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1 **Generalized Sub-Gaussian processes: theory and application to hydrogeological and**
2 **geochemical data**

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10 **Key Points:**

- 11 • We develop the theoretical formulation of the Generalized Sub-Gaussian (GSG) model
12 for a general distributional form of the subordinator.
- 13 • The GSG formulations tested on laboratory- and field-scale data effectively capture the
14 observed scale dependence of increments' statistics.
- 15 • Our formulations can improve flexibility and accuracy of the GSG model, supporting its
16 applicability to a wide range of data.
- 17

18 Abstract

19 We start from the well-documented scale dependence displayed by the probability distribution and
20 associated statistical moments of a variety of hydrogeological and soil science variables and their
21 spatial or temporal increments. These features can be captured by a Generalized Sub-Gaussian
22 (GSG) model, according to which a given variable, Y , is subordinated to a (typically spatially
23 correlated) Gaussian random field, G , through a subordinator, U . This study extends the theoretical
24 framework originally proposed by Riva et al. (2015a) to include the possibility of selecting a
25 general form of the subordinator distribution, thus enhancing the flexibility of the GSG framework
26 for data interpretation and modeling. Analytical expressions for the GSG process associated with
27 (i) lognormal, (ii) Pareto, and (iii) Gamma subordinator distributions are then derived. We
28 demonstrate the ability of the GSG modeling framework to capture the way key features of the
29 statistics associated with two datasets transition across scales. The latter correspond to variables
30 which are typical of a geochemical and a hydrogeological setting, i.e., (i) data characterizing the
31 micrometer-scale surface roughness of a crystal of calcite, collected within a laboratory-scale
32 setting under given environmental conditions inducing mineral dissolution; and (ii) a vertical
33 distribution of decimeter-scale porosity data, collected along a deep km-scale borehole within a
34 sandstone formation and typically used in hydrogeological and geophysical characterization of
35 aquifer systems. The theoretical developments and the successful applications of the approach we
36 propose provide a unique framework within which one can interpret a broad range of scaling
37 behaviors displayed by a variety of Earth and environmental variables in various scenarios.

38 Plain Language Summary

39 Characterization of hydrogeological and geochemical systems aims at assessing the heterogeneity
40 and scale dependency exhibited by their attributes and the associated key statistics. It has been
41 shown that complex scaling features documented for the statistics of a wide range of Earth,
42 environmental (and several other) variables and their spatial/temporal increments can be captured
43 through a Generalized Sub-Gaussian (GSG) model. The latter relies on the subordination of a
44 Gaussian random field through a subordinator. This study extends the theoretical framework
45 originally proposed for the GSG model to include multiple choices of the subordinator distribution.
46 We provide the theoretical formulation and discuss the main features of the GSG model resulting
47 from (i) a general form of the subordinator and (ii) three selected distributional forms. We show

48 the effectiveness of the GSG modeling framework for the interpretation of real data encompassing
49 a considerably wide range of scales by analyzing (i) a set of surface topography (roughness) data
50 collected on a calcite sample in a laboratory-scale geochemical setting; and (ii) a field-scale
51 distribution of porosity data, collected along a deep borehole within a sandstone formation.

52 **1 Introduction**

53 Geostatistical models adopted for the interpretation of key features of spatial heterogeneity
54 of quantities related to subsurface flow and transport processes consider available observations of
55 a variable of interest as samples from a random field with a given distribution. Analyses of a wide
56 collection of datasets of hydrogeological attributes, including, *e.g.*, (log) hydraulic conductivity
57 and permeability (Liu & Molz, 1997; Meerschaert et al., 2004; Painter, 2001; Painter, 1996, Riva
58 et al., 2013a, 2013b; Siena et al., 2012, 2019), electrical resistivity (Painter, 2001), and neutron
59 porosity (Guadagnini et al., 2015; Riva et al., 2015a) observations, clearly document the
60 occurrence of distinct non-Gaussian features characterizing their distributions. Notably, it has been
61 shown that spatial increments, $\Delta Y(\mathbf{s}) = Y(\mathbf{x} + \mathbf{s}) - Y(\mathbf{x})$, evaluated over separation distance (or lag)
62 \mathbf{s} (\mathbf{x} being a position vector) of a given quantity Y are characterized by distributions displaying
63 peaks that become sharper and tails that tend to become heavier with decreasing lag. A similar
64 behavior, corresponding to distributions transitioning from heavy tailed at small lags to seemingly-
65 Gaussian at increased lags, is documented by analyses of a variety of spatial and/or temporal
66 increments of environmental data, including sediment transport processes (*e.g.*, Ganti et al., 2009)
67 and fully developed turbulence (Boffetta et al., 2008) as well as datasets of Earth, environmental
68 and several other variables (see Neuman et al., 2013 and references therein). Such a scale
69 dependence is directly imprinted to the associated (statistical) moments of increment distributions.

70 All of these evidences suggest that modeling the (spatially correlated) variability of Y
71 through a Gaussian model is not generally warranted. With specific reference to the spatial
72 variability of hydrogeologic quantities, a number of studies evidence that the heterogeneity of
73 natural aquifers is generally more complex than what can be captured through a Gaussian model
74 (*e.g.*, among others, with reference to hydraulic conductivity, Gómez-Hernández & Wen, 1998;
75 Haslauer et al., 2012; Mariethoz et al., 2010; Xu & Gómez-Hernández, 2015 and references
76 therein).

77 In this context, it is also noted that attributes/properties of porous media that at some scale
 78 can be considered as composed by distinct facies/regions, each corresponding to a given material
 79 characterized by an internal degree of heterogeneity, could be represented through multi-modal
 80 distributions (see, *e.g.*, Desbarats, 1990; Lu & Zhang, 2002; Rubin, 1995; Russo, 2002, 2010;
 81 Winter et al., 2003 and references therein). The latter are representative of a conceptual (and
 82 mathematical) model that views the otherwise composite nature of the system as a unique
 83 continuum at the given scale of observation, natural variability within each region being
 84 characterized by a statistical behavior of the kind described above.

85 Riva et al. (2015a, b) show that the above illustrated scale-dependent behavior of the
 86 probability density function (pdf) of ΔY can be captured through a Generalized Sub-Gaussian
 87 (GSG) model. This theoretical framework relies on the idea that the spatial random field
 88 $Y(\mathbf{x}) = \langle Y \rangle + Y'(\mathbf{x})$, $\langle Y \rangle$ and $Y'(\mathbf{x})$ being respectively the ensemble mean and a local zero-mean
 89 fluctuation, can be interpreted through the following model

$$90 \quad Y'(\mathbf{x}) = U(\mathbf{x})G(\mathbf{x}). \quad (1)$$

91 Here, $G(\mathbf{x})$ is a zero-mean, Gaussian random field and $U(\mathbf{x})$ is a so-called *subordinator*,
 92 independent of G , consisting of statistically independent identically distributed (iid) non-negative
 93 random variables. The underlying Gaussian random field generally (but not necessarily) displays
 94 a multi-scale nature which can be captured, for example, through a geostatistical description based
 95 on a Truncated Power Variogram model (*e.g.*, Di Federico & Neuman, 1997; Neuman & Di
 96 Federico, 2003).

97 As opposed to mathematical models based on multifractals (*e.g.*, Boffetta et al., 2008; Frisch,
 98 2016; Lovejoy & Schertzer, 1995; Mandelbrot, 1974; Monin & Yaglom, 1975; Veneziano et al.,
 99 2006) or fractional Laplace approaches (*e.g.*, Kozubowski et al., 2006; Kozubowski et al., 2013;
 100 Meerschaert et al., 2004), which have been employed to mimic the above-mentioned pattern of
 101 increment frequency distributions, the GSG model enables one to represent jointly within a unique
 102 theoretical framework the documented behavior (as described by probability distributions and/or
 103 moments) of a quantity and its incremental values.

104 Riva et al. (2015a) provide the first analytical formulation of the GSG model, illustrating
 105 that the characteristic scaling behavior of the increments results from the decay of the correlation
 106 function of the underlying Gaussian random field with increasing lag. Riva et al. (2015b) illustrate

107 an approach for the generation of unconditional random realizations of statistically isotropic or
108 anisotropic GSG fields in multiple dimensions. Panzeri et al. (2016) develop an algorithm for the
109 generation of GSG fields conditional to a given set of data. Siena et al. (2019) rely on the GSG
110 model for the interpretation of the spatial variability of a set of air permeability data collected
111 along a core of limestone. Guadagnini et al. (2018) present a 9-step procedure for the detection of
112 GSG signatures in a given dataset. Notably, theoretical developments and applications to date rest
113 solely on a lognormal distribution of U characterized by a single parameter, thus limiting the range
114 of possible applications of the GSG model.

115 The present study focuses on a generalization of the GSG framework by extending the
116 formulation of Riva et al. (2015a) to include a generic subordinator. This enables us to enhance
117 the flexibility of the model for data interpretation and modeling by taking into account specific
118 features exhibited by the way statistics of a given dataset transition across scales. We then
119 demonstrate the applicability of the general theoretical framework by considering a (i) two-
120 parameter lognormal; (ii) Pareto; and (iii) Gamma distributional form of U and developing the
121 ensuing analytical expressions for the GSG process. We analyze in this context two datasets
122 associated with differing processes and observation scales. The first application includes direct
123 observations of μm -scale surface topography (or roughness) of mm-scale calcite crystals resulting
124 from induced mineral dissolution. Calcite is a common mineral in the Earth's crust and is
125 characterized by significant dynamics of its surface, depending on environmental conditions (*e.g.*
126 Fischer et al., 2012; Jordan & Rammensee, 1998; Noiriél et al., 2009, 2020). Acquisition of the
127 type of data we consider is subject to increased interest to characterize micro-scale geochemical
128 processes deriving from interactions taking place at fluid-rock interfaces (*e.g.*, Bouissonnié et al.,
129 2018; Pollet-Villard et al., 2016a, b and references therein). While the possibility of acquiring
130 these direct observations is continuously enhanced through the use of modern atomic force
131 microscopy and vertical scanning interferometry, statistical analyses of available datasets are still
132 limited to standard variography (Pollet-Villard et al., 2016a). As an additional test-bed, we analyze
133 a vertical profile of neutron porosity data, collected along a deep borehole in a sandstone formation
134 and encompassing a vertical depth of about 1 km at a 15-cm resolution (Dashtian et al., 2011). As
135 these types of data are routinely available in (hydro)geological and geophysical subsurface
136 exploration, they constitute a remarkable dataset to assess the applicability of our statistical scaling
137 framework at such scales.

138 The work is structured as follows. Section 2 illustrates the key features of the GSG model
 139 and describes a moment-based method of inference of model parameters. The detailed original
 140 analytical formulation of the GSG model associated with a generic subordinator and the ensuing
 141 derivations for the three subordinators here considered are provided in Appendix A and B,
 142 respectively. In Section 3 we compare the performance of these three alternative GSG models for
 143 the interpretation of the two datasets illustrated above. Concluding remarks are provided in Section
 144 4.

145 2 Generalized Sub-Gaussian model

146 2.1 Theoretical framework

147 Zero-mean fluctuations, Y' , at two spatial locations, \mathbf{x}_1 and \mathbf{x}_2 , can be expressed as

$$148 Y'(\mathbf{x}_i) = U(\mathbf{x}_i)G(\mathbf{x}_i) = Y'_i = U_i G_i, \quad \text{with } i = 1, 2. \quad (2)$$

149 The bivariate pdf of Y'_1 and Y'_2 is (Riva et al., 2015a)

$$150 f_{Y'_1, Y'_2}(y'_1, y'_2) = \int_0^\infty \int_0^\infty f_{U_1}(u_1) f_{U_2}(u_2) f_{G_1 G_2} \left(\frac{y'_1}{u_1}, \frac{y'_2}{u_2} \right) \frac{du_2}{u_2} \frac{du_1}{u_1}, \quad (3)$$

151 where $f_{U_i}(u_i)$ is the pdf of U_i and $f_{G_1 G_2}$ is the bivariate pdf of G_1 and G_2 , given by

$$152 f_{G_1 G_2} \left(\frac{y'_1}{u_1}, \frac{y'_2}{u_2} \right) = \frac{e^{-\frac{1}{2\sigma_G^2(1-\rho_G^2)} \left(\frac{y_1'^2}{u_1^2} + \frac{y_2'^2}{u_2^2} - 2\rho_G \frac{y'_1 y'_2}{u_1 u_2} \right)}}{2\pi\sigma_G^2 \sqrt{1-\rho_G^2}}, \quad (4)$$

153 where, σ_G^2 is the variance of G and ρ_G is the correlation coefficient between G_1 and G_2 , which
 154 typically decreases as the separation distance (or lag) $s = |\mathbf{x}_1 - \mathbf{x}_2|$ increases. Starting from Riva et
 155 al. (2015a), who developed the analytical framework for the specific case of a single-parameter
 156 lognormal subordinator, we provide in Appendix A an original theoretical formulation of the GSG
 157 model considering a generic distributional form of U . It is worth noting that, regardless the
 158 distributional form of U , the variogram of Y' is always characterized by a nugget effect (see Eq.
 159 A14), rendered by the product of the variance of G and the variance of U . This result implies that
 160 nugget effects, which are typically considered to appear due to variability of Y' at scales smaller
 161 than the sampling interval and/or to measurement errors, may in fact be (at least in part) considered
 162 as a symptom of non-Gaussianity of the type embedded in the GSG theoretical framework.

163 The general framework introduced in Appendix A encompasses multiple possible
 164 formulations of the GSG model: in this context, we evaluate three possible alternative models for
 165 U , corresponding to a lognormal, Pareto, or Gamma distribution. Each of these models is
 166 characterized by $N_P = 2$ parameters, respectively controlling the shape (*shape* parameter) and the
 167 spreading (*scale* parameter) of the pdfs of the ensuing GSG formulation for Y' . Hereinafter, we
 168 denote the latter as LN-GSG, P-GSG, and Γ -GSG for the lognormal, Pareto, and Gamma
 169 subordinator, respectively. The theoretical formulation of each of these GSG models is provided
 170 in Appendix B.

171 Equations (A7) indicate that the pdfs $f_{\Delta Y}$ of incremental values (ΔY) corresponding to
 172 differing lags depend on (i) σ_G^2 and the N_P parameters of U ; and (ii) ρ_G . While the former
 173 parameters are constant for all lags, the correlation function of G is lag-dependent, thus imprinting
 174 a scaling behavior, i.e., an intrinsic variability with lag, to the shape of the pdf of incremental
 175 values of Y' , independent of the GSG model considered. This feature is clearly illustrated in
 176 Figures 1a-c, where we depict $f_{\Delta Y}$ for selected values of the three GSG model parameters
 177 (analytical expressions being collected in Eq. (B8)) and three values of ρ_G corresponding to short,
 178 intermediate and large lags. The pattern associated with the behavior of peaks and tails of the pdfs
 179 of Y' and ΔY can be described quantitatively by analyzing their standardized kurtosis, $\kappa_{Y'}$ (see
 180 Eq. (A6)) and $\kappa_{\Delta Y}$ (see Eq. (A11)), respectively, deviations from Gaussianity being clearly
 181 revealed by the excess kurtosis, $\kappa_{Y'} - 3$ and $\kappa_{\Delta Y} - 3$. As these quantities increase, the peak of the
 182 pdf of Y' or ΔY grows sharper and the associated tails become heavier. Figures 1d-f depict the
 183 excess kurtosis of Y' , as well as of ΔY , as a function of ρ_G for selected values of the shape
 184 parameter α (for the LN-GSG model, Fig. 1d), a (for the P-GSG model, Fig. 1e), and k (for the Γ -
 185 GSG model, Fig. 1f). Inspection of these figures, together with Eqs. (B7) and (B11), suggests that
 186 for all GSG models (i) $\kappa_{Y'} - 3$ and $\kappa_{\Delta Y} - 3$ do not depend on the scale parameter of the subordinator
 187 and on the variance of G ; (ii) for a given value of ρ_G , $\kappa_{Y'} - 3$ and $\kappa_{\Delta Y} - 3$ increase as the shape
 188 parameter of U decreases; (iii) for a given value of the shape parameter, $\kappa_{\Delta Y} - 3$ increases as ρ_G
 189 increases (or, equivalently, as lag decreases), i.e., the pdfs of ΔY transition with lag. One can note
 190 that, in all cases, $\kappa_{\Delta Y} - 3$ exceeds zero by a significant margin at small lags (i.e., as $\rho_G \rightarrow 1$), even

191 for the largest values of the shape parameter of U considered. With reference to the LN-GSG
 192 model, Figure 1d and Eqs. (B7) and (B11) highlight that there is a threshold value of the shape
 193 parameter, corresponding to $\alpha_T = 2 - \sqrt{\ln 3} \approx 0.95$, such that (i) for $\alpha < \alpha_T$, the pdfs of ΔY are
 194 characterized by lower peaks and lighter tails than those of Y' at all lags; while (ii) for $\alpha > \alpha_T$,
 195 $\kappa_{\Delta Y} - 3$ is higher/lower than $\kappa_{Y'} - 3$ at small/large lags (see also Riva et al., 2015a). An analogous
 196 behavior is exhibited by the results associated with the Γ -GSG model (Fig. 1f), the threshold value,
 197 k_T , of the shape parameter being equal to 1.0. Otherwise, one can demonstrate analytically (see
 198 also Fig. 1e) that $\kappa_{\Delta Y} - 3$ is always larger than $\kappa_{Y'} - 3$ at small lags for the P-GSG model,
 199 regardless the value of the shape parameter a . Besides, the range of values of ρ_G for which (
 200 $\kappa_{\Delta Y} - 3) > (\kappa_{Y'} - 3)$ (i.e., the range of lags where the pdfs of the increments display sharper peaks
 201 and heavier tails than the pdf of Y') tends to increase as a decreases.

202 2.2 Parameter estimation methods

203 The Method of Moment (MOM) is a straightforward way to infer model parameters from a
 204 dataset. Here, we illustrate two approaches to estimate model parameters through MOM. These
 205 are respectively based on (i) sample statistics of the parent variable (Method A) and (ii) sample
 206 statistics of both the parent variable and the incremental data at multiple lags (Method B). Sections
 207 2.2.1 and 2.2.2 examine merits and drawbacks of these methods.

208 2.2.1 Parameter estimation Method A

209 Method A (henceforth denoted as MOM_A) relies on the marginal frequency distribution
 210 and associated moments of Y' . Estimates of GSG model parameters are obtained by replacing
 211 $\langle Y'^2 \rangle$ and $\langle Y'^4 \rangle$ in Eqs. (A3) and (A6) with their sample counterparts, $M_2^{Y'}$ and $M_4^{Y'}$, inferred
 212 from data. The shape parameter of U for the LN-GSG, P-GSG, and Γ -GSG models can be
 213 estimated by making use of Eq. (B7) as

$$\frac{M_4^{Y'}}{3(M_2^{Y'})^2} = \begin{cases} e^{4(2-\alpha)^2}, & \text{with } \alpha < 2 \quad \text{for LN-GSG} \\ \frac{(a-2)^2}{a(a-4)}, & \text{with } a > 4 \quad \text{for P-GSG} \\ 1 + \frac{4k+6}{k(1+k)}, & \text{with } k > 0 \quad \text{for } \Gamma\text{-GSG} \end{cases} \quad (5)$$

Then, by making use of Eq. (B5), one can estimate the product between σ_G and the scale parameter of U (e.g., e^μ , b , and θ , for LN-GSG, P-GSG and Γ -GSG model, respectively) as

$$\sigma_G^2 e^{2\mu} = \frac{M_2^{Y'}}{e^{2(2-\alpha)^2}} \quad \text{for LN-GSG,} \quad (6)$$

$$\sigma_G^2 b^2 = M_2^{Y'} \frac{a-2}{a} \quad \text{for P-GSG,} \quad (7)$$

$$\sigma_G^2 \theta^2 = \frac{M_2^{Y'}}{k(1+k)} \quad \text{for } \Gamma\text{-GSG.} \quad (8)$$

It is noted that the analytical expressions of the marginal pdf (as well as its statistical moments) of Y' for all GSG models are characterized by the scale parameter of the subordinator being always coupled with the scale parameter, σ_G , of the underlying Gaussian process (see Eqs. (B4) - (B6)). It then follows that the set of Eqs.(5)-(8) fully determines $f_{Y'}(y')$, i.e., it is not necessary to estimate σ_G and the scale parameter of U independently to determine $f_{Y'}(y')$. As an additional remark, it is noted that one cannot estimate ρ_G with the methodology here implemented. As such, its application, while straightforward, does not allow ascertaining the degree of spatial correlation of the random field Y' .

2.2.2 Parameter estimation Method B

Method B (henceforth denoted as MOM_B) yields estimates of GSG model parameters by relying jointly on sample statistics of Y' and ΔY . For any given lag, one replaces $\langle Y'^2 \rangle$, $\langle \Delta Y^2 \rangle$ and $\langle \Delta Y^4 \rangle$ in Eqs. (A3), (A8), and (A9) by their sample counterparts $M_2^{Y'}$, $M_2^{\Delta Y}$, and $M_4^{\Delta Y}$, respectively. Making use of Eqs. (B5), (B9) and (B11), the resulting systems of equations are

$$\begin{cases}
 M_2^{Y'} = \sigma_G^2 e^{2\mu} e^{2(2-\alpha)^2} \\
 \frac{M_2^{\Delta Y}}{2M_2^{Y'}} = 1 - \frac{\rho_G}{e^{(2-\alpha)^2}} \\
 \frac{M_4^{\Delta Y}}{(M_2^{\Delta Y})^2} = 3e^{2(2-\alpha)^2} \left\{ 1 + \frac{1}{2} \left(\frac{e^{2(2-\alpha)^2} - 1}{e^{(2-\alpha)^2} - \rho_G} \right)^2 \right\}
 \end{cases}
 \quad \text{for LN-GSG,} \quad (9)$$

$$\begin{cases}
 M_2^{Y'} = \sigma_G^2 b^2 \frac{a}{a-2} \\
 \frac{M_2^{\Delta Y}}{2M_2^{Y'}} = 1 - \frac{a(a-2)\rho_G}{(a-1)^2} \\
 \frac{M_4^{\Delta Y}}{(M_2^{\Delta Y})^2} = \frac{3}{2} \left[\frac{1}{a(a-4)} - \frac{4\rho_G}{(a-1)(a-3)} + \frac{1+2\rho_G^2}{(a-2)^2} \right] \left[\frac{1}{(a-2)} - \frac{a\rho_G}{(a-1)^2} \right]^{-2}
 \end{cases}
 \quad \text{for P-GSG,} \quad (10)$$

$$\begin{cases}
 M_2^{Y'} = \sigma_G^2 \theta^2 k(1+k) \\
 \frac{M_2^{\Delta Y}}{2M_2^{Y'}} = 1 - \frac{k}{1+k} \rho_G \\
 \frac{M_4^{\Delta Y}}{(M_2^{\Delta Y})^2} = 3 \left[1 + \frac{1}{k} + \frac{(k+1)(2+k+\rho_G^2 k - 2k\rho_G)}{k(k+1-\rho_G k)^2} \right]
 \end{cases}
 \quad \text{for } \Gamma\text{-GSG.} \quad (11)$$

236 Equations (9)-(11) allow estimating all parameters characterizing the joint pdf of ΔY , i.e.,
 237 (i) the product of the scale parameters of G and U , (ii) the shape parameter of U , and (iii) the
 238 correlation coefficient ρ_G , which enables us to diagnose the dependence on lag of increment
 239 statistics. We further note that relying on the joint use of Y' and ΔY data is recommended because
 240 it leads to an (often considerably) augmented set of data upon which sample moments are
 241 evaluated, thus improving the accuracy of the estimates. This approach yields a set of three
 242 parameter estimates for each investigated lag. Riva et al. (2015a) document that MOM_B provides
 243 results of similar quality to those one could obtain upon relying on parameter estimation through
 244 analyzing incremental data at various lags via Maximum Likelihood (ML). This element, together
 245 with the high computational demand associated with ML, leads us to rely on MOM_B for the
 246 purpose of our analyses.

247 According to our theoretical framework (see Section 2.1), we expect that values of the shape
248 parameter and of the product between the scale parameters of U and G remain (approximately)
249 constant with lag. It then follows that an additional benefit of relying on MOM_B, as opposed to
250 MOM_A, is that it enables one to assess the consistency of the parameter estimation results with
251 these theoretical requirements. As already noted for the pdf of Y' , the pdf of ΔY (as well as its
252 statistical moments) for all GSG models is also characterized by the scale parameter of U being
253 always coupled with σ_G (see Eqs. (B8) - (B10)). Therefore, the inability to provide unique
254 estimates of the scale parameters of U and G (while only their product is estimated) does not
255 hamper the use of the results of the analysis for further applications, typically involving
256 generations of a collection of realizations of a given random field to be employed in the context of
257 studies on flow and transport processes in a Monte Carlo framework.

258 **3 Application to laboratory- and field- scale datasets**

259 The three alternative GSG models illustrated in Section 2 and in Appendix B are here
260 considered for the characterization of the spatial variability of two datasets. These are selected to
261 represent two differing observation scales, i.e., a laboratory- and a field- scale setting. Both
262 systems are characterized by the availability of a considerable amount of observations, which is
263 achievable with modern measurement techniques, and are therefore well suited for the analysis.

264 3.1 Micrometer-scale topography of a millimeter-scale calcite sample resulting from 265 mineral dissolution

266 The first dataset we consider (hereinafter denoted as Dataset 1) comprises direct observations
267 of surface topography collected on a (104) calcite cleavage plan. While calcite is the main rock-
268 forming mineral of limestones and has a key role in a variety of geological and biological systems,
269 its surface is characterized by remarkable dynamics when put in contact with aqueous fluid, which
270 are still not completely characterized, the (104) surface plane being very common in natural
271 settings. The sample consisted of a ~5mm-sized single crystal of calcite polished through a multi-
272 step abrasive sequence. The initial arithmetic roughness of the surface was on the order of 50 nm.
273 The sample was introduced in a mixed-flow reactor set-up. The crystal was subject to reaction for
274 8 days at room temperature and at a saturation index with respect to calcite of 0.8, corresponding
275 to conditions where dissolution occurs while the nucleation of etch pits is thermodynamically

276 impossible. Measurements of surface topography, (x, y) being spatial coordinates in the horizontal
 277 plane, are collected by means of a vertical scanning interferometer (Zygo NewView 7300) with a
 278 vertical resolution of 3 nm, on a two-dimensional grid of $N_1 = 250 \times 250 = 62500$ cells, with lateral
 279 resolution $dl = 2.2 \mu\text{m}$. Additional details of the experimental set-up and procedure are offered in
 280 Bouissonnié et al. (2018). The surface is characterized by a slight curvature, resulting from the
 281 preliminary polishing of the sample. Mean-removed topography data, Y' , have been obtained by
 282 subtracting the best-fitting quadratic surface from the measurements. Figure 2a depicts the spatial
 283 distribution of Y' , the sample standard deviation being equal to $\sqrt{M_2^{Y'}} = 0.21 \mu\text{m}$.

284 3.2 Field-scale neutron porosity data

285 Dataset 2 is a collection of neutron porosity data sampled from a (km-scale) deep vertical
 286 borehole in southwestern Iran. The data are part of a wider dataset comprising multiple wells, some
 287 of which have been recently analyzed by Dashtian et al. (2011), Riva et al. (2015a), and
 288 Guadagnini et al. (2015). The borehole considered here is drilled in the Ahwaz field (see Dashtian
 289 et al., 2011), where oil and natural gas are produced from a sandstone formation. A large number
 290 ($N_2 = 6949$) of neutron porosity data collected at a uniform distance of $dz = 15 \text{ cm}$ is available.
 291 The one-dimensional profile of mean-removed porosity data is depicted in Fig. 2b, the associated
 292 sample standard deviation being equal to $\sqrt{M_2^{Y'}} = 8.35\%$.

293 3.3 Results and discussion

294 Figures 3a and 3b depict sample pdfs of Y' for Dataset 1 and 2, respectively. Depictions are
 295 provided in linear and semi-logarithmic scales for ease of analysis. A slight bimodality and
 296 asymmetry are exhibited by the pdf of porosity observations in Dataset 2, the pdf of surface
 297 topography (Dataset 1) being left-skewed. A qualitative comparison (based on visual inspection)
 298 between each of these sample pdfs and a normal distribution with the corresponding variance, $M_2^{Y'}$
 299 , (also included in the figures) suggests deviation from Gaussianity for both variables. This
 300 qualitative result is also confirmed quantitatively by the outcomes of formal (Shapiro-Wilk,
 301 Kolmogorov-Smirnov, and Anderson-Darling) tests performed on randomly-sampled subsets of
 302 data, which reject the Gaussian model at a significance level of 0.05 for both datasets.

303 We compute sample statistics of incremental data, ΔY , evaluated (i) along all directions in
 304 the x - y plane for Dataset 1 and (ii) along the z axis for Dataset 2. The pdfs of ΔY at three diverse
 305 lags ($s = 1, 5$, and $50 dl$ for Dataset 1; and $s = 5, 50$, and $250 dz$ for Dataset 2) are depicted in Figs.
 306 4a and 4b, respectively. As a term of comparison, corresponding normal distributions with the
 307 same variance are juxtaposed to the increment pdfs. These results illustrate that sample pdfs of
 308 increments (i) exhibit the characteristic scale dependence mentioned in Section 1; and (ii)
 309 progressively tend to distributions with lower peaks and lighter tails, resembling the Gaussian
 310 distribution as lag increases, this feature being particularly evident for Dataset 2.

311 Figures 5a and 5b depict the dependence of sample values of $\kappa_{\Delta Y} - 3$ on lag for Dataset 1
 312 and Dataset 2, respectively, dashed horizontal lines denoting values of excess kurtosis of the parent
 313 variable Y' . For both sets, incremental data excess kurtosis is significantly larger than zero at small
 314 lags. Excess kurtosis (EK) of (omnidirectional) incremental data associated with Dataset 1
 315 decreases rapidly as lag increases and tends to attain a quite stable value of ≈ 3.5 at large lags.
 316 Otherwise, values of EK for Dataset 2 tend to consistently decrease across the whole range of lags
 317 considered, attaining values smaller than 1 (i.e., approaching a Gaussian distribution, consistent
 318 with the qualitative result depicted in Fig. 4b) from $s = 400 dz$.

319 To provide an appraisal of the accuracy associated with the sample estimates of EK, we apply
 320 a standard bootstrapping technique (Efron, 1992) to each set of incremental data. This procedure
 321 relies on sampling (with replacement) from a collection of ΔY data related to a given lag a total
 322 of m (here we set $m = 10,000$) sets, each characterized by the same number of elements of the
 323 original collection of ΔY . The same procedure is then repeated for all lags considered. Figures 5a
 324 and 5b depict the 95%-confidence intervals, CI, associated with the estimates of EK at four
 325 representative lags. Uncertainties associated with EK estimates are (in general) negligible.
 326 Therefore, we consider the observed overall decrease of EK with the lag to be significant for both
 327 datasets. We note that $\kappa_{\Delta Y} - 3 > \kappa_{Y'} - 3$ at small lags for Dataset 1 (Fig. 5a), implying that
 328 frequency distributions of ΔY exhibit sharper peaks and heavier tails than does that of Y' , whereas
 329 the opposite behavior is documented at large lags. Otherwise, $\kappa_{\Delta Y} - 3 > \kappa_{Y'} - 3$ over the whole
 330 range of lags considered for Dataset 2 (Fig. 5b). Considering the type of analyses documented in
 331 Figs. 1d-f, the behavior observed for both datasets is consistent with our theoretical models for (i)
 332 $0.95 < a < 2$ in the case of LN-GSG; (ii) $a > 4$ for P-GSG, and (iii) $k > 1$ for Γ -GSG.

333 Estimates of (i) the shape parameter and (ii) the product of the scale parameters of U and G
 334 (henceforth denoted only as *global scale parameter* for conciseness) obtained via MOM_A and
 335 MOM_B for each GSG model formulation are depicted in Fig. 6 (Dataset 1) and Fig. 7 (Dataset
 336 2) as a function of normalized lag. These results are complemented by Table 1 where we list
 337 parameter estimates obtained via MOM_A, together with mean and coefficient of variation (cv)
 338 evaluated over all lags of MOM_B estimates, obtained for all GSG model formulations and both
 339 datasets.

340 Considering Dataset 1, results obtained via MOM_B for LN-GSG (i.e., α in Fig. 6a and
 341 $e^{\theta}\sigma_G$ in Fig. 6d) and P-GSG (i.e., a in Fig. 6b and $b\sigma_G$ in Fig. 6e) do not vary appreciably with
 342 lag (cv \approx 2-3%), consistent with our theoretical framework. Otherwise, MOM_B estimates of k
 343 and $\theta\sigma_G$ (Figs. 6c and 6f, respectively) associated with Γ -GSG are characterized by stronger
 344 oscillations around an average value, as indicated by larger values of the corresponding coefficient
 345 of variation, as compared to the other models. Nevertheless, values of cv range between 18% (for
 346 the shape parameter) and 22% (for the global scale parameter), which (also in view of ubiquitously
 347 present experimental uncertainties) can still be considered as a good approximation of the
 348 constraints associated with theoretical requirements. Figure 6 and Table 1 also document that
 349 MOM_A estimates are consistent with their counterparts obtained via MOM_B for all models.

350 Results for Dataset 2 (Fig. 7) obtained through MOM_B generally reveal more pronounced
 351 oscillations around a constant value and larger values of cv than those observed for Dataset 1, in
 352 particular considering the Γ -GSG model. We remark that the two considered datasets are
 353 associated with differing dimensionalities (Dataset 1 and Dataset 2 being two- and one-
 354 dimensional, respectively) and considering that $N_1 / N_2 \approx 9$, statistics of incremental data for
 355 Dataset 2 are evaluated on a much smaller sample of data as compared to Dataset 1. We regard
 356 this as the main reason related to the (slightly) increased deviations from the expected theoretical
 357 pattern.

358 We rely on the bootstrapping procedure mentioned above to evaluate the uncertainty
 359 associated with the GSG parameter estimates obtained via MOM_B. Figures 6-7 include
 360 depictions of the 95% CIs related to the GSG parameter estimates evaluated at four representative
 361 lags. The width of these intervals is in general very limited. The results obtained via MOM_A (see
 362 Table 1 and dashed lines in Fig. 7) tend to overestimate all parameters, as compared to their

363 MOM_B-based counterparts (except for $\theta\sigma_G$ in Fig. 7f), a notable discrepancy between the two
 364 estimation methods being observed for the shape parameters of P-GSG (Fig. 7b) and Γ -GSG (Fig.
 365 7c).

366 Results collected in Table 1 also evidence that estimates of the shape parameter stemming
 367 from the application of each GSG model to Dataset 1 are smaller than their counterparts related to
 368 Dataset 2. This finding is indicative of a stronger non-Gaussian signature in the former data set, a
 369 behavior that can also be inferred from the increased values of excess kurtosis exhibited by Dataset
 370 1 (see Figs. 5a and 5b).

371 Figures 8a and 8b depict estimates of ρ_G as a function of lag obtained for Dataset 1 and 2,
 372 respectively. These results show that the correlation function of the underlying Gaussian process
 373 is quite insensitive to the choice of subordinator adopted in the GSG model, in particular
 374 considering Dataset 1. Figure 8b suggests that the width of the 95% CIs for Dataset 2 is particularly
 375 wide in the range of lags where the results associated with the three models do not overlap. This
 376 observation suggests that differences observed between ρ_G estimates obtained with the three GSG
 377 models may not be particularly significant in this dataset and can be related to effects of the limited
 378 size of this sample. This result (i) is in agreement with the theoretical framework according to
 379 which the subordinator should be statistically independent of G and (ii) suggests that the
 380 correlation structure provided by the underlying Gaussian process can be considered as a
 381 distinctive signature of the system.

382 Figure 9 depicts sample pdfs of the parent variables (Figs. 9a, 9c) and their increments (Figs.
 383 9b, 9d) corresponding to two separation lags included in Fig. 4 and presented here for the sake of
 384 comparison against theoretical pdfs corresponding to the various GSG models considered. In these
 385 plots, $f_{Y'}$ and $f_{\Delta Y}$ associated with GSG models are evaluated respectively on the basis of (i)
 386 parameters estimated via MOM_A and (ii) the mean values of shape and global scale parameters
 387 obtained via MOM_B, $f_{\Delta Y}$ also including the lag dependent parameter, ρ_G , computed with
 388 MOM_B and depicted in Fig. 8. From a qualitative comparison between Figs. 9a-d and Figs. 3 and
 389 4, it can be appreciated that all GSG models are generally in better agreement with the target
 390 sample pdfs than the Gaussian model. The degree of similarity between sample and analytical pdfs
 391 is quantified through the Kullback-Leibler (KL) divergence (Kullback & Leibler, 1951), D_{KL} . The
 392 latter is a measure of the information lost when a given distribution is used to approximate a target

393 one. As such, smaller values of D_{KL} are associated with reduced loss of information. Considering
394 the pdf of Y' , (i) for Dataset 1 we obtain $D_{KL} = 0.048$ (for LN-GSG), 0.013 (for P-GSG), and
395 0.071 (for Γ -GSG), thus suggesting P-GSG as the best among the models considered; (ii)
396 $D_{KL} \approx 0.068$ for Dataset 2, regardless the subordinator employed. This latter outcome is consistent
397 with Fig. 9c, where all GSG pdf are virtually overlapping. Therefore, when considering Dataset 2
398 the sole analysis of the parent data population does not allow discriminating between alternative
399 GSG models. We finally evaluate D_{KL} between sample and theoretical pdfs of incremental data for
400 diverse lags. Figure 10a depicts D_{KL} versus lag for Dataset 1, Fig. 10b showing a corresponding
401 depiction for Dataset 2. These results highlight that, considering Dataset 1, the P-GSG model
402 provides the highest degree of similarity between sample and theoretical pdfs of increments at
403 almost all lags ($s > 25 dz$), and is consistent with the results obtained for the parent variable as well
404 as with those collected in Table 1 and Fig. 6. Considering Dataset 2, Fig. 10b suggests that the
405 three models provide results of similar quality for lags $s > 200 dz$, a feature that can also be noted
406 from the almost overlapping analytical results depicted in Fig. 9d for $s = 250 dz$. Otherwise, LN-
407 GSG and Γ -GSG outperform P-SGS in the range $0 < s < 100 dz$. This observation, in conjunction
408 with the analysis performed in Fig. 7 and Table 1, leads to favoring LN-GSG for Dataset 2.

409 Overall, our results support the ability of the GSG model to provide a theoretical
410 interpretation of characteristic features associated with the statistics of both investigated datasets.
411 We note that having at our disposal these tools forms the basis to achieve the overarching goal to
412 quantify the way one can transfer the key statistics of a variable (and its increments) across scales,
413 with direct implications on uncertainty quantification. With reference to the spatial distribution of
414 surface roughness, these results constitute an important step to bridge across characterizations of
415 reactive phenomena at microscopic and laboratory scales. In this context, there is documented and
416 growing interest in the application of statistical methods (Fischer et al., 2012; Lüttge et al., 2013;
417 Pollet-Villard et al., 2016; Trindade Pedrosa et al., 2019) to firmly ground the multiscale nature of
418 such processes on rigorous theoretical bases. The quality of our results is encouraging to promote
419 further studies targeting statistically-based descriptions of the temporal evolution of the surface
420 topography of calcite minerals subject to precipitation/dissolution processes acting at diverse
421 scales. We envision addressing this objective in the future by coupling our theoretical approach
422 with direct in situ observations through, e.g., time-lapse nanoscale imaging. In this context,

423 characterizing porosity of natural porous media has the clear potential to link geochemical
424 processes acting at small scales with descriptions of flow and transport at scales compatible with
425 a continuum description of the system. Hydraulic conductivity is intimately related to porosity. As
426 mentioned in the Introduction, statistics of its spatial increments have also been documented to
427 display a behavior consistent with what we have observed here for porosity. These concepts have
428 already been employed in the context of preliminary analytical and numerical studies of flow and
429 transport in porous media associated with such a statistical description by Riva et al. (2017) and
430 Libera et al. (2017).

431 **4 Concluding remarks**

432 We extend the Generalized Sub-Gaussian (GSG) stochastic model proposed by Riva et al.
433 (2015a) by providing theoretical formulations of the GSG for a generic subordinator U . Properties
434 of such an extended and more general model are analyzed and alternative formulations of the GSG
435 model, derived for three selected subordinator forms, are considered to interpret observations
436 associated with two datasets: (i) a set of observations characterizing the surface-roughness
437 resulting from the dissolution of a crystal of calcite, collected in a geochemical laboratory-scale
438 setting under given environmental conditions (Dataset 1); and (ii) a field-scale spatial distribution
439 of porosity data, collected along a deep borehole within a sandstone formation (Dataset 2). Our
440 study leads to the following key conclusions.

- 441 1. For any subordinator type associated with the GSG, the analytical formulation of
442 standardized kurtosis, $\kappa_{Y'}$ and $\kappa_{\Delta Y}$, governing the behavior of peaks and tails of the pdf of
443 Y' and ΔY , respectively, does not depend on scale parameters of U and G . Values of $\kappa_{Y'}$,
444 and $\kappa_{\Delta Y}$ increase as the shape parameter of U decreases, $\kappa_{\Delta Y}$ decreasing as the separation
445 distance (or lag) at which increments are evaluated increases. Thus, GSG models are suitable
446 to capturing the extensively documented peculiar features of Earth and environmental
447 variable whose distributions transition from heavy tailed at small lags to seemingly-Gaussian
448 at increased lags.
- 449 2. The proposed theoretical framework successfully captures the main features of the
450 distributions of the variables analyzed as well as their spatial increments. Results of
451 statistical analyses performed on both datasets are consistent with theoretical expectations:
452 (i) estimates of shape and (global) scale parameters of the GSG models are nearly constant

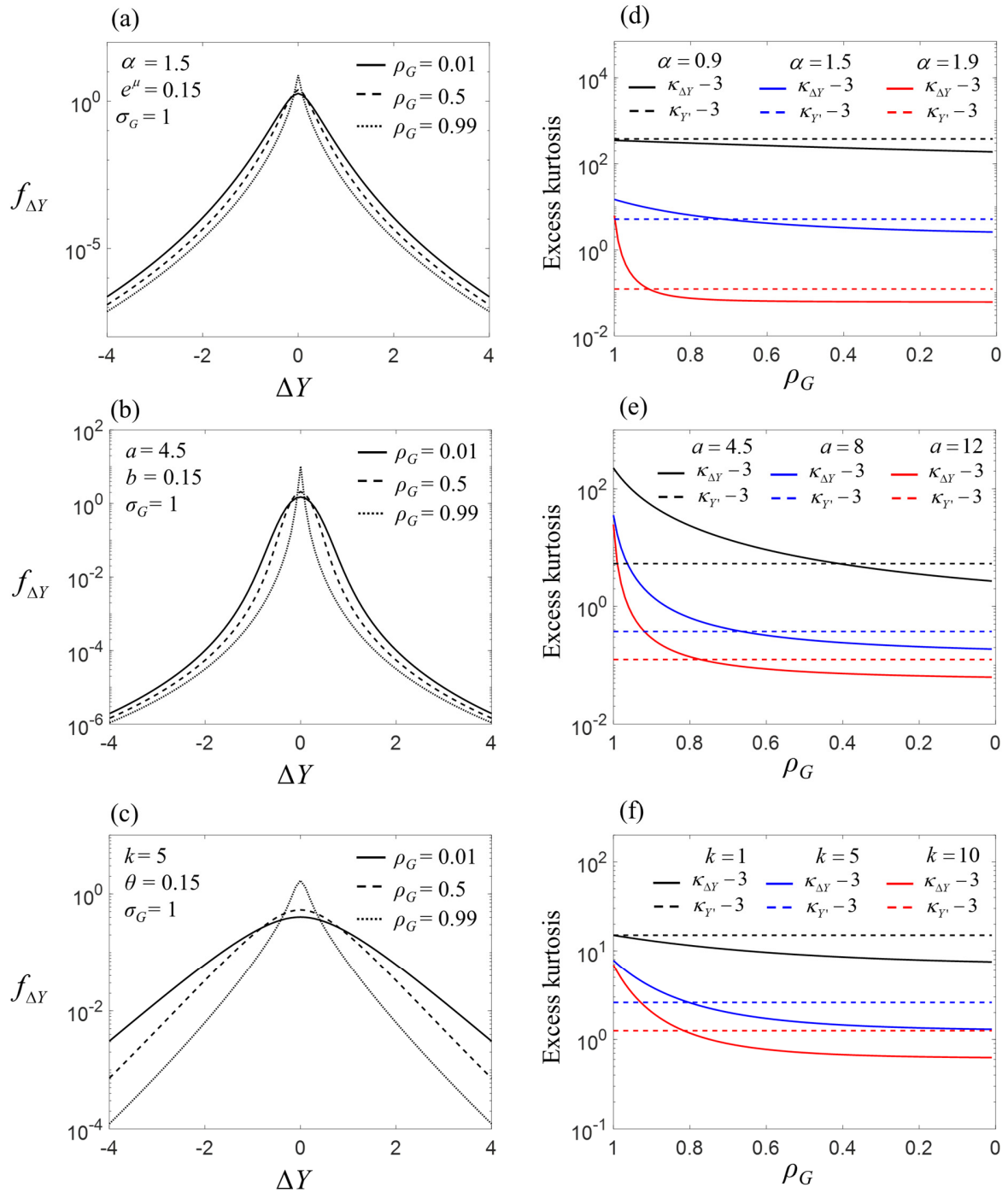
453 with lag; (ii) the correlation coefficient (ρ_G) of the underlying Gaussian process decreases
454 as lag increases, according to a trend that is almost insensitive to the type of subordinator
455 considered. The latter results suggest that the correlation structure provided by the Gaussian
456 process underlying the GSG field can be considered as a distinctive signature of the system
457 behavior.

458 3. The Kullback-Leibler (KL) divergence is adopted to evaluate degree of similarity between
459 theoretical (i.e., based on the various GSG model formulations) and sample Y' and ΔY pdfs
460 in each dataset. Our results indicate that the implementation of multiple subordinators within
461 the GSG framework can enhance the flexibility of the model and improve the accuracy of
462 the interpretation of statistical behavior of a given dataset.

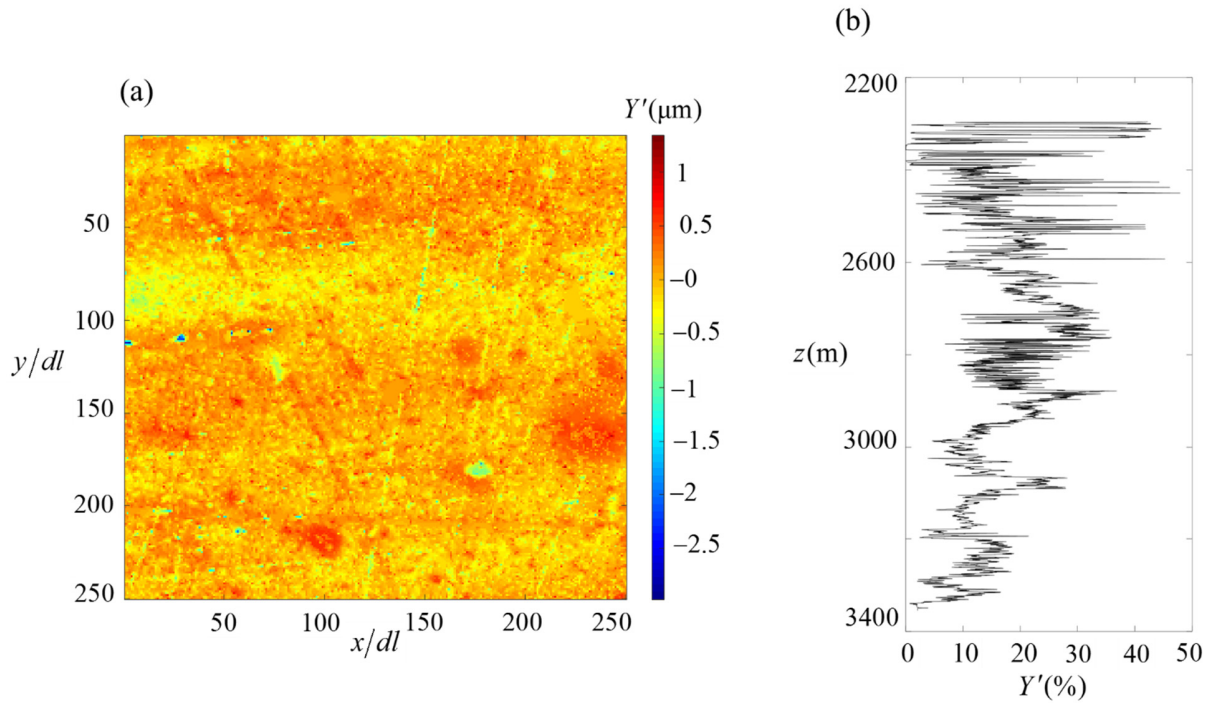
463 The approach and theoretical developments we propose provide a unique framework within
464 which one can interpret a broad range of scaling behaviors displayed by a variety of Earth and
465 environmental variables in various settings. The successful demonstration we present imbues us
466 with confidence about research applications targeting hydrogeological and geochemical scenarios
467 upon leveraging on modern experimental investigation techniques leading to characterize natural
468 systems across a diverse range of scales. These include, for example, further experiments and
469 theoretical analyses devoted to the assessment of micro-scale reaction rates taking place at rock-
470 liquid interfaces.

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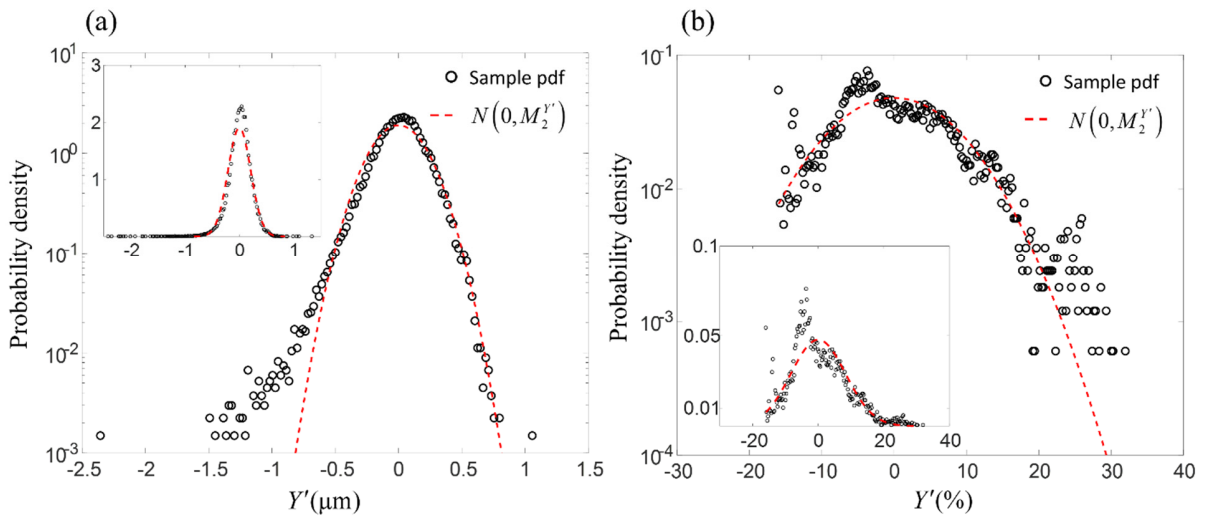


473
 474 **Figure 1.** Probability density function, $f_{\Delta Y}$, computed for three values of ρ_G according to Eq.
 475 (B8) for the (a) LN-GSG; (b) P-GSG, and (c) Γ -GSG model, respectively. Excess kurtosis of (i)
 476 Y' (dashed lines) and (ii) ΔY as a function of ρ_G (solid curves) for three selected values of the
 477 shape parameter (d) α , (e) a , and (f) k for the LN-GSG, P-GSG, and Γ -GSG model, respectively.



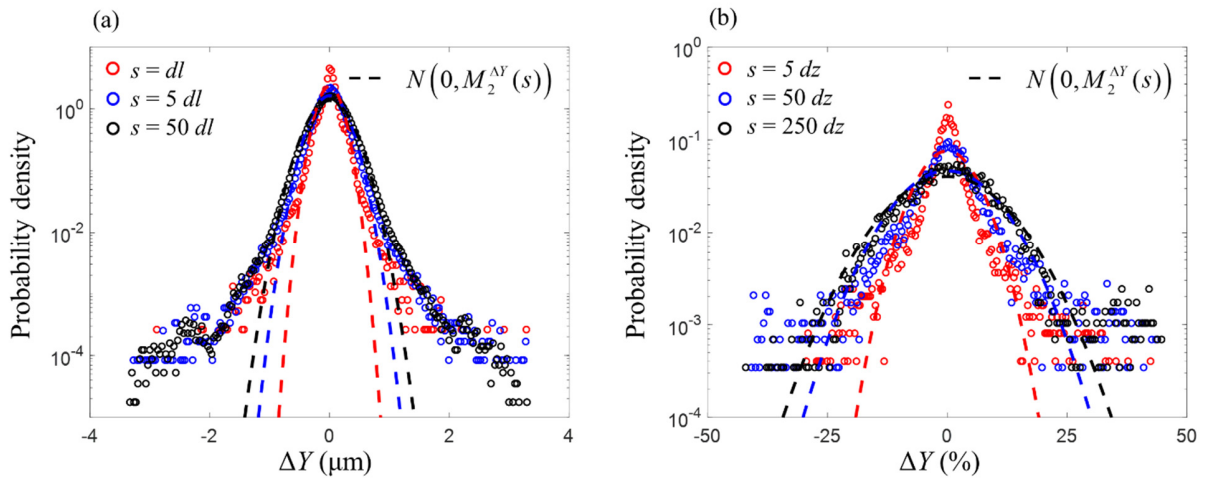
478

479 **Figure 2.** Spatial distributions of zero-mean fluctuations, Y' , of (a) (micrometer-scale) surface
 480 topography data (Dataset 1); and (b) (decimeter-scale) neutron-porosity data (Dataset 2).
 481



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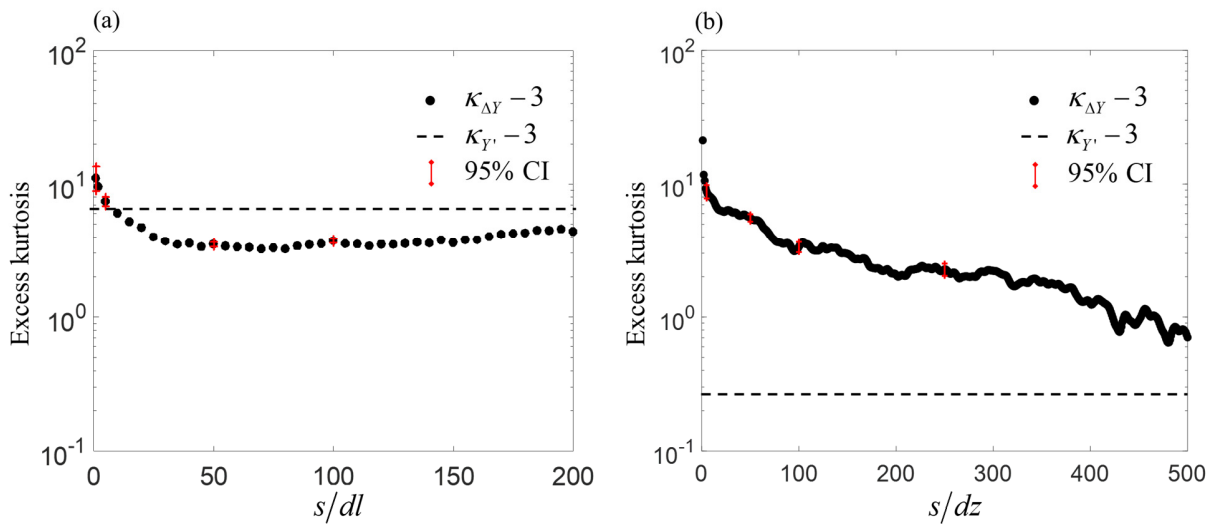
483 **Figure 3.** Sample pdf of Y' data obtained for (a) Dataset 1 and (b) Dataset 2 on semi-logarithmic
 484 and arithmetic (inset) scales. Gaussian pdfs with variance equal $M_2^{Y'}$ are also shown (dashed
 485 curves).



486

487 **Figure 4.** Sample pdfs of incremental data, ΔY , at three lags: (a) $s/dl = 1, 5,$ and 50 for Dataset 1
 488 and (b) $s/dz = 5, 50,$ and 250 for Dataset 2. Gaussian pdfs with variance equal to that of the data
 489 sample are also shown (dashed curves).

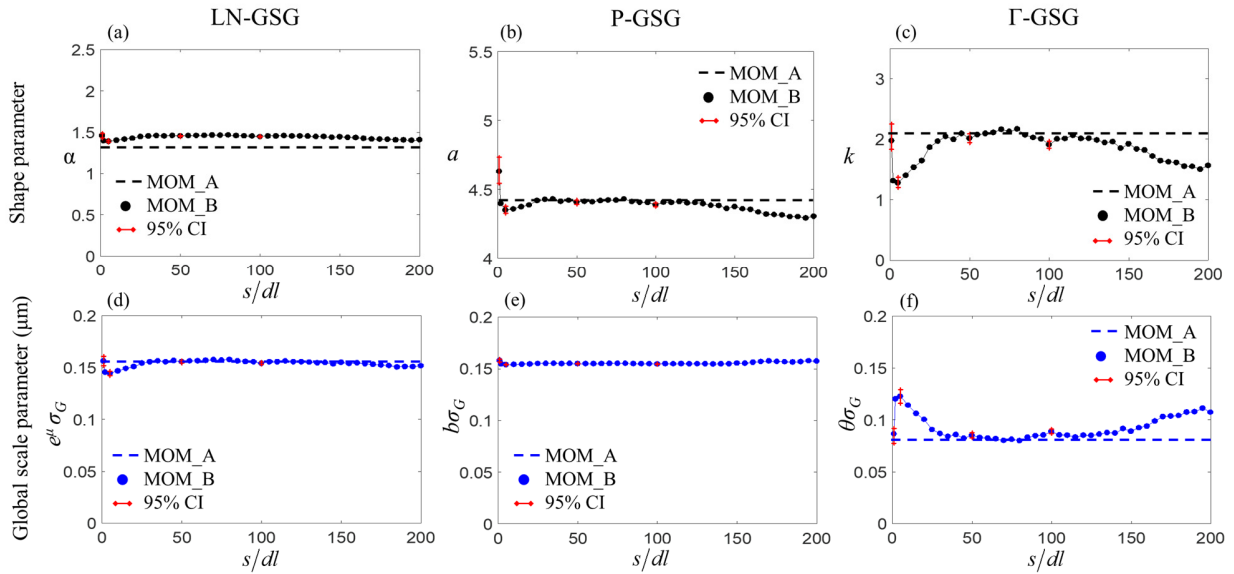
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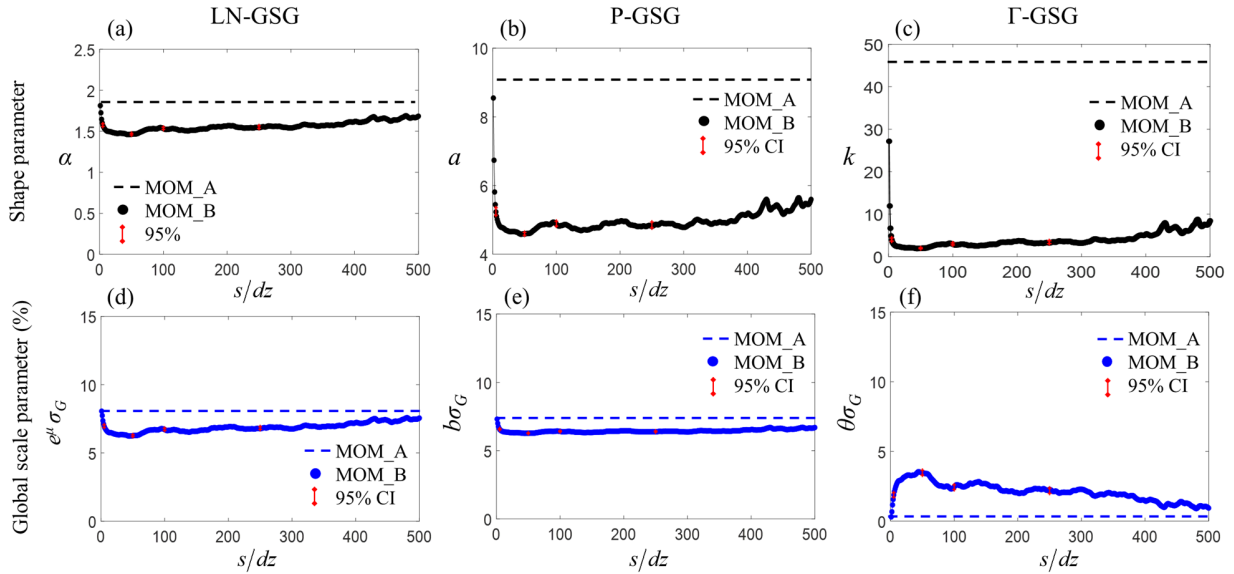
492 **Figure 5.** Excess kurtosis of (i) Y' (dashed horizontal lines); and (ii) ΔY (symbols) versus
 493 (normalized) lag for (a) Dataset 1 and (b) Dataset 2; 95% CIs (evaluated through bootstrapping)
 494 associated with EK estimates at four lags are reported in red.

495



496

497 **Figure 6.** Dataset 1: Estimates of shape (a, b, c) and global scale (d, e, f) parameters of the GSG
 498 model formulations obtained via MOM_A (dashed horizontal lines) and MOM_B (symbols)
 499 versus (normalized) lag; 95% CIs (evaluated through bootstrapping) associated with MOM_B
 500 estimates at four lags are reported in red.



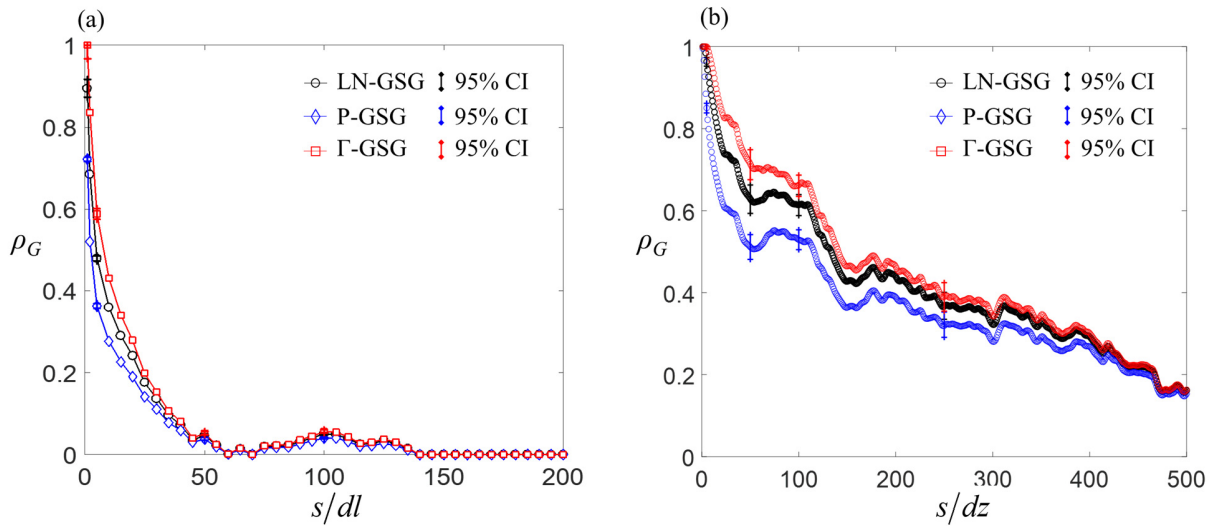
501

502 **Figure 7.** Dataset 2: Estimates of shape (a, b, c) and global scale (d, e, f) parameters of the GSG
 503 model formulations obtained via MOM_A (dashed horizontal lines) and MOM_B (symbols)
 504 versus (normalized) lag; 95% CIs (evaluated through bootstrapping) associated with MOM_B
 505 estimates at four lags are reported in red.

506

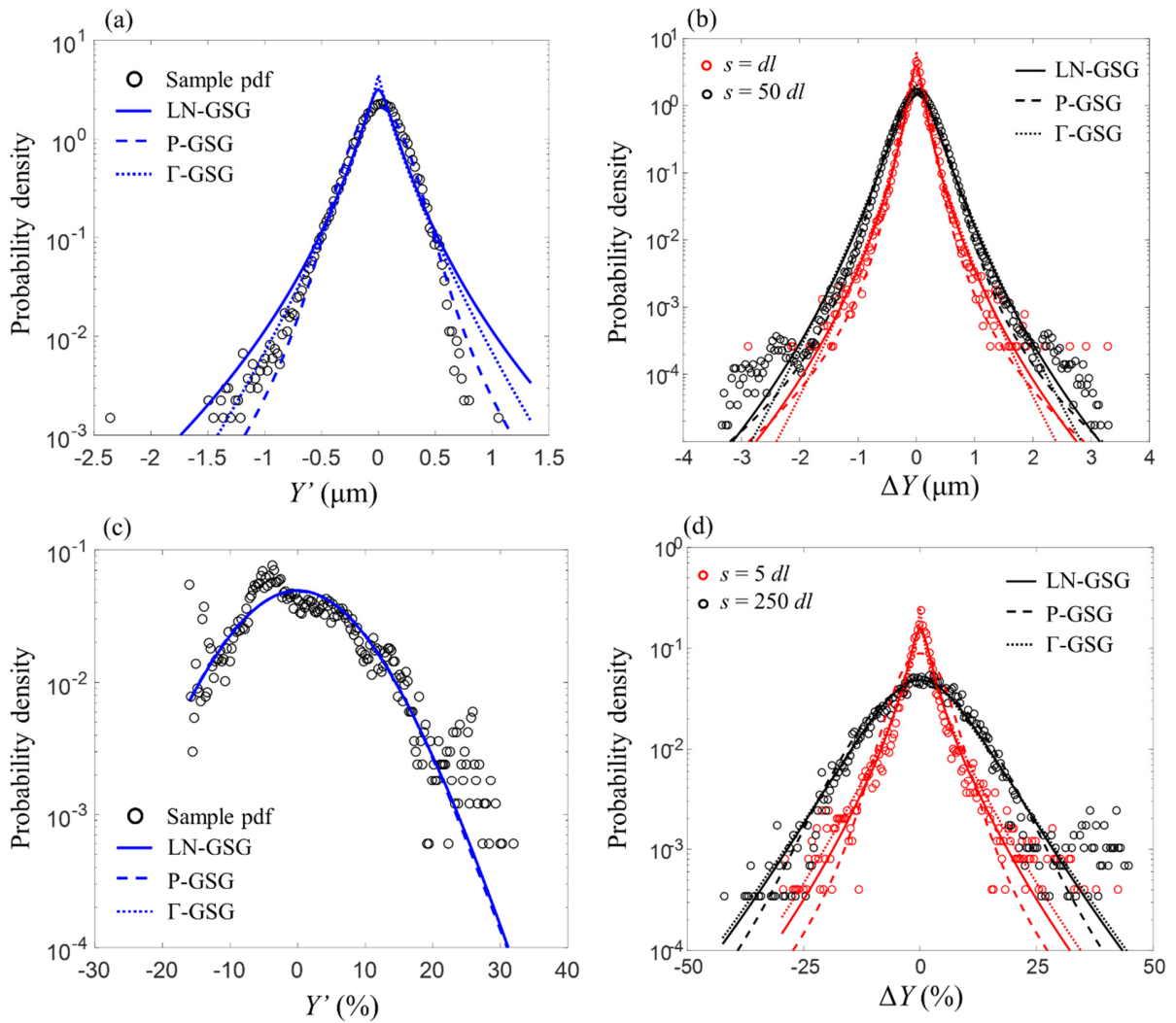
		Dataset 1			Dataset 2		
		LN-GSG	P-GSG	Γ -GSG	LN-GSG	P-GSG	Γ -GSG
Shape parameter	MOM_A	1.34	4.14	2.10	1.86	9.05	46.12
	MOM_B (mean)	1.43	4.36	1.76	1.56	4.96	3.96
	MOM_B (cv)	0.02	0.02	0.18	0.04	0.06	0.48
Global scale parameter	MOM_A	0.16	0.16	0.08	8.18	7.37	0.18
	MOM_B (mean)	0.15	0.16	0.11	6.88	6.44	2.12
	MOM_B (cv)	0.02	0.03	0.22	0.05	0.02	0.30

507 **Table 1.** Parameter estimates obtained via MOM_A; mean and coefficient of variation (cv)
 508 evaluated over all lags of MOM_B estimates obtained for all tested GSG model formulations and
 509 both datasets.
 510



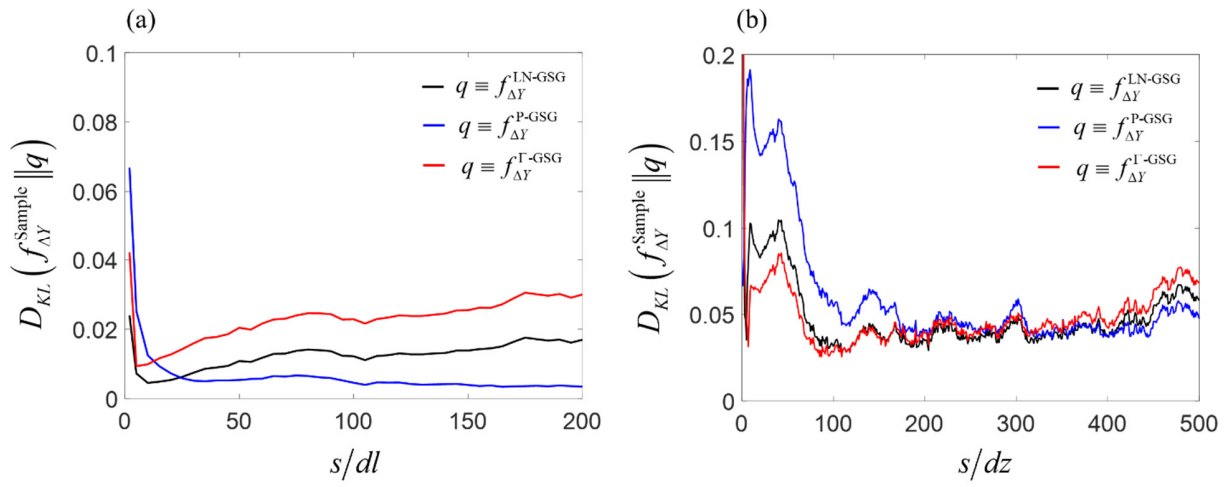
511 **Figure 8.** Estimates of ρ_G obtained via MOM_B for all tested GSG model formulations versus
 512 (normalized) lag for (a) Dataset 1 and (b) Dataset 2; 95% CIs (evaluated through bootstrapping)
 513 associated with ρ_G estimates at four lags are also shown.
 514

515



516

517 **Figure 9.** Sample pdfs (symbols) of Y' and of ΔY at two selected lags for Dataset 1 (a, b) and
 518 Dataset 2 (c, d). Theoretical distributions $f_{Y'}$ and $f_{\Delta Y}$ (Eqs. (B4) and (B8)) computed by using
 519 model parameters estimated via MOM_A (a, c) and MOM_B (b, d) are also shown.



520

521 **Figure 10.** Kullback-Leibler divergence, D_{KL} , between sample and analytical pdfs of ΔY versus
 522 (normalized) lag for (a) Dataset 1 and (b) Dataset 2.

523

524

525 **Appendix A: Analytical formulation of the GSG model for a general distributional form of**
 526 **the subordinator**

527 The theoretical framework of the GSG model is here presented considering a general
 528 distributional form of the subordinator. We do so by deriving the analytical expressions of (i) pdf,
 529 statistical moments and standardized kurtosis of the parent variable Y' ; (ii) pdf, statistical moments
 530 and standardized kurtosis of increments, ΔY , as a function of separation lag; (iii) covariance and
 531 variogram functions as well as integral scale of Y' .

532 Substituting Eq. (4) into Eq. (3) yields

$$533 \quad f_{Y'_1, Y'_2}(y'_1, y'_2) = \frac{1}{2\pi\sigma_G^2\sqrt{1-\rho_G^2}} \int_0^\infty \int_0^\infty f_{U_1}(u_1) f_{U_2}(u_2) e^{-\frac{1}{2\sigma_G^2(1-\rho_G^2)}\left(\frac{y_1'^2+y_2'^2}{u_1^2+u_2^2}-2\rho_G\frac{y_1'y_2'}{u_1u_2}\right)} \frac{du_2}{u_2} \frac{du_1}{u_1}. \quad (\text{A1})$$

534 The marginal pdf of Y' can then be obtained from Eq. (A1) as

$$535 \quad f_{Y'}(y') = \int_{-\infty}^\infty f_{Y'_1, Y'_2}(y'_1, y'_2 = y') dy'_1 = \frac{1}{\sqrt{2\pi}\sigma_G} \int_0^\infty f_U(u) e^{-\frac{1}{2\sigma_G^2} \frac{y'^2}{u^2}} \frac{du}{u}. \quad (\text{A2})$$

536 All odd-order statistical moments of Y' identically vanish, whereas variance, kurtosis and
 537 (in general) q -th even order moments can be respectively expressed as

$$538 \quad \langle Y'^2 \rangle = \int_{-\infty}^{+\infty} y'^2 f_{Y'}(y') dy' = \sigma_G^2 \int_0^\infty u^2 f_U(u) du = \sigma_G^2 \langle U^2 \rangle, \quad (\text{A3})$$

$$539 \quad \langle Y'^4 \rangle = \int_{-\infty}^{+\infty} y'^4 f_{Y'}(y') dy' = 3\sigma_G^4 \int_0^\infty u^4 f_U(u) du = 3\sigma_G^4 \langle U^4 \rangle, \quad (\text{A4})$$

$$540 \quad \langle Y'^q \rangle = \int_{-\infty}^{+\infty} y'^q f_{Y'}(y') dy' = \frac{2^{\frac{q}{2}}}{\sqrt{\pi}} \Gamma\left(\frac{1+q}{2}\right) \sigma_G^q \int_0^\infty u^q f_U(u) du = \langle G^q \rangle \langle U^q \rangle. \quad (\text{A5})$$

541 The standardized kurtosis of Y' is then given by

$$542 \quad \kappa_{Y'} = \frac{\langle Y'^4 \rangle}{\langle Y'^2 \rangle^2} = \frac{3\langle U^4 \rangle}{\langle U^2 \rangle^2} \quad (\text{A6})$$

543 and depends only on the subordinator (and not on G).

544 The pdf of incremental values, $\Delta Y(\mathbf{s}) = Y'(\mathbf{x} + \mathbf{s}) - Y'(\mathbf{x})$, can be evaluated as

$$545 \quad f_{\Delta Y}(\Delta y) = \int_{-\infty}^{\infty} f_{Y_1 Y_2}(\Delta y + y'_2, y'_2) dy'_2 = \frac{1}{\sqrt{2\pi}\sigma_G} \int_0^{\infty} \int_0^{\infty} f_{U_1}(u_1) f_{U_2}(u_2) \frac{e^{-\frac{\Delta y^2}{2\sigma_G^2 r^2}}}{r} du_2 du_1, \quad (\text{A7})$$

546 with $r = \sqrt{u_1^2 + u_2^2 - 2\rho_G u_1 u_2}$. Odd-order moments of ΔY are identically zero, whereas variance,
547 kurtosis, and moments of even order q can be respectively expressed as

$$548 \quad \langle \Delta Y^2 \rangle = \sigma_G^2 \int_0^{\infty} \int_0^{\infty} r^2 f_{U_1}(u_1) f_{U_2}(u_2) du_2 du_1 = 2\sigma_G^2 [\langle U^2 \rangle - \langle U \rangle^2 \rho_G], \quad (\text{A8})$$

$$549 \quad \langle \Delta Y^4 \rangle = 3\sigma_G^4 \int_0^{\infty} \int_0^{\infty} r^4 f_{U_1}(u_1) f_{U_2}(u_2) du_2 du_1 = 6\sigma_G^4 [\langle U^4 \rangle - 4\langle U^3 \rangle \langle U \rangle \rho_G + \langle U^2 \rangle^2 (1 + 2\rho_G^2)], \quad (\text{A9})$$

$$550 \quad \langle \Delta Y^q \rangle = \frac{\sigma_G^q}{\sqrt{\pi}} 2^{\frac{q}{2}} \Gamma\left(\frac{q+1}{2}\right) \int_0^{\infty} \int_0^{\infty} r^q f_{U_1}(u_1) f_{U_2}(u_2) du_2 du_1$$

$$= \sum_{k=0}^q (-1)^{q-k} \binom{q}{k} \langle U^k \rangle \langle U^{q-k} \rangle \langle G(x)^k G(x+s)^{q-k} \rangle. \quad (\text{A10})$$

551 The standardized kurtosis of ΔY is derived from Eqs. (A8) and (A9) as

$$552 \quad \kappa_{\Delta Y} = \frac{\langle \Delta Y^4 \rangle}{\langle \Delta Y^2 \rangle^2} = \frac{3}{2} \frac{\langle U^4 \rangle - 4\langle U^3 \rangle \langle U \rangle \rho_G + \langle U^2 \rangle^2 (1 + 2\rho_G^2)}{[\langle U^2 \rangle - \langle U \rangle^2 \rho_G]^2}, \quad (\text{A11})$$

553 the latter depending on the subordinator and on the correlation coefficient ρ_G (but not on σ_G^2).

554 The Covariance of Y' between two points \mathbf{x}_1 and \mathbf{x}_2 is

$$555 \quad C_{Y'}(\mathbf{x}_1, \mathbf{x}_2) = \langle Y'(\mathbf{x}_1) Y'(\mathbf{x}_2) \rangle$$

$$= \langle U(\mathbf{x}_1) U(\mathbf{x}_2) \rangle \langle G(\mathbf{x}_1) G(\mathbf{x}_2) \rangle = \langle U(\mathbf{x}_1) U(\mathbf{x}_2) \rangle \sigma_G^2 \rho_G(\mathbf{x}_1, \mathbf{x}_2). \quad (\text{A12})$$

556 From Eq. (A12), one derives

$$557 \quad C_{Y'}(0) = \sigma_{Y'}^2 = \langle U^2 \rangle \sigma_G^2, \quad C_{Y'}(s > 0) = \langle U \rangle^2 \sigma_G^2 \rho_G. \quad (\text{A13})$$

558 Note that according to Eq. (A13) the covariance $C_{Y'}$ of the Sub-Gaussian field is discontinuous at
559 the origin, i.e., at $s = 0$, thus exhibiting a nugget effect. The variogram of Y' can be evaluated from
560 Eq. (A8) as

$$561 \quad \gamma_{Y'} = \frac{\langle \Delta Y^2 \rangle}{2} = \sigma_G^2 [\langle U^2 \rangle - \langle U \rangle^2 \rho_G] = \sigma_G^2 [\sigma_U^2 + \langle U \rangle^2 - \langle U \rangle^2 \rho_G] = \sigma_G^2 \sigma_U^2 + \langle U \rangle^2 \gamma_G \quad (\text{A14})$$

562 and is characterized by a nugget effect, quantified by $\sigma_G^2 \sigma_U^2$, $\gamma_G = \sigma_G^2 (1 - \rho_G)$ and σ_U^2 being the
 563 variogram of G and the variance of U , respectively.

564 The integral scale of Y' can be obtained by making use of Eq.(A12) as

$$565 \quad I_{Y'} = \frac{\langle U \rangle^2}{\langle U^2 \rangle} I_G = \frac{\langle U \rangle^2}{\sigma_U^2 + \langle U \rangle^2} I_G, \quad (\text{A15})$$

566 so that one can recognize that $0 < I_{Y'} < I_G$, independent of the type of subordinator considered. An
 567 increase of σ_U^2 results in a decrease of the (integral) correlation scale of Y' .

568

569 **Appendix B: GSG formulation for lognormal, Pareto, and Gamma distribution of U**

570 Here, we consider U_1 and U_2 to be described by (i) a lognormal distribution,

571 $U_i \sim \ln N(\mu, (2 - \alpha)^2)$, (ii) a Pareto distribution, $U_i \sim PD(a, b)$, and (iii) a Gamma distribution,

572 $U_i \sim \Gamma(k, \theta)$, i.e.,

$$573 \quad f_{U_i}(u_i) = \frac{e^{-\frac{(\ln u_i - \mu)^2}{2(2 - \alpha)^2}}}{\sqrt{2\pi} u_i (2 - \alpha)} \quad \text{with } \alpha < 2; u_i \in (0, +\infty) \quad (\text{B1a})$$

$$574 \quad f_{U_i}(u_i) = \frac{ab^a}{u_i^{a+1}} \quad \text{with } a > 0; b > 0; u_i \in [b, +\infty) \quad (\text{B1b})$$

$$575 \quad f_{U_i}(u_i) = \frac{u_i^{k-1} e^{-\frac{u_i}{\theta}}}{\Gamma(k) \theta^k} \quad \text{with } k > 0; \theta > 0; u_i \in (0, +\infty); \Gamma(k) = \int_0^{\infty} x^{k-1} e^{-x} dx \quad (\text{B1c})$$

576 here, $i = 1, 2$; α , a , and k are shape parameters, while e^μ , b , and θ are scale parameters. Note that
 577 the exponential distribution can be obtained from Eq. (B1c) by setting $k = 1$.

578 The q -th order raw moment of U is

$$579 \quad \langle U^q \rangle = \begin{cases} e^{q\mu + \frac{q^2}{2}(2-\alpha)^2} & \text{if } U_i \sim \ln N(\mu, (2-\alpha)^2), \\ \frac{ab^q}{a-q} & \text{if } U_i \sim PD(a, b), \\ \frac{\Gamma(k+q)\theta^q}{\Gamma(k)} & \text{if } U_i \sim \Gamma(k, \theta), \end{cases} \quad (B2)$$

580 the variance being equal to

$$581 \quad \sigma_U^2 = \begin{cases} e^{2\mu + (2-\alpha)^2} (e^{(2-\alpha)^2} - 1) & \text{if } U_i \sim \ln N(\mu, (2-\alpha)^2), \\ ab^2 / [(a-1)^2(a-2)] & \text{if } U_i \sim PD(a, b), \\ k\theta^2 & \text{if } U_i \sim \Gamma(k, \theta). \end{cases} \quad (B3)$$

582 As specified in Section 2.2, the application of Method of Moment (MOM) requires U_i to
583 have finite raw moments up to order $q = 2N_P$ (thus implying $a > 4$ in (B2)).

584 Substituting Eq. (B1) into Eq. (A2) yields the following marginal pdf of Y'

$$585 \quad f_{Y'}(y) = \begin{cases} \frac{1}{2\pi \sigma_G e^\mu (2-\alpha)} \int_0^\infty e^{-\frac{1}{2} \left[\left(\frac{\ln u}{2-\alpha} \right)^2 + \left(\frac{y}{\sigma_G e^\mu u} \right)^2 \right]} \frac{du}{u^2} & \text{for LN-GSG,} \\ \frac{a}{\sqrt{2\pi} (b\sigma_G)^{-a}} \int_{b\sigma_G}^\infty e^{-\frac{1}{2} \left(\frac{y}{u} \right)^2} \frac{du}{u^{a+2}} & \text{for P-GSG,} \\ \frac{1}{\sqrt{2\pi} \sigma_G \theta \Gamma(k)} \int_0^\infty u^{k-2} e^{-\left[u + \frac{y^2}{2\sigma_G^2 \theta^2 u^2} \right]} du & \text{for } \Gamma\text{-GSG.} \end{cases} \quad (B4)$$

586 Note that LN-GSG coincides with a normal-lognormal distribution (NLN) when $\mu = 0$. The
587 latter has been shown to well represent some financial (Clark, 1973) and environmental
588 (Guadagnini et al., 2015) data. Making use of Eqs. (A3) - (A4) and (B4), variance, kurtosis and
589 standardized kurtosis of Y' are respectively given by

$$590 \quad \langle Y'^2 \rangle = \sigma_G^2 \begin{cases} e^{2\mu} e^{2(2-\alpha)^2} & \text{for LN-GSG,} \\ b^2 \frac{a}{a-2} & \text{for P-GSG,} \\ \theta^2 k(1+k) & \text{for } \Gamma\text{-GSG.} \end{cases} \quad (B5)$$

$$591 \quad \langle Y^4 \rangle = 3\sigma_G^2 \begin{cases} e^{4\mu} e^{8(2-\alpha)^2} & \text{for LN-GSG,} \\ b^4 \frac{a}{a-4} & \text{for P-GSG,} \\ \theta^4 \prod_{i=0}^3 (k+i) & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B6)$$

$$592 \quad \kappa_{Y'} = 3 \begin{cases} e^{4(2-\alpha)^2} & \text{for LN-GSG,} \\ \frac{(a-2)^2}{a(a-4)} & \text{for P-GSG,} \\ 1 + \frac{4k+6}{k(k+1)} & \text{for } \Gamma\text{-GSG.} \end{cases} \quad (B7)$$

593 Substituting Eq. (B1) into Eq. (A7) yields the following expressions for the pdf of ΔY

$$594 \quad f_{\Delta Y}(\Delta y) = \frac{1}{\sqrt{2}} \begin{cases} \frac{\sqrt{\pi}}{2\pi^2 (2-\alpha)^2 \sigma_G e^\mu} \int_0^\infty \int_0^\infty e^{-\frac{1}{2} \left[\frac{1}{(2-\alpha)^2} (\ln^2 u_1 + \ln^2 u_2) + \frac{\Delta y^2}{\sigma_G^2 e^{2\mu} r^2} \right]} \frac{du_2 du_1}{u_2 u_1 r} & \text{for LN-GSG,} \\ \frac{a^2}{\sqrt{\pi} (b\sigma_G)^{-2a}} \int_0^\infty \int_0^\infty e^{-\frac{1}{2} \frac{\Delta y^2}{r^2}} \frac{du_2 du_1}{u_1^{a+1} u_2^{a+1} r} & \text{for P-GSG,} \\ \frac{1}{\sqrt{\pi} (\sigma_G \theta)^{2k} \Gamma^2(k)} \int_0^\infty \int_0^\infty (u_1 u_2)^{k-1} e^{-\left[\frac{1}{\sigma_G \theta} (u_1 + u_2) + \frac{1}{r^2} \frac{\Delta y^2}{2} \right]} \frac{du_2 du_1}{r} & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B8)$$

595 with $r = \sqrt{u_1^2 + u_2^2 - 2\rho_G u_1 u_2}$.

596 Making use of Eqs. (A8)-(A11) and (B2), variance, kurtosis and standardized kurtosis of ΔY are
597 respectively given by

$$598 \quad \langle \Delta Y^2 \rangle = 2\sigma_G^2 \begin{cases} e^{2\mu} e^{(2-\alpha)^2} \left[e^{(2-\alpha)^2} - \rho_G \right] & \text{for LN-GSG,} \\ b^2 a \left[\frac{1}{a-2} - \frac{a\rho_G}{(a-1)^2} \right] & \text{for P-GSG,} \\ k\theta^2 \left[1+k - \rho_G k \right] & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B9)$$

$$\begin{aligned}
 \langle \Delta Y^4 \rangle = 6\sigma_G^4 & \begin{cases} e^{4\mu} e^{4(2-\alpha)^2} \left[1 + e^{4(2-\alpha)^2} - 4e^{(2-\alpha)^2} \rho_G + 2\rho_G^2 \right] & \text{for LN-GSG,} \\ b^4 a^2 \left[\frac{1}{a(a-4)} - \frac{4\rho_G}{(a-1)(a-3)} + \frac{(1+2\rho_G^2)}{(a-2)^2} \right] & \text{for P-GSG,} \\ 2\theta^4 k^2 (k+1) \left[\frac{3}{k} + k + 3 + \rho_G (\rho_G (k+1) - 2(k+2)) \right] & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B10)
 \end{aligned}$$

$$\begin{aligned}
 \kappa_{\Delta Y} = 3 & \begin{cases} e^{2(2-\alpha)^2} \left\{ 1 + \frac{1}{2} \left(\frac{e^{2(2-\alpha)^2} - 1}{e^{(2-\alpha)^2} - \rho_G} \right)^2 \right\} & \text{for LN-GSG,} \\ \frac{1}{2} \left[\frac{1}{a(a-4)} - \frac{4\rho_G}{(a-1)(a-3)} + \frac{1+2\rho_G^2}{(a-2)^2} \right] \left[\frac{1}{(a-2)} - \frac{a\rho_G}{(a-1)^2} \right]^{-2} & \text{for P-GSG,} \\ 1 + \frac{1}{k} + \frac{(k+1)(2+k+\rho_G^2 k - 2k\rho_G)}{k(k+1-\rho_G k)^2} & \text{for } \Gamma\text{-GSG.} \end{cases} \quad (B11)
 \end{aligned}$$

601 The variogram, $\gamma_{Y'}$, covariance, $C_{Y'}(s)$, for $s > 0$ (note that $C_{Y'}(s = 0)$ coincides with $\langle Y'^2 \rangle$)
 602 evaluated in Eq. (B5)) and integral scale, $I_{Y'}$, of Y' , can be derived from Eqs. (A13) – (A15) and
 603 (B2) as

$$\begin{aligned}
 \gamma_{Y'} = & \begin{cases} \sigma_G^2 e^{2\mu+(2-\alpha)^2} \left(e^{(2-\alpha)^2} - 1 \right) + e^{2\mu+(2-\alpha)^2} \gamma_G & \text{for LN-GSG,} \\ \sigma_G^2 \frac{b^2 a}{(a-1)^2 (a-2)} + \frac{a^2 b^2}{(a-1)^2} \gamma_G & \text{for P-GSG,} \\ \sigma_G^2 k \theta^2 + k^2 \theta^2 \gamma_G & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B12)
 \end{aligned}$$

$$\begin{aligned}
 C_{Y'}(s > 0) = \sigma_G^2 \rho_G & \begin{cases} e^{2\mu+(2-\alpha)^2} & \text{for LN-GSG,} \\ \frac{a^2 b^2}{(a-1)^2} & \text{for P-GSG,} \\ k^2 \theta^2 & \text{for } \Gamma\text{-GSG,} \end{cases} \quad (B13)
 \end{aligned}$$

$$I_{Y'} = I_G \begin{cases} e^{-(2-\alpha)^2} & \text{for LN-GSG,} \\ a \frac{a-2}{(a-1)^2} & \text{for P-GSG,} \\ 1 - \frac{1}{1+k} & \text{for } \Gamma\text{-GSG.} \end{cases} \quad (\text{B14})$$

607 It is thus seen that when the pdf of U tends to the Dirac delta function (i.e., when $\alpha \rightarrow 2$ for
 608 LN-GSG; $a \rightarrow \infty$ for P-GSG; or $k \rightarrow \infty$ for Γ -GSG), then $I_{Y'} \rightarrow I_G$. Otherwise, $I_{Y'}$ is smaller than
 609 I_G (regardless the subordinator adopted), while never vanishing. The range of values which can
 610 be undertaken by $I_{Y'}$ depends on the type of subordinator employed and on the threshold values of
 611 the shape parameters (see Section 2.1). The broadest range of variability of $I_{Y'}$ is associated with
 612 the LN-GSG, where $0.33 < I_{Y'} / I_G < 1$. Otherwise, the smallest interval is obtained through P-
 613 GSG, where $0.89 < I_{Y'} / I_G < 1$, Γ -GSG being associated with $0.5 < I_{Y'} / I_G < 1$.

614 **Data**

615 Datasets are available at: [https://data.mendeley.com/datasets/trdgwfwsvn/draft?a=ee55e214-](https://data.mendeley.com/datasets/trdgwfwsvn/draft?a=ee55e214-386a-48f4-88bf-461c5ddaf7ec)
 616 [386a-48f4-88bf-461c5ddaf7ec](https://data.mendeley.com/datasets/trdgwfwsvn/draft?a=ee55e214-386a-48f4-88bf-461c5ddaf7ec)

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