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Linear Response for a Family of Self-Consistent Transfer Operators

Fanni M. Sélley*, Matteo Tanzi†

Abstract

We study a system of globally coupled uniformly expanding circle maps in the thermodynamic limit. The state of the system is described by a probability density and its evolution is given by the action of a nonlinear operator called the self-consistent transfer operator. Self-consistency is understood in the following sense: if φ corresponds to the system's state, then the evolution of φ is given by the application of the transfer operator of the circle map F_φ . This is a φ -dependent map describing the dynamics of a single unit in the finite system where the interaction coming from all the nodes is replaced with the average of the interaction term with respect to the density φ . Assuming some level of smoothness of the coupled circle maps and of the coupling, we prove that when the coupling strength is sufficiently small, the system has a unique smooth stable state. We then show that this stable state is continuously differentiable as a function of the coupling strength. Finally, we prove that the derivative satisfies a linear response formula which at zero simplifies to a formula reminiscent to the one obtained for the linear transfer operators of perturbed expanding circle maps.

1 Introduction

A fundamental question in the theory of dynamical systems is how the statistical properties of a system change when it is subjected to perturbation. A system exhibits linear response if the invariant measure depends smoothly on the perturbation, and an expression for the derivative of the invariant measure (in the strong or in the dual sense) is called a linear response formula.

The rigorous study of linear response dates back to Ruelle, who proved linear response of uniformly hyperbolic Axiom A systems [29, 30, 31]. Similar results can be proved in some non-uniformly hyperbolic [12, 37] or non-uniformly expanding cases such as intermittent maps [2, 9, 21] and piecewise expanding unimodal maps [7, 8]. But caution is required, as examples for the lack of linear response are also well known [1, 3, 5, 6] (see [4] for a more comprehensive collection of references).

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Deterministic systems subjected to random perturbations are also known to exhibit linear response in some cases [14, 16]. In addition to theoretical works, linear response theory is successfully applied in geophysics, in particular climate science [24, 27].

In this paper we rigorously prove a linear response formula for all-to-all weakly coupled uniformly expanding maps in the thermodynamic limit. To the best of our knowledge, there are few results available on linear response of complex systems with multiple interacting units. A program focusing on the control of statistical properties of extended systems by exploiting linear response was outlined by Mackay [26]. This question was further pursued in the case of a single dynamical unit defined by a uniformly expanding circle map [15] and by a smooth map of a compact smooth Riemannian manifold [20]. In case of coupled map lattices, the smooth dependence of the SRB measure and the linear response formula for was studied by Jiang and de la Llave [18, 17]. Wormell and Gottwald [35, 36] recently showcased numerical examples and some rigorous arguments for all-to-all coupled maps with mean-field interaction showing that it is possible for a high dimensional compound system to exhibit linear response, even if its units do not.

Our setup is based on the model studied by Fernandez [13]. The finite size model is described by M coordinates (x_1, \dots, x_M) with $x_i \in \mathbb{T}$, whose evolution is given by

$$x_i(t+1) = f \circ \Phi_\varepsilon(x_i(t); x_1(t), \dots, x_M(t)) \quad (1)$$

where $\Phi_\varepsilon : \mathbb{T} \times \mathbb{T}^M \rightarrow \mathbb{T}$,

$$\Phi_\varepsilon(x_i; x_1, \dots, x_M) = x_i + \frac{\varepsilon}{M} \sum_{j=1}^M h(x_i, x_j), \quad (2)$$

is a smooth mean-field all-to-all coupling between the coordinates (ε is the coupling strength*), and $f : \mathbb{T} \rightarrow \mathbb{T}$ is a sufficiently smooth map describing the uncoupled dynamics.

When taking the thermodynamic limit, $M \rightarrow \infty$, one can replace the finite dimensional state space \mathbb{T}^M with the infinite-dimensional state space $\mathbb{T}^\mathbb{N}$, and the coupling term in (2) with

$$\lim_{M \rightarrow \infty} \frac{\varepsilon}{M} \sum_{j=1}^M h(x_i, x_j).$$

The limit above might not exist. Let's assume that for the initial condition $(x_i(0))_{i=1}^\infty \in \mathbb{T}^\mathbb{N}$ there is a probability measure μ such that

$$\lim_{M \rightarrow \infty} \frac{1}{M} \sum_{i=1}^M \delta_{x_i(0)} = \mu \quad (3)$$

where the limit is with respect to the weak convergence of measures [†]. This assumption

*For example, when $\varepsilon = 0$, Φ_0 equals the identity and the evolution of any coordinate i equals the uncoupled dynamics f .

[†]Notice that the existence of the limit and value of the limit in (3) depends, in general, on the ordering of the coordinates.

Given a sequence of i.i.d. random variables $\{X_i\}_{i \in \mathbb{N}}$ such that the law of X_i is given by μ , then almost every realization of the sequence $\{X_i\}_{i \in \mathbb{N}}$ satisfies (3). This is a consequence of the strong Law of Large Numbers and the separability of continuous functions from S^1 to \mathbb{R} with respect to the topology induced by the sup norm.

tion ensures that for any continuous h ,

$$\lim_{M \rightarrow \infty} \frac{\varepsilon}{M} \sum_{j=1}^M h(x_i(0), x_j(0)) = \varepsilon \int h(x_i(0), y) d\mu(y),$$

and the time evolution of each coordinate can be written as:

$$x_i(1) = f \circ \Phi_{\varepsilon, \mu}(x_i(0)) \quad \text{where} \quad \Phi_{\varepsilon, \mu}(x) = x + \varepsilon \int h(x, y) d\mu(y).$$

Calling f_* and $(\Phi_{\varepsilon, \mu})_*$ the push-forwards of f and $\Phi_{\varepsilon, \mu}$,

$$\lim_{M \rightarrow \infty} \frac{1}{M} \sum_{i=1}^M \delta_{x_i(1)} = f_* (\Phi_{\varepsilon, \mu})_*(\mu)$$

where the convergence is weak and follows from the continuity of the push forwards (which only requires the continuity of f and $\Phi_{\varepsilon, \mu}$). In case μ has a density φ with respect to a suitable reference measure (Lebesgue, for instance) it is sufficient to study the evolution of the density. By the above argument, this is given by a transfer operator that depends on the probability measure, thus φ itself (for a precise definition, see (7)). This nonlinear application is called a *self-consistent transfer operator*. *This will be the object under study in this paper.*

Fernandez [13] studied the finite setting in the case of $h(x, y) = g(x - y)$ with $g : \mathbb{R} \rightarrow \mathbb{R}$ defined to be a specific discontinuous, but piecewise affine function. Studying the thermodynamic limit of his setup in case of ε sufficiently small, Bálint and Sélley [33] first showed that when $f(x) = 2x \bmod 1$, the constant function is the only smooth fixed density. Expanding these results, in [10], they considered a general uniformly expanding smooth circle map f , and showed that the invariant density is unique and a Lipschitz continuous function of the coupling strength, with higher regularity being obstructed by the discontinuity of the function h .

In this paper we present an approach based on convex cones of functions to study self-consistent operators for general uniformly expanding f and general h satisfying certain smoothness conditions. We show that in the weak coupling regime (ε small), the self-consistent operator has a unique fixed smooth density that depends in a differentiable way on the coupling strength. Furthermore, we compute a linear response formula for the density.

We believe that assuming stronger smoothness conditions on h and f , higher smoothness of the invariant density in terms of the coupling could be also achieved and a quadratic response formula could be deduced, reminiscent of the formula obtained in [16]. Another, more difficult, open question is whether one could apply ideas similar to the one in this paper to treat the case where h is not smooth. This would require the definition of some suitable Banach spaces and cones playing analogous roles to the spaces and cones of smooth functions that we consider here.

1.1 Setup

Function spaces and norms Let S^1 denote the unit circle and denote the Lebesgue (Haar) measure on S^1 by λ . Throughout this paper, we consider the representation of S^1 as the unit interval $[0, 1]$ with 0 and 1 identified. We denote by

$L^1(S^1, \mathbb{R})$ the real valued functions on S^1 for which

$$\int_{S^1} |f| d\lambda < \infty$$

holds. Let $C^0(S^1, S^1)$ denote the space of continuous mappings of S^1 , endowed with the norm

$$\|f\|_{C^0} = \max_{x \in S^1} |f(x)|.$$

For integers $k \geq 1$, let $C^k(S^1, S^1)$ denote the space of k -times continuously differentiable mappings of S^1 , endowed with the norm

$$\|f\|_{C^k} = \max_{x \in S^1} |f(x)| + \sum_{\ell=1}^k \max_{x \in S^1} |f^{(\ell)}(x)|.$$

Uncoupled Dynamics Let $f \in C^k(S^1, S^1)$ (with $k > 1$ to be specified later) be a uniformly expanding circle map, i.e. there exists an $\omega > 1$ such that $|f'(x)| > \omega$ for all $x \in S^1$, and denote by P its transfer operator with respect to the Lebesgue measure. That is, $P : L^1(S^1, \mathbb{R}) \rightarrow L^1(S^1, \mathbb{R})$ is defined as

$$P\varphi(x) = \sum_{y \in f^{-1}(x)} \frac{\varphi(y)}{|f'(y)|}, \quad \varphi \in L^1(S^1, \mathbb{R}). \quad (4)$$

Mean-Field Coupling Given $\varepsilon_0 > 0$, $t \in [-\varepsilon_0, \varepsilon_0]$, $h \in C^k(S^1 \times S^1, \mathbb{R})$ (with $k > 1$ to be specified later), and $\psi \in L^1(S^1, \mathbb{R}_0^+)$ [‡] with $\int_{S^1} \psi = 1$, define $\Phi_{t,\psi} : S^1 \rightarrow S^1$

$$\Phi_{t,\psi}(x) = x + t \int_{S^1} h(x, y) \psi(y) dy.$$

We assume that $\varepsilon_0 > 0$ is sufficiently small so that $\Phi_{t,\psi}$ is a diffeomorphism such that $\Phi'_{t,\psi} > 0$, and we denote by $Q_{t,\psi}$ its transfer operator. More precisely, similarly to (4), the operator $Q_{t,\psi} : L^1(S^1, \mathbb{R}) \rightarrow L^1(S^1, \mathbb{R})$ is defined as

$$Q_{t,\psi}\varphi(x) = \frac{\varphi}{\Phi'_{t,\psi}} \circ \Phi_{t,\psi}^{-1}(x), \quad \varphi \in L^1(S^1, \mathbb{R}). \quad (5)$$

The evolution of the density of states ψ under the mean-field coupling is then given by the nonlinear application of Q_t defined as $Q_t(\psi) := Q_{t,\psi}\psi$. In the following we use the short-hand notation $A_\psi(x) = \int_{S^1} h(x, y) \psi(y) dy$, so we write $\Phi_{t,\psi}(x) = x + tA_\psi(x)$. Thus substituting $\varphi = \psi$ in (5) gives the following expression for Q_t :

$$Q_t\psi(x) = \frac{\psi}{1 + tA'_\psi} \circ \Phi_{t,\psi}^{-1}(x). \quad (6)$$

Coupled Dynamics Define $F_{t,\psi} := f \circ \Phi_{t,\psi}$, and let $\mathcal{L}_{t,\psi} := PQ_{t,\psi}$ denote its transfer operator. Given $t \in [-\varepsilon_0, \varepsilon_0]$, define the self-consistent transfer operator $\mathcal{L}_t : L^1(S^1, \mathbb{R}) \rightarrow L^1(S^1, \mathbb{R})$ as

$$\mathcal{L}_t := PQ_t \quad (7)$$

which describes the evolution of the all-to-all coupled system in the thermodynamic limit.

[‡] \mathbb{R}_0^+ is the set containing the positive real numbers together with 0.

Smooth densities Throughout the paper we are going to work with the restrictions of P, Q_t and \mathcal{L}_t to some appropriate spaces of smoothly differentiable functions. These restrictions are well-defined: on the one hand, $\varphi \in C^\ell$ implies $P\varphi \in C^\ell$ for $\ell < k$ supposing $f \in C^k(S^1, S^1)$. On the other, if $h \in C^{k'}(S^1 \times S^1, \mathbb{R})$ then $\Phi_{t,\psi} \in C^{k'}(S^1, S^1)$ for sufficiently small values of t , thus $Q_t\varphi \in C^\ell$ for $\ell < k'$.

By continuity, f is an N -fold covering map of S^1 for some $N \in \mathbb{N}$. We make a technical assumption on f ensuring that the expansion is large enough and the distortion is small enough, more precisely

$$N \left(\frac{\max |f''|}{\omega^3} + \frac{1}{\omega^2} \right) < 1. \quad (8)$$

Note that for example the map $f(x) = kx \bmod 1$ (where k is an integer) and moderate perturbations satisfy this condition. We remark that we could exclude this condition by working with higher iterates of the self-consistent transfer operator. However, this would complicate the calculations a great deal and we will assume (8) for the sake of clarity of the presentation. From now on, $K > 0$ will denote an unspecified positive number that depends on h and bounds of its derivatives only, while $\tilde{K} > 0$ is another unspecified constant depending also on f , unless otherwise specified.

1.2 Summary of the results and outline of the paper

The first theorem of this paper claims that for t sufficiently small, the system has a unique equilibrium density $\rho(t)$, and any sufficiently smooth initial state will evolve towards $\rho(t)$.

Theorem 1.1. *Assume $f \in C^2(S^1, S^1)$ and $h \in C^2(S^1 \times S^1, \mathbb{R})$. Then, there is $\varepsilon_0 > 0$ sufficiently small such that for any $t \in [-\varepsilon_0, \varepsilon_0]$, \mathcal{L}_t has a unique fixed non-negative probability density $\rho(t) \in \text{Lip}(S^1, \mathbb{R}^+)$. Furthermore, there exist constants $\tilde{K} > 0$ and $\Lambda < 1$ such that for any non-negative $\varphi \in \text{Lip}(S^1, \mathbb{R})$ with $\int \varphi = 1$ it holds that*

$$\|\mathcal{L}_t^n \varphi - \rho(t)\|_{C^0} \leq \tilde{K} \Lambda^n \quad \text{for all } n \in \mathbb{N}.$$

Remark 1.1. A very similar result was stated by [19, Theorem 4]. However the coupling mechanism there is different in the sense that an integral average is computed from the current system statistic which feeds into the dynamics as a constant. Our coupling is more complicated: it varies from point to point due to the x -dependence of the function h . However, the results and the methods that can be used to prove them are essentially the same. We view this statement as a preparatory result for the next theorem, our main result.

The second theorem claims that taking f and h sufficiently smooth, the dependence of $\rho(t)$ on t is C^1 in a strong sense.

Theorem 1.2. *Assume $f \in C^5(S^1, S^1)$ and $h \in C^5(S^1 \times S^1, \mathbb{R})$. Then*

- (i) *there is an $\varepsilon_1 > 0$, $\varepsilon_1 \leq \varepsilon_0$ such that $\rho(t) \in C^3(S^1, \mathbb{R})$ for all $t \in [-\varepsilon_1, \varepsilon_1]$;*
- (ii) *there is an $\varepsilon_2 > 0$, $\varepsilon_2 \leq \varepsilon_1$ such that $t \mapsto \rho(t)$ belongs to $C^1([-\varepsilon_2, \varepsilon_2], C^1(S^1, \mathbb{R}))$;*

(iii) there is an $\varepsilon_3 > 0$, $\varepsilon_3 \leq \varepsilon_2$ such that the following linear response formula holds for $|\hat{t}| < \varepsilon_3$:

$$\partial_t \rho(t)|_{t=\hat{t}} = -(1 - P_{\hat{t}} + \hat{t}PK_{\hat{t}})^{-1}PK_{\hat{t}}(\rho(\hat{t})),$$

where $P_{\hat{t}} = PQ_{\hat{t}, \rho(\hat{t})}$ and the operator $K_{\hat{t}}$ is given by (10).

Remark 1.2. Both theorems depend intimately on the fact that when the coupling strength is sufficiently weak, the operators satisfy uniform Lasota-Yorke type inequalities implying uniform spectral properties. However, standard perturbation theory arguments are insufficient to obtain the results above due to the nonlinear character of the operators \mathcal{L}_t . Getting around the nonlinearity of the operators is where the main novelty of our approach lies.

We break down the proof of the above theorems into several propositions listed below and proved in the following sections.

In Section 2 we study existence, uniqueness and regularity of the fixed point for the self-consistent transfer operator \mathcal{L}_t (7). To this end, we consider the cone of log-Lipschitz functions

$$\mathcal{V}_a := \left\{ \varphi : S^1 \rightarrow \mathbb{R}^+ : \frac{\varphi(x)}{\varphi(y)} \leq \exp(a|x - y|) \right\},$$

where $|x - y|$ denotes the Euclidean distance between x and y on the unit circle, and show that for large enough a , \mathcal{L}_t keeps \mathcal{V}_a invariant. Most importantly, we can also show that \mathcal{L}_t is a contraction with respect to the Hilbert metric θ_a (see (11) below) on \mathcal{V}_a .

Remark 1.3. Cones of functions and the Hilbert metric are common tools to prove uniqueness of the invariant density for some linear transfer operators [23]. The standard approach uses the fact that a linear operator mapping a cone with finite diameter strictly inside itself is a contraction in the Hilbert metric. Unlike the standard case, \mathcal{L}_t is a nonlinear operator, and one has to prove explicitly that \mathcal{L}_t is a contraction.

Proposition 1.1. *There is $a_0 > 0$ sufficiently large such that for $a > a_0$, there is $\varepsilon_0 > 0$ sufficiently small such that for $t \in [-\varepsilon_0, \varepsilon_0]$, \mathcal{L}_t keeps \mathcal{V}_a invariant. Furthermore, there are $a > 0$, and $\varepsilon_0 > 0$ small such that \mathcal{L}_t is a contraction on $(\mathcal{V}_a, \theta_a)$ for any $t \in [-\varepsilon_0, \varepsilon_0]$.*

The proof of Theorem 1.1 then follows by standard arguments showing that under repeated application of \mathcal{L}_t , all the elements in the cone \mathcal{V}_a converge exponentially fast to the fixed point $\rho(t)$ (in the C^0 norm).

To show that the C^5 -smoothness of f and h implies that $\rho(t) \in C^3(S^1, \mathbb{R})$, define the following set for the quadruple $\mathbf{C} = (C_1, C_2, C_3, C_4) \in (\mathbb{R}^+)^4$:

$$\mathcal{C}_{\mathbf{C}} := \left\{ \varphi \in C^4(S^1, \mathbb{R}^+) : \int \varphi = 1, \left\| \frac{d^i \varphi}{dx^i} \right\|_{C^0} < C_i \ i = 1, \dots, 4 \right\}.$$

We show that

Proposition 1.2. *There is $\varepsilon_1 > 0$ and $C_1, C_2, C_3, C_4 > 0$ such that for $t \in [-\varepsilon_1, \varepsilon_1]$*

$$\mathcal{L}_t \mathcal{C}_{\mathbf{C}} \subset \mathcal{C}_{\mathbf{C}}.$$

The existence of a fixed point in C^3 is then implied by the fact that \mathcal{L}_t is continuous (Lemma 2.3 below), \mathcal{C}_C is relatively compact in C^3 , and the Schauder fixed-point theorem. Since one can find $\varphi \in \mathcal{C}_C \cap \mathcal{V}_a$ (for example take $\varphi = 1$), the fixed point $\rho \in \mathcal{V}_a$, whose existence was argued above, must belong to the closure of \mathcal{C}_C and therefore is in C^3 , that will prove point (i) of Theorem 1.2.

In Section 3 we show that, when f and h are C^5 , the curve of fixed points $t \mapsto \rho(t)$ found above is C^1 from $[-\varepsilon_2, \varepsilon_2]$ to $C^1(S^1, \mathbb{R})$ for some $\varepsilon_2 > 0$.

Remark 1.4. In the case of a one-parameter family $\{f_t\}_t$ of sufficiently smooth uniformly expanding maps, the strong differentiability in t follows from the spectral properties of the associated transfer operators and the smoothness of $t \mapsto f_t$ [4]. Here, since the one-parameter family depends on the fixed density itself, i.e. $\rho(t)$ is the fixed density of $f_t = F_{t, \rho(t)}$ (self-consistent relation), the usual arguments do not apply since there is no a priori knowledge of the regularity of $t \mapsto f_t$ that has to be established first.

This is achieved by showing that the set of curves

$$\hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2} := \left\{ \gamma : [-\varepsilon_2, \varepsilon_2] \rightarrow \mathcal{C}_C \text{ s.t. } \begin{array}{l} (t, x) \mapsto \gamma(t)(x) \text{ is } C^2 \\ \sup_t \left\| \frac{d}{dt} \gamma(t) \right\|_{C^3} < K_1, \sup_t \left\| \frac{d^2}{dt^2} \gamma(t) \right\|_{C^2} < K_2 \end{array} \right\}$$

is invariant for suitably chosen K_1 and K_2 greater than zero under the mapping

$$(\mathcal{L}\gamma)(t) = \mathcal{L}_t(\gamma(t)) \text{ for all } t \in [-\varepsilon_2, \varepsilon_2]$$

naturally defined by the family of operators on a curve.

Proposition 1.3. *Choose $C_1, C_2, C_3, C_4 > 0$ so that Proposition 1.2 holds. Then there is $\varepsilon_2 > 0$ sufficiently small and $K_1, K_2 > 0$ such that*

$$\mathcal{L}\hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2} \subset \hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2}.$$

Since $\hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2}$ is relatively compact in $C^1([-\varepsilon_2, \varepsilon_2], C^1(S^1, \mathbb{R}))$, for any $\gamma \in \hat{\mathcal{C}}$, there is an increasing sequence $\{n_k\}_{k \in \mathbb{N}}$ and $\bar{\gamma} \in C^1([-\varepsilon_2, \varepsilon_2], C^2(S^1, \mathbb{R}))$ such that $\mathcal{L}^{n_k} \gamma \rightarrow \bar{\gamma}$, and since by Proposition (1.3) $\mathcal{L}_t^{n_k}(\gamma(t)) \rightarrow \rho(t)$ for every t , $\bar{\gamma}(t) = \rho(t)$. This proves point (ii) of Theorem 1.2.

In Section 4, we prove point (iii) of Theorem 1.2 by calculating a linear response formula. To ease the notation, we denote $P_t := \mathcal{L}_{t, \rho(t)}$.

Proposition 1.4. *We have the formula*

$$\partial_t \rho|_{t=0} = -(1 - P)^{-1} P(\rho(0) A_{\rho(0)})', \quad (9)$$

where ∂_t is understood as $\partial_t : C^1([-\varepsilon_3, \varepsilon_3], C^1(S^1, \mathbb{R})) \mapsto C^0([-\varepsilon_3, \varepsilon_3], C^1(S^1, \mathbb{R}))$. More generally, for any $\hat{t} \in (-\varepsilon_3, \varepsilon_3)$ where $\varepsilon_3 \leq \varepsilon_2$ as in Proposition 1.3,

$$\partial_t \rho|_{t=\hat{t}} = -(1 - P_{\hat{t}} + \hat{t} P K_{\hat{t}})^{-1} P(K_{\hat{t}}(\rho(\hat{t})))$$

where $P_{\hat{t}} = P Q_{\hat{t}, \rho(\hat{t})}$ and

$$K_{\hat{t}}(g) = \left(\left(\frac{\rho(\hat{t})}{\Phi'_{\hat{t}, \rho(\hat{t})}} A_g \right) \circ \Phi_{\hat{t}, \rho(\hat{t})}^{-1} \right)', \quad g \in C^1(S^1, \mathbb{R}). \quad (10)$$

The formula in (9) is reminiscent of the one obtained in [4, Theorem 2.2] for perturbations of expanding circle maps. This is not surprising, since for t approaching 0 the operator is very close to linear.

At the end of the paper, Appendix A gathers some elementary formulae used throughout our calculations.

2 Invariant Densities for the Self-Consistent Transfer Operators

We now show that the fixed point found above is unique using an invariant cone argument and the Hilbert metric. We remind the reader of our definition

$$\mathcal{V}_a := \left\{ \varphi : S^1 \rightarrow \mathbb{R}^+ : \frac{\varphi(x)}{\varphi(y)} \leq \exp(a|x-y|) \right\}.$$

and of the Hilbert metric

$$\theta_a(\varphi, \psi) = \log \frac{1}{\alpha_a(\varphi, \psi)\alpha_a(\psi, \varphi)}$$

where

$$\alpha_a(\psi, \varphi) = \inf \left\{ \frac{\varphi(x)}{\psi(x)}, \frac{e^{a|x-y|}\varphi(x) - \varphi(y)}{e^{a|x-y|}\psi(x) - \psi(y)} : x \neq y \right\}. \quad (11)$$

It is a known important fact that P is a contraction with respect to this metric (see e.g. [23, 34]). In the following we show that for suitable values of t and $a > 0$, the operator \mathcal{L}_t is also a contraction. The lemmas below gather a few useful inequalities.

Lemma 2.1. *There is $\varepsilon_0 > 0$ sufficiently small such that for every $a > 0$, every $t \in (-\varepsilon_0, \varepsilon_0)$, and $\varphi, \psi \in C^0$ with $\int \varphi = \int \psi = 1$,*

$$\left| \Phi_{t,\psi}^{-1}(x) - \Phi_{t,\varphi}^{-1}(x) \right| \leq K|t|\|\psi - \varphi\|_{C^0} \quad (12)$$

Proof. Notice that

$$\begin{aligned} |\Phi_{t,\psi}(x) - \Phi_{t,\varphi}(x)| &= |t| \left| \int \partial_1 h(x, y)(\psi(y) - \varphi(y)) \right| \\ &\leq |t|K\|\psi - \varphi\|_{C^0} \end{aligned}$$

By the mean value theorem, for any x

$$\begin{aligned} \min |\Phi'_{t,\psi}| &\leq \frac{|\Phi_{t,\psi}(\Phi_{t,\psi}^{-1}(x)) - \Phi_{t,\psi}(\Phi_{t,\varphi}^{-1}(x))|}{|\Phi_{t,\varphi}^{-1}(x) - \Phi_{t,\psi}^{-1}(x)|} \\ &= \frac{|\Phi_{t,\varphi}(\Phi_{t,\varphi}^{-1}(x)) - \Phi_{t,\psi}(\Phi_{t,\varphi}^{-1}(x))|}{|\Phi_{t,\varphi}^{-1}(x) - \Phi_{t,\psi}^{-1}(x)|} \\ &\leq \frac{K|t|\|\psi - \varphi\|_{C^0}}{|\Phi_{t,\varphi}^{-1}(x) - \Phi_{t,\psi}^{-1}(x)|}. \end{aligned}$$

Since by assumption $\min |\Phi'_{t,\varphi}| \geq 1 - K|t| > 0$ for any φ , (12) follows. \square

Lemma 2.2. *There is $\varepsilon_0 > 0$ sufficiently small such that for every $a > 0$ there is $K_a > 0$ such that for every $t \in (-\varepsilon_0, \varepsilon_0)$ and $\varphi, \psi \in \mathcal{V}_a$ with $\int \varphi = \int \psi = 1$, the following holds:*

$$\left| \Phi_{t,\psi}^{-1}(x) - \Phi_{t,\varphi}^{-1}(x) \right| \leq K_a |t| \max \left| 1 - \frac{\varphi}{\psi} \right| \quad (13)$$

$$\left| \Phi_{t,\varphi}^{-1}(x) - \Phi_{t,\varphi}^{-1}(y) \right| \leq (1 + K|t|)|x - y| \quad (14)$$

$$\max \left| 1 - \frac{\varphi}{\psi} \right| \geq K_a \log \alpha(\varphi, \psi) \quad (15)$$

Proof. The proof of (13) follows along the same lines of the proof of Lemma 2.1 after noticing that

$$\begin{aligned} \left| \int \partial_1 h(x, y)(\psi(y) - \varphi(y)) \right| &\leq \max \psi \cdot K \cdot \max \left| 1 - \frac{\varphi}{\psi} \right| \\ &\leq e^{\frac{a}{2}} K \cdot \max \left| 1 - \frac{\varphi}{\psi} \right|. \end{aligned}$$

Since $\psi \in \mathcal{V}_a$ and has integral equal to one, $\max \psi \leq e^{\frac{a}{2}}$. To prove (14), notice that for t sufficiently small $|(\Phi_{t,\varphi}^{-1})'| < 1 + Kt$. To prove (15), pick $x_0 \in \mathbb{T}$ such that $\max |1 - \frac{\varphi}{\psi}| = 1 - \frac{\varphi}{\psi}(x_0)$. Then the inequality is implied by

$$\log \alpha(\varphi, \psi) \leq \log \left(\frac{\varphi}{\psi}(x_0) \right) = \log \left[1 - \left(1 - \frac{\varphi}{\psi}(x_0) \right) \right] \leq K_a^{-1} \left(1 - \frac{\varphi}{\psi}(x_0) \right)$$

as φ/ψ is lower bounded by $e^{-\frac{a}{2}}$. Analogous conclusion holds if $\max_{x \in \mathbb{T}} |1 - \frac{\varphi}{\psi}(x)| = \frac{\varphi}{\psi}(x_0) - 1$ for some $x_0 \in \mathbb{T}$. \square

We are now ready to prove Proposition 1.1.

Proof of Proposition 1.1. Step 1: There is $\lambda < 1$ such that $\mathcal{L}_t \mathcal{V}_a \subset \mathcal{V}_{\lambda a}$

First of all, we show that there is K such that $Q_t \mathcal{V}_a \subset \mathcal{V}_{a+K(a+1)|t|}$ for t sufficiently small and $a > 0$. In fact, we show more. Pick $\varepsilon_0 > 0$ such that $1 - |t| \max |\partial_1 h| > 0$ for $t \in [-\varepsilon_0, \varepsilon_0]$. For any $\varphi, \psi \in \mathcal{V}_a$ we have

$$\frac{\varphi \circ \Phi_{t,\psi}^{-1}(x)}{\varphi \circ \Phi_{t,\psi}^{-1}(y)} \leq e^{\alpha(1+K|t|)|x-y|} \quad \text{and} \quad \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(y)}{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(x)} \leq e^{K|t||x-y|},$$

where the bound on the first expression follows from $\varphi \in \mathcal{V}_a$ and (14), while that on the second expression follows from the smoothness of h (in particular K depends on the norms of the derivatives of h). Thus

$$\frac{Q_{t,\psi} \varphi(x)}{Q_{t,\psi} \varphi(y)} = \frac{\varphi \circ \Phi_{t,\psi}^{-1}(x)}{\varphi \circ \Phi_{t,\psi}^{-1}(y)} \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(y)}{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(x)} \leq e^{(a+(a+1)K|t|)|x-y|}.$$

Analogously, it is easy to check (by noticing that $\left| \frac{f'(x)}{f'(y)} \right| \leq e^{\frac{1}{\omega} \max |f''||x-y|}$) that if $\varphi \in \mathcal{V}_a$ then $P\varphi \in \mathcal{V}_{a/\omega + \max |f''|/\omega^2}$.

Picking

$$a_0 := \frac{\max |f''|}{\omega(\omega - 1)}$$

one has that for any $a > a_0$,

$$a/\omega + \max |f''|/\omega^2 < a.$$

Therefore, for any $a > a_0$ there is $\varepsilon_0 > 0$ and a_1, a_2 satisfying

$$a < a + (a + 1)K|\varepsilon_0| =: a_1 < a_2 < \omega \left(a - \frac{\max |f''|}{\omega^2} \right).$$

With this choice of a_1 and a_2 we have: $Q_t \mathcal{V}_a \subset \mathcal{V}_{a_1} \subset \mathcal{V}_{a_2}$ for $t \in (-\varepsilon_0, \varepsilon_0)$, and $P \mathcal{V}_{a_2} \subset \mathcal{V} \left(\frac{a_2}{\omega} + \frac{\max |f''|}{\omega^2} \right) \subset \mathcal{V}_a$. More precisely, $\mathcal{L}_t \mathcal{V}_a \subset \mathcal{V}_{\lambda a}$ with $\lambda < 1$ depending on a . [§]

From now on, $a > a_0$ is considered to be fixed.

Step 2: for t small, $\mathcal{L}_t : \mathcal{V}_a \rightarrow \mathcal{V}_a$ is a contraction with respect to θ_a

Let us start by showing how assuming

$$\theta_{a_2}(Q_t \varphi, Q_t \psi) \leq [1 + \mathcal{O}(|t|)] \theta_a(\varphi, \psi) \quad (16)$$

yields the proof of the claim in Step 2. Recall that with the choice of a and a_2 as in Step 1, $P : (\mathcal{V}_{a_2}, \theta_{a_2}) \rightarrow (\mathcal{V}_a, \theta_a)$ is a contraction as implied by the linearity of P and the fact that $\mathcal{V}_a \subset \mathcal{V}_{a_2}$. Call $\sigma < 1$ the contraction rate, i.e. $\theta_a(P\varphi, P\psi) \leq \sigma \theta_{a_2}(\varphi, \psi)$ for any $\varphi, \psi \in \mathcal{V}_{a_2}$. This and (16) imply that $\theta_a(\mathcal{L}_t \varphi, \mathcal{L}_t \psi) \leq \sigma [1 + \mathcal{O}(|t|)] \theta_a(\varphi, \psi)$. Picking $|t|$ sufficiently small, equation (16) implies that \mathcal{L}_t is also a contraction.

Now we proceed with the proof of (16). Notice that by the triangle inequality

$$\theta_{a_2}(Q_t \varphi, Q_t \psi) \leq \theta_{a_2}(Q_{t,\varphi} \varphi, Q_{t,\varphi} \psi) + \theta_{a_2}(Q_{t,\varphi} \psi, Q_{t,\psi} \psi).$$

For the first term, the operator $Q_{t,\varphi}$ is linear and maps the cone \mathcal{V}_a to \mathcal{V}_{a_1} . \mathcal{V}_{a_1} is strictly inside the cone \mathcal{V}_{a_2} . Since $Q_{t,\varphi}(\mathcal{V}_a)$ is strictly contained in \mathcal{V}_{a_2} and its diameter in $(\mathcal{V}_{a_2}, \theta_{a_2})$ is finite [34, Proposition 2.5], the mapping $Q_{t,\psi} : (\mathcal{V}_a, \theta_a) \rightarrow (\mathcal{V}_{a_2}, \theta_{a_2})$ is a contraction [34]. Therefore,

$$\theta_{a_2}(Q_{t,\varphi} \varphi, Q_{t,\varphi} \psi) \leq \theta_a(\varphi, \psi).$$

For the second term, computations are lengthier and occupy the rest of the proof. Recall that

$$\theta_{a_2}(Q_{t,\varphi} \psi, Q_{t,\psi} \psi) = \log \frac{1}{\alpha_{a_2}(\varphi, \psi) \alpha_{a_2}(\psi, \varphi)}$$

where α is the infimum of two expressions as defined in (11). For the first expression in the definition of α_{a_2} :

$$\begin{aligned} \frac{Q_{t,\varphi} \psi(x)}{Q_{t,\psi} \psi(x)} &= \frac{\psi \circ \Phi_{t,\varphi}^{-1}(x)}{\psi \circ \Phi_{t,\psi}^{-1}(x)} \cdot \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(x)}{(1 + t\partial_x A_\varphi) \circ \Phi_{t,\varphi}^{-1}(x)} = \\ &= \frac{\psi \circ \Phi_{t,\varphi}^{-1}(x)}{\psi \circ \Phi_{t,\psi}^{-1}(x)} \cdot \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(x)}{(1 + t\partial_x A_\psi) \circ \Phi_{t,\varphi}^{-1}(x)} \cdot \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\varphi}^{-1}(x)}{(1 + t\partial_x A_\varphi) \circ \Phi_{t,\varphi}^{-1}(x)} \end{aligned}$$

[§]If \mathcal{L}_t were a linear operator, this would be enough to conclude that \mathcal{L}_t is a contraction with respect to the Hilbert metric. Since it is nonlinear, we need to show it directly by estimating the Hilbert distance.

The bound on the first two factors follows from $\psi, \varphi \in \mathcal{V}_a$, $(1 + tA_\psi) \in \mathcal{V}_{K|t|}$ (again K depends on bounds on h and its derivatives), and from inequality (13) that give

$$\frac{\psi \circ \Phi_{t,\varphi}^{-1}(x)}{\psi \circ \Phi_{t,\psi}^{-1}(x)} \cdot \frac{(1 + t\partial_x A_\psi) \circ \Phi_{t,\psi}^{-1}(x)}{(1 + t\partial_x A_\varphi) \circ \Phi_{t,\varphi}^{-1}(x)} \geq e^{-(a+K|t|)|\Phi_{t,\varphi}^{-1}(x) - \Phi_{t,\psi}^{-1}(x)|} \geq e^{-aK_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|}.$$

For the third factor

$$\begin{aligned} \frac{(1 + t\partial_x A_\psi)(x)}{(1 + t\partial_x A_\varphi)(x)} &= \frac{1 + t \int \partial_1 h(x, y) \psi(y) dy}{1 + t \int \partial_1 h(x, y) \varphi(y) dy} \\ &\geq \frac{1 + t \int \partial_1 h(x, y) \varphi(y) dy + t \left(\min\left\{\frac{\psi}{\varphi}\right\} - 1 \right) \int |\partial_1 h(x, y)| \varphi(y) dy}{1 + t \int \partial_1 h(x, y) \varphi(y) dy} \\ &\geq 1 - K|t| \max\left|\frac{\psi}{\varphi} - 1\right| \\ &\geq \alpha_a(\varphi, \psi)^{K|t|}. \end{aligned} \tag{17}$$

Putting all the estimates together and using (15) one obtains

$$\frac{Q_{t,\varphi}\psi(x)}{Q_{t,\psi}\psi(x)} \geq \alpha_a(\varphi, \psi)^{K_a|t|}.$$

where $K_a > 0$ indicates a constant depending on h and on a only.

For the second expression in the definition of α we need to estimate

$$\frac{e^{a_2|x-y|} Q_{t,\varphi}\psi(x) - Q_{t,\varphi}\psi(y)}{e^{a_2|x-y|} Q_{t,\psi}\psi(x) - Q_{t,\psi}\psi(y)} = \frac{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\varphi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\varphi}^{-1}(y)}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\psi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\psi)} \circ \Phi_{t,\psi}^{-1}(y)} = A \cdot B$$

where

$$\begin{aligned} A &:= \frac{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\varphi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\varphi}^{-1}(y)}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y)} \\ B &:= \frac{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y)}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\psi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\psi)} \circ \Phi_{t,\psi}^{-1}(y)} \end{aligned}$$

$$\begin{aligned} A &\geq \frac{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) e^{-a_1|\Phi_{t,\psi}^{-1}(x) - \Phi_{t,\varphi}^{-1}(x)|} - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y) e^{a_1|\Phi_{t,\psi}^{-1}(y) - \Phi_{t,\varphi}^{-1}(y)|}}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y)} \\ &\geq \frac{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) e^{-K'_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|} - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y) e^{K'_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|}}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y)} \end{aligned} \tag{18}$$

$$\geq e^{K'_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|} \left(1 + \frac{e^{-2K'_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|} - 1}{e^{a_2|x-y|} \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(x) - \frac{\psi}{(1+tA_\varphi)} \circ \Phi_{t,\psi}^{-1}(y)} \right) \tag{19}$$

$$\geq e^{K'_a|t| \max\left|1 - \frac{\varphi}{\psi}\right|} \geq \alpha_a(\varphi, \psi)^{K_a|t|} \tag{20}$$

where (18) follows from (13), (20) follows from the fact that the denominator in (19) is bounded since $\frac{\varphi}{(1+tA\varphi)} \circ \Phi_{t,\psi}^{-1}$ belongs to \mathcal{V}_{a_1} with $a_1 < a_2$, and K_a , K'_a and K''_a are constants depending on a . The factor B can be bounded in a similar way:

$$B \geq \alpha_a(\varphi, \psi)^{K_a|t|}.$$

Putting together all the previous estimates one obtains

$$\alpha_{a_2}(Q_{t,\varphi}\psi, Q_{t,\psi}\psi) \geq \alpha_a(\varphi, \psi)^{2K_a|t|},$$

and analogously

$$\alpha_{a_2}(Q_{t,\psi}\psi, Q_{t,\varphi}\psi) \geq \alpha_a(\psi, \varphi)^{2K_a|t|}$$

which implies

$$\theta_{a_2}(Q_{t,\varphi}\psi, Q_{t,\psi}\psi) = \log \frac{1}{\alpha_{a_2}(Q_{t,\varphi}\psi, Q_{t,\psi}\psi)\alpha_{a_2}(Q_{t,\psi}\psi, Q_{t,\varphi}\psi)} \leq 2K_a|t|\theta_a(\varphi, \psi).$$

□

We are now ready to conclude the proof of Theorem 1.1

Proof of Theorem 1.1. The proof can be finished by standard arguments. The existence of an invariant density $\rho(t)$ in $\text{Lip}(S^1, \mathbb{R}^+)$ follows for example by the argument presented in [34, Section 2.3] and the exponential convergence of $\mathcal{L}_t^n \varphi$ to $\rho(t)$ in the C^0 norm, hence uniqueness of $\rho(t)$ can be argued as in [34, Section 2.4]. □

From now on let's assume that $f \in C^5(S^1, S^1)$ and $h \in C^5(S^1 \times S^1, \mathbb{R})$, and prove point (i) of Theorem 1.2, i.e. we show that in this case $\rho(t) \in C^3(S^1, \mathbb{R}^+)$. Recall that $\|\cdot\|_{C^0}$ denotes the supremum norm on $C^0(S^1, \mathbb{R})$, and

$$\mathcal{C}_C := \left\{ \varphi \in C^4(S^1, \mathbb{R}^+) : \int \varphi = 1, \left\| \frac{d^i \varphi}{dx^i} \right\|_{C^0} < C_i \ i = 1, \dots, 4 \right\}.$$

Proof of Proposition 1.2. Let $t \in [-\varepsilon_0, \varepsilon_0]$ and $\varphi \in C^4(S^1, \mathbb{R}^+)$. We remind the reader that $F_{t,\varphi}$ is defined as $F_{t,\varphi} = f \circ \Phi_{t,\varphi}$. By (43)–(45) of the Appendix and some elementary bounds, we obtain

$$|F'_{t,\varphi}| \geq \omega(1 - K_0|t|), \tag{21}$$

$$|F''_{t,\varphi}| \leq c((1 + K_0|t|)^2 + K_1)|t|, \tag{22}$$

$$|F'''_{t,\varphi}| \leq c((1 + K_0|t|)^3 + (1 + K_0|t|)K_1|t| + K_2|t|), \tag{23}$$

where $K_i = \max |h^{(i)}|$ and $c = \max\{\max |f'''|, 3 \max |f''|, \max |f'|\}$.

We now choose the constants C_i , $i = 1, \dots, 4$, starting with the choice of C_1 . We denote the inverse branches of $F_{t,\varphi}$ by $F_{t,\varphi,i}^{-1}$, $i = 1, \dots, N$.

$$\begin{aligned} \|(\mathcal{L}_t \varphi)'\|_{C^0} &= \left\| \left(\sum_{i=1}^N \frac{\varphi}{F'_{t,\varphi}} \circ F_{t,\varphi,i}^{-1} \right)' \right\|_{C^0} = \left\| \sum_{i=1}^N \left(\frac{\varphi'}{(F'_{t,\varphi})^2} - \frac{\varphi F''_{t,\varphi}}{(F'_{t,\varphi})^3} \right) \circ F_{t,\varphi,i}^{-1} \right\|_{C^0} \\ &\leq N \left\| \frac{\varphi'}{(F'_{t,\varphi})^2} - \frac{\varphi F''_{t,\varphi}}{(F'_{t,\varphi})^3} \right\|_{C^0} \\ &\leq N \left(\frac{1}{\omega^2(1 - K_0t)^2} \|\varphi'\|_{C^0} + \frac{\max |f''|(1 + K_0t)^2 + cK_1t}{\omega^3(1 - K_0t)^3} \|\varphi\|_{C^0} \right) \end{aligned}$$

by the bounds (21)–(22). Since φ is continuous and has unit integral, we have $\|\varphi\|_{C^0} \leq 1 + \|\varphi'\|_{C^0}$ giving us

$$\begin{aligned} \|(\mathcal{L}_t\varphi)'\|_{C^0} &\leq N \left(\frac{1}{\omega^2(1-K_0t)^2} + \frac{\max|f''|(1+Kt)^2 + cK_1t}{\omega^3(1-K_0t)^3} \right) \|\varphi'\|_{C^0} \\ &\quad + \frac{N(\max|f''|(1+K_0t)^2 + cK_1t)}{\omega^3(1-K_0t)^3} \\ &=: \sigma_1\|\varphi'\|_{C^0} + R_1. \end{aligned} \tag{24}$$

By the assumption (8) on f , we can choose t small enough (by potentially decreasing ε_0) such that $\sigma_1 < 1$. Choose C_1 such that

$$\frac{R_1}{1-\sigma_1} \leq C_1,$$

and then $\|\varphi'\|_{C^0} \leq C_1$ implies $\|(\mathcal{L}_t\varphi)'\|_{C^0} \leq C_1$.

We move on to the choice of C_2 .

$$\begin{aligned} \|(\mathcal{L}_t\varphi)''\|_{C^0} &= \left\| \left(\sum_{i=1}^N \frac{\varphi}{F'_{t,\varphi}} \circ F_{t,\varphi}^{-1} \right)'' \right\|_{C^0} \\ &\leq N \left\| \frac{\varphi''}{(F'_{t,\varphi})^3} - 3 \frac{\varphi' F''_{t,\varphi}}{(F'_{t,\varphi})^4} - \frac{\varphi F'''_{t,\varphi}}{(F'_{t,\varphi})^4} + 3 \frac{\varphi (F''_{t,\varphi})^2}{(F'_{t,\varphi})^5} \right\|_{C^0} \\ &\leq N \frac{1}{\omega^3(1-K_0t)^3} \|\varphi''\|_{C^0} + N \left\| 3 \frac{F''_{t,\varphi}}{(F'_{t,\varphi})^4} + \frac{F'''_{t,\varphi}}{(F'_{t,\varphi})^4} - 3 \frac{(F''_{t,\varphi})^2}{(F'_{t,\varphi})^5} \right\|_{C^0} \|\varphi'\|_{C^0} \\ &\quad + N \left\| \frac{F'''_{t,\varphi}}{(F'_{t,\varphi})^4} - 3 \frac{(F''_{t,\varphi})^2}{(F'_{t,\varphi})^5} \right\|_{C^0} \\ &=: \sigma_2\|\varphi''\|_{C^0} + R_2^{(1)}\|\varphi'\|_{C^0} + R_2^{(2)}. \end{aligned} \tag{25}$$

For sufficiently small t we have $\sigma_2 < 1$ by assumption (8) and by (21)–(23) $R_2^{(1)}$, $R_2^{(2)}$ are bounded uniformly in φ and in t if t is in some fixed neighborhood of zero, so choosing

$$\frac{R_2^{(1)}C_1 + R_2^{(2)}}{1-\sigma_2} \leq C_2$$

will be sufficient.

Similar calculations can be done for the third and fourth derivative. We can conclude that

$$\begin{aligned} \|(\mathcal{L}_t\varphi)'''\|_{C^0} &\leq \sigma_3\|\varphi'''\|_{C^0} + R_3^{(1)}\|\varphi''\|_{C^0} + R_3^{(2)}\|\varphi'\|_{C^0} + R_3^{(3)} \\ \|(\mathcal{L}_t\varphi)''''\|_{C^0} &\leq \sigma_4\|\varphi''''\|_{C^0} + R_4^{(1)}\|\varphi'''\|_{C^0} + R_4^{(2)}\|\varphi''\|_{C^0} + R_4^{(3)}\|\varphi'\|_{C^0} + R_4^{(4)} \end{aligned}$$

where $\sigma_3, \sigma_4 < 1$ for sufficiently small t and $R_i^{(j)}$ only depend on the derivatives of f and h and on t (such that they are bounded uniformly in t if t is in some fixed

neighborhood of zero). Choosing

$$\begin{aligned} \frac{R_3^{(1)}C_2 + R_3^{(2)}C_1 + R_3^{(3)}}{1 - \sigma_3} &\leq C_3 \\ \frac{R_4^{(1)}C_3 + R_4^{(2)}C_2 + R_4^{(3)}C_1 + R_4^{(4)}}{1 - \sigma_4} &\leq C_4 \end{aligned}$$

will be sufficient for our purposes. \square

Lemma 2.3. *If $f \in C^5(S^1, S^1)$ and $h \in C^5(S^1 \times S^1, \mathbb{R})$, then for $1 \leq k \leq 3$, $Q_t : \mathcal{C}_C \rightarrow \mathcal{C}_C$ is continuous with respect to the C^k distance.*

Proof. Fix $k \in \{1, 2, 3\}$. For any $\varphi, \psi \in C^k(S^1, \mathbb{R})$, using expression (6) for Q_t one obtains

$$\begin{aligned} Q_t\varphi - Q_t\psi &= \frac{\varphi - \psi}{1 + tA'_\varphi} \circ \Phi_{t,\varphi}^{-1}(x) + \\ &+ \left[\frac{\psi}{1 + tA'_\varphi} - \frac{\psi}{1 + tA'_\psi} \right] \circ \Phi_{t,\varphi}^{-1}(x) + \\ &+ \frac{\psi}{1 + tA'_\psi} \circ \Phi_{t,\varphi}^{-1}(x) - \frac{\psi}{1 + tA'_\psi} \circ \Phi_{t,\psi}^{-1}(x). \end{aligned}$$

For t sufficiently small $\frac{1}{1+tA'_\varphi}$ and $\Phi_{t,\varphi}^{-1}$ belong to C^4 . Then it follows that

$$\begin{aligned} \left\| \frac{\varphi - \psi}{1 + tA'_\varphi} \circ \Phi_{t,\varphi}^{-1}(x) \right\|_{C^k} &\leq M \left\| \frac{1}{1 + tA'_\varphi} \circ \Phi_{t,\varphi}^{-1}(x) \right\|_{C^k} \|(\varphi - \psi) \circ \Phi_{t,\varphi}^{-1}(x)\|_{C^k} \\ &\leq M' \left\| \frac{1}{1 + tA'_\varphi} \right\|_{C^k} \left(\max_{\ell=1,\dots,k} \|\Phi_{t,\varphi}^{-1}\|_{C^k}^\ell \right)^2 \|\varphi - \psi\|_{C^k} \\ &\leq \tilde{K} \|\varphi - \psi\|_{C^k}, \end{aligned}$$

where the first two inequalities hold by (42) and (47) of the Appendix, respectively. The third inequality can be proved by using (52) of the Appendix and some standard bounds on the derivatives of $\Phi_{t,\varphi}$. To bound the second term, notice that

$$\frac{1}{1 + tA'_\varphi} - \frac{1}{1 + tA'_\psi} = t \frac{A'_\psi - A'_\varphi}{(1 + tA'_\varphi)(1 + tA'_\psi)}$$

and $\|A'_\psi - A'_\varphi\|_{C^k} \leq K\|\psi - \varphi\|_{C^0}$. Then similarly to the calculation above,

$$\begin{aligned} &\left\| \left[\frac{\psi}{1 + tA'_\varphi} - \frac{\psi}{1 + tA'_\psi} \right] \circ \Phi_{t,\varphi}^{-1}(x) \right\|_{C^k} \\ &\leq M \left\| t \frac{\psi}{(1 + tA'_\varphi)(1 + tA'_\psi)} \circ \Phi_{t,\varphi}^{-1}(x) \right\|_{C^k} \|(\varphi - \psi) \circ \Phi_{t,\varphi}^{-1}(x)\|_{C^k} \\ &\leq M' \left\| t \frac{\psi}{(1 + tA'_\varphi)(1 + tA'_\psi)} \right\|_{C^k} \left(\max_{\ell=1,\dots,k} \|\Phi_{t,\varphi}^{-1}\|_{C^k}^\ell \right)^2 \|\varphi - \psi\|_{C^k} \\ &\leq \tilde{K}_C \|\varphi - \psi\|_{C^k}, \end{aligned}$$

where the last inequality also exploits that $\psi \in \mathcal{C}_{\mathbf{C}}$. Using (38)–(41) of the Appendix we can compute that there exists an explicit polynomial function p such that

$$\begin{aligned} & \left\| \frac{\psi}{1+tA'_\psi} \circ \Phi_{t,\varphi}^{-1}(x) - \frac{\psi}{1+tA'_\psi} \circ \Phi_{t,\psi}^{-1}(x) \right\|_{C^k} \\ & \leq p \left(\left\| \frac{\psi}{1+tA'_\psi} \right\|_{C^{k+1}}, \|\Phi_{t,\varphi}^{-1}\|_{C^k}, \|\Phi_{t,\psi}^{-1}\|_{C^k} \right) \|\varphi - \psi\|_{C^k}. \end{aligned}$$

Since $\frac{\psi}{1+tA'_\psi} \in C^4$, simple calculations as before give us

$$\left\| \frac{\psi}{1+tA'_\psi} \circ \Phi_{t,\varphi}^{-1}(x) - \frac{\psi}{1+tA'_\psi} \circ \Phi_{t,\psi}^{-1}(x) \right\|_{C^k} \leq \tilde{K}_{\mathbf{C}} \|\varphi - \psi\|_{C^k}.$$

By triangle inequality one can bound $\|Q_t\varphi - Q_t\psi\|_{C^k}$ and the thesis follows. \square

We can finally prove that $\rho(t) \in C^3(S^1, \mathbb{R})$.

Proof of Theorem 1.2 Point (i). $\mathcal{C}_{\mathbf{C}} \subset C^4(S^1, \mathbb{R})$ is convex and relatively compact in $C^3(S^1, \mathbb{R})$. The invariance of the set $\mathcal{C}_{\mathbf{C}}$ under the action of \mathcal{L}_t , as claimed in Proposition 1.2 and the continuity of \mathcal{L}_t restricted to $\mathcal{C}_{\mathbf{C}}$ as stated in Lemma 2.3 implies by the Schauder fixed point theorem that \mathcal{L}_t has a fixed density in $C^3(S^1, \mathbb{R})$. By Theorem 1.1 this density must be the unique Lipschitz density $\rho(t)$. \square

3 Smooth Dependence of the Invariant Densities on the Coupling Strength

Given $t \in [-\varepsilon_1, \varepsilon_1]$, the relation $t \mapsto \rho(t)$ defines a curve of fixed densities, that is, for each t the density $\rho(t)$ is such that $\mathcal{L}_t\rho(t) = \rho(t)$. Recall the definition of $\hat{\mathcal{C}}_{\varepsilon, K_1, K_2}$

$$\hat{\mathcal{C}}_{\varepsilon, K_1, K_2} := \left\{ \gamma : [-\varepsilon, \varepsilon] \rightarrow \mathcal{C}_{\mathbf{C}} \text{ s.t. } \begin{array}{l} (t,x) \mapsto \gamma(t)(x) \text{ is } C^2 \\ \sup_t \|\frac{d}{dt}\gamma(t)\|_{C^3} < K_1, \sup_t \|\frac{d^2}{dt^2}\gamma(t)\|_{C^2} < K_2 \end{array} \right\}$$

and the operator \mathcal{L}

$$(\mathcal{L}\gamma)(t) = \mathcal{L}_t(\gamma(t)) \text{ for all } t \in [-\varepsilon, \varepsilon]$$

for $\gamma \in \hat{\mathcal{C}}_{\varepsilon, K_1, K_2}$. The strategy consists of (i) showing that there are ε_2 , K_1 , and K_2 such that \mathcal{L} keeps $\hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2}$ invariant, and (ii) using Proposition 1.1 to show that this implies that $t \mapsto \rho(t)$ is C^1 .

Proof of Proposition 1.3. Let us start by noticing that the implicit function theorem implies that the function $g : [-\varepsilon_1, \varepsilon_1] \times S^1 \rightarrow S^1$, defined as $g(t, y) := (\Phi_{t,\varphi})^{-1}(y)$, is differentiable and C^5 . Furthermore, its C^5 norm can be uniformly bounded with respect to φ [¶]. Let $\varepsilon \leq \varepsilon_1$ and pick $\gamma \in \hat{\mathcal{C}}_{\varepsilon, K_1, K_2}$.

$$\frac{d}{dt}(\mathcal{L}\gamma)(t) = \left(\frac{d}{ds}\mathcal{L}_s \right)_{s=t} \gamma(t) + \lim_{\delta \rightarrow 0} \frac{\mathcal{L}_t\gamma(t+\delta) - \mathcal{L}_t\gamma(t)}{\delta}. \quad (26)$$

[¶]This is a consequence of the fact that the function to which we apply the implicit function theorem, $F(t, y, x) = \Phi_{t,\varphi}(x) - y$, has C^5 norm uniformly bounded varying φ .

For the first term in the right hand side of (26)

$$\begin{aligned}
\left(\frac{d}{ds}\mathcal{L}_s\right)_{s=t} \varphi &= P \left(\lim_{\delta \rightarrow 0} \frac{Q_{t+\delta} - Q_t}{\delta} \varphi \right) \\
&= P \frac{d}{ds} \left(\frac{\varphi}{\partial_x \Phi_{s,\varphi}} \circ \Phi_{s,\varphi}^{-1} \right)_{s=t} \\
&= P \left(\frac{\varphi'}{\partial_x \Phi_{t,\varphi}} \circ \Phi_{t,\varphi}^{-1} \partial_t \Phi_{t,\varphi}^{-1} - \frac{\varphi}{(\partial_x \Phi_{t,\varphi})^2} \circ \Phi_{t,\varphi}^{-1} \partial_t [\partial_x \Phi_{t,\varphi} \circ \Phi_{t,\varphi}^{-1}] \right) \\
&= P \left(\frac{\varphi'}{\partial_x \Phi_{t,\varphi}} \circ \Phi_{t,\varphi}^{-1} \partial_t \Phi_{t,\varphi}^{-1} \right. \\
&\quad \left. - \frac{\varphi}{(\partial_x \Phi_{t,\varphi})^2} \circ \Phi_{t,\varphi}^{-1} [\partial_t \partial_x \Phi_{t,\varphi} \circ \Phi_{t,\varphi}^{-1} + \partial_x^2 \Phi_{t,\varphi} \circ \Phi_{t,\varphi}^{-1} \partial_t \Phi_{t,\varphi}^{-1}] \right). \quad (27)
\end{aligned}$$

Therefore, putting $\varphi = \gamma(t)$ and recalling that $\gamma(t) \in \mathcal{C}_C$, for t sufficiently small all the functions above have bounded norm in $C^3(S^1, \mathbb{R})$, and therefore there is a constant \tilde{K} depending on f, h and C such that

$$\left\| \left(\frac{d}{ds} \mathcal{L}_s \right)_{s=t} \gamma(t) \right\|_{C^3} \leq \tilde{K}. \quad (28)$$

The second term in the RHS of (26) equals

$$\begin{aligned}
\Xi(t) &:= P \lim_{\delta \rightarrow 0} \frac{Q_t \gamma(t+\delta) - Q_t \gamma(t)}{\delta} \\
&= P \lim_{\delta \rightarrow 0} \frac{1}{\delta} \left[\frac{\gamma(t+\delta)}{1 + tA'_{\gamma(t+\delta)}} \circ \Phi_{t,\gamma(t+\delta)}^{-1} - \frac{\gamma(t)}{1 + tA'_{\gamma(t)}} \circ \Phi_{t,\gamma(t)}^{-1} \right] \\
&= P \lim_{\delta \rightarrow 0} \frac{\gamma(t+\delta) - \gamma(t)}{\delta(1 + A'_{\gamma(t+\delta)})} \circ \Phi_{t,\gamma(t+\delta)}^{-1} \\
&\quad + P \lim_{\delta \rightarrow 0} \frac{1}{\delta} \left[\gamma(t) \left(\frac{1}{1 + tA'_{\gamma(t+\delta)}} - \frac{1}{1 + tA'_{\gamma(t)}} \right) \right] \circ \Phi_{t,\gamma(t+\delta)}^{-1} \\
&\quad + P \lim_{\delta \rightarrow 0} \frac{1}{\delta} \left[\frac{\gamma(t)}{1 + tA'_{\gamma(t)}} \circ \Phi_{t,\gamma(t+\delta)}^{-1} - \frac{\gamma(t)}{1 + tA'_{\gamma(t)}} \circ \Phi_{t,\gamma(t)}^{-1} \right],
\end{aligned}$$

giving

$$\begin{aligned}
\Xi(t) &= \\
&= PQ_{t,\gamma(t)} \frac{d}{dt} \gamma(t) \quad (29)
\end{aligned}$$

$$+ P \left[\left(\gamma(t) \frac{d}{ds} \left(\frac{1}{1 + tA'_{\gamma(s)}} \right)_{s=t} \right) \circ \Phi_{t,\gamma(t)}^{-1} + \partial_x (Q_{t,\gamma(t)} \gamma(t)) \frac{d}{ds} \left(\Phi_{t,\gamma(s)}^{-1} \right)_{s=t} \right] \quad (30)$$

For the first term of the sum, one can argue in a way similar to the proof of Proposition 1.2 and obtain the inequality

$$\left\| PQ_{t,\gamma(t)} \frac{d}{dt} \gamma(t) \right\|_{C^3} \leq \sigma \left\| \frac{d}{dt} \gamma(t) \right\|_{C^3} + \tilde{K} \quad (31)$$

with $\sigma < 1$ and $\tilde{K} > 0$ (in general different from the \tilde{K} above). To see this, notice that in the referenced calculations we use uniform bounds for $F_{t,\varphi}^{(k)}$ in φ , only exploiting its unit integral. Furthermore, since $\int \frac{d}{dt}\gamma(t) = 0$, the bound $\|\varphi\|_{C^0} \leq 1 + \|\varphi'\|_{C^0}$ simplifies to $\|\frac{d}{dt}\gamma(t)\|_{C^0} \leq \|(\frac{d}{dt}\gamma(t))'\|_{C^0}$.

The regularity assumptions allow to deduce that the C^3 norm of the other terms is uniformly bounded. With computations similar to those in the proof of Proposition 1.2, one can deduce there are $\sigma < 1$ (by possibly decreasing ε_1) and $\tilde{K} > 0$ such that $\|\Xi\|_{C^3} \leq \sigma\|\frac{d}{dt}\gamma\|_{C^3} + \tilde{K}$. Putting this estimate together with the one in (28) we get

$$\begin{aligned} \left\| \frac{d}{dt}(\mathcal{L}\gamma)(t) \right\|_{C^3} &\leq \left\| \left(\frac{d}{ds}\mathcal{L}_s \right)_{s=t} \gamma(t) \right\|_{C^3} + \|\Xi(t)\|_{C^3} \\ &\leq \sigma \left\| \frac{d}{dt}\gamma(t) \right\|_{C^3} + \tilde{K}. \end{aligned}$$

We conclude that there is a K_1 sufficiently large such that if $\|\frac{d}{dt}\gamma(t)\|_{C^3} < K_1$ for all t values sufficiently small, $\|\frac{d}{dt}(\mathcal{L}\gamma)(t)\|_{C^3} \leq K_1$.

Now, assuming that $\|\frac{d}{dt}\gamma\|_{C^3} < K_1$, we carry out similar computations in order to bound $\|\frac{d^2}{dt^2}(\mathcal{L}\gamma)(t)\|_{C^2}$.

$$\frac{d^2}{dt^2}(\mathcal{L}\gamma)(t) = \frac{d}{dt} \left(\left(\frac{d}{ds}\mathcal{L}_s \right)_{s=t} \gamma(t) \right) + \frac{d}{dt}\Xi(t)$$

The first term corresponds to the derivative with respect to t of (27) when $\varphi = \gamma(t)$, which in turn is a combination of the functions

$$\begin{aligned} &\partial_x \left(\frac{\gamma(t)}{\partial_x \Phi_{t,\gamma(t)}} \right) \circ \Phi_{t,\gamma(t)}^{-1}, \quad \partial_s \Phi_{s,\gamma(t)}^{-1}|_{s=t}, \\ &\frac{\gamma(t)}{(\partial_x \Phi_{t,\gamma(t)})^2} \circ \Phi_{t,\gamma(t)}^{-1}, \quad \partial_s \partial_x \Phi_{s,\gamma(t)}|_{s=t} \circ \Phi_{t,\gamma(t)}^{-1}, \quad \partial_x^2 \Phi_{t,\gamma(t)} \circ \Phi_{t,\gamma(t)}^{-1}. \end{aligned}$$

With the given regularity hypotheses all the derivatives with respect to t of the above terms are continuous in the variable x . In particular, for the first term notice that the assumed regularity of all the terms implies

$$\frac{d}{dt} \partial_x \left(\frac{\gamma(t)}{\partial_x \Phi_{t,\gamma(t)}} \right) = \partial_x \frac{d}{dt} \left(\frac{\gamma(t)}{\partial_x \Phi_{t,\gamma(t)}} \right).$$

We conclude that there exists a $\tilde{K} > 0$ such that $\|\frac{d}{dt} \left(\left(\frac{d}{ds}\mathcal{L}_s \right)_{s=t} \gamma(t) \right)\|_{C^2} < \tilde{K}$.

As for $\frac{d}{dt}\Xi(t)$, consider the decomposition of $\Xi(t)$ given by (29)–(30). The derivative of $PQ_{t,\gamma(t)}\frac{d}{dt}\gamma(t)$ is

$$\begin{aligned} \frac{d}{ds} \left(PQ_{s,\gamma(s)} \frac{d}{ds}\gamma(s) \right)_{s=t} &= P \lim_{\delta \rightarrow 0} \frac{(Q_{t+\delta,\gamma(t+\delta)} - Q_{t,\gamma(t+\delta)}) \frac{d}{dt}\gamma(t+\delta)}{\delta} \\ &\quad + P \lim_{\delta \rightarrow 0} \frac{(Q_{t,\gamma(t+\delta)} - Q_{t,\gamma(t)}) \frac{d}{dt}\gamma(t+\delta)}{\delta} \\ &\quad + P \lim_{\delta \rightarrow 0} \frac{Q_{t,\gamma(t)} \left(\frac{d}{dt}\gamma(t+\delta) - \frac{d}{dt}\gamma(t) \right)}{\delta}. \end{aligned}$$

The first term in the sum on the RHS can be treated with computations similar to those leading to (27), and equals

$$\begin{aligned} & P \lim_{\delta \rightarrow 0} \frac{(Q_{t+\delta, \gamma(t)} - Q_{t, \gamma(t)}) \frac{d}{dt} \gamma(t)}{\delta} \\ &= P \left(\partial_x \left(\frac{\frac{d\gamma(t)}{dt}}{\partial_x \Phi_{t, \gamma(t)}} \right) \circ \Phi_{t, \gamma(t)}^{-1} \partial_s \Phi_{s, \gamma(t)}^{-1} \Big|_{s=t} \right. \\ & \quad \left. - \frac{\frac{d\gamma(t)}{dt}}{(\partial_x \Phi_{t, \gamma(t)})^2} \circ \Phi_{t, \gamma(t)}^{-1} \left[\partial_s \partial_x \Phi_{s, \gamma(t)} \Big|_{s=t} \circ \Phi_{t, \gamma(t)}^{-1} + \partial_x^2 \Phi_{t, \gamma(t)} \circ \Phi_{t, \gamma(t)}^{-1} \partial_s \Phi_{s, \gamma(t)}^{-1} \Big|_{s=t} \right] \right) \end{aligned}$$

which can be verified to have uniformly bounded C^2 norm.

The second term can be expressed as

$$P \lim_{\delta \rightarrow 0} \frac{(Q_{t, \gamma(t+\delta)} - Q_{t, \gamma(t)}) \frac{d}{dt} \gamma(t+\delta)}{\delta} = \frac{d}{ds} \left(\frac{\frac{d\gamma(t)}{dt}}{\partial_x \Phi_{t, \gamma(s)}} \circ \Phi_{t, \gamma(s)}^{-1} \right)_{s=t} \quad (32)$$

and noticing that

$$\frac{d}{ds} \frac{1}{\partial_x \Phi_{t, \gamma(s)}} = - \frac{1}{(\partial_x \Phi_{t, \gamma(s)})^2} \frac{d}{ds} \partial_x \Phi_{t, \gamma(s)}, \quad \frac{d}{ds} \left(\frac{d\gamma(t)}{dt} \right), \quad \text{and } \partial_s \Phi_{t, \gamma(s)}^{-1}$$

have uniformly bounded C^2 norm, it follows from straightforward computations that the whole of (32) has bounded norm. \parallel

For the third term we can find an inequality with a similar argument that leads to (31):

$$\left\| P Q_{t, \gamma(t)} \frac{d^2}{dt^2} \gamma(t) \right\|_{C^2} \leq \sigma \left\| \frac{d^2}{dt^2} \gamma(t) \right\|_{C^2} + \tilde{K}$$

For what concerns the derivative of (30), it involves the same terms listed above. Most of the factors appearing have been previously treated apart from

$$\frac{d}{ds} \left(\frac{1}{1 + t A'_{\gamma(s)}} \right)_{s=t} \quad \text{and} \quad \partial_x (Q_{t, \gamma(t)} \gamma(t)).$$

For the first one is easy to see that all the derivatives in t of $1 + t A'_{\gamma(s)}$ have bounded C^2 norm in x . For the second one, it follows from previous computations that $Q_{t, \gamma(t)} \gamma(t)$ is at least $C^1((-\varepsilon, \varepsilon), S^1)$ and this implies

$$\frac{d}{dt} \partial_x Q_{t, \gamma(t)} \gamma(t) = \partial_x \frac{d}{dt} Q_{t, \gamma(t)} \gamma(t).$$

$\frac{d}{dt} Q_{t, \gamma(t)} \gamma(t)$ can be shown to have uniformly bounded C^3 norm in the variable x with the same computations used to control (26), and therefore $\|\partial_x \frac{d}{dt} Q_{t, \gamma(t)} \gamma(t)\|_{C^2} \leq \tilde{K}$. Therefore,

$$\left\| \frac{d^2}{dt^2} (\mathcal{L}\gamma)(t) \right\|_{C^2} \leq \sigma \left\| \frac{d^2}{dt^2} \gamma(t) \right\|_{C^2} + \tilde{K}$$

so we can find K_2 sufficiently large so that $\|\frac{d}{dt} \gamma(t)\|_{C^3} < K_1$, with K_1 as above, and $\|\frac{d^2}{dt^2} \gamma(t)\|_{C^2} < K_2$ implies $\|\frac{d^2}{dt^2} (\mathcal{L}\gamma)(t)\|_{C^2} < K_2$. \square

\parallel Notice that in the expression above the factor $\partial_x \left(\frac{d\gamma}{dt} \right)$ appears, and it is for this kinds of factors that the choice of controlling two different norms for $\frac{d\gamma}{dt}$ and $\frac{d^2\gamma}{dt^2}$ becomes crucial.

Proof of Theorem 1.2 Point (ii). Given any $\gamma \in \hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2}$, it follows from Proposition 1.3 that $\{\mathcal{L}^n \gamma\}_{n \in \mathbb{N}_0} \in \hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2}$. Since $\hat{\mathcal{C}}_{\varepsilon_2, K_1, K_2} \subset C^2([-\varepsilon_2, \varepsilon_2], C^2(S^1, \mathbb{R}))$ is relatively compact in $C^1([-\varepsilon_2, \varepsilon_2], C^1(S^1, \mathbb{R}))$ **, there is a strictly increasing sequence n_k and $\bar{\gamma} \in C^1([-\varepsilon_2, \varepsilon_2], C^1(S^1, \mathbb{R}))$ such that $\mathcal{L}^{n_k} \gamma \rightarrow \bar{\gamma}$ in the C^1 topology. In particular, for every t , $\|\mathcal{L}^{n_k} \gamma(t) - \bar{\gamma}(t)\|_{C^0} \rightarrow 0$ for $k \rightarrow \infty$. However, from Proposition 1.1 we know that for every t , $\|\mathcal{L}_t^n \gamma(t) - \rho(t)\|_{C^0} \rightarrow 0$ for $n \rightarrow \infty$ from which we deduce that $\bar{\gamma}(t) = \rho(t)$, and therefore $t \mapsto \rho(t)$ belongs to $C^1([-\varepsilon_2, \varepsilon_2], C^1(S^1, \mathbb{R}))$. \square

4 Linear Response Formula

In this section we derive a linear response formula for the curve $t \mapsto \rho(t)$ of fixed densities. We proceed to prove the proposition in an analogous way to [4, Theorem 2.2] and apply perturbation arguments to deal with the self-consistency.

Lemma 4.1. *Considering $\{P_t\}_{t \in [-\varepsilon_2, \varepsilon_2]}$ as bounded operators from $C^2(S^1, \mathbb{R})$ to $C^1(S^1, \mathbb{R})$, we claim that the map $t \mapsto P_t$ is differentiable at $\hat{t} \in (-\varepsilon_2, \varepsilon_2)$ and*

$$\partial_t P_t|_{t=\hat{t}}(\varphi) = -P \left(\left(\frac{\varphi}{\Phi'_{\hat{t}, \rho(\hat{t})}} (A_{\rho(\hat{t})} + \hat{t} A_{\partial_t \rho(t)|_{t=\hat{t}}}) \right) \circ \Phi_{\hat{t}, \rho(\hat{t})}^{-1} \right)'.$$

Proof. Recall that $P_t = P Q_{t, \rho(t)}$ which implies that

$$\partial_t P_t = P \partial_t Q_{t, \rho(t)}. \quad (33)$$

Recall also that for $t \in [-\varepsilon_0, \varepsilon_0]$, $\Phi_{t, \rho(t)}$ is a diffeomorphism of S^1 and $t \mapsto \Phi_{t, \rho(t)}$ is C^2 . Let $\varphi \in C^2$.

$$\begin{aligned} & Q_{t, \rho(t)}(\varphi) - Q_{\hat{t}, \rho(\hat{t})}(\varphi) = \\ &= \frac{\varphi}{\Phi'_{t, \rho(t)}} \circ \Phi_{t, \rho(t)}^{-1} - \frac{\varphi}{\Phi'_{\hat{t}, \rho(\hat{t})}} \circ \Phi_{\hat{t}, \rho(\hat{t})}^{-1} \\ &= -(t - \hat{t}) (\partial_t \Phi_{t, \rho(t)}|_{t=t^*} \circ \Phi_{t^*, \rho(t^*)}^{-1})' \left(\frac{\varphi}{\Phi'_{t^*, \rho(t^*)}} \circ \Phi_{t^*, \rho(t^*)}^{-1} \right) \\ &\quad - (t - \hat{t}) \partial_t \Phi_{t, \rho(t)}|_{t=t^*} \circ \Phi_{t^*, \rho(t^*)}^{-1} \left(\frac{\varphi'}{(\Phi'_{t^*, \rho(t^*)})^2} \circ \Phi_{t^*, \rho(t^*)}^{-1} - \frac{\varphi \cdot \Phi''_{t^*, \rho(t^*)}}{(\Phi'_{t^*, \rho(t^*)})^3} \circ \Phi_{t^*, \rho(t^*)}^{-1} \right). \end{aligned}$$

for some $t^* \in (t, \hat{t})$. Notice that

$$\begin{aligned} & (\partial_t \Phi_{t, \rho(t)}|_{t=t^*} \circ \Phi_{t^*, \rho(t^*)}^{-1})' \left(\frac{\varphi}{\Phi'_{t^*, \rho(t^*)}} \circ \Phi_{t^*, \rho(t^*)}^{-1} \right) \\ &= \left(\frac{\varphi}{(\Phi'_{t^*, \rho(t^*)})^2} \circ \Phi_{t^*, \rho(t^*)}^{-1} \right) \partial_t \Phi'_{t, \rho(t)}|_{t=t^*} \circ \Phi_{t^*, \rho(t^*)}^{-1} \end{aligned}$$

**Recall that $C^2(S^1, \mathbb{R})$ is compactly embedded in $C^1(S^1, \mathbb{R})$, and thus given a bounded ball $B' \subset C^2(S^1, \mathbb{R})$ and any $\delta > 0$ there are $\mathcal{F}_\delta := \{\varphi_i\}_{i=1}^{N_\delta} \subset C^1(S^1, \mathbb{R})$ such that the union of the $\delta/2$ C^1 -balls centered at the functions in \mathcal{F}_δ covers B' . Fix a bounded ball $B \subset C^2([-\varepsilon_2, \varepsilon_2], C^2(S^1, \mathbb{R}))$, there is $B' \subset C^2(S^1, \mathbb{R})$ such that $\gamma([-\varepsilon_2, \varepsilon_2]) \subset B'$ for any $\gamma \in B$. Now fix $\delta > 0$ and pick points $\{t_i\}_{i=1}^M$ partitioning $[-\varepsilon_2, \varepsilon_2]$ such that $\|\gamma(s_1) - \gamma(s_2)\|_{C^0} < \delta/2$ whenever $s_1, s_2 \in [t_i, t_{i+1}]$ for some i . Then, for any $(i_1, \dots, i_M) \in \{1, \dots, N\}^M$ fix a C^1 function $\gamma_{(i_1, \dots, i_M)} : [-\varepsilon_2, \varepsilon_2] \rightarrow C^1(S^1, \mathbb{R})$ such that $\gamma_{(i_1, \dots, i_M)}(t_j) = \varphi_{i_j}$. It is then easy to check that for every $\gamma \in B$ there is $(i_1, \dots, i_M) \in \{1, \dots, N\}^M$ such that γ is at C^1 distance no more than δ from $\gamma_{(i_1, \dots, i_M)}$.

implying that

$$\begin{aligned}
\partial_t Q_{t,\rho(t)}|_{t=\hat{t}}(\varphi) &= - \left(\frac{\varphi}{(\Phi'_{\hat{t},\rho(\hat{t})})^2} \partial_t \Phi'_{t,\rho(t)}|_{t=\hat{t}} \right) \circ \Phi_{\hat{t},\rho(\hat{t})}^{-1} \\
&\quad - \left(\left(\frac{\varphi'}{(\Phi'_{\hat{t},\rho(\hat{t})})^2} - \frac{\varphi \cdot \Phi''_{\hat{t},\rho(\hat{t})}}{(\Phi'_{\hat{t},\rho(\hat{t})})^3} \right) \cdot \partial_t \Phi_{t,\rho(t)}|_{t=\hat{t}} \right) \circ \Phi_{\hat{t},\rho(\hat{t})}^{-1} \\
&= - \left(\left(\frac{\varphi}{\Phi'_{\hat{t},\rho(\hat{t})}} (\partial_t \Phi_{t,\rho(t)})_{t=\hat{t}} \right) \circ \Phi_{\hat{t},\rho(\hat{t})}^{-1} \right)'
\end{aligned}$$

and by (33)

$$\begin{aligned}
\partial_t P_t|_{t=\hat{t}}(\varphi) &= -P \left(\left(\frac{\varphi}{\Phi'_{\hat{t},\rho(\hat{t})}} (\partial_t \Phi_{t,\rho(t)})_{t=\hat{t}} \right) \circ \Phi_{\hat{t},\rho(\hat{t})}^{-1} \right)' \\
&= -P \left(\left(\frac{\varphi}{\Phi'_{\hat{t},\rho(\hat{t})}} (A_{\rho(\hat{t})} + \hat{t} A_{\partial_t \rho(t)|_{t=\hat{t}}}) \right) \circ \Phi_{\hat{t},\rho(\hat{t})}^{-1} \right)'.
\end{aligned}$$

□

Notice that this lemma implies that

$$\partial_t P_t|_{t=\hat{t}}\rho(\hat{t}) = -P[K_{\hat{t}}(\rho(\hat{t})) + \hat{t}K_{\hat{t}}(\partial_t \rho|_{t=\hat{t}})],$$

where $K_{\hat{t}}$ is given by (10).

Proof of Proposition 1.4. Recall that P_t is the transfer operator of $f \circ \Phi_{t,\rho(t)}$ which is an at least C^3 transformation of S^1 . Furthermore, there is $K > 0$ such that $|\Phi'_{t,\rho(t)}| > 1 - Kt$ and thus $f \circ \Phi_{t,\rho(t)}$ is uniformly expanding for any t with $|t| < \varepsilon$ for some sufficiently small ε . Standard results on uniformly expanding maps (see [11]) imply that P_t has a spectral gap in $C^2(S^1, \mathbb{R})$. In particular, 1 is an isolated and simple eigenvalue and $\mathbb{R}\rho(t)$ is the corresponding eigenspace. Because of the uniform Lasota-Yorke type inequalities (24)–(25), a uniform bound can be given on the spectral gap of all P_t such that t is small enough. This implies that we can find a positively oriented curve γ on the complex plane around 1 such that no other element of the spectrum of any P_t is contained inside of it, and the projection formula

$$\rho(t) = \frac{1}{2\pi i} \int_{\gamma} (z - P_t)^{-1} \varphi(z) dz \quad (34)$$

holds for any $\varphi \in C^2$ such that $\int_{S^1} \varphi = 1$. One can prove this using the decomposition of $\mathcal{L}_{t,\rho(t)}$ as stated in [11, Theorem 7.1.1] (exploiting the fact that $F_{t,\rho(t)}$ is C^2 and uniformly expanding which properties by [22] imply that $F_{t,\rho(t)}$ mixing) and a residue computation.

For $z \in \gamma$ we have

$$(z - P_{\hat{t}})^{-1} - (z - P_t)^{-1} = (z - P_{\hat{t}})^{-1} (P_{\hat{t}} - P_t) (z - P_t)^{-1} \quad (35)$$

where we view $(z - P_t)^{-1}$ as an operator acting on $C^2(S^1, \mathbb{R})$, $(P_{\hat{t}} - P_t)$ as an operator from $C^2(S^1, \mathbb{R})$ to $C^1(S^1, \mathbb{R})$ and $(z - P_{\hat{t}})^{-1}$ as acting on $C^1(S^1, \mathbb{R})$.

Letting $t \rightarrow \hat{t}$ in (35) by Lemma 4.1 we have that

$$\partial_t(z - P_t)^{-1}|_{t=\hat{t}} = (z - P_{\hat{t}})^{-1}\partial_t P_t|_{t=\hat{t}}(z - P_{\hat{t}})^{-1}. \quad (36)$$

Substituting $\varphi = \rho(\hat{t})$ in (34) and differentiating both sides we get

$$\begin{aligned} \partial_t \rho(t)|_{t=\hat{t}} &= \frac{1}{2\pi i} \int_{\gamma} (z - P_{\hat{t}})^{-1} \partial_t P_t|_{t=0} (z - P_{\hat{t}})^{-1} \rho(\hat{t})(z) dz \\ &= \frac{1}{2\pi i} \int_{\gamma} (z - P_{\hat{t}})^{-1} \partial_t P_t|_{t=\hat{t}} \rho(\hat{t})(z) \frac{1}{z-1} dz \\ &= -\frac{1}{2\pi i} \int_{\gamma} (z - P_{\hat{t}})^{-1} P(\mathbf{K}_{\hat{t}}(\rho(\hat{t})) + \hat{t}\mathbf{K}_{\hat{t}}(\partial_t \rho(t)|_{t=\hat{t}}))(z) \frac{1}{z-1} dz \\ &= -(1 - P_{\hat{t}})^{-1} P(\mathbf{K}_{\hat{t}}(\rho(\hat{t})) + \hat{t}\mathbf{K}_{\hat{t}}(\partial_t \rho(t)|_{t=\hat{t}})), \end{aligned} \quad (37)$$

where the last step is a residue computation.

Thus we can write

$$(1 - P_{\hat{t}} + \hat{t}PK_{\hat{t}})\partial_t \rho(t)|_{t=\hat{t}} = -PK_{\hat{t}}(\rho(\hat{t})).$$

It is a well-known fact that invertible operators form an open set in the space of bounded linear operators between normed spaces. More precisely, if an operator T is invertible, then any operator S such that $\|T - S\| \leq \|T^{-1}\|^{-1}$ is also invertible. Since $1 - P_{\hat{t}}$ is invertible as an operator from $C^2(S^1, \mathbb{R})$ to $C^1(S^1, \mathbb{R})$, the invertibility of $1 - P_{\hat{t}} + \hat{t}PK_{\hat{t}}$ follows if we are able to choose \hat{t} so small that

$$\|\hat{t}PK_{\hat{t}}\|_{C^2 \rightarrow C^1} \leq \|(1 - P_{\hat{t}})^{-1}\|_{C^2 \rightarrow C^1}^{-1} \Leftrightarrow \hat{t} \leq \frac{1}{\|PK_{\hat{t}}\|_{C^2 \rightarrow C^1} \|(1 - P_{\hat{t}})^{-1}\|_{C^2 \rightarrow C^1}}.$$

For this it suffices to show that $t \mapsto \|PK_t\|_{C^2 \rightarrow C^1} \|(1 - P_t)^{-1}\|_{C^2 \rightarrow C^1}$ is uniformly bounded in a small interval around zero. The norm $\|(1 - P_t)^{-1}\|_{C^2 \rightarrow C^1}$ is finite for all sufficiently small t (because of the above mentioned spectral gap), so it can be uniformly bounded in an interval $[-\varepsilon_2, \varepsilon_2]$. As for K_t , we can compute that

$$\begin{aligned} \sup_{x \in S^1} |K_t \varphi(x)| &= \\ &= \sup_{x \in S^1} \left| \left(\left(\frac{\rho(t)}{\Phi'_{t,\rho(t)}} A_{\varphi} \right) \circ \Phi_{t,\rho(t)}^{-1}(x) \right)' \right| \\ &= \sup_{x \in S^1} \left| \frac{\rho'(t) A_{\varphi} \Phi'_{t,\rho(t)} + \rho(t) A'_{\varphi} \Phi'_{t,\rho(t)} - \rho(t) A_{\varphi} \Phi''_{t,\rho(t)}}{(\Phi'_{t,\rho(t)})^2} \circ \Phi_{t,\rho(t)}^{-1}(x) \cdot (\Phi_{t,\rho(t)}^{-1})'(x) \right| \\ &\leq \frac{1}{(1 - tK)^3} \left(K(1 + tK) \sup_{x \in S^1} |\rho'(t)| + (K(1 + tK) + tK^2) \sup_{x \in S^1} |\rho(t)| \right) \end{aligned}$$

and a similar formula for $\sup_{x \in S^1} |(K_t \varphi)'(x)|$ depending smoothly on ρ, ρ', ρ'' and t . Since $(t, x) \mapsto \rho(t)(x)$ is C^2 , $\|K_t\|_{C^2 \rightarrow C^1}$ is uniformly bounded for $t \in [-\varepsilon_2, \varepsilon_2]$.

Thus we can choose ε_3 further decreasing ε_2 such that $1 - P_{\hat{t}} + \hat{t}PK_{\hat{t}}$ is invertible for $t \in [-\varepsilon_3, \varepsilon_3]$ and

$$\partial_t \rho(t)|_{t=\hat{t}} = -(1 - P_{\hat{t}} + \hat{t}PK_{\hat{t}})^{-1} PK_{\hat{t}}(\rho(\hat{t})).$$

□

A Appendix

In this appendix we recall some simple formulae on derivatives that have been used extensively throughout the document. First of all higher derivatives of a product: Assume $\Phi, \Psi \in C^4$, then

$$(\Phi\Psi)' = \Phi'\Psi + \Phi\Psi' \quad (38)$$

$$(\Phi\Psi)'' = \Phi''\Psi + 2\Phi'\Psi' + \Phi\Psi'' \quad (39)$$

$$(\Phi\Psi)''' = \Phi'''\Psi + 3\Phi''\Psi' + 3\Psi'\Phi'' + \Psi\Phi''' \quad (40)$$

$$(\Phi\Psi)'''' = \Phi''''\Psi + 4\Phi'''\Psi' + 6\Phi''\Psi'' + 4\Phi'\Psi''' + \Psi\Phi'''' \quad (41)$$

which implies that there is a constant $M > 0$ such that

$$\|\Phi\Psi\|_{C^k} \leq M\|\Phi\|_{C^k}\|\Psi\|_{C^k}, \quad k = 1, \dots, 4. \quad (42)$$

Higher derivatives of a composition: Assume $\Phi, \Psi \in C^4$, then

$$(\Phi \circ \Psi)' = \Phi' \circ \Psi\Psi' \quad (43)$$

$$(\Phi \circ \Psi)'' = \Phi'' \circ \Psi(\Psi')^2 + \Phi' \circ \Psi\Psi'' \quad (44)$$

$$(\Phi \circ \Psi)''' = \Phi''' \circ \Psi(\Psi')^3 + 2\Phi'' \circ \Psi\Psi'\Psi'' + \Phi'' \circ \Psi\Psi''\Psi' + \Phi' \circ \Psi\Psi''' \quad (45)$$

$$\begin{aligned} (\Phi \circ \Psi)'''' &= \Phi'''' \circ \Psi(\Psi')^4 + \Phi''' \circ \Psi [3(\Psi')^2\Psi'' + \Psi''(\Psi')^2 + 2(\Psi')^2\Psi'''] + \\ &\quad + \Phi'' \circ \Psi [3(\Psi'')^2 + 4\Psi'\Psi'''] + \Phi' \circ \Psi(\Psi''''). \end{aligned} \quad (46)$$

which implies that there is a numerical constant M such that

$$\|\Phi \circ \Psi\|_{C^k} \leq M\|\Phi\|_{C^k} \cdot \max_{\ell=1, \dots, k} \|\Psi\|_{C^k}^\ell, \quad k = 1, \dots, 4. \quad (47)$$

We also give the formulae for the higher derivatives of the inverse function. Let $\Phi \in C^4$ be invertible and such that $\min |\Phi'| > 0$, then

$$(\Phi^{-1})' = \frac{1}{\Phi'} \circ \Phi^{-1} \quad (48)$$

$$(\Phi^{-1})'' = -\frac{\Phi''}{(\Phi')^2} \circ \Phi^{-1}(\Phi^{-1})' = -\frac{\Phi''}{(\Phi')^3} \circ \Phi^{-1} \quad (49)$$

$$\begin{aligned} (\Phi^{-1})''' &= \left[2\frac{(\Phi'')^2}{(\Phi')^3} - \frac{\Phi'''}{(\Phi')^2} \right] \circ \Phi^{-1}(\Phi^{-1})' - \frac{\Phi''}{(\Phi')^2} \circ \Phi^{-1}(\Phi^{-1})'' \\ &= \left[2\frac{(\Phi'')^2}{(\Phi')^4} - \frac{\Phi'''}{(\Phi')^3} - \frac{(\Phi'')^2}{(\Phi')^5} \right] \circ \Phi^{-1} \end{aligned} \quad (50)$$

$$(\Phi^{-1})'''' = \left[4\frac{\Phi''\Phi'''}{(\Phi')^5} - 8\frac{(\Phi'')^3}{(\Phi')^6} - \frac{\Phi''''}{(\Phi')^4} + 3\frac{\Phi'''\Phi''}{(\Phi')^5} - 2\frac{\Phi'''}{(\Phi')^6} + 5\frac{(\Phi'')^3}{(\Phi')^7} \right] \circ \Phi^{-1} \quad (51)$$

and so there exists an $M > 0$ such that

$$\|\Phi^{-1}\|_{C^k} \leq M \cdot \max_{n=1, \dots, 7} (\min |\Phi'|)^{-n} \cdot \max_{\ell=1, \dots, 3} \|\Phi\|_{C^k}^\ell, \quad k = 1, \dots, 4. \quad (52)$$

References

- [1] ASPENBERG, M., BALADI, V., LEPPÄNEN, J., AND PERSSON, T. On the fractional susceptibility function of piecewise expanding maps. *arXiv preprint arXiv:1910.00369* (2019).
- [2] BAHOUN, W., AND SAUSSOL, B. Linear response in the intermittent family: differentiation in a weighted C^0 -norm. *Discrete and Continuous Dynamical Systems (12)*, 36 (2016), 6657–6668.
- [3] BALADI, V. On the susceptibility function of piecewise expanding interval maps. *Communications in Mathematical Physics* 275, 3 (2007), 839–859.
- [4] BALADI, V. Linear response, or else. Proceedings of the International Congress of Mathematicians-Seoul. Vol. III. Invited lectures, 2014.
- [5] BALADI, V., BENEDICKS, M., AND SCHNELLMANN, D. Whitney–Hölder continuity of the SRB measure for transversal families of smooth unimodal maps. *Inventiones mathematicae* 201, 3 (2015), 773–844.
- [6] BALADI, V., KUNA, T., AND LUCARINI, V. Linear and fractional response for the SRB measure of smooth hyperbolic attractors and discontinuous observables. *Nonlinearity* 30, 3 (2017), 1204.
- [7] BALADI, V., AND SMANIA, D. Linear response formula for piecewise expanding unimodal maps. *Nonlinearity* 21, 4 (2008), 677.
- [8] BALADI, V., AND SMANIA, D. Alternative proofs of linear response for piecewise expanding unimodal maps. *Ergodic Theory and Dynamical Systems* 30, 1 (2010), 1–20.
- [9] BALADI, V., AND TODD, M. Linear response for intermittent maps. *Communications in Mathematical Physics* 347, 3 (2016), 857–874.
- [10] BÁLINT, P., KELLER, G., SÉLLEY, F. M., AND TÓTH, I. P. Synchronization versus stability of the invariant distribution for a class of globally coupled maps. *Nonlinearity* 31, 8 (2018), 3770.
- [11] BOYARSKY, A., AND GORA, P. *Laws of chaos: invariant measures and dynamical systems in one dimension*. Springer Science & Business Media, 2012.
- [12] DOLGOPYAT, D. On differentiability of SRB states for partially hyperbolic systems. *Inventiones mathematicae* 155, 2 (2004), 389–449.
- [13] FERNANDEZ, B. Breaking of ergodicity in expanding systems of globally coupled piecewise affine circle maps. *Journal of Statistical Physics* 154, 4 (2014), 999–1029.
- [14] GALATOLO, S., AND GIULIETTI, P. A linear response for dynamical systems with additive noise. *Nonlinearity* 32, 6 (2019), 2269.
- [15] GALATOLO, S., AND POLLICOTT, M. Controlling the statistical properties of expanding maps. *Nonlinearity* 30, 7 (2017), 2737.

- [16] GALATOLO, S., AND SEDRO, J. Quadratic response of random and deterministic dynamical systems. *Chaos: An Interdisciplinary Journal of Nonlinear Science* 30, 2 (2020), 023113.
- [17] JIANG, M., AND DE LA LLAVE, R. Smooth dependence of thermodynamic limits of SRB-measures. *Communications in Mathematical Physics* 211, 2 (2000), 303–333.
- [18] JIANG, M., AND DE LA LLAVE, R. Linear response function for coupled hyperbolic attractors. *Communications in Mathematical Physics* 261, 2 (2006), 379–404.
- [19] KELLER, G. An ergodic theoretic approach to mean field coupled maps. In *Fractal Geometry and Stochastics II*. Springer, 2000, pp. 183–208.
- [20] KLOECKNER, B. The linear request problem. *Proceedings of the American Mathematical Society* 146, 7 (2018), 2953–2962.
- [21] KOREPANOV, A. Linear response for intermittent maps with summable and nonsummable decay of correlations. *Nonlinearity* 29, 6 (2016), 1735.
- [22] KRZYŻEWSKI, K., AND SZLENK, W. On invariant measures for expanding differentiable mappings. *Studia Mathematica* 33, 1 (1969), 83–92.
- [23] LIVERANI, C. Decay of correlations. *Annals of Mathematics-Second Series* 142, 2 (1995), 239–302.
- [24] LUCARINI, V., BLENDER, R., HERBERT, C., RAGONE, F., PASCALE, S., AND WOUTERS, J. Mathematical and physical ideas for climate science. *Reviews of Geophysics* 52, 4 (2014), 809–859.
- [25] LUCARINI, V., AND SARNO, S. A statistical mechanical approach for the computation of the climatic response to general forcings. *Nonlinear Processes in Geophysics* 18, 1 (2011), 7–28.
- [26] MACKAY, R. Management of complex dynamical systems. *Nonlinearity* 31, 2 (2018), R52.
- [27] MAJDA, A. J. Challenges in climate science and contemporary applied mathematics. *Communications on Pure and Applied Mathematics* 65, 7 (2012), 920–948.
- [28] RAGONE, F., LUCARINI, V., AND LUNKEIT, F. A new framework for climate sensitivity and prediction: a modelling perspective. *Climate Dynamics* 46, 5-6 (2016), 1459–1471.
- [29] RUELLE, D. Differentiation of SRB states. *Communications in Mathematical Physics* 187, 1 (1997), 227–241.
- [30] RUELLE, D. General linear response formula in statistical mechanics, and the fluctuation-dissipation theorem far from equilibrium. *Physics Letters A* 245, 3-4 (1998), 220–224.

- [31] RUELLE, D. A review of linear response theory for general differentiable dynamical systems. *Nonlinearity* 22, 4 (2009), 855.
- [32] RUELLE, D. Structure and f -dependence of the acim for a unimodal map f of Misiurewicz type. *Communications in Mathematical Physics* 287, 3 (2009), 1039–1070.
- [33] SÉLLEY, F., AND BÁLINT, P. Mean-field coupling of identical expanding circle maps. *Journal of Statistical Physics* 164, 4 (2016), 858–889.
- [34] VIANA, M. Stochastic dynamics of deterministic systems. Lecture Notes XXI Braz. Math. Colloq. IMPA Rio de Janeiro, 1997.
- [35] WORMELL, C. L., AND GOTTWALD, G. A. On the validity of linear response theory in high-dimensional deterministic dynamical systems. *Journal of Statistical Physics* 172, 6 (2018), 1479–1498.
- [36] WORMELL, C. L., AND GOTTWALD, G. A. Linear response for macroscopic observables in high-dimensional systems. *Chaos: An Interdisciplinary Journal of Nonlinear Science* 29, 11 (2019), 113127.
- [37] ZHANG, Z. On the smooth dependence of SRB measures for partially hyperbolic systems. *Communications in Mathematical Physics* 358, 1 (2018), 45–79.