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Impact of channel length on the operation of junctionless transistors with substrate-biasing

Dae-Young Jeon, Mireille Mouis, *Senior Member, IEEE*, Sylvain Barraud and Gérard Ghibaudo, *Fellow, IEEE*

Abstract—Junctionless transistors (JLTs) have promising advantages such as structural simplicity without PN-junctions and bulk conduction-based operation for the realization of advanced CMOS technologies. Here, channel length dependence on the operation of JLTs with substrate-biasing (V_{gb}) was investigated in detail. Parasitic series resistance (R_{sd}) was noticeably decreased as increasing V_{gb} . In addition, transconductance (g_m) , its derivative (dg_m/dV_{gf}) and on-drain current (I_{on}) in a short-channel JLT were significantly affected by the V_{gb} -modulated R_{sd} with charge coupling effects. This work provides important information for better understanding and true estimation of intrinsic JLT performance, for practical applications based on polycrystalline Si, III-V semiconductors and TMDs nano-materials as well as advanced logic devices.

Index Terms—Junctionless transistors (JLTs), substrate-biasing effect, channel length dependence, series resistance (R_{sd}), electrical parameters modulated by substrate bias.

I. INTRODUCTION

tructural simplicity without PN-junctions of junctionless transistors (JLTs) is the most promising advantage for the realization of sub-5 nm complementary metal-oxide-semiconductor (CMOS) technologies [1-5].

The junctionless architecture with uniformly doped source (S), channel and drain (D) enables no need of precise gate-alignment and abrupt S/D junction formation with a low-thermal-budget [3-5]. Bulk conduction-based operation in JLTs also results in better reliability, lower 1/f noise and less mobility degradation against transverse electric-field (E-field)

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[6-9]. In addition, JLTs in a variety of materials, including polycrystalline silicon (Si), III-V semiconductors and 2-dimentional (2D) transition metal dichalcogenides (TMDs), have been reported for high performance thin-film transistors (TFTs), 3D vertical transistors, sensors and multi-functional devices with excellent short-channel characteristics [10-14].

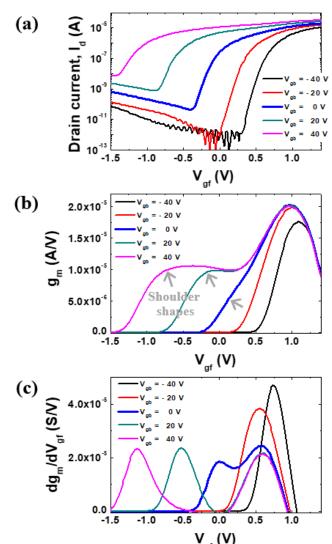


Fig. 1. (a) Transfer curves, (b) transconductance (g_m) and (c) derivative of g_m (dg_m/dV_{gf}) measured from a very long channel JLT of $L_m=10~\mu m$ with varying $V_{gb}.$ The device was biased in the linear operation-regime with drain voltage $V_d=50~mV.$

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Silicon-on-insulator (SOI) technology based JLTs can be operated by both front-gate and back-gate [15,16]. In-depth studies about substrate-biasing effects and charge coupling between front and back interfaces are very important for optimizing performance and better understanding operation mechanism of JLTs.

In this paper, channel length dependent substrate-biasing effects on the operation of JLTs were investigated in detail. Parasitic series resistance (R_{sd}) of JLTs was noticeably modulated by the substrate bias. Transconductance (g_m), its derivative (dg_m/dV_{gf}) and effective mobility in a short-channel JLT were significantly affected by the substrate-bias-controlled R_{sd} with charge coupling effects between front and back interfaces.

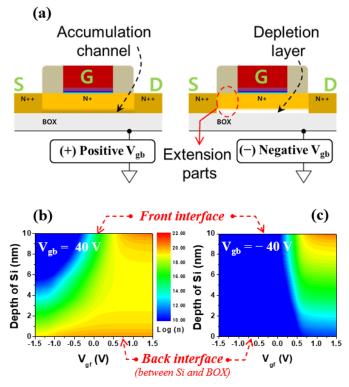


Fig. 2. (a) Schematic architecture of JLT describing the accumulated charge carriers at back interface for positive V_{gb} , and the depletion of charge carriers due to negative V_{gb} . Simulated distribution of electron carrier concentration per unit volume (n) along the depth of device Si varying V_{gf} with different V_{gb} , (b) $V_{gb} = 40 \text{ V}$ and (c) $V_{gb} = -40 \text{ V}$.

II. EXPERIMENTS

JLTs with device Si \approx 10 nm and high-k dielectric \approx 1.2 nm were fabricated on (100) SOI wafers with thick buried oxide (BOX) \approx 145 nm. Channel of JLTs was highly doped with

phosphorus, and effective doping concentration was estimated as $\approx 5 \times 10^{18} / \text{cm}^3$. An additional S/D implantation was also performed to reduce access resistance. Further detail information with respect to fabrication and material parameters was described in previous works [17]. Electrical characteristics were measured by using HP 4155A semiconductor parameter analyzer, and simulated results were obtained by FlexPDE software, which is a numerical solver based on scripted finite element models.

III. DISCUSSION

The bold blue plots as shown in Fig.1 show transfer curve, transconductance (g_m) and derivative of g_m (dg_m/dV_{gf}) , where V_{gf} means front-gate voltage), obtained from a very long channel JLT ($L_m = 10 \mu m$) with back-gate voltage (substrate bias) $V_{gb} = 0V$. Shoulder shape on the g_m and two clear peaks on the $dg_{\text{m}}\!/dV_{\text{gf}}$ in JLTs represent unique electrical properties based on both bulk neutral and surface accumulation conductions, which is very different from the operation of conventional inversion-mode (IM) transistors [8,17]. The first (left-hand side) and second peak (right-hand side) on dg_m/dV_{gf} denote conduction threshold voltage (V_{th}) and flat-band voltage (V_{fb}), i.e. threshold of accumulation conduction channel, respectively. The shoulder shape on g_m is also strongly associated with the amount of charge carriers on the operation at $V_{gf} = V_{fb}$ condition. As V_{gb} decreased, the shoulder shrunk in Fig. 1(b) and the two peaks merged as a single peak in Fig. 1(c), due to the effectively reduced bulk channel thickness resulted from the negative V_{gb}-induced depletion of free charge carriers at back interface between device Si and BOX [18]. On the other hand, bulk neutral channel of JLTs can be early turned on, and additional accumulation channel is newly created at the back interface with a positive V_{gb}. Those gave rise to the pronounced shoulder shape in Fig. 1(b) and left-shift of the first peak in Fig. 1(c), as V_{gb} increased.

Figure 2 shows the simulated substrate-biasing characteristics of JLTs. The distribution of electron carrier concentration per unit volume (n) along the depth of device Si was visualized by varying V_{gf} with different V_{gb} . One can clearly identified the shrunk bulk channel with negative V_{gb} , and the early turned-on

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bulk channel with the newly accumulated carriers for positive V_{gb} . Indeed, a depletion layer in blue color was observed at back interface for $V_{gb} = -40 \text{ V}$ in Fig. 2(c), while an additional accumulation layer in red or orange color was created at the back interface for $V_{gb} = 40 \text{ V}$ in Fig. 2(b).

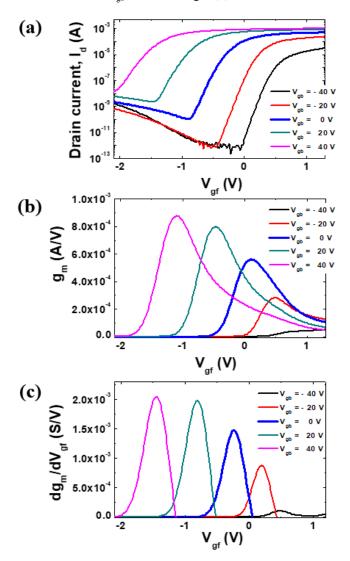


Fig. 3. (a) Transfer curves, (b) transconductance (g_m) and (c) derivative of g_m (dg_m/dV_{gf}) measured from a very short channel JLT of L_m = 40 nm with varying V_{gb} . The device was biased in the linear operation-regime with drain voltage V_d = 50 mV.

Completely different substrate-biasing behaviors were investigated in JLTs with a very short channel length ($L_m = 40$ nm), as shown in Fig. 3. A shoulder shape was not observed any more in Fig. 3(b), and only a single peak was shown in Fig. 3(c), since the parasitic series resistance (R_{sd}) considerably affects the short-channel JLTs characteristics [19,20]. In addition, the on-drain current (I_{on} , $V_{gf} > 1.0$ V) in Fig. 3(a) was

dramatically changed by varying V_{gb} , as compared to that of the long channel JLT in Fig. 1(a). The g_m in Fig. 3(b) was also getting noticeably degraded as decreasing V_{gb} . The variation of V_{gb} could have an effect on the value of R_{sd} , and the R_{sd} modulated by V_{gb} could affect overall results related to substrate-biasing in the short channel JLTs.

 V_{gb} controlled R_{sd} of JLTs was estimated by using transfer length method (TLM) in Fig. 4. Total resistance (R_{tot}) including JLT channel resistance (R_{ch}) and R_{sd} can be calculated as [21]:

$$R_{tot} = R_{ch} + R_{sd} = \rho_{ch} \times \frac{L_m - \Delta L}{W_m} + R_{sd}$$
 (1)

where $\rho_{ch},~\Delta L$ and W_m are channel sheet resistance, channel length reduction and on-mask channel width, respectively. The clear common intersection point of the linear fitted lines on the accumulation operation regime $(V_{gf} > V_{fb})$ successfully determined $R_{sd} \approx 104~\Omega$ for $V_{gb} = 0V,$ as shown in Fig. 4(a).

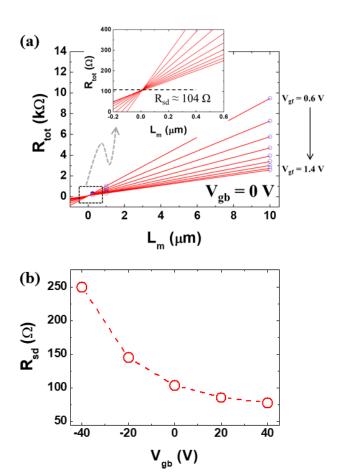


Fig. 4. (a) Total resistance (R_{tot}) versus channel lengths ($L_m = 250$ nm, 1 μm and 10 μm) with varying V_{gf} , for determination of series resistance (R_{sd}). The inset shows a clear common intersection of the linear fitted lines. (b) Extracted R_{sd} modulated by V_{gb} .

Finally, V_{gb} dependent R_{sd} was plotted in Fig. 4(b) after the same procedure for R_{sd} extraction with different V_{gb} from -40 V to 40 V. Indeed, R_{sd} was decreasing for positive V_{gb} , while it was increasing when applying negative V_{gb} . Accumulated charge carriers at back interface for positive V_{gb} reduces resistance of the extension parts (R_{ext}), as illustrated in Fig. 2(a). R_{ext} can be also enhanced by the depletion of charge carriers due to negative V_{gb} . Similar trend of the substrate-bias affected R_{sd} have been also reported in SOI based transistors with lightly doped drain (LDD) [22].

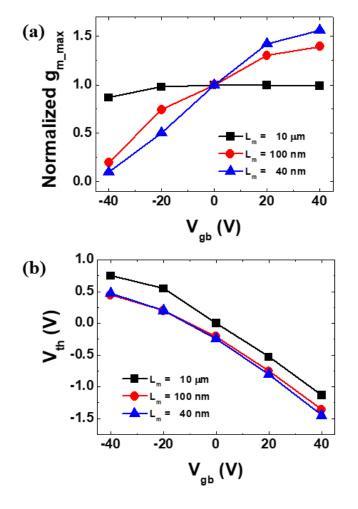


Fig. 5. (a) Normalized g_{m_max} [i.e. maximum- $g_m(V_{gb})$ /maximum- $gm(V_{gb}=0\ V)$] versus V_{gb} , with different JLT lengths. (b) Threshold voltage (V_{th}) as a function of V_{gb} with different JLT lengths, the V_{th} was extracted from the first peak in the dg_m/dV_{gf} plots.

Figure 5(a) shows the normalized g_{m_max} [i.e. maximum- $g_m(V_{gb})$ /maximum- $g_m(V_{gb} = 0 \text{ V})$] versus V_{gb} , with different JLT lengths. It should be noted that the trend of g_m is closely related to characteristics of field-effect mobility

behaviors [23]. There was a slight decrease in the normalized g_{m_max} of long length JLT (L $_m$ = 10 $\mu m)$ as decreasing $V_{gb}.$ This feature is attributed to the lower mobility at the top interface, since the high-k dielectric (HfSiON) for top-gate insulator has poorer interface-properties than that in pure thermal oxide (SiO₂) of BOX at the back interface [15-17]. However, the variation of V_{gb} significantly changed the normalized g_{m_max} in JLTs with $L_m = 40$ nm. The increased value of R_{sd} at a negative $V_{\mbox{\scriptsize gb}}$ caused the dramatically enhanced degradation of $g_{\mbox{\scriptsize m}}$ and field-effect mobility in the short channel JLTs. Moreover, the threshold voltage (V_{th}) extracted from the first peak in the dg_m/dV_{gf} plots was plotted as a function of V_{gb} as shown in Fig. 5(b). As decreasing length of JLTs with a same V_{gb} , V_{th} was reduced due to short-channel effects [17]. Whereas the variation rate of V_{th} against V_{gb} changes was very similar regardless of JLT lengths. Interestingly, this means that substrate-biasing impact on the operation near V_{th} regime is independent of channel-length variation in JLTs.

IV. CONCLUSIONS

The operation of junctionless transistors (JLTs) with substrate-biasing (V_{gb}) has been investigated, and its channel length dependence was also discussed in detail. A shoulder shape on transconductance (g_m) and clear two peaks on derivative of g_m (dg_m/dV_{gf}) were observed in a long-channel $(L_{\text{m}}$ = 10 $\mu\text{m})$ JLT with V_{gb} = 0V, since JLTs have both bulk neutral and surface accumulation conduction channels. The shoulder shape shrunk and the two peaks merged as a single peak when applying negative $V_{\rm gb}$. For positive $V_{\rm gb}$, the shoulder shape was pronounced and the first peak on the $dg_m/dV_{\rm gf}$ moved to the left-hand side. Completely different features (no shoulder shape on g_m and only a single peak on dg_m/dV_{gf}) were investigated in a very short-length ($L_m = 40 \text{ nm}$) JLT due to parasitic series resistance (R_{sd}) effect. The R_{sd} was noticeably decreased as increasing V_{gb}, and this gave rise to the dramatical changes on-drain current (Ion) and field-effect mobility of the short-channel JLTs against the variation of V_{gb}. However, substrate-biasing impact on the operation near V_{th} regime was similar regardless of JLT lengths. Further calibration with the experimental data and stronger coupling

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effects in ultra-thin body and buried oxide (UTBB) structures will be our future works [24,25].

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