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Low-power High Gain Observers for Wake Flow Rebuild

Javeria Ahmed, Matthieu Fruchard, Estelle Courtial and Youssoufi Touré

Abstract—This paper deals with the observation of the wake flow behind a bluff body based on reduced order models. Two Galerkin models with different truncation orders are considered to capture the intermittent character of the vortex shedding. The system output is the oscillation amplitude obtained thanks to pression measurements on the sides and on the back of the bluff body. For flow control purposes, the state estimation has to be efficient and robust to noises and dynamic changes. A high gain observer is modified i) by extending the state with the introduction of fictitious outputs in order to avoid an observability singularity and to ensure the existence of the inverse jacobian matrix and ii) by integrating a high gain parameter whose power is limited to 2 instead of n to address the sensitivity issue of the standard high gain observer. Simulation results illustrate the efficiency of the proposed observers and highlight promising applications for flow control at high Reynolds numbers.

I. INTRODUCTION

A. Context and motivation

The flow at the rear of a bluff body is a common occurrence such as the flows at the back of an airplane, a submarine or a vehicle. The turbulent wakes downstream the bluff body can be very complex, exhibiting coherent structures with different scales. The strongly intermittent behavior of the flow depends on the Reynolds number. At high Reynolds numbers, vortex shedding occurs in the wake, leading to a significant pressure drop on the rear surface of the body. This phenomenon gives rise to structural vibrations, acoustic noise and increases the drag force. The control of vortex shedding is then of major interest for engineering applications. Active control in fluid mechanics is a promising and challenging research area. Promising because lower energy consumption and better performances are expected by decreasing the drag force. Challenging because observation and control require both real-time tractable and robust models with respect to dynamics changes in the wake flow [1], [2]. Fluid flows are governed by the Navier-Stokes equations known for their complexity (strong nonlinearities, high dimensionality and time-delays). This kind of modeling can not be used when real-time applications are targeted [3]. Linear models are unable to recover unstabilities characterizing the wake flow [4], [5]. Within this context, reduced order models approaches have been widely used to address the extraction of the coherent structures dynamics in fluid flows: Proper Orthogonal Decomposition [6] and more recently Dynamic

Mode Decomposition [7], [8] using projection of the Navier-Stokes equations onto a modal basis. Galerkin models have proven to be a good compromise between simplicity and robustness. They are based on a Karhunen-Loève expansion around the steady or time-averaged Navier-Stokes solution u_s . The unsteady flow $u(x,t)$ is then approximated by an orthonormal Galerkin series:

$$u(x,t) \approx u_s(x) + \sum_{i=1}^{2N} a_i(t)u_i(x), \quad (1)$$

where $u_i(x)$ denote the orthonormal spatial modes and $a_i(t)$ are the temporal modes amplitudes. For time periodic flows, structures are described by pairs of modes, acting as modal oscillators. Since the first two modes ($N = 1$) convey 95% of the turbulent kinetic energy, is it necessary to truncate the equation (1) at higher order? However, one mode plays a crucial role to robustify the low-order Galerkin model with respect to Reynolds range and transient dynamics [9], [10]: it's the shift mode, noted a_Δ , modeling the energy flow between the mean flow and the vortex shedding. In the sequel, two Galerkin models are considered: a minimal Galerkin representation [9] using the dominant von Karman modes a_1, a_2 and the shift mode a_Δ (the decay rate $\mu > 0$ is supposed to be constant):

$$\begin{pmatrix} \dot{a}_1 \\ \dot{a}_2 \\ \dot{a}_\Delta \\ \dot{\mu} \end{pmatrix} = \begin{pmatrix} (\mu - a_\Delta)a_1 - a_2 \\ (\mu - a_\Delta)a_2 + a_1 \\ a_1^2 + a_2^2 - a_\Delta \\ 0 \end{pmatrix} \quad (2)$$

and a more refined Galerkin representation including higher damped ($\mu' > 0$) modes (a_3, a_4) interacting with the first two ones through the shift mode:

$$\begin{pmatrix} \dot{a}_1 \\ \dot{a}_2 \\ \dot{a}_3 \\ \dot{a}_4 \\ \dot{a}_\Delta \\ \dot{\mu} \end{pmatrix} = \begin{pmatrix} (\mu - a_\Delta)a_1 - a_2 \\ (\mu - a_\Delta)a_2 + a_1 \\ (-\mu' - a_\Delta)a_3 - a_4 \\ (-\mu' - a_\Delta)a_4 + a_3 \\ a_1^2 + a_2^2 + a_3^2 + a_4^2 - a_\Delta \\ 0 \end{pmatrix} \quad (3)$$

The latter model, referred as a hybrid model [9], [11] between the empirical Galerkin model (2) and higher dimensional invariant-manifold model, is expected to be more precise. Considering the oscillation amplitudes $A_i = a_{2i-1}^2 + a_{2i}^2$, the above models become two systems of order 3 and 4 respectively. With the notation $x_{min} = (A_1 \ a_\Delta \ \mu) \in \mathcal{X} \subset \mathbb{R}^3$, system (2) becomes:

$$\mathcal{S}_{min} : \dot{x}_{min} = \begin{pmatrix} 2x_1(x_3 - x_2) \\ x_1 - x_2 \\ 0 \end{pmatrix} = f(x_{min}) \quad (4)$$

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and using the state vector $x_{hyb} = (A_1 \ A_2 \ a_\Delta \ \mu) \in \mathcal{X} \subset \mathbb{R}^4$, system (3) becomes

$$\mathcal{S}_{hyb} : \dot{x}_{hyb} = \begin{pmatrix} 2x_1(x_4 - x_3) \\ 2x_2(-\mu' - x_3) \\ x_1 + x_2 - x_3 \\ 0 \end{pmatrix} = f(x_{hyb}) \quad (5)$$

For the sake of simplicity, the state vector x will be used for both x_{min} and x_{hyb} in the sequel. The only state component accessible to measurement is the oscillation amplitude A_1 . So the output is $y = h(x) = x_1$ in both cases.

B. Problem statement

Rebuilding the unmeasured states is mandatory to *i)* understand the macro-physics at work in the wake, i.e. the dominant structure dynamics, and *ii)* synthesize control laws aiming at stabilizing the wake around some desired trajectory. A first prerequisite is to study the conditions guaranteeing the observability of the system.

Proposition 1 *Systems \mathcal{S}_{min} and \mathcal{S}_{hyb} respectively given by (4) and (5) are observable on $\Omega_{min} = \{x \in \mathcal{X} \subset \mathbb{R}^3 : |x_1| > \varepsilon\}$ and $\Omega_{hyb} = \{x \in \mathcal{X} \subset \mathbb{R}^4 : |x_1| > \varepsilon, |x_2 + x_3 + \mu'| > \varepsilon'\}$ for any $\varepsilon, \varepsilon' > 0$, respectively.*

Proof: Let \mathcal{O} denote the smallest vector space containing h and closed under the Lie derivative \mathcal{L}_f . A system is said weakly observable if $\dim d\mathcal{O}(x) = \dim \mathcal{X}$ where $d\mathcal{O} = \text{Span}\{d\tau, \tau \in \mathcal{O}\}$ is the observability co-distribution.

Computing the differentials of the successive Lie derivatives of the output $y = h(x)$ along the system vector field f leads to the observability co-distributions $d\mathcal{O}_{min}$ and $d\mathcal{O}_{hyb}$ for systems \mathcal{S}_{min} and \mathcal{S}_{hyb} respectively, whose determinants are:

$$\begin{aligned} |d\mathcal{O}_{min}| &= -4x_1^2 \\ |d\mathcal{O}_{hyb}| &= -16x_1^3(x_2 + x_3 + \mu'). \end{aligned} \quad (6)$$

An observability singularity appears in particular when $x_1 = 0$, whereas the control aims at forcing the oscillation amplitude to zero. The designed observer should be able to handle this singularity in order to synthesize control laws based on the estimated states.

Among the observer approaches, the High Gain Observer (HGO) is well-adapted to estimate nonlinear system states and exploits a fast error convergence depending on the high gain parameter. However, it has some flaws which need to be addressed properly. The power of its high gain parameter increases up to the dimension of the system causing some sensitivity issues. Besides, a non-singular diffeomorphism between the system expressed in natural coordinates and the system in its normal form is required.

C. Contribution

The present paper proposes a solution for the aforementioned issues. The first issue is addressed by limiting the power of the high gain parameter. As suggested by [12] for high dimensional physical systems, the proposed observer limits the high gain parameter power to two instead of n at

the expense of the increased observer dimension from n to $2n-2$. In order to deal with the singularity issue that happens during the inversion of the jacobian matrix to get back the observer in original coordinates, the HGO is modified by extending the state with the introduction of fictitious outputs [13]. In this way, we avoid the singularity by making the system observable and ensuring the existence of the inverse jacobian matrix everywhere in space \mathbb{R}^n .

This paper is organized as follows: we briefly recall the design of high gain observers in section 2. Then the design of the proposed observer is detailed in section 3. Section 4 is devoted to the simulation results of the observation of the wake flow behind a bluff body. The observer efficiency for the two reduced order models are compared. Finally concluding remarks and on-going work are discussed in section 5.

II. RECALL ON HIGH GAIN OBSERVERS

High gain observers [14], [15] are widely used since their arbitrarily fast exponential decay and practical stability arbitrary bound, in case of unmodeled dynamics, make possible to synthesize real-time output feedback stabilizing control laws for systems with few state measurements. We briefly recall the standard high gain observer, the low-power high gain observer and the observer in natural coordinates.

Proposition 2 (High gain observer) *Consider the nonlinear system $\dot{x} = f(x), y = h(x)$ fulfilling the observability condition with $x \in \mathcal{X} \subset \mathbb{R}^n, y \in \mathbb{R}$. There exists an injective immersion ϕ from $\mathcal{X} \subset \mathbb{R}^n$ to \mathbb{R}^m with $m \geq n$ such that the nonlinear system can be rewritten in the observability normal form with $z = \phi(x)$ as:*

$$\dot{z} = A_m z + B_m F(z), y = C_m z = z_1 \quad (7)$$

where the matrices A_m, B_m and C_m are in prime form: A_m is a $(m \times m)$ square matrix whose non null entry is the unitary superdiagonal, $B_m(C_m^T)$ is a m -dimensional column vector whose non null entry is its unitary last (first) component.

Then a high gain observer is given by

$$\dot{\hat{z}} = A_m \hat{z} + B_m F_s(\hat{z}) + \Delta_m K(y - C_m \hat{z}) \quad (8)$$

with $\Delta_m = \text{diag}(\theta, \dots, \theta^m)$ where the high gain parameter $\theta > \theta_0 \geq 1$ is related to the Lipschitz constant of F_s , F_s is a saturated function of F and K is chosen such that the matrix $(A_m - KC_m)$ is Hurwitz.

Remark 1 *Two well-known limitations of the high gain observer are related to the increasing powers of the high gain θ and to the inversion of ϕ . First, as dimension increases, so does θ^m , and it may result in an innovation term that dominates the observer dynamics and an increased sensitivity to noise on the last states, especially at high frequencies. Second, since the observer is built in the z -space, recovering an estimation useful for control purposes often requires to go back to the natural coordinates x , meaning that ϕ has to be inverted. This inversion is in general not easy, especially if $m > n$.*

In a recent study [12], the authors proposed a new observer of dimension $(2m-2)$ using a high gain limited to power 2 instead of designing an observer of dimension m with a gain that grows up to power m . The limited gain to power 2 reduces the numerical implementation problems of the classic high gain observer and also reduces the sensitivity to noise measurements.

Proposition 3 (Low-power high gain observer) *Under assumptions and notations of Proposition 2, let*

$$\begin{aligned}\dot{\zeta}_i &= A_2\zeta_i + N\zeta_{i+1} + \Delta_2 K_i e_i, \quad i \geq 1 \\ &\vdots \\ \dot{\zeta}_{m-1} &= A_2\zeta_{m-1} + B_2 F_s(\zeta) + \Delta_2 K_{m-1} e_{m-1}\end{aligned}\quad (9)$$

with $\zeta_i \in \mathbb{R}^2$, $\zeta_{0,2} = y$ given by convention, $e_i = \zeta_{i-1,2} - \zeta_{i,1}$ and $K_i = (k_{i1} \quad k_{i2})^T$.

Let $N = (0_{2,1} \quad B_2)$, $E_i = (-K_i \quad C_2^T)$ and $Q_i = (0_{2,1} \quad K_i)$. Choosing the high gain $\theta > \theta_0$ as in Proposition 2 and gains K_i such that the matrix M

$$M = \begin{pmatrix} E_1 & N & 0 & \dots & \dots & 0 \\ Q_2 & E_2 & N & \ddots & & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & Q_i & E_i & N & \ddots \\ \vdots & & \ddots & \ddots & \ddots & \vdots \\ \vdots & & & \ddots & Q_{m-2} & E_{m-2} & N \\ 0 & \dots & \dots & 0 & Q_{m-1} & E_{m-1} \end{pmatrix}\quad (10)$$

is Hurwitz then

$$\dot{\zeta} = \Upsilon(\zeta, y)\quad (11)$$

with Υ given by (9) provides a low-power high gain observer for system (7). We can use either $\hat{z} = P\zeta$ or $\hat{z}' = P'\zeta$ with the block diagonal projectors P given by I_2 and $(m-1)$ matrices B_2^T , and P' given by $(m-1)$ matrices C_2 and I_2 .

Remark 2 *Despite this dimension increase, the low-power (Δ_2 depends only on θ, θ^2) high gain observer advantage is twofold: innovation term predominance is limited and sensitivity to output noise is consequently reduced, especially at high frequencies which is of particular interest considering wake flow dynamics.*

Both aforementioned observers make use of a generically nonlinear mapping $\phi : x \mapsto z$ whose inversion is required to infer the observed state \hat{x} in the original coordinates of the system dynamics. Such an inversion can be obtained through optimization but often at the price of costly computations. Furthermore, the non unicity of the inverse when $m > n$ may raise numerical issues. An alternative approach is to build the observer in the natural coordinates x as suggested initially by [16].

Proposition 4 (Observer in natural coordinates) *If the transformation $\phi : \mathcal{X} \subset \mathbb{R}^n \rightarrow \Pi \subset \mathbb{R}^q$ mapping the original system $\dot{x} = f(x)$ to a target system admitting an observer in the form $\dot{\pi} = \varphi(\pi, \hat{x}, y)$ is a diffeomorphism, then an observer in natural coordinates is given by:*

$$\dot{\hat{x}} = \frac{\partial \phi}{\partial x}(\hat{x})^{-1} \varphi(\phi(\hat{x}), \hat{x}, y).\quad (12)$$

Remark 3 *Obviously, building an observer in natural coordinates is more difficult when $q = m > n$ using a high gain observer [17], or when $q = 2(m-1)$ for low-power high gain observer, even if $m = n$.*

III. MAIN RESULTS

To address the observability singularity avoidance, we follow the core idea initiated in [18] and then enhanced and structured in [13]. We follow a three steps procedure. The first step consists in augmenting the initial n dimensional domain of the mapping ϕ to equal the dimension $q = 2(m-1)$ of the low-power observer. The resulting augmented mapping ϕ_a is now a diffeomorphism on a subset of \mathbb{R}^q . We then proceed to an extension of ϕ_a into an injective immersion ϕ_e on \mathbb{R}^q using fictitious outputs to artificially remove the observability singularity. The third step lies in the augmentation of ϕ_e into a global diffeomorphism ϕ where we augment the natural coordinates by the coordinate $\tau \subset \mathbb{R}^{N-n}$ so that $\phi(x, 0) = \phi_e(x)$.

A. A meaningful example

Let us illustrate the previous section results as well as the observability singularity issue on the simple empirical Galerkin system (4). First we define $\phi \in \mathcal{C}^1(\Omega_{min} \subset \mathcal{X}, \mathbb{R}^3)$ using

$$z = \phi(x) = \begin{pmatrix} x_1 \\ 2x_1(x_3 - x_2) \\ 2x_1(2(x_3 - x_2)^2 - x_1 + x_2) \end{pmatrix}\quad (13)$$

so ϕ maps (4) on the observability canonical form (7) with $m = 3$ and $F_3(x) = 2x_1[4(x_3 - x_2)^3 - 2(x_3 - x_2)(4x_1 - 3x_2) + x_1 - x_2]$. One can build an observer using Proposition 2 and get back to natural coordinates through the explicit expression of ϕ^{-1} :

$$\hat{x} = \phi^{-1}(\hat{z}) = \begin{pmatrix} \hat{z}_1 \\ \frac{\hat{z}_1 \hat{z}_3 - \hat{z}_2^2 + 2\hat{z}_1^3}{2\hat{z}_1^2} \\ \frac{\hat{z}_1 \hat{z}_2 + \hat{z}_1 \hat{z}_3 - \hat{z}_2^2 + 2\hat{z}_1^3}{2\hat{z}_1^2} \end{pmatrix}\quad (14)$$

which is undefined at the singularity $\hat{z}_1 = \hat{x}_1 = 0$. One can also get an observer in the natural coordinates using Propositions 2 and 4:

$$\dot{\hat{x}} = \frac{\partial \phi}{\partial x}(\hat{x})^{-1} (A_3 \phi(\hat{x}) + B_3 F_s(\hat{x}) + \Delta_3 K(y - C_3 \hat{x}))\quad (15)$$

as long as ϕ is a diffeomorphism. Yet we know from Proposition 1 that ϕ is singular at $\hat{x}_1 = 0$ and there is no guarantee that the observer will not leave Ω_{min} .

Now using Propositions 2 and 3, we get a low-power high gain observer

$$\dot{\zeta} = \varphi(\zeta, \hat{x}, y) = \begin{pmatrix} \zeta_{12} + \theta k_{11}(y - \zeta_{11}) \\ \zeta_{22} + \theta^2 k_{21}(y - \zeta_{11}) \\ \zeta_{22} + \theta k_{12}(\zeta_{12} - \zeta_{21}) \\ F_{3s}(\zeta) + \theta^2 k_{22}(\zeta_{12} - \zeta_{21}) \end{pmatrix} \quad (16)$$

and we can get back to the natural coordinates using either $\hat{z} = P\zeta$ or $\hat{z}' = P'\zeta$ combined with (14). Note that such an observer exhibits the same redundancy as in using the low-power high gain observer since we can use either $\hat{x} = P\hat{x}_a$ or $\hat{x}' = P'\hat{x}_a$. This observer also faces the problem linked with the observability singularity at $\hat{x}_1 = 0$.

B. Augmentation of the mapping ϕ

Since the low-power high gain observer underlies a mapping ϕ from $x \in \mathbb{R}^n$ to $\zeta \in \mathbb{R}^q$ with $q = 2(m-1) > m \geq n$, the mapping is not surjective for $m > 2$. A first step is to augment the mapping into a diffeomorphism by augmenting the mapping domain. Since the state of the observer of Proposition 3 is redundant in the form (using $\hat{z} = P\zeta$ and $\hat{z}' = P'\zeta$)

$$\zeta = \begin{pmatrix} \zeta_1 \\ \zeta_2 \\ \vdots \\ \zeta_{n-2} \\ \zeta_{n-1} \end{pmatrix} = \begin{pmatrix} \begin{pmatrix} \zeta_{11} \\ \zeta_{12} \\ \zeta_{21} \\ \zeta_{22} \end{pmatrix} \\ \vdots \\ \begin{pmatrix} \zeta_{n-21} \\ \zeta_{n-22} \\ \zeta_{n-11} \\ \zeta_{n-12} \end{pmatrix} \end{pmatrix} = \begin{pmatrix} \begin{pmatrix} \hat{z}_1 \\ \hat{z}_2 \\ \hat{z}_3 \\ \vdots \\ \hat{z}_{n-1} \\ \hat{z}_n \end{pmatrix} \\ \vdots \\ \begin{pmatrix} \hat{z}'_{n-2} \\ \hat{z}'_{n-1} \\ \hat{z}'_n \end{pmatrix} \end{pmatrix}, \quad (17)$$

a natural way of addressing a diffeomorphism augmentation is to merge the two states \hat{z} and \hat{z}' defining a new augmented state $\hat{z}_a = (\hat{z}_1 \ \hat{z}_2 \ \hat{z}'_2 \ \hat{z}_3 \ \dots \ \hat{z}'_{n-1} \ \hat{z}_n)^T$ associated with a natural coordinate augmented state $x_a = (x_1 \ x_2 \ x'_2 \ x_3 \ \dots \ x'_{n-1} \ x_n)^T \in \mathbb{R}^q$.

Proposition 5 For systems \mathcal{S}_{min} and \mathcal{S}_{hyb} , defining respectively the natural coordinate augmented state x_a and ϕ_a as

$$\phi_a : x_a = \begin{pmatrix} x_1 \\ x_2 \\ x'_2 \\ x_3 \end{pmatrix} \mapsto z_a = \begin{pmatrix} x_1 \\ 2x_1(x_3 - x_2) \\ 2x_1(x_3 - x'_2) \\ 2x_1(2(x_3 - x_2)^2 - x_1 + x_2) \end{pmatrix} \quad (18)$$

and $\phi_a : x_a = (x_1 \ x_2 \ x'_2 \ x_3 \ x'_3 \ x_4)^T \mapsto z_a$ given by

$$z_a = \begin{pmatrix} x_1 \\ 2x_1(x_4 - x_3) \\ 2x_1(x_4 - x'_3) \\ 4x_1(x_4 - x_3)^2 - 2x_1(x_1 + x_2 - x_3) \\ 4x_1(x_4 - x_3)^2 - 2x_1(x_1 + x'_2 - x_3) \\ 2x_1 \left(4(x_4 - x_3)^3 + (x_1 + x_2 - x_3) - 2x_2(-\mu' - x_3) \right) \\ -2(x_4 - x_3)(3(x_1 + x_2 - x_3) - x_1(x_4 - x_3)) \end{pmatrix} \quad (19)$$

ϕ_a is a diffeomorphism on $\Omega_{a,min} = \{x_a \in \mathbb{R}^4 : |x_1| > \varepsilon\}$, respectively on $\Omega_{a,hyb} = \{x_a \in \mathbb{R}^6 : |x_1| > \varepsilon, |x_2 + x_3 + \mu'| > \varepsilon'\}$.

Proof: To lighten notations, one can write

$$\frac{\partial \phi_a}{\partial x_a}(x_a) = \begin{pmatrix} 1 & 0_{1,q-1} \\ * & J \end{pmatrix}. \quad (20)$$

For system \mathcal{S}_{min} , differentiating ϕ_a leads to

$$J_{min} = 2 \begin{pmatrix} -x_1 & 0 & x_1 \\ 0 & -x_1 & x_1 \\ x_1 - 4x_1q_0 & 0 & 4x_1q_0 \end{pmatrix} \quad (21)$$

with $q_0 = x_3 - x_2$, so we have

$$\left| \frac{\partial \phi_a}{\partial x_a}(x_a) \right| = x_1^3$$

and ϕ_a is a diffeomorphism on $\Omega_{a,min}$.

Similarly for the system \mathcal{S}_{hyb} , differentiating ϕ_a leads to

$$J_{hyb} = 2 \begin{pmatrix} 0 & 0 & -x_1 & 0 & x_1 \\ 0 & 0 & 0 & -x_1 & x_1 \\ -x_1 & 0 & -4x_1q_1 + x_1 & 0 & 4x_1q_1 \\ 0 & -x_1 & -4x_1q_1 + x_1 & 0 & 4x_1q_1 \\ x_1q_2 & 0 & x_1q_3 & 0 & x_1q_4 \end{pmatrix} \quad (22)$$

with $q_1 = x_4 - x_3$, $q_2 = 1 - (-\mu' - x_3) - 6q_1$, $q_3 = -12q_1^2 + 6q_1 + 8x_2 + 8x_1 - 6x_3 - 1$ and $q_4 = 12q_1^2 - 2(4x_1 + 3x_2 - 3x_3)$. The determinant is given by

$$\left| \frac{\partial \phi_a}{\partial x_a}(x_a) \right| = x_1^5(x_2 + x_3 + \mu'),$$

so ϕ_a is a diffeomorphism on $\Omega_{a,hyb}$. \blacksquare

Invertibility of the jacobian matrix of ϕ_a is now possible as long as the estimated augmented state remains away from the observability singularity. Yet there is no guarantee that the estimated state stays in this set. The next step addresses this concern.

C. Extension of ϕ_a into an injective immersion on \mathbb{R}^q

To avoid losing any access to some unmeasured states because of the observability rank loss pinpointed in Proposition 1, for instance the loss of information about x_2 and x_3 when $x_1 = 0$ for system \mathcal{S}_{min} , a way to get around the observability singularity is to provide fictitious outputs aiming at getting information about these states around the singularity.

Proposition 6 For systems \mathcal{S}_{min} and \mathcal{S}_{hyb} , adding respectively fictitious outputs $y_f \in \mathbb{R}^p$

$$y_f(x) = \begin{pmatrix} \rho(x)x_2 \\ \rho(x)x'_2 \\ \rho(x)x_3 \end{pmatrix}, \quad y_f(x) = \begin{pmatrix} \rho(x)x_2 \\ \rho(x)x'_2 \\ \rho(x)x_3 \\ \rho(x)x'_3 \\ \rho(x)x_4 \\ \rho'(x) \end{pmatrix} \quad (23)$$

, where ρ and ρ' are differentiable functions that satisfy for all positive δ :

$$\begin{aligned} \rho(x_1) &: \quad \rho(x_1) + \delta x_1^2 > 0, \forall x_1 \in \mathbb{R} \\ \rho'(x_2, x_3) &: \quad \frac{\partial \rho'}{\partial x_2} + \frac{\partial \rho'}{\partial x_3} + x_2 + x_3 + \mu' > 0, \end{aligned} \quad (24)$$

to the state z_a results in an extended state $z_e = \begin{pmatrix} z_a \\ y_f \end{pmatrix}$. Extending the transformation ϕ_e mapping $x_a \in \mathbb{R}^q$ to $z_e \in \mathbb{R}^{q+p}$ results in an injective immersion on the domain \mathbb{R}^q .

Proof: The Jacobian matrix of ϕ_e is given by

$$\frac{\partial \phi_e}{\partial x_a}(x_a) = \begin{pmatrix} 1 & 0_{1,q-1} \\ * & J \\ * & J_f \end{pmatrix} \quad (25)$$

where J is J_{min} given by (21) (J_{hyb} given by (22)) for system \mathcal{S}_{min} (\mathcal{S}_{hyb}) respectively, and J_f is given by

$$J_{f,min} = \rho I_3, \quad J_{f,hyb} = \begin{pmatrix} \rho I_5 & & & & \\ \rho'_2 & 0 & \rho'_3 & 0 & 0 \end{pmatrix} \quad (26)$$

where $\rho'_i = \frac{\partial \rho'}{\partial x_i}$. Since by construction $2x_1^2 + \rho(x) > 0$ and $x_2 + x_3 + \mu' + \rho'_2(x) + \rho'_3(x) > 0$, it is straightforward that $\frac{\partial \phi_e}{\partial x_e}(x_e)$ has a constant rank q for any $x_a \in \mathbb{R}^q$. ■

D. Augmentation of ϕ_e into a global diffeomorphism

The observability singularity has been solved, but using Proposition 4 requires a diffeomorphism to get an observer in natural coordinates. This last step consists in augmenting the immersion ϕ_e defined on the full domain \mathbb{R}^q to \mathbb{R}^{q+p} into a global diffeomorphism ϕ on \mathbb{R}^{q+p} , which requires an augmentation of the natural coordinates $x_a \in \mathbb{R}^q$ to $\mathbf{x} = \begin{pmatrix} x_a \\ \tau \end{pmatrix} \in \mathbb{R}^{q+p}$. We follow here the Jacobian completion initially proposed in [18].

Proposition 7 Let $\tau \in \mathbb{R}^p$ and $\mathbf{x} = \begin{pmatrix} x_a \\ \tau \end{pmatrix} \in \mathbb{R}^{q+p}$. Defining

$$\phi : \mathbf{x} \mapsto \mathbf{z} = z_e + \Psi(x_a)\tau = \phi_e(x_a) + \Psi(x_a)\tau \quad (27)$$

with matrix $\Psi(x_a)$, the completion of the jacobian matrix $\frac{\partial \phi_e}{\partial x_a}$:

$$\frac{\partial \phi}{\partial \mathbf{x}} = \begin{pmatrix} \frac{\partial \phi_e}{\partial x_a} + \frac{\partial \Psi}{\partial x_a} & \Psi(x_a) \end{pmatrix} \quad (28)$$

is fulfilled for systems \mathcal{S}_{min} and \mathcal{S}_{hyb} using

$$\Psi_{min}(x_a) = \begin{pmatrix} 0_{1,3} \\ H_1 \\ -B_3^T \\ x_{1,3} \end{pmatrix} \text{ with } H_1 = \begin{pmatrix} I_2 & -1_{2,1} \end{pmatrix} \quad (29)$$

and

$$\Psi_{hyb}(x_a) = \begin{pmatrix} 0_{1,4} & 0 \\ H_2 & 0_{4,1} \\ -B_5^T & -2x_1 \\ x_{1,5} & 0_{5,1} \\ 0_{1,4} & 1 \end{pmatrix} \text{ with } H_2 = \begin{pmatrix} 0_2 & I_2 & -1_{2,1} \\ I_2 & 0_2 & -1_{2,1} \end{pmatrix}. \quad (30)$$

Then ϕ defines a global diffeomorphism.

Proof: Since Ψ depends only on x_1 , differentiating ϕ using (25) leads to

$$\frac{\partial \phi}{\partial \mathbf{x}}(\mathbf{x}) = \begin{pmatrix} 1 & 0_{1,q-1} & \Psi(x_1) \\ * & J & \\ * & J_f & \end{pmatrix} \quad (31)$$

Due to expressions (29) and (30), it follows that

$$\begin{aligned} \left| \frac{\partial \phi}{\partial \mathbf{x}}(\mathbf{x}) \right| &= (\rho + 2x_1^2)^3 \\ \left| \frac{\partial \phi}{\partial \mathbf{x}}(\mathbf{x}) \right| &= (\rho + 2x_1^2)^4 (2x_1^2(\rho'_2 + \rho'_3 + 2(x_2 + x_3 + \mu')) + \rho) \end{aligned} \quad (32)$$

for \mathcal{S}_{min} and \mathcal{S}_{hyb} , respectively. By construction of functions ρ and ρ' , these Jacobian determinants are strictly positive for all $\mathbf{x} \in \mathbb{R}^{q+p}$. ■

E. Low-power high gain observer in natural coordinates

Proposition 8 The dynamic system

$$\dot{\hat{\mathbf{x}}} = \left(\frac{\partial \phi}{\partial \mathbf{x}}(\hat{\mathbf{x}}) \right)^{-1} \begin{pmatrix} \Upsilon(\phi(\hat{\mathbf{x}}), y) \\ -K_p \phi_p(\hat{\mathbf{x}}) \end{pmatrix} \quad (33)$$

is a converging observer for system \mathcal{S}_{min} (\mathcal{S}_{hyb}) in their natural coordinates on any bounded subset of \mathbb{R}^{q+p} with $(q, p) = (4, 3)$ (with $(q, p) = (6, 6)$). Υ is defined by Proposition 3, ϕ is defined by Propositions 5, 6 and 7, K_p is a positive definite diagonal matrix of dimension p and the notation ϕ_p stands for the last p components of ϕ .

Proof: It can be checked that $\phi^{-1} \in \mathcal{C}^1$, hence it is L -Lipschitz on any compact of \mathbb{R}^{q+p} . It follows that, using $\mathbf{z}^* = \phi(x_a, 0) = \phi_e(x_a)$:

$$\begin{aligned} \|x_a(t) - \hat{x}_a(t)\| + \|\tau(t)\| &\leq \sqrt{2} \left\| \begin{pmatrix} x_a - \hat{x}_a \\ \tau \end{pmatrix} \right\| \\ &\leq \sqrt{2} \left\| \begin{pmatrix} x_a \\ 0 \end{pmatrix} - \begin{pmatrix} \hat{x}_a \\ \tau \end{pmatrix} \right\| \\ &\leq \sqrt{2} \|\phi^{-1}(\mathbf{z}^*) - \phi^{-1}(\hat{\mathbf{z}})\| \\ &\leq L\sqrt{2} \|\mathbf{z}^* - \hat{\mathbf{z}}\| \\ &\leq L\sqrt{2} \|\phi_e(x_a) - \phi(\hat{x}_a, \tau)\|. \end{aligned} \quad (34)$$

Since ϕ_e is injective and (9) defines a converging observer from Proposition 3, it follows that

$$\lim_{t \rightarrow \infty} \|x_a(t) - \hat{x}_a(t)\| + \|\hat{\tau}(t)\| = 0. \quad (35)$$

IV. SIMULATION RESULTS

x_0	$(0 \ 0.51 \ 0.8)^T$
\hat{x}_0	$(0.1 \ 0.1 \ 0.2)^T$
K_1, K_2	$\begin{pmatrix} 0.18 \\ 0.1 \end{pmatrix}, \begin{pmatrix} 0.18 \\ 0.02 \end{pmatrix}$
K_p	I_3
ε	0.5
θ	2

TABLE I

INITIAL CONDITIONS AND PARAMETERS FOR SYSTEM \mathcal{S}_{min}

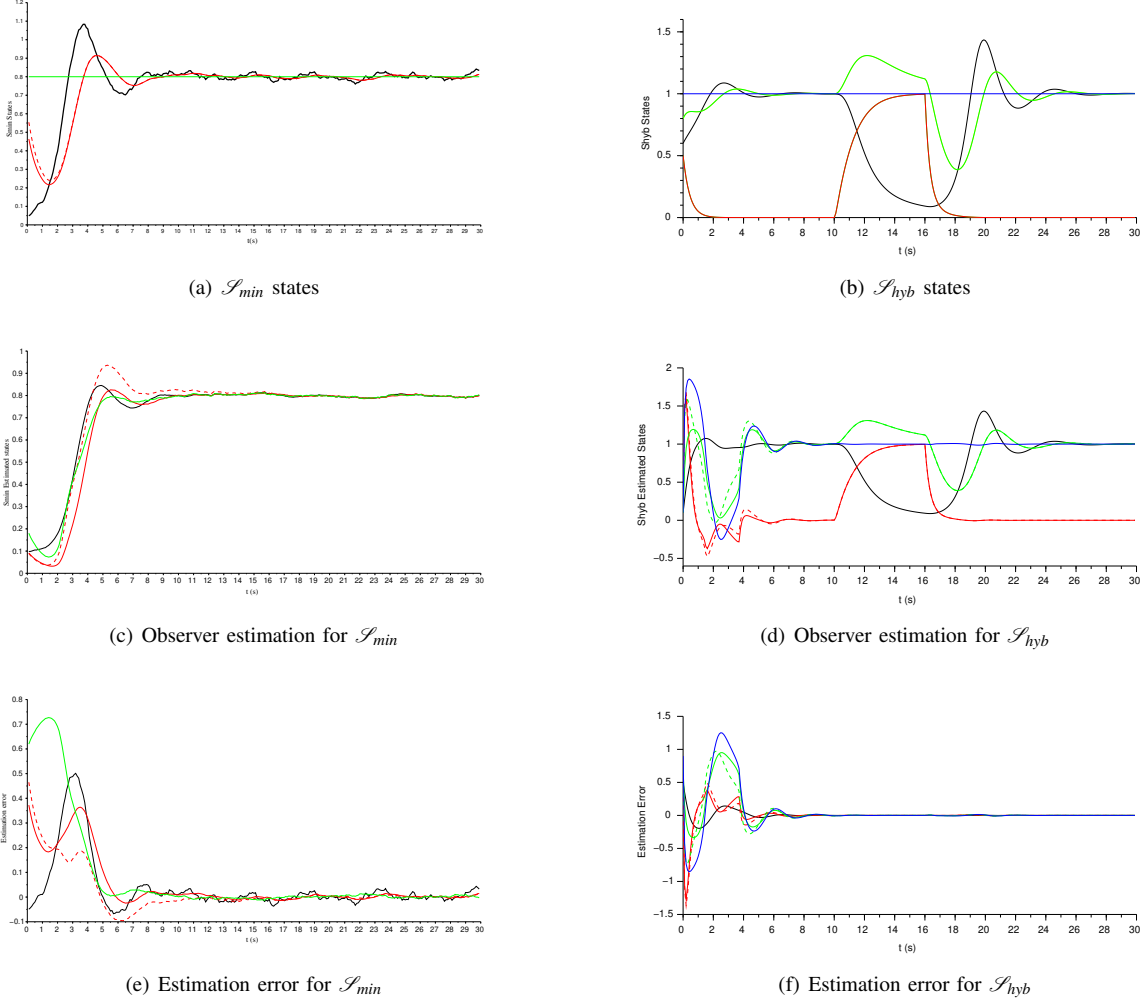


Fig. 1. Estimation using the low-power high gain observer in natural coordinates. From top to bottom, on the subfigures: (a)-(b) the states of the physical systems \mathcal{S}_{min} (\mathcal{S}_{hyb}): x_1, x_2, x_3 (and x_4) are depicted in black, red, green and blue solid lines; (c)-(d) state estimation $\hat{x}_a^T = (\hat{x}_1 \hat{x}_2 \hat{x}_2' \dots \hat{x}_4)$ is plotted with \hat{x}_1 in black solid line, (\hat{x}_2, \hat{x}_2') in red solid and dash dotted lines, (\hat{x}_3, \hat{x}_3') in green solid and dash dotted lines, and \hat{x}_4 in blue; (e)-(f) the estimation error is drawn using the previous color convention. Due to the low-power high gain observer redundancy, there are two estimates \hat{x}_i and \hat{x}_i' for $2 \leq i \leq n-1$.

Simulations are performed using the low-power high gain observer in natural coordinates given by Proposition 8. The first simulation is led on system \mathcal{S}_{min} and illustrates the robustness with respect to noise. The x_1 dynamics and the output y are corrupted by the addition of zero mean Gaussian white noises of standard deviations 0.2 and 0.1 respectively. The former noise is intended to model turbulence and impact of other unmodeled dynamics on the main mode x_1 . System \mathcal{S}_{min} given by (4) has been simulated with the initial conditions and observer parameters given in Table I. Fictitious outputs required to get a full-ranked Jacobian matrix in Proposition 6 are defined using

$$\rho : x \mapsto \max(0, \varepsilon^2 - x_1^2)^2. \quad (36)$$

Despite the initial condition $x_1(0) = 0$, the proposed observer is well defined and exhibits a good convergence property irrespective of the singularity. However, some peaking phenomenon can be noticed at the beginning of the simulation on Figure 1(e) because of the inappropriate

choice for initial conditions that "activates" the fictitious outputs. It implies changes in the dynamics of the extension τ , and in turn on the estimates since the observer aims at making $\phi(\hat{x}) = \phi_e(\hat{x}_a) + \Psi(\hat{x}_a)\tau$ converge, whilst $\phi_e(\hat{x}_a)$ is a better representation of the estimates. Then the error asymptotically converges to zero and the role of fictitious outputs vanishes outside the $\hat{x}_1 > |\varepsilon|$, hence displaying the effectiveness of the proposed observer, particularly with respect to noise filtering.

Simulation parameters for system \mathcal{S}_{hyb} described by (5) are given in Table II. As highlighted in Proposition 1, the system \mathcal{S}_{hyb} can face an additional observability singularity for $x_2 + x_3 + \mu' = 0$, so we use functions ρ defined by (36) and $\rho'(x) = (1/(x_2 + x_3 + \mu') - 1/\varepsilon')^2$ in order to maintain $\hat{x}_2 + \hat{x}_3 + \mu' > 0$. As it can be noticed on Figure 1(b), the initial dynamics of \mathcal{S}_{hyb} are driven by the growth rate parameter $x_4 = \mu$, here assumed to be a constant. \mathcal{S}_{hyb} behaves as a double oscillator whose main mode amplitude

x_0	$(0.6 \ 0.5 \ 0.8 \ 1)^T$
\hat{x}_0	$(0.1 \ 0.1 \ 0.1 \ 0.1)^T$
K_1, K_2, K_3	$\begin{pmatrix} 1.8 \\ 1 \end{pmatrix}, \begin{pmatrix} 2.4 \\ 0.4 \end{pmatrix}, \begin{pmatrix} 2 \\ 0.2 \end{pmatrix}$
K_p	0.1I ₆
θ	2
$\varepsilon, \varepsilon'$	0.1, 0.5

TABLE II

INITIAL CONDITIONS AND PARAMETERS FOR SYSTEM \mathcal{S}_{hyb}

x_1 asymptotically follows a transcritical bifurcation, whilst the enslaved mode amplitude x_2 is damped thanks to the term $(-\mu')$. During the time interval [10s; 16s], a control input is applied on the dynamics of the second mode as suggested in [11] in order to excite this oscillator. As expected the coupling –through the shift mode x_3 in the third equation of (5)– results in an energy flow from the main mode x_1 to the mode x_2 , forcing x_1 to go to zero. During the transient phase, Figure 1(d) illustrates the convergence of the state estimates to the real states despite initial values close to zero. As in the first simulation, avoiding the singularity induces some peaking: it is expected since the observer is then built on a system that does not represent exactly the physical model. The estimation tracking is ensured as x_1 goes to zero due to the control input.

V. CONCLUSIONS AND FUTURE WORKS

We have proposed an observer capable of being sufficiently robust with respect to noises and efficient to rebuild the system states while avoiding singularities as shown by the simulation results. The modified low-power high gain observer has provided a solution to many challenges : *i*) a rebuild of turbulent wake flow from an unique measured output based on reduced order models, *ii*) an observability singularity avoidance, *iii*) a robustness in regard to high frequency noises and sudden changes of the dynamics. It would be interesting to study if the redundancy induced by the low-power high gain observer may be exploited to reduce the proposed observer dimension. The on-going works are on one hand, the experimental validation of the observer in real conditions (the experimental setup is being fully equipped by sensors and actuators) to get a better knowledge of physical phenomena and, on the other hand, the design of feedback laws based on the estimated state to control the flow in order to reduce the drag force. Both perspectives are of great interest for applications in fluid mechanics.

VI. ACKNOWLEDGMENTS

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