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Rotational spectroscopy of imidazole: Accurate spectroscopic information for three vibrationally excited states and the heavy-atom isotopologues up to 295 GHz

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Abstract

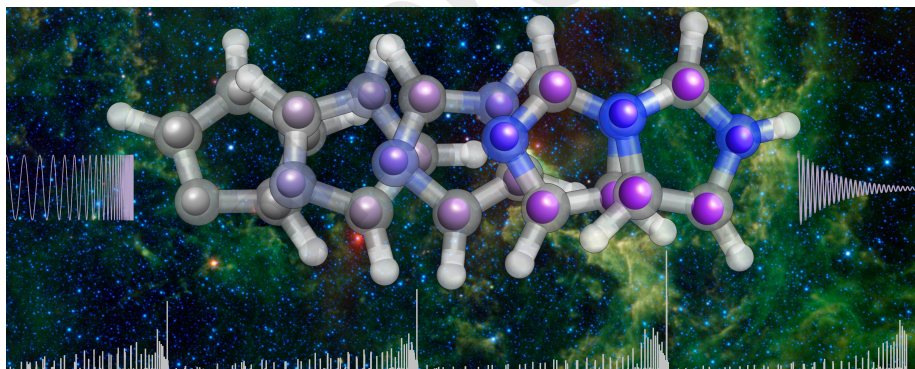
We report our analysis of the pure rotational spectra of low-lying vibrationally excited states and heavy-atom rare isotopologues of imidazole. To facilitate searches for imidazole in the interstellar medium, we previously described the analysis of the rotational spectrum of the imidazole main isotopologue across the 2-295 GHz range, and we extend this analysis here. Structure optimisation and anharmonic frequency calculations were performed to aid the spectral analysis. Three vibrationally excited states of imidazole were assigned in our room-temperature spectra, with energies up to approximately 670 cm^{-1} above the vibronic ground state. The vibrational states could act as temperature probes in warmer star-forming regions. The ^{13}C and ^{15}N heavy-atom isotopologues were assigned, which allowed for the structure of imidazole in the gas phase to be determined. Structural comparisons are drawn between the related heterocyclic molecules hydantoin and imidazolidine. An experimental gas-phase structure of the former is also determined in this work. The potential detection of the isotopologues could help deduce formation pathways towards imidazole and other nitrogen-containing cyclic compounds in interstellar space.

Keywords

astrochemistry, imidazole, isotopologues, molecular structure, rotational spectroscopy, vibrationally excited states

Highlights

- Imidazole was studied between 2-295 GHz; hydantoin was studied between 18-26 GHz.
- Three vibrationally excited states of imidazole were observed and assigned.
- Isotopologue analysis allowed gas-phase structure determination for both molecules.
- Imidazole, hydantoin, and imidazolidine structures and vibrations are compared.
- Precise rotational constants and line lists are delivered for all assigned species.

Graphical Abstract

1. Introduction

Activities to fully characterise the chemical inventory of the interstellar medium (ISM) are key drivers for laboratory astrochemistry, with over 200 molecules having been identified in the ISM so far [1]; an up-to-date list can be found at the Cologne Database for Molecular Spectroscopy¹. In particular, the increase in sensitivity of observational facilities such as the Atacama Large Millimeter/submillimeter Array (ALMA) is facilitating the detection of more complex molecules. The availability of accurate rest frequencies and precise rotational constants for astrochemically relevant molecules, such as complex organic molecules (COMs), must follow these advancements in observational astronomy. These experimentally determined data sets are necessary for searches for these molecules in the ISM, and we present here such data that will allow searches for vibrationally excited states and isotopologues of imidazole (*c*-NHCHN(CH)₂) to be performed in observational spectra.

Imidazole is an interesting molecule for astrochemical studies due to the confirmed presence of N-heterocycles on meteorites (see Ref. [2] and references therein), suggesting conditions exist in space for its formation in substantial abundance. N-heterocycles are also prevalent in biologically important motifs. For example, imidazole is present in the purine bases of DNA and RNA, as well as in the structures of the amino acid histidine and the neurotransmitter histamine. Despite a number of cyclic compounds already being observed in the ISM, imidazole has yet to be detected. Indeed, following on from an unsuccessful search for imidazole by the authors of Ref. [3], our analysis of the ground state rotational spectrum was accompanied with an observational search towards the molecular core Sagittarius B2(N2) with the EMOCA line survey [4], which resulted in a non-detection [5]. We had also previously suggested a new search towards TMC-1, where benzonitrile was detected in 2018 [6]. The presence of 1-cyano-1,3-cyclopentadiene [7] and 1- and 2-cyanonaphthalene [8] towards the same source was recently reported. These detections, as concluded in Ref. [7], hint towards a reservoir of ring molecules in TMC-1 that defies the current chemical models, and this could prove promising in the search for imidazole.

Vibrationally excited states and isotopologues are important considerations when analysing observational data sets [9]. Indeed, as noted in Ref. [10], these species are believed to be the source of many unidentified lines. The percentage of unidentified lines can be as much as 70% in very deep observations at ALMA Bands 3 (84-116 GHz) and 6 (211-275 GHz) [11]. Further, the 2018 NASA Laboratory Astrophysics Workshop identified “spectral studies producing rotational line lists, especially focusing on vibrational states, [and] isotopologues as one of the highest-priority recommendations [12]. The experimentally determined rotational constants and line lists, like those we report here, facilitate searches for these species in the ISM.

Vibrationally excited states are important as indicators of the physical conditions of the regions they are detected in (see, for example, Ref. [13] for a comparison between vibrational and local thermal equilibrium temperatures), and they can thus aid in the characterisation of, for example, hot core regions. The lines generated by vibrational states of already-detected molecules can also be considered as ‘weeds’ [14] or a hindrance to identifying new COMs in observational data sets. Their elimination from

¹<https://cdms.astro.uni-koeln.de/cdms/porta1/>

astronomical spectra is an important step in processing complex, confusion-limited spectra. Knowledge of these lines can draw attention to cases of blended transitions and prevent misassignments, particularly in broadband, unbiased observational campaigns.

From an astrochemical perspective, the search for and detection of the ^{13}C and ^{15}N isotopologues in regions where imidazole is detected could be of importance for elucidating formation pathways of the parent species. Furuya *et al.* demonstrated with a gas-grain chemical network model that isotopic ratios can be used to deduce likely routes to interstellar molecules. In their study, they showed that the $^{12}\text{C}/^{13}\text{C}$ ratio for CCH and CCS differed from the elemental $^{12}\text{C}/^{13}\text{C}$ ratio, and the ratio was dependent on whether the molecule was formed from the carbon ion or from carbon monoxide [15]. One could postulate that constructing similar models with the nitrogen atom/ion or other small nitrogen-containing species (some examples are discussed below) could allow for the deduction of likely formation pathways for imidazole. Any future detection of imidazole isotopologues, made possible by the high-level line lists provided in this study, would provide the interstellar abundances and $^{14}\text{N}/^{15}\text{N}$ ratios required for such an analysis to be performed.

Isotopologues are also important species from a spectroscopic and structural point of view. With the experimentally determined rotational constants of the singly-substituted isotopologues, a substitution structure using Kraitchman's equations can be determined [16]. Blackman *et al.* presented an analysis of the rotational spectrum of imidazole and N-deutero imidazole, which focussed on the characterisation of the quadrupole hyperfine structure [17]. Building on this, and their previous study in Ref. [18], Christen *et al.* provided a set of rotational constants for the molecule's parent species, as well as the D, ^{13}C , and ^{15}N isotopologues determined from data up to 40 GHz [19]. This allowed a comprehensive structure of imidazole in the gas phase to be constructed, which is re-evaluated here based on our extended data set.

Imidazole was previously studied with an astrochemical motivation by Schwell *et al.* [20], where the authors studied the fragmentation patterns of imidazole with photoionisation mass spectrometry. The survival of imidazole, and the other nitrogen-containing cyclic molecules studied, was discussed. In summary, the authors of Ref. [20] conclude that protection from vacuum ultraviolet (VUV) radiation is required for detectable quantities of the studied N-bearing molecules to survive in the ISM. The molecules would likely undergo dissociative ionisation if irradiated with VUV radiation. It was therefore suggested that searches for imidazole be done in hot molecular cores and in young circumstellar disks, where such shielded conditions are likely to exist. If a search for the imidazole ground state in such an environment would be successful, knowledge of the rotational parameters for vibrational states would be important for further searches.

A number of previous studies have looked at the incorporation of smaller molecules, such as HCN, into larger ring structures. For example, acetylene polymerisation with HCN in the gas-phase has been found to produce nitrogen-containing polycyclic aromatic structures [21, 22, 23], which could have interesting implications for interstellar imidazole synthesis. A further example of this was included in Ref. [24], where a number of imidazole-like structures were presented in an interstellar synthetic pathway towards adenine. The pathways start from familiar small molecules - HCN, HCCN,

NH_2CN , and CN^- - and one could imagine termination of the reaction at any of the relevant steps could result in imidazole being formed.

In this paper, we present the analysis of the rotational spectra of the imidazole vibrationally excited states and heavy-atom isotopologues up to 295 GHz. The resulting rotational constants, quadrupole coupling constants, and high order distortion constants have allowed us to accurately predict the rotational line lists for these species in the millimeter-wave regime and beyond. This data is essential for observational searches. Further, the motions of the assigned vibrational states were revealed from the results of anharmonic frequency calculations. The isotopologue analysis allowed us to characterise the gas-phase structure of imidazole's ring motif and compare its structural parameters to related molecules, including hydantoin ($c\text{-NHC(O)NHCH}_2\text{C(O)}$), whose heavy-atom isotopologues were also experimentally examined in this work.

2. Experimental Details and Theoretical Methods

The experimental work was conducted across a number of high-resolution broadband chirped-pulse Fourier transform rotational spectrometers at the Deutsches Elektronen-Synchrotron (DESY, Germany) and at the University of Virginia (USA). The work covers the frequency ranges 2-8, 12-15.5, 18-26, 75-110, and 260-295 GHz. The experiments reported on here were performed for the analysis of the imidazole main isotopologue as presented in Ref. [5], where there was no analysis presented on the vibrationally excited states nor the isotopologues. The methodology has therefore been described in detail previously, and a short summary follows here for information.

The Hamburg COMPACT spectrometer was used to record the spectrum of imidazole across the 2-8 GHz and 12-15.5 GHz regions [25, 26, 27]. For these measurements, solid imidazole (purchased from Sigma Aldrich, $\geq 99\%$ purity, and used without further treatment) was placed in a reservoir nozzle, which was heated to 90 °C. Neon, at a pressure of 3 bar, was used as a backing gas, and the imidazole vapour was supersonically expanded into the vacuum chamber. The molecules then interacted with the incident radiation, which is created by an arbitrary waveform generator (AWG) and amplified by a 300 W travelling wavetube amplifier (for the 2-8 GHz experiment) or a 50 W solid state amplifier (for the 12-15.5 GHz experiment) before being transmitted into the chamber. The spectrometer uses the fast frame method, whereby each molecular pulse is probed by eight chirps spanning the frequency regime of the experiment. The relaxation of the induced macroscopic dipole moment is collected in the form of a free induction decay (FID). For the 2-8 GHz experiment, 1.6 million FIDs (each recorded for 40 μs) were co-added and fast Fourier transformed into the frequency domain. For the 12-15.5 GHz experiment, 1.2 million FIDs (each recorded for 20 μs) were co-added and fast Fourier transformed.

The spectrum across the 18-26 GHz region was recorded using a segmented chirped-pulse Fourier transform microwave spectrometer [27]. The spectrometer combines the segmented approach and the multi-train method. Here, the segmented approach involves splitting the spectrum into 800 MHz segments, and the spectrum across the whole bandwidth is obtained by concatenating the segments together [28]. The multi-train method involves utilising multiple pulse trains per molecular pulse [27], which is similar to the fast frame approach exploited in the Hamburg COMPACT spectrometer. The 1.5 μs excitation pulses are written by an AWG, and they cover the region

7-3 GHz. The pulses are then frequency up-converted with a number of mixing and doubling stages. The FIDs were collected, frequency down-converted, then recorded for 10 μ s. Segment concatenation yielded the spectrum across the whole region, which consisted of 1.8 million co-added FIDs. Full details of the design and characterisation of the instrument are given in Ref. [27]. The imidazole sample was placed in a heatable reservoir nozzle similar to that used in the experiments performed on the COMPACT. Three bar backing pressure of neon was used as the carrier gas. The rotational spectrum of hydantoin (also purchased from Sigma Aldrich, >99% purity, and used without further treatment) was also recorded in the 18-26 GHz region. This was performed in extension to the work in Ref. [29] and in order to obtain an experimental gas-phase structure that could be used for structural comparisons between imidazole and hydantoin (see Section 3). For the hydantoin experiment, the sample was placed in an internal reservoir nozzle, heated to 185 °C, and supersonically expanded into the vacuum chamber with 3 bar of neon as the backing gas. The spectrum consists of 3.5 million co-added FIDs. The frequency accuracy achieved in both the COMPACT and the 18-26 GHz spectrometer was in the order of 25 kHz.

The spectrum in the 75-110 GHz region was recorded in a room temperature flow cell set-up with a BrightSpec, Inc. W-band spectrometer. This spectrometer also utilises the segmented approach, and the 35 GHz of bandwidth is covered in segments of 30 MHz, as implemented in the instrument's High Dynamic Range (HDR) mode [30]. The excitation pulses are created by an AWG in the 1.5-2.3 GHz range, and the millimeter-wave radiation is achieved with an up-conversion stage and an active multiplier chain (AMC, x6). The excitation pulses were 500 ns in duration, and after collection and frequency down-conversion, the FIDs were recorded for 4 μ s. A total of 500,000 FIDs were co-added in the W-band experiment. The imidazole sample was heated to 90 °C in an external reservoir, and the vapours were pulsed into the sample chamber without a backing gas.

The 260-295 GHz spectrum was recorded at the University of Virginia. The instrument has been previously described in detail [28, 31, 32]. The operation of the 260-295 GHz spectrometer also follows the segmented operating regime. The AWG produces a pulse train in the region 2-4 GHz, which is then frequency up-converted by mixing and an AMC to the 260-295 GHz range. The spectrum is split into fifty segments of 720 MHz bandwidth, where each 720 MHz is made up of thirty 24 MHz chirps in a similar fashion to the HDR mode of the W-band spectrometer. The pulse duration was 100 ns, and the FIDs, after collection and frequency down-conversion, were recorded for 1.9 μ s. The experiment was performed with imidazole crystals, which were sublimed, and the vapours were allowed to enter the sample chamber, where a pressure of 7 μ bar was maintained in a slow-flow configuration. The number of co-added FIDs in this case was 10,000, which equates to approximately three minutes of measurement time. The frequency accuracy achieved in both the W-band spectrometer and the 260-295 GHz spectrometer was about 35 kHz.

The assignment of the imidazole parent species was presented and discussed in our previous work [5]. Using the same experimental data, we were able to assign the rotational spectra of the singly-substituted heavy isotopologues in natural abundance. This was achieved by combining and refining the accurately determined nitrogen (^{14}N) quadrupole coupling constants and high order centrifugal distortion constants from Ref.

[5] with the A , B , and C rotational constants determined in Ref. [19]. A global fit was performed for the individual isotopologues, which involved fitting the rotational constants and the centrifugal distortion constants to the assigned line frequencies across all of the frequency regimes studied in this work. Further, the nuclear quadrupole coupling constants were fit only to transitions which exhibited hyperfine splitting (HFS) - in this work, this included the spectra up to 26 GHz. This global fit procedure was also used in the analysis of the rotational spectrum of the parent imidazole isotopologue [5].

The spectra were initially fit using the PGOPHER programme [33]. Global fits were obtained using Pickett's SPFIT/SPCAT suite of programmes [34]. The data were fit to Watson's S -reduction rotational Hamiltonian in the III' representation with centrifugal distortion corrections and hyperfine splitting considerations. The substitution structure analysis was performed using the KRA programme [35], and the EVAL programme [35] was used to determine structural parameters and their associated errors; the effective structure analysis was performed using the STRFIT programme [36].

Extensive *ab initio* calculations have been performed to derive theoretical predictions for many spectroscopic and structural parameters of imidazole and its heavy-atom isotopic variants. The computational treatment included complete basis set extrapolation (CBS) by using several calculations carried out at both the Hartree-Fock and the CCSD(T) levels of theory in conjunction with basis sets up to cc-pV6Z and a composite approach for modelling core-valence and full electron correlation contributions. Anharmonicity corrections to the vibrational fundamentals and to vibrationally-averaged rotational constants have been computed at the MP2/cc-pVQZ level (hybrid approach). With regards to the calculation of the sextic centrifugal distortion constants, the fc-CCSD level of theory in conjunction with the cc-pVTZ basis set has been chosen due to its good balance between accuracy and computational time (see, for example, Ref [37]) with respect to the more demanding fc-CCSD(T). All of these calculations have been carried out using the CFOUR suite of programmes². Full details can be found in the preceding publication and the references therein [5].

Structure optimisation calculations were performed at the MP2/aug-cc-pVTZ level of theory as implemented in Gaussian09 [38] for hydantoin and imidazolidine. The hydantoin calculations supported the assignment of the isotopologues and the subsequent gas-phase structure determination. The imidazolidine calculations supplied theoretical parameters for the structural comparisons discussed later.

3. Results and Discussions

3.1 Vibrationally Excited States

The experiments in the 75-110 GHz and 260-295 GHz ranges were performed in room temperature flow cells, and thus the rotational spectra of low-lying vibrationally excited states could be observed and assigned. We present here the assignment of three of these vibrationally excited states: ν_{21} , ν_{20} , and ν_{19} . The assignment process was supported by

²CFOUR, Coupled-Cluster techniques for Computational Chemistry, a quantum-chemical programme package written by J.F. Stanton, J. Gauss, L. Cheng, M.E. Harding, D.A. Matthews, P.G. Szalay et al., and the integral packages MOLECULE (J. Almlöf and P.R. Taylor), PROPS (P.R. Taylor), ABACUS (T. Helgaker, H.J. Aa.Jensen, P. Jrgensen, and J. Olsen), and ECP routines by A. V.Mitin and C. van Wullen. For the current version, see <http://www.cfour.de>.

theoretical calculations of the corresponding fundamental anharmonic frequencies and vibrational contributions to the rotational constants (see Section 2). The experimentally determined rotational constants are detailed in Table 1. We observe a very good agreement between the experimental and theoretical rotational constants. It should be noted that the number of lines assigned per state varies, mainly due to the decrease in signal-to-noise ratio (SNR) with increasing vibrational state energy, which means that some weaker lines observable for lower energy states blend into the noise for higher energy states. The sextic centrifugal distortion constants were required, as was the case for the vibronic ground state [5] and isotopologues (see below), to obtain the best possible fit.

The motions associated with the assigned fundamental modes could also be inferred from anharmonic frequency calculations. The ν_{21} vibration is a stretch of the N(1)-H bond, resulting in an out-of-plane displacement of the hydrogen atom. The ν_{20} and ν_{19} vibrations are wagging motions along the C(4)-C(5)-N(1) and the N(1)-C(2)-N(3) sections of the ring, respectively, which manifest as the out-of-plane displacement of the hydrogen atoms connected to these heavy atoms. Table 1 also contains depictions of these fundamental modes. Parts of the experimental spectra are shown in Figure 1, and they are compared to the simulated spectra based on the fitted rotational constants for the assigned vibrational states. The spectra of these states, like those of the imidazole vibronic ground state and isotopologues in the same frequency region, are characterised by an intense Q-branch ($\Delta J=0$) ladder, which culminates with a set of R-branch ($\Delta J=+1$) transitions [5].

The rigidity of the imidazole ring motif may be hinted towards by the relatively high energy of the vibrationally excited states. The lowest energy vibrational state, ν_{21} , has $\Delta E^{exp}=513\text{ cm}^{-1}$ relative to the imidazole ground state [39]. In contrast, imidazolidine's lowest energy vibrational state, ν_{33} , is only $\Delta E^a=71\text{ cm}^{-1}$ above the ground state at MP2/aug-cc-pVTZ, pointing towards the increased flexibility of the fully saturated structure.

The vibrationally excited states of a number of other aromatic molecules have been characterised, and one can gain valuable information from comparisons between such systems. The energies of the vibrationally excited states of the nitrogen-containing aromatic molecule indole have previously been reported. The energies of the ν_{42} and ν_{41} fundamental modes were determined as 207 and 241 cm^{-1} , respectively [40, 41], and the rotational spectra of these states were analysed recently in Ref. [42]. The ν_{42} state is an out-of-plane butterfly motion, which deforms the ring structure [40]. These states are substantially lower in energy compared to the lowest vibrational states of imidazole. The first indole vibrational motion analogous to an imidazole motion is found about 740 cm^{-1} above the ground state, which is a C-H bending motion [43].

An important implication is the effect that the vibrational energies of imidazole may have on the probabilities of astronomical detection. The difference in the SNR between the experimental spectra of the imidazole ground state and the lowest energy vibrationally excited state, ν_{21} , is a factor of 10 at approximately 300 K. This factor matches well with the assumed Boltzmann distribution at 300 K, which suggests a population difference between the vibronic ground state and the ν_{21} state of a factor of twelve. For hydantoin, a structurally related molecule to imidazole, for which a structural comparison is laid out in the next section, the ratio of the SNR of the ground state

to the first vibrationally excited state (140 cm^{-1} above the ground state) was found to be 1:0.6 [44]. This intensity ratio was reported for the experiment at 420 K [44], and the Boltzmann population difference at that temperature is a factor of 1.5. This reinforces experimentally the higher relative energies of the imidazole vibrationally excited states. The detection of the imidazole vibrational states will therefore require searches in regions with relatively higher temperatures, where the states are more substantially populated. This can also be seen in the Boltzmann distributions at the lower temperatures. For example, in the Sagittarius B2 region N1S, where a rotational temperature of approximately 160 K has been previously determined [45], the ground state to ν_{21} state population difference is a ratio of 100:1. The Boltzmann population of the vibrationally excited states is virtually zero at the low temperatures of TMC-1.

Table 1: The experimentally determined rotational constants of the observed imidazole vibrational states, which are tabulated alongside the theoretical rotational constants (obtained by CCSD(T) calculations extrapolated to the CBS limit and augmented by vibrational corrections computed at the fc-MP2/cc-pVQZ level of theory (hybrid approach); see Section 2 for details). Values contained in square brackets were kept fixed to the ground state values determined experimentally in Ref. [5]. The harmonic and anharmonic energies of each state relative to the ground state computed with the hybrid approach are also given. The experimental values ΔE^{exp} are taken from Ref. [39]. The structures show the associated motions of the states (see text for full descriptions). The figures depict the optimised structure of imidazole superimposed with the extreme positions of the atoms during the vibration. The fits include transitions assigned in the 75-110 and 260-295 GHz regions.

| Mode assigned | ν_{21} | | ν_{20} | | ν_{19} | |
|-----------------------------------|----------------|---------|----------------|---------|----------------|---------|
| | Experiment | Theory | Experiment | Theory | Experiment | Theory |
| $\Delta E^h / \text{cm}^{-1}$ | | 512.3 | | 643.2 | | 678.7 |
| $\Delta E^a / \text{cm}^{-1}$ | | 498.6 | | 628.0 | | 663.2 |
| $\Delta E^{exp} / \text{cm}^{-1}$ | 513 | | 626 | | 668 | |
| A / MHz | 9705.93562(85) | 9704.54 | 9695.8368(15) | 9693.12 | 9703.4261(18) | 9698.51 |
| B / MHz | 9356.37255(79) | 9359.01 | 9352.6093(15) | 9356.95 | 9344.6215(14) | 9347.66 |
| C / MHz | 4771.23603(40) | 4772.53 | 4773.41203(66) | 4773.12 | 4773.48988(67) | 4773.72 |
| D_J / kHz | 3.26871(40) | | 3.22412(84) | | 3.21492(71) | |
| D_{JK} / kHz | -5.15120(53) | | -5.0235(12) | | -4.98702(63) | |
| D_K / kHz | 2.23314(36) | | 2.14345(85) | | 2.11768(43) | |
| d_1 / kHz | 0.01875(31) | | 0.0109(21) | | 0.0512(15) | |
| d_2 / kHz | 0.01779(69) | | 0.0310(51) | | 0.0148(27) | |
| H_J / mHz | [1.367] | | [1.367] | | [1.367] | |
| H_{JK} / mHz | [-6.148] | | [-6.148] | | [-6.148] | |
| H_{KJ} / mHz | [-10.15] | | [-10.15] | | [-10.15] | |
| H_K / mHz | [-2.813] | | [-2.813] | | [-2.813] | |
| h_1 / mHz | [0.0459] | | [0.0459] | | [0.0459] | |
| h_2 / mHz | [-0.757] | | [-0.757] | | [-0.757] | |
| h_3 / mHz | [-0.0147] | | [-0.0147] | | [-0.0147] | |
| $L_{JKK} / \mu\text{Hz}$ | [-1.98] | | [-1.98] | | [-1.98] | |
| N_{lines} | 234 | | 119 | | 115 | |
| σ / kHz | 20.6 | | 17.9 | | 27.4 | |

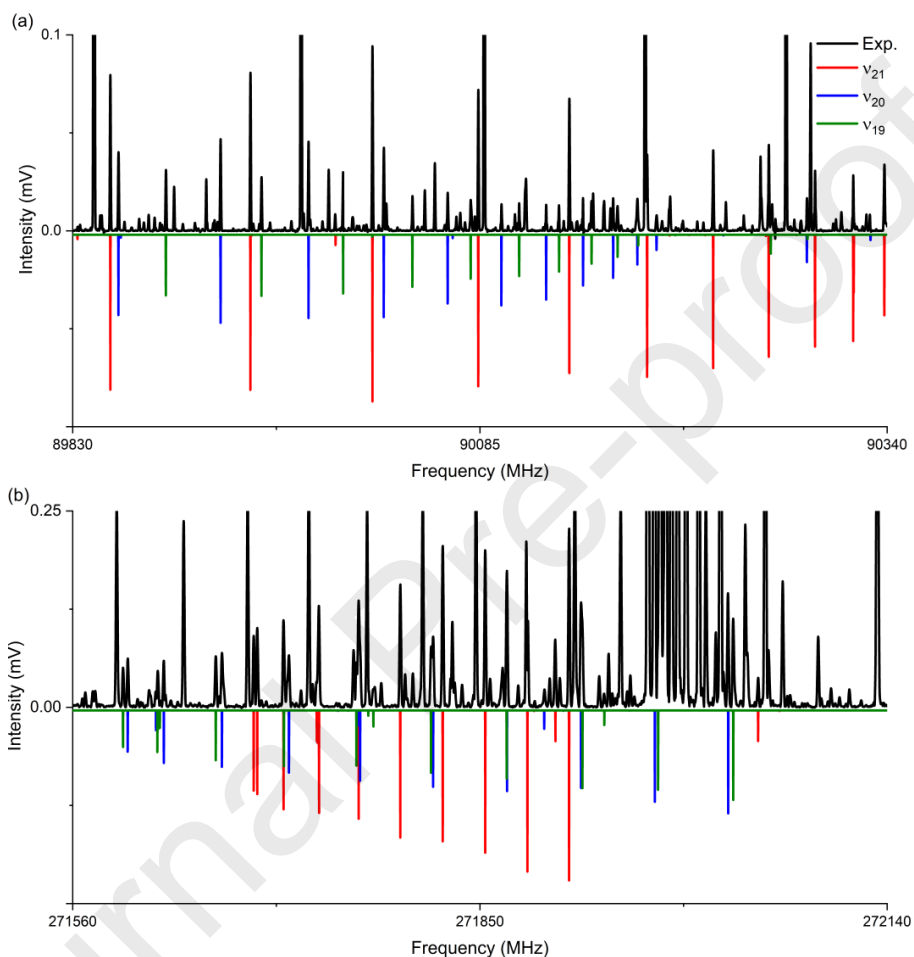


Figure 1: Parts of the rotational spectrum of imidazole (black) with the simulated spectra based on the experimentally determined rotational constants for the three assigned vibrationally excited states at 300 K (ν_{21} is shown in red, ν_{20} is shown in blue, and ν_{19} is shown in green). (a) A zoom-in of the W-band spectrum, showing a number of assigned Q-branch transitions. (b) Assigned R-branch transitions of the normal modes in the 260-295 GHz region. The simulated spectra intensities have been normalised to the experimental spectrum for clarity.

3.2 Heavy-atom Isotopologues

Experimental values for the ground state rotational constants of many isotopic variants of imidazole (three ^{13}C s, two ^{15}N s, four Ds, and one 1,2- D_2) were already available in the literature [19]. In addition to these rotational constants, we used theoretically computed values for the quartic and sextic distortion constants (see Section 2) as the initial input for our spectral analyses. The spectral assignments of the singly-substituted ^{13}C - and ^{15}N -isotopologues were thus accomplished in a straightforward way. We noted that, up to moderately low J levels ($J \leq 25$), the centrifugal distortion effects are small enough to be considered almost isotopically invariant, and the set of parameters determined for the parent species works equally well in modelling the rotational energy trends. As the value of J increased ($25 \leq J \leq 50$), the use of isotopologue-specific centrifugal distortion constants becomes critical to achieve a satisfactory fit of the measured transition frequencies. The small d_1 and d_2 constants, in particular, display significant changes in magnitude among isotopologues. For the $^{13}\text{C}(2)$ and $^{13}\text{C}(5)$ isotopologues, the sign also changed with respect to the parent isotopologue's distortion constants. While this effect is predicted by the *ab initio* results (see Table 2), it is a less commonly observed phenomenon upon isotopic substitution of a heavy atom, and the behaviour is worth a more detailed treatment. The distortion constants are a linear combination of the values $\tau'_{\alpha\beta\gamma\delta}$, which are introduced as an energy correction term in the Hamiltonian of the rigid rotor to account for distortion [46]. The six $\tau'_{\alpha\beta\gamma\delta}$ values relevant for the determination of the quartic centrifugal distortion constants are τ'_{aaaa} , τ'_{bbbb} , τ'_{cccc} , τ'_{aabb} , τ'_{aacc} , and τ'_{bbcc} . The underlying theory and relevant equations used to determine these values can be found in Ref. [46].

The calculated $\tau'_{\alpha\beta\gamma\delta}$ for all the heavy-atom isotopologues (see Table 2) do not show significant deflection from the values obtained for the parent isotopologue. These $\tau'_{\alpha\beta\gamma\delta}$ values can be linearly combined to give the quartic centrifugal distortion constants. Upon isotopic substitution, the inertial axes can shift, which can result in a small change in the moment of inertia elements in the inverse moment of inertia matrix. This can lead to a small change in the $\tau'_{\alpha\beta\gamma\delta}$ values. Since the distortion constants are the linear combination of these $\tau'_{\alpha\beta\gamma\delta}$ values, a change in magnitude and sign of the quartic centrifugal distortion constants may be possible, as is the case for imidazole, specifically seen in d_1 and d_2 . The Supplementary Material contains a comparison between the values of τ'_{aaaa} and τ'_{bbbb} determined in this work and those determined in Ref. [19]. The analysis performed in this work increases the precision of these values by two orders of magnitude. The phenomenon of a change of magnitude or sign of the quartic centrifugal distortion constants upon isotopic substitution is typically more pronounced in lighter molecules such as water, H_2O , and hydrogen sulphide, H_2S [46], and ring structures, such as 1,3,4-oxadiazoles, $\text{C}_2\text{H}_2\text{N}_2\text{O}$ [47], and pyridine, $\text{C}_5\text{H}_5\text{N}$ [48].

Table 2: A comparison of the values determined for τ'_{aaaa} , τ'_{bbbb} , τ'_{cccc} , τ'_{aabb} , τ'_{aacc} , and τ'_{bbcc} for the imidazole main isotopologue and the five main heavy-atom isotopologues.

| $\tau'_{aff\%b}$ | Main isotopologue | $^{15}\text{N}(1)$ | $^{13}\text{C}(2)$ | $^{15}\text{N}(3)$ | $^{13}\text{C}(4)$ | $^{13}\text{C}(5)$ |
|----------------------|-------------------|--------------------|--------------------|--------------------|--------------------|--------------------|
| τ'_{aaaa} / kHz | -13.11972(62) | -13.23030(201) | -12.50712(269) | -13.14752(807) | -13.20630(142) | -12.83071(185) |
| τ'_{bbbb} / kHz | -12.80252(62) | -12.16344(201) | -13.12808(269) | -12.39952(807) | -12.08770(142) | -12.90560(185) |
| τ'_{cccc} / kHz | -1.3897(12) | -1.3539(42) | -1.3433(33) | -1.440(351) | -1.3434(33) | -1.3472(39) |
| τ'_{aabb} / kHz | -13.4288(38) | -13.4092(130) | -12.6863(112) | -12.9389(1049) | -13.3388(102) | -12.5681(121) |
| τ'_{aacc} / kHz | -2.6800(26) | -2.6421(89) | -2.5414(68) | -2.5420(749) | -2.7176(69) | -2.5151(82) |
| τ'_{bbcc} / kHz | -2.8729(24) | -2.7659(80) | -2.8293(65) | -2.8148(662) | -2.6596(65) | -2.8590(75) |

Table 3 shows the experimentally determined constants and compares them to those obtained theoretically. Figures 2 and 3 show the experimental spectrum compared with the simulated spectra of the fitted rotational constants for the assigned isotopologues. Note that Figure 2 shows a section of the low frequency data set (12-15.5 GHz), where the spectra of the isotopologues exhibit HFS due to the presence of the ^{14}N nucleus/nuclei. Figure 3 shows a section of the high frequency data set (260-295 GHz), where this splitting is no longer present due to the collapse at the relatively high principal rotational energy level, J (transitions assigned in the 260-295 GHz region have J values $17 \leq J \leq 52$). The parameters provided in Table 3 are a result of fitting the rotational constants and centrifugal distortion constants to all the assigned line frequencies and fitting quadrupole coupling constants only to line frequencies in the 2 - 26 GHz region, as described in the previous section.

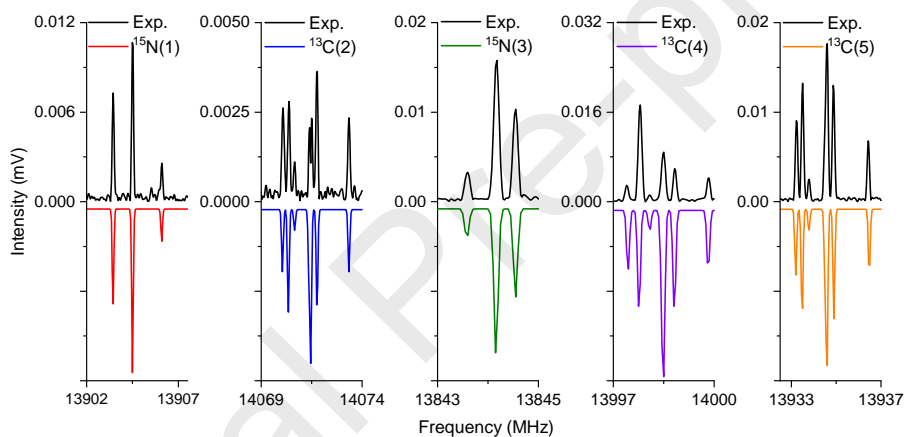


Figure 2: A zoom-in of the assigned spectra of the singly-substituted isotopologues of imidazole in the frequency range 12-15.5 GHz, where the $1_{0,1}-0_{0,0}$ transitions show HFS. The experimental spectra are shown in black, and the simulated spectra of the fitted rotational constants at 3 K are shown in colour. The atom labelling scheme is detailed in Table 1 and Figure 4. The simulated spectra intensities have been normalised to the experimental spectrum for clarity.

Table 3: A comparison of the experimental rotational constants for the singly-substituted heavy-atom isotopologues of imidazole determined in this work and those theoretically calculated. The theoretical rotational constants were obtained by CCSD(T) calculations extrapolated to the CBS limit and augmented by vibrational corrections computed at the fc-MP2/cc-pVQZ level of theory, the quartic centrifugal distortion constants were calculated at the fc-CCSD(T)/cc-pVQZ level of theory, and the sextic centrifugal distortion constants were calculated at the fc-CCSD/cc-pVTZ level of theory; see Section 2. The values contained in square brackets were kept fixed to theoretical values. Note that the value of L_{JKK} was fixed to the value obtained from the main isotopologue analysis in Ref. [5]. The fits from this work include transitions assigned in the 2-8, 12-15.5, 18-26, 75-110, and 260-295 GHz regions. The positions of each isotopic substitution can be found in Table 1 and Figure 4.

| | ¹⁵ N(1) | | ¹³ C(2) | | ¹⁵ N(3) | |
|---|--------------------|-----------|--------------------|-----------|--------------------|-----------|
| | Experiment | Theory | Experiment | Theory | Experiment | Theory |
| <i>A</i> / MHz | 9695.27280(60) | 9693.2567 | 9522.2481(12) | 9516.0036 | 9721.5291(17) | 9716.6844 |
| <i>B</i> / MHz | 9188.15572(59) | 9191.9578 | 9353.9384(11) | 9362.2887 | 9135.7459(12) | 9142.5408 |
| <i>C</i> / MHz | 4716.17271(55) | 4716.3799 | 4717.38619(35) | 4717.6589 | 4708.48056(66) | 4708.8368 |
| <i>D_J</i> / kHz | 3.21549(45) | 3.1632 | 3.19784(45) | 3.1510 | 3.2036(11) | 3.1459 |
| <i>D_{JK}</i> / kHz | -5.06600(40) | -4.9784 | -5.05961(57) | -4.9705 | -5.0516(15) | -4.9506 |
| <i>D_K</i> / kHz | 2.18898(44) | 2.1506 | 2.19760(40) | 2.1527 | 2.1840(12) | 2.1389 |
| <i>d₁</i> / kHz | 0.066679(100) | 0.0657 | -0.03881(25) | -0.0310 | 0.04675(65) | 0.0458 |
| <i>d₂</i> / kHz | 0.020636(47) | 0.0209 | [-0.0033] | -0.0033 | 0.00511(54) | 0.0060 |
| <i>H_J</i> / mHz | [1.3169] | 1.3169 | [1.3326] | 1.3326 | [1.3100] | 1.3100 |
| <i>H_{JK}</i> / mHz | [-5.3526] | -5.3526 | [-5.3844] | -5.3844 | [-5.3218] | -5.3218 |
| <i>H_{KJ}</i> / mHz | [6.7607] | 6.7607 | [6.7795] | 6.7795 | [6.7177] | 6.7177 |
| <i>H_K</i> / mHz | [-2.7233] | -2.7233 | [-2.7263] | -2.7263 | [-2.7045] | -2.7045 |
| <i>h₁</i> / mHz | [0.0488] | 0.0488 | [0.0178] | 0.0178 | [0.0414] | 0.0414 |
| <i>h₂</i> / mHz | [-0.0443] | -0.0443 | [-0.0267] | -0.0267 | [-0.0292] | -0.0292 |
| <i>h₃</i> / mHz | [-0.0061] | -0.0061 | [-0.0103] | -0.0103 | [-0.0129] | -0.0129 |
| <i>L_{JKK}</i> / μHz | -1.98 | - | -1.98 | - | -1.98 | - |
| $\chi_{cc}^{(1)}$ / MHz | - | - | -2.608(35) | -2.666 | -2.592(24) | -2.666 |
| $\chi_{bb}^{(1)}-\chi_{aa}^{(1)}$ / MHz | - | - | [0] | -0.071 | [0] | -0.071 |
| $\chi_{cc}^{(3)}$ / MHz | 2.265(39) | 2.21 | 2.203(47) | 2.21 | - | - |
| $\chi_{bb}^{(3)}-\chi_{aa}^{(3)}$ / MHz | 4.812(56) | 5.74 | 5.720(68) | 5.74 | - | - |
| No. of lines | 248 | - | 399 | - | 244 | - |
| σ / kHz | 32.5 | - | 33.8 | - | 32.7 | - |
| | ¹³ C(4) | | ¹³ C(5) | | | |
| | Experiment | Theory | Experiment | Theory | | |
| <i>A</i> / MHz | 9573.81156(53) | 9576.1789 | 9632.88401(69) | 9626.3298 | | |
| <i>B</i> / MHz | 9285.89570(41) | 9286.6274 | 9223.79544(49) | 9232.0577 | | |
| <i>C</i> / MHz | 4712.53945(32) | 4712.9821 | 4710.66720(35) | 4710.9362 | | |
| <i>D_J</i> / kHz | 3.20500(31) | 3.1559 | 3.19777(34) | 3.1456 | | |
| <i>D_{JK}</i> / kHz | -5.06574(35) | -4.9750 | -5.04990(41) | -4.9574 | | |
| <i>D_K</i> / kHz | 2.19660(29) | 2.1524 | 2.18893(41) | 2.1454 | | |
| <i>d₁</i> / kHz | 0.069913(77) | 0.0649 | -0.00468(15) | -0.0031 | | |
| <i>d₂</i> / kHz | 0.021625(40) | 0.0227 | -0.009635(45) | -0.0078 | | |
| <i>H_J</i> / mHz | [1.3210] | 1.3210 | [1.3259] | 1.3259 | | |
| <i>H_{JK}</i> / mHz | [-5.3628] | -5.3628 | [-5.3645] | -5.3645 | | |
| <i>H_{KJ}</i> / mHz | [6.7713] | 6.7713 | [6.7576] | 6.7576 | | |
| <i>H_K</i> / mHz | [-2.7278] | -2.7278 | [-2.7175] | -2.7175 | | |
| <i>h₁</i> / mHz | [0.0414] | 0.0414 | [0.0341] | 0.0341 | | |
| <i>h₂</i> / mHz | [-0.0284] | -0.0284 | [-0.0037] | -0.0037 | | |
| <i>h₃</i> / mHz | [-0.0014] | -0.0014 | [-0.0051] | -0.0051 | | |
| <i>L_{JKK}</i> / μHz | -1.98 | - | -1.98 | - | | |
| $\chi_{cc}^{(1)}$ / MHz | -2.566(21) | -2.666 | -2.541(30) | -2.666 | | |
| $\chi_{bb}^{(1)}-\chi_{aa}^{(1)}$ / MHz | [0] | -0.071 | [0] | -0.071 | | |
| $\chi_{cc}^{(3)}$ / MHz | 2.281(27) | 2.21 | 2.171(29) | 2.21 | | |
| $\chi_{bb}^{(3)}-\chi_{aa}^{(3)}$ / MHz | 3.028(40) | 5.74 | [5.74] | 5.74 | | |
| No. of lines | 422 | - | 360 | - | | |
| σ / kHz | 25.9 | - | 27.6 | - | | |

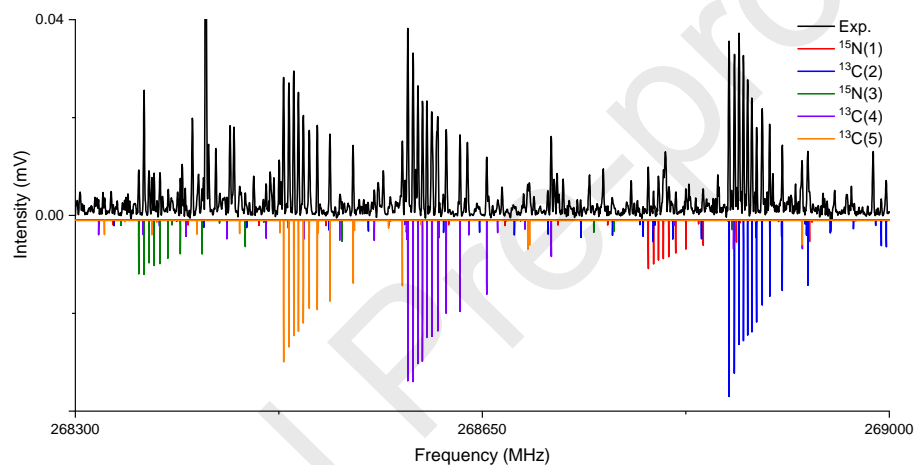


Figure 3: A zoom-in of the 260-295 GHz region, showing the assignments of the singly-substituted isotopologues of imidazole, which do not exhibit hyperfine structure. The experimental spectra are shown in black, and the simulated spectra of the fitted rotational constants at 300 K are shown in colour. The atom labelling scheme is detailed in Table 1 and Figure 4. The simulated spectra intensities have been normalised to the experimental spectrum for clarity.

3.3 Molecular Structures

The assigned isotopologues allowed us to determine the structure of imidazole in the gas phase. This was done to obtain both the substitution (r_s) and the effective (r_0) structures. These structures were determined by taking advantage of the presence of heavy-atom isotopologues in natural abundance, each of which has a distinct rotational spectrum. The r_s structure analysis involves applying Kraitchman's equations to obtain the coordinates of the substituted heavy atom. The r_0 structure was determined by a least-squares fit, which calculated the bond lengths and angles that best reproduced the rotational constants of the main and heavy-atom isotopologues. Figure 4 shows an overlay of the theoretical structure optimised using the coupled-cluster composite method [5], shown in transparent colour, with the experimentally determined atom positions from the r_s structure analysis, shown as solid spheres. The experimentally determined bond distances between the nuclei that make up the ring structure of imidazole and the experimentally determined internal angles of the five-membered ring are labelled. In Tables 4 and 5, these structural parameters are compared to the parameters presented in Refs. [5] and [19]. It can be seen that the substitution structure obtained from this analysis shows little deviation from that reported in Ref. [19], as can be expected from the good agreement of the rotational constants of each of the isotopologues. Indeed, whilst the experimental structural parameters - the bond lengths and internal bond angles - are of comparable quality to those in Ref. [19], the rotational constants are more precise by up to two significant figures / orders of magnitude.³ Our extended line lists and more precise rotational constants, which were derived from experimental data with an increased frequency range, are nevertheless important, especially for astrophysical applications at higher frequencies, with which our extended data sets overlap.

Table 4: A comparison of imidazole bond lengths, in Å, determined in this work and by Christen *et al.* from the r_s and r_0 structure analysis [19], and those from the structure optimised using the coupled-cluster composite method [5] and the MP2/aug-cc-pVTZ level of theory.

| Bond length / Å | r_s | | r_0 | | Theory | |
|-----------------|-----------|-----------------------------|-----------|-----------------------------|----------------------------|---------------|
| | This work | Christen <i>et al.</i> [19] | This work | Christen <i>et al.</i> [19] | Giuliano <i>et al.</i> [5] | This work MP2 |
| N(1)-C(2) | 1.365(6) | 1.3643 | 1.367(6) | 1.3704 | 1.347 | 1.362 |
| C(2)-N(3) | 1.321(7) | 1.3135 | 1.315(6) | 1.3164 | 1.286 | 1.322 |
| N(3)-C(4) | 1.373(7) | 1.3822 | 1.384(7) | 1.3875 | 1.370 | 1.373 |
| C(4)-C(5) | 1.370(3) | 1.3638 | 1.370(6) | 1.3676 | 1.347 | 1.377 |
| N(1)-C(5) | 1.375(3) | 1.3774 | 1.383(7) | 1.3788 | 1.368 | 1.372 |

This structure of imidazole can be compared directly to the structures of related molecules. Indeed, in Ref. [19], the imidazole structure was discussed in comparison with those of oxazole, pyrrole, and pyrazole. To add to these, and to put the structure of imidazole (c -NHCHN(CH₂)₂) in context with rings of the same elemental composition and connectivity, we present a comparison to the fully saturated ring analogues hydantoin (c -NHC(O)NHCH₂C(O)) and imidazolidine (c -NHCH₂NH(CH₂)₂).

³The error in the structural parameters was not reported directly in Ref. [19], but we have reprocessed their data in order to calculate the associated errors in order to facilitate a thorough comparison. The relevant data can be found in the Supplementary Material.

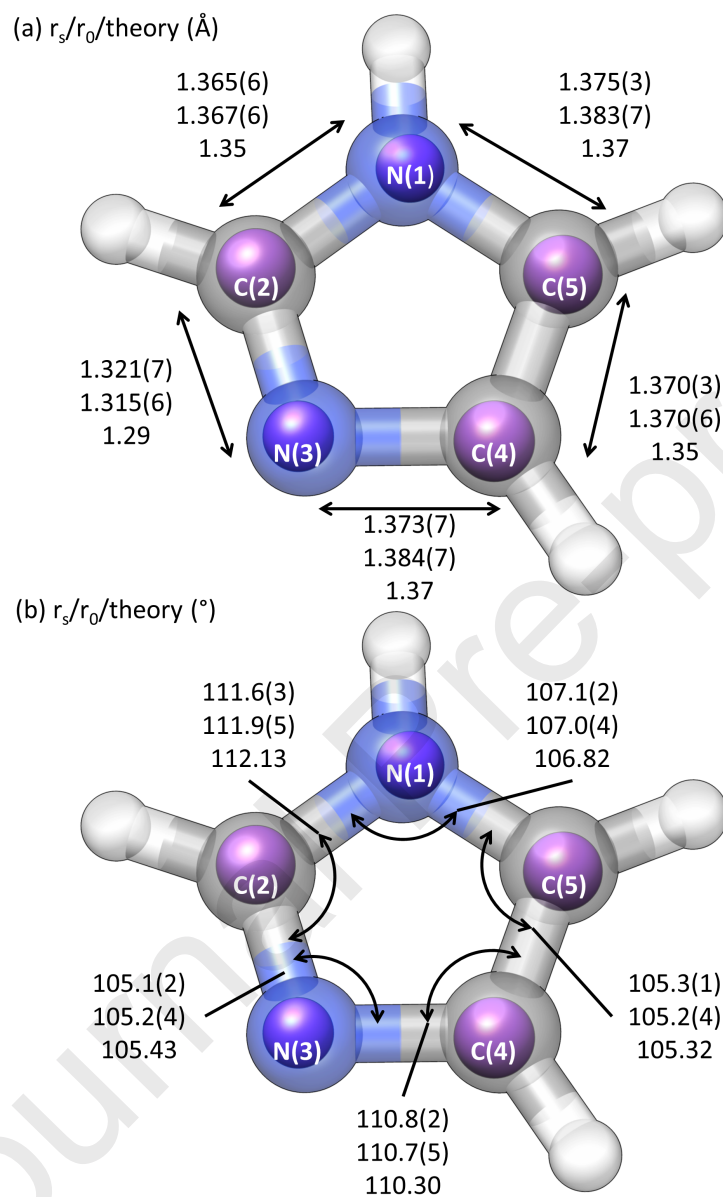
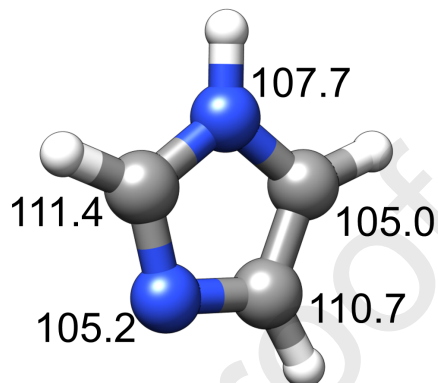
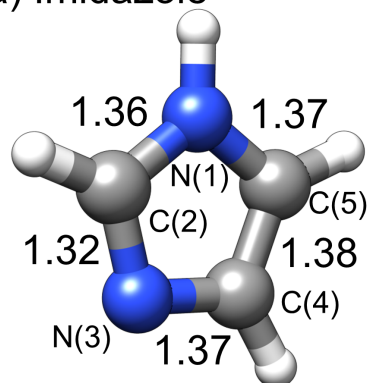
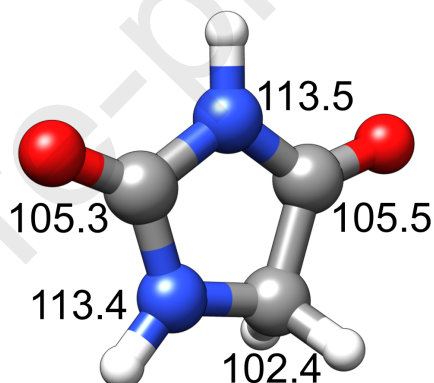
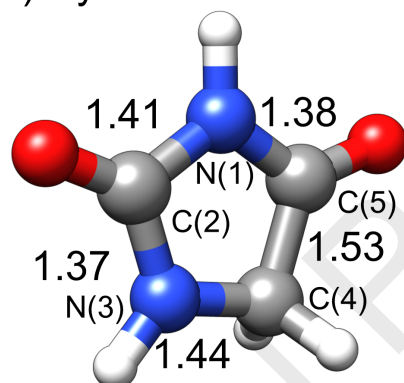


Figure 4: The theoretical structure of imidazole, optimised using the coupled-cluster composite method [5] and shown in transparent colour, is compared to the experimentally determined atom positions from the r_s structure analysis, shown as solid spheres. Three sets of values are shown, which are from the r_s structure analysis, the r_0 structure analysis, and from the coupled-cluster composite method. The bond lengths between the heavy atoms of the ring are labelled in (a), and the internal angles of the ring are given in (b).

(a) Imidazole



(b) Hydantoin



(c) Imidazolidine

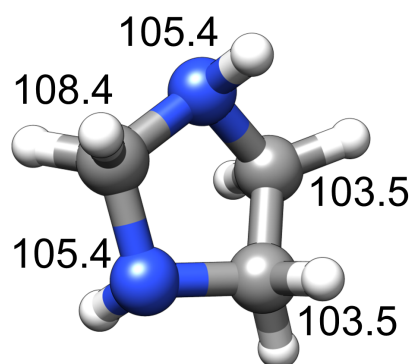
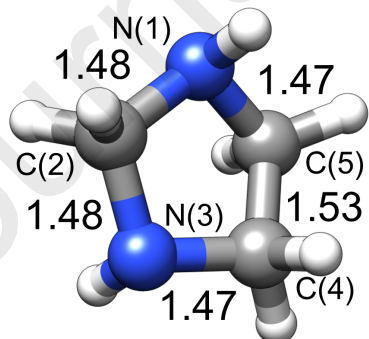


Figure 5: The structures of (a) imidazole, (b) hydantoin, and (c) imidazolidine. The bond lengths between the ring atoms are shown, in Å, on the left side, while the internal angles are marked at the atom of the angle vertex, in degrees ($^{\circ}$), on the right side. The values are all from structure optimisation calculations at the MP2/aug-cc-pVTZ level of theory.

Table 5: A comparison of the internal bond angles of the imidazole ring, in degrees ($^{\circ}$), determined in this work and by Christen *et al.* from the r_s and r_0 structure analysis [19], and those from the structure optimised using the coupled-cluster composite method [5] and the MP2/aug-cc-pVTZ level of theory.

| Bond angle / $^{\circ}$ | r_s | | r_0 | | Theory | |
|-------------------------|-----------|-----------------------------|-----------|-----------------------------|----------------------------|---------------|
| | This work | Christen <i>et al.</i> [19] | This work | Christen <i>et al.</i> [19] | Giuliano <i>et al.</i> [5] | This work MP2 |
| N(1)-C(2)-N(3) | 111.6(3) | 111.99 | 111.9(5) | 111.85 | 112.1 | 111.4 |
| C(2)-N(3)-C(4) | 105.1(2) | 104.93 | 105.2(4) | 105.08 | 105.4 | 105.2 |
| N(3)-C(4)-C(5) | 110.8(2) | 110.69 | 110.7(5) | 110.50 | 110.3 | 110.7 |
| C(4)-C(5)-N(1) | 105.3(1) | 105.48 | 105.2(4) | 105.65 | 105.3 | 105.0 |
| C(5)-N(1)-C(2) | 107.1(2) | 106.90 | 107.0(4) | 106.93 | 106.8 | 107.7 |

Figure 5 shows the ring structures of the three molecules along with the relevant bond lengths and internal angles. The values are taken from structure optimisation calculations at the MP2/aug-cc-pVTZ level of theory, which were performed for this work, in particular because of the absence of the experimental gas-phase structure for imidazolidine in the literature. The main difference between the structures is the degree of saturation of the ring bonds, and this impacts the bond lengths. The C(2)-N(3) bond illustrates this clearly. In the fully aromatic system, imidazole, where the bond is formally a carbon-nitrogen double bond, the experimental bond length is 1.321(7) Å. This can be compared to the theoretical value for this bond length, which is 1.32 Å at the MP2/aug-cc-pVTZ level of theory. This demonstrates the excellent agreement between the MP2 calculation and the r_s structure (for most parameters, a better agreement than between the experimental structure and the structure from the coupled-cluster method), which we can exploit for further comparisons. The C(2)-N(3) bond increases to 1.37 Å in hydantoin, in which the bond has some double bond character due to resonance effects. From an experimental point of view, as can be seen in Tables 6 and 7, the bond distance in hydantoin⁴ is 1.359(10) Å. This is also essentially the same as the theoretical value. The C(2)-N(3) bond is the longest in the fully saturated imidazolidine, where it is a carbon-nitrogen single bond, at 1.48 Å. The C(5)-N(1) bond in hydantoin also has some double bond character, and this bond, at 1.38 Å, is similar in length to the corresponding bond in imidazole (1.37 Å). The same bond in imidazolidine is longer, at 1.47 Å. The comparison of the experimental gas-phase structures shows a slight elongation in the hydantoin C(5)-N(1) bond (1.401(5) Å) when compared to the same imidazole bond (1.375(3) Å).

Another consequence of unsaturation and aromaticity is structure planarity. This can be characterised by the dihedral angles between the atoms in the ring, where Θ_1 is the dihedral angle N(1)-C(2)-N(3)-C(4), Θ_2 is C(2)-N(3)-C(4)-C(5), Θ_3 is N(3)-C(4)-C(5)-N(1), Θ_4 is C(4)-C(5)-N(1)-C(2), and Θ_5 is C(5)-N(1)-C(2)-N(3). In the fully saturated, non-aromatic imidazolidine, the ring structure is non-planar. The dihedral angles in imidazolidine are: $\Theta_1 = 11.1^{\circ}$, $\Theta_2 = -27.8^{\circ}$, $\Theta_3 = 34.5^{\circ}$, $\Theta_4 = -27.8^{\circ}$, and $\Theta_5 = 11.1^{\circ}$. These values are in contrast to those of the other two structures, where the dihedral angles are essentially zero for both imidazole and hydantoin. This planar structure is expected for imidazole because of its aromatic nature and for hydantoin

⁴A full account of the assigned transitions and the experimentally determined rotational constants and heavy-atom coordinates for hydantoin can be found in the Supplementary Material.

Table 6: A comparison of hydantoin bond lengths, in Å, determined from the r_s structure analysis and the MP2/aug-cc-pVTZ structure optimisation calculation. The atom labelling scheme is detailed in Figure 5.

| Bond length / Å | r_s | Theory |
|-----------------|-----------|--------|
| N(1)-C(2) | 1.389(6) | 1.407 |
| C(2)-N(3) | 1.359(10) | 1.366 |
| N(3)-C(4) | 1.445(3) | 1.444 |
| C(4)-C(5) | 1.516(9) | 1.526 |
| N(1)-C(5) | 1.401(5) | 1.374 |

Table 7: A comparison of the internal bond angles of the hydantoin ring, in degrees ($^\circ$), determined from the r_s structure analysis and the MP2/aug-cc-pVTZ structure optimisation calculation. The atom labelling scheme is detailed in Figure 5.

| Bond angles / $^\circ$ | r_s | Theory |
|------------------------|----------|--------|
| N(1)-C(2)-N(3) | 107.0(3) | 105.3 |
| C(2)-N(3)-C(4) | 112.7(2) | 113.4 |
| N(3)-C(4)-C(5) | 102.7(2) | 102.4 |
| C(4)-C(5)-N(1) | 105.4(2) | 105.5 |
| C(5)-N(1)-C(2) | 112.2(5) | 113.5 |

because of the carbonyl groups introducing resonance effects. The planarity can be explained by the hybridisation of the atoms in the ring motif. The sp^2 -hybridised molecular orbitals in the imidazole ring locks the planar structure, whilst the sp^3 -hybridised molecular orbitals in imidazolidine force some deviation from planarity.

The planarity of the imidazole ring can also be revealed by considering the second moments of inertia, P_{ii} , along the three axes, $i = a, b, c$. The values of P_{ii} can indicate the extension of mass along the i axis. The rotational constants, moments of inertia, and second moments of inertia for the imidazole normal isotopologue (reported in Ref. [5]) and the rare isotopologues (reported here) are detailed in Table 8, and the values were converted using the protocols contained in Ref. [49]. The tabulated values of P_{cc} , which were all calculated using the experimentally determined rotational constants, are around zero. Therefore, the mass of the imidazole ring atoms lie on the ab plane; there is almost zero extension of mass along the c axis. The fact that the P_{cc} values for the isotopologues show little-to-no change when compared to the value for the normal isotopologue further shows that the substituted atoms lie on the ab plane, reinforcing the conclusion of planarity.

Further structural information can be retrieved from the values of the nuclear quadrupole coupling constants, χ_i , which are an indication of the electronic environment of the respective nuclei. In Ref. [50], the authors concluded that the hydantoin molecule contains two pyrrolic nitrogen nuclei. That is, the lone pairs of electrons on the nitrogen nuclei are within the aromatic ring system. The electrons are, therefore, unavailable for bonding to an external electrophile, as is the case in pyrrole. This was signified by the negative χ_{cc} values $-4.7338(70)$ and $-3.4635(75)$ MHz for N(1) and N(3), respectively

Table 8: The experimental rotational constants (A, B, C), moments of inertia (I_a , I_b , I_c), and second moments of inertia (P_{aa} , P_{bb} , P_{cc}) for the normal isotopologue and heavy-atom isotopologues of imidazole.

| | Rotational constant (MHz) | Moment of inertia (amu.Å ²) | Second moment of inertia (amu.Å ²) |
|----------------------|---------------------------|---|--|
| Main isotopologue | | | |
| A / I_a / P_{aa} | 9725.30651 | 51.96535497 | 53.9272361 |
| B / I_b / P_{bb} | 9373.99203 | 53.91289041 | 51.97970066 |
| C / I_c / P_{cc} | 4771.91599 | 105.9069368 | -0.014345693 |
| ¹⁵ N(1) | | | |
| A / I_a / P_{aa} | 9695.27145 | 52.12633887 | 55.01783889 |
| B / I_b / P_{bb} | 9188.15484 | 55.00331827 | 52.14085948 |
| C / I_c / P_{cc} | 4716.17342 | 107.1586984 | -0.014520619 |
| ¹³ C(2) | | | |
| A / I_a / P_{aa} | 9522.243 | 53.07352532 | 54.0430351 |
| B / I_b / P_{bb} | 9353.9408 | 54.02845879 | 53.08810164 |
| C / I_c / P_{cc} | 4717.38675 | 107.1311367 | -0.014576318 |
| ¹⁵ N(3) | | | |
| A / I_a / P_{aa} | 9721.5279 | 51.98555311 | 55.33353593 |
| B / I_b / P_{bb} | 9135.7461 | 55.31885403 | 52.00023501 |
| C / I_c / P_{cc} | 4708.48085 | 107.3337709 | -0.014681892 |
| ¹³ C(4) | | | |
| A / I_a / P_{aa} | 9573.8098 | 52.78765879 | 54.43902448 |
| B / I_b / P_{bb} | 9285.8951 | 54.42437154 | 52.80231172 |
| C / I_c / P_{cc} | 4712.53924 | 107.2413362 | -0.014652938 |
| ¹³ C(5) | | | |
| A / I_a / P_{aa} | 9632.88165 | 52.46394831 | 54.80539971 |
| B / I_b / P_{bb} | 9223.79437 | 54.79079267 | 52.47855535 |
| C / I_c / P_{cc} | 4710.66717 | 107.2839551 | -0.014607039 |

[50]. For the vibronic ground state of imidazole, the χ_{cc} values indicate that the N(1) nucleus is pyrrolic and the N(3) nucleus is pyridinic. In the latter case, the lone pair lies outside of the ring structure, and it is available for bonding to an external electrophile, as is the case in pyridine. The χ_{cc} values for the parent species of imidazole were determined to be $-2.5592(62)$ and $2.2293(66)$ MHz for N(1) and N(3), respectively [5]. This can also be seen in Table 3, where all of the values determined for $\chi_{cc}^{(1)}$ for the isotopologues are negative, and the $\chi_{cc}^{(3)}$ values are all positive.

The experimental analysis up to 295 GHz for the vibrational states and the isotopologues allows for the prediction of the rotational spectra of these species to be accurate well beyond the laboratory frequency range. The most intense lines (to $\log_{10}(\text{Intensity}) = -5.0$ nm²MHz [34]) up to 600 GHz, for example, have frequency accuracies well below 1 MHz, meaning they can be relied upon for astrophysical searches.

4. Conclusions

The broadband rotational spectrum of imidazole was recorded in the regions 2-8, 12-15.5, 18-26, 75-110, and 260-295 GHz. The analysis of the parent species spectrum was previously reported [5], and the investigation is extended here to include three low-lying vibrationally excited states and the singly-substituted heavy-atom isotopologues.

The low-lying vibrational states assigned in the room-temperature data can be designated as an N-H bond stretch (the fundamental mode ν_{21}) and wagging motions along the C-C-N and N-C-N sections of the ring (the fundamental modes ν_{20} and ν_{19}). Their relatively high energies reinforce the notion that the aromatic ring structure is rigid, especially when compared to the energies of the vibrationally excited states of the more flexible molecule imidazolidine. The experimental signal levels and predicted Boltzmann populations of the imidazole vibrational states compared to the ground state suggest searches for these species could be successful in warm regions of the ISM.

The gas-phase structure of imidazole was determined. When compared to the previous structure from Ref. [19], there was excellent agreement, with changes in the bond lengths (within 0.67%) and bond angles (within 0.35%) in the r_s structures being negligible. The rotational spectrum of hydantoin was recorded in the 18-26 GHz region, and an experimental gas-phase structure was also obtained for this heterocycle. The effect of the aromatic character of the imidazole ring was also explored, with comparisons being drawn to the experimental structure of hydantoin and the theoretical structure of the related molecule imidazolidine. Aromaticity explains the decrease in lengths of the ring bonds in imidazole when compared to the other cyclic molecules.

Searches for all of these species in the ISM are facilitated with the line lists, rotational constants, centrifugal distortion constants, and nuclear quadrupole coupling constants that result from our comprehensive analysis.

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