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# A Virgo Environmental Survey Tracing Ionised Gas Emission (VESTIGE)

## IV. A tail of ionised gas in the merger remnant NGC 4424

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### ABSTRACT

We observed the late-type peculiar galaxy NGC 4424 during the Virgo Environmental Survey Tracing Galaxy Evolution (VESTIGE), a blind narrow-band  $H\alpha$ + $[NII]$  imaging survey of the Virgo cluster carried out with MegaCam at the Canada-French-Hawaii Telescope (CFHT). The presence of a  $\sim 110$  kpc (in projected distance) HI tail in the southern direction indicates that this galaxy is undergoing a ram pressure stripping event. The deep narrow-band image revealed a low surface brightness ( $\Sigma(H\alpha) \approx 4 \times 10^{-18}$  erg s<sup>-1</sup> cm<sup>-2</sup> arcsec<sup>-2</sup>) ionised gas tail  $\sim 10$  kpc in length extending from the centre of the galaxy to the north-west, thus in the direction opposite to the HI tail. *Chandra* and XMM X-rays data do not show a compact source in the nucleus or an extended tail of hot gas, while IFU spectroscopy (MUSE) indicates that the gas is photo-ionised in the inner regions and shock-ionised in the outer parts. Medium-resolution (MUSE) and high-resolution (Fabry-Perot) IFU spectroscopy confirms that the ionised gas is kinematically decoupled from the stellar component and indicates the presence of two kinematically distinct structures in the stellar disc. The analysis of the SED of the galaxy indicates that the activity of star formation was totally quenched in the outer disc  $\sim 250$ – $280$  Myr ago, while only reduced by  $\sim 80\%$  in the central regions. All this observational evidence suggests that NGC 4424 is the remnant of an unequal-mass merger that occurred  $\lesssim 500$  Myr ago when the galaxy was already a member of the Virgo cluster, and is now undergoing a ram pressure stripping event that has removed the gas and quenched the activity of star formation in the outer disc. The tail of ionised gas probably results from the outflow produced by a central starburst fed by the collapse of gas induced by the merging episode. This outflow is sufficiently powerful to overcome the ram pressure induced by the intracluster medium on the disc of the galaxy crossing the cluster. This analysis thus suggests that feedback can participate in the quenching process of galaxies in high-density regions.

**Key words.** galaxies: individual: NGC 4424 – galaxies: clusters: general – galaxies: clusters: individual: Virgo – galaxies: evolution – galaxies: interactions – galaxies: ISM

## 1. Introduction

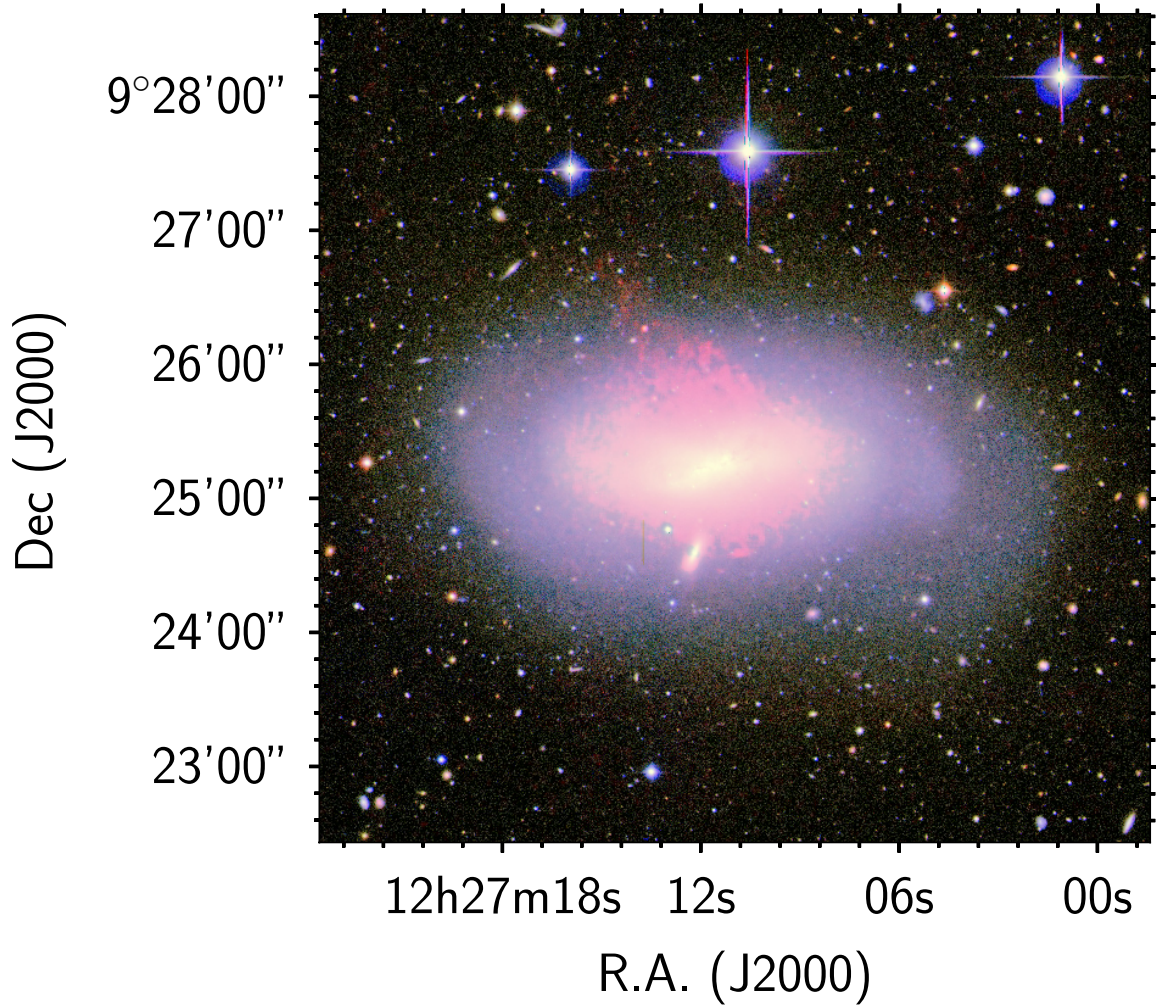
It is well established that the environment plays a major role in shaping galaxy evolution. Rich clusters of galaxies in the local (Dressler 1980; Whitmore et al. 1993) and intermediate-redshift Universe (Dressler et al. 1997) are dominated by a quiescent galaxy population mainly composed of ellipticals and lenticulars. At the same time, compared to their field counterparts, star forming systems in rich environments are generally deprived of their interstellar medium (ISM) in its different phases, from the atomic (Haynes et al. 1984; Solanes et al. 2001; Vollmer et al. 2001a; Gavazzi et al. 2005, 2006b, 2013) and molecular gas (Fumagalli et al. 2009; Boselli et al. 2014b) to the interstellar dust (Cortese et al. 2010, 2012a). As a consequence, their star formation could be significantly reduced (Kennicutt 1983; Gavazzi et al. 1998, 2002, 2006a, 2013; Boselli et al. 2014c, 2015).

Different physical processes have been proposed in the literature to explain the differences observed, as reviewed in

Boselli & Gavazzi (2006, 2014). These processes can be broadly divided in two main families: those related to the gravitational perturbation that a galaxy can suffer in a dense environment (galaxy–galaxy interactions, Merritt 1983, galaxy-cluster interactions, Byrd & Valtonen 1990, galaxy harassment, Moore et al. 1998), and those due to the interaction of the galaxy ISM with the hot ( $T \approx 10^7$ – $10^8$  K) and dense ( $\rho_{\text{ICM}} \approx 10^{-3}$  cm<sup>-3</sup>; Sarazin 1986) intracluster medium (ICM; ram pressure, Gunn & Gott 1972, thermal evaporation, Cowie & Songaila 1977, viscous stripping, Nulsen 1982, starvation, Larson et al. 1980).

The identification of the dominant mechanism responsible for the morphological transformation of galaxies in dense environments is one of the major challenges of modern extragalactic astronomy. It is expected that the relative importance of the aforementioned processes varies with the mass of galaxies and of the overdense regions, and thus with cosmic time given that these structures are growing with time. The identification of the principal process can be carried on by comparing complete samples of galaxies in different environments, from voids to rich clusters, and at different redshifts, with the prediction of cosmological simulations and semi-analytic models of galaxy evolution. It can also be accomplished by comparing

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**Fig. 1.** Pseudo-colour image of NGC 4424 obtained combining the NGVS (Ferrarese et al. 2012) optical  $u$  and  $g$  in the blue channel, the  $r$  and NB in the green, and the  $i$  and the continuum-subtracted  $H\alpha$  in the red. At the assumed distance of the galaxy (16.5 Mpc), 1 arcmin = 4.8 kpc. Stellar shells are not detected down to the NGVS surface brightness limit of  $\mu_g = 29$  mag arcsec $^{-2}$ .

the detailed multifrequency observations of representative objects to the predictions of models and simulations especially tailored to take into account the effects of the different perturbations. This second approach has been successfully applied to a large number of galaxies in nearby clusters where the sensitivity and the angular resolution of multifrequency observations allows us to compare in a coherent study the properties of the different galaxy components (e.g. stellar populations, dust, gas in its different phases, magnetic fields) to the predictions of models. Typical examples are the studies of several galaxies in the Virgo cluster (Kenney & Koopmann 1999; Kenney & Yale 2002; Boselli et al. 2005, 2006; Kenney et al. 2004, 2008, 2014; Vollmer et al. 1999, 2000, 2004a,b, 2005a,b, 2006, 2008a,b, 2009, 2012; Fumagalli et al. 2011; Abramson & Kenney 2014; Abramson et al. 2011, 2016; Jáchym et al. 2013), Coma cluster (Vollmer et al. 2001b; Kenney et al. 2015), A1367 (Gavazzi et al. 1995, 2001a,b; Consolandi et al. 2017), and Norma cluster (Sun et al. 2006, 2007, 2010; Fumagalli et al. 2014; Jáchym et al. 2014; Fossati et al. 2016), or the jellyfish galaxies of the GASP survey (Poggianti et al. 2017). The analysis of these representative objects has been important not only to understand the very nature of the physical mechanisms in act, but also to see that different processes can jointly participate in and compete in the ongoing transformation (Gavazzi et al. 2001a,b,

2003a; Vollmer 2003; Vollmer et al. 2005b; Cortese et al. 2006; Fritz et al. 2017).

NGC 4424, the galaxy analysed in this work, is a typical example of an object undergoing a violent transformation due to different mechanisms (Fig. 1 and Table 1). This star forming system is located at  $\approx 0.9$  Mpc ( $\approx 0.6R_{200}$ ) projected distance from M87 and M49, the two dominant substructures of the Virgo cluster, but it is likely a member of the main cluster A given its distance of  $\approx 15.5$ – $15.8$  Mpc (Munari et al. 2013; Hatt et al. 2018)<sup>1</sup>. The first analysis on optical data done by Kenney et al. (1996) revealed a peculiar morphology with the presence of shells suggesting that this object has recently undergone a merging event. This scenario was later confirmed by high-resolution spectroscopic and CO observations by Cortés et al. (2006, 2015). The presence of a 110 kpc tail in projected length of HI gas has been interpreted as clear evidence that the galaxy is now undergoing a ram pressure stripping event (Chung et al. 2007, 2009; Sorgho et al. 2017). Radio continuum observations also indicate the presence of two polarised cones perpendicular to the disc

<sup>1</sup> Consistently with other VESTIGE works, throughout this work we assume for NGC 4424 a distance of 16.5 Mpc, as for all galaxies within the Virgo cluster A substructure (Gavazzi et al. 1999; Mei et al. 2007).

**Table 1.** Properties of the galaxy NGC 4424 (VCC 979).

Variable	Value	Ref.
Type	SB(s)a: HII	1
$cz$	438 km s <sup>-1</sup>	2
$M_{\text{star}}^a$	10 <sup>10.17</sup> $M_{\odot}$	3
$M(\text{HI})$	10 <sup>8.35</sup> $M_{\odot}$	2
$M(\text{H}_2)$	10 <sup>8.57</sup> $M_{\odot}$	2
Distance <sup>b</sup>	16.5 Mpc	4,5,6,7
incl	62°	8
Proj. distance from M 87	0.9 Mpc	T.W.
$\log f(\text{H}\alpha + [\text{NII}])$	-12.12 ± 0.02 erg s <sup>-1</sup> cm <sup>-2</sup>	T.W.
SFR <sup>a</sup>	0.25 $M_{\odot}$ yr <sup>-1</sup>	T.W.

**Notes.** <sup>(a)</sup> $M_{\text{star}}$  and SFR are derived assuming a [Chabrier \(2003\)](#) IMF and the [Kennicutt \(1998\)](#) calibration; <sup>(b)</sup>assumed distance for the determination of the physical parameters.

**References:** (1) [NED](#); (2) [Boselli et al. \(2014a\)](#); (3) [Boselli et al. \(2015\)](#); (4) [Mei et al. \(2007\)](#); (5) [Gavazzi et al. \(1999\)](#); (6) [Blakeslee et al. \(2009\)](#); (7) [Cantiello et al. \(2018\)](#); (8) [Cortese et al. \(2012b\)](#).

of the galaxy, probably due to a nuclear outflow ([Vollmer et al. 2013](#)).

To study the effects of the environment on galaxy evolution, we are undertaking Virgo Environmental Survey Tracing Ionised Gas Emission (VESTIGE), a deep blind H $\alpha$ + $[\text{NII}]^2$  narrow-band (NB) imaging survey of the Virgo cluster at the CFHT ([Boselli et al. 2018a](#); Paper I). One of the purposes of this survey is to identify galaxies undergoing perturbations with the hostile cluster environment through the observation of peculiar morphologies in the ionised gas component. Deep blind observations of galaxies in nearby clusters, now made possible by the advent of large panoramic detectors coupled with 4–8 m class telescopes, have indicated H $\alpha$  NB imaging as one of the most efficient techniques for this purpose ([Gavazzi et al. 2001a](#); [Yoshida et al. 2002](#); [Yagi et al. 2007, 2010, 2017](#); [Sun et al. 2007](#); [Kenney et al. 2008](#); [Fossati et al. 2012](#); [Zhang et al. 2013](#); [Boselli et al. 2016a](#)). The new H $\alpha$  image of NGC 4424 taken as part of the VESTIGE survey reveals the presence of spectacular features in the ionised gas distribution. We use these new data, combined with *Chandra* and XMM imaging and high-resolution 2D Fabry-Perot spectroscopic data gathered during follow-up observations and integral field spectroscopic data from VLT/MUSE to further analyse the kinematical and physical properties of this intriguing object. The observations and data reduction are described in Sect. 2, while the description of the multifrequency data used in the analysis is given in Sect. 3. The derivation of the physical and kinematical parameters is given in Sect. 4, while the discussion and conclusions follow in Sects. 5 and 6.

## 2. Observations and data reduction

### 2.1. Narrow-band imaging

The galaxy NGC 4424 was observed during VESTIGE pilot observations in April 2016. The observations were carried out using MegaCam at the CFHT coupled with the new NB filter

MP9603 ( $\lambda_c = 6590 \text{ \AA}$ ;  $\Delta\lambda = 104 \text{ \AA}$ ) perfectly suited for the observations of Virgo cluster galaxies. At the redshift of the galaxy ( $cz = 438 \text{ km s}^{-1}$ ) the filter encompasses the H $\alpha$  line ( $\lambda = 6563 \text{ \AA}$ ), where the transmissivity is  $T = 92\%$ , and the two  $[\text{NII}]$  lines ( $\lambda = 6548, 6583 \text{ \AA}$ ). MegaCam is composed of 40 CCDs with a pixel scale of  $0.187 \text{ arcsec pixel}^{-1}$ . As for the other pilot observations ([Boselli et al. 2016a, 2018b](#); [Fossati et al. 2018](#)), the galaxy was observed using the pointing-macro QSO LDP-CCD7 especially designed for the Elixir-LSB data reduction pipeline. This macro is composed of seven different overlapping pointings over a region of  $40 \times 30 \text{ arcmin}^2$ . The determination of the stellar continuum is secured via a similar observation in the  $r$ -band filter. The integration time for each single pointing within the macro was of 660 s in the NB filter and 66 s in the  $r$ -band. The macro was run twice, thus the total integration time is of 924 s in  $r$ -band and 9240 s in NB. The typical sensitivity of the images is similar to that reached during the VESTIGE survey, i.e.  $f(\text{H}\alpha) \sim 4 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2}$  for point sources ( $5\sigma$ ) and  $\Sigma(\text{H}\alpha) \sim 2 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  for extended sources ( $1\sigma$  after smoothing at 3 arcsec resolution; see Paper I for details). The images are of excellent quality, with a median image quality of 0.63 arcsec in the  $r$ -band and 0.62 arcsec in the NB filter.

As for the other VESTIGE observations, the MegaCam images were reduced using Elixir-LSB ([Ferrarese et al. 2012](#)), a data reduction pipeline especially designed to detect the diffuse emission of extended low surface brightness features associated with perturbed cluster galaxies. This pipeline efficiently removes any contribution of the scattered light in the different science frames as demonstrated by the analysis of broadband images taken during the the Next Generation Virgo Cluster Survey (NGVS; [Ferrarese et al. 2012](#); [Mihos et al. 2015](#)). Pilot observations of the galaxy NGC 4569 ([Boselli et al. 2016a](#)) have shown that Elixir-LSB is perfectly suited for the NB frames whenever the images are background dominated, such as those under discussion. The photometric calibration in both filters was done following the procedures given in [Fossati et al. \(in prep.\)](#), with a typical uncertainty of 0.02–0.03 mag in both bands. A further check on the spectroscopic calibration of the images was done by comparing the emission in the NB filter to that obtained by MUSE on the  $\sim 1 \times 1 \text{ arcsec}^2$  gas emitting region. This comparison gives consistent results within 3.5%. The astrometry of each single image was corrected using the MegaPipe pipeline ([Gwyn 2008](#)) before stacking.

### 2.2. X-rays imaging

*Chandra* observation of NGC 4424 were taken on April 17, 2017, with the Advanced CCD Imaging Spectrometer (ACIS). ACIS chip S3 is on the aimpoint (ObsID: 19408; PI: Soria). The image has no background flares and the clean time for the S3 chip is 14.6 ks. There were two XMM observations of NGC 4424. The first (0651790101; PI: Sun) was taken on June 13, 2010, but was completely ruined by background flares. The second (0802580201; PI: Sun) was gathered on December 5, 2017, with an integration time of 34 ks. After removing the background flares, the clean time is 14.1 ks for MOS1, 13.8 ks for MOS2, and 9.2 ks for PN. Because of the accelerated build-up of contamination on the optical blocking filter of ACIS/*Chandra*, the XMM data are much more sensitive than the *Chandra* data below 1.5 keV. Following [Sun et al. \(2010\)](#) and [Ge et al. \(2018\)](#), the X-ray data were reduced using the standard procedures with CIAO 4.10 and CALDB 4.7.9 (*Chandra*), and SAS 17.0.0 (XMM).

<sup>2</sup> Hereafter we refer to the H $\alpha$ + $[\text{NII}]$  band simply as H $\alpha$ , unless otherwise stated.

### 2.3. MUSE spectroscopy

The MUSE data were taken as part of the 097.D-0408(A) and 098.D-0115(A) programmes of the AMUSING supernova survey (Galbany et al. 2018). Observations were carried out using the Wide Field Mode in the spectral range 4800–9300 Å with a spectral resolution of 2.55 Å, corresponding to  $\sim 50 \text{ km s}^{-1}$  at H $\alpha$  as described in Galbany et al. (2016) and Krühler et al. (2017). The data were gathered during two different runs, one in May 2016 with four dithered exposures 625 s long, and centred at  $\sim 18 \text{ arcsec}$  on the east of the nucleus of the galaxy, and the other in January 2017, again with four dithered exposures of 701 s, centred at  $\sim 75 \text{ arcsec}$  on the east of the nucleus. To secure a correct sampling of the sky on this extended galaxy, two exposures were taken with the telescope pointed on an empty region close to the galaxy. The observations were carried out in mostly clear sky conditions. The typical sensitivity of MUSE to low and extended surface brightness features at H $\alpha$  is  $\Sigma(\text{H}\alpha) \sim 4 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ .

The MUSE data were reduced using a combination of in-house python procedures and the MUSE pipeline (Weilbacher et al. 2014), as first presented in Fumagalli et al. (2014) and later improved in Fossati et al. (2016) and Consolandi et al. (2017). These procedures perform the cube reconstruction using the daytime calibrations, and later improve the illumination spatial uniformity. The skylines subtraction is achieved with the Zurich Atmospheric Purge (ZAP, V1.0, Soto et al. 2016). Individual exposures are projected on a common output grid and are then combined with mean statistics. Lastly, the cube is smoothed by  $10 \times 10$  spaxels in the spatial direction to improve the signal-to-noise ratio (S/N) of the spectra without compromising the image quality.

The stellar continuum is modelled and subtracted from the data cube using the GANDALF code (Sarzi et al. 2006), which uses the penalised pixel-fitting code (pPXF, Cappellari & Emsellem 2004). We only fit spaxels with a S/N per spectral channel greater than 5 and we used the MILES spectral library (Vazdekis et al. 2010), and we record the stellar kinematics from the best fit spectra. We fit the emission lines from the continuum subtracted cube using the KUBEVIZ software as detailed in Fossati et al. (2016). After fitting all the spaxels independently we flag as good fits those with a  $S/N > 5$  for the H $\alpha$  flux. In the present paper we also ran a new feature of the KUBEVIZ code that iteratively attempts to improve the unreliable fits using initial guesses obtained from the nearby good spaxels. With this method we obtained good fits for  $\sim 2\%$  more spaxels than with the original fitting method.

Finally, we find no ionised gas emission in the eastern data cube and we find that the stellar kinematics fits are highly uncertain in this cube due to the very low surface brightness of the stellar continuum. For these reasons we only use the central field in the following analysis.

### 2.4. Fabry-Perot spectroscopy

High-resolution 3D spectroscopic observations of NGC 4424 were carried out using two different Fabry-Perot interferometers, SAM-FP on the 4.2 m SOAR telescope located at Cerro Pachón (Chile), and GHASP on the 1.93 m Observatoire de Haute Provence (OHP) telescope (France). SAM-FP, a Fabry-Perot mounted inside the SOAR telescope Adaptive-optics Module, is the only imaging Fabry-Perot interferometer which uses a laser-assisted ground layer adaptive optic over a  $3 \times 3 \text{ arcmin}^2$  field-of-view (Mendes de Oliveira et al. 2017). SAM-FP is equipped

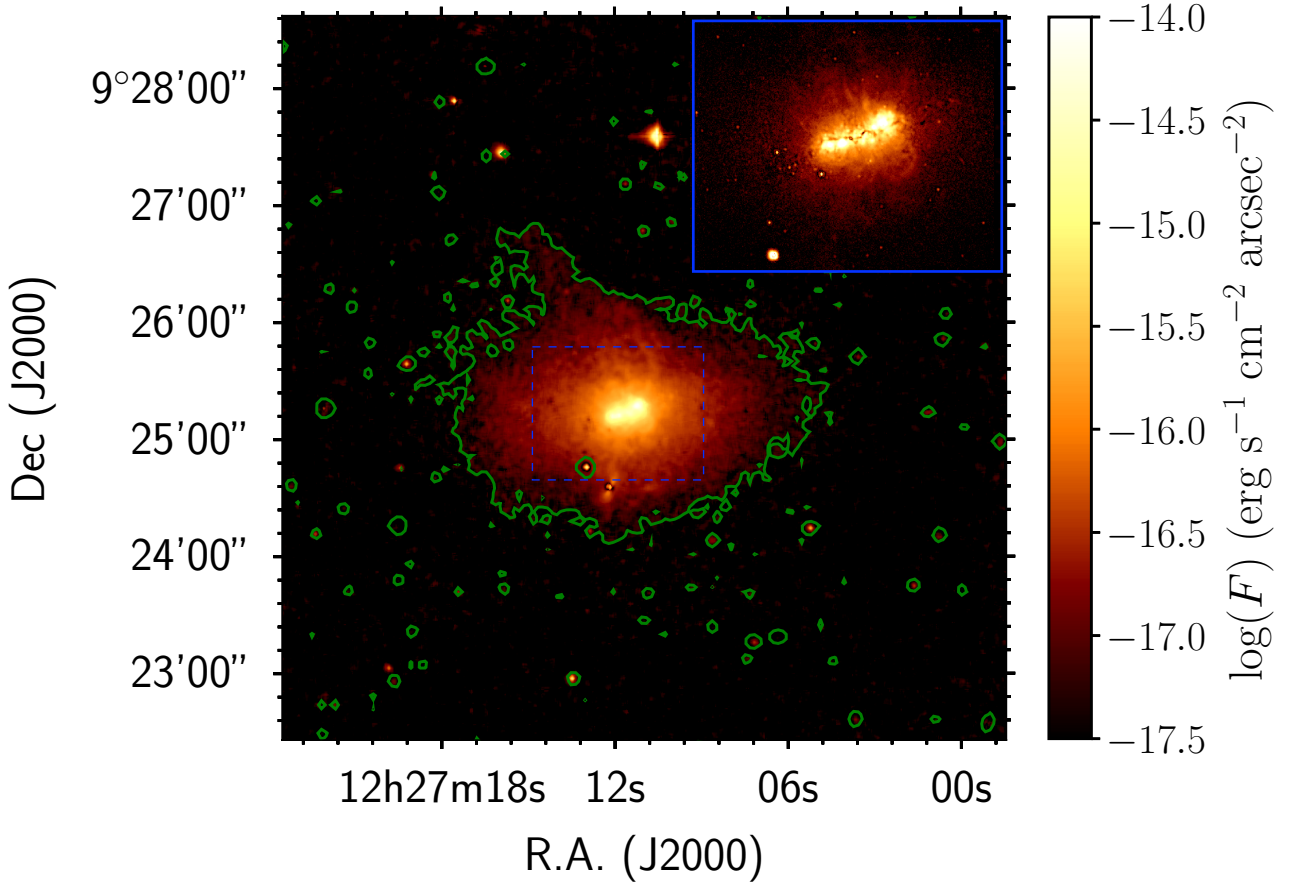
with a single e2v CCD with  $4096 \times 4112$  pixels with a  $2 \times 2$  binned pixel of size  $0.18 \times 0.18 \text{ arcsec}^2$ . The observations were carried out in April 2017 during the commissioning of the instrument. The galaxy was observed with a 1 h exposure in dark sky under photometric conditions. Thanks to the Ground Layer Adaptive Optics (GLAO) module, the mean image quality is 0.55 arcsec. GHASP is a focal reducer including a Fabry-Perot providing a  $5.8 \times 5.8 \text{ arcmin}^2$  field-of-view equipped with a  $512 \times 512$  Imaging Photon Counting System (IPCS) with a pixel scale of  $0.68 \times 0.68 \text{ arcsec}^2$  (Gach et al. 2002; Garrido et al. 2005). The galaxy was observed in dark sky time during photometric conditions, with a total exposure time of 4 h and a typical seeing of 3 arcsec. The free spectral range of the two Fabry-Perot instruments,  $492 \text{ km s}^{-1}$  and  $378 \text{ km s}^{-1}$  for SAM-FP and GHASP, were scanned through 36 and 32 channels, respectively. This provides a comparable spectral resolution of  $R \simeq 12000$  at H $\alpha$ . For both instruments, suitable  $15 \text{ \AA}$  FWHM interference filters were used to select the spectral range of interest. The Ne I 6598.95 Å emission line was used for wavelength calibration. The line detection limits for the diffuse emission of both instruments after Voronoi binning is  $\sim 2.5 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  at OHP and  $\sim 5 \times 10^{-17} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  at Cerro Pachón.

Fabry-Perot data were reduced following the standard procedures fully described in Daigle et al. (2006) and Epinat et al. (2008). Night-sky line subtraction was applied after data cube construction and wavelength calibrations. Using adaptive binning techniques based on 2D Voronoi tessellations applied to the 3D data cubes, we derived Voronoi H $\alpha$  data cubes from which, among others, H $\alpha$  maps, radial velocity, and velocity dispersion fields are computed. With the spatial adaptive binning technique, a bin accretes new pixels until it reaches an a priori criterion, which was fixed here to a  $S/N = 7$  for the emission line within the bin for both sets of observations. The main advantage of this technique is that the best spatial resolution is obtained at each point of the galaxy and that no spatial contamination between the bins is possible, each bin corresponds to the same S/N and therefore has the same statistical weight.

## 3. Multifrequency data

### 3.1. Data

The galaxy NGC 4424 is located close to the core of Virgo and has been observed during several multifrequency surveys of the cluster. Deep GALEX ultraviolet images in the far-ultraviolet (FUV;  $\lambda_c = 1539 \text{ \AA}$ ; integration time = 2008 s) and near-ultraviolet (NUV;  $\lambda_c = 2316 \text{ \AA}$ ; integration time = 2009 s) bands were gathered during the GUViCS survey of the Virgo cluster (Boselli et al. 2011). The galaxy was also observed in the *ugiz* broad bands during the NGVS (Ferrarese et al. 2012), in the near-infrared *H* band (Boselli et al. 2000), in the mid-infrared by WISE (Wright et al. 2010), in the four IRAC bands by *Spitzer* (Ciesla et al. 2014), and in the far-infrared by *Spitzer* with MIPS at 24, 70, and  $160 \mu\text{m}$  (Bendo et al. 2012) and by *Herschel* with PACS at 100 and  $160 \mu\text{m}$  (Cortese et al. 2014) and SPIRE at 250, 350, and  $500 \mu\text{m}$  (Ciesla et al. 2012) as part of the *Herschel* Reference Survey (HRS; Boselli et al. 2010) and the *Herschel* Virgo Cluster Survey (HeViCS; Davies et al. 2010). The images of the galaxy in different photometric bands, including the continuum-subtracted NB H $\alpha$ , the FUV, NUV, *u*, IRAC1  $3.6 \mu\text{m}$ , and the MIPS  $24 \mu\text{m}$  bands, are shown in Fig. 3. High-quality images of the galaxy in the visible were also obtained with the WFC3 on the HST by Silverman et al. (2012; Fig. 4). HI VLA, KAT-7, and WSRT data of NGC 4424 are available



**Fig. 2.** Continuum-subtracted H $\alpha$  image of NGC 4424 with contours at  $\Sigma(\text{H}\alpha) = 4 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ . A prominent tail is visible in the north-east direction extending from the ionised gas halo. The inset shows a contrasted image of the central regions (here on an arbitrary surface brightness scale) with prominent filaments of ionised gas.

from Chung et al. (2009) and Sorgho et al. (2017). VLA radio continuum data at 6 cm and 20 cm, including polarisation, are also available (Vollmer et al. 2013).

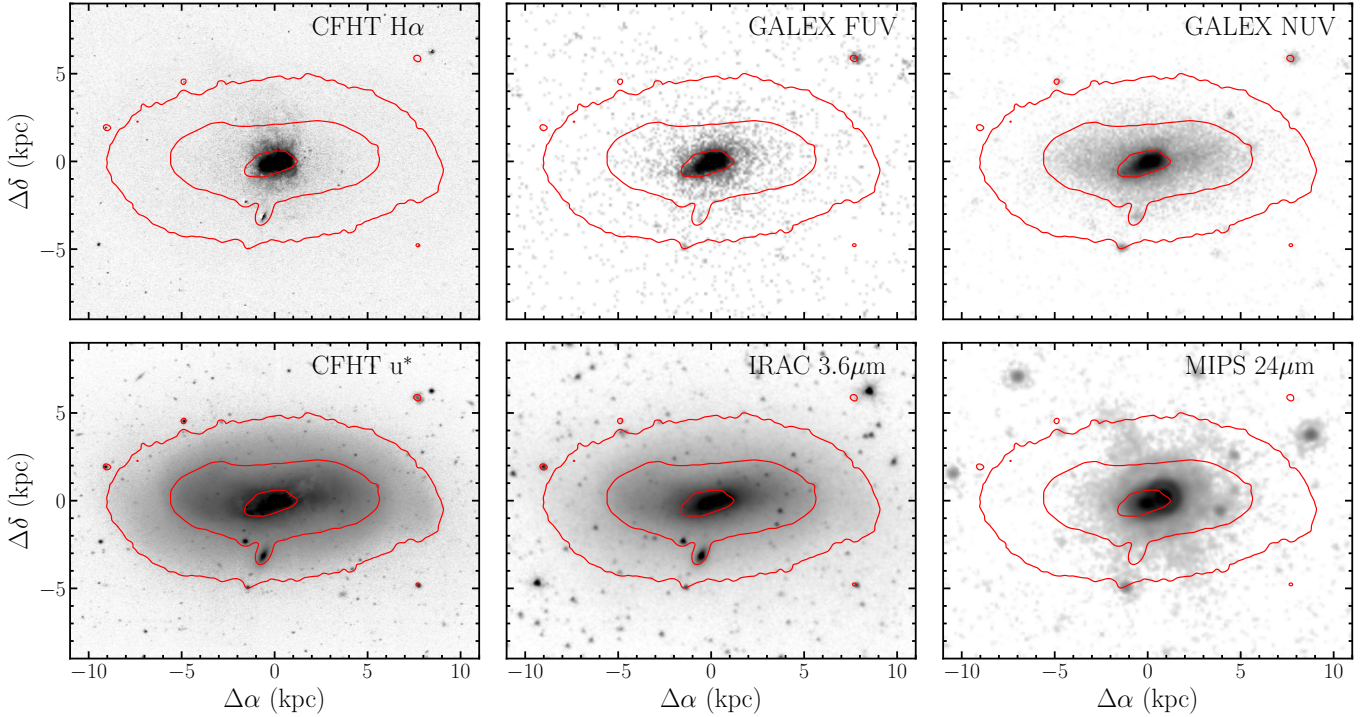
### 3.2. Image analysis

A first glance at the new images presented in this paper clearly shows the peculiar morphology of the galaxy at all wavelengths, although with some remarkable differences with respect to what was found in previous works. First of all, the very deep NGVS images (Figs. 1 and 3) do not show any of the fine structures generally interpreted as a proof of a recent merging event (shells, rings, plumes, tails) in the stellar component in the outer regions (e.g. Duc et al. 2011). We recall that these NGVS images were gathered and reduced using Elixir-LSB, a pipeline especially tailored to detect extended low surface brightness features (Ferrarese et al. 2012; Paudel et al. 2017). Fine structures in the stellar component are not detected down to a surface brightness limit of  $\mu_g \approx 29 \text{ mag arcsec}^{-2}$ . The excellent HST image of the inner region not only confirms a perturbed morphology, but also shows plumes and filaments of dust in absorption located along the major axis of the galaxy or perpendicular to it (Fig. 4). A peculiar distribution of the dust emission is also revealed for the first time by the  $24 \mu\text{m}$  MIPS/Spitzer image (Fig. 3). This figure shows a  $\approx 5 \text{ kpc}$  filament (projected length) extending from the nucleus to the north-north-east. The  $3.6 \mu\text{m}$  IRAC/Spitzer image (Fig. 3), sensitive to the distribution of the old stellar population, shows an elongated feature slightly misaligned with respect to

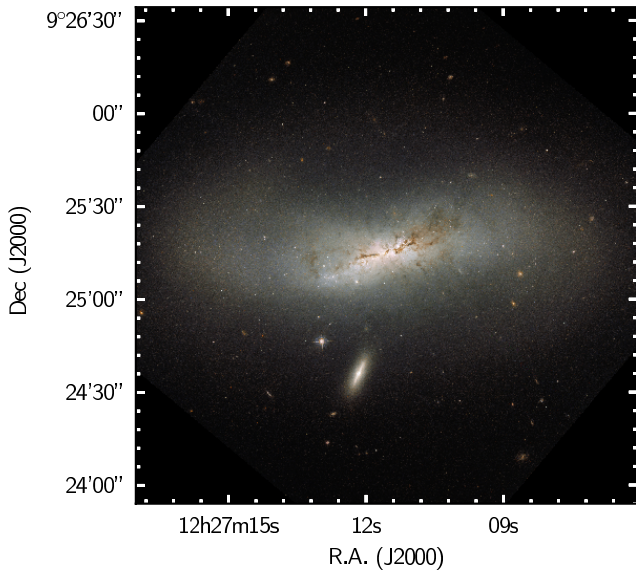
the major axis, with a prominent compact structure at the western edge and another in the centre.

Thanks to their excellent quality in terms of sensitivity and angular resolution, the very deep VESTIGE continuum-subtracted NB H $\alpha$  image shows several features never observed before:

- (1) a spectacular tail of ionised gas in the north-east  $\approx 10 \text{ kpc}$  long (projected length) extending from the central star forming region (Figs. 1 and 2). This feature, which corresponds to the one detected at  $24 \mu\text{m}$  (Fig. 3) and with a typical surface brightness of  $\Sigma(\text{H}\alpha) \approx 4 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ , is diffuse and does not contain clumpy spots indicating the presence of HII regions. The star formation activity is present only in the central  $\approx 3 \text{ kpc}^2$  region and is mainly located along the thin elongated structure visible in the near-infrared, and in particular at its edges. There are, however, a few blobs of star formation south to the nucleus outside the stellar disc;
- (2) filaments of ionised gas perpendicular to the star forming bar extending up to  $\approx 1.5 \text{ kpc}$ ;
- (3) a low surface brightness ( $\Sigma(\text{H}\alpha) \approx 4 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ ) diffuse halo extending along the major axis of the galaxy up to  $\approx 6\text{--}7 \text{ kpc}$  from the nucleus. The lack of any structured feature suggests that this emission is due to the diffuse gas ionised by the radiation escaping from the nuclear starburst or that it is gas shock-ionised by the ram pressure stripping event (see Sect. 4.1.2). The activity of star formation in the outer disc is thus completely quenched.



**Fig. 3.** Multifrequency images on the same scale of the galaxy NGC 4424. *Upper panels, from left to right:* continuum-subtracted  $H\alpha$  (VESTIGE), FUV (GALEX), NUV (GALEX); *lower panels:*  $u$  (NGVS), IRAC  $3.6\mu\text{m}$  (*Spitzer*), MIPS  $24\mu\text{m}$  (*Spitzer*). Red contours are drawn from the  $u$ -band image at 22nd, 24th, and 26th mag arcsec $^{-2}$  and highlight the truncation of the disc at different wavelengths.



**Fig. 4.** HST colour image of the galaxy NGC 4424, showing a perturbed morphology at an unequal sensitivity and angular resolution, with plumes and filaments of dust in absorption located along the major axis of the galaxy or perpendicular to it.

## 4. Physical parameters

### 4.1. Narrow-band imaging

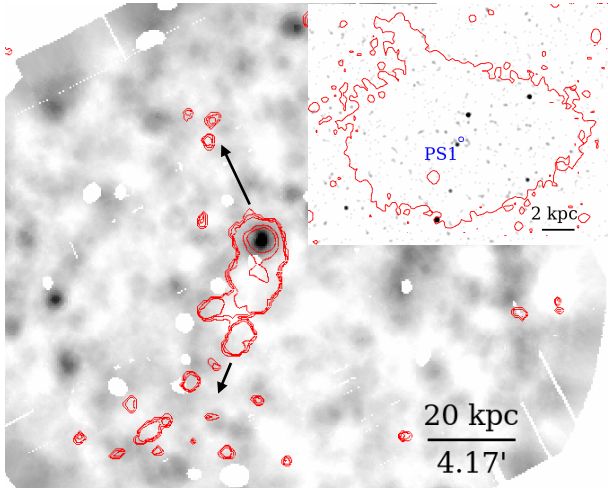
#### 4.1.1. Galaxy

The VESTIGE data can be used to estimate the total  $H\alpha$  luminosity of the galaxy and of the tail of ionised gas detected in the north-eastern direction. The total  $H\alpha$  flux of the galaxy including its extended tail is  $\log f(H\alpha + [\text{NII}]) = -12.12 \pm$

$0.02 \text{ erg s}^{-1} \text{ cm}^{-2}$ , where the uncertainty on the flux was measured after deriving the uncertainty from the sky background on an aperture similar to the size of the galaxy in 100 regions randomly located on the frame around the target. Assuming a typical  $[\text{NII}]\lambda 6583 \text{ \AA}/H\alpha = 0.36$  and a Balmer decrement  $H\alpha/H\beta = 4.45$  as measured from the MUSE data in the central field (see Sect. 4.3), the total extinction-corrected  $H\alpha$  luminosity of the galaxy is  $L(H\alpha) = 4.50 (\pm 0.21) \times 10^{40} \text{ erg s}^{-1}$ , which corresponds to a  $\text{SFR} = 0.25 \pm 0.04 M_{\odot} \text{ yr}^{-1}$  assuming the escape fraction and the fraction of ionising photons absorbed by dust before ionising the gas to be zero (Boselli et al. 2009), the Kennicutt (1998)  $\text{SFR} - L(H\alpha)$  calibration, and a Chabrier (2003) IMF.

#### 4.1.2. Diffuse tail

We also measured the total  $H\alpha$  emission in the tail of ionised gas within a box of size  $77 \times 32 \text{ arcsec}^2$ ,  $\text{PA} = -30 \text{ deg}$ , centred at  $\text{RA} = 12:27:13.56$ ;  $\text{Dec} = +9:26:17.1$  obtaining  $\log f(H\alpha + [\text{NII}]) = -13.85 \pm 0.19 \text{ erg s}^{-1} \text{ cm}^{-2}$ . The lack of IFU data in this low surface brightness region prevents us from having an accurate estimate of the  $[\text{NII}]$  contribution and of the Balmer decrement. The BPT diagram analysed in Sect. 4.3 indicates that the  $[\text{NII}]\lambda 6583 \text{ \AA}/H\alpha$  ratio increases in the outer regions of the galaxy where it reaches  $[\text{NII}]\lambda 6583 \text{ \AA}/H\alpha \approx 1.0$  probably because of the shock-ionised gas. The MUSE map of the Balmer decrement indicates that in these outer regions  $H\alpha/H\beta \approx 4.5$  ( $A(H\alpha) \approx 1.2 \text{ mag}$ ). Assuming that these ratios are also valid in the diffuse extraplanar tail, where the gas is supposedly shock-ionised and dust might be present as indicated by the extended  $24\mu\text{m}$  emission, we estimate that the total  $H\alpha$  luminosity of the northern diffuse component is  $L(H\alpha) \approx 5.34 (\pm 2.34) \times 10^{38} \text{ erg s}^{-1}$ . Using relation (2) in Boselli et al. (2016a), and assuming a filling factor  $f = 0.1$  and a gas distributed within a cylinder of

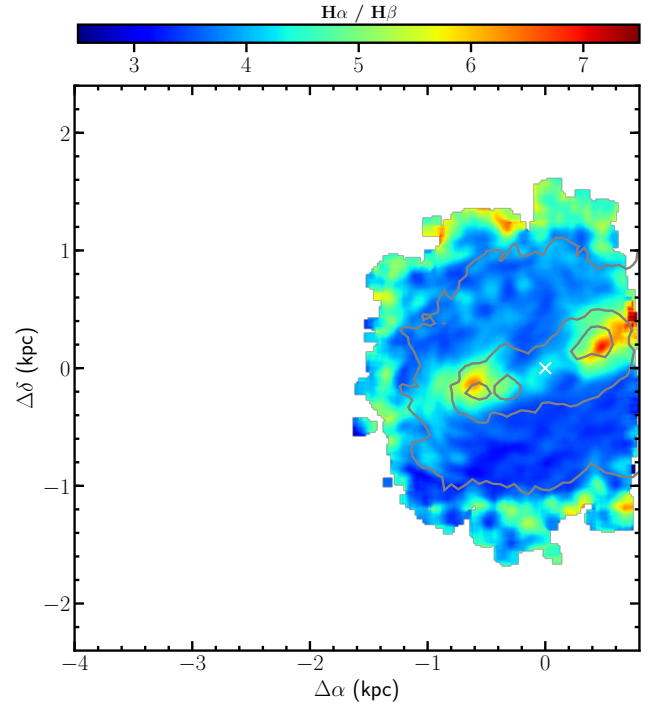


**Fig. 5.** Adaptively smoothed 0.4–1.3 keV XMM image of the NGC 4424 with the HI contours (from Chung et al. 2009) superposed. Bright X-ray point sources outside of the galaxy are masked. No soft X-ray enhancement is observed in the position of the HI tail. The two arrows at the north and south of the galaxy point towards the centre of Virgo cluster A and B, respectively. Their length is proportional to the distance to M 87 and M 49. The zoomed-in region at the upper right shows the *Chandra* 0.5–7 keV image of the galaxy, with the  $H\alpha$  contours (as in Fig. 2) superposed. The small blue circle with a radius of 2'' shows the putative nuclear position without a corresponding X-ray source. An obscured X-ray source 4.9'' to the south-east of the nucleus (PS1) is detected by *Chandra*.

diameter ( $\approx 2.6$  kpc) and height ( $\approx 7$  kpc) similar to the deprojected rectangular box used for the flux extraction in the tail, we obtain a rough estimate of the typical gas density in the tail,  $n_e \approx 7 \times 10^{-2} \text{ cm}^{-3}$ , with a corresponding total ionised gas mass  $M_{\text{tail}}(H\alpha) \approx 6.0 \times 10^6 M_\odot$ , to be compared to the HI mass in the tail,  $M_{\text{tail}}(\text{HI}) \approx 5.1 \times 10^7 M_\odot$  (Sorgho et al. 2017). We can also estimate the typical recombination timescale for the diffuse gas by using Eq. (5) in Boselli et al. (2016a):  $t_{\text{rec}} \approx 1.4$  Myrs.

#### 4.2. X-ray imaging

Figure 5 shows the XMM and *Chandra* images of NGC 4424. The putative nucleus is undetected by *Chandra*. Assuming a power law model with a photon index of 1.7, the  $3\sigma$  upper limit of the 2–10 keV luminosity is  $5.6 \times 10^{37} \text{ erg s}^{-1}$  for an absorption column density of  $1.7 \times 10^{20} \text{ cm}^{-2}$  (Galactic absorption). As a significant amount of intrinsic absorption is expected in NGC 4424, the  $3\sigma$  upper limit becomes  $6.6 \times 10^{37} \text{ erg s}^{-1}$  and  $1.5 \times 10^{38} \text{ erg s}^{-1}$  for an absorption column density of 10 times and 100 times greater than the Galactic value, respectively. The expected super massive black hole (SMBH) mass of NGC 4424 derived using the Marconi & Hunt (2003) relation is  $\sim 2.4 \times 10^7 M_\odot$ . Thus, in NGC 4424 the SMBH is quiescent and has an Eddington ratio of  $< 10^{-6}$ . There is indeed an X-ray source 4.9'' to the south-east of the nucleus (PS1 in Fig. 5). It is only detected in 2–7 keV, suggesting that it is obscured. Its X-ray hardness ratio is consistent with an absorption column density greater than  $2.5 \times 10^{22} \text{ cm}^{-2}$ . The 2–10 keV luminosity is  $4.3 \times 10^{38} \text{ erg s}^{-1}$  for an absorption column density of  $2.5 \times 10^{22} \text{ cm}^{-2}$  and  $1.1 \times 10^{39} \text{ erg s}^{-1}$  for an absorption column density of  $10^{23} \text{ cm}^{-2}$ . If it is associated with NGC 4424, this point source can be an X-ray binary (almost an ultraluminous X-ray source in that case) or the nucleus of a merged galaxy.



**Fig. 6.** Spaxel distribution of the Balmer decrement  $H\alpha/H\beta$  derived from the MUSE pixels with  $S/N > 5$ . The cross indicates the position of the nucleus, the grey contours the distribution of the  $H\alpha$  emitting gas (contour levels are  $\Sigma(H\alpha) = 5 \times 10^{-17} - 10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ ).

As shown in Fig. 5, the HI tail is undetected in X-rays. The  $3\sigma$  upper limit of the 0.4–1.3 keV emission in the HI tail can be estimated from the XMM data, and is  $4.1 \times 10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcmin}^{-2}$  and  $4.8 \times 10^{-16} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcmin}^{-2}$  for an assumed gas temperature of 0.6 keV and 0.3 keV, respectively. The resulting X-ray bolometric luminosity of the tail is thus at least 370–700 times smaller than that of ESO 137-001 (Sun et al. 2010). Although the mixing of the cold gas stripped from the galaxy with the hot ICM enhances the soft X-rays emission within the tail, the ambient pressure, which in the position of NGC 4424 is  $\sim 100$  smaller than around ESO 137-001 (see Sect. 5.4), is not sufficient to make the tail detectable with the existing X-ray observations (Tonnesen et al. 2011).

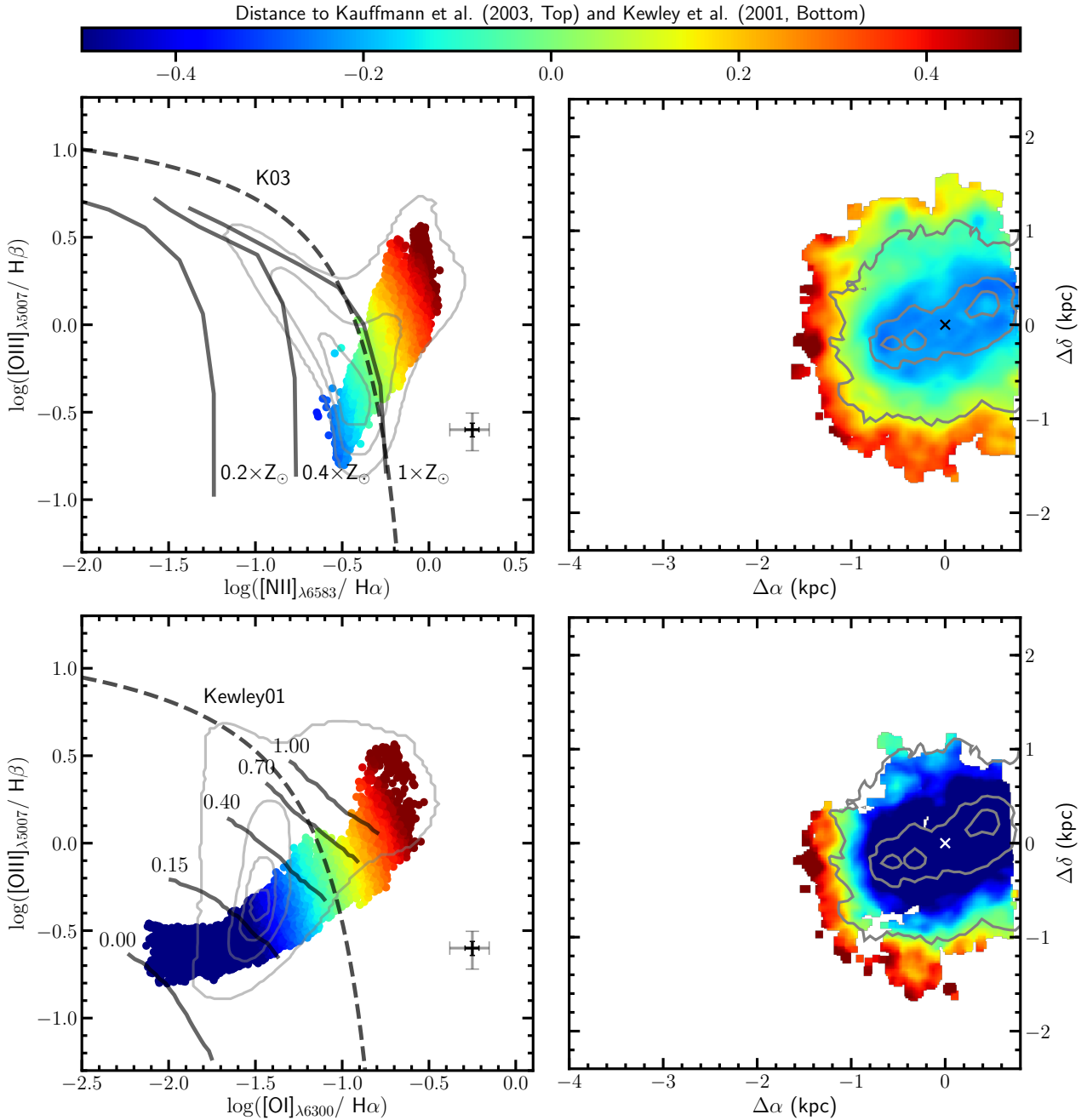
#### 4.3. Spectroscopy

##### 4.3.1. Line ratios

The MUSE data can be used to derive the typical properties of the ISM. The Balmer decrement shown in Fig. 6 traces the dust attenuation in the Balmer lines. The  $H\alpha$ -to- $H\beta$  ratio ranges from  $\approx 7$  in the western peak of  $H\alpha$  emission ( $A(H\alpha) \approx 2.3$  mag) to  $\approx 3$  in the diffuse gas ( $A(H\alpha) \approx 0.1$  mag). It is worth noting that the eastern peak of  $H\alpha$  emission does not match the peak of dust attenuation which is located further out ( $\approx 0.3$  kpc) along the elongated structure dominating the ionised gas emission.

The  $[\text{SII}]\lambda 6716/6731$  line ratio, a direct tracer of the gas density, is always fairly high, with the lower values found on the two bright  $H\alpha$  nuclei, where  $[\text{SII}]\lambda 6716/6731 \approx 1.2$ –1.3, corresponding to gas densities  $n_e \approx 150 \text{ cm}^{-3}$ , while  $[\text{SII}]\lambda 6716/6731 \approx 1.4$ –1.5 in the diffuse gas ( $n_e \lesssim 30 \text{ cm}^{-3}$ ) (Osterbrock et al. 2006; Proxauf et al. 2014). Line measurements can also be used to derive the metallicity of the gas. Using the calibration of Curti et al. (2017) based on the  $[\text{OIII}]\lambda 4959/H\beta$ ,



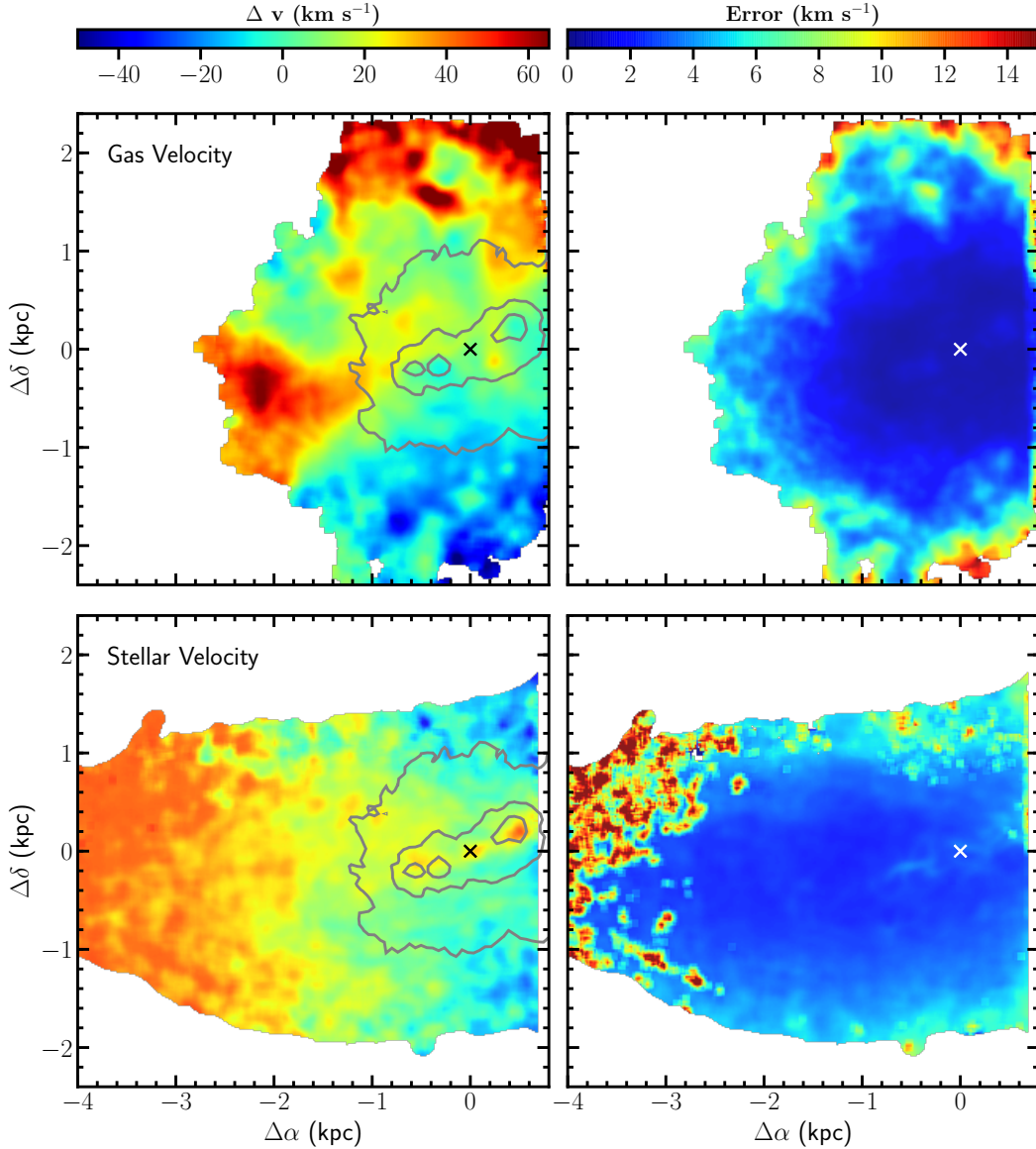


**Fig. 7.** Line diagnostic BPT diagrams  $\log([\text{OIII}]_{\lambda 5007}/\text{H}\beta)$  vs.  $\log([\text{NII}]_{\lambda 6583}/\text{H}\alpha)$  (*upper left panel*) and  $\log([\text{OIII}]_{\lambda 5007}/\text{H}\beta)$  vs.  $\log([\text{OI}]_{\lambda 6300}/\text{H}\alpha)$  (*lower left panel*) derived from MUSE pixels with  $S/N > 5$ . The different points are colour-coded according to their distance from the dashed lines defined by [Kauffmann et al. \(2003; upper panels\)](#) and [Kewley et al. \(2001; lower panels\)](#) to separate AGN from star forming regions. The grey contours show the distribution of a random sample of nuclear spectra of SDSS galaxies in the redshift range 0.01–1 and stellar mass  $10^9$ – $10^{11}$ . The solid thick lines in the *upper left panel* show the expected behaviour of star forming regions as derived from the photo-ionisation models of [Kewley et al. \(2001\)](#) for three different metallicities (0.2, 0.4,  $1 Z_{\odot}$ ). *Lower left panel*: shock models of [Rich et al. \(2011\)](#) for increasing shock fraction (from left to right) in a twice solar gas. *Right panels*: distribution of the same points colour-coded as in the left panels over the disc of the galaxy. The cross shows the centre of the galaxy, the grey contours the distribution of the  $\text{H}\alpha$  emitting gas.

$[\text{NII}]$ , and  $\text{H}\alpha$  lines, we obtain typical metallicities ranging from  $12 + \log \text{O}/\text{H} \approx 8.65$  in the outer gas to  $12 + \log \text{O}/\text{H} \approx 8.8$  in the inner regions. These values are comparable to those derived for galaxies of similar stellar mass (e.g. [Hughes et al. 2013](#)).

The spectroscopic data can also be used to derive several line diagnostic diagrams such as those proposed by [Baldwin et al. \(1981; BPT diagrams\)](#). These diagrams, shown in Fig. 7, can be used to identify the dominant ionising source of the gas. Figure 7

shows that the gas is photo-ionised by the young stars in the inner regions, while it is shock-ionised in the outer ones. The high cut in the  $S/N$  ratio in all the emission lines used in the two diagrams ( $S/N > 5$ ) and the agreement in the  $[\text{OIII}]/\text{H}\beta$  versus  $[\text{NII}]/\text{H}\alpha$ ,  $[\text{OIII}]/\text{H}\beta$  versus  $[\text{OI}]/\text{H}\alpha$ , and  $[\text{OIII}]/\text{H}\beta$  versus  $[\text{SII}]/\text{H}\alpha$  (not shown here) BPT diagrams secure this result against systematic effects in the data. Figures 7 and 5 do not show the presence of any hard ionising source near the nucleus, suggesting a very low



**Fig. 8.** Velocity field (left panels) and error maps (right panels) of the gaseous (upper panels) and stellar (lower panels) components derived from the MUSE data. The black and white crosses indicate the position of the nucleus, the grey contours the distribution of the H $\alpha$  emitting gas.

accretion rate despite the abundant cold material in the nuclear region (Cortés et al. 2006).

#### 4.3.2. Stellar kinematics

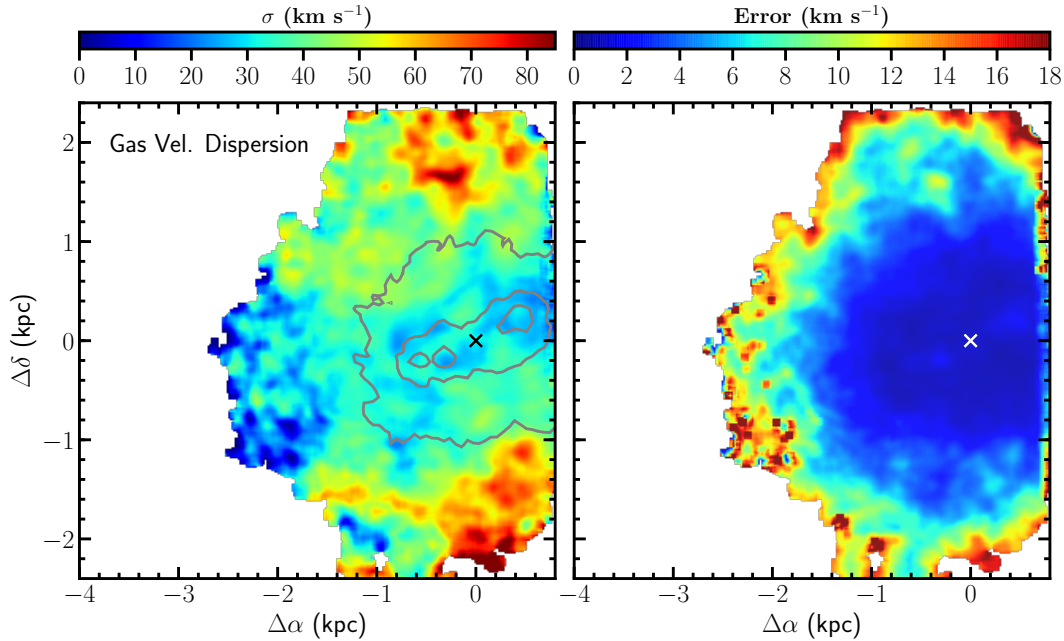
Figure 8 shows the stellar kinematics derived using the MUSE data (lower left panel). The stellar velocity field of NGC 4424, although showing a gradient from the nucleus to the eastern regions<sup>3</sup> of  $\approx 40 \text{ km s}^{-1}$ , does not have the typical shape of a rotating disc for a spiral galaxy inclined  $\sim 60$  deg on the plane of the sky with a steep gradient in the inner regions. Furthermore, this rotation ( $vel \approx 85 \text{ km s}^{-1}$  when corrected for inclination) is significantly smaller than that expected for a rotating system of similar stellar mass ( $M_{\text{star}} \approx 10^{10.17} M_{\odot}$ ;  $vel \approx 200 \text{ km s}^{-1}$ ) as indeed noticed by Kenney et al. (1996), Rubin et al. (1999), Coccato et al. (2005), and Cortés et al. (2015). The superior

<sup>3</sup> The velocity does not seem to increase in the eastern disc mapped by the second field of MUSE. The low  $S/N$  in this outer region prevents us from deriving an accurate stellar velocity field.

quality of the MUSE data allows us to reconstruct the stellar velocity field down to one arcsec angular resolution ( $\approx 80 \text{ pc}$ ). This stellar velocity field shows a kinematically decoupled structure on the elongated feature emitting in H $\alpha$  and in the nucleus of the galaxy not seen in previous data. This feature suggests that the stellar component associated with the nuclear starburst has a different origin than the outer disc.

#### 4.3.3. Gas kinematics

Figure 8 also shows the velocity of the gas derived using the MUSE data (upper panel), while its velocity dispersion is shown in Fig. 9. The MUSE data suggest that the gas is rotating at large scales along the major axis as the stellar component, although with a slightly steeper gradient. We expect that the same rotation of the two components is present on the western side where MUSE data are unfortunately unavailable, as indeed suggested by the CO velocity field of Cortés et al. (2006). As for the stellar component, in the inner region we do not see any steep gradi-



**Fig. 9.** Velocity dispersion (*left panel*) and error maps (*right panel*) of the gaseous component derived from the MUSE data. The cross indicates the position of the nucleus, the grey contours the distribution of the H $\alpha$  emitting gas.

ent in the gas component along the major axis, as expected in a massive unperturbed system. The two dominant H $\alpha$  emitting regions located to the east and west of the stellar nucleus are at approximately the same line-of-sight velocity ( $\approx -10 \text{ km s}^{-1}$ , see Fig. 10), suggesting that the elongated structure to which they belong is not a rotating bar. Between these two H $\alpha$  emitting regions there is a rather narrow (50 pc) north–south filament in the velocity field displaying a higher mean line-of-sight velocity ( $\approx 30 \pm 10 \text{ km s}^{-1}$ ) surrounded by velocities closer to those of the dominant H $\alpha$  blobs than to that of the filament. This filament is also present in the MUSE data. The kinematical data suggest that this high-velocity filament might be in a different plane than the low-velocity component embedding the two dominant H $\alpha$  emitting regions.

Figures 8 and 10 consistently indicate that there is a systematic difference in velocity between the northern ( $\approx +50 \text{ km s}^{-1}$ ) and the southern regions ( $\approx -30 \text{ km s}^{-1}$ ), suggesting the presence of a second gas component. Figure 9 also indicates that in these outer regions the velocity dispersion of the gas is fairly high ( $\approx 80 \text{ km s}^{-1}$ ). This second component might be gas rotating along the minor axis (Cocato et al. 2005), as well as gas infalling or outflowing from the nucleus, depending on the orientation of the galaxy disc on the plane of the sky. The line-of-sight overlapping of the minor axis component on the major axis gas could explain why in the inner regions we do not observe a velocity gradient along the major axis. The high spectral resolution Fabry-Perot data ( $\approx 5 \text{ km s}^{-1}$ ) do not show multiple profiles, suggesting that the two gaseous components are well mixed within this inner region.

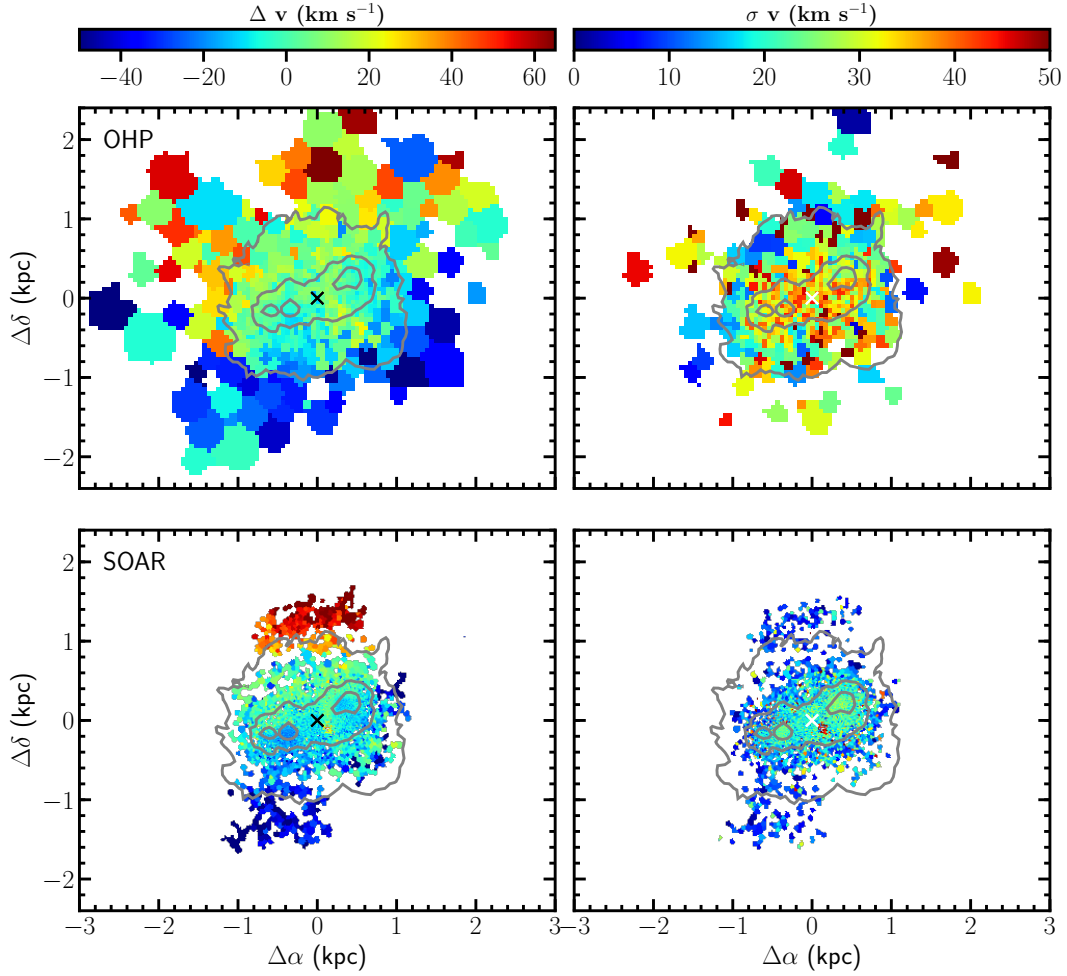
## 5. Discussion

### 5.1. SED fitting

The H $\alpha$  image of NGC 4424 obtained during the VESTIGE survey clearly indicates that the activity of star formation is present only within the inner  $\sim 1.5 \times 1.5 \text{ kpc}^2$  region, while it is totally quenched in the outer disc. Figure 3 also shows that the colour

gradually increases with galactic radius (the size of the stellar disc increases with wavelength). The lack of diffuse dust in the outer regions, shown by the  $24 \mu\text{m}$  image, indicates that the radial reddening of the disc is not due to dust attenuation, but is instead produced by the ageing of the stellar population. The VLA HI data suggest that the atomic gas in the outer disc was removed during a ram pressure episode. The lack of gas has induced the decrease in star formation activity. We use the spectro-photometric data to measure the epoch when the activity was quenched by the interaction with the ICM. This is done by dating the last episode of star formation in the outer disc. To this end, we fitted the FUV to far-IR SED of the galaxy extracted in three distinct  $0.5 \times 0.5 \text{ arcmin}^2$  regions ( $2.5 \times 2.5 \text{ kpc}^2$ ) located along the major axis of the galaxy, as depicted in Fig. 11. These regions were defined to be representative of the inner star forming region and of the outer quenched disc. Fluxes and uncertainties were measured within these regions using the same procedures described in Fossati et al. (2018) and Boselli et al. (2018b).

To quantify the time elapsed since the truncation of the star formation activity we followed the same procedure given in Boselli et al. (2016b), i.e. we used an abruptly truncated star formation law and the following free parameters: the rotational velocity of the galaxy, the quenching age QA (epoch when the quenching of the star formation activity started to occur), and the quenching factor QF (QF = 0 for unperturbed SFR, QF = 1 for a totally quenched SFR). The SED of the galaxy was fitted with CIGALE (Noll et al. 2009; Boquien et al. 2018) using Bruzual & Charlot (2003) stellar population models derived with a Chabrier IMF for the stellar continuum and the Draine & Li (2007) models for the dust emission. To represent the observed SED, we used the two GALEX FUV and NUV bands, the optical NGVS and VESTIGE *ugriz* bands, a Calar Alto near-infrared observation in the *H* band, the four IRAC bands, the MIPS  $24 \mu\text{m}$  image, the PACS images at  $100$  and  $160 \mu\text{m}$ , and the SPIRE  $250$  and  $350 \mu\text{m}$  bands (17 photometric bands). We excluded from the analysis the WISE bands, the MIPS bands at longer wavelengths, and the SPIRE  $500 \mu\text{m}$  band for their lower sensitivity and/or lower angular resolution to avoid any possible contamina-

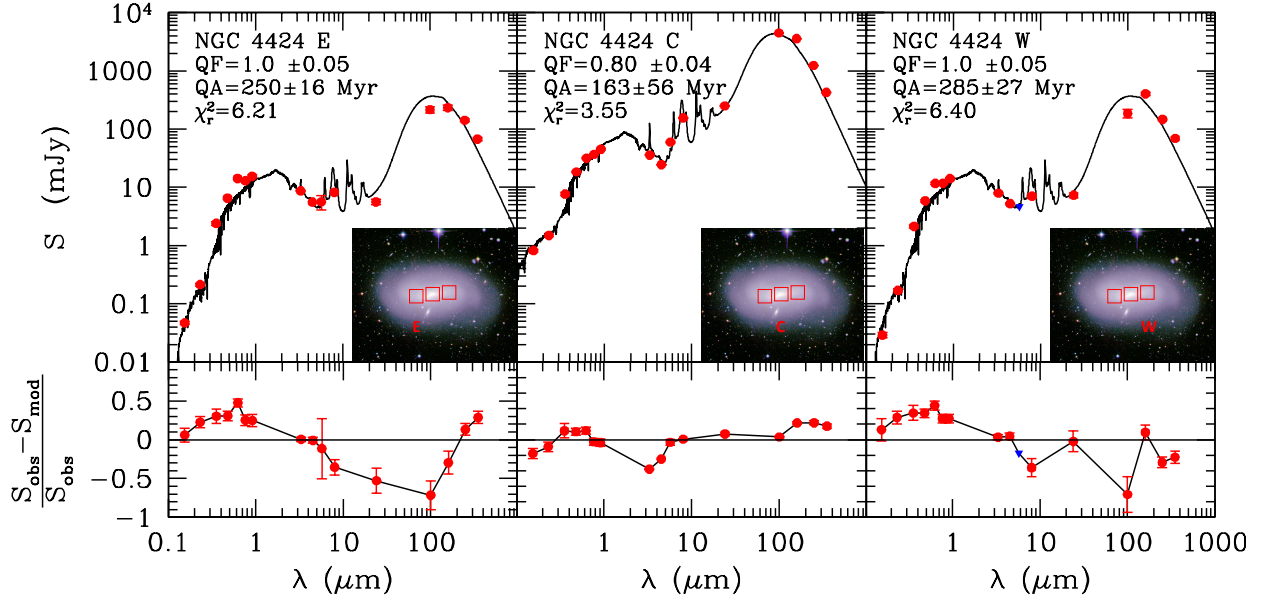


**Fig. 10.** Velocity field (*left panels*) and velocity dispersion (*right panels*) of the gaseous component as derived from the OHP (*upper panels*) and SOAR (*lower panels*) Fabry-Perot observations. The black and white crosses indicate the position of the nucleus, the grey contours the distribution of the  $H\alpha$  emitting gas.

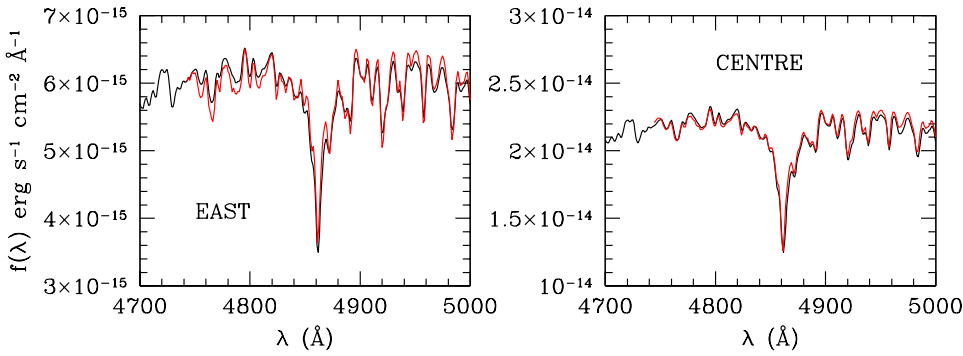
tion within the selected regions. Given their size, we did not apply any aperture correction. As in [Boselli et al. \(2016b\)](#) we included the  $H\alpha$  emission by measuring the number of ionising photons from the VESTIGE image. The observed  $H\alpha + [\text{NII}]$  flux measured within the central region was corrected for dust attenuation and  $[\text{NII}]$  contamination using the mean values derived from the MUSE data (see Sect. 4.3). We also included in the fit the age-sensitive  $H\beta$  absorption line ([Poggianti & Barbaro 1997](#)) by measuring its contribution within a pseudo-filter as in [Boselli et al. \(2016b\)](#). This was done using a light-weight mean spectrum derived from the MUSE data cube in the central and eastern regions.

The results of the fit are shown in Figs. 11 and 12 and summarised in Table 2. The quality of the fit is satisfactory in the two outer regions, and good in the centre of the galaxy. The best fit models perfectly match the age-sensitive  $H\beta$  absorption line in the eastern and central regions where MUSE data are available. The fitted model indicates that the activity of star formation was totally quenched ( $\text{QF} = 1$ ) in the outer disc  $\approx 250\text{--}280$  Myr ago. These numbers are consistent with those derived by [Crowl & Kenney \(2008\)](#) using Balmer absorption indices extracted from IFU spectroscopy or the GALEX FUV-NUV colour index in the outer disc in the eastern direction ( $275 \pm 75$  Myr), although with a much better accuracy thanks to our larger number of photometric (17) and spectroscopic (2)

bands and to the more refined SED fitting technique. To test the reliability of the output parameters of the fit, and in particular of the quenching ages and quenching factors given above, following [Boselli et al. \(2016b\)](#) we created a mock catalogue by introducing some extra noise into the observed data, randomly distributed according to a Gaussian curve of dispersion similar to the median error in each band, and re-fitted the data using the SED models. The output of the fit on the mock catalogue gives again  $\text{QF} = 1$  and QA values consistent within  $\approx 6\%$ , confirming that the determination of these parameters is robust. The fit also suggests that, in the central region, the activity of star formation was reduced by 79% approximately 180 Myr ago. The mock analysis suggests that while in this region QF is fairly well constrained, the uncertainty on the quenching age QA is very large. Furthermore, this value should be considered with caution given that the parametrised star formation history used to trace the evolution of NGC 4424 is not optimal for the inner region, which has probably suffered from two competitive mechanisms, a quenching episode (during the ram pressure stripping event) and a starburst (during the merging event) which might last 10–100 Myr ([Cortijo-Ferrero et al. 2017](#)). This very peculiar and changing star formation history is reproducible with difficulty using a simple SED fitting technique like the one adopted for the outer disc. To conclude, the overall star formation activity of NGC 4424 was significantly reduced after the gas stripping



**Fig. 11.** Far-UV to far-IR SED of the three selected regions along the disc of NGC 4424 (from east to west). In each panel the observational data and their error bars are indicated with red filled dots, the upper limit with a blue triangle; the black solid line shows the best fit model. The inset shows the colour *gri* RGB image of the galaxy; overlaid are the red boxes indicating the three fields used to study the SED with the CIGALE fitting code.



**Fig. 12.** Best model fit obtained with CIGALE (black) is compared to the MUSE spectrum (red) in the  $H\beta$  line after emission is removed using GANDALF. The best model has been shifted on the y-axis to match the observed spectrum.

**Table 2.** Results of the SED fitting

Region	RA(J2000)	Dec	QF	QA (Myr)	SFR ( $M_{\odot} \text{ yr}^{-1}$ )
West <sup>a</sup>	12:27:09.001	+9:25:19.06	1	$281 \pm 27$	0
Centre	12:27:11.478	+9:25:15.27	$0.79 \pm 0.04$	$179 \pm 68$	$0.13 \pm 0.02$
East	12:27:13.955	+9:25:10.70	1	$251 \pm 14$	0

**Notes.** <sup>(a)</sup>Fit done without the  $H\beta$  pseudo filter.

episode, moving the galaxy below the main sequence relation of unperturbed objects (Boselli et al. 2016b).

## 5.2. Ram pressure stripping

The presence of a long ( $\sim 110$  kpc in projected distance) symmetric HI tail of gas is clear evidence that the galaxy is undergoing a ram pressure stripping event (Chung et al. 2007; Sorgho et al. 2017). The relative orientation of the galaxy and the position of the tail on the plane of the sky combined with a relatively low recessional velocity with respect to the mean velocity of the cluster ( $\langle \text{vel}_{\text{cluster A}} - \text{vel}_{\text{N4424}} \rangle \sim 500 \text{ km s}^{-1}$ ) also indicate that the stripping event is occurring almost face-on,

making it particularly efficient at removing the gas from the outer regions producing truncated gaseous discs (Quilis et al. 2000; Roediger & Brüggén 2006, 2007; Tonnesen & Bryan 2009). Given its accurate distance estimate of  $\approx 15.5\text{--}16$  Mpc (Munari et al. 2013; Cantiello et al. 2018; Hatt et al. 2018), the galaxy is now in the foreground of the cluster (16.5 Mpc). It has thus probably already crossed the southern envelopes of Virgo cluster A coming from the south-eastern region behind the cluster and moving in the north-west direction in front of it. The lack of gas quenches the activity of star formation outside-in, also producing truncated discs in the young stellar populations (Boselli et al. 2006; Fossati et al. 2018). The SED fitting analysis presented in the previous section indicates that the quenching

episode is relatively recent ( $\approx 250\text{--}280$  Myr ago). This timescale can be compared with the time necessary for the galaxy to travel within the cluster to form an HI tail of  $\approx 110$  kpc. Assuming the typical velocity dispersion measured within Virgo cluster A ( $\sigma_{\text{cluster A}} = 800 \text{ km s}^{-1}$ ; Boselli et al. 2014c), this timescale is  $t_{\text{HI}} \approx 140$  Myr, just a factor of  $\sim 2$  shorter than the typical quenching time. These numbers are fairly consistent considering that (i) the total extension of the HI tail can be underestimated for sensitivity reasons and (ii) the cold atomic gas, once stripped and in contact with the hot intracluster medium can change phase becoming ionised gas or hot plasma via heat conduction, as indeed observed in the majority of cluster galaxies (e.g. Boselli & Gavazzi 2014; Boselli et al. 2016a; Yagi et al. 2017; Gavazzi et al. 2018). Within the stellar disc the gas can be shocked-ionised by the stripping process, as observed in the outer regions of the MUSE field (Fig. 7). The significant difference between the HI mass in the tail measured by Sorgho et al. (2017;  $M_{\text{tail}}(\text{HI}) \approx 5.1 \times 10^7 M_{\odot}$ ) and that of the stripped gas derived from the HI-deficiency parameter ( $\text{HI} - \text{def} = 0.98$ ; Boselli et al. 2014a, thus  $M_{\text{stripped}}(\text{HI}) \approx 2 \times 10^9 M_{\odot}$ ) supports this interpretation.

### 5.3. Merging

The morphological and kinematical properties of NGC 4424 also suggest that the galaxy has undergone a more violent gravitational perturbation (Kenney et al. 1996; Cortés et al. 2006, 2008, 2015). Although the presence of fine structures witnessing an old major merging event is not confirmed in our NGVS data, the perturbed and decoupled kinematics of both the gas and stellar components unquestionably indicate an unrelaxed object. In this regard, we recall the decoupled kinematics of the stellar component along the elongated star forming structure with respect to the slowly rotating stellar disc revealed by MUSE. Simulations suggest that unequal-mass merger remnants might result in systems with a spiral-like morphology but elliptical-like kinematics (Bournaud et al. 2004), characteristics similar to those observed in NGC 4424: a grand-design spiral morphology (Fig. 1) and a low rotational velocity compared to the velocity dispersion of both the gas and the stellar components. The fact that the gas is still turbulent and not stabilised over the disc suggests that the merging event is relatively recent ( $t_{\text{merging}} \lesssim 500$  Myr). Considering that the typical crossing time of Virgo is  $\approx 1.7$  Gyr (Boselli & Gavazzi 2006), this means that the merging episode occurred once the galaxy was already within the Virgo cluster. We recall that in massive clusters such as Virgo ( $M_{\text{Virgo}} \approx 2\text{--}4 \times 10^{14} M_{\odot}$ , Paper I, and references therein) the high-velocity dispersion ( $\sigma_{\text{Virgo}} \approx 800 \text{ km s}^{-1}$ ) makes merging events among intermediate-mass galaxies very unlikely (Boselli & Gavazzi 2006).

### 5.4. The nature of the northern ionised gas tail

The presence of a tail of ionised gas in the northern direction, thus in the direction opposite to the direction of the HI gas tail, is hard to explain in a simple ram pressure stripping scenario. The kinematics of the gas over the plane of the disc, as mentioned in Sect. 4.3, can be interpreted as being due to (i) a rotating gas ring, (ii) infalling, or (iii) outflowing gas. It is hard to discriminate between these different scenarios, in particular because the peculiar morphology of the galaxy prevents us from determining the orientation of the stellar disc on the plane of the sky. There are, however, a few arguments suggesting that this ionised gas tail results from an outflow:

- (a) The morphology of the inner regions is similar to that observed in other starburst galaxies such as M82 or the Circinus galaxy: the  $\text{H}\alpha$  emitting gas shows prominent filaments perpendicular to the stellar disc (Fig. 2) and dust filamentary structures (Fig. 4) in the inner regions, while extended tails in  $\text{H}\alpha$  (Fig. 2) and in the mid-IR (Fig. 3) on large scales (up to 10 kpc) are also present in these iconic objects (Elmouttie et al. 1998a,b; Shopbell & Bland-Hawthorn 1998; Alton et al. 1999; Wilson et al. 2000; Engelbracht et al. 2006).
- (b) The velocity dispersion of the gas increases from the centre of the galaxy along the minor axis, reaching values as high as  $\approx 80 \text{ km s}^{-1}$  at  $\approx 2$  kpc from the disc (e.g. Cresci et al. 2017; López-Cobá et al. 2017).
- (c) The galaxy has two extended radio continuum lobes and a magnetic field perpendicular to the stellar disc, interpreted by Vollmer et al. (2013) as evidence of outflow.
- (d) The flat radio continuum spectral index ( $\alpha = -0.26$ ; Vollmer et al. 2013) and the BPT diagram (Fig. 7) consistently indicate that the source of energy is a nuclear starburst. The energy of this starburst, although now limited ( $\text{SFR} = 0.25 M_{\odot} \text{ yr}^{-1}$ ), could have been sufficient to overcome the gravitational forces also thanks to the perturbed gravitational potential well of the galaxy and to eject matter in the direction perpendicular to the disc. Following Boselli et al. (2016a) we can make a rough estimate of the total energy (potential and kinetic) of the outflow,  $E_{\text{out}} \approx 1.7 \times 10^{55} \text{ erg s}^{-1}$ . This energy can be produced by  $1.7 \times 10^4$  supernovae, or a larger number if we assume a more realistic 1–10% energy transfer efficiency, associated with a nuclear star cluster of stellar mass  $\gtrsim 3.7 \times 10^6 M_{\odot}$  and to a star formation rate  $\gtrsim 3.7 M_{\odot} \text{ yr}^{-1}$ , a star formation rate comparable to that in the galaxy before the quenching episode (Sect. 5.1). As remarked in Boselli et al. (2016a), however, ram pressure can produce low-density superbubble holes and can thus supply extra energy to the outflow through Kelvin-Helmholtz instabilities (Roediger & Hensler 2005) and viscous stripping (Roediger & Brüggén 2008).
- (e) The presence of a nuclear outflow might decouple the rotation of the stellar disc from the kinematics of the gas which becomes turbulent, as indeed observed in blue compact galaxies (Östlin et al. 2004).
- (f) The recombination time of the gas within the tail is very low ( $t_{\text{rec}} \approx 1.4$  Myrs). It is thus unlikely that it remains ionised if it is gas orbiting along the minor axis or if it is infalling gas given the lack of any ionising source in these regions. In the case of an outflow, the gas can be ionised by the central starburst or by the diffuse hot gas emitting in X-rays via heat conduction. The BPT diagram shown in Fig. 7, however, indicates that at the periphery of the central star forming region detected by MUSE the gas is mainly shock-ionised probably by the ram pressure episode.
- (g) The outflow has sufficient kinetic energy to overcome the external pressure of the ICM ( $p_{\text{RP}}$ ) if

$$p_{\text{out}} \geq p_{\text{RP}}, \quad (1)$$

where  $p_{\text{out}}$  is the pressure within the outflow. A rough estimate of  $p_{\text{out}}$  can be derived as

$$p_{\text{out}} = \rho_{\text{out}} v_{\text{out}}^2 \approx \frac{M_{\text{out}}}{V_{\text{out}}} v_{\text{out}}^2, \quad (2)$$

where  $V_{\text{out}}$  is the volume,  $M_{\text{out}}$  the mass, and  $v_{\text{out}}$  the velocity of the outflow. As described in Sect. 4.1.2, the mass of the

outflow is  $M_{\text{out}} = 6 \times 10^6 M_{\odot}$  and its volume  $V_{\text{out}} = 3.9 \times 10^{66} \text{ cm}^3$ . We do not have any estimate of the velocity of the gas up to 10 kpc from the galaxy disc, but we can extrapolate it from the MUSE data (Sect. 4.3.2, Fig. 8): the gas velocity gradually increases from  $0 \text{ km s}^{-1}$  to  $\frac{50 \text{ km s}^{-1}}{\cos(i)}$  at  $\approx 2 \text{ kpc}$  from the galaxy disc, and might thus reach  $vel_{\text{out}} \approx \frac{250 \text{ km s}^{-1}}{\cos(i)} \approx 500 \text{ km s}^{-1}$  at the edge of the observed outflow (10 kpc). This velocity is close to the escape velocity of the galaxy, which we estimate as in Boselli et al. (2018b) at  $v_e \approx 700 \text{ km s}^{-1}$ . The mean pressure exerted by the outflow on the surrounding medium should thus be  $p_{\text{out}} \approx 7.5 \times 10^{-12} \text{ dyn cm}^{-2}$ . Considering the typical velocity of the galaxy within Virgo cluster A ( $vel = 800 \text{ km s}^{-1}$ , Boselli et al. 2014c), and estimating the density of the ICM in the position of NGC 4424 (at a physical distance of 1.3 Mpc from M 87) using *Planck* and *Suzaku* data from Simionescu et al. (2017), we obtain  $n_e(r_{\text{N4424}}) = 3 \times 10^{-5} \text{ cm}^{-3}$  and  $p_{\text{RP}} = 6.2 \times 10^{-13} \text{ dyn cm}^{-2}$ .

The accurate distance estimate (15.5–16.0 Mpc) and its relative velocity of  $\approx -500 \text{ km s}^{-1}$  with respect to M 87 place the galaxy in the foreground of the cluster and suggest that it has probably crossed the periphery of Virgo coming from the south-east and moving towards the north-west. If the  $H\alpha$  tail is an outflow, the upper edge of the disc is the one closest to us. The  $H\alpha$  tail, like the HI tail, is bent to the east because of the motion of the galaxy along the line of sight. The line-of-sight velocity of the HI gas in the tail is  $vel_{\text{tail}}(\text{HI}) = 453 \text{ km s}^{-1}$ , slightly higher than that of the HI gas over the stellar disc ( $vel_{\text{disc}}(\text{HI}) = 433 \text{ km s}^{-1}$ ; Sorgho et al. 2017), again suggesting that the gas was lost while the galaxy was approaching from the cluster background, as has indeed been observed in other ram pressure stripped galaxies (Fumagalli et al. 2014). Remarkably, the ionised gas in the  $z$ -plane in the south of the galaxy disc, as observed by MUSE and by the Fabry-Perot observations, has a lower line-of-sight velocity than the galaxy, again suggesting that the outflow overcomes the ram pressure stripping process also in this southern region.

The new result of this analysis is that despite a moderate nuclear star formation activity probably induced by the infall of fresh gas accreted after a recent merging event, the galactic winds produced in the star forming regions are able to overcome the ram pressure stripping process and drag the gas out from the disc of the galaxy from the nuclear regions even in the direction opposite to that of the motion within the cluster. This picture is fairly consistent with the prediction of cosmological simulations and semi-analytic models in a starvation scenario (e.g. Larson et al. 1980), where the removal of the hot halo once the galaxy becomes a satellite of a larger dark matter structure makes the feedback process more efficient in removing the cold gas from the galaxy disc and quenching the activity of star formation (e.g. De Lucia 2011). It must be noted, however, that as in NGC 4569 (Boselli et al. 2016a) the amount of ionised gas ejected by feedback is between one and two orders of magnitude lower than that removed by the ram pressure mechanism.

## 6. Conclusion

NGC 4424 is a peculiar galaxy located at  $\sim 0.6R_{200}$  from the Virgo cluster centre. New very deep NB  $H\alpha$  imaging data acquired during the VESTIGE survey have revealed the presence of a  $\sim 10 \text{ kpc}$  tail (projected distance) of ionised gas perpendicular to the stellar disc. This newly discovered tail is located in the direction opposite to the  $\sim 110 \text{ kpc}$  long HI tail observed

at 21 cm, and is generally interpreted as clear evidence of an ongoing ram pressure stripping event. The detailed analysis of this object based on multifrequency data covering the whole electromagnetic spectrum, from the X-rays to the radio continuum, including high- and medium-resolution Fabry-Perot and MUSE spectroscopic data, consistently indicate that the galaxy is indeed being subjected to an ongoing ram pressure stripping event able to remove the gas from the outer disc and quench the activity of star formation on timescales of  $\sim 250\text{--}280 \text{ Myr}$ . The galaxy, however, has also been perturbed by a recent merging event. This gravitational perturbation, as suggested by simulations, might have contributed to making the ram pressure stripping event more efficient than in other cluster galaxies. The tail of ionised gas observed in the opposite direction than the HI tail probably results from a nuclear outflow fed by the collapse of the gas in the central regions after the merging episode. Despite a moderate nuclear star formation activity, the outflow produced by the starburst is able to overcome the ram pressure stripping force and contribute to the exhaustion of the gas in the inner regions. This observational result is the first evidence that the feedback process generally invoked by cosmological simulations and semi-analytic models (e.g. De Lucia 2011) in a starvation scenario (e.g. Larson et al. 1980) might contribute to remove gas from galaxies once they become satellites of a larger dark matter structure. As observed in other cluster objects, however, the amount of gas expelled by feedback seems to be only a small fraction of that removed via ram pressure stripping, thus suggesting that feedback plays only a minor role in quenching the activity of star formation of galaxies in high-density environments.

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