, which for a long time have been recognized as a prime messenger to indirectly detect DM annihilation or decay. Gamma-rays are not deflected by magnetic fields and therefore point back to their sources. Among the most promising DM targets observed by ground-based imaging atmospheric Cerenkov telescopes (IACTs) such as H.E.S.S. are the Galactic center (Abdallah et al. 2016, 2018a) and nearby dwarf galaxies (Aharonian et al. 2008a; Abramowski et al. 2011, 2014; Abdalla et al. 2018b).

Other compelling and complementary DM targets for IACTs are DM subhalos populating galactic halos (see, e.g., Kamionkowski et al. 2010). The observed universe today is believed to have formed hierarchically with the smallest structures first. DM particles first collapse into gravitationally bound systems that later merge to form the first subhalos, which subsequently form more massive ones. The merging history leads to DM halos massive enough to retain gas and trigger star formation and give rise to the

galaxies we observe today. However, most of the subhalos remain completely dark. Assuming that DM is made of WIMPs, they could shine in gamma-rays. The annihilation process in subhalos could be frequent enough to be detectable by IACTs provided that the WIMPs are sufficiently massive. If DM subhalos are made of WIMPs, 10^{-4} to $10^{10} M_\odot$ mass subhalos are expected to lie in DM halos of Milky Way– (MW-) sized galaxies (Diemand et al. 2008; Springel et al. 2008b), with the most massive of them ($\gtrsim 10^8 M_\odot$) hosting dwarf galaxies.

DM subhalos are predicted to be compact and concentrated and are not expected to harbor conventional astrophysical high-energy emitters. Provided they are close and/or massive enough, DM annihilations in these objects could produce gamma-ray fluxes detectable with satellite and ground-based experiments such as IACTs (Calcano-Roldan & Moore 2000; Tasitsiomi & Olinto 2002; Stoehr et al. 2003; Koushiappas et al. 2004). However, their actual location in the galaxy is not known. Their search can be performed using all-sky gamma-ray observations (see, for instance, Diemand et al. 2007) such as with the Large Area Telescope (LAT) instrument onboard the Fermi satellite (see, for instance, Berlin & Hooper 2014) or wide-field surveys carried out with IACTs (see, for instance, Aharonian et al. 2008b; Brun et al. 2011).

All-sky Fermi-LAT observations revealed a population of sources that lack associated signals from observations at other wavelengths (Ajello et al. 2017; Abdollahi et al. 2020). These sources are therefore classified as unidentified Fermi objects (UFOs). If the DM particle mass lies below 100 GeV, some UFOs detected by Fermi-LAT could be potentially described

³⁹ Corresponding authors: D. Glawion, D. Malyshev, A. Montanari, E. Moulin; contact.hess@hess-experiment.eu

by DM models (Belikov et al. 2012; Zechlin et al. 2012; Bertoni et al. 2015, 2016; Calore et al. 2017; Coronado-Blazquez et al. 2019b). Identifying some of the UFOs as DM subhalos requires, however, higher ($\gtrsim 100$ GeV) DM masses given their hard gamma-ray spectra in the energy range of a few tens to 100 GeV. Such objects are therefore excellent targets for IACTs to perform searches for tera-electron-volt DM subhalos. In 2018 and 2019, the H.E.S.S. collaboration carried out an observational campaign for a selection of the most promising UFOs in order to probe their potential tera-electron-volt-mass DM-induced emission.

The paper is organized as follows. Section 2 presents the expected DM-induced gamma-ray signals from Galactic subhalos. Section 3 describes the selection procedure of UFOs as DM subhalo candidates relevant for H.E.S.S. observations along with the Fermi-LAT data analysis of the selected UFOs as DM subhalo candidates. In Section 4, the H.E.S.S. observations and data analysis of the selected UFOs are presented. The constraints from H.E.S.S. observations on DM-induced emission models for the UFOs are derived in Section 5. Sections 6 and 7 are devoted to the discussion of the results obtained in this work and to a general summary, respectively.

2. DM Annihilation Signals

2.1. Expected Gamma-Ray Flux from DM Annihilation

The energy-differential gamma-ray flux expected from the self-annihilation of Majorana DM particles of mass m_{DM} can be expressed as

$$\frac{d\Phi_\gamma(E_\gamma, \Delta\Omega)}{dE_\gamma} = \frac{\langle\sigma v\rangle}{8\pi m_{\text{DM}}^2} \sum_f \text{BR}_f \frac{dN^f}{dE_\gamma} J(\Delta\Omega),$$

$$\text{with } J(\Delta\Omega) = \int_{\Delta\Omega} \int_{\text{los}} \rho^2(s(r, \theta)) ds d\Omega. \quad (1)$$

where $\langle\sigma v\rangle$ is the thermally averaged velocity-weighted annihilation cross section and $\sum_f \text{BR}_f dN^f/dE_\gamma$ is the sum of the annihilation spectra dN^f/dE_γ per annihilation in the final states f with associated branching ratios BR_f . The expected DM annihilation signal consists of a continuum spectrum of gamma-rays extending up to the DM mass and possibly a line-like feature close to the DM mass. The former contribution arises from the hadronization and/or decay of quarks, heavy leptons, and gauge bosons involved in the annihilation process. The latter comes from the direct annihilation into γX with $X = \gamma, h, Z$ or a non-Standard Model neutral particle, providing a spectral line at an energy $E_\gamma = m_{\text{DM}}[1 - (m_X/2m_{\text{DM}})^2]$. When the DM particles self-annihilate into charged particles, gamma-rays are produced via processes involving virtual internal bremsstrahlung and final state radiation. These processes provide an additional bump-like feature that peaks at an energy close to the DM mass.

The term $J(\Delta\Omega)$, hereafter referred to the J factor, corresponds to the integration of the square of the DM density over the line of sight s and solid angle $\Delta\Omega$. As opposed to objects with measured stellar dynamics like dwarf galaxies, UFOs have unknown distances to Earth and their J factors cannot be derived from stellar kinematics.

2.2. Expected Subhalo J -factor Distribution in the MW

Cosmological N -body simulations (see, for instance, Diemand et al. 2008; Springel et al. 2008b) predict MW-sized galaxy halos to harbor as of yet unmerged DM substructures called galactic subhalos. These simulations make robust predictions on the slope and normalization of the subhalo mass function defined as $d\ln N/d\ln M \propto M^{-\alpha_m}$, with a slope $\alpha_m \simeq 1.9$ for MW-like galaxies (see, for instance, Diemand et al. 2008; Springel et al. 2008a; Gao et al. 2012; Fiacconi et al. 2017). Galactic subhalos are not expected to host conventional high-energy astrophysical sources, and a large number density of subhalos in MW-like galaxies with high DM concentrations is expected. Only close by and massive enough subhalos are expected to provide detectable gamma-ray signals.

The cosmological simulations provide the abundance of the resolved subhalos, their radial distribution, and structural properties such as their mass and concentration and suggest a DM distribution in the subhalos following a cuspy density profile that can be well described by Navarro–Frenk–White (NFW) (Navarro et al. 1997) or Einasto (Springel et al. 2008c) parameterizations. The limited spatial resolution of the current simulations can only probe the slope of the radial DM density distribution in a subhalo for the most massive subhalos. Assuming a parameterization of the DM density distribution in subhalos, the distribution of J factors of the galactic subhalo population, dN/dJ , can be derived. The cumulative J -factor distribution, $N(J) \equiv N(\geq J)$, is defined as the number of subhalos with a J factor higher than or equal to specified.

In order to derive the J -factor distribution of DM subhalos in the MW, the CLUMPY code v3.0.0 (Charbonnier et al. 2012; Bonnivard et al. 2016; Hütten et al. 2019a) is used. One thousand simulations of an MW-like galaxy are performed with a smooth NFW (Navarro et al. 1997) DM main halo profile with the parameters corresponding to the best-fit NFW parameters presented in a recent study of DM distribution in the MW (Cautun et al. 2020). For each simulation, the subhalo parameters were chosen similar to the ones used in Hütten et al. (2016) for the ‘‘HIGH’’ model. The power-law slope of the subhalo mass function is chosen to be $\alpha_m = 1.9$ (Diemand et al. 2008); the number of objects between 10^8 and $10^{10} M_\odot$ is taken as $N_{\text{calib}} = 300$ following Springel et al. (2008b); and the subhalo mass-concentration relation follows the distance-dependent prescription of Molin e et al. (2017). From each simulation, the Galactic coordinates of all subhalos and their J factors integrated in circular regions with 0.1° radius around the centers of gravity of the subhalos are derived.

The cumulative J -factor distribution $N(\geq J)$ is shown in the upper panel of Figure 1 for subhalos located at Galactic latitudes $|b| > 5^\circ$. The solid red curve shows the averaged distribution computed from all of the realizations, and the shaded region shows the formal 1σ statistical dispersion calculated over all simulated MW-like galaxies. In the lower panel of Figure 1, the red dashed curve illustrates the probability of finding in any simulation at least one subhalo with a J factor higher than specified. The blue dotted curve corresponds to the probability of finding three or more subhalos. The horizontal black dashed line gives the 5% probability: the probability on average of finding one or more subhalos with $J \geq 3 \times 10^{20} \text{ GeV}^2 \text{ cm}^{-5}$ (three or more subhalos with J factors $J \geq 1 \times 10^{20} \text{ GeV}^2 \text{ cm}^{-5}$) is only about 5%.

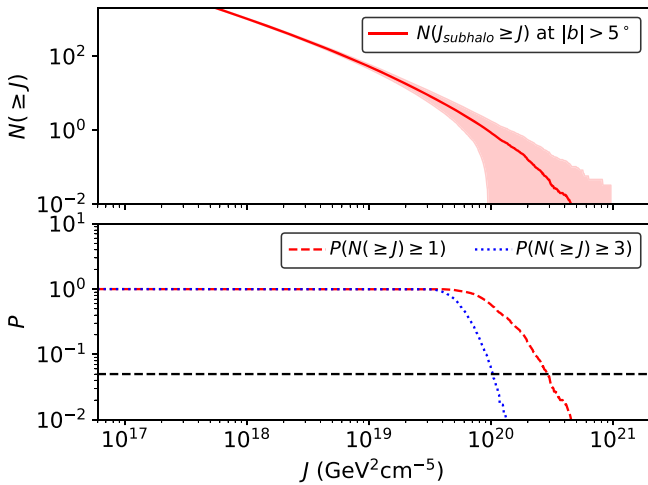


Figure 1. Top panel: Cumulative J -factor distribution, $N(\geq J)$, of an MW-like subhalo population. The number of subhalos with a J factor exceeding a given value is plotted (red solid curve). The red-shaded band corresponds to the 1σ statistical uncertainty obtained from 1000 simulations of the subhalo population for MW-like galaxies. Bottom panel: Probability P of finding at least one subhalo with a J factor higher than specified (red dashed line). The blue dotted line presents the same probability for at least three subhalos. The horizontal black dashed line shows the 5% probability. See text for more details.

3. Fermi-LAT Unidentified Sources as DM Subhalo Candidates

A hypothetical gamma-ray emission from a DM subhalo would show up in the all-sky gamma-ray surveys (Kamionkowski et al. 2010) as unidentified objects, that is, detected with the Fermi satellite but without counterparts at any other wavelengths. A smoking-gun signature for DM detection is a very distinct energy cutoff close to the DM particle mass, assuming the two-body annihilation process taking place almost at rest. For sufficiently large DM particle masses, that is, above a few hundred gigaelectron-volts, the energy cutoff could be too high in energy to be measurable by the LAT instrument onboard the Fermi satellite within a reasonable observation time. The combination of Fermi-LAT and IACT observations is therefore mandatory for DM subhalo searches with unidentified sources detected by the Fermi-LAT.

3.1. Selection Procedure for H.E.S.S. Observations

In order to determine the best candidates for H.E.S.S. observations among the unidentified Fermi-LAT sources, a thorough selection has been performed in the Third Catalog of High-Energy Fermi-LAT Sources (3FHL) (Ajello et al. 2017), which comprises pointlike sources detected above 10 GeV. The source selection requires the unidentified sources to be steady, that is, to not show flux variability over time according to the 3FHL catalog,⁴⁰ exhibiting a hard power-law spectral index ($\Gamma < 2$), as expected for DM-induced signals for DM masses above 100 GeV with no obvious conventional counterpart at other wavelengths. The multiwavelength (MWL) search for possible counterparts is based on the Fermi-LAT source coordinates for each source, and counterparts are searched in catalogs of MWL facilities (XMM-Newton, ROSAT, SUZAKU, Compton

Gamma Ray Observatory, Chandra, Swift, Wilkinson Microwave Anisotropy Probe [WMAP], RXTE, NuStar, Sloan Digital Sky Survey, Planck, Wide-field Infrared Survey Explorer, Hubble Space Telescope) assuming a searched radius around the source determined by the position uncertainty derived in the 3FHL catalog. The unidentified Fermi-LAT sources are usually quite faint, and only the ones sufficiently far away from the Galactic plane are considered in order to avoid challenging background modeling connected to Galactic plane diffuse emission in the Fermi-LAT energy range. The selected sources do not lie in a complex astrophysical environment, that is, they are relatively isolated with no high-energy gamma-ray emission within about 1° .⁴¹ In addition, a maximum zenith angle of 45° for H.E.S.S. observations is required in order to obtain low-energy thresholds. The selection criteria were applied on the 3FHL source catalog and are summarized in Table 1. The small number of suitable DM subhalo candidates obtained by a straightforward selection confirms that the observation of a selection of UFOs is a viable DM search strategy for targeted observations performed by IACTs. Only four UFOs were eventually observed due to H.E.S.S. observational program scheduling constraints. The characteristics of the UFOs selected for observations with the H.E.S.S. telescopes are summarized in Table 2.

3.2. Fermi-LAT Data Analysis on the Selected Objects

Fermi-LAT data selected for the analysis presented in this paper were collected over a time span of more than 12 yr (2008 August–2020 October). The latest available `fermitools` v.2.0.0 with P8R3_V3 response functions (CLEAN photon class)⁴² are used. The initial stage of the analysis aims at determining the energy spectra of UFO sources. In order to extract the differential energy spectrum for each object, the standard binned likelihood analysis of a 14° radius region around each of the considered objects is performed for a set of eight log-equal energy bins in the range 0.1–1000 GeV. The spectral analysis is based on the fitting of a spatial and spectral model of the sky region around the source of interest to the data. The model of the region includes (1) all sources from the most recent 4FGL-DR2 catalog (Abdollahi et al. 2020) within the 14° radius region around UFO position and (2) components for isotropic and Galactic diffuse emissions given by the standard spatial and spectral templates `iso_P8R3_CLEAN_V2_v1.txt` and `gll_iem_v07.fits`, respectively. For the fitting procedure, the spectral models of these sources are selected according to the 4FGL catalog with all parameters except normalization fixed to the catalog values. In addition, 4FGL sources up to 10° beyond the considered region of interest (ROI) are included in the model, with all of their parameters fixed to the catalog values in order to reduce the bias connected to a possible presence of bright sources outside the considered region and effects connected to the poor point-spread function (PSF) of the LAT at low (~ 0.1 GeV) energies. The UFO spectra are modeled by a pure power-law function with the slope defined from a broad energy range fit.

Following the recommendation of the Fermi-LAT collaboration, our analysis is performed with energy dispersion handling enabled. We additionally checked that the TS maps of the

⁴⁰ While the criterium on the variability provides steady candidates as expected for DM sources, Fermi-LAT photon properties at the highest energies have been checked. None of them could be attributed to flaring of nearby Fermi-LAT sources.

⁴¹ The closest 3FHL source for 3FHL J1915.2-1323 is at $0^\circ 8$, while for the other UFOs, the closest source is at distance higher than 1.7° .

⁴² See [description of Fermi-LAT response functions](#).

Table 1
Criteria Applied to 3FHL Catalog for the Selection of DM Subhalo Candidates

Criteria	Numbers of sources
Without association	178
Far enough from the Galactic plane, cut in Galactic latitude of $ b > 5^\circ$	126
Nonvariable, cut in variability index (No. of Bayesian blocks in var. analysis) equal to 1	125
Maximum zenith angle at H.E.S.S. site of 45°	83
Follow a simple power law with significance for curvature $< 3\sigma$	83
Hard spectrum, cut in spectral index below 2	18
No MWL counterparts	6

Note. For the MWL counterpart search, individual search radii were used ($\sim 2-4'$) based on the uncertainty of the Fermi position quoted in the 3FHL. The following list of MWL facilities was checked: XMM-Newton, ROSAT, SUZAKU, Compton Gamma Ray Observatory, Chandra, Swift, WMAP, RXTE, NuStar, Sloan Digital Sky Survey, Planck, Wide-field Infrared Survey Explorer, Hubble Space Telescope.

Table 2
Properties of the Selected UFOs Together with Their Spectral Parameters

Name	RA	Dec.	TS for $E \geq 10$ GeV	Position uncertainty (arcmin)	Pivot energy (GeV)	Spectral energy distribution at pivot energy (10^{-13} TeV cm^{-2} s^{-1})	Power-law index	$\Delta\chi^2$	E_{cut} (GeV)
	(degrees)	(degrees)							
3FHL J0929.2-4110	142.3345	-41.1833	36	2.4	0.39	0.12 ± 0.01	1.37 ± 0.07	0.15	>33
3FHL J1915.2-1323 [†]	288.8182	-13.3916	23	3.0	62.8	2.1 ± 0.9	1.5 ± 0.4	0.05	>35
3FHL J2030.2-5037	307.5901	-50.6344	40	2.6	6.3	1.9 ± 0.3	1.85 ± 0.1	0.40	>67
3FHL J2104.5+2117 ^{a,b}	316.1226	21.2831	58	2.2	1.56	5.3 ± 0.5	2.22 ± 0.06	0.02	>85

Note. The second and third columns provide the R.A.-decl. coordinates of the UFOs along with their test statistic (TS) values for energies above 10 GeV in the fourth column. The fifth column gives their position uncertainty. Pivot energy, spectral energy distribution at the pivot energy, and best-fit power-law spectral index are given in the sixth, seventh, and eighth columns, respectively. The ninth column provides the $\Delta\chi^2$ value between a pure power law and a power law with exponential cutoff fit to the data. The last column gives the 95% confidence level (C.L.) lower limit on a possible energy cutoff in the energy spectrum. The 3FHL J1915.2-1323 source marked with a dagger ([†]) is detected only above 10 GeV. For this source, the spectral index, pivot energy, differential flux, and $\Delta\chi^2$ value are given for this energy band. For the other sources, these quantities are given for energies higher than 0.1 GeV.

^a The spectral index in the 3FHL catalog is 1.8 (Ajello et al. 2017).

^b 3FHL J2104.5.2117 was recently associated in the 4FGL catalog (Abdollahi et al. 2020) with an active galactic nucleus with a probability of 0.4.

considered regions do not show significant residuals between the data and the model; see Appendix A for more details. We conclude that the considered models well describe the UFO sources' regions.

Table 2 summarizes the analysis results for each UFO. Given the available photon statistics in the Fermi-LAT data set of the selected objects, a power-law spectrum with a (super-) exponential energy cutoff in the tera-electron-volt energy range, as expected from DM-induced emissions (Belikov & Silk 2013), and a pure power-law emission cannot be significantly discriminated; see Table 2 for a corresponding change $\Delta\chi^2$ between these models. The last column of Table 2 summarizes 95% C.L. Fermi-LAT lower limits on the energy cutoff, defined as the energy at which $\Delta\chi^2$ changes by 2.71 between power-law and exponential energy cutoff power-law models.

3.3. DM Models for the Selected Sources

Figure 2 shows DM annihilation models together with Fermi-LAT flux measurements. Model predictions for DM masses of 1 and 10 TeV, respectively, are plotted separately for the W^+W^- and $\tau^+\tau^-$ annihilation channels. Some DM models are able to qualitatively describe the observed gamma-ray flux from the selected UFOs. For instance, the predictions shown for $m_{\text{DM}} = 1$ TeV describe well the Fermi-LAT data for 3FHL J0929.2-4110 in the W^+W^- and $\tau^+\tau^-$ annihilation channels. In contrast, for 3FHL J2030.2-5037, the predictions for $m_{\text{DM}} = 1$ TeV in the $\tau^+\tau^-$ annihilation channel provide a poor

description of the emission, while a better one is obtained in the W^+W^- one.

In order to assess quantitatively the above statements on viable DM-induced emission models, the spectrum of each UFO is explicitly modeled with a DM annihilation-induced spectral template.⁴³ For a given m_{DM} and annihilation channel, the model is characterized only by the overall normalization of the spectra given by $\langle\sigma v\rangle J$. To identify the range of viable parameters for DM annihilation, a scan over a large range of $\langle\sigma v\rangle J$ is performed.

A TS is defined as a difference between best-fit log-likelihood functions for models with no DM emission (\mathcal{L}_0 , “background-only” hypothesis) and the model (\mathcal{L}), which includes the UFO source described by the corresponding parameter $\langle\sigma v\rangle J$: $TS = -2 \log(\mathcal{L}/\mathcal{L}_0)$ (Mattox et al. 1996).⁴⁴ Negative TS values correspond to the detection of the source, that is, adding a source with a corresponding parameter improves the fit compared with the background-only hypothesis.

The left panels of Figure 3 illustrate the results for a single UFO 3FHL J0929.2-4110 for the W^+W^- (top) and $\tau^+\tau^-$ (bottom) annihilation channels, respectively. For each $\langle\sigma v\rangle J$, the color scale shows the TS values. Assuming that the TS follows a χ^2 distribution, a TS equal to -9 , (resp. -25) would

⁴³ Provided within fermitools as `DMFitFunction` based on Jeltema & Profumo (2008).

⁴⁴ The TS value for a source with N -parametric (spectral) model follows a χ^2 distribution with N degrees of freedom in the high statistic limit (Wilks 1938).

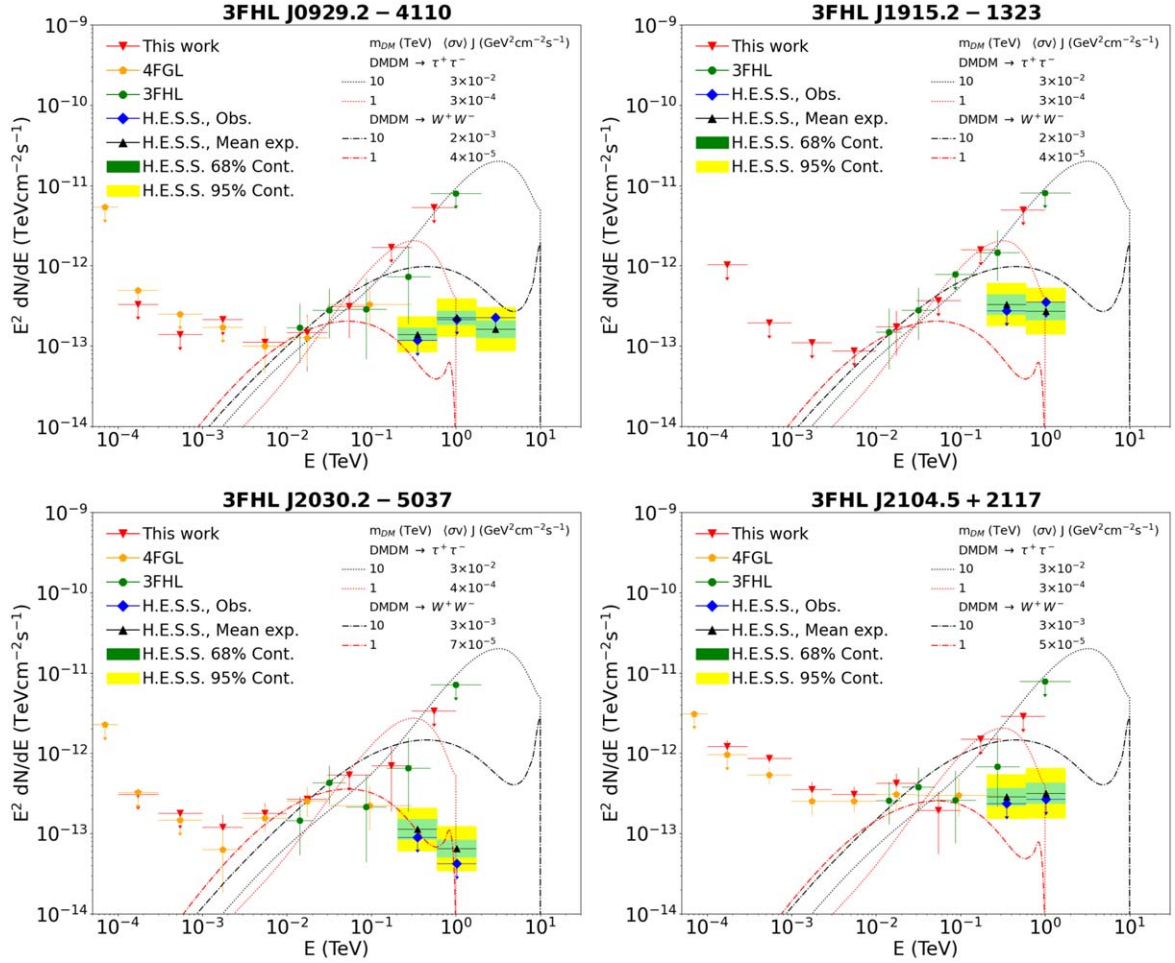


Figure 2. Spectral energy distributions of the selected UFOs observed with Fermi-LAT and H.E.S.S. for 3FHL J0929.2-4110 (top left), 3FHL J1915.2-1323 (top right), 3FHL J2030.2-5037 (bottom left), and 3FHL J2104.5+2117 (bottom right), respectively. The differential flux points computed in this work from the Fermi-LAT data set (red dots) and taken from the 4FGL (orange dots) and from the 3FHL (green dots) catalogs (Ajello et al. 2017; Abdollahi et al. 2020) are shown with the vertical and horizontal error bars corresponding to the 1σ statistical errors and the bin size, respectively. Upper limit points (red, orange, and green arrows) are given at 95% C.L. The observed flux upper limits from H.E.S.S. observations (blue arrows) are plotted at 95% C.L., together with the mean expected flux upper limits (black) and the 1σ (green) and 2σ (yellow) containment bands. Overlaid are theoretical DM-induced fluxes for 1 and 10 TeV DM masses in the W^+W^- (dashed-dotted lines) and $\tau^+\tau^-$ (dotted lines) annihilation channels, respectively.

correspond to a 3σ (resp. 5σ) detection for 1 degree of freedom. The dashed cyan and orange lines show the detection region that corresponds to the improvement of TS by -9 and -25 , respectively. The right panels in Figure 3 present the results for the analysis of the combined data sets of the three selected UFOs (3FHL J0929.2-4110, 3FHL J1915.2-1323, and 3FHL J2030.2-5037) obtained through the combination of the log-likelihood profiles from individual objects, for the W^+W^- (top) and $\tau^+\tau^-$ (bottom) annihilation channels, respectively.

4. H.E.S.S. Observations and Analysis

4.1. Data Analysis

H.E.S.S. is an array of five IACTs located in the Khomas Highland in Namibia, at an altitude of 1800 m. The array is composed of four 12 m diameter telescopes (CT1-4) and a fifth 28 m diameter telescope (CT5) at the middle of the array. The observations presented here were performed in 2018 and 2019 with the full five-telescope array for the selection of unidentified Fermi-LAT objects presented in Table 2. They were carried out in the *wobble* mode, where the telescope pointing direction is offset from the nominal target position by

an angle between $0^\circ.5$ and $0^\circ.7$. The observations for the data analysis are selected according to the standard run selection criteria (Aharonian et al. 2006). After the calibration of raw shower images recorded in the camera, the reconstruction of the direction and energy of the gamma-ray events is performed with a template-fitting technique (de Naurois & Rolland 2009) in which the recorded images are compared to precalculated showers computed from a semianalytical model. An energy resolution of 10% and an angular resolution of $0^\circ.06$ at 68% containment radius for gamma-ray energies above 200 GeV are achieved. The results described here have been cross-checked with an independent calibration and analysis chain yielding compatible results (Parsons & Hinton 2014).

The selected UFOs are assumed to be pointlike sources according to the PSF of Fermi-LAT, which reaches $\sim 0^\circ.1$ above 100 GeV. Given the H.E.S.S. PSF, the ROI, hereafter referred to as the ON source region, is therefore defined as for pointlike-emission searches for H.E.S.S., and the ROI is taken as a disk of $0^\circ.12$ radius. The residual background is measured in OFF regions according to the *MultipleOff* technique (Aharonian et al. 2006). For each telescope pointing position, the OFF regions are defined at the same distance from the pointing position as for the ON

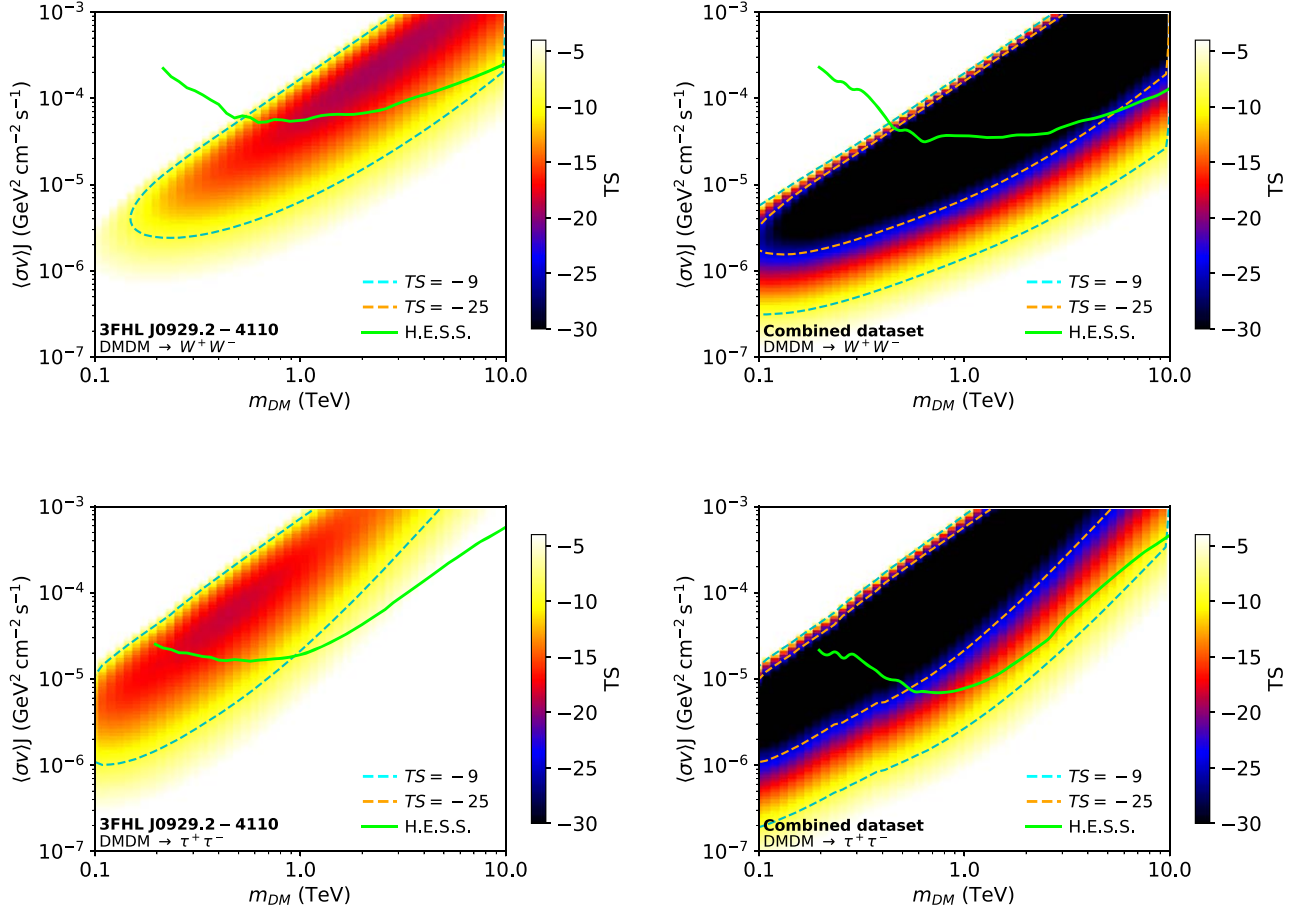


Figure 3. Contours of TS computed from Fermi-LAT data sets on the 3FHL J0929.2-4110 (left panels) and the combined UFO data sets (right panels). The contours are given in the $(\langle\sigma v\rangle, m_{DM})$ plane for the W^+W^- (upper panels) and $\tau^+\tau^-$ (lower panels) annihilation channels. The cyan and orange dashed lines show the -9 and -25 TS contours. Overlaid (solid green line) are H.E.S.S. upper limits displayed at 95% C.L. Contours of TS for 3FHL J1915.2-1323, 3FHL J2030.2-5037, and 3FHL J2104.5+2117 are shown in Figure 7 in Appendix B.

Table 3
H.E.S.S. Data Analysis Results for Each UFO

Name	Live time (hours)	Mean zenith angle (degrees)	N_{ON} (counts)	N_{OFF} (counts)	$\bar{\alpha}$	Significance (σ)
3FHL J0929.2-4110	27.4	29.0	424	5884	13.9	0.1
3FHL J1915.2-1323	3.6	19.4	87	1181	13.9	0.2
3FHL J2030.2-5037	9.8	31.3	160	2192	13.9	0.1
3FHL J2104.5+2117	6.8	46.7	73	853	13.9	1.1

Note. The second and third columns give the live time and mean zenith angle of the H.E.S.S. observations, respectively. Count numbers measured in the ON and OFF regions are given in the fourth and fifth columns, respectively, with the α parameter averaged over all observations, $\bar{\alpha}$, given in the sixth column. The seventh column provides the measured excess significance between the ON and OFF counts.

region, which leads to identical acceptances in the ON and OFF regions. An excluded region defined as a disk of radius equal to twice the ON region radius is used in order to avoid any leakage of the searched signal into the OFF regions. The α parameter is defined as the ratio between the solid angle size of the OFF and ON regions by $\alpha = \Delta\Omega_{OFF}/\Delta\Omega_{ON}$. The excess significance in the ROI is computed following the statistical approach of Li & Ma (1983). Table 3 summarizes for each UFO the live time, the mean zenith angle of the observations, the ON and OFF counts, the α parameter averaged over all the observations, and the excess significance in the ROI. No significant gamma-ray excess is found, neither in the ON source region nor anywhere else in the field of view.

Since no significant excess is found from any of the selected UFOs, energy-differential observed flux upper limits are computed at 95% C.L. assuming the best-fit power-law spectral index derived from the Fermi-LAT data analysis. Expectations for the energy-differential flux upper limits are computed from 100 independent Poisson realizations of the measured OFF counts for the ON and OFF regions. From the realizations, mean and standard deviation values are extracted and used to compute the 68% and 95% containment bands. The differential flux upper limits are shown in Figure 2 with an energy binning of 0.5 dex. The expectations and the containment bands for the upper limits shown in the figure are computed including 25% systematic uncertainty. Assuming a power-law spectral index

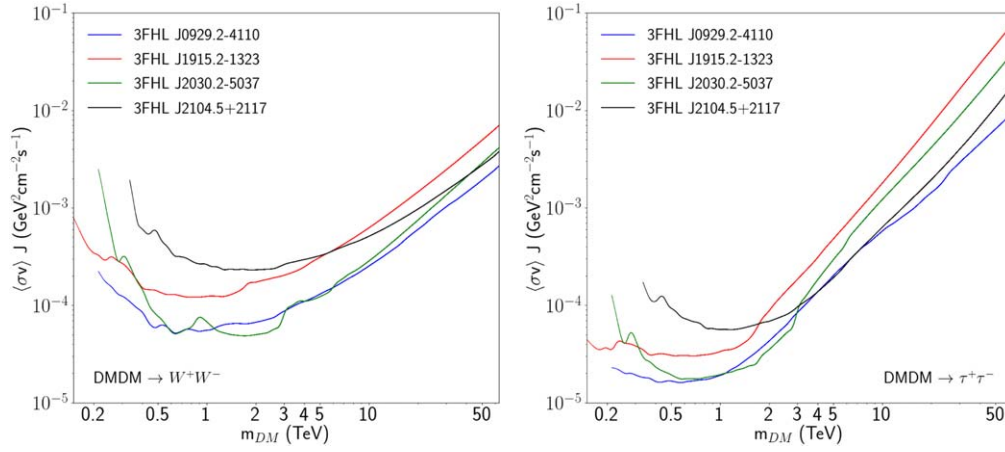


Figure 4. The 95% C.L. upper limits on the product of the annihilation cross section $\langle\sigma v\rangle$, and the J factor as a function of the DM mass m_{DM} in the W^+W^- (left panel) and $\tau^+\tau^-$ (right panel) annihilation channels for 3FHL J0929.2-4110 (blue line), 3FHL J1915.2-1323 (red line), 3FHL J2030.2-5037 (green line), and 3FHL J2104.5+2117 (black line), respectively.

of 2 would change the differential flux upper limits by less than 6%.

4.2. Upper Limit Computation for the DM Emission

A binned Poisson maximum likelihood analysis is performed to search for spectral features expected from DM annihilation signals with respect to the background. For each UFO, the H.E.S.S. energy range is divided into 62 logarithmically spaced bins from 100 GeV up to 70 TeV. For a given DM mass and annihilation channel, the Poisson likelihood function in the energy bin i can be written as

$$\mathcal{L}_i(N_i^S, N_i^B | N_{ON,i}, N_{OFF,i}, \alpha) = \frac{(N_i^S + N_i^B)^{N_{ON,i}} e^{-(N_i^S + N_i^B)}}{N_{ON,i}!} \times \frac{(N_i^{S'} + \alpha N_i^B)^{N_{OFF,i}} e^{-(N_i^{S'} + \alpha N_i^B)}}{N_{OFF,i}!}. \quad (2)$$

The variables $N_{ON,i}$ and $N_{OFF,i}$ stand for the number of measured events in the ON and OFF regions, respectively; N_i^B is the expected number of background events in the ON region; and N_i^S and $N_i^{S'}$ are the expected number of DM signal events in the ON and OFF regions, respectively. They are computed by folding the expected theoretical DM flux given in Equation (1) with the energy-dependent acceptance and energy resolution of H.E.S.S. for the considered data set. The term dN/dE_γ^f in Equation (1) is extracted from Cirelli et al. (2011) for each assumed DM mass and annihilation channel. The energy resolution of H.E.S.S. is represented by a Gaussian function with a width of $\sigma_E/E = 10\%$ above 200 GeV. The UFOs are also pointlike sources for H.E.S.S.; therefore no leakage is expected in the background region and $N_i^{S'}$ is taken to $N_i^{S'} \equiv 0$. The likelihood function for a given object \mathcal{L} is defined as $\mathcal{L} = \prod_i \mathcal{L}_i$.

Since no significant excess is found in any of the selected UFOs by H.E.S.S., upper limits can be derived assuming UFOs are DM-induced gamma-ray emitters from a likelihood ratio TS given by

$$TS = -2 \log \frac{\mathcal{L}(N_i^S(\langle\sigma v\rangle J), \widehat{N}_i^B(\langle\sigma v\rangle J) | N_{ON}, N_{OFF}, \alpha)}{\mathcal{L}(\widehat{N}_i^S(\langle\sigma v\rangle J), \widehat{N}_i^B | N_{ON}, N_{OFF}, \alpha)}. \quad (3)$$

The variable \widehat{N}_i^B is obtained through a conditional maximization, achieved by solving $d\mathcal{L}/dN_i^B = 0$; \widehat{N}_i^S and \widehat{N}_i^B are computed using an unconditional maximization. Following the procedure defined in Cowan et al. (2011), upper limits are computed assuming a positive signal, that is, $\widehat{N}_i^S > 0$. The $\langle\sigma v\rangle J$ value for which the TS value is equal to 2.71 provides the one-sided 95% C.L. upper limit on the quantity $\langle\sigma v\rangle J$.

The hypothesis that all UFOs are indeed DM subhalos, but too faint to be detected as such in the tera-electron-volt energy range with the given exposure, can be tested when the individual data sets are combined. If no significant overall excess is found in the combined data set, combined upper limits on $\langle\sigma v\rangle J$ can be derived versus the DM mass assuming J to be an averaged J factor. In Equation (3), the likelihood function is replaced by the combined likelihood expressed as $\mathcal{L}_{\text{comb}} = \prod_{j=1}^{N_{\text{targets}}} \mathcal{L}_j$, where \mathcal{L}_j is the likelihood of the target j .

5. Results

No significant excess is measured in any of the H.E.S.S. data sets of the selected UFOs. The 95% C.L. upper limits on $\langle\sigma v\rangle J$ are derived versus the DM mass using Equation (3). Figure 4 shows for each UFO the upper limits as a function of the DM mass for the W^+W^- and $\tau^+\tau^-$ annihilation channels, respectively. In most of the DM mass range, the strongest constraints are reached for 3FHL J0929.2-4110 observations. For a 1 TeV DM mass, the constraints $\langle\sigma v\rangle J = 5.5 \times 10^{-5} \text{ GeV}^2 \text{ cm}^{-2} \text{ s}^{-1}$ and $1.9 \times 10^{-5} \text{ GeV}^2 \text{ cm}^{-2} \text{ s}^{-1}$ in the W^+W^- and $\tau^+\tau^-$ annihilation channels, respectively, for 3FHL J0929.2-4110.

The combined analysis of the four H.E.S.S. data sets does not show any significant excess. Therefore, the UFO data sets are combined and upper limits on $\langle\sigma v\rangle J$ are computed. Given its possible association with an active galactic nucleus, the source 3FHL J2104.5+2117 is removed to provide conservative combined upper limits. The right panel of Figure 3 shows

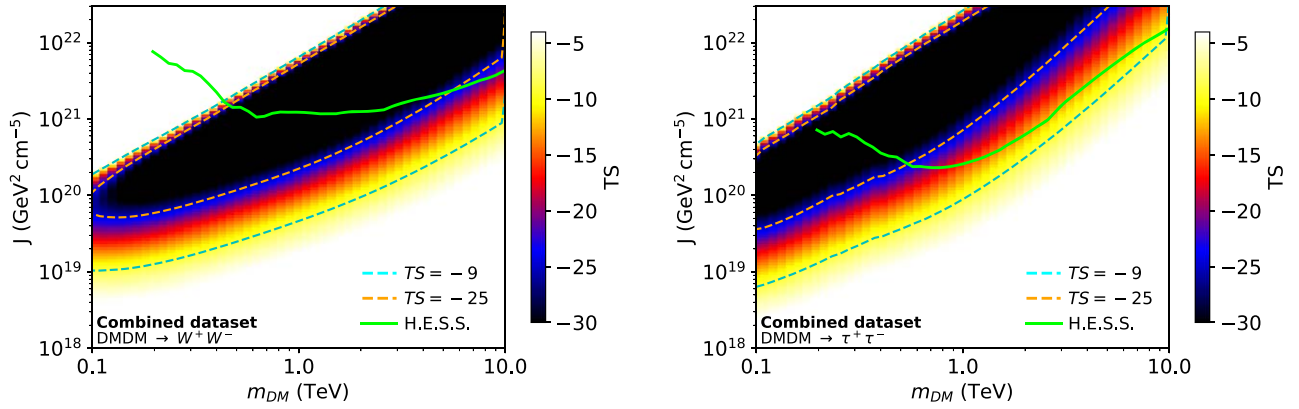


Figure 5. Contours of TS computed from the Fermi-LAT combined UFO data sets. The contours are given in the (J, m_{DM}) plane for the W^+W^- (left panel) and $\tau^+\tau^-$ (right panel) annihilation channel, assuming the $\langle\sigma v\rangle$ value expected for thermal WIMPs. The cyan and orange dashed lines show the -9 and -25 TS contours. Overlaid (solid green line) are the H.E.S.S. 95% C.L. upper limits from the combined UFO data sets.

the combined 95% C.L. upper limits on $\langle\sigma v\rangle J$ as a function of the DM mass for the W^+W^- and $\tau^+\tau^-$ annihilation channels, respectively. The analysis of the combined data sets allows for an improvement of about 10% and 20% for 1 TeV DM mass in the W^+W^- and $\tau^+\tau^-$ annihilation channels, respectively, with respect to the most constraining upper limit from the individual UFO data sets. The combined limits reach 3.7×10^{-5} and 8.1×10^{-6} $\text{GeV}^2 \text{cm}^{-2} \text{s}^{-1}$ in the W^+W^- and $\tau^+\tau^-$ channels, respectively, for a 1 TeV DM mass.

In order to derive the J -factor values required to explain the UFO emission in terms of DM models, the value of the annihilation cross section expected for thermal WIMPs ($\langle\sigma v\rangle_{\text{th}} \simeq 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$) is used (Steigman et al. 2012). The J -factor upper limits for the DM models of the UFOs as a function of the DM mass are given at 95% C.L. in Figure 5. For a 1 TeV DM mass in the W^+W^- channel, the J -factor values are constrained to be between 2.4×10^{20} and 1.3×10^{21} $\text{GeV}^2 \text{cm}^{-5}$ for DM models with $\text{TS} \leq -25$ (which corresponds to $\geq 5\sigma$ confidence interval assuming TS follows χ^2 distribution). For a DM mass of 10 TeV in the W^+W^- channel, all the J -factor values for DM models with $\text{TS} \leq -25$ are ruled out at 95% C.L. by the H.E.S.S. constraints. In the $\tau^+\tau^-$ channels, the H.E.S.S. constraints are even stronger. For 300 GeV DM mass, the allowed J -factor values are between 1.4×10^{20} and 5.9×10^{20} $\text{GeV}^2 \text{cm}^{-5}$ for $\text{TS} \leq -25$ DM models. The H.E.S.S. upper limits restrict the J factors to lie in the range 6.1×10^{19} – 2.0×10^{21} $\text{GeV}^2 \text{cm}^{-5}$ and the masses to lie between 0.2 and 6 TeV in the W^+W^- channel. For the $\tau^+\tau^-$ channel, the J factors lie in the range 7.0×10^{19} – 7.1×10^{20} $\text{GeV}^2 \text{cm}^{-5}$ and the masses lie between 0.2 and 0.5 TeV.

Using predictions of N -body cosmological simulations, the number of subhalos with a J factor higher than a given value for an MW-like galaxy can be extracted as displayed in Figure 1. The probability of having at least three subhalos with a J factor higher than 10^{20} $\text{GeV}^2 \text{cm}^{-5}$ is below 5% (Figure 1, blue dotted line). According to this prediction, the interpretation of the UFO emissions in terms of DM particle annihilations in Galactic DM subhalos can be further constrained from Figure 5 to $m_{\text{DM}} \lesssim 1$ TeV for W^+W^- and $m_{\text{DM}} \lesssim 0.3$ TeV for $\tau^+\tau^-$ channels.

6. Discussion

A substantial number of UFOs may emit gamma-rays from DM annihilations in subhalos. However, some of them could

be active galactic nuclei or other types of galaxies that so far lack detection at other wavelengths. Alternative astrophysical interpretations of UFOs as pulsars or low-luminosity globular clusters hosting millisecond pulsars (Mirabal et al. 2016) may be less plausible since typical gamma-ray spectra for these types of objects are characterized by an energy cutoff at energies of a few giga-electron-volts.

The cumulative J -factor distribution is in very good agreement with the results of Hütten et al. (2016) for the “HIGH” model intended to predict the highest possible number of subhalos in a typical MW-like galaxy. The real number of DM subhalos can be an order of magnitude smaller as shown for the predictions in the “LOW” model of Hütten et al. (2016). The choice of the number of subhalos of masses between 10^8 and $10^{10} M_{\odot}$ of $N_{\text{calib}} = 300$ is motivated by the output of DM-only simulations Springel et al. (2008b). Baryon feedback can significantly reduce this value, up to a factor of two (Mollitor et al. 2015; Sawala et al. 2016). This would make the highest J -factor values even more unlikely. As discussed in Coronado-Blázquez et al. (2019a), subhalos with the highest J factors should appear as extended sources for Fermi-LAT given its PSF of about 0.1° above 10 GeV. However, even these brightest DM subhalos would produce faint gamma-ray sources whose spatial extension would be challenging to measure for Fermi-LAT. On the simulation front, further work is likely needed to use predictions for subhalo angular sizes in MW-like galaxies to definitely rule out pointlike UFOs as potential DM subhalos.

The DM density distribution of the Galactic halo is assumed here to follow an NFW parameterization. Incorporating hydrodynamics and baryon feedback in cosmological simulations tends to soften the inner cusp of the DM profiles in MW-like galaxies, leading to a flattening of order 1 kpc (see, for instance, Chan et al. 2015). However, these predictions for the expected DM distribution have large uncertainties due to the effects of baryonic physics. The resolution limit of the simulations at sufficiently small distances also becomes pertinent. Alternative Galactic mass models can be used to describe subhalo parameters for MW-like galaxies (Catena & Ullio 2010; McMillan 2011; Stref & Laval 2017; McMillan 2017). Using different Galactic mass models, the subhalo luminosity functions derived in Stref & Laval (2017) provide compatible results. Considering a core profile would make the high DM mass exclusion of the DM models for the UFO emission even stronger. DM subhalos could also have cored profiles with lower DM concentration, which would make them more prone to tidal disruption. This would therefore both

decrease the normalization of the J -factor distribution function and shift it to lower values. The J -factor distribution prediction is based on DM-only Via Lactea-II simulations using WMAP cosmology. The simulations based on the most recent cosmological results from the Planck mission with added baryonic physics could somewhat change the predicted properties of the MW subhalos. The effect of baryon feedback and tidal effects induced by both DM and baryons is likely to alter the DM concentration of subhalos (Despali & Vegetti 2017; Stref & Lavalle 2017). Details of the modeling of tidal disruption of Galactic DM subhalos on the brightest subhalo can be found in Di Mauro et al. (2020). Therefore, the cumulative J -factor distribution would be shifted to lower values. In this case, the DM-induced interpretation of the UFO emission would be even more constrained given that the probability of finding high J -factor values would be even smaller. The presence of baryons affects both the amplitude and the slope of the DM halo and subhalo mass functions. It can reduce the slope by a few percent in the 10^6 – $10^9 M_\odot$ subhalo mass range (Benson 2020), which can also alter the rate of subhalos with large J factors. No cut is applied to the maximal value of the subhalo mass when computing the cumulative J -factor distribution. While subhalos with masses of about $10^7 M_\odot$ or higher may be able to host star formation and actually be dwarf galaxies, simulations including hydrodynamics and feedback physics in addition to the gravitational effects for the expected DM distribution in both the main halo and its subhalos such as in Zhu et al. (2016) show that a significant fraction of subhalos with masses of about $10^9 M_\odot$ is found to host no stars. For subhalo masses larger than 10^7 – $10^8 M_\odot$, a noticeable fraction of them can start triggering star formation and indeed form a faint dwarf galaxy. Arguably these masses critically depend on the baryonic physics implementation in the simulations and its associated feedback. Considering subhalos with masses lower than $10^7 M_\odot$ implies the probability of having at least one subhalo with $J \geq 3 \times 10^{20} \text{ GeV}^2 \text{ cm}^{-5}$ to be about 0.3%, therefore a factor of about 16 lower than in the case without mass cut.

The interpretation of UFOs as DM subhalos of tera-electron-volt-mass scale thermal WIMPs requires J factors in excess of a few $10^{20} \text{ GeV}^2 \text{ cm}^{-5}$. Such J -factor values are only occasionally obtained in N -body simulations of MW-type galaxies. The highest subhalo J -factor values are usually subject to a large statistical variance. The precise value of the brightest subhalos can be subject to a large uncertainty, implying a factor of 10 uncertainty on the J -factor value for $J \gtrsim 10^{20} \text{ GeV}^2 \text{ cm}^{-5}$ in the “HIGH” model (Hütten et al. 2019b).⁴⁵ Previous studies using Fermi-LAT data only searched for UFOs as promising DM subhalo candidates for DM masses below 100 GeV (see, for instance, Bertoni et al. 2015, 2016; Coronado-Blázquez et al. 2019b). For instance, in the study of Coronado-Blázquez et al. (2019a), masses of a few tens of giga-electron-volts are ruled out for canonical thermal WIMPs. The present analysis searches for UFOs as DM subhalos for the uncharted tera-electron-volt-mass thermal WIMP models.

The maximal value of subhalo J factors obtained in simulations is model dependent and can be increased even in comparison with the optimistic estimate of the J -factor

distribution considered here as discussed below. The normalization of the subhalo mass function assumes usually about 10% of the total DM halo content to be in the form of subhalos, with the total DM halo density being normalized to the DM density at the location of the Sun $\rho(r_\odot) = \rho_\odot = 0.39 \text{ GeV cm}^{-3}$. The precise value is subject to uncertainties of a factor of 1.5–2 (Read 2014). Varying the input parameters of the simulations in the relevant ranges of interest, such as $\rho_\odot = 0.6 \text{ GeV cm}^{-3}$ and the scale radius of the main DM halo $r_s = 25 \text{ kpc}$, would increase the highest J -factor values by a factor of two. The predicted cumulative J -factor distribution would therefore probe higher J -factor values. Considering substructures in galactic subhalos (see, for instance, Hiroshima et al. 2018) results in higher expected J -factor values for the Galactic subhalo population with typical increase factors of a few, therefore shifting the cumulative J -factor distribution to higher values.

The abovementioned large systematic uncertainties in the prediction of the J -factor distribution weaken significantly the constraints from cosmological simulations, which makes them comparable to or weaker than the H.E.S.S. constraints in, for example, the $\tau^+\tau^-$ channel. This makes the model-independent H.E.S.S. constraints the only relevant ones for robust interpretation of UFOs as Galactic subhalos of annihilating DM.

7. Summary

In this work, a straightforward filtering of the unidentified sources in the 3FHL point-source catalog has been performed using selection cuts to identify the most promising DM subhalo candidates for DM masses above a few hundred giga-electronvolts. The data sets for the four UFOs were collected with the Fermi satellite in a 12 yr observation period. Using H.E.S.S. observations, no significant signal has been detected from any of the selected UFOs. From H.E.S.S. flux upper limits, DM models describing the UFO emissions with high significance are strongly constrained in the tera-electron-volt DM mass range for different annihilation channels. Assuming thermally produced DM particles, the DM models for the UFO emissions require high J -factor values. When the model-dependent predictions from N -body simulations of the MW-like subhalo population are taken into account, the required high J -factor values for the DM models explaining the UFOs as Galactic subhalos are unlikely. This can point toward interpretation of the UFOs as subhalos of relatively light WIMPs with masses $m_{\text{DM}} \lesssim 0.3 \text{ TeV}$, for which somewhat lower J factors are preferred. However, this could be in tension with constraints on thermal WIMPs from dwarf galaxy observations by Fermi-LAT (Ackermann et al. 2015; Albert et al. 2017).

The support of the Namibian authorities and of the University of Namibia in facilitating the construction and operation of H.E.S.S. is gratefully acknowledged, as is the support by the German Ministry for Education and Research (BMBF), the Max Planck Society, the German Research Foundation (DFG), the Helmholtz Association, the Alexander von Humboldt Foundation, the French Ministry of Higher Education, Research and Innovation, the Centre National de la Recherche Scientifique (CNRS/IN2P3 and CNRS/INSU), the Commissariat à l’énergie atomique et aux énergies alternatives (CEA), the U.K. Science and Technology Facilities Council (STFC), the Knut and Alice Wallenberg Foundation, the National Science Centre, Poland grant No. 2016/22/M/ST9/00382, the

⁴⁵ Using the “LOW” model for the predictions of the J -factor distribution would make the interpretation of UFOs as DM subhalos even more unlikely given that the probability to get high J -factor values would be lowered compared with what it would be in the “HIGH” model.

South African Department of Science and Technology and National Research Foundation, the University of Namibia, the National Commission on Research, Science & Technology of Namibia (NCRST), the Austrian Federal Ministry of Education, Science and Research and the Austrian Science Fund (FWF), the Australian Research Council (ARC), the Japan Society for the Promotion of Science, and the University of Amsterdam.

We appreciate the excellent work of the technical support staff in Berlin, Zeuthen, Heidelberg, Palaiseau, Paris, Saclay, Tübingen, and Namibia in the construction and operation of the equipment. This work benefitted from services provided by the H.E.S.S. Virtual Organisation, supported by the national resource providers of the EGI Federation.

Appendix A

TS Maps for the UFO Sources from FERMI-LAT Data Sets

In this section we present the results for the search of possible sources not accounted for by the considered model based on the 4FGL catalog of the UFOs vicinity. We built TS maps for $5^\circ \times 5^\circ$ region centered at the position of each UFO. These maps illustrate the significance ($\sim\sqrt{TS}$) above the background model of a test point-like source with a power-law spectrum characterized by a free normalization and slope fixed to -2 , in each pixel.

For the background model, where the corresponding UFO source is removed, we consider the same spatial/spectral models as used for the analysis of Fermi-LAT data described in

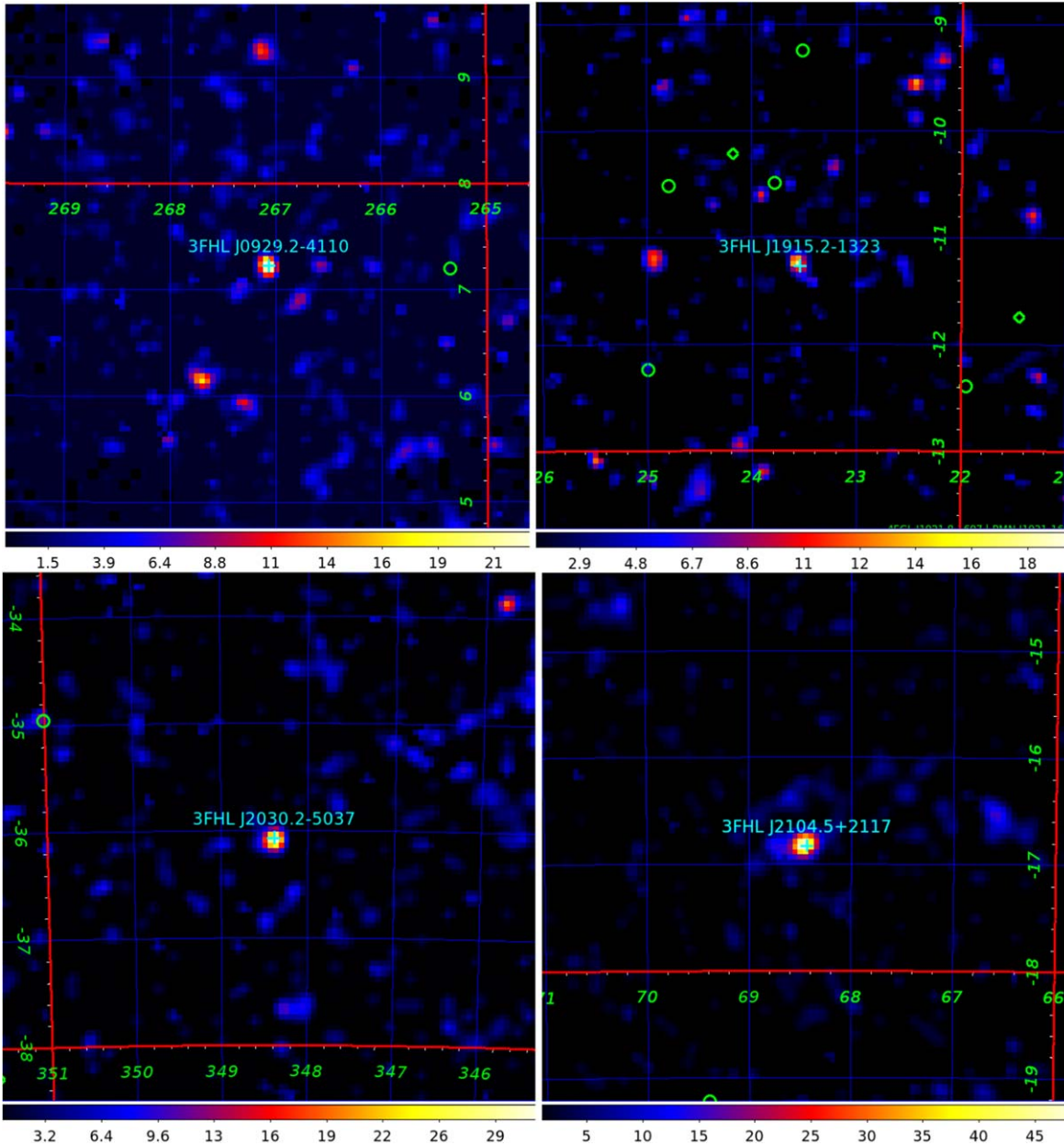


Figure 6. TS maps for the $5^\circ \times 5^\circ$ region around each of the considered UFOs for energies above 10 GeV in Galactic coordinates with pixel size of 0.05° . The color scale indicates the values of the TS. No smoothing to the map is applied. Cyan crosses show the position of the corresponding UFO source. Green crosses show positions of nearby 4FGL sources included in the background model.

Section 3.2. Such a choice of the background model allows us to check the presence of unaccounted for sources close to UFOs' positions as well as verify the point-like spatial morphology of the emission associated with the UFO.

Figure 6 show the TS maps for all four UFO sources for energies above 10 GeV in Galactic coordinates with a pixel size of $0^{\circ}05$. No smoothing is applied to the maps. Cyan crosses show the positions of UFO sources. Green circles indicate the position of the nearby 4FGL sources included in the background model. The maps indicate the absence of significant residuals, which could affect the Fermi-LAT analysis of the UFOs.

Appendix B

Contours of TS from Fermi-LAT data sets on the 3FHL J1915.2-1323, 3FHL J2030.2-5037, and 3FHL J2104.5+2117 for the W^+W^- and $\tau^+\tau^-$ annihilation channels, respectively

In Figure 7 we show the contour of TS from Fermi-LAT datasets and the 95% C.L. upper limits on J as a function of the DM mass for 3FHL J1915.2-1323, 3FHL J2030.2-5037 and 3FHL J2104.5+2117. Contours and upper limits for the W^+W^- and $\tau^+\tau^-$ annihilation channels are shown in the left and right panels, respectively.

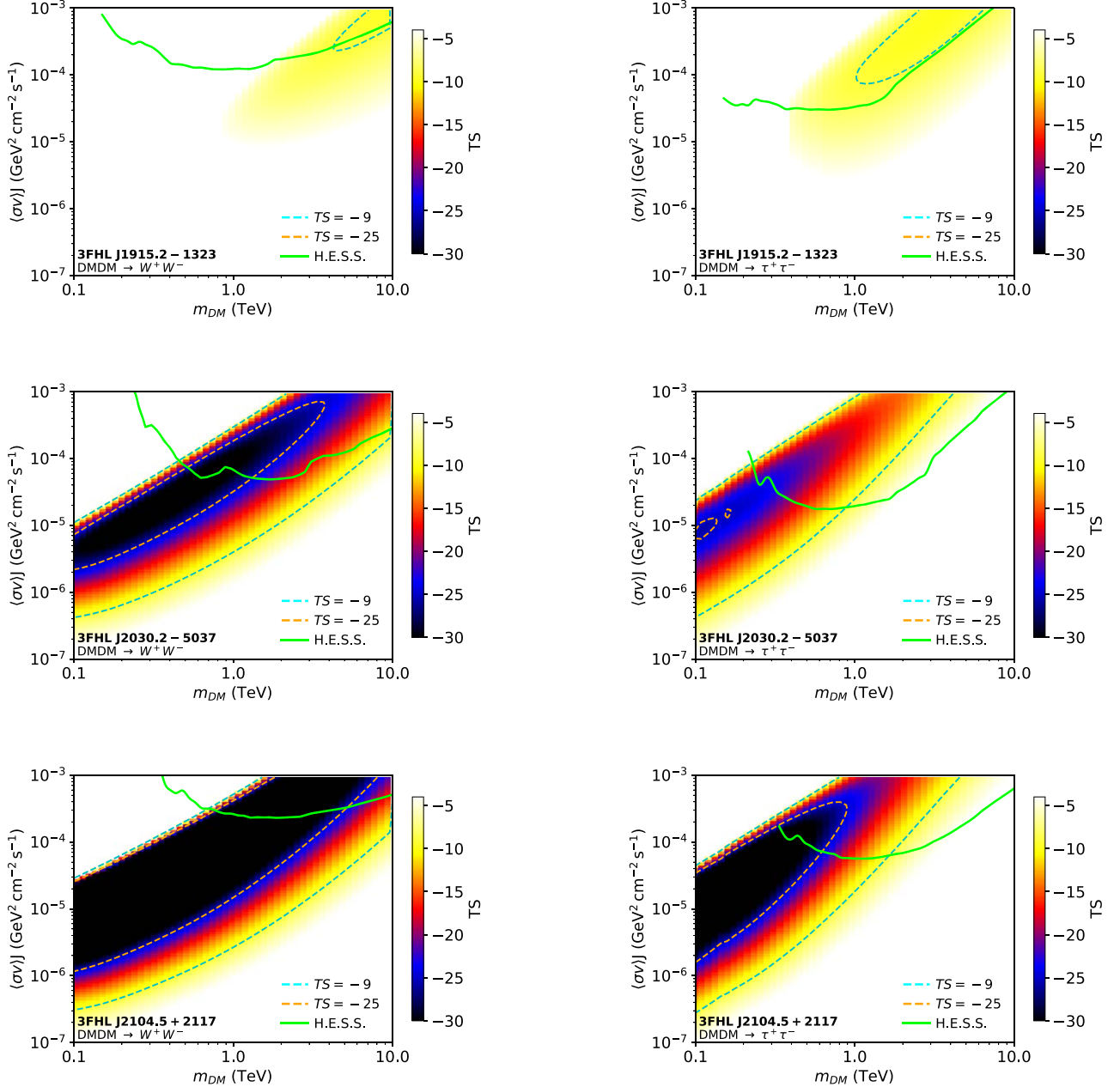


Figure 7. Contours of TS computed from Fermi-LAT data sets on 3FHL J1915.2-1323 (upper panels), 3FHL J2030.2-5037 (middle panels), and 3FHL J2104.5+2117 (lower panels). The contours are given in the $(\langle\sigma v\rangle, m_{\text{DM}})$ plane for the W^+W^- (left panels) and $\tau^+\tau^-$ (right panels) annihilation channels. The cyan and orange dashed lines show the -9 and -25 TS contours. Overlaid (solid green line) are the H.E.S.S. constraints displayed at 95% C.L.

ORCID iDs

- F. Aharonian <https://orcid.org/0000-0003-1157-3915>
 T. Armstrong <https://orcid.org/0000-0001-5067-2620>
 H. Ashkar <https://orcid.org/0000-0002-2153-1818>
 M. Backes <https://orcid.org/0000-0002-9326-6400>
 V. Baghmanyant <https://orcid.org/0000-0003-0477-1614>
 V. Barbosa-Martins <https://orcid.org/0000-0002-5085-8828>
 A. Barnacka <https://orcid.org/0000-0001-5655-4158>
 M. Barnard <https://orcid.org/0000-0003-1720-7959>
 D. Berge <https://orcid.org/0000-0002-2918-1824>
 K. Bernlöhr <https://orcid.org/0000-0001-8065-3252>
 M. Böttcher <https://orcid.org/0000-0002-8434-5692>
 C. Boisson <https://orcid.org/0000-0001-5893-1797>
 J. Bolmont <https://orcid.org/0000-0003-4739-8389>
 M. Breuhaus <https://orcid.org/0000-0003-0268-5122>
 R. Brose <https://orcid.org/0000-0002-8312-6930>
 F. Brun <https://orcid.org/0000-0003-0770-9007>
 T. Bylund <https://orcid.org/0000-0003-2946-1313>
 S. Caroff <https://orcid.org/0000-0002-1103-130X>
 S. Casanova <https://orcid.org/0000-0002-6144-9122>
 A. Chen <https://orcid.org/0000-0001-6425-5692>
 G. Cotter <https://orcid.org/0000-0002-9975-1829>
 J. Damascene Mbarubucyeye <https://orcid.org/0000-0002-4991-6576>
 J. Davies <https://orcid.org/0000-0002-2394-4720>
 L. Dreyer <https://orcid.org/0000-0002-4971-3672>
 G. Fichet de Clairfontaine <https://orcid.org/0000-0003-1143-3883>
 G. Fontaine <https://orcid.org/0000-0002-2357-1012>
 S. Funk <https://orcid.org/0000-0002-2012-0080>
 G. Giavitto <https://orcid.org/0000-0002-7629-6499>
 D. Glawion <https://orcid.org/0000-0003-4865-7696>
 J. F. Glicenstein <https://orcid.org/0000-0003-2581-1742>
 M.-H. Grondin <https://orcid.org/0000-0002-8383-251X>
 T. L. Holch <https://orcid.org/0000-0003-4865-7696>
 M. Jamrozy <https://orcid.org/0000-0002-0870-7778>
 V. Joshi <https://orcid.org/0000-0003-4467-3621>
 U. Katz <https://orcid.org/0000-0002-7063-4418>
 D. Khangulyan <https://orcid.org/0000-0002-7576-7869>
 B. Khélifi <https://orcid.org/0000-0001-6876-5577>
 Nu. Komin <https://orcid.org/0000-0003-3280-0582>
 R. Konno <https://orcid.org/0000-0003-1892-2356>
 D. Kostunin <https://orcid.org/0000-0002-0487-0076>
 M. Kreter <https://orcid.org/0000-0002-3092-3506>
 A. Kundu <https://orcid.org/0000-0003-2128-1414>
 M. Lemoine-Goumard <https://orcid.org/0000-0002-4462-3686>
 J.-P. Lenain <https://orcid.org/0000-0001-7284-9220>
 F. Leuschner <https://orcid.org/0000-0001-9037-0272>
 J. Mackey <https://orcid.org/0000-0002-5449-6131>
 D. Malyshev <https://orcid.org/0000-0001-9689-2194>
 D. Malyshev <https://orcid.org/0000-0002-9102-4854>
 V. Marandon <https://orcid.org/0000-0001-9077-4058>
 A. Mares <https://orcid.org/0000-0003-3345-7800>
 G. Martí-Devesa <https://orcid.org/0000-0003-0766-6473>
 R. Marx <https://orcid.org/0000-0002-6557-4924>
 A. Mitchell <https://orcid.org/0000-0003-3631-5648>
 R. Moderski <https://orcid.org/0000-0002-8663-3882>
 L. Mohrmann <https://orcid.org/0000-0002-9667-8654>
 A. Montanari <https://orcid.org/0000-0002-3620-0173>
 P. Morris <https://orcid.org/0000-0002-8533-8232>
 E. Moulin <https://orcid.org/0000-0003-4007-0145>
 J. Muller <https://orcid.org/0000-0003-0004-4110>
 T. Murach <https://orcid.org/0000-0003-1128-5008>
 H. Ndiyavala <https://orcid.org/0000-0001-9279-1775>
 J. Niemiec <https://orcid.org/0000-0001-6036-8569>
 P. O'Brien <https://orcid.org/0000-0002-5128-1899>
 S. Ohm <https://orcid.org/0000-0002-3474-2243>
 L. Olivera-Nieto <https://orcid.org/0000-0002-9105-0518>
 M. Ostrowski <https://orcid.org/0000-0002-9199-7031>
 S. Panny <https://orcid.org/0000-0001-5770-3805>
 R. D. Parsons <https://orcid.org/0000-0003-3457-9308>
 G. Peron <https://orcid.org/0000-0003-3255-0077>
 V. Poireau <https://orcid.org/0000-0002-4768-0256>
 M. Punch <https://orcid.org/0000-0002-4710-2165>
 P. Reichherzer <https://orcid.org/0000-0003-4513-8241>
 A. Reimer <https://orcid.org/0000-0001-8604-7077>
 O. Reimer <https://orcid.org/0000-0001-6953-1385>
 F. Rieger <https://orcid.org/0000-0003-1334-2993>
 C. Romoli <https://orcid.org/0000-0003-2541-4499>
 G. Rowell <https://orcid.org/0000-0002-9516-1581>
 B. Rudak <https://orcid.org/0000-0003-0452-3805>
 H. Rueda Ricarte <https://orcid.org/0000-0001-9833-7637>
 E. Ruiz-Velasco <https://orcid.org/0000-0001-6939-7825>
 A. Santangelo <https://orcid.org/0000-0003-4187-9560>
 M. Sasaki <https://orcid.org/0000-0001-5302-1866>
 F. Schüssler <https://orcid.org/0000-0003-1500-6571>
 H. M. Schutte <https://orcid.org/0000-0002-1769-5617>
 U. Schwanke <https://orcid.org/0000-0002-1229-278X>
 M. Senniappan <https://orcid.org/0000-0001-6734-7699>
 J. N. S. Shapopi <https://orcid.org/0000-0002-7130-9270>
 S. Spencer <https://orcid.org/0000-0001-5516-1205>
 S. Steinmassl <https://orcid.org/0000-0002-2865-8563>
 A. M. Taylor <https://orcid.org/0000-0001-9473-4785>
 R. Terrier <https://orcid.org/0000-0002-8219-4667>
 C. van Eldik <https://orcid.org/0000-0001-9669-645X>
 C. Venter <https://orcid.org/0000-0002-2666-4812>
 H. J. Völk <https://orcid.org/0000-0003-2386-8067>
 S. J. Wagner <https://orcid.org/0000-0002-7474-6062>
 A. Wierzycholska <https://orcid.org/0000-0003-4472-7204>
 M. Zacharias <https://orcid.org/0000-0001-5801-3945>
 D. Zargaryan <https://orcid.org/0000-0002-2876-6433>
 A. A. Zdziarski <https://orcid.org/0000-0002-0333-2452>
 S. J. Zhu <https://orcid.org/0000-0002-6468-8292>
 S. Zouari <https://orcid.org/0000-0002-5333-2004>
 N. Żywucka <https://orcid.org/0000-0003-2644-6441>

References

- Abdallah, H., Abramowski, A., Aharonian, F., et al. (H.E.S.S. Collaboration) 2016, *PhRvL*, **117**, 111301
 Abdallah, H., Abramowski, A., Aharonian, F., et al. (H.E.S.S. Collaboration) 2018a, *PhRvL*, **120**, 201101 H
 Abdalla, H., Aharonian, F., Benkhali, F., et al. (H.E.S.S. Collaboration) 2018b, *JCAP*, **11**, 037
 Abdollahi, S., Acero, F., Ackermann, M., et al. (Fermi-LAT Collaboration) 2020, *ApJS*, **247**, 33
 Abramowski, A., Acero, F., Aharonian, F., et al. (H.E.S.S. Collaboration) 2011, *Aph*, **34**, 608
 Abramowski, A., Aharonian, F., Benkhali, F., et al. (H.E.S.S. Collaboration) 2014, *PhRvD*, **90**, 112012
 Ackermann, M., Albert, A., Anderson, B., et al. (Fermi-LAT Collaboration) 2015, *PhRvL*, **115**, 231301
 Adam, R., Ade, P. A. R., Aghanim, N., et al. (Planck Collaboration) 2016, *A&A*, **594**, A1
 Aharonian, F., Akhperjanian, A. G., Bazer-Bachi, A. R., et al. (H.E.S.S. Collaboration) 2006, *A&A*, **457**, 899

- Aharonian, F., Akhperjanian, A. G., Bazer-Bachi, A. R., et al. (H.E.S.S. Collaboration) 2008a, *Aph*, **29**, 55
- Aharonian, F., Akhperjanian, A. G., Barres de Almeida, U., et al. (H.E.S.S. Collaboration) 2008b, *PhRvD*, **78**, 072008
- Ajello, M., Atwood, W. B., Baldini, L., et al. (Fermi-LAT Collaboration) 2017, *ApJS*, **232**, 18
- Albert, A., Anderson, B., Bechtol, K., et al. (Fermi-LAT and DES Collaborations) 2017, *ApJ*, **834**, 110
- Belikov, A. V., Hooper, D., & Buckley, M. R. 2012, *PhRvD*, **86**, 043504
- Belikov, A. V., & Silk, J. 2013, *PhRvL*, **111**, 071302
- Benson, A. J. 2020, *MNRAS*, **493**, 1268
- Berlin, A., & Hooper, D. 2014, *PhRvD*, **89**, 016014
- Bertoni, B., Hooper, D., & Linden, T. 2015, *JCAP*, **12**, 035
- Bertoni, B., Hooper, D., & Linden, T. 2016, *JCAP*, **05**, 049
- Bonnivard, V., Hütten, M., Nezri, E., et al. 2016, *CoPhC*, **200**, 336
- Brun, P., Moulin, E., Diemand, J., & Glicenstein, J.-F. 2011, *PhRvD*, **83**, 015003
- Calcano-Roldan, C., & Moore, B. 2000, *PhRvD*, **62**, 123005
- Calore, F., De Romeri, V., Di Mauro, M., Donato, F., & Marinacci, F. 2017, *PhRvD*, **96**, 063009
- Catena, R., & Ullio, P. 2010, *JCAP*, **08**, 004
- Cautun, M., Benítez-Llambay, A., Deason, A. J., et al. 2020, *MNRAS*, **494**, 4291
- Chan, T. K., Kereš, D., Oñorbe, J., et al. 2015, *MNRAS*, **454**, 2981
- Charbonnier, A., Combet, C., & Maurin, D. 2012, *CoPhC*, **183**, 656
- Cirelli, M., Corcella, G., Hektor, A., et al. 2011, *JCAP*, **1103**, 051
- Coronado-Blázquez, J., Sánchez-Conde, M. A., Di Mauro, M., et al. 2019a, *JCAP*, **11**, 045
- Coronado-Blázquez, J., Sánchez-Conde, M. A., Domínguez, A., et al. 2019b, *JCAP*, **07**, 020
- Cowan, G., Cranmer, K., Gross, E., & Vitells, O. 2011, *EPJC*, **71**, 1554
- de Naurois, M., & Rolland, L. 2009, *Aph*, **32**, 231
- Despali, G., & Vegetti, S. 2017, *MNRAS*, **469**, 1997
- Di Mauro, M., Stref, M., & Calore, F. 2020, *PhRvD*, **102**, 103010
- Diemand, J., Kuhlen, M., & Madau, P. 2007, *ApJ*, **657**, 262
- Diemand, J., Kuhlen, M., Madau, P., et al. 2008, *Natur*, **454**, 735
- Fiacconi, D., Mayer, L., Madau, P., et al. 2017, *MNRAS*, **467**, 4080
- Gao, L., Navarro, J., Frenk, C., et al. 2012, *MNRAS*, **425**, 2169
- Hiroshima, N., Ando, S., & Ishiyama, T. 2018, *PhRvD*, **97**, 123002
- Hütten, M., Combet, C., Maier, G., & Maurin, D. 2016, *JCAP*, **2016**, 047
- Hütten, M., Combet, C., & Maurin, D. 2019a, *CoPhC*, **235**, 336
- Hütten, M., Stref, M., Combet, C., Lavalle, J., & Maurin, D. 2019b, *Galax*, **7**, 60
- Jeltema, T. E., & Profumo, S. 2008, *JCAP*, **2008**, 003
- Kamionkowski, M., Koushiappas, S. M., & Kuhlen, M. 2010, *PhRvD*, **81**, 043532
- Koushiappas, S. M., Zentner, A. R., & Walker, T. P. 2004, *PhRvD*, **69**, 043501
- Li, T. P., & Ma, Y. Q. 1983, *ApJ*, **272**, 317
- Mattox, J. R., Bertsch, D. L., Chiang, J., et al. 1996, *ApJ*, **461**, 396
- McMillan, P. J. 2011, *MNRAS*, **414**, 2446
- McMillan, P. J. 2017, *MNRAS*, **465**, 76
- Mirabal, N., Charles, E., Ferrara, E., et al. 2016, *ApJ*, **825**, 69
- Moliné, Á., Sánchez-Conde, M. A., Palomares-Ruiz, S., & Prada, F. 2017, *MNRAS*, **466**, 4974
- Mollitor, P., Nezri, E., & Teyssier, R. 2015, *MNRAS*, **447**, 1353
- Navarro, J. F., Frenk, C. S., & White, S. D. M. 1997, *ApJ*, **490**, 493
- Parsons, R. D., & Hinton, J. A. 2014, *Aph*, **56**, 26
- Read, J. 2014, *JPhG*, **41**, 063101
- Sawala, T., Frenk, C. S., Fattahi, A., et al. 2016, *MNRAS*, **457**, 1931
- Springel, V., Wang, J., Vogelsberger, M., et al. 2008b, *MNRAS*, **391**, 1685
- Springel, V., White, S. D. M., Frenk, C. S., et al. 2008a, *Natur*, **456**, 73
- Springel, V., White, S. D. M., Frenk, C. S., et al. 2008c, *Natur*, **456**, 73
- Steigman, G., Dasgupta, B., & Beacom, J. F. 2012, *PhRvD*, **86**, 023506
- Stoehr, F., White, S. D. M., Springel, V., Tormen, G., & Yoshida, N. 2003, *MNRAS*, **345**, 1313
- Stref, M., & Lavalle, J. 2017, *PhRvD*, **95**, 063003
- Tasitsiomi, A., & Olinto, A. V. 2002, *PhRvD*, **66**, 083006
- Wilks, S. S. 1938, *Annals Math. Statist.*, **9**, 60
- Zechlin, H. S., Fernandes, M. V., Elsaesser, D., & Horns, D. 2012, *A&A*, **538**, A93
- Zhu, Q., Marinacci, F., Maji, M., et al. 2016, *MNRAS*, **458**, 1559