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1 **Experimental evidence for crustal control over seismic fault**
2 **segmentation**

3

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11

12 **ABSTRACT**

13 **Strike-slip faults are described as continuous structures, while they are actually**
14 **formed of successive segments separated by geometrical complexities. Although**
15 **this along-strike segmentation is known to affect the overall dynamics of**
16 **earthquakes, the physical processes governing the scale of this segmentation**
17 **remain unclear. Here, we use analogue models to investigate the structural**
18 **development of strike-slip faults and the physical parameters controlling**
19 **segmentation. We show that length of fault segments is regular along strike, and**
20 **scales linearly with the thickness of the brittle material. Variations of the**
21 **rheological properties only have minor effects on the scaling relationship. Ratios**
22 **between the segment length and the brittle material thickness are similar for co-**
23 **seismic ruptures and sandbox experiments. This supports a model where crustal**
24 **seismogenic thickness controls fault geometry. Finally, we show that the**

25 **geometrical complexity acquired during the strike-slip fault formation withstands**
26 **cumulative displacement. Thus, inherited complexity impedes the formation of an**
27 **ever straighter fault, and might control the length of earthquakes ruptures.**

28

29 INTRODUCTION

30 Strike-slip faults are discontinuous features formed by a succession of segments,
31 separated by geometrical complexities, such as bends, or relay zones (Fig. 1.a, S1)(e.g.
32 Segall and Pollard, 1980). Although the size of these complexities is highly variable, even
33 the smaller ones can significantly affect initiation, propagation, and termination of
34 earthquakes (e.g. King and Nabelek, 1985; Wesnousky, 2006; Klinger et al., 2005;
35 Manighetti et al., 2007), implying their persistence at seismogenic depth (e.g. Schwartz
36 and Coppersmith, 1984; Wei et al., 2011; Schwartz, 2018). Fault discontinuities also
37 modulate amplitude and style of co-seismic surface deformation (Klinger et al., 2006;
38 Vallée et al., 2008), and might stop ruptures propagation if larger than 5 km (e.g.
39 Wesnousky, 2006). Recently, numerical simulations of rupture scenarios have confirmed
40 the strong dependency of earthquake rupture models on the geometry of the fault
41 system (e.g. Aochi and Ulrich, 2015; Lozos, 2016; Klinger et al., 2018). As a consequence,
42 understanding the genesis of the segmentation and its scaling is critical for physics-
43 based earthquake modeling, as well as for hazard assessment, since complexities impact
44 the earthquake rupture length, and thus its magnitude.

45 Based on experimental, seismological, and geomorphological studies, fault segmentation
46 was first described as fractal (e.g. Aviles et al., 1987; Okubo and Aki, 1987). Alternatively,
47 it was suggested that faults are composed of a fixed number of segments, implying a
48 segment length varying between faults (e.g. Manighetti et al., 2015).

49 Conversely, studies based on co-seismic rupture maps have argued in favor of a constant
50 segment length of about $18 \text{ km} \pm 5 \text{ km}$, independent of the regional tectonic setting or
51 the earthquake magnitude, suggesting an external controlling factor, which is assumed
52 to be the seismogenic crust thickness (e.g. Bilham and Williams, 1985; Klinger,
53 2010)(Fig. 1.a). However, detailed co-seismic rupture maps are available only for a
54 limited number of earthquakes, preventing a reliable statistical analysis. The lack of data
55 combined with our limited knowledge of the fault segmentation process prevent us from
56 discriminating between these models.

57 To tackle the origin of strike-slip fault segmentation, we followed an analogue modeling
58 approach, which provides the complete geometrical evolution of the fault structure
59 while accumulating slip. Moreover, a large number of experiments can be conducted,
60 allowing for a thorough exploration of the impact of the geometrical and rheological
61 parameters.

62 **ANALOGUE MODELING OF STRIKE-SLIP FAULTS**

63 Numerous analogue studies already focused on the formation and final geometry of
64 strike-slip faults systems (Dooley and Schreurs, 2012 and references therein). In
65 experiments with sand overlying a straight basal dislocation, the first structures to
66 appear are the Riedel shears (R-shears)(Riedel, 1929; Naylor et al., 1986), propagating
67 upward (e.g. Ueta et al., 2000, Cambonie et al., 2018), oriented at the surface at an angle
68 of $\varphi_{\text{int}}/2$ relative to strike-slip basal fault (φ_{int} , angle of internal friction of the material).
69 This orientation is due to the helicoidal shape of the Riedels (Fig 1,b) (Naylor et al.,
70 1986; Mandl, 1987). The next structures to appear are first, synthetic shears (S-shear)
71 that only slightly deviate from the strike of the basal fault, and finally P-shears, which
72 cross the basal fault in the opposite direction. Eventually, coalescence of these different

73 shears leads to the formation of an anastomosed fault zone (e.g. Naylor et al., 1986;
74 Richard et al., 1995).

75 Here, we investigate if the first structures to appear during the formation of strike-slip
76 faults can be responsible for the long-term segmentation. As these faults result from the
77 linkage of the S-shears, the initial discontinuities between them at the locations of the
78 Riedels might persist and give rise to geometrical complexities. Therefore, we consider
79 the distance between two adjacent Riedels as equivalent to a fault segment. Although
80 dependency of the length of Riedels on the thickness of the frictional material had
81 already been suspected, it was never indisputably demonstrated (Tchalenko, 1970;
82 Atmaoui et al., 2006, Hatem et al., 2017). Thus, we first evaluate the influence of the
83 material thickness on the inter-Riedels distance, and then the impact of intrinsic
84 properties of the material.

85 We used an analogue model set-up known as Riedel experiments in a 1.2m x 0.8m box
86 (detailed set-up in Supp. Mat., Fig. S2). The model reproduces brittle deformation in a
87 non-cohesive material during the onset of a strike-slip fault in a homogeneous sandpack
88 above a straight-basal fault. To test the impact of the frictional properties on the fault
89 geometry, we used four different Aeolian sands with different internal frictions (Sand 1,
90 $\varphi_{\text{int}}=43.7^\circ$; Sand 2, $\varphi_{\text{int}}=35.6^\circ$; Sand 3, $\varphi_{\text{int}}=33.4^\circ$; Sand 4, $\varphi_{\text{int}}=22.1^\circ$) and variable basal
91 materials, such as PVC ($\varphi_{\text{b}}=13^\circ$) Alkor-foil ($\varphi_{\text{b}}=18^\circ$) and sand paper ($\varphi_{\text{b}}>43,7^\circ$). Some of
92 these sands are poured and have a rate-hardening behavior (sand 2, 3 and 4), while sand
93 1 is sedimented and rate-weakening (see Supp. Mat. for material properties). A camera
94 placed above the box records the deformation every 0.5mm of basal displacement to
95 generate orthoimages of the sandbox surface, with a resolution of the order of the grain
96 size. Moreover, to assess how the shear is accommodated by the different structures, we
97 computed the incremental displacement field from optical images correlation using

98 Micmac software (Rosu et al., 2015). The set-up is intentionally simple to limit the
99 number of free parameters and to ensure a good control on the boundary conditions of
100 the experiment. The aim of these experiments is not to reproduce details of a realistic
101 strike-slip fault evolution, but to retrieve relationships between fault-characteristic
102 patterns and properties of material.

103 For each sandpack thickness (T), three parameters were extracted from surface images
104 for each Riedel: the inter-Riedel distances (S) measured parallel to the trace of the basal
105 fault, the Riedels length (L), and the angle (α) formed by the Riedel with the direction of
106 the basal fault (Fig. 1.c). We ran repeated experiments with thicknesses varying from 2
107 to 5 cm (Table S1). Only the experiments with thicknesses lower or equal to 5cm were
108 free of edge effects (Fig. S3). The minimum thickness is constrained by the ability of the
109 sand to develop the 3D helicoidal shape of Riedel shears. When sand is thinner than 1.5
110 cm, instead of the Riedel shears, a continuous fault directly appears at the surface (Fig.
111 S4).

112 To provide robust correlations between thickness, material properties, and fault
113 geometry, we need repeated experimental measurements. As several Riedels are visible
114 in each experiment, we checked that we could use them as independent measurements
115 to improve statistical significance (see verification in Supp. Mat.). In the following, for
116 clarity's sake, only average quantities per experiment are presented and discussed
117 (individual measurements in Supp. Mat. Fig. S7-S11).

118 **SCALING RELATIONSHIPS DERIVED FROM SANDBOX**

119 Surface images of the experiments for different thicknesses (T) show that the deformed
120 area widens and the inter-Riedel distance (S) increases for thicker sandpack (Fig. 2.a).

121 Indeed, systematic measurements of (S) for different (T) show a linear correlation
122 between those parameters (Fig. 2.b), with a scale factor of about 2.8 for Sand 1.

123 To test the potential effect of the material properties, we first performed experiments
124 with sands of different internal friction. The three sets of experiments exhibit the same
125 trend: (S) increases with (T)(Fig. 2.c, S7), and the measurements distributions are well
126 modeled by linear relationships. Nevertheless, depending on the internal friction, the
127 slope of the best-fit lines changes, indicating a partial control of the material properties
128 on the size of the structures. The slope is 1.43 for the smaller friction ($\phi_{\text{int}}=33.4^\circ$, Sand
129 3), 1,61 for the intermediate friction ($\phi_{\text{int}}=35.6^\circ$, Sand 2), and 2.76 for the highest one
130 ($\phi_{\text{int}}=43.7^\circ$, Sand 1). As Sand 2 and 3 present relatively similar relationships (Fig. 2.c)
131 and are both strain-hardening, we cannot exclude from our data that the discrepancy
132 between Sand 1 and other sands could be due to its strain-weakening properties instead
133 of its internal friction (Fig. S11). Nevertheless, as Riedel shears form when shear stress
134 reaches the static friction (Tchalenko, 1970, Ritter et al. 2018), we think that the
135 internal friction, of the sands controls the relation between S and T, and not its strain
136 hardening or weakening property.

137 Using different materials to test the impact of the internal friction might also change the
138 basal friction between the different experiments. To assess a possible trade-off between
139 these two parameters, we conducted three series of experiments using Sand 1 with
140 different basal materials (Fig. 2.d, S7). For the three tested basal frictions we observed
141 linear relations between (S) and (T). The slopes of the regressions are quite similar, and
142 only the intercepts change slightly (Fig. 2.d). The associated change of x-axis intercept
143 corresponds to the minimal thickness necessary to develop the Riedels-helicoid. This
144 observation suggests that the basal friction impacts the formation of the helicoid,

145 probably by changing the width of the deformed area at depth (see Supp. Mat.
146 mechanical implications).

147 We then examined the effect of the cohesion on the inter-Riedel distance (S), using non-
148 cohesive (Sand 3) and strongly cohesive (Sand 4) sands (Fig. 2.e, S8). Again, (S)
149 increases linearly with (T). The two linear regressions present the same slope but
150 shifted intercepts, with a larger value for the highly cohesive sand (of about 1.5 cm).
151 However, this difference may not be significant in view of the small number of
152 experiments associated to a high dispersion of the Sand 4 measurements, interpreted as
153 resulting from heterogeneities in the sand mixture due to its electro-static charge (see
154 Supp. Mat.). We consider that cohesion only marginally affects the fault geometry, as (S)
155 is only slightly affected whereas all parameters vary and uncertainties are significant.

156 **COMPARISON OF EXPERIMENTAL RESULTS WITH EARTHQUAKE SURFACE-** 157 **RUPTURES GEOMETRY**

158 This new set of experiments demonstrates unambiguously that the inter-Riedel distance
159 (S) is directly controlled by the thickness of the sandpack (T)(Fig. 2), while the material
160 properties have secondary effects. The internal friction controls the scaling factor of the
161 linear relation between (S) and (T), whereas the basal friction and the cohesion only
162 marginally impact this relation. The width and rheology of the basal deformation zone,
163 which are not addressed here, might also affect the size of structures (Hatem et al. 2017,
164 Zuza et al. 2017, Yang et al. 2020).

165 According to field studies, strike-slip fault segments, measured between two successive
166 relay zones, have similar lengths independently of the local tectonic setting (Klinger,
167 2010). Thus, it was suggested that the average length of fault segments scales with the
168 thickness of the seismogenic crust (Bilham and Williams, 1985; Klinger, 2010), with a

169 ratio close to 1. To compare our experiments with natural observations, we selected
170 sands 2 and 3, both non-dilatant and with frictions within the same range as standard
171 upper crust materials ($\mu \sim 0.7$) (e.g. Byerlee, 1978).

172 For comparison, the inter-Riedel distance and the length of fault segments were
173 normalized respectively by the thickness of sand and of seismogenic crust, which is
174 derived from the literature (details in Supp. Mat.). The average normalized length for
175 natural fault segments is 1.16, and for sands 2 and 3 it is 1.17 and 1.3, which are
176 indistinguishable from the natural case (Fig. 3). Moreover, individual ratio values are all
177 in a narrow range ($\sigma = 0.43$), which indicates that such similarity is not coincidental. This
178 supports the hypothesis that the length of continental strike-slip segments is directly
179 controlled by the thickness of the seismogenic crust.

180 Although this approach does not capture the effects of a dynamic rupture, it sheds light
181 upon the way faults sections grow to eventually coalesce into an apparent continuous
182 structure. Interestingly, our experiments show that even when the cumulative
183 displacement during the experiment starts to be large enough to dismantle the Riedels
184 and to allow coalescence of successive S-shears into a longer fault, the initial geometric
185 discontinuity related to the location of the Riedel is preserved (Fig. 4.a). Riedels act as
186 inherited faults influencing the location of subsequent deformation. In the case of strike-
187 slip faults, such discontinuities will turn into relay zones, either compressional or
188 extensional. Indeed, in nature faults form in crust that already went through multiple
189 stages of deformation and that displays local geological variation. Thus the simple
190 scenario of faults structured only according to the thickness of the brittle crust is likely
191 modulated by inherited crustal heterogeneities, resulting in some relay zones being
192 larger than others. Sandbox experiments can also present such variability with larger
193 complexities randomly appearing (Fig. S14), suggesting that the initial orientation of the

194 S-shears might also influence the size of the relay zones. Other studies have proposed
195 that the off-fault deformation of earthquakes could be related to the reactivation of these
196 Riedels (Hatem et al. 2017), although we could not observe such reactivation in our
197 experiments. Moreover, in nature evidence of the initial Riedels structures is tenuous, as
198 they accommodate a very short amount of deformation before being passively offset by
199 the main active fault, which explains their weak morphological expression at the surface.
200 These results show that the first stages of deformation determine the long-term scaling
201 of the of strike-slip faults segmentation: Riedels create a regular pattern of geometrical
202 complexities, which is inherited throughout the fault history (Fig. 4.b). This pattern of
203 segmentation does not seem to experience significant changes even after the Riedels
204 have ceased to be active, as we keep seeing it on the coseismic-rupture map. Thus, this
205 questions the paradigm that faults become even more linear while accumulating
206 displacement through time, and thus prone to larger earthquakes. Instead, strike-slip
207 fault geometry may retain a constant level of complexity, which controls the size of
208 earthquakes.

209

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315

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321

322 **FIGURES**

323 Figure 1: **a.** Surface rupture map of the 1992 Landers earthquake, California.
324 Segmentation (black lines), according to Klinger (2010). **b.** Evolutionary scheme of the
325 formation of strike-slip faults, seen from the surface, and a 3D sketch to illustrate the
326 helicoidal-shape of the Riedels. **c.** Top view of the sandbox, with the measured
327 parameters, S: inter-Riedel distance, L: Riedel length, α : angle between the Riedel and the
328 basement fault.

329 Figure 2: Comparison between inter-Riedel distance (S) and sandpack thickness (T) for
330 different frictional and cohesive properties of the sand. **a.** Pictures of the top of the sandbox

331 after ~11-12 mm of displacement for experiments with a 2, 3 and 4 cm-thick sandpack. **b., c.,**
332 **d., e.** Average inter-Riedel distance (S) per experiment versus sand-pack thickness (T) for
333 Sand 1 (b.); for three different sands with different internal friction (**c**); for three different
334 basal materials to study the basal friction effect (**d**); and for two sands with different cohesion
335 (**e**). Note that for clarity's sake, for Sand 1 on graphics **c** and **d**, we represent the average
336 value by thickness rather than by experiment.

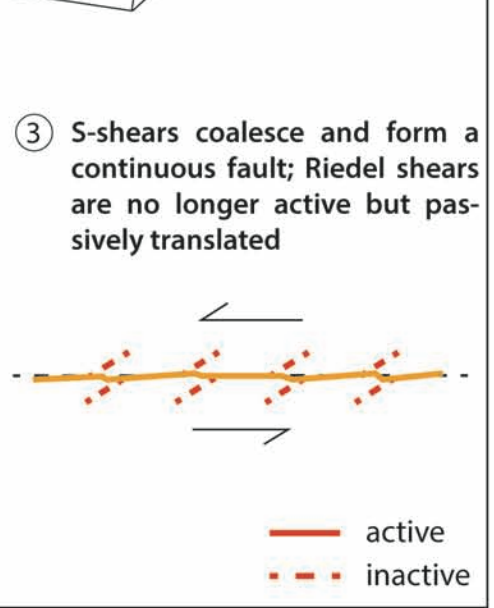
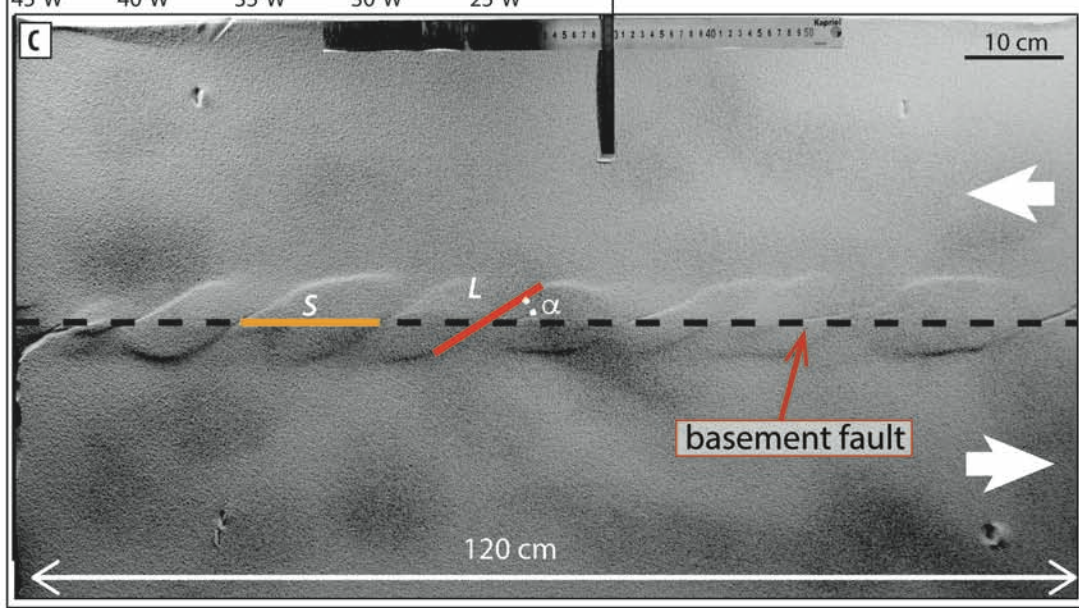
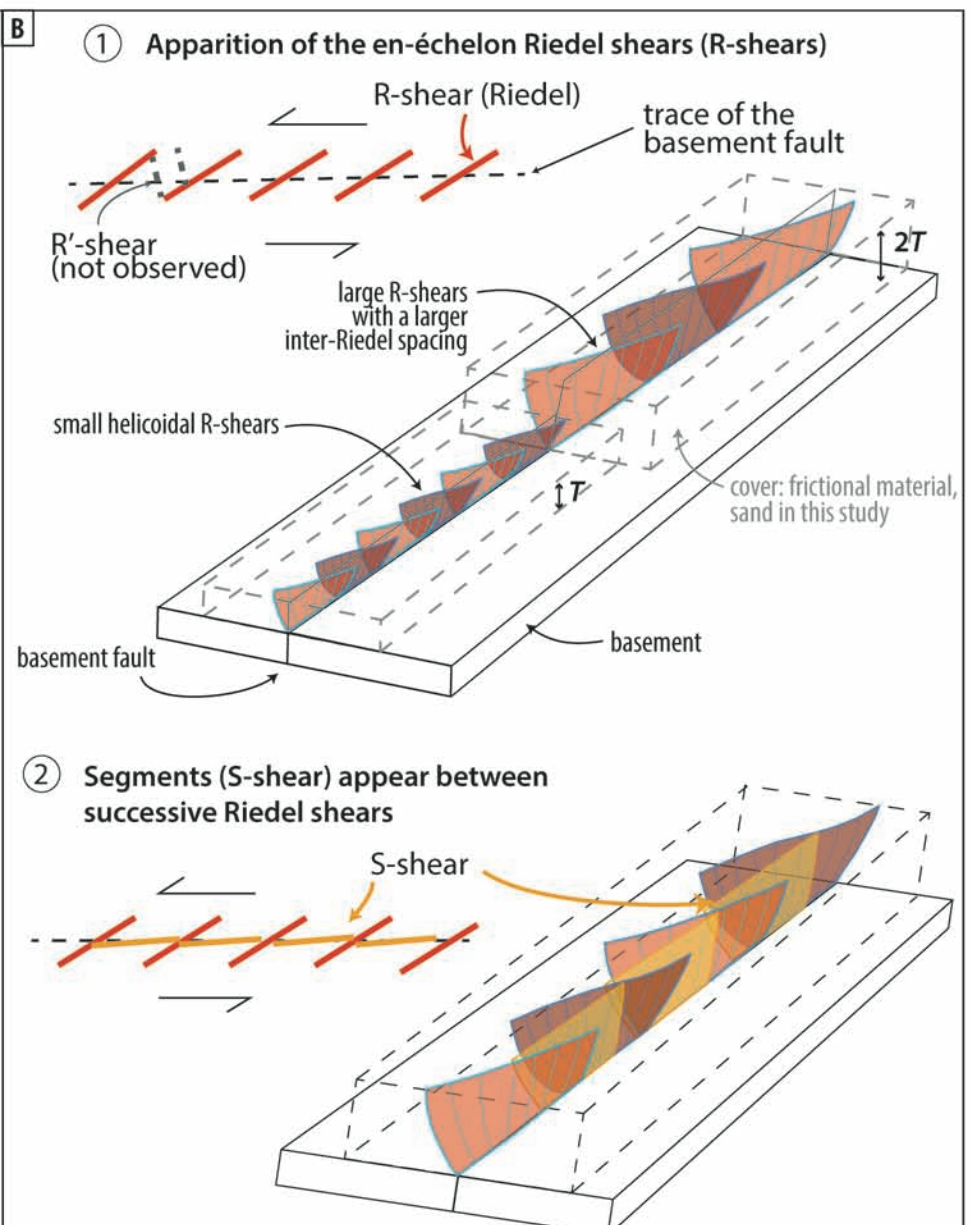
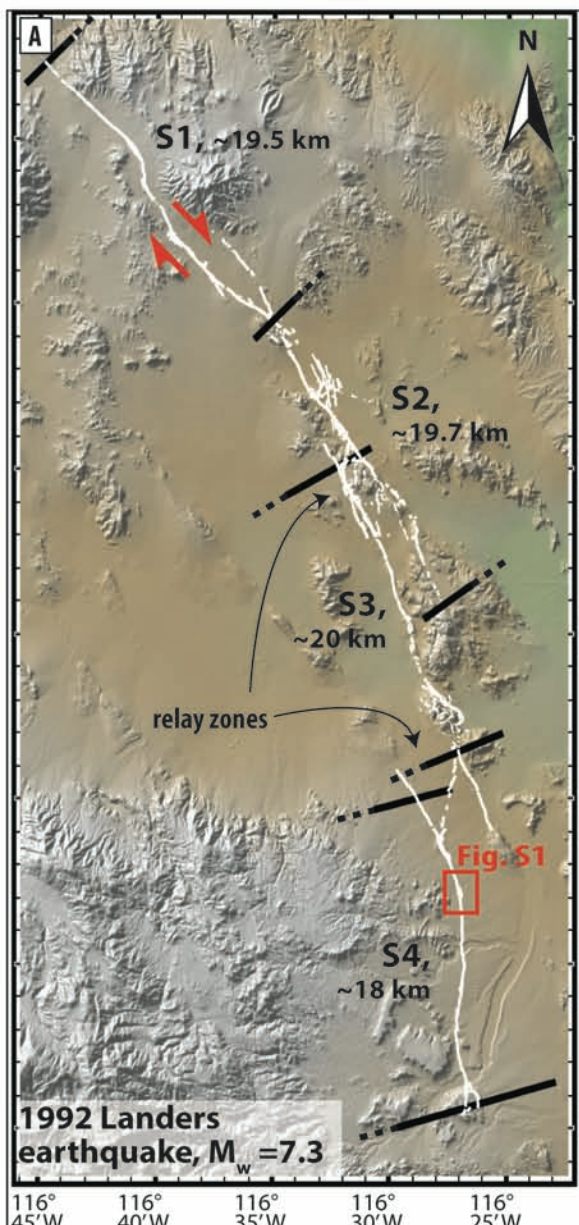
337 Figure 3 : Normalised length of individual fault segments for several continental strike-
338 slip earthquakes (Klinger, 2010, Lauer et al., 2018) and for individual analogue-fault
339 segments. Sand 1 is shown only for comparison. The length is normalized to inferred
340 seismogenic thickness or sand thickness accordingly. Values of seismogenic crust
341 thicknesses are estimated from the literature (see Supp. Mat.)

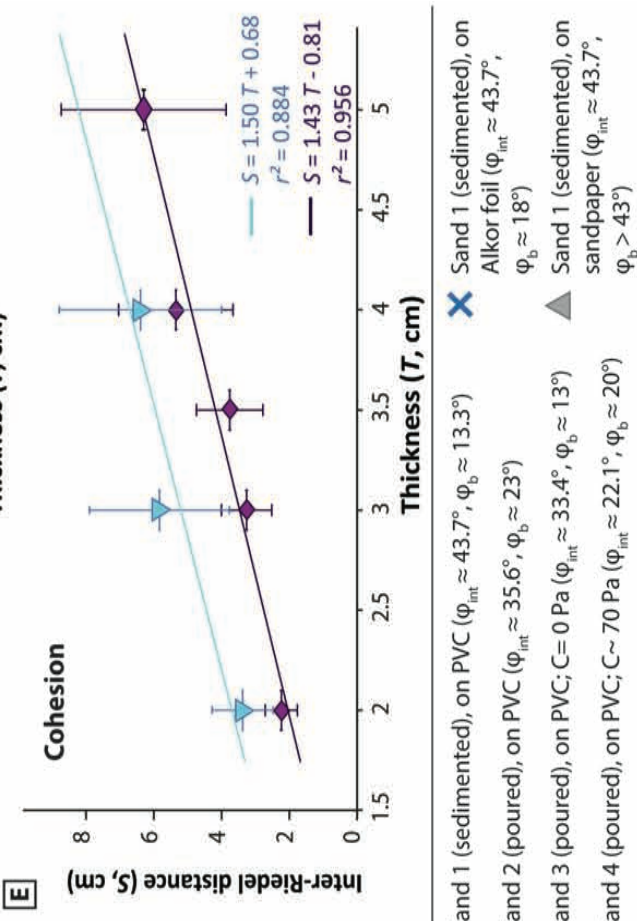
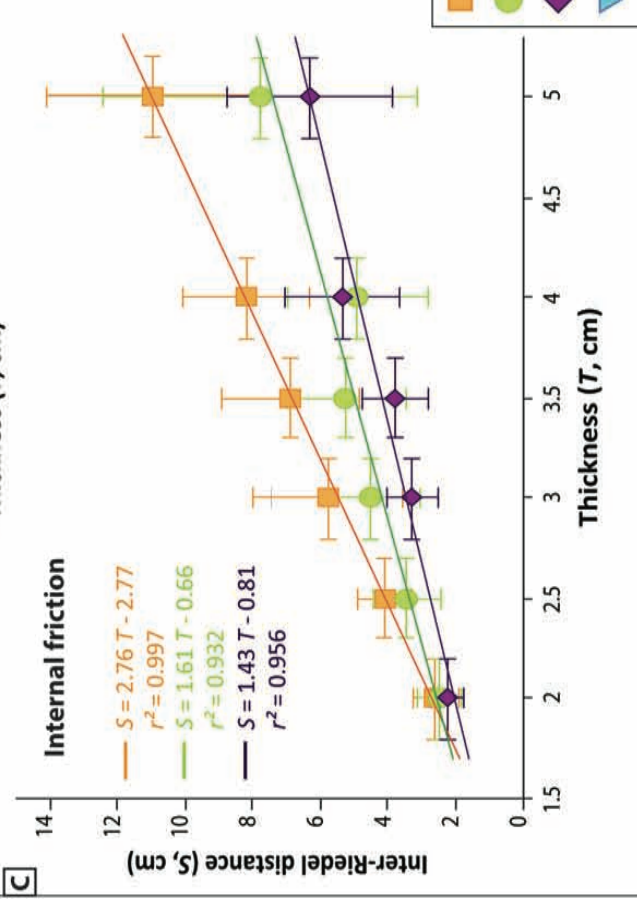
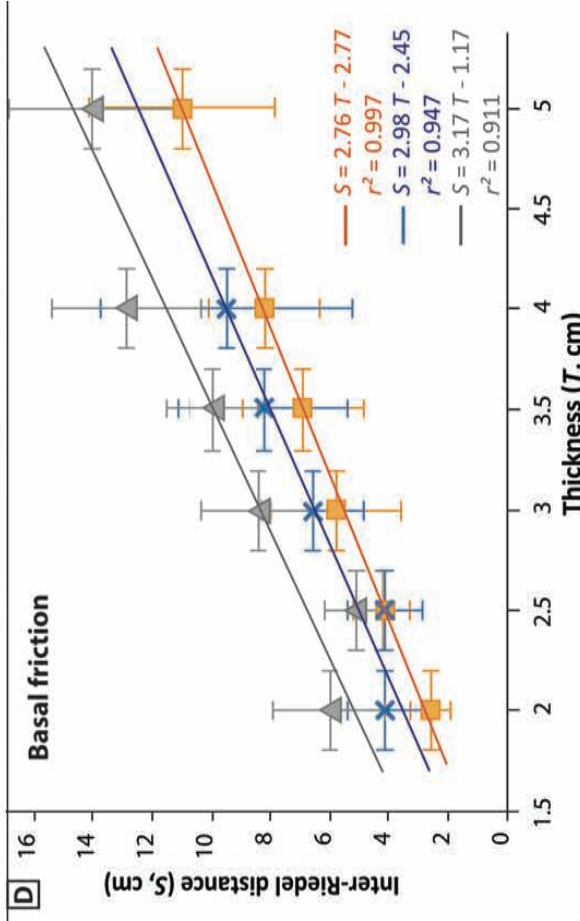
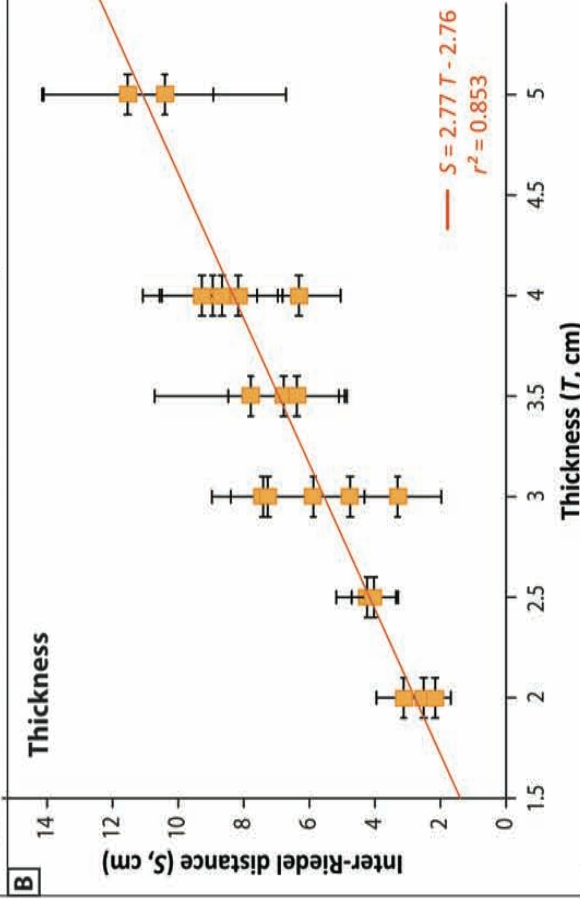
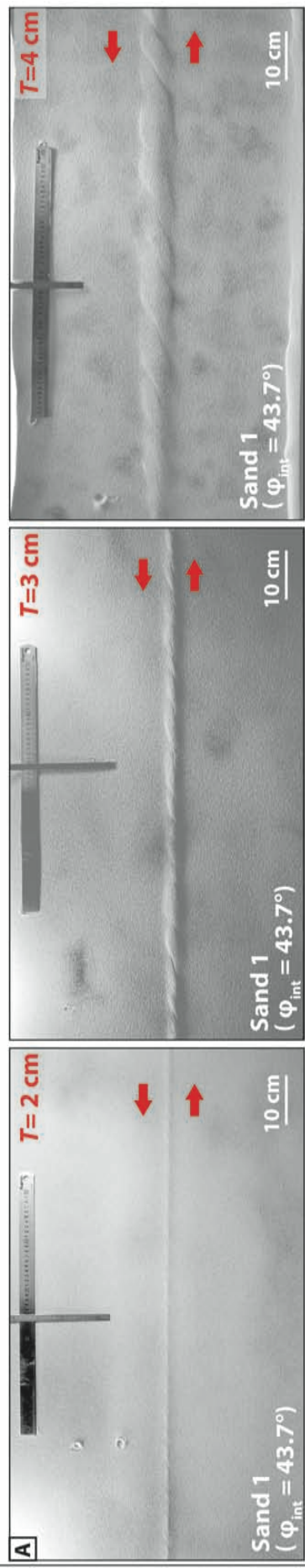
342 Figure 4: **a.** Curl of the incremental displacement field for a 4 cm-thick experiment (T=4cm)
343 at three different stages of the deformation (d=7.5mm, d=10mm, d=13mm). The curl is
344 derived from the correlation image between two successive pictures of the sandbox surface. It
345 highlights the location of shear. The location of the different active faults (R-shear, S-shear) is
346 highlighted. **b.** Scheme of the 3D-structures of a strike-slip fault at the crustal scale.

347

348 Figure 3 : Normalised length of individual fault segments for several continental strike-
349 slip earthquakes (Klinger, 2010, Lauer et al., 2018) and for individual analogue-fault
350 segments (see Fig. S13 for comparison with sands 1 and 4). The length is normalized to
351 inferred seismogenic thickness or sand thickness accordingly. Values of seismogenic
352 crust thicknesses are estimated from the literature (see Supp. Mat.)

353





- Sand 1 (sedimented), on PVC ($\phi_{int} \approx 43.7^\circ$, $\phi_b \approx 13.3^\circ$)
- Sand 2 (poured), on PVC ($\phi_{int} \approx 35.6^\circ$, $\phi_b \approx 23^\circ$)
- ◆ Sand 3 (poured), on PVC; C=0 Pa ($\phi_{int} \approx 33.4^\circ$, $\phi_b \approx 13^\circ$)
- ▲ Sand 4 (poured), on PVC; C~70 Pa ($\phi_{int} \approx 22.1^\circ$, $\phi_b \approx 20^\circ$)
- ◆ Sand 1 (sedimented), on Alkor foil ($\phi_{int} \approx 43.7^\circ$, $\phi_b \approx 18^\circ$)
- ▲ Sand 1 (sedimented), on sandpaper ($\phi_{int} \approx 43.7^\circ$, $\phi_b > 43^\circ$)

