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CURVATURE FORMULA FOR DIRECT IMAGES OF TWISTED RELATIVE CANONICAL BUNDLES ENDOWED WITH A SINGULAR METRIC

by

Junyan Cao, Henri Guenancia & Mihai Păun

Dedicated to Ahmed Zeriahi, on the occasion of his retirement

Abstract. — In this note, we obtain various formulas for the curvature of the L^2 metric on the direct image of the relative canonical bundle twisted by a holomorphic line bundle endowed with a positively curved metric with analytic singularities, generalizing some of Berndtsson’s seminal results in the smooth case. When the twist is assumed to be relatively big, we further provide a very explicit lower bound for the curvature of the L^2 metric.

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1. Introduction

Let $p : \mathcal{X} \rightarrow D$ be a smooth, proper fibration from a $(n + 1)$ -dimensional Kähler manifold \mathcal{X} onto the unit disk $D \subset \mathbb{C}$, and let (L, h_L) be a holomorphic line bundle endowed with a possibly singular hermitian metric h_L assumed to be positively curved (i.e. when $i\Theta_{h_L}(L) \geq 0$ in the sense of currents). Then, the positivity properties of the direct image sheaves

$$\mathcal{F} := p_*((K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L))$$

endowed with the L^2 metric $h_{\mathcal{F}}$ are well-known, cf e.g. [Ber09, PT18, Pău18, DNWZ20] among many others. Moreover, when h_L is *smooth*, we have at hand explicit formulas obtained by Berndtsson [Ber09, Ber11] that compute the curvature of the L^2 metric on the direct image sheaf above.

In this article we are aiming at the generalisation of Berndtsson’s curvature formulas in case where the metric h_L has relatively simple singularities, e.g. *analytic singularities*. As for our main motivation, we recall that all the results in connection to the famous

Itaka conjecture (predicting the subadditivity of Kodaira dimension for algebraic fiber spaces) are relying on the properties of the sheaves $\mathcal{F} = \mathcal{F}_m$ in which the bundle L is of type $(m-1)K_{\mathcal{X}/D} + mB$ for some $m \geq 1$ and an effective divisor B on \mathcal{X} such that the pair (X, B) is klt. If $\det(\mathcal{F}_m)$ is either big or trivial for some $m \gg 0$, then the Itaka conjecture is already established. We expect that in the “intermediate” cases a better understanding of the geometric properties of \mathcal{F}_m will be required, in particular the structure of the flat directions of the curvature tensor and their interaction with the variation of the complex structure of the fibers of the family p . Our main result in this framework states as follows.

Theorem A. — *Let $p : \mathcal{X} \rightarrow D$ and $(L, h_L) \rightarrow \mathcal{X}$ as above, and let $u \in H^0(D, \mathcal{F})$. We assume that*

- *The metric h_L has analytic singularities and $i\Theta_{h_L}(L) \geq 0$ in the sense of currents.*
- *The section u is flat with respect to $h_{\mathcal{F}}$.*

Set $E := \{h_L = \infty\}$. Then, there exists a continuous L^2 -integrable representative \mathbf{u} of u defined on the restriction $\mathcal{X}^ \setminus E$ of the family p to some punctured disk D^* such that*

$$\frac{\bar{\partial}\mathbf{u}}{dt} \Big|_{X_t \setminus E} = 0$$

for any $t \in D^$ and*

$$(1.1) \quad D'\mathbf{u} = 0, \quad \Theta_{h_L}(L) \wedge \mathbf{u} = 0$$

on $\mathcal{X}^ \setminus E$. Here $\mathcal{X}^* := p^{-1}(D^*)$ and \mathbf{u} is L^2 with respect to h_L and a Poincaré type metric cf. Section 3.*

By “punctured disk” in the previous statement we mean that $D \setminus D^*$ is a discrete set, possibly empty. By L^2 , we mean locally L^2 with respect to the base D^* . The meaning of $\frac{\bar{\partial}\mathbf{u}}{dt}$ is provided in (3.13) below.

Finally, a representative \mathbf{u} of u (in the sense of Berndtsson) is a section of $(K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L)$ over \mathcal{X}^* such that for any $t \in D^*$, we have $\mathbf{u}|_{X_t} = u_t$ under the canonical identification $\left((K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L)\right)|_{X_t} \simeq (K_{X_t} + L|_{X_t}) \otimes \mathcal{I}(h_L|_{X_t})$, cf Lemma 2.2 and the comments below. In other words, \mathbf{u} is an L -valued $(n, 0)$ -form on \mathcal{X} which coincides with u_t on X_t .

The result we are next mentioning concerns the case of a twisting line bundle L which is p -big. It is then expected that the strict positivity of (L, h_L) is inducing stronger positivity properties of the curvature of the direct image than in the general case of a semi-positively curved L . This is confirmed by the following statement, which is a version of [Ber11, Thm 1.2].

Theorem B. — *Let $p : \mathcal{X} \rightarrow D$ be a smooth projective fibration and let $(L, h_L) \rightarrow \mathcal{X}$ be a line bundle such that*

- *h_L has analytic singularities and $i\Theta_{h_L}(L) \geq 0$ in the sense of currents.*
- *For any $t \in D$, the absolutely continuous part $\omega_L := (i\Theta_{h_L}(L))_{\text{ac}}$ satisfies $\int_{X_t} \omega_L^n > 0$.*

Then there exists a punctured disk $D^* \subset D$ such that for any $u \in H^0(D, \mathcal{F})$ we have the following inequality

$$(1.2) \quad \langle \Theta_{h_{\mathcal{F}}}(\mathcal{F})u, u \rangle_t \geq c_n \int_{X_t} c(\omega_L)u \wedge \bar{u}e^{-\phi_t}$$

for any $t \in D^*$.

In the statement above we identify $\Theta_{h_{\mathcal{F}}}(\mathcal{F})$ with an endomorphism of \mathcal{F} by "dividing" with $idt \wedge d\bar{t}$. Moreover, $c_n = (-1)^{\frac{n^2}{2}}$ is the usual unimodular constant. We denote by $c(\omega_L) := \frac{\omega_L^{n+1}}{\omega_L^n \wedge idt \wedge d\bar{t}}$ the geodesic curvature associated to ω_L , cf. Definition 2.3 for a precise definition in the degenerate case.

Actually we can provide some details about the punctured disk D^* in Theorem B. Under the hypothesis of this result, it turns out that the L^2 metric $h_{\mathcal{F}}$ is smooth in a complement of a discrete subset of D . We will show that the formula (1.2) is valid for points $t \in D$ in the neighborhood of which the metric $h_{\mathcal{F}}$ is smooth, and such that $\mathcal{F}_t = H^0(X_t, (K_{X_t} + L) \otimes \mathcal{I}(h_L|_{X_t}))$, cf Remark 5.5.

As it was kindly recalled to us by P. Naumann, many of the results in the current article overlap with his work [Nau20]. Let us emphasize that his results are stated for a metric h_L on L with a very specific type of singularities, which implies among other things that L is relatively ample. In particular, Theorem B may not be obtained as a consequence of those results. The readers interested in these topics can benefit of complementary points of view, especially in the treatment of the Hodge-theoretic part of the current article and [Nau20], respectively.

• Strategy of the proof

Roughly speaking, the idea of the proof of Theorem A and B respectively is as follows: we endow the complement $\mathcal{X} \setminus E$ with a complete metric of Poincaré type and proceed by taking advantage of what is known in the compact case, combined with the existence of families of cut-off functions specific to the complete setting. There are however quite a few difficulties along the way. Probably the most severe stems from the Hodge decomposition in the complete case: the image of the usual operators $\bar{\partial}$ and $\bar{\partial}^*$ may not be closed. We show in Section 3.2 that at least in bi-degree $(n, 1)$ this is the case, cf. Theorem 3.6, as consequence of the fact that the background metric has Poincaré singularities.

In order to construct the form u in Theorem A, we start with a representative of u given by the contraction with the lifting V of $\frac{\partial}{\partial t}$ with respect to a Poincaré metric ω_E . It turns out that this specific representative has all the desired properties needed to fit into the L^2 -theory. Then we "correct" it: this is possible by the flatness hypothesis, and it boils down to solving a fiberwise $\bar{\partial}^*$ -equation. It is both in the resolution of this equation as well as in the study of the regularity of the resulting solution that Theorem 3.6 is used. Another important ingredient of the proof is Proposition 4.1, which gives a general curvature formula for $(\mathcal{F}, h_{\mathcal{F}})$ when h_L has e.g. analytic singularities. It provides a rather wide generalisation of a result due to Berndtsson.

As for Theorem B, the starting point is the fact that the positivity properties of (L, h_L) allow us to construct a family of Poincaré metric $(\omega_\varepsilon)_{\varepsilon>0}$ on $\mathcal{X} \setminus E$. Then, the representatives \mathbf{u}_ε of u - obtained as above as the contraction with the lifting V_ε of $\frac{\partial}{\partial t}$ with respect to ω_ε - enjoy a special property that allows us to extract the desired inequality from the general curvature formula from Proposition 4.1 and a limiting argument when ε approaches zero. Although the use of this special representative goes back to Berndtson, several new analytic inputs are required to deal with the present singular situation.

• **Organization of the paper**

- In Section 2, we introduce our set-up, notation and main objects of study (the L^2 metric on \mathcal{F} , the geodesic curvature).
- In Section 3, we review two aspects of Poincaré metrics: first, the integrability properties of representatives \mathbf{u} of sections u of \mathcal{F} constructed via such metrics (Lemma 3.3) and then, we investigate the closedness of the image of the operators $\bar{\partial}, \bar{\partial}^*$ on a hermitian line bundle with analytic singularities (Theorem 3.6).
- In Section 4, we establish a general curvature formula (Proposition 4.1). This allows us to find very special representatives of flat sections of \mathcal{F} (Theorem 4.6), leading to the proof of Theorem A.
- In Section 5, we analyze the relatively big case in the "snc situation" (Theorem 5.1), from which we then deduce Theorem B.

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It is our privilege to dedicate this article to our friend and colleague Ahmed Zeriahi, with our admiration for his outstanding mathematical achievements and wishing him a very happy and active retirement!

2. Set-up and notation

The set of assumptions we need for our results to hold is the following.

Set-up 2.1. — Let $p : \mathcal{X} \rightarrow D$ be a smooth, proper fibration from a $(n + 1)$ -dimensional Kähler manifold \mathcal{X} onto the unit disk $D \subset \mathbb{C}$, and let (L, h_L) be a holomorphic line bundle endowed with a possibly singular hermitian metric h_L .

We assume that there exists a divisor $E = E_1 + \dots + E_N$ whose support is contained in the total space \mathcal{X} of p . such that the following requirements are fulfilled.

- (A.1) For every $t \in D$ the divisor $E + X_t$ has simple normal crossings. Let $\Omega \subset \mathcal{X}$ be a coordinate subset on \mathcal{X} . We take $(z_1, \dots, z_n, t = z_{n+1})$ a coordinate system

on Ω such that the last one z_{n+1} corresponds to the map p itself and such that $z_1 \dots z_p = 0$ (for some index $p \leq N$) is the local equation of $E \cap \Omega$.

(A.2) The metric h_L has generalised analytic singularities along E ; i.e. its local weights φ_L on Ω can be written as

$$\varphi_L \equiv \sum_{i=1}^p a_i \log |z_i|^2 - \sum_I b_I \log \left(\phi_I(z) - \log \left(\prod_{i \in I} |z_i|^{2k_i} \right) \right)$$

modulo C^∞ functions, where a_i, b_I are positive real numbers, k_i are positive integers and $(\phi_I)_I$ are smooth functions on Ω . The set of indexes in the second sum coincides with the non-empty subsets of $\{1, \dots, p\}$.

(A.3) The Chern curvature of (L, h_L) satisfies

$$i\Theta_{h_L}(L) \geq 0$$

in the sense of currents on \mathcal{X} .

We then set

$$\mathcal{F} := p_*((K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L))$$

and assume that that this vector bundle on D has positive rank. As a consequence of the previous requirements (A.1)-(A.3), we have the following statement.

Lemma 2.2. — *Under the assumptions (A.1)-(A.3), we have*

$$\mathcal{F}_t = H^0(X_t, (K_{X_t} + L) \otimes \mathcal{I}(h_L|_{X_t}))$$

for every $t \in D$. Moreover, the canonical L^2 metric (cf. Notation 2.4) on \mathcal{F} is non-singular.

Proof. — We first remark that \mathcal{F} is indeed locally free given that it is torsion-free and $D \subset \mathbb{C}$ is a disk.

The fibers of \mathcal{F} are indeed identified with $H^0(X_t, (K_{X_t} + L) \otimes \mathcal{I}(h_L|_{X_t}))$ because of the transversality hypothesis (A.1), combined with the type of singularities we are allowing for h_L in (A.2). The point is that a holomorphic function f defined on the coordinate subset Ω belongs to $\mathcal{I}(h_L)$ exactly when the restriction $f|_{\Omega \cap X_t}$ belongs to the ideal $\mathcal{I}(h_L|_{X_t})$. On the other hand the Kähler version of Ohsawa-Takegoshi theorem [Cao17] implies that any element of $H^0(X_t, (K_{X_t} + L) \otimes \mathcal{I}(h_L|_{X_t}))$ extends to \mathcal{X} –it is at this point that the hypothesis (A.3) plays a crucial role.

Concerning the smoothness of the L^2 -metric on \mathcal{F} , we can use partitions of unity to reduce to checking that integrals of the form $\int_{\Omega \cap X_t} |f_t|^2 e^{-\varphi_L}$ vary smoothly with t , where $f_t = f|_{X_t}$ for some $f \in \mathcal{I}(h_L)|_\Omega$ and φ_L is given by the expression in (A.2). Now it is clear there that all derivatives in the t, \bar{t} variables of φ_L are bounded, so that the result follows from general smoothness results for integrals depending on a parameter. \square

• *A few comments about the conditions (A.1)-(A.2).*

The point we want to make here is that the transversality requirements in (A.1)-(A.2) can be obtained starting from a quite general context.

We consider $p : \mathcal{X} \rightarrow D$ a proper fibration from a $(n+1)$ -dimensional Kähler manifold \mathcal{X} onto the unit disk $D \subset \mathbb{C}$, and let (L, h_L) be a holomorphic line bundle endowed

with a possibly singular hermitian metric h_L . We assume that (A.3) holds true, and that the singularities of h_L are of the form

$$(2.1) \quad \varphi_L \equiv \sum_{i=1}^p a_i \log |f_i|^2 - \sum_{i=1}^p b_i \log (\tau_i - \log |g_i|^2)$$

modulo C^∞ functions, where a_i, b_i are positive real numbers, f_i, g_i are holomorphic, and τ_i are smooth.

If (A.1)-(A.2) are not satisfied for $p : \mathcal{X} \rightarrow D$, then one can consider a log resolution $\pi : \mathcal{X}' \rightarrow \mathcal{X}$ of $(\mathcal{X}, \mathcal{I}_Z)$ where Z is the singular set of h_L . Set $p' := p \circ \pi : \mathcal{X}' \rightarrow D$, $(L', h_{L'}) := (\pi^* L, \pi^* h_L)$, E is the reduced divisor induced by $\pi^{-1}(Z)$. It is immediate that

$$\pi_*((K_{\mathcal{X}'/D} + L') \otimes \mathcal{I}(h_{L'})) = (K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L)$$

so that, in particular, $\pi_*((K_{\mathcal{X}'/D} + L') \otimes \mathcal{I}(h_{L'})) = \mathcal{F}$.

The map p' may not be smooth anymore (e.g. some components of E may be irreducible components of fibers of p'). Set D_{reg} to be the Zariski open set of regular values of p' , $D_1 := D_{\text{reg}} \cap D$, $\mathcal{X}_1 := p'^{-1}(D_1)$, $p_1 := p'|_{\mathcal{X}_1}$, $(L_1, h_{L_1}) := (L', h_{L'})|_{\mathcal{X}_1}$. Then, the triplet (p_1, L_1, h_{L_1}) satisfies the assumptions (A.1)-(A.3).

In conclusion, starting with a map p as above and singular metric h_L as in (2.1), we can use our results on the family p_1 restricted to some punctured disk $D^* \subset D$.

• *The geodesic curvature in a degenerate setting.*

Let $p : \mathcal{X} \rightarrow D$ be a smooth, proper fibration where \mathcal{X} is a Kähler manifold of dimension $n + 1$. Let ω be a closed positive $(1, 1)$ -current on \mathcal{X} such that ω is smooth on a non-empty Zariski open subset $\mathcal{X}^\circ \subset \mathcal{X}$. Let $\omega_{\mathcal{X}}$ be a Kähler metric on \mathcal{X} .

Definition 2.3. — *The geodesic curvature $c(\omega)$ of ω on \mathcal{X}° is defined by*

$$c(\omega) := \lim_{\varepsilon \rightarrow 0} c(\omega + \varepsilon \omega_{\mathcal{X}}) = \lim_{\varepsilon \rightarrow 0} \frac{(\omega + \varepsilon \omega_{\mathcal{X}})^{n+1}}{(\omega + \varepsilon \omega_{\mathcal{X}})^n \wedge \text{idt} \wedge \bar{d}\bar{f}}.$$

A few explanations are in order.

First, if ω is relatively Kähler on \mathcal{X} , one recovers the usual definition.

Next, it is easy to observe that $c(\omega + \varepsilon \omega_{\mathcal{X}}) = 1/\|dt\|_{\omega + \varepsilon \omega_{\mathcal{X}}}^2$. In particular, that non-negative quantity is non-increasing when ε decreases towards 0, hence it admits a limit. By the same token, one can see that the limit is independent of the choice of the Kähler metric $\omega_{\mathcal{X}}$.

Finally, if $t \in D$ is such that $X_t^\circ := X_t \cap \mathcal{X}^\circ$ is dense in X_t , then one defines $c(\omega)$ on the whole X_t by extending it by zero across $X_t \setminus X_t^\circ$. Note that if the absolutely continuous part of ω satisfies $\omega_{\text{ac}} \leq C \omega_{\mathcal{X}}$ on X_t for some constant $C > 0$, then $c(\omega)$ is a bounded function on X_t (this follows e.g. from the inequality $c(\omega) \leq \|\frac{\partial}{\partial \bar{t}}\|_{\omega}^2$ for a set of coordinates $(z_1, \dots, z_n, z_{n+1} = t)$ such that $p(z) = t$). In particular, the integral $\int_{X_t} c(\omega) \omega_{\mathcal{X}}^n$ is finite.

Notation 2.4. — *In the Set-up 2.1 above:*

• We set $\mathcal{X}^\circ := \mathcal{X} \setminus E$, $X_t^\circ := X_t \cap \mathcal{X}^\circ$, $L_t := L|_{X_t}$, $h_{L_t} := h_L|_{X_t}$.

· We use interchangeably h_L and $e^{-\phi_L}$; when working in a trivializing chart of L , we will denote by φ_L the local weight of h_L . The $(1,0)$ -part of the Chern connection of (L, h_L) over $\mathcal{X} \setminus E$ is denoted by D' .

· Under assumption (A.1), we will write $E := \sum_{i=1}^N E_i$ for the decomposition of E into its (smooth) irreducible components. Next, let s_i be a section of $\mathcal{O}_{\mathcal{X}}(E_i)$ that cuts out E_i , and let h_i be a smooth hermitian metric on $\mathcal{O}_{\mathcal{X}}(E_i)$. In the following, $|s_i|^2$ stands for $|s_i|_{h_i}^2$, and we assume that $|s_i|^2 < e^{-1}$.

· We will interchangeably denote by $\|\cdot\|$ or $h_{\mathcal{F}}$ the L^2 metric on \mathcal{F} ; i.e. if $u \in \mathcal{F}_t = H^0(X_t, (K_{X_t} + L_t) \otimes \mathcal{I}(h_{L_t}))$, then $\|u\|^2 := c_n \int_{X_t} u \wedge \bar{u} e^{-\phi_{L_t}}$ with $c_n = (-1)^{\frac{n^2}{2}}$. Lemma 2.2 ensures that the L^2 metric is *smooth* on D . We denote by ∇ the $(1,0)$ part of the Chern connection of $(\mathcal{F}, \|\cdot\|)$ on D .

3. A few technicalities about Poincaré type metrics

Throughout this section we adopt Set-up 2.1 and use the Notation 2.4. Let ω be a fixed Kähler metric on \mathcal{X} , and let

$$\omega_E := \omega + dd^c \left[- \sum_{i=1}^N \log \log \frac{1}{|s_i|^2} \right] \quad \text{on } \mathcal{X}^\circ$$

be a metric with Poincaré singularities along E . Thanks to (A.1) we infer that $\omega_E|_{X_t^\circ}$ is a complete Kähler metric on X_t° with Poincaré singularities along $E \cap X_t$ for each $t \in D$.

In the next subsections we will be concerned with the following two main themes.

Let V be the horizontal lift of $\frac{\partial}{\partial t}$ with respect to the Poincaré-type metric ω_E . We estimate the size of its coefficients near the singularity divisor E , and then show that the representatives of direct images constructed by using V have the expected L^2 properties allowing us to use them in the computation of the curvature of \mathcal{F} . This is the content of Subsection 3.1.

In Subsection 3.2 we establish a few important properties of the L^2 -Hodge decomposition for $(n,1)$ -forms with values in (L, h_L) , where the background metric is (X, ω_E) . The main result here is that the image of $\bar{\partial}^*$ is closed, cf Theorem 3.6, a very useful result *per se* and for the next sections of this paper as well.

3.1. Estimates of the background metric, lifting. — We choose local coordinates $(z_1, \dots, z_n, z_{n+1} = t)$ on \mathcal{X} such that $p(z, t) = t$ —as in (A.1)—and in which ω_E is locally given by

$$\omega_E = g_{t\bar{t}} idt \wedge d\bar{t} + \sum_{\alpha=1}^n (g_{\alpha\bar{t}} idz_\alpha \wedge d\bar{t} + g_{t\bar{\alpha}} idt \wedge d\bar{z}_\alpha) + \sum_{\alpha, \beta=1}^n g_{\alpha\bar{\beta}} idz_\alpha \wedge d\bar{z}_\beta$$

By the estimates in [Gue14, §4.2] the coefficients of ω_E are can be written as follows

$$(3.1) \quad g_{\alpha\bar{\beta}} = g_{\alpha\bar{\beta}}^0 + \frac{\delta_\alpha \cdot \delta_{\alpha\beta}}{|z_\alpha|^2 \log^2 |z_\alpha|^2} + \frac{\delta_\alpha A_\alpha}{z_\alpha \log^2 |z_\alpha|^2} + \frac{\delta_\beta B_\beta}{\bar{z}_\beta \log^2 |z_\beta|^2} + \sum_{\gamma=1}^p \frac{C_\gamma}{\log |z_\gamma|^2}$$

where $A_\alpha, B_\beta, C_\gamma, g_{\alpha\bar{\beta}}^0$ are smooth functions on Ω . We use the notation

$$\delta_\alpha = \begin{cases} 1 & \text{if } \alpha \in \{1, \dots, p\} \\ 0 & \text{else.} \end{cases}$$

and $\delta_{\alpha\beta}$ is the usual Kronecker symbol.

In order to present the computations to follow in a reasonably simple way, we introduce for $\alpha = 1, \dots, n$ the functions

$$(3.2) \quad f_\alpha(z) := \begin{cases} w_\alpha := z_\alpha \log |z_\alpha|^2 & \text{if } \alpha \in \{1, \dots, p\} \\ 1 & \text{else} \end{cases}$$

or $f_\alpha(z) = \delta_\alpha w_\alpha + 1 - \delta_\alpha$ more concisely. The coefficients of the metric ω_E can be written as

$$(3.3) \quad g_{\alpha\bar{\beta}} = \frac{1}{f_\alpha \bar{f}_\beta} \Psi_{\alpha\bar{\beta}}(z'', w, \rho).$$

The notations we are using in (3.3) are:

- (a) $\alpha, \beta \in \{1, \dots, n\}$.
- (b) $\Psi_{\alpha\bar{\beta}}$ is a smooth function defined in the neighborhood of $0 \in \mathbb{C}^{n-p} \times \mathbb{C}^p \times \mathbb{C}^p$.
- (c) $w := (w_1, \dots, w_p)$ (cf. (3.2)) and $z'' := (z_{p+1}, \dots, z_n)$.
- (d) For $\alpha = 1, \dots, p$ we introduce $\rho_\alpha := \frac{1}{\log |z_\alpha|^2}$.

To justify the smoothness of $\psi_{\alpha\bar{\beta}}$, we keep in mind that $f_\alpha = \delta_\alpha w_\alpha + 1 - \delta_\alpha$ is a smooth function of w and that if $\alpha \in \{1, \dots, p\}$, $z_\alpha = w_\alpha \rho_\alpha$ is a smooth function of (w, ρ) . Smoothness of Ψ now follows once one rewrites (3.1) as

$$f_\alpha \bar{f}_\beta g_{\alpha\bar{\beta}} = f_\alpha \bar{f}_\beta g_{\alpha\bar{\beta}}^0 + \delta_\alpha \delta_{\alpha\beta} + \delta_\alpha A_\alpha \rho_\alpha \bar{f}_\beta + \delta_\beta B_\beta \rho_\beta f_\alpha + \sum_{\gamma=1}^p C_\gamma \rho_\gamma f_\alpha \bar{f}_\beta.$$

Given that ω_E is a metric, the functions $\Psi_{\alpha\bar{\beta}}$ are not arbitrary (since the matrix $(\Psi_{\alpha\bar{\beta}})_{\alpha\bar{\beta}}$ is definite positive at each point). We simply want to emphasize in (3.3) the general shape of the coefficients, which will be useful in the statement that follows.

Lemma 3.1. — *The following estimates hold true:*

- (i) $\det(g) = \prod_{\alpha=1}^n |f_\alpha(z)|^{-2} (c + \Psi(z'', w, \rho))$, where the c inside the parentheses is a strictly positive constant, and Ψ is smooth such that $\Psi(0) = 0$.
- (ii) For each pair of indexes α, β we have

$$g^{\bar{\beta}\alpha} = f_\alpha(z) \bar{f}_\beta(z) \Psi^{\alpha\bar{\beta}}(z'', w, \rho),$$

where $g^{\bar{\beta}\alpha}$ are the coefficients of the inverse of $(g_{\alpha\bar{\beta}})$.

Proof. — Both statements above are obtained by a direct calculation, using the expression (3.3) of the coefficients. We skip the straightforward details. \square

Thanks to Lemma 3.1 it is easy to infer the following useful estimates: for each set of indexes α, β, q, r we have

$$(3.4) \quad \begin{aligned} \frac{\partial g^{\bar{\beta}\alpha}}{\partial z_q}(z) &= \left(\delta_{\alpha q} \delta_\alpha (1 + \log |z_\alpha|^2) \bar{f}_\beta(z) + \delta_{\beta q} \delta_\beta f_\alpha(z) \frac{\bar{z}_q}{z_q} \right) \Psi_{\alpha\bar{\beta}}(z'', w, \rho) \\ &\quad + (1 + \delta_q (\log |z_q|^2 - 1)) f_\alpha(z) \bar{f}_\beta(z) \Psi_{\alpha\bar{\beta}}(z'', w, \rho) + \frac{\delta_q}{z_q \log^2 |z_q|^2} f_\alpha(z) \bar{f}_\beta(z) \Psi_{\alpha\bar{\beta}}(z'', w, \rho) \end{aligned}$$

as well as

$$(3.5) \quad \begin{aligned} \frac{\partial^2 g^{\bar{\beta}\alpha}}{\partial z_q \partial \bar{z}_r}(z) &= \mathcal{O}(1) \frac{\partial}{\partial \bar{z}_r} \left(\delta_{\alpha q} \delta_\alpha (1 + \log |z_\alpha|^2) \bar{f}_\beta(z) + \delta_{\beta q} \delta_\beta f_\alpha(z) \frac{\bar{z}_q}{z_q} \right) \\ &\quad + \mathcal{O}(1) \left(1 + \frac{\delta_r}{\bar{z}_r \log |z_r|^2} \right) \left(\delta_{\alpha q} \delta_\alpha (1 + \log |z_\alpha|^2) \bar{f}_\beta(z) + \delta_{\beta q} \delta_\beta f_\alpha(z) \frac{\bar{z}_q}{z_q} \right) \\ &\quad + \mathcal{O}(1) \frac{\delta_q \delta_{rq}}{|z_q|^2 \log^3 |z_q|^2} f_\alpha(z) \bar{f}_\beta(z) \\ &\quad + \mathcal{O}(1) \left(1 + \frac{\delta_q}{z_q \log |z_q|^2} \right) \left(1 + \frac{\delta_r}{\bar{z}_r \log |z_r|^2} \right) f_\alpha(z) \bar{f}_\beta(z) \\ &\quad + \mathcal{O}(1) \left(1 + \frac{\delta_q}{z_q \log |z_q|^2} \right) \left(\delta_{\beta r} \delta_\beta (1 + \log |z_\beta|^2) f_\alpha(z) + \delta_{\alpha r} \delta_\alpha \bar{f}_\beta(z) \frac{z_r}{\bar{z}_r} \right) \end{aligned}$$

Again, in the relations (3.4)–(3.5) we are using $\Psi_{\alpha\bar{\beta}}$ and $\mathcal{O}(1)$ as generic notation, these functions are allowed to change from one line to another, the point is that they are of the same type. The verification of formula (3.5) is immediate, one simply takes the derivative in (3.4).

3.1.1. Horizontal lift, μ and η . — One can define the lift V of $\frac{\partial}{\partial t}$ with respect to the Poincaré type metric ω_E , cf. [Siu86, Ber11, Sch12]. It is a vector field of type $(1, 0)$ on \mathcal{X}° such that dp maps it to $\frac{\partial}{\partial t}$ pointwise on \mathcal{X}° and which is ω_E -orthogonal to $T^{1,0}X_t$ for any t . In local coordinates, one has the following formula

$$(3.6) \quad V = \frac{\partial}{\partial t} - \sum_{\alpha, \beta} g^{\bar{\beta}\alpha} g_{t\bar{\beta}} \frac{\partial}{\partial z_\alpha}.$$

Let u be a holomorphic section of $p_*(\mathcal{O}(K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L))$. One can choose an arbitrary representative U_0 of u , this is an L -valued $(n, 0)$ -form on \mathcal{X} which coincides with u_t on X_t . Now, let

$$(3.7) \quad \mathbf{u} := V \lrcorner (dt \wedge U_0).$$

One can write locally (using the previous system of coordinates):

$$U_0 \wedge dt = a(z, t) dt \wedge dz_1 \wedge \dots \wedge dz_n$$

where $a(z, t)$ is holomorphic with values in L . We have an explicit formula:

$$(3.8) \quad \mathbf{u} = a(z, t) \left(dz_1 \wedge \dots \wedge dz_n - \sum_{\alpha, \beta} (-1)^\alpha g^{\bar{\beta}\alpha} g_{t\bar{\beta}} dt \wedge dz_1 \wedge \dots \wedge \widehat{dz_\alpha} \wedge \dots \wedge dz_n \right).$$

By construction, we have

$$(3.9) \quad dt \wedge \mathbf{u} = dt \wedge U_0.$$

Therefore, although \mathbf{u} is only well defined on \mathcal{X}° , $\mathbf{u} \wedge dt$ can be extended as a smooth form on \mathcal{X} .

Remark 3.2. — The representative \mathbf{u} in (3.8) has the following interesting property

$$(3.10) \quad \mathbf{u} \wedge \omega_E|_\Omega = a(z, t) g_{t\bar{t}} dt \wedge \bar{d}\bar{t} \wedge dz_1 \wedge \dots \wedge dz_n.$$

In particular, we have $\frac{\mathbf{u} \wedge \omega_E}{dt} \Big|_{X_t} = 0$ for every t .

Moreover, as U_0 is a smooth representative of u on \mathcal{X} and u is a holomorphic form, we have $\bar{\partial}(U_0 \wedge dt) = 0$ on \mathcal{X} . Combining with (3.9), we know that $\bar{\partial}\mathbf{u} \wedge dt = 0$ on \mathcal{X}° . As a consequence, we can find a smooth $(n-1, 1)$ -form η on \mathcal{X}° such that

$$(3.11) \quad \bar{\partial}\mathbf{u} = dt \wedge \eta \quad \text{on } \mathcal{X}^\circ$$

and one has

$$(3.12) \quad \begin{aligned} D'_{\text{loc}} \mathbf{u} &= \partial\mathbf{u} - \partial\varphi_L \wedge \mathbf{u} \\ &= dt \wedge \mu \end{aligned}$$

for some $(n, 0)$ -form μ on \mathcal{X}° . The forms

$$(3.13) \quad \frac{\bar{\partial}\mathbf{u}}{dt} \Big|_{X_t^\circ} := \eta|_{X_t^\circ}, \quad \text{and} \quad \frac{D'\mathbf{u}}{dt} \Big|_{X_t^\circ} := \mu|_{X_t^\circ}$$

are well-defined and do not depend on the choice of η satisfying (3.11) (resp. μ satisfying (3.12)).

3.1.2. Estimates for η and μ . — In this section our goal is to establish the following statement.

Lemma 3.3. — *We consider the forms $\eta|_{X_t}, \mu|_{X_t}$ induced by the representative \mathbf{u} constructed in the previous section in (3.13). Then $\eta|_{X_t}, \mu|_{X_t}$ as well as $\bar{\partial}\mu|_{X_t}$ are L^2 with respect to $(\omega_E, e^{-\phi_L})$.*

Moreover, \mathbf{u}, η, μ and $\bar{\partial}\mu$ are also in $L^2(\mathcal{X}^\circ)$ with respect to $(\omega_E, e^{-\phi_L})$, say up to shrinking D .

Proof. — This is routine: we can easily obtain the explicit expression of η and μ , and then we simply evaluate their respective L^2 norms by using the estimates (3.4)–(3.5). We detail to some extent the calculations next.

First of all, we have

$$\eta|_{X_t} = a(z, t) \sum_{\alpha, \beta, r} (-1)^{\alpha+1} (g_{\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta}} + g^{\bar{\beta}\alpha} g_{t\bar{\beta}, \bar{r}}) dz_1 \wedge \dots \wedge \widehat{dz_\alpha} \wedge \dots \wedge dz_n \wedge d\bar{z}_r.$$

where we use the notation $g_{r\bar{r}}^{\bar{\beta}\alpha} := \frac{\partial g^{\bar{\beta}\alpha}}{\partial \bar{z}_r}$.

We consider first the quantity

$$(3.14) \quad \left| g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right|^2.$$

Thanks to the equality (3.4), up to a constant it is smaller than

$$(3.15) \quad \frac{1}{|f_\beta|^2} \left(\delta_{\alpha r} \delta_\alpha \log^2 |z_\alpha|^2 |f_\beta|^2 + \delta_{\beta r} \delta_\beta |f_\alpha|^2 + \left(1 - \delta_r + \frac{\delta_r}{|z_r|^2 \log^2 |z_r|^2} \right) |f_\alpha f_\beta|^2 \right)$$

which simplifies to

$$|f_\alpha|^2 + \delta_r \left[\delta_{\alpha r} \log^2 |z_r|^2 + \delta_{\beta r} \frac{|f_\alpha|^2}{|f_r|^2} + \left(\frac{1}{|z_r|^2 \log^2 |z_r|^2} - 1 \right) |f_\alpha|^2 \right].$$

Since

$$\left| dz_1 \wedge \dots \wedge \widehat{dz_\alpha} \wedge \dots \wedge dz_n \wedge d\bar{z}_r \right|_{\omega_E}^2 dV_{\omega_E} \lesssim \frac{|f_r|^2}{|f_\alpha|^2} dV_\omega$$

we eventually find

$$(3.16) \quad \left| g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right|^2 \left| dz_1 \wedge \dots \wedge \widehat{dz_\alpha} \wedge \dots \wedge dz_n \wedge d\bar{z}_r \right|_{\omega_E}^2 dV_{\omega_E} \leq C(1 + \delta_r \delta_{\alpha r} \log^2 |z_r|^2) dV_\omega.$$

The term

$$(3.17) \quad \left| g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\bar{r}} \right|^2$$

is bounded by

$$(3.18) \quad \delta_{\beta r} \delta_\beta (1 + \log^2 |z_\beta|^2) \frac{|f_\alpha|^2}{|f_\beta|^2} + \left(1 - \delta_r + \frac{\delta_r}{|w_r|^2} \right) |f_\alpha|^2$$

and we see that the same thing as before occurs, i.e.

$$(3.19) \quad \left| g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\bar{r}} \right|^2 \left| dz_1 \wedge \dots \wedge \widehat{dz_\alpha} \wedge \dots \wedge dz_n \wedge d\bar{z}_r \right|_{\omega_E}^2 dV_{\omega_E} \leq C(1 + \delta_r \log^2 |z_r|^2) dV_\omega.$$

Thus, the restriction of η to any fiber of the family $p : \mathcal{X} \rightarrow D$ is L^2 with respect to $(\omega_E, e^{-\varphi_L})$ since the holomorphic function $a(z, t)$ belongs to the multiplier ideal sheaf defined by h_L . Indeed, setting $\nu := |a|^2 dV_\omega$, one has $e^{-\varphi_L} \in L^{1+\varepsilon}(\nu)$ for some $\varepsilon > 0$ (this is easily checked since φ_L has analytic singularities and $e^{-\varphi_L} \in L^1(\nu)$) while $\log |z_r|^2 \in L^p(\nu)$ for all $p > 0$, so that Hölder inequality shows the claim.

The local expression of the form $\mu|_{X_t}$ is obtained by restricting $\frac{D'\mathbf{u}}{dt}$ to the fiber X_t ; it reads as

$$(3.20) \quad \begin{aligned} \frac{\mu|_{X_t}}{dz} &= a_{,t} - a \sum_{\alpha,\beta} (g_{r,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} + g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha}) - \sum_{\alpha,\beta} a_{,\alpha} g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \\ &\quad - a(z, t) \varphi_{L,t} + a(z, t) \sum_{\alpha,\beta} \varphi_{L,\alpha} g_{r\bar{r}}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \end{aligned}$$

where $dz := dz_1 \wedge \cdots \wedge dz_n$.

By our transversality conditions, the function a_t is still L^2 with respect to h_L . The term $g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} + g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha}$ is treated as we did for (3.14) and (3.17), with the exception that the indexes r and β coincide (and the type of the form is different). We have up to some constant

$$(3.21) \quad \left| g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right| \leq \left((1 - \log |z_\alpha|^2) |f_\beta| + \delta_{\alpha\beta} \delta_\alpha |f_\alpha| + \left[(1 - \delta_\alpha) + \frac{\delta_\alpha}{|z_\alpha| \log^2 |z_\alpha|^2} \right] \cdot |f_\alpha f_\beta| \right) \cdot \frac{1}{|f_\beta|} \\ \lesssim 1 - \log |z_\alpha|^2$$

and we can bound this term as before. The second term satisfies $g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha} = g_{,\alpha}^{\bar{\beta}\alpha} g_{\alpha\bar{\beta},t}$ since ω_E is Kähler, hence it is bounded.

Next, we have $\left| g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right|^2 \leq C |z_\alpha|^2 \log^2 |z_\alpha|$, and also that

$$(3.22) \quad |a_{,\alpha}|^2 |z_\alpha|^2 \log^2 |z_\alpha|^2 e^{-\varphi_L} = \mathcal{O}(|a|^2 \log^2 |z_\alpha|^2 e^{-\varphi_L}) \in L^1$$

for any $\alpha = 1, \dots, p$ again by the transversality/ L^2 conditions we impose to $a(z, t)$, so the third term in (3.20) is in L^2 . A similar argument applies to the second line of (3.20).

The last part of the proof of our lemma concerns $\bar{\partial}\mu$; the computations are using (3.5). We will only discuss the term

$$(3.23) \quad \frac{\partial}{\partial \bar{z}_r} \left(g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} + g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha} \right) dz \wedge d\bar{z}_r$$

since for (3.23) the computations are the most involved. The reason why we are able to conserve the L^2 property is that the partial derivative with respect to \bar{z}_r will induce a new term of order $\mathcal{O}(1/|z_r|)$ if $r \leq p$, and its square will be compensated by $|d\bar{z}_r|_{\omega_E}^2$. As for the computations: the singularities induced by $g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}}$ are bounded by the following quantity

$$(3.24) \quad \left| g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right| \leq \delta_\alpha \delta_{r\alpha} \frac{1}{|z_r|} + \frac{\delta_\alpha \delta_{\beta r}}{|f_\beta|} \log \frac{1}{|z_\alpha|^2} \log \frac{1}{|z_\beta|^2} + \delta_\alpha \delta_{r\alpha\beta} \frac{1}{|f_\beta| |z_\alpha|^2} \\ + \frac{\delta_r}{|w_r|} \left(\delta_\alpha \log \frac{1}{|z_\alpha|^2} + \delta_\alpha \delta_{\alpha\beta} \right) \\ + \frac{\delta_\alpha \delta_r}{|w_\alpha w_r|} |f_\alpha| + \frac{\delta_\alpha}{|w_\alpha|} \left(\delta_\beta \delta_{\beta r} \log \frac{1}{|z_\beta|^2} \frac{|f_\alpha|}{|f_\beta|} + \delta_\alpha \delta_{\alpha r} \right)$$

from which we see that the first part of (3.23) is L^2 . The remaining terms are

$$(3.25) \quad \left(g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\bar{r}} + g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha} + g_{,\alpha}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha\bar{r}} \right) dz \wedge d\bar{z}_r$$

for which one could use the fact that the metric ω_E is Kähler and so we have

$$(3.26) \quad g_{t\bar{\beta},\alpha} = g_{\alpha\bar{\beta},t}, \quad g_{t\bar{\beta},\alpha\bar{r}} = g_{\alpha\bar{\beta},t\bar{r}}.$$

The equalities (3.26) are simplifying a bit the calculations, since the derivative with respect to t does not increase at all the order of the singularity.

For the first term of (3.25), we have, up to a multiplicative constant

$$\begin{aligned}
 \left| g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\bar{r}} \right| &\leq \left| \delta_\alpha (1 - \log |z_\alpha|^2) \bar{f}_\beta + \delta_\alpha \delta_{\alpha\beta} f_\alpha + \left((1 - \delta_\alpha) + \frac{\delta_\alpha}{z_\alpha \log^2 |z_\alpha|^2} \right) f_\alpha \bar{f}_\beta \right| \times \\
 &\quad \left| \delta_r \delta_{r\beta} \frac{1}{|z_r|^2 \log^2 |z_r|^2} + \frac{1}{\bar{f}_\beta} \left(1 + \frac{\delta_r}{\bar{z}_r \log^2 |z_r|^2} \right) \right| \\
 &\leq \left| (1 - \delta_\alpha \log |z_\alpha|) \cdot \left(1 + \frac{\delta_r}{|z_r| (-\log |z_r|^2)} \right) \right|
 \end{aligned}$$

In particular, we get

$$\left| g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\bar{r}} \right| \cdot |d\bar{z}_r|_{\omega_E} \lesssim 1 - \delta_\alpha \log |z_\alpha|$$

and we are done with this term as before.

For the second term of (3.25), we have, using (3.26)

$$\begin{aligned}
 \left| g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{\alpha\bar{\beta},t} \right| &\leq \frac{1}{|f_\alpha f_\beta|} \cdot \left(\delta_{\alpha r} \delta_\alpha (1 - \log |z_\alpha|^2) |f_\beta| + \delta_{\beta r} \delta_\beta |f_\alpha| \right. \\
 &\quad \left. + \left[(1 - \delta_r) + \delta_r (-\log |z_r|^2) + \frac{1}{|z_r| \log^2 |z_r|^2} \right] |f_\alpha f_\beta| \right) \\
 &\leq 1 + \delta_r \frac{-\log |z_r|^2}{|z_r| (-\log |z_r|^2)}.
 \end{aligned}$$

In particular, we get

$$\left| g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{\alpha\bar{\beta},t} \right| \cdot |d\bar{z}_r|_{\omega_E} \lesssim 1 - \delta_r \log |z_r|$$

and we are done.

As for the last term of (3.25), we use (3.26) and (3.5) to see that the expansion of $g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha\bar{r}}$ will only involve terms like

$$\psi_{\alpha\bar{\beta},\bar{r}'} \frac{\partial_{\bar{r}} f_\alpha}{f_\alpha}, \frac{\partial_{\bar{r}} \bar{f}_\beta}{\bar{f}_\beta}$$

which are respectively of order

$$\frac{\delta_r}{z_r \log |z_r|^2}, \frac{\delta_r \delta_{\alpha r}}{z_r \log |z_r|^2}, \frac{\delta_r \delta_{\beta r}}{\bar{z}_r}.$$

All in all, we find

$$\left| g_{\alpha\bar{\alpha}}^{\bar{\beta}\alpha} g_{t\bar{\beta},\alpha\bar{r}} \right| \cdot |d\bar{z}_r|_{\omega_E} \lesssim 1 - \delta_r \log |z_r|$$

and this is the end of the main part of the proof.

The integrability of \mathbf{u}, η, μ on \mathcal{X}° follows directly from the estimates we have obtained above. Concerning $\bar{\partial}\mu$ there is one additional term given by $\frac{\partial}{\partial \bar{t}}$ of the expression in (3.20). This is however harmless: given the shape of the coefficients $(g_{\alpha\bar{\beta}})$ (i.e. the

transversality conditions), the additional anti-holomorphic derivative with respect to t induces no further singularity and the estimates e.g. for the term

$$\frac{\partial}{\partial \bar{t}} \left(g^{\bar{\beta}\alpha} g_{t\bar{\beta}} \right)$$

will be completely identical to those already obtained $g^{\bar{\beta}\alpha} g_{t\bar{\beta}}$. We leave the details to the interested reader. \square

Remark 3.4. — Using quasi-coordinates adapted to the Poincaré metric ω_E (cf. e.g. [CY75, Kob84, TY87]), we can prove easily that η and its derivatives are in L^2 . However, that argument cannot be applied to μ because of the singularity in the Chern connection of (L, h_L) .

3.2. A few results from L^2 Hodge theory. — We recall briefly a few results of L^2 -Hodge theory for a complete manifold endowed with a Poincaré type metric, following closely [CP20]. We are in the following setting.

Let X be a n -dimensional compact Kähler manifold, and let (L, h_L) be a line bundle endowed with a (singular) metric $h_L = e^{-\phi_L}$ such that

- h_L has analytic singularities;
- Its Chern curvature satisfies $i\Theta_{h_L}(L) \geq 0$ in the sense of currents.

We consider a modification $\pi : \widehat{X} \rightarrow X$ of X such that the support of the singularities of $\varphi_L \circ \pi$ is a simple normal crossing divisor E . As usual, we can construct π such that its restriction to $\widehat{X} \setminus E$ is a biholomorphism. Then

$$(3.27) \quad \varphi_L \circ \pi|_{\Omega} \equiv \sum_{\alpha=1}^p e_{\alpha} \log |z_{\alpha}|^2$$

modulo a smooth function. Here $\Omega \subset \widehat{X}$ is a coordinate chart, and $(z_{\alpha})_{\alpha=1, \dots, n}$ are coordinates such that $E \cap \Omega = (z_1 \dots z_p = 0)$.

Let $\widehat{\omega}_E$ be a complete Kähler metric on $\widehat{X} \setminus E$, with Poincaré singularities along E , and let

$$(3.28) \quad \omega_E := \pi_{\star}(\widehat{\omega}_E)$$

be the direct image metric. We note that in this way (X°, ω_E) becomes a complete Kähler manifold, where $X^{\circ} := X \setminus (h_L = \infty)$.

Remark 3.5. — If u is a L -valued $(p, 0)$ -form on X° which is L^2 with respect to ω_E , then it is also L^2 with respect to an euclidean metric on \widehat{X} (or X , too). Therefore, if u is holomorphic, then it extends holomorphically to X and more generally any smooth compactification of X° .

The main goal of this section is to establish the following decomposition theorem, which is a slight generalization of the corresponding result in [CP20].

Theorem 3.6. — Consider a line bundle $(L, h_L) \rightarrow X$ endowed with a metric h_L with analytic singularities, as well as the corresponding complete Kähler manifold (X°, ω_E) , cf. (3.28). If $i\Theta_{h_L}(L) \geq 0$ on X , we have the following Hodge decomposition

$$L^2_{n,1}(X^\circ, L) = \mathcal{H}_{n,1}(X^\circ, L) \oplus \text{Im } \bar{\partial} \oplus \text{Im } \bar{\partial}^*.$$

Here $\mathcal{H}_{n,1}(X^\circ, L)$ is the space of L^2 Δ'' -harmonic $(n, 1)$ -forms.

The proof we will present here is almost contained in the aforementioned reference, in which the case $\Theta_{h_L}(L) = 0$ is treated. The few additional arguments together with the precise references we are using are explained below.

We start by recalling the following result.

Lemma 3.7. — There exists a family of smooth functions $(\mu_\varepsilon)_{\varepsilon>0}$ with the following properties.

- (a) For each $\varepsilon > 0$, the function μ_ε has compact support in X° , and $0 \leq \mu_\varepsilon \leq 1$.
- (b) The sets $(\mu_\varepsilon = 1)$ are providing an exhaustion of X° .
- (c) There exists a positive constant $C > 0$ independent of ε such that we have

$$\sup_{X^\circ} (|\partial\mu_\varepsilon|_{\omega_E}^2 + |\bar{\partial}\mu_\varepsilon|_{\omega_E}^2) \leq C.$$

We have also the Poincaré type inequality for the $\bar{\partial}$ -operator acting on $(p, 0)$ -forms.

Proposition 3.8. — [CP20] Let $(\Omega_j)_{j=1,\dots,N}$ be a finite union of coordinate sets of \widehat{X} covering E , and let \widehat{U} be any open subset contained in their union and $U := \pi(\widehat{U})$. Let τ be a $(p, 0)$ -form with compact support in a set $U \setminus \pi(E) \subset X$ and values in (L, h_L) . Then we have

$$(3.29) \quad \frac{1}{C} \int_U |\tau|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} \leq \int_U |\bar{\partial}\tau|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E}$$

where C is a positive numerical constant.

We emphasize that the constant C in (3.29) only depends on the distortion between the model Poincaré metric on Ω_j with singularities on E and the global metric $\widehat{\omega}_E$ restricted to Ω_j . Another important observation is that by using the cut-off function μ_ε in Lemma 3.7, we infer that (3.29) holds in fact for any L^2 -bounded form with compact support in U .

- Quick recap around the Bochner-Kodaira-Nakano formula.

We recall the following formula, which is central in complex differential geometry

$$(3.30) \quad \Delta'' = \Delta' + [i\Theta_{h_L}(L), \Lambda_{\omega_E}]$$

where $\Delta'' = \bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial}$ and $\Delta' = D'D'^* + D'^*D'$ where D' is the $(1, 0)$ -part of the Chern connection on (L, h_L) . Let us also recall the well-known fact that the self-adjoint operator

$$A := [i\Theta_{h_L}(L), \Lambda_{\omega_E}]$$

is semi-positive when acting on (n, q) forms, for any $0 \leq q \leq n$ as long as $i\Theta_{h_L}(L) \geq 0$. An immediate consequence of (3.30) is that for a L^2 -integrable form u with values in L of any type in the domains of Δ' and Δ'' , we have

$$(3.31) \quad \|\bar{\partial}u\|_{L^2}^2 + \|\bar{\partial}^*u\|_{L^2}^2 = \|D'u\|_{L^2}^2 + \|D'^*u\|_{L^2}^2 + \int_{X^\circ} \langle Au, u \rangle dV_{\omega_E}$$

where $\|\cdot\|_{L^2}$ (resp. $\langle \cdot, \cdot \rangle$) denotes the L^2 -norm (resp. pointwise hermitian product) taken with respect to (h_L, ω_E) , cf e.g. [BDIP02, Proposition 12.2 c)]. Here the main point is that in case of complete manifolds, the forms with compact support are dense in the domains of the operators $\bar{\partial}$ and $\bar{\partial}^*$ (unlike e.g. in the case of bounded domains in \mathbb{C}^n , where further boundary requirements for forms are to be imposed so that Bochner formula holds true).

Let $\star : \Lambda^{p,q} T_{X^\circ}^* \rightarrow \Lambda^{n-q, n-p} T_{X^\circ}^*$ be the Hodge star with respect to ω_E ; we introduce for any integer $0 \leq p \leq n$ the space

$$(3.32) \quad H^{(p)} := \{F \in H^0(X^\circ, \Omega_{X^\circ}^p \otimes L) \cap L^2; \int_{X^\circ} \langle A \star F, \star F \rangle dV_{\omega_E} = 0\}$$

and we can observe by Bochner formula that for a L^2 integrable, L -valued $(p,0)$ -form F , one has

$$(3.33) \quad \Delta''(\star F) = 0 \iff \Delta'(\star F) = 0 \text{ and } \int_{X^\circ} \langle A \star F, \star F \rangle dV_{\omega_E} = 0 \iff F \in H^{(p)}.$$

Indeed, the holomorphicity of F follows from the identity $\bar{\partial}F = -\star D'^*(\star F)$.

The proof of Theorem 3.6, which we give below, makes use of the following proposition which is the $\bar{\partial}$ -version of the Poincaré inequality established in [Auv17].

Proposition 3.9. — *Let $p \leq n$ be an integer. There exists a positive constant $C > 0$ such that the following inequality holds*

$$(3.34) \quad \int_{X^\circ} |u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} \leq C \left(\int_{X^\circ} |\bar{\partial}u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} + \int_{X^\circ} \langle A \star u, \star u \rangle dV_{\omega_E} \right)$$

for any L -valued form u of type $(p,0)$ which belongs to the domain of $\bar{\partial}$ and which is orthogonal to the space $H^{(p)}$ defined by (3.32). Here \star is the Hodge star operator with respect to the metric ω_E .

Proof of Proposition 3.9. — If a positive constant as in (3.34) does not exist, then we obtain a sequence u_j of L -valued forms of type $(p,0)$ orthogonal to $H^{(p)}$ such that

$$(3.35) \quad \int_{X^\circ} |u_j|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} = 1, \quad \lim_j \int_{X^\circ} |\bar{\partial}u_j|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} = 0$$

and

$$\lim_j \int_{X^\circ} \langle A \star u_j, \star u_j \rangle dV_{\omega_E} = 0.$$

It follows that the weak limit u_∞ of (u_j) is holomorphic and belongs to $H^{(p)}$. On the other hand, each u_j is perpendicular to $H^{(p)}$, so it follows that u_∞ is equal to zero.

Let us first show that the weak convergence $u_i \rightharpoonup u_\infty$ also takes place in $L_{\text{loc}}^2(X^\circ)$. To that purpose, let us pick a small Stein open subset $U \Subset X^\circ$. By solving the $\bar{\partial}$ -equation on U , we can find w_j such that $\bar{\partial}w_j = \bar{\partial}u_j$ on U and $\int_U |w_j|^2 \rightarrow 0$. Therefore $u_j - w_j$ is holomorphic on U and converges weakly, hence strongly to $u_\infty|_U$. In particular u_j converges to u_∞ in L^2 on U . As $u_\infty = 0$, we have

$$(3.36) \quad u_j|_K \rightarrow 0$$

in L^2 for any compact subset $K \subset X^\circ$.

The last step in the proof is to notice that the considerations above contradict the fact that the L^2 norm of each u_j is equal to one. This is not quite immediate, but is precisely as the end of the proof of Lemma 1.10 in [Auv17], so we will not reproduce it here. The idea is however very clear: in the notation of Proposition 3.8, we choose V small enough so that it admits a cut-off function χ with small gradient with respect to ω_E . Then, we decompose each u_j as $u_j = \chi u_j + (1 - \chi)u_j$. Then the L^2 norm of χu_j is small by (3.36). The L^2 norm of $(1 - \chi)u_j$ is equally small by (3.29), and this is how we reach a contradiction. \square

We have the following direct consequences of Proposition 3.9.

Corollary 3.10. — *There exists a positive constant $C > 0$ such that the following inequality holds*

$$(3.37) \quad \int_{X^\circ} |u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} \leq C \left(\int_{X^\circ} |\bar{\partial}u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} \right)$$

for any L -valued form u of type $(n, 0)$ which belongs to the domain of $\bar{\partial}$ and which is orthogonal to the kernel of $\bar{\partial}$.

Proof. — This follows immediately from Proposition 3.9 combined with the observation that the curvature operator A is equal to zero in bi-degree $(n, 0)$. \square

The next statement shows that in bi-degree $(n, 2)$ the image of the operator $\bar{\partial}^*$ is closed.

Corollary 3.11. — *There exists a positive constant $C > 0$ such that the following holds true. Let v be a L -valued form of type $(n, 2)$. We assume that v is L^2 , in the domain of $\bar{\partial}$ and orthogonal to the kernel of the operator $\bar{\partial}^*$. Then we have*

$$(3.38) \quad \int_{X^\circ} |v|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} \leq C \int_{X^\circ} |\bar{\partial}^*v|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E}.$$

Proof. — Let us first observe that the Hodge star $u := \star v$, of type $(n - 2, 0)$, is orthogonal to $H^{(n-2)}$. This can be seen as follows. Let us pick $F \in H^{(n-2)}$; it follows from (3.33) that we have $\bar{\partial}^*(\star F) = 0$. In other words, $\star F \in \text{Ker } \bar{\partial}^*$. We thus have

$$\int_{X^\circ} \langle u, F \rangle dV_{\omega_E} = \int_{X^\circ} \langle v, \star F \rangle dV_{\omega_E} = 0.$$

Applying Bochner formula (3.31) to v and using the facts that $\bar{\partial}v = 0$ (since v is orthogonal to $\text{Ker } \bar{\partial}^*$) and that $\bar{\partial}^*u = 0$ for degree reasons, we get

$$(3.39) \quad \|\bar{\partial}^*v\|_{L^2}^2 = \|\bar{\partial}u\|_{L^2}^2 + \int_{X^\circ} \langle A \star u, \star u \rangle dV_{\omega_E}$$

This proves the corollary by applying Proposition 3.9. \square

We discuss next the relative version of the previous estimates. Let $p : \mathcal{X} \rightarrow D$ and (L, h_L) be the family of manifolds and the line bundle, respectively fixed in the previous section. We assume that

$$(3.40) \quad D \ni t \mapsto \dim(\text{Ker}(\Delta_t'')) \quad \text{is constant}$$

where the Laplace operator Δ_t'' is the one acting on L^2 $(n,1)$ -forms with respect to (ω_E, h_L) .

The next result is a consequence of the proof of Proposition 3.9.

Corollary 3.12. — *Under the additional assumptions (3.40) and (A.1), there exists a constant $C > 0$ independent of t such that*

$$(3.41) \quad \int_{X_t^c} |u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} \leq C \left(\int_{X_t^c} |\bar{\partial}u|_{\omega_E}^2 e^{-\phi} dV_{\omega_E} + \int_{X_t^c} \langle A \star u, \star u \rangle dV_{\omega_E} \right)$$

for all L^2 forms u orthogonal to the space $H_t^{(p)}$ defined in (3.32) on the fiber X_t .

Here the constant C is uniform in the sense that for any subset U of compact support in D , we can find a constant C depending on U such that (3.41) is satisfied for any $t \in U$

Proof. — We first show that every form F_0 on the central fiber which is in the space $H_0^{(p)}$ can be written as limit of $F_{t_i} \in H_{t_i}^{(p)}$. This is of course well-known in the compact case, but we include a proof here since we could not find a reference fitting in our context. Let $(F_t)_{t \in D^*}$ any family of L -valued holomorphic p -forms on the fibers above the pointed disk D^* such that

$$(3.42) \quad \int_{X_t^c} |F_t|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} = 1.$$

Then we can definitely extract a limit F_∞ on the central fiber X_0 , but in principle it could happen that $F_\infty \equiv 0$ is identically zero. Such assumption would lead however to a contradiction, as follows.

We write locally on a coordinate chart Ω for \mathcal{X}

$$(3.43) \quad F_t|_\Omega = \sum f_I dz_I \otimes e_L,$$

where the coefficients f_I are holomorphic, and of course depending on t . We can assume that the multiplier ideal sheaf of h_L is trivial, given the transversality conditions that we have imposed (we can simply divide F_t with the corresponding sections). If the weak limit of F_t is zero, we can certainly extract a limit in strong sense, because the L^2 norm with respect to a smooth metric is smaller than the L^2 norm with respect to Poincaré metric, cf. also Remark 3.5.

In this case, the sup norm of the coefficients f_I above converges to zero as $t \rightarrow 0$. Since the Poincaré metric we are using has uniformly bounded volume, the equality (3.42) will not be satisfied as soon as $t \ll 1$.

We now take an orthonormal basis $(F_{t,j})$ of the space $H_t^{(p)}$ (this is obtained by the \star_t of an orthonormal basis for the $\text{Ker}(\Delta_t'')$, for example). The previous considerations will allow us to construct by extraction an orthonormal family $(F_{\infty,j})$ in $H_0^{(p)}$; this will be a basis because of dimension considerations.

We argue by contradiction and assume that the smallest constant C_t for which (3.34) holds true to for the fiber X_t tends to infinity when $t \rightarrow 0$. Then we get u_i on X_{t_i} such

that u_i is orthogonal to the space $H_{t_i}^{(p)}$ and such that

$$(3.44) \quad \int_{X_{t_i}} |u_i|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} = 1, \quad \lim_i \int_{X_{t_i}} |\bar{\partial} u_i|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} = 0.$$

Then the limit (computed using any C^∞ trivialization of the family)

$$u_0 := \lim_{i \rightarrow 0} u_i$$

is still orthogonal to $H_0^{(p)}$: this is exactly where the previous considerations are needed. The rest of the proof of the corollary follows the arguments already given for Proposition 3.9, so we simply skip it. \square

Now we can prove Theorem 3.6.

Proof of Theorem 3.6. — This statement is almost contained in [Dem12, chapter VIII, pages 367-370]. Indeed, in the context of complete manifolds one has the following decomposition

$$(3.45) \quad L_{n,1}^2(X^\circ, L) = \mathcal{H}_{n,1}(X^\circ, L) \oplus \overline{\text{Im } \bar{\partial}} \oplus \overline{\text{Im } \bar{\partial}^*}.$$

We also know (see *loc. cit.*) that the adjoints $\bar{\partial}^*$ and D'^* in the sense of von Neumann coincide with the formal adjoints of $\bar{\partial}$ and D' respectively.

It remains to show that the range of the $\bar{\partial}$ and $\bar{\partial}^*$ -operators are closed with respect to the L^2 topology. In our set-up, this is a consequence of the particular shape of the metric ω_E at infinity (i.e. near the support of $\pi(E)$): we are simply using the inequalities (3.37) and (3.38). The former shows that the image of $\bar{\partial}$ is closed, and the latter does the same for $\bar{\partial}^*$. \square

We finish this section with the following result (relying on the decomposition theorem obtained above), identifying the L^2 -integrable Δ'' -harmonic forms of bi-degree $(n, 1)$ on (X°, ω_E, h_L) with the vector space $H^1(X, K_X \otimes L \otimes \mathcal{I}(h_L))$ - which is independent of ω_E .

Proposition 3.13. — *In the setting of Theorem 3.6, we have a natural isomorphism*

$$\mathcal{H}_{n,1}(X^\circ, L) \xrightarrow{\simeq} H^1(X, K_X \otimes L \otimes \mathcal{I}(h_L))$$

where $\mathcal{H}_{n,1}(X^\circ, L)$ is the space of L^2 integrable, Δ'' -harmonic $(n, 1)$ -forms on X° .

Proof. — We proceed in several steps.

- *Step 1. Reduction to the snc case.*

The first observation is that the statement is invariant by blow-up whose centers lie on $X \setminus X^\circ$. It is obvious for the LHS while it follows from the usual formula $\pi_*(K_{X'} \otimes \pi^*L \otimes \mathcal{I}(\pi^*h_L)) \simeq K_X \otimes L \otimes \mathcal{I}(h_L)$ as well as Grauert-Riemenschneider vanishing $R^1\pi_*(K_{X'} \otimes \pi^*L \otimes \mathcal{I}(\pi^*h_L)) = 0$ (see [Mat16, Cor. 1.5]) valid for any modification $\pi : X' \rightarrow X$. So from now on, we assume that the singular locus of h_L is an snc divisor. In the following, we pick a finite Stein covering $(U_i)_{i \in I}$ of X .

- *Step 2. Statement of the claim to solve the $\bar{\partial}$ -equation.*

Our main tool in the proof will be the following estimate

Claim 3.14. — Let v be a $(n, 1)$ -form on X° with values in (L, h_L) , and such that

$$\Delta''v = 0, \quad \int_X |v|_{\omega_E}^2 e^{-\varphi_L} dV_{\omega_E} < \infty.$$

Then for each coordinate set $\Omega \subset X$ there exists an $(n, 0)$ -form u on Ω such that

$$(3.46) \quad \bar{\partial}u = v, \quad \int_\Omega |u|_{\omega_E}^2 e^{-\varphi_L} dV_{\omega_E} < \infty,$$

For bi-degree reasons, the $(n, 0)$ -form u in (3.46) is L^2 with respect to h_L (independently of any background metric). We postpone the proof of the claim for the moment and we will use it in order to prove Proposition 3.13.

• *Step 3. The map "harmonic to cohomology".*

We first construct an application

$$(3.47) \quad \Phi : \mathcal{H}_{n,1}(X^\circ, L) \longrightarrow \check{H}^1(X, K_X \otimes L \otimes \mathcal{I}(h_L))$$

as follows. Let $f \in \mathcal{H}_{n,1}(X^\circ, L)$; by definition we have $\Delta''f = 0$. Therefore, one can solve on each $U_i^\circ := U_i \cap X^\circ$ the equation $\bar{\partial}u_i = f$ where u_i is an L -valued $(n, 0)$ -form on U_i° satisfying the condition (3.46). In particular, the form $u_{ij} := u_i - u_j$ is a holomorphic L -valued n -form on U_{ij}° such that

$$\int_{U_{ij}^\circ} |u_{ij}|^2 e^{-\varphi_L} \leq 2 \int_{U_{ij}^\circ} |f|_{\omega_E}^2 e^{-\varphi_L} dV_{\omega_E}$$

It follows that u_{ij} extends holomorphically across E as a section of $K_X \otimes L \otimes \mathcal{I}(h_L)$ on U_{ij} and therefore it defines a 1-cocycle of the latter sheaf. It is straightforward to check that the class

$$\Phi(f) := \{(u_{ij})_{i,j \in I}\} \in \check{H}^1(X, K_X \otimes L \otimes \mathcal{I}(h_L))$$

is independent of the choice of the L^2 -integrable form u_i solving $\bar{\partial}u_i = f$.

• *Step 4. The map "cohomology to harmonic".*

Next, we have a natural morphism

$$(3.48) \quad \Psi : \check{H}^1(X, K_X \otimes L \otimes \mathcal{I}(h_L)) \longrightarrow \mathcal{H}_{n,1}(X^\circ, L)$$

Indeed, given a cocycle $v := (v_{ij})_{i,j \in I}$ and a partition of unity $(\theta_k)_{k \in K}$ we use the Leray isomorphism and consider as usual the L -valued $(n, 0)$ -form

$$(3.49) \quad \tau_k := \sum_{i \in I} \theta_i v_{ki} \quad \text{on } U_k$$

and then the local L -valued forms of type $(n, 1)$

$$\bar{\partial}\tau_k$$

are gluing on overlapping sets. Let β_v be the resulting form. We have

$$(3.50) \quad \bar{\partial}\beta_v = 0 \quad \text{and} \quad \beta_v \in L^2,$$

where the second property in (3.50) is due to the fact that $\omega_E \geq \omega$. Under the canonical decomposition

$$\text{Ker } \bar{\partial} = \mathcal{H}_{n,1} \oplus^\perp \text{Im } \bar{\partial}$$

from Theorem 3.6, we define $\Psi(v)$ to be the orthogonal projection of β_v onto $\mathcal{H}_{n,1}$. It is clear that Ψ above is well-defined: if $v_{ij} = v_i - v_j$, then $\tau_k - v_k$ is a global, L^2 form and our β_v is exact and therefore its projection onto the kernel of Δ'' is zero.

• *Step 5. Compatibility of the maps.*

We are left to showing that the maps Φ and Ψ in (3.47) and (3.48) are inverse to each other. Let $f \in \mathcal{H}_{n,1}$, $u_i \in L^2$ such that $\bar{\partial}u_i = f$ on U_i° and $u = (u_{ij})$. Then on U_k° , one has

$$\beta_u - f = \bar{\partial}\left(\sum_i \theta_i u_{ki} - u_k\right) = \bar{\partial}\left(-\sum_{i \in I} \theta_i u_i\right)$$

and that last form is globally exact in X° and L^2 , hence $\Psi(\Phi(f)) = f$.

In the other direction, let $v := (v_{ij})_{i,j \in I}$ be a cocycle and let us write $\beta_v = \Psi(v) + \bar{\partial}w$ for some L^2 -integrable $(n,0)$ -form w . On U_k , one has $\Psi(v) = \bar{\partial}(\tau_k - w)$ so that $\Phi(\Psi(v))$ is represented by the cocycle $(\tau_i - \tau_j)_{i,j \in I} = v$.

• *Step 6. Proof of Claim 3.14.*

In order to complete the proof of Proposition 3.13, we need to prove the Claim 3.14 that we used in the course of the proof.

By (3.33), the form $\star v$ is holomorphic and its restriction to a coordinate subset Ω can be written as

$$(3.51) \quad \star v|_\Omega = \sum (-1)^{i-1} \alpha_i \widehat{d}z_i \otimes e_L, \quad \sum_j \int_\Omega \frac{|\alpha_j|^2}{|f_j|^2} e^{-\varphi_L} d\lambda < \infty$$

where the α_i are holomorphic on Ω and f_j is as in (3.2).

Then we have

$$(3.52) \quad v|_\Omega = (-1)^n \sum_{i,k} \alpha_i g_{i\bar{k}} dz \wedge d\bar{z}_k \otimes e_L$$

where $g_{i\bar{k}}$ are the coefficients of the metric ω_E . The construction of the metric at the beginning shows that

$$(3.53) \quad g_{i\bar{k}} = \frac{\partial^2}{\partial z_i \partial \bar{z}_k} \left(\phi - \sum_j \log \log \frac{1}{|s_j|^2} \right)$$

where ϕ is a local potential for the smooth metric ω . Therefore we can get a primitive for $v|_\Omega$ by defining

$$(3.54) \quad u = \sum_i \alpha_i \frac{\partial}{\partial z_i} \left(\phi - \sum_j \log \log \frac{1}{|s_j|^2} \right) dz \otimes e_L.$$

By equality (3.53) it verifies $\bar{\partial}u = v$ and in is also in L^2 as one can see by an direct explicit computation combined with the second inequality in (3.51).

The proof of Proposition 3.13 is now complete. \square

4. Curvature formulas and applications

In this section, we use the Set-up 2.1. We also borrow the Notation 2.4 for the L^2 metric denoted by $h_{\mathcal{F}}$ on the direct image bundle $\mathcal{F} = p_*(\mathcal{O}(K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L))$ induced by $e^{-\phi_L}$.

Let $u \in H^0(D, \mathcal{F})$ and let \mathbf{u} be a $(n, 0)$ -form on \mathcal{X}° representing u . Thanks to (3.9), for any smooth function $f(t)$ with compact support in D , we have

$$(4.1) \quad \int_D \|u\|_{h_{\mathcal{F}}}^2 \cdot dd^c f(t) = c_n \int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge dd^c f(t).$$

Recall that $h_{\mathcal{F}}$ is smooth by Lemma 2.2.

The aim of this section is to generalize formulas [Ber09, (4.4), (4.8)] to our singular setting, cf Proposition 4.1 and Proposition 4.5.

4.1. A general curvature formula. — In this context we establish the following general formula, which generalise the corresponding result in [Ber09, (4.4)].

Proposition 4.1. — *Let \mathbf{u} be a continuous representative of u such that:*

- (i) $\mathbf{u}, D'\mathbf{u}$ and $\bar{\partial}(D'\mathbf{u})$ are L^2 on \mathcal{X}° with respect to ω_E, h_L ,
- (ii) $\bar{\partial}\mathbf{u} \wedge \overline{\bar{\partial}\mathbf{u}}$ is L^1 on \mathcal{X}° with respect to ω_E, h_L .

Then the following formula holds true

$$(4.2) \quad \partial\bar{\partial}\|u\|_{h_{\mathcal{F}}}^2 = c_n \left[-p_*((\Theta_{h_L}(L))_{ac} \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L}) + (-1)^n p_*(D'\mathbf{u} \wedge \overline{D'\mathbf{u}} e^{-\phi_L}) \right. \\ \left. + (-1)^n p_*(\bar{\partial}\mathbf{u} \wedge \overline{\bar{\partial}\mathbf{u}} e^{-\phi_L}) \right]$$

Here $(\Theta_{h_L}(L))_{ac}$ is the absolutely continuous part of the current $\Theta_{h_L}(L)$.

Remark 4.2. — Here we merely require the L^1 integrability of $\bar{\partial}\mathbf{u} \wedge \overline{\bar{\partial}\mathbf{u}}$ and not the L^2 -integrability of $\bar{\partial}\mathbf{u}$. The reason is that for later application in Theorem 4.6, we could only obtain the former condition. It is not clear whether the term $\bar{\partial}\mathbf{u}$ in Theorem 4.6 is L^2 .

The proof of Proposition 4.1 will require a few preliminary computations and will be given on page 24 below. First, we start with the following result legitimizing integration by parts.

Lemma 4.3. — *If \mathbf{u} and $D'\mathbf{u}$ are L^2 on \mathcal{X}° , we have*

$$\int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\partial\bar{\partial}f(t) = - \int_{\mathcal{X}^\circ} D'\mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\bar{\partial}f(t).$$

Proof. — Let ψ_ε be the cut-off function in Lemma 3.7. Since \mathbf{u} is L^2 bounded with respect to ϕ_L and ω_E , we have

$$\int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\partial\bar{\partial}f(t) = \lim_{\varepsilon \rightarrow 0} \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\partial\bar{\partial}f(t).$$

An integration by parts yields

$$(4.3) \quad \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\partial\bar{\partial}f(t) = - \int_{\mathcal{X}^\circ} i\partial\psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \bar{\partial}f(t) \\ - \int_{\mathcal{X}^\circ} \psi_\varepsilon \wedge D'\mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\bar{\partial}f(t) - (-1)^n \int_{\mathcal{X}^\circ} \psi_\varepsilon \wedge \mathbf{u} \wedge \overline{\partial\mathbf{u}} e^{-\phi_L} i\bar{\partial}f(t).$$

Since \mathbf{u} is a representative of a holomorphic section u , we know by (3.11) that $\bar{\partial}\mathbf{u} = dt \wedge \eta$, hence

$$(4.4) \quad \bar{\partial}\mathbf{u} \wedge dt = 0$$

and the third term of RHS of (4.3) vanishes.

The first term of RHS of (4.3) tends to 0 because \mathbf{u} is assumed to be globally L^2 -integrable. Similarly, we see that the second term of RHS of (4.3) tends to

$$- \int_{\mathcal{X}^\circ} D'\mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} i\bar{\partial}f(t).$$

The lemma is thus proved. \square

As a corollary of Lemma 4.3 above, we can compute the Chern connection of $(\mathcal{F}, h_{\mathcal{F}})$ as follows.

Corollary 4.4. — *Let $u \in H^0(D, \mathcal{F})$ and let \mathbf{u} be a smooth representative of u . We have*

$$\nabla u = P(\mu)dt$$

where

- ∇ is the $(1,0)$ -part of the Chern connection on $(\mathcal{F}, h_{\mathcal{F}})$.
- μ is defined by $D'\mathbf{u} = dt \wedge \mu$, cf. (3.12), and $\mu|_{X_t}$ only depends on u .
- $P(\mu)$ is the fiberwise projection onto $H^0(X_t, (K_{X_t} + L_t) \otimes \mathcal{I}(h_{L_t}))$ with respect to the L^2 -norm.

Proof. — Let ∇ be the $(1,0)$ -part of the Chern connection of $(\mathcal{F}, h_{\mathcal{F}})$. Then we have

$$\nabla u = \sigma \otimes dt,$$

where $\sigma = \frac{\nabla u}{dt} \in \mathcal{C}^\infty(D, \mathcal{F})$. Let u, v be two holomorphic sections of \mathcal{F} and let f be a smooth function with compact support in D . Since v is a holomorphic, we have

$$\int_D \langle u, v \rangle i\partial\bar{\partial}f(t) = \int_D \langle \nabla u, v \rangle \wedge i\bar{\partial}f(t).$$

Let \mathbf{u} and \mathbf{v} be the representatives of u and v respectively given by (3.7). The argument already used in Lemma 4.3 shows that we have

$$\int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{v}} e^{-\phi_L} i\partial\bar{\partial}f(t) = \int_{\mathcal{X}^\circ} D'\mathbf{u} \wedge \bar{\mathbf{v}} e^{-\phi_L} i\bar{\partial}f(t).$$

Here D' is, as before, the Chern connection on $(L \rightarrow \mathcal{X}^\circ, h_L)$. As a consequence, we have

$$\int_D \langle \nabla u, v \rangle \wedge i\bar{\partial}f(t) = \int_{\mathcal{X}^\circ} D'\mathbf{u} \wedge \bar{\mathbf{v}} e^{-\phi_L} i\bar{\partial}f(t).$$

Since we can choose f on the base D arbitrarily, we infer

$$\begin{aligned} \int_{X_t^\circ} \langle \sigma_t, v_t \rangle &\stackrel{\text{def}}{=} \int_{X_t^\circ} \left\langle \frac{\nabla u}{dt}, v \right\rangle_t = \int_{X_t^\circ} \frac{D'u}{dt} \wedge \bar{v} e^{-\phi_L} \\ &= \int_{X_t^\circ} \mu \wedge \bar{v}_t e^{-\phi_L} \\ &= \int_{X_t^\circ} P(\mu|_{X_t}) \wedge \bar{v}_t e^{-\phi_L} \end{aligned}$$

for each $t \in D$.

As the above holds for any holomorphic section v , we obtain thus

$$\nabla u = P(\mu)dt$$

on D . □

We can now complete the proof of Proposition 4.1.

Proof of Proposition 4.1. — Let $f \in \mathcal{C}_c^\infty(D)$. By (4.4), we have

$$\int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} \bar{\partial} f(t) = 0$$

for every ε . By integration by parts, we obtain

$$(4.5) \quad \int_{\mathcal{X}^\circ} \bar{\partial} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t) + \int_{\mathcal{X}^\circ} \psi_\varepsilon \bar{\partial} \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t) + (-1)^n \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{D}' \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t) = 0.$$

For the first term of (4.5), by integration by parts again, we have

$$\begin{aligned} (-1)^n \int_{\mathcal{X}^\circ} \bar{\partial} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t) &= - \int_{\mathcal{X}^\circ} \bar{\partial} \psi_\varepsilon \wedge D' \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} f(t) + \int_{\mathcal{X}^\circ} \partial \bar{\partial} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} f(t) \\ &\quad - \int_{\mathcal{X}^\circ} \bar{\partial} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \partial f(t). \end{aligned}$$

Recall that $d\psi_\varepsilon$ and $dd^c\psi_\varepsilon$ are uniformly bounded with respect to ω_E and converge to zero pointwise. Since \mathbf{u} and $D'\mathbf{u}$ are L^2 by assumption, we see from Lebesgue dominated convergence theorem that the RHS tends to 0. Therefore the first term of (4.5) tends to 0.

Since $\bar{\partial} \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}}$ is L^1 , the second term of (4.5) tends to $\int_{\mathcal{X}^\circ} \bar{\partial} \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t)$. We obtain thus

$$(4.6) \quad \int_{\mathcal{X}^\circ} \bar{\partial} \mathbf{u} \wedge \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t) = (-1)^{n-1} \lim_{\varepsilon \rightarrow 0} \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{D}' \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} f(t).$$

We complete in what follows the proof of the proposition. We have

$$\begin{aligned} \int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \bar{\partial} \partial f(t) &= \lim_{\varepsilon \rightarrow 0} \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \bar{\partial} \partial f(t) \\ &= - \lim_{\varepsilon \rightarrow 0} \left[\int_{\mathcal{X}^\circ} \bar{\partial} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \partial f(t) + \int_{\mathcal{X}^\circ} \psi_\varepsilon \wedge \bar{\partial} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \partial f(t) \right. \\ &\quad \left. + (-1)^n \int_{\mathcal{X}^\circ} \psi_\varepsilon \wedge \mathbf{u} \wedge \bar{D}' \bar{\partial} \bar{\mathbf{u}} e^{-\phi_L} \wedge \partial f(t) \right] \end{aligned}$$

Note that the first term tends to 0 since \mathbf{u} is L^2 . The second term vanishes because of (4.4). Then we have

$$\int_{\mathcal{X}^\circ} \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} \wedge \bar{\partial} \partial f(t) = (-1)^{n-1} \lim_{\varepsilon \rightarrow 0} \int_{\mathcal{X}^\circ} \psi_\varepsilon \wedge \mathbf{u} \wedge \overline{D' \mathbf{u}} e^{-\phi_L} \wedge \partial f(t).$$

Applying again integration by part, the RHS above becomes

$$(4.7) \quad \lim_{\varepsilon \rightarrow 0} (-1)^{n-1} \int_{\mathcal{X}^\circ} \partial \psi_\varepsilon \wedge \mathbf{u} \wedge \overline{D' \mathbf{u}} e^{-\phi_L} f(t) + (-1)^{n-1} \int_{\mathcal{X}^\circ} \psi_\varepsilon D' \mathbf{u} \wedge \overline{D' \mathbf{u}} e^{-\phi_L} f(t) - \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \overline{\partial D' \mathbf{u}} e^{-\phi_L} f(t).$$

As \mathbf{u} and $D' \mathbf{u}$ are L^2 , the first term of (4.7) tends to 0, and the second term of (4.7) tends to $\int_{\mathcal{X}^\circ} D' \mathbf{u} \wedge \overline{D' \mathbf{u}} e^{-\phi_L} f(t)$. For the third term, as $\bar{\partial} D' \mathbf{u} = \Theta_{h_L}(L) - D' \bar{\partial} \mathbf{u}$, we have

$$(4.8) \quad \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \overline{\partial D' \mathbf{u}} e^{-\phi_L} f(t) = - \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \overline{D' \bar{\partial} \mathbf{u}} e^{-\phi_L} f(t) - \int_{\mathcal{X}^\circ} \psi_\varepsilon \Theta_{h_L}(L) \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} f(t).$$

Combining with (4.6), we obtain

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathcal{X}^\circ} \psi_\varepsilon \mathbf{u} \wedge \overline{\partial D' \mathbf{u}} e^{-\phi_L} f(t) = (-1)^n \int_{\mathcal{X}^\circ} \bar{\partial} \mathbf{u} \wedge \overline{\partial \mathbf{u}} e^{-\phi_L} f(t) - \int_{\mathcal{X}^\circ} \Theta_{h_L}(L) \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} f(t).$$

All three terms of the RHS of (4.7) have now been calculated and the sum is just

$$(-1)^{n-1} \int_{\mathcal{X}^\circ} D' \mathbf{u} \wedge \overline{D' \mathbf{u}} e^{-\phi_L} f(t) + (-1)^{n-1} \int_{\mathcal{X}^\circ} \bar{\partial} \mathbf{u} \wedge \overline{\partial \mathbf{u}} e^{-\phi_L} f(t) + \int_{\mathcal{X}^\circ} \Theta_{h_L}(L) \wedge \mathbf{u} \wedge \bar{\mathbf{u}} e^{-\phi_L} f(t).$$

The proposition is thus proved. \square

4.2. A characterization of flat sections. — Now for applications, we need to generalize [Ber09, Prop 4.2] and formula [Ber09, (4.8)] to our singular setting. Following the argument of [Ber09, Prop 4.2], we have the following.

Proposition 4.5. — *We assume that the coefficients b_i in the Set-up condition (A.2) are equal to zero. Let u be a holomorphic section of \mathcal{F} on D such that $\nabla u(0) = 0$. Then u can be represented by a smooth $(n, 0)$ -form \mathbf{u} on \mathcal{X}° , L^2 with respect to h_L, ω_E , such that*

$$\bar{\partial} \mathbf{u} = dt \wedge \eta$$

for some L^2 -form η which is primitive (with respect to $\omega_E|_{X^\circ}$) on X° , and

$$D' \mathbf{u} = dt \wedge \mu$$

for some μ satisfying $\mu|_{X^\circ} = 0$. Here $X^\circ := X_0 \cap \mathcal{X}^\circ$ is on the central fiber.

Proof. — Let \mathbf{u} be the representative constructed in (3.7). We have

$$D' \mathbf{u} = dt \wedge \mu \quad \bar{\partial} \mathbf{u} = dt \wedge \eta.$$

Then $\mu|_{X^\circ}$ is orthogonal to the space of L^2 -holomorphic section by Corollary 4.4.

By Remark 3.2 our representative \mathbf{u} has the following property

$$(4.9) \quad \mathbf{u} \wedge \omega_E = dt \wedge d\bar{t} \wedge u_1$$

for some $(n, 0)$ -form u_1 on \mathcal{X}° . It follows that we have

$$(4.10) \quad \eta \wedge \omega_E = 0$$

on each fiber X_t .

Moreover, as $\mu|_{X^\circ}$ is orthogonal to $\text{Ker } \bar{\partial}$, Theorem 3.6 shows that $\mu|_{X^\circ}$ is $\bar{\partial}^{\star_0}$ -exact, i.e., there exists a $\bar{\partial}$ -closed L^2 -form β_0 on X° such that

$$\bar{\partial}^{\star_0} \beta_0 = \mu|_{X^\circ}.$$

Let $\tilde{\beta}_0$ be an arbitrary (globally L^2) extension of $\star_0 \beta_0$. Then $\mathbf{u} - dt \wedge \tilde{\beta}_0$ is the representative we are looking for. \square

The result above produces a representative enjoying nice properties in restriction to the central fiber; more precisely, the fact that $\eta|_{X^\circ}$ is primitive ensure that the third term in the RHS of the curvature formula (4.2) has a (positive) sign at $t = 0$.

In order to generalize that to each fiber, we consider the case where $u \in H^0(D, \mathcal{F})$ is a flat section with respect to $h_{\mathcal{F}}$. For that purpose, we introduce an additional cohomological assumption.

In the Set-up 2.1, assume that the coefficients b_I appearing in (A.2) vanish. That is to say, h_L has analytic singularities in the usual sense. We let Δ_t'' be the Laplace operator on L^2 -integrable $(n, 1)$ -forms with values in L on X_t° , taken with respect to ω_E, h_L . Let us consider the following assumption.

(A.4) The dimension $\dim \text{Ker } \Delta_t''$ is independent of $t \in D$.

Note that by Bochner formula, we already know that $\dim \text{Ker } \Delta_t'' < +\infty$. Indeed, if α is a Δ_t'' -harmonic L^2 $(n, 1)$ -form on X_t° , then (3.31) shows that α is Δ_t' -harmonic and $(D')^* \alpha = 0$. In particular, $\star \alpha$ is a L^2 -holomorphic section. Thanks to Remark 3.5, we get an injection $\text{Ker } \Delta_t'' \hookrightarrow H^0(X_t, \Omega_{X_t}^{n-1} \otimes L_t)$. In particular, the former space is finite-dimensional.

The main result of this subsection states as follows.

Theorem 4.6. — *In the Set-up 2.1, assume that h_L has analytic singularities and that the condition (A.4) above is satisfied.*

Let $u \in H^0(D, \mathcal{F})$ be a flat section with respect to $h_{\mathcal{F}}$. Then, we can find a continuous $(n, 0)$ -form \mathbf{u} on $\mathcal{X} \setminus E$ representing u such that

- (i) \mathbf{u} is L^2 and $D' \mathbf{u} = 0$,
- (ii) $\eta|_{X_t^\circ} = 0$ for any $t \in D$, and $\bar{\partial} \mathbf{u} \wedge \overline{\bar{\partial} \mathbf{u}} = 0$, where η is –as usual– given by $\bar{\partial} \mathbf{u} = dt \wedge \eta$.
Moreover, the equality

$$(4.11) \quad \Theta_{h_L}(L) \wedge \mathbf{u} = 0$$

holds true point-wise on $\mathcal{X} \setminus E$.

Remark 4.7. — Let us collect a few remarks about the theorem.

- (a) The content of Theorem 4.6 is clear: it "converts" the abstract data $\bar{\partial} u = 0$ and $\nabla u = 0$ into an effective result.
- (b) The identity (4.11) is equivalent to saying that the hermitian metric induced by $i\Theta_{h_L}(L)$ on $\Lambda^n T_{\mathcal{X}^\circ}^*$ has \mathbf{u} in its kernel.

Proof. — As in the proof of Proposition 4.5, we start with a representative \mathbf{u} given by (3.7) (i.e. constructed via the contraction with the canonical lifting of $\frac{\partial}{\partial t}$ with respect ω_E).

Since u is flat on D , we have $D'\mathbf{u} = dt \wedge \mu$ where $\mu|_{X_t^\circ}$ is L^2 and $\bar{\partial}^{\star_t}$ -exact for every $t \in D$. Therefore we can solve the $\bar{\partial}^{\star_t}$ -equation fiberwise, namely there exists a unique L^2 -form β_t on X_t° such that β_t is orthogonal to the $\text{Ker } \bar{\partial}^{\star_t}$ and such that

$$\bar{\partial}^{\star_t} \beta_t = \mu|_{X_t^\circ}.$$

By taking the $\bar{\partial}$ in both side and taking into account the fact that β_t is orthogonal to the $\text{Ker } \bar{\partial}^{\star_t}$, we obtain

$$(4.12) \quad \Delta_t'' \beta_t = \bar{\partial}(\mu|_{X_t^\circ}) \text{ on } X_t^\circ \quad \text{and} \quad \beta_t \perp \text{Ker } \Delta_t''.$$

By analogy to the compact case it is expected that the minimal solution of a Δ_t'' equation varies smoothly provided that $\dim \text{Ker } \Delta_t''$ is constant. We partly confirm this expectation in Proposition 4.8 by showing that it is continuous; for the moment, we will admit this fact and finish the proof of the theorem.

We set

$$\mathbf{u}_1 := \mathbf{u} - dt \wedge (\star_t \beta_t).$$

It is a continuous, fiberwise smooth form on \mathcal{X}° and we show now that \mathbf{u}_1 is a representative for which the points (i)-(ii) above are satisfied.

By construction, $D'\mathbf{u}_1 = 0$ on \mathcal{X}° . By (4.18) of Proposition 4.8 below, the L^2 -norm of β_t is smaller than the L^2 -norm of $\bar{\partial}\mu|_{X_t^\circ}$, i.e.,

$$\|\beta_t\|_{L^2} \leq C \|\bar{\partial}\mu|_{X_t^\circ}\|_{L^2}$$

for some constant C independent of t . Moreover, we recall the estimates in Lemma 3.3: $\bar{\partial}\mu|_{X_t^\circ}$ is uniformly L^2 -bounded. Therefore $dt \wedge (\star_t \beta_t)$ is L^2 and so our representative \mathbf{u}_1 is L^2 .

We have $\bar{\partial}\mathbf{u}_1 = dt \wedge (\eta + \bar{\partial}(\star_t \beta_t))$ and since $\bar{\partial}(\star_t \beta_t) \wedge \omega_E = \bar{\partial}\beta_t = 0$, it follows that

$$(4.13) \quad \frac{\bar{\partial}\mathbf{u}_1}{dt} \Big|_{X_t} \wedge \omega_E = 0.$$

In order to use Proposition 4.1, we show next that we have $\bar{\partial}\mathbf{u}_1 \wedge \overline{\bar{\partial}\mathbf{u}_1} \in L^1$. To this end, we write

$$(4.14) \quad \bar{\partial}\mathbf{u}_1 \wedge \overline{\bar{\partial}\mathbf{u}_1} = dt \wedge \eta \wedge \overline{dt \wedge \eta} + dt \wedge \bar{\partial}(\star_t \beta_t) \wedge \overline{dt \wedge \eta}$$

$$(4.15) \quad + dt \wedge \eta \wedge \overline{dt \wedge \bar{\partial}(\star_t \beta_t)} + dt \wedge \bar{\partial}(\star_t \beta_t) \wedge \overline{dt \wedge \bar{\partial}(\star_t \beta_t)}.$$

By the estimates in Lemma 3.3, η is L^2 . Then the first term of RHS of (4.14) is L^1 . Degree considerations show that we have

$$(4.16) \quad dt \wedge \bar{\partial}(\star_t \beta_t) \wedge \overline{dt \wedge \eta} = dt \wedge \bar{\partial}_t(\star_t \beta_t) \wedge \overline{dt \wedge \eta},$$

where $\bar{\partial}_t$ is the $\bar{\partial}$ -operator on X_t . Since $\Delta_t'' \beta_t = \bar{\partial}\mu$ and β_t is of degree $(n, 1)$, Bochner formula shows that the L^2 -norm of $\bar{\partial}_t(\star_t \beta_t)$ is equal to the L^2 norm of the form $(D')^{\star_t} \beta_t$

(this is due to the fact that β_t is $\bar{\partial}$ -closed), which in turn is bounded by the L^2 -norm of $\bar{\partial}\mu$. Once again, the estimates provided by Lemma 3.3 show that the L^2 -norm of $\bar{\partial}\mu|_{X_t}$ is bounded uniformly with respect to t . It follows that $dt \wedge \bar{\partial}_t(\star_t \beta_t)$ is L^2 .

Therefore $dt \wedge \bar{\partial}(\star_t \beta_t) \wedge \overline{dt} \wedge \bar{\eta}$ is L^1 -bounded by using (4.16). The same type of arguments show that the two terms in (4.15) are also L^1 .

We apply Proposition 4.1 for the representative \mathbf{u}_1 of u . The flatness of u imply that

$$c_n \int_{\mathcal{X}^\circ} ((-1)^n \bar{\partial} \mathbf{u}_1 \wedge \overline{\bar{\partial} \mathbf{u}_1} + i\Theta_{h_L}(L) \wedge \mathbf{u}_1 \wedge \overline{\mathbf{u}_1}) e^{-\phi_L} = 0.$$

Thanks to the assumption $i\Theta_{h_L}(L) \geq 0$ combined with (4.13), both two terms in the integral are semi-positive, and modulo the continuity of the family (β_t) our result is proved. \square

The proof of the following result is very similar to the familiar situation in which the couple of metrics (ω_E, h_L) are non-singular. We provide a complete argument because we were unable to find a reference.

Proposition 4.8. — *The minimal solution β_t in (4.12) varies continuously with respect to t .*

Proof. — We have divided our proof in a few steps.

4.2.1. Step 1. — Let (u_t) be a family of L -valued, L^2 forms of $(n, 1)$ -type on the fibers of p , such that we have

$$(4.17) \quad \Delta_t'' v_t = u_t$$

on the fiber X_t . If moreover we assume that each v_t is perpendicular to $\text{Ker } \Delta_t''$, then we claim that

$$(4.18) \quad \int_{X_t} |v_t|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E} \leq C \int_{X_t} |u_t|_{\omega_E}^2 e^{-\phi_L} dV_{\omega_E}$$

for some constant C uniform with respect to t .

Indeed, our claim follows instantly from Corollary 3.12 and (3.33) applied to $u := \star_t v_t$.

4.2.2. Step 2. — Let $\lambda \in \mathbb{C}$ such that $0 < |\lambda| \ll 1$ –we will make this precise in a moment. We claim that the operator

$$A_{\lambda,t} := \lambda - \Delta_t''$$

is invertible, which we show by proving that the equation $A_{\lambda,t} v = u$ admits a solution v , as soon as u is in L^2 . This can be seen via the usual Riesz representation theorem, as follows.

We define on $L_{n,1}^2$ the functional

$$I(\phi) = \int_{X_t} \langle u, \phi \rangle e^{-\phi_L} dV_{\omega_E}$$

We write $u = u_1 + u_2$ and $\phi = \phi_1 + \phi_2$ according to the decomposition $L_{n,1}^2 = \text{Ker } \Delta_t'' \oplus (\text{Ker } \Delta_t'')^\perp$. Then we have

$$I(\phi) = \int_{X_t} \langle u_1, \phi_1 \rangle e^{-\phi_L} dV_{\omega_E} + \int_{X_t} \langle u_2, \phi_2 \rangle e^{-\phi_L} dV_{\omega_E}$$

and then the squared absolute value of the second integral is smaller than

$$(4.19) \quad \int_{X_t} |\Delta_t'' \phi|^2 e^{-\phi_L} dV_{\omega_E}$$

up to a uniform constant, by Step 1. Therefore, we get

$$(4.20) \quad |I(\phi)|^2 \lesssim \int_{X_t} |\phi_1|^2 e^{-\phi_L} dV_{\omega_E} + \int_{X_t} |\Delta_t'' \phi|^2 e^{-\phi_L} dV_{\omega_E}.$$

Since

$$\left| \int_{X_t} \langle \phi, \Delta_t'' \phi \rangle dV_{\omega_E} \right|^2 \leq \int_{X_t} |\phi_2|^2 e^{-\phi_L} dV_{\omega_E} \cdot \int_{X_t} |\Delta_t'' \phi|^2 e^{-\phi_L} dV_{\omega_E}$$

we see that from (4.20) that

$$(4.21) \quad |I(\phi)|^2 \leq C_{u,\lambda} \int_{X_t} |\lambda \phi - \Delta_t'' \phi|^2 e^{-\phi_L} dV_{\omega_E}$$

as we see from the previous step, provided that $|\lambda| \leq \frac{1}{2C}$, where C is the constant in (4.18). Moreover, $C_{u,\lambda}$ is of the form $C_\lambda \cdot \|u\|_{L^2}^2$.

Taking $\phi = u$ in the identity above, we see that $A_{\lambda,t} := \lambda - \Delta_t''$ is injective. Moreover, the functional

$$\begin{aligned} J : \text{Im } A_{\lambda,t} &\longrightarrow \mathbb{C} \\ A_{\lambda,t} \phi &\longmapsto I(\phi) \end{aligned}$$

is well-defined and continuous by (4.21). In particular, it extends to $F = \overline{\text{Im } A_{\lambda,t}}$ and Riesz theorem provides us with an element $v \in F$ satisfying

$$(4.22) \quad \forall \psi \in F, \quad J(\psi) = \int_{X_t} \langle v, \psi \rangle e^{-\phi_L} dV_{\omega_E} \quad \text{and} \quad \|v\|_{L^2} \leq C_\lambda \|u\|_{L^2}^2.$$

The equality $J(A_{\lambda,t} \phi) = I(\phi)$ for any ϕ in L^2 shows that $A_{\lambda,t} v = u$. This concludes this step.

4.2.3. Step 3. — Let $\lambda \in \mathbb{C}$ as in the previous step, and let u_t be a continuous L^2 -family. We show that $v_t := (\lambda - \Delta_t'')^{-1} u_t$ is continuous with respect to t (with respect to the L^2 -norm).

It would be sufficient to check the continuity at one point $0 \in D$. For any $\varepsilon > 0$ we define the form $v_{0,\varepsilon} := \mu_\varepsilon v_0$ with compact support in $X_0 \setminus E$. We then have

$$\|v_{0,\varepsilon} - v_0\|_{L^2} \leq \varepsilon, \quad \|(\lambda - \Delta_0'') v_{0,\varepsilon} - (\lambda - \Delta_0'') v_0\|_{L^2} \leq \varepsilon$$

by the properties of $(\mu_\varepsilon)_{\varepsilon>0}$.

We next construct a smooth extension v_ε of $v_{0,\varepsilon}$ as follows. Let $(\Omega_i)_{i \in I}$ be a finite covering of $p^{-1}(\frac{1}{2}D)$ by coordinate charts, and let $(\theta_i)_{i \in I}$ be a partition of unity subordinate to this covering. The L -valued form

$$v_\varepsilon := \mu_\varepsilon \sum_i \theta_i(z, t) v_{0,i}$$

extends $v_{0,\varepsilon}$ and it is compactly supported in $\mathcal{X} \setminus E$. Here we denote by $v_{0,i}$ is the local expression of $v_{0,\varepsilon}|_{X_0 \cap \Omega_i}$, extended trivially to Ω_i (note that this is still L^2).

Since the metrics ω_E, h_L are smooth in $\mathcal{X} \setminus E$,

$$u_{\varepsilon,t} := (\lambda - \Delta_t'')(v_\varepsilon|_{X_t})$$

is a smooth L^2 -family.

By the second part in (4.22) we have

$$\|v_t - v_\varepsilon|_{X_t}\|_{L^2} \leq C\|u_t - u_{\varepsilon,t}\|_{L^2}.$$

As u_t and $u_{\varepsilon,t}$ are continuous with respect to t , we infer that we have $\|v_t - v_\varepsilon|_{X_t}\|_{L^2} \leq C_\varepsilon \mathcal{O}(|t|) + C\|u_0 - u_{\varepsilon,0}\|_{L^2}$. It follows that we have

$$\|v_t - v_\varepsilon|_{X_t}\|_{L^2} \leq o(1) + C\varepsilon$$

as $|t| \rightarrow 0$. The -small- quantity $o(1)$ here depends on ε , but since by construction the family $v_\varepsilon|_t$ is continuous with respect to t and its continuity modulus is independent of ε we infer that v_t is continuous at 0.

Now we define the operator $\text{Pr}_t := \int_{\lambda \in \Gamma} (\lambda - \Delta_t'')^{-1} d\lambda$ where Γ is a small circle centered at 0. We have proved above that Pr_t is continuous with respect to t . Moreover, Pr_t coincides with the orthogonal projection onto $\text{Ker } \Delta_t''$: we postpone the proof of this claim for the moment, see the Remark 4.9 below.

4.2.4. Step 4. — This is the main step in the proof of the proposition. Let β_t be the $\bar{\partial}^*$ -solution on X_t in question. Then

$$\Delta_t'' \beta_t = \bar{\partial}(\mu|_{X_t^\circ}).$$

By the estimates in Proposition 4.5, the RHS is L^2 and continuous with respect to t . Let s_t be a continuous L^2 -family (continuous with respect to t) such that $s_0 = \beta_0$. Then we have

$$\lambda_0 s_t - \Delta_t'' \beta_t = \lambda_0 s_t - \bar{\partial}(\mu|_{X_t^\circ})$$

for every t , where $0 < |\lambda_0| \ll 1$ is fixed, small enough as in Step 2. By Step 3, we can find a continuous family γ_t such that

$$\lambda_0 s_t - \Delta_t'' \beta_t = \lambda_0 \gamma_t - \Delta_t'' \gamma_t$$

for every t . Then $\Delta_t''(\beta_t - \gamma_t^\perp) = \lambda_0(s_t - \gamma_t)$, where $\gamma_t^\perp := \gamma_t - \text{Pr}_t \gamma_t$ is the projection onto $(\text{Ker } \Delta_t'')^\perp$.

Now β_t is orthogonal to $\text{Ker } \Delta_t''$ by construction. Then $\beta_t - \gamma_t^\perp$ is orthogonal to $\text{Ker } \Delta_t''$. By Step 1 we thus have

$$\|\beta_t - \gamma_t^\perp\|_{L^2} \leq C\|s_t - \gamma_t\|_{L^2}.$$

Note that $s_0 = \gamma_0$ and s_t and γ_t are continuous, then $\|\beta_t - \gamma_t^\perp\|_{L^2} = o(1)$. By Step 4, γ_t^\perp is continuous, therefore β_t is continuous at 0.

Summing up, the continuity with respect to the L^2 norm of $(\beta_t)_{t \in D}$ is established.

4.2.5. Step 5. — We show here that the form

$$(4.23) \quad dt \wedge \star_t \beta_t$$

induced by the family $(\beta_t)_{t \in D}$ in (4.12) is continuous on $\mathcal{X} \setminus E$. This is a consequence of the fact that the family of operators $(\Delta_t'')_{t \in D}$ is smooth and it has a smooth variation when restricted to a compact subset $K \subset \mathcal{X} \setminus E$, combined with the continuity property established in Proposition 4.8.

Let $\Omega \Subset X_0 \setminus E$ be a small coordinate chart. We can interpret the $(\Delta_t'')_{t \in D}$ as family of operators on the forms defined on Ω , since p is locally trivial. Then we have

$$(4.24) \quad \Delta_0'' \beta_t = \bar{\partial} \mu_t + (\Delta_0'' - \Delta_t'')(\beta_t)$$

from which it follows that

$$(4.25) \quad \Delta_0''(\beta_t - \beta_0) = \bar{\partial} \mu_t - \bar{\partial} \mu_0 + (\Delta_0'' - \Delta_t'')(\beta_t).$$

The equality (4.25) combined with the usual a-priori estimates for the elliptic operators imply that

$$(4.26) \quad \|\beta_t - \beta_0\|_{W^2}^2 \leq C(\|\bar{\partial} \mu_t - \bar{\partial} \mu_0\|_{L^2}^2 + \|\beta_t - \beta_0\|_{L^2}^2) + \delta_t \|\beta_t\|_{W^2}^2$$

where $\delta_t \rightarrow 0$ as $t \rightarrow 0$. We infer that

$$(4.27) \quad \|\beta_t - \beta_0\|_{W^2}^2 \leq \delta_t$$

for some (other) function δ_t tending to zero.

The usual boot-strapping method implies that $\lim_{t \rightarrow 0} \beta_t = \beta_0$ smoothly on any compact subset in $\mathcal{X} \setminus E$. In global terms this translates as

$$dt \wedge \star_t \beta_t$$

is a continuous $(n, 0)$ -form on $\mathcal{X} \setminus E$, so our lemma is proved. \square

Remark 4.9. — For the sake of completeness, we provide the details for the fact that the linear operator $\text{Pr}_t := \int_{\lambda \in \Gamma} (\lambda - \Delta_t'')^{-1} d\lambda$ is the orthogonal projection onto $\text{Ker } \Delta_t''$. Let \mathcal{H}_t be the Hilbert space $(\text{Ker } \Delta_t'')^\perp$. We need to prove two points:

- (i) $\text{Pr}_t u = 0$ for any $u \in \mathcal{H}_t$;
- (ii) $\text{Pr}_t u = u$ for any $u \in \text{Ker } \Delta_t''$.

For the first point, we have the following equality

$$(4.28) \quad (\lambda - \Delta_t'')^{-1} = - \sum_{k \geq 0} \lambda^k \mathcal{G}_t^{k+1} \quad \text{on } \mathcal{H}_t$$

where \mathcal{G}_t is the inverse of the operator Δ_t'' restricted to \mathcal{H}_t . The sum in (4.28) is indeed convergent (for the operator norm), given the estimates (4.18) and the fact that λ belongs to the circle Γ of small enough radius.

It follows that we can exchange integration/sum and then we have

$$(4.29) \quad \int_{\lambda \in \Gamma} (\lambda - \Delta_t'')^{-1} u d\lambda = - \sum_{k \geq 0} \mathcal{G}_t^{k+1}(u) \int_{\lambda \in \Gamma} \lambda^k d\lambda$$

and this shows that $\text{Pr}_t(u) = 0$ for any $u \in \mathcal{H}_t$.

For the second point, let $u \in \text{Ker } \Delta_t''$. For any $\alpha \in \text{Ker } \Delta_t''$, we have

$$\left\langle \frac{u}{\lambda}, \alpha \right\rangle = \left\langle \frac{(\lambda - \Delta_t'')(\lambda - \Delta_t'')^{-1}u}{\lambda}, \alpha \right\rangle = \langle (\lambda - \Delta_t'')^{-1}u, \alpha \rangle.$$

Therefore

$$\langle u, \alpha \rangle = \int_{\lambda \in \Gamma} \left\langle \frac{u}{\lambda}, \alpha \right\rangle d\lambda = \int_{\lambda \in \Gamma} \langle (\lambda - \Delta_t'')^{-1}u, \alpha \rangle d\lambda.$$

Then $u - \text{Pr}_t u$ is orthogonal to $\text{Ker } \Delta_t''$. On the other hand, thanks to the equality

$$\Delta_t'' \circ (\lambda - \Delta_t'')^{-1}(u) = (\lambda - \Delta_t'')^{-1} \circ \Delta_t'' u = 0,$$

we know that $\text{Pr}_t u \in \text{Ker } \Delta_t''$. Therefore $\text{Pr}_t u = u$.

Remark 4.10. — Actually the form β_t can be obtained as usually via an integral formula,

$$(4.30) \quad \beta_t = - \int_{\lambda \in \Gamma} \frac{1}{\lambda} (\lambda - \Delta_t'')^{-1} (\bar{\partial}\mu|_{X_t^c})$$

which gives the hope that its variation with respect to t is actually smooth. This can probably be obtained along the same lines as in [Kod86, Thm 7.5] modulo the fact that in the present situation, we have to deal with the additional difficulty induced by the fact that we are working with singular metrics ω_E and h_L .

We can now end this section by providing a proof of Theorem A.

Proof of Theorem A. — Up to shrinking D to a punctured disk $D_1 \subset D$, one may assume that the assumptions (A.1)-(A.3) are satisfied (with $b_I = 0$ for each I), cf. Section 2.

Next, there exists another punctured disk $D_2 \subset D_1$ such that the coherent sheaf $R^1 p_*(K_{\mathcal{X}/D} \otimes L \otimes \mathcal{I}(h_L))$ is locally free and commutes with base change; i.e. its fiber at $t \in D_2$ is given by $H^1(X_t, K_{X_t} \otimes L|_{X_t} \otimes \mathcal{I}(h_L|_{X_t}))$ and the dimension of the latter is independent of $t \in D_2$. Thanks to Proposition 3.13, the dimension of the space of harmonic $(n, 1)$ -forms, i.e. $\dim \text{Ker } (\Delta_t'')$, is independent of $t \in D_2$. In other words, the condition (A.4) is satisfied over D_2 .

Theorem A is now a direct consequence of Theorem 4.6. \square

5. A lower bound for the curvature in case of a -relatively- big twist

Let $p : \mathcal{X} \rightarrow D$ be a smooth, projective family, and let $L \rightarrow \mathcal{X}$ be a line bundle endowed with a metric $h_L = e^{-\phi_L}$ satisfying the following requirements.

(B.1) There exist a smooth, semi-positive real $(1,1)$ -form ω_L as well as an effective \mathbb{R} -divisor E_0 on \mathcal{X} such that

$$i\Theta_{h_L}(L) = \omega_L + [E_0]$$

where we denote by $[E_0]$ the current of integration associated to the \mathbb{R} -divisor E_0 .

(B.2) ω_L is relatively Kähler, i.e., $\omega_L|_{X_t} > 0$ for every t .

(B.3) The support of the divisor $E := \text{Supp}(E_0)$ is snc, and transverse to the fibers of p .

Let

$$c(\phi_L) := \frac{\omega_L^{n+1}}{\omega_L^n \wedge \text{idt} \wedge \bar{d}\bar{t}}$$

be the so-called geodesic curvature associated to ω_L .

Our goal here is to establish the following result.

Theorem 5.1. — *Under the assumptions (B.1)-(B.3) above, let $(\mathcal{F}, h_{\mathcal{F}})$ be the direct image bundle $p_*((K_{\mathcal{X}/D} + L) \otimes \mathcal{I}(h_L))$ endowed with the L^2 metric. Then for every $u \in H^0(D, \mathcal{F})$ and every $t \in D$, the following inequality holds*

$$(5.1) \quad \langle \Theta_{h_{\mathcal{F}}}(\mathcal{F})u, u \rangle_t \geq c_n \int_{X_t} c(\phi_L)u \wedge \bar{u}e^{-\phi_L}$$

where we identify $\Theta_{h_{\mathcal{F}}}(\mathcal{F})$ with an endomorphism of \mathcal{F} by "dividing" with $\text{idt} \wedge \bar{d}\bar{t}$.

Prior to providing the arguments for Theorem 5.1 we propose here the following problem.

Question 5.2. — *We assume that \mathcal{Y} is a foliation on a Kähler manifold Z . In which cases the bundle $K_{\mathcal{Y}} + L$ admits a positively curved metric? That is to say, is there some analogue of the Bergman metric on twisted relative canonical bundles in the more general contexts of a foliation? If yes, can we equally obtain a lower bound of the curvature form?*

Proof of Theorem 5.1. — The idea of our proof is to construct an approximation of the metric h_L so that the resulting absolutely continuous part of the associated curvature form has Poincaré singularities along the support of E . Then we can use the curvature formulas we have obtained in the previous sections, and finally conclude by a limit argument.

5.0.1. Approximation of the metric. — Let $\varepsilon > 0$ be a (small) positive real number. We introduce the form

$$(5.2) \quad \omega_\varepsilon := \omega_L - \varepsilon \sum_{i \in I} dd^c \log \log \frac{1}{|s_i|^2}$$

where I is the set of irreducible components of E , and s_i cut outs exactly one of these for a given $i \in I$, following notation in Section 2. We note that ω_ε is positive and has Poincaré singularities along E as soon as the metrics h_i used to measure the norm of s_i are suitably scaled, which is what we assume from now on.

Next, we introduce the following weight on L

$$(5.3) \quad \phi_\varepsilon := \phi_L - \varepsilon \sum_{i \in I} \log \log \frac{1}{|s_i|^2}.$$

Clearly, ϕ_ε has generalized analytic singularities in the sense of (A.2) in Set-up 2.1 and it satisfies $dd^c \phi_\varepsilon = \omega_\varepsilon$.

The properties of $(\phi_\varepsilon)_\varepsilon$ are collected in the following statement.

Lemma 5.3. — *Let $(\Omega, (z_1, \dots, z_n, t = z_{n+1}))$ be a coordinate system on \mathcal{X} adapted to the pair (\mathcal{X}, E) as in Set-up 2.1. Then the following hold.*

(i) *The geodesic curvature $c(\phi_\varepsilon)$ is uniformly bounded from above.*

(ii) We have $\lim_{\varepsilon \rightarrow 0} c(\phi_\varepsilon) = c(\phi_L)$ point-wise on $\mathcal{X} \setminus E$.

(iii) For every ε small enough, the multiplier ideal sheaf of $h_\varepsilon := e^{-\phi_\varepsilon}$ coincides with $\mathcal{I}(h_L)$. Moreover, the induced L^2 metric, say H_ε on the direct image is smooth, and it converges to $h_{\mathcal{F}}$ as $\varepsilon \rightarrow 0$.

Proof. — For the point (i), we write $c(\omega_\varepsilon) = 1/\|dt\|_{\omega_\varepsilon}^2$ and the result follows from the transversality conditions (B.3) and e.g. the estimates for the coefficients provided in (3.1).

The point (ii) follows easily from the local smooth convergence $\phi_\varepsilon \rightarrow \phi_L$ on $\mathcal{X} \setminus E$ combined with the positivity requirement (B.2).

As for the third point (iii), the smoothness of H_ε and its convergence to $h_{\mathcal{F}}$ is a consequence of the transversality assumption (B.3) by the same arguments as for Lemma 2.2. As for the statement about multiplier ideal sheaves, one has clearly $\mathcal{I}(\phi_\varepsilon) \subset \mathcal{I}(\phi_L)$ while the reverse inclusion is an easy consequence of (B.1)-(B.3). \square

5.0.2. Application of the curvature formula. — We consider u a local holomorphic section of the bundle \mathcal{F} , and let \mathbf{u}_ε be the representative of u constructed in (3.7), by using the contraction with the vector field V_ε associated to the metric ω_ε .

Let

$$(5.4) \quad \bar{\partial}\mathbf{u}_\varepsilon = dt \wedge \eta_\varepsilon, \quad D'\mathbf{u}_\varepsilon = dt \wedge \mu_\varepsilon$$

where $D' = D'_\varepsilon$ is the Chern connection corresponding to (L, h_ε) . Moreover we have

$$\omega_\varepsilon \wedge \mathbf{u}_\varepsilon \wedge \bar{\mathbf{u}}_\varepsilon = c(\phi_\varepsilon) \mathbf{u}_\varepsilon \wedge \bar{\mathbf{u}}_\varepsilon \wedge p^*(dt \wedge d\bar{t}) \quad \text{on } \mathcal{X} \setminus E$$

by [Ber11, Lem 4.2]. Proposition 4.1 then gives

$$(5.5) \quad -\frac{\partial^2}{\partial t \partial \bar{t}}(\|u\|_{H_\varepsilon}^2) = c_n \int_{X_t} c(\phi_\varepsilon) \mathbf{u}_\varepsilon \wedge \bar{\mathbf{u}}_\varepsilon e^{-\phi_\varepsilon} + \int_{X_t} |\eta_\varepsilon|^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon} - \int_{X_t} |\mu_\varepsilon|^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon},$$

since η_ε is primitive on fibers of p . We discuss next the terms which occur in (5.5).

• The LHS of (5.5) is equal to

$$(5.6) \quad \langle \Theta_{H_\varepsilon}(\mathcal{F})u, u \rangle - \|P(\mu_\varepsilon)\|^2$$

by the usual formula of the Hessian of the norm of a holomorphic section of a vector bundle. Then (5.5) becomes

$$(5.7) \quad \langle \Theta_{H_\varepsilon}(\mathcal{F})u, u \rangle_t = c_n \int_{X_t} c(\phi_\varepsilon) \mathbf{u}_\varepsilon \wedge \bar{\mathbf{u}}_\varepsilon e^{-\phi_\varepsilon} + \int_{X_t} |\eta_\varepsilon|_{\omega_\varepsilon}^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon} - \int_{X_t} |\mu_\varepsilon^\perp|_{\omega_\varepsilon}^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon},$$

where $\mu_\varepsilon = P(\mu_\varepsilon) + \mu_\varepsilon^\perp$ is the L^2 decomposition of μ_ε according to the $\text{Ker } \bar{\partial}$ and its orthogonal.

• As observed in [Ber11, Lem 4.4], we have

$$(5.8) \quad \bar{\partial}\mu_\varepsilon = D'\eta_\varepsilon$$

—even if the curvature is not zero!—, and actually μ_ε^\perp is the solution of (5.8) whose L^2 norm is minimal. By [CP20, Thm 1.6], we have the precise estimate

$$(5.9) \quad \int_{X_t} |\mu_\varepsilon^\perp|_{\omega_\varepsilon}^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon} \leq \int_{X_t} |\eta_\varepsilon|_{\omega_\varepsilon}^2 e^{-\phi_\varepsilon} dV_{\omega_\varepsilon}$$

and then we get

$$(5.10) \quad \langle \Theta_{H_\varepsilon}(\mathcal{F})u, u \rangle_t \geq c_n \int_{X_t} c(\phi_\varepsilon)u \wedge \bar{u}e^{-\phi_\varepsilon}.$$

as consequence of (5.7).

• The last step in our proof is to notice that as the parameter ε approaches zero, the inequality (5.10) implies

$$(5.11) \quad \langle \Theta_{h_{\mathcal{F}}}(\mathcal{F})u, u \rangle_t \geq c_n \int_{X_t} c(\phi_L)u \wedge \bar{u}e^{-\phi_L}.$$

Indeed, we are using Lemma 5.3 for the LHS of (5.10) and Lemma 5.3 (i) combined with dominated convergence theorem for the RHS. Theorem 5.1 is proved. \square

In the last lines, we now explain how to deduce Theorem B from Theorem 5.1 above.

Proof of Theorem B. — We start by making the observation that if $\pi : \mathcal{X}' \rightarrow \mathcal{X}$ is a proper birational morphism inducing birational morphisms $X'_t \rightarrow X_t$, then one has $\int_{X'_t} c(\phi'_L)u' \wedge \bar{u}'e^{-\phi'_L} = \int_{X_t} c(\phi_L)u \wedge \bar{u}e^{-\phi_L}$, with the self-explanatory notation.

Therefore, by blowing up \mathcal{X} and restricting the family to a punctured disk $D_1 \subset D$, one can from now on assume that the conditions (B.1) and (B.3) are satisfied.

Now, one has to show that one can further assume that condition (B.2) is satisfied. This is a bit more involved and can be shown as follows.

Since $\int_{X_t} \omega_L^n > 0$ and ω_L is smooth, it follows from e.g. [Bou02] that $[\omega_L]$ is p -big. In particular, there exists a punctured disk $D_2 \subset D_1$, an effective, horizontal \mathbb{R} -divisor F and an ample \mathbb{R} -line bundle A on \mathcal{X} such that

$$(5.12) \quad [\omega_L] = A + F \quad \text{in } H^{1,1}(\mathcal{X}, \mathbb{R}).$$

After blowing-up once again and restricting to a smaller punctured disk $D_3 \subset D_2$, one can assume without loss of generality that $E + F$ is snc and transverse to the fiber. Of course, the pull-back of A is not ample anymore, but there exists an effective divisor G contained in the exceptional locus of the blow-up such that $A - G$ is ample. All in all, one will assume from now on that one has a decomposition (5.12) where A is ample and $E + F$ is snc and transverse to the fibers.

We pick a strictly psh smooth weight ϕ_A on A and set ϕ_E (resp. ϕ_F) for the singular psh weight on the corresponding \mathbb{R} -divisor.

For $\delta > 0$, we introduce the psh weight ϕ_δ on L defined by

$$\phi_\delta = (1 - \delta)\phi_L + \delta(\phi_A + \phi_F + \phi_E).$$

Clearly, ϕ_δ has analytic singularities along the divisor $E + F$ and $(dd^c \phi_\delta)_{\text{ac}}$ is a relative Kähler metric for any $\delta > 0$. That is, the metric $h_{L,\delta} := e^{-\phi_\delta}$ satisfies (B.2).

Thanks to Theorem 5.1, the proof of Theorem B will be complete once we show the following

Claim 5.4. — *With the notation above, one has*

- (i) $\mathcal{I}(\phi_\delta) = \mathcal{I}(\phi_L)$ for δ small enough.
- (ii) The L^2 metric H_δ induced by $h_{L,\delta}$ on \mathcal{F} is smooth and converges smoothly to $h_{\mathcal{F}}$ when $\delta \rightarrow 0$.

(iii) For any $t \in D_3$ and $u \in \mathcal{F}_t$, one has

$$\lim_{\delta \rightarrow 0} \int_{X_t} c(\phi_\delta) u \wedge \bar{u} e^{-\phi_\delta} = \int_{X_t} c(\phi_L) u \wedge \bar{u} e^{-\phi_L}.$$

Proof of Claim 5.4. — Since $\phi_L - \phi_E$ is smooth (its curvature is nothing but ω_L), we have $\mathcal{I}(\phi_L) = \mathcal{I}(\phi_E)$ and $\mathcal{I}(\phi_\delta) = \mathcal{I}(\phi_E + \delta\phi_F)$, which coincides with $\mathcal{I}(\phi_E)$ when δ is small enough. This shows (i).

The item (ii) can be proved along the same lines as Lemma 2.2, using the fact that $E + F$ is snc and transverse to the fibers.

As for item (iii), we have pointwise convergence $c(\phi_\delta) \rightarrow c(\phi_L)$ on a Zariski open set of each X_t , $t \in D_3$, cf. Definition 2.3. Moreover, the Kähler metric $(dd^c \phi_\delta)_{ac}$ on \mathcal{X} is uniformly bounded above by a fixed Kähler metric on \mathcal{X} (for instance, $\omega_L + dd^c \phi_A$). In particular $c(\phi_\delta)$ is uniformly bounded above (say over compact subsets of D_3) and one can apply Lebesgue dominated convergence theorem to conclude. \square

The proof of Theorem B is now complete. \square

Remark 5.5. — The following limit argument shows that we can take $D^* \subset D$ to be the set of $t \in D$ such that the following hold:

- the metric $h_{\mathcal{F}}$ is smooth locally near t ;
- the fiber \mathcal{F}_t coincides with $H^0(X_t, (K_{X_t} + L) \otimes \mathcal{I}(h_L|_{X_t}))$.

Let $0 \in D$ be a point which satisfies these requirements. Let $U \Subset \mathcal{X} \setminus (h_L = \infty)$ be any open subset of \mathcal{X} whose closure does not meet the singular locus of the metric h_L . Then we have

$$(5.13) \quad \int_{U \cap X_0} c(\phi_L) u \wedge \bar{u} e^{-\phi_L} = \lim_{t \rightarrow 0} \int_{U \cap X_t} c(\phi_L) u \wedge \bar{u} e^{-\phi_L}$$

since all the objects involved are non-singular.

The next observation is that since $h_{\mathcal{F}}$ is smooth near 0 -by assumption-, the function $t \mapsto \langle \Theta_{h_{\mathcal{F}}}(\mathcal{F})u, u \rangle_t$ is thus continuous at 0. Theorem B combined with (5.13) and the positivity of $c(\phi)$ shows that we have

$$\int_{U \cap X_0} c(\phi_L) u \wedge \bar{u} e^{-\phi_L} \leq \langle \Theta_{h_{\mathcal{F}}}(\mathcal{F})u, u \rangle(0).$$

It follows that the estimate in (5.1) of Theorem B extends across $0 \in D$ as well.

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