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G. Griso: **griso@ljll.math.upmc.fr**
Sorbonne Université, CNRS, Université de Paris, Laboratoire Jacques-Louis Lions (LJLL), F-75005 Paris, France,
L. Khilkova: **larysa.khilkova@itwm.fraunhofer.de**, J. Orlik: **julia.orlik@itwm.fraunhofer.de**,
O. Sivak: **olena.sivak@itwm.fraunhofer.de**,
Department SMS, Fraunhofer ITWM, 1 Fraunhofer Platz, 67663 Kaiserslautern, Germany

Abstract

In this paper, we study the asymptotic behavior of an ε-periodic 3D stable structure made of beams of circular cross-section of radius r when the periodicity parameter ε and the ratio r/ε simultaneously tend to 0. The analysis is performed within the frame of linear elasticity theory and it is based on the known decomposition of the beam displacements into a beam centerline displacement, a small rotation of the cross-sections and a warping (the deformation of the cross-sections). This decomposition allows to obtain Korn type inequalities. We introduce two unfolding operators, one for the homogenization of the set of beam centerlines and another for the dimension reduction of the beams. The limit homogenized problem is still a linear elastic, second order PDE.

Keywords: linear elasticity, homogenization, stable structure, periodic beam structure, periodic unfolding method, dimension reduction, Korn inequalities.
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1 Introduction

The aim of this work is to study the asymptotic behavior of an ε-periodic 3D stable structure made of "thin" beams of circular cross-section of radius r when the periodicity parameter ε tends to 0, in the framework of the linear elasticity. By "thin", we mean that the radius r of the beams is much smaller than the periodicity parameter ε and that we deal with the case where ε and r/ε simultaneously tend to 0.

It is well known to engineers that for wire trusses, lattices made of very thin beams, bending dominates the stretching-compression. A contrario, if the same structures are made of thick beams the stretching-compression dominates. This is what several mathematical studies of recent decades have obtained for periodic structures made of beams. For such structures, from the mathematical point of view, this means that the processes of homogenization and dimension reduction do not commute (see the pioneer works [5, 11, 12] and also [1, 6, 8, 24, 25, 27, 28, 31]). Our aim is to investigate between these extreme cases. More precisely, we consider the case for which the ratios $\mathrm{diam}(\Omega)/\varepsilon$ and ε/r are of the same order (Ω is the $3D$ domain covered by the beam structure). In Sections 5 and following, we show that the ratio r/ε^2 and its limit $\kappa \in [0, +\infty]$ play an important role in the estimates and the asymptotic behaviors. It worth to notice that in our analysis, $\kappa = 0$ also corresponds to the case where first the dimension reduction is done and then the homogenization, while $\kappa = +\infty$ is for the vice-versa case. In the convergences (7.12) of Theorem 7.2, we show that the rescaled global displacement depends on κ. If $\kappa \in (0, +\infty)$, its limit is a combination of a global displacement (a pure stretching-compression) and a local bending; if $\kappa = +\infty$ it is just a global displacement and if $\kappa = 0$ it is a local bending.

Our analysis relies on a displacement decomposition for a single beam introduced in [13, 14, 15]. According to those studies, a beam displacement is the sum of an elementary displacement and a warping. The elementary displacement has two components. The first one is the displacement of the beam centerline while the second stands for the small rotation of the beam cross-sections (see [13, 15]). This decomposition has been extended for structures made of a large number of beams in [14] (see [4] for the structures made of beams in the nonlinear elasticity framework). Here, similar displacement decompositions are obtained, these decompositions are used for stable beam structures (see Lemma 4.2) and then for periodic 3D stable structures made of beams. It is important

to note that estimate $(4.5)_1$ is the key point of this paper. It characterizes the stable structures. In a forthcoming paper, we will investigate the unstable and auxetic 3D periodic structures made of beams and we will see that all the estimates of Lemma 4.2 will remain except $(4.5)_1$. These decompositions allow to obtain Korn type inequalities as well as relevant estimates of the centerline displacements.

To study the asymptotic behavior of periodic stable structures and derive limit problem we use the periodic unfolding method introduced in [9] and then developed in [10]. This method has been applied to a large number of different types of problems. We mention only a few of them which deal with periodic structures in the framework of the linear elasticity (see [3, 16, 17, 18, 19, 20, 21, 26]). As general references on the theory of beams or structures made of beams, we refer to [2, 7, 22, 23, 29, 30].

The paper is organized as follows. Section 2 introduces structures made of segments and remind properties of Sobolev spaces defined on these structures. Furthermore, in this section we give a simple definition of stable and unstable structures and present several examples. In Section 3 we remind known results concerning the decomposition of a beam displacement into an elementary displacement and a warping. This section also gives estimates with respect to the L^2-norm of the strain tensor of the terms appearing in the decomposition. In Section 4 we extend the results of the previous section to structures made of beams. Complete estimates of our decomposition terms and Korn-type inequalities are obtained for stable structures.

In Section 5 we deal with an ε-periodic stable 3D structure made of r-thin beams, $\mathcal{S}_{\varepsilon,r}$. For this structure we introduce a linearized elasticity problem and specify the assumptions on the applied forces. Using results from the previous section we decompose every displacement of $\mathcal{S}_{\varepsilon,r}$ as the sum of an elementary displacement and a warping and provide estimates of the terms of this decomposition. The scaling of the applied forces are given with respect to ε and r. That leads to an upper bound for the L^2-norm of the strain tensor of the solution of the elasticity problem of order 1.

In Section 6 we introduce different types of unfolding operators, mainly one for the centerline beams and another for the cross-sections. This last operator concerns the dimension reduction. Several results on these operators are given in this section and Appendix C.

Section 7, deals with the asymptotic behavior of a sequence of displacements and their strain tensors. Then, in Section 8, in order to obtain the limit unfolded problem we split it into three problems: the first involving the limit warpings (these fields are concentrated in the cross-sections, this step corresponds mainly to the process of dimension reduction), the second involving the local extensional and inextensional limit displacements posed on the skeleton structure and the third involving the macroscopic limit displacement posed in the homogeneous domain Ω.

In Section 9 we complete this analysis by giving the homogenized limit problem (Theorem 9.1). We obtain a linear elasticity problem with constant coefficients calculated using the correctors.

In Section 10 we apply the previously obtained results in the case where the periodic 3D beam structure is made of isotropic and homogeneous material. We present an approximation to the solution of the linearized elasticity problem which can be explicitly computed using the solution of the homogenized problem.

In the Appendix we give the most technical results.

2 Geometric setting

2.1 Structures made of segments

In this paper we consider structures made up of a large number of segments.

Definition 2.1. *Let* $\mathcal{S} = \bigcup_{\ell=1}^{m} \gamma_\ell$, $\gamma_\ell \doteq [\mathbf{A}^\ell, \mathbf{B}^\ell]$, *be a set of segments and* $\mathcal{K}$ *the set of the extremities of these segments.*

$\mathcal{S}$ *is a* ***structure*** *if*

- $\mathcal{S}$ *is nonincluded in a plane,*
- $\mathcal{S}$ *is connected,*
- *a common point to two segments is a common extremity of these segments,*
- *if an element of* $\mathcal{K}$ *belongs to only two segments then the directions of these segments are noncollinear,*
- *for every segment* γ_ℓ *we denote* $\mathbf{t}_1^\ell$ *a unit vector in the direction of* γ_ℓ, $\ell \in \{1, \dots, m\}$.

We denote $\mathbf{t}_1$ the field belonging to $L^\infty(\mathcal{S})^3$ defined by

$$\mathbf{t}_1 = \mathbf{t}_1^\ell \quad \text{a.e. in } \gamma_\ell, \ \ \ell \in \{1, \dots, m\}.$$

The segment $\gamma_\ell \subset \mathcal{S}$ of length l_ℓ is parameterized by $S_1 \in [0, l_\ell]$, $\ell \in \{1, \dots, m\}$

$$\gamma_\ell \doteq [\mathbf{A}^\ell, \mathbf{B}^\ell] \doteq \big\{\mathbf{A}^\ell + S_1 \mathbf{t}_1^\ell \in \mathbb{R}^3 \mid \ S_1 \in [0, l_\ell]\big\}, \quad (\mathbf{A}^\ell, \mathbf{B}^\ell) \in \mathcal{K}^2.$$

The running point of $\mathcal{S}$ is denoted $\mathbf{S}$. For all $\mathbf{S} \in \gamma_\ell$ one has $\mathbf{S} = \mathbf{A}^\ell + S_1 \mathbf{t}_1^\ell$, $S_1 \in [0, l_\ell]$, $\ell \in \{1, \dots, m\}$.

2.2 Some reminders on the Sobolev spaces $L^p(\mathcal{S})$ and $H^1(\mathcal{S})$

A measurable function Φ defined on $\mathcal{S}$ belongs to $L^p(\mathcal{S})$, $p \in [1, +\infty]$, if for every segment $\gamma_\ell \subset \mathcal{S}$, one has $\Phi_{|\gamma_\ell} \in L^p(\gamma_\ell)$, $\ell \in \{1, \dots, m\}$.
For every $\Phi \in L^1(\mathcal{S})$ define

$$\int_{\mathcal{S}} \Phi(\mathbf{S}) d\mathbf{S} \doteq \sum_{\ell=1}^{m} \int_0^{l_\ell} \Phi(\mathbf{A}^\ell + S_1 \mathbf{t}_1^\ell) dS_1.$$

Observe that the right-hand side of the above equality does not depend on the choice of a unit vector in the directions of the segments. The space $L^2(\mathcal{S})$ is endowed with the norm

$$\|\psi\|_{L^2(\mathcal{S})} \doteq \sqrt{\int_{\mathcal{S}} |\psi(\mathbf{S})|^2 \, d\mathbf{S}}, \qquad \forall\, \psi \in L^2(\mathcal{S}).$$

Set

$$H^1(\mathcal{S}) \doteq \Big\{\psi \in C(\mathcal{S}) \mid \psi_{|\gamma_\ell} \in H^1(\gamma_\ell), \quad \ell \in \{1, \dots, m\}\Big\},$$

where $C(\mathcal{S})$ is the set of continuous functions on $\mathcal{S}$.
For every $\phi \in H^1(\mathcal{S})$ denote

$$\frac{d\phi}{d\mathbf{S}}(\mathbf{S}) \doteq \frac{d\phi}{dS_1}(\mathbf{A}^\ell + S_1 \mathbf{t}_1^\ell) \qquad \text{for a.e. } \mathbf{S} = \mathbf{A}^\ell + S_1 \mathbf{t}_1^\ell, \quad S_1 \in (0, l_\ell), \ \ \ell \in \{1, \dots, m\}. \tag{2.1}$$

We endow $H^1(\mathcal{S})$ with the norm

$$\|\psi\|_{H^1(\mathcal{S})} \doteq \sqrt{\|\psi\|^2_{L^2(\mathcal{S})} + \Big\|\frac{d\psi}{d\mathbf{S}}\Big\|^2_{L^2(\mathcal{S})}}, \qquad \forall\, \psi \in H^1(\mathcal{S}).$$

2.3 Stable structures

The space of all rigid displacements is denoted by $\mathbf{R}$

$$\mathbf{R} \doteq \Big\{\mathbf{r} \in C^1(\mathbb{R}^3) \mid \mathbf{r}(x) = \mathbf{a} + \mathbf{b} \wedge x, \ \ \forall x \in \mathbb{R}^3, \ \ (\mathbf{a}, \mathbf{b}) \in \mathbb{R}^3 \times \mathbb{R}^3\Big\}.$$

We define the space $\mathbf{U}_\mathcal{S}$ as follows:

$$\mathbf{U}_\mathcal{S} \doteq \Big\{U \in C(\mathcal{S})^3 \mid \ \text{for every segment } \gamma_\ell \subset \mathcal{S}, \ U_{|\gamma_\ell} \text{ is an affine function}, \ \ell \in \{1, \dots, m\}\Big\}.$$

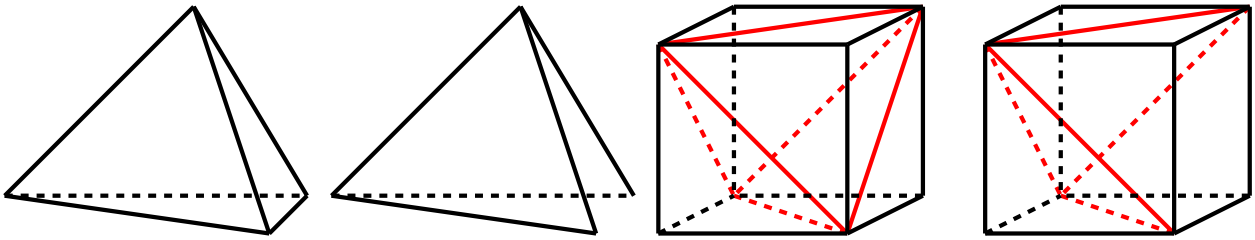
Definition 2.2. *A structure $\mathcal{S}$ is a **stable** structure if*

$$\forall\, U \in \mathbf{U}_\mathcal{S}, \quad \frac{dU}{d\mathbf{S}} \cdot \mathbf{t}_1 = 0 \implies U \in \mathbf{R}.$$

*If the above condition is not satisfied, $\mathcal{S}$ is an **unstable** structure.*

Remark 2.1.

1. *The structure made of the edges of a tetrahedron is stable (see Fig.1.a). If we remove one edge then the structure becomes unstable (see Fig.1.b).*
2. *The structure made of 12 edges and 6 diagonals of the faces of a cube is stable (see Fig.1.c). If we remove one diagonal then the structure becomes unstable (see Fig.1.d).*



a. Stable tetrahedron *b. Unstable tetrahedron* *c. Stable cube* *d. Unstable cube*

Figure 1: Stable and unstable structures

We equip $\mathbf{U}_{\mathcal{S}}$ with the following bilinear form:

$$< \Phi, \Psi >_1 = \int_{\mathcal{S}} \frac{\partial \Phi}{\partial S_1} \cdot \mathbf{t}_1 \, \frac{\partial \Psi}{\partial S_1} \cdot \mathbf{t}_1 \, d\mathbf{S}, \qquad \forall \, (\Phi, \Psi) \in \mathbf{U}_{\mathcal{S}} \times \mathbf{U}_{\mathcal{S}} \tag{2.2}$$

and the associated semi-norm

$$\|U\|_{\mathcal{S}} \doteq \sqrt{< U, U >_1} = \left\| \frac{dU}{d\mathbf{S}} \cdot \mathbf{t}_1 \right\|_{L^2(\mathcal{S})}, \qquad \forall \, U \in \mathbf{U}_{\mathcal{S}}. \tag{2.3}$$

Lemma 2.1. *Let $\mathcal{S}$ be a stable structure. There exists a constant C, which depends on $\mathcal{S}$, such that for every U in $\mathbf{U}_{\mathcal{S}}$ there exists $\mathbf{r} \in \mathbf{R}$ such that*

$$\|U - \mathbf{r}\|_{H^1(\mathcal{S})} \leq C \|U\|_{\mathcal{S}}. \tag{2.4}$$

Proof. Let $\mathbf{R}^{\perp}$ be the orthonormal of $\mathbf{R}$ in $\mathbf{U}_{\mathcal{S}}$ for the scalar product

$$< \Phi, \Psi >= \int_{\mathcal{S}} \Phi \cdot \Psi \, d\mathbf{S}, \qquad \forall \, (\Phi, \Psi) \in \mathbf{U}_{\mathcal{S}} \times \mathbf{U}_{\mathcal{S}}.$$

If U belongs to $\mathbf{R}^{\perp}$ and satisfies $\|U\|_{\mathcal{S}} = 0$ then, since $\mathcal{S}$ is a stable structure, U belongs to $\mathbf{R}$. Therefore U is equal to 0. The semi-norm $\|\cdot\|_{\mathcal{S}}$ is a norm on the space $\mathbf{R}^{\perp}$. Since $\mathbf{R}^{\perp}$ is a finite dimensional vector space, all the norms are equivalent. Thus (2.4) is proved. □

3 Decomposition of beam displacements

In this section, we remind some results concerning the decomposition of a beam displacement. These results will be used later and can be found in [15]. For the sake of simplicity these results are formulated for the beam $B_{l,\mathfrak{r}} \doteq (0, l) \times D_{\mathfrak{r}}$ whose cross-sections are disc of radius $\mathfrak{r}$ ($\mathfrak{r} \leq l$). The beam is referred to the orthonormal frame $(O; \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ ($\mathbf{e}_1$ is the direction of the centerline). In this frame the running point is denoted $x = (x_1, x_2, x_3)$.

Any displacement $u \in H^1(B_{l,\mathfrak{r}})^3$ of the beam $B_{l,\mathfrak{r}}$ is uniquely decomposed as follows

$$u = U^e + \overline{u} \tag{3.1}$$

where U^e is called elementary displacement and it stands for the displacement of the centerline of the beam and the small rotation of the cross-section at every point of the centerline

$$U^e(x) = \mathcal{U}(x_1) + \mathcal{R}(x_1) \wedge (x_2 \mathbf{e}_2 + x_3 \mathbf{e}_3), \quad \text{for a.e. } x = (x_1, x_2, x_3) \in B_{l,\mathfrak{r}}. \tag{3.2}$$

$\mathcal{U} = (\mathcal{U}_1, \mathcal{U}_2, \mathcal{U}_3)$ and $\mathcal{R} = (\mathcal{R}_1, \mathcal{R}_2, \mathcal{R}_3)$ belong to $H^1(0, l)^3$. The residual displacement $\overline{u} \in H^1(B_{l,\mathfrak{r}})^3$ is the warping (the deformation of the cross-sections), it satisfies (for more details see [15])

$$\int_{D_{\mathfrak{r}}} \overline{u}(x) \, dx_2 dx_3 = \int_{D_{\mathfrak{r}}} \overline{u}(x) \wedge (x_2 \mathbf{e}_2 + x_3 \mathbf{e}_3) \, dx_2 dx_3 = 0 \quad \text{for a.e. } x_1 \in (0, l). \tag{3.3}$$

Taking into account the decomposition (3.1) and the representation for the elementary displacement given by (3.2) the strain tensor $e(u)$ has the following form:

$$\begin{aligned} e(u) &= e(U^e) + e(\overline{u}) \\ &= \begin{pmatrix} \frac{d\mathcal{U}_1}{dx_1} - x_2\frac{d\mathcal{R}_3}{dx_1} + x_3\frac{d\mathcal{R}_2}{dx_1} & \frac{1}{2}\Big[\Big(\frac{d\mathcal{U}_2}{dx_1} - \mathcal{R}_3\Big) - x_3\frac{d\mathcal{R}_1}{dx_1}\Big] & \frac{1}{2}\Big[\Big(\frac{d\mathcal{U}_3}{dx_1} + \mathcal{R}_2\Big) + x_2\frac{d\mathcal{R}_1}{dx_1}\Big] \\ \frac{1}{2}\Big[\Big(\frac{d\mathcal{U}_2}{dx_1} - \mathcal{R}_3\Big) - x_3\frac{d\mathcal{R}_1}{dx_1}\Big] & 0 & 0 \\ \frac{1}{2}\Big[\Big(\frac{d\mathcal{U}_3}{dx_1} + \mathcal{R}_2\Big) + x_2\frac{d\mathcal{R}_1}{dx_1}\Big] & 0 & 0 \end{pmatrix} + e(\overline{u}). \end{aligned} \tag{3.4}$$

Below is a lemma proven in [13, 15]. It gives estimates for the warping and the terms from U^e in the above strain tensor (3.4).

Lemma 3.1. *Let u be in $H^1(B_{l,\mathfrak{r}})^3$ decomposed as (3.1)-(3.2)-(3.3). The following estimates hold:*

$$\begin{aligned} &\|\overline{u}\|_{L^2(B_{l,\mathfrak{r}})} \le C\mathfrak{r}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \qquad \|\nabla\overline{u}\|_{L^2(B_{l,\mathfrak{r}})} \le C\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \\ &\Big\|\frac{d\mathcal{R}}{dx_1}\Big\|_{L^2(0,l)} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \quad \Big\|\frac{d\mathcal{U}}{dx_1} - \mathcal{R}\wedge\mathbf{e}_1\Big\|_{L^2(0,l)} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}. \end{aligned} \tag{3.5}$$

The constants are independent of l and $\mathfrak{r} \le l$.

The function $\mathcal{U}$, defined in (3.1), is decomposed into the sum of two functions $\mathcal{U}^h$ and $\overline{\mathcal{U}}$, where $\mathcal{U}^h$ coincides with $\mathcal{U}$ in the extremities of the centerline and is affine between them, $\overline{\mathcal{U}} = \mathcal{U} - \mathcal{U}^h$ is the residual part, i.e.

$$\mathcal{U}^h(x_1) = \frac{(l-x_1)}{l}\mathcal{U}(0) + \frac{x_1}{l}\mathcal{U}(l).$$

In the same way the function $\mathcal{R}$, defined in (3.1), is decomposed into the sum of two functions $\mathcal{R}^h$ and $\overline{\mathcal{R}}$. It is obvious, but important to note that

$$\overline{\mathcal{U}}(0) = \overline{\mathcal{U}}(l) = 0, \quad \overline{\mathcal{R}}(0) = \overline{\mathcal{R}}(l) = 0.$$

Lemma 3.2. *The following estimates hold:*

$$\begin{aligned} &\Big\|\frac{d\overline{\mathcal{R}}}{dx_1}\Big\|_{L^2(0,l)} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \quad \|\overline{\mathcal{R}}\|_{L^2(0,l)} \le \frac{Cl}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})} \\ &\Big\|\frac{d\overline{\mathcal{U}}}{dx_1}\cdot\mathbf{e}_1\Big\|_{L^2(0,l)} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \quad \Big\|\frac{d\overline{\mathcal{U}}}{dx_1}\Big\|_{L^2(0,l)} \le C\frac{l}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \\ &\|\overline{\mathcal{U}}\cdot\mathbf{e}_1\|_{L^2(0,l)} \le C\frac{l}{\mathfrak{r}}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \qquad \|\overline{\mathcal{U}}\|_{L^2(0,l)} \le C\frac{l^2}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \\ &\Big\|\frac{d\mathcal{U}^h}{dx_1} - \mathcal{R}^h\wedge\mathbf{e}_1\Big\|_{L^2(0,l)} \le \frac{Cl}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}, \quad \Big\|\frac{d\mathcal{R}^h}{dx_1}\Big\|_{L^2(0,l)} + \frac{1}{\mathfrak{r}}\Big\|\frac{d\mathcal{U}^h}{dx_1}\cdot\mathbf{e}_1\Big\|_{L^2(0,l)} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})}. \end{aligned} \tag{3.6}$$

The constants do not depend on l and $\mathfrak{r}$.

Proof. Since $\dfrac{d\mathcal{R}^h}{dx_1}$ and $\dfrac{d\overline{\mathcal{U}}}{dx_1} - (\mathcal{R} - m(\mathcal{R}))\wedge\mathbf{e}_1$ ($m(\mathcal{R}) = \dfrac{1}{l}\displaystyle\int_0^l \mathcal{R}(t)\,dt$) are constant on $(0,l)$, one gets

$$\begin{aligned} &\Big\|\frac{d\mathcal{R}^h}{dx_1}\Big\|^2_{L^2(0,l)} + \Big\|\frac{d\overline{\mathcal{R}}}{dx_1}\Big\|^2_{L^2(0,l)} = \Big\|\frac{d\mathcal{R}}{dx_1}\Big\|^2_{L^2(0,l)} \le \frac{C}{\mathfrak{r}^4}\|e(u)\|^2_{L^2(B_{l,\mathfrak{r}})}, \\ &\Big\|\frac{d\mathcal{U}^h}{dx_1} - m(\mathcal{R})\wedge\mathbf{e}_1\Big\|^2_{L^2(0,l)} + \Big\|\frac{d\overline{\mathcal{U}}}{dx_1} - (\mathcal{R}-m(\mathcal{R}))\wedge\mathbf{e}_1\Big\|^2_{L^2(0,l)} = \Big\|\frac{d\mathcal{U}}{dx_1} - \mathcal{R}\wedge\mathbf{e}_1\Big\|^2_{L^2(0,l)} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|^2_{L^2(B_{l,\mathfrak{r}})}. \end{aligned}$$

Then, the Poincaré and the Poincaré-Wirtinger inequalities together with the above estimates yield

$$\|\mathcal{R} - \mathcal{R}^h\|_{L^2(0,l)} = \|\overline{\mathcal{R}}\|_{L^2(0,l)} \le C\frac{l}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})} \quad \text{and} \quad \|\mathcal{R} - m(\mathcal{R})\|_{L^2(0,l)} \le C\frac{l}{\mathfrak{r}^2}\|e(u)\|_{L^2(B_{l,\mathfrak{r}})},$$

from which we derive the other estimates in (3.6). □

4 Decomposition of the displacements of a beam structure

From now on, $\mathcal{S}$ is a stable structure.

The beam structure $\mathcal{S}_{1,\mathfrak{r}}$ is defined as follows:

$$\mathcal{S}_{1,\mathfrak{r}} = \big\{x \in \mathbb{R}^3 \mid \mathrm{dist}(x,\mathcal{S}) < \mathfrak{r}\big\}.$$

For $\ell \in \{1,\dots,m\}$, denote $\mathcal{P}_{\ell,\mathfrak{r}}$ the straight beam with centerline $\gamma_\ell = [\mathbf{A}^\ell, \mathbf{B}^\ell]$ and reference cross-section the disk $D_\mathfrak{r} \doteq D(O,\mathfrak{r})$ of radius $\mathfrak{r}$, $0 < \mathfrak{r} \le l_\ell$ (the disk D_1 for simplicity will be denoted D). The straight beam $\mathcal{P}_{\ell,\mathfrak{r}}$ is referred to the orthonormal frame $(\mathbf{A}^\ell; \mathbf{t}_1^\ell, \mathbf{t}_2^\ell, \mathbf{t}_3^\ell)$

$$\mathcal{P}_{\ell,\mathfrak{r}} = \big\{x \in \mathbb{R}^3 \mid x = \mathbf{A}^\ell + S_1\mathbf{t}_1^\ell + S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell, \quad (S_1,S_2,S_3) \in (0,l_\ell)\times D_\mathfrak{r}\big\}. \tag{4.1}$$

By definition, the whole structure $\mathcal{S}_{1,\mathfrak{r}}$ contains the straight beams $\mathcal{P}_{\ell,\mathfrak{r}}$, $\ell \in \{1,\dots,m\}$ and the balls of radius $\mathfrak{r}$ centered in the points of $\mathcal{K}$, more precisely one has

$$\mathcal{S}_{1,\mathfrak{r}} \doteq \Big(\bigcup_{A\in\mathcal{K}} B(A,\mathfrak{r})\Big) \cup \Big(\bigcup_{\ell=1}^{m} \mathcal{P}_{\ell,\mathfrak{r}}\Big).$$

The set of junction domains is denoted by $\mathcal{J}_\mathfrak{r}$. There exists c_0 which only depends on $\mathcal{S}$ such that

$$\mathcal{J}_\mathfrak{r} \subset \bigcup_{A\in\mathcal{K}} B(A, c_0\mathfrak{r}).$$

The set $\mathcal{J}_\mathfrak{r}$ is defined in such a way that $\mathcal{S}_{1,\mathfrak{r}} \setminus \overline{\mathcal{J}}_\mathfrak{r}$ only consists of disjoint straight beams.

Definition 4.1. *An elementary beam-structure displacement is a displacement U^e belonging to $H^1(\mathcal{S}_{1,\mathfrak{r}})^3$ whose restriction to each beam is an elementary displacement and whose restriction to each junction is a rigid displacement*

$$\begin{aligned}
&U^e(x) = \mathcal{U}(\mathbf{A}^\ell + S_1\mathbf{t}_1^\ell) + \mathcal{R}(\mathbf{A}^\ell + S_1\mathbf{t}_1^\ell) \wedge (S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell),\\
&\qquad \textit{for a.e. } x = \mathbf{A}^\ell + S_1\mathbf{t}_1^\ell + S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell \in \mathcal{P}_{\ell,\mathfrak{r}},\ (S_1,S_2,S_3) \in (0,l_\ell)\times D_\mathfrak{r},\ \ell \in \{1,\dots,m\},\\
&U^e(x) = \mathcal{U}(A) + \mathcal{R}(A) \wedge (x - A), \quad \textit{for a.e. } x \in B(A,r), \quad \textit{for all } A \in \mathcal{K}
\end{aligned}$$

with $\mathcal{U}$ and $\mathcal{R}$ in $H^1(\mathcal{S})^3$.

In [14] it is shown that every displacement $u \in H^1(\mathcal{S}_{1,\mathfrak{r}})^3$ can be decomposed as

$$u = U^e + \overline{u},$$

where U^e is an elementary beam-structure displacement and where $\overline{u} \in H^1(\mathcal{S}_{1,\mathfrak{r}})^3$ is the warping. Here, the pair $(U^e, \overline{u})$ is not uniquely determined. Furthermore, the warping satisfies the conditions (3.3) "outside" the domain $\mathcal{J}_\mathfrak{r}$ (see [14, 15]), more precisely, one has ($\ell \in \{1,\dots,m\}$)

$$\begin{aligned}
&\int_{D_r} \overline{u}(S_1,S_2,S_3)\, dS_2dS_3 = 0,\\
&\int_{D_r} \overline{u}(S_1,S_2,S_3) \wedge (S_2\mathbf{e}_2 + S_3\mathbf{e}_3)\, dS_2dS_3 = 0,
\end{aligned} \qquad \text{for a.e. } S_1 \in (2c_0\mathfrak{r}, l_\ell - 2c_0\mathfrak{r}). \tag{4.2}$$

The following lemma is proved in [14, Lemma 3.4]:

Lemma 4.1. *Let u be in $H^1(\mathcal{S}_{1,\mathfrak{r}})^3$. There exists a decomposition of u, $u = U^e + \overline{u}$ for which U^e is an elementary beam-structure displacement. The terms of this decomposition satisfy*

$$\begin{aligned}
&\|\overline{u}\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})} \le C\mathfrak{r}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}, \quad \|\nabla\overline{u}\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})} \le C\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})},\\
&\Big\|\frac{d\mathcal{R}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}, \quad \Big\|\frac{d\mathcal{U}}{d\mathbf{S}} - \mathcal{R}\wedge\mathbf{t}_1\Big\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}.
\end{aligned} \tag{4.3}$$

The constants do not depend on $\mathfrak{r}$.

Here, again we split the field $\mathcal{U}$ into the sum of two fields $\mathcal{U}^h$ and $\overline{\mathcal{U}}$, where $\mathcal{U}^h$ coincides with $\mathcal{U}$ in the nodes of $\mathcal{S}$ and is affine between two contiguous nodes and $\overline{\mathcal{U}} = \mathcal{U} - \mathcal{U}^h$ is the residual part.
In the same way the fields $\mathcal{R}^h$ and $\overline{\mathcal{R}}$ are introduced. The field $\mathcal{U}^h$ describes the displacement of the nodes, i.e. the global behavior of the structure, whereas $\overline{\mathcal{U}}$ stands for the local displacement of the beams.
By construction the fields $\mathcal{U}^h$ and $\mathcal{R}^h$ belong to $\mathbf{U}_{\mathcal{S}}$. Furthermore one has

Lemma 4.2. *For every $u \in H^1(\mathcal{S}_{1,\mathfrak{r}})^3$ the following estimates hold:*

$$\begin{aligned}
&\left\|\frac{d\overline{\mathcal{R}}}{d\mathbf{S}}\right\|_{L^2(\mathcal{S})} + \left\|\overline{\mathcal{R}}\right\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})},\\
&\left\|\frac{d\overline{\mathcal{U}}}{d\mathbf{S}}\cdot\mathbf{t}_1\right\|_{L^2(\mathcal{S})} + \left\|\overline{\mathcal{U}}\cdot\mathbf{t}_1\right\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})},\\
&\left\|\frac{d\overline{\mathcal{U}}}{d\mathbf{S}}\right\|_{L^2(\mathcal{S})} + \left\|\overline{\mathcal{U}}\right\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})},\\
&\left\|\frac{d\mathcal{U}^h}{d\mathbf{S}} - \mathcal{R}^h\wedge\mathbf{t}_1\right\|_{L^2(\mathcal{S})} + \left\|\frac{d\mathcal{R}^h}{d\mathbf{S}}\right\|_{L^2(\mathcal{S})} + \frac{1}{\mathfrak{r}}\left\|\frac{d\mathcal{U}^h}{d\mathbf{S}}\cdot\mathbf{t}_1\right\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}.
\end{aligned} \tag{4.4}$$

Moreover, since $\mathcal{S}$ is a stable structure, there exists a rigid displacement $\mathbf{r} \in \mathbf{R}$, $(\mathbf{r}(x) = \mathbf{a} + \mathbf{b}\wedge x)$, such that

$$\|\mathcal{U}^h - \mathbf{r}\|_{H^1(\mathcal{S})} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}, \qquad \|\mathcal{R}^h - \mathbf{b}\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}. \tag{4.5}$$

The constants do not depend on $\mathfrak{r}$.

Proof. Estimates (4.4) are the immediate consequences of the Lemmas 3.2 and 4.1. Since $\mathcal{S}$ is a stable structure, Lemma 2.1 and again (4.4) yield a rigid displacement $\mathbf{r} \in \mathbf{R}$ $(\mathbf{r}(x) = \mathbf{a} + \mathbf{b}\wedge x)$ such that $(4.5)_1$ holds.
Besides, from the Poincaré-Wirtinger inequality and $(4.4)_4$, there exists $\widetilde{\mathbf{b}} \in \mathbb{R}^3$ such that

$$\|\mathcal{R}^h - \widetilde{\mathbf{b}}\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}.$$

The constant does not depend on $\mathfrak{r}$. Then, $(4.5)_1$ and the above estimate give

$$\|(\mathbf{b} - \widetilde{\mathbf{b}})\wedge\mathbf{t}_1\|_{L^2(\mathcal{S})} \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}.$$

Since the structure has more than two segments with non-collinear directions, this yields

$$|\mathbf{b} - \widetilde{\mathbf{b}}| \le \frac{C}{\mathfrak{r}^2}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}.$$

Hence, $(4.5)_2$ is proved. □

Let $\mathcal{S}$ be a stable structure such that $\mathcal{S}\cup(\mathcal{S}+\mathbf{e}_1)$ is a stable structure. For every displacement $u \in H^1(\mathcal{S}_{1,\mathfrak{r}}\cup(\mathcal{S}_{1,\mathfrak{r}}+\mathbf{e}_1))^3$, Lemma 4.2 gives two rigid displacements $\mathbf{r}_0$, $\mathbf{r}_1$ such that

$$\begin{aligned}
&\mathbf{r}_0(x) = \mathbf{a}_0 + \mathbf{b}_0\wedge(x-G), \quad \mathbf{r}_1(x) = \mathbf{a}_1 + \mathbf{b}_1\wedge(x-G-\mathbf{e}_1) \qquad \forall x\in\mathbb{R}^3,\\
&\|\mathcal{U}^h - \mathbf{r}_0\|_{H^1(\mathcal{S})} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}})}, \quad \|\mathcal{U}^h - \mathbf{r}_1\|_{H^1(\mathcal{S}+\mathbf{e}_1)} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}+\mathbf{e}_1)},
\end{aligned} \tag{4.6}$$

where G is the center of mass of $\mathcal{S}$.

Lemma 4.3. *Let $\mathcal{S}$ be a stable structure such that $\mathcal{S}\cup(\mathcal{S}+\mathbf{e}_1)$ is also a stable structure. The following estimate holds:*

$$\|\mathbf{r}_1 - \mathbf{r}_0\|_{H^1(\mathcal{S}\cup(\mathcal{S}+\mathbf{e}_1))} \le \frac{C}{\mathfrak{r}}\|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}}\cup(\mathcal{S}_{1,\mathfrak{r}}+\mathbf{e}_1))}. \tag{4.7}$$

The constant does not depend on $\mathfrak{r}$.

Proof. From Lemma 4.2, there exists a rigid displacement $\mathbf{r}$ such that

$$\mathbf{r}(x) = \mathbf{a} + \mathbf{b} \wedge (x - G - \mathbf{e}_1/2) \qquad \forall\, x \in \mathbb{R}^3,$$
$$\|\mathcal{U}^h - \mathbf{r}\|_{H^1(\mathcal{S}\cup(\mathcal{S}+\mathbf{e}_1))} \leq \frac{C}{\mathfrak{r}} \|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}}\cup(\mathcal{S}_{1,\mathfrak{r}}+\mathbf{e}_1))}.$$

The constant does not depend on $\mathfrak{r}$. Hence

$$\|\mathbf{r} - \mathbf{r}_0\|_{H^1(\mathcal{S})} \leq \frac{C}{\mathfrak{r}} \|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}}\cup(\mathcal{S}_{1,\mathfrak{r}}+\mathbf{e}_1))}, \qquad \|\mathbf{r} - \mathbf{r}_1\|_{H^1(\mathcal{S}+\mathbf{e}_1)} \leq \frac{C}{\mathfrak{r}} \|e(u)\|_{L^2(\mathcal{S}_{1,\mathfrak{r}}\cup(\mathcal{S}_{1,\mathfrak{r}}+\mathbf{e}_1))}.$$

The above estimates yield (4.7) since in $\mathbf{R}$ the norms $\|\cdot\|_{H^1(\mathcal{S})}$, $\|\cdot\|_{H^1(\mathcal{S}+\mathbf{e}_1)}$ and $\|\cdot\|_{H^1(\mathcal{S}\cup(\mathcal{S}+\mathbf{e}_1))}$ are equivalent. □

5 A periodic beam structure as 3D-like domain

From now on, in all the estimates, we denote by C a strictly positive constant which does not depend on ε and r.

5.1 Notations and statement of the problem

Below we consider periodic structures $\mathcal{S}$ included in a closed parallelotope.

Definition 5.1. *A structure $\mathcal{S}$ is a* 3*D*-**periodic structure** *if for every $i \in \{1,2,3\}$ the set $\mathcal{S} \cup \big(\mathcal{S} + \mathbf{e}_i\big)$ is a structure in the sense of Definition 2.1.*

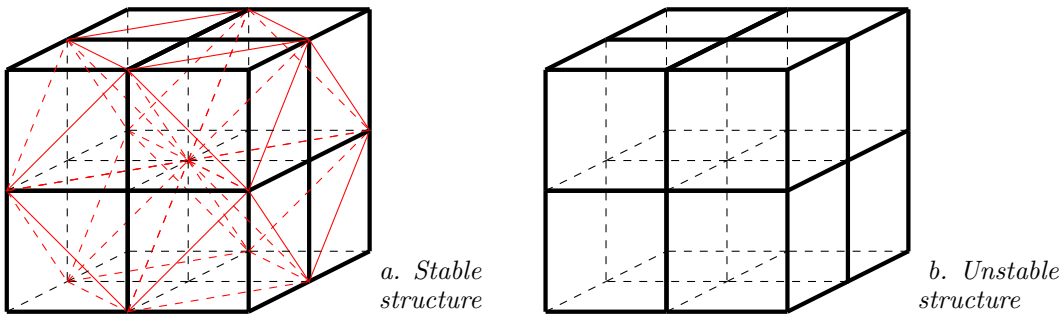
Definition 5.2. *A* 3*D*-**periodic structure** *$\mathcal{S}$ is a* 3*D*-**periodic stable structure** *(briefly* **3-PSS***) if $\mathcal{S}$ and $\mathcal{S} \cup \big(\mathcal{S} + \mathbf{e}_i\big)$, $i \in \{1,2,3\}$, are stable structures in the sense of Definition 2.2.*

Remark 5.1.

1. *The structure made of 12 edges and 6 diagonals of the faces of a cube is a 3D-periodic stable structure (Fig. 2.a).*
2. *The structure made of 12 edges of a cube is not a 3D-periodic stable structure (Fig. 2.b).*

Let Ω be a bounded domain in $\mathbb{R}^3$ with a Lipschitz boundary and Γ be a subset of $\partial\Omega$ with nonnull measure. We assume that there exists an open set Ω' with a Lipschitz boundary such that $\Omega \subset \Omega'$ and $\Omega' \cap \partial\Omega = \Gamma$.
Denote

- $\Omega_1 \doteq \big\{x \in \mathbb{R}^N \mid \mathrm{dist}(x,\Omega) < 1\big\}$, $\Omega_\varepsilon^{int} = \big\{x \in \Omega \mid \mathrm{dist}(x,\partial\Omega) > 2\sqrt{3}\varepsilon\big\}$,
- $Y \doteq (0,1)^3$,
- $G = (1/2, 1/2, 1/2)$ the center of mass of Y,
- $\mathcal{S}$ a 3-periodic structure included in $\overline{Y}$,



a. Stable structure *b. Unstable structure*

Figure 2: 3D-periodic stable and unstable structures

- $\Xi_\varepsilon \doteq \big\{\xi \in \mathbb{Z}^3 \ \mid \ (\varepsilon\xi + \varepsilon Y) \cap \Omega \neq \emptyset\big\}$, $\widetilde{\Xi}_\varepsilon \doteq \big\{\xi \in \mathbb{Z}^3 \ \mid \ (\varepsilon\xi + \varepsilon Y) \subset \Omega\big\}$
- $\Xi_\varepsilon^{int} \doteq \big\{\xi \in \mathbb{Z}^3 \ \mid \ (\varepsilon\xi + \varepsilon Y) \subset \Omega_\varepsilon^{int}\big\}$,
- $\Xi'_\varepsilon \doteq \big\{\xi \in \mathbb{Z}^3 \ \mid \ (\varepsilon\xi + \varepsilon Y) \cap \Omega' \neq \emptyset\big\}$,
- $\widehat{\Xi}_\varepsilon \doteq \big\{\xi \in \Xi_\varepsilon \mid \text{all the vertices of } \xi + \overline{Y} \text{ belong to } \Xi_\varepsilon\big\}$,
- $\Xi_{\varepsilon,i} \doteq \big\{\xi \in \Xi_\varepsilon \mid \xi + \mathbf{e}_i \in \Xi_\varepsilon\big\}$, $i \in \{1,2,3\}$,
- $\Omega_\varepsilon \doteq \text{interior}\Big(\bigcup_{\xi\in\Xi_\varepsilon} (\varepsilon\xi + \varepsilon\overline{Y})\Big)$, $\widehat{\Omega}_\varepsilon \doteq \text{interior}\Big(\bigcup_{\xi\in\widehat{\Xi}_\varepsilon} (\varepsilon\xi + \varepsilon\overline{Y})\Big)$, $\Omega'_\varepsilon \doteq \text{interior}\Big(\bigcup_{\xi\in\Xi'_\varepsilon} (\varepsilon\xi + \varepsilon\overline{Y})\Big)$
- $\widehat{\Omega}_\varepsilon^{int} \doteq \text{interior}\Big(\bigcup_{\xi\in\Xi_\varepsilon^{int}} (\varepsilon\xi + \varepsilon\overline{Y})\Big)$, $\widetilde{\Omega}_\varepsilon \doteq \text{interior}\Big(\bigcup_{\xi\in\widetilde{\Xi}_\varepsilon} (\varepsilon\xi + \varepsilon\overline{Y})\Big)$.

One has

$$\Xi_\varepsilon^{int} \subset \widehat{\Xi}_\varepsilon \subset \bigcap_{i=1}^{3} \Xi_{\varepsilon,i} \subset \bigcup_{i=1}^{3} \Xi_{\varepsilon,i} = \Xi_\varepsilon.$$

The open sets Ω_ε, Ω'_ε, $\widehat{\Omega}_\varepsilon$, $\widehat{\Omega}_\varepsilon^{int}$ and Ω_ε^{int} are connected. Moreover, the following inclusions hold

$$\widehat{\Omega}_\varepsilon^{int} \subset \Omega_\varepsilon^{int} \subset \Omega \subset \Omega_\varepsilon \subset \Omega'_\varepsilon, \qquad \widehat{\Omega}_\varepsilon^{int} \subset \Omega_\varepsilon^{int} \subset \widehat{\Omega}_\varepsilon \subset \Omega_\varepsilon.$$

Set

$$\mathcal{S}_\varepsilon \doteq \bigcup_{\xi\in\Xi_\varepsilon} \big(\varepsilon\xi + \varepsilon\mathcal{S}\big), \qquad \mathcal{S}_{\varepsilon,r} \doteq \big\{x \in \mathbb{R}^3 \mid \text{dist}(x, \mathcal{S}_\varepsilon) < r\big\},$$
$$\mathcal{S}'_\varepsilon \doteq \bigcup_{\xi\in\Xi'_\varepsilon} \big(\varepsilon\xi + \varepsilon\mathcal{S}\big), \qquad \mathcal{S}'_{\varepsilon,r} \doteq \big\{x \in \mathbb{R}^3 \mid \text{dist}(x, \mathcal{S}'_\varepsilon) < r\big\},$$
$$\mathcal{K}_\varepsilon \doteq \bigcup_{\xi\in\Xi_\varepsilon} \big(\varepsilon\xi + \varepsilon\mathcal{K}\big).$$

The running point of $\mathcal{S}_\varepsilon$ is denoted $\mathbf{s}$.

Let $\mathcal{S}_{\varepsilon,r}$ be a beam structure consisting of balls of radius r centered on the points of $\mathcal{K}_\varepsilon$ and beams, whose cross-sections are discs of radius r and their centerlines are the segments of $\mathcal{S}_\varepsilon$

$$\mathcal{P}^\xi_{\varepsilon\ell,r} \doteq \varepsilon\xi + \varepsilon\mathcal{P}_{\ell,\mathfrak{r}}, \quad \ell \in \{1,\dots,m\}, \quad \mathfrak{r} = r/\varepsilon,$$
$$\mathcal{S}_{\varepsilon,r} \doteq \Big(\bigcup_{A\in\mathcal{K}_\varepsilon} B(A,r)\Big) \cup \Big(\bigcup_{\xi\in\Xi_\varepsilon} \bigcup_{\ell=1}^{m} \mathcal{P}^\xi_{\varepsilon\ell,r}\Big).$$

The parametrization of the beam $\mathcal{P}^\xi_{\varepsilon\ell,r}$ ($\ell \in \{1,\dots,m\}$) is given by (see (4.1))

$$x = \varepsilon\xi + \varepsilon\mathbf{A}^\ell + s_1\mathbf{t}_1^\ell + s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell, \quad (s_1,s_2,s_3) \in (0, \varepsilon l_\ell) \times D_r.$$

The junction domains (the common parts of the beams) is denoted $\mathcal{J}_{\varepsilon,r}$. One has

$$\bigcup_{A\in\mathcal{K}_\varepsilon} B(A,r) \subset \mathcal{J}_{\varepsilon,r} \subset \bigcup_{A\in\mathcal{K}_\varepsilon} B(A,c_0 r). \tag{5.1}$$

The structure $\mathcal{S}_{\varepsilon,r}$ is included in Ω_ε.

The space of all admissible displacements is denoted $\mathbf{V}_{\varepsilon,r}$

$$\mathbf{V}_{\varepsilon,r} = \big\{u \in H^1(\mathcal{S}_{\varepsilon,r})^3 \ \mid \ \exists u' \in H^1(\mathcal{S}'_{\varepsilon,r})^3 \text{ such that } u'_{|\mathcal{S}_{\varepsilon,r}} = u \text{ and } u' = 0 \text{ in } \mathcal{S}'_{\varepsilon,r} \setminus \overline{\mathcal{S}_{\varepsilon,r}}\big\}.$$

It means that the displacements belonging to $\mathbf{V}_{\varepsilon,r}$ "vanish" on a part $\Gamma_{\varepsilon,r}$ included in $\partial\mathcal{S}_{\varepsilon,r} \cap \partial\Omega$.
We assume that $\mathcal{S}_{\varepsilon,r}$ is made of isotropic and homogeneous material.

For a displacement $u \in \mathbf{V}_{\varepsilon,r}$, we denote by e the strain tensor (or symmetric gradient)

$$e(u) \doteq \frac{1}{2}\Big(\nabla u + (\nabla u)^T\Big), \quad e_{ij}(u) \doteq \frac{1}{2}\Big(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i}\Big). \tag{5.2}$$

We have two coordinate systems. The first one is the global Cartesian system (x_1, x_2, x_3) and is related to the frame $(O; \mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$. The second one is the local coordinate system (s_1, s_2, s_3) defined for every beam and related to the frame $(\varepsilon\xi + \varepsilon\mathbf{A}_2^\ell; \mathbf{t}_1^\ell, \mathbf{t}_2^\ell, \mathbf{t}_3^\ell)$, $\ell \in \{1, \dots, m\}$. The orthonormal transformation matrix from the basis $(\mathbf{t}_1^\ell, \mathbf{t}_2^\ell, \mathbf{t}_3^\ell)$ to the basis $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ is $\mathbf{T}^\ell = \big(\mathbf{t}_1^\ell \mid \mathbf{t}_2^\ell \mid \mathbf{t}_3^\ell\big)$, this matrix belongs to $SO(3)$.

Hence, for every displacement $v \in H^1(\mathcal{P}^\xi_{\varepsilon\ell,r})$ a straightforward calculation gives

$$\begin{aligned} e(v) =& \frac{1}{2}\Big(\nabla_x v + \big(\nabla_x v\big)^T\Big) = \frac{1}{2}\mathbf{T}^\ell\Big(\nabla_s v + \big(\nabla_s v\big)^T\Big)(\mathbf{T}^\ell)^T = \frac{1}{2}\mathbf{T}^\ell\, e_s(v)\, (\mathbf{T}^\ell)^T \\ e_s(v) =& \begin{pmatrix} \frac{\partial v}{\partial s_1}\cdot\mathbf{t}_1^\ell & \frac{1}{2}\big(\frac{\partial v}{\partial s_2}\cdot\mathbf{t}_1^\ell + \frac{\partial v}{\partial s_1}\cdot\mathbf{t}_2^\ell\big) & \frac{1}{2}\big(\frac{\partial v}{\partial s_3}\cdot\mathbf{t}_1^\ell + \frac{\partial v}{\partial s_1}\cdot\mathbf{t}_3^\ell\big) \\ * & \frac{\partial v}{\partial s_2}\cdot\mathbf{t}_2^\ell & \frac{1}{2}\big(\frac{\partial v}{\partial s_3}\cdot\mathbf{t}_2^\ell + \frac{\partial v}{\partial s_2}\cdot\mathbf{t}_3^\ell\big) \\ * & * & \frac{\partial v}{\partial s_3}\cdot\mathbf{t}_3^\ell \end{pmatrix}. \end{aligned} \tag{5.3}$$

Let $a^{\varepsilon,r}_{ijkl} \in L^\infty(\mathcal{S}_{\varepsilon,r}), (i,j,k,l) \in \{1,2,3\}^4$, be the components of the elasticity tensor. These functions satisfy the usual symmetry and positivity conditions

- $a^{\varepsilon,r}_{ijkl} = a^{\varepsilon,r}_{jikl} = a^{\varepsilon,r}_{klij}$ a.e. in $\mathcal{S}_{\varepsilon,r}$;
- for any $\tau \in M^3_s$, where M^3_s is the space of 3×3 symmetric matrices, there exists $C_0 > 0$ (independent of ε and r) such that
$$a^{\varepsilon,r}_{ijkl}\tau_{ij}\tau_{kl} \geq C_0\tau_{ij}\tau_{ij} \quad \text{a.e. in } \mathcal{S}_{\varepsilon,r}. \tag{5.4}$$

The coefficients a^ε_{ijkl} are given via the functions $a_{ijkl} \in L^\infty(\mathcal{S}\times D)$

$$\begin{aligned} &a^{\varepsilon,r}_{ijkl}(x) = a^{\varepsilon,r}_{ijkl}(\varepsilon\xi + \varepsilon A^\ell + s_1\mathbf{t}_1^\ell + s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell) = a_{ijkl}\Big(A^\ell + \frac{s_1}{\varepsilon}, \frac{s_2}{r}, \frac{s_3}{r}\Big) \\ &\text{for a.e. } x = \varepsilon\xi + \varepsilon A^\ell + s_1\mathbf{t}_1^\ell + s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell \text{ in } \mathcal{P}^\xi_{\ell,r}, \quad \ell \in \{1,\dots,m\}, \quad \xi \in \Xi_\varepsilon. \end{aligned} \tag{5.5}$$

The constitutive law for the material occupying the domain $\mathcal{S}_{\varepsilon,r}$ is given by the relation between the linearized strain tensor and the stress tensor

$$\sigma_{ij}(u) \doteq a^{\varepsilon,r}_{ijkl}\, e_{s,kl}(u), \quad \forall\, u \in \mathbf{V}_{\varepsilon,r}. \tag{5.6}$$

The unknown displacement u_ε[1] $: \mathcal{S}_{\varepsilon,r} \to \mathbb{R}^3$ is the solution to the linearized elasticity system:

$$\begin{cases} \nabla\cdot\sigma(u_\varepsilon) = -f_\varepsilon & \text{in} \quad \mathcal{S}_{\varepsilon,r}, \\ u_\varepsilon = 0 & \text{on} \quad \Gamma_{\varepsilon,r}\cap\partial\mathcal{S}_{\varepsilon,r}, \\ \sigma(u_\varepsilon)\cdot\nu_\varepsilon = 0 & \text{on} \quad \partial\mathcal{S}_{\varepsilon,r}\setminus\Gamma_{\varepsilon,r}, \end{cases} \tag{5.7}$$

where ν_ε is the outward normal vector to $\partial\mathcal{S}_{\varepsilon,r}\setminus\Gamma$, f_ε is the density of volume forces.

The variational formulation of problem (5.7) is

$$\begin{cases} \text{Find } u_\varepsilon \in \mathbf{V}_{\varepsilon,r} \text{ such that,} \\ \displaystyle\int_{\mathcal{S}_{\varepsilon,r}} \sigma(u_\varepsilon) : e(v)\, dx = \int_{\mathcal{S}_{\varepsilon,r}} f_\varepsilon\cdot v\, dx, \qquad \forall\, v \in \mathbf{V}_{\varepsilon,r}. \end{cases} \tag{5.8}$$

5.2 Final decomposition of the displacements of a periodic beam stable structure as a 3D-like domain

Let u be a displacement belonging to $\mathbf{V}_{\varepsilon,r}$. As proved in [14], we can decompose u as the sum of an elementary displacement and a warping.

The decompositions introduced in Section 4, the estimates of Lemma 4.2 lead to the following estimates:

[1]Of course, the solution to this problem depends on ε and r, but for simplicity, we omit the index r. The same holds for the applied forces f_ε and for every function which in fact depends on both indexes.

Lemma 5.1. *For every $u \in \mathbf{V}_{\varepsilon,r}$ the following estimates hold:*

$$\begin{aligned}
&\big\|\overline{u}\big\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le Cr\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad \big\|\nabla\overline{u}\big\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{d\mathcal{R}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \Big\|\frac{d\mathcal{U}}{d\mathbf{s}} - \mathcal{R}\wedge\mathbf{t}_1\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned} \tag{5.9}$$

Moreover, one has

$$\begin{aligned}
&\Big\|\frac{d\overline{\mathcal{R}}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \|\overline{\mathcal{R}}\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{d\overline{\mathcal{U}}}{d\mathbf{s}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \big\|\overline{\mathcal{U}}\cdot\mathbf{t}_1\big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{d\overline{\mathcal{U}}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \big\|\overline{\mathcal{U}}\big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon^2}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{d\mathcal{U}^h}{d\mathbf{S}} - \mathcal{R}^h\wedge\mathbf{t}_1\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \Big\|\frac{d\mathcal{R}^h}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} + \frac{1}{r}\Big\|\frac{d\mathcal{U}^h}{d\mathbf{s}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned} \tag{5.10}$$

Proof. We apply Lemma 4.2 to the structure $\varepsilon(\xi+\mathcal{S}_{1,\mathfrak{r}})$. Replacing $\mathfrak{r}$ by $\dfrac{r}{\varepsilon}$ and then summing over all $\xi\in\Xi_\varepsilon$ give the estimates (5.9) and (5.10). □

Let u be in $H^1(\mathcal{S}_{\varepsilon,r})^3$. In Lemma 4.2 replace $\mathcal{S}_{1,\mathfrak{r}}$ by $\varepsilon(\xi+\mathcal{S}_{r/\varepsilon})$, with $\xi\in\Xi_\varepsilon$, and let $\mathbf{r}_{\varepsilon\xi}$ be a rigid displacement given by this lemma

$$\mathbf{r}_{\varepsilon\xi}(x) = \mathbf{a}(\varepsilon\xi) + \mathbf{b}(\varepsilon\xi)\wedge(x-\varepsilon G-\varepsilon\xi), \quad \forall\, x\in\mathbb{R}^3.$$

One has

$$\begin{cases}
\|\mathcal{U}^h - \mathbf{r}_{\varepsilon\xi}\|_{L^2(\varepsilon(\xi+\mathcal{S}))} \le \dfrac{C\varepsilon}{r}\|e(u)\|_{L^2(\varepsilon(\xi+\mathcal{S}_{r/\varepsilon}))},\\
\Big\|\dfrac{d\mathcal{U}^h}{d\mathbf{s}} - \mathbf{b}(\varepsilon\xi)\wedge\mathbf{t}_1\Big\|_{L^2(\varepsilon(\xi+\mathcal{S}))} \le \dfrac{C}{r}\|e(u)\|_{L^2(\varepsilon(\xi+\mathcal{S}_{r/\varepsilon}))}
\end{cases} \tag{5.11}$$

and

$$\|\mathcal{R}^h - \mathbf{b}(\varepsilon\xi)\|_{L^2(\varepsilon(\xi+\mathcal{S}))} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\varepsilon(\xi+\mathcal{S}_{r/\varepsilon}))}. \tag{5.12}$$

Recall that if ξ belongs to $\Xi_{\varepsilon,i}$, the domains $\varepsilon(\xi+\mathcal{S}_{r/\varepsilon})$ and $\varepsilon(\xi+\mathbf{e}_i+\mathcal{S}_{r/\varepsilon})$, $i\in\{1,2,3\}$, are included in $\mathcal{S}_{\varepsilon,r}$. Then, applying estimates (4.7) in Lemma 4.3 to the structure $\varepsilon(\xi+\mathcal{S}_{r/\varepsilon})$ we obtain

$$\begin{cases}
\displaystyle\sum_{\imath=1}^{3}\sum_{\xi\in\Xi_{\varepsilon,i}} |\mathbf{b}(\varepsilon\xi+\varepsilon\mathbf{e}_i) - \mathbf{b}(\varepsilon\xi)|^2\varepsilon^3 \le C\frac{\varepsilon^2}{r^2}\|e(u)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})},\\
\displaystyle\sum_{i=1}^{3}\sum_{\xi\in\Xi_{\varepsilon,i}} |\mathbf{a}(\varepsilon\xi+\varepsilon\mathbf{e}_i) - \mathbf{a}(\varepsilon\xi) - \varepsilon\mathbf{b}(\varepsilon\xi+\varepsilon\mathbf{e}_i)\wedge\mathbf{e}_i|^2\varepsilon^3 \le C\frac{\varepsilon^4}{r^2}\|e(u)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{cases} \tag{5.13}$$

Set

$$\boldsymbol{\mathcal{U}}(\varepsilon\xi) = \mathbf{a}(\varepsilon\xi), \quad \boldsymbol{\mathcal{R}}(\varepsilon\xi) = \mathbf{b}(\varepsilon\xi), \quad \text{for every } \ \xi\in\Xi_\varepsilon.$$

Now, define

- $\boldsymbol{\mathcal{U}}$ (resp. $\boldsymbol{\mathcal{R}}$) in the cell $\varepsilon(\xi+\overline{Y})$, $\xi\in\widehat{\Xi}_\varepsilon$, as the Q_1 interpolate of its values on the vertices of this parallelotope.

$$\boldsymbol{\mathcal{U}},\ \boldsymbol{\mathcal{R}}\in W^{1,\infty}(\widehat{\Omega}_\varepsilon)^3,$$

- $\mathbf{a}$ (resp. $\mathbf{b}$) as a piecewise constant function, equals to $\mathbf{a}(\varepsilon\xi)$ (resp. $\mathbf{b}(\varepsilon\xi)$) in the cell $\varepsilon(\xi+Y)$, $\xi\in\Xi_\varepsilon$.

$$\mathbf{a},\ \mathbf{b}\in L^\infty(\Omega_\varepsilon)^3. \tag{5.14}$$

We remind the following classical results ([10, Lemmas 5.22 and 5.35] and [16, Lemmas 5.2 and 5.3]):

Lemma 5.2. *Let Ω be a bounded domain in $\mathbb{R}^N$ with Lipschitz boundary. There exists $\delta_0 > 0$ such that for all $\delta \in (0, \delta_0]$ the sets $\Omega_\delta^{int} = \big\{x \in \Omega \mid dist(x, \partial\Omega) > \delta\big\}$ are uniformly Lipschitz.*

Lemma 5.3. *Let Ψ be a function defined on Ξ_ε and extended using the classical Q_1 interpolation procedure in a function denoted Ψ and belonging to $W^{1,\infty}(\widehat{\Omega}_\varepsilon)$ then we have*

$$\varepsilon^3 \sum_{\xi\in\Xi_\varepsilon^{int}} |\Psi(\xi)|^2 \le \|\Psi\|^2_{L^2(\Omega_\varepsilon^{int})}, \qquad \sum_{\xi\in\Xi_\varepsilon} |\Psi(\xi)|^2 \le C\Big(\sum_{\xi\in\Xi_\varepsilon^{int}} |\Psi(\xi)|^2 + \sum_{i=1}^{3}\sum_{\xi\in\Xi_{\varepsilon,i}} |\Psi(\xi+\mathbf{e}_i)-\Psi(\xi)|^2\Big). \tag{5.15}$$

Proposition 5.1. *Let $\mathcal{S}$ be a* **3-PSS**. *For every displacement $u \in H^1(\mathcal{S}_{\varepsilon,r})^3$, one has*

$$\begin{aligned}
&\|\nabla\boldsymbol{\mathcal{R}}\|_{L^2(\Omega_\varepsilon^{int})} \le \frac{C}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_i} - \boldsymbol{\mathcal{R}}\wedge\mathbf{e}_i\Big\|_{L^2(\Omega_\varepsilon^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad i\in\{1,2,3\},\\
&\big\|e(\boldsymbol{\mathcal{U}})\big\|_{L^2(\Omega_\varepsilon^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned} \tag{5.16}$$

Moreover, there exists a rigid displacement $\mathbf{r}$ such that

$$\|\boldsymbol{\mathcal{U}} - \mathbf{r}\|_{H^1(\Omega_\varepsilon^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}. \tag{5.17}$$

Proof. The estimates $(5.13)_{1,2}$ and Lemma 5.3 yield

$$\begin{aligned}
&\|\nabla\boldsymbol{\mathcal{R}}\|_{L^2(\Omega_\varepsilon^{int})} \le \|\nabla\boldsymbol{\mathcal{R}}\|_{L^2(\widehat{\Omega}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\Big\|\frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_i} - \boldsymbol{\mathcal{R}}\wedge\mathbf{e}_i\Big\|_{L^2(\Omega_\varepsilon^{int})} \le \Big\|\frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_i} - \boldsymbol{\mathcal{R}}\wedge\mathbf{e}_i\Big\|_{L^2(\widehat{\Omega}_\varepsilon)} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad i\in\{1,2,3\}.
\end{aligned}$$

And $(5.16)_{1,2}$ are proved. From which we get

$$\Big\|\frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_i}\cdot\mathbf{e}_j + \frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_j}\cdot\mathbf{e}_i\Big\|_{L^2(\Omega_\varepsilon^{int})} \le \Big\|\frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_i}\cdot\mathbf{e}_j + \frac{\partial\boldsymbol{\mathcal{U}}}{\partial x_j}\cdot\mathbf{e}_i\Big\|_{L^2(\widehat{\Omega}_\varepsilon)} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \quad \forall\,(i,j)\in\{1,2,3\}^2,$$

which also read $(5.16)_3$. Lemma 5.2 allows to apply the 3D-Korn inequality in the domain Ω_ε^{int} using estimate $(5.16)_3$. That gives (5.17). □

Proposition 5.2. *Let $\mathcal{S}$ be a* **3-PSS**. *For every u in $\mathbf{V}_{\varepsilon,r}$, the following estimates of the elementary displacement holds:*

$$\begin{aligned}
&\|\mathcal{U}\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r}\Big(1+\frac{\varepsilon^2}{r}\Big)\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, && \Big\|\frac{d\mathcal{U}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\|\mathcal{R}\|_{L^2(\mathcal{S}_\varepsilon)} + \varepsilon\Big\|\frac{d\mathcal{R}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, &&\\
&\|U^e\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\Big(1+\frac{\varepsilon^2}{r}\Big)\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, && \|\nabla U^e\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned} \tag{5.18}$$

Moreover, one has the Korn type inequalities

$$\|u\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\Big(1+\frac{\varepsilon^2}{r}\Big)\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad \|\nabla u\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}; \tag{5.19}$$

Proof. This proposition is a consequence of Proposition 5.1 and two lemmas postponed in Appendix A. □

5.3 Assumptions on the applied forces

We distinguish two types of applied forces. The first ones are applied in the beams (between the junctions) and the second ones are applied in the junctions.

⋆ *The applied forces $\mathbf{f}_\varepsilon$ in the set of beams* $\displaystyle\bigcup_{\xi\in\Xi_\varepsilon}\bigcup_{\ell=1}^{m}\mathcal{P}^\xi_{\varepsilon\ell,r}$.

For simplicity, we choose these applied forces constant in the cross-sections and equal to

$$\mathbf{f}_\varepsilon = \frac{\varepsilon}{r+\varepsilon^2}\mathbf{f}_{|\mathcal{S}_\varepsilon} \quad \text{a.e. in } \bigcup_{\xi\in\Xi_\varepsilon}\bigcup_{\ell=1}^{m}\mathcal{P}^\xi_{\varepsilon\ell,r}$$

$\star$ *The applied forces $F_{r,\mathcal{K}_\varepsilon}$ in the junctions.*
These forces are defined in the balls centered in the nodes with radius r

$$F_{r,\mathcal{K}_\varepsilon} = \sum_{A\in\mathcal{K}_\varepsilon}\frac{\varepsilon^2}{r^2}F(A)\mathbf{1}_{B(A,r)} + \sum_{A\in\mathcal{K}_\varepsilon}\frac{\varepsilon}{r^3}G(A)\wedge(x-A)\mathbf{1}_{B(A,r)},$$

Lemma 5.4. *Taking the applied forces as*

$$f_\varepsilon = \sum_{A\in\mathcal{K}_\varepsilon}\Big[\frac{\varepsilon^2}{r^2}F(A)+\frac{\varepsilon}{r^3}G(A)\wedge(x-A)\Big]\mathbf{1}_{B(A,r)} + \frac{\varepsilon}{r+\varepsilon^2}\mathbf{f}_{|\mathcal{S}_\varepsilon}\mathbf{1}_{\cup_{\xi\in\Xi_\varepsilon}\cup_{\ell=1}^m\mathcal{P}^\xi_{\varepsilon\ell,r}}, \tag{5.20}$$

where $(\mathbf{f},\, F,\, G)\in\big(C(\overline{\Omega})^3\big)^3$ and where $\mathbf{1}_{\mathcal{O}}$ is the characteristic function of the set $\mathcal{O}$, we obtain

$$\Big|\int_{\mathcal{S}_{\varepsilon,r}} f_\varepsilon\cdot u\,dx\Big| \le C\big(\|\mathbf{f}\|_{L^\infty(\Omega)}+\|F\|_{L^\infty(\Omega)}+\|G\|_{L^\infty(\Omega)}\big)\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\quad \forall\, u\in\mathbf{V}_{\varepsilon,r}. \tag{5.21}$$

Proof. The proof is postponed in Appendix B. □

As a consequence of the above lemma one obtains

Proposition 5.3. *The solution u_ε to the problem* (5.8) *satisfies*

$$\|e(u_\varepsilon)\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\big(\|\mathbf{f}\|_{L^\infty(\Omega)}+\|F\|_{L^\infty(\Omega)}+\|G\|_{L^\infty(\Omega)}\big). \tag{5.22}$$

Proof. In order to obtain apriori estimate of u_ε, we test (5.8) with $v=u_\varepsilon$. From (5.21), we obtain

$$\|e(u_\varepsilon)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\big(\|\mathbf{f}\|_{L^\infty(\Omega)}+\|F\|_{L^\infty(\Omega)}+\|G\|_{L^\infty(\Omega)}\big)\|e(u_\varepsilon)\|_{L^2(\mathcal{S}_{\varepsilon,r})},$$

which leads to (5.22). □

6 The unfolding operators

The classical unfolding operator $\mathcal{T}_\varepsilon$ is developed in [9, 10]. Here, we will use similar operators $\mathcal{T}^{ext}_\varepsilon$, $\mathcal{T}^{\mathcal{S}}_\varepsilon$, $\mathcal{T}^{b,\ell}_\varepsilon$ in the context of the domains Ω_ε, $\mathcal{S}_\varepsilon$ and $\mathcal{S}_{\varepsilon,r}$.

Definition 6.1 (Classical unfolding-operator)**.** *For a measurable function ϕ on Ω, the unfolding operator $\mathcal{T}_\varepsilon$ is defined as follows:*

$$\mathcal{T}_\varepsilon(\phi)(x,y) = \phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon y\Big) \quad \text{for a.e. } (x,y)\in\widetilde{\Omega}_\varepsilon\times Y,$$
$$\mathcal{T}_\varepsilon(\phi)(x,y) = 0 \qquad\qquad \text{for a.e. } (x,y)\in\big(\Omega\setminus\overline{\widetilde{\Omega}_\varepsilon}\big)\times Y.$$

Definition 6.2 (Unfolding-operator)**.** *For a measurable function ϕ on Ω_ε, the unfolding operator $\mathcal{T}^{ext}_\varepsilon$ is defined as follows:*

$$\mathcal{T}^{ext}_\varepsilon(\phi)(x,y) = \phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon y\Big) \quad \text{for a.e. } (x,y)\in\Omega_\varepsilon\times Y.$$

Lemma 6.1. *Let ϕ be in $L^p(\Omega_\varepsilon)$, $p\in[1,+\infty)$. One has*

$$\|\mathcal{T}^{ext}_\varepsilon(\phi)-\mathcal{T}_\varepsilon(\phi)\|_{L^p(\Omega\times Y)} \le \|\phi\|_{L^p(\Omega^{bl}_\varepsilon)} \tag{6.1}$$

where

$$\Omega^{bl}_\varepsilon \doteq \big\{x\in\Omega_\varepsilon \mid dist(x,\partial\Omega)\le\varepsilon\sqrt{3}\big\}.$$

Proof. Inequality (6.1) is an immediate consequence of the definitions of these operators. □

As a consequence of the above lemma, the properties of the operator $\mathcal{T}_\varepsilon^{ext}$ are similar to those of the classical unfolding operator $\mathcal{T}_\varepsilon$. For the main properties of the unfolding operator $\mathcal{T}_\varepsilon$, we refer the reader to [10, Chapter 1]. Below, we introduce two new unfolding operators. The first one is used for the centerlines of beams and the second one is used for the small beams (it concerns the reduction of dimension).

In the definitions below, $\varepsilon\Big[\frac{x}{\varepsilon}\Big]$ represents a macroscopic coordinate (the same coordinate for all the points in the cell $\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon Y$) while $\mathbf{S}$ is the coordinate of a point belonging to $\mathcal{S}$. Hence, $\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{S}$ represents the coordinate of a point belonging to $\mathcal{S}_\varepsilon$. In order to get a map $(x,\mathbf{S})\longmapsto\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{S}$ almost one to one, we have to restrict the set $\mathcal{S}$. This is why from now on, to introduce the unfolding operator, in lieu of $\mathcal{S}$ we consider the set

$$\mathcal{S}\cap[0,1)^3.$$

For simplicity we still refer to it as $\mathcal{S}$. The set of new nodes is always denoted $\mathcal{K}$ and the number of beams of $\mathcal{S}$ is still denoted m.

Definition 6.3 (Centerlines unfolding)**.** *For a measurable function ϕ on $\mathcal{S}_\varepsilon$, the unfolding operator $\mathcal{T}_\varepsilon^{\mathcal{S}}$ is defined as follows:*

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)(x,\mathbf{S})=\phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{S}\Big)\quad \text{for a.e. } (x,\mathbf{S})\in\Omega_\varepsilon\times\mathcal{S}.$$

Definition 6.4 (Beams unfolding)**.** *For a measurable function u on $\mathcal{S}_{\varepsilon,r}$, the unfolding operator $\mathcal{T}_\varepsilon^{b,\ell}$ is defined as follows ($\ell\in\{1,\dots,m\}$):*

$$\mathcal{T}_\varepsilon^{b,\ell}(u)(x,\widehat{S})=u\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell\Big)\quad \text{for a.e. } (x,\widehat{S})\in\Omega_\varepsilon\times(0,l_\ell)\times D,$$

where $\widehat{S}=(S_1,S_2,S_3)$, $\mathbf{A}^\ell$ is an extremity of the segment $\gamma_\ell\subset\mathcal{S}$ and $D=D_1$ is the disc of radius 1.

Let ϕ be measurable on $\mathcal{S}_\varepsilon$, one has

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)(x,\mathbf{S})=\phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{S}\Big)=\phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell\Big)=\mathcal{T}_\varepsilon^{b,\ell}(\phi)(x,\widehat{(S_1,0,0)})\quad\text{for a.e. } (x,S_1)\in\Omega_\varepsilon\times(0,l_\ell).$$

Lemma 6.2 (Properties of the operators $\mathcal{T}_\varepsilon^{\mathcal{S}}$ and $\mathcal{T}_\varepsilon^{b,\ell}$)**.**
For every $\phi\in L^1(\mathcal{S}_\varepsilon)$

$$\int_{\Omega_\varepsilon\times\mathcal{S}}\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)(x,\mathbf{S})\,d\mathbf{S}\,dx=\varepsilon^2\int_{\mathcal{S}_\varepsilon}\phi(x)\,dx.\tag{6.2}$$

For every $\phi\in L^2(\mathcal{S}_\varepsilon)$

$$\|\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)\|_{L^2(\Omega_\varepsilon\times\mathcal{S})}=\varepsilon\|\phi\|_{L^2(\mathcal{S}_\varepsilon)}.\tag{6.3}$$

For every ϕ in $H^1(\mathcal{S}_\varepsilon)$

$$\frac{\partial\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)}{\partial\mathbf{S}}(x,\mathbf{S})=\varepsilon\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi}{d\mathbf{s}}\Big)(x,\mathbf{S})\qquad\text{for a.e. } (x,\mathbf{S})\in\Omega_\varepsilon\times\mathcal{S}.\tag{6.4}$$

For every ψ in $L^2(\mathcal{S}_{\varepsilon,r})$

$$\big\|\mathcal{T}_\varepsilon^{b,\ell}(\psi)\big\|_{L^2(\Omega_\varepsilon\times\gamma_\ell\times D)}\le C\frac{\varepsilon}{r}\|\psi\|_{L^2(\mathcal{S}_{\varepsilon,r})}\qquad\text{for all }\ell\in\{1,\dots,m\}.\tag{6.5}$$

For every ψ in $L^1(\mathcal{S}_{\varepsilon,r})$

$$\Big|\sum_{\ell=1}^m\int_{\Omega_\varepsilon\times\gamma_\ell\times D}\frac{r^2}{\varepsilon^2}\mathcal{T}_\varepsilon^{b,\ell}(\psi)(x,\widehat{S})\,dx\,d\widehat{S}-\int_{\mathcal{S}_{\varepsilon,r}}\psi(x)\,dx\Big|\le C\|\psi\|_{L^1(\mathcal{J}_{\varepsilon,r})}.\tag{6.6}$$

The constant only depends on $\mathcal{S}$.
For every u in $H^1(\mathcal{S}_{\varepsilon,r})$ ($j\in\{2,3\}$ and $\ell\in\{1,\dots,m\}$)

$$\begin{aligned}&\varepsilon\mathcal{T}_\varepsilon^{b,\ell}(\nabla u)(x,\widehat{S})\cdot\mathbf{t}_1^\ell=\frac{\partial\mathcal{T}_\varepsilon^{b,\ell}(u)}{\partial S_1}(x,\widehat{S}),\\ &r\mathcal{T}_\varepsilon^{b,\ell}(\nabla u)(x,\widehat{S})\cdot\mathbf{t}_j^\ell=\frac{\partial\mathcal{T}_\varepsilon^{b,\ell}(u)}{\partial S_j}(x,\widehat{S}),\end{aligned}\qquad\text{for a.e. } (x,\widehat{S})\in\Omega_\varepsilon\times(0,l_\ell)\times D.\tag{6.7}$$

Proof. We prove (6.2) and (6.3). Let ϕ be in $L^1(\mathcal{S}_\varepsilon)$

$$\int_{\Omega_\varepsilon\times\mathcal{S}}\mathcal{T}_\varepsilon^S(\phi)(x,\mathbf{S})\,d\mathbf{S}\,dx=\sum_{\ell=1}^m\int_{\Omega_\varepsilon\times\gamma_\ell}\mathcal{T}_\varepsilon^S(\phi)(x,\mathbf{A}^\ell+S_1\mathbf{t}_1^\ell)\,dx\,dS_1=\sum_{\ell=1}^m\sum_{\xi\in\Xi_\varepsilon}|\varepsilon\xi+\varepsilon Y|\int_0^{l_\ell}\phi(\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon t)dt$$
$$=\sum_{\ell=1}^m\sum_{\xi\in\Xi_\varepsilon}\varepsilon^3\int_0^{l_\ell}\phi(\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon t)dt=\varepsilon^2\int_{\mathcal{S}_\varepsilon}\phi(x)\,dx.$$

We prove (6.6). For $u\in L^1(\mathcal{S}_{\varepsilon,r})$ we have

$$\begin{aligned}\int_{\Omega_\varepsilon\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(u)(x,\widehat{S})\,dx\,d\widehat{S}&=\sum_{\xi\in\Xi_\varepsilon}\int_{(\varepsilon\xi+\varepsilon Y)\times\gamma_\ell\times D}u\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big]+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell\Big)\,dx\,d\widehat{S}\\&=\sum_{\xi\in\Xi_\varepsilon}\int_{(\varepsilon\xi+\varepsilon Y)\times\gamma_\ell\times D}u\Big(\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell\Big)\,dx\,d\widehat{S}\\&=\sum_{\xi\in\Xi_\varepsilon}|\varepsilon\xi+\varepsilon Y|\int_{\gamma_\ell\times D}u\Big(\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell\Big)\,d\widehat{S}\\&=\varepsilon^3\sum_{\xi\in\Xi_\varepsilon}\int_{\gamma_\ell\times D}u\Big(\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell\Big)\,d\widehat{S}.\end{aligned}$$

Now, replacing $\varepsilon\xi+\varepsilon\mathbf{A}^\ell+\varepsilon S_1\mathbf{t}_1^\ell+rS_2\mathbf{t}_2^\ell+rS_3\mathbf{t}_3^\ell$ by x and taking into account that the matrix $(\mathbf{t}_1^\ell|\mathbf{t}_2^\ell|\mathbf{t}_3^\ell)$ belongs to $SO(3)$, we obtain

$$\int_{\Omega_\varepsilon\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(u)(x,\widehat{S})\,dx\,d\widehat{S}=\frac{\varepsilon^2}{r^2}\sum_{\xi\in\Xi_\varepsilon}\int_{(\varepsilon\xi+\varepsilon\mathcal{P}_{\ell,r/\varepsilon})}u(x)\,dx=\frac{\varepsilon^2}{r^2}\sum_{\xi\in\Xi_\varepsilon}\int_{\mathcal{P}^\xi_{\varepsilon\ell,r}}u(x)\,dx$$

and (6.6) follows.
Properties (6.4)-(6.7) are direct consequences of the definitions of the unfolding operators. □

Corollary 6.1. *For every ϕ in $L^2(\mathcal{S}_\varepsilon)$, $\ell\in\{1,\dots,m\}$*

$$\|\mathcal{T}_\varepsilon^{b,\ell}(\phi)\|_{L^2(\Omega\times\gamma_\ell\times D)}\le C\varepsilon\|\phi\|_{L^2(\mathcal{S}_\varepsilon)}.\tag{6.8}$$

From now on, every function belonging to $L^p(\Omega)$ ($p\in[1,+\infty]$) will be extended by 0 in $\Omega_\varepsilon\setminus\overline{\Omega}$.

Denote $Q_1(Y)$ the subspace of $W^{1,\infty}(Y)$ containing the functions which are the Q_1 interpolations of their values at the vertices of the parallelotope $\overline{Y}$.

Lemma 6.3. *For every Φ in $W^{1,\infty}(\Omega_\varepsilon)$ satisfying*

$$\mathcal{T}_\varepsilon^{ext}(\Phi)\in L^\infty(\Omega;Q_1(Y)).\tag{6.9}$$

Then $\Phi_{|\mathcal{S}_\varepsilon}$ belongs to $W^{1,\infty}(\mathcal{S}_\varepsilon)$ and it satisfies

$$\begin{aligned}&\big\|\Phi_{|\mathcal{S}_\varepsilon}\big\|_{L^2(\mathcal{S}_\varepsilon)}\le\frac{C}{\varepsilon}\|\Phi\|_{L^2(\Omega_\varepsilon)},\\&\frac{d\Phi_{|\mathcal{S}_\varepsilon}}{d\mathbf{s}}=\nabla\Phi_\varepsilon\cdot\mathbf{t}_1\quad a.e.\ in\ \mathcal{S}_\varepsilon\quad and\quad\Big\|\frac{d\Phi_{|\mathcal{S}_\varepsilon}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)}\le\frac{C}{\varepsilon}\|\nabla\Phi\|_{L^2(\Omega_\varepsilon)}.\end{aligned}\tag{6.10}$$

Let $\{\Phi_\varepsilon\}_\varepsilon$ be a sequence of functions belonging to $W^{1,\infty}(\Omega_\varepsilon)$ satisfying (6.9) and

$$\|\Phi_\varepsilon\|_{L^2(\Omega_\varepsilon)}\le C\tag{6.11}$$

then, up to a subsequence of $\{\varepsilon\}$, there exists $\Phi\in L^2(\Omega)$ such that

$$\begin{aligned}&\Phi_\varepsilon\rightharpoonup\Phi\quad weakly\ in\ L^2(\Omega),\\&\mathcal{T}_\varepsilon^{ext}(\Phi_\varepsilon)\rightharpoonup\Phi\quad weakly\ in\ L^2(\Omega;Q_1(Y)),\\&\mathcal{T}_\varepsilon^{ext}(\Phi_\varepsilon)_{|\Omega\times\mathcal{S}}=\mathcal{T}_\varepsilon^{\mathcal{S}}(\Phi_\varepsilon)\rightharpoonup\Phi\quad weakly\ in\ L^2(\Omega;H^1(\mathcal{S})).\end{aligned}\tag{6.12}$$

Moreover, if one also has

$$\|\nabla\Phi_\varepsilon\|_{L^2(\Omega_\varepsilon)} \leq C$$

then Φ *belongs to* $H^1(\Omega)$ *and*

$$\begin{aligned}
&\Phi_\varepsilon \rightharpoonup \Phi \quad \textit{weakly in } H^1(\Omega),\\
&\mathcal{T}_\varepsilon^{ext}(\nabla\Phi_\varepsilon) \rightharpoonup \nabla\Phi \quad \textit{weakly in } L^2(\Omega\times Y)^3,\\
&\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\Phi_\varepsilon}{d\mathbf{s}}\Big) = \mathcal{T}_\varepsilon^{ext}(\nabla\Phi_\varepsilon\cdot\mathbf{t}_1)_{|\Omega\times\mathcal{S}} \rightharpoonup \nabla\Phi\cdot\mathbf{t}_1 \quad \textit{weakly in } L^2(\Omega\times\mathcal{S}).
\end{aligned} \tag{6.13}$$

Proof. The proof is given in Appendix C. □

First convergence results for sequences in $H^1(\mathcal{S}_\varepsilon)$.

Lemma 6.4. *Let* $\{\phi_\varepsilon\}_\varepsilon$ *be a sequence of functions belonging to* $H^1(\mathcal{S}_\varepsilon)$ *satisfying*

$$\|\phi_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon)} + \varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}.$$

Then, up to a subsequence, there exists $\widehat{\phi}\in L^2(\Omega; H^1_{per}(\mathcal{S}))$ *such that*

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon) \rightharpoonup \widehat{\phi} \quad \textit{weakly in } \ L^2(\Omega; H^1(\mathcal{S})). \tag{6.14}$$

If we only have

$$\|\phi_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon\cap\Omega_\varepsilon^{int})} + \varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon\cap\Omega_\varepsilon^{int})} \leq \frac{C}{\varepsilon},$$

then, up to a subsequence, there exists $\widehat{\phi}\in L^2(\Omega; H^1_{per}(\mathcal{S}))$ *such that*

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \rightharpoonup \widehat{\phi} \quad \textit{weakly in } \ L^2(\Omega; H^1(\mathcal{S})). \tag{6.15}$$

Proof. The proof is postponed in Appendix C. □

Definition 6.5. *The local average operator* $\mathcal{M}_\varepsilon^*$ *is defined from* $L^2(\mathcal{S}_\varepsilon)$ *to* $L^2(\Omega_\varepsilon)$ *as*

$$\mathcal{M}_\varepsilon^*(\phi)(x) = \frac{1}{|\mathcal{S}|}\int_{\mathcal{S}} \mathcal{T}_\varepsilon^{\mathcal{S}}(\phi)(x,\mathbf{S})\,d\mathbf{S}, \qquad \textit{for a.e. } x\in\Omega_\varepsilon.$$

By convention the value of $\mathcal{M}_\varepsilon^*(\phi)$ *on the cell* $\varepsilon(\xi+Y)$ *is simply denoted* $\mathcal{M}_\varepsilon^*(\phi)(\varepsilon\xi)$.

A second lemma for sequences in $H^1(\mathcal{S}_\varepsilon)$.

Lemma 6.5. *Let* $\{\phi_\varepsilon\}_\varepsilon$ *be a sequence of functions belonging to* $H^1(\mathcal{S}_\varepsilon)$ *satisfying*

$$\|\phi_\varepsilon\|_{H^1(\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}. \tag{6.16}$$

Then, up to a subsequence, there exists $(\Phi,\widehat{\phi})\in H^1(\Omega)\times L^2(\Omega; H^1_{per}(\mathcal{S}))$ *such that*

$$\begin{aligned}
&\mathcal{T}_\varepsilon^{\mathcal{S}}\big(\phi_\varepsilon\big)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \longrightarrow \Phi \quad \textit{strongly in } \ L^2(\Omega; H^1(\mathcal{S})),\\
&\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \rightharpoonup \nabla\Phi\cdot\mathbf{t}_1 + \frac{\partial\widehat{\phi}}{\partial\mathbf{S}} \quad \textit{weakly in } \ L^2(\Omega\times\mathcal{S}).
\end{aligned} \tag{6.17}$$

Proof. The proof is postponed in Appendix C. □

Denote

$$H^1_\Gamma(\Omega) \doteq \big\{\phi\in H^1(\Omega)\mid \phi = 0 \text{ on } \Gamma\big\}.$$

Corollary 6.2. *Let $\{\phi_\varepsilon\}_\varepsilon$ be a sequence of functions belonging to $H^1(\mathcal{S}_\varepsilon)^3 \cap \mathbf{V}_{\varepsilon,r}$ and satisfying the following*

$$\|\phi_\varepsilon\|_{H^1(\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}.$$

Then, up to a subsequence, there exists $(\Phi, \widehat{\phi}) \in H^1_\Gamma(\Omega)^3 \times L^2(\Omega; H^1_{per}(\mathcal{S}))^3$ such that

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon'^{int}\times\mathcal{S}} \longrightarrow \Phi \quad \textit{strongly in } \ L^2(\Omega; H^1(\mathcal{S}))^3,$$
$$\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big)\mathbf{1}_{\widehat{\Omega}_\varepsilon'^{int}\times\mathcal{S}} \rightharpoonup \nabla\Phi\cdot\mathbf{t}_1 + \frac{\partial\widehat{\phi}}{\partial\mathbf{S}} \quad \textit{weakly in } \ L^2(\Omega\times\mathcal{S})^3.$$

Proof. Since $\{\phi_\varepsilon\}_\varepsilon$ belongs to $V_{\varepsilon,r}$, these functions equal to 0 in $\mathcal{S}'_\varepsilon \setminus \mathcal{S}_\varepsilon$. Applying Lemma 6.5 with $\mathcal{S}'_\varepsilon$ instead $\mathcal{S}_\varepsilon$ and with Ω' instead Ω give the result. □

7 Asymptotic behaviors

7.1 Asymptotic behavior of a sequence of displacements

From now on, we assume that r is a function of ε satisfying the following conditions:

$$\lim_{\varepsilon\to 0}\frac{r}{\varepsilon} = 0, \qquad \lim_{\varepsilon\to 0}\frac{r}{\varepsilon^2} = \kappa \in [0,+\infty]. \tag{7.1}$$

In addition, every field appearing in the decomposition introduced in the previous sections will be denoted with only the index ε.

In this section we consider a sequence $\{u_\varepsilon\}_\varepsilon$ of displacements belonging to $\mathbf{V}_{\varepsilon,r}$ and satisfying

$$\|e(u_\varepsilon)\|_{L^2(\mathcal{S}_{\varepsilon,r})} \leq C.$$

Theorem 7.1. *For a subsequence of $\{\varepsilon\}$, still denoted $\{\varepsilon\}$, one has*
(i) there exist $\mathcal{U} \in H^1_\Gamma(\Omega)^3$, $\overline{\mathcal{U}} \in L^2(\Omega; H^1_{per}(\mathcal{S}))^3$ such that $S \longmapsto \overline{\mathcal{U}}(\cdot,S)\wedge\mathbf{t}_1$ is an affine function on every segment of $\mathcal{S}$ and the following convergences hold:

$$\begin{aligned}
&\frac{r}{\varepsilon}\boldsymbol{\mathcal{U}}_\varepsilon\mathbf{1}_{\Omega_\varepsilon^{int}} \rightharpoonup \mathcal{U} \ \textit{ weakly in } \ L^2(\Omega)^3,\\
&\frac{r}{\varepsilon}\nabla\boldsymbol{\mathcal{U}}_\varepsilon\mathbf{1}_{\Omega_\varepsilon^{int}} \rightharpoonup \nabla\mathcal{U} \ \textit{ weakly in } \ L^2(\Omega)^9,\\
&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\big(\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1\big) \rightharpoonup \mathcal{U} \quad \textit{weakly in} \quad L^2(\Omega; H^1(\mathcal{S}))^3,\\
&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\big(\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1\big)\Big) \rightharpoonup \nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}} \quad \textit{weakly in} \quad L^2(\Omega\times\mathcal{S})^3,\\
&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\mathcal{U}_\varepsilon}{d\mathbf{s}}\Big)\cdot\mathbf{t}_1 \rightharpoonup (e(\mathcal{U})\,\mathbf{t}_1)\cdot\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 \quad \textit{weakly in} \quad L^2(\Omega\times\mathcal{S}),
\end{aligned} \tag{7.2}$$

where $e(\mathcal{U})$ is the symmetric gradient of the displacement $\mathcal{U}$
(ii) there exists $\widehat{\mathcal{U}} \in L^2(\Omega; H^1_{per}(\mathcal{S}))^3$ such that $\widehat{\mathcal{U}}_{|\gamma_\ell} \in L^2(\Omega; H^1_0(\gamma_\ell)\cap H^2(\gamma_\ell))^3$, $\widehat{\mathcal{U}}_{|\gamma_\ell}\cdot\mathbf{t}_1^\ell = 0$, $\ell\in\{1,\dots,m\}$ and

$$\frac{r^2}{\varepsilon^3}\mathcal{T}_\varepsilon^{\mathcal{S}}\big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1\big) \rightharpoonup \widehat{\mathcal{U}} \quad \textit{weakly in} \quad L^2(\Omega; H^1(\mathcal{S}))^3, \tag{7.3}$$

(iii) there exists $\mathcal{Z} \in L^2(\Omega\times\mathcal{S})^3$ such that

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\mathcal{U}_\varepsilon}{d\mathbf{s}} - \mathcal{R}_\varepsilon\wedge\mathbf{t}_1\Big) \rightharpoonup \nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}} + \mathcal{Z} \quad \textit{weakly in} \quad L^2(\Omega\times\mathcal{S})^3, \tag{7.4}$$

(iv) there exists $\widehat{\mathcal{R}} \in L^2(\Omega; H^1_{per}(\mathcal{S}))^3$ such that

$$\frac{r^2}{\varepsilon^2}\mathcal{T}_\varepsilon^{\mathcal{S}}(\mathcal{R}_\varepsilon) \rightharpoonup \widehat{\mathcal{R}} \quad \textit{weakly in} \quad L^2(\Omega; H^1(\mathcal{S}))^3 \tag{7.5}$$

and

$$\widehat{\mathcal{R}} \wedge \mathbf{t}_1 = \frac{\partial \widehat{\mathcal{U}}}{\partial \mathbf{S}}, \tag{7.6}$$

(v) there exists $\overline{u} \in L^2(\Omega \times \mathcal{S}; H^1(D))^3$ *such that*

$$\begin{aligned} &\frac{1}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon) \rightharpoonup \overline{u} \quad \textit{weakly in} \quad L^2(\Omega \times \gamma_\ell; H^1(D))^3, \\ &\frac{r}{\varepsilon^2}\frac{\partial}{\partial S_1}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon) \rightharpoonup 0 \quad \textit{weakly in} \quad L^2(\Omega \times \gamma_\ell \times D)^3. \end{aligned} \tag{7.7}$$

Proof. Below, every convergence is up to a subsequence of $\{\varepsilon\}$ still denoted $\{\varepsilon\}$.
(i) From Lemma A.1 and Proposition 5.3 we have the following estimates:

$$\frac{r}{\varepsilon}\|\boldsymbol{\mathcal{U}}_\varepsilon\|_{H^1(\Omega_\varepsilon'^{int})} \leq C. \tag{7.8}$$

Lemma 5.1 in [16] gives a field $\mathcal{U} \in H^1_\Gamma(\Omega)^3$ such that $(7.2)_{1,2}$ hold.
From the estimates (5.10) and (A.2) one obtains

$$\|\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\|_{H^1(\mathcal{S}_\varepsilon)} \leq \frac{C}{r}.$$

Hence, the convergences $(7.2)_{3,4}$ are the consequences of Corollary 6.2.
Since

$$\frac{d\mathcal{U}_\varepsilon}{d\mathbf{s}} \cdot \mathbf{t}_1 = \frac{d}{d\mathbf{s}}\big(\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\big) \cdot \mathbf{t}_1,$$

the convergence $(7.2)_5$ holds (observe that $(\nabla \mathcal{U}\, \mathbf{t}_1) \cdot \mathbf{t}_1 = (e(\mathcal{U})\, \mathbf{t}_1) \cdot \mathbf{t}_1$).
(ii) From (5.10), (5.22), (6.1) and the fact that by construction $\overline{\mathcal{U}}_{\varepsilon|\gamma_\ell}(0) = \overline{\mathcal{U}}_{\varepsilon|\gamma_\ell}(\varepsilon l_\ell) = 0$, we obtain

$$\frac{r^2}{\varepsilon^3}\big\|\mathcal{T}_\varepsilon^{\mathcal{S}}\big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\big)\big\|_{L^2(\Omega; H^1_0\gamma_\ell)} \leq C.$$

Thus, up to a subsequence, there exists $\widehat{\mathcal{U}} \in L^2(\Omega; H^1(\mathcal{S}))^3$ such that $\widehat{\mathcal{U}}_{|\gamma_\ell} \in L^2(\Omega; H^1_0(\gamma_\ell))^3$, $\widehat{\mathcal{U}}_{|\gamma_\ell} \cdot \mathbf{t}_1^\ell = 0$, $\ell \in \{1, \ldots, m\}$ and convergence $(7.3)_1$ holds.
(iii) Estimates $(5.9)_4$-(5.10) and (6.2) yield

$$\Big\|\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\big) - \mathcal{R}_\varepsilon \wedge \mathbf{t}_1\Big)\Big\|_{L^2(\Omega \times \mathcal{S})} \leq C\frac{\varepsilon}{r}. \tag{7.9}$$

Then, there exists a field $\mathcal{Z} \in L^2(\Omega \times \mathcal{S})^3$ such that

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\big) - \mathcal{R}_\varepsilon \wedge \mathbf{t}_1\Big) \rightharpoonup \mathcal{Z} \quad \text{weakly in} \quad L^2(\Omega \times \mathcal{S})^3$$

and by $(7.2)_4$ we have

$$\begin{aligned} \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\mathcal{U}_\varepsilon}{d\mathbf{s}} - \mathcal{R}_\varepsilon \wedge \mathbf{t}_1\Big) &= \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\Big(\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\Big)\Big) + \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\Big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\Big) - \mathcal{R}_\varepsilon \wedge \mathbf{t}_1\Big) \\ &\rightharpoonup \nabla \mathcal{U}\, \mathbf{t}_1 + \frac{\partial \overline{\mathcal{U}}}{\partial \mathbf{S}} + \mathcal{Z} \quad \text{weakly in} \quad L^2(\Omega \times \mathcal{S})^3. \end{aligned}$$

(iv) Estimate $(5.18)_2$ gives

$$\|\mathcal{R}_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon)} + \varepsilon\Big\|\frac{d\mathcal{R}_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq C\frac{\varepsilon}{r^2}.$$

Thus, up to a subsequence, there exists a function $\widehat{\mathcal{R}} \in L^2(\Omega; H^1_{per}(\mathcal{S}))^3$ (see Lemma 6.4) such that (7.5) holds.
On the one hand, from (7.9) we have

$$\frac{r^2}{\varepsilon^2}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{s}}\big(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon \cdot \mathbf{t}_1)\mathbf{t}_1\big) - \mathcal{R}_\varepsilon \wedge \mathbf{t}_1\Big) \longrightarrow 0 \quad \text{strongly in} \quad L^2(\Omega \times \mathcal{S})^3.$$

On the other hand from convergences $(7.3)_1$, (7.5) we obtain

$$\frac{r^2}{\varepsilon^2}\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d}{d\mathbf{S}}(\overline{\mathcal{U}}_\varepsilon-(\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1)-\mathcal{R}_\varepsilon\wedge\mathbf{t}_1\Big)\rightharpoonup\frac{\partial\widehat{\mathcal{U}}}{d\mathbf{S}}-\widehat{\mathcal{R}}\wedge\mathbf{t}_1\quad\text{weakly in}\quad L^2(\Omega\times\mathcal{S})^3.$$

Hence, we obtain (7.6) and

$$\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\wedge\mathbf{t}_1=\frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\ \text{ a.e. in }\ \Omega\times\mathcal{S}. \tag{7.10}$$

Then $\widehat{\mathcal{U}}_{|\gamma_\ell}\in L^2(\Omega;H_0^1(\gamma_\ell)\cap H^2(\gamma_\ell))^3$.
(v) Taking into account $(5.9)_{1,2}$, $(6.7)_2$ and (6.5) for $j=2,3$, $\ell\in\{1,\dots,m\}$, we have

$$\|\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon)\|_{L^2(\Omega\times\gamma_\ell\times D)}+\Big\|\frac{\partial}{\partial S_j}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon)\Big\|_{L^2(\Omega\times\gamma_\ell\times D)}\le C\varepsilon.$$

Hence, up to a subsequence, there exists $\overline{u}\in L^2(\Omega\times\mathcal{S};H^1(D))^3$ such that $(7.7)_1$ holds.
In order to show convergence $(7.7)_2$, note that from $(5.9)_2$, $(6.7)_1$ and (6.5) it follows

$$\frac{r}{\varepsilon^2}\Big\|\frac{\partial}{\partial S_1}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon)\Big\|_{L^2(\Omega\times\gamma_\ell\times D)}\le C.$$

Therefore, convergence $(7.7)_2$ is proved, since

$$\frac{r}{\varepsilon^2}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon)\longrightarrow 0\quad\text{strongly in}\quad L^2(\Omega\times\gamma_\ell;H^1(D))^3.$$ □

Remark 7.1. *Due to* (4.2)*, the warping* $\overline{u}$ *satisfies*

$$\begin{aligned}&\int_D\overline{u}(\cdot,S_2,S_3)\,dS_2dS_3=0,\\&\int_D\overline{u}(\cdot,S_2,S_3)\wedge(S_2\mathbf{t}_2^\ell+S_3\mathbf{t}_3^\ell)\,dS_2dS_3=0,\end{aligned}\qquad a.e.\ in\ \ \Omega\times\gamma_\ell,\quad\forall\ell\in\{1,\dots,m\}. \tag{7.11}$$

Denote

$$\boldsymbol{\mathcal{D}}_{Ex}\doteq\Big\{\overline{\mathcal{A}}\in H^1_{per,0}(\mathcal{S})^3\ \mid\ \overline{\mathcal{A}}\wedge\mathbf{t}_1\text{ is an affine function on every segment }\gamma_\ell,\ \ell\in\{1,\dots,m\}\Big\},$$

$$\boldsymbol{\mathcal{D}}_{In}\doteq\Big\{(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\in H^1_{per}(\mathcal{S})^3\times H^1_{per}(\mathcal{S})^3\ \mid\ \frac{d\widehat{\mathcal{A}}}{d\mathbf{S}}=\widehat{\mathcal{B}}\wedge\mathbf{t}_1,\ \ \widehat{\mathcal{A}}=0\text{ on all the nodes of }\mathcal{S}\Big\}.$$

The field $\overline{\mathcal{U}}$ is in $L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex})$ while the pair $(\widehat{\mathcal{U}},\widehat{\mathcal{R}})$ belongs to $L^2(\Omega;\boldsymbol{\mathcal{D}}_{In})$. It worth to notice that a field $\overline{\mathcal{A}}$ belonging to $H^1_{per,0}(\mathcal{S})^3$ is a local extensional displacement if and only if

$$\int_{\mathcal{S}}\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\frac{d\widehat{\mathcal{A}}}{d\mathbf{S}}d\mathbf{S}=0$$

for all $\widehat{\mathcal{A}}\in H^1_{per}(\mathcal{S})^3$ which is the first component of an element belonging to $\boldsymbol{\mathcal{D}}_{In}$.
We endow $\boldsymbol{\mathcal{D}}_{Ex}$ (resp. $\boldsymbol{\mathcal{D}}_{In}$) with the semi-norm

$$\|\overline{\mathcal{A}}\|_{\mathcal{S}}\doteq\Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})},\qquad(\text{resp.}\quad\|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{\boldsymbol{\mathcal{D}}_{In}}\doteq\Big\|\frac{d\widehat{\mathcal{B}}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})}).$$

Lemma 7.1. *On* $\boldsymbol{\mathcal{D}}_{Ex}$ *the semi-norm* $\|\cdot\|_{\mathcal{S}}$ *is a norm equivalent to the norm of* $H^1(\mathcal{S})^3$*. On* $\boldsymbol{\mathcal{D}}_{In}$ *the semi-norm* $\|(\cdot,\cdot)\|_{\boldsymbol{\mathcal{D}}_{In}}$ *is a norm equivalent to the norm of* $H^1(\mathcal{S})^3\times H^1(\mathcal{S})^3$*.*

Proof. The proof is given in Appendix D. □

7.2 Asymptotic behavior of the strain tensor

For every $\mathcal{V} \in H^1_\Gamma(\Omega)^3$, $(\overline{\mathcal{V}}, \widehat{\mathcal{V}}, \widehat{\mathcal{B}}) \in L^2(\Omega; \boldsymbol{\mathcal{D}}_{Ex} \times \boldsymbol{\mathcal{D}}_{In})$ and $\widetilde{v} \in L^2(\Omega \times \mathcal{S}; H^1(D))^3$ we define the symmetric tensors $\mathcal{E}$, $\mathcal{E}_\mathcal{S}$, $\mathcal{E}_D$ by

$$\mathcal{E}(\mathcal{V}) \doteq \begin{pmatrix} \big(e(\mathcal{V})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \qquad \mathcal{E}_S(\overline{\mathcal{V}}, \widehat{\mathcal{V}}, \widehat{\mathcal{B}}) \doteq \begin{pmatrix} \frac{\partial \overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_1 - \frac{\partial^2 \widehat{\mathcal{V}}}{\partial S_1^2}\cdot\Big(S_2\,\mathbf{t}_2 + S_3\,\mathbf{t}_3\Big) & * & * \\ -\frac{S_3}{2}\frac{\partial \widehat{\mathcal{B}}}{\partial S_1}\cdot\mathbf{t}_1 & 0 & 0 \\ \frac{S_2}{2}\frac{\partial \widehat{\mathcal{B}}}{\partial S_1}\cdot\mathbf{t}_1 & 0 & 0 \end{pmatrix},$$

$$\mathcal{E}_D(\widetilde{v}) \doteq \begin{pmatrix} 0 & \frac{1}{2}\frac{\partial \widetilde{v}}{\partial S_2}\cdot\mathbf{t}_1 & \frac{1}{2}\frac{\partial \widetilde{v}}{\partial S_3}\cdot\mathbf{t}_1 \\ * & \frac{\partial \widetilde{v}}{\partial S_2}\cdot\mathbf{t}_2 & \frac{1}{2}\frac{\partial \widetilde{v}}{\partial S_3}\cdot\mathbf{t}_2 + \frac{1}{2}\frac{\partial \widetilde{v}}{\partial S_2}\cdot\mathbf{t}_3 \\ * & * & \frac{\partial \widetilde{v}}{\partial S_3}\cdot\mathbf{t}_3 \end{pmatrix} \qquad \text{a.e. in } \Omega\times\mathcal{S}\times D.$$

Theorem 7.2. *Let u_ε be the solution to (5.8). There exist a subsequence of $\{\varepsilon\}$, still denoted $\{\varepsilon\}$, and $\mathcal{U} \in H^1_\Gamma(\Omega)^3$, $(\overline{\mathcal{U}}, \widehat{\mathcal{U}}, \widehat{\mathcal{R}}) \in L^2(\Omega; \boldsymbol{\mathcal{D}}_{Ex} \times \boldsymbol{\mathcal{D}}_{In})$ and $\widetilde{u} \in L^2(\Omega \times \mathcal{S}; H^1(D))^3$ such that the following convergences hold ($\ell \in \{1, \dots, m\}$):*

$$\begin{cases} \dfrac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(u_\varepsilon) \rightharpoonup \mathcal{U} + \dfrac{1}{\kappa}\widehat{\mathcal{U}} \ \text{ weakly in } L^2(\Omega\times\gamma_\ell; H^1(D))^3, & \text{if } \kappa \in (0, +\infty], \\ \dfrac{r^2}{\varepsilon^3}\mathcal{T}_\varepsilon^{b,\ell}(u_\varepsilon) \rightharpoonup \widehat{\mathcal{U}} \ \text{ weakly in } L^2(\Omega\times\gamma_\ell; H^1(D))^3, & \text{if } \kappa = 0 \end{cases} \tag{7.12}$$

and

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon)) \rightharpoonup \mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}}, \widehat{\mathcal{U}}, \widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u}) \quad \text{weakly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times 3}. \tag{7.13}$$

Proof. Below, we give the asymptotic behavior of the sequence $\{\mathcal{T}_\varepsilon^{b,\ell}(u_\varepsilon)\}$ as $\varepsilon \to 0$ and $r/\varepsilon \to 0$. One has

$$\mathcal{T}_\varepsilon^{b,\ell}(u_\varepsilon) = \mathcal{T}_\varepsilon^{b,\ell}(U_\varepsilon^e) + \mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon).$$

From $(7.7)_1$ we have ($\ell \in \{1, \dots, m\}$)

$$\frac{1}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(\overline{u}_\varepsilon) \rightharpoonup \overline{u} \quad \text{weakly in} \quad L^2(\Omega\times\gamma_\ell; H^1(D))^3.$$

From Definition 4.1 we have ($\ell \in \{1, \dots, m\}$)

$$\mathcal{T}_\varepsilon^{b,\ell}(U_\varepsilon^e) = \mathcal{T}_\varepsilon^{\mathcal{S}}(\mathcal{U}_\varepsilon^h + (\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1) + \mathcal{T}_\varepsilon^{\mathcal{S}}(\overline{\mathcal{U}}_\varepsilon - (\overline{\mathcal{U}}_\varepsilon\cdot\mathbf{t}_1)\mathbf{t}_1) + r\mathcal{T}_\varepsilon^{\mathcal{S}}(\mathcal{R}_\varepsilon) \wedge (S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell), \ \text{ a.e. in } \Omega\times\gamma_\ell\times D.$$

The convergences $(7.2)_3$, (7.3), (7.5) yield

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(U_\varepsilon^e) \longrightarrow \mathcal{U} + \frac{1}{\kappa}\widehat{\mathcal{U}} \ \text{ weakly in } L^2(\Omega\times\gamma_\ell; H^1(D))^3, \quad \text{if } \kappa \in (0, +\infty].$$

if $\kappa = 0$, from (7.3) we obtain

$$\frac{r^2}{\varepsilon^3}\mathcal{T}_\varepsilon^{b,\ell}(U_\varepsilon^e) \to \widehat{\mathcal{U}} \ \text{ weakly in } L^2(\Omega\times\gamma_\ell; H^1(D))^3.$$

Hence, the convergences (7.12) hold.

Now we consider the asymptotic behavior of the strain tensors $\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon))$

$$\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon)) = \mathcal{T}_\varepsilon^{b,\ell}(e_s(\overline{u}_\varepsilon)) + \mathcal{T}_\varepsilon^{b,\ell}(e_s(U_e^e)).$$

From (7.7), we obtain ($\ell \in [1, ..., m]$)

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(\overline{u}_\varepsilon)) \rightharpoonup \mathcal{E}_D(\overline{u}) \quad \text{weakly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times 3}.$$

Next from the convergences $(7.2)_4$, $(7.3)_2$, (7.5) and (7.6) we obtain

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(U_\varepsilon^e)) \rightharpoonup \begin{pmatrix} \Big(\nabla\mathcal{U}\,\mathbf{t}_1^\ell + \frac{\partial \overline{\mathcal{U}}}{\partial S_1}\Big)\cdot\mathbf{t}_1^\ell - \frac{\partial^2 \widehat{\mathcal{U}}}{\partial S_1^2}\cdot\Big(S_2\cdot\mathbf{t}_2^\ell + S_3\cdot\mathbf{t}_3^\ell\Big) & * & * \\ \frac{1}{2}\big(\nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial \overline{\mathcal{U}}}{\partial S_1} + \mathcal{Z}\big)\cdot\mathbf{t}_2^\ell - \frac{S_3}{2}\frac{\partial \widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1^\ell & 0 & 0 \\ \frac{1}{2}\big(\nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial \overline{\mathcal{U}}}{\partial S_1} + \mathcal{Z}\big)\cdot\mathbf{t}_3^\ell + \frac{S_2}{2}\frac{\partial \widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1^\ell & 0 & 0 \end{pmatrix} \quad \text{weakly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times 3}.$$

We set

$$\widetilde{u} = \overline{u} + S_2\Big(\big(\nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial S_1} + \mathcal{Z}\big)\cdot\mathbf{t}_2\Big)\,\mathbf{t}_1 + S_3\Big(\big(\nabla\mathcal{U}\,\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial S_1} + \mathcal{Z}\big)\cdot\mathbf{t}_3\Big)\,\mathbf{t}_1 \qquad \text{a.e. in } \Omega\times\mathcal{S}\times D.$$

Hence, one has

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon)) \rightharpoonup \mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u}) \quad \text{weakly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times3}$$

and (7.13) holds. □

Denote

$$\begin{aligned}\boldsymbol{\mathcal{D}}_w = \Big\{(\widetilde{w}_1,\widetilde{w}_2,\widetilde{w}_3)\in H^1(D)^3 \mid & \int_D \big(S_3\widetilde{w}_2(S_2,S_3) - S_2\widetilde{w}_3(S_2,S_3)\big)\,dS_2dS_3 = 0,\\ & \int_D \widetilde{w}_i(S_2,S_3)\,dS_2dS_3 = 0,\ i\in\{1,2,3\}\Big\}.\end{aligned} \tag{7.14}$$

Thanks to the conditions (7.11) satisfied by $\overline{u}$ and the definition of $\widetilde{u}$, one obtains

$$\widetilde{u} = (\widetilde{u}\cdot\mathbf{t}_1)\mathbf{t}_1 + (\widetilde{u}\cdot\mathbf{t}_2)\mathbf{t}_2 + (\widetilde{u}\cdot\mathbf{t}_3)\mathbf{t}_3 \text{ is such that } \big(\widetilde{u}\cdot\mathbf{t}_1,\widetilde{u}\cdot\mathbf{t}_2,\widetilde{u}\cdot\mathbf{t}_3\big)\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w). \tag{7.15}$$

For the sake of simplicity, if $\widetilde{v}$ belongs to $L^2(\Omega\times\mathcal{S};H^1(D)^3)$ and is such that

$$\widetilde{v} = (\widetilde{v}\cdot\mathbf{t}_1)\mathbf{t}_1 + (\widetilde{v}\cdot\mathbf{t}_2)\mathbf{t}_2 + (\widetilde{v}\cdot\mathbf{t}_3)\mathbf{t}_3 \quad \textit{satisfies} \quad \big(\widetilde{v}\cdot\mathbf{t}_1,\widetilde{v}\cdot\mathbf{t}_2,\widetilde{v}\cdot\mathbf{t}_3\big)\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w)$$

we will write that $\widetilde{v}$ belongs to $L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w)$.

8 The limit unfolded problem

To obtain the limit unfolded problem, we will choose test displacements v in $\mathbf{V}_{\varepsilon,r}$ which vanish in the junction domain $\mathcal{J}_{\varepsilon,r}$ or which are equal to rigid displacements in $\mathcal{J}_{\varepsilon,r}$. In doing so, we will have

$$\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(v)\,dx = \sum_{\ell=1}^m\frac{r^2}{\varepsilon^2}\int_{\Omega\times\gamma_\ell\times D} a_{ijkl}\mathcal{T}_\varepsilon^{b,\ell}(e_{s,ij}(u_\varepsilon))\,\mathcal{T}_\varepsilon^{b,\ell}(e_{s,kl}(v))\,dxd\widehat{S}.$$

The step-by-step construction of the unfolded limit problem (8.12) is considered in Lemmas 8.1, 8.2, 8.3.

Lemma 8.1 (The limit problem involving the limit warping)**.** *For every $\ell\in\{1,\dots,m\}$ one has*

$$\int_{\Omega\times\gamma_\ell\times D} a_{ijkl}\left(\mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\right)_{ij}\left(\mathcal{E}_D(\widetilde{v})\right)_{kl}\,d\widehat{S}\,dx = 0, \qquad \forall\,\widetilde{v}\in L^2(\Omega\times\gamma_\ell;H^1(D))^3. \tag{8.1}$$

Proof. Set

$$\begin{aligned}&\widetilde{v}_{\varepsilon,r}(x) = \varepsilon W(\varepsilon\xi+\varepsilon A^\ell)\,V\Big(\frac{s_1}{\varepsilon}\Big)\varphi\Big(\frac{s_2}{r},\frac{s_3}{r}\Big)\\ &\text{for a.e. } x = \varepsilon\xi+\varepsilon A^\ell + s_1\mathbf{t}_1^\ell + s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell,\ (s_1,s_2,s_3)\in(0,\varepsilon l_\ell)\times D_r,\ \xi\in\Xi_\varepsilon,\end{aligned} \tag{8.2}$$

where $W\in\mathcal{D}(\Omega)$, $V\in\mathcal{D}(\gamma_\ell)$ and $\varphi\in H^1(D)^3$, $\ell\in\{1,\dots,m\}$. Since V belongs to $\mathcal{D}(\gamma_\ell)$ and r/ε goes to 0, the support of the above test-displacement is only included in the beams whose centerline is $\varepsilon\xi+\varepsilon\gamma_\ell$. Moreover, this displacement vanishes in the neighborhood of the extremities of this beam, it means that this displacement vanishes in the junction domain $\mathcal{J}_{\varepsilon,r}$.

One has

$$e_s(\widetilde{v}_{\varepsilon,r}) = \frac{\varepsilon}{r}W(\varepsilon\xi+\varepsilon A^\ell)\begin{pmatrix}\frac{r}{\varepsilon}\frac{dV}{dS_1}\varphi\cdot\mathbf{t}_1^\ell & \frac12\big(V\frac{\partial\varphi}{\partial S_2}\cdot\mathbf{t}_1^\ell + \frac{r}{\varepsilon}\frac{dV}{dS_1}\varphi\cdot\mathbf{t}_2^\ell\big) & \frac12\big(V\frac{\partial\varphi}{\partial S_3}\cdot\mathbf{t}_1^\ell + \frac{r}{\varepsilon}\frac{dV}{dS_1}\varphi\cdot\mathbf{t}_3^\ell\big)\\ * & V\frac{\partial\varphi}{\partial S_2}\cdot\mathbf{t}_2^\ell & \frac12\big(V\frac{\partial\varphi}{\partial S_3}\cdot\mathbf{t}_2^\ell + V\frac{\partial\varphi}{\partial S_2}\cdot\mathbf{t}_3^\ell\big)\\ * & * & V\frac{\partial\varphi}{\partial S_3}\cdot\mathbf{t}_3^\ell\end{pmatrix} \tag{8.3}$$

We apply the unfolding operator $\mathcal{T}_\varepsilon^{b,\ell}$ and pass to the limit, this gives

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(\widetilde{v}_{\varepsilon,r})) \longrightarrow W\,V\,\mathcal{E}_D(\varphi) \quad \text{strongly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times3}. \tag{8.4}$$

Hence

$$\int_{\mathcal{S}_{\varepsilon,r}} \sigma(u_\varepsilon) : e(\widetilde{v}_{\varepsilon,r})\,dx = \int_{\Omega\times\gamma_\ell\times D} \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(\sigma_s(u_\varepsilon)) : \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(\widetilde{v}_{\varepsilon,r}))\,dx$$
$$\to \int_{\Omega\times\gamma_\ell\times D} a_{ijkl}\big(\mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\big)_{ij} W\,V\,\big(\mathcal{E}_D(\varphi)\big)_{kl}\,dx\,dS.$$

Using (5.20) and then unfolding and passing to the limit yield

$$\Big|\int_{\mathcal{S}_{\varepsilon,r}} \mathbf{f}_\varepsilon\cdot\widetilde{v}_{\varepsilon,r}\,dx\Big| = \Big|\frac{r^2}{\varepsilon^2}\int_{\Omega\times\gamma_\ell\times D} \mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f}_\varepsilon)\cdot\mathcal{T}_\varepsilon^{b,\ell}(\widetilde{v}_{\varepsilon,r})\,dx\,d\widehat{S}\Big|$$
$$\le C\frac{r^2}{\varepsilon^2}\cdot\frac{\varepsilon^2}{r+\varepsilon^2}\|\mathbf{f}\|_{L^\infty(\Omega)}\|W\|_{L^\infty(\Omega)}\|V\|_{L^\infty(\gamma_\ell)}\|\varphi\|_{L^2(D)} \longrightarrow 0.$$

The above convergences lead to

$$\int_{\Omega\times\gamma_\ell\times D} a_{ijkl}\big(\mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\big)_{ij} W\,V\big(\mathcal{E}_D(\varphi)\big)_{kl}\,dx\,d\widehat{S} = 0.$$

Finally, since the space $\mathcal{D}(\Omega)\otimes\mathcal{D}(\gamma_\ell)\otimes H^1(D)^3$ is dense in $L^2(\Omega\times\gamma_\ell; H^1(D))^3$ we obtain (8.1). □

Lemma 8.2 (The limit problem involving the extensional and inextensional limit displacements)**.** *One has*

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}\times D} a_{ijkl}\big(\mathcal{E}(\mathcal{U}) + \mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\big)_{kl}\,dx\,d\widehat{S}\\ =&\frac{4\pi}{5}\int_\Omega G\cdot\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(\cdot,A)\Big)\,dx + \frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\widehat{\mathcal{V}}(\cdot,\mathbf{S})\,d\mathbf{S}\,dx, \qquad \forall\,(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}).\end{aligned} \tag{8.5}$$

Proof. Let ϕ be in $\mathcal{D}(\Omega)$ and $(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})$ in $\boldsymbol{\mathcal{D}}_{EX}\times\boldsymbol{\mathcal{D}}_{In}$ such that $\overline{\mathcal{V}}$ and $(\widehat{\mathcal{V}},\widehat{\mathcal{B}})$ are constant in the neighborhood of every node of $\mathcal{S}$.

Step 1. The test displacement.

Set

$$\overline{\mathcal{V}}_{\varepsilon,r} \doteq \phi_{\varepsilon,r}\overline{\mathcal{V}}\Big(\frac{\cdot}{\varepsilon}\Big), \qquad \widehat{\mathcal{V}}_{\varepsilon,r} \doteq \varepsilon\phi_{\varepsilon,r}\widehat{\mathcal{V}}\Big(\frac{\cdot}{\varepsilon}\Big), \qquad \widehat{\mathcal{B}}_{\varepsilon,r} \doteq \phi_{\varepsilon,r}\widehat{\mathcal{B}}\Big(\frac{\cdot}{\varepsilon}\Big)$$

where $\phi_{\varepsilon,r}$ is defined in Appendix F. Since the above fields are constant in the neighborhood of every node of $\mathcal{S}_\varepsilon$, this allows to extend them in functions belonging to $H^1(\mathcal{S}_{\varepsilon,r})$. Hence, these functions are constant in the cross-sections and in the neighborhood of every node. We remind (see Appendix F)

$$\begin{aligned}\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_{\varepsilon,r}) &\longrightarrow \phi \quad \text{strongly in} \quad L^2(\Omega\times\mathcal{S}),\\ \varepsilon\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi_{\varepsilon,r}}{d\mathbf{s}}\Big),\ \ \varepsilon^2\mathcal{T}_\varepsilon^{b,\ell}\Big(\frac{d^2\phi_{\varepsilon,r}}{d\mathbf{s}^2}\Big) &\longrightarrow 0 \quad \text{strongly in} \quad L^2(\Omega\times\mathcal{S}).\end{aligned} \tag{8.6}$$

We define $v_{\varepsilon,r}$ in the beam whose centerline is $\varepsilon\xi+\varepsilon\gamma_\ell$, $\ell\in\{1,\dots,m\}$ by

$$v_{\varepsilon,r}(x) = \frac{\varepsilon^2}{r}\overline{\mathcal{V}}_{\varepsilon,r}\big(\varepsilon\xi+\varepsilon A^\ell+s_1\big) + \frac{\varepsilon^2}{r^2}\widehat{\mathcal{V}}_{\varepsilon,r}\big(\varepsilon\xi+\varepsilon A^\ell+s_1\big) + \frac{\varepsilon^2}{r^2}\widehat{\mathcal{B}}_{\varepsilon,r}\big(\varepsilon\xi+\varepsilon A^\ell+s_1\big)\wedge(s_2\mathbf{t}_2^\ell+s_3\mathbf{t}^\ell) + \widetilde{v}_{\varepsilon,r}(x),$$
$$\widetilde{v}_{\varepsilon,r}(x) = -\frac{\varepsilon^3}{r^2}\frac{d\phi_{\varepsilon,r}}{ds_1}\big(\varepsilon\xi+\varepsilon A^\ell+s_1\big)\widehat{\mathcal{V}}\Big(A^\ell+\frac{s_1}{\varepsilon}\Big)\cdot\big(s_2\mathbf{t}_2^\ell+s_3\mathbf{t}_3^\ell\big)\mathbf{t}_1^\ell,$$
$$\text{for a.e. } x=\varepsilon\xi+\varepsilon A^\ell+s_1\mathbf{t}_1^\ell+s_2\mathbf{t}_2^\ell+s_3\mathbf{t}_3^\ell,\ (s_1,s_2,s_3)\in(0,\varepsilon l_\ell)\times D_r,\ \ell\in\{1,\dots,m\},\ \ \xi\in\Xi_\varepsilon.$$

Observe that for every x in $B(\varepsilon\xi+\varepsilon A^\ell,c_0r)\cap\mathcal{S}_{\varepsilon,r}$ one has

$$v_{\varepsilon,r}(x) = \phi\big(\varepsilon\xi+\varepsilon A^\ell\big)\Big[\frac{\varepsilon^2}{r}\overline{\mathcal{V}}\big(A^\ell\big)+\frac{\varepsilon^2}{r^2}\widehat{\mathcal{B}}\big(A^\ell\big)\wedge\big(x-\varepsilon\xi-\varepsilon A^\ell\big)\Big].$$

Hence, $v_{\varepsilon,r}$ is a rigid displacement in $B(\varepsilon\xi+\varepsilon A^\ell,c_0r)\cap\mathcal{S}_{\varepsilon,r}$. This test displacement belongs to $\mathbf{V}_{\varepsilon,r}$.

Step 2. Limit of the LHS.

One has

$$\begin{aligned}
\frac{\partial v_{\varepsilon,r}}{\partial s_1} &= \frac{\varepsilon^2}{r}\frac{\partial \overline{\mathcal{V}}_{\varepsilon,r}}{\partial s_1} + \frac{\varepsilon^2}{r^2}\frac{\partial \widehat{\mathcal{V}}_{\varepsilon,r}}{\partial s_1} + \frac{\varepsilon^2}{r^2}\frac{\partial \widehat{\mathcal{B}}_{\varepsilon,r}}{\partial s_1} \wedge (s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell) + \frac{\partial \widetilde{v}_{\varepsilon,r}}{\partial s_1},\\
\frac{\partial v_{\varepsilon,r}}{\partial s_1} &= \frac{\varepsilon^2}{r}\frac{d\phi_{\varepsilon,r}}{ds_1}\overline{\mathcal{V}}\Big(\frac{\cdot}{\varepsilon}\Big) + \frac{\varepsilon}{r}\phi_{\varepsilon,r}\frac{d\overline{\mathcal{V}}}{dS_1}\Big(\frac{\cdot}{\varepsilon}\Big) + \frac{\varepsilon^3}{r^2}\frac{d\phi_{\varepsilon,r}}{ds_1}\widehat{\mathcal{V}}\Big(\frac{\cdot}{\varepsilon}\Big) + \frac{\varepsilon^2}{r^2}\phi_{\varepsilon,r}\frac{d\widehat{\mathcal{V}}}{dS_1}\Big(\frac{\cdot}{\varepsilon}\Big)\\
&\quad + \frac{\varepsilon}{r^2}\Big(\varepsilon\frac{d\phi_{\varepsilon,r}}{ds_1}\widehat{\mathcal{B}}\Big(\frac{\cdot}{\varepsilon}\Big) + \phi_{\varepsilon,r}\frac{d\widehat{\mathcal{B}}}{dS_1}\Big(\frac{\cdot}{\varepsilon}\Big)\Big) \wedge (s_2\mathbf{t}_2^\ell + s_3\mathbf{t}_3^\ell) + \frac{\partial \widetilde{v}_{\varepsilon,r}}{\partial s_1},\\
\frac{\partial v_{\varepsilon,r}}{\partial s_2} &= \frac{\varepsilon^2}{r^2}\phi_{\varepsilon,r}\widehat{\mathcal{B}}\Big(\frac{\cdot}{\varepsilon}\Big) \wedge \mathbf{t}_2^\ell + \frac{\partial \widetilde{v}_{\varepsilon,r}}{\partial s_2}, \qquad \frac{\partial v_\varepsilon^\ell}{\partial s_3} = \frac{\varepsilon^2}{r^2}\phi_{\varepsilon,r}\widehat{\mathcal{B}}\Big(\frac{\cdot}{\varepsilon}\Big) \wedge \mathbf{t}_3^\ell + \frac{\partial \widetilde{v}_{\varepsilon,r}}{\partial s_3}.
\end{aligned}$$

Observe that $\dfrac{\partial v_{\varepsilon,r}}{\partial s_2}\cdot\mathbf{t}_2 = \dfrac{\partial v_{\varepsilon,r}}{\partial s_3}\cdot\mathbf{t}_3 = \dfrac{\partial v_{\varepsilon,r}}{\partial s_2}\cdot\mathbf{t}_3 + \dfrac{\partial v_{\varepsilon,r}}{\partial s_3}\cdot\mathbf{t}_2 = 0$ and by definition of $(\widehat{\mathcal{V}},\widehat{\mathcal{B}}) \in \boldsymbol{\mathcal{D}}_{In}$, one has $\widehat{\mathcal{V}}\cdot\mathbf{t}_1 = 0$. The convergences (8.6) yield

$$\begin{aligned}
&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}\Big(\frac{\partial v_{\varepsilon,r}}{\partial s_1}\cdot\mathbf{t}_1^\ell\Big) \longrightarrow \frac{\partial \overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_1^\ell - \frac{\partial^2\widehat{\mathcal{V}}}{\partial S_1^2}\cdot(S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell) \qquad \text{strongly in} \quad L^2(\Omega\times\gamma_\ell\times D),\\
&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}\Big(\frac{\partial \widetilde{v}_{\varepsilon,r}}{\partial s_1}\Big) \longrightarrow 0 \qquad\qquad\qquad\qquad\qquad\qquad \text{strongly in} \quad L^2(\Omega\times\gamma_\ell\times D)^3.
\end{aligned}$$

The presence of $\widetilde{v}_{\varepsilon,r}$ in the test displacement is just to eliminate $\dfrac{\varepsilon^3}{r^2}\dfrac{d\phi_{\varepsilon,r}}{ds_1}\widehat{\mathcal{V}}\Big(\dfrac{\cdot}{\varepsilon}\Big)\cdot\mathbf{t}_\alpha^\ell$ in $\dfrac{\partial v_{\varepsilon,r}}{\partial s_i}\cdot\mathbf{t}_1^\ell + \dfrac{\partial v_{\varepsilon,r}}{\partial s_1}\cdot\mathbf{t}_i^\ell$, $i\in\{2,3\}$. Then, again using the convergences (8.6), we obtain

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(v_{\varepsilon,r})) \longrightarrow \phi\begin{pmatrix} \frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_1^\ell - \frac{\partial^2\widehat{\mathcal{V}}}{\partial S_1^2}\cdot(S_2\mathbf{t}_2^\ell + S_3\mathbf{t}_3^\ell) & * & * \\ \frac{1}{2}\big(\frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_2^\ell - S_3\frac{\partial\widehat{\mathcal{B}}}{\partial S_1}\cdot\mathbf{t}_1^\ell\big) & 0 & * \\ \frac{1}{2}\big(\frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_3^\ell + S_2\frac{\partial\widehat{\mathcal{B}}}{\partial S_1}\cdot\mathbf{t}_1^\ell\big) & 0 & 0 \end{pmatrix} \quad \text{strongly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times 3}.$$

Hence,

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(v_{\varepsilon,r})) \longrightarrow \phi\big(\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}}) + \mathcal{E}_D(\widetilde{\overline{v}})\big) \quad \text{strongly in} \quad L^2(\Omega\times\gamma_\ell\times D)^{3\times 3} \tag{8.7}$$

where

$$\widetilde{\overline{v}} = S_2\Big(\frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_2\Big)\mathbf{t}_1 + S_3\Big(\frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_3\Big)\mathbf{t}_1.$$

Unfolding the left-hand side of (5.8) and passing to the limit give

$$\begin{aligned}
&\int_{\mathcal{S}_{\varepsilon,r}} \sigma(u_\varepsilon) : e(v_{\varepsilon,r})\,dx = \sum_{\ell=1}^m \int_{\Omega\times\gamma_\ell\times D} \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(\sigma_s(u_\varepsilon)) : \frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(v_{\varepsilon,r}))\,dx\\
&\rightharpoonup \int_{\Omega\times\mathcal{S}\times D} a_{ijkl}\big(\mathcal{E}(\mathcal{U}) + \mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\big)_{ij}\phi\big(\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}}) + \mathcal{E}_D(\widetilde{\overline{v}})\big)_{kl}\,dx\,d\widehat{S}.
\end{aligned}$$

Step 3. Limit of the RHS.
Now, we consider the right-hand side of (5.8)

$$\int_{\mathcal{S}_{\varepsilon,r}} f_\varepsilon\cdot v_{\varepsilon,r}\,dx = \sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)} F_{r,\mathcal{K}_\varepsilon}\cdot v_{\varepsilon,r}\,dx + \int_{\mathcal{S}_{\varepsilon,r}} \mathbf{f}_\varepsilon\cdot v_{\varepsilon,r}\,dx. \tag{8.8}$$

Let's take the first term in the right-hand side of (8.8). Taking into account the symmetries of the ball $B(\varepsilon\xi+\varepsilon A^\ell,r)$ and the fact that $\displaystyle\int_{B(O,r)}|x|^2\,dx = \frac{4\pi r^5}{5}$. After a straightforward calculation, one obtains

$$\begin{aligned}
\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)} F_{r,\mathcal{K}_\varepsilon}\cdot v_{\varepsilon,r}\,dx &= \sum_{A^\ell\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}\int_{B(\varepsilon\xi+\varepsilon A^\ell,r)}\Big[\frac{\varepsilon^2}{r^2}F(\varepsilon\xi+\varepsilon A^\ell) + \frac{\varepsilon}{r^3}G(\varepsilon\xi+\varepsilon A^\ell)\wedge(x-(\varepsilon\xi-\varepsilon A^\ell)\Big]\\
&= \frac{4\pi}{3}\varepsilon^4\sum_{A^\ell\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}\phi\big(\varepsilon\xi+\varepsilon A^\ell\big)F(\varepsilon\xi+\varepsilon A^\ell)\cdot\overline{\mathcal{V}}\big(A^\ell\big) + \frac{4\pi}{5}\varepsilon^3\sum_{A^\ell\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}\phi\big(\varepsilon\xi+\varepsilon A^\ell\big)G(\varepsilon\xi+\varepsilon A^\ell)\cdot\widehat{\mathcal{B}}\big(A^\ell\big).
\end{aligned}$$

Since $|Y|=1$, one has

$$\sum_{A^\ell\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}\varepsilon^3\phi\big(\varepsilon\xi+\varepsilon A^\ell\big)F(\varepsilon\xi+\varepsilon A^\ell)\cdot\overline{\mathcal{V}}\big(A^\ell\big)\longrightarrow\int_\Omega F\cdot\phi\Big(\sum_{A\in\mathcal{K}}\overline{\mathcal{V}}\big(A\big)\Big)\,dx$$
$$\sum_{A^\ell\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}\varepsilon^3\phi\big(\varepsilon\xi+\varepsilon A^\ell\big)G(\varepsilon\xi+\varepsilon A^\ell)\cdot\widehat{\mathcal{B}}\big(A^\ell\big)\longrightarrow\int_\Omega G\cdot\phi\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}\big(A\big)\Big)\,dx.$$

Hence,

$$\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}F_{r,\mathcal{K}_\varepsilon}\cdot v_{\varepsilon,r}\,dx\longrightarrow\frac{4\pi}{5}\int_\Omega G\cdot\phi\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}\big(A\big)\Big)\,dx.\tag{8.9}$$

Now, we take the second term in the right-hand side of (8.8).

Due to (6.6), we only need to consider $\dfrac{r^2}{\varepsilon^2}\displaystyle\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f}_\varepsilon)\cdot\mathcal{T}_\varepsilon^{b,\ell}(v_{\varepsilon,r})\,dx\,d\widehat{S}$. One has

$$\begin{aligned}&\frac{r^2}{\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f}_\varepsilon)\cdot\mathcal{T}_\varepsilon^{b,\ell}(v_{\varepsilon,r})\,dx\,d\widehat{S}\\ =&\frac{r^2}{\varepsilon^2}\frac{\varepsilon}{r+\varepsilon^2}\frac{\varepsilon^2}{r^2}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f})\cdot\Big[r\mathcal{T}_\varepsilon^{b,\ell}(\overline{\mathcal{V}}_{\varepsilon,r})+\mathcal{T}_\varepsilon^{b,\ell}(\widehat{\mathcal{V}}_{\varepsilon,r})+r\mathcal{T}_\varepsilon^{b,\ell}(\widehat{\mathcal{B}}_{\varepsilon,r})\wedge(S_2\mathbf{t}_2^\ell+S_3\mathbf{t}_3^\ell)\Big]\,dx\,d\widehat{S}.\end{aligned}$$

Assumptions (7.1) and convergence $(8.6)_1$ lead to

$$\begin{aligned}&\frac{\varepsilon^2}{r+\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f})\cdot\frac{1}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(\widehat{\mathcal{V}}_{\varepsilon,r})\,dx\,d\widehat{S}\longrightarrow\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}(x)\cdot\phi(x)\,\widehat{\mathcal{V}}(\mathbf{S})\,dx\,d\mathbf{S}\\ &\frac{\varepsilon r}{r+\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f})\cdot\mathcal{T}_\varepsilon^{b,\ell}(\overline{\mathcal{V}}_{\varepsilon,r})\,d\widehat{S}\,dx\longrightarrow0,\\ &\frac{\varepsilon r}{r+\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f}_\varepsilon)\cdot\big[\mathcal{T}_\varepsilon^{b,\ell}(\widehat{\mathcal{B}}_{\varepsilon,r})\wedge(S_2\mathbf{t}_2^\ell+S_3\mathbf{t}_3^\ell)\big]\,dx\,d\widehat{S}\longrightarrow0.\end{aligned}$$

Hence,

$$\int_{\mathcal{S}_{\varepsilon,r}}f_\varepsilon\cdot v_{\varepsilon,r}\,dx\longrightarrow\frac{4\pi}{5}\int_\Omega G\cdot\phi\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}\big(A\big)\Big)\,dx+\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\phi\,\widehat{\mathcal{V}}(\mathbf{S})\,dx\,d\mathbf{S}.$$

Lemma E.2 and the density of $\mathcal{D}(\Omega)\otimes\boldsymbol{\mathcal{D}}_{Ex}$ in $L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex})$ and $\mathcal{D}(\Omega)\otimes\boldsymbol{\mathcal{D}}_{In}$ in $L^2(\Omega;\boldsymbol{\mathcal{D}}_{In})$ lead to

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{\overline{v}})\big)_{kl}\,dx\,d\widehat{S}\\ =&\frac{4\pi}{5}\int_\Omega G\cdot\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}\big(\cdot,A\big)\Big)\,dx+\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\widehat{\mathcal{V}}(\cdot,\mathbf{S})\,dx\,d\mathbf{S},\qquad\forall\,(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}).\end{aligned}$$

Besides, since $\widetilde{\overline{v}}$ belongs to $L^2(\Omega\times\mathcal{S};H^1(D))^3$ equality (8.1) together with the one above yield (8.5). □

Lemma 8.3 (The limit problem involving the macroscopic limit displacement)**.** *One has*

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{V})\big)_{kl}\,dx\,d\widehat{S}\\ &=\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx,\qquad\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\end{aligned}\tag{8.10}$$

where $|\mathcal{K}|$ is the number of points of $\mathcal{K}$ and $\mathcal{S}$ the measure of $\mathcal{S}$.

[2]Here, by convention $\dfrac{+\infty}{1+\infty}=1$.

Proof. Step 1. Limit of the LHS of (5.8).

Let $\mathcal{V}$ be in $\mathcal{D}(\mathbb{R}^3)^3$ such that $\mathcal{V}=0$ in $\Omega'\setminus\overline{\Omega}$. We define $\mathcal{V}_{\varepsilon,r}$ using F. This function is extended as in Step 1 of the proof of Lemma 8.2. Set

$$v_{\varepsilon,r}=\frac{\varepsilon}{r}\mathcal{V}_{\varepsilon,r}\in\mathbf{V}_{\varepsilon,r}.$$

We have

$$\begin{aligned}&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(v_{\varepsilon,r})\longrightarrow\mathcal{V}\quad\text{strongly in}\quad L^2(\Omega\times\gamma_\ell\times D)^3,\\ \text{and}\quad&\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(v_{\varepsilon,r}))\longrightarrow\begin{pmatrix}(\nabla\mathcal{V}\,\mathbf{t}_1^\ell)\cdot\mathbf{t}_1^\ell & * & *\\ \frac12(\nabla\mathcal{V}\,\mathbf{t}_1^\ell)\cdot\mathbf{t}_2^\ell & 0 & 0\\ \frac12(\nabla\mathcal{V}\,\mathbf{t}_1^\ell)\cdot\mathbf{t}_3^\ell & 0 & 0\end{pmatrix}=\mathcal{E}(\mathcal{V})+\mathcal{E}_D(\widetilde{\widetilde{v}})\quad\text{strongly in}\quad L^2(\Omega\times\gamma_\ell\times D)^{3\times3}\end{aligned}\tag{8.11}$$

where

$$\widetilde{\widetilde{v}}=S_2\big(\big(\nabla\mathcal{V}\,\mathbf{t}_1\big)\cdot\mathbf{t}_2\big)\,\mathbf{t}_1+S_3\big(\big(\nabla\mathcal{V}\,\mathbf{t}_1\big)\cdot\mathbf{t}_3\big)\,\mathbf{t}_1,\qquad\text{a.e. in }\ \Omega\times\mathcal{S}\times D.$$

Convergence (8.11) leads to

$$\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(\mathcal{V})\,dx\longrightarrow\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{V})+\mathcal{E}_D(\widetilde{\widetilde{v}})\big)_{kl}\,dx\,d\widehat{S}.$$

Step 2. Limit of the RHS.

Now we consider the right-hand side of (5.8). By (5.20), firstly we have

$$\begin{aligned}\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}F_{r,\mathcal{K}_\varepsilon}\cdot v_{\varepsilon,r}\,dx&=\sum_{\mathcal{A}\in\mathcal{K}_\varepsilon}\int_{B(\mathcal{A},r)}\Big(\frac{\varepsilon^2}{r^2}F(\mathcal{A})+\frac{\varepsilon}{r^3}G(\mathcal{A})\wedge(x-\mathcal{A})\Big)\cdot\frac{\varepsilon}{r}\mathcal{V}(\mathcal{A})\,dx\\&=\frac{4\pi}{3}\sum_{A\in\mathcal{K}}\sum_{\xi\in\Xi_\varepsilon}F(\varepsilon\xi+\varepsilon A)\cdot\mathcal{V}(\varepsilon\xi+\varepsilon A)\varepsilon^3\longrightarrow\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx\end{aligned}$$

and secondly, due to (6.6), we pass to the limit in

$$\frac{r^2}{\varepsilon^2}\frac{\varepsilon}{r+\varepsilon^2}\frac{\varepsilon}{r}\sum_{\ell=1}^m\int_{\Omega\times\gamma_\ell\times D}\mathcal{T}_\varepsilon^{b,\ell}(\mathbf{f}_\varepsilon)\cdot\mathcal{T}_\varepsilon^{b,\ell}(\mathcal{V}_{\varepsilon,r})\,dxdS\longrightarrow\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx.$$

Hence

$$\int_{\mathcal{S}_{\varepsilon,r}}f_\varepsilon\cdot v_{\varepsilon,r}=\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}F_{r,\mathcal{K}_\varepsilon}\cdot v_{\varepsilon,r}\,dx+\int_{\mathcal{S}_{\varepsilon,r}}\mathbf{f}_\varepsilon\cdot v_{\varepsilon,r}\,dx\longrightarrow\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx.$$

Since the set of functions belonging to $\mathcal{D}(\mathbb{R}^3)^3$ and vanishing in $\Omega'\setminus\Omega$ is dense in $H^1_\Gamma(\Omega)^3$, we obtain

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{V})+\mathcal{E}_D(\widetilde{\widetilde{v}})\big)_{kl}dxd\widehat{S}\\&=\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx,\qquad\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\end{aligned}$$

Taking into account that $\widetilde{\widetilde{v}}$ belongs to $L^2(\Omega\times\mathcal{S};H^1(D))^3$ and using (8.1), equality (8.10) is proved. □

Theorem 8.1 (The unfolded limit problem)**.** *Let u_ε be the solution to* (5.8)*. There exist $\mathcal{U}\in H^1_\Gamma(\Omega)^3$, $(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In})$ and $\widetilde{u}\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w)$ such that $\big(\mathcal{U},\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}},\widetilde{u}\big)$ is the solution to the following unfolded problem:*

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{V})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{v})\big)_{kl}\,dx\,d\widehat{S}\\&=\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{4\pi}{5}\int_\Omega G\cdot\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(\cdot,A)\Big)\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx+\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\widehat{\mathcal{V}}(\cdot,\mathbf{S})\,dx\,d\mathbf{S}\\&\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\quad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}),\quad\forall\widetilde{v}\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w).\end{aligned}\tag{8.12}$$

Moreover, the following convergences hold ($\ell\in\{1,\dots,m\}$):

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon))\longrightarrow\mathcal{E}(\mathcal{U})+\mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\quad\textit{strongly in}\quad L^2(\Omega\times\gamma_\ell\times D)^{3\times3}.\tag{8.13}$$

Denote

$$\mathbf{M}^{11}=\begin{pmatrix}1&0&0\\0&0&0\\0&0&0\end{pmatrix},\quad \mathbf{M}^{22}=\begin{pmatrix}0&0&0\\0&1&0\\0&0&0\end{pmatrix},\quad \mathbf{M}^{33}=\begin{pmatrix}0&0&0\\0&0&0\\0&0&1\end{pmatrix},$$
$$\mathbf{M}^{12}=\mathbf{M}^{21}=\frac{1}{2}\begin{pmatrix}0&1&0\\1&0&0\\0&0&0\end{pmatrix},\quad \mathbf{M}^{13}=\mathbf{M}^{31}=\frac{1}{2}\begin{pmatrix}0&0&1\\0&0&0\\1&0&0\end{pmatrix},\quad \mathbf{M}^{23}=\mathbf{M}^{32}=\frac{1}{2}\begin{pmatrix}0&0&0\\0&0&1\\0&1&0\end{pmatrix}.$$

Proof. From Lemmas 8.1, 8.2, 8.3 we obtain that $(\mathcal{U},\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}},\widetilde{u})$ satisfies (8.12) for every test function $\mathcal{V}\in H^1_\Gamma(\Omega)^3$, $(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In})$ and $\widetilde{w}\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w)\subset L^2(\Omega\times\mathcal{S};H^1(D))^3$.
The coercivity of this problem is given by Lemma G.2. Since the problem (8.12) admits a unique solution, the whole sequences in Theorems 7.1, 7.2 and (8.13) converge to their limits.

Now, we prove the strong convergence (8.13). First, observe that due to the inclusion of $\mathcal{J}_{\varepsilon,r}$ in $\bigcup_{A\in\mathcal{K}_\varepsilon} B(A,c_0r)$ given by (5.1), the portions of beams which correspond to $S_1\in(2c_0r,l_\ell-2c_0r)$ are all disjoint. Furthermore, since $\sigma(u_\varepsilon):e(u_\varepsilon)$ is non-negative, one has

$$\frac{r^2}{\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times(0,l_\ell)\times D}\mathcal{T}_\varepsilon^{b,\ell}\big(\sigma_s(u_\varepsilon)\big):\mathcal{T}_\varepsilon^{b,\ell}\big(e_s(u_\varepsilon)\big)\mathbf{1}_{(2c_0r,l_\ell-2c_0r)}\,dx\,d\widehat{S}\le\liminf_{\varepsilon\to0}\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(u_\varepsilon)\,dx.$$

From (7.13) and the fact that r goes to 0, one obtains ($\ell\in\{1,\dots,m\}$)

$$\frac{r}{\varepsilon}\mathcal{T}_\varepsilon^{b,\ell}(e_s(u_\varepsilon))\mathbf{1}_{(2c_0r,l_\ell-2c_0r)}\rightharpoonup\mathcal{E}(\mathcal{U})+\mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\quad\text{weakly in}\quad L^2(\Omega\times\gamma_\ell\times D)^{3\times3}.$$

Hence, choosing u_ε as a test function in (5.8) and using a weak lower semi-continuity of convex functionals, one has

$$\begin{aligned}
&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{u})\big)_{kl}\,dx\,d\widehat{S}\\
\le&\liminf_{\varepsilon\to0}\frac{r^2}{\varepsilon^2}\sum_{\ell=1}^m\int_{\Omega\times(0,l_\ell)\times D}\mathcal{T}_\varepsilon^{b,\ell}\big(a^\varepsilon_{ijkl}\big)\mathcal{T}_\varepsilon^{b,\ell}\big(e_{s,ij}(u_\varepsilon)\big)\mathcal{T}_\varepsilon^{b,\ell}\big(e_{s,kl}(u_\varepsilon)\big)\mathbf{1}_{(2c_0r,l_\ell-2c_0r)}\,dx\,d\widehat{S}\\
\le&\liminf_{\varepsilon\to0}\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(u_\varepsilon)\,dx\le\limsup_{\varepsilon\to0}\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(u_\varepsilon)\,dx=\limsup_{\varepsilon\to0}\int_{\mathcal{S}_{\varepsilon,r}}f_\varepsilon\cdot u_\varepsilon dx\\
=&\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{U}\,dx+\frac{4\pi}{5}\int_\Omega G\cdot\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{R}}(\cdot,A)\Big)\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{U}\,dx+\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\widehat{\mathcal{U}}(\cdot,\mathbf{S})\,dx\,d\mathbf{S},\\
=&\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{u})\big)_{kl}\,dx\,d\widehat{S}.
\end{aligned}$$

Thus, all inequalities above are equalities and

$$\lim_{\varepsilon\to0}\int_{\mathcal{S}_{\varepsilon,r}}\sigma(u_\varepsilon):e(u_\varepsilon)\,dx=\int_{\Omega\times\mathcal{S}\times D}a_{ijkl}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})+\mathcal{E}_D(\widetilde{u})\big)_{ij}\big(\mathcal{E}(\mathcal{U})+\mathcal{E}_\mathcal{S}(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{u})\big)_{kl}\,dx\,d\widehat{S},$$

which in turn leads to the strong convergence (8.13). □

9 The homogenized problem

9.1 Expression of the warping $\widetilde{u}$

In this subsection we give the expression of the warping $\widetilde{u}$ in terms of the macroscopic displacement $\mathcal{U}$ and the microscopic fields $\overline{\mathcal{U}}$, $\widehat{\mathcal{U}}$, $\widehat{\mathcal{R}}$.

To this end, we use the variational formulation (8.1). For every $\ell\in\{1,\dots,m\}$ one has

$$\int_D a_{ijkl}\left(\mathcal{E}_D(\widetilde{u})\right)_{ij}\left(\mathcal{E}_D(\widetilde{v})\right)_{kl}dS_2dS_3=-\int_D a_{ijkl}\left(\mathcal{E}(\mathcal{U})+\mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})\right)_{ij}\left(\mathcal{E}_D(\widetilde{v})\right)_{kl}dS_2dS_3,\ \text{ a.e. in }\Omega\times\gamma_\ell,\ \forall\,\widetilde{v}\in\boldsymbol{\mathcal{D}}_w.$$

This shows that $\widetilde{u}$ can be expressed in terms of the elements of the tensors $\mathcal{E}$ and $\mathcal{E}_{\mathcal{S}}$.
We write

$$\mathcal{E}(\mathcal{U})+\mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{V}},\widehat{\mathcal{R}})=\begin{pmatrix}\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1 & 0 & 0\\ 0 & 0 & 0\\ 0 & 0 & 0\end{pmatrix}+\begin{pmatrix}\frac{\partial\overline{\mathcal{U}}}{\partial S_1}\cdot\mathbf{t}_1-\frac{\partial^2\widehat{\mathcal{U}}}{\partial S_1^2}\cdot\Big(S_2\,\mathbf{t}_2+S_3\,\mathbf{t}_3\Big) & * & *\\ -\frac{S_3}{2}\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1 & 0 & 0\\ \frac{S_2}{2}\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1 & 0 & 0\end{pmatrix}$$

$$\begin{aligned}&=\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{U}}}{\partial S_1}\cdot\mathbf{t}_1\Big)\mathbf{M}^{11}-\sum_{\alpha=2}^{3}\frac{\partial^2\widehat{\mathcal{U}}}{\partial S_1^2}\cdot\mathbf{t}_\alpha\,S_\alpha\mathbf{M}^{11}+\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1\Big(S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\Big)\\ &=\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{U}}}{\partial S_1}\cdot\mathbf{t}_1\Big)\mathbf{M}^{11}+\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_1\Big(S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\Big)+\Big(\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_2\,S_3-\frac{\partial\widehat{\mathcal{R}}}{\partial S_1}\cdot\mathbf{t}_3\,S_2\Big)\mathbf{M}^{11}\\ &\quad\text{a.e. in}\quad\Omega\times\mathcal{S}\times D.\end{aligned}\tag{9.1}$$

Now, we introduce 4 correctors which are the solutions to the following cell problems:

$$\begin{aligned}&\int_D a_{ijkl}(\mathbf{S},\cdot)\,\big(\mathcal{E}_D(\widetilde{\chi}_1)(\mathbf{S},\cdot)+\mathbf{M}^{11}\big)_{ij}\big(\mathcal{E}_D(\widetilde{v})\big)_{kl}\,dS_2dS_3=0,\\ &\int_D a_{ijkl}(\mathbf{S},\cdot)\,\big(\mathcal{E}_D(\widetilde{\chi}_2)(\mathbf{S},\cdot)+S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\big)_{ij}\big(\mathcal{E}_D(\widetilde{v})\big)_{kl}\,dS_2dS_3=0,\\ &\int_D a_{ijkl}(\mathbf{S},\cdot)\,\big(\mathcal{E}_D(\widetilde{\chi}_3)(\mathbf{S},\cdot)+S_3\mathbf{M}^{11}\big)_{ij}\big(\mathcal{E}_D(\widetilde{v})\big)_{kl}\,dS_2dS_3=0,\\ &\int_D a_{ijkl}(\mathbf{S},\cdot)\,\big(\mathcal{E}_D(\widetilde{\chi}_4)(\mathbf{S},\cdot)-S_2\mathbf{M}^{11}\big)_{ij}\big(\mathcal{E}_D(\widetilde{v})\big)_{kl}\,dS_2dS_3=0,\end{aligned}\qquad\text{for a.e. }\mathbf{S}\text{ in }\mathcal{S},\quad\forall\,\widetilde{v}\in\boldsymbol{\mathcal{D}}_w.\tag{9.2}$$

Since a_{ijkl}'s belong to $L^\infty(\mathcal{S}\times D)$, then $\widetilde{\chi}_q\in L^\infty(\mathcal{S};\boldsymbol{\mathcal{D}}_w)$, $q\in\{1,\dots,4\}$.
Hence, we have

$$\widetilde{u}=\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big)\,\widetilde{\chi}_1+\sum_{q=1}^{3}\frac{\partial\widetilde{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_q\,\widetilde{\chi}_{q+1}\quad\text{a.e. in }\Omega\times\mathcal{S}\times D.$$

9.2 Expression of the microscopic fields $\overline{\mathcal{U}}$, $\widehat{\mathcal{U}}$, $\widehat{\mathcal{R}}$

In this subsection we give the expression of the microscopic fields $\overline{\mathcal{U}}$, $\widehat{\mathcal{U}}$, $\widehat{\mathcal{R}}$ in terms of the macroscopic displacement $\mathcal{U}$. To this end, as before, we use the variational formulation (8.12).
Thus, taking $\mathcal{V}=0$, $\widetilde{v}=0$ in (8.12), then replacing $\widetilde{u}$ by its expression, using the following equality:

$$\mathcal{E}_D(\widetilde{u})=\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big)\,\mathcal{E}_D(\widetilde{\chi}_1)+\sum_{q=1}^{3}\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_q\,\mathcal{E}_D(\widetilde{\chi}_{q+1})\quad\text{a.e. in }\Omega\times\mathcal{S}\times D$$

together with (9.1) give

$$\begin{aligned}&\int_{\mathcal{S}}\Big(\int_D a_{ijkl}\Big[\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big)\big(\mathcal{E}_D(\widetilde{\chi}_1)+\mathbf{M}^{11}\big)+\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big(\mathcal{E}_D(\widetilde{\chi}_2)+S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\Big)\\ &+\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_2\Big(\mathcal{E}_D(\widetilde{\chi}_3)+S_3\mathbf{M}^{11}\Big)+\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_3\Big(\mathcal{E}_D(\widetilde{\chi}_4)-S_2\mathbf{M}^{11}\Big)\Big]_{ij}\Big[\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\Big]_{kl}\,dS_2dS_3\Big)d\mathbf{S}\\ &=\frac{4\pi}{5}G\cdot\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(A)+\frac{\pi}{1+\kappa}\mathbf{f}\cdot\int_{\mathcal{S}}\widehat{\mathcal{V}}(\mathbf{S})\,d\mathbf{S},\qquad\text{a.e. in }\Omega,\qquad\forall\,(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}.\end{aligned}\tag{9.3}$$

We write

$$\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})=\Big(\frac{\partial\overline{\mathcal{V}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big)\mathbf{M}^{11}+\frac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big(S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\Big)+\Big(\frac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_2\,S_3-\frac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_3\,S_2\Big)\mathbf{M}^{11}$$

and the variational problem (9.3) has the following form:

$$\begin{aligned}&\int_{\mathcal{S}}\mathfrak{A}\,\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{U}}\\ \cdots\\ \widehat{\mathcal{R}}\end{pmatrix}\cdot\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}=-\int_{\mathcal{S}}\mathfrak{A}\begin{pmatrix}(e(\mathcal{U})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}\cdot\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}\\ &+\frac{4\pi}{5}\,G\cdot\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(A)+\frac{\pi}{1+\kappa}\mathbf{f}\cdot\int_{\mathcal{S}}\widehat{\mathcal{V}}(\mathbf{S})\,d\mathbf{S},\quad\text{a.e. in }\Omega,\quad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}\end{aligned}\tag{9.4}$$

where the symmetric matrix $\mathfrak{A}$ belongs to $L^\infty(\mathcal{S})^{4\times4}$.

Here, the column $\dfrac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}$ *stands for the column* $\Big(\dfrac{\partial\overline{\mathcal{V}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\ \ \dfrac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\ \ \dfrac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_2\ \ \dfrac{\partial\widehat{\mathcal{B}}}{\partial\mathbf{S}}\cdot\mathbf{t}_3\Big)^T$*, while the column* $\begin{pmatrix}(e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}$ *stands for* $\big((e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\ \ 0\ \ 0\ \ 0\big)^T$.

Matrix $\mathfrak{A}$ satisfies

$$\begin{aligned}\forall\zeta\in\mathbb{R}^4,\qquad\mathfrak{A}\,\zeta\cdot\zeta=\int_D a_{ijkl}\Big[&\zeta_1\big(\mathcal{E}_D(\widetilde{\chi}_1)+\mathbf{M}^{11}\big)+\zeta_2\left(\mathcal{E}_D(\widetilde{\chi}_2)+S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\right)\\ &+\zeta_3\left(\mathcal{E}_D(\widetilde{\chi}_3)+S_3\mathbf{M}^{11}\right)+\zeta_4\left(\mathcal{E}_D(\widetilde{\chi}_4)-S_2\mathbf{M}^{11}\right)\Big]_{ij}\\ &\times\Big[\zeta_1\big(\mathcal{E}_D(\widetilde{\chi}_1)+\mathbf{M}^{11}\big)+\zeta_2\left(\mathcal{E}_D(\widetilde{\chi}_2)+S_2\mathbf{M}^{13}-S_3\mathbf{M}^{12}\right)\\ &+\zeta_3\left(\mathcal{E}_D(\widetilde{\chi}_3)+S_3\mathbf{M}^{11}\right)+\zeta_4\left(\mathcal{E}_D(\widetilde{\chi}_4)-S_2\mathbf{M}^{11}\right)\Big]_{kl}dS_2dS_3\end{aligned}\qquad\text{a.e. in }\mathcal{S}$$

since $\widetilde{\chi}_q$'s verify (9.2).

At this step, the unfolded problem becomes

$$\begin{aligned}&\int_{\Omega\times\mathcal{S}}\mathfrak{A}\left[\begin{pmatrix}(e(\mathcal{U})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{U}}\\ \cdots\\ \widehat{\mathcal{R}}\end{pmatrix}\right]\cdot\left[\begin{pmatrix}(e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}\right]dx\,d\mathbf{S}\\ =&\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{4\pi}{5}\int_\Omega G\cdot\Big(\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(\cdot,A)\Big)\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx+\frac{\pi}{1+\kappa}\int_{\Omega\times\mathcal{S}}\mathbf{f}\cdot\widehat{\mathcal{V}}(\cdot,\mathbf{S})\,dx\,d\mathbf{S},\\ &\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\quad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}).\end{aligned}\tag{9.5}$$

Now, we introduce 12 correctors

$$\chi^{ij}\doteq(\overline{\chi}^{ij},\widehat{\chi}^{ij},\widehat{\boldsymbol{\chi}}^{ij}),\quad\chi^q\doteq(\overline{\chi}^q,\widehat{\chi}^q,\widehat{\boldsymbol{\chi}}^q)\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\qquad(i,j)\in\{1,2,3\}^2,\qquad q\in\{1,\dots,6\}.$$

They are the solutions to the following variational problems:

$$\begin{aligned}&\chi^{ij}\doteq(\overline{\chi}^{ij},\widehat{\chi}^{ij},\widehat{\boldsymbol{\chi}}^{ij})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\\ &\int_{\mathcal{S}}\mathfrak{A}\,\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\cdot\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}=-\int_{\mathcal{S}}\mathfrak{A}\begin{pmatrix}(\mathbf{M}^{ij}\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}\cdot\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}\qquad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\\ &\chi^q\doteq(\overline{\chi}^q,\widehat{\chi}^q,\widehat{\boldsymbol{\chi}}^q)\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\qquad q\in\{1,2,3\},\\ &\int_{\mathcal{S}}\mathfrak{A}\,\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\chi}^{q}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{q}\end{pmatrix}\cdot\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}=\mathbf{e}_q\cdot\sum_{A\in\mathcal{K}}\widehat{\mathcal{B}}(A)\qquad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\\ &\chi^{q+3}\doteq(\overline{\chi}^{q+3},\widehat{\chi}^{q+3},\widehat{\boldsymbol{\chi}}^{q+3})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\qquad q\in\{1,2,3\},\\ &\int_{\mathcal{S}}\mathfrak{A}\,\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\chi}^{q+3}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{q+3}\end{pmatrix}\cdot\frac{d}{d\mathbf{S}}\begin{pmatrix}\overline{\mathcal{V}}\\ \cdots\\ \widehat{\mathcal{B}}\end{pmatrix}d\mathbf{S}=\mathbf{e}_q\cdot\int_{\mathcal{S}}\widehat{\mathcal{V}}\,d\mathbf{S},\qquad\forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In},\end{aligned}\tag{9.6}$$

where $\mathbf{e}_1=\big(1\ \ 0\ \ 0\big)^T$, $\mathbf{e}_2=\big(0\ \ 1\ \ 0\big)^T$ and $\mathbf{e}_3=\big(0\ \ 0\ \ 1\big)^T$. Note that $\chi^{ij}=\chi^{ji}$.

Hence, one has

$$(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) = \sum_{i,j=1}^{3} e_{ij}(\mathcal{U})\,\chi^{ij} + \frac{4\pi}{5}\sum_{q=1}^{3} G_q\chi^q + \frac{\pi}{1+\kappa}\sum_{q=1}^{3}\mathbf{f}_q\chi^{q+3}, \tag{9.7}$$

where $G=\sum_{q=1}^{3} G_q\mathbf{e}_q,\ \mathbf{f}=\sum_{q=1}^{3}\mathbf{f}_q\mathbf{e}_q$.

In problem (9.5), we replace $(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}})$ by (9.7) and we choose $(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})=(0,0,0)$. That gives

$$\begin{aligned}
&\int_{\Omega\times\mathcal{S}} \mathfrak{A}\left[\begin{pmatrix}(e(\mathcal{U})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} + \sum_{i,j=1}^{3} e_{ij}(\mathcal{U})\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right]\cdot\begin{pmatrix}(e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} dx\,d\mathbf{S}\\
&\quad= -\int_\Omega\left(\frac{4\pi}{5}\sum_{q=1}^{3} G_q\left[\int_{\mathcal{S}}\mathfrak{A}\,\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{q}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{q}\end{pmatrix}\cdot\begin{pmatrix}(e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} d\mathbf{S}\right]\right.\\
&\qquad\left.+\frac{\pi}{1+\kappa}\sum_{q=1}^{3}\mathbf{f}_q\left[\int_{\mathcal{S}}\mathfrak{A}\,\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{q+3}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{q+3}\end{pmatrix}\cdot\begin{pmatrix}(e(\mathcal{V})\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} d\mathbf{S}\right]\right) dx\\
&\qquad+\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx,\qquad\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3.
\end{aligned}\tag{9.8}$$

Now, taking into account the definition of the corrector $\chi^{ij}\doteq(\overline{\chi}^{ij},\widehat{\chi}^{ij},\widehat{\boldsymbol{\chi}}^{ij})$, the left-hand side becomes

$$\int_\Omega \mathfrak{B}^{hom}\big(e(\mathcal{U}),e(\mathcal{V})\big)\,dx,$$

where $\mathfrak{B}^{hom}$ is a symmetric bilinear form associated to the definite positive quadratic form

$$\begin{aligned}
\mathfrak{B}^{hom}(\zeta,\zeta) &= \int_{\mathcal{S}}\mathfrak{A}\left[\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right]\cdot\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} d\mathbf{S}\\
&= \int_{\mathcal{S}}\mathfrak{A}\left[\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right]\cdot\left[\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right] d\mathbf{S}\\
&= \mathfrak{b}^{hom}_{ijkl}\,\zeta_{ij}\zeta_{kl}
\end{aligned}\tag{9.9}$$

for every 3×3 symmetric matrix ζ.

Write $\zeta=\sum_{i,j=1}^{3}\zeta_{ij}\mathbf{M}^{ij}$. Hence,

$$\mathfrak{b}^{hom}_{ijkl} = \int_{\mathcal{S}}\mathfrak{A}\left[\begin{pmatrix}(\mathbf{M}^{ij}\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right]\cdot\left[\begin{pmatrix}(\mathbf{M}^{kl}\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{kl}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{kl}\end{pmatrix}\right] d\mathbf{S}. \tag{9.10}$$

Now, we simplify the right-hand side of (9.8). Set

$$\mathfrak{c}^{hom}_{ijq} = \int_{\mathcal{S}}\mathfrak{A}\,\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{q}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{q}\end{pmatrix}\cdot\left[\begin{pmatrix}(\mathbf{M}^{ij}\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix}+\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right] d\mathbf{S},\qquad (i,j,q)\in\{1,2,3\}^2\times\{1,\dots,6\}. \tag{9.11}$$

Thus, the limit field $\mathcal{U}\in H^1_\Gamma(\Omega)^3$ is the solution to the homogenized problem

$$\begin{aligned}
\int_\Omega \mathfrak{b}^{hom}_{ijkl}\,e_{ij}(\mathcal{U})\,e_{kl}(\mathcal{V})\,dx = &-\frac{4\pi}{5}\sum_{q=1}^{3}\int_\Omega G_q\,\mathfrak{c}^{hom}_{ijq}\,e_{ij}(\mathcal{V})\,dx+\frac{\pi}{1+\kappa}\sum_{q=1}^{3}\int_\Omega\mathbf{f}_q\,\mathfrak{c}^{hom}_{ijq+3}\,e_{ij}(\mathcal{V})\,dx\\
&+\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx+\frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega\mathbf{f}\cdot\mathcal{V}\,dx,\qquad\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3.
\end{aligned}\tag{9.12}$$

Lemma 9.1. *The components of the homogenized elasticity tensor $\mathfrak{b}_{ijkl}\in\mathbb{R}$ satisfy the usual symmetry and positivity conditions*

- $\mathfrak{b}^{hom}_{ijkl} = \mathfrak{b}^{hom}_{jikl} = \mathfrak{b}^{hom}_{klij}$;
- *there exists* $C_0^* > 0$ *such that for every* 3×3 *symmetric matrix, one has*

$$\mathfrak{B}^{hom}(\zeta,\zeta) = \mathfrak{b}^{hom}_{ijkl}\zeta_{ij}\zeta_{kl} \geq C_0^*|\zeta|^2.$$

Proof. By definition of the $\mathfrak{b}^{hom}_{ijkl}$'s, the symmetry of matrices $M^{ij} = M^{ji}$ and correctors $\chi^{ij} = \chi^{ji}$ we obtain the symmetries of the $\mathfrak{b}^{hom}_{ijkl}$'s.
From equality (9.9), Lemma G.3 and estimate (G.4) we have

$$\begin{aligned}\mathfrak{B}^{hom}(\zeta,\zeta) &= \int_{\mathcal{S}} \mathfrak{A}\left[\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} + \sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right]\cdot\left[\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} + \sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right] d\mathbf{S}\\ &\geq C_0'\int_{\mathcal{S}}\left|\begin{pmatrix}(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1\\ \cdots\\ 0\end{pmatrix} + \sum_{i,j=1}^{3}\zeta_{ij}\frac{\partial}{\partial\mathbf{S}}\begin{pmatrix}\overline{\chi}^{ij}\\ \cdots\\ \widehat{\boldsymbol{\chi}}^{ij}\end{pmatrix}\right|^2 d\mathbf{S} \geq C_0^*|\zeta|^2.\end{aligned}$$

□

Theorem 9.1 (The homogenized limit problem)**.** *The limit field* $\mathcal{U} \in H^1_\Gamma(\Omega)^3$ *is the unique solution to the homogenized problem*

$$\begin{aligned}\int_\Omega \mathfrak{b}^{hom}_{ijkl}\,e_{ij}(\mathcal{U})\,e_{kl}(\mathcal{V})\,dx = &-\frac{4\pi}{5}\sum_{q=1}^{3}\int_\Omega G_q\,\mathfrak{c}^{hom}_{ijq}\,e_{ij}(\mathcal{V})\,dx + \frac{\pi}{1+\kappa}\sum_{q=1}^{3}\int_\Omega \mathbf{f}_q\,\mathfrak{c}^{hom}_{ijq+3}\,e_{ij}(\mathcal{V})\,dx\\ &+\frac{4\pi|\mathcal{K}|}{3}\int_\Omega F\cdot\mathcal{V}\,dx + \frac{\kappa|\mathcal{S}|}{1+\kappa}\int_\Omega \mathbf{f}\cdot\mathcal{V}\,dx, \qquad \forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\end{aligned} \tag{9.13}$$

where the $\mathfrak{b}^{hom}_{ijkl}$ *are given by* (9.10) *and the* $\mathfrak{c}^{hom}_{ijq}$ *by* (9.11).

10 The case of an isotropic and homogeneous material

We consider an isotropic and homogeneous material for which the relation between the linearized strain tensor and the stress tensor is given as follows

$$\sigma(u) = \lambda\,\mathrm{Tr}(e(u))\,\mathbf{I}_3 + 2\mu\,e(u), \tag{10.1}$$

where $\mathbf{I}_3$ is the unit 3×3 matrix and λ, μ are the material Lamé constants.
The correctors $\widetilde{\chi}_q \in L^\infty(\mathcal{S};\boldsymbol{\mathcal{D}}_w)$, $q\in\{1,2,3,4\}$, have the following form (see [13])

$$\begin{aligned}&\widetilde{\chi}_1(\cdot,S_2,S_3) = -\nu\big(S_2\mathbf{t}_2 + S_3\mathbf{t}_3\big), && \widetilde{\chi}_2(\cdot,S_2,S_3) = 0,\\ &\widetilde{\chi}_3(\cdot,S_2,S_3) = \nu\Big(-S_2S_3\mathbf{t}_2 + \frac{S_3^2-S_2^2}{2}\mathbf{t}_3\Big), && \widetilde{\chi}_4(\cdot,S_2,S_3) = \nu\Big(\frac{S_2^2-S_3^2}{2}\mathbf{t}_2 + S_2S_3\mathbf{t}_3\Big),\end{aligned}$$

where $\nu = \dfrac{\lambda}{2(\mu+\lambda)}$ is the Poisson coefficient.
Due to the symmetries of the elasticity coefficients and cross-sections, we have immediately

$$\widetilde{\chi}_3(\cdot,S_2,S_3) = -\widetilde{\chi}_4(\cdot,S_3,S_2).$$

Hence, we obtain

$$\begin{aligned}\widetilde{u} = \nu\Big[-\Big(\big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1\Big)\big(S_2\mathbf{t}_2+S_3\mathbf{t}_3\big) + \frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\mathbf{t}_2\Big(\frac{S_2^2-S_3^2}{2}\mathbf{t}_2 + S_2S_3\mathbf{t}_3\Big)\\ +\frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\mathbf{t}_3\Big(S_2S_3\mathbf{t}_2 + \frac{S_3^2-S_2^2}{2}\mathbf{t}_3\Big)\Big] \quad \text{a.e. in } \Omega\times\mathcal{S}\times D.\end{aligned} \tag{10.2}$$

The matrix $\mathfrak{A}$ becomes

$$\mathfrak{A} = \begin{pmatrix} \pi E & 0 & 0 & 0 \\ 0 & \frac{\pi}{2}\mu & 0 & 0 \\ 0 & 0 & \frac{\pi}{4}E & 0 \\ 0 & 0 & 0 & \frac{\pi}{4}E \end{pmatrix}, \tag{10.3}$$

where $E = \dfrac{\mu(3\lambda+2\mu)}{\lambda+\mu}$ is the Young's modulus.

The correctors $\chi^{ij} \doteq \big(\overline{\chi}^{ij}, \widehat{\chi}^{ij}, \widehat{\boldsymbol{\chi}}^{ij}\big) \in L^\infty\big(\Omega; \boldsymbol{\mathcal{D}}_{Ex} \times \boldsymbol{\mathcal{D}}_{In}\big)$, $(i,j) \in \{1,2,3\}^2$.
These correctors are the solutions to the variational problems $(9.6)_1$. Hence, by virtue of (10.3), we have

$$\int_{\mathcal{S}} \Big(\mathfrak{a}_{1,1} \frac{d\overline{\chi}^{ij}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, \frac{d\overline{\mathcal{V}}}{d\mathbf{S}} \cdot \mathbf{t}_1 + \sum_{q=1}^{3} \mathfrak{a}_{q+1,q+1} \frac{d\widehat{\boldsymbol{\chi}}^{ij}}{d\mathbf{S}} \cdot \mathbf{t}_q \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_q \Big) d\mathbf{S} = - \int_{\mathcal{S}} \mathfrak{a}_{1,1} (\mathbf{M}^{ij}\, \mathbf{t}_1) \cdot \mathbf{t}_1 \frac{d\overline{\mathcal{V}}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, d\mathbf{S}. \tag{10.4}$$

Choosing the function $\big(0, \widehat{\chi}^{ij}, \widehat{\boldsymbol{\chi}}^{ij}\big)$ as a test function we obtain

$$\int_{\mathcal{S}} \sum_{q=1}^{3} \mathfrak{a}_{q+1,q+1} \frac{d\widehat{\boldsymbol{\chi}}^{ij}}{d\mathbf{S}} \cdot \mathbf{t}_q \, \frac{d\widehat{\boldsymbol{\chi}}^{ij}}{d\mathbf{S}} \cdot \mathbf{t}_q \, d\mathbf{S} = 0.$$

Hence, for every $(i,j) \in \{1,2,3\}^2$ one has $\big(\widehat{\chi}^{ij}, \widehat{\boldsymbol{\chi}}^{ij}\big) = (0,0)$.
Let ℓ be in $\{1, \dots, m\}$ and $\phi \in H_0^1(\gamma_\ell)$. Consider the test function $\overline{\mathcal{V}} \in \boldsymbol{\mathcal{D}}_{Ex}$ defined by

$$\overline{\mathcal{V}} = \begin{cases} \phi\, \mathbf{t}_1^\ell & \text{on } \ \gamma_\ell, \\ 0 & \text{on the other segments of } \mathcal{S}. \end{cases}$$

That gives

$$\int_{\gamma_\ell} \frac{d\overline{\chi}^{ij}}{dS_1} \cdot \mathbf{t}_1^\ell \, \frac{d\phi}{dS_1} \, d\mathbf{S} = - \int_{\gamma_\ell} (\mathbf{M}^{ij}\, \mathbf{t}_1^\ell) \cdot \mathbf{t}_1^\ell \frac{d\phi}{dS_1} \, d\mathbf{S}$$

and then

$$\frac{d^2\overline{\chi}^{ij}}{dS_1^2} \cdot \mathbf{t}_1^\ell = 0 \quad \text{in } \ H^{-1}(\gamma_\ell).$$

It means that $\overline{\chi}^{ij} \cdot \mathbf{t}_1$ is affine on every segment of $\mathcal{S}$. The function $\overline{\chi}^{ij}$ belongs to $\mathbf{U}_{\mathcal{S}}$. Set

$$\mathbf{U}_{\mathcal{S},per,0} \doteq \mathbf{U}_{\mathcal{S}} \cap H^1_{per,0}(\mathcal{S})^3.$$

For every $(i,j) \in \{1,2,3\}^2$ one has

$$\overline{\chi}^{ij} \in \mathbf{U}_{\mathcal{S},per,0}.$$

Denote $\overline{\mathbf{M}}^{ij}$ the restriction to $\mathcal{S}$ of the linear field $x \in \mathbb{R}^3 \longmapsto \mathbf{M}^{ij} x \in \mathbb{R}^3$. It belongs to $\mathbf{U}_{\mathcal{S}}$. Problem (10.4) becomes

$$\int_{\mathcal{S}} \frac{d}{d\mathbf{S}} \big(\overline{\chi}^{ij} + \overline{\mathbf{M}}^{ij}\big) \cdot \mathbf{t}_1 \, \frac{d\overline{\mathcal{V}}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, d\mathbf{S} = 0, \qquad \forall\, \overline{\mathcal{V}} \in \mathbf{U}_{\mathcal{S},per,0}. \tag{10.5}$$

The corrector $\overline{\chi}^{ij}$ is the projection on $\mathbf{U}_{\mathcal{S},per,0}$ of the field $\overline{\mathbf{M}}^{ij} \in \mathbf{U}_{\mathcal{S}}$ for the scalar product $< \cdot, \cdot >_1$ (see (2.2) and Lemma 2.1).

The correctors: $\chi^q \doteq \big(\overline{\chi}^q, \widehat{\chi}^q, \widehat{\boldsymbol{\chi}}^q\big) \in L^\infty\big(\Omega; \boldsymbol{\mathcal{D}}_{Ex} \times \boldsymbol{\mathcal{D}}_{In}\big)$, $q \in \{1,2,3\}$.
They are the solution to the following variational problems $(9.6)_2$. Hence, by virtue (10.3), we have

$$\int_{\mathcal{S}} \Big(\mathfrak{a}_{1,1} \frac{d\overline{\chi}^{q}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, \frac{d\overline{\mathcal{V}}}{d\mathbf{S}} \cdot \mathbf{t}_1 + \sum_{i=1}^{3} \mathfrak{a}_{i+1,i+1} \frac{d\widehat{\boldsymbol{\chi}}^{q}}{d\mathbf{S}} \cdot \mathbf{t}_i \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_i \Big) d\mathbf{S} = \mathbf{e}_q \cdot \sum_{A \in \mathcal{K}} \widehat{\mathcal{B}}(A). \tag{10.6}$$

Choosing the function $\big(\overline{\chi}^q, 0, 0\big)$ as a test function we obtain

$$\int_{\mathcal{S}} \frac{d\overline{\chi}^{q}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, \frac{d\overline{\chi}^{q}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, d\mathbf{S} = 0.$$

Hence, for every $q \in \{1,2,3\}$ one has $\overline{\chi}^q = 0$, since this function belongs to $\boldsymbol{\mathcal{D}}_{Ex}$.
Let ℓ be in $\{1,\dots,m\}$ and $\phi_1 \in H_0^1(\gamma_\ell)$, $\phi_2,\ \phi_3 \in H_0^2(\gamma_\ell)$. Consider the test function defined by

$$\widehat{\mathcal{V}} = \begin{cases} \phi_2 \mathbf{t}_2^\ell + \phi_3 \mathbf{t}_3^\ell \quad \text{on } \gamma_\ell, \\ 0 \quad \text{on the other segments of } \mathcal{S}, \end{cases} \qquad \widehat{\mathcal{B}} = \begin{cases} \phi_1 \mathbf{t}_1^\ell - \dfrac{d\widehat{\mathcal{V}}}{d\mathbf{S}} \wedge \mathbf{t}_1^\ell \qquad \text{on } \gamma_\ell, \\ 0 \quad \text{on the other segments of } \mathcal{S}. \end{cases} \tag{10.7}$$

The couple $(\widehat{\mathcal{V}}, \widehat{\mathcal{B}})$ belongs to $\boldsymbol{\mathcal{D}}_{In}$. Choosing this couple as a test function in (10.5) leads to

$$\frac{d^2\widehat{\boldsymbol{\chi}}^q}{dS_1^2} \cdot \mathbf{t}_1^\ell = 0 \quad \text{in } \ H^{-1}(\gamma_\ell), \qquad \frac{d^3\widehat{\boldsymbol{\chi}}^q}{dS_1^3} \cdot \mathbf{t}_2^\ell = \frac{d^3\widehat{\boldsymbol{\chi}}^q}{dS_1^3} \cdot \mathbf{t}_3^\ell = 0 \quad \text{in } \ H^{-2}(\gamma_\ell).$$

Hence, for every $\ell \in \{1,\dots,m\}$ $\widehat{\boldsymbol{\chi}}^q \cdot \mathbf{t}_1^\ell$ is an affine function on γ_ℓ, while $\widehat{\boldsymbol{\chi}}^q \cdot \mathbf{t}_2^\ell$ and $\widehat{\boldsymbol{\chi}}^q \cdot \mathbf{t}_3^\ell$ are polynomial functions of degree less than 2 on γ_ℓ. A straightforward calculation gives the restriction of $\widehat{\boldsymbol{\chi}}^q$ to the segment γ_ℓ ($S_1 \in [0, l_\ell]$)

$$\begin{aligned} \widehat{\boldsymbol{\chi}}^q(S_1) = \widehat{\boldsymbol{\chi}}^q(A)\Big(1 - \frac{S_1}{l_\ell}\Big) + \widehat{\boldsymbol{\chi}}^q(B)\frac{S_1}{l_\ell} \\ - 3\Big(\big(\widehat{\boldsymbol{\chi}}^q(A) + \widehat{\boldsymbol{\chi}}^q(B)\big) - \big(\widehat{\boldsymbol{\chi}}^q(A) + \widehat{\boldsymbol{\chi}}^q(B)\big) \cdot \mathbf{t}_1^\ell \ \cdot \mathbf{t}_1^\ell\Big)\frac{S_1}{l_\ell}\Big(1 - \frac{S_1}{l_\ell}\Big) \end{aligned} \tag{10.8}$$

since $\displaystyle\int_{\gamma_\ell} \widehat{\boldsymbol{\chi}}^q \cdot \mathbf{t}_2^\ell \, d\mathbf{S} = \int_{\gamma_\ell} \widehat{\boldsymbol{\chi}}^q \cdot \mathbf{t}_3^\ell \, d\mathbf{S} = 0$. Then, an integration gives

$$\widehat{\chi}^q(S_1) = \widehat{\boldsymbol{\chi}}^q(A) \wedge \mathbf{t}_1^\ell \, S_1\Big(1 - \frac{S_1}{2l_\ell}\Big) + \widehat{\boldsymbol{\chi}}^q(B) \wedge \mathbf{t}_1^\ell \frac{S_1^2}{2l_\ell} - 3\Big(\widehat{\boldsymbol{\chi}}^q(A) + \widehat{\boldsymbol{\chi}}^q(B)\Big) \wedge \mathbf{t}_1^\ell \frac{S_1^2}{l_\ell}\Big(\frac{1}{2} - \frac{S_1}{3l_\ell}\Big), \tag{10.9}$$

since $\widehat{\chi}^q(A) = \widehat{\chi}^q(B) = 0$.

The correctors: $\chi^{q+3} \doteq \big(\overline{\chi}^{q+3}, \widehat{\chi}^{q+3}, \widehat{\boldsymbol{\chi}}^{q+3}\big) \in L^\infty\big(\Omega; \boldsymbol{\mathcal{D}}_{Ex} \times \boldsymbol{\mathcal{D}}_{In}\big)$, $q \in \{1,2,3\}$**.**
They are the solution to the variational problems $(9.6)_3$. Hence by virtue (10.3) we have

$$\int_{\mathcal{S}} \Big(\mathfrak{a}_{1,1} \frac{d\overline{\chi}^{q+3}}{d\mathbf{S}} \cdot \mathbf{t}_1 \, \frac{d\overline{\mathcal{V}}}{d\mathbf{S}} \cdot \mathbf{t}_1 + \sum_{i=1}^{3} \mathfrak{a}_{i+1,i+1} \frac{d\widehat{\boldsymbol{\chi}}^{q+3}}{d\mathbf{S}} \cdot \mathbf{t}_i \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_i\Big) d\mathbf{S} = \mathbf{e}_q \cdot \int_{\mathcal{S}} \widehat{\mathcal{V}} \, d\mathbf{S}. \tag{10.10}$$

As in the previous case, for every $q \in \{1,2,3\}$ one obtains $\overline{\chi}^{q+3} = 0$.
Again, we consider the test function defined by (10.7). That leads to ($\ell \in \{1,\dots,m\}$)

$$\frac{d^2\widehat{\boldsymbol{\chi}}^{q+3}}{dS_1^2} \cdot \mathbf{t}_1^\ell = 0 \text{ in } \ H^{-1}(\gamma_\ell), \qquad \frac{d^3\widehat{\boldsymbol{\chi}}^{q+3}}{dS_1^3} \cdot \mathbf{t}_2^\ell = \frac{4}{\pi E}\mathbf{e}_q \cdot \mathbf{t}_2^\ell, \quad \frac{d^3\widehat{\boldsymbol{\chi}}^{q+3}}{dS_1^3} \cdot \mathbf{t}_3^\ell = \frac{4}{\pi E}\mathbf{e}_q \cdot \mathbf{t}_3^\ell \quad \text{in } \ H^{-2}(\gamma_\ell). \tag{10.11}$$

Hence, for every $\ell \in \{1,\dots,m\}$, the restriction of $\widehat{\boldsymbol{\chi}}^{q+3}$ to the segment γ_ℓ is ($S_1 \in [0, l_\ell]$)

$$\begin{aligned} \widehat{\boldsymbol{\chi}}^{q+3}(S_1) = \widehat{\boldsymbol{\chi}}^{q+3}(A)\Big(1 - \frac{S_1}{l_\ell}\Big) + \widehat{\boldsymbol{\chi}}^{q+3}(B)\frac{S_1}{l_\ell} - 3\Big(\big(\widehat{\boldsymbol{\chi}}^{q+3}(A) + \widehat{\boldsymbol{\chi}}^{q+3}(B)\big) \\ - \big(\widehat{\boldsymbol{\chi}}^{q+3}(A) + \widehat{\boldsymbol{\chi}}^{q+3}(B)\big) \cdot \mathbf{t}_1^\ell \, \mathbf{t}_1^\ell\Big)\frac{S_1}{l_\ell}\Big(1 - \frac{S_1}{l_\ell}\Big) + \frac{2l_\ell^3}{3\pi E}\big(\mathbf{e}_q - \mathbf{e}_q \cdot \mathbf{t}_1^\ell \, \mathbf{t}_1^\ell\big)\frac{S_1}{l_\ell}\Big(\frac{1}{2} - \frac{S_1}{l_\ell}\Big)\Big(1 - \frac{S_1}{l_\ell}\Big). \end{aligned} \tag{10.12}$$

Then, integrating leads to

$$\begin{aligned} \widehat{\chi}^{q+3}(S_1) = \widehat{\boldsymbol{\chi}}^{q+3}(A) \wedge \mathbf{t}_1^\ell \, S_1\Big(1 - \frac{S_1}{2l_\ell}\Big) + \widehat{\boldsymbol{\chi}}^{q+3}(B) \wedge \mathbf{t}_1^\ell \frac{S_1^2}{2l_\ell} \\ - 3\Big(\widehat{\boldsymbol{\chi}}^{q+3}(A) + \widehat{\boldsymbol{\chi}}^{q+3}(B)\Big) \wedge \mathbf{t}_1^\ell \frac{S_1^2}{l_\ell}\Big(\frac{1}{2} - \frac{S_1}{3l_\ell}\Big) - \frac{l_\ell^4}{6\pi E}\frac{S_1^2}{l_\ell^2}\Big(1 - \frac{S_1}{l_\ell}\Big)^2 \mathbf{e}_q \wedge \mathbf{t}_1^\ell. \end{aligned} \tag{10.13}$$

The last step allows us to reduce the corrector problems $(9.6)_{2,3}$ to the algebraic equations with respect to the unknown vector of nodal values. Denote $\mathbf{E}_q$ the function belonging to $H^1_{per,0}(\mathcal{S})^3$ and defined by ($\ell \in \{1,\dots,m\}$)

$$\mathbf{E}_q(S_1) = \frac{l_\ell^4}{6\pi E}\frac{S_1^2}{l_\ell^2}\Big(1 - \frac{S_1}{l_\ell}\Big)^2 \mathbf{e}_q \wedge \mathbf{t}_1 \quad \text{on } \ \gamma_\ell, \ \ (S_1 \in [0, l_\ell]).$$

Set

$$\mathcal{P}_{per}(\mathcal{S}) = \Big\{ \widehat{\mathcal{B}} \in H^1_{per}(\mathcal{S})^3 \mid \widehat{\mathcal{B}}(S_1) = \widehat{\mathcal{B}}(A)\Big(1 - \frac{S_1}{l_\ell}\Big) + \widehat{\mathcal{B}}(B)\frac{S_1}{l_\ell} \\ - 3\Big[\big(\widehat{\mathcal{B}}(A) + \widehat{\mathcal{B}}(B)\big) - \big(\big(\widehat{\mathcal{B}}(A) + \widehat{\mathcal{B}}(B)\big)\cdot \mathbf{t}_1^\ell\big)\mathbf{t}_1^\ell\Big]\frac{S_1}{l_\ell}\Big(1 - \frac{S_1}{l_\ell}\Big), \\ \text{on } \ \gamma_\ell = [A,B], \quad S_1 \in [0, l_\ell], \ \ \ell \in \{1,\dots,m\}\Big\}.$$

So $\widehat{\boldsymbol{\chi}}^q$, $(q \in \{1,2,3\})$, belongs to $\mathcal{P}_{per}(\mathcal{S})$ and solves the discrete problem

$$\widehat{\boldsymbol{\chi}}^q \in \mathcal{P}_{per}(\mathcal{S}), \quad q \in \{1,2,3\}, \\ \int_{\mathcal{S}} \Big(\sum_{i=1}^3 \mathfrak{a}_{i+1,i+1} \frac{d\widehat{\boldsymbol{\chi}}^q}{d\mathbf{S}} \cdot \mathbf{t}_i \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_i \Big) d\mathbf{S} = \mathbf{e}_q \cdot \sum_{A \in \mathcal{K}} \widehat{\mathcal{B}}(A), \quad \forall\, \mathcal{B} \in \mathcal{P}_{per}(\mathcal{S}). \tag{10.14}$$

Similarly $\widehat{\boldsymbol{\chi}}^{q+3} + \dfrac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1$, $(q \in \{1,2,3\})$, belongs to $\mathcal{P}_{per}(\mathcal{S})$ and solves the discrete problem

$$\widehat{\boldsymbol{\chi}}^{q+3} + \frac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1 \in \mathcal{P}_{per}(\mathcal{S}), \quad q \in \{1,2,3\}, \\ \int_{\mathcal{S}} \sum_{i=1}^3 \mathfrak{a}_{i+1,i+1} \frac{d}{d\mathbf{S}} \Big(\widehat{\boldsymbol{\chi}}^{q+3} + \frac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1 \Big) \cdot \mathbf{t}_i \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_i \, d\mathbf{S} = \int_{\mathcal{S}} \mathbf{e}_q \cdot \widehat{\mathcal{V}} \, d\mathbf{S} \\ + \frac{\pi E}{4} \int_{\mathcal{S}} \Big(\frac{d^2\mathbf{E}_q}{d\mathbf{S}^2} \wedge \mathbf{t}_1 \Big) \cdot \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \, d\mathbf{S}, \qquad \forall\, \mathcal{B} \in \mathcal{P}_{per}(\mathcal{S}). \tag{10.15}$$

One has

$$\frac{\pi E}{4} \int_{\mathcal{S}} \Big(\frac{d^2\mathbf{E}_q}{d\mathbf{S}^2} \wedge \mathbf{t}_1 \Big) \cdot \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} d\mathbf{S} = \sum_{\ell=1}^m \frac{l_\ell^4}{24} \int_{\gamma_\ell} \frac{d^2}{dS_1^2} \Big(\frac{S_1^2}{l_\ell^2} \Big(1 - \frac{S_1}{l_\ell} \Big)^2 \Big) (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \, d\mathbf{S} \\ = \sum_{\ell=1}^m \frac{l_\ell^2}{12} (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \big(\widehat{\mathcal{B}}(\mathbf{B}) - \widehat{\mathcal{B}}(\mathbf{A}) \big) - \sum_{\ell=1}^m \frac{l_\ell^4}{24} \int_{\gamma_\ell} \frac{d^3}{dS_1^3} \Big(\frac{S_1^2}{l_\ell^2} \Big(1 - \frac{S_1}{l_\ell} \Big)^2 \Big) (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \widehat{\mathcal{B}} \, d\mathbf{S} \\ = \sum_{\ell=1}^m \frac{l_\ell^2}{12} (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \big(\widehat{\mathcal{B}}(\mathbf{B}) - \widehat{\mathcal{B}}(\mathbf{A}) \big) - \sum_{\ell=1}^m \int_{\gamma_\ell} \Big(S_1 - \frac{l_\ell}{2} \Big) (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \widehat{\mathcal{B}} \, d\mathbf{S}$$

and

$$\int_{\mathcal{S}} \mathbf{e}_q \cdot \widehat{\mathcal{V}} \, d\mathbf{S} = \sum_{\ell=1}^m \int_{\gamma_\ell} \mathbf{e}_q \cdot \widehat{\mathcal{V}} \, d\mathbf{S} = - \sum_{\ell=1}^m \int_{\gamma_\ell} \Big(S_1 - \frac{l_\ell}{2} \Big) \mathbf{e}_q \cdot \frac{d\widehat{\mathcal{V}}}{d\mathbf{S}} \, d\mathbf{S} = - \sum_{\ell=1}^m \int_{\gamma_\ell} \Big(S_1 - \frac{l_\ell}{2} \Big) \mathbf{e}_q \cdot \big(\widehat{\mathcal{B}} \wedge \mathbf{t}_1 \big) \, d\mathbf{S} \\ = \sum_{\ell=1}^m \int_{\gamma_\ell} \Big(S_1 - \frac{l_\ell}{2} \Big) (\mathbf{e}_q \wedge \mathbf{t}_1) \cdot \widehat{\mathcal{B}} \, d\mathbf{S}.$$

Hence $\widehat{\boldsymbol{\chi}}^{q+3} + \dfrac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1$, $(q \in \{1,2,3\})$ are solutions of the discrete problem

$$\widehat{\boldsymbol{\chi}}^{q+3} + \frac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1 \in \mathcal{P}_{per}(\mathcal{S}), \quad q \in \{1,2,3\}, \\ \int_{\mathcal{S}} \sum_{i=1}^3 \mathfrak{a}_{i+1,i+1} \frac{d}{d\mathbf{S}} \Big(\widehat{\boldsymbol{\chi}}^{q+3} + \frac{d\mathbf{E}_q}{d\mathbf{S}} \wedge \mathbf{t}_1 \Big) \cdot \mathbf{t}_i \, \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \cdot \mathbf{t}_i \, d\mathbf{S} = - \sum_{\ell=1}^m \frac{l_\ell^2}{12} \Big(\int_{\gamma_\ell} \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \wedge \mathbf{t}_1 d\mathbf{S} \Big) \cdot \mathbf{e}_q \\ \forall\, \mathcal{B} \in \mathcal{P}_{per}(\mathcal{S}). \tag{10.16}$$

11 Conclusion

We conclude, that for our ε-periodic r-thin structure, the solution to the linearized elasticity problem (5.7) (in the strong), or (5.8) (in the weak/variational form) can be reconstructed in the following form:

$$u_\varepsilon(x) \approx \frac{\varepsilon}{r} \mathcal{U}(x) + \frac{4\pi}{5} \frac{\varepsilon^3}{r^2} \sum_{q=1}^3 G_q(x) \, \widehat{\chi}^q\Big(\Big\{ \frac{x}{\varepsilon} \Big\}_{|\mathcal{S}} \Big) + \frac{\varepsilon^5}{r^2(\varepsilon^2 + r)} \pi \sum_{q=1}^3 \mathbf{f}_q(x) \, \widehat{\chi}^{q+3}\Big(\Big\{ \frac{x}{\varepsilon} \Big\}_{|\mathcal{S}} \Big) + O\Big(\frac{\varepsilon^2}{r} \Big), \\ \text{for a.e. } x \in \mathcal{S}_{\varepsilon,r}. \tag{11.1}$$

From Proposition 7.2 we have

$$e_s(u_\varepsilon) \approx \frac{\varepsilon}{r}\Big(\mathcal{E}(\mathcal{U}) + \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) + \mathcal{E}_D(\widetilde{u})\Big).$$

$$\mathcal{E}(\mathcal{U}) \doteq \begin{pmatrix} \big(e(\mathcal{U})\,\mathbf{t}_1\big)\cdot\mathbf{t}_1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \qquad \mathcal{E}_S(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) \doteq \begin{pmatrix} \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 - \frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\Big(S_2\,\mathbf{t}_2 + S_3\,\mathbf{t}_3\Big) & * & * \\ -\frac{S_3}{2}\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 & 0 & 0 \\ \frac{S_2}{2}\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 & 0 & 0 \end{pmatrix},$$

$$\mathcal{E}_D(\widetilde{u}) \doteq \begin{pmatrix} 0 & \frac{1}{2}\frac{\partial\widetilde{u}}{\partial S_2}\cdot\mathbf{t}_1 & \frac{1}{2}\frac{\partial\widetilde{u}}{\partial S_3}\cdot\mathbf{t}_1 \\ * & \frac{\partial\widetilde{u}}{\partial S_2}\cdot\mathbf{t}_2 & \frac{1}{2}\frac{\partial\widetilde{u}}{\partial S_3}\cdot\mathbf{t}_2 + \frac{1}{2}\frac{\partial\widetilde{u}}{\partial S_2}\cdot\mathbf{t}_3 \\ * & * & \frac{\partial\widetilde{u}}{\partial S_3}\cdot\mathbf{t}_3 \end{pmatrix} \qquad \text{a.e. in } \Omega\times\mathcal{S}\times D,$$

where

$$(\overline{\mathcal{U}},\widehat{\mathcal{U}},\widehat{\mathcal{R}}) = \sum_{i,j=1}^{3} e_{ij}(\mathcal{U})\,(\overline{\chi}^{ij},0,0) + \frac{4\pi}{5}\sum_{q=1}^{3} G_q\big(0,\widehat{\chi}^q,\widehat{\boldsymbol{\chi}}^q\big) + \frac{\pi}{1+\kappa}\sum_{q=1}^{3}\mathbf{f}_q\big(0,\widehat{\chi}^{q+3},\widehat{\boldsymbol{\chi}}^{q+3}\big).$$

and

$$\widetilde{u} = \nu\Big[\frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\mathbf{t}_2\Big(\frac{S_2^2-S_3^2}{2}\mathbf{t}_2 + S_2S_3\mathbf{t}_3\Big) + \frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\mathbf{t}_3\Big(S_2S_3\mathbf{t}_2 + \frac{S_3^2-S_2^2}{2}\mathbf{t}_3\Big)\Big].$$

The strain tensor in the global coordinates can be obtained using (5.3). Then, we can reconstruct the local stress field for $\mathcal{P}^{\xi}_{\varepsilon\ell,r}$ beam as follows

$$\sigma_s(u_\varepsilon) \approx \frac{\varepsilon}{r}\begin{pmatrix} E\Big(\big(e(\mathcal{U})\,\mathbf{t}_1 + \frac{\partial\overline{\mathcal{U}}}{\partial\mathbf{S}}\big)\cdot\mathbf{t}_1 - \frac{\partial^2\widehat{\mathcal{U}}}{\partial\mathbf{S}^2}\cdot\big(S_2\,\mathbf{t}_2 + S_3\,\mathbf{t}_3\big)\Big) & * & * \\ -S_3\mu\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 & 0 & * \\ S_2\mu\frac{\partial\widehat{\mathcal{R}}}{\partial\mathbf{S}}\cdot\mathbf{t}_1 & 0 & 0 \end{pmatrix} \qquad \text{a.e. in } \Omega\times\mathcal{S}\times D. \tag{11.2}$$

A Proof of Proposition 5.2

Lemma A.1. *Let $\mathcal{S}$ be a* **3-PSS**. *For every u in $\mathbf{V}_{\varepsilon,r}$, one has*

$$\begin{aligned} &\|\boldsymbol{\mathcal{U}}\|_{L^2(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad &&\|\nabla\boldsymbol{\mathcal{U}}\|_{L^2(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \\ &\|\boldsymbol{\mathcal{R}}\|_{L^2(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad &&\|\nabla\boldsymbol{\mathcal{R}}\|_{L^2(\Omega_\varepsilon'^{int})} \le C\frac{1}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}. \end{aligned} \tag{A.1}$$

Proof. Since u belongs to $\mathbf{V}_{\varepsilon,r}$, by definition, it is equal to 0 in $\mathcal{S}'_{\varepsilon,r}\setminus\overline{\mathcal{S}_{\varepsilon,r}}$. Then, there exists a rigid displacement $\mathbf{r}'(x) = \mathbf{a}' + \mathbf{b}'\wedge x$, $(\mathbf{a}',\mathbf{b}')\in\mathbb{R}^3\times\mathbb{R}^3$ such that (using (5.17) with Ω' instead of Ω)

$$\|\boldsymbol{\mathcal{U}} - \mathbf{r}'\|_{H^1(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.$$

Let $\mathcal{O}$ be an open set satisfying $\mathcal{O}$ strictly included in $\Omega'\setminus\overline{\Omega}$.
If ε is small enough then $\mathcal{O}_\varepsilon = \big\{x\in\mathbb{R}^3 \mid \mathrm{dist}(x,\partial\mathcal{O}) < 2\sqrt{3}\varepsilon\big\} \subset \Omega_\varepsilon'^{int}\setminus\overline{\Omega_\varepsilon^{int}}$. As a consequence $\boldsymbol{\mathcal{U}} = \boldsymbol{\mathcal{R}} = 0$ a.e. in $\mathcal{O}$. Hence,

$$\|\mathbf{r}'\|_{H^1(\mathcal{O})} \le \|\boldsymbol{\mathcal{U}} - \mathbf{r}'\|_{H^1(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})} \quad\Longrightarrow\quad |\mathbf{a}'| + |\mathbf{b}'| \le C\|\mathbf{r}'\|_{H^1(\mathcal{O})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.$$

The constants do not depend on ε and r. Therefore,

$$\|\mathbf{r}'\|_{H^1(\Omega_\varepsilon'^{int})} \le C_0(|\mathbf{a}'| + |\mathbf{b}'|) \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},$$

where the constant C_0 only depends on the volume and diameter of Ω'. Finally,

$$\|\boldsymbol{\mathcal{U}}\|_{H^1(\Omega_\varepsilon^{int})} \le \|\boldsymbol{\mathcal{U}}\|_{H^1(\Omega_\varepsilon'^{int})} \le \|\boldsymbol{\mathcal{U}} - \mathbf{r}'\|_{H^1(\Omega_\varepsilon'^{int})} + \|\mathbf{r}'\|_{H^1(\Omega_\varepsilon'^{int})} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}$$

and $(\mathrm{A}.1)_{1,2}$ are proved. Estimates $(\mathrm{A}.1)_{3,4}$ follow from $(\mathrm{A}.1)_{1,2}$ and $(5.16)_{1,2}$. □

Lemma A.2. *Let $\mathcal{S}$ be a* **3-PSS**. *One has (see* (5.14) *for* $\mathbf{a}$ *and* $\mathbf{b}$*)*

$$\begin{aligned}
&\|\mathbf{a}\|_{L^2(\Omega_\varepsilon)} + \|\mathbf{b}\|_{L^2(\Omega_\varepsilon)} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\|\boldsymbol{\mathcal{U}} - \mathbf{a}\|_{L^2(\Omega_\varepsilon^{int})} \le C\frac{\varepsilon^2}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad \|\boldsymbol{\mathcal{R}} - \mathbf{b}\|_{L^2(\Omega_\varepsilon^{int})} \le \frac{C\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\|\mathcal{U}^h - \mathbf{a}\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad \Big\|\frac{d\mathcal{U}^h}{d\mathbf{s}} - \mathbf{b}\wedge\mathbf{t}_1\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})},\\
&\|\mathcal{U}^h\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{1}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}, \qquad \Big\|\frac{d\mathcal{U}^h}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \frac{C}{r}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned} \tag{A.2}$$

Proof. From estimates $(5.13)_1$, (5.15), $(A.1)_3$ and the definition of $\boldsymbol{\mathcal{R}}$ we obtain

$$\sum_{\xi\in\Xi_\varepsilon} |\mathbf{b}(\varepsilon\xi)|^2\varepsilon^3 \le C\frac{\varepsilon^2}{r^2}\|e(u)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})}. \tag{A.3}$$

Then, from the above estimate and $(5.13)_2$ we obtain

$$\sum_{i=1}^3\sum_{\xi\in\Xi_{\varepsilon,i}} |\mathbf{a}(\varepsilon\xi+\varepsilon\mathbf{e}_i) - \mathbf{a}(\varepsilon\xi)|^2\varepsilon^3 \le C\frac{\varepsilon^4}{r^2}\|e(u)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})}, \tag{A.4}$$

which in turn with (5.15), $(A.1)_1$ and the definition of $\boldsymbol{\mathcal{U}}$ lead to

$$\sum_{\xi\in\Xi_\varepsilon} |\mathbf{a}(\varepsilon\xi)|^2\varepsilon^3 \le C\frac{\varepsilon^2}{r^2}\|e(u)\|^2_{L^2(\mathcal{S}_{\varepsilon,r})}. \tag{A.5}$$

Hence we have $(A.2)_1$. Estimate $(A.2)_2$ is the consequences of $(A.2)_1$, (A.4) and the definition of $\boldsymbol{\mathcal{U}}$ while $(A.2)_3$ follows from $(5.13)_1$ and the definition of $\boldsymbol{\mathcal{R}}$.
Estimate $(5.11)_1$ yields

$$\|\mathcal{U}^h - \mathbf{a}(\varepsilon\xi)\|^2_{L^2(\varepsilon(\xi+\mathcal{S}))} \le \frac{C\varepsilon^2}{r^2}\|e(u)\|^2_{L^2(\varepsilon(\xi+\mathcal{S}_{r/\varepsilon}))} + C|\mathbf{b}(\varepsilon\xi)|^2\varepsilon^3.$$

Summing up over all $\xi\in\Xi_\varepsilon$ the above inequality, using $(5.11)_2$ and applying (A.3) give $(A.2)_{4,5}$.
Inequalities $(A.2)_{6,7}$ are the immediate consequences of $(A.2)_{1,4,5}$. □

Proof of Proposition 5.2: Since $\mathcal{U} = \mathcal{U}^h + \overline{\mathcal{U}}$ we have

$$\|\mathcal{U}\|_{L^2(\mathcal{S}_\varepsilon)} \le \|\mathcal{U}^h\|_{L^2(\mathcal{S}_\varepsilon)} + \|\overline{\mathcal{U}}\|_{L^2(\mathcal{S}_\varepsilon)}, \qquad \Big\|\frac{d\mathcal{U}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \le \Big\|\frac{d\mathcal{U}^h}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} + \Big\|\frac{d\overline{\mathcal{U}}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)}.$$

From the estimates of Lemmas 5.1-A.2 we obtain $(5.18)_{1,2}$. Estimate (5.12) yields

$$\|\mathcal{R}^h\|^2_{L^2(\varepsilon(\xi+\mathcal{S}))} \le \frac{C\varepsilon^2}{r^4}\|e(u)\|^2_{L^2(\varepsilon(\xi+\mathcal{S}_{r/\varepsilon}))} + C|\mathbf{b}(\varepsilon\xi)|^2\varepsilon^2.$$

Summing up over all $\xi\in\Xi_\varepsilon$ and applying (A.3) give

$$\|\mathcal{R}^h\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{\varepsilon}{r^2}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.$$

Then, this inequality together with the estimates (5.10) yield $(5.18)_2$.
From Definition 4.1 we have

$$\begin{aligned}
&\|U^e\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le Cr\|\mathcal{U}\|_{L^2(\mathcal{S}_\varepsilon)} + Cr^2\|\mathcal{R}\|_{L^2(\mathcal{S}_\varepsilon)},\\
&\|\nabla U^e\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le C\Big(r\Big\|\frac{d\mathcal{U}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} + r^2\Big\|\frac{d\mathcal{R}}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} + r\|\mathcal{R}\|_{L^2(\mathcal{S}_\varepsilon)}\Big).
\end{aligned}$$

Then, the estimates $(5.18)_{1,2}$ and $(5.18)_2$ lead to $(5.18)_{3,4}$.

B The applied forces

First, note that the number of elements in $\mathcal{K}_\varepsilon$, which is denoted by $|\mathcal{K}_\varepsilon|$ is less than

$$|\mathcal{K}_\varepsilon| \le \frac{|\mathcal{K}| mes |\Omega_1|}{\varepsilon^3}, \tag{B.1}$$

where $|\mathcal{K}|$ is the number of elements in $\mathcal{K}$.

Proof of Lemma 5.4. Let u be in $\mathbf{V}_{\varepsilon,r}$. By the estimates of Proposition 5.2, we have

$$\Big| \sum_{\xi\in\Xi_\varepsilon} \int_{\mathcal{P}^\xi_{\varepsilon\ell,r}} \mathbf{f}_\varepsilon \cdot u \, dx \Big| \le C\|\mathbf{f}\|_{L^\infty(\Omega)} \|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}. \tag{B.2}$$

Now, taking into account that for every node $A \in \mathcal{K}_\varepsilon$ the following decomposition holds:

$$u(x) = \mathcal{U}(A) + \mathcal{R}(A) \wedge (x - A) + \overline{u}, \quad \text{for a.e. } x \in B(A,r),$$

we have

$$\begin{aligned}
&\int_{\mathcal{S}_{\varepsilon,r}} \Big(\sum_{A\in\mathcal{K}_\varepsilon} F(A)\mathbf{1}_{B(A,r)} + \sum_{A\in\mathcal{K}_\varepsilon} G(A)\wedge(x-A)\mathbf{1}_{B(A,r)} \Big) \cdot u \, dx \\
&= \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} F(A)\cdot\mathcal{U}(A)\,dx + \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} F(A)\cdot\big(\mathcal{R}(A)\wedge(x-A)\big)\,dx \\
&+ \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} \big(G(A)\wedge(x-A)\big)\cdot\mathcal{U}(A)\,dx + \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} \big(G(A)\wedge(x-A)\big)\cdot\big(\mathcal{R}(A)\wedge(x-A)\big)\,dx \\
&+ \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} F(A)\cdot\overline{u}\,dx + \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} \big(G(A)\wedge(x-A)\big)\cdot\overline{u}\,dx.
\end{aligned} \tag{B.3}$$

Let us estimate every integral in (B.3) separately. Due to the symmetries of the ball we have

$$\int_{B(A,r)} F(A)\cdot\big(\mathcal{R}(A)\wedge(x-A)\big)\,dx = 0, \qquad \int_{B(A,r)} \big(G(A)\wedge(x-A)\big)\cdot\mathcal{U}(A)\,dx = 0, \qquad \forall\, A\in\mathcal{K}.$$

Thus, the second and third terms in the right-hand side of (B.3) vanish. Then, using the Cauchy-Schwarz inequality, $(5.9)_1$ and (B.1), the last two integrals in (B.3) are estimated as follows

$$\begin{aligned}
&\Big| \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} F(A)\cdot\overline{u}\,dx \Big| \le \Big(\sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} |F(A)|^2\,dx \Big)^{1/2} \Big(\int_{\mathcal{S}_{\varepsilon,r}} |\overline{u}|^2\,dx \Big)^{1/2} \\
&\le C\Big(\frac{r}{\varepsilon}\Big)^{3/2} \|F\|_{L^\infty(\Omega)} \|\overline{u}\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le Cr\Big(\frac{r}{\varepsilon}\Big)^{3/2} \|F\|_{L^\infty(\Omega)} \|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}
\end{aligned}$$

and

$$\begin{aligned}
&\Big| \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} \big(G(A)\wedge(x-A)\big)\cdot\overline{u}\,dx \Big| \le \Big(\sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} r^2|G(A)|^2\,dx \Big)^{1/2} \Big(\int_{\mathcal{S}_{\varepsilon,r}} |\overline{u}|^2\,dx \Big)^{1/2} \\
&\le Cr\Big(\frac{r}{\varepsilon}\Big)^{3/2} \|G\|_{L^\infty(\Omega)} \|\overline{u}\|_{L^2(\mathcal{S}_{\varepsilon,r})} \le Cr^2\Big(\frac{r}{\varepsilon}\Big)^{3/2} \|G\|_{L^\infty(\Omega)} \|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned}$$

Since $\mathcal{U}^h(A) = \mathcal{U}(A)$ and $\mathcal{R}^h(A) = \mathcal{R}(A)$, then using the fact that $\mathcal{U}^h$, $\mathcal{R}^h$ are affine functions between two contiguous nodes

$$\sum_{A\in\mathcal{K}_\varepsilon} |\mathcal{U}^h(A)|^2\varepsilon \le C\|\mathcal{U}^h\|^2_{L^2(\mathcal{S}_\varepsilon)}, \quad \sum_{A\in\mathcal{K}_\varepsilon} |\mathcal{R}^h(A)|^2\varepsilon \le C\|\mathcal{R}^h\|^2_{L^2(\mathcal{S}_\varepsilon)}. \tag{B.4}$$

Then, the remaining two integrals in the right-hand side of (B.3) are estimated using (B.1), $(A.2)_6$, $(5.18)_2$ and (B.4)

$$\begin{aligned}
&\Big| \sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} F(A)\cdot\mathcal{U}(A)\,dx \Big| \le \Big(\sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} |F(A)|^2\,dx \Big)^{1/2} \Big(\sum_{A\in\mathcal{K}_\varepsilon} \int_{B(A,r)} |\mathcal{U}(A)|^2\,dx \Big)^{1/2} \\
&\le C\frac{r^{3/2}}{\varepsilon^{3/2}} \|F\|_{L^\infty(\Omega)} \Big(\frac{r^3}{\varepsilon} \sum_{A\in\mathcal{K}_\varepsilon} |\mathcal{U}^h(A)|^2\varepsilon \Big)^{1/2} \le C\frac{r^3}{\varepsilon^2} \|F\|_{L^\infty(\Omega)} \|\mathcal{U}^h\|_{L^2(\mathcal{S}_\varepsilon)} \le C\frac{r^2}{\varepsilon^2} \|F\|_{L^\infty(\Omega)} \|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}
\end{aligned}$$

and

$$\begin{aligned}
&\Big|\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}\big(G(A)\wedge(x-A)\big)\cdot\big(\mathcal{R}(A)\wedge(x-A)\big)\,dx\Big|\\
&\le\Big(\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}|G(A)\wedge(x-A)|^2\,dx\Big)^{1/2}\Big(\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}|\mathcal{R}(A)\wedge(x-A)|^2\,dx\Big)^{1/2}\\
&\le\Big(\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}r^2|G(A)|^2\,dx\Big)^{1/2}\Big(\sum_{A\in\mathcal{K}_\varepsilon}\int_{B(A,r)}r^2|\mathcal{R}(A)|^2\,dx\Big)^{1/2}\\
&\le C\frac{r^{5/2}}{\varepsilon^{3/2}}\|G\|_{L^\infty(\Omega)}\Big(\frac{r^5}{\varepsilon}\sum_{A\in\mathcal{K}_\varepsilon}|\mathcal{R}^h(A)|^2\varepsilon\Big)^{1/2}\le C\frac{r^5}{\varepsilon^2}\|G\|_{L^\infty(\Omega)}\|\mathcal{R}^h\|_{L^2(\mathcal{S}_\varepsilon)}\le C\frac{r^3}{\varepsilon}\|G\|_{L^\infty(\Omega)}\|e(u)\|_{L^2(\mathcal{S}_{\varepsilon,r})}.
\end{aligned}$$

The above estimates, those of Lemma A.2 and the fact that $r\le\varepsilon$ end the proof of Lemma 5.4.

C Unfolding method results

Proof of Lemma 6.3: Since Φ_ε belongs to $W^{1,\infty}(\Omega_\varepsilon)$ then $\Phi_{\varepsilon|_{\mathcal{S}_\varepsilon}}$ is in $W^{1,\infty}(\mathcal{S}_\varepsilon)$. Taking into account that $x=\varepsilon\xi+\varepsilon\mathbf{A}^\ell+s\mathbf{t}_1$ in $\mathcal{S}_\varepsilon$, we have equality $(6.10)_2$.

Since $Q_1(Y)$ is a finite dimensional vector space, there exist two strictly positive constants c and C such that for every $\Psi\in Q_1(Y)$

$$c\big\|\Psi_{|\mathcal{S}}\big\|_{L^2(\mathcal{S})}\le\|\Psi\|_{L^2(Y)}\le C\big\|\Psi_{|\mathcal{S}}\big\|_{L^2(\mathcal{S})},\qquad c\Big\|\frac{d\Psi_{|\mathcal{S}}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})}\le\|\nabla\Psi\|_{L^2(Y)}\le C\Big\|\frac{d\Psi_{|\mathcal{S}}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})}.$$

Now, for every $\Phi\in W^{1,\infty}(\Omega_\varepsilon)$ satisfying (6.9), after ε-scaling, we obtain

$$\begin{aligned}
c\varepsilon^2\big\|\Phi_{|\mathcal{S}_\varepsilon}\big\|^2_{L^2((\varepsilon\xi+\varepsilon\overline{Y})\cap\mathcal{S}_\varepsilon)}&\le\|\Phi\|^2_{L^2(\varepsilon\xi+\varepsilon Y)}\le C\varepsilon^2\big\|\Phi_{|\mathcal{S}_\varepsilon}\big\|^2_{L^2((\varepsilon\xi+\varepsilon\overline{Y})\cap\mathcal{S}_\varepsilon)},\\
c\varepsilon^2\Big\|\frac{d\Phi_{|\mathcal{S}_\varepsilon}}{d\mathbf{s}}\Big\|^2_{L^2((\varepsilon\xi+\varepsilon\overline{Y})\cap\mathcal{S}_\varepsilon)}&\le\|\nabla\Phi\|^2_{L^2(\varepsilon\xi+\varepsilon Y)}\le C\varepsilon^2\Big\|\frac{d\Phi_{|\mathcal{S}_\varepsilon}}{d\mathbf{s}}\Big\|^2_{L^2((\varepsilon\xi+\varepsilon\overline{Y})\cap\mathcal{S}_\varepsilon)},
\end{aligned}\qquad\forall\xi\in\Xi_\varepsilon.$$

Summing up all these inequalities for all $\xi\in\Xi_\varepsilon$ yields $(6.10)_{1,3}$.

Now, suppose that the sequence $\{\Phi_\varepsilon\}_\varepsilon$ of functions belonging to $W^{1,\infty}(\Omega_\varepsilon)$ satisfies (6.11). Then, up to a sequence of ε, there exists $\Phi\in L^2(\Omega)$ such that $(6.12)_1$ holds and furthermore due to (6.3) (see also [9, Theorem 3.6]), one has

$$\mathcal{T}_\varepsilon^{ext}(\Phi_\varepsilon)\rightharpoonup\Phi\quad\text{weakly in } L^2(\Omega\times Y).$$

But, taking into account (6.9), we have the convergence $(6.12)_2$ which implies $(6.12)_3$, since the embedding $Q_1(Y)\subset H^1(\mathcal{S})$ is continuous.

Moreover, if $\|\Phi_\varepsilon\|_{H^1(\Omega_\varepsilon)}\le C$ then, up to a sequence of ε, there exists $\Phi\in H^1(\Omega)$ such that $(6.13)_1$ holds. In the same way as [9, Theorem 3.6], we obtain convergence $(6.13)_2$, from which, taking into account $(6.10)_1$, we have convergence $(6.13)_3$.

Proof of Lemma 6.4: Using the properties of the unfolding operator $\mathcal{T}_\varepsilon^{\mathcal{S}}$ (6.3)-(6.4) and the estimates for ϕ_ε, we obtain

$$\|\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\|_{L^2(\Omega_\varepsilon\times\mathcal{S})}=\varepsilon\|\phi_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon)}\le C,$$

and

$$\Big\|\frac{\partial\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)}{\partial\mathbf{S}}\Big\|_{L^2(\Omega_\varepsilon\times\mathcal{S})}=\Big\|\varepsilon\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big)\Big\|_{L^2(\Omega_\varepsilon\times\mathcal{S})}=\varepsilon^2\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)}\le C.$$

Thus,

$$\|\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\|_{L^2(\Omega;H^1(\mathcal{S}))}\le\|\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\|_{L^2(\Omega_\varepsilon;H^1(\mathcal{S}))}\le C.$$

Hence, up to a subsequence ε, there exists $\widehat{\phi}\in L^2(\Omega;H^1(\mathcal{S}))$ such that (6.14) holds.

In order to prove of (6.15), first observe that $\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}})$ belongs to $L^2(\Omega;H^1(\mathcal{S}))$ and

$$\|\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}}\|_{L^2(\Omega;H^1(\mathcal{S}))}\le\varepsilon\Big(\|\phi_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon\cap\Omega_\varepsilon^{int})}+\varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon\cap\Omega_\varepsilon^{int})}\Big)\le C.$$

And, up to a subsequence of $\{\varepsilon\}$, there exists $\widehat{\phi} \in L^2(\Omega; H^1(\mathcal{S})$ such that (6.15) holds.

In both cases, the periodicity of $\widehat{\phi}$ is obtained proceeding in the same way as to prove [10, Theorem 4.28].

Proof of Lemma 6.5: The Poincaré-Wirtinger inequality gives a constant such that

$$\forall\, \psi \in H^1(\mathcal{S}), \qquad \Big\|\psi - \frac{1}{|\mathcal{S}|}\int_{\mathcal{S}} \psi\, d\mathbf{S}\Big\|_{L^2(\mathcal{S})} \leq C\Big\|\frac{d\psi}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})}.$$

We apply the above inequality with the function $\psi(S) = \phi_\varepsilon(\varepsilon\xi + \varepsilon S)$, $\xi \in \Xi_\varepsilon$. Then, after summation over $\xi \in \Xi_\varepsilon$, that yields

$$\|\phi_\varepsilon - \mathcal{M}^*_\varepsilon(\phi_\varepsilon)\|_{L^2(\mathcal{S}_\varepsilon)} \leq C\varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)}. \tag{C.1}$$

Now, consider the function Φ_ε defined in the cell $\varepsilon(\xi + \overline{Y})$, $\xi \in \widehat{\Xi}_\varepsilon$, as the Q_1 interpolation of $\mathcal{M}^*_\varepsilon(\phi)(\varepsilon\xi)$ on the vertices of this parallelotope. One has

$$\Phi_\varepsilon \in W^{1,\infty}(\widehat{\Omega}_\varepsilon).$$

Observe that Φ_ε also belongs to $W^{1,\infty}(\Omega^{int}_\varepsilon)$. Proceeding as in [10, Chapter 4] we obtain the following estimates:

$$\begin{aligned}
&\|\Phi_\varepsilon\|_{L^2(\widehat{\Omega}_\varepsilon)} \leq C\varepsilon\|\phi_\varepsilon\|_{L^2(\mathcal{S}_\varepsilon)}, \quad \|\nabla\Phi_\varepsilon\|_{L^2(\widehat{\Omega}_\varepsilon)} \leq C\varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)},\\
&\|\Phi_\varepsilon - \mathcal{M}^*_\varepsilon(\phi_\varepsilon)\|_{L^2(\widehat{\Omega}_\varepsilon)} \leq C\varepsilon^2\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)}.
\end{aligned} \tag{C.2}$$

Therefore,

$$\|\Phi_\varepsilon\|_{H^1(\Omega^{int}_\varepsilon)} \leq \|\Phi_\varepsilon\|_{H^1(\widehat{\Omega}_\varepsilon)} \leq C.$$

Lemma 5.1 in [16] gives $\Phi \in H^1(\Omega)$ such that (up to a subsequence)

$$\begin{aligned}
&\Phi_\varepsilon \mathbf{1}_{\Omega^{int}_\varepsilon} \longrightarrow \Phi \quad \text{strongly in} \quad L^2(\Omega),\\
&\nabla\Phi_\varepsilon \mathbf{1}_{\Omega^{int}_\varepsilon} \rightharpoonup \nabla\Phi \quad \text{weakly in} \quad L^2(\Omega)^3.
\end{aligned} \tag{C.3}$$

Besides, by definition of Φ_ε, $\Phi_{\varepsilon|\widehat{\Omega}_\varepsilon\cap\mathcal{S}_\varepsilon}$ belongs to $W^{1,\infty}(\widehat{\Omega}_\varepsilon \cap \mathcal{S}_\varepsilon)$ and Lemma 6.3 gives

$$\frac{d\Phi_{\varepsilon|\mathcal{S}_\varepsilon}}{d\mathbf{s}} = \nabla\Phi_\varepsilon \cdot \mathbf{t}_1 \quad \text{a.e. in } \widehat{\Omega}_\varepsilon \cap \mathcal{S}_\varepsilon, \quad \Big\|\frac{d\Phi_{\varepsilon|\mathcal{S}_\varepsilon}}{d\mathbf{s}}\Big\|_{L^2(\widehat{\Omega}_\varepsilon\cap\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}\|\nabla\Phi_\varepsilon\|_{L^2(\widehat{\Omega}_\varepsilon)}. \tag{C.4}$$

Hence,

$$\Big\|\frac{d\Phi_\varepsilon}{d\mathbf{s}} - \frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\widehat{\Omega}_\varepsilon\cap\mathcal{S}_\varepsilon)} \leq \Big\|\frac{d\Phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\widehat{\Omega}_\varepsilon\cap\mathcal{S}_\varepsilon)} + \Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}\|\nabla\Phi_\varepsilon\|_{L^2(\widehat{\Omega}_\varepsilon)} + \Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq C\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq \frac{C}{\varepsilon}.$$

By (6.16), (C.1)-(C.2), we obtain

$$\|\Phi_\varepsilon - \phi_\varepsilon\|_{L^2(\widehat{\Omega}_\varepsilon\cap\mathcal{S}_\varepsilon)} \leq C\varepsilon\Big\|\frac{d\phi_\varepsilon}{d\mathbf{s}}\Big\|_{L^2(\mathcal{S}_\varepsilon)} \leq C.$$

Therefore, Lemma 6.4 gives a function $\widehat{\phi} \in L^2(\Omega; H^1_{per}(\mathcal{S}))$ such that (up to a subsequence)

$$\frac{1}{\varepsilon}\mathcal{T}^{\mathcal{S}}_\varepsilon(\phi_\varepsilon - \Phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}^{int}_\varepsilon\times\mathcal{S}} \rightharpoonup \widehat{\phi} \quad \text{weakly in } \ L^2(\Omega; H^1(\mathcal{S})). \tag{C.5}$$

Due to estimate (C.4), there exist a subsequence of $\{\varepsilon\}$ and $F \in L^2(\mathcal{S})$ such that

$$\mathcal{T}^{\mathcal{S}}_\varepsilon\Big(\frac{d\Phi_\varepsilon}{d\mathbf{s}}\Big)\mathbf{1}_{\widehat{\Omega}_\varepsilon\times\mathcal{S}} \rightharpoonup F \quad \text{weakly in } L^2(\Omega\times\mathcal{S}).$$

Let $\mathcal{O}$ be an open set strictly included in Ω. If ε is small enough, one has

$$\mathcal{O} \subset \widehat{\Omega}^{int}_\varepsilon \subset \widehat{\Omega}_\varepsilon.$$

Applying Lemma 6.3 in the context of the open set $\mathcal{O}$ leads to (up to a subsequence)

$$\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\Phi_\varepsilon}{d\mathbf{s}}\Big)\mathbf{1}_{\mathcal{O}\times\mathcal{S}} = \mathcal{T}_\varepsilon^{ext}(\nabla\Phi_\varepsilon\cdot\mathbf{t}_1)\mathbf{1}_{\mathcal{O}\times\mathcal{S}} \rightharpoonup \nabla\Phi\cdot\mathbf{t}_1 \quad \text{weakly in } L^2(\mathcal{O}\times\mathcal{S}).$$

Hence,

$$F = \nabla\Phi\cdot\mathbf{t}_1 \quad \text{a.e. in } \mathcal{O}\times\mathcal{S}.$$

As a consequence $F = \nabla\Phi\cdot\mathbf{t}_1$ a.e. in $\Omega\times\mathcal{S}$ and (6.17) are proved. Now, from (C.3) and (C.5) we obtain

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\Phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \longrightarrow \Phi \quad \text{strongly in} \quad L^2(\Omega;H^1(\mathcal{S})),$$
$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon-\Phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \longrightarrow 0 \quad \text{strongly in} \quad L^2(\Omega;H^1(\mathcal{S})).$$

Hence,

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_\varepsilon)\mathbf{1}_{\widehat{\Omega}_\varepsilon^{int}\times\mathcal{S}} \longrightarrow \Phi \quad \text{strongly in} \quad L^2(\Omega;H^1(\mathcal{S})).$$

D Proof of Lemma 7.1

Step 1. We show that the semi-norm $\|\cdot\|_{\mathcal{S}}$ is a norm in $\boldsymbol{\mathcal{D}}_{Ex}$.

Indeed, if $\Big\|\dfrac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})} = 0$ then $\overline{\mathcal{A}}$ is a rigid displacement (remind that $\mathcal{S}$ is a stable structure). The periodicity of $\overline{\mathcal{A}}$ implies that $\overline{\mathcal{A}}$ is a constant field. Since the mean value of $\overline{\mathcal{A}}$ is equal to zero then $\overline{\mathcal{A}} = 0$. Hence, the semi-norm $\|\cdot\|_{\mathcal{S}}$ is a norm in $\boldsymbol{\mathcal{D}}_{Ex}$.

Step 2. We show that the norm $\|\cdot\|_{\mathcal{S}}$ is equivalent to the norm $\|\cdot\|_{H^1(\mathcal{S})}$.

First, we have

$$\forall\,\overline{\mathcal{A}}\in\boldsymbol{\mathcal{D}}_{Ex}, \qquad \|\overline{\mathcal{A}}\|_{\mathcal{S}} = \Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})} \le \Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})} \le \|\overline{\mathcal{A}}\|_{H^1(\mathcal{S})}.$$

The map

$$\overline{\mathcal{A}}\in\boldsymbol{\mathcal{D}}_{Ex} \longmapsto \overline{\mathcal{A}}_{\rm Aff}\in H^1_{per}(\mathcal{S})^3\cap\mathbf{U}_{\mathcal{S}},$$

where $\overline{\mathcal{A}}_{\rm Aff}$ is defined by

$$\overline{\mathcal{A}}_{\rm Aff}(A) = \overline{\mathcal{A}}(A) \qquad \forall A\in\mathcal{K}.$$

Lemma 2.1 claims that there exists a rigid displacement $\mathbf{r}$ such that

$$\|\overline{\mathcal{A}}_{\rm Aff}-\mathbf{r}\|_{H^1(\mathcal{S})} \le C\|\overline{\mathcal{A}}_{\rm Aff}\|_{\mathcal{S}} = C\Big\|\frac{d\overline{\mathcal{A}}_{\rm Aff}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}.$$

Since $\mathcal{S}$ is a 3-periodic structure and $\overline{\mathcal{A}}_{\rm Aff}$ is a periodic function, we can choose $\mathbf{r}$ constant. Hence,

$$\Big\|\frac{d\overline{\mathcal{A}}_{\rm Aff}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})} \le C\Big\|\frac{d\overline{\mathcal{A}}_{\rm Aff}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}. \tag{D.1}$$

The function $\overline{\mathcal{A}}-\overline{\mathcal{A}}_{\rm Aff}$ vanishes on all the nodes. Therefore by the definitions of the functions $\overline{\mathcal{A}}$ and $\overline{\mathcal{A}}_{\rm Aff}$ we obtain

$$\big(\overline{\mathcal{A}}-\overline{\mathcal{A}}_{\rm Aff}\big)\wedge\mathbf{t}_1 \text{ is an affine function on all the segments } \gamma_\ell,\ \ell\in\{1,\dots,m\}\ .$$

Hence,

$$\Big\|\frac{d\big(\overline{\mathcal{A}}-\overline{\mathcal{A}}_{\rm Aff}\big)}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|^2_{L^2(\mathcal{S})} + \Big\|\frac{d\overline{\mathcal{A}}_{\rm Aff}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|^2_{L^2(\mathcal{S})} = \Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|^2_{L^2(\mathcal{S})} \tag{D.2}$$

and, therefore,

$$\Big\|\frac{d\big(\overline{\mathcal{A}}-\overline{\mathcal{A}}_{\rm Aff}\big)}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})} = \Big\|\frac{d\big(\overline{\mathcal{A}}-\overline{\mathcal{A}}_{\rm Aff}\big)}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})} \le \Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}. \tag{D.3}$$

As a consequence of (D.1)-(D.3), one obtains

$$\forall\overline{\mathcal{A}}\in\boldsymbol{\mathcal{D}}_{Ex}, \qquad \Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\Big\|_{L^2(\mathcal{S})} \le C\Big\|\frac{d\overline{\mathcal{A}}}{d\mathbf{S}}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}.$$

Remind that since $\overline{\mathcal{A}}$ belongs to $H^1_{per,0}(\mathcal{S})^3$, the Poincaré-Wirtinger inequality gives

$$\forall \overline{\mathcal{A}} \in \boldsymbol{\mathcal{D}}_{Ex}, \qquad \|\overline{\mathcal{A}}\|_{L^2(\mathcal{S})} \leq C \Big\| \frac{d\overline{\mathcal{A}}}{d\mathbf{S}} \Big\|_{L^2(\mathcal{S})} \leq C \Big\| \frac{d\overline{\mathcal{A}}}{d\mathbf{S}} \cdot \mathbf{t}_1 \Big\|_{L^2(\mathcal{S})}.$$

Thus,

$$\forall \overline{\mathcal{A}} \in \boldsymbol{\mathcal{D}}_{Ex}, \qquad \|\overline{\mathcal{A}}\|_{H^1(\mathcal{S})} \leq C \Big\| \frac{d\overline{\mathcal{A}}}{d\mathbf{S}} \cdot \mathbf{t}_1 \Big\|_{L^2(\mathcal{S})} = C \|\overline{\mathcal{A}}\|_{\mathcal{S}}.$$

Both norms are equivalent.

Step 3. We show that the semi-norm $\|(\cdot,\cdot)\|_{\boldsymbol{\mathcal{D}}_{In}}$ is a norm in $\boldsymbol{\mathcal{D}}_{In}$.

Indeed, if $\Big\| \dfrac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \Big\|_{L^2(\mathcal{S})} = 0$, then $\widehat{\mathcal{B}}$ is a constant field. Remind that $\widehat{\mathcal{A}}$ vanishes on all the nodes, therefore one has $\widehat{\mathcal{B}} \wedge \mathbf{t}_1 = 0$ in $\mathcal{S}$. Since every node is a common extremity of at least two segments with non-collinear direction, then $\widehat{\mathcal{B}}$ vanishes on every node and thus $\widehat{\mathcal{B}} = 0$ in $\mathcal{S}$. Hence, we have $\widehat{\mathcal{A}} = 0$ on $\mathcal{S}$ and the semi-norm $\|(\cdot,\cdot)\|_{\boldsymbol{\mathcal{D}}_{In}}$ is a norm in $\boldsymbol{\mathcal{D}}_{In}$.

Step 4. We show that the norm $\|(\cdot,\cdot)\|_{\boldsymbol{\mathcal{D}}_{In}}$ is equivalent to the norm $\|(\cdot,\cdot)\|_{H^1(\mathcal{S})\times H^1(\mathcal{S})}$.

First, we have

$$\forall (\widehat{\mathcal{A}},\widehat{\mathcal{B}}) \in \boldsymbol{\mathcal{D}}_{In}, \qquad \|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{\boldsymbol{\mathcal{D}}_{In}} \doteq \Big\| \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \Big\|_{L^2(\mathcal{S})} \leq \|\widehat{\mathcal{A}}\|_{H^1(\mathcal{S})} + \|\widehat{\mathcal{B}}\|_{H^1(\mathcal{S})} = \|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{H^1(\mathcal{S})\times H^1(\mathcal{S})}.$$

We prove by contradiction that there exists a constant C strictly positive such that

$$\forall (\widehat{\mathcal{A}},\widehat{\mathcal{B}}) \in \boldsymbol{\mathcal{D}}_{In}, \qquad \|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{H^1(\mathcal{S})\times H^1(\mathcal{S})} \leq C \|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{\boldsymbol{\mathcal{D}}_{In}}.$$

Suppose that such constant does not exist, then for every $n \geq 1$, there exists $(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n) \in \boldsymbol{\mathcal{D}}_{In}$ such that

$$\|(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n)\|_{H^1(\mathcal{S})\times H^1(\mathcal{S})} = 1 \qquad \text{and} \qquad \|(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n)\|_{\boldsymbol{\mathcal{D}}_{In}} \leq \frac{1}{n}.$$

Thus, there exists a subsequence, still denoted n, such that

$$(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n) \rightharpoonup (\widehat{\mathcal{A}},\widehat{\mathcal{B}}) \ \text{ weakly in } \ \big(H^1_{per}(\mathcal{S})^3 \times H^1_{per}(\mathcal{S})^3\big) \cap \boldsymbol{\mathcal{D}}_{In}.$$

Then, one has

$$\|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{\boldsymbol{\mathcal{D}}_{In}} = \Big\| \frac{d\widehat{\mathcal{B}}}{d\mathbf{S}} \Big\|_{L^2(\mathcal{S})} \leq \liminf_{n\to+\infty} \Big\| \frac{d\widehat{\mathcal{B}}_n}{d\mathbf{S}} \Big\|_{L^2(\mathcal{S})} = \liminf_{n\to+\infty} \|(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n)\|_{\boldsymbol{\mathcal{D}}_{In}} = 0.$$

Hence, $\|(\widehat{\mathcal{A}},\widehat{\mathcal{B}})\|_{\boldsymbol{\mathcal{D}}_{In}} = 0$ which implies $(\widehat{\mathcal{A}},\widehat{\mathcal{B}}) = (0,0)$. As a consequence of the above convergences, the Sobolev embedding and the definition of $\boldsymbol{\mathcal{D}}_{In}$ we obtain

$$\widehat{\mathcal{B}}_n \longrightarrow 0 \ \text{ strongly in } \ H^1_{per}(\mathcal{S})^3,$$
$$\text{and then} \quad \widehat{\mathcal{A}}_n \longrightarrow 0 \ \text{ strongly in } \ H^1_{per}(\mathcal{S})^3.$$

Finally $\|(\widehat{\mathcal{A}}_n,\widehat{\mathcal{B}}_n)\|_{H^1(\mathcal{S})\times H^1(\mathcal{S})} \to 0$ which gives us a contradiction.

E Density results

Let $\mathfrak{r}$ and a be two constants such that $0 < 4\mathfrak{r} < a$.

Lemma E.1. *For every ϕ in $H^1(0,a)$, we define $\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi) \in H^1(0,a)$ by*

$$\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)(t) = \begin{cases} \phi(0) & \text{if } t \in [0,\mathfrak{r}], \\ \big(\phi(2\mathfrak{r}) - \phi(0)\big)\dfrac{t-\mathfrak{r}}{\mathfrak{r}} + \phi(0) + A(t-\mathfrak{r})(t-2\mathfrak{r}) & \text{if } t \in [\mathfrak{r},2\mathfrak{r}], \\ \phi(t) & \text{if } t \in [2\mathfrak{r},a-2\mathfrak{r}], \\ \big(\phi(a) - \phi(a-2\mathfrak{r})\big)\dfrac{t-a+2\mathfrak{r}}{\mathfrak{r}} + \phi(a) + B(t-a+\mathfrak{r})(t-a+2\mathfrak{r}) & \text{if } t \in [a-2\mathfrak{r},a-\mathfrak{r}], \\ \phi(a) & \text{if } t \in [a-\mathfrak{r},a], \end{cases}$$

where A and B are determined by the equalities

$$\int_0^{2\mathfrak{r}} \boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)(t)\,dt = \int_0^{2\mathfrak{r}} \phi(t)\,dt, \qquad \int_{a-2\mathfrak{r}}^{a} \boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)(t)\,dt = \int_{a-2\mathfrak{r}}^{a} \phi(t)\,dt. \tag{E.1}$$

Then one has

$$\|\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)'\|_{L^2(0,2\mathfrak{r})} + \|\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)'\|_{L^2(a-2\mathfrak{r},a)} \leq C\big(\|\phi'\|_{L^2(0,2\mathfrak{r})} + \|\phi'\|_{L^2(a-2\mathfrak{r},a)}\big). \tag{E.2}$$

The constant does not depend on a and $\mathfrak{r}$.

Proof. From $(\text{E.1})_1$ we have

$$A = \frac{6}{\mathfrak{r}^3}\Big(\frac{\mathfrak{r}}{2}(\phi(2\mathfrak{r})-\phi(0)) + \int_0^{2\mathfrak{r}} (\phi(0)-\phi(t))\,dt\Big) = \frac{6}{\mathfrak{r}^3}\Big(\frac{\mathfrak{r}}{2}\int_0^{2\mathfrak{r}} \phi'(\tilde{t})\,d\tilde{t} - \int_0^{2\mathfrak{r}}\int_0^{t} \phi'(\tilde{t})\,d\tilde{t}\,dt\Big).$$

Then, using the Cauchy-Schwartz inequality we obtain

$$|A|^2 \leq \frac{36}{\mathfrak{r}^4}\Big(\mathfrak{r}\|\phi'\|^2_{L^2(0,2r)} + \int_0^{2\mathfrak{r}}\int_0^{t} |\phi'(\tilde{t})|^2\,d\tilde{t}\,dt\Big) \leq \frac{36}{\mathfrak{r}^4}\Big(\mathfrak{r}\|\phi'\|^2_{L^2(0,2r)} + 2\mathfrak{r}\|\phi'\|^2_{L^2(0,2r)}\Big) = \frac{108}{\mathfrak{r}^3}\|\phi'\|^2_{L^2(0,2r)}.$$

By definition of $\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}$ and again using the Cauchy-Schwartz inequality we have

$$\begin{aligned}
\|\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)'\|^2_{L^2(0,2\mathfrak{r})} &= \int_{\mathfrak{r}}^{2\mathfrak{r}} \Big(\frac{\phi(2\mathfrak{r})-\phi(0)}{\mathfrak{r}} + A(t-\mathfrak{r}) + A(t-2\mathfrak{r})\Big)^2\,dt \\
&\leq \frac{3}{\mathfrak{r}^2}\int_{\mathfrak{r}}^{2\mathfrak{r}} (\phi(2\mathfrak{r})-\phi(0))^2\,dt + 3A^2\int_{\mathfrak{r}}^{2\mathfrak{r}} (t-\mathfrak{r})^2\,dt + 3A^2\int_{\mathfrak{r}}^{2\mathfrak{r}} (t-2\mathfrak{r})^2\,dt \\
&= \frac{3}{\mathfrak{r}}\Big(\int_0^{2\mathfrak{r}} \phi'(\tilde{t})\,d\tilde{t}\Big)^2\,dt + 2A^2\mathfrak{r}^3 \leq \frac{3}{\mathfrak{r}}\Big(\int_0^{2\mathfrak{r}} 1^2\,dt\Big)\Big(\int_0^{2\mathfrak{r}} |\phi'(\tilde{t})|^2\,d\tilde{t}\Big) + 204\|\phi'\|^2_{L^2(0,2r)} \\
&\leq 6\|\phi'\|^2_{L^2(0,2\mathfrak{r})} + 204\|\phi'\|^2_{L^2(0,2r)} = 210\|\phi'\|^2_{L^2(0,2\mathfrak{r})}.
\end{aligned}$$

In the same way we obtain

$$\|\boldsymbol{\mathcal{F}}_{\mathfrak{r},a}(\phi)'\|^2_{L^2(a-2\mathfrak{r},a)} \leq 210\|\phi'\|^2_{L^2(a-2\mathfrak{r},a)}$$

and (E.2) holds. □

Let $\mathcal{S}$ be a $3D$-periodic structure. For every $\mathfrak{r}$ satisfying (remind that l_ℓ is the length of the segment $\gamma_\ell \subset \mathcal{S}$)

$$0 < 4\mathfrak{r} < \min_{\ell\in\{1,\dots,m\}} l_\ell$$

we define the map $\boldsymbol{\mathcal{F}}_{\mathfrak{r}}$ from $H^1(\mathcal{S})$ into $H^1(\mathcal{S})$ by

$$\forall \phi \in H^1(\mathcal{S}), \quad \forall \gamma_\ell \subset \mathcal{S}, \qquad \boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\phi)_{|\gamma_\ell} = \boldsymbol{\mathcal{F}}_{\mathfrak{r},l_\ell}(\phi_{|\gamma_\ell}).$$

Lemma E.2. *$\boldsymbol{\mathcal{F}}_{\mathfrak{r}}$ is a linear and continuous map from $H^1_{per}(\mathcal{S})$ into $H^1_{per}(\mathcal{S})$. We have*

$$\forall \phi \in H^1_{per}(\mathcal{S}), \qquad \boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\phi) \longrightarrow \phi \qquad \text{strongly in } \ H^1_{per}(\mathcal{S}). \tag{E.3}$$

For every $(\widehat{\mathcal{A}},\widehat{\mathcal{B}}) \in \boldsymbol{\mathcal{D}}_{In}$, we define $\widehat{\mathcal{A}}_{\mathfrak{r}} \in H^1_{per}(\mathcal{S})^3$ by

$$\frac{d\widehat{\mathcal{A}}_{\mathfrak{r}}}{d\mathbf{s}} = \boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\widehat{\mathcal{B}}) \wedge \mathbf{t}_1, \qquad \widehat{\mathcal{A}}_{\mathfrak{r}} = 0 \text{ on all the nodes of } \mathcal{S}. \tag{E.4}$$

The couple $\big(\widehat{\mathcal{A}}_{\mathfrak{r}}, \boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\widehat{\mathcal{B}})\big)$ belongs to $\boldsymbol{\mathcal{D}}_{In}$ and we have

$$\begin{aligned}
\boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\widehat{\mathcal{B}}) &\longrightarrow \widehat{\mathcal{B}} \qquad \text{strongly in } \ H^1_{per}(\mathcal{S})^3, \\
\widehat{\mathcal{A}}_{\mathfrak{r}} &\longrightarrow \widehat{\mathcal{A}} \qquad \text{strongly in } \ H^1_{per}(\mathcal{S})^3.
\end{aligned} \tag{E.5}$$

Proof. First observe that for every $\phi \in H^1(\mathcal{S})$, the function $\boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\phi) - \phi$ vanishes on every nodes. As a consequence, $\boldsymbol{\mathcal{F}}_{\mathfrak{r}}$ maps $H^1_{per}(\mathcal{S})$ into $H^1_{per}(\mathcal{S})$. Then (E.3) follows from Lemma E.1.

Due to properties (E.1) of $\boldsymbol{\mathcal{F}}_{\mathfrak{r}}(\widehat{\mathcal{B}})$, the function $\widehat{\mathcal{A}}_{\mathfrak{r}}$ is well define by (E.4). Then, the convergences (E.5) are the immediate consequences of Lemma E.1. □

F The test functions $\phi_{\varepsilon,r}$ and $\mathcal{V}_{\varepsilon,r}$

Let ϕ be in $\mathcal{D}(\Omega)$. We define the field $\phi_{\varepsilon,r}$ belonging to $W^{1,\infty}(\mathcal{S}_\varepsilon)$ as follows:
for every $\xi \in \Xi_\varepsilon$ and every $\gamma_\ell = [\mathbf{A}^\ell, \mathbf{A}^\ell + l_\ell \mathbf{t}_1^\ell] \in \mathcal{S}$ we set

$$\phi_{\varepsilon,r}(\mathbf{s}) = \begin{cases} \phi(\varepsilon\xi + \varepsilon\mathbf{A}^\ell) & \text{in } [-c_0 r, c_0 r], \\ \mathbf{\Phi}_{\varepsilon,r}(s_1) & \text{if } s_1 \in [c_0 r, \varepsilon l_\ell - c_0 r], \\ \phi(\varepsilon\xi + \varepsilon\mathbf{B}^\ell) & \text{in } [\varepsilon l_\ell - c_0 r, \varepsilon l_\ell + c_0 r], \end{cases}$$

where $\mathbf{\Phi}_{\varepsilon,r}$ is a polynomial function of degree less than 3 with respect to s_1 such that

$$\mathbf{\Phi}_{\varepsilon,r}(c_0 r) = \phi(\varepsilon\xi + \varepsilon\mathbf{A}^\ell), \qquad \mathbf{\Phi}_{\varepsilon,r}(\varepsilon l_\ell - c_0 r) = \phi(\varepsilon\xi + \varepsilon\mathbf{B}^\ell), \qquad \frac{d\mathbf{\Phi}_{\varepsilon,r}}{ds_1}(c_0 r) = \frac{d\mathbf{\Phi}_{\varepsilon,r}}{ds_1}(\varepsilon l_\ell - c_0 r) = 0.$$

By construction $\phi_{\varepsilon,r|_{\varepsilon\xi+\varepsilon\gamma_\ell}}$ belongs to $W^{2,\infty}(\varepsilon\xi + \varepsilon\gamma_\ell)$. We easily check that

$$\mathcal{T}_\varepsilon^{\mathcal{S}}(\phi_{\varepsilon,r}) \longrightarrow \phi \quad \text{strongly in} \quad L^2(\Omega' \times \mathcal{S}'),$$
$$\varepsilon \mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\phi_{\varepsilon,r}}{d\mathbf{s}}\Big), \ \varepsilon^2 \mathcal{T}_\varepsilon^{b,\ell}\Big(\frac{d^2\phi_{\varepsilon,r}}{d\mathbf{s}^2}\Big) \longrightarrow 0 \quad \text{strongly in} \quad L^2(\Omega' \times \mathcal{S}').$$

Let $\mathcal{V}$ be in $\mathcal{D}(\mathbb{R}^3)$ such that $\mathcal{V} = 0$ in $\Omega' \setminus \overline{\Omega}$. We define the field $\phi_{\varepsilon,r}$ belonging to $W^{1,\infty}(\mathcal{S}'_\varepsilon)$ as follows:
for every $\xi \in \Xi_\varepsilon$ and every $\gamma_\ell = [\mathbf{A}^\ell, \mathbf{A}^\ell + l_\ell \mathbf{t}_1^\ell] \in \mathcal{S}$ we set

$$\mathcal{V}_{\varepsilon,r}(\mathbf{s}) = \begin{cases} \mathcal{V}(\varepsilon\xi + \varepsilon\mathbf{A}^\ell) & \text{in } [-c_0 r, c_0 r], \\ \mathcal{V}(\varepsilon\xi + \varepsilon\mathbf{A}^\ell) + \big(\mathcal{V}(\varepsilon\xi + \varepsilon\mathbf{A}^\ell + \varepsilon\mathbf{t}_1^\ell) - \mathcal{V}(\varepsilon\xi + \varepsilon\mathbf{A}^\ell)\big)\dfrac{s_1 - c_0 r}{\varepsilon l_\ell - 2c_0 r} & \text{if } s_1 \in [c_0 r, \varepsilon l_\ell - c_0 r], \\ \mathcal{V}(\varepsilon\xi + \varepsilon\mathbf{B}^\ell) & \text{in } [\varepsilon l_\ell - c_0 r, \varepsilon l_\ell + c_0 r], \end{cases}$$

We easily check that

$$\mathcal{T}_\varepsilon^{\mathcal{S}}\Big(\frac{d\mathcal{V}_{\varepsilon,r}}{d\mathbf{s}}\Big) \longrightarrow \nabla\mathcal{V} \cdot \mathbf{t}_1 \quad \text{strongly in} \quad L^2(\Omega' \times \mathcal{S}').$$

G Coercivity results.

Lemma G.1. *For every $\widetilde{v} = \big(\widetilde{v}_1, \widetilde{v}_2, \widetilde{v}_3\big) \in \boldsymbol{\mathcal{D}}_w$ and every $\zeta \in \mathbb{R}^4$, one has*

$$\int_D \big|\mathcal{E}_D(\widetilde{v}) + \mathbf{M}_\zeta\big|^2 dS_2 dS_3 = \pi\Big(|\zeta_1|^2 + \frac{1}{8}\big(|\zeta_3|^2 + |\zeta_4|^2\big) + \frac{1}{16}|\zeta_2|^2\Big) + \frac{1}{4}\|\nabla\widetilde{v}_1\|^2_{L^2(D)} + \sum_{j,k=2}^{3} \|e_{jk}(\widetilde{v})\|^2_{L^2(D)}, \quad \text{(G.1)}$$

where $\mathbf{M}_\zeta = \big(\zeta_1 + S_3\zeta_3 - S_2\zeta_4\big)\mathbf{M}^{11} - S_3\mathbf{M}^{12} + S_2\mathbf{M}^{13}$.
Moreover, there exists a strictly positive constant C such that

$$|\zeta|^2 + \|\widetilde{v}\|^2_{H^1(D)} \leq C \int_D \big|\mathcal{E}_D(\widetilde{v}) + \mathbf{M}_\zeta\big|^2 dS_2 dS_3. \quad \text{(G.2)}$$

Proof. A direct calculation gives

$$\int_D \big|\mathcal{E}_D(\widetilde{v}) + \mathbf{M}_\zeta\big|^2 dS_2 dS_3 = \int_D \big(\zeta_1 + S_3\zeta_3 - S_2\zeta_4\big)^2 dS_2 dS_3 + \frac{1}{4}\int_D \Big(\frac{\partial\widetilde{v}_1}{\partial S_2} - S_3\zeta_2\Big)^2 dS_2 dS_3$$
$$+ \frac{1}{4}\int_D \Big(\frac{\partial\widetilde{v}_1}{\partial S_3} + S_2\zeta_2\Big)^2 dS_2 dS_3 + \sum_{j,k=2}^{3} \|e_{jk}(\widetilde{v})\|^2_{L^2(\Omega)}.$$

Observe that

$$\int_D \Big(- S_3\frac{\partial\widetilde{v}_1}{\partial S_2} + S_2\frac{\partial\widetilde{v}_1}{\partial S_3}\Big) dS_2 dS_3 = 0.$$

Hence, we obtain (G.1). Then (G.2) follows from the definition of $\boldsymbol{\mathcal{D}}_w$, the Poincaré-Wirtinger and the Korn inequality. □

Lemma G.2. *There exists a strictly positive constant C such that*

$$\begin{aligned}&\forall\,\mathcal{V}\in H^1_\Gamma(\Omega)^3,\quad \forall(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\in L^2(\Omega;\boldsymbol{\mathcal{D}}_{Ex}\times\boldsymbol{\mathcal{D}}_{In}),\quad \forall\widetilde{v}\in L^2(\Omega\times\mathcal{S};\boldsymbol{\mathcal{D}}_w),\\ &\|\mathcal{V}\|^2_{H^1(\Omega)}+\|(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})\|^2_{L^2(\Omega;H^1(\mathcal{S}))}+\|\widetilde{v}\|^2_{L^2(\Omega;H^1(D))}\le C\int_{\Omega\times\mathcal{S}\times D}\left|\mathcal{E}(\mathcal{V})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{v})\right|^2 dx\,d\widehat{S}.\end{aligned}\tag{G.3}$$

Proof. Step 1. A preliminary result.
Let ζ be a 3×3 symmetric matrix. Consider the displacements $\overline{\mathcal{W}}\in\boldsymbol{\mathcal{D}}_{Ex}$ and $\mathbb{V}_\zeta(x)=\zeta\,x$, $x\in\mathbb{R}^3$. The restriction of $\mathbb{V}_\zeta$ to $\mathcal{S}$ belongs to $\mathbf{U}_{\mathcal{S}}$ and one has

$$\frac{d\mathbb{V}_\zeta}{dS_1}\cdot\mathbf{t}_1=(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1,\qquad(\zeta\,\mathbf{t}_1)\cdot\mathbf{t}_1+\frac{\partial\overline{\mathcal{W}}}{\partial S_1}\cdot\mathbf{t}_1=\frac{\partial(\mathbb{V}_\zeta+\overline{\mathcal{W}})}{\partial S_1}\cdot\mathbf{t}_1.$$

As in Step 2 of Lemma 7.1 and since the structure $\mathcal{S}$ is stable, we obtain a rigid displacement $\mathbf{r}$ such that

$$\|\mathbb{V}_\zeta+\overline{\mathcal{W}}-\mathbf{r}\|_{H^1(\mathcal{S})}\le C\Big\|\frac{\partial(\mathbb{V}_\zeta+\overline{\mathcal{W}})}{\partial S_1}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}.$$

Remind that $\mathcal{S}$ is also a 3-periodic structure. Therefore, comparing the values of $\mathbb{V}_\zeta+\overline{\mathcal{W}}-\mathbf{r}$ on the opposite faces of $Y\cap\mathcal{S}$ gives

$$|\zeta-\mathbf{b}|\le C\|\mathbb{V}_\zeta+\overline{\mathcal{W}}-\mathbf{r}\|_{H^1(\mathcal{S})},$$

where $\mathbf{b}=\nabla\mathbf{r}$ is a 3×3 antisymmetric matrix. Hence,

$$|\zeta|+|\mathbf{b}|\le C\Big\|\frac{\partial(\mathbb{V}_\zeta+\overline{\mathcal{W}})}{\partial S_1}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}.$$

Since $\overline{W}$ belongs to $H^1_{per,0}(\mathcal{S})^3$, we obtain

$$|\zeta|+\|\overline{\mathcal{W}}\|_{H^1(\mathcal{S})}\le C\Big\|\frac{\partial(\mathbb{V}_\zeta+\overline{\mathcal{W}})}{\partial S_1}\cdot\mathbf{t}_1\Big\|_{L^2(\mathcal{S})}.\tag{G.4}$$

Step 2. Inequality (G.2) leads to

$$\begin{aligned}\int_{\Omega\times\mathcal{S}}\Big|(e(\mathcal{V})\,\mathbf{t}_1)\mathbf{t}_1+\frac{\partial\overline{\mathcal{V}}}{\partial S_1}\cdot\mathbf{t}_1\Big|^2 d\mathbf{S}\,dx+\|(\widehat{\mathcal{V}},\widehat{\mathcal{B}})\|^2_{L^2(\Omega;\boldsymbol{\mathcal{D}}_{In})}+\|\widetilde{v}\|^2_{L^2(\Omega;H^1(D))}&\\ \le C\int_{\Omega\times\mathcal{S}\times D}\left|\mathcal{E}(\mathcal{V})+\mathcal{E}_{\mathcal{S}}(\overline{\mathcal{V}},\widehat{\mathcal{V}},\widehat{\mathcal{B}})+\mathcal{E}_D(\widetilde{v})\right|^2 d\widehat{S}\,dx.&\end{aligned}$$

Then, the estimate (G.4) and Lemma 7.1 give (G.3). □

Lemma G.3. *There exists $C_0'>0$ which does not depend on the variable $\mathbf{S}$, such that*

$$\forall\zeta\in\mathbb{R}^4,\qquad\mathfrak{A}\,\zeta\cdot\zeta\ge C_0'|\zeta|^2\qquad\text{a.e. in }\mathcal{S}.\tag{G.5}$$

Proof. Set $\widetilde{\chi}_\zeta=\sum_{q=1}^4\zeta_q\widetilde{\chi}_q$. By (5.4) one obtains

$$\mathfrak{A}\,\zeta\cdot\zeta\ge C_0\int_D\left|\mathcal{E}_D\big(\widetilde{\chi}_\zeta\big)+\mathbf{M}_\zeta\right|^2 dS_2dS_3,\qquad\text{a.e. in }\mathcal{S},$$

Lemma G.1 yields $\displaystyle\int_D\left|\mathcal{E}_D\big(\widetilde{\chi}_\zeta\big)+\mathbf{M}_\zeta\right|^2 dS_2dS_3\ge\pi\Big(|\zeta_1|^2+\frac{1}{8}\big(|\zeta_3|^2+|\zeta_4|^2\big)+\frac{1}{16}|\zeta_2|^2\Big)$. Thus, (G.5) is proved. □

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