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1 **Rheology of natural sediments and its influence on the settling of dropstones**
2 **in marine hemipelagic sediment**

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7

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9 **Key Points:**

- 10 • We obtain a new dataset of rheology measurements from naturally occurring hemipelagic
11 sediment and a terrestrial mud volcano
- 12 • Comparison with existing measurements indicates that particle composition influences
13 yield stress and consistency.
- 14 • Most pumice clasts recovered from IODP 340 are stratigraphically in place and thus date
15 the occurrence of eruptions.
16

17 **Abstract**

18 We investigate the rheology of naturally occurring hemipelagic marine sediment and compare
19 measurements to another naturally occurring sediment from a terrestrial mud volcano and
20 literature values. The hemipelagic marine sediment, collected by IODP 340, has a median grain
21 size of 5.5 microns, is poorly sorted, and contains 31% clay, including smectite. The yield
22 stresses and consistency are calculated by applying a range of shear stresses and shear-rates
23 using a cone-and-plate rheometer. A Herschel-Bulkley model is fit to measured shear stresses
24 and shear rates to calculate the yield stress and consistency. These measurements are performed
25 at a range of particle concentrations and show that the hemipelagic sediment has a yield stress at
26 particle concentrations as low as 10%. Increasing particle concentration increases the yield stress
27 and consistency. We apply our results to show that natural pumice clasts need to have a radius
28 greater than about 1 cm in order to settle through hemipelagic sediment on the sea floor. Most
29 recovered pumice clasts from IODP 340 are thus preserved in the same horizon in which they
30 were deposited.

31

32 **1 Introduction**

33 The rheological properties of sediment are influenced by multiple factors including
34 particle concentration, size, size distribution, composition and pore-water salinity (Imran et al.,
35 2001). Models for sediment rheology are often based on laboratory characterization of idealized
36 kaolinite and montmorillonite mixtures (e.g. Maar et al., 2001; Baas and Best, 2002; Talling,
37 2013). These idealized mixtures differ from naturally occurring submarine sediment in particle
38 size and distribution as well as composition. However, there are few studies of the rheology of
39 natural submarine deposits (e.g., Imran et al., 2001; Locat et al., 2004; Jeong, 2013; Menapace et
40 al., 2019).

41 The composition of the sediment can affect its rheology and cohesive sediments exhibit
42 complex behaviors that are often difficult to characterize (Yang et al., 2014). When materials
43 contain high clay fractions, strong attractive colloidal forces can lead to thixotropic behavior,
44 that is, the viscosity of the material depends on the shear history (Santolo et al., 2005).
45 Thixotropy is caused by a time dependent microstructural evolution of a material. When a shear
46 stress or strain is applied to the material the viscosity decreases (Barnes, 1997). In sediment, this
47 decrease in viscosity could be caused by the particles developing a preferred orientation or by the
48 breakdown of flocs and aggregates which release entrapped fluid and increases the lubrication
49 between the particles (Barnes, 1997; Tregger et al., 2010). However, this evolution is reversible:
50 the longer a material is at rest, the more the microstructure rebuilds until it reaches its initial state
51 again.

52 In addition to the clay fraction, the type of clay can also influence rheology due to
53 differences in cohesion. Smectite and montmorillonite have high cohesion compared to kaolinite
54 and illite. The difference in cohesion affects how easily the microstructure can be broken down.
55 Naturally occurring sediments differ from idealized mixtures because they contain a range of
56 sediment sizes, fossils, minerals and clays.

57 Thixotropy is also influenced by salinity and particle concentration. At high salinity,
58 greater than 0.35 g/l to 0.55 g/l, the inter-particle forces are strong enough to favor fast
59 reconstruction of the microstructure, thus making the material strongly thixotropic (Perret et al.,
60 1996). Particle concentration also affects the thixotropic behavior of a material. At low particle
61 concentrations, there are fewer interactions between the particles and thixotropic effects are less
62 significant (Santolo et al., 2012).

63 The rheological behavior of sediment has implications for the initiation and flow of
64 submarine sediment-laden density currents. Additionally, seafloor sediment preserves a record of
65 volcanic eruptions in tephra and crypto-tephra layers including isolated cm-sized pumice clasts
66 that record the occurrence of additional volcanic eruptions (le Friant et al., 2013, 2015; Palmer et
67 al., 2016; Jutzeler et al., 2016). These eruptions can be dated using the surrounding sediments but
68 this assumes that the clasts are at the same stratigraphic location and did not settle through
69 underlying sediment upon deposition.

70 In this study we analyze naturally occurring hemipelagic sediment collected from the
71 seafloor adjacent to Lesser Antilles Volcanic Arc. To determine yield stress and consistency as a
72 function of particle concentration a suite of rheometric tests are utilized. The yield stress and
73 consistency of a material depend strongly on the particle concentration. As particle concentration
74 increases, there is an increase in particle-particle interactions, which causes an increase in the
75 friction and thus increases the yield stress and consistency (Coussot and Piau, 1994). We thus
76 use a range of particle concentrations to develop a better understanding of how these natural
77 sediments flow. The same tests are performed on an additional naturally occurring sediment from
78 the Davis-Schrimp mud volcanoes for further comparison. Additionally, our results are
79 compared to previously published measurements on other materials in order to better identify the
80 factors that influence rheology. For illustrative purposes, we use our results to determine the size
81 of particles that would settle through the hemipelagic sediment on the seafloor and hence what
82 size particles would be preserved in situ within marine sediments. We use these results to assess
83 whether volcanic dropstones collected by IODP 340 are preserved in situ.

84

85 **2 Methods**

86 The hemipelagic sediment was collected offshore of the Lesser Antilles volcanic arc from
87 a water depth of 2745 meters. The Lesser Antilles volcanic arc is located at the eastern border of
88 the Caribbean plate. The proximity of our sampling location to volcanic islands, specifically the
89 active volcano Montagne Pelée, causes our sample to contain volcanoclastic sediment as well as
90 marine fossils (le Friant et al., 2008). The sample, 340 U1400C 3H 4, was collected during IODP
91 340 from 26 meters below the sea floor from a sedimentary unit that includes multiple landslide
92 deposits (le Friant et al., 2013; Lafuerza et al., 2015; Brunet et al., 2016); see supplementary
93 Figure S1a for a location map. For comparison, another natural mud from the Davis-Schrimpf
94 mud volcanoes located next to the Salton Sea in the Imperial Valley, California is analyzed (map
95 in supplementary Figure S1b). The Salton Sea mud volcanoes are driven by the ascent of CO₂
96 from underlying geothermal systems and the mud recirculates within the subsurface plumbing
97 system (Svenson et al., 2009).

98

99 **2.1 Sample Properties**

100 Size distribution is measured using a Sequoia Portable XR Laser Diffraction Particle Size
101 Analyzer. The cumulative size distribution for both samples is shown in Figure 1a. The median
102 grain size for the hemipelagic sediment is 5.5 microns. Grain density was measured as 2700
103 kg/m³ on the IODP ship using helium pycnometry (le Friant et al., 2013).

104 Particle concentration, the amount of solids in the total mixture, is given by

$$105 \quad \phi = \frac{V_s}{V_t} \quad (1)$$

106 where V_s is the volume of solids, and V_t is the total volume. We measure ϕ by weighing a known
107 volume of material. Uncertainty in ϕ is dominated by uncertainty in measured volume. The
108 density of the hemipelagic sediment was confirmed to be 2700 kg/m³.

109 Figure 1b,c shows three dimensional x-ray microtomography reconstructions of the
110 samples imaged at the Advanced Light Source 8.3.2 beamline, Lawrence Berkeley National
111 Laboratory, with a resolution of 1.3 microns/pixel. The hemipelagic marine sediment is poorly
112 sorted and contains large voids (Figure 1b). Additionally, the sediment also contains foraminifera
113 fossils and crystal fragments, such as hornblende, quartz and muscovite, that are large compared
114 to particles in the surrounding matrix. The Salton Sea sediment has a more uniform size
115 distribution by comparison (Figure 1c).

116 Using x-ray diffraction, we determine that the hemipelagic sediment is composed of 20%
117 feldspar and plagioclase, 19% calcite, 16% quartz, 7% aragonite, and 31% clay. The Salton Sea
118 sediment is composed of 41% quartz, 3% dolomite, 16% plagioclase, 8% orthoclase and 26%
119 clay. The clay in both samples contains kaolinite, illite, and montmorillonite while the
120 hemipelagic sediment also contains smectite.

121

122 **3 Rheology Measurement Methods**

123 We measure the rheology using a HAAKE Rheoscope 1. To ensure uniform strain-rate
124 through the sample a cone and plate geometry, with a 4° cone angle and 60 mm diameter plate, is
125 used. We apply either a controlled angular velocity, and hence a shear rate, and then measure the
126 torque required to maintain this rate, or alternatively apply a torque and measure the resulting
127 deformation.

128 In order to prevent slippage, both the cone and the plate are covered in ANSI 150 grit
129 sand paper (average particle size of 100 microns). The gap size between the plate and the center
130 of the cone was 0.142 mm after adhesion of sand paper. Each test utilizes 8ml of sample. To
131 prevent jamming of particles within the gap, the mud is sieved to remove particles larger than

132 100 microns, which comprise less than a few percent of the volume of the sediment. The
133 volumetrically dominant fine particles have a larger controlling influence on the yield stress, so
134 we do not expect that the small number of fossils removed from the sample would alter the
135 rheology (Yu et al., 2013). In order to prevent temperature fluctuations from affecting the
136 measurements, the temperature of the mud is kept constant at 20°C using a Thermo Scientific
137 Haake DC30-K20 Digital Bath. At this temperature water viscosity is 1.05 mPa s compared to
138 1.57 mPa s at a seafloor temperature of 4 °C.

139 The hemipelagic sediment is from a marine setting so we control salinity of the water by
140 adding artificial seawater since the salinity of the water affects suspension rheology (Jeong,
141 2010). 35 grams of Instant Ocean salt mixture is added to 1 liter of water to create salinity
142 similar to that of the Caribbean Sea.

143 The thixotropic nature of the sediment led us to focus on the static yield stress, above
144 which the material will start to flow, and the dynamic yield stress, below which the material will
145 no longer flow. Two separate tests are performed: the first to quantify the static yield and the
146 second to determine the dynamic yield and consistency.

147

148 **3.1 Flow Characteristics**

149 In order to characterize rheology we consider an idealized constitutive model. The most
150 frequently used model to describe naturally occurring sediments and muds is the Herschel-
151 Bulkley model (Huang and Garcia, 1998):

$$152 \tau = \tau_y + K\dot{\gamma}^n \quad (2)$$

153 where τ is shear stress, τ_y is the yield stress, $\dot{\gamma}$ is strain rate, K is consistency, and n is the flow
154 index. We adopt the Herschel-Bulkley model, but note that models without a yield stress have

155 also been used to model mud rheology especially for low strain rates and high viscosities (e.g.,
156 Menapace et al., 2019).

157 During the measurements, it became apparent that both materials did not simply follow
158 the Herschel-Bulkley model (Supplementary Figure S2) as has been documented previously
159 (e.g., Yang et al., 2014). Different stresses are measured when applying increasing and
160 decreasing strain rates, indicating that the properties are history dependent, one signature of
161 thixotropic rheology (Moller et al., 2009) (Supplementary Figure 2b). This leads to the static
162 yield stress, which is associated with the start of flow, being higher than the dynamic yield stress,
163 below which the material will stop flowing, due to the static yield stress having to breakdown the
164 microstructure that has formed while the material was at rest. The microstructure breakdown is
165 reversible; if the material is left at rest, the yield stress will increase again (Moller et al., 2009).
166 There are models that quantify the evolving microstructure (e.g., Dullaert and Mewis, 2006; de
167 Souza Mendes, 2009), however for this study we focus on the static yield stress and the dynamic
168 yield stress. Consistency is analogous to viscosity, and is also quantified prior to the cessation of
169 flow.

170

171 **3.2 Static Yield**

172 Thixotropic fluids have a viscosity that decreases over time as they are sheared. To
173 mitigate the variations in shear history, samples are allowed to rest for ten minutes between each
174 test (Van Kessel and Blom, 1998). After running hysteresis tests at varying time intervals
175 between tests, the material recovers most of its strength after 10 minutes. Ten minutes limits the
176 amount of evaporation of water and settling of particles so as to not affect measurements. Static
177 yield is measured using a controlled stress mode while the sample was maintained at a constant

178 temperature. For this controlled stress mode, an increasing shear stress, τ , or torque, is applied
179 and the resulting deformation is measured (Supplementary Figure S2a). Shear stress is increased
180 linearly in 70 steps from 0 Pa to 1500 Pa, for 10 seconds each step. Initially, strain increases
181 linearly with stress, however with further increases of stress the strain grows exponentially fast
182 and is immeasurable after a certain shear stress. The location of this exponential increase varied
183 with each test depending on the water content. To calculate static yield stress a Herschel-Bulkley
184 model is fit to the data from increasing portion of the test, providing τ_y , K , and n .

185

186 **3.3 Dynamic Yield**

187 In order to determine the dynamic yield stress and the consistency, a shear strain rate, $\dot{\gamma}$,
188 is applied and the torque, or shear stress, τ , required to maintain this strain rate is measured.
189 Strain rate is increased in 40 steps uniformly spaced in the log of strain rate from 0.01 s^{-1} to 1.5 s^{-1} ,
190 and subsequently decreased in the same 40 steps, with each step lasting 15 seconds
191 (Supplementary Figure S3 shows results for two different particle concentrations). To determine
192 the dynamic yield stress and consistency, the decreasing portion of the test is fit with a Bingham
193 model.

194

195 **4 Results**

196 Figure 2a shows the relationship between particle concentration (%) and yield stress (Pa)
197 for our samples compared to other studies. For the hemipelagic sediment and Salton Sea mud,
198 the filled circles show static yield stress and the open circles show dynamic yield stress. Table 1
199 summarizes the results for our study and Supplementary Table 1 summarizes comparative
200 studies. In our samples the dynamic yield stress is always lower than the static yield stress, with

201 the exception of the Salton Sea mud at a concentration of 40.5%. This exception could be the
 202 result of sedimentation within the sample that allowed the material to start flowing at a lower
 203 static yield stress. In addition to measuring the static yield using a controlled stress mode, we can
 204 estimate the static yield from the inflection point of the increasing strain rate portion of the
 205 dynamic yield tests (Santolo et al., 2012), and the two separately determined yield stresses are
 206 similar.

207 A yield stress becomes apparent in the hemipelagic sediment at a particle concentration
 208 of ~10% and increases with increasing particle concentration. Figure 2a also shows the yield
 209 stress as a function of particle concentration measured in other studies (Coussot and Piau, 1994;
 210 Coussot et al., 1996; Huang and Garcia, 1998; Matel et al., 2005; Remaitre et al., 2005; Maciel et
 211 al., 2009; Santolo et al., 2010; Blasio et al., 2010; Manga and Bonini, 2012; Jeong et al., 2013;
 212 Yang et al., 2014; Tran et al., 2014). In Figure 2, non-natural sediments are represented with
 213 squares.

214 To determine the yield stress as a function of particle concentration, the hemipelagic data
 215 is fit with the model developed by Mueller et al. (2009):

$$216 \quad \tau_y = \tau^* \left(\left(1 - \frac{\phi}{\phi_m} \right)^{-2} - 1 \right) \quad (3)$$

217 where τ^* is a fitting parameter, ϕ is the particle concentration, ϕ_m is the maximum particle
 218 concentration. The fit of the Mueller et al. (2009) model to the hemipelagic static yield data is
 219 shown in Supplementary Figure 4. We find $\tau^* = 34.9 \pm 6.9$ Pa and $\phi_m = 0.23 \pm 0.004$.

220 Consistency also increases with increasing particle concentration. Following Mader et al.
 221 (2013), the hemipelagic data is fit to the model:

$$222 \quad K_r = \left(1 - \frac{\phi}{\phi_m} \right)^{-2} \quad (4)$$

223 where $K_r = K/\mu_0$ where μ_0 is viscosity of the suspending fluid (Pa•s) (Supplementary Figure 5).
224 We find $\phi_m = 0.25 \pm 0.02$ Pa and $\mu_0 = 0.14 \pm 0.05$ Pa s.

225

226 **5 Discussion**

227 Both the static and dynamic yield tests indicate that the hemipelagic sediment is
228 thixotropic. In the dynamic tests, there is a significant difference between the increasing strain
229 rate section and the decreasing strain rate section (Supplement 2b). This indicates that the
230 material has a history dependence in its rheology and hence its microstructure. The higher the
231 particle concentration, the more dramatic this hysteresis. In the static yield tests, the material
232 initially behaves like a linear elastic material until a specific stress is overcome and then the
233 strain exponentially increases until the material reaches another critical point and flows too fast
234 for the rheometer to measure (Supplement 2a). Coussot et al. (1992) similarly found that the
235 response of their clay-rich material is essentially elastic until a critical point is reached and the
236 flow becomes unstable and fast. This behavior indicates that the material is initially structured,
237 but when the apparent yield stress is overcome the sediment becomes unstructured and starts to
238 flow at increasing velocities. Alignment of particles with the flow causes the material to flow
239 faster when a shear is applied (Barnes, 1997). Additionally, the static yield stress is significantly
240 larger than the dynamic yield stress, which is consistent with thixotropic behavior (Figure 2).

241 The scatter within our data could be caused by the history dependence of our samples due
242 to slightly different handling of the material prior to testing. We tried to mitigate this effect by
243 allowing the samples to rest between tests and by loading the samples into the rheometer in the
244 same way. The scatter could additionally be caused by heterogeneities within the natural sample
245 as each data point was obtained from a different aliquot of the mud.

246 The range of particle concentrations over which a yield stress can be measured for the
247 hemipelagic sediment is approximately between 10-20%. Below a concentration of 10% the
248 particles settle creating a muddy fluid top layer and particle rich layer underneath. Comparing
249 the results with other studies that used sediment suspensions, the hemipelagic marine sediment is
250 able to develop a static stress at comparatively low particle concentrations (Figure 2a).
251 Additionally, other studies that measured yield stresses at low particle concentrations were
252 kaolinite and water mixtures produced in the lab (Figure 2a – laboratory produced mixtures are
253 represented with squares).

254 The clay content in the hemipelagic sediment is in the higher range compared to the other
255 studies shown in Figure 2a, although not all studies record clay content. The hemipelagic
256 sediment has approximately 31% clay, which could be one of the contributing factors to why
257 there is a yield stresses at such low particle concentrations. Huang and García (1998) measure a
258 yield stress at lower concentration, approximately 7%, but do not record the clay content of the
259 sample they used. The higher clay content increases electrostatic forces between the clay
260 particles. Yu et al. (2013) found that increasing the clay content from approximately 10% to 40%
261 clay increased the yield stress by more than an order of magnitude, consistent with Jeong (2013)
262 who reported the highest clay content and similarly large yield stresses. While both the
263 hemipelagic sediment and Salton Sea mud have kaolinite, illite, montmorillonite, the
264 hemipelagic sediment also contains smectite, which is highly cohesive. The addition of this clay
265 type could also contribute to the higher yield stress and more thixotropic behavior compared to
266 the Salton Sea mud. The yield stress of the hemipelagic sediment is also more than an order of
267 magnitude larger than kaolinite suspensions with similar particle concentrations suggesting that

268 the presence of smectite is largely responsible for the higher yield stress (in Figure 2, compare
269 Huang and Garcia (1998) and Maciel et al. (2009) data with the present measurements).

270 The salinity of the hemipelagic marine sediment, at 35 g/l, is well above the salinity that
271 Perret et al. (1996) identify as influencing rheology. Above this salinity, there is fast
272 reconstruction of the microstructure, and below the material is likely to favor slower
273 reconstruction. Yang et al. (2014), however, found no effect on rheology of salinity in the range
274 of fresh water to 50 g/L salinity. Using the method outlined by Perret et al. (1996) to determine
275 the thixotropic nature of a material, the area between the hemipelagic sediment thixotropy test
276 curves is larger than that of the Salton Sea sediment, indicating that the hemipelagic material is
277 slower to recover back to its initial state. This indicates that the clay content has a larger effect
278 than salinity for the two samples we considered.

279 Materials with high clay contents have a slow rate of reconstruction, increasing their
280 recovery time (Barnes 1997). This is also the case for the two samples we considered. While the
281 Salton Sea mud has a significant portion of clay, the median particle size of the hemipelagic
282 sediment is much smaller at 5.5 microns, compared to the Salton Sea sediment 35.7 microns,
283 which could account for the difference in recovery times.

284

285 **5.1 Particle Settling through seafloor sediment**

286 Hemipelagic sediment often contains isolated “dropstones” delivered by melting icebergs
287 or volcanic eruptions; in the case of sediments offshore the Lesser Antilles volcanic arc, these are
288 predominately pumice clasts, deposited from rafts of pumice or fall out. Whether these clasts are
289 preserved within their deposit or subsequently sink multiple clast diameters through underlying
290 layers depends on their size and weight. In order to determine whether a larger particle with

291 radius R_0 can settle within mud, we approximate the clast as a completely submerged sphere.

292 The force is

$$293 \quad F = \frac{4}{3}\pi(\rho_s - \rho)R_0^3g \quad (5)$$

294 where ρ_s is the density of the sphere (kg/m^3), ρ is the density of the mud, R_0 is the radius of the
 295 sphere (m), and g is gravity (m/s^2). The density of the sediment depends on the particle
 296 concentration (C_v), $\rho=1800\cdot C_v+1000 \text{ kg/m}^3$. The sphere will cause yielding of the mud and sink
 297 when the yield stress parameter

$$298 \quad Y_g = \frac{2\tau_y\pi R_0^2}{F} \quad (6)$$

299 exceeds a critical value of 0.143 (Beris et al., 1985). Pumice is porous, between 60-80%
 300 porosity, and at the bottom of the seafloor these pores will be filled with water. The mean
 301 density was calculated assuming 60% porosity and using a glass density of 2700 kg/m^3 . We use a
 302 dynamic yield stress because we assume the clasts are initially moving when they reach the
 303 seafloor. Using equation (6) and equation (3) to determine the yield stress as a function of
 304 particle concentration, we calculate the relationship between the particle concentration and the
 305 radius of a clasts that would settle through the sediment (Figure 3). Below this radius, clasts
 306 would not settle but instead be preserved in situ. Using an average mud concentration at the
 307 seafloor of 10% (Hamilton, 1976) a pumice clast would have to have a radius larger than about
 308 1 centimeter to settle through uncompacted hemipelagic sediment. Most of the various volcanic
 309 dropstones collected during IODP 340 (le Friant et al., 2013; Jutzeler et al., 2016) are thus in
 310 place and their stratigraphic locations record the time of eruptions.

311

312 **6 Conclusion**

313 By studying the rheology of a naturally occurring mud we can better discern some of the
314 influencing factors in rheology and thixotropic behavior. Particle concentration affects the yield
315 stress and consistency, with increasing particle concentrations leading to increasing yield stresses
316 and consistency. As particle concentration increases, the thixotropic behavior also increases and
317 the material becomes more history dependent. The particle concentration of the material also
318 dictates the sizes of clasts that can settle through seafloor sediment. Pumice clasts would need
319 radii greater than ~1 cm to settle through the hemipelagic sediment, indicating that most pumice
320 clasts collected in IODP 340 are stratigraphically in place. Finally, the high smectite clay content
321 in the hemipelagic sediment leads to more attractive colloidal forces in the sample and appears to
322 influence the yield stress and enhance the thixotropic nature of the material.

323

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337

338 **8 References**

339 Baas, J. H., Best, J. L. 2002, Turbulence modulation in clay-rich sediment-laden flows and some
340 implications for sediment deposition. *Journal of Sedimentary Research*, Vol 72, No. 3.

341

342 Barnes, H. A. 1997, Thixotropy — a review. *Journal of Non-Newtonian Fluid Mechanics*, Vol.
343 70, pp. 1-33.

344 Beris, A. N., Tsamopoulos, J. A., Armstrong, R. C., Brown, R. A. 1985, Creeping motion of a
345 sphere through a Bingham plastic. *Journal of Fluid Mechanics*, Vol. 158, pp. 219-244.

346

347 Blasio, F. V., Breien, H., Elverhøi, A. 2011, Modelling a cohesive-frictional debris flow: an
348 experimental, theoretical and field-based study. *Earth Surface Processes and Landforms*,
349 Vol. 36, pp. 753-766.

350

351 Brunet, M., A. Le Friant, G. Boudon, S. Lafuerza, P. Talling, M. Hornbach, O. Ishizuka, E.
352 Lebas, H. Guyard, and IODP Expedition 340 Science Party (2016), Composition,
353 geometry, and emplacement dynamics of a large volcanic island landslide offshore
354 Martinique: From volcano flank- collapse to seafloor sediment failure?, *Geochem.*
355 *Geophys. Geosyst.*, 17-3, 699– 724, doi:10.1002/2015GC006034.

356

357 Coussot, P., Leonov, A. I., Piau, J. M. 1992, Rheologic modeling and peculiar properties of
358 some debris flows, *Proceedings of International Symposium on Erosion, Debris Flows and*
359 *Environment in Mountain Regions*, I.A.H.S. Publication 209, pp. 207-216.

360

361 Coussot, P., Piau, J. M. 1994, On the behavior of fine mud suspensions. *Rheologica Acta*, Vol.
362 33, pp. 175-184.

363

364 Coussot, P., Proust, S., Ancey, C. 1996, Rheological interpretation of deposits of yield stress
365 Fluids. *Journal of Non-Newtonian Fluid Mechanics*, Vol. 66, pp. 55-70.

366

367 de Souza Mendes, P. R. 2009, Modeling the thixotropic behavior of structured fluids. *Journal of*
368 *Non-Newtonian Fluid Mechanics*, Vol. 164, pp. 66-75.

369

370 Dullaert, K., Mewis, J. 2006, A structural kinetics model for thixotropy. *Journal of Non-*
371 *Newtonian Fluid Mechanics*, Vol. 139, pp. 21-30.

372

373 Hamilton, E. L. 1976, Variations of density and porosity with depth in deep-sea sediments.
374 *Journal of Sedimentary Petrology*, Vol. 46, pp. 280-300.

375

376 Huang, X., García, M. H. 1998, A Herschel-Bulkley model for mud flow down a slope. *Journal*
377 *of Fluid Mechanics*, Vol. 374, pp. 305-333.

378

- 379 Imran, J., Parker, G., Locat, J., Lee, H. 2001, 1D Numerical model of muddy subaqueous and
 380 subaerial debris flows. *Journal of Hydraulic Engineering*, Vol. 127, pp. 959-968.
 381
- 382 Jeong, S. W. 2010, Grain size dependent rheology on the mobility of debris flows. *Geoscience*
 383 *Journal*, Vol. 14, pp. 359-369.
 384
- 385 Jeong, S. W. 2013, Determining the viscosity and yield surface of marine sediments using
 386 modified Bingham models. *Geosciences Journal*, Vol. 17, pp. 241-247.
 387
- 388 Jutzeler, M., M. Manga, J.D.L. White, P.J. Talling, A.A. Proussevitch, S.F.L. Watt, M. Cassidy,
 389 R.N. Taylor, A. Le Friant, and O. Ishizuka (2016) Submarine deposits from pumice-rich
 390 pyroclastic density currents dispersing over water: an outstanding example from offshore
 391 Montserrat (IODP 340), *Bulletin of the Geological Society of America*, doi:
 392 dx.doi.org/20.1130/B31448.1
- 393 Lafuerza S., A. Le Friant, M. Manga, G. Boudon, B. Villemant, N. Stroncik, B. Voight, M.
 394 Hornbach, O. Ishizuka and the Expedition 340 Scientific Party. 2014, Geomechanical
 395 characterizations of submarine volcano flank sediments, Martinique, Lesse Antilles Arc.
 396 S. Krastel et al., (Eds.), Sumarine mass movements and consequences, *Advances in*
 397 *Natural and Technological Hazards Research*, Springer Inten. Publishing, Switzerland,
 398 2014, 37, 73-81, doi 10.1007/978-3-319-00972-8_7.
 399
- 400 Le Friant, A., Lock, E.J., Hart, M.B., Boudon, G., Sparks, R.S.J., Leng, M.J., Smart, C.W.,
 401 Komorowski, J.C., Deplus, C., Fisher, J.K. 2008, Late Pleistocene tephrochronology of
 402 marine sediments adjacent to Montserrat, Lesser Antilles volcanic arc. *Journal of rthe*
 403 *Geological Society*, London, (London, U. K.), Vol. 165, pp. 279–289.
 404
- 405 Le Friant, A., Ishizuka, O., Stroncik, N. A., and the Expedition 340 Scientists 2013, Lesser
 406 Antilles Volcanism and Landslides, Integrated Ocean Drilling Program Expedition 340
 407 Preliminary Report, doi:10.2204.iodp.pr.340.2012.
 408
- 409 Le Friant, A., Ishizuka, O. , Boudon, G. , Palmer, M.R. , Talling, P. , Villemant, B. , Adachi, T.,
 410 Aljhdali, M., Breikreuz, C., Brunet, M., Caron, B., Coussens, M., Deplus, C., Endo, D.,
 411 Feuillet, N., Fraas, A.J., Fujinawa, A., Hart, M. B., Hatfield, R.G., Hornbach, M.,
 412 Jutzeler, M., Kataoka, K. S., Komorowski, J-C., Lebas, E., Lafuerza, S., Maeno, F.,
 413 Manga, M., Martínez-Colón, M., McCanta, M., Morgan, S., Saito, T., Slagle, A., Sparks,
 414 S., Stinton, A., Stroncik, N., Subramanyam, K. S.V., Tamura, Y., Trofimovs, J., Voight,
 415 B., Wall-Palmer, D., Wang, F., Watt, S.F.L. 2015, Submarine record of volcanic island
 416 construction and collapse in the Lesser Antilles arc: First scientific drilling of submarine
 417 volcanic island landslides by IODP Expedition 340. *Geochem., Geophys., Geosyst*, 16, 2,
 418 420-442.
 419
- 420 Locat, J, Lee, H.J., Locat, P., Imran, J. 2004, Numerical analysis of the mobility of the Palos
 421 Verdes debris avalanche, California, and its implication for the generation of tsunamis.
 422 *Marine Geology*, Vol. 203, pp. 269-280.
 423

- 424 Maciel, G. F., Santos, H. K., Ferreira, F. O. 2009, Rheological Analysis of water clay
425 compositions in order to investigate mudflows developing in canals. *Journal of the Brazil*
426 *Society of Mechanical Science and Engineer*, 31, 64-74.
427
- 428 Mader, H.M., Llewelin, E. W., Mueller, S. P. 2013, The rheology of two-phase magmas: A
429 review and analysis. *Journal of Volcanology and Geothermal research*, 257, 135-158.
430
- 431 Malet, J.P., Laigle, D., Remaitre, A., Maquaire, O. 2005, Triggering condition and mobility of
432 debris flows associated to complex earthflow. *Geomorphology*, 66, 215-235.
433
- 434 Manga, M., Bonini, M. 2012, Large historical eruptions at subaerial mud volcanoes, Italy.
435 *Natural Hazards and Earth System Science*, Vol. 12, pp. 3377-3386.
436
- 437 Menapace, W., Tangunan, D., Maas, M., Williams, T., Kopf, A. 2019, Rheology and
438 biostratigraphy of the Maraian serpentinite muds unravel mud volcano evolution. *Journal*
439 *of Geophysical Research*, Vol. 124, doi.org/10.1029/2019JB018265.
440
- 441 Marr, J. G., Harff, P.A., Shanmugam, G., Parker, G. 2001. Experiments on subaqueous sandy
442 gravity flows: The role of clay and water content in flow dynamics and depositional
443 structures. *GSA Bulletin*, Vol. 113, No. 11, 1377-1386.
444
- 445 Moller, P., Fall, A., Chikkadi, V., Derks, D., Bonn, D. 2009, An attempt to categorize yield
446 stress fluid behavior. *Philosophical Transactions of the Royal Society*, Vol. 367, pp.
447 5139-5155.
448
- 449 Mueller, S., Llewelin, E. W., Mader, H. M. 2009, The rheology of suspensions of solid
450 particles. *Proceedings of the Royal Society*, Vol. 446, pp. 1201-1228.
451
- 452 Palmer, M., Hatter S.J., Gernon T.M., Taylor, R.N., Cassidy M., Johnson P., Le Friant A.,
453 Ishizuka, O., 2016, Discovery of a large 2.4 Ma Plinian eruption of Basse-Terre,
454 Guadeloupe, from the marine sediment record. *Geology*, 44-2, 123-126,
455 doi:10.1130/G37193.1
456
- 457 Perret, D., Locat, J., Martignoni, P. 1996, Thixotropic behavior during shear of a fine-grained
458 mud from Eastern Canada. *Engineering Geology*, Vol. 43, pp. 31-44.
459
- 460 Remaitre, Al., Matel, J.P., Maquaire, O., Ancey, C., Locat, J. 2005, Flow behavior and runout
461 modeling of a complex debris flow in a clay-shale basin. *Earth Surface Processes and*
462 *Landforms*, 30, 479-488.
463
- 464 Santolo, A. S., Evangelista, A. 2005, Some observations on the prediction of the dynamic
465 parameters of debris flows in pyroclastic deposits in the Campania region of Italy.
466 *Natural Hazards*, Vol. 50, pp. 605-622.
467
- 468 Santolo, A. S., Pellegrino, A. M., Evangelista, A. 2010, Experimental study on the rheologic
469 behavior of debris flow. *Natural Hazards Earth System Science*, Vol. 10, pp. 2507-2514.

- 470
471 Santolo, A. S., Pellegrino, A. M., Evangelista, A., Coussot, P. 2012, Rheologic behavior of
472 reconstituted pyroclastic debris flow. *Geotechnique*, Vol. 62, pp. 19-27.
473
- 474 Svensen, H., Hammer, O., Mazzini, A., Onderdonk, N., Polteau S., Planke, S., Podladchikov,
475 Y.Y. 2009, Dynamics of hydrothermal seeps from the Salton Sea geothermal system
476 (California, USA) constrained by temperature monitoring and time series analysis.
477 *Journal of Geophysical Research*, Vol. 114, B09201, doi:10.1029/2008JB006247.
478
- 479 Talling, P. J. 2013, Hybrid submarine flows comprising turbidity current and cohesive debris
480 flow: Deposits, theoretical and experimental analyses and generalized models,
481 *Geosphere*, Vol. 9, No. 3, doi: 10.1130/GES00793.
482
- 483 Tran, A., Rudolph, M.L., Manga, M. 2014, Bubble mobility in mud and yield stress fluids,
484 submitted to *Journal of Volcanology and Geothermal Research*.
485
- 486 Tregger, N.A., Pakula, M. E., Shah, S. P. 2010, Influence of clays on the rheology of cement
487 pastes. *Cement and Concrete Research*, Vol. 40, pp. 384-391.
488
- 489 Van Kessel T. and Blom C. 1998, Rheology of cohesive sediments: comparison between anatural
490 and an artificial mud. *Journal of Hydraulic Research*, Vol. 36, pp. 591–612
- 491 Yang, W.Y., Yu, G.-L, Tan, S., Wang, H.K. 2014, Rheological properties of dense natural
492 cohesive sediments subject to shear loadings. *International Journal of Sediment Research*,
493 Vol. 29, pp. 454-470.
494
- 495 Yu, B., Ma, Y., Qi, X. 2013, Experimental study on the influence of clay mineral on the yield
496 stress of debris flows. *Journal of Hydraulic Engineering*, Vol. 139, pp. 364-373.
497
498

499 **Figure 1.** a) The mean cumulative size distribution of the sediments. The hemipelagic sediment,
500 shown in blue, is the average of four tests, and the Salton Sea sediment, shown in red, is the
501 average of two tests. b) 3-D microtomography reconstructions of the materials. Reconstruction of
502 the hemipelagic sediment, which is poorly sorted and contains large fossils. C) Reconstruction of
503 the Salton Sea sediment, which is much more uniform in size than the hemipelagic sediment.
504

505 **Figure 2.** a) Yield stress as a function of particle concentration from multiple different studies,
506 b) consistency as a function of particle concentration from multiple different studies. Clay
507 contents shown in numbers next to the respective data when reported. For our measurements,
508 open symbols show static yield stress and filled symbols show dynamic yield stress. All data in
509 Table 1 and Supplementary Table 1.
510

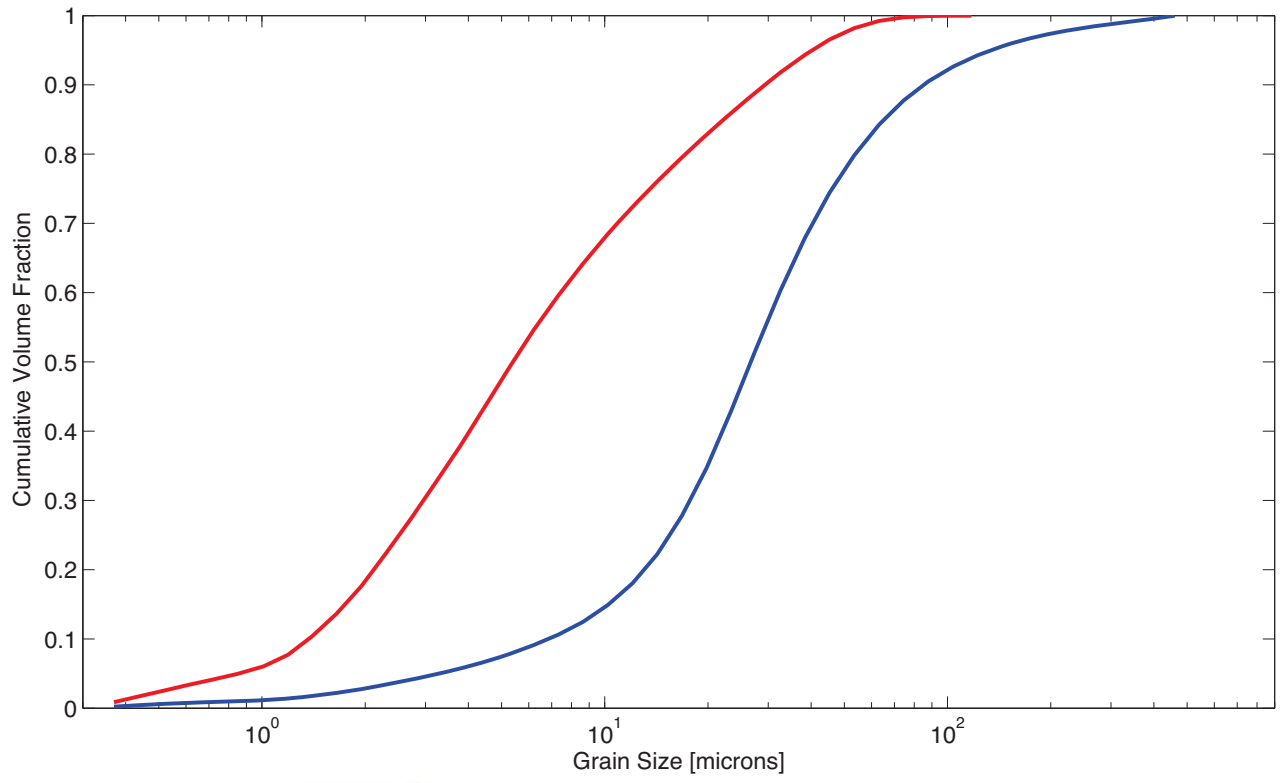
511 **Table 1.** The hemipelagic sediment and Salton Sea sediment measured in this study. Includes the
512 parameters used to fit the Herschel-Bulkley model. Data plotted in Figure 2.
513

514 **Figure 3.** Particle concentration of hemipelagic sediment and the clast radius that would exceed
515 the yield stress parameter defined by Beris et al. (1985). For radii above the blue curve, clasts

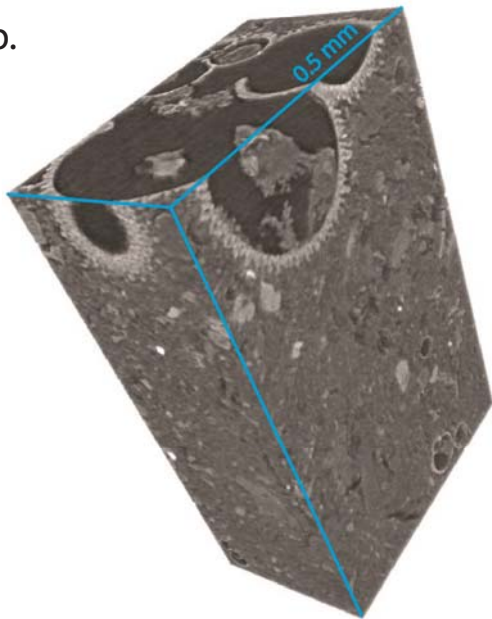
516 will settle through the hemipelagic marine sediment. Clasts are approximated as spheres of
517 density 1.65 g/cm^3 .
518

Figure 1.

a.



b.



c.

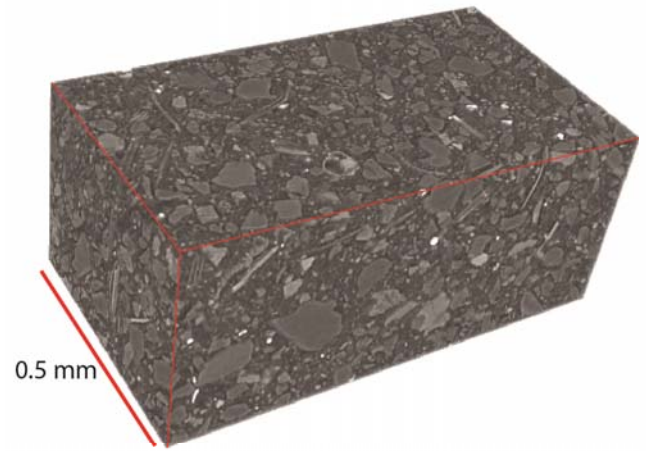
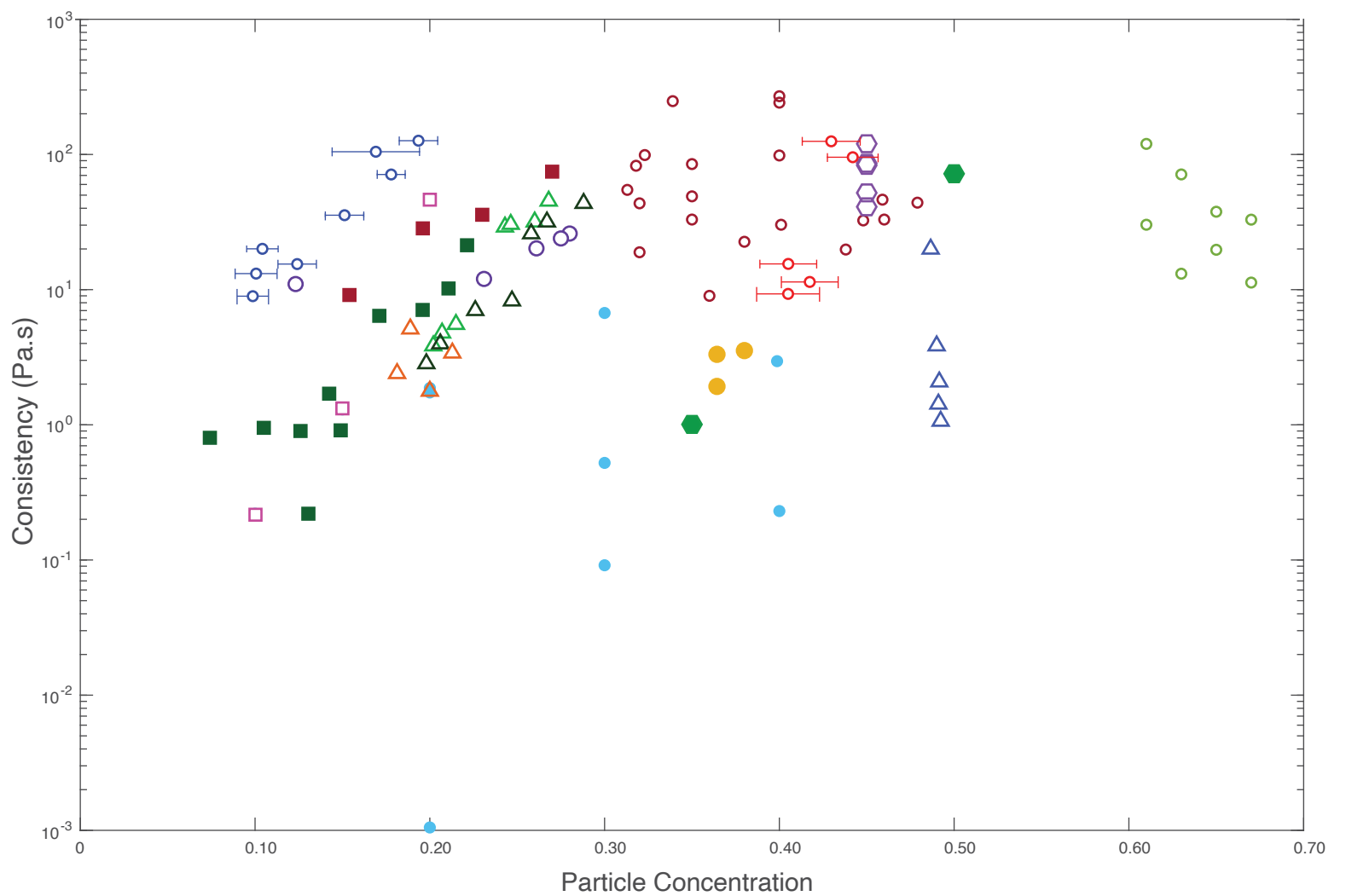
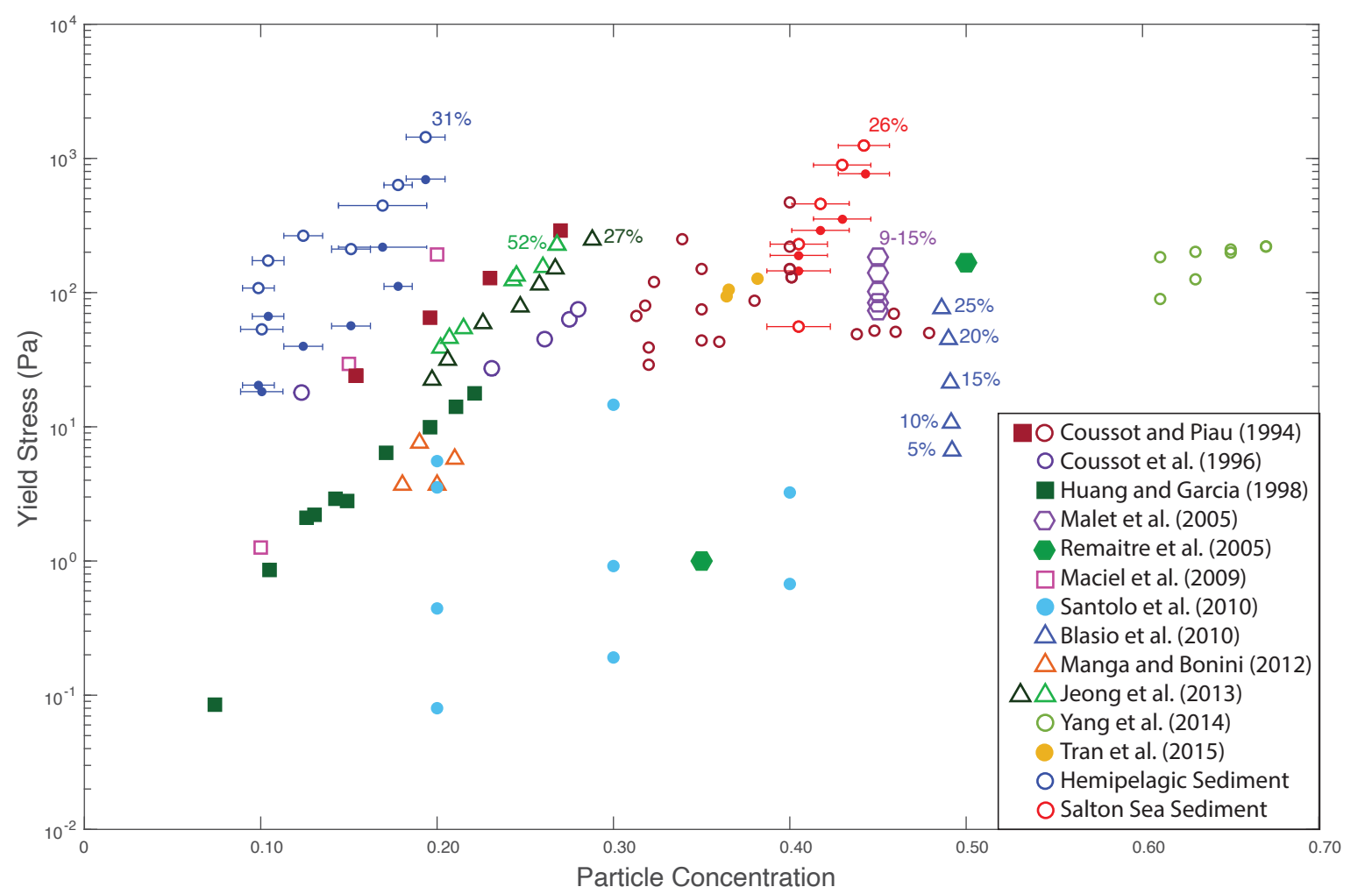


Figure 2.



Sample	Particle Concentration (%)	Dynamic Yield (Pa)	Consistency (Pa·s)	Static Yield (Pa)
Hemipelagic Marine Sediment	19.4 ± 1.1%	701.9 ± 2.6	126.4 ± 5.8	1440.9 ± 1.5
	17.8 ± 0.8%	111.4 ± 2.8	71.1 ± 4.1	634.2 ± 17.4
	16.9 ± 2.5%	218.1 ± 2.1	104.8 ± 3.5	445.3 ± 13.9
	12.4 ± 1.1%	39.8 ± 0.2	15.5 ± 0.6	265.3 ± 3.2
	15.1 ± 1.1%	56.3 ± 0.7	35.5 ± 2.2	211.1 ± 9.4
	10.4 ± 0.9%	66.5 ± 0.3	20.0 ± 0.9	173.1 ± 5.6
	9.9 ± 0.9%	20.5 ± 0.6	8.9 ± 1.4	108.1 ± 7.3
	10.1 ± 1.2%	18.3 ± 0.5	13.1 ± 1.0	53.3 ± 3.9
Salton Sea Mud Volcano	44.2 ± 1.45%	769.1 ± 3.0	95.4 ± 22.3	1247.8 ± 2.6
	42.9 ± 1.6%	353.6 ± 0.7	125.1 ± 1.9	893.4 ± 0.2
	41.7 ± 1.6%	291.1 ± 0.8	11.4 ± 4.7	458.3 ± 0.9
	40.5 ± 1.6%	189.4 ± 0.9	15.5 ± 8.2	229.7 ± 11.1
	40.5 ± 1.8%	145.3 ± 0.5	9.3 ± 4.9	55.8 ± 20.1

Density (g/cm ³)	% Clay	n
1.34	31	3.5
1.31	31	3.5
1.29	31	3.5
1.22	31	3.5
1.27	31	4.5
1.18	31	3.5
1.17	31	3.6
1.18	31	3.5
1.76	26	3.5
1.73	26	3.5
1.7	26	3.5
1.67	26	3.5
1.68	26	3.5

Figure 3.

