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1	Transition Periods in the Diurnally-Varying Atmospheric Boundary
2	Layer Over Land
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20	Abstract
21	The atmospheric boundary layer undergoes transitions between stable and convective states. Over
22	land, in undisturbed conditions, these transitions occur daily in the morning and late afternoon or
23	early evening. Though less well studied and presenting more challenges than the fully stable and

fully convective states, such transitions have been the subject of growing interest over the last few decades. During transitions, all forcings are weak, and few simplifications are possible. Factors such as terrain, radiation, advection, and subsidence can seldom be safely neglected. In this paper, we review research on transitions over recent decades, with an emphasis on work published in *Boundary-Layer Meteorology*. The review is brief and inevitably reflects the interests and views of the authors.

30 Keywords Afternoon transition · Evening transition · Morning transition · Boundary-layer field
 31 studies · Boundary-layer simulation

32 **1 Introduction**

33 Over land under clear skies and in the absence of any synoptic disturbance, the atmospheric 34 boundary layer (ABL) is often deep and strongly turbulent during the day, shallow and weakly 35 turbulent at night. In between, the ABL undergoes morning and afternoon/evening transitions. 36 Those transitions have been the subject of considerable research in the last 25 years, much of which 37 has been published in *Boundary-Layer Meteorology*. Figure 1 is a schematic showing a large-eddy simulation (LES) of the diurnal cycle of an idealized mid-latitude ABL, see the caption for more 38 39 details. Shown are the times of sunrise, crossover (change of sign of the surface buoyancy flux), 40 onset of convective turbulence at a height of 200 m, rising of the entrainment zone and growth of 41 the mixed layer during the day, cessation of growth and eventual shrinkage during the afternoon, 42 and the times of crossover and sunset during the evening transition, followed by the growth of the 43 stable boundary layer (SBL), and the subsidence of the capping inversion overnight. These features are discussed in more detail below. We do not provide precise definitions of the transitions 44 because they themselves evolve through the literature reviewed. Each section addresses definitions 45 as appropriate. 46

47 Motivations for the study of transitions are various. For air-quality forecasts and analyses, the 48 timing of changes in the depth of mixing with respect to time-varying emissions is important. The 49 evening transition is key to the evolution of the low-level jet (Smith et al. 2019), which often 50 influences severe weather. Fog and frost forecasting depend on understanding the evening 51 transition (Boutle et al. 2018; Price 2019).



Fig. 1 A schematic illustration of the diurnal cycle of the ABL derived from a large-eddy simulation. **a** The evolution of the vertical gradient of the horizontally-averaged potential temperature, θ , with height above the surface. Vertical white lines indicate the times of sunrise, crossover and onset during the morning transition, while vertical black lines show the times of crossover and sunset during the evening transition. The light grey line marks the top of the stable

56 boundary layer, defined as the height at which the linearly-extrapolated horizontally-averaged buoyancy flux becomes 57 zero. The dashed magenta line marks the height of the elevated maximum of the gradient of θ . The crimson dashed line 58 indicates the height of the minimum of the buoyancy flux during the period when the boundary layer is unstable. 59 Crimson dotted lines mark the upper and lower boundaries of the entrainment zone during the same period, the former 60 defined as the lowest height above that of the minimum of the buoyancy flux at which the absolute value 61 of the buoyancy flux first falls below 5% of the surface value, and the latter as the lowest height at which the buoyancy 62 flux changes sign. The figure also shows the rising of the entrainment zone through the morning, the development of 63 the well-mixed layer, and, after the evening transition, the development of the nocturnal stable boundary layer, the 64 radiative stabilization of the residual layer, and the subsidence of the capping inversion. **b** The evolution of the net 65 downward shortwave flux at the surface (left axis) and the surface buoyancy flux (right axis). c The evolution of the 66 square root of the variance of the vertical velocity with height. Lines as in \mathbf{a} and also showing the pre-residual layer.

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Transitions are complex because many processes are roughly equally important, and the simplifications and approximations that work well for the fully convective or the fully stable boundary layer do not apply. In budget analyses, more terms must be retained, often including direct radiative heating of the air, mixing due to shear, and advection. Stationarity may not apply. Full descriptions of the surface energy balance and of the energy balance at the top of the ABL are not included in this review, reflecting the contents of the literature we are considering.

Most of our increased understanding of the transition periods has been enabled by ground-based remote sensing measurements (Wilczak et al. 1996), since measurements at the surface provide only a limited perspective. The ability to view continuously the full depth of the boundary layer and especially its upper region is critical to understanding transitions. Traditional tall-tower measurements and observations taken by aircraft have also contributed.

While this should not be mistaken for a full description of the diverse possibilities, considering the diurnal cycle is an improvement over looking only at the fully convective and fully stable states. Throughout this review, we remain mindful of ways in which the real atmosphere frequently departs from the idealization. Horizontal inhomogeneity of the surface, land use, soil and vegetation properties has strong impacts. Coasts create land and sea breezes. Terrain induces mountain and valley breeze systems. The daytime ABL is not always strongly turbulent, and the nocturnal boundary layer is not always shallow.

The depth of the boundary layer is an ambiguous but important variable in discussions of the ABL 86 87 in general and transition periods in particular. When the convective boundary layer (CBL) is 88 topped by a well-defined increase in virtual potential temperature, its depth is fairly precisely 89 measured by any of several definitions (e.g. turbulence intensity, virtual-temperature profile, 90 virtual-temperature-flux profiles, or Richardson number) (LeMone et al. 2013). The depth of the 91 stable boundary layer is fundamentally more difficult to define, since the turbulence-intensity 92 profile may have a variety of shapes. Clouds add another level of complexity. Measurements of 93 the depth of an aerosol layer, for example by lidar, can be useful but must be carefully examined to 94 avoid mistaking the residual layer for the current boundary layer (Hogan et al. 2009; Tucker et al. 95 2009). The preferred technique is to measure the profile of vertical velocity variance, but this 96 approach can produce ambiguous results because of gravity waves near the boundary-layer top 97 (Lothon et al. 2009). Whether the ABL is stable or convective, some definitions can be readily 98 derived from NWP or LES model output, but rarely or with difficulty in observations. Throughout 99 this review, we will use the terms "boundary-layer depth" and "boundary-layer top" as if they were 100 well-defined, without tediously acknowledging all of this complexity every time the term comes 101 up.

The review proceeds as follows: First, the morning and afternoon/evening transitions as seen in observations are described in Sects. 2 and 3. In Sect. 4, key processes involved are discussed, including entrainment, radiation, advection and subsidence, length and time scales, and clouds. Sections 5 and 6 describe the current understanding provided by LES and mesoscale numerical weather prediction (NWP) models respectively, and the challenges that continue to arise in the modelling of the diurnal cycle. Finally, we provide an outlook and thoughts on needs for ongoing research.

109 **2** The Morning Transition in Observations

110 The role of mixing from above the near-surface layer relative to heating from the surface was an 111 important question in early research on the morning transition. Data from the Cabauw tower in the 112 Netherlands were used by Driedonks (1982) to test analytical models of the early-morning 113 boundary-layer growth. He recognized that entrainment was important in many cases, and his 114 preferred model included a term proportional to the surface stress, based on theory provided by 115 Tennekes (1973). Aircraft data in Australia were used by Coulman (1978) to test predictions of 116 morning boundary-layer growth, with mixed results. He found that entrainment was probably 117 important, and that estimating some parameters with suficient precision was not always possible, 118 particularly the inversion temperature gradient. Several expressions were provided by Garratt 119 (1992) along with a useful figure showing the change in potential-temperature profile for a 120 transition driven by surface heating vs. shear. A large dataset was examined by Angevine et al. 121 (2001) using data from the Cabauw tower (van Ulden et al. 1996) and radar wind-profiler data from 122 the Cabauw site and from the Flatland (Bondville, Illinois) site (Angevine et al. 1998). They 123 defined two milestones in course of the transition, "crossover" (when the surface buoyancy flux 124 first becomes positive) and "onset" (when convective turbulence is first detected at 200 m above

125 the ground). Sunrise was used as a reference time (Fig. 1). The most interesting discovery from 126 this study was that the time required to erode the stable layer from below, using the measured heat 127 flux and ignoring entrainment, was much longer than the time over which the layer was actually 128 eroded. In other words, the surface heat flux was insufficient to warm the stable layer in the 129 The conclusion was, and remains, that much of the warming arises from observed time. 130 entrainment at the top of the eroding stable layer. The data showed a situation more like Garratt's 131 Fig. 6.7b for a shear-driven transition. This was foreshadowed by Wyngaard (1983) in a brief 132 remark indicating that only a small heat flux was required to switch from a stable to a convective 133 layer.

134 We must caution that the definition of entrainment used in the above papers is imprecise. 135 Formally, entrainment is "The process by which turbulent fluid within a mixed layer incorporates 136 adjacent fluid that is nonturbulent, or much less turbulent." (American Meteorological Society 137 Glossary of Meteorology). However, in the context of the morning transition, within a statically 138 stable layer, the turbulent state of the air at each level is often not known. Some of the air that is 139 mixed to produce a near-neutrally-stratified layer may already have been turbulent and therefore 140 formally part of the boundary layer. The unknown turbulent state makes the formal definition 141 difficult to apply rigorously.

142 It should be noted that the measurements used by Angevine et al. (2001) were influenced by a 143 general effort in previous years to understand and exploit the capabilities of ground-based remote 144 sensing instruments (e.g. (Angevine et al. 1994; Beyrich et al. 1995; Carter et al. 1995).

145 A six-year dataset from the rural resarch site at Cardington in central England was used by 146 Lapworth (2006) to examine the morning and evening transitions, with an emphasis on wind speed 147 and forecasting rules. He confirmed the finding that most of the warming during the morning transition was due to mixing from above. Before crossover, the warming is better described as eddy diffusion rather than entrainment, since the boundary-layer top may not be rising. The nearsurface air temperature is controlled by the balance between radiative cooling and turbulent mixing, which, under weak to moderate stability, is continuous. As the radiative cooling decreases, the turbulent heat flux from above warms the near-surface layer, before the surface heat flux becomes positive (upward).

154 Figure 2 shows observed soundings during a morning transition at Cardington. Entrainment is 155 evidenced by cooling in the upper part of the growing boundary layer in the progression of 156 observed soundings, especially between 0400 and 0700 UTC for example. Profiles from an 157 operational NWP model are also shown. The simulated boundary layer is consistently too cool and 158 too shallow, probably because of insufficient entrainment. Advection is important in most 159 individual cases, but is not observed, so attribution of the changes in profiles due to entrainment is 160 suggestive rather than conclusive; see Angevine et al. (2001) for more cases and discussion of their 161 variations.

An interesting feature of the morning transition is a peak in the time series of near-surface specific humidity that is sometimes observed, e.g. (LeMone et al. 2002). When transpiration by vegetation becomes significant while the mixed layer is still shallow, the humidity may exhibit a local maximum, and then decline when the mixed layer grows. The morning peak is not seen in the Flatland or Cabauw data of Angevine et al. (2001), but was seen at Cabauw on the day used by Holtslag et al. (1995), and seems therefore to depend on the season and vegetation type. This has been further studied for the afternoon/evening transition by Blumberg et al. (2019).

169 It is notable that multi-day averages are necessary to make understandable behaviour emerge from170 data. Individual days exhibit a wide range of behaviour, dominated by advection, which is hardly

ever negligible (Baas et al. 2010), but sometimes averages out over multiple days (Betts et al. 1996). Even mild terrain can produce phenomena that overwhelm column-wise analyses, as was found by Lenschow et al. (1979). Some influence of heterogeneity cannot be ruled out in the Flatland and Cabauw cases, although the Flatland site in particular is extremely uniform for a land site.

Further development of observational equipment and techniques continues to add to our knowledge of the morning transition. Wildmann et al. (2015) conducted a study with a small remotely-piloted aircraft, emphasizing the scaling behaviour of variances compared with theory, and they include a thorough discussion of uncertainties and challenges.

180 **3** Afternoon and Evening Transition in Observations

181 From a surface perspective, the boundary layer changes from unstable to stable in the "evening," 182 some time near sunset, but looking at the boundary-layer top, changes happen much earlier in the 183 day. This has led to a lack of clarity and consensus in terminology. In this review, we refer to the 184 afternoon or evening transition according to the terms used by the cited papers. Many of these 185 terms are illustrated in Fig. 1. In general, "afternoon transition" denotes the period when the surface 186 buoyancy flux decreases (from midday to sunset), and "evening transition" refers to the period near 187 sunset including the change of sign of the surface buoyancy flux. A sunset-relative frame of 188 reference was used by Wingo et al. (2015). In complex terrain, local sunset is a key time, 189 especially when slopes are steep (Nadeau et al. 2011). In Fig. 1, we have indicated sunset and 190 crossover, newly defined here for the evening transition as the time when the surface buoyancy flux 191 changes sign from upward to downward, the opposite of morning crossover.

192 Mahrt (1981) made an early observational study of the "early evening" transition based on data 193 from the Wangara (Australia) study, emphasizing the behavior of the wind and the low-level jet. 194 Mahrt observed that the wind changes began well before sunset. Many studies of the evening 195 transition consider slightly or moderately complex terrain. Acevedo et al. (2001) and Mahrt (2017) 196 emphasized the evolution of temperature and wind through three phases. The cooling is small at the 197 start (before zero flux), greater in the second phase (around zero net radiation), and decreases again 198 after sunset. The wind decelerates, changes to calm or local drainage, and evolves toward regional 199 drainage or synoptic flow, as also shown by LeMone et al. (2003), Busse et al. (2012), Román-200 Cascón et al. (2015), and Nilsson et al. (2016a). Those phases all depend on the characteristics of 201 the studied sites. A peak in the time series of specific humidity may also be present, and is 202 dependent on the season, location, and crop behavior (Blumberg et al. 2019).

Recent works have highlighted new features like the counter-gradient flux (Blay-Carreras et al. 2014a), caused by a delay between the times when the surface flux changes sign and when the temperature gradient changes sign.. It has been shown to be partly related to the eddy turnover time, and dependent on the vegetation cover and height (Jensen et al. 2016). Also the lifted temperature minimum is worth mentioning (Blay-Carreras et al. 2015), that is, a minimum of temperature at 0.1-1 m above the surface.

A detailed study by Grant (1997) used data from tethered balloon equipment at the Cardington site to describe the latter part of the transition period, after the surface buoyancy flux had changed to negative (downward). Grant showed profiles and spectra to explore the applicability of proposed scaling frameworks during the evening transition. Further data from the same site and similar equipment were used by Edwards et al. (2006). Observations of turbulent kinetic energy near the surface were used by Nadeau et al. (2011), who fitted a simple analytical model which approximately showed a t^2 power-law decay of turbulent kinetic energy (TKE) early in the transition, and t^{-6} later (*t* being time).

217 Observations through the depth of the boundary layer in the afternoon and evening were rare until 218 ground-based remote sensing instruments became available. The first full analysis was undertaken 219 by Grimsdell et al. (2002) using wind-profiling radar observations in conjunction with surface-flux 220 measurements. They found that the height of the boundary layer decreases in the afternoon, 221 starting some hours before sunset. The explanation is this: as solar heating of the surface begins to 222 decline, and the air continues to warm, thermals become less vigorous. Since the bulk of the 223 boundary layer is roughly adiabatic, the weaker thermals can still reach the boundary-layer top at 224 first. Meanwhile, several other processes are at work, and all of these tend to reduce the height of 225 Subsidence directly lowers the boundary-layer top. Subsidence, shear-driven thermals. 226 entrainment, and direct radiative heating all act to stabilize the upper part of the boundary layer. In 227 the case of subsidence, regime selection is at work: If large-scale upward motion is present, clouds 228 and precipitation often form, and these days are seldom considered for further case study analysis. 229 Direct measures of turbulence show the decreasing top height, while indirect measures, particularly 230 aerosol laver height, may not. Since the boundary layer is defined as being in recent contact with 231 the surface (Stull 1988; LeMone et al. 2019), the direct turbulence measurement is preferred 232 (Tucker et al. 2009; Barlow et al. 2011; Harvey et al. 2013; Manninen et al. 2018).

Lidar observations of the boundary layer over London, UK covering several full diurnal cycles were published by Barlow et al. (2011). The urban boundary layer has additional complexity not seen in simpler situations, but some of its characteristics are similar. Long lidar datasets from two sites were analyzed by Manninen et al. (2018). 237 The Boundary Layer Late Afternoon and Sunset Turbulence (BLLAST) field study was focused on 238 these transitions. It provided a comprehensive set of measurements by all types of instruments and 239 platforms. Unlike most earlier studies, the BLLAST study embraced the challenge of complex 240 terrain and heterogeneous land use. Lothon et al. (2014) provides an overview, early results, and a 241 rich bibliography. The project started with careful definitions, and in addition to measurements, 242 incorporated theory and modelling. Frequent and continuous observations through the whole depth 243 of the boundary layer allowed more complete analysis than had previously been possible. The 244 decay of TKE in a vertically resolved sense, and the ways in which turbulent length scales change, 245 were particular foci.

246 In terms of the scaling of turbulence within the boundary layer, the BLLAST study results from 247 aircraft (Lothon et al. 2014; Darbieu et al. 2015) and near-surface measurements (Nilsson et al. 248 2016a; Nilsson et al. 2016b) indicated that statistics of turbulence on days with weak shear are 249 quasi-stationary until late afternoon. Shear delays the decay, or even prevents decay if strong 250 enough (Pino et al. 2006). While the decay remains slow enough for statistics to be quasi-251 stationary, the turbulence structure does not evolve substantially, and remains scaled with the 252 boundary-layer depth. Later in the afternoon, the decay accelerates and the turbulence structure 253 departs from CBL scaling. The decay was found to start from the top (Darbieu et al. 2015), 254 although this strongly depends on the shear at surface and at boundary layer top. Consistent with 255 this, Nilsson et al. (2016b) used a simple model to identify a region of reduced turbulence that 256 appears earlier in the upper part of the boundary layer, before the stable layer develops, and expands downward with time. They described this as a "pre-residual" layer. Scaling by the 257 258 buoyancy flux clearly becomes questionable as it approaches zero.

259 4 Processes of Boundary-Layer Transitions

260 The afternoon transition is nearly unforced; all contributing processes are weak. The morning 261 transition is forced by the rising sun, but this forcing is surprisingly weak as well. Many processes influence the boundary layer through the diurnal cycle, and the key challenge to understanding is 262 263 the need to include all relevant processes. Turbulent eddy sizes and time scales change. 264 Entrainment at the boundary-layer top and mixing of entrained air play important roles. Advection 265 is almost always important when examining individual days and hours, and is difficult to estimate. 266 Direct radiative heating or cooling of the air (as opposed to heating or cooling of the surface) is 267 often small, but cannot be safely neglected during transitions. When fluxes or turbulent energy are 268 changing rapidly, methods of analysis that neglect time rates of change may fail. Clouds may 269 affect the profiles of mean and turbulent quantities, especially near the top of the boundary layer. 270 This section describes these processes in more detail.

271 4.1 Eddy Length and Time Scales

Time scales are fundamental to transitions, in fact to the idea of transition itself. Some relevant time scales include the eddy turnover time, the time required for the boundary-layer height to change by a given fraction, and the time required for the surface buoyancy flux to change by a given fraction (or to reverse sign). Around midday, the eddy turnover time (10–20 min) is much shorter than the time scale of the evolution of the boundary-layer height, which changes only slowly. In the morning growth phase, however, the boundary layer may grow as rapidly as a thermal rises (Sorbjan 1997).

The time scale for turbulence to decay is important in the afternoon and evening, and depends on the eddy size (Pino et al. 2006; Lothon et al. 2014). Larger eddies decay more slowly, leading to an effective increase in the length scale above the surface layer. 282 The impact of variation in surface fluxes on time scales in the CBL was explored by Van Driel et 283 al. (2011) in an idealized framework. They made several interesting findings: the time scale for 284 changes in the boundary-layer depth is very long, of order 10 h, because it is controlled by the 285 entrainment velocity. A phase lag between the changing surface flux and the mixed-layer height is 286 due to these differing time scales. For rapidly varying fluxes, the production and dissipation of 287 TKE become important. Their simulations included subsidence and radiative cooling specifically 288 tuned to produce a steady boundary-layer depth and temperature. The application of these results 289 to transitions is challenging, because, for example, the entrainment velocity during the morning 290 transition can be very large (shortening the time scale). . Subsidence and other processes that are 291 roughly in balance at midday, allowing the boundary-layer depth to remain approximately constant, 292 become unbalanced when the surface flux permanently weakens in the afternoon, resulting in rapid 293 changes of the boundary-layer depth. Nonetheless, the concept of competing time scales and 294 transfer functions introduced by Van Driel et al. (2011) is a promising framework for future 295 exploration.

Another important time scale from a practical point of view is the averaging time required in processing turbulence measurements, whether from a sonic anemometer or a remote sensor. If important changes happen during the averaging time, the measurements are hard to interpret (Mahrt 1981; Grant 1997). On the other hand, averaging over too short a time gives noisy and excessively uncertain results, and can lead to systematic underestimation of fluxes.

Length scales and their variation are also vital to understanding boundary-layer behaviour. With the dataset from the BLLAST study, Darbieu et al. (2015) found that the length scales of turbulence remain scaled with the boundary-layer depth until late afternoon (more precisely until about 0.75 times the decay-forcing time scale). Later, they observed a sharp decrease of the length scale in the surface layer, and an increase of the length scale above. The changes occurred first at the top of the
boundary layer, although this depends on shear.

307 4.2 Entrainment

As described in sections 2 and 3 above and sections 5 and 6 below, entrainment at the boundary-308 layer top is critical to the transitions. In the morning, most of the warming during the transition is 309 310 due to mixing from above, either from within the upper part of the boundary layer or from 311 entrainment. In the afternoon, entrainment is one of the processes that causes warming in the upper 312 part of the boundary layer and thereby reduces the depth of mixing. Entrainment is a complex topic that deserves a whole review by itself; see Angevine (2008) for further discussion and 313 314 references. It is particularly worth noting that entrainment relationships depend strongly on the 315 interpretive framework, e.g. Liu et al. (2017). Figure 2 shows the importance of entrainment in the 316 morning transition, and insufficient entrainment appears to be the cause of the cool and shallow 317 bias in the morning boundary layer in the UKV NWP model.



Fig. 2 Profiles of potential temperature at the Cardington field site during the clear-sky morning of 19 July 2016.
Radiosonde ascents launched on each hour are shown as solid lines, together with tower sensors (square points) at 2 m,
10 m, 25 m and 50 m; horizontal bars on the mast data show variations of +/- 5 mins centred on each hour. Model data
(dotted lines) are the instantaneous hourly forecast profiles from the Met Office UKV operational 1.5 km grid spacing
model, initialized at 0300 UTC.

A nearly adiabatic residual layer often exists above the nocturnal boundary layer. In the morning, when the growing boundary layer reaches the base of the residual layer, its growth rate increases until the residual-layer top is reached (Deardorff 1974; Sorbjan 1996; Blay-Carreras et al. 2014b). It is useful to note, however, that the residual layer is not simply the remains of the previous day's mixed layer at the same location; because ofdifferential advection, the mixed layer found in the morning above a certain site was likely formed somewhere else.

331 Experience with setting up cases for single-column model (SCM) and LES comparisons has shown 332 that advection is rarely negligible, even at nearly ideal sites. Measurements of advection are 333 difficult or impossible. The best that can be done is to estimate advection terms from observations at levels that are thought not to be affected by turbulent mixing, or to use output from mesoscale 334 335 models (Barr et al. 1997; Angevine 1999). Estimates of subsidence are at least as problematic. 336 Efforts involving many dropsondes have succeeded in producing subsidence measurements for the 337 simplest possible situation over water (Bony et al. 2018). It should be noted, however, that 338 comparisons between LES and other types of models using the same forcings are not subject to 339 these uncertainties.

340 Baas et al. (2010) discuss the issue of specifying forcing (advection) for SCM studies. They 341 demonstrate strong impacts on the formation of the low-level jet (an evening-transition issue) from 342 different plausible specifications of advection. These differences make it difficult to determine 343 which vertical mixing formulation is best. They also show that a composite case, constructed by averaging several similar cases, removes some of the "noise" and allows for clearer conclusions. 344 345 Different choices of forcing, all from plausible sources, had a strong impact on shallow cumulus 346 simulations shown by Angevine et al. (2018). Averaging over many cases was done by Angevine et 347 al. (2001) to understand the morning transition. Systematic advection can mislead the analyst 348 despite averaging. For example, in the central United States, the prevailing wind direction is 349 southerly (warm and moist) in summer and northerly in winter, so advection cannot be expected to 350 average out within a season.

The effects of subsidence (Blay-Carreras et al. 2014b) and advection (Pietersen et al. 2015) were studied for some BLLAST cases. At the BLLAST site, there was a persistent diurnal pattern of subsidence, induced by the nearby mountain range and accentuated by moist convection over themountains.

355 4.4 Radiation

356 Vertical turbulent heat fluxes driven by net radiation at the surface play a substantial role in the 357 diurnal evolution of the boundary layer, but direct atmospheric radiative heating is also relevant. 358 Radiative cooling at cloud top can generate turbulence and affect the entrainment rate (Stevens 359 2002; de Lozar et al. 2015). Clear-sky radiative processes have received less attention, but can play 360 a significant role in transitions. Longwave radiative cooling erodes the capping inversion and 361 accelerates the growth of the boundary layer through the morning transition (Edwards et al. 2014). Radiative cooling near the surface during the evening transition stabilizes the developing stable 362 363 boundary layer and suppresses turbulence (Garratt et al. 1981; Ha et al. 2003). In near-calm 364 conditions radiation can be the dominant process determining the temperature profile (Savijarvi 2006; Edwards 2009). Longwave radiative cooling rates vary on short length scales near the 365 366 surface, so poor vertical resolution can produce incorrect results in either observations or models 367 (Ha et al. 2003; Edwards 2009; Steeneveld et al. 2010).

Aerosols in the boundary layer may have large effects on radiative heating or cooling when the aerosol loading is significant or the aerosols are strongly absorbing. Aerosol layers within or above the boundary layer reduce the net radiation reaching the surface (Barbaro et al. 2014; Liu et al. 2019), thereby affecting transition timing. Aerosols within the boundary layer or the residual layer can directly heat those layers enough to affect timing as well. The effects are complex and subtle because the sun is low in the sky during transitions. Barbaro et al. (2014) give a detailed discussion of these effects, supported by model studies of two observed cases.

376 Clouds are very common at the top of the boundary layer. The University of Washington Cloud 377 Atlas (https://atmos.uw.edu/CloudMap/index.html; accessed 15 March 2019) shows that 26% of 378 land and 55% of the oceans is covered with low cloud on an annual average basis. Most of these 379 clouds are in the boundary layer. Nonetheless, strongly cloud-influenced conditions are commonly 380 avoided or selected against in boundary-layer field research. During the morning transition, 381 Angevine et al. (2001) noted that the most extreme day-to-day variations in timing were found on 382 days with cloud. In their data from the Flatland site, cloud-free days showed more rapid increases 383 in temperature than cloudy days, but no systematic differences were found between cloud-free and 384 cloud-influenced days at Cabauw. In the afternoon, Grimsdell et al. (2002) found little systematic 385 influence of clouds. Over water, where the diurnal cycle is damped, Pennell et al. (1974) and 386 LeMone (1978) found little influence of clouds on subcloud-layer TKE budgets. Brown et al. 387 (2002) found that subcloud-layer buoyancy-flux profiles were insensitive to cloud cover. In all of 388 these analyses, regime selection probably prevented strong cloud influences.

An interesting and complex but not uncommon case is presented by Hogan et al. (2009). A thin stratocumulus layer in the early morning is replaced by a convective boundary layer with shallow cumulus. Radiative cooling in the stratocumulus layer produces mixing, which must affect the timing of the morning transition. Similar cases probably arise commonly, and deserve further exploration. Clouds and their effects are also seen in the data of Harvey et al. (2013) and Manninen et al. (2018).

5 Transitions in large-eddy simulations and direct numerical simulations

396 Large-eddy simulations of transitions are challenging because they require both fine resolution and 397 large domains, and are thus computationally expensive (Beare 2008). Several LES studies of 398 transitions have been done, however. Beare (2008) produced simulations of the morning transition 399 by starting with a small, finely resolved domain and expanding to a larger, more coarsely resolved 400 domain partway through the transition. These simulations started at the time of heat-flux 401 crossover, so they describe the late phase of the transition. Beare found a shallow convective 402 boundary layer with penetrating, entraining updrafts under the remnant stable layer, which he 403 named the "mixed CBL-SBL state". The morning transition in the third GEWEX Atmospheric 404 Boundary Layer Studies (GABLS3) case was simulated by Edwards et al. (2014), who 405 demonstrated the effect of radiation in weakening the capping inversion, thus accelerating the 406 growth of the CBL.

407 Evening-transition LESs have a longer history. It seems to have been assumed that the resolution 408 suitable for convective boundary layers was also sufficient to represent the transition. Nieuwstadt 409 et al. (1986) used an abrupt cut-off of the surface heat flux. Sorbian (1997) introduced and showed 410 the importance of a smoother reduction, which introduces a second timescale. Pino et al. (2006) 411 used LES to show that horizontal length scales become larger during the afternoon transition. Beare 412 et al. (2006) may have been the first to simulate the evening transition at fine resolution, with more 413 realistic fluxes. Early simulations had simplified representations of non-turbulent processes such 414 as radiation, which was shown to be important in light winds by Edwards (2009). Few LES models 415 have an interactive land surface, and the choice of surface boundary condition (flux or temperature) 416 has been controversial. Details of initialization and forcing were found to be important, e.g. Kumar 417 et al. (2010). Observed surface temperature was used by Rizza et al. (2013) to study the evening 418 transition, and they found TKE-decay power-law exponents of -2 early and -6 late in the

transition, similar to those observed by Nadeau et al. (2011). El Guernaoui et al. (2019) explored
the TKE decay in the context of interacting time scales.

421 Both morning and afternoon transitions were simulated by Conzemius et al. (2008) for one day in 422 Oklahoma. In their LES model, they simulated very large relative contributions of shear-driven 423 entrainment to the total warming during the morning and afternoon. These are caused by the 424 realistic nature of the case. Profiles and TKE budgets are shown. Basu et al. (2008) simulated a 425 diurnal cycle from Wangara, emphasizing the low-level jet that appeared after the evening 426 transition. Their evening transition was abrupt, probably due to the lack of radiation and/or 427 subsidence in their setup. Full diurnal cycles were also simulated by Sharma et al. (2017), again 428 finding an abrupt evening transition without a representation of radiation and with no subsidence. 429 Both morning and evening transitions were included in the LES of Blay-Carreras et al. (2014b), 430 who emphasized the importance of the profiles above the growing boundary layer in the morning 431 and the forcing by subsidence in the afternoon. Darbieu et al. (2015) used a BLLAST case to set 432 the surface flux and initial conditions, and NWP model outputs to account for total advection 433 (including subsidence). Although the amount of TKE was different than in observations, and the 434 real conditions were much more complex than in the LES, the way that the TKE and the turbulence 435 structure evolved were very similar to the observations, showing a certain universality of the decay, 436 well reproduced by LES.

Figure 1 illustrates the evolving sophistication of coupled LES models. Itwas performed using the large-eddy Met Office-NERC Cloud model (MONC) (Brown et al. 2015) coupled to the Joint UK Land Environment Simulator (JULES) land-surface model (Best et al. 2011) and the Suite of Community Radiative Transfer codes based on Edwards and Slingo (SOCRATES) radiation code (Manners et al. 2015). The specifications of the simulation were chosen for visual clarity in the 442 figure. The latitude was set as 57° N and a geostrophic wind speed of 8 m s⁻¹ was prescribed. The 443 land surface was specified as short grass and the initial atmospheric profiles of temperature and 444 specific humidity were chosen to avoid the formation of cloud. The period shown begins on 2 445 March.

446 Direct numerical simulation (DNS) has recently become practical for application to atmospheric 447 problems. Mellado et al. (2016) simulated a free-encroachment case relevant to the late morning 448 transition. Fedorovich et al. (2017) performed DNS simulations relevant to the formation of the 449 low-level jet at the end of the evening transition.

450 6 Transitions in Mesoscale Numerical Models

451 Operational NWP models simulate transitions and the diurnal cycle regularly over large areas. 452 These models inherently simulate subsidence and advection. Radiation, boundary-layer physics, 453 clouds, and interaction of the atmosphere with the surface are parametrized. Climate-model 454 simulations likewise include many diurnal cycles under all kinds of conditions. However, research 455 specifically targeting transitions in mesoscale models is relatively rare in the literature. The 456 GEWEX Atmospheric Boundary Layer Study (GABLS) is a long-running project to improve the 457 simulation of the boundary layer in models and has proceeded through increasingly realistic 458 intercomparisons. Two of the three GABLS studies looked at diurnal cycles over land, and provide 459 the most complete view of mesoscale model performance during the transitions. The results show 460 considerable room for improvement. An overview of GABLS is given by Holtslag et al. (2013).

461 For GABLS2, Svensson et al. (2011) simulated a diurnal cycle from the CASES-99 experiment in
462 Kansas, USA. Many single-column models participated. Their behaviour during the transitions
463 was diverse. One conclusion was that choosing the most realistic or most useful boundary

464 conditions and forcings is both difficult and critical. Specified surface temperatures were used as 465 the bottom boundary condition, and it became clear that the fluxes produced by the different 466 models for the same atmospheric and surface temperatures diverged. Heat flux during the morning 467 transition was too small and the boundary layer grew too slowly. This may have been due to too 468 little entrainment (see above), which also led to wind speeds that were too low. Different kinds of 469 boundary-layer scheme in the models had different kinds of errors, but no type of scheme was 470 clearly better for the whole diurnal cycle.

471 The GABLS3 study (Bosveld et al. 2014) benefited from the lessons learned in GABLS2. The case 472 was derived from measurements at Cabauw, The Netherlands. The single-column models included 473 land-surface and radiation modules, removing the need for specified lower boundary conditions. 474 Several of the conclusions of GABLS3 involved the land-surface formulations, which dominated 475 the transitions and stable boundary-layer behaviour. Another important conclusion was the 476 importance of longwave radiation from within the stable boundary layer (see above). Evening-477 transition timing varied by 1.5 h among the models. In the morning, the simulated time of 478 crossover from downward to upward surface flux varied by 2 hours between models, and the 479 interval from crossover to onset of turbulence at 200 m was about half an hour too long.

Using a single-column model, Edwards et al. (2006) simulated two evening-transition cases, as observed at Cardington Observed surface temperature was prescribed to the model as a lower boundary condition. They emphasized the importance of specifying realistic advection. In both cases, the simulated boundary layer was too shallow during the transition. They found limited sensitivity to the resolution or the parameterization of turbulence in the stable boundary layer. At the moderate wind speeds in these cases, longwave radiative cooling was a significant, but not the dominant process in the development of the nocturnal boundary layer. Edwards (2009) described a 487 case of very light winds (about 1 m s⁻¹) and strong surface cooling and showed that radiative 488 cooling to the surface was the dominant process determining the evolution of the near-surface air 489 temperature through the transition. It was noted that the assumptions underlying standard surface-490 layer similarity theory failed during the transition, leading to cold biases in operational forecasts in 491 such conditions.

Extensive observations of temperature profiles were examined by van Hooijdonk et al. (2017) to study the growth rate of the surface-based temperature inversion, and compared with bulk and single-column models of varying complexity. They found that the simplest models could not reproduce the observations, because history effects were important. This reinforces the idea that the boundary layer is not in equilibrium during some parts of the evening transition.

497 Mesoscale modelling for the BLLAST campaign also pointed out the importance of the land 498 surface. Lothon et al. (2014) and Couvreux et al. (2016) found variations of the vegetation among 499 different measurement sites to be an important contribution to variations in afternoon transition 500 timing. Angevine et al. (2014) devised a technique for making soil moisture match a limited-area 501 model better to produce more realistic simulations. Couvreux et al. (2016) found that a finer-502 resolution NWP model (around 2 km grid spacing) resolves the vertical structures better than the 503 coarser models (around 10 km), in particular a well-defined maximum of the water-vapour mixing 504 ratio that arises during the evening in the upper part of the residual layer due to fine-scale 505 advection. The finer- resolution model also provides a good simulation of the diurnal variability of 506 the turbulent kinetic energy, including its late afternoon decay, while the coarser model still shows 507 the right order of magnitude.



Fig. 3 (Top) surface and minimum sensible heat flux from a single-column simulation with the TEMF scheme and two LES models for 6 June 2015 at the ARM SGP site in Oklahoma. Potential temperature flux is shown for SCM and LES1, virtual potential temperature flux for LES2. See text for details. (Bottom) Ratio of minimum flux to surface flux, one way of calculating the entrainment ratio.

512 Diagnosing entrainment in mesoscale or NWP models is difficult because they rarely produce flux-513 profile output by default. One exception is the Total Energy Mass Flux (TEMF) boundary-layer 514 scheme available in the Weather Research and Forecast model (WRF) (Angevine et al. 2018). 515 Figure 3 shows the diurnal cycle of the surface kinematic heat flux and the minimum flux in the 516 profile within the boundary layer for 6 June 2015 at the Atmospheric Radiation Measurement 517 Southern Great Plains site in Oklahoma. Along with the single-column model running TEMF, two 518 LESs of that day are also shown. The ratio of entrainment flux to surface flux is commonly used in 519 parameterizations, even though it is strictly applicable only to pure free convection driven from the 520 surface, and different methods of calculation yield different ratios (Conzemius et al. 2006; Liu et al. 521 2017). With those caveats in mind, the lower portion of Fig. 3 shows the entrainment ratio through 522 the diurnal cycle. TEMF has large ratios of the magnitude of entrainment flux to that of the surface 523 flux early and late in the day, consistent with some of the observational studies. TEMF also has a 524 relatively large ratio in midday. Breaks in the curve are due to limited vertical resolution in the 525 SCM. Both LESs show some of the same tendency to larger ratios in the morning and afternoon 526 transition periods. LES1 is from the U.S. Department of Energy LES Atmospheric Radiation Measurement Symbiotic Simulation and Observation Workflow (LASSO) archive, simulation code 527 528 0030 of 6 June 2015 (Gustafson et al. 2016a; Gustafson et al. 2016b). It has a 30-m vertical grid 529 spacing. LES2 was created by Glenn et al. (2019) using the System for Atmospheric Modeling 530 (SAM), with a vertical grid spacing 20 m above 600 m.

531 7 Outlook and Research Needs

532 The progress in understanding ABL transitions over land in the last several decades has laid a firm 533 foundation for future research. Much could be learned from better-resolved LES modeling through 534 transitions, including all relevant processes, such as interactive radiation and an interactive ground 535 surface. These processes are already represented in mesoscale models. Better understanding of 536 cloudy boundary layers could also come from LES, and from mining existing datasets for cases that 537 may have been discarded from previous analyses. Entrainment is still an open topic, especially its 538 variation with wind shear, aerosol, and surface inhomogeneity. Non-column effects, such as 539 terrain, coasts, and surface heterogeneity, must be considered as model grid sizes decrease. The 540 proper representation of transitioning boundary layers in the grey zone or terra icognita (Wyngaard 541 2004; Efstathiou et al. 2016) requires exploration. The full energy balance at the surface and at the 542 top of the boundary layer should be taken into account, including soil, canopy, and vegetation 543 contributions to heat and moisture. Evolving measurement techniques can also contribute. Because 544 real transitional boundary layers are influenced by even small surface or terrain variations, multi-545 scale networks of observations have the most promise to provide useful insight.

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550 **References**

- Acevedo OC, Fitzjarrald DR (2001) The early evening surface-layer transition: Temporal and
 spatial variability. J Atmos Sci. 58(17):2650-2667.
- Angevine WM (1999) Entrainment results including advection and case studies from the Flatland
 boundary layer experiments. J Geophys Res. 104(D26):30947-30963.
- Angevine WM (2008) Transitional, entraining, cloudy, and coastal boundary layers. Acta
 Geophys. 56:2-20 doi:10.2478/s11600-11007-10035-11601.
- Angevine WM, Bazile E, Legain D, Pino D (2014) Land surface spinup for episodic modeling.
 Atmos Chem Phys. 14(4):8165-8172.
- Angevine WM, Grimsdell AW, Hartten LM, Delany AC (1998) The Flatland boundary-layer
 experiments. Bull Am Meteorol Soc. 79:419-431.
- Angevine WM, Klein Baltink H, Bosveld FC (2001) Observations of the morning transition of the
 convective boundary layer. Boundary-Layer Meteorol. 101:209-227.

- Angevine WM, Olson J, Kenyon JS, Gustafson WI, Endo S, Sušelj K, Turner DD (2018) Shallow
 cumulus in WRF parameterizations evaluated against LASSO large-eddy simulations. Mon
 Weather Rev. 146:4303-4322.
- Angevine WM, White AB, Avery SK (1994) Boundary-layer depth and entrainment zone
 characterization with a boundary-layer profiler. Boundary-Layer Meteorol. 68:375-385.
- Baas P, Bosveld FC, Lenderink G, van Meijgaard E, Holtslag AAM (2010) How to design single column model experiments for comparison with observed nocturnal low-level jets. Q J R
 Meteorol Soc. 136(648):671-684.
- Barbaro E, de Arellano JV-G, Ouwersloot HG, Schröter JS, Donovan DP, Krol MC (2014)
 Aerosols in the convective boundary layer: Shortwave radiation effects on the coupled landatmosphere system. J Geophys Res Atmos. 119(10):5845-5863.
- Barlow JF, Dunbar TM, Nemitz EG, Wood CR, Gallagher MW, Davies F, O'Connor EJ, Harrison
 RM (2011) Boundary layer dynamics over London, UK, as observed using Doppler lidar
 during REPARTEE-II. Atmos Chem Phys. 11:2111-2125.
- 577 Barr AG, Betts AK (1997) Radiosonde boundary-layer budgets above a boreal forest. J Geophys
 578 Res. 102:29,205-229,212.
- Basu S, Vinuesa J-F, Swift A (2008) Dynamic LES modeling of a diurnal cycle. J Appl Meteorol
 Climatol. 47:1156-1174.
- Beare RJ (2008) The Role of Shear in the Morning Transition Boundary Layer. Boundary-Layer
 Meteorol. 129(3):395-410.
- Beare RJ, Edwards JM, Lapworth AJ (2006) Simulation of the observed evening transition and
 nocturnal boundary layers: Large-eddy simulation. Q J R Meteorol Soc. 132:81-99.
- Best MJ, Pryor M, Clark DB, Rooney GG, Essery RLH, Ménard CB, Edwards JM, Hendry MA,
 Porson A, Gedney N, Mercado LM, Sitch S, Blyth E, Boucher O, Cox PM, Grimmond
 CSB, Harding RJ (2011) The Joint UK Land Environment Simulator (JULES), model
 description Part 1: Energy and water fluxes. Geosci Model Dev. 4(3):677-699.
- 589 Betts AK, Barr AG (1996) First International Satellite Land Surface Climatology Field Experiment 590 1987 sonde budget revisited. J Geophys Res. 101(D18):23,285-223,288.
- Beyrich F, Goersdorf U (1995) Composing the diurnal cycle of mixing height from simultaneous
 sodar and wind profiler measurements. Boundary-Layer Meteorol. 76(999):387-394.
- Blay-Carreras E, Pardyjak ER, Pino D, Alexander DC, Lohou F, Lothon M (2014a)
 Countergradient heat flux observations during the evening transition period. Atmos Chem
 Phys. 14(17):9077-9085.
- Blay-Carreras E, Pardyjak ER, Pino D, Hoch SW, Cuxart J, Martínez D, Reuder J (2015) Lifted
 temperature minimum during the atmospheric evening transition. Atmos Chem Phys.
 15(12):6981-6991.
- Blay-Carreras E, Pino D, Vilà-Guerau de Arellano J, van de Boer A, De Coster O, Darbieu C,
 Hartogensis O, Lohou F, Lothon M, Pietersen H (2014b) Role of the residual layer and
 large-scale subsidence on the development and evolution of the convective boundary layer.
 Atmos Chem Phys. 14(9):4515-4530.
- Blumberg WG, Turner DD, Cavallo SM, Gao J, Basara J, Shapiro A (2019) An Analysis of the
 Processes Affecting Rapid Near-Surface Water Vapor Increases during the Afternoon to
 Evening Transition in Oklahoma. J Appl Meteorol Climatol. 58(10):2217-2234.
- Bony S, Stevens B (2018) Measuring Area-Averaged Vertical Motions with Dropsondes. J Atmos
 Sci. 76(3):767-783.
- Bosveld F, Baas P, Steeneveld G-J, Holtslag AM, Angevine W, Bazile E, de Bruijn EF, Deacu D,
 Edwards J, Ek M, Larson V, Pleim J, Raschendorfer M, Svensson G (2014) The Third

610	GABLS Intercomparison Case for Evaluation Studies of Boundary-Layer Models. Part B:
611	Results and Process Understanding. Boundary-Layer Meteorol. 152(2):157-187.
612	Boutle I, Price J, Kudzotsa I, Kokkola H, Romakkaniemi S (2018) Aerosol-fog interaction and the
613	transition to well-mixed radiation fog. Atmos Chem Phys. 18:7827-7840.
614	Brown AR, Cederwall RT, Chlond A, Duynkerke PG, Golaz J-C, Khairoutdinov M, Lewellen DC,
615	Lock AP, Macvean MK, Moeng C-H, Neggers RAJ, Siebesma AP, Stevens B (2002)
616	Large-eddy simulation of the diurnal cycle of shallow cumulus convection over land. Q J R
617	Meteorol Soc. 128:1075-1093.
618	Brown N, Weiland M, Hill A, Shipway B, Maynard C, Allen T, Rezny M (2015) A highly scalable
619	Met Office NERC cloud model. Proceedings of the 3rd International Conference on
620	Exascale Applications and Software, Edinburgh, UK2015, City, pp. 132-137.
621	Busse J, Knupp K (2012) Observed Characteristics of the Afternoon–Evening Boundary Layer
622	Transition Based on Sodar and Surface Data. J Appl Meteorol Climatol. 51(3):571-582.
623	Carter DA, Gage KS, Ecklund WL, Angevine WM, Johnston PE, Riddle AC, Wilson J, Williams
624	CR (1995) Developments in UHF lower tropospheric wind profiling at NOAA's Aeronomy
625	Laboratory. Radio Science. 30(4):977-1001.
626	Conzemius R. Fedorovich E (2006) Dynamics of sheared convective boundary layer entrainment.
627	Part II: Evaluation of bulk model predictions of entrainment flux. J Atmos Sci. 63:1179-
628	1199.
629	Conzemius RJ, Fedorovich E (2008) A case study of convective boundary layer development
630	during IHOP-2002: Numerical simulations compared to observations. Mon Weather Rev.
631	136:2305-2320.
632	Coulman CE (1978) Boundary-layer evolution and nocturnal inversion dispersal - part I.
633	Boundary-Layer Meteorol. 14:471-491.
634	Couvreux F, Bazile E, Canut G, Seity Y, Lothon M, Lohou F, Guichard F, Nilsson E (2016)
635	Boundary-layer turbulent processes and mesoscale variability represented by numerical
636	weather prediction models during the BLLAST campaign. Atmos Chem Phys. 16(14):8983-
637	9002.
638	Darbieu C, Lohou F, Lothon M, Vilà-Guerau de Arellano J, Couvreux F, Durand P, Pino D, Patton
639	EG, Nilsson E, Blay-Carreras E, Gioli B (2015) Turbulence vertical structure of the
640	boundary layer during the afternoon transition. Atmos Chem Phys. 15(17):10071-10086.
641	de Lozar A, Mellado JP (2015) Mixing driven by radiative and evaporative cooling at the
642	stratocumulus top. J Atmos Sci. 72:4681-4700.
643	Deardorff JW (1974) Three-dimensional numerical study of the height and mean structure of a
644	heated planetary boundary layer. Boundary-Layer Meteorol. 7(81-106).
645	Driedonks AGM (1982) Models and observations of the growth of the atmospheric boundary
646	layer. Boundary-Layer Meteorol. 23:283-306.
647	Edwards JM (2009) Radiative processes in the stable boundary layer: Part II. The development of
648	the nocturnal boundary layer. Boundary-Layer Meteorol. 131:127-146.
649	Edwards JM, Basu S, Bosveld FC, Holtslag AAM (2014) The impact of radiation on the GABLS3
650	large-eddy simulation through the night and during the morning transition. Boundary-Layer
651	Meteorol. 152:189-211.
652	Edwards JM, Beare RJ, Lapworth AJ (2006) Simulation of the observed evening transition and
653	nocturnal boundary layers: Single-column modelling. Q J R Meteorol Soc. 132:61-80.
654	Etstathiou GA, Beare RJ, Osborne S, Lock A (2016) Grey zone simulations of the morning
622	convective boundary layer development. J Geophys Res. 121:4/69-4/82.

- El Guernaoui O, Reuder J, Esau I, Wolf T, Maronga B (2019) Scaling the Decay of Turbulence
 Kinetic Energy in the Free-Convective Boundary Layer. Boundary-Layer Meteorol.
 173(1):79-97.
- Fedorovich E, Gibbs JA, Shapiro A (2017) Numerical study of nocturnal low-level jets over gently
 sloping terrain. J Atmos Sci. 74:2813-2834.
- 661 Garratt JR (1992) The atmospheric boundary layer. Cambridge University Press.
- 662 Garratt JR, Brost RA (1981) Radiative cooling rates within and above the nocturnal boundary
 663 layer. J Atmos Sci. 38:2730-2746.
- 664 Glenn IB, Feingold G, Gristey JJ, Yamaguchi T (2019) Quantification of the radiative effect of 665 aerosol-cloud-interactions in shallow continental cumulus clouds. J Atmos Sci. submitted.
- 666 Grant ALM (1997) An observational study of the evening transition boundary layer. Q J R
 667 Meteorol Soc. 123:657-677.
- 668 Grimsdell AW, Angevine WM (2002) Observations of the afternoon transition of the convective
 669 boundary layer. J Appl Meteorol. 41(1):3-11.
- Gustafson WI, Vogelmann AM, Cheng X, Endo S, Krishna B, Li Z, Toto T, Xiao H: 2016a,
 Description of the LASSO Alpha 1 Release, DOE Atmospheric Radiation Measurement
 Research Facility
- Gustafson WI, Vogelmann AM, Cheng X, Endo SK, B., Li Z, Toto T, Xiao H: 2016b, LASSO
 alpha 1 data bundles. 36° 36' 18.0" N, 97° 29' 6.0" W: Southern Great Plains Central
 Facility (C1), ARM Data Archive: Oak Ridge, Tennessee, USA
- Ha K-J, Mahrt L (2003) Radiative and turbulent fluxes in the nocturnal boundary layer. Tellus.
 55A:317-327.
- Harvey NJ, Hogan RJ, Dacre HF (2013) A method to diagnose boundary-layer type using Doppler
 lidar. Q J R Meteorol Soc. 139:1681–1693.
- Hogan RJ, Grant ALM, Illingworth AJ, Pearson GN, O'Connor EJ (2009) Vertical velocity
 variance and skewness in clear and cloud-topped boundary layers as revealed by Doppler
 lidar. Q J R Meteorol Soc. 135:635-643.
- Holtslag AAM, Svensson G, Baas P, Basu S, Beare RJ, Beljaars A, Bosveld FC, Cuxart J, Lindvall
 J, Steeneveld GJ, Tjernstroem M, van de Wiel BJH (2013) Stable atmospheric boundary
 layers and diurnal cycles: Challenges for weather and climate models. Bull Am Meteorol
 Soc. 94:1691-1706.
- Holtslag AAM, van Meijgaard E, De Rooy WC (1995) A comparison of boundary layer diffusion
 schemes in unstable conditions over land. Boundary-Layer Meteorol. 76:69-95.
- Jensen DD, Nadeau DF, Hoch SW, Pardyjak ER (2016) Observations of Near-Surface Heat-Flux
 and Temperature Profiles Through the Early Evening Transition over Contrasting Surfaces.
 Boundary-Layer Meteorol. 159(3):567-587.
- Kumar V, Svensson G, Holtslag AAM, Meneveau C, Parlange MB (2010) Impact of surface flux
 formulations and geostrophic forcing on large-eddy simulations of diurnal atmospheric
 boundary layer flow. J Appl Meteorol Climatol. 49:1496-1516.
- Lapworth A (2006) The Morning Transition of the Nocturnal Boundary Layer. Boundary-Layer
 Meteorol. 119(3):501-526.
- LeMone MA (1978) The Marine Boundary Layer, Workshop on the Planetary Boundary Layer.
 AMS, Boston.
- LeMone MA, Angevine WM, Bretherton CS, Chen F, Dudhia J, Fedorovich E, Katsaros KB,
- Lenschow DH, Mahrt L, Patton EG, Sun J, Tjernstroem M, Weil J (2019) 100 Years of
 Progress in Boundary-Layer Meteorology, AMS Centennial Monograph. AMS, Boston.

- LeMone MA, Grossman RL, McMillen RT, Liou K-N, Ou SC, McKeen S, Angevine W, Ikeda K,
 Chen F (2002) Cases-97: Late-morning warming and moistening of the convective
 boundary layer over the Walnut River watershed. Boundary-Layer Meteorol. 104:1-52.
- LeMone MA, Ikeda K, Grossman RL, Rotach MW (2003) Horizontal Variability of 2-m
 Temperature at Night during CASES-97. J Atmos Sci. 60(20):2431-2449.
- LeMone MA, Tewari M, Chen F, Dudhia J (2013) Objectively determined fair-weather CBL
 depths in the ARW-WRF model and their comparison to CASES-97 observations. Mon
 Weather Rev. 141:30-54.
- Lenschow DH, Stankov BB, Mahrt L (1979) The rapid morning boundary-layer transition. J
 Atmos Sci. 36:2108-2124.
- Liu C, Fedorovich E, Huang J (2017) Revisiting entrainment relationships for shear-free and
 sheared convective boundary layers through large-eddy simulations. Q J R Meteorol Soc.
 144:2182-2195.
- Liu C, Fedorovich E, Huang J, Hu X-M, Wang Y, Lee X (2019) Impact of aerosol shortwave
 radiative heating on entrainment in the atmospheric boundary layer: A large-eddy
 simulation study. J Atmos Sci. 76:785-799.
- Lothon M, Lenschow DH, Mayor SD (2009) Doppler Lidar Measurements of Vertical Velocity
 Spectra in the Convective Planetary Boundary Layer. Boundary-Layer Meteorol.
 132(2):205-226.
- Lothon M, Lohou F, Pino D, Couvreux F, Pardyjak ER, Reuder J, Vilà-Guerau de Arellano J,
 Durand P, Hartogensis O, Legain D, Augustin P, Gioli B, Lenschow DH, Faloona I, Yagüe
 C, Alexander DC, Angevine WM, Bargain E, Barrié J, Bazile E, Bezombes Y, BlayCarreras E, van de Boer A, Boichard JL, Bourdon A, Butet A, Campistron B, de Coster O,
 Cuxart J, Dabas A, Darbieu C, Deboudt K, Delbarre H, Derrien S, Flament P, Fourmentin
 M, Garai A, Gibert F, Graf A, Groebner J, Guichard F, Jiménez MA, Jonassen M, van den
- Kroonenberg A, Magliulo V, Martin S, Martinez D, Mastrorillo L, Moene AF, Molinos F,
 Moulin E, Pietersen HP, Piguet B, Pique E, Román-Cascón C, Rufin-Soler C, Saïd F,
 Sastre-Marugán M, Seity Y, Steeneveld GJ, Toscano P, Traullé O, Tzanos D, Wacker S,
- Wildmann N, Zaldei A (2014) The BLLAST field experiment: Boundary-Layer Late
 Afternoon and Sunset Turbulence. Atmos Chem Phys. 14(20):10931-10960.
- 732 Mahrt L (1981) The early evening boundary layer transition. Q J R Meteorol Soc. 107:329-343.
- 733 Mahrt L (2017) The near-surface evening transition. Q J R Meteorol Soc. 143:2940-2948.
- Manners J, Edwards JM, Hill P, Thelen JC: 2015, SOCRATES (Suite Of Community RAdiative
 Transfer codes based on Edwards and Slingo) Technical Guide, UK Met Office
- Manninen AJ, Marke T, Tuononen M, O'Connor EJ (2018) Atmospheric boundary layer
 classification with Doppler lidar. J Geophys Res. 123:8172-8189.
- Mellado JP, van Heerwarden CC, Garcia JR (2016) Near-surface effects of free atmosphere
 stratification in free convection. Boundary-Layer Meteorol. 159:69-95.
- Nadeau DF, Pardyjak E, Higgins CW, Fernando HJS, Parlange MB (2011) A simple model for the
 afternoon and early evening decay of convective turbulence over different land surfaces.
 Boundary-Layer Meteorol. 141:301-324.
- Nieuwstadt FTM, Brost RA (1986) The decay of convective turbulence. J Atmos Sci. 43(6):532546.
- Nilsson E, Lohou F, Lothon M, Pardyjak E, Mahrt L, Darbieu C (2016a) Turbulence kinetic
 energy budget during the afternoon transition Part 1: Observed surface TKE budget and
 boundary layer description for 10 intensive observation period days. Atmos Chem Phys.
 16(14):8849-8872.

749 Nilsson E, Lothon M, Lohou F, Pardyjak E, Hartogensis O, Darbieu C (2016b) Turbulence kinetic 750 energy budget during the afternoon transition – Part 2: A simple TKE model. Atmos Chem 751 Phys. 16(14):8873-8898. 752 Pennell WT, LeMone MA (1974) An experimental study of turbulence structure in the fair-753 weather trade wind boundary layer. J Atmos Sci. 31:1308-1323. 754 Pietersen HP, Vilà-Guerau de Arellano J, Augustin P, van de Boer A, de Coster O, Delbarre H, 755 Durand P, Fourmentin M, Gioli B, Hartogensis O, Lohou F, Lothon M, Ouwersloot HG, 756 Pino D, Reuder J (2015) Study of a prototypical convective boundary layer observed 757 during BLLAST: contributions by large-scale forcings. Atmos Chem Phys. 15(8):4241-758 4257. 759 Pino D, Jonker HJJ, vilà-Guerau de Arellano J, Dosio A (2006) Role of shear and the inversion 760 strength during sunset turbulence over land: characteristic length scales. Boundary-Layer 761 Meteorol. 121:537-556. Price J (2019) On the formation and development of radiation fog: An observational study. 762 763 Boundary-Layer Meteorol. 172:167-197. 764 Rizza U, Miglietta MM, Degrazia GA, Acevedo OC, Marques Filho EP (2013) Sunset decay of the 765 convective turbulence with large-eddy simulation under realistic conditions. Physica A. 766 392:4481-4490. 767 Román-Cascón C, Yagüe C, Mahrt L, Sastre M, Steeneveld GJ, Pardyjak E, van de Boer A, 768 Hartogensis O (2015) Interactions among drainage flows, gravity waves and turbulence: a 769 BLLAST case study. Atmos Chem Phys. 15(15):9031-9047. 770 Savijarvi H (2006) Radiative and turbulent heating rates in the clear-air boundary layer. Q J R 771 Meteorol Soc. 132:147-161. 772 Sharma V, Parlange MB, Calaf M (2017) Perturbations to the spatial and temporal characteristics 773 of the diurnally-varying atmospheric boundary layer due to an extensive wind farm. 774 Boundary-Layer Meteorol. 162:255-282. 775 Smith EN, Gebauer JG, Klein PM, Fedorovich E, Gibbs JA (2019) The Great Plains low-level jet 776 during PECAN: Observed and simulated characteristics. Mon Weather Rev. 147:1845-777 1869. 778 Sorbjan Z (1996) Effects caused by varying the strength of the capping inversion based on a large 779 eddy simulation model of the shear-free convective boundary layer. J Atmos Sci. 53:2015-780 2024. 781 Sorbjan Z (1997) Decay of convective turbulence revisited. Boundary-Layer Meteorol. 82:501-782 515. 783 Steeneveld GJ, Wokke MJJ, Groot Zwaaftink CD, Pijlman S, Heusinkveld BG, Jacobs AFG, 784 Holtslag AAM (2010) Observations of the radiation divergence in the surface layer and its 785 implication for its parameterization in numerical weather prediction models. J Geophys 786 Res. 115:D06107. 787 Stevens B (2002) Entrainment in stratocumulus-topped mixed layers. Q J R Meteorol Soc. 788 128:2663-2690. 789 Stull RB (1988) An introduction to boundary layer meteorology. Kluwer Academic. 790 Svensson G, Holtslag AAM, Kumar V, Mauritsen T, Steeneveld GJ, Angevine WM, Bazile E, 791 Beljaars A, de Bruijn EIF, Cheng A, Conangla L, Cuxart J, Ek M, Falk MJ, Freedman F, 792 Kitagawa H, Larson VE, Lock A, Mailhot J, Masson V, Park S, Pleim J, Söderberg S, 793 Weng W, Zampieri M (2011) Evaluation of the diurnal cycle in the atmospheric boundary 794 layer over land as represented by a variety of single column models - the second GABLS 795 experiment. Boundary-Laver Meteorol. 140:177-206.

- Tennekes H (1973) A model for the dynamics of the inversion above a convective boundary layer.
 J Atmos Sci. 30:558-567.
- Tucker SC, Senff CJ, Weickmann A, Brewer WA, Banta RM, Sandberg SJ, Law D, Hardesty RM
 (2009) Doppler lidar estimation of mixing height using turbulence, shear, and aerosol
 profiles. J Atmos Ocean Technol. 26:673-688.
- Van Driel R, Jonker HJJ (2011) Convective boundary layers driven by nonstationary surface heat
 fluxes. J Atmos Sci. 68:727-738.
- van Hooijdonk IGS, Clercx JH, Abraham C, Holdsworth AM, Monahan AH, Vignon E, Moene
 AF, Baas P, van de Wiel BJH (2017) Near-surface temperature inversion growth rate
 during the onset of the stable boundary layer. J Atmos Sci. 74:3433-3449.
- van Ulden AP, Wieringa J (1996) Atmospheric boundary layer research at Cabauw. Boundary Layer Meteorol. 78:39-69.
- Wilczak JM, Gossard EE, Neff WD, Eberhard WL (1996) Ground-based remote sensing of the
 atmospheric boundary layer: 25 years of progress. Boundary-Layer Meteorol. 78:321-349.
- Wildmann N, Rau GA, Bange J (2015) Observations of the Early Morning Boundary-Layer
 Transition with Small Remotely-Piloted Aircraft. Boundary-Layer Meteorol. 157(3):345-373.
- Wingo SM, Knupp KR (2015) Multi-platform Observations Characterizing the Afternoon-toEvening Transition of the Planetary Boundary Layer in Northern Alabama, USA.
 Boundary-Layer Meteorol. 155(1):29-53.
- 816 Wyngaard JC (1983) Lectures on the planetary boundary layer. In: Lilly DK and Gal-Chen T 817 (eds.), Mesoscale Meteorology - Theories, Observations, and Models. D. Reidel, Dordrecht.
- 818 Wyngaard JC (2004) Toward numerical modeling in the "Terra Incognita". J Atmos Sci. 61:1816-
- 819

1826.