



**HAL**  
open science

## Engineering the magnetocaloric properties of PrVO<sub>3</sub> epitaxial oxide thin films by strain effects

H. Bouhani, A. Endichi, D. Kumar, O. Copie, H. Zaari, A. David, A. Fouchet, Wilfrid Prellier, O. Mounkachi, M. Balli, et al.

► **To cite this version:**

H. Bouhani, A. Endichi, D. Kumar, O. Copie, H. Zaari, et al.. Engineering the magnetocaloric properties of PrVO<sub>3</sub> epitaxial oxide thin films by strain effects. *Applied Physics Letters*, 2020, 117 (7), pp.072402. 10.1063/5.0021031 . hal-03014237v2

**HAL Id: hal-03014237**

**<https://hal.science/hal-03014237v2>**




Submitted on 20 Nov 2020

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Engineering the magnetocaloric properties of $\text{PrVO}_3$ epitaxial oxide thin films by strain effects F

Cite as: Appl. Phys. Lett. **117**, 072402 (2020); <https://doi.org/10.1063/5.0021031>  
 Submitted: 06 July 2020 . Accepted: 31 July 2020 . Published Online: 18 August 2020

 H. Bouhani,  A. Endichi, D. Kumar, O. Copie, H. Zaari, A. David,  A. Fouchet,  W. Prellier,  O. Mounkachi, M. Balli, A. Benyoussef, A. El Kenz, and  S. Mangin

## COLLECTIONS

F This paper was selected as Featured



View Online



Export Citation



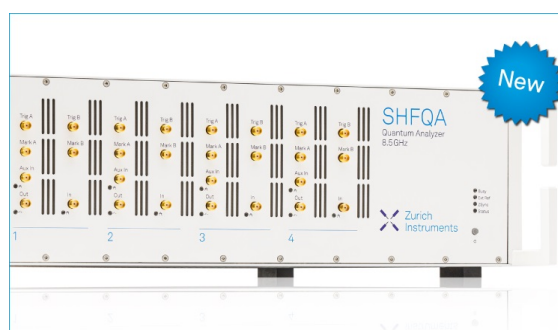
CrossMark

## ARTICLES YOU MAY BE INTERESTED IN

[A four-state magnetic tunnel junction switchable with spin-orbit torques](#)  
 Applied Physics Letters **117**, 072404 (2020); <https://doi.org/10.1063/5.0014771>

[In-plane ferroelectricity and enhanced Curie temperature in perovskite  \$\text{BaTiO}\_3\$  epitaxial thin films](#)  
 Applied Physics Letters **117**, 072902 (2020); <https://doi.org/10.1063/5.0013484>

[Room temperature infrared detectors made of  \$\text{PbTe}/\text{CdTe}\$  multilayer composite](#)  
 Applied Physics Letters **117**, 072102 (2020); <https://doi.org/10.1063/5.0018686>



## Your Qubits. Measured.

Meet the next generation of quantum analyzers

- Readout for up to 64 qubits
- Operation at up to 8.5 GHz, mixer-calibration-free
- Signal optimization with minimal latency

Find out more



# Engineering the magnetocaloric properties of $\text{PrVO}_3$ epitaxial oxide thin films by strain effects

Cite as: Appl. Phys. Lett. **117**, 072402 (2020); doi: [10.1063/5.0021031](https://doi.org/10.1063/5.0021031)

Submitted: 6 July 2020 · Accepted: 31 July 2020 ·

Published Online: 18 August 2020







View Online



Export Citation



CrossMark

H. Bouhani,<sup>1,2,a)</sup>  A. Endichi,<sup>1,2</sup>  D. Kumar,<sup>3</sup> O. Copie,<sup>1</sup> H. Zaari,<sup>2</sup> A. David,<sup>3</sup> A. Fouchet,<sup>3</sup>  W. Prellier,<sup>3</sup>   
O. Mounkachi,<sup>2</sup>  M. Balli,<sup>4</sup> A. Benyoussef,<sup>2</sup> A. El Kenz,<sup>2</sup> and S. Mangin<sup>1</sup> 

## AFFILIATIONS

<sup>1</sup>Université de Lorraine, CNRS, Institut Jean Lamour, F-54000 Nancy, France

<sup>2</sup>Laboratory of Condensed Matter and Interdisciplinary Sciences (LaMCSci), Faculty of Science, Mohammed V University, 1014 Rabat, Morocco

<sup>3</sup>Normandie University ENSICAEN UNICAEN CNRS, CRISMAT 6 Boulevard Maréchal Juin, F-14050 Caen, Cedex 4, France

<sup>4</sup>AMEEC Team, LERMA, ESIE, International University of Rabat, Parc Technopolis, Rocade de Rabat-Salé, 11100, Morocco

<sup>a)</sup>Author to whom correspondence should be addressed: [hamza.bouhani@univ-lorraine.fr](mailto:hamza.bouhani@univ-lorraine.fr)

## ABSTRACT

Combining multiple degrees of freedom in strongly correlated materials such as transition-metal oxides would lead to fascinating magnetic and magnetocaloric features. Herein, the strain effects are used to markedly tailor the magnetic and magnetocaloric properties of  $\text{PrVO}_3$  thin films. The selection of an appropriate thickness and substrate enables us to dramatically decrease the coercive magnetic field from 2.4 T previously observed in sintered  $\text{PVO}_3$  bulk to 0.05 T for compressive thin films making from the  $\text{PrVO}_3$  compound a nearly soft magnet. This is associated with a marked enhancement of the magnetic moment and the magnetocaloric effect that reaches unusual maximum values of roughly  $4.86 \mu_B$  and  $56.8 \text{ J/kg K}$  with the magnetic field change of 6 T applied in the sample plane in the cryogenic temperature range (3 K), respectively. This work strongly suggests that taking advantage of different degrees of freedom and the exploitation of multiple instabilities in a nanoscale regime is a promising strategy for unveiling unexpected phases accompanied by a large magnetocaloric effect in oxides.

Published under license by AIP Publishing. <https://doi.org/10.1063/5.0021031>

The interplay between several degrees of freedom in complex functional materials gained a lot of interest due to its potential to enhance the caloric effects in new alternative cooling technologies such as magnetic cooling.<sup>1–11</sup> The latter, which is based on the magnetocaloric effect (MCE), is an emergent, innovating, and potentially low carbon technology. The MCE, which is an intrinsic property of certain magnetic materials, results in a change of their thermal state when subjected to an external magnetic field. Currently, the gadolinium metal (Gd) is the magnetic material used in the vast majority of magnetic cooling prototypes, mainly due to its magnetic phase transition taking place at 294 K, leading to excellent magnetocaloric properties close to room temperature.<sup>9</sup> However, this metal presents multiple disadvantages such as its easy oxidation as well as the limitation of its working temperature range close to room temperature. In addition, gadolinium cannot be used in large scale applications because of its high cost. These issues have motivated the scientists to search for cheaper, safe, and performant magnetocaloric materials under moderate magnetic fields including intermetallic and oxides.<sup>1–11</sup> The magnetocaloric effect in manganite-based perovskites exhibiting multiferroic

behaviors has become an interesting topic because of the potential application of these oxides in some specific applications such as the liquefaction of hydrogen and space industry.<sup>10,11</sup> This sort of compounds fulfills the necessary conditions for practical applications as they unveil a strong chemical stability, high electrical resistivity, low hysteresis, and mechanical stability.<sup>1</sup> In contrast, the magnetocaloric potential of the  $\text{RVO}_3$  vanadates (R = rare earth) has not yet been explored except for the bulk  $\text{HoVO}_3$ .<sup>12</sup> However, perovskite-type vanadium oxides  $\text{RVO}_3$  display a great variety of phase transitions associated with a series of charge and spin and orbital ordering phenomena, making them interesting candidates from a magnetocaloric point of view.<sup>12</sup>

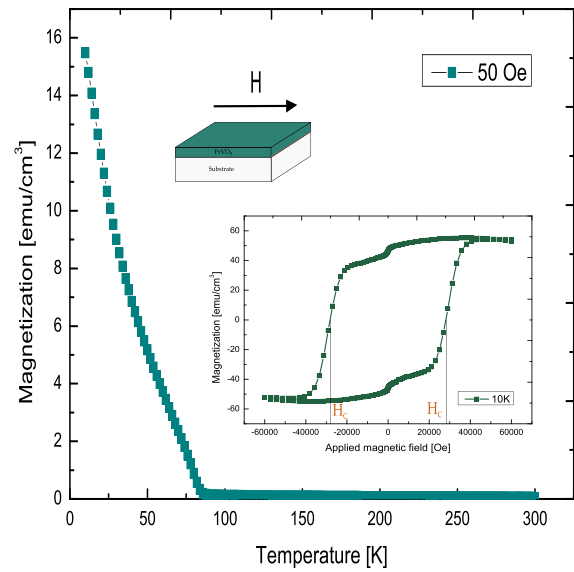
Today's research activities on magnetocaloric thin films attract widespread interest due to their potential integration in miniaturized electronic devices.<sup>13,14</sup> This particularly motivated us to investigate the MCE in  $\text{RVO}_3$  thin films since their behavior strongly depends on the cooperative nature of the Jahn–Teller distortion, making them sensitive to strain effects including the induced biaxial strain due to the lattice mismatch between the substrate and the film. Such structural

effects tend to play an important role in tuning the film properties.<sup>15,16</sup> In this work, we mainly focus on exploring and tuning the magnetic and magnetocaloric properties of high quality epitaxial PrVO<sub>3</sub> (PVO) thin films by applying a compressive strain via a proper choice of substrate.

The epitaxial PVO films were grown by pulsed-laser deposition (PLD) on two different cubic substrates, namely, (001)-oriented (La,Sr)(Al,Ta)O<sub>3</sub> (LSAT) and (001)-oriented SrTiO<sub>3</sub> (STO). Their in- and out-of-plane lattice parameters are found to be 3.868 Å and 3.95 Å with a thickness of 41.7 nm for the LSAT film, while those of the SrTiO<sub>3</sub> film are 3.905 Å and 3.92 Å with a thickness of 100 nm. The lattice parameters have been determined by performing x-ray diffraction measurements; see Refs. 21 and 23 for details. The film deposition was carried out with a KrF excimer laser ( $\lambda = 248$  nm) with a repetition rate of 2 Hz and a laser fluence of  $\sim 2$  J/cm<sup>2</sup> focusing on stoichiometric ceramic targets. The thin films were deposited at an optimum growth temperature ( $T_{GROWTH}$ ) of 650 °C and under an oxygen partial pressure ( $P_{GROWTH}$ ) of  $10^{-6}$  bar. In ZFC measurements, our sample was cooled to the desired temperature under no magnetic field, and then data were collected while heating under the magnetic field. For the FC process, the sample was cooled in the presence of an external magnetic field to the desired temperature. We used two procedures, field cooled cooling (FCC) and field cooled warming (FCW) where data are collected during the cooling and heating processes, respectively. Magnetization was performed by using a Quantum Design SQUID-VSM. Each hysteresis loop was measured after a 150 K excursion above the Néel temperature and corrected by subtracting the diamagnetic contribution arising from the substrate.

To investigate the strain effect on the magnetic properties of PrVO<sub>3</sub> thin films, the magnetization dependence on temperature was measured for PVO/SrTiO<sub>3</sub> under an applied magnetic field of 50 Oe as shown in Fig. 1. The hysteresis loop is also obtained at 10 K in magnetic fields changing between  $-6$  and  $6$  T (see the inset of Fig. 1). Magnetic measurements indicate hard-ferromagnetic behavior below 80 K, which is similar to that reported previously for bulk PVO.<sup>17–19</sup> In fact, the S-shape of magnetization depicts a metamagnetic transition, which is defined as the transition between antiferromagnetic (AF) and ferromagnetic (F) configurations of spins under the effect of magnetic fields or temperature change.<sup>20</sup> The intrinsic coercivity is  $\sim 2.8$  T, while the remanence magnetization is  $\sim 48$  emu/cm<sup>3</sup>. The magnetization saturation is found to be  $\sim 54$  emu/cm<sup>3</sup> being equivalent to only  $0.291 \mu_B/f.u.$  at 10 K.

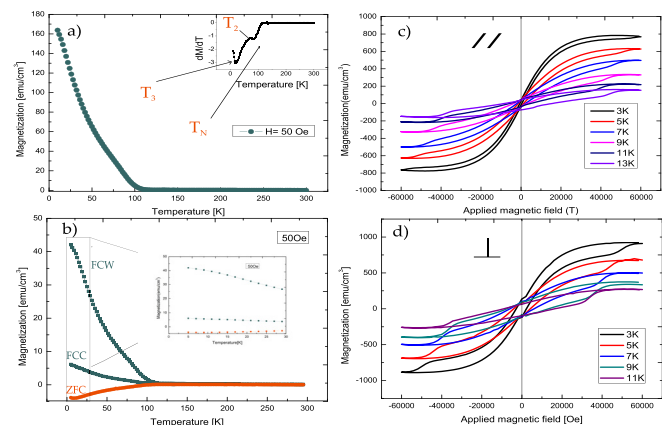
The presence of a soft component can be seen at a magnetic field of  $\sim 0.2$  T indicated by the shape of the  $M$  vs  $H$  loop caused by the magnetic field-induced transition, which is absent at higher temperatures as observed earlier.<sup>21</sup> The high coercivity may arise from the pinning mechanism due to the microstructure as well as the film unit cell orientation compared to the in-plane and out-of plane magnetic field directions. On the other hand, we observed a reduction of  $T_N$  (see Fig. S2) compared to bulk PVO ( $T_N \approx 140$  K), which could be explained by the oxygen vacancy-induced film lattice distortion.<sup>22</sup> XRD reveals that the pseudo cubic volume of PVO unit cell when deposited on SrTiO<sub>3</sub> ( $\approx 60.91 \text{ \AA}^3$ ) is larger than its equivalent of the bulk ( $\approx 58.86 \text{ \AA}^3$ ).<sup>23</sup> As a result, the volume expansion decreases the transfer integral, which tends to reduce the neighbor exchange interactions as the magnetic interactions in this system are governed by superexchange mechanisms. It is worth mentioning that no significant



**FIG. 1.** Magnetization dependence of temperature for the PrVO<sub>3</sub> film on the SrTiO<sub>3</sub> substrate performed under an in-plane applied magnetic field of 50 Oe. The inset displays the magnetic hysteresis loops measured at 10 K after subtracting the diamagnetic contribution of the substrate and the holder.

magnetic anisotropy is observed when comparing the performed magnetic measurements under magnetic fields applied in and out of the sample plane. This can be related to the crystallographic orientation as a strong perpendicular magnetic anisotropy is revealed when the substrate orientation is changed from (001)- to (111)- or (110)-oriented SrTiO<sub>3</sub>.<sup>24</sup>

Figures 2(c) and 2(d) display some selected isothermal magnetization curves for two different orientations measured up to 6 T of a



**FIG. 2.** (a) Temperature dependence of magnetization of PrVO<sub>3</sub> on the LSAT thin film under an in-plane applied magnetic field of 50 Oe. Inset: differentiation of the temperature-dependent magnetization. (b) Temperature dependence of magnetization in zero field-cooling (ZFC), field cooled cooling (FCC), and field cooled warming (FCW) conditions. (c) and (d) Some selected magnetic isotherms in the temperature range of 3–13 K for PVO/LSAT collected in-plane (c) and out-of-plane (d) applied magnetic fields.

PVO film deposited on a (001)-oriented LSAT substrate. As shown, the magnetization saturation is markedly enhanced when compared to PVO/SrTiO<sub>3</sub>, reaching about  $\simeq 900$  emu/cm<sup>3</sup> and  $\simeq 785$  emu/cm<sup>3</sup> at 3 K for H $\perp$  [Fig. 2(d)] and H// [Fig. 2(c)], respectively. At 10 K, the corresponding magnetization saturations are about  $\simeq 402$  emu/cm<sup>3</sup> and  $\simeq 305$  emu/cm<sup>3</sup>, respectively. These values are much larger than those of PVO/SrTiO<sub>3</sub> as can be clearly seen from the inset of Fig. 1. In addition, the coercive field is largely reduced to attain about 0.3 T and 1.1 T for hysteretic loops performed in plane and out of plane (Fig. 2), respectively. More surprisingly, the coercive field decreases dramatically at 3 K reaching only 0.05 T for magnetic fields applied within the film plane as shown in Fig. 2(c). This markedly contrasts with the conventional magnets in which usually the thermal excitations lead to the reduction of coercivity. This contrast could be attributed to FM and AFM couplings and/or the spin and orbital transitions usually leading to stair-like hysteresis observed in bulk PrVO<sub>3</sub>.<sup>18</sup> The enhancement of coercivity when heating may also be attributed to the stress-induced magnetic anisotropy due to the relaxation of the surface stress reported for various magnetic thin films.<sup>24</sup> The temperature dependence of magnetization at an in-plane applied magnetic field of 50 Oe is displayed in Fig. 2(a). As can be clearly observed, a sharp decrease in magnetization at low temperature and a magnetic transition from paramagnetic (PM) to an antiferromagnetic (AFM) phase transition occur at  $T_N = 125$  K. Such a transition is attributed to the beginning of a G-type spin ordering (G-SO).<sup>17</sup> The differentiation of the temperature-dependent magnetization is displayed in the inset of Fig. 2(a) where two additional magnetic transitions take place at  $T_2 \simeq 20$  K and  $T_3 \simeq 80$  K. These transitions were absent in bulk PVO,<sup>25</sup> but they were recently reported in strained PVO films<sup>17</sup> and in doped Pr<sub>1-x</sub>Ca<sub>x</sub>VO<sub>3</sub> compounds.<sup>26</sup> Upon cooling down to 3 K, the plot of the first derivative of the magnetization temperature dependence exhibits a minimum at very low temperature, which could be explained by the polarization of the praseodymium magnetic moments. The newly established order is due to the fact that the antiferromagnetic vanadium sublattice produces an exchange field that results in a ferrimagnetic structure of Pr sublattice under cooling as already observed in other vanadates.<sup>26</sup> This is supported by the presence of a soft component at temperatures below 20 K [Figs. 2(c) and 2(d)] and also by the fact that the magnetization saturation reaches 4.86 and 5.54  $\mu_B$  at 3 K when a magnetic field is applied in and out-of-plane, respectively. On the other hand, these findings inform us on the contribution of both Pr<sup>3+</sup> and V<sup>3+</sup> ions to the whole magnetization since the theoretical saturated moments of free Pr<sup>3+</sup> and V<sup>3+</sup> ions are 3.22 and 2.12  $\mu_B$ , suggesting that all the praseodymium and vanadium moments are fully aligned parallel to the magnetic field.

The ZFC, FCW, and FCC curves were measured in the 50 Oe field applied in the sample plane from 3 to 300 K for PVO/LSAT as shown in Fig. 2(b). The bifurcation between FC and ZFC magnetizations indicates an intrinsic disorder and irreversibility, which is the characteristic of a complex system. This difference reflects the impact of the anisotropy on the shapes of ZFC and FC curves below the ordering temperature since the coercivity is related to the magnetic anisotropy. The latter plays an important role in determining the magnetization at a given field strength during both the ZFC and FC processes since it aligns the spins in a preferred direction. During the ZFC process, the spins are locked in random directions since no magnetic field is applied while cooling the thin films to the desired temperature.

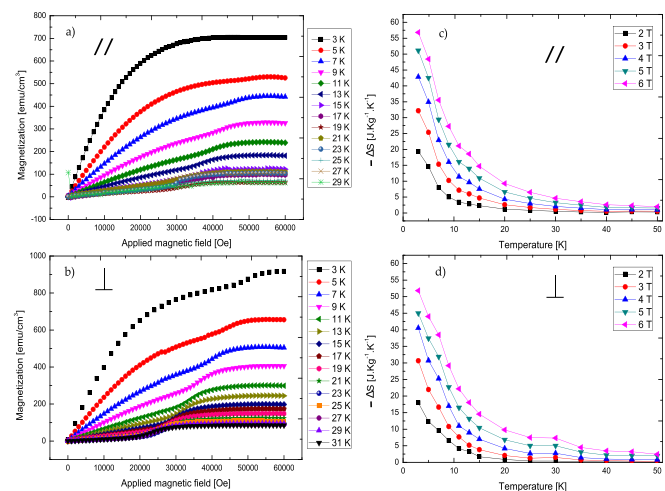
When a small magnetic field is applied at temperatures far below  $T_N$  and as the system is anisotropic,<sup>27</sup> the magnetization decreases to reach negative values, indicating a possible competition between antiferromagnetic interactions, a characteristic that is observed in orthovanadate RVO<sub>3</sub> compounds.<sup>28</sup> A small negative trapped field in the sample space as well as the coercivity could be responsible for the negative magnetization.<sup>29</sup> During the FC process, the PVO film is cooled under the application of a magnetic field. Therefore, the spins will be aligned in a specific direction depending on the strength of the applied magnetic field. Consequently,  $M_{FC}$  continuously increases below  $T_N$  as the temperature decreases.

Magnetic isotherms collected under magnetic fields going from 0 up to 6 T at different temperatures are shown in Figs. 3(a) and 3(b) for the PVO/LSAT films. Except the isothermal magnetization at 3 K, which shows a typical behavior of a ferromagnetic material, all the other isotherms follow a sharp increase when the magnetic field is below 30 kOe, indicating a field induced first order metamagnetic transition from the AFM to FM state as a result of the strong competition between Pr 4f and V 3d spins.<sup>21</sup> Such a competition often leads to a giant MCE in strongly correlated materials.<sup>30-32</sup> A similar behavior is found in the corresponding Arrott plots ( $M^2$  vs  $H/M$ ), confirming the first order nature of the transition according to the Banerjee<sup>33</sup> criterion as the curves show negative slopes at some points (not shown here).

The large field-induced metamagnetic transition in PVO/LSAT films and the soft component in  $M$  vs  $H$  below 30 K are a clear indication of a possible giant magnetic entropy change. In order to explore the magnetocaloric effect in PVO films, the magnetic field-induced entropy change  $-\Delta S_M$  was calculated from magnetic isotherms by using the well-known Maxwell relation (MR) given as follows:<sup>1</sup>

$$\Delta S_M = \int_0^H \left( \frac{\partial M}{\partial T} \right)_H dH.$$

We are aware that the utilization of the Maxwell relation to evaluate entropy changes in materials showing large hysteretic effects could



**FIG. 3.** Magnetic and MCE properties of PrVO<sub>3</sub> on LSAT. (a) Magnetic isotherms in the temperature range of 3–31 K with a step of 2 K under in-plane (a) and out-of-plane (b) magnetic fields. (c) and (d) Temperature dependence of the magnetic entropy change of PrVO<sub>3</sub>/LSAT under some selected magnetic fields applied within (c) and out-of-plane (d).

lead to spurious values as already demonstrated by one of the present paper authors.<sup>34</sup> However, the Maxwell relation can also be used to reasonably evaluate the MCE in terms of the entropy change even in first-order phase transition materials, provided that the remaining magnetization from previous isotherms is suppressed via a thermal loop.<sup>35</sup> On the other hand, it has been demonstrated recently that the magnetocaloric effect in multiferroics could be evaluated perfectly via the Maxwell relation.<sup>36</sup> In this way, it has been particularly found that the deduced entropy change from magnetization measurements of  $\text{EuTiO}_3$  fits perfectly with that obtained from specific heat data.<sup>36</sup> For more information about the impact of hysteretic phenomena on the MCE, we refer the interested reader to Refs. 34 and 35.

In our case, magnetic isotherms of Figs. 3(a) and 3(b) are used to calculate the entropy change exhibited by the PVO films on LSAT. However, since the entropy change is proportional to the area between two successive isotherms,  $\Delta S_M$  was directly calculated without subtracting the magnetic contribution arising from the substrate as already done in the case of  $\text{La}_2\text{NiMnO}_6$  thin films.<sup>37</sup> For PVO/STO, it was difficult to calculate the MCE because of the overlap between  $M$  vs  $H$  curves as well as the very low (negligible) magnetization (see the [supplementary material](#)) at low temperatures.

The temperature dependence of the magnetic entropy change unveils larger values at very low temperature for PVO/LSAT films.  $-\Delta S_M$  reaches roughly a maximum value of  $56.7 \text{ J kg}^{-1} \text{ K}^{-1}$  for a magnetic field changing from 0 to 6 T applied in the sample plane [Fig. 3(a)], being about 63% of its theoretical limit given by  $R \cdot \ln(2J + 1)$ . In a similar field change applied out of plane [Fig. 3(d)],  $-\Delta S_M$  is slightly lower and found to be about  $52.7 \text{ J kg}^{-1} \text{ K}^{-1}$ . Also, the magnetic entropy change shows a large magnetocaloric effect under relatively low magnetic fields, which can be easily reached via permanent magnets. In the magnetic field change of 2 T applied within and out of thin film plane,  $-\Delta S_M$  reaches 19.5 and  $16.3 \text{ J kg}^{-1} \text{ K}^{-1}$ , respectively. As shown in Figs. 3(c) and 3(d), a large magnetocaloric effect can be induced below the AFM transition temperature. This suggests that a major part of the contribution to the MCE comes from the praseodymium 4f spins. A list of some relevant magnetocaloric materials working in the cryogenic temperature range are given in Table I for comparison. As shown, the exhibited  $\Delta S_M$  by the strained PVO film is significantly larger than its equivalent reported for several rare-earth metal transition oxides, making the PVO films potential candidates for low temperature magnetic refrigeration.

To sum up, we have investigated the magnetic and magnetocaloric properties of PVO films grown by pulsed laser deposition, in view of their potential application in cryogenic magnetic cooling. The obtained results reveal that the magnetic and magnetocaloric properties of PVO compounds can be easily tailored by using the thin film

**TABLE I.** Maximum magnetic entropy change  $\Delta S_M$  shown by  $\text{PrVO}_3$  deposited on the LSAT substrate compared to some relevant cryomagnetocaloric compounds. SC means single crystal.

Materials	T(K)	$\Delta H$ (T)	$\Delta S(\text{J/kg K})$	References
PVO/LSAT ( $//-\perp$ )	3	6 T	56.7–52.7	Present work
$\text{GdFeO}_3$ (SC)	2.5	6 T	43.1	38
$\text{EuTiO}_3$	5.6	5 T	42.4	36
$\text{HoVO}_3$ (SC)	15	7 T	17.2	12

approach. Particularly, the coercive magnetic field was dramatically decreased making from the PVO compound a nearly soft magnet. Accordingly, a giant MCE is exhibited by PVO thin films grown on LSAT substrates at low temperatures, pointing out the great impact of strain effects and the competition between AFM and FM exchange interactions. These findings would open the way for the implementation of PVO thin films in some specific applications such as on-chip magnetic micro-refrigeration and sensor technology. Our result not only suggests that epitaxial PVO thin films present non-negligible potential for refrigeration at cryogenic temperatures but may also pave the way for new applications taking, for example, advantage of the possibility to tailor their magnetic coercivity. However, we are aware that the reported entropy change values in PVO thin films are too large when compared to the best magnetocaloric materials working in a similar working temperature range. The necessary was done to reasonably evaluate the MCE in terms of the entropy change by considering the impact of hysteretic phenomena. In order to accurately estimate  $\Delta S$ , the measurements of specific heat in equilibrium conditions are highly required. This point will be certainly addressed in the future.

See the [supplementary material](#) for in-plane MH loops for the PVO/STO film as well as its differentiation of the temperature-dependent magnetization and higher temperature MH loops for the PVO/LSAT film with the magnetic field applied in- and out-of-plane.

This work was supported by the PHC Toubkal 17/49 Project, the French PIA project “Lorraine Université d’Excellence” Reference No. ANR-15-IDEX-04-LUE, and the Institut Carnot ICEEL. M. Balli highly appreciates the financial support from the International University of Rabat.

## DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## REFERENCES

- M. Balli, S. Jandl, P. Fournier, and A. Kedous-Lebouc, *Appl. Phys. Rev.* **4**, 021305 (2017).
- S. Murakami and N. Nagaosa, *Phys. Rev. Lett.* **90**, 197201 (2003).
- Y. Tokura, *Rep. Prog. Phys.* **69**, 797 (2006).
- Y. Taguchi, H. Sakai, and D. Choudhury, *Adv. Mater.* **29**, 1606144 (2017).
- D. Choudhury, T. Suzuki, D. Okuyama, D. Morikawa, K. Kato, M. Takata, K. Kobayashi, R. Kumai, H. Nakao, Y. Murakami, M. Bremholm, B. B. Iversen, T. Arima, Y. Tokura, and Y. Taguchi, *Phys. Rev. B* **89**, 104427 (2014).
- N. H. Dung, Z. Q. Ou, L. Caron, L. Zhang, D. T. Cam Thanh, G. A. de Wijs, R. A. de Groot, K. H. J. Buschow, and E. Brück, *Adv. Energy Mater.* **1**, 1215 (2011).
- M. S. Song, K. K. Cho, B. Y. Kang, S. B. Lee, and B. K. Cho, *Sci. Rep.* **10**, 803 (2020).
- X. Moya, S. Kar-Narayan, and N. D. Mathur, *Nat. Mater.* **13**, 439 (2014).
- G. V. Brown, *J. Appl. Phys.* **47**, 3673 (1976).
- M. Balli, S. Jandl, P. Fournier, and M. M. Gospodinov, *Appl. Phys. Lett.* **104**, 232402 (2014).
- M. H. Phan and S. C. Yu, *J. Magn. Magn. Mater.* **308**, 325 (2007).
- M. Balli, B. Roberge, S. Jandl, P. Fournier, T. T. M. Palstra, and A. A. Nugroho, *J. Appl. Phys.* **118**, 073903 (2015).
- X. Moya, L. E. Hueso, F. Maccherozzi, A. I. Tovstolytkin, D. I. Podyalovskii, C. Ducati, L. C. Phillips, M. Ghidini, O. Hovorka, A. Berger, M. E. Vickers, E. Defay, S. S. Dhési, and N. D. Mathur, *Nat. Mater.* **12**, 52 (2013).
- Y. Liu, C. Phillips, R. Mattana, M. Bibes, A. Barthélémy, and B. Dkhil, *Nat. Commun.* **7**, 11614 (2016).

- <sup>15</sup>R. A. Rao, D. Lavric, T. K. Nath, C. B. Eom, L. Wu, and F. Tsui, *Appl. Phys. Lett.* **73**, 3294 (1998).
- <sup>16</sup>M. Chen, S. Bao, Y. Zhang, Y. Wang, Y. Liang, J. Wu, T. Huang, L. Wu, P. Yu, J. Zhu, Y. Lin, J. Ma, C. W. Nan, A. J. Jacobson, and C. Chen, *Appl. Phys. Lett.* **115**, 081903 (2019).
- <sup>17</sup>O. Copie, J. Varignon, H. Rotella, G. Steciuk, P. Boullay, A. Pautrat, A. David, B. Mercey, P. Ghosez, and W. Prellier, *Adv. Mater.* **29**, 1604112 (2017).
- <sup>18</sup>L. D. Tung, *Phys. Rev. B* **72**, 054414 (2005).
- <sup>19</sup>F. Wang, J. Zhang, P. Yuan, Q. Yan, and P. Zhang, *J. Phys.: Condens. Matter.* **12**, 3037 (2000).
- <sup>20</sup>N. Anuniwat, M. Ding, S. J. Poon, S. A. Wolf, and J. Lu, *J. Appl. Phys.* **113**, 43905 (2013).
- <sup>21</sup>D. Kumar, A. David, A. Fouchet, A. Pautrat, J. Varignon, C. U. Jung, U. Lüders, B. Domengès, O. Copie, P. Ghosez, and W. Prellier, *Phys. Rev. B* **99**, 224405 (2019).
- <sup>22</sup>S. Chikazumi, *Physics of Ferromagnetism* (Oxford University Press Inc., New York, 1991).
- <sup>23</sup>O. Copie, H. Rotella, P. Boullay, M. Morales, A. Pautrat, P.-E. Janolin, I. C. Infante, D. Pravathana, U. Lüders, and W. Prellier, *J. Phys.: Condens. Matter.* **25**, 492201 (2013).
- <sup>24</sup>D. Kumar, "Thin film growth by combinatorial epitaxy for electronic and energy applications," Ph.D. thesis (Normandie University, 2019).
- <sup>25</sup>S. Miyasaka, Y. Okimoto, M. Iwama, and Y. Tokura, *Phys. Rev. B* **68**, 100406(R) (2003).
- <sup>26</sup>M. Reehuis, C. Ulrich, P. M. Abdala, P. Pattison, G. Khaliullin, J. Fujioka, S. Miyasaka, Y. Tokura, and B. Keimer, *Phys. Rev. B* **94**, 104436 (2016).
- <sup>27</sup>P. A. Joy, P. S. Anil Kumar, and S. K. Date, *J. Phys.: Condens. Matter.* **10**, 11049–11054 (1998).
- <sup>28</sup>L. D. Tung, M. R. Lees, G. Balakrishnan, and D. M. Paul, *Phys. Rev. B* **75**, 104404 (2007).
- <sup>29</sup>N. Kumar and A. Sundaresan, *Solid State Commun.* **150**, 1162 (2010).
- <sup>30</sup>A. Midya, P. Mandal, S. Das, S. Banerjee, L. S. S. Chandra, V. Ganesan, and S. R. Barman, *Appl. Phys. Lett.* **96**, 142514 (2010).
- <sup>31</sup>M. Shao, S. Cao, S. Yuan, J. Shang, B. Kang, B. Lu, and J. Zhang, *Appl. Phys. Lett.* **100**, 222404 (2012).
- <sup>32</sup>A. Midya, N. Khan, D. Bhoi, and P. Manda, *Appl. Phys. Lett.* **103**, 092402 (2013).
- <sup>33</sup>B. K. Banerjee, *Phys. Lett.* **12**, 16 (1964).
- <sup>34</sup>M. Balli, D. Fruchart, D. Gignoux, and R. Zach, *Appl. Phys. Lett.* **95**, 072509 (2009).
- <sup>35</sup>L. Caron, Z. Q. Oub, T. T. Nguyen, D. T. Cam Thanh, O. Tegus, and E. Bruck, *J. Magn. Magn. Mater.* **321**, 3559 (2009).
- <sup>36</sup>A. Midya, P. Mandal, K. Rubi, R. Chen, J. S. Wang, R. Mahendiran, G. Lorusso, and M. Evangelisti, *Phys. Rev. B* **93**, 094422 (2016).
- <sup>37</sup>D. Matte, M. de Lafontaine, A. Ouellet, M. Balli, and P. Fournier, *Phys. Rev. Appl.* **9**, 054042 (2018).
- <sup>38</sup>M. Das, S. Roy, and P. Mandal, *Phys. Rev. B* **96**, 174405 (2017).