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Pilar Sango Solanas, Kevin Tse Ve Koon, Hélène Ratiney, Fabien Millioz, Cyrielle Caussy, et al.. Harmonic wideband simultaneous dual-frequency MR Elastography. NMR in Biomedicine, 2021, 34 (2), pp.e4442. 10.1002/nbm.4442 . hal-03007661

HAL Id: hal-03007661

<https://hal.science/hal-03007661>

Submitted on 4 Jan 2021

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Harmonic Wideband simultaneous dual-frequency MR Elastography

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Word count: 4407

Key words (3): MR Elastography, multifrequency, phase offset undersampling

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List of abbreviations

A_{exc} , excitation shear wave amplitude

f_{MEG} , motion encoding gradient oscillating frequency

f_{exc} , excitation shear wave frequency

f_s , sampling frequency

G' , shear storage modulus

G'' , shear loss modulus

G^* , complex shear modulus

MEG, motion encoding gradient

MRE, Magnetic Resonance Elastography

NAFLD, Nonalcoholic fatty liver disease

N_G , number of motion encoding gradient cycles

PNR, Phase-to-Noise Ratio

q , fractional encoding ratio

RARE, Rapid Acquisition with Refocused Echoes

ROI, Region of interest

SNR, Signal-to-Noise Ratio.

SUMMARY ABSTRACT

Magnetic resonance elastography (MRE) is used to non-invasively quantify viscoelastic properties of tissues based on the measurement of propagation characteristics of shear waves. Because some of these viscoelastic parameters show a frequency dependence, multifrequency analysis allows us to measure the wave propagation dispersion, leading to a better characterization of tissue properties. Conventionally, motion encoding gradients (MEGs) oscillating at the same frequency as the mechanical excitation encode motion. Hence, multifrequency data is usually obtained by sequentially repeating monochromatic wave excitations experiments at different frequencies. The result is that the total acquisition time is multiplied by a factor corresponding to the number of repetitions of monofrequency experiments, which is a major limitation of multifrequency MRE. In order to make it more accessible, a novel single-shot harmonic wideband dual-frequency MRE method is proposed. Two superposed shear waves of different frequencies are simultaneously generated and propagate in a sample. Trapezoidal oscillating MEGs are used to encode mechanical vibrations having frequencies that are an odd multiple of the MEG frequency. The number of phase offsets is optimized to reduce the acquisition time. For this purpose, a sampling method not respecting the Shannon theorem is used to produce a controlled temporal aliasing that allows us to encode both frequencies without any additional examination time. Phantom experiments were run to compare conventional monofrequency MRE with the single-shot dual-frequency MRE method and showed excellent agreement between the reconstructed shear storage moduli G' . In addition, dual-frequency MRE yielded an increased signal-to-noise ratio compared with conventional monofrequency MRE acquisitions when encoding the high frequency component. The novel proposed multifrequency MRE method could be applied to simultaneously acquire more than two frequency components, reducing examination time. Further studies are needed to confirm its applicability in preclinical and clinical models.

INTRODUCTION

The development of a disease is associated with morphological and structural changes in the affected organs. Various imaging techniques enable to obtain images of organs as well as the non-invasive assessment of their different physical characteristics. Mechanical properties of these tissues are of particular interest as they undergo appreciable modifications which can be quantified¹, enabling evaluation of the stage of a disease or assessment of longitudinal changes such as progression or regression of a disease stage. Elastography techniques using ultrasound² or MRI³ are used to non-invasively assess the stiffness of tissues by providing information on viscoelastic properties. Specifically, magnetic resonance elastography³ (MRE) quantifies, based on viscoelastic models, the mechanical properties of tissues from measurements of the shear wave propagation through an organ.

Several studies have already shown that MRE is a relevant non-invasive tool in clinical practice for characterizing liver fibrosis disease evolution by providing comparable results to liver biopsy⁴⁻⁷ without being limited to a small number of samples. Actually, MRE is a recognized non-invasive modality to assess liver fibrosis stage in patients with nonalcoholic fatty liver disease (NAFLD). Indeed, MRE generally outperforms ultrasound modalities for the staging of liver fibrosis in patients with NAFLD^{6,8}. Clinical researches have also exploited the valuable tissue mechanical characterization technique of MRE to target other organs with different applications⁹ such as the diagnosis of cardiac abnormalities¹⁰⁻¹², breast cancer¹³⁻¹⁶ and brain viscoelastic properties measurement¹⁷⁻²⁰. However, the utility of MRE in clinical practice is limited by its long acquisition duration²¹.

The MRE biomarker currently used to assess the severity of the different mentioned diseases is the shear storage modulus (G') or more rarely the complex shear modulus (G^*), both parameters are frequency dependent and represent elastic and viscous characteristics of tissues. The frequency dispersion of G' over a wide dynamic range can be evaluated to detect microstructural²²⁻²⁴ changes, providing thus a more detailed mechanical characterization of tissues^{22,25}. It was demonstrated that viscoelastic power-law behavior of living soft tissue is altered by disease^{26,27} and presently, multifrequency MRE is one of the few imaging-based techniques able to measure the dispersion of the complex shear modulus of soft tissue *in vivo*.

In conventional MRE, oscillating motion encoding gradients (MEGs) encode the cyclic motion of nuclear spins when both MEGs and spins are oscillating at the same frequency^{3,21}. But there are also other motion encoding strategies in which oscillating MEGs are not needed^{28–32}.

Conventionally, to encode low excitation frequencies, long MEGs are used. This implies an increase in the TE and consequently, the signal to noise ratio (SNR) is degraded especially for short T2 tissues. High excitation frequencies can overcome this limitation. However, they are not appropriate for soft deep tissues for which shear waves are quickly attenuated. Each acquisition provides the mechanical information of a single excitation frequency and multifrequency results have to be obtained by repeating the experiences with monochromatic wave excitations at different frequencies. This yields very long total acquisition times, which is one of the major limitations of the technique in routine clinical practice. Therefore, reducing its duration would ultimately improve its efficiency, reduce its financial cost and thus increase its availability in the future.

In the last decades some multifrequency MRE methods have been developed. To this aim, the resulting spectral sensitivity of a time-limited sinusoidal MEGs has been exploited: for a given MEG, motion of spins oscillating at a frequency contained in the MEG's sensitivity bandwidth can be encoded.

The study by Rump *et al.*³³ initially proposed a fractional encoding MRE method in which only a fraction of one vibration cycle per TR is used for motion sensitization, defined by the fractional encoding ratio ($q = \frac{f_{exc}}{f_{MEG}}$) where f_{MEG} and f_{exc} are the frequencies of the MEG and the mechanical excitation respectively. In the liver, they used a 100 Hz MEG to encode a 76 Hz excitation frequency ($q = 0.76$). This approach demonstrated that low frequencies can be encoded with shorter TEs yielding higher SNR. From then on, several studies have exploited this strategy for different applications^{10,26,34–37}. For instance, Klatt *et al.*³⁴ and Asbach *et al.*³⁵ demonstrated the high consistency and reproducibility of MRE in human liver when using a 50 Hz sinusoidal MEG to simultaneously encode four superimposed sinusoidal waveforms with frequencies of 25.0, 37.5, 50.0, and 62.5 Hz ($q = 0.5, 0.75, 1, 1.25$). And more recently, Garteiser *et al.*³⁷ introduced a MRE sequence that enabled the acquisition of multiple frequencies on several slices in the same examination. These approaches proved the feasibility of fractional encoding to simultaneously acquire several mechanical excitation frequencies to extend the dynamic range of a single MRE

examination instead of being monochromatic and therefore increasing the acquisition speed.

Despite these advantages, fractional encoding is usually limited *in vivo* to relatively large targeted organs where low frequencies are needed. Current methods are thus not suitable for the encoding of motion in small organs where higher frequencies are required. In view of carrying out multifrequency MRE and extracting the subsequent power laws, accessing data covering much higher frequency ranges would represent a substantial added value. However, this strategy is confronted to two major limitations. The first limitation, as previously mentioned, is linked to the higher attenuation of shear waves at higher frequencies, which therefore limits their penetration depth and the region of interest. The second limitation is related to hardware specifications of MR scanners. Indeed, gradient amplitudes and switching rates are limited which is particularly true in clinical systems. Thus, even if we were able to generate and propagate higher frequencies shear waves, encoding them in the phase of the MR signal would be challenging.

Here, we first propose a harmonic wideband MRE method, demonstrated by applying low oscillating MEG frequencies to encode higher frequencies, usually hard to be reached by the gradient system of MRI scanners. In this way, we overcome the limitation due to the hardware of MR scanner although a remaining challenge is the propagation of high frequencies through tissues. Then, this strategy is combined with a dual-frequency mechanical excitation, yielding the simultaneous encoding and acquisition of both frequency motions. This allows exploring a wider range of frequencies that may improve the performance of the power-law fit analysis of viscoelastic parameters frequency dispersion. Finally, the number of phase offset is optimized in order to still reduce the total acquisition time. For this purpose, a non-respecting Shannon sampling method is used to produce a controlled temporal aliasing that allows to encode both frequencies in a minimal acquisition time.

The main objectives of this study were first, to demonstrate the feasibility of the novel single-shot wideband dual-frequency method in phantom experiments and second, to test the phase offset undersampling method based on temporal aliasing, allowing to significantly reduce the total acquisition time.

THEORETICAL BACKGROUND

Motion encoding and viscoelastic parameters reconstruction

MRE is a phase-contrast motion encoding method. An oscillating gradient $\vec{G}(t)$ encodes the cyclic motion of a spin $\vec{r}(t)$ at a given frequency into the MRI signal phase ϕ ³:

$$\phi = \gamma \int_{t_0}^{t_0+T_{enc}} \vec{G}(t) \cdot \vec{r}(t) dt, \quad (1)$$

where γ is the gyromagnetic ratio and T_{enc} denotes the application time of MEG cycles.

The cumulated phase reaches maximum values when the frequency of motion is equal to the MEG frequency. However even for sinusoidal MEGs, this cumulated phase is still appreciable if the motion frequency belongs to the sensitivity band of the MEG as shown in Figure 1a. This is the underlying principle of fractional encoding. This sensitivity band has a periodical behavior when using a trapezoidal-shaped MEG since the latter can be decomposed as a Fourier series that is as a sum of odd-harmonics: non-resonant sensitivity bands appear around the odd-multiples of the MEG frequency ($f_{exc} = (2n+1) \cdot f_{MEG}$ with $n=0,1,2,\dots$) as can be seen in Figure 1b.

Equation [1] was used prior to the experiments to numerically calculate the cumulated encoded phase at the driving frequency used. We considered a shear excitation frequency f_{exc} varying between 0 and 3000 Hz and a single period of 300 Hz MEG with either a sinusoidal (Figure 1a) or a trapezoidal shape (Figure 1b). Here, we can see the sensitivity bands around 300, 900, 1500 and 2100 Hz that appear when using a 300 Hz trapezoidal MEG.

In our experiments, we considered only plane waves and only one direction was sensitized with the MEGs. The propagating wave was represented by a scalar wave field $u(x,y,t)$ that was proportional to the unwrapped accumulated phase (φ_{unw}). The displacement in each pixel was calculated as follow:

$$u(x, y, t) = C \cdot \varphi_{unw}(x, y, t), \quad (2)$$

where C is a conversion factor depending on the MEGs amplitude, the excitation wave frequency and amplitude. The number of acquired phase offsets gave the size of the temporal dimension.

Afterwards, a temporal Fourier transform was applied to the wave field in order to work in the frequency domain $U(x,y,\omega)$. In the case of dual-frequency acquisitions, the two harmonics corresponding to both excitation frequencies are kept.

The Helmholtz equation inversion applied to a spin harmonically oscillating at a given frequency is expressed by

$$\rho\omega^2 U(x,y,\omega) + G^*(\omega)\Delta U(x,y,\omega) = 0, \quad (3)$$

where $G^*(\omega)$ and ρ denote the complex shear modulus and the density of the tissue respectively. From this equation we can derive the expression of $G^*(\omega)$.

The complex shear modulus can be expressed by $G^* = G' + iG''$, where G' represents the storage modulus, which is related to the tissue elasticity and G'' the loss modulus, related to the viscosity of the medium.

In our results, we focused on the value of the parameter G' but G'' were also calculated. Values for different excitation frequencies and with different acquisition modes (mono and dual-frequency acquisitions) are compared.

Undersampling

According to the Nyquist-Shannon sampling theorem, $f_s > 2 \cdot f_{max}$, where f_s is the sampling frequency and f_{max} is the maximum frequency contained in the acquired signal.

For a correct sampling of a monochromatic wave, at least three sampling steps equally spaced over a period of the signal wave are needed. But in fact, in conventional MRE, four equally spaced phase offsets are usually acquired in order to assure the acquisitions of maxima and minima phase points in the shear wave time evolution. This implies that to do repetitive monofrequency acquisitions at n different mechanical excitation frequencies, the total acquisition time would be n times greater.

However, in our specific case of a two components wave motion with an odd harmonic relationship between them ($f_2 = 3f_1$), a small number of phase offsets could be able to encode both frequencies because of the temporal aliasing that appears when the Shannon theorem is not respected to sample polychromatic waves. This aliasing phenomenon is shown in Figure 2. Here, we considered a polychromatic wave resulting from the summation of two shear waves of $f_1 = 300$ Hz and $f_2 = 900$ Hz and we realized a sampling

with a sampling period $T_s = T_1/N = 1/(N f_1)$ (this means that we acquire N equally-spaced time points over one period of the fundamental frequency shear wave). As the spectrum of the original signal (represented as a block in dashed lines) is replicated every f_s , when the sampling frequency respects the Nyquist criterion (N should be greater than 6 since the second component $f_2 = 3f_1$), there is no overlap between adjacent blocks (Figure 2a). When the Nyquist criterion is not respected, the blocks overlap leading to contamination from components of adjacent blocks (Figure 2b-c). For instance, in Figure 2b, both frequencies overlap in the same acquired points meaning the two components cannot be separated. However, as can be seen in Figure 2c, due to the fact that the second sought component is an odd harmonic of the first one, even in the case of 5 sampling steps, the resulting aliasing doesn't prevent separation of the two components since the aliased higher component corresponds to a different reconstructed point than the lower component in Fourier space. Therefore, such a chosen configuration with odd harmonics in addition to a fundamental component produces a controlled frequency aliasing meaning that it is possible to encode simultaneously several odd multiples of the fundamental frequency with a smaller number of phase offsets. In general, to simultaneously encode n different shear wave frequencies, $2n + 1$ phase offsets would be needed.

MATERIALS AND METHODS

MRE setup

All acquisitions were performed on a preclinical 4.7 T MRI scanner (BioSpec Bruker system, Ettlingen, Germany), with a 60 mm inner diameter volume coil in transmit/receive mode.

Experiments were carried out on a uniform phantom composed of 75%-standard plastisol³⁸ (Plastileurre Standard, Bricoleurre, France) and 25% softener (Di-2-ethylhexyl adipate). Relaxation times measured using a RARE sequence with different TEs and TRs gave $[T_1, T_2] = [300, 35]$ ms. Although the phantom used here could not be characterized from a mechanical point of view by a complementary method, in previous studies we compared values of storage and loss moduli obtained both with MRE and

High-Frequency Rheometry for different plastisol phantoms, confirming that the complex modulus of plastisol phantoms can be reliably assessed using MRE³⁸.

A piezoelectric actuator (CEDRAT Technologies) was used as mechanical transducer for motion generation. A dual output wave generator was used to produce two harmonic signals of frequencies 300 and 900 Hz. They were then summed and fed to the actuator thus resulting in the latter generating waves at both frequencies.

MRE acquisition

The MRE sequence used was based on a conventional RARE sequence. Inserted MEGs may have variable sinusoidal or trapezoidal shapes with adjustable frequencies and amplitudes. Two acquisitions with inversion of the MEG polarity were made for phase images subtraction to reduce static phase errors. Oscillating MEGs were synchronized with the mechanical excitation. The wave generator was triggered at the beginning of each TR in order to reach a steady-state wave propagation during the MEGs application period.

Fixed acquisition parameters were axial orientation, TE/TR=16.67/1000 ms, single slice of 1 mm thickness, field of view = 4x4 cm², matrix size =128 x 128, RARE factor = 4.

The number of MEG cycles (N_G) and the amplitude of the shear waves (A_{exc}) were adjusted according to the excitation frequency. In order to account for the reduction of the encoding efficiency of the MEG at 900 Hz, the amplitude of this shear wave component was tripled with respect to the 300 Hz component. The amplitude of MEGs was always the same (250 mT/m). A single direction was sensitized with the MEGs. Further details of the acquisitions are summarized in Table 1.

In our protocol, several acquisitions were realized: monochromatic acquisitions synchronized on the fundamental frequency with $f_{exc}=f_{MEG}$ (acquisitions 1 and 2), monochromatic acquisitions synchronized on the first odd harmonic with $f_{exc}=3 \cdot f_{MEG}$ (3,4) and a dual-frequency acquisition with $f_{MEG}=\min(f_{exc})$ (5). Numbers in brackets correspond to the acquisition numbers in Table 1. During the dual-frequency acquisition (Figure 3), the oscillating trapezoidal gradients were fixed to the lowest frequency and allowed encoding of both frequencies due to the sensitive motion encoding wideband of MEGs, as explained previously in the MRE subchapter. According to the

proposed phase offset undersampling method, five equally spaced phase offset of the dual-frequency acquisition were obtained.

To avoid artifacts during the viscoelastic parameter reconstruction, the lowest shear wave frequency was chosen so that the corresponding wavelength was smaller than the phantom dimension in the propagation direction. The pair of frequencies 300 and 900 Hz was considered in our work.

Image analysis

Elastograms reconstruction

Raw phase images were unwrapped with a quality guided path following phase unwrapping algorithm³⁹. The proportional 2D displacement fields were temporally Fourier transformed. Then, each frequency was reconstructed from a different harmonic. In the multifrequency acquisition, the second and the fourth harmonic, which correspond to the fundamental frequency (f_1) and to the high frequency (f_3) respectively, were kept for the elastograms reconstruction.

Then, a spatial fourth-order 2D Butterworth filter with lower and upper thresholds of 0.25 and 2.5 cm^{-1} at 300 Hz and 1 and 3.5 cm^{-1} at 900 Hz was applied to filter and preserve the frequencies of interest as much as possible^{4,21}. The spatial filter limits were chosen based on simulations of estimated wavelength values with different cut-offs filters values (data not shown).

Finally, the Helmholtz equation inversion^{40,41} (Equation 3) was applied to the displacement fields to calculate the shear storage modulus G' and the loss modulus G'' at each frequency assuming a density of 1000 kg/m^3 .

A 3x3 median filter was applied to the elastograms. A ROI (about 1600 pixels) was manually drawn in the center of the phantom, excluding the phantom borders in order to avoid reconstruction artifacts. The mean and standard deviation of G' and G'' were calculated inside the ROI. The same ROI was used for all acquisitions and excitation frequencies.

SNR and PNR analysis

Signal-to-noise ratio (SNR) and Phase-to-Noise ratio (PNR) were extracted from magnitude and phase images respectively. SNR was calculated as the ratio of the mean

signal of a ROI in the phantom to the standard deviation in the background. PNR was calculated as the product of the SNR and the phase encoding Δ ($PNR = SNR \times \Delta$). The phase encoding Δ is defined as the difference of the maximal and minimal phase encoded along a single line parallel to the wave propagation direction (Shown in Figure 4 at right). The PNR is calculated to evaluate the motion encoding efficiency of the different acquisition modes and image quality through SNR. TE and TR were always the same in order to acquire all scans under the same conditions. Note that all SNR were measured on the same regions of interest and the same phase offset of the wave propagation and line were taken for PNR computation so as to have a reliable estimation of the phase encoding.

RESULTS

The ability of the 300 Hz trapezoidal MEG to encode the 900 Hz excitation frequency was experimentally demonstrated as seen in Figure 4 and compared to the quasi-null encoded phase of a sinusoidal one.

The phase image of the dual-frequency excitation is showed in Figure 5, as well as phase images obtained from the inversion of the two separate frequency components after temporal Fourier transform. Both vibrations are correctly encoded. Likewise, the G' elastograms are computed for each excitation frequency. Both maps present some edge artifacts resulting from the viscoelastic parameters reconstruction method but they are otherwise quite uniform and conform to the homogeneous composition of the phantom.

The different acquisition modes are compared in terms of SNR, Δ and PNR in Figure 6. Only the acquisitions with trapezoidal MEGs are considered here. Dual-frequency results present the highest values of SNR yielding the best values of PNR. PNRs were 74.3, 67.4, 77.0, 82.7 rad (indicated respectively for acquisitions 1-3 and 5). Especially, we find an improved motion encoding efficiency of the highest frequency when using the 300 Hz MEG in both mono and dual-frequency acquisitions, compared to the conventional 900 Hz MEG monofrequency acquisition.

The mean values and standard deviation of G' obtained for both excitations frequencies by using the different acquisition modes are presented in Figure 7. Storage modulus values were 20.2 ± 2 , 20.3 ± 2.2 kPa at 300 Hz (in acquisitions 1 and 5 respectively) and

23.6 ± 3.7 , 23.9 ± 2 and 23.6 ± 1.5 kPa at 900 Hz (in acquisitions 2, 3 and 5). Loss modulus values were 6.9 ± 2.6 , 6.5 ± 2.8 kPa at 300 Hz (in acquisitions 1 and 5 respectively) and 8.3 ± 1.6 , 7.9 ± 3.8 and 7.7 ± 1.4 kPa at 900 Hz (in acquisitions 2, 3 and 5). Results obtained were comparable between mono and dual-frequency acquisitions. The undersampling method shows also consistent results.

DISCUSSION

This study demonstrates the feasibility of the novel single-shot harmonic wideband dual-frequency MRE method. The simultaneous excitation of two distant frequencies (300-900 Hz) shear waves combined to the sensitive harmonic wideband of trapezoidal motion encoding gradients enabled the mechanical characterization of tissues without any additional acquisition time. In addition, this study demonstrates that the proposed undersampling method, based on a controlled frequency aliasing, allows to reduce the examination time up to 38% of the time required in conventional MRE.

Although a better mechanical characterization of tissues can be obtained with multifrequency MRE, it is mostly limited by the long acquisition times required²¹ to repeat monofrequency acquisitions. To increase the speed of multifrequency MRE acquisitions, other investigators proposed different methods to do simultaneous multifrequency MRE. First approaches utilized the hypothesis that spectrally selective gradient waveforms could be designed in order to detect multiple harmonics of the spin motion simultaneously⁴². Manduca *et al.*⁴³ exploited this property and found excellent agreements of shear stiffness results between mono and multifrequency acquisitions. Similarly to our study, Manduca *et al.* used two frequencies farther apart that may better determine dispersion effects. However, this method was restricted in terms of SNR due to a limitation of the gradient waveform amplitude. The major advantage of our proposed method is that SNR and PNR of dual-frequency acquisitions were not reduced and were even slightly higher than those found with conventional monofrequency MRE as shown in Figure 7.

Currently, the MRE technique commonly used to simultaneously encode several driving frequencies is the fractional encoding and is based on the sensitivity band of motion encoding gradients. It was introduced by Rump *et al.*³³ and it has been exploited in several

studies^{26,34,35,37}. These studies have applied the fractional encoding strategy *in vivo* to encode low frequency shear waves in large targeted organs (liver, brain). The key difference with our study rests on the fractional encoding ratios (between 0.3 and 1.25 *versus* 3 in our case) while reducing the examination time by using a non-classical sampling method.

In the present work, the harmonic wideband sensitivity of trapezoidal MEGs enables to encode high excitation frequencies (>500 Hz in preclinical experiments) with lower MEG frequencies. This is an interesting property since high oscillating frequencies are not obtainable by most gradient systems due to their low switching rate of the oscillating polarity, especially in clinical MRI scanners. Moreover, even in preclinical setups, fast oscillating gradients are more prone to inducing image imperfections due to hardware limitations. This could explain why in our results we found a lower SNR and PNR when encoding the 900 Hz wave with a 900 Hz MEG compared to the efficacy of the 300 Hz MEG (Figure 6). Likewise, this could have an effect in the viscoelastic parameters reconstruction and could explain the higher standard deviation obtained for this acquisition (Figure 7). One solution could be to reduce the amplitude of MEGs but this would lead to low cumulated phases encoding, which is a very frequent imposed compromise. The benefit of using high frequencies is that they enable a better power-law fit. Recent preclinical studies have investigated the effect of the considered range of frequencies on the power law fit. Riek *et al.*⁴⁴ found a good match between the fit with low and high frequency ranges on *ex vivo* experiments. However, Reiter *et al.*⁴⁵ found an increasing correlation of elasticity with the degree of liver fibrosis when using high mechanical excitation frequencies. In this context, our method could be useful to fill the whole spectrum with few acquisitions in order to reduce the total acquisition time.

In addition, by accessing the information of two different vibrations, we have no penalty in the acquisition time since the proposed sampling method allows to drastically reduce the number of phase offsets to be acquired. Such undersampling, based on a controlled frequency aliasing, allows to simultaneously encode several odd harmonics of the fundamental frequency. Generally, in order to simultaneously encode n different shear wave frequencies with odd harmonics relationship, the acquisition of $2n + 1$ phase offsets is required. This could be easily implemented to extend the dynamic of MRE to more than two frequencies in future experiments. Our proposed method presents a higher acceleration speed of multifrequency MRE acquisitions when compared with current

existent methods. For instance, Garteiser *et al.*³⁷ acquire 8 phase offsets to encode 3 different frequencies. Considering that usually 4 phase offset are acquired to encode a monofrequency shear wave, the gain ratio would be 33%. On the other hand, with our proposed method, it would be possible to encode 3 different frequencies with 7 phase offsets so the gain would be 42%. Therefore, the gain ratio of our proposed method is slightly better.

The challenging limitation of our study is the increased damping of shear waves at higher frequencies that results in limiting the penetration depth and the region of interest. In addition, we could not reach very high frequencies since the mechanical actuator in our experiment could not support frequencies higher than 1 kHz. Obtaining a plane wave pattern propagating through the sample was an important condition to respect in order to avoid possible artifacts in the viscoelastic parameter reconstruction.

The fast attenuation of high frequency shear waves through the tissues combined to the long TEs obtained with the low frequency MEGS, makes complicated the application of our method for the actual targeted organs with MRE (liver, brain) in a clinical context. Nevertheless, since our method enables to encode lower wavelengths (~ 5 mm), it might be useful in a clinical context for a novel MRE targeted organ: the characterization of thin superficial tissues such as the skin⁴⁶. Also our proposed method could be interesting for stiff tissues such as the cartilage^{47,48} that require much higher vibration frequencies than soft tissues.

The current study is the proof of concept of a novel multifrequency MRE method in a preclinical context. The experiments were carried on a homogenous phantom to realize a first validation of the proposed method. Since we can obtain very different wavelengths with our dual-frequency method, it might be useful for heterogeneous phantoms in order to provide a better visualization of interfaces between softer and stiffer parts of the phantom⁴⁹. However, the attenuation of the high frequency component would still be a limitation.

Further work is needed to explore the mentioned clinical and preclinical applications of the proposed method.

CONCLUSION

In this study, we present an alternative technique to the standard simultaneous multifrequency motion encoding: using harmonic widebands of motion encoding gradients, it is possible to significantly increase the frequency of shear waves used in MRE without increasing the burden on gradient systems. In addition, the application of a non-respecting Shannon sampling, which produces a controlled temporal aliasing, enables also a notable decrease in acquisition times of multifrequency MRE. Both combinations could ultimately be used to increase the speed of multifrequency MRE acquisitions and further work is needed to explore the clinical and preclinical applications of the proposed method.

ACKNOWLEDGEMENTS

This work was performed within the framework of the LABEX PRIMES of University of Lyon (ANR-11-LABX-0063), within the program "Investissements d'Avenir" (ANR-11-IDEX-0007) operated by the French National Research Agency (ANR). MR experiments were acquired on the PILoT platform, member of the France Life Imaging infrastructure (ANR-11-INBS-0006).

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TABLES

#Acquisition	Number of phase offsets	Acquisition time	f_{exc} (Hz)	A_{exc} (V)	f_{MEG} (Hz)	N_G	MEG shape
1	4	2min 8s	300	1	300	1	trapezoidal
2	4	2min 8s	900	1	900	3	trapezoidal
3	4	2min 8s	900	3	300	1	trapezoidal
4	4	2min 8s	900	3	300	1	sinusoidal
5	5	2min 40s	300 : 900	1 : 3	300	1	trapezoidal

Table 1. Experimental MRE parameters

FIGURES

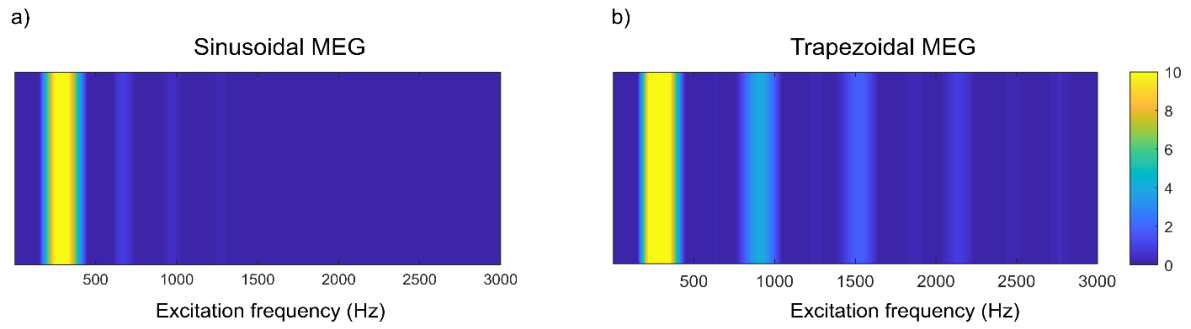


Figure 1. Accumulated phase with one cycle of 300 Hz MEG for various excitation frequencies and with different oscillating gradient shapes. A, Sinusoidal shaped; B, trapezoidal shaped. Values are in arbitrary units and were scaled from 0 to 10 for a better visualization. The periodical behavior of the sensitivity bands is visible when using a trapezoidal MEG.

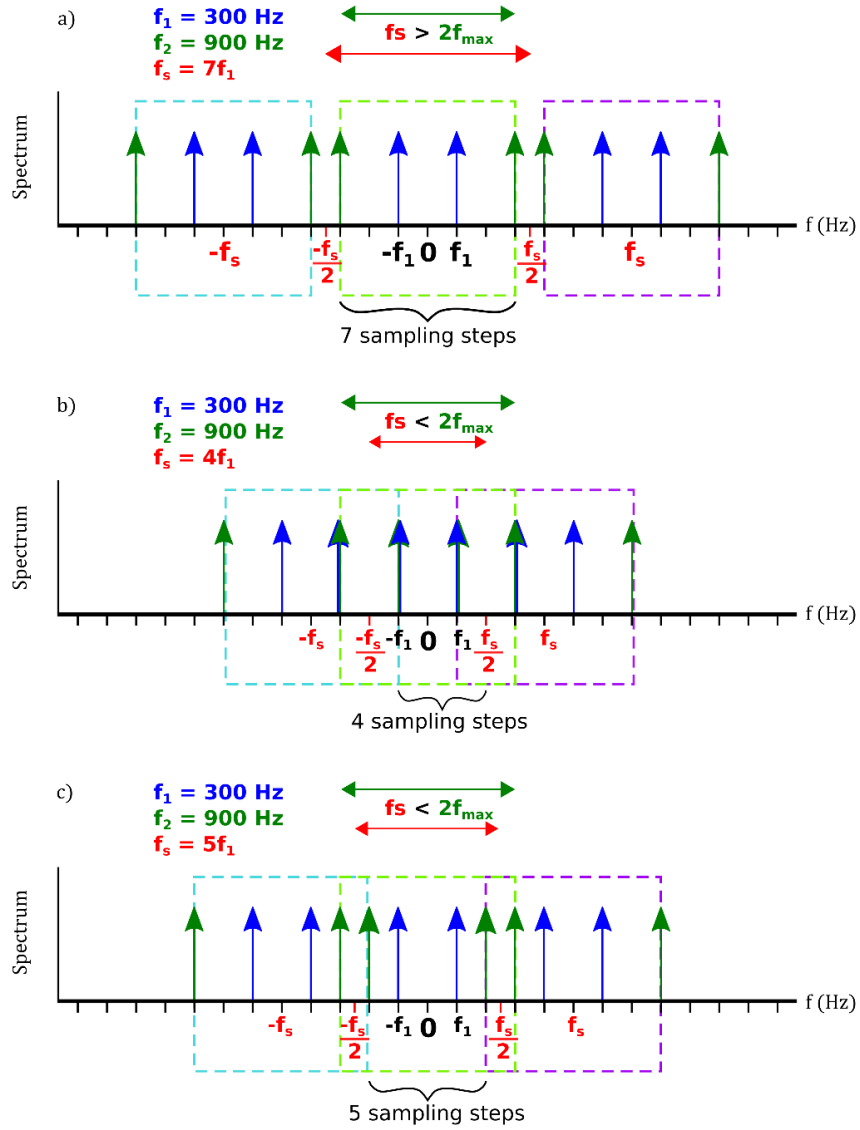


Figure 2. Different representations of the spectra obtained in the frequency domain after temporal Fourier transform of sampled points at a sampling frequency f_s . The original signal was a dual-frequency wave (f_1, f_2) equal to (300, 900) Hz. A, The spectra when $f_s = 7 f_1$, respecting the Shannon theorem. B, C, The spectra when a sampling not respecting the Shannon theorem is applied, producing a temporal aliasing. In B with $f_s = 4 f_1$ the temporal aliasing produces an overlap of the two frequency components at the same acquired points. In C with $f_s = 5 f_1$ the temporal aliasing is controlled in order to obtain both frequency components at different acquired points, enabling separation of the two components. Green and blue arrows represent the Dirac delta functions of the 300 and 900 Hz components, respectively.

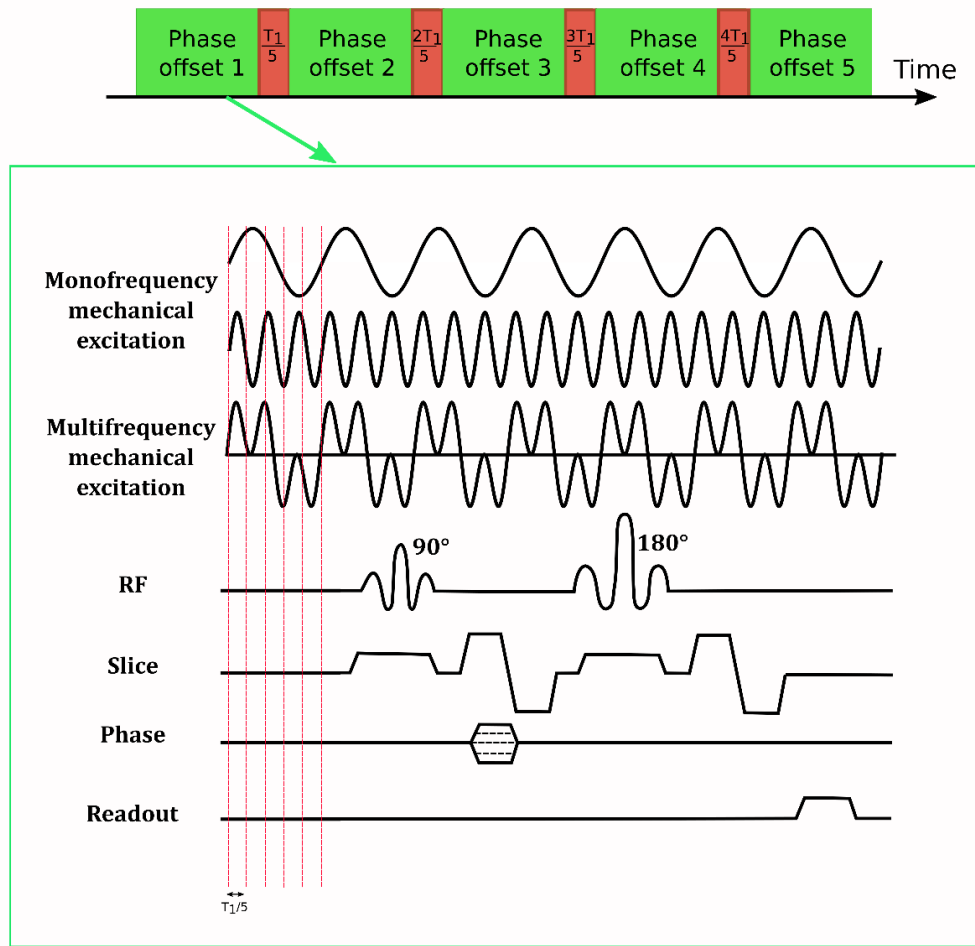
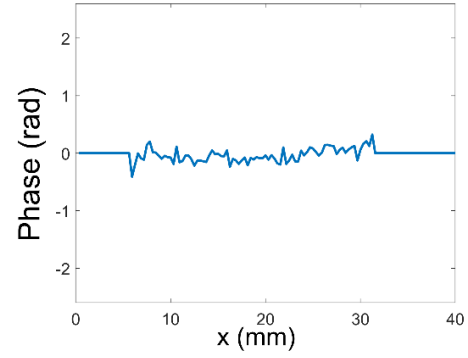
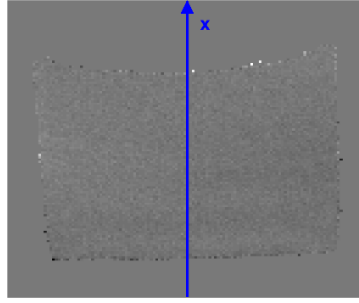


Figure 3. Chronogram of the multifrequency acquisition pulse sequence. The period of the lowest-frequency mechanical vibration is denoted as T_1 . After the complete image of a phase offset has been acquired (green square), a delay of $T_1/5$ (red square) is added to perform the acquisition for the next phase offset image.

Sinusoidal
300 Hz MEG



Trapezoidal
300 Hz MEG

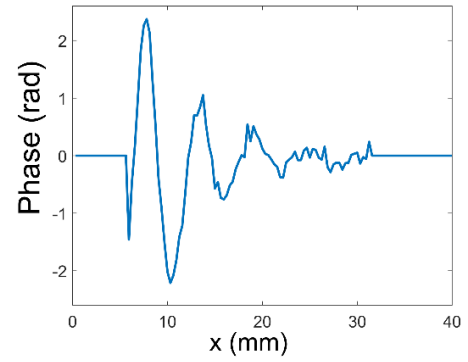
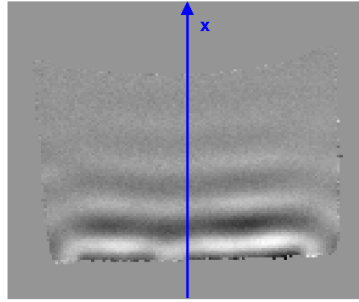


Figure 4. Wave propagation visualization on experimentally acquired phase images (left), and profiles following the blue arrow (right), when a 300 Hz MEG is applied (top, sinusoidal shaped; bottom, trapezoidal shaped) with a 900 Hz shear wave excitation, in radians. The motion is encoded when using a trapezoidal MEG.

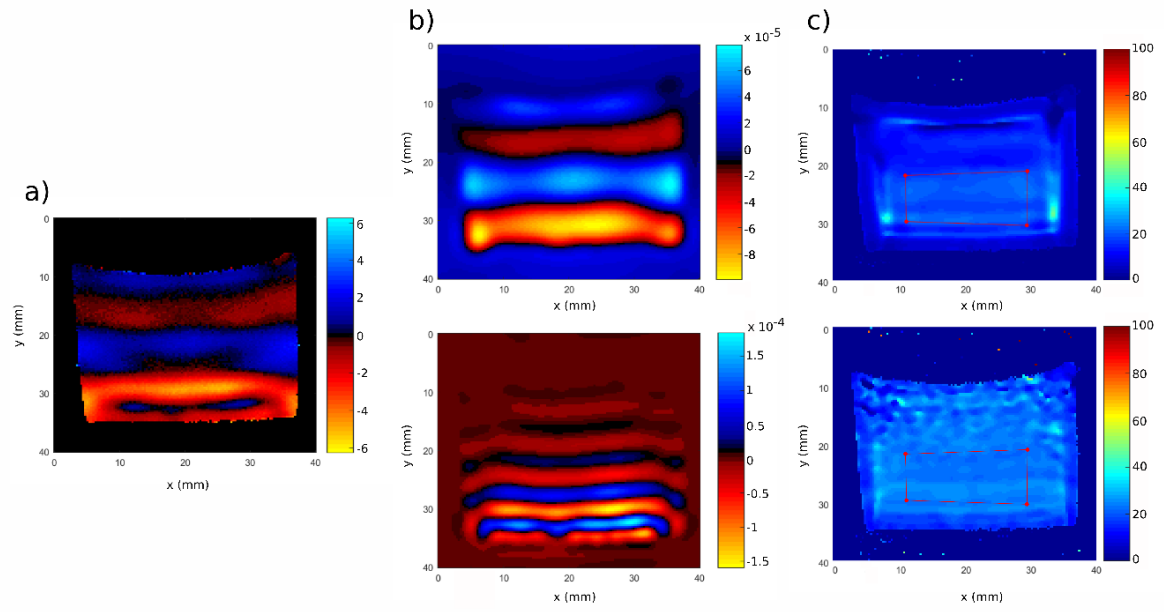


Figure 5. A, Phase image of the composite excitation (rad). B, Phase images (a.u.) and C, G' elastograms (kPa) of every extracted frequency component: top, 300 Hz; bottom, 900 Hz. The ROI used for the computation of the shear storage modulus mean and standard deviations is shown with red lines on the maps in C.

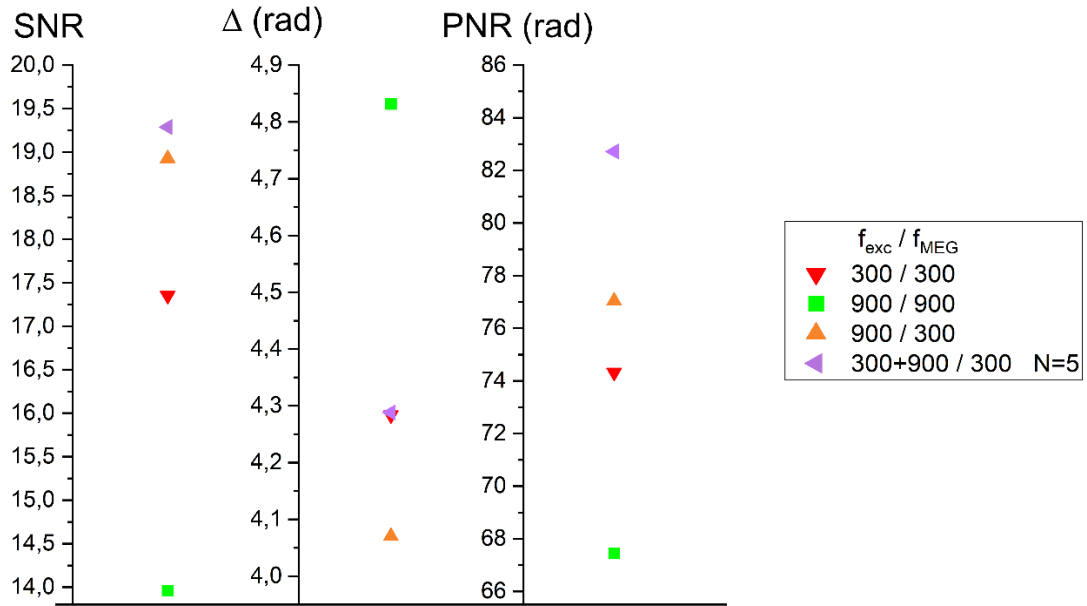


Figure 6. Plot of the SNR, Δ and PNR values obtained with the different acquisition modes. SNRs are in arbitrary units, Δ and PNRs in radians. Note that each parameter has its own vertical axis. Only the acquisitions with trapezoidal MEGs are considered here. In the legend: conventional monofrequency acquisition (\blacktriangledown , \blacksquare), harmonic monofrequency acquisition (\blacktriangle) and dual-frequency acquisition (\blacktriangleleft).

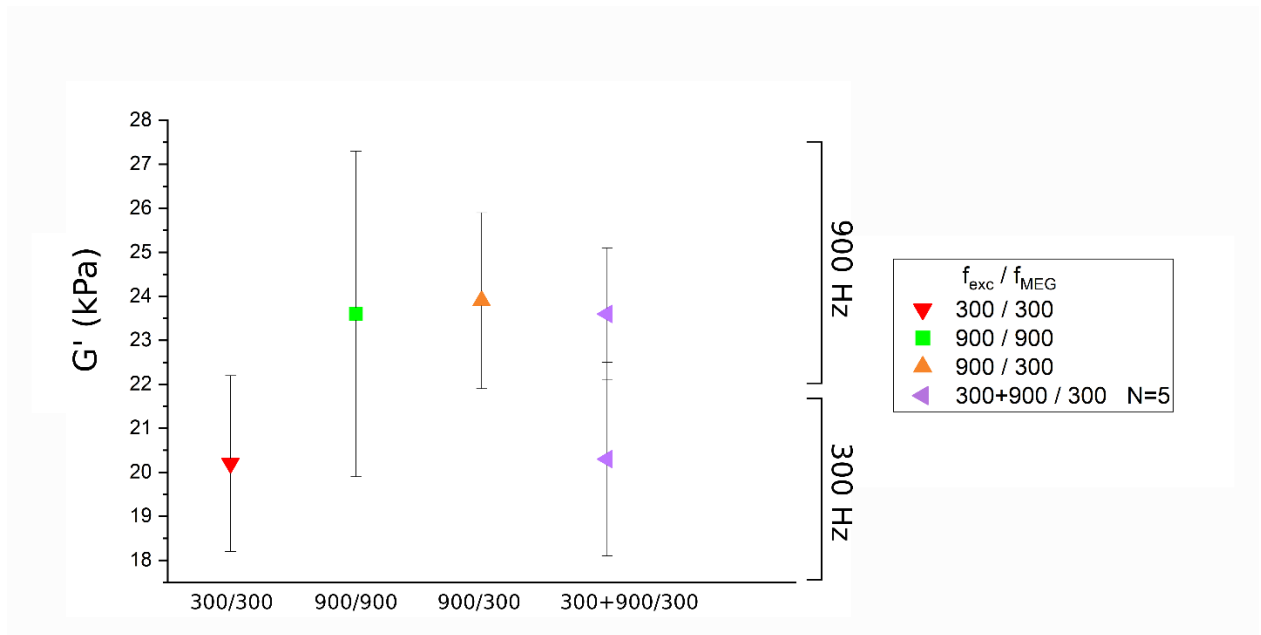


Figure 7. Mean values, with standard errors displayed as error bars, of the reconstructed shear storage modulus G' for both excitation frequencies (300 and 900 Hz) obtained with the different acquisition modes. Only the acquisitions with trapezoidal MEGs are considered here. An excellent agreement of the G' values is found for all cases.