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Sylvain David, D. Chateau, H.-J. Chang, L.H. Karlsson, M.V. Bondar, et al.. High-Performance Optical Power Limiting Filters at Telecommunication Wavelengths: When Aza-BODIPY Dyes Bond to Sol-Gel Materials. *Journal of Physical Chemistry C*, 2020, 124 (44), pp.24344-24350. 10.1021/acs.jpcc.0c08006 . hal-03103459

HAL Id: hal-03103459

<https://hal.science/hal-03103459v1>

Submitted on 26 May 2023

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High performance optical power limiting filters at telecommunication wavelengths: when aza-BODIPY dyes bond to sol-gel materials.

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KEYWORDS: *Optical limiting; nonlinear-optics; Aza-BODIPY dyes; two-photon absorption; sol-gel material.*

ABSTRACT: An optical power limiter (OPL) device can change its optical properties depending on the input light intensity: it is ideally transparent at low incident power and presents a decrease in transmission when the incident intensity increases. This phenomenon can be used for protection of optical sensors. This paper presents the design of a functionalized aza-BODIPY chromophore, combining good nonlinear two-photon absorption and excited state absorption in the short wave infra-red range (1300-1600 nm). The aza-BODIPY chromophores also show excellent thermal and photochemical stabilities. Their covalent grafting to a sol-gel matrix lead to the formation of class-II sol-gel materials that allows for extraordinary doping concentrations of up to 280 mol.L⁻¹ (40 wt%), which results in highly efficient OPL performances at 1550 nm.

An optical power limiting (OPL) filter has the capability of modifying its global transmittance depending on the input light intensity.¹⁻³ Nowadays, the tremendous development of eye-safe LIDAR (Light Detection And Ranging) systems working at telecommunication wavelengths (1550 nm), for applications in aeronautics (telemetry), but also in defense (active imaging) and civilian transportation (autonomous vehicles guiding systems) triggered a rapidly growing demand for efficient OPL devices at this wavelength, but more generally in the SWIR (Short Wave Infra-Red) band. For field applications, the OPL filter should fulfill a set of requirements: (i) a broad spectral efficiency in the spectral range-of-interest (ROI = 1300-1600 nm), (ii) a high transmission in the off-state at low incident light intensity (minimum 85%), (iii) a low transmission in the on-state at high incident light intensity (below the damage threshold of the optical detector), (iv) a very fast response time and (v) a good

environmental stability in order to be adapted to existing optical devices.³

As a general point of view, the OPL process can be driven by nonlinear optical phenomena,³⁻⁴ such as nonlinear refraction,⁵⁻⁷ nonlinear scattering,⁸⁻⁹ reverse saturable absorption¹⁰⁻¹³ and, two- (multi)-photon absorption (2PA). Consequently, the process is triggered by the incident light itself, which induces an instantaneous nonlinear response within the material (sub-fs timescale). This self-activated instantaneous response is a key advantage for practical applications. It has been widely demonstrated that OPL in the visible spectral range can be strongly enhanced if an excited-state absorption (ESA) phenomenon overlaps temporally and spectrally with the 2PA.^{3, 10, 14-18} However, in spite of the rapid development since 2005 of two-photon dyes absorbing in the SWIR spectral range,¹⁹⁻²⁴ OPL at 1300-1600 nm is a field of research still in its infancy. Therefore, few reports describing OPL

experiments at telecommunication wavelengths have appeared; so far, polymethine dyes,²⁵⁻²⁶ functionalized porphyrins²⁷ and aza-BODIPY dyes in a solvent solution have been presented.²⁸⁻²⁹ In all cases, a 2PA induced ESA photophysical process is responsible for the OPL mechanism.³⁰ However, if solid-state materials operating in the visible wavelengths are well-known,³¹ up to now, only one solid-state material for OPL in the SWIR has been described in the literature.³² This material, reported by our group is a sol-gel glass doped with the functionalized aza-BODIPY chromophore **1** (**Erreur ! Source du renvoi introuvable.**).³² Despite the encouraging results obtained with this first material, the OPL response remained far from the technical requirements for practical use in real optical systems. In particular, the dye loading, which is one of the key parameters for improving the OPL properties was limited to 8 wt% (approx. 85 mM), consequently limiting the performances. This limitation is explained by the poor compatibility between the dyes and the matrix in such host-guest approach, where the matrix-dye interface strength is weak. Above this concentration the gel, during its contraction, expelled the dye. An aggregation of the dye inside the matrix is also observed which causes a drop in the linear transmission.

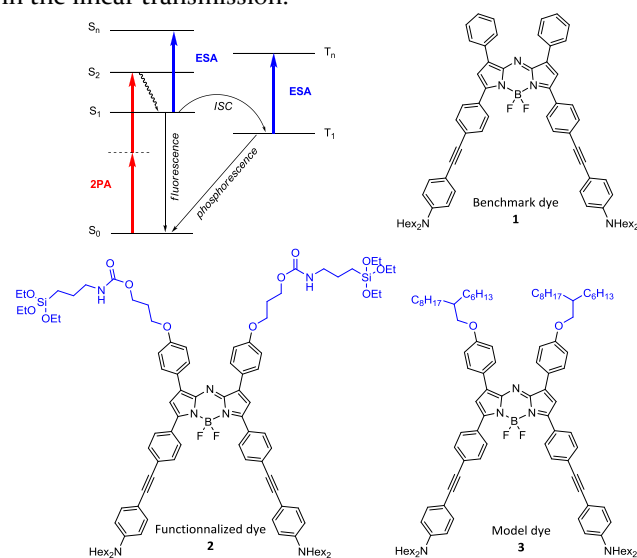


Figure 1. Perrin-Jablonski diagram with the different photophysical processes involved in a molecular based OPL. Structure of the benchmark dye **1** and the functionalized and model chromophores **2** and **3**.

In order to prevent such phase segregation during the process, an alternative strategy is known, increasing the interface strength through covalent bonding between the hybrid matrix and the molecular host, here the dye.³³ In contrast to weak interaction systems, named Class-I sol-gel materials, where the guest dye is only dispersed in the matrix, Class-II refers to materials where the dye is covalently grafted to the host matrix. Such alternative has been previously investigated for OPL hybrid materials in the visible range. It brought improvements but in most cases with still limited optical performances.^{15, 34-37} These results led us to turn to the so-called class-II sol-gel hybrid materials with covalent interactions between matrix and molecular host. This process usually permits

to achieve high doping concentrations without separation of the dye.¹⁵ However, in some cases, the rigidifying of the dye structure induced by the grafting on the matrix can be counter-productive, as the solid-state constraints disrupts the conjugated skeleton hence reducing the nonlinear optical properties.³⁸ Class-II materials for OPL applications in the visible range have been reported using platinum acetylide complexes as the nonlinear doping chromophore. A high concentration was achieved (*i.e.* 120 mM) without any perturbation of the optical properties of the dye.¹⁵ This procedure was therefore extended to aza-BODIPY dye. **1** was functionalized with triethoxysilane moieties, through the alkoxy-functionalization of upper 1,7 positions (molecule **2**, **Erreur ! Source du renvoi introuvable.**). The effect of such alkoxy-donating group functionalization on the photophysical properties of the dye was thoroughly investigated on the model chromophore **3** and compared to **1** both theoretically and experimentally. Finally, class-II materials with improved doping ratio were prepared for OPL properties studies.

The synthesis of compounds **2** and **3** follows the classical route for symmetrical aza-BODIPY dyes, as already described for the benchmark compound **1** (Schemes S₁ and S₂).³⁸ In the case of **2**, an intermediate aza-BODIPY featuring hydroxyl-propanol (-O-(CH₂)₃-OH) group at the top positions is prepared and the tri-ethoxysilane reactive function was introduced at the last step by reaction with 3-(triethoxysilyl)propyl isocyanate in the presence of dibutyltindilaurate as catalyst. Synthetic details and complete characterization are given in the Supporting Information (SI).

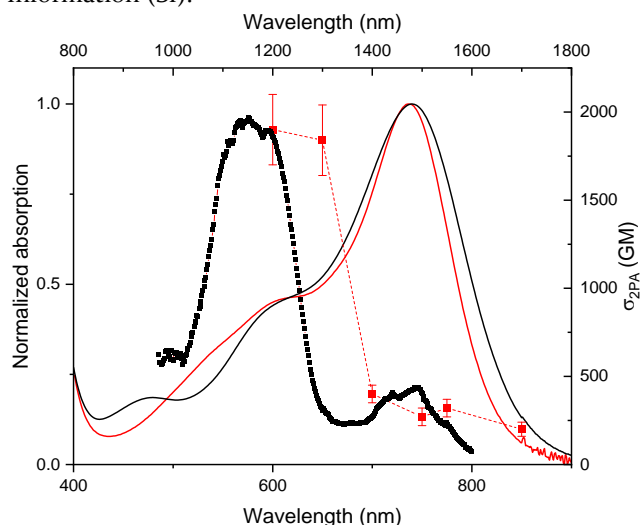


Figure 2. Comparison of the linear and nonlinear two-photon absorption of **1** (black) and **3** (red) in DCM solution. Two-photon absorption spectra of **1** and **3** were measured by two-photon excited fluorescence in CCl₄ (▪) and open-aperture Z-scan measurements in DCM (▪), respectively.

Due to the moisture sensitivity of the triethoxysilane group in molecule **2**, all photophysical measurements were performed in DCM on the model compound **3** which has simple alkoxy substituents (**Erreur ! Source du renvoi introuvable.** and Table 1). The measurements, performed in dichloromethane (DCM) solution, were compared with the reference dye **1**. Both benchmark dye **1** and molecule **3** show a very similar linear absorption spectrum with

respective absorption maxima in DCM at 740 and 736 nm and a slight narrowing of the lower energy transition for **3**. The second transition in the 500-600 nm region is slightly modified as well. If the shoulder at ca. 600 nm, attributed to a transition having mainly Charge Transfer (CT) character in the benchmark chromophore is identical for the dye **3**, it seems that a second band around 525 nm superimposes with the latter band, causing the broadening of the spectrum. All these minor modifications in linear absorption indicate that upper alkoxy donating groups have weak impact on the electronic transitions. In addition, both compounds are not emissive in DCM²⁹ though the emission can be restored in apolar solvents such as toluene (Figure S1). The emission profiles are almost identical but dye **3** exhibits an improved quantum yield and a longer lifetime (Table 1).

Table 1. Linear and nonlinear optical properties of **1** and **3** in dichloromethane. λ_{\max} and λ_{em} are the maximal absorption and emission wavelength, ϵ the extinction coefficient, ϕ and τ the radiative quantum yield and lifetime, $\sigma_{2\text{PA}}$ and σ_{ESA} the 2PA and ESA cross-sections at 1550 nm.

Molecule	1	3
λ_{\max} [nm]	740	736
ϵ [$\text{M}^{-1}\cdot\text{cm}^{-1}$]	57000	65000
$\lambda_{\text{em}}^{\text{a}}$ [nm]	826	804
$\phi^{\text{a,b}}$	0.04	0.11
τ^{a} [ns]	0.7	1.4
$\sigma_{2\text{PA}@1550\text{ nm}}$ [GM]	236 ^c	320 ± 50 ^c
$\sigma_{\text{ESA}@1550\text{ nm}}$ [10^{-16} cm^2]	0.84 ^d	1.75 ^e

a) in toluene; b) using aza-BODIPY dyes as a reference³⁹ ($\phi = 0.36$ in CHCl_3); c) using TPEF in CCl_4 ; ²⁹ d) using pump-probe experiment in DCM; ²⁹ e) using Z-scan measurements in DCM.

To obtain more insights into the excited states involved in electronic transitions of these dyes, we performed first principle calculations, using an approach validated previously for a similar class of dyes^{29, 40} (see the SI for details). To make these calculations feasible, the long alkyl chains of the dihexylamino and decylhexylalcoxy fragments have been replaced by methyl groups (simplified dyes **1**th and **3**th, respectively depicted Figure S3) and the main results are displayed in Figure 3 and Table 2. The lowest energy transition S_0-S_1 , was calculated at 611 and 610 nm respectively for **1**th and **3**th. The fact that the computed absorption wavelength is significantly smaller than the experimental value is the logical consequence of the neglect of vibronic couplings and the cyanine nature of the transition where the electronic density changes are mainly localized on the aza-BODIPY core, the peripheral groups at the top or the bottom of the molecule playing a minor role for both systems. As previously observed for **1**,²⁹ the excited state representations (Figure 3) show that the lowest transition is mainly localized on the aza-BODIPY core. There is

nevertheless a small CT from the ethynyl groups towards the core.²⁹ Moderate values of transferred charge (q_{CT} 0.48 and 0.44 e) over a distance (d_{CT}) of 2.26 and 1.63 Å, for **1**th and **3**th respectively are in good agreement with this assumption. As a consequence, the calculated 2PA cross section are modest. In contrast, the three higher energy transitions S_0-S_i ($i = 2, 3$, and 4) present stronger CT character from the bottom donating group (for **1**th and **3**th) to the aza-BODIPY core, with an additional short-range contribution of the top substituent for **3**th in the case of S_0-S_2 . As expected, all these transitions present larger q_{CT} values spread over a longer d_{CT} distance compare to S_0-S_1 (Table 2) and do contribute more efficiently to the 2PA cross-section.

Table 2. Theoretical (vertical) absorption wavelength (λ), one-photon absorption oscillator strength (f), CT distance and charge (d_{CT} and q_{CT}), as well as 2PA cross-section ($\sigma_{2\text{PA}}$) for the four lowest-lying states of the simplified dyes **1**th and **3**th.

dyes	S_0-S_i	λ [nm]	f (1PA)	d_{CT} [Å]	q_{CT} [e]	$\sigma_{2\text{PA}}$ [GM]
1 th	S_1	611	1.02	2.26	0.48	175
	S_2	431	0.99	4.37	1.08	3760
	S_3	417	0.01	0.60	0.68	1160
	S_4	394	0.00	3.90	0.88	1340
3 th	S_1	610	1.06	1.63	0.44	95
	S_2	447	0.02	1.41	0.69	1030
	S_3	430	1.29	1.96	0.79	1850
	S_4	403	0.03	2.58	0.78	3410

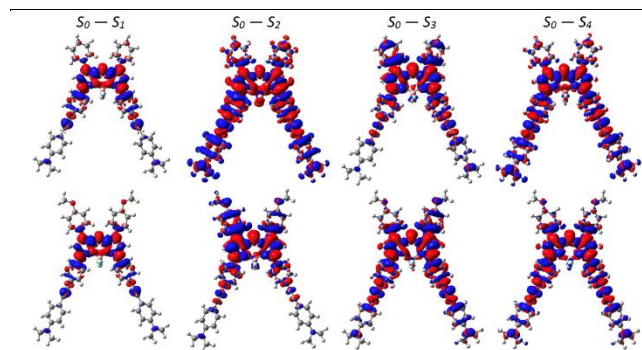


Figure 3. Electron density difference plots for the four lowest-lying excited states of **1**th (top) and **3**th (bottom). The red (blue) zones represent regions of increase (decrease) of the density upon electronic transition. A contour threshold of 8×10^{-4} au is used.

The 2PA properties of **3** were determined by open aperture Z-scan in DCM solutions using a fs-optical parametric amplifier pumped by a Ti:Sapphire source with a typical pulse duration of 150 fs (Table 1). The experimental nonlinear transmittance curve obtained (Figure S2) was fitted using the complete propagation equation (eq 1) taking into account both 2PA and ESA phenomena,

$$\frac{dI}{dz} = -\alpha_2 I^2 - \sigma_{\text{ESA}} N_e I \quad (1)$$

where I represents the light irradiance, z the propagation axis, α_2 the 2PA coefficient, σ_{ESA} the ESA cross section and N_e describes the first excited-state population after 2PA absorption. Using this model, both 2PA and ESA cross-sections could be determined simultaneously for a given wavelength using the following equation (2)

$$\sigma_{2\text{PA}} = \frac{\hbar\omega\alpha_2}{N} \quad (2)$$

where N represents the molecular density and $\hbar\omega$ the photon energy at the operating wavelength. Typical Z-scans at 1550 nm for **3** and related fitting curves are depicted in Figure S2 leading to a 2PA cross section of 320 ± 50 GM and an ESA cross section of $1.75 \pm 0.3 \cdot 10^{-16}$ cm² (see Table S1 for other wavelengths). These values are comparable with that obtained for the benchmark molecule **1** ($\sigma_{2\text{PA}} = 236$ GM and $\sigma_{\text{ESA}} = 0.84 \cdot 10^{-16}$ cm² in Table 1) keeping in mind that these latter data have been obtained using different and independent techniques, *i.e.* two-photon excited fluorescence and NIR pump-probe techniques for 2PA and ESA measurements, respectively.²⁹ Considering experimental uncertainties for both techniques, it appears that molecules **1** and **3** present comparable 2PA and ESA in the 1300-1600 nm range (Figure 2).

For both systems, the weak 2PA band overlaps with the intense linear “cyanine” type transition (Figure 2) in agreement with the calculated $\sigma_{2\text{PA}}$ of 175 and 95 GM respectively for simplified molecules **1'** and **3'** (Table 2); the origin of nonlinear properties can be ascribed to the weak CT character of this absorption transition, as mentioned above. In addition, an intense 2PA band around 2000 GM, has been measured at higher energies matching with the shoulder of the linear absorption spectra, (Figure 2) corresponding to the overlap of the S_0-S_i ($i = 2, 3,$ and 4) with stronger CT character. These experimental data are further corroborated by the high $\sigma_{2\text{PA}}$ calculated values (Table 2) (3760 and 3410 GM) for S_0-S_2 and S_0-S_4 transitions of **1'** and **3'** respectively. It must be noted that, as previously observed,²⁹ calculations show for molecules **1**th and **3**th the existence of an intense 2PA transition corresponding to a forbidden 1PA transition (S_0-S_3 and S_0-S_4 for **1**th and S_0-S_2 and S_0-S_4 for **3**th, respectively). All these close energy transitions may contribute to the overall 2PA spectrum.

Finally, the thermal and photochemical stability of dye **3** was evaluated before incorporation into the matrix. Thermogravimetric analysis show a Td¹⁰ value of 375°C (Figure S4). This value is very high for an organic dye⁴¹ and fully compatible with hybrid materials synthesis and practical application. On the other hand; the photochemical stability of **3** was evaluated in air-saturated DCM with the absorption method described previously in

detail.⁴² The photodecomposition quantum yield, ϕ_{ph} , of solutions of **3** at ≈ 635 nm, performed with a continuous wave (CW) diode laser with an average irradiance ≈ 500 mW/cm², was determined to 1.10^{-8} . This level of photostability is very high, comparable with the best laser dyes.⁴³⁻⁴⁵

Having checked that the model chromophore **3** fulfills all the molecular prerequisites for being used as the active dye in OPL materials, the synthesis of hybrid sol-gel materials containing molecule **2** featuring reactive triethoxysilane moieties was undertaken. The general synthesis protocol consisted in mixing compound **2** with a sol composed of a 1:1 mixture of VTMOs (vinyltrimethoxysilane) and MTEOS (methyltriethoxysilane). Modifications have been made compared to the previously published protocol³² to account for class-II materials specificities. The overall hydrolysis-condensation kinetics were controlled by adding gelation agents like *N,N*-Dimethylaminopropyl-trimethoxysilane (DMATMS) and tetrabutylammonium fluoride (TBAF) in catalytic amounts. The catalyst amount was adapted to the quantity of dyes added to the sol-gel matrix. The condensation ratio of the prepared matrix is extremely high, leaving very few available silanols for the condensation of the added dyes with the silica network. This favors the dyes self-condensation and further aggregation and optical diffusion in the solid material. In order to provide available silanol groups for the reaction with the dyes, a low amount of water and fresh VTMOs monomer were added.

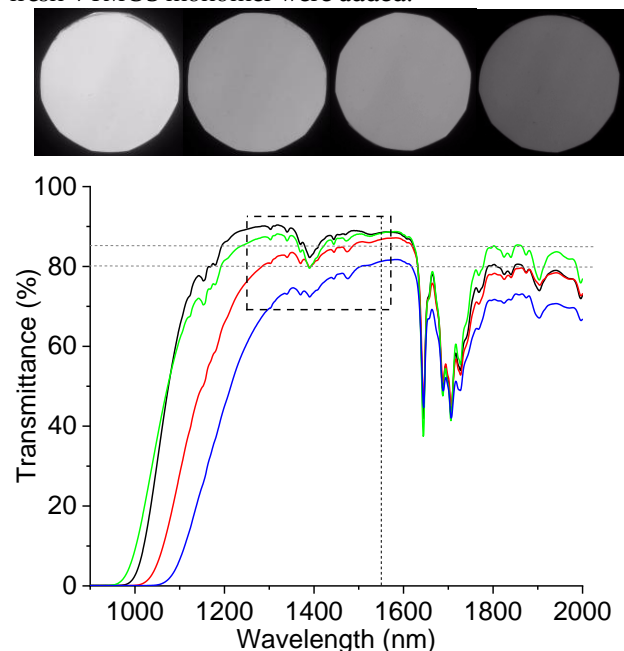


Figure 4. (Top) NIR picture of xerogel **MRef**, **M10**, **M20**, **M40**, after polishing (from left to right). The increased darkening is due to red-shift of the cut-off absorption wavelength of the dye with incremental concentration and also to the lower sensitivity of the camera over 1200 nm. All four images were taken with similar exposure. (bottom) Transmittance of materials **MRef** (black), **M10** (green), **M20** (red) and **M40** (blue) after polishing. Dotted box represents the ROI, horizontal dotted lines represent the transmittance at 80 and 85% and vertical dotted line is 1550 nm wavelength used for OPL experiments.

With this strategy, no dye aggregation or phase separation is observed during the preparation. The xerogel was readily obtained after drying at 45°C for 2 days as shiny dark disks (1.5 cm diameter, 2 mm thickness). The optical qualities of these materials were qualitatively estimated by a home-made device able to evaluate qualitatively their SWIR transparency and homogeneity (see the SI, Figure S5, S6 and Figure 4). The presence of TBAF was mandatory for obtaining xerogels with good optical quality (Figure S6).

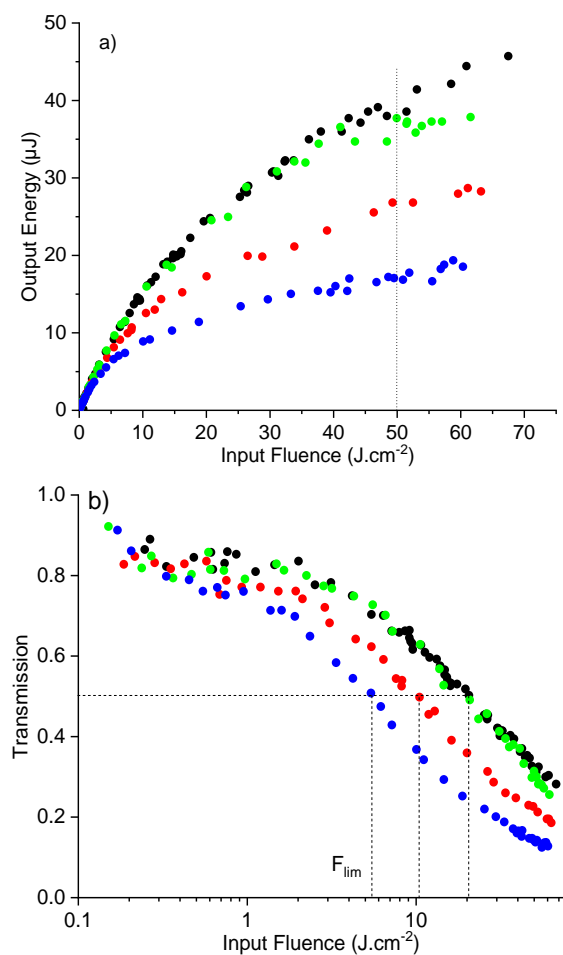


Figure 5. Optical limiting performances at 1550 nm for materials **M10** (●), **M20** (●), **M40** (●). For comparison, the reference doped material **MRef** (●) was measured under the same conditions. Two representations are proposed: (top) Out-put energy vs incident fluence and (bottom) transmittance vs incident fluence (in log scale). The F_{lim} thresholds are indicated by dotted lines.

This optimized protocol was used for the preparation of a series of xerogels (**M10**, **M20**, **M40**) containing increased concentration of aza-BODIPY dye **2** (10, 20 and 40 wt% corresponding to 70, 140 and 280 mol.L⁻¹, respectively, Table S1). As a reference, we used the already described **MRef** material,³² with a 5 wt% doping concentration of dye **1** (approx. 56 mM). The final materials were mechanically polished, yielding 1 mm thick shiny black disks of optical quality and perfectly parallel faces.

Interestingly, all studied materials are very homogeneous without any apparent defect (dark spots, cracks...Figure 4). In addition, transmittance spectra were measured and compared with reference host-guest material **MRef** (Erreur ! Source du renvoi introuvable.. Material **M10** has a transmittance very close to that of **MRef**. The increase in concentration caused a red-shift in cutoff wavelength on materials **M20** and **M40** corresponding to the red tail of the chromophore absorption. The linear transmission, in the 1600-1800 nm region, decreased due to matrix absorption. However, in the ROI (1300-1600 nm) and more particularly at 1550 nm, the wavelength of interest, all materials have an excellent transmission ranging from 81% to 88%, which matches the 85 % requirement.

The optical limiting experiments were performed at 1550 nm using 5 ns laser pulses and a f/5 system (see SI for description) and the experimental results are reported in Figure 5. Material **M10** shows a very similar OPL behavior as **MRef**, which agrees with the comparable dye concentration in both materials. As expected the OPL performance improves with an increase in doping concentration. In order to compare the materials OPL performances further, the limiting threshold, F_{lim} , defined as the fluence at which the transmittance reaches 50% is used. This value is 21 J.cm⁻² for both materials **MRef** and **M10**, 10 J.cm⁻² for **M20** and 5 J.cm⁻² for **M40**. This indicates a 4-fold increase of the material performances compared to the reference. We also observe a two-fold decrease in F_{lim} with a two-fold increase in doping concentration. The linear stepwise increase of the OPL performances of these sol-gel materials therefore indicates that the nonlinear properties of the dye is not affected by an increase in concentration nor by the covalent bonding to the matrix. The materials were thoroughly inspected after the OPL experiments, under visible and IR light, and no trace of laser induced damages was observed for the applied experimental conditions. This is in accordance with the observed high photochemical stability of the materials.

In conclusion, this article describes the potential use of class-II sol-gel materials, containing substituted aza-BODIPY dyes, for OPL applications at telecommunication wavelengths. The presented preparation method allows for reaching high chromophore concentrations and avoiding aggregation issues, thanks to the covalent grafting of the chromophore to the matrix. More importantly is the fact that the nonlinear properties are preserved. The resulting xerogels are homogeneous with high linear transparency in the SWIR spectral range. Their remarkable mechanical properties allowed preparation of prototype with high optical quality after polishing. Finally, the materials showed a significant improvement of the OPL performances with the concentration increase. This is accompanied with a damage resistance to high fluence laser irradiation. The presented results, therefore, constitute a significant step towards high performance optical limiters for applications in SWIR.

ASSOCIATED CONTENT

Supporting Information. Synthetic procedure and chemical characterisation of dyes **2** and **3**, sol gel material preparation protocols, complete description of the spectroscopy, imaging and OPL set-ups are described in the Supporting Information. This material is available free of charge via the Internet at <http://pubs.acs.org>.”

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£ both authors contribute equally to the work. The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

ACKNOWLEDGMENT

Authors acknowledge Drs. Eric Van Stryland and David Hagan and the CREOL Nonlinear Optics Laboratories for their collaboration and assistance for the determination of the two-photon cross-section. Authors thank Thales LAS for financial support and for the grant of SD. The Swedish Defence Research Agency (FOI) is also thanked for financial support. D.J. is indebted to the CCIPL computational center installed in Nantes for (very) generous allocation of cpu time. Hampus Lundén is acknowledged for his contribution regarding the OPL set-up.

ABBREVIATIONS

SWIR, short wavelength infra-red; OPL, optical power limiting; DCM, dichloromethane.

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