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1 **THE BOUSSINESQ SYSTEM WITH NON-SMOOTH BOUNDARY**
2 **CONDITIONS : EXISTENCE, RELAXATION AND TOPOLOGY**
3 **OPTIMIZATION.**

4 ALEXANDRE VIEIRA* AND PIERRE-HENRI COCQUET*†

5 **Abstract.** In this paper, we tackle a topology optimization problem which consists in finding
6 the optimal shape of a solid located inside a fluid that minimizes a given cost function. The motion
7 of the fluid is modeled thanks to the Boussinesq system which involves the unsteady Navier-Stokes
8 equation coupled to a heat equation. In order to cover several models presented in the literature, we
9 choose a non-smooth formulation for the outlet boundary conditions and an optimization parameter
10 of bounded variations. This paper aims at proving existence of solutions to the resulting equations,
11 along with the study of a relaxation scheme of the non-smooth conditions. A second part covers
12 the topology optimization problem itself for which we proved the existence of optimal solutions and
13 provides the definition of first order necessary optimality conditions.

14 **Key words.** Non-smooth boundary conditions, topology optimization, relaxation scheme, di-
15 rectional do-nothing boundary conditions

16 **AMS subject classifications.** 49K20, 49Q10, 76D03, 76D55

17 **1. Introduction.**

18 *Directional do-nothing conditions.* For many engineering applications, simula-
19 tions of flows coupled with the temperature are useful for predicting the behaviour
20 of physical designs before their manufacture, reducing the cost of the development
21 of new products. The relevance of the model and the adequacy with the experiment
22 therefore become important [16, 41, 46]. In this paper, we choose to model the flow
23 with the Boussinesq system which involves the Navier-Stokes equations coupled with
24 an energy equation. In most mathematical papers analyzing this model [8, 27, 47],
25 homogeneous Dirichlet boundary conditions are considered on the whole boundary.
26 This simplifies the mathematical analysis of the incompressible Navier-Stokes equa-
27 tion since the non-linear term vanishes after integrating by part hence simplifying the
28 derivation of a priori estimates [7, 21, 26, 47].

29 However, several applications use different boundary conditions that model inlet,
30 no-slip and outlet conditions [1]. Unlike the inlet and the no-slip conditions, the
31 outlet conditions are more subject to modelling choices. A popular choice consists in
32 using a do-nothing outlet condition (see e.g. [25, 34, 48]) which naturally comes from
33 integration by parts when defining a weak formulation of the Navier-Stokes equations.
34 However, since this outlet condition can not deal with re-entering flows, several papers
35 use a non-smooth outlet boundary conditions for their numerical simulations (see e.g.
36 [5, 23]). A focus on non-smooth outflow conditions when the temperature appears
37 can be found in [12, 23, 42, 43].

38 In particular, directional do-nothing (DDN) boundary conditions are non-smooth
39 boundary conditions that become popular. The idea is originally described in [13],
40 and several other mathematical studies followed [5, 9, 11]. These conditions were
41 considered especially for turbulent flows. In this situation, the flow may alternatively

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42 exit and re-enter the domain. These directional boundary conditions tries to capture
43 this phenomenon, while limiting the reflection. It is worth noting that other boundary
44 conditions can be used, namely the so-called local/global Bernouilli boundary condi-
45 tions [12, 23, 43]. The latter implies the do-nothing boundary condition is satisfied
46 for exiting fluid and that both the normal velocity gradient and the total pressure
47 vanish for re-entering fluid. Nevertheless, in this paper, we are going to used non-
48 smooth DDN boundary condition since they are easier to impose though a variational
49 formulation.

50 Concerning the mathematical study of Boussinesq system with directional do-
51 nothing conditions, the literature is rather scarce. To the best of our knowledge, we
52 only found [6, 15], where the steady case is studied in depth, but the unsteady case
53 only presents limited results. Indeed, while [15, p. 16, Theorem 3.2] gives existence
54 and uniqueness of a weak solution with additional regularity to the Boussinesq system
55 involving non-smooth boundary conditions at the inlet, it requires the source terms
56 and the physical constants (e.g. Reynolds, Grashof numbers) to be small enough. We
57 emphasize that these limitations comes from the proof which relies on a fixed-point
58 strategy. The first aim of this paper will then be to fill that gap by proving existence
59 and, in a two-dimensional setting, uniqueness of solutions for the unsteady Boussinesq
60 system with non-smooth DDN boundary condition on the outlet.

61 *Topology optimization.* On top of the previous considerations, this paper aims at
62 using these equations in a topology optimization (TO) framework. In fluid mechanics,
63 the term *topology optimization* refers to the problem of finding the shape of a solid
64 located inside a fluid that either minimizes or maximizes a given physical effect. There
65 exist various mathematical methods to deal with such problems that fall into the
66 class of PDE-constrained optimization, such as the topological asymptotic expansion
67 [3, 14, 40] or the shape optimization method [24, 38, 39]. In this paper, we choose to
68 locate the solid thanks to a penalization term added in the unsteady Navier-Stokes
69 equations, as exposed in [4]. However, the binary function introduced in [4] is usually
70 replaced by a smooth approximation, referred as *interpolation function* [43], in order
71 to be used in gradient-based optimization algorithms. We refer to the review papers
72 [1, 22] for many references that deal with numerical resolution of TO problems applied
73 to several different physical settings. However, as noted in [1, Section 4.7], most
74 problems tackling topology optimization for flows only focus on steady flows, and
75 time-dependant approaches are still rare. Furthermore, to the best of our knowledge,
76 no paper is dedicated to the mathematical study of unsteady TO problems involving
77 DDN boundary conditions, even though they are already used in numerical studies
78 [12, 23, 42, 43]. Therefore, a second goal of this paper will be to prove existence
79 of optimal solution to a TO problem involving Boussinesq system with non-smooth
80 DDN boundary conditions at the outlet.

81 *First order optimality conditions.* As hinted above, a gradient based method is
82 often used in order to compute an optimal solution of a TO problem. However, the
83 introduction of the non-smooth DDN boundary conditions implies that the control-
84 to-state mapping is no longer differentiable. The literature presents several ways to
85 deal with such PDE-constrained optimization problems. Most focus on elliptic equa-
86 tions, using subdifferential calculus [17, 30, 19] or as the limit of relaxation schemes
87 [18, 35, 45]. We may also cite [37] for a semilinear parabolic case and [49] which
88 involves the Maxwell equations. We emphasize that using directly a subdifferential
89 approach presents several drawbacks: the subdifferential of composite functions may
90 be hardly computed, and the result may be hardly enlightening nor used [17]. We will
91 therefore use a differentiable relaxation approach, as studied in [45]. First, we will

92 be able to use standard first order necessary optimality conditions since the relaxed
 93 control-to-state mapping will be smooth. A convergence analysis will let us design
 94 necessary optimality condition for the non-smooth problem. Secondly, we find this
 95 approach more enlightening, as it may be used as a numerical scheme for solving the
 96 TO problem.

97 **1.1. Problem settings.** Let $\Omega \subset \mathbb{R}^d$, $d \in \{2, 3\}$ be a bounded open set with
 98 Lipschitz boundary whose outward unitary normal is \mathbf{n} . We assume the fluid occupies
 99 a region $\Omega_f \subset \Omega$ and that a solid is defined by a region Ω_s such that $\Omega = \Omega_f \cup \Omega_s$.
 100 The penalized Boussinesq approximation (see e.g. [43] for the steady case) of the
 101 Navier-Stokes equations coupled to convective heat transfer reads:

$$\begin{aligned}
 & \nabla \cdot \mathbf{u} = 0, \\
 & \partial_t \theta + \nabla \cdot (\mathbf{u}\theta) - \nabla \cdot (Ck(\alpha)\nabla\theta) = 0, \quad \text{a.e. in } \Omega \\
 & \partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla)\mathbf{u} - A\Delta\mathbf{u} + \nabla p - B\theta e_y + h(\alpha)\mathbf{u} = f, \\
 & \mathbf{u}(0) = u_0(\alpha), \quad \theta(0) = \theta_0(\alpha),
 \end{aligned}
 \tag{1.1}$$

103 where \mathbf{u} denotes the velocity of the fluid, p the pressure and θ the temperature (all
 104 dimensionless), $u_0(\alpha), \theta_0(\alpha)$ are initial conditions. In (1.1), $A = \text{Re}^{-1}$ with Re being
 105 the Reynolds number, $B = \text{Ri}$ is the Richardson number and $C = (\text{RePr})^{-1}$ where Pr
 106 is the Prandtl number. In a topology optimization problem, it is classical to introduce
 107 a function $\alpha : x \in \Omega \mapsto \alpha(x) \in \mathbb{R}^+$ as optimization parameter (see e.g. [1, 22]). The
 108 function $h(\alpha)$ then penalizes the flow in order to mimic the presence of solid:

- 109 • if $h \equiv 0$, then one retrieves the classical Boussinesq approximation.
- 110 • if, for some $s > s_0$ and large enough α_{\max} , $h : s \in [0, \alpha_{\max}] \mapsto h(s) \in [0, \alpha_{\max}]$
 111 is a smooth function such that $h(s) \approx 0$ for $s \leq s_0$ and $h(s) \approx \alpha_{\max}$ for $s \geq s_0$,
 112 one retrieves the formulations used in topology optimization [1, 8, 43]. In the
 113 sequel, we work in this setting since we wish to study a TO problem.

114 Since the classical Boussinesq problem is retrieved when $h(\alpha) = 0$, the fluid zones
 115 $\Omega_f \subset \Omega$ and the solid ones $\Omega_s \subset \Omega$ can be defined as $\Omega_s := \{x \in \Omega \mid \alpha(x) < s_0\}$, $\Omega_f :=$
 116 $\{x \in \Omega \mid \alpha(x) > s_0\}$, where $\alpha_{\max} > 0$ is large enough to ensure the velocity \mathbf{u} is small
 117 enough for the Ω_s above to be considered as a solid. The function $k(\alpha) : x \in \Omega \mapsto$
 118 $k(\alpha(x))$ is the dimensionless diffusivity defined as $k(\alpha)|_{\Omega_f} = 1$ and $k(\alpha)|_{\Omega_s} = k_s/k_f$
 119 with k_s and k_f are respectively the diffusivities of the solid and the fluid. We also
 120 assume that k is a smooth regularization of $(k_s/k_f)\mathbf{1}_{\Omega_s} + \mathbf{1}_{\Omega_f}$. In this framework, α
 121 is thus defined as a parameter function, which will let us control the distribution of
 122 the solid in Ω .

123 Let us now specify the boundary conditions. Assume $\partial\Omega = \Gamma$ is Lipschitz and we
 124 split it in three parts: $\Gamma = \Gamma_w \cup \Gamma_{\text{in}} \cup \Gamma_{\text{out}}$. Here, Γ_w are the walls, Γ_{in} the inlet/entrance
 125 and Γ_{out} is the exit/outlet of the computational domain. As exposed above, we would
 126 like to rigorously study a non-smooth outlet boundary condition. Inspired by [13],
 127 the following formulation tries to encapsulate these different approaches. Let β be
 128 a function defined on Γ_{out} and define: $\forall x \in \mathbb{R} : x^+ = \text{pos}(x) = \max(0, x), x^- =$
 129 $\text{neg}(x) = \max(0, -x), x = x^+ - x^-$. On top of (1.1), we impose the following boundary
 130 conditions:

$$\begin{aligned}
 & \underline{\text{On } \Gamma_{\text{in}}} : \quad \mathbf{u} = \mathbf{u}_{\text{in}}, \quad \theta = 0, \\
 & \underline{\text{On } \Gamma_w} : \quad \mathbf{u} = 0, \quad Ck\partial_n\theta = \phi, \\
 & \underline{\text{On } \Gamma_{\text{out}}} : \quad A\partial_n\mathbf{u} - \mathbf{n}p = A\partial_n\mathbf{u}^{\text{ref}} - \mathbf{n}p^{\text{ref}} - \frac{1}{2}(\mathbf{u} \cdot \mathbf{n})^-(\mathbf{u} - \mathbf{u}^{\text{ref}}), \\
 & \quad Ck\partial_n\theta + \beta(\mathbf{u} \cdot \mathbf{n})^-\theta = 0,
 \end{aligned}
 \tag{1.2}$$

132 with $\phi \in L^2(0, T; L^2(\Gamma_w))$, $f \in L^2(0, T; H^{-1}(\Omega))$, $\mathbf{u}_{\text{in}} \in L^2(0, T; H_{00}^{1/2}(\Gamma_{\text{in}}))$, \mathbf{n} de-
 133 notes the normal vector to the boundary, $\partial_n = \mathbf{n} \cdot \nabla$ and $(\mathbf{u}^{\text{ref}}, p^{\text{ref}})$ denotes a reference
 134 solution. As stated in [29], this nonlinear condition is physically meaningful: if the
 135 flow is outward, we impose the constraint coming from the selected reference flow ;
 136 if it is inward, we need to control the increase of energy, so, according to Bernoulli's
 137 principle, we add a term that is quadratic with respect to velocity.

138 To define a weak formulation of (1.1)-(1.2), we introduce $\mathcal{V}^u = \{\mathbf{u} \in \mathcal{C}^\infty(\Omega; \mathbb{R}^d);$
 139 $\mathbf{u}|_{\Gamma_{\text{in}} \cup \Gamma_w} = 0\}$, and define V^u (resp H^u) as the closure of \mathcal{V}^u in $(H^1(\Omega))^d$ (resp. in
 140 $(L^2(\Omega))^d$). Similarly, we define $\mathcal{V}^\theta = \{\theta \in \mathcal{C}^\infty(\Omega; \mathbb{R}); \theta|_{\Gamma_{\text{in}}} = 0\}$, and V^θ and H^θ
 141 accordingly. A weak formulation of (1.1)-(1.2) then reads as:

$$142 \quad \int_{\Omega} \partial_t \theta \varphi - \int_{\Omega} \theta \mathbf{u} \cdot \nabla \varphi + \int_{\Omega} Ck \nabla \theta \cdot \nabla \varphi + \int_{\Gamma} (\theta(\mathbf{u} \cdot \mathbf{n}) - Ck \partial_n \theta) \varphi = 0,$$

143 for all $\varphi \in H^\theta$. However, from (1.2), we have:

$$\begin{aligned} \int_{\Gamma} (\theta(\mathbf{u} \cdot \mathbf{n}) - Ck \partial_n \theta) \varphi &= - \int_{\Gamma_w} \phi \varphi + \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta(\mathbf{u} \cdot \mathbf{n})^-) \theta \varphi \\ &\quad - \int_{\Gamma_{\text{out}}} (\beta \theta(\mathbf{u} \cdot \mathbf{n})^- + Ck \partial_n \theta) \varphi \\ 144 \quad &= - \int_{\Gamma_w} \phi \varphi + \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta(\mathbf{u} \cdot \mathbf{n})^-) \theta \varphi. \end{aligned}$$

145 Therefore:

(WF.1)

$$146 \quad \int_{\Omega} \partial_t \theta \varphi - \int_{\Omega} \theta \mathbf{u} \cdot \nabla \varphi + \int_{\Omega} Ck \nabla \theta \cdot \nabla \varphi + \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta(\mathbf{u} \cdot \mathbf{n})^-) \theta \varphi = \int_{\Gamma_w} \phi \varphi.$$

147 Doing similar computations with the Navier-Stokes system yield:

$$150 \quad (\text{WF.2}) \quad \int_{\Omega} q \nabla \cdot \mathbf{u} = 0, \quad \forall q \in L^2(\Omega),$$

151

$$\begin{aligned} &\int_{\Omega} (\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u}) \cdot \Psi + A \int_{\Omega} \nabla \mathbf{u} : \nabla \Psi - \int_{\Omega} B \theta e_y \cdot \Psi \\ 152 \quad (\text{WF.3}) \quad &- \int_{\Omega} p \nabla \cdot \Psi + \int_{\Omega} h \mathbf{u} \cdot \Psi + \frac{1}{2} \int_{\Gamma_{\text{out}}} (\mathbf{u} \cdot \mathbf{n})^- (\mathbf{u} - \mathbf{u}^{\text{ref}}) \cdot \Psi \\ &= \int_{\Omega} f \cdot \Psi + \int_{\Gamma_{\text{out}}} (A \partial_n \mathbf{u}^{\text{ref}} - \mathbf{n} p^{\text{ref}}) \cdot \Psi, \end{aligned}$$

153 for all $\Psi \in H^u$.

154 **1.2. The topology optimization problem.** A goal of this paper is to analyze
 155 the next topology optimization problem

$$156 \quad (\text{OPT}) \quad \begin{aligned} &\min \mathcal{J}(\alpha, \mathbf{u}, \theta, p) \\ &\text{s.t.} \quad \begin{cases} (\mathbf{u}, \theta, p) \text{ solution of (WF) parametrized by } \alpha, \\ \alpha \in \mathcal{U}_{\text{ad}}, \end{cases} \end{aligned}$$

157 where \mathcal{J} is a given cost function and, for some $\kappa > 0$, we set $\mathcal{U}_{\text{ad}} = \{\alpha \in \text{BV}(\Omega)$
 158 $: 0 \leq \alpha(x) \leq \alpha_{\text{max}} \text{ a.e. on } \Omega, |D\alpha|(\Omega) \leq \kappa\}$. $\text{BV}(\Omega)$ stands for functions of bounded

159 variations, as exposed in [2]. We recall that the weak-* convergence in $BV(\Omega)$ is
 160 defined as follows [2]: $(\alpha_\varepsilon)_\varepsilon \subset BV(\Omega)$ weakly-* converges to $\alpha \in BV(\Omega)$ if (α_ε)
 161 strongly converges to α in $L^1(\Omega)$ and $(D\alpha_\varepsilon)$ weakly-* converges to $D\alpha$ in Ω , meaning:

$$162 \quad \lim_{\varepsilon \rightarrow +\infty} \int_{\Omega} \nu dD\alpha_\varepsilon = \int_{\Omega} \nu dD\alpha, \quad \forall \nu \in C_0(\Omega),$$

163 where $C_0(\Omega)$ denotes the closure, in the sup norm, of the set of real continuous
 164 functions with compact support over Ω . We choose \mathcal{U}_{ad} as a subset of $BV(\Omega)$ since
 165 it is a nice way to approximate piecewise constant functions, which is close to the
 166 desired solid distribution.

167 It is classical for these problems to compute first order optimality conditions
 168 (see e.g. [33, 44]). This approach needs smoothness of the control-to-state mapping.
 169 However, the presence of the non-differentiable function $\text{neg}(x) = x^-$ makes this
 170 approach impossible. Therefore, we adopt a smoothing approach, as studied in [35,
 171 45], and we approximate the neg function with a C^1 positive approximation, denoted
 172 neg_ε . We suppose this approximation satisfies the following assumptions:

173 **(A1)** $\forall x \in \mathbb{R}, \text{neg}_\varepsilon(x) \geq \text{neg}(x)$.

174 **(A2)** $\forall x \in \mathbb{R}, 0 \leq \text{neg}'_\varepsilon(x) \leq 1$.

175 **(A3)** neg_ε converges to neg uniformly over \mathbb{R} .

176 **(A4)** for every $\delta > 0$, the sequence $(\text{neg}'_\varepsilon)_{\varepsilon > 0}$ converges uniformly to 1 on $[\delta, +\infty)$
 177 and uniformly to 0 on $(-\infty, -\delta]$ as $\varepsilon \rightarrow +\infty$.

178 As presented in [45], we may choose:

$$179 \quad (1.3) \quad \text{neg}_\varepsilon(x) = \begin{cases} x^- & \text{if } |x| \geq \frac{1}{2\varepsilon}, \\ \frac{1}{2} \left(\frac{1}{2} - \varepsilon x\right)^3 \left(\frac{3}{2\varepsilon} + x\right) & \text{if } |x| < \frac{1}{2\varepsilon}. \end{cases}$$

180 We thus redefine (WF) with an approximation of neg , which gives:

$$181 \quad (\text{WFe.1}) \quad \begin{aligned} & \int_{\Omega} (\partial_t \mathbf{u}_\varepsilon + (\mathbf{u}_\varepsilon \cdot \nabla) \mathbf{u}_\varepsilon) \cdot \Psi + A \int_{\Omega} \nabla \mathbf{u}_\varepsilon : \nabla \Psi - \int_{\Omega} B \theta e_y \cdot \Psi \\ & + \int_{\Omega} h(\alpha_\varepsilon) \mathbf{u}_\varepsilon \cdot \Psi - \int_{\Omega} p_\varepsilon \nabla \cdot \Psi + \frac{1}{2} \int_{\Gamma_{out}} \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) (\mathbf{u}_\varepsilon - \mathbf{u}_\varepsilon^{\text{ref}}) \cdot \Psi \\ & = \int_{\Omega} f \cdot \Psi - \int_{\Gamma_{out}} (A \partial_n \mathbf{u}^{\text{ref}} + \mathbf{n} p^{\text{ref}}) \cdot \Psi, \end{aligned}$$

182

$$183 \quad (\text{WFe.2}) \quad \int_{\Omega} \mathbf{u}_\varepsilon \cdot \nabla q = 0,$$

184

$$185 \quad (\text{WFe.3}) \quad \begin{aligned} & \int_{\Omega} \partial_t \theta_\varepsilon \varphi - \int_{\Omega} \theta_\varepsilon \mathbf{u}_\varepsilon \cdot \nabla \varphi + \int_{\Omega} Ck \nabla \theta_\varepsilon \cdot \nabla \varphi \\ & + \int_{\Gamma_{out}} ((\mathbf{u}_\varepsilon \cdot \mathbf{n}) + \beta \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})) \theta_\varepsilon \varphi = \int_{\Gamma_w} \phi \varphi, \end{aligned}$$

186 for all $(\Psi, \varphi, q) \in H^u \times H^\theta \times L^2(\Omega)$. We then define the approximate optimal control
 187 problem:

$$188 \quad (\text{OPTe}) \quad \begin{aligned} & \min \mathcal{J}(\alpha_\varepsilon, \mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon) \\ & \text{s.t.} \quad \begin{cases} (\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon) \text{ solution of (WFe) parametrized by } \alpha_\varepsilon, \\ \alpha_\varepsilon \in \mathcal{U}_{ad}. \end{cases} \end{aligned}$$

189 As it will be made clear later, the control-to-state mapping in (WFe) is smooth, which
 190 will let us derive first order conditions.

191 **1.3. Plan of the paper.** The rest of this introduction is dedicated to the pre-
 192 sentation of some notations used in this article and some important results of the
 193 literature. The core of this paper is organized in two sections. First, we will prove the
 194 existence of solutions to (WFe), which will let us prove, with a compactness argument,
 195 the existence of solutions to (WF). We then focus on the two dimensional case, where
 196 we prove uniqueness of the solutions along with stronger convergence results. This
 197 is an extension of the work done by [13], where only the pressure and the velocity
 198 where considered, and to [6, 15], where the steady case was studied in depth, but the
 199 results concerning the unsteady case were obtained using restrictive assumptions. We
 200 then study the approximate optimal control problem (OPTe), for which we will derive
 201 first order conditions. We conclude this paper with the convergence of the optimality
 202 conditions of (OPTe), which let us design first order conditions of (OPT).

203 *Notations.* We denote by $a \lesssim b$ if there exists a constant $C(\Omega) > 0$ depending
 204 only on Ω such that $a \leq C(\Omega)b$. Denote:

- 205 • $\mathcal{A} : V^u \rightarrow (V^u)'$ defined by $\langle \mathcal{A}\mathbf{u}, \mathbf{v} \rangle_{(V^u)', V^u} = A \int_{\Omega} \nabla \mathbf{u} : \nabla \mathbf{v}$,
- 206 • $\mathcal{B} : V^u \times V^u \rightarrow (V^u)'$ defined by $\langle \mathcal{B}(\mathbf{u}, \mathbf{v}), \mathbf{w} \rangle_{(V^u)', V^u} = \int_{\Omega} (\mathbf{u} \cdot \nabla) \mathbf{v} \cdot \mathbf{w}$,
- 207 • $\mathcal{T} : V^{\theta} \rightarrow (V^u)'$ defined by $\langle \mathcal{T}\theta, \mathbf{v} \rangle_{(V^u)', V^u} = \int_{\Omega} B\theta e_y \cdot \mathbf{v}$,
- 208 • $\mathcal{P} : L^2(\Omega) \rightarrow (V^u)'$ defined by $\langle \mathcal{P}p, \mathbf{w} \rangle_{(V^u)', V^u} = \int_{\Omega} p \nabla \cdot \mathbf{w}$,
- 209 • $\mathcal{N} : V^u \times V^u \rightarrow (V^u)'$ defined by $\langle \mathcal{N}(\mathbf{u}, \mathbf{v}), \mathbf{w} \rangle_{(V^u)', V^u} = \int_{\Gamma_{\text{out}}} \text{neg}(\mathbf{u} \cdot \mathbf{n})(\mathbf{v} \cdot \mathbf{w})$,
- 210 • $\mathcal{N}_{\varepsilon} : V^u \times V^u \rightarrow (V^u)'$ given by $\langle \mathcal{N}_{\varepsilon}(\mathbf{u}, \mathbf{v}), \mathbf{w} \rangle_{(V^u)', V^u} = \int_{\Gamma_{\text{out}}} \text{neg}_{\varepsilon}(\mathbf{u} \cdot \mathbf{n})(\mathbf{v} \cdot \mathbf{w})$.
- 211 • $\mathcal{C}(\alpha) : V^{\theta} \rightarrow (V^{\theta})'$ defined by $\langle \mathcal{C}(\alpha)\theta, \varphi \rangle_{(V^{\theta})', V^{\theta}} = \int_{\Omega} Ck(\alpha) \nabla \theta \cdot \nabla \varphi$,
- 212 • $\mathcal{D} : V^u \times V^{\theta} \rightarrow (V^{\theta})'$ defined by $\langle \mathcal{D}(\mathbf{u}, \theta), \varphi \rangle_{(V^{\theta})', V^{\theta}} = \int_{\Omega} \theta \mathbf{u} \cdot \nabla \varphi$,
- 213 • $\mathcal{M} : V^u \times V^{\theta} \rightarrow (V^{\theta})'$ defined by $\langle \mathcal{M}(\mathbf{u}, \theta), \varphi \rangle_{(V^{\theta})', V^{\theta}} = \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta \text{neg}(\mathbf{u} \cdot \mathbf{n})) \theta \varphi$,
- 214 • $\mathcal{M}_{\varepsilon} : V^u \times V^{\theta} \rightarrow (V^{\theta})'$ defined by $\langle \mathcal{M}_{\varepsilon}(\mathbf{u}, \theta), \varphi \rangle_{(V^{\theta})', V^{\theta}} = \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta \text{neg}_{\varepsilon}(\mathbf{u} \cdot \mathbf{n})) \theta \varphi$,

215 By a slight abuse of notation, we will still denote by σ^{ref} the element of $(V^u)'$ defined
 216 by $\langle \sigma^{\text{ref}}, \mathbf{w} \rangle_{(V^u)', V^u} = \int_{\Gamma_{\text{out}}} (A \partial_n \mathbf{u}^{\text{ref}} - p^{\text{ref}} \mathbf{n}) \cdot \mathbf{w}$, $h(\alpha) : V^u \rightarrow (V^u)'$ the function
 217 defined by $\langle h(\alpha)\mathbf{u}, \mathbf{v} \rangle_{(V^u)', V^u} = \int_{\Omega} h(\alpha) \mathbf{u} \cdot \mathbf{v}$, and ϕ the element of $(V^{\theta})'$ defined by
 218 $\langle \phi, \varphi \rangle_{(V^{\theta})', V^{\theta}} = \int_{\Gamma_{\text{out}}} \phi \varphi$.

219 *Results from the literature.* We now recall two results from the literature that will
 220 be heavily used throughout this paper.

221 **PROPOSITION 1.1.** ([10, Proposition III.2.35]) *Let Ω be a Lipschitz domain of*
 222 *\mathbb{R}^d with compact boundary. Let $p \in [1, +\infty]$ and $q \in [p, p^*]$, where p^* is the critical*
 223 *exponent associated with p , defined as:*

$$224 \begin{cases} \frac{1}{p^*} = \frac{1}{p} - \frac{1}{d} & \text{for } p < d, \\ p^* \in [1, +\infty[& \text{for } p = d, \\ p^* = +\infty & \text{for } p > d. \end{cases}$$

225 *Then, there exists a positive constant C such that, for any $u \in W^{1,p}(\Omega)$:*

$$226 \|u\|_{L^q(\Omega)} \leq C \|u\|_{L^p(\Omega)}^{1 + \frac{d}{q} - \frac{d}{p}} \|u\|_{W^{1,p}(\Omega)}^{\frac{d}{p} - \frac{d}{q}}.$$

230 PROPOSITION 1.2. ([10, Theorem III.2.36]) Let Ω be a Lipschitz domain of \mathbb{R}^d
 231 with compact boundary, and $1 < p < d$. Then for any $r \in \left[p, \frac{p(d-1)}{d-p} \right]$, there exists a
 232 positive constant C such that, for any $u \in W^{1,p}(\Omega)$:

$$233 \quad \|u|_{\partial\Omega}\|_{L^r(\partial\Omega)} \leq C \|u\|_{L^p(\Omega)}^{1-\frac{d}{p}+\frac{d-1}{r}} \|u\|_{W^{1,p}(\Omega)}^{\frac{d}{p}-\frac{d-1}{r}}.$$

234 In the case $p = d$, the previous result holds true for any $r \in [p, +\infty[$.

235 **2. Existence of solutions.** In this section, we will focus on proving the existence
 236 of solutions to (WFe) and prove their convergence toward the ones of (WF).

237 We make the following assumptions throughout this paper:

238 ASSUMPTIONS 2.1. • The source term $f \in L^2(0, T; H^{-1}(\Omega))$.

239 • $(\mathbf{u}^{ref}, p^{ref})$ are such that:

$$240 \quad \begin{cases} \mathbf{u}^{ref} \in L^r(0, T; (H^1(\Omega))^d) \cap L^\infty(0, T; (L^2(\Omega))^d) \\ \text{with } r = 2 \text{ if } d = 2 \text{ and } r = 4 \text{ if } d = 3, \\ \nabla \cdot \mathbf{u}^{ref} = 0, \\ \partial_t \mathbf{u}^{ref} \in L^2(0, T; (L^2(\Omega))^d), \\ \mathbf{u}^{ref} = \mathbf{u}_{in} \text{ on } \Gamma_{in}. \end{cases}$$

- 241 • There exists k_{min} such that $k(x) \geq k_{min} > 0$ and $h(x) \geq 0$ for a.e. $x \in \Omega$.
- 242 • The initial condition \mathbf{u}_0 (resp. θ_0) is a Fréchet-differentiable function from
- 243 \mathcal{U}_{ad} to V^u (resp. V^θ). Furthermore, for all $\alpha \in \mathcal{U}_{ad}$, $\mathbf{u}_0(\alpha)|_{\Gamma_{in}} = \mathbf{u}_{in}(0)$,
- 244 $\mathbf{u}_0(\alpha)|_{\Gamma_w} = 0$, $\nabla \cdot \mathbf{u}_0(\alpha) = 0$ and $\theta_0(\alpha)|_{\Gamma_{in}} = 0$.
- 245 • $\beta \in L^\infty(0, T; L^\infty(\Gamma_{out}))$ such that $\beta(t, x) \geq \frac{1}{2}$, for a.e. $(t, x) \in [0, T] \times \Gamma_{out}$.

246 **2.1. Existence in dimension 2 or 3.** In this part, we work with a fixed $\varepsilon > 0$
 247 and a given α_ε in \mathcal{U}_{ad} .

248 In order to prove the existence of solutions to (WFe), we follow the classical Fedeo-
 249 Galerkin method, as used in [13, 36, 47]. By construction, V^u and V^θ are separable.
 250 Therefore, both admit a countable Hilbert basis $(w_k^u)_k$ and $(w_k^\theta)_k$. Let us construct
 251 an approximate problem, which will converge to a solution of the original problem
 252 (WFe). Denote by V_n^u (resp. V_n^θ) the space spanned by $(w_k^u)_{k \leq n}$ (resp. $(w_k^\theta)_{k \leq n}$).

253 We consider the following Galerkin approximated problem:

254 find $t \mapsto \mathbf{v}_n(t) \in V_n^u$, $t \mapsto p_n(t) \in L^2(\Omega)$ and $t \mapsto \theta_n(t) \in V_n^\theta$ such that, defining
 255 $\mathbf{u}_n = \mathbf{v}_n + \mathbf{u}^{ref}$, $(\mathbf{u}_n, p_n, \theta_n)$ satisfy (WFe) for all $t \in [0, T]$ and for all $(\Psi, q, \varphi) \in$
 256 $V_n^u \times L^2(\Omega) \times V_n^\theta$.

257 As done in [47], we prove that such $(\mathbf{u}_n, \theta_n, p_n)$ exist. We now prove that these
 258 solutions are bounded with respect to n and ε :

259 PROPOSITION 2.2. There exist positive constants c_1^θ , c_2^θ , c_1^Ψ and c_2^Ψ , independent
 260 of ε and n , such that:

$$261 \quad (2.1) \quad \sup_{[0, T]} \|\theta_n\|_{L^2(\Omega)} \leq c_1^\theta, \quad 265 \quad (2.3) \quad \sup_{[0, T]} \|\mathbf{v}_n\|_{L^2(\Omega)} \leq c_1^\Psi,$$

$$262 \quad (2.2) \quad \int_0^T \|\nabla \theta_n\|_{L^2(\Omega)}^2 \leq c_2^\theta, \quad 266 \quad (2.4) \quad \int_0^T \|\nabla \mathbf{v}_n\|_{L^2(\Omega)}^2 \leq c_2^\Psi.$$

268 *Proof.* Taking $\varphi_n = \theta_n$ in (WFe.1) and integrating by part give:

$$269 \quad \begin{aligned} & \frac{d}{dt} \|\theta_n\|_{L^2(\Omega)}^2 - \frac{1}{2} \int_{\Gamma_{\text{out}}} \theta_n^2 (\mathbf{u}_n \cdot \mathbf{n}) + \int_{\Omega} Ck |\nabla \theta_n|^2 \\ & + \int_{\Gamma_{\text{out}}} ((\mathbf{u}_n \cdot \mathbf{n}) + \beta \text{neg}_{\varepsilon} (\mathbf{u}_n \cdot \mathbf{n})) \theta_n^2 = \int_{\Gamma_w} \phi \theta_n. \end{aligned}$$

270 Since $\beta \geq \frac{1}{2}$ and using assumption (A1), one has on Γ_{out} :

$$271 \quad \begin{aligned} ((\mathbf{u}_n \cdot \mathbf{n}) + \beta \text{neg}_{\varepsilon} (\mathbf{u}_n \cdot \mathbf{n})) \theta_n^2 - \frac{1}{2} (\mathbf{u}_n \cdot \mathbf{n}) \theta_n^2 & \geq \frac{1}{2} ((\mathbf{u}_n \cdot \mathbf{n}) + \text{neg}_{\varepsilon} (\mathbf{u}_n \cdot \mathbf{n})) \theta_n^2 \\ & \geq \frac{1}{2} (\mathbf{u}_n \cdot \mathbf{n})^+ \theta_n^2 \geq 0. \end{aligned}$$

272 Therefore: $\frac{d}{dt} \|\theta_n\|_{L^2(\Omega)}^2 + Ck_{\min} \|\nabla \theta_n\|_{L^2(\Omega)}^2 \leq \|\phi\|_{L^2(\Gamma_w)} \|\theta_n\|_{L^2(\Gamma_w)}$. Using continuity
273 of the trace operator and Young's inequality, one proves that there exists a positive
274 constant $C(\Omega)$ such that, for any $\nu > 0$:

$$275 \quad \frac{d}{dt} \|\theta_n\|_{L^2(\Omega)}^2 + Ck_{\min} \|\nabla \theta_n\|_{L^2(\Omega)}^2 \leq \frac{1}{2\nu} \|\phi\|_{L^2(\Gamma_w)}^2 + \frac{C(\Omega)\nu}{2} (\|\theta_n\|_{L^2(\Omega)}^2 + \|\nabla \theta_n\|_{L^2(\Omega)}^2).$$

276 Taking ν small enough, we are left with:

$$277 \quad \frac{d}{dt} \|\theta_n\|_{L^2(\Omega)}^2 \leq \frac{1}{2\nu} \|\phi\|_{L^2(\Gamma_w)}^2 + \frac{C(\Omega)\nu}{2} \|\theta_n\|_{L^2(\Omega)}^2.$$

278 Integrating this equation and using Gronwall's lemma then give (2.1) and (2.2).

279 Now, take $\Psi_n = \mathbf{v}_n$ in (WFe.3). After some calculations, one gets:

$$280 \quad \begin{aligned} & \frac{d}{dt} |\mathbf{v}_n|^2 + A |\nabla \mathbf{v}_n|^2 + \frac{1}{2} \int_{\Gamma_{\text{out}}} \text{neg}_{\varepsilon} (\mathbf{u}_n \cdot \mathbf{n}) |\mathbf{v}_n|^2 + \int_{\Omega} h |\mathbf{v}_n|^2 \\ & = \int_{\Omega} f_{\theta} \cdot \mathbf{v}_n - \int_{\Omega} \partial_t \mathbf{u}^{\text{ref}} \cdot \mathbf{v}_n - A \int_{\Omega} \nabla \mathbf{u}^{\text{ref}} : \nabla \mathbf{v}_n + \int_{\Omega} (\mathbf{u}_n \cdot \nabla) \mathbf{v}_n \cdot \mathbf{u}^{\text{ref}} \\ & - \int_{\Omega} h \mathbf{u}^{\text{ref}} \cdot \mathbf{v}_n + \int_{\Gamma_{\text{out}}} (A \partial_n \mathbf{u}^{\text{ref}} - \mathbf{n} p^{\text{ref}}) \mathbf{v}_n \end{aligned}$$

281 where $f_{\theta} = f + B\theta_n e_y$. First, using (2.2), one has $\|f_{\theta}\|_{(H^u)'} \leq \|f\|_{(H^u)'} + Bc_1^{\theta}$.
282 Secondly, using (A1) gives that $\int_{\Gamma_{\text{out}}} \text{neg}_{\varepsilon} (\mathbf{u}_n \cdot \mathbf{n}) |\mathbf{v}_n|^2 \geq 0$ and following then the
283 same pattern of proof as in [13, Proposition 2], one proves (2.3) and (2.4). \square

284 Following [47, 10], we need to bound the fractional derivatives of the solution in
285 order to prove some convergence results. For any real-valued function f defined on
286 $[0, T]$, define by \tilde{f} the extension by 0 of f to the whole real line \mathbb{R} , and by $\mathcal{F}(\tilde{f})$
287 the Fourier transform of \tilde{f} , which we define as: $\mathcal{F}(\tilde{f})(\tau) = \int_{\mathbb{R}} \tilde{f}(t) e^{-it\tau} dt$. Using the
288 Hausdorff-Young inequality [10, Theorem II.5.20] we can prove the

289 **PROPOSITION 2.3.** *For all $\sigma \in [0, \frac{1}{6}]$, there exists a constant $C(\sigma) > 0$ indepen-*
290 *dent of ε and n such that:*

$$291 \quad (2.5) \quad \int_{\mathbb{R}} |\tau|^{2\sigma} \left\| \mathcal{F}(\tilde{\theta}_n) \right\|_{(L^2(\Omega))^d}^2 \leq C(\sigma),$$

292

$$293 \quad (2.6) \quad \int_{\mathbb{R}} |\tau|^{2\sigma} \left\| \mathcal{F}(\tilde{\mathbf{u}}_n) \right\|_{L^2(\Omega)}^2 \leq C(\sigma).$$

294 *Proof.* We emphasize that (2.6) is proved if (2.5) holds by using [10, Proposition
 295 VII.1.3] by replacing f by $f_\theta = f + B\theta e_y$. The proof of (2.5) consists in adapting the
 296 one of [10, Proposition VII.1.3] and is thus omitted. \square

297 Combining the two previous results, we can now prove the following existence
 298 theorem for (WFe).

299 **THEOREM 2.4.** *For all $(\mathbf{v}_0, \theta_0) \in H^u \times H^\theta$ and all $T > 0$, there exists $\mathbf{v}_\varepsilon \in$
 300 $L^\infty(0, T; H^u) \cap L^2(0, T; V^u)$, $\theta_\varepsilon \in L^\infty(0, T; H^\theta) \cap L^2(0, T; V^\theta)$ and $p_\varepsilon \in W^{-1, \infty}(0, T;$
 301 $L^2(\Omega))$ solution of (WFe) such that, defining $\mathbf{u}_0 = \mathbf{v}_0 + \mathbf{u}^{ref}(0)$ and $\mathbf{u}_\varepsilon = \mathbf{v}_\varepsilon + \mathbf{u}^{ref}$,
 302 one has for all $(\Psi, \varphi) \in V^u \times V^\theta$ such that $\nabla \cdot \Psi = 0$: $(\int_\Omega \mathbf{u}_\varepsilon \cdot \Psi)(0) = \int_\Omega \mathbf{u}_0 \cdot \Psi$,
 303 $(\int_\Omega \theta_\varepsilon \varphi)(0) = \int_\Omega \theta_0 \varphi$. Moreover, one has $\mathbf{v}'_\varepsilon = \frac{d\mathbf{v}_\varepsilon}{dt} \in L^{\frac{4}{3}}(0, T; (V^u)')$ and $\theta'_\varepsilon \in$
 304 $L^2(0, T; (V^\theta)')$.*

305 *Proof.* The proof of existence is similar to part (iv) of the proof of [47, Theorem
 306 3.1] and the proof of [10, Proposition VII.1.4], where estimates (2.1)-(2.4) and (2.5)-
 307 (2.6) are used in a compactness argument.

308 We only add the proof that (\mathbf{u}_n, θ_n) converges to a solution of (WFe.1). Using
 309 (2.5) and [47, Theorem 2.2], one shows that, up to a subsequence, θ_n strongly
 310 converges to an element θ_ε of $L^2(0, T; H^\theta)$. The only technical points which needs more
 311 detail are the non-linear terms in (WFe.1). Using the strong convergence of \mathbf{u}_n to \mathbf{u}_ε
 312 in $L^2(0, T; H^u)$ proved in [47, Eq (3.41)], one proves that $(\theta_n \mathbf{u}_n)$ strongly converges
 313 to $\theta_\varepsilon \mathbf{u}_\varepsilon$ in $L^1(0, T; L^2(\Omega))$. Furthermore, notice that:

$$\begin{aligned} \int_0^T \|(\mathbf{u}_n \cdot \mathbf{n})\theta_n\|_{L^{\frac{4}{3}}(\Gamma)}^{\frac{4}{3}} &\leq \int_0^T \|\mathbf{u}_n\|_{L^{\frac{4}{3}}(\Gamma)}^{\frac{4}{3}} \|\theta_n\|_{L^{\frac{4}{3}}(\Gamma)}^{\frac{4}{3}} \\ &\leq C \int_0^T \|\mathbf{u}_n\|_{L^2(\Omega)}^{\frac{1}{3}} \|\theta_n\|_{L^2(\Omega)}^{\frac{1}{3}} \|\mathbf{u}_n\|_{H^1(\Omega)} \|\theta_n\|_{H^1(\Omega)} \\ &\leq C \|\mathbf{u}_n\|_{L^\infty(0, T; L^2(\Omega))}^{\frac{1}{3}} \|\theta_n\|_{L^\infty(0, T; L^2(\Omega))}^{\frac{1}{3}} \\ &\quad \|\mathbf{u}_n\|_{L^2(0, T; H^1(\Omega))} \|\theta_n\|_{L^2(0, T; H^1(\Omega))}. \end{aligned}$$

315 This inequality together with (2.1)-(2.4) proves that $((\mathbf{u}_n \cdot \mathbf{n})\theta_n)_n$ is bounded in
 316 $L^{\frac{4}{3}}(0, T; L^{\frac{4}{3}}(\Gamma))$, which is reflexive. Therefore, it proves that, up to a subsequence,
 317 there exists a weak limit κ_1 in $L^{\frac{4}{3}}(0, T; L^{\frac{4}{3}}(\Gamma))$ of $((\mathbf{u}_n \cdot \mathbf{n})\theta_n)_n$. A simple adapta-
 318 tion of the above reasoning proves that $(\text{neg}_\varepsilon(\mathbf{u}_n \cdot \mathbf{n})\theta_n)_n$ weakly converges to some
 319 κ_2 in $L^{\frac{4}{3}}(0, T; L^{\frac{4}{3}}(\Gamma))$. Using the strong convergence of θ_n in $L^2(0, T; L^2(\Omega))$, [10,
 320 Proposition II.2.12] implies that:

$$321 \quad ((\mathbf{u}_n \cdot \mathbf{n}) + \beta \text{neg}_\varepsilon(\mathbf{u}_n \cdot \mathbf{n}))\theta_n \rightharpoonup ((\mathbf{u}_\varepsilon \cdot \mathbf{n}) + \beta \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}))\theta_\varepsilon \text{ in } L^{\frac{4}{3}}(0, T; L^1(\Gamma))$$

322 obtained using the continuity of $x \in \mathbb{R} \mapsto \text{neg}_\varepsilon(x)$. By uniqueness of the limit in
 323 the sense of distribution, we can identify $\kappa_1 + \beta\kappa_2$ with $((\mathbf{u}_\varepsilon \cdot \mathbf{n}) + \beta \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}))\theta_\varepsilon$.
 324 Therefore, $(\mathbf{u}_\varepsilon, \theta_\varepsilon)$ is a solution of (WFe.1).

325 The convergence of the weak derivative with respect to time of \mathbf{v}_ε in $L^{\frac{4}{3}}(0, T;$
 326 $(V^u)')$ is proved in [10, Proposition V.1.3]. Concerning the weak derivative with
 327 respect to time of θ_ε , it follows immediately from the fact that differentiation with
 328 respect to time is continuous in the sense of distribution. Existence of the pressure
 329 p_ε follows from [10, Chapter V]. \square

330 We now use the existence of solutions to the approximate problem (WFe) to prove
 331 existence of solutions to the limit problem (WF), along with the convergence of the
 332 approximate solutions to the solutions of (WF).

333 THEOREM 2.5. Let $(\alpha_\varepsilon) \subset \mathcal{U}_{ad}$ and $\alpha \in \mathcal{U}_{ad}$ such that $\alpha_\varepsilon \overset{*}{\rightharpoonup} \alpha$ in BV. Define by
 334 $(\mathbf{v}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ a solution of (WFe) parametrized by α_ε , and define $\mathbf{u}_\varepsilon = \mathbf{v}_\varepsilon + \mathbf{u}^{\text{ref}}$. Then,
 335 there exists $(\mathbf{v}, \theta, p) \in L^\infty(0, T; H^u) \cap L^2(0, T; V^u) \times L^\infty(0, T, H^\theta) \cap L^2(0, T; V^\theta) \times$
 336 $W^{-1, \infty}(0, T; L^2(\Omega))$ such that, defining $\mathbf{u} = \mathbf{v} + \mathbf{u}^{\text{ref}}$, up to a subsequence, we have

- 337 • $\mathbf{u}_\varepsilon \overset{*}{\rightharpoonup} \mathbf{u}$ in $L^\infty(0, T; H^u)$ and $\theta_\varepsilon \overset{*}{\rightharpoonup} \theta$ in $L^\infty(0, T; H^\theta)$,
- 338 • $\mathbf{u}_\varepsilon \rightharpoonup \mathbf{u}$ in $L^2(0, T; V^u)$ and in $L^2(0, T; (L^6(\Omega)))$,
- 339 • $\theta_\varepsilon \rightharpoonup \theta$ in $L^2(0, T; V^\theta)$ and in $L^2(0, T; (L^6(\Omega)))$,
- 340 • $\mathbf{u}_\varepsilon \rightharpoonup \mathbf{u}$ in $L^4(0, T; (L^2(\Gamma))^d)$ and $\theta_\varepsilon \rightharpoonup \theta$ in $L^4(0, T; (L^2(\Gamma)))$,
- 341 • $\mathbf{u}_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} \mathbf{u}$ in $L^2(0, T; (L^2(\Omega))^d)$ and $\theta_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} \theta$ in $L^2(0, T; (L^2(\Omega)))$,
- 342 • $\mathbf{u}_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} \mathbf{u}$ in $L^2(0, T; (L^2(\Gamma))^d)$ and $\theta_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} \theta$ in $L^2(0, T; (L^2(\Gamma)))$,
- 343 • $p_\varepsilon \rightharpoonup p$ in $L^{\frac{4}{3}}(0, T; L^2(\Omega))$.

344 Furthermore, (\mathbf{v}, θ, p) is a solution to (WF) parametrized by α .

345 *Proof.* Using (2.1)-(2.4) and (2.5)-(2.6), we prove that there exists \mathbf{u} and θ such
 346 that all the convergences above are verified in the same manner as in [10, Propo-
 347 sition VII.1.4]. Let us prove first that \mathbf{u} is a solution of (WF.3) parametrized
 348 by α and θ . With the same pattern of proof as in Theorem 2.4, one proves im-
 349 mediately that $(\mathbf{u}_\varepsilon \cdot \nabla) \mathbf{u}_\varepsilon \rightharpoonup (\mathbf{u} \cdot \nabla) \mathbf{u}$ in $L^1(0, T; (L^1(\Omega))^d)$, and $(\mathbf{u}_\varepsilon \cdot \mathbf{n}) \mathbf{u}^{\text{ref}} \rightharpoonup$
 350 $(\mathbf{u} \cdot \mathbf{n}) \mathbf{u}^{\text{ref}}$ in $L^4(0, T; (L^{\frac{4}{3}}(\Gamma))^d)$. Regarding the penalization term:

$$351 \quad \begin{aligned} \|h(\alpha_\varepsilon) \mathbf{u}_\varepsilon - h(\alpha) \mathbf{u}\|_{L^2(0, T; L^2(\Omega)^d)}^2 &\lesssim \|h\|_\infty^2 \|\mathbf{u}_\varepsilon - \mathbf{u}\|_{L^2(0, T; L^2(\Omega)^d)}^2 \\ &\quad + \int_0^T \int_\Omega (h(\alpha_\varepsilon) - h(\alpha))^2 |\mathbf{u}|^2. \end{aligned}$$

352 Since $\alpha_\varepsilon \rightarrow \alpha$ strongly in $L^1(\Omega)$, $h(\alpha_\varepsilon) \rightarrow h(\alpha)$ pointwise in Ω up to a subsequence
 353 (which is not relabeled). Lebesgue dominated convergence theorem then implies:
 354 $h(\alpha_\varepsilon) \mathbf{u}_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} h(\alpha) \mathbf{u}$ in $L^2(0, T; (L^2(\Omega))^d)$.

355 Concerning the boundary terms, we only consider the term with the approxima-
 356 tion of the neg function. First, we claim that there exists γ such that $\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})(\mathbf{u}_\varepsilon +$
 357 $\mathbf{u}^{\text{ref}}) \rightharpoonup \gamma$ in $L^{\frac{4}{3}}(0, T; L^{\frac{4}{3}}(\Gamma)^d)$. Notice that, for ε large enough and using the proper-
 358 ties of the neg approximation, we have:

$$(2.7) \quad \begin{aligned} \int_0^T \|\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})(\mathbf{u}_\varepsilon + \mathbf{u}^{\text{ref}})\|_{L^{\frac{4}{3}}(\Gamma)}^{\frac{4}{3}} &\lesssim \int_0^T \left(\|\mathbf{u}_\varepsilon \cdot \mathbf{n} + 1\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} \|\mathbf{u}_\varepsilon + \mathbf{u}^{\text{ref}}\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} \right) \\ &\lesssim \int_0^T \left(\|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} + C \right) \|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} \\ &\quad + \int_0^T \left(\|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} + C \right) \|\mathbf{u}^{\text{ref}}\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{4}{3}} \\ &\lesssim \int_0^T \|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{8}{3}} + 2 \left(\int_0^T \|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{8}{3}} \right)^{\frac{1}{2}} + C. \end{aligned}$$

360 In addition, from Proposition 1.2, we have $\|\mathbf{u}_\varepsilon\|_{L^{\frac{8}{3}}(\Gamma)}^{\frac{8}{3}} \lesssim \|\mathbf{u}_\varepsilon\|_{L^2(\Omega)}^{\frac{8}{3}} \|\mathbf{u}_\varepsilon\|_{H^1(\Omega)}$. Since
 361 \mathbf{u}_ε is bounded in $L^\infty(0, T; (L^2(\Omega))^d)$ and in $L^2(0, T; (H^1(\Omega))^d)$ as proved in Propo-
 362 sition 2.2, we see that $\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})(\mathbf{u}_\varepsilon + \mathbf{u}^{\text{ref}})$ is bounded in $L^{\frac{4}{3}}(0, T; L^{\frac{4}{3}}(\Gamma)^d)$ in-
 363 dependently of ε . Since this Banach space is reflexive, it proves the claimed weak
 364 convergence. Let us now prove that γ can be identified with $(\mathbf{u} \cdot \mathbf{n})^-(\mathbf{u} + \mathbf{u}^{\text{ref}})$.

365 First, since $\mathbf{u}_\varepsilon \rightarrow \mathbf{u}$ strongly in $L^2(0, T; L^2(\Gamma)^d)$ and $\text{neg}_\varepsilon(\cdot) \rightarrow (\cdot)^-$ uniformly,
 366 one proves that $\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) \rightarrow (\mathbf{u} \cdot \mathbf{n})^-$ strongly in $L^2(0, T; L^2(\Gamma))$. Then, the
 367 weak convergence of \mathbf{u}_ε in $L^4(0, T; L^2(\Gamma)^d)$ and [10, Proposition II.2.12] implies that
 368 $\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})(\mathbf{u}_\varepsilon + \mathbf{u}^{\text{ref}}) \rightharpoonup (\mathbf{u} \cdot \mathbf{n})^-(\mathbf{u} + \mathbf{u}^{\text{ref}})$ weakly in $L^{\frac{4}{3}}(0, T; L^1(\Gamma)^d)$. Using [10,
 369 Proposition II.2.9], we argue that $\gamma = (\mathbf{u} \cdot \mathbf{n})^-(\mathbf{u} + \mathbf{u}^{\text{ref}})$.

370 Regarding p_ε , we use an inf-sup condition as the one introduced in the proof of
 371 [28, Theorem 5.1, eq. (5.14)], which states that

$$372 \quad (2.8) \quad \|p\|_{L^2(\Omega)} \lesssim \sup_{\Psi \in V} \frac{\int_{\Omega} p \nabla \cdot \Psi}{\|\Psi\|_{H^1(\Omega)}}.$$

373 Therefore, using (WFe.3), one shows that:

$$374 \quad \|p_\varepsilon\|_{L^2(\Omega)} \lesssim \|\partial_t \mathbf{u}_\varepsilon\|_{V'} + \|\mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon)\|_{V'} + \|\mathcal{A}\mathbf{u}_\varepsilon\|_{V'} + \|h(\alpha)\mathbf{u}_\varepsilon\|_{V'} + \|\mathcal{T}\theta\|_{V'} \\ + \|\mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}})\|_{V'} + \|f\|_{V'} + \|\sigma^{\text{ref}}\|_{V'}.$$

375 We now bound each term depending on ε :

- 376 • Since the Stokes operator is continuous, $\|\mathcal{A}\mathbf{u}_\varepsilon\|_{V'} \lesssim \|\mathbf{u}_\varepsilon\|_{H^1(\Omega)}$ and therefore,
 377 $\mathcal{A}\mathbf{u}_\varepsilon$ is bounded in $L^2(0, T; V')$.
- 378 • Using [10, Eq. (V.3)], we prove that $\|\mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon)\|_{V'} \lesssim \|\mathbf{u}_\varepsilon\|_{L^2(\Omega)}^{2-\frac{d}{2}} \|\mathbf{u}_\varepsilon\|_{H^1(\Omega)}^{\frac{d}{2}}$,
 379 which in turn shows that $\mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon)$ is bounded in $L^{\frac{4}{d}}(0, T; V')$.
- 380 • Obviously, $\|h(\alpha)\mathbf{u}_\varepsilon\|_{V'} \leq \|h\|_\infty \|\mathbf{u}_\varepsilon\|_{L^2(\Omega)}$ and therefore, $h(\alpha)\mathbf{u}_\varepsilon$ is bounded
 381 in $L^\infty(0, T; V')$.

382 We are left with the boundary term. Let $0 \neq \Psi \in V$. In a similar manner as before
 383 and using Proposition 1.2, there exists a constant $C > 0$ such that:

$$384 \quad \frac{1}{\|\Psi\|_{H^1(\Omega)}} \int_{\Gamma_{\text{out}}} |\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})(\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \Psi| \lesssim \left(\|\mathbf{u}_\varepsilon\|_{L^2(\Omega)}^{\frac{3-d}{4}} \|\mathbf{u}_\varepsilon\|_{H^1(\Omega)}^{\frac{d-1}{4}} + C \right)^2.$$

385 As proved before, \mathbf{u}_ε is bounded in $L^\infty(0, T; L^2(\Omega)) \cap L^2(0, T; H^1(\Omega))$. Therefore, \mathbf{u}_ε
 386 is also bounded in $L^{2-\frac{2}{d}}(0, T; H^1(\Omega))$. Taking the supremum over Ψ , this proves that
 387 $\mathcal{N}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon)$ is bounded in $L^{\frac{4}{d}}(0, T; V')$. Finally, in a similar fashion as in [10, Lemma
 388 V.1.6], the above bounds prove that $\partial_t \mathbf{u}_\varepsilon$ is bounded in $L^{\frac{4}{d}}(0, T; V')$. These bounds
 389 prove that (p_ε) is bounded in $L^{\frac{4}{d}}(0, T; L^2(\Omega))$, and therefore (p_ε) weakly converges to
 390 some p in $L^{\frac{4}{d}}(0, T; L^2(\Omega))$.

391 Concerning θ , the convergence is largely proved in the same way as in Theorem 2.4.
 392 The only difference concerns the convergence of $\text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})\theta_\varepsilon$ to $(\mathbf{u} \cdot \mathbf{n})^-\theta$, which
 393 is proved in the same manner as (2.7). All these convergence results let us say that
 394 (\mathbf{u}, θ, p) is a solution to (WF) in the distribution sense. \square

395 **2.2. Further results in dimension 2.** It is notably known that the solution
 396 of the Navier-Stokes equations with homogeneous Dirichlet boundary conditions are
 397 unique in dimension 2. We prove here that uniqueness still holds with the boundary
 398 conditions (1.2). We only sketch the proof.

399 **PROPOSITION 2.6.** *Let $d = 2$. Then, the solution $(\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ of (WFe) is unique.*

400 *Sketch of proof* First of all, note that uniqueness of $(\mathbf{u}_\varepsilon, \theta_\varepsilon)$ implies the uniqueness of
 401 p_ε via the De Rham Theorem [10, Theorem IV.2.4 and Chapter V].

402 Let $(\mathbf{u}_{\varepsilon_1}, \theta_{\varepsilon_1})$ and $(\mathbf{u}_{\varepsilon_2}, \theta_{\varepsilon_2})$ be two solutions of (WF.1)-(WF.3). Define $\mathbf{u} = \mathbf{v} =$
 403 $\mathbf{v}_{\varepsilon_1} - \mathbf{v}_{\varepsilon_2}$ and $\theta = \theta_{\varepsilon_1} - \theta_{\varepsilon_2}$. Slightly adapting the proof in [10, Section VII.1.2.5],

404 one proves that:

$$405 \quad (2.9) \quad \frac{d}{dt} |\mathbf{v}|_{L^2(\Omega)}^2 + A |\nabla \mathbf{v}|_{L^2(\Omega)}^2 \lesssim g^v(t) |\mathbf{v}|_{L^2(\Omega)}^2 + B |\theta|_{L^2(\Omega)}^2 + \nu^v |\nabla \mathbf{v}|_{L^2(\Omega)}^2$$

406 where ν^v is a positive constant and g^v is a function in $L^1([0, T])$.

407 Testing the differential equation verified by θ with θ proves that:

$$408 \quad \begin{aligned} & \frac{d}{dt} |\theta|_{L^2(\Omega)}^2 + 2C \int_{\Omega} k |\nabla \theta|^2 + \int_{\Gamma_{\text{out}}} \theta^2 \left(\frac{1}{2} (\mathbf{u}_{\varepsilon_1} \cdot \mathbf{n}) + \beta \text{neg}_{\varepsilon} (\mathbf{u}_{\varepsilon_1} \cdot \mathbf{n}) \right) \\ & = - \int_{\Gamma_{\text{out}}} \left(\beta (\text{neg}_{\varepsilon} (\mathbf{u}_{\varepsilon_1} \cdot \mathbf{n}) - \text{neg}_{\varepsilon} (\mathbf{u}_{\varepsilon_2} \cdot \mathbf{n})) + \frac{1}{2} (\mathbf{u} \cdot \mathbf{n}) \right) \theta_{\varepsilon_2} \theta. \end{aligned}$$

409 With a similar proof as the one of [Proposition 2.2](#), we can prove that, on Γ_{out} ,
410 $\theta^2 \left(\frac{1}{2} (\mathbf{u}_1 \cdot \mathbf{n}) + \beta \text{neg}_{\varepsilon} (\mathbf{u}_1 \cdot \mathbf{n}) \right) \geq 0$. Therefore, using [\(A3\)](#), one has:

$$411 \quad (2.10) \quad \frac{d}{dt} |\theta|_{L^2(\Omega)}^2 + 2C \int_{\Omega} k |\nabla \theta|^2 \lesssim \left(|\beta|_{L^\infty(\Gamma_{\text{out}})} + \frac{1}{2} \right) |\mathbf{u} \cdot \mathbf{n}|_{L^3(\Gamma_{\text{out}})} |\theta_{\varepsilon_2}|_{L^3(\Gamma_{\text{out}})} |\theta|_{L^3(\Gamma_{\text{out}})}.$$

412 Using Sobolev embeddings and Young inequality, we prove:

$$413 \quad \begin{aligned} & \left(|\beta|_{L^\infty(\Gamma_{\text{out}})} + \frac{1}{2} \right) |\mathbf{u} \cdot \mathbf{n}|_{L^3(\Gamma_{\text{out}})} |\theta_{\varepsilon_2}|_{L^3(\Gamma_{\text{out}})} |\theta|_{L^3(\Gamma_{\text{out}})} \\ & \lesssim \left(|\beta|_{L^\infty(\Gamma_{\text{out}})} + \frac{1}{2} \right)^3 \frac{|\theta_{\varepsilon_2}|_{L^2(\Omega)} |\nabla \theta_{\varepsilon_2}|_{L^2(\Omega)}^2}{2(\nu^\theta)^3} (|\mathbf{u}|_{L^2(\Omega)}^2 + |\theta|_{L^2(\Omega)}^2) \\ & \quad + \frac{(\nu^\theta)^{\frac{3}{2}}}{2} \left(|\nabla \mathbf{u}|_{L^2(\Omega)}^2 + |\nabla \theta|_{L^2(\Omega)}^2 \right), \end{aligned}$$

414 where ν^θ is a positive constant. Therefore, summing [\(2.9\)](#) and [\(2.10\)](#) gives $\frac{d}{dt} (|\mathbf{u}|_{L^2(\Omega)}^2 + |\theta|_{L^2(\Omega)}^2) \lesssim \max(g_1^v, g^\theta) (|\mathbf{u}|_{L^2(\Omega)}^2 + |\theta|_{L^2(\Omega)}^2)$, with g_1^v and g^θ integrable. Therefore,
415 applying Gronwall's lemma and noticing that $|\mathbf{u}(0)|_{L^2(\Omega)}^2 + |\theta(0)|_{L^2(\Omega)}^2 = 0$, one shows
416 that $\mathbf{u} = 0$ and $\theta = 0$. \square

417 Note that we may also prove that, for $d = 2$, the solution (\mathbf{u}, θ, p) of [\(WF\)](#) is
418 unique. We can also state stronger convergence (compared to the ones stated in
419 [Theorem 2.5](#)) in dimension 2. These results will be useful in the analysis of the
420 optimisation problems.
421

422 **COROLLARY 2.7.** *Suppose $d = 2$. Under the assumptions of [Theorem 2.5](#), $\mathbf{u}_\varepsilon \rightarrow$
423 \mathbf{u} strongly in $L^\infty(0, T; L^2(\Omega)^2)$, $\nabla \mathbf{u}_\varepsilon \rightarrow \nabla \mathbf{u}$ strongly in $L^2(0, T; L^2(\Omega)^2)$, $\theta_\varepsilon \rightarrow \theta$
424 strongly in $L^\infty(0, T; L^2(\Omega))$, $\nabla \theta_\varepsilon \rightarrow \nabla \theta$ strongly in $L^2(0, T; L^2(\Omega))$ and $p_\varepsilon \rightarrow p$
425 strongly in $L^2(0, T; L^2(\Omega))$.*

426 *Proof.* Denote $\bar{\mathbf{u}} = \mathbf{u} - \mathbf{u}_\varepsilon$, $\bar{\theta} = \theta - \theta_\varepsilon$ and $\bar{p} = p_\varepsilon - p$. The variational formulation
427 verified by $(\bar{\mathbf{u}}, \bar{\theta}, \bar{p})$ reads as: for all $\Psi \in V^u$:

$$428 \quad (2.11a) \quad \begin{aligned} & -\langle \mathcal{P}\bar{p}, \Psi \rangle = \langle \partial_t \bar{\mathbf{u}} + \mathcal{A}\bar{\mathbf{u}} + h(\alpha) \bar{\mathbf{u}}, \Psi \rangle_{(V^u)', V^u} + \langle (h(\alpha) - h(\alpha_\varepsilon)) \mathbf{u}_\varepsilon, \Psi \rangle_{(V^u)', V^u} + \\ & \quad \langle \mathcal{B}(\mathbf{u}, \mathbf{u}) - \mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon), \Psi \rangle_{(V^u)', V^u} + \langle \mathcal{T}\bar{\theta}, \Psi \rangle_{(V^u)', V^u} \\ & \quad \frac{1}{2} \langle \mathcal{N}(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}}) - \mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}), \Psi \rangle_{(V^u)', V^u}, \end{aligned}$$

429

$$430 \quad (2.11b) \quad 0 = \langle \nabla \cdot \bar{\mathbf{u}}, q \rangle_{L^2(\Omega)},$$

431

$$\begin{aligned}
0 &= \langle \partial_t \bar{\theta}, \varphi \rangle_{(V^\theta)', V^\theta} - \langle \mathcal{D}(\mathbf{u}, \bar{\theta}) + \mathcal{D}(\bar{\mathbf{u}}, \theta_\varepsilon), \varphi \rangle_{(V^\theta)', V^\theta} \\
&+ \langle (\mathcal{C}(\alpha) - \mathcal{C}(\alpha_\varepsilon))\theta + \mathcal{C}(\alpha_\varepsilon)\bar{\theta}, \varphi \rangle_{(V^\theta)', V^\theta} \\
&+ \langle \mathcal{M}(\mathbf{u}, \theta) + \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon, \theta_\varepsilon), \varphi \rangle_{(V^\theta)', V^\theta}.
\end{aligned}
\tag{2.11c}$$

433 The following inequalities, valid for $d = 2$, will be useful throughout this proof:

434 • As proved in [10, Eq. (V.5)]:

$$\begin{aligned}
&(2.12) \\
&\frac{\langle \mathcal{B}(\mathbf{u}, \mathbf{u}) - \mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon), \Psi \rangle_{(V^u)', V^u}}{\|\Psi\|_{H^1(\Omega)}} = \frac{\langle \mathcal{B}(\bar{\mathbf{u}}, \mathbf{u}) + \mathcal{B}(\mathbf{u}_\varepsilon, \bar{\mathbf{u}}), \Psi \rangle_{(V^u)', V^u}}{\|\Psi\|_{H^1(\Omega)}} \\
&\lesssim (\|\bar{\mathbf{u}}\|_{L^2(\Omega)} \|\mathbf{u}\|_{H^1(\Omega)} + \|\mathbf{u}_\varepsilon\|_{L^2(\Omega)} \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}).
\end{aligned}
\tag{2.12}$$

436 • Concerning the boundary term in (2.11a):

$$\begin{aligned}
&\langle \mathcal{N}(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}}) - \mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}), \Psi \rangle_{(V^u)', V^u} = \\
&\langle \mathcal{N}(\mathbf{u}, \bar{\mathbf{u}}), \Psi \rangle_{(V^u)', V^u} + \int_{\Gamma_{\text{out}}} (\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}))(\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \Psi.
\end{aligned}
\tag{2.11a}$$

438 We now deal with each term separately. Concerning the first term, Young's
439 inequality and Proposition 1.1 imply:

$$\langle \mathcal{N}(\mathbf{u}, \bar{\mathbf{u}}), \Psi \rangle_{(V^u)', V^u} \lesssim \|\bar{\mathbf{u}}\|_{H^1(\Omega)} \|\Psi\|_{H^1(\Omega)}.
\tag{2.13}$$

441 Owing to the Lipschitz behavior of the neg function, and the uniform con-
442 vergence of neg_ε toward neg (see (A3)), there exists $C_\varepsilon > 0$ such that:

$$\begin{aligned}
&\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) = \text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}(\mathbf{u}_\varepsilon \cdot \mathbf{n}) \\
&+ \text{neg}(\mathbf{u}_\varepsilon \cdot \mathbf{n}) - \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) \\
&\leq |\bar{\mathbf{u}} \cdot \mathbf{n}| + C_\varepsilon
\end{aligned}
\tag{2.14}$$

444 where $C_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} 0$. Therefore, using Proposition 1.2, we infer:

$$\begin{aligned}
&\int_{\Gamma_{\text{out}}} (\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}))(\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \Psi \\
&\lesssim \|\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})\|_{L^4(\Gamma)} \|\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}\|_{L^2(\Gamma)} \|\Psi\|_{L^4(\Gamma)} \\
&\lesssim \left(\|\bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{3}{4}} + C_\varepsilon \right) \|\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}\|_{L^2(\Omega)}^{\frac{1}{2}} \|\nabla(\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}})\|_{L^2(\Omega)}^{\frac{1}{2}} \\
&\quad \times \|\Psi\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \Psi\|_{L^2(\Omega)}^{\frac{3}{4}}.
\end{aligned}
\tag{2.15}$$

446 • The inequality proved in Proposition 1.1 shows that:

$$\int_{\Omega} \theta_\varepsilon \bar{\mathbf{u}} \cdot \nabla \bar{\theta} \lesssim \|\theta_\varepsilon\|_{H^1(\Omega)} \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{1}{2}} \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{1}{2}} \|\nabla \bar{\theta}\|_{L^2(\Omega)}.
\tag{2.16}$$

448 • One will need also to bound the terms involving \mathbf{u} and θ on the boundary.
449 Using once again Proposition 1.1, one shows directly that:

$$\int_{\Gamma_{\text{out}}} (\bar{\mathbf{u}} \cdot \mathbf{n}) \theta_\varepsilon \bar{\theta} \lesssim \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{3}{4}} \|\theta_\varepsilon\|_{L^2(\Omega)}^{\frac{1}{2}} \|\nabla \theta_\varepsilon\|_{L^2(\Omega)}^{\frac{1}{2}} \|\bar{\theta}\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \bar{\theta}\|_{L^2(\Omega)}^{\frac{3}{4}}.
\tag{2.17}$$

451 With the same technique as for (2.15), one proves:

(2.18)

$$452 \int_{\Gamma_{\text{out}}} (\text{neg}(\bar{\mathbf{u}} \cdot \mathbf{n}) - \text{neg}_{\varepsilon}(\mathbf{u}_{\varepsilon} \cdot \mathbf{n})) \theta_{\varepsilon} \bar{\theta} \lesssim \left(\|\bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{3}{4}} + C_{\varepsilon} \right) \|\theta_{\varepsilon}\|_{L^2(\Omega)}^{\frac{1}{2}} \\ \|\nabla \theta_{\varepsilon}\|_{L^2(\Omega)}^{\frac{1}{2}} \|\bar{\theta}\|_{L^2(\Omega)}^{\frac{1}{4}} \|\nabla \bar{\theta}\|_{L^2(\Omega)}^{\frac{3}{4}}.$$

453 Since $d = 2$, one has $\bar{\mathbf{u}}' \in L^2(0, T; (V^u)')$ and we may choose $\Psi = \bar{\mathbf{u}}(t)$ for fixed t
454 in (2.11a). Using the fact that $\nabla \cdot \Psi = 0$ in this case, and after rearranging the terms,
455 we obtain:

$$\begin{aligned} & \frac{d}{dt} \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + 2A \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + 2 \int_{\Omega} h(\alpha) |\bar{\mathbf{u}}|^2 + \int_{\Gamma_{\text{out}}} \text{pos}(\mathbf{u} \cdot \mathbf{n}) |\bar{\mathbf{u}}|^2 = \\ & - 2 \langle (h(\alpha) - h(\alpha_{\varepsilon})) \mathbf{u}_{\varepsilon}, \bar{\mathbf{u}} \rangle_{(V^u)', V^u} - \langle \mathcal{B}(\bar{\mathbf{u}}, \mathbf{u}), \bar{\mathbf{u}} \rangle_{(V^u)', V^u} + \langle \mathcal{B}(\bar{\mathbf{u}}, \bar{\mathbf{u}}), \mathbf{u} \rangle_{(V^u)', V^u} \\ 456 & - \int_{\Omega} B \bar{\theta} e_y \cdot \bar{\mathbf{u}} - \int_{\Gamma_{\text{out}}} (\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_{\varepsilon}(\mathbf{u}_{\varepsilon} \cdot \mathbf{n})) (\mathbf{u}_{\varepsilon} - \mathbf{u}^{\text{ref}}) \cdot \bar{\mathbf{u}} \\ & - \int_{\Gamma_{\text{out}}} (\bar{\mathbf{u}} \cdot \mathbf{n}) \mathbf{u}_{\varepsilon} \cdot \bar{\mathbf{u}} = 0. \end{aligned}$$

457 Therefore, (2.12), (2.15), Proposition 1.2 and Young's inequality imply there ex-
458 ists $C_1 > 0$ independent of ε such that:

$$459 \frac{d}{dt} \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + C_1 \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \lesssim \|\bar{\theta}\|_{L^2(\Omega)}^2 + 2 \int_{\Omega} |h(\alpha) - h(\alpha_{\varepsilon})|^2 |\mathbf{u}_{\varepsilon}|^2 \\ + g_1^u \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + (g_2^u)^{\frac{4}{5}} \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^{\frac{2}{5}},$$

460 where $g_1^u = 1 + \|\mathbf{u}\|_{H^1(\Omega)}^2 + \|\mathbf{u}\|_{L^2(\Omega)}^2 \|\nabla \mathbf{u}\|_{L^2(\Omega)}^2 + 2 \|\mathbf{u}_{\varepsilon}\|_{L^2(\Omega)} \|\nabla \mathbf{u}_{\varepsilon}\|_{L^2(\Omega)}^2 +$
461 $\|\mathbf{u}^{\text{ref}}\|_{L^2(\Omega)} \|\mathbf{u}^{\text{ref}}\|_{H^1(\Omega)}^2$ and $g_2^u = C_{\varepsilon}^2 \|\mathbf{u}_{\varepsilon} - \mathbf{u}^{\text{ref}}\|_{L^2(\Omega)} \|\nabla(\mathbf{u}_{\varepsilon} - \mathbf{u}^{\text{ref}})\|_{L^2(\Omega)}$.

462 Using once again Young's inequality, one has:

$$463 \frac{d}{dt} \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + C_1 \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \lesssim \|\bar{\theta}\|_{L^2(\Omega)}^2 + (1 + g_1^u) \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \\ + 2 \int_{\Omega} |h(\alpha) - h(\alpha_{\varepsilon})|^2 |\mathbf{u}_{\varepsilon}|^2 + g_2^u.$$

464 We now move back to (2.11c) and choose $\varphi = \bar{\theta}$, which gives, after some manip-
465 ulation:

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\bar{\theta}\|_{L^2(\Omega)}^2 + C \int_{\Omega} k(\alpha_{\varepsilon}) |\nabla \bar{\theta}|^2 + \int_{\Gamma_{\text{out}}} \left(\frac{1}{2} (\mathbf{u} \cdot \mathbf{n}) + \beta \text{neg}(\mathbf{u} \cdot \mathbf{n}) \right) \bar{\theta}^2 \\ 466 & = \int_{\Omega} \theta_{\varepsilon} \bar{\mathbf{u}} \cdot \nabla \bar{\theta} - C \int_{\Omega} (k(\alpha) - k(\alpha_{\varepsilon})) \nabla \theta \cdot \nabla \bar{\theta} \\ & - \int_{\Gamma_{\text{out}}} [((\bar{\mathbf{u}} \cdot \mathbf{n}) + \beta (\text{neg}(\mathbf{u} \cdot \mathbf{n}) - \text{neg}_{\varepsilon}(\mathbf{u}_{\varepsilon} \cdot \mathbf{n})))] \theta_{\varepsilon} \bar{\theta}. \end{aligned}$$

467 As shown in Proposition 2.2, $\int_{\Gamma_{\text{out}}} \left(\frac{1}{2} (\mathbf{u} \cdot \mathbf{n}) + \beta \text{neg}(\mathbf{u} \cdot \mathbf{n}) \right) \bar{\theta}^2$ is positive. There-
468 fore, using (2.17), (2.18), Proposition 1.2 and Young's inequality, one proves that
469 there exists $C_3 > 0, C_4 > 0$, such that:

$$470 \frac{d}{dt} \|\bar{\theta}\|_{L^2(\Omega)}^2 + C_3 \|\nabla \bar{\theta}\|_{L^2(\Omega)}^2 \lesssim \|\theta_{\varepsilon}\|_{L^2(\Omega)}^2 \|\nabla \theta_{\varepsilon}\|_{L^2(\Omega)}^2 \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + C_4 \|\nabla \bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \\ + \left(C \int_{\Omega} (k(\alpha) - k(\alpha_{\varepsilon}))^2 |\nabla \theta|^2 \right) + g_1^{\theta} \|\bar{\theta}\|_{L^2(\Omega)}^2 + g_2^{\theta},$$

471 where $g_1^\theta = 1 + \|\theta_\varepsilon\|_{L^2(\Omega)}^2 \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}^2$, $g_2^\theta = C_\varepsilon^2 \|\theta_\varepsilon\|_{L^2(\Omega)} \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}$.

472 Summing (2.19) and (2.20) and choosing C_4 small enough, there exists $C^* > 0$
 473 such that:

$$\begin{aligned}
 & \frac{d}{dt} (\|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + \|\bar{\theta}\|_{L^2(\Omega)}^2) + C^* (\|\nabla\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + \|\nabla\bar{\theta}\|_{L^2(\Omega)}^2) \lesssim g_2^u + g_2^\theta \\
 474 \quad (2.21) \quad & + (1 + \|\theta_\varepsilon\|_{L^2(\Omega)}^2 \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}^2 + g_1^u) \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + (g_1^\theta + 1) \|\bar{\theta}\|_{L^2(\Omega)}^2 \\
 & + \int_\Omega (k(\alpha) - k(\alpha_\varepsilon))^2 |\nabla\theta|^2 + \int_\Omega |h(\alpha) - h(\alpha_\varepsilon)|^2 |\mathbf{u}_\varepsilon|^2.
 \end{aligned}$$

475 We now introduce the following functions

$$\begin{aligned}
 476 \quad a_\varepsilon^u &= (1 + \|\theta_\varepsilon\|_{L^2(\Omega)}^2 \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}^2 + g_1^u), & b_\varepsilon^u &= \int_\Omega |h(\alpha) - h(\alpha_\varepsilon)|^2 |\mathbf{u}_\varepsilon|^2 + g_2^u, \\
 477 \quad a_\varepsilon^\theta &= (1 + g_1^\theta), & b_\varepsilon^\theta &= \int_\Omega (k(\alpha) - k(\alpha_\varepsilon))^2 |\nabla\theta|^2 + g_2^\theta.
 \end{aligned}$$

478 Since \mathbf{u} and \mathbf{u}_ε both belong to $L^2(0, T; H^1(\Omega)^2) \cap L^\infty(0, T; L^2(\Omega)^2)$ (the same holds
 479 for θ and θ_ε), a_ε^u , b_ε^u , a_ε^θ and b_ε^θ are integrable, and so are $a_\varepsilon = \max(a_\varepsilon^u, a_\varepsilon^\theta)$ and $b_\varepsilon =$
 480 $b_\varepsilon^u + b_\varepsilon^\theta$. Grönwall's lemma proves that for all $t \in [0, T]$, $\|\bar{\mathbf{u}}(t)\|_{L^2(\Omega)}^2 + \|\bar{\theta}(t)\|_{L^2(\Omega)}^2 \leq$
 481 $\left(\int_0^t b_\varepsilon(s) ds\right) \exp\left(\int_0^t a_\varepsilon(s) ds\right)$. Since $a_\varepsilon \geq 0$ and $b_\varepsilon \geq 0$, $t \mapsto \left(\int_0^t b_\varepsilon(s) ds\right)$ and $t \mapsto$
 482 $\exp\left(\int_0^t a_\varepsilon(s) ds\right)$ are non-decreasing and we have

$$483 \quad (2.22) \quad \sup_{t \in [0, T]} (\|\bar{\mathbf{u}}(t)\|_{L^2(\Omega)} + \|\bar{\theta}(t)\|_{L^2(\Omega)}) \leq \left(\int_0^T b_\varepsilon(s) ds\right)^{\frac{1}{2}} \exp\left(\frac{1}{2} \int_0^T a_\varepsilon(s) ds\right).$$

484 Since, on one hand, $\alpha_\varepsilon \rightarrow \alpha$ in $L^1(\Omega)$ and α_ε is independent of time, and on the other
 485 hand, $\mathbf{u}_\varepsilon \rightarrow \mathbf{u}$ strongly in $L^2(0, T; L^2(\Omega))$, Lebesgue's dominated convergence gives a
 486 subsequence (ε_k) such that:

$$487 \quad (2.23) \quad \int_0^T \int_\Omega |h(\alpha) - h(\alpha_{\varepsilon_k})|^2 |\mathbf{u}_\varepsilon|^2 \xrightarrow{k \rightarrow +\infty} 0, \quad \int_0^T \int_\Omega |k(\alpha) - k(\alpha_{\varepsilon_k})|^2 |\nabla\theta|^2 \xrightarrow{k \rightarrow +\infty} 0.$$

488 Notice that, owing to the convergence of \mathbf{u}_ε and θ_ε , $\|\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}\|_{L^2(\Omega)} \|\nabla(\mathbf{u}_\varepsilon -$
 489 $\mathbf{u}^{\text{ref}})\|_{L^2(\Omega)}$ and $\|\theta_\varepsilon\|_{L^2(\Omega)} \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}$ are bounded w.r.t ε in $L^1([0, T])$. Therefore,
 490 since $C_\varepsilon \xrightarrow{\varepsilon \rightarrow +\infty} 0$, it proves that $\int_0^T (g_2^u + g_2^\theta) \xrightarrow{\varepsilon_k \rightarrow +\infty} 0$. Gathering the previous
 491 convergence results then prove that $\int_0^T b_{\varepsilon_k}(s) ds \xrightarrow{k \rightarrow +\infty} 0$. In addition, thanks to the
 492 convergence proved in [Theorem 2.5](#), we show that $\int_0^T a_\varepsilon(s) ds$ is bounded w.r.t. ε .
 493 Therefore, it proves that $\|\mathbf{u} - \mathbf{u}_{\varepsilon_k}\|_{L^\infty(0, T; L^2(\Omega))} + \|\theta - \theta_{\varepsilon_k}\|_{L^\infty(0, T; L^2(\Omega))} \xrightarrow{k \rightarrow +\infty} 0$.

494 We now move back to (2.21). We integrate each side of the inequality:

$$\begin{aligned}
 495 \quad & \int_0^T \|\nabla\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 + \|\nabla\bar{\theta}\|_{L^2(\Omega)}^2 \lesssim F_\varepsilon^{u, \theta} + \int_0^T (g_1^u + \|\theta_\varepsilon\|_{L^2(\Omega)}^2 \|\nabla\theta_\varepsilon\|_{L^2(\Omega)}^2 + 1) \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \\
 & + \int_0^T (g_1^\theta + 1) \|\bar{\theta}\|_{L^2(\Omega)}^2,
 \end{aligned}$$

496 with $F_\varepsilon^{u, \theta} = \|\mathbf{u}_0(\alpha_\varepsilon) - \mathbf{u}_0(\alpha)\|_{L^2(\Omega)}^2 + \|\theta_0(\alpha_\varepsilon) - \theta_0(\alpha)\|_{L^2(\Omega)}^2 + \int_0^T (g_2^u + g_2^\theta)$
 497 $+ \int_0^T \int_\Omega |k(\alpha) - k(\alpha_\varepsilon)|^2 |\nabla\theta|^2 + \int_0^T \int_\Omega |h(\alpha) - h(\alpha_\varepsilon)|^2 |\mathbf{u}_\varepsilon|^2$. From [Assumptions 2.1](#),

498 the initial conditions are continuous with respect to α and thus the two first terms
499 above goes to 0 as $\varepsilon \rightarrow +\infty$. The third, fourth and fifth terms have been already
500 treated (see (2.23)). Concerning the two last terms, notice that due to the con-
501 vergence of \mathbf{u}_ε , one has $\|\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}\|_{L^2(\Omega)} \|\nabla(\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}})\|_{L^2(\Omega)}$ bounded w.r.t ε in
502 $L^1([0, T])$. The main problem concerns the term $\int_0^T \|\mathbf{u}_\varepsilon\|_{L^2(\Omega)} \|\nabla \mathbf{u}_\varepsilon\|_{L^2(\Omega)}^2 \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2$.
503 Since $\mathbf{u}_\varepsilon \in L^2(0, T; L^2(\Omega)^2)$, we only need to deal with $\|\nabla \mathbf{u}_\varepsilon\|_{L^2(\Omega)}^2 \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2$. How-
504 ever, as proved in Theorem 2.5, up to a subsequence, $\nabla \mathbf{u}_\varepsilon$ weakly converges to $\nabla \mathbf{u}$
505 in $L^2(0, T, L^2(\Omega))$ and $\|\mathbf{u} - \mathbf{u}_\varepsilon\|_{L^2(\Omega)}^2 \rightarrow 0$ in $L^\infty([0, T])$. Concerning the other terms
506 in g_1^u , they are all independent of ε , and we mainly use the fact that $\|\bar{\mathbf{u}}\|_{L^2(\Omega)} \rightarrow 0$
507 in $L^2([0, T])$. We may do the same proof concerning $\int_0^T \|\theta_\varepsilon\|_{L^2(\Omega)} \|\nabla \theta_\varepsilon\|_{L^2(\Omega)}^2 \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2$
508 and $\int_0^T \|\theta_\varepsilon\|_{L^2(\Omega)} \|\nabla \theta_\varepsilon\|_{L^2(\Omega)}^2 \|\bar{\theta}\|_{L^2(\Omega)}^2$.
509 Therefore, $\int_0^T (1 + \|\theta_{\varepsilon_k}\|_{L^2(\Omega)} \|\nabla \theta_{\varepsilon_k}\|_{L^2(\Omega)}^2 + g_1) \|\bar{\mathbf{u}}\|_{L^2(\Omega)}^2 \xrightarrow{\varepsilon_k \rightarrow +\infty} 0$ and $\int_0^T (g_1^\theta +$
510 $1) \|\bar{\theta}\|_{L^2(\Omega)}^2 \xrightarrow{\varepsilon_k \rightarrow +\infty} 0$. It eventually proves that $\|\nabla(\mathbf{u} - \mathbf{u}_{\varepsilon_k})\|_{L^2(0, T; L^2(\Omega))} + \|\nabla(\theta -$
511 $\theta_{\varepsilon_k})\|_{L^2(0, T; L^2(\Omega))} \xrightarrow{k \rightarrow +\infty} 0$.

512 Concerning the pressure, we use once again the inf-sup condition (2.8) on \bar{p} , which
513 proves that:

$$\begin{aligned}
& \|\bar{p}\|_{L^2(\Omega)} \lesssim \|\partial_t \bar{\mathbf{u}}\|_{V'} + \|\bar{\mathbf{u}}\|_{L^2(\Omega)} + \|h\| \|\bar{\mathbf{u}}\|_{L^2(\Omega)} + \|(h(\alpha) - h(\alpha_\varepsilon))\mathbf{u}_\varepsilon\|_{V'} + \\
& \|\mathcal{B}(\mathbf{u}, \mathbf{u}) - \mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon)\|_{V'} + \|\mathcal{T}\bar{\theta}\|_{V'} + \\
& \frac{1}{2} \|\mathcal{N}(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}}) - \mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}})\|_{V'}.
\end{aligned}
\tag{2.24}$$

515 Most of the terms in the right hand side were proved to converge strongly to 0 in
516 $L^2([0, T])$.

- 517 • In the same pattern of proof as in (2.23), one proves easily that $(h(\alpha) -$
518 $h(\alpha_{\varepsilon_k}))\mathbf{u}_{\varepsilon_k}$ converges strongly to 0 in $L^2(0, T; V')$
- 519 • Using (2.12), one proves that $\mathcal{B}(\mathbf{u}, \mathbf{u}) - \mathcal{B}(\mathbf{u}_{\varepsilon_k}, \mathbf{u}_{\varepsilon_k})$ strongly converges to 0
520 in $L^2(0, T; V')$.
- 521 • Summing (2.13) and (2.15) prove that: $\|\mathcal{N}(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}}) - \mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}})\|_{V'}$
522 $\lesssim \|\bar{\mathbf{u}}\|_{H^1(\Omega)} + C_\varepsilon$. Since $\mathbf{u}_{\varepsilon_k} \rightarrow \mathbf{u}$ strongly in $L^2(0, T; H^1(\Omega))$ and $C_\varepsilon \rightarrow 0$,
523 this proves that $\mathcal{N}(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}}) - \mathcal{N}_\varepsilon(\mathbf{u}_{\varepsilon_k}, \mathbf{u}_{\varepsilon_k} - \mathbf{u}^{\text{ref}})$ converges strongly to 0
524 in $L^2(0, T; V')$.
- 525 • Finally, in a similar fashion as in [10, Lemma V.1.6], the above bounds prove
526 that $\partial_t \bar{\mathbf{u}} \rightarrow 0$ strongly in $L^2(0, T; V')$.

527 Therefore, all the convergence results above prove that, up to a subsequence, $p_\varepsilon \rightarrow p$
528 strongly in $L^2(0, T; L^2(\Omega))$. Owing to Urysohn's subsequence principle and the
529 uniqueness of the solution to (WF), we actually obtain that the whole sequence
530 $(\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ strongly converges toward (\mathbf{u}, θ, p) . \square

531 **3. Optimal control and necessary conditions.** We now begin the analysis of
532 the optimal control problems (OPT) and (OPTe). Let us detail first some assumptions
533 made on the objective functional:

- 534 ASSUMPTIONS 3.1. • For $d = 2$, \mathcal{J} is lower semi-continuous with respect to
535 the (weak-*, strong, strong, strong) topology of $\mathcal{U}_{\text{ad}} \times L^2(0, T; V^u) \times L^2(0, T;$
536 $V^\theta) \times L^2(0, T; L^2(\Omega))$.
537 • In dimension 3, \mathcal{J} is either lower semi-continuous with respect to the (weak-
538 *, strong, strong, weak) topology of $\mathcal{U}_{\text{ad}} \times L^2(0, T; H^u) \times L^2(0, T; H^\theta) \times$

539 $L^2(0, T; L^2(\Omega))$, or lower semi-continuous with respect to the (weak-*, weak,
540 weak, weak) topology of $\mathcal{U}_{ad} \times L^2(0, T; V^u) \times L^2(0, T; V^\theta) \times L^2(0, T; L^2(\Omega))$.

541 The existence of solutions to (OPTe) and (OPT) is rather classical and we refer
542 for instance to [20, 31, 33]. We state a first result that let us see that a solution of
543 (OPT) can be approximated by (OPTe).

544 **THEOREM 3.2.** *Assume Assumptions 3.1 is verified. Let $(\alpha_\varepsilon^*, \mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ be a glob-*
545 *ally optimal solution of (OPTe). Then $(\alpha_\varepsilon^*) \subset \mathcal{U}_{ad}$ is a bounded sequence. Further-*
546 *more, there exists $(\alpha^*, \mathbf{u}^*, \theta^*, p^*) \in \mathcal{U}_{ad} \times L^2(0, T; V^u) \times L^2(0, T; V^\theta) \times L^2(0, T; L^2(\Omega))$*
547 *such that a subsequence of $(\alpha_\varepsilon^*, \mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ converges to $(\alpha^*, \mathbf{u}^*, \theta^*, p^*)$ in the topol-*
548 *ogy of Assumptions 3.1, and for all $(\alpha, \mathbf{u}, \theta, p)$ in $\mathcal{U}_{ad} \times L^2(0, T; V^u) \times L^2(0, T; V^\theta) \times$*
549 *$L^2(0, T; L^2(\Omega))$: $\mathcal{J}(\alpha^*, \mathbf{u}^*, \theta^*, p^*) \leq \mathcal{J}(\alpha, \mathbf{u}, \theta, p)$. Hence, any accumulation point of*
550 *$(\alpha_\varepsilon^*, \mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon)$ is a globally optimal solution of (OPT).*

551 *Proof.* The proof can be adapted from [20, Theorem 15] or [31, Theorem 3]. \square

552 However, the fact that this only concerns global solutions may appear restrictive.
553 Under an additional assumption, we can state a slightly stronger result.

554 **COROLLARY 3.3.** *Assume Assumptions 3.1 hold. Let α^* be a local strict solution*
555 *of (OPT), meaning that there exists $\rho > 0$ such that $\mathcal{J}(\alpha^*, \mathbf{u}^*, \theta^*, p^*) < \mathcal{J}(\alpha, \mathbf{u}, \theta, p)$*
556 *for all α such that $\|\alpha^* - \alpha\|_{BV} < \rho$. Then, there exists a family of local solution (α_ε^*)*
557 *of (OPTe) such that (α_ε^*) converges weak-* to α^* .*

558 *Proof.* Similar to [35, Theorem 3.14]. \square

559 **3.1. First order necessary conditions for (OPTe).** From now on, we set
560 $d = 2$, in order to have uniqueness of solution of (WFe). We make the following
561 assumption on the cost function:

562 **ASSUMPTIONS 3.4.** *Assume $d = 2$ and \mathcal{J} is Fréchet-differentiable.*

563 We define the sets $W^u(0, T) = \{\mathbf{u} \in L^2(0, T; (V^u)); \partial_t \mathbf{u} \in L^2(0, T; (V^u)')\}$, and
564 $W^\theta(0, T) = \{\theta \in L^2(0, T; (V^\theta)); \partial_t \theta \in L^2(0, T; (V^\theta)')\}$. Write, in $(V^u)' \times (V^\theta)' \times$
565 $L^2(0, T; L^2(\Omega))$, the equation (WFe) as $e(\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon, \alpha_\varepsilon) = 0$, where $e : W^u(0, T) \times$
566 $W^\theta(0, T) \times L^2(0, T; L^2(\Omega)) \times \mathcal{U}_{ad} \rightarrow L^2(0, T; (V^u)') \times L^2(0, T; L^2(\Omega)) \times L^2(0, T; (V^\theta)')$
567 $\times H^u \times H^\theta$ is defined as:

$$568 \quad e(\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon, \alpha_\varepsilon) = \begin{pmatrix} \partial_t \mathbf{u}_\varepsilon + \mathcal{A}\mathbf{u}_\varepsilon + \mathcal{B}(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon) + h(\alpha_\varepsilon)\mathbf{u}_\varepsilon + \mathcal{P}p \\ \quad \quad \quad + \frac{1}{2}\mathcal{N}_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) - f - \sigma^{\text{ref}} \\ \mathcal{P}^*\mathbf{u}_\varepsilon \\ \partial_t \theta_\varepsilon - \mathcal{D}(\mathbf{u}_\varepsilon, \theta_\varepsilon) + \mathcal{C}(\alpha_\varepsilon)\theta_\varepsilon + \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon, \theta_\varepsilon) - \phi \\ \quad \quad \quad \mathbf{u}_\varepsilon(0, \cdot) - \mathbf{u}_0(\alpha_\varepsilon) \\ \quad \quad \quad \theta_\varepsilon(0, \cdot) - \theta_0(\alpha_\varepsilon) \end{pmatrix}.$$

569 The operators \mathcal{N}_ε and \mathcal{M}_ε are Fréchet differentiable with the same smoothness as
570 the approximation neg_ε . Their derivatives with respect to \mathbf{u}_ε are denoted by $d_u \mathcal{N}_\varepsilon :$
571 $W^u(0, T)^2 \rightarrow \mathcal{L}(W(0, T), L^2(0, T; (V^u)'))$, $d_u \mathcal{M}_\varepsilon : W^u(0, T) \times W^\theta(0, T) \rightarrow$
572 $\mathcal{L}(W^u(0, T), L^2(0, T; (V^\theta)'))$, defined by:

$$573 \quad d_u \mathcal{N}_\varepsilon(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}})\mathbf{v} = \mathcal{N}_\varepsilon(\mathbf{u}, \mathbf{v}) + \mathcal{N}'(\mathbf{u}, \mathbf{u} - \mathbf{u}^{\text{ref}})\mathbf{v},$$

574

$$575 \quad \langle d_u \mathcal{M}_\varepsilon(\mathbf{u}, \theta)\mathbf{v}, \varphi \rangle_{(V^\theta)', V^\theta} = \int_{\Gamma_{\text{out}}} (1 + \beta \text{neg}'_\varepsilon(\mathbf{u} \cdot \mathbf{n})) (\mathbf{v} \cdot \mathbf{n}) \theta \varphi,$$

609 where $(e'_{\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon}(\alpha_\varepsilon))^*$ denotes the adjoint operator of $e'_{\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon}(\alpha_\varepsilon)$.

610 After some calculations, equation (3.2) is equivalent to solve, for all $(\mathbf{v}, \ell, q) \in$
 611 $W^u(0, T) \times W^\theta(0, T) \times L^2(0, T; L^2(\Omega))$, the following variational problem:

$$\begin{aligned}
 & \langle -\partial_t \lambda_\varepsilon^{\mathbf{u}} + \mathcal{A} \lambda_\varepsilon^{\mathbf{u}} + (\nabla \mathbf{u}_\varepsilon)^\top \lambda_\varepsilon^{\mathbf{u}} + \mathcal{B}(\mathbf{u}_\varepsilon, \lambda_\varepsilon^{\mathbf{u}}) + h(\alpha_\varepsilon) \lambda_\varepsilon^{\mathbf{u}} + \mathcal{P} \lambda_\varepsilon^p - \mathcal{D}_1(\theta_\varepsilon) \lambda_\varepsilon^\theta \\
 & \quad + \mathcal{N}'_\varepsilon(\mathbf{u}_\varepsilon, \lambda_\varepsilon^{\mathbf{u}}) + (\mathcal{N}'_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}))^* \lambda_\varepsilon^{\mathbf{u}} \\
 & \quad + \mathcal{N}'_2(\mathbf{u}_\varepsilon, \lambda_\varepsilon^{\mathbf{u}}) + (d_u \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon, \theta_\varepsilon))^* \lambda_\varepsilon^\theta, \mathbf{v} \rangle_{W(0, T)', W(0, T)} \\
 & \quad + \langle \mathbf{v}(0, \cdot), \lambda_\varepsilon^{u_0} \rangle_H \\
 612 \quad (3.3) \quad & = \langle g_1, \mathbf{v} \rangle_{W^u(0, T)', W^u(0, T)}, \\
 & \langle \mathcal{P}^* \lambda_\varepsilon^{\mathbf{u}}, q \rangle_{L^2(0, T; L^2(\Omega))} = \langle g_2, q \rangle_{L^2(0, T; L^2(\Omega))} \\
 & \langle -\partial_t \lambda_\varepsilon^\theta + \mathcal{T}^* \lambda_\varepsilon^{\mathbf{u}} + \mathcal{C}(\alpha_\varepsilon) \lambda_\varepsilon^\theta - \mathcal{D}_2(\mathbf{u}_\varepsilon) \lambda_\varepsilon^\theta + \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon)^* \lambda_\varepsilon^\theta, \ell \rangle_{W^\theta(0, T)', W^\theta(0, T)} \\
 & = \langle g_3, \ell \rangle_{W^\theta(0, T)', W^\theta(0, T)}
 \end{aligned}$$

613 where $\langle \mathcal{D}(\theta, \mathbf{u}), \varphi \rangle = \langle \mathcal{D}_1(\theta) \varphi, \mathbf{u} \rangle = \langle \mathcal{D}_2(\mathbf{u}) \varphi, \theta \rangle$, $\langle \mathcal{M}_\varepsilon(\mathbf{u}) \theta, \varphi \rangle = \langle \mathcal{M}_\varepsilon(\mathbf{u}) \varphi, \theta \rangle$
 614 $= \int_{\Gamma_{\text{out}}} ((\mathbf{u} \cdot \mathbf{n}) + \beta \text{neg}_\varepsilon(\mathbf{u} \cdot \mathbf{n})) \theta \varphi$, and $\langle \mathcal{N}'_2(\mathbf{u}_\varepsilon, \lambda_\varepsilon^{\mathbf{u}}), \mathbf{v} \rangle = \int_{\Gamma_{\text{out}}} (\mathbf{u}_\varepsilon \cdot \mathbf{n}) \lambda_\varepsilon^{\mathbf{u}} \cdot \mathbf{v}$. This
 615 equation, in turn, is the weak formulation of:

$$\begin{aligned}
 & -\partial_t \lambda_\varepsilon^{\mathbf{u}} - \Delta \lambda_\varepsilon^{\mathbf{u}} + h(\alpha_\varepsilon) \lambda_\varepsilon^{\mathbf{u}} + \nabla \lambda_\varepsilon^p + (\nabla \mathbf{u}_\varepsilon)^\top \lambda_\varepsilon^{\mathbf{u}} - (\mathbf{u}_\varepsilon \cdot \nabla) \lambda_\varepsilon^{\mathbf{u}} - \theta_\varepsilon \nabla \lambda_\varepsilon^\theta = g_1 \\
 & \nabla \cdot \lambda_\varepsilon^{\mathbf{u}} = g_2 \\
 & -\partial_t \lambda_\varepsilon^\theta + B \lambda_\varepsilon^{\mathbf{u}} \cdot e_y - \nabla \cdot (Ck(\alpha_\varepsilon) \nabla \lambda_\varepsilon^\theta) - \nabla \cdot (\mathbf{u}_\varepsilon \lambda_\varepsilon^\theta) = g_3 \\
 & \lambda_\varepsilon^{\mathbf{u}}|_{\Gamma_w \cup \Gamma_{\text{in}}} = 0, \\
 & \lambda_\varepsilon^\theta|_{\Gamma_{\text{in}}} = 0, \\
 616 \quad (3.4a) \quad & \partial_n \lambda_\varepsilon^\theta|_{\Gamma_w} = 0, \\
 & A \partial_n \lambda_\varepsilon^{\mathbf{u}} - \mathbf{n} \lambda_\varepsilon^p|_{\Gamma_{\text{out}}} = \left(\frac{1}{2} \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) + (\mathbf{u}_\varepsilon \cdot \mathbf{n}) \right) \lambda_\varepsilon^{\mathbf{u}} + (1 + \beta \mu_\varepsilon) \theta_\varepsilon \lambda_\varepsilon^\theta \mathbf{n} \\
 & \quad + \frac{1}{2} \mu_\varepsilon ((\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \lambda_\varepsilon^{\mathbf{u}}) \mathbf{n}, \\
 & Ck(\alpha_\varepsilon) \partial_n \lambda_\varepsilon^\theta + \beta \lambda_\varepsilon^\theta \text{neg}_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})|_{\Gamma_{\text{out}}} = 0 \\
 & \lambda_\varepsilon^{\mathbf{u}}(T) = 0, \lambda_\varepsilon^\theta(T) = 0,
 \end{aligned}$$

617

$$618 \quad (3.4b) \quad \mu_\varepsilon = \text{neg}'_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n})$$

619 and, as shown in a similar fashion in [32], $\lambda_\varepsilon^{u_0} = \lambda_\varepsilon^{\mathbf{u}}(0, \cdot)$, $\lambda_\varepsilon^{\theta_0} = \lambda_\varepsilon^\theta(0, \cdot)$.

620 An other consequence of Theorem 3.6 is that we can apply [33, Corollary 1.3]
 621 which states that at any local solution $(\alpha_\varepsilon^*, \mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*)$ of (OPTe), the following opti-
 622 mality conditions hold:

623 **THEOREM 3.7.** *Let α_ε^* be an optimal solution of (OPTe) with associated states*
 624 *$(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*)$. Then there exists adjoint states $(\lambda_\varepsilon^{\mathbf{u}}, \lambda_\varepsilon^\theta, \lambda_\varepsilon^p) \in L^2(0, T; V^u) \times L^2(0, T; V^\theta)$*
 625 *$\times L^2(0, T; L^2(\Omega))$ such that, denoting $(\lambda_\varepsilon^{u_0}, \lambda_\varepsilon^{\theta_0}) = (\lambda_\varepsilon^{\mathbf{u}}(0, \cdot), \lambda_\varepsilon^\theta(0, \cdot))$ and $\Lambda_\varepsilon = (\lambda_\varepsilon^{\mathbf{u}}, \lambda_\varepsilon^\theta,$
 626 $\lambda_\varepsilon^p, \lambda_\varepsilon^{u_0}, \lambda_\varepsilon^{\theta_0})$:*

$$\begin{aligned}
 & e(\alpha_\varepsilon^*, \mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*) = 0, \\
 & \mathcal{J}'_{\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*}(\alpha_\varepsilon^*) + (e_{\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*}(\alpha_\varepsilon^*))' \Lambda_\varepsilon = 0, \\
 627 \quad (3.5) \quad & \left\langle \mathcal{J}'_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*) + (e_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*))' \Lambda_\varepsilon, \alpha - \alpha_\varepsilon^* \right\rangle_{\mathcal{U}'_{ad}, \mathcal{U}_{ad}} \geq 0, \quad \forall \alpha \in \mathcal{U}_{ad}, \\
 & \alpha_\varepsilon \in \mathcal{U}_{ad}.
 \end{aligned}$$

628 REMARK 3.8. As stated in [33, Eq. (1.89)], since e and \mathcal{J} are Fréchet differ-
629 entiable, the mapping $\alpha_\varepsilon \mapsto \tilde{\mathcal{J}}(\alpha_\varepsilon) = \mathcal{J}(\alpha_\varepsilon, \mathbf{u}_\varepsilon, p_\varepsilon)$ is Fréchet differentiable, and
630 $\tilde{\mathcal{J}}'(\alpha_\varepsilon) = \mathcal{J}'_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*) + (e_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*))^* \Lambda_\varepsilon$, which reads as:

$$631 \quad (e_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*))^* \Lambda_\varepsilon = \int_0^T (h'(\alpha_\varepsilon) \mathbf{u}_\varepsilon \cdot \lambda_\varepsilon^{\mathbf{u}} + Ck'(\alpha_\varepsilon) \nabla \theta_\varepsilon \cdot \nabla \lambda_\varepsilon^\theta) \\ + \mathbf{u}'_0(\alpha_\varepsilon) \cdot \lambda_\varepsilon^{u_0} + \theta'_0(\alpha_\varepsilon) \lambda_\varepsilon^{\theta_0}.$$

632 **3.2. Limit adjoint system.** To conclude this paper, we will now study the con-
633 vergence of the adjoint states $(\lambda_\varepsilon^{\mathbf{u}}, \lambda_\varepsilon^\theta, \lambda_\varepsilon^p)$ to functions $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p)$. The only trouble
634 concerns the multiplier μ_ε defined in (3.4b). We will prove that at the limit, μ is
635 defined thanks to the convex-hull of the Heaviside function $H : \mathbb{R} \rightarrow [0, 1]$, defined as:

$$636 \quad (3.6) \quad H(u) = \begin{cases} \{0\} & \text{if } u < 0, \\ \{1\} & \text{if } u > 0, \\ [0, 1] & \text{if } u = 0. \end{cases}$$

637 As we will prove in this section, these limit adjoint states $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p)$ let us define
638 necessary conditions of optimality for the unrelaxed problem (OPT).

639 LEMMA 3.9. Let $(\alpha_\varepsilon) \subset \mathcal{U}_{ad}$ and $\alpha \in \mathcal{U}_{ad}$ such that $\alpha_\varepsilon \xrightarrow{*} \alpha$. Define by $(\lambda_\varepsilon^{\mathbf{u}}, \lambda_\varepsilon^\theta, \lambda_\varepsilon^p)$
640 a weak solution of (3.4) parametrized by α_ε . Then, there exists $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p) \in L^\infty(0, T;$
641 $H^u) \cap L^2(0, T; V^u) \times L^\infty(0, T, H^\theta) \cap L^2(0, T; V^\theta) \times L^\infty(0, T; L^2(\Omega))$ such that, up to
642 a subsequence:

- 643 • $\lambda_\varepsilon^{\mathbf{u}} \rightarrow \lambda^{\mathbf{u}}$ in $L^\infty(0, T; (L^2(\Omega))^2)$ and $\lambda_\varepsilon^\theta \rightarrow \lambda^\theta$ in $L^\infty(0, T; L^2(\Omega))$,
- 644 • $\lambda_\varepsilon^{\mathbf{u}} \xrightarrow{\varepsilon \rightarrow +\infty} \lambda^{\mathbf{u}}$ in $L^2(0, T; (H^1(\Omega))^2)$ and $\lambda_\varepsilon^\theta \xrightarrow{\varepsilon \rightarrow +\infty} \lambda^\theta$ in $L^2(0, T; (H^1(\Omega)))$,
- 645 • $\lambda_\varepsilon^{\mathbf{u}} \xrightarrow{\varepsilon \rightarrow +\infty} \lambda^{\mathbf{u}}$ in $L^2(0, T; (L^2(\Gamma))^2)$ and $\lambda_\varepsilon^\theta \xrightarrow{\varepsilon \rightarrow +\infty} \lambda^\theta$ in $L^2(0, T; (L^2(\Gamma)))$,
- 646 • $\lambda_\varepsilon^p \rightarrow \lambda^p$ in $L^2(0, T; L^2(\Omega))$.

647 Furthermore, there exists $\mu \in L^\infty(0, T; L^\infty(\Gamma_{out}))$ defined by $-\mu(x) \in H(-\mathbf{u}(x) \cdot$
648 $\mathbf{n}(x))$ a.e. in Γ_{out} such that $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p)$ is a weak solution to (3.4a) parametrized by
649 α and μ .

650 *Proof.* The proof is very similar to the ones presented in section 2.

- 651 • In a similar manner as for Proposition 2.2 and Proposition 2.3, one shows
652 that, for all $\sigma \in [0, \frac{1}{6}]$, there exist constants $c_\lambda^\theta(\sigma)$ and $c_\lambda^u(\sigma)$, independent of
653 ε , such that:

$$654 \quad \sup_{[0, T]} \|\lambda_\varepsilon^{\mathbf{u}}\|_{L^2(\Omega)} + \int_0^T \|\nabla \lambda_\varepsilon^{\mathbf{u}}\|_{L^2(\Omega)} + \int_{\mathbb{R}} |\tau|^{2\sigma} \left\| \mathcal{F} \left(\widetilde{\lambda_\varepsilon^{\mathbf{u}}} \right) \right\|_{L^2(\Omega)} d\tau \leq c_\lambda^u(\sigma),$$

655

$$656 \quad \sup_{[0, T]} \|\lambda_\varepsilon^\theta\|_{L^2(\Omega)} + \int_0^T \|\nabla \lambda_\varepsilon^\theta\|_{L^2(\Omega)} + \int_{\mathbb{R}} |\tau|^{2\sigma} \left\| \mathcal{F} \left(\widetilde{\lambda_\varepsilon^\theta} \right) \right\|_{L^2(\Omega)} d\tau \leq c_\lambda^\theta(\sigma).$$

- 657 • These bounds prove a weaker set of convergence in the same manner as in
658 Theorem 2.5. Since once again, we set $d = 2$, one proves the strong conver-
659 gence stated above as in Corollary 2.7.

660 We only need to prove that $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p)$ is a weak solution to (3.4a). The terms
661 $\langle (\mathcal{N}'_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}))^* \lambda_\varepsilon^{\mathbf{u}}, \mathbf{v} \rangle$ and $\langle (d_u \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon, \theta_\varepsilon))^* \lambda_\varepsilon^\theta, \mathbf{v} \rangle_{W^u(0, T)', W^u(0, T)}$ need a more thor-
662 ough examination. We start with the first term for which we have

$$663 \quad \langle (\mathcal{N}'_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}))^* \lambda_\varepsilon^{\mathbf{u}}, \mathbf{v} \rangle_{W^u} = \int_0^T \int_{\Gamma_{out}} \text{neg}'_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) ((\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \lambda_\varepsilon^{\mathbf{u}}) \mathbf{n} \cdot \mathbf{v}.$$

664 Thanks to assumptions **(A3)** and **(A4)** and $\mathbf{u}_\varepsilon \rightarrow \mathbf{u}$ pointwise a.e. in Γ (due to
665 strong convergence in $L^2([0, T] \times \Gamma)$), there is a subsequence (not relabeled) such that
666 $\text{neg}'_\varepsilon(\mathbf{u}_\varepsilon \cdot \mathbf{n}) \xrightarrow{*} \mu$ in $L^\infty(0, T; L^\infty(\Gamma_{\text{out}}))$, and such that $0 \leq \mu \leq 1$ a.e. in Γ_{out} and

667 $\mu = 1$ a.e. in $\{x \in \Gamma_{\text{out}}, \mathbf{u}(x) \cdot \mathbf{n}(x) > 0\}$, $\mu = 0$ a.e. in $\{x \in \Gamma_{\text{out}}, \mathbf{u}(x) \cdot \mathbf{n}(x) < 0\}$.

668 Furthermore, due to the convergence presented above, $((\mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}) \cdot \lambda_\varepsilon^{\mathbf{u}}) \rightarrow ((\mathbf{u} - \mathbf{u}^{\text{ref}}) \cdot \lambda^{\mathbf{u}})$
669 $\lambda^{\mathbf{u}}$ in $L^1(0, T; L^1(\Gamma_{\text{out}}))$. Therefore, it proves that:

$$670 \quad \langle (\mathcal{N}'_\varepsilon(\mathbf{u}_\varepsilon, \mathbf{u}_\varepsilon - \mathbf{u}^{\text{ref}}))^* \lambda_\varepsilon^{\mathbf{u}}, \mathbf{v} \rangle_{W^u(0, T)', W^u(0, T)} \rightarrow \int_0^T \int_{\Gamma_{\text{out}}} \mu ((\mathbf{u} - \mathbf{u}^{\text{ref}}) \cdot \lambda^{\mathbf{u}}) \mathbf{n} \cdot \mathbf{v}.$$

671 Similarly, one proves that:

$$672 \quad \langle (d_u \mathcal{M}_\varepsilon(\mathbf{u}_\varepsilon, \theta_\varepsilon))^* \lambda_\varepsilon^\theta, \mathbf{v} \rangle_{W^u(0, T)', W^u(0, T)} \rightarrow \int_0^T \int_{\Gamma_{\text{out}}} (1 + \beta\mu) (\mathbf{v} \cdot \mathbf{n}) \theta \lambda^\theta.$$

673 All other terms in **(3.3)** are easily proved to converge in the same manner as in
674 **Theorem 2.5**. Therefore, $(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p)$ is a weak solution to **(3.4a)** parametrized by α
675 and μ . \square

676 We may now prove the final result of this paper ; namely the necessary optimality
677 conditions of **(OPT)**.

678 **THEOREM 3.10.** *Let α^* be an optimal solution of **(OPT)** with associated state*
679 *$\mathbf{u}^*, \theta^*, p^*$. Then there exist a multiplier $\mu \in L^\infty(0, T; L^\infty(\Gamma_{\text{out}}))$ and adjoint states*
680 *$(\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p) \in L^2(0, T; V^u) \times L^2(0, T; V^\theta) \times L^2(0, T; L^2(\Omega))$ solution of **(3.4a)** such*
681 *that, denoting $(\lambda^{\mathbf{u}_0}, \lambda^{\theta_0}) = (\lambda^{\mathbf{u}}(0, \cdot), \lambda^\theta(0, \cdot))$ and $\Lambda = (\lambda^{\mathbf{u}}, \lambda^\theta, \lambda^p, \lambda^{\mathbf{u}_0}, \lambda^{\theta_0})$:*

$$682 \quad \langle \mathcal{J}'_{\alpha^*}(\mathbf{u}^*, \theta^*, p^*) + (e_{\alpha^*}(\mathbf{u}^*, \theta^*, p^*))^* \Lambda, \alpha - \alpha^* \rangle_{\mathcal{U}'_{ad}, \mathcal{U}_{ad}} \geq 0, \quad \forall \alpha \in \mathcal{U}_{ad}.$$

683 *Proof.* The proof follows the lines of [17, Theorem 4.4]. Denote by S_ε the solution
684 operator which to α associates the solution of the relaxed equations **(WFe)** and by S
685 the solution operator which to α associates the solution of **(WF)**. For some $\rho > 0$,
686 consider the auxiliary optimal control problem:

$$687 \quad (3.7) \quad \begin{aligned} \min F_\varepsilon(\alpha_\varepsilon) &= \mathcal{J}(\alpha_\varepsilon, \mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon) + \frac{1}{2} \|\alpha^* - \alpha_\varepsilon\|_{L^2(\Omega)}^2 \\ \text{s.t.} \quad &\begin{cases} (\mathbf{u}_\varepsilon, \theta_\varepsilon, p_\varepsilon) = S_\varepsilon(\alpha_\varepsilon), \\ \alpha_\varepsilon \in \mathcal{U}_{ad}, \\ \|\alpha_\varepsilon - \alpha^*\|_{L^2(\Omega)} \leq \rho. \end{cases} \end{aligned}$$

688 Since α_ε and α^* are both in \mathcal{U}_{ad} , they are both bounded in $L^\infty(\Omega)$ and therefore,
689 $\|\alpha^* - \alpha_\varepsilon\|_{L^2(\Omega)}$ is well defined. It is classical to show that **(3.7)** admits a global
690 minimizer $\alpha_\varepsilon^* \in \mathcal{U}_{ad}$.

691 Using **(2.22)** and **(2.24)** (but with $\alpha_\varepsilon \equiv \alpha$), one proves that (in the norm of the
692 topology from **Assumptions 3.1** with $d = 2$):

$$693 \quad (3.8) \quad \|S(\alpha) - S_\varepsilon(\alpha)\| \lesssim C_\varepsilon, \quad \forall \alpha \in \mathcal{U}_{ad},$$

694 where C_ε has been defined in **(2.14)**.

695 Note that due to the Fréchet-differentiability of \mathcal{J} supposed in **Assumptions 3.4**
696 and **(3.8)**, it holds, for ε large enough:

$$697 \quad |\mathcal{J}(\alpha, S(\alpha)) - \mathcal{J}(\alpha, S_\varepsilon(\alpha))| \lesssim C_\varepsilon, \quad \forall \alpha \in \mathcal{U}_{ad}, \quad \|\alpha - \alpha^*\| \leq \rho.$$

698 We obtain as a consequence that $F_\varepsilon(\alpha^*) \lesssim C_\varepsilon + \mathcal{J}(\alpha^*, S(\alpha^*))$, and:

$$699 \quad F_\varepsilon(\alpha) \gtrsim -C_\varepsilon + \mathcal{J}(\alpha^*, S(\alpha^*)) + \frac{1}{2} \|\alpha - \alpha^*\|_{L^2(\Omega)}^2, \quad \forall \alpha \in \mathcal{U}_{\text{ad}}, \quad \|\alpha - \alpha^*\|_{L^2(\Omega)} \leq \rho.$$

700 Therefore, for all $\alpha \in \mathcal{U}_{\text{ad}}$ such that $\|\alpha - \alpha^*\|_{L^2(\Omega)} \leq \rho$:

$$701 \quad F_\varepsilon(\alpha^*) \lesssim C_\varepsilon + \mathcal{J}(\alpha^*, S(\alpha^*)) \lesssim C_\varepsilon + \mathcal{J}(\alpha, S(\alpha)) \lesssim 2C_\varepsilon + F_\varepsilon(\alpha).$$

702 Hence, for some constant C' , and denoting $C'_\varepsilon = C' C_\varepsilon$, one has the implication:

$$703 \quad \forall \alpha \in \mathcal{U}_{\text{ad}}, \quad 2C'_\varepsilon < \frac{1}{2} \|\alpha - \alpha^*\|_{L^2(\Omega)}^2 \leq \frac{1}{2} \rho^2 \implies F_\varepsilon(\alpha^*) < F_\varepsilon(\alpha).$$

704 One has therefore the following necessary condition of optimality:

$$705 \quad (3.9) \quad \|\alpha_\varepsilon^* - \alpha^*\|_{L^2(\Omega)} \leq \sqrt{4C'_\varepsilon}.$$

706 Hence, for ε large enough, α_ε^* is in the ρ -ball around α^* ; therefore, α_ε^* is a local
707 solution of (OPTe). Using [Theorem 3.7](#), one then proves that there exists adjoint
708 states $(\lambda_\varepsilon^{\mathbf{u}}, \lambda_\varepsilon^\theta, \lambda_\varepsilon^p)$ solution of (3.4a) such that, for all $\alpha \in \mathcal{U}_{\text{ad}}$:

$$709 \quad (3.10) \quad \left\langle \mathcal{J}'_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*) + (e_{\alpha_\varepsilon^*}(\mathbf{u}_\varepsilon^*, \theta_\varepsilon^*, p_\varepsilon^*))' \Lambda_\varepsilon, \alpha - \alpha_\varepsilon^* \right\rangle_{\mathcal{U}'_{\text{ad}}, \mathcal{U}_{\text{ad}}} + \langle \alpha_\varepsilon^* - \alpha^*, \alpha - \alpha_\varepsilon^* \rangle_{L^2(\Omega)} \geq 0.$$

710 From (3.9), one has $\alpha_\varepsilon^* \rightarrow \alpha^*$ strongly in $L^2(\Omega)$, and therefore, in $L^1(\Omega)$. Since
711 $(\alpha_\varepsilon^* - \alpha^*)_\varepsilon \subset \mathcal{U}_{\text{ad}}$, one has also $(\alpha_\varepsilon^* - \alpha^*)_\varepsilon$ bounded in $BV(\Omega)$. Hence, $\alpha_\varepsilon^* \xrightarrow{*} \alpha^*$ in
712 \mathcal{U}_{ad} . Using then [Corollary 2.7](#), [Assumptions 3.1](#) and [Lemma 3.9](#), we can pass to the
713 limit in (3.10), which concludes this proof. \square

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