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TIME DEPENDENT QUANTUM PERTURBATIONS UNIFORM IN THE SEMICLASSICAL REGIME

FRANÇOIS GOLSE AND THIERRY PAUL

ABSTRACT. We present a time dependent quantum perturbation result, uniform in the Planck constant for potential whose gradient is bounded a.e.. We show also that the classical limit of the perturbed quantum dynamics remains in a tubular neighbourhood of the classical unperturbed one, the size of this neighbourhood being of the order of the square root of the size of the perturbation. We treat both Schrödinger and von Neumann-Heisenberg equations.

in memory of Arthur Wightman

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1. INTRODUCTION

Perturbation theory has a very special status in Quantum Mechanics. On one side, it is responsible to most of its more spectacular success, from atomic to nuclear physics. On the other side, it has a very peculiar epistemological status: it was while he was working with Max Born [B25] on the Bohr-Sommerfeld quantization of celestial perturbations series, as explicitly stated by Poincaré in his famous “Mémoires” [P1892], that Heisenberg went to the idea of replacing the commutative algebra of convolution - corresponding to multiple multiplications of Fourier series appearing in computations on action-angles variables - by the famous noncommutative algebra of matrices [H25].

After quantum mechanics was truly settled, perturbation theory took a completely different form, in the paradigm of functional analysis “à la Kato” and appeared then mostly in the framework of the so-called Rayleigh-Schrödinger series. A kind of paradox is that it took a long time to link back the Raileigh-Schrödinger series to the “original” formalism of quantization of, say, Birkhoff series [B28],

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though, in the mean time, the latter continued to be extensively used for applied purpose e.g. in heavy chemical computations.

It seems that Arthur Wightman proposed to several PhD students to work on this problem. One of the difficulty is that, starting with the second term of the Raileigh-Schrödinger expansion,

$$E_i^2 = \sum_k \frac{\langle \psi_i^0, V \psi_k^0 \rangle \langle \psi_k^0, V \psi_i^0 \rangle}{E_i^0 - E_k^0},$$

appear poles in the Planck constant, for example when the unperturbed eigenvalues are the one of the harmonic oscillator $E_i^0 = (i + \frac{1}{2})\hbar$. Although this pole disappears at the classical limit $\hbar \rightarrow 0$ because the sum $\sum_k \frac{\langle \psi_i^0, V \psi_k^0 \rangle \langle \psi_k^0, V \psi_i^0 \rangle}{i-k}$ vanishes in this limit for parity reasons, controlling all the terms of the series remained for years a task considered as unachievable.

To our knowledge, the first proof on the convergence term by term of the Raileigh-Schrödinger expansion to the quantized Birkhoff one, for perturbations of non-resonant harmonic oscillators, was given in [G87], by implementing the perturbation procedure in the so-called Bargman representation (see also [D91] for an implementation in the framework of the Lie method). The reader interested in this subject can also consult [P16, P162] for a proof (also for general non harmonic unperturbed Hamiltonians) in a generalization of Écalle's mould theory and [NPST18] for a link between Raileigh-Schrödinger expansion and Hopf algebras.

When one considers time dependent perturbation theory, i.e. comparison between two quantum evolution associated to two "close" Hamiltonians H, H' , the situation is more difficult. The simple Duhamel formula

$$e^{-i\frac{tH}{\hbar}} - e^{-i\frac{tH'}{\hbar}} = \frac{1}{i\hbar} \int_0^t e^{-i\frac{(t-s)H}{\hbar}} (H' - H) e^{-i\frac{sH'}{\hbar}} ds$$

shows clearly that a pole in the Planck constant is again involved. But to our knowledge, no combinatorics or normal form can help to remove it in general and one si usually reduced to the trivial estimate

$$\|e^{-i\frac{tH}{\hbar}} - e^{-i\frac{tH'}{\hbar}}\| \leq t \frac{\|H' - H\|}{\hbar}$$

valid for, e.g. any Schatten norm, the operator, Hilbert-Schmidt or trace norm for example.

In the present paper, we will get rid of this pole in \hbar phenomenon by estimating the difference between two quantum evolutions (in a weak topology consisting in tracing against a set of test observables) in two forms:

- one linear in the norm of the difference of the Hamiltonians plus a term vanishing with \hbar
- the other proportional to the norm of the difference of the Hamiltonians to the power 1/3 and independent of \hbar .

The proofs of our results, Sections 5 and ??, will be using the framework of the von Neumann-Heisenberg equation for density operators D ,

$$\partial_t D = \frac{1}{i\hbar} [D, H],$$

but our results, Theorem 2.1 and Corollary 2.2, will be first presented for pure states, Section 2, that is when $D = |\psi\rangle\langle\psi|$, in which case it reduces to the usual Schrödinger one (modulo a global phase of the wave function)

$$i\hbar\partial_t\psi = H\psi.$$

The mixed state situation will be treated in Section 4, Theorem 4.1.

Our results will need very low regularity of the perturbed potential, namely the boundness of its gradient, and of the unperturbed one, Lipschitz continuity of its gradient. In this situation, the classical underlying dynamics is well posed for the unperturbed Hamiltonian, but not for the perturbed one. To our knowledge, the classical limit for pure state in this perturbed situation is unknown. We show in Section 3 that the limit as $\hbar \rightarrow 0$ of the Wigner function of the wave function at time t is close to the one of the initial state pushed forward by the unperturbed classical flow, Theorem 3.1.

2. MAIN RESULT

For $\lambda, \mu \in [0, 1]$, let us consider the quantum Hamiltonian

$$\mathcal{H}_0^{\lambda, \mu} = \mathcal{H}_0 := -\frac{1}{2}\hbar^2\Delta_x + \frac{\lambda}{2}|x|^2 + \mu V$$

on $\mathfrak{H} := L^2(\mathbf{R}^d)$. Here $V \equiv V(x) \in \mathbf{R}$ such that $V \in C^{1,1}(\mathbf{R}^d)$. For any other real potential $U \in W^{1,\infty}(\mathbf{R}^d)$, we define, for $\epsilon \in [0, 1]$

$$\mathcal{H}_\epsilon^{\lambda, \mu} = \mathcal{H}_\epsilon := -\frac{1}{2}\hbar^2\Delta_x + \frac{\lambda}{2}|x|^2 + \mu V + \epsilon U.$$

Henceforth we denote

$$\mathcal{H} := \mathcal{H}_0^{1,0} = -\frac{1}{2}\hbar^2\Delta_x + \frac{1}{2}|x|^2 \text{ (harmonic oscillator)}$$

$$\mathcal{D}(\mathfrak{H}) := \{R \in \mathcal{L}^1(\mathfrak{H}) \text{ s.t. } R = R^* \geq 0 \text{ and } \text{trace}_{\mathfrak{H}}(R) = 1\} \text{ (density operators),}$$

$$\mathcal{D}_2(\mathfrak{H}) := \{R \in \mathcal{D}(\mathfrak{H}) \text{ s.t. } \text{trace}_{\mathfrak{H}}(R^{1/2}\mathcal{H}R^{1/2}) < \infty\} \text{ (finite second moments).}$$

For $\psi \in \mathfrak{H}$, we define

$$(1) \quad \Delta(\psi) := \sqrt{(\psi, (x - (\psi, x\psi))^2 + (-i\hbar\nabla_x - (\psi, -i\hbar\nabla_x\psi))^2\psi)}$$

Note that the Heisenberg inequalities

$$\sqrt{(\psi, (x_k - (\psi, x_k\psi))^2, \psi)} \cdot \sqrt{(\psi, (-i\hbar\nabla_{x_k} - (\psi, -i\hbar\nabla_{x_k}\psi))^2\psi)} \geq \hbar/2, \quad k = 1, \dots, d,$$

imply that

$$\Delta(\psi) \geq \sqrt{2d\hbar}.$$

On $\mathcal{D}(\mathfrak{H})$ we define the following distance

$$(2) \quad \mathbf{d}(R, S) := \sup_{\substack{|\alpha|, |\beta| \leq 2[\frac{d}{4}] + 3 \\ \|\mathcal{D}_{i\hbar\nabla}^\alpha \mathcal{D}_x^\beta F\|_1 \leq 1}} |\text{trace}(F(R - S))|,$$

where $\mathcal{D}_A = \frac{1}{i\hbar}[A, \cdot]$ for each (possibly unbounded) self-adjoint operator A on \mathfrak{H} and $\|\cdot\|_1$ is the trace norm on $\mathcal{D}(\mathfrak{H})$. The fact that \mathbf{d} is a distance has been proved in [GJP20, Appendix A].

By abuse of notation, we will call for $\psi, \varphi \in \mathfrak{H}$

$$\mathbf{d}(\psi, \varphi) := \mathbf{d}(|\psi\rangle\langle\psi|, |\varphi\rangle\langle\varphi|)$$

Consider the family of Schrödinger equations, for $\epsilon \in [0, 1]$,

$$(3) \quad i\hbar\partial_t\psi_\epsilon(t) = \mathcal{H}_\epsilon, \psi_\epsilon(t) \quad \psi_\epsilon(0) = \psi_\epsilon^{in} \in H^2(\mathbf{R}^d),$$

Theorem 2.1. *Let ψ_ϵ^{in} satisfy the following hypothesis:*

$$(4) \quad \Delta(\psi^{in}) = O(\sqrt{\hbar}).$$

Then, for every t ,

$$\mathbf{d}(\psi_0(t), \psi_\epsilon(t))^2 \leq C(t)\epsilon + D(t)\hbar,$$

where $C(t), D(t)$, given by (19),(20), satisfy

$$\begin{aligned} C(t) &= \frac{e^{|t|(1-\lambda+\mu\text{Lip}(\nabla V))} - 1}{1 - \lambda + \mu\text{Lip}(\nabla V)} C_{(\psi_\epsilon^{in}, \mathcal{H}\psi_\epsilon^{in}), \|V\|_\infty, \|U\|_\infty, \|\nabla U\|_\infty} < \infty \\ D(t) &= e^{|t|(1-\lambda+\mu\text{Lip}(\nabla V))} D_d < \infty \end{aligned}$$

The following result gives an upper bound independent of \hbar .

Corollary 2.2.

$$\mathbf{d}(\psi_0(t), \psi_\epsilon(t)) \leq E(t)\epsilon^{\frac{1}{3}},$$

with

$$E(t) = \min(\sqrt{C(t) + D(t)}, 2|t|\|U\|_\infty^2).$$

Note that when $\lambda = 1$, $\text{Lip}(\nabla V) = 0$ (perturbation of the harmonic oscillator), $C(t)$ increases linearly in time and $D(t)$ is independent of time and $C(t), E(t)$ increase linearly in time.

Remark 2.3. *Other choices than the hypothesis (i) are possible, that we didn't mention for sake of clarity of the main statements. For example*

$$(i)' \quad \Delta(\psi^{in}) = O(\epsilon):$$

in this case the statement of both Theorem 2.1 and Corollary 2.2 remain the same with a slight change of the constants $C(t), D(t)$.

$$(i)'' \quad \Delta(\psi^{in}) = O(\hbar^\alpha), \quad 0 \leq \alpha < 1:$$

in this case the statement of both Theorem 2.1 and Corollary 2.2 become

$$\begin{aligned} \mathbf{d}(\psi_0(t), \psi_\epsilon(t))^2 &\leq C(t)\epsilon + D'(t)\hbar^\alpha \\ \mathbf{d}(\psi_0(t), \psi_\epsilon(t)) &\leq E'(t)\epsilon^{\frac{\alpha}{\alpha+2}} \end{aligned}$$

for constants $C(t), D(t), E(t)$ easily computable from the proofs of Section 5.

Let us finish this section by some topological remarks, inspired by [GP18, Section 4]. The distance \mathbf{d} defines a weak topology, very different a priori of the usual strong topologies associated to Hilbert spaces in quantum mechanics. Nevertheless, it seems to us better adapted to the semiclassical approximation for the following reason.

Let us consider two coherent states pinned up at two points $z_1 = (p_1, q_1), z_2 = (p_2, q_2)$ of phase -space $T^*\mathbf{R}^d$: $\psi_{z_j}(x) = (\pi\hbar)^{-d/4} e^{-i\frac{p_j \cdot x}{\hbar}} e^{-\frac{(x-q_j)^2}{2\hbar}}$, $j = 1, 2$.

An easy computation shows that

$$\|\psi_{z_1} - \psi_{z_2}\|_{L^2(\mathbf{R}^d)}^2 = \|\psi_{z_1}\rangle\langle\psi_{z_1}| - |\psi_{z_2}\rangle\langle\psi_{z_2}|\|_{Hilbert-Schmidt}^2 = 1 - e^{-\frac{(z_1 - z_2)^2}{2\hbar}}$$

so that, as $\hbar \rightarrow 0$,

$$\begin{aligned} \|\psi_{z_1} - \psi_{z_2}\|_{L^2(\mathbf{R}^d)}^2 &= \|\psi_{z_1}\rangle\langle\psi_{z_1}| - |\psi_{z_2}\rangle\langle\psi_{z_2}|\|_{Hilbert-Schmidt}^2 &= 0 & \text{if } z_1 = z_2 \\ & & \rightarrow 1 & \forall z_1 \neq z_2. \end{aligned}$$

In other words, the Lebesgue or Schatten norms behave for small values of \hbar as the discrete topology, the one which only discriminate points.

At the contrary, \mathbf{d} is much more sensitive to the localization on phase space as shows our next result, proven in Section 7 below.

Proposition 2.4. *For any bounded convex domain $\Omega \subset \mathbf{R}^{2d}$, there exists $C_\Omega > 0$ such that, for any $z_1, z_2 \in \Omega$,*

$$C_\Omega |z_1 - z_2| - \hbar \leq \mathbf{d}(\psi_{z_1}, \psi_{z_2}) \leq \sqrt{|z_1 - z_2|^2 + 2d\hbar} + C_d \hbar,$$

where C_d is defined in Lemma 5.3 Section 5 below.

3. APPLICATIONS TO THE CLASSICAL LIMIT

The estimates provided by the results of the two preceding sections do not require ∇U to be Lipschitz continuous — in other words, the classical dynamics underlying the quantum dynamics generated by \mathcal{H}_ϵ , $\epsilon > 0$, may fail to satisfy the assumptions of the Cauchy-Lipschitz theorem.

Let us recall that one way to look at the transition from quantum to classical dynamics as $\hbar \rightarrow 0$ is to associate to a quantum (pure or mixed) state, namely a positive trace one operator R^\hbar on \mathfrak{H} (density operator) with integral kernel $r^\hbar(x, x')$, e.g. a pure state $R^\hbar = |\psi^\hbar\rangle\langle\psi^\hbar|$, for any vector ψ^\hbar in \mathfrak{H} the so-called Wigner transform defined on phase-space by (with a slight abuse of notation again)

$$(5) \quad W_\hbar[R^\hbar](x, \xi) := \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} e^{-i\xi \cdot y} r^\hbar\left(x + \frac{1}{2}\hbar y, x - \frac{1}{2}\hbar y\right) dy$$

$$(6) \quad W_\hbar[|\psi^\hbar\rangle\langle\psi^\hbar|](x, \xi) := \frac{1}{(2\pi)^d} \int_{\mathbf{R}^d} e^{-i\xi \cdot y} \overline{\psi^\hbar\left(x + \frac{1}{2}\hbar y\right)} \psi^\hbar\left(x - \frac{1}{2}\hbar y\right) dy.$$

An easy computation shows that $W_\hbar[|\psi^\hbar\rangle\langle\psi^\hbar|]$ is linked to ψ^\hbar by the two following marginal properties

$$(7) \quad \int_{\mathbf{R}^d} W_\hbar[|\psi^\hbar\rangle\langle\psi^\hbar|](x, \xi) d\xi = |\psi^\hbar(x)|^2$$

$$(8) \quad \int_{\mathbf{R}^d} W_\hbar[|\psi^\hbar\rangle\langle\psi^\hbar|](x, \xi) dx = |\widehat{\psi^\hbar}(p)|^2, \quad \widehat{\psi^\hbar}(p) := \int_{\mathbf{R}^d} e^{-ip \cdot x/\hbar} \psi^\hbar(x) \frac{dx}{(2\pi\hbar)^{d/2}}$$

It has been proved, see e.g. [LP93], that, under the tightness conditions

$$(9) \quad \lim_{R \rightarrow +\infty} \sup_{\hbar \in (0,1)} \int_{\mathbf{R}^d \setminus B_R^{(d)}} r^\hbar(x, x) dx = 0,$$

$$(10) \quad \lim_{R \rightarrow +\infty} \sup_{\hbar \in (0,1)} \frac{1}{(2\pi\hbar)^d} \int_{\mathbf{R}^d \setminus B_R^{(d)}} \mathcal{F}r^\hbar\left(\frac{p}{\hbar}, \frac{p}{\hbar}\right) dp = 0,$$

where $B_R^{(d)}$ is the ball of radius R in \mathbf{R}^d and \mathcal{F} is the Fourier transform on \mathbf{R}^{2d} , the family of Wigner functions $W_\hbar[R_\hbar]$ converges weakly, after extraction of a subsequence of values of \hbar to $W_0 \in \mathcal{P}(\mathbf{R}^{2d})$, the space of probability measures on \mathbf{R}^{2d} . W_0 is called the Wigner measure of the family R_\hbar .

As it is quite standard, we will omit to mention the extraction of subsequences, together with the explicit dependence of states in the Planck constant, and we will just write, when it doesn't create any confusion,

$$\lim_{\hbar \rightarrow 0} W_\hbar[R] = W_0.$$

Note that when $R = |\psi\rangle\langle\psi|$ is a pure state, (9)-(10) reads

$$(11) \quad \lim_{R \rightarrow +\infty} \sup_{\hbar \in (0,1)} \int_{\mathbf{R}^d \setminus B_R^{(d)}} |\psi(x)|^2 dx = 0,$$

$$(12) \quad \lim_{R \rightarrow +\infty} \sup_{\hbar \in (0,1)} \int_{\mathbf{R}^d \setminus B_R^{(d)}} |\widehat{\psi}(p)|^2 dp = 0$$

Considering the quantum Hamiltonian \mathcal{H}_ϵ , the expected underlying classical dynamics is the one driven by the Liouville equation

$$(13) \quad \partial_t \rho = \left\{ \frac{1}{2}(p^2 + \lambda q^2) + \mu V(q) + \epsilon U, \rho \right\}, \quad \rho^t|_{t=0} = \rho^{in}$$

where $\{.,.\}$ is the Poisson bracket on the symplectic manifold $T^*\mathbf{R}^d \sim \mathbf{R}^{2d}$.

When $\epsilon = 0$, the Hamiltonian vector field of Hamiltonian $\frac{p^2 + \lambda q^2}{2} + \mu V(q)$ is Lipschitz continuous $C_0([0, T]; \mathcal{P}(\mathbf{R}^{2d}))$. Moreover, it was proven, [LP93], that R_\hbar^t is tight for any $t \in \mathbf{R}$ and $W_0^t := \lim_{\hbar \rightarrow 0} W_\hbar[R_\hbar(t)] = \rho^t$ solving (13) with $\rho^{in} = W_0$.

When $\epsilon > 0$, the Liouville equation (13) exits the Cauchy-Lipschitz category: the associated Hamiltonian vector field might not have a unique characteristic out of every point of the phase-space. Nevertheless, as shown in [AFFGP10, Theorem 6.1], (13) is still well posed in $L^+([0, T]; L^1(\mathbf{R}^{2d}) \cap L^\infty(\mathbf{R}^{2d}))$, and it was proven in [FLP13] (after [AFFGP10], that the Wigner function $W_\hbar[R_\hbar(t)]$ of the solution of the von Neumann equation

$$i\hbar \partial_t R_\epsilon(t) = [\mathcal{H}_\epsilon, R_\epsilon(t)], \quad R_\epsilon(0) = R_\epsilon^{in}$$

tends weakly to the solution of (13), under certain conditions on R_ϵ^{in} .

Unfortunately, these conditions exclude definitively pure states, as, for example, one of them impose that $\|R_\epsilon^{in}\| = O(\hbar^d)$ and, to our knowledge, nothing is known concerning the dynamics of the (possible) limit of $W_\hbar[\psi_\epsilon(t)]$ as $\hbar \rightarrow 0$ where $\psi_\epsilon(t)$ solves the Schrödinger equation (15).

Our next result will show that such a limit remains $\sqrt{\epsilon}$ -close to the pushforward of the Wigner measure of the initial condition by the flow of the unperturbed classical Hamiltonian.

Theorem 3.1. *Let $R_\epsilon(t)$ be the solution of the Schrödinger equation (15) with R_ϵ^{in} satisfying*

$$\Delta(R_\epsilon^{in}) = O(\sqrt{\hbar}).$$

Let R_ϵ^{in} be tight, in the sense that it satisfies (9)-(10), so that $W_\hbar[R_\epsilon^{in}] \rightarrow W_0^{in}$ as $\hbar \rightarrow 0$, its Wigner measure.

Then, for any $t \in \mathbf{R}$,

(1) the family $R_\epsilon(t)$ is tight, so that

$$W_\hbar[R_\epsilon(t)] \rightarrow W_0(t) \in \mathcal{P}(\mathbf{R}^{2d}) \text{ as } \hbar \rightarrow 0$$

(2) $W_0(t)$ is $\sqrt{\epsilon}$ -close to $\Phi^t \# W_0^{in}$, where Φ^t is the flow of Hamiltonian

$$\frac{p^2 + (1-\lambda)q^2}{2} + \mu V(q),$$

in the sense that

$$\int_{\mathbf{R}^d} \sup_{q \in \mathbf{R}^d} |\mathcal{F}_p f(q, z)| dz \leq 1 \quad \left| \int_{\mathbf{R}^{2d}} f(q, p) (W_0(t) - \Phi^t \# W_0^{in})(q, p) dq dp \right| \leq 2^{-d} \sqrt{C(t)} \sqrt{\epsilon},$$

$$\max_{|\alpha|+|\beta| \leq 2[d/4]+3} \|\partial_q^\alpha \partial_p^\beta f\|_{L^\infty(\mathbf{R}^{2d})} \leq 1$$

where $\mathcal{F}_p f(q, z) := \int_{\mathbf{R}^d} e^{-ip \cdot z} f(q, p) dp$ and $C(t)$ is as in Theorem 2.1 after replacing $(\psi_\epsilon^{in}, \mathcal{H}\psi_\epsilon^{in})$ by $\int_{\mathbf{R}^{2d}} (\frac{p^2+q^2}{2}) W_0(dp dq)$.

(3) in particular, $W_0(t)$ is weakly $\sqrt{\epsilon}$ -close to $\Phi^t \# W_0^{in}$ in the sense of distribution as, for all test functions $\varphi \in \mathcal{S}(\mathbf{R}^{2d})$,

$$\left| \int_{\mathbf{R}^{2d}} \varphi(q, p) (W_0(t) - \Phi^t \# W_0^{in})(dp, dq) \right| \leq C \varphi(t) \sqrt{\epsilon}$$

with

$$C_\varphi(t) = \max \left(\int_{\mathbf{R}^d} \sup_{q \in \mathbf{R}^d} |\mathcal{F}_p f(q, z)| dz, \max_{|\alpha|+|\beta| \leq 2[d/4]+3} \|\partial_q^\alpha \partial_p^\beta f\|_{L^\infty(\mathbf{R}^{2d})} \right) 2^{-d} \sqrt{C(t)}.$$

Proof. (1) the propagation of tightness is proved as follows.

Let $\chi \in C^\infty(\mathbf{R}^d)$, $0 \leq \chi \leq 1$ such that $\chi(x) = 0$ if $|x| < 1/2$ and $\chi(x) = 1$ if $|x| > 1$, and define $\chi_R(x) := \chi(x/R)$. Obviously

$$\int_{\mathbf{R}^d \setminus B_R^{(d)}} |\psi_\epsilon(t)(x)|^2 dx \leq \int_{\mathbf{R}^d} \chi_R(x) |\psi_\epsilon(t)(x)|^2 dx$$

Moreover, for some $C > 0$, $\|\nabla \chi_R\|_\infty, \|\Delta \chi_R\|_\infty \leq C/R^2$, and

$$-\frac{i}{\hbar} [\chi_R, \mathcal{H}_\epsilon] = -\frac{i}{\hbar} [\chi_R, -\frac{\hbar^2}{2} \Delta] = -i\hbar (\frac{1}{2} \Delta \chi_R - i \nabla \chi_R \cdot \nabla).$$

Therefore

$$\begin{aligned} \partial_t \int_{\mathbf{R}^d} \chi_R(x) |\psi_\epsilon(t)(x)|^2 dx &= -i\hbar \int_{\mathbf{R}^d} \bar{\psi}_\epsilon(t)(x) ((\frac{1}{2} \Delta \chi_R - i \nabla \chi_R \cdot \nabla) \psi_\epsilon(t))(x) dx \\ &= \int_{\mathbf{R}^d} (-i\hbar \frac{1}{2} \Delta \chi_R(x) |\psi_\epsilon(t)|^2 \\ &\quad + \bar{\psi}_\epsilon(t)(x) \nabla \chi_R(x) \cdot (-i\hbar \nabla \psi_\epsilon(t)(x))) dx, \end{aligned}$$

so that

$$\begin{aligned}
\partial_t \int_{\mathbf{R}^d} \chi_R(x) |\psi_\epsilon(t)(x)|^2 dx &\leq \hbar \frac{C}{2R^2} + \frac{C}{R} \| -i\hbar \nabla \psi_\epsilon(t) \|_{L^2(\mathbf{R}^d)} \\
&= \hbar \frac{C}{2R^2} + 2 \frac{C}{R} (\psi_\epsilon(t), \mathcal{H}_0^{0,0} \psi_\epsilon(t))_{L^2(\mathbf{R}^d)} \\
&\leq \hbar \frac{C}{2R^2} + 2 \frac{C}{R} ((\psi_\epsilon(t), \mathcal{H}_\epsilon^{\lambda,\mu} \psi_\epsilon(t))_{L^2(\mathbf{R}^d)} + \mu \|V\|_\infty + \epsilon \|vd\|_\infty) \\
&= \hbar \frac{C}{2R^2} + 2 \frac{C}{R} ((\psi_\epsilon^{in}, \mathcal{H}_\epsilon^{\lambda,\mu} \psi_\epsilon^{in})_{L^2(\mathbf{R}^d)} + \mu \|V\|_\infty + \epsilon \|vd\|_\infty)
\end{aligned}$$

and finally, for $t \in [0, T]$

$$\begin{aligned}
\int_{\mathbf{R}^d} \chi_R(x) |\psi_\epsilon(t)(x)|^2 dx &\leq \int_{\mathbf{R}^d} \chi_R(x) |\psi_\epsilon^{in}(x)|^2 dx \\
&\quad + \left(\hbar \frac{C}{2R^2} + 2 \frac{C}{R} ((\psi_\epsilon^{in}, \mathcal{H}_\epsilon^{\lambda,\mu} \psi_\epsilon^{in})_{L^2(\mathbf{R}^d)} + \mu \|V\|_\infty + \epsilon \|vd\|_\infty) \right) T.
\end{aligned}$$

Therefore $\psi_\epsilon(t)$ satisfies (12) as soon as ψ_ϵ^{in} does.

Finally, let us remark that

$$\begin{aligned}
\int_{\mathbf{R}^d \setminus B_R^{(d)}} |\widehat{\psi}(p)|^2 dp &\leq \frac{1}{R^2} \int_{\mathbf{R}^d} p^2 |\widehat{\psi}(p)|^2 dp \\
(14) \qquad \qquad \qquad &= \frac{2}{R^2} (\psi_\epsilon(t), \mathcal{H}_0^{0,0} \psi_\epsilon(t))_{L^2(\mathbf{R}^d)}
\end{aligned}$$

and one concludes the same way.

- (2) one knows from [LP93] that the convergence of Wigner functions to Wigner measure as $\hbar \rightarrow 0$ happens in the dual of the set of test functions f on \mathbf{R}^{2d} satisfying

$$\int_{\mathbf{R}^d} \sup_{q \in \mathbf{R}^d} |\mathcal{F}_p f(q, z)| dz < \infty.$$

Since $V \in C^{1,1}$ one knows that, for such a test function,

$$\lim_{\hbar \rightarrow 0} \int_{\mathbf{R}^{2d}} f(x, \xi) W_\hbar[\psi_0(t)](x, \xi) dx d\xi = \int_{\mathbf{R}^{2d}} f(x, \xi) \Phi^t \# W_0(x, \xi) dx d\xi.$$

On the other site, we have the slight variant of Theorem 2.1, proven also in Section 5.

Proposition 3.2. *Let δ be defined by (17) below. Then*

$$\delta(W_\hbar[\psi_0(t)], W_\hbar[\psi_\epsilon(t)]) \leq 2^{-d} \sqrt{C(t)\epsilon + D(t)\hbar},$$

where $C(t), D(t)$ are the constants defined in Theorem 2.1.

Proposition 3.2 tells us that, for any f satisfying

$$\max_{|\alpha|+|\beta| \leq 2[d/4]+3} \|\partial_q^\alpha \partial_p^\beta f\|_{L^\infty(\mathbf{R}^{2d})} \leq 1,$$

$$\left| \int_{\mathbf{R}^{2d}} f(q, p) (W_\hbar[\psi_\epsilon(t)] - W_\hbar[\psi_0(t)])(dp dq) \right| \leq 2^{-d} \sqrt{C(t)\epsilon + D(t)\hbar}$$

Hence for any f satisfying

$$\int_{\mathbf{R}^d} \sup_{q \in \mathbf{R}^d} |\mathcal{F}_p f(q, z)| dz \leq 1, \quad \max_{|\alpha|+|\beta| \leq 2[d/4]+3} \|\partial_q^\alpha \partial_p^\beta f\|_{L^\infty(\mathbf{R}^{2d})} \leq 1,$$

we have

$$\begin{aligned} & \left| \int_{\mathbf{R}^{2d}} f(q, p) (W_h[\psi_\epsilon(t)] - \Phi^t \# W_0(x, \xi)) (dpdq) \right| \\ & \leq 2^{-d} \sqrt{C(t)\epsilon + D(t)\hbar} + \left| \int_{\mathbf{R}^{2d}} f(q, p) (W_h[\psi_0(t)] - \Phi^t \# W_0(x, \xi)) (dpdq) \right| \end{aligned}$$

and we conclude by taking the supremum on the functions f and the limit $\hbar \rightarrow 0$ on both sides.

(3) the proof is obvious by homogeneity. \square

4. THE CASE OF MIXED STATES

Consider the family of von Neumann equations, for $\epsilon \in [0, 1]$,

$$(15) \quad i\hbar \partial_t R_\epsilon(t) = [\mathcal{H}_\epsilon, R_\epsilon(t)], \quad R_\epsilon(0) = R_\epsilon^{in} \in \mathcal{D}(\mathfrak{H}),$$

Let $R_\epsilon^{in} = R_0^{in} = R \in \mathcal{D}_2(\mathfrak{H})$ satisfy one of the five following hypothesis:

- (i) $\Delta(R^{in}) = O(\sqrt{\hbar})$ where the standard deviation $\Delta(R)$ is defined in Lemma 5.4 below.
- (ii) $\sqrt{R^{in}}$ satisfies, for some $C > 0$,

$$\sup_{|\beta_1|, \dots, |\beta_d| \leq 7} \left| \prod_{m=1}^d D_{(x, \xi)}^{\beta_m} W_h[\sqrt{R^{in}}](x, \xi) \right| \leq \frac{C(2\pi\hbar)^{-\frac{d}{2}}}{((\xi^2 + x^2)^2 + d)^{\frac{10}{4} + 3\epsilon}} \quad \forall (x, \xi) \in \mathbf{R}^{2d}.$$

where $W_h[\sqrt{R^{in}}]$ is the Wigner transform of R^{in}

- (iii) $\sqrt{R^{in}}$ satisfies, for $C > 0$ and all $j \in \mathbf{N}^d$,
 - (a) $|(H_i, \sqrt{R^{in}} H_j)| \leq C(2\pi\hbar)^{\frac{d}{2}} \prod_{1 \leq l \leq d} |\hbar j_l + \frac{1}{2}|^{-\frac{3}{4} - \epsilon} (|i_l - j_l| + 1)^{-2 - \epsilon}$,
 - (b) $\sup_{O \in \Omega_1} |(H_i, \frac{1}{i\hbar} [O, \sqrt{R^{in}}] H_j)| \leq C(2\pi\hbar)^{\frac{d}{2}} \prod_{1 \leq l \leq d} |\hbar j_l + \frac{1}{2}|^{-\frac{1}{2} - \epsilon} (|i_l - j_l| + 1)^{-1 - \epsilon}$,
 where $\Omega_1 = \{y_j, \pm \hbar \partial_{y_j} \text{ on } L^2(\mathbf{R}^d, dy), j = 1, \dots, d\}$ and the H_j s are the semiclassical Hermite functions..

(iv) R^{in} is a Töplitz operator.

(v) there exist a Töplitz operator T_F such that $T_F^{-\frac{1}{2}} \sqrt{R^{in}}$, $\sqrt{R^{in}} T_F^{-\frac{1}{2}}$ and $T_F^{-\frac{1}{2}} \sqrt{R^{in}} \frac{1}{i\hbar} [O, \sqrt{R^{in}} T_F^{-\frac{1}{2}}]$, $O \in \Omega_1$ are bounded on $L^2(\mathbf{R}^d)$.

Theorem 4.1. *For every $t \geq 0$,*

$$\mathbf{d}(R_0(t), R_\epsilon(t))^2 \leq C(t)\epsilon + D(t)\hbar,$$

where $C(t), D(t)$ are the same given by (19), (20), satisfy

$$\begin{aligned} C(t) &= \frac{e^{|t|(1-\lambda+\mu \text{Lip}(\nabla V))} - 1}{1 - \lambda + \mu \text{Lip}(\nabla V)} C_{\|\mathcal{H}R^{in}\|_1, \|V\|_\infty, \|U\|_\infty, \|\nabla U\|_\infty} < \infty \\ D(t) &= e^{|t|(1-\lambda+\mu \text{Lip}(\nabla V))} D_d < \infty \end{aligned}$$

Corollary 4.2.

$$\mathbf{d}(R_0(t), R_\epsilon(t)) \leq E(t)\epsilon^{\frac{1}{3}},$$

In Theorem 4.1 and Corollary 4.2, the constants $C(t), D(t), E(t)$ are the same as in Theorem 2.1 and Corollary 2.2 after replacing $(\sigma^{in}, \mathcal{H}\psi^{in})$ by $\|\mathcal{H}R^{in}\|_1$ in $C(t)$.

5. PROOF OF THEOREMS 2.1 AND 4.1, AND PROPOSITION 3.2

For all $R, S \in \mathcal{D}(\mathfrak{H})$, we denote by $\mathcal{C}(R, S)$ the set of couplings of R and S , i.e. $\mathcal{C}(R, S) := \{Q \in \mathcal{D}(\mathfrak{H} \otimes \mathfrak{H}) \text{ s.t. } \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}((A \otimes I + I \otimes B)Q) = \text{trace}_{\mathfrak{H}}(AR + BS)\}$. We recall the definition of the pseudo-distance MK_h (see Definition 2.2 in [GMP16]).

Definition 5.1. For each $R, S \in \mathcal{D}_2(\mathfrak{H})$,

$$MK_h(R, S) := \inf_{Q \in \mathcal{C}(R, S)} \sqrt{\text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q^{1/2} C Q^{1/2})},$$

where

$$C := \sum_{j=1}^d ((q_j \otimes I - I \otimes q_j)^2 - \hbar^2 (\partial_{q_j} \otimes I - I \otimes \partial_{q_j})^2).$$

Theorem 2.1 and Proposition 3.2 are a consequence of the following inequality, which controls the continuous dependence of the solution to the von Neumann equation in terms of the initial data and on the potential.

Theorem 5.2. Let $R_\epsilon^{in} \in \mathcal{D}_2(\mathfrak{H})$ and $R_\epsilon(t)$ be the solution of (15), $\epsilon \in [0, 1]$. Then, for each $t \in \mathbf{R}$, $\epsilon \in [0, 1]$, one has

$$\begin{aligned} MK_h(R_0(t), R_\epsilon(t))^2 &\leq e^{|\Lambda(\nabla V)|} MK_h(R_0^{in}, R_\epsilon^{in})^2 \\ &+ \epsilon \frac{e^{|\Lambda(\nabla V)|} - 1}{\Lambda(\nabla V)} \|\nabla U\|_{L^\infty(\mathbf{R}^d)} \sqrt{\text{trace}((R_0^{in})^{1/2} \mathcal{H}(R_0^{in})^{1/2}) + 2\mu \|V\|_{L^\infty(\mathbf{R}^d)}} \\ &+ \epsilon \frac{e^{|\Lambda(\nabla V)|} - 1}{\Lambda(\nabla V)} \|\nabla U\|_{L^\infty(\mathbf{R}^d)} \sqrt{\text{trace}((R_\epsilon^{in})^{1/2} \mathcal{H}(R_\epsilon^{in})^{1/2}) + 2(\mu \|V\|_{L^\infty(\mathbf{R}^d)} + \epsilon \|U\|_{L^\infty(\mathbf{R}^d)})}, \end{aligned}$$

where

$$\Lambda(\nabla V) = 1 - \lambda + \mu \text{Lip}(\nabla V).$$

Note that when $\lambda = 1$, $\Lambda(\nabla V) = \mu \text{Lip}(\nabla V)$ and in the inequality above, the function

$$z \mapsto \frac{e^z - 1}{z}$$

is extended by continuity at $z = 0$.

Proof. In order to lighten the formulas we will use the following notations

$$(16) \quad V_1 = \mu V, \quad V_2 = \mu V + \epsilon U,$$

so that

$$\mathcal{H}_0 = \mathcal{H}_0^{\lambda, 0} + V_1 \text{ and } \mathcal{H}_\epsilon = \mathcal{H}_0^{\lambda, 0} + V_2.$$

Let $Q^{in} \in \mathcal{C}(R_0^{in}, R_\epsilon^{in})$, and let Q be the solution of the von Neumann equation

$$i\hbar \partial_t Q = [(\mathcal{H}_\epsilon^{\lambda, 0} + V_1) \otimes I + I \otimes (\mathcal{H}_\epsilon^{\lambda, 0} + V_2), Q], \quad Q|_{t=0} = Q^{in}.$$

Then

$$Q(t) \in \mathcal{C}(R_0(t), R_\epsilon(t)), \quad \text{for each } t \geq 0$$

(see for instance Lemma 5.1 in [GMP16]).

Next we compute

$$\begin{aligned} &\frac{d}{dt} \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \\ &= \frac{i}{\hbar} \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} [(\mathcal{H}_\lambda + V_1) \otimes I + I \otimes (\mathcal{H}_\lambda + V_2), C] Q(t)^{1/2}). \end{aligned}$$

One finds that

$$\frac{i}{\hbar}[-\frac{1}{2}\hbar^2(\Delta \otimes I + I \otimes \Delta), C] = \sum_{j=1}^d (q_j \otimes I - I \otimes q_j) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})),$$

while

$$\frac{i}{\hbar}[(V_1 \otimes I + I \otimes V_2), C] = -\sum_{j=1}^d (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_2) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})),$$

with the notation

$$A \vee B := AB + BA.$$

In particular

$$\frac{i}{\hbar}[\frac{1}{2}(|q|^2 \otimes I + I \otimes |q|^2), C] = -\sum_{j=1}^d (q_j \otimes I - I \otimes q_j) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})).$$

See [GMP16] on p. 190. Hence

$$\begin{aligned} & \frac{d}{dt} \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \\ & - (1 - \lambda) \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}} \left(Q(t)^{1/2} \sum_{j=1}^d (q_j \otimes I - I \otimes q_j) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) Q(t)^{1/2} \right) \\ & = \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}} \left(Q(t)^{1/2} \sum_{j=1}^d (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_2) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) Q(t)^{1/2} \right) \\ & = \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}} \left(Q(t)^{1/2} \sum_{j=1}^d (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_1) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) Q(t)^{1/2} \right) \\ & \quad + \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}} \left(Q(t)^{1/2} \sum_{j=1}^d (I \otimes \partial_{q_j} (V_1 - V_2)) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) Q(t)^{1/2} \right) \\ & =: \tau_1 + \tau_2. \end{aligned}$$

At this point, we recall the elementary operator inequality

$$A^* B + B^* A \leq A^* A + B^* B.$$

Therefore,

$$\sum_{j=1}^d (q_j \otimes I - I \otimes q_j) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) \leq C$$

and, for each $\ell > \text{Lip}(\nabla V_1)^{1/2}$, one has

$$\begin{aligned} & \sum_{j=1}^d (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_1) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) \\ & = \sum_{j=1}^d \frac{1}{\ell} (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_1) \vee \ell (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) \\ & \leq \sum_{j=1}^d \left(\frac{\text{Lip}(\nabla V_1)^2}{\ell^2} (q_j \otimes I - I \otimes q_j)^2 + \ell^2 (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j}))^2 \right). \end{aligned}$$

Letting $\ell \rightarrow \text{Lip}(\nabla V)^{1/2}$ shows that

$$\sum_{j=1}^d (\partial_{q_j} V_1 \otimes I - I \otimes \partial_{q_j} V_1) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j})) \leq \text{Lip}(\nabla V_1)C,$$

so that

$$\tau_1 \leq \text{Lip}(\nabla V_1) \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}).$$

For the term τ_2 , we simply use the Cauchy-Schwarz inequality:

$$\begin{aligned} \tau_2 &= \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}} \left(Q(t)^{1/2} \sum_{j=1}^d (I \otimes \partial_{q_j} (V_1 - V_2) \vee (-i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j}))) Q(t)^{1/2} \right) \\ &\leq \sum_{j=1}^d \|Q(t)^{1/2} (I \otimes \partial_{q_j} (V_1 - V_2))\|_2 \| -i\hbar(\partial_{q_j} \otimes I - I \otimes \partial_{q_j}) Q(t)^{1/2} \|_2 \\ &\leq \sum_{j=1}^d \|Q(t)^{1/2} (I \otimes \partial_{q_j} (V_1 - V_2))\|_2 \|(-i\hbar \partial_{q_j} \otimes I) Q(t)^{1/2}\|_2 \\ &\quad + \sum_{j=1}^d \|Q(t)^{1/2} (I \otimes \partial_{q_j} (V_1 - V_2))\|_2 + \|(I \otimes (-i\hbar \partial_{q_j})) Q(t)^{1/2}\|_2 \\ &= \tau_{21} + \tau_{22}. \end{aligned}$$

Now

$$\begin{aligned} \tau_{21} &\leq \left(\sum_{j=1}^d \|Q(t)^{1/2} (I \otimes \partial_{q_j} (V_1 - V_2))\|_2^2 \right)^{1/2} \left(\sum_{j=1}^d \|(-i\hbar \partial_{q_j} \otimes I) Q(t)^{1/2}\|_2^2 \right)^{1/2} \\ &= \left(\sum_{j=1}^d \text{trace} (Q(t)^{1/2} (I \otimes \partial_{q_j} (V_1 - V_2))^2 Q(t)^{1/2}) \right)^{1/2} \\ &\quad \times \left(\text{trace} (Q(t)^{1/2} (-\hbar^2 \Delta \otimes I) Q(t)^{1/2}) \right)^{1/2} \end{aligned}$$

so that

$$\tau_{21} \leq \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \text{trace}(R_0(t)^{1/2} \mathcal{H}_\lambda R_0(t)^{1/2})^{1/2},$$

and likewise

$$\tau_{22} \leq \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \text{trace}(R_\epsilon(t)^{1/2} \mathcal{H}_\lambda R_\epsilon(t)^{1/2})^{1/2}.$$

Summarizing, we have proved that

$$\begin{aligned} \frac{d}{dt} \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) &\leq \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \\ &\quad + \text{Lip}(\nabla V_1) \text{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \\ &\quad + \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} (\text{trace}(R_0(t)^{1/2} \mathcal{H}_\lambda R_0(t)^{1/2})^{1/2} + \text{trace}(R_\epsilon(t)^{1/2} \mathcal{H}_\lambda R_\epsilon(t)^{1/2})^{1/2}). \end{aligned}$$

On the other hand, since

$$i\hbar \partial_t R_j(t) = [\mathcal{H}_\lambda + V_j, R_j(t)],$$

one has

$$\text{trace}(R_j(t)^{1/2} (\mathcal{H}_\lambda + V_j) R_j(t)^{1/2}) = \text{trace}((R_j^{in})^{1/2} (\mathcal{H}_\lambda + V_j) (R_j^{in})^{1/2})$$

so that

$$\text{trace}(R_j(t)^{1/2} \mathcal{H}_\lambda R_j(t)^{1/2}) \leq \text{trace}((R_j^{in})^{1/2} \mathcal{H}_\lambda (R_j^{in})^{1/2}) + 2\|V_j\|_{L^\infty(\mathbf{R}^d)}.$$

Hence

$$\begin{aligned} \frac{d}{dt} \operatorname{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) &\leq (1 - \lambda + \operatorname{Lip}(\nabla V_1)) \operatorname{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \\ &\quad + \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_0^{in})^{1/2} \mathcal{H}_\lambda (R_0^{in})^{1/2}) + 2\|V_1\|_{L^\infty(\mathbf{R}^d)}} \\ &\quad + \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_\epsilon^{in})^{1/2} \mathcal{H}_\lambda (R_\epsilon^{in})^{1/2}) + 2\|V_2\|_{L^\infty(\mathbf{R}^d)}}. \end{aligned}$$

By Gronwall's inequality, choosing Q^{in} to be an optimal coupling of R_0^{in} and R_ϵ^{in} , one finds that, denoting $\Lambda(\nabla V_1) = 1 - \lambda + \operatorname{Lip}(\nabla V_1)$,

$$\begin{aligned} MK_h(R_0(t), R_\epsilon(t))^2 &\leq \operatorname{trace}_{\mathfrak{H} \otimes \mathfrak{H}}(Q(t)^{1/2} C Q(t)^{1/2}) \leq e^{t\Lambda(\nabla V_1)} MK_h(R_0^{in}, R_\epsilon^{in})^2 \\ &\quad + \frac{e^{t\Lambda(\nabla V_1)} - 1}{\Lambda(\nabla V_1)} \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_0^{in})^{1/2} \mathcal{H}_\lambda (R_0^{in})^{1/2}) + 2\|V_1\|_{L^\infty(\mathbf{R}^d)}} \\ &\quad + \frac{e^{t\Lambda(\nabla V_1)} - 1}{\Lambda(\nabla V_1)} \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_\epsilon^{in})^{1/2} \mathcal{H}_\lambda (R_\epsilon^{in})^{1/2}) + 2\|V_2\|_{L^\infty(\mathbf{R}^d)}}, \end{aligned}$$

which is the desired inequality by coming back to V, U through (16) and using

$$\|\mu V + \epsilon U\| \leq \mu \|V\| + \epsilon \|U\|.$$

□

Proof of Theorem 2.1 and Proposition 3.2. Observe that the estimate above is uniform in \hbar — more precisely, the moduli of continuity in the initial data and in the potential are independent of \hbar . Of course, the pseudo-distance MK_h itself is not independent of \hbar . For $R, S \in \mathcal{D}(\mathfrak{H})$ let us define

$$(17) \quad \delta(W_h[R], W_h[S]) := \sup_{\substack{\max_{|\alpha|+|\beta| \leq 2[d/4]+3} \|\partial_q^\alpha \partial_p^\beta f\|_{L^\infty(\mathbf{R}^{2d})} \leq 1}} \left| \int_{\mathbf{R}^{2d}} f(q, p) (W_h[R], W_h[S])^* q, p) dq dp \right|$$

Lemma 5.3. *For any $R, S \in \mathcal{D}_2(\mathfrak{H})$,*

$$\mathbf{d}(R, S) \leq 2^d \delta(W_h[R], W_h[S]) \leq 2^d (MK_h(R, S) + C_d \hbar) \quad \text{with } C_d = (1 + \frac{\gamma_d}{\sqrt{\pi}}) 2d,$$

where $\gamma_d \leq \frac{d^{3/4} (192e^{-\frac{1}{4}} \pi^{-\frac{5}{4}})^d}{4e^{\frac{1}{4}}} (d^d)^{11/4}$ is the constant appearing in the Calderon-Vaillancourt theorem (see Appendix C in [GP20a]).

Proof. The proof consists in applying Theorem A.7 in [GJP20] and Theorem 2.3 (2) in [GMP16]. □

Lemma 5.3 shows that the inequality above implies that

$$\begin{aligned} &(2^{-d} \mathbf{d}(R_0(t), R_\epsilon(t)) - C_d \hbar)^2 \\ &(\delta(W_h[R_0(t)], W_h[R_\epsilon(t)]) - C_d \hbar)^2 \\ &\leq e^{t\Lambda(\epsilon \nabla V_1)} MK_h(R_0^{in}, R_\epsilon^{in})^2 + 2d\hbar \\ &+ \epsilon \frac{e^{t\Lambda(\epsilon \nabla V_1)} - 1}{\Lambda(\nabla V_1)} \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_0^{in})^{1/2} \mathcal{H}_\lambda (R_0^{in})^{1/2}) + 2\epsilon \|V_1\|_{L^\infty(\mathbf{R}^d)}} \\ &+ \epsilon \frac{e^{t\Lambda(\epsilon \nabla V_1)} - 1}{\Lambda(\nabla V_1)} \|\nabla(V_1 - V_2)\|_{L^\infty(\mathbf{R}^d)} \sqrt{\operatorname{trace}((R_\epsilon^{in})^{1/2} \mathcal{H}_\lambda (R_\epsilon^{in})^{1/2}) + 2\epsilon \|V_2\|_{L^\infty(\mathbf{R}^d)}} \\ &:= e^{t\Lambda(\nabla V_1)} MK_h(R_0^{in}, R_\epsilon^{in})^2 + 2d\hbar + \gamma(t)\epsilon. \end{aligned}$$

Therefore Theorem 2.1 and Proposition 3.2 are proven as soon as

$$(18) \quad MK_{\hbar}(R^{in}, R^{in})^2 \leq D' \hbar$$

since then

$$\begin{aligned} \mathbf{d}(R_0(t), R_{\epsilon}(t)) &\leq 2^d \delta(W_{\hbar}[R_0(t)], W_{\hbar}[R_{\epsilon}(t)]) \\ &\leq \sqrt{2^{2d}(e^{|\mu| \text{Lip } V} D' + 2d)\hbar + 2^{2d}\gamma(t)\epsilon + 2^d C_d \hbar} \\ &\leq \sqrt{C(t)\epsilon + D(t)\hbar} \end{aligned}$$

with

$$(19) \quad C(t) = 2^{2d}\gamma(t)$$

$$(20) \quad D(t) = e^{|\mu| \text{Lip } V} 2^{2d} 2 \max(D' + 2d, C_d^2)$$

Lemma 5.4. *For $R \in \mathcal{D}_2(\mathfrak{H})$ let*

$$\Delta(R) := \sqrt{\text{trace}\left(R^{\frac{1}{2}}\left((x - \text{trace}\left(R^{\frac{1}{2}}xR^{\frac{1}{2}}\right))^2 + (-i\nabla_x - \text{trace}\left(R^{\frac{1}{2}}(-i\nabla_x)R^{\frac{1}{2}}\right))^2\right)R^{\frac{1}{2}}\right)}.$$

Note that when R is a pure state $R|\psi\rangle\langle\psi|$, $\Delta(R)$ is, modulo a slight abuse of notation, the same as in the definition (1).

Then

$$MK_{\hbar}(R, R) \leq \sqrt{2}\Delta(R).$$

Moreover, for any $\psi \in \mathfrak{H}$,

$$MK_{\hbar}(|\psi\rangle\langle\psi|, |\psi\rangle\langle\psi|) = \sqrt{2}\Delta(|\psi\rangle\langle\psi|).$$

Proof. The proof consists in remarking that $R \otimes R$ is indeed a coupling between R and itself. Therefore

$$\begin{aligned} MK_{\hbar}(R)^2 &\leq \text{trace}\left(\left(R^{\frac{1}{2}} \otimes R^{\frac{1}{2}}\right)C\left(R^{\frac{1}{2}} \otimes R^{\frac{1}{2}}\right)\right) \\ &= \text{trace}\left(\left(R^{\frac{1}{2}} \otimes R^{\frac{1}{2}}\right)\left(x^2 \otimes I + I \otimes x^2 - 2x \otimes x\right)\left(R^{\frac{1}{2}} \otimes R^{\frac{1}{2}}\right)\right) \\ &\quad + \text{same with } x \leftrightarrow -i\hbar\nabla_x \\ &= \text{trace}\left(R^{\frac{1}{2}}x^2R^{\frac{1}{2}} + R^{\frac{1}{2}}x^2R^{\frac{1}{2}} - 2R^{\frac{1}{2}}xR^{\frac{1}{2}} \cdot R^{\frac{1}{2}}xR^{\frac{1}{2}}\right) \\ &\quad + \text{same with } x \leftrightarrow -i\hbar\nabla_x \\ &= 2 \text{trace}\left(R^{\frac{1}{2}}x^2R^{\frac{1}{2}} - \left(R^{\frac{1}{2}}xR^{\frac{1}{2}}\right)^2\right) \\ &\quad + \text{same with } x \leftrightarrow -i\hbar\nabla_x \\ &= 2\Delta(R). \end{aligned}$$

The equality is proven the same way, after Lemma 2.1 (ii) in [GP20b] which stipulates that the only coupling between $|\psi\rangle\langle\psi|$ and itself is $|\psi\rangle\langle\psi| \otimes |\psi\rangle\langle\psi|$. \square

This proves Theorem 2.1, and Theorem 4.1 when R^{in} satisfies hypothesis (i).

If the initial data R^{in} is e Töplitz operator, specifically if

$$R^{in} = \text{OP}_{\hbar}^T[(2\pi\hbar)^d \mu^{in}],$$

with the notation of [GMP16] one can go further and apply Theorem 2.3 (1) in [GMP16]:

$$MK_{\hbar}(R^{in}, R^{in}) \leq 2d\hbar,$$

so that (18) is again satisfied and Theorem 4.1 is proven when R^{in} satisfies the hypothesis (iv).

The proof in the case of the hypothesis (ii), (iii) and (v) follows directly the first inequality of Theorem 8.1 in [GP20b] with $R = S = R^{in}$, together with item (I) (through the Corollary of Theorem 3.1 in [GP20a]) and item (II) of Theorem 4.1 in [GP20a], with $\mu(\hbar) = \mu'(\hbar) = C$, $\nu(\hbar) = \nu'(\hbar) = C\sqrt{\hbar}$ and $\tau(\hbar) = \hbar$, respectively.

Indeed, [GP20b, Theorem 8.1, (iii) first inequality] stipulates that, for all density matrix R ,

$$MK_{\hbar}(R^{in}, R^{in}) \leq 2E_{\hbar}(\widetilde{W}_{\hbar}[R^{in}], R^{in})$$

where E_{\hbar} is a semiquantum pseudometric whose knowledge of the definition [GP17, Definition 2.2] is not strictly necessary for our purpose here since Theorem 4.1 in [GP20a] shows that, when $\mu(\hbar) = \mu'(\hbar) = C$, $\nu(\hbar) = \nu'(\hbar) = C\sqrt{\hbar}$ and $\tau(\hbar) = \hbar$,

$$E_{\hbar}(\widetilde{W}_{\hbar}[R^{in}], R^{in}) = O(\sqrt{\hbar}).$$

Hence

$$MK_{\hbar}(R^{in}, R^{in})^2 \leq D'\hbar$$

for some constant D' explicitly recoverable from [GP20a, Theorem 4.1].

This completes the proof of Theorems 2.1 and 4.1. \square

6. PROOF OF COROLLARIES 2.2 AND 4.2

Let us first derive the easy standard following estimate.

Proposition 6.1. *For every $t \in \mathbf{R}$,*

$$\|R_0(t) - R_{\epsilon}(t)\|_1 \leq 2t \frac{\epsilon}{\hbar} \|U\|_{\infty}$$

Proof. The solution of (15) is explicitly given by

$$R_{\epsilon}(t) = e^{-i\frac{t\mathcal{H}_{\epsilon}}{\hbar}} R^{in} e^{i\frac{t\mathcal{H}_{\epsilon}}{\hbar}}$$

Therefore, one easily shows that

$$R_0(t) - R_{\epsilon}(t) = \frac{1}{i\hbar} \int_0^t e^{-i\frac{(t-s)\mathcal{H}_{\epsilon}}{\hbar}} [\epsilon U, R_0(t)] e^{i\frac{(t-s)\mathcal{H}_{\epsilon}}{\hbar}}.$$

and the result follows by

$$\|[\epsilon U, R_0(t)]\|_1 \leq 2\epsilon \|U\| \|R_0(t)\|_1.$$

\square

Corollaries 2.2, 4.2 will follow by interpolation between Theorem 5.2 and Proposition 6.1, through the following inequality.

Lemma 6.2. *For any $R, S \in \mathcal{D}_2(\mathfrak{H})$,*

$$\mathbf{d}(R, S) \leq 2^d \|R - S\|_1$$

Proof. This is Theorem A7 in [GJP20], item (2) \square

Therefore, by Theorem 2.1 and Lemma 6.2 we get, for each $\hbar, \epsilon \in [0, 1]$,

$$\mathbf{d}(R_1(t), R_2(t)) \leq \min(\sqrt{C(t)\epsilon + D(t)\hbar}, 2|t| \frac{\epsilon}{\hbar} \|U\|_{\infty}).$$

Obviously, $\min(\sqrt{C(t)\epsilon + D(t)\hbar}, 2|t| \frac{\epsilon}{\hbar} \|U\|_{\infty}) \leq \min(\sqrt{C(t) + D(t)}, 2|t| \|U\|_{\infty}) \epsilon^{\frac{1}{3}}$ for $\epsilon, \hbar \leq 1$, since, when $\hbar \leq \epsilon^{\frac{2}{3}}$, $\hbar, \epsilon \leq \epsilon^{\frac{2}{3}}$, and, when $\hbar \geq \epsilon^{\frac{2}{3}}$, $\frac{\epsilon}{\hbar} \leq \epsilon^{\frac{1}{3}}$. The Corollary is proved.

7. PROOF OF PROPOSITION 2.4

The upper bound is given simply by Lemma 5.3 and the following inequality, proved in [GP18, Section 4]

$$MK_{\hbar}(|\psi_{z_1}\rangle\langle\psi_{z_1}|, |\psi_{z_2}\rangle\langle\psi_{z_2}|)^2 \leq |z_1 - z_2|^2 + 2d\hbar.$$

For the lower bound, we will pick up a test operator in the form of a Töplitz operator with symbol $f \geq 0$:

$$F := \text{OP}_{\hbar}^T[(2\pi\hbar)^d f],$$

One easily verifies that $\text{trace } F = \int_{\mathbf{R}^{2d}} f(q, p) dq dp$.

Moreover, see [GJP20, Appendix B],

$$\begin{aligned} \frac{1}{i\hbar}[F, -i\hbar\partial_{x_j}] &= \text{OP}_{\hbar}^T[(2\pi\hbar)^d \partial_{p_j} f], \quad j = 1, \dots, d \\ \frac{1}{i\hbar}[F, x_j] &= \text{OP}_{\hbar}^T[(2\pi\hbar)^d (-\partial_{q_j} f)], \quad j = 1, \dots, d \\ \|F\|_1 &\leq \|f\|_{L^1(\mathbf{R}^{2d})} \end{aligned}$$

Therefore, it is easy to construct functions f such that $F := \text{OP}_{\hbar}^T[(2\pi\hbar)^d f]$ satisfies the constraints of the maximization problem in the definition of \mathbf{d} .

Moreover, denoting $z = (q, p) \in \mathbf{R}^{2d}$,

$$\begin{aligned} \mathbf{d}(|\psi_{z_1}\rangle\langle\psi_{z_1}|, |\psi_{z_2}\rangle\langle\psi_{z_2}|) &\geq |\text{trace}(F(|\psi_{z_1}\rangle\langle\psi_{z_1}| - |\psi_{z_2}\rangle\langle\psi_{z_2}|))| \\ &= |\langle\psi_{z_1}|F|\psi_{z_1}\rangle - \langle\psi_{z_2}|F|\psi_{z_2}\rangle| \\ &= \left| \int_{\mathbf{R}^{2d}} f(z) (|\langle\psi_z|\psi_{z_1}\rangle|^2 - |\langle\psi_z|\psi_{z_2}\rangle|^2) \frac{dq dp}{(2\pi\hbar)^d} \right| \\ &= \left| \int_{\mathbf{R}^{2d}} f(z) (e^{-\frac{|z-z_1|^2}{2\hbar}} - e^{-\frac{|z-z_2|^2}{2\hbar}}) \frac{dq dp}{(2\pi\hbar)^d} \right| \\ &\geq |f(z_1) - f(z_2)| \\ &\quad - \frac{2}{e} \hbar \max_{\|\alpha\|, \|\beta\| \leq 2} \int_{\mathbf{R}^{2d}} |\partial_q^\alpha \partial_p^\beta f(q, p)| dq dp. \end{aligned}$$

Let us suppose now that $f \in \mathcal{S}(\mathbf{R}^{2d})$ and f is convex in a convex domain containing z_1, z_2 . Then, one can certainly rescale, translate and rotate f such that

- $f(z_1) - f(z_2) \geq \nabla f(z_2) \cdot (z_1 - z_2)$ convexity
- $|f(z_1) - f(z_2)| \geq C|z_1 - z_2|$, $C > 0$ rotation and translation
- $\max_{\|\alpha\|, \|\beta\| \leq 2[\frac{d}{4}]+3} \|\mathcal{D}_{-i\hbar\nabla}^\alpha \mathcal{D}_x^\beta F\|_1 \leq \max_{\|\alpha\|, \|\beta\| \leq 2} \|\partial_q^\alpha \partial_p^\beta f\|_{L^1(\mathbf{R}^{2d})} \leq 1$ rescaling.

Hence

$$C|z_1 - z_2| - \hbar \leq \mathbf{d}(|\psi_{z_1}\rangle\langle\psi_{z_1}|, |\psi_{z_2}\rangle\langle\psi_{z_2}|).$$

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