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1 **Acoustical properties of an immersed Corner-Cube Retroreflector alone**
2 **and behind screen for ultrasonic telemetry applications**

3

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17

18 **Abstract:**

19 Ultrasonic telemetry measurements consist in remotely detecting and locating an
20 object. To maximize the signal-to-noise ratio, a target may be used, positioned at a
21 reference point. In this framework, the ultrasonic reflective characteristics of a corner-
22 cube retroreflector (CCR) are investigated. The most interesting property of a CCR is its
23 ability to fully reverse an incoming wave in the same direction under certain
24 conditions. Theoretical developments are performed in order to understand its
25 acoustic behaviour, and experimentations are made in various configurations: CCR
26 alone in water, and behind an immersed plate that acts as a screen, with normal and
27 non-normal incidence. The results highlight its strong performance. Moreover, the
28 study of two other couples of CCR material and surrounding fluid underlines the
29 relevance of considering the acoustic properties of each medium, as they have a strong
30 influence on the acoustic response of the CCR.

31

32

33 **Keywords:** Corner-Cube Retroreflector; trihedral reflector; acoustical properties;
34 ultrasonic telemetry; telemetry behind screen(s)

35

36 **1. Introduction and state of the art**

37 The concept of sodium-cooled fast reactor (SFR) is being studied in France for the 4th
38 generation of nuclear power plants [1]. The improvement of in-service inspection and
39 repair (ISI&R) is a major issue in the case of nuclear power plants. Ultrasonic solutions
40 are particularly adapted to this harsh environment, and are therefore being studied for
41 the inspection of the main vessel and its internal components. In this context, different
42 methods to perform ultrasonic NDT and NDE are being investigated by CEA and its
43 partners [2].

44 Localisation and positioning of an object by using ultrasounds require an appropriate
45 relative orientation between the object and the ultrasonic beam. In pulse-echo mode,
46 sufficient specular energy has to travel back towards the transducer to be analysed
47 and interpreted. This condition may fail to occur for example in the case of a
48 misaligned or nonplanar object. A solution studied here consists in using of a specific
49 target part attached to the object, this target being designed to reflect well towards
50 the emitter. The Corner-Cube Retroreflector (CCR), also named retroreflector, corner
51 reflector or trihedral rectangular reflector in the literature, is an efficient target found
52 to reflect incoming waves back to the emitter even when these incoming waves do not
53 strike it perpendicularly (over a certain angular range, about $\pm 30^\circ$). Retroreflectors are
54 common in everyday life, in the field of optics [3,4] (for example as security reflectors
55 on bikes), or for advanced technological applications, such as satellite and lunar laser
56 ranging systems [5,6], or else with radar waves for geophysical observations and
57 monitoring [7–11].

58

59 In the field of ultrasound, CCR targets remain rather unknown despite their obvious
60 efficiency. A few authors [12–18] propose various ultrasonic reflectors (conical, linear,
61 trihedral etc.) for very specific applications, such as transducer characterisation or
62 determination of misalignment of a transducer. For telemetry applications, CCR seems
63 the best candidate because not only does it send back a lot of energy even if
64 misaligned but also it provides an easy-to-detect spatial reference. However, no
65 extensive study was really performed on the acoustical properties of a CCR alone or
66 behind screens.

67 A CCR acts as a planar reflector (waves are reflected parallel to the incidence direction)
68 even off “normal incidence”, and the reflection can be total under certain conditions.
69 The reflected amplitude level depends on the angular incidence, on the alignment of
70 the transducer with the CCR inner vertex and on the two media (CCR material and
71 surrounding liquid). It is more or less equivalent to a corner effect, but in 3D. The main
72 developments (theoretical and experimental) are performed here in water, but the
73 ultimate objective is telemetry applications in liquid sodium. Thus to try to get closer
74 to this case, a CCR made of copper is also studied. This paper aims at reviewing and
75 enhancing knowledge of the acoustical properties of a CCR, from a theoretical and
76 experimental point of view, in the general framework of ultrasonic telemetry on an
77 object which may have moved out of position and potentially hidden by one (or
78 several) immersed plate(s). In this case the target is fixed on the object of interest.

79

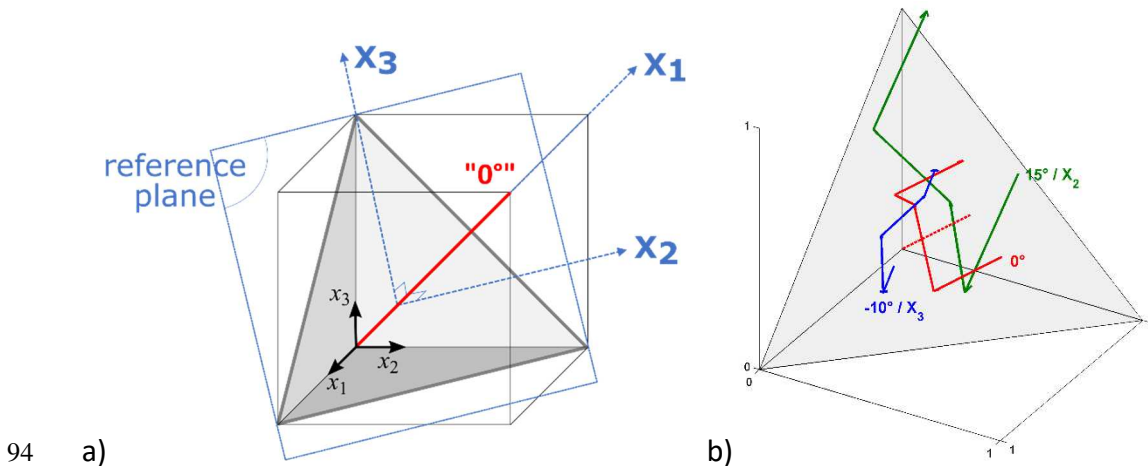
80 Section 2 presents the theoretical background and results. Then experimentations in
81 various configurations are detailed and analysed in section 3. Finally, in section 4 we

82 study the influence of different acoustic properties of the CCR material and the
83 surrounding fluid.

84

85 **2. Theoretical background and modelling**

86 A CCR consists of three mutually orthogonal planes (isosceles right triangular faces)
87 joining the inner vertex as schematized in Figure 1a. In the following, the “incidence
88 angle” on the CCR designates the incidence relative to the reference plane (X_2X_3) which
89 is the cutting plane of the cube the CCR is virtually obtained from. An incident ray, non-
90 aligned with the vertex, will reflect successively on the three faces (with a different
91 local incidence on each face) and will result in a final reflected ray parallel to the initial
92 incident one. Figure 1b shows examples of incident rays and their successive
93 reflections with respect to Snell’s law.



94 a) Definition of coordinate systems, and normal incidence, referred to as “0°”

95 and b) examples of successive reflections of a ray on the faces with various incidences

96

97
98 Moreover, after the three reflections, the times-of-flight of all the rays are the same in
99 the plane parallel to (X_2X_3) containing the ray starting point, equal to the time-of-flight

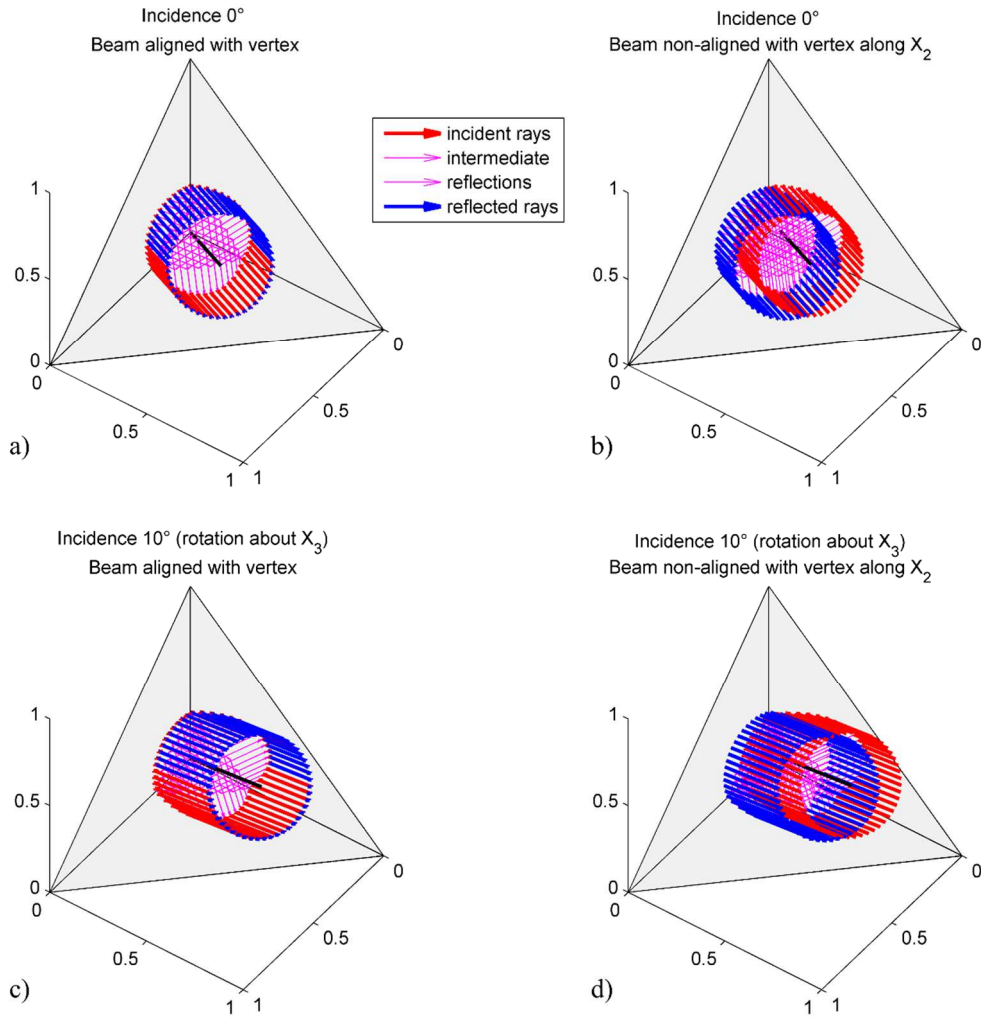
100 of the travel on the X_1 -axis impacting the inner corner of the target. Then the
101 calculation of the distance from the time-of-flight will provide the distance between
102 the transducer and the inner corner of the target.

103 Note that there are ineffective (shadow) areas of the CCR that reflect an incident ray
104 only partially (one or two reflections instead of three). For normal incidence, such
105 areas are the three triangles sizing a third of the CCR side located at the three
106 endpoints of the CCR [7,11].

107

108 If we consider now the contour of the incident beam (modelled by a set of rays
109 cylindrically arranged), it can be easily known whether the reflected beam will merge
110 with the incident beam or not. As shown by the examples in Figure 2, whatever the
111 incidence angle, if the reflected beam exists (in terms of amplitude, see below), it will
112 merge with the incident beam only if the latter is aligned with the inner vertex of the
113 CCR. Otherwise the reflected beam is shifted from the initial incident beam location,
114 the angle of propagation remaining the same.

115



116

117 Figure 2: Calculation of cylindrical beam reflection as a function of the angle of incidence and
 118 the alignment with the inner vertex (the total length of the arrows of one ray travel
 119 corresponds to a single time-of-flight)

120

121 The overall reflected amplitude depends on the successive local reflections on the CCR
 122 faces. If we consider an incident direction noted \vec{U}_{inc} (normalized vector) in the
 123 coordinate system $(X_1X_2X_3)$, it becomes \vec{u}_{inc} in the coordinate system $(x_1x_2x_3)$:

124

$$\vec{u}_{inc} = M\vec{U}_{inc}$$

125 where $M = \begin{bmatrix} \sqrt{3}/3 & -\sqrt{2}/2 & -\sqrt{6}/6 \\ \sqrt{3}/3 & \sqrt{2}/2 & -\sqrt{6}/6 \\ \sqrt{3}/3 & 0 & \sqrt{6}/3 \end{bmatrix}$ is the transformation matrix between the two

126 coordinate systems.

127 The local incidence angles on each face are thus given by:

128
$$\alpha_{(x_1x_2)} = \cos^{-1} \left(-\vec{u}_{inc} \cdot \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \right)$$

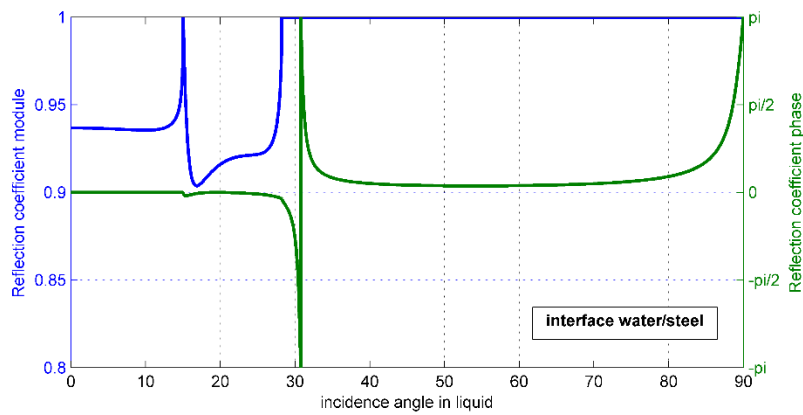
129
$$\alpha_{(x_1x_3)} = \cos^{-1} \left(-\vec{u}_{inc} \cdot \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} \right)$$

130
$$\alpha_{(x_2x_3)} = \cos^{-1} \left(-\vec{u}_{inc} \cdot \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \right)$$

131

132 At normal incidence for example (i.e. $\vec{U}_{inc} = -\vec{X}_1$), the three local incidence angles are
133 equal, and their value is about 54.7°. The theoretical reflection coefficient [19,20] in
134 the case of water as surrounding liquid and stainless steel as CCR metal is displayed in
135 Figure 3. The two critical angles are 14.9° for the disappearance of longitudinal waves
136 and 28.2° for the disappearance of transversal waves, and conversion into Rayleigh
137 wave occurs at 30.7°. It can therefore be concluded that local reflection on one face is
138 total at 54.7° incidence and thus global reflection at “normal incidence” on the CCR is
139 total. Therefore, in water, at normal incidence, a steel CCR is more efficient as a
140 reflector than a steel plate, provided that the CCR is larger than the beam.

141



142

143

Figure 3: Reflection coefficient (module and phase) at the water/steel interface

144

as a function of the incidence angle

145

146

Moreover if we consider arbitrary incidence on a geometrically infinite CCR, the

147

calculation of local incidence angles on each face (see Figure 4) coupled with the

148

reflection coefficient of Figure 3 makes it possible to conclude that:

149

- For incidence angles about the X_3 -axis from -36° to $+36^\circ$ (incidence in X_1X_2 plane),

150

the global reflection is always total, except in the cases of Rayleigh wave

151

generation at -31.1° and $+31.1^\circ$. For incidence angles between 36° and 40° , some

152

energy is transferred into the metal thanks to its conversion into transversal

153

waves. The global reflection will thus be weaker.

154

- For incidence angles about the X_2 -axis from -40° to $+26.7^\circ$ (incidence in X_1X_3

155

plane), the global reflection is always total, except in the cases of Rayleigh wave

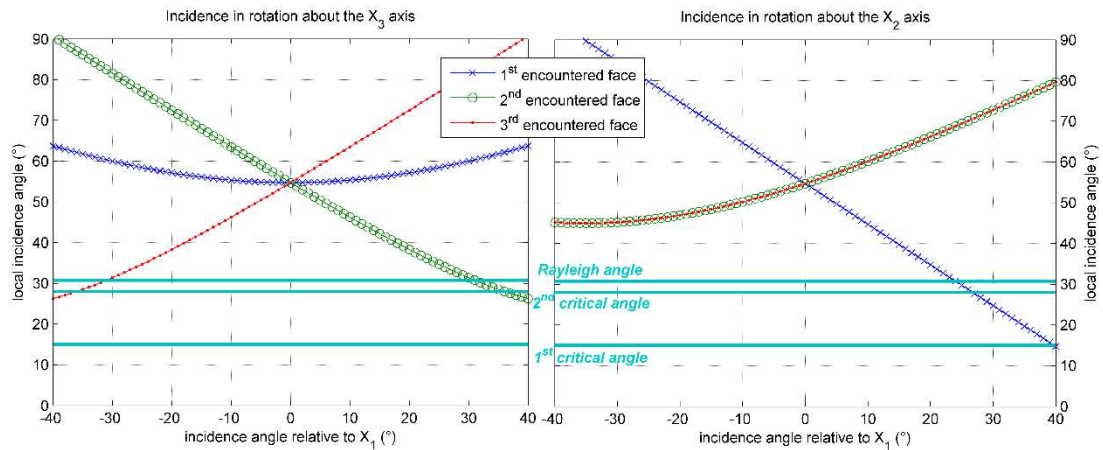
156

generation at $+24^\circ$. Beyond 26.8° a part of the incident energy is converted into

157

transversal waves in the steel.

158



159

160 Figure 4: Local incidence angles on the faces as a function of the initial incidence on CCR

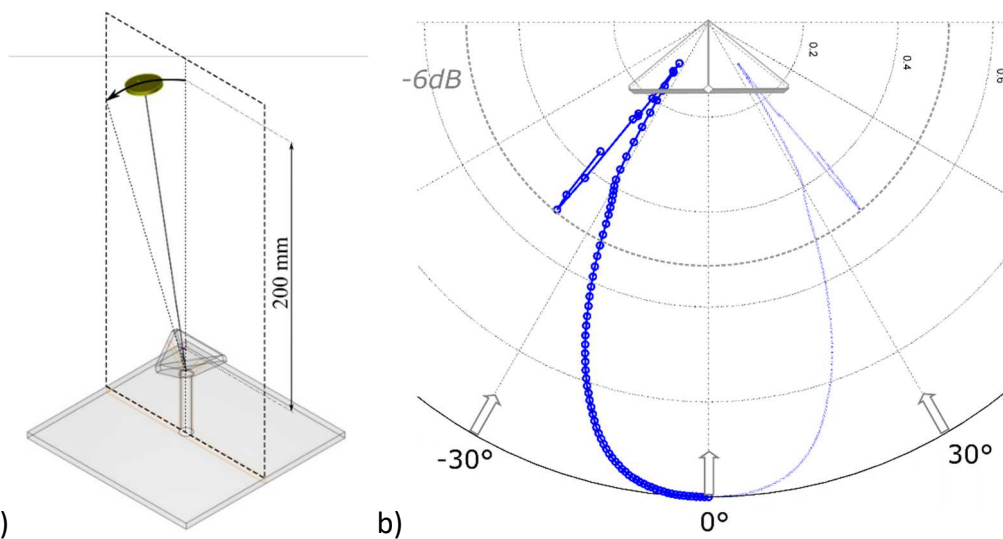
161

162 Three physical phenomena can reduce the amplitude of an arbitrary incident beam
 163 reflected from a CCR: 1) energy is converted into Rayleigh waves on one face, 2) the
 164 bulk reflection on one face is not total, which happens rather rarely in the case of steel
 165 in water, or 3) due to the finite dimensions of the CCR regardless of the beam
 166 diameter, a part of the energy is lost (the beam reaches ineffective areas and is
 167 reflected only once or twice, or a part of the incident beam strikes off the target area,
 168 or the reflected beam is shifted and therefore not entirely incident onto the
 169 transducer).

170

171 These phenomena are taken into account in our model using the CIVA expertise
 172 platform [21]. The CCR is imported from CAD software. A rotational scan in echo mode,
 173 centred on the inner CCR vertex, is computed (Figure 5.a), and the maximum reflected
 174 amplitude is plotted as a function of the incidence angle (Figure 5.b). The maximum
 175 reflected amplitude occurs at 0°, and it decreases slowly to reach a minimum at 36°.

176 In our case of computation of echoes from an immersed target, the CIVA model is
 177 invalid in the near vicinity of the longitudinal, transversal and Rayleigh critical angles.
 178 Indeed this model based on the Kirchhoff approximation [22,23] does not take into
 179 account head waves or Rayleigh waves and uses the plane wave reflection coefficient
 180 which is discontinuous in terms of magnitude or phase at these critical angles.
 181 Nevertheless, to simulate the response of flaws inside a solid specimen, CIVA mimics
 182 some head waves using its semi-analytical PTD (Physical Theory of Diffraction) model
 183 [24,25] and head/Rayleigh waves are henceforth correctly modelled in the CIVA hybrid
 184 ECHO model [26] using the Finite Elements method for flaw modelling.
 185 The CIVA result presented below will be confronted with experimentation in the next
 186 section.



187 a) b)
 188 Figure 5: Rotational scan about X_3 on CIVA: (a) configuration and (b) polar plot of the maximum
 189 received amplitude
 190

191 **3. Experimentations for various configurations**

192 Experimental measurements are performed in immersion, in water at room
193 temperature as a first step. The 43mm side stainless steel CCR is fixed on a plate (see
194 Figure 6). The ultrasonic emitter-receiver is a flat transducer, 1"-diameter, broadband,
195 centred at 2.25MHz. It is first set so as to be aligned with the inner vertex of the CCR at
196 0° incidence, at about 200mm from the cutting plane (X_2X_3) of the CCR.



197

198

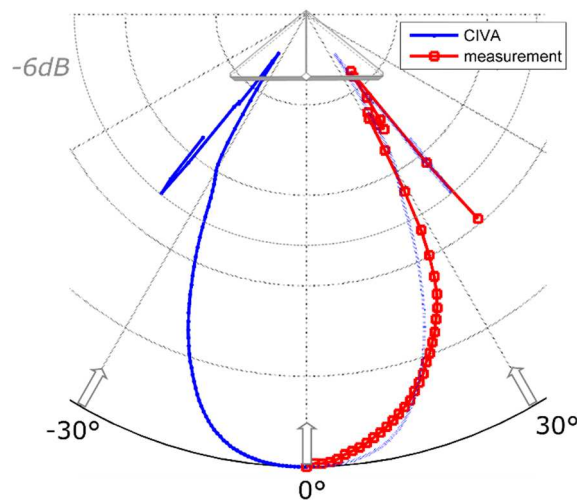
Figure 6: Stainless steel CCR and experimental setup

199

200 3.1. Validation of modelling

201 Bscan type acquisitions (1D scanning) are made along the X_2 axis for various incidence
202 angles of the beam, from 0° to 40° by steps of 1°. For each angle, the position at which
203 the incident beam is aligned with the vertex of the CCR corresponds to the maximum
204 reflected amplitude because there is no beam shift in this position (see Figure 2).

205



206 Figure 7: Polar plot of the maximum reflected amplitude: comparison between CIVA
207 simulations and experimentation

208
209 The maximum reflected amplitude value is then saved and plotted as a function of the
210 incidence angle in Figure 7 along with CIVA results. The agreement between simulation
211 and experimentation is quite good (except at the Rayleigh angle, as previously
212 explained). This will allow the use of CIVA in more complex cases as will be seen in
213 section 4.

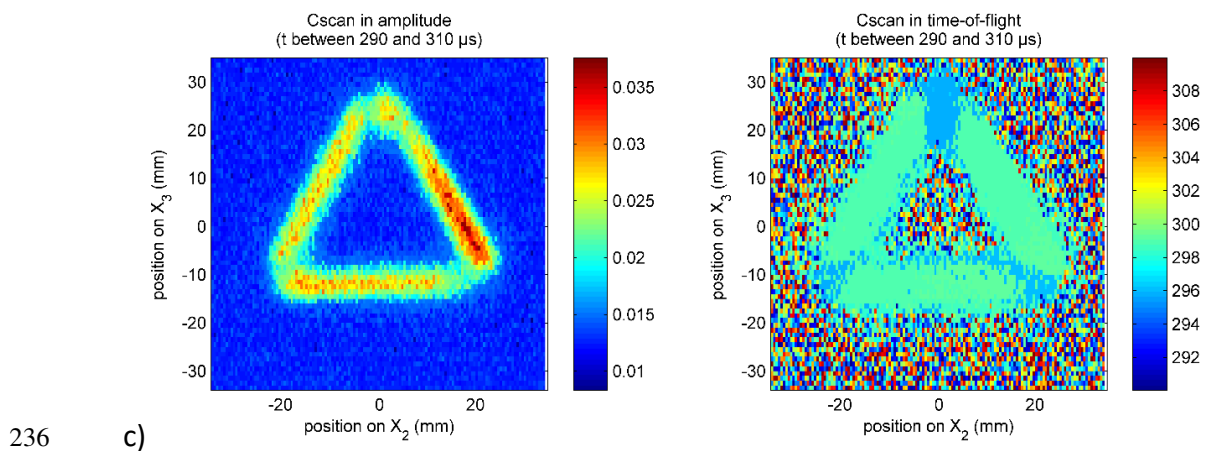
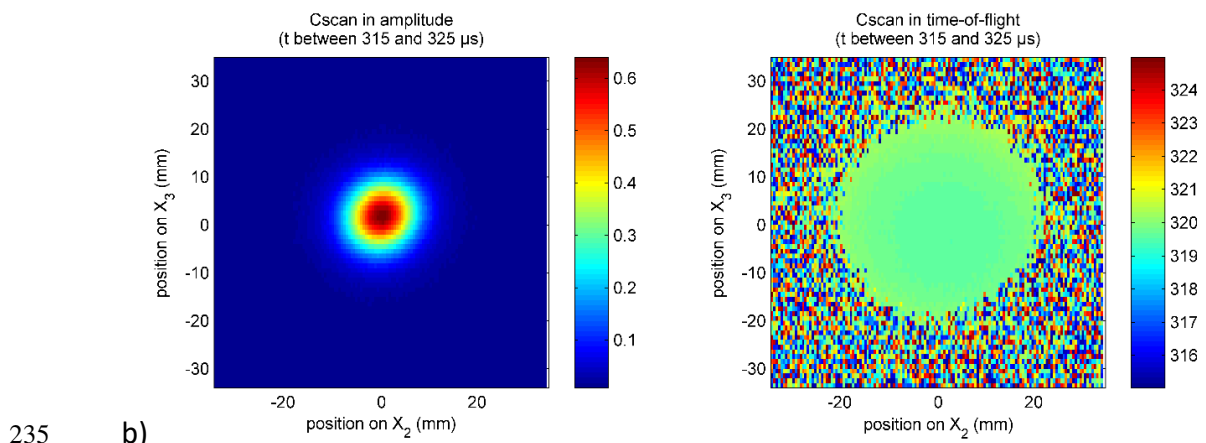
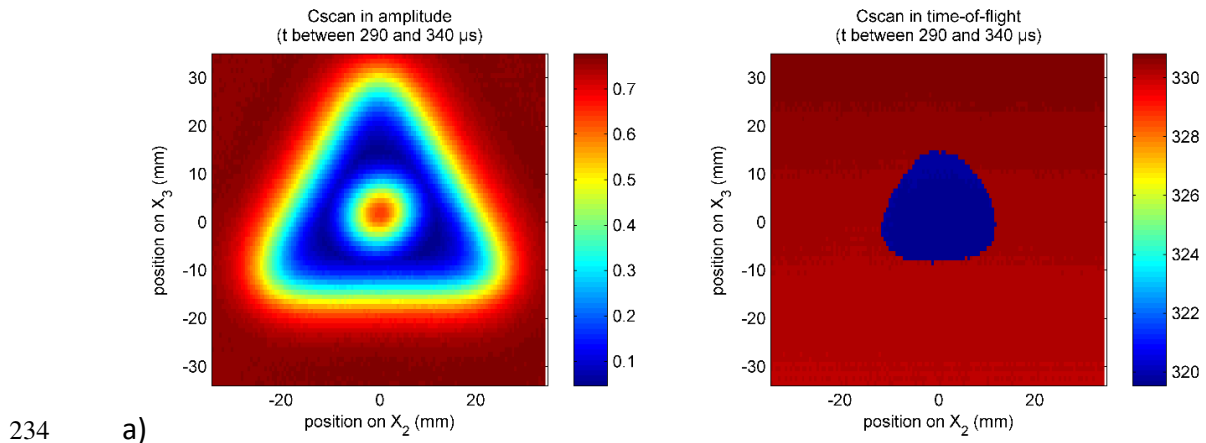
214

215 **3.2. Ultrasonic field reflected by a CCR alone**

216 Cscan type acquisitions (2D scanning) are performed over the area of the target,
217 parallel to (X_2X_3) , in order to analyse how it reflects the beam. Results at 0° incidence
218 (amplitude and time-of-flight) are displayed in Figure 8 for different choices of time-
219 domain windowing: a global window over all the received echoes (including those
220 reflected on the plate carrying the CCR), a window focusing around the time-of-flight
221 from the inner vertex of the CCR, and a window centred on the weak echoes coming
222 back from its edges.

223 Global reflections from the CCR and the supporting plate are plotted in Figure 8.a. The
224 amplitude image (left) shows up the position of the CCR on the plate and confirms that
225 the maximal amplitude from the CCR is registered when the acoustic beam is aligned
226 with the CCR inner vertex. As expected, the corresponding time-of-flight (right) is
227 constant and is equivalent to the direct path to the CCR inner vertex. The tightened
228 windowing on the CCR inner vertex echoes only (Figure 8.b) highlights the field sent
229 back by the CCR and the constant time-of-flight. This treatment makes it possible to

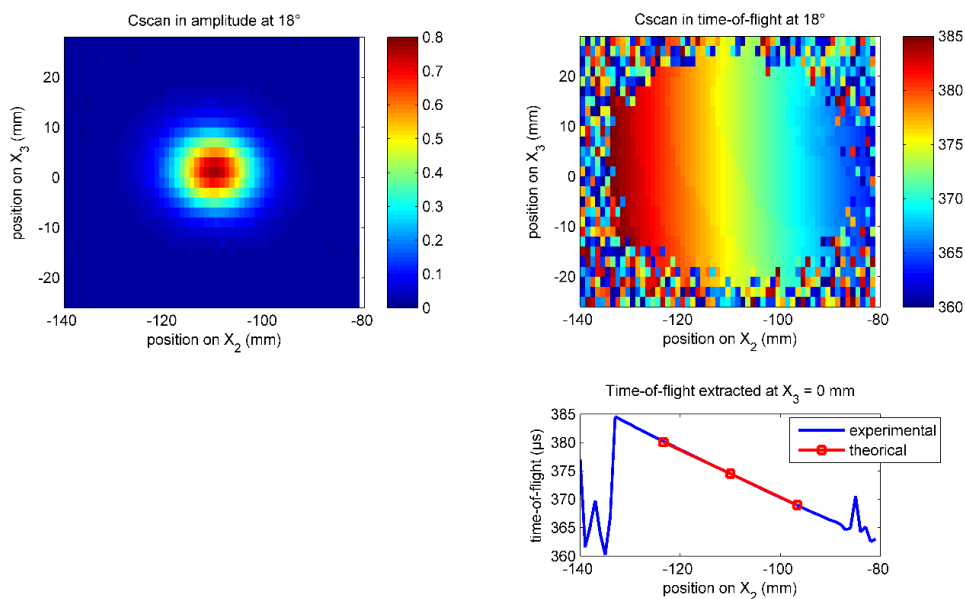
230 check the proper angular positioning of the target against the transducer. As for the
 231 last windowing (Figure 8.c) around echoes from the edges, its plot also confirms that
 232 the CCR is well positioned. However, the amplitudes are weak because the echoes are
 233 generated by scattering (or diffraction).



237 Figure 8: Normal incidence Cscan images in amplitude (left) and time-of-flight (right):
 238 a) global windowing, b) windowing on CCR inner corner echoes, and c) on CCR edges

239

240 These acquisitions are made for various incidence angles (transducer tilted around X_3 ,
 241 and scanned plane always parallel to (X_2X_3)). At oblique incidence, the time-of-flight is
 242 not constant anymore but evolves linearly as a function of the inclination of the
 243 transducer with respect to the scanned plane. The edges are not visible anymore. An
 244 example of such acquisition is shown in Figure 9. The maximum of amplitude is again
 245 reached when the acoustic beam is aligned with the CCR vertex.



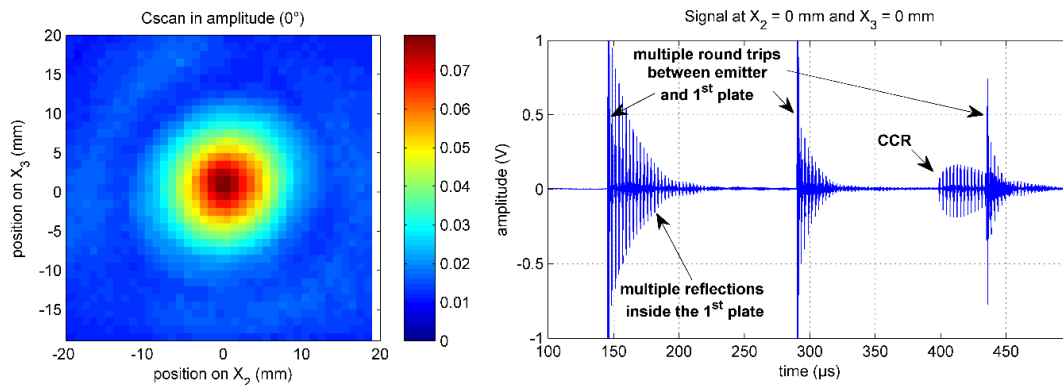
246

247 Figure 9: Example of Cscan acquisition at oblique incidence (18°):
 248 in amplitude (left) and in time-of-flight(right)

249

250 **3.3. Ultrasonic response of a CCR behind a plate**

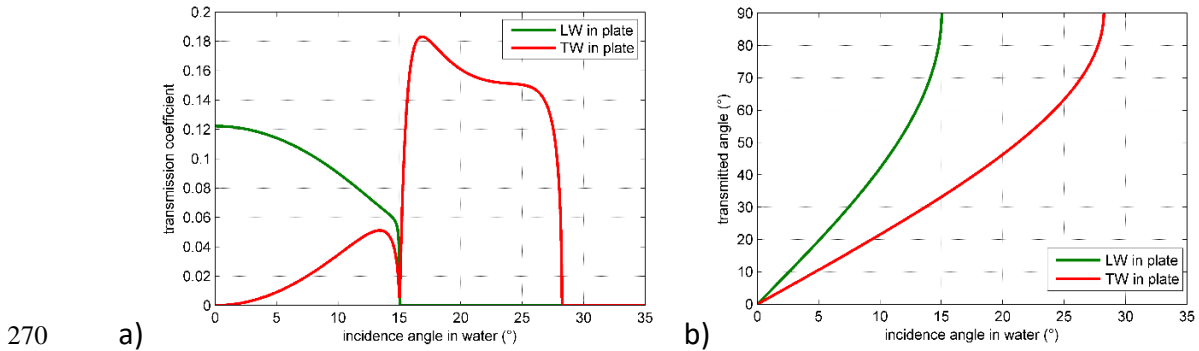
251 The purpose of this study is to perform telemetry on an object equipped with a CCR
 252 and located behind a plate. This 7.8 mm thick 316L stainless steel plate is positioned
 253 parallel to the scanning plane (X_2X_3).
 254 At normal incidence, each time the beam passes through the plate, it loses about 90%
 255 of its energy in reflection. However, the CCR is still well identified, as shown in Figure
 256 10, provided that its echoes do not overlap other echoes, such as those resulting from
 257 the multiple round trips between the transducer and the plate hiding the target.



258
 259 **Figure 10: Acquisition at normal incidence on the CCR behind a plate (Cscan with narrow time**
 260 **window focused on the first echo from the CCR, between 398.5 and 400 μ s)**

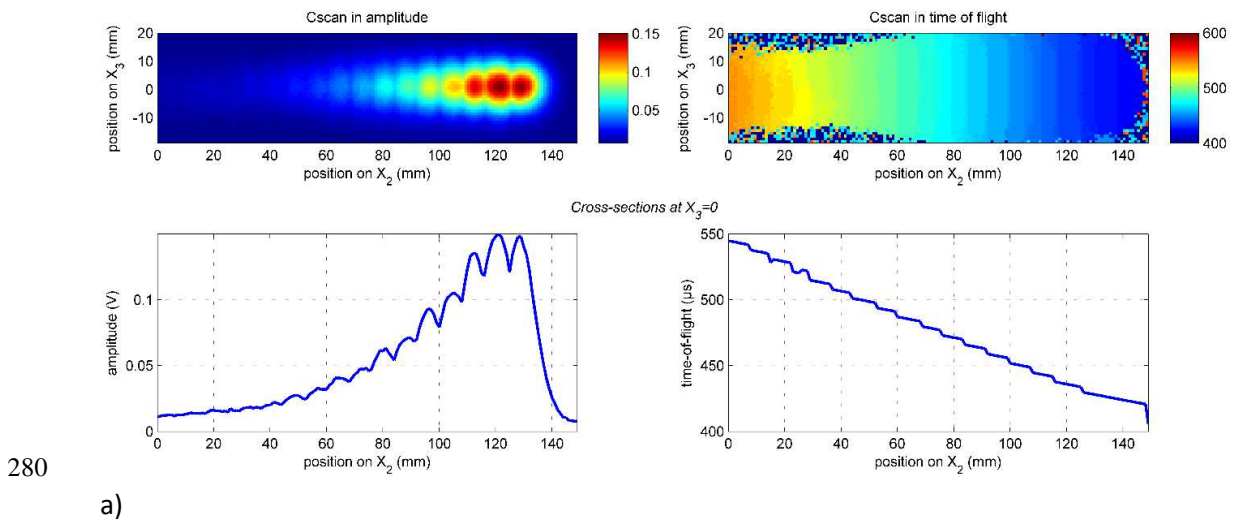
261
 262 Thus, positioning the transducer at oblique incidence will eliminate the problem of
 263 multiple echoes. Moreover, the energy transmitted through the first plate can be
 264 increased if the incidence angle is well chosen. Therefore, the transmission coefficients
 265 through a steel plate immersed in water are calculated and plotted in Figure 11. The
 266 maximum of transmission is reached for an incidence around 17° , generating shear
 267 waves at about 40° in the plate. We fix here the incidence angle at 18° , because it

268 corresponds to the maximum of transmission measured experimentally using our
 269 setup.



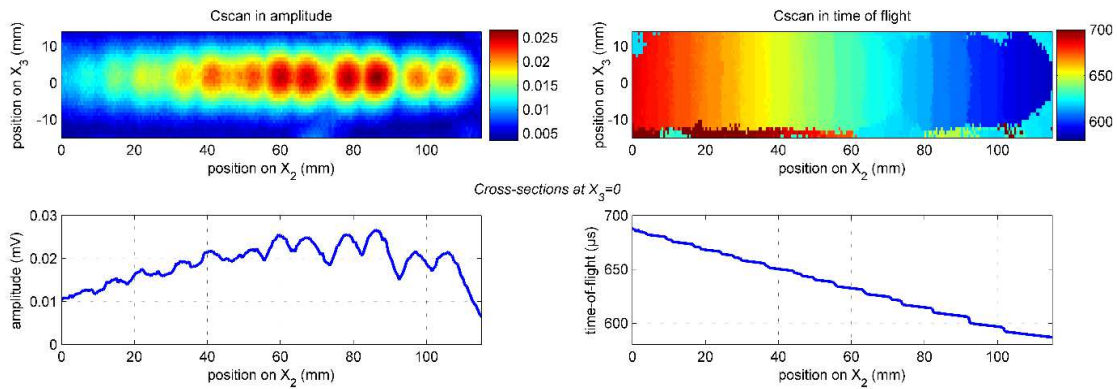
270 a)
 271 Figure 11: a) Transmission coefficient (separate contribution of longitudinal and transverse
 272 waves) through a steel plate immersed in water. b) Angle of propagation in the plate, as a
 273 function of the incidence angle

274 Resulting Cscan images (amplitude and time-of-flight) are plotted in the upper part of
 275 the Figure 12a. Again, the CCR is clearly identifiable and its ultrasonic image is visible
 276 with the repetition corresponding to the successive oblique reflections in the plate. As
 277 previously the time-of-flight evolves linearly with the displacement of the transducer,
 278 and exhibits steps corresponding to the increase in the number of round trips in the
 279 plate.



280

a)



281

b)

282

Figure 12: Acquisition at oblique incidence (18°) on the CCR (a) behind one plate, and (b)

283

behind two plates

284

285 Signal-to-noise ratios (SNR) were estimated on the acquired maximal signal in each

286 case (at normal and oblique incidence, alone and behind one and two screens), and

287 orders of magnitude are indicated in Table 1. The similar experiments conducted with

288 two screens (parallel plates) behind the target show that the target is not visible at

289 normal incidence but clearly identifiable at oblique incidence (see Figure 12b). All

290 these results clearly highlight how improved the back echo amplitude is when a CCR is

291 adding on the object of interest for detection purposes and telemetry applications.

292

Configuration	CCR alone (normal incidence)	CCR behind a plate, normal incidence	CCR behind a plate, oblique incidence (18°)	CCR behind two plates, normal incidence	CCR behind two plates, oblique incidence (18°)
SNR / inner corner	more than 100	about 8 (when non-	about 50	/	about 3

		overlapping)			
SNR / edges	about 2	/	/	/	/

293 Table 1: Signal-to-noise ratio **estimated** of the back echoes from the CCR according to
294 configuration

295

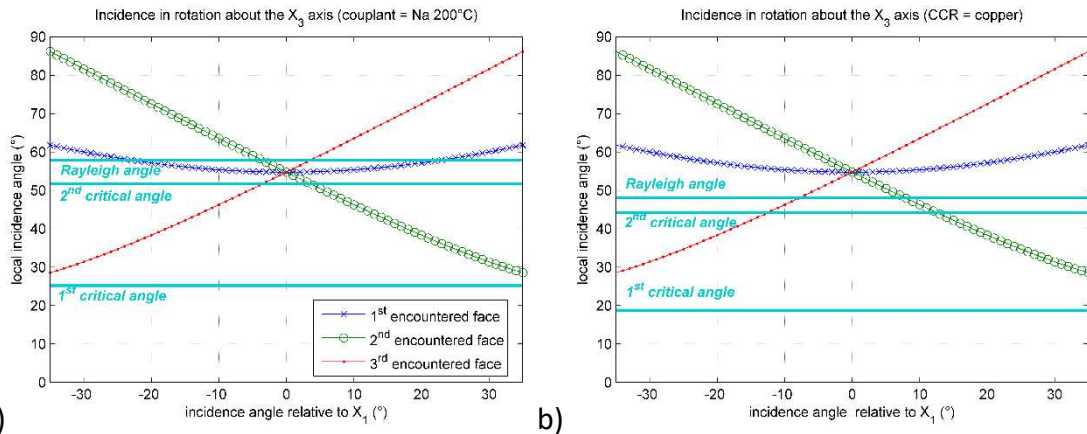
296 **4. Discussion on the change of materials: stainless steel in liquid sodium and copper**
297 **in water**

298 The case of a stainless steel CCR immersed in water is well adapted for detection and
299 telemetry operations in 2nd and 3rd generation nuclear power plants, but not for the 4th
300 generation reactors, in which liquid sodium is used as coolant [27]. Ultrasound
301 propagates faster in liquid sodium than in water (2470 m.s⁻¹ in sodium at 200°C [28]),
302 and this shifts all the critical angles up: the first and second critical angles are now
303 25.5° and 51.7°, and the Rayleigh angle is now 58°, instead of 14.9°, 28.2° and 31.5°.

304 This means that the angular range with total reflection on the CCR is null because no
305 incidence angle allows total reflection locally on all the three faces (cf. Figure 13.a).
306 This leads to a modelled reflection diagram (Figure 14.a) very different from the one
307 for the previous case. The global reflection is no longer maximal at normal incidence,
308 and in particular there are two local minima, at 0° and 10°.

309 Validation tests are not easy to perform in liquid sodium, as its handling requires
310 special safety precautions. That is why we searched for another (common) material
311 whose behaviour in water would be similar to that of steel in liquid sodium. Then we
312 calculated critical angles and reflection coefficients for various materials and chose
313 copper. In fact, the critical angles for copper in water are close to what they are for

314 steel in liquid sodium (18.5°, 40.4° and 43.9°), as shown in Figure 13.b. A narrow range
 315 of incidence angle leading to total reflection remains, from about -8° to 8°.



316 a) b)
 317 Figure 13: Local incidence angles on the faces as a function of initial incidence on CCR

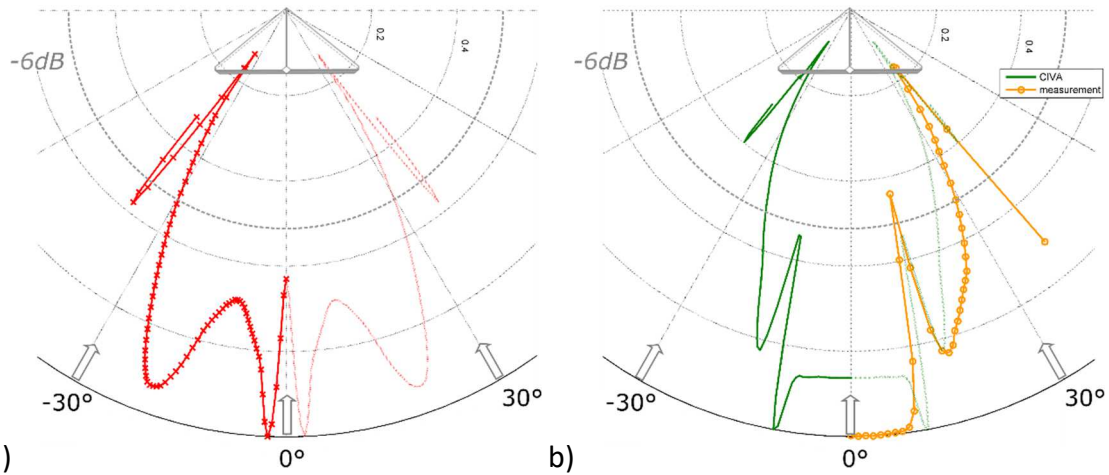
318 (a) for the case of steel in liquid sodium, and (b) for the case of copper in water

319

320 The modelled reflection diagram is plotted in Figure 14.b (green curve). This diagram
 321 also exhibits a particular behaviour, very different from the case of steel in water. Local
 322 minima are found at 0° and 13°, and maxima at 11° and 15°. A CCR made of copper
 323 was produced and studied by ultrasound using the same procedure as previously
 324 (section 3.1). Experimental results (orange curve) are represented along with modelled
 325 results. They are in quite good agreement: the local minimum at 13° and the local
 326 maximum at 15° are well recovered. However, the maximum at 11° is experimentally
 327 not observed: this incidence angle being in the vicinity of the second critical angle
 328 locally on the faces (see Figure 13b), it is outside the limits of validity of CIVA, as
 329 mentioned earlier.

330 These results show that one can always find large angular ranges where high energy is
 331 sent back to the transducer. These ranges are reduced compared with those for the

332 case of steel in water. The knowledge of the media considered for both CCR and
333 surrounding fluid is thus predominant for the overall acoustic reflexion from an
334 immersed CCR.
335



337 Figure 14: Polar plot of the (normalized) maximum reflected amplitude computed with CIVA
338 (a) for the case of steel in liquid sodium, and (b) for the case of copper in water,
339 with comparison with experimentation

340

341 5. Conclusion

342 The Corner Cube Retroreflector is studied in the framework of ultrasonic telemetry. Its
343 ability to fully reverse an incoming wave in the same direction after three reflections is
344 key to performing telemetry operations (i.e. position measurement) at non-normal
345 incidence (within the angular range of $\pm 30^\circ$) on an immersed object. To demonstrate
346 the feasibility of telemetry in various configurations, theoretical developments,
347 modelling and experimentations have been made. It appears possible to locate the
348 steel CCR even behind one and two plates immersed in water, because the amplitude
349 of the back echo is clearly sufficient: SNR equals about 100 for CCR alone, about 50 for

350 CCR behind a plate, and about 3 for CCR behind two plates. When the CCR is behind
351 screens, the oblique incidence is chosen to maximize the energy transmitted through
352 the screens.

353 We have also underlined the relevance of considering the acoustic properties of each
354 medium (CCR material and surrounding fluid). They affect the local coefficients of
355 reflection on the CCR faces and thus the global reflection of the target.

356 Future work will deal with the ability to measure not only the position but also the tilt
357 of a new compact target.

358

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366

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450 **Figures Captions**

451 Figure 1: a) Definition of coordinate system, and normal incidence, referred to as “0°” and b)
452 examples of successive reflections of a ray on the faces with various incidences

453 Figure 2: Calculation of cylindrical beam reflection as a function of the angle of incidence and
454 the alignment with the inner vertex (the total length of the arrows of one ray travel
455 corresponds to a single time-of-flight)

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457 of the incidence angle

458 Figure 4: Local incidence angles on the faces as a function of the initial incidence on CCR

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465 windowing, b) windowing on CCR inner corner echoes, and c) on CCR edges

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467 time-of-flight(right)

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469 window focused on the first echo from the CCR, between 398.5 and 400 μ s)

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471 waves) through a steel plate immersed in water. b) Angle of propagation in the plate, as a
472 function of the incidence angle

473 Figure 12: Acquisition at oblique incidence (18°) on the CCR behind a plate

474 Figure 13: Local incidence angles on the faces as a function of initial incidence on CCR (a) for
475 the case of steel in liquid sodium, and(b) for the case of copper in water

476 Figure 14: Polar plot of the maximum reflected amplitude computed with CIVA (a) for the case
477 of steel in liquid sodium, and (b) for the case of copper in water, with comparison with
478 experimentation

479

480 **Tables Captions**

481 Table 1: Signal-to-noise ratio of the back echoes from the CCR according to configuration

482