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Wigner Matrix Formalism for phase modulated signals

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Laser beam can carry multi-scales properties in space and time that impact the beam quality. The study of their evolution along complex optical sequences is of crucial interest, particularly in high intensity laser chains. For such analysis, results obtained with standard numerical calculus are strongly dependent of the sampling. In this paper, we develop an analytic model for a sinusoidal phase modulation inside a sequence of first order optics elements based on the Wigner matrix formalism. A Bessel decomposition of the Wigner function gives pseudo-Wigner functions that obey to the general ABCD matrix law transformation without approximations and sampling considerations. Applied to a Gaussian beam, explicit expressions are obtained for the projections of the Wigner function in the sub-spaces and gives a powerful tool for laser beam analysis. The formalism is established in the spatial and temporal domain and can be used to evaluate the impact of the phase noise to the beam properties and is not limited to small modulation depths. In a sake of illustration, the model is applied to the Talbot effect with the analysis of the propagation in the spatial and phase-space domain. A comparison with full numerical calculations evidences the high accuracy of the analytic model that retrieves all the features of the diffracted beam.

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1. INTRODUCTION

Wigner Distribution Function (*WDF*) [1] has been successfully used to study the propagation of optical beams through first-order optical systems [2]. *WDF* is a signal representation in space (time) and frequency simultaneously and this forms an intermediate signal description between pure space (time) representation (electric field) and pure frequency representation (spectrum).

ABCD matrix formalism is widely used technique to evaluate the output beam properties of laser chains including beam propagation, imaging and focusing with lenses for high energy lasers [3,4], and also stretching and compression with gratings for high power lasers [5].

Accurate numerical modelling of such laser chains containing multi-scale considerations in space and time requires important computing resources and is subject to numerical imprecisions.

In this paper we propose a general model obtained by coupling the Wigner approach [2] with the *ABCD* matrix formalism without approximations and errors induced by numerical sampling. Moreover, it allows to extract significant physical parameters from a specific laser chain. We show how the association of these two powerful techniques allows to model the impact of a coherent and incoherent phase noise on the optical properties of the beam. The general model is valid in the spatio-temporal domain and the phase noise can be added to a perfect beam and inserted in any *ABCD* matrix sequence. This allows to model imperfections of any optical components inside the laser chain and to evaluate its impact in the output beam quality.

The coherent phase noise modeling is based on the Bessel decomposition of the sinusoidal phase modulation [6,7], by analogy

with the decomposition of the cross-spectral density function for Bragg diffraction [8].

This decomposition gives pseudo-Wigner functions that obey to the general Wigner *ABCD* matrix transformation. Studying the evolution of the elementary pseudo-Wigner function through *ABCD* transformation provide a complete description of the beam by coherent addition of the pseudo-Wigner function obtained after the *ABCD* transformation. The analytical approach allows to model the impact of high frequency phase modulation in the propagation of large laser beam. Let us notice that the Bessel decomposition is not limited to small modulation of the phase noise. The formalism presented in this work can be also used to model strongly modulated signals in both spatial and time domains. Strongly modulated signals turn out to be of the highest importance in high energy laser chains. The phase of signal can be modulated in the time domain for avoiding Brillouin and Raman scattering [9] and for beam smoothing [10,11] in high energy laser chains.

The contents of the paper are as follows. In section 2 we analyze the decomposition of the Wigner function of a spatially phase modulated signal with sinusoidal modulation at the frequency $K_d/2\pi$, and can be viewed as a coherent analysis of interacting beams. We show that the Wigner function can be written as the sum of pseudo-Wigner functions using the Bessel decomposition of the phase modulation function. These pseudo-Wigner functions obey to the standard linear *ABCD* matrix transformation. In section 3, we establish a matrix formalism for a phase modulated signal in the general case where the phase mask which creates the phase modulation is positioned anywhere in a sequence of linear optical components. In section 4, for comparison, we develop an equivalent matrix formalism for the Gaussian-Schell Model (*GSM*), which is the Wigner model for incoherent or partially coherent

beams. For the incoherent phase noise treatment, Gaussian-Schell model (*GSM*) has been intensively used to model partially coherent beams [12,13] and is well adapted to model imperfections of optical components. A matrix formulation is proposed to transform in a sequence a coherent beam to an incoherent beam. In section 5, we apply the model to Gaussian beams. In this case, the Wigner functions and their projections in the different sub-spaces are described with analytical expressions featuring a Gaussian formulation, avoiding imprecisions induced by numerical integration. In this part, the Wigner projection in the spatio-temporal domain is described, and gives the exact formulation of intensity $I(x,t)$, fluence $F(x)$ and Power $P(t)$ of the pulse. In section 6, we apply the Wigner Bessel decomposition to the analysis of the Talbot effect. The Bessel decomposition in power series for five contributions, respecting the energy conservation, allows an explicit formulation for the Talbot effect in the spatial domain for finite size beams. The first application is the analysis of the Talbot effect in the phase space domain (x,k_x) with dimensionless parameters which allows to observe invariant quantities. A comparison with the *GSM* beam is proposed and can be easily extended to the 2D phase space domain (x,y,k_x,k_y) . A second illustration is the analysis of the Talbot effect in the spatial domain, which corresponds to a transformation of a phase modulation to an amplitude modulation with the propagation. Such properties cannot be obtained with *GSM* model which only treats the evolution of average quantities while the Bessel decomposition gives interferometric information.

2. WIGNER FORMALISM FOR PHASE MODULATED SIGNALS

A. Wigner formalism

Let us consider, in the spatio-temporal domain, an electric field $E(x,t)$. The corresponding Wigner function can be written as [2] with the convention of positive argument for Fourier transformation:

$$W_0(x,t,k_x,\omega) = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} E\left(x + \frac{1}{2}x_1, t + \frac{1}{2}t_1\right) E^*\left(x - \frac{1}{2}x_1, t - \frac{1}{2}t_1\right) \times \exp[ix_1k_x + it_1\omega] dx_1 dt_1 \quad (1)$$

For linear transformations, the *ABCD* matrix law gives [14]:

$$W^{Seq}(x,t,k_x,\omega) = W_0(X(x,t,k_x,\omega), T(x,t,k_x,\omega), K_x(x,t,k_x,\omega), \Omega(x,t,k_x,\omega)) \quad (2)$$

where the functions X , T , K_x , and Ω express the transformation of the arguments of the Wigner function with *ABCD* matrix. This allows considering only the arguments transformation instead of the function transformation.

For Gaussian beams, the Wigner function and its projections into the different subspaces are Gaussian (see Part 5). For a standard initial Gaussian beam of spatial and temporal size Δx and Δt , and a total energy E_i :

$$E_0(x,t) = \sqrt{\frac{2E_i}{\Delta x \Delta t \pi}} \exp\left[-\frac{x^2}{\Delta x^2} - \frac{t^2}{\Delta t^2}\right] \quad (3)$$

The Wigner function is [15]:

$$W_0(x,t,k_x,\omega) = \frac{E_i}{\pi^2} \exp\left[-\frac{2}{\Delta x^2}x^2 - \frac{2}{\Delta t^2}t^2 - \frac{\Delta x^2}{2}k_x^2 - \frac{\Delta t^2}{2}\omega^2\right] \quad (4)$$

B. Sinusoidally phase modulated electric field and Bessel decomposition

Let us consider now, in the spatio-temporal domain, an electric field $E_0(x,t)$ interacting with a spatial sinusoidal phase mask with of period $2\pi/K_d$ and modulation depth m . The resulting field can be expressed as a sum of an infinite Bessel decomposition:

$$E(x,t) = E_0(x,t) \exp[im \sin(K_d x)] = E_0(x,t) \sum_{n=-\infty}^{n=+\infty} J_n(m) \exp[inK_d x] \quad (5)$$

where J_n represents the Bessel function of the first kind of order n . The zero order contribution ($n=0$) is proportional to the field component transmitted without perturbations. We can reduce the analysis to the field Kernel $G_n(x,t)$:

$$G_n(x,t) = E_0(x,t) \exp[inK_d x] \quad (6)$$

to describe the electric field:

$$E(x,t) = \sum_{n=-\infty}^{n=+\infty} J_n(m) G_n(x,t) \quad (7)$$

For the field Kernel $G_n(x,t)$, the equivalent of the Wigner function is:

$$W_{n,l}(x,t,k_x,\omega) = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} G_n\left(x + \frac{1}{2}x_1, t + \frac{1}{2}t_1\right) G_l^*\left(x - \frac{1}{2}x_1, t - \frac{1}{2}t_1\right) \times \exp[ix_1k_x + it_1\omega] dx_1 dt_1 \quad (8)$$

This function is not rigorously a Wigner function because it corresponds to a conjugation of the electric field Kernels with different orders n and l , and it is a complex function. Let us describe two properties of this function:

$$W_{n,l}(x,t,k_x,\omega) = W_{l,n}^*(x,t,k_x,\omega) \quad (9)$$

and:

$$W(x,t,k_x,\omega) = \sum_{n=-\infty}^{n=+\infty} \sum_{l=-\infty}^{l=+\infty} J_n(m) J_l(m) W_{n,l}(x,t,k_x,\omega). \quad (10)$$

The fact that $\overline{W_{n,l}(x,t,k_x,\omega)} = W_{l,n}^*(x,t,k_x,\omega)$ allows us to prove that W is real:

$$W(x,t,k_x,\omega) = \sum_{n=-\infty}^{n=+\infty} \sum_{l=-\infty}^{l=+\infty} J_n(m) J_l(m) \text{Re}(W_{n,l}(x,t,k_x,\omega)) \quad (11)$$

In the following parts, we will reduce our analysis to the pseudo-Wigner function $W_{n,l}(x,t,k_x,\omega)$, which include self-Wigner terms ($n=l$), and cross-Wigner terms ($n \neq l$) [16].

C. Wigner function for Gaussian beam

The Wigner formulation with *ABCD* matrix formalism applied to Gaussian beams gives specific relations between the electric field and the Wigner function [17,18]. For a phase-modulated signal applied to a Gaussian beam, the pseudo-Wigner function is:

$$W_{n,l}(x,t,k_x,\omega) = \frac{E_i}{\pi^2} \times \exp\left[-\frac{2}{\Delta x^2}x^2 - \frac{2}{\Delta t^2}t^2 - \frac{\Delta x^2}{2}\left(k_x + \left(\frac{n+l}{2}\right)K_d\right)^2 - \frac{\Delta t^2}{2}\omega^2\right] \times \exp[i(n-l)K_d x] \quad (12)$$

The spatial phase modulation induces a shift in the k_x space for the amplitude. This expression can be then reduced to:

$$W_{n,l}(x,t,k_x,\omega) = W_0\left(x,t,k_x + \left(\frac{n+l}{2}\right)K_d, \omega\right) \exp[i(n-l)K_d x] \quad (13)$$

or, in an equivalent representation:

$$W_{n,l}(x,t,k_x,\omega) = W_0(x,t,k_x,\omega) \times \exp\left[-\frac{\Delta x^2}{2}\left(\frac{n+l}{2}\right)K_d\left(2k_x + \left(\frac{n+l}{2}\right)K_d\right)\right] \exp[i(n-l)K_d x] \quad (14)$$

It shows that the pseudo-Wigner can be written as:

- the product of the Wigner function of the initial field containing modulated term with an exponential function containing linear phase term in x .

- or the product between the Wigner function of the initial field, not modulated, with an exponential function of linear amplitude term in k_x and a linear phase term in x .

In the following part, we study the evolution of the Wigner function with the matrix formalism and show that we can always write the Wigner function in those two representations.

3. MATRIX FORMALISM

In this section, we aim at describing separately the different elementary transformations of the pseudo-Wigner functions needed for describing a complete $ABCD$ sequence transformations and phase modulation. In the proposed model, the phase modulation can be positioned anywhere on the sequence, it is the reason why we consider an $ABCD$ transformation M_0 followed by a phase modulation, and followed by a second $ABCD$ transformation M_1 .

A. Matrix formalism for a phase modulated signal

The arguments of the initial Wigner function are denoted into the form of a column matrix:

$$\mathbf{r} = \begin{pmatrix} x \\ t \\ k_x \\ \omega \end{pmatrix}, \quad (15)$$

The Wigner expression of the field is then given by [2]:

$$W_0(\mathbf{r}) = \frac{E_t}{\pi^2 \sqrt{\det(\mathbf{Q}^{-1})}} \exp(-\mathbf{r}^T \mathbf{Q} \mathbf{r}), \quad (16)$$

where T denotes the transpose of the matrix column \mathbf{r} and where the matrix \mathbf{Q} represents the Curvature Matrix [19], which is also equivalent to the inverse of the 2nd order matrix moment of the field [20]. \mathbf{Q} is real symmetric, positive definite and symplectic [17].

Integration over the four arguments (x,t,k_x,ω) gives the total energy E_t of the beam:

$$E_t = \int W_0(\mathbf{r}) d\mathbf{r}, \quad (17)$$

with $\det(\mathbf{Q}^{-1}) = 1$.

If the phase of the initial field is modulated, the pseudo-Wigner functions are given by:

$$W_{n,l}(\mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-(\mathbf{r} + \mathbf{S}^{n,l})^T \mathbf{Q} (\mathbf{r} + \mathbf{S}^{n,l})\right) \exp\left(i(\mathbf{D}^{n,l})^T \mathbf{r}\right), \quad (18)$$

with:

$$\begin{cases} \mathbf{S}^{n,l} = \left(\frac{n+l}{2}\right)K_d \mathbf{I}_3, \\ \mathbf{D}^{n,l} = (n-l)K_d \mathbf{I}_1 \end{cases}, \quad (19)$$

and,

$$\mathbf{I}_3 = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \quad \mathbf{I}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}. \quad (20)$$

It can be noticed that the phase modulation induces an additional linear phase term and a modification of the amplitude of the pseudo Wigner function. The matrix formalism is very convenient since it allows to address also phase modulation in the time domain or even both spatial and time domains considering the right choice of $\mathbf{S}^{n,l}$, $\mathbf{D}^{n,l}$, and \mathbf{I} column matrices.

$W_{n,l}$ are self ($n=l$) or cross ($n \neq l$) Wigner terms. The self-Wigner terms are real ($\mathbf{D}_{n,l}=0$) and they transport the energy of the beam on the even ($n+l=2n$) contributions. The cross-Wigner terms have a non-null phase complement ($\mathbf{D}_{n,l} \neq 0$). They don't carry energy if the number of fringes is significant in the field envelope (i.e. $\cos[(n-l)K_d x] \approx 0$ in mean value). The $(n+l)$ contribution of the pseudo-Wigner terms can be even or odd, but the odd contributions are composed with only cross-terms without energy. It will be illustrated in Part 6 in Fig. 1 and Fig. 2.

B. ABCD transformation of an initially phase modulated field

Let us now consider the impact of an $ABCD$ transformation M_1 on the pseudo-Wigner function. If M_1 represents the 4×4 $ABCD$ matrix describing the transformation of the argument \mathbf{r} through a first order optical system, then, the evolution of the pseudo-Wigner function after the transformation is given by:

$$W_{n,l}^{M_1}(\mathbf{r}) = W_{n,l}(M_1 \mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-(\mathbf{M}_1 \mathbf{r} + \mathbf{S}^{n,l})^T \mathbf{Q} (\mathbf{M}_1 \mathbf{r} + \mathbf{S}^{n,l})\right) \exp\left(i(\mathbf{D}^{n,l})^T \mathbf{M}_1 \mathbf{r}\right). \quad (21)$$

This transformation corresponds to the change:

$$\mathbf{r} \rightarrow \mathbf{M}_1 \mathbf{r}. \quad (22)$$

The $ABCD$ transformation law is applied only to the column matrix \mathbf{r} and not to the $\mathbf{S}^{n,l}$ because $ABCD$ law acts only on quadratic terms for Gaussian beams. For first order optical systems, $ABCD$ matrix are symplectic [7,14].

C. ABCD transformation for a field without modulation to the phase mask component

We now consider the case of the $ABCD$ transformation M_0 before the phase modulation, i.e. for a beam without phase modulation. In this case the transformation through $ABCD$ matrix M_0 is:

$$W_0^{M_0}(\mathbf{r}) = W_0(M_0 \mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-(\mathbf{M}_0 \mathbf{r})^T \mathbf{Q} (\mathbf{M}_0 \mathbf{r})\right). \quad (23)$$

If, after this $ABCD$ transformation, the field is modulated in phase, the pseudo Wigner function becomes:

$$W_{n,l}^{M_0+\varphi(x)}(\mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-(\mathbf{M}_0(\mathbf{r} + \mathbf{S}^{n,l}))^T \mathbf{Q} (\mathbf{M}_0(\mathbf{r} + \mathbf{S}^{n,l}))\right) \times \exp\left(i(\mathbf{D}^{n,l})^T \mathbf{r}\right). \quad (24)$$

An important remark has to be done at this stage: if the field is transformed with an $ABCD$ sequence followed by a phase modulation, the \mathbf{M}_0 matrix transformation acts both on \mathbf{r} and $\mathbf{S}^{n,l}$ components.

D. Phase mask in a middle of an ABCD sequence

We now consider the complete sequence of a phase mask located between two 4×4 \mathbf{M}_0 and \mathbf{M}_1 matrices. By combining the results derived in the last two sections, we obtain the following pseudo-Wigner function of the complete sequence:

$$W_{n,l}^{Seq}(\mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-\left(\mathbf{M}_0(\mathbf{M}_1\mathbf{r} + \mathbf{S}^{n,l})\right)^T \mathbf{Q}\left(\mathbf{M}_0(\mathbf{M}_1\mathbf{r} + \mathbf{S}^{n,l})\right)\right) \times \exp\left(i\left(\mathbf{D}^{n,l}\right)^T \mathbf{M}_1\mathbf{r}\right), \quad (25)$$

which is also equivalent to:

$$W_{n,l}^{Seq}(\mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-\left(\mathbf{M}_0\mathbf{M}_1\mathbf{r} + \mathbf{M}_0\mathbf{S}^{n,l}\right)^T \mathbf{Q}\left(\mathbf{M}_0\mathbf{M}_1\mathbf{r} + \mathbf{M}_0\mathbf{S}^{n,l}\right)\right) \times \exp\left(i\left(\mathbf{D}^{n,l}\right)^T \mathbf{M}_1\mathbf{r}\right). \quad (26)$$

This general expression can be used for any linear $ABCD$ transformation, including propagation and focusing in the spatial domain, dispersion in the temporal domain, but also the spatio-temporal coupling *via* diffraction gratings stretching and compression of the beam.

An alternative formulation can also be found by writing:

$$W_{n,l}^{Seq}(\mathbf{r}) = W_0^{Seq}(\mathbf{r}) \exp\left(A_{n,l}(\mathbf{r}) + i\Phi_{n,l}(\mathbf{r})\right), \quad (27)$$

with:

$$\begin{cases} W_0^{Seq}(\mathbf{r}) = W_0(\mathbf{M}_0\mathbf{M}_1\mathbf{r}) = \frac{E_t}{\pi^2} \exp\left(-\left(\mathbf{M}_0\mathbf{M}_1\mathbf{r}\right)^T \mathbf{Q}\left(\mathbf{M}_0\mathbf{M}_1\mathbf{r}\right)\right) \\ A_{n,l}(\mathbf{r}) = -\left(\frac{n+l}{2}\right) K_d \\ \quad \times \left[\left(\mathbf{M}_0\mathbf{I}_3\right)^T \mathbf{Q}\left(\mathbf{M}_0\mathbf{M}_1\mathbf{r}\right) + \left(\mathbf{M}_0\mathbf{M}_1\mathbf{r}\right)^T \mathbf{Q}\left(\mathbf{M}_0\mathbf{I}_3\right) \right] \\ \quad - \left(\frac{n+l}{2}\right)^2 K_d^2 \left(\mathbf{M}_0\mathbf{I}_3\right)^T \mathbf{Q}\left(\mathbf{M}_0\mathbf{I}_3\right) \\ \Phi_{n,l}(\mathbf{r}) = \left(\mathbf{D}^{n,l}\right)^T \mathbf{M}_1\mathbf{r} = (n-l)K_d \left(\mathbf{I}_1\right)^T \mathbf{M}_1\mathbf{r} \end{cases} \quad (28)$$

It shows that the complete Wigner function can be written as the product of the Wigner function without phase modulation with a sum of amplitude and phase terms, as the formulation established in section 2.C:

$$W_{n,l}^{Seq}(\mathbf{r}) = W_0^{Seq}(\mathbf{r}) \sum_{n=-\infty}^{n=+\infty} \sum_{l=-\infty}^{l=+\infty} J_n(m) J_l(m) \exp\left(A_{n,l}(\mathbf{r}) + i\Phi_{n,l}(\mathbf{r})\right). \quad (29)$$

E. Discussion

The formulations (26) or (27) can be used for modelling and visualizing the impact of the phase modulation on the middle of complex optical sequences and are directly applicable to high laser chains [3-5] containing not only propagation, but lenses, gratings etc... Compared to classical numerical calculation (FFT based algorithms), (26) or (27) takes the advantage to carry multiscale properties ($1/K_d \ll \Delta x$ for example) without approximations and errors induced by numerical sampling.

The Bessel decomposition presented here can be extended to any periodic signal by replacing Bessel J_n and J_l coefficients by the general Fourier series decomposition (plane waves decomposition). In the same way, the amplitude envelope of the electric field $E_0(x,t)$ can be decomposed in sum of off-axis Gaussian fields (see for example [21] for

supergaussian decomposition). Decomposition of periodic signals and amplitude envelope on Gaussian basis are very similar on the Wigner domain:

- an amplitude shift in k_x (respectively ω in the time domain) for periodic signals, or in x (respectively t) for the envelope decomposition.
- a phase shift in x (respectively t) for periodic signals, or in k_x (respectively ω) for the envelope decomposition.

In the other hand, a matrix representation of the Wigner formalism allows to represent the evolution of the second order moments of an arbitrary beam. Let us point out that the study of arbitrary beams limited to their second order moments reduces to a Gaussian representation.

To obtain the physical observables (intensity, fluence, power etc...), it is necessary to project the Wigner function in the sub-spaces. At this step, it is obtained by numerical integration of formulations (26) or (27). In Part 5, we propose a new step consisting of an analytical integration of the Wigner function without approximation. It then gives exact formulations of the physical observables that can be directly used.

4. COMPARISON WITH THE GAUSSIAN SCHELL-MODEL (GSM)

A. Gaussian-Schell Model

The correlation function of an incoherent beam in the spatial and temporal domain can take the general form [22]:

$$\Gamma_{GSM}(x_1, t_1, x_2, t_2) = \frac{E_t}{2\pi\Delta x\Delta t} \times \exp\left[-\frac{x_1^2 + x_2^2}{\Delta x^2} - \frac{(x_1 - x_2)^2}{2\delta_x^2} - \frac{t_1^2 + t_2^2}{\Delta t^2} - \frac{(t_1 - t_2)^2}{2\delta_t^2}\right] \quad (30)$$

where Δx and Δt are respectively the spatial and temporal dimensions of the beam and the parameters δ_x et δ_t describe respectively the spatial and temporal coherence characteristics of the beam. If $\Delta x/\delta_x$ and $\Delta t/\delta_t \ll 1$, the beam becomes a perfect coherent Gaussian beam, in the other cases, it describes a partially coherent beam.

The correlation function can also be rearranged under a matrix form [23]. For such correlation function, the Wigner function is written as:

$$W_{GSM}(x, t, k_x, \omega) = \frac{1}{4\pi^2} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \Gamma_{GSM}\left(x + \frac{1}{2}x_1, t + \frac{1}{2}t_1, x - \frac{1}{2}x_1, t - \frac{1}{2}t_1\right) \exp\left[ix_1 k_x + it_1 \omega\right] dx_1 dt_1. \quad (31)$$

A matrix formulation of the Wigner function is [24]:

$$W_{GSM}(\mathbf{r}) = E_t \frac{pq}{\pi^2} \exp\left(-\mathbf{r}^T \mathbf{Q}_{GSM} \mathbf{r}\right), \quad (32)$$

with:

$$\mathbf{Q}_{GSM} = \begin{pmatrix} \frac{2}{\Delta x^2} & 0 & 0 & 0 \\ 0 & \frac{2}{\Delta t^2} & 0 & 0 \\ 0 & 0 & p^2 \frac{\Delta x^2}{2} & 0 \\ 0 & 0 & 0 & q^2 \frac{\Delta t^2}{2} \end{pmatrix} \quad (33)$$

$$p = \frac{1}{\sqrt{1 + \frac{\Delta x^2}{\delta_x^2}}}$$

$$q = \frac{1}{\sqrt{1 + \frac{\Delta t^2}{\delta_t^2}}}$$

In such representation, the parameters p and q will vary from 0 (incoherent field) to 1 (coherent field). Values between 0 and 1 will give a description of partially coherent field. For highly incoherent field, we will have:

$$\begin{cases} p \approx \frac{\delta_x}{\Delta x} \\ q \approx \frac{\delta_t}{\Delta t} \end{cases}, \quad (34)$$

and ($p, q < 1$).

The evolution of the Wigner function through an $ABCD$ optical system described by the 4×4 matrix \mathbf{M} is then:

$$W_{GSM}(\mathbf{r}) = E_t \frac{pq}{\pi^2} \exp\left(-(\mathbf{M}\mathbf{r})^T \mathbf{Q}_{GSM}(\mathbf{M}\mathbf{r})\right). \quad (35)$$

In such case, the spatial and temporal dimensions Δx and Δt in \mathbf{Q}_{GSM} matrix should be taken at the phase mask position.

B. Matrix transformation of a standard beam to a GSM beam

In order to take into account the incoherent noise given by an optical component in the middle of an $ABCD$ sequence matrix, we need to transform a part P_m of a standard Gaussian coherent beam to an incoherent beam, both in spatial and temporal domain. To find the transformation law from a coherent beam to an incoherent beam we should find the matrix expression \mathbf{M}_{GSM} obeying the equation:

$$\begin{aligned} W_{GSM}(\mathbf{r}) &= E_t \frac{pq}{\pi^2} \exp\left(-\mathbf{r}^T \mathbf{Q}_{GSM} \mathbf{r}\right) \\ &= E_t \frac{pq}{\pi^2} \exp\left(-(\mathbf{M}_{GSM} \mathbf{r})^T \mathbf{Q}_0(\mathbf{M}_{GSM} \mathbf{r})\right), \end{aligned} \quad (36)$$

with the expression of the coherent beam matrix \mathbf{Q}_0 :

$$\mathbf{Q}_0 = \begin{pmatrix} 2/\Delta x^2 & 0 & 0 & 0 \\ 0 & 2/\Delta t^2 & 0 & 0 \\ 0 & 0 & \Delta x^2/2 & 0 \\ 0 & 0 & 0 & \Delta t^2/2 \end{pmatrix}. \quad (37)$$

For a standard Gaussian initial field, we find the expression for the matrix \mathbf{M}_{GSM} in a general case:

$$\mathbf{M}_{GSM} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & q \end{pmatrix}. \quad (38)$$

For a perfect optical component (no noise), we obtain the identity matrix. In the general case, the determinant of the matrix \mathbf{M}_{GSM} is not equal to 1 and should be taken into account in a total sequence for the normalization of the Wigner function. A realistic beam can be written as the sum of an ideal beam and a noise beam:

$$W(\mathbf{r}) = (1 - P_m^2)W_0(\mathbf{r}) + P_m^2 W_{GSM}(\mathbf{r}), \quad (39)$$

where P_m is the modulation depth. If the noise is created at the beginning, the evolution of the Wigner function after a matrix sequence \mathbf{M} is:

$$W(\mathbf{r}) = \frac{E_t}{\pi^2} \left\{ (1 - P_m^2) \exp\left(-(\mathbf{M}\mathbf{r})^T \mathbf{Q}_0(\mathbf{M}\mathbf{r})\right) + P_m^2 pq \exp\left(-(\mathbf{M}\mathbf{r})^T \mathbf{Q}_{GSM}(\mathbf{M}\mathbf{r})\right) \right\}. \quad (40)$$

If the noise is created after a first matrix sequence \mathbf{M}_0 , then after a second matrix sequence \mathbf{M}_1 the Wigner function is:

$$W(\mathbf{r}) = \frac{E_t}{\pi^2} \left\{ (1 - P_m^2) \exp\left(-(\mathbf{M}\mathbf{r})^T \mathbf{Q}_0(\mathbf{M}\mathbf{r})\right) + P_m^2 pq \exp\left(-(\mathbf{M}_0 \mathbf{M}_{GSM} \mathbf{M}_1 \mathbf{r})^T \mathbf{Q}_0(\mathbf{M}_0 \mathbf{M}_{GSM} \mathbf{M}_1 \mathbf{r})\right) \right\} \quad (41)$$

C. Comparison with the Bessel decomposition

The Wigner formalism for partially coherent beam (41) takes the same aspect as the one expressed for the Bessel decomposition (26) or (27) and the two formalism can be compared and combined together to model a more complex beam or a more complex sequence containing both coherent and partially coherent contributions.

If the phase decomposition is given by $N \times N$ spectral contributions (10) or (29), the Wigner formulation of GSM beam is lighter. It allows simplify the impact of phase modulation complexity but is giving only partial information of the beam behavior.

5. ANALYTIC INTEGRATION OF OBSERVABLES

A. Motivation

This step is important to give an explicit formulation of the Wigner function and the marginals [15] (intensity, fluence and spectrum observables) by the direct analytical integration of the Wigner function over the different sub-spaces and without approximations.

Compared to Part 3, it gives exact and ready-made additional formulations without the errors induced by the sampling of numerical integration, which could be non-negligible in the case of multi-scale process.

Because all the results are analytical, they can be used to extract physical properties of the beam. Nevertheless, specific properties like invariant quantities depend strongly of the optical sequence and initial properties of the beam and cannot be generalized. However, this step gives a common tool, independent of the optical sequence, which is a foundation for further and deeper analysis, by following the same approach as illustrated in Part 6 with a specific application.

B. Injected beam characteristics

Let us consider an initial field described by a Gaussian function in the spatio-temporal domain. The pseudo-Wigner functions therefore become Gaussian like their projections in every sub-spaces. Let us consider the case where a single phase mask is inserted in an $ABCD$ sequence described with a matrix $\mathbf{M}_0 = [p_{ij}]$ before the phase mask and a

matrix $M_I=[q_{ij}]$ after the phase mask with a total $ABCD$ sequence $M=M_0M_I=[S_{ij}]$, with $ij=[1;4]$.

The diagonal curvature matrix of a standard Gaussian field can be cast:

$$Q_0 = \begin{pmatrix} 2/\Delta x^2 & 0 & 0 & 0 \\ 0 & 2/\Delta t^2 & 0 & 0 \\ 0 & 0 & \Delta x^2/2 & 0 \\ 0 & 0 & 0 & \Delta t^2/2 \end{pmatrix}. \quad (42)$$

For arbitrary field, we must consider the general curvature matrix Q that can always be diagonalized through a symmetric transformation with symplectic matrix S : $Q=S^T Q_0 S$ (Williamson's theorem [17]).

C. Steps for Wigner calculation and projections in the sub-spaces

The input parameters describe the optical system with the 4×4 matrix coefficients p_{ij} and q_{ij} . The sinusoidal phase modulation is characterized by its modulation depth m , its spatial frequency $K_d/2\pi$ and also, the number of terms selected in the Bessel decomposition. Projections of the pseudo Wigner function $W_{n,l}$ in the sub-spaces are given by [14]:

$$\begin{cases} I_{n,l}(x,t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} W_{n,l}(x,t,k_x,\omega) dk_x d\omega \\ F_{n,l}(x) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} W_{n,l}(x,t,k_x,\omega) dk_x d\omega dt = \int_{-\infty}^{\infty} I_{n,l}(x,t) dt \\ P_{n,l}(t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} W_{n,l}(x,t,k_x,\omega) dk_x d\omega dx = \int_{-\infty}^{\infty} I_{n,l}(x,t) dx \end{cases} \quad (43)$$

These expressions are complex, but their summation over n and l are real and positive.

In a first step, the pseudo-Wigner function is written for the complete process. It is a Gaussian function characterized by the polynomial coefficients $C_{ij}^{W_0}$, A_j^L , Φ_j^L , and A_0 (Appendixes A and B).

The part of the Wigner function which is not affected by the phase modulation is:

$$W_0^{Seq}(x,t,k_x,\omega) = \exp \left[\begin{array}{l} C_{xx}^{W_0} x^2 + C_{tt}^{W_0} t^2 + C_{k_x k_x}^{W_0} k_x^2 + C_{\omega \omega}^{W_0} \omega^2 \\ + C_{xk_x}^{W_0} x k_x + C_{xt}^{W_0} x t + C_{x\omega}^{W_0} x \omega \\ + C_{k_x t}^{W_0} k_x t + C_{k_x \omega}^{W_0} k_x \omega + C_{t\omega}^{W_0} t \omega \end{array} \right] \quad (44)$$

where the $C_{ij}^{W_0}$ terms depend only of the 4×4 matrix coefficients S_{ij} . They represent the quadratic part in (x,t,k_x,ω) arguments of the Wigner function.

The complete pseudo-Wigner function becomes:

$$W_{n,l}^{Seq}(x,t,k_x,\omega) = W_0^{Seq}(x,t,k_x,\omega) \times \exp \left[\begin{array}{l} (n+l)K_d(A_x^L x + A_t^L t + A_{k_x}^L k_x + A_{\omega}^L \omega) \\ + i(n-l)K_d(\Phi_x^L x + \Phi_t^L t + \Phi_{k_x}^L k_x + \Phi_{\omega}^L \omega) \\ + (n+l)^2 K_d^2 A_0 \end{array} \right] \quad (45)$$

where the A_j^L and Φ_j^L terms are respectively the real and imaginary parts of the linear terms in (x,t,k_x,ω) . They depend of the matrix coefficients p_{ij} and q_{ij} . For the real part, they act as a shift of the (x,t,k_x,ω) coordinate system. The A_0 term is the real and negative constant term function of p_{ij} only. The real constant term is important because it plays the role of attenuation term of the contribution. All these terms do not depend on the modulation depth m , frequency K_d , integers n and l but only on the matrix parameters p_{ij} , q_{ij} and initial beam parameters defined by Q_0 .

The complete Wigner function is then given by the summation of the pseudo-Wigner functions with a ponderation by the first kind Bessel function product term $J_n(m)J_l(m)$.

In a second step, we calculate the projection of the pseudo-Wigner function in the selected sub-space by integration over the appropriate arguments. Because integration of Gaussian functions gives Gaussian function, the general expression in the sub-space projection is characterized by the coefficients $C_{ij}^{I_0}$, A_j^L , Φ_j^L , A_0^S , A_0^D and Φ^D (see appendices C and D for details), which have, in the general case, similar dependence properties as the $C_{ij}^{W_0}$, A_j^L , Φ_j^L , and A_0 coefficients. For the quadratic part of the intensity, not affected by the phase modulation, we have:

$$I_0^{Seq}(x,t) = \exp \left[C_{xx}^{I_0} x^2 + C_{tt}^{I_0} t^2 + C_{xt}^{I_0} x t \right], \quad (46)$$

and the complete pseudo-intensity is then:

$$I_{n,l}^{Seq}(x,t) = I_0^{Seq}(x,t) \exp \left[\begin{array}{l} (n+l)K_d \left[A_{I_x}^L x + A_{I_t}^L t \right] \\ + i(n-l)K_d \left[\Phi_{I_x}^L x + \Phi_{I_t}^L t \right] \\ + K_d^2 \left[(n+l)^2 A_{I_0}^S + (n-l)^2 A_{I_0}^D \right] \\ + i(n^2 - l^2) K_d^2 \Phi_0^D \end{array} \right] \quad (47)$$

The complete intensity function is given by the summation of the pseudo-intensity functions with a ponderation by the first kind Bessel function product term $J_n(m)J_l(m)$.

This general approach is valid for all other sub-spaces like phase sub-spaces (x,k_x) . But the (x,t) sub-space choice is of great interest for the spatio-temporal coupling analysis of the impact of the spatial noise on the focal spot of the stretched and compressed pulses [25].

The last step consists of the projection of the (x,t) sub-space of both (x) and (t) spaces which give us respectively the pseudo-fluence and the pseudo-power of the beam (appendices E and F).

$$\begin{cases} F_{n,l}^{Seq}(x) = \exp \left[\begin{array}{l} C_x^{F_0} x^2 + K_d \left[(n+l)A_x^L + i(n-l)\Phi_x^L \right] x \\ + K_d^2 \left[(n+l)^2 A_{F_0}^S + (n-l)^2 A_{F_0}^D + i(n^2 - l^2)\Phi_0^F \right] \end{array} \right] \\ P_{n,l}^{Seq}(t) = \exp \left[\begin{array}{l} C_t^{P_0} t^2 + K_d \left[(n+l)A_t^L + i(n-l)\Phi_t^L \right] t \\ + K_d^2 \left[(n+l)^2 A_{P_0}^S + (n-l)^2 A_{P_0}^D + i(n^2 - l^2)\Phi_0^P \right] \end{array} \right] \end{cases} \quad (48)$$

All the described pseudo-functions depend on the sum $(n+l)$ and the difference $(n-l)$.

D. Reduction of the number of significant terms

The number of significant terms (n,l) in the sum in Eq. 10 depends on the modulation depth m . The development of the Bessel functions in power of m , in relation to the number of the limited (n,l) pairs of integers gives a good approximation of the Wigner function and its projections in the sub-spaces.

A sinusoidal phase modulation $\phi(x)=m \sin(K_d x)$ corresponds, by analogy with the instantaneous frequency in the time domain [26], to the spectrum:

$$K(x) = \frac{d\phi(x)}{dx} = m K_d \cos(K_d x) \quad (49)$$

with a maximum spectral excursion $\pm m K_d$ corresponding to a total spectrum broadening $\Delta K = 2m K_d$. So, the complete spectrum will be well described with at least $|n_l|, |l_l| \approx m$ spectral components.

6. APPLICATION TO THE TALBOT EFFECT

A. Wigner function and projections in the sub-spaces

Let us illustrate the sinusoidal phase mask in an elementary propagation sequence. It corresponds to the following transformation matrix according to Part 5.B, where z is the propagation distance and $k_0=2\pi/\lambda$ where λ is the central wavelength:

$$\begin{aligned} \mathbf{M}_\theta &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\ \mathbf{M}_l = \mathbf{M} &= \begin{pmatrix} 1 & 0 & -\frac{z}{k_0} & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \end{aligned} \quad (50)$$

The following pure quadratic terms C_{ij}^{W0} for the Gaussian Wigner function are unchanged through the transformation:

$$\begin{aligned} C_{xx}^{W_0} &= -\frac{2}{\Delta x^2} \\ C_{uu}^{W_0} &= -\frac{2}{\Delta t^2}, \\ C_{\omega\omega}^{W_0} &= -\frac{\Delta t^2}{2} \end{aligned} \quad (51)$$

Only one term $C_{k_x k_x}^{W0}$ is modified with an additional term in z :

$$C_{k_x k_x}^{W_0} = -\frac{\Delta x^2}{2} - \frac{2}{\Delta x^2} \times \frac{z^2}{k_0^2}. \quad (52)$$

All the cross-terms are null except $C_{xk_x}^{W0}$ and all the linear terms are null except:

$$\begin{aligned} C_{xk_x}^{W_0} &= \frac{2}{\Delta x^2} \times \frac{2z}{k_0} \\ A_{k_x}^L &= -\frac{\Delta x^2}{2}, \\ \Phi_x^L &= 1, \\ \Phi_{k_x}^L &= -\frac{z}{k_0} \end{aligned} \quad (53)$$

with the constant term:

$$A_0 = -\frac{\Delta x^2}{8}. \quad (54)$$

B. Expressions with dimensionless parameters

We introduce the dimensionless parameters $X=x/\Delta x$, $K_x=k_x\Delta x$, $u=K_d\Delta x$, and $Z=z/z_c$ with $z_c=k_0\Delta x^2$, to obtain the following Wigner function (after integration over t and ω , and neglecting the normalization terms):

$$W_{n,l}^{Seq}(X, K_x) = \exp \left[-2(X - ZK_x)^2 - \frac{1}{2} \left(K_x + \left(\frac{n+l}{2} \right) u \right)^2 + i(n-l)u(X - ZK_x) \right]. \quad (55)$$

The parameter $u/2\pi$ represents the number of fringes in the beam and z_c corresponds to the distance where diffraction effects on the spatial envelope of the beam are observed. Integration over the X variable gives the spatial spectrum of the beam, which is independent of the normalized variable Z :

$$W_{n,l}^{Seq}(K_x) = \exp \left[-\frac{1}{2} \left(K_x + \left(\frac{n+l}{2} \right) u \right)^2 - \frac{1}{8} (n-l)^2 u^2 \right]. \quad (56)$$

Each spectral contribution is a Gaussian centered around $K_x = -(n+l)/2u$ with the attenuation factor $\exp(-(n-l)^2 u^2/8)$. Integration of $W(X, K_x)$ over the K_x variable provides the fluence of the beam:

$$\begin{aligned} F_{n,l}^{Seq}(X) &= \exp \left[-\frac{2 \left[\left(X + \left(\frac{n+l}{2} \right) u \right)^2 + \frac{1}{4} (n-l)^2 u^2 Z^2 \right]}{1 + 4Z^2} \right] \\ &\times \exp \left[i(n-l)u \frac{\left(X + \left(\frac{n+l}{2} \right) u \right)}{1 + 4Z^2} \right]. \end{aligned} \quad (57)$$

Identification with the terms in Part 5.C gives:

$$\begin{aligned} C_{F_0}^L &\approx -\frac{2}{\Delta x^2} \times \frac{1}{1 + 4Z^2} \\ A_F^L &\approx -\frac{2Z}{1 + 4Z^2} \\ \Phi_F^L &\approx \frac{1}{1 + 4Z^2} \\ A_{F_0}^S &= A_{F_0}^D = -\frac{\Delta x^2}{2} \times \frac{Z^2}{1 + 4Z^2} \\ \Phi_0^F &\approx \frac{\Delta x^2}{2} \times \frac{Z}{1 + 4Z^2} \end{aligned} \quad (58)$$

Neglecting the diffraction effects ($4Z^2 \ll 1$) and assuming $uZ \ll 1$, provides a simplified formulation of the fluence:

$$\begin{aligned} F(X) &= \exp[-2X^2] \\ &\times \sum_{n=-\infty}^{+\infty} \sum_{l=-\infty}^{+\infty} J_n(m) J_l(m) \cos \left[(n-l)u \left(X + \left(\frac{n+l}{2} \right) u \right) \right]. \end{aligned} \quad (59)$$

Introducing the Talbot distance $z_t = 4\pi k_0 / K_d^2 = 4\pi z_c / u^2$ and the dimensionless parameter $Z_t = z/z_c$ allows us to get the identity $z/z_c = Z/Z_t$ and to obtain the usual expression of the Talbot effect [27]:

$$\begin{aligned} F(X) &= \exp[-2X^2] \\ &\times \sum_{n=-\infty}^{+\infty} \sum_{l=-\infty}^{+\infty} J_n(m) J_l(m) \cos \left[(n-l) \left(uX + 2\pi \frac{2Z}{Z_t} \left(\frac{n+l}{2} \right) \right) \right]. \end{aligned} \quad (60)$$

C. Explicit representation for five contributions

The expansion order of $J_n(m)$ and $J_l(m)$ must be adapted to the maximum value of $(n+l)$ to avoid pure imaginary residual terms and to respect the energy conservation. For five contributions, $(n+l) \in [-2, 2]$, the polynomial expansion of the Bessel function is limited to the second order in m . It gives:

$$\begin{aligned} J_0(m) &\approx 1 - \frac{1}{4} m^2 \\ J_1(m) &= -J_{-1}(m) \approx \frac{1}{2} m \\ J_2(m) &= J_{-2}(m) \approx \frac{1}{8} m^2 \end{aligned} \quad (61)$$

The different contributions (Tab.1) are given by the value of $n+l$ with $|n+l| < 2$ in respect to the second order development of $J_n(m)$. It

corresponds to the different spectral contributions centered at $K_x=(n+l)u/2$. According to Part 5.D, the development for five contributions is valid for $m \ll 2$. A practical analysis (Fig. 2) shows a good agreement with numerical calculation up to $m=0.5$. Higher values of modulation depth m need more than five contributions.

For these selected couples (n,l) , we have:

$$\sum_n \sum_l J_n J_l^* = 1. \quad (62)$$

For the energy conservation, the total Wigner function will be given by the sum of the five contributions:

$$W^{Seq}(X, K_x) = W_{n+l=0}(X, K_x) + W_{n+l=1}(X, K_x) + W_{n+l=-1}(X, K_x) + W_{n+l=2}(X, K_x) + W_{n+l=-2}(X, K_x) \quad (63)$$

Table 1. Wigner decomposition for five contributions.

(n,l)	$ J_n J_l $	Resulting Wigner function
$(0,0)$	$1 - \frac{m^2}{2}$	$W_{n+l=0}(X, K_x) =$
$(-1,1)$	$-\frac{m^2}{4}$	$\exp[-2(X - ZK_x)^2] \exp\left[-\frac{1}{2}K_x^2\right]$
$(1,-1)$	$-\frac{m^2}{4}$	$\times \left\{ 1 - \frac{m^2}{2} (1 + \cos[2u(X - ZK_x)]) \right\}$
$(0,1)$	$\frac{m}{2}$	$W_{n+l=1}(X, K_x) =$
$(1,0)$	$\frac{m}{2}$	$m \exp[-2(X - ZK_x)^2]$
$(-1,0)$	$-\frac{m}{2}$	$\times \exp\left[-\frac{1}{2}\left(K_x + \frac{u}{2}\right)^2\right] \cos[u(X - ZK_x)]$
$(0,-1)$	$-\frac{m}{2}$	$W_{n+l=-1}(X, K_x) =$
$(1,-1)$	$-\frac{m}{2}$	$-m \exp[-2(X - ZK_x)^2]$
$(0,-1)$	$-\frac{m}{2}$	$\times \exp\left[-\frac{1}{2}\left(K_x - \frac{u}{2}\right)^2\right] \cos[u(X - ZK_x)]$
$(1,1)$	$\frac{m^2}{4}$	$W_{n+l=2}(X, K_x) =$
$(0,2)$	$\frac{m^2}{8}$	$\frac{m^2}{4} \exp[-2(X - ZK_x)^2]$
$(2,0)$	$\frac{m^2}{8}$	$\times \exp\left[-\frac{1}{2}(K_x + u)^2\right]$
$(-1,-1)$	$\frac{m^2}{4}$	$\times \{1 + \cos[2u(X - ZK_x)]\}$
$(0,-2)$	$\frac{m^2}{8}$	$W_{n+l=-2}(X, K_x) =$
$(-2,0)$	$\frac{m^2}{8}$	$\frac{m^2}{4} \exp[-2(X - ZK_x)^2]$
		$\times \exp\left[-\frac{1}{2}(K_x - u)^2\right]$
		$\{1 + \cos[2u(X - ZK_x)]\}$

D. Phase-space domain analysis

The Wigner decomposition in Table 1 can be separated in:

- a common factor to the five contributions which gives the Gaussian envelope in X , shifted of a value $X_0 = ZK_x$,
- a contribution in $\exp(-K_x^2/2)$ which represents the contribution centered at $K_x=0$, with modulated (frequency $2K_d$) and not modulated parts,

- a modulated term in $m \cos(u(X - ZK_x))$ in factor of the two lateral contributions, centered in $(K_x = \pm u/2)$,
- Gaussian contributions in m^2 centered $(K_x = \pm u)$ with modulated (frequency $2K_d$) and not modulated parts.

Figure 1 represents the projection of the Wigner function in the phase space domain with $u=50$ for equivalent propagation distances of $Z=0$, $Z=0.1Z_t$, $Z=Z_t/4$ and $Z=Z_t/2$ where Z_t is the normalized Talbot distance $Z_t = z_t/z_c = 4\pi/\mu^2$ with $z_t = 4\pi k_0/K_d^2$. For a better visualization, we choose a modulation depth $m=0.5$.

On the left part and the bottom of each graphic, we have represented respectively the functions $W(K_x)$ and $W(X)=F(X)$ of the $W(X, K_x)$ function integrated over X and K_x respectively. In the dimensionless phase-space (X, K_x) , the Wigner function is composed of one central contribution (centered in $K_x=0$) and of two lateral contributions at frequency $K_d/2$, spatially modulated along X with alternatively positive and negative values [28,29]. The thin dashed lines represent the linear curves $X=0$ and $X=2\pi/u$. The thick dashed line represents the displacement of the maxima of the Wigner phase-space function on the lateral contributions. As mentioned in Part 3.A, the odd contributions $K_x = \pm u/2$ are composed of cross-terms only and have no energy while even contributions $K_x = 0, \pm u$ are a combination of self and cross-Wigner terms with energy given by the self-terms only.

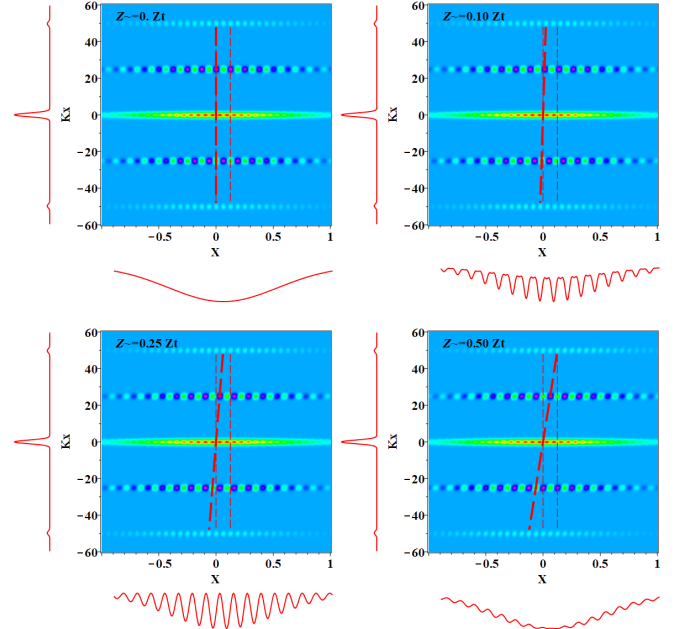


Fig. 1. Phase-space representation of the Talbot effect with five contributions for Gaussian beams for $u=50$ and $m=0.5$. On the left of each figure is represented the spatial spectrum (integration over X), and on the bottom, is represented the spatial profile (integration over K_x). The case $Z=0$ represents the beam spatially-phase modulated with no propagation (no spatial fluence modulations). The case $Z=0.1Z_t$ represents a propagation over a fraction of the Talbot distance. The case $Z=0.25Z_t$ presents a maximal spatial modulation, when the shift of the periodic is half period. When the shift is $Z=0.5Z_t$, modulations are minimized but not cancelled because the finite size of the beam.

For $Z=0$, the sum of the two lateral contributions, obtained by integration on K_x variable, is equal to zero because they are exactly in opposition of phase. The effect of the propagation along z is to shift along X the lateral contributions, but in opposite directions.

For $Z=Z_t/4$ modulation are in phase and the integration along K_x of the two lateral contributions gives an extremum of modulation along X .

For $Z=Z_i/2$ the modulations are again in opposition and the sum of the lateral contributions becomes zero again. Nevertheless, modulations are not cancelled due to the finite size of the envelope of the beam. The propagation along z implies a displacement of the lateral contributions and also an inclination of each envelope contribution.

If we operate the change of variable $(X, K_x) \rightarrow (X', K_x)$ with:

$$X' = X - ZK_x, \quad (64)$$

the Wigner phase-space function becomes invariant with Z in this new coordinate system.

E. Spectral analysis

Integration of the Wigner function in the spatial domain gives the spatial spectrum of the beam (without the common normalization factor of each contribution):

$$\begin{cases} W_{n+l=0}(K_x) = \exp\left[-\frac{1}{2}K_x^2\right] \left(1 - \frac{m^2}{2} \left(1 + \exp\left[-\frac{u^2}{2}\right]\right)\right) \\ W_{n+l=1}(K_x) = m \exp\left[-\frac{u^2}{8}\right] \exp\left[-\frac{1}{2}\left(K_x + \frac{u}{2}\right)^2\right] \\ W_{n+l=-1}(K_x) = -m \exp\left[-\frac{u^2}{8}\right] \exp\left[-\frac{1}{2}\left(K_x - \frac{u}{2}\right)^2\right] \\ W_{n+l=2}(K_x) = \frac{m^2}{4} \left\{1 + \exp\left[-\frac{u^2}{2}\right]\right\} \exp\left[-\frac{1}{2}(K_x + u)^2\right] \\ W_{n+l=-2}(K_x) = \frac{m^2}{4} \left\{1 + \exp\left[-\frac{u^2}{2}\right]\right\} \exp\left[-\frac{1}{2}(K_x - u)^2\right] \end{cases} \quad (65)$$

The spatial spectrum is independent of the normalized propagation distance Z . Assuming $u^2 \gg 1$ (i.e. more than one fringe in the beam), we get $\exp(-u^2) \approx 0$ and terms in m disappear because they correspond to odd terms. The expression reduces to the three Gaussian contributions centered at $K_x=0$ (signal without modulations, proportional to $W_0(K_x)$) and $K_x=\pm u$ (noise), with a weight in m^2 . This result was expected because a convergent description of the beam requires five contributions, with a development at the second order with the modulation depth m (terms in m^2):

$$W^{Seq}(K_x) \approx \left\{ \begin{array}{l} \exp\left[-\frac{1}{2}K_x^2\right] \left(1 - \frac{m^2}{2}\right) \\ + \frac{m^2}{4} \left\{ \exp\left[-\frac{1}{2}(K_x - u)^2\right] + \exp\left[-\frac{1}{2}(K_x + u)^2\right] \right\} \end{array} \right\} \quad (66)$$

This expression gives a good description of the spatial spectrum of the beam with central and lateral contributions.

F. Fluence expression for five contributions

Integration of the Wigner function over the spectral coordinate K_x gives the fluence of the beam. From Table 1, integration of the five contributions of the Wigner function, without approximations and without normalization factors, gives the five contributions of the fluence function:

$$\begin{cases} F_{n+l=0}(X) = \exp\left[-\frac{2X^2}{1+4Z^2}\right] \\ \quad \times \left\{1 - \frac{m^2}{2} \left(1 + \exp\left[-\frac{2u^2Z^2}{1+4Z^2}\right] \cos\left(\frac{2Xu}{1+4Z^2}\right)\right)\right\} \\ F_{n+l=1}(X) = m \exp\left[-\frac{2\left(X + \frac{u}{2}Z\right)^2 + \frac{u^2}{2}Z^2}{1+4Z^2}\right] \cos\left(\frac{u\left(X + \frac{u}{2}Z\right)}{1+4Z^2}\right) \\ F_{n+l=-1}(X) = -m \exp\left[-\frac{2\left(X - \frac{u}{2}Z\right)^2 + \frac{u^2}{2}Z^2}{1+4Z^2}\right] \cos\left(\frac{u\left(X - \frac{u}{2}Z\right)}{1+4Z^2}\right) \\ F_{n+l=2}(X) = \frac{m^2}{4} \exp\left[-\frac{2(X+uZ)^2}{1+4Z^2}\right] \\ \quad \times \left\{1 + \exp\left[-\frac{2u^2Z^2}{1+4Z^2}\right] \cos\left[\frac{2u(X+uZ)}{1+4Z^2}\right]\right\} \\ F_{n+l=-2}(X) = \frac{m^2}{4} \exp\left[-\frac{2(X-uZ)^2}{1+4Z^2}\right] \\ \quad \times \left\{1 + \exp\left[-\frac{2u^2Z^2}{1+4Z^2}\right] \cos\left[\frac{2u(X-uZ)}{1+4Z^2}\right]\right\} \end{cases} \quad (67)$$

with a total fluence:

$$F^{Seq}(X) = F_{n+l=0}(X) + F_{n+l=1}(X) + F_{n+l=-1}(X) + F_{n+l=2}(X) + F_{n+l=-2}(X). \quad (68)$$

This expression contains all the information concerning the Talbot effect coupled with finite pupil and diffraction effects. The fluence formulation consists of three Gaussian waves centered at $X_0=0, \pm uZ$ interacting together with different behavior in different areas [30].

Figure 2 illustrates the evolution of the fluence $F(X)$ with Z (at the left) obtained with the decomposition in Table 2 and the numerical calculation (Beam Propagation Method with 2048×200 points [31]) for a Gaussian beam evolution with Z (at the right) after a phase mask $\Phi(X) = m \sin(u)$ for $u=50$ and $m=0.5$. For the numerical calculus, we bring $\Delta x = 0.03m$ and $K_0 = 1666m^{-1}$ in the (x, z) coordinates and translate it in (X, Z) coordinates for the comparison with the Bessel decomposition.

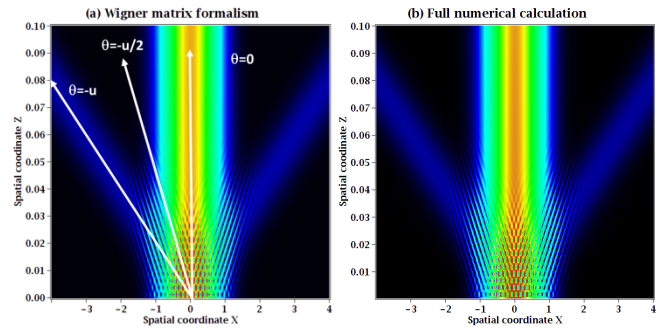


Fig. 2. Evolution of the normalized fluence with Z for a Gaussian beam: a) Bessel decomposition with five contributions and b) numerical calculus (Beam Propagation Method) for $u=50$ and $m=0.5$. Normalized fluence in linear scale.

We can observe in Figure 2 that the Wigner matrix formulation catches all the optical features of the phase modulated beam and matches very accurately with the full numerical calculation. For larger modulation depths, more contributions in the Bessel expansion will need to be taken into account. In relation with Fig. 1, we identify the even contributions characterized by the three beams carrying the energy in the directions $\theta=0,\pm u$, while the odd contributions can be observed in the directions $\theta =\pm u/2$ and characterized by fringes direction.

For an infinite beam, and neglecting the diffraction effects ($4Z^2 \ll 1$), the spatial evolution of the beam reduces to the standard Talbot formulation:

$$F(X) = 1 - 2m \sin(uX) \sin\left(\frac{u^2 Z}{2}\right) = 1 - 2m \sin(uX) \sin\left(2\pi \frac{Z}{Z_t}\right). \quad (69)$$

G. Comparison between Bessel decomposition and GSM beam

Let us consider a beam with a noise in phase and amplitude. This noise is characterized by a modulation depth P_m^2 and a coherence parameter p (see section 4A) given by a *GSM* model. The total Wigner function is given by:

$$W(x, k_x) = (1 - P_m^2) W_0(x, k_x) + P_m^2 W_{GSM}(x, k_x), \quad (70)$$

with a weight $(1 - P_m^2)$ on W_0 contribution to preserve the total energy E_t of the beam. With the dimensionless parameters X , K_x and Z introduced in part 6B, we have, for the propagation of the noise over a distance Z [11]:

$$W_{GSM}(X, K_x) = \frac{E_t p}{\pi} \exp\left[-2(X - ZK_x)^2 - \frac{1}{2} p^2 K_x^2\right]. \quad (71)$$

The *GSM* beam corresponds to a Gaussian statistical distribution $W(K_x)$ of the noise in the spatial frequency domain and is independent of the propagation distance Z :

$$W_{GSM}(K_x) = \int_{-\infty}^{+\infty} W_{GSM}(X, K_x) dX = \frac{E_t p \Delta x}{\sqrt{2\pi}} \exp\left[-\frac{1}{2} p^2 K_x^2\right]. \quad (72)$$

In the spatial domain:

$$W_{GSM}(X) = \int_{-\infty}^{+\infty} W_{GSM}(X, K_x) dK_x = \sqrt{\frac{2}{\pi}} \frac{E_t}{\sqrt{1 + \frac{4Z^2}{p^2}}} \exp\left[-\frac{2X^2}{1 + \frac{4Z^2}{p^2}}\right]. \quad (73)$$

When neglecting diffraction ($Z^2 \ll p^2/4$), $W_{GSM}(X)$ becomes also independent of z :

$$W_{GSM}(X) = \int_{-\infty}^{+\infty} W_{GSM}(X, K_x) dK_x \approx \sqrt{\frac{2}{\pi}} E_t \exp[-2X^2]. \quad (74)$$

To compare the Bessel decomposition with the *GSM* beam, we need to add a sufficient number of sinusoidal phase masks with different frequencies u . For each frequency u , the noise part is given by the two lateral frequencies in the spectrum:

$$W_{noise}(K_x, u, m) = W(K_x, u, m) - \left(1 - \frac{m^2}{2}\right) W_0(K_x). \quad (75)$$

The weight of the spectrum with no modulation ($1 - m^2/2$) in the Bessel decomposition is to compare with the weight of the spectrum with no modulations in the *GSM* model ($1 - P_m^2$) and gives the following relation for one frequency u :

$$P_m^2 = \frac{m^2}{2}. \quad (76)$$

The expression in phase-space domain is then, for N noise contributions at frequency $u[i]$ and depth modulation $m(u[i])$:

$$W_{noise}(X, K_x) = \sum_{i=1}^N W(X, K_x, u[i], m(u[i])) - W_0(X, K_x) \sum_{i=1}^N \left(1 - \frac{m(u[i])^2}{2}\right). \quad (77)$$

with the relation between the *GSM* model and multi-frequencies Bessel decomposition:

$$P_m^2 \approx \sum_{i=1}^N \frac{m(u[i])^2}{2}. \quad (78)$$

We need to have, for each frequency $u[i]$, a modulation depth following the same distribution noise as the *GSM* beam:

$$m(u[i]) = P_0 \exp\left[-\frac{1}{2} p^2 u[i]^2\right]. \quad (79)$$

where the $u[i]$ frequencies are randomly distributed over the frequency domain. For each frequency, we describe a phase modulation $\sin(u[i]X + \Phi[i])$ where $\Phi[i]$ is randomly distributed between 0 and 2π . For a given factor P_0 , equaling the noise energy for the *GSM* beam and Bessel decomposition gives the value of P_m . For the *GSM* beam, we choose $E_t=1$ and $p=0.05$, which gives a total spectrum width of $\Delta K_x=47$ (FWHM). The representation in Figure 3 is made with $N=10$ different frequencies for the Bessel decomposition (Tab 2.):

Table 2. Random choice following the Gaussian statistics of the GSM beam for the corresponding values of $m(u)$

$u[i]$ (m^{-1})	3.1	21.6	32.9	33.5	34.0	43.5	46.0	53.48	57.7	59.6
$m(u[i])$ (x100)	9.93	7.46	5.08	4.95	4.48	3.05	2.65	1.67	1.24	1.08
$\Phi[i]$ (rad)	5.69	2.22	3.76	0.455	2.06	0.27	1.61	4.79	1.24	5.26

For $P_0=0.1$, we find $P_m=0.11$ with the parameters of Table 2. Figure 3 shows the Wigner function in the dimensionless (X, K_x) phase-space domain, in log scale, for the Bessel decomposition (positive part only) and the *GSM* model, for a propagation distance $Z=0$ (upper part) and $Z=2Z_t$ (two times the dimensionless Talbot distance Z_t) for $u=50$ (lower part). The low frequency distribution allows the observation of the substructures of the Bessel decomposition. If the *GSM* beam is a good model for describing the envelope of the beam (e.g. the rotation of the envelope with z), it cannot model fine local behaviors.

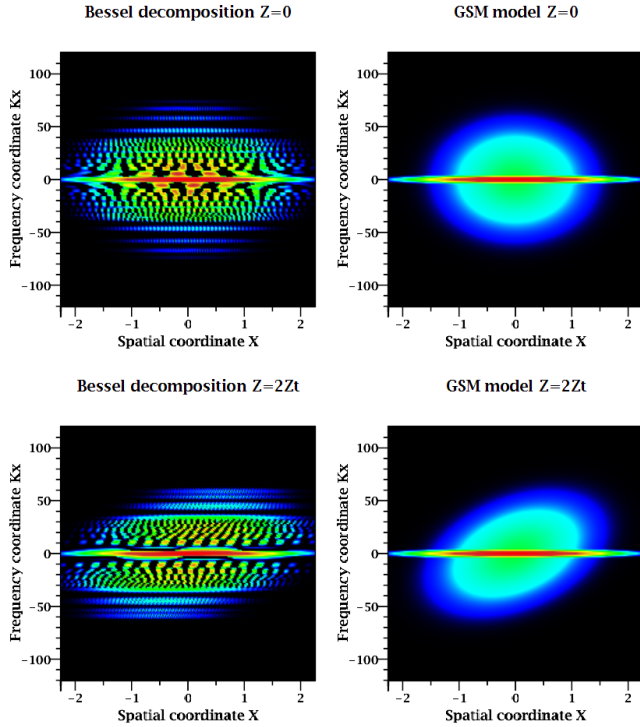


Fig. 3. Wigner representation of the Wigner function $W(x, k_x)$ in the phase-space for 10 different spatial frequencies K_d randomly given with the modulation depth $m(K_d)$ following a Gaussian distribution given by the GSM model (log scale).

Figure 4(a) shows the corresponding spatial envelope when integrating the *GSM* beam over k_x (blue), compared with the Bessel decomposition (dash red). For $Z=0$ the envelopes are superimposed, but for $z=2Z_t$ (Z_t is given for $u=50$ i.e. at the bottom of the total distribution spectrum $\Delta K_x=47$ FWHM), the Bessel decomposition gives spatial modulations which are not predicted by the *GSM* model, because the mean value of the modulations remains the same.

Figure 4(b) shows the frequency spectrum given when integrating the *GSM* beam over the x coordinate (in blue), compared with the Bessel decomposition (dash red) in log scale. The central contribution corresponds to the part of the signal which is not disturbed by the noise. In both cases, the spectrum is invariant with Z .

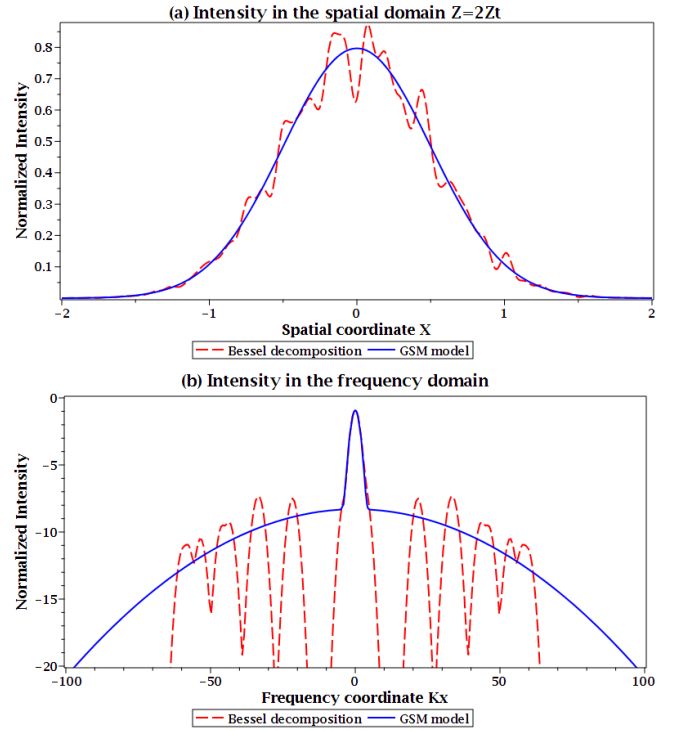


Fig. 4. (a). Spatial envelope of the beam obtained by the integration of the Wigner functions over the spatial frequencies (full blue line: GSM model; dashed red line: Bessel decomposition). (b). Spatial spectra obtained by the integration of the Wigner functions over the spatial coordinate (full blue line: GSM model; dashed red line: Bessel decomposition).

7. CONCLUSIONS

To conclude, we have developed an analytic Wigner matrix formalism for describing the propagation of optical beams modulated in phase. We have coupled this formalism with a complete *ABCD* matrix sequence to apply this model to complex optical chains. A matrix Wigner formalism allows to represent the evolution of the second order moments of an arbitrary beam.

We illustrated the interest and versatility of this model by considering the textbook case of the Talbot effect. In particular, we proved that the Wigner matrix model is able to retrieve all the features of the diffracted beam and a comparison with full numerical calculations evidenced its accuracy. Let us emphasize that modeling the propagation of phase modulated beam, with a finite size envelope, is very challenging since it is time consuming and requires important computing resources. This analytic formalism was shown to retrieve the same results as those obtained with full numerical calculation. In a second step, we compared this model with another Wigner model, the so-called Gaussian Shell model and showed that the Wigner Matrix method retrieves more features of the diffracted beam, in particular, oscillations of the envelope of the beam in both spatial and frequency coordinates. This model can be easily extended to phase modulations in the three spatial dimensions, with an interest in polarized beams and coupling effects between the spatial components.

Besides its interest in the spatial domain, this formalism will find many applications for describing beam properties in the time domain. Coupling description in both spatial and time domains will be of particular interest to model the spatio-temporal coupling in laser chains. For example, in *CPA* systems, taking into account the spatio-temporal coupling in the stretcher and the compressor allow to

evaluate the amplitude of spatial modulations in the near field (for optical damages) and the spatial and temporal contrast in the far field (focal plane).

Let us emphasize that this model is not limited to weak modulation depths and that strong modulations can be modeled at the price of an expansion of the Bessel functions to higher orders. This formalism turns out to be of high interest for modeling the spatio-temporal coupling in *LMJ* (Laser Mégajoules) including the strong modulation depth of the phase plate on the spatial domain ($\Delta x \approx 1m$, $m \approx 1rad$, $K_d \approx 1mm^{-1}$) and the strong modulation depth in the temporal domain ($\Delta t \approx 1ns$, $m \approx 7rad$, $f = 14GHz$) can help to analyze the smoothing properties at the focal plane.

APPENDIX

A. Polynomial coefficients of the general Gaussian-Wigner formulation without modulation

$$\begin{aligned} \mathbf{M}_0 &= \begin{pmatrix} p_{11} & p_{12} & p_{13} & p_{14} \\ p_{21} & p_{22} & p_{23} & p_{24} \\ p_{31} & p_{32} & p_{33} & p_{34} \\ p_{41} & p_{42} & p_{43} & p_{44} \end{pmatrix}, \\ \mathbf{M}_1 &= \begin{pmatrix} q_{11} & q_{12} & q_{13} & q_{14} \\ q_{21} & q_{22} & q_{23} & q_{24} \\ q_{31} & q_{32} & q_{33} & q_{34} \\ q_{41} & q_{42} & q_{43} & q_{44} \end{pmatrix}, \end{aligned} \quad (\text{A1})$$

with a total *ABCD* sequence

$$\mathbf{M} = \mathbf{M}_0 \mathbf{M}_1 = \begin{pmatrix} s_{11} & s_{12} & s_{13} & s_{14} \\ s_{21} & s_{22} & s_{23} & s_{24} \\ s_{31} & s_{32} & s_{33} & s_{34} \\ s_{41} & s_{42} & s_{43} & s_{44} \end{pmatrix}. \quad (\text{A2})$$

$$W_0^{Seq}(x, t, k_x, \omega) = \exp \left[\begin{aligned} & C_{xx}^{W_0} x^2 + C_{tt}^{W_0} t^2 + C_{k_x k_x}^{W_0} k_x^2 + C_{\omega \omega}^{W_0} \omega^2 \\ & + C_{xk_x}^{W_0} x k_x + C_{xt}^{W_0} x t + C_{x\omega}^{W_0} x \omega \\ & + C_{k_x t}^{W_0} k_x t + C_{k_x \omega}^{W_0} k_x \omega + C_{t\omega}^{W_0} t \omega \end{aligned} \right]. \quad (\text{A3})$$

The C_{ij} ($ij=x,t,k_x,\omega$) polynomial coefficients of the general Gaussian-Wigner formulation without modulation are related to the matrix coefficients s_{ij} ($ij=1..4$) of the total *ABCD* matrix sequence \mathbf{M} and the matrix coefficients of \mathbf{Q}_0 (functions of Δx and Δt):

1. Pure quadratic terms in (x, t, k_x, ω)

$$\begin{aligned} C_{xx}^{W_0} &= - \left(\frac{2}{\Delta x^2} s_{11}^2 + \frac{2}{\Delta t^2} s_{21}^2 + \frac{\Delta x^2}{2} s_{31}^2 + \frac{\Delta t^2}{2} s_{41}^2 \right) \\ C_{tt}^{W_0} &= - \left(\frac{2}{\Delta x^2} s_{12}^2 + \frac{2}{\Delta t^2} s_{22}^2 + \frac{\Delta x^2}{2} s_{32}^2 + \frac{\Delta t^2}{2} s_{42}^2 \right) \\ C_{k_x k_x}^{W_0} &= - \left(\frac{2}{\Delta x^2} s_{13}^2 + \frac{2}{\Delta t^2} s_{23}^2 + \frac{\Delta x^2}{2} s_{33}^2 + \frac{\Delta t^2}{2} s_{43}^2 \right) \\ C_{\omega \omega}^{W_0} &= - \left(\frac{2}{\Delta x^2} s_{14}^2 + \frac{2}{\Delta t^2} s_{24}^2 + \frac{\Delta x^2}{2} s_{34}^2 + \frac{\Delta t^2}{2} s_{44}^2 \right) \end{aligned} \quad (\text{A4})$$

2. Cross-terms

$$\begin{aligned} C_{xt}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{12} s_{11} + \frac{2}{\Delta t^2} s_{22} s_{21} + \frac{\Delta x^2}{2} s_{32} s_{31} + \frac{\Delta t^2}{2} s_{42} s_{41} \right) \\ C_{k_x k_x}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{13} s_{11} + \frac{2}{\Delta t^2} s_{23} s_{21} + \frac{\Delta x^2}{2} s_{33} s_{31} + \frac{\Delta t^2}{2} s_{43} s_{41} \right) \\ C_{x\omega}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{14} s_{11} + \frac{2}{\Delta t^2} s_{24} s_{21} + \frac{\Delta x^2}{2} s_{34} s_{31} + \frac{\Delta t^2}{2} s_{44} s_{41} \right) \\ C_{k_x t}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{13} s_{12} + \frac{2}{\Delta t^2} s_{23} s_{22} + \frac{\Delta x^2}{2} s_{33} s_{32} + \frac{\Delta t^2}{2} s_{43} s_{42} \right) \\ C_{t\omega}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{14} s_{12} + \frac{2}{\Delta t^2} s_{24} s_{22} + \frac{\Delta x^2}{2} s_{34} s_{32} + \frac{\Delta t^2}{2} s_{44} s_{42} \right) \\ C_{k_x \omega}^{W_0} &= -2 \left(\frac{2}{\Delta x^2} s_{14} s_{13} + \frac{2}{\Delta t^2} s_{24} s_{23} + \frac{\Delta x^2}{2} s_{34} s_{33} + \frac{\Delta t^2}{2} s_{44} s_{43} \right) \end{aligned} \quad (\text{A5})$$

These coupling coefficients are real but not necessary positives.

B. Polynomial coefficient of the complementary part of the Wigner function

The part of the Wigner function which is affected by the phase modulation takes the following general form:

$$\begin{aligned} W_{n,l}^{Seq}(x, t, k_x, \omega) &= W_0^{Seq}(x, t, k_x, \omega) \\ &\times \exp \left[\begin{aligned} & (n+l) K_d (A_x^L x + A_t^L t + A_{k_x}^L k_x + A_\omega^L \omega) \\ & + i(n-l) K_d (\Phi_x^L x + \Phi_t^L t + \Phi_{k_x}^L k_x + \Phi_\omega^L \omega) \\ & + (n+l)^2 K_d^2 A_0 \end{aligned} \right]. \end{aligned} \quad (\text{B1})$$

The A^L and Φ^L ($j=x,t,k_x,\omega$) polynomial coefficients of the complementary part of the Wigner function due to the phase modulation are related to the matrix coefficients s_{ij} and p_{ij} ($ij=1..4$) *ABCD* matrix sequences \mathbf{M} and \mathbf{M}_0 and the matrix coefficients of \mathbf{Q}_0 :

1. A^L coefficients of the real linear terms in (x, t, k_x, ω) :

$$\begin{aligned} A_x^L &= - \left(\frac{2}{\Delta x^2} p_{13} s_{11} + \frac{2}{\Delta t^2} p_{23} s_{21} + \frac{\Delta x^2}{2} p_{33} s_{31} + \frac{\Delta t^2}{2} p_{43} s_{41} \right) \\ A_t^L &= - \left(\frac{2}{\Delta x^2} p_{13} s_{12} + \frac{2}{\Delta t^2} p_{23} s_{22} + \frac{\Delta x^2}{2} p_{33} s_{32} + \frac{\Delta t^2}{2} p_{43} s_{42} \right) \\ A_{k_x}^L &= - \left(\frac{2}{\Delta x^2} p_{13} s_{13} + \frac{2}{\Delta t^2} p_{23} s_{23} + \frac{\Delta x^2}{2} p_{33} s_{33} + \frac{\Delta t^2}{2} p_{43} s_{43} \right) \\ A_\omega^L &= - \left(\frac{2}{\Delta x^2} p_{13} s_{14} + \frac{2}{\Delta t^2} p_{23} s_{24} + \frac{\Delta x^2}{2} p_{33} s_{34} + \frac{\Delta t^2}{2} p_{43} s_{44} \right) \end{aligned} \quad (\text{B2})$$

2. Φ^L coefficients of the imaginary linear terms in (x, t, k_x, ω) :

$$\begin{aligned} \Phi_x^L &= q_{11} \\ \Phi_t^L &= q_{12} \\ \Phi_{k_x}^L &= q_{13} \\ \Phi_\omega^L &= q_{14} \end{aligned} \quad (\text{B3})$$

3. Coefficient of the constant term:

$$A_0 = - \frac{1}{4} \left(\frac{2}{\Delta x^2} p_{13}^2 + \frac{2}{\Delta t^2} p_{23}^2 + \frac{\Delta x^2}{2} p_{33}^2 + \frac{\Delta t^2}{2} p_{43}^2 \right). \quad (\text{B4})$$

From these explicit coefficients, we can obtain the Wigner projection in all the sub-spaces.

C. Polynomial coefficients of the Intensity without modulation

$$I_0^{Seq}(x,t) = \exp\left[C_{xx}^{l_0}x^2 + C_{tt}^{l_0}t^2 + C_{xt}^{l_0}xt\right], \quad (\text{C1})$$

with:

$$\begin{cases} C_{xx}^{l_0}, C_{tt}^{l_0} < 0 \\ 4C_{xx}^{l_0}C_{tt}^{l_0} - (C_{xt}^{l_0})^2 > 0 \end{cases} \quad (\text{C2})$$

The expressions of the pure quadratic terms $C_{xx}^{l_0}$ and $C_{tt}^{l_0}$ and the cross-term $C_{xt}^{l_0}$ of the intensity $I_0^{Seq}(x,y)$ of the field, not affected by the phase modulation, are related to the C_{ij}^{W0} coefficients:

$$\begin{cases} C_{xx}^{l_0} = \frac{1}{g^2} \left[4C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} C_{xx}^{W_0} + C_{x\omega}^{W_0} C_{xk_x}^{W_0} C_{k_x\omega}^{W_0} - C_{xx}^{W_0} (C_{k_x\omega}^{W_0})^2 - C_{\omega\omega}^{W_0} (C_{xk_x}^{W_0})^2 - C_{k_x k_x}^{W_0} (C_{x\omega}^{W_0})^2 \right] \\ C_{tt}^{l_0} = \frac{1}{g^2} \left[4C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} C_{tt}^{W_0} + C_{t\omega}^{W_0} C_{tk_x}^{W_0} C_{k_x\omega}^{W_0} - C_{tt}^{W_0} (C_{k_x\omega}^{W_0})^2 - C_{\omega\omega}^{W_0} (C_{tk_x}^{W_0})^2 - C_{k_x k_x}^{W_0} (C_{t\omega}^{W_0})^2 \right] \\ C_{xt}^{l_0} = \frac{1}{g^2} \left[-4C_{xt}^{W_0} C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} - C_{t\omega}^{W_0} C_{xk_x}^{W_0} C_{k_x\omega}^{W_0} + C_{xt}^{W_0} (C_{k_x\omega}^{W_0})^2 + 2C_{xk_x}^{W_0} C_{k_x\omega}^{W_0} C_{\omega\omega}^{W_0} + 2C_{x\omega}^{W_0} C_{k_x k_x}^{W_0} C_{t\omega}^{W_0} - C_{x\omega}^{W_0} C_{k_x\omega}^{W_0} C_{k_x k_x}^{W_0} \right] \end{cases} \quad (\text{C3})$$

with:

$$g = \sqrt{4C_{\omega\omega}^{W_0} C_{k_x k_x}^{W_0} - (C_{k_x\omega}^{W_0})^2}, \quad (\text{C4})$$

C_{xx}^{W0} and C_{tt}^{W0} are negative, and C_{xt}^{W0} is real with:

$$4C_{\omega\omega}^{W_0} C_{k_x k_x}^{W_0} > (C_{k_x\omega}^{W_0})^2. \quad (\text{C5})$$

D. Polynomial coefficient of the complementary part of the intensity

$$I_{n,l}^{Seq}(x,t) = I_0^{Seq}(x,t) \exp \left[\begin{aligned} & (n+l)K_d [A_{lx}^L x + A_{lt}^L t] \\ & + i(n-l)K_d [\Phi_{lx}^L x + \Phi_{lt}^L t] \\ & + K_d^2 [(n+l)^2 A_{F0}^S + (n-l)^2 A_{F0}^D] \\ & + i(n^2 - l^2)K_d^2 \Phi_0^I \end{aligned} \right]. \quad (\text{D1})$$

For the real coefficients of the linear terms in (x,t) , we have:

$$\begin{cases} A_{lx}^L = \frac{1}{g^2} \left[C_{k_x\omega}^{W_0} A_{lx}^L C_{x\omega}^{W_0} + C_{k_x\omega}^{W_0} A_{lx}^L C_{xk_x}^{W_0} - 2C_{k_x k_x}^{W_0} A_{lx}^L C_{x\omega}^{W_0} + 4C_{k_x k_x}^{W_0} A_{lx}^L C_{\omega\omega}^{W_0} - 2C_{\omega\omega}^{W_0} A_{lx}^L C_{xk_x}^{W_0} - (C_{k_x\omega}^{W_0})^2 A_{lx}^L \right] \\ \Phi_{lx}^L = \frac{1}{g^2} \left[\Phi_{lx}^L C_{k_x\omega}^{W_0} C_{x\omega}^{W_0} + \Phi_{lx}^L C_{k_x\omega}^{W_0} C_{xk_x}^{W_0} - 2\Phi_{lx}^L C_{\omega\omega}^{W_0} C_{xk_x}^{W_0} + 4\Phi_{lx}^L C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} - 2\Phi_{lx}^L C_{\omega\omega}^{W_0} C_{x\omega}^{W_0} - (C_{k_x\omega}^{W_0})^2 \Phi_{lx}^L \right] \end{cases} \quad (\text{D2})$$

for the coefficients proportional to x , and:

$$\begin{cases} A_{lt}^L = \frac{1}{g^2} \left[C_{k_x\omega}^{W_0} A_{lt}^L C_{t\omega}^{W_0} + C_{k_x\omega}^{W_0} A_{lt}^L C_{tk_x}^{W_0} - 2C_{k_x k_x}^{W_0} A_{lt}^L C_{t\omega}^{W_0} + 4C_{k_x k_x}^{W_0} A_{lt}^L C_{\omega\omega}^{W_0} - 2C_{\omega\omega}^{W_0} A_{lt}^L C_{tk_x}^{W_0} - (C_{k_x\omega}^{W_0})^2 A_{lt}^L \right] \\ \Phi_{lt}^L = \frac{1}{g^2} \left[\Phi_{lt}^L C_{k_x\omega}^{W_0} C_{t\omega}^{W_0} + \Phi_{lt}^L C_{k_x\omega}^{W_0} C_{tk_x}^{W_0} - 2\Phi_{lt}^L C_{\omega\omega}^{W_0} C_{tk_x}^{W_0} + 4\Phi_{lt}^L C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} - 2\Phi_{lt}^L C_{\omega\omega}^{W_0} C_{t\omega}^{W_0} - (C_{k_x\omega}^{W_0})^2 \Phi_{lt}^L \right] \end{cases} \quad (\text{D3})$$

for the coefficients proportional to t . The constant terms are respectively:

$$\begin{cases} A_{F0}^S = \frac{1}{g^2} \left[\left(A_{lx}^L A_{lx}^L C_{k_x\omega}^{W_0} - (A_{lx}^L)^2 C_{k_x k_x}^{W_0} - (A_{lx}^L)^2 C_{\omega\omega}^{W_0} \right) + 4A_0 \left(4C_{k_x k_x}^{W_0} C_{\omega\omega}^{W_0} - (C_{k_x\omega}^{W_0})^2 \right) \right] \\ A_{F0}^D = \frac{1}{g^2} \left[(\Phi_{lx}^L)^2 C_{\omega\omega}^{W_0} - \Phi_{lx}^L \Phi_{lx}^L C_{k_x\omega}^{W_0} + (\Phi_{lx}^L)^2 C_{k_x k_x}^{W_0} \right] \\ \Phi_0^I = \frac{1}{g^2} \left[\Phi_{lx}^L A_{lx}^L C_{k_x\omega}^{W_0} + \Phi_{lx}^L A_{lx}^L C_{k_x\omega}^{W_0} - 2\Phi_{lx}^L A_{lx}^L C_{\omega\omega}^{W_0} - 2\Phi_{lx}^L A_{lx}^L C_{k_x k_x}^{W_0} \right] \end{cases} \quad (\text{D4})$$

E. Fluence coefficients

$$F_{n,l}^{Seq}(x) = \exp \left[\begin{aligned} & C_{F0}^l x^2 + K_d ((n+l)A_F^L) \\ & + i(n-l)\Phi_F^L x \\ & + K_d^2 ((n+l)^2 A_{F0}^S + (n-l)^2 A_{F0}^D) \\ & + i(n^2 - l^2)\Phi_0^F \end{aligned} \right], \quad (\text{E1})$$

with:

$$\begin{cases} C_{F0}^l = C_{xx}^{l_0} - \frac{(C_{xt}^{l_0})^2}{C_{tt}^{l_0}} \\ A_F^L = A_{lx}^L - A_{lt}^L \frac{C_{xt}^{l_0}}{C_{tt}^{l_0}} \\ \Phi_F^L = \Phi_{lx}^L - \Phi_{lt}^L \frac{C_{xt}^{l_0}}{C_{tt}^{l_0}} \\ A_{F0}^S = A_{F0}^S - \frac{(A_{lx}^L)^2}{4C_{tt}^{l_0}} \\ A_{F0}^D = A_{F0}^D + \frac{(\Phi_{lx}^L)^2}{4C_{tt}^{l_0}} \\ \Phi_0^F = \Phi_0^I - \frac{A_{lx}^L \Phi_{lx}^L}{2C_{tt}^{l_0}} \end{cases} \quad (\text{E2})$$

F. Power coefficients

$$P_{n,l}^{Seq}(t) = \exp \left[\begin{aligned} & C_{P0}^l t^2 + K_d ((n+l)A_P^L) \\ & + i(n-l)\Phi_P^L t \\ & + K_d^2 ((n+l)^2 A_{P0}^S + (n-l)^2 A_{P0}^D) \\ & + i(n^2 - l^2)\Phi_0^P \end{aligned} \right], \quad (\text{F1})$$

with:

$$\begin{cases} C_{P0}^l = C_{tt}^{l_0} - \frac{(C_{xt}^{l_0})^2}{C_{xx}^{l_0}} \\ A_P^L = A_{lt}^L - A_{lx}^L \frac{C_{xt}^{l_0}}{C_{xx}^{l_0}} \\ \Phi_P^L = \Phi_{lt}^L - \Phi_{lx}^L \frac{C_{xt}^{l_0}}{C_{xx}^{l_0}} \\ A_{P0}^S = A_{P0}^S - \frac{(A_{lx}^L)^2}{4C_{xx}^{l_0}} \\ A_{P0}^D = A_{P0}^D + \frac{(\Phi_{lx}^L)^2}{4C_{xx}^{l_0}} \\ \Phi_0^P = \Phi_0^I - \frac{A_{lx}^L \Phi_{lx}^L}{2C_{xx}^{l_0}} \end{cases} \quad (\text{F2})$$

References

1. E. Wigner, "On the quantum correction for thermodynamic equilibrium," *Phys. Rev.* **40**, 749 – 759 (1932)
2. M. J. Bastiaans, "Wigner distribution function and its application to first-order optics," *J. Opt. Soc. Am.* **69**(12), 1710 – 1716 (1979).
3. M. Nicolaizeau and J.L. Miquel, "LMJ status: Fifth bundle commissioning and PW class laser coupling," *Proc. SPIE* **10898**, 1089802 (2019).
4. M. L. Spaeth et al., "Description of the NIF Laser," *Fus. Sci. and Techn.* **69**, 25 – 145 (2016).
5. C. N. Danson, C. Haefner, J. Bromage, T. Butcher, J.-C. Chanteloup, E. A. Chowdhury, A. Galvanauskas, L. A. Gizzi, J. Hein, D. Hillier, N. W. Hopps, Y. Kato, E. A. Khazanov, R. Kodama, G. Korn, R. Li, Y. Li, J. Limpert, J. Ma, C. H. Nam, D. Neely, D. Papadopoulos, R. R. Penman, L. Qian, J. J. Rocca, A. A. Shaykin, C. W. Siders, C. Spindloe, S. Szatmari, R. M. G. M. Trines, J. Zhu, P. Zhu, J. D. Zuegel, "Petawatt and exawatt class lasers worldwide," *High Power Laser Science and Engineering* **7**, 1 – 54 (2019).
6. F. Frei, A. Galler and T. Feuer, "Space-time coupling in femtosecond pulse shaping and its effects on coherent control," *J. Chem. Phys.* **130**, 034302 (2009).
7. T. Alieva and J. Bastiaans, "Properties of linear canonical integral transformation," *J. Opt. Soc. Am. A* **24**(11), 3658 – 3665 (2007).
8. P. Vahimaa and J. Turunen, "Bragg diffraction of spatially partially coherent fields," *J. Opt. Soc. Am. A* **14**(1), 54 – 59 (1997).
9. J. R. Murray, J. Ray Smith, R. B. Ehrlich, D. T. Kyzazis, C. E. Thompson, T. L. Weiland, and R. B. Wilcox, "Experimental observation and suppression of transverse stimulated Brillouin scattering in large optical components," *J. Opt. Soc. Am. B* **6**(12), 2402 – 2411 (1989).
10. J. E. Rothenberg, "Comparison of beam-smoothing methods for direct-drive inertial confinement fusion," *J. Opt. Soc. Am. B* **14**(7), 1664 – 1671 (1997).
11. J. Garnier, L. Videau, C. Guedard, and A. Migus, "Statistical analysis for beam smoothing and some applications," *J. Opt. Soc. Am. A* **14**(8), 1928 – 1937 (1997).
12. R. Simon, E.C.G. Sudarshan and N. Mukunda, "Generalized rays in first order optics: transformation properties of Gaussian-Schell model fields," *Phys. Rev. A* **29**, 3273 – 3279 (1984).
13. D. Dragoman, "Wigner distribution function for Gaussian-Schell beams in complex matrix optical systems," *Appl. Opt.* **34**(18), 3352 – 3357 (1995).
14. J. Paye and A. Migus, "Space-time Wigner functions and their application to the analysis of a pulse shaper," *J. Opt. Soc. Am. B* **12**(8), 1480 – 1490 (1995).
15. M. A. Alonso, "Wigner functions in optics: describing beams as ray bundles and pulses as particles ensembles," *Advances in Optics and Photonics* **3**, 272 – 365 (2011).
16. H. Harold Szu, J. A. Blodgett, "Wigner distribution and ambiguity function," *AIP conference Proceedings* **65**, 355 – 381 (1980).
17. R. Simon, E.C.G. Sudarshan and N. Mukunda, "Gaussian pure states in quantum mechanics and the symplectic group," *Phys. Rev. A* **37**, 3028 – 3038 (1988).
18. R. Simon, E.C.G. Sudarshan and N. Mukunda, "Gaussian-Wigner distributions in quantum mechanics and optics," *Phys. Rev. A* **36**, 3868 – 3880 (1987).
19. G. A. Deschamps, "Ray techniques in electromagnetics," *Proc. IEEE* **60**, 1022 – 1035 (1972).
20. M. J. Bastiaans, "Second order-moments of the Wigner distribution function in first order optical systems," *OPTIK* **88**(4), 163 – 168 (1991).
21. A. A. Tovar, "Propagation of flat-topped multi-Gaussian laser beams," *J. Opt. Soc. Am. B* **18**(8), 1897 – 1904 (2001)
22. L. G. Wang, Q. Lin, H. Chen and S. Y. Zhu, "Propagation of partially coherent pulses beams in the spatiotemporal domain," *Phys. Rev. E* **67**, 056613-1-7 (2003).
23. M. Yao, Y. Cai, O. Korotkova, Q. Lin and Z. Wang, "Spatio-temporal coupling of random electromagnetic pulses interacting with reflecting gratings," *Opt. Express* **18**(21), 22503 – 22514 (2010).
24. M.J. Baastians, "ABCD law for partially coherent gaussian light propagating through first order optical systems," *Optical and Quantum Electronics* **24**(21), S1011 – S1 19 (1992).
25. J. Bromage, C. Dorner and R. K. Jungquist, "Temporal contrast degradation at the focus of ultrafast pulses from high-frequency spectral phase modulation," *J. Opt. Soc. Am. B* **29**(5), 1125 – 1135 (2012).
26. C. Finot, F. Chaussard, S. Broscolo, "Impact of a temporal sinusoidal phase modulation on the optical spectrum," *European Journal of Physics, European Physical Society* **39**, 055303 (2018).
27. V. Arrizon and J. Ojeda-Castañeda, "Irradiance at Fresnel planes of a phase grating," *J. Opt. Soc. Am. A* **9**(10), 1801 – 1806 (1992).
28. M. Testorf, B. Hennelly, and J. Ojeda-Castañeda, "Phase-space optics, fundamentals and applications," McGraw-Hill (2010).
29. M. Testorf and J. Ojeda-Castañeda, "Fractional Talbot effect: analysis in phase space," *J. Opt. Soc. Am. B* **13**(1), 119 – 125 (1996).
30. S. Bouillet, F. Audo, S. Fréville, L. Eupherte, C. Rouyer and J. Daurios, "Optical diffraction interpretation: an alternative to interferometers," *SPIE* **9575**, 9575A-1 (2015)
31. O. Morice, "Miró: complete modelling and software for pulse amplification and propagation in high power laser systems," *Opt. Eng.* **42**, 1530-1541(2003)