



HAL
open science

Local microscopic behavior for 2D Coulomb gases

Thomas Leblé

► **To cite this version:**

| Thomas Leblé. Local microscopic behavior for 2D Coulomb gases. 2020. hal-03040308

HAL Id: hal-03040308

<https://hal.science/hal-03040308>

Preprint submitted on 4 Dec 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Local microscopic behavior for 2D Coulomb gases

Thomas Leblé*

October 10, 2016

Abstract

The study of two-dimensional Coulomb gases lies at the interface of statistical physics and non-Hermitian random matrix theory. In this paper we give a large deviation principle (LDP) for the empirical fields obtained, under the canonical Gibbs measure, by zooming around a point in the bulk of the equilibrium measure, up to the finest averaging scale $N^{-1/2+\varepsilon}$. The rate function is given by the sum of the “renormalized energy” of Serfaty *et al.* weighted by the inverse temperature, and of the specific relative entropy. We deduce a local law which quantifies the convergence of the empirical measures of the particles to the equilibrium measure, up to the finest scale.

1 Introduction

1.1 General setting

We consider a system of N classical charged point particles in the Euclidean space \mathbb{R}^2 with pairwise logarithmic interaction, in a confining potential V , and associate to any N -tuple $\vec{X}_N = (x_1, \dots, x_N)$ the energy

$$(1.1) \quad h_N(\vec{X}_N) := \sum_{1 \leq i \neq j \leq N} -\log |x_i - x_j| + N \sum_{i=1}^N V(x_i), \quad x_i \in \mathbb{R}^2.$$

We only impose mild conditions on the potential V (see Assumption 1).

For any value of the *inverse temperature* parameter $\beta > 0$ we consider the associated N -point Gibbs measure, which is absolutely continuous with respect to the Lebesgue measure on $(\mathbb{R}^2)^N$ with a density given by

$$(1.2) \quad d\mathbb{P}_N^\beta(\vec{X}_N) := \frac{1}{Z_{N,\beta}} e^{-\frac{1}{2}\beta h_N(\vec{X}_N)} d\vec{X}_N,$$

where we denote a N -tuple of points by $\vec{X}_N = (x_1, \dots, x_N)$ and $d\vec{X}_N := dx_1 \dots dx_N$. The constant $Z_{N,\beta}$ is a normalizing constant, also called the *partition function*, so that the total mass of \mathbb{P}_N^β is 1. The factor $\frac{1}{2}$ in front of the physical inverse temperature parameter β is chosen in order to match the usual convention in random matrix theory. Let us also emphasize the fact that the external potential scales like N , which corresponds to a mean-field/small temperature regime.

*Sorbonne Universités, UPMC Univ. Paris 06 and CNRS, UMR 7598, Laboratoire Jacques-Louis Lions, F-75005, Paris. E-mail leble@ann.jussieu.fr

Motivations. The model described by (1.1) and (1.2) is known in statistical physics as a *two-dimensional Coulomb gas*, *two-dimensional log-gas* or *two-dimensional one-component plasma*, we refer e.g. to [AJ81], [JLM93], [SM76] for a physical treatment of its main properties.

When $\beta = 2$ and V is quadratic, the probability measure (1.2) coincides with the joint law of eigenvalues of a non-Hermitian matrix model known as the *complex Ginibre ensemble*, which is obtained by sampling a $N \times N$ matrix whose coefficients are (properly normalized) i.i.d. complex Gaussians, see [Gin65]. For $\beta = 2$, more general potentials can be considered, which are associated to “random normal matrices” (see e.g. [AHM15]). Systems of particles with a logarithmic interaction as in (1.1), called *log-gases*, have also (and mostly) been studied on the real line, motivated by their link with Hermitian random matrix theory. We refer to [For10] for a survey of the connection between log-gases and random matrix theory, and in particular to [For10, Chap.15] for the two-dimensional (non-Hermitian) case.

The Ginibre case (and the case $\beta = 2$ in general) has the special property that the point process associated to \mathbb{P}_N^β becomes *determinantal*, which allows for an exact computation of many interesting quantities, e.g. the n -point correlation functions. The existence of a matrix model also allows for local laws at the microscopic scale as in [BYY14a, BYY14b]. In the present paper we rather work with general $\beta > 0$ and potential V , thus dealing with what could be called *two-dimensional β -ensembles* by analogy with the one-dimensional β -ensembles which generalize the laws of eigenvalues of random Hermitian matrices (see e.g. [DE02]). The microscopic behavior of one-dimensional β -ensembles has been recently investigated in [BEY12, BEY14] and we aim at a similar understanding in the two-dimensional case.

First-order results: the macroscopic behavior. Let us first recall some results about the macroscopic behavior of the particle system as $N \rightarrow \infty$.

If the potential V has some regularity and grows fast enough at infinity (see Assumption 1) there is an associated *equilibrium measure* μ_{eq} , such that the sequence $\{\mu_N\}_N$ (where $\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{x_i}$ denotes the *empirical measure* of the points) converges almost surely to μ_{eq} . Moreover the law of $\{\mu_N\}_N$ satisfies a Large Deviation Principle (LDP) at speed $\frac{\beta}{2}N^2$ on the space $\mathcal{P}(\mathbb{R}^2)$ of probability measures, with good rate function given by

$$(1.3) \quad I(\mu) := \iint -\log|x-y|d\mu(x)d\mu(y) + \int V(x)d\mu(x),$$

(see [PH98], and [BAZ98] for a closely related result). This characterizes the *first-order* or *macroscopic* behavior of the interacting particle system. Typically, as N becomes large, the N points x_1, \dots, x_N arrange themselves according to the probability density $d\mu_{\text{eq}}$, which has compact support Σ . Events that deviate from this prediction occur only with \mathbb{P}_N^β -probability of order $\exp(-N^2)$. We refer to [Ser15, Chap.2] and the references therein for a detailed exposition.

Microscopic behavior with macroscopic average. In this section we summarize the main result of [LS15], which describes the behavior as $N \rightarrow \infty$ of a *microscopic* quantity obtained through a *macroscopic average*.

Let \mathcal{X} be the set of locally finite point configurations in \mathbb{R}^2 , endowed with the topology of vague convergence, and let us denote by $\mathcal{P}(\mathcal{X})$ the set of Borel probability measures on \mathcal{X} i.e. the set of random point processes on \mathbb{R}^2 (we refer to Section 2.2 for more details).

In [LS15] (following the line of work [SS12], [SS15b], [SS15a], [RS15], [PS15]) S. Serfaty and the author have investigated the microscopic behavior of the system by making a statement on the point processes arising when zooming in by a factor $N^{1/2}$ (which is the typical inter-particle distance) and averaging over translations in a way that we now briefly present.

For any N -tuple \vec{X}_N , let $x'_i = N^{1/2}x_i$, $\nu'_N := \sum_{i=1}^N \delta_{x'_i}$, let $\Sigma' := N^{1/2}\Sigma$ denote the support of μ_{eq} after rescaling, and let i_N be the map $i_N : (\mathbb{R}^2)^N \rightarrow \mathcal{P}(\mathcal{X})$ defined by

$$(1.4) \quad i_N(\vec{X}_N) := \frac{1}{|\Sigma'|} \int_{\Sigma'} \delta_{\theta_{z'} \cdot \nu'_N} dz',$$

where $\theta_{z'} \cdot$ denotes the action of translation by $z' \in \Sigma'$, and where δ is the Dirac mass. The map i_N transforms a N -tuple of points into the data of all the blown-up point configurations obtained by zooming in by a factor $N^{1/2}$ around any $z \in \Sigma$. Such quantities are called *empirical fields*. Empirical fields (in our case, those could also be called ‘‘Gibbs random fields’’), appear in mathematical statistical physics as natural objects to encode the microscopic behavior of the systems, see e.g. [FO88, Oll88, Geo93].

We let $\mathfrak{P}_{N,\beta}$ be the push-forward of \mathbb{P}_N^β by i_N . The main result of [LS15] gives a large deviation principle for $\{\mathfrak{P}_{N,\beta}\}_N$ at speed N , on the space of stationary random point processes. The rate function on this subset of $\mathcal{P}(\mathcal{X})$ is given by

$$(1.5) \quad \mathcal{F}_\beta(P) := \frac{\beta}{2} \mathbf{E}_P[\mathbb{W}] + \text{ent}[P|\mathbf{\Pi}^1],$$

where \mathbb{W} is an energy functional which will be defined later, \mathbf{E}_P denotes the expectation under P , and $\text{ent}[P|\mathbf{\Pi}^1]$ is the *specific relative entropy* of P with respect to the Poisson point process of intensity 1 in \mathbb{R}^2 (see Section 2.5). In (1.5) the energy term is weighted by the inverse temperature β , it could easily be re-written following the usual convention of weighting the entropy term by $T = \frac{1}{\beta}$. A similar free energy functional is derived in a fairly general setting in [GZ93, Theorem 3.1], *if one knows a priori how to identify the limit energy term*, which turns out to be the central difficulty.

This LDP characterizes the microscopic behavior only in an *averaged* way, because of the average over translations in the definition of i_N . In fact (this is still a consequence of [LS15, Theorem 1]) this description can be enhanced by replacing the average over translations in Σ' by an average over translations in arbitrary small macroscopic regions (seen in blown-up scale), for example the square $C(z'_0, \varepsilon N^{1/2})$, where $\varepsilon > 0$ is fixed and $z'_0 = N^{1/2}z_0$ for some z_0 in the interior of Σ (the bulk). Let us emphasize that the average still takes place at the macroscopic scale $N^{1/2}$. This is done in [LS15] by considering ‘‘tagged’’ empirical fields which are elements of $\mathcal{P}(\Sigma \times \mathcal{X})$ keeping track of the point around which the configuration has been zoomed, thus allowing for a macroscopic localization.

Finally, let us emphasize that proving a LDP implies that the law of observable concentrates around minimizers of the rate function. If uniqueness of the minimizer is unknown (which is, so far, the case in our setting), the LDP does not characterize the limit *in law* of the observable.

Microscopic behavior with mesoscopic averages. The goal of this paper is to push further the analysis of [LS15] at finer scales and to consider *mesoscopic* versions of the map i_N . In other terms we look at the empirical fields obtained by averaging over translations in $C(z'_0, N^\delta)$ for $0 < \delta < 1/2$, and we obtain a LDP at speed $N^{2\delta}$ with essentially the same rate function as above. It is crucial to average over a relatively large set and although one might hope to go down to even finer scales (e.g. $O(\log^k N)$ for k large enough) we do not expect a similar result to hold for a strictly speaking *microscopic* average at scale $O(1)$ (in blown-up coordinates).

The first-order results show that the empirical measure $\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{x_i}$ converges to the equilibrium measure μ_{eq} almost surely. As a consequence of our analysis we get a ‘‘local law’’ (borrowing the terminology of [BYY14a] and [TV15, Theorem 20]) which implies that μ_N and μ_{eq} are close at small scales with very high probability.

The paper [BBNY15], which appeared simultaneously, also proves a local law for two-dimensional Coulomb gases. It does not study the large deviations of the empirical fields but obtain better (in fact, optimal) bounds on the errors (it can be understood as a “moderate deviations” result at all scales, whereas our local law only deals with the large deviations scale). The proof of [BBNY15] relies on a multiscale/bootstrap analysis, with very different methods. First, they establish a local law with error bounds comparable to ours: in that step, the “screening” method of the present paper is replaced by a fine understanding of the properties of potential-theoretic equilibrium measures with perturbed external fields. Then, the technology of loop equations is used to obtain optimal bounds on the error. In a forthcoming paper [LS], we recover such moderate deviations (as well as a central limit theorem for fluctuations), by different means.

1.2 Preliminary notation and definitions

1.2.1 Notation

For $R > 0$ we denote by C_R the square $[-R/2, R/2]^2$ and by $C(z, R)$ the translate of C_R by $z \in \mathbb{R}^2$. We denote by $D(p, r)$ the disk of center p and radius $r > 0$. If N is fixed and $\vec{X}_N \in (\mathbb{R}^2)^N$ we denote by $\nu_N := \sum_{i=1}^N \delta_{x_i}$ and $\nu'_N := \sum_{i=1}^N \delta_{x'_i}$ (where $x'_i = N^{1/2}x_i$).

Let $0 < \delta < 1/2$. We say that an event \mathcal{A} occurs with δ -overwhelming probability if

$$\limsup_{N \rightarrow \infty} N^{-2\delta} \log \mathbb{P}_N^\beta(\mathcal{A}^c) = -\infty,$$

where \mathcal{A}^c is the complement of \mathcal{A} . In particular, for any event \mathcal{B} , if \mathcal{A} occurs with δ -overwhelming probability we have

$$\limsup_{N \rightarrow \infty} N^{-2\delta} \log \mathbb{P}_N^\beta(\mathcal{B}) = \limsup_{N \rightarrow \infty} N^{-2\delta} \log \mathbb{P}_N^\beta(\mathcal{B} \cap \mathcal{A}),$$

and the same goes for the \liminf . In other terms, when evaluating probabilities of (logarithmic) order $N^{2\delta}$ we may restrict ourselves to the intersection with any event of δ -overwhelming probability.

If $\{a_N\}_N, \{b_N\}_N$ are two sequences of non-negative real numbers, we will write $a_N \preceq b_N$ if there exists $C > 0$ such that $a_N \leq Cb_N$ (\mathbb{P}_N^β -a.s. if the numbers are random), and we will write $a_N \preceq_\delta b_N$ if there exists $C > 0$ such that $a_N \leq Cb_N$ with δ -overwhelming probability.

We will write $a_N \ll N^\delta$ if there exists $\tau > 0$ such that $a_N \leq N^{\delta-\tau}$ (\mathbb{P}_N^β -a.s. if the numbers are random) and $a_N \ll_{\delta'} N^\delta$ if there exists $\tau > 0$ such that $a_N \leq N^{\delta-\tau}$ with δ' -overwhelming probability.

1.2.2 Equilibrium measure and splitting of the energy

Under mild hypotheses on V (see Assumption 1) it is known (see e.g. [ST97, Chap.1]) that there exists a probability measure μ_{eq} with compact support Σ which is the unique minimizer of I (as in (1.3)) over $\mathcal{P}(\mathbb{R}^2)$ (the set of probability measures). Defining ζ as

$$(1.6) \quad \zeta(x) := \int_{\mathbb{R}^2} -\log|x-y|d\mu_{\text{eq}}(x) + \frac{V}{2} - \left(\iint_{\mathbb{R}^2} -\log|x-y|d\mu_{\text{eq}}(x)d\mu_{\text{eq}}(y) - \frac{1}{2} \int_{\mathbb{R}^2} Vd\mu_{\text{eq}} \right),$$

we have $\zeta \geq 0$ quasi-everywhere (q.e.) in \mathbb{R}^2 and $\zeta = 0$ q.e. on Σ , and in fact this characterizes μ_{eq} uniquely, see [Fro35]. If $N \geq 1$ is fixed we let $\mu'_{\text{eq}}(x) := \mu_{\text{eq}}(xN^{-1/2})$.

If \mathcal{C} is a finite point configuration we define the second-order energy functional $w_N(\mathcal{C})$ as

$$(1.7) \quad w_N(\mathcal{C}) := \iint_{\Delta^c} -\log|x-y|(d\mathcal{C} - d\mu'_{\text{eq}})(x)(d\mathcal{C} - d\mu'_{\text{eq}})(y),$$

where Δ^c denotes the complement of the diagonal Δ . It computes the electrostatic interaction of the electric system made of the point charges in \mathcal{C} and a negatively charged background of density μ'_{eq} , without the infinite self-interactions of the point charges.

Let $\tilde{\zeta}(\mathcal{C}) := \int \zeta d\mathcal{C}$. It was proven in [SS15b] (see also [Ser15, Chap.3]) that the following *exact splitting formula* holds:

Lemma 1.1. *For any $N \geq 1$ and any $\vec{X}_N \in (\mathbb{R}^2)^N$ we have, with I as in (1.3)*

$$(1.8) \quad h_N(\vec{X}_N) = N^2 I(\mu_{\text{eq}}) - \frac{N \log N}{2} + w_N(\nu'_N) + 2N\tilde{\zeta}(\nu_N).$$

We may thus re-write the Gibbs measure \mathbb{P}_N^β as

$$(1.9) \quad d\mathbb{P}_N^\beta(\vec{X}_N) = \frac{1}{K_{N,\beta}} e^{-\frac{1}{2}\beta(w_N(\nu'_N) + 2N\tilde{\zeta}(\nu_N))} d\vec{X}_N,$$

where $K_{N,\beta}$ is a new normalizing constant. The exponent $(w_N(\nu'_N) + 2N\tilde{\zeta}(\nu_N))$ is expected to be typically of order N , and it was proven in [LS15, Cor. 1.5] that $\log K_{N,\beta} = -N \min \bar{\mathcal{F}}_\beta + o(N)$, where $\bar{\mathcal{F}}_\beta$ is closely related to the free energy functional \mathcal{F}_β mentioned in (1.5).

1.2.3 Energy and entropy

Renormalized energy. In [LS15], following [SS12, SS15b, RS15, PS15], an energy functional is defined at the level of random stationary point processes (see also [Ser15, Chap.3-6]), which is the Γ -limit of $\frac{1}{N}w_N$ as $N \rightarrow \infty$. We will define it precisely in Section 2.4 and we denote it by \mathbb{W}_m (where $m \geq 0$ is a parameter - the notation differs slightly from that of [LS15] where it corresponds to $\tilde{\mathbb{W}}_m$). It can be thought of as the infinite-volume limit of (1.7) and as a way of computing the interaction energy of an infinite configuration of point charges \mathcal{C} together with a negatively charged background of constant density m .

Specific relative entropy. For any $m \geq 0$ we let $\mathbf{\Pi}^m$ be the law of a Poisson point process of intensity m in \mathbb{R}^2 . Let P be a stationary random point process on \mathbb{R}^2 . The relative specific entropy $\text{ent}[P|\mathbf{\Pi}^m]$ of P with respect to $\mathbf{\Pi}^m$ is defined by

$$(1.10) \quad \text{ent}[P|\mathbf{\Pi}^m] := \lim_{R \rightarrow \infty} R^{-2} \text{Ent} \left(P|_{C_R} | \mathbf{\Pi}|_{C_R}^m \right),$$

where $P|_{C_R}$ denotes the random point process induced in C_R , and $\text{Ent}(\cdot|\cdot)$ denotes the usual relative entropy (or Kullback-Leibler divergence) of two probability measures defined on the same probability space. A sub-additivity argument ensures that the limit in (1.10) exists, see e.g. [GZ93]. We take the appropriate sign convention for the entropy so that it is non-negative: if μ, ν are two probability measures defined on the same space we let $\text{Ent}(\mu|\nu) := \int \log \frac{d\mu}{d\nu} d\mu$ if μ is absolutely continuous with respect to ν and $+\infty$ otherwise. For more details we refer to Section 2.5.

1.2.4 Good control on the energy

In this paragraph we define the notion of “good control at scale δ ”, which expresses the fact that our particle system has good properties in any square of sidelength N^δ (after blow-up). The assumption that good control at scale δ holds will be a key point in order to prove the LDP at slightly smaller scales. Moreover we will see that the “good control” assumption can be bootstrapped, i.e. good control at scale δ implies good control at scale δ_1 for $\delta_1 < \delta$ large enough.

Let us first introduce the local electric field E^{loc} and its truncation E_η^{loc} , we will come back to these definitions in more detail in Section 2.3. If \vec{X}_N is a N -tuple of points in \mathbb{R}^2 , for any $0 < \eta < 1$ we denote by $\nu'_{N,\eta}$ the measure $\nu'_{N,\eta} := \sum_{i=1}^N \delta_{x'_i}^{(\eta)}$, where $\delta_{x'_i}^{(\eta)}$ denotes the uniform probability measure on the circle of center x'_i and radius η . We let E^{loc} be the associated “local electric field”

$$(1.11) \quad E^{\text{loc}}(x) := (-\nabla \log) * (\nu'_N - \mu'_{\text{eq}})$$

and E_η^{loc} its truncation at scale η , defined by

$$(1.12) \quad E_\eta^{\text{loc}}(x) := (-\nabla \log) * (\nu'_{N,\eta} - \mu'_{\text{eq}}).$$

Finally, we denote by $\overset{\circ}{\Sigma}$ the interior of Σ .

Definition 1.2. *For any $0 < \delta \leq \frac{1}{2}$ we say that a good control at scale δ holds if for any $z_0 \in \overset{\circ}{\Sigma}$ and any $0 < \delta_1 < \delta$ we have with δ_1 -overwhelming probability:*

1. *The number of (blown-up) points $\mathcal{N}_\delta^{z_0}$ in the square $C(z'_0, N^\delta)$ is of order $N^{2\delta}$*

$$(1.13) \quad \mathcal{N}_\delta^{z_0} \preceq_{\delta_1} N^{2\delta}.$$

2. *For any $0 < \eta < 1$ we have*

$$(1.14) \quad \int_{C_\delta^{z_0}} |E_\eta^{\text{loc}}|^2 + \mathcal{N}_\delta^{z_0} \log \eta \preceq_{\delta_1} N^{2\delta},$$

which expresses the fact that the energy in the square $C(z'_0, N^\delta)$ (after blow-up) is of order $N^{2\delta}$.

Let us emphasize that (1.13), (1.14) control quantities at scale δ by looking at probabilities at scale $\delta_1 < \delta$. This is mostly a technical choice, which allows to simplify some arguments below.

1.3 Rate function

Let us define the *local* rate function as

$$(1.15) \quad \mathcal{F}_\beta^m(P) := \frac{\beta}{2} \mathbb{W}_m(P) + \text{ent}[P|\mathbf{\Pi}^m].$$

It is a *good rate function* because both terms are good rate functions (see e.g. [LS15, Lemma 4.1]).

For any $m > 0$ we let $\mathcal{P}_{s,m}(\mathcal{X})$ be the set of random stationary point processes of intensity m . Let us define a scaling map $\sigma_m : \mathcal{P}_s(\mathcal{X}) \rightarrow \mathcal{P}_s(\mathcal{X})$ such that $\sigma_m(P)$ is the push-forward of P by $\mathcal{C} \mapsto m^{-1/2}\mathcal{C}$. It is easy to see that σ_m induces a bijection from $\mathcal{P}_{s,m}(\mathcal{X})$ to $\mathcal{P}_{s,1}(\mathcal{X})$ for any $m > 0$.

It is proven (see [LS15, Def. 2.4, Lemma 4.2]) that

Lemma 1.3. *The map σ_m induces a bijection between the minimizers of \mathcal{F}_β^m over $\mathcal{P}_{s,m}(\mathcal{X})$ and the minimizers of \mathcal{F}_β^1 over $\mathcal{P}_{s,1}(\mathcal{X})$.*

Let us observe that Lemma 1.3 crucially uses the scaling properties of the logarithmic kernel. Replacing the logarithmic interaction by, e.g. a Riesz potential $|x - y|^s$ (as studied in [LS15]) gives rise to a more intricate situation where the effect of the equilibrium measure may be felt at microscopic scale, not only through a scaling factor.

We may now state our main results.

1.4 Statement of the results

If $z_0 \in \mathring{\Sigma}$ and $0 < \delta < 1/2$ are fixed, let us define the map $i_{N,\delta}^{z_0} : \mathcal{X} \rightarrow \mathcal{P}(\mathcal{X})$ as

$$(1.16) \quad i_{N,\delta}^{z_0}(\mathcal{C}) := N^{-2\delta} \int_{C(z'_0, N^\delta)} \delta_{\theta_{z'_0}(\mathcal{C} \cap C(z'_0, N^\delta))} dz'.$$

Such quantities are called empirical fields. We denote by $\mathfrak{P}_{N,\beta,\delta}^{z_0}$ the law of the push-forward of \mathbb{P}_N^β by the map $i_{N,\delta}^{z_0}$ - in other words: the empirical field observed around z_0 by averaging at the mesoscopic scale N^δ . Finally we denote by m_{eq} the density of μ_{eq} (see Assumption 1).

Theorem 1. *For any $0 < \delta < 1/2$, for any z_0 in $\mathring{\Sigma}$, the sequence $\{\mathfrak{P}_{N,\beta,\delta}^{z_0}\}_N$ obeys a large deviation principle at speed $N^{2\delta}$ with good rate function $(\mathcal{F}_\beta^{m_{\text{eq}}(z_0)} - \min \mathcal{F}_\beta^{m_{\text{eq}}(z_0)})$.*

Moreover a good control at scale δ holds in the sense of Definition 1.2.

Theorem 1 tells us in particular that the behavior around $z_0 \in \mathring{\Sigma}$ depends on V only through the value $m_{\text{eq}}(z_0)$, and in view of Lemma 1.3 it has only the effect of scaling the configurations. Conditionally to the uniqueness of minimizers of \mathcal{F}_β , it would yield another example of the *universality* phenomenon: the small scale behavior of the particle system is essentially independent of the choice of V (however, since uniqueness of minimizers has not yet been proven, we can only rigorously speak of *universality of the large deviations functional*).

The first consequence of Theorem 1 is a bound on the discrepancy i.e. the difference between the number of points of ν'_N in a given square and the mass given by μ'_{eq} .

Corollary 1.4. *Let $z_0 \in \mathring{\Sigma}$, let $0 < \delta < 1/2$ and $\delta_1 \in (\delta/2, \delta)$. We have*

$$(1.17) \quad \left| \int_{C(z'_0, N^{\delta_1})} d\nu'_N - d\mu'_{\text{eq}} \right| \ll_{\delta_1} N^{\frac{4}{3}\delta}.$$

For $\delta_1 < \delta$ close to δ , the bound $N^{\frac{4}{3}\delta}$ on the difference is much smaller than the typical value of each term, of order $N^{2\delta_1}$. It allows us to prove a *local law* in the following sense:

Corollary 1.5. *Let $z_0 \in \mathring{\Sigma}$ and $0 < \delta \leq 1/2$ be fixed. Let f be a C^1 function (which may depend on N) such that f is supported in $C(z'_0, N^\delta)$. Then for any $\delta/2 < \delta_1 < \delta$ we have*

$$(1.18) \quad N^{-2\delta} \left| \int_{C(z'_0, N^\delta)} f(d\nu'_n - d\mu'_{\text{eq}}) \right| \preceq_{\delta_1} \|\nabla f\|_\infty N^{\delta_1} + \|f\|_\infty N^{-2\delta/3}.$$

In particular if $f(z) = \tilde{f}(N^{-\delta}(z - z'_0))$ for some compactly supported C^1 function \tilde{f} then $\|\nabla f\|_\infty \preceq N^{-\delta}$ and $\|f\|_\infty \preceq 1$, thus we get

$$N^{-2\delta} \left| \int_{C(z'_0, N^\delta)} f(d\nu'_n - d\mu'_{\text{eq}}) \right| \ll_{\delta_1} 1.$$

Comments and open questions. In the statement of the results we restrict ourselves to the following setting: we first pick a point z_0 in the interior of Σ (called the *bulk*) and then look at the point process in $C(z', N^\delta)$ with $N^{-1/2}z' = z_0$. A careful inspection of the proof shows that we might have taken z' depending on N more finely, e.g. by considering a sequence z' with $N^{-1/2}z' \rightarrow z_0 \in \overset{\circ}{\Sigma}$, while keeping the same conclusions. It does not seem possible to take $z' \rightarrow z_0 \in \partial\Sigma$ (the “edge case”) in general because the density m_{eq} may vanish near the boundary - however, this does not happen in the standard example of the quadratic potential, in which the density is constant up to the boundary of the support. Our analysis might be done in the edge case at a scale $\delta \geq \delta_c$ depending on the speed at which $m_{\text{eq}}(z)$ vanishes, but we do not pursue this goal here.

The minimizers of the rate function are unknown in general, however it is proven in [LS15, Corollary 1.4] that the Ginibre point process minimizes \mathcal{F}_β^1 over $\mathcal{P}_{s,1}(\mathcal{X})$ for $\beta = 2$. We do not know whether uniqueness of the minimizers holds for $\beta = 2$, nor for any value of β . Uniqueness of the minimizers for some $\beta > 0$ would imply that the empirical fields have a limit in law as $N \rightarrow \infty$, which would heuristically correspond to some “ β -Ginibre” random point process. In that case, our results shows that the hypothetical convergence

$$\text{Empirical field averaged at scale } N^\delta \rightarrow \beta\text{-Ginibre}$$

holds at arbitrarily fine scales $\delta > 0$, which would hint at the convergence in law of the non-averaged point process ν'_N to the conjectural β -Ginibre point process.

Another open question is the behavior of the minimizers as $\beta \rightarrow \infty$ (the low-temperature limit). The *crystallization conjecture* (see e.g. [BL15] for a review) predicts that the minimum of \mathbb{W}_1 on $\mathcal{P}_{s,1}(\mathcal{X})$ is (uniquely) attained by the random stationary point process associated to the triangular lattice. In the high-temperature limit $\beta \rightarrow 0$, it is proven in [Leb15, Theorem 2] that minimizers of \mathcal{F}_β^1 converge (in a strong sense) to $\mathbf{\Pi}^1$.

The result of [LS15] and most of the methods used in this paper are valid in a broader setting than the two-dimensional, logarithmic case, in particular we could think of treating the $1d$ log-gas (i.e. the β -ensembles). It turns out that an adaptation of the present method in the one-dimensional case allows one to improve the result of [LS15] to finer, mesoscopic scales, however, we have been unable so far to go down to the finest scale $N^{-1+\varepsilon}$ and we hope to return on this question in a subsequent work.

1.5 Plan of the paper and sketch of the proof

In Section 2 we introduce some notation and we give the definitions of the main objects used throughout the paper, as well as their key properties. In Section 3 we gather preliminary results about the energy w_N and we prove the main technical tool, called the “screening lemma”. Section 4 is devoted to the proof of a LDP upper bound and Section 4 to the lower bound. We combine these two steps to prove Theorem 1 in Section 6, together with Corollary 1.4 and 1.5. Section 7 is devoted to intermediate results which we postpone there.

Let us now sketch how the proof of Theorem 1 goes. The basic idea is a bootstrap argument, we find that there exists $t < 1$ such that

$$(1.19) \quad \text{Good control at scale } \delta \longrightarrow \begin{cases} \text{Large deviations at scale } \delta_1 \\ \text{Good control at scale } \delta_1 \end{cases} \quad \text{for all } t\delta \leq \delta_1 < \delta.$$

Once good control at scale $\delta = 1/2$ is established, Theorem 1 follows. A similar bootstrap argument was used in [RNS15] for studying the minimizers of w_N (which corresponds to the $\beta = +\infty$, or zero temperature case).

The main obstruction to obtaining LDP for empirical fields in our context is the non-locality of the energy (1.7): due to the long-range nature of the interactions, it is hard to localize the energy in a given square in such a way that it only depends of the point configuration in this square. Another way of seeing it is that $E^{\text{loc}}(x)$ depends *a priori* on the whole configuration \vec{X}_N and not only on the points close to x .

To prove (1.19) we rely on the following steps: let $z_0 \in \mathring{\Sigma}$ be fixed. For the sake of simplicity let us assume that $w_N(\vec{X}_N) = \frac{1}{2\pi} \int_{\mathbb{R}^2} |E^{\text{loc}}|^2$, where E^{loc} is the local electric field defined in Section 1.2.4 (see also Section 2.3).

1. For any \vec{X}_N , we split the energy $w_N(\vec{X}_N)$ as

$$\int_{\mathbb{R}^2} |E^{\text{loc}}|^2 = \int_{C(z'_0, N^{\delta_1})} |E^{\text{loc}}|^2 + \int_{C(z'_0, N^{\delta_1})^c} |E^{\text{loc}}|^2,$$

and we split \vec{X}_N as $\vec{X}^{\text{in}} + \vec{X}^{\text{out}}$ where \vec{X}^{in} is the point configuration in $C(z'_0, N^{\delta_1})$ and \vec{X}^{out} is the point configuration in $C(z'_0, N^{\delta_1})^c$.

2. We define $F^{\text{in}}(\vec{X}^{\text{in}})$ (resp. $F^{\text{out}}(\vec{X}^{\text{out}})$) as the minimal energy of an electric field associated to \vec{X}^{in} (resp. \vec{X}^{out}), see (4.13) (resp. (4.10)). We thus have

$$w_N(\vec{X}_N) \geq F^{\text{in}}(\vec{X}^{\text{in}}) + F^{\text{out}}(\vec{X}^{\text{out}}).$$

The two terms in the right-hand side become independent (they depend from two distinct sets of variables).

3. Inserting the previous inequality into the expression of the Gibbs measure (1.9) we obtain, for any event \mathcal{A} “concerning” \vec{X}^{in}

$$(1.20) \quad \mathbb{P}_N^\beta(\mathcal{A}) \leq \frac{1}{K_{N,\beta}} \left(\int_{\mathcal{A}} e^{-\frac{1}{2}\beta F^{\text{in}}(\vec{X}^{\text{in}})} d\vec{X}^{\text{in}} \right) \left(\int e^{-\frac{1}{2}\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + 2N\zeta(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right).$$

This can be used to prove a first LDP upper bound (taking $\mathcal{A} = \{i_{N,\delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)\}$) or a first “good control” estimate (taking $\mathcal{A} = \{F^{\text{in}}(\vec{X}^{\text{in}}) \gg N^{2\delta_1}\}$).

4. Then we need to prove that (1.20) is sharp (at scale δ_1). Given \vec{X}^{out} and \vec{X}^{in} , it amounts to be able to reconstruct (a family of) point configurations $\vec{X}_N \approx \vec{X}^{\text{out}} + \vec{X}^{\text{in}}$ such that $w_N(\vec{X}_N) \leq F^{\text{in}}(\vec{X}^{\text{in}}) + F^{\text{out}}(\vec{X}^{\text{out}}) + o(N^{2\delta_1})$. This is where the screening procedure is used: we modify \vec{X}^{out} and the associated electric field a little bit (this procedure follows the line of work [SS15b, SS15a, RS15, PS15] and is called *screening* for reasons that will appear later) so that we may glue together \vec{X}^{in} and the new \vec{X}^{out} and create an electric field compatible with the new (slightly modified) point configuration \vec{X}_N . It is then a general fact that $w_N(\vec{X}_N)$ (the energy of the *local* electric field associated to \vec{X}_N) is the smallest energy in a wide class of compatible electric fields. Let us emphasize that in this step, the good control on the energy is the key ingredient which allows to bootstrap the argument. Indeed, the screening procedure requires to find good boundaries where the energy is controlled, and the existence of such boundaries is ensured by the good control at a larger scale.

In particular, proving a partial converse to (1.20) allows us to estimate the “local partition function”

$$\frac{1}{K_{N,\beta}} \left(\int e^{-\frac{1}{2}\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + 2N\zeta(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right),$$

and also to show a LDP lower bound. Combined with the estimates of the previous step, it proves (1.19).

Acknowledgements. The author would like to thank his PhD supervisor, Sylvia Serfaty, for helpful comments on this work, as well as the anonymous referees for many useful corrections and comments.

2 Notation, assumptions and main definitions

2.1 Assumption on the potential

Assumption 1. *The potential V is such that*

1. V is lower semi-continuous (l.s.c.) and bounded below.
2. The set $\{x \in \mathbb{R}^2 \mid V(x) < \infty\}$ has positive logarithmic capacity.
3. We have $\lim_{|x| \rightarrow \infty} \frac{V(x)}{2} - \log|x| = +\infty$.

These first three conditions ensure that the equilibrium measure μ_{eq} is well-defined and has compact support Σ . Furthermore we ask that the measure μ_{eq} has a density m_{eq} which is κ -Hölder in Σ , for some $0 < \kappa \leq 1$

$$(2.1) \quad |m_{\text{eq}}(x) - m_{\text{eq}}(y)| \leq |x - y|^\kappa.$$

If V is C^2 , it is known that μ_{eq} is absolutely continuous with respect to the Lebesgue measure on \mathbb{R}^2 and its density coincides with ΔV in Σ . Thus in particular (2.1) is satisfied as soon as V is $C^{2,\kappa}$. Let us observe that the third assumption (that V is “strongly confining”) could be slightly relaxed into

$$\liminf_{|x| \rightarrow \infty} \frac{V(x)}{2} - \log|x| > -\infty,$$

i.e. V is only “weakly confining”, in which case the support Σ might not be compact (see [Har12] for a proof of the first-order LDP in this case). We believe that Theorem 1 should extend to the non-compact case as well, since it is really a *local* result, but we do not pursue this goal here.

2.2 Point configurations and point processes

Point configurations. If B is a Borel set of \mathbb{R}^2 we denote by $\mathcal{X}(B)$ the set of locally finite point configurations in B or equivalently the set of non-negative, purely atomic Radon measures on B giving an integer mass to singletons. We will often write \mathcal{C} for $\sum_{p \in \mathcal{C}} \delta_p$. We endow the set $\mathcal{X} := \mathcal{X}(\mathbb{R}^2)$ (and the sets $\mathcal{X}(B)$ for B Borel) with the topology induced by the topology of weak convergence of Radon measure (also known as vague convergence or convergence against compactly supported continuous functions), these topologies are metrizable (e.g. take the 1-Wasserstein distance for B compact, and a convergent serie of such distances over an increasing sequence of compact domains if B is not - in fact we will only encounter the cases $B = \mathbb{R}^2$ or $B = \text{some square}$) and we fix a compatible distance $d_{\mathcal{X}}$.

Volume of configurations. Let B be a Borel set of \mathbb{R}^2 . For any $N \geq 1$, let \sim_N be the equivalence relation on B^N defined as $(x_1, \dots, x_N) \sim_N (y_1, \dots, y_N)$ if and only if there exists a permutation $\sigma \in \mathfrak{S}_N$ (the symmetric group on N elements) such that $x_i = y_{\sigma(i)}$ for $i = 1, \dots, N$. We denote by B^N/\mathfrak{S}_N the quotient set and by π_N the canonical projection $B^N \rightarrow B^N/\mathfrak{S}_N$. The set of finite point configurations in B can be identified to $\{\emptyset\} \cup \bigcup_{N=1}^{+\infty} B^N/\mathfrak{S}_N$.

If $\mathcal{A} \subset B^N/\mathfrak{S}_N$ we define $\hat{\mathcal{A}} \subset B^N$ as $\hat{\mathcal{A}} := \bigcup_{\mathcal{C} \in \mathcal{A}} \mathcal{C}$. It is easy to see that $\hat{\mathcal{A}}$ is the largest subset of B^N such that the (direct) image of $\hat{\mathcal{A}}$ by π_N is \mathcal{A} .

We will call “the volume of \mathcal{A} ” and write (with a slight abuse of notation) $\mathbf{Leb}^{\otimes N}(\mathcal{A})$ the quantity $\mathbf{Leb}^{\otimes N}(\hat{\mathcal{A}})$ (where \mathbf{Leb} is the normalized Lebesgue measure on \mathbb{R}^2).

Random point process. A random point process is a probability measure on \mathcal{X} . We denote by $\mathcal{P}_s(\mathcal{X})$ the set of stationary random point processes i.e. those which are invariant under (push-forward by) the natural action of \mathbb{R}^2 on \mathcal{X} by translations. We endow $\mathcal{P}_s(\mathcal{X})$ with the topology of weak convergence of probability measures, and we fix a compatible distance $d_{\mathcal{P}(\mathcal{X})}$, e.g. the 1-Wasserstein distance. Throughout the text we will denote by $B(P, \varepsilon)$ the closed ball of center P and radius ε for $d_{\mathcal{P}(\mathcal{X})}$.

2.3 Electric systems and electric fields

Finite electric system. We will call an “electric system” a couple (\mathcal{C}, μ) where \mathcal{C} is a point configuration and μ is a non-negative measurable bounded function in \mathbb{R}^2 . One should think of \mathcal{C} as a set of point charges (carrying a +1 charge) and μ as the density of (negative) charges in the background. We say that the system is finite if \mathcal{C} is finite and μ is compactly supported. We say that the system is neutral if it is finite and $\int_{\mathbb{R}^2} d\mathcal{C} = \int_{\mathbb{R}^2} \mu(x) dx$.

Electric fields. Let $1 < p < 2$ be fixed. We define the set of electric fields Elec as the set of vector fields in $L_{\text{loc}}^p(\mathbb{R}^2, \mathbb{R}^2)$ such that

$$(2.2) \quad -\operatorname{div} E = 2\pi(\mathcal{C} - \mu) \text{ in } \mathbb{R}^2$$

for some electric system (\mathcal{C}, μ) . When (2.2) holds we say that E is compatible with (\mathcal{C}, μ) in \mathbb{R}^2 and we denote it by $E \in \text{Elec}(\mathcal{C}, \mu)$. If K is a compact subset of \mathbb{R}^2 with piecewise C^1 boundary we let $\text{Elec}(\mathcal{C}, \mu, K)$ be the set of electric fields which are compatible with (\mathcal{C}, μ) in K i.e. such that

$$-\operatorname{div} E = 2\pi(\mathcal{C} - \mu) \text{ in } K.$$

We denote by Elec^0 the set of decaying electric fields, such that $E(z) = O(|z|^{-2})$ as $|z| \rightarrow \infty$. We let $\text{Elec}^0(\mathcal{C}, \mu, K)$ be the set of electric fields which are compatible with \mathcal{C}, μ in K and decay. As an example of decaying electric fields, one could consider the local electric field (defined in (1.11)) associated to a neutral (hence finite) electric system.

Local electric fields. If (\mathcal{C}, μ) is a finite electric system there is a natural compatible electric field, namely the local electric field defined as

$$E^{\text{loc}} := -\nabla \log * (\mathcal{C} - \mu).$$

We also define the “local electric potential”

$$H^{\text{loc}} := -\log * (\mathcal{C} - \mu).$$

The scalar field H^{loc} corresponds physically to the electrostatic potential generated by the point charges of \mathcal{C} together with a background of “density” μ . The vector field E^{loc} can be thought of as the associated electrostatic field. It is easy to see that E^{loc} fails to be in L_{loc}^2 because it blows up like $|x|^{-1}$ near each point of \mathcal{C} , however E^{loc} is in $L_{\text{loc}}^p(\mathbb{R}^2, \mathbb{R}^2)$ for any $1 < p < 2$.

Truncation procedure. The renormalization procedure of [RS15, PS15] uses a truncation of the singularities which we now recall. We define the truncated Coulomb kernel as follows: for $0 < \eta < 1$ and $x \in \mathbb{R}^2$, let $f_\eta(x) = (-\log|x| + \log \eta)_+$. If (\mathcal{C}, μ) is an electric system and $E \in \text{Elec}(\mathcal{C}, \mu)$ we let

$$(2.3) \quad E_\eta(X) := E(X) - \sum_{p \in \mathcal{C}} \nabla f_\eta(X - p).$$

If E is the local field associated to some finite electric system (\mathcal{C}, μ) , then E_η is given by

$$E_\eta = -\log * \left(\sum_{p \in \mathcal{C}} \delta_p^{(\eta)} - d\mu \right),$$

where $\delta_p^{(\eta)}$ is the uniform measure on the circle of center p and radius η .

2.4 Renormalized energy

For finite point configurations. It follows from the definition, and the fact that $-\log$ is (up to an additive constant) the Coulomb kernel in dimension 2, that $-\Delta H^{\text{loc}} = 2\pi(\mathcal{C} - \mu)$, where H^{loc} denotes the local electric potential associated to a finite electric system (\mathcal{C}, μ) . We may thus observe that $w_N(\mathcal{C})$ (defined in (1.7)) can be written

$$w_N(\mathcal{C}) \approx -\frac{1}{2\pi} \int H^{\text{loc}} \Delta H^{\text{loc}}$$

(up to diagonal terms). Using $E^{\text{loc}} = \nabla H^{\text{loc}}$ and integrating by parts we obtain heuristically $w_N(\mathcal{C}) = \int_{\mathbb{R}^2} |E^{\text{loc}}|^2$. However, this computation does not make sense because E^{loc} fails to be in L^2 around each point charge (and indeed the diagonal is excluded in the definition of w_N). The correct way of giving a sense to “ $w_N(\mathcal{C}) \approx \int_{\mathbb{R}^2} |E^{\text{loc}}|^2$ ” is to use a *renormalization* procedure, using the truncation at scale η defined above. The following is proven in [SS15b].

Lemma 2.1. *For any $N \geq 1$ and any $\vec{X}_N \in (\mathbb{R}^2)^N$, we have*

$$(2.4) \quad w_N(\nu'_N) = \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \int_{\mathbb{R}^2} (|E_\eta^{\text{loc}}|^2 + N \log \eta).$$

Moreover w_N is bounded below on $(\mathbb{R}^2)^N$ by $O(N)$.

For infinite electric fields. Let (\mathcal{C}, μ) be an electric system and $E \in \text{Elec}(\mathcal{C}, \mu)$. We let $\mathcal{W}_\eta(E)$ be

$$\mathcal{W}_\eta(E) := \frac{1}{2\pi} \limsup_{R \rightarrow \infty} R^{-2} \int_{C_R} (|E_\eta|^2 + \mu \log \eta).$$

The renormalized energy of E is then defined as $\mathcal{W}(E) := \limsup_{\eta \rightarrow 0} \mathcal{W}_\eta(E)$.

For (random) infinite point configurations. If (\mathcal{C}, μ) is an electric system with μ constant equal to some $m > 0$ we define

$$(2.5) \quad \mathbb{W}_m(\mathcal{C}) = \inf_{E \in \text{Elec}(\mathcal{C}, \mu)} \mathcal{W}(E).$$

Similarly if P is a random point process we let $\mathbb{W}_m(P) = \mathbf{E}_P[\mathbb{W}_m]$ for any $m > 0$.

The following lower semi-continuity result was proven in [LS15, Lemma 4.1].

Lemma 2.2. *For any $m > 0$ the map $P \mapsto \mathbb{W}_m(P)$ is lower semi-continuous on $\mathcal{P}_s(\mathcal{X})$. Moreover its sub-level sets are compact. In particular, \mathbb{W}_m is bounded below.*

2.5 Specific relative entropy

For any $P \in \mathcal{P}_s(\mathcal{X})$, the specific relative entropy $\text{ent}[P|\mathbf{\Pi}^m]$ is defined as in (1.10) (see e.g. [GZ93, (2.12)]).

Lemma 2.3. *For any $m \geq 0$ the map $P \mapsto \text{ent}[P|\mathbf{\Pi}^m]$ is well-defined on $\mathcal{P}_s(\mathcal{X})$, it is affine lower semi-continuous and its sub-level sets are compact.*

Proof. We refer to [RAS09, Chap. 6] for a proof of these statements. □

Large deviations for the reference measure. If we throw $\approx mR^2$ points uniformly and independently in some large square C_R , we expect the associated empirical field to look like a Poisson point process. The following proposition states that the large deviations are governed by the specific relative entropy defined in (1.10). It can be understood as the “process-level” (or “type III”, in the language of large deviations) extension of the classical Sanov large deviation principle for empirical measure. It consists mainly of an adaptation of [GZ93, Theorem 3.1] (with $F = 0$) to our context, while taking care of the additional terms induced by changing the density of points.

Proposition 2.4. *Let P be in $\mathcal{P}_{s,m}(\mathcal{X})$ for some $m > 0$. Let $\{N_R\}_{R>0}$ be such that*

$$N_R = mR^2(1 + o(1)).$$

Let us denote by i_R the “empirical field” map with average in C_R i.e. (similarly to (1.4), (1.16))

$$i_R(\mathcal{C}) := \frac{1}{R^2} \int_{C_R} \delta_{\theta_z \cdot (\mathcal{C} \cap C_R)} dz.$$

We have

$$(2.6) \quad \lim_{\varepsilon \rightarrow 0} \lim_{R \rightarrow \infty} \frac{1}{R^2} \log \int_{(C_R)^{N_R}} \mathbf{1}_{i_R(x_1, \dots, x_{N_R}) \in B(P, \varepsilon)} dx_1 \dots dx_{N_R} = -(\text{ent}[P|\mathbf{\Pi}^m] - m + m \log m).$$

We postpone the proof to Section 7.1. Henceforth we let, for any $z \in \Sigma$

$$(2.7) \quad c_z := m_{\text{eq}}(z) - m_{\text{eq}}(z) \log m_{\text{eq}}(z).$$

3 Preliminary considerations on the energy

3.1 Monotonicity estimates

Almost monotonicity in η of the local energy. The next lemma expresses the fact that the limit $\eta \rightarrow 0$ in (2.4) is almost monotonous.

Lemma 3.1. *Let (\mathcal{C}, μ) be a neutral electric system with N points and E^{loc} be the associated local electric field. We have, for any $0 < \eta < \eta_1 < 1$,*

$$(3.1) \quad \left(\int_{\mathbb{R}^2} |E_\eta^{\text{loc}}|^2 + N \log \eta \right) - \left(\int_{\mathbb{R}^2} |E_{\eta_1}^{\text{loc}}|^2 + N \log \eta_1 \right) \succeq -N \|\mu\|_\infty \eta_1.$$

Proof. This is [PS15, Lemma 2.3]. □

Let us note that, integrating by parts, we may re-write $\int_{\mathbb{R}^2} |E_\eta^{\text{loc}}|^2$ as $\int_{\mathbb{R}^2} -H_\eta^{\text{loc}} \Delta H_\eta^{\text{loc}}$ and (3.1) is really a monotonicity estimate for $(\int_{\mathbb{R}^2} -H_\eta^{\text{loc}} \Delta H_\eta^{\text{loc}} + N \log \eta)$ as η varies.

A localized monotonicity estimate.

Lemma 3.2. *Let (\mathcal{C}, μ) be an electric system and E^{loc} be the associated local electric field. Let $R_2 > 10$ and let N^{in} be the number of points of \mathcal{C} in C_{R_2} . We let also N^{bou} be the number of points of \mathcal{C} in $C_{R_2} \setminus C_{R_2-5}$. We have for any $0 < \eta_1 < \eta_0 < 1$,*

$$(3.2) \quad \left(\int_{C_{R_2}} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) - \left(\int_{C_{R_2}} |E_{\eta_0}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_0 \right) \succeq -N^{\text{in}} \|\mu\|_{\infty} \eta_0 \\ + N^{\text{bou}} \log \eta_1 + (\log \eta_1 - 1) \int_{C_{R_2} \setminus C_{R_2-5}} \left(|E_{\eta_0}^{\text{loc}}|^2 + |E_{\eta_1}^{\text{loc}}|^2 \right).$$

Proof. It follows from the proof of [PS15, Lemma 2.4], see e.g. [PS15, Equation 2.29] and the one immediatly after. \square

Almost monotonicity with no points near the boundary.

Lemma 3.3. *Let (\mathcal{C}, μ) be an electric system in \mathbb{R}^2 and E^{loc} be the associated local electric field. Let $0 < R_2$ and let $0 < \eta_1 < 1$ be such that the smeared out charges at scale η_1 do not intersect ∂C_{R_2} i.e. $\bigcup_{p \in \mathcal{C}} B(p, \eta_1) \cap \partial C_{R_2} = \emptyset$. Let us denote by N^{in} the number of points in C_{R_2} . Then we have for any $\eta \leq \eta_1$*

$$(3.3) \quad \left(\int_{C_{R_2}} |E_{\eta}^{\text{loc}}|^2 + N^{\text{in}} \log \eta \right) - \left(\int_{C_{R_2}} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) \succeq -N^{\text{in}} \eta_1 \|\mu\|_{\infty}$$

Proof. We postpone the proof to Section 7.2. \square

3.2 Discrepancy estimates

Lemma 3.4. *Let $N \geq 1$, and let \mathcal{C} be a finite point configuration in \mathbb{R}^2 . Let E be a gradient electric field in $\text{Elec}(\mathcal{C}, \mu'_{\text{eq}})$. For any $R > 0$, let \mathcal{D}_R be the discrepancy $\mathcal{D}_R := \int_{C_R} (d\mathcal{C} - d\mu'_{\text{eq}})$ in C_R .*

For any $\eta \in (0, 1)$ we have

$$(3.4) \quad \mathcal{D}_R^2 \min \left(1, \frac{\mathcal{D}_R}{R^2} \right) \preceq \int_{C_{2R}} |E_{\eta}|^2.$$

Proof. This follows from [RS15, Lemma 3.8]. \square

As a corollary, we see that if a good control holds at scale δ , then the discrepancies are controlled at smaller scales.

Lemma 3.5. *Let $0 < \delta \leq \frac{1}{2}$ and let us assume that a good control holds at scale δ . Then for any $R \in (\frac{1}{2}N^{\delta_1}, 2N^{\delta_1})$ with $\frac{\delta}{2} < \delta_1 < \delta$ we have*

$$(3.5) \quad |\mathcal{D}_R| \ll_{\delta_1} N^{\frac{4}{3}\delta}.$$

Proof. Let us apply Lemma 3.4 with $\eta = 1/2$, taking E to be the local electric field E^{loc} . It yields

$$\mathcal{D}_R^2 \min \left(1, \frac{\mathcal{D}_R}{R^2} \right) \preceq \int_{C_{2R}} |E_{1/2}^{\text{loc}}|^2.$$

Using the good control on the energy at scale δ (in the sense of Definition 1.2) we have

$$\int_{C_{2R}} |E_{1/2}^{\text{loc}}|^2 + \mathcal{N}_{2R} \log(1/2) \preceq N^{2\delta},$$

where \mathcal{N}_{2R} (the number of points in C_{2R}) is itself $\preceq_{\delta_1} N^{2\delta}$. We thus obtain that

$$\mathcal{D}_R^2 \min\left(1, \frac{\mathcal{D}_R}{R^2}\right) \preceq_{\delta_1} N^{2\delta}.$$

Then, elementary considerations imply that if $\frac{\delta}{2} < \delta_1 < \delta$, we have $\mathcal{D}_R^3 \preceq_{\delta_1} N^{2(\delta_1+\delta)}$ which yields (3.5). \square

Application: number of points in a square.

Lemma 3.6. *Let $0 < \delta \leq \frac{1}{2}$ and let us assume that a good control holds at scale δ . Then we have, for any $R \in (\frac{1}{2}N^{\delta_1}, 2N^{\delta_1})$ with $\frac{2}{3}\delta < \delta_1 < \delta$, and for any $z_0 \in \dot{\Sigma}$, letting $\mathcal{N}_R^{z_0} := \int_{C(z'_0, R)} d\mathcal{C}$*

$$(3.6) \quad |\mathcal{N}_R^{z_0} - m_{\text{eq}}(z_0)R^2| \ll_{\delta_1} N^{2\delta_1}.$$

Proof. We have by definition

$$\int_{C(z'_0, R)} d\mu'_{\text{eq}} = \int_{C(0, R)} m_{\text{eq}}(z_0 + tN^{-1/d}) dt.$$

Using the Hölder assumption (2.1) we get

$$\left| \int_{C(z'_0, R)} d\mu'_{\text{eq}} - R^2 m_{\text{eq}}(z_0) \right| \preceq N^{2\delta_1 + \kappa(\delta_1 - \frac{1}{2})}.$$

Since $\kappa > 0$ and $\delta_1 < 1/2$ we get

$$\left| \int_{C(z'_0, R)} d\mu'_{\text{eq}} - R^2 m_{\text{eq}}(z_0) \right| \ll N^{2\delta_1},$$

while Lemma 3.5 yields

$$\left| \mathcal{N}_R^{z_0} - \int_{C(z'_0, R)} d\mu'_{\text{eq}} \right| \ll_{\delta_1} N^{\frac{4}{3}\delta}.$$

Combining these two inequalities we see that if $\delta_1 > \frac{2}{3}\delta$ then (3.6) holds. \square

3.3 Minimality of local energy against decaying fields

Lemma 3.7. *Let $K \subset \mathbb{R}^2$ be a compact set with piecewise C^1 boundary. Let (\mathcal{C}, μ) be a neutral electric system in K . Let E^{loc} be the local electric field associated to (\mathcal{C}, μ) and let $E \in \text{Elec}^0(\mathcal{C}, \mu)$. For any $0 < \eta < 1$ we have*

$$(3.7) \quad \int_{\mathbb{R}^2} |E_\eta^{\text{loc}}|^2 \leq \int_{\mathbb{R}^2} |E_\eta|^2.$$

The proof is very similar to that of [LS15, Lemma 3.12] and we postpone it to Section 7.

3.4 The screening lemma

The following result is strongly inspired from the “screening lemmas” of [SS15b, Prop. 6.4], [PS15, Prop 6.1] and [LS15, Prop 5.2]. Our setting is slightly simpler because we have ruled out the possibility of having point charges close to the boundary ∂C_{R_2} . A difference is that here we screen the electric field “from the inside” (a similar procedure is used in [RNS15]) whereas the aforementioned screening results were constructing a field E such that $E \equiv 0$ outside a certain hypercube. Another difference is that in the present lemma we really need to deal with a variable background μ .

Let us briefly explain the purpose of Lemma 3.8. Let μ be a fixed background, defined on a large square C_{R_2} . We start with a point configuration and an electric field associated to it (and to the negative background given by μ), that we can think of as living in $C_{R_2} \setminus C_{R_1}$, and we want to fill C_{R_1} with some other point configuration and associated electric field. To connect both electric fields, we need to check that their normal components coincide along the boundary of C_{R_1} (otherwise divergence is created). A way to ensure that is to have them both equal to 0 on ∂C_{R_1} . We thus need to change the initial point configuration a little bit so that we can find an associated electric field whose normal component vanishes on ∂C_{R_1} . The next lemma shows that under some assumptions, and especially if the energy control (3.9) holds, we may slightly perturb the point configuration and pass from a given electric field E to a new one E^{tran} which is “screened” in that (3.12) holds, and whose energy is close to that of E (as expressed in (3.13)). In Section 4 we will ensure that, by construction, the assumptions of Lemma 3.8 are satisfied.

Lemma 3.8. *There exists $C > 0$ universal such that the following holds.*

Let $z \in \mathbb{R}^2$ and for any N let (\mathcal{C}, μ) be an electric system in \mathbb{R}^2 , let $0 < \delta_3 < \delta_2 < \delta_1 < 1/2$ and let R_1, R_2 positive be such that, letting $\eta_1 := N^{-10}$, we have

1. $R_2 \in [N^{\delta_1} + N^{\delta_2}, N^{\delta_1} + 2N^{\delta_2}]$ and $R_1 \in [R_2 - 3N^{\delta_3}, R_2 - 2N^{\delta_3}]$.
2. *The smeared out charges at scale η_1 do not intersect $\partial C(z, R_2)$, i.e.*

$$\bigcup_{p \in \mathcal{C}}^N D(p, \eta_1) \cap \partial C(z, R_2) = \emptyset.$$

3. N^{gen} is an integer, where $N^{\text{gen}} := \int_{C(z, R_1)} d\mu$.
4. Letting $N^{\text{mid}} := \int_{C(z, R_2) \setminus C(z, R_1)} d\mathcal{C}$, it holds $N^{\text{mid}} \ll N^{2\delta_1}$.

Let us assume that μ satisfies $0 < \underline{m} \leq \mu \leq \bar{m}$ on $C_{R_2} \setminus C_{R_1}$ and that furthermore there exists $C_\mu > 0$ such that

$$(3.8) \quad \forall (x, y) \in (C_{R_2} \setminus C_{R_1})^2, |\mu(x) - \mu(y)| \leq C_\mu |x - y|^\kappa,$$

where κ is as in (2.1). Let E be in $\text{Elec}(\mathcal{C}, \mu, \mathbb{R}^2 \setminus C_{R_1})$ and let $M := \int_{\partial C_{R_2}} |E_\eta|^2$.

If the following inequality is satisfied

$$(3.9) \quad M \leq C \min(\underline{m}^2, 1) N^{3\delta_3},$$

then there exists a measurable family A_N^{tran} of point configurations such that for any $\mathcal{C}^{\text{tran}} \in A_N^{\text{tran}}$

1. *The configuration $\mathcal{C}^{\text{tran}}$ is supported in $C_{R_2} \setminus C_{R_1}$.*
2. *The configuration $\mathcal{C}^{\text{tran}}$ has N^{tran} points where (with \vec{n} the unit normal vector)*

$$(3.10) \quad N^{\text{tran}} := \int_{\partial C_{R_2}} E_{\eta_1} \cdot \vec{n} - \int_{C_{R_2} \setminus C_{R_1}} d\mu.$$

3. The points of $\mathcal{C}^{\text{tran}}$ are well-separated from each other and from the boundaries

$$(3.11) \quad \min_{p_1 \neq p_2 \in \mathcal{C}^{\text{tran}}} |p_1 - p_2| \succeq \bar{m}^{-1/2}, \quad \min_{p \in \mathcal{C}^{\text{tran}}} \text{dist}(p, \partial C_{R_2} \cup \partial C_{R_1}) \succeq \bar{m}^{-1/2}.$$

4. There exists an electric field $E^{\text{tran}} \in \text{Elec}(\mathcal{C}^{\text{tran}}, \mu, C_{R_2} \setminus C_{R_1})$ such that

(a) We have

$$(3.12) \quad E_{\eta_1}^{\text{tran}} \cdot \vec{n} = \begin{cases} E_{\eta_1} \cdot \vec{n} & \text{on } \partial C_{R_2} \\ 0 & \text{on } \partial C_{R_1} \end{cases}.$$

(b) The energy of E^{tran} is bounded by

$$(3.13) \quad \int_{C_{R_2} \setminus C_{R_1}} |E_{\eta_1}^{\text{tran}}|^2 \preceq N^{\delta_3} M + N^{\delta_1 + 3\delta_3} C_\mu N^{\kappa\delta_3} + N^{\delta_1 + \delta_3} \log N.$$

Moreover the volume of A_N^{tran} is bounded below as follows

$$(3.14) \quad \log \mathbf{Leb}^{\otimes N^{\text{tran}}}(A_N^{\text{tran}}) \succeq -(\bar{m}N^{\delta_1 + \delta_3} + N^{\delta_1} + M) \log \bar{m}.$$

Proof. We set $l = N^{\delta_3}$.

Step 1. Subdividing the domain. We claim that we may split $C_{R_2} \setminus C_{R_1}$ into a finite family of rectangles $\{H_i\}_{i \in I}$ with sidelengths in $[l/2, 2l]$ such that letting $\bar{m}_i := \frac{1}{|H_i|} \int_{H_i} d\mu$ and

$$(3.15) \quad m_i := \bar{m}_i + \frac{1}{|H_i|} \frac{1}{2\pi} \int_{\partial C_{R_2} \cap \partial H_i} E_\eta \cdot \vec{n},$$

we have for any $i \in I$

$$(3.16) \quad |m_i - \bar{m}_i| < \frac{1}{2} \underline{m}, \quad m_i |H_i| \in \mathbb{N}.$$

The fact that we may split $C_{R_2} \setminus C_{R_1}$ into a finite family of rectangles $\{H_i\}_{i \in I}$ with sidelengths in $[l/2, 2l]$ is elementary. Using Cauchy-Schwarz's inequality we see that

$$\left| \frac{1}{|H_i|} \int_{\partial C_{R_2} \cap \partial H_i} E_\eta \cdot \vec{n} \right| \preceq \frac{1}{l^{3/2}} \left(\int_{\partial C_{R_2}} |E_\eta|^2 \right)^{1/2} = M^{1/2} l^{-3/2}$$

hence assuming (3.9) (with C large enough) we have $|m_i - \bar{m}_i| < \frac{1}{2} \underline{m}$ for any tiling of $C_{R_2} \setminus C_{R_1}$ by rectangles of sidelengths $\in [l/2, 2l]$. It remains to show that we may obtain a tiling such that $m_i |H_i| \in \mathbb{N}$. We have by definition

$$m_i |H_i| = \int_{H_i} d\mu + \frac{1}{2\pi} \int_{\partial C_{R_2} \cap \partial H_i} E_\eta \cdot \vec{n}.$$

Increasing the sidelengths of H_i by t (with $t \leq l/10$) increases $\int_{H_i} d\mu$ by a quantity $\succeq \underline{m} t$ whereas it changes $\int_{\partial C_{R_2} \cap \partial H_i} E_\eta \cdot \vec{n}$ by a quantity $\preceq \sqrt{tM}$. We thus see that if (3.9) holds, up to modifying the boundaries of H_i a little bit (e.g. changing the sidelengths by a quantity less than $l/10$) we can ensure that each $m_i |H_i| \in \mathbb{N}$.

We may then subdivide further each rectangle H_i into a finite family of rectangles $\{R_\alpha\}_{\alpha \in I_i}$ which all have an area m_i^{-1} and sidelengths bounded above and below by universal constants

times $m_i^{-1/2}$ (for a proof of this fact, see [PS15, Lemma 6.3.]). Let us emphasize that since μ is bounded above we have $m_i^{-1} \preceq 1$.

Step 2. *Defining the transition field and configuration.* For $i \in I$ we let $E^{(1,i)} := \nabla h^{(1,i)}$ where $h^{(1,i)}$ is the mean-zero solution to

$$-\Delta h^{(1,i)} = 2\pi(m_i - \mu) \text{ in } H_i,$$

such that $\nabla h^{(1,i)} \cdot \vec{n} = -E_{\eta_1} \cdot \vec{n}$ on each side of H_i which is contained in ∂C_{R_2} and $\nabla h^{(1,i)} \cdot \vec{n} = 0$ on the other sides.

We also let, for any $\alpha \in I_i$, $h^{(2,\alpha)}$ be the mean-zero solution to

$$-\Delta h^{(2,\alpha)} = 2\pi(\delta_{p_\alpha} - m_i) \text{ in } R_\alpha, \nabla h^{(2,\alpha)} \cdot \vec{n} = 0 \text{ on } \partial R_\alpha,$$

where p_α is the center of R_α . We define $E^{(2,i)}$ as $E^{(2,i)} := \sum_{\alpha \in I_i} \nabla h^{(2,\alpha)}$. Finally we define the *transition field* E^{tran} as $E^{\text{tran}} := \sum_{i \in I} E^{(1,i)} + E^{(2,i)}$, and the *transition configuration* $\mathcal{C}^{\text{tran}}$ as $\mathcal{C}^{\text{tran}} := \sum_{i \in I} \sum_{\alpha \in I_i} \delta_{p_\alpha}$. It is easy to see that

$$(3.17) \quad -\operatorname{div}(E^{\text{tran}}) = 2\pi(\mathcal{C}^{\text{tran}} - \mu) \text{ in } C_{R_2} \setminus C_{R_1}, \quad E_{\eta_1}^{\text{tran}} \cdot \vec{n} = \begin{cases} E_{\eta_1} \cdot \vec{n} & \text{on } \partial C_{R_2} \\ 0 & \text{on } \partial C_{R_1} \end{cases}.$$

In particular (3.10) and (3.12) hold.

Step 3. *Controlling the energy.* For any $i \in I$ the energy of $E^{(1,i)}$ can be bounded using Lemma 7.1 as follows

$$\int_{H_i} |E^{(1,i)}|^2 \preceq l \int_{\partial H_i \cap \partial C_{R_2}} |E_\eta|^2 + l^4 \|\mu - \bar{m}_i\|_{L^\infty(H_i)}^2,$$

and using the Hölder assumption (3.8) on μ we have $\|\mu - \bar{m}_i\|_{L^\infty(H_i)}^2 \preceq C_\mu l^\kappa$ hence

$$\int_{H_i} |E^{(1,i)}|^2 \preceq l \int_{\partial H_i \cap \partial C_{R_2}} |E_\eta|^2 + C_\mu l^{4+\kappa}.$$

For any $\alpha \in I_i$ we also have, again by standard estimates

$$\int_{R_\alpha} |\nabla h_{\eta_1}^{(2,\alpha)}|^2 \preceq -\log \eta_1 \preceq \log N \text{ by choice of } \eta_1.$$

The number of rectangles R_α for $\alpha \in I_i$, $i \in I$ is bounded by the volume of $C_{R_2} \setminus C_{R_1}$ hence is $\preceq N^{\delta_1 + \delta_3}$. We deduce that

$$\int_{C_{R_2} \setminus C_{R_1}} |E_{\eta_1}^{\text{tran}}|^2 \preceq l \int_{\partial C_{R_2}} |E_\eta|^2 + \#I C_\mu l^{4+\kappa} + N^{\delta_1 + \delta_3} \log N,$$

where $\#I$ denotes the cardinality of I . We may observe that $\#I \preceq N^{\delta_1 + \delta_3} l^{-2}$ and get (using that $l = N^{\delta_3}$)

$$\int_{C_{R_2} \setminus C_{R_1}} |E_{\eta_1}^{\text{tran}}|^2 \preceq N^{\delta_3} M + N^{\delta_1 + 3\delta_3} C_\mu N^{\kappa\delta_3} + N^{\delta_1 + \delta_3} \log N.$$

which yields (3.13).

Step 4. *Constructing a family.* As was observed in [PS15, Remark 6.7], [LS15, Proposition 5.2], we may actually construct a whole family of configurations $\mathcal{C}^{\text{tran}}$ and associated electric fields E^{tran} such that (3.17) and (3.13) hold. Indeed, since $\mu \leq \bar{m}$ the sidelengths of R_α are

$\asymp \bar{m}^{-1/2}$ hence we may move each of the points p_α (for $\alpha \in I_i, i \in I$) arbitrarily within a disk of radius $\frac{1}{10}\bar{m}^{-1/2}$ and proceed as above (so that (3.11) is conserved). This creates a volume of configurations of order $\bar{m}^{-N^{\text{tran}}}$. To get (3.14) it suffices to observe that

$$N^{\text{tran}} \leq \bar{m}N^{\delta_1+\delta_3} + N^{\delta_1} + M,$$

which follows from (3.10) by applying the Cauchy-Schwarz inequality and using the fact that $R_2^2 - R_1^2 \leq N^{\delta_1+\delta_3}$. \square

4 A large deviation upper bound

Given a point z_0 and a scale δ_1 we localize the energy $w_N(\vec{X}_N)$ by splitting it between the “interior part” (which, roughly speaking, corresponds to the L^2 -norm of E^{loc} on $C(z'_0, N^{\delta_1})$) and the “exterior part”, then we replace both quantities by smaller ones which are independent (this corresponds to separating variables). Roughly speaking, it allows us to localize the problem on $C(z'_0, N^{\delta_1})$ by deriving a “local energy” (which corresponds to the interior part) and a “local partition function” (the exponential sum of contributions of the exterior part). This lower bound on the energy will be complemented by a matching upper bound in Section 5.

4.1 Definition of good interior and exterior boundaries and energies

The decomposition between interior and exterior part will be done at the boundary of some “good” square, not much larger than $C(z'_0, N^{\delta_1})$. We give the definition of good exterior and interior boundaries, with an abuse of notation which is discussed in the next subsection.

Definition 4.1 (Exterior boundary). *Let $1/2 > \delta > \delta_1 > \delta_2 > 0$ and $\eta_0 > 0$ be fixed, let $N \geq 1$ and let $\eta_1 := N^{-10}$. Let $z_0 \in \tilde{\Sigma}$, let $R_2 > 0$ and let \vec{X}^{out} be a point configuration in $C(z'_0, R_2)^c$.*

Exterior fields. *We say that E is in $\text{Elec}^{\text{out}}(\vec{X}^{\text{out}})$ if E is in $\text{Elec}^0(\mathcal{C}, \mu'_{\text{eq}}, \mathbb{R}^2)$ for some point configuration \mathcal{C} with N points such that $\mathcal{C} = \vec{X}^{\text{out}}$ on $C(z, R_2)^c$.*

Good exterior boundary. *Let $E \in \text{Elec}^{\text{out}}(\vec{X}^{\text{out}})$. We say that $\partial C(z'_0, R_2)$ is a good exterior boundary for E if the following conditions are satisfied:*

1. *We have*

$$(4.1) \quad R_2 \in [N^{\delta_1} + N^{\delta_2}, N^{\delta_1} + 2N^{\delta_2}].$$

2. *The smeared out charges at scale η_1 do not intersect $\partial C(z'_0, R_2)$*

$$(4.2) \quad \bigcup_{p \in \mathcal{C}} D(p, \eta_1) \cap \partial C(z'_0, R_2) = \emptyset.$$

3. *The energy near $\partial C(z'_0, R_2)$ is controlled as follows*

$$(4.3) \quad \int_{C(z'_0, R_2) \setminus C(z'_0, R_2-5)} |E_{\eta_0}|^2 \leq N^{2\delta-\delta_2} |\log \eta_0|,$$

$$(4.4) \quad \int_{C(z'_0, R_2) \setminus C(z'_0, R_2-5)} |E_{\eta_1}|^2 \leq N^{2\delta-\delta_2} \log N,$$

$$(4.5) \quad \int_{\partial C(z'_0, R_2)} |E_{\eta_1}|^2 \leq N^{2\delta-\delta_2} (\log N)^2.$$

4. We have, letting $N^{\text{bou}} := \int_{C(z'_0, R_2) \setminus C(z'_0, R_2 - 5)} d\mathcal{C}$

$$(4.6) \quad N^{\text{bou}} \ll N^{2\delta_1}.$$

5. We have, letting $N^{\text{in}} := \int_{C(z'_0, R_2)} d\mathcal{C}$

$$(4.7) \quad N^{\text{in}} \leq N^{2\delta_1}.$$

6. We have, letting

$$(4.8) \quad \mathcal{N}_{\delta_1}^{z_0} := \int_{C(z'_0, N^{\delta_1})} d\mathcal{C}$$

that the following holds

$$(4.9) \quad N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0} \ll N^{2\delta_1} \text{ and } |\mathcal{N}_{\delta_1}^{z_0} - m_{\text{eq}}(z_0)N^{2\delta_1}| \ll N^{2\delta_1}.$$

This very last inequality does not depend on R_2 , but it is convenient to include it in the definition of the exterior boundary.

Best exterior energy. Let N^{out} denote the number of points of \vec{X}^{out} in $C(z'_0, R_2)^c$. We define $F^{\text{out}}(\vec{X}^{\text{out}})$ as

$$(4.10) \quad F^{\text{out}}(\vec{X}^{\text{out}}) := \frac{1}{2\pi} \min_E \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)^c} |E_\eta|^2 + N^{\text{out}} \log \eta \right)$$

where the min is taken over the set of electric fields E satisfying $E \in \text{Elec}^{\text{out}}(\vec{X}^{\text{out}})$ and such that $\partial C(z'_0, R_2)$ is a good exterior boundary for E (if there is no such field we set $F^{\text{out}}(\vec{X}^{\text{out}}) = +\infty$). The minimum is achieved because on the one hand $E \mapsto |E_\eta|^2$ is coercive for the weak L^p_{loc} topology, and on the other hand the L^2 norm is coercive and lower semi-continuous for the weak L^2 topology.

Definition 4.2 (Interior boundary). Let $1/2 > \delta_1 > \delta_2 > \delta_3 > 0$ and $\eta_0 > 0$ be fixed. Let $N \geq 1$ and let $z_0 \in \hat{\Sigma}$, let $R_1 > 0$, let $R_2 > 0$ such that (4.1) holds and let \vec{X}^{in} be a point configuration in $C(z'_0, R_2)$.

Interior fields. Let $E \in \text{Elec}(\mathbb{R}^2)$. We say that E is in $\text{Elec}^{\text{in}}(\vec{X}^{\text{in}})$ if E is in $\text{Elec}^0(C, \mu'_{\text{eq}}, \mathbb{R}^2)$ for some $C \in \mathcal{X}(\mathbb{R}^2)$ such that $C = \vec{X}^{\text{in}}$ on $C(z'_0, R_2)$.

Good interior boundary. We say that $\partial C(z'_0, R_1)$ is a good interior boundary for \vec{X}^{in} if

1. We have

$$(4.11) \quad R_1 \in [R_2 - 3N^{\delta_3}, R_2 - 2N^{\delta_3}].$$

2. N^{gen} is an integer, where $N^{\text{gen}} := \int_{C(z'_0, R_1)} d\mu'_{\text{eq}}$.

3. Letting $N^{\text{mid}} := \int_{C(z'_0, R_2) \setminus C(z'_0, R_1)} d\mathcal{C}$, it holds

$$(4.12) \quad N^{\text{mid}} \ll N^{2\delta_1}.$$

Best interior energy. Let N^{in} denote the number of points of \vec{X}^{in} in $C(z'_0, R_2)$. For any $0 < \eta_0 < 1$ we define $F_{\eta_0}^{\text{in}}$ as

$$(4.13) \quad F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}}) := \frac{1}{2\pi} \min_E \left(\int_{C(z'_0, R_2)} |E_{\eta_0}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_0 \right),$$

where the minimum is taken over the set of electric fields E such that $E \in \text{Elec}^{\text{in}}(\vec{X}^{\text{in}})$ (if $\text{Elec}^{\text{in}}(\vec{X}^{\text{in}})$ is empty we set $F^{\text{in}}(\vec{X}^{\text{in}}) = +\infty$).

4.2 Finding good boundaries

The conditions (4.3), (4.4), (4.5), (4.6), (4.7), (4.9), (4.12) are *asymptotic* as $N \rightarrow \infty$, in particular they do not make sense for a finite N (nonetheless, (4.1), (4.2) and (4.11) do). Strictly speaking one thus has to consider sequences $R_2 = R_2(N)$ and $R_1 = R_1(N)$.

Lemma 4.3. *Let $1/2 \geq \delta > \delta_1 > \delta_2 > \delta_3 > 0$ and $\eta_0 > 0$ be fixed, with $\delta_1 > 2\delta/3$. Let $z_0 \in \mathring{\Sigma}$ and $\eta_0 > 0$ be fixed. Assume that good control at scale δ holds and let $\delta_1 \in (\delta/2, \delta)$. With δ_1 -overwhelming probability, there exists R_1, R_2 such that, letting $\vec{X}^{\text{in}} = \vec{X}'_N \cap C(z'_0, R_2)$ and $\vec{X}^{\text{out}} = \vec{X}'_N \cap C(z'_0, R_2)^c$, we have*

1. $\partial C(z'_0, R_2)$ is a good exterior boundary for $\vec{X}^{\text{out}}, E^{\text{loc}}$.
2. $\partial C(z'_0, R_1)$ is a good interior boundary for \vec{X}^{in} .

Proof of Lemma 4.3. First we look for a good exterior boundary $\partial C(z'_0, R_2)$. The good control at scale δ implies that

$$(4.14) \quad \int_{C(z'_0, N^\delta)} |E_{\eta_1}^{\text{loc}}|^2 \preceq_{\delta_1} N^{2\delta} (1 + |\log \eta_1|) \preceq N^{2\delta} \log N,$$

$$(4.15) \quad \int_{C(z'_0, N^\delta)} |E_{\eta_0}^{\text{loc}}|^2 \preceq_{\delta_1} N^{2\delta} (1 + |\log \eta_0|) \preceq N^{2\delta} |\log \eta_0|.$$

In view of (4.14), (4.15), by the pigeonhole principle, we may find (with δ_1 -overwhelming probability) an interval $[R - 10, R + 10]$ included in $[N^{\delta_1} + N^{\delta_2}, N^{\delta_1} + 2N^{\delta_2}]$ such that

$$(4.16) \quad \int_{C(z'_0, R+10) \setminus C(z'_0, R-10)} (|E_{\eta_1}^{\text{loc}}|^2 + |E_{\eta_0}^{\text{loc}}|^2) \preceq N^{2\delta - \delta_2} (\log N + |\log \eta_0|).$$

We may find N^2 disjoint strips of width $2N^{-2}(\log N)^{-1}$ in $[R - 8, R + 8]$. In view of (4.16), there are at least $N^2/2$ such strips on which the integral of $|E_{\eta_1}^{\text{loc}}|^2 + |E_{\eta_0}^{\text{loc}}|^2$ is $\preceq N^{2\delta - \delta_2 - 2} (\log N + |\log \eta_0|)$. On the other hand there are at most N point charges, thus since $\eta_1 = N^{-10}$ by Dirichlet's drawer principle we may moreover assume that no smeared out charge (at scale η_1) intersects the strips (more precisely, there is at least one such strip, and in fact many more). Finally a mean value argument on one of these strips shows that we may find R_2 such that (4.5) and (4.2) holds. By (4.16) we also have (4.3) and (4.4).

Next, we look for a good interior boundary $\partial C(z'_0, R_1)$. Since z_0 is in the interior of Σ , the density m'_{eq} is bounded above and below by positive constants on $C(z'_0, N^\delta)$ (for N large enough) and thus the derivative of $R \mapsto \int_{C(z'_0, R)} dm'_{\text{eq}}$ is bounded above and below by (positive) constants times N^{δ_1} on $C(z'_0, 2N^{\delta_1})$. Hence we may find $R_1 \in [R_2 - 3N^{\delta_3}, R_2 - 2N^{\delta_3}]$ such that $\int_{C(z'_0, R_1)} d\mu'_{\text{eq}}$ is an integer, hence the first two points of the definition of a good interior boundary are satisfied.

We have (4.7) with δ_1 -overwhelming probability according to the good control at scale δ . Since $\delta_1 > 2\delta/3$, the discrepancy estimates of Lemma 3.6 imply that, up to an error $\ll_{\delta_1} N^{2\delta_1}$ we have

$$\begin{aligned} N^{\text{bou}} &= m_{\text{eq}}(z_0)(R_2^2 - (R_2 - 10)^2) \preceq_{\delta_1} N^{\delta_1} \ll N^{2\delta_1} \\ N^{\text{tran}} &= m_{\text{eq}}(z_0)(R_2^2 - R_1^2) \preceq_{\delta_1} N^{\delta_1 + \delta_2} \ll N^{2\delta_1}, \end{aligned}$$

which proves (4.6) and (4.12). We obtain (4.9) with similar arguments. \square

4.3 A first LDP upper bound

Let us recall that $F_{\eta_0}^{\text{in}}, F^{\text{out}}$ are defined in (4.13), (4.10).

Proposition 4.4. *Let $1/2 \geq \delta > \delta_1 > \delta_2 > \delta_3 > 0$ be fixed with $\delta_1 > 2\delta/3$ and $2\delta - \delta_2 < \delta_1$. Let $z_0 \in \dot{\Sigma}$ and $0 < \eta_0 < 1$ be fixed. Assume that a good control holds at scale δ .*

For any $P \in \mathcal{P}_s(\mathcal{X})$ and any $\varepsilon > 0$ we have

$$(4.17) \quad \log \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(B(P,\varepsilon)) \leq -\log K_{N,\beta} \\ + \log \max_{R_1, R_2, N^{\text{out}}} \binom{N}{N^{\text{out}}} \left(\int_{i_{N,\delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P,\varepsilon)} e^{-\frac{1}{2}\beta F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}})} d\vec{X}^{\text{in}} \right) \left(\int e^{-\frac{1}{2}\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right) \\ + N^{2\delta_1} O(\eta_0).$$

Let us first give some clarifications about (4.17). The max on R_1, R_2, N^{out} is restricted to the set of $\{R_1, R_2\}$ such that (4.1) and (4.11) hold, with N^{out} between 1 and N . Once N^{out} is fixed, we let $d\vec{X}^{\text{in}} = dx_1 \dots dx_{N^{\text{in}}}$ and $d\vec{X}^{\text{out}} = dx_1 \dots dx_{N^{\text{out}}}$, with $N^{\text{out}} + N^{\text{in}} = N$. The binomial coefficient originates from the fact that the particles are initially labeled, but for commodity we prefer to think e.g. that the N^{in} first ones are in the ‘‘interior part’’ and that the N^{out} last ones are in the ‘‘exterior part’’. This induces a simple combinatorial factor.

Proof. Using the definition of $\mathfrak{P}_{N,\beta,\delta_1}^{z_0}$ and of the Gibbs measure \mathbb{P}_N^β we have

$$(4.18) \quad \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(B(P,\varepsilon)) = \frac{1}{K_{N,\beta}} \int_{(i_{N,\delta_1}^{z_0})^{-1}(B(P,\varepsilon))} e^{-\frac{1}{2}\beta(w_N(\nu'_N) + N\tilde{\zeta}(\nu_N))} d\vec{X}_N.$$

Step 1. *Finding good boundaries.* We apply Lemma 4.3. With δ_1 -overwhelming probability we obtain R_1, R_2 such that (4.1) and (4.11) hold, and such that, letting $\vec{X}^{\text{in}} = \vec{X}'_N \cap C(z'_0, R_2)$ and $\vec{X}^{\text{out}} = \vec{X}'_N \cap C(z'_0, R_2)^c$, we have

1. $\partial C(z'_0, R_2)$ is a good exterior boundary for $\vec{X}^{\text{out}}, E^{\text{loc}}$.
2. $\partial C(z'_0, R_1)$ is a good interior boundary for \vec{X}^{in} .

In order to prove (4.17) we may restrict ourselves, in the right-hand side of (4.18), to any event of δ_1 -overwhelming probability, and we will henceforth assume that good boundaries exist. **Step**

2. *Splitting the energy w_N .* For any \vec{X}'_N , let R_1, R_2 be as above. We have, using (2.4)

$$(4.19) \quad w_N(\nu'_N) = \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)} |E_\eta^{\text{loc}}|^2 + N^{\text{in}} \log \eta \right) + \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)^c} |E_\eta^{\text{loc}}|^2 + N^{\text{out}} \log \eta \right).$$

Since (4.2) holds we may apply Lemma 3.3 to $(\vec{X}'_N, \mu'_{\text{eq}})$ and R_2, η_1 as above. Since $\eta_1 = N^{-10}$, it yields

$$\lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)} |E_\eta^{\text{loc}}|^2 + N^{\text{in}} \log \eta \right) \geq \left(\int_{C(z'_0, R_2)} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) + o(1), \text{ as } N \rightarrow \infty.$$

We may then apply Lemma 3.2 to $(\nu'_N, \mu'_{\text{eq}})$ and R_2 , with η_0, η_1 as above. The number N^{in} of points in $C(z'_0, R_2)$ is controlled by (4.7), the number N^{bou} of points near the boundary is controlled by (4.6) and the energy near the boundary is controlled by (4.3) and (4.4). We obtain

$$\left(\int_{C(z'_0, R_2)} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) - \left(\int_{C(z'_0, R_2)} |E_{\eta_0}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_0 \right) \\ \geq -N^{2\delta_1} \eta_0 - N^{\text{bou}} \log N - N^{2\delta - \delta_2} (\log N)^2,$$

which may be simplified (assuming that $2\delta - \delta_2 < \delta_1$, which will be later ensured by the choice (5.4)) as

$$(4.20) \quad \left(\int_{C(z'_0, R_2)} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) - \left(\int_{C(z'_0, R_2)} |E_{\eta_0}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_0 \right) \succeq -N^{2\delta_1} \eta_0.$$

Using Definition 4.2 we thus get

$$(4.21) \quad \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)} |E_{\eta}^{\text{loc}}|^2 + N^{\text{in}} \log \eta \right) - F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}}) \succeq -\eta_0 N^{2\delta_1}.$$

On the other hand, we may write, using Definition 4.1

$$(4.22) \quad \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)^c} |E_{\eta}^{\text{loc}}|^2 + N^{\text{out}} \log \eta \right) \geq F^{\text{out}}(\vec{X}^{\text{out}}).$$

Combining (4.19), (4.21), (4.22) and inserting them into (4.18) yields (4.17). More precisely, we may split the domain of integration in (4.18) according to the value of R_1, R_2 and N^{out} , and then (up to a binomial coefficient as in (4.17)) we may assume that \vec{X}^{in} consists of the N^{in} first particles and \vec{X}^{out} of the N^{out} last ones, and the constraint $(i_{N, \delta_1}^{z_0})^{-1}(B(P, \varepsilon))$ on the domain of integration only concerns the points of \vec{X}^{in} . Inserting the splitting of the energy and the lower bounds on each term as described above leads then to (4.17). \square

4.4 Good control upper bound

Lemma 4.5. *Let $1/2 \geq \delta > \delta_1 > \delta_2 > \delta_3 > 0$ be fixed with $\delta_1 > 2\delta/3$. Let $z_0 \in \mathring{\Sigma}$ and $0 < \eta_0 < 1$ be fixed. Assume that a good control holds at scale δ .*

Let us denote by \mathcal{E}_M the event

$$(4.23) \quad \mathcal{E}_M := \left\{ \int_{C(z'_0, N^{\delta_1})} |E_{\eta_0}^{\text{loc}}|^2 + \mathcal{N}_{\delta_1}^{z_0} \log \eta_0 \geq 2\pi N^{2\delta_1} M \right\}.$$

We have

$$(4.24) \quad \log \mathbb{P}_N^\beta(\mathcal{E}_M) \leq -\log K_{N, \beta} - \frac{\beta}{2} M \\ + \log \max_{R_1, R_2, N^{\text{out}}} \binom{N}{N^{\text{out}}} \left(\int e^{-\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right) + N^{2\delta_1} O(\eta_0).$$

Proof. We follow the same steps as in the proof of Proposition 4.4, replacing the event $B(P, \varepsilon)$ by \mathcal{E}_M .

Let us write

$$\left(\int_{C(z'_0, R_2)} |E_{\eta_0}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_0 \right) \geq \left(\int_{C(z'_0, R_1)} |E_{\eta_0}^{\text{loc}}|^2 + \mathcal{N}_{\delta_1}^{z_0} \log \eta_0 \right) + (N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0}) \log \eta_0$$

Using (4.20) and the definition of \mathcal{E}_M we get, conditionally to \mathcal{E}_M

$$\left(\int_{C(z'_0, R_2)} |E_{\eta_1}^{\text{loc}}|^2 + N^{\text{in}} \log \eta_1 \right) \geq N^{2\delta_1} M + (N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0}) \log \eta_0 + O(N^{2\delta_1}).$$

Since we have $|N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0}|_{\delta_1} \ll N^{2\delta_1}$ we deduce that (conditionally to \mathcal{E}_M), using Lemma 3.3

$$(4.25) \quad \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)} |E_\eta^{\text{loc}}|^2 + N^{\text{in}} \log \eta \right) \geq N^{2\delta_1} M + O(N^{2\delta_1}).$$

Using (4.25) instead of (4.21) and arguing as in the proof of Proposition 4.4 gives (4.24). \square

4.5 Large deviation upper bound for the interior part

In (4.17) we have separated the Gibbs measure into its interior and exterior parts. In the next lemma, we give a large deviation upper bound for the interior part, namely

$$\int_{i_{N, \delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)} e^{-\frac{1}{2}\beta F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}})} d\vec{X}^{\text{in}}.$$

Up to technical details, this is a classical application of Varadhan's lemma (see e.g. [DZ10, Theorem 4.3.1]): the large deviations for the reference point process (without the exponential term) are known from Proposition 2.4, and on the other hand the lower semi-continuity of the energy near a random stationary point process P implies that $F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}}) \geq \mathbb{W}_{m_{\text{eq}}(z_0)}(P)$ on $\{i_{N, \delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)\}$, hence we obtain

Lemma 4.6. *Let $1/2 \geq \delta > \delta_1 > \delta_2 > \delta_3 > 0$ be fixed with $\delta_1 > 2\delta/3$, and let us assume that good control at scale δ holds. We have, for any R_1, R_2 satisfying (4.1), (4.11), and any N^{in}*

$$(4.26) \quad \limsup_{\eta_0 \rightarrow 0, \varepsilon \rightarrow 0, N \rightarrow \infty} N^{-2\delta_1} \log \int_{i_{N, \delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)} e^{-\frac{1}{2}\beta F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}})} d\vec{X}^{\text{in}} \leq -\mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0},$$

where c_{z_0} is the constant (depending only on $m_{\text{eq}}(z_0)$) defined in (2.7) and which appears in Proposition 2.4.

Let us emphasize that $\limsup_{\eta_0 \rightarrow 0, \varepsilon \rightarrow 0, N \rightarrow \infty}$ is a shorthand for taking the limit

$$\limsup_{\eta_0 \rightarrow 0} \limsup_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty},$$

and these limits do not commute in general.

Proof. Step 1. Lower semi-continuity of the energy. We claim that

$$(4.27) \quad \liminf_{\eta_0 \rightarrow 0, \varepsilon \rightarrow 0, N \rightarrow \infty} N^{-2\delta_1} \inf \left\{ F_{\eta_0}^{\text{in}}(\vec{X}^{\text{in}}), \vec{X}^{\text{in}} \in \mathcal{X}(C(z'_0, R_2)), i_{N, \delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon) \right\} \\ \geq \mathbb{W}_{m_{\text{eq}}(z_0)}(P).$$

To prove (4.27), let $E = E(N)$ be a minimizing sequence in (4.27), let $\mathcal{C} = \mathcal{C}(N)$ be the associated point configuration in $C(z'_0, N^{\delta_1})$ and let us define

$$P_N^{\text{elec}} := N^{-2\delta_1} \int_{C(z'_0, N^{\delta_1})} \delta_{\theta'_z \cdot E} dz', \quad P_N := N^{-2\delta_1} \int_{C(z'_0, N^{\delta_1})} \delta_{\theta'_z \cdot \mathcal{C}} dz'.$$

We have, for any $m \geq 1$,

$$\mathbf{E}_{P_N^{\text{elec}}} \left[\frac{1}{|\mathcal{C}_m|} \int_{\mathcal{C}_m} |E_{\eta_0}|^2 \right] \leq N^{-2\delta_1} \int_{C(z'_0, R_2)} |E_{\eta_0}|^2,$$

which proves that the push-forward of P_N^{elec} by $E \mapsto E_{\eta_0}$ is tight in $L^2(C_m, \mathbb{R}^2)$ for the weak topology, whereas the sequence $\{P_N^{\text{elec}}\}$ itself is tight in $L^p(C_m, \mathbb{R}^2)$ (indeed when η is fixed, the L^2 -norm of E_η controls the L^p -norm of E , this follows easily from Hölder's inequality, see [RS15, Lemma 3.9]). On the other hand, the sequence of random point processes $\{P_N\}$ is also tight because the expectation of the number of points in any square is bounded and, up to subsequence extraction, it converges to some $Q \in B(P, \varepsilon)$. Denoting by Q^{elec} a limit point of $\{P_N^{\text{elec}}\}$ it is not hard to see that Q^{elec} is compatible with Q , and by lower semi-continuity of the L^2 -norm with respect to weak convergence we have

$$\frac{1}{|C_m|} \int_{C_m} |E_{\eta_0}|^2 dQ^{\text{elec}} \leq \liminf_{N \rightarrow \infty} N^{-2\delta_1} \int_{C(z'_0, R_2)} |E_{\eta_0}|^2.$$

Up to applying a standard diagonal extraction procedure we may assume that it holds for any $m \geq 1$, hence

$$\mathbf{E}_{Q^{\text{elec}}} [|E_{\eta_0}|^2] \leq \liminf_{N \rightarrow \infty} N^{-2\delta_1} \int_{C(z'_0, R_2)} |E_{\eta_0}|^2.$$

Letting $\varepsilon \rightarrow 0$ and arguing as above concerning the tightness of Q^{elec} and Q (in L^p_{loc} and in \mathcal{X}) and for the lower semi-continuity of the norm with respect to weak convergence we obtain

$$\mathbf{E}_{P^{\text{elec}}} [|E_{\eta_0}|^2] \leq \liminf_{N \rightarrow \infty} N^{-2\delta_1} \int_{C(z'_0, N^{\delta_1})} |E_{\eta_0}|^2$$

where P^{elec} is some random electric field compatible with P . Letting $\eta_0 \rightarrow 0$, using (4.9) and the definition of $\mathbb{W}_{m_{\text{eq}}(z_0)}(P)$ we thus obtain

$$\mathbb{W}_{m_{\text{eq}}(z_0)}(P) \leq \lim_{\eta_0 \rightarrow 0, \varepsilon \rightarrow 0} \liminf_{N \rightarrow \infty} N^{-2\delta_1} \left(\int_{C(z'_0, N^{\delta_1})} |E_{\eta_0}|^2 + \mathcal{N}_{\delta_1}^{z_0} \log \eta_0 \right).$$

Step 2. Large deviations without interactions. We claim that, on the other hand

$$(4.28) \quad \limsup_{\varepsilon \rightarrow 0, N \rightarrow \infty} N^{-2\delta_1} \log \int_{i_{N, \delta_1}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)} \mathbf{1}_{C(z'_0, R_2)}(\vec{X}^{\text{in}}) d\vec{X}^{\text{in}} \leq -\text{ent}[P | \mathbf{\Pi}^{m_{\text{eq}}(z_0)}] + c_{z_0},$$

where c_{z_0} is as in (2.7).

The configuration \vec{X}^{in} has N^{in} points. Among them, $N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0}$ are not affected by the constraint $i_{N, \delta}^{z_0}(\vec{X}^{\text{in}}) \in B(P, \varepsilon)$ because they belong to $C(z'_0, N^{\delta_1})^c$, and are free to move in $C(z'_0, R_2) \setminus C(z'_0, N^{\delta_1})$, but we know from (4.9) that $N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0} \ll N^{2\delta_1}$, thus the volume contribution of these points is negligible because $N^{-2\delta_1} \log |C(z'_0, R_2) \setminus C(z'_0, N^{\delta_1})|^{N^{\text{in}} - \mathcal{N}_{\delta_1}^{z_0}} = o(1)$. On the other hand we know, using (4.9), that $\mathcal{N}_{\delta_1}^{z_0} \sim_{N \rightarrow \infty} m_{\text{eq}}(z_0) N^{2\delta_1}$ and then (4.28) follows from Proposition 2.4.

Step 3. Conclusion. Combining (4.27) and (4.28) yields (4.26). \square

4.6 A second LDP upper bound

Combining Proposition 4.4 and Lemma 4.6 and letting $\eta_0 \rightarrow 0$ we obtain

Proposition 4.7.

$$(4.29) \quad \lim_{\varepsilon \rightarrow 0} \limsup_{N \rightarrow \infty} N^{-2\delta_1} \log \mathfrak{P}_{N, \beta, \delta_1}^{z_0}(B(P, \varepsilon)) \leq -\mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0} \\ + \limsup_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N, z, \delta_1}^\beta$$

where c_{z_0} is as in (2.7) and we let K_{N,z,δ_1}^β be such that

$$(4.30) \quad \log K_{N,z,\delta_1}^\beta := -\log K_{N,\beta} + \log \int_{R_1, R_2} \sum_{N^{\text{out}}} \binom{N}{N^{\text{out}}} \int e^{-\frac{1}{2}\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}}.$$

5 Large deviation lower bound

In this section, we derive a converse estimate to (4.29), by showing that splitting the energy as in Proposition 4.4 is essentially sharp as far as probabilities of order $\exp(-N^{2\delta_1})$ are concerned.

5.1 Generating microstates

In the next lemma, we recall a tool which was introduced in [LS15]. Given a stationary point process P and a large square C_R , Lemma 5.1 can be thought of as a way of generating a family of point configurations in C_R whose empirical field is close to P , whose interaction energy is close to the renormalized energy of P , and such that the volume of the family is optimal in view of the specific relative entropy of P .

Lemma 5.1. *Let $z_0 \in \mathring{\Sigma}$ and $0 < \delta_2 < \delta_1 < 1/2$ be fixed. Let $P \in \mathcal{P}_{s, m_{\text{eq}}(z_0)}(\mathcal{X})$ such that $\mathbb{W}_{m_{\text{eq}}(z_0)}(P)$ and $\text{ent}[P|\mathbf{\Pi}^{m_{\text{eq}}(z_0)}]$ are finite.*

For any $N \geq 1$, let $R_1 > 0$ be such that $R_1 \in (N^{\delta_1}, N^{\delta_1} + 2N^{\delta_2})$ and $N^{\text{gen}} := \int_{C(z'_0, R_1)} dm'_{\text{eq}}$ is an integer. Moreover let us assume that $N^{\text{gen}} \sim_{N \rightarrow \infty} m_{\text{eq}}(z_0)N^{2\delta_1}$.

Then there exists a family A_N^{int} of point configurations in $C(z'_0, R_1)$ such that the following properties hold for any $\mathcal{C}^{\text{int}} \in A_N^{\text{int}}$:

1. *The configuration \mathcal{C}^{int} has N^{gen} points in $C(z'_0, R_1)$.*
2. *The continuous average of \mathcal{C}^{int} on $C(z'_0, R_1)$ is close to P , i.e.*

$$(5.1) \quad i_{N, \delta_1}^{z_0}(\mathcal{C}^{\text{int}}) \in B(P, o(1)) \text{ as } N \rightarrow \infty.$$

3. *There exists an electric field $E^{\text{int}} \in \text{Elec}(\mathcal{C}^{\text{int}}, \mu'_{\text{eq}}, C(z'_0, R_1))$ such that*
 - (a) *$E^{\text{int}} \cdot \vec{n} = 0$ on $\partial C(z'_0, R_1)$, where \vec{n} is the unit normal vector.*
 - (b) *The energy of E^{int} is controlled by $\mathbb{W}_{m_{\text{eq}}(z_0)}(P)$: as $N \rightarrow \infty$ we have*

$$(5.2) \quad \frac{1}{2\pi} N^{-2\delta_1} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_1)} |E_\eta^{\text{int}}|^2 + N^{\text{gen}} \log \eta \right) \leq \mathbb{W}_{m_{\text{eq}}(z_0)}(P) + o(1)$$

uniformly on A_N^{int} .

4. *The (logarithmic) volume of the family is close to the relative specific entropy of P (plus constant terms depending on $m_{\text{eq}}(z_0)$)*

$$(5.3) \quad N^{-2\delta_1} \log \mathbf{Leb}^{\otimes N^{\text{gen}}}(A_N^{\text{int}}) \geq -\text{ent}[P|\mathbf{\Pi}^{m_{\text{eq}}(z_0)}] + c_{z_0} + o(1), \text{ as } N \rightarrow \infty.$$

Proof. This follows from the analysis of [LS15, Section 6]. Let us sketch the main steps here.

We fix $R > 0$ and we tile $C(z'_0, R_1)$ by squares of sidelength $\approx R$. We let $\{C_i\}_{i \in I}$ be this collection of squares and x_i be the center of C_i . We sample a point configuration \mathcal{C} in $C(z'_0, R_1)$ according to the law $\mathbf{B}_{N^{\text{gen}}}$ of a Bernoulli point process with N^{gen} points (which is nothing but throwing N^{gen} independent points uniformly in $C(z'_0, R_1)$), and we decompose \mathcal{C} as $\mathcal{C} = \sum_{i \in I} \mathcal{C}_i$ where $\mathcal{C}_i := \mathcal{C} \cap C_i$ is the point configuration in C_i . We form two points processes, the continuous average $M_1 := i_{N, \delta_1}^{z_0}(\mathcal{C})$ and the discrete average $M_2 := \frac{1}{\#I} \sum_{i \in I} \delta_{C_i}$. Classical large deviations

arguments (similar to Section 2.5) show that both M_1 and M_2 belong to $B(P, \varepsilon)$ with probability $\approx \exp(-N^{2\delta_1}(\text{ent}[P|\mathbf{\Pi}^{\mu_{\text{eq}}(z_0)}] - c_{z_0}))$.

Then we apply to each point configuration \mathcal{C}_i the ‘‘screening-then-regularization’’ procedure of [LS15, Section 5]. The screening procedure is similar in spirit to the one described in Lemma 3.8, except that here we change \mathcal{C}_i to $\mathcal{C}_i^{\text{scr}}$ by modifying the configuration only in a thin layer near $\partial\mathcal{C}_i$ and we construct an electric field E_i^{scr} compatible with $\mathcal{C}_i^{\text{scr}}$ and which is screened outside \mathcal{C}_i (whereas in Lemma 3.8 we rather ‘‘screen the configurations from the inside’’). By gluing the fields E_i^{scr} together we define E^{scr} which is compatible with $\mathcal{C}^{\text{scr}} := \sum_{i \in I} \mathcal{C}_i^{\text{scr}}$. The next task is to ‘‘regularize’’ the point configurations, which means to separate the pairs of points which are too close from each other. This changes \mathcal{C}^{scr} into \mathcal{C}^{mod} (which is very much like \mathcal{C}^{scr}) and E^{scr} into an electric field E^{mod} which is still screened outside $C(z'_0, R_1)$.

The energy of E^{scr} can be seen to satisfy, for any $0 < \eta < 1$

$$\int_{C(z'_0, R_1)} |E_\eta^{\text{scr}}|^2 + N^{\text{gen}} \log \eta = \sum_{i \in I} \int_{\mathcal{C}_i} |(E_i^{\text{scr}})_\eta|^2 + N^{\text{gen}} \log \eta,$$

and a certain continuity property of the energy shows that the right-hand side is smaller than $N^{2\delta_1}(\mathbb{W}_{\text{meq}(z_0)}(P) + o(1))$ often enough. Passing from E^{scr} to E^{mod} does not affect this estimate, on the contrary the regularization procedure allows to bound the difference between the truncated energy $\int_{C(z'_0, R_1)} |E_\eta^{\text{mod}}|^2 + N^{\text{gen}} \log \eta$ and its limit as $\eta \rightarrow 0$. \square

5.2 Choice of the deltas

In order to apply Lemma 3.8 (the ‘‘screening lemma’’) at scale δ , we need to find good boundaries as in Section 4.1, which requires to set $\delta_3, \delta_2, \delta_1$ with $\delta_3 < \delta_2 < \delta_1 < \delta$, and satisfying some inequalities dispersed among the various statements. Then the bootstrap argument allows one to pass from the scale δ to the scale $\delta_1 < \delta$. We thus need to set up a way of finding $\delta_3, \delta_2, \delta_1$ such that the required inequalities are satisfied at each step, and such that the multiscale analysis can be pushed up to $N^{o(1)}$. This is the goal of the next lemma.

Lemma 5.2. *Let $\gamma := \sqrt{\frac{1+\kappa/2}{1+\kappa/3}}$, with κ as in (2.1). Since $0 < \kappa \leq 1$ we have $1 < \gamma \leq \frac{3\sqrt{2}}{4} \approx 1.06$. Let also $\alpha := \frac{\gamma-1}{1-\frac{1}{3}}$, we have $\alpha \in (0, 1)$.*

For any $0 < \delta \leq 1/2$ and any $\delta_1, \delta_2, \delta_3$ such that

$$(5.4) \quad \delta > \delta_1 > \max\left(\frac{3}{4}\delta, \delta\frac{1-\alpha}{1-\alpha^2}, \delta(1+\kappa/2) - \kappa/2\right), \delta_3 = \frac{1}{3}\delta\gamma, \quad \delta_2 = \alpha^2\delta_3 + (1-\alpha^2)\delta_1,$$

we have $0 < \delta_3 < \delta_2 < \delta_1 < \delta$ and the following inequalities are satisfied

$$\begin{aligned} \delta_1 &> \frac{2}{3}\delta \\ 3\delta_3 &> \delta \\ \delta_1 + 3\delta_3 + \kappa(\delta_3 - 1/2) &< 2\delta_1 \\ 2\delta &< \delta_2 + 3\delta_3 \\ 2\delta - \delta_2 &< 2\delta_1. \end{aligned}$$

Moreover, if we consider the lower bound on δ_1 as a function $f(\delta)$, i.e.

$$f(\delta) := \max\left(\frac{3}{4}\delta, \delta\frac{1-\alpha}{1-\alpha^2}, \delta(1+\kappa/2) - \kappa/2\right)$$

we have $f^{\circ k}(\delta) \rightarrow 0$ as $k \rightarrow \infty$, where $f^{\circ k}$ denotes the k -th iteration of f .

Proof. It is clear that $\delta_3 > 0$. From the fact that $\delta_1 \geq \delta(1 + \kappa/2) - \kappa/2$ and the expression of δ_3 we get $\delta_1 + 3\delta_3 + \kappa(\delta_3 - 1/2) < 2\delta_1$. Since $\kappa > 0$ we have $3\delta_3 > \delta$. Since $\kappa \leq 1$ we also have $\delta_3 \leq \frac{1}{3}\sqrt{\frac{9}{8}}\delta < \frac{3}{4}\delta$ hence $\delta_3 < \delta_1$. On the other hand from the definition of δ_1 it is clear that $\frac{2}{3}\delta < \delta_1 < \delta$ because $\kappa > 0$ and $\delta \leq 1/2$. Since $\delta_1 > \delta\frac{1-\alpha}{1-\alpha^2}$ and since $\alpha < \alpha^2$, the inequality $2\delta < \delta_2 + 3\delta_3$ follows from checking that $\delta(2 - (1 + \alpha/3)\gamma) \leq (1 - \alpha)\delta$. The inequality $2\delta < \delta_2 + 3\delta_3$ implies the last one, $2\delta - \delta_2 < 2\delta_1$, because it can be easily checked that $3\delta_3 = \delta\gamma < \frac{6}{4}\delta < 2\delta_1$.

Finally, we may observe that $f(\delta) \leq \max(\frac{3}{4}\delta, \frac{1-\alpha}{1-\alpha^2}\delta, \delta - \kappa/4)$ hence $f^{ok}(\delta)$ is decreasing and tends to 0 as $k \rightarrow \infty$. \square

5.3 A LDP lower bound

We use Lemma 5.1 and the screening result of Lemma 3.8 to prove a first LDP lower bound.

Proposition 5.3. *Let $0 < \delta \leq 1/2$ and $z_0 \in \mathring{\Sigma}$ be fixed. Assume that a good control holds at scale δ and let us fix $\delta_1, \delta_2, \delta_3$ as in (5.4). Let us recall that c_{z_0} is the constant defined in (2.7) which appears in (5.3).*

For any $P \in \mathcal{P}_{s, m_{\text{eq}}(z_0)}(\mathcal{X})$ we have

$$(5.5) \quad \log \mathfrak{P}_{N, \beta, \delta_1}^{z_0}(B(P, \varepsilon)) \geq -N^{2\delta_1} \mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + N^{2\delta_1} c_{z_0} - \log K_{N, \beta} \\ + \max_{R_1, R_2, N^{\text{out}}} \log \binom{N}{N^{\text{out}}} \left(\int e^{-\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right) + o(N^{2\delta_1}).$$

The maximum $\max_{R_1, R_2, N^{\text{out}}}$ is taken among $\{R_1, R_2\}$ satisfying (4.1) and (4.11) and with N^{out} between 1 and N .

Proof. By definition of $\mathfrak{P}_{N, \beta, \delta_1}^{z_0}$ and \mathbb{P}_N^β we may forget about the partition function and prove

$$(5.6) \quad \log \int_{(i_{N, \delta_1}^{z_0})^{-1}(B(P, \varepsilon))} e^{-\beta(w_N(\nu'_N) + N\tilde{\zeta}(\nu_N))} d\vec{X}_N \geq -N^{2\delta_1} \mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + N^{2\delta_1} c_{z_0} \\ + \max_{R_1, R_2, N^{\text{out}}} \log \binom{N}{N^{\text{out}}} \left(\int e^{-\beta(F^{\text{out}}(\vec{X}^{\text{out}}) + N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right) + o(N^{2\delta_1}).$$

Let R_1, R_2, N^{out} be fixed. Let \vec{X}^{out} be a finite point configuration in $C(z'_0, R_2)^c$ such that $F^{\text{out}}(\vec{X}^{\text{out}})$ is finite. Let $E \in \text{Elec}^{\text{out}}(\vec{X}^{\text{out}})$ be a minimizer in the definition of F^{out} . We claim that there exists a set A^{tot} of N -tuples \vec{X}_N such that $\vec{X}_N = \vec{X}^{\text{out}}$ on $C(z'_0, R_2)^c$, that the energy is controlled uniformly on A^{tot} as follows

$$w_N(\nu'_N) \leq -(\mathbb{W}_{m_{\text{eq}}(z_0)}(P) + F^{\text{out}}(\vec{X}^{\text{out}})) + o(N^{2\delta_1}),$$

and that the volume of A^{tot} is almost optimal

$$\log \text{Leb}^{\otimes N}(A^{\text{tot}}) \geq -N^{2\delta_1} \text{ent}[P | \mathbf{\Pi}^{m_{\text{eq}}(z_0)}] + N^{2\delta_1} c_{z_0} + \log \binom{N}{N^{\text{out}}} + o(N^{2\delta_1}).$$

Step 1. Screening E . We may apply Lemma 3.8 to the point configuration \vec{X}^{out} and the electric field E , with $\mu = \mu'_{\text{eq}}$. Let us check that the assumptions of Lemma 3.8 are satisfied:

1. The first condition on R_1, R_2 is satisfied by assumption (see (4.1) and (4.11)).
2. Since $F^{\text{out}}(\vec{X}^{\text{out}})$ is finite, (4.2) holds i.e. the smeared out charges at scale η_1 do not intersect $\partial C(z'_0, R_2)$.

3. The third and fourth condition on $N^{\text{gen}}, N^{\text{mid}}$ follow from the fact that R_1 is a good interior boundary.

From Assumption 1 we know that $\mu'_{\text{eq}} \leq 1$, and since we are blowing-up the configuration around $z_0 \in \mathring{\Sigma}$ the density m'_{eq} is bounded below on $C(z'_0, N^\delta)$ by some $\underline{m} > 0$ depending on z_0 . We deduce from (2.1) that $|\mu'_{\text{eq}}(x) - \mu'_{\text{eq}}(y)| \leq N^{-\kappa/2}|x - y|^\kappa$, hence we may chose $C_\mu = N^{-\kappa/2}$ in (3.8).

By definition of a good exterior boundary (see (4.5)) we have $\int_{\partial C(z'_0, R_2)} |E_{\eta_1}|^2 \leq N^{2\delta - \delta_2} \log^2 N$, thus (3.9) is satisfied (for N large enough) as long as $2\delta < \delta_2 + 3\delta_3$ (which is ensured by the choice (5.4)).

We obtain a family A_N^{tran} of point configurations such that the conclusions of Lemma 3.8 hold.

Step 2. Generating microstates. Now we apply Lemma 5.1 with R_1 as above and obtain a family A_N^{int} of point configurations with N^{gen} points in $C(z'_0, R_1)$ together with screened electric fields E^{int} such that (5.1), (5.2), (5.3) are satisfied.

Step 3. Gluing pieces together and bounding the energy. For any $\mathcal{C}^{\text{tran}} \in A_N^{\text{tran}}$ and $\mathcal{C}^{\text{int}} \in A_N^{\text{int}}$ we form the configuration

$$\mathcal{C}^{\text{tot}} := \mathcal{C}^{\text{tran}} + \mathcal{C}^{\text{int}} + \vec{X}^{\text{out}}.$$

It is easy to check that \mathcal{C}^{tot} always has N points. Indeed we know that

- By Lemma 5.1, \mathcal{C}^{int} always has $\int_{C(z'_0, R_1)} d\mu'_{\text{eq}}$ points.
- By construction, $\mathcal{C}^{\text{tran}}$ has $N^{\text{tran}} = \int_{\partial C(z'_0, R_2)} E_{\eta_1} \cdot \vec{n} - \int_{C(z'_0, R_2) \setminus C(z'_0, R_1)} d\mu'_{\text{eq}}$ points.
- By integrating the compatibility relation of E and \vec{X}^{out} , we get

$$N^{\text{out}} = N - \int_{\partial C(z'_0, R_2)} E_{\eta_1} \cdot \vec{n} + \int_{C(z'_0, R_2)} d\mu'_{\text{eq}}.$$

If E^{tran} and E^{int} are the electric fields associated to $\mathcal{C}^{\text{tran}}$ and \mathcal{C}^{int} we also define

$$E^{\text{tot}} := E^{\text{tran}} + E^{\text{int}} + E\mathbf{1}_{C(z'_0, R_2)^c}$$

By construction the normal derivatives of E^{tran} and E^{int} coincide on $\partial C(z'_0, R_1)$ (they both vanish), and the normal derivatives of E_η^{tran} and E_η coincide on $\partial C(z'_0, R_2)$ for any $\eta \leq \eta_1$ (they coincide for η_1 by construction, but since there are no points at distance $\leq \eta_1$ of $\partial C(z'_0, R_2)$ the value of the fields E_η and E_η^{tran} on $\partial C(z'_0, R_2)$ do not depend on η for $\eta \leq \eta_1$). Thus E^{tot} satisfies

1. $-\text{div } E^{\text{tot}} = 2\pi(\mathcal{C}^{\text{tot}} - \mu'_{\text{eq}})$ in \mathbb{R}^2
2. E^{tot} coincides with E on $C(z'_0, R_2)^c$. In particular E^{tot} belongs to Elec^0 , as E does.
3. The energy of E^{tot} is bounded as follows

$$(5.7) \quad \lim_{\eta \rightarrow 0} \left(\int_{\mathbb{R}^2} |E_\eta^{\text{tot}}|^2 + N \log \eta \right) \leq F^{\text{out}}(\vec{X}^{\text{out}}) + N^{2\delta_1} \mathbb{W}_{m_{\text{eq}}(z_0)}(P) + o(N^{2\delta_1}) \text{ as } N \rightarrow \infty.$$

To show (5.7), let us split the energy of E^{tot} as

$$(5.8) \quad \int_{\mathbb{R}^2} |E_\eta^{\text{tot}}|^2 + N \log \eta = \left(\int_{C(z'_0, R_2)^c} |E_\eta|^2 + N^{\text{out}} \log \eta \right) + \left(\int_{C(z'_0, R_2) \setminus C(z'_0, R_1)} |E_\eta^{\text{tran}}|^2 + N^{\text{tran}} \log \eta \right) + \left(\int_{C(z'_0, R_1)} |E_\eta^{\text{int}}|^2 + N^{\text{gen}} \log \eta \right)$$

By definition of $F^{\text{out}}(\vec{X}^{\text{out}})$ and by the choice of E we have

$$(5.9) \quad \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2)^c} |E_\eta|^2 + N^{\text{out}} \log \eta \right) = F^{\text{out}}(\vec{X}^{\text{out}}).$$

In view of (5.2) we have

$$(5.10) \quad \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_1)} |E_\eta^{\text{int}}|^2 + N^{\text{gen}} \log \eta \right) \leq N^{2\delta_1} \mathbb{W}_{m_{\text{eq}}(z_0)}(P) + o(N^{2\delta_1}).$$

Finally, the conclusions of Lemma 3.8 combined with the control (4.5) and the fact that $C_\mu = N^{-\kappa/2}$ ensure that

$$\begin{aligned} \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2) \setminus C(z'_0, R_1)} |E_\eta^{\text{tran}}|^2 + N^{\text{tran}} \log \eta \right) &\leq N^{\delta_3} N^{2\delta - \delta_2} (\log N)^2 + N^{\delta_1 + 3\delta_3} N^{\kappa(\delta_3 - 1/2)} \\ &\quad + N^{\delta_1 + \delta_3} \log N. \end{aligned}$$

The choice of $\delta_1, \delta_2, \delta_3$ as in (5.4) yields

$$(5.11) \quad \lim_{\eta \rightarrow 0} \left(\int_{C(z'_0, R_2) \setminus C(z'_0, R_1)} |E_\eta^{\text{tran}}|^2 + N^{\text{tran}} \log \eta \right) \ll N^{2\delta_1}.$$

Inserting (5.9), (5.10) and (5.11) into (5.8) yields (5.7).

Now, using the minimality of local energy as stated in Lemma 3.7 and the formula (2.4) we conclude that

$$(5.12) \quad w_N(\mathcal{C}^{\text{tot}}) \leq \frac{1}{2\pi} \lim_{\eta \rightarrow 0} \left(\int_{\mathbb{R}^2} |E_\eta^{\text{tot}}|^2 + N \log \eta \right) \leq F^{\text{out}}(\vec{X}^{\text{out}}) + N^{2\delta_1} \mathbb{W}_{m_{\text{eq}}(z_0)}(P) + o(N^{2\delta_1}) \text{ as } N \rightarrow \infty.$$

Step 4. Volume considerations. For any \vec{X}^{out} we let $A^{\text{tot}}(\vec{X}^{\text{out}})$ be the set of point configurations \mathcal{C}^{tot} obtained as above. Now, let \mathcal{A} be a measurable set of finite point configurations \vec{X}^{out} with N^{out} points in $C(z'_0, R_2)^c$ such that $F^{\text{out}}(\vec{X}^{\text{out}})$ is finite for all $\vec{X}^{\text{out}} \in \mathcal{A}$. We let \mathcal{A}^{tot} be

$$\mathcal{A}^{\text{tot}} := \bigcup_{\vec{X}^{\text{out}} \in \mathcal{A}} A^{\text{tot}}(\vec{X}^{\text{out}}),$$

Using the volume estimate (3.14) we obtain, with the choice (5.4)

$$\log \mathbf{Leb}^{\otimes N}(\mathcal{A}^{\text{tot}}) \geq \log \mathbf{Leb}^{\otimes N^{\text{out}}}(\mathcal{A}) + \log \mathbf{Leb}^{\otimes N^{\text{gen}}}(A_N^{\text{int}}) + o(N^{2\delta_1}) + \log \binom{N}{N^{\text{out}} N^{\text{tran}} N^{\text{gen}}},$$

where the last term denotes a multinomial coefficient. This coefficient comes from the fact that the elements of a point configuration are not labeled. In particular in any $\mathcal{C}^{\text{tot}} \in \mathcal{A}^{\text{tot}}$ we cannot determine which points belong to the interior, exterior or “transition” part. When enumerating the total point configurations we can thus multiply the volume of the resulting “ordered” configurations by $\binom{N}{N^{\text{out}} N^{\text{tran}} N^{\text{gen}}}$. Using (5.3) and a straightforward combinatorial inequality $\binom{a}{bcd} \geq \binom{a}{b}$ yields

$$(5.13) \quad \log \mathbf{Leb}^{\otimes N}(\mathcal{A}^{\text{tot}}) \geq \log \mathbf{Leb}^{\otimes N^{\text{out}}}(\mathcal{A}) - N^{2\delta_1} \text{Ent}[P | \mathbf{\Pi}^{m_{\text{eq}}(z_0)}] + N^{2\delta_1} c_{z_0} + o(N^{2\delta_1}) + \log \binom{N}{N^{\text{out}}}.$$

This proves the claim made before Step 1.

Step 5. Conclusion. Combining (5.12) and (5.13), we obtain that

$$\begin{aligned} \log \int_{(i_{N,\delta_1}^{z_0})^{-1}(B(P,\varepsilon))} e^{-\frac{1}{2}\beta(w_N(\nu'_N)+N\tilde{\zeta}(\nu_N))} d\vec{X}_N \\ \geq \log \binom{N}{N^{\text{out}}} \left(\int e^{-\frac{1}{2}\beta(F^{\text{out}}(\vec{X}^{\text{out}})+N\tilde{\zeta}(\vec{X}^{\text{out}}))} d\vec{X}^{\text{out}} \right) \\ - N^{2\delta_1} \left(\text{Ent}[P|\mathbf{\Pi}^{\text{meq}}(z_0)] + \frac{1}{2}\beta\mathbb{W}_{\text{meq}}(z_0)(P) \right) + N^{2\delta_1}c_{z_0} + o(N^{2\delta_1}), \end{aligned}$$

for any choice of R_1, R_2 and N^{out} as in the definitions 4.1 and 4.2. It yields (5.6). \square

6 Conclusion

6.1 Proof of Theorem 1

6.1.1 Good control at macroscopic scale

The next lemma is used to initiate the multiscale analysis by proving that good control of the energy holds at the largest scale $\delta = \frac{1}{2}$. A slightly stronger version, with deviations at speed N and not $N^{1-o(1)}$, can be deduced from the analysis [SS15b] (see e.g. Theorem 5.2).

Lemma 6.1. *Good control holds at scale $\delta = \frac{1}{2}$.*

Proof. Let $z_0 \in \mathring{\Sigma}$. Using (1.9) we see that

$$\log \mathbb{P}_N^\beta(w_N(\vec{X}_N) \geq NM) \leq -\log K_{N,\beta} + \log \int e^{-\frac{\beta}{2}(NM+N\tilde{\zeta}(\vec{X}_N))} d\vec{X}_N.$$

Since w_N is bounded below by $O(N)$ (see Lemma 2.1) we have $-\log K_{N,\beta} + \log \left(\int e^{-\frac{\beta}{2}N\tilde{\zeta}(\vec{X}_N)} \right) = O(N)$ (which may depend on β), hence we get

$$\log \mathbb{P}_N^\beta \left(w_N(\vec{X}_N) \geq M \right) \leq N \left(-\frac{\beta}{2}M + O(1) \right).$$

We deduce that for M_0 large enough, $\limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}_N^\beta(w_N(\vec{X}_N) \geq M_0) < 0$, which in particular implies that $\limsup_{N \rightarrow \infty} \frac{1}{N^{2\delta}} \log \mathbb{P}_N^\beta(w_N(\vec{X}_N) \geq M_0) < 0$ for any $\delta < \frac{1}{2}$, thus we have $w_N(\vec{X}_N) \preceq N$ with δ -overwhelming probability. Using (2.4) and Lemma 3.1 we get for any $\eta \in (0, 1)$ that

$$\lim_{\eta \rightarrow 0} \int_{\mathbb{R}^2} |E_\eta^{\text{loc}}|^2 + N \log \eta \preceq_\delta N,$$

which yields (1.14), and we deduce (1.13) from the discrepancy estimates of Lemma 3.4. \square

6.1.2 Exponential tightness

Lemma 6.2. *For any $0 < \delta \leq 1/2$, if good control holds at scale δ then $\mathfrak{P}_{N,\beta,\delta_1}^{z_0}$ is exponentially tight (at speed $N^{2\delta_1}$) for any $z_0 \in \mathring{\Sigma}$ and $\frac{2}{3}\delta < \delta_1 < \delta$.*

Proof. Let $z_0 \in \mathring{\Sigma}$ and $\frac{2}{3}\delta < \delta_1 < \delta$ be fixed. The good control at scale δ , combined with Lemma 3.6, implies that there exists $C > 0$ such that the number of points in $C(z'_0, N^{\delta_1})$ is bounded above by $CN^{2\delta_1}$ with δ_1 -overwhelming probability. It implies that $\mathfrak{P}_{N,\beta,\delta_1}^{z_0}$ is concentrated on the compact subset

$$\{P \in \mathcal{P}(\mathcal{X}), \mathbf{E}_P[\mathcal{N}(0, R)] \leq CR^2 \quad \forall R > 0\},$$

with δ_1 -overwhelming probability, which ensures exponential tightness at speed $N^{2\delta_1}$. \square

6.1.3 Proof of the theorem

Proof of Theorem 1.

Step 1. *Good control \implies LDP.*

In this first step we claim that if a good control holds at scale δ , then the LDP of Theorem 1 holds for δ_1 as in (5.4).

Indeed, comparing the right-hand side of (5.5) with the definition (4.30) of K_{N,z,δ_1}^β we see that $\lim_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N,z,\delta_1}^\beta$ exists and that the following weak large deviation principle holds (with c_{z_0} as in (2.7))

$$(6.1) \quad \lim_{\varepsilon \rightarrow 0} \lim_{N \rightarrow \infty} N^{-2\delta_1} \log \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(B(P, \varepsilon)) = -\mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0} - \lim_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N,z,\delta_1}^\beta,$$

for any $P \in \mathcal{P}_{s, m_{\text{eq}}(z_0)}(\mathcal{X})$, hence since $\mathbb{W}_{m_{\text{eq}}(z_0)}(P) = +\infty$ as soon as P is not of intensity $m_{\text{eq}}(z_0)$, we may write (6.1) for any $P \in \mathcal{P}_s(\mathcal{X})$. By exponential tightness we obtain a full large deviation inequality: for any measurable $A \subset \mathcal{P}_s(\mathcal{X})$ it holds

$$(6.2) \quad -\inf_{P \in A} \mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0} - \lim_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N,z,\delta_1}^\beta \\ \leq \liminf_{N \rightarrow \infty} N^{-2\delta_1} \log \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(A) \leq \limsup_{N \rightarrow \infty} N^{-2\delta_1} \log \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(A) \\ \leq -\inf_{P \in A} \mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0} - \lim_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N,z,\delta_1}^\beta.$$

In particular, taking $A = \mathcal{P}_s(\mathcal{X})$ we obtain

$$(6.3) \quad \lim_{N \rightarrow \infty} N^{-2\delta_1} \log K_{N,z,\delta_1}^\beta = -\inf \mathcal{F}_\beta^{m_{\text{eq}}(z_0)}(P) + c_{z_0},$$

and inserting (6.3) into (6.2) yields the LDP for $\{\mathfrak{P}_{N,\beta,\delta_1}^{z_0}\}_N$ as stated in Theorem 1.

Step 2. *Good control \implies good control.*

We now claim that if a good control holds at scale δ , then it holds at scale δ_1 with δ_1 as in (5.4). Combining the ‘‘good control upper bound’’ of Lemma 4.5 and the lower bound estimates which yield (6.3) we deduce that

$$(6.4) \quad \log \mathfrak{P}_{N,\beta,\delta_1}^{z_0}(\mathcal{E}_M) \preceq -MN^{2\delta_1} + O(N^{2\delta_1}),$$

where \mathcal{E}_M is as in (4.23). In particular it implies that $\int_{C(z'_0, N^{\delta_1})} |E_{\eta_0}^{\text{loc}}|^2 + \mathcal{N}_{\delta_1}^{z_0} \log \eta_0 \preceq_{\delta'} N^{2\delta_1}$ for any $\delta' < \delta_1$ and any $\eta_0 \in (0, 1)$. We also have $\mathcal{N}_{\delta_1}^{z_0} \preceq_{\delta_1} N^{2\delta_1}$ (since it was proven in Lemma 4.3 that (4.9) holds with δ_1 -overwhelming probability) hence in particular $\mathcal{N}_{\delta_1}^{z_0} \preceq_{\delta'} N^{2\delta_1}$ for $\delta' < \delta_1$.

Step 3. *Conclusion.*

Combining both steps with the initialization of Lemma 6.1 and the conclusions of Lemma 5.2 yields the proof of Theorem 1. \square

6.2 Proof of Corollary 1.4

Proof. We simply combine the fact that a good control holds at any scale $0 < \delta \leq 1/2$ (which follows from Theorem 1) with Lemma 3.5. \square

6.3 Proof of Corollary 1.5

Proof. We may split $C(z'_0, N^\delta)$ into a family $\{C_i\}_{i \in I}$ of squares of sidelength $\approx N^{\delta_1}$, with $\#I \approx N^{2(\delta - \delta_1)}$. For any $i \in I$ we have, letting z_i be the center of C_i and \mathcal{D}_i the discrepancy in C_i

$$\int_{C_i} f(d\nu'_N - d\mu'_{\text{eq}}) = \mathcal{D}_i f(z_i) + \int_{C_i} (f(z) - f(z_i))(d\nu'_N - d\mu'_{\text{eq}}).$$

Since good control holds at scale δ we have $|\mathcal{D}_i| \preceq_{\delta_1} N^{4\delta/3}$ (from the discrepancy estimates of Lemma 3.5) and $\int_{C(z'_0, N^\delta)} d\nu'_N \preceq_{\delta_1} N^{2\delta}$. On the other hand $\int_{C(z'_0, N^\delta)} d\mu'_{\text{eq}} \preceq N^{2\delta}$ (because μ'_{eq} is bounded above). Moreover the mean value theorem yields $|f(z) - f(z_i)| \leq N^{\delta_1} \|\nabla f\|_\infty$. We thus have

$$\left| \sum_{i \in I} \int_{C_i} f(d\nu'_N - d\mu'_{\text{eq}}) \right| \leq N^{2(\delta - \delta_1)} N^{4\delta/3} \|f\|_\infty + N^{2\delta} N^{\delta_1} \|\nabla f\|_\infty,$$

hence we see that

$$N^{-2\delta} \left| \int_{C(z'_0, N^\delta)} f(d\nu'_n - d\mu'_{\text{eq}}) \right| \preceq_{\delta_1} \|\nabla f\|_\infty N^{\delta_1} + \|f\|_\infty N^{-2\delta/3},$$

which concludes the proof since $\delta_1 < \delta$. \square

7 Additional proofs

7.1 Proof of Proposition 2.4

Proof. First, let us treat the case $m = 1$. It follows from [GZ93, Theorem 3.1] (with $F = 0$) that, if $\mathbf{\Pi}_R^1$ denotes the law of a Poisson point process of intensity 1 in C_R , we have for any P stationary

$$(7.1) \quad \lim_{\varepsilon \rightarrow 0} \lim_{R \rightarrow \infty} \frac{1}{R^2} \log \mathbf{\Pi}_R^1(i_R(\mathcal{C}) \in B(P, \varepsilon)) = -\text{ent}[P|\mathbf{\Pi}^1].$$

It is not hard to extend (7.1) to the case where the points are distributed according to a Bernoulli point process with R^2 points in C_R , i.e. the number of points is set to R^2 (here and in the rest of the proof, for simplicity we do not take integer values). Describing the Bernoulli point process in an explicit fashion yields

$$(7.2) \quad \lim_{\varepsilon \rightarrow 0} \lim_{R \rightarrow \infty} \frac{1}{R^2} \frac{R^2!}{(R^2)^{R^2}} \int_{(C_R)^{R^2}} \mathbf{1}_{i_R(x_1, \dots, x_{R^2}) \in B(P, \varepsilon)} dx_1 \dots dx_{R^2} = -\text{ent}[P|\mathbf{\Pi}^1],$$

which we re-write (using Stirling's formula) as

$$(7.3) \quad \int_{(C_R)^{R^2}} \mathbf{1}_{i_R(x_1, \dots, x_{R^2}) \in B(P, \varepsilon)} dx_1 \dots dx_{R^2} = -R^2 (\text{ent}[P|\mathbf{\Pi}^1] - 1 + o(1)),$$

where the $o(1)$ goes to zero as $R \rightarrow \infty, \varepsilon \rightarrow 0$ (in this order). It proves Proposition 2.4 in the case $m = 1$ (the fact that the number of points is not exactly R^2 but $R^2(1 + o(1))$ is not important).

In order to treat the general $m > 0$ case, let us change variables by setting $y_i = m^{-1/2}x_i$. We have

$$(7.4) \quad \int_{(C_R)^{R^2}} \mathbf{1}_{i_R(x_1, \dots, x_{R^2}) \in B(P, \varepsilon)} dx_1 \dots dx_{R^2} \\ = m^{R^2} \int_{(C_{m^{-1/2}R})^{R^2}} \mathbf{1}_{i_R(m^{1/2}y_1, \dots, m^{1/2}y_{R^2}) \in B(P, \varepsilon)} dy_1 \dots dy_{R^2}.$$

One can check that having $i_R(m^{1/2}y_1, \dots, m^{1/2}y_{R^2})$ close to P is equivalent to having $i_R(y_1, \dots, y_{R^2})$ close to $\sigma_m^{-1}(P)$ (which has density m , with σ_m defined in Section (1.3)). On the other hand, let us set $S = m^{-1/2}R$. Combining (7.3) and (7.4) and sending S (or R) to $+\infty$ and then ε to 0 we obtain

$$\int_{(C_S)^{mS^2}} \mathbf{1}_{i_R(y_1, \dots, y_{R^2}) \in B(\sigma_m^{-1}(P), \varepsilon)} dy_1 \dots dy_{R^2} = -mS^2 (\text{ent}[P|\mathbf{\Pi}^1] - 1 + \log m + o(1)).$$

Writing $Q = \sigma_m^{-1}(P)$, it is not difficult to check that

$$\text{ment}[P|\mathbf{\Pi}^1] = \text{ent}[Q|\mathbf{\Pi}^m],$$

e.g. by a simple scaling argument (see also [LS15, Lemma 4.2] together with a change of reference measure in the classical relative entropy). We thus obtain

$$\int_{(C_S)^{mS^2}} \mathbf{1}_{i_R(y_1, \dots, y_{R^2}) \in B(Q, \varepsilon)} dy_1 \dots dy_{R^2} = -S^2 (\text{ent}[Q|\mathbf{\Pi}^m] - m + m \log m + o(1)),$$

which proves the proposition. \square

7.2 Proof of Lemma 3.3

Proof. We may decompose E^{loc} as $E^{\text{in}} + E^{\text{out}}$ where E^{in} is the local electric field generated by the electric system *inside* C_{R_2} and E^{out} is the local electric field generated by the electric system *outside* C_{R_2} . We have

$$\int_{C_{R_2}} |E_\eta^{\text{loc}}|^2 = \int_{C_{R_2}} |E_\eta^{\text{in}}|^2 + \int_{C_{R_2}} |E_\eta^{\text{out}}|^2 + 2 \int_{C_{R_2}} E_\eta^{\text{in}} \cdot E_\eta^{\text{out}}.$$

Since the charges outside C_{R_2} are at distance at least η_1 from ∂C_{R_2} we may replace E_η^{out} by E^{out} in the previous identity (in fact we have $E_\eta^{\text{out}} = E^{\text{out}}$ on C_R for $\eta \leq \eta_1$). Integrating by parts we obtain

$$\int_{C_{R_2}} |E_\eta^{\text{loc}}|^2 = \int_{C_{R_2}} -H_\eta^{\text{in}} \Delta H_\eta^{\text{in}} + \int_{\partial C_{R_2}} H_\eta^{\text{in}} E_\eta \cdot \vec{n} + \int_{\partial C_{R_2}} H_\eta^{\text{in}} E^{\text{out}} \cdot \vec{n}$$

(up to additive terms which do not depend on $\eta \leq \eta_1$), where H^{in} is the local electric potential generated by the electric system inside C_{R_2} . By assumption we have $H_\eta^{\text{in}} = H_{\eta_1}^{\text{in}}$ and $E_\eta^{\text{in}} = E_{\eta_1}^{\text{in}}$ on ∂C_{R_2} . Finally we see that

$$\int_{C_{R_2}} |E_\eta^{\text{loc}}|^2 - |E_{\eta_1}^{\text{loc}}|^2 = - \int_{\mathbb{R}^2} (H_\eta^{\text{in}} \Delta H_\eta^{\text{in}} - H_{\eta_1}^{\text{in}} \Delta H_{\eta_1}^{\text{in}})$$

for any $\eta \leq \eta_1$, and (3.3) is obtained as Lemma 3.1 (*cf.* the remark after the statement of Lemma 3.1). \square

7.3 Proof of Lemma 3.7

Proof. The neutrality of the system implies that the local electric potential H^{loc} decays like $|x|^{-1}$ as $|x| \rightarrow \infty$ in \mathbb{R}^2 and E^{loc} decreases like $|x|^{-2}$. This can be seen by using Taylor's formula and the fact that the electric system is compactly supported

$$(7.5) \quad \begin{aligned} H^{\text{loc}}(x) &= - \int \log|x-y|(d\mathcal{C} - d\mu)(y) \\ &= - \log|x| \int (d\mathcal{C} - d\mu)(y) + O\left(\frac{1}{|x|}\right) \left(\int d\mathcal{C}(y) + d\mu(y) \right) = O(|x|^{-1}), \end{aligned}$$

and a similar argument goes for E^{loc} .

If the right-hand side of (3.7) is infinite then there is nothing to prove. If it is finite, given $M > 1$ and letting χ_M be a smooth nonnegative function equal to 1 in C_M and 0 at distance ≥ 1 from C_M , we may write

$$(7.6) \quad \begin{aligned} \int_{\mathbb{R}^2} \chi_M |E_\eta|^2 &= \int_{\mathbb{R}^2} \chi_M |E_\eta - E_\eta^{\text{loc}}|^2 + \int_{\mathbb{R}^2} \chi_M |E_\eta^{\text{loc}}|^2 + 2 \int_{\mathbb{R}^2} \chi_M (E_\eta - E_\eta^{\text{loc}}) \cdot E_\eta^{\text{loc}} \\ &\geq \int_{\mathbb{R}^2} \chi_M |E_\eta^{\text{loc}}|^2 + 2 \int_{\mathbb{R}^2} \chi_M (E_\eta - E_\eta^{\text{loc}}) \cdot (\nabla H_\eta^{\text{loc}}) \\ &= \int_{\mathbb{R}^2} \chi_M |E_\eta^{\text{loc}}|^2 + 2 \int_{\mathbb{R}^2} H_\eta^{\text{loc}} (E_\eta - E_\eta^{\text{loc}}) \cdot \nabla \chi_M, \end{aligned}$$

where we have integrated by parts and we have used the fact that E, E^{loc} are compatible with the same configuration (hence $\text{div}(E_\eta - E_\eta^{\text{loc}}) = 0$). Letting $M \rightarrow \infty$, the last term tends to 0 by finiteness of the right-hand side of (3.7), the decay properties of H^{loc} and E^{loc} and the decay assumption on E . \square

7.4 Auxiliary estimate for screening

Lemma 7.1. *Let $l > 0$ and let H be a rectangle of \mathbb{R}^2 with sidelengths in $(l/2, 3l/2)$. Let $g \in L^2(\partial H)$ and let m be a function on H of average $m_0 := \frac{1}{|H|} \int_H m$ such that*

$$(7.7) \quad -2\pi m_0 |H| = \int_{\partial H} g.$$

Then there exists a solution h to $-\Delta h = 2\pi m$ in H with $\nabla h \cdot \vec{n} = 0$ on ∂H satisfying

$$(7.8) \quad \int_H |\nabla h|^2 \leq l \int_{\partial H} |g|^2 + l^4 \|m - m_0\|_{L^\infty(H)}^2.$$

Proof. A solution exists thanks to the compatibility condition (7.7). We may split h as $h_1 + h_2$ where h_1 solves

$$-\Delta h_1 = 2\pi m_0 \text{ in } H, \quad \nabla h_1 \cdot \vec{n} = g \text{ on } \partial H,$$

and h_2 is the mean zero solution to

$$-\Delta h_2 = 2\pi(m - m_0) \text{ in } H, \quad \nabla h_2 \cdot \vec{n} = 0 \text{ on } \partial H.$$

In view of [RS15, Lemma 5.8] we may find h_1 satisfying

$$(7.9) \quad \int_H |\nabla h_1|^2 \leq l \int_{\partial H} |g|^2.$$

We also claim that

$$(7.10) \quad \int_H |\nabla h_2|^2 \leq l^4 \|m - m_0\|_{L^\infty(H)}^2.$$

Indeed it is easy to check that (7.10) holds when $l = 1$, and the general case follows by a scaling argument.

Combining (7.9) and (7.10) concludes the proof. \square

References

- [AHM15] Y. Ameur, H. Hedenmalm, and N. Makarov. Random normal matrices and Ward identities. *Ann. Probab.*, 43(3):1157–1201, 2015.
- [AJ81] A. Alastuey and B. Jancovici. On the classical two-dimensional one-component Coulomb plasma. *Journal de Physique*, 42(1):1–12, 1981.
- [BAZ98] G. Ben Arous and O. Zeitouni. Large deviations from the circular law. *ESAIM: Probability and Statistics*, 2:123–134, 1998.
- [BBNY15] R. Bauerschmidt, P. Bourgade, M. Nikula, and H.-T. Yau. Local density for two-dimensional one-component plasma. <http://arxiv.org/abs/1510.02074>, 10 2015.
- [BEY12] P. Bourgade, L. Erdős, and H.-T. Yau. Bulk universality of general β -ensembles with non-convex potential. *J. Math. Phys.*, 53(9):095221, 19, 2012.
- [BEY14] P. Bourgade, L. Erdős, and H.-T. Yau. Universality of general β -ensembles. *Duke Math. J.*, 163(6):1127–1190, 2014.
- [BL15] X. Blanc and M. Lewin. The crystallization conjecture: A review. <http://arxiv.org/abs/1504.01153>, 2015.
- [BYY14a] P. Bourgade, H.-T. Yau, and J. Yin. Local circular law for random matrices. *Probab. Theory Related Fields*, 159(3-4):545–595, 2014.
- [BYY14b] P. Bourgade, H.-T. Yau, and J. Yin. The local circular law II: the edge case. *Probab. Theory Related Fields*, 159(3-4):619–660, 2014.
- [DE02] I. Dumitriu and A. Edelman. Matrix models for beta ensembles. *J. Math. Phys.*, 43:5830–5847, 2002.
- [DZ10] A. Dembo and O. Zeitouni. *Large deviations techniques and applications*, volume 38 of *Stochastic Modelling and Applied Probability*. Springer-Verlag, Berlin, 2010. Corrected reprint of the second (1998) edition.
- [FO88] H. Föllmer and S. Orey. Large deviations for the empirical field of a gibbs measure. *The Annals of Probability*, pages 961–977, 1988.
- [For10] P. J. Forrester. *Log-gases and random matrices*, volume 34 of *London Mathematical Society Monographs Series*. Princeton University Press, Princeton, NJ, 2010.
- [Fro35] O. Frostman. Potentiel d'équilibre et capacité des ensembles avec quelques applications à la théorie des fonctions. *Meddelanden Mat. Sem. Univ. Lund* 3, 115 s, 1935.

- [Geo93] H.-O. Georgii. Large deviations and maximum entropy principle for interacting random fields on \mathbb{Z}^d . *The Annals of Probability*, pages 1845–1875, 1993.
- [Gin65] J. Ginibre. Statistical ensembles of complex, quaternion, and real matrices. *J. Mathematical Phys*, 6:440–449, 1965.
- [GZ93] H.-O. Georgii and H. Zessin. Large deviations and the maximum entropy principle for marked point random fields. *Probab. Theory Related Fields*, 96(2):177–204, 1993.
- [Har12] A. Hardy. A note on large deviations for 2D Coulomb gas with weakly confining potential. *Electron. Commun. Probab.*, 17:no. 19, 12, 2012.
- [JLM93] B. Jancovici, J. Lebowitz, and G. Manificat. Large charge fluctuations in classical Coulomb systems. *J. Statist. Phys*, 72(3-4):773–7, 1993.
- [Leb15] T. Leblé. Logarithmic, Coulomb and Riesz energy of point processes. <http://arxiv.org/abs/1509.05253>, 2015.
- [LS] T. Leblé and S. Serfaty. Fluctuations of two-dimensional coulomb gases. In preparation.
- [LS15] T. Leblé and S. Serfaty. Large Deviation Principle for empirical fields of Log and Riesz gases. <http://arxiv.org/abs/1502.02970>, 2015.
- [Oll88] S. Olla. Large deviations for gibbs random fields. *Probability Theory and related fields*, 77(3):343–357, 1988.
- [PH98] D. Petz and F. Hiai. Logarithmic energy as an entropy functional. *Contemporary Mathematics*, 217:205–221, 1998.
- [PS15] M. Petrache and S. Serfaty. Next order asymptotics and renormalized energy for Riesz interactions. *Journal of the Institute of Mathematics of Jussieu*, FirstView:1–69, 5 2015.
- [RAS09] F. Rassoul-Agha and T. Seppäläinen. *A course on large deviation theory with an introduction to Gibbs measures*, volume 162 of *Graduate Studies in Mathematics*. American Mathematical Society, 2015 edition, 2009.
- [RNS15] S. Rota Nodari and S. Serfaty. Renormalized energy equidistribution and local charge balance in 2D Coulomb systems. *Int. Math. Res. Not. IMRN*, (11):3035–3093, 2015.
- [RS15] N. Rougerie and S. Serfaty. Higher-dimensional Coulomb gases and renormalized energy functionals. *Communications on Pure and Applied Mathematics*, 2015.
- [Ser15] S. Serfaty. *Coulomb Gases and Ginzburg-Landau Vortices*. Zurich Lecture Notes in Mathematics, 2015.
- [SM76] R. Sari and D. Merlini. On the ν -dimensional one-component classical plasma: the thermodynamic limit problem revisited. *J. Statist. Phys.*, 14(2):91–100, 1976.
- [SS12] E. Sandier and S. Serfaty. From the Ginzburg-Landau model to vortex lattice problems. *Comm. Math. Phys.*, 313:635–743, 2012.
- [SS15a] E. Sandier and S. Serfaty. 1D log gases and the renormalized energy: crystallization at vanishing temperature. *Probab. Theory Related Fields*, 162(3-4):795–846, 2015.

- [SS15b] E. Sandier and S. Serfaty. 2D Coulomb gases and the renormalized energy. *Ann. Probab.*, 43(4):2026–2083, 2015.
- [ST97] E. B. Saff and V. Totik. *Logarithmic Potentials with External Fields*. Grundlehren der mathematischen Wissenschaften **316**, Springer-Verlag, Berlin, 1997.
- [TV15] T. Tao and V. Vu. Random matrices: universality of local spectral statistics of non-Hermitian matrices. *Ann. Probab.*, 43(2):782–874, 2015.