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# Rotational spectroscopy in the $v_3=v_6=1$ and $v_6=2$ vibrational states of $\text{CH}^{35}\text{Cl}_3$

Adina CEAUSU – VELCESCU\*, Petr PRACNA†, Roman A. MOTIYENKO, Laurent MARGULES§

*\*Université de Perpignan, Laboratoire de Mathématiques et Physique, 52 Avenue Paul Alduy, 66860 Perpignan Cedex, France*

*† University of Chemistry and Technology, Department of Analytical Chemistry, Technická 5, 166 28 Prague 6, Czech Republic*

*§ Univ. Lille, CNRS, UMR 8523 - PhLAM - Physique des Lasers Atomes et Molécules, F-59000 Lille, France*

Proofs to: Dr. Adina CEAUSU-VELCESCU

Université de Perpignan,  
Laboratoire de Mathématiques et Physique,  
52 Avenue Paul Alduy, 66860 Perpignan Cedex France  
Tel. (0033)4 68 66 22 19  
Fax: (0033)4 68 66 22 34  
E-mail address: [adina@univ-perp.fr](mailto:adina@univ-perp.fr)

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## Abstract

The overtone  $\nu_6=2$  ( $\sim 520$   $\text{cm}^{-1}$ ) and combination  $\nu_3=\nu_6=1$  ( $\sim 627$   $\text{cm}^{-1}$ ) vibrational levels of chloroform  $\text{CH}^{35}\text{Cl}_3$  have been studied by millimeter wave spectroscopy, in the regions 150-330 and 360-660 GHz, corresponding to  $J=22-49$  and 55-99. Like the  $\nu_6=1$  fundamental vibration, the  $\nu_3=\nu_6=1$  combination level has an accidentally small value of  $C-B-C\zeta \sim -0.0029$   $\text{cm}^{-1}$ . Hence, a global  $q_{22}$  resonance ( $\Delta k=\Delta l=\pm 2$ ) connects the  $(k-1, l_6=-1)$  and  $(k+1, l_6=+1)$  sublevels. As a consequence, several resonant  $q_{12}$  ( $\Delta k=\pm 1, \Delta l=\mp 2$ ) crossings occur, which enabled the precise determination of the  $q_{12}$  interaction parameter for this state. Altogether, 1182 rotational transitions in the  $\nu_3=\nu_6=1$  combination level have been measured and fitted, with a standard deviation of 41.2 kHz.

For the  $\nu_6=2$  ( $A_1+E$ ) overtone level, the combined effects of the  $q_{22}$  and  $q_{12}$   $l$ -type interactions induce resonant crossings between the  $(k, l_6=0)/(k+1, l_6=-2)$  pairs of levels, for a large variety of  $k$  values. Altogether, 1463 rotational transitions in the  $(l_6=0)$   $A_1$  sublevel and 2924 rotational transitions in the  $(l_6=\pm 2)$   $E$  sublevels, together with nine perturbation-allowed transitions induced by the  $q_{12}$  ( $\Delta k=\pm 1, \Delta l=\mp 2$ ) interaction, have been measured and fitted, with a global standard deviation of 41.4 kHz. This resulted in the first accurate determination of the molecular constants of this overtone level.

## 1. Introduction

$\text{CHCl}_3$  (chloroform, trichloromethane) is a clear, colorless, highly volatile, non-flammable liquid at room temperature. It is widely used as an extraction solvent for fats, oils, greases, rubber, resins, lacquers, artificial silk manufacture, gums and adhesives. Chloroform is also employed as an industrial solvent in the extraction and purification of some antibiotics, alkaloids, vitamins and flavors [1]. This explains why most of the oldest investigations of its vibrational spectrum were done in the liquid phase.

Chloroform may be hazardous to the environment, as it can be transported in air as vapor and dissolved in water. In the air, chloroform eventually breaks down to other toxic chemicals, such as phosgene and hydrogen chloride.

Gas-phase spectroscopy is an invaluable tool for monitoring atmospheric pollutants [2]. Unfortunately, with several low-lying excited states, small rotational constants, and four isotopologues simultaneously present, the gas-phase high-resolution infrared spectra of natural chloroform are almost impossible to interpret. The use of (expensive) isotopically pure samples seems thus to be the only possibility.

Rotational spectra of the symmetric species, in the ground vibrational state and the low-lying vibrational states  $\nu_2=1$  ( $676\text{ cm}^{-1}$ ),  $\nu_3=1$  ( $367\text{ cm}^{-1}$ ), and  $\nu_6=1$  ( $260\text{ cm}^{-1}$ ), were studied by Carpenter *et al.* [3]. The ground vibrational state was also studied using millimeter and submillimeter-wave spectroscopy by Cazzoli *et al.* [4], who obtained very accurate centrifugal distortion constants. For the main isotopologue  $\text{CH}^{35}\text{Cl}_3$ , these authors also resolved the splitting of the  $K=3$  transitions, for  $J \geq 45$ , and obtained thus a value of the  $h_3$  ( $\Delta k = \pm 6$ ) interaction parameter in the ground state. Colmont *et al.* [5] analyzed the ground state rotational spectra of various isotopic species of chloroform and obtained empirical experimental and *ab initio* structures. Later, Demaison [6] recalculated the *ab initio* equilibrium structure and derived also, as a check, a semi-experimental equilibrium structure.

Rovibrational studies were performed for all fundamental bands of  $\text{CH}^{35}\text{Cl}_3$  [7, 8, 9, 10, 11] including recently the  $\nu_6$  fundamental vibration, located near  $260\text{ cm}^{-1}$  [12]. The latter study, which merged rovibrational wavenumbers of the  $\nu_6$  band and rotational data from a previous work of Margulès *et al.* [13] in a global least-squares fit, comprises also a first accurate determination of the axial ground state constants  $C_0$ ,  $D_K^0$ , and  $H_K^0$ . Unfortunately, the infrared study of the  $2\nu_6$  lowest overtone levels was not possible so far due to very low intensities of both the overtone and hot bands reaching this level.

The aim of the present study is thus to provide a first accurate parameter set for the  $\nu_6=2$  overtone levels of  $\text{CH}^{35}\text{Cl}_3$ , near  $520\text{ cm}^{-1}$ , through the analysis of the rotational spectrum, which has been recorded on a natural sample. In fact, contrarily to the infrared vibration-rotation spectra, where the high congestion imposes the use of rather expensive isotopically pure samples, the rotational (millimeter-wave) spectrum of a natural sample may present several distinct regions where the contributions of the different isotopologues and vibrational levels of a given isotopologue occur.

The present study includes also a detailed analysis of the rotational spectrum of  $\text{CH}^{35}\text{Cl}_3$  in the  $\nu_3=\nu_6=1$  combination level, near  $627\text{ cm}^{-1}$ . This combination vibration has been already analyzed through a high-resolution study of the IR  $\nu_3+\nu_6-\nu_6$  hot band, in the  $\nu_3$  region, near  $367\text{ cm}^{-1}$  [9]. In this study, Pietilä *et al.* assigned 199 peaks of the above mentioned hot band, the wavenumbers of which were fitted using a fourth-order polynomial in  $m^l$ . However, the two components ( $l_6=-1$  and  $l_6=+1$ ) of this bi-parallel band were treated separately, the strong  $l(2,2)$  interaction, present both in the lower and the upper levels, being absorbed into effective  $B(\pm l)$  constants. The authors further remarked that the resonance was almost exact, since the difference  $B(-l)-B(+l)=237(3)\times 10^{-6}\text{ cm}^{-1}$  almost equals  $4|q_{22}|=243.87(4)\times 10^{-6}\text{ cm}^{-1}$ . Thus, the aim of the present study is to provide a more accurate treatment of this combination level with explicit inclusion of the  $l$ -type interactions, resulting in accurate molecular parameters, necessary to further studies of merged IR and rotational data.

## 2. Experimental details

The measurements in the frequency range under investigation (150-660 GHz) were performed on a chloroform sample with natural abundance, using the Lille spectrometer [14]. The absorption cell was a stainless-steel tube (6 cm diameter, 220 cm long). Measurements were conducted at room temperature and at a sample pressure of about 30 Pa. The frequency ranges 150-330 and 360-660 GHz were covered with various active and passive frequency multipliers from VDI Inc. and an Agilent synthesizer (12.5-18.25 GHz) was used as the source of radiation. In order to improve the signal sensitivity, the sources were frequency modulated at 30.5 kHz and lock-in detection with the second harmonic was used. Typical estimated uncertainties for measured line frequencies are 30 to 50 kHz, depending on the observed S/N ratio and the frequency range.

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<sup>1</sup> Where, as usual,  $m = -J$  for a  $P$  transition ( $\Delta J = -1$ ) and  $J+1$  for a  $R$  transition ( $\Delta J = +1$ ).

Despite the rather low pressure used in the experiment, we could not achieve Doppler-limited linewidths, for at least two reasons: (i) The hyperfine structure (HFS), which is still apparent for rotational lines with intermediate  $J$ /high  $K$  values. Many of the observed lines are thus partially or completely unresolved hyperfine structure components, even at Doppler linewidth. (ii) The increase of the frequency modulation depth during the experiment, which was necessary for improving the S/N ratio, for a better observation of transitions pertaining to the excited vibrational states. This increase resulted in additional broadening, giving for instance linewidths of the order of 1 MHz in the intermediate  $J$ -values region (compared to an estimated Doppler linewidth of the order of 0.3 MHz). Consequently, broader intermediate- $J$ /low- $K$  lines, with no apparent HFS, were included in the least-squares fit with higher uncertainties, whereas weaker, intermediate- $J$ /high- $K$  lines, showing partly resolved HFS, were excluded from the fit (as explained later).

### 3. The rotational spectrum

The degenerate fundamental vibration  $\nu_6=1$  ( $\text{CCl}_3$  asymmetric bending vibration), located at  $260.6 \text{ cm}^{-1}$ , is well separated from other vibrational levels and could be therefore treated as isolated [12, 13]. This is also the case of the  $\nu_6=2$  ( $A_1+E$  symmetry) and  $\nu_3=\nu_6=1$  (E) levels, which are separated by more than  $100 \text{ cm}^{-1}$ .

We started assignments with the rotational transitions in the  $\nu_3=\nu_6=1$  level. Like the  $\nu_6=1$  fundamental vibration, the  $\nu_3=\nu_6=1$  combination level is characterized by an accidentally small value of  $|C - B - C\zeta| \approx 88 \text{ MHz}$ . This combination of constants determines the separation of **unperturbed**  $(K-1, l_6=-1)/(K+1, l_6=+1)$  levels coupled by the  $q_{22}$   $l$ -type resonance [13], which is roughly given by  $\Delta E \approx 4(C - B - C\zeta)K \approx -0.012K$  (in  $\text{cm}^{-1}$  units). This separation increases only slowly with  $K$ , whereas the  $l(2,2)$  coupling matrix elements rapidly increase with  $J$ . This confers to the  $q_{22}$  interaction a globally resonant character, which is for vibrational levels with  $(C - B - C\zeta) \approx B$  'reserved' to the  $kl=+1$  sublevels only. The  $q_{22}$  resonance produces an effective contribution to the  $B$  rotational constant  $\Delta B \approx \pm 2|q_{22}|$ , with the upper sign corresponding to the higher level in the  $(K+1, l_6=+1) / (K-1, l_6=-1)$  pairs of levels in resonance. Thus, the  $K$ -clusters become wide separated, with the  $l_6=+1$  lines pushed down to lower frequencies and those of  $l_6=-1$  up to higher frequencies. For a given  $J$ , the  $l(2,2)$  resonance term decreases with  $k$ , whereas the unperturbed energies difference increases only slightly with  $k$ . Hence, the absolute magnitude of the energy shift decreases with  $k$ , for a given  $J$ . The, otherwise degenerate,  $|kl=+1, J+1\rangle \leftarrow$

$|kl=+1, J\rangle$  transitions are symmetrically split by about  $\pm 4|q_{22}|(J+1)$ , with the  $A$ -component at lower frequencies. For  $kl\neq+1$ , due to the negative sign of  $(C-B-C\zeta)$ , the  $l_6=+1$  lines are pushed down to lower frequencies, whereas the  $l_6=-1$  lines are pushed up to higher frequencies. The displacement decreases with  $k$ , for a given  $J$  cluster. The combined effects of  $D_{JK} < 0$  and of the  $q_{22}$  resonance make thus the lower  $l_6=+1$  cluster degraded towards higher frequencies, for all the values of  $J$ . The  $l_6=-1$  clusters are degraded towards lower frequencies for  $J \leq 35$ ; a bandhead is observed for moderate  $J$  values (Fig. 1), whereas, for higher  $J$  values, these clusters become also degraded towards higher frequencies (Fig. 2).

Assignment of the  $v_3=v_6=1$ ,  $l_6 = -1$  and  $+1$ , clusters of rotational lines, which much resemble those of  $v_6=1$ , was thus straightforward, despite their weakness as shown in Fig. 1. A total of 1182 lines of the  $l_6=\pm 1$  components have been assigned and fitted.

The pattern of rotational transitions in the  $v_6=2$  overtone level is quite different from that of the fundamental level  $v_6=1$ . The  $l(2, 2)$  resonance connects the triplets of levels  $(K+2, l_6=+2)/(K, l_6=0)/(K-2, l_6=-2)$ , the relative positions of which are mainly determined by the  $g_{66}$  anharmonicity parameter and the combination of the  $C\zeta$  and  $(C-B)$  parameters, as follows:  $E(K \pm 2, l_6 = \pm 2) - E(K, l_6 = 0) \approx \pm 4(C - B - C\zeta)(K \pm 1) + 4(g_{66} - C\zeta) = \mp 0.012(K \pm 1) + 0.550$  (in  $\text{cm}^{-1}$  units).

This puts the  $(K, l_6=0)$  levels below their  $(K-2, l_6=-2)$  resonance partners for all values of  $K$ . Also, the  $(K, l_6=0)$  levels lie below the  $(K+2, l_6=+2)$  ones for  $K < 45$ .

The appearance of the  $K$ -clusters of rotational lines, for a given  $J$  value, in the  $v_6=2$  level, can be easily understood again by considering the case of exact resonance, which concerns the  $|0^0\rangle$  and  $|\pm 2^{\pm 2}\rangle$  levels, the latter being degenerate. The Wang transformation yields an  $A_+$  block of interacting levels, symmetrically displaced with respect to the mean value of their unperturbed energies, and a  $|2^2, A_- \rangle$  unperturbed level. Hence, the  $(2^2, A_-, J+1) \leftarrow (2^2, A_-, J)$  transition is unperturbed. Three clusters can thus be observed, for each value of  $J$  (Fig. 1): a lower-frequency one, which corresponds to  $l=0$ , a middle one, composed of  $l=+2$  transitions, which 'follows' the  $(2^2, A_-, J+1) \leftarrow (2^2, A_-, J)$  transition, and a higher-frequency one, composed of  $l=-2$  transitions and 'following' the  $(2^2, A_+, J+1) \leftarrow (2^2, A_+, J)$  transition. This latter cluster contains also the  $(1^2)$  and  $(0^2)$  transitions. The effect of the  $l(2,2)$  resonance on the appearance and behavior of the  $l=0, +2$ , and  $-2$   $J$ -clusters of lines is by the way nicely depicted in Fig. 5 of Ref. [15], for the  $v_{10}=2$  rotational spectrum of  $\text{CF}_3\text{CCH}$ .

The lower frequency clusters, corresponding to  $l_6=0$  and composed of regularly spaced lines, degraded to higher frequencies, were clearly observed in the spectrum, distinct from the much stronger rotational transitions of the vibrational ground state and the fundamental  $v_6=1$  level, not being obscured by them (Figs. 1 and 2). Assignment of the  $l_6=0$  transitions at moderate  $J$ -values (Fig. 1), facilitated by the intensity alternation due to nuclear spin statistics, was the starting point of the analysis, yielding more reasonable values of the  $B^{(0)}$  and  $B^{(2)}$  rotational constants, which were used for further assignments. At this stage, the AABS package, designed by Kisiel *et al.* for broadband rotational spectroscopy [16], has been extensively used. Finally, the middle and the upper clusters, corresponding to  $l_6=+2$  and  $l_6=-2$ , respectively, were also assigned. Like before, intensity alternation due to nuclear spin statistics and the Loomis-Wood representation of the spectrum in the AABS package were extremely helpful in progressive extension to higher  $J$  values. A total of 1463 lines of the  $l_6=0$  component and 2924 lines of the  $l_6=\pm 2$  components ( $22 \leq J'' \leq 49$  and  $55 \leq J'' \leq 99$ ) have been thus assigned and fitted.

With the energies of the  $l_6=0$  levels being pushed down by the  $q_{22}$  resonance, these fall into local resonances with both  $l_6=+2$  and  $l_6=-2$  levels due to other vibration-rotation interactions, such as  $q_{12}$  ( $\Delta k=\pm 1, \Delta l=\mp 2$ ), but also higher-order  $f_{42}$  ( $\Delta k=\pm 4, \Delta l=\mp 2$ ) and  $f_{52}$  ( $\Delta k=\pm 5, \Delta l=\pm 2$ ) interactions. In addition to this, resonant crossings between the  $l_6=+2$  and  $l_6=-2$  levels due to the  $f_{14}$  interaction occur for several low  $K$  values, as shown in Figure 3 with the reduced energy plot of the whole system of rotational levels of  $v_6=2$ . Nine perturbation-allowed transitions, induced by the  $q_{12}$  interaction, were also included in the fit.

#### 4. Results and discussion

For fitting the experimental frequencies, a conventional Hamiltonian of a symmetric-top molecule has been used, the matrix elements of which are given in Appendix A. For the least-squares fitting, using the SIMFIT program [17], the data were given statistical weights proportional to the inverse square of their estimated uncertainties.

The present work is based on a huge set of rotational data. Part of the observed transitions would be further split by the hyperfine quadrupole interactions due to the three chlorine nuclei with nuclear spin  $3/2$ . The detailed description of the tensorial coupling techniques allowing to construct the nuclear quadrupole part of the Hamiltonian, for any number of nuclei of spin  $\geq 1$ , can be found for instance in Refs. [18, 19]. This tensorial coupling mechanism is implemented in the SPFIT/SPCAT program package [20], which was



already used for studying the HFS of the  $J=1\leftarrow 0$  and  $J=2\leftarrow 1$  transitions in the ground vibrational state of  $\text{CHCl}_3$  [21]. This structure is quite complex, extending over more than 50 MHz and consisting of a great number of measurable components. The analysis performed by the authors of Ref. [21] allowed fitting all determinable components of the inertial nuclear splitting tensors for the chlorine nuclei. On the contrary, there are only few studies dealing with HFS of rotational transitions in degenerate excited states of symmetric top molecules [3, 22]. For the  $\nu_6=1$  state of  $\text{CHCl}_3$ , Carpenter *et al.* [3] constrained the  $eQq_{ZZ}$  and  $\eta_Q$  parameters to the corresponding ground state values.

In our current analysis we could in principle attempt either at assigning the resolved HFS and fit them explicitly with the SPFIT program, or reducing the resolved HFS to ‘hyperfine-free’ rotational frequencies, using e.g. the HFS parameters constrained to their ground vibrational state values, and fit them with the use of SIMFIT program [17]. Our moderate  $J$ -values ( $22 \leq J'' \leq 35$ ) spectrum clearly shows hyperfine quartets in the vibrational ground,  $\nu_3=1$ , and  $\nu_6=1$  states, for high  $K$  values, as, for a given  $J$ , the quadrupole splitting of lines increases with  $K$ . At the same time, intensities are decreasing with increasing  $K$ ; this makes the quartets of the, much weaker,  $\nu_6=2$  band, more difficult to distinguish, as they are often overlapped by other, stronger, transitions. For such  $J$ -values, HFS structure is already partly collapsing, even though still visible. This means, however, that the reliable assignments of the hyperfine components for using them in either of the mentioned fitting approaches, remains somewhat problematic. Therefore we decided to exclude from the data set the lines corresponding to these intermediate  $J$  / high  $K$ -values, showing partly resolved hyperfine structure, and those only broadened due to hyperfine interaction give lower weights (i.e. higher uncertainties of 100-200 kHz) and fit them with the SIMFIT program.

Despite this reduction of the data set, frequent overlaps with much stronger transitions in the lower lying levels and of course vanishing intensities for growing  $J$  values, we were still able to collect over 1000 transitions in the two  $l$ -sublevels of the  $\nu_3=\nu_6=1$  vibrational state and over 4000 transitions in the three  $l$ -sublevels of the  $\nu_6=2$  vibrational state, which we believe to be a sound basis for determination of parameters of the centrifugal distortion expansion and all relevant off-diagonal interaction terms of the effective rotational Hamiltonian of a symmetric top molecule.

#### 4.1. The $v_3=v_6=1$ level

In the  $v_3=v_6=1$  vibrational level, situation is almost the same as that described for  $v_6=1$  [13]: the  $q_{22}$  resonance pushes the pairs of interacting levels  $(k, l_6=-1)/(k+2, l_6=+1)$  apart so that they are brought into local resonances due to the  $q_{12}$   $l$ -type interaction. We were able to detect and assign transitions on both sides of the crossings for the pairs of levels  $1^{+1}/2^{-1}$  (at  $J=43/44$ ) and  $3^{+1}/4^{-1}$  (at  $J=67/68$ ) only: the  $2^{+1}/3^{-1}$  crossing, occurring at  $J=56/57$ , is in a spectral range not recorded by our spectrometer, whereas the crossings  $4^{+1}/5^{-1}$ ,  $5^{+1}/6^{-1}$ , and so on, occur at  $J$  values which are too high for detecting the corresponding transitions.

The fits were performed with two reduced forms of the effective rotational Hamiltonian, the  $Q$  and  $D$  reductions [23]. With 14 refined parameters (out of 22), these reduction schemes provided an identical reproduction of the data, with a standard deviation of 41.2 kHz (Table 1). We have to notice that the  $C$ ,  $D_K$ , and  $\eta_K$  coefficients of the  $K$ -dependent part of the rotational energies were kept fixed, in order to avoid too high correlation between these parameters and  $C\zeta$ , the latter parameter being refined.

The unitary equivalence checking of the  $Q$  and  $D$  reductions, up to 3<sup>rd</sup> order is shown in Table 2. The results of this Table can be compared to those presented for the  $v_6=1$  level, given in Table 3 of Ref. [12]. Surprisingly, the discrepancies noticed for  $v_6=1$ , for instance that between  $-2q_{12}^D/F$  and  $\frac{1}{2}d_t^Q/q_{22}$ , seem to be less important for the  $v_3=v_6=1$  level: for the latter level, the above mentioned quantities agree within  $2\sigma$ . The 3<sup>rd</sup> order combinations also show quite good agreement, except the  $\Delta^{Q,D}\eta_J$  difference, which is statistically not defined.

#### 4.2. The $v_6=2$ level

For each value of the rotational quantum number  $J$ , rotational levels are grouped into three-level-systems with the same value of  $G=|k-l|$ . Several kinds of level crossings take place, but only two of them result into resonances due either to  $\Delta(k-l) = 3$  or to  $\Delta(k-l) = 6$  interactions, which have strong, local effects on energy levels with a large variety of  $G$  values.

First, the  $q_{22}$  resonance pushes the pairs of levels  $(k, l_6=0)/(k-2, l_6=-2)$  apart, so that local resonances due to the  $q_{12}$  interaction ( $\Delta G=\pm 3$ ) occur (Fig. 3). We have observed such crossings between  $(k, l_6=0)/(k+1, l_6=-2)$  levels for a large variety of  $k$  values and were thus able to refine not only the main interaction parameter,  $q_{12}$ , but also the parameter of the  $l$ -dependent term,  $q_{12}^l$ , which is of the same order of magnitude as  $q_{12}$  [24].

Second, the  $(k, l_6=0)/(k+4, l_6=-2)$  levels are interacting through a  $f_{42}$   $l$ -type interaction ( $\Delta k=\pm 4, \Delta l=\mp 2$ ). Level crossings occur for  $J=60$  in the  $0^0/4^{-2}$  pair, 71 in the  $1^0/5^{-2}$  pair, 81 in the  $2^0/6^{-2}$  pair and so on. Unfortunately, only the first crossing occurs at a  $J$ -value which is low enough for observing transitions with a reasonable intensity.

Another interesting phenomenon, due to the strong  $l(2,2)$  resonance and the resulting mixing of states, occurs: for each  $G$  value,  $G \geq 45$ , the two lower levels in the three-level-system (i.e.  $(k, l_6=0)$  and  $(k+2, l_6=+2)$ ) change their relative positions, though they do not cross in a conventional sense. This phenomenon has been already described for the  $v_{10}=2$  state of 1,1,1-trifluoropropane by Wötzel *et al.* [25]. Thus, for a certain value of  $J$ , the mixing of the  $|k, l_6 = 0\rangle/|k + 2, l_6 = +2\rangle$  wavefunctions is so strong that an unambiguous labeling of the levels is not possible. For higher values of  $J$ , the labeling becomes again unambiguous, but with reversed ordering of the levels.

The approach employed in the present work for fitting the frequencies of the assigned transitions is that proposed by Sarka and Harder [24], who developed, for an isolated  $v_t=2$  level of a  $C_{3v}$  symmetric-top molecule, five reduced forms of the effective Hamiltonian. The common feature of these reduced forms is that from the four lower order  $\Delta(k-l)=3$  interactions, two have to be constrained and the other two refined, but the simultaneous constraint of the  $k$ - and  $l$ -dependent terms of the  $(2,-1)$  interaction is not allowed. Therefore, the following constraints have been suggested up to fifth order [24]: all parameters of the  $K$ -dependent terms of off-diagonal interactions ( $f_{22}^K, f_{12}^K, f_{12}^{lK}, f_{22}^{lK}, \varepsilon_K, d_t^K$ ) and the parameter  $f_{14}$  of the  $(4,1)$  interaction should be fixed to zero. In the present case, all these constraints have been applied except for  $f_{22}^K$  (see below). Indeed, it has been shown by Watson *et al.* [23] that the constraint  $\tilde{f}_{22}^K = 0$  should be avoided if  $|C - B - C\zeta|$  is accidentally too small. This is the case for the  $2\nu_6$  vibration of  $\text{CHCl}_3$  (and situation was of course the same for  $\nu_6$  [12]), where  $C - B - C\zeta \approx -92 \text{ MHz}$ . For accidentally small values of  $(C - B - C\zeta)$ , one can alternatively use one of the following constraints:  $\tilde{\tau}_J = 0$ , or  $\tilde{\tau}_{JK} = 0$ , or  $\tilde{\tau}_K = 0$ .

The reduction schemes employed in the present work correspond to the  $Q$ ,  $D$ , and  $L$  reductions described in Ref. [24]. As pointed out in this Reference, the parameters  $\tilde{f}_{14}, \tilde{f}_{12}^{lJ}, \tilde{f}_{12}^{lK}, \tilde{f}_{22}^{lK}$  depend on the transformation parameters  $o, p, r$  and therefore three of them should be constrained. The constraints suggested in the above mentioned Reference are thus  $\tilde{f}_{14} = 0$ ;  $\tilde{f}_{12}^{lK} = 0$ ;  $\tilde{f}_{22}^{lK} = 0$ , the  $\tilde{f}_{12}^{lJ}$  parameter being refined. In the present case, the  $f_{24}$  interaction

parameter ( $\Delta k = \pm 2$ ,  $\Delta l = \mp 4$ ) was not refined in the final stage of the calculations but kept fixed to a value obtained in a preliminary stage (see footnote b of Table 3).

With the final choice of 28 refined parameters, the data could be reproduced within error limits comparable to the estimated experimental accuracy, of 50 kHz. The global standard deviation is thus of 41.4 kHz for the three reductions  $D$ ,  $Q$ , and  $L$  (Table 3).

The unitary equivalence of the three parameter sets, up to third order, was also checked and the results are presented in Table 4. For this purpose, the following relations, derived for the lowest order parameters of the  $\Delta(k-l)=3$  interactions, have been employed [26]

$$-\frac{2q_{12}^D}{F} = \frac{d_t^Q}{2q_{22}} \quad \frac{d_t^L}{2q_{22}} = \frac{2q_{12}^{L,D}}{G}$$

$$q_{12}^{L,D} = q_{12}^{L,Q} + \frac{Gd_t^Q}{4q_{22}} \quad q_{12}^L = q_{12}^D + \frac{Fq_{12}^{L,D}}{G}$$

As shown in Table 4, the above equalities are satisfactorily, though not perfectly, fulfilled. Good fulfillment of the unitary equivalence of parameters occurs also for the third-order constants  $\eta_J$  and  $\eta_K$ , for which the following equalities can be established [26]:

$$\Delta^{Q,D}\eta_J = -\Delta^{Q,D}\eta_K = -\frac{8(q_{12}^D)^2}{F} = -\frac{(d_t^Q)^2 F}{2q_{22}^2}$$

$$\Delta^{Q,L}\eta_J = -\Delta^{Q,L}\eta_K = -\frac{8(q_{12}^L)^2}{F} = -\frac{8(q_{12}^{L,Q})^2 F}{G^2}$$

$$\Delta^{L,D}\eta_J = -\Delta^{L,D}\eta_K = \frac{8[(q_{12}^L)^2 - (q_{12}^D)^2]}{F}$$

In these relations, the symbol  $\Delta^{X,Y}$  implies the difference between the values of a parameter in the  $X$  and  $Y$  reductions ( $X, Y = Q, D$ , or  $L$ ). By the way, the refinement of the  $f_{22}^{LK}$  parameter has a noticeable effect on the third-order parameters  $\eta_J$  and  $\eta_K$ , thus improving the fulfillment of the above mentioned equalities.

We have not shown the unitary equivalence relations for the fourth-order parameters  $f_{22}^J$  and  $f_{42}$ , as suggested for instance in Ref. [26], because the corresponding differences are only poorly defined. It would be however interesting to check, even only partly, the fourth order contributions to the rotational constants  $B$  and  $C$ . The dependence of these rotational constants on the transformation parameter  $s$  writes ([26], corrected in [27]):

$$B_{v,l}(s) = B_{v,l} + 4 \left[ \frac{1}{2}(v^2 - l^2) + v \right] \left( q_{12}s + \frac{1}{4}Fs^2 \right) - 4(v^2 + 2v - 3l^2) \left( q_{12}s - \frac{1}{4}Gs^2 \right)$$

$$C_{v,l}(s) = C_{v,l} - 8 \left[ \frac{1}{2}(v^2 - l^2) + v \right] \left( q_{12}s + \frac{1}{4}Fs^2 \right)$$

From these relations, one can derive for a  $v_6=2$  state the following equalities:

$$\Delta^{Q,D}B^{(0)} = -16 \left[ \frac{(q_{12}^D)^2}{F} + 2 \frac{(q_{12}^{l,D})^2 - (q_{12}^{l,Q})^2}{G} \right]$$

$$\Delta^{Q,D}B^{(2)} = -8 \left[ \frac{(q_{12}^D)^2}{F} - 2 \frac{(q_{12}^{l,D})^2 - (q_{12}^{l,Q})^2}{G} \right]$$

as well as

$$\frac{1}{2} \Delta^{Q,D}C^{(0)} = \Delta^{Q,D}C^{(2)} = 16 \frac{(q_{12}^D)^2}{F}$$

For the  $C$  constants, checking of the above equalities is meaningless, the uncertainties being larger than the experimental differences. For the  $\Delta B$  values however, good fulfillment is observed, even though uncertainties of the experimental differences are of the same order of magnitude as the differences themselves. We obtain for instance for the  $l_6=0$  sublevel,  $(\Delta^{Q,D}B^{(0)})_{exp} = 5.27(188) \times 10^{-9} \text{ cm}^{-1}$  and  $(\Delta^{Q,D}B^{(0)})_{th} = 4.701(62) \times 10^{-9} \text{ cm}^{-1}$ . For the  $l_6=\pm 2$  sublevels, we obtain  $(\Delta^{Q,D}B^{(2)})_{exp} = -1.97(116) \times 10^{-9} \text{ cm}^{-1}$  and  $(\Delta^{Q,D}B^{(2)})_{th} = -1.740(31) \times 10^{-9} \text{ cm}^{-1}$ .

As said previously, the experimental frequencies could be reproduced with a global standard deviation of 41.4 kHz. It should be mentioned, however, that we have not succeeded in reproducing in such a quantitative manner the  $(k=\pm 1, l_6=\mp 2)$  and  $(k=\pm 5, l_6=\pm 2)$  series, as well as the low- $J$   $(k=\pm 3, l_6=0)$  transitions, which were excluded from the final fits.

These series connect levels pertaining to the nine-level system  $(0^0, A_+)$ ,  $(3^0, A_+/A_-)$ ,  $(1^2, A_+/A_-)$ ,  $(2^2, A_+/A_-)$  and  $(5^2, A_+/A_-)$  in which the higher-order diagonal splitting terms  $f_{24}$  and  $h_3$  compete with the, much stronger,  $q_{22}$  and  $q_{12}$  resonances. These latter resonances act globally all over the system of rotational levels of  $v_6=2$ , the  $q_{12}$  term producing a series of resonant level crossings. This is illustrated in Fig. 3 and also in a broader  $J/K$  span in the reduced energy diagram included in the Supplementary material.

The interplay of all these interactions produces strongly irregular patterns of the  $A_+/A_-$  splittings of the above mentioned series. The independent determination of the parameters of the three  $\Delta(k-l) = 6$  interactions (i.e.  $f_{42}$ ,  $f_{24}$ , and  $h_3$ ) would be probably enabled by the observation and inclusion in the fit of transitions with different selection rules, e.g. direct  $l$ -type resonance transitions, like achieved in analogous studies of the vibrational state  $v_{10}=2$  of  $\text{CF}_3\text{CCH}$  [25] or the vibrational state  $v_6=2$  of  $\text{SiHF}_3$  [26]. This was unfortunately not the case in the present work as such measurements would have required a completely different

experimental setup, which we did not have at our disposal. As a consequence, the  $f_{42}$  and  $h_3$  parameters could be determined, but may still remain correlated. The  $f_{24}$  parameter could not be determined and was kept fixed in the final stage of the calculations (as described earlier on page 10).

Nevertheless, the exclusion of the above mentioned low- $K$  series from the final fits did not deteriorate the consistency of the diagonal and off-diagonal parameters of the  $v_6=2$  level with those of  $v_6=1$ . This is shown in Table 5, for the reduction  $Q$ , by comparing the  $q_{22}$ ,  $f_{22}^J$ ,  $d_t$ , and  $f_{42}$  off-diagonal parameters of the two levels. Inspection of this Table shows an excellent agreement for  $q_{22}$ ,  $f_{22}^J$ , and  $d_t$  (within less than 2%), whereas the  $f_{42}$  values differ by about 5%. The same difference is obtained for the  $f_{42}$  values in the reduction  $D$ . Situation is somewhat worse for the off-diagonal parameters of  $v_3=v_6=1$ , some of which present larger differences with the corresponding  $v_6=1$  values (Table 5). In this case, excellent agreement is observed for  $q_{22}$  and  $f_{42}$  (within less than 2%), whereas the  $f_{22}^J$  values differ by about 5%, and the  $d_t$  values by about 14%.

Supplementary data associated with this article can be found in the online version, at doi: ...

## 5. Conclusions

The present work reports on the first high-resolution study of the  $v_6=2$  ( $A_1+E$ ) and  $v_3=v_6=1$  ( $E$ ) vibrational states of the  $\text{CH}^{35}\text{Cl}_3$  molecule through the analysis of the corresponding rotational spectra. The reproduction of the rotational data pertaining to  $v_3=v_6=1$  level is fully quantitative. This is also the case for the majority of the data pertaining to the  $v_6=2$  overtone level, except the  $(\pm 3,0)$ ,  $(\pm 1,\mp 2)$  and  $(\pm 5,\pm 2)$  series, where unusually large splittings occur.

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## Appendix A. Definition of the matrix elements of the effective vibration-rotational Hamiltonian

The diagonal matrix elements were defined conventionally as

$$\begin{aligned}
E_{vr}(J, k, l_t) = & E_v + B_v J(J+1) + (C_v - B_v) k^2 - D_v^v J^2(J+1)^2 - D_{JK}^v J(J+1) k^2 - D_K^v k^4 \\
& + H_J^v J^3(J+1)^3 + H_{JK}^v J^2(J+1)^2 k^2 + H_{KJ}^v J(J+1) k^4 + H_K^v k^6 \\
& + L_J^v J^4(J+1)^4 + L_{JK}^v J^3(J+1)^3 k^2 + L_{JK}^v J^2(J+1)^2 k^4 + L_{JKK}^v J(J+1) k^6 + L_K^v k^8 \\
& + \left[ -2(C\zeta_t)_v + \eta_{tJ} J(J+1) + \eta_{tK} k^2 + \tau_{tJ} J^2(J+1)^2 \right. \\
& + \tau_{tJK} J(J+1) k^2 + \tau_{tK} k^4 + \sigma_{tJ} J^3(J+1)^3 \\
& \left. + \sigma_{tJK} J^2(J+1)^2 k^2 + \sigma_{tKJ} J(J+1) k^4 + \sigma_{tK} k^6 \right] k l_t
\end{aligned} \tag{A1}$$

The following off-diagonal matrix elements were considered

a)  $l$ -type operators

$$\begin{aligned}
& \langle v_t^{l_t \pm 2}; J, k \pm 2 | \mathbf{H}_{22} + \mathbf{H}_{24} + \mathbf{H}_{26} | v_t^{l_t}; J, k \rangle \\
& = \left[ (v_t \mp l_t)(v_t \pm l_t + 2) \right]^{1/2} \left\{ q_{22} + f_{22}^J J(J+1) + f_{22}^K \left[ k^2 + (k \pm 2)^2 \right] \right. \\
& \left. + f_{22}^{JJ} J^2(J+1)^2 + f_{22}^{JK} J(J+1) \left[ k^2 + (k \pm 2)^2 \right] + f_{22}^{KK} \left[ k^4 + (k \pm 2)^4 \right] \right\} F_2^\pm(J, k)
\end{aligned} \tag{A2}$$

$$\begin{aligned}
& \langle v_t^{l_t \pm 2}; J, k \mp 1 | \mathbf{H}_{22} + \mathbf{H}_{24} | v_t^{l_t}; J, k \rangle \\
& = \left[ (v_t \mp l_t)(v_t \pm l_t + 2) \right]^{1/2} \left\{ \left[ q_{12} + f_{12}^J J(J+1) \right] (2k \mp 1) + f_{12}^K \left[ k^3 + (k \mp 1)^3 \right] \right\} F_1^\mp(J, k)
\end{aligned} \tag{A3}$$

$$\begin{aligned}
& \langle v_t^{l_t \pm 2}; J, k \mp 1 | \mathbf{H}_{41} + \mathbf{H}_{43} | v_t^{l_t}; J, k \rangle \\
& = \left[ (v_t \mp l_t)(v_t \pm l_t + 2) \right]^{1/2} \left\{ q_{12}^l + f_{12}^{lJ} J(J+1) + f_{12}^{lK} \left[ k^2 + (k \mp 1)^2 \right] \right\} (2l_t \pm 2) F_1^\mp(J, k)
\end{aligned} \tag{A4}$$

$$\begin{aligned}
& \langle v_t^{l_t \pm 2}; J, k \mp 4 | \mathbf{H}_{24} + \mathbf{H}_{26} | v_t^{l_t}; J, k \rangle = \left[ (v_t \mp l_t)(v_t \pm l_t + 2) \right]^{1/2} \left\{ f_{42} + f_{42}^J J(J+1) \right. \\
& \left. + f_{42}^K \left[ k^2 + (k \mp 4)^2 \right] \right\} F_4^\mp(J, k)
\end{aligned} \tag{A5}$$

$$\langle v_t^{l_t \pm 4}; J, k \mp 2 | \mathbf{H}_{42} | v_t^{l_t}; J, k \rangle = \frac{f_{24}}{4} \left[ (v_t \mp l_t - 2)(v_t \mp l_t)(v_t \pm l_t + 2)(v_t \pm l_t + 4) \right]^{1/2} F_2^\mp(J, k) \tag{A6}$$

$$\begin{aligned} \langle v_i^{l_i \pm 4}; J, k \pm 1 | \mathbf{H}_{42} | v_i^{l_i}; J, k \rangle &= \frac{f_{14}}{4} (2k \pm 1) [(v_i \mp l_i - 2)(v_i \mp l_i) \\ &\times (v_i \pm l_i + 2)(v_i \pm l_i + 4)]^{1/2} F_1^\pm(J, k) \end{aligned} \quad (\text{A7})$$

b)  $\Delta k = \pm 3$  operators

$$\langle v_i^{l_i}; J, k \pm 3 | \mathbf{H}_{23} + \mathbf{H}_{25} | v_i^{l_i}; J, k \rangle = \left\{ d_i + d_i^J J(J+1) + d_i^K [k^2 + (k \pm 3)^2] \right\} l_i F_3^\pm(J, k) \quad (\text{A8})$$

$$\langle v_i^{l_i}; J, k \pm 3 | \mathbf{H}_{04} + \mathbf{H}_{06} | v_i^{l_i}; J, k \rangle = \left\{ [\varepsilon + \varepsilon_J J(J+1)](2k \pm 3) + \varepsilon_K [k^2 + (k \pm 3)^2] \right\} F_3^\pm(J, k) \quad (\text{A9})$$

c)  $\Delta k = \pm 6$  operator

$$\langle v_i^{l_i}; J, k \pm 6 | \mathbf{H}_{06} + \mathbf{H}_{08} | v_i^{l_i}; J, k \rangle = \left\{ h_3 + h_3^J J(J+1) + h_3^K [k^2 + (k \pm 6)^2] \right\} F_6^\pm(J, k) \quad (\text{A10})$$

The notation of the matrix elements of the rotational shifting operators was taken as

$$F_n^\pm(J, k) = \prod_{i=1}^n [J(J+1) - (k \pm i \mp 1)(k \pm i)]^{1/2} \quad (\text{A11})$$



## Figure captions

**Figure 1.** The rotational spectrum of chloroform in the region of  $J=26\leftarrow 25$  transitions. Upper panel: experimental spectrum (natural sample). Lower panel: simulated spectrum, showing several  $J$ -clusters of the  $\text{CH}^{35}\text{Cl}_3$  isotopologue, as follows: ground vibrational state – yellow;  $\nu_3=1$  – gray;  $\nu_6=1$  – blue;  $\nu_3=\nu_6=1$  – black;  $\nu_6=2$  – red.

**Figure 2.** The rotational spectrum of chloroform in the region of  $J=61\leftarrow 62$  transitions. Upper panel: experimental spectrum (natural sample). Lower panel: simulated spectrum, showing several  $J$ -clusters of the  $\text{CH}^{35}\text{Cl}_3$  isotopologue, with the same color code as for Figure 1. The lines marked with an asterisk are ground state rotational transitions of the  $\text{CH}^{35}\text{Cl}_2^{37}\text{Cl}$  isotopologue.

**Figure 3.** Reduced energies (pure  $J$ -dependent terms not included) against  $J$ , for the  $\nu_6=2$  levels. Positions of resonant crossings due to  $q_{12}$ ,  $f_{42}$ , and  $f_{14}$  interactions are indicated. Dashed arrows indicate interaction between only one of either the  $A_+$  or  $A_-$  sublevels.

**Figure 4.** Hyperfine splittings of the rotational transitions in the  $\nu_6=1$  level. The  $J=34$ ,  $l_6=-1$  cluster is shown; the simulated spectrum (lower panel) does not account for these splittings. Several, much weaker,  $\nu_6=2$  transitions, are indicated by an asterisk.

## Table captions

**Table 1.** Molecular parameters of the  $\nu_3=\nu_6=1$  level of  $\text{CH}^{35}\text{Cl}_3$ .

**Table 2.** Demonstration of unitary equivalence between the  $D$  and  $Q$  reductions, for the  $\nu_3=\nu_6=1$  level.

**Table 3.** Molecular parameters of the  $\nu_6=2$  level of  $\text{CH}^{35}\text{Cl}_3$ .

**Table 4.** Demonstration of unitary equivalence of the different parameter sets for the  $\nu_6=2$  level (2<sup>nd</sup> and 3<sup>rd</sup> order).

**Table 5.** Parameters of the (2,2), (3,0) and (4,-2) interactions (in MHz units) in the vibrational states  $\nu_6=1$ ,  $\nu_3=\nu_6=1$ , and  $\nu_6=2$  of  $\text{CH}^{35}\text{Cl}_3$  (reduction  $Q$ ).

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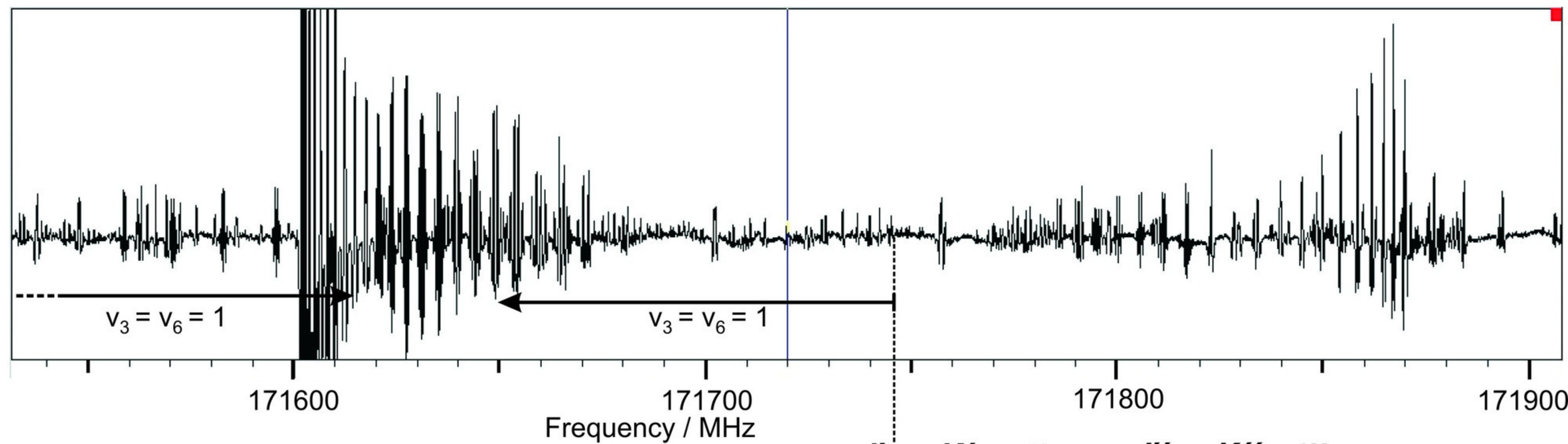
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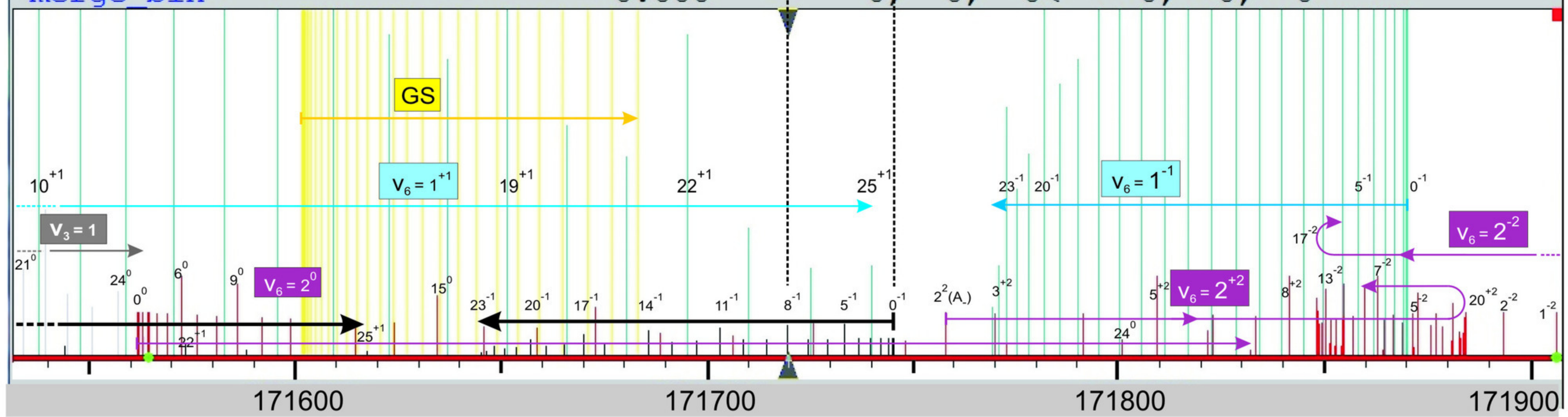
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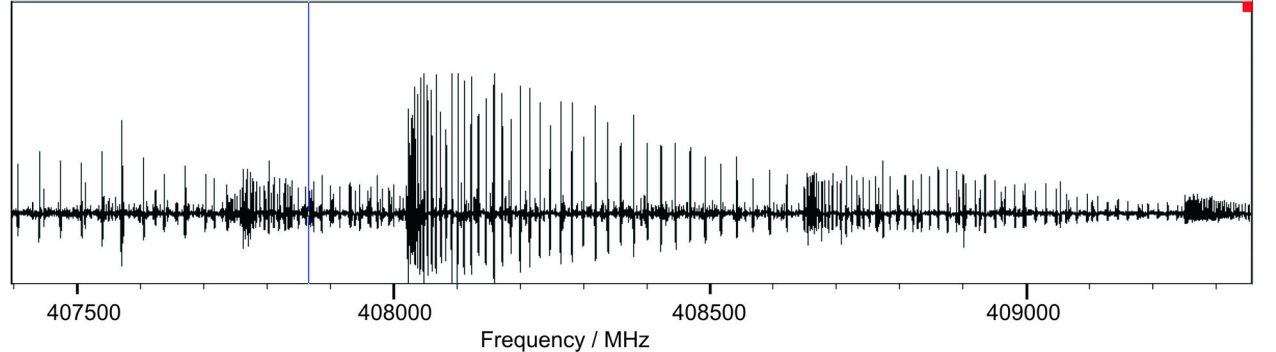
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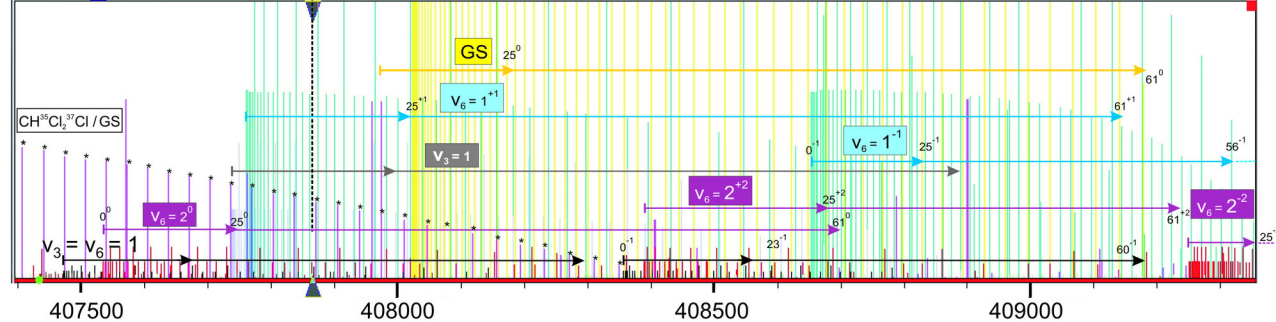


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merge bin			0.000	0,	8,	0 <--	25,	8,	-1		
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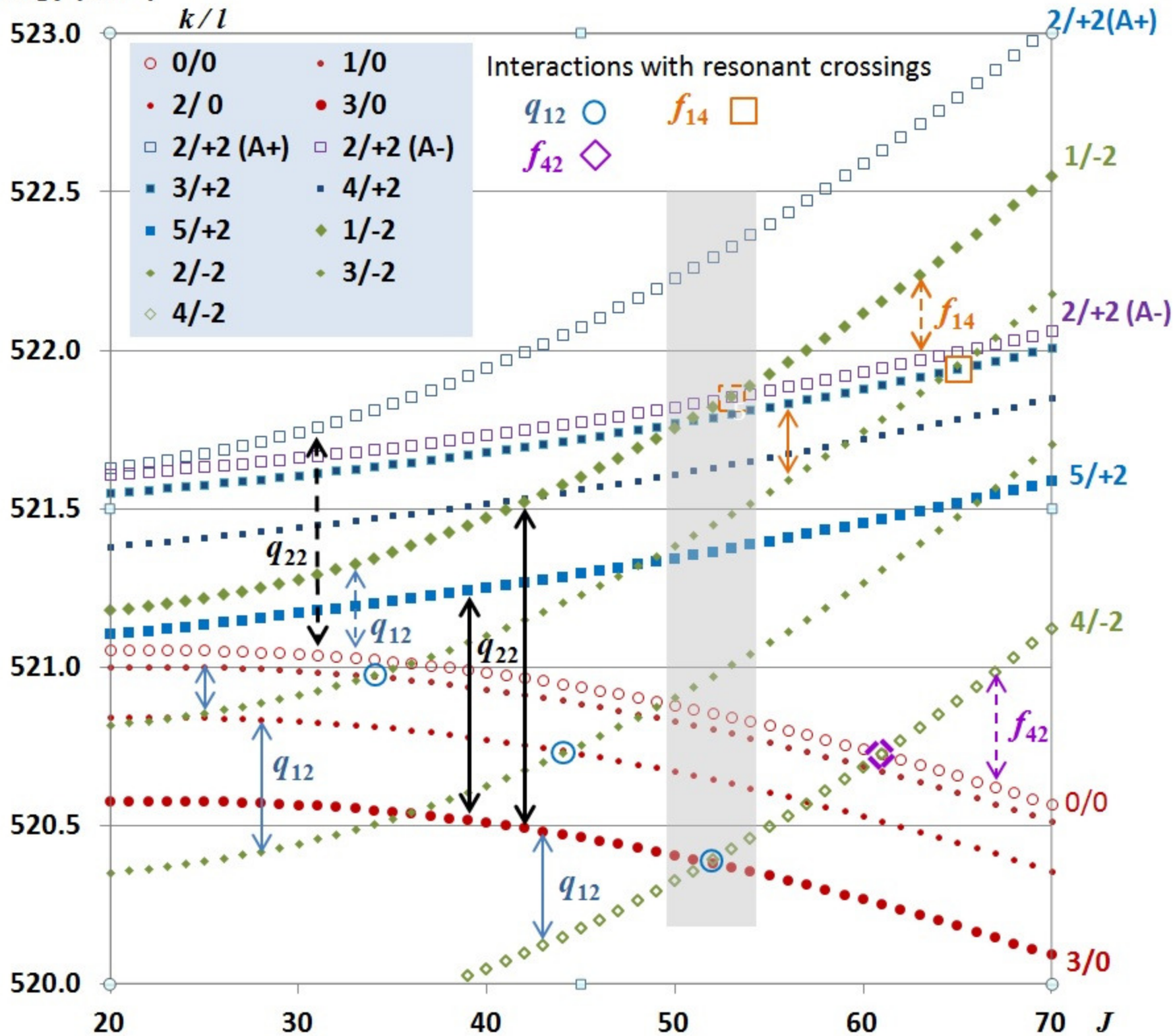


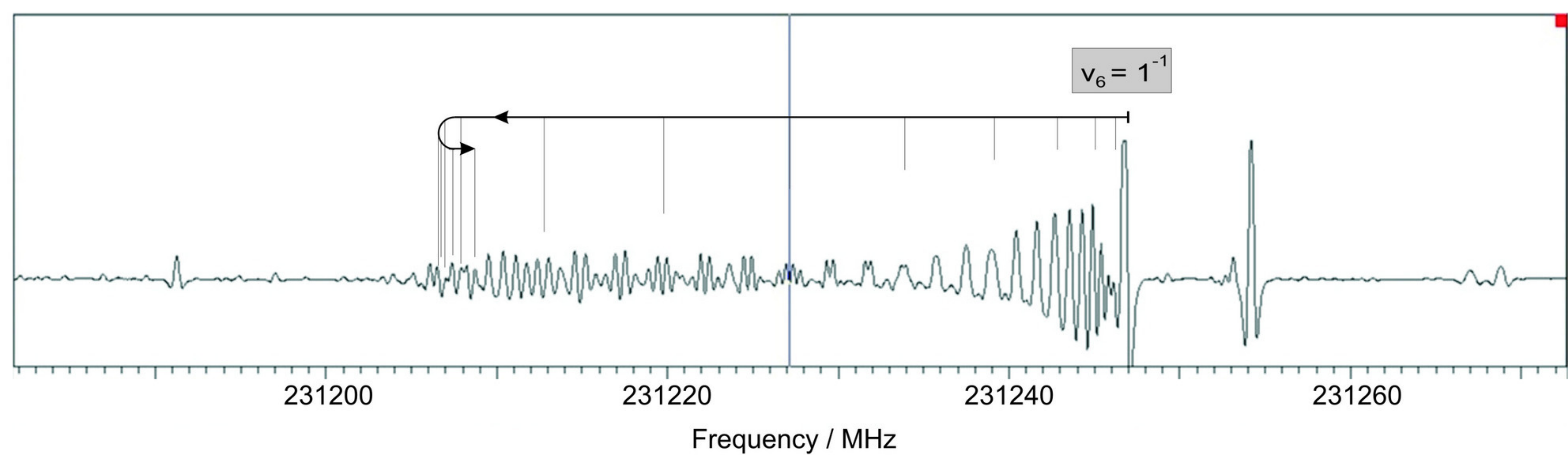


Line	1 / 1:	407867.7890	1.45E-01	<i>J'</i>	<i>K'</i>	<i>l'</i> ←	<i>J''</i>	<i>K''</i>	<i>l''</i>		
			0.000	62	34	1 <--	61	34	1	■	.R 0, 0
	all2_bin			0	0	0 <--	0	0	0		

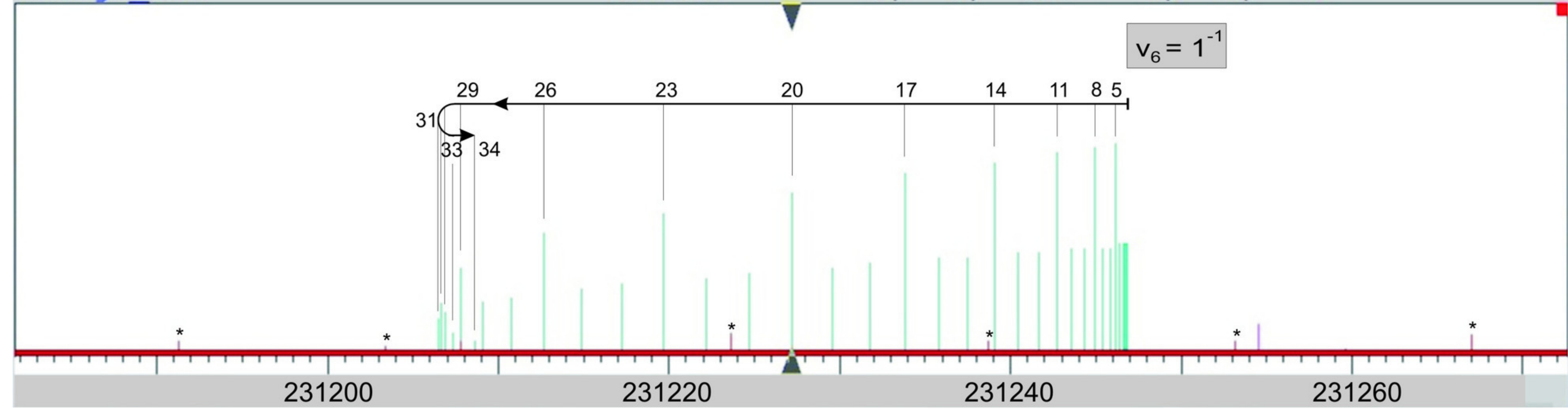




Energy (cm<sup>-1</sup>)



Line 1/ 1: 231227.1648 8.68E+00 35, 20, -1<-- 34, 20, -1 merge\_bin 0.000 0, 0, 0<-- 0, 0, 0 .R 0, 0





**Table 1. Molecular parameters of the  $v_3=v_6=1$  level of  $\text{CH}^{35}\text{Cl}_3$** 

Parameter	Reduction $D$	Reduction $Q$
$B$ (MHz)	3301.228 336 (55)	3301.228 335 (55)
$C$ (MHz)	1711.524 214 <sup>a</sup>	1711.524 214 <sup>a</sup>
$D_J$ (kHz)	1.514 619 (13)	1.514 618 (13)
$D_{JK}$ (kHz)	-2.527 72 (22)	-2.527 71 (21)
$D_K$ (kHz)	0.834 58 <sup>b</sup>	0.834 58 <sup>b</sup>
$H_J$ (mHz)	1.256 7 (11)	1.256 6 (11)
$H_{JK}$ (mHz)	-5.092 (22)	-5.089 (22)
$H_{KJ}$ (mHz)	6.96 (22)	6.94 (22)
$H_K$ (mHz)	-2.435 <sup>b</sup>	-2.435 <sup>b</sup>
$L_J$ (nHz)	-2.538 <sup>b</sup>	-2.538 <sup>b</sup>
$C\zeta$ (MHz)	-1502.060 8 (84)	-1502.062 4 (84)
$\eta_J$ (kHz)	-4.407 (20)	-4.392 (20)
$\eta_K$ (kHz)	3.267 5 <sup>b</sup>	3.264 9 <sup>b</sup>
$q_{22}$ (MHz)	-1.813 161 (18)	-1.813 154 (18)
$f_{22}^J$ (Hz)	4.051 4 (21)	4.052 9 (21)
$f_{22}^K$ (Hz)	-2.981 (19)	-2.975 (19)
$f_{22}^{JJ}$ ( $\mu\text{Hz}$ )	-8.677 <sup>b</sup>	-8.677 <sup>b</sup>
$f_{22}^{JK}$ ( $\mu\text{Hz}$ )	23.058 <sup>b</sup>	23.058 <sup>b</sup>
$q_{12}$ (kHz)	-82.07 (17)	0.
$d_r$ (kHz)	0.	0.129 03 (27)
$f_{42}$ (Hz)	0.490 0 (59)	0.493 5 (59)
$h_3$ (mHz)	0.155 5 (10)	0.155 6 (10)
$\sigma$ (kHz)	41.2	41.2

<sup>a</sup> Constrained to a value obtained from the GS,  $v_3=1$ , and  $v_6=1$  values (Refs. [9] and [12]).

<sup>b</sup> Constrained to the corresponding  $v_6=1$  value quoted in [12].

Numbers in parentheses are standard deviations in units of the last digit quoted.

**Table 2. Demonstration of the unitary equivalence between the  $D$  and  $Q$  reductions, for the  $v_3=v_6=1$  level**

Equation <sup>a</sup>	Unit	Verification
$-2 q_{12}^D / F$	$10^{-5}$	-3.573 2 (75)
$\frac{1}{2} d_t^Q / q_{22}$	$10^{-5}$	-3.558 0 (75)
$-8( q_{12}^D)^2 / F$	$10^{-10} \text{ cm}^{-1}$	3.913 (16)
$-\frac{1}{2} F( d_t^Q)^2 / q_{22}^2$	$10^{-10} \text{ cm}^{-1}$	3.880 (16)
$(\eta_J^Q - \eta_J^D)$	$10^{-10} \text{ cm}^{-1}$	5.1 (94)
$-(\eta_K^Q - \eta_K^D)$	$10^{-10} \text{ cm}^{-1}$	0.89 <sup>b</sup>

<sup>a</sup>  $F = C - B + 2C\zeta$ .

<sup>b</sup> Constrained to the value obtained in Ref. [12] for the  $v_6=1$  level.

**Table 3. Molecular parameters of the  $v_6=2$  states of  $\text{CH}^{35}\text{Cl}_3$**

Parameter		Reduction D	Reduction Q	Reduction L	
$l_6=0$	$B$	MHz	3305.138 203 (80)	3305.138 361 (80)	3305.138 006 (80)
	$C$	MHz	1710.636 4 (37)	1710.636 3 (37)	1710.636 4 (37)
	$D_J$	kHz	1.520 418 (16)	1.520 426 (16)	1.520 405 (16)
	$D_{JK}$	kHz	-2.545 53 (14)	-2.545 54 (14)	-2.545 54 (14)
	$D_K$	kHz	0.834 58 <sup>a</sup>	0.834 58 <sup>a</sup>	0.834 58 <sup>a</sup>
	$H_J$	mHz	1.284 09 (96)	1.284 16 (96)	1.284 00 (96)
	$H_{JK}$	mHz	-5.172 8 (71)	-5.172 9 (71)	-5.173 5 (71)
	$H_{KJ}$	mHz	6.902 (20)	6.903 (20)	6.901 (20)
	$H_K$	mHz	- 2.434 67 <sup>a</sup>	- 2.434 67 <sup>a</sup>	- 2.434 67 <sup>a</sup>
	$L_J$	nHz	- 2.538 <sup>a</sup>	- 2.538 <sup>a</sup>	- 2.538 <sup>a</sup>
$l_6=\pm 2$	$4g_{ll}$	MHz	10468.65 (17)	10468.64 (17)	10468.67 (17)
	$C$	MHz	1710.665 5 (33)	1710.665 5 (33)	1710.665 5 (33)
	$D_J$	kHz	1.520 277 (10)	1.520 273 (10)	1.520 283 (10)
	$D_{JK}$	kHz	-2.545 540 (79)	-2.545 535 (79)	-2.545 541 (79)
	$D_K$	kHz	0.835 04 (20)	0.835 03 (20)	0.835 06 (20)
	$H_J$	mHz	1.293 13 (60)	1.293 09 (60)	1.293 17 (60)
	$H_{JK}$	mHz	-5.163 6 (41)	-5.163 4 (41)	-5.163 5 (41)
	$H_{KJ}$	mHz	6.774 0 (87)	6.773 4 (87)	6.775 2 (87)
	$H_K$	mHz	- 2.434 67 <sup>a</sup>	- 2.434 67 <sup>a</sup>	- 2.434 67 <sup>a</sup>
	$L_J$	nHz	- 2.538 <sup>a</sup>	- 2.538 <sup>a</sup>	- 2.538 <sup>a</sup>
Off-diagonal parameters	$C\zeta$	MHz	-1502.373 4 (31)	-1502.373 0 (31)	-1502.374 2 (31)
	$\eta_J$	kHz	-4.951 4 (18)	-4.942 4 (18)	-5.001 7 (18)
	$\eta_K$	kHz	3.238 3 (23)	3.229 3 (23)	3.288 3 (23)
	$q_{12}$	kHz	-72.507 (70)	0.	-184.78 (49)
	$q_{12}^I$	kHz	91.50 (40)	150.58 (40)	0.
Off-diagonal parameters	$q_{12}^{I,J}$	Hz	-6.541 (88)	-6.602 (88)	-6.445 (87)
	$q_{22}$	MHz	-1.829 779 5 (77)	-1.829 780 3 (76)	-1.829 778 7 (77)
	$f_{22}^J$	Hz	3.861 4 (14)	3.863 3 (14)	3.865 5 (14)

$f_{22}^K$	Hz	-3.028 6 (19)	-3.029 4 (19)	-3.031 9 (19)
$f_{22}^J$	$\mu$ Hz	-8.738 (82)	-8.743 (82)	-8.716 (82)
$f_{22}^{JK}$	$\mu$ Hz	23.011 <sup>a</sup>	23.011 <sup>a</sup>	23.011 <sup>a</sup>
$d_t$	kHz	0.	0.114 89 (11)	-0.177 90 (77)
$f_{42}$	Hz	0.506 69 (45)	0.509 51 (44)	0.510 15 (45)
$f_{24}$	kHz	-0.661 <sup>b</sup>	-0.661 <sup>b</sup>	-0.661 <sup>b</sup>
$h_3$	mHz	0.154 14 (99)	0.153 88 (99)	0.154 50 (99)
$\sigma$	kHz	41.4	41.4	41.4

<sup>a</sup> Constrained to the  $v_6=1$  values obtained in Ref. [12].

<sup>b</sup> Constrained to a value obtained in a preliminary fit, which still included the ( $k=1, l_6=-2$ ) lines. These lines were excluded from the final fit (see text).

Numbers in parentheses are standard deviations in units of the last digit quoted.

**Table 4**

**Demonstration of unitary equivalence of the different parameter sets for the  $v_6=2$  level (2<sup>nd</sup> and 3<sup>rd</sup> order)**

Equation <sup>a</sup>	Unit	Verification	Equation <sup>a</sup>	Unit	Verification
$-2 q_{12}^D / F$	$10^{-5}$	-3.153 0 (30)	$2 q_{12}^{L,D} / G$	$10^{-5}$	4.903 (21)
$= \frac{1}{2} d_t^Q / q_{22}$	$10^{-5}$	-3.139 4 (30)	$= \frac{1}{2} d_t^L / q_{22}$	$10^{-5}$	4.861 (21)
$q_{12}^{L,D}$	$10^{-6} \text{ cm}^{-1}$	3.052 (13)	$q_{12}^L$	$10^{-6} \text{ cm}^{-1}$	-6.164 (16)
$= q_{12}^{L,Q} + \frac{G d_t^Q}{4 q_{22}}$	$10^{-6} \text{ cm}^{-1}$	3.069 (13)	$= q_{12}^D + \frac{F}{G} q_{12}^{L,D}$	$10^{-6} \text{ cm}^{-1}$	-6.180 (16)
$-8 (q_{12}^D)^2 / F$	$10^{-10} \text{ cm}^{-1}$	3.050 4 (59)	$-8 (q_{12}^L)^2 / F$	$10^{-9} \text{ cm}^{-1}$	1.981 (11)
$= -\frac{1}{2} F (d_t^Q)^2 / q_{22}^2$	$10^{-10} \text{ cm}^{-1}$	3.024 1 (58)	$= -8 F (q_{12}^{L,Q})^2 / G^2$	$10^{-9} \text{ cm}^{-1}$	1.998 (11)
$= (\eta_J^Q - \eta_J^D)$	$10^{-10} \text{ cm}^{-1}$	3.02 (84)	$= (\eta_J^Q - \eta_J^L)$	$10^{-9} \text{ cm}^{-1}$	1.980 (85)
$= -(\eta_K^Q - \eta_K^D)$	$10^{-10} \text{ cm}^{-1}$	2.99 (106)	$= -(\eta_K^Q - \eta_K^L)$	$10^{-9} \text{ cm}^{-1}$	1.966 (107)

<sup>a</sup>  $F = C - B + 2C\zeta$ ;  $G = C\zeta + 2g_u$ .

**Table 5. Parameters of the (2,2), (3,0) and (4,-2) interactions (in MHz units) in the vibrational states  $v_6=1$ ,  $v_3=v_6=1$ , and  $v_6=2$  of  $\text{CH}^{35}\text{Cl}_3$  (reduction  $Q$ )**

Parameter	$v_6=1^a$	$v_3=v_6=1$	$v_6=2$
$q_{22}$	-1.827 752 (19)	-1.813 154 (18)	-1.829 780 3 (76)
$f_{22}^J \times 10^6$	3.855 2 (43)	4.052 9 (21)	3.863 3 (14)
$d_i \times 10^4$	1.127 14 (15)	1.290 3 (27)	1.148 89 (111)
$f_{42} \times 10^7$	4.873 6 (60)	4.935 (59)	5.095 1 (44)

<sup>a</sup> Parameters taken from [12].

