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Intrinsic timescales of variability in a marine plankton model

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Abstract

Phytoplankton are known to exhibit temporal variability in biomass and community composition. While physically driven sources of variability have been studied extensively, ecosystems often exhibit complicated intrinsic dynamics that are not as well understood. As a first step towards assessing the contribution of this intrinsic variability to the total variability in the ocean, we examine the temporal scales of intrinsic variability in a marine plankton model suitable for use in climate model projections. Our rationale is that a better understanding of the time scales over which intrinsic variability manifests could help in the attribution of observed variability. Our model includes multiple phytoplankton, dissolved inorganic nutrients, and zooplankton and supports two oscillatory mechanisms: "R-oscillations", corresponding to patterns of species succession and associated with changes in resources, and "Z-oscillations", corresponding to changes in total phytoplankton biomass due to predator-prey interactions.

Over a wide range of model parameters, we found that while Z-oscillations typically occurred on time scales not exceeding 60 days, R-oscillations ranged from roughly 100 to 900 days under predation-free conditions, and R-oscillations occurred on longer time scales when interacting with Z-oscillations. Thus the two kinds of oscillations can be easily distinguished. At high grazing rates, we identified aperiodic cases where the dominant period never resolved, with distinct regimes emerging over decadal (or longer) time scales. These chaotic regime shifts are likely highly dependent on the model parameters and structure. More work must be done to understand how these oscillations interact with physical forcings.

Keywords: Intrinsic variability, phytoplankton, predator-prey dynamics

1 1. Introduction

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Phytoplankton communities exhibit different temporal modes of variability both in total biomass and community structure. Understanding the drivers of this natural variability is a prerequisite for confidently identifying changes in the ecosystem attributable to anthropogenic forcing. Ocean color observations of chlorophyll-a (a proxy for phytoplankton biomass), for example, have revealed variability at intraseasonal (weeks to months), seasonal, and interannual (> 1 year) time scales over most open ocean regions (Behrenfeld et al., 2006; Martinez et al., 2009; Resplandy et al., 2009; Vantrepotte and Mélin, 2011; Thomalla et al.,

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2011; Demarcq et al., 2012; Mayot et al., 2017; Salgado-Hernanz et al., 2019; Keerthi et al., 2020; Huisman et al., 2006). In-situ measurements of phytoplankton species abundances, as well as observations from ocean color remote sensing (which has been used to infer phytoplankton functional types and sizes), reveal that community structure varies over a similar range of time scales to chlorophyll-a (Alvain et al., 2008; d'Ovidio et al., 2010; Demarcq et al., 2012; Rousseaux and Gregg, 2015; Mayot et al., 2017; Dakos et al., 2009). This wide range of time variability in phytoplankton communities is also seen in global biogeochemical models of the ocean (Aumont et al., 2018; Dutkiewicz et al., 2019).

Much of the variability found in marine phytoplankton communities can be attributed to variability in $\mathbf{15}$ the physical environment in which they evolve. A typical example is the spring phytoplankton bloom, which 16 is driven by seasonal changes in light and vertical mixing (Platt et al., 2009; Sathyendranath et al., 2015, and 17 $\mathbf{18}$ references therein), and which explains the largest part of the time variability (Demarcq et al., 2012). At 19 the intraseasonal timescale, variability can be driven by intraseasonal basin-scale climate modes such as the Madden-Julian oscillation (Resplandy et al., 2009), synoptic atmospheric forcing due to storms (Fauchereau $\mathbf{20}$ $\mathbf{21}$ et al., 2011; Carranza and Gille, 2014) and tropical cyclones (Menkes et al., 2016), or by oceanic mesoscale and submesoscale processes acting at even smaller spatial scales (Mahadevan et al., 2012; Lévy et al., 2018; $\mathbf{22}$ Poggiale et al., 2013; Keerthi et al., 2020). At lower frequencies (interannual to multidecadal timescales), part $\mathbf{23}$ $\mathbf{24}$ of the variability is forced by large-scale, low-frequency climate modes such as the Pacific Decadal Oscillation, the El Nino Southern Oscillation or the Atlantic Multidecadal Oscillation (Chavez et al., 2011, and references $\mathbf{25}$ 26 therein).

 $\mathbf{27}$ However, phytoplankton also exhibit complex dynamics even in the absence of external variability. Empir-28 ical evidence for such intrinsic variability is difficult in the natural environment because the external forcing $\mathbf{29}$ is always present. However, it can be seen in dedicated laboratory experiments where the experimental con-30 ditions are artificially held constant. For example, Fussmann et al. (2000) have demonstrated the potential $\mathbf{31}$ for predator-prey oscillations in a culture with one nutrient, one phytoplankton species and one zooplankton $\mathbf{32}$ species in a continuous flow-through chemostat. More recently, predator-prey oscillations have been observed 33 to persist over several predator-prey cycles in a controlled culture of planktonic rotifers and unicellular green 34 algae in the laboratory (Blasius et al., 2019; Hastings, 2020). Benincà et al. (2008) conducted a long-term (> 8 years) laboratory mesoscosm experiment with a more complex plankton food web (two nutrients and ten $\mathbf{35}$ 36 functional groups) that exhibited striking fluctuations of different periodicities (from 15 days to 225 days). 37 attributable only to interactions between species in the food web.

Phytoplankton intrinsic variability is a recurrent topic in theoretical ecology. Limit cycles — and some-38 times chaos (Benincà et al., 2015) — in phytoplankton populations are emergent properties of mathematical 39 models of interspecies interactions. These models, usually constructed to represent a system in a chemostat 40 under constant conditions (Gothlich and Oschlies, 2015), have shown that internal variability may result 41from two fairly different processes: competition between different phytoplankton species for limiting re-42sources (Tilman, 1977; Huisman and Weissing, 1999, 2001) and predator-prev interactions (Gilpin, 1979; $\mathbf{43}$ $\mathbf{44}$ Edwards and Brindley, 1996). The zero-dimensional food web models developed for these ecological studies $\mathbf{45}$ are similar to the ones used for global biogeochemical studies, where the plankton equations are solved in three dimensions and are embedded into an ocean circulation model. We might therefore expect internal 46

47 variability to manifest in these forced 3D models.

In the natural environment, it is difficult to disentangle internal variability from forced variability. An $\mathbf{48}$ improved understanding of phytoplankton intrinsic variability is a prerequisite to confidently identify the 49 phytoplankton trends attributable to physical forcings, which are likely to dominate in the open ocean (and 50in 3D models of the open ocean) but might not explain all of the observed variability. As a first step towards $\mathbf{51}$ assessing how much intrinsic variability in phytoplankton biomass and community structure contributes to 52the total variability in the ocean, our objective is to examine the temporal scales of intrinsic variability in a $\mathbf{53}$ marine plankton model. A better understanding of the time scales over which intrinsic variability manifests $\mathbf{54}$ could help in the attribution of observed variability. 55

In this study, we focus on the two types of intrinsic variability mentioned previously: those due to 56 57 competition between phytoplankton for multiple resources, which we call R-oscillations, and those due to $\mathbf{58}$ predator-prey interactions, which we call Z-oscillations. These two kinds of oscillations have been studied in $\mathbf{59}$ isolation, but never together to our knowledge. Here we ask specifically: (1) what are the time scales for each 60 type of oscillation, (2) how do they interact, and (3) what is their sensitivity to a realistic range of model 61 parameters? To address these questions, we have developed a moderately-complex ecosystem model that supports both kinds of intrinsic variability, and we use it to explore the time scales of variability that emerge 62 63 under constant environmental forcing in a chemostat system. The closest analogue of our chemostat scenario to a physical regime is a low-latitude eastern boundary upwelling region, such as the Benguela (Messié et al., 64 2009). High nutrient delivery via upwelling, combined with relatively low seasonality due to its proximity to 65 66 the equator, provide an optimal environment for any existing mechanisms of intrinsic variability to persist. 67 The present paper can be considered the first entry in a two-part study; the second step will be to investigate 68 how these intrinsic variations manifest under time-varying external forcings.

69 A key goal in constructing our model is to bridge the gap between highly idealized conceptual mod-70 els focused on specific ecological interactions, and the much more complex biogeochemical components of 71Earth-system models used for global climate simulations. The latter class of models achieved a milestone $\mathbf{72}$ with Fasham et al. (1990), who developed a nitrogen-based biogeochemical model used in a General Circula-73 tion Model for the North Atlantic Ocean (Sarmiento et al., 1993), and Aumont et al. (2003), who investigated 74 the effect of multi-nutrient limitation on the biological carbon pump. Building from these first biogeochemical $\mathbf{75}$ models, the PISCES model (Aumont et al., 2015) was incorporated into the IPSL-CM5 earth system model 76 as part of CMIP5 (Bopp et al., 2013). The even-more-complex Darwin ecosystem model (Follows et al., 2007) 77 was used by Prowe et al. (2014) to examine the effect of phytoplankton diversity on primary production in 78 response to environmental changes. Wherever possible, we follow common choices in the class of models described above in determining model representations of processes and parameter values. 79

Section 2 introduces the model that we will use, including the equations and the numerical integration scheme (a simplified version of the equations is also presented in Appendix A). Section 3 presents four example simulations: two that exhibit each type of oscillation in isolation, and two that demonstrate the complexity of the full model. Section 4 explores the sensitivity of the results to key parameters involved in the phytoplankton growth and decay rates, using multidimensional sweeps over model parameters adding up to 6400 separate simulations. The results demonstrate the emergent regimes arising from the underlying oscillations. Finally in section 5 we discuss the implications of these results, in particular how the emergent
 frequencies depend on the presence of and interaction between different kinds of oscillations. We discuss how
 the time scales of intrinsic variability that emerges compares with time scales of external forcings in the open
 ocean. Ideas regarding future experiments to address this topic are also presented.

90 2. A chemostat-like model for studying intrinsic variability

Here we develop a chemostat-like model, with a continuous inflow of nutrients and without variations 91 92 in light or temperature. This is arguably the simplest system that mimics a physical setting — something 93 roughly like a low-latitude Eastern boundary upwelling system (see e.g. Messié et al. (2009)), with a high rate of nutrient influx from the deep ocean and relatively low seasonality — yet permits only variability driven 94 95 by internal dynamics. The model shares common terms with other biogeochemical models, and all ecological parameters are assigned reasonable values based on the literature. While our model does support oscillatory 96 97 behavior that is not common in other models, we believe that these features are reasonably explained by the **98** general complexity of ecosystem dynamics.

To represent the dynamics discussed in the introduction, the model requires multiple nutrient, phytoplankton and zooplankton compartments. Huisman and Weissing (2001) have shown that R-oscillations start emerging when the number of species and the number of resources is equal or larger than three. And of course, Z-oscillations require the presence of at least one phytoplankton and one zooplankton species. As is common in biogeochemistry models of this kind (e.g. Dutkiewicz et al., 2009; Aumont et al., 2015), we include two size classes for phytoplankton: small species (gleaners) that thrive in nutrient-poor environments, and larger species (opportunists) that dominate in nutrient-rich environments.

106 Given these constraints, we therefore include two phytoplankton size classes with three phytoplankton each, for six phytoplankton in total: P_1^s through P_3^s are small species, while P_1^l through P_3^l are large 107108 species. The different phytoplankton compartments within the same size class represent distinct species (e.g. prochlorococcus and synechococcus are both picophytoplankton), rather than members of the same species 109with different ecological properties. Two species of zooplankton are also included: Z^s is a microzooplankton 110 species that consumes small phytoplankton, while Z^{l} is a larger mesozooplankton species that consumes 111 both large phytoplankton and microzooplankton. The model includes three dissolved inorganic resource 112compartments R_j , representative of nitrogen (R_1 , μ mol/m³ N), phosphorus (R_2 , μ mol/m³ PO₄), and iron 113 $(R_3, \mu \text{mol/m}^3 \text{ Fe})$. We restrict our model to three nutrients because it is the minimum number required to 114produce R-oscillations. Iron is chosen over silicate because the need for silicate is specific to large diatoms. 115116 Also, for simplicity, the model does not distinguish between different forms of nitrogen such as nitrate and am-117 monium. Phytoplankton and zooplankton compartments have units μ mol C, with stoichiometric conversions for the phytoplankton compartments described below. The model is configured in a chemostat setting, with 118relaxation of resource R_j toward target values S_j at a dilution rate τ . Biomass is removed through mortality 119 120and via outflux at the same dilution rate τ ; explicit remineralization and settling are omitted. Phytoplankton 121mortality is linear as is typical in these models (Cropp and Norbury, 2009). Quadratic zooplankton mortality 122parameterizes density-dependent losses such as viral infection, and also acts as a closure term for predation

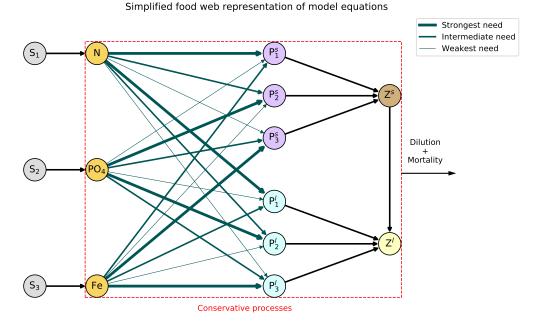


Figure 1: A food web representation of the model equations. Phytoplankton are controlled from the bottom-up by externally-supplied nutrients and from the top-down by zooplankton. The widths of the arrows pointing to a given phytoplankton depend on its relative needs for each nutrient. *Conservative processes*, or exchanges that do not result in any loss of mass from the system, are contained within the dashed red box. That is, without the deep nutrient supply; natural phytoplankton and zooplankton mortality terms; and loss of biomass through dilution, the system would exhibit biomass conservation.

- by species unrepresented in this model (Aumont et al., 2015). A model schematic is shown in Figure 1. Asimplified version of this model is provided in Appendix A.
- **125** 2.1. Model equations
- **126** The model equations are

$$\frac{dR_j}{dt} = \tau \left(S_j - R_j\right) - \sum_{i=1}^3 C_{ji} \left[\mu^s \min_j \frac{R_j}{K_{ji}^s + R_j} P_i^s + \mu^l \min_j \frac{R_j}{K_{ji}^l + R_j} P_i^l \right],\tag{1a}$$

$$\frac{dP_i^s}{dt} = -\tau P_i^s + \left[\mu^s \min_j \frac{R_j}{K_{ji}^s + R_j} - m_P - \frac{g^s Z^s}{K_Z + P_{\text{tot}}^s} \right] P_i^s,$$
(1b)

$$\frac{dP_i^l}{dt} = -\tau P_i^l + \left[\mu^l \min_j \frac{R_j}{K_{ji}^l + R_j} - m_P - \frac{g^l Z^l}{K_Z + Z^s + P_{\text{tot}}^l} \right] P_i^l,$$
(1c)

$$\frac{dZ^{s}}{dt} = -\tau Z^{s} + \frac{g^{s} Z^{s}}{K_{Z} + P_{\text{tot}}^{s}} P_{\text{tot}}^{s} - \left[\frac{g^{l} Z^{l}}{K_{Z} + Z^{s} + P_{\text{tot}}^{l}} + m_{Z} Z^{s}\right] Z^{s}$$
(1d)

$$\frac{dZ^{l}}{dt} = -\tau Z^{l} + \frac{g^{l} Z^{l}}{K_{Z} + Z^{s} + P_{\text{tot}}^{l}} \left(Z^{s} + P_{\text{tot}}^{l} \right) - m_{Z} (Z^{l})^{2}$$
(1e)

127 where i, j = 1, 2, 3, and $P_{\text{tot}}^{s,l} = \sum_{i=1}^{3} P_i^{s,l}$.

C

128 The phytoplankton growth rates are equal to a nutrient limitation factor scaled by the maximum growth 129 rate μ . Nutrient limitation is set dynamically by the smallest Monod growth factor, which depends on the 130 values of K_{ji} . Large phytoplankton species (opportunists) grow faster and require more nutrient than small 131 phytoplankton species (gleaners), allowing the former to thrive in high-nutrient environments, and the latter

Parameter	Default Value	Units	Description
τ	0.04	day^{-1}	Nutrient restoring rate ¹
S_{j}	18.1, 1.21, 1.21×10^{-3}	$[R_j]$	Deep ocean nutrient source values ¹
μ^s,μ^l	0.308, 0.616	day^{-1}	Maximum phytoplankton growth rates^2
m_P	0.1	day^{-1}	Phytoplankton mortality $rate^2$
r_j	$0.15, 0.01, 1 \times 10^{-5}$	$\frac{[R_j]}{\mu \mathrm{mol/m}^3 \mathrm{C}}$	Nutrient : Carbon ratios ³
c_{\min}, c_{\max}	0.9, 1.3		Stoichiometric scaling factors (Eq. 3)
$\kappa_j^s = \frac{1}{4}\kappa_j^l$	$0.15, 0.01, 1 \times 10^{-5}$	$[R_j]$	Nutrient uptake half-saturation $constants^2$
k_{\min}, k_{\max}	0.7, 1.1		Nutrient half-saturation scaling factors (Eq. 4)
g^s, g^l	1.5, 0.5	day^{-1}	Grazing rates ⁴
K_Z	10	$\mu { m mol}/{ m m^3}~{ m C}$	Grazing half-saturation $constant^4$
m_Z	0.015	$(\mu { m mol}/{ m m}^3~{ m C} \cdot { m day})^{-1}$	Zooplankton mortality rate ⁵

¹ Typical for the Benguela upwelling system (Messié et al., 2009) ² Follows et al. (2007) ³ $r_1 = N:C, r_2 = PO_4:C, r_3 = Fe:C$ (Redfield, 1934) ⁴ Aumont et al. (2015) ⁵ Messié and Chavez (2017)

Table 1: Definitions and default values for model parameters. Subscript j = 1, 2, 3. The notation $[R_j]$ means "the units of R_j ", which are $[R_1] = \mu \text{mol N}$, $[R_2] = \mu \text{mol PO}_4$, and $[R_3] = \mu \text{mol Fe}$. References for default values are given where appropriate in footnotes.

132 to overtake competitors in low-nutrient environments. This is modeled by setting $\mu^l = 2\mu^s$ (see Table 1 for 133 values) and $K_{ji}^l = 4K_{ji}^s$.

134 2.2. Nutrient source, stoichiometry, and uptake limitation

135 We set the nutrient source relaxation rate $\tau = 0.04 \text{ days}^{-1}$ and deep nitrogen concentration value $S_1 =$ 136 18.1 µmol N, which are characteristic of, for example, the Benguela upwelling system (Messié et al., 2009). 137 The other deep nutrient concentrations are derived from S_1 through the Redfield ratios r_i (Table 1), i.e.

$$\frac{S_1}{r_1} = \frac{S_2}{r_2} = \frac{S_3}{r_3}.$$
(2)

 $\mathbf{138}$ The phytoplankton stoichiometric coefficients C_{ji} in the nutrient uptake term convert increases in biomass 139 to corresponding losses in the inorganic nutrient pools; large and small phytoplankton share the same stoichiometry. The Redfield ratio is an average of the stoichiometry found across a variety of phyotoplankton 140species (Klausmeier et al., 2004), and so we include some variation in the stoichiometry for each modeled 141species. The stoichiometry matrix C_{ji} is constructed such that for a given nutrient and phytoplankton group, 142one species has exactly the Redfield value, another species has a slightly larger value controlled by $c_{\rm max} > 1$, 143and the third species has a slightly smaller value controlled by $c_{\min} < 1$. Dividing each row by r_j , the matrix 144145is

$$\frac{C_{ji}}{r_j} = \begin{pmatrix} c_{\max} & 1 & c_{\min} \\ c_{\min} & c_{\max} & 1 \\ 1 & c_{\min} & c_{\max} \end{pmatrix}$$
(3)

146 So for example, P_1 , P_2 , and P_3 have the largest, intermediate, and smallest stoichiometric values for nitrogen, 147 respectively. This is reflected in Figure 1, with the thickness of the lines between the phytoplankton and 148 nutrients corresponding to the magnitudes of the stoichiometric coefficients.

149 Half-saturation constants for a given nutrient and a specific size class (gleaners, say) are similar between 150 species in that size class. Denoting $\kappa_j^{s,l}$ as the typical half-saturation value for resource j and size class s or l151 (See Table 1 for values), the half-saturation matrix has a similar structure to the stoichiometric matrix, with

$$\frac{K_{ji}^{s,l}}{\kappa_j^{s,l}} = \begin{pmatrix} 1 & k_{\min} & k_{\max} \\ k_{\max} & 1 & k_{\min} \\ k_{\min} & k_{\max} & 1 \end{pmatrix}$$
(4)

- where $k_{\text{max}} > 1$ and $k_{\text{min}} < 1$ designate the entries with the largest and smallest coefficients per row (species), respectively.
- **154** 2.3. Model numerics and spinup

The model is written in Python 3.6 and integrated in time with the SciPy odeint library (Millman and Aivazis, 2011). This module acts as a wrapper for ODEPACK's LSODA solver that automatically switches
between Adams (non-stiff) and BDF (stiff) integration methods when appropriate (Hindmarsh, 1982; Petzold, 1983). The time series for each of the model compartments is stored in intervals of one day.

All simulations were run for 200 years, with the exception of one 500 year simulation to illustrate alonger-term phenomenon, depicted in Figure 6. We have observed that the spinup time for these simulations

is typically around 1-2 years. But because some of our simulations have very long periods of oscillation —
and may even exhibit aperiodic behavior — we took a conservative approach and only used the last 150 years
of each simulation for analysis.

164 3. Example simulations: Resource oscillations, predator-prey oscillations and their interactions

165 The model described above is designed to capture two primary mechanisms of intrinsic variability: competition between phytoplankton for resources (R-oscillations), and predator-prey interactions (Z-oscillations). 166 167In this section we consider four basic example simulations. The first two are designed to understand the 168 nature of each oscillation independently. The NoPred simulation removes both zooplankton, in order to 169 highlight how resource competition can give rise to coexistence amongst phytoplankton in a similar size 170class. The SimpleCom (for "simple competition") case retains only two phytoplankton species, one from each of the two size classes, thus allowing only predator-prey interactions. The third case, called *WeakPred*, is an 171172example of the full system that uses parameters given in Table 1. In this case, the two types of oscillations 173are both present, but remain distinct enough to isolate in their time-series and frequency spectra. The last 174case, termed StrongPred, employs increased grazing rates; the result is a simulation in which the types of oscillations interact in complex ways, producing time series that exhibit broadband chaotic behavior. 175

176 *3.1. Case NoPred: Isolated R-oscillations*

The first example, which we call *NoPred*, removes predators from the model by disabling all terms that contain Z^s and Z^l . Figure 2 shows time series and frequency spectra for the nutrients and small phytoplankton species. The spectra were computed using the discrete Fourier transform (DFT) after centering the data around zero and multiplying the remainder by a Hamming window. In the caption and throughout the paper, we use the term "dominant period" to denote the reciprocal of the frequency with the largest amplitude in the spectrum. If the dominant period is equal to the length of the series, then we say that the series has not equilibrated and we do not assign a dominant period.

184For this particular parameter regime, the large phytoplankton are not shown because they die out via competitive exclusion. The large phytoplankton have higher growth rates, but larger half-saturation coefficients, 185186which means they require greater ambient nutrient concentrations to compete successfully with gleaners; an increase in the nutrient input rate τ or deep nutrient concentrations S_j , for example, would lead to the 187 elimination of small phytoplankton instead. Peaks in the individual small phytoplankton alternate roughly 188189 every 90 days. Every time an individual species peaks, one also sees peaks in both the total phytoplankton 190 P_{tot}^s and in the nutrient compartments. The correspondence between a peak in PO₄ and a peak in P_2^s (as well as a peak in the total phytoplankton) is highlighted by the vertical red lines. The same correspondence 191 occurs between P_s^1 and N, and between P_s^3 and Fe. The oscillations of total phytoplankton biomass are small 192 compared with fluctuations in the individual compartments. 193

These R-oscillations are well described by the *resource-ratio hypothesis* of species succession (Tilman, 195 1985). Such oscillations depend on two factors: differences in the phytoplankton stoichiometric coefficients 196 C_{ji}/r_j (3) across the nutrient pools, and differences in the half-saturation coefficients K_{ji}/κ_j (4) between 197 species. The driving mechanism behind R-oscillations can be summarized in two steps:

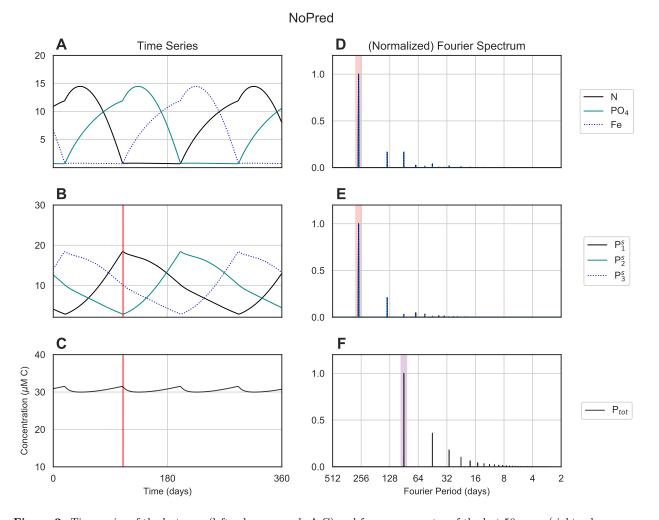


Figure 2: Time series of the last year (left column, panels A-C) and frequency spectra of the last 50 years (right column, panels D-F) from a 200 year simulation of the *NoPred* (isolated R-oscillations) scenario. The rows correspond to (top) nutrients, (middle) surviving phytoplankton (large phytoplankton do not survive in this scenario), and (bottom) total phytoplankton. See legends for curve identification. In panel B, the red vertical line corresponds to a peak in P_1^s , while the blue lines correspond to peaks in P_1^l . The purple lines in panel C correspond to peaks in the small and large zooplankton. These peak indicators were all copied into panel D to demonstrate that zooplankton peaks follow those of the phytoplankton. The red shading in panels D and E indicates the dominant periods for the phytoplankton and nutrients (all share the same peak at approx. 273 days), and the purple shading in panel F indicates the dominant period for the total phytoplankton (approx. 91 days).

- 198 1. The current dominant phytoplankton depletes the nutrient for which it has the largest C_{ji}/r_j , at a rate 199 proportional to c_{max} .
- 200 2. The phytoplankton with the largest growth rate at low values of the newly-depleted nutrient will become 201 the next dominant phytoplankton. By construction, this phytoplankton is the one with the minimum 202 half-saturation constant K_{ji}/κ_j (proportional to k_{\min}) for this given nutrient.

 $\mathbf{203}$ Let us say that P_1^s is dominant. For this phytoplankton, the largest nutrient need (largest C_{ji}/r_j) is for 204nitrogen. Thus, P_1^s depletes nitrogen first as it grows. This sets up P_2^s , which has the largest growth rate at $\mathbf{205}$ low values of nitrogen (lowest K_{ji}/κ_j) compared with its competitors, to become the next dominant species. 206 This two-step process repeats as P_2^s depletes phosphorus and P_3^s is primed to thrive when this nutrient is 207 limiting. Finally, P_3^s depletes iron, allowing P_1^s to thrive and completing a full cycle. In this case, c_{\max} $\mathbf{208}$ and k_{\min} are two key parameters that regulate this oscillatory mechanism. Increasing c_{\max} accelerates the 209 depletion of what will become the next limiting nutrient in the system. Decreasing k_{\min} enhances the growth 210 rate of the next dominant species, resulting in a faster invasion.

With the chosen parameter set, individual phytoplankton oscillations take about 273 days to complete, while total phytoplankton biomass cycles take one third as long (91 days) due to symmetries in the oscillations of the individual compartments; these symmetries emerge because the stoichiometric matrix (3) and halfsaturation matrix (4) are circulant. Also, note that fluctuations in the total phytoplankton concentrations are relatively small compared to oscillations in the individual compartments. This is because R-oscillations primarily represent variability at the species level.

217 *3.2. Case SimpleCom: Isolated Z-oscillations*

 The second example, which we call *SimpleCom*, is obtained by including both zooplankton, but just one large and one small phytoplankton with the same stoichiometric coefficients: P_1^s and P_1^l . In this case, nitrogen will always be the limiting nutrient, so the stoichiometric diversity that generates oscillations among competing species is removed, while predator-prey cycles are left intact.

222While R-oscillations primarily represent variability at the species level, Z-oscillations are a form of total biomass variability that occurs when zooplankton growth rates are sufficiently large. They induce significant 223224fluctuations in both the total phytoplankton and individual species concentrations. In panel G of Figure 3 we 225see that peaks in the zooplankton time series always follow peaks in the phytoplankton series. As is typical of a predator-prey model, zooplankton concentrations evolve in the direction of phytoplankton biomass changes, 226227while phytoplankton concentrations evolve inversely to zooplankton biomass changes. In this case, because of the zooplankton size selectivity, the small and large phytoplankton alternate dominance. The large and $\mathbf{228}$ 229small zooplankton alternate dominance as well, with each following peaks in their target prey. But small 230zooplankton are restricted to small bursts following peaks in small phytoplankton, as large zooplankton $\mathbf{231}$ exhibit top-down control on their biomass. The dominant period of the total phytoplankton in this scenario is about 40 days. The dominant period of P_1^s is about 81 days, while the double-peaked large phytoplankton $\mathbf{232}$ 233has a dominant period of about 27 days.

SimpleCom

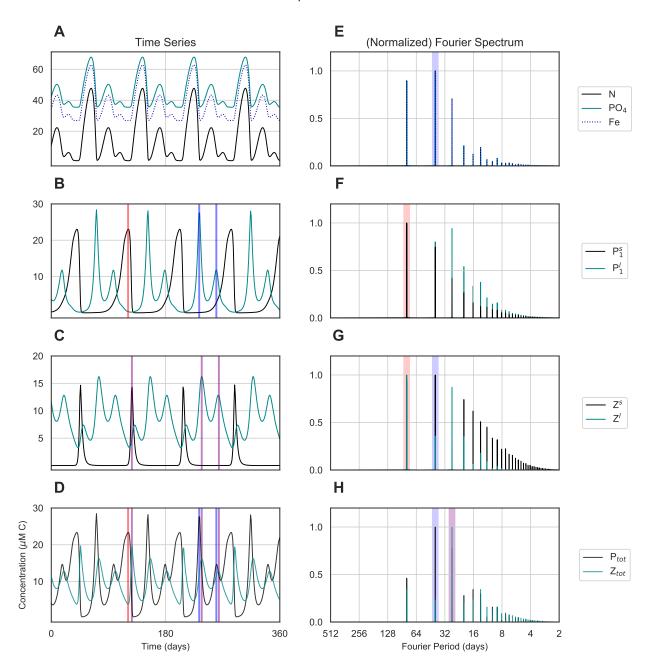


Figure 3: Time series of the last year (left column, panels A-E) and frequency spectra of the last 50 years (right column, panels F-J) from a 200 year simulation of the *SimpleCom* (isolated Z-oscillations) experiment. The rows correspond to (top) nutrients, (second) phytoplankton, (third) zooplankton, and (bottom) total phytoplankton and total zooplankton. See legends for curve identification. The red vertical line in panels B and C corresponds to a peak in P_1^s . In the spectra plots, dominant periods are indicated by: red shadings for individual phytoplankton and large zooplankton (approx. 81 days); blue shadings for nutrients, total phytoplankton, and small zooplankton (approx. 40 days); and purple shadings for nutrients and total phytoplankton (approx. 27 days).

WeakPred

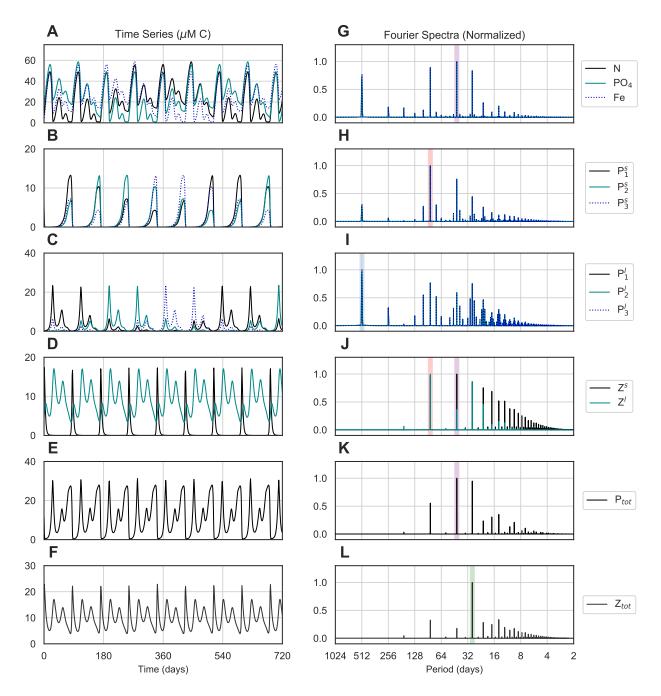


Figure 4: The time series (first column) and frequency spectra (second column) for phytoplankton, nutrient, and zooplankton compartments from a 200 year run of the *WeakPred* simulation. See legends for curve identification. The last two years of the time series are shown, and the last 150 years of the simulation were used to compute the spectra. In the spectra plots, dominant periods are indicated by: red shadings for the small phytoplankton and large zooplankton (approx. 85 days); blue shadings for the large phytoplankton (approx. 509 days); green shadings for the total zooplankton (approx. 28 days); and purple shadings for the total phytoplankton and small zooplankton (approx. 43 days).

234 3.3. Case WeakPred: The full system with distinct R- and Z-oscillations

Here we consider an example simulation in which all compartments of the full model equations (1) are active, using the parameters given in Table 1. The results can be seen in Figure 4. From the plots of N, P_1^s and P_1^l , we see that both R- and Z-oscillations are present, and can be easily distinguished. The low frequency peak in panel I, with a period of approximately 514 days, is indicative of R-oscillations: analogous species P_1^s and P_1^l evolve together, while nitrogen is depleted and restored. A similar pattern exists between phosphate, P_2^s and P_2^l , and amongst iron, P_3^s and P_3^l .

By contrast, the high frequency patterns of alternating dominance between the small and large phytoplankton resemble those of the *SimpleCom* simulation; the large phytoplankton exhibit a double-peaked pattern, as before, while the small phytoplankton are single-peaked, with low concentrations of biomass away from the peaks. And like in the *SimpleCom* simulation, peaks in the zooplankton series follow their phytoplankton counterparts. The dominant period of the total phytoplankton, whose peaks reflect those of the alternating small and large phytoplankton, is 28 days.

247 3.4. Case StrongPred: The full system with chaotic interactions

In this last example, the grazing pressure is increased relative to the WeakPred simulation, with grazing $\mathbf{248}$ rates g^s, g^l set to 1.15 times the values given in Table 1. The result is a simulation that exhibits more 249 $\mathbf{250}$ complex behavior, as shown by the time series and frequency spectra in Figure 5. The dominant period of $\mathbf{251}$ the total phytoplankton and zooplankton (purple vertical shading in spectra plots) is roughly 31 days, and 252the dominant period of the nutrient and small zooplankton (green vertical shading) is roughly 64 days. There 253are several properties that distinguish this scenario from the *WeakPred* case. First, the spectrum is much $\mathbf{254}$ more broadband, with dominant periods appearing more like clusters than individual peaks. Second, the time series exhibits low frequency variations in amplitude that are not present in WeakPred: an entire WeakPred 255cycle is visible within a two year window, whereas that is not the case for the *StrongPred* simulation. Third, in 256257WeakPred (and in SimpleCom), Z-oscillations take place between analogous small and large phytoplankton.

 $\mathbf{258}$ For the individual phytoplankton, the dominant period never resolves. The reason for this becomes clear 259in Figure 6, which shows the last 50 years of an extended 500 year StrongPred run. The series exhibits 260 heteroclinic cycles of increasing periodicity until roughly 35 years into the plotted series (indicated by the $\mathbf{261}$ orange vertical line), where there is a marked shift in the phytoplankton and nutrient series. From this $\mathbf{262}$ point until the end of the run, P_l^l is present at a low concentration and continues to oscillate with the other species. This regime shift indicates an aperiodicity in the series, making the Fourier decomposition a less 263 useful tool for understanding the nature of the oscillations. This is confirmed by looking at the frequency 264 $\mathbf{265}$ spectra in panels D, E, and F (taken from the last 400 years of the 200 year run). We see that the spectrum $\mathbf{266}$ is quite broadband compared with the narrow high frequency peaks observed in previous simulations, shown 267in Figures 4 and 5.

268 4. Intrinsic time scales of community change

269 In this section we use large sets of simulations to sample swaths of parameter space and explore the nature270 of the emergent frequencies. Specifically, we vary a subset of parameters involved in the phytoplankton growth

StrongPred

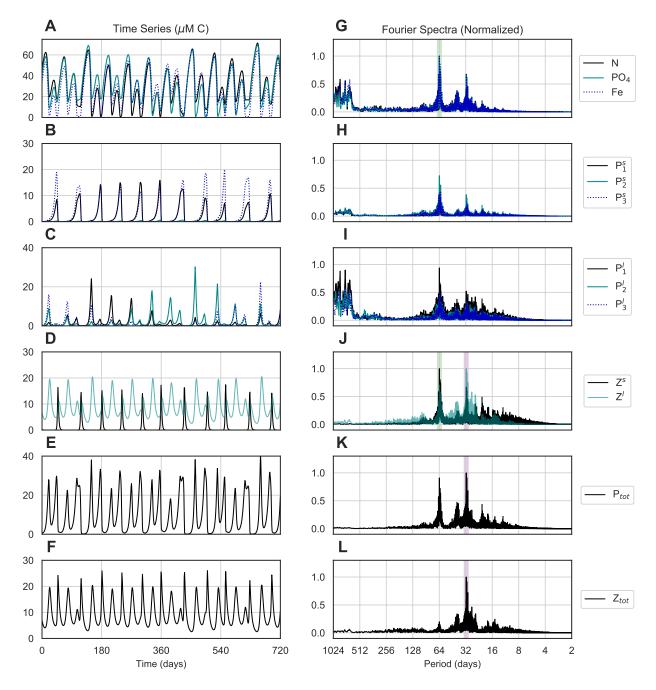


Figure 5: The time series and frequency spectra for the *StrongPred* simulation, with the same layout as for Figure 4. Two dominant periods occur at approximately 64 days for nutrients and small zooplankton (green shading) and 31 days for total phytoplankton and large/total zooplankton (purple shading). But there is also a low frequency component that does not fully resolve, as indicated by the lack of a dominant period in the individual phytoplankton spectra.

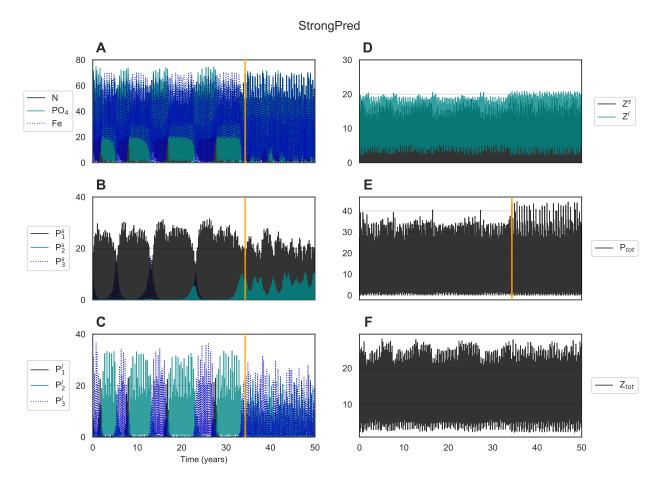


Figure 6: The last 50 years of an extended 500 year *StrongPred* simulation, highlighting a notable regime shift indicated by the orange vertical line. This shift is preceded by clear heteroclinic cycles, characterized by oscillations whose periods lengthen with each iteration, in the nutrient and phytoplankton compartments.

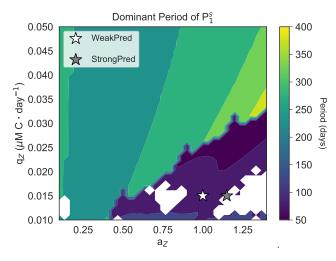


Figure 7: Dominant period for P_1^s as a function of zooplankton mortality m_Z and grazing rate scaling factor a_Z . There are three primary regimes: domination by slow R-oscillations; domination by fast Z-oscillations; and aperiodic regimes where the dominant period never resolves (shown in white).

and decay rates, i.e. grazing rates, zooplankton mortality rate, the nondimensional scaling factor c_{max} for the stoichiometry matrix (3), and the nondimensional scaling factor k_{\min} for the half-saturation matrix (4).

273 We first focus on controlling Z-oscillations by varying the quadratic zooplankton mortality rate m_Z , and 274 the grazing rates g^s and g^l . We used the same multiplicative coefficient a_Z to scale the large and the small 275 grazing rate as $a_Z g^s$ and $a_Z g^l$ (e.g. the *StrongPred* example used $a_Z = 1.15$). The parameter range used for 276 these sweeps is

$$0.01 \le m_Z \le 0.05$$
 and $0.1 \le a_Z \le 1.5$, (5)

where recall that the values for m_Z have units μ mol C day⁻¹. For each sweep, we assigned 40 possible linearly-spaced values per parameter, for a total of 1600 simulations. Each simulation was run for 200 years and the last 150 were used for analysis.

Figure 7 shows the dominant period of P_1^s as a function of m_Z and a_Z . There are three primary regimes: 280 281 domination by slow R-oscillations; domination by fast Z-oscillations; and regimes where a dominant period never resolves over the 150 year period considered (shown in white). In the regime where R-oscillations $\mathbf{282}$ 283 dominate, the dominant period tends to increase with the grazing pressure. When zooplankton reduce phytoplankton biomass, they also reduce the rate at which the phytoplankton deplete nutrients. This increases 284the time it takes a new nutrient to become limiting and a new species to become dominant. Furthermore, $\mathbf{285}$ the next dominant species invades at a slower rate because phytoplankton growth rates are inhibited. Thus, 286287grazing interferes with the species succession mechanism.

We can better understand how enhanced grazing affects species succession by examining the dominant period over a range of saturation and stoichiometric coefficients. The time-scale of R-oscillations depends on the value of the parameters that control these oscillations, i.e. in our case c_{max} and k_{min} (refer to section 3.1 for an explanation). To examine how these key parameters impact the frequency of oscillations, we started from the default set of parameters in Table 1 and varied them over a range that respects their constraints:

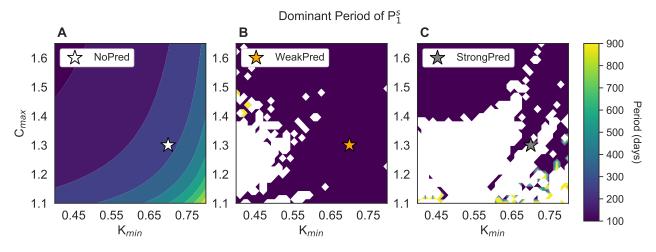


Figure 8: Dominant period of P_1^s for sweeps over c_{\max} and k_{\min} , in order of increasing grazing pressure from left to right. White regions indicate cases where no dominant period emerges, due to aperiodicity. As the grazing pressure increases, the dominant period is less likely to resolve over the considered range of parameters.

293 $k_{\min} < 1 \text{ and } c_{\max} > 1$:

$$0.4 \le k_{\min} \le 0.8$$
 and $1.1 \le c_{\max} \le 1.65.$ (6)

294 The range chosen is just large enough to clearly illustrate the dependence on the parameters over a range of295 time scales, from intraseasonal to interannual.

Again, we assigned 40 possible linearly-spaced values for each of the sweep parameters, k_{\min} and c_{\max} , 296 and consider three combinations for the grazing parameters (m_Z, a_Z) : (0.05, 1) (weak grazing); (0.015, 1) 297 298 (intermediate grazing); and (0.015, 1.15) (strong grazing). We ran each simulation for 200 years and used the last 150 years for analysis. Figure 8 shows the dominant period of P_1^s for the three sweeps, in order 299 300 of increasing grazing pressure, from left to right. Panel A shows that under weak grazing, the expected 301 relationship between the sweep parameters and the succession time holds: larger values of $c_{\rm max}$ and smaller 302 values of k_{\min} lead to faster succession times. In the other two panels, the regions in white indicate that a 303 dominant period never resolves due to aperiodicity. Most of the parameter space in panel C, where grazing 304 is strongest, has dominant periods that never resolve.

305 5. Discussion and Conclusion

306 Partitioning the observed variability in planktonic ecosystems and community composition between in-307 trinsic and external sources is a difficult task. But a reasonable first step, which we have taken in this study, is to model one kind of variability (intrinsic, in our case) and the time scales on which it occurs. To do 308 309 so, we developed and explored a zero-dimensional ecosystem model that describes a range of possible interactions between phytoplankton, nutrients, and zooplankton. The model supports two kinds of oscillations, 310 which have been previously described in the literature and which we call R-oscillations and Z-oscillations. 311 R-oscillations describe patterns of species succession, while Z-oscillations are predator-prev limit cycles of 312 total biomass. Previous studies have incorporated one type of oscillation or the other into their models, but 313 we are not aware of any that have studied them simultaneously. And there is a dearth of knowledge on the 314time scales of these oscillations in a realistic open ocean system. 315

316 We found that the species succession time for R-oscillations depends on the phytoplankton growth rates at high and low nutrient concentrations. In our model, these depend strongly on two parameters: c_{\max} , which 317 controls the depletion rate of the nutrient in greatest abundance; and k_{\min} , which controls the growth rate 318 of the next phytoplankton species in line to invade and dominate the system. Over the range of parameters 319 320 in (6), we saw R-oscillations in the individual phytoplankton time series that ranged from roughly 100 to 900 days under predation-free conditions. Predator-prev Z-oscillations, on the other hand, consistently occurred 321 322 on time scales not exceeding 60 days. When the two kinds of oscillations interacted, R-oscillations became 323 even slower due to a top-down control on the rate of phytoplankton growth/nutrient depletion. And in the 324 aperiodic scenarios where the dominant period never resolves, distinct regimes emerged over decadal (or 325 longer) time scales.

326 The specific periods of oscillation and frequency of chaotic regime shifts are likely highly dependent 327 on the model parameters and structure. However, the wide range of time scales on which the community 328 structure and biomass varies is consistent with observations from laboratory data. The mesocosm experiment 329 from Benincà et al. (2008) mentioned in our introduction (Section 1) contains multiple time scales of variability 330 in species composition and biomass, and exhibits chaotic tendencies. While the structure of the food web 331 is different from our modeling experiment (it contains only two nutrients but additional compartments for 332 bacteria, detritus and detritivores), the results suggest that non-equilibrium dynamics play a key role in 333 shaping ecosystems. Our model examines this same idea via a particular set of mechanisms (namely R- and $\mathbf{334}$ Z-oscillations).

335 Though our model is suitable for this study in many ways, it also has some clear limitations. First, 336 only six phytoplankton species are included with two feeding strategies represented (small gleaners and large 337 opportunists), whereas there are tens of thousands of species in the real ocean with a continuum of sizes and 338 optimal growth conditions. Second, we assume that any intrinsic variability reflected in the data must be 339 attributed to one or two prescribed mechanisms. By design, we have omitted contributions to variability from 340 stochasticity in the environment, which may affect the evolution of biomass and community composition in 341 a nontrivial way, even in systems with minimal seasonality. We leave these avenues of investigation open for 342 future studies.

From the results above, we determined that intrinsic oscillations can occur on broad range of time scales. 343 344 These oscillations are not tied to any physical phenomena, such as seasonal changes in light or interannual changes in environmental parameters, but they share overlap in their frequency ranges. Thus, we believe that 345346 phytoplankton may exhibit complex responses to many external sources of variability. The degree to which 347 the nature of these responses diverges from the external variability will likely depend on the strength of the forcings and the region of interest. At highly seasonal latitudes, R- and Z-oscillations would be inhibited 348 by winter mixing, which reduces predator-prey encounters and the amount of available light and nutrient. 349350 Phytoplankton evolution may therefore be phase locked to seasonal changes, as reflected in SeaWiFS data from the North Atlantic, for example in (Behrenfeld, 2010). The idea is also reflected in (Winder and Cloern, 3512010), which contrasts a strong annual bloom for an ecosystem in the Wadden Sea (North Atlantic) with one 352353 in Lake Hatchineha (Florida), whose blooms are irregular and occur on shorter time scales.

354 In contrast, coastal upwelling "hotspots" are characterized by a persistent inflow of nutrients from the

355 deep and elevated predatory pressure — conditions that are conducive to domination by intrinsic oscillations. Evidence of intrinsic long-term variability in coastal communities is also reflected in phytoplankton time 356 series. Bode et al. (2015) found from the RADIALES dataset, collected off the northwest coast of the 357 Iberian peninsula, that phytoplankton variability tends to be stochastic. Trends emerge over interannual and 358359 decadal time scales that do not correlate meaningfully with physical properties such as temperature and and stratification. Cianelli et al. (2017) also concluded that the structure of coastal ecosystems in the Gulf of 360 361 Naples is primarily influenced by intrinsic factors, with physical contributions aiding biodiversity by diluting 362 ecosystems and preventing dominance by a handful of species.

Another way to frame the above discussion is to ask how the amplitude of intrinsic variability compares to that of external variability under a variety of physical conditions. The oscillations in this study were quite large for both individual compartments (R-oscillations) and in total biomass (Z-oscillations). We expect that the presence of external oscillations will reduce the amplitude of the intrinsic contributions. But the degree to which this happens will depend on the strength of the relevant forcings. We intend to explore this open question in a follow-up study.

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374 Appendix A. Simplified Model

Here, we present a simplified ecosystem model based on the equations outlined in Section 2. We include two nutrients R_1 and R_2 ; two small phytoplankton P_1 and P_2 ; and one zooplankton Z. We use the same symbols for our parameters, but we remove superscripts and abandon the concept of size classes.

The equations are

$$\frac{dR_1}{dt} = \tau \left(S_1 - R_1\right) - \mu \left[C_{11} \min\left\{\frac{R_1}{K_{11} + R_1}, \frac{R_2}{K_{21} + R_2}\right\} P_1 + C_{12} \min\left\{\frac{R_1}{K_{12} + R_1}, \frac{R_2}{K_{22} + R_2}\right\} P_2\right],\tag{A.1a}$$

$$\frac{dR_2}{dt} = \tau \left(S_2 - R_2\right) - \mu \left[C_{21} \min\left\{\frac{R_1}{K_{11} + R_1}, \frac{R_2}{K_{21} + R_2}\right\} P_1 + C_{22} \min\left\{\frac{R_1}{K_{12} + R_1}, \frac{R_2}{K_{22} + R_2}\right\} P_2\right],\tag{A.1b}$$

$$\frac{dP_1}{dt} = -\tau P_1 + \left[\mu \min\left\{\frac{R_1}{K_{11} + R_1}, \frac{R_2}{K_{21} + R_2}\right\} - m_P - \frac{gZ}{K_Z + P_1 + P_2}\right] P_1,\tag{A.1c}$$

$$\frac{dP_2}{dt} = -\tau P_2 + \left[\mu \min\left\{\frac{R_1}{K_{12} + R_1}, \frac{R_2}{K_{22} + R_2}\right\} - m_P - \frac{gZ}{K_Z + P_1 + P_2}\right] P_2,\tag{A.1d}$$

$$\frac{dZ}{dt} = -\tau Z + \frac{gZ}{K_Z + P_1 + P_2} (P_1 + P_2) - m_Z Z^2$$
(A.1e)

378 There are two nutrient currencies R_1 and R_2 , and we can convert from phytoplankton units (carbon) to 379 nutrient units via multiplication by the stoichiometric coefficients C_{ji} . The stoichiometric coefficients for Z,

- which we may call C_j^Z , do not appear in the equations because zooplankton do not introduce a loss term to the nutrient equations.
- **382** We can examine the evolution of ϕ_j , or the total amount of R_j in the system

$$\phi_1 = R_1 + C_{11}P_1 + C_{12}P_2 + C_1^Z Z \tag{A.2a}$$

$$\phi_2 = R_2 + C_{21}P_1 + C_{22}P_2 + C_2^Z Z \tag{A.2b}$$

383 If we differentiate ϕ_j in time, we have a balance between: nutrient inflow, outflow from all compartments, 384 and phytoplankton/zooplankton mortality

$$\frac{d\phi_1}{dt} = \tau(S_1 - \phi_1) - m_P(C_{11}P_1 + C_{12}P_2) - C_1^Z m_Z^2$$
(A.3a)

$$\frac{d\phi_2}{dt} = \tau (S_2 - \phi_2) - m_P (C_{21}P_1 + C_{22}P_2) - C_2^Z m_Z^2$$
(A.3b)

385 Note that these equations do not depend on growth or uptake terms, as these processes are completely386 balanced and cancel upon adding up the right-hand sides of the equations.

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