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Dependency of the Dimethyldihydropyrene Photochromic Properties on the Number of Pyridinium Electron-Withdrawing Groups

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In honor of Olivier Poizat

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Supporting information for this article is given via a link at the end of the document.

Abstract: Photochromic dimethyldihydropyrenes substituted with electron-withdrawing pyridinium groups have shown an increase of photo-induced ring-opening efficiency and a light sensitivity that is red shifted relative to the unsubstituted compound. However, a recently synthesized tetrapyrindinium derivative showed a considerable decrease of the photo-isomerization quantum yield relative to the monopyridinium and bispyridinium derivatives. We provide a rationale for this unexpected photochemical behavior based on the comparative theoretical investigations of the relevant excited states of these systems. In particular, we found that the nature and order of the lowest two excited states depend on the number of pyridinium groups and on the symmetry of the system. While the lowest S_1 excited state is photo-active in the monopyridinium and bispyridinium derivatives, the photo-isomerizing state is S_2 in the reference unsubstituted compound and both S_1 and S_2 lead to isomerization in the tetrapyrindinium derivative, albeit with a low efficiency. In the latter derivative, the photo-isomerization is hindered by the particular S_1/S_2 conical intersection topology.

Introduction

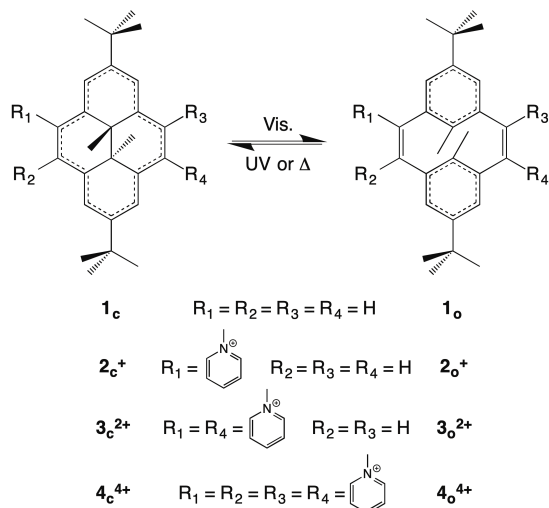
Besides being a clean and inexhaustible energy source, light is now envisioned as a powerful way to control the physical properties of materials, for applications in micro and nanotechnologies. Among them, molecular systems exhibiting switchable optical properties, emission and excitation fluorescence and electronic communication by light are among the most promising topics in the field of nanosciences.^[1]

During the last decade, the dimethyldihydropyrene (DHP) – cyclophanediene (CPD) photochromic couple has gained an increasing interest because of its original behavior.^[2] In particular it belongs to the class of negative photochromic systems. Indeed, this system switches by irradiation with visible light between a highly conjugated closed-ring isomer (DHP) and a less aromatic

open-ring counterpart (CPD). The reverse process from the CPD to the DHP isomer is achieved by UV-irradiation or heating (Scheme 1).^[3]

The chemical modification of the DHP photochromic core by one or two electron-withdrawing pyridinium substituents ($2c^+$ and $3c^{2+}$; Scheme 1) has recently been studied and the photo-isomerization quantum yield has been improved by at least one order of magnitude relative to the unsubstituted reference compound $1c$.^[4] Indeed, the ring-opening quantum yield increases from 0.006^[5] in $1c$ to 0.040 and 0.042^[4] in $2c^+$ and $3c^{2+}$, respectively. This improved photoswitching behavior was rationalized by theoretical calculations: in the reference compound, the photo-isomerization can be triggered by blue light irradiation ($\lambda > 480$ nm),^[5] populating a S_2 zwitterionic excited state dominated by a HOMO→LUMO transition.^[6] This state is photo-active (*i.e.*, state that leads to photo-isomerization) but is largely depopulated by internal conversion because of the presence of a lower S_1 covalent excited state, which is unreactive (*i.e.*, does not lead to isomerization).^[6] In the mono and bispyridinium substituted DHPs, the S_1 and S_2 excited states gain charge transfer characters and the S_1 excited state becomes the photo-active state upon red light irradiation ($\lambda > 630$ nm) with an improved efficiency.^[4]

The tetrapyrindinium DHP derivative ($4c^{4+}$; Scheme 1) was later synthesized and showed very different electrochemical and photochemical behaviors compared to its bispyridinium counterpart.^[7] On the one hand, $4c^{4+}$ undergoes an efficient ring-opening isomerization leading to $4o$ upon electrochemical reduction, whereas $3c^{2+}$ does not produce its corresponding opening isomer $3o$ when reduced. Thus, the use of pyridinium redox-active units can provide a strategy for the electrochemically triggered switching processes of DHPs. On the other hand, the photo-isomerization quantum yield of $4c^{4+}$ was measured with a 300-fold decrease compared to $3c^{2+}$ (and is also lower than the parent $1c$), which is a surprising and yet unexplained observation.



Scheme 1. DHP (closed form) – CPD (open form) isomerization and various pyridinium-appended systems studied.

In this study, in order to provide a rationale for the unexpectedly poor efficiency of the tetrapyrindinium DHP derivative $4_{c^{4+}}$, we investigate its photoswitching behavior using theoretical calculations based on density functional theory (DFT) and time-dependent DFT (TD-DFT) and we compare our results with those of the mono 2_{c^+} and bispyridinium $3_{c^{2+}}$ derivatives. First, the computed ground-state structure of $4_{c^{4+}}$ is compared with the experimental X-ray data and the thermal isomerization pathway is investigated through the characterization of the biradicaloid transition state between the closed- and open-ring isomers and its associated energy barrier. Second, we investigate the potential energy landscape of the first two excited states S_1 and S_2 along the initial photo-isomerization pathway and we show how it differs from that of 2_{c^+} and $3_{c^{2+}}$. Finally, we propose a possible explanation for the change of behavior observed in $4_{c^{4+}}$ based on steric and electronic factors.

Results and Discussion

Ground-state properties. Single crystals of $4_{c^{4+}}$ suitable for X-ray analysis have been obtained by slow diffusion of diethyl ether into an acetonitrile solution of the compound. $4_{c^{4+}}$ crystallizes in the space group $C2/c$ (monoclinic system). The X-ray structure (Figure 1 and Table S1 in Supporting Information) reveals that the length of the central C–C bond is 1.522 Å, which is similar to values observed for other DHP derivatives. The distance is also in good agreement with the computed value of 1.545 Å for the isolated system using DFT (Table 1). In this experimental structure, the four pyridinium units adopt a twisted arrangement, with dihedral angles with respect to the main DHP plane of 64° for two pyridinium and 76° for the other two. The DFT structure of the isolated model system (see Computational Section and Table S2 in Supporting Information) belongs to the C_{2h} point group and the torsion angle with respect to the DHP core is 64° for the four pyridinium groups. This torsion is more pronounced than that of the corresponding monopyridinium- and bispyridinium-appended DHP for which torsion angles of 34° and 42° were computed, respectively (Table 1).^[4] Note also that the introduction of

pyridinium groups slightly breaks the planarity of the DHP core by a few degrees compared to the reference compound 1_c (Table 1).

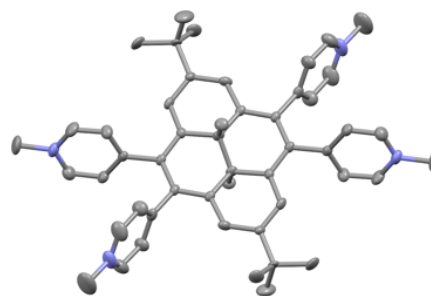


Figure 1. Molecular structure of $4_{c^{4+}}$. Ellipsoids are set at 50% probability; hydrogen atoms and anions have been omitted for clarity.

Table 1. Geometrical parameters^[a] characterizing the ground-state structures of the studied DHP derivatives.

DHPs	q(Å)	ϕ (°)	τ (°)
1_c	1.543	2.4	–
2_{c^+}	1.547	6.2	34
$3_{c^{2+}}$	1.549	7.9	42
$4_{c^{4+}}$	1.545	7.7	64

[a] q is the transannular bond distance, ϕ is an average dihedral angle measuring the DHP planarity and τ is the torsion angle of the pyridinium groups with respect to the DHP core. Taken from Ref. [12] for 1_c and from Ref. [4] for $3_{c^{2+}}$.

DFT calculations confirm that the DHP isomers are more stable than their respective CPD forms by about 20 kcal·mol⁻¹ (Table 2). Interestingly, all CPDs exhibit similar thermal isomerization barriers of about 18 to 20 kcal·mol⁻¹, values in agreement with those reported in the literature for other related DHPs.^[8,9] We can thus predict that the thermal stability of all these CPDs is rather similar. Note that the transition states involved have a strong biradical nature, as indicated by their heavily spin-contaminated wave function ($\langle S^2 \rangle \approx 1$). This is consistent with previous studies of other DHP/CPD couples.^[10]

Table 2. Computed relative energies between closed- and open-ring isomers and potential energy barriers for the thermal cyclization process.

Systems	$\Delta E_{\text{DHP/CPD}}$ (kcal/mol)	$\Delta E_{\text{CPD} \rightarrow \text{DHP}}^\ddagger$ (kcal/mol)
1	22.1	20.0
2^+	20.6	18.9
3^{2+}	17.1	18.9
4^+	20.5	17.9

Excited-state properties. The absorption spectrum of $4_{c^{4+}}$ recorded in acetonitrile at 3·10⁻⁵ M (Figure 2) shows four broad absorption bands at 672, 496, 398 and 334 nm. Based on TD-DFT calculations (Table 3 and Figure S1 in Supporting Information), these bands can be attributed to π - π^* transitions

localized mainly within the DHP core but with some charge transfer character towards the electron-withdrawing pyridinium groups, as shown by the natural transition orbitals^[11] involved in some of these electronic transitions.

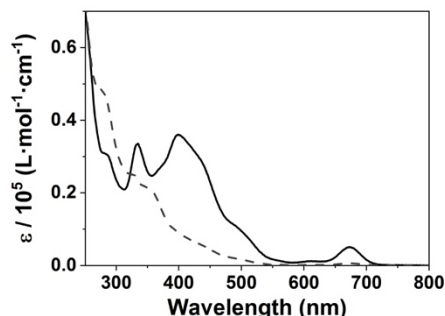
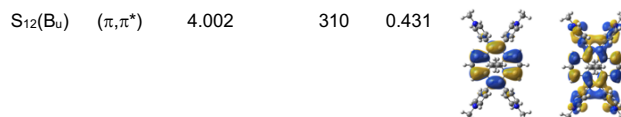


Figure 2. UV-Vis absorption spectra of $4c^{4+}$ (full line) and $4o^{4+}$ (dashed line) recorded in acetonitrile at $3 \cdot 10^{-5}$ M.

Table 3. UV/Vis absorption spectral data and TD-DFT results of $4c^{4+}$ in acetonitrile.

State	Nature	Excitation Energy (eV)	λ (nm)	$f^{[a]}$	NTOs ^[b]
$S_1(A_u)$	(π, π^*)	2.052	604	0.057	
$S_2(B_u)$	(π, π^*)	2.403	516	0.004	
$S_3(A_u)$	(π, π^*)	2.986	415	1.063	
$S_4(B_u)$	(π, π^*)	3.234	383	0.502	
$S_7(A_u)$	(π, π^*)	3.512	353	0.026	
$S_9(B_u)$	(π, π^*)	3.573	347	0.087	



[a] Oscillator strength. [b] Main pair of natural transition orbitals.

Upon visible light irradiation ($\lambda > 630$ nm or $\lambda > 480$ nm), $4c^{4+}$ can be converted into its open-ring isomer $4o^{4+}$ with a low quantum yield. The C-C central bond opening and the concomitant loss of aromaticity are revealed by a decrease of the absorption bands in the visible region and the emergence of a new set of bands in the UV-region (Figure 2). As published previously,^[4] the introduction of one or two pyridinium units (in *trans* position of the DHP core; Scheme 1) enhances the photo-induced ring-opening efficiency compared to the reference $1c$ system by one order of magnitude. In addition, the photo-isomerization in $2c^+$ and $3c^{2+}$ can be triggered by exciting their S_1 state with red light,^[4] whereas $1c$ needs to be excited in its S_2 state with blue light (no isomerization from S_1).^[5] Quantum chemical calculations brought a simple rationalization for this behavior: in the reference compound $1c$, the DHP \rightarrow CPD photo-isomerization is largely inefficient because of the depopulation of the photo-active $S_2(\pi, \pi^*)$ zwitterionic (L_a) state by the lower $S_1(\pi, \pi^*)$ covalent (L_b) state. While vibrational relaxation in the $S_2(\pi, \pi^*)$ state leads to a CPD precursor (denoted CPD* herein) characterized by an elongated transannular C-C bond and a loss of planarity of the DHP core relative to the ground-state structure, the $S_1(\pi, \pi^*)$ relaxes to an excited-state minimum (denoted DHP*) with a similar structure as ground-state DHP (see Table 4). From DHP*, the system deactivates back to the initial ground-state DHP either radiatively or non-radiatively. Thus, the $S_1(\pi, \pi^*)$ L_b state does not lead to CPD isomerization.^[6,12,13] In the mono $2c^+$ and bispyridinium $3c^{2+}$ DHP compounds, the S_1 and S_2 electronic excited states gain a partial charge transfer character from the DHP core to the electron-withdrawing pyridinium units. Furthermore, the S_1 state becomes the photo-active state, as this state leads directly to CPD* by vibrational relaxation on S_1 . No population quenching by a lower excited state is present and irradiation at low energy (red light) can be used to trigger the DHP \rightarrow CPD photo-isomerization.^[4,8,9] Note that CPD* is only an excited-state intermediate towards isomerization and that further vibronic couplings and non-adiabatic processes are involved in the nonradiative decay to S_0 on the way to the CPD isomer.^[6]

Table 4. Main geometrical parameters^[a] characterizing the minima on the S_0 , S_1 and S_2 potential energy surface of the studied DHPs.

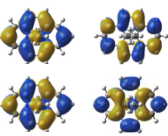
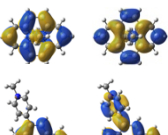
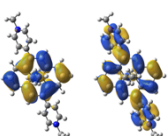
	$1c$		$2c^+$		$3c^{2+}$		$4c^{4+}$	
	q(Å)	$\phi(^{\circ})$	q(Å)	$\phi(^{\circ})$	q(Å)	$\phi(^{\circ})$	q(Å)	$\phi(^{\circ})$
S_0^{\min}	1.543	2.4	1.547	6.2	1.549	7.9	1.545	7.7
S_1^{\min}	1.552	4.1	1.555	9.9	1.568	14.4	1.554	11.3
S_2^{\min}	1.608	12.3	1.552	7.6	1.558	10.8	1.613 ^[b]	21.1 ^[b]

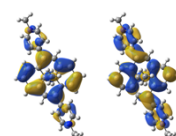
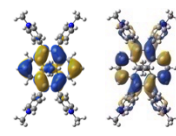
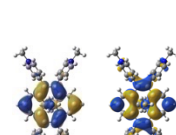
[a] q is the transannular bond distance and ϕ is an average dihedral angle measuring the DHP planarity. Taken from Ref. [12] for $1c$ and from Ref. [4] for $3c^{2+}$. [b] Values corresponding to a S_1 saddle point structure (no S_2 minimum for $4c^{4+}$; see Figure 3).

Remarkably, the quantum yield for the photo-induced ring-opening isomerization of $4c^{4+}$ is $\Phi_{c/o}=1.4 \cdot 10^{-4}$,^[7] determined by actinometry ($\lambda=480$ nm) using 2,7-di-tert-butyl-trans-10b,10c-dimethyl-10b,10c-dihydropyrene as standard, showing a significant 300-fold decrease compared to the mono and bispyridinium DHPs. Such a drastic change of photoswitching efficiency necessarily coincides with a substantial change of the photochemical pathway. In other words, important modifications in the potential energy landscape must occur such that the formation of CPD* is impeded in $4c^{4+}$. In order to rationalize this unexpected behavior, TD-DFT calculations were performed to investigate the initial photochemical pathway.

The first remarkable feature concerns the change of nature of the two lowest excited states S_1 and S_2 between the bispyridinium and tetrapyrindinium DHPs (Table 5). In fact, the S_1 and S_2 excited states of $4c^{4+}$ correlate directly with the S_1 L_b and S_2 L_a states of the reference $1c$, as shown by the direct comparison of the natural transition orbitals (NTOs) involved in the S_1 and S_2 states of these two DHPs. A consequence is that the associated potential energy surfaces present a similar topology between $1c$ and $4c^{4+}$, but very different from that of $3c^{2+}$. In the bispyridinium DHP $3c^{2+}$, the S_1 potential energy surface presents a minimum corresponding to CPD* and the relaxation on this surface leads the system directly to this minimum without any surface crossing involved.^[4] In the tetrapyrindinium DHP $4c^{4+}$, the S_1 minimum does not correspond to CPD* but to DHP* (*i.e.*, the L_b state minimum of similar structure as ground-state DHP), as illustrated by the excited-state geometrical distortions (Table 4). The CPD* structure is associated with a noticeable elongation of the transannular C–C bond by 0.07 Å and an increase of the non-planarity of the DHP core of 13.4° with respect to the ground-state structure. CPD* is thus diabatically linked to the S_2 L_a state and is reached from S_2 through a peaked S_2/S_1 conical intersection, as shown in Figure 3. This is similar to what was found in the reference DHP $1c$.^[6,12] The main difference is that CPD* is a first-order saddle point on the S_1 potential energy surface in $4c^{4+}$, while it is a minimum in $1c$. This is reminiscent of a second-order Jahn-Teller effect, where the derivative coupling between S_1 and S_2 is responsible for the negative curvature at CPD*. Thus, CPD* is an unstable structure in the tetrapyrindinium DHP $4c^{4+}$, and relaxation from CPD* leads preferentially to DHP* (Figure 3). This peculiar potential energy landscape accounts for the remarkably low DHP→CPD photoswitching efficiency in this compound (even lower than that of the reference DHP $1c$ by a factor 40).

Table 5. Comparison of the first two electronic transitions in various DHPs in acetonitrile.

System	State	Excitation Energy (eV)	λ (nm)	$f^{[a]}$	NTOs ^[b]
$1c$	$S_1(A_u)$	2.189	567	0.003	
	$S_2(B_u)$	2.567	483	0.030	
$3c^{2+}$	$S_1(A_u)$	1.968	630	0.080	

$S_2(A_u)$	2.284	543	0.238		
$4c^{4+}$	$S_1(A_u)$	2.052	604	0.057	
$S_2(B_u)$	2.403	516	0.004		

[a] Oscillator strength. [b] Main pair of natural transition orbitals.

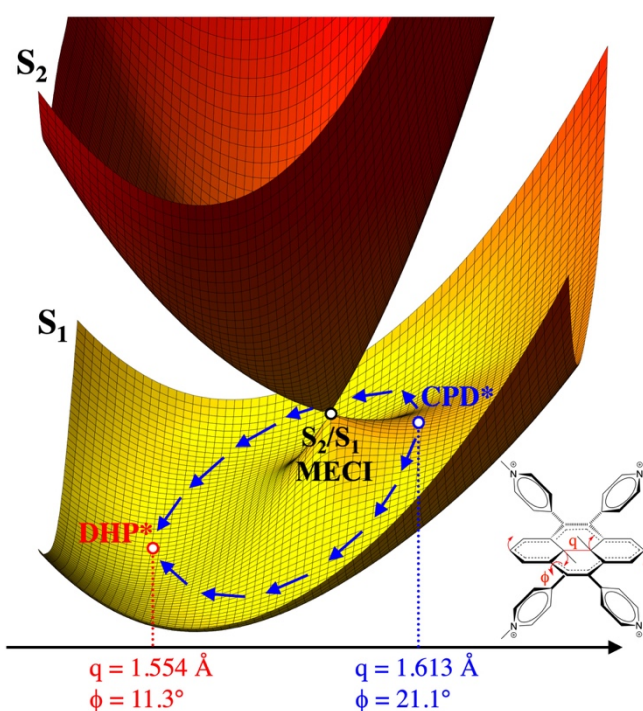


Figure 3. Illustration of the S_1 and S_2 potential energy landscape in $4c^{4+}$. DHP* is a minimum on the S_1 potential energy surface, while CPD* is a first-order saddle point diabatically connected to S_2 via the S_2/S_1 minimum energy conical intersection (MECI). Blue arrows indicate the preferential relaxation from CPD* to DHP*. The transannular central C–C bond length (q) and the average twisting angle between the pyridinium and the DHP core (ϕ) are indicated at the DHP* and CPD* structures.

It is noteworthy that the effect of the four pyridinium units cancel out, while the dissymmetry brought by one or two pyridinium units alters the excited states in favor of the DHP→CPD photoisomerization. In the following, we try to provide some possible explanation for this surprising effect. The first factor that comes to mind is a steric one. As mentioned before, the ground-state structure of $4c^{4+}$ mainly differs from that of $2c^+$ and $3c^{2+}$ by the torsion angle of the pyridinium groups. Because of the proximity and bulkiness of these groups, the decrease of the steric repulsion requires that the pyridinium rotates further relative to the

quasi-planar DHP core in $4c^{4+}$. As a result, the torsion angle ϕ increases by 30° and 22° , with respect to that in $2c^+$ and $3c^{2+}$. Thus, the pyridinium are less conjugated to the DHP moiety and the charge transfer character in the excited states is expected to decrease. If this was the case, one would expect that the excited states of $4c^{4+}$ would resemble those of unsubstituted $1c$, with weak electronic contribution from the pyridinium groups. Comparing the NTOs associated with the S_1 and S_2 excited states of $1c$, and $4c^{4+}$ (Table 5), this scenario seems plausible as almost no charge transfer to the pyridinium is observed in $4c^{4+}$, whereas this charge transfer is pronounced in $3c^{2+}$. Figure 4 represents for $4c^{4+}$ the S_1 and S_2 potential energy profiles along the pyridinium torsion angle ϕ following a relaxed scan on S_0 . The representation of the NTOs for these two excited states at the ground-state geometry (*i.e.*, at $\phi = 64^\circ$) and at a ϕ angle of 40° (*i.e.*, similar to the value of ϕ for the S_0 structure of $3c^{2+}$) confirms that the charge transfer character increases as the torsion angle decreases. However, the order of the electronic states is unchanged along the torsion coordinate. Similarly, increasing the pyridinium torsion in $3c^{2+}$ leads to a decrease of the charge transfer character as expected, but no electronic state crossing is observed (Figure S2 in Supporting Information). In other words, even when inducing the charge transfer between the pyridinium groups and the DHP core upon reducing the torsion angle in $4c^{4+}$, the 1B_u state remains S_2 . Thus, although the torsion controls the degree of charge transfer, it does not account for the order of the 1A_u and 1B_u electronic states.

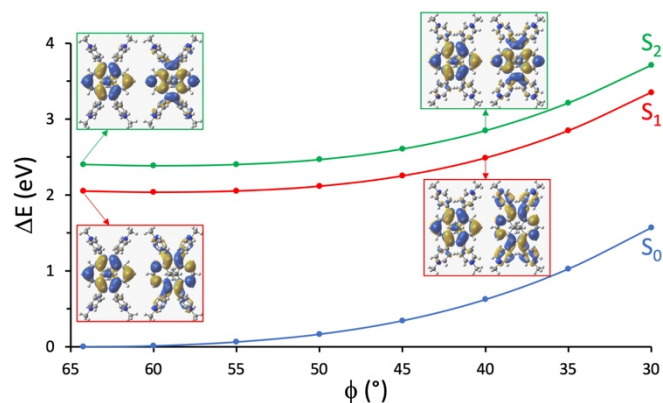


Figure 4. Potential energy profiles of the S_0 , S_1 and S_2 electronic states of $4c^{4+}$ along the pyridinium torsion angle ϕ . Illustration of the NTOs associated with the electronic transition to S_1 and S_2 are shown at two different torsion angles.

In order to explain why the 1B_u photo-isomerizing state is S_2 in $1c$ and $4c^{4+}$, whereas the correlated state becomes S_1 in $2c^+$ and $3c^{2+}$, we need to consider electronic factors and symmetry consideration. The orbitals involved in the single-excitation transition of the S_2 state in $1c$ and $4c^{4+}$ belong respectively to the a_g and b_u irreducible representation of the C_{2h} point group. These orbitals match with the pair of NTOs associated with the $S_2(B_u)$ electronic state shown in Table 5 for these two DHPs. The a_g orbital (HOMO) has a bonding character along the C–C transannular bond, whereas the b_u orbital (LUMO in $1c$ and LUMO+1 in $4c^{4+}$) has an anti-bonding character for this bond (Figure S3 in Supporting Information). Upon symmetry reduction to the C_i point group in $3c^{2+}$, these orbitals belong respectively to the a_g and a_u irreducible representation, meaning that this state now becomes an A_u electronic state. Of course, the $S_1(A_u)$ state of $1c$ and $4c^{4+}$, which is dominated by an $a_g \rightarrow a_u$ transition^[13]

(where a_u has a non-bonding character along the C–C transannular bond; Figure S3 in Supporting Information), remains an A_u state upon symmetry reduction to the C_i point group in $3c^{2+}$. Thus, the two lowest excited states S_1 and S_2 both belong to the same irreducible representation in $3c^{2+}$ (a similar statement is valid for $2c^+$ for which no symmetry exists). The $S_1(A_u)$ and $S_2(A_u)$ can therefore mix in C_i symmetry. This mixing is clearly happening through orbital mixing. Indeed, the second NTO of the $S_1(A_u)$ state in $3c^{2+}$ appears to be a mixture of the second NTOs of the $S_1(A_u)$ and $S_2(B_u)$ states in the more symmetric $1c$ and $4c^{4+}$ molecules. In other words, this S_1 state in $3c^{2+}$ correlates with a mixture of the $S_1(A_u)$ and $S_2(B_u)$ states of $1c$ or $4c^{4+}$. A structural manifestation of this fact is that the excited-state deformations in the S_1 state of $2c^+$ and of $3c^{2+}$ are lying in between those of the S_1 and S_2 states of $1c$ and $4c^{4+}$: for example, the increase of the ϕ angle with respect to the ground state is about 7° in the relaxed S_1 state of $3c^{2+}$, whereas it is respectively 2° and 10° in the S_1 and S_2 states of $1c$ (Table 4). This is also consistent with the fact that the excited a_u LUMO orbital in $3c^{2+}$ has a less pronounced anti-bonding character along the central C–C bond than the b_u orbital in $4c^{4+}$ (Figure S3 in Supporting Information).

A more delicate question to address concerns the reason for the photo-activity of the S_1 state in $4c^{4+}$. As explained above, this state correlates with the S_1 covalent L_b state of $1c$, which does not lead to photo-isomerization in the unsubstituted compound. Several reasons could come to mind. The first one is that the $S_0 \rightarrow S_1$ transition in $4c^{4+}$ is more efficient upon red light irradiation compared to $1c$, with a ratio of ca. 20 in the oscillator strengths in favor of $4c^{4+}$ (Table S3 in Supporting Information). A second argument is that the S_1 potential energy well is deeper in $1c$ compared to the one in $4c^{4+}$. The energy difference between DHP* and CPD* reduces from $4.9 \text{ kcal.mol}^{-1}$ in $1c$ to $3.4 \text{ kcal.mol}^{-1}$ in $4c^{4+}$ giving the latter system more chances to escape from the unreactive DHP* excited-state intermediate to the photoreactive CPD* excited-state species. Another reason could be that this state may undergo a more efficient radiationless deactivation back to the initial ground state in $1c$ than in $4c^{4+}$. In other words, the S_0/S_1 conical intersection could be more difficult to access in $4c^{4+}$ than in $1c$ possibly due to the bulkiness of the pyridinium groups. However, this hypothesis cannot be substantiated, as the photophysical nonradiative decay path from DHP* to DHP is still not well understood.^[6]

Conclusion

We investigated theoretically and experimentally the photo-isomerizing capability of a tetrapyridinium-appended derivative $4c^{4+}$ of the dimethyldihydropyrene photochromic system $1c$. The structures and electronic properties of the ground and lowest electronic excited states of $4c^{4+}$ are described. This derivative displays a surprisingly low quantum yield for the photo-induced ring-opening isomerization from DHP to CPD with a 300-fold decrease relative to that of the mono $2c^+$ and bispyridinium $3c^{2+}$ DHPs.

The comparison of the electronic structures of the first two singlet excited states S_1 and S_2 between these pyridinium-appended derivatives and the reference dimethyldihydropyrene compound shows that $1c$ and $4c^{4+}$ share the same set of excited states, with both S_1 and S_2 being photo-isomerizing states in $4c^{4+}$, albeit with a very low efficiency. The analysis of the topological features of

the S_1 and S_2 potential energy surfaces predict the transient excited-state precursor for the CPD formation (CPD*) to be unstable and short-lived in $4c^{4+}$ because it corresponds to a saddle-point structure on the S_1 potential energy surface in the vicinity of an easily accessible S_1/S_2 conical intersection. Thus, after an efficient $S_2 \rightarrow S_1$ nonradiative decay to CPD*, the system will relax preferentially to the lower S_1 DHP* minimum, an unreactive excited-state species. However, $4c^{4+}$ can still be isomerized into $4o^{4+}$ by direct excitation to S_1 with red light irradiation, something not achievable in $1c$. This could possibly be explained by a more efficient $S_0 \rightarrow S_1$ transition and a shallower DHP* potential energy well, allowing the system to reach the reactive CPD* region.

In the mono $2c^+$ and bispyridinium $3c^{2+}$ DHPs, the situation is different as the reduction of symmetry allows for orbital and state mixing. As a result, the S_1 state in these systems can be viewed as a mixture of the $S_1(A_u)$ and $S_2(B_u)$ states of $1c$ or $4c^{4+}$, producing the CPD* directly on the lowest S_1 excited state. This study thus shows that charge transfer to the pyridinium groups is not the only reason for the increased isomerization efficiency but that symmetry reduction or symmetry breaking also plays a crucial role.

Computational and Experimental Section

Computational Section. DFT has been used to perform calculations of closed-shell and open-shell species for isolated model systems of the reference compound (**1**), of the monopyridinium- (**2**), bispyridinium- (**3**) and tetrapyrindinium-appended (**4**) DHP derivatives in the ground electronic state. All the model systems were obtained by simply replacing the bulky *tert*-butyl groups by hydrogen atoms in order to reduce the computational cost. Geometry optimizations were carried out for the closed-ring and open-ring isomers of all the compounds and for the transition states connecting them. These transition states correspond to open-shell singlet (biradical) species and were computed using broken-symmetry DFT calculations. To account for the spin contamination, spin-projected energies were computed with the approximated spin-correction procedure proposed by Yamaguchi and coworkers.^[14] All these geometry optimizations were performed with the B3LYP functional.^[15] Excited-state geometry optimizations and potential energy profile calculations were performed at the TD-DFT level using the CAM-B3LYP functional.^[16] DHP* and CPD* were obtained by geometry optimization of the $S_1(A_u)$ and $S_2(B_u)$ states, respectively. The schematic representation of the S_1 and S_2 potential energy surfaces shown in Figure 3 results from the energy and structural data along these geometry optimizations and from the numerical frequency calculations performed at these stationary points. The absorption spectrum of $4c^{4+}$ was computed in acetonitrile using the polarizable continuum model (PCM). Natural transition orbitals (NTOs)^[11] for relevant electronic transitions were generated in order to analyze the nature of the excited states involved. The triple- ζ quality basis set including polarization functions 6-311G(d,p) was used throughout.^[17] All calculations were performed with Gaussian 09^[18] using the default convergence criteria for energies and gradients along with an ultrafine integration grid. All the optimized Cartesian coordinates and energies are collected in Table S2 in Supporting Information.

Experimental Section. The synthesis of the tetrapyrindinium-appended dimethyldihydropyrene $4c^{4+}$ is described in ref. [7]. Irradiation experiments in solution were performed under inert atmosphere. The visible irradiations for making the isomerization of the closed-ring (DHP) isomers to their corresponding open-ring (CPD) forms were carried out with a Xe-Hg lamp (500 W), using a $\lambda > 630$ nm cut-off filter and the samples were placed in an ice/water bath in order to limit the reverse thermal reaction. Single crystal X-ray diffraction data have been collected at 200 K on a Bruker AXS Enraf-Nonius Kappa CCD diffractometer using the Mo- K_α radiation from

an Incoatec microsource. Intensity data have been collected for Lorentz-polarization effects and absorption with the Olex² software. Structures solutions and refinements have been performed with Olex² by direct methods. All non-hydrogen atoms have been refined by full matrix least squares with anisotropic thermal parameters. Hydrogen atoms have been introduced at calculated positions as riding atoms. Crystallographic structures have been drawn using the Mercury 3.1 software. A summary of the data collection and structure refinements is provided in Table S1 in Supporting Information. Absorption spectra of $[4c(\text{PF}_6)_4]$ and of its opening isomer were recorded in acetonitrile solution using a Varian Cary 50 Scan UV-visible spectrophotometer equipped with a temperature controller unit. Absorption measurements were carried out in quartz cells with optical path length of 1 cm.

Acknowledgements

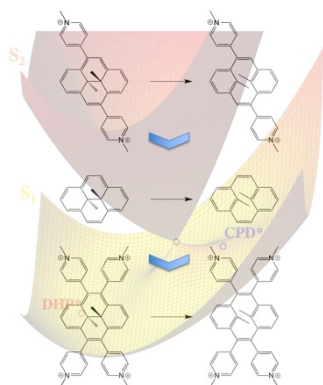
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Keywords: Computational Chemistry • Conical Intersection • Excited States • Photochromism • Photo-isomerization

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Entry for the Table of Contents



While substitution with one or two pyridinium electron-withdrawing groups improves the dimethyldihydropyrene (DHP) photoswitchability, the use of four pyridinium groups has a surprising detrimental effect. We report a theoretical study comparing the excited-state properties of these DHP compounds. A rationale for this unexpected behavior is provided based on topological features of the potential energy surfaces and symmetry arguments.

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