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# Absorbent Porous Paper Reveals How Earthquakes Could be Mitigated

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4 **Absorbent porous paper reveals**  
5 **how earthquakes could be mitigated**

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9 **Key Points:**

- 10 • From an energetic point of view, absorbent porous paper can be an ideal, low-cost  
11 surrogate rock material for studying induced seismicity.  
12 • Segmentation of faults and sequential fluid injection in each segment can mitigate  
13 potential earthquakes by at least one order of magnitude.  
14 • We show that fault segmentation, segment-activation rate and stress state predominantly  
15 control the result of applied injection strategies.

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## Abstract

Earthquakes nucleate when large amounts of elastic energy, stored in the earth's crust, are suddenly released due to abrupt sliding over a fault. Fluid injections can reactivate existing seismogenic faults and induce/trigger earthquakes by increasing fluid pressure. Here we develop an analogous experimental system of simultaneously loaded and wetted absorbent porous paper to quantify theoretically the process of wetting-induced earthquakes. This strategy allows us to gradually release the stored energy by provoking low intensity tremors. We identify the key parameters that control the outcome of the applied injection strategy, which include the initial stress state, fault segmentation, and segment-activation rate. Subsequent injections, initiated at high stress levels, can drive the system faster towards its instability point, nucleating a large earthquake. Starting at low stress levels, however, they can reduce the magnitude of the natural event by at least one unit.

## Plain Language Summary

Understanding natural and anthropogenic seismicity is a major scientific challenge. Here we present a novel analogue fault model using absorbent porous paper, which gives new insights on earthquake mitigation. When scaled to in-situ conditions, the porous paper model represents a natural seismic rupture of magnitude  $M_w = 5.9$ . By progressively wetting it, we simulate fluid injections in the earth's crust and draw analogies to large-scale industrial projects. In our experiments, each injection is accompanied by tremors, which progressively release energy and modify the energy budget of the system. Without precise knowledge of the fault properties, we risk driving the system faster towards an unexpected large seismic event. However, provided that the model's key parameters - fault segmentation, segment-activation rate, and stress state - are well known or controlled, the natural rupture can be mitigated by at least one unit. We expect that these results will facilitate risk reduction in current fluid injection projects and inspire earthquake mitigation strategies for real tectonic faults.

## 1 Introduction

It is well recognized today that humans can cause earthquakes (Raleigh et al., 1976; McGarr et al., 2002; Guglielmi et al., 2015; Foulger et al., 2018). Examples of anthropogenic seismicity involve earthquakes provoked by large artificial water reservoirs such as dams (Gupta, 2002), mining (Li et al., 2007), underground nuclear explosions (Hamilton et al., 1972) or by fluid injections in the earth's crust (Ellsworth, 2013; Rubinstein & Mahani, 2015; Garagash & Germanovich, 2012; Schultz et al., 2020). The latter type of anthropogenic seismicity is of particular interest, due to the numerous ongoing industrial applications (Rubinstein & Mahani, 2015; Hosseini et al., 2018). These anthropogenic seismic events could shift the Gutenberg and Richter (1954) power law, which describes the relationship between the total number of earthquakes and their magnitudes, towards smaller events.

This idea is similar to the one traditionally used for reducing the risk of large snow avalanches. Snow avalanches follow the same frequency-magnitude distributions as natural earthquakes do (Birkeland & Landry, 2002). Nowadays, it is common practice to avoid large avalanches by provoking smaller ones. Similarly, large earthquakes could be probably mitigated by inducing low intensity tremors. This is what is called here earthquake mitigation. This notion was first mentioned by Raleigh et al. (1976), but has not been explored further since.

Surrogate materials can effectively substitute in-situ rock and gouge materials (Rosenau et al., 2017). Some examples of such materials are sandpaper (King, 1975), cardboard (Heslot et al., 1994), pasta (Knuth & Marone, 2007), steel (Popov et al., 2012), hydrogel (Latour et al., 2013) and puffed rice (Einav & Guillard, 2018). Here we use absorbent porous paper as an analogue fault material to explore earthquake mitigation. We show

that it can be an ideal low-cost surrogate material for reproducing earthquake-like instabilities in the laboratory. Wetting the absorbent paper allows us not only to weaken the material but also to induce small instabilities. By deriving adequate scaling laws (see supporting information §6), we can simulate fluid injections and study the transition from seismic (unstable, sudden) to aseismic (stable, creep-like) slip of an ideal fault (Stefanou, 2019; Cappa et al., 2019).

## 2 Paper-quakes vs. earth-quakes

Consider an isolated planar dip-slip fault, as depicted in Figure 1a, with a length of  $L_{ac}^{real} = 6.5$  km dipped at an angle of  $60^\circ$  (see also supporting information §5). Based on its properties and assuming a square rupture area ( $A^{real} \approx (L_{ac}^{real})^2$ ), this fault leads to an earthquake of  $M_w \approx 6$  (see supporting information §6). The energy budget is contained in a single sheet of absorbent paper which, consequently, can represent this fault (see also supporting information §1-2) if strained as shown in Figure 1b. Long range interactions with other faults are not considered herein.

The energetic equivalence of the paper analogue with a real fault can be observed in Figures 1c-e. The progressive accumulation of elastic energy in the rocks surrounding the fault zone is taken into account by a spring attached at the one end of the sheet (Figures 1d-e). The apparent stiffness of the spring is chosen to represent the real system. A constant slow velocity is applied at the extremity of this spring, simulating the slow far-field tectonic loading. Finally, paper has a non-negligible softening branch that simulates the critical slip distance  $d_c$  of faults (Scholz, 2002) (Figure 1f).

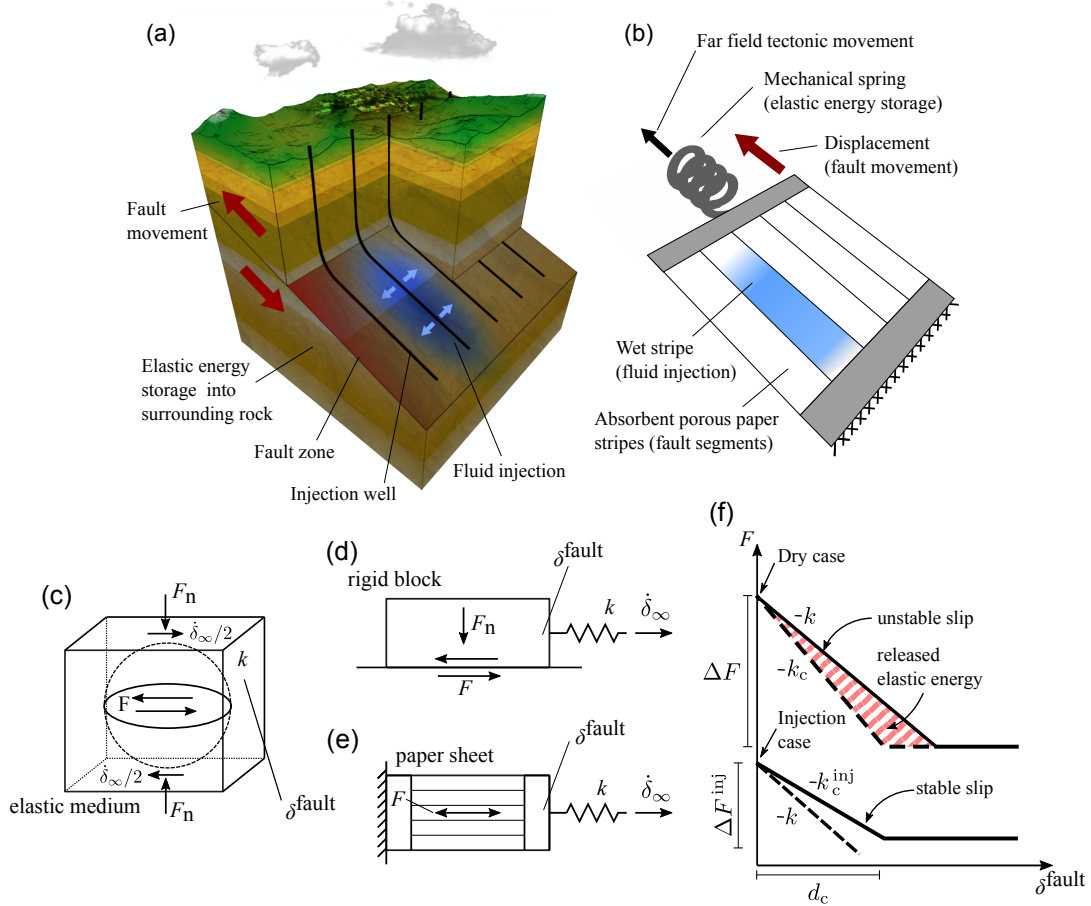
We can also simulate fluid injections by simply wetting the paper sheet. In a fault system, the apparent friction drops when fluids under pressure are injected into the fault zone due to the decrease of the effective normal stress. Similarly, porous paper shows a noticeable stress drop when it is wetted due to the reduction of its strength (Figures 1f and 2c). The ratio of the shear stress drop of real faults to the strength weakening of porous paper is defined here by the scaling factor  $a$ , which is a free parameter in our model (see also supporting information §6). Finally, opposite to wetting, healing could be considered by drying the paper stripes, which could potentially lead to repeatedly growing slip events. However, this is out of the scope of the present work, which focuses on earthquake mitigation by fault reactivation.

Notice that the fluid diffusion process which takes places during injections in wellbores can be considered in our analogue model in two ways, by water absorption of porous paper and by progressive wetting of many isolated stripes (see Figures 1b and e and Section 4). Assuring that tectonic loading is much slower (see Section 5 for more details) than diffusion, only the latter way is examined here.

The configuration shown in Figures 1b and e leads to a sudden release of the elastic energy upon rupture, in the same way as the energy stored in the rocks surrounding the fault zone is released during an earthquake. Note that typical failure modes II and/or III that take place during seismic slip in faults are represented here by a mode I failure of the porous paper. These systems are equivalent in terms of energy budget (Nussbaum & Ruina, 1987), provided that appropriate scaling laws are applied. By using these scaling laws (see supporting information §6) and measuring the elastic energy  $E_R^{paper}$  that accompanies paper failure in the surrogate system, one can estimate the earthquake magnitude  $M_w$  of the real system as follows:

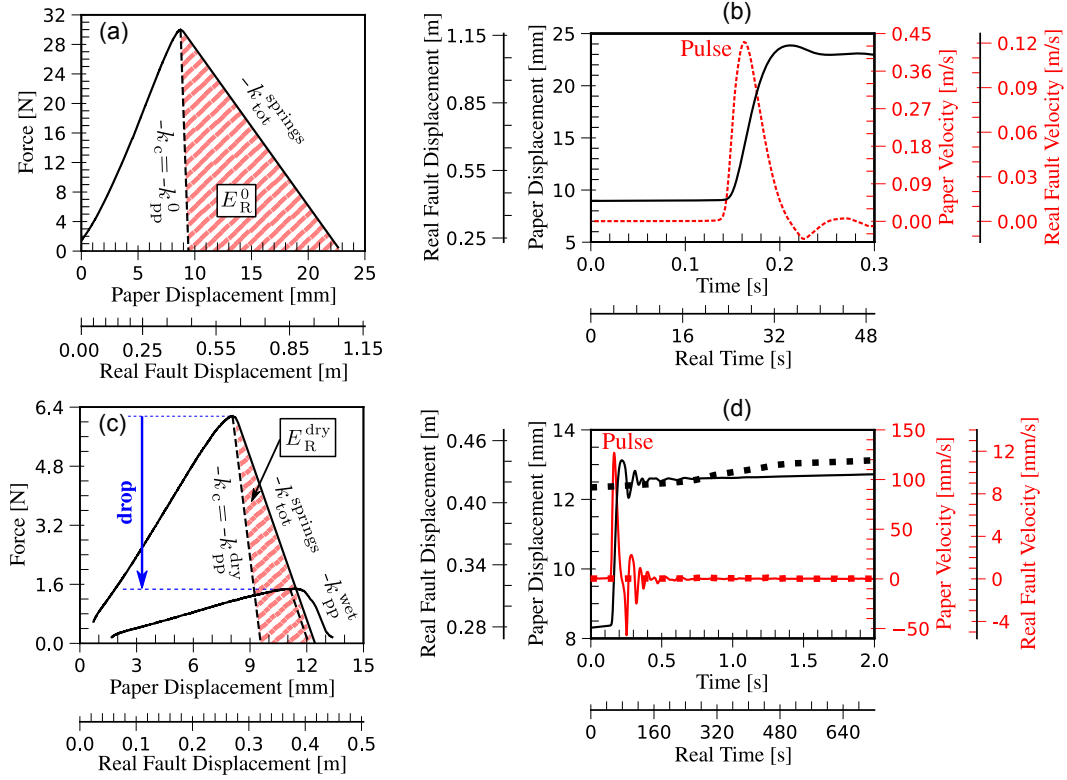
$$M_w = \frac{2}{3} \log_{10} (E_R^{paper}) + 6.36 \quad (E_R^{paper} \text{ in Nm}) \quad (1)$$

In Figure 2a, we present the force-displacement evolution of a paper sheet that is put under tension as described above. In the beginning, the porous paper shows a linear



**Figure 1.** (a) Simplified dip-slip fault where the fault zone area is divided in five segments parallel to the slip direction. Five wells allow to inject fluid into each segment independently. (b) Surrogate laboratory experiment consisting of five parallel stripes of absorbent paper, held together at their ends by clamps. Far-field movement is applied through a mechanical spring. The paper stripes can be wetted individually. (c) Pre-existing fault within an elastic medium of stiffness  $k$ , loaded by a very low velocity,  $\dot{\delta}_\infty$ . Normal forces  $F_n$  act on the fault, while frictional forces  $F$  are generated along the interface (Reid, 1910). (d) Spring-slider model (Scholz, 2002), composed of a rigid block, which is pulled through a spring over a rough surface. (e) Paper-spring model, where a stripe of paper is pulled through an elastic spring. (f) Schematic force-displacement diagram in which the post-peak properties are modified due to fluid injection on the fault (wetting in the porous paper case), allowing transition from unstable to stable slip.

113 behavior up to its peak strength. Then, a fracture appears in the paper sheet and the  
 114 spring is unloaded abruptly. The energy that is released during the unloading is equal  
 115 to  $E_R^{paper} = E_R^0 = 0.201 \text{ Nm}$  (hatched area), corresponding to an equivalent earthquake  
 116 of  $M_w = 5.9$  according to Eq. (1). This dynamic instability can also be observed by  
 117 the velocity pulse shown in Figure 2b, illustrating the analogy between paper-quakes and  
 118 earth-quakes (Kanamori & Brodsky, 2004) (see also supporting information Movie S1).



**Figure 2.** (a) Force-paper displacement diagram of a paper sheet experiment, showing paper behavior without a spring (dashed line) and with a spring attached to one end (solid line). (b) Pulse-like velocity (dotted line) of the dynamic paper sheet experiment, with corresponding slip (solid line). (c) Force-paper displacement diagram of representative dry and wet experiments, for a single paper stripe. Post-peak slope reduction ( $k_{pp}^{dry} > k_{pp}^{wet}$ ) and force drop are observed. In this example, a dynamic instability happens in the dry case ( $k_{tot}^{springs} < k_{pp}^{dry}$ ) releasing energy  $E_R^{dry}$ . In the wet case no instability occurs ( $k_{tot}^{springs} > k_{pp}^{wet}$ ). (d) Paper displacement (black) and velocity (red) evolution with time for dry (solid lines) and wet (dotted lines). The stabilizing effect of wetting is apparent. The additional axes scale the respective quantities to the real fault case (cf. supporting information §6).

### 119 3 From seismic to aseismic rupture

120 A seismic rupture, i.e. an earthquake, is a dynamic instability that happens when  
 121 the (elastic) unloading of the rocks surrounding the fault zone cannot be counterbalanced  
 122 by fault friction. A necessary but not sufficient condition for the occurrence of this instability  
 123 is the reactivation of the fault. We say that a fault is reactivated when the shear stress  
 124 on the fault area is high enough for frictional slip to take place. However, this slip can  
 125 be slow (aseismic) or sudden and abrupt (seismic), depending on the amount of slip weakening  
 126 (Figure 1f). It can be shown that the condition for sudden, unstable slip is (Dieterich,  
 127 1979; Scholz, 2002; Kanamori & Brodsky, 2004; Stefanou, 2019):

$$k < k_c = \frac{\Delta F}{d_c} \quad (2)$$

128 where  $k$  is the apparent stiffness of the rocks surrounding the fault zone for the real system,  
 129 and  $k_c$  is the critical stiffness. In a real scenario, the drop of shear force is  $\Delta F = A^{\text{real}} \Delta \tau^{\text{real}}$ ,

130 where  $\Delta\tau^{\text{real}}$  is the apparent shear stress drop of the fault zone. The stiffness  $k$  is proportional  
 131 to the effective elastic shear modulus of the surrounding rocks  $G^{\text{real}}$ , and inversely proportional  
 132 to the fault length  $L_{\text{ac}}^{\text{real}}$ , i.e.  $k \propto G^{\text{real}}/L_{\text{ac}}^{\text{real}}$ . Moreover, according to Coulomb friction  
 133  $\tau = \mu\sigma'_{\text{n}}$ , where  $\sigma'_{\text{n}} = \sigma_{\text{n}} - p_{\text{w}}$  is the effective normal stress,  $\sigma_{\text{n}}$  is the total normal  
 134 stress, which is a fraction of the overburden load, depending on the tectonic setting, and  
 135  $p_{\text{w}}$  is the fluid pressure.

136 Therefore, fluid injections have a double effect. On the one hand, they can reactivate  
 137 a fault by increasing  $p_{\text{w}}$ , reduce friction and promote frictional slip (Cappa et al., 2019).  
 138 On the other hand, they can reduce  $k_c$  as they make the post-peak slope less steep (Scuderi  
 139 et al., 2017; Lockner et al., 1991) (Figure 1f). Consequently, in an earthquake mitigation  
 140 attempt, one could adjust fluid pressure in such a way to avoid sudden, seismic slip and  
 141 assure stable, creep-like rupture (Stefanou, 2019).

#### 142 **4 Earthquake mitigation by fault segmentation and fluid injections**

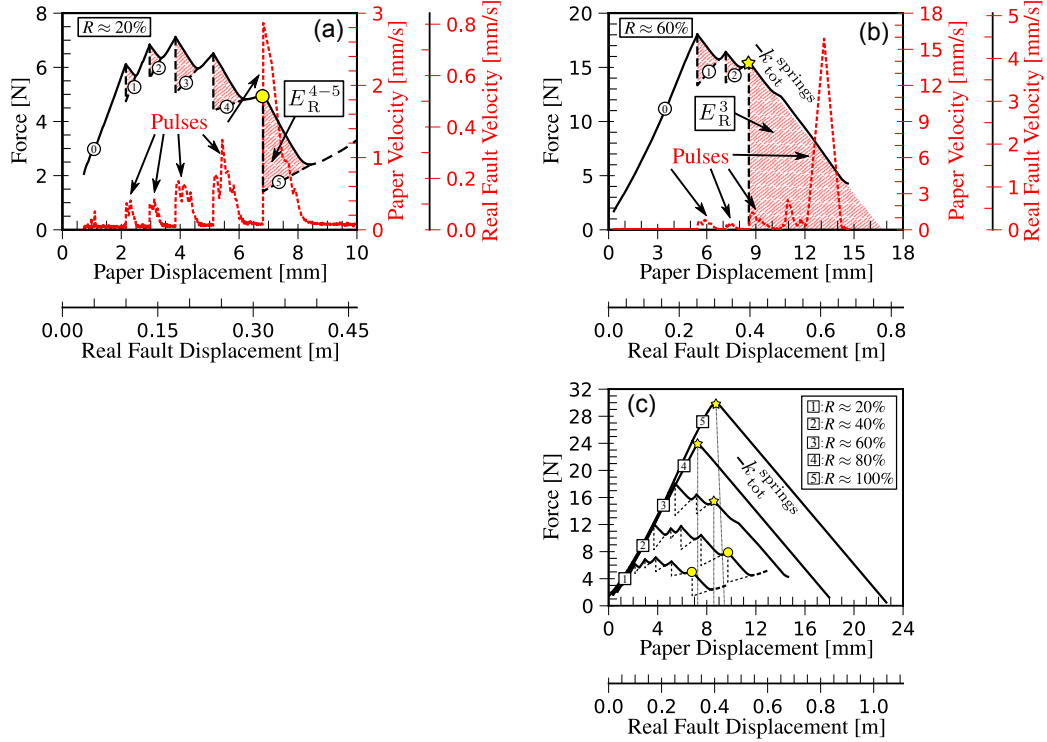
143 Controlling the fluid pressure simultaneously, across an entire fault of several kilometers,  
 144 seems impossible with current technologies. Yet, we could imagine to divide the potential  
 145 rupture area into several zones and inject fluids through a network of well-bores (Figure  
 146 1a). In this way, the energy stored in the system could be in theory released gradually,  
 147 mitigating the maximum earthquake magnitude.

148 The rupture area of our fault scenario is divided into five segments and so is the  
 149 paper sheet, as shown in Figures 1a-b. In order to have a better understanding of the  
 150 behavior of the segmented paper sheet (Figure 1b) under wetting, we tested first a single  
 151 segment (Figure S4a). In Figure 2c, we show the transition from seismic rupture to an  
 152 aseismic one by wetting. While the dry sample fails suddenly, liberating energy  $E_{\text{R}}^{\text{dry}} =$   
 153  $9.2 \text{ Nmm}$  (hatched area), the wet sample fails progressively with  $E_{\text{R}}^{\text{wet}} \approx 0$ , i.e. aseismically  
 154 (see also supporting information Movie S2).

155 Figure 2d corroborates the aseismic failure of the single stripe when wet. In particular,  
 156 the time-profiles of displacement and velocity are presented and compared for both dry  
 157 and wet samples. While the dry sample slips abruptly, the wet sample reaches the same  
 158 displacement in an almost constant, slow velocity, which is two orders of magnitude smaller  
 159 than the peak velocity of the dry case (see also supporting information Movie S2). If we  
 160 apply our scaling laws on the experimental data (see supporting information §6), the dry  
 161 single-stripe test gives an earthquake of magnitude  $M_{\text{w}} = 4.5$ , which is nucleated in  
 162 a normal fault with a length of  $L_{\text{ac}}^{\text{real}} = 1.3 \text{ km}$ .

163 Focusing on our surrogate experiment, five stripes (see supporting information Figure  
 164 S4b) are put in tension and are wetted subsequently to simulate multiple discrete fluid  
 165 injections. A certain amount of fluid is injected each time assuring fast saturation all over  
 166 the respective fault segment. Until the completion of the injection strategy, no more fluid  
 167 is injected on an already saturated segment. The injection program starts before reaching  
 168 instability, at a stress level  $R = 20\%$  (Figure 3a), where  $R$  is the ratio of the tensile force  
 169 at the first injection over the maximum tensile resistance. Wetting one stripe is performed  
 170 approximately every 1 mm of total displacement (segment-activation rate = 3 stripes/min).  
 171 This displacement corresponds to  $\sim 5 \text{ cm}$  of average slip over the real fault zone (see also  
 172 supporting information §4 and Movies S1 and S3).

173 Each fluid injection is accompanied by an instantaneous stress drop and stress redistribution  
 174 over the intact, dry porous paper stripes (Harris, 1998; Cappa et al., 2019) (Figure 3a).  
 175 In this setup, the redistribution is quasi-uniform and may not capture phenomena related  
 176 to real fault geometries. Due to these stress drops, energy is released abruptly, corresponding  
 177 to the triggering of small dynamic events (hatched areas). Ideally, we would like to minimize  
 178 their magnitudes or assure aseismic slip after each injection. The maximum magnitude



**Figure 3.** Wetting scenarios starting at  $R = 20\%$ , (a), and  $R = 60\%$ , (b). Force (equivalent friction, solid lines) and velocity (equivalent slip-rate, dotted lines) evolve with displacement (equivalent slip). Circled numbers depict the number of wetted paper stripes (reactivated fault segments). In (a), the maximum earthquake event takes place at the fifth injection (filled circle) and releases  $E_R^{4-5}$ . In (b), the sample fails dynamically at the third injection (star) and all remaining elastic energy is released abruptly (rupture of the entire fault area outside the injection area). (c) Force-paper displacement for scenarios started at different stress levels  $R$ . The star indicates that global failure occurs before wetting all the stripes (i.e. rupture outside the injection region).

179 among these smaller events characterizes the effectiveness of the earthquake mitigation  
 180 strategy.

181 Figure 3a shows the energy release and the developed velocities during this injection  
 182 program in our experiment. In the same plot we present also the corresponding magnitudes  
 183 of displacement and velocity that would develop in the real fault system. A sequence of  
 184 dynamic events are triggered with magnitudes:  $M_w^{0-1} = 3.9$ ,  $M_w^{1-2} = 3.9$ ,  $M_w^{2-3} =$   
 185  $4.3$ ,  $M_w^{3-4} = 4.4$  and  $M_w^{4-5} = 4.7$ . It is worth emphasizing that the released energy  
 186 increases with subsequent slip and it is maximum after the last injection ( $E_R^{4-5} = 2.94$  Nmm),  
 187 which is 68 times smaller than the energy of the large natural event. Therefore, we were  
 188 able to mitigate the initial natural earthquake event of magnitude  $M_w = 5.9$  to five  
 189 smaller earthquake events, whose maximum magnitude (filled circle in Figure 3a) is  $M_w =$   
 190  $4.7$ . In terms of velocities (dotted line in Figure 3a), after each injection, we observe a  
 191 distinct pulse corresponding to the released dynamic energy.

192 The system behaved differently when the injections started at a stress level ratio  
 193  $R = 60\%$  (Figure 3b, see also supporting information Movie S4). In this case, three dynamic  
 194 events are triggered with magnitudes:  $M_w^{0-1} = 4.7$ ,  $M_w^{1-2} = 4.5$  and  $M_w^3 = 5.5$ . A



195 large event followed the third injection (filled star in Figure 3b), leading to seismic rupture.  
 196 Energetically, this event ( $M_w^3 = 5.5$ ) is equivalent to the natural earthquake event ( $M_w =$   
 197  $5.9$ , see also Figures 2a-b).

198 In Figure 3c, we show the energy release for injections started at different stress  
 199 level ratios with the same segment-activation rate as before. In the case of  $R \approx 100\%$ ,  
 200 no injection can be carried out, as any tiny perturbation leads directly to an earthquake  
 201 nucleation (filled star) of magnitude  $M_w^0 = 5.9$  (Figures 2a-b). At  $R = 40\%$ , all five  
 202 injections are accomplished as in  $R = 20\%$ . These injections result in a series of five  
 203 induced earthquakes where their maximum magnitude (filled circle) is  $M_w^{4-5} = 4.8$ . When  
 204  $R = 80\%$ , though, a dynamic rupture occurs (filled star) after the first injection leading  
 205 to an earthquake event of  $M_w^1 = 5.8$ , similar to the case of  $R = 60\%$  (Figure 3b).

206 The experimental observations show clearly that as the stress level at the initiation  
 207 of the injection process becomes smaller, so does the magnitude of the subsequent events.  
 208 In other words, a sequence of earthquakes could be triggered, showing a maximum magnitude  
 209 which is at least one order of magnitude smaller than the natural earthquake.

210 Notice that our study is in agreement with the findings of the modeling work of  
 211 van der Elst et al. (2016). As we can observe in Figure 3c, the natural earthquake event  
 212 (curve 5) is the maximum event that can be nucleated in a fault in terms of energy, no  
 213 matter how much fluid volume has been injected into the fault and no matter when the  
 214 injection commences.

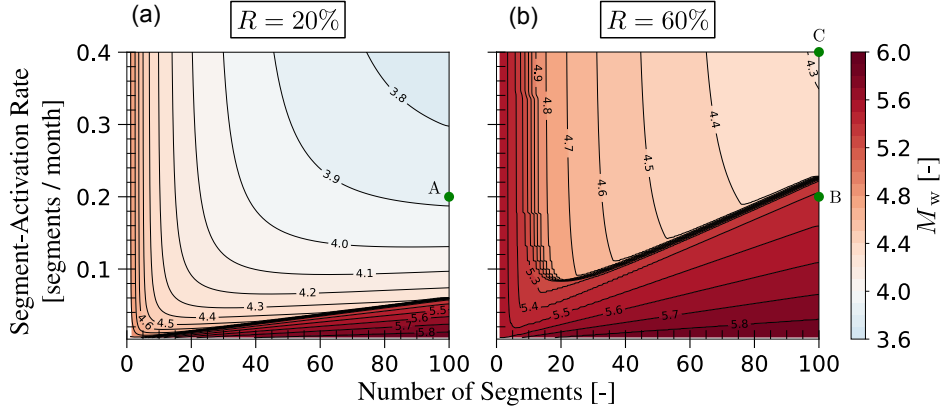
## 215 5 Discussion

216 The proposed analogy between uniaxial tension experiments of absorbent porous  
 217 paper and a real fault system is based on energy considerations. Our model is a pertinent  
 218 example for building understanding regarding possible mitigation of the earthquake phenomenon  
 219 (for a synthesis of the main assumptions and limitations we refer to the supporting information  
 220 §7). Intuitively, our approach could help to limit anthropogenic seismicity (Shapiro et  
 221 al., 2013) during fluid injections in the earth's crust, in parallel with the Traffic Light  
 222 System (TLS) method (Bommer et al., 2006; Edwards et al., 2015) used in deep geothermal  
 223 projects.

224 It is worth emphasizing that, according to our experiments, preceding small seismic  
 225 events do not guarantee the avoidance of large ones. Even though the released energy  
 226 in induced events is always smaller than the natural event, subsequent injections can drive  
 227 the system faster towards its instability point, provoking a large event (Figures 3b-c).  
 228 These events ( $R = 60\%$  and  $80\%$  in our experiments) would correspond to anthropogenic  
 229 seismicity when significant amounts of fluid are injected in the earth's crust, close to critically  
 230 stressed tectonic faults. On the other hand, we showed scenarios of progressive wetting,  
 231 where the maximum seismic moment could be reduced by one order of magnitude (Figures  
 232 3a and c). An important aspect, necessary for the possible mitigation in this scenario,  
 233 is the relatively low initial stress level ( $R = 20\%$  and  $40\%$  in our experiments). Therefore,  
 234 the earthquake mitigation strategy can succeed only if the injection process starts at relatively  
 235 low stress levels.

236 Besides the in-situ stress level, the paper experiments uncovered two additional factors  
 237 which govern the magnitude of induced events. These factors are the segment-activation  
 238 rate, which expresses the number of segments that are wetted per unit of time, and the  
 239 number of segments that our samples are divided into. In order to explore the response  
 240 of the fault system under the variation of these two additional factors, we use a multi  
 241 spring-slider model (see supporting information §3-4).

242 The segment-activation rate can be seen as the rate under which we force the system  
 243 to release its internal (potential/elastic) energy. According to our experiments and this



**Figure 4.** The maximum magnitude is plotted in function of the segment-activation rate, number of fault segments and two different stress levels, based on our fault scenario: (a)  $R = 20\%$  and (b)  $R = 60\%$ .

244 model, this rate has to be fast enough to outpace the progressive energy build-up due  
 245 to the far-field tectonic displacements. When our segment-activation rate is fast enough,  
 246 the system releases its internal energy and the large event is avoided. It seems that this  
 247 is the case in many industrial projects that involve injections of large amounts of fluids  
 248 in the earth's crust (McGarr, 2014; Guglielmi et al., 2015; Grigoli et al., 2018).

249 According to Hosseini et al. (2018), fluid boundary conditions play an important  
 250 role in induced seismicity. In our experimental work, fluid boundary conditions are expressed  
 251 through segmentation by adjusting plastic (impermeable) barriers between the stripes  
 252 (see also supporting information §4). Each injection leads to only one seismic event resulting  
 253 in a linear relation between the number of cumulative events in the fault and the time,  
 254 presuming constant segment-activation rate. Eventually, the number of segments, that  
 255 a fault has been divided into, depicts the maximum number of stimulated events.

256 In the absence of impermeable barriers, high number of segments would represent  
 257 better the physical reality, where the distribution of injection pressure does not occur  
 258 instantly over a region (paper segment), but follows a diffusion process (Bhattacharya  
 259 & Viesca, 2019). The diffusion process due to fluid injection in faults is equivalent to the  
 260 progressive wetting of paper stripes. If we assume, for instance, a high-permeable damage  
 261 fault zone with hydraulic diffusivity of the surrounding rocks of the order of  $10^{-1}$  to  $10^1$  m<sup>2</sup>/s  
 262 (values taken from Lim et al., 2020), the time it takes for each segment ( $1.3 \times 6.5$  km<sup>2</sup>)  
 263 to be saturated after fluid injection ranges between 10 to 980 days, respectively. In Figure  
 264 3a, where the experimental results of the proposed injection strategy are presented, the  
 265 segment-activation rate is 3 stripes/min which corresponds to 0.01 segments/month in  
 266 a real case scenario. Therefore, the above evaluated diffusion time is sufficient enough  
 267 in order to assure fast saturation of each segment.

268 Figure 4 presents the computed magnitude of the maximum earthquake event that  
 269 would occur, as a function of segment-activation rate and number of segments. For instance,  
 270 under a given rate of one injection per five months and for  $R = 20\%$ , one obtains a maximum  
 271  $M_w \approx 3.9$ , when segmenting the fault into 100 parts (point A, Figure 4a). For  $R =$   
 272  $60\%$ , though, the injection program leads to a large event of  $M_w \approx 5.4$ , close to the  
 273 natural one of  $M_w \approx 5.9$  (point B, Figure 4b). Doubling the rate from one to two injections  
 274 per five months reduces the maximum event to  $M_w \approx 4.3$  (point C, Figure 4b).

275 Despite the numerous uncertainties in the properties of the earth's crust (Cornet,  
 276 2019), our experiments (Figure 3c) and model (Figure 4) reveal the strong dependency  
 277 between the outcome of the injection strategy and the three aforementioned key parameters  
 278 (initial stress level, segment-activation rate and segmentation). By adequately controlling  
 279 these parameters, we managed to artificially reduce the stored elastic energy in an analogue  
 280 tectonic fault. However, in practice these parameters are hard to control and other strategies  
 281 based on the mathematical theory of control could provide rigorous alternatives (Stefanou,  
 282 2019, 2020).

## 283 Data Availability Statement

284 The raw data from all the performed experiments of this study are publicly available in  
 285 Zenodo repository in Tzortzopoulos et al. (2020).

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