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On biases in displacement estimation for image registration, with a focus on photomechanics

Frédéric Sur*, Benoît Blaysat†, Michel Grédiac†

Abstract

Image registration under small displacements is the keystone of several image analysis tasks such as optical flow estimation, stereoscopic imaging, or full-field displacement estimation in photomechanics. A popular approach consists in locally modeling the displacement field between two images by a parametric transformation and performing least-squares estimation afterwards. This procedure is known as “digital image correlation” (DIC) in several domains as in photomechanics. The present article is part of this approach. First, the estimated displacement is shown to be impaired by biases related to the interpolation scheme needed to reach subpixel accuracy, the image gradient distribution, as well as the difference between the hypothesized parametric transformation and the true displacement. A quantitative estimation of the difference between the estimated value and the actual one is of importance in application domains such as stereoscopy or photomechanics, which have metrological concerns. Second, we question the extent to which these biases could be eliminated or reduced. We also present numerical assessments of our predictive formula in the context of photomechanics. Software codes are freely available to reproduce our results. Although this paper is focused on a particular application field, namely photomechanics, it is relevant to various scientific areas concerned by image registration.

Keywords: Image registration, digital image correlation, fattening effect, Savitzky-Golay filtering, experimental mechanics, stereo-imaging, optical flow

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1 Introduction

Two images \mathcal{I} and \mathcal{I}' of a scene being given, image registration consists in estimating a bidimensional displacement field ϕ which permits to map pixels from one image to the corresponding pixels in the other image. Assuming that corresponding pixels have the same intensity, for any pixel \mathbf{x} , the following relation thus holds: $\mathcal{I}(\mathbf{x}) = \mathcal{I}'(\mathbf{x} + \phi(\mathbf{x}))$. When small displacements are sought (typically below one pixel), the displacement field is often estimated locally by imposing the preceding relation over corresponding subsets of the domains of \mathcal{I} and \mathcal{I}' . While image registration is the very first step of many image analysis tasks, the case of small amplitude displacements is of particular interest in optical flow estimation [11, 12], disparity estimation in stereoscopic imaging [9, 31, 33, 34], or extensometry in experimental mechanics [20, 43], to mention some representative examples. Although we will focus our presentation on displacement estimation in photomechanics, which is a branch of experimental solid mechanics dedicated to full-field measurements from images, many aspects of the problem discussed in this paper are of general interest. However, in this latest application domain, no occlusions are present, in contrast to stereoscopic imaging or optical flow estimation.

1.1 Problem statement

In photomechanics, digital image correlation (DIC) [20, 43] is of prime interest to measure displacement fields on the surface or in the bulk of materials subjected to thermo-mechanical loads. The term DIC equally refers to methods which are based on cross-correlation (CC), on sum of squared differences (SSD), or on normalized CC or SSD as well [27]. DIC methods are based on two images \mathcal{I} and \mathcal{I}' of the surface of the specimen taken before and after deformation, respectively. The specimen shall be marked beforehand with a contrasted random pattern, called a speckle pattern, as shown in Section 3. The aim is to retrieve the displacement field \mathbf{u} such that for any \mathbf{x} , $\mathcal{I}(\mathbf{x}) = \mathcal{I}'(\mathbf{x} + \mathbf{u}(\mathbf{x}))$. Strain fields are then deduced by differentiation. These fields are used to observe various phenomena, which occur on the surface of deformed specimens, and which are revealed by displacement or strain heterogeneities. They can eventually be used to identify parameters governing constitutive equations, which are then used to design structural components [19]. It is therefore of prime importance to retrieve displacement and strain maps affected by the lowest possible measurement errors, so that the resulting identified parameters are themselves as reliable as possible. The displacement field usually features values well below one pixel. The so-called *local subset-based DIC* consists in registering \mathcal{I}' on \mathcal{I} by optimizing some criterion as the above-mentioned CC or SSD over subsets of the image domains. These criteria being equivalent under mild assumptions [27], we focus on the sum of squared differences (SSD), which is also used in block matching for stereoscopic imaging [31]. Since \mathcal{I} and \mathcal{I}' are known only at integer pixel coordinates and since subpixel accuracy of \mathbf{u} is sought, interpolation is required. In order to estimate the displacement \mathbf{u} at a pixel \mathbf{x} , the

following SSD, defined over a subset $\Omega_{\mathbf{x}}$ of the image of the specimen surface centered at \mathbf{x} , is minimized with respect to a displacement $\phi^{\mathbf{x}}$ defined over $\Omega_{\mathbf{x}}$:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \phi^{\mathbf{x}}(\mathbf{x}_i)) \right)^2 \quad (1)$$

Here $\tilde{\mathcal{I}}'$ denotes a continuous interpolation of \mathcal{I}' (\mathcal{I}' being sampled at integer coordinates), and $\Omega_{\mathbf{x}}$ is a set of M pixels $(\mathbf{x}_1, \dots, \mathbf{x}_M)$. Minimization is performed with respect to a displacement field $\phi^{\mathbf{x}}$ expected to approximate the actual unknown displacement \mathbf{u} over $\Omega_{\mathbf{x}}$. We do not consider a weighted SSD criterion, but the formulas in the remainder of the paper would easily adapt.

An estimation $\phi(\mathbf{x})$ of the displacement field over the whole specimen surface is eventually obtained by taking, at any pixel \mathbf{x} , the value of $\phi^{\mathbf{x}}$ at the center of the subset $\Omega_{\mathbf{x}}$.

Displacement estimation is an ill-posed problem because of under-determination. Indeed, at each pixel, a bidimensional displacement must be retrieved. Moreover, information is lost from the component of the displacement orthogonal to the image gradient. This is the so-called aperture problem [11, 12]. Consequently, in experimental mechanics (and in general applications as well [7]), the displacement $\phi^{\mathbf{x}}$ is usually sought as the linear combination of N shape (or basis) functions $(\phi_j)_{1 \leq j \leq N}$, such that the parameters $(\lambda_j)_{1 \leq j \leq N}$ minimize

$$\text{SSD}(\lambda_1, \dots, \lambda_N) = \sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}' \left(\mathbf{x}_i + \sum_{j=1}^N \lambda_j \phi_j(\mathbf{x}_i) \right) \right)^2 \quad (2)$$

It may be noted that classic ways to deal with the aperture problem are either to consider a constant displacement $\phi^{\mathbf{x}}$ over $\Omega_{\mathbf{x}}$ and a first-order Taylor expansion of the preceding equation, giving fast approaches *à la* Lucas-Kanade [25], or to consider global smoothing with regularization constraints, giving approaches *à la* Horn-Schunck [21]. In experimental mechanics, it is often preferred to numerically minimize SSD because an accurate estimation of ϕ is required and setting the hyperparameter involved by regularization is not an easy task, which rules out such approaches.

From now on, we assume that $\mathbf{\Lambda} = (\lambda_1, \dots, \lambda_N)^T$ minimizes the expression of SSD given by Equation 2, and $\phi^{\mathbf{x}} = \sum_{j=1}^N \lambda_j \phi_j$ denotes the corresponding local displacement field. For example, as recalled in [5], zero-order shape functions are such that $N = 2$ and $\phi_1(\mathbf{x}) = (1 \ 0)^T$, $\phi_2(\mathbf{x}) = (0 \ 1)^T$, giving a constant $\phi^{\mathbf{x}}$ over $\Omega_{\mathbf{x}}$. First-order shape functions are such that $N = 6$ and functions $\phi_3(\mathbf{x}) = (x \ 0)^T$, $\phi_4(\mathbf{x}) = (0 \ x)^T$, $\phi_5(\mathbf{x}) = (y \ 0)^T$, $\phi_6(\mathbf{x}) = (0 \ y)^T$ are added to ϕ_1 and ϕ_2 . Second-order shape functions are such that $N = 12$ and embed the following additional functions: $\phi_7(\mathbf{x}) = (x^2, 0)^T$, $\phi_8(\mathbf{x}) = (xy, 0)^T$, $\phi_9(\mathbf{x}) = (y^2, 0)^T$, $\phi_{10}(\mathbf{x}) = (0, x^2)^T$, $\phi_{11}(\mathbf{x}) = (0, xy)^T$, $\phi_{12}(\mathbf{x}) = (0, y^2)^T$.

The origin of the axis is often the center of the subset $\Omega_{\mathbf{x}}$, so that the displacement estimated at \mathbf{x} by minimizing the SSD criterion over $\Omega_{\mathbf{x}}$ is given by $\phi(\mathbf{x}) = \phi^{\mathbf{x}}(0, 0) = (\lambda_1, \lambda_2)^T$ for any order. Moreover, one can see that first- and second-order shape functions satisfy $\partial \phi^{\mathbf{x}} / \partial x(0, 0) = (\lambda_3, \lambda_4)^T$ and $\partial \phi^{\mathbf{x}} / \partial y(0, 0) = (\lambda_5, \lambda_6)^T$.

The displacement field is supposed to be smooth enough so that approximating \mathbf{u} by ϕ makes sense. In particular, no occlusions (an object hide the part of another object in stereoscopy) or missing parts (as a mechanical fracture) are allowed. Displacement is locally invertible if the mapping from \mathbf{x} to $\mathbf{x} + \phi(\mathbf{x})$ has a non-singular Jacobian matrix, thanks to the inverse function theorem. In the case of polynomial shape functions, the determinant of this matrix at the center of the subset is $(1 + \lambda_3)(1 + \lambda_6) - \lambda_4\lambda_5 = 1 + \lambda_3 + \lambda_6 + \lambda_3\lambda_6 - \lambda_4\lambda_5$. Of course, zero-order shape functions (corresponding to simple translations) give invertible displacement since, in this case, $\lambda_i = 0$ for $i \in \{3, 4, 5, 6\}$ and the determinant is non-zero. Higher-order shape functions also yield locally invertible displacement fields if the λ_i 's are small with respect to 1, which holds true under the classic small strain hypothesis in photomechanics.

1.2 Motivation

The motivation of our work is the observation that, in photomechanics, displacement fields are often impaired by spurious small fluctuations. These fluctuations are often believed, in this community, to be caused by sensor noise, interpolation bias, or the numerical scheme minimizing the DIC criterion. Nevertheless, recent works consider the random marking on the surface of the specimen [14, 24, 45], in addition to the aforementioned causes, the authors of [10] coining the term *pattern-induced bias*. From a different perspective, it is known that the local parametric estimation of a displacement field is intrinsically biased, the retrieved displacement being merely the convolution of the true displacement by a Savitzky-Golay filter [36]. In this latter approach, specimen marking does not play any role, which is contradictory with the presence of a pattern-induced bias. The main motivation of this study is thus to see if this contradiction can be removed.

1.3 Contributions

The goal of the present paper is to provide a theoretical basis for these different viewpoints and thus to show that they can coexist. More precisely, we shall give the relation between the retrieved Λ (thus, the retrieved displacement ϕ) and the actual unknown displacement \mathbf{u} over any subset $\Omega_{\mathbf{x}}$. By revisiting papers dealing with disparity estimation in stereo-imaging [1, 9, 31, 34], we show in Section 2 that the estimation of Λ is mainly the sum of two terms, namely a term depending on the unknown displacement \mathbf{u} and the approximation induced by the shape functions, and a term caused by the interpolation error. We propose a characterization of these terms, which actually depend on the gradient distribution over the imaged specimen. This unifies the points of view of stereoscopy (fattening effect) and photomechanics (pattern-induced bias). In Section 3, numerical experiments validate the theoretical developments and assess their limitations. To limit the size of the paper, we discuss how sensor noise propagates from the images to the estimated displacement field in a separate companion research report [42].

1.4 Related work

Stereoscopy and photomechanics share the same objective, namely estimating displacement fields very accurately. The proposed contribution echoes different papers from these research areas.

First, the proposed approach relies on the calculation presented in [31] (and, to some extent, in [9]) in the context of disparity estimation in stereo-imaging. In this research field, disparity plays the role of the displacement field considered in the present paper. On the one hand, disparity may be estimated at any pixel as a constant monodimensional displacement between small patches extracted from the stereo image pair. Indeed, disparity is collinear with a given direction, since images are rectified so that the epipolar lines are parallel to each other. In this context, the authors of [9, 31] give predictive formulas for quantifying the fattening effect, i.e., the bias in disparity estimation caused by image gradient distribution. On the other hand, in experimental mechanics, local non-constant displacements are estimated as linear combinations of shape functions and they are not constrained to be collinear with a given direction. The main difference with disparity is that displacement fields are usually smooth and have tiny fluctuations, the strain components (defined from the partial derivatives of the displacement) being below 10^{-2} for many materials and load intensity. Gradient distribution is also likely to affect displacement estimation, giving pattern-induced bias. A first attempt at characterizing this bias in photomechanics is available in [24].

In addition to pattern-induced bias or fattening effect, the interpolation scheme required for registration is certainly a source of error. It is a common assumption in stereo-imaging that the input images satisfy Shannon-Nyquist sampling conditions, cf. [9, 31, 34]. This involves that continuous images can be perfectly interpolated from the Fourier coefficients without any interpolation bias. However, it is mentioned in [29, 40] that aliasing, although hardly noticeable to the naked eye on the raw input images, may strongly affect the estimated displacement and strain fields in photomechanics. This motivates the use of bilinear or bicubic interpolation schemes in this field [43] instead of Fourier interpolation. The drawback is that $\tilde{\mathcal{I}}'$ matches \mathcal{I}' only at integer pixel coordinates: non-integer pixel coordinates are thus affected by interpolation error, as numerically illustrated in [6]. Seeing interpolation as a convolution filter, the authors of [35] have proposed a characterization of the interpolation-induced bias giving rise to the famous “S-shape function”. The authors of [38] reduce the interpolation bias by sampling the subset Ω at non-integer pixels.

As mentioned in the introduction, the most popular approach in experimental mechanics is certainly to locally approximate the displacement field by a linear combination of shape functions, which may in turn undermatch the true displacement. This is another source of error discussed in [36], where it is shown that the retrieved displacement field is the convolution of the true displacement with a Savitzky-Golay filter characterized by the order of the shape functions and the size of the subset. This characterization is used in [14] and [46] to define some metrological parameters, and in [16] to restore displacement fields

through a dedicated deconvolution procedure.

1.5 Notation and reminder

In what follows, $\langle \cdot, \cdot \rangle$ and $|\cdot|$ denote the bidimensional Euclidean product and norm, respectively. The gradient of any 2D function is denoted by ∇ . We identify any vector with the corresponding column matrix, in boldface letters. The transpose of any matrix A is denoted by A^T .

We denote by \mathcal{O} Landau’s “big-O” for a variable tending to 0. We recall that if \mathbf{f} and \mathbf{g} are 2D-valued functions, $\mathbf{f} = \mathcal{O}(|\mathbf{g}|)$ means that for some $K > 0$, any small enough \mathbf{x} satisfies $|\mathbf{f}(\mathbf{x})| \leq K|\mathbf{g}(\mathbf{x})|$. If $\mathbf{f} = \mathcal{O}(|\mathbf{g}|)$ and A is some 2-column matrix, then $A\mathbf{f} = \mathcal{O}(|\mathbf{g}|)$ and in particular $\langle \mathbf{h}, \mathbf{f} \rangle = \mathcal{O}(|\mathbf{g}|)$ for any constant 2-D vector \mathbf{h} .

In the remainder of this paper, we will make use of partial derivatives of interpolated 2-D functions. While some interpolation schemes, such as bilinear interpolation, do not provide us with derivatives at integer pixel coordinates, it should be noted that derivatives of the interpolated functions will be calculated at points like $\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)$, \mathbf{x}_i being an integer pixel coordinate. In most situations, $\boldsymbol{\phi}(\mathbf{x}_i)$ has a subpixel value, thus $\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)$ has a non-integer value and the derivatives are well defined.

2 Estimating displacements by minimizing SSD

This section gives a closed-form expression of the displacement $\boldsymbol{\phi}$ estimated by minimizing the SSD criterion (defined in Equation 2) as a function of the actual unknown displacement \mathbf{u} . The aim is to emphasize the role of systematic errors caused by image texture and interpolation scheme. Non-noisy images are considered in this section. Images affected by signal-dependent noise are investigated in a separate research report [42].

Since the present section deals with the minimization problem at a given pixel \mathbf{x} and in order to simplify notations, we do not write the index \mathbf{x} to remind that the subset is centered at a pixel \mathbf{x} , and we simply write Ω and $\boldsymbol{\phi}$ instead of $\Omega_{\mathbf{x}}$ and $\boldsymbol{\phi}^{\mathbf{x}}$.

2.1 Relation between retrieved and actual displacement fields

The goal here is to express $\boldsymbol{\Lambda}$ as a function of the images \mathcal{I} and \mathcal{I}' , of the interpolated continuous image $\tilde{\mathcal{I}}$, and of the unknown displacement field \mathbf{u} .

By definition, $\boldsymbol{\Lambda} = (\lambda_1, \dots, \lambda_N)^T$ is a stationary point of the SSD given by Equation 2. For any $j \in \{1, \dots, N\}$, taking the derivative with respect to λ_j thus gives:

$$\sum_{\mathbf{x}_i \in \Omega} \left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) \right) \cdot \left\langle \nabla \tilde{\mathcal{I}}(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)), \boldsymbol{\phi}_j(\mathbf{x}_i) \right\rangle = 0 \quad (3)$$

Let $\boldsymbol{\delta}$ be the difference between the unknown displacement field and the retrieved one, such that for any $\mathbf{x}_i \in \Omega$, $\boldsymbol{\delta}(\mathbf{x}_i) = \mathbf{u}(\mathbf{x}_i) - \boldsymbol{\phi}(\mathbf{x}_i)$. It quantifies the undermatching of shape

functions mentioned in [36].

First,

$$\begin{aligned} \mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) &= (\mathcal{I}(\mathbf{x}_i) - \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i))) \\ &\quad + \left(\mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) \right) \end{aligned} \quad (4)$$

By definition of \mathbf{u} ,

$$\mathcal{I}(\mathbf{x}_i) - \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) = \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) - \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) \quad (5)$$

A first-order Taylor series expansion allows writing:

$$\mathcal{I}(\mathbf{x}_i) - \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\delta}(\mathbf{x}_i) \rangle + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2) \quad (6)$$

Denoting by $\mathbf{D}\mathcal{I}' = \mathcal{I}' - \tilde{\mathcal{I}}'$ the interpolation error (equal to 0 at integer coordinates), we conclude from Equation 4 that

$$\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\delta}(\mathbf{x}_i) \rangle + \mathbf{D}\mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2) \quad (7)$$

Second, another Taylor series expansion gives:

$$\begin{aligned} \nabla \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) - \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) &= -\boldsymbol{\delta}(\mathbf{x}_i)^T \mathbf{H}\mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) \boldsymbol{\delta}(\mathbf{x}_i) + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^3) \\ &= \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2) \end{aligned} \quad (8)$$

where $\mathbf{H}\mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i))$ is the Hessian matrix of \mathcal{I}' at $\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)$.

Consequently,

$$\begin{aligned} \nabla \tilde{\mathcal{I}}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) &= \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) + \left(\nabla \tilde{\mathcal{I}}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) - \nabla \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) \right) \\ &\quad + \left(\nabla \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) - \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) \right) \\ &= \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) - \mathbf{D}\nabla \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2) \end{aligned} \quad (9)$$

where $\mathbf{D}\nabla \mathcal{I}' = \nabla \mathcal{I}' - \nabla \tilde{\mathcal{I}}'$ is the gradient interpolation error.

Plugging Equations 7 and 9 into Equation 3 gives:

$$\begin{aligned} \sum_{\mathbf{x}_i \in \Omega} \left(\langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\delta}(\mathbf{x}_i) \rangle + \mathbf{D}\mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2) \right) \\ \times \left\langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) - \mathbf{D}\nabla \mathcal{I}'(\mathbf{x}_i + \boldsymbol{\phi}(\mathbf{x}_i)) + \mathcal{O}(|\boldsymbol{\delta}(\mathbf{x}_i)|^2), \boldsymbol{\phi}_j(\mathbf{x}_i) \right\rangle = 0 \end{aligned} \quad (10)$$

For any $i \in \{1, \dots, M\}$ and $j \in \{1, \dots, N\}$, let $L_{i,j}^{\mathbf{u}} = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \phi_j(\mathbf{x}_i) \rangle$, such that $L^{\mathbf{u}}$ is a $M \times N$ matrix. We eventually obtain:

$$\begin{aligned} & \sum_{\mathbf{x}_i \in \Omega} \left(L_{i,j}^{\mathbf{u}} \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \delta(\mathbf{x}_i) \rangle + L_{i,j}^{\mathbf{u}} \mathbf{D}\mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i)) + \mathcal{O}(|\delta(\mathbf{x}_i)|^2) \right. \\ & \quad + \mathcal{O}(|\mathbf{D}\nabla \mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i))| \cdot |\delta(\mathbf{x}_i)|) + \mathcal{O}(|\mathbf{D}\mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i))| \cdot |\mathbf{D}\nabla \mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i))|) \\ & \quad \left. + \mathcal{O}(|\mathbf{D}\mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i))| \cdot |\delta(\mathbf{x}_i)|) \right) = 0 \quad (11) \end{aligned}$$

or in a simpler way:

$$\begin{aligned} & \sum_{\mathbf{x}_i \in \Omega} L_{i,j}^{\mathbf{u}} \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \delta(\mathbf{x}_i) \rangle + \sum_{\mathbf{x}_i \in \Omega} L_{i,j}^{\mathbf{u}} \mathbf{D}\mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i)) \\ & \quad + \mathcal{O}(|\delta|^2) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\mathbf{D}\nabla \mathcal{I}'|) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\delta|) + \mathcal{O}(|\mathbf{D}\nabla \mathcal{I}'| \cdot |\delta|) = 0 \quad (12) \end{aligned}$$

where $\mathcal{O}(|\mathbf{f}|)$ denotes $\sum_{\mathbf{x}_i \in \Omega} (\mathcal{O}|\mathbf{f}(\mathbf{x}_i)|)$.

Since $\delta(\mathbf{x}_i) = \mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i) = \mathbf{u}(\mathbf{x}_i) - \sum_{k=1}^N \lambda_k \phi_k(\mathbf{x}_i)$, the preceding equation gives:

$$\begin{aligned} & \sum_{\mathbf{x}_i \in \Omega} L_{i,j}^{\mathbf{u}} \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \mathbf{u}(\mathbf{x}_i) \rangle - \sum_{\mathbf{x}_i \in \Omega} L_{i,j}^{\mathbf{u}} \sum_{k=1}^N \lambda_k L_{i,k}^{\mathbf{u}} + \sum_{\mathbf{x}_i \in \Omega} L_{i,j}^{\mathbf{u}} \mathbf{D}\mathcal{I}'(\mathbf{x}_i + \phi(\mathbf{x}_i)) \\ & \quad + \mathcal{O}(|\delta|^2) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\mathbf{D}\nabla \mathcal{I}'|) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\delta|) + \mathcal{O}(|\mathbf{D}\nabla \mathcal{I}'| \cdot |\delta|) = 0 \quad (13) \end{aligned}$$

If we denote by \mathbf{G} the vector of components $\mathbf{G}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \mathbf{u}(\mathbf{x}_i) \rangle$ for any $i \in \{1, \dots, M\}$, we obtain with Equation 13 the following matrix relation:

$$(L^{\mathbf{u}})^T L^{\mathbf{u}} \mathbf{\Lambda} = (L^{\mathbf{u}})^T \mathbf{G} + (L^{\mathbf{u}})^T \mathbf{D}\mathcal{I}' + \mathcal{O}(|\delta|^2) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\mathbf{D}\nabla \mathcal{I}'|) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\delta|) \quad (14)$$

Matrix $(L^{\mathbf{u}})^T L^{\mathbf{u}}$ is invertible as soon as the columns of $L^{\mathbf{u}}$ are linearly independent (it is a Gramian matrix), which holds if the ϕ_j form a valid basis and if the gradient is not equal to zero over the whole subset Ω . We assume in the following that these assumptions hold.

We have finally demonstrated the following theorem.

Theorem 1 *If $\mathbf{\Lambda}$ minimizes the SSD criterion of Equation 2 registering image \mathcal{I}' over \mathcal{I} , then the following equality holds:*

$$\begin{aligned} \mathbf{\Lambda} = & ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{G} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D}\mathcal{I}' \\ & + \mathcal{O}(|\delta|^2) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\mathbf{D}\nabla \mathcal{I}'|) + \mathcal{O}(|\mathbf{D}\mathcal{I}'| \cdot |\delta|) \quad (15) \end{aligned}$$

where \mathbf{u} is the actual displacement field, and for any $i \in \{1, \dots, M\}$ and $j \in \{1, \dots, N\}$, the following notations hold: $L_{i,j}^{\mathbf{u}} = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \phi_j(\mathbf{x}_i) \rangle$, $\mathbf{G}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \mathbf{u}(\mathbf{x}_i) \rangle$, $\mathbf{D}\mathcal{I}' = \mathcal{I}' - \tilde{\mathcal{I}}'$, $\mathbf{D}\nabla \mathcal{I}' = \nabla \mathcal{I}' - \nabla \tilde{\mathcal{I}}'$, $\tilde{\mathcal{I}}'$ denotes the interpolation of \mathcal{I}' at non-integer coordinates, and $\delta = \mathbf{u} - \phi$.

With this presentation, it is not straightforward to see how the estimation of $\mathbf{\Lambda}$ is affected by a systematic error.

Let us now decompose (for instance, in the least-squares sense) the unknown, true displacement \mathbf{u} over the basis functions (ϕ_k) . We introduce the N -dimensional vector $\mathbf{\Lambda}^{\mathbf{u}} = (\lambda_1^{\mathbf{u}}, \dots, \lambda_N^{\mathbf{u}})^T$ such that $\mathbf{u} = \sum_{k=1}^N \lambda_k^{\mathbf{u}} \phi_k + \boldsymbol{\epsilon}$ where $\boldsymbol{\epsilon}$ is the residual in the decomposition of \mathbf{u} over the basis $(\phi_k)_{k \in \{1, \dots, N\}}$. Consequently, for any $i \in \{1, \dots, M\}$,

$$\mathbf{G}_i = \sum_{k=1}^N \lambda_k^{\mathbf{u}} \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \phi_k(\mathbf{x}_i) \rangle + \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\epsilon}(\mathbf{x}_i) \rangle \quad (16)$$

If we denote by \mathbf{E} the vector of components $\mathbf{E}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\epsilon}(\mathbf{x}_i) \rangle$, we obtain $\mathbf{G} = L^{\mathbf{u}} \mathbf{\Lambda} + \mathbf{E}$. Consequently, Theorem 1 gives the following corollary.

Corollary 1 *If $\mathbf{\Lambda}$ minimizes the SSD criterion of Equation 2 registering image \mathcal{I}' over \mathcal{I} , then the following equality holds:*

$$\begin{aligned} \mathbf{\Lambda} = \mathbf{\Lambda}^{\mathbf{u}} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{E} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D} \mathcal{I}' \\ + \mathcal{O}(|\boldsymbol{\delta}|^2) + \mathcal{O}(|\mathbf{D} \mathcal{I}'| \cdot |\mathbf{D} \nabla \mathcal{I}'|) + \mathcal{O}(|\mathbf{D} \mathcal{I}'| \cdot |\boldsymbol{\delta}|) \end{aligned} \quad (17)$$

with, in addition to the notations of Theorem 1, $\mathbf{E}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\epsilon}(\mathbf{x}_i) \rangle$ for any $i \in \{1, \dots, M\}$, where $\boldsymbol{\epsilon}$ is the component of \mathbf{u} out of the basis (ϕ_k) .

2.2 Discussion

Skipping second-order terms from Theorem 1 gives:

$$\mathbf{\Lambda} = ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{G} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D} \mathcal{I}' \quad (18)$$

and from Corollary 1:

$$\mathbf{\Lambda} = \mathbf{\Lambda}^{\mathbf{u}} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{E} + ((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D} \mathcal{I}' \quad (19)$$

It is easy to see that if interpolation is perfect ($\mathbf{D} \mathcal{I}' = \mathbf{0}$), and if the function basis is expressive enough to perfectly represent the true displacement \mathbf{u} ($\mathbf{E} = \mathbf{0}$) or at least if the degree of the polynomial basis is large enough ($\mathbf{E} \simeq \mathbf{0}$), then $\mathbf{\Lambda}$ is equal to the sought $\mathbf{\Lambda}^{\mathbf{u}}$.

These assumptions are, however, optimistic. They are discussed in the remainder of this section.

2.2.1 Interpolation bias

In Equation 18 the term $((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D} \mathcal{I}'$ quantifies the effect of subpixel interpolation error. Since the underlying image \mathcal{I}' is unknown, it is not possible to bound a priori

this error, except if some additional information is available. For instance, well-sampled images satisfying the Shannon-Nyquist condition give no interpolation error when using Fourier interpolation as in [31, 44] and in this case, $\mathbf{D}\mathcal{I}' = 0$. As discussed in the introduction, real images from photomechanical experiments probably do not satisfy these hypotheses.

It should be noted that most image interpolation schemes (such as Fourier, bilinear, bicubic, or Lanczos interpolations) are actually linear, in the sense that any interpolated value is a weighted mean of image values at integer pixel coordinates [13, 23]. For any $\mathbf{x} \in \mathbb{R}^2$, there exists a row matrix $P(\mathbf{x})$ of size M such that the interpolation of \mathcal{I}' at \mathbf{x} is given by $\tilde{\mathcal{I}}'(\mathbf{x}) = P(\mathbf{x})\mathcal{I}'$, where \mathcal{I}' momentarily denotes the matrix of image values reshaped as a column vector. For instance, row-vector $P(\mathbf{x})$ has four non-zero entries in bilinear interpolation, sixteen in bicubic or bicubic spline interpolation [13, 23], which are the most popular schemes in photomechanics.

The interpolation bias thus writes:

$$((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T (I - P(\mathbf{x})) \mathcal{I}' \quad (20)$$

It has been derived by another approach in [2] (refining [30, 43]), and experimentally assessed in [3]. It can be noted that the image gradient is involved in $L^{\mathbf{u}}$ and affects the interpolation bias. While interpolation error may be neglected (either by assuming the Shannon-Nyquist condition to be satisfied or by using high-order interpolation schemes), the next section deals with the so-called ‘‘undermatched subset shape functions’’ which potentially gives pattern-induced bias.

2.2.2 Undermatched shape functions and pattern-induced bias

In Equation 18, the term $((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{G}$ links $\mathbf{\Lambda}$ with the gradient of \mathcal{I}' and the actual displacement \mathbf{u} . Let us recall that, for any $i \in \{1, \dots, M\}$ and $j \in \{1, \dots, N\}$,

$$L_{i,j}^{\mathbf{u}} = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \phi_j(\mathbf{x}_i) \rangle \quad (21)$$

$$\mathbf{G}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \mathbf{u}(\mathbf{x}_i) \rangle. \quad (22)$$

and

$$\mathbf{E}_i = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \boldsymbol{\epsilon}(\mathbf{x}_i) \rangle. \quad (23)$$

We can see that we find again the well-known *aperture problem*: only the displacement component collinear with the image gradient plays a role in \mathbf{G}_i . Moreover, Theorem 1 (Corollary 1, respectively) shows that each component (the error on each component, respectively) of $\mathbf{\Lambda}$ is a weighted mean of \mathbf{G}_i (of the undermatching error \mathbf{E} , respectively), the weights being proportional to the squared components of the gradient along the shape functions (here, $((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1}$ acts as a normalization).

For didactic purposes, we explicit the relation in a simple, yet realistic case, where displacements is sought as a pure translation along the x -direction. This situation corresponds to disparity estimation in stereo-imaging with rectified images. We have $\mathbf{u}(\mathbf{x}_i) = (u(\mathbf{x}_i) \ 0)^T$, $N = 1$ and $\phi_1(\mathbf{x}) = (1 \ 0)^T$. In this case, $L^{\mathbf{u}} = (\mathcal{I}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)))_{1 \leq i \leq M}$ is a column-vector, and:

$$((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{G} = \frac{\sum_{i=1}^M (\mathcal{I}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)))^2 u(\mathbf{x}_i)}{\sum_{i=1}^M (\mathcal{I}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)))^2} \quad (24)$$

where \mathcal{I}'_x denotes the partial derivative of \mathcal{I}' along direction x .

In this case, the term related to the interpolation error simplifies into:

$$((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{D}\mathcal{I}' = \frac{\sum_{i=1}^M \mathcal{I}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)) \mathbf{D}\mathcal{I}'(\mathbf{x}_i)}{\sum_{i=1}^M (\mathcal{I}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)))^2} \quad (25)$$

When $\mathbf{u} = (u, 0)$ is a constant displacement, Equation 24 simplifies to u . This observation is consistent with Corollary 1 since in this case, the displacement is perfectly represented by the basis function ϕ_1 as, for any \mathbf{x} , $\mathbf{u}(\mathbf{x}) = u \phi_1(\mathbf{x})$.

It should be noted that Equation 24 is exactly the relation given in [1, 31]. When the sought displacement is not constant over Ω , it is a weighted mean of the actual displacement, the weights being the squared image derivatives. This justifies the well-known *fattening effect* in stereo-imaging (also called *adhesion effect* in [9]): points lying in the neighborhood of edges have a disparity essentially governed by the edge points, which have large gradient values. Fattening effects can be seen for example in [31, 34]: foreground objects (which have a larger disparity than the background) appear fatter than they are because background pixels near their edges inherit their disparity.

In experimental mechanics the relevant formula is given by Theorem 1. The estimated displacement field given by $\mathbf{\Lambda}$ is biased since it is a weighted sum of the actual displacement at the pixels belonging to the subset, the weights increasing with the squared gradient at the pixels (more precisely, with the squared component of the gradient collinear with the shape functions). In particular, the displacement returned at the center of the considered subset is essentially the average of the actual displacement taken at high-gradient points, even if the actual displacement is different at the considered point. Although this corresponds to the intuition, we are not aware of earlier papers giving a rigorous description of the phenomenon in the context of experimental mechanics. As a consequence, we keep on calling this phenomenon “pattern-induced bias” after [10], since the expression “fattening effect” does not seem to be adequate in experimental mechanics, nothing really becoming “fatter”.

Note that the relation between $\mathbf{\Lambda}$ and \mathbf{G} in Theorem 1 can be seen as a generalized convolution, with a spatially varying kernel.

2.3 From pattern-induced bias to Savitzky-Golay filtering

The authors of [36] claim that the estimated displacement is simply the convolution product between the true displacement and the Savitzky-Golay (SG) kernel (a low-pass filter [32, 26]), causing the so-called matching bias [14]. The SG filter only depends on the degree of the polynomial shape functions and on the size of the considered subset Ω [32]. In particular, it does not depend on the gradient of the underlying image. Moreover, the discussion of Section 2.2.2 concludes that the relation between the estimated displacement and the true one can be seen as a convolution with a spatially varying kernel. The claim of [36], backed by results from [14] or [16], therefore seems to contradict Theorem 1 and pattern-induced bias discussed in the preceding section. In the remainder of this section, we explain how these two viewpoints can be accommodated.

2.3.1 The case of stationary random patterns

To establish the relation with the SG filter, the ground hypothesis of [36] is that, with our notations, ϕ minimizes

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} |\mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i)|^2 \quad (26)$$

over a subset $\Omega_{\mathbf{x}}$ centered at pixel \mathbf{x} , for any \mathbf{x} . Let us write for a while the x - and y -components of ϕ as:

$$\phi_x(\mathbf{x}) = \sum_{0 \leq i+j \leq d} \alpha_{i,j} x^i y^j \quad (27)$$

$$\phi_y(\mathbf{x}) = \sum_{0 \leq i+j \leq d} \beta_{i,j} x^i y^j \quad (28)$$

where d is the maximum degree of the shape functions. Thus, $\lambda_1 = \alpha_{0,0}$, $\lambda_2 = \beta_{0,0}$, $\lambda_3 = \alpha_{1,0}$, $\lambda_4 = \beta_{1,0}$, $\lambda_5 = \alpha_{0,1}$, $\lambda_6 = \beta_{0,1}$, etc.

With $J_{\mathbf{x}} = (1, x, \dots, x^d, y, \dots, yx^{d-1}, y^2, \dots, y^2x^{d-2}, \dots, y^d)^T$, $\alpha = (\alpha_{0,0}, \alpha_{1,0}, \dots, \alpha_{d,0}, \alpha_{1,0}, \dots, \alpha_{1,d-1}, \alpha_{2,0}, \dots, \alpha_{2,d-2}, \dots, \alpha_{d,0})^T$, and $\beta = (\beta_{0,0}, \beta_{1,0}, \dots, \beta_{d,0}, \beta_{1,0}, \dots, \beta_{1,d-1}, \beta_{2,0}, \dots, \beta_{2,d-2}, \dots, \beta_{d,0})^T$, Equation 26 writes:

$$\sum_{\mathbf{x}_i} (u_1(\mathbf{x}_i) - J_{\mathbf{x}_i}^T \alpha)^2 + (u_2(\mathbf{x}_i) - J_{\mathbf{x}_i}^T \beta)^2 \quad (29)$$

Consequently, minimizing Equation 26 with respect to α and β amounts to solving the normal equations $J^T \alpha = u_1$ and $J^T \beta = u_2$ where matrix J collects all column vectors $J_{\mathbf{x}}$. In other words, $\alpha = (JJ^T)^{-1} J u_1$ and $\beta = (JJ^T)^{-1} J u_2$. Since the displacement estimated over the subset Ω is simply $(\lambda_1, \lambda_2) = (\alpha_{0,0}, \beta_{0,0})$, each component of the retrieved displacement is the convolution of the components of \mathbf{u} by a Savitzky-Golay [32] filter of order equal to the degree of the shape functions and of support given by the dimensions of Ω . With first-

and second-degree shape functions, the partial derivatives of the displacement field being given by $(\lambda_3, \lambda_4, \lambda_5, \lambda_6)$, they are also given by SG filters, as explained in [26, 32].

The justification that the sought ϕ minimizes Equation 26 is only based on heuristic arguments in [36]. We shall see that the very specific nature of the speckle patterns used in experimental mechanics permits to justify this claim.

Since $\mathcal{I}(\mathbf{x}_i) = \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i))$, if we neglect interpolation error and identify $\tilde{\mathcal{I}}'$ with \mathcal{I}' , we obtain the following first-order approximation as in Equation 6:

$$\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \phi(\mathbf{x}_i)) = \langle \nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i)), \mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i) \rangle \quad (30)$$

Provided this first-order approximation holds (for instance because an initial guess of the solution is available), minimizing the SSD criterion thus amounts to minimizing:

$$\sum_{\mathbf{x}_i \in \Omega} g_i^2 \cos^2(\theta_i) |\mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i)|^2 \quad (31)$$

with g_i the norm of $\nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i))$ and θ_i the angle between $\nabla \mathcal{I}'(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i))$ and the error $\mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i)$. As can be seen, this is not Equation 26. It can be noted that the component of $\mathbf{u} - \phi$ orthogonal to the gradient of \mathcal{I}' does not play any role, which is consistent with the preceding discussion about the aperture problem. This was also mentioned in [24].

As explained in the introduction, specimens tested in experimental mechanics are marked with random speckle patterns which can be modeled as stationary textures. We can thus safely assume that the gradient norms $(g_i)_{\mathbf{x}_i \in \Omega}$ are identically distributed random variables, as well as the angles $(\theta_i)_{\mathbf{x}_i \in \Omega}$, and that at each pixel \mathbf{x}_i , g_i and θ_i are independent. Nevertheless, these random variables are spatially correlated. A common assumption is that spatial correlation vanishes with the distance between pixels as in natural images [37].

Let $X_i = g_i^2 \cos^2(\theta_i) |\mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i)|^2$ for any $1 \leq i \leq M$. The classic law of large numbers does not hold here because of spatial correlations. However, generalizations such as Bernstein's weak law of large numbers [8, Ex. 254 p. 67] still hold. Assuming that the variance of the X_i is bounded, i.e., there exists $c > 0$ such that $\text{Var}(X_i) \leq c$, and assuming also that spatial correlations vanish with distance, i.e., $\text{Cov}(X_i, X_j) \rightarrow 0$ when $|\mathbf{x}_i - \mathbf{x}_j| \rightarrow +\infty$, we obtain¹:

$$\begin{aligned} \text{Var} \left(\sum_{i=1}^M X_i \right) &= \sum_{i=1}^M \text{Var}(X_i) + 2 \sum_{i=1}^M \sum_{j \in [1, M] \setminus V_i} \text{Cov}(X_i, X_j) \\ &\quad + 2 \sum_{i=1}^M \sum_{j \in [1, M] \cap V_i} \text{Cov}(X_i, X_j) \quad (32) \end{aligned}$$

¹The proposed calculation is adapted from <https://math.stackexchange.com/questions/245327/>

where, for any $\varepsilon > 0$, V_i is the set of indices j such that for any i and $j \in V_i$, $\text{Cov}(X_i, X_j) \leq \varepsilon$. Here, $B \setminus A$ denotes the relative complement of a set A with respect to B . Let N_i be the cardinality of $\mathbb{N} \setminus V_i$ which is a finite set since $\text{Cov}(X_i, X_j) \rightarrow 0$ when $|\mathbf{x}_i - \mathbf{x}_j| \rightarrow +\infty$, and $N = \max_i N_i$.

Since for any i, j , $\text{Cov}(X_i, X_j) \leq c$ by Cauchy-Schwartz inequality, the following upper bound holds:

$$\text{Var} \left(\sum_{i=1}^M X_i \right) \leq Mc + 2MNc + 2M^2\varepsilon \quad (33)$$

Thus,

$$\text{Var} \left(\frac{1}{M} \sum_{i=1}^M X_i \right) \leq c/M + 2Nc/M + 2\varepsilon \quad (34)$$

Chebyshev's inequality implies that the random variable $\frac{1}{M} \sum_{i=1}^M X_i - \frac{1}{M} \sum_{i=1}^M E(X_i)$ tends to 0 in probability as $M \rightarrow +\infty$.

In other words, this justifies that minimizing Equation 31 amounts to minimizing

$$\sum_{i=1}^M E(X_i) = E(g^2)E(\cos^2(\theta)) \sum_{\mathbf{x}_i \in \Omega} |\mathbf{u}(\mathbf{x}_i) - \phi(\mathbf{x}_i)|^2 \quad (35)$$

as soon as the size M of the domain Ω is “large enough”. Equation 34 states that, for a given M , the approximation is as tight as the bound c on the variance is small, or as the spatial correlation of the image gradients quickly vanishes, giving a small N .

2.3.2 Towards an optimum pattern with respect to pattern-induced bias?

The result of the preceding section can be interpreted as follows: minimizing Equation 31 to estimate ϕ is similar to minimizing Equation 26, provided that the size M of the subset Ω is large enough. From Equation 34, this is all the more valid as the speckle pattern is fine (giving quickly vanishing spatial correlations, thus a small N for a given ε) and as the variance of the X_i is small (giving a small c). Since $\text{Var}(X_i)$ is proportional to $\text{Var}(g_i^2 \cos^2(\theta_i)) = E(g_i^4)E(\cos^4(\theta_i)) - E^2(g_i^2)E^2(\cos^2(\theta_i)) = 3E(g_i^4)/8 - E^2(g_i^2)/4$ (assuming that the θ_i are uniformly distributed in the interval $(0, 2\pi)$, which is sound with an isotropic speckle pattern, we indeed obtain $E(\cos^2(\theta)) = 1/2$ and $E(\cos^4(\theta)) = 3/8$). As a consequence,

$$\text{Var}(g_i^2 \cos^2(\theta_i)) = 3\text{Var}(g_i^2)/8 + E^2(g_i^2)/8 \quad (36)$$

A fine pattern minimizing this quantity should have lower pattern-induced bias. One can see that a concentrated gradient distribution with a low gradient average value is of interest. However, such a speckle pattern is still to be designed.

2.3.3 Link with a fattening-free criterion in stereoscopic imaging

Disparity in stereo-imaging being a 1-D displacement, $\theta_i = 0$ in Equation 31 and $g_i = \widetilde{\mathcal{I}}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i))$ where $\widetilde{\mathcal{I}}'_x$ denotes the partial derivative of $\widetilde{\mathcal{I}}$ along the epipolar line.

This motivates the authors of [1] to estimate ϕ by minimizing the following weighted SSD criterion, denoted $\widetilde{\text{SSD}}$ in what follows:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \frac{\left(\mathcal{I}(\mathbf{x}_i) - \widetilde{\mathcal{I}}'(\mathbf{x}_i + \phi(x))\right)^2}{\max(|\widetilde{\mathcal{I}}'_x(\mathbf{x}_i + \mathbf{u}(\mathbf{x}_i))|^2, \kappa)} \quad (37)$$

where $\kappa > 0$ avoids divisions by zero.

In stereo-imaging, ϕ is sought as a constant displacement. The solution of $\widetilde{\text{SSD}}$ minimizes $\sum_i |\mathbf{u}(\mathbf{x}_i) - \phi|^2$ and does not depend on the image gradient. It is shown in [1] that estimating ϕ by minimizing $\widetilde{\text{SSD}}$ at any pixel gives an estimated displacement at \mathbf{x} which is the mean of the $\mathbf{u}(\mathbf{x}_i)$ over $\Omega_{\mathbf{x}}$. Interestingly, this is consistent with the preceding section, since a zero-order Savitzky-Golay filter is a simple moving average with a kernel constant over its domain.

3 Numerical assessment

The goal of this section is to provide the reader with illustrative didactic experiments, and to assess the validity of the predictive formulas given by Theorem 1 and Corollary 1. We also discuss to what extent PIB can be eliminated or decreased. In the proposed numerical assessments, the real displacement fields are known, and we are able to compare the estimated displacement to this ground truth. It should be noted that the numerical scheme actually used to minimize the SSD criterion over the subsets (Equation 2) (see, e.g., [28]) is not important here. Nevertheless, the stopping criterion must be set carefully so that the stationarity assumption, which is the ground of Section 2, is valid. In practice, we use the Broyden–Fletcher–Goldfarb–Shanno (BFGS) quasi-Newton method with a cubic line search implemented in MATLAB’s `fminunc` function.

Since the predictive formulas are based on the ideal, continuous images \mathcal{I} and \mathcal{I}' whose derivatives are required, we shall first consider in Section 3.1 images and deformation fields expressed as simple closed-form expressions. We also assess the effect of gray-level quantization. Nevertheless, the true derivatives of quantized images are, of course, unknown. Section 3.2 deals with synthetic speckle images which mimic real images used in experimental solid mechanics.

The numerical experiments proposed in this section can be reproduced with datasets and MATLAB codes available at the following URL:

<https://members.loria.fr/FSur/software/PIB/>

The interested reader can also easily modify the parameters and datasets for further investigations.

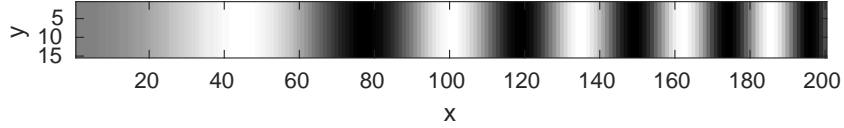


Figure 1: Synthetic image in deformed state defined by Equation 39 (image in reference state defined by Equation 38 is not discernible to the naked eye).

3.1 Image pairs given by a closed-form expression

In this section, we make use of images defined by closed-form expressions in order to assess the proposed formulas in a controlled experimental setting.

3.1.1 Displacement estimation

We define the image \mathcal{I} of the reference state and the image \mathcal{I}' of the deformed state at any pixel of coordinates (x, y) as sine waves by the following equations:

$$\mathcal{I}(x, y) = 2^{b-1}(1 + \gamma \sin(2\pi(x + u(x))^{1.5}/p^{1.5})) \quad (38)$$

$$\mathcal{I}'(x, y) = 2^{b-1}(1 + \gamma \sin(2\pi x^{1.5}/p^{1.5})) \quad (39)$$

where (x, y) spans a 200×15 pixel domain, $b = 8$ (so that the gray level of both images spans an 8-bit range), $\gamma = 0.9$ is the contrast, $p = 50$ pixels governs the varying period of the sine wave, and $u(x)$ is the ground-truth displacement field, supposed to be restricted along the x -axis. See Figure 1. Note that in this section, images are not quantized over b bits.

We seek for a constant displacement field $(\phi, 0)$ on each subset $\Omega_{\mathbf{x}}$ of size 15×15 pixels² distributed along the x -axis. The size of the subset is chosen in accordance with the characteristic scale of the modulated sine wave giving images \mathcal{I} and \mathcal{I}' . Such a constant displacement corresponds to zero-order shape functions. We therefore calculate at any abscissa x the quantity $\phi(x) = (\phi(x), 0)$ minimizing the SSD criterion:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \phi(x)) \right)^2. \quad (40)$$

The numerical assessment in the present section deals with a 1-D displacement along the x -axis in an image which varies only along the x -axis. The gradient is thus always collinear with the displacement. Consequently, the aperture problem manifests itself only because of vanishing image gradients: there is no loss of information orthogonally to the gradient as in the general 2-D case.

The 1-D displacement considered here is the case of interest of stereoscopy. It is possible to also implement the SSD criterion of [1], recalled in Section 2.3.3 above. We set the value

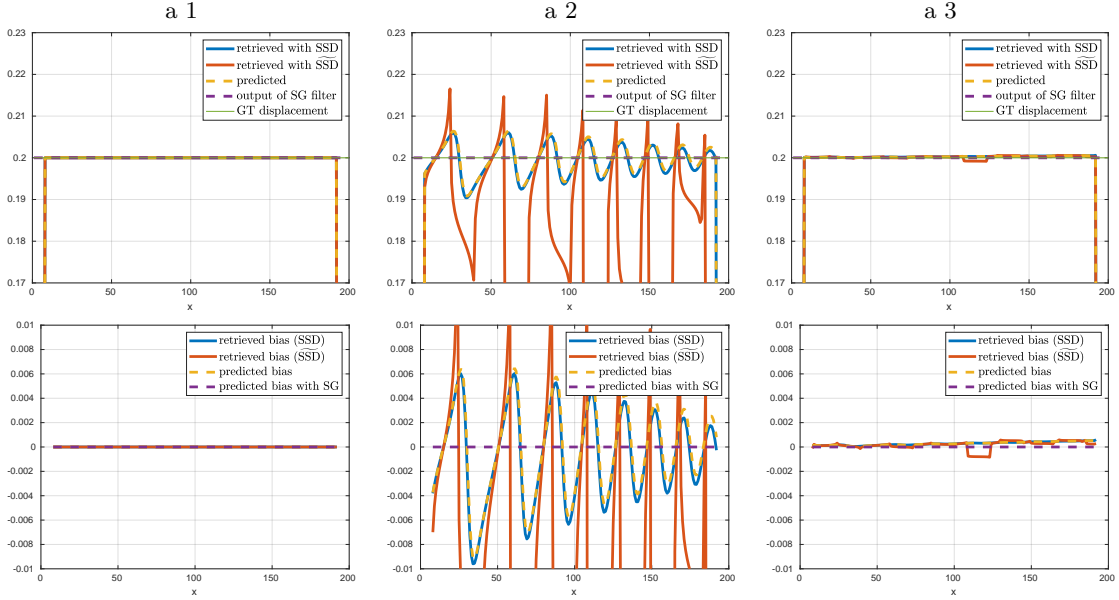


Figure 2: For displacement a) (constant displacement of 0.2 pixel): 1) closed-form expression (first column), 2) bilinear interpolation (second column), and 3) bicubic interpolation (third column). In each of these three cases, the first row depicts, superimposed on the same graph: the ground truth u , the displacement $\phi(x)$ estimated from SSD, the displacement $\widetilde{\phi}(x)$ estimated from \widetilde{SSD} , the displacement predicted by Theorem 1, and Savitzky-Golay filtering of u . In each of the three cases, the second row depicts, superimposed on the same graph : the differences (bias estimations) between u on the one hand, and displacement retrieved with SSD, predicted displacement, displacement retrieved with \widetilde{SSD} , output of the SG filter on the other hand.

of κ to achieve the best trade-off between reducing the fattening effect and numerical stability, see Section 3.2.5.

In the SSD and \widetilde{SSD} criteria, $\widetilde{\mathcal{I}}$ denotes a continuous image. Since we use images given by closed-form expressions, image values at non-integer pixels are available. Real experiments require, however, to interpolate images. In this illustrative experiment, we consider three possibilities: 1) using Equation 39 which allows avoiding any interpolation scheme (in this case we also use the closed-form expression of the gradient in the quasi-Newton scheme), 2) using bilinear interpolation, or 3) using bicubic interpolation.

We also consider two displacement fields: a) a constant $u(x) = 0.2$ pixel, b) a low-frequency sine wave $u(x) = 0.2 \sin(2\pi x/q)$ where $q = 90$ pixels. We also discuss in [42] the case of a high-frequency sine wave (not shown here). In all cases, the largest displacement

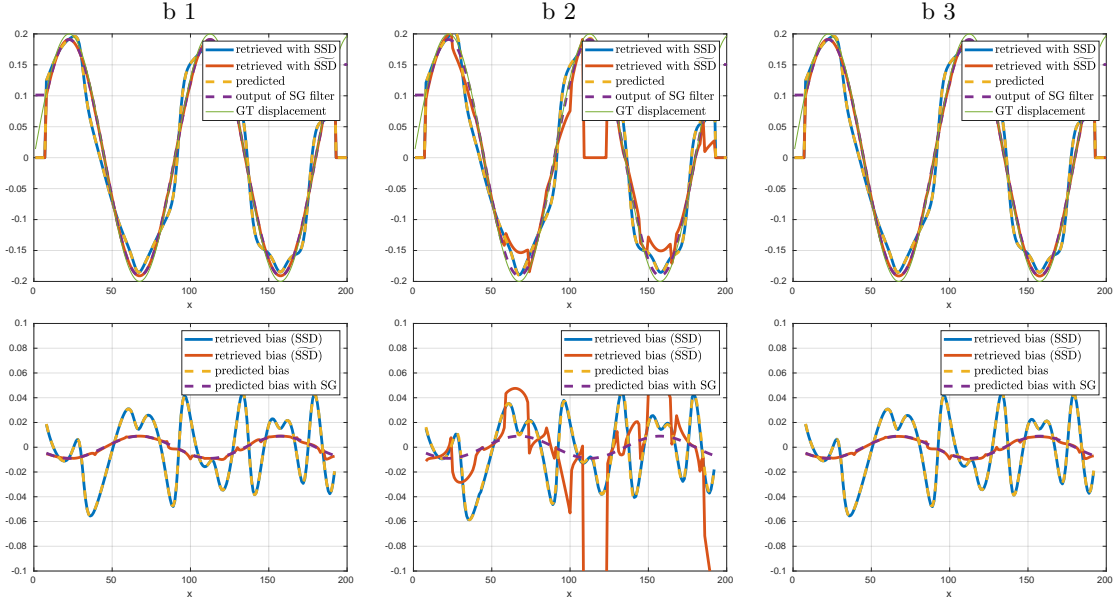


Figure 3: For displacement b) (sine wave of amplitude 0.2 pixel and period 90 pixels): 1) closed-form expression (first column), 2) bilinear interpolation (second column), and 3) bicubic interpolation (third column). In each of these three cases, the first row depicts, superimposed on the same graph: the ground truth u , the displacement $\phi(x)$ estimated from SSD, the displacement $\tilde{\phi}(x)$ estimated from \tilde{SSD} , the displacement predicted by Theorem 1, and Savitzky-Golay filtering of u . In each of the three cases, the second row depicts, superimposed on the same graph: the differences (bias estimations) between u on the one hand, and displacement retrieved with SSD, predicted displacement, displacement retrieved with \tilde{SSD} , output of the SG filter on the other hand.

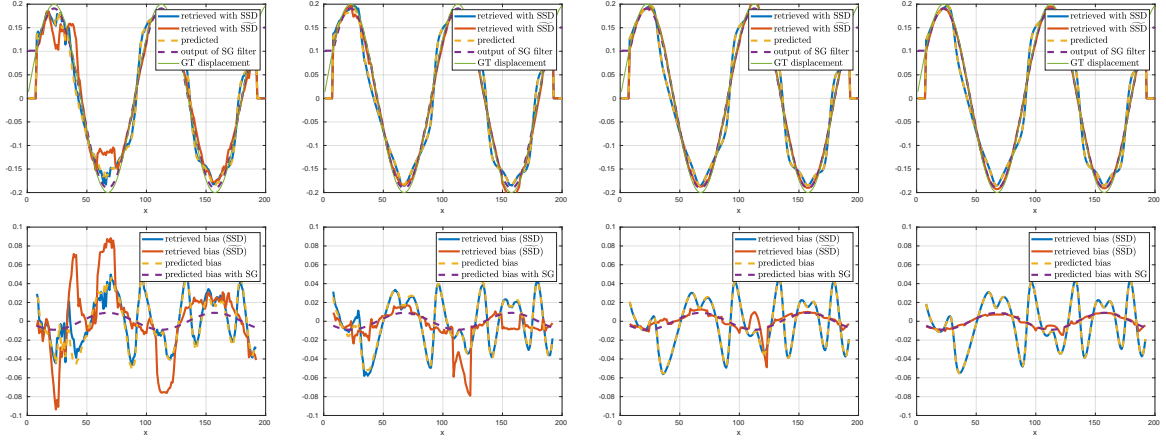


Figure 4: Effect of gray-level quantization. From left to right: 8-bit images, 10-bit images, 12-bit images, 14-bit images. First row: displacements. Second row: biases (difference between the predicted or retrieved displacements and the ground truth displacement).

value is 0.2 pixel.

Figures 2 and 3 show various plots. Each of these figures permits discussing displacement fields on the top and biases (systematic errors) on the bottom, as a function of the interpolation scheme (from left to right: closed-form expression, bilinear interpolation, and bicubic interpolation). Concerning the displacement, we plot the ground truth u (thin green line), the displacement $\phi(x)$ retrieved by minimizing the SSD criterion (blue), the displacement $\tilde{\phi}(x)$ retrieved from the modified criterion $\widetilde{\text{SSD}}$ defined in Equation 37 (red), the predicted displacement (cf. Theorem 1, yellow) and the filtering of the ground-truth displacement by the Savitzky-Golay filter of order 0 and frame length 15 (purple). Concerning the biases, we plot the differences between the ground truth displacement on the one hand, and the retrieved displacement with SSD, the predicted displacement, the retrieved displacement with $\widetilde{\text{SSD}}$, and the output of the SG filter on the other hand.

Concerning the constant displacement a) in Figure 2, since it can be represented with zero-order shape functions, the error ϵ (thus \mathbf{E}) in Corollary 1 is null: no marking bias should be noticed. With the closed-form expression of image \mathcal{I}' (case 1), the interpolation error $\mathbf{D}\mathcal{I}'$ is also null. It can be seen, indeed, that all curves are superimposed in this case. With bilinear interpolation (case 2), Theorem 1 predicts an error term caused by the interpolation error. We can see that the retrieved displacement fits well the prediction, the blue and yellow curves being superimposed. We can also see that correcting the marking bias (which is, here, non-existent) significantly amplifies the interpolation error, giving the erratic red curve. When looking closely at the curves showing the biases, we can notice a small difference between the retrieved bias and the predicted bias, which could probably be

explained by higher-order error terms than the first-order terms of the present calculation. With bicubic interpolation (case 3), the interpolation error is very small, all the curves being close to each other. With the constant displacement considered here, we can see that interpolation still causes a very small drift in the estimation, giving the increasing bias. Correcting the marking bias with $\widetilde{\text{SSD}}$ does not give an erratic red curve but a few spurious estimations can be seen. Interpolation errors seem to be amplified.

Concerning the low-frequency sine displacement b) in Figure 3, it cannot be represented by zero-order shape functions, thus marking bias should affect the retrieved displacement, the error term \mathbf{E} being non-null in Theorem 1. We can see that with closed-form expression (case 1), thus no interpolation error, the retrieved displacement fits perfectly the prediction, the blue and yellow curves being superimposed. The marking bias causes departures from the ground truth displacement, whose amplitude is governed by the gradient of the underlying image, giving a rather chaotic bias curve. The marking bias is perfectly removed with the $\widetilde{\text{SSD}}$ criterion: the retrieved displacement indeed fits the output of the SG filter, as predicted by the theory. With bilinear interpolation (case 2), an additional interpolation error is predicted by Theorem 1, although it is difficult to see a difference between the blue and yellow curves. The retrieved displacement globally fits the prediction. With bicubic interpolation (case 3), the curves are close to case 1. When using the $\widetilde{\text{SSD}}$ criterion, we can see that cases 1 and 3 fit well the predicted Savitzky-Golay filtering (green and red curves are superimposed), but as for displacement a), bilinear interpolation error (case 2) is amplified, giving an erratic red curve.

3.1.2 Effect of quantization

In the preceding section, image intensity is not quantized. We now perform experiments with quantized images in order to illustrate its impact on the predictive formulas. Bicubic interpolation is used in order to minimize interpolation bias, as illustrated earlier. We perform the same experiment as the one described in Figure 3.

Figure 4 shows the result of displacement estimation by minimizing SSD and $\widetilde{\text{SSD}}$. The former estimation should match the predicted one, and the latter should match the output of SG filter. We can see that the proposed predictive formulas are quite accurate as soon as quantization is performed over 10 bits (the yellow and blue curves are superimposed), and that retrieving the output of the SG filter with $\widetilde{\text{SSD}}$ requires to quantize image intensity over 12 bits (so that red and purple curves are superimposed).

3.2 Speckle patterns: assessing the prediction for pattern-induced bias

The previous section shows numerical assessments with smooth images and simple ground-truth displacement fields. The present section presents numerical assessments based on speckle pattern images corresponding to typical usage cases in experimental mechanics.

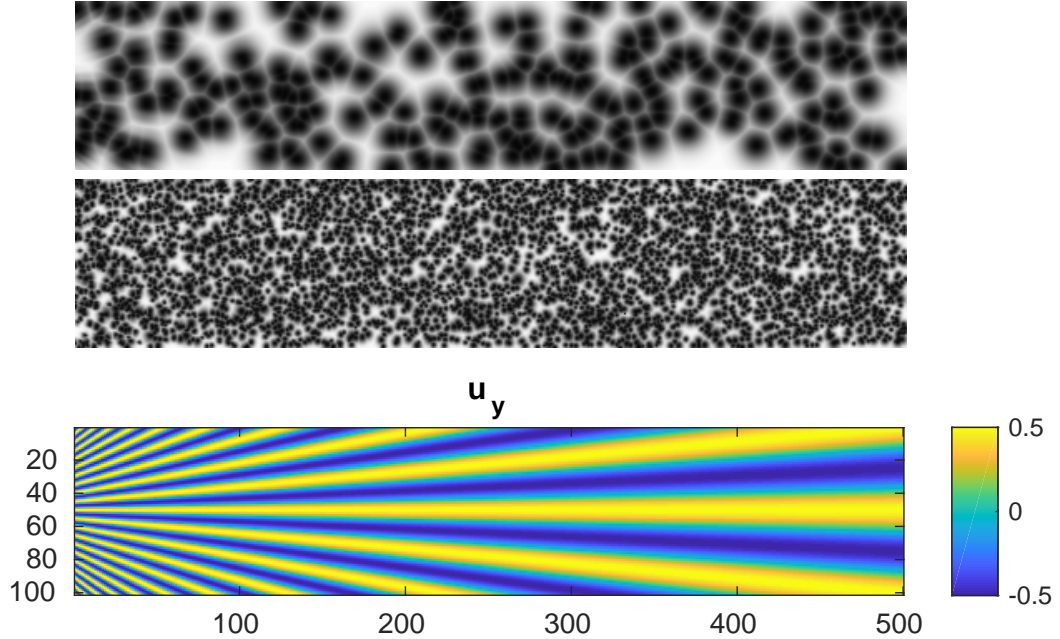


Figure 5: Two speckle images (size: 500×100 pixels) and y -component u_y of the prescribed displacement field. There is no displacement along the x -direction ($u_x = 0$). The first image shows the large speckle pattern discussed in Section 3.2.2, the second one the fine speckle pattern of Section 3.2.3.

BSpeckleRender software² [41] is used to render synthetic speckle images of size 500×100 pixels: one image corresponds to a reference state and another one corresponds to a deformed state. Speckles are designed to mimic real patterns such as the ones that can be seen in [22, 38, 39] for instance. The displacement field $\mathbf{u} = (u_x, u_y)$ is given by the following closed formula:

$$\begin{cases} u_x = 0 \\ u_y = 0.5 \cos\left(2\pi \frac{y-50}{5+45x/500}\right) \end{cases} \quad (41)$$

The speckle pattern is deformed along the y -direction, the prescribed displacement being a sine wave along the x -direction of maximum amplitude of 0.5 pixel, whose period ranges from 5 pixels ($x = 0$) to 50 pixels ($x = 500$). Such a deformation field is relevant in order to

²We use BSpeckleRender.b which renders speckle images with patterns of intensity at pixel \mathbf{x} varying as $\exp(-4|\mathbf{x}-\mathbf{x}_0|^2/R^2)$ with a center \mathbf{x}_0 given by a Poisson point process and a random radius R , instead of random black disks over a white background, so that the image gradient is a smooth function. MATLAB software code is available at the following URL: <https://members.loria.fr/FSur/software/BSpeckleRender/>

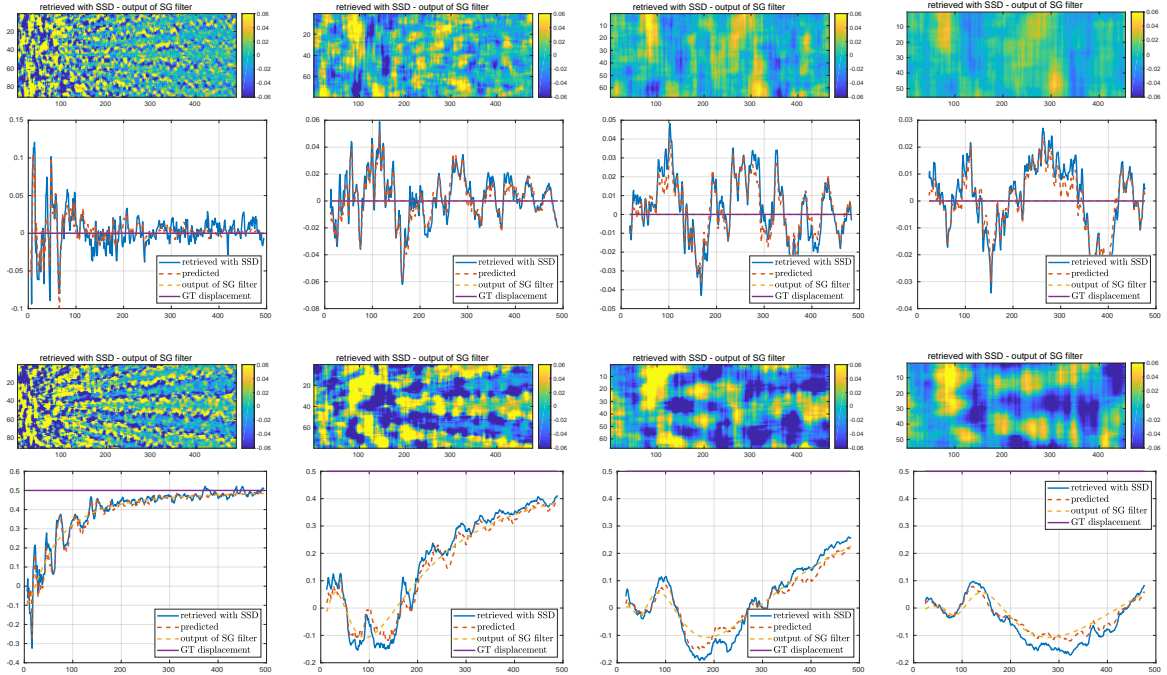


Figure 6: Fine speckle, zero-order shape functions. From left to right: 7×7 , 19×19 , 31×31 , and 43×43 subset Ω . On the top, displacement following x -axis; on the bottom, displacement following y -axis (difference between retrieved displacement field and output of SG filter, and cross-section along the line $y = 50$).

highlight the frequency response of the SG filter, see [10, 16, 17, 18] for instance. A 12-bit quantization is used, following the prescription of Section 3.1.2. Note that these images are smaller than in assessment datasets used in some recent works [14, 15, 17, 18] because of the time needed to render 12-bit images, and in order to facilitate the reading of the graphs. Because of the discrete nature of the rendered speckle images, interpolation is also required. Bicubic interpolation is used, following Section 3.1.1. Consequently, estimation error is likely to be only caused by the $((L^{\mathbf{u}})^T L^{\mathbf{u}})^{-1} (L^{\mathbf{u}})^T \mathbf{E}$ term in Corollary 1.

Figure 5 shows two speckle images and the imposed (ground truth) displacement field. Images (before and after deformation) with the large speckle pattern are discussed in Section 3.2.2. Section 3.2.3 concerns images with the fine speckle pattern. The goal here is to compare the displacement field retrieved with SSD minimization to its counterpart given by the predictive formula for the pattern-induced bias (PIB). For didactic purpose, several values are tested for both the subset size and the order of the shape functions. In the remainder of this section, we show the displacement maps (in each direction) retrieved by SSD minimization, the maps predicted by Theorem 1, the ground-truth (GT) displace-

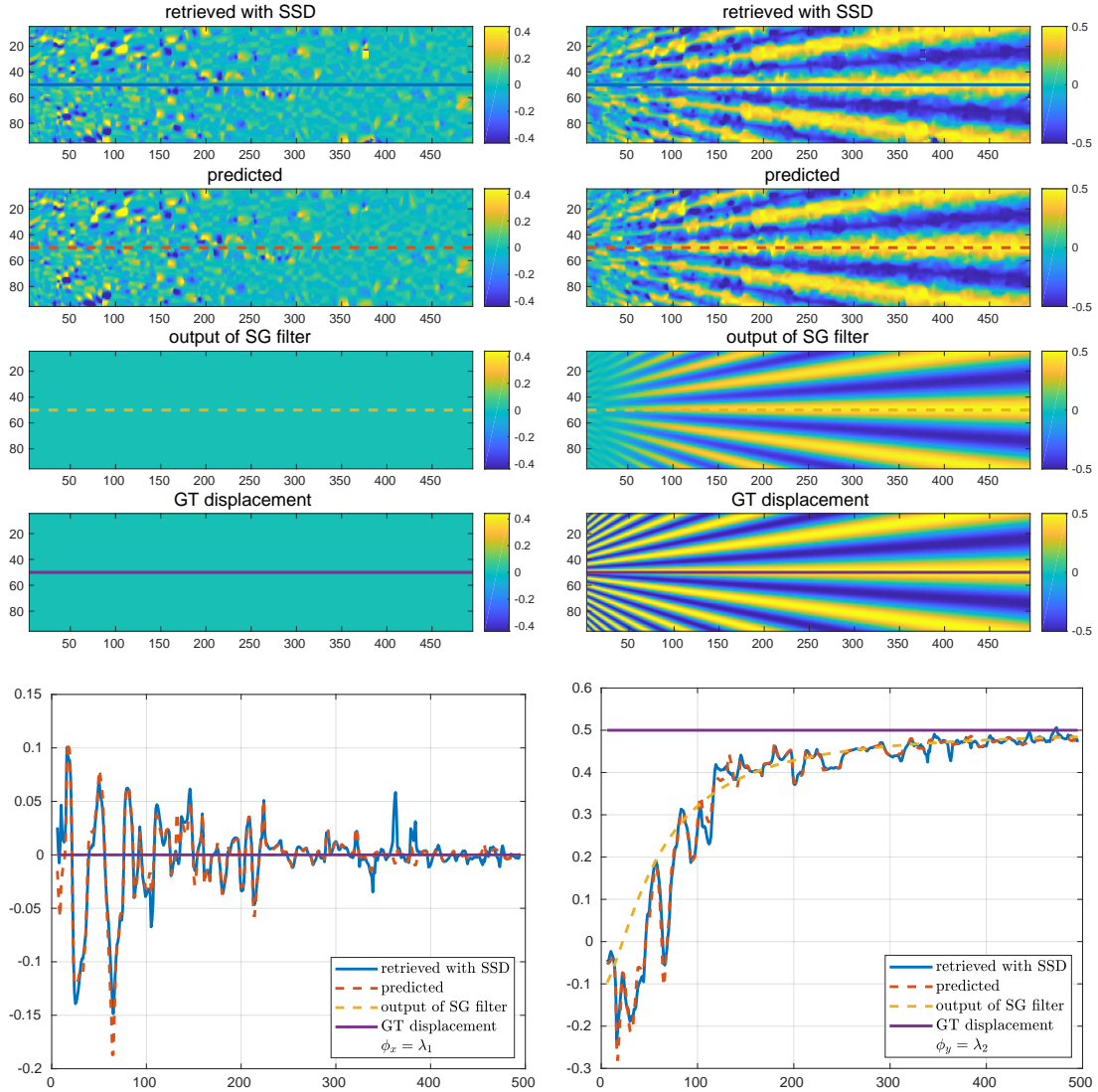


Figure 7: Large speckle, 7×7 subset Ω and 0-order shape functions. Displacement along x - (left) and y -axis (right), and corresponding cross-section plots.

ment, and the output of the Savitzky-Golay filtering of the GT displacement. We also show a cross-section plot of the displacement field along its middle-line $y = 50$ where GT displacement is constant, either null along the x -direction, or equal to 0.5 pixel along the y -direction.

3.2.1 Influence of the subset size

Figure 6 shows the evolution of the difference between the retrieved displacement and the output of the Savitzky-Golay filter as a function of the subset size, for zero-order shape functions used in stereoscopy. The results (not shown) for first-order shape functions are similar. The fine speckle pattern images are used here. As discussed in the following sections, the predictive formulas for the retrieved displacement are quite accurate. These figures illustrate the convergence of the retrieved displacement towards the output of the SG filter, as predicted in Section 2.3. It also illustrates that PIB shows a large amplitude over areas where the gradient of the displacement has a large value (here, on the left of the u_y displacement field). We can also see that large subsets give very smooth displacement fields, the corresponding lowpass SG filter having the same support as the subset (see [14, 36]). Such large subsets are not used in practice. Small subsets are required to avoid a large Savitzky-Golay smoothing, but they also give a potentially large PIB.

3.2.2 Large speckle pattern

Figure 7 shows the results for a subset Ω of size 7×7 and zero-order shape functions (that is, a constant displacement is estimated over each subset). We can see that the seemingly random fluctuations in the retrieved displacement fields are actually caused by PIB (thus not by sensor noise) and are well predicted by our formulas. In particular, the cross-section plots of the displacement maps confirm that the seemingly random fluctuations of the estimated displacement along the SG filtering of the GT displacement are caused by PIB, the retrieved and predicted curves being superimposed. Interestingly, despite the null GT displacement along the x -direction, the x -component of the retrieved displacement is still affected by PIB. We can see that the amplitude of the PIB is quite large compared to the true displacement. It can be noticed that PIB is large where displacement gradient is large (that is, on the left-hand side of the displacement field). Since the retrieved displacement is a weighted mean of the true displacement in the considered subset, PIB is indeed likely to be larger if the displacement strongly varies within the subset.

Figure 8 shows results for 13×13 subsets and first-order shape functions. It turns out that higher-order shape functions require larger subsets, since numerical issues affect smaller subsets (not shown here). With first-order shape functions, $\mathbf{\Lambda}$ has 6 components. The first two components correspond to the displacement $\phi(x)$ at the center of the subset $\Omega_{\mathbf{x}}$ and can thus be compared to the convolution of the true displacement with SG filter. The four other components correspond to partial derivatives of the displacement

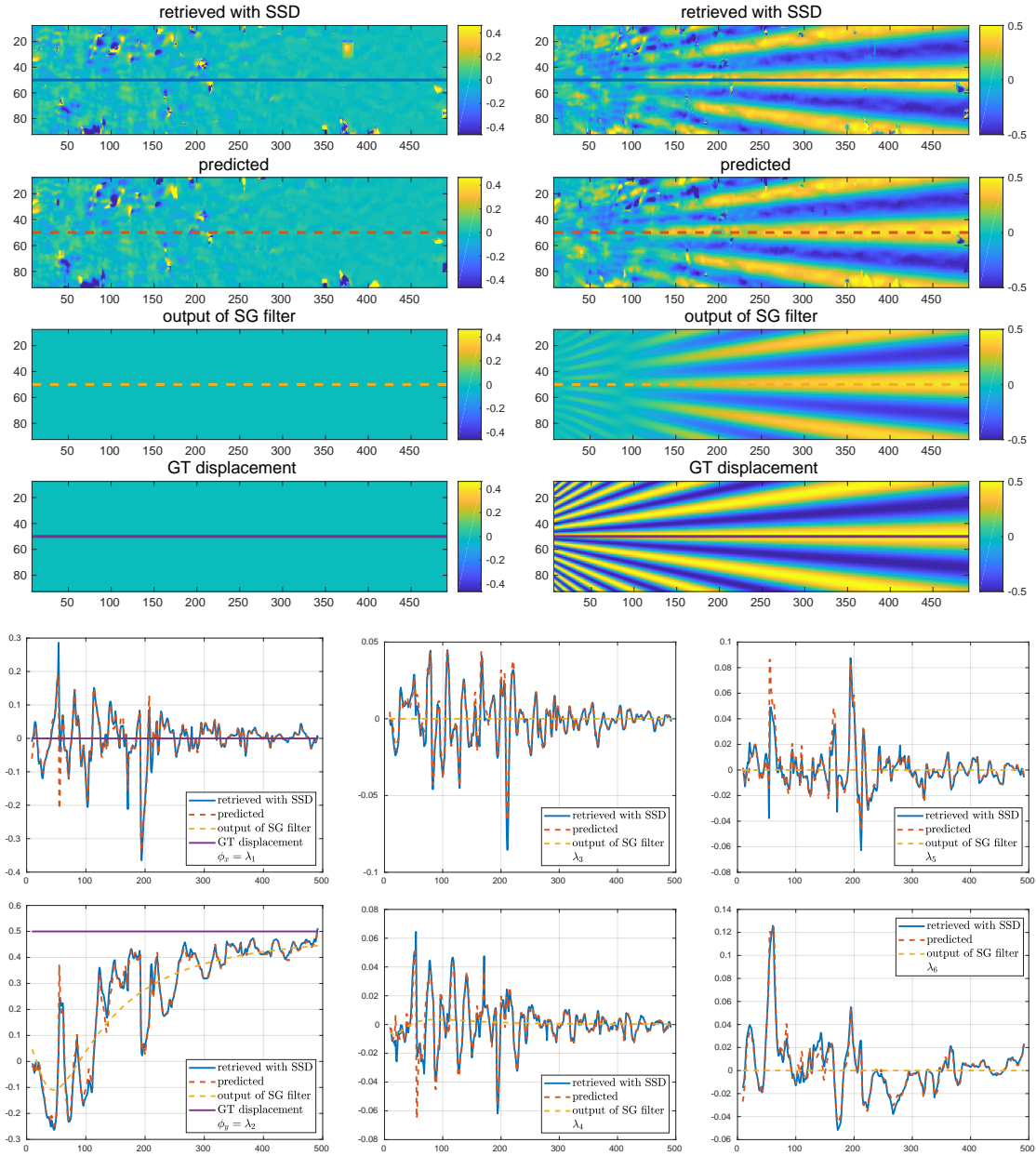


Figure 8: Large speckle, 13×13 subset Ω and first-order shape functions. The cross-section plots show, from left to right and top to bottom, λ_1 , λ_3 , λ_5 , λ_2 , λ_4 , λ_6 , such that the x -component of the displacement over a subset (giving fields shown on the upper left) is $\lambda_1 + \lambda_3x + \lambda_5y$ and the y -component (on the upper right) is $\lambda_2 + \lambda_4x + \lambda_6y$.

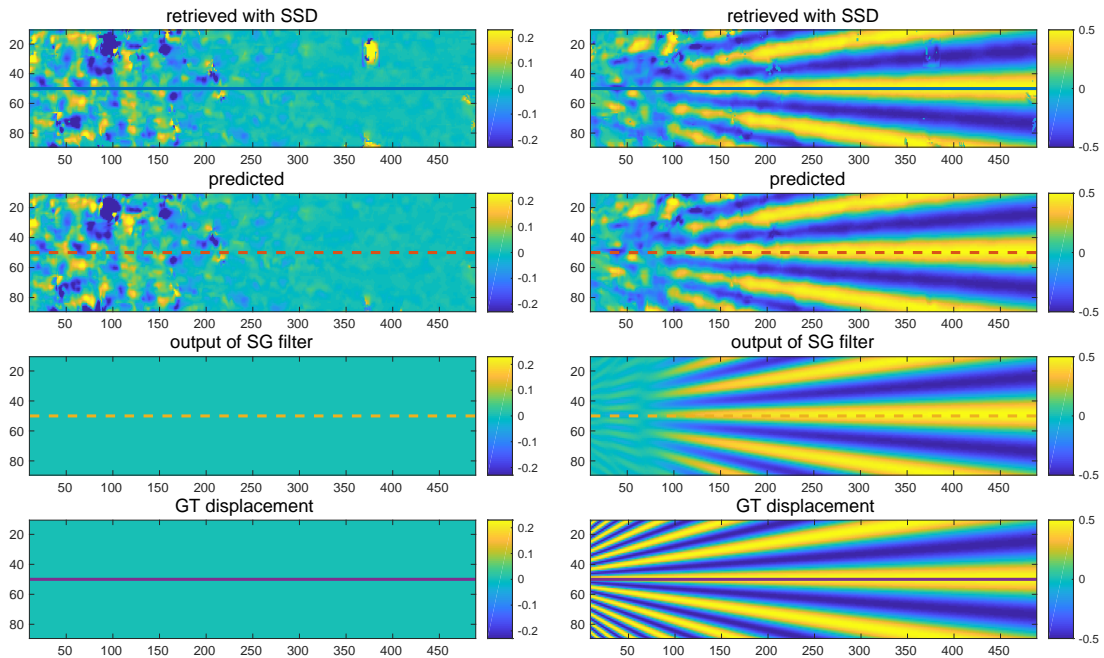


Figure 9: Large speckle, 19×19 subset Ω and second-order shape functions. Displacement along x - (left) and y -axis (right).

field. They also correspond to the output of (other) SG filters, as recalled in Section 2.3³. In all cases, we can see fluctuations of the estimated displacement around the output of the SG filter, these fluctuations being due to PIB. Because of the transfer function of the SG filter, the retrieved displacement field vanishes for certain values. It has even a wrong sign, as discussed in detail in [14]. It is quite surprising that the PIB give such large spurious measurements.

Figures 9 and 10 show results for 19×19 subsets and second-order shape functions. Even if second-order shape functions are rarely used in commercial DIC software programs for experimental mechanics, it can be seen that predictive formulas are still valid in this case.

3.2.3 Fine speckle pattern

Figures 11 to 14 show the results of the same experiments as in the preceding section, but with the fine speckle pattern. Although such a pattern makes it difficult to reliably

³It should be noted that these four components are the derivatives of the local displacement field ϕ^x estimated over Ω_x . In photomechanics, the derivatives of the displacement field, which are related to strain components, are rather computed from the derivatives of the global displacement ϕ .

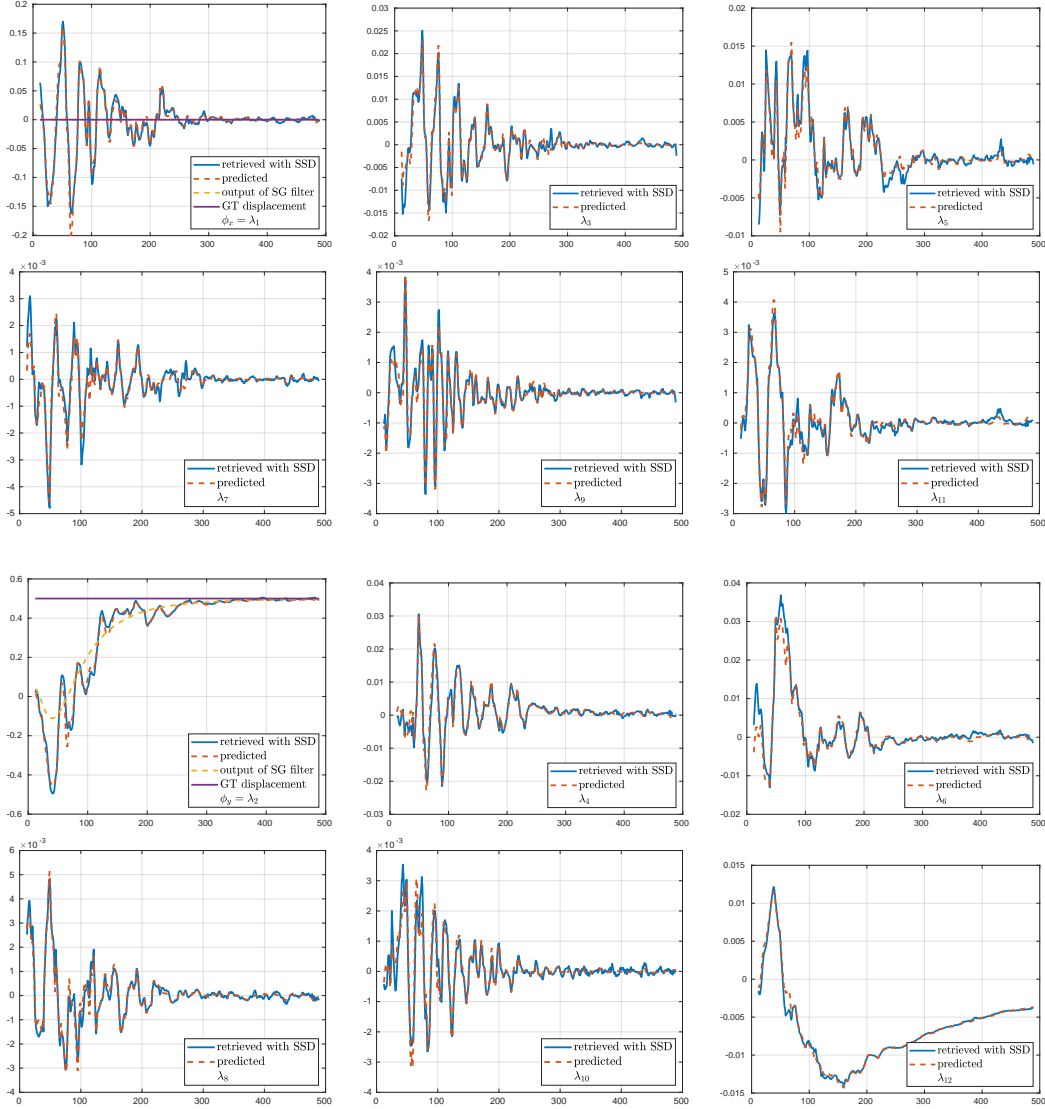


Figure 10: Large speckle, 19×19 subset Ω and second-order shape functions. From left to right and top to bottom: $\lambda_1, \lambda_3, \lambda_5, \lambda_7, \lambda_9, \lambda_{11}$ concerning x -displacement, and $\lambda_2, \lambda_4, \lambda_6, \lambda_8, \lambda_{10}, \lambda_{12}$ concerning y -displacement.

estimate the image gradients, we can see that PIB is still accurately predicted. With the same shape functions and subset sizes, we can see that the fine speckle pattern gives a PIB with a smaller amplitude, as discussed in Section 2.3.

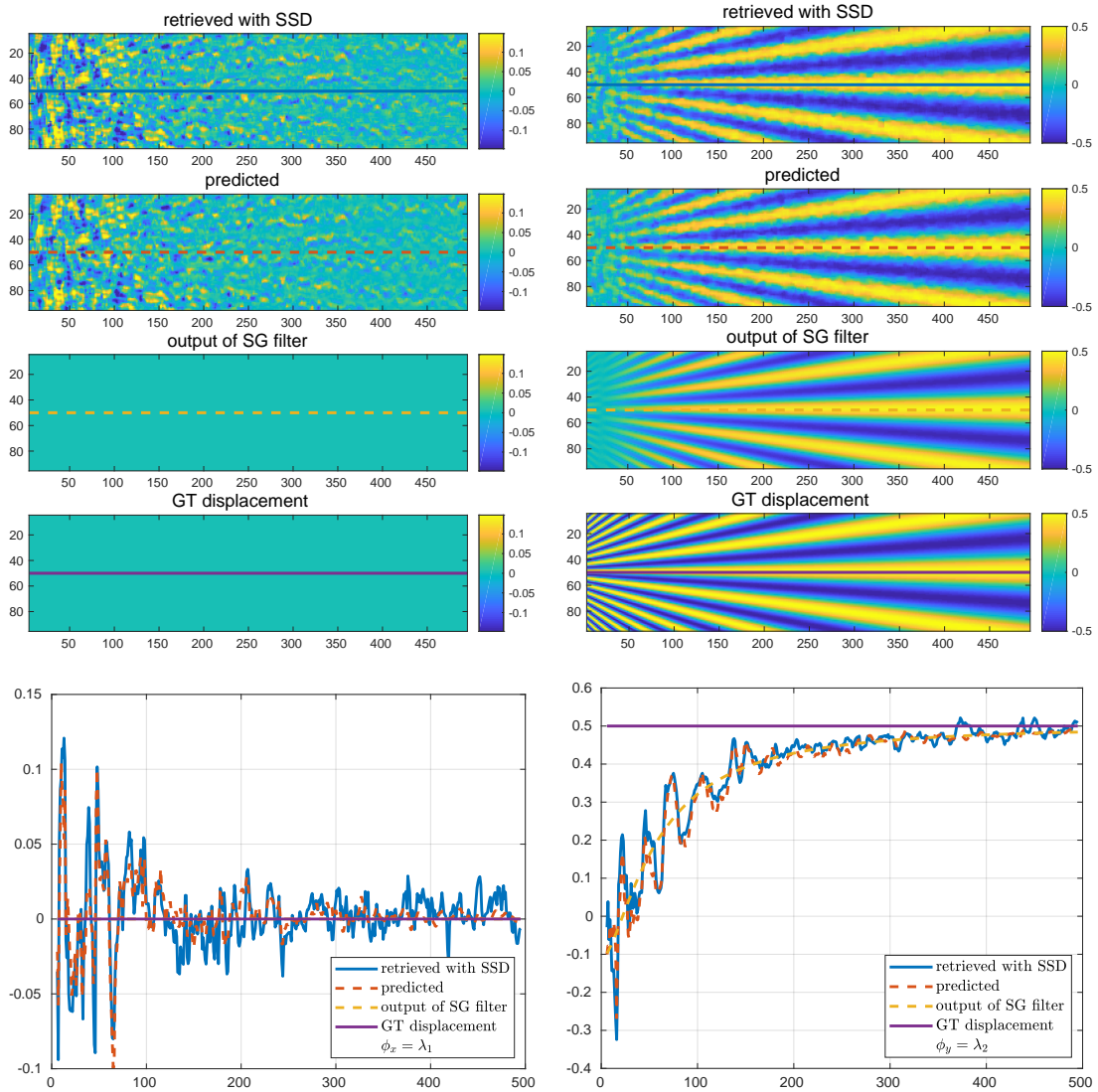


Figure 11: Fine speckle, 7×7 subset Ω and 0-order shape functions. Displacement maps along x - (left) and y -axis (right), and cross-section plots.

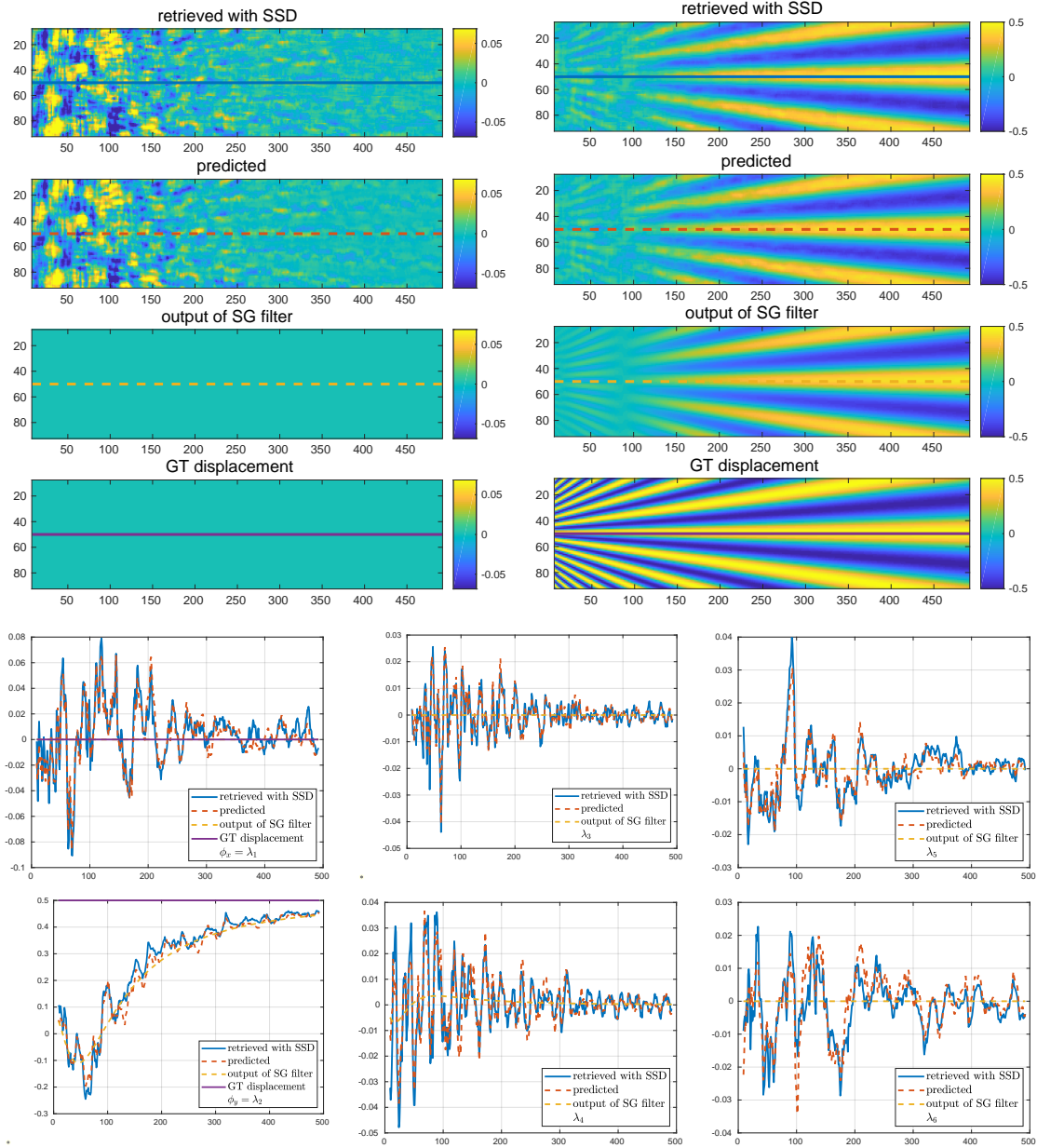


Figure 12: Fine speckle, 13×13 subset Ω and first-order shape functions. Displacement maps along x - (left) and y -axis (right). Cross-section plots, from left to right and top to bottom: λ_1 , λ_3 , λ_5 , λ_2 , λ_4 , λ_6 .

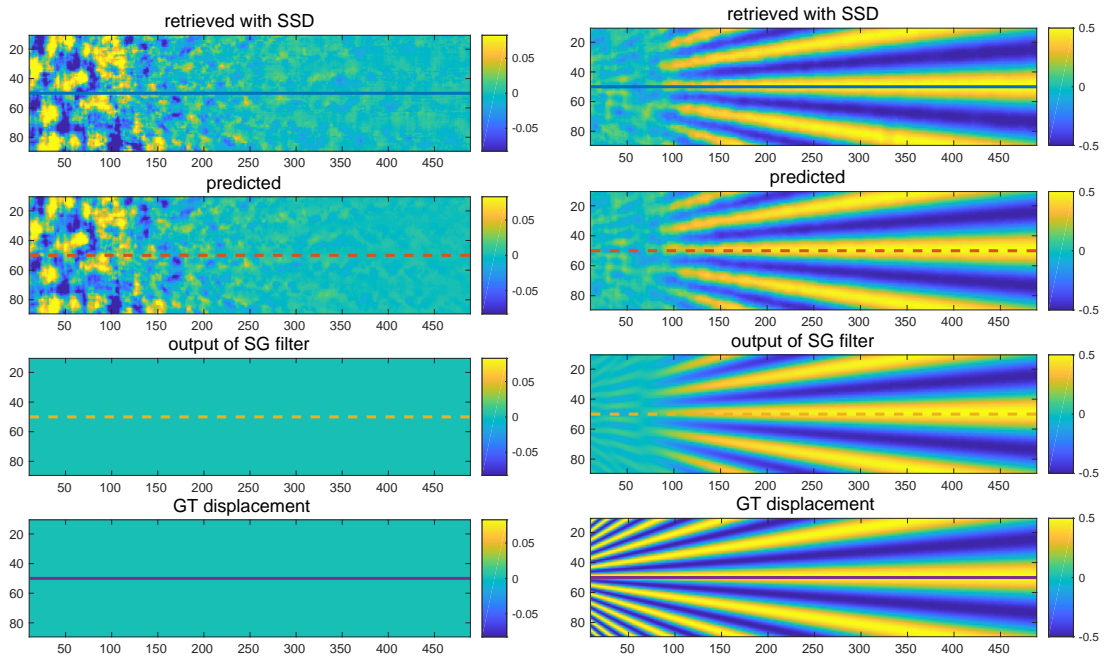


Figure 13: Fine speckle, 19×19 subset Ω and second-order shape functions. Displacement maps along x - (left) and y -axis (right).

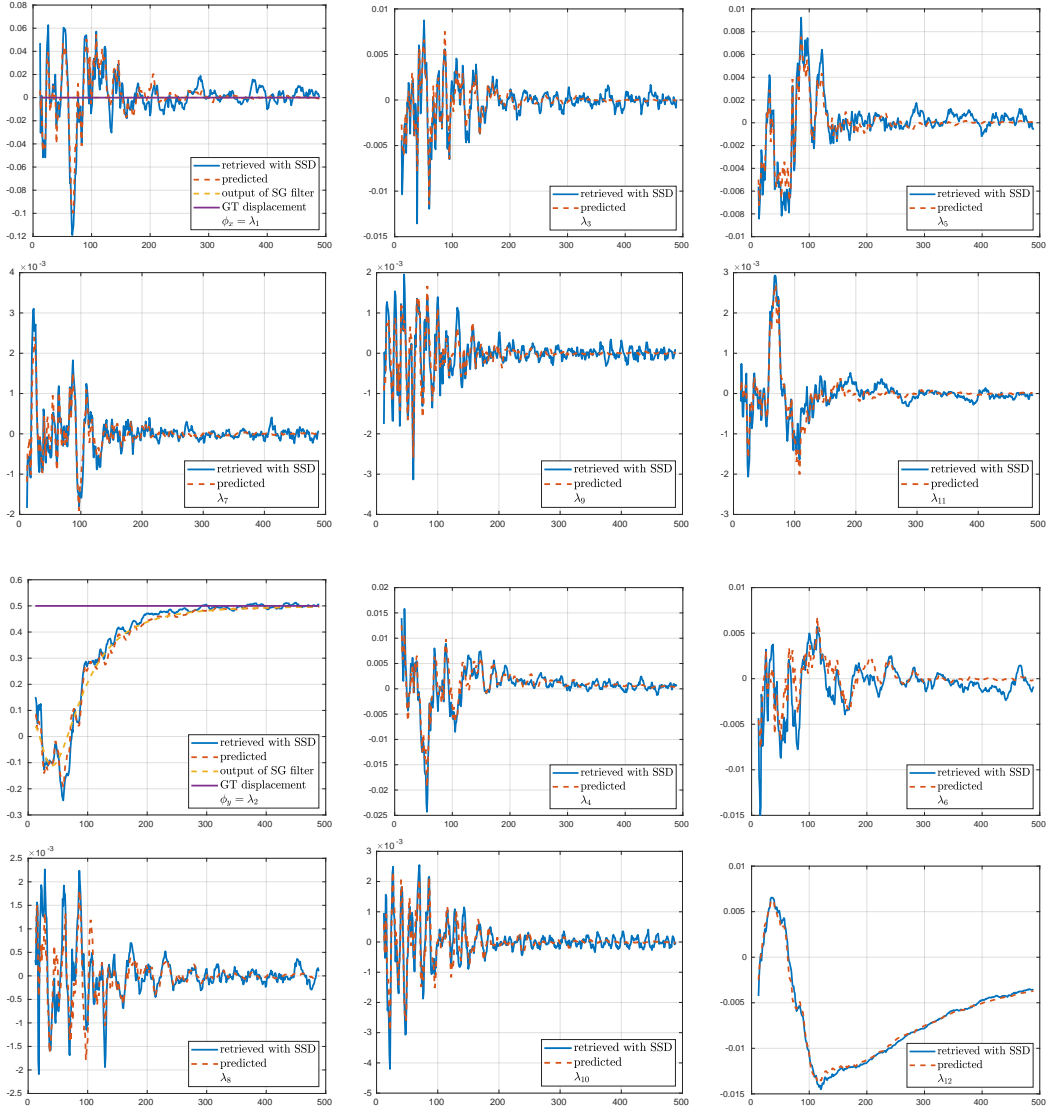


Figure 14: Fine speckle, 19×19 subset Ω and second-order shape functions. From left to right and top to bottom: $\lambda_1, \lambda_3, \lambda_5, \lambda_7, \lambda_9, \lambda_{11}$ concerning x -displacement, and $\lambda_2, \lambda_4, \lambda_6, \lambda_8, \lambda_{10}, \lambda_{12}$ concerning y -displacement.

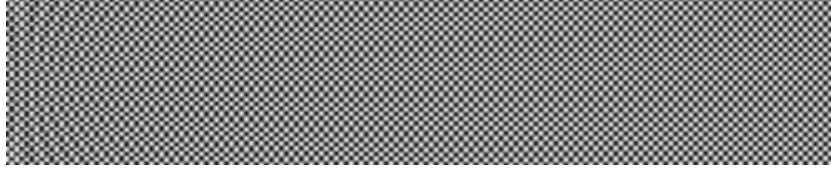


Figure 15: Checkerboard pattern.

3.2.4 A remark on checkerboard patterns

While random speckle patterns have a huge popularity in photomechanics, very recent papers show that checkerboard patterns give less random noise in the retrieved displacement, see [4, 18]. The reason is that the average gradient norm within a subset has larger values than with any typical random speckle patterns. Figure 15 shows a checkerboard pattern of pitch equal to 6 (that is, it is made of juxtaposed black and white squares of width 3 pixels). Such a pattern has been used in the experimental assessment of [15, 18]. In these papers, less spurious fluctuations have been observed in the displacement with checkerboard than with random speckle.

This is confirmed by Figures 16 and 17 which show estimation of the $\mathbf{\Lambda}$ parameters with checkerboards patterns, with the same settings as in Figures 7-8 with large speckle and 11-12 with fine speckle. While predictive formulas are not fully satisfied here, the curves being not perfectly superimposed, it can be noted that the amplitude of the spurious fluctuations is much smaller than with speckle patterns. The amplitude of the high-frequency spurious displacement on the x -component of the displacement is less than $5 \cdot 10^{-3}$, an order of magnitude smaller than with speckles. This high-frequency phenomenon is caused by aliasing; it cannot be seen with large checkerboard pitches (not shown here). Moreover, with such a fine pattern, the numerical estimation of the gradient needed in the predictive formulas is certainly not consistent. While spurious fluctuations cannot be seen, contrary to the fluctuations caused by random speckle patterns, the PIB still plays a role, giving displacement fields which do not fit the output of the SG filter, as can be seen in the cross-section plots in Figures 16 and 17. The difference caused by PIB reaches, however, values smaller than with random speckle patterns.

This experiment illustrates that a periodic and fine pattern such as a checkerboard gives a smaller difference with the output of the SG filter than the classic speckle patterns.

3.2.5 Towards a PIB-free SSD criterion?

This section discusses the extent to which it is possible to get rid of PIB, after Blanchet et al.'s approach [1]. We can see from Section 2.3.3 that if ϕ minimizes the following GT-SSD

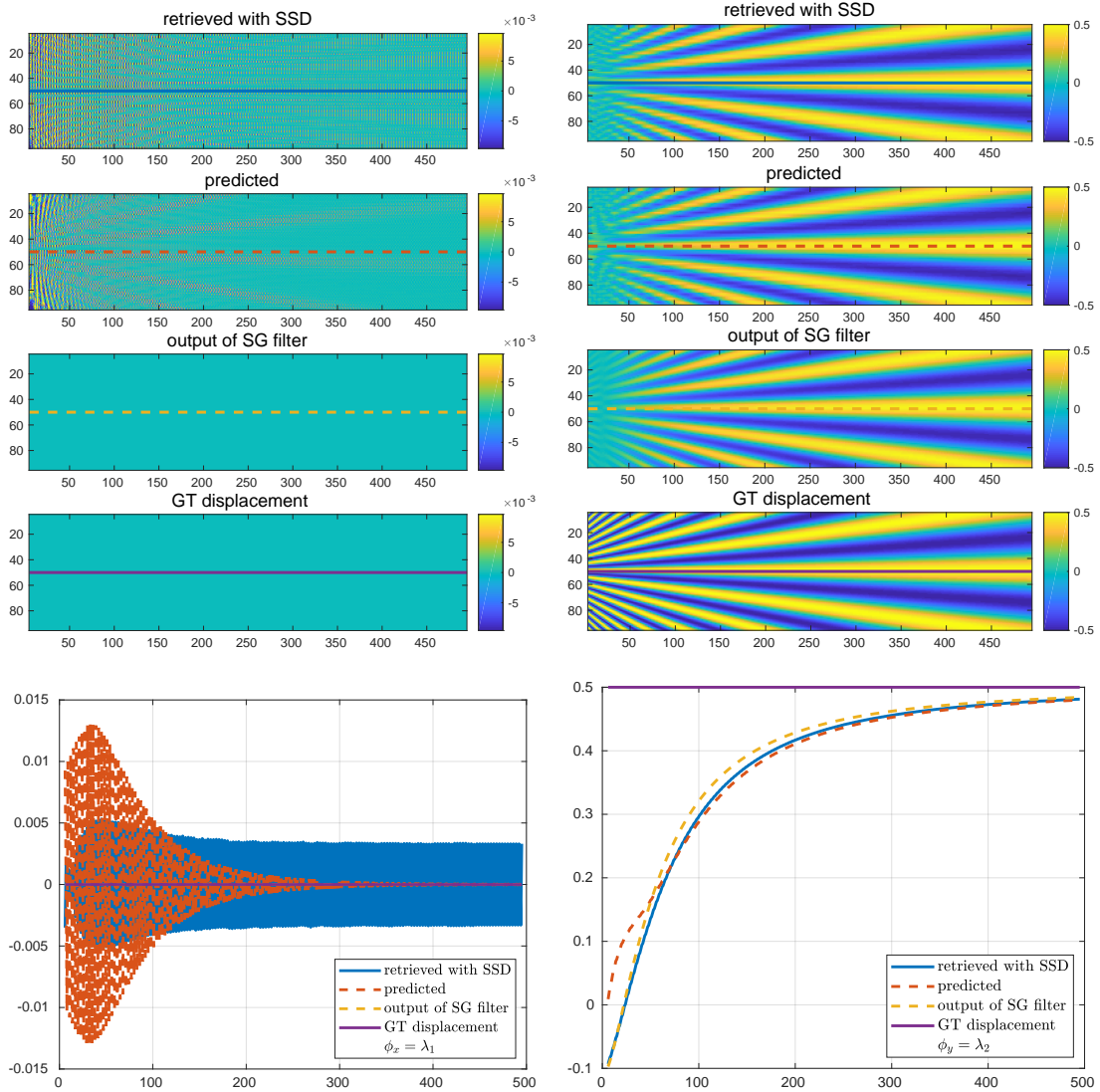


Figure 16: Checkerboard, 7×7 subset Ω and 0-order shape functions. Displacement maps along x - (left) and y -axis (right), and cross-section plots.

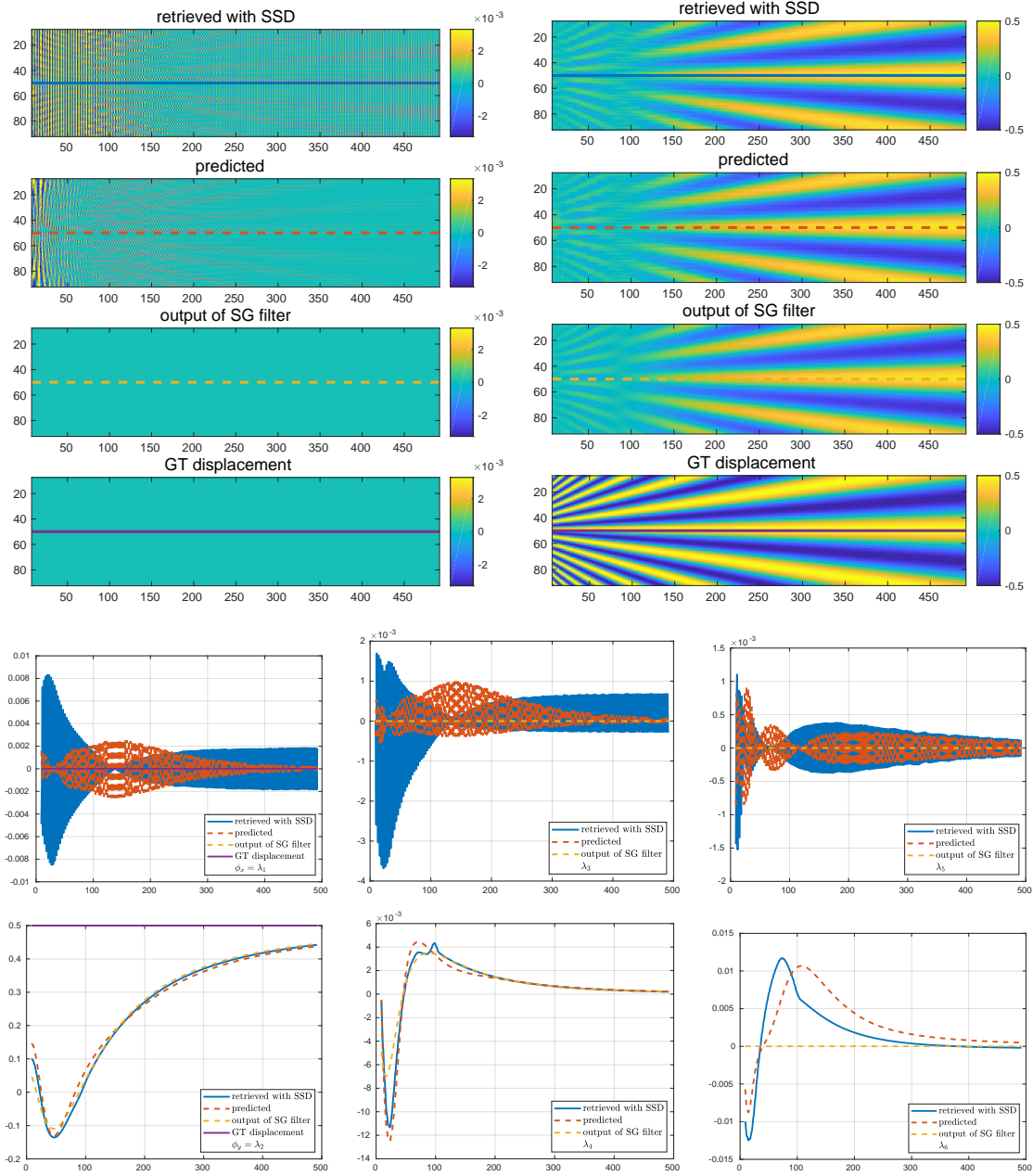


Figure 17: Checkerboard, 13×13 subset Ω and first-order shape functions. Displacement maps along x - (left) and y -axis (right). Cross-section plots, from left to right and top to bottom: $\lambda_1, \lambda_3, \lambda_5, \lambda_2, \lambda_4, \lambda_6$.

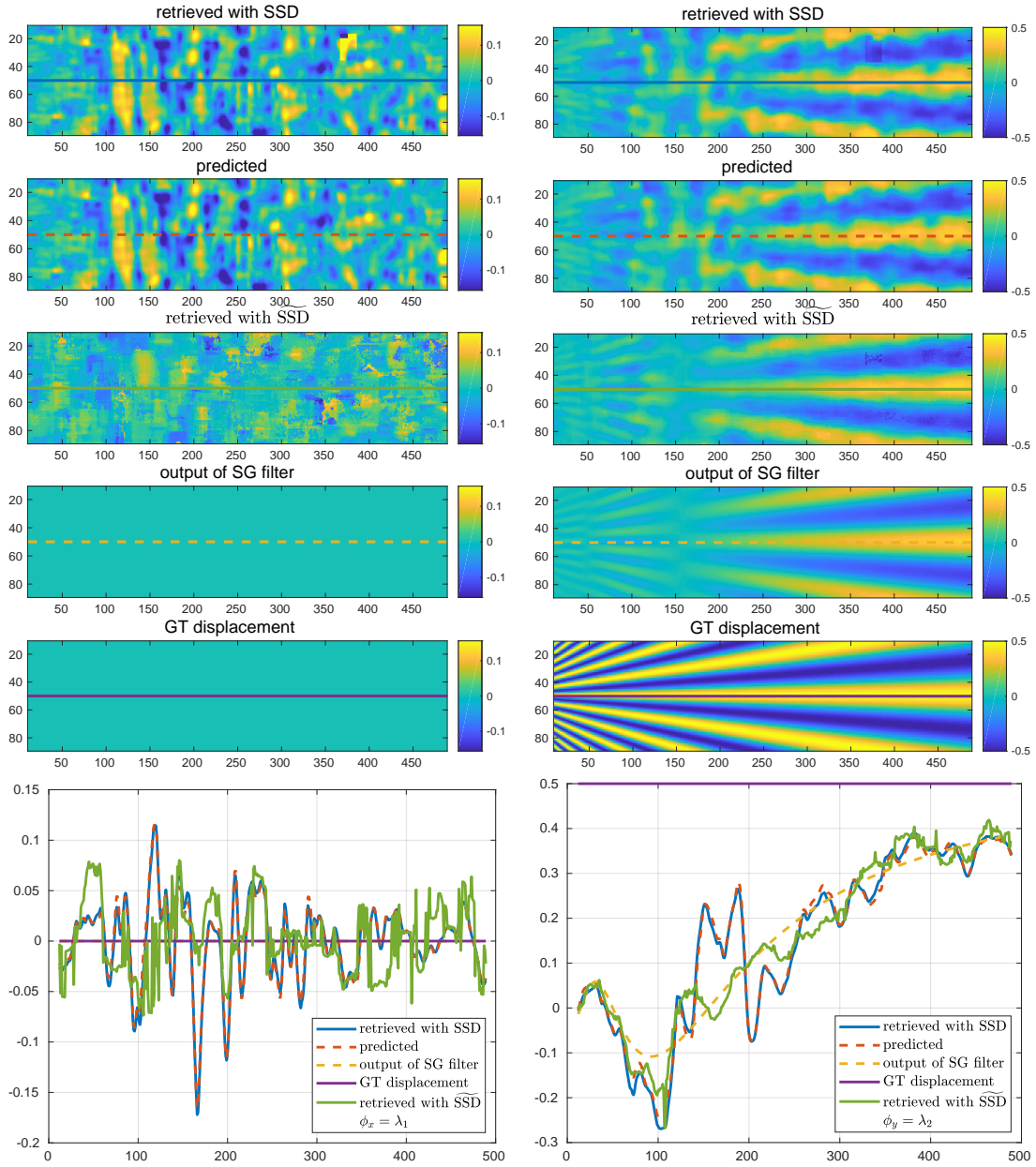


Figure 18: Large speckle, 19×19 subset Ω , zero-order shape functions, and displacement estimation (along x - on the left and y -axis on the right) with the GT-SSD criterion.

criterion:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \frac{1}{\max(g_i^2 \cos^2(\theta_i), \kappa)} \left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \phi(x)) \right)^2 \quad (42)$$

(with the notations of Section 2.3.3), then it is also the least-squares estimate of \mathbf{u} over Ω . In this section, $\kappa = 10^{-2}$ for the large speckle pattern, and $\kappa = 10^{-4}$ for the fine one.

Figure 18 is representative of results obtained with the large speckle pattern. We can see that the displacement retrieved with this criterion is much less impaired by the spurious fluctuations caused by PIB. Here, it is possible to estimate $g_i^2 \cos^2(\theta_i)$ since the GT displacement is known. However, Figure 19 shows that this approach is much less efficient with the fine speckle pattern. The reason is that the GT-SSD criterion requires an estimation of the gradient g_i and of the angle θ_i , which is less accurate with the fine speckle than with the smooth softly-varying speckle pattern. For the same reason, similar results are obtained with the checkerboard pattern discussed in Section 3.2.4.

We now discuss two approaches to PIB-free estimation which do not require the knowledge of the true displacement \mathbf{u} . In the present experiment, \mathbf{u} is actually a 1-D displacement along the x -axis. It is thus possible to use Blanchet et al.'s approach to fattening-free block matching, and estimate ϕ over Ω_x by minimizing (again with MATLAB's `fminunc` function) the following 1D-SSD criterion:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \frac{\left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}'(\mathbf{x}_i + \phi(x)) \right)^2}{\max(|\tilde{\mathcal{I}}'_x(\mathbf{x}_i + \phi(x))|^2, \kappa)} \quad (43)$$

where $\tilde{\mathcal{I}}'_x$ denotes the partial derivative along x of the interpolated image $\tilde{\mathcal{I}}$, an initial guess of ϕ being given by the classic SSD criterion.

The result is shown in Figures 20 and 21 (to be compared to Figures 18 and 19). PIB significantly decreases: 1D-SSD gives a displacement that roughly follows the output of the SG filter.

Nevertheless, realistic displacements in experimental mechanics are bidimensional. We also introduce and test the following 2D-SSD criterion:

$$\sum_{\mathbf{x}_i \in \Omega_{\mathbf{x}}} \frac{\left(\mathcal{I}(\mathbf{x}_i) - \tilde{\mathcal{I}}(\mathbf{x}_i + \phi(x)) \right)^2}{\max(|\nabla \tilde{\mathcal{I}}(\mathbf{x}_i + \phi(x))|^2, \kappa)} \quad (44)$$

Figure 22 is representative of results generally obtained: this approach only marginally allows us to decrease PIB, contrary to the case of 1D-SSD. A slight decreasing of the amplitude of PIB can be noticed. As noted earlier in the 1-D case, displacement information is lost only at points where the derivative of the image vanishes. On the contrary, in the 2-D case the aperture problems manifests itself at any pixel since the component of the displacement orthogonal to the gradient always vanishes. This probably leads here to an incomplete correction of the criterion.

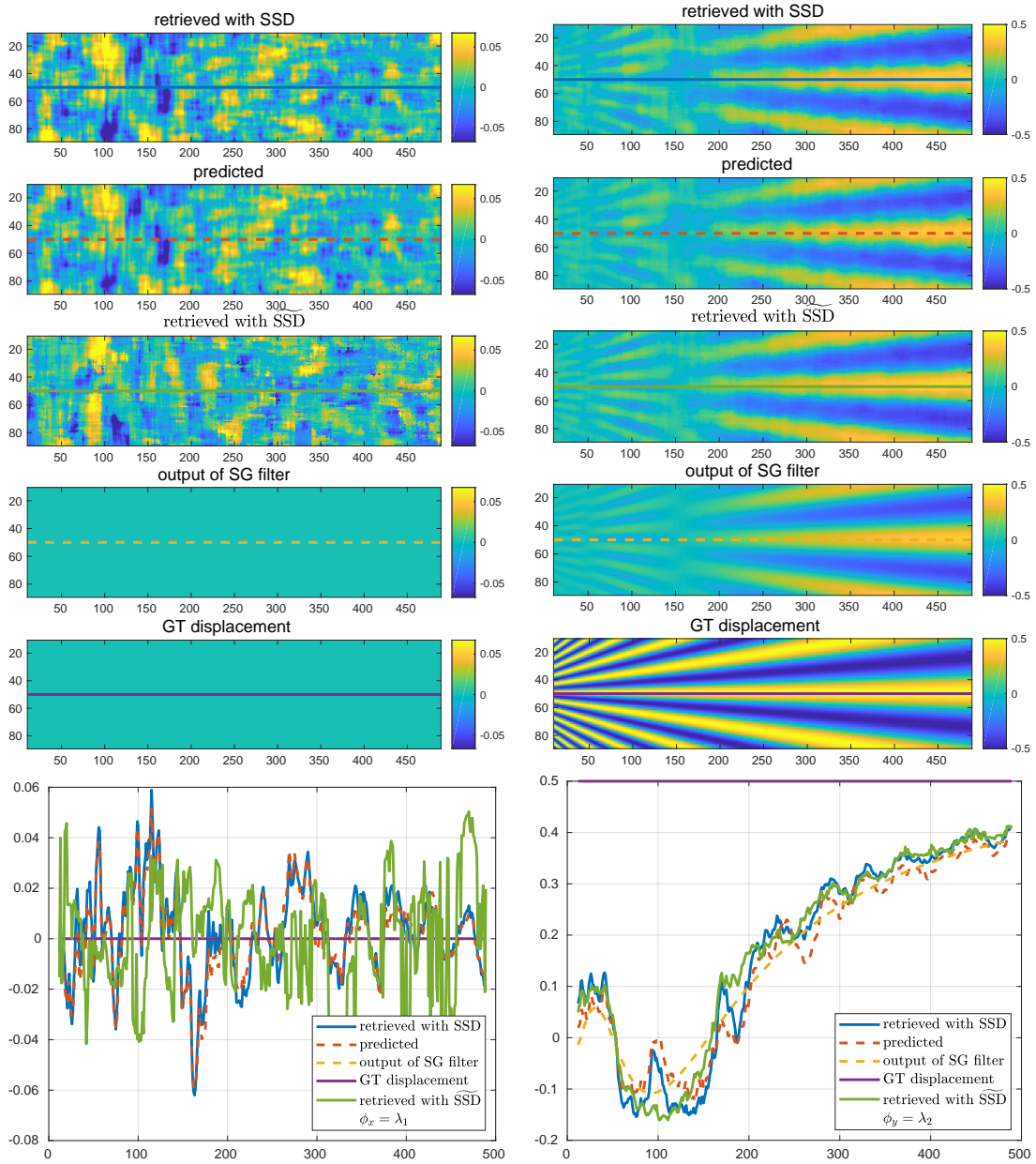


Figure 19: Fine speckle, 19×19 subset Ω , zero-order shape functions, and displacement estimation (along x - on the left and y -axis on the right) with the GT-SSD criterion.

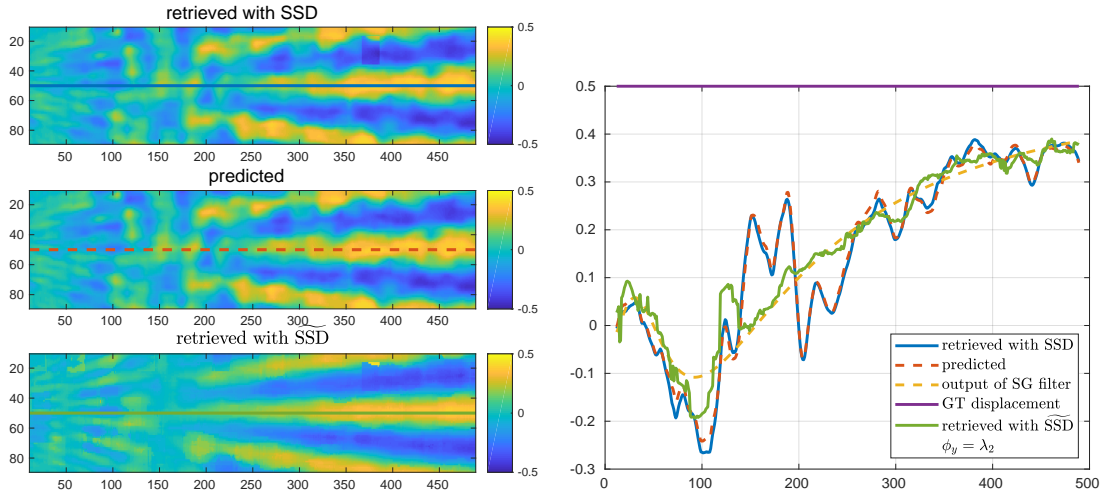


Figure 20: Large speckle, 19×19 subset Ω , zero-order shape functions, and displacement estimation (along x -axis) with the 1D-SSD criterion. To be compared to Figure 18.

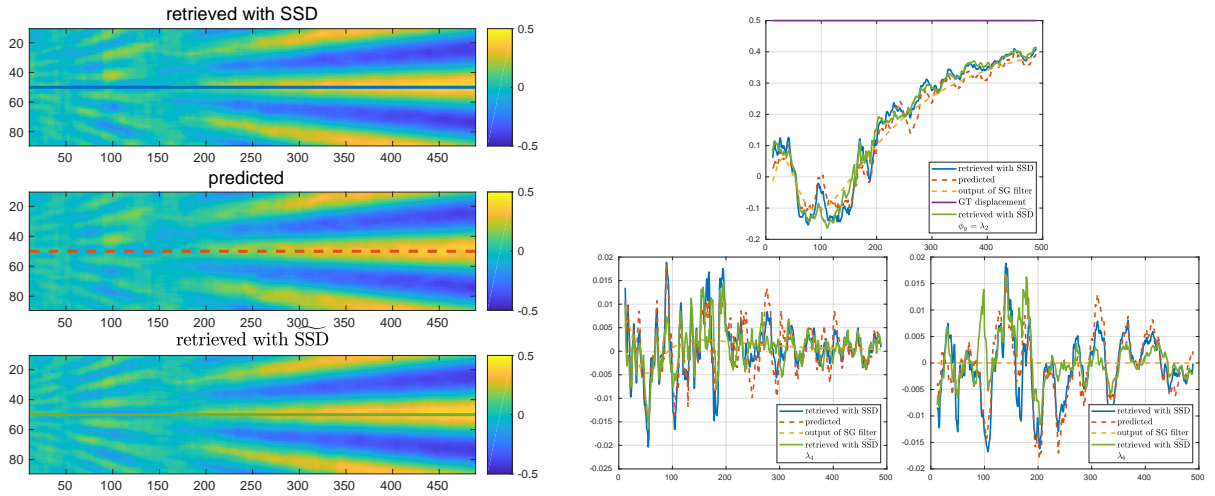


Figure 21: Fine speckle, 19×19 subset Ω , first-order shape functions, and displacement estimation (along x -axis) with the 1D-SSD criterion. To be compared to Figure 19.

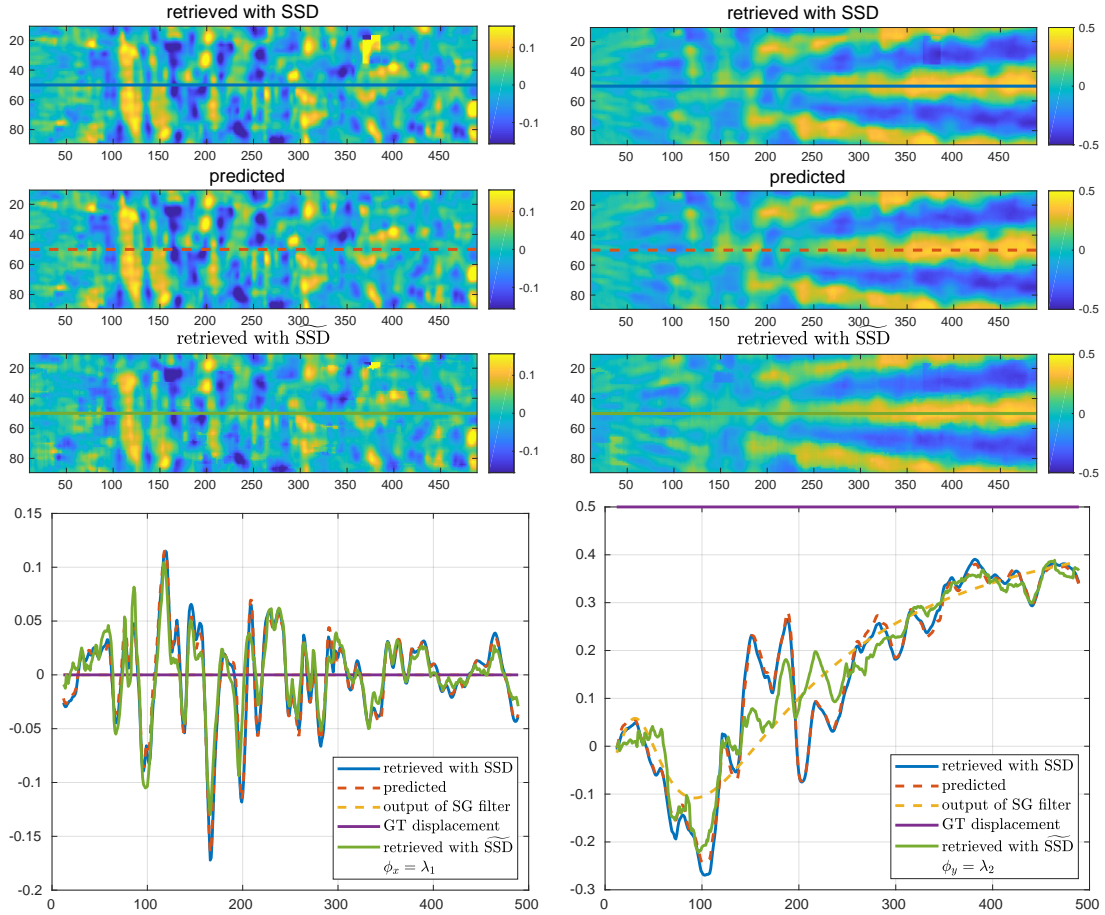


Figure 22: Large speckle, 19×19 subset Ω , zero-order shape functions, and displacement estimation (along x - on the left and y -axis on the right) with the 2D-SSD criterion.

4 Conclusion and open questions

This paper discusses several bias sources in image registration with local parametric estimation via a sum of squared differences criterion, with a focus on photomechanics. A predictive formula is proposed in Theorem 1 for biases caused by interpolation and under-matched shape functions. These sources of errors all depend on the gradient distribution in the underlying images. In particular, the pattern-induced bias (PIB), known as fattening effect in stereoscopic imaging, is caused by the gradient distribution and by the difference between the true displacement field and its local approximation by shape functions. In addition to these biases, the retrieved displacement is affected by spatially correlated random fluctuations caused by sensor noise propagation. This point is discussed in [42].

Several results from the literature are extended or presented in a unifying way. In Section 2.3, we have also completed the contribution of [36] by establishing a rigorous link between the estimated displacement and the true displacement through the Savitzky-Golay (SG) filter, whose parameters depends on the order of the shape functions and on the size of the analysis subset. The link holds because of the very random nature of speckle images.

A numerical assessment of the predictive formulas is discussed as well. First, we have noticed that bicubic interpolation gives biases well below PIB in amplitude, in contrast to bilinear interpolation. Second, it is shown that PIB may have a large amplitude, either if zero-order (used in stereo-imaging) or first and second-order (used in DIC for experimental mechanics applications) shape functions are used. This bias term gives fluctuations around the output of the SG filter of the true displacement. It is striking to note that these fluctuations may have an amplitude of twice the true displacement. As mentioned in Section 2.3, the effect of PIB is equivalent to a spatially varying convolution. A fine understanding of this question would permit to go beyond the deconvolution procedure that is proposed in [16] in order to reduce the measurement bias.

In experimental mechanics applications, defining a marking pattern which is optimal with respect to relevant metrological criteria still remains an open question. Some guidelines are given in Section 3.2.5. In addition, first results discussed in [18] show that checkerboards give a smaller measurement bias than classic speckle patterns used in DIC, and suggest that checkerboards indeed give lower PIB than random speckle patterns. This was numerically verified and illustrated in Section 3.2.4.

Besides, the set of shape functions used to parameterize the local displacement also plays a role: we have shown that PIB involves the scalar product of the image gradient and the difference between the true and the retrieved displacements. It could be of interest to select shape functions minimizing PIB, by improving statistical criteria such as the ones discussed in [7].

Concerning modifications of the sum of squared differences criterion permitting to get rid of pattern-induced bias, we have verified that the approach proposed by Blanchet et al. [1] is an effective method for 1-D displacements met in stereo-imaging. Our numerical experiments show that it is also a valid approach with higher-order shape functions than the

constant local displacement considered in [1]. We have shown that the resulting fattening-free estimation approximates the output of a SG filter, which generalizes [1]. However, the case of 2-D displacements is much more complicated because of the aperture problem, which discards displacement information at any pixel. A PIB-free estimation is still to be designed in this case.

Conflict of interest

The authors declare that they have no conflict of interest.

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