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A Mountaineering Strategy to Excited States: Highly-Accurate Oscillator Strengths and Dipole Moments of Small Molecules

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Abstract

This work presents a series of highly-accurate excited-state properties obtained using high-order coupled-cluster (CC) calculations performed with a series of diffuse containing basis sets, as well as extensive comparisons with experimental values. Indeed, we have computed both the main ground-to-excited transition property, the oscillator strength, as well as the ground- and excited-state dipole moments, considering thirteen small molecules (hydridoboron, hydrogen chloride, water, hydrogen sulfide, boron fluoride, carbon monoxide, dinitrogen, ethylene, formaldehyde, thioformaldehyde, nitroxyl, fluorocarbene, and silylidene). We systematically include corrections up to the quintuple (CCSDTQP) in the CC expansion and extrapolate to the complete basis set limit. When comparisons with experimental measurements are possible, that is, when a number of consistent experimental data can be found, theory typically provides values falling within the experimental error bar for the excited-state properties. Besides completing our previous studies focussed on transition energies (*J. Chem. Theory Comput.* **14** (2018) 4360–4379, *ibid.* **15** (2019) 1939–1956, *ibid.* **16** (2020) 1711–1741, and *ibid.* **16** (2020) 3720–3736), this work also provides ultra-accurate dipoles and oscillator strengths that could be employed for future theoretical benchmarks.

1. INTRODUCTION

In the formidable quest aiming at reaching high accuracy in the modeling of electronically excited states (ESs), the primary focus has been set on vertical excitation energies,[1] defined as the difference in total energies between a given ES and its corresponding ground state (GS) at fixed geometry (typically the ground-state equilibrium geometry). The main reasons for this choice are, on the one hand, the availability of many theoretical models for computing total ES energies, and, on the other hand, the fact that obtaining accurate transition energies is generally viewed as a prerequisite for further ES investigations. Along the years, more and more accurate vertical transition energies have become available. Probably the most illustrative example is provided by the well-known Thiel set,[2–4] encompassing 223 values obtained at the CC3 and CASPT2 levels, as well as our recent efforts to obtain chemically-accurate vertical energies within ±0.03 eV of the full configuration interaction (FCI) limit for more than 400 ES.[5–8] Whilst such sets are obviously useful to benchmark lower-order methods, they remain nevertheless intrinsically limited by two factors. First, vertical transition energies remain inaccessible experimentally (in the vast majority of the cases), preventing direct comparisons with measurements. To circumvent this issue, several groups have turned their attention to 0-0 energies,[9–23] defined as the differences between the ES and GS energies determined at their respective minimum and corrected for zero-point vibrational effects, because 0-0 energies allow straightforward theory-experiment comparisons. To compute 0-0 energies, one must however determine ES geometries and vibrations, which limits the number of methods that can be applied. This explains why compromise "hybrid" protocols employing different levels of theory for the structural and energetic parameters are popular in this particular field.[12,18,19,21] Interestingly, the accuracy of the underlying geometries has been shown to be rather irrelevant,[14,22] indicating that benchmarks of 0-0 energies still mainly assess the quality of the (adiabatic) transition energies. Second, vertical transition energies do provide a very partial characterization of the ES as one typically needs to determine transition probabilities (oscillator strengths, f) as well as ES properties (such as structures and dipole moments, μ) to attain a better grasp of the photophysics. On a more theoretical viewpoint, it is also known that a method can provide accurate transition energies while failing to deliver accurate ES properties,[24–26] indicating that benchmarks relying on transition energies as unique gauge can lead to incomplete and/or erroneous conclusions regarding the strengths and weaknesses of a given theoretical model.

In the above-defined framework, let us now summarize the efforts that have been made to define accurate reference ES properties for significant sets of compounds.

First, for ES geometries, which are not our focus here, several ensembles of structures have been reported at CASPT2,[27] SAC-CI,[28,29] Mk-MRCCSD,[30] QMC,[31] and

CC2[32] levels. To the best of our knowledge, the most extensive dataset of accurate ES structural parameters remains the one defined by some of us during the last few years: it contains bond lengths, valence angles, and torsion angles determined at the CC3 and CASPT2 levels with *aug*-cc-pVTZ for several dozens of small organic compounds.[22,24,33–35] In this framework, it is also worth mentioning the works of Szalay's group on the shape of potential energy surfaces, in which CCSDT references are defined,[25,26,36] as well as the publications of Olivucci's group focussing on the topology of conical intersections, in which the results of high-level multi-reference calculations are available.[37,38]

Second, for transition properties, the most accurate reference values we are aware of (for a significant set of transitions) are: i) the CC3/TZVP oscillator strengths obtained by Kannar and Szalay[39] for the Thiel set (this work also includes 15 CCSDT/TZVP oscillator strengths) that can be compared to CASPT2/TZVP values obtained earlier for the same set;[2] ii) the numerous CC3/*aug*-cc-pVTZ oscillator strengths determined to complete our FCI energy calculations;[5,7,8] and iii) the transition dipole moments computed at the CC3 or ADC(3) levels for 15 molecules by Robinson.[40] There are also earlier sets of oscillator strengths computed at lower levels of theory, *e.g.*, CCSD,[41] SAC-CI,[42] and CC2.[43,44] The typical error associated with these approaches (with respect to FCI) remains unknown at this stage. To the best of our knowledge, a significant set of FCI-quality oscillator strengths has yet to be published.

Third, theoretical datasets for accurate ES dipoles are apparently even scarcer. The most advanced sets we are aware of contain: i) CASPT2/TZVP,[2] and CC2/*aug*-cc-pVTZ reference values[43] determined for Thiel's set; ii) CC2/*aug*-cc-pVQZ results for small- and medium-sized molecules;[45] and iii) CC2/*aug*-cc-pVTZ data for "real-life" organic dyes.[46] Comparisons with experimental ES dipoles performed in Ref. 45 yield an error of ca. 0.2 Debye (D) for both ADC(2) and CC2, a value that can be viewed as acceptable, but is nevertheless far from chemical accuracy and not on par with the precision provided by state-of-the-art approaches for GS dipole moments.[47] Of course, one can find specific works focussing on a small number of ES dipoles and using high level(s) of theory (see our Result Section for references). However, such specific studies preclude valuable statistical conclusions.

Globally, these previous studies have demonstrated that both f and μ^{ES} are much more basis set sensitive than transition energies and geometries. Indeed, a specific challenge comes from intensity-borrowing effects that can vastly change the properties of two close-lying ES of the same symmetry while their energies remain almost unaffected. When a change of basis set slightly tunes the energy gap between two ESs, it might simultaneously drastically affect the properties. This highlights that properties are much harder to accurately estimate than energies, not only for implementation or computational reasons, but also due to more fundamental aspects.

Another important fact is that, in contrast to vertical transition energies, both oscillator strengths and ES dipole moments are accessible experimentally, so comparisons between theoretical and experimental measurements are, in principle, possible. However, as we pointed out previously for geometries,[33] such comparisons are generally far from straightforward. For the oscillator strengths, the two main measurement techniques are electron impact and optical spectroscopies. As nicely summarized elsewhere,[48] the former typically requires extrapolations of the cross-sections measured at various momenta/angles (such extrapolations are not error-free), whereas the latter can be plagued by saturation and interaction effects (yielding to underestimation of the actual f value). In any case, the so-called "electronic" or "optical" f, which is of interest here, is not directly measurable, and post-processing of the experimental raw data (lifetimes or cross sections) is needed to access its "experimental" value.[48,49] The measurements of ES dipoles are also cumbersome.[50] The typical strategy is to investigate solvatofluorochromism, that is, to measure the shift of the emission wavelength in a series of solvents of various polarities, and to fit the results with the Lippert-Mataga equation within an Onsager-like interaction model. Such approach is obviously not a direct gas-phase measurement and comes with significant assumptions, making the final error often too large to allow benchmarking.[51] Alternatively, one can also measure μ^{ES} directly in gas phase by studying how external electric fields tune the position and shape of the vibronic peaks (the so-called Stark effect).[50,51] Only the latter approach can be considered as sufficiently robust for reference purposes. Yet, such Stark effect measurements do typically provide an "adiabatic" version of μ^{ES}, that is, the ES dipole measured at the ES equilibrium geometry. If a non-trifling structural reorganization takes place after absorption, such as in formaldehyde, this means that the value of μ^{ES} computed at the GS equilibrium structure has again no experimental counterpart. In addition, the direction of the ES dipole moment cannot be determined directly from the measurements of the Stark effect.

For all these reasons, it is of interest to define a set of coherent near-FCI oscillator strengths and excited-state dipole moments. Indeed, the oscillator strengths are directly related to the transition dipole moments, whereas the dipole moments can be seen as a measure of the quality of the total (GS and ES) densities, so that these data go beyond the "simple" characterization of vertical transition energies. This contribution therefore aims at tackling this ambitious objective for a set of small molecules similar to the one treated in our original *Mountaineering* paper.[5] To this end, we take advantage of the CC hierarchy, going from CCSD to CCSDTQP for both properties, in combination with increasingly large atomic basis sets including one or two sets of diffuse basis functions. Whilst such calculations provide "definite answers", they remain at the limit of today's computational capabilities and are achievable for compact molecules only, which stands as a clear limit of the present contribution. Nevertheless, providing highly-accurate numbers for properties directly related to the quality of the transition and the ES density is, we believe, useful. Very recently, Hait and Head-Gordon have used, in a density-functional theory (DFT) context, the GS dipole moment as a metric to estimate the quality of the GS density given by many exchange-correlation functionals.[47] As these authors nicely stated: the dipole *"is perhaps the simplest observable that captures errors in the underlying density, (...*

and is*) a relevant density derived quantity to examine for DFA testing and development"*. The extension to ESs (and time-dependent DFT) is obviously natural, but such task is difficult due to the lack of definite reference values. In this framework, we also note some very recent efforts for computing ES first-order properties with the many-body expansion FCI (MBE-FCI) approach.[52]

2. COMPUTATIONAL METHODS

2.1 Geometries and basis sets

All our geometries are obtained at the CC3/*aug*-cc-pVTZ[53,54] without using the frozen-core (FC) approximation (*i.e*, we correlate all the electrons). These geometries are given in the Supporting Information (SI). Note that several structures come from previous works,[5,8] and we used this level of theory here for consistency. New optimizations have been achieved with the DALTON 2017[55] and CFOUR 2.1[56] codes applying default parameters in both cases.

As in our previous works,[5–8] we consider the diffuse-containing Pople's 6-31+G(d) and Dunning's (d-)*aug*-cc-pVXZ (X = D, T, Q, and 5) atomic basis sets in all ES calculations.

In contrast with our previous studies, in which the complete basis set limit (CBS) could be obtained by a brute-force approach, f and $\mu^{\rm ES}$ converge, in some cases, slower than the energy with respect to the basis set size. Therefore, we have performed CBS extrapolation by applying the well-known Helgaker formula,[57]

$$P_{\rm CBS} = \frac{P_{X-1}\,(X-1)^3 - P_X X^3}{(X-1)^3 - X^3}, \qquad (1)$$

in which X equals 2, 3, 4, . . . for D, T, Q, . . . in the Dunning series and P is the property under investigation. Interestingly, this approach was successfully used for GS dipole moments.[58] In practice, we performed four extrapolations using both the CCSD and CCSDT results obtained with the pairs of largest singly- and doubly-augmented basis sets accessible of at least triple-ζ size, i.e., i) CCSD/*aug*-cc-pVQZ and *aug*-cc-pV5Z; CCSDT/*aug*-cc-pVTZ and *aug*-cc-pVQZ; CCSD/d-*aug*-cc-pVTZ and d-*aug*-cc-pVQZ; and CCSDT/d-*aug*-cc-pVTZ and d-*aug*-cc-pVQZ for molecules in Sections 3.1–3.5; ii) CCSD/*aug*-cc-pVTZ and *aug*-cc-pVQZ; and CCSD/d-*aug*-cc-pVTZ and d-*aug*-cc-pVQZ for molecules in Sections 3.6–3.9. In such a way, we can estimate the extrapolation error, and provide error bars for the CBS values, although such error bar is likely underestimated for the latter set of compounds. These CCSD and/or CCSDT CBS values are next used to correct the properties obtained at higher levels (*e.g.*, CCSDTQ) with a finite basis using the approach described below.

2.2 Reference calculations

We have chosen to use the MRCC (2017 and 2019) program,[59,60] for performing our CC calculations, as this code allows to set up an arbitrary CC expansion order. We therefore use the CCSD,[61–65] CCSDT,[66–70] CCSDTQ,[71–74] and CCSDTQP[72–74] hierarchy for energies, oscillator strengths, GS and ES dipoles. All these values have been obtained within the FC approximation. The interested reader may find discussions about the impact of this approximation and the importance of core correlation functions for transition energies in some of our previous works.[22,75] For CCSD, we performed several test calculations with GAUSSIAN 16,[76] Q-CHEM 5.2,[77] DALTON 2017,[55] and e^T 1.0[78] and we could not detect any significant discrepancy with respect to the MRCC results. At this stage, it is important to stress that all these calculations rely on the so-called linear-response (LR) formalism,[73,79,80] so that while the same transition energies would be obtained with the equation-of-motion (EOM) approach, the properties would be different. However, it is known that the two formalisms become equivalent when the CC wave function becomes exact. As we strive here to be as close as possible from the FCI limit, we trust that our theoretical best estimates (TBEs) are not significantly affected by the selection of the LR implementation. At the CCSD level, we also provide a comparison between the EOM and LR oscillator strength values obtained with various codes in Table S1 in the SI. The differences found between the EOM and LR formalisms are very small. The interested reader may also find extensive comparisons of oscillator strengths determined within the two formalisms elsewhere.[81] We also note that all our oscillator strengths are given in the length gauge, the most commonly applied gauge, but again this choice is likely irrelevant when one is targeting near-exact values. The interested reader can find comparisons between CC oscillator strengths determined with the length, velocity, and mixed length-velocity gauges on small compounds elsewhere.[82] As expected the impact of the gauge was found to decrease when increasing the order of the CC expansion, especially when triples are included. Finally, we report in the Tables below the so-called *orbital-relaxed* dipoles, which are more accurate than the so-called *orbital-unrelaxed* dipoles in which the impact of the external field on the orbitals is neglected. Details on various approaches and their implementations for correlated first-order properties can be found elsewhere.[80,83,84]

Beyond the basis set extrapolation discussed above, we also define TBEs in the following. To this end, we use an incremental strategy for the transition energies and dipoles, *e.g.*, at the *aug*-cc-pVTZ level,

$$\begin{aligned} P({\rm TBE}) &= P({\rm Low/AVTZ}) \\ &\quad + P({\rm High/AVDZ}) - P({\rm Low/AVDZ}), \end{aligned} \qquad (2)$$

where Low and High denote, *e.g.*, CCSDTQ and CCSDTQP (see footnotes in the Tables for specific details). For oscillator strengths, we applied the corresponding multiplicative approach, *e.g.*,

$$f({\rm TBE}) = f({\rm Low/AVTZ})\frac{f({\rm High/AVDZ})}{f({\rm Low/AVDZ})}. \qquad (3)$$

Such incremental strategy that used a double-ζ result to estimate "Q" or "P" effects is commonly employed in the CC literature.[73,85–89]

3. RESULTS AND DISCUSSION

Below, we discuss individual molecules going up on size progressively. Concerning literature references, we do not

Table 1: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, oscillator strengths f, and excited-state dipole moments, $\mu^{\rm ES}_{\rm vert}$ and $\mu^{\rm ES}_{\rm adia}$, determined for BH (GS and ES geometries) and HCl (GS geometry). Transition energies are in eV and dipoles in D.

		BH					HCl			
		$^1\Sigma^+$	$^1\Pi$ (Val)				$^1\Sigma^+$	$^1\Pi$ (Val)		
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$	$\mu^{\rm ES}_{\rm adia}$	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$
aug-cc-pVDZ	CCSD	1.389	2.970	0.051	0.530	0.513	1.147	7.862	0.066	-2.773
	CCSDT	1.371	2.946	0.049	0.543	0.527	1.131	7.815	0.065	-2.745
	CCSDTQ	1.370	2.947	0.049	0.545	0.528	1.130	7.822	0.065	-2.728
	CCSDTQP	1.370	2.947	0.049	0.545	0.528	1.130	7.823	0.065	-2.727
aug-cc-pVTZ	CCSD	1.433	2.928	0.050	0.550	0.534	1.097	7.906	0.056	-2.526
	CCSDT	1.410	2.900	0.048	0.558	0.541	1.085	7.834	0.055	-2.515
	CCSDTQ	1.409	2.901	0.048	0.559	0.542	1.084	7.837	0.055	-2.502
aug-cc-pVQZ	CCSD	1.440	2.918	0.050	0.555	0.538	1.111	7.954	0.051	-2.410
	CCSDT	1.416	2.890	0.048	0.561	0.544	1.098	7.880	0.050	-2.410
aug-cc-pV5Z	CCSD	1.443	2.915	0.050	0.556	0.540	1.109	7.961	0.048	-2.336
d-*aug*-cc-pVDZ	CCSD	1.388	2.969	0.051	0.519	0.503	1.135	7.836	0.064	-2.697
	CCSDT	1.370	2.945	0.049	0.533	0.516	1.119	7.787	0.063	-2.670
d-*aug*-cc-pVTZ	CCSD	1.432	2.927	0.049	0.550	0.534	1.096	7.894	0.055	-2.491
	CCSDT	1.409	2.900	0.048	0.558	0.541	1.084	7.822	0.054	-2.480
d-*aug*-cc-pVQZ	CCSD	1.440	2.918	0.050	0.555	0.538	1.111	7.949	0.050	-2.405
	CCSDT	1.416	2.890	0.048	0.561	0.544	1.098	7.876	0.050	-2.406
aug-cc-pVTZ	TBE[a]	1.409	2.901	0.048	0.559	0.542	1.084	7.837	0.055	-2.501
CBS	TBE[b]	1.42±0.00	2.88±0.01	0.048±0.001	0.57±0.01	0.55±0.01	1.11±0.01	7.91±0.01	0.046±0.001	-2.32±0.01
Lit.	Th.	1.425[c]	2.944[d]			0.57[e]	1.106[f]	7.94[g]	0.081[g]	
	Exp.	1.27±0.21[h]		0.044[i]		0.58±0.04[h]		8.05[j]	0.051[j]	
									0.042±0.004[k]	
									0.052±0.006[l]	

[a] Computed using CCSDTQ/*aug*-cc-pVTZ values and CCSDTQP/*aug*-cc-pVDZ corrections. Note that CCSDTQ is equivalent to FCI for BH; [b] See Computational Methods section; [c] Average between the MR-ACPF/*aug*-cc-pCV7Z(i) values obtained for r = 1.220 and 1.225 Å in Ref. 90; [d] FCI/*aug*-cc-pVDZ value of Ref. 54; [e] CC2/*aug*-cc-pVQZ result from Ref. 45; [f] CCSD(T)/CBS value from Ref. 47; [g] CISDTQ/*aug*-cc-pCVQZ value from Ref. 91. A f value of 0.071 is also reported in Table II of the same work; [h] Stark (emission) measurements of Ref. 92; [i] f_{00} obtained from laser-induced fluorescence in Ref. 93. There is no major contributions from other bands according to this work (see references therein for previous experimental values). A slightly older experiment (Ref. 94) reports a smaller estimate of 0.045±0.02; [j] Absorption values from Ref. 95 (ΔE corresponds to the maximum of absorption); [k] EELS value from Ref. 96; [l] HR-EELS value from Ref. 97.

intend to provide an exhaustive list of all previous works for each system considered here, but rather to highlight the studies and comparisons that we have found valuable for the present work.

3.1 BH and HCl

Let us start by a tiny compound, BH. Our results are listed in Table 1 and although the size of this molecule seems ridiculous (only 4 valence electrons), some valuable conclusions can be obtained. We note that $\mu^{\rm GS}$ is slightly too large with CCSD (irrespective of the basis set) but the CC convergence is fast and CCSDT is obviously sufficient. Our TBE/CBS of 1.42 D for $\mu^{\rm GS}$ is equivalent to a recent ultra-accurate estimate,[90] and also falls within the error bar of the only available experimental value we are aware of: 1.27±0.21 D.[92] Such large error bar is explainable: the experiment relied on an analysis of the emission spectrum.[92] It is quite obvious that the theoretical estimate is more trustworthy in this specific case, indicating that previous error analyses based on the 1.27 D value likely significantly exaggerated the ADC(2) and CC2 overestimations,[45] but underestimated the QMC error[98] for $\mu^{\rm GS}$. For the lowest transition energy in BH, our values fit with previous FCI calculations,[54] and one again notices rather quick convergence of $\Delta E_{\rm vert}$ with respect to both basis set size and CC order. For the oscillator strength, there is also an astonishing stability of the values, as the considered ES is well separated from higher-lying ones of the same spatial symmetry. Our TBE is close to the most recent measurement (of f_{00}) we could find.[93] For $\mu^{\rm ES}$, one notes a very small decrease of the amplitude between the GS and ES geometries, which is a logical consequence of the tiny geometrical relaxation (+0.018 Å), and our TBE is within the rather small experimental error bar (see bottom of Table 1). Finally, as can be seen in Table S2, these results are not affected by the FC approximation, e.g., the difference between CCSDTQ(FC)/*aug*-cc-pVTZ and CCSDTQ(Full)/*aug*-cc-pCVTZ is 0.001 for f and 0.002 D for the dipole moments.

For HCl, the GS dipole moment does not cause any specific challenge and our TBE is equivalent to the one recently reported by the Head-Gordon group.[47] As expected,[5] CCSDTQ provides converged $\Delta E_{\rm vert}$ and this holds for both f and $\mu^{\rm ES}$. When increasing the size of the basis set, one sees a significant decrease of the oscillator strength (ca. -25% from *aug*-cc-pVDZ to *aug*-cc-pV5Z) and of the ES dipole (-16% for the same basis pair), whereas the transition energy varies by 1% only. Nevertheless, the CBS extrapolations are stable for all investigated properties. For HCl, it is difficult to obtain a very accurate experimental f value, in part due to the mixing of the $^1\Pi$ and $^3\Pi$ states.[95] There is therefore a

broad range of measured values (see bottom of Table 1 as well as Table 2 in Ref. 97), and our TBE/CBS of 0.046 is compatible with the two most recent experiments. In contrast, we could not find any experimental μ^{ES} estimate, which is a logical consequence of the dissociative character of the lowest singlet ES of HCl.

3.2 H_2O and H_2S

The results obtained for water and hydrogen sulfide are listed in Tables 2 and 3, respectively. The CCSDT/*aug*-cc-pVTZ estimate of the ground-state dipole of water is already within 0.01 D of both previous CASPT2[99] and experimental values.[101] For all transition energies, as discussed in our earlier work,[5] the calculations are converged with CCSDTQ, but one needs a rather large basis set (especially in terms of diffuse functions) to be chemically accurate, which is quite usual for Rydberg transitions in small compounds. Nevertheless, CCSDT/*aug*-cc-pVTZ delivers ΔE_{vert} with an error of 1–3% only as compared to the most accurate estimates (see bottom of Table 2) for all three transitions. For the lowest B_1 excitation, the oscillator strength f varies rather mildly with the selected level of theory and basis set, although one notices a general decreasing trend when improving the method. Our TBE of 0.052±0.001 falls in the error bar of a recent experiment,[103] and is only slightly larger than the previous most accurate TBE we are aware of.[48] More exhaustive lists of additional theoretical and experimental values can be found in Table IV of Ref. 103 and Table 6 of Ref. 48 for the lowest transition. The reported f values in these tables are in the range 0.041–0.060. The magnitude of μ^{ES} for this B_1 ES increases significantly with the basis set size, and a difference of -0.21 D exists between our *aug*-cc-pVTZ and CBS TBEs. The latter compares well with an earlier CASPT2 value.[99] This trend is even exacerbated for the μ^{ES} value associated with the A_2 state that changes by -0.56 D from *aug*-cc-pVTZ to CBS, whereas no significant changes can be noticed between CCSDT and CCSDTQ. For the lowest A_1 ES, there is a significant mixing with a close-lying state of the same symmetry, and the addition of a second set of diffuse is mandatory to obtain reasonable estimates of the oscillator strength (which is much too large with *aug*-cc-pVTZ). Our TBE/CBS f value of 0.062 could still be slightly too large, but the experimental values range from 0.041 to 0.073 (see Table 3 in Ref. 100). In contrast to what we found for the two lower-lying ES, μ^{ES} for the A_1 ES is not very sensitive to the basis set size, with a difference of -0.09 D only between the *aug*-cc-pVTZ and CBS TBEs. We can already conclude, from the data of Table 2, that the basis set required to reach accurate estimates not only differs from one ES to another, but might also be very different, for a given ES, in oscillator strengths and ES dipoles. On a brighter note, the improvements brought by the Q and P excitations are rather limited, meaning that CCSDT seems already sufficient (for water at least).

As can be seen in Table 3, the ground-state dipole moment of hydrogen sulfide remains almost unchanged when increasing the CC order or the size of the basis set: all estimates fall in a tight window: 1.00±0.03 D. Our TBE/CBS of 0.99 D is very close from the experimental value of 0.974±0.005 D, although we use a theoretical geometry. For the lowest $^1A_2$ transition, the excitation energies are almost converged with CCSDT, but a rather large basis set is required, like in water. Our TBE/CBS of 6.10±0.03 eV agrees with previous high-order estimates.[5,99,105] To the best of our knowledge there are no accurate experimental estimates for this dark state. While the methodological effects remain firmly under control for the ES dipole, CCSDT being again sufficient, the basis set effects are huge: at the CCSDT level, we have μ^{ES} = 0.13 D with *aug*-cc-pVDZ but more than four times larger (0.65 D) with d-*aug*-cc-pVQZ. Our TBE/CBS of 0.72±0.02 D is slightly larger than the best previous estimate we have found (0.65 D).[99] The transition energy to the second ES, of B_1 symmetry, is also basis set sensitive, although in that case larger bases yield larger (and not smaller) transition energies, as shown in Table 3. Our TBE/CBS of 6.33±0.01 eV is close to previous estimates,[5,99,105] and also consistent with measurements,[107] although we recall that such comparisons should be made with care. The dipole moment of the B_1 state is very large and relatively insensitive to the basis set as compared to its A_2 counterpart. Our TBE/CBS of -1.74±0.02 D is very close to an earlier CASPT2 estimate.[99] Finally, the computed oscillator strength is rather methodologically insensitive and our best estimate of 0.060 is within 0.010 of the two most recent measurements we could found,[108,109] and compares favorably with two (rather old) theoretical estimates: 0.081[110] and 0.075.[111]

3.3 BF

For this diatomic, the lowest Π ES behaves rather nicely, and one notices in Table 4 that we could reach very stable estimates for all investigated properties, CCSDT/*aug*-cc-pVTZ delivering already sufficiently accurate values. It is noteworthy, that for any given basis set, CCSD underestimates μ^{GS} (μ^{ES}) by ca. 4% (25%), highlighting the difficulty posed by ES for "simple" methods. The experimental value of μ^{GS} was measured to be 0.5±0.2 D,[112] a value that was suggested to be too low more than five scores ago.[113] Indeed, the present TBE/CBS of 0.80 D is significantly above the upper limit of the experimental error bar, but in good agreement with an earlier MRCI+Q value (0.84 D) obtained with a very large basis set.[114] Our best estimate for ΔE_{vert}, 6.39 eV, is slightly larger than a rather old CIPSI value of 6.329 eV,[115] whereas unfortunately, Ref. 114 lists adiabatic energies only. For the oscillator strength, the TBE of Table 4 is larger than previous experimental[116] and theoretical[113] data, but the uniformity of our estimates gives confidence in their quality. Eventually, we predict a significant drop in polarity when going from the GS to the ES, with a μ^{ES} value of ca. 0.27 D. This value is much larger than the MRCI+Q μ^{ES} reported by Magoulas and coworkers (0.01 D),[114] but the latter is obtained on the ES rather than the GS geometry.

3.4 CO

For C=O, we considered six different excited states of various nature, and our results are listed in Table 5. For the three lowest transitions of valence character, one notices that CCSDT is again sufficient, with very small corrections brought by the Q and P excitations, e.g., the Q-induced changes attain ca. ±0.01–0.02 D only for μ^{ES}. For these three transitions, the basis set effects are also firmly under control for all prop-

Table 2: Ground-state dipole moment μ^{GS}, vertical transition energies ΔE_{vert}, oscillator strengths f, and excited state dipole moments $\mu^{\mathrm{ES}}_{\mathrm{vert}}$ determined for H_2O (GS geometry). See caption of Table 1 for details.

		$^1A_1$	$^1B_1$(Ryd, $n \to 3s$)			$^1A_2$(Ryd, $n \to 3p$)		$^1A_1$(Ryd, $n \to 3s$)		
Basis	Method	μ^{GS}	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	ΔE_{vert}	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$
aug-cc-pVDZ	CCSD	1.870	7.447	0.057	-1.404	9.213	-0.936	9.861	0.103	-1.095
	CCSDT	1.849	7.497	0.058	-1.420	9.279	-0.978	9.903	0.104	-1.149
	CCSDTQ	1.848	7.529	0.058	-1.415	9.313	-0.974	9.937	0.105	-1.143
	CCSDTQP	1.848	7.522	0.058	-1.414	9.318	-0.972	9.941	0.105	-1.141
aug-cc-pVTZ	CCSD	1.864	7.597	0.053	-1.549	9.361	-1.056	9.957	0.098	-1.163
	CCSDT	1.842	7.591	0.054	-1.565	9.368	-1.110	9.949	0.100	-1.221
	CCSDTQ	1.840	7.620	0.054	-1.559	9.401	-1.107	9.981	0.100	-1.215
aug-cc-pVQZ	CCSD	1.873	7.660	0.052	-1.655	9.422	-1.237	10.004	0.095	-1.226
	CCSDT	1.850	7.637	0.053	-1.667	9.410	-1.294	9.980	0.097	-1.286
aug-cc-pV5Z	CCSD	1.876	7.683	0.051	-1.726	9.444	-1.309	10.010	0.089	-1.229
d-*aug*-cc-pVDZ	CCSD	1.861	7.429	0.052	-1.759	9.179	-1.658	9.731	0.050	-0.884
	CCSDT	1.841	7.479	0.053	-1.754	9.244	-1.703	9.792	0.057	-1.021
d-*aug*-cc-pVTZ	CCSD	1.865	7.592	0.051	-1.764	9.348	-1.626	9.869	0.057	-1.154
	CCSDT	1.843	7.586	0.052	-1.767	9.353	-1.681	9.872	0.062	-1.232
d-*aug*-cc-pVQZ	CCSD	1.874	7.659	0.051	-1.765	9.416	-1.622	9.932	0.058	-1.228
	CCSDT	1.851	7.636	0.052	-1.770	9.403	-1.678	9.917	0.063	-1.295
aug-cc-pVTZ	TBE[a]	1.840	7.614	0.054	-1.558	9.405	-1.106	9.985	0.100	-1.213
CBS	TBE[b]	1.86±0.01	7.71±0.02	0.052±0.001	-1.77±0.04	9.49±0.02	-1.67±0.01	9.99±0.01	0.062±0.002	-1.30±0.03
Lit.	Th.	1.853[c]	7.66[c]	0.054[d]	-1.787[c]	9.42[c]	-1.682[c]	9.97[e]	0.100[d]	
			7.70[e]	0.049[f]		9.47[e]				
			7.71[f]					9.92[g]	0.055[g]	
	Exp.	1.850[h]	7.41[i]	0.046±0.007[j]		9.20[i]		9.67[i]	0.051[k]	

[a,b] See corresponding footnotes in Table 1; [c] CASPT2/d-*aug*-cc-pVQZ values from Ref. 99; [d] LR-CC3/*aug*-cc-pVTZ value from Ref. 5; [e] Basis set corrected exFCI/*aug*-cc-pVQZ values from Ref. 5; [f] 2FVCAS/MR-CI/CBS values (f in length gauge) from Ref. 48; [g] 2FVCAS/MR-CI/d-*aug*-cc-pV5Z values (f in length gauge) from Ref. 100; [h] Average of the three (very close) experimental values reported in Table 1 of Ref. 101; [i] Energy loss experiment from Ref. 102; [j] Electron impact from Ref. 103; [k] Electron impact from Ref. 104 integrated by Borges in Ref. 100.

Table 3: Ground-state dipole moment μ^{GS}, vertical transition energies ΔE_{vert}, oscillator strengths f, and excited state dipole moments μ^{ES} determined for H_2S (GS geometry). See caption of Table 1 for details.

		$^1A_1$	$^1A_2$(Ryd, $n \to 4p$)		$^1B_1$(Ryd, $n \to 4s$)		
Basis	Method	μ^{GS}	ΔE_{vert}	μ^{ES}	ΔE_{vert}	f	μ^{ES}
aug-cc-pVDZ	CCSD	1.031	6.343	0.113	6.141	0.068	-1.983
	CCSDT	1.016	6.286	0.131	6.098	0.067	-1.946
	CCSDTQ	1.015	6.286	0.137	6.103	0.067	-1.934
	CCSDTQP	1.015	6.286	0.137	6.103	0.067	-1.933
aug-cc-pVTZ	CCSD	0.990	6.246	0.503	6.295	0.064	-1.893
	CCSDT	0.978	6.185	0.496	6.237	0.063	-1.875
	CCSDTQ	0.977	6.181	0.498	6.238	0.063	-1.866
aug-cc-pVQZ	CCSD	1.001	6.212	0.650	6.349	0.062	-1.822
	CCSDT	0.990	6.153	0.636	6.288	0.061	-1.815
aug-cc-pV5Z	CCSD	0.998	6.177	0.691	6.368	0.061	-1.794
d-*aug*-cc-pVDZ	CCSD	1.017	6.297	0.445	6.130	0.065	-1.811
	CCSDT	1.002	6.241	0.458	6.086	0.065	-1.776
d-*aug*-cc-pVTZ	CCSD	0.989	6.228	0.597	6.292	0.062	-1.774
	CCSDT	0.977	6.167	0.587	6.234	0.062	-1.761
d-*aug*-cc-pVQZ	CCSD	1.001	6.206	0.668	6.347	0.061	-1.758
	CCSDT	0.989	6.147	0.653	6.286	0.061	-1.755
aug-cc-pVTZ	TBE[a]	0.977	6.181	0.498	6.238	0.063	-1.865
CBS	TBE[b]	0.99±0.01	6.10±0.03	0.72±0.01	6.33±0.01	0.060±0.001	-1.74±0.02
Lit.	Th.	0.989[c]	6.12[c];6.10[d];6.10[e]	0.653[c]	6.27[c];6.29[d];6.33[e]	0.063[f]	-1.733[c]
	Exp.	0.974±0.005[g]			6.326[h]	0.0542[i];0.0547[j]	

[a,b] See corresponding footnotes in Table 1; [c] CASPT2/d-*aug*-cc-pVQZ values from Ref. 99; [d] Basis set corrected exFCI/*aug*-cc-pVQZ values from Ref. 5; [e] SS-RASPT2/ANO-RCC-VTZP+diffuse values from Ref. 105; [f] LR-CC3/*aug*-cc-pVTZ value from Ref. 5; [g] From Ref. 106; [h] From Ref. 107 (see Table 9 of this work); [i] From Ref. 108, obtained by integrating the experimental absorption spectrum in the 5.2–7.7 eV region: [j] From Ref. 109, obtained by integrating the experimental absorption spectrum in the 5.2–7.7 eV region.

Table 4: Ground-state dipole moment μ^{GS}, vertical transition energies ΔE_{vert} , oscillator strengths f, and excited state dipole moments μ^{ES} determined for BF (GS geometry). See caption of Table 1 for details.

		$^1A_1$	$^1\Pi$(Val, $\sigma \rightarrow \pi^\star$)		
Basis	Method	μ^{GS}	ΔE_{vert}	f	μ^{ES}
aug-cc-pVDZ	CCSD	0.832	6.534	0.479	0.240
	CCSDT	0.861	6.491	0.475	0.311
	CCSDTQ	0.861	6.486	0.474	0.316
	CCSDTQP	0.860	6.485	0.474	0.316
aug-cc-pVTZ	CCSD	0.794	6.464	0.475	0.222
	CCSDT	0.824	6.423	0.469	0.293
	CCSDTQ	0.824	6.417	0.468	0.300
aug-cc-pVQZ	CCSD	0.783	6.449	0.475	0.207
	CCSDT	0.812	6.411	0.468	0.279
aug-cc-pV5Z	CCSD	0.782	6.443	0.475	0.202
d-*aug*-cc-pVDZ	CCSD	0.822	6.521	0.478	0.262
	CCSDT	0.850	6.477	0.474	0.330
d-*aug*-cc-pVTZ	CCSD	0.794	6.459	0.476	0.237
	CCSDT	0.824	6.419	0.469	0.308
d-*aug*-cc-pVQZ	CCSD	0.784	6.448	0.475	0.213
	CCSDT	0.813	6.409	0.468	0.285
aug-cc-pVTZ	TBE[a]	0.824	6.417	0.468	0.299
CBS	TBE[b]	0.80±0.01	6.39±0.01	0.467±0.001	0.27±0.01
Lit.	Th.	0.84[c]	6.329[d]	0.30[e]	0.01[c]
	Exp.	0.5±0.2[f]		0.40[g]	

[a,b] See corresponding footnotes in Table 1; [c] MRCI+Q/*aug*-cc-pV6Z values from Ref. 114, the ES dipole was obtained on the corresponding ES geometry; [d] CIPSI/(11s7p4d3f)/[6s4p4d3f] estimate from Ref. 115; [e] Sum of the MRCI f value determined for GS $v'' = 0$, from Table 5 of Ref. 113; [f] From Ref. 112; [g] f_{00} value from Ref. 116.

erties, and, both f and μ^{ES} are not significantly affected by the second set of diffuse orbitals and *aug*-cc-pVQZ is likely sufficiently large to obtain accurate estimates. As a consequence, it is rather straightforward to get stable CBS extrapolations. The moderate impact brought by the quadruples seems to pertain for the three higher-lying Rydberg ESs of carbon monoxide, except for the dipole moment of the second $^1\Sigma^+$ transition for which μ^{ES} increases by +0.111 D when going from CCSDT/*aug*-cc-pVTZ to CCSDTQ/*aug*-cc-pVTZ. As expected, much larger basis sets are needed to obtain converged properties for the Rydberg ESs than for their valence counterparts. Indeed, when going from *aug*-cc-pVTZ to d-*aug*-cc-pVQZ, the CCSDT ES dipoles of the three lowest Rydberg states drastically change from -3.89 to -1.81 D, from +6.09 to +4.34 D, and from +0.14 to -1.44 D, respectively. In the same time, the CCSDT value of f for the two significantly dipole-allowed transitions decrease by -36% when considering the same basis set pair. Obviously such behavior makes the error bar obtained for our TBE/CBS estimates non-negligible (see below). In contrast, and as already noted above, the transition energies of these Rydberg ESs can be effectively estimated with high accuracy.

When comparing with previously reported experimental or theoretical estimates, one should keep in mind that the properties of carbon monoxide are strongly affected by the bond length,[118] and that we used a theoretical geometry with a slightly too elongated double bond (2.142 bohr *versus* 2.132 bohr experimentally). Our TBE/CBS for μ^{GS} is 0.10 D, which is slightly too low as compared to both the experimental measurement[116] and the theoretical calculations performed at the experimental geometry.[47,123] For the lowest and well-studied Π ES, the present 8.46±0.01 eV ΔE_{vert} value is close to our recent FCI estimate (8.48 eV)[5] as well as to the experimental value (8.51 eV) deduced with the help of the measured spectroscopic constants and the reconstruction of the potential energy surfaces.[124] The computed value of $f = 0.165$ is in very reasonable agreement with the estimates obtained by dipole (e,e) and (γ,γ) spectroscopies: 0.181[119] and 0.194,[121] respectively. We refer the interested readers to Tables 3 and 4 of Ref. 119, Table 6 of Ref. 125, and Table II of Ref. 126 for more complete lists of experimental and theoretical estimates of the oscillator strength. There is a quite significant elongation of the double bond in the lowest ES. As a consequence the value of μ^{ES} determined on the GS (-0.19±0.01 D) and ES (-0.39±0.01 D) structures do differ significantly. The two Stark effect measurements we are aware of yield $\mu^{ES} = -0.15 \pm 0.05$ D (analysis of the 0-0 band)[120] and $\mu^{ES} = -0.34 \pm 0.01$ D (two-photon LIF),[122] which are also somehow inconsistent, yet of the same order of magnitude as our GS and ES values, respectively. For the two other valence transitions, of Σ^- and Δ symmetries, the TBE/CBS for ΔE_{vert} show small errors (1–2 %) as compared to the "experimental" values.[49] We could not find any experimental estimate of the ES dipoles for these states, likely because they are dark, and we believe that the data listed in Table 5 stand as the most accurate ES dipole values proposed to date. They indicate that the dipole moments of these two ESs are parallel to the one of the GS but have much larger amplitudes. For the three Rydberg transitions considered, our best estimates of the transition energies are fully compatible with the measurements. The lowest Rydberg ES has a very low oscillator strength, a result consistent with previous theoretical[5,49] and experimental[119] investigations. More interestingly, its dipole moment is large and negative. The two available measurements return dipoles of -1.60±0.15[120] and -1.95±0.03 D,[122] whereas the strong basis set effects bring uncertainty to our TBE/CBS (-2.49±0.59 D). The lower bound (-1.90 D) fits the latest experimental value quite well. We have also computed the ES dipole of this state at its equilibrium geometry, but the changes are very limited (see Table S4 in the SI). The second Rydberg Σ^+ transition is much more dipole allowed than the first, and our TBE/CBS of 0.146±0.016 for the oscillator strength is consistent with the most recent measurement (0.136).[121] The dipole moment of this ES is large, positive, and rather unaffected by structural relaxation effects (Table S4 in the SI), our extrapolated value of 4.94±0.47 D agreeing well with the two experimental estimates of 4.52±0.35[120] and 4.50±0.07 D.[122] Eventually, for the highest ES of CO considered herein, the extrapolated f value perfectly fits the measurements.[119,121] We provide, as far as we know, the first estimate of its dipole moment, which is rather small and negative (see bottom of Table 5).

3.5 N_2

The understanding of the nature of the ES in the highly symmetric N_2 molecule is certainly challenging from both a experimental and a theoretical point of view.[127,128,130] Of particular relevance for the present work is the nature and

Table 5: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, oscillator strengths f, and excited state dipole moments, $\mu^{\rm ES}_{\rm vert}$ and $\mu^{\rm ES}_{\rm adia}$, determined for various transitions in CO at its GS geometry. The transition energies are in eV and the dipoles in D.

		$^1\Sigma^+$		$^1\Pi(\text{Val}, n \to \pi^\star)$			$^1\Sigma^-(\text{Val}, \pi \to \pi^\star)$		$^1\Delta(\text{Val}, \pi \to \pi^\star)$	
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$	$\mu^{\rm ES}_{\rm adia}$	$\Delta E_{\rm vert}$	$\mu^{\rm ES}_{\rm vert}$	$\Delta E_{\rm vert}$	$\mu^{\rm ES}_{\rm vert}$
aug-cc-pVDZ	CCSD	0.078	8.671	0.167	-0.157	-0.438	10.096	1.574	10.210	1.405
	CCSDT	0.121	8.574	0.173	-0.079	-0.331	10.062	1.597	10.178	1.432
	CCSDTQ	0.130	8.563	0.174	-0.072	-0.317	10.057	1.613	10.169	1.446
	CCSDTQP	0.132	8.561	0.175	-0.069	-0.312	10.057	1.615	10.168	1.448
aug-cc-pVTZ	CCSD	0.051	8.587	0.161	-0.227	-0.478	9.986	1.577	10.123	1.413
	CCSDT	0.104	8.492	0.164	-0.137	-0.358	9.940	1.595	10.076	1.432
	CCSDTQ	0.113	8.480	0.166	-0.129	-0.344	9.932	1.612	10.066	1.449
aug-cc-pVQZ	CCSD	0.039	8.574	0.160	-0.264	-0.511	9.992	1.567	10.127	1.402
	CCSDT	0.094	8.480	0.163	-0.169	-0.385	9.940	1.586	10.073	1.421
aug-cc-pV5Z	CCSD	0.037	8.571	0.160	-0.280	-0.523			10.130	1.398
d-*aug*-cc-pVDZ	CCSD	0.085	8.663	0.167	-0.149	-0.442	10.087	1.561	10.199	1.386
	CCSDT	0.128	8.565	0.173	-0.069	-0.334	10.053	1.584	10.167	1.413
d-*aug*-cc-pVTZ	CCSD	0.053	8.582	0.160	-0.232	-0.483	9.983	1.568	10.120	1.401
	CCSDT	0.106	8.487	0.164	-0.141	-0.364	9.937	1.586	10.073	1.421
d-*aug*-cc-pVQZ	CCSD	0.040	8.572	0.160	-0.269	-0.513	9.992	1.563	10.126	1.397
	CCSDT	0.094	8.478	0.163	-0.173	-0.389	9.940	1.581	10.073	1.416
aug-cc-pVTZ	TBE[a]	0.115	8.478	0.166	-0.126	-0.339	9.932	1.614	10.065	1.450
CBS	TBE[b]	0.097±0.002	8.46±0.01	0.165±0.001	-0.19±0.01	-0.39±0.01	9.94±0.01	1.60±0.01	10.07±0.01	1.43±0.01
Lit.	Th.	0.1172[c]	8.541[d]	0.121[e]	-0.135[f]		10.045[d]		10.182[d]	
		0.091[f]	8.48[g]	0.168[g]		-0.05[h];-0.19[i]	9.98[g]		10.10[g]	
	Exp.	0.122[j]	8.51[k]	0.181[l]		-0.15±0.05[m]	9.88[k]		10.23[k]	
				0.194[n]		-0.335±0.013[o]				

		$^1\Sigma^+$(Ryd)			$^1\Sigma^+$(Ryd)			$^1\Pi$(Ryd)		
Basis	Method	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$
aug-cc-pVDZ	CCSD	11.171	0.003	-4.134	11.710	0.248	6.323	11.973	0.132	0.192
	CCSDT	10.944	0.001	-4.519	11.518	0.240	6.699	11.767	0.124	0.149
	CCSDTQ	10.926	0.000	-4.572	11.510	0.238	6.768	11.758	0.122	0.143
	CCSDTQP	10.919	0.000	-4.592	11.506	0.238	6.792	11.753	0.121	0.141
aug-cc-pVTZ	CCSD	11.222	0.008	-3.376	11.751	0.208	5.600	11.960	0.115	0.195
	CCSDT	10.987	0.004	-3.894	11.540	0.203	6.094	11.737	0.110	0.138
	CCSDTQ	10.963	0.003	-3.994	11.523	0.202	6.205	11.720	0.108	0.126
aug-cc-pVQZ	CCSD	11.190	0.010	-2.792	11.733	0.183	5.097	11.916	0.103	0.084
	CCSDT	10.954	0.006	-3.350	11.514	0.180	5.619	11.687	0.098	0.007
aug-cc-pV5Z	CCSD	11.133	0.012	-2.077	11.691	0.160	4.614	11.851	0.091	-0.158
d-*aug*-cc-pVDZ	CCSD	10.795	0.009	-1.462	11.393	0.134	3.498	11.535	0.067	-1.957
	CCSDT	10.569	0.006	-1.981	11.175	0.129	3.927	11.313	0.061	-2.146
d-*aug*-cc-pVTZ	CCSD	10.960	0.010	-1.337	11.567	0.134	3.814	11.700	0.074	-1.449
	CCSDT	10.726	0.007	-1.855	11.340	0.130	4.216	11.468	0.069	-1.648
d-*aug*-cc-pVQZ	CCSD	11.010	0.011	-1.298	11.617	0.134	3.946	11.749	0.075	-1.233
	CCSDT	10.774	0.007	-1.810	11.388	0.130	4.337	11.516	0.070	-1.437
aug-cc-pVTZ	TBE[a]	10.956	0.003	-4.014	11.519	0.202	6.229	11.715	0.107	0.124
CBS	TBE[b]	10.84±0.06	0.005±0.001	-2.49±0.59	11.43±0.04	0.146±0.016	4.94±0.47	11.58±0.05	0.078±0.009	-0.70±0.60
Lit.	Th.	10.983[d]	0.029[e]	-2.79[i]		0.133[e]	5.34[i]			
		10.80[g]	0.003[g]		11.42[g]	0.200[g]		11.55[g]	0.106[g]	
	Exp.	10.78[k]	0.009[l]	-1.60±0.15[m]	11.40[k]	0.121[l]	4.52±0.35[m]	11.53[k]	0.074[l]	
				-1.95±0.03[o]		0.136[n]	4.50±0.07[o]		0.074[n]	

[a,b] See corresponding footnotes in Table 1; [c] CCSD(T)/CBS value from Ref. 47; [d] CCSDT/PVTZ+ values from Ref. 68; [e] SOPPA (with "all corrections") from Ref. 49; [f] FCI/cc-pVDZ from Ref. 117; [g] From Ref. 5: the values od $\Delta E_{\rm vert}$ are basis set corrected exFCI/*aug*-cc-pVQZ values whereas f are LR-CC3/*aug*-cc-pVTZ values. Note that a factor of two linked to the degeneracy was incorrectly omitted in this earlier work for the oscillator strengths of the two Π transitions; [h] CC2/AVQZ result from Ref. 45; [i] From Ref. 118; [j] From Ref. 116; [k] Vertical values estimated in Ref. 49 on the basis of the experimental spectroscopic data of Ref. 116; [l] Dipole (e,e) spectroscopy of Ref. 119. For the higher ES, we give the contributions given for the two vibrationnal quanta; [m] Stark measurements of the 0-0 bands from Ref. 120. Note that the value for the lowest ES is considered as an upper limit in this work. The sign of the dipole is assumed to be consistent with theory; [n] Dipole (γ,γ) measurements from Ref. 121, summing over the different vibrational contributions; [o] Two-photon laser induced fluorescence spectroscopy measurement of Stark effect from Ref. 122. The sign of the dipole is assumed to be consistent with theory.

Table 6: Vertical transition energies ΔE_{vert} and oscillator strengths f determined for N_2 (GS geometry). See caption of Table 1 for details.

		$^1\Pi_u$ (Ryd)		$^1\Sigma_u^+$ (Ryd)		$^1\Pi_u$ (Ryd)		$^1\Pi_u$ (Ryd)	
Basis	Method	ΔE_{vert}	f	ΔE_{vert}	f	ΔE_{vert}	f	ΔE_{vert}	f
aug-cc-pVDZ	CCSD	13.451	0.531	13.250	0.311	13.765	0.014	14.497	0.148
	CCSDT	13.174	0.469	13.131	0.334	13.591	0.020	14.228	0.163
	CCSDTQ	13.131	0.458	13.109	0.337	13.560	0.027	14.221	0.164
	CCSDTQP	13.127	0.457	13.107	0.338	13.558	0.028	14.216	0.164
aug-cc-pVTZ	CCSD	13.422	0.439	13.264	0.263	13.674	0.053	14.307	0.136
	CCSDT	13.140	0.435	13.118	0.281	13.455	0.008	14.034	0.148
	CCSDTQ	13.095	0.424	13.090	0.285	13.419	0.014	14.014	0.149
aug-cc-pVQZ	CCSD	13.354	0.357	13.242	0.242	13.638	0.113	14.216	0.134
	CCSDT	13.108	0.422	13.088	0.258	13.372	0.000	13.935	0.145
aug-cc-pV5Z	CCSD	13.235	0.266	13.195	0.219	13.621	0.181	14.100	0.126
	CCSDT	13.037	0.356	13.039		13.306	0.044	13.816	0.138
d-*aug*-cc-pVDZ	CCSD	12.784	0.170	12.822	0.172	13.640	0.101	13.537	0.234
	CCSDT	12.669	0.198	12.712	0.183	13.262	0.256	13.300	0.014
d-*aug*-cc-pVTZ	CCSD	12.978	0.178	13.026	0.181	13.599	0.314	13.682	0.007
	CCSDT	12.827	0.224	12.885	0.193	13.257	0.168	13.396	0.073
d-*aug*-cc-pVQZ	CCSD	13.039	0.181	13.090	0.184	13.601	0.276	13.730	0.043
	CCSDT	12.876	0.234	12.939	0.196	13.256	0.460	13.443	0.082
aug-cc-pVTZ	TBE[a]	13.090	0.423	13.088	0.286	13.417	0.015	14.009	0.148
CBS	TBE[b]	12.83±0.08	0.22±0.06[c]	12.96±0.01	0.207±0.005	13.27±0.07	[d]	13.57±0.12	0.10±0.03[c]
Lit.	Th.	12.73[e]	0.458[e]	12.95[e]	0.296[e]	13.27[e]	0.000[e]		
			0.091[f];0.063[g]		0.65[f];0.221[g]		0.32[f];0.091[g]		0.083[g]
	Exp.	12.78[h]	0.243[i]	12.96[h]	0.279[i]	13.10[h]	0.145[i]		0.080[i]
		12.90[j]		12.98[j]	0.223[k]	13.24[j]		13.63[j]	

[a,b] See corresponding footnotes in Table 1; [c] The extrapolation is very challenging due to the mixing, tentative values; [d] Too unstable to report any reasonable CBS estimate; [e] From Ref. 5: energies are basis set corrected exFCI/*aug*-cc-pVQZ values and the oscillator strengths are LR-CC3/*aug*-cc-pVTZ values. Note that a factor of two linked to the degeneracy was incorrectly omitted in this earlier work for the Π_u transitions; [f] RPA values in length gauge from Ref. 127; [g] SOPA values from Ref. 128; [h] Experimental vertical values given in Ref. 127 deduced from spectroscopic data of Ref. 116; [i] Integrated intensities from electron scattering of Ref. 129. [j] Experimental vertical values given in Ref. 130 deduced from spectroscopic data of Ref. 116; [k] Integrated electron impact induced emission intensities of Ref. 131.

relative ordering of the three lowest $^1\Pi_u$ states. Here, we have classified them following the nature of the underlying MOs,[132] and characterized all of them as Rydberg, although alternative yet reasonable assignments can be found in the literature (*e.g.*, see Tables I and II in Ref. 128). Our results are given in Table 6. The convergence with respect to the CC excitation order is rather quick for both energies and oscillator strengths, with nevertheless non-negligible contributions from the quadruples for the transition energies, a likely consequence of the presence of a triple bond. If basis set effects follow the expected trends for the energies of these high-lying ES (*i.e.*, a decrease of ΔE_{vert} when increasing the basis set size, and a significant impact of the second set of diffuse functions), the changes in the oscillator strength of the $^1\Pi_u$ states when enlarging the basis set are dramatic. As an illustration, the oscillator strength of the second $^1\Pi_u$ transition is 0.000 at the CCSDT/*aug*-cc-pVQZ level but 0.460 at the CCSDT/d-*aug*-cc-pVQZ level. If the extreme sensitivity of N_2's oscillator strengths to the computational setup is known for years (*e.g.*, see Table 6 in Ref. 127), it remains a very striking example of the state mixing nightmare. As a consequence, while we could obtain solid TBE/*aug*-cc-pVTZ values in all cases, the extrapolation to CBS of the f values of the $^1\Pi_u$ transitions is clearly problematic. In comparison, while the basis set effects are far from being negligible for the $^1\Sigma_u^+$ ES, it is definitely possible to establish robust CBS f values for this transition.

For the vertical excitation energies, our TBE/CBS of 12.83, 13.27, and 13.57 eV for the three lowest $^1\Pi_u$ transitions and of 12.96 eV for the lowest $^1\Sigma_u^+$ transition are all in reasonable agreement with current state-of-the-art values. Indeed, as can be seen from the bottom of Table 6, these values are close to the experimental vertical data deduced elsewhere,[130] as well as to our recent FCI/CBS estimates.[5] When turning to the oscillator strengths, there is quite a diversity in the experimentally-measured values (see Table 8 of Ref. 129 for integrated values, and Table 1 of Ref. 133 for individual vibronic contributions from various earlier works). For the $^1\Sigma_u^+$ transition, several estimates are available and they show a wide range on both the theoretical and experimental sides (see Table 3 of Ref. 134), and our TBE/CBS f value of 0.207 is reasonably in line with the most recent electron impact value that we have found (0.223),[131] yet significantly smaller than a value obtained by electron scattering (0.279).[129] For the first and third $^1\Pi_u$ transitions, our estimates of the oscillator strength come with large error bars (0.22±0.06 and 0.10±0.03, respectively), which nevertheless cover the most recent experimental values (0.243 and 0.080),[129] an outcome that we found satisfying. As stated above, no reasonable estimate could be obtained for the oscillator strengths of the remaining transition. It is likely that high-level multi-reference calculations would be welcome in this specific case.

3.6 Ethylene

The ESs of this model π-conjugated hydrocarbon have puzzled chemists for years, in particular the relative ordering and nature of the two lowest singlet ESs considered here (Table

7). As can be seen, the oscillator strength of the lowest Rydberg state of B_{3u} symmetry is much too large with 6-31+G(d) but rapidly converges with the Dunning series, the addition of a second set of diffuse functions playing no significant role. In fact the f values for this transition converges faster when climbing the methodology ladder than the ΔE_{vert} values. For the valence $\pi \rightarrow \pi^\star$ B_{1u} transition, CCSD seems to slightly overestimate the oscillator strength and the convergence with basis set size is nearly reached with *aug*-cc-pVTZ. Our best estimates are 7.42 eV (f=0.076) and 7.90 eV (f=0.338), for the Rydberg and valence transitions, respectively.

Table 7: Vertical transition energies ΔE_{vert} and oscillator strengths f of ethylene (GS geometry). See caption of Table 1 for details.

		$^1B_{3u}$ (Ryd)		$^1B_{1u}$ (Val)	
Basis	Method	ΔE_{vert}	f	ΔE_{vert}	f
6-31+G(d)	CCSD	7.814	0.152	8.275	0.380
	CCSDT	7.725	0.151	8.152	0.365
	CCSDTQ	7.722	0.150	8.137	0.364
	CCSDTQP	7.722	0.150	8.135	0.364
aug-cc-pVDZ	CCSD	7.323	0.080	8.035	0.365
	CCSDT	7.294	0.080	7.944	0.352
	CCSDTQ	7.303	0.080	7.932	0.351
aug-cc-pVTZ	CCSD	7.417	0.078	8.020	0.362
	CCSDT	7.365	0.078	7.918	0.346
aug-cc-pVQZ	CCSD	7.451	0.078	8.023	0.360
d-*aug*-cc-pVDZ	CCSD	7.301	0.078	8.008	0.345
	CCSDT	7.273	0.078	7.920	0.336
d-*aug*-cc-pVTZ	CCSD	7.409	0.077	8.009	0.353
	CCSDT	7.357	0.077	7.908	0.339
d-*aug*-cc-pVQZ	CCSD	7.446	0.077	8.017	0.355
aug-cc-pVTZ	TBE[a]	7.374	0.078	7.905	0.345
CBS	TBE[b]	7.42±0.02	0.076±0.001	7.90±0.01	0.338±0.005
Lit.	Th.	7.45[c]	0.069[c]	8.00[c]	0.333[c]
		7.43[d]	0.078[d]	7.92[d]	0.348[d]
	Exp.	7.11[e]	~0.04[f]	7.60[e]	0.34[g]

[a] CCSDT/*aug*-cc-pVTZ values corrected for Q effects using *aug*-cc-pVDZ and for P effects using 6-31+G(d); [b] See Computational section; [c] From Ref. 88: best composite theory for energies, icCAS(12/15)-CI/VDZ+ for the oscillator strengths of the Rydberg transition and icINO(12/16)-CI/VDZ+ extrapolated to FCI for the oscillator strengths of the valence transition; [d] FCI/CBS (for transition energies) and CC3/*aug*-cc-pVTZ (for f) values from Ref. 5; [e] Experimental values collected in Ref. 135 (see the discussions in Refs. 2, 88, and 136); [f] From Ref. 137 (see text); [g] Vacuum absorption from Ref. 138.

On the theoretical side, the most advanced theoretical study of ethylene's ES likely remains the 2014 investigation of Feller *et al.*[88] With their best composite theory, these authors reported transition energies of 7.45 eV and 8.00 eV with respective f values of 0.069 and 0.333 (see footnotes of Table 7 for details). One can also compare to the values of Ref. 5: FCI/CBS estimates of 7.43 and 7.92 eV with CC3/AVTZ f values of 0.078 and 0.346, respectively. Older works report oscillator strengths of 0.389 for the valence state at the CCSDT/TZVP level,[39] and of 0.078 and 0.358 with CCSD.[139] The experimental measurements of the oscillator strength do not allow to attribute values to individual transitions due to strong overlapping.[140] The generally used experimental reference values are 0.34 or 0.29 for the valence transition, both estimates being obtained from measurements of the vacuum absorption spectrum performed in 1953[138] and 1955,[141] respectively. However, a more recent dipole (e,e) spectroscopy study suggests that the originally measured oscillator strengths in the 7.4–8.0 eV regions are probably too low by ca. 10–15%.[140] Our TBE of 0.338 is therefore reasonably in line with the current experimental knowledge. For the Rydberg ES, the only experimental estimate we are aware of has been reported in 1969 as a *"total f perhaps about 0.04"*.[137] Given the consistency of all theoretical estimates, it seems rather reasonable to state that our current TBE of 0.076 is significantly more trustworthy. In ethylene, in contrast to N_2, theory has clearly the edge because the considered transitions have different spatial symmetries.

3.7 Formaldehyde and thioformaldehyde

The ESs of formaldehyde have been extensively studied before with almost all possible theoretical approaches,[2,5,45,74,142–155] and we have considered here two valence and three Rydberg states. Our results are collected in Table 8 for these fives ESs. The hallmark $n \rightarrow \pi^\star$ transition behaves nicely from a theoretical point of view, and the convergences of both ΔE_{vert} and μ^{ES} with respect to the basis set size are very quick, with a negligible effect of the second set of diffuse functions. Likewise, the corrections brought by the Q and P terms in the CC expansion are rather negligible, and CCSDT/*aug*-cc-pVTZ values are trustworthy. The higher-lying $\pi \rightarrow \pi^\star$ valence ES is significantly more challenging as, on the one hand, the second set of diffuse functions significantly increases the transition energy (by ca. +0.2 eV), decreases the oscillator strength (by roughly 10%), and greatly amplifies the ES dipole (by a factor of 2 or 3), whereas, on the other hand, going from CCSDT to CCSDTQ yields a non-negligible drop of the computed dipole. Extrapolation to the CBS limit is therefore uneasy for the latter property. The oscillator strengths of the three Rydberg ESs are all relatively small, but their absolute and relative amplitudes are fairly independent on the selected level of theory and basis sets, though double augmentation induces a small decrease of the magnitude of the oscillator strength. The value of μ^{ES} for the lowest-lying Rydberg ES cannot be adequately described with the selected Pople basis set, but are easy to extrapolate using Dunning's series. For the second (third) Rydberg ES, all tested approaches agree on the rather small (moderate) amplitude for μ^{ES}, but the basis set effects are quite drastic. For instance, considering the higher-lying Rydberg ES, the CCSDT μ^{ES} value is -0.06 D with 6-31+G(d), -2.19 D with *aug*-cc-pVDZ, and -0.37 D with d-*aug*-cc-pVTZ. Clearly it is challenging to get a definitive CBS estimate.

A very accurate measurement of μ^{GS} for formaldehyde is available at 2.3321±0.0005 D (molecular beam electric resonance spectroscopy),[159] and our theoretical TBE of 2.41±0.01 D seems slightly too large, but is in very good agreement with previous CCSD(T)/CBS (2.393 D)[47] and CCSD(T)/*aug*-cc-pVQZ (2.382 D) estimates.[166] As early as 2004, Hirata proposed a CCSDT/*aug*-cc-pVDZ value of μ^{GS} at 2.33 D,[74] right on the experimental spot, but an experimental geometry was used and the orbital relaxation effects neglected, which might have induced a very slight error compensation. For the hallmark lowest $n \rightarrow \pi^\star$ transition, one can find several experimental estimates of μ^{ES}:

Table 8: Ground-state dipole moment μ^{GS}, vertical transition energies ΔE_{vert}, oscillator strengths f, and excited state dipole moments, $\mu^{\mathrm{ES}}_{\mathrm{vert}}$ and $\mu^{\mathrm{ES}}_{\mathrm{adia}}$, determined of formaldehyde. See caption of Table 1 for details.

		$^1A_1$	$^1A_2$(Val, $n \to \pi^\star$)			$^1B_2$(Ryd, $n \to 3s$)			
Basis	Method	μ^{GS}	ΔE_{vert}	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	$\mu^{\mathrm{ES}}_{\mathrm{adia}}$	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	$\mu^{\mathrm{ES}}_{\mathrm{adia}}$
6-31+G(d)	CCSD	2.584	4.031	1.710	1.870	7.238	0.017	-0.901	-0.479
	CCSDT	2.529	4.012	1.649	1.790	7.232	0.021	-1.459	-0.990
	CCSDTQ	2.518	4.022	1.629	1.751	7.279	0.020	-1.364	-0.906
	CCSDTQP	2.517	4.023	1.627	1.746	7.287	0.020	-1.342	-0.885
aug-cc-pVDZ	CCSD	2.427	4.020	1.397	1.602	7.043	0.018	-2.078	-1.634
	CCSDT	2.368	3.986	1.337	1.524	7.040	0.020	-2.457	-1.976
	CCSDTQ	2.356	3.997	1.319	1.486	7.091	0.020	-2.403	-1.932
aug-cc-pVTZ	CCSD	2.457	4.013	1.416	1.620	7.231	0.018	-1.929	-1.476
	CCSDT	2.389	3.954	1.346	1.534	7.165	0.020	-2.379	-1.873
aug-cc-pVQZ	CCSD	2.475	4.024	1.436	1.635	7.296	0.018	-1.861	-1.406
d-*aug*-cc-pVDZ	CCSD	2.417	4.012	1.386	1.596	7.027	0.016	-1.839	-1.410
	CCSDT	2.359	3.978	1.327	1.518	7.024	0.019	-2.282	-1.803
d-*aug*-cc-pVTZ	CCSD	2.454	4.011	1.418	1.619	7.224	0.017	-1.751	-1.309
	CCSDT	2.387	3.952	1.347	1.533	7.158	0.020	-2.248	-1.743
d-*aug*-cc-pVQZ	CCSD	2.475	4.023	1.439	1.636	7.192	0.018	-1.747	-1.299
aug-cc-pVTZ	TBE[a]	2.375	3.966	1.325	1.491	7.225	0.020	-2.302	-1.809
CBS	TBE[b]	2.41±0.01	3.99±0.01	1.36±0.01	1.52±0.01	7.34±0.01	0.020±0.001	-2.15±0.03	-1.66±0.03
Lit.	Th.	2.393[c]	3.98[d]	1.33[e]	1.48[e]	7.12[d]	0.018[d];0.025[e]	-2.52[e]	-3.45[f]
		2.33[e]	3.97[g]		1.73[h]	7.30[g]	0.021[g];0.018[i]		
		2.44[j]	3.98[j]	1.46[j]					
	Exp.	2.332[k]	4.07[l]		1.53±0.11[m]	7.11[k]	0.028[n]		-0.33±0.16[o]
					1.47[p]		0.041[q]		

		$^1B_2$(Ryd, $n \to 3p$)			$^1A_1$(Ryd, $n \to 3p$)			$^1A_1$(Val, $\pi \to \pi^\star$)		
Basis	Method	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$	ΔE_{vert}	f	$\mu^{\mathrm{ES}}_{\mathrm{vert}}$
6-31+G(d)	CCSD	7.994	0.042	0.199	8.282	0.060	0.082	10.042	0.183	2.057
	CCSDT	8.007	0.040	0.625	8.295	0.058	-0.062	9.829	0.163	1.737
	CCSDTQ	8.045	0.040	0.504	8.341	0.058	-0.076	9.779	0.159	1.670
	CCSDTQP	8.051	0.041	0.483	8.350	0.058	-0.074	9.776	0.158	1.663
aug-cc-pVDZ	CCSD	7.993	0.044	0.414	8.052	0.058	-2.058	9.752	0.157	1.470
	CCSDT	8.002	0.042	0.631	8.068	0.057	-2.190	9.588	0.147	1.356
	CCSDTQ	8.045	0.042	0.557	8.119	0.057	-2.202	9.544	0.143	1.292
aug-cc-pVTZ	CCSD	8.120	0.040	0.319	8.210	0.054	-1.438	9.670	0.139	1.503
	CCSDT	8.070	0.038	0.549	8.164	0.052	-1.617	9.488	0.131	1.372
aug-cc-pVQZ	CCSD	8.153	0.039	0.292	8.267	0.053	-1.145	9.647	0.125	1.388
d-*aug*-cc-pVDZ	CCSD	7.834	0.036	0.062	7.962	0.049	-0.019	9.885	0.126	3.672
	CCSDT	7.846	0.035	0.280	7.983	0.049	-0.301	9.748	0.119	3.553
d-*aug*-cc-pVTZ	CCSD	8.026	0.035	-0.104	8.169	0.050	-0.044	9.862	0.124	3.783
	CCSDT	7.978	0.033	0.117	8.126	0.049	-0.369	9.719	0.108	4.347
d-*aug*-cc-pVQZ	CCSD	8.095	0.035	-0.115	8.240	0.050	-0.055	9.869	0.123	3.918
aug-cc-pVTZ	TBE[a]	8.119	0.039	0.454	8.224	0.052	-1.628	9.441	0.127	1.301
CBS	TBE[b]	8.16±0.02	0.035±0.002	0.21±0.20	8.28±0.04	0.050±0.001	-0.69±0.43	9.52±0.12	0.107±0.002	2.46±1.36
Lit.	Th.	7.94[d]	0.040[d];0.041[e]	0.85[e]	8.16[d]	0.043[d];0.058[e]	-2.16[e]	9.83[d]	0.100[d]	
		8.14[g]	0.037[g];0.035[i]		8.27[g]	0.052[g];0.050[i]		9.26[g]	0.135[g];0.093[i]	
	Exp.	7.97[l]	0.017[n]		8.14[l]	0.032[n]				
			0.018[q]			0.061[q]				

[a,b] See corresponding footnotes in Table 7; [c] CCSD(T)/CBS value from Ref. 47; [d] MR-AQCC-LRT calculations from Ref. 156; [e] CCSDT/*aug*-cc-pVDZ results from Ref. 74; [f] MRDCI value from Ref. 157; [g] exFCI/*aug*-cc-pVTZ transition energies corrected for basis set effects up to d-*aug*-cc-pVFZ and LR-CC3/*aug*-cc-pVTZ for f from Ref. 5; [h] CC2/*aug*-cc-pVQZ figure from Ref. 45; [i] EOM-CCSD from Ref. 158; [j] CASPT2/TZVP values from Ref. 2; [k] Electric resonance spectroscopy from Ref. 159; [l] Various experimental sources collected in Ref. 135; [m] Stark effect measurement on lineshapes from Ref. 160; [n] EELS values from Ref. 161; [o] Values measured from polarized electrochromism reported in Refs. 162,163; [p] Stark effect from quantum beat spectroscopy from Ref. 164; [q] Dipole (e,e) spectroscopy from Ref. 165.

Table 9: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, oscillator strengths f, and excited state dipole moments, $\mu^{\rm ES}_{\rm vert}$ and $\mu^{\rm ES}_{\rm adia}$, determined of thioformaldehyde. See caption of Table 1 for details.

		$^1A_1$	$^1A_2$(Val, $n \rightarrow \pi^\star$)			$^1B_2$(Ryd, $n \rightarrow 4s$)				$^1A_1$(Val, $\pi \rightarrow \pi^\star$)		
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	$\mu^{\rm ES}_{\rm vert}$	$\mu^{\rm ES}_{\rm adia}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$	$\mu^{\rm ES}_{\rm adia}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$
6-31+G(d)	CCSD	1.747	2.302	0.933	0.968	5.937	0.019	-3.378	-3.070	6.961	0.261	1.964
	CCSDT	1.733	2.244	0.948	0.990	5.875	0.018	-3.570	-3.257	6.790	0.223	1.745
	CCSDTQ	1.720	2.246	0.919	0.947	5.890	0.019	-3.563	-3.253	6.713	0.191	1.334
	CCSDTQP	1.719	2.247	0.917	0.943	5.893	0.019	-3.562	-3.252	6.708	0.189	1.305
aug-cc-pVDZ	CCSD	1.742	2.325	0.851	0.873	5.841	0.012	-3.938	-3.693	6.749	0.251	1.875
	CCSDT	1.716	2.253	0.870	0.898	5.796	0.011	-4.140	-3.885	6.597	0.182	1.753
	CCSDTQ	1.704	2.255	0.848	0.864	5.817	0.011	-4.134	-3.885	6.512	0.152	1.264
aug-cc-pVTZ	CCSD	1.737	2.291	0.848	0.870	5.970	0.014	-3.374	-3.149	6.633	0.206	2.379
	CCSDT	1.706	2.207	0.865	0.890	5.900	0.012	-3.665	-3.422	6.467	0.163	1.698
aug-cc-pVQZ	CCSD	1.759	2.296	0.869	0.890	6.018	0.014	-3.162	-2.944	6.607	0.200	2.385
d-*aug*-cc-pVDZ	CCSD	1.735	2.324	0.858	0.877	5.804	0.015	-3.258	-3.055	6.678	0.165	-0.038
	CCSDT	1.709	2.252	0.875	0.901	5.761	0.014	-3.495	-3.281	6.577	0.191	1.749
d-*aug*-cc-pVTZ	CCSD	1.737	2.291	0.851	0.872	5.958	0.015	-3.050	-2.847	6.627	0.206	2.376
	CCSDT	1.707	2.207	0.867	0.892	5.888	0.013	-3.379	-3.154	6.463	0.162	2.193
d-*aug*-cc-pVQZ	CCSD	1.759	2.296	0.871	0.892	6.012	0.014	-2.999	-2.794	6.604	0.199	2.383
aug-cc-pVTZ	TBE[a]	1.694	2.210	0.840	0.851	5.923	0.013	-3.658	-3.421	6.377	0.135	1.179
CBS	TBE[b]	1.73±0.01	2.22±0.01	0.088±0.01	0.089±0.01	6.01±0.01	0.013±0.001	-3.27±0.03	-3.05±0.02	6.33±0.01	0.13±0.01	1.20±0.01
Lit.	Th.	1.700[c]	2.20[d]		0.96[e]	5.99[d]	0.012[f]			6.34[g]	0.178[f]	
		1.72[e]										
	Exp.	1.649[h]	2.033[i]		0.850±0.002[j]	5.841[i]			-2.2±0.3[k]	6.60[i]		
		1.647[l]			0.815±0.020[m]							

[a,b] See corresponding footnotes in Table 7; [c] CCSD(T)/*aug*-cc-pVQZ from Ref. 166; [d] exFCI/*aug*-cc-pVTZ transition energies corrected for basis set effects up to d-*aug*-cc-pVQZ from Ref. 5; [e] CC2/*aug*-cc-pVQZ from Ref. 45; [f] LR-CC3/*aug*-cc-pVTZ f from Ref. 5; [g] CCSDTQ/*aug*-cc-pVDZ transition energy corrected for basis set effects up to d-*aug*-cc-pVQZ from Ref. 5; [h] Molecular beam electric resonance value from Ref. 159; [i] 0-0 energies listed in Table 13 of Ref. 167; [j] Microwave-optical double resonance measurements of Stark effect from Ref. 168; [k] Stark effect on the absorption spectrum from Ref. 169; [l] Stark effect measurement on the microwave spectra from Ref. 170; [m] Intermodulated fluorescence of Stark effect from Ref. 171.

1.48±0.07,[172] 1.56±0.07,[173] 1.4±0.1 D,[174] 1.53±0.11D,[160] and 1.47 D.[164] Somehow surprisingly, the most recent value obtained by Stark quantum beat spectroscopy has hardly been considered as reference in theoretical works than the *maximal* measured value of 1.56 D. On the theory side, one can highlight two significant earlier contributions (on the ES geometry): 1.48 D (CCSDT/*aug*-cc-pVDZ)[74] and 1.73 D (CC2/*aug*-cc-pVQZ).[45] We somehow reconcile these earlier results by using both large basis sets and high CC levels and considering both geometries, leading to a TBE/CBS of 1.52±0.01 D for the adiabatic value, right at the center of the experimental cloud. It is noteworthy that the geometrical relaxation induces a non-negligible increase of the magnitude of the dipole moment for the A_2 ($^1A''$) ES. We indeed found a TBE/CBS value of 1.36±0.02 D for the GS geometry. For this $\mu^{\rm ES}_{\rm vert}$, earlier estimates include the 1.46 D (CASPT2/TZVP)[2] and 1.38 D (CC2/*aug*-cc-pVTZ).[43] For the lowest B_2 Rydberg transition, we are aware of one experiment only (polarized electrochromism), leading to an ES dipole of -0.33±0.16 D.[162,163] While theory does confirm the sign change, it returns a much larger amplitude for the dipole with -2.15±0.03 D (GS structure) or -1.66±0.03 D (ES geometry). At the CCSDT/*aug*-cc-pVDZ level, Hirata reported -2.52 D (vertical),[74] likely the best previous estimate. This significant discrepancy between theory and experiment was previously attributed to an (experimental) mixing between the two lowest-lying B_2 transitions.[157] It seems reasonable to state that theory has the edge in this case. For the higher-lying ES, no Stark effect measurement are available, and our values are very likely more accurate than the previous ones reported at the CCSDT level but with a rather small basis set.[74] Nevertheless, the CBS extrapolation is uneasy and large error bars are obtained for all these high-lying ESs. The oscillator strengths of the three lowest Rydberg transitions, $B_2(n \rightarrow 3s)$, $B_2(n \rightarrow 3p)$, and $A_1(n \rightarrow 3s)$ have been respectively measured as 0.038, 0.017±0.02, and 0.038± 0.04 (absorption spectroscopy),[175] 0.028, 0.017, and 0.032 (EELS),[161] 0.032, 0.019, and 0.036 (absorption),[176] and 0.041, 0.028, and 0.061 [dipole (e,e) spectroscopy].[165] Although the orders of magnitude are consistent with the present calculations, the theoretical values do not follow the same ranking as they yield f values of 0.020, 0.035, and 0.050. Such discrepancy has been attributed by other groups to the difficulty of assigning individual vibronic bands to a specific electronic transition in the experimental spectra.[74,177] Interestingly, our current values are agreeing very well with previous theoretical estimates, that returned 0.018, 0.040, and 0.043 (MR-AQCC-LRT),[156] 0.025, 0.041, and 0.058 (CCSDT),[74] 0.018, 0.035, and 0.050 (EOM-CCSD),[158] and 0.021, 0.037, and 0.052 (CC3).[5] For the brighter $\pi \rightarrow \pi^\star$ transition, we are not aware of experimental f values, but theoretical values reported in previous works are of the order of 0.1: 0.100,[156] 0.093,[158] and 0.135,[5] and our current TBE of 0.107±0.002 lies in the middle of these earlier data. In the original Thiel benchmark, the next A_1 ES with a larger f was actually considered.[2]

In thioformaldehyde (Table 9), one notes relatively stable $\Delta E_{\rm vert}$ and $\mu^{\rm ES}$ for the lowest dipole-forbidden A_2 transition: the convergence is rather fast with respect to both CC expansion and basis set size, so that we can safely report accurate TBE/CBS for both the GS and ES structures. The change of ES dipole between the two geometries is limited as well, contrasting with formaldehyde. This is because there is no puckering effect in thioformaldehyde's A_2 ES: the true minimum belongs to the C_{2v} point group.[33,178,179] For the first Rydberg transition (B_2) the impact of the basis set size is logically more pronounced with, *e.g.*, a +0.76 D change between the CCSDT/*aug*-cc-pVTZ and CCSDT/d-*aug*-cc-pVTZ dipoles, making the TBE/CBS extrapolation uncertainty larger than for the A_2 ES. The difference between the values of $\mu^{\rm ES}$ determined at the GS and ES equilibrium geometries are also larger for the Rydberg excitation than for the lowest transition, despite the planarity of all geometries. When selecting Dunning's basis sets, the weak oscillator strength of the B_2 transition always falls in a rather tight window (0.011–0.015). The valence $A_1(\pi \rightarrow \pi^\star)$ transition is clearly no cakewalk: not only the enlargement of the basis set yields significant changes of $\mu^{\rm ES}$ (*e.g.*, +0.49 D from CCSDT/*aug*-cc-pVTZ to CCSDT/d-*aug*-cc-pVTZ), but, in addition, the impact of the quadruples in the CC expansion becomes significant: the Q term induces a drop of $\mu^{\rm ES}$ by ca. -0.40 D and a decrease of f by ca. 15%. For this transition CCSD is clearly insufficient to obtain accurate ES properties. A chemical understanding of the underlying reasons for this large Q effect in the A_1 ES of thioformaldehyde would likely require an in-depth analysis of the various densities determined at various levels of theory, which is beyond our scope.

The value of $\mu^{\rm GS}$ in thioformaldehyde was measured very accurately: 1.6491 ± 0.0004 D,[159] and our TBE of 1.73±0.01 D is slightly higher. Same comment applies to a previous CCSD(T)/*aug*-cc-pVQZ value of (1.700 D)[166] and a CC2/*aug*-cc-pVQZ value of 1.72 D.[45] Like in formaldehyde, one can find a series of Stark measurements relying on various spectroscopic techniques for the lowest $^1A_2$ ES. Quite a range of magnitudes have been reported for $\mu^{\rm ES}$: 0.79±0.04 D (absorption spectroscopy),[180] 0.838±0.008 D (laser-induced fluorescence excitation),[181] 0.850±0.002 D (microwave-optical double resonance),[168] and 0.815±0.020 D (intermodulated fluorescence).[171] Obviously, the error bars of these measurements are not overlapping, but the latter work warns that values between 0.77 and 0.93 D can be obtained.[171] Our TBEs of 0.88±0.01 D (GS geometry) and 0.89±0.01 D (ES geometry) are therefore obviously compatible with the experimental measures. The only previous wave function-based TBE we are aware of is the CC2/*aug*-cc-pVQZ value of 0.96 D estimated by Hellweg (ES geometry),[45] that appears approximately 0.10 D too large. For the second ES (of Rydberg nature), we know only one measurement of the Stark effect (on the absorption spectrum) that led to a $\mu^{\rm ES}$ value of -2.2 ± 0.3 D.[169] Theory clearly confirms the flip of the dipole as compared to the GS, but our TBEs are significantly larger than this experimental value, irrespective of the considered geometry: -3.27 (GS geometry) and -3.05 D (ES geometry). Given the significant basis set dependence of $\mu^{\rm ES}$ of this state, one clearly needs to be cautious but it is nevertheless likely that the experimental value of -2.2 D is too low. To the best of our knowledge, there is no previous published value of $\mu^{\rm ES}$ for the A_1 ES. Concerning the oscillator strengths, the previous TBEs are likely our CC3/*aug*-cc-pVTZ values: 0.012 (B_1) and 0.178 (A_1),[5] which are consistent with the new values listed in Table 9. On the "experimental side", an estimate of 0.38 was proposed for the valence transition,[182] but it is based on a empirical ratio of 10 compared to an earlier estimate of the oscillator strength for the corresponding Rydberg ES of formaldehyde.[175] We trust that our current TBEs are significantly more accurate.

3.8 Nitroxyl and fluorocarbene

Table 10 provides the dipole moments and transition energies of the lowest ES of HNO. Although this transition is not strictly forbidden by symmetry, all methods return very low f values (< 0.001). Thus, we have not bothered reporting the values of the oscillator strength.

Table 10: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, and excited state dipole moments, $\mu^{\rm ES}_{\rm vert}$ and $\mu^{\rm ES}_{\rm adia}$, determined of nitroxyl. We report the norm of the dipoles. See caption of Table 1 for details.

		$^1A'$	$^1A''$(Val, $n \rightarrow \pi^\star$)		
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	$\mu^{\rm ES}_{\rm vert}$	$\mu^{\rm ES}_{\rm adia}$
6-31+G(d)	CCSD	1.902	1.802	1.982	2.111
	CCSDT	1.876	1.797	1.948	2.076
	CCSDTQ	1.869	1.799	1.938	2.063
	CCSDTQP	1.868	1.800	1.927	2.062
aug-cc-pVDZ	CCSD	1.701	1.779	1.719	1.840
	CCSDT	1.667	1.767	1.681	1.799
	CCSDTQ	1.658	1.770	1.670	1.785
aug-cc-pVTZ	CCSD	1.722	1.756	1.727	1.850
	CCSDT	1.683	1.737	1.688	1.807
aug-cc-pVQZ	CCSD	1.735	1.753	1.735	1.859
d-*aug*-cc-pVDZ	CCSD	1.695	1.778	1.709	1.831
	CCSDT	1.661	1.766	1.671	1.790
d-*aug*-cc-pVTZ	CCSD	1.720	1.755	1.724	1.847
	CCSDT	1.681	1.737	1.685	1.804
d-*aug*-cc-pVQZ	CCSD	1.735	1.753	1.679	1.795
aug-cc-pVTZ	TBE^a	1.674	1.740	1.676	1.791
CBS	TBE^b	1.69±0.01	1.73±0.01	1.69±0.01	1.80±0.01
Lit.	Th.	1.654^c	1.74^d		
	Exp.	1.67±0.03^e	1.63^f		1.69±0.01^g
		1.62±0.02^h			

^{a,b}See corresponding footnotes in Table 7; ^{c}CCSD(T)/CBS value from Ref. 47; ^{d}Unpublished exFCI/*aug*-cc-pVTZ value (from our groups); ^{e}From microwave spectroscopy (Ref. 183); ^{f}0-0 energy from Ref. 184 and references therein. ^{g}Mircowave optical double resonance value from Ref. 185 ^{h}From Stark effects measurements (Ref. 186).

As can be seen in Table 10, the convergences with respect to the CC excitation order and basis size are rather fast: quadruples tune the dipole par ∼-0.01 D only, and basis set extension beyond triple-ζ is unnecessary. In other words CCSDT/*aug*-cc-pVTZ provides very accurate estimates and the CBS extrapolations come with small error bars. One notes that the geometrical relaxation of the ES increases the predicted dipole by ca. +0.11 D for all methods. Our TBE/CBS for $\mu^{\rm GS}$, 1.69±0.01 D, is slightly above the Hai and Head-Gordon value (1.654 D), but the two available experiments

also show discrepancies larger than the reported uncertainties (see bottom of Table 10). For the vertical transition energy, our TBE/*aug*-cc-pVTZ is the same as the result of a CIPSI calculation performed with the same basis set, and logically exceeds the experimental 0-0 energy. For the ES dipole we are aware of two experiments,[184,185] but the former investigated μ_a (one of the two dipole components) only. The most recent experiment yields a total $\mu^{\rm ES}$ of 1.69±0.01 D,[185] which indicates a very slight increase as compared to the GS dipole. Our $\mu^{\rm ES}_{\rm vert}$ (1.69±0.01 D) and $\mu^{\rm ES}_{\rm adia}$ (1.80±0.01 D) values apparently slightly undershoots and overestimates the measured *change* of dipole moment. Again, the final call on the relative accuracy of theory and experiment is hard to make.

Table 11: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, oscillator strengths f, and excited state dipole moments, $\mu^{\rm ES}_{\rm vert}$ and $\mu^{\rm ES}_{\rm adia}$, determined of fluorocarbene. We report the norm of the dipoles. See caption of Table 1 for details.

		$^1A'$		$^1A''$	
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$
6-31+G(d)	CCSD	1.572	2.581	0.009	1.316
	CCSDT	1.552	2.573	0.009	1.287
	CCSDTQ	1.549	2.577	0.009	1.282
	CCSDTQP	1.549	2.578	0.009	1.282
aug-cc-pVDZ	CCSD	1.451	2.541	0.007	0.991
	CCSDT	1.430	2.529	0.006	0.970
	CCSDTQ	1.428	2.534	0.006	0.965
aug-cc-pVTZ	CCSD	1.465	2.507	0.006	0.991
	CCSDT	1.441	2.493	0.006	0.969
aug-cc-pVQZ	CCSD	1.468	2.500	0.006	0.994
d-*aug*-cc-pVDZ	CCSD	1.445	2.536	0.006	0.978
	CCSDT	1.425	2.524	0.006	0.958
d-*aug*-cc-pVTZ	CCSD	1.461	2.505	0.006	0.987
	CCSDT	1.437	2.491	0.006	0.965
d-*aug*-cc-pVQZ	CCSD	1.468	2.500	0.006	0.994
aug-cc-pVTZ	TBE[a]	1.438	2.498	0.006	0.964
CBS	TBE[b]	1.44±0.01	2.48±0.01	0.006±0.001	0.97±0.01
Lit.	Th.	1.426[c]	2.49[d]	0.006[e]	
	Exp.	1.403[f]	2.14[g]		

[a,b] See corresponding footnotes in Table 7; [c] CISD value from Ref. 187; [d] exFCI/*aug*-cc-pVTZ transition energies corrected for basis set effects up to *aug*-cc-pV5Z from Ref. 8; [e] LR-CC3/*aug*-cc-pVTZ f from Ref. 8; [f] From Stark effects measurements (Ref. 186); [g] Experiment 0-0 energy from Ref. 188.

Table 11 provides the dipole moments, oscillator strength, and transition energies for the smallest halocarbene, HCF, a system isoelectronic to the previous one. Although we note a small oscillation of the GS dipole and transition energies going from CCSD to CCSDT and CCSDTQ, it is obvious that the convergence with respect to the CC order is fast. Likewise, basis set effects are moderate in the Dunning series, whereas the use of Pople's basis set yields grossly overestimated oscillator strengths, and ES dipole moments. In short reaching accurate values is not problematic. Our TBE/CBS for $\mu^{\rm GS}$, 1.44±0.01 D, is very close to an earlier CISD estimate, (1.43 D)[187] and both are slightly larger than the most recent measurement we are aware of (1.40 D).[189] As for nitroxyl, our TBE for the vertical transition energy is equivalent to the result of a recent CIPSI calculation, and both are logically larger than the experimental 0-0 energy. The small oscillator strength determined here is also the same as our CC3/*aug*-cc-pVTZ value. We could not find previous estimates of the ES dipole in the literature, and our calculations yield a decrease of ca. 50% as compared to the ground state value, which contrasts with the very similar values obtained for the two states of HNO.

3.9 Silylidene

Let us finish our tour by considering silylidene, $H_2C{=}Si$, a small original molecule presenting two well-defined low-lying Rydberg ES.[191–193] Our results are collected in Table 12. The values of $\mu^{\rm GS}$ are small in magnitude for all methods and one notes that CCSD significantly overestimates the dipole whereas one needs at least a triple-ζ basis set to obtain reasonable data. Our TBE of 0.16 D for the GS dipole is close to the only previous theoretical value we found.[190] There is, to the best of our knowledge, no experimental measurement available. The values of $\Delta E_{\rm vert}$ of the two lowest transitions are insensitive to the addition of quadruples in the CC series and they converge quite well with respect to the basis set size. The current TBEs are equal to the ones we obtained earlier applying a different strategy to reach the FCI limit,[7] and they remain slightly larger than the experimental 0-0 energies.[191,192] For the two ESs, we disclose here the two first estimates of the dipole moments. For the lowest excitation, the dipole clearly flips direction as compared to the GS, which contrasts with thioformaldehyde, and also becomes much larger than the GS dipole with a trustworthy TBE of -1.92±0.01 D. For the second ES, all methods predict relatively small $\mu^{\rm ES}_{\rm vert}$ values, with the CCSD and CCSDT signs sometimes in disagreement for a given basis set. Although the addition of a second set of diffuse orbitals has a quite small effect, the convergence with the size of the basis is quite slow, and we can only state that the final dipole should be almost null, though its sign remains unknown. Finally, one notes that 6-31+G(d) provides too large oscillator strengths for the second ES, but that the stability is otherwise remarkable. One can likely be confident in the proposed TBE value (0.034).

4. CONCLUSIONS

In this contribution, we have considered 30 singlet excited states in thirteen small molecules and strived to obtain oscillator strengths and dipole moments as accurate as possible. To this end, we have performed a series of CC calculations going from (LR) CCSD to CCSDTQP using a large panel of basis sets containing one or two sets of diffuse functions. In all cases, we have obtained FCI/*aug*-cc-pVTZ quality properties, as well as estimates of the corresponding FCI/CBS values, the latter coming with quite large uncertainties in some cases. While FCI results do obviously yield rather definitive answers, we treated only small molecules here with computationally expensive methods, so that the transferability of this strategy to larger compounds is indeed limited. Regarding the CC expansion, we found that the correction brought by the P term is always negligible, whereas the impact of Q is often rather small, although some exceptions

Table 12: Ground-state dipole moment $\mu^{\rm GS}$, vertical transition energies $\Delta E_{\rm vert}$, oscillator strengths f, and excited state dipole moments $\mu^{\rm ES}_{\rm vert}$ determined of silylidene (GS geometry). See caption of Table 1 for details.

		$^1A_1$	$^1A_2$ (Ryd		$^1B_2$ (Ryd)		
Basis	Method	$\mu^{\rm GS}$	$\Delta E_{\rm vert}$	$\mu^{\rm ES}_{\rm vert}$	$\Delta E_{\rm vert}$	f	$\mu^{\rm ES}_{\rm vert}$
6-31+G(d)	CCSD	0.091	2.254	-1.845	3.966	0.045	-0.080
	CCSDT	0.028	2.107	-1.891	3.874	0.042	-0.237
	CCSDTQ	0.019	2.101	-1.909	3.876	0.042	-0.256
	CCSDTQP	0.018	2.101	-1.911	3.877	0.042	-0.259
aug-cc-pVDZ	CCSD	0.181	2.289	-1.836	3.875	0.036	0.162
	CCSDT	0.115	2.146	-1.889	3.795	0.034	0.005
	CCSDTQ	0.105	2.140	-1.905	3.798	0.034	-0.012
aug-cc-pVTZ	CCSD	0.235	2.286	-1.851	3.877	0.036	0.161
	CCSDT	0.153	2.128	-1.905	3.779	0.033	-0.018
aug-cc-pVQZ	CCSD	0.249	2.301	-1.848	3.891	0.037	0.186
d-*aug*-cc-pVDZ	CCSD	0.187	2.290	-1.831	3.876	0.036	0.164
	CCSDT	0.121	2.146	-1.885	3.796	0.034	0.007
d-*aug*-cc-pVTZ	CCSD	0.236	2.287	-1.848	3.877	0.036	0.166
	CCSDT	0.155	2.130	-1.902	3.779	0.033	-0.014
d-*aug*-cc-pVQZ	CCSD	0.249	2.301	-1.847	3.890	0.037	0.186
aug-cc-pVTZ	TBE[a]	0.142	2.122	-1.924	3.783	0.034	-0.039
CBS	TBE[b]	0.16±0.01	2.15±0.00	-1.92±0.01	3.81±0.01	0.034±0.01	0.00±0.04
Lit.	Th.	0.144[c]	2.12[d]		3.80[d]	0.033[e]	
	Exp.		1.88[f]		3.63[g]		

[a,b] See corresponding footnotes in Table 7; [c] CCSD(T)/cc-pVCQZ from Ref. 190; [d] exFCI/*aug*-cc-pVTZ transition energies corrected for basis set effects up to *aug*-cc-pV5Z from Ref. 8; [e] LR-CC3/*aug*-cc-pVTZ f from Ref. 8; [f] 0-0 energy from Ref. 191; [g] 0-0 energy from Ref. 192.

to the latter statement have been observed for the considered set. For instance, a reduction of the oscillator strength and ES dipole by -15% and -0.40 D, respectively, is observed for the valence $\pi \rightarrow \pi^\star$ transition of thioformaldehyde. More problematic is the convergence of the computed properties while increasing the size of the atomic basis set: this convergence can go from very smooth to erratic. For some states, huge differences between the results obtained with simply- and doubly-augmented basis sets are indeed found. The oscillator strengths determined for the three close-lying Π_u excited states of dinitrogen being a typical example of this problem caused by state mixing. All in all, when choices have to be made, it seems a better option to use CCSDT with a very large basis set rather than CCSDTQ with a smaller one when one wishes to perform comparisons with experiment, which as explained in the Introduction is always a challenging task. We have also found several examples herein in which one property, e.g., the oscillator strength, is rather independent from the selected basis set, whereas another, e.g., the ES dipole, is not. One must therefore carefully check the basis set convergence and dependence for all considered states and properties.

Despite these challenges, it is certainly noteworthy that for the vast majority of the properties studied here, we could not only establish the most accurate theoretical estimates available to date, but also obtain values that are compatible with the experimental knowledge when these are accessible, which is not always the case even for the small molecules considered here. Theory sometimes deliver smaller error bars than the corresponding experimental data. This is certainly the case for the smallest compounds treated here, for which very large basis sets could be employed. It should be stressed that the measurements of both f and $\mu^{\rm ES}$ are difficult, so that depending on the experimental techniques, apparently incompatible results are quite commonly reported and that the role of theory is likely critical. At this stage, it might be important to recall that we did not used any experimental input, as even our geometries are theoretically determined. The present effort is thus truly *ab initio*. However, it only provides an idealized picture as we did not aimed at modeling vibronic effects.

As we expected at the beginning of the study, getting the right answer for the right reason in the context of ES properties is certainly more challenging than for the corresponding energies. It is therefore our hope that the reference values reported here will be useful benchmarks and will stimulate further studies in both the theoretical and experimental communities.

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SUPPORTING INFORMATION AVAILABLE

Extra data and Cartesian coordinates. The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/doi.

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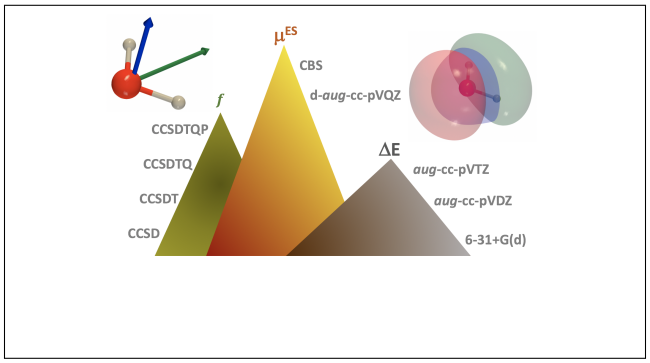
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