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► **To cite this version:**

Benjamin Fildier, W D Collins, C Muller. Changing skewness of the rain distribution with warming, with and without self-aggregation. *Journal of Advances in Modeling Earth Systems*, In press. hal-03023335

**HAL Id: hal-03023335**

**<https://hal.science/hal-03023335>**

Submitted on 25 Nov 2020

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# 1 **Changing skewness of the rain distribution with warming, with** 2 **and without self-aggregation**

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## 7 **Key Points:**

- 8 • In aggregated RCE, the mean rain fall rate is larger, but its relative increase with  
9 warming is similar to that of disorganized RCE
- 10 • Rainrates are sensitive to the strength of aggregation and are enhanced when aggre-  
11 gation feedbacks are combined
- 12 • In the presence of aggregation, extreme rain can increase faster than Clausius-  
13 Clapeyron because of increasing precipitation efficiency

**Abstract**

We investigate how mesoscale circulations associated with convective aggregation can modulate the sensitivity of the hydrologic cycle to warming. We quantify changes in the full distribution of rain across radiative-convective equilibrium states in a cloud-resolving model. For a given SST, the shift in mean rainfall between disorganized and organized states is associated with a shift in atmospheric radiative cooling, and is analogous to the effect of a 4K SST increase. With rising temperatures, the increase in mean rain rate is insensitive to the presence of organization, while extremes intensify faster in the aggregated state, leading to a faster amplification in the skewness of rain.

Overall, heavy rain is enhanced by 20-30% when convection aggregates, and its sensitivity to warming shows an excess of 2.5%/K beyond the Clausius-Clapeyron scaling. However, nonlinear behaviors are observed under aggregation. First, radiative- and surface-flux aggregation feedbacks have multiplicative effects on extremes, illustrating a non-trivial sensitivity to the degree of organization. Second, alternating super- and sub-Clausius-Clapeyron regimes in extreme rainfall are found as a function of SST, corresponding to varying contributions in the thermodynamic, dynamic, and precipitation efficiency increases in different SST ranges.

The potential for mesoscale circulations in amplifying the hydrologic cycle is established. However these nonlinear distortions question the quantitative relevance of idealized self-aggregation. This calls for a systematic investigation of universal relationships which capture the coupling between global energetics, aggregation feedbacks and local convection, and which could hold across domain configurations, surface boundary conditions, microphysics and turbulence schemes.

**Plain Language Summary**

Convective aggregation, or organization, is known to affect the spatial distribution of clouds, the wind circulation and the intensity of rain as a result of feedbacks that couple convective processes, radiative transfer in the atmosphere and energy fluxes from the Earth's surface. We investigate how the hydrologic cycle responds to warming in various conditions of forcing and aggregation feedbacks in a hierarchy of idealized simulations, and provide a fine characterization of the statistical distribution of rain in order to connect its modes of change to the physical drivers involved in aggregation. The complex behav-

45   ior of the rain distribution in these simulations feeds a discussion on the use of idealized  
46   experiments to investigate convective organization and on their relevance to understand  
47   future changes in the hydrologic cycle.

## 1 Introduction

The strength of the global hydrologic cycle can be compactly summarized by the global statistical distribution of rain. The properties of this distribution emerge from the interplay between a variety of atmospheric processes, from the large-scale energy budget and atmospheric general circulation to short convective processes. Its skewness, or “unevenness”, is related to the relative importance between dry areas, drizzle and more active regimes of convection [Pendergrass and Knutti, 2018]. As the climate warms, the breadth of the distribution is stretched further by the tiered increases in mean and extreme precipitation [Pendergrass and Hartmann, 2014a], driven respectively by large-scale changes in atmospheric radiative cooling and the shorter-scale response of convection to atmospheric moistening (and in particular the Clausius-Clapeyron formula) [Allen and Ingram, 2002]. Emulating the relevant atmospheric radiative and convective processes with simple scaling approximations has enabled linking specific processes to the statistical properties of rain and their change with warming [O’Gorman and Schneider, 2009; Pfahl et al., 2017; Fildier et al., 2018], identifying sources of uncertainty in climate models [O’Gorman and Schneider, 2009; Fildier and Collins, 2015] and comparing the behavior of different modeling strategies [Fildier, 2019]. In particular, one common deficiency to all modeling frameworks is the scale separation between the large-scale circulation and small-scale convective processes due to limitations in computing capabilities. The resulting inability of models to represent the full continuum of scales responsible for the spatial organization of convection raises the central question motivating this work: as the climate warms, how could changes in the mesoscale circulations involved in the spatial structure of rain affect its global statistics?

Warmer air holds more water vapor, and that tends to cause an increase in extreme precipitation intensities at the 6-7%/K rate of the Clausius-Clapeyron (CC) formula [O’Gorman and Muller, 2010]. Departures from this thermodynamic reference has been noted when convection is parameterized in climate models [Pall et al., 2006; O’Gorman and Schneider, 2009; Fildier et al., 2017], as well as on short time scales for regional extremes with respect to local temperatures [Lenderink et al., 2017; Loriaux et al., 2013]. However, in cases of radiative-convective equilibrium (RCE, when atmospheric radiative cooling is fully balanced by latent convective heating in the absence of lateral inflow or outflow), the CC-scaling seems to hold consistently at the convective scale with respect to global temperature [Muller et al., 2011; Romps, 2011; Fildier et al., 2017; Muller and Takayabu,

81 2020]. In the latter case, mesoscale circulations are unresolved, since the simulated con-  
82 vection is either in a disorganized state (random, or pop-up convection) or superparame-  
83 terized [*Khairoutdinov and Randall, 2001; Grabowski, 2001*]. Hence the question of the  
84 effects of mesoscale circulations on extreme rainfall in a changing climate remains open.

85 The approach chosen here is to use a cloud-resolving model (CRM) in a mesoscale  
86 RCE domain that exhibits spatial organization in the form of convective self-aggregation  
87 (see, e.g., *Wing et al. [2017]* for a review). This idealized setup is a specific case of organ-  
88 ization, since convective organization can also encompass more general features observed  
89 such as mesoscale convective systems [*Maddox, 1980*] or patterns of shallow convection  
90 [*Feingold et al., 2010; Bony et al., 2020*]. More generally, organization can also occur in  
91 GCMs [*Coppin and Bony, 2015*], with or without convective parameterization, and under  
92 various conditions of forcing [*Wing et al., 2017*]. Under homogeneous boundary condi-  
93 tions, convection can also organize in squall lines due to imposed wind shear [*Muller,*  
94 *2013*]. Although changing the degree of organization (squall lines in that study) can lead  
95 to up to a doubling of extreme rainfall rates, for a given degree of organization the inten-  
96 sification of precipitation extremes with warming remains similar to the Clausius-Clapeyron  
97 theoretical expectation, close to about 7%/K with warming. Here instead, we will focus on  
98 the amplification of precipitation extremes when convection spontaneously self-aggregates,  
99 without any large-scale forcing or shear. In this modeling setup, a mesoscale circulation  
100 spontaneously develops within the RCE domain, and this circulation is explicitly coupled  
101 to the resolved convective processes.

102 Here, the internal circulations of interest emerge spontaneously in response to in-  
103 ternal feedbacks which drive and maintain the system into a lower energy state [*Emanuel*  
104 *et al., 2014*]. Self-aggregation feedbacks involve, in particular, the spatial heterogeneities  
105 in longwave radiative fluxes [*Muller and Held, 2012*] and wind-induced evaporation [*Brether-*  
106 *ton et al., 2005*]. In this case, it takes the form of a moist patch where convection is ac-  
107 tive, surrounded by a much drier region which lets the system lose more energy radia-  
108 tively to space. This organized RCE state corresponds to different mean climate properties  
109 and climate sensitivity than the disorganized RCE state [*Mauritsen and Stevens, 2015*], as  
110 well as heterogeneous thermodynamic properties that could have substantial effects on the  
111 characteristics of the hydrologic cycle [*Tobin et al., 2013; Tan et al., 2015*].

112 The potential role of self-aggregation on precipitation extremes in the framework  
113 of RCE has been highlighted by several studies. At the coarse resolutions of GCMs, *Pen-*  
114 *dergrass et al.* [2016] show that precipitation extremes could increase faster with warming  
115 than the CC rate because of a changing degree of organization, and because aggregation  
116 may be more likely to occur above a critical SST threshold [*Held et al.*, 1993; *Emanuel*  
117 *et al.*, 2014; *Wing and Emanuel*, 2014]. *Bao et al.* [2017] further confirmed this diagnostic  
118 by showing that the ratio of extreme-rainfall-increase to mean-rainfall-increases is greater  
119 for larger degrees of organization. However, using an idealized CRM and fixed SST, *Bao*  
120 *and Sherwood* [2018] noticed that the statistical distribution of instantaneous precipitation  
121 does not change with the degree of aggregation because increases in precipitation effi-  
122 ciency are compensated by a reduced updraft speed of condensing parcels. This behavior  
123 appears independent from the microphysics scheme chosen in their study. These apparent  
124 contradictions raise the question: can self-aggregation lead to stronger mean and extreme  
125 precipitation, and can it amplify their increase with global warming?

126 We perform a series of CRM simulations in radiative-convective equilibrium (RCE)  
127 with fixed SSTs between 300 K and 308 K. Following most earlier studies of self-aggregation,  
128 for simplicity we neglect the Earth's rotation. This is a reasonable approximation for the  
129 study of deep tropical clouds, as the Coriolis parameter is small at low latitudes. The ra-  
130 diative and surface-flux feedbacks are alternately turned on or off to constrain the system  
131 in organized and disorganized states, to vary the strength of aggregation, and to remove  
132 some methodological differences for more adequate comparison with previous studies. The  
133 general procedure and simulations are introduced in section 2. Section 3 quantifies the  
134 separate roles of convective organization and surface temperatures on the distribution of  
135 rain and describes the different behavior of organized and disorganized precipitation with  
136 warming. Section 4 investigates which mechanisms involved in self-aggregation affect the  
137 strength of mean and extreme rain, in particular the circulations induced and reinforced  
138 by the radiative and surface-flux feedbacks. Section 5 investigates the thermodynamic and  
139 dynamic response of extreme events themselves in order to explain their sensitivity to the  
140 circulation. Because the relevance and strengthening of convective organization in future  
141 climates is an active area of research, we discuss the sensitivity of these results to the sim-  
142 ulation design, and argue for future investigations of how the coupling between global cli-  
143 mate and local convection can be modulated by these mesoscale circulations (section 6).

## 144 2 Methodology

### 145 2.1 Numerical experiments

146 We perform a series of experiments using the System for Atmospheric Modeling  
 147 (SAM) version 6.10.10 [Khairoutdinov and Randall, 2003]. This anelastic CRM uses as  
 148 prognostic variables liquid and ice static energy and non-precipitating and precipitating  
 149 water. Outputs are saved on hourly averages. All runs analyzed use a 4km grid resolu-  
 150 tion and a 1024-km square domain with doubly-periodic horizontal boundary conditions.  
 151 The vertical grid has 64 levels with a resolution of 500m at the tropopause and 50m at  
 152 the surface; in the upper levels a sponge layer is added to prevent gravity wave reflection.  
 153 The following parameterizations are used: a 1.5-order subgrid-scale TKE closure for tur-  
 154 bulent processes, the native 1-moment scheme for microphysics, and the CAM3 radiation  
 155 scheme [Collins *et al.*, 2006]. Surface fluxes are calculated from Monin-Obukhov similar-  
 156 ity theory. The model is forced with a fixed solar constant of 650 W/m<sup>2</sup> at a 50.0° angle,  
 157 typical mean insolation of equatorial regions, and with uniform SSTs of 300, 302, 304 and  
 158 306 K. It is run to RCE before analysis.

159 The experiments performed in this paper are summarized in table 1 at each SST.  
 160 The reference simulations are the first two lines. The organized run  $O_{\text{ref}}$  has interactive  
 161 surface fluxes and interactive, locally-computed radiation to allow spontaneous aggrega-  
 162 tion due to the surface and radiative feedbacks. In the reference disorganized run  $D_{\text{ref}}$ , the  
 163 radiative feedback is removed by homogenizing the radiative heating rates  $Q$  at each time  
 164 step in the horizontal dimensions similarly to Muller and Held [2012], while the surface  
 165 fluxes are kept interactive to allow the simulated atmosphere to reach mass and energy  
 166 equilibrium. Both  $O_{\text{ref}}$  and  $D_{\text{ref}}$  are initialized from a disorganized RCE state obtained on  
 167 a smaller square domain 128-km wide, used for spin up at each SST. Because the equi-  
 168 libration time scale of atmospheric humidity is about 40 days for disorganized runs and  
 169 80 days for organized simulations,  $D_{\text{ref}}$  is run for 100 days and  $O_{\text{ref}}$  for 150 days, and the  
 170 last 50 days of each run is used for analysis. Comparing  $O_{\text{ref}}$  and  $D_{\text{ref}}$  for different SSTs  
 171 allows us to quantify the overall effect of organization on extreme rainfall.



172

**Table 1.** Simulations run performed for SST  $\in \{300, 302, 304, 306, 308\}$  K.

| Name  | Radiative heating rates $Q$   | Surface fluxes $\mathcal{F}$               | Expected RCE state | Duration |
|---|---|--|--------------------|----------|
| $\mathcal{O}_{\text{ref}}$                            | interactive, heterogenous   | interactive, heterogenous                  | organized          | 150 days |
| $\mathcal{D}_{\text{ref}}$                            | interactive, homogenized  | interactive, heterogenous                  | disorganized       | 100 days |
| $\mathcal{O}(\mathcal{F}_{\mathcal{O}_{\text{ref}}})$ | interactive, heterogenous   | prescribed from $\mathcal{O}_{\text{ref}}$ | organized          | 100 days |
| $\mathcal{O}(\mathcal{F}_{\mathcal{D}_{\text{ref}}})$ | interactive, heterogenous   | prescribed from $\mathcal{D}_{\text{ref}}$ | organized          | 100 days |
| $\mathcal{D}(\mathcal{Q}_{\mathcal{O}_{\text{ref}}})$ | radiative heating profile<br>prescribed from $\mathcal{O}_{\text{ref}}$ | interactive, heterogenous                  | disorganized       | 100 days |

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## 2.2 Separate effects of aggregation: circulation-reinforcing feedbacks and mean climate shift

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The additional experiments shown on Table 1 are designed to separate the distinct roles that organization can have on the strength of mean and extreme rainfall intensities: “circulation effects”, associated with the horizontal heterogeneities maintained by the radiative feedback and further strengthened by the surface flux feedback, and “mean-climate” effects, associated with shifts in the domain-averaged atmospheric radiative cooling and surface fluxes. Indeed, convective organization is associated with the development of a large-scale circulation, with low-level divergence from dry regions and low-level convergence into moist regions [Muller and Held, 2012]. This large-scale circulation is believed to be driven by differential radiative heating rates between moist and dry regions [Muller and Bony, 2015]. Convective organization is also accompanied by enhanced domain-mean outgoing longwave radiative cooling to space [Wing and Emanuel, 2014], as the free troposphere dries in the subsiding environment. This entails a stronger domain-averaged radiative cooling and by energy conservation a larger domain-averaged surface enthalpy flux.

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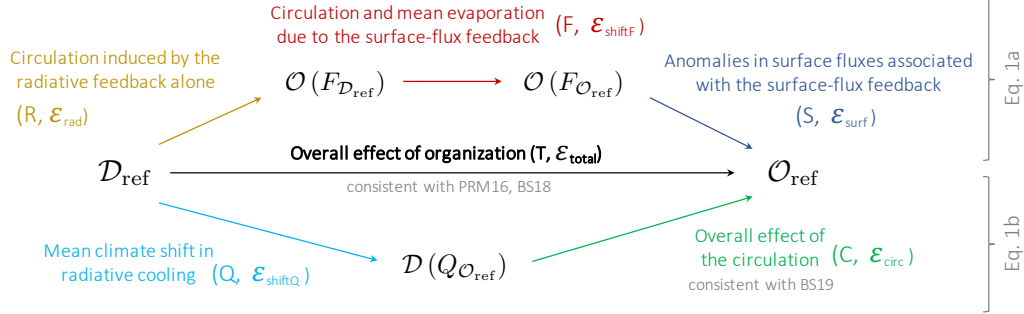
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Overall, this distinction between heterogeneities and domain-averages is artificial because these elements likely interact in a nonlinear fashion. But this exercise will help to emphasize the role of the spatial heterogeneities on the strength of extremes, as well as to highlight that changes in mean climate state induced by organization mainly affect the domain-mean rainfall rather than the extremes. Figure 1 explains which pairing of exper-

194 iments can be used to estimate the role of these individual processes on the intensity and  
 195 change in extreme rainfall. These three simulations are initialized from the end state of  
 196  $O_{\text{ref}}$  and run for 100 days to achieve a robust steady state.



197 **Figure 1.** Summary of the experiments and their relation to individual mechanisms. The same color coding  
 198 is used in later figures. Arrows are labeled with  $\mathcal{E}$  symbols for defining enhancement factors (equation 1) and  
 199 letters for plotting them (Figure 7).

200 The two pathways drawn on Figure 1 can be decomposed into a product of enhance-  
 201 ment factors  $\mathcal{E}$  for precipitation, as follows:

$$202 \quad P(O_{\text{ref}}) = \underbrace{\frac{P(O_{\text{ref}})}{P(O(F_{O_{\text{ref}}}))}}_{\mathcal{E}_{\text{surf}}} \times \underbrace{\frac{P(O(F_{O_{\text{ref}}}))}{P(O(F_{D_{\text{ref}}}))}}_{\mathcal{E}_{\text{shift}F}} \times \underbrace{\frac{P(O(F_{D_{\text{ref}}}))}{P(D_{\text{ref}})}}_{\mathcal{E}_{\text{rad}}} \times P(D_{\text{ref}}) \quad (1a)$$

$$203 \quad \text{and } P(O_{\text{ref}}) = \underbrace{\frac{P(O_{\text{ref}})}{P(D(Q_{O_{\text{ref}}}))}}_{\mathcal{E}_{\text{circ}}} \times \underbrace{\frac{P(D(Q_{O_{\text{ref}}}))}{P(D_{\text{ref}})}}_{\mathcal{E}_{\text{shift}Q}} \times P(D_{\text{ref}}) \quad (1b)$$

204 where  $P(X)$  is the precipitation statistic of interest (the mean, or an extreme percentile)  
 205 for simulation  $X$ . Simulations  $O(F_{O_{\text{ref}}})$  and  $O(F_{D_{\text{ref}}})$  use fixed sensible and latent heat  
 206 fluxes, but active radiation to allow organization to persist. These prescribed surface fluxes  
 207 are diagnosed from the end states of  $O_{\text{ref}}$  and  $D_{\text{ref}}$  respectively, as averages in space across  
 208 the domain and in time over the last 50 days of simulation.  
 209

210 The comparison proceeds as follows. The organized  $O(F_{D_{\text{ref}}})$  and the disorganized  
 211  $D_{\text{ref}}$  have the same mean surface fluxes, which emphasizes the role of the circulation in-  
 212 duced by the radiative feedback alone on the intensity of extremes. The only difference  
 213 between  $O_{\text{ref}}$  and  $O(F_{O_{\text{ref}}})$  are the surface-flux spatial structure, so that their differences  
 214 highlights how the circulation is reinforced by the spatial structure of the surface-flux  
 215 feedback. Differences between  $O(F_{O_{\text{ref}}})$  and  $O(F_{D_{\text{ref}}})$  represents to first order the mean

216 shift in the atmospheric energy budget between the disorganized and organized states,  
 217 which is apparent in the surface fluxes themselves in steady state.

218 The last simulation  $\mathcal{D}(Q_{O_{\text{ref}}})$  uses a prescribed radiative cooling profile chosen as  
 219 the domain mean profile at equilibrium in the  $O_{\text{ref}}$  simulation, and interactive surface  
 220 fluxes to let the system free to reach RCE. This leads to a disorganized state constrained  
 221 in a similar manner as *Bao and Sherwood* [2018] and facilitates comparison with their re-  
 222 sults. By comparing  $\mathcal{D}(Q_{O_{\text{ref}}})$  with  $\mathcal{D}_{\text{ref}}$ , it also allows to isolate the shift in mean atmo-  
 223 spheric radiative cooling that is induced by organization: it represents the effect of mean  
 224 climate shift on precipitation intensities when convection remains disorganized, that is,  
 225 without any adjustment in the circulation. Instead, differences between  $\mathcal{D}(Q_{O_{\text{ref}}})$  and  $O_{\text{ref}}$   
 226 represent the full effect of the circulation independently from changes in the mean climate.

227 The relationships between these simulations and their physical interpretation intro-  
 228 duced above are summarized in Figure 2. The main effect of organization (central arrow)  
 229 will be described in section 3 and the upper and lower pathways shown in the figure will  
 230 be further analyzed and commented in section 4.

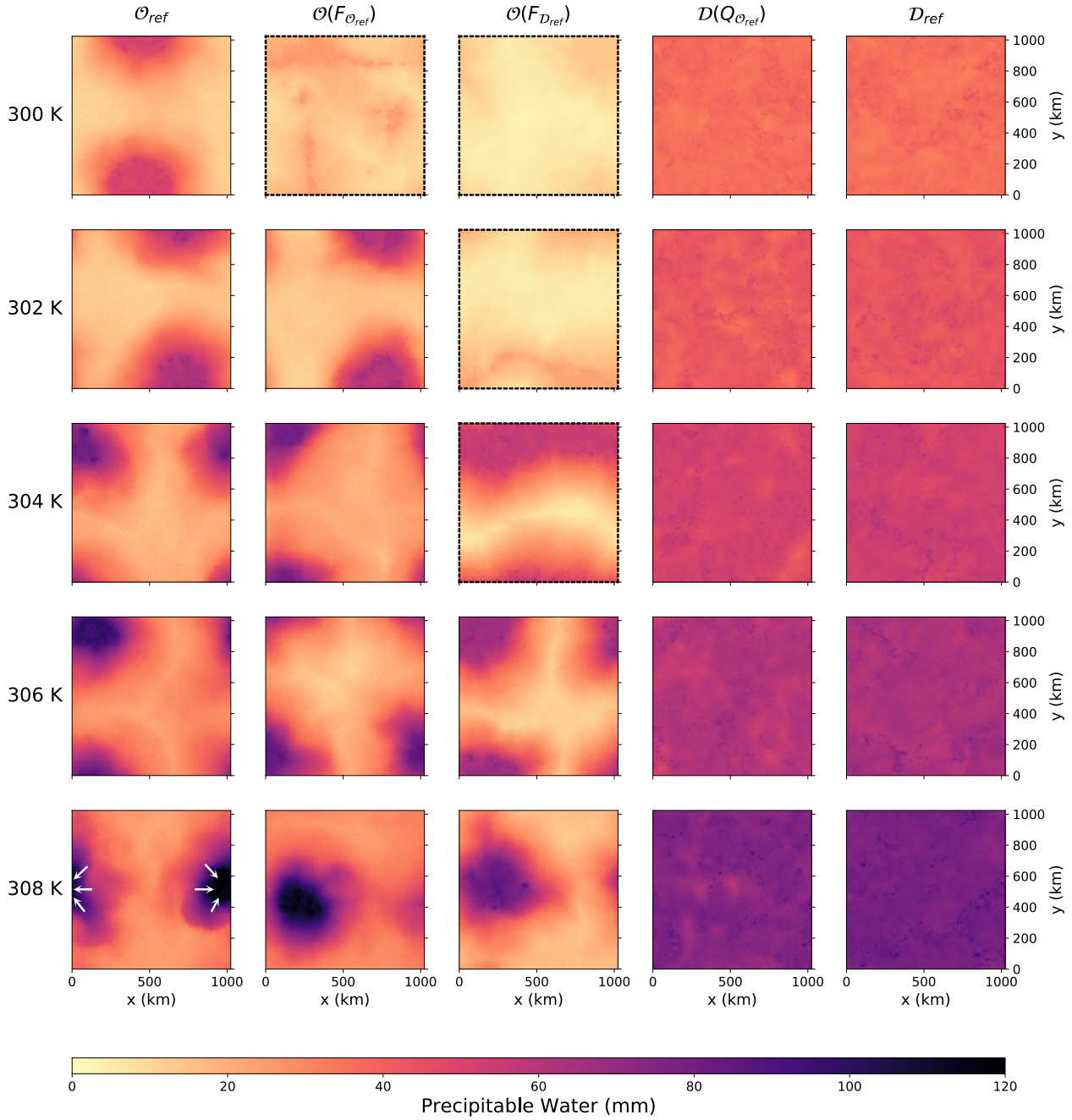
### 231 **2.3 Methodological limitations**

232 Figure 2 shows 2D snapshots of precipitable water  $PW$  at the end of the simulation  
 233 for all simulation types and all SSTs. In some cases with fixed surface fluxes at the lowest  
 234 SSTs, represented with a dashed frame on Figure 2, the system cannot maintain its aggre-  
 235 gated states. These runs display oscillations in the spatial pattern of convection, where the  
 236 convectively-active regions alternate between a small circular shapes and elongated stripes,  
 237 and this behavior gradually leads to a strong drying of the entire domain (not shown). Be-  
 238 cause of the lack of robustness of the organized state, these runs will not be analyzed in  
 239 this article.

## 244 **3 Organized precipitation is heavier and intensifies faster with warming**

### 245 **3.1 Acceleration of the hydrologic cycle amplified by organization**

246 We first investigate differences between the organized and disorganized simulations  
 247 about precipitation statistics and their change with warming, using the 300K to 304K  
 248 range as a reference (Figure 3).



240 **Figure 2.** Snapshots of precipitable water for all SSTs (rows) and all simulation types (columns) at the end  
 241 of the simulation. Simulation details are provided in Table 1. Arrows are hand-drawn on one of the simula-  
 242 tions to indicate the direction of the low-level circulation. Simulations shown in dashed frames are discarded  
 243 from the analysis because their final states do not stay robustly aggregated.

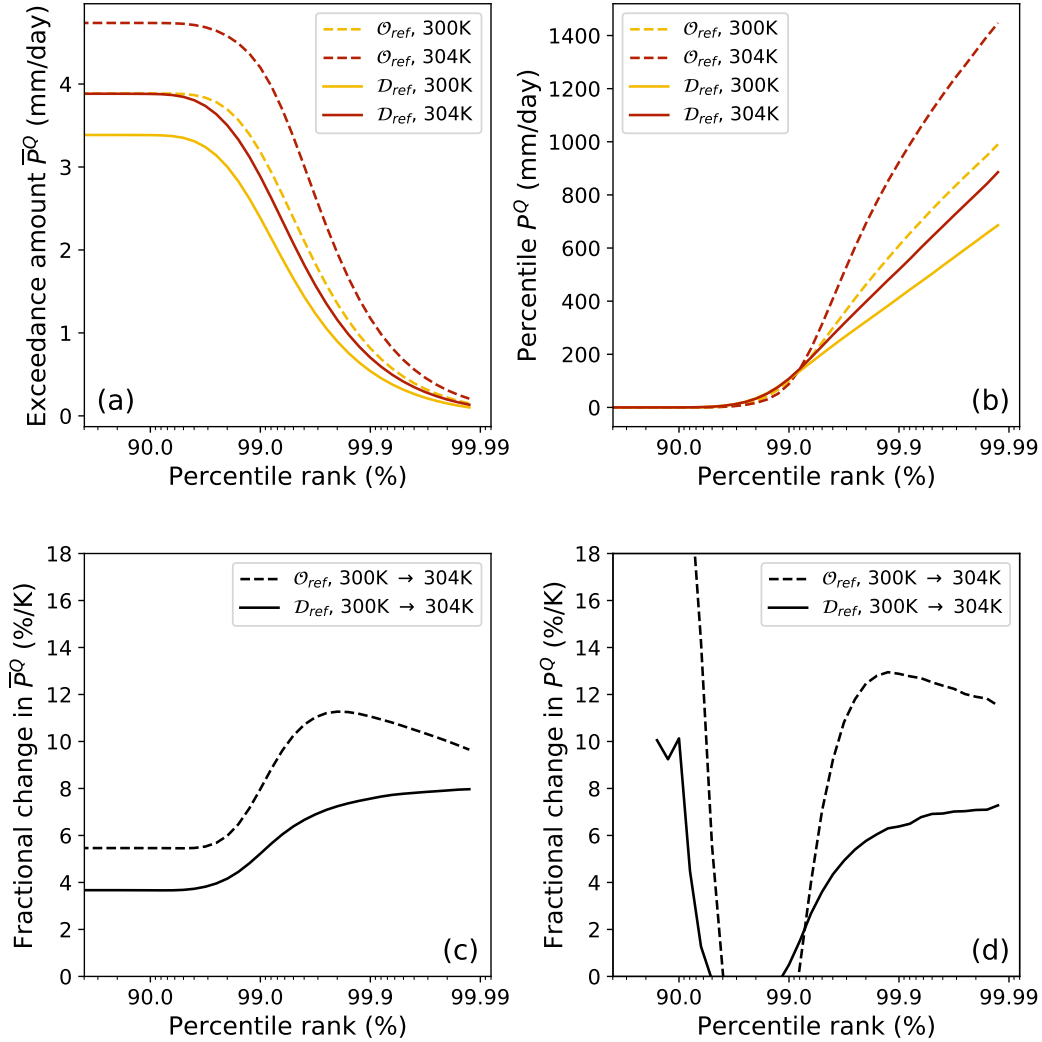
### 3.1.1 Simultaneous characterization of mean and extreme rainfall

Characteristics of the hydrologic cycle are quantified by calculating the full distribution of rainfall intensities and rainfall amounts. Rainfall intensities, or percentiles, are expressed in  $\text{kg/m}^2/\text{s}$  or  $\text{mm/day}$  and are calculated on hourly time scales at each percentile of the distribution of rain, including wet and dry points. Percentiles  $P^Q$  are chosen on an inverse-logarithmic scale ( $\dots, 90^{\text{th}}, 99^{\text{th}}, 99.9^{\text{th}}, \dots$ ) with 10 bins regularly spaced per decade, to zoom on the extremes. This metric is traditionally used as a way to quantify extreme precipitation [Pall *et al.*, 2006; Muller and O’Gorman, 2011; Fildier *et al.*, 2018] that facilitates interpreting rain intensities in terms of return times or frequencies of occurrence ( $\dots, 1/10, 1/100, 1/1000, \dots$ ). The second metric is rainfall exceedance amount  $\bar{P}^Q$ , also used by Pendergrass and Knutti [2018]: it has the same units as  $P^Q$  and is calculated by summing the rainfall intensities above each percentile rank  $Q$ :

$$\bar{P}^Q \equiv \int_Q^1 P^{Q'} f(Q') dQ' \quad (2)$$

where  $f(Q)$  is the frequency of points in the  $Q^{\text{th}}$  percentile bin. This integrated approach enables a more comprehensive characterization of the hydrologic cycle by quantifying mean and extreme rainfall simultaneously. At the lowest percentiles,  $\bar{P}^0$  corresponds to the domain mean rainfall and is controlled by the domain mean atmospheric energy budget, while at the largest percentiles with  $Q \rightarrow 1$ ,  $\bar{P}^Q$  is the mean water falling inside the most intense events that are controlled by convective processes. As a result,  $\bar{P}^Q/\bar{P}^0$  represents the fraction of global rain which falls in the form of extremes.

The distributions of exceedance amounts  $\bar{P}^Q$  (Figure 3a) reveal several general features of the changes in the hydrologic cycle induced by convective organization and SSTs. Two modes of change in the distribution of rain can be observed on these curves (with a similar physical interpretation as Pendergrass and Hartmann [2014a], although the statistical representation differs). First, a horizontal shift of  $\bar{P}^Q$  towards the right (higher percentile rank  $Q$ ) shows that rain becomes more extreme, meaning that given amounts of rain fall within fewer precipitation events. Second, a vertical shift of  $\bar{P}^Q$  at fixed  $Q$  is characteristic of the increase in domain-mean rainfall that is consistent with changes in the atmospheric energy budget [Pendergrass and Hartmann, 2014a; Fildier and Collins, 2015]. As shown on Figure 3a, these two shifts can either occur from changes in SSTs or in the degree of convective organization. When fractional changes in  $\bar{P}^Q$  are calculated for all percentiles (Figure 3c), it also becomes apparent that convective organization could be si-



269 **Figure 3.** Comparison of statistical distribution of precipitation for the organized (solid) and disorganized  
 270 (dashed) experiments at 300K (yellow) and 304K (red). The statistics are shown on individual climates (upper  
 271 row) as well as their fractional changes with warming (lower row), for percentiles precipitation intensity  $P^Q$   
 272 (left panels) and exceedance amounts  $\bar{P}^Q$  (right panels).

285 multaneously associated with larger rates of increase in domain mean rainfall (left end of  
 286 the curve) and extreme rainfall (right end) as surface temperatures rise.

287 Focusing on the response of extreme rain intensities, it appears  $P^Q$  can be amplified  
 288 by convective organization or SST (Figure 3b) and that the aggregation-driven enhance-  
 289 ment of extremes is exacerbated with warming (Figure 3d). Note that the magnitudes of  
 290  $P^Q$  are large because they are computed on scales of 1h and 4km; 1000 mm/day approx-  
 291 imately corresponds to 40 mm/hr, which is realistic for hourly extremes. Consistent with

292 the CC scaling, extreme disorganized precipitation intensities increase at about 6-7%/K,  
 293 whereas organized extremes increase faster with warming, at 12.5%/K for this specific pair  
 294 of SSTs (Figure 3d).

### 3.1.2 Apparent link between the skewness of the distribution and the degree of ag- 295 gregation 296

297 The skewness of the distribution of rain, shown by the sharpness of the curves on  
 298 Figure 3a, expresses the unevenness of precipitation [Pendergrass and Knutti, 2018]. In  
 299 our simulations, it appears that this unevenness is amplified by convective organization,  
 300 regardless of its effect on the mean climate.

301 It so happens that domain-mean rainfall rates are identical to within 0.1% in the  
 302  $O_{\text{ref}}^{300K}$  and  $D_{\text{ref}}^{304K}$  simulations (see Figure 3a) as a result of similar radiative cooling and  
 303 mean surface fluxes in these two simulations. This implies that at fixed SST, the reorga-  
 304 nization of convection induces a mean climate shift equivalent to a 4K warming for disor-  
 305 ganized rainfall. This equivalence does not hold for extreme precipitation, as can be seen  
 306 by comparing the distributions of rain intensities and exceedance amounts between  $O_{\text{ref}}^{300K}$   
 307 and  $D_{\text{ref}}^{304K}$ . So, for equivalent mean climates, an organized state corresponds to a stronger  
 308 skewness of the distribution of rain towards more intense extreme events. Later sections  
 309 will confirm that this skewness is likely associated with differences in the circulation and  
 310 in thermodynamic heterogeneities.

311 Following this reasoning, the enhancement of extremes by self-aggregation circula-  
 312 tions alone can be estimated as:

$$313 \frac{P^{99.9}(O_{\text{ref}}^{300}) - P^{99.9}(D_{\text{ref}}^{300})}{P^{99.9}(D_{\text{ref}}^{300})} = 46.7\% > 27.6\% = \frac{P^{99.9}(D_{\text{ref}}^{304}) - P^{99.9}(D_{\text{ref}}^{300})}{P^{99.9}(D_{\text{ref}}^{300})} \quad (3)$$

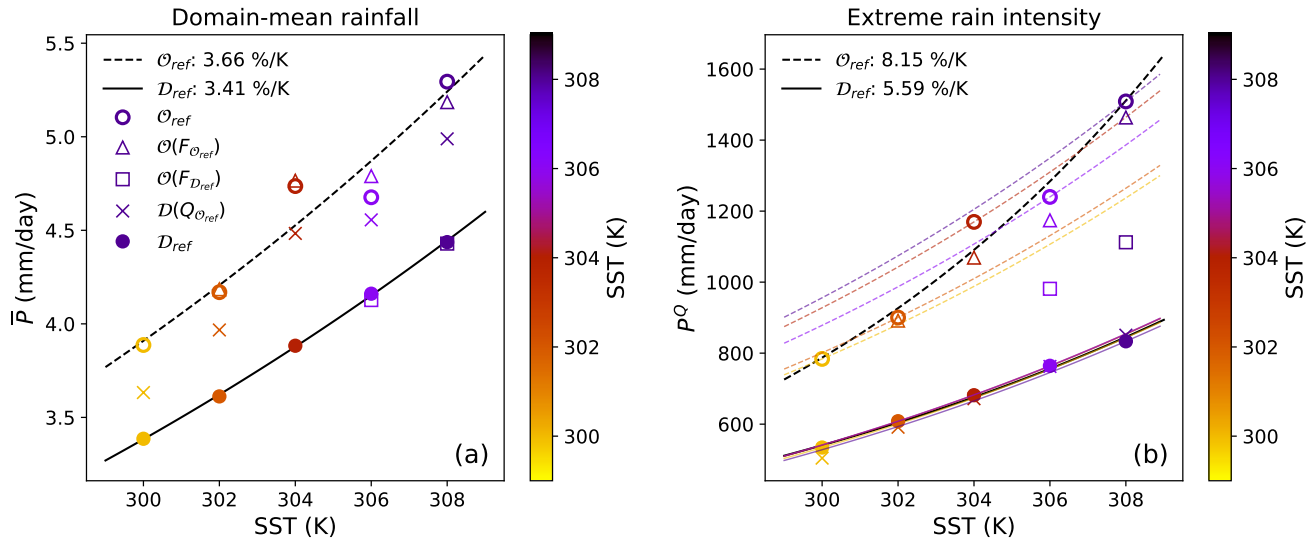
314 This “4K-warming equivalence” also applies for higher SSTs:  $O_{\text{ref}}^{302}$  and  $O_{\text{ref}}^{304}$  have similar  
 315 domain mean rainfall as  $D_{\text{ref}}^{306}$  and  $D_{\text{ref}}^{308}$  respectively. As a consequence, the same exercise  
 316 can be done to quantify how changes in the circulation alone affect the fractional increase  
 317 in extreme rainfall. Assuming that simulations  $O_{\text{ref}}^T$  and  $D_{\text{ref}}^{T+4K}$  have similar mean cli-  
 318 mates, the amplification of the fractional increase in extremes by the circulation alone is

$$319 \delta P^{99.9}(O_{\text{ref}})|_{300 \rightarrow 302} - \delta P^{99.9}(D_{\text{ref}})|_{304 \rightarrow 306} = 1.35\%/K \quad \text{at } 300 \text{ K.} \quad (4)$$

320 More robust estimates of the amplification of extremes are calculated next (sec-  
 321 tion 3.2), followed by a more detailed attribution of this amplification to effects to aggre-  
 322 gation feedbacks (sections 4 and 5).

### 3.2 Trends in mean and extreme rain, and alternating regimes

324 Mean and extreme rainfall intensities are shown on Figure 4 across the complete  
 325 range of SSTs. Each dot shows the value of mean and extreme rain for each simulation  
 326 analyzed at each SST and the black curves represent exponential fits to the  $O_{ref}$  and  $D_{ref}$   
 327 simulations. The large offset between the two exponential fits on both graphs shows that  
 328 convective organization induces a clear amplification of mean and extreme rainfall in fixed  
 329 climates. This systematic effect from organization on mean rain was not always found in  
 330 previous studies [*Craig and Mack, 2013*]. Mean rainfall is enhanced by a fixed fraction  
 331 ( $\approx 15\%$ ) which leads to global-warming trends of similar magnitude for  $O_{ref}$  (3.7%) and  
 332  $D_{ref}$  (3.4%). Section 4 further attributes this mean rainfall enhancement to the mean cli-  
 333 mate shift induced by aggregation.



334 **Figure 4.** Strength of mean precipitation (left) and precipitation extremes  $P^{99.9}$  (right) as a function of SST  
 335 for all simulation types. Four simulations are excluded, as explained in section 2.2 and displayed in Figure 2.  
 336 Curves show the expected Clausius-Clapeyron increase of precipitation extremes that corresponds to the  
 337 reference  $D_{ref}$  (solid) and  $O_{ref}$  (dashed) simulations at each SST (color scale).



338 In contrast to mean precipitation, the global-warming trend of extreme rainfall in-  
 339 tensities is accelerated in the case of convective organization, leading to super-CC rates of  
 340 increase. Figure 4b shows in color the trends that would be consistent with the Clausius-  
 341 Clapeyron formula, for each point  $(SST, P^{\circ})$  in the reference simulations  $\mathcal{O}_{\text{ref}}$  and  $\mathcal{D}_{\text{ref}}$ .  
 342 Here the trends are calculated using the Buck’s formula for accuracy [Buck A. L., 1981; A.  
 343 L. Buck, 1996]:

$$344 \quad P^{99.9}(T) \propto q_v^*(T) \propto \exp\left(\left(18.678 - \frac{T}{234.5}\right) \frac{T}{257.14 + T}\right). \quad (5)$$

345 For disorganized convection, the CC curves collapse onto each other, indicating a robust  
 346 agreement with CC. The dashed curves, representing CC trends for organized precipi-  
 347 tation, do not lie onto each other, indicating departures from the CC scaling. While the  
 348 overall contribution of organization is a positive departure from CC of about 2.5%/K,  
 349 these simulations exhibit several regimes, further detailed in section 5. Notably, super-CC  
 350 increases occur below 304K and above 306K, while sub-CC increases occur in between,  
 351 demonstrating a nonlinear global-warming behavior in the dynamics of extreme events due  
 352 to changes in organization.

353 Section 4 emphasizes the multiplicity of effects that self-aggregation can have on  
 354 the overall offset and trends in mean and extreme rainfall. Section 5 will investigate more  
 355 specifically the reason for the overall superCC trend in extremes and the three separate  
 356 regimes that appear as SST rises.

#### 357 **4 Shifts in mean climate and circulation strengthening**

358 In this section we will use the intermediate simulations  $\mathcal{O}(\mathcal{F}_{\mathcal{O}_{\text{ref}}})$ ,  $\mathcal{O}(\mathcal{F}_{\mathcal{D}_{\text{ref}}})$  and  
 359  $\mathcal{D}(\mathcal{Q}_{\mathcal{O}_{\text{ref}}})$  in an attempt to discuss the role of aggregation on mean and extreme rainfall  
 360 via two general mechanisms, using the decomposition shown on Figure 1:

- 361 (a) *shift in the mean climate state*: changes in the atmospheric radiative cooling can be  
 362 measured by comparing  $\mathcal{D}_{\text{ref}}$  and  $\mathcal{D}(\mathcal{Q}_{\mathcal{O}_{\text{ref}}})$  (light blue arrow), and shifts in mean sur-  
 363 face fluxes can be measured by comparing  $\mathcal{O}(\mathcal{F}_{\mathcal{O}_{\text{ref}}})$  and  $\mathcal{O}(\mathcal{F}_{\mathcal{D}_{\text{ref}}})$  (red arrow);  
 364 (b) *changes in the circulation* affect the local thermodynamic environment in which con-  
 365 vective clouds can form.  $\mathcal{D}(\mathcal{Q}_{\mathcal{O}_{\text{ref}}}) \rightarrow \mathcal{O}_{\text{ref}}$  (green) quantifies the total contribution from  
 366 changes in the circulation,  $\mathcal{D}_{\text{ref}} \rightarrow \mathcal{O}(\mathcal{F}_{\mathcal{D}_{\text{ref}}})$  (yellow) quantifies the role of the circu-  
 367 lation induced by the radiative feedback alone, and  $\mathcal{O}(\mathcal{F}_{\mathcal{D}_{\text{ref}}}) \rightarrow \mathcal{O}(\mathcal{F}_{\mathcal{O}_{\text{ref}}})$  (red) and

368  $O(\mathcal{F}_{O_{\text{ref}}}) \rightarrow O_{\text{ref}}$  (dark blue) both quantify the reinforcement of the circulation by the  
 369 surface-flux feedback.

370 Interpreting each contribution separately is obviously difficult, because these mechanisms  
 371 are coupled. In particular, the red arrow ( $O(\mathcal{F}_{\mathcal{D}_{\text{ref}}}) \rightarrow O(\mathcal{F}_{O_{\text{ref}}})$ ) can be understood as a  
 372 change in the mean climate, because it corresponds to a mean increase in sensible and  
 373 latent heat fluxes, but also as a change in the circulation, because surface fluxes in  $O_{\text{ref}}$  are  
 374 larger in response to the stronger surface winds caused by the radiative and surface flux  
 375 feedbacks. In addition, there is no unidirectional causality running from the strength of  
 376 the circulation to the strength of convection, because one adjusts to changes in the other.  
 377 However, this decomposition will associate changes in precipitation preferentially to shifts  
 378 in mean climate or in the circulation, and will also reveal the importance of the simulation  
 379 design for the distribution of rain.

380 We first aim at explaining the offsets between the exponential fits to  $\mathcal{D}_{\text{ref}}$  and  $O_{\text{ref}}$   
 381 shown in Figures 4a and 4b. They are calculated as

$$382 \quad P(\mathcal{S}) \approx \exp(\delta_{\mathcal{S}}T + \beta_{\mathcal{S}}) \quad (6)$$

383 where  $\delta_{\mathcal{S}}$  is the overall fractional change in  $P \in \{\bar{P}, P^Q\}$  for simulation  $\mathcal{S}$ . Recalling that  
 384 the pathways drawn on Figure 1 can be decomposed into a product of enhancement factors  
 385  $\mathcal{E}$  (Equation 1), the enhancement factor between two simulations  $\mathcal{S}$  and  $\mathcal{S}'$  can be written

$$386 \quad \mathcal{E}_{\mathcal{S} \rightarrow \mathcal{S}'} = \exp((\delta_{\mathcal{S}'} - \delta_{\mathcal{S}})T + \beta_{\mathcal{S}'} - \beta_{\mathcal{S}}) \quad (7)$$

387 The enhancement factors are reported in Table 2. In the case of mean precipita-  
 388 tion, the enhancement  $\mathcal{E}$  is roughly independent of  $T$ , for its fractional change is similar  
 389 in all simulations:  $\delta \approx 3.5\%/K$  shown on Figure 4. The mean total enhancement (the first  
 390 row of the table, denoted by  $\bar{P}$ ) is  $\mathcal{E}_{\text{total}} = 1.17$ , which corresponds to a 17% amplifica-  
 391 tion. In the first pathway (equation (1a)), it is mostly explained by the increase in mean  
 392 surface fluxes  $\mathcal{E}_{\text{shift}\mathcal{F}} = 1.19$  while the two other enhancing components are negligible  
 393  $\mathcal{E}_{\text{rad}} = 1.00$  and  $\mathcal{E}_{\text{surf}} = 0.99$ . This pathway simply highlights the consistency of a closed  
 394 water budget in steady state, so that changes in mean precipitation must match changes  
 395 in evaporation from the surface. In the second pathway (equation (1b)), the shift in mean  
 396 radiative cooling alone cannot explain the mean rainfall shift (i.e.,  $\mathcal{E}_{\text{shift}Q} = 1.11$  is less  
 397 than  $\mathcal{E}_{\text{total}} = 1.17$ ), because it is partly compensated by an increase in surface sensible  
 398 heat fluxes in simulation  $\mathcal{D}(Q_{O_{\text{ref}}})$  which dims the latent heat response (not shown). The

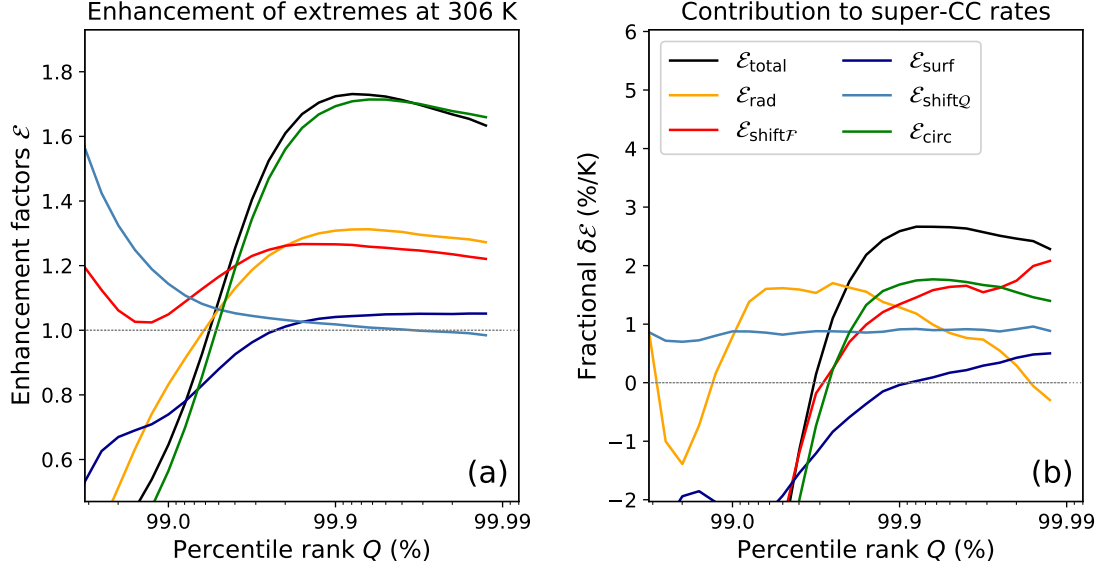
399 last term, corresponding to the experimental setup of previous studies [*Bao and Sherwood,*  
 400 2018], shows a complementary enhancement ( $\mathcal{E}_{\text{circ}} = 1.05$ ). Because  $\mathcal{D}(Q_{O_{\text{ref}}})$  and  $O_{\text{ref}}$   
 401 have equal mean radiative cooling, this last term can be interpreted as the effect of the  
 402 circulation on the surface enthalpy fluxes, via stronger surface winds and the drier near-  
 403 surface air in the subsidence region.

404 **Table 2.** Enhancing factors  $\mathcal{E}$  (see Fig. 1) to mean ( $\bar{P}$ ) and extreme precipitation averaged between per-  
 405 centiles 99.9 and 99.99 ( $\delta P^{\text{ext}}$ ), and contributions to the departures from Clausius-Clapeyron ( $d_{\text{CC}}$ ):

|  | 1st pathway (eq. 1a)             |                                |   | 2nd pathway (eq. 1b)            |                                   |                                 |
|--|----------------------------------|--------------------------------|---|---------------------------------|-----------------------------------|---------------------------------|
|  | $\mathcal{E}_{\text{total}}$ (T) | $\mathcal{E}_{\text{rad}}$ (R) | $\mathcal{E}_{\text{shift}\mathcal{F}}$ (F) | $\mathcal{E}_{\text{surf}}$ (S) | $\mathcal{E}_{\text{shift}Q}$ (Q) | $\mathcal{E}_{\text{circ}}$ (C) |
| $\bar{P} : \mathcal{E}$ (all SSTs)             | 1.17                             | 0.99                           | 1.19  | 1.00                            | 1.11                              | 1.05                            |
| $P^{\text{ext}} : \mathcal{E}$ (304-308K)      | 1.72                             | 1.27                           | 1.27  | 1.06                            | 1.00                              | 1.71                            |
| $d_{\text{CC}} : \delta\mathcal{E}$ (304-308K) | 2.57                             | 0.73                           | 1.62  | 0.22                            | 0.91                              | 1.66                            |

406 At fixed SST, these contributions can also be calculated for the enhancement of pre-  
 407 cipitation extremes, as shown on Figure 5a for the largest percentiles at 306 K and by the  
 408 second row of Table 1 labeled  $P^{\text{ext}}$ .  $\mathcal{E}_{\text{shift}Q}$  is close to 1, which suggests that for disorga-  
 409 nized convection, changes in domain mean radiative cooling has little effect on precipita-  
 410 tion extremes at this particular SST. It confirms recent results showing that radiative cool-  
 411 ing mostly affects weaker rain rates [*Chua et al., 2019*] and leaves the extremes unchanged  
 412 despite the adjustment of the mean surface fluxes. As a consequence, the response of ex-  
 413 tremes to self-aggregation at fixed SST can be attributed to the overall circulation and its  
 414 effect on surface fluxes:  $\mathcal{E}_{\text{circ}} \approx \mathcal{E}_{\text{total}} = \frac{P^Q(O_{\text{ref}})}{P^Q(\mathcal{D}_{\text{ref}})}$ . On the upper pathway, the spatial redis-  
 415 tribution of surface enthalpy fluxes ( $\mathcal{E}_{\text{surf}}$ ) only enhances  $P^Q$  by a few percent. The overall  
 416 contribution is evenly split between the circulation induced by the radiative feedback and  
 417 the mean effect of the surface-flux feedback:  $\mathcal{E}_{\text{rad}} \approx \mathcal{E}_{\text{shift}\mathcal{F}} \approx 1.3$ .

422 The decomposition is then applied to the super-CC rates of increase in  $P^Q$  (Fig-  
 423 ure 5b and the 3rd row of Table 2 labeled  $d_{\text{CC}}$ ). Departures  $d_{\text{CC}}$  from the CC scaling can  
 424 be decomposed as a sum of contributions coming from changes in the enhancing factors.  
 425 Denoting the fractional change in precipitation extremes for a given warming gap  $\Delta T$  by



418 **Figure 5.** Enhancement of precipitation intensities  $P^Q$  by the circulation (green), induced by the radiative  
 419 feedback (yellow), reinforced by the surface-flux feedback (dark blue and red), as well as the negligible effect  
 420 of the larger radiative cooling (light blue) in a steady climate at 306K (left, multiplicative contributions); and  
 421 corresponding contributions to the fractional changes  $\delta P^Q$  (right, additive contributions).

426  $\delta P^Q \equiv \Delta P^Q / P^Q / \Delta T$ , departures from CC can be written as

427 
$$d_{\text{CC}} \equiv \delta P^Q (O_{\text{ref}}) - \delta P^Q (D_{\text{ref}}) = \delta \mathcal{E}_{\text{surf}} + \delta \mathcal{E}_{\text{shift}\mathcal{F}} + \delta \mathcal{E}_{\text{rad}} \quad (8a)$$

428 
$$= \delta \mathcal{E}_{\text{circ}} + \delta \mathcal{E}_{\text{shift}Q}. \quad (8b)$$

429

430 where likewise  $\delta \mathcal{E}_i \equiv \Delta \mathcal{E}_i / \mathcal{E}_i / \Delta T$ . Each term is independent of SST and can be calculated  
 431 as a difference in fractional changes for pairs of simulations:  $\delta \mathcal{E}_{S \rightarrow S'} = \delta_{S'} - \delta_S$ . Sur-  
 432 prisingly, the increase in mean radiative cooling  $\delta \mathcal{E}_{\text{shift}Q}$  does bring a small contribution  
 433 of just under 1%/K to the superCC rate (0.91%/K in Table 2), despite its small role on  
 434 extreme precipitation in fixed climates. This occurs because the latent heat flux increases  
 435 more rapidly than the sensible flux in the  $\mathcal{D} (Q_{O_{\text{ref}}})$  simulation as a function of SST, re-  
 436 inforcing convection (not shown). The remaining contribution comes from the overall  
 437 circulation and the additional surface evaporation that it causes ( $\delta \mathcal{E}_{\text{circ}} \approx 1.5\%/K$ ). On  
 438 the upper pathway, most of the change is attributable to the mean enhancement in evap-  
 439 oration due to the surface-flux feedback ( $\delta \mathcal{E}_{\text{shift}\mathcal{F}} \approx 1.6\%/K$ ). Changes in the radiatively-  
 440 driven circulation have its strongest contribution between the 99th and 99.9th percentiles  
 441 ( $\delta \mathcal{E}_{\text{rad}} \approx 1.7\%/K$ ) and affects less the heaviest rain events ( $\approx 0.7\%/K$ ).

442 Testing further the sensitivity of mean and extreme rainfall to specific self-aggregation  
 443 mechanisms would likely require a more physically-based analytic framework which con-  
 444 nects the strength of the circulation to the strength of precipitation extremes. In our case,  
 445 feedback mechanisms maintain the circulation and affect the mean climate simultaneously,  
 446 and their effect on rain statistics cannot be easily separated. In particular, the surface-flux  
 447 feedback alone is not sufficient to maintain organization [Holloway and Woolnough, 2016],  
 448 as seen in the  $\mathcal{D}(Q_{O_{\text{ref}}})$  experiments, so its effectiveness to strengthen the circulation can-  
 449 not be physically separated from the effectiveness of the radiative feedback to organize  
 450 convection in the first place.

451 Despite the difficulty to tease apart these contributions, two main conclusions can be  
 452 drawn:

- 453 • Even for a fixed mean climate (identical values of SST, mean surface fluxes and ra-  
 454 diative cooling), the presence of a mesoscale circulation alone can induce heavier  
 455 precipitation extremes, thereby causing a stronger skewness of the distribution of  
 456 rain. This result indicates that even the rarest and heaviest rain events can be un-  
 457 derstood as an integral component of the circulation. This is seen in this section by  
 458 the significant enhancement  $\mathcal{E}_{\text{rad}}$ , although an enhancement of similar magnitude  
 459 occurs in response to the stronger surface fluxes.
- 460 • The superCC trend in precipitation extremes is mainly associated with the mean  
 461 enhancement of surface fluxes, indicating that the surface-flux feedback plays a ma-  
 462 jor role in shaping the distribution of rain through its possible effects on convective  
 463 strength. This result suggests that organized precipitation extremes could be highly  
 464 sensitive to changes in the large-scale winds in the real atmosphere because of their  
 465 effects on surface evaporation.

466 Section 5 will investigate how a reinforcement of the circulation can amplify the  
 467 heaviest rain intensities.

## 468 **5 The dynamics of disorganized and organized extremes**

### 469 **5.1 Thermodynamic, dynamic and precipitation efficiency contributions**

470 We now seek to understand the cause for the overall super-CC trend in precipita-  
 471 tion extremes  $P^{\text{Q}}$  across steady RCE states at different SSTs as well as the reasons for the  
 472 super- and sub-CC regimes seen in section 3.2 in the  $O_{\text{ref}}$  simulation. To do so, we use

473 a scaling formula to approximate the largest rainfall percentiles  $P^Q$  from the average dy-  
 474 namic and thermodynamic characteristics of extreme events [O’Gorman and Schneider,  
 475 2009]. This expression approximates the condensation rate at each level in the troposphere  
 476 as the vertical advection of saturation specific humidity  $q_v^*$  at speed  $w$  along a moist adia-  
 477 bat  $\theta_e^*$ , and integrates it along the vertical to estimate the surface precipitation rate:

$$478 \quad P^Q \approx -\varepsilon \int_0^{z^T} w^Q \left. \frac{\partial q_v^*}{\partial z} \right|_{\theta_e^{*Q}} \rho dz. \quad (9)$$

479 Superscript  $Q$  denotes that variables have been composited at the locations of extreme  
 480 events: for all surface precipitation rates counted as the  $Q$ th percentile of the distribu-  
 481 tion, the profiles are sampled 1h early and averaged across rainfall events in order to re-  
 482 construct the approximate conditions in which the corresponding convective clouds were  
 483 formed. The 1h time scale happens to match the characteristic timescale of convective  
 484 updrafts and corresponds to the output time step. Coefficient  $\varepsilon$  can be interpreted as a  
 485 proxy for the *precipitation efficiency of extremes*: as such,  $\varepsilon$  approximates the fraction  
 486 of condensed water that reaches the surface, while  $1 - \varepsilon$  corresponds to the fraction of  
 487 cloud water that mixes in the environment.  $\varepsilon$  is calculated as a tuning coefficient from a  
 488 least-square fit between true percentiles  $P^Q$  and approximated percentiles  $\int w^Q \partial_z q_v^*|_Q$  be-  
 489 tween the 99.9th and 99.99th percentile ranks. Although  $\varepsilon$  is an efficiency coefficient, it  
 490 remains a tuning parameter and could potentially exceed 1. It can be affected by sampling  
 491 issues and its variability can reveal a limited explanatory power of the approximation for-  
 492 mula [Fildier et al., 2018] but it also embeds additional processes such as entrainment of  
 493 dry air and subtle differences in precipitation microphysics that could affect precipitation  
 494 extremes [Muller, 2013].

495 This formulation facilitates decomposing the fractional changes  $\delta P^Q$  into a sum of  
 496 independent contributions coming from changes in the dynamic and thermodynamic prop-  
 497 erties of extreme rainfall events. Denoting the pressure-weighted integral across the tropo-  
 498 sphere  $\langle X \rangle = \int_{p_s}^{p^T} X \frac{dp}{g} \approx - \int_0^{z^T} X \rho dz$ , this scaling approximation can be rewritten as

$$499 \quad P^Q \approx \underbrace{\varepsilon \langle w^Q \rangle}_M \underbrace{\left\langle \frac{w^Q}{\langle w^Q \rangle} \left( \left. \frac{\partial q_v^*}{\partial z} \right|_{\theta_e^{*Q}} \right) \right\rangle}_\Gamma = \varepsilon M \Gamma \quad (10)$$

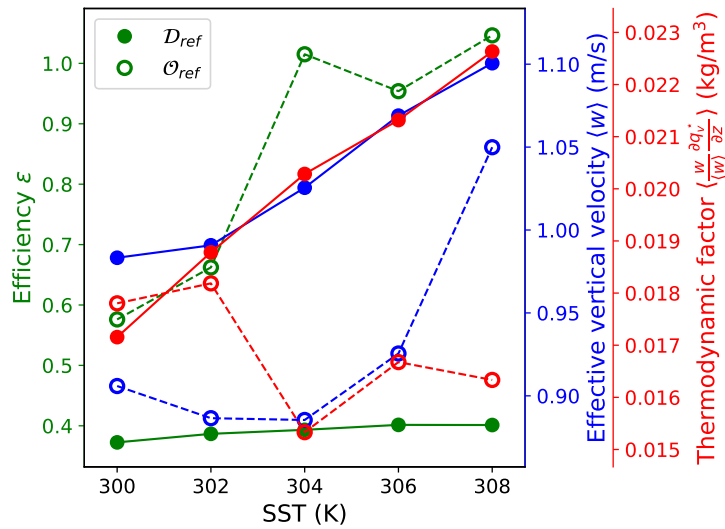
500 where  $M$  represents the vertically-integrated mass flux across the depth of the troposphere,  
 501 such that  $M/\langle 1 \rangle$  represents an effective cloud updraft velocity. In turn  $\Gamma \times \langle 1 \rangle$  is a thermo-  
 502 dynamic term representing an effective amount of moisture available for condensation, in  
 503  $\text{kg/m}^3$ .

504 In the case of disorganized convection, the fractional increase in extremes  $P^Q$  can be  
 505 decomposed in a sum of fractional contributions expressed in %/K, because the terms  $\varepsilon$ ,  
 506  $M$  and  $\Gamma$  all follow exponential increases with warming:

$$507 \quad \delta P^Q \approx \underbrace{\delta \varepsilon}_{\substack{\text{Contribution from} \\ \text{changes in:} \\ \text{precipitation} \\ \text{efficiency}}} + \underbrace{\delta M}_{\substack{\text{mass} \\ \text{flux}}} + \underbrace{\delta \Gamma}_{\text{thermodynamics}} \quad (11)$$

508 These individual contributions, precipitation efficiency  $\varepsilon$ , pressure-weighted mean vertical  
 509 velocity  $M/\langle 1 \rangle$ , in m/s, and the remaining thermodynamic term  $\Gamma \times \langle 1 \rangle$  which represents  
 510 the moisture available for condensation in  $\text{kg/m}^3$ , are shown on Figure 6 for each SST for  
 511 simulations  $\mathcal{D}_{\text{ref}}$  and  $\mathcal{O}_{\text{ref}}$ .

512 For the reference disorganized case  $\mathcal{D}_{\text{ref}}$ , precipitation efficiency is constant around  
 513 0.4, and the CC increase is explained by a joint and steady increase in the thermodynamic  
 514 and dynamic components. Indeed, the CC increase of 5.6%/K results from an increase in  
 515 the thermodynamic term  $\Gamma$  of +3.5%/K and increase in mass flux of +1.5%/K, while pre-  
 516 cipitation efficiency is roughly constant with warming (+0.8%/K). These numbers are con-  
 517 sistent in relative magnitude with previous analyses of disorganized extremes [e.g. *Romps*,  
 518 2011].



519 **Figure 6.** Precipitation efficiency (green), and the dynamic (blue) and thermodynamic (red) components of  
 520 extremes rainfall  $P^Q$  for simulations  $\mathcal{D}_{\text{ref}}$  and  $\mathcal{O}_{\text{ref}}$ .

521 In  $\mathcal{O}_{\text{ref}}$ , this decomposition allows to investigate the three regimes identified in sec-  
 522 tion 3.2 (superCC-subCC-superCC). The superCC rates seem to be explained by changes

523 in precipitation efficiency and updraft speeds. First, the large and fast increase in  $\varepsilon$  ex-  
 524 plains the first superCC regime between 300K and 304K, and it stabilizes around 1. Then,  
 525 the delayed and large increase in vertical velocities at high SSTs seems to strongly con-  
 526 tribute to the last superCC regime between 306K and 308 K. The gradual decrease in  $\Gamma$   
 527 and the delayed amplification of the dynamic component at higher SSTs can be related  
 528 to changes in vertical velocity profiles (Figure 7b): the profile gradually shifts towards  
 529 higher altitudes, causing a reduction in  $\Gamma$ , and vertical velocities strongly increase in the  
 530 upper troposphere at high SST. This explains the simultaneous enhancement of  $M$ , but  
 531 understanding what actually sets the shape of vertical velocity profiles remains an open  
 532 question.

533 Explaining why such nonlinear dynamics occur would likely require a closer anal-  
 534 ysis of the interactions between precipitation efficiency, mixing with environmental air,  
 535 and buoyancy. Although these results should be interpreted with care, it is unlikely that  
 536 this behavior be an artifact of the scaling approximation: this formula closely captures the  
 537 shape of the distribution tail at different SSTs, with correlation coefficients above 0.95 in  
 538 all cases between the 99.9th and 99.99th percentiles, and the approximate magnitude of  
 539  $\varepsilon$  is realistic for disorganized convection [Lutsko and Cronin, 2018]. However, limitations  
 540 of the approximation formula could lead to overestimates of precipitation efficiency, and  
 541 sampling uncertainties could have caused small fluctuations in  $\varepsilon$  and the subCC increase  
 542 in extremes between 304K and 306K.

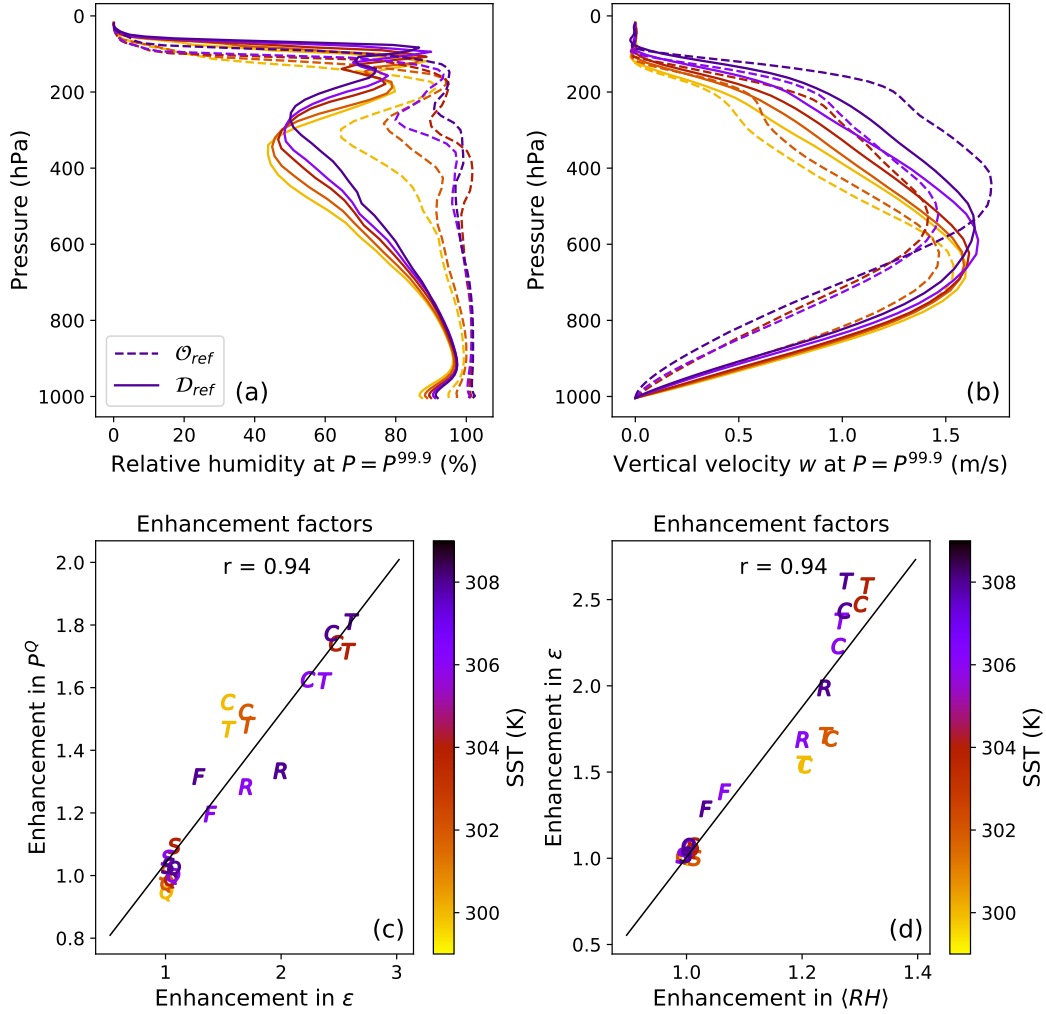
## 543 **5.2 Amplification of extremes by larger precipitation efficiencies**

544 We now make a stronger case for the possible role of precipitation efficiency in am-  
 545 plifying the strength of extremes. Figure 7a shows the vertical profile of relative humidity  
 546 at the location of precipitation extremes for the reference disorganized and organized sim-  
 547 ulations. Convective aggregation appears to bring the environmental air close to 100%  
 548 relative humidity in the moist environment in which deep convective clouds form. This  
 549 lower environmental saturation deficit could have a doubly enhancing effect on the surface  
 550 precipitation rate: it would increase condensation, or conversion efficiency, by reducing  
 551 the dilution of cloud parcels into their environment, and could increase sedimentation ef-  
 552 ficiency by reducing the re-evaporation of condensates [Langhans et al., 2015; Lutsko and  
 553 Cronin, 2018]



554            Figures 7c,d show the correspondence between the amplification in extreme pre-  
555            cipitation  $P^Q$ , in  $\varepsilon$  and in vertically-integrated relative humidity  $\langle RH \rangle$  at the location of  
556            extremes, for all SSTs and all pairs of simulations as explained on Figure 1. The large  
557            correlation on panel (c) confirms the importance of fluctuations in precipitation efficiency  
558            for the strength of precipitation extremes. The large correlation on panel (d) suggests that  
559            precipitation efficiency closely relates to the degree of saturation in the moist environment  
560            where extreme events occur. This amplification is related to the strength of aggregation,  
561            as this effect is more pronounced for simulations that combine several aggregation feed-  
562            backs (for instance, enhancements  $C$  and  $R$  on Figures 7c,d correspond to the enhance-  
563            ments  $\mathcal{E}_{\text{total}}$  and  $\mathcal{E}_{\text{rad}}$  respectively, introduced in section 4).

564            These results highlight the key role played by the precipitation efficiency in modu-  
565            lating precipitation extremes with organization. Further investigation of the sensitivity of  
566            efficiency to mesoscale organization is desirable to achieve accurate rainfall predictions in  
567            a warming climate.



568 **Figure 7.** (a-b) Relative humidity and vertical velocity profiles at location of precipitation extremes. (c-d)  
 569 Correlations between the enhancement in precipitation extremes  $P^Q$ , precipitation efficiency  $\varepsilon$  and vertically  
 570 integrated relative humidity  $\langle RH \rangle$  for the pairs of simulations defined in Figure 1 (letters), at all SSTs (colors).

571 **6 Discussion**

572 The present article summarizes the behavior of mean and extreme precipitation in  
 573 conditions of disorganized and organized convection in large 3D domains. The analysis is  
 574 based upon a diversity of forcing strategies and application of a large range of SSTs in order  
 575 to highlight the strong sensitivity of precipitation in steady RCE states to the strength  
 576 of the internal circulation and the global atmospheric energy budget. This section dis-  
 577 cusses additional lessons drawn from the variety of forcings and presents some remaining  
 578 modeling simplifications that can affect the scope of these results.

## 579 6.1 Connection between mean climate, circulations and convection

580 Our analysis assumes that the response of mean precipitation, which is driven by  
 581 the global impact of organization on the energy balance, is distinct and independent from  
 582 the response of extreme precipitation, which is driven by the strength of the circulation  
 583 through local changes in saturation deficit and precipitation efficiency. Even though this  
 584 separation appears valid to first order, changes in mean climate properties are tightly con-  
 585 nected to the horizontal structure of convection itself. In the case of convective aggrega-  
 586 tion, this complicates the attribution of causality and blurs the traditional distinction be-  
 587 tween the sensitivity of mean and extreme precipitation to warming.

588 In particular, understanding how the shape of the rainfall distribution and the inten-  
 589 sity of rainfall extremes respond to changes in the atmospheric energy budget remains an  
 590 open question. In lack of a better theory, *Pendergrass and Hartmann [2014b]* make the  
 591 objective assumption that the energetic constraint on mean precipitation changes applies  
 592 uniformly across all percentiles of the distribution of rain amounts. This uniformity as-  
 593 sumption has later been disproved. *Chua et al. [2019]* show that doubling radiative cool-  
 594 ing at fixed SST affects weak precipitation rates while leaving the heaviest rain rates un-  
 595 changed. In contrast, *Thackeray et al. [2018]* show that changes in mean precipitation  
 596 do correlate with changes in heavy precipitation rates across GCMs, suggesting that the  
 597 strength of precipitation extremes is not actually independent from the domain-mean atmo-  
 598 spheric radiative cooling.

599 In the case of disorganized convection, our simulations tend to confirm *Chua et al.*  
 600 [2019]’s result because extremes are unchanged between the  $\mathcal{D}_{\text{ref}}$  and  $\mathcal{D}(Q_{O_{\text{ref}}})$  simula-  
 601 tions at fixed SST (Figure 4). But in the case of organized convection, changing the mean  
 602 climate at fixed SST does affect the extremes, as seen by the enhancement  $\mathcal{E}_{\text{shift}\mathcal{F}} > 1$   
 603 on Figure 5. In this example, both the domain-mean atmospheric radiative cooling and  
 604 the circulation adjust to the increased surface fluxes, which leads to a simultaneous re-  
 605 inforcement of mean and extreme precipitation rates. In particular, departures from their  
 606 exponential fits can be seen at the same SSTs for  $\overline{P}$  and  $P^{\text{Q}}$  on Figure 4a,b.

607 Overall, a plausible link exists between mean and extreme rainfall in given climates,  
 608 and this work suggests that the connection possibly lies in mesoscale circulation changes.  
 609 As a result, these mesoscale circulations could also act as an additional degree of freedom

610 on their rate of change with warming as well as on the increased skewness of the rain dis-  
611 tribution.

## 612 **6.2 Sources of methodological uncertainty**

613 Further analysis is required to test the relationships between mean/extreme precip-  
614 itation and convective organization in more realistic representations of the tropical atmo-  
615 sphere. In particular, it seems necessary to quantify on which scales these relationships  
616 hold and validate the physical relevance of the superCC and subCC regimes identified.  
617 Indeed, several implicit methodological assumptions could bias our interpretation.

618 First, the limited domain size is suspected to amplify the degree of convective or-  
619 ganization [*Cronin and Wing, 2017*], which could cause an artificial amplification of heavy  
620 precipitation intensities. On larger tropical domains than we employed, longwave radiative  
621 fluxes would tend to stretch the size of moist patches [*Beucler and Cronin, 2019*], which  
622 could damp the strong amplification in relative humidity that we find in our study. This  
623 could prevent the superCC rates that arise from increases in precipitation efficiencies, or  
624 transpose this behavior to higher SSTs. An improved characterization of the relationship  
625 between strength of self-aggregation feedbacks and domain size seems necessary to quan-  
626 tify the sensitivity of rain intensities to specific modeling choices.

627 The second issue resides in the coarse 4km resolutions used. They prevent the ad-  
628 equate representation of low clouds and tend to bias the system towards an excessively  
629 dry free troposphere in the subsiding regions [*Holloway et al., 2017*]. This could lead  
630 to an overestimate of mean rainfall increases due to the model's inability to resolve the  
631 absorption of upwelling shortwave radiation reflected by stratocumulus. In addition to  
632 its role on turbulent mixing as well as rain reevaporation and cold pools [*Jeevanjee and*  
633 *Romps, 2013*], the coarse resolutions could bias the strength of self-aggregation feed-  
634 backs by modifying the low-level circulation induced by longwave radiative cooling in  
635 low-clouds [*Muller and Held, 2012*], thus affecting the thermodynamic environment in  
636 which extreme events occur.

637 Third, the timescale required for these experiments to reach equilibrium is substan-  
638 tially larger than the lifetime of typical MCSs [*Houze, 2004*]. Even though the effect of  
639 self-aggregation on the largest percentiles of the distribution does appear gradually during  
640 the transition between disorganized and organized states (not shown), the actual impor-

641 tance of aggregation for precipitation rates in shorter-lived convective systems could be  
642 smaller, and deserves further investigation.

643 A fourth kind of methodological limitation is the quantification of precipitation ex-  
644 tremes themselves. In particular, the distribution of rain calculated over the entire domain  
645 actually depends on the relative occurrence of different precipitation regimes. As a result,  
646 organized RCE states over idealized small domains likely exhibit different rainfall distri-  
647 butions and stronger rainfall extremes than larger simulation domains with realistic occur-  
648 rences of deep and shallow convective systems, even for the same conditions of forcing  
649 and in a similar mean climate.

650 In addition to these methodological considerations, additional uncertainties arise  
651 from the magnitude and type of forcing conditions, in addition to the choice of closure  
652 required for unresolved processes. Winds, turbulence and cloud microphysics could affect  
653 extremes indirectly through the strength of self-aggregation feedbacks and convective mix-  
654 ing, but also directly, by affecting convective dynamics, mixing, and condensation. *Bao*  
655 *and Sherwood* [2018] also report changes in precipitation efficiency with convective or-  
656 ganization, but connect them to a microphysical response rather than changes in the local  
657 environment of convective updrafts (we note though that both can be related, as for in-  
658 stance changes in low-level relative humidity can affect the evaporation of rain and thus  
659 precipitation efficiency). They document changes in condensate species with a reduction  
660 in graupel production and argue that it enhances precipitation efficiency while reducing  
661 buoyancy and updraft velocity. Further analysis is desirable to test their results in SAM  
662 with two-moment microphysics schemes and under a large range of SSTs and forcing con-  
663 ditions.

664 Besides, *Lane and Moncrieff* [2015] and *Moncrieff and Lane* [2015] showed that the  
665 propagation, vertical structure and organizational properties of MCSs depend on a balance  
666 between the strength of cold pools, convective inhibition and subcloud layer saturation  
667 levels, which largely varies with large-scale wind shear conditions. Such analyses have not  
668 yet been extended to the context of self-aggregation in RCE, so that the sensitivity of self-  
669 aggregation feedbacks, precipitation efficiency and rainfall extremes to wind shear is still  
670 largely unknown.

## 7 Conclusion

Convective organization provides a framework to study the interaction between atmospheric circulations that occur on long time scales, and convective processes that occur on short time scales. Our simulations show that these mesoscale circulations can directly change the spatial distribution of thermodynamic variables such as moist static energy and relative humidity and affect local convective processes and the statistics of rain. The reverse interaction, namely the effect of local convective activity on the circulation strength or on the redistribution of moisture [Romps, 2014, e.g.], is outside the scope of the present study but is of importance when estimating the strengthening of organization in future climates.

This work provides evidence that a stronger degree of convective organization in warmer climates can be associated with a faster increase in extreme rainfall intensities than what disorganized convection indicates, with an excess relative to the Clausius-Clapeyron scaling by 2.5%/K in this modeling setup. These departures correspond to variations in the strength of the organization and concomitant mesoscale circulation, here simply modified by turning on and off the radiative and surface flux feedbacks. When convection is organized, heavier extreme rainfall intensities strongly correlate with larger precipitation efficiencies: this likely results from the enhanced air saturation in the moistest areas of the domain and a reduced ability of cloud parcels to mix with dry air as they form.

Because this coupling between convection and the large-scale circulation cannot be resolved on the coarse GCM grids, it could be an important source of bias for the intensification of extreme rain events in current climate models. The magnitude of self-aggregation feedbacks in future climates has not been quantified yet because of the sensitivity of aggregation to model parameterizations and differences across simulation designs. Some of the processes from which these uncertainties originate can actually affect convection and precipitation directly, such as cloud microphysics and turbulent processes. As a result, improved understanding of future changes in self-aggregation and precipitation extremes could be gained simultaneously by focusing on the sensitivity of these two processes to the model formulations for cloud microphysics, turbulence, surface enthalpy fluxes and radiative transfer.

This work has demonstrated a changing behavior of precipitation extremes as a function of SST via multiplicative effects of self-aggregation feedbacks through a change in

703 the strength of aggregation. In section 6, it is proposed that the mesoscale circulations  
704 could possibly modulate the response of mean and extreme rain to global warming simul-  
705 taneously. These nonlinear behaviors raise fundamental questions about the design of ide-  
706 alized experiments and the nature of their connection to the real atmosphere. Can these  
707 idealized model configurations (with self-aggregation and without) be interpreted quantita-  
708 tively to estimate shifts in the hydrologic cycle? What correspondence can be achieved to  
709 map these results onto realistic estimates? *Tan et al.* [2015] show that a large uncertainty  
710 in current estimates of changing precipitation extremes with warming comes from lack of  
711 knowledge of how convective organization will change with warming. So, improved fun-  
712 damental understanding of convective organization and its sensitivity to warming is hence  
713 an area of priority for climate model development to achieve accurate rainfall projections  
714 in a warming climate.

## 715 **Acknowledgments**

716 The authors gratefully acknowledge diverse funding agencies and resources used for  
717 this work. It was initiated as part of a PhD thesis [*Fildier, 2019*] under the CASCADE  
718 project, supported by the Director, Office of Science, Office of Biological and Environ-  
719 mental Research of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231  
720 and DE-SC0012548. It was conducted as part of the Regional and Global Climate Mod-  
721 eling Program and used the computational cluster resource Cori provided by the IT Divi-  
722 sion at the Lawrence Berkeley National Laboratory. This work was later completed using  
723 funding from the European Research Council (ERC) under the European Union’s Hori-  
724 zon 2020 research and innovation programme (Project CLUSTER, grant agreement No  
725 805041).

726 The Python scripts developed for this analysis are freely available at [https://](https://github.com/bfildier/Fildier2020)  
727 [github.com/bfildier/Fildier2020](https://github.com/bfildier/Fildier2020). Because of the large volume of model outputs,  
728 the simulation results are available from the authors upon request ([benjamin.fildier@lmd.ens.fr](mailto:benjamin.fildier@lmd.ens.fr))  
729 and are archived at the National Energy Research Scientific Computing Center (NERSC).

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