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# HAL Id: hal-03208707 https://hal.science/hal-03208707v1

Submitted on 26 Apr 2021

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Assessing the effective elastic properties of the tendon-to-bone insertion: A multiscale modeling approach

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#### Abstract

The interphase joining tendon to bone plays the crucial role of integrating soft to hard tissues, by effectively transferring stresses across two tissues displaying a mismatch in mechanical properties of nearly two orders of magnitude. The outstanding mechanical properties of this interphase are attributed to its complex hierarchical structure, especially by means of competing gradients in mineral content and collagen fibers organization at different length scales. The goal of this study is to develop a multiscale model to describe how the tendon-to-bone insertion derives its overall mechanical behavior. To this end, the effective anisotropic stiffness tensor of the interphase is predicted by modeling its elastic response at different scales, spanning from the nanostructural to the mesostructural levels, using continuum micromechanics methods. The results obtained at a lower scale serve as inputs for the modeling at a higher scale. The obtained predictions are in good agreement with stochastic finite element simulations and experimental trends reported in literature. Such model has implication for the design of bioinspired bi-materials that display the functionally graded properties of the tendon-to-bone insertion.

Keywords: Tendon-to-bone insertion, Homogenization, Continuum micromechanics, Biological interphase, Functionally graded material, Partially mineralized tissue

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#### 1. Introduction

The interaction between soft and hard tissues is essential to ensure good mobility to the musculoskeletal system. Specifically, the integration between tendon (or ligament) and bone occurs through a particular tissue interphase called *enthesis*, which derives from the ancient Greek word referring to the insertion. According to their musculoskeletal site of insertion, entheses can be characterized as fibrous or fibrocartilaginous [1]. Fibrous (or indirect) inertions are usually found where tendons and ligaments attach to the shaft of long bones, such as the insertion of the deltoid tendon into the humerus or the tibial insertion of the medial collateral ligament. Fibrocartilaginous (or direct) insertions are more common and present at the bony attachments of the rotator cuff, the anterior cruciate ligament and the 10 Achilles tendons. Fibrocartilaginous insertions are generally divided into four distinct re-11 gions with different compositions, functions, and biomechanical properties: (1) the tendon, 12 which consists of aligned type I collagen fibers with a small amount of non-collagenous 13 proteins (NCPs), including proteoglycans and the remaining volume filled with water [2]; 14 (2) a non-mineralized fibrocartilaginous region, which is an avascular zone consisting of 15 unmineralized types I, II and III collagen fibers and proteoglycan aggrecan [3]; (3) a miner-16 alized fibrocartilaginous region, which is mostly characterized by partially mineralized type 17 II collagen fibers, the proteoglycan aggrecan being produced by fibrocartilage cells or fibro-18 chondrocytes between bundles of collagen fibers [4]; and (4) the bone, which is a composite 19 material that mainly consists of highly mineralized type I collagen [5]. 20

Although these different tissue regions are compositionally distinct, they are structurally continuous. Moreover, it is acknowledged that the mechanical environment of the enthesis is complex and heterogeneous, as it is subject to tensile, compressive and shear forces [6], in order to ensure an effective transfer of the mechanical stresses across the two surroundings tissues, which show a mismatch in mechanical properties of nearly two orders of magnitude. Three main features contribute to the mechanical effectiveness of the enthesis: (1) its geometry is characterized by a shallow attachment angle and an optimized shape of the outward splay, which improves its resistance to mechanical loads [7]; (2) the interdigitation of the

bone within the tendon, described as a wave-like structure, allows a gain in toughness, as
well as a better distribution of stresses on the fraction of tissue implied in this mechanism
[8]; and (3) a compositional and structural gradient, through a gradual increase in mineral
content together with a corresponding reorganization of the collagen fibers from the tendon
to the bone [9]. The combination of these multiscale factors is hypothesized to lead to an
intermediate zone more compliant than either tendon or bone, thus ensuring an effective
distribution of mechanical stresses, which allows reducing the risk of rupture [10, 11].

Musculoskeletal injuries involving tendon rupture typically require surgical reattachment 36 back to bone. Such operations exhibit a high rate of failure because the enthesis rarely 37 regenerates and does not recover its mechanical properties. This clinical burden is mainly 38 due to the lack of understanding of the healthy attachment, as well as of the damaging 39 process that may occur at this interphase [12]. To date, studies on the tendon-to-bone 40 insertion have mainly focused on the rehabilitation methods to be implemented after a 41 tendon operation, as well as on the histology of the enthesis [13]. The targeted parameters 42 generally included compositional and structural features at lower length scales (e.q., collagen 43 types, organization and morphology of collagen fibers, degree of mineralization), but seldom the mechanical properties at the tissue scale. Indeed, the direct mechanical measurement of 45 the tendon-to-bone insertion has been shown to be very complicated, particularly because 46 of its heterogeneity and small dimensions (typically a few  $\mu$ m to some mm according to the 47 insertion site) [10, 14]. The excellent mechanical properties of the enthesis are attributed to its complex hierarchical structure. Nonetheless, understanding how this hierarchy and the 49 resulting mechanical properties at the different scales affect its overall mechanical behavior 50 and allow such an effective stress transfer across this interphase represents a considerable 51 challenge. To address this issue, tools such as multiscale modeling could help predicting 52 the effective stiffness properties of the enthesis at the tissue level, as well as studying the 53 individual impact of lower-scale features on its overall mechanical behavior at higher scales. 54 Regardless of the considered medium –tendon, enthesis or bone–, all these regions sub-55 56

higher scales [15]. Several analytical and computational models have been proposed to infer the mechanical properties of bone at different scales, in particular using homogenization 59 methods [16–19]. In the same way, a few studies have been conducted on the hierarchical 60 modeling of (mineralized) tendon [20, 21]. However, little information concerning the multi-61 scale modeling of the tendon-to-bone insertion is available in the literature. Biomechanical 62 models were initially proposed by considering the collagen-mineral interactions and the col-63 lagen fiber architecture at the microscale. In a pioneering study, Genin et al. [22] estimated 64 the stiffness of partially mineralized fibers by randomly distributing minerals among un-65 mineralized cylindrical collagen fibrils using Monte Carlo simulations, which subsequently allowed determining the isotropic mechanical properties of the partially mineralized tissue at 67 the macroscale. Later, the same group introduced a progressive stiffening model for collagen, 68 by accounting for different sequences of mineral accumulation across spatial hierarchy levels 69 spanning from the fibril to the tissue level, to estimate the mechanical properties of par-70 tially mineralized collagen tissue [23]. In a subsequent study, they presented an estimate for 71 composites containing multiple classes of aligned ellipsoidal inclusions with a relatively high 72 volume fraction, which demonstrated its potential application to the graded tissue inter-73 phase at the tendon-to-bone attachment [24]. More recently, the micromechanical behavior 74 of tendon-to-bone structures was studied and compared with a composite beam-bending 75 model [25]. The beam was modeled as a three-phase composite in series, in which the 76 modulus of the transitional graded interphase was defined as a linear extrapolation between 77 those of the mineralized and unmineralized regions. All these models provided a general 78 understanding on the mechanical properties of the insertion site (i.e., average Young mod-79 ulus) at different structural levels. However, to the best of our knowledge, these models 80 only partially accounted for the anisotropic nature of the interphase at the different scales, 81 and the transition between different hierarchies from the nanoscale to the mesoscale was 82 not achieved. Moreover, the impact of the individual nanoscale properties on the effective 83 stiffness tensor at the mesoscale was not addressed so far. 84

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ing its elastic response at different scales spanning from the nano- to the mesostructural levels. To this end, a continuum micromechanics approach was proposed by distinguishing 88 six homogenization steps at four different scales. First, we derived the effective stiffness 89 properties of the enthesis at the mesoscale by accounting for gradients in mineral content 90 and collagen fibers organization across the interphase at lower scales. Second, we carried out 91 a parametric study to evaluate the impact of the mechanical properties of each nanoscale 92 component on the overall mechanical behavior of the insertion site. Third, based on the 93 knowledge of the mechanical properties of the two surroundings tissues at the mesoscale, an optimization procedure was solved to identify the nanoscale properties that yield the best 95 agreement between the model predictions and ultrasonic data reported in the literature. 96 Numerical results show that the effective mechanical properties predicted by our multiscale 97 model are in good agreement with earlier modeling approaches and experimentally observed trends. Moreover, our model is straighforward and could be readily extended to account for 99 damage characteristics of the interphase as a footprint of clinical burden. 100

The paper is structured as follows: Sec. 2 recalls the fundamentals of micromechanics-based models and introduces the proposed multiscale modeling of the enthesis. The model parameters, *i.e.*, the mechanical properties, volume fractions and gradients of the tendon-to-bone insertion's components, are then described in Sec. 3. Finally, the numerical results are presented in Sec. 4 and discussed in the light of the literature in Sec. 5.

#### 106 2. Methods

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In this section, a multiscale model of the tendon-to-bone insertion based on a hierarchical modeling approach that consists of successive homogenization steps from the nanostructural to the mesostructural levels is proposed. This multiscale model relies on several hypotheses:

(1) whereas mineral grading can occur at any length scale, it was incorporated here at the nanoscale [26]; (2) the gradient in collagen fibers organization occured at the microscale [22];

(3) type I collagen was assumed to be the main organic component across the interphase;

and (4) the mechanical properties of the elementary nanoscale components (collagen, hydroxyapatite (HA), water and NCPs), which were used as input data in our model, were all

s assumed to have a linear elastic and isotropic behavior.

#### 2.1. Fundamentals of micromechanics-based models

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Continuum micromechanics allows estimating the effective stiffness properties of heteroge-117 neous materials [27], which can generally be derived by solving the matrix-inclusion problem 118 according to Eshelby [28]. A basic concept of continuum micromechanics consists in selecting 119 a representative volume element (RVE) in the structure. Based on the separation-of-scales requirement, the characteristic length  $\ell$  of this RVE must be considerably larger than the 121 dimension of the heterogeneities d within the RVE and implicitly smaller than the charac-122 teristic length of the structure  $\mathcal{L}$ , so that  $d \ll \ell \ll \mathcal{L}$ . Useful information concerns volume 123 fractions and elastic properties of phases, the existence of one connected matrix phase in which one or several inclusion phases with different shapes are embedded or the disordered 125 arrangement of all phases. Hence, the homogenized stiffness tensor of such RVE,  $\mathbb{C}_{hom}$ , can 126 be stated as in [29]

$$\mathbb{C}_{\text{hom}} = \sum_{r=1}^{n} f_r \mathbb{C}_r : \left[ \mathbb{I} + \mathbb{P}_r^0 : (\mathbb{C}_r - \mathbb{C}^0) \right]^{-1} : \left\{ \sum_{s=1}^{n} f_s \left[ \mathbb{I} + \mathbb{P}_s^0 : (\mathbb{C}_s - \mathbb{C}^0) \right]^{-1} \right\}^{-1}$$
(1)

where: denotes the double contracted product of two tensors, n is the number of phases in the RVEs,  $f_r$  denotes the volume fraction of the phase r,  $\mathbb{I}$  is the fourth-order identity tensor,  $\mathbb{C}_r$  is the stiffness tensor of phase r,  $\mathbb{C}^0$  is the stiffness tensor of the matrix phase, and  $\mathbb{P}^0$  is the so-called Hill tensor of phase r embedded in the matrix phase, whose form depends on the shape of the inclusion and on the stiffness tensor of the surrounding matrix  $\mathbb{C}^0$ . The Hill tensor can be related to the Eshelby tensor as  $\mathbb{P}^0 = \mathbb{S}^0_{\mathrm{Esh}} : (\mathbb{C}^0)^{-1}$ . The components of the Hill tensors for the different inclusion shapes considered in this study are provided in Appendix A.

Several approaches for determining  $\mathbb{C}^0$  have been proposed in the literature. Among these, two methods are particularly relevant to our context. First, the Mori-Tanaka scheme can be applied when the inclusion phase consists of small particles surrounded by a continuous homogeneous matrix [30]. In such a case, the matrix phase can be directly identified as the actual matrix phase (*i.e.*,  $\mathbb{C}^0 = \mathbb{C}_{\text{matrix}}$ ). Second, when no matrix phase can be

clearly identified, it is convenient to apply the self-consistent scheme [31]. The underlying hypothesis of this second method consists in estimating the homogenized properties of the effective medium by locating each inhomogeneity in the homogeneous effective medium (*i.e.*,  $\mathbb{C}^0 = \mathbb{C}_{hom}$ ).

#### 2.2. Multiscale modeling of the tendon-to-bone insertion

According to earlier studies applied to partially mineralized tissues, such as bone [17, 32, 33] and mineralized tendon [20], we hierarchically modeled here the mechanical behavior of the 147 tendon-to-bone insertion by considering four relevant hierarchical levels compound of six 148 homogenization steps (see Fig. 1), and by assuming gradients in mineral content and colla-149 gen fibers organization across the interphase between the two surroundings tissues [22, 23]. 150 Across the enthesis, mineral crystals replace the water by filling the gap zones between 151 collagen thereafter spreading through the fibrils surface [34]. Extrapolating this considera-152 tion to our modeling approach, the fibril array was considered as a hydrated collagen fibril 153 at the nanoscale (i) without mineral particles in the tendon region, (ii) with a functional 154 grading of mineral across the insertion, and (iii) as a partially mineralized collagen fibril 155 in the bone region. The gradient of mineralization was introduced in the model by means 156 of a normalized mineral volume fraction,  $\phi(x)$ , defined as the ratio of the overall mineral 157 volume fraction  $f_{\text{mineral}}(x)$  at a normalized position along the insertion, x, to the overall volume fraction of mineral at the bone level,  $f_{\text{mineral}}^{\text{bone}}$ , so that  $\phi(x) = f_{\text{mineral}}(x)/f_{\text{mineral}}^{\text{bone}}$ . The 159 normalized position x across the interphase was thus x = 0 at the end of the tendon (i.e., 160  $\phi(x) = 0$ ) and x = 1 at the beginning of the bone (i.e.,  $\phi(x) = 1$ ), so that  $\phi(x) \approx x$ . 161 Moreover, the gradient in collagen fibers organization was accounted for by considering the 162 angular deviation of the collagen fibers distribution at a normalized position x along the 163 insertion [22]. At each scale, the homogenized stiffness tensor of the enthesis,  $\mathbb{C}_{hom}(x)$ , was 164 therefore related to a precise description of the geometric and mechanical properties of each 165 of the subscales at any normalized position x across the interphase.

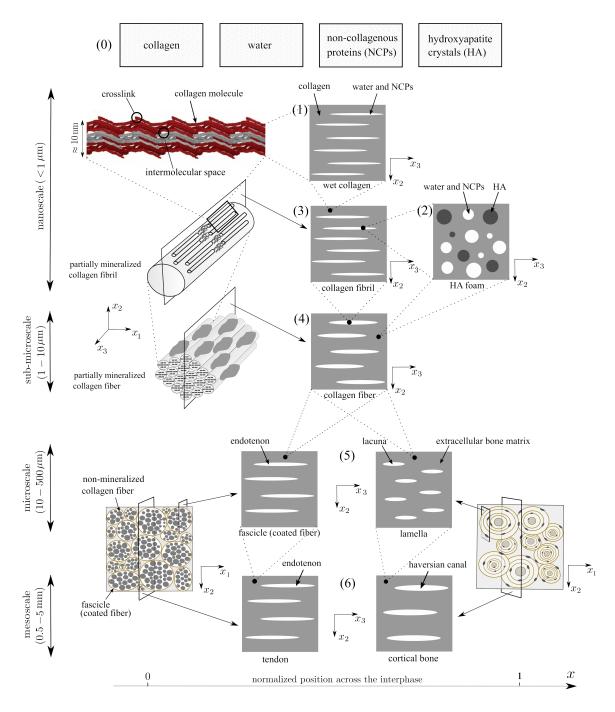


Figure 1: Successive steps of our multiscale modeling approach: (0) elementary nanoscale components; (1) wet collagen composite formed by a collagen molecules matrix containing water and NCPs-filled pores (image adapted from [35]); (2) HA foam formed by HA crystals containing water and NCPs-filled pores; (3) partially mineralized collagen fibril made up from wet collagen and intrafibrillar HA foam; (4) partially mineralized collagen fibre formed by partially mineralized collagen fibrils and extrafibrillar HA foam; (5) lamella and unit fascicle built up from a partially mineralized collagen fiber matrix holding lacunae cavities and endotenon, respectively; and (6) cortical bone and tendon formed by a continuous bone matrix (extravascular bone) and a bundle of parallel fascicles matrix containing haversian canals and endotenon, respectively.

#### $A.\ Nanostructural\ level$

At the nanostructural level, collagen molecules are attached to each other by crosslinks, 168 which transfer the stresses from one collagen molecule to another [35]. Water initially fills the 169 spaces between collagen molecules. During the mineralization process, mineral crystals are 170 nucleated in the gap zones, thus progressively replacing the water. They then further grow 171 into the overlap zones during development and may be distributed throughout the collagen 172 molecules to form a partially mineralized collagen fibril [36]. To mimic this physiological process, three homogenization steps were used to model the enthesis at the nanostructural 174 level: (1) combining collagen together with water and NCPs, (2) combining HA crystals 175 with water and NCPs, and (3) combining the collagen composite from step 1 with the HA foam from step 2 to form a partially mineralized collagen fibril. 177

178 Step 1: wet collagen composite

The wet collagen composite was modeled as a contiguous cross-linked collagen molecules matrix with holes hosting water and NCPs, where the holes were considered as cylindrical inclusions [33, 37]. Based on the assumption that a continuum approach can be applied at the nanostructural level, the Mori-Tanaka scheme was used to estimate the effective stiffness tensor of the composite,  $\mathbb{C}_{wc}$ , at this scale as

$$\mathbb{C}_{wc} = (1 - f_{wp})\mathbb{C}_{col} + f_{wp}\mathbb{C}_{wp} : \left[\mathbb{I} + \mathbb{P}_{wp}^{col}(\mathbb{C}_{wp} - \mathbb{C}_{col})\right]^{-1} \\
: \left\{ (1 - f_{wp})\mathbb{I} + f_{wp}\left[\mathbb{I} + \mathbb{P}_{wp}^{col} : (\mathbb{C}_{wp} - \mathbb{C}_{col})\right]^{-1} \right\}^{-1}, \tag{2}$$

where subscripts "wc", "col" and "wp" are referred to the wet collagen composite, collagen, and water-protein mixture, respectively.

186 Step 2: hydroxyapatite foam

At the scale of a few hundred nanometers, HA crystals were assumed to be interpenetrated by intercrystalline space filled with water and NCPs [19]. Since no well-defined matrix phase can be identified at this scale, it results convenient to apply the self-consistent scheme to estimate the effective stiffness tensor of the HA foam,  $\mathbb{C}_{Hw}(x)$ , which depends upon the mineral gradient across the insertion. Both the mineral phase and the water-protein mixture were here assumed to be spherical in shape with a linear elastic and isotropic behavior.

Hence,

$$\mathbb{C}_{\mathrm{Hw}}(x) = \sum_{r=1}^{2} f_r(x) \mathbb{c}_r : \left[ \mathbb{I} + \mathbb{P}_r^0 : (\mathbb{c}_r - \mathbb{C}_{\mathrm{Hw}}(x)) \right]^{-1} : \left\{ \sum_{s=1}^{2} f_s(x) \left[ \mathbb{I} + \mathbb{P}_s^0 : (\mathbb{c}_s - \mathbb{C}_{\mathrm{Hw}}(x)) \right]^{-1} \right\}^{-1},$$
(3)

with  $r, s \in [HA, wp]$ . The subscript "Hw" denotes the HA foam. The detailed calculation of the stiffness coefficients is provided in Appendix B.

196 Step 3: partially mineralized collagen fibril

The collagen fibril structure was modeled as a composite consisting of a wet collagen composite matrix reinforced with aligned needle-shaped interfibrillar minerals, which were assumed to be distributed along the main axis of the collagen fibril [38]. Again, the classical
Mori-Tanaka method was applied to determine the effective stiffness tensor of a partially
mineralized collagen fibril,  $\mathbb{C}_{fib}(x)$ , as

$$\mathbb{C}_{\text{fib}}(x) = (1 - f_{\text{Hw}}) \mathbb{c}_{\text{wc}} + f_{\text{Hw}} \mathbb{c}_{\text{Hw}}(x) : \left[ \mathbb{I} + \mathbb{P}_{\text{Hw}}^{\text{wc}} : (\mathbb{c}_{\text{Hw}}(x) - \mathbb{c}_{\text{wc}}) \right]^{-1} \\
: \left\{ (1 - f_{\text{Hw}}) \mathbb{I} + f_{\text{Hw}} \left[ \mathbb{I} + \mathbb{P}_{\text{Hw}}^{\text{wc}} : (\mathbb{c}_{\text{Hw}}(x) - \mathbb{c}_{\text{wc}}) \right]^{-1} \right\}^{-1}, \tag{4}$$

where  $\mathbb{C}_{wc}$  and  $\mathbb{C}_{Hw}(x)$  are the effective stiffness tensors resulting from the application of steps 1 and 2, respectively.

204 B. Sub-microstructural level – partially mineralized collagen fiber

At the scale of several micrometers (step 4), randomly dispersed extrafibrillar minerals with different sizes and shapes strongly adhere to the outer parts of the collagen fibrils [34]. During the mineralization, these extrafibrillar minerals are located around the fibril surface and then start growing along the main fibril axis, therefore forming a sort of reinforcing structure around the fibril. Consequently, a partially mineralized collagen fiber can be modeled as an extrafibrillar HA foam matrix pervaded by a cylindrical collagen fibril inclusion. In this way, the effective stiffness tensor of the partially mineralized collagen fiber,  $\mathbb{C}_{fbr}(x)$ , can be

estimated using the Mori-Tanaka scheme as

$$\mathbb{C}_{\text{fbr}}(x) = (1 - f_{\text{fib}}) \mathbb{c}_{\text{Hw}}(x) + f_{\text{fib}} \mathbb{c}_{\text{fib}}(x) : \left[ \mathbb{I} + \mathbb{P}_{\text{fib}}^{\text{Hw}}(x) : \left( \mathbb{c}_{\text{fib}}(x) - \mathbb{c}_{\text{Hw}}(x) \right) \right]^{-1} \\
: \left\{ (1 - f_{\text{fib}}) \mathbb{I} + f_{\text{fib}} \left[ \mathbb{I} + \mathbb{P}_{\text{fib}}^{\text{Hw}}(x) : \left( \mathbb{c}_{\text{fib}}(x) - \mathbb{c}_{\text{Hw}}(x) \right) \right]^{-1} \right\}^{-1}, \tag{5}$$

where  $c_{Hw}(x)$  and  $c_{fib}(x)$  are the effective stiffness tensors resulting from the application of steps 2 and 3, respectively.

Moving from the tendon to bone, apart from the gradual increase in mineral content (already accounted for in step 2), the insertion site also shows a gradual decrease in the organization of collagen fibers. At the tendon level, collagen fibers are aligned and parallelly-oriented along the  $x_3$ -axis. These then start bending and intercrossing along the insertion, and become less organized near the bone [9]. Consequently, the effective stiffness tensor of an individual partially mineralized collagen fiber,  $\mathbb{C}_{fbr}(x)$ , can be related to the microstructural level (*i.e.*, fiber bundle or pattern) by averaging over the fibers orientation at each position x of the insertion. Hence,

$$\bar{\mathbb{C}}_{fbr}(x) = \int_{S} \left( \mathbb{R}(\boldsymbol{n}) \mathbb{C}_{fbr}(x) \mathbb{R}^{T}(\boldsymbol{n}) \right) p(\boldsymbol{n}, s(x)) dS$$
 (6)

where the superscript T denotes the transpose operator;  $p(\boldsymbol{n}, s(x))$  is an axisymmetric and spherical probability density function for finding a fiber whose axis is parallel to the unit vector  $\boldsymbol{n}$ , with s(x) the angular deviation at a position x with respect to the mean fiber direction  $x_3$ ;  $\mathbb{R}(\boldsymbol{n})$  rotates a tensor from this coordinate system into the global coordinate system; and  $\boldsymbol{n}$  includes all directions over the unit sphere S [22]. Note that the resulting stiffness tensor of a partially mineralized fibers pattern,  $\bar{\mathbb{C}}_{fbr}(x)$ , retains the symmetry class of that of an individual collagen fiber, but results more compliant along the  $x_3$ -direction.

230 C. Microstructural level – partially mineralized fibers pattern

At the microstructural level (step 5), the structure and composition of tendon and bone are no longer identical, and therefore lead to a different approach for modeling a partially mineralized fibers pattern. At the tendon level, a bundle of aligned collagen fibers forms a unit fascicle, which represents the basic unit of a tendon. A thin connective tissue called

endotenon surrounds each collagen fiber and binds the fibers together [39]. Since the self-235 consistent and Mori-Tanaka methods are limited to an inclusion associated with a low volume 236 fraction, the interactions between neighboring fibers cannot be neglected anymore [24]. To 237 circumvent this issue, the so-called "inverse Mori-Tanaka" scheme was applied here simply 238 by inverting the matrix and inclusion phases [17]. In this regard, by considering the con-239 nected bundle of collagen fibers as a matrix and endotenon as a cylindrical inclusion, the 240 effective stiffness tensor of a unit fascicle can be estimated. At the bone level, a lamella 241 containing cavities filled with osteocytes called *lacunae* is enclosed by the continuous bone 242 matrix holding a bundle of preferentially oriented collagen fibers. In addition, the main axis of lacunae was assumed to be oriented along the longitudinal direction of a lamella [40]. 244 Therefore, similarly to the tendon region, the effective stiffness tensor of a lamella can be 245 estimated by considering the lacuna as a cylindrical inclusion coated with a continuous bone matrix. Hence, the effective stiffness tensor of a partially mineralized fibers pattern at the microstructural level,  $\mathbb{C}_{\text{mic}}(x)$ , was estimated by considering  $\mathbb{C}^0 \equiv \mathbb{C}_{\text{end}} \equiv \mathbb{C}_{\text{lac}} \equiv \mathbb{C}_{\text{wp}}$ , so 248 that 249

$$\mathbb{C}_{\text{mic}}(x) = f_{\text{fbr}}(x)\bar{\mathbb{c}}_{\text{fbr}}(x) + (1 - f_{\text{fbr}}(x))\mathbb{c}_{\text{wp}} : \left[\mathbb{I} + \bar{\mathbb{P}}_{\text{wp}}^{\text{fbr}}(x) : (\mathbb{c}_{\text{wp}} - \bar{\mathbb{c}}_{\text{fbr}}(x))\right]^{-1} \\
: \left\{ f_{\text{fbr}}(x)\mathbb{I} + (1 - f_{\text{fbr}}(x))\left[\mathbb{I} + \bar{\mathbb{P}}_{\text{wp}}^{\text{fbr}}(x) : (\mathbb{c}_{\text{wp}} - \bar{\mathbb{c}}_{\text{fbr}}(x))\right]^{-1} \right\}^{-1}.$$
(7)

In this way, the main difference between tendon and bone at this scale was accounted for by considering a linearly varying volume fraction across the interphase (i.e.,  $1 - f_{\text{fbr}}(x)$  with  $x \in \{0,1\}$ ) for their corresponding inclusion (endotenon for x = 0 or lacuna for x = 1).

253 D. Mesostructural level – partially mineralized tissue

Similarly to the microstructural level, a different modeling approach was used for the tendon and bone to account for their differences in structure and composition at the mesostructural level (step 6). On the one hand, a bundle of parallel fascicles with an interfascicular matrix (endotenon) forms the tendon [41]. Considering a bundle of parallel fascicles as a matrix and endotenon as a cylindrical inclusion, the effective stiffness tensor of tendon can be estimated. On the other hand, cortical bone hosts cylindrical pores called haversian canals, which contain the bone's nerve and blood supplies being embedded into the osteonal lamella matrix (extravascular bone) [37]. The effective stiffness tensor of cortical bone can thus be estimated by considering haversian canals as a cylindrical inclusion coated with a continuous bone matrix.

As in step 5, the effective stiffness tensors of a partially mineralized tissue,  $\mathbb{C}_{\text{meso}}(x)$ , can be estimated by considering  $\mathbb{C}^0 \equiv \mathbb{C}_{\text{end}} \equiv \mathbb{C}_{\text{hav}} \equiv \mathbb{C}_{\text{wp}}$ ,

$$\mathbb{C}_{\text{meso}}(x) = f_{\text{mic}}(x)\mathbb{c}_{\text{mic}}(x) + (1 - f_{\text{mic}}(x))\mathbb{c}_{\text{wp}} : \left[\mathbb{I} + \mathbb{P}_{\text{wp}}^{\text{mic}}(x) : (\mathbb{c}_{\text{wp}} - \mathbb{c}_{\text{mic}}(x))\right]^{-1} \\
: \left\{ f_{\text{mic}}(x)\mathbb{I} + (1 - f_{\text{mic}}(x)) \left[\mathbb{I} + \mathbb{P}_{\text{wp}}^{\text{mic}}(x) : (\mathbb{c}_{\text{wp}} - \mathbb{c}_{\text{mic}}(x))\right]^{-1} \right\}^{-1}.$$
(8)

Again, the variation across the interphase was here driven by the a linearly varying volume fraction,  $1 - f_{\rm mic}(x)$ , for their corresponding inclusion only.

#### 268 3. Model parameters

Like for any other hierarchical material, the knowledge of the mechanical properties and volume fractions of the tendon-to-bone insertion's elementary components, along with the competing gradients in collagen fibers organization and mineral content, are required to assess its overall mechanical behavior. A broad range of values for the mechanical properties (e.g., Young's modulus E and Poisson's ratio  $\nu$ ) of the nanoscale components has been reported in the literature, based on experimental testing facilities or molecular dynamics (MD) simulations.

Earlier theoretical studies provided values of around 2 GPa for the Young's modulus of 276 collagen, but recent MD simulations reported higher values for the Young's modulus of a 277 single collagen molecule, ranging from 7 to 19 GPa [33]. The Poisson's ratio of collagen 278 has typically been chosen equal to 0.28 to deliver a Poisson's ratio of around 0.35 for the 279 collagen-water composite (recall step 1) [18, 17]. Furthermore, HA minerals were typically 280 assumed to have a linear elastic and isotropic behavior, with a Young's modulus ranging 281 from 63 to 165 GPa, and a Poisson's ratio of around 0.23 [33]. To date, little is known about 282 the mechanical properties of NCPs. Nevertheless, water-NCPs composite has generally 283 been assumed to have isotropic properties with a Young's modulus ranging from 1.3 MPa to 0.7 GPa, and a Poisson's ratio ranging from 0.4999 to 0.45, thus encompassing typical 285

values for water (*i.e.*, a nearly incompressible material with a bulk modulus of 2.2 GPa) and soft polymers [17]. The mechanical properties range of each nanoscale component used as input values in our multiscale model is summarized in Tab. 1.

Material	Young's modulus, $E$ (GPa)	Poisson's ratio, $\nu$ (–)
Collagen molecules	1.2 - 9.6	0.23 - 0.33
HA minerals	63 - 165	0.18 - 0.28
Water-NCPs composite	0.0013 - 0.7	0.4999 - 0.45

Table 1: Mechanical properties range of the nanoscale components used for modeling the mesoscale effective stiffness tensor of the tendon-to-bone insertion in this study.

A volume fraction of mineral  $f_{\rm HA}(x)$  varying linearly from 0 to 65% was considered (see Fig. 2a), based on the experimentally observed mineral concentration gradient at the tendon-to-bone insertion [26]. Moreover, the collagen fibers distribution s(x) was chosen according to experimental data [42], which reported that the angular deviation values for the tendon-to-bone insertion reached a peak in the non-mineralized fibrocartilaginous region and varied little beyond it (see Fig. 2b)

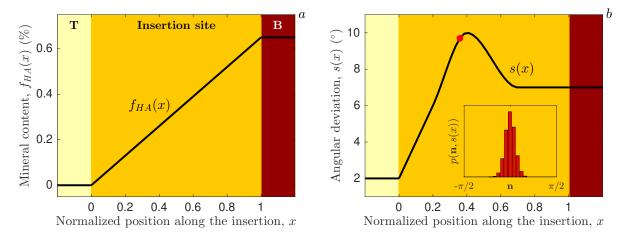


Figure 2: Spatially varying properties of the interphase from the tendon (T) to bone (B) based on competing gradients in (a) mineral content and (b) collagen fibers organization. The insert displays the distribution of collagen fibers near the region of largest fibers disorganization (red dot). These trends are based upon data reported in Refs. [22, 24].

The volume fractions of each phase, along with the different homogenization methods 295 applied at each step, are summarized in Tab. 2. At the nanostructural level, the volume 296 fraction for the water-protein mixture,  $f_{\rm wp}$ , has been taken according to [17]. At the nanos-297 tructural and sub-microstructural levels, we used the values for a mineralization inside and 298 outside the fibrilsas a function of the total mineral content, i.e.,  $f_{\text{Hw}}$  and  $1 - f_{\text{fib}}$ , following 299 the work of Nikolov and Raabe [17]. In this way, these concurrent variations lead to an 300 overall volume fraction of mineral varying from 0 to 52% between tendon and bone at the 301 mesoscale, which is in good agreement with values reported in the literature [16, 17]. Fur-302 thermore, to account for the difference in structure and composition of the tendon and bone at the micro- and mesostructural levels, volume fractions that vary linearly with respect to 304 the normalized position x across the interphase, i.e.,  $1 - f_{\text{fbr}}(x)$  and  $1 - f_{\text{mic}}(x)$ , have been 305 considered for their respective inclusions at each of these two levels [16, 20].

Hierarchical levels	Steps	Homogenization schemes		Phases	Volume fractions (%)	Relations	
Nano	1	Mori-Tanaka	Inclusion Matrix	Water-NCPs Collagen	$f_{\rm wp} = 35$	Eq. (2)	
	2	Self-consistent	Phase 1 Phase 2	Water-NCPs HA	$f_{\rm HA}(x) = 65x$	Eq. (3)	
	3	Mori-Tanaka	Inclusion Matrix	Intra-HA foam Wet collagen	$f_{\rm Hw} = 43$	Eq. (4)	
Sub-micro	4	Mori-Tanaka	Inclusion Matrix	Collagen fibril Extra-HA foam	$f_{\mathrm{fib}} = 73$	Eq. (5)	
Micro	5	Inverse Mori-Tanaka	Inclusion Matrix	Endotenon – Lacunae Collagen fiber	$f_{\rm fbr}(x) = 8x + 90$	Eqs. (6–7)	
Meso	6	Inverse Mori-Tanaka	Inclusion Matrix	Endotenon – Haversian Fascicle – Lamella	$f_{\rm mic}(x) = 8x + 90$	Eq. (8)	

Table 2: Volume fractions of each phase along with the different homogenization schemes applied at each step of the multiscale modeling.

#### 4. Numerical results

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The effective stiffness properties of the tendon-to-bone insertion from the nanostructural to
the mesostructural levels were estimated using the six homogenization steps described in
Fig. 1. Three analyses were conducted to rate the performance of the proposed multiscale
model. First, several characteristics of the effective stiffness coefficients obtained at the
mesostructural level (step 6) were assessed. Second, a parametric study was performed to
evaluate the impact of the input parameters on the resulting effective stiffness tensor. Third,
an optimization procedure was proposed to identify the model parameters that yielded the
best agreement between modeled and measured effective stiffness tensors.

## 316 4.1. Effective stiffness tensor at the mesostructural level, $\mathbb{C}_{meso}(x)$

Figure 3 depicts the effective stiffness coefficients of the tendon-to-bone insertion obtained at the mesostructural level, together with the Voigt-Reuss bounds calculated according to [43], using the average mechanical properties from Tab. 1. Note that Voigt notation is adopted to define the coefficients of the fourth-order stiffness tensor and that the  $x_3$ -axis denotes the longitudinal axis, which is aligned with the mean direction of the collagen fibers distribution (recall step 4 of Fig. 1).

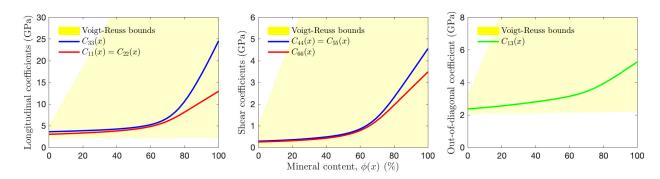


Figure 3: Effective stiffness coefficients (in Voigt notation) of the tendon-to-bone insertion obtained at the mesostructural level (continuous lines), along with the Voigt-Reuss bounds (yellow area), using the average mechanical properties from Tab. 1.

First, it is worth pointing out that our model led to a transversely isotropic elastic behavior for the insertion site at the mesoscale, where the stiffness coefficients along the fibers

(i.e.,  $C_{33}(x)$  and  $C_{44}(x)$ ) were consistently larger than those normal to the fibers (i.e.,  $C_{11}(x)$ and  $C_{66}(x)$ ), thus retaining the symmetry class of the two surroundings tissues. Recalling 326 that the main constituent properties at the nanoscale were all considered as linear elastic and 327 isotropic, the anisotropy at the different scales was initially triggered by the consideration of 328 geometric features (i.e., the shape of the inclusion) in step 1. Second, the derived effective 329 stiffness tensors at the extremities of the insertion, i.e.,  $\mathbb{C}_{\text{meso}}(x=0)$  and  $\mathbb{C}_{\text{meso}}(x=1)$  from 330 Eq. (8), were in good agreement with experimental values reported in the literature for 331 the Achilles tendon [44, 45] and cortical bone, measured at the tibia mid-diaphysis [46] and 332 femoral diaphysis [47]. Third, the derived stiffness profile across the interphase was nonlinear, displaying a smooth stiffness increase in the non-mineralized fibrocartilaginous region 334  $(\phi(x) < 50\%)$  followed by a sudden rise in stiffness in the mineralized fibrocartilaginous 335 region ( $\phi(x) > 50\%$ ). The fact that an increase in mineral accumulation within collagen fibers can provide significant stiffening of the enthesis, but only for concentration of minerals above a certain percolation threshold (i.e.,  $\phi(x) \approx 60\%$ ), corroborates earlier results and 338 confirms that our model can account for effects of nonuniform mineral accumulation [22]. 339 Finally, it should be noted that the obtained stiffness coefficients all fall within the lower and upper limits defined by the Voigt-Reuss bounds. 341

#### 342 4.2. Parametric study

The predicted output values of our multiscale model strongly depend on the mechanical prop-343 erties of the nanoscale components, which showed a rather high dispersion among different 344 studies (recall Tab. 1). We thus performed a parametric study to assess the impact of these 345 input parameters on the values of the resulting effective stiffness tensor at the mesostruc-346 tural level. The obtained bounds for each of the five independent stiffness coefficients are 347 depicted in Fig. 4. As can be observed, the mechanical properties of the water-NCPs com-348 posite inclusion were mostly responsible for relative variations of the shear coefficients (i.e., 349  $C_{44}(x)$  and  $C_{66}(x)$ ) in the low-mineralized region, whereas they barely affected the longitu-350 dinal coefficients. The mechanical properties of the collagen molecules matrix significantly 351 affected all stiffness coefficients over the entire tendon-to-bone transition. As expected, the mechanical properties of HA minerals were only responsible for variations of the stiffness coefficients in the highly mineralized region.

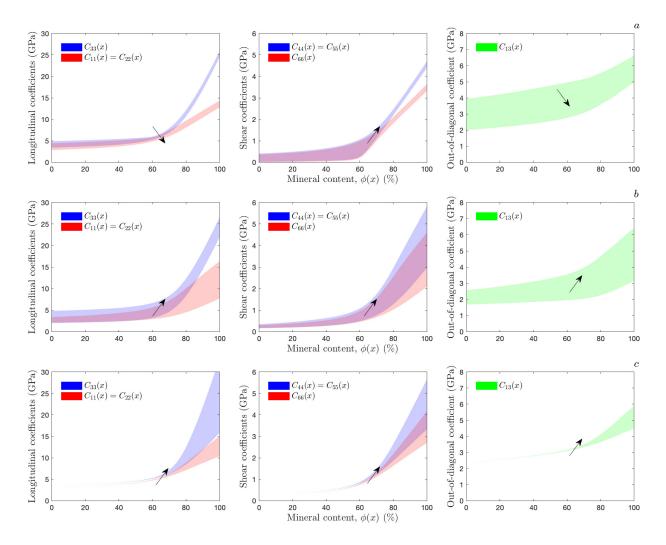


Figure 4: Stiffness coefficients bounds of the tendon-to-bone insertion at the mesostructural level for a variation of the mechanical properties of the nanoscale components: (a) water-NCPs composite; (b) collagen molecules; and (c) HA minerals. Black arrows indicate the resulting trends in stiffness with respect to increasing nanoscale properties.

#### 4.3. Identification procedure

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To further investigate the sensitivity of the effective stiffness tensor across the interphase, which, in turn, reflects the ability of our model to predict experimentally observed data, we proposed an optimization procedure to identify the model parameters  $\theta$  (*i.e.*, the six nanoscale properties from Tab. 1) that yielded the best agreement between modeled and measured effective stiffness tensors at the mesostructural scale, denoted by  $\mathbb{C}_{\text{meso}}(\boldsymbol{\theta};x)$  and  $\mathbb{C}^{\text{exp}}(x)$ , respectively. Since experimental data for the interphase at the tissue scale are not available in the literature, the optimization was conducted on data from the two surrounding tissues only, *i.e.*,  $\mathbb{C}^{\text{exp}}(x=0)$  and  $\mathbb{C}^{\text{exp}}(x=1)$ , which are summarized in Tab. 3.

Material	Position	Stiffness tensor $\mathbb{C}^{\exp}(x)$				D. C	
		$C_{11}^{\mathrm{exp}}$	$C_{33}^{\mathrm{exp}}$	$C_{13}^{\mathrm{exp}}$	$C_{55}^{\mathrm{exp}}$	$C_{66}^{\mathrm{exp}}$	Reference
Tendon	x = 0	3.08	4.51	3.10	0.04	0.02	[44]
Cortical bone	x = 1	14.79	26.64	6.31	5.52	3.65	[46]

Table 3: Reference stiffness coefficients (in GPa) for tendon and bone at the mesostructural scale.

Formally, the optimal model parameters  $\hat{\boldsymbol{\theta}}$  result from the minimization of an objective function  $F(\boldsymbol{\theta})$  in a least-squares sense as,

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$$\hat{\boldsymbol{\theta}} = \arg\min_{\boldsymbol{\theta} \in \Theta} F(\boldsymbol{\theta}) \quad \text{with} \quad F(\boldsymbol{\theta}) = \frac{1}{N} \sum_{n=1}^{N} \left( \frac{C_n^{\text{exp}} - C_n(\boldsymbol{\theta})}{C_n^{\text{exp}}} \right)^2, \tag{9}$$

where  $C^{\text{exp}}$  and  $C(\theta)$  are vectors that contain the measured and modeled stiffness coefficients, respectively, N is the number of stiffness coefficients, and  $\Theta$  denotes the bounds of the 367 model parameters  $\theta$ , which were taken according to Tab. 1. Genetic algorithms were applied to solve Eq. (9) owing to their ability in finding a near global solution for non-convex multi-369 dimensional objective functions, without the need of an accurate initial guess for the model 370 parameters [48]. For this set of experimental data, the identification procedure delivered the 371 following estimates for the model parameters,  $\hat{\boldsymbol{\theta}} = [0.0153 \ 0.4991 \ 7.9 \ 0.29 \ 125.5 \ 0.23]$ , which 372 pairwise correspond to the Young's moduli (GPa) and Poisson's ratios of the water-NCPs 373 composite, collagen molecules and HA minerals, respectively. Figure 5 depicts the optimal 374 matching between the modeled and measured effective stiffness tensors at the mesostructural scale. An excellent agreement can be observed between the modeled and measured 376 longitudinal stiffness coefficients, with relative errors that are less than 5\%. However, the 377 agreement for the shear and out-of-diagonal stiffness coefficients was moderate for the tendon region, thus indicating that our model lacks some mechanical features that may account for a stronger anisotropy in the transverse direction.

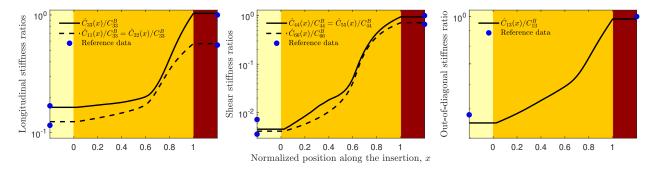


Figure 5: Optimal matching between the modeled and measured effective stiffness tensors at the mesostructural scale. Results were normalized with respect to the stiffness coefficients of bone (Tab. 3).

#### 5. Discussion

The tendon-to-bone insertion plays a crucial role in the musculoskeletal system, achieving an effective transfer of mechanical stresses across two tissues displaying a substantial gap in mechanical properties of nearly two orders of magnitude. An accurate modeling accounting for the mutiscale and composite nature of this interphase is essential to deepen our understanding of complex biological interphases and has potential applications both for clinical purposes and for the development of biomimetic strategies in engineering. In this study, we proposed a multiscale model spanning four different hierarchical levels, from the nanoto the mesoscale, connected based on homogeneization procedures. Starting from the mechanical properties of the elementary nanoscale components, this multiscale modeling strategy allowed deriving the effective stiffness tensor of the enthesis at the mesoscale. Since an analytical formulation was used at every step, the computational cost of modeling was very low.

The main findings from this study were as follows: (1) the effective stiffness tensor across the interphase was found to be transversely isotropic at the mesostructural level (Fig. 3), thus retaining the symmetry class of the two surrounding tissues (*i.e.*, tendon and bone); (2) the mechanical properties of the elementary nanoscale components were shown to impact

the stiffness coefficients of the interphase at the mesoscale (Fig. 4), and were subsequently optimized, by solving an inverse identification procedure, to match available ultrasonic data for the two surroundings tissues (Fig. 5); and (3) our modeling results for the interphase were in qualitative agreement with stochastic finite-element estimates [22, 23] and reported experimental values for the insertional zones of human meniscal attachments into underlying bone [49], thus supporting the hypothesis that the tendon-to-bone insertion can be seen as a continuous functionally graded material.

Indeed, in accordance with these earlier models, our results indicated that the competing 405 gradients in mineral concentration (at the nanoscale) and collagen fibers organization (at the microscale) are important factors in determining the effective mechanical behavior of 407 the tendon-to-bone insertion at the mesoscale. On the one hand, the linear increase in 408 mineral content caused a stiffening of the interphase that became significant beyond a certain 409 percolation threshold. On the other hand, the decreasing collagen fiber organization across 410 the interphase led to a reduced tissue stiffness along the main fibers direction. The existence 411 of a region that is more compliant than either tendon or bone (continuous line in Fig. 6) 412 was the result of these competing gradients. Moreover, it can be shown that the dominant factors driving the width and depth of this compliant region are the onset of mineralization 414 and the angular deviation from the main collagen fibers direction (dashed line in Fig. 6). 415

Despite these promising results, mismatch observed between the model predictions and 416 experimental data (Fig. 5) could rise from simplifying hypotheses that were adopted at 417 different stages of the modeling. First, the transition between different hierarchies from the 418 nanoscale to the mesoscale is continuous rather than discrete in real biological structures [33]. 419 However, a limited number of length scales was accounted for in the modeling, by assuming 420 the existence of a RVE at each scale and each normalized position across the interphase, 421 in order to fulfill the separation-of-scales requirements (i.e., features at a previous scale are 422 much smaller than those at the next scale). This is not necessarily correct for the tendon-to-423 bone insertion because the features at a previous scale may not be strictly infinitesimal with 424 respect to a larger scale. Although several continuum micromechanics approaches have been 425 used to model partially mineralized tissues with a certain success [16–20, 29, 50], their use 426

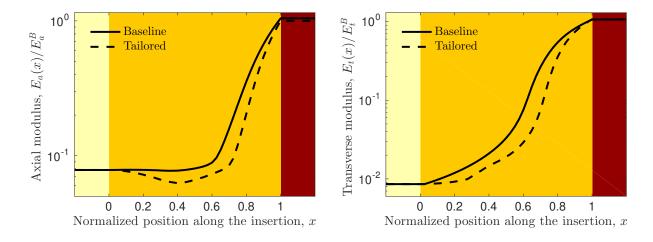


Figure 6: Mesostructural level estimates of the longitudinal elastic moduli. The baseline (continuous line) displays the axial and transverse moduli obtained from the optimal stiffness coefficients depicted in Fig. 5, whereas the tailored case (dashed line) was obtained by delaying the onset of mineralization (to  $x \approx 0.2$ ) and reducing the collagen fibers organization (angular deviation increased to  $5^{\circ}$ ).

at the nanoscale has sparked important discussions among the homogenization community 427 and is open to debate [51]. For instance, the nano-sized dimensions of mineralized tissue 428 components as well as their spatial arrangements and interactions motivated alternative 429 approaches based on discrete atomistic simulations, specifically molecular dynamics [52]. 430 A second challenging issue was the selection of the mechanical properties of the nanoscale 431 components, i.e., Young's moduli and Poisson's ratios. In this study, all nanoscale com-432 ponents (i.e., water-NCPs composite, collagen molecules and HA crystals) were assumed 433 to have a linear elastic and isotropic behavior. Wide range of values for these nanoscale 434 properties have been reported in the literature (recall Tab. 1) and different choices of such 435 properties may lead to very different results. For instance, atomistic modeling approach of 436 collagen molecules revealed that collagen may have a highly nonlinear viscoelastic behav-437 ior [53]. Likewise, the geometry and properties of HA crystals are still subject to debate. 438 This includes the shape of the mineral particles, which have been considered both as nee-439 dles [38, 54] or plates [5, 34]. Some evidences also shown that HA crystals may be oriented 440 randomly around the mineralized collagen fibrils [55], that their size can change during

the mineralization from the tendon to bone [56], or that their mechanical properties are anisotropic [57]. Third, only a few studies differentiated between intra- and extrafibrillar mineral volume fractions. Indeed, the mineralization of fibril arrays, in particular the ratio 444 between extra- and intrafibrillar mineral content, has sparked important discussions. Some 445 authors mentioned a high intrafibrillar mineral content [58, 59], whereas some others evidenced that the mineral content is mainly located in the extrafibrillar matrix [60, 61]. More 447 recent contributions underlined that mineralization pathways in bone are neither exclusively 448 intrafibrillar nor extrafibrillar [62], but rather form a continuous cross-fibrillar phase [38], 449 but this feature remains largely unexplored for the tendon-to-bone insertion [23]. Fourth, the profiles of the competing gradients in mineral content and collagen fibers orientation 451 were selected according to earlier modeling strategies [22, 23], themselves based on Raman 452 spectroscopy [26] and polarized light microscopy [42]. Nevertheless, recent experimental 453 developments for measuring these finely tuned gradients across the interphase are likely 454 to provide different (nonlinear) profiles [63, 64], which could be used as well to feed our 455 modeling approach. Fifth, based on the considered biological phases and their interactions 456 at each scale [34, 56], different methods have been used to estimate the effective stiffness tensor at each homogenization step: The Mori-Tanaka method (steps 1, 3 and 4), the self-458 consistent method (step 2), and the "inverse Mori-Tanaka" scheme (steps 5 and 6). It is 459 commonly acknwledged that the Mori-Tanaka method shows limited performance for prob-460 lems involving high volume fraction of inclusions over around 35% [65], since the inclusions' 461 distribution and interactions are not accounted for properly. Therefore, the attributions of 462 matrix and inclusions roles in our modeling steps 3 and 4 are open to question, so that other 463 homogenization methods could have been used. For instance, a rigorous approach has been proposed for composites containing multiple classes of aligned ellipsoidal inclusions with a 465 relatively high volume fraction [24], although this scheme can violate the Hashin-Shtrikman 466 bounds at low volume fraction for certain anisotropy of the phases. Lastly, our modeling ap-467 proach contains many steps, each of which should be validated experimentally in the future. 468 To date, little information is available in the literature concerning the experimental multi-469 scale evaluation of the mechanical properties at the tendon-to-bone attachment. Current 470

imaging modalities have been used to map the local content of the main constituents across
the interphase [26, 63, 10], but a direct link with the corresponding variations in mechanical properties is somehow still missing. Ongoing developments of experimental techniques
for measuring the local mechanical properties of the interphase at different scales, such as
nanoscale dynamic mechanical analysis [66] or scanning acoustic microscopy [67], are expected to provide means of comparing our numerically obtained predictions.

All these uncertainties along with some other model parameters ignored in this work, such 477 as the gradient in protein content [68], the different collagen types across the interphase [4] 478 or the unraveling of tendon fibers into smaller interphase fibrils [10], make the multiscale modeling of the tendon-to-bone insertion a challenging problem with much potential for fu-480 ture works. As such, the modeling process proposed here is straightforwardly extendable to 481 account for damage characteristics as a footprint of clinical burden that occur at the tendon-482 to-bone insertion, which may include for instance the multiscale effects of unloading (e.g., 483 stiffening of the HA foam at the nanoscale or fibers misalignment at the microscale) [56]. Be-484 sides, such model paves the way to the design of bioinspired bi-materials [69] that display the 485 functionally graded properties of the enthesis, and their characterization using quantitative 486 ultrasound [70, 71]. 487

#### 488 6. Conclusion

We modeled the tendon-to-bone insertion as a hierarchical composite material and pre-489 dicted its effective anisotropic stiffness tensor at the tissue scale. Our multiscale analysis 490 involved a bottom-up approach, starting from the nanostructural level (partially mineralized 491 collagen fibril) and then moving up the scales through the sub-microstructural level (par-492 tially mineralized fiber), the microstructural level (pattern of partially mineralized fibers), 493 up to the mesostructural level (partially mineralized tissue level). Our modeling results sup-494 ported the hypothesis that the tendon-to-bone insertion can be considered as a continuous 495 functionally graded material with respect to its mineral concentration and collagen fibers 496 organization, thus confirming earlier results obtained using phenomenological models (i.e., stochastic finite element simulations) and experimentally reported trends.

## 499 Acknowledgement

- This work was partially funded by the BEST-AMUS project (IIN program, CNRS-INSIS),
- 501 the "Support for research for newly appointed Associate Professors" and the "Bonus Qualité
- 502 Recherche" (Faculté des Sciences et Technologie, Université Paris-Est Créteil).

# 503 Appendix A. Hill tensor $\mathbb{P}^0$

Appendix A.1. Hill tensor for a cylindrical inclusion in a transversely isotropic medium. The non-zero components of the Hill tensor  $\mathbb{P}^0$  for a cylindrical inclusion embedded in a transversely isotropic matrix of stiffness  $\mathbb{C}^0$  are given according to [72] using Voigt notation,  $x_3$  being the axis of rotational symmetry,

$$P_{11}^{0} = \frac{1}{8} \frac{(C_{22}^{0} + C_{66}^{0}) + 2C_{66}^{0}}{C_{22}^{0}C_{66}^{0}}, \tag{A.1}$$

$$P_{22}^0 = \frac{1}{8} \frac{(C_{22}^0 + 3C_{66}^0)}{C_{22}^0 C_{66}^0}, \tag{A.2}$$

$$P_{12}^0 = \frac{1}{8} \frac{C_{66}^0 - C_{22}^0}{C_{22}^0 C_{66}^0} \,, \tag{A.3}$$

$$P_{66}^0 = \frac{1}{2} \frac{C_{22}^0 + C_{66}^0}{C_{22}^0 C_{66}^0}, \tag{A.4}$$

$$P_{44}^0 = \frac{1}{2C_{44}^0} \,, \tag{A.5}$$

$$P_{55}^0 = P_{44}^0 \,. \tag{A.6}$$

Appendix A.2. Hill tensors for a spherical inclusion in an isotropic medium

Considering the case of a spherical inclusion embedded in an isotropic matrix of stiffness  $\mathbb{C}^0$  [72], the components of the Hill tensor  $\mathbb{P}^0$  now read as

$$P_{11}^{0} = \frac{7C_{44}^{0} + 2C_{12}^{0}}{15C_{44}^{0}(C_{12}^{0} + 2C_{44}^{0})}, \tag{A.7}$$

$$P_{12}^{0} = \frac{C_{44}^{0} + C_{12}^{0}}{-15C_{44}^{0}(C_{12}^{0} + 2C_{44}^{0})}, \tag{A.8}$$

$$P_{44}^{0} = \frac{2(3C_{12}^{0} + 8C_{44}^{0})}{15C_{44}^{0}(C_{12}^{0} + 2C_{44}^{0})},$$
(A.9)

$$P_{22}^0 = P_{33}^0 = P_{11}^0, \ P_{55}^0 = P_{66}^0 = P_{44}^0, \ P_{13}^0 = P_{23}^0 = P_{12}^0.$$
 (A.10)

# Appendix B. Effective stiffness tensor of the hydroxyapatite foam $\mathbb{C}_{\mathrm{Hw}}(x)$

The self-consistent scheme for two interpenetrating (spherical) inclusion phases with a linear elastic and isotropic behavior can be solved according to [32]. In such a case, the nonlinear system of equations (3) can be substituted by a system of two nonlinear equations (note that the spatial variable x was omitted here for sake of clarity),

$$\frac{f_{\rm HA}(K_{\rm HA} - K_{\rm Hw})}{1 + \alpha_{\rm Hw}(K_{\rm HA} - K_{\rm Hw})/K_{\rm Hw}} + \frac{(1 - f_{\rm HA})(K_{\rm wp} - K_{\rm Hw})}{1 + \alpha_{\rm Hw}(K_{\rm wp} - K_{\rm Hw})/K_{\rm Hw}} = 0,$$
(B.1)

$$\frac{f_{\rm HA}(G_{\rm HA} - G_{\rm Hw})}{1 + \beta_{\rm Hw}(G_{\rm HA} - G_{\rm Hw})/G_{\rm Hw}} + \frac{(1 - f_{\rm HA})(G_{\rm wp} - G_{\rm Hw})}{1 + \beta_{\rm Hw}(G_{\rm wp} - G_{\rm Hw})/G_{\rm Hw}} = 0,$$
(B.2)

where the two unknowns  $K_{\text{Hw}}$  and  $G_{\text{Hw}}$  denote the bulk and shear moduli of the HA foam, respectively. The parameters  $\alpha_{\text{Hw}}$  and  $\beta_{\text{Hw}}$  are defined as

$$\alpha_{\text{Hw}} = \frac{3K_{\text{Hw}}}{3K_{\text{Hw}} + 4G_{\text{Hw}}}, \quad \beta_{\text{Hw}} = \frac{6(K_{\text{Hw}} + 2G_{\text{Hw}})}{5(3K_{\text{Hw}} + 4G_{\text{Hw}})}.$$
 (B.3)

Solving the aforementioned system yields the following components for the stiffness tensor of the hydroxyapatite foam,

$$C_{11}^{\text{Hw}} = K_{\text{Hw}} + \frac{4}{3}G_{\text{Hw}},$$
 (B.4)

$$C_{12}^{\text{Hw}} = K_{\text{Hw}} - \frac{2}{3}G_{\text{Hw}},$$
 (B.5)

$$C_{44}^{\text{Hw}} = G_{\text{Hw}},$$
 (B.6)

$$C_{22}^{\text{Hw}} = C_{33}^{\text{Hw}} = C_{11}^{\text{Hw}}, \quad C_{13}^{\text{Hw}} = C_{23}^{\text{Hw}} = C_{12}^{\text{Hw}}, \quad C_{55}^{\text{Hw}} = C_{66}^{\text{Hw}} = C_{44}^{\text{Hw}}.$$
 (B.7)

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