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# The new nitrogen dioxide (NO<sub>2</sub>) linelist in the GEISA database and first identification of the $\nu_1+2\nu_3-\nu_3$ band of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>

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## ► To cite this version:

Agnès Perrin, L. Manceron, J.-M. Flaud, F. Kwabia-Tchana, Raymond Armante, et al.. The new nitrogen dioxide (NO<sub>2</sub>) linelist in the GEISA database and first identification of the  $\nu_1+2\nu_3-\nu_3$  band of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>. *Journal of Molecular Spectroscopy*, 2020, pp.111394. 10.1016/j.jms.2020.111394. hal-03086660

**HAL Id: hal-03086660**

**<https://hal.science/hal-03086660>**

Submitted on 22 Dec 2020

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## Journal Pre-proofs

The new nitrogen dioxide (NO<sub>2</sub>) linelist in the GEISA database and first identification of the  $\nu_1+2\nu_3-\nu_3$  band of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>

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PII: S0022-2852(20)30162-4  
DOI: <https://doi.org/10.1016/j.jms.2020.111394>  
Reference: YJMSP 111394

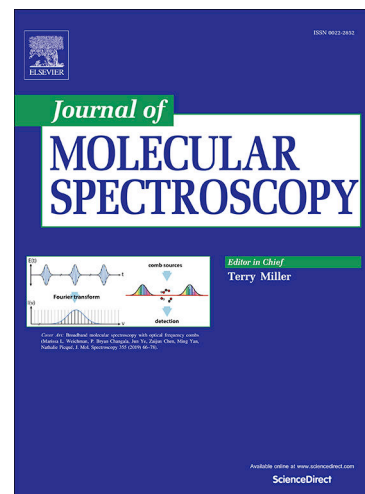
To appear in: *Journal of Molecular Spectroscopy*

Received Date: 21 September 2020  
Revised Date: 30 October 2020  
Accepted Date: 21 November 2020

Please cite this article as: A. Perrin, L. Manceron, J.-M. Flaud, F. Kwabia-Tchana, R. Armante, P. Roy, D. Doizi, The new nitrogen dioxide (NO<sub>2</sub>) linelist in the GEISA database and first identification of the  $\nu_1+2\nu_3-\nu_3$  band of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>, *Journal of Molecular Spectroscopy* (2020), doi: <https://doi.org/10.1016/j.jms.2020.111394>

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The new nitrogen dioxide (NO<sub>2</sub>) linelist in the GEISA database; first  
identification of the  $\nu_1+2\nu_3-\nu_3$  band of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>.

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Number of tables: 5

Number of figures: 11

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*Keywords:* Nitrogen dioxide; <sup>14</sup>N<sup>16</sup>O<sub>2</sub>; Fourier transform spectroscopy; Electron spin-rotation resonances; Vibration-rotation resonances; HITRAN; GEISA;

### Abstract

We have generated new lists of line position, line intensity and line shape parameters of nitrogen dioxide ( $^{14}\text{N}^{16}\text{O}_2$  and  $^{15}\text{N}^{16}\text{O}_2$ ), here labeled as “GEISA-19”, which have been included in the (Gestion et Etude des Informations Spectroscopiques Atmosphériques) GEISA database (<https://geisa.aeris-data.fr/>). Except for the far infrared and the 13.3  $\mu\text{m}$  regions, all spectral regions of the 1153- 4775  $\text{cm}^{-1}$  spectral domain are significantly modified by this major update of the GEISA linelist. For the 6.2  $\mu\text{m}$  and 3.4  $\mu\text{m}$  spectral regions, which correspond to the strongest absorption of  $\text{NO}_2$ , we proceed to a complete replacement of the lists for the first hot bands,  $\nu_2 + \nu_3 - \nu_2$  and  $\nu_1 + \nu_2 + \nu_3 - \nu_2$ , respectively, and to the inclusion, whenever possible, of higher order hot bands involving the (1,0,0), (0,2,0) and (0,0,1), (1,1,0), (2,0,0) or (0,0,2) states as lower states. Also, the  $\nu_1 + \nu_3$  linelist was improved for high rotational quantum numbers and the  $\nu_3$  and  $\nu_1 + \nu_3$  bands for  $^{15}\text{N}^{16}\text{O}_2$ , which is the second most abundant isotopologue of  $\text{NO}_2$ , are now included in the database. Finally several weak cold bands in the 2.2 to 4.9  $\mu\text{m}$  region were added for the first time to the GEISA linelist. These new vibration rotation transitions were generated using existing literature data or making use of experimental data extracted from high resolution Fourier transform spectra recorded at SOLEIL for the purpose of this study. One outcome of this study was the first identification of the  $\nu_1 + 2\nu_3 - \nu_3$  hot band, leading to the first determination of the (1,0,2) energy level parameters. Also, an improved set of parameters was derived for the (0,1,1) state. The validation of the GEISA-19 linelist was performed through a detailed comparison at 296K between computed and observed Fourier transform laboratory spectra. Also, the consistency, from one band to another, of the energy levels values was carefully checked. Finally inter-comparisons and verifications were performed using the recent versions of the HITRAN (<https://hitran.org/>) and HITEMP databases [R.J. Hargreaves, I. E. Gordon, L. S. Rothman, S. A. Tashkun, V. I. Perevalov, A. A. Lukashvskaya, S. N. Yurchenko, J. Tennyson, and H. S.P. Müller. *J. Quant. Spectrosc. Radiat. Transf.* **232** 35 (2019)]. Our conclusions are that, at 296K, GEISA-19 is of better quality than HITRAN2016-updated or HITEMP in the overall 1153- 4775  $\text{cm}^{-1}$  spectral region. As compared to its previous version, this new linelist will lead to an improved quality of the  $\text{NO}_2$  retrievals that will be performed for the future IASI-NG (Infrared Atmospheric Sounding Interferometer New Generation) satellite instrument (<https://iasi-ng.cnes.fr/fr>). However, contrary to HITEMP, GEISA-19 which does not include transitions involving high rotational quantum numbers or belonging to very high order hot bands cannot be used for hot temperature conditions.

## I. Introduction

Because of its important role in the photochemistry of the atmosphere, nitrogen dioxide has been the subject of numerous spectroscopic studies. These led to the generation of linelists which are now included in the GEISA [1], HITRAN [2], and HITEMP [3] databases.

Until very recently (2016), the linelists present in the 2016 version of HITRAN [2] and in the 2015 version of GEISA [1] for nitrogen dioxide did not differ significantly. This 2015 version of the NO<sub>2</sub> linelist, which concerns only the <sup>14</sup>N<sup>16</sup>O<sub>2</sub> isotopomer, will be noted as “HITRAN-GEISA-15” in the rest of the text. For the <sup>14</sup>N<sup>16</sup>O<sub>2</sub> species, four spectral regions are considered in HITRAN-GEISA-15 which correspond to the microwave to far infrared region (pure rotation within the (0,0,0) state), the 13.3 μm (ν<sub>2</sub> band), the 6.2 μm (the strong ν<sub>3</sub> band and the weaker ν<sub>1</sub> and 2ν<sub>2</sub> dark resonating bands) and 3.4 μm (the rather strong ν<sub>1</sub>+ν<sub>3</sub> band and the weaker ν<sub>1</sub>+2ν<sub>2</sub> dark resonating band) regions, respectively. Together with these cold bands, the HITRAN-GEISA-15 list includes in all four regions lines from the first hot bands which involve the (0,1,0) state as lower vibrational state. These linelists were generated at different time, and with different accuracy, and as it will be discussed in this paper, HITRAN-GEISA-15 presents clear deficiencies.

The most recent version of the HITEMP (“HIGH-TEMPerature molecular spectroscopic database”) linelist includes NO<sub>2</sub> for the first time [3]. The aim of HITEMP is to model gas phase spectra for high-temperature applications. For NO<sub>2</sub>, a composite HITEMP linelist was generated in the 0- 4775 cm<sup>-1</sup> by extending the current HITRAN linelist [2] (here labelled as HITRAN2016-updated<sup>1</sup>) using inputs from the recent “Nitrogen Dioxide Spectroscopic Databank” (NDS-1000) [4,5] line list. These additional inputs concern all spectral regions except the pure rotations bands (microwave to far infrared). For the vibrational transitions already considered in HITRAN, the NDS-1000 has provided extension of the current list up to the higher N (N≤100) and K<sub>a</sub> values. HITEMP also includes linelists for several cold and hot

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<sup>1</sup> In this paper, the reference date for the « HITRAN2016-updated” is September 2020.

bands<sup>2</sup> which, up to now, were missing in HITRAN-GEISA-15. Up to now, HITEMP does not include linelists for the  $^{15}\text{N}^{16}\text{O}_2$  (second) isotope species of nitrogen dioxide.

Subsequently, the “HITRAN2016-updated”, the 2020 version of HITRAN, was created [6]. The 1153-4775  $\text{cm}^{-1}$  spectral range now include lines from several weaker vibration-rotation bands present in HITEMP which were not considered previously in HITRAN-GEISA-15 [1,2]. More recently, the spectral range for  $\text{NO}_2$  has been extended to include transitions between 5720-8000  $\text{cm}^{-1}$ , with the addition of several combination and overtone bands [7-&15]. Finally, HITRAN2016-updated now includes line parameters for the  $\nu_3$  band of  $^{15}\text{N}^{16}\text{O}_2$  [16] which is the most abundant daughter isotopologue of  $\text{NO}_2$ .

Therefore, as compared to the previous version of HITRAN [2], both HITRAN2016-updated and HITEMP [3] databases were significantly extended. However, the existing deficiencies in HITRAN-GEISA-15 were not corrected.

The purpose of the present work was to generate a new version of GEISA linelist for  $\text{NO}_2$ , here identified as ‘GEISA-19’, which now includes  $\text{NO}_2$  lines parameters for the 0 – 4775.3  $\text{cm}^{-1}$  region.<sup>3</sup> As far as the line positions and intensities are concerned, this update does not concern the microwave to far infrared region, while all other spectral regions took benefit of new parameters generated during this work.

For the 6.2  $\mu\text{m}$  and 3.4  $\mu\text{m}$  regions which correspond to the strongest  $\text{NO}_2$  infrared absorption, we process to the replacement of the cold and of «first» hot bands, and GEISA-19 now includes, when possible, line parameters from higher order hot bands. We also implement in GEISA-19 linelists for several weak cold bands that absorb in the 2000- 4775  $\text{cm}^{-1}$  spectral range. In addition, GEISA-19 includes linelists for the  $\nu_3$  and  $\nu_1+\nu_3$  bands for the  $^{15}\text{N}^{16}\text{O}_2$  isotopic species.

Therefore, it is clear that the GEISA-19 linelist differs from HITRAN-GEISA-15, HITRAN2016-updated [6], and HITEMP [3].

The present work was performed using, as input to our calculations, not only the spectroscopic data available in the literature but also additional data issued from two recent spectroscopic

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<sup>2</sup> In the rest of the paper, we will designate as “first”, “second”, “third” and “fourth” hot bands the vibrational bands involving, as lower vibrational state, the (0,1,0) state, one of the states belonging to the first triad  $\{(1,0,0),(0,2,0),(0,0,1)\}$ , second triad  $\{(1,1,0),(0,3,0),(0,1,1)\}$ , or first hexade  $\{(2,0,0),(1,2,0),(1,0,1), (0,4,0), \{(0,2,1), (0,0,2)\}$  of interacting states of  $\text{NO}_2$ , respectively.

<sup>3</sup> This work does not concern 4800-8000  $\text{cm}^{-1}$  spectral range which is now considered in HITRAN16-updated.

studies (Refs. [17] and [18]), or obtained from the analysis of two new FTS spectra recorded during this work.

Before going into detail description of the work, we will define the vibrational and rotational quantum numbers and the energy levels used in this paper. Then we will describe the theoretical models used for the computations.

## **II General considerations**

### **II-A The NO<sub>2</sub> vibrational and rotational quantum numbers**

Nitrogen dioxide (<sup>14</sup>N<sup>16</sup>O<sub>2</sub>) is an asymmetric rotor, with three vibrational modes,  $\omega_1$ ,  $\omega_2$ , which are symmetric and  $\omega_3$  which is antisymmetric, and the vibrational modes are noted as ( $v_1, v_2, v_3$ ). As a matter of consequence, the vibrational states are grouped in polyads of interacting states. When necessary, the rovibrational interactions were accounted for in order to reproduce the measured line positions and intensities, and the present computations were performed polyad by polyad.

This molecule possesses an unpaired electron, and the electron spin-rotation interaction causes a doublet structure which is usually observable in the infrared. In addition, the hyperfine structure due to the I=1 nuclear spin is observable in the microwave, far infrared and 13.3  $\mu\text{m}$  regions [19, 20]. This work deals with the computation of line positions and intensities in the 1000- 4775.3  $\text{cm}^{-1}$  spectral range. Therefore, only four (N,  $K_a$ ,  $K_c$ , and J) rotational quantum numbers are considered for the upper and lower energy levels. The N,  $K_a$  and  $K_c$  are the main rotational quantum numbers. Because the electronic spin-rotation interaction has to be considered, the J quantum number (associated to  $\mathbf{J}=\mathbf{N}+\mathbf{S}$  (with  $S=\pm 1/2$ )) is indicated by a “+” and “-“ flag.

### **II B Definition of the upper $E_U$ and lower $E_L$ state energy levels.**

Each given vibration spin-rotation transition located at a  $\sigma$  position is then identified by its  $(v_1^U, v_2^U, v_3^U) [N^U, K_a^U, K_c^U, J^U]$  and  $(v_1^L, v_2^L, v_3^L) [N^L, K_a^L, K_c^L, J^L]$  upper and lower state vibration and rotational quantum numbers. For each considered transition, the HITRAN-GEISA-15, HITRAN2016-updated and HITEMP databases include the value of the lower state energy,  $E_L$ , while the upper level energy,  $E_U$ , is easily computable from the position,  $\sigma$  of the

considered transition ( $E_U = E_L + \sigma$ ). Within these databases one should expect that the individual values of  $E_L$  (resp. of  $E_U$ ) should be consistent from one given vibration – transition band to the other one, provided that they involve the same lower  $(v_1^L, v_2^L, v_3^L)$   $[N^L, K_a^L, K_c^L, J^L]$  (resp. same upper  $(v_1^U, v_2^U, v_3^U)$   $[N^U, K_a^U, K_c^U, J^U]$ ) vibrational spin-rotational level.

### **III Line position, line intensity and line shape parameters**

During this work the linelists for numerous vibration – rotation bands were generated in the 1153- 4775.4  $\text{cm}^{-1}$  spectral region, and the theoretical models used to compute the line positions and intensities account both for the spin- rotation resonances and for various vibration – rotational resonances.

#### **III-A Line positions**

For the computation of the energy levels, the Hamiltonian matrix involves in its  $v$ -diagonal blocks the sum of Watson’s type rotational [21] and of electronic spin rotational interaction operators [22]. These  $v$ - type operators are written for an A-type reduction and in an F representation. The  $(v_1, v_2, v_3) \Leftrightarrow (v_1', v_2', v_3')$   $v$ - off diagonal operators are Fermi-type or C-type Coriolis operators, for  $|\Delta v_3| = \text{even}$  and  $|\Delta v_3| = \text{odd}$ , respectively. Starting with the first triad of interacting states  $\{(1,0,0), (0,2,0), (0,0,1)\}$ , strong second order C-type Coriolis resonances are coupling together the spin-rotational levels of the  $(v_1, v_2, v_3 \pm 1)$  and  $(v_1, v_2 \pm 2, v_3)$  vibrational states. In addition, a first order C-type Coriolis resonance may have to be accounted for the  $(v_1 \pm 1, v_2, v_3 \pm 1)$  and  $(v_1, v_2, v_3)$  resonating energy levels involving rather high  $K_a$  values [23, 18].

In the rest of the paper, the generic term “energy level parameters” will designate the set of vibrational energies, rotational, spin- rotation, and coupling constants which are associated to a given block of interacting vibration states of  $\text{NO}_2$ .

Several of the computations performed during this study were done using the energy level parameters quoted in the literature, and the form of the Hamiltonian matrices used during these calculations are explicitly described in the attached papers. However, during this work we could make use of new FTS spectra recorded at 6.3  $\mu\text{m}$  and 4.8  $\mu\text{m}$  during this work and of a spectrum recorded at 3.4  $\mu\text{m}$  in Ref. [18] which allowed a detailed investigation of several hot bands. Using these results, we could perform an update of the energy level parameters for the  $\{(0,3,0), (0,1,1)\}$  interacting states and generate the first existing set of energy level parameters



for the  $\{(1,2,1),(1,0,2)\}$  interacting states. The Hamiltonian matrices used during these two computations are given in Table 1.

### III-B Line intensities

A detailed description of the method which is used to compute the NO<sub>2</sub> line intensities was given in Refs. [23, 24]. The intensity,  $k_{\tilde{\nu}}^N(T)$ , of a line of a pure <sup>14</sup>N<sup>16</sup>O<sub>2</sub> isotopic sample is given (in cm<sup>-1</sup>/(molecule×cm<sup>-2</sup>)) by :

$$k_{\tilde{\nu}}^N(T) = a_{\text{iso}} \frac{8\pi^3 \tilde{\nu}}{4\pi\epsilon_0 3hc} \exp\left(-\frac{E_L}{kT}\right) \left(1 - \exp\left(-\frac{\tilde{\nu}}{kT}\right)\right) \frac{g_L}{Z(T)} R_L^U \quad \text{Eq.(1)}$$

In this expression,  $a_{\text{iso}}$  is the isotopic concentration of the considered isotopic species ( $a_{14}=0.991616$  and  $a_{15}=0.00364$  for <sup>14</sup>N<sup>16</sup>O<sub>2</sub> and <sup>15</sup>N<sup>16</sup>O<sub>2</sub>, respectively [25]).  $\tilde{\nu} = (E_U - E_L) / hc$  is the wavenumber of the transition, and  $E_L$  and  $E_U$  are the energies of lower and upper levels of the transition (in cm<sup>-1</sup>). The total partition function,  $Z(T) = Z_{\text{vib}}(T) \times Z_{\text{rot}}(T)$ , is  $Z_{14}(296 \text{ K}) = 13618$  for <sup>14</sup>N<sup>16</sup>O<sub>2</sub> which includes the nuclear spin contribution ( $g_{\text{Nuclear}}^{14} = 2I + 1 = 3$ ) (resp.  $Z_{15}(296 \text{ K})=9324.7$ , and  $g_{\text{Nuclear}}^{15} = 2I + 1 = 2$ , for <sup>15</sup>N<sup>16</sup>O<sub>2</sub>).

Finally,  $R_L^U$  is the square of the matrix element of transformed transition moment operator for the considered band. Let us remind that, as for line positions, the intensity computations account of the vibration-rotation resonances, of the centrifugal distortion effects, and of the electron-spin rotation resonances.

For the line intensities of the strong  $\nu_3$  band at 6.2  $\mu\text{m}$  and for the rather strong  $\nu_1+\nu_3$  at 2.3  $\mu\text{m}$  of <sup>14</sup>N<sup>16</sup>O<sub>2</sub> and <sup>15</sup>N<sup>16</sup>O<sub>2</sub> we used the expansion of the transition moment parameter given in Ref. [23] and Ref. [26], respectively. For the transition moment operators of the hot bands associated to  $\nu_3$  and  $\nu_1+\nu_3$  bands, at 6.2  $\mu\text{m}$  and 2.3  $\mu\text{m}$ , we used the harmonic approximation, which leads to:

$${}_{(v_1, v_2, v_3+1)}(v_1, v_2, v_3) \mu'_Z = \sqrt{v_3} \times {}_{(0,0,0)}(0,0,1) \mu'_Z \quad \text{Eq. (2-a)}$$

and

$${}_{(v_1, v_2, v_3)(v_1+1, v_2, v_3+1)}\mu'_Z = \sqrt{v_1 \times v_3} \times {}_{(0,0,0)(1,0,1)}\mu'_Z \quad \text{Eq.( 2-b)}$$

During this work, several weak cold bands were also considered for GEISA-19, and details on the line intensity parameters that we used for these bands are given in the text.

### III-C Line shape parameters

In Ref. [27], Benner *et al.* measured a very complete set of accurate line shape parameters for the  $v_3$  band of  $\text{NO}_2$ . Following these results, polynomial expansions in ‘ $m$ ’ ( $m=N_{\text{lower}}$  for P or Q lines, and  $m=N_{\text{upper}}$  for the R transitions) were proposed for the air-broadening linewidths ( $\gamma_{\text{Air}}$ ) and their associated “ $n$ ”- temperature dependences and for the pressure shift ( $\delta^0$ ) parameters.

The polynomial expressions achieved for  $\gamma_{\text{Air}}$  and “ $n$ ” in Ref. [27] for the  $v_3$  band were implemented for all vibrational bands in the GEISA-19 database. This means that we assume that the vibrational dependence of  $\gamma_{\text{Air}}$  and “ $n$ ” is negligible. More precisely, the air broadening linewidths ( $\gamma_{\text{Air}}$ ) were computed using the ‘ $m$ ’ polynomial expressions that are quoted for  $K_a=0$  to  $K_a=8$  in Figs. 5 of Ref. [27], and for larger  $K_a$  values ( $K_a > 8$ ) we used the  $K_a=8$  expression. These polynomial expressions can lead, for large “ $m$ ” values, to unreasonably weak values for  $\gamma_{\text{Air}}$ : in such conditions, we fixed  $\gamma_{\text{Air}}$  to a default value ( $\text{Default}\gamma_{\text{Air}}=0.06 \text{ cm}^{-1}/\text{atm}$ ) which corresponds to the asymptotic behaviour of the plots presents on Figs. 5 of Ref. [27]. For the air broadening associated “ $n$ ”- temperature dependences, we used the expression quoted in Fig. 8 of Ref. [27].

As far as the  $\delta^0$  pressure shifts are concerned, we used the “ $m$ ”- expansions given on Fig. 9 of Ref. [27] for the  $v_3$  bands of the  $^{14}\text{NO}_2$  and  $^{15}\text{NO}_2$  species, while for other vibrational bands, the pressure shifts were set to a (default) zero value. This is because it was established that the results on pressure shifts for the  $v_3$  band are not transferable to any other vibrational band [28]. Finally, the default value  $\gamma_{\text{Self}}=0.095 \text{ cm}^{-1}/\text{atm}$  [29] was set up for the self-broadening parameters.

This strategy clearly differs clearly from the one which was adopted in HITRAN2016- updated and HITEMP, where a vibrational dependence of the air- broadening and self linewidths, of the “ $n$ ”- temperature dependence of air- broadening linewidth, and of the pressure

line shift parameters was computed (see Ref. [3] and Ref. [5] and references therein), starting from the experimental results of Ref. [27].

Therefore, our computed line shape parameters differ from those quoted in HITEMP and HITRAN2016-updated.

#### IV. Experimental details

Three high-resolution absorption spectra (FTS12, FTS3 and FTS3bis) of nitrogen dioxide were recorded on the Bruker IFS125HR Fourier transform spectrometer of the AILES Beamline at Synchrotron SOLEIL, coupled to the newly developed corrosive gas multipass cell [30] set to a 10.88 m path length. For the spectrum in the  $\nu_1 + \nu_2$  region, (FTS12) the cell path length was set to 10.88 m, the interferometer was equipped with a Si/CaF<sub>2</sub> beam splitter and an InSb detector. The spectral resolution was chosen to give an apparatus function ( $0.0028 \text{ cm}^{-1}$ ) smaller than the Doppler width (ca.  $0.0048 \text{ cm}^{-1}$ ) in the considered spectral domain. The instrument was operated with a 1.3mm diameter entrance aperture and a quartz-halogen source, as the synchrotron source presents no advantage at this resolution in this spectral domain. For the second spectrum, in the  $\nu_3$  region, (FTS3 and FTS3bis), the cell path length was set to 2.72m and a Ge/KBr beamsplitter, a home-made 4K-cooled HgCdTe with a 1180-1880  $\text{cm}^{-1}$  cold band pass filter and 1.2 mm aperture [31] were used. The instrument was operated with a 1.3mm aperture and at  $0.002 \text{ cm}^{-1}$  resolution, well below the Doppler width. The spectrometer was evacuated to about  $5 \times 10^{-5}$  hPa in order to minimize H<sub>2</sub>O and CO<sub>2</sub> absorptions. Spectra were ratioed against single channel background spectra of the empty cell, recorded at a resolution of  $0.04 \text{ cm}^{-1}$  in order to ensure the best possible signal-to-noise in the ratioed spectrum. For the Fourier transform, a Mertz-phase correction with  $2 \text{ cm}^{-1}$  phase resolution, a zero-filling factor of 2 and no apodization (boxcar option) were applied to the averaged interferograms (880 scans). The spectra were calibrated with residual CO<sub>2</sub> or H<sub>2</sub>O lines observed in the spectra with their wavenumbers taken from HITRAN database [2]. The standard deviation after calibration with well isolated reference gas lines is  $0.00005 \text{ cm}^{-1}(1\sigma)$ . Thus, the estimated frequency accuracy of our measured lines is thus close to the lines reported accuracy ( $0.0001 \text{ cm}^{-1}$ ), estimated to be  $0.0002 \text{ cm}^{-1}$ .

The NO<sub>2</sub> gas bottle used (Sigma-Aldrich, France 99.5%) was found to contain NO, N<sub>2</sub>O and other impurities at a much higher level than the stated purity. It was then first purified following the standard procedure [23] by pumping on the frozen solid at about 200K until the bluish colour due to N<sub>2</sub>O<sub>3</sub> disappeared. This eliminated most of the impurities. We thus added a further step by letting 5 mmoles of the gas mixture react with about 0.5 mmole of ozone, prepared separately from 99.999% pure O<sub>2</sub>. The remaining ozone and oxygen were removed by pumping above the mixture re-cooled to about 210K. This successfully removed the NO and N<sub>2</sub>O traces. The total pressures were measured using a Pfeiffer 10 hPa capacitive gauge. A small contamination due to CO<sub>2</sub> remained visible, but could be estimated to about 0.2% of the gas sample, using IR integrated intensity measurements.

In the near infrared,  $\nu_1+\nu_2$  region, the sample gas pressure used was  $5.1 \pm 0.1$  hPa. Assuming a negligible contribution of foreign gases, the total pressure can be attributed to the mixing of the monomer (NO<sub>2</sub>) and dimer (N<sub>2</sub>O<sub>4</sub>) forms of nitrogen dioxide. These two forms exist in equilibrium, according to the equation ( $2\text{NO}_2 \leftrightarrow \text{N}_2\text{O}_4$ ), with:

$$P^2(\text{NO}_2) = K_p P(\text{N}_2\text{O}_4) \quad \text{Eq. (3)}$$

where  $K_p$  is the equilibrium constant between NO<sub>2</sub> and N<sub>2</sub>O<sub>4</sub>, whose value depends on temperature [32] (at 296K,  $K_p = 123.3$  hPa)  $P(\text{NO}_2)$  and  $P(\text{N}_2\text{O}_4)$  are the partial pressures of the monomer and dimer, respectively. In such conditions, the monomer (NO<sub>2</sub>) and dimer (N<sub>2</sub>O<sub>4</sub>) partial pressures can be estimated in the at about  $P(\text{NO}_2) \approx 4.9 \pm 0.1$  hPa,  $P(\text{N}_2\text{O}_4) \leq 0.2$  hPa. For the second (FTS3) and third (FTS3bis) spectra, in the  $\nu_3$  mid-infrared range, the sample gas pressure was much lower, (0.20 hPa and 0.62 hPa, respectively) and although N<sub>2</sub>O<sub>4</sub> is detectable, due to the long optical path length its partial pressure can be estimated to be quite low, on the order of 0.001 hPa and 0.003 hPa, respectively.

Fig. (A1) and (A2) give overviews of the FT3 and FT12 spectra. While saturated in its central part, the FTS3 and FRS3bis spectra were used in its P branch to identify lines belonging to hot bands and test the quality of the GEISA-19, HITRAN2016-updated, and HITEMP databases. The FTS12 spectrum corresponds to the B-type  $\nu_1+\nu_2$  band.

Also, we used FTS spectra, or portions of spectra, recorded during previous NO<sub>2</sub> investigations. Table 2 summarizes the information on these spectra (FTS13, FTS23, FTS33, FTS113 and FTS333) which are associated to the investigation of the  $2\nu_1$  and  $\nu_1+\nu_3$  bands [18],

the  $\nu_2+\nu_3$  band [29], the  $2\nu_2+\nu_3$ ,  $4\nu_2$ , and  $2\nu_3$  [34] interacting bands, and the  $2\nu_1+\nu_3$  and  $3\nu_3$  bands [35]. All spectra were recorded on Bruker instruments, type IFS125HR (FTS12, FTS3, FTS3bis, FTS13) or FTS120HR (other spectra), at a resolution ranging as  $R=0.002 - 0.003 \text{ cm}^{-1}$ , and at room temperature conditions.

Some additional comments concern spectra FTS113 and FTS333 which were used to perform the most recent line positions and (relative) line intensities investigations for the  $2\nu_1+\nu_3$  and  $3\nu_3$  bands of  $\text{NO}_2$  [35]. Hargreaves et al. [3] pointed out that the experimental intensities measured for these two bands in Ref. [35] lead to cross sections which are, on the average, 3.5 times weaker than those delivered at  $25^\circ\text{C}$  by the Pacific Northwest National Laboratory (PNNL) [36]. At the present time we decided to adopt the strategy suggested by Hargreaves et al. [3], this means that in GEISA-19, the intensities of the  $2\nu_1+\nu_3$  and  $3\nu_3$  bands were calibrated on the absolute scale using the PNNL cross sections.

### **V The GEISA-19 database in the far infrared and $13.3 \mu\text{m}$ region, and status of the energy parameters for the $(0,0,0)$ , $(0,1,0)$ and $\{(1,0,0),(0,2,0),(0,0,1)\}$ states of $^{14}\text{N}^{16}\text{O}_2$ :**

#### V-A The microwave to far infrared region and the $13.3 \mu\text{m}$ regions.

In HITRAN-GEISA-15 the microwave and far infrared region, involves rotational transitions within the  $(0,0,0)$  and  $(0,1,0)$  vibrational state [19,20, 37]. The  $13.3 \mu\text{m}$  region considers the lines belonging to the  $\nu_2$  band [20] and its associated  $2\nu_2-\nu_2$  «first» hot band [33, 20]. The details on the line intensities calculations, models and parameters, are given in Refs. [19] [20] and [33]. As far as the line positions and intensities, these microwave to far-infrared and  $13.3 \mu\text{m}$  linelists were not updated during the present work because we do not have any new spectroscopic data or a new model which encourage us to perform such modification.

#### V-B Status of the $(0,0,0)$ , $(0,1,0)$ and $\{(1,0,0),(0,2,0),(0,0,1)\}$ energy levels parameters

At this level, it is necessary to justify the choice of the energy level parameters used during this work for the computation of energy levels of the  $(0,0,0)$ ,  $(0,1,0)$  vibrational states and for those of the first triad of  $\{(1,0,0),(0,2,0),(0,0,1)\}$  interacting states.

*The (0,0,0), (0,1,0) vibrational states*

During the present update of GEISA-19 we used the (0,0,0) and (0,1,0) energy level parameters quoted in Ref. [37] and [20], respectively, for the whole set of line position computations. These parameters were derived from combined fits of experimental data extracted from high resolution Fourier transform infrared spectra and from microwave measurements. Indeed those (0,0,0) and (0,1,0) energy level parameters are significantly more accurate than those quoted in Ref. [38]: at that time the investigation was performed using infrared grating spectra or diode laser measurements, and the spin-rotation resonances were accounted only through second order perturbation.

*The  $\{(1,0,0),(0,2,0),(0,0,1)\}$  interacting states:*

As far as the  $\{(1,0,0),(0,2,0),(0,0,1)\}$  energy levels are concerned, the  $2\nu_2-\nu_2$  linelist at 13.3  $\mu\text{m}$  is not consistent with the  $\nu_1$ ,  $2\nu_2$ , and  $\nu_3$  linelist at 6.2  $\mu\text{m}$ . This is because two different sets of energy level parameters, Ref. [33] at 13.3  $\mu\text{m}$ , and Ref. [23] and 6.2  $\mu\text{m}$ , were used at the computation of the upper state energy levels. Historically, a first set of energy level parameters for the first triad of  $^{14}\text{NO}_2$ , was obtained in Ref. [23] using as input to the least squares fit calculation experimental energy levels for the  $\{(1,0,0),(0,2,0),(0,0,1)\}$  interacting states that were obtained during the investigation of high-resolution FTS spectra at 6.2  $\mu\text{m}$ . At that time, the identification of the  $\nu_1$  and  $\nu_3$  bands was very extended, while for  $2\nu_2$ , only several series were observed, which involves for the (0,2,0) vibration state only the  $K_a=0-2$  and  $K_a=6$  rotational quantum numbers. The additional missing information for the  $K_a=3,4$  and 5 series of (0,2,0) could be filled through the latter identification of the  $2\nu_2-\nu_2$  hot band in high resolution spectra recorded at 13.3  $\mu\text{m}$  [33]. Therefore, in Ref. [33] the energy levels parameters for the first triad of  $^{14}\text{NO}_2$  were refined during a least squares fit which combined together the experimental spin-rotation energy levels of the (0,0,1), (1,0,0) and (0,2,0) states originating from both analyses, at 6.2  $\mu\text{m}$  [23], and at 13.3  $\mu\text{m}$  [33]. This justify the choice of the parameters of Ref. [33] rather than those of Ref. [23] for the computations the (0,0,1), (1,0,0) and (0,2,0) upper or lower state energy levels in the whole GEISA-19 database.

**VI The GEISA-19 database in the 1153-4775  $\text{cm}^{-1}$  spectral region**

For wavenumbers greater than 1153  $\text{cm}^{-1}$ , GEISA-19 was largely modified as compared to its previous version. First, we updated the lists for the main cold bands and the first hot bands

which were the only originally present in HITRAN-GEISA-15 at 6.2 and 3.4  $\mu\text{m}$ . Also, we add new linelists which concern (i) higher order hot bands in the 6.2  $\mu\text{m}$  and 3.4  $\mu\text{m}$  regions (ii) several weaker bands in the 2.3 to 4.9  $\mu\text{m}$  region, and (iii) the  $\nu_3$  and  $\nu_1+\nu_3$  bands for the  $^{15}\text{N}^{16}\text{O}_2$  isotopic species of nitrogen dioxide.

For GEISA-19, we tried to generate accurate vibration-rotation bands each time this proved to be “technically possible”. The limitation concerns the quality of line position and line intensity parameters that we had at our disposal. For example, the very weak  $\nu_1+2\nu_3$  weak cold band at 2.2  $\mu\text{m}$ , present in HITRAN-2016 and HITEMP, is absent in GEISA-19 because we do not have at our disposal intensity parameters for this band. Contrary to HITEMP, GEISA-19 does not include lines involving very high rotational quantum numbers, and numerous higher order hot bands are missing.

#### VI-A Updates of the lists included in HITRAN-GEISA-15

##### *V-A-1 The cold bands in the 6.2 $\mu\text{m}$ and 3.4 $\mu\text{m}$ regions for $^{14}\text{N}^{16}\text{O}_2$ .*

The 6.2  $\mu\text{m}$  bands

For GEISA-19 a new linelist was generated for the  $\nu_1$ ,  $2\nu_2$ , and  $\nu_3$  interacting bands at 6.2  $\mu\text{m}$  using the energy level parameters of Ref. [33]. Let us mention that the  $\nu_3$  and  $\nu_1$  bands are only weakly affected in this new calculation as compared to the previous linelist [39]. In fact, only the very weak  $K_a=3-5$  series, or those with  $K_a \geq 7$  of the dark  $2\nu_2$  band have their positions significantly modified.

The 3.4  $\mu\text{m}$  bands:

The recent update of the  $\nu_1+\nu_3$  band at 3.4  $\mu\text{m}$  was presented in details in Ref. [18]. Let us recall that this recent analysis was unable to assign  $\nu_1+\nu_3$  lines up to  $N=75$  and  $K_a=14$  (instead of  $N=55$  and  $K_a=10$ , in Ref. [26]). In addition, a new investigation of the  $2\nu_1$  band was performed. The energy level computation involves six interacting states  $\{(2,0,0),(1,2,0),(1,0,1),(0,4,0),(0,2,1),(0,0,2)\}$ , instead of two  $\{(1,2,0),(1,0,1)\}$  previously [26]. Also, detailed line intensity measurements and calculations were performed for the  $2\nu_1$  band [18]. Indeed, this band exhibits an “unusual” intensity behaviour for a B-type band, with, for example, a very weak P-branch. The progress achieved for the  $\nu_1+\nu_3$  and  $2\nu_1$  linelists [18] as compared to their counterparts in HITRAN2016-updated [6] and HITEMP [3] are resumed shortly later in the text.



*VI-A-2 The “first” hot bands in the 6.2  $\mu\text{m}$  and 3.4  $\mu\text{m}$  regions for  $^{14}\text{N}^{16}\text{O}_2$ .*

We will present the updates performed for the “first” hot bands associated to the strongest absorbing bands of  $\text{NO}_2$  in the 6.2  $\mu\text{m}$  and 3.4  $\mu\text{m}$  regions.

The 6.2  $\mu\text{m}$  bands and update of the energy level parameters for the  $\{(0,3,0),(0,1,1)\}$  interacting states.

As far as the line positions are concerned, the list for the  $\nu_2+\nu_3-\nu_2$  hot band present in HITRAN-GEISA-15 was generated using for the  $\{(0,3,0),(0,1,1)\}$  upper and  $(0,1,0)$  lower states the energy level parameters quoted in Ref. [38]. Let us remind that these parameters were achieved using a composite set of experimental data (grating and diode laser spectra). In addition, the spin-rotation resonances were only accounted for through a second order treatment. This  $\nu_2+\nu_3-\nu_2$  list, which is unchanged in HITRAN2016-updated, was only extended in HITEMP up to higher values of the  $N$  and  $K_a$  rotational quantum numbers.

For the preparation of GEISA-19 we generated a new version of the linelist for this hot band. In a preliminary step, the line positions were computed using the upper and lower state energy level parameters of Refs. [29] and [20] for the  $\{(0,3,0),(0,1,1)\}$  and  $(0,1,0)$  states, respectively. The comparison between calculated spectrum and the observed one (the FTS3 or FTS3bis spectra in the P branch region) proved that this preliminary  $\nu_2+\nu_3-\nu_2$  linelist was already more precise than the existing one in HITRAN-GEISA-15, although not perfect. Indeed, we notice that the  $\nu_2+\nu_3-\nu_2$  transitions involving high  $K_a$  values ( $11 \geq K_a \geq 9$ ) differ from their predicted positions. We used these  $\nu_2+\nu_3-\nu_2$  new assignments to generate an additional set of 37 experimental spin-rotational energy levels for the  $(0,1,1)$  state by adding the  $\nu_2+\nu_3-\nu_2$  line positions to the calculated  $(0,1,0)$  state energy levels [20]. These additional set of  $(0,1,1)$  levels were combined with those already measured during the investigation of the  $\nu_2+\nu_3$  band [29], and these data were introduced in a least squares fit calculation to get refined values for the  $\{(0,3,0),(0,1,1)\}$  interacting states parameters. These constants which are given, together with their associated uncertainties in Table 3, were used to generate the final linelist for the  $\nu_2+\nu_3-\nu_2$  band which is quoted in GEISA-19.

The list of assigned  $\nu_2+\nu_3-\nu_2$  transitions involving high  $K_a$  values ( $11 \geq K_a \geq 9$ ), and the results of the energy level computation for the  $\{(0,3,0),(0,1,1)\}$  state are given in the Supplementary data of this article.



The 3.4  $\mu\text{m}$  bands:

In HITRAN2016-updated the list for the  $\nu_1+\nu_2+\nu_3-\nu_2$  band, located at 2888.191  $\text{cm}^{-1}$ , coincides with HITRAN-GEISA-15's. This list, which is restricted to the 2850.082 - 2917.363  $\text{cm}^{-1}$  spectral region, was generated using for the energy positions in the upper  $\{(1,2,0),(1,1,1)\}$  and lower (0,1,0) states the parameters quoted in Ref. [40] and [38], respectively. As for its  $\nu_2+\nu_3-\nu_2$  counterpart at 6.2  $\mu\text{m}$ , this list is not very accurate for the same reasons (low quality of the theoretical model and of the experimental data used as input to the calculation). For HITEMP the  $\nu_1+\nu_2+\nu_3-\nu_2$  linelist was extended to include lines involving higher N and Ka values.

For GEISA-19, the generation of the new linelist for the  $\nu_1+\nu_2+\nu_3-\nu_2$  band was performed using for the line positions the upper  $\{(2,1,0),(1,3,0),(1,1,1)\}$  and lower (0,1,0) energy level parameters collected in Ref. [17] and [20], respectively.

#### VI-B New input in GEISA-19: higher order hot band in the 6.2 $\mu\text{m}$ and 3.4 $\mu\text{m}$ regions

##### *VI-B-1 Higher order hot bands in the 3.4 $\mu\text{m}$ region, and first investigation of the $\nu_1+2\nu_3-\nu_3$ hot band:*

In the 3.4  $\mu\text{m}$  region, the three stronger second order hot bands are the  $2\nu_1+\nu_3-\nu_1$ ,  $\nu_1+2\nu_3-\nu_3$ , and  $\nu_1+2\nu_2+\nu_3-2\nu_2$  hot bands, located at (about) 2860.144, 2844.218 and 2870.76  $\text{cm}^{-1}$ , with a relative intensity at 296K of 0.32%, 0.077%, and 0.065%, respectively as compared to the cold  $\nu_1+\nu_3$  band (located at 2906.074  $\text{cm}^{-1}$ )

It was rather easy to generate a linelist for the  $2\nu_1+\nu_3-\nu_1$  band using for the computation of the upper and lower state levels the energy level parameters of Refs [35] and [33], respectively. In this way, lines of the  $2\nu_1+\nu_3-\nu_1$  hot band could be identified in the FTS13 spectrum within strong lines of the P branch of the  $\nu_1+\nu_3$  band

On the other hand, the situation for the  $\nu_1+2\nu_3-\nu_3$  and  $\nu_1+2\nu_2+\nu_3-2\nu_2$  hot bands was not so simple since, to our knowledge, the very weak  $\nu_1+2\nu_3$  and  $\nu_1+2\nu_2+\nu_3$  bands were never investigated at high or medium resolution. The only existing information concerns the band centers which were measured by laser fluorescence spectroscopy [41] at 4461.07 and 4369.10  $\text{cm}^{-1}$  for the (1,0,2) and (1,2,1) vibrational states.

Therefore, the right method was to try to identify these hot bands in the congested FTS13 spectrum which was recorded for a large (pressure  $\times$  path length) product (see Ref. [18]),

because numerous lines were unidentified at that time. To solve this problem, we generated predictive line lists for the  $\nu_1+2\nu_3-\nu_3$  and  $\nu_1+2\nu_2+\nu_3-2\nu_2$  bands.

We were unfortunately not able to identify the  $\nu_1+2\nu_2+\nu_3-2\nu_2$  band which is located in a congested part of the FTS13 spectrum. Fortunately, we were luckier for the  $\nu_1+2\nu_3-\nu_3$  hot band, as the line prediction helped us to perform several assignments. The calculated (0,0,1) state energy levels [33] were added to the first observed  $\nu_1+2\nu_3-\nu_3$  line positions to get a preliminary set of (1,0,2) experimental spin-rotational energy levels. These upper state levels were then introduced in a least squares fit calculation to get refined values for the  $\{(1,0,2)(1,2,1)\}$  energy level parameters, and therefore perform new assignments, and then improve the accuracy of the parameters, allowing then further new assignments. This iterative process was performed until no further assignments of  $\nu_1+2\nu_3-\nu_3$  lines could be possible. Altogether, about 550 lines (resolved or unresolved spin doublets) were assigned with  $K_a$  and  $N$  values such that  $(0 \leq K_a \leq 8)$   $(0 \leq N \leq 45)$ , leading to the determination of 394 energy levels of the (1,0,2) state.

The final set of vibrational energies, spin rotation, rotational and interacting constants for the  $\{(1,0,2)(1,2,1)\}$  interacting states are collected in Table 3. During this computation, the  $E_v$  (vibrational energy) of the (1,2,1) was maintained at the value proposed by Delon and Jost [41] ( $E_{121} = 4369.10 \text{ cm}^{-1}$ ). The  $E_v$  value of the (1,0,2),  $E_{102} = 4460.8648 \pm 0.0003 \text{ cm}^{-1}$  is in reasonable agreement with the laser fluorescence value proposed by Delon and Jost [41] ( $E_{102} = 4461.07 \text{ cm}^{-1}$ ). We also notice that the  $K_a=4$  levels of the (1,0,2) state are resonating with  $K_a=5$  levels of (1,2,1), and that the mixing of the energy levels grows up to 44% for  $N=44$ . We used these energy level parameters together with the first triad constants of Ref. [33] to generate the line positions of the  $\nu_1+2\nu_3-\nu_3$  hot band. The associated intensities were computed using the method described previously.

The complete list of assignments and the results of this calculation are provided as supplementary data of this article.

#### *VI-B-2 Higher order hot bands in the 6.2 $\mu\text{m}$ region*

In the 6.2  $\mu\text{m}$  region, the second order hot bands which lead to relative strong contributions are those associated to  $|\Delta\nu_3|=1$  i.e. the  $2\nu_2+\nu_3-2\nu_2$ ,  $2\nu_3-\nu_3$  and  $\nu_1+\nu_3-\nu_1$  hot bands located at 1594.14, 1584.59 and 1586.95  $\text{cm}^{-1}$ . The line positions were computed using

the energy level parameters of Ref. [18] and of Ref. [33] for the upper and lower states respectively.

Whenever possible, we generated also higher order hot bands. This was possible in the 6.2  $\mu\text{m}$  region for the  $\nu_1+\nu_2+\nu_3-(\nu_1+\nu_2)$ ,  $2\nu_1+\nu_3-2\nu_1$ ,  $\nu_1+2\nu_3-(\nu_1+\nu_3)$  and  $3\nu_3-2\nu_3$  hot bands. For the upper state energy levels we used the energy level parameters of the (1,1,1) (Table 3 of Ref. [17]), (2,0,1), and (0,0,3) (Table 3 and 4 of Ref. [35]), and for (1,0,2) those in Table 3 of this work, while for the (1,1,0) and  $\{(2,0,0),(1,2,0),(1,0,1), (0,0,2), (0,4,0),(0,0,2)\}$  states, we used the parameters quoted in Table 1 of Ref. [42] and Table 4 of Ref. [18], respectively.

#### VI-C New input in GEISA: weak cold bands in the 2.2 to 4.9 $\mu\text{m}$ spectral regions :

The new versions of HITRAN2016-updated [6] and HITEMP [3] databases include linelists for numerous weak bands which were not considered in HITRAN-GEISA-15. For several of these rather dark bands, this is now the case for GEISA-19.

□ The  $2\nu_1$  band centered at 2627.34  $\text{cm}^{-1}$  was already the subject of an extensive study (positions and intensities) in Ref. [18], leading to a linelist now included in GEISA-19.

□ GEISA-19 includes linelists from the  $\nu_1+\nu_2$  band [42],  $2\nu_2+\nu_3$  and  $2\nu_3$  bands [34,14], and  $2\nu_1+\nu_3$  and  $3\nu_3$  bands [35], located at 2063.12, 3092.48, 3201.44, 4179.94, and 4754.21  $\text{cm}^{-1}$ , respectively together with some lines associated to their dark interacting bands. For these bands we used, without modifications, the energy level parameters and the expansion of the transition moment operators which are quoted in the original papers. The only exception concerns the intensities of the  $2\nu_1+\nu_3$  and  $3\nu_3$  bands which were calibrated relative to the PNNL cross sections for  $\text{NO}_2$  [36].

□ GEISA-19 includes linelists from the  $\nu_2+\nu_3$  band located at 2355.15  $\text{cm}^{-1}$ . For this computation we used the improved energy level parameters achieved for the  $\{(0,3,0),(0,2,1)\}$  interacting states during this work (see Table 3) for the computation of the (0,1,1) upper levels, and the expansion of the  $\nu_2+\nu_3$  transition parameter quoted in Ref. [29] for the  $\nu_2+\nu_3$  intensities.

□ The very weak  $\nu_1+\nu_2+\nu_3$  band located at 3637.848  $\text{cm}^{-1}$  [17] is not included in HITEMP or HITRAN2016-updated. The FTS123 Fourier transform spectrum used in Ref. [17] was recorded without purification of the  $\text{NO}_2$  sample. So, the presence of impurities prevents to use this spectrum for line intensities measurements. However, Fig.1 in Ref. [17] shows that the band intensity pattern of the  $\nu_1+\nu_2+\nu_3$  band does not differ from what is expected for a

classical A –type band of nitrogen dioxide, without any significant distortion effect. Therefore, only the first order term in the expansion of the  $\nu_1+\nu_2+\nu_3$  transition moment operator is necessary to compute the line intensities for this band. Its value was estimated to be :

$${}^{111,000}\mu_1' \approx 0.27 \times 10^{-2} \text{ D } (\pm 20 \%)^4, \quad \text{Eq.(4)}$$

using the PNNL cross section data [36]. The uncertainty quoted in Eq. (4) is set up according to the noise level ( $\sim 5.E-22 \text{ cm}^2/\text{molecule}$ ) in the PNNL cross section data.

#### VI-D New inputs in GEISA-19: The $\nu_3$ and $\nu_1+\nu_3$ bands for ${}^{15}\text{N}{}^{16}\text{O}_2$ :

The linelists for the  $\nu_3$  and  $\nu_1+\nu_3$  bands of  ${}^{15}\text{N}{}^{16}\text{O}_2$ , located at 1582.1029 and 2858.7088  $\text{cm}^{-1}$ ,<sup>5</sup> respectively, were included in GEISA-19. We used the line positions and intensities gathered in Refs. [16], and [43], respectively. The computed partition functions for  ${}^{15}\text{N}{}^{16}\text{O}_2$  at different temperatures (T=50K to T=500K) are also provided as Supplementary data of the present article.

#### VI-E Final description of the GEISA-19 database in the 1153- 4775 $\text{cm}^{-1}$ spectral region:

The GEISA-19 database for nitrogen dioxide is available on the freely accessible <https://geisa.aeris-data.fr/> website. Also, in the 1153-4775  $\text{cm}^{-1}$  spectral range, which is the object of this study, this linelist is provided as supplementary data of this work. Table 4 gives a global description of this linelist, with, for each vibrational band and isotopic species ( ${}^{14}\text{N}{}^{16}\text{O}_2$  and  ${}^{15}\text{N}{}^{16}\text{O}_2$ ), the vibrational assignment, the number of lines, an approximate band center, the wavenumber range (in  $\text{cm}^{-1}$ ), the band intensity and the minimum and maximum values of the individual line intensities intensity (in  $\text{cm}^{-1}/(\text{molecule.cm}^{-2})$  at 296K), the maximum N values and the range for the  $K_a$  rotational quantum numbers in the upper vibrational state.

<sup>4</sup> in Debye (1 D =  $3.33564 \times 10^{-30} \text{ C m}$ ),

<sup>5</sup> The actual band center of the  $\nu_1+\nu_3$  band is 2858.7088  $\text{cm}^{-1}$  instead of 2858.7077  $\text{cm}^{-1}$  in Ref. [43]. Also the  $Z_{15}(296\text{K})=9324.7$  value was used during the intensities computations for the  $\nu_1+\nu_3$  band, instead of the erroneous value “ $Z_{15}(296\text{K})=13987$ ” which is quoted in Ref. [43].

It is necessary to estimate the accuracy attached to the computed line positions and intensities. In GEISA-19, each of vibration – rotation linelist was generated through calculations. It is obvious that extrapolations far away from the range of the lower or upper states experimental data used during this computation can be uncertain. To document this point, Table 4 provides for each vibrational band the maximum N value and the  $K_a$  range for the upper and lower states experimental energy levels used during the computation.

One can estimate at about 0.001- 0.003  $\text{cm}^{-1}$  the accuracy of the line position for transitions involving rotational quantum numbers within the range of the observed energy levels. For transitions outside these limits, the line position accuracy is decreasing with increasing rotational quantum

As far as the line intensities are concerned, the information describing the quality and accuracy of these data are provided in the literature articles which are referred in this paper.

## **VII-Validations of the new GEISA-19 database.**

Two different types of validations for the GEISA-19, HITEMP and HITRAN2016-updated linelists were performed. First the existing laboratory FTS spectra (see Table 2) were used to check these linelists through inter-comparison between observed and calculated spectra. In a second step, we checked, within these databases, the internal consistency of the computed energy levels in term of the Rydberg–Ritz combination principle.

### VII-A Validation with the observed laboratory spectra at 296K.

To validate GEISA-19 we compared the calculated and observed spectra using the laboratory FTS spectra at our disposal (see Table 2). The same intercomparison was performed using the HITEMP or HITRAN2016-updated linelist.

#### *VII-A-1 The 6.2 $\mu\text{m}$ region:*

We used the FTS3 and FTS3bis spectra recorded during this work to validate our computations.

Figures 1 to 5 detail the computations for several vibrational bands which were computed for GEISA-19 during this work.

Figure 1 presents a portion of the spectrum in the 1557  $\text{cm}^{-1}$  region where lines from the  $\nu_2 + \nu_3 - \nu_2$  hot band are clearly observable.

Figure 2, 3, and 4 present a portion of the same spectrum in the 1555, 1561 and 1572  $\text{cm}^{-1}$  spectral ranges where lines from the P branches of the  $\nu_1+\nu_3-\nu_1$ ,  $2\nu_3-\nu_3$  and  $2\nu_2+\nu_3-2\nu_2$  hot bands are clearly visible.

Finally, Figure 5 present a portion of the P- branch of the  $\nu_3$  band in the 1559.1  $\text{cm}^{-1}$  region. In this case a global comparison GEISA-19 versus HITEMP and HITRAN2016-updated is performed.

#### *VII-A-2 The 2.3 $\mu\text{m}$ region:*

For this validation we used the FTS13 spectrum recorded in Ref. [18].

Figures 6, 7, and 8 give portions of FT13 spectrum in the 2838.8  $\text{cm}^{-1}$ , 2821.6  $\text{cm}^{-1}$  and 2834.8  $\text{cm}^{-1}$  spectral ranges, and the comparisons focus specifically on the  $\nu_1+\nu_2+\nu_3-\nu_2$ ,  $\nu_1+2\nu_3-\nu_3$  and  $2\nu_1+\nu_3-\nu_1$  hot bands, respectively.

Finally, Figure 9 presents a portion of the 2820  $\text{cm}^{-1}$  spectral region. In this region, lines from the  $\nu_1+\nu_3$  band of  $^{15}\text{N}^{16}\text{O}_2$  are also observable, together with numerous lines from the hot bands from the main species. The comparison of GEISA-19 with HITEMP and HITRAN2016-updated is performed. Note than in HITRAN2016-updated the linelist for the  $\nu_1+\nu_2+\nu_3-\nu_2$  hot band is restricted to the 2850.08 -2917.36  $\text{cm}^{-1}$  spectral region.

#### *VII-A-3 The weak bands:*

During this work several weak bands were added to the GEISA linelist for  $\text{NO}_2$ .

Fig. 10 presents a portion of the FTS13 spectrum in the 2814.4  $\text{cm}^{-1}$  spectral range, where lines from the R branch of the  $2\nu_1$  band are easily observable within from those from the P- branches of the  $\nu_1+\nu_3$  band and associated hot bands. The agreement (positions and intensities) between the observed and calculated spectra is excellent when using GEISA-20 both for the  $2\nu_1$  lines and for the  $\nu_1+\nu_3$  band (and for the associated hot bands), while this is not the case for HITEMP. For this figure, Table TA1 (Supp-5 in the supplementary data) provides the assignments, line positions and line intensities for several  $^{14}\text{N}^{16}\text{O}_2$  lines. As far as the  $\nu_1+\nu_3$ ,  $2\nu_1+\nu_3-\nu_1$ , and  $\nu_1+2\nu_3-\nu_3$  (associated hot) bands are concerned, the intensities in HITEMP and

GEISA-20 are in good agreement, while for the  $2\nu_1$  lines, the intensities in HITEMP are about two times weaker than in GEISA-20. This is also evidenced on Fig. A9 which presents an overview of the  $2\nu_1$  band. Also Fig. A9 in the archive (Supp-5 in the supplementary data) presents an overview of the 2500- 3000  $\text{cm}^{-1}$  region of the FTS13 spectrum which corresponds to the  $2\nu_1$  and  $\nu_1+\nu_3$  bands of  $^{14}\text{N}^{16}\text{O}_2$ , it is clear that that HITEMP underestimates the  $2\nu_1$  band line intensities by a factor of about two, and that the strange intensity pattern of the  $2\nu_1$  band is not accounted for.

For the other weak “cold” bands, the figures presenting intercomparisons are accessible in Supp-5 of the supplementary data of this paper.

Figures A3, A4, A5, A6, A7 and A8 present inter-comparisons for the  $\nu_1+\nu_2$  band near 2090.7  $\text{cm}^{-1}$  (FTS12),  $\nu_2+\nu_3$  band near 2365.2  $\text{cm}^{-1}$  (FTS23),  $2\nu_2+\nu_3$  band near 3069.68  $\text{cm}^{-1}$  (FTS223),  $\nu_1+\nu_2+\nu_3$  band near 3650  $\text{cm}^{-1}$  (FTS111),  $2\nu_1+\nu_3$  band near 4155.2  $\text{cm}^{-1}$  (FTS113) and  $3\nu_3$  band near 4774.8  $\text{cm}^{-1}$  (FTS333).

Except for  $\nu_1+\nu_2+\nu_3$ , all these weak bands are included in HITRAN2016-updated and HITEMP.

#### *VII-A-4 Conclusion for the intercomparison between observed and calculated spectra:*

We notice that the agreement between the observed and calculated spectra is better in all spectral regions when using GEISA-19 rather than with HITRAN2016-updated or HITEMP.

However, the line position differences between databases depend on the considered vibrational transitions.

The  $\nu_3$  band was not significantly updated. For the  $\nu_1+\nu_3$  band, the most significant discrepancies (more than  $\sim 0.01 \text{ cm}^{-1}$ ) concern lines involving high rotational quantum numbers ( $K_a=0$  and  $N \geq 69$ ;  $K_a=5$  and  $N \geq 55$ ; or  $K_a=10$  and  $N \geq 47$ ) [18].

For vibrational transitions involving only excitations in the  $\nu_3$  vibrational mode, like the  $3\nu_3$  weak cold band at 2.1  $\mu\text{m}$ , the  $2\nu_3-\nu_3$  and  $3\nu_3-2\nu_3$  hot bands at 6.2  $\mu\text{m}$ , the differences are in general weaker than for other type of vibrational transitions. However, differences larger than  $\sim 0.005 \text{ cm}^{-1}$  concern the lines involving  $N$  and  $K_a$  quantum numbers such that  $K_a=1$  and  $N \geq 35$  of  $3\nu_3-2\nu_3$ ;  $K_a=1$  and  $N \geq 49$  for  $2\nu_3-\nu_3$  and  $2\nu_3$ );  $K_a=1$  and  $K_a=2$  and  $N \geq 41$  for  $3\nu_3$ .

More significant discrepancies (larger than  $\sim 0.01 \text{ cm}^{-1}$ ) are observed for the following series ( $K_a=1$  and  $N \geq 23$  of  $2\nu_1+\nu_3$ ;  $K_a \geq 5$  and  $N \geq (50-K_a)$  or  $K_a \geq 9$  of  $\nu_2+\nu_3-\nu_2$ ;  $K_a \geq 5$  and  $N \geq$



( $50-K_a$ ) or  $K_a \geq 9$  of  $\nu_2+\nu_3-\nu_2$ ; for  $K_a \leq 8$  and  $N \geq (42-K_a)$  or  $K_a \geq 9$  for  $\nu_1+\nu_2+\nu_3-\nu_2$ ). Let us remind that the  $\nu_2+\nu_3-\nu_2$  and  $\nu_1+\nu_2+\nu_3-\nu_2$  hot bands give rise to the strongest  $\text{NO}_2$  infrared signature in the 6.3  $\mu\text{m}$  and 3.4  $\mu\text{m}$  region after their associated  $\nu_3$  and  $\nu_1+\nu_3$  associated cold bands.

As expected, larger line position differences (larger than  $\sim 0.02 \text{ cm}^{-1}$ ) between the databases concern bands involving the (2,0,0) or (1,0,2) vibrational states in the upper or lower state of the vibrational transition. However it is clear that these bands are extremely weak, except for hot bands ( $2\nu_1+\nu_3-2\nu_1$ ,  $\nu_1+2\nu_3-2\nu_3$ , etc...) for  $\text{NO}_2$  spectra in very very high temperature conditions.

#### VII-B Consistency of the GEISA-19 database for the $E_L$ or $E_U$ as energy levels values (Rydberg–Ritz combination principle).

Within the HITRAN-GEISA-15, HITRAN2016-updated or HITEMP databases one should expect that the individual values of  $E_L$  (resp. of  $E_U$ ) lower (resp. upper) state energy level values to be consistent from one given vibration–transition band to the other one, provided that they involve the same lower ( $\nu_1^L, \nu_2^L, \nu_3^L$ )  $[N^L, K_a^L, K_c^L, J^L]$  (resp. same upper ( $\nu_1^U, \nu_2^U, \nu_3^U$ )  $[N^U, K_a^U, K_c^U, J^U]$ ) vibrational spin-rotational energy level. However, this not the case because these linelists were generated by merging numerous unrelated data sources. Table 5 provides several examples of vibration spin-rotation transitions for which this is not the case in HITEMP or HITRAN2016-updated. For example, the computed positions of the (0,2,0) [21,3,19] spin rotational levels are different in the  $2\nu_2$ ,  $2\nu_2-\nu_2$ , and  $2\nu_2+\nu_3-2\nu_2$  bands. This is the case also for the (1,1,0) [42,0,42] or (1,1,0) [33, 1, 33] levels in the  $\nu_1+\nu_2$  and  $\nu_1+\nu_2-\nu_2$  bands, for the (0,1,1) [31, 9, 22], (0,1,1) [35, 8,27], (0,1,1) [45, 4, 41], or (0,1,1) [27, 10, 17], in the  $\nu_2+\nu_3$  and  $\nu_2+\nu_3-\nu_2$  bands, or for the (0,2,1) [20,3,18] spin rotational levels in the  $2\nu_2+\nu_3$ ,  $2\nu_2+\nu_3-\nu_2$ , and  $2\nu_2+\nu_3-2\nu_2$  bands. This problem was solved in GEISA-19, and to validate these corrections, Table 5 provides, each time it was possible, the observed line positions which were either observed during previous investigations (see the literature) or identified during the present work.



### VII-C Comparison with the PNNL cross-sections.

In Fig. 11, we compared the cross-sections computed using the HITEMP and GEISA-19 linelists with the PNNL NO<sub>2</sub> cross sections [36]. Although remaining absorption due to the N<sub>2</sub>O<sub>4</sub> dimer are quite observable (in the 1263, 1751, 2630 and 2971 cm<sup>-1</sup> spectral regions) in the PNNL spectra, we can notice the following points.

- In GEISA-19, the 2ν<sub>1</sub> band is significantly stronger than in HITEMP, and this is in agreement with the PNNL data.
- GEISA-19 includes data for the ν<sub>1</sub>+ν<sub>2</sub>+ν<sub>3</sub> band at 3637 cm<sup>-1</sup>, while this linelist is absent in HITEMP.
- HITEMP includes data for the ν<sub>1</sub>+2ν<sub>3</sub> band at ~4500 cm<sup>-1</sup>, while this linelist is absent in GEISA-19. This is because we do not have reliable intensity data for this very weak band which, in Fig. 11 is masked in the noise level (about 5×10<sup>-22</sup> cm<sup>2</sup>/molecule) of the PNNL cross-section.

### **Conclusion**

During this work, a new version of the line by energy level parameters (line positions, intensities and shape) for nitrogen dioxide was generated for the 0- 4700 cm<sup>-1</sup> spectral range and implemented in the (Gestion et Etude des Informations Spectroscopiques Atmosphériques) GEISA database (<https://geisa.aeris-data.fr/>). Except for the far infrared and 13.3 μm regions all spectral regions are significantly affected by this major update.

For the main (<sup>14</sup>N<sup>16</sup>O<sub>2</sub>) isotope species, the updated lists involve the main cold bands (ν<sub>1</sub>, 2ν<sub>2</sub> and ν<sub>3</sub> and ν<sub>1</sub>+ν<sub>3</sub>, respectively) together with their associated «first» hot bands, and, whenever possible, higher order hot bands. One of the outputs of this work was the first identification of the ν<sub>1</sub>+2ν<sub>3</sub>-ν<sub>3</sub> hot band, leading to the first determination of the energy level parameters for the (1,0,2) vibrational state of <sup>14</sup>N<sup>16</sup>O<sub>2</sub>. Also, the quality of the (0,1,1) energy level parameters was improved during the investigation of the ν<sub>2</sub>+ν<sub>3</sub>-ν<sub>2</sub> hot band. In addition, the linelists for several weak cold bands were implemented in GEISA-19 for the first time. For <sup>15</sup>N<sup>16</sup>O<sub>2</sub>, which is the most abundant daughter isotopologue of nitrogen dioxide, we included in GEISA-19 the line parameters for the ν<sub>3</sub> and ν<sub>1</sub>+ν<sub>3</sub> bands.

These new lists, were validated using high resolution Fourier transform spectra recorded at SOLEIL at 296K or during previous laboratory studies. It appeared that, as compared to HITRAN2016-updated and HITEMP, GEISA-19 leads to better agreements between calculated and observed spectra in all spectral regions. In term of line position calculations, GEISA-19 proved to be more consistent than HITRAN2016-updated and HITEMP. However, GEISA-19, which does not include transitions involving very high rotational quantum numbers or belonging to higher order hot bands, is not suitable for uses at high temperatures.

### **Acknowledgments**

The authors acknowledge support from synchrotron SOLEIL (project 99200018). This work was supported by the French National program LEFE « Les Enveloppes Fluides et l'Environnement » through the « MEANSPECFORUM » project. It is also funded by the French National program ANR (ANR-19-CE29-0013) through the « QUASARS » project. Also, we thank the reviewer for his very valuable comments.

### **APPENDIX-A: description of the supplementary data**

**Supp-1:** The GEISA-19 linelist in the 1153-4775  $\text{cm}^{-1}$  spectral range (in the HITRAN's type format [2]).

**Supp-2:** Results of the calculation of the  $\{(0,1,1),(0,3,0)\}$  energy levels. List of assignments for the  $\nu_2+\nu_3-\nu_2$  hot band. List of assignments for the  $(0,1,1)-(0,1,0)$  hot band.

**Supp-3:** Results of the calculation of the  $\{(1,0,2),(1,2,1)\}$  energy levels, and list of assignments for the  $\nu_1+2\nu_3-\nu_3$  hot band. List of assignments for the  $(1,0,2)-(0,0,1)$  hot band.

**Supp-4:** Partition function for  $^{15}\text{N}^{16}\text{O}_2$ .

**Supp-5:** Table TA1 and list of nine figures (A1-A9).

Table TA1 :

Comparison of the line positions and intensities in the 2814-2815  $\text{cm}^{-1}$  spectral range of the FTS13 spectrum (see Fig. 10)

Fig A1 : Overview of the FTS3 spectrum ( $\nu_3$  band spectral region)

Fig A2: Overview of the  $\nu_1+\nu_2$  band (FTS12).

Fig A3: Portion of the R branch of the  $\nu_1+\nu_2$  band in the 2090 - 2093  $\text{cm}^{-1}$  region (FTS12).

We give some  $[N, K_a, K_c]$  ( $J=N\pm 1/2$ ) rotational assignments for the  $\nu_1+\nu_2$  band. A comparison between the model provided by GEISA-19 and HITRAN2016-updated (or HITEMP) is given.

Fig. A4 : Portion of the R branch of the  $\nu_2+\nu_3$  band in the 2365.2  $\text{cm}^{-1}$  region (FTS23). In the bottom part of the figure, the quoted  $[N, K_a, K_c]$  (and eventually J) assignments are for the (0,1,1) upper state energy levels. A comparison between the model provided by GEISA-19 and HITRAN2016-updated (or HITEMP) is given.

Fig. A5 : Portion of the P- branch of the  $2\nu_2+\nu_3$  band near 3079.8  $\text{cm}^{-1}$  (FTS223). In the upper part of the figure, the quoted  $[N, K_a, K_c]$  (and eventually J) assignments are for the (0,2,1) upper energy levels. The GEISA-19 and HITEMP (or HITRAN2016 -updated) linelists are compared.

Fig. A6: Portion of the R- branch of the  $\nu_1+\nu_2+\nu_3$  band of  $\text{NO}_2$  in the 3650  $\text{cm}^{-1}$  region (FTS111 spectrum). Several  $N'$  assignments in the upper state are given for,  $[N, K_a = 0, K_c = N] = [N-1, 0, N-1]$  (triangles). This band is not considered in the HITEMP or HITRAN2016-updated databases.

Fig. A7 : Portion of the R- branch of the  $2\nu_1+\nu_3$  band near 4155.2  $\text{cm}^{-1}$  (FTS113). In the bottom part of the figure, the quoted  $[N, K_a, K_c]$  (and eventually J) assignments are for the (2,0,1) upper energy levels of the  $2\nu_1+\nu_3$  band. Due to a local resonance the  $[23,3,21]$ - $[24,0,24]$  transition of the dark  $2\nu_1+2\nu_2$  band are observable on the spectrum. The GEISA-19 and HITEMP (or HITRAN2016-updated) linelists are compared.

Fig. A8 : Portion of the R- branch of the  $3\nu_3$  band near  $4774.8\text{ cm}^{-1}$  (FTS333). In the bottom part of the figure, the quoted  $[N, K_a, K_c]$  (and  $J$ , eventually) assignments are for the  $(0,0,3)$  upper energy levels of the  $3\nu_3$  band. The GEISA-19 and HITEMP linelists are compared.

Fig. A9 : Overview of the FTS13 spectrum in the  $2500\text{-}3000\text{ cm}^{-1}$  spectral region of the  $\text{NO}_2$  absorption region. For  $\text{NO}_2$ , this spectral region corresponds to the  $2\nu_1$  band and to the  $\nu_1+\nu_3$  band and its associated hot bands. In this spectral region absorptions due the  $\nu_1+\nu_{11}$  and  $\nu_5+\nu_{11}$  bands of the  $\text{N}_2\text{O}_4$  dimer are quite observable at  $2630$  and  $2971\text{ cm}^{-1}$ , respectively. For  $\text{NO}_2$ , the comparison of the line by line model generated using the GEISA-19 and HITEMP databases are also presented. As noticed in Ref. [18] the  $2\nu_1$  (B-type) band exhibits an “anusual” intensity behaviour, with a very weak P- branch. In HITEMP, the line intensities of the  $2\nu_1$  band are underestimated by a factor of about two, and the strange intensity pattern is not accounted for.

## References

- [1] N. Jacquinet-Husson, R. Armante, N.A. Scott, A. Chédin, L. Crépeau, C. Boutammine, A. Bouhdaoui, C. Crevoisier, V. Capelle, C. Boone, N. Poulet-Crovisier, A. Barbe, D.C. Benner, V. Boudon, L.R. Brown, J. Buldyreva, A. Campargue, L.H. Coudert, V.M. Devi, M.J. Down, B.J. Drouin, A. Fayt, C. Fittschen, J.-M. Flaud, R.R. Gamache, J.J. Harrison, C. Hill C, Ø. Hodnebrog, S.M. Hu, D. Jacquemart, A. Jolly, E. Jiménez, N. Lavrentieva, A.W. Liu, L. Lodi, O.M. Lyulin, S.T. Massie, S. Mikhailenko, H.S.P. Müller, O.V. Naumenko, A. Nikitin, C.J. Nielsen, J. Orphal, V. Perevalov, A. Perrin, E. Polovtseva, A. Predoi-Cross, M. Rotger, A.A. Ruth, Y. Shanshan, K. Sung, S. Tashkun, J. Tennyson, V.G. Tyuterev, J. Vander Auwera, B. Voronin, and A. Makie. *J. Mol. Spectrosc.* **327** (2016) 31-72.
- [2] I.E. Gordon, Rothman LS, Hill C, Kochanov RV, Tana Y, Bernath PF, Birk M, Boudon V, Campargue A, Chance KV, B.J. Drouin, J.-M. Flaud, R.R. Gamache, J.T. Hodges, D. Jacquemart, V.I. Perevalov, A. Perrin, Shine KP, Smith MA, Tennyson J, Toon GC, Tran H, Tyuterev VG, A. Barbe, A.G. Császár, V.M. Devi, T. Furtenbacher, J.J. Harrison, J.M. Hartmann, A. Jolly, T.J. Johnson, T. Karman, I. Kleiner, A.A. Kyuberis, J. Loos, O.M. Lyulin, S.T. Massie, S.N. Mikhailenko, N. Moazzen-Ahmadi, H. Müller, O.V. Naumenko, A.V. Nikitin, O.L. Polyansky, M. Rey, M. Rotger, S.W. Sharpe, K. Sung E. Starikova, S.A. Tashkun, J. Vander Auwera, G. Wagner, J. Wilzewski, P. Wcisło, S. Yuh, E.J. Zak. *J. Quant. Spectrosc. Radiat. Transf.* **203**, (2017) 3-69.
- [3] R.J. Hargreaves, I. E. Gordon, L. S. Rothman, S. A. Tashkun, V. I. Perevalov, A. A. Lukashetskaya, S. N. Yurchenko, J. Tennyson, and H. S.P. Müller. *J. Quant. Spectrosc. Radiat. Transf.* **232** (2019) 35-53.
- [4] A.A. Lukashetskaya, O.M. Lyulin, A. Perrin, and V. I. Perevalov, *Atmospheric and Oceanic Optics*, **28** (2015) 216-231.
- [5] A.A. Lukashetskaya, N.N. Lavrentieva, A.C. Dudaryonok, and V.I. Perevalov. *J. Quant. Spectrosc. Radiat. Transf.* **184** (2016) 205-217 ; *ibid.* **202** (2017) 37-37.

- [6] High-resolution transmission molecular absorption database (HITRAN).  
[<https://hitran.org/>]
- [7] A.A.Lukashevskaya, O.V.Naumenko, S.Kassi, A.Campargue. *J. Quant. Spectrosc. Radiat. Transfer.* 338 (2017) 91-96.
- [8] A.A.Lukashevskaya, S.Kassi, A.Campargue, V.I.Perevalov. *J. Quant. Spectrosc. Radiat. Transfer.* 202 (2017) 302-307.
- [9] A.A.Lukashevskaya, S.Kassi, A.Campargue, V.I.Perevalov. *J. Quant. Spectrosc. Radiat. Transfer.* 200 (2017) 17-24.
- [10] A.A.Lukashevskaya, D.Mondelain, A.Campargue, V.I.Perevalov. *J. Quant. Spectrosc. Radiat. Transfer.* 219 (2018) 393-398.
- [11] O.V.Naumenko, A.A.Lukashevskaya, S.Kassi, S.Béguier, A.Campargue. *J. Quant. Spectrosc. Radiat. Transfer.* 232 (2019) 146-151.
- [12] A.Perrin, S.Kassi and A.Campargue. *J. Quant. Spectrosc. Radiat. Transf.* 111, (2010) 2246–2255.
- [13] D. Mondelain, A.Perrin, S. Kassi, A. Campargue.. *J. Quant. Spectrosc. Radiat. Transf.* 113 (2012) 1058-1065.
- [14] A.A. Lukashevskaya, O.V. Naumenko, A. Perrin, D. Mondelain, S. Kassi and A. Campargue, , *J. Quant. Spectrosc. Rad. Trans.*, **130**, (2013) 249-259.
- [15] A.A. Lukashevskaya, O.V. Naumenko, D. Mondelain, S.Kassi, *J. Quant. Spectrosc. Radiat. Transf.* 177 (2016) 225-233.
- [16] A. Perrin, G. Toon, J. Orphal. *J. Quant. Spectrosc. Radiat. Transfer.* 154 (2015) 91–97

- [17] F. Gueye, F. Kwabia Tchana, X. Landsheere, and A. Perrin. *J. Quant. Spectrosc. Radiat. Transf.* 138 (2014) 60-69.
- [18] A. Perrin, L. Manceron, and F. Kwabia Tchana (2020). *Mol. Phys.* 118 (2020) DOI: 10.1080/00268976.2019.1711235.
- [19] A. Perrin, J.-M. Flaud, C. Camy-Peyret, B. Carli, and M. Carlotti. *Mol. Phys.* 63 (1988) 791-910.
- [20] A. Perrin, C. Camy-Peyret, J.-M. Flaud, and J. Kauppinen. *J. Mol. Spectrosc.* 130 (1988) 168-182.
- [21] J.K.G. Watson, in: J.R. Durig (Ed.), *Vibrational Spectra and Structure*, vol. 6, Elsevier, Amsterdam, 6 (1997) 1–89.
- [22] I. C. Bowater, J. M. Brown, and A. Carrington, F.R.S. *Proc. B. Soc. Lond. A.* 333 (1973) 265-288.
- [23] A. Perrin, J.-M. Flaud, C. Camy-Peyret, A.-M. Vasserot, G. Guelachvili, A. Goldman, F.J. Murcray, and R.D. Blatherwick. *J. Mol. Spectrosc.* 154 (1992) 391-406.
- [24] A. Perrin, M. Ndao, and L. Manceron. *J. Quant. Spectrosc. Radiat. Transf.* 200 (2017) 12-16.
- [25] De Bievre P, Holden NE, Barnes IL. *J. Phys. Chem. Ref. Data* 13 (1984) 809–891.
- [26] J.-Y. Mandin, V. Dana, A. Perrin, J.-M. Flaud, C. Camy-Peyret, L. Régalia, and A. Barbe, *J. Mol. Spectrosc.* 181 (1997) 379-388.
- [27] D. C. Benner, T.A. Blake, L.R. Brown, V. Malathy Devi, M.A.H. Smith, and R.A. Toth. *J. Mol. Spectrosc.* 228 (2004) 593-619.
- [28] J.M.Hartmann (private communication).

- [29] A. Perrin, J.-M. Flaud, C. Camy-Peyret, D. Hurtmans, and M. Herman and G. Guelachvili. *J. Mol. Spectrosc.* 168 (1994) 54-66.
- [30] S. Reymond-Laruinaz, M. Faye, V. Boudon, D. Doizi, and L. Manceron, *J. Mol. Spectrosc.* 336 (2017) 29-35.
- [31] M. Faye, M. Bordessoule, B. Kanouté, J.-B. Brubach, P. Roy, L. Manceron, *Rev. Sci. Inst.* 87, 063119 (2016).
- [32] K. Schofield, Evaluated chemical kinetic rate constants for various gas phase reactions. *J. Phys. Chem. Ref. Data.* 2 (1973) 25-84.
- [33] A. Perrin, J.-M. Flaud, C. Camy-Peyret, A. Goldman, J.F. Murcray, R.D. Blatherwick, and C.P. Rinsland. *J. Mol. Spectrosc.* 160, 456–463 (1993).
- [34] A. Perrin, J.-M. Flaud, C. Camy-Peyret, D. Hurtmans, D. Hurtmans, and M. Herman. *J. Mol. Spectrosc.* 177 (1996) 58-65.
- [35] T.M. Stephen, A. Goldman, A. Perrin, J.-M. Flaud, F. Keller, and C.P. Rinsland. *J. Mol. Spectrosc.* 201 (2000) 134-142.
- [36] S. Sharpe, T. Johnson, R. Sams, P. Chu, G. Rhoderick, P. Johnson, *Appl. Spectrosc.* 58 (2004) 1452– 1461.
- [37] N. Semmoud-Monnanteuil, J.-M. Colmont, A. Perrin, J.-M. Flaud, and C. Camy-Peyret, *J. Mol. Spectrosc.* 134 (1989) 176 -182.
- [38] C. Camy-Peyret, J.-M. Flaud, A. Perrin, and K. Narahari-Rao, *J. Mol. Spectrosc.* 95 (1982) 72-79.
- [39] A. Perrin, J.-M. Flaud, A. Goldman, C. Camy-Peyret, W.J. Lafferty, Ph. Arcas, and C.P. Rinsland, *J. Quant. Spectrosc. Radiat. Transfer.* 60 (1998) 839- 850.
- [40] A. Perrin, J.-M. Flaud, C. Camy-Peyret, *Infrared Phys.* 22 (1982) 343-348.



[41] A.Delon, and R.Jost, J. Chem. Phys. 95 (1991) 5686-5700.  
<https://doi.org/10.1063/1.461617>

[42] A.Perrin, J.-M.Flaud, C.Camy-Peyret, A. N’Gom, R. M’Biaké, H. Gbaguidi, and G. Guelachvili, J. Mol. Spectrosc. 171 (1995) 354-357.

[43] S. Miljanic, A. Perrin, and J. Orphal, J. Mol. Spectrosc. 242 (2007) 176–181.

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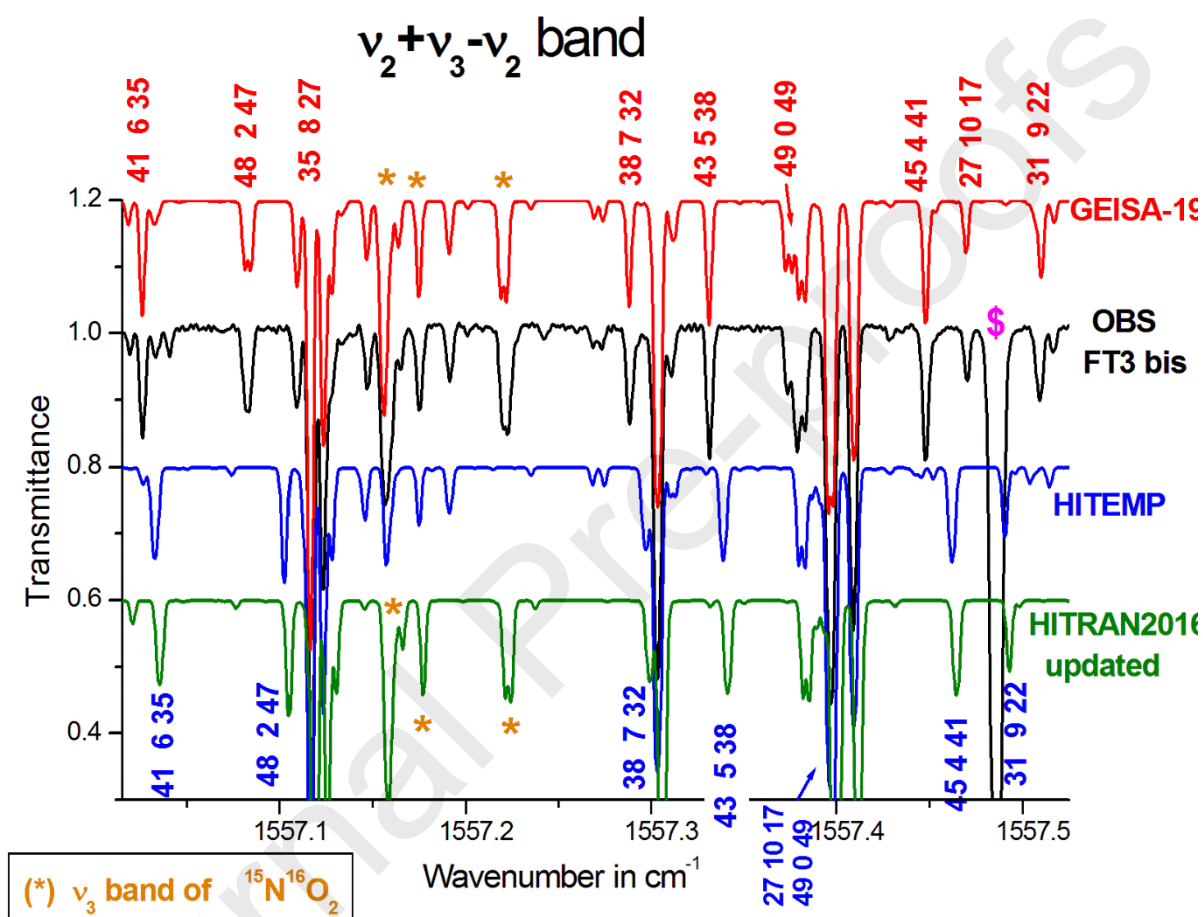


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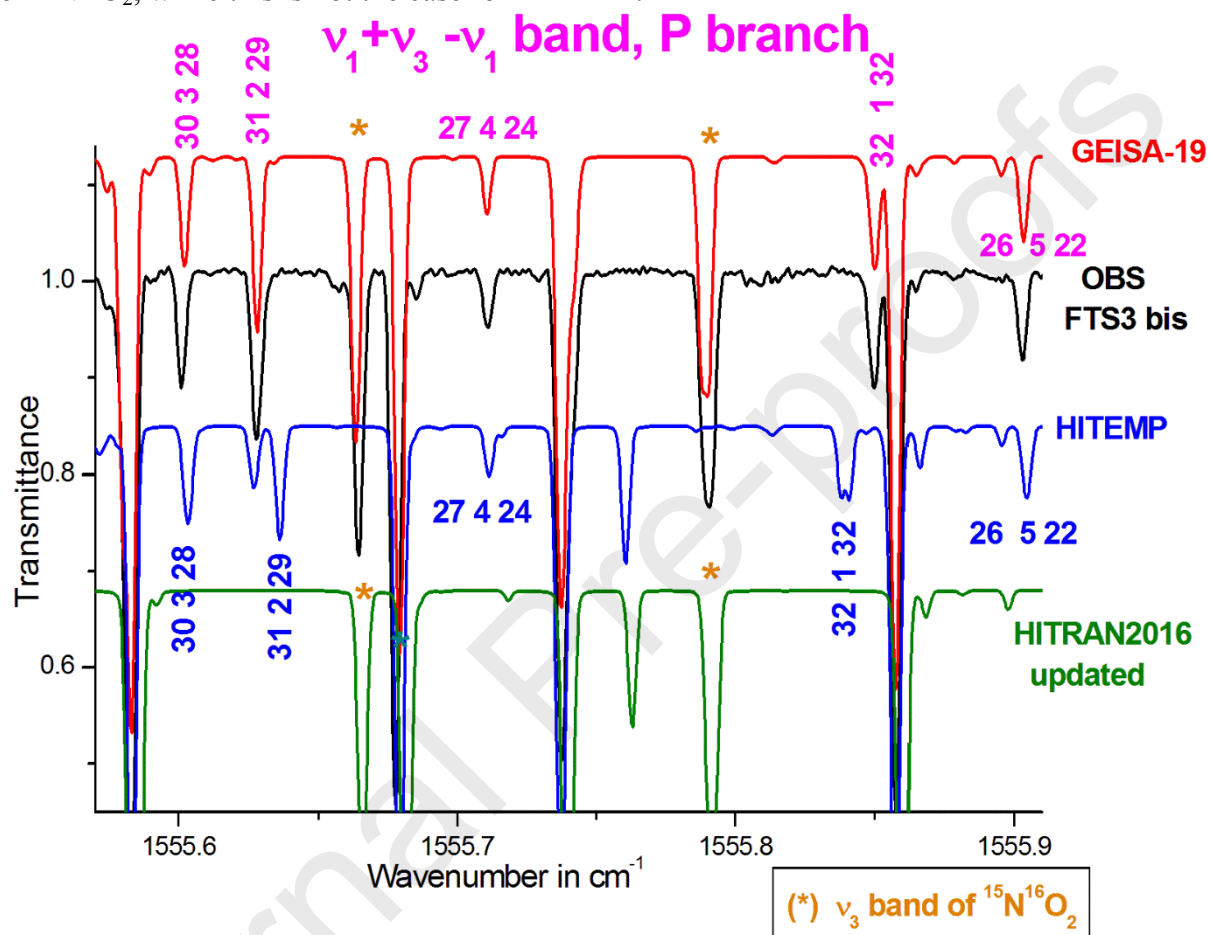


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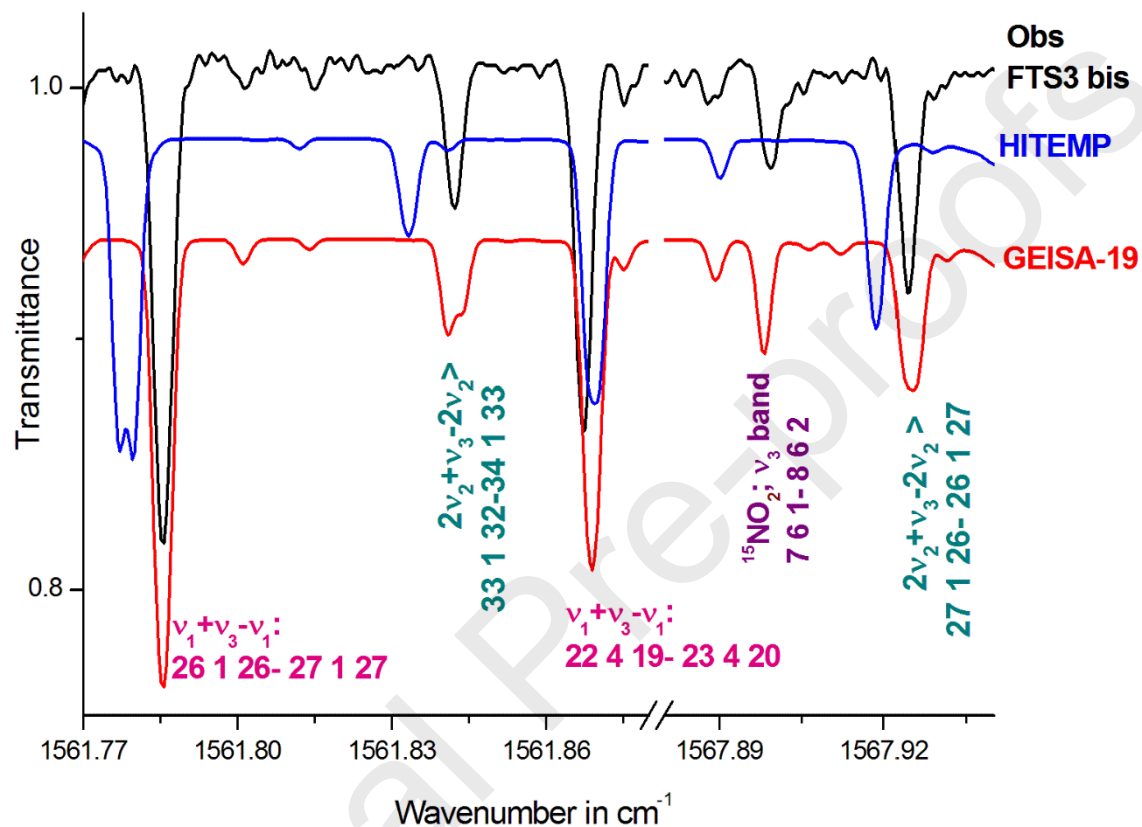


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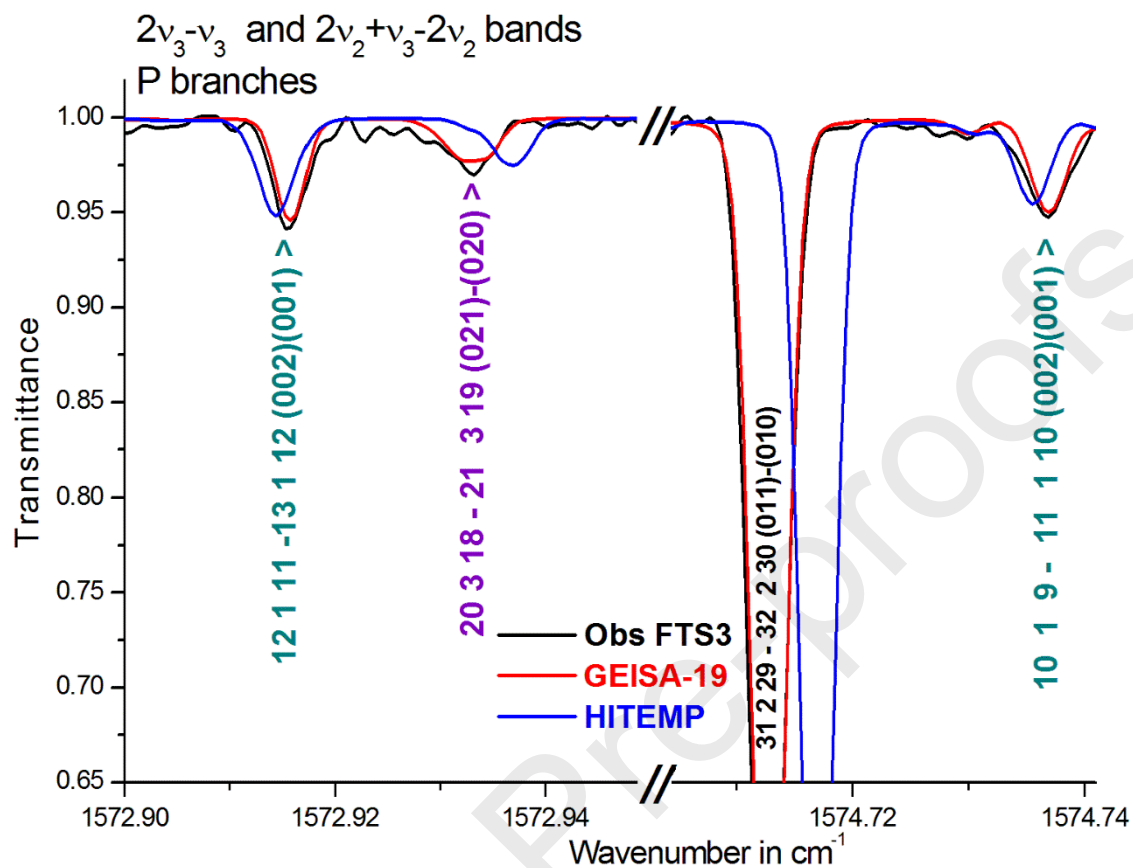


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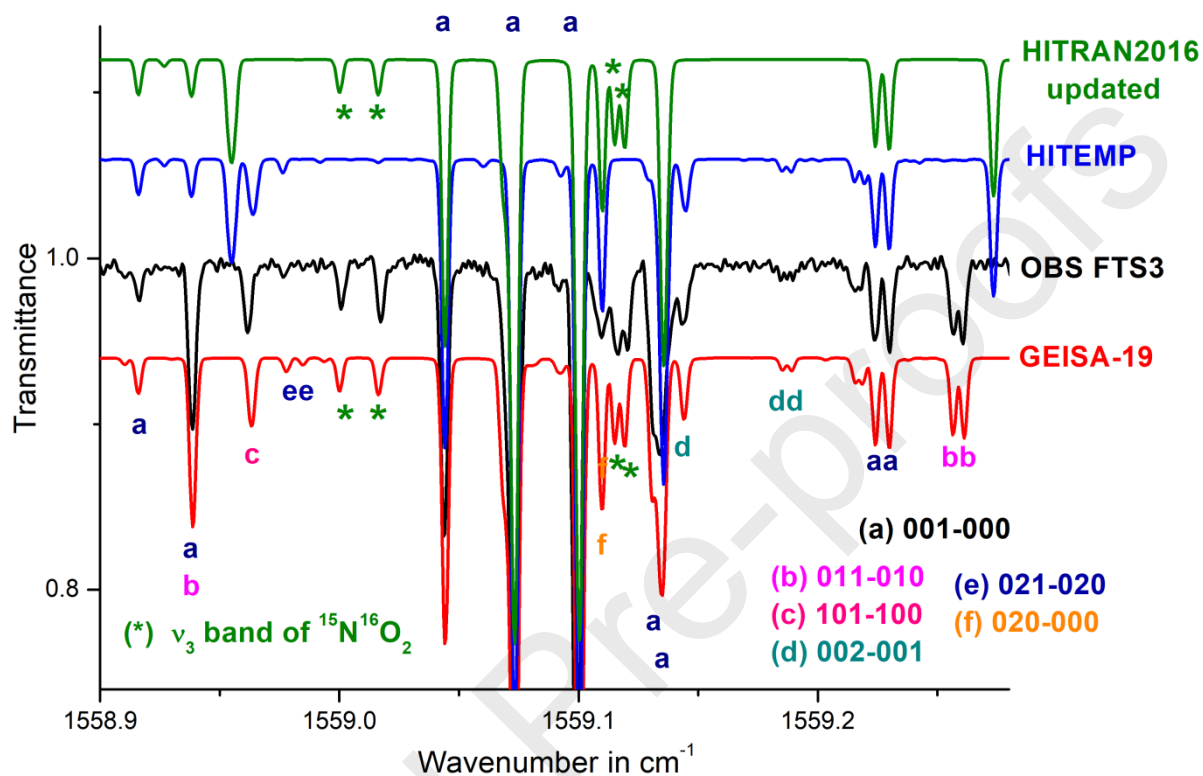


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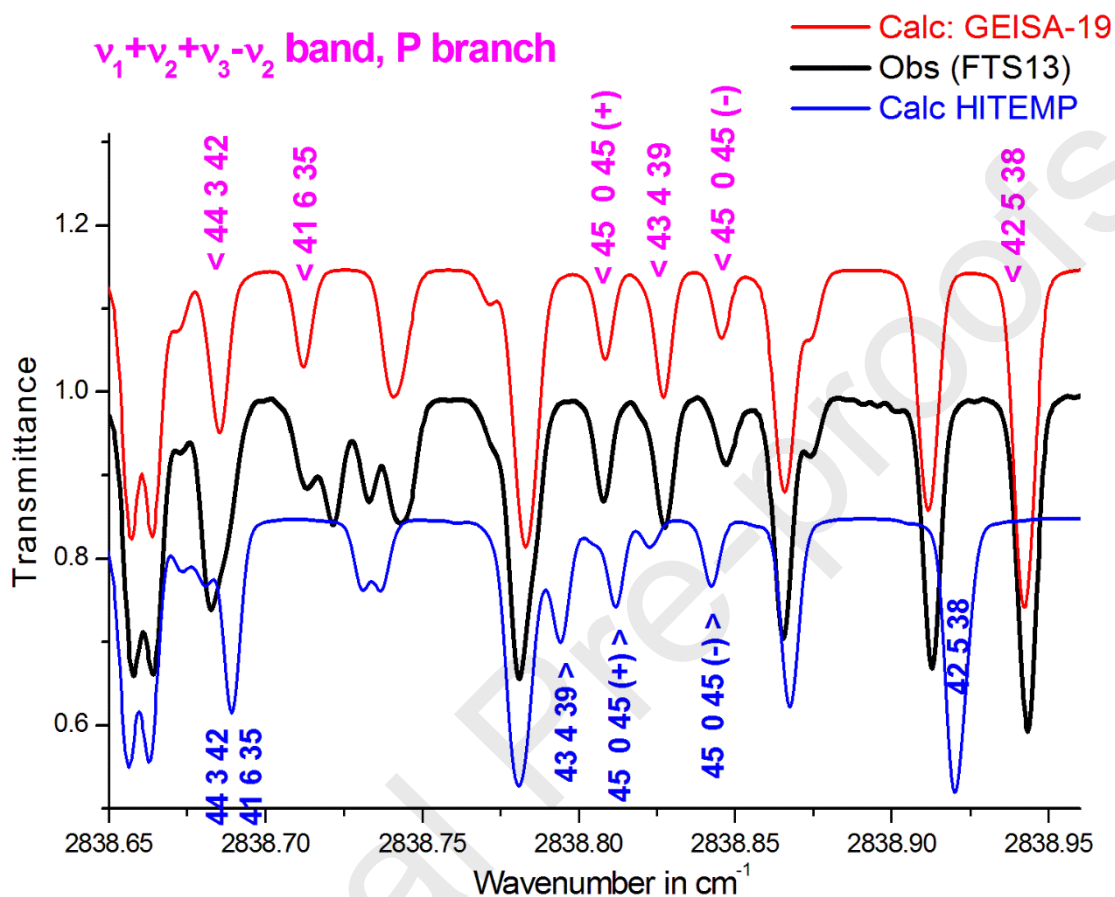


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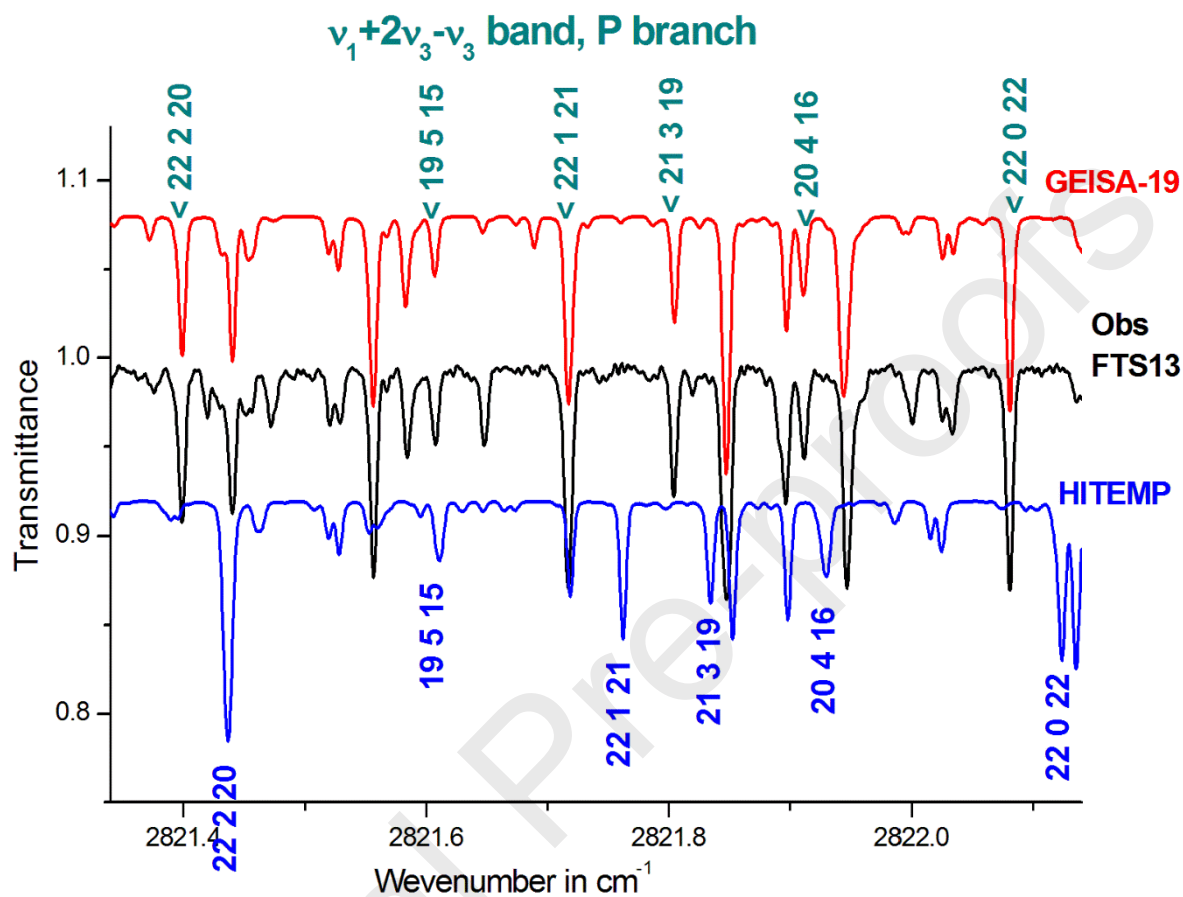




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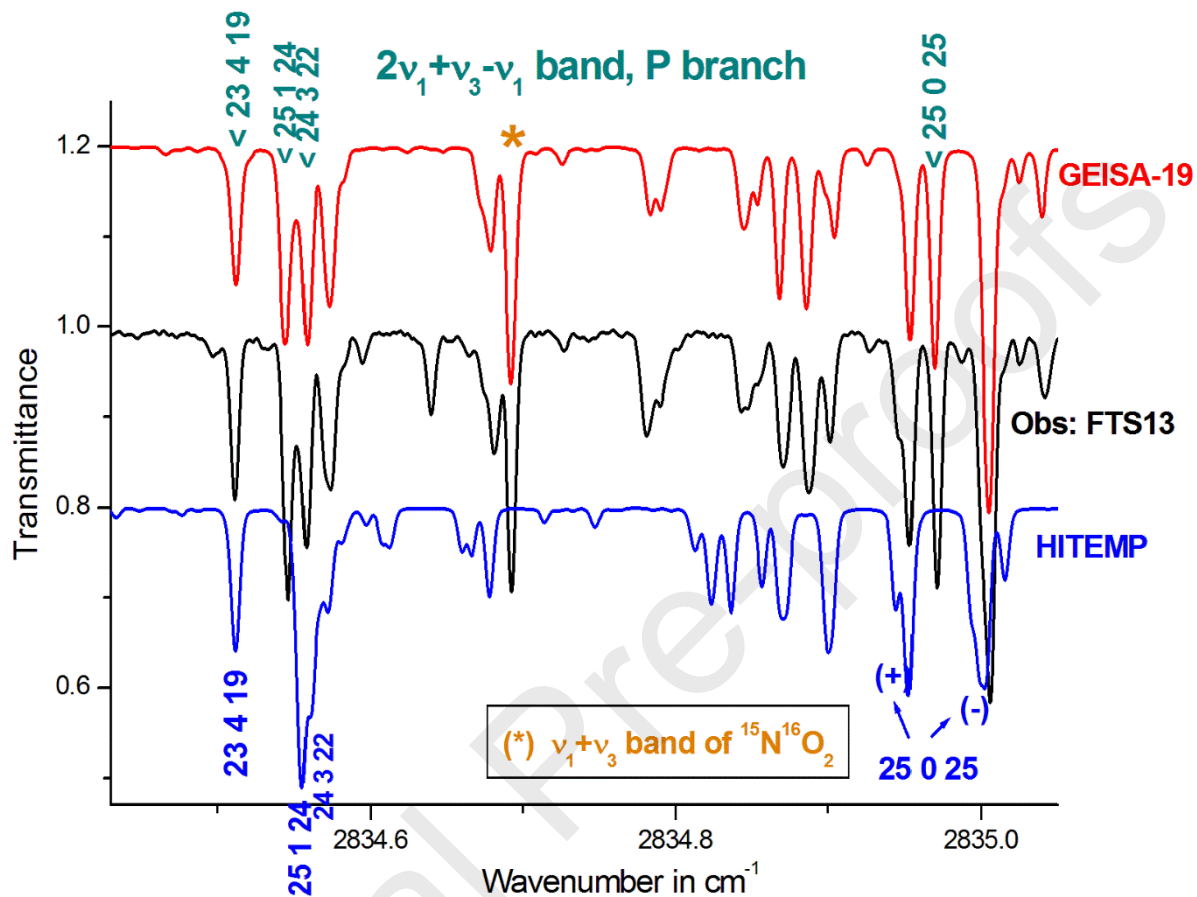


Fig 9: Portion of the P-branch of the  $\nu_1+\nu_3$  band in the  $2820\text{ cm}^{-1}$  spectral range (FTS13 spectrum). In this region, lines from the  $\nu_1+\nu_3$  band of  $^{15}\text{N}^{16}\text{O}_2$  are also observable, together with numerous lines from the hot bands from the main species. The model of the observed spectrum performed by GEISA-19, HITRAN2016-updated and HITEMP are compared. In this region, HITRAN2016-updated includes only lines from the  $\nu_1+\nu_3$  band of  $^{14}\text{N}^{16}\text{O}_2$ .

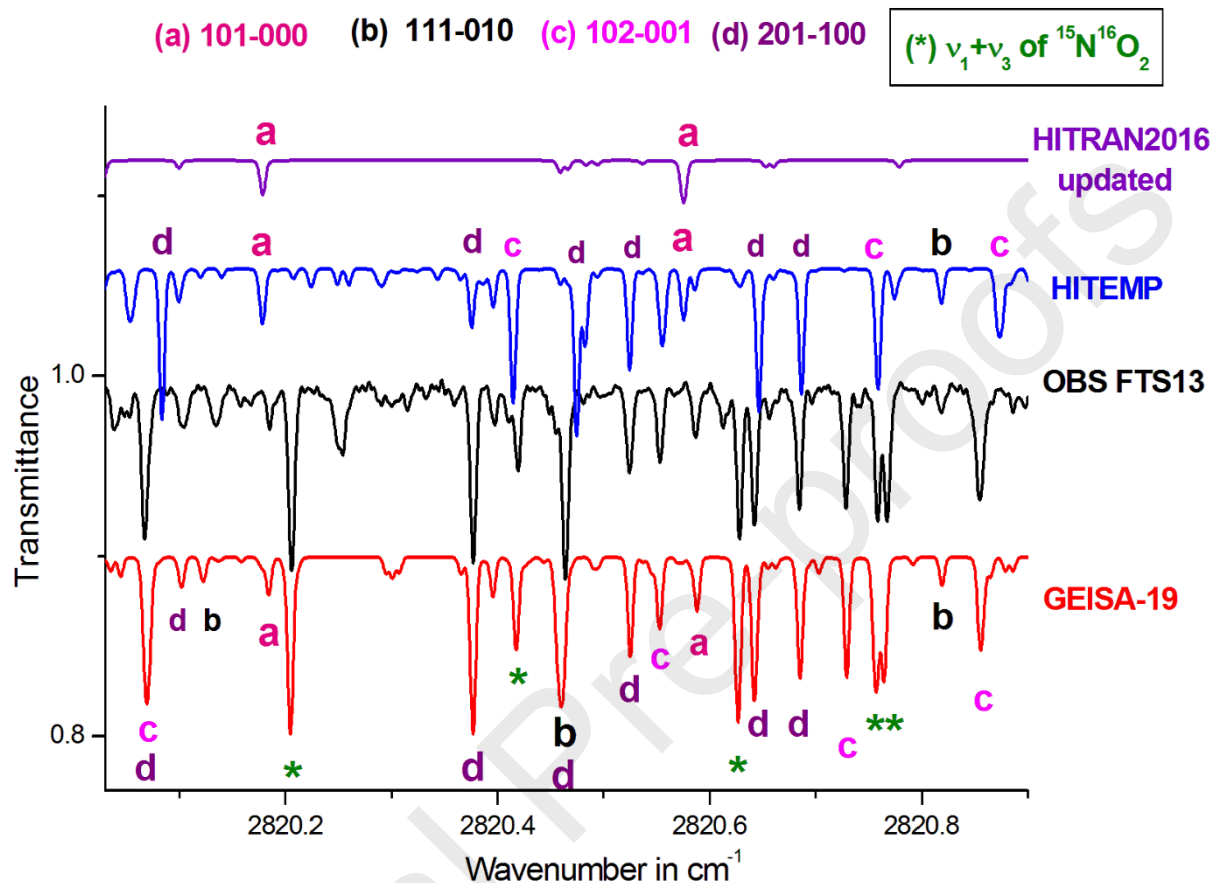


Fig. 10 : Portion of the P branch of  $\nu_1+\nu_3$  band and of the R branch of the the  $2\nu_1$  band in the 2814.4  $\text{cm}^{-1}$  regions (FTS13). The upper and lower traces compare the observed spectrum to line by line models performed using the HITEMP and GEISA-19 linelists, respectively. On the “GEISA-19” trace, the assignments for two (split) doublets of the  $2\nu_1$  band are also given, and compared to their corresponding (unsplit doublets) predictions in HITEMP. Lines from the P branches of the  $2\nu_1+\nu_3-\nu_1$  and the  $\nu_1+2\nu_3-\nu_3$  hot bands of  $^{14}\text{N}^{16}\text{O}_2$  and of the  $\nu_1+\nu_3$  band of the  $^{15}\text{N}^{16}\text{O}_2$  isotopic species species are also indicated. More detailed assignments are provided in Table TA1 of “Supp-5” of the Supplementary data.

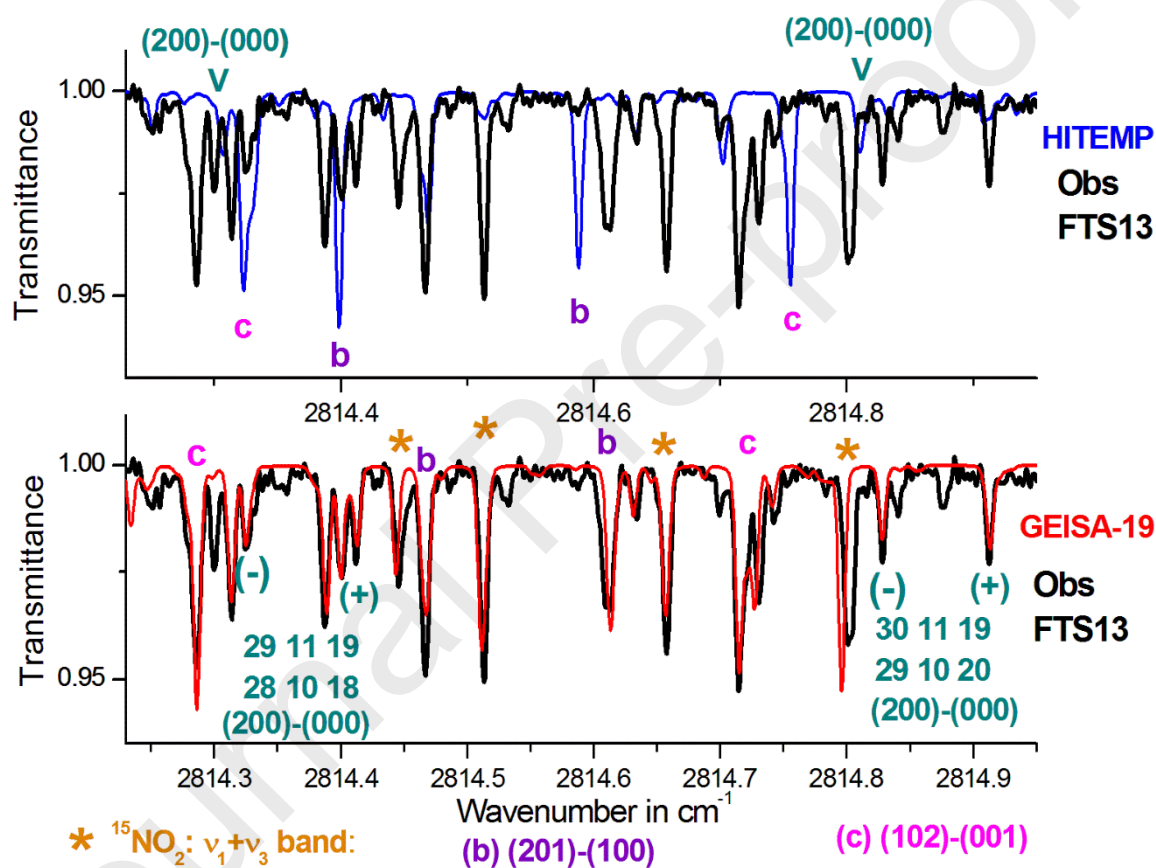
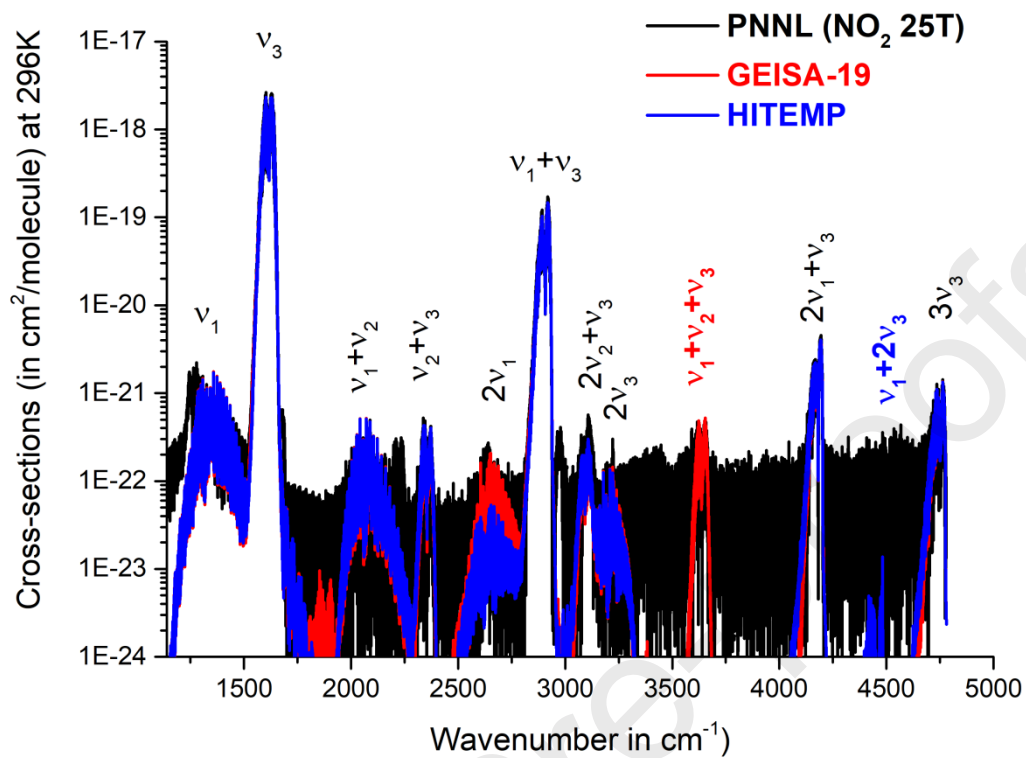


Fig 11: The PNNL cross sections for NO<sub>2</sub> (NO2\_25T). Comparison with the computed cross sections using the GEISA-19 and HITEMP databases.



**Table 1**  
**Hamiltonian matrix**

	$(v'_1, v'_2, v'_3)$	$(v_1, v_2, v_3)$
$(v'_1, v'_2, v'_3)$	W+SR	C
$(v_1, v_2, v_3)$	C	W+SR

With  $(v'_1, v'_2, v'_3) = (0, 3, 0)$  and  $(v_1, v_2, v_3) = (0, 1, 1)$  ;  
or  $(v'_1, v'_2, v'_3) = (1, 0, 2)$  and  $(v_1, v_2, v_3) = (1, 2, 1)$ .

**Example of a “W+SR” Watson “W” and spin-rotational “SR” v-diagonal block:**

		$(v_1, v_2, v_3)$	
		$N = J-1/2$	$N = J+1/2$
$(v_1, v_2, v_3)$	$N = J-1/2$	$W_{vv} + SR_{vv}$	$SR_{vv}$
	$N = J+1/2$	$SR_{vv}$	$W_{vv} + SR_{vv}$

**Example of a Coriolis “C” vibrational off diagonal block:**

		$(v_1, v_2, v_3)$	
		$N = J-1/2$	$N = J+1/2$
$(v'_1, v'_2, v'_3)$	$N = J-1/2$	$C_{v'v}$	
	$N = J+1/2$		$C_{v'v}$

$W_{vv}$ : Watson’s A-type I<sup>r</sup> representation Hamiltonian.

$$\begin{aligned}
 W_{vv} = & E_v + \left[ A^v - \frac{1}{2}(B^v + C^v) \right] N_z^2 + \frac{1}{2}(B^v + C^v) N^2 + \frac{1}{2}(B^v - C^v) N_{xy}^2 \\
 & - \Delta_K^v N_z^4 - \Delta_{JK}^v N_z^2 N^2 - \Delta_J^v (N^2)^2 - \delta_K^v \{ N_z^2, N_{xy}^2 \} - 2\delta_J^v N_{xy}^2 N^2 \\
 & + H_K^v N_z^6 + H_{KJ}^v N_z^4 N^2 + H_{JK}^v N_z^2 (N^2)^2 + H_J^v (N^2)^3 \\
 & + h_K^v \{ N_z^4, N_{xy}^2 \} + h_{KJ}^v \{ N_z^2, N_{xy}^2 \} N^2 + 2h_J^v N_{xy}^2 (N^2)^2 + \dots
 \end{aligned}$$

$SR_{vv}$  : Electron-spin rotation interactions.

$$SR_{vv} = \varepsilon_{aa}^v S_a N_a + \varepsilon_{bb}^v S_b N_b + \varepsilon_{cc}^v S_c N_c +$$

$$\begin{aligned} & \Delta_N^{vS} N^2 (\mathbf{N} \cdot \mathbf{S}) + \frac{1}{2} \Delta_{NK}^{vS} \left\{ N^2 N_z S_z + S_z N_z N^2 \right\} + \Delta_{KN}^{vS} (\mathbf{N} \cdot \mathbf{S}) \\ & + \Delta_K^{vS} N_z^3 S_z + \delta_N^{vS} (N_+^2 + N_-^2) (\mathbf{N} \cdot \mathbf{S}) + \frac{1}{2} \delta_K^{vS} \left\{ (N_+^2 + N_-^2) N_z S_z + N_z S_z (N_+^2 + N_-^2) \right\} \end{aligned}$$

$C_{v'v}$  : Coriolis-type interactions (for  $|\Delta v_3| = \text{odd}$ ):

$$C_{v'v}^C = h_{v'v}^{1C} i N_y + h_{v'v}^{2C} (N_-^3 - N_+^3) + \dots$$

with:  $N_{xy}^2 = N_x^2 - N_y^2$ ,  $N_{\pm} = N_x \mp i N_y$ , and  $\{X, Y\} = XY + YX$

cc: complex conjugate

**Table 2**

FTS spectrum	Laboratory	Main band	Spectral range in $\text{cm}^{-1}$	$\text{NO}_2$ pressure in hPa	Path length in m.	Ref.
FTS12	AILES on SOLEIL	$\nu_1+\nu_2$	1958-2289	$\approx 4.9$	10.88	This work
FTS3	AILES on SOLEIL	$\nu_3$	1540-1665	0.2	2.72	This work
FTS3bis	AILES on SOLEIL	$\nu_3$	1540-1665	0.62	2.72	This work
FTS13	AILES on SOLEIL	$2\nu_1$ and $\nu_1+\nu_3$	2520-3000	5.1	10.88	Ref. [18]
FTS23	ULB-Brussels	$\nu_2+\nu_3$	2270-2400	2.90	48.8	Ref. [29]
FTS33	ULB-Brussels	$2\nu_2+\nu_3$ , $4\nu_2$ , and $2\nu_3$	3000-3400	2.93	40.8	Ref. [34]
FTS123	University Paris Est	$\nu_1+\nu_2+\nu_3$	3578-3680	$\sim 3$	41.6	Ref. [17]
FTS113	University of Denver	$2\nu_1+\nu_3$	4125-4210	???	7.2	Ref.[35]
FTS333	University of Denver	$3\nu_3$	4680-4780	???	7.2	Ref.[35]

All measurements were done at room temperature ( $\sim 296\text{K}$ ).

\$ In Ref. [35] the partial pressure of the  $\text{NO}_2$  monomer was estimated at 10 hPa. This value may be incorrect (see text).

**Table 3**  
Vibrational energies, spin-rotational and rotational constants:

	(000) <sup>#</sup>	(0,3,0)	(0,1,1)
$E_v$		2246.04 <sup>a</sup>	2355.14980(8)
A	8.00235469	9.2178873(640)	8.1261153(290)
B	0.433706798	0.433080(130)	0.430878518(290)
C	0.410442540	0.407802(140)	0.406866093(170)
$\Delta_K$	$0.26878757 \times 10^{-2}$	#	$0.3216229(600) \times 10^{-2}$
$\Delta_{KN}$	$-0.196822 \times 10^{-4}$	#	$-0.2457511(860) \times 10^{-4}$
$\Delta_N$	$0.2992447 \times 10^{-6}$	#	$0.306505(110) \times 10^{-6}$
$\delta_K$	$0.40547 \times 10^{-5}$	#	$0.57572(420) \times 10^{-5}$
$\delta_N$	$0.3192774 \times 10^{-7}$	#	$0.335862(400) \times 10^{-7}$
$H_K$	$0.303157 \times 10^{-5}$	#	$0.428664(940) \times 10^{-5}$
$H_{KN}$	$-0.270439 \times 10^{-7}$	#	$-0.24547(270) \times 10^{-7}$
$H_N$	$0.2866 \times 10^{-12}$	#	$0.5250(220) \times 10^{-12}$
$h_K$	$0.29297 \times 10^{-7}$	#	$0.46349(580) \times 10^{-7}$
$L_K$	$-0.51104 \times 10^{-8}$	#	$-0.72675(450) \times 10^{-8}$
$L_{KKN}$	$0.35117 \times 10^{-10}$	#	$1.8141(240) \times 10^{-10}$
$\epsilon_{aa}^v$	0.180353006	#	0.1886686(710)
$\epsilon_{bb}^v$	$0.257833 \times 10^{-3}$	#	$0.25761(510) \times 10^{-3}$
$\epsilon_{cc}^v$	$-0.3178107 \times 10^{-2}$	#	$-0.32570(490) \times 10^{-2}$
$\Delta_K^{vS}$	$-0.17606 \times 10^{-3}$	#	$-0.20700(130) \times 10^{-3}$



Vibrational energies, spin-rotational and rotational constants:

	(1,0,2)	(1,2,1)
$E_v$	4460.8648(2)	4369.10 <sup>a</sup>
A	7.62253802(380)	8.1261153(290)
B	0.42593600(170)	0.430878518(290)
C	0.4020917(150)	0.406866093(170)
$\Delta_K$	$0.252081(180)\times 10^{-2}$	#
$\Delta_{KN}$	$-0.251017(260)\times 10^{-4}$	#
$\Delta_N$	$0.309119(520)\times 10^{-6}$	#
$\delta_K$	$0.3879(430)\times 10^{-5}$	#
$\delta_N$	$0.3245(460)\times 10^{-7}$	#
$H_K$	$0.29843(210)\times 10^{-5}$	#
$H_{KN}$	$-0.2042(340)\times 10^{-8}$	#
$H_N$	#	#
$h_K$	#	#
$L_K$	#	#
$L_{KKN}$	#	#
$\epsilon_{aa}^v$	0.165066(270)	#
$\epsilon_{bb}^v$	$0.2528(240)\times 10^{-3}$	#
$\epsilon_{cc}^v$	$-0.32199(230)\times 10^{-2}$	#
$\Delta_K^{vS}$	#	#

Higher order spin-rotational and rotational constants for all states :

	Rotational constants <sup>#</sup>	Spin-rotational constants <sup>#</sup>	
$H_{NK}$	$0.2995\times 10^{-10}$	$\Delta_{KN}^{vS} + \Delta_{NK}^{vS}$	$0.6005\times 10^{-6}$
$h_{KN}$	$-0.3637\times 10^{-10}$	$\Delta_N^{vS}$	$0.6322\times 10^{-9}$
$h_N$	$0.1057\times 10^{-12}$	$\Delta_{NK}^{vS}$	$0.1678\times 10^{-5}$

$L_{KN}$	$0.12158 \times 10^{-12}$	$\delta_K^{vS}$	$0.3769 \times 10^{-6}$
$P_K$	$0.867 \times 10^{-11}$	$\delta_N^{vS}$	$0.244 \times 10^{-9}$
$Q_K$	$-0.839 \times 10^{-14}$	$H_K^S$	$0.29673 \times 10^{-6}$
		$L_K^S$	$-0.3569 \times 10^{-9}$

---

***Coriolis resonances***

---

Operator	<b>(0,3,0)-(0,1,1)</b>	<b>(2,1,0)-(2,0,1)</b>
$i N_y$	0.046100(350)	0.0517441(430)

Note: The results are in  $\text{cm}^{-1}$  and the quoted errors correspond to one standard deviation. The spin-rotational constants, the higher order rotational constants together with constants marked with # were held fixed to the ground state values [37]. <sup>a</sup> From Ref. [41].

Table 4  
Description of the GEISA-19 database (1153- 4775.4 cm<sup>-1</sup>)  
The <sup>14</sup>N<sup>16</sup>O<sub>2</sub> linelist:

V'	V''	Nb	Centre	S_min	S_max	Band Int	Int <sub>Min</sub>	Int <sub>Max</sub>	Calculation		Observation	
									N Max	K <sub>a</sub> range	N Max	K <sub>a</sub> range
101	001	966	1289.89	1224.94	1607.29	0.185E- 22	0.10E- 25	0.59E- 25	52	0-8	75	0-14
100	000	8145	1319.79	1153.01	1664.24	0.721E- 19	0.10E- 24	0.17E- 21	79	0-16	56	0-15
101	020	2	1408.40	1508.14	1615.42	0.204E- 25	0.10E- 25	0.10E- 25	65	8	60	1-6
220	200	5	1473.24	1530.94	1582.85	0.285E- 24	0.14E- 25	0.11E- 24	44	2-3	0	0
021	001	239	1475.63	1533.53	1621.31	0.450E- 21	0.11E- 24	0.49E- 23	53	4-8	49	0-6
130	110	78	1483.98	1533.70	1605.43	0.132E- 23	0.10E- 25	0.73E- 25	43	3-5	0	0
120	100	520	1485.81	1455.08	1652.23	0.230E- 22	0.10E- 25	0.16E- 24	65	0-7	0	0
040	020	256	1494.66	1545.26	1631.78	0.140E- 20	0.10E- 24	0.20E- 22	52	3-5	41	K <sub>a</sub> =5
030	010	1927	1496.39	1429.25	1690.68	0.761E- 21	0.20E- 25	0.16E- 22	71	0-8	0	0
020	000	4426	1498.34	1369.96	1800.30	0.645E- 19	0.10E- 24	0.61E- 21	82	0-14	62	0-6
201	200	1659	1552.60	1503.90	1586.05	0.150E- 21	0.10E- 25	0.36E- 24	51	0-10	55	0-11
003	002	850	1552.77	1516.38	1581.14	0.244E- 22	0.10E- 25	0.66E- 25	40	0-8	60	0-8
102	101	1350	1554.12	1510.73	1586.95	0.765E- 22	0.10E- 25	0.18E- 24	47	0-9	45	0-8
111	110	3073	1574.72	1509.97	1614.74	0.238E- 20	0.10E- 25	0.56E- 23	61	0-12	55	0-10
002	001	3147	1584.59	1517.86	1626.60	0.432E- 19	0.10E- 24	0.99E- 22	65	0-10	60	0-8

101	100	4383	1586.95	1504.16	1667.57	0.909E-19	0.10E-25	0.21E-21	65	0-10	56	0-10
021	020	3032	1594.14	1526.69	1635.95	0.351E-19	0.10E-24	0.88E-22	65	0-10	49	0-6
011	010	6947	1605.50	1505.60	1752.36	0.145E-17	0.20E-25	0.34E-20	71	0-13	63	0-10
001	000	10856	1616.85	1429.63	1835.38	0.569E-16	0.10E-24	0.13E-18	90	0-18	81	0-15
002	020	243	1703.10	1525.53	1616.92	0.314E-21	0.10E-24	0.36E-23	55	4-5	55	4-5
002	100	880	1881.65	1754.08	1953.09	0.186E-21	0.10E-24	0.72E-24	53	4-5	49	0-6
110	000	7593	2063.12	1921.88	2346.16	0.165E-19	0.10E-25	0.41E-22	65	0-14	61	0-10
011	000	2551	2355.15	2289.98	2395.03	0.850E-20	0.50E-25	0.23E-22	62	0-10	63	0-10
200	000	5982	2627.34	2457.95	2921.80	0.839E-20	0.50E-25	0.23E-22	71	0-16	62	0-12
130	010	293	2797.45	2826.76	2924.89	0.671E-22	0.51E-25	0.23E-23	47	3-5	0	0
120	000	1531	2805.60	2740.02	3023.06	0.853E-21	0.50E-25	0.53E-23	77	0-9	61	3-5
102	001	1836	2844.01	2787.39	2872.69	0.217E-20	0.10E-24	0.51E-23	53	0-10	45	0-8
201	100	3245	2860.14	2790.47	2888.91	0.923E-20	0.50E-26	0.22E-22	52	0-10	55	0-11
111	010	3904	2888.19	2811.08	2919.85	0.734E-19	0.51E-25	0.17E-21	61	0-12	55	0-12
101	000	8731	2906.74	2688.75	3148.79	0.287E-17	0.50E-25	0.66E-20	80	0-16	75	0-14
040	000	194	2993.00	3040.81	3133.67	0.257E-21	0.11E-24	0.37E-23	49	3-5	41	K <sub>a</sub> =5
021	000	2512	3092.48	3032.58	3283.84	0.650E-20	0.10E-24	0.15E-22	58	0-12	49	6
002	000	3597	3201.44	3019.48	3322.09	0.377E-20	0.10E-24	0.14E-22	65	0-10	60	0-8
210	000	1018	3364.57	3291.70	3658.07	0.215E-22	0.10E-25	0.15E-24	52	0-9	0	0
130	000	275	3547.10	3581.21	3678.91	0.110E-22	0.10E-25	0.36E-24	46	3-5	0	0
111	000	3999	3637.84	3563.65	3690.60	0.123E-19	0.10E-25	0.28E-22	61	0-15	55	0-12
220	000	333	4100.58	4100.56	4207.21	0.191E-21	0.50E-25	0.58E-22	63	2-7	N=23	3
201	000	4234	4179.94	4087.83	4230.50	0.790E-19	0.50E-25	0.18E-21	66	0-13	55	0-11
022	000	133	4656.34	4702.20	4768.40	0.145E-22	0.50E-25	0.18E-24	41	5-6	0	0
003	000	3647	4754.21	4624.75	4775.31	0.304E-19	0.50E-25	0.70E-22	66	0-12	57	0-10

The <sup>15</sup>N<sup>16</sup>O<sub>2</sub> linelist

									Calculation		Observation	
V <sub>U</sub>	V <sub>L</sub>	Nb	Band center	S <sub>Min</sub>	S <sub>Max</sub>	Band Int	Int <sub>Min</sub>	Int <sub>Max</sub>	N <sub>Max</sub>	K <sub>a</sub> range	N <sub>Max</sub>	K <sub>a</sub> range
001	000	5860	1582.10	1500.73	1660.45	2.039E-19	2.00E-26	4.56E-22	77	0-15	70	0-13
120	000	157	2761.21	2819.67	2889.29	0.162E-22	0.21E-25	0.21E-24	42	5-5	0	0
101	000	3640	2858.71	2784.88	2890.81	1.036E-20	3.64E-28	2.34E-23	61	0-15	54	0-10

Total : 118249 lines

Caption :

V<sub>U</sub> V<sub>L</sub> : (v<sub>1</sub>,v<sub>2</sub>,v<sub>3</sub>) vibrational assignments

Nb : number of lines

Band center: approximate band center (in  $\text{cm}^{-1}$ )

SMin Smax= Wavenumber range (in  $\text{cm}^{-1}$ )

Band int, Int<sub>Min</sub>, Int<sub>Max</sub>: band intensity, minimum and maximum value of the line strength (in  $\text{cm}^{-1}/(\text{molecule}\cdot\text{cm}^{-2})$  at 296K.

Calculation: N<sub>Max</sub>:, K<sub>a</sub> range: maximum value of N, minimum and maximum value of K<sub>a</sub> for the upper state in the calculated linelist

Observation: N<sub>Max</sub>:, K<sub>a</sub> range: maximum value of N, minimum and maximum value of K<sub>a</sub> for the upper and lower experimental energy levels used for the computation.

## Table 4

(see the separate file)

Description of the GEISA-19database

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Table 5

N K <sub>a</sub> K <sub>c</sub> J - N K <sub>a</sub> K <sub>c</sub> J	vib	Obs	§	TW			HIT		
				Sigma	E <sub>L</sub>	E <sub>U</sub>	Sigma	E <sub>L</sub>	E <sub>U</sub>
42 0 42 + 41 1 41 +	110 000	2090.6336 <sup>(2)</sup>	G	2090.6330	722.2130	<u>2812.8460</u>	2090.6366	722.2130	<u>2812.8496</u>
42 0 42 - 41 1 41 -	110 000	2090.6336 <sup>(2)</sup>	G	2090.6345	722.3139	<u>2812.9484</u>	2090.6353	722.3139	<u>2812.9492</u>
42 0 42 + 41 1 41 +	110 010			1341.8081	1471.0379	<u>2812.8460</u>	1341.8173	1471.0399	<u>2812.8572</u>
42 0 42 - 41 1 41 -	110 010			1341.8076	1471.1408	<u>2812.9484</u>	1341.8134	1471.1389	<u>2812.9523</u>
33 1 33 + 32 0 32 +	110 000	2090.6234 <sup>(2)</sup>	G	2090.6234	443.0513	<u>2533.6747</u>	2090.6230	443.0513	<u>2533.6743</u>
33 1 33 - 32 0 32 -	110 000	2090.6408 <sup>(2)</sup>	G	2090.6410	443.1170	<u>2533.7580</u>	2090.6380	443.1170	<u>2533.7550</u>
33 1 33 + 32 0 32 +	110 010			1341.5283	1192.1464	<u>2533.6747</u>	1341.5309	1192.1609	<u>2533.6918</u>
33 1 33 - 32 0 32 -	110 010			1341.5444	1192.2136	<u>2533.7580</u>	1341.5438	1192.2109	<u>2533.7547</u>
35 8 27 + 36 8 28 +	011 010	1557.1093 <sup>(2)</sup>	G	1557.1085	1808.5312	<u>3365.6397</u>	1557.1130	1808.5329	<u>3365.6459</u>
35 8 27 - 36 8 28 -	011 010	1557.1093 <sup>(2)</sup>	G	1557.1095	1808.2940	<u>3365.4035</u>	1557.1150	1808.2999	<u>3365.4149</u>
35 8 27 + 36 8 28 +	011 000		G	2327.4023	1038.2374	<u>3365.6396</u>	2327.4101	1038.2374	<u>3365.6475</u>
35 8 27 - 36 8 28 -	011 000		G	2327.3855	1038.0178	<u>3365.4033</u>	2327.3941	1038.0178	<u>3365.4119</u>
45 4 41 + 46 4 42 +	011 010	1557.4479 <sup>(2)</sup> )	G	1557.4488	1787.5838	<u>3345.0326</u>	1557.4630	1787.5839	<u>3345.0469</u>
45 4 41 - 46 4 42 -	011 010	1557.4479 <sup>(2)</sup> )	G	1557.4476	1787.5899	<u>3345.0375</u>	1557.4610	1787.5729	<u>3345.0339</u>
45 4 41 + 46 4 42 +	011 000		G	2311.9175	1033.1151	<u>3345.0326</u>	2311.9282	1033.1151	<u>3345.0433</u>
45 4 41 - 46 4 42 -	011 000		G	2311.9143	1033.1233	<u>3345.0376</u>	2311.9252	1033.1233	<u>3345.0485</u>
27 10 17 + 28 10 18 +	011 010	1557.4704 <sup>(2)</sup> )	G	1557.4698	1858.4606	<u>3415.9304</u>	1557.3910	1858.1159	<u>3415.5029</u>
27 10 17 - 28 10 18 -	011 010	1557.4704 <sup>(2)</sup> )	G	1557.4698	1857.9318	<u>3415.4017</u>	1557.3870	1858.6289	<u>3416.0199</u>
27 10 17 + 28 10 18 +	011 000		G	2337.9880	1077.9428	<u>3415.9309</u>	2338.1092	1077.4496	<u>3415.5588</u>
27 10 17 - 28 10 18 -	011 000		G	2337.9523	1077.4496	<u>3415.4019</u>	2338.1409	1077.9428	<u>3416.0837</u>
31 9 22 + 32 9 23 +	011 010	1557.5094 <sup>(2)</sup> )	G	1557.5097	1820.1611	<u>3377.6708</u>	1557.4900	1820.2009	<u>3377.6909</u>
31 9 22 - 32 9 23 -	011 010	1557.5094 <sup>(2)</sup> )	G	1557.5107	1819.7996	<u>3377.3103</u>	1557.4910	1819.8479	<u>3377.3389</u>
31 9 22 + 32 9 23 +	011 000		G	2332.7906	1044.8802	<u>3377.6708</u>	2332.8372	1044.8802	<u>3377.7174</u>
31 9 22 - 32 9 23 -	011 000		G	2332.7661	1044.5441	<u>3377.3102</u>	2332.8146	1044.5440	<u>3377.3586</u>

21 3 19 + 20 2 18 +	020 010	809.0824 <sup>(1)</sup>	G	809.0826	958.7949	<u>1767.8775</u>	809.0826	958.7949	<u>1767.8775</u>
21 3 19- 20 2 18 -	020 010	809.0445 <sup>(1)</sup>	G	809.0438	958.7871	<u>1767.8309</u>	809.0438	958.7871	<u>1767.8309</u>
21 3 19 + 20 2 18 +	020 000		G	1560.0222	207.8553	<u>1767.8775</u>	1560.0103	207.8553	<u>1767.8656</u>
21 3 19- 20 2 18 -	020 000		G	1559.9817	207.8493	<u>1767.8310</u>	1559.9695	207.8493	<u>1767.8188</u>
20 3 18 + 21 3 19 +	021 020	1572.9332 <sup>(2)</sup> )	G	1572.9346	<u>1767.8774</u>	<u>3340.8120</u>	1572.9378	<u>1767.8793</u>	<u>3340.8171</u>
20 3 18 - 21 3 19 -	021 020	1572.9332 <sup>(2)</sup> )	G	1572.9322	<u>1767.8310</u>	<u>3340.7632</u>	1572.9359	<u>1767.8313</u>	<u>3340.7672</u>
20 3 18 + 21 3 19 +	021 010			2324.9823	<b>1015.8297</b>	<u>3340.8120</u>	2324.9839	<b>1015.8349</b>	<u>3340.8188</u>
20 3 18 - 21 3 19 -	021 010			2324.9750	<b>1015.7882</b>	<u>3340.7632</u>	2324.9771	<b>1015.7949</b>	<u>3340.7720</u>
20 3 18 + 21 3 19 +	021 000	3077.7153 <sup>(3)</sup>	G	3077.7153	263.0968	<u>3340.8121</u>	3077.7192	263.0968	<u>3340.8160</u>
20 3 18 - 21 3 19 -	021 000	3077.7036 <sup>(3)</sup>	G	3077.7033	263.0598	<u>3340.7631</u>	3077.7078	263.0598	<u>3340.7676</u>
23 0 23 + 24 0 24 +	201 200		G	1530.9994	2876.7510	<u>4407.7504</u>	1531.0200	2876.7552	<u>4407.7752</u>
23 0 23 - 24 0 24 -	201 200		G	1531.0282	2876.7998	<u>4407.8279</u>	1531.0171	2876.7993	<u>4407.8164</u>
23 0 23 + 24 0 24 +	201 000	4155.4091 <sup>(2)</sup> )	G	4155.4106	252.3398	<u>4407.7504</u>	4155.4340	252.3398	<u>4407.7738</u>
23 0 23 - 24 0 24 -	201 000	4155.4467 <sup>(2)</sup> )	G	4155.4439	252.3841	<u>4407.8280</u>	4155.4439	252.3841	<u>4407.8280</u>

Caption :

N K<sub>a</sub> K<sub>c</sub> J - N K<sub>a</sub> K<sub>c</sub> J: Rotational assignments

Vib: vibrational transition

Obs: (when possible) Observed line positions: <sup>(1)</sup> Line assigned in Ref. [29]; <sup>(2)</sup> on spectrum FTS3, FTS12, or FTS113 (see Fig. 1, Fig. 4, Fig. A3 or Fig. A7 for example). <sup>(3)</sup> Lines assigned in Ref. [34].

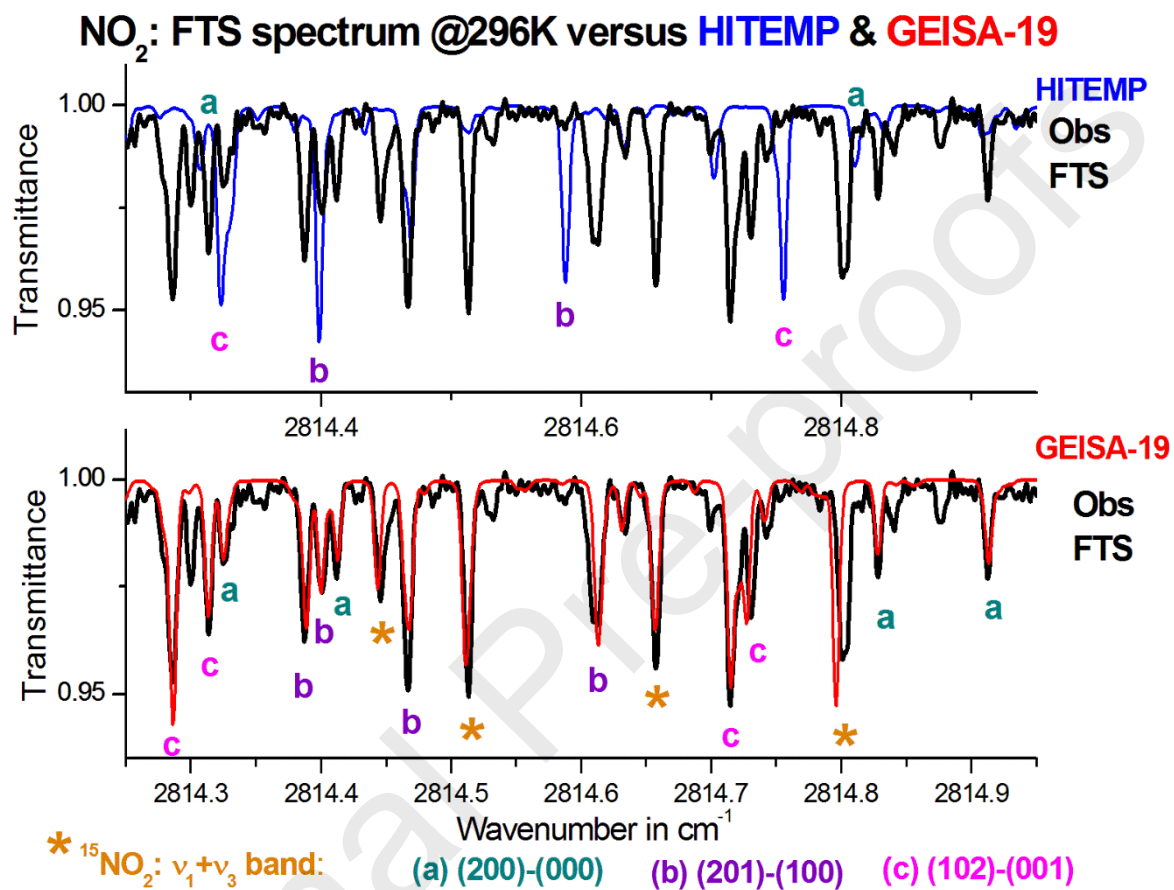
TW: This work.

\$: a "G" indicates the lines which are quoted in GEISA-19.

HIT: HITEMP database.

This work: results of the present work. G-19 if included in GEISA-19.

E<sub>U</sub> and E<sub>L</sub>: upper and lower state energy level





### Highlights

New lists of line position, line intensity and line shape parameters of nitrogen dioxide (GEISA-19) for GEISA (<https://geisa.aeris-data.fr/>).dioxide.  $^{14}\text{N}^{16}\text{O}_2$  and  $^{15}\text{N}^{16}\text{O}_2$ . 1153-4775  $\text{cm}^{-1}$ . Validation with high resolution Fourier transform spectra. Inter-comparison with HITEMP and HITRAN2016-updated. First identification of the  $\nu_1+2\nu_3-\nu_3$  hot band. First determination of the (1,0,2) energy level parameters and an improved set of parameters for the (0,1,1) vibrational state.

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