



HAL
open science

ON THE CAUCHY PROBLEM FOR A DERIVATIVE NONLINEAR SCHRÖDINGER EQUATION WITH NONVANISHING BOUNDARY CONDITIONS

Phan van Tin

► **To cite this version:**

Phan van Tin. ON THE CAUCHY PROBLEM FOR A DERIVATIVE NONLINEAR SCHRÖDINGER EQUATION WITH NONVANISHING BOUNDARY CONDITIONS. *Evolution Equations and Control Theory (EECT)*, 2021, 10.3934/eect.2021028 . hal-03083232v3

HAL Id: hal-03083232

<https://hal.science/hal-03083232v3>

Submitted on 27 Sep 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

ON THE CAUCHY PROBLEM FOR A DERIVATIVE
NONLINEAR SCHRÖDINGER EQUATION WITH
NONVANISHING BOUNDARY CONDITIONS

PHAN VAN TIN

Institut de Mathématiques de Toulouse, UMR5219
Université de Toulouse, CNRS
UPS IMT, F-31062 Toulouse Cedex 9, France

(Communicated by Remi Carles)

ABSTRACT. In this paper we consider the Schrödinger equation with nonlinear derivative term. Our goal is to initiate the study of this equation with non vanishing boundary conditions. We obtain the local well posedness for the Cauchy problem on Zhidkov spaces $X^k(\mathbb{R})$ and in $\phi + H^k(\mathbb{R})$. Moreover, we prove the existence of conservation laws by using localizing functions. Finally, we give explicit formulas for stationary solutions on Zhidkov spaces.

CONTENTS

1. Introduction	2
2. Local existence in Zhidkov spaces	6
2.1. Preliminaries on Zhidkov spaces	6
2.2. From the equation to the system	6
2.3. Resolution of the system	8
2.4. Preservation of the differential identity	9
2.5. From the system to the equation	12
3. Results on the space $\phi + H^k(\mathbb{R})$ for $\phi \in X^k(\mathbb{R})$	13
3.1. The local well posedness on $\phi + H^2(\mathbb{R})$	13
3.2. The local well posedness on $\phi + H^1(\mathbb{R})$	17
4. Conservation of the mass, the energy and the momentum	21
4.1. Conservation of mass	21
4.2. Conservation of energy	22
4.3. Conservation of momentum	24
5. Stationary solutions	25
Acknowledgments	30
REFERENCES	30

2020 *Mathematics Subject Classification.* 35Q55, 35A01.

Key words and phrases. Nonlinear derivative Schrödinger equations, Cauchy problem, nonvanishing boundary condition.

1. **Introduction.** We are interested in the Cauchy problem for the following derivative nonlinear Schrödinger equation with nonvanishing boundary conditions:

$$\begin{cases} i\partial_t u + \partial^2 u = -iu^2 \partial \bar{u}, \\ u(0) = u_0, \end{cases} \quad (1.1)$$

where $u : \mathbb{R}_t \times \mathbb{R}_x \rightarrow \mathbb{C}$, $\partial = \partial_x$ denotes derivative in space and ∂_t denotes derivative in time.

Our attention was drawn to this equation by the work of Hayashi and Ozawa [10] concerning the more general nonlinear Schrödinger equation

$$\begin{cases} i\partial_t u + \partial^2 u = i\lambda|u|^2 \partial u + i\mu u^2 \partial \bar{u} + f(u), \\ u(0) = u_0. \end{cases} \quad (1.2)$$

When $\lambda = 0$, $\mu = -1$, $f \equiv 0$, then (1.2) reduces to (1.1). This type of equation is usually referred to as *derivative nonlinear Schrödinger equations*. It may appear in various areas of physics, e.g. in Plasma Physics for the propagation of Alfvén waves [13, 15].

Under Dirichlet boundary conditions in space, the Cauchy problem for (1.1) has been solved in [10]: local well-posedness holds in $H^1(\mathbb{R})$, i.e. for any $u_0 \in H^1(\mathbb{R})$ there exists a unique solution $u \in C(I, H^1(\mathbb{R}))$ of (1.1) on a maximal interval of time I . Moreover, we have continuous dependence with respect to the initial data, blow-up at the ends of the time interval of existence I if I is bounded and conservation of energy, mass and momentum.

The main difficulty is the appearance of the derivative term $-iu^2 \partial_x$. We cannot use the classical contraction method for this type of nonlinear Schrödinger equations. In [10] Hayashi and Ozawa use the Gauge transform to establish the equivalence of the local well-posedness between the equation (1.2) and a system of equations without derivative terms. By studying the Cauchy problem for this system, they obtain the associated results for (1.2). In [9], Hayashi and Ozawa construct a sequence of solutions of approximated equations and prove that this sequence is converging to a solution of (1.2), obtaining this way the local well-posedness of (1.2). The approximation method has also been used by Tsutsumi and Fukuda in [16, 17]. The difference between [9] and [16, 17] lies in the way of constructing the approximate equation. In [9], the authors use approximation on the non-linear term, whereas in [16, 17] the authors use approximation on the linear operator.

To our knowledge, the Cauchy problem for (1.1) has not been studied under non-zero boundary conditions, and our goal in this paper is to initiate this study. Note that non-zero boundary conditions on the whole space are much rarely considered in the literature around nonlinear dispersive equations than Dirichlet boundary conditions. In the case of the nonlinear Schrödinger equation with power-type nonlinearity, we refer to the works of Gérard [7, 8] for local well-posedness in the energy space and to the works of Gallo [5] and Zhidkov [18] for local well-posedness in Zhidkov spaces (see Section 2.1 for the definition of Zhidkov spaces) and Gallo [6] for local well-posedness in $u_0 + H^1(\mathbb{R})$. In this paper, using the method of Hayashi and Ozawa as in [10] on the Zhidkov-space $X^k(\mathbb{R})$, ($k \geq 4$) and in the space $\phi + H^k(\mathbb{R})$ ($k = 1, 2$) for ϕ in a Zhidkov space, we obtain the existence, uniqueness and continuous dependence on the initial data of solutions of (1.1) in these spaces. Using the transform

$$v = \partial u + \frac{i}{2}|u|^2 u, \quad (1.3)$$

we see that if u is a solution of (1.1) then (u, v) is a solution of a system of two equations without derivative terms. It is easy to obtain the local wellposedness of this system on Zhidkov spaces. The main difficulty is how to obtain a solution of (1.1) from a solution of the system. Actually, we must prove that the relation (1.3) is conserved in time. The main difference in our setting with the setting in [11] is that we work on Zhidkov spaces instead of the space of localized functions $H^1(\mathbb{R})$. Our first main result is the following.

Theorem 1.1. *Let $u_0 \in X^4(\mathbb{R})$. Then there exists a unique maximal solution of (1.1) $u \in C((T_{min}, T_{max}), X^4(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^2(\mathbb{R}))$. Moreover, u satisfies the two following properties.*

- Blow-up alternative. *If $T_{max} < \infty$ (resp. $T_{min} > -\infty$) then*

$$\lim_{t \rightarrow T_{max} \text{ (resp. } T_{min})} \|u(t)\|_{X^2} = \infty.$$

- Continuity with respect to the initial data. *If $u_0^n \in X^4(\mathbb{R})$ is such that $u_0^n \rightarrow u_0$ in $X^4(\mathbb{R})$ then for any subinterval $[T_1, T_2] \subset (T_{min}, T_{max})$ the associated solutions of equation (1.1) (u^n) satisfy*

$$\lim_{n \rightarrow \infty} \|u^n - u\|_{L^\infty([T_1, T_2], X^4)} = 0.$$

To obtain the local wellposedness on $\phi + H^k(\mathbb{R})$ for ϕ in Zhidkov spaces $X^l(\mathbb{R})$. First, we use the transform $v = \partial u + \frac{i}{2}|u|^2u$. We see that if $u \in \phi + H^k(\mathbb{R})$ then $v \in \frac{i}{2}|\phi|^2\phi + H^{k-1}(\mathbb{R})$. This motivates us to define $\tilde{u} = u - \phi$ and $\tilde{v} = v - \frac{i}{2}|\phi|^2\phi$. We have

$$\tilde{v} = \partial \tilde{u} + \frac{i}{2}(|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2\phi) + \partial \phi. \quad (1.4)$$

We see that if u is a solution of (1.1) then (\tilde{u}, \tilde{v}) is a solution of a system of two equations without the derivative terms. For technical reasons, we will need some regularity on ϕ . With a solution of the system in hand, we want to obtain a solution of (1.1). In practice, we need to prove that the relation (1.4) is conserved in time. Our main second result is the following.

Theorem 1.2. *Let $\phi \in X^4(\mathbb{R})$ and $u_0 \in \phi + H^2(\mathbb{R})$. Then the problem (1.1) has a unique maximal solution $u \in C((T_{min}, T_{max}), \phi + H^2(\mathbb{R}))$ which is differentiable as a function of $C((T_{min}, T_{max}), \phi + L^2(\mathbb{R}))$ and such that $u_t \in C((T_{min}, T_{max}), L^2(\mathbb{R}))$. Moreover u satisfies the following properties.*

- (1) *Blow-up alternative: If $T_{max} < \infty$ (resp. $T_{min} > -\infty$) then*

$$\lim_{t \rightarrow T_{max} \text{ (resp. } T_{min})} (\|u(t) - \phi\|_{H^2(\mathbb{R})}) = \infty.$$

- (2) *Continuous dependence on initial data: If $(u_0^n) \subset \phi + H^2(\mathbb{R})$ is such that $\|u_0^n - u_0\|_{H^2} \rightarrow 0$ as $n \rightarrow \infty$ then for all $[T_1, T_2] \subset (T_{min}, T_{max})$ the associated solutions (u^n) of (1.1) satisfy*

$$\lim_{n \rightarrow \infty} \|u^n - u\|_{L^\infty([T_1, T_2], H^2)} = 0.$$

In the less regular space $\phi + H^1(\mathbb{R})$, we obtain the local well posedness under a smallness condition on the initial data. Our third main result is the following.

Theorem 1.3. *Let $\phi \in X^4(\mathbb{R})$ such that $\|\partial \phi\|_{L^2}$ is small enough, $u_0 \in \phi + H^1(\mathbb{R})$ such that $\|u_0 - \phi\|_{H^1(\mathbb{R})}$ is small enough. There exist $T > 0$ and a unique solution u of (1.1) such that*

$$u - \phi \in C([-T, T], H^1(\mathbb{R})) \cap L^4([-T, T], W^{1, \infty}(\mathbb{R})).$$

In the proof of Theorem 1.3, the main difference with the case $\phi + H^2(\mathbb{R})$ is that we use Strichartz estimates to prove the contractivity of a map on $L^\infty([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$. In the case of a general nonlinear term (as in (1.2)), our method is not working. The main reason is that we do not have a proper transform to give a system without derivative terms. Moreover, our method is not working if the initial data lies on $X^1(\mathbb{R})$. It is because when we study the system of equations, we would have to study it on $L^\infty(\mathbb{R})$, but we know that the Schrödinger group is not bounded from $L^\infty(\mathbb{R})$ to $L^\infty(\mathbb{R})$. Thus, the local wellposedness on less regular spaces is a difficult problem for nonlinear derivative Schrödinger equations.

To prove the conservation laws of (1.1), we need to use a localizing function, which is necessary for integrals to be well defined. Indeed, to obtain the conservation of the energy, using (1.1), at least formally, we have

$$\partial_t(|\partial u|^2) = \partial_x(F(u)) + \partial_t(G(u)),$$

for functions F and G which will be defined later. The important thing is that when u is not in $H^1(\mathbb{R})$, there are some terms in $G(u)$ which do not belong to $L^1(\mathbb{R})$, hence, it is impossible to integrate the two sides as in the usual case. However, we can use a localizing function to deal with this problem. Similarly, we use the localizing function to prove the conservation of the mass and the momentum. The localizing function χ is defined as follows

$$\chi \in C^1(\mathbb{R}) \text{ and even, } \text{supp}\chi \subset [-2, 2], \text{ and } \chi = 1 \text{ on } [-1, 1]. \quad (1.5)$$

For all $a \in \mathbb{R}$ and $R > 0$, we define

$$\chi_{a,R}(x) = \chi\left(\frac{x-a}{R}\right) = \chi\left(\frac{|x-a|}{R}\right). \quad (1.6)$$

To prove the conservation of mass, we use the similar notations as in [4, section 7]

$$m^+(u) = \inf_{a \in \mathbb{R}} \limsup_{R \rightarrow \infty} \int_{\mathbb{R}} (|u|^2 - q_0^2) \chi_{a,R} dx,$$

$$m^-(u) = \sup_{a \in \mathbb{R}} \liminf_{R \rightarrow \infty} \int_{\mathbb{R}} (|u|^2 - q_0^2) \chi_{a,R} dx.$$

If u is such that $m^+(u) = m^-(u)$ we define *generalized mass* as

$$m(u) \equiv m^+(u) = m^-(u).$$

Especially, for $a = 0$ we define

$$\chi_R(x) = \chi\left(\frac{x}{R}\right). \quad (1.7)$$

Our fourth main result is the following.

Theorem 1.4. *Let $q_0 \in \mathbb{R}$ be a constant and $u_0 \in q_0 + H^2(\mathbb{R})$ and $u \in C((T_{min}, T_{max}), q_0 + H^2(\mathbb{R}))$ be the associated solution of (1.1) given by Theorem 1.2. Then, we have*

$$E(u) := \int_{\mathbb{R}} |\partial u|^2 dx + \frac{1}{2} \mathcal{I}m \int_{\mathbb{R}} (|u|^2 \bar{u} - q_0^3) \partial u dx$$

$$+ \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - |q_0|^2)^2 (|u|^2 + 2|q_0|^2) dx = E(u_0), \quad (1.8)$$

$$P(u) := \frac{1}{2} \mathcal{I}m \int_{\mathbb{R}} (u - q_0) \partial \bar{u} dx - \int_{\mathbb{R}} \frac{1}{4} (|u|^2 - |q_0|^2)^2 dx = P(u_0), \quad (1.9)$$

for all $t \in (T_{min}, T_{max})$. Moreover, u satisfies $m^+(u(t)) = m^+(u_0)$ (respectively $m^-(u(t)) = m^-(u_0)$). In particular, if u_0 has finite generalized mass then the generalized mass is conserved by the flow, that is $m(u(t)) = m(u_0)$.

Remark 1.5. When $q_0 = 0$, we recover the classical conservation of mass, energy and momentum as usually defined.

In the classical Schrödinger equation, there are special solutions which are called *standing waves*. There are many works on standing waves (see e.g [12], [2] and the references therein). In [18], Zhidkov shows that there are two types of bounded solitary waves possessing limits as $x \rightarrow \pm\infty$. These are monotone solutions and solutions which have precisely one extreme point. They are called *kinks* and *soliton-like solutions*, respectively. In [18], Zhidkov studied the stability of kinks of classical Schrödinger equations. In [1], the authors have studied the stability of kinks in the energy space. To our knowledge, all these solitary waves are in Zhidkov spaces i.e the Zhidkov space is largest space we know to find special solutions. We want to investigate stationary solutions of (1.1) in Zhidkov spaces. Before stating the next main result, we need the following definition:

Definition 1.6. The stationary solutions of (1.1) are functions $\phi \in X^2(\mathbb{R})$ satisfying

$$\phi_{xx} + i\phi^2\bar{\phi}_x = 0. \quad (1.10)$$

In [14], the authors proved the existence of periodic traveling waves of a derivative nonlinear Schrödinger equation using a skillful changes of variables. In this paper, we use a similar changes of variables as in [14] to prove the existence and uniqueness of stationary solution of (5.2) on $X^2(\mathbb{R})$. Our fifth main result is the following.

Theorem 1.7. *Let ϕ be a stationary solution of (1.1) (see Definition 1.6). The followings is true:*

(1) *If ϕ is not a constant function and satisfies*

$$\inf_{x \in \mathbb{R}} |\phi(x)| > 0$$

then ϕ is of the form $e^{i\theta}\sqrt{k}$ where

$$k(x) = 2\sqrt{B} + \frac{-1}{\sqrt{\frac{5}{72B}} \cosh(2\sqrt{B}(x - x_0)) + \frac{5}{12\sqrt{B}}},$$

$$\theta = \theta_0 - \int_x^\infty \left(\frac{B}{k(y)} - \frac{k(y)}{4} \right) dy,$$

for some constants $\theta_0, x_0 \in \mathbb{R}$, $B > 0$.

(2) *If ϕ is a stationary solution of (1.1) such that $\phi(\infty) = 0$ then $\phi \equiv 0$ on \mathbb{R} .*

Remark 1.8. We have classified stationary solutions of (1.1) for the functions which are vanishing at infinity, and for the functions which are not vanishing on $\overline{\mathbb{R}}$. One question still unanswered is the class of stationary solutions of (1.1) vanishing at a point in \mathbb{R} .

This paper is organized as follows. In Section 2, we give the proof of local well posedness of solution of (1.1) on Zhidkov spaces. In Section 3, we prove the local well posedness on $\phi + H^2(\mathbb{R})$ and $\phi + H^1(\mathbb{R})$, for $\phi \in X^4(\mathbb{R})$ a given function. In Section 4, we give the proof of conservation laws when the initial data is in $q_0 + H^2(\mathbb{R})$, for a given constant $q_0 \in \mathbb{R}$. Finally, in Section 5, we have some results on stationary solutions of (1.1) on Zhidkov spaces.

Notation. In this paper, we will use in the following notation L for the linear part of the Schrödinger equation, that is

$$L = i\partial_t + \partial^2.$$

Moreover, C denotes various positive constants and $C(R)$ denotes constants depending on R .

2. Local existence in Zhidkov spaces. In this section, we give the proof of Theorem 1.1.

2.1. Preliminaries on Zhidkov spaces. Before presenting our main results, we give some preliminaries. We start by recalling the definition of Zhidkov spaces, which were introduced by Peter Zhidkov in his pioneering works on Schrödinger equations with non-zero boundary conditions (see [18] and the references therein).

Definition 2.1. Let $k \in \mathbb{N}$, $k \geq 1$. The Zhidkov space $X^k(\mathbb{R})$ is defined by

$$X^k(\mathbb{R}) = \{u \in L^\infty(\mathbb{R}) : \partial u \in H^{k-1}(\mathbb{R})\}.$$

It is a Banach space when endowed with the norm

$$\|\cdot\|_{X^k} = \|\cdot\|_{L^\infty} + \sum_{\alpha=1}^k \|\partial^\alpha \cdot\|_{L^2}.$$

It was proved by Gallo [5, Theorem 3.1 and Theorem 3.2] that the Schrödinger operator defines a group on Zhidkov spaces. More precisely, we have the following result.

Proposition 2.2. Let $k \geq 1$ and $u_0 \in X^k(\mathbb{R})$. For $t \in \mathbb{R}$ and $x \in \mathbb{R}$, the quantity

$$S(t)u_0(x) := \begin{cases} e^{-i\pi/4}\pi^{-1/2} \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}} e^{(i-\varepsilon)z^2} u_0(x + 2\sqrt{t}z) dz & \text{if } t \geq 0, \\ e^{i\pi/4}\pi^{-1/2} \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}} e^{(-i-\varepsilon)z^2} u_0(x + 2\sqrt{-t}z) dz & \text{if } t \leq 0. \end{cases} \quad (2.1)$$

is well-defined and S defines a strongly continuous group on $X^k(\mathbb{R})$. For all $u_0 \in X^k(\mathbb{R})$ and $t \in \mathbb{R}$ we have

$$\|S(t)u_0\|_{X^k} \leq C(k)(1 + |t|^{1/4})\|u_0\|_{X^k}.$$

The generator of the group $(S(t))|_{t \in \mathbb{R}}$ on $X^k(\mathbb{R})$ is $i\partial^2$ and its domain is $X^{k+2}(\mathbb{R})$.

Remark 2.3. Since, for all $\phi \in X^k(\mathbb{R})$, we have $\phi + H^k(\mathbb{R}) \subset X^k(\mathbb{R})$, the uniqueness of solution in $X^k(\mathbb{R})$ implies the uniqueness of solution in $\phi + H^k(\mathbb{R})$, and the existence of solution in $\phi + H^k(\mathbb{R})$ implies the existence of solution in $X^k(\mathbb{R})$.

2.2. From the equation to the system. The equation (1.1) contains a spatial derivative of u in the nonlinear part, which makes it difficult to work with. In the following proposition, we indicate how to eliminate the derivative in the nonlinearity by introducing an auxiliary function and converting the equation into a system.

Proposition 2.4. Let $k \geq 2$. Given $u \in X^k(\mathbb{R})$, we define v by

$$v = \partial u + \frac{i}{2}|u|^2 u. \quad (2.2)$$

Hence, $v \in X^{k-1}(\mathbb{R})$. Furthermore, if u satisfies the equation (1.1), then the couple (u, v) verifies the system

$$\begin{cases} Lu = P_1(u, v), \\ Lv = P_2(u, v), \end{cases} \quad (2.3)$$

where P_1 and P_2 are given by

$$\begin{aligned} P_1(u, v) &= -iu^2\bar{v} + \frac{1}{2}|u|^4u, \\ P_2(u, v) &= i\bar{u}v^2 + \frac{3}{2}|u|^4v + u^2|u|^2\bar{v}. \end{aligned} \quad (2.4)$$

Proof. Let u be a solution of (1.1) and v be defined by (2.2). Then we have

$$Lu = -iu^2\partial\bar{u} = -iu^2\left(\bar{v} + \frac{i}{2}(|u|^2\bar{u})\right) = -iu^2\bar{v} + \frac{1}{2}|u|^4u,$$

which gives us the first equation in (2.3).

On the other hand, since L and ∂ commute and u solves (1.1), we have

$$\begin{aligned} Lv &= \partial(Lu) + \frac{i}{2}L(|u|^2u) = \partial(-iu^2\partial\bar{u}) + \frac{i}{2}L(|u|^2u) \\ &= -i(u^2\partial^2\bar{u} + 2u|\partial u|^2) + \frac{i}{2}L(|u|^2u). \end{aligned} \quad (2.5)$$

Using

$$L(uv) = L(u)v + uL(v) + 2\partial u\partial v, \quad L(\bar{u}) = -\overline{Lu} + 2\partial^2\bar{u}, \quad (2.6)$$

we have

$$\begin{aligned} L(|u|^2u) &= L(u^2\bar{u}) = L(u^2)\bar{u} + u^2L(\bar{u}) + 2\partial(u^2)\partial\bar{u} \\ &= (2L(u)u + 2(\partial u)^2)\bar{u} + u^2(-\overline{Lu} + 2\partial^2\bar{u}) + 4u|\partial u|^2 \\ &= 2L(u)|u|^2 + 2\bar{u}(\partial u)^2 + 2u^2\partial^2\bar{u} - u^2\overline{Lu} + 4u|\partial u|^2. \end{aligned} \quad (2.7)$$

We now recall that u verifies (1.1) to obtain

$$\frac{i}{2}L(|u|^2u) = u^2\partial\bar{u}|u|^2 + i\bar{u}(\partial u)^2 + iu^2\partial^2\bar{u} + \frac{1}{2}\partial u|u|^4 + 2iu|\partial u|^2. \quad (2.8)$$

Substituting in (2.5), we get

$$\begin{aligned} Lv &= -i(u^2\partial^2\bar{u} + 2u|\partial u|^2) + u^2\partial\bar{u}|u|^2 + i\bar{u}(\partial u)^2 + iu^2\partial^2\bar{u} + \frac{1}{2}\partial u|u|^4 + 2iu|\partial u|^2, \\ &= u^2\partial\bar{u}|u|^2 + i\bar{u}(\partial u)^2 + \frac{1}{2}\partial u|u|^4. \end{aligned}$$

Observe here that the second order derivatives of u have vanished and only first order derivatives remain. Therefore, using the expression of v given in (2.2) to substitute ∂u , we obtain by direct calculations

$$Lv = i\bar{u}v^2 + \frac{3}{2}|u|^4v + u^2|u|^2\bar{v},$$

which gives us the second equation in (2.3). \square

2.3. Resolution of the system. We now establish the local well-posedness of the system (2.3) in Zhidkov spaces.

Proposition 2.5. *Let $k \geq 3$, and $(u_0, v_0) \in X^k(\mathbb{R}) \times X^k(\mathbb{R})$. There exist $T_{min} < 0$, $T_{max} > 0$ and a unique maximal solution (u, v) of system (2.3) such that $(u, v) \in C((T_{min}, T_{max}), X^k(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^{k-2}(\mathbb{R}))$. Furthermore the following properties are satisfied.*

- Blow-up alternative. *If $T_{max} < \infty$ (resp. $T_{min} > -\infty$) then*

$$\lim_{t \rightarrow T_{max} \text{ (resp. } T_{min})} (\|u(t)\|_{X^1} + \|v(t)\|_{X^1}) = \infty.$$

- Continuity with respect to the initial data. *If $(u_0^n, v_0^n) \in X^k \times X^k$ is such that*

$$\|u_0^n - u_0\|_{X^k} + \|v_0^n - v_0\|_{X^k} \rightarrow 0$$

then for any subinterval $[T_1, T_2] \subset (T_{min}, T_{max})$ the associated solution (u^n, v^n) of (2.3) satisfies

$$\lim_{n \rightarrow \infty} (\|u^n - u\|_{L^\infty([T_1, T_2], X^k)} + \|v^n - v\|_{L^\infty([T_1, T_2], X^k)}) = 0.$$

Proof. Consider the operator $A : D(A) \subset X^{k-2}(\mathbb{R}) \rightarrow X^{k-2}(\mathbb{R})$ defined by $A = i\partial^2$ with domain $D(A) = X^k(\mathbb{R})$. From Proposition 2.2 we know that the operator A is the generator of the Schrödinger group $S(t)$ on $X^{k-2}(\mathbb{R})$. From classical arguments (see [3, Lemma 4.1.1 and Corollary 4.1.8]) the couple $(u, v) \in C((T_{min}, T_{max}), X^k(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^{k-2}(\mathbb{R}))$ solves (2.3) if and only if the couple $(u, v) \in C((T_{min}, T_{max}), X^k(\mathbb{R}))$ solves

$$\begin{cases} (u, v) = S(t)(u_0, v_0) - i \int_0^t S(t-s)P(u, v)(s)ds, \\ u(0) = u_0 \in X^k(\mathbb{R}), v(0) = v_0 \in X^k(\mathbb{R}), \end{cases} \quad (2.9)$$

where $S(t)(u, v) := (S(t)u, S(t)v)$, $P(u, v) = (P_1(u, v), P_2(u, v))$ and P_1 and P_2 are defined in (2.4). Consider P as a map from $X^k(\mathbb{R}) \times X^k(\mathbb{R})$ into $X^k(\mathbb{R}) \times X^k(\mathbb{R})$. Since P_1 and P_2 are polynomial in u and v , the map P is Lipschitz continuous on bounded sets of $X^k(\mathbb{R}) \times X^k(\mathbb{R})$. Since (see [3, Theorem 4.3.4 and Theorem 4.3.7]), there exists unique maximal solution $(u, v) \in C((T_{min}, T_{max}), X^k(\mathbb{R}) \times X^k(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^{k-2}(\mathbb{R}) \times X^{k-2}(\mathbb{R}))$ of system (2.3). Moreover, (u, v) satisfy blow-up alternative continuous dependence on initial data in $X^k(\mathbb{R}) \times X^k(\mathbb{R})$. It remains to prove the blow-up alternative in $X^1(\mathbb{R}) \times X^1(\mathbb{R})$. We use the similar arguments as in [18, Proof of Theorem 1.2.4]. For each $1 \leq s \leq k-1$, since the map P is Lipschitz continuous on bounded sets of $X^s(\mathbb{R}) \times X^s(\mathbb{R})$, there exists T_{smin} and T_{smax} such that (u, v) is the maximal $X^s(\mathbb{R}) \times X^s(\mathbb{R})$ solution of system (2.9) on (T_{smin}, T_{smax}) and (u, v) satisfy:

$$\lim_{t \rightarrow T_{smax} \text{ (resp. } T_{smin})} (\|u(t)\|_{X^s} + \|v(t)\|_{X^s}) = \infty.$$

It is sufficient to prove that $T_{1max} = T_{max}$ and $T_{1min} = T_{min}$. We have

$$T_{1max} \geq T_{2max} \geq \dots \geq T_{(k-1)max} \geq T_{max}.$$

We first prove $T_{1max} = T_{2max}$. Assume $T_{1max} > T_{2max}$. For $t \in [0, T_{2max}]$, since (2.9) we have

$$\begin{aligned} & \|u\|_{X^2} + \|v\|_{X^2} \\ & \leq \|u_0\|_{X^2} + \|v_0\|_{X^2} + \max_{t \in [0, T_{2max}]} (\|u\|_{X^1} + \|v\|_{X^1} + 1)^4 \int_0^t (\|u(s)\|_{X^2} \\ & \quad + \|v(s)\|_{X^2}) ds. \end{aligned}$$

By Gronwall's inequality in integral form we obtain

$$\sup_{t \in [0, T_{2max}]} (\|u\|_{X^2} + \|v\|_{X^2}) < \infty.$$

This contradicts to blow-up alternative of (u, v) in $X^2(\mathbb{R}) \times X^2(\mathbb{R})$. Thus, $T_{1max} = T_{2max}$. By apply many times this arguments we obtain $T_{1max} = T_{max}$ and by similar arguments we have $T_{1min} = T_{min}$. This completes the proof of Proposition 2.5. \square

2.4. Preservation of the differential identity. The following proposition establishes the link from (2.3) to (1.1) by showing preservation along the time evolution of the differential identity

$$v_0 = \partial u_0 + \frac{i}{2}|u_0|^2 u_0.$$

Proposition 2.6. *Let $u_0, v_0 \in X^3(\mathbb{R})$ be such that*

$$v_0 = \partial u_0 + \frac{i}{2}u_0|u_0|^2.$$

Then the associated solution $(u, v) \in C((-T_{min}, T_{max}), X^3(\mathbb{R}) \times X^3(\mathbb{R}))$ obtained in Proposition 2.5 satisfies for all $t \in (-T_{min}, T_{max})$ the differential identity

$$v = \partial u + \frac{i}{2}|u|^2 u.$$

Proof. Given $(u, v) \in C((-T_{min}, T_{max}), X^3(\mathbb{R}) \times X^3(\mathbb{R}))$ the solution of (2.3) obtained in Proposition 2.5, we define

$$w = \partial u + \frac{i}{2}|u|^2 u.$$

Our goal will be to show that $w = v$. We first have

$$\begin{aligned} Lu &= -iu^2 \bar{v} + \frac{1}{2}|u|^4 u \\ &= -iu^2(\bar{v} - \bar{w}) - iu^2 \bar{w} + \frac{1}{2}|u|^4 u \\ &= -iu^2(\bar{v} - \bar{w}) - iu^2 \partial \bar{u}. \end{aligned}$$

Applying L to w and using (2.7) and the expression previously obtained for Lu , we get

$$\begin{aligned} Lw &= \partial(Lu) + \frac{i}{2}L(|u|^2 u) \\ &= \partial(Lu) + \frac{i}{2}(2Lu|u|^2 + 2\bar{u}(\partial u)^2 + 2u^2 \partial^2 \bar{u} - u^2 \overline{Lu} + 4u|\partial u|^2) \\ &= \partial(-iu^2(\bar{v} - \bar{w}) - iu^2 \partial \bar{u}) \\ &+ \frac{i}{2}\left(2(-iu^2 \partial \bar{u})|u|^2 + 2\bar{u}(\partial u)^2 - u^2 \overline{(-iu^2 \partial \bar{u})} + 2u^2 \partial^2 \bar{u} + 4u|\partial u|^2\right) \end{aligned}$$

$$\begin{aligned}
& + \frac{i}{2} \left[2(-iu^2(\bar{v} - \bar{w}))|u|^2 - u^2 \overline{(-iu^2(\bar{v} - \bar{w}))} \right] \\
& = \left(-i\partial(u^2(\bar{v} - \bar{w})) + u^2|u|^2(\bar{v} - \bar{w}) + \frac{1}{2}|u|^4(v - w) \right) \\
& + \left(-i\partial(u^2\bar{u}) + u^2\partial\bar{u}|u|^2 + i\bar{u}(\partial u)^2 + \frac{1}{2}|u|^4\partial u + iu^2\partial^2\bar{u} + 2iu|\partial u|^2 \right) \\
& =: I_1 + I_2.
\end{aligned}$$

As in the proof of Proposition 2.4, we obtain

$$I_2 = i\bar{u}w^2 + \frac{3}{2}|u|^4w + |u|^2u^2\bar{w}.$$

Furthermore

$$\begin{aligned}
I_1 & = \partial(-iu^2(\bar{v} - \bar{w})) + u^2|u|^2(\bar{v} - \bar{w}) + \frac{1}{2}|u|^4(v - w) \\
& = -iu^2\partial(\bar{v} - \bar{w}) - 2iu\partial u(\bar{v} - \bar{w}) + u^2|u|^2(\bar{v} - \bar{w}) + \frac{1}{2}|u|^4(v - w).
\end{aligned}$$

It follows that

$$Lw - Lv = I_1 + (I_2 - Lv) \quad (2.10)$$

$$= I_1 + i\bar{u}(w - v)(w + v) + \frac{3}{2}|u|^4(w - v) + |u|^2u^2(\bar{w} - \bar{v}) \quad (2.11)$$

$$= (w - v)A_1 + (\bar{w} - \bar{v})A_2 - iu^2\partial(\bar{v} - \bar{w}), \quad (2.12)$$

where A_1 and A_2 are polynomials of degree at most 4 in u , ∂u , v , ∂v and their complex conjugates. Hence,

$$(Lw - Lv)(\bar{w} - \bar{v}) = |w - v|^2A_1 + (\bar{w} - \bar{v})^2A_2 - iu^2\frac{\partial(\bar{v} - \bar{w})^2}{2} := K, \quad (2.13)$$

where K is a polynomial of degree at most 6 in u , v , w , ∂u , ∂v , ∂w and their complex conjugates. Remembering that $L = i\partial_t + \partial^2$, and taking imaginary part in the two sides of (2.13) we obtain

$$\frac{1}{2}\partial_t|w - v|^2 + \mathcal{I}m(\partial((\partial w - \partial v)(\bar{w} - \bar{v}))) = \mathcal{I}m(K). \quad (2.14)$$

Let $\chi : \mathbb{R} \rightarrow \mathbb{R}$ be a cut-off function such that

$$\begin{aligned}
& \chi \in C^1(\mathbb{R}), \quad \text{supp}(\chi) \subset [-2, 2], \\
& \chi \equiv 1 \text{ on } (-1, 1), \quad 0 \leq \chi \leq 1, \quad |\chi'(x)|^2 \lesssim \chi(x) \text{ for all } x \in \mathbb{R}.
\end{aligned}$$

For each $n \in \mathbb{N}$, define

$$\chi_n(x) = \chi\left(\frac{x}{n}\right).$$

Multiplying both sides of (2.14) by χ_n and integrating in space we obtain

$$\frac{1}{2}\partial_t\|(w - v)\sqrt{\chi_n}\|_{L^2}^2 + \int_{\mathbb{R}} \mathcal{I}m(\partial((\partial w - \partial v)(\bar{w} - \bar{v})))\chi_n dx = \int_{\mathbb{R}} \mathcal{I}m(K)\chi_n dx. \quad (2.15)$$

For the right hand side, we have

$$\begin{aligned}
& \int_{\mathbb{R}} \mathcal{I}m(K)\chi_n dx \\
& = \mathcal{I}m \int_{\mathbb{R}} |w - v|^2 A_1 \chi_n dx + \mathcal{I}m \int_{\mathbb{R}} (\bar{w} - \bar{v})^2 A_2 \chi_n dx - \mathcal{I}m \int_{\mathbb{R}} iu^2 \frac{\partial((\bar{v} - \bar{w})^2)}{2} \chi_n dx,
\end{aligned}$$

and therefore

$$\begin{aligned} & \left| \int_{\mathbb{R}} \mathcal{I}m(K)\chi_n dx \right| \\ & \leq \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 (\|A_1\|_{L^\infty} + \|A_2\|_{L^\infty}) + \frac{1}{2} \left| \int_{\mathbb{R}} u^2 \partial((\bar{v}-\bar{w})^2)\chi_n dx \right|. \end{aligned}$$

We now fix some arbitrary interval $[-T_1, T_2]$ such that $0 \in [-T_1, T_2] \subset (-T_{min}, T^{max})$ in which we will be working from now on, and we set

$$R = \|u\|_{L^\infty([T_1, T_2], X^3)} + \|v\|_{L^\infty([T_1, T_2], X^3)}.$$

From the fact that A_1 and A_2 are polynomials in $u, \partial u, v, \partial v$ of degree at most 4, for all $t \in [T_1, T_2]$ we have

$$\|A_1\|_{L^\infty} + \|A_2\|_{L^\infty} \leq C(R).$$

It follows that

$$\begin{aligned} & \left| \int_{\mathbb{R}} \mathcal{I}m(K)\chi_n dx \right| \\ & \leq \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 C(R) + \frac{1}{2} \left| \int_{\mathbb{R}} (\bar{v}-\bar{w})^2 (\partial(u^2)\chi_n + u^2\partial\chi_n) dx \right|. \end{aligned}$$

By definition of χ we have

$$\begin{aligned} & |\partial(u^2)\chi_n| \leq C(R)\chi_n, \\ & |u^2\partial\chi_n| \leq |u^2| \frac{1}{n} \left| \chi' \left(\frac{\cdot}{n} \right) \right| \leq \frac{1}{n} C(R) \sqrt{\chi \left(\frac{\cdot}{n} \right)} \leq C(R) \frac{1}{n} \sqrt{\chi_n(\cdot)}. \end{aligned}$$

Hence,

$$\begin{aligned} \left| \int_{\mathbb{R}} \mathcal{I}m(K)\chi_n dx \right| & \leq \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 C(R) + \frac{C(R)}{n} \left| \int_{\mathbb{R}} (\bar{v}-\bar{w})^2 \sqrt{\chi_n} dx \right| \\ & \leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{C(R)^2}{n} \int_{\mathbb{R}} |v-w| \sqrt{\chi_n} dx \\ & \leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{C(R)^2}{n} \int_{-2n}^{2n} |v-w| \sqrt{\chi_n} dx \\ & \leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{C(R)^2}{n} \left(\int_{-2n}^{2n} (|v-w| \sqrt{\chi_n})^2 dx \right)^{\frac{1}{2}} \left(\int_{-2n}^{2n} dx \right)^{\frac{1}{2}} \\ & \leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{2C(R)^2}{\sqrt{n}} \|(w-v)\sqrt{\chi_n}\|_{L^2}. \end{aligned} \quad (2.16)$$

In addition, we have

$$\begin{aligned} \left| \int_{\mathbb{R}} \mathcal{I}m(\partial((\partial w - \partial v)(\bar{w} - \bar{v}))\chi_n) dx \right| & = \left| \int_{\mathbb{R}} \mathcal{I}m(((\partial w - \partial v)(\bar{w} - \bar{v}))\chi_n') dx \right| \\ & = \left| \int_{\mathbb{R}} \mathcal{I}m \left((\partial w - \partial v)(\bar{w} - \bar{v}) \frac{1}{n} \chi' \left(\frac{x}{n} \right) \right) dx \right| \\ & \leq \int_{\mathbb{R}} |\partial w - \partial v| |w-v| \frac{1}{n} \sqrt{\chi_n} dx \\ & \leq \frac{1}{n} \|\partial w - \partial v\|_{L^2} \|(w-v)\sqrt{\chi_n}\|_{L^2} \\ & \leq \frac{C(R)}{n} \|(w-v)\sqrt{\chi_n}\|_{L^2}. \end{aligned} \quad (2.17)$$

From (2.15), (2.16), (2.17) we obtain that

$$\partial_t \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 \leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{C(R)}{\sqrt{n}} \|(w-v)\sqrt{\chi_n}\|_{L^2} \quad (2.18)$$

$$\leq C(R) \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 + \frac{C(R)}{\sqrt{n}} \quad (2.19)$$

where we have used the Cauchy inequality $|x| \leq \frac{|x|^2+1}{2}$. Define the function $g : [-T_1, T_2]$ by

$$g = \|(w-v)\sqrt{\chi_n}\|_{L^2}^2.$$

Then by definition of w we have $g(t=0) = 0$. Furthermore, from (2.19) we have

$$\partial_t g \leq C(R)g + \frac{C(R)}{\sqrt{n}}.$$

By Gronwall inequality for all $t \in [-T_1, T_2]$ we have

$$g \leq \frac{C(R)}{\sqrt{n}} \exp(C(R)(T_2 + T_1)) \leq \frac{C(R)}{\sqrt{n}}. \quad (2.20)$$

Assume by contradiction that there exist t and x such that

$$w(t, x) \neq v(t, x).$$

By continuity of v and w , there exists $\varepsilon > 0$ such that (for $n > |x|$) we have

$$g(t) = \|(w-v)\sqrt{\chi_n}\|_{L^2}^2 > \varepsilon.$$

Since $\varepsilon > 0$ is independant of n , we obtain a contradiction with (2.20) when n is large enough. Therefore for all t and x , we have

$$v(t, x) = w(t, x),$$

which concludes the proof. \square

2.5. From the system to the equation. With Proposition 2.6 in hand, we give the proof of Theorem 1.1.

Proof of Theorem 1.1. We start by defining v_0 by

$$v_0 = \partial u_0 + \frac{i}{2}|u_0|^2 u_0 \in X^3(\mathbb{R}).$$

From Proposition 2.5 there exists a unique maximal solution $(u, v) \in C((T_{min}, T_{max}), X^3(\mathbb{R}) \times X^3(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^1(\mathbb{R}) \times X^1(\mathbb{R}))$ of the system (2.3) associated with (u_0, v_0) . From Proposition 2.6, for all $t \in (T_{min}, T_{max})$ we have

$$v = \partial u + \frac{i}{2}|u|^2 u. \quad (2.21)$$

It follows that

$$Lu = -iu^2 \bar{v} + \frac{1}{2}|u|^4 u = -iu^2 \partial \bar{u},$$

and therefore u is a solution of (1.1) on (T_{min}, T_{max}) . Furthermore

$$u \in C((T_{min}, T_{max}), X^3(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^1(\mathbb{R})).$$

To obtain the desired regularity on u , we observe that, since v has the same regularity as u , and verifies (2.21), we have

$$\partial u = v - \frac{i}{2}|u|^2 u \in C((T_{min}, T_{max}), X^3(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^1(\mathbb{R}))$$

This implies that

$$u \in C((T_{min}, T_{max}), X^4(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^2(\mathbb{R})).$$

This proves the existence part of the result. Uniqueness is a direct consequence from Proposition 2.4 and Proposition 2.5.

To prove the blow-up alternative, assume that $T_{max} < \infty$. Then from Proposition 2.5 we have

$$\lim_{t \rightarrow T_{max}} (\|u(t)\|_{X^1(\mathbb{R})} + \|v(t)\|_{X^1(\mathbb{R})}) = \infty$$

On the other hand, since (2.21) we obtain

$$\lim_{t \rightarrow T_{max}} (\|u(t)\|_{X^1(\mathbb{R})} + \|\partial u(t)\|_{X^1(\mathbb{R})}) = \infty.$$

It follows that

$$\lim_{t \rightarrow T_{max}} \|u(t)\|_{X^2(\mathbb{R})} = \infty.$$

Finally, we establish the continuity with respect to the initial data. Take a subinterval $[T_1, T_2] \subset (T_{min}, T_{max})$, and a sequence $(u_0^n) \in X^4(\mathbb{R})$ such that $u_0^n \rightarrow u_0$ in X^4 . Let u_n be the solution of (1.1) associated with u_0^n and define v_n by

$$v_n = \partial u_n + \frac{i}{2}|u_n|^2 u_n. \quad (2.22)$$

By Proposition 2.5 the couple (u_n, v_n) is the unique maximal solution of system (2.3) in

$$C((T_{min}, T_{max}), X^3(\mathbb{R}) \times X^3(\mathbb{R})) \cap C^1((T_{min}, T_{max}), X^1(\mathbb{R}) \times X^1(\mathbb{R})).$$

Moreover, we have

$$\lim_{n \rightarrow +\infty} (\|u_n - u\|_{L^\infty([T_1, T_2], X^3)} + \|v_n - v\|_{L^\infty([T_1, T_2], X^3)}) = 0 \quad (2.23)$$

Since v and v_n verify the differential identity (2.22), we have

$$\partial(u_n - u) = (v_n - v) - \frac{i}{2}(|u_n|^2 u_n - |u|^2 u).$$

Therefore we have

$$\lim_{n \rightarrow +\infty} \|u_n - u\|_{L^\infty([T_1, T_2], X^4)} = 0,$$

which completes the proof. \square

3. Results on the space $\phi + H^k(\mathbb{R})$ for $\phi \in X^k(\mathbb{R})$. In this section, we give the proof of Theorem 1.2 and Theorem 1.3. For $k \geq 1$, let $\phi \in X^k(\mathbb{R})$.

3.1. The local well posedness on $\phi + H^2(\mathbb{R})$.

3.1.1. *From the equation to the system.* Define

$$v = \partial u + \frac{i}{2}|u|^2 u. \quad (3.1)$$

Since Proposition 2.4, if u solves (1.1) then (u, v) solves the following system:

$$\begin{cases} Lu = -iu^2 \bar{v} + \frac{1}{2}|u|^4 u, \\ Lv = i\bar{u}v^2 + \frac{3}{2}|u|^4 v + u^2|u|^2 \bar{v}, \\ u(0) = u_0, \\ v(0) = v_0 := \partial u_0 + \frac{i}{2}|u_0|^2 u_0. \end{cases} \quad (3.2)$$

Let $\phi \in X^4(\mathbb{R})$. Define $\tilde{u} = u - \phi$, $\tilde{v} = v - \frac{i}{2}|\phi|^2\phi$. We have if u solves (1.1) then (\tilde{u}, \tilde{v}) solves:

$$\begin{cases} L\tilde{u} = Q_1(\tilde{u}, \tilde{v}, \phi), \\ L\tilde{v} = Q_2(\tilde{u}, \tilde{v}, \phi), \\ \tilde{u}(0) = \tilde{u}_0 := u_0 - \phi, \\ \tilde{v}(0) = \tilde{v}_0 := v_0 - \frac{i}{2}|\phi|^2\phi, \end{cases} \quad (3.3)$$

where

$$Q_1(\tilde{u}, \tilde{v}, \phi) = -i(\tilde{u} + \phi)^2 \left(\bar{\tilde{v}} - \frac{i}{2}|\phi|^2\bar{\phi} \right) + \frac{1}{2}|\tilde{u} + \phi|^4(\tilde{u} + \phi) - L(\phi), \quad (3.4)$$

$$\begin{aligned} Q_2(\tilde{u}, \tilde{v}, \phi) &= i(\bar{\tilde{u}} + \bar{\phi}) \left(\tilde{v} + \frac{i}{2}|\phi|^2\phi \right)^2 + \frac{3}{2}|\tilde{u} + \phi|^4 \left(\tilde{v} + \frac{i}{2}|\phi|^2\phi \right) \\ &\quad + (\tilde{u} + \phi)^2|\tilde{u} + \phi|^2 \left(\bar{\tilde{v}} - \frac{i}{2}|\phi|^2\bar{\phi} \right) - \frac{i}{2}L(|\phi|^2\phi). \end{aligned} \quad (3.5)$$

3.1.2. Resolution of the system. Let $k \geq 1$. We note that if $\phi \in X^{k+2}$ then $Q_1 : (\tilde{u}, \tilde{v}) \rightarrow Q_1(\tilde{u}, \tilde{v}, \phi)$ and $Q_2 : (\tilde{u}, \tilde{v}) \rightarrow Q_2(\tilde{u}, \tilde{v}, \phi)$ defined as in (3.4) and (3.5) are Lipschitz continuous on bounded set of $H^k(\mathbb{R}) \times H^k(\mathbb{R})$. By similar arguments to the one used for the proof of Proposition 2.5, we obtain the following local well-posedness result:

Proposition 3.1. *Let $k \geq 1$, $\phi \in X^{k+2}$, $\tilde{u}_0, \tilde{v}_0 \in H^k(\mathbb{R})$. There exist $T_{min} < 0$, $T_{max} > 0$ and a unique maximal solution (\tilde{u}, \tilde{v}) of the system (3.3) such that $\tilde{u}, \tilde{v} \in C((T_{min}, T_{max}), H^k(\mathbb{R})) \cap C^1((T_{min}, T_{max}), H^{k-2}(\mathbb{R}))$. Furthermore the following properties are satisfied.*

- Blow-up alternative. *If $T_{max} < \infty$ (resp. $T_{min} > -\infty$) then*

$$\lim_{t \rightarrow T_{max} \text{ (resp. } T_{min})} (\|\tilde{u}\|_{H^k} + \|\tilde{v}\|_{H^k}) = \infty.$$

- Continuity with respect to the initial data. *If $\tilde{u}_0^n, \tilde{v}_0^n \in H^k(\mathbb{R})$ are such that*

$$\|\tilde{u}_0^n - \tilde{u}_0\|_{H^k} + \|\tilde{v}_0^n - \tilde{v}_0\|_{H^k} \rightarrow 0$$

then for any subinterval $[T_1, T_2] \subset (T_{min}, T_{max})$ the associated solution $(\tilde{u}^n, \tilde{v}^n)$ of (3.3) satisfies

$$\lim_{n \rightarrow +\infty} (\|\tilde{u}^n - \tilde{u}\|_{L^\infty([T_1, T_2], H^k)} + \|\tilde{v}^n - \tilde{v}\|_{L^\infty([T_1, T_2], H^k)}) = 0.$$

3.1.3. Preservation of a differential identity. Let $(\tilde{u}_0, \tilde{v}_0)$ be defined as in section 3.1.1. By elementary calculation we have

$$\tilde{w}_0 = \partial\tilde{u}_0 + \frac{i}{2}(|\tilde{u}_0 + \phi|^2(\tilde{u}_0 + \phi) - |\phi|^2\phi) + \partial\phi. \quad (3.6)$$

We have the following results:

Proposition 3.2. *Let $\phi \in X^4(\mathbb{R})$ and $\tilde{u}_0, \tilde{v}_0 \in H^2(\mathbb{R})$ satisfy (3.6). Then the associated solution (\tilde{u}, \tilde{v}) obtained in Proposition 3.1 also satisfy (3.6) for all $t \in (T_{min}, T_{max})$.*

Proof. We define

$$\tilde{w} = \partial\tilde{u} + \frac{i}{2}(|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2\phi) + \partial\phi. \quad (3.7)$$

Set $u = \tilde{u} + \phi$, $v = \tilde{v} + \frac{i}{2}|\phi|^2\phi$, $w = \tilde{w} + \frac{i}{2}|\phi|^2\phi$. We have

$$w = \partial u + \frac{i}{2}|u|^2u. \quad (3.8)$$

Since (\tilde{u}, \tilde{v}) is a solution of (3.3), we have (u, v) is a solution of (3.2). We have

$$Lu = -iu^2(\bar{v} - \bar{w}) + H,$$

where H defined by

$$H = -iu^2\bar{w} + \frac{1}{2}|u|^4u.$$

By using (2.7) and the previously expression obtained for Lu , we get

$$\begin{aligned} Lw &= \partial(Lu) + \frac{i}{2}L(|u|^2u) \\ &= \partial(Lu) + \frac{i}{2}\left(2L(u)|u|^2 + 2\bar{u}(\partial u)^2 + 2u^2\partial^2\bar{u} - u^2\overline{L(u)} + 4u|\partial u|^2\right) \\ &= \partial(-iu^2(\bar{v} - \bar{w})) + \partial H \\ &+ i\left(H|u|^2 - iu^2|u|^2(\bar{v} - \bar{w}) + \bar{u}(\partial u)^2 + u^2\partial^2\bar{u} - \frac{1}{2}u^2(i\bar{u}^2(v - w) + \bar{H}) + 2u|\partial u|^2\right) \\ &= -i\partial(u^2(\bar{v} - \bar{w})) + u^2|u|^2(\bar{v} - \bar{w}) + \frac{1}{2}|u|^4(v - w) + K, \end{aligned}$$

where K is defined by

$$K = \partial H + iH|u|^2 + i\bar{u}(\partial u)^2 + iu^2\partial^2\bar{u} - \frac{i}{2}u^2\bar{H} + 2iu|\partial u|^2.$$

Using (3.8) to replace the term ∂u in K and remark that the role of w is the same the one of v as in Proposition 2.4, we have

$$K = i\bar{u}w^2 + \frac{3}{2}|u|^4w + u^2|u|^2\bar{w}.$$

Thus,

$$\begin{aligned} Lw - Lv &= -i\partial(u^2(\bar{v} - \bar{w})) + u^2|u|^2(\bar{v} - \bar{w}) + \frac{1}{2}|u|^4(v - w) + (K - L(v)) \\ &= -i\partial(u^2(\bar{v} - \bar{w})) + u^2|u|^2(\bar{v} - \bar{w}) \\ &+ \frac{1}{2}|u|^4(v - w) + i\bar{u}(w^2 - v^2) + \frac{3}{2}|u|^4(w - v) + u^2|u|^2(\bar{w} - \bar{v}) \\ &= -iu^2\partial(\bar{v} - \bar{w}) + A(v - w) + B(\bar{v} - \bar{w}), \end{aligned}$$

where

$$\begin{aligned} A &:= -|u|^4 - i\bar{u}(v + w), \\ B &:= -2iu\partial u = -2iu\left(w - \frac{i}{2}|u|^2u\right) = -2iuv - |u|^2u^2. \end{aligned}$$

This implies that

$$L(\tilde{w} - \tilde{v}) = -i(\tilde{u} + \phi)^2\partial(\bar{\tilde{v}} - \bar{\tilde{w}}) + A(\tilde{v} - \tilde{w}) + B(\bar{\tilde{v}} - \bar{\tilde{w}}). \quad (3.9)$$

Multiplying both sides of (3.9) by $\bar{\tilde{w}} - \bar{\tilde{v}}$, taking the imaginary part, and integrating over space with integration by part for the first term of right hand side of (3.9), we obtain

$$\frac{d}{dt}\|\tilde{w} - \tilde{v}\|_{L^2}^2 \lesssim (\|\tilde{u} + \phi\|_{L^\infty}\|\partial\tilde{u} + \partial\phi\|_{L^\infty} + \|A\|_{L^\infty} + \|B\|_{L^\infty})\|\tilde{w} - \tilde{v}\|_{L^2}^2.$$

By Grönwall's inequality we obtain

$$\begin{aligned} & \|\tilde{w} - \tilde{v}\|_{L^2}^2 \\ & \leq \|\tilde{w}(0) - \tilde{v}(0)\|_{L^2}^2 \times \exp\left(C \int_0^t (\|\tilde{u} + \phi\|_{L^\infty} \|\partial\tilde{u} + \partial\phi\|_{L^\infty} + \|A\|_{L^\infty} + \|B\|_{L^\infty}) ds\right). \end{aligned}$$

Using the fact that $\tilde{w}(0) = \tilde{v}(0)$, we obtain $\tilde{w} = \tilde{v}$, for all t . This implies that

$$\tilde{v} = \partial\tilde{u} + \frac{i}{2}(|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2\phi) + \partial\phi.$$

This completes the proof of Proposition 3.2. \square

3.1.4. *From the system to the equation.* Now, we finish the proof of Theorem 1.2.

Proof of Theorem 1.2. Let $\phi \in X^4(\mathbb{R})$ and $u_0 \in \phi + H^2(\mathbb{R})$. We define $v_0 \in X^1(\mathbb{R})$, $\tilde{u}_0 \in H^2(\mathbb{R})$ and $\tilde{v}_0 \in H^1(\mathbb{R})$ in the following way:

$$v_0 = \partial u_0 + \frac{i}{2}u_0|u_0|^2, \quad \tilde{u}_0 = u_0 - \phi, \quad \text{and} \quad \tilde{v}_0 = v_0 - \frac{i}{2}|\phi|^2\phi.$$

We have

$$\tilde{v}_0 = \partial\tilde{u}_0 + \frac{i}{2}(|\tilde{u}_0 + \phi|^2(\tilde{u}_0 + \phi) - |\phi|^2\phi) + \partial\phi.$$

From Proposition 3.1 there exists a unique maximal solution $(\tilde{u}, \tilde{v}) \in C((T_{min}, T_{max}), H^1(\mathbb{R})) \cap C^1((T_{min}, T_{max}), H^{-1}(\mathbb{R}))$ of (3.3). Let $\tilde{u}_0^n \in H^3(\mathbb{R})$ be such that

$$\|\tilde{u}_0^n - \tilde{u}_0\|_{H^2(\mathbb{R})} \rightarrow 0$$

as $n \rightarrow \infty$. Define $\tilde{v}_0^n \in H^2(\mathbb{R})$ by

$$\tilde{v}_0^n = \partial\tilde{u}_0^n + \frac{i}{2}(|\tilde{u}_0^n + \phi|^2(\tilde{u}_0^n + \phi) - |\phi|^2\phi) + \partial\phi.$$

From Proposition 3.1, there exists a unique solution maximal solution.

$$\tilde{u}^n, \tilde{v}^n \in C((T_{min}^n, T_{max}^n), H^2(\mathbb{R})) \cap C^1((T_{min}^n, T_{max}^n), L^2(\mathbb{R}))$$

of the system (3.3). Let $[T_1, T_2] \subset (T_{min}, T_{max})$ be any closed interval. From [3, proposition 4.3.7], for $n \geq N_0$ large enough, we have $[T_1, T_2] \subset (T_{min}^n, T_{max}^n)$. By Proposition 3.2, for $n \geq N_0$, $t \in [T_1, T_2]$, we have

$$\tilde{v}^n = \partial\tilde{u}^n + \frac{i}{2}(|\tilde{u}^n + \phi|^2(\tilde{u}^n + \phi) - |\phi|^2\phi) + \partial\phi.$$

By Proposition 3.1, we have

$$\lim_{n \rightarrow \infty} \sup_{t \in [T_1, T_2]} (\|\tilde{u}^n(t) - \tilde{u}(t)\|_{H^1(\mathbb{R})} + \|\tilde{v}^n(t) - \tilde{v}(t)\|_{H^1(\mathbb{R})}) \rightarrow 0.$$

We obtain that for all $t \in [T_1, T_2]$, and then for all $t \in (T_{min}, T_{max})$:

$$\tilde{v} = \partial\tilde{u} + \frac{i}{2}(|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2\phi) + \partial\phi.$$

This follows that

$$\partial\tilde{u} \in C((T_{min}, T_{max}), H^1(\mathbb{R})) \cap C^1((T_{min}, T_{max}), H^{-1}(\mathbb{R})).$$

Hence we have

$$\tilde{u} \in C((T_{min}, T_{max}), H^2(\mathbb{R})) \cap C^1((T_{min}, T_{max}), L^2(\mathbb{R})).$$

Define $u = \phi + \tilde{u}$ and define v by

$$v = \tilde{v} + \frac{i}{2}|\phi|^2\phi = \partial u + \frac{i}{2}|u|^2u.$$

Since (\tilde{u}, \tilde{v}) solves (3.3), we have (u, v) solves (3.2). Therefore, $u \in \phi + C((T_{min}, T_{max}), H^2(\mathbb{R})) \cap C^1((T_{min}, T_{max}), L^2(\mathbb{R}))$ solves:

$$Lu = -iu^2\bar{v} + \frac{1}{2}|u|^4u = -iu^2\partial\bar{u}.$$

This establishes the existence of a solution to (1.1). To prove uniqueness, assume that $U \in \phi + C((T_{min}, T_{max}), H^2(\mathbb{R})) \cap C^1((T_{min}, T_{max}), L^2(\mathbb{R}))$ is another solution of (1.1). Set $V = \partial U + \frac{i}{2}|U|^2U$, and $\tilde{U} = U - \phi$, $\tilde{V} = V - \frac{i}{2}|\phi|^2\phi$. Thus, $(\tilde{U}, \tilde{V}) \in C((T_{min}, T_{max}), H^1(\mathbb{R})) \cap C^1((T_{min}, T_{max}), H^{-1}(\mathbb{R}))$ is a solution of (3.3). By the uniqueness statement in Proposition 3.1, we obtain $\tilde{U} = \tilde{u}$. Hence, $u = U$, which proves uniqueness. The blow-up alternative and continuity with respect to the initial data are proved using similar arguments as in the proof of Theorem 1.1. This completes the proof of Theorem 1.2. \square

3.2. The local well posedness on $\phi + H^1(\mathbb{R})$. In this section, we give the proof of Theorem 1.3, using the method of Hayashi and Ozawa [11]. As in Section 3.1.1, we work with the system (3.3).

3.2.1. Resolution of the system. Since we are working in the less regular space $\phi + H^1(\mathbb{R})$, we cannot use Proposition 3.1. Instead, we establish the following result using Strichartz estimate.

Proposition 3.3. *Consider the system (3.3). Let $\phi \in X^2(\mathbb{R})$, $\tilde{u}_0, \tilde{v}_0 \in L^2(\mathbb{R})$. There exists $R > 0$ such that if $\|\tilde{u}_0\|_{L^2} + \|\tilde{v}_0\|_{L^2} < R$ then there exist $T > 0$ and a unique solution (\tilde{u}, \tilde{v}) of the system (3.3) verifying*

$$\tilde{u}, \tilde{v} \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R})).$$

Moreover, we have the following continuous dependence on initial data property: If $(\tilde{u}_0^n, \tilde{v}_0^n) \in L^2(\mathbb{R}) \times L^2(\mathbb{R})$ is a sequence such that $\|\tilde{u}_0^n - \tilde{u}_0\|_2 + \|\tilde{v}_0^n - \tilde{v}_0\|_2 \rightarrow 0$ then for n large enough we have $\|\tilde{u}_0^n\|_2 + \|\tilde{v}_0^n\|_2 < R$ and the associated solutions $(\tilde{u}^n, \tilde{v}^n)$ satisfy:

$$\|\tilde{u}^n - \tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|\tilde{v}^n - \tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty} \rightarrow 0,$$

where we have used the following notation:

$$L^\infty L^2 = L^\infty([-T, T], L^2(\mathbb{R})), \quad L^4 L^\infty = L^4([-T, T], L^\infty(\mathbb{R}))$$

and the norm on $L^\infty L^2 \cap L^4 L^\infty$ is defined, as usual for the intersection of two Banach spaces, as the sum of the norms on each space.

Proof. Let Q_1, Q_2 be defined as in system (3.3). By direct computations, we have

$$Q_1(\tilde{u}, \tilde{v}, \phi) = -i(\tilde{u} + \phi)^2\bar{\tilde{v}} - \frac{1}{2}|\phi|^2\bar{\phi}(\tilde{u}^2 + 2\tilde{u}\phi) + \frac{1}{2}(|\tilde{u} + \phi|^4 - |\phi|^4)\tilde{u} - \partial^2\phi, \quad (3.10)$$

$$\begin{aligned} Q_2(\tilde{u}, \tilde{v}, \phi) &= i\bar{\tilde{u}} \left(\tilde{v} + \frac{i}{2}|\phi|^2\phi \right)^2 + i\bar{\phi} \left[\left(\tilde{v} + \frac{i}{2}|\phi|^2\phi \right)^2 - \left(\frac{i}{2}|\phi|^2\phi \right)^2 \right] + \frac{3}{2}|\tilde{u} + \phi|^4\tilde{v}, \\ &+ \frac{3}{4}i|\phi|^2\phi(|\tilde{u} + \phi|^4 - |\phi|^4) + \bar{\tilde{v}}(\tilde{u} + \phi)|\tilde{u} + \phi|^2 \\ &- \frac{i}{2}|\phi|^2\bar{\phi}((\tilde{u} + \phi)^2|\tilde{u} + \phi|^2 - |\phi|^2\phi^2) - \frac{i}{2}\partial^2(|\phi|^2\phi). \end{aligned} \quad (3.11)$$

Thus,

$$\begin{aligned}
& |Q_1(\tilde{u}, \tilde{v}, \phi)| \\
& \lesssim |\tilde{v}|(|\tilde{u}|^2 + |\phi|^2) + |\phi|^3|\tilde{u}|^2 + |\phi|^4|\tilde{u}| + (|\tilde{u}|^5 + |\phi|^4|\tilde{u}|) + |\phi|(|\tilde{u}|^4 + |\tilde{u}||\phi|^3) + |\partial^2\phi| \\
& \lesssim |\tilde{v}||\tilde{u}|^2 + |\tilde{v}||\phi|^2 + |\tilde{u}|^5 + |\tilde{u}||\phi|^4 + |\partial^2\phi|, \\
& |Q_2(\tilde{u}, \tilde{v}, \phi)| \\
& \lesssim |\tilde{u}|(|\tilde{v}|^2 + |\phi|^6) + |\phi|(|\tilde{v}|^2 + |\tilde{v}||\phi|^3) + |\tilde{v}|(|\tilde{u}|^4 + |\phi|^4) \\
& \quad + |\phi|^3(|\tilde{u}|^4 + |\tilde{u}||\phi|^3) + |\tilde{v}|(|\tilde{u}|^3 + |\phi|^3) + |\phi|^3(|\tilde{u}|^4 + |\phi|^3|\tilde{u}|) + |\partial^2(|\phi|^2\phi)| \\
& \lesssim |\tilde{u}||\tilde{v}|^2 + |\tilde{u}||\phi|^6 + |\phi||\tilde{v}|^2 + |\phi|^4|\tilde{v}| + |\tilde{u}|^4|\tilde{v}| + |\phi|^3|\tilde{u}|^4 \\
& \quad + |\tilde{u}|^3|\tilde{v}| + |\phi|^3|\tilde{v}| + |\partial^2(|\phi|^2\phi)|.
\end{aligned}$$

Consider the following problem

$$(\tilde{u}, \tilde{v}) = S(t)(\tilde{u}_0, \tilde{v}_0) - i \int_0^t S(t-s)Q(\tilde{u}, \tilde{v}, \phi) ds \quad (3.12)$$

where $Q = (Q_1, Q_2)$. Let

$$\Phi(\tilde{u}, \tilde{v}) = S(t)(\tilde{u}_0, \tilde{v}_0) - i \int_0^t S(t-s)Q ds.$$

Assume that $\|\tilde{u}_0\|_{L^2(\mathbb{R})} + \|\tilde{v}_0\|_{L^2(\mathbb{R})} \leq \frac{R}{4}$ for $R > 0$ small enough. For $T > 0$ we define the space $X_{T,R}$ by

$$\begin{aligned}
X_{T,R} = \{ & (\tilde{u}, \tilde{v}) \in (C([-T, T], L^2(\mathbb{R}))) \\
& \cap L^4([-T, T], L^\infty(\mathbb{R}))\}^2 : \|\tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|\tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty} \leq R \}.
\end{aligned}$$

We are going to prove that for R, T small enough the map Φ is a contraction from $X_{T,R}$ to itself.

We first prove that for R, T small enough, Φ maps $X_{T,R}$ into $X_{T,R}$. Let $(\tilde{u}, \tilde{v}) \in X_{T,R}$. By Strichartz estimates we have

$$\begin{aligned}
\|\Phi(\tilde{u}, \tilde{v})\|_{(L^\infty L^2 \cap L^4 L^\infty)^2} & \lesssim \|(\tilde{u}_0, \tilde{v}_0)\|_{L^2 \times L^2} + \|Q\|_{L^1 L^2 \times L^1 L^2}, \\
& \lesssim \frac{R}{4} + (\|Q_1\|_{L^1 L^2} + \|Q_2\|_{L^1 L^2}).
\end{aligned}$$

We have

$$\begin{aligned}
\|Q_1\|_{L^1 L^2} & \lesssim \| |\tilde{u}|^2 \tilde{v} \|_{L^1 L^2} + \| |\tilde{v}| |\phi|^2 \|_{L^1 L^2} + \| |\tilde{u}|^5 \|_{L^1 L^2} + \| \partial^2 \phi \|_{L^1 L^2} \\
& \lesssim \| \tilde{v} \|_{L^2 L^2} \| \tilde{u} \|_{L^4 L^\infty}^2 + \| \tilde{v} \|_{L^2 L^2} \| |\phi|^2 \|_{L^4 L^\infty} + \| \tilde{u} \|_{L^4 L^\infty}^4 \| \tilde{u} \|_{L^\infty L^2} + \| \partial^2 \phi \|_{L^1 L^2} \\
& \lesssim (2T)^{\frac{1}{2}} \| \tilde{v} \|_{L^\infty L^2} \| \tilde{u} \|_{L^4 L^\infty}^2 + (2T)^{\frac{1}{2}} \| \tilde{v} \|_{L^\infty L^2} \| \phi \|_{L^\infty} (2T)^{\frac{1}{4}} \\
& \quad + \| \tilde{u} \|_{L^4 L^\infty}^4 \| \tilde{u} \|_{L^\infty L^2} + \| \partial^2 \phi \|_{L^2} (2T) \\
& \lesssim (2T)^{\frac{1}{2}} R^3 + (2T)^{\frac{3}{4}} \| \phi \|_{L^\infty} R + R^5 + (2T) \| \phi \|_{X^2} < \frac{R}{4}.
\end{aligned}$$

for T, R small enough. Similarly, we also have

$$\|\tilde{Q}_2\|_{L^1 L^2} < \frac{R}{4}$$

for T, R small enough. Therefore, for T, R small enough, we have

$$\|\Phi(\tilde{u}, \tilde{v})\|_{(L^\infty L^2 \cap L^4 L^\infty)^2} < \frac{3R}{4} < R.$$

Hence, Φ maps from $X_{T,R}$ into itself.

We now show that for T, R small enough, the map Φ is a contraction from $X_{T,R}$ to itself.

Indeed, let $(u_1, v_1), (u_2, v_2) \in X_{T,R}$. By Strichartz estimates we have

$$\begin{aligned} & \|\Phi(u_1, v_1) - \Phi(u_2, v_2)\|_{(L^\infty L^2 \cap L^4 L^\infty)^2} \\ &= \left\| \int_0^t S(t-s) (Q(u_1, v_1) - Q(u_2, v_2)) ds \right\|_{(L^\infty L^2 \cap L^4 L^\infty)^2}, \\ &\lesssim \|Q_1(u_1, v_1) - Q_1(u_2, v_2)\|_{L^1 L^2} + \|Q_2(u_1, v_1) - Q_2(u_2, v_2)\|_{L^1 L^2}. \end{aligned}$$

Using the same kind of arguments as before we obtain that Φ is a contraction on $X_{T,R}$. Therefore, using the Banach fixed-point theorem, there exist $T > 0$ and a unique solution $(\tilde{u}, \tilde{v}) \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$ of the problem (3.12). As above, we see that if $h, k \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$ then $Q_1(h, k, \phi), Q_2(h, k, \phi) \in L^1([-T, T], L^2(\mathbb{R}))$. By [3, Proposition 4.1.9], $(\tilde{u}, \tilde{v}) \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$ solves (3.12) if and only if (\tilde{u}, \tilde{v}) solves (3.3). Thus, we proved the existence of a solution of (3.3). The uniqueness of solution of (3.3) is obtained by the uniqueness of solution of (3.12).

It remains to prove the continuous dependence on initial data. Assume that $(u_0^n, v_0^n) \in L^2(\mathbb{R}) \times L^2(\mathbb{R})$ is such that

$$\|u_0^n - \tilde{u}_0\|_{L^2(\mathbb{R})} + \|v_0^n - \tilde{v}_0\|_{L^2(\mathbb{R})} \rightarrow 0,$$

as $n \rightarrow \infty$. In particular, for n large enough, we have

$$\|u_0^n\|_{L^2(\mathbb{R})} + \|v_0^n\|_{L^2(\mathbb{R})} < R.$$

There exists a unique maximal solution (u^n, v^n) of system (3.3), and we may assume that for n large enough, (u^n, v^n) is defined on $[-T, T]$. Assume that T small enough such that

$$\|\tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|\tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty} + \sup_n (\|u^n\|_{L^\infty L^2 \cap L^4 L^\infty} + \|v^n\|_{L^\infty L^2 \cap L^4 L^\infty}) \leq 2R. \quad (3.13)$$

We have (\tilde{u}, \tilde{v}) is a solution of the following system

$$(\tilde{u}, \tilde{v}) = S(t)(\tilde{u}_0, \tilde{v}_0) - i \int_0^t S(t-s)(Q_1(\tilde{u}, \tilde{v}, \phi), Q_2(\tilde{u}, \tilde{v}, \phi)).$$

Similarly, (u^n, v^n) are solutions of the following system

$$(u^n, v^n) = S(t)(u_0^n, v_0^n) - i \int_0^t S(t-s)(Q_1(u^n, v^n, \phi), Q_2(u^n, v^n, \phi)).$$

Hence,

$$\begin{aligned} & (u^n - \tilde{u}, v^n - \tilde{v}) \\ &= S(t)(u_0^n - \tilde{u}_0, v_0^n - \tilde{v}_0) \\ &- i \int_0^t S(t-s)(Q_1(\tilde{u}, \tilde{v}, \phi) - Q_1(u^n, v^n, \phi), Q_2(\tilde{u}, \tilde{v}, \phi) - Q_2(u^n, v^n, \phi)). \end{aligned}$$

Using Strichartz estimates and (3.13), for all $t \in [-T, T]$ and R, T small enough, we have

$$\begin{aligned}
& \|u^n - \tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|v^n - \tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty} \\
& \lesssim \|u_0^n - \tilde{u}_0\|_{L^2} + \|v_0^n - \tilde{v}_0\|_{L^2} \\
& + \|Q_1(\tilde{u}, \tilde{v}, \phi) - Q_1(u^n, v^n, \phi)\|_{L^1 L^2} + \|Q_2(\tilde{u}, \tilde{v}, \phi) - Q_2(u^n, v^n, \phi)\|_{L^1 L^2} \\
& \lesssim \|u_0^n - \tilde{u}_0\|_{L^2} + \|v_0^n - \tilde{v}_0\|_{L^2} \\
& + R(\|u^n - \tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|v^n - \tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty}).
\end{aligned}$$

For $R < \frac{1}{2}$ small enough, we have

$$\frac{1}{2}(\|u^n - \tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|v^n - \tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty}) \leq \|\tilde{u}_0 - u_0^n\|_{L^2(\mathbb{R})} + \|\tilde{v}_0 - v_0^n\|_{L^2(\mathbb{R})}.$$

Letting $n \rightarrow +\infty$ we obtain the desired result. \square

3.2.2. From the system to the equation. Now, we finish the proof of Theorem 1.3.

Proof of Theorem 1.3. Let $\phi \in X^4(\mathbb{R})$ be such that $\|\partial\phi\|_{L^2}$ is small enough. Let $u_0 \in \phi + H^1(\mathbb{R})$ be such that $\|u_0 - \phi\|_{H^1}$ is small enough. Set $v_0 = \partial u_0 + \frac{i}{2}|u_0|^2 u_0$, $\tilde{u}_0 = u_0 - \phi$ and $\tilde{v}_0 = v_0 - \frac{i}{2}|\phi|^2 \phi$. We have

$$\tilde{v}_0 = \partial \tilde{u}_0 + \frac{i}{2}(|\tilde{u}_0 + \phi|^2(\tilde{u}_0 + \phi) - |\phi|^2 \phi) + \partial \phi.$$

Furthermore, $\tilde{u}_0 \in H^1(\mathbb{R})$, $\tilde{v}_0 \in L^2(\mathbb{R})$ satisfy:

$$\|\tilde{u}_0\|_{L^2(\mathbb{R})} + \|\tilde{v}_0\|_{L^2(\mathbb{R})} \lesssim \|\tilde{u}_0\|_{H^1(\mathbb{R})} + \|\partial\phi\|_{L^2},$$

which is small enough by the assumption. By Proposition 3.3, there exist $T > 0$ and a unique solution $(\tilde{u}, \tilde{v}) \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$ of the system (3.3). Let $u_0^n \in H^3(\mathbb{R})$ satisfy $\|u_0^n - \tilde{u}_0\|_{H^1(\mathbb{R})} \rightarrow 0$ as $n \rightarrow +\infty$. Set

$$v_0^n = \partial u_0^n + \frac{i}{2}(|u_0^n + \phi|^2(u_0^n + \phi) - |\phi|^2 \phi) + \partial \phi.$$

Let (u^n, v^n) be the $H^2(\mathbb{R})$ solution of the system (3.3) obtained by Proposition 3.1 with data (u_0^n, v_0^n) . By Proposition 3.2 we have

$$v^n = \partial u^n + \frac{i}{2}(|u^n + \phi|^2(u^n + \phi) - |\phi|^2 \phi) + \partial \phi. \quad (3.14)$$

Furthermore,

$$\|u_0^n - \tilde{u}_0\|_{L^2(\mathbb{R})} + \|v_0^n - \tilde{v}_0\|_{L^2(\mathbb{R})} \rightarrow 0.$$

From the continuous dependence on the initial data obtained in Proposition 3.3, $(u^n, v^n), (\tilde{u}, \tilde{v})$ are solutions of the system (3.3) on $[-T, T]$ for n large enough, and

$$\|u^n - \tilde{u}\|_{L^\infty L^2 \cap L^4 L^\infty} + \|v^n - \tilde{v}\|_{L^\infty L^2 \cap L^4 L^\infty} \rightarrow 0$$

as $n \rightarrow \infty$. Letting $n \rightarrow \infty$ on the two sides of (3.14), we obtain for all $t \in [-T, T]$

$$\tilde{v} = \partial \tilde{u} + \frac{i}{2}(|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2 \phi) + \partial \phi, \quad (3.15)$$

which makes sense in $H^{-1}(\mathbb{R})$. From (3.15) we see that $\partial \tilde{u} \in C([-T, T], L^2(\mathbb{R}))$ and (3.15) makes sense in $L^2(\mathbb{R})$. Then $\tilde{u} \in C([-T, T], H^1(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$. By the Sobolev embedding of $H^1(\mathbb{R})$ in $L^\infty(\mathbb{R})$ we obtain that

$$\begin{aligned}
\| |\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2 \phi \|_{L^4 L^\infty} & \lesssim \| |\tilde{u}|^3 \|_{L^4 L^\infty} + \| |\tilde{u}| |\phi|^2 \|_{L^4 L^\infty} \\
& < \| \tilde{u} \|_{L^4 L^\infty} \| \tilde{u} \|_{L^\infty L^\infty} + \| \tilde{u} \|_{L^4 L^\infty} \| \phi \|_{L^\infty L^\infty}^2 < \infty.
\end{aligned}$$

Hence, $|\tilde{u} + \phi|^2(\tilde{u} + \phi) - |\phi|^2\phi \in L^4L^\infty$. From (3.15) we obtain that $\partial\tilde{u} \in L^4L^\infty$ which implies $\tilde{u} \in L^4([-T, T], W^{1,\infty}(\mathbb{R}))$. Set $u = \tilde{u} + \phi$, $v = \tilde{v} + \frac{i}{2}|\phi|^2\phi$, then $u - \phi \in C([-T, T], H^1(\mathbb{R})) \cap L^4([-T, T], W^{1,\infty}(\mathbb{R}))$ and $v - \frac{i}{2}|\phi|^2\phi \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$. Moreover,

$$v = \partial u + \frac{i}{2}|u|^2u.$$

Since (u, v) solves (3.2), we have

$$Lu = -iu^2\bar{v} + \frac{1}{2}|u|^4u = -iu^2\partial\bar{u}.$$

The existence of a solution of the equation (1.1) follows. To prove the uniqueness property, assume that $U \in C([-T, T], \phi + H^1(\mathbb{R})) \cap L^4([-T, T], \phi + W^{1,\infty}(\mathbb{R}))$ is another solution of the equation (1.1). Set $V = \partial U + \frac{i}{2}|U|^2U$ and $\tilde{U} = U - \phi$, $\tilde{V} = V - \frac{i}{2}|\phi|^2\phi$. Hence $\tilde{U} \in C([-T, T], H^1(\mathbb{R})) \cap L^4([-T, T], W^{1,\infty}(\mathbb{R}))$ and $\tilde{V} \in C([-T, T], L^2(\mathbb{R})) \cap L^4([-T, T], L^\infty(\mathbb{R}))$. Moreover, (\tilde{U}, \tilde{V}) is a solution of the system (3.3). By the uniqueness of solutions of (3.3), we obtain that $\tilde{U} = \tilde{u}$. Hence, $u = U$, which completes the proof. \square

4. Conservation of the mass, the energy and the momentum. In this section, we prove Theorem 1.4. Let $q_0 \in \mathbb{R}$ and $u \in q_0 + H^2(\mathbb{R})$ be a solution of (1.1). Let χ and χ_R be the functions defined as in (1.5) and (1.7). We have

$$\|\chi'_R\|_{L^2(\mathbb{R})} = \left(\int_{\mathbb{R}} \frac{1}{R^2} \left(\chi' \left(\frac{x}{R} \right) \right)^2 dx \right)^{\frac{1}{2}} = \frac{1}{R^{\frac{1}{2}}} \|\chi'\|_{L^2(\mathbb{R})} \rightarrow 0 \text{ as } R \rightarrow \infty. \quad (4.1)$$

Similarly, for each $a \in \mathbb{R}$, we have

$$\|\chi'_{a,R}\|_{L^2(\mathbb{R})} \rightarrow 0 \text{ as } R \rightarrow \infty. \quad (4.2)$$

By the continuous dependence on initial data property of solution, we can assume that $u_0 \in q_0 + H^3(\mathbb{R})$, so that

$$u \in C((T_{min}, T_{max}), q_0 + H^3(\mathbb{R})).$$

It is enough to prove conservation of generalized mass, conservation of energy (1.8) and conservation of momentum (1.9) for any closed interval $[T_0, T_1] \in (T_{min}, T_{max})$. Let $T_0 < 0$, $T_1 > 0$ be such that $[T_0, T_1] \subset (T_{min}, T_{max})$. Let $M > 0$ be defined by

$$M = \sup_{t \in [T_0, T_1]} \|u - q_0\|_{H^3(\mathbb{R})}.$$

4.1. Conservation of mass. Multiplying both sides of (1.1) by \bar{u} and taking imaginary part to obtain

$$\mathcal{R}e(u_t\bar{u}) + \mathcal{I}m(\partial^2 u\bar{u}) + \mathcal{R}e(|u|^2 u \partial\bar{u}) = 0.$$

This implies that

$$\begin{aligned} 0 &= \frac{1}{2} \partial_t(|u|^2) + \partial(\mathcal{I}m(\partial u\bar{u})) + \frac{1}{4} \partial(|u|^4) \\ &= \frac{1}{2} \partial_t(|u|^2 - q_0^2) + \partial(\mathcal{I}m(\partial u\bar{u})) + \frac{1}{4} \partial(|u|^4 - q_0^4). \end{aligned}$$

By multiplying both sides by χ_R , integrating on space, and integrating by part we have

$$\begin{aligned} 0 &= \partial_t \int_{\mathbb{R}} \frac{1}{2}(|u|^2 - q_0^2)\chi_R dx - \int_{\mathbb{R}} \mathcal{I}m(\partial u \bar{u})\chi'_R - \int_{\mathbb{R}} \frac{(|u|^4 - q_0^4)}{4}\chi'_R dx \\ &= \partial_t \int_{\mathbb{R}} \frac{1}{2}(|u|^2 - q_0^2)\chi_R dx - \int_{\mathbb{R}} \left(\mathcal{I}m(\partial u \bar{u}) + \frac{1}{4}(|u|^4 - q_0^4) \right) \chi'_R dx. \end{aligned} \quad (4.3)$$

Denote the second term of (4.3) by K , using (4.1), we have

$$|K| \leq \| \mathcal{I}m(\partial u \bar{u}) + \frac{1}{4}(|u|^4 - q_0^4) \|_{L^2} \| \chi'_R \|_{L^2} \lesssim C(M) \frac{1}{R^{\frac{1}{2}}} \rightarrow 0 \text{ as } R \rightarrow \infty.$$

Thus, by integrating from 0 to t and taking R to infinity we obtain

$$\lim_{R \rightarrow \infty} \left(\int_{\mathbb{R}} \frac{1}{2}(|u|^2 - q_0^2)\chi_R dx - \int_{\mathbb{R}} \frac{1}{2}(|u_0|^2 - q_0^2)\chi_R dx \right) = 0. \quad (4.4)$$

Similarly, for each $a \in \mathbb{R}$, we have

$$\lim_{R \rightarrow \infty} \left(\int_{\mathbb{R}} \frac{1}{2}(|u|^2 - q_0^2)\chi_{a,R} dx - \int_{\mathbb{R}} \frac{1}{2}(|u_0|^2 - q_0^2)\chi_{a,R} dx \right) = 0. \quad (4.5)$$

as $R \rightarrow \infty$. This implies that $m^+(u(t))$ and $m^-(u(t))$ are conserved in time. In particular, if $m^+(u_0) = m^-(u_0) = m(u_0)$ then $m^+(u(t)) = m^-(u(t)) = m(u(t)) = m(u_0)$. This completes the proof of conservation of mass.

4.2. Conservation of energy. Now, we prove the conservation of the energy. Since u solves (1.1), after elementary calculations, we have

$$\begin{aligned} \partial_t(|\partial u|^2) &= \partial(2\mathcal{R}e(\partial u \partial_t \bar{u}) + \mathcal{R}e(u^2(\partial \bar{u})^2) - |\partial u|^2|u|^2 - |u|^4\mathcal{I}m(\bar{u}\partial u)) \\ &\quad + |u|^4\partial\mathcal{I}m(\bar{u}\partial u) + 2\mathcal{I}m(|u|^2\partial \bar{u} u_t). \end{aligned} \quad (4.6)$$

Recall that we have

$$\partial\mathcal{I}m(\partial u \bar{u}) = -\frac{1}{2}\partial_t(|u|^2) - \frac{1}{4}\partial(|u|^4). \quad (4.7)$$

Furthermore,

$$\partial_t\mathcal{I}m(|u|^2 u \partial \bar{u}) = 4\mathcal{I}m(u_t |u|^2 \partial \bar{u}) + \partial\mathcal{I}m(|u|^2 u \partial_t \bar{u}).$$

Thus,

$$2\mathcal{I}m(|u|^2 u_t \partial \bar{u}) = \frac{1}{2}(\partial_t\mathcal{I}m(|u|^2 u \partial \bar{u}) - \partial\mathcal{I}m(|u|^2 u \partial_t \bar{u})). \quad (4.8)$$

Combining (4.6), (4.7) and (4.8) we obtain

$$\begin{aligned} &\partial_t(|\partial u|^2) \\ &= \partial \left(2\mathcal{R}e(\partial u \partial_t \bar{u}) + \mathcal{R}e(u^2(\partial \bar{u})^2) - |\partial u|^2|u|^2 - |u|^4\mathcal{I}m(\partial u \bar{u}) - \frac{1}{2}\mathcal{I}m(|u|^2 u \partial_t \bar{u}) \right) \\ &\quad + \frac{1}{2}\partial_t\mathcal{I}m(|u|^2 u \partial \bar{u}) - \frac{1}{8}\partial(|u|^8) - \frac{1}{6}\partial_t(|u|^6). \end{aligned}$$

Hence,

$$\begin{aligned} &\partial_t \left(|\partial u|^2 - \frac{1}{2}\mathcal{I}m(|u|^2 u - q_0^3)\partial \bar{u} + \frac{1}{6}(|u|^6 - q_0^6) \right) \\ &= \partial(2\mathcal{R}e(\partial u \partial_t \bar{u}) + \mathcal{R}e(u^2(\partial \bar{u})^2) - |\partial u|^2|u|^2 - |u|^4\mathcal{I}m(\partial u \bar{u})) \end{aligned}$$

$$-\frac{1}{2}\mathcal{I}m(|u|^2u\partial_t\bar{u}) - \frac{1}{8}(|u|^8 - q_0^8) + \frac{1}{2}q_0^3\mathcal{I}m\partial_t\overline{(u - q_0)}.$$

Multiplying both sides by χ_R , integrating in space and integrating by part we obtain

$$\begin{aligned} & \partial_t \int_{\mathbb{R}} \left(|\partial u|^2 - \frac{1}{2}\mathcal{I}m(|u|^2u - q_0^3)\partial\bar{u} + \frac{1}{6}(|u|^6 - q_0^6) \right) \chi_R dx \\ &= - \int_{\mathbb{R}} \chi'_R \left(2\mathcal{R}e(\partial u\partial_t\bar{u}) + \mathcal{R}e(u^2(\partial\bar{u})^2) - |u|^2|\partial u|^2 - |u|^4\mathcal{I}m(\partial u\bar{u}) \right. \\ & \quad \left. - \frac{1}{2}\mathcal{I}m(|u|^2u\partial_t\bar{u}) - \frac{1}{8}(|u|^8 - q_0^8) \right) dx \\ & \quad - \frac{q_0^3}{2}\mathcal{I}m\partial_t \int_{\mathbb{R}} \overline{(u - q_0)}\chi'_R dx. \end{aligned}$$

Integrating from 0 to t we obtain

$$\int_{\mathbb{R}} \left(|\partial u|^2 - \frac{1}{2}\mathcal{I}m(|u|^2u - q_0^3)\partial\bar{u} + \frac{1}{6}(|u|^6 - q_0^6) \right) \chi_R dx \quad (4.9)$$

$$- \int_{\mathbb{R}} \left(|\partial u_0|^2 - \frac{1}{2}\mathcal{I}m(|u_0|^2u_0 - q_0^3)\partial\bar{u}_0 + \frac{1}{6}(|u_0|^6 - q_0^6) \right) \chi_R dx \quad (4.10)$$

$$\begin{aligned} &= \int_0^t \int_{\mathbb{R}} \chi'_R \left(2\mathcal{R}e(\partial u\partial_t\bar{u}) + \mathcal{R}e(u^2(\partial\bar{u})^2) - |u|^2|\partial u|^2 - |u|^4\mathcal{I}m(\partial u\bar{u}) \right. \\ & \quad \left. - \frac{1}{2}\mathcal{I}m(|u|^2u\partial_t\bar{u}) - \frac{1}{8}(|u|^8 - q_0^8) \right) dx ds \quad (4.11) \end{aligned}$$

$$- \frac{q_0^3}{2} \left(\mathcal{I}m \int_{\mathbb{R}} \overline{(u - q_0)}\chi'_R dx - \mathcal{I}m \int_{\mathbb{R}} \overline{(u_0 - q_0)}\chi'_R dx \right) . \quad (4.12)$$

Denoting the term (4.11) by A_R , using (4.1), we have

$$\begin{aligned} |A_R| &\leq \|\chi'_R\|_{L^2} \|2\mathcal{R}e(\partial u\partial_t\bar{u}) + \mathcal{R}e(u^2(\partial\bar{u})^2) - |u|^2|\partial u|^2 - |u|^4\mathcal{I}m(\partial u\bar{u}) \\ & \quad - \frac{1}{2}\mathcal{I}m(|u|^2u\partial_t\bar{u}) - \frac{1}{8}(|u|^8 - q_0^8)\|_{L^2} \\ &\lesssim C(M)\|\chi'_R\|_{L^2} \rightarrow 0 \text{ as } R \rightarrow \infty. \end{aligned} \quad (4.13)$$

Moreover, using (4.1) again, we have

$$\left| \mathcal{I}m \int_{\mathbb{R}} \overline{(u - q_0)}\chi'_R dx \right| \leq \|u - q_0\|_{L^2} \|\chi'_R\|_{L^2} \lesssim C(M)\|\chi'_R\|_{L^2} \rightarrow 0 \text{ as } R \rightarrow \infty. \quad (4.14)$$

$$\left| \mathcal{I}m \int_{\mathbb{R}} \overline{(u_0 - q_0)}\chi'_R dx \right| \leq \|u_0 - q_0\|_{L^2} \|\chi'_R\|_{L^2} \lesssim C(M)\|\chi'_R\|_{L^2} \rightarrow 0 \text{ as } R \rightarrow \infty. \quad (4.15)$$

To deal with the term (4.9), we need to divide it into two terms. First, using $u \in q_0 + H^3(\mathbb{R})$, as $R \rightarrow \infty$, we have

$$\int_{\mathbb{R}} \left(|\partial u|^2 - \frac{1}{2}\mathcal{I}m(|u|^2u - q_0^3)\partial\bar{u} \right) \chi_R dx \rightarrow \int_{\mathbb{R}} \left(|\partial u|^2 - \frac{1}{2}\mathcal{I}m(|u|^2u - q_0^3)\partial\bar{u} \right) dx. \quad (4.16)$$

Second, by easy calculations, we have

$$\frac{1}{6} \int_{\mathbb{R}} (|u|^6 - q_0^6)\chi_R - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^6 - q_0^6)\chi_R dx$$

$$\begin{aligned}
&= \frac{1}{6} \int_{\mathbb{R}} [(|u|^2 - q_0^2)(|u|^4 + q_0^2|u|^2 - 2q_0^4) + 3q_0^4(|u|^2 - q_0^2)] \chi_R dx \\
&\quad - \frac{1}{6} \int_{\mathbb{R}} [(|u_0|^2 - q_0^2)(|u_0|^4 + q_0^2|u_0|^2 - 2q_0^4) + 3q_0^4(|u_0|^2 - q_0^2)] \chi_R dx \\
&= \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - q_0^2)^2(|u|^2 + 2q_0^2) \chi_R dx - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^2 - q_0^2)^2(|u_0|^2 + 2q_0^2) \chi_R dx \quad (4.17)
\end{aligned}$$

$$+ \frac{q_0^4}{2} \int_{\mathbb{R}} (|u|^2 - q_0^2) \chi_R dx - \frac{q_0^4}{2} \int_{\mathbb{R}} (|u_0|^2 - q_0^2) \chi_R dx. \quad (4.18)$$

Denote the term (4.17) by B_R , we have

$$\begin{aligned}
B_R &\rightarrow \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - q_0^2)(|u|^4 + q_0^2|u|^2 - 2q_0^4) dx \\
&\quad - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^2 - q_0^2)(|u_0|^4 + q_0^2|u_0|^2 - 2q_0^4) dx \text{ as } R \rightarrow +\infty. \quad (4.19)
\end{aligned}$$

The term (4.18) converges to 0 as $R \rightarrow \infty$ by (4.4). Finally, we have

$$\begin{aligned}
&\lim_{R \rightarrow \infty} \left(\frac{1}{6} \int_{\mathbb{R}} (|u|^6 - q_0^6) \chi_R dx - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^6 - q_0^6) \chi_R dx \right) \\
&= \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - q_0^2)^2(|u|^2 + 2q_0^2) dx - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^2 - q_0^2)^2(|u_0|^2 + 2q_0^2) dx. \quad (4.20)
\end{aligned}$$

Combining (4.20) and (4.16) we have

$$\begin{aligned}
&\lim_{R \rightarrow \infty} (\text{the term (4.9)} - \text{the term (4.10)}) \\
&= \int_{\mathbb{R}} |\partial u|^2 - \frac{1}{2} \mathcal{I}m(|u|^2 u - q_0^3) \partial \bar{u}) dx + \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - q_0^2)^2(|u|^2 + 2q_0^2) dx \\
&\quad - \int_{\mathbb{R}} |\partial u_0|^2 - \frac{1}{2} \mathcal{I}m(|u_0|^2 u_0 - q_0^3) \partial \bar{u}_0) dx - \frac{1}{6} \int_{\mathbb{R}} (|u_0|^2 - q_0^2)^2(|u_0|^2 + 2q_0^2) dx \quad (4.21)
\end{aligned}$$

Combining (4.9)-(4.15), (4.21), we have

$$\begin{aligned}
&\int_{\mathbb{R}} |\partial u|^2 - \frac{1}{2} \mathcal{I}m(|u|^2 u - q_0^3) \partial \bar{u}) dx + \frac{1}{6} \int_{\mathbb{R}} (|u|^2 - q_0^2)^2(|u|^2 + 2q_0^2) dx \\
&= \int_{\mathbb{R}} |\partial u_0|^2 - \frac{1}{2} \mathcal{I}m(|u_0|^2 u_0 - q_0^3) \partial \bar{u}_0) dx + \frac{1}{6} \int_{\mathbb{R}} (|u_0|^2 - q_0^2)^2(|u_0|^2 + 2q_0^2) dx.
\end{aligned}$$

This implies (1.8).

4.3. Conservation of momentum. Now, we prove (1.9). Multiplying both sides of (1.1) by $-\partial \bar{u}$ and taking real part we obtain

$$\begin{aligned}
0 &= -\mathcal{R}e(iu_t \partial \bar{u} + \partial^2 u \partial \bar{u} + iu^2 (\partial \bar{u})^2) \\
&= \mathcal{I}m(u_t \partial \bar{u}) + \mathcal{I}m(u^2 (\partial \bar{u})^2) - \frac{1}{2} \partial (|\partial u|^2). \quad (4.22)
\end{aligned}$$

Moreover, by elementary calculation, we have

$$\partial_t \mathcal{I}m(u \partial \bar{u}) = 2\mathcal{I}m(u_t \partial \bar{u}) + \partial \mathcal{I}m(u \partial_t \bar{u}).$$

Replacing $\mathcal{I}m(u_t \partial \bar{u}) = \frac{1}{2} (\partial_t \mathcal{I}m(u \partial \bar{u}) - \partial \mathcal{I}m(u \partial_t \bar{u}))$ in (4.22), we obtain that

$$0 = \left(\frac{1}{2} \partial_t \mathcal{I}m(u \partial \bar{u}) - \frac{1}{2} \partial \mathcal{I}m(u \partial_t \bar{u}) \right) + 2\mathcal{R}e(u \partial \bar{u}) \mathcal{I}m(u \partial \bar{u}) - \frac{1}{2} \partial (|\partial u|^2)$$

$$= \partial_t \left[\frac{1}{2} \mathcal{I}m(u \partial \bar{u}) - \frac{1}{4} (|u|^4 - q_0^4) \right] + \partial \left[\mathcal{I}m(|u|^2 u \partial \bar{u}) - \frac{1}{2} |\partial u|^2 - \frac{1}{6} (|u|^6 - q_0^6) \right].$$

Multiply both sides by χ_R , integrating on space and integrating by part, we have

$$\begin{aligned} 0 &= \partial_t \int_{\mathbb{R}} \left[\frac{1}{2} \mathcal{I}m(u \partial \bar{u}) - \frac{1}{4} (|u|^4 - q_0^4) \right] \chi_R dx \\ &\quad - \int_{\mathbb{R}} \left[\mathcal{I}m(|u|^2 u \partial \bar{u}) - \frac{1}{2} |\partial u|^2 - \frac{1}{6} (|u|^6 - q_0^6) \right] \chi'_R dx \\ &= \partial_t \int_{\mathbb{R}} \left[\frac{1}{2} \mathcal{I}m(u \partial \bar{u}) - \frac{1}{4} (|u|^2 - q_0^2)^2 - \frac{1}{2} q_0^2 (|u|^2 - q_0^2) \right] \chi_R dx \\ &\quad - \int_{\mathbb{R}} \left[\mathcal{I}m(|u|^2 u \partial \bar{u}) - \frac{1}{2} |\partial u|^2 - \frac{1}{6} (|u|^6 - q_0^6) \right] \chi'_R dx. \end{aligned} \quad (4.23)$$

Denoting the second term of (4.23) by D_R , we have

$$\begin{aligned} |D_R| &\leq \left\| \mathcal{I}m(|u|^2 u \partial \bar{u}) - \frac{1}{2} |\partial u|^2 - \frac{1}{6} (|u|^6 - q_0^6) \right\|_{L^2} \|\chi'_R\|_{L^2} \\ &\lesssim C(M) \|\chi'_R\|_{L^2} \rightarrow 0 \text{ as } R \rightarrow \infty. \end{aligned} \quad (4.24)$$

Integrating from 0 to t the two sides of (4.23) and taking R to infinity, using (4.24) and (4.4), we have

$$\int_{\mathbb{R}} \left[\frac{1}{2} \mathcal{I}m(u \partial \bar{u}) - \frac{1}{4} (|u|^2 - q_0^2)^2 \right] dx = \int_{\mathbb{R}} \left[\frac{1}{2} \mathcal{I}m(u_0 \partial \bar{u}_0) - \frac{1}{4} (|u_0|^2 - q_0^2)^2 \right] dx. \quad (4.25)$$

We thus obtain the conservation of momentum, which completes the proof of Theorem 1.4.

5. Stationary solutions. In this section, we give the proof of Theorem 1.7. To convenience for readers, we first introduce a fundamental lemma which is a classical version of the Cauchy-Lipschitz theorem:

Lemma 5.1. *Let $C_1, C_2 \in \mathbb{R}$ and $f : \mathbb{R} \rightarrow \mathbb{R}$ be a C^1 function. There exists a unique real valued C^2 local solution of following equation*

$$\begin{cases} u_{xx} = f(u), \\ u(0) = C_1, \\ u_x(0) = C_2. \end{cases} \quad (5.1)$$

Remark 5.2. Let $C_1, C_2 \in \mathbb{C}$ and f be considered as C^1 function from \mathbb{R}^2 to \mathbb{R}^2 . By using Picard's uniqueness and existence theorem for system equations, we obtain the existence and uniqueness of complex valued solution for (5.1). However, the Lemma 5.1 is sufficient for our analysis in this paper.

Now, we give the proof of Theorem 1.7. We use the similar of variable changing as in [14, Proof of Proposition 1.1].

Proof of Theorem 1.7. Let ϕ be a nonconstant solution of (1.10) such that $m = \inf_{x \in \mathbb{R}} |\phi(x)| > 0$. From (1.10), we have $\phi \in X^3(\mathbb{R})$. Using the assumptions on ϕ we can write ϕ as

$$\phi(x) = R(x) e^{i\theta(x)}$$

where $R > 0$ and $R, \theta \in \mathcal{C}^2(\mathbb{R})$ are real-valued functions. We have

$$\begin{aligned}\phi_x &= e^{i\theta}(R_x + i\theta_x R), \\ \phi_{xx} &= e^{i\theta}(R_{xx} + 2iR_x\theta_x + iR\theta_{xx} - R\theta_x^2).\end{aligned}$$

Hence, since ϕ satisfies (1.10) we obtain

$$0 = (R_{xx} - R\theta_x^2 + R^3\theta_x) + i(2R_x\theta_x + R\theta_{xx} + R^2R_x).$$

This is equivalent to

$$0 = R_{xx} - R\theta_x^2 + R^3\theta_x, \quad (5.2)$$

$$0 = 2R_x\theta_x + R\theta_{xx} + R^2R_x. \quad (5.3)$$

The equation (5.3) is equivalent to

$$0 = \partial_x \left(R^2\theta_x + \frac{1}{4}R^4 \right).$$

Hence there exists $B \in \mathbb{R}$ such that

$$B = R^2\theta_x + \frac{1}{4}R^4. \quad (5.4)$$

This implies

$$\theta_x = \frac{B}{R^2} - \frac{R^2}{4}. \quad (5.5)$$

Substituting the above equality in (5.2) we obtain

$$\begin{aligned}0 &= R_{xx} - R \left(\frac{B}{R^2} - \frac{R^2}{4} \right)^2 + R^3 \left(\frac{B}{R^2} - \frac{R^2}{4} \right) \\ &= R_{xx} - \frac{B^2}{R^3} - \frac{5R^5}{16} + \frac{3BR}{2}.\end{aligned} \quad (5.6)$$

We prove that the set $V = \{x \in \mathbb{R} : R_x(x) \neq 0\}$ is dense in \mathbb{R} . Indeed, assume there exists $x \in \mathbb{R} \setminus \bar{V}$. Thus, there exists ε such that $B(x, \varepsilon) \in \mathbb{R} \setminus \bar{V}$. It implies that for all $y \in B(x, \varepsilon)$, we have $R_x(y) = 0$ so $R \equiv C_0$ on $B(x, \varepsilon)$ for some constants C_0 . Let $x_0 \in B(x, \varepsilon)$ then $R(x_0) = C_0$ and $R_x(x_0) = 0$. By Lemma 5.1, $R \equiv C_0$. By (5.5), θ_x is constant. Thus, $\phi(x)$ is of form $Ce^{i\alpha x}$, for some constants $C, \alpha \in \mathbb{R}$. If $\alpha = 0$, ϕ is a constant and if $\alpha \neq 0$ ϕ is not in $X^1(\mathbb{R})$, which contradicts the assumption of ϕ . From (5.6), we have

$$\begin{aligned}0 &= R_x \left(R_{xx} - \frac{B^2}{R^3} - \frac{5R^5}{16} + \frac{3BR}{2} \right) \\ &= \frac{d}{dx} \left[\frac{1}{2}R_x^2 + \frac{B^2}{2R^2} - \frac{5}{96}R^6 + \frac{3B}{4}R^2 \right].\end{aligned}$$

Hence there exists $a \in \mathbb{R}$ such that

$$a = \frac{1}{2}R_x^2 + \frac{B^2}{2R^2} - \frac{5}{96}R^6 + \frac{3B}{4}R^2.$$

This is equivalent to

$$\begin{aligned}0 &= R_x^2 R^2 + B^2 - \frac{5}{48}R^8 + \frac{3B}{2}R^4 - 2aR^2 \\ &= \frac{1}{4}[(R^2)_x]^2 + B^2 - \frac{5}{48}R^8 + \frac{3B}{2}R^4 - 2aR^2.\end{aligned}$$

Set $k = R^2$. We have

$$0 = \frac{1}{4}k_x^2 + B^2 - \frac{5}{48}k^4 + \frac{3B}{2}k^2 - 2ak. \quad (5.7)$$

Differentiating the two sides of (5.7) we have

$$0 = k_x \left(\frac{k_{xx}}{2} - \frac{5}{12}k^3 + 3Bk - 2a \right)$$

On the other hand, since $k_x = 2R_x R \neq 0$ for a.e x in \mathbb{R} , we obtain the following equation for a.e x in \mathbb{R} , hence, by continuity of k , it is true for all x in \mathbb{R} :

$$0 = \frac{k_{xx}}{2} - \frac{5}{12}k^3 + 3Bk - 2a. \quad (5.8)$$

Now, using Lemma 5.3 we have $k - 2\sqrt{B} \in H^3(\mathbb{R})$. Combining with (5.8) we obtain $a = \frac{4B\sqrt{B}}{3}$. Set $h = k - 2\sqrt{B}$. Then from (5.8) $h \in H^3(\mathbb{R})$ solves

$$\begin{cases} 0 = h_{xx} - \frac{5}{6}h^3 - 5\sqrt{B}h^2 - 4Bh, \\ h > -2\sqrt{B}, \end{cases} \quad (5.9)$$

Since $h \in H^3(\mathbb{R})$, there exists $x_0 \in \mathbb{R}$ such that $h_x(x_0) = 0$. Indeed, if h_x does not change sign on \mathbb{R} then $|h(-\infty)| > 0$ or $|h(\infty)| > 0$. This contradicts to $h \in H^3(\mathbb{R})$. Multiplying both sides of (5.9) by h_x we obtain

$$0 = \frac{1}{2}\partial_x(h_x^2) - \frac{5}{24}\partial_x(h^4) - \frac{5\sqrt{B}}{3}\partial_x(h^3) - 2B\partial_x(h^2).$$

Since $h \in H^3(\mathbb{R})$ we have $h(\infty) = h_x(\infty) = 0$ and hence,

$$\frac{1}{2}(h_x)^2 = \frac{5}{24}h^4 + \frac{5\sqrt{B}}{3}h^3 + 2Bh^2. \quad (5.10)$$

Using $h_x(x_0) = 0$, since (5.10), we have $h(x_0) = 0$ or $h(x_0) = \frac{4}{5}(-5 \pm \sqrt{10})\sqrt{B}$. If $h(x_0) = 0$ then by using Lemma 5.1, we have $h \equiv 0$, this is a contradiction. Since $h > -2\sqrt{B}$, we obtain $h(x_0) = \frac{4}{5}(-5 + \sqrt{10})\sqrt{B}$. Define $v(x) = h(x + x_0)$. We have

$$\begin{cases} 0 = v_{xx} - \frac{5}{6}v^3 - 5\sqrt{B}v^2 - 4Bv, \\ v(0) = \frac{4}{5}(-5 + \sqrt{10})\sqrt{B}, \\ v_x(0) = 0. \end{cases} \quad (5.11)$$

Using Lemma 5.1, there exists a unique solution v of (5.11). Moreover, we can check that the following function is a solution of (5.11):

$$v(x) = \frac{-1}{\sqrt{\frac{5}{72B} \cosh(2\sqrt{B}x) + \frac{5}{12\sqrt{B}}}}.$$

Hence,

$$h(x) = \frac{-1}{\sqrt{\frac{5}{72B} \cosh(2\sqrt{B}(x - x_0)) + \frac{5}{12\sqrt{B}}}}$$

This implies

$$k = 2\sqrt{B} + h = 2\sqrt{B} + \frac{-1}{\sqrt{\frac{5}{72B} \cosh(2\sqrt{B}(x - x_0)) + \frac{5}{12\sqrt{B}}}}.$$

Furthermore, using $\theta_x = \frac{B}{k} - \frac{k}{4}$, there exists $\theta_0 \in \mathbb{R}$ such that

$$\theta(x) = \theta_0 - \int_x^\infty \left(\frac{B}{k} - \frac{k}{4} \right) dy.$$

Now, assume that ϕ is a solution of (1.10) such that $\phi(\infty) = 0$. We prove $\phi \equiv 0$ on \mathbb{R} . Multiplying both sides of (1.10) by $\bar{\phi}$ then taking the imaginary part we obtain

$$\partial_x \mathcal{I}m(\phi_x \bar{\phi}) + \frac{1}{4} \partial_x (|\phi|^4) = 0$$

On the other hand, $\phi(\infty) = \phi_x(\infty) = 0$ then on \mathbb{R} we have

$$\mathcal{I}m(\phi_x \bar{\phi}) + \frac{1}{4} |\phi|^4 = 0. \quad (5.12)$$

If there exists y_0 such that $\phi_x(y_0) = 0$ then from (5.12) we have $\phi(y_0) = 0$. By the uniqueness of Cauchy problem we obtain $\phi \equiv 0$ on \mathbb{R} . Otherwise, ϕ_x does not vanish on \mathbb{R} . From now on, we will consider this case. Multiplying both sides of (1.10) by $\bar{\phi}_x$ then taking the real part, we have

$$\begin{aligned} 0 &= \mathcal{R}e(\phi_{xx} \bar{\phi}_x) - \mathcal{I}m(\phi^2 \bar{\phi}_x^2) \\ &= \frac{1}{2} \frac{d}{dx} |\phi_x|^2 - 2 \mathcal{R}e(\phi \bar{\phi}_x) \mathcal{I}m(\phi \bar{\phi}_x) \\ &= \frac{1}{2} \frac{d}{dx} |\phi_x|^2 - \partial_x (|\phi|^2) \frac{1}{4} |\phi|^4 \\ &= \frac{d}{dx} \left(\frac{1}{2} |\phi_x|^2 - \frac{1}{12} |\phi|^6 \right). \end{aligned}$$

This implies that

$$|\phi_x|^2 - \frac{1}{6} |\phi|^6 = 0.$$

Hence, since ϕ_x is non vanishing, ϕ is also non vanishing on \mathbb{R} . We can write $\phi = \rho e^{i\theta}$ for $\rho > 0$, $\rho, \theta \in \mathcal{C}^2(\mathbb{R})$. Similar to (5.2) we have

$$0 = -\rho \theta_x^2 + \rho_{xx} + \rho^3 \theta_x. \quad (5.13)$$

Replacing $\phi = \rho e^{i\theta}$ in (5.12) we have

$$0 = \rho^2 \theta_x + \frac{1}{4} \rho^4.$$

Then $\theta_x = -\frac{1}{4} \rho^2$, replacing this equality in (5.13) we obtain

$$0 = \rho_{xx} - \frac{5}{16} \rho^5.$$

Multiplying both sides of the above equality by ρ_x we obtain

$$0 = \rho_{xx} \rho_x - \frac{5}{16} \rho^5 \rho_x = \frac{d}{dx} \left(\frac{1}{2} \rho_x^2 - \frac{5}{96} \rho^6 \right).$$

Hence,

$$0 = \rho_x^2 - \frac{5}{48} \rho^6.$$

Moreover, ϕ is non vanishing on \mathbb{R} then $\rho > 0$ and then ρ_x is not change sign on \mathbb{R} . If $\rho_x > 0$ then since $\rho(\infty) = 0$ we have $\rho < 0$ on \mathbb{R} , a contradiction. Hence, $\rho_x < 0$

and $\rho_x = -\sqrt{\frac{5}{48}}\rho^3$. From this we easily check that

$$\rho^2(x) = \frac{1}{\rho(0)^2 + \sqrt{5/12}x},$$

which implies the contradiction, for the right hand side is not a continuous function on \mathbb{R} . This completes the proof. \square

Lemma 5.3. *Let $B > 0$ be the constant given as the above. The following is true:*

$$k - 2\sqrt{B} \in L^2(\mathbb{R}), \quad k \in X^3(\mathbb{R}).$$

Proof. Using $\phi \in L^\infty(\mathbb{R})$ we obtain $k \in L^\infty(\mathbb{R})$. On the other hand, since $\phi \in X^3(\mathbb{R})$, we have $\phi_x \in L^2(\mathbb{R})$, $\phi_{xx} \in L^2(\mathbb{R})$ and it easy to see that

$$\begin{aligned} |\phi_x|^2 &= \frac{k_x^2}{4k} + k\theta_x^2 \in L^1(\mathbb{R}), \\ |\phi_{xx}|^2 &= \left| \frac{k_x\theta_x}{\sqrt{k}} + \theta_{xx}\sqrt{k} \right|^2 + \left| \frac{k_{xx}}{2\sqrt{k}} - \sqrt{k}\theta_x^2 - \frac{k_x^2}{4k\sqrt{k}} \right|^2 \in L^1(\mathbb{R}). \end{aligned}$$

This implies

$$\begin{aligned} \frac{k_x}{2\sqrt{k}} &\in L^2(\mathbb{R}) \text{ and } \sqrt{k}\theta_x \in L^2(\mathbb{R}) \\ \frac{k_x\theta_x}{\sqrt{k}} + \theta_{xx}\sqrt{k} &\in L^2(\mathbb{R}) \text{ and } \frac{k_{xx}}{2\sqrt{k}} - \sqrt{k}\theta_x^2 - \frac{k_x^2}{4k\sqrt{k}} \in L^2(\mathbb{R}). \end{aligned}$$

Using $\sqrt{m} < k < \|k\|_{L^\infty}$, $\theta_x = \frac{4B-k^2}{4k} \in L^\infty(\mathbb{R})$, $k_x = 2RR_x \in L^\infty$ (indeed $|\phi_x|^2 = |R_x|^2 + |R\theta_x|^2 \in L^\infty(\mathbb{R})$) we have

$$\begin{aligned} k_x &\in L^2 \text{ and } \theta_x \in L^2, \\ \theta_{xx} &\in L^2 \text{ and } k_{xx} \in L^2. \end{aligned}$$

By using $\theta_x = \frac{4B-k^2}{4k} \in L^2(\mathbb{R})$, we have $4B - k^2 \in L^2(\mathbb{R})$. Thus, $B \geq 0$ and $2\sqrt{B} - k \in L^2(\mathbb{R})$. If $B = 0$ then $k \in L^2(\mathbb{R})$, hence, $R \in L^2(\mathbb{R})$. Which contradicts to the assumption $m > 0$. Thus, $B > 0$. It remains to prove that $k_{xxx} \in L^2(\mathbb{R})$. Indeed, from $\phi_{xxx} \in L^2(\mathbb{R})$ we have

$$|\phi_{xxx}|^2 = |\theta_{xxx}\sqrt{k} + \mathcal{M}|^2 + \left| \frac{k_{xxx}}{2\sqrt{k}} + \mathcal{N} \right|^2 \in L^1(\mathbb{R}) \quad (5.14)$$

where \mathcal{M}, \mathcal{N} are functions of $\theta, \theta_x, \theta_{xx}, k, k_x, k_{xx}$. We can easily check that $\mathcal{M}, \mathcal{N} \in L^2(\mathbb{R})$. Hence, from (5.14) and the facts that $\theta_x \in H^1(\mathbb{R})$, $k \in X^2(\mathbb{R})$, k bounded from below we obtain $\theta_{xxx}, k_{xxx} \in L^2(\mathbb{R})$. This implies the desired results. \square

From now on, we will denote ϕ_B is the stationary solution of (1.10) given by Theorem 1.7 with $\theta_0 = 0$. We have

$$\phi_B = e^{i\theta}\sqrt{k}, \quad (5.15)$$

$$k(x) = 2\sqrt{B} + \frac{-1}{\sqrt{\frac{5}{72B}} \cosh(2\sqrt{B}x) + \frac{5}{12\sqrt{B}}}, \quad (5.16)$$

$$\theta(x) = -\int_x^\infty \frac{B}{k(y)} - \frac{k(y)}{4} dy. \quad (5.17)$$

We have the following asymptotic properties for ϕ_B at ∞ .

Proposition 5.4. *Let $B > 0$ and ϕ_B be kink solution of (1.1). Then for $x > 0$, we have*

$$|\phi_B - \sqrt{2\sqrt{B}}| \lesssim e^{-\sqrt{B}x}.$$

As consequence ϕ_B converges to $\sqrt{2\sqrt{B}}$ as x tends to ∞ and there exists limit of ϕ_B as x tends to $-\infty$.

Proof. Since (5.16) we have

$$|k - 2\sqrt{B}| \lesssim e^{-2\sqrt{B}x}.$$

Hence, for all $x \in \mathbb{R}$ we have

$$|\phi_B(x) - \sqrt{2\sqrt{B}}| \lesssim |e^{i\theta(x)}\sqrt{k(x)} - \sqrt{k(x)}| + |\sqrt{k(x)} - \sqrt{2\sqrt{B}}| \quad (5.18)$$

$$\lesssim \|k\|_{L^\infty}^{\frac{1}{2}} |e^{i\theta(x)} - 1| + e^{-\sqrt{B}x} \quad (5.19)$$

Moreover, for $x > 0$, we have

$$\begin{aligned} |e^{i\theta(x)} - 1| &\leq |\theta(x)| \leq \int_x^\infty \left| \frac{B}{k} - \frac{k}{4} \right| dx \\ &\leq \int_x^\infty \left| \frac{B}{k} - \frac{\sqrt{B}}{2} \right| + \left| \frac{\sqrt{B}}{2} - \frac{k}{4} \right| dx \\ &\lesssim \int_x^\infty |k - 2\sqrt{B}| dx \lesssim \int_x^\infty e^{-2\sqrt{B}x} dx \lesssim e^{-2\sqrt{B}x}. \end{aligned}$$

Combining with (5.19) we obtain

$$|\phi_B(x) - \sqrt{2\sqrt{B}}| \lesssim e^{-\sqrt{B}x}.$$

As consequence ϕ_B converges to $\sqrt{2\sqrt{B}}$ as x tends to ∞ . Since (5.16), we have $|k - 2\sqrt{B}| \in L^1(\mathbb{R})$ and $k > \left(2 - \frac{1}{\frac{5}{12} + \sqrt{\frac{5}{72}}}\right)\sqrt{B}$. Thus, $\frac{B}{k} - \frac{k}{4} = \frac{4B - k^2}{4k} \in L^1(\mathbb{R})$.

Hence since (5.17) we have

$$\lim_{x \rightarrow -\infty} \theta(x) = - \int_{-\infty}^\infty \left(\frac{B}{k(y)} - \frac{k(y)}{4} \right) dy.$$

Hence,

$$\lim_{x \rightarrow -\infty} \phi_B(x) = \exp \left(-i \int_{-\infty}^\infty \left(\frac{B}{k(y)} - \frac{k(y)}{4} \right) dy \right) \sqrt{2\sqrt{B}}.$$

This completes the proof. \square

Acknowledgments. The author wishes to thank Prof. Stefan Le Coz for his guidance and encouragement.

REFERENCES

- [1] F. Béthuel, P. Gravejat and D. Smets, [Stability in the energy space for chains of solitons of the one-dimensional Gross-Pitaevskii equation](#), *Ann. Inst. Fourier (Grenoble)*, **64** (2014), 19–70.
- [2] T. Cazenave, [Semilinear Schrödinger Equations](#), volume 10 of *Courant Lecture Notes in Mathematics*, New York University, Courant Institute of Mathematical Sciences, New York, 2003.

- [3] T. Cazenave and A. Haraux, *An Introduction to Semilinear Evolution Equations*, volume 13 of *Oxford Lecture Series in Mathematics and its Applications*, The Clarendon Press Oxford University Press, New York, 1998.
- [4] A. de Laire, [Global well-posedness for a nonlocal Gross-Pitaevskii equation with non-zero condition at infinity](#), *Comm. Partial Differential Equations*, **35** (2010), 2021–2058.
- [5] C. Gallo, Schrödinger group on Zhidkov spaces, *Adv. Differential Equations*, **9** (2004), 509–538.
- [6] C. Gallo, [The Cauchy problem for defocusing nonlinear Schrödinger equations with non-vanishing initial data at infinity](#), *Comm. Partial Differential Equations*, **33** (2008), 729–771.
- [7] P. Gérard, [The Cauchy problem for the Gross-Pitaevskii equation](#), *Ann. Inst. H. Poincaré Anal. Non Linéaire*, **23** (2006), 765–779.
- [8] P. Gérard, [The Gross-Pitaevskii equation in the energy space](#), In *Stationary and Time Dependent Gross-Pitaevskii Equations*, volume 473 of *Contemp. Math.*, Amer. Math. Soc., Providence, RI, 2008, pages 129–148.
- [9] M. Hayashi and T. Ozawa, [Well-posedness for a generalized derivative nonlinear Schrödinger equation](#), *J. Differential Equations*, **261** (2016), 5424–5445.
- [10] N. Hayashi and T. Ozawa, [On the derivative nonlinear Schrödinger equation](#), *Phys. D*, **55** (1992), 14–36.
- [11] N. Hayashi and T. Ozawa, [Finite energy solutions of nonlinear Schrödinger equations of derivative type](#), *SIAM J. Math. Anal.*, **25** (1994), 1488–1503.
- [12] S. Le Coz, Standing waves in nonlinear Schrödinger equations, In *Analytical and Numerical Aspects of Partial Differential Equations*, Walter de Gruyter, Berlin, 2009, pages 151–192.
- [13] E. Mjølhus, [On the modulational instability of hydromagnetic waves parallel to the magnetic field](#), *Journal of Plasma Physics*, **16** (1976), 321–334.
- [14] M. Murai, K. Sakamoto and S. Yotsutani, [Representation formula for traveling waves to a derivative nonlinear Schrödinger equation with the periodic boundary condition](#), *Discrete Contin. Dyn. Syst.*, (Dynamical systems, differential equations and applications. 10th AIMS Conference. Suppl.), 2015, 878–900.
- [15] C. Sulem and P.-L. Sulem, *The Nonlinear Schrödinger Equation*, volume 139 of *Applied Mathematical Sciences*, Springer-Verlag, New York, 1999. Self-focusing and wave collapse.
- [16] M. Tsutsumi and I. Fukuda, On solutions of the derivative nonlinear Schrödinger equation. Existence and uniqueness theorem, *Funkcial. Ekvac.*, **23** (1980), 259–277.
- [17] M. Tsutsumi and I. Fukuda, On solutions of the derivative nonlinear Schrödinger equation. II, *Funkcial. Ekvac.*, **24** (1981), 85–94.
- [18] P. E. Zhidkov, *Korteweg-de Vries and Nonlinear Schrödinger Equations: Qualitative Theory*, volume 1756 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 2001.

Received January 2021; 1st revision April 2021; final revision April 2021.

E-mail address: van-tin.phan@univ-tlse3.fr