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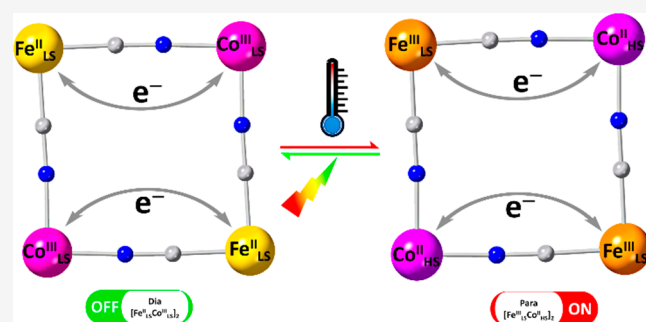
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# Two-Step Thermoinduced Metal-to-Metal Electron Transfer and ON/OFF Photoswitching in a Molecular $[\text{Fe}_2\text{Co}_2]$ Square Complex

Sujit Kamilya, Subrata Ghosh, Yanling Li, Pierre Dechambenoit, Mathieu Rouzières, Rodrigue Lescouëzec, Sakshi Mehta, and Abhishake Mondal\*

**ABSTRACT:** A cyanide-bridged  $[\text{Fe}_2\text{Co}_2]$  molecular square complex,  $\{[\text{Fe}(\text{Tp})(\text{CN})_3]_2[\text{Co}(\text{L})_2]_2\}(\text{BF}_4)_2 \cdot 2\text{CH}_3\text{CN} \cdot 6\text{H}_2\text{O}$  [**1**; Tp = hydrotris(pyrazol-1-yl)borate and L = bis(1-ethyl-1H-imidazol-2-yl)ketone], has been synthesized and characterized fully by single-crystal X-ray diffraction, (photo)magnetic measurements, optical reflectivity, and other physical measurements. **1** exhibits a two-step metal-to-metal electron-transfer (MMET)-induced spin transition accompanied by thermal hysteresis ( $T_{1/2\uparrow} = 332$  and  $407$  K and  $T_{1/2\downarrow} = 320$  and  $405$  K, respectively), converting the low-temperature diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}$  ground state into the high-temperature paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}$  state. Additionally, **1** displays reversible photoinduced MMET under light irradiation (ON mode using  $808$  nm laser light and OFF mode using  $532$  nm laser light), as confirmed by optical reflectivity and (photo)magnetic measurements. The photoinduced paramagnetic metastable state relaxes back to the diamagnetic ground state at  $91$  K ( $T_{\text{LIESST}} = 91$  K). Astonishingly, **1** also exhibits a  $27$  K wide light-induced thermal hysteresis below  $100$  K. The overall results show that **1** is a multistimuli-responsive bistable material that exhibits reversible switching between the diamagnetic state,  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}$ , and the paramagnetic state,  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}$ , under the application of temperature and light.



## INTRODUCTION

Magnetic materials and switches have attracted great interest in the field of material sciences from the last 2 decades because of huge potential applications in future electronic devices.<sup>1</sup> Bistable molecules that exhibit ON/OFF switching of their magnetization with the application of various external perturbations, e.g., temperature,<sup>2</sup> pressure,<sup>3</sup> light,<sup>4</sup> electric field,<sup>5</sup> etc., are the best candidates for information storage,<sup>6</sup> switches,<sup>7</sup> and sensor devices<sup>8</sup> at a molecular level. So far, materials based on cyanide-bridged Prussian Blue (PB) and its analogues (PBAs) have shown their potential and also become an eminent example in this growing field with its exciting reversible switching of (photo)physical properties. It has been started with the revolutionary example of metal-to-metal electron transfer (MMET) occurring between iron and cobalt centers via a cyanide bridge in a  $[\text{FeCo}]$  PBA,  $\text{K}_{0.2}\text{Co}_{1.4}[\text{Fe}(\text{CN})_6] \cdot 6.9\text{H}_2\text{O}$ , where the low-temperature  $\{\text{Fe}^{\text{II}}_{\text{LS}}(S=0)\text{-CN-Co}^{\text{III}}_{\text{LS}}(S=0)\}$  diamagnetic ground state changes into the metastable  $\{\text{Fe}^{\text{III}}_{\text{LS}}(S=1/2)\text{-CN-Co}^{\text{II}}_{\text{HS}}(S=3/2)\}$  paramagnetic state via light irradiation at  $5$  K.<sup>9</sup> In 2004, Dunbar et al. synthesized the first-ever thermally induced MMET in a discrete cyanide-bridged  $[\text{Co}_3\text{Fe}_2]$  complex,<sup>10</sup> and later in 2011, the (photo)magnetic effect of this complex was studied by Clérac and Dunbar et al.<sup>11</sup> After that, various molecular model complexes of  $[\text{FeCo}]$  PBAs; e.g., molecular cubes,<sup>12,13</sup>

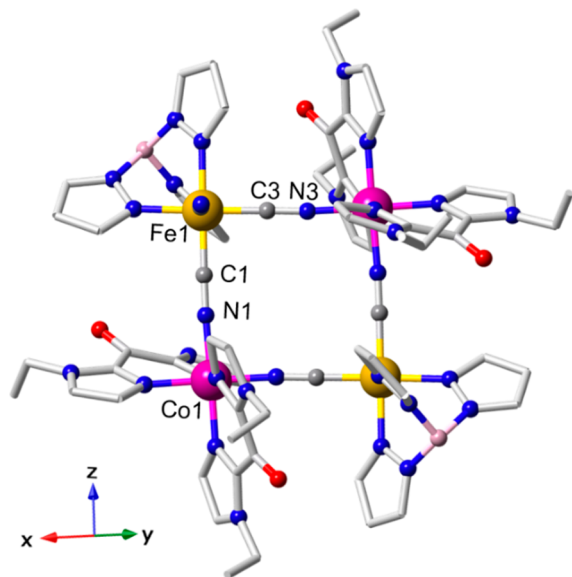
squares,<sup>14–21</sup> and dinuclear units<sup>22,23</sup> have been reported by different research groups, including us, which exhibit thermo- and/or photoinduced MMET.<sup>24,25</sup> Additionally, a few  $[\text{FeCo}]$  systems have also been investigated where the photoinduced metastable state exhibits a slow relaxation of magnetization, i.e., single molecule magnets,<sup>26,27</sup> while some 1D  $[\text{FeCo}]$  chains exhibit interesting single chain magnet properties in the photoinduced state.<sup>28–31</sup> All of these stimuli-responsive materials open a new horizon in the field of electronic device technology with their huge potential.<sup>32</sup> The most challenging task down the lane is to synthesize a multiresponsive material that will exhibit bistability at ambient conditions. To the best of our knowledge, until now there is no such example of the  $\{\text{Fe}_2\text{Co}_2\}$  molecular square complex that exhibits thermoinduced two-step MMET, accompanied by thermal hysteresis and ON/OFF photoswitching and the light-induced thermal hysteresis effect (LITH).

In this aspect, herein we present our recent works with a significant focus on the development and a better understanding of the role of the redox potential along with intra- and intermolecular interactions associated with the stimuli-responsive reversible electron-transfer properties of a  $\{\text{Fe}_2\text{Co}_2\}$  molecular square complex. Applying building block as a ligand strategy, we have used here the cyanide building block  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  [Tp = hydrotris(pyrazol-1-yl)borate] to prepare the molecular model complex of [FeCo] PBAs, i.e., the  $\{\text{Fe}_2\text{Co}_2\}$  square complex  $\{[\text{Fe}(\text{Tp})(\text{CN})_3]_2[\text{Co}(\text{L})_2]_2\}(\text{BF}_4)_2\cdot 2\text{CH}_3\text{CN}\cdot 6\text{H}_2\text{O}$  [**1**; L = bis(1-ethyl-1*H*-imidazol-2-yl)ketone] and explored the structural, (photo)magnetic, and optical properties during the MMET process.

## RESULTS AND DISCUSSION

**Synthesis.** The reaction of  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  with a solution of  $\text{Co}(\text{BF}_4)_2\cdot 6\text{H}_2\text{O}$  and the L ligand in a mixture of acetonitrile/water (9:1) afforded a green reaction mixture. Slow evaporation of the reaction mixture gave analytically pure green crystals of **1** in good yield. The purity of the complex has been studied by elemental analyses and powder X-ray diffraction (PXRD) measurements (Figure S1). The thermal stability of **1** was investigated by thermogravimetric analysis (TGA) in the temperature range of 300–600 K with a scan rate of 5 K  $\text{min}^{-1}$  under a nitrogen atmosphere (Figure S2). **1** first undergoes a weight loss between 300 and 353 K (ca. 2.7%), corresponding to a partial release of interstitial solvent molecules. After that, **1** shows thermal stability up to 450 K. The partial solvent loss is probably the reason for the slight deviation in the experimental PXRD data.

**Structure Analysis.** Single-crystal X-ray structure analysis was performed on suitable single crystal of **1** at 120 K (Figure 1). The complex crystallizes in the triclinic  $\overline{P}1$  space group ( $Z = 1$ ; Table S1) and consists of a  $\{\text{Fe}_2\text{Co}_2\}$  square structural motif where the iron and cobalt centers occupy an alternate corner of the square and are bridged by cyanide ions. Each iron



**Figure 1.** Perspective view of the cyanide-bridged  $\{\text{Fe}_2\text{Co}_2\}$  square unit of **1**. Solvent molecules, hydrogen atoms, and counteranions are omitted for clarity. Color code: Fe, yellow; Co, purple; C, gray; N, blue; O, red; B, light pink.

center is in a distorted-octahedral geometry with a  $\text{FeC}_3\text{N}_3$  coordination environment that can best be described by using the continuous shape measures (CSHM) program<sup>33</sup> (Table S2), where the iron center is coordinated to three nitrogen-donor atoms of the Tp ligand and three carbon-donor atoms from two bridging and one terminal cyanide ligands.

The coordination geometry around each cobalt center is best described as a slightly distorted octahedron with surrounding  $\text{CoN}_6$  (Table S2), in which the cobalt center is coordinated to four nitrogen-donor atoms of two L ligands and two nitrogen-donor atoms of two bridging cyanide ligands. The overall structure of **1** consists of a dicationic  $\{[\text{Fe}(\text{Tp})(\text{CN})_3]_2[\text{Co}(\text{L})_2]_2\}^{2+}$  square unit, two  $\text{BF}_4^-$  counteranions, and interstitial water and acetonitrile solvent molecules (Figure S3). Estimations of their numbers are 6 and 2, respectively, but their exact number must be taken with care because some of them had to be squeezed (see the Supporting Information).<sup>34</sup> In **1**, the Fe–C and Fe–N bond distances are in the ranges of 1.860(3)–1.905(4) and 1.994(3)–1.999(3) Å, respectively (Table S3), which are in agreement with a low-spin (LS) iron(II) state in a  $\text{FeC}_3\text{N}_3$  distorted octahedral geometry.<sup>12–20,24</sup>

The average Co–N bond distance is 1.912 Å, lying in the range expected for the LS cobalt(III) ion in a  $\text{CoN}_6$  distorted octahedral coordination environment, where typical Co–N bond distances of LS cobalt(III) and high-spin (HS) cobalt(II) are  $\sim 1.9$  and  $\sim 2.1$  Å, respectively.<sup>12–20,35–38</sup> The Fe–CN–Co edges are almost identical with distances of 4.892(1) and 4.888(1) Å, and angles at the corners are close to the orthogonality [ $\text{C1–Fe1–C3} = 88.71(14)^\circ$  and  $\text{N1–Co1–N3} = 91.73(12)^\circ$ ; Table S4]. The cyanide bridges are slightly more bent on the cobalt side [ $\text{Co1–N1–C1} = 178.0(3)^\circ$  and  $\text{Co1–N3–C3} = 174.9(3)^\circ$ ] than on the iron side, which remains close to linearity [ $\text{Fe1–C1–N1} = 178.3(3)^\circ$  and  $\text{Fe1–C3–N3} = 179.3(3)^\circ$ ]. All of these geometric features indicate the presence of a  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{–CN–Co}^{\text{III}}_{\text{LS}}\}_2$  square moiety in **1** at 120 K. The high-temperature single-crystal X-ray diffraction (SCXRD) data collection of **1** at 300 K or above has remained unsuccessful even after several attempts because the crystals lose crystallinity as a result of interstitial solvent loss. In **1**, the nonbridging cyanide groups are involved in hydrogen-bonding interaction with the interstitial water molecules (Figure S4). Weak supramolecular interactions (C–H $\cdots$ N and C–H $\cdots$ F) have been observed between the molecular square motif, acetonitrile solvent molecules, and  $\text{BF}_4^-$  counteranions in **1** (Figure S4). In addition, various C–H $\cdots$ N and C–H $\cdots$ C interactions between the square motifs form a supramolecular chain in **1** (Figure S5). The  $\{\text{Fe}_2\text{Co}_2\}$  square units in **1** are isolated quite well from each other, with the shortest intermolecular metal–metal distance of 8.759(1) Å (Fe $\cdots$ Co).

**Spectroscopic Studies.** IR spectroscopy has been considered to be an important tool in characterizing the electronic states and bridging modes of cyanide-bridged metal complexes because the cyanide stretching frequencies  $\nu_{\text{CN}}$  are sensitive to the oxidation states as well as the spin states of the cyanide-bridged metal complexes.<sup>12–20,24,25,37,39–41</sup> Also,  $\nu_{\text{BH}}$  and  $\nu_{\text{CO}}$  of the Tp and L ligands, respectively, are significantly sensitive to the electronic nature (both spin and oxidation states) of the coordinated metal ions. Therefore, IR measurements were carried out on the L ligand, building block  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$ , L-coordinated cobalt complex  $[\text{Co}(\text{L})_3]^{2+}$ , and complex **1** at 300 K (Figures S6–S8). The IR spectrum of  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  exhibits character-

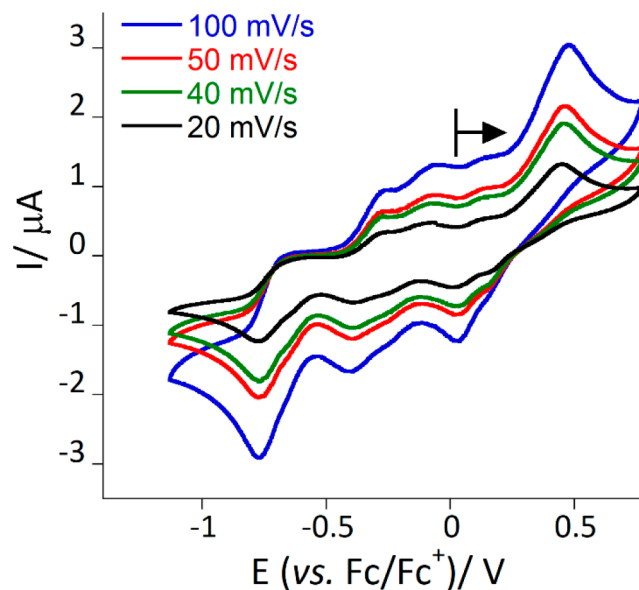
istic  $\nu_{\text{CN}}$  and  $\nu_{\text{BH}}$  bands at 2117 and 2510  $\text{cm}^{-1}$ , respectively (Figure S7). At 300 K, the IR spectrum of **1** shows three  $\nu_{\text{CN}}$  bands at 2127, 2112, and 2065  $\text{cm}^{-1}$  (Figure S8). The first two bands are attributed to the bridging  $\nu_{\text{CN}}$  band of the  $\{\text{Fe}^{\text{II}}_{\text{LS}}-\text{CN}-\text{Co}^{\text{III}}_{\text{LS}}\}$  unit, while the latter band can be ascribed to the  $\nu_{\text{CN}}$  band of terminal  $\text{Fe}^{\text{II}}_{\text{LS}}-\text{CN}$  of the  $[\text{Fe}^{\text{II}}(\text{Tp})(\text{CN})_3]^{2-}$  unit.<sup>12–20,22–25,37,39–41</sup> The characteristic  $\nu_{\text{BH}}$  band was observed at 2468  $\text{cm}^{-1}$ , which is significantly lower than that found in the  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  (2510  $\text{cm}^{-1}$ ) building block and very similar to that reported for the  $[\text{Fe}^{\text{II}}(\text{Tp})(\text{CN})_3]^{2-}$  unit, suggesting the presence of an iron(II) oxidation state.<sup>12–20,22–25,37,39–41</sup> The peak at 1663  $\text{cm}^{-1}$  corresponds to the  $\nu_{\text{CO}}$  band of the coordinated L ligand, which is significantly shifted to higher energy in comparison to the free ligand (1630  $\text{cm}^{-1}$ ) and HS  $[\text{Co}(\text{L})_3]^{2+}$  complex (1641  $\text{cm}^{-1}$ ).<sup>12</sup> Additionally, the presence of four characteristic peaks in the range of 1030–1115  $\text{cm}^{-1}$  indicates the presence of  $\text{BF}_4^-$  counteranions in **1**. All of these results suggest the presence of the  $\{\text{Fe}^{\text{II}}_{\text{LS}}-\text{CN}-\text{Co}^{\text{III}}_{\text{LS}}\}$  square moiety in **1** at 300 K.

Variable-temperature IR measurements were performed on **1** in the temperature range of 300–420 K in both heating and cooling modes with 10 K intervals (Figures S9–S12). Upon an increase in the temperature, the intensities of the three  $\nu_{\text{CN}}$  bands of **1** were decreased with the appearance of three new bands at 2169, 2163, and 2120  $\text{cm}^{-1}$  (Figure S10). The former two bands lie in the range expected for the bridging  $\nu_{\text{CN}}$  band of the  $\{\text{Fe}^{\text{III}}_{\text{LS}}-\text{CN}-\text{Co}^{\text{II}}_{\text{HS}}\}$  unit.<sup>14</sup> The latter value is ascribed to the  $\nu_{\text{CN}}$  band of free terminal  $\text{Fe}^{\text{III}}_{\text{LS}}-\text{CN}$  of the  $[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]^-$  unit.<sup>14</sup> In addition, a significant change has also been observed in the  $\nu_{\text{BH}}$  absorption band of the coordinated Tp ligand. Upon an increase in the temperature, the intensity of the  $\nu_{\text{BH}}$  band decreases with the appearance of a new band at 2510  $\text{cm}^{-1}$ , which is similar to the value observed in  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  (Figure S7). Interestingly, an important variation has also been noticed in the  $\nu_{\text{CO}}$  absorption band of the L ligand coordinated to cobalt centers. A decrease in the  $\nu_{\text{CO}}$  absorption band at 1667  $\text{cm}^{-1}$  is observed with the formation of a new peak at 1646  $\text{cm}^{-1}$ , which lies in the range observed in the HS  $[\text{Co}(\text{L})_3]^{2+}$  complex (1641  $\text{cm}^{-1}$ ; Figure S6). These changes in  $\nu_{\text{CN}}$ ,  $\nu_{\text{BH}}$ , and  $\nu_{\text{CO}}$  are reversible upon cooling (Figures S10–S12). The overall results clearly suggest a reversible thermal conversion of the  $\{\text{Fe}^{\text{II}}_{\text{LS}}-\text{CN}-\text{Co}^{\text{III}}_{\text{LS}}\}$  pair to the  $\{\text{Fe}^{\text{III}}_{\text{LS}}-\text{CN}-\text{Co}^{\text{II}}_{\text{HS}}\}$  pair occurs in **1** via the intramolecular MMET process.<sup>12–20,22–25,37,39–41</sup>

UV/vis/near-IR spectroscopic studies were carried out on the building block  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$ ,  $[\text{Co}^{\text{II}}(\text{L})_2(\text{NCMe})_2]^{2+}$ , and complex **1** in both the solid and solution states at room temperature. The UV/vis spectra of  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  in acetonitrile show a ligand-to-metal charge-transfer (LMCT) transition band centered at around 408 nm, while a LMCT band centered at around 330 nm was observed for  $[\text{Co}^{\text{II}}(\text{L})_2(\text{NCMe})_2]^{2+}$  (Figure S13).<sup>15</sup> For **1**, the spectrum exhibits two sharp bands centered at around 337 and 420 nm with a broad band centered at around 695 nm (Figure S14). The former two bands are ascribed to the metal-to-ligand charge-transfer (MLCT) transition band coming from the  $[\text{Co}^{\text{II}}(\text{L})_2(\text{NCMe})_2]^{2+}$  unit and the LMCT band coming from the  $[\text{Fe}^{\text{II}}(\text{Tp})(\text{CN})_3]^{2-}$  unit, respectively, while the latter band is corresponding to a metal-to-metal charge-transfer (MMCT) transition band coming from the  $\text{Fe}^{\text{II}}_{\text{LS}} \rightarrow \text{Co}^{\text{III}}_{\text{LS}}$  unit (Figure S15).<sup>12–20,22–25,37,39–41</sup> The UV/vis

spectra of **1** in different solvents reveal an absorption spectrum similar to that in an acetonitrile solution (Figure S16). The solid-state spectrum of **1** displays an absorption band very similar to that of the solution-state spectrum (Figure S17), suggesting that the complex retained its identity in both the solid and solution states.

**Electrochemical Studies.** The electrochemical properties of **1**, along with  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  and  $[\text{Co}(\text{L})_2(\text{NCMe})_2]^{2+}$ , were investigated by cyclic voltammetry (CV), square-wave voltammetry, and differential-pulse voltammetry in an acetonitrile solution containing respective samples ( $\sim 1$  mM) with 0.1 M  $(\text{TBA})\text{PF}_6$  as an electrolyte at 300 K (Figures 2, S18, and S19 and Table S5). The cyclic



**Figure 2.** Cyclic voltammograms of **1** in acetonitrile containing 0.1 M  $(\text{TBA})\text{PF}_6$  as an electrolyte with various scan rates at 300 K. Arrows indicate an open-circuit potential along with the direction of the potential sweep.

voltammogram of  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  exhibits a quasi-reversible reduction for  $\text{Fe}^{\text{III}}$  to  $\text{Fe}^{\text{II}}$  at around  $-0.81$  V versus ferrocene/ferrocenium ( $\text{Fc}/\text{Fc}^+$ ), while a slow quasi-reversible oxidation process at around  $0.22$  V versus  $\text{Fc}/\text{Fc}^+$  was observed for  $\text{Co}^{\text{II}}$  to  $\text{Co}^{\text{III}}$  in  $[\text{Co}(\text{L})_2(\text{NCMe})_2]^{2+}$  (Figure S18).

The cyclic voltammograms of **1** reveal a quasi-reversible oxidation process at around  $0.08$  V versus  $\text{Fc}/\text{Fc}^+$ , corresponding to oxidation of the iron(II) center ( $\{\text{Fe}^{\text{II}}_2\text{Co}^{\text{III}}_2\}^{2+}/\{\text{Fe}^{\text{III}}\text{Fe}^{\text{II}}\text{Co}^{\text{III}}_2\}^{3+}$  along with an irreversible oxidation peak at around  $0.47$  V versus  $\text{Fc}/\text{Fc}^+$  ( $E_{\text{pc}}$ ; Figure 2).<sup>12–20</sup> The latter value is similar to that observed for  $[\text{Co}(\text{L})_2(\text{NCMe})_2]^{2+}$  (Figure S18, right). Apart from oxidation, the cyclic voltammograms also exhibit two quasi-reversible reductions for the cobalt(III) centers at around  $-0.15$  and  $-0.34$  V versus  $\text{Fc}/\text{Fc}^+$  ( $\{\text{Fe}^{\text{II}}_2\text{Co}^{\text{III}}_2\}^{2+}/\{\text{Fe}^{\text{II}}_2\text{Co}^{\text{II}}\text{Co}^{\text{III}}_2\}^+$  and  $\{\text{Fe}^{\text{II}}_2\text{Co}^{\text{III}}\text{Co}^{\text{II}}\}^+/\{\text{Fe}^{\text{II}}_2\text{Co}^{\text{II}}_2\}$ ). Also, the voltammograms show a quasi-reversible reduction process at around  $-0.76$  V versus  $\text{Fc}/\text{Fc}^+$  in **1**, which is slightly lower than the value observed for the iron(III) building block, i.e.,  $(\text{TBA})[\text{Fe}^{\text{III}}(\text{Tp})(\text{CN})_3]\cdot\text{H}_2\text{O}$  (Figure S18, left). These results are also consistent with the square-wave and differential-pulse voltammograms (Figure S19). The Gibbs free-energy difference ( $\Delta G_{\text{HL}}^0$ ) between the paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\}_2\{\text{Co}^{\text{II}}_{\text{HS}}\}_2$



and diamagnetic  $\{(\text{Fe}^{\text{II}}_{\text{LS}})_2(\text{Co}^{\text{III}}_{\text{LS}})_2\}$  species correlates with the difference between the redox potentials of the iron and cobalt centers ( $\Delta E = E_{\text{Co}^{\text{III}}/\text{Co}^{\text{II}}} - E_{\text{Fe}^{\text{II}}/\text{Fe}^{\text{III}}}$ ).<sup>12–20</sup> The measured  $|\Delta E|$  value for **1** is 0.23 V versus  $\text{Fc}/\text{Fc}^+$ , belonging in the range observed for similar  $\{\text{Fe}_2\text{Co}_2\}$  square complexes.<sup>12–20</sup> The corresponding  $\Delta G_{\text{HL}}^0$  value of **1** suggests the possibilities of entropy-driven stimuli-responsive reversible electron transfer between the donor and acceptor in **1**.

**Differential Scanning Calorimetry (DSC).** DSC studies were carried out for **1** to explore the possibilities of the thermally induced MMET process, which typically follows a first-order phase transition, and to estimate the associated changes in the thermodynamic parameters, i.e., enthalpy and entropy.<sup>24,25</sup> DSC measurements were performed in the temperature range of 300–420 K, with different sweep rates varying from 15 to 5  $\text{K min}^{-1}$  (Figure 3). DSC studies of **1**

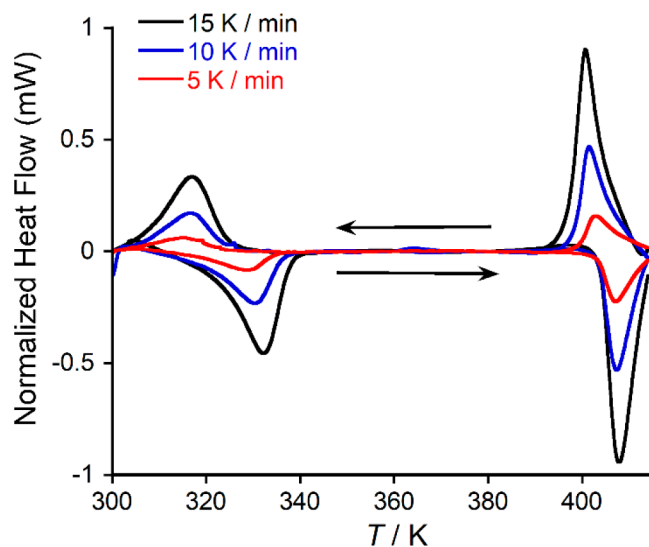


Figure 3. DSC plots of **1** shown between 300 and 420 K at different sweep rates.

show two reversible peaks at around 310–340 and 400–410 K, suggesting the first-order phase transition for **1**. Upon heating/cooling with a sweep rate of 10  $\text{K min}^{-1}$ , two sets of reversible endothermic/exothermic peaks were observed with a peak maximum at 332/316 K and 407/400 K, respectively. These values suggest the presence of a two-step transition in **1**, which is consistent with the magnetic studies (vide infra). All of the thermodynamic parameters have been calculated, and the corresponding  $\Delta H$  values were estimated to be 4.68 ( $T_{1/2}^{\uparrow}$ )/5.06 ( $T_{1/2}^{\downarrow}$ )  $\text{kJ mol}^{-1}$  and 8.78 ( $T_{1/2}^{\uparrow}$ )/7.62 ( $T_{1/2}^{\downarrow}$ )  $\text{kJ mol}^{-1}$ , respectively. The associated  $\Delta S$  values were calculated to be 14.17 ( $T_{1/2}^{\uparrow}$ )/15.98 ( $T_{1/2}^{\downarrow}$ ) and 21.56 ( $T_{1/2}^{\uparrow}$ )/19.01 ( $T_{1/2}^{\downarrow}$ )  $\text{J K}^{-1} \text{mol}^{-1}$ , respectively. These  $\Delta H$  and  $\Delta S$  values match well with the obtained values for other MMET complexes having stepwise transitions.<sup>24,25</sup>

**Magnetic Studies.** The temperature dependence of the magnetic susceptibility measurement was performed on polycrystalline samples of **1** in the temperature range of 2–420 K at a 1000 Oe external magnetic field (Figure 4).

For **1**, the measured  $\chi T$  values ( $\chi$  is the magnetic susceptibility equal to  $M/H$  per  $\{\text{Fe}_2\text{Co}_2\}$  unit) in the temperature range of 15–300 K are nearly constant with a value of 0.30  $\text{cm}^3 \text{mol}^{-1} \text{K}$ , which is slightly higher than the value expected for two LS iron(II) ions ( $S = 0$ ) and two LS

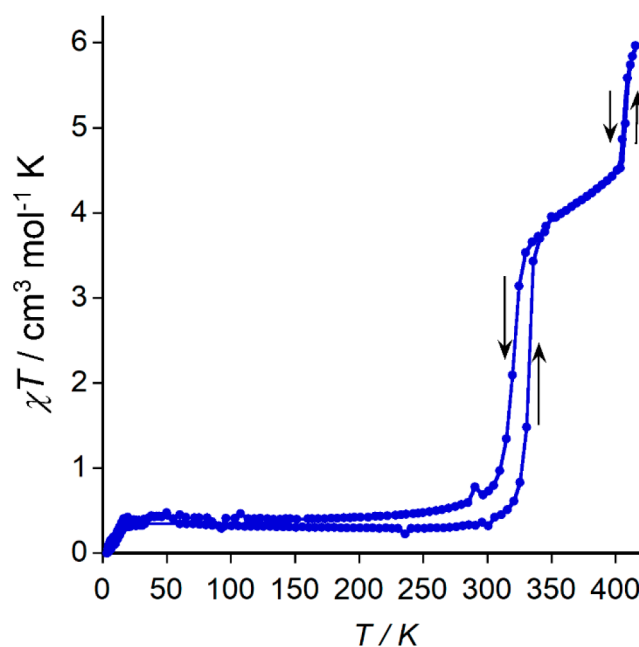
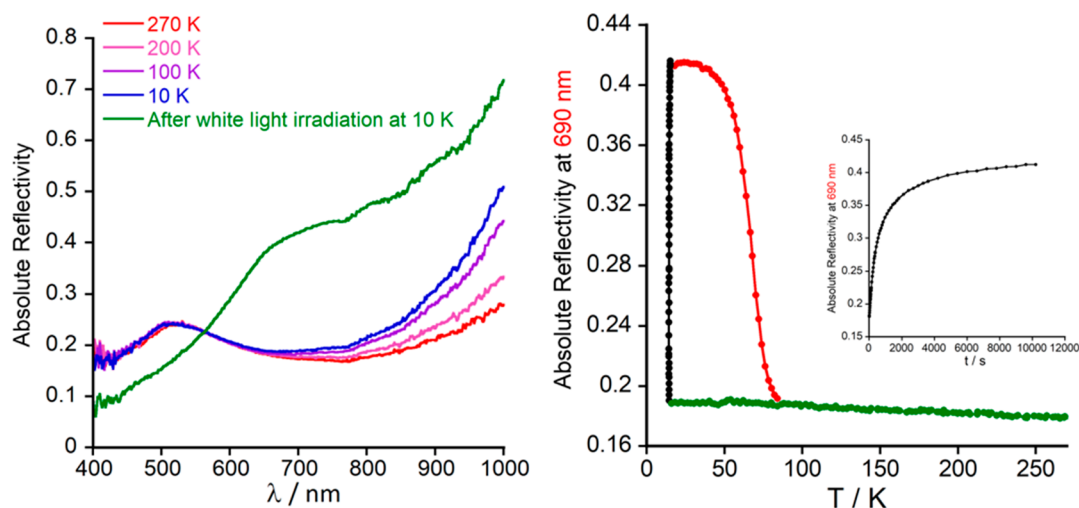


Figure 4. Temperature dependence of  $\chi T$  for **1** under a 1000 Oe external magnetic field with a sweep rate of 2  $\text{K min}^{-1}$ .

cobalt(III) ions ( $S = 0$ ). The slightly positive  $\chi T$  value might be due to the presence of a small fraction of the residual paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}(S=1/2) - \text{Co}^{\text{II}}_{\text{HS}}(S=3/2)\}$  state. With increasing temperature, one interesting phenomenon has been observed starting from 310 K. The  $\chi T$  value increases abruptly with a transition temperature  $T_{1/2}(1) \uparrow$  at 332 K to reach a plateau with values of 4.02–4.00  $\text{cm}^3 \text{mol}^{-1} \text{K}$  at 350 K. This  $\chi T$  value is higher than the value obtained for one LS iron(III) ion ( $S = 1/2$ ;  $g \approx 2.6$ )<sup>14,16,17,42</sup> and one HS cobalt(II) ion ( $S = 3/2$ ;  $g \approx 2.4$ )<sup>14,16,17</sup> both with significant orbital contributions, suggesting an intermediate phase with the  $\{\text{Fe}^{\text{II}}_{\text{LS}}(S=0) - \text{Co}^{\text{III}}_{\text{LS}}(S=0)\}$  and  $\{\text{Fe}^{\text{III}}_{\text{LS}}(S=1/2) - \text{Co}^{\text{II}}_{\text{HS}}(S=3/2)\}$  pairs in a ratio close to 1:1 generated via MMET between the iron and cobalt centers.<sup>14,16,17,24,25</sup> Upon a further increase in the temperature, the  $\chi T$  value increases sharply with a transition temperature  $T_{1/2}(2) \uparrow = 407$  K to reach a value of 5.98  $\text{cm}^3 \text{mol}^{-1} \text{K}$  at 420 K. This  $\chi T$  value is slightly lower than the values observed for two LS iron(III) ions ( $S = 1/2$ ;  $g \approx 2.6$ ) and two HS cobalt(II) ions ( $S = 3/2$ ;  $g \approx 2.4$ ) with significant orbital contributions, which can be ascribed by the nonsaturation nature of the curve measured up to this temperature.<sup>14,16,17,20,24,25</sup>

This indicates that the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}(S=0) - \text{Co}^{\text{III}}_{\text{LS}}(S=0)\}_2$  ground state of **1** converts to the high-temperature paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}(S=1/2) - \text{Co}^{\text{II}}_{\text{HS}}(S=3/2)\}_2$  state via a two-step MMET associated with spin transition on the cobalt center ( $\text{Co}^{\text{III}}_{\text{LS}} \rightarrow \text{Co}^{\text{II}}_{\text{HS}}$ ) upon an increase in the temperature. Because of the instability or possible decomposition of complex **1** at high temperature, these measurements do not show a plateau at 420 K, but the transition is almost fully described. During the cooling mode with a sweep rate of 2  $\text{K min}^{-1}$ , a narrow thermal hysteresis of 2 K at  $T_{1/2}(2) \downarrow = 405$  K and a wide thermal hysteresis of 12 K at  $T_{1/2}(1) \downarrow = 320$  K have also been observed for **1**. The small decrease of the  $\chi T$  values below 15 K is probably due to the combined effect coming from the intramolecular antiferromagnetic interactions or the anisotropy resulting from spin–



**Figure 5.** Left: Comparison of the surface reflectivity spectra in the temperature range of 270–10 K (red to blue) in the dark and after 160 min of white-light irradiation (green). Right: Thermal evolution of the absolute reflectivity signal plotted at 690 nm in cooling mode in the dark (270–10 K, at 4 K min<sup>-1</sup>, green), during white-light irradiation (0.6 mW cm<sup>-2</sup>, at 10 K, black), and in heating mode (10–270 K, at 0.4 K min<sup>-1</sup>, red) of **1**. Inset: Time evolution of the absolute reflectivity during white-light irradiation.

orbit coupling of HS cobalt(II) ions from the residual paramagnetic state.

The field-dependent magnetization studies were performed for **1** from 0 to 5 T at 2, 5, and 10 K (Figure S21, left). At 2 K, the magnetization value is 0.60 μ<sub>B</sub> for **1** at 5 T, confirming the presence of {Fe<sup>II</sup><sub>LS</sub>(S=0)–Co<sup>III</sup><sub>LS</sub>(S=0)} pairs for **1**, and the small magnetization value is coming from the residual paramagnetic {Fe<sup>III</sup><sub>LS</sub>(S=1/2)–Co<sup>I</sup><sub>HS</sub>(S=3/2)} state.<sup>14,16,17,20,24,25,35</sup> In addition, *M* versus *H*/*T* measurements exhibit nonsuperposition of the data on a single master curve (Figure S21, right). The nonsaturation and nonsuperposition of the magnetization curves for **1** suggest the presence of magnetic anisotropy due to the residual paramagnetic {Fe<sup>III</sup><sub>LS</sub>(S=1/2)–Co<sup>I</sup><sub>HS</sub>(S=3/2)} state.

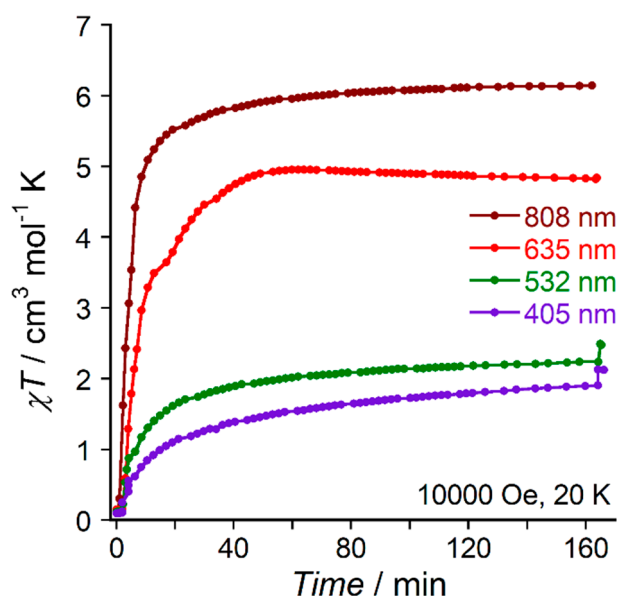
**Optical Studies.** Solid-state optical reflectivity measurements for complex **1** were carried out in the temperature range of 270–10 K with a sweep rate of 4 K min<sup>-1</sup> in both cooling and heating modes. No significant changes were observed in the evolution of the absolute reflectivity as a function of the temperature. After, the sample was cooled in the dark to 10 K, and the optical reflectivity of **1** was measured during irradiation with white light (power 0.6 mW cm<sup>-2</sup>) for 2 h. For **1**, the spectrum resulting upon irradiation at 10 K shows a clear difference in the absolute reflectivity of ca. 500 nm and also in the range of 650–800 nm (Figure 5, left), the region expected for MMET between the iron and cobalt centers. This suggests the presence of a photoinduced MMET and the formation of the metastable paramagnetic {Fe<sup>III</sup><sub>LS</sub>–CN–Co<sup>I</sup><sub>HS</sub>}<sub>2</sub> state of complex **1** at low temperature.<sup>14–20</sup>

The evolution of the absolute reflectivity at 690 nm versus time of irradiation indicates a rapid increase before it tends to saturation, and the major part of the photoexcitation was observed after 90 min, reaching a value of 0.42 (from 0.18) after 160 min of light irradiation (Figure 5, right). After white-light irradiation, the thermal stability of the photoinduced metastable state has also been studied in the dark by measuring the absolute reflectivity as a function of the temperature (sweep rate 0.4 K min<sup>-1</sup>). The reflectivity value was constant up to 60 K, and then the value starts to decrease as the photoinduced metastable state thermally relaxed back to the

diamagnetic ground state where complete relaxation was observed at 85 K (*T*<sub>LIESST</sub>; Figure 5, right), as further confirmed by photomagnetic measurements (vide infra). The absolute reflectivity data obtained beyond this temperature exhibited the same behavior as that in the dark. In addition to the above experiments, a set of 14 light-emitting diodes (LEDs; 385–1050 nm) were also used to conduct the photoexcitation experiments at 10 K to find the most efficient irradiation wavelength for MMET (Figure S22). All of 14 LEDs show some photoeffect, i.e., an increase in the absolute reflectivity value, in agreement with the efficiency of the white light, where the maximum efficiency was observed for the 850 nm region (Figure S23, left). Consequently, 850 nm LED irradiation has been applied for photoexcitation at 10 K. A more complete conversion to the photoinduced excited state can be reached from LED irradiation (for 1 h with 80 mW cm<sup>-2</sup>) than from white-light irradiation. After 850 nm LED light irradiation, the thermal stability of the photoinduced metastable state has also been studied in the dark by measuring the absolute reflectivity as a function of time (Figure S23, right). It is worth noting that the photoinduced metastable state has a long lifetime and was relatively stable at this temperature. Additionally, thermal relaxation of the photoinduced excited state to the ground state has been observed (*T*<sub>LIESST</sub> = 90 K; sweep rate 0.4 K min<sup>-1</sup>), a little higher than that of white-light irradiation (Figure S24). Moreover, the effect of the LED power of 850 nm LED light was also studied (Figure S25), which indicates that 80 mW cm<sup>-2</sup> brings the best ratio conversion by power. It is worth mentioning that the higher LED power does not significantly improve the conversion.

**Photomagnetic Studies. ON Switching Mode.** The photosensitivity of **1** was further investigated by measuring the magnetic properties under diode laser light irradiation at 20 K (measured *χT* as a function of time) to explore the light-induced MMET between iron and cobalt centers. Complex **1** was irradiated at 20 K and a 10000 Oe external magnetic dc field with four different diode laser sources (405, 532, 635, and 808 nm with powers ranging from 5 to 12 mW cm<sup>-2</sup>), covering the visible region. An obvious increase of the *χT* value was observed for each laser source (photomagnetic effect);

however, the maximum effect in the  $\chi T$  value was obtained for 808 nm light (Figure 6), which is also consistent with the reflectivity and UV/vis studies (MMET from the  $\text{Fe}^{\text{II}}_{\text{LS}} \rightarrow \text{Co}^{\text{III}}_{\text{LS}}$  unit, vide supra).

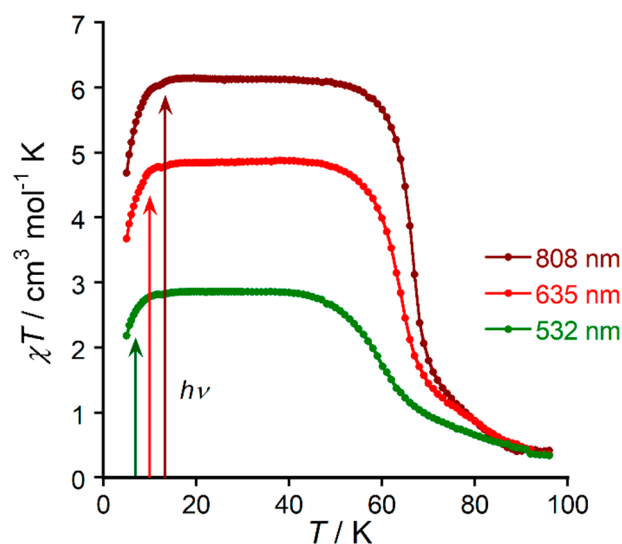


**Figure 6.**  $\chi T$  versus time curves for **1** under 405, 532, 635, and 808 nm laser light irradiation at 20 K and 10000 Oe field.

Under 808 nm light irradiation, the  $\chi T$  value increases rapidly to reach a saturation value of  $6.14 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$  after 170 min. This value lies in the expected range observed in similar  $\{\text{Fe}_2\text{Co}_2\}$  model complexes,<sup>14–20,24,25</sup> signifying complete photoinduced conversion of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state to the metastable paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state at 20 K. It is worth mentioning that around 25–70% a paramagnetic state can be achieved after irradiation with 405, 532, and 635 nm laser light. Field-dependent magnetization was studied for **1** after irradiation with all four wavelengths (808, 635, 532, and 405 nm) from 0 to 5 T at 2 K (Figure S26). At 2 K and 5 T, the magnetization value was  $5.91 \mu_{\text{B}}$  for **1** after irradiation with 808 nm, confirming also the presence of the  $\{\text{Fe}^{\text{III}}_{\text{LS}}(S=1/2)\text{-Co}^{\text{II}}_{\text{HS}}(S=3/2)\}_2$  state at this temperature for **1**.<sup>14–20,24,25</sup> Moreover, all of these results are consistent with complete photoconversion of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state to the metastable paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state for **1**.

To investigate the thermal stability and relaxation of the photoinduced metastable  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state, the temperature dependence of the  $\chi T$  value was measured (sweep rate  $0.4 \text{ K min}^{-1}$ ) in the dark after it reached saturation (ca. 160 min at 808 nm; Figure S27). In the dark, in the temperature range of 20–2 K, the  $\chi T$  value decreases sharply to reach a minimum from  $6.20$  to  $4.60 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ , indicating magnetic anisotropy coming from the cobalt(II) side and/or intermolecular antiferromagnetic interactions.<sup>14–20,24,25</sup> Upon an increase in the temperature from 20 to 70 K,  $\chi T$  values remain almost constant, reaching a plateau value of  $6.15 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$ , with such a small increase of the  $\chi T$  value at low temperature suggesting intramolecular ferromagnetic interaction between the  $\text{Fe}^{\text{III}}_{\text{LS}}$  and  $\text{Co}^{\text{II}}_{\text{HS}}$  centers in a photogenerated  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  metastable state of **1**.<sup>14–20,24,25</sup> Upon a

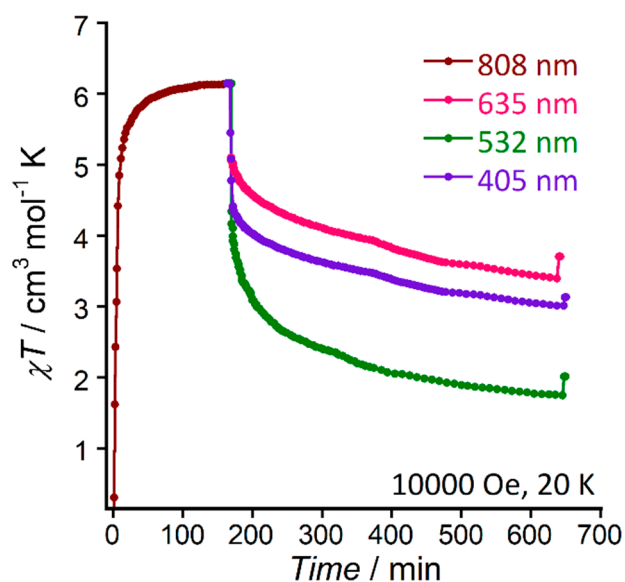
further increase in the temperature, the photoinduced paramagnetic metastable state relaxes back to the thermodynamic diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state at 91 K, which is consistent with optical reflectivity studies (vide supra). Similar experiments have been carried out after irradiation with the other two laser lights e.g., 635 and 532 nm (sweep rate  $0.4 \text{ K min}^{-1}$ ). It is worth noting that relaxation of the photoinduced metastable state to the diamagnetic ground state happened nearly at the same temperature, 91 K (Figure 7).



**Figure 7.** Temperature dependence of  $\chi T$  of **1** after irradiation with 532, 635, and 808 nm laser light at 10000 Oe field (sweep rate  $0.4 \text{ K min}^{-1}$ ).

**OFF Switching Mode.** Photomagnetic studies were further continued to explore the possibility of observing the reverse photoeffect, which converts the photoinduced metastable paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state to the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state via light-induced MMET at the same temperature compared to the thermal effect. After the saturation  $\chi T$  values are reached using 808 nm laser light irradiation (ca. 160 min), the evolution of  $\chi T$  as a function of time was measured under three different laser sources (405, 532, and 635 nm) at 20 K and 10000 Oe (Figure 8). Remarkably, all three wavelengths produced a significant reduction in the  $\chi T$  values, suggesting depopulation of the photoinduced metastable paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state to the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state.<sup>17</sup> The maximum difference was observed upon irradiation with 532 nm laser light (matches well with the reflectivity and UV/vis data; MMET from the  $\text{Co}^{\text{II}}_{\text{HS}} \rightarrow \text{Fe}^{\text{III}}_{\text{LS}}$  unit), where the  $\chi T$  value decreases from  $6.14$  to  $1.75 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$  after 475 min, indicating 75% recovery of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  ground state at 20 K.

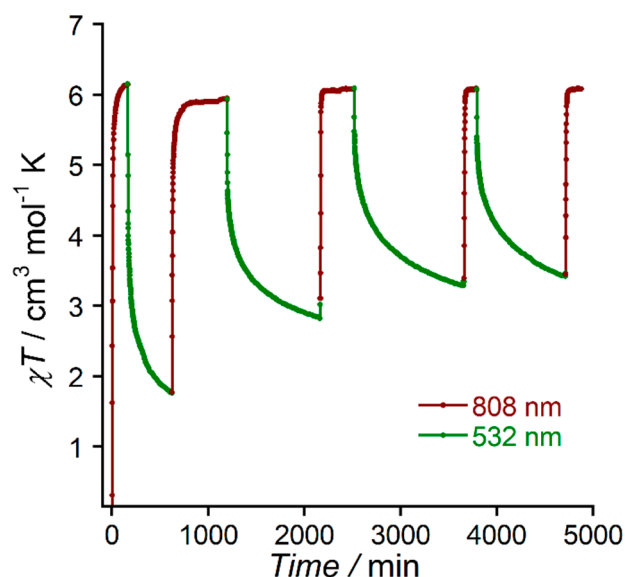
The incomplete conversion is consistent with the results found during the above ON switching mode of photomagnetic measurements, where all laser sources show a significant increase of the  $\chi T$  values (populating the paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state). Indeed, it is also observed that the most effective laser source in the ON switching mode is less effective in the OFF switching mode. In addition, for a particular wavelength, an increase and a decrease in the  $\chi T$  values of photoinduced states show a significant difference in



**Figure 8.**  $\chi T$  versus time curves for **1** under 405, 532, 635, and 808 nm laser light irradiation at 20 K and 10000 Oe field.

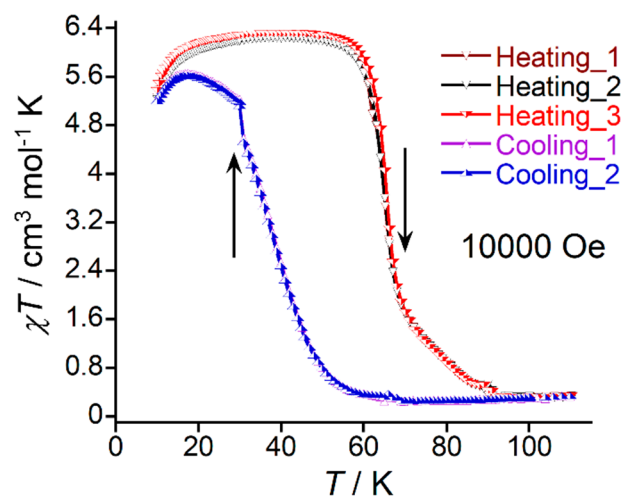
the ON and OFF switching modes, respectively. For example, in the case of the 405 nm laser, a ca. 28% conversion of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  state to the paramagnetic metastable  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  state in the ON switching mode has been observed, while a 51% reduction of the paramagnetic state to the diamagnetic ground state was observed in the OFF switching mode. This can be ascribed to a larger sample size, the thickness of the sample rather than a light penetration length, and faster kinetics in the ON switching mode in comparison to the OFF switching mode.<sup>17</sup> It is worth noting that the photoconversion rate in the OFF switching mode is slower than that in the ON switching mode for complex **1**, and a similar effect has been observed in another  $[\text{Fe}_2\text{Co}_2]$  complex.<sup>17,25</sup>

**ON/OFF Cycling.** The photoreversibility studies were performed on **1** by measuring  $\chi T$  versus time under several cycles of successive irradiation with 808 nm (most effective light for the ON mode) and 532 nm (most effective light for the OFF mode) lasers for ca. 200 and 1000 min, respectively, at 20 K under 10000 Oe field. Figure 9 represents variation of the  $\chi T$  value with time after successive irradiation cycles. The individual curves do not superimpose each other; however, they are very similar. Besides, the first irradiation curve in the ON mode appears to be slightly different from the following ones. The  $\chi T$  values obtained after successive irradiation at 808 nm are almost constant, while a gradual decrease in the  $\chi T$  value has been noticed after successive 532 nm laser light irradiation cycles, which can be ascribed to an aging effect. After 808 nm light irradiation on the diamagnetic ground state, the  $\chi T$  value increases from 0.12 to 6.14  $\text{cm}^3 \text{mol}^{-1} \text{K}$ , whereas after 532 nm light irradiation on the photoexcited metastable state, the  $\chi T$  value decreases from 6.14 to 1.76  $\text{cm}^3 \text{mol}^{-1} \text{K}$ . These results suggest that a reversible ON/OFF switching of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  and paramagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  states of **1** can be achieved by 808 and 532 nm successive irradiation over numerous cycles at 20 K. In addition, field-dependent magnetization was studied for **1** under cycles of successive irradiation at 808 and 532 nm, from 0 to 5 T at 2 K (Figure S28). These results are consistent with the above findings.



**Figure 9.** Plot of  $\chi T$  versus time under cycles of successive irradiation with 808 nm ( $6 \text{ mW cm}^{-2}$ ) and 532 nm ( $10 \text{ mW cm}^{-2}$ ) at 20 K and 10000 Oe field of **1**.

**Light-Induced Thermal Hysteresis Effect.** Given the reversibility of thermal relaxation of the photoinduced metastable state and the possibility of observing the photoinduced memory effect, the photomagnetic measurement was further investigated following the temperature dependence of the  $\chi T$  product under continuous 808 nm laser irradiation during both heating and cooling modes in the temperature range of 10–120 K (sweep rate  $0.4 \text{ K min}^{-1}$ ). Remarkably, an interesting hysteresis width of 27 K has been observed in **1**, which is unique in MMET systems, termed the light-induced thermal hysteresis (LITH) effect<sup>43</sup> (Figure 10). First, the sample was irradiated with a 808 nm laser at 20 K for ca. 160 min under a 10000 Oe magnetic field to reach complete conversion of the diamagnetic  $\{\text{Fe}^{\text{II}}_{\text{LS}}\text{-CN-Co}^{\text{III}}_{\text{LS}}\}_2$  state to the photomagnetic  $\{\text{Fe}^{\text{III}}_{\text{LS}}\text{-CN-Co}^{\text{II}}_{\text{HS}}\}_2$  metastable state, and the 808 nm light was continuously switched ON during the measurement process of thermal relaxation ( $\chi T$  versus  $T$ )



**Figure 10.**  $\chi T$  versus  $T$  curves of **1** under 808 nm light irradiation during cooling as well as warming with a sweep rate of  $0.4 \text{ K min}^{-1}$  at a 10000 Oe field.



compared to the general thermal relaxation measurements (Figures 7 and S27).

Upon an increase in the temperature (sweep rate 0.4 K min<sup>-1</sup>), the measured  $\chi T$  value (6.30 cm<sup>3</sup> mol<sup>-1</sup> K) at 20 K remained almost constant up to 60 K and then decreased smoothly to reach a value of 0.33 cm<sup>3</sup> mol<sup>-1</sup> K at 100 K. In a successive cooling mode (808 nm light is switched ON), the  $\chi T$  value remained almost constant up to 70 K (0.27 cm<sup>3</sup> mol<sup>-1</sup> K) and then increased gradually to reach a plateau with 5.67 cm<sup>3</sup> mol<sup>-1</sup> K at 20–15 K. Upon further cooling,  $\chi T$  reached a value of 5.00 cm<sup>3</sup> mol<sup>-1</sup> K at 10 K (lower than the heating mode). The reversibility of the LITH effect was also confirmed by measuring  $\chi T$  versus  $T$  in numerous heating and cooling cycles while switching ON the 808 nm laser light (Figure 10). The temperature dependence of the  $\chi T$  value (switching ON the 808 nm laser) exhibited an effect similar to that of the OFF switching mode (Figure S27) upon an increase of the temperature from 10 to 120 K (measuring thermal relaxation, 0.4 K min<sup>-1</sup>). When the sample was cooled to 10 K while switching ON the 808 nm laser (0.4 K min<sup>-1</sup>), the temperature dependence of the  $\chi T$  value in both curves, i.e., cooling and heating modes (Figure 10, red and blue curves), followed different paths. It is worth noting that at 60 K the  $\chi T$  value in the heating mode was 6.30 cm<sup>3</sup> mol<sup>-1</sup> K, whereas in the cooling mode, the  $\chi T$  value was 0.75 cm<sup>3</sup> mol<sup>-1</sup> K (switching ON the 808 nm laser), which indicates the LITH effect. While the temperature is being increased in the ON switching mode, a competition between photoexcitation of the diamagnetic state to the paramagnetic state and thermal relaxation of the photoinduced metastable paramagnetic state to the diamagnetic ground state is occurring. When the temperature is higher, the thermal relaxation process predominates over the photoexcitation effect, whereas at a lower temperature, the photoexcitation effect wins over thermal relaxation, responsible for this interesting LITH effect.

## CONCLUSION

We have presented a new molecular model complex of [FeCo] PBAs, i.e., a [Fe<sub>2</sub>Co<sub>2</sub>] square complex. X-ray structural analysis, (photo)magnetic, spectroscopic, and other physical measurements revealed that the redox potential of the building block, solvents, and capping ligand (Tp) played a major role in their physical properties. The [Fe<sub>2</sub>Co<sub>2</sub>] complex exhibits reversible two-step thermo- and photoinduced MMET between the iron and cobalt centers, converting the diamagnetic phase (Fe<sup>II</sup><sub>LS</sub>-CN-Co<sup>III</sup><sub>LS</sub>) to the paramagnetic phase (Fe<sup>III</sup><sub>LS</sub>-CN-Co<sup>II</sup><sub>HS</sub>). Again, electrochemical studies suggested that the redox potential difference between the iron and cobalt units in the complex determines the electron-transfer probability. If the difference in the redox potential of the donor and acceptor units (iron and cobalt units) is too high or low, it will stabilize only one state (either the diamagnetic or paramagnetic state); when the redox potential difference is comparable with  $k_B T$  (order of magnitude of  $\Delta G$ ), a reversible MMET is possible upon application of external perturbation, i.e., for complex 1. All interesting magnetic properties of the square complex can be explained by considering the different electronic states of the metal ions coming from the substituent group of capping ligands. 1 exhibits a reversible switching property, indicating that the energy gaps between the frontier orbitals of the iron and cobalt centers are close to each other, comparable with  $k_B T$ . The magnetic (thermo- and photoinduced) bistability probably arises from MMET, which induces a reversible

conversion between the diamagnetic phase {Fe<sup>II</sup><sub>LS</sub>-CN-Co<sup>III</sup><sub>LS</sub>}<sub>2</sub> and the paramagnetic phase {Fe<sup>III</sup><sub>LS</sub>-CN-Co<sup>II</sup><sub>HS</sub>}<sub>2</sub>. On the basis of the above data, diamagnetic-paramagnetic reversible switching of 1 is possible in two ways. So, the stimuli-responsive nature of 1 may pave a way in several new directions to rationally design materials with novel functionalities for applications as multiresponsive materials for future technological applications. Studies of the impact of substituents in the building block as well as in the capping ligand on MMET with detailed characterization of the photoinduced excited state using SCXRD and PXRD analyses are currently ongoing in our research group.

## EXPERIMENTAL SECTION

Detailed experimental procedures, including materials, crystal structure analyses, physical methods, (photo)magnetic measurements, and optical studies, are described in the Supporting Information.

**Synthesis of {[Fe(Tp)(CN)<sub>3</sub>]<sub>2</sub>[Co(L)<sub>2</sub>]<sub>2</sub>}(BF<sub>4</sub>)<sub>2</sub>·2CH<sub>3</sub>CN·6H<sub>2</sub>O (1).** The reaction of Co(BF<sub>4</sub>)<sub>2</sub>·6H<sub>2</sub>O (34 mg, 0.1 mmol) and L (44 mg, 0.2 mmol) in a mixture of solvent [acetonitrile and water mixture (9:1)] (10 mL) produces a brownish-pink solution. After 10 min of stirring, the resulting solution was added dropwise to a solution of (TBA)[Fe<sup>III</sup>(Tp)(CN)<sub>3</sub>]-H<sub>2</sub>O (62.4 mg, 0.1 mmol) in an acetonitrile and water mixture (9:1) (15 mL). The resulting dark-green solution indicates MMET between the iron(III) and cobalt(II) centers. The reaction mixture was filtered and kept for crystallization by slow evaporation. After 10 days, dark-green rectangular plate-type crystals of 1 were obtained. Anal. Calcd for C<sub>72</sub>H<sub>94</sub>Co<sub>2</sub>Fe<sub>2</sub>N<sub>36</sub>O<sub>10</sub>B<sub>4</sub>F<sub>8</sub> (MW 2048.63 g mol<sup>-1</sup>): C, 42.41; H, 4.63; N, 24.61. Found: C, 41.78; H, 4.90; N, 24.27. IR (KBr, cm<sup>-1</sup>): 3423, 3154, 3123, 3106, 2983, 2940, 2473, 2463, 2127, 2110, 2065, 1664, 1486, 1426, 1311, 1211, 1186, 1113, 1086, 1046, 980, 902, 802, 760, 715, 624, 507, 421.

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